

Threshold production of top-quark pairs at NNNLO

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based on work in collaboration with
M. Beneke, Y. Kiyo, P. Marquard,
A. Penin, T. Rauh, D. Seidel, M. Steinhauser

Threshold Production

relative velocity of quark-antiquark pair is small:

$$\sqrt{s} = E + 2m_Q \approx 2m_Q \quad \Rightarrow \quad v = \sqrt{\frac{E}{m_Q}} \ll 1; \quad v \sim \alpha_s(m_Q v)$$

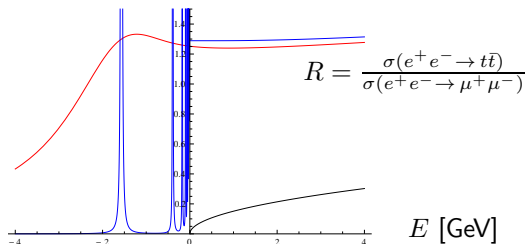
- multi-scale problem: mass m_Q , momentum $m_Q v$, energy $m_Q v^2$
- perturbation theory breaks down due to terms proportional to $\frac{\alpha_s}{v}$
 \rightsquigarrow Coulomb resummation
- formation of bound states below threshold

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 \rightsquigarrow Coulomb resummation
- formation of bound states below threshold
- $b\bar{b}$: bound-state resonances
- $t\bar{t}$: large width prevents existence of bound states



[Thacker, Lepage; Lepage, Magnea, Nakhleh, Magnea, Hornbostel; Bodwin, Braaten, Lepage]

[Pineda, Soto; Beneke, Signer, Smirnov; Brambilla, Pineda, Soto, Vairo]

scale hierarchy: $m_t \gg m_tv \gg m_tv^2 \gg \Lambda_{\text{QCD}}$

QCD

full theory

Effective Theory Setup

[Thacker, Lepage; Lepage, Magnea, Nakhleh, Magnea, Hornbostel; Bodwin, Braaten, Lepage]
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integrate out hard modes: $k^0 \sim k^i \sim m_t$
hard subgraphs become point-like vertices
 \rightsquigarrow hard matching coefficients

NRQCD

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NRQCD



integrate out soft modes: $k^0 \sim k^i \sim m_tv$
soft subgraphs become instantaneous, non-local vertices
 \rightsquigarrow potentials

pNRQCD

contains potential quarks and ultrasoft gluons

Calculating the Cross Section

use optical theorem:

$$\int d\Phi_{\text{PS}} \left| \text{Im} \left[\text{Diagram with wavy line and green arrows} \right] \right|^2 \sim \text{Im} \left[\text{Diagram with wavy line and green loop} \right] \sim \text{Im} \Pi(q)$$

compute polarisation function in pNRQCD: $\Pi(q) \sim c_v^2 G(\vec{0}, \vec{0}; E)$

- compute **hard matching coefficients**
- compute **potential**
- solve Schrödinger equation
- compute ultrasoft corrections

power counting:

$$\sum_n \left(\frac{\alpha_s}{v} \right)^n \times \left\{ 1; \underbrace{\alpha_s, v}_{\text{NLO}}; \underbrace{\alpha_s^2, \alpha_s v, v^2}_{\text{NNLO}}; \underbrace{\alpha_s^3, \alpha_s^2 v, \alpha_s v^2, v^3}_{\text{NNNLO}}; \dots \right\}$$

- NNLO

[Hoang, Teubner; Melnikov, Yelkhovsky; Yakovlev; Beneke, Signer, Smirnov;
Nagano, Ota, Sumino; Penin, Pivovarov]

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 - potential contributions [Beneke, Kiyo, Schuller; Kniehl, Penin, Smirnov, Steinhauser]
 - 3-loop static potential [Anzai, Kiyo, Sumino; Smirnov, Smirnov, Steinhauser]
 - ultrasoft corrections [Beneke, Kiyo, Penin]
 - 3-loop contribution to c_v [Marquard, JP, Seidel, Steinhauser]
 - P-wave contribution [Penin, Pivovarov; Beneke, JP, Rauh]

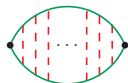
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- NNLL
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 - pNRQCD [Pineda, Signer]
- electroweak corrections [Guth, Kühn; Hoang, Reißer; Eiras, Steinhauser; Kiyo, Seidel, Steinhauser]
- finite-width effects [Fadin, Khoze; Hoang, Reißer, Ruiz-Femenía; Beneke, Jantzen, Ruiz-Femenía; Penin, JP]
- Higgs contribution [Harlander, Ježabek, Kühn; Eiras, Steinhauser; Beneke, JP, Rauh]

Computing the Green's Function

Coulomb Green's function at the origin:

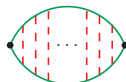
$$G_C(E) = \langle 0 | \frac{1}{H_0 - E} | 0 \rangle = \langle 0 | G^{(0)} | 0 \rangle$$



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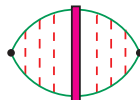
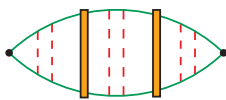
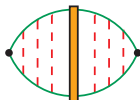
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corrections due to higher order potential can be computed perturbatively:

$$\begin{aligned} \delta G(E) &= \langle 0 | \frac{1}{H_0 + \delta V - E} | 0 \rangle - G_C(E) \\ &= -\langle 0 | G^{(0)} \delta V^{(1)} G^{(0)} | 0 \rangle \\ &\quad + \langle 0 | G^{(0)} \delta V^{(1)} G^{(0)} \delta V^{(1)} G^{(0)} | 0 \rangle - \langle 0 | G^{(0)} \delta V^{(2)} G^{(0)} | 0 \rangle + \dots \end{aligned}$$



Extracting Divergences

- divergent diagrams have to be calculated in $D = 4 - 2\epsilon$
- Coulomb Green function only known in $D = 4$
- adding Coulomb exchanges decreases degree of divergence

↪ subtract divergent graphs in $D = 4$ and calculate them in $D = 4 - 2\epsilon$

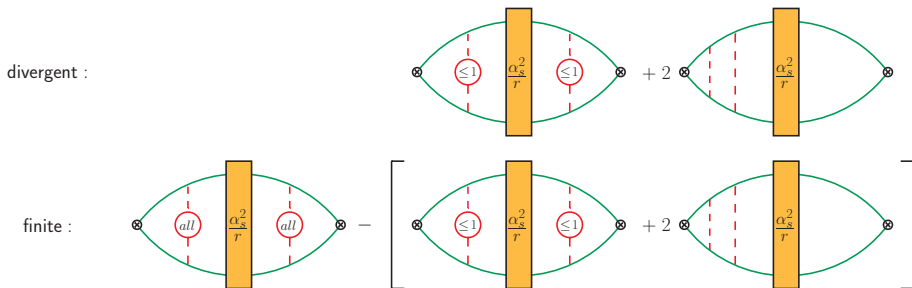
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example: NLO P-wave contribution

[Beneke, JP, Rauh, 2013]



- $E \rightarrow E + i\Gamma_t$ at LO $\rightsquigarrow \Gamma_t/\epsilon$ at NNLO
- finite-width divergences lead to scale dependence
- dependence does not decrease in higher orders
- divergences are cancelled by non-resonant contributions

example: LO P-wave Green function

[Beneke, JP, Rauh, 2013]

$$G_0^P(E) = \frac{m_t^4}{4\pi} \left\{ \frac{\alpha_s C_F}{2} \left[\frac{1}{2\epsilon} - \ln \left(\frac{-4m_t E}{\mu_w^2} \right) + 4 \right] \frac{E}{m_t} + \dots \right\}$$

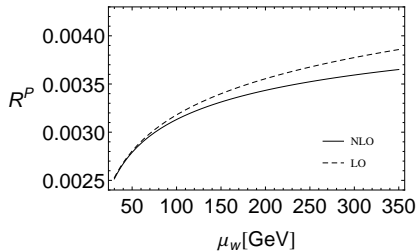
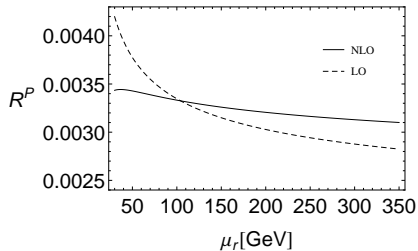
$$\rightsquigarrow \text{Im}G_0^P(E + i\Gamma_t) = \frac{m_t^4}{4\pi} \left\{ \frac{\alpha_s C_F}{2} \left[\frac{1}{2\epsilon} - \ln \left(\frac{4m_t \sqrt{E^2 + \Gamma_t^2}}{\mu_w^2} \right) + 4 \right] \frac{\Gamma_t}{m_t} + \dots \right\}$$

Finite-Width Divergences

- $E \rightarrow E + i\Gamma_t$ at LO $\rightsquigarrow \Gamma_t/\epsilon$ at NNLO
- finite-width divergences lead to scale dependence
- dependence does not decrease in higher orders
- divergences are cancelled by non-resonant contributions

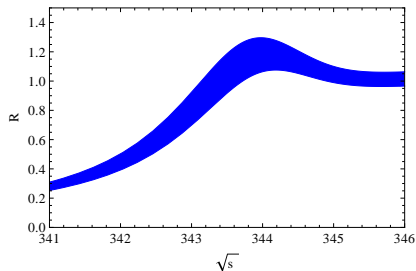
example: P-wave contribution close to threshold

[Beneke, JP, Rauh, 2013]

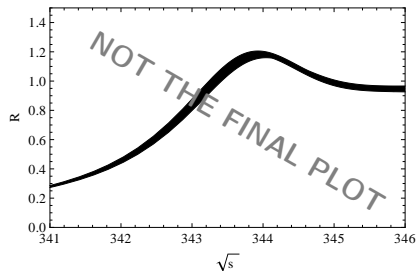


(effect is less pronounced for S-wave)

NNLO:

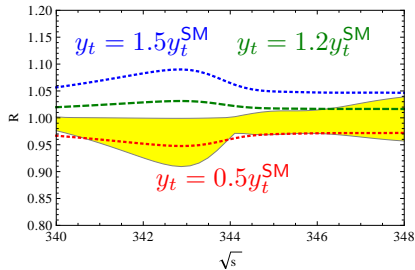
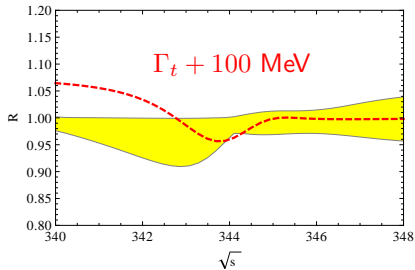
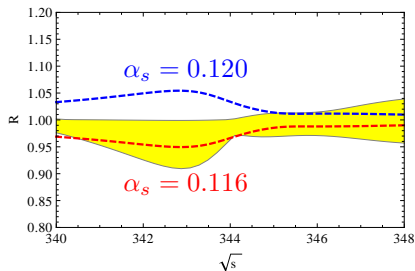
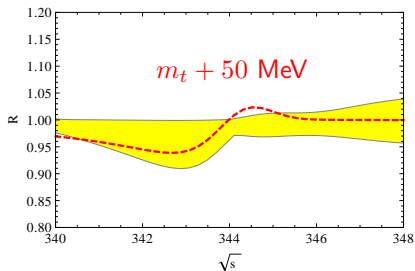


NNNLO:



PS scheme: $m_t^{\text{PS}} = 171.5$ GeV, $\Gamma_t = 1.33$ GeV, $\alpha_s(M_Z) = 0.118$
 $\mu = 50 - 350$ GeV

Parameter Dependence (preliminary)



- threshold production allows for very precise and theoretically clean determination of top-quark mass
- very high theoretical precision is required
- NNNLO pNRQCD calculation is almost done

lots of work still to be done, e.g.:

- electroweak/finite-width corrections
- combination of fixed-order and resummed calculations