Non-SUSY BSM Theory

Markus Luty UC Davis

Reasons to believe in new physics accessible to LHC:

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"The SUSY train is late." — G. Altarelli, 2001

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Example: $m_{\tilde{t}} \sim 100 \text{ TeV}$ \Rightarrow threshold correction $\Delta m_h^2 \sim 10^5 \times (125 \text{ GeV})^2$

S. Chang, ML, unpublished

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Compute loop effects in terms of observable quantities \Rightarrow no quadratic dependence on heavy masses.

Example: loop corrections to $hh \rightarrow hh$

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Optical theorem: Im $= \underbrace{\frac{1}{t}}_{-\frac{1}{t}} + \cdots$ $= \underbrace{\frac{y_t^2 m_h^2}{s}}_{-\frac{y_t^2 m_h^2}{s}} - y_t^4$

Use to compute loop corrections to $hh \rightarrow hh$:

$$t \to 0$$
: $\mathcal{A}(s) = \frac{1}{\pi} \int_0^\infty ds' \frac{\operatorname{Im} \mathcal{A}(s')}{s' - s} = \log \operatorname{UV} \operatorname{divergent}$

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 \Rightarrow can write once-subtracted dispersion relation

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My conclusion: think harder!

The standard model is incomplete.

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Gauge unification, ν masses, flavor, ... have simplest explanation in terms of physics at scales \gg TeV.

* If the result holds up.

Naturalness requires a mechanism to prevent physics at high scales from contributing to m_h^2 .

$$\Delta m_h^2 = \cdots \sim \sqrt[]{\lambda^2 m_\chi^2} \sim \frac{\lambda^2 m_\chi^2}{16\pi^2}$$

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$$\Delta m_h^2 = \cdots \sim \sum_{h=1}^{N} \sim \frac{\lambda \Lambda^2}{16\pi^2}$$

Y

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Motivates $\Lambda \sim \text{TeV}$

 $m_h = 125 \text{ GeV}$ $\frac{g_{hVV}}{g_{hVV}^{(SM)}} = 1 + O(10\%)$

⇒ Higgs VEV dominates electroweak symmetry breaking

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Georgi, Kaplan, 1984 : Contino, Nomura, Pomarol, 2003 :



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v = f (Technicolor)



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 \Rightarrow precision Higgs coupling measurements directly probe tuning in PNGB Higgs models

Top quark is a potential additional source of tuning:



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BMSSM Higgs

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Increase Higgs quartic:

- F terms (NMSSM, ...)
- *D* terms (new gauge interactions)

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Another possibility:

• Higgs tadpole from "auxiliary" Higgs sector

Azatov, Galloway, ML, 2012



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If auxiliary sector is at a strong conformal fixed point, SUSY breaking triggers confinement and electroweak symmetry breaking at the TeV scale.

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$$V_{\text{eff}} \simeq m_H^2 H^{\dagger} H + \kappa H^{\dagger} e^{i\Pi/f} \begin{pmatrix} 0\\ f \end{pmatrix} + \text{h.c.}$$

$$m_H^2 > 0 \quad \Rightarrow \quad v = \frac{\kappa_J}{m_H^2}$$

"induced EWSB"

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Signals:

• Non-minimal Higgs signals

A → Zh, tt, ττ H → WW, hh

 $\rho \rightarrow WW, AA, \ldots$

Chang, ML, Salvioni, Tsai, to appear

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Induced EWSB

Galloway, ML, Tsai, Zhao, 2014

Can also construct perturbative calculable models



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General message: SUSY naturalness motivates BMSSM Higgs

It exists. It is BSM.

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 \Rightarrow search for effective interactions coupling SM with SM

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Bai, Fox, Harnik, 2010 Beltran, Hooper, Kolb, Krusberg, Tait, 2010

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⇒ monojet signal



Generalized to mono-X searches ...

Chang, Hutchinson, Edezath, ML 2013 An, Wang, Zhang, 2013 Bai, Berger, 2013

Relic abundance motivates *renormalizable* couplings to SM.

$$\Omega_{\rm DM} \sim 0.1 \left(\frac{\sigma_{\rm ann}}{\rm pb}\right)^{-1}$$
 $\sigma_{\rm ann} \sim \frac{g^4}{m_{\rm DM}^2} \sim \rm pb$ for $m_{\rm DM} \sim {\rm TeV}$

The "WIMP miracle"

Chang, Hutchinson, Edezath, ML 2013 An, Wang, Zhang, 2013 Bai, Berger, 2013

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Fix coupling by requiring correct relic abundance \Rightarrow parameterized by m_{χ} , m_Q

Majorana fermion dark matter



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