Non-SUSY BSM Theory

Markus Luty UC Davis

Reasons to believe in new physics accessible to LHC:

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"The SUSY train is late." — G. Altarelli, 2001

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 m_h^2 *h* **=** 44848354663004959003564458711382292 GeV² -44848354663004959003564458711366667 GeV² $= (125 \text{ GeV})^2$

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Example: $m_{\tilde{t}} \sim 100$ TeV \Rightarrow threshold correction $\Delta m_h^2 \sim 10^5 \times (125 \text{ GeV})^2$

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Compute loop effects in terms of observable quantities **⇒** no quadratic dependence on heavy masses.

Example: loop corrections to $hh \rightarrow hh$

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Im Optical theorem: Im $\left(\begin{array}{ccc} \bullet & \bullet & \bullet \\ \bullet & \bullet & \bullet \end{array} \right)$ =

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= + + *···* $\overbrace{}^{x }$ **⇠** *y*2 *^t ^m*² *h s* \longrightarrow $\sim y_t^4$ *t* Im *t* ¯*t* **Optical theorem:**

Use to compute loop corrections to $hh \rightarrow hh$:

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⇒ can write once-subtracted dispersion relation

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My conclusion: think harder!

The standard model is incomplete.

• Not UV complete

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Gauge unification, ν masses, flavor, ... have *simplest* explanation in terms of physics at scales \gg TeV.

* If the result holds up.

Naturalness requires a mechanism to prevent physics at high scales from contributing to m_h^2 .

$$
\Delta m_h^2 = -\sqrt{\frac{X}{\sqrt{16\pi^2}}} \approx \frac{\lambda^2 m_X^2}{16\pi^2}
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SUSY: $\Delta m_h^2 = -4 \sum_{\substack{n=-\infty \\ n \text{ odd}}}^X$ + \cdots $\left(\sum_{\substack{n=-\infty \\ n \text{ odd}}}^X (m_x^2 - m_{\tilde{x}}^2)\right)$

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Two mechanisms:

SUSY:
$$
\Delta m_h^2 = -2\pi \sum_{\dot{x}} m_{\dot{x}}^2 + \frac{1}{2} m_{\dot{x}}^2 + \cdots \approx \frac{\lambda^2}{16\pi^2} (m_x^2 - m_{\dot{x}}^2)
$$

 $\tilde{\mathsf{X}}$

Compositeness:
$$
\Delta m_h^2 = -4. \sum_{n=1}^{\infty} m_m^2
$$
 $\sim \frac{\lambda h^2}{16\pi^2}$

 \boldsymbol{X}

 Λ = scale of form factors

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 $\tilde{\mathbf{v}}$

Compositeness:
$$
\Delta m_h^2 = -\mu \sum_{n=1}^{N} \frac{\lambda h^2}{16\pi^2}
$$

 \checkmark

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SUSY:
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\Delta m_h^2 = -2 \sum_{\substack{\lambda \vdots \\ \lambda \vdots \\
$$

 \sim

$$
mpositeness: \quad \Delta m_h^2 = -\sqrt{2} \quad \sqrt{2} \quad \sqrt{2} \quad \sqrt{16 \pi^2}
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SUSY:
$$
\Delta m_h^2 = -4 \sum_{\text{min}}^{\text{max}} \sqrt{m_x^2 + m_y^2 + \cdots + m_z^2}
$$

 \sim

Compositeness:
$$
\Delta m_h^2 =
$$
 $\therefore \qquad \qquad \searrow$ $\qquad \qquad \frac{\lambda \Lambda^2}{16 \pi^2}$

 $\sqrt{ }$

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Motivates $\Lambda \sim \text{TeV}$

 $m_h = 125$ GeV *ghVV g*(SM) *hVV* **=** 1 **+** *O***(**10%**)**

⇒ Higgs VEV dominates electroweak symmetry breaking

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Georgi, Kaplan, 1984 Contino, Nomura, Pomarol, 2003
.
. . .

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 $v = f$ (Technicolor)

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⇒ precision Higgs coupling measurements directly probe tuning in PNGB Higgs models

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BMSSM Higgs

Naturalness ⇒ Higgs mass too small in MSSM

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Another possibility:

• Higgs tadpole from "auxiliary" Higgs sector

Azatov, Galloway, ML, 2012

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$$
V_{\text{eff}} \simeq m_H^2 H^{\dagger} H + \kappa H^{\dagger} e^{i\Pi/f} \begin{pmatrix} 0 \\ f \end{pmatrix} + \text{h.c.}
$$

$$
m_H^2 > 0
$$
 \Rightarrow $v = \frac{N_H}{m_H^2}$

"induced EWSB"

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Yukawa couplings to elementary Higgs

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Yukawa couplings to elementary Higgs

Signals:

• Non-minimal Higgs signals

 $A \rightarrow Zh$, tt, $\tau\tau$

 $H \rightarrow WW, hh$ $\rho \rightarrow W W$, AA, . . .

Chang, ML, Salvioni, Tsai, to appear \sim 120 (350) GeV in the model of Section 3. The dot-dashed black line corresponds Chang, ML, Salvioni, isal, to appear

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Galloway, ML, Tsai, Zhao, 2014

Can also construct perturbative calculable models

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General message: SUSY naturalness motivates BMSSM Higgs

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⇒ search for effective interactions coupling SM with SM

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Generalized to mono-X searches *...*

Chang, Hutchinson, Edezath, ML 2013 An, Wang, Zhang, 2013 Bai, Berger, 2013

Relic abundance motivates renormalizable couplings to SM.

$$
\Omega_{DM} \sim 0.1 \left(\frac{\sigma_{\text{ann}}}{pb} \right)^{-1} \qquad \sigma_{\text{ann}} \sim \frac{g^4}{m_{DM}^2} \sim pb \qquad \text{for } m_{DM} \sim \text{TeV}
$$

The "WIMP miracle"

Chang, Hutchinson, Edezath, ML 2013 An, Wang, Zhang, 2013 Bai, Berger, 2013

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Relic abundance motivates renormalizable couplings to SM.

Fix coupling by requiring correct relic abundance \Rightarrow parameterized by m_X , m_Q

Majorana fermion dark matter

Conclusions

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