



Beam Instrumentation

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CAS
Introductory level course on
Accelerator Physics

Introduction

- Beam Instrumentation is a very wide subject; with a large range of technologies and fields involved, including:
 - Accelerator physics
 - understand the beam parameters to be measured
 - distinguish beam effects from sensor effects
 - Particle physics and detector physics
 - understand the interaction of the beam with the sensor
 - RF technology
 - Optics
 - Mechanics
 - Electronics
 - Analogue signal treatment
 - Low noise amplifiers
 - High frequency analogue electronics
 - Digital signal processing
 - Digital electronics for data readout
 - Software engineering
 - Front-end and Application Software

Introduction, cont'd

- Harsh environment:
 - Radiation (single event effects, radiation ageing, activation)
 - Many sources of measurement noise and background
 - Place readout close to detector, but → radiation
 - RF heating by the beam
 - Accessibility and maintenance
 - Sometimes: cryogenic temperatures
 - Mostly: must operate in vacuum and be UHV compatible

Introduction, cont'd

- Aim: assist in commissioning, tuning and operating the accelerator and to improve performance → see tomorrow
- Some instrument classifications:
 - **LINAC and transport lines:** Single pass, can have separate measurement lines ↔ **Synchrotron:** multi pass
 - **Total Beam Energy** (beam particles x particle energy) low ↔ high
 - **Non-intercepting** ↔ **Intercepting** ↔ **Destructive** (often depending on beam energy)
- **In this presentation:**
 - Explain working principles of some of the most important instruments
 - Give indication on achievable performance
 - Give selected examples from operating machines and current developments

Resources and References

- Peter Forck: *Lecture on Beam Instrumentation and Diagnostics* at the Joint University Accelerator School (JUAS), see also the extended Bibliography.
<http://www-bd.gsi.de/conf/juas/juas.html>
- CERN Accelerator Schools (CAS):
<http://cas.web.cern.ch/cas/CAS%20Welcome/Previous%20Schools.htm> and
http://cas.web.cern.ch/cas/CAS_Proceedings.html
 - Rhodri Jones and Hermann Schmickler: *Introduction to Beam Instrumentation and Diagnostics*, CERN-2006-002.
 - Daniel Brandt (Ed.), 2008 CAS on *Beam Diagnostics for Accelerators*, Dourdan, CERN-2009-005 (2009).
 - Heribert Koziol, *Beam Diagnostic for Accelerators*, Loutraki, Greece (2000), CERN/PS 2001-012 (DR), see also extended Bibliography.
- Jacques Bosser (Ed.), *Beam Instrumentation*, CERN-PE-ED 001-92, Rev. 1994

Measured Quantities

- **Beam intensity**
- Ideally: 6D phase space of the beam
- Real measurements: mean values and 1D-projection, some 2D-projections
 - **Transverse position** (mean x, y) → **trajectory and orbit**
 - **Transverse profile**
 - Bunch length, **bunch shape**
 - Mean momentum and momentum spread
 - **Emittance and 2D phase space reconstruction (transverse and longitudinal)**
 - Beam halo measurements
- **Tune**, chromaticity, coupling, beta function, dispersion
- **Beam Losses**
- Polarisation
- Luminosity

Classification of Selected Devices

- Different devices (techniques) to measure the same quantity ↔
Same device to measure different quantities

PROPERTY MEASURED →	Intensity/charge	tr. Position	tr. Size/shape	tr. Emittance	Effect on beam		
					N	+	D
Current transformers	●				x		
Pick-ups	●	●			x		
Faraday cup	●						x
Secondary emission monitors	●	●	●	●	x	x	
Wire scanners		●	●	●	x		
Scintillator screens		●	●		x	x	
OTR screen		●	●		x	x	
Residual-gas profile monitors		●	●	●	x		
Beam loss monitors					x		

Effect on beam depends on circumstances
 N none
 - slight
 + perturbing
 D destructive

- Different Labs have different names for the same device!

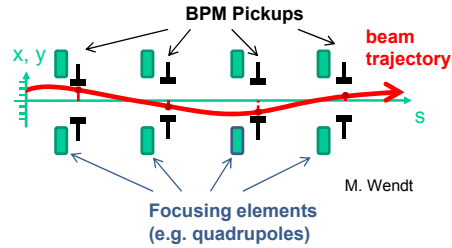
Beam Position Monitors

Capacitive Pick-Ups for Bunched Beams

- Among the most numerous instruments

- **Measurements:**

- Transverse beam position (typically next to focusing elements)
 - Beam trajectory or closed orbit
 - injection oscillations
- Tune and lattice function in synchrotrons

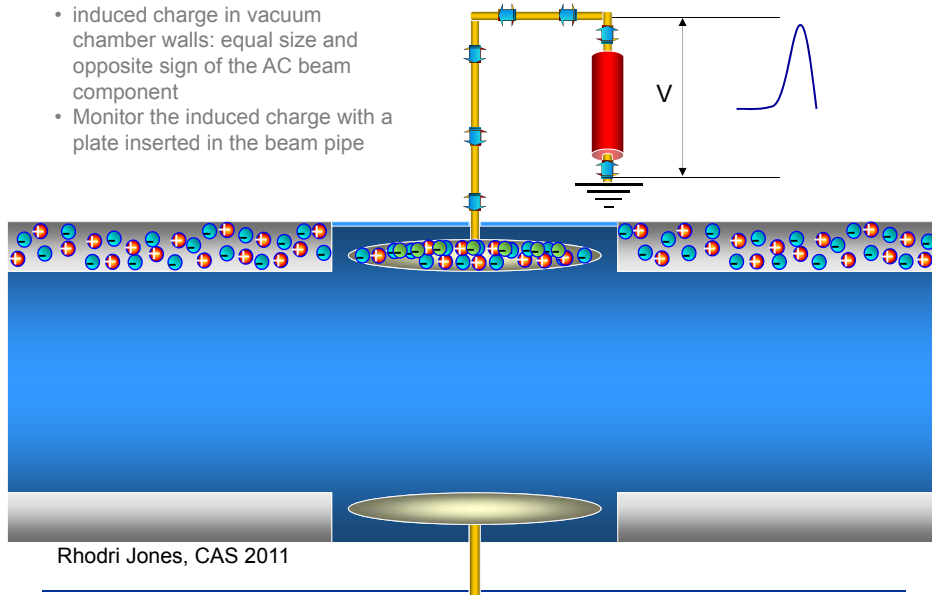


- **Working principle:**

- Image current in vacuum chamber walls: equal size and opposite sign of the AC beam component
- Monitor the induced charge with a plate inserted in the beam pipe
- **Adapt BPM electronics integration time**
 - Single-bunch ↔ multi-bunch
 - turn-by-turn (single pass) ↔ multi-turn average

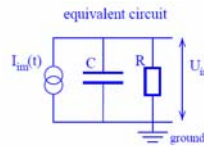
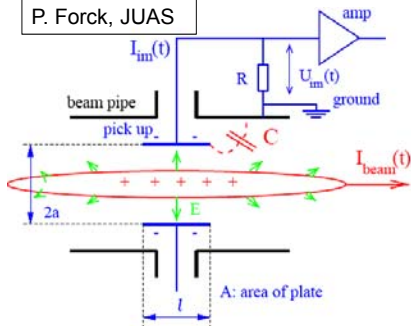
Capacitive Pick-Up – The Principle

- induced charge in vacuum chamber walls: equal size and opposite sign of the AC beam component
- Monitor the induced charge with a plate inserted in the beam pipe



Schematics and Simplified Equivalent Circuit

P. Forck, JUAS



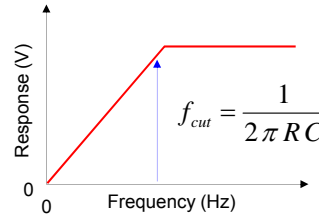
$$U_{im}(\omega) = \frac{R}{1 + i\omega RC} \cdot I_{im}(\omega)$$

$$U_{im}(\omega) = \frac{A}{2\pi a} \cdot \frac{1}{\beta c} \cdot \frac{1}{C} \cdot \frac{i\omega RC}{1 + i\omega RC} \cdot I_{beam}(\omega)$$

$$\equiv Z_t(\omega, \beta) \cdot I_{beam}(\omega)$$

Z_t ... longitudinal transfer impedance

⇒ High pass characteristics with a cut-off frequency, f_{cut}



$$I_{im} = \frac{A}{2\pi a l} \left(-\frac{l}{\beta c} \frac{dI_{beam}}{dt} \right) = \frac{A}{2\pi a l} \frac{1}{\beta c} i\omega I_{beam}(\omega)$$

frequency domain: $I_{beam} = I_0 e^{-i\omega t}$

- U_{im} ... voltage measured due to image current
- R ... amplifier input resistor
- ω ... frequency
- βc ... beam velocity

Transfer Impedance

$$|Z_t| = \frac{A}{2\pi a} \frac{1}{\beta c} \frac{1}{C} \frac{\omega/\omega_{cut}}{\sqrt{1 + \omega^2/\omega_{cut}^2}}$$

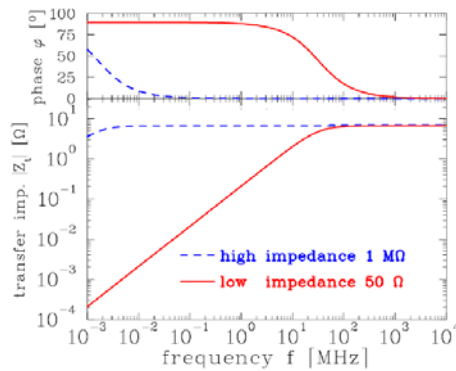
$$f_{cut} = \frac{1}{2\pi RC}$$

$R=50 \Omega \Rightarrow f_{cut} = 32 \text{ MHz}$

$R=1 \text{ M}\Omega \Rightarrow f_{cut} = 1.6 \text{ kHz}$

- large signal at lower frequencies
⇒ **high impedance**
- smooth signal transmission
⇒ **50 Ω**

P. Forck, JUAS



$l = 10 \text{ cm}$ long cylindrical pick-up with a capacitance of $C = 100 \text{ pF}$ and an ion velocity of $\beta = 50\%$

Operational Regime Determines Signal Shape

$$Z_t \propto \frac{i\omega/\omega_{cut}}{1 + i\omega/\omega_{cut}} \quad \frac{dI_{beam}}{dt} = -i\omega I_{beam} \quad U_{im}(\omega) = Z_t(\omega)I_{beam}(\omega)$$

▪ $\omega \ll \omega_{cut}$: $U_{im}(t) \propto \frac{dI_{beam}}{dt}$, phase shift $90^\circ \rightarrow$ derivative signal

▪ $\omega \gg \omega_{cut}$: $U_{im}(t) \propto I_{beam}(t)$, no phase shift

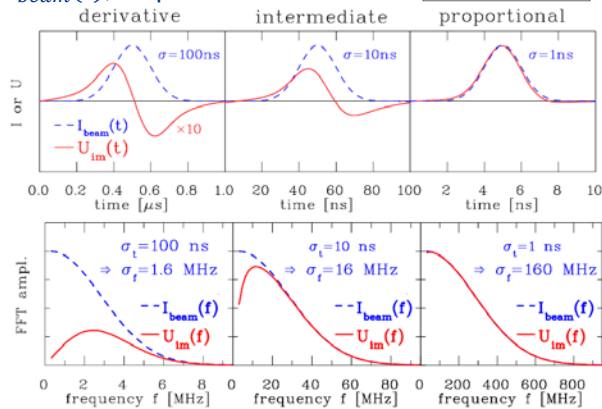
P. Forck, JUAS

$l = 10$ cm long cylindrical pick-up with a capacitance of $C = 100$ pF and an ion velocity of $\beta = 0.5$,

$R = 50 \Omega$

$f_{cut} = 32$ MHz

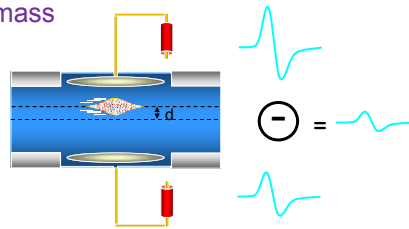
Operational regime depends on the bunch length!



Beam Position

- Signal on each plate is proportional to the beam intensity
- The difference signal (ΔU), top - bottom, or left - right, is proportional to the position of the beam center of mass
- Normalization to the sum signal (ΣU) gives the position:

$$x = \frac{1}{S_x(\omega, x, y)} \cdot \frac{\Delta U}{\Sigma U} \quad (S_x \dots \text{position sensitivity})$$

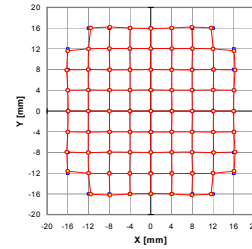
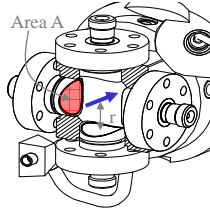


- The difference signal (ΔU) is normally at least a factor 10 lower than the sum signal (ΣU)
- Difficult to do electronically without some of the intensity information leaking through
- When looking for small differences this leakage can dominate the measurement
- Resolution for typical apertures:
 - \approx tens μm turn-by-turn
 - \approx μm multi-turn resolution

Example: Button Pick-up

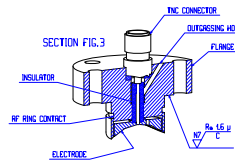
- ✓ Low cost ⇒ most popular
- ✗ Non-linear
 - requires correction algorithm when beam is off-centre

$$X = 2.30 \cdot 10^{-5} X_1^3 + 3.70 \cdot 10^{-5} X_1^3 + 1.035 X_1 + 7.53 \cdot 10^{-6} X_1^2 Y_1^2 + 1.53 \cdot 10^{-5} X_1 Y_1^4$$



R. Jones, CAS

LHC buttons



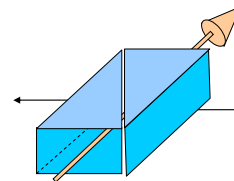
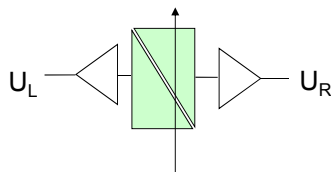
$$f_{cut} = \frac{1}{2\pi RC} = \frac{1}{2\pi \times 50\Omega \times 8pF} = 400MHz$$

$$Z_{loc} = \frac{A}{(2\pi a) \times c \times C_e} = \frac{\pi \times (12mm)^2}{(2\pi \times 24.5mm) \times c \times (8pF)} = 1.2\Omega$$

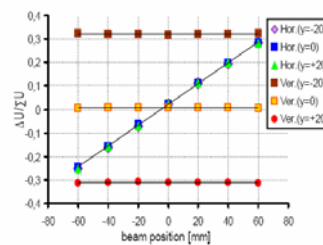
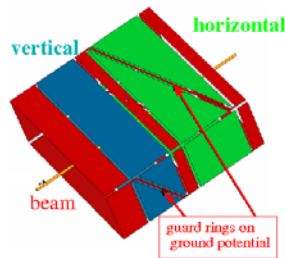
Shoebbox Pick-up

- ΔU gives linear position reading (no geometric correction)
- Condition: Linear cut: projection on the measurement plane must be linear:

$$X \propto \frac{U_L - U_R}{U_L + U_R} = \frac{\Delta U}{\Sigma U}$$

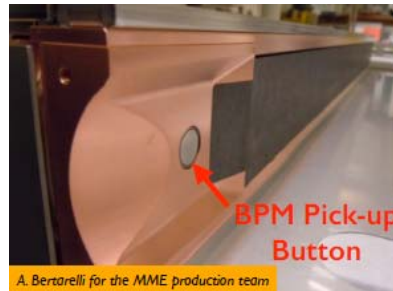


- Various geometries have been built, example from GSI optimization study (P.Kowina et al., DIPAC 2005)



New LHC Collimators with Integrated BPMs

- Beam-based setup currently with BLM signal → time consuming
- Tighter tolerances will be required for future LHC operation
- BPM integrated in the tapered end of the collimator jaws (10.6mm retraction from jaw surface)
 - Drastically reduce set-up time
 - Allow constant monitoring of beam position to jaw position
- **Successfully tested in the SPS**
(D. Wollmann, HB2012)
 - $25\ \mu\text{m}$ difference to BLM setup
- believed to be dominated by the BLM setup method
 - single pass (transfer line):
$90\ \mu\text{m}$ rms
 - no disturbance observed from protons hitting the jaws or from shower particles

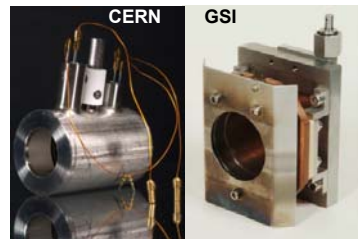
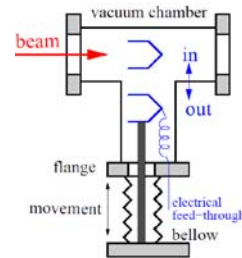


Beam Current

Faraday Cup

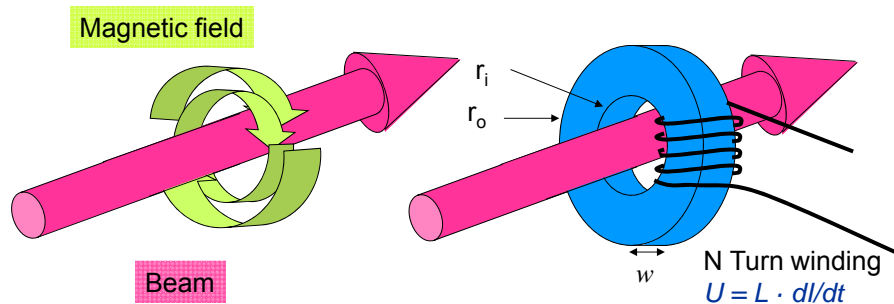
- Measurement of the beam's **electrical charges**
 - Low energies only
 - Particles are stopped in the device
→ **Destructive**
 - Sensitive to low currents: down to 1 pA can be measured
 - Creation of secondary electrons of low energy (below 20 eV)
 - Repelling electrode with some 100 V polarization voltage pushes secondary electrons back onto the electrode
 - Absolute accuracy:
 - $\approx 1\%$ (some monitors reach 0.1%)

Faraday Cup at GSI LINAC, P. Forck, JUAS



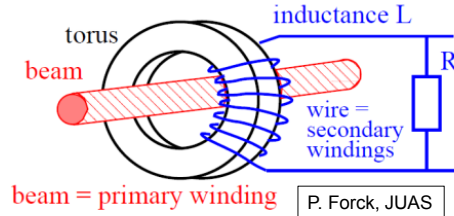
Beam Current Transformer (BCT)

- Measurement of the **magnetic field** of the beam
- **Non-interceptive**
- Independent on beam energy
- Beam as primary winding of a transformer



U. Raich, CAS

Current Transformers



P. Forck, JUAS

$$\text{Beam current } I_{\text{Beam}} = \frac{e N_q}{t} = \frac{e N_q \beta c}{w}$$

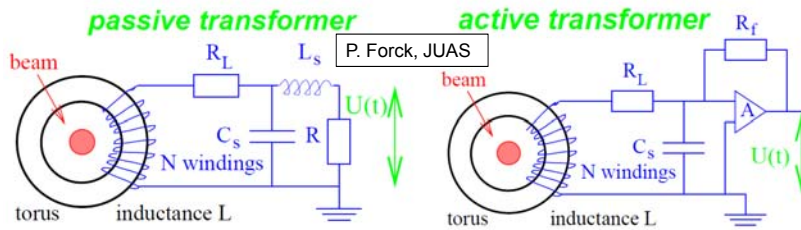
Transformer Inductance

$$L = \frac{\mu_0 \mu_r}{2\pi} w N^2 \ln \frac{r_0}{r_i}$$

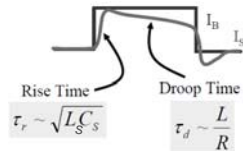
- Magnetic field of the beam is very low (Example: $1 \mu\text{A}$, $r = 10\text{cm} \Rightarrow 2 \text{ pT}$; compared to earth magnetic field of $\approx 50 \mu\text{T}$)

- Aim of the Torus:
 - Capture magnetic field lines with cores of high relative permeability
 - Signal strength nearly independent of beam position.
 - (CoFe based amorphous alloy Vitrovac: $\mu_r = 10^5$)

Adapt Droop Time with Active Transformer



P. Forck, JUAS



Bunch trains:



- Equal areas
- Baseline shift proportional to intensity

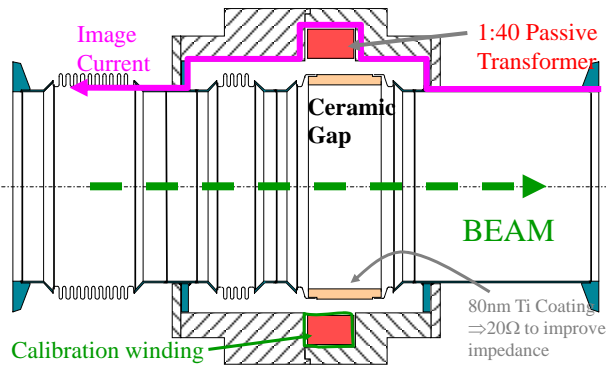
H. Koziol, CAS

- Use a trans-impedance amplifier (current-to-voltage converter) for observation of beam pulses $> 10 \mu\text{s}$, e.g. at pulsed LINAC
- Droop time constants of up to 1s
- Longer rise times as well (to reduce high frequency noise of the amplifier)

$$\tau_d = \frac{L}{R_f/A + R_L} \approx \frac{L}{R_L}$$

Transformer Housing

- Image current passing outside of the transformer torus
- High permeability material shields the transformer against external magnetic fields



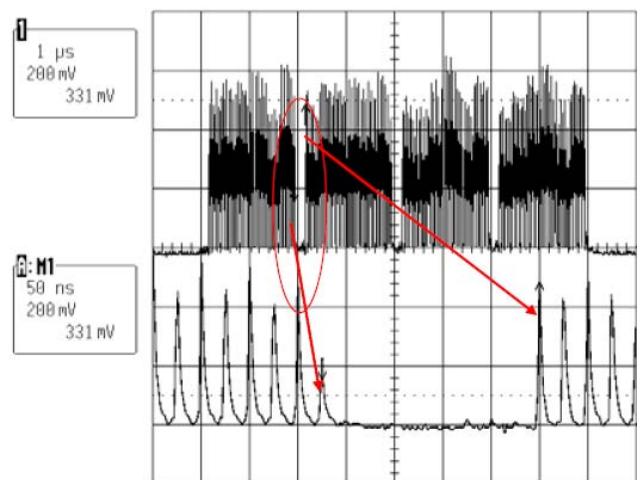
CERN SPS Fast Beam Current Transformer (FBCT)

R. Jones, DIPAC'03

500 MHz Bandwidth; Low droop (< 0.2%/ms)

CERN FBCT Readings of LHC Type Beams in the SPS

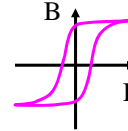
- 4 batches each containing 72 bunches separated by 25 ns



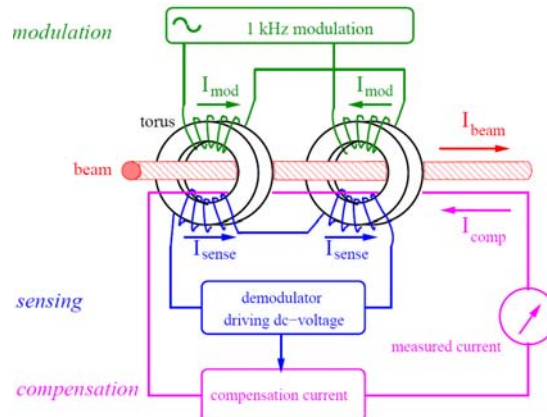
R. Jones, DIPAC'03

DCCT: DC Beam Current Transformer

- DC current $dB/dt = 0 \Rightarrow$ no voltage induced
- Use two **identical** toroids
- Take advantage of non-linear magnetisation curve



- **Modulation** of opposite sign drives toroids into saturation
- **Sense windings** measure the modulation signal
 - Signals from the two toroids cancel each other as long as there is no beam

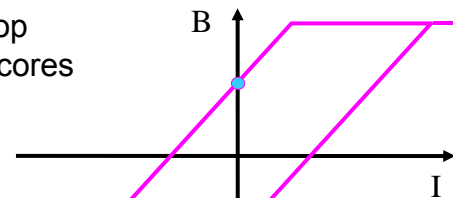


P. Forck, JUAS

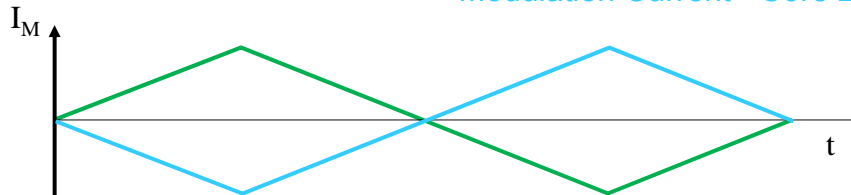
DCCT Principle – Case 1: No Beam

R. Jones, CAS

Hysteresis loop of modulator cores

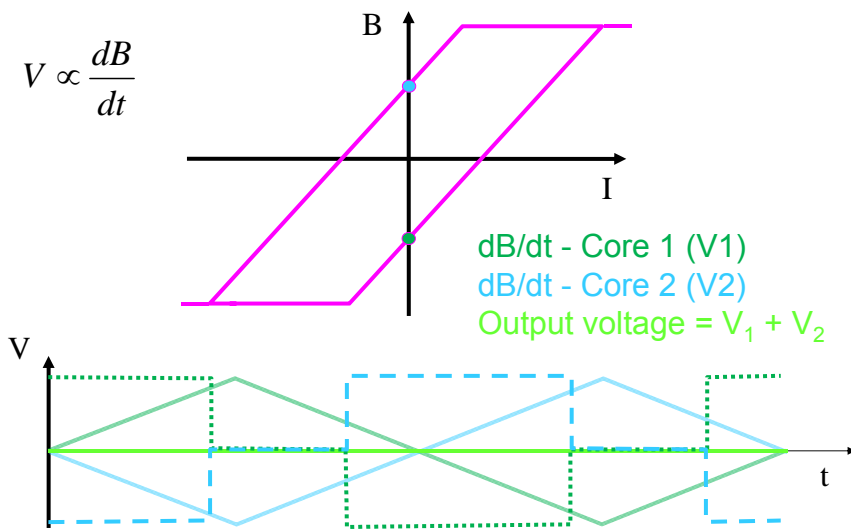


Modulation Current - Core 1
Modulation Current - Core 2



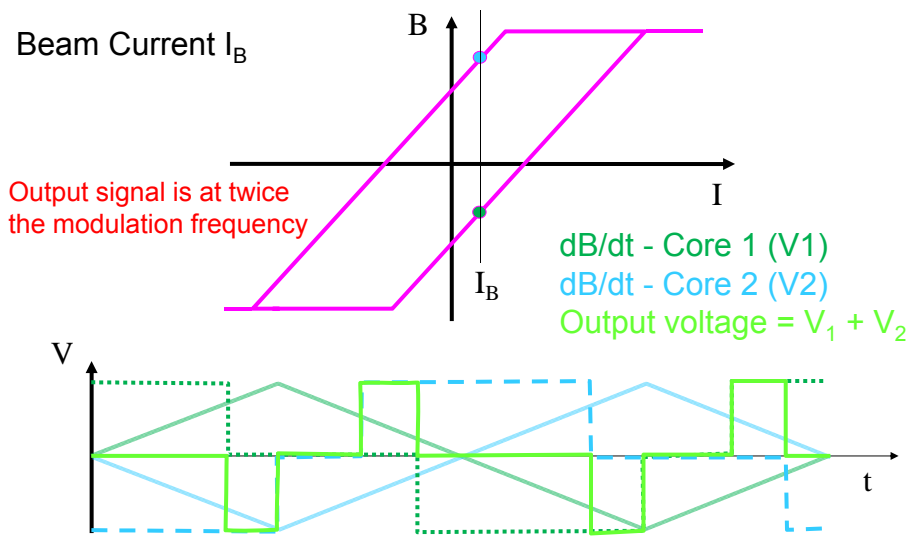
DCCT Principle – Case 1: No Beam

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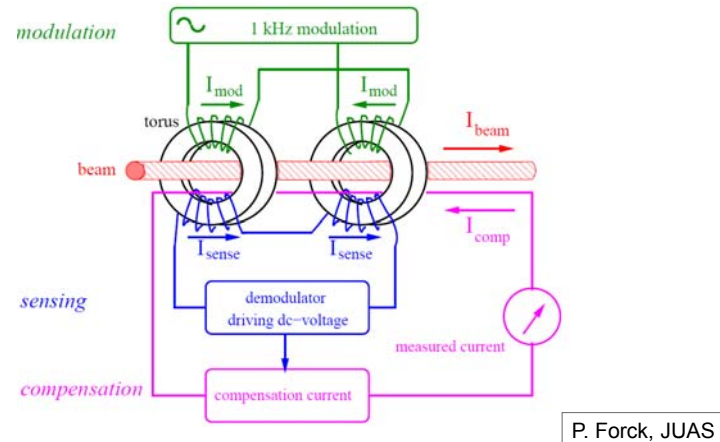
DCCT Principle – Case 2: With Beam

R. Jones, CAS



DCCT in the “Zero Flux” Scheme

- The length of the pulses is a measure for the beam current
- **Zero-flux scheme:** compensate for the beam current and measure the magnitude of the **compensation current**



Performance

- **Achievable performance Fast Beam Current Transformers (FBCT):**
 - Absolute accuracy: 1%
 - Reproducibility / relative precision: 0.1%
 - Dynamic range: 10^3 (10^4)
- **Performance LHC DC Beam Current Transformers (DCCT):**
 - Absolute accuracy: 0.2%
 - Noise floor: $2 \mu\text{A}$
 - Dynamic range: 10^6 ($\mu\text{A} - 1\text{A}$)

Transverse Profile

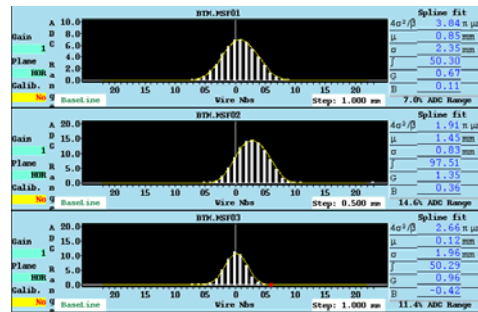
Overview - Beam Profile measurement

- Methods which intercept the beam with matter:
 - Secondary emission (SEM) grids
 - Screens
 - Wire scanners
 - more or less perturbing to the beam
 - Energies/intensity threshold for safe operation
 - Material damage (e.g. wire sublimation, breakage)
 - Radiation to other machine components (e.g. quenching of superconducting magnets)
- (Quasi) Non-Invasive Methods:
 - Synchrotron light monitors
 - Rest Gas Ionisation monitors
 - Luminescence monitors
 - Laser wire scanner
 - Electron beam scanner
 - Gas curtains, gas pencil beams
 - Beam Gas Vertex Detector

SEM grids and wire scanners:
Used as reference measurement for the other methods

Secondary Emission (SEM) Grids

- When the beam passes through a wire, secondary electrons are emitted, proportional to beam intensity
- The current flowing back onto the wires is measured using one amplifier/ADC chain for each wire
- Clearing field removes liberated electrons
- Problem: thermal emission
- Very high sensitivity, semi-transparent
- Good absolute measurement
- Spatial resolution limited by wire spacing to $\approx 0.25\text{mm}$
- Dynamic range: $\approx 10^6$



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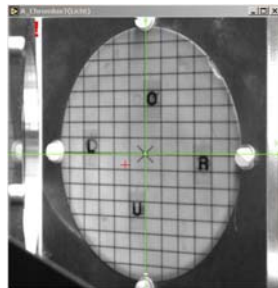
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Scintillation Screens

- Typically for setting-up with low intensities, thick screens (mm) \rightarrow emittance blow-up
- Workshop in 2011 at GSI to look at resolution possible with various screen materials: <http://www-bd.gsi.de/ssabd/home.htm>
- Sensitivities of different materials vary by orders of magnitudes



Abbreviation	Material	Activator	max. emission	decay time
Quartz	SiO ₂	none	470 nm	< 10 ns
	CsI	Tl	550 nm	1 μs
Chromolux	Al ₂ O ₃	Cr	700 nm	100 ms
YAG	Y ₃ Al ₅ O ₁₂	Ce	550 nm	0.2 μs
	Li glass	Ce	400 nm	0.1 μs
P11	ZnS	Ag	450 nm	3 ms
P43	Gd ₂ O ₂ S	Tb	545 nm	1 ms
P46	Y ₃ Al ₅ O ₁₂	Ce	530 nm	0.3 μs
P47	Y ₂ Si ₅ O ₅	Ce&Tb	400 nm	100 ns

Approximate values for inorganic scintillators

P. Forck, JUAS

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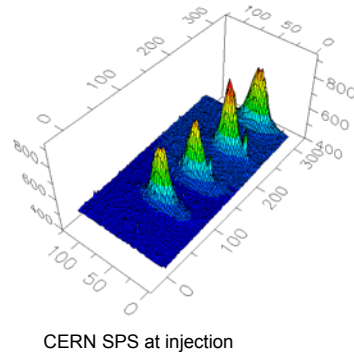
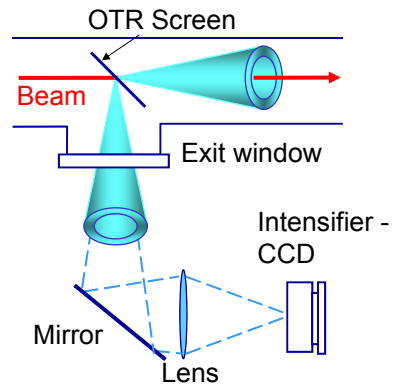
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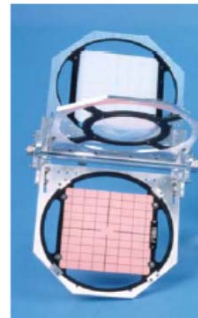
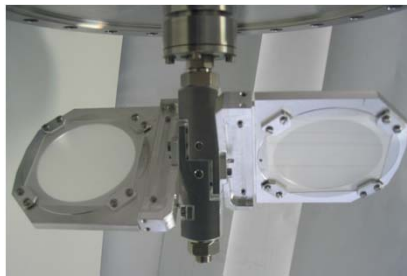
Optical Transition Radiation (OTR) Screens

- Radiation emitted when a charged particle beam goes through the interface of two media with different dielectric constants
- Surface phenomenon allows the use of very thin screens ($\geq 0.25 \mu\text{m}$)
- Much less intercepting, but requires higher intensity



Beam Profile Monitoring Using Screens

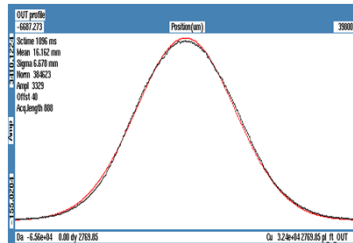
- Combine several screens in one housing e.g.
 - Al_2O_3 scintillation screen for setting-up with low intensity
 - Thin ($\approx 10\mu\text{m}$) Ti OTR screen for high intensity measurements
 - Carbon OTR screen for very high intensity operation



- Cameras:
 - CCD cameras are radiation sensitive
 - Analogue VIDICON camera can be used with high radiation

Wire Scanners

- A thin wire (down to 10 μm) is moved across the beam
 - Has to move fast to avoid excessive heating of the wire
 - Rotational scanner up to 10 m/s with special pneumatic mechanism (linear scanners slower)
- Detection
 - Secondary particle shower detected outside the vacuum chamber e.g. using a scintillator/photo-multiplier assembly
 - Secondary emission current detected as for SEM grids
- Correlating wire position with detected signal gives the beam profile
 - Wire vibrations limit position resolution
- Less invasive than screen or SEM grids



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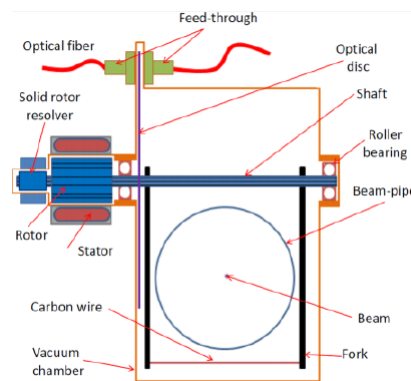
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New Wire Scanner being developed at CERN

- Design goals:
 - Spatial resolution of few μm (using high resolution angular position sensor)
 - Dynamic range: 10^4
 - Usage of sensor with large dynamic (diamond)
 - Automatic electronic switching of gain ranges
 - Minimize fork and wire deformations
 - Study of dynamic behavior of fork/wire system
 - Vibration mode optimized acceleration profile
- Current Wire Scanners at CERN:
 - Dynamic range 100; accuracy 5-10%; spatial resolution 50 μm (linear type) and 200 μm (rotational)



B. Dehning

Eva Barbara Holzer

CAS intr. Level course on Accelerator Physics

September, 2014

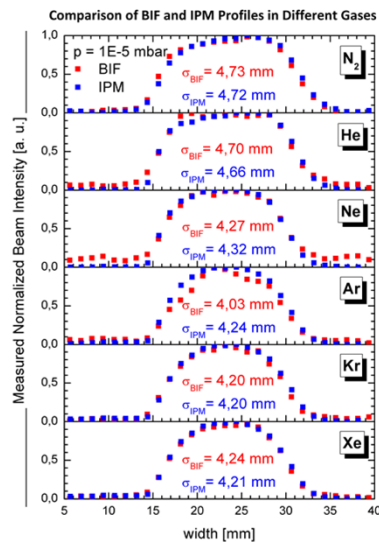
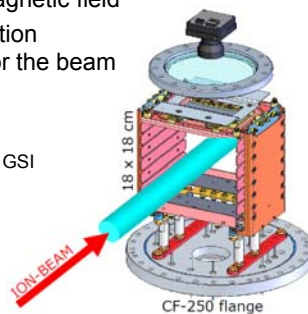
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(Quasi) Non-Invasive Beam Size Measurement

IPM (Ionization Profile Monitors)

- Residual Gas Ionisation
- dynamic range: up to 10^3
- ≈ 10 times more sensitive than Luminescence
- Image broadening due to space charge
- More complicated to build
 - High voltage
 - Guiding magnetic field
 - Compensation magnets for the beam

T. Giacomini et al., GSI



M.Schwickert, P.Forck, F.Becker, GSI

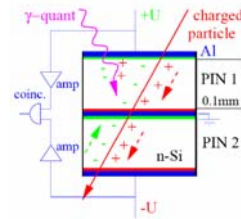
Beam Loss Measurement for Protection and Diagnostics

Common types of monitors

- Long ionisation chamber (charge detection)
 - Up to several km of gas filled hollow coaxial cables
 - Longitudinal position information by arrival time measurement
 - e.g. SLAC – 8m position resolution (30ns) over 3.5km cable length
 - Dynamic range of up to 10^4
- Cherenkov fibres
 - Time resolution 1 ns
 - Minimal space requirement
 - Insensitive to gamma background, E and B fields
 - Radiation hard (depending on type)
 - Combination fiber / readout can adapt to a wide dose range
 - Dynamic range 10^4 seems feasible

Common types of monitors cont'd

- **Short ionisation chamber** (charge detection)
 - Typically gas filled with many metallic electrodes and kV bias
 - Speed limited by ion collection time - tens of microseconds
 - Dynamic range of up to 10^8
- **PIN photodiode** (count detection)
 - Detect charged particle
 - Insensitive to photons from synchrotron radiation due to coincidence counting in two back-to-back mounted PIN diodes (K. Wittenburg, DESY)
 - Count rate proportional to beam loss
 - Speed limited by integration time
 - Dynamic range of up to 10^9
- **Scintillators plus photo-multipliers**
- etc. ...



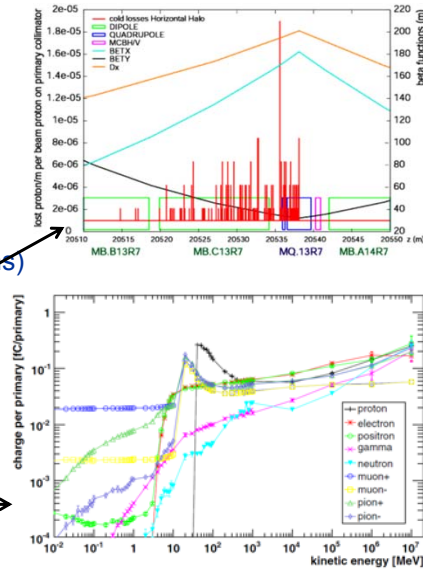
LHC BLM System

- **Main purpose: prevent damage and quench**
- 3600 Ionization chambers
- Beam abort thresholds:
 - 12 integration intervals: $40\mu\text{s}$ to 84s (32 energy levels)
 - 1.5 Million threshold values
- Each monitor aborts beam
 - One of 12 integration intervals over threshold
 - Internal test failed
- **Requirements and Challenges**
 - High Dependability (Reliability, Availability, Safety)
 - Threshold precision (factor 2)
 - Reaction time 1-2 turns ($100 - 200 \mu\text{s}$)
 - Dynamic range: 10^8 (at $40\mu\text{s}$ 10^5 achieved – 10^6 planned)
 - Radiation hard: currently at CERN development of kGy radiation hard readout to avoid noise from long cables



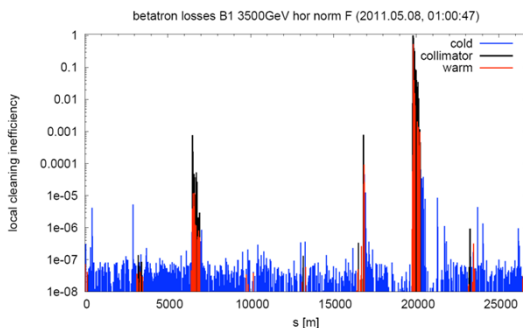
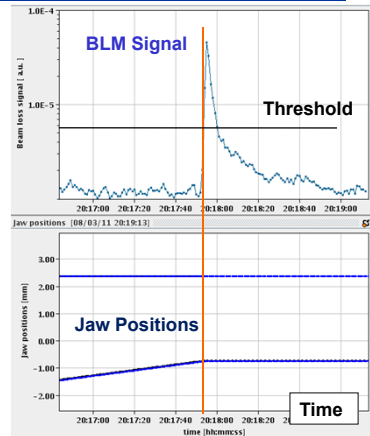
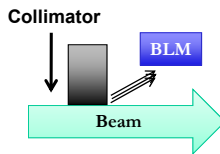
Threshold Determination

- Typically: Applied threshold set to 30% of the magnet quench level
- Relate the BLM signal to the:
 - Number of locally lost beam particles
 - Deposited energy in the machine
 - Quench and damage levels
- Extensive simulations during system design (and experimental verifications)
 - Proton loss locations (MAD-X, SIXTRACK)
 - Hadronic showers through magnets (GEANT, FLUKA)
 - Magnet quench levels as function of p energy and loss duration (SPQR)
 - Chamber response to the mixed radiation field (GEANT, FLUKA, GARFIELD)
- experimental verifications and beam tests in the LHC



Set-up and validation of collimation performance

- Find the beam center with each collimator jaw by stepping the jaw towards the beam and observing the BLM signal



'loss map': losses along the ring normalized to the losses at the primary collimator: performance verification

Diamond Detectors

- Fast and sensitive
- Small and radiation hard
- Used in LHC to distinguish bunch by bunch losses
- Dynamic range of monitor: 10^9
- Temporal resolution: few ns
- Investigations now ongoing to see if they can work in cryogenic conditions

