

Initial state of the HIC: thermalization and isotropization

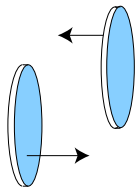
Yan Zhu



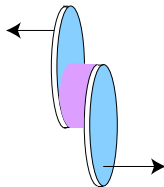
Particle Physics Day, 30 Oct. 2015, Helsinki

Where are we at?

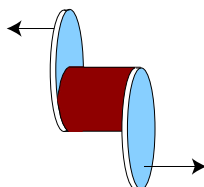
Lorentz contracted nuclei



Pre-thermal plasma



Locally thermalised plasma



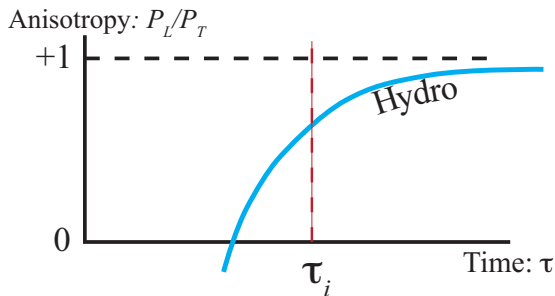
- Soft physics of HIC described by relativistic hydrodynamics

$$\partial_\mu T^{\mu\nu} = 0$$

- Gradient expansion around local thermal equilibrium

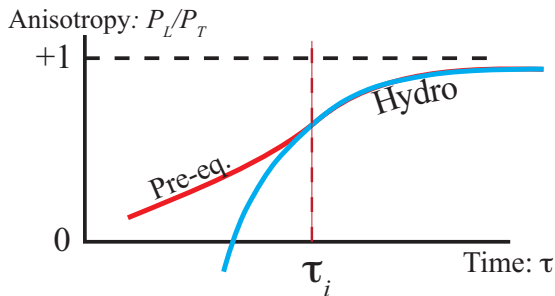
$$T^{\mu\nu} = T_{\text{Thermal equilibrium}}^{\mu\nu} - \eta(\epsilon)\sigma^{\mu\nu} - \zeta(\epsilon)\{g^{\mu\nu} + u^\mu u^\nu\}(\nabla \cdot u) + \dots$$

Where are we at?



- Strong anisotropy $P_L/P_T \ll 1$, sign of large corrections
- At early times *pre-equilibrium* evolution
- Hydro simulations start at *intialization time* τ_i

Where are we at?

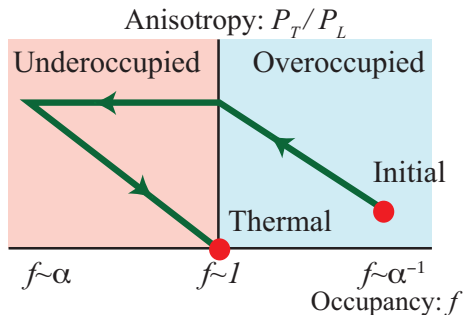


- If prethermal evolution converges smoothly to hydro, independence of unphysical τ_i
- Explicit example: Strong coupling $\mathcal{N} = 4$ SYM

Chesler, Yaffe PRL 106 (2011) 021601; van der Schee et al. PRL 111 (2013) 22, 222302, arXiv:1507.08195

This has proven to be challenging in QCD, even at weak coupling

Bottom-up thermalization at weak coupling



- Color Glass Condensate: Initial condition overoccupied

McLerran, Venugopalan PRD49 (1994) 2233-2241 , PRD49 (1994) 3352-3355 ; Gelis et. al Int.J.Mod.Phys. E16 (2007) 2595-2637 , Ann.Rev.Nucl.Part.Sci. 60 (2010) 463-489

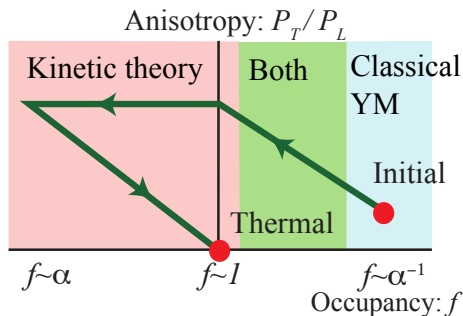
$$f(Q_s) \sim 1/\alpha_s, \quad Q_s \sim 2\text{GeV}$$

- Expansion makes system underoccupied before thermalizing

Baier et al Phys.Lett. B502 (2001) 51-58; Kurkela, Moore JHEP 1111 (2011) 120

$$f(Q_s) \ll 1$$

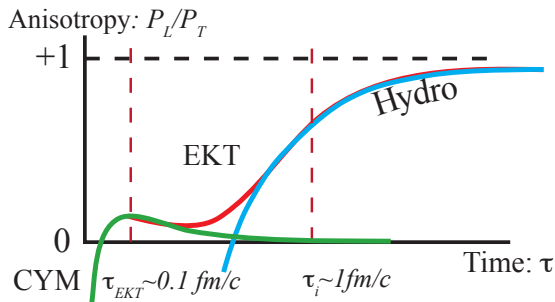
Bottom-up thermalization at weak coupling



- Degrees of freedom:
 - $f \gg 1$: Classical Yang-Mills theory (CYM)
 - $f \ll 1/\alpha_s$: (Semi-)classical particles, Eff. Kinetic Theory (EKT)
- Transmutation of fields to particles: Field-particle duality
Son, Mueller PLB582 (2004) 279-287; Jeon PRC72 (2005) 014907; Mathieu et al EPJ. C74 (2014) 2873 ; Kurkela, Moore PRD89 (2014) 7, 074036

$$1 \ll f \ll 1/\alpha_s$$

Strategy at weak coupling



Strategy: Switch from CYM to EKT at τ_{EKT} ,

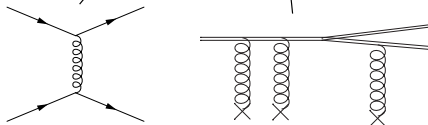
$$1 \ll f \ll 1/\alpha_s$$

From EKT to hydro at τ_i ,

$$P_L/P_T \sim 1$$

Effective kinetic theory of Arnold, Moore, Yaffe

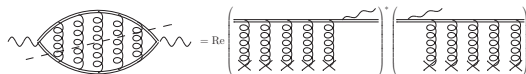
JHEP 0301 (2003) 030

$$\frac{df}{dt} = -C_{2 \leftrightarrow 2}[f] - C_{1 \leftrightarrow 2}[f]$$


The diagram on the left represents the $C_{2 \leftrightarrow 2}$ process, showing two incoming particles and two outgoing particles connected by a vertical gluon line. The diagram on the right represents the $C_{1 \leftrightarrow 2}$ process, showing a single incoming particle and two outgoing particles connected by two vertical gluon lines, with a horizontal line representing a soft gluon exchange.

- Soft and collinear divergences lead to nontrivial matrix elements

soft: screening, Hard-loop; collinear: LPM, ladder resum



The diagram shows a resummed matrix element for the $C_{2 \leftrightarrow 2}$ process. It consists of a central loop diagram with two incoming and two outgoing particles, connected by a vertical gluon line. This is followed by an equals sign and the real part of a complex expression. The expression is the real part of the product of two terms in large parentheses. The first term is a ladder diagram with two incoming and two outgoing particles, connected by a vertical gluon line, and a horizontal line representing a soft gluon exchange. The second term is a similar ladder diagram with two incoming and two outgoing particles, connected by a vertical gluon line, and a horizontal line representing a soft gluon exchange.

- No free parameters; LO accurate in the $\alpha_s \rightarrow 0$, $\alpha_s f \rightarrow 0$ limit.
- Used for LO transport coefficients in QCD, jet energy loss

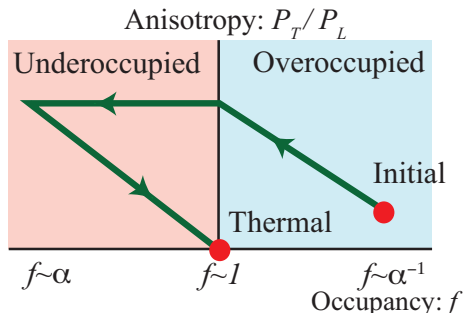
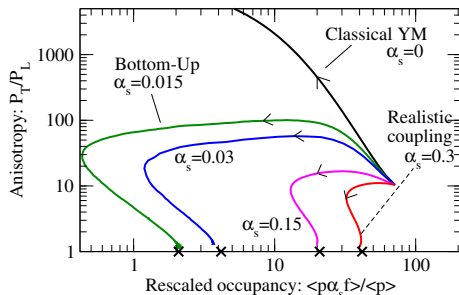
Arnold et al. JHEP 0305 (2003) 051; Moore, York PRD79 (2009) 054011; Ghiglieri, Teaney 1502.03730; Kurkela, Wiedemann PLB740 (2015) 172-178; Iancu, Wu 1506.07871

- Now also available in NLO $\mathcal{O}(\sqrt{\alpha_s})$

Ghiglieri 1509.07773

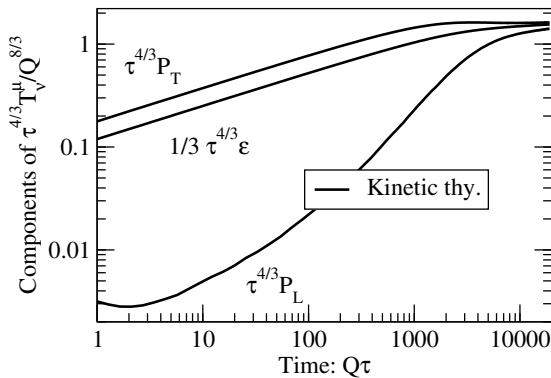
Route to equilibrium in EKT

Kurkela, YZ, PRL in press



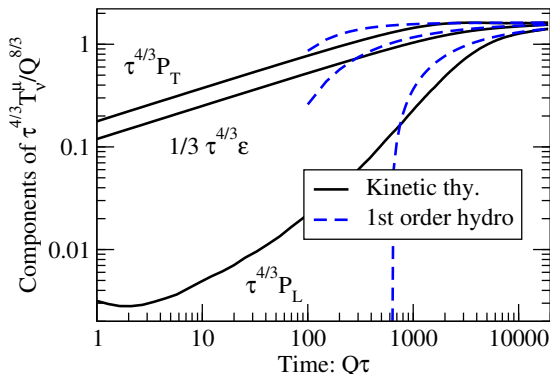
- Initial condition ($f \sim 1/\alpha_s$) from classical field theory calculation
Lappi PLB703 (2011) 325-330
- In the classical limit ($\alpha_s \rightarrow 0, \alpha_s f$ fixed), no thermalization
- At small values of couplings, clear Bottom-Up behaviour
- Features become less defined as α_s grows

$$\alpha_s = 0.03$$



- Kinetic theory converges to hydro smoothly and automatically

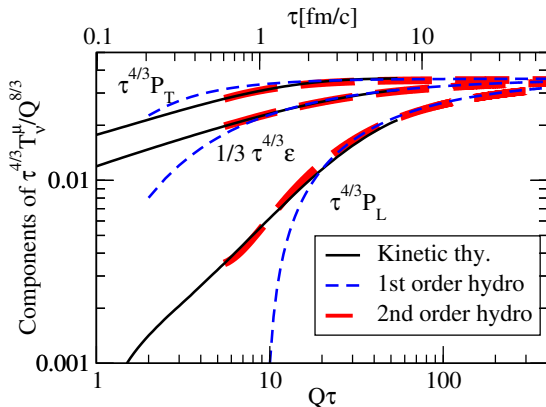
$$\alpha_s = 0.03$$



- Kinetic theory converges to hydro smoothly and automatically
- Hydro prediction fixed by perturbative η/s

Arnold et al. JHEP 0305 (2003) 051

$$\alpha_s = 0.3$$



- For realistic couplings, hydrodynamics reached around $\lesssim 1\text{fm}/c$.
- Hydro seems to give a good description even when $P_L/P_T \sim 1/5$

Where are we going?

- Combination of classical Yang-Mills simulations and effective kinetic theory allows to follow the time evolution from highly occupied initial condition to thermal equilibrium.
- Weak coupling thermalization extrapolated to realistic couplings shows agreement with hydro around

$$\tau_1 \sim 1\text{fm}/c$$

- Unified description of soft and hard physics: hydro, jets, etc.