

OPEN ISSUES AND TOOLS FOR SINGLE TOP

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CONTENTS

- ✿ Some publicity about *MadGraph5_aMC@NLO* to be officially released soon (next week?)...
- ✿ When to use NLO+PS?
- ✿ Four and five-flavour schemes
- ✿ Non-resonant contributions to single top
- ✿ $Wt+ttbar$ combined
- ✿ $W+b$ background

MADGRAPH5_AMC@NLO EXAMPLE

	LO	NLO
Start the python interface	<code>./bin/mg5</code>	<code>./bin/mg5</code>
Generate a process	<code>generate p p > t t~ w+</code>	<code>generate p p > t t~ w+ [QCD]</code>
write the process to disk	<code>output MY_TTW_PROC</code>	<code>output MY_TTW_NLOPROC</code>
start the event generation	<code>launch</code> unweighted events	<code>launch</code> unweighted events up to a sign
Accuracy	LO + PS	NLO + PS

WHICH PROCESSES CAN I GENERATE AT NLO?

- ✱ Any SM process you are interested in...
 - ✱ The only thing to keep in mind is that NLO processes are a lot more complicated than the corresponding LO process (so, do not try to do $t\bar{t} + 3\text{jets}$ @ NLO...)

WHICH PROCESSES CAN I GENERATE AT NLO?

- ✱ $p p > t t^{\sim}$ [QCD]
- ✱ $p p > t t^{\sim} j$ [QCD]
- ✱ $p p > t t^{\sim} h$ [QCD]
- ✱ $p p > t t^{\sim} z$ [QCD]
- ✱ $p p > t t^{\sim} t t^{\sim}$ [QCD]
- ✱ $p p > t t^{\sim} w^+ a$ [QCD]
- ✱ $p p > t t^{\sim} b b^{\sim}$ [QCD]

✱ ...

- ✱ $p p > t b^{\sim}$ [QCD]
- ✱ $p p > t j \$\$ w^+$ [QCD]
- ✱ $p p > t b^{\sim} j \$\$ w^+$ [QCD]
- ✱ $p p > t b^{\sim} h \$\$ w^+$ [QCD]
- ✱ $p p > t j z \$\$ w^+$ [QCD]
- ✱ $p p > t j a \$\$ w^+$ [QCD]
- ✱ ...

- ✱ All possible with aMC@NLO on a multi-core desktop machine (or a small computer cluster)

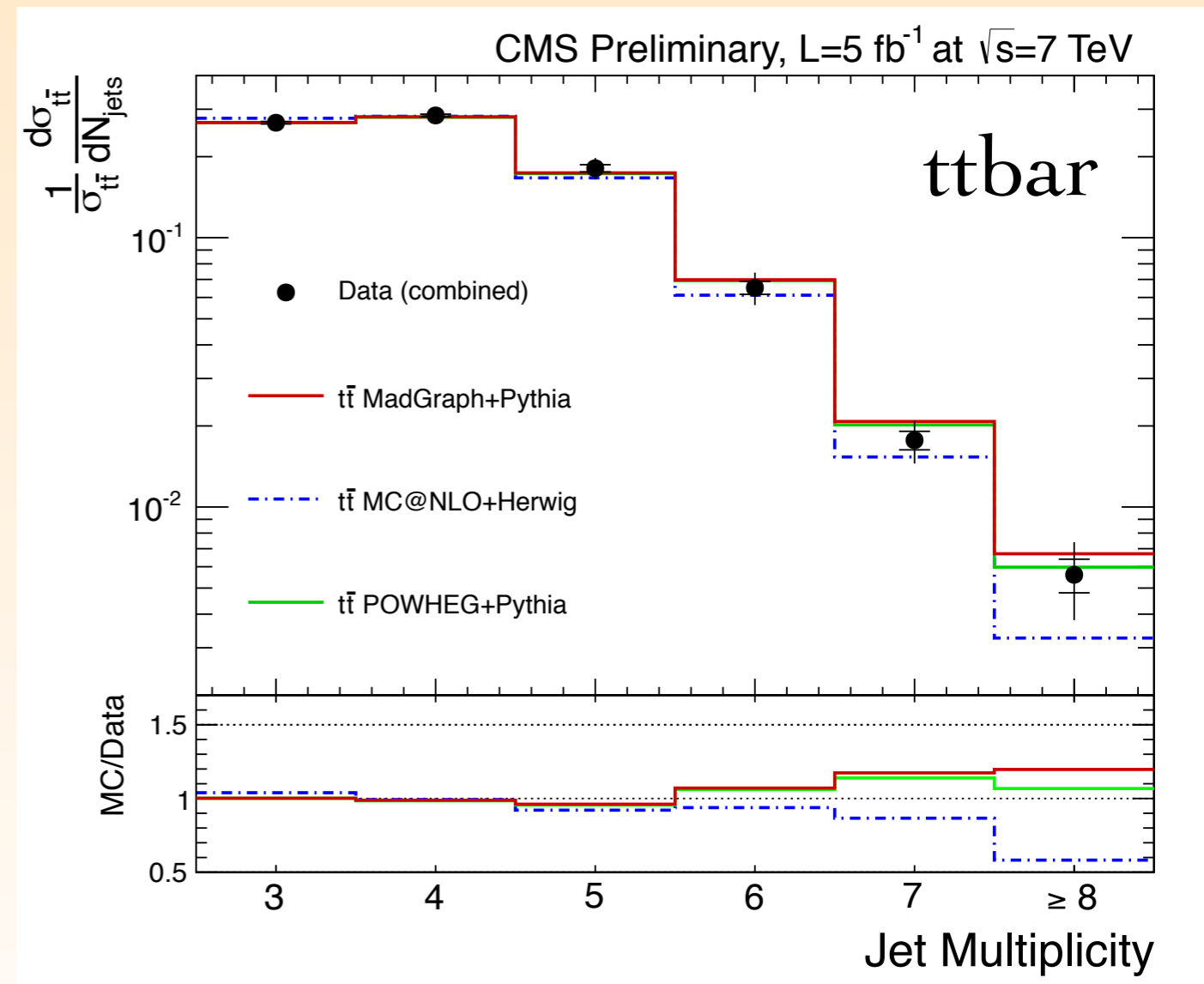
“ $\$\$ w^+$ ” means to forbid W^+ -bosons to appear in the s-channel; in general not gauge invariant, so use with extreme care

FURTHER SUBTLETY

- ✿ There is another class of processes for which QCD NLO corrections cannot simply be included
- ✿ In general, QCD NLO corrections to an EW process are of the same order as EW NLO corrections to a QCD process
 - ✿ In general, one cannot disentangle them and both need to be included for consistency
 - ✿ But NLO EW corrections are usually not available; at least not in the codes that aim for full automation of NLO+PS (aMC@NLO, Sherpa, ..)
 - ✿ implementation in aMC@NLO has started
- ✿ Possibly a problem for single top production because it is an EW process, but there is no equivalent QCD process so it's okay in most cases (we'll see an example later where some tricks were needed)

WHEN TO USE NLO+PS

- ✿ It makes only sense to NLO corrections if the observable is sensitive to those corrections.
- ✿ Otherwise it is just a waste of time and one might as well use stand-alone Pythia or Herwig

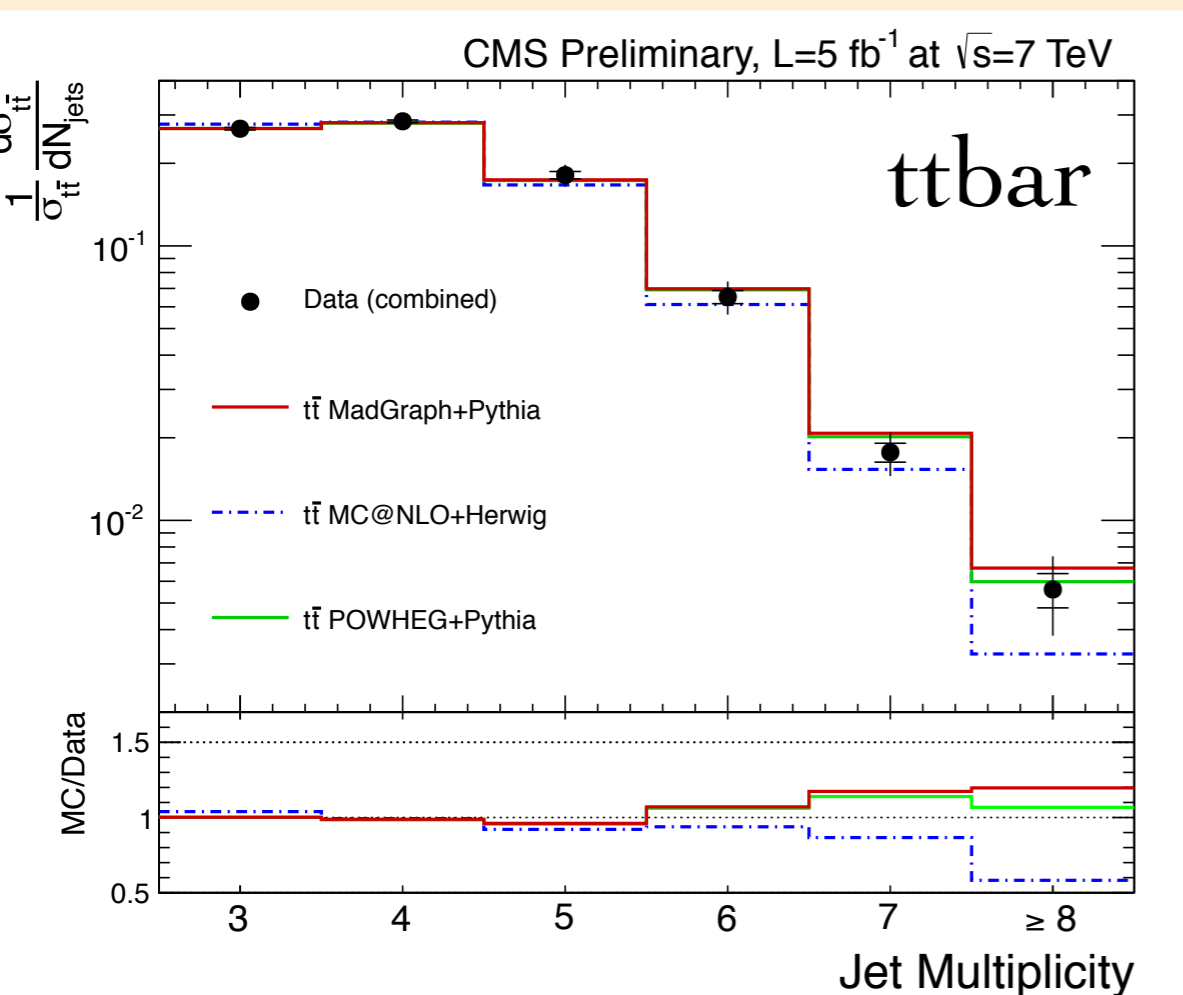


UNCERTAINTY ESTIMATES

- ✱ For observables that **have NLO precision**, the theory/generator uncertainties can be estimated by
 - ✱ Independent renormalisation and factorisation scale variations
 - ✱ PDF error sets (preferably following the PDF4LHC agreement)
 - ✱ Matching an NLO computation to at least 2 different parton showers
 - ✱ These PDF and scale variations can be obtained via reweighting in aMC@NLO and POWHEG, not yet possible in Sherpa.
- ✱ For observables that **do not have NLO precision**, further uncertainties are coming from the shower starting scale (“Power” or “Wimpy” shower). Currently these cannot be approximated with the (a)MC@NLO program, but not really relevant because why use an NLO+PS computation for these observables in the first place?
They can be estimated more correctly in the NLO Sherpa program.

NLO CORRECTIONS TO HIGHER MULTIPLICITIES

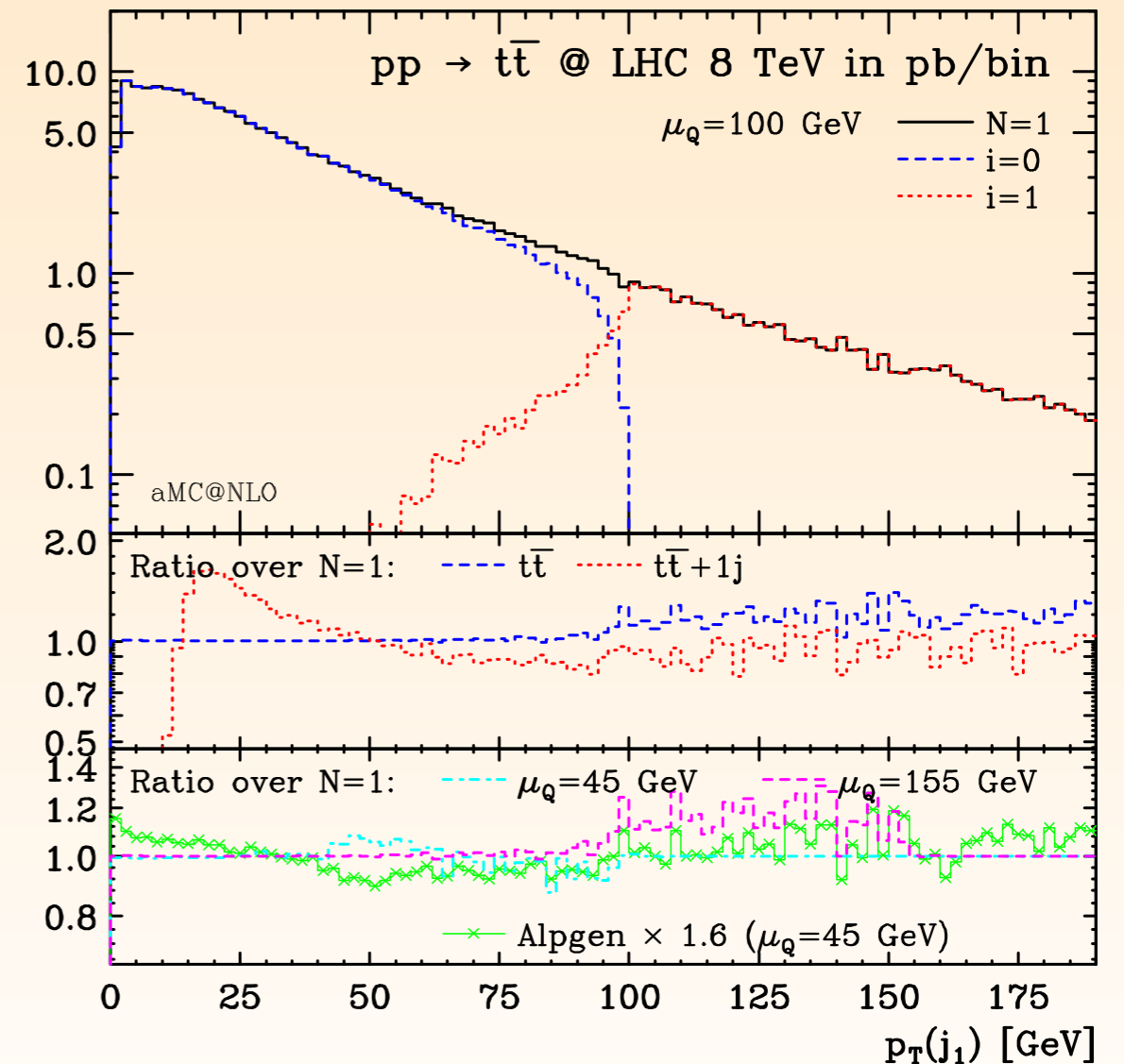
- For this example/observable there is a straight-forward way of also getting (some of) the higher multiplicity bins at NLO accuracy:



- Compute NLO corrections to $t\bar{t} + \text{jet}$ or $t\bar{t} + 2\text{jets}$ and match that calculation to the shower.
- This requires a generation cut, otherwise the process diverges, and it should be checked that the final results are independent from the generation cut
- This will also work for single top+jets or any other processes

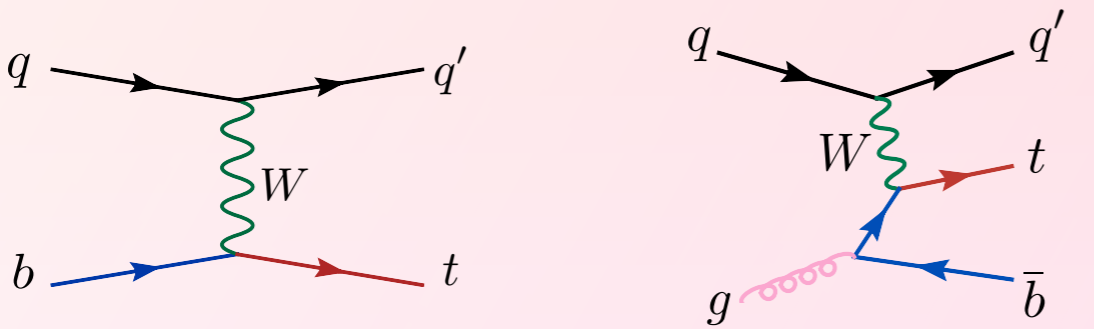
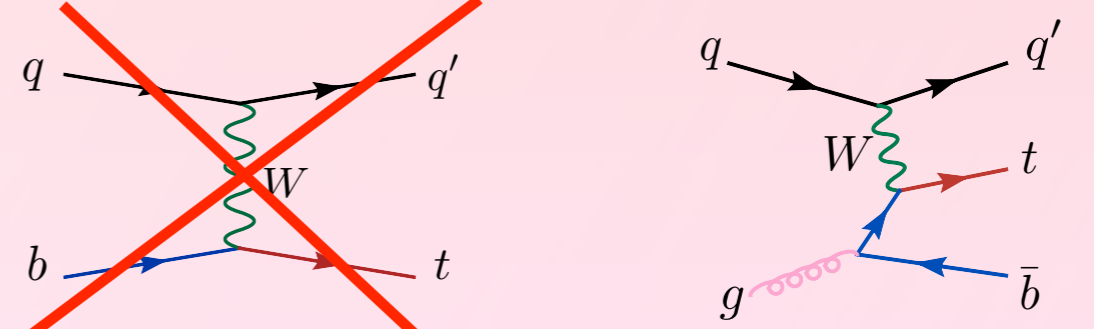
EXTENDING NLO ACCURACY OVER MORE THAN ONE MULTIPLICITY

- Using a special procedure ('**FxFx**' or '**MEPS@NLO**' or '**MINLO**' or '**UNLOPS**' or '**some other acronym**') NLO accuracy can be extended over more than one multiplicity.
- This only works for processes that do not have matrix element jets at the lowest multiplicity Born matrix elements (ttbar is okay, single top isn't)



4- & 5-FLAVOR SCHEMES

2 ways of making predictions

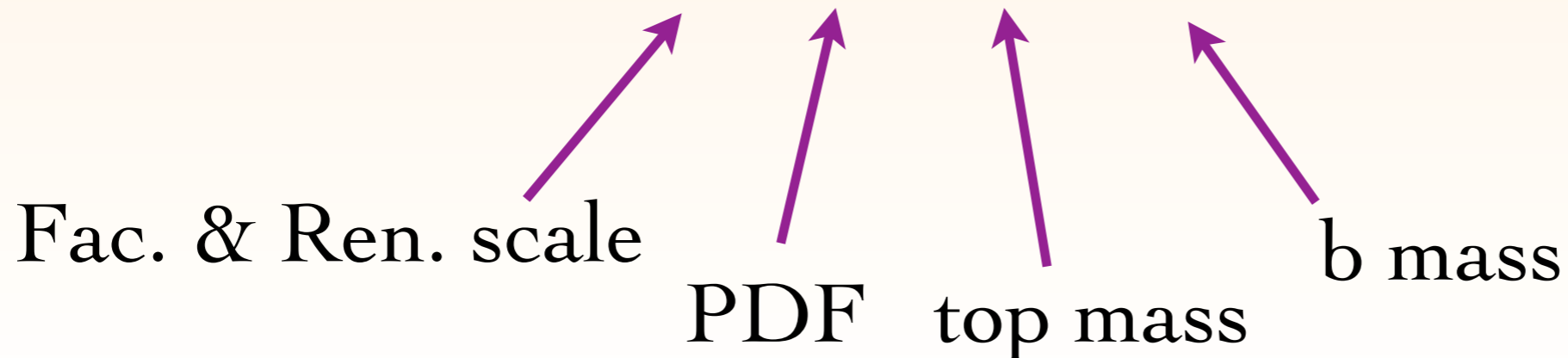
5 flavour scheme	4 flavour scheme
massless b	massive b
PDF includes initial state b quarks	No b quarks in PDF
Log[m_b/μ_F] resummed in PDF	Finite terms correctly included
Simpler calculation	More involved prediction
 <p>leading order (contribution to) NLO</p>	 <p>Does not exist (part of) leading order</p>

Descriptions are equivalent when including all orders in perturbation theory

UNCERTAINTIES

- ✱ Estimate of the theory uncertainty:
 - ✱ independent variation of renormalization and factorization scales by a factor 2
 - ✱ 44 eigenvector CTEQ6.6 PDF's
 - ✱ Top mass: 172 ± 1.7 GeV (sorry, it's a bit of an old table!)
 - ✱ Bottom mass: 4.5 ± 0.2 GeV
- ✱ Uncertainties from scales slightly larger in 4F (as expected)
- ✱ Other sources are comparable in size
- ✱ Don't forget the uncertainty from the bottom quark mass uncertainty!

$\sigma_{t\text{-ch}}^{\text{NLO}}(t + \bar{t})$	$2 \rightarrow 2$ (pb)					$2 \rightarrow 3$ (pb)				
Tevatron Run II	1.96	+0.05 -0.01	+0.20 -0.16	+0.06 -0.06	+0.05 -0.05	1.87	+0.16 -0.21	+0.18 -0.15	+0.06 -0.06	+0.04 -0.04
LHC (7 TeV)	62.6	+1.1 -0.5	+1.4 -1.6	+1.1 -1.1	+1.1 -1.1	59.4	+2.1 -3.4	+1.4 -1.4	+1.0 -1.0	+1.3 -1.2
LHC (14 TeV)	244	+5 -4	+5 -6	+3 -3	+4 -4	234	+7 -9	+5 -5	+3 -3	+4 -4

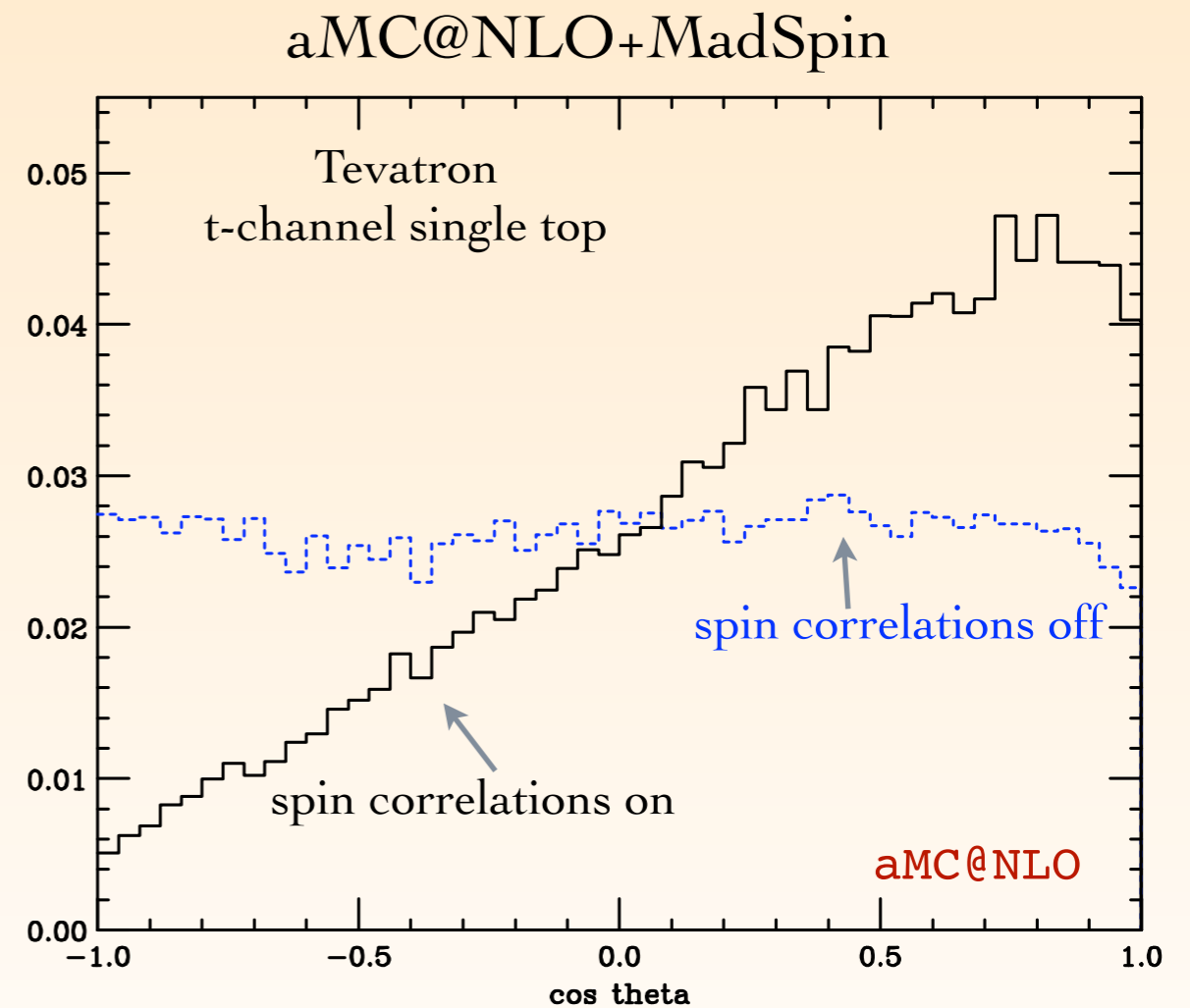


T-CHANNEL SINGLE TOP: “BEST PRESCRIPTION”

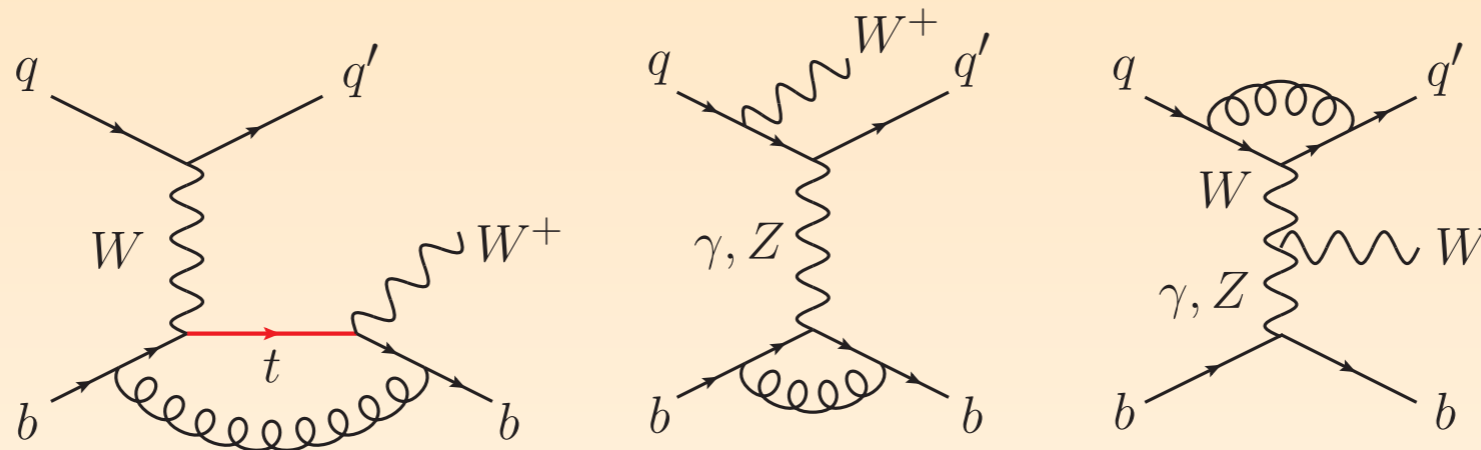
- ✱ Use the NLO+PS 4-flavour scheme predictions for the kinematics, but normalised to the NLO 5-flavour scheme with an overall factor
 - ✱ All kinematics of the top, the light jet and the **“spectator b-quark”** are NLO correct because of the 4 flavour scheme
 - ✱ Overall normalisation is improved because log's are resummed in the PDF in the 5 flavour scheme

SPIN CORRELATIONS

- ✿ In aMC@NLO, spin correlations can be included using MadSpin
 - ✿ Uses only tree-level information
 - ✿ Includes some off-shell effects
- ✿ Also POWHEG events can be decayed with MadSpin this way (when not directly available in POWHEG, i.e. single top 4 flavour), although some MadGraph-style information should be added to the header of the event file to tell MadSpin the process definition and all that.



OFF-SHELL SINGLE TOP



- ✿ To include the full NLO spin correlations, as well as all the all off-shell effects, the complete set of diagrams are needed
- ✿ Not really possible for single top in the 4 flavour scheme, because of the EW nature of the single top process, QCD corrections mix with the EW corrections to a corresponding QCD process

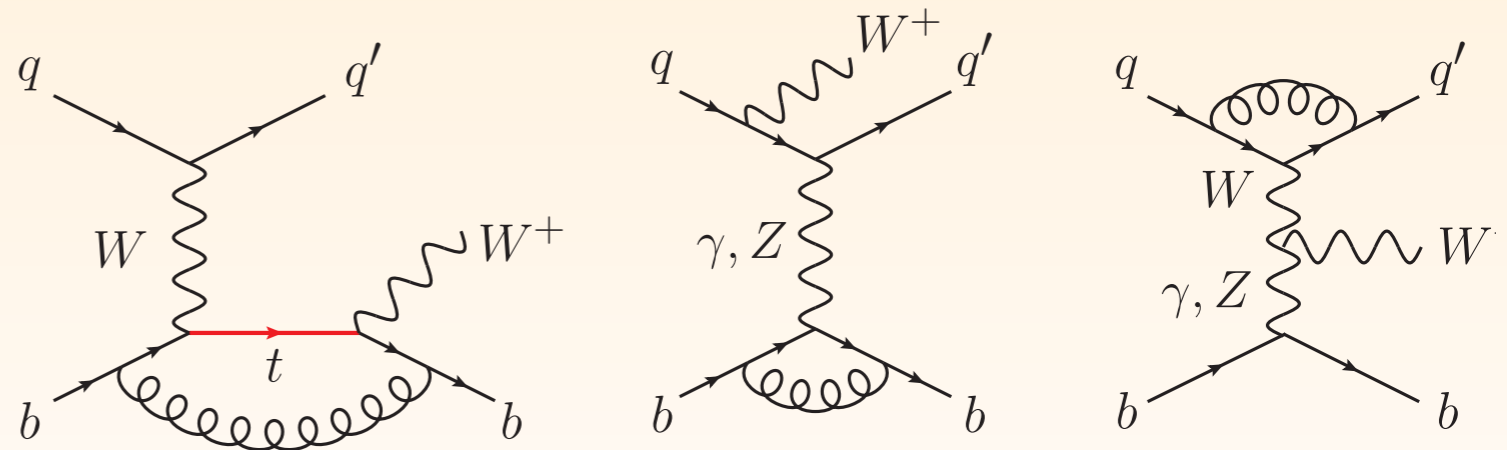
PROCESS DEFINITION

Papanastasiou, RF, Frixione, Hirschi, Maltoni arXiv:1305.7088

- ✱ Particular care needed in process definition: $pp \rightarrow W^+ J_b J_{\text{light}} + X$
- ✱ Single top production is EW process. In general QCD corrections to an EW process cannot be disentangled from EW corrections to a QCD process
 - ✱ Need to use **diagonal CKM matrix** (at least for the 3rd generation) to prevent interference from QCD Born with EW Born. Only possible for 5F scheme calculations, not possible in the 4F scheme.

- ✱ Setting $V_{tb}=1$ also allows for **separation of t and s-channel contributions** (which results in easier comparison with literature)

- ✱ With decays, the process is not finite: **need a cut on the b-jet**

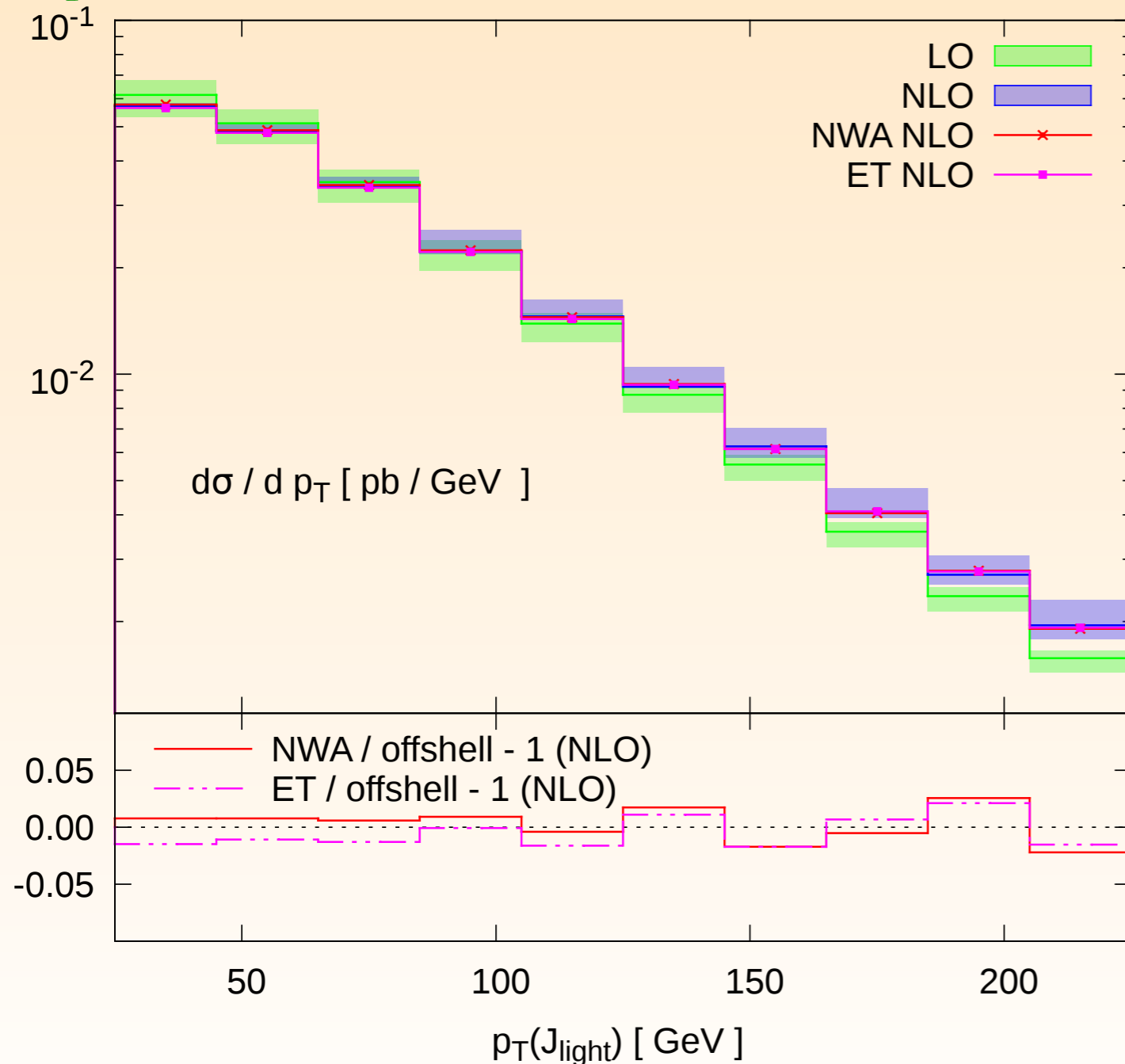


- ✱ s-channel could be included in our approach but requires that J_b does not include two b quarks

- ✱ With $V_{tb}=1$, the NLO corrections to the t-channel process are finite and well defined, provided that J_b has a non-zero transverse momentum

COMPARISON BETWEEN APPROXIMATIONS

Papanastasiou, RF, Frixione, Hirschi, Maltoni arXiv:1305.7088

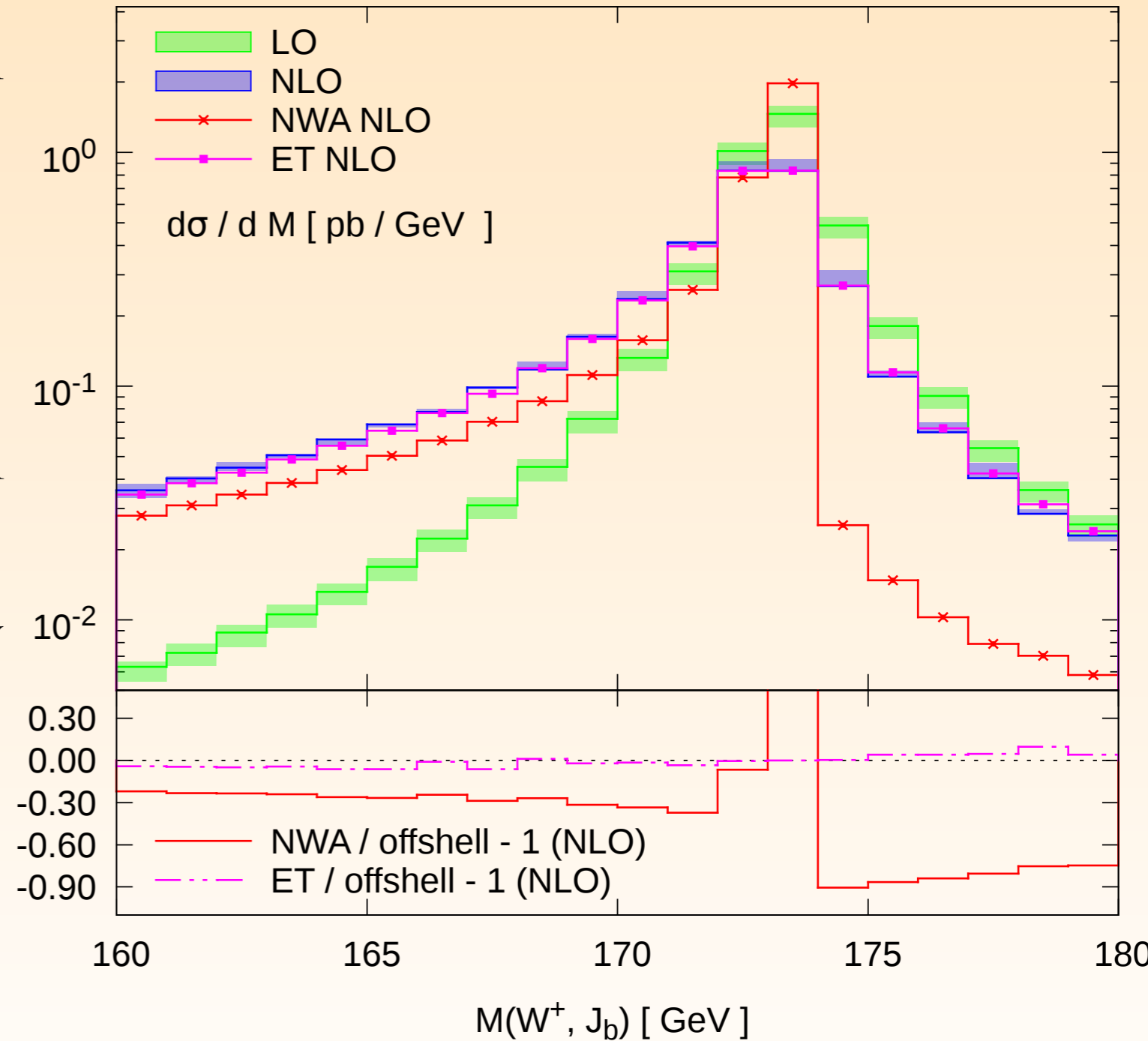


- ✱ Comparison of LO, NLO, NLO in the narrow width approximation (NWA) [MCFM] and the effective theory (ET) approach [Falgari, Mellor, Signer arXiv:1007.0893; + Giannuzzi arXiv:1102.5267] that includes the leading contributions beyond the NWA
- ✱ For observables that are integrated over the top resonance peak, differences between NWA, ET and full NLO are rather small:
 - ✱ No visible effects in shape of the transverse momentum of the light jet

“TOP INVARIANT MASS”: $M(W^+, J_B)$

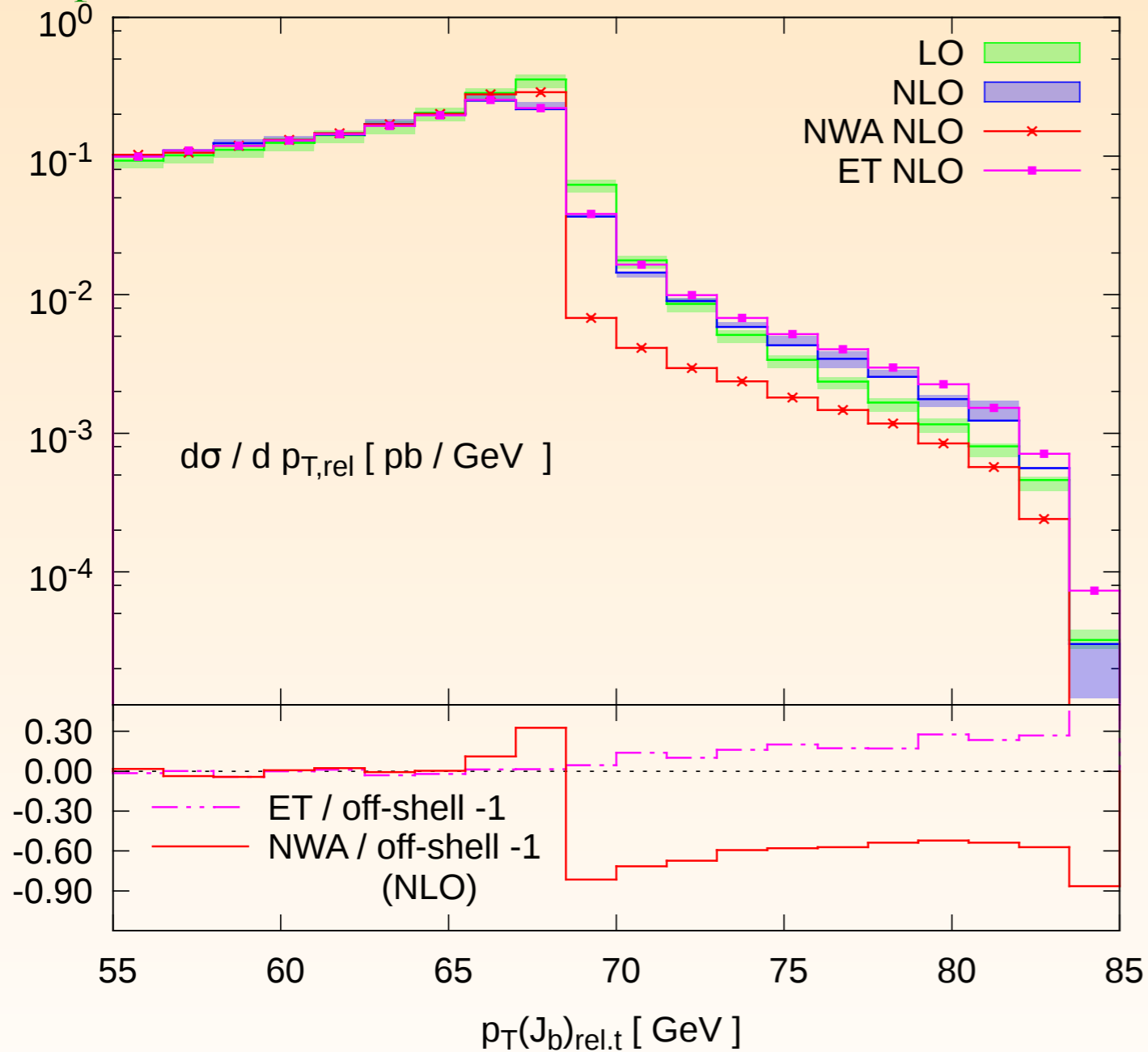
Papanastasiou, RF, Frixione, Hirschi, Maltoni arXiv:1305.7088

- ✱ Invariant mass of the W-boson and the b-jet
- ✱ LO prediction greatly undershoots the NLO results below the peak: no radiation from the b-jet
 - ✱ Including parton shower should get most of this right
- ✱ Above peak, NWA undershoots by a long way: no off-shell effects are included
- ✱ Effective theory approach results in an excellent description over the whole range plotted here



RELATIVE $p_T(J_B)$ TO THE (RECONSTRUCTED) TOP QUARK DIRECTION

Papanastasiou, RF, Frixione, Hirschi, Maltoni arXiv:1305.7088



- ✱ At LO with on-shell top quarks, this distribution has a kinematic cut-off at $p_T(J_b) = (m_t^2 - m_W^2) / (2m_t)$
- ✱ In the NWA this cut-off largely remains
- ✱ Again ET approach does a very good job, apart in the far tail. This is a sign that subleading Γ_t contributions become important

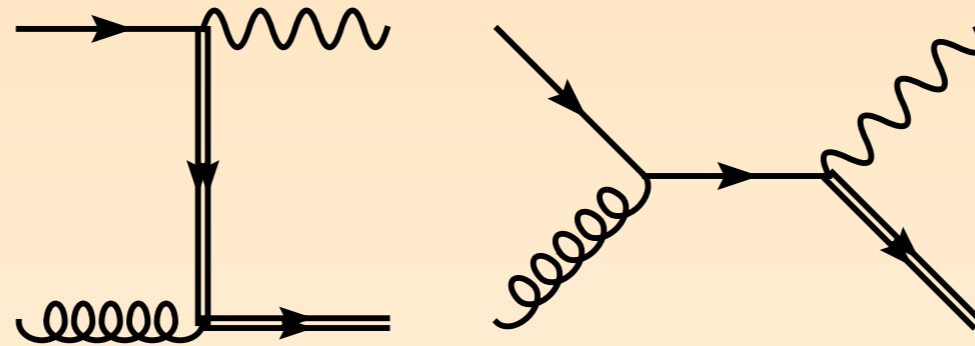
NARROW WIDTH APPROXIMATION

- ✱ Most observables are well-described by the NWA
- ✱ However, always keep this approximation in mind: care must be taken when observables as $M(W^+, J_b)$ or $p_T(J_b)_{rel,t}$ are used (e.g. in template fitting).
- ✱ Effective theory approach does an excellent job close to the resonance, but will ultimately fail as well
 - ✱ Excellent for top quarks as signals
 - ✱ When top quarks are background, and on-shell contributions are removed with cuts, relative enhancement of non-resonant contributions

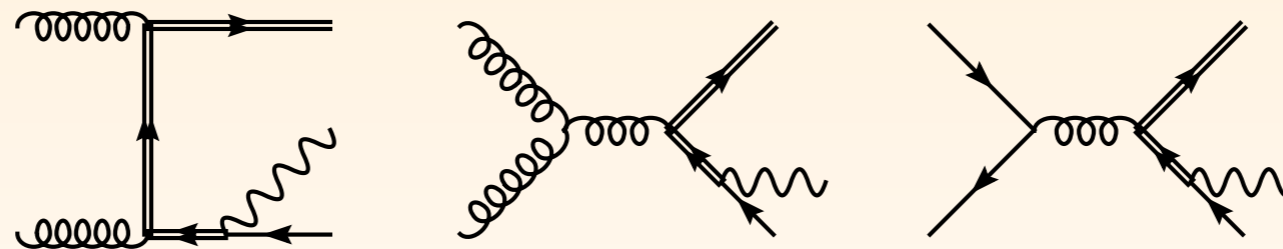
ADDING PARTON SHOWER

- ✱ For processes with intermediate top quark resonances it is not straight-forward to include parton shower effects
- ✱ Possible double counting from radiation from those top quarks and real-emission corrections. In the default MC@NLO scheme this double counting is not subtracted
- ✱ Possibly simple solution with excellent approximation is to forbid the parton shower to include radiation from the top quarks
- ✱ Under investigation...

WT-CHANNEL



- Biggest issues in single top modeling are coming from the W-boson associated single top channel
- At NLO, the real-emission diagrams have a contribution from top pair production



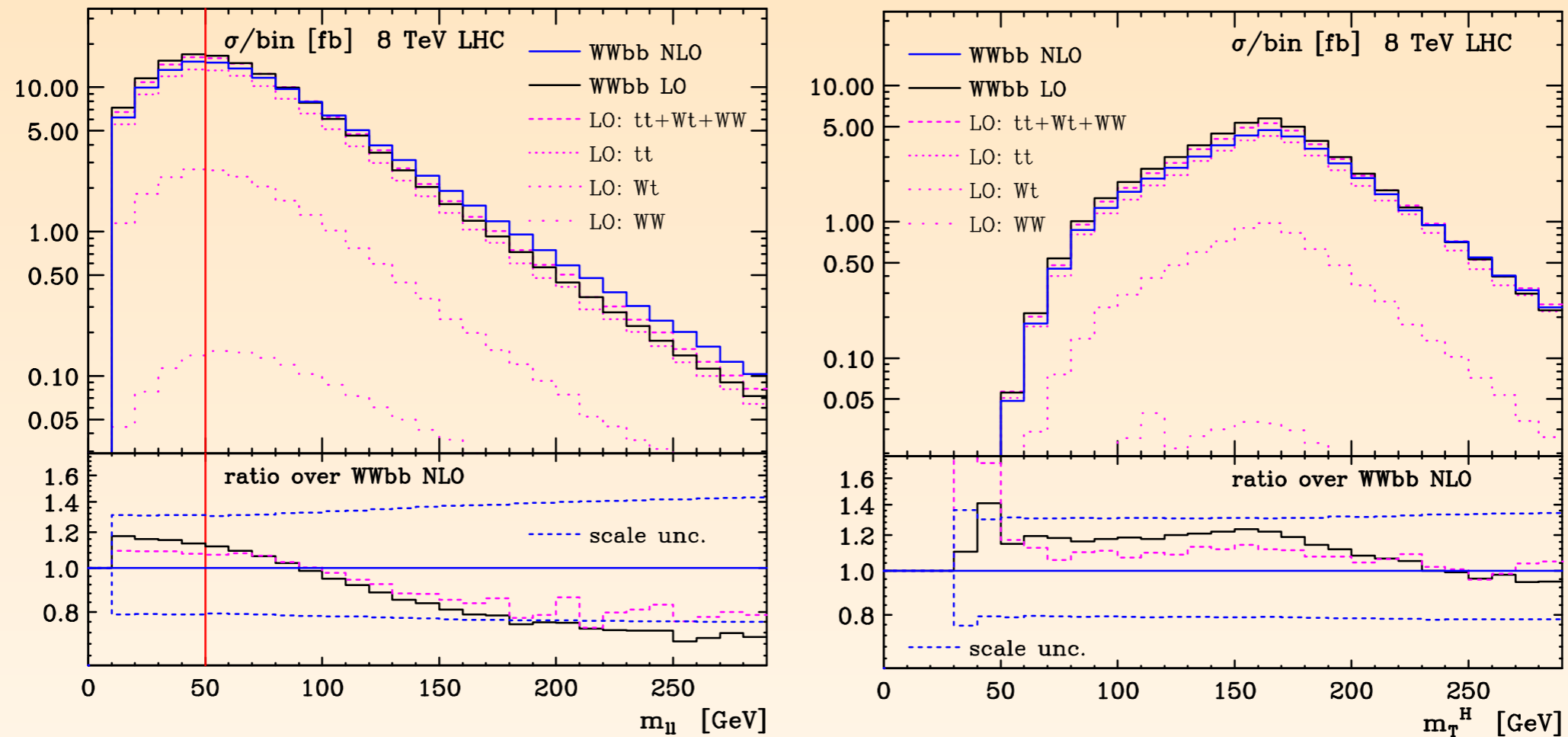
- “Perturbation theory breaks down”: the full NLO corrections to Wt production are much larger than the Born, because they receive a contribution from LO top pair production
- DR/DS scheme has been developed to remove/subtract them

DR/DS SCHEME

- ✱ Remove double resonant ($t\bar{t}$ -like) contributions at the level of amplitude (DR) or matrix elements (DS)
- ✱ Difference is the interference between Wt and $t\bar{t}$ production
- ✱ Both descriptions are formally not gauge invariant
- ✱ Wt with DR/DS available in MC@NLO and POWHEG
- ✱ Not available in aMC@NLO (not so easy to automate), i.e. not worth it for Wt production only...
 - ✱ Wt is the only example in the SM where this is needed, but many BSM processes have the same problems
 - ✱ When BSM can be done in an automated way, this is something that we'll start thinking about more carefully

WWBB AT NLO

RF arXiv:1311.4893

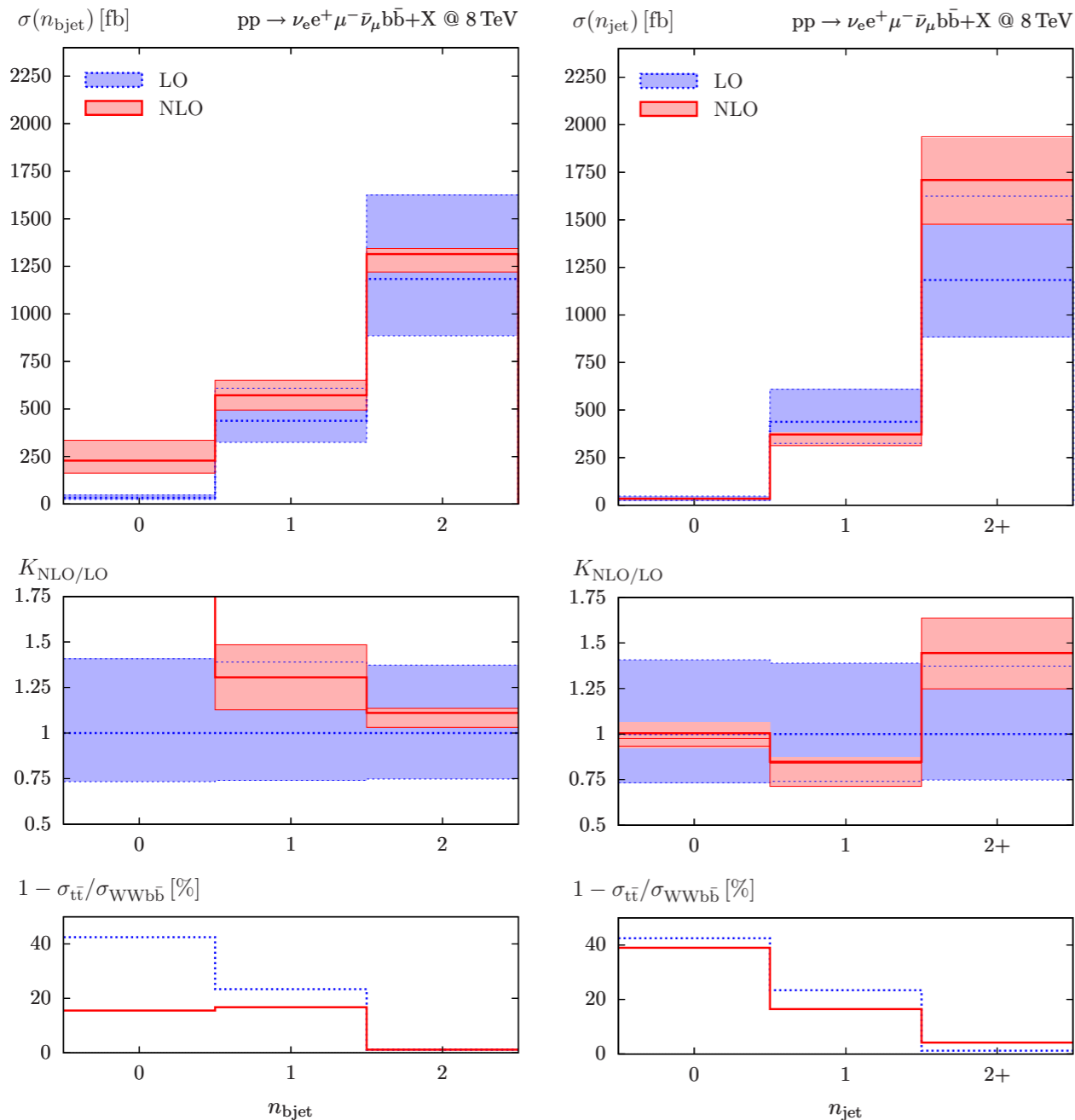


- ✱ This has recently been achieved and applied to the $H \rightarrow WW^*$ measurement channel in the one-jet bin, requiring a single jet
- ✱ Long list of cuts to suppress backgrounds
- ✱ For these observables, NLO scale dependence remains large

SIMILAR STUDY UNDERWAY BY THE ZURICH GROUP

Numerical Results for the 8TeV LHC Integrated cross sections at NLO QCD

Multiplicity of (b)jets with $p_T \geq 30$ GeV



First “off-shell” insights into **exclusive 0- and 1-(b)jet bins** of the combined process $t\bar{t} + tW$

(b)jet multiplicity

- large K factors in 0-bjet and 1-bjet bins due to light-jet emission,
- tW contribution to 0-bjet and 1-bjet bins reduced at NLO.

(generic) jet multiplicity

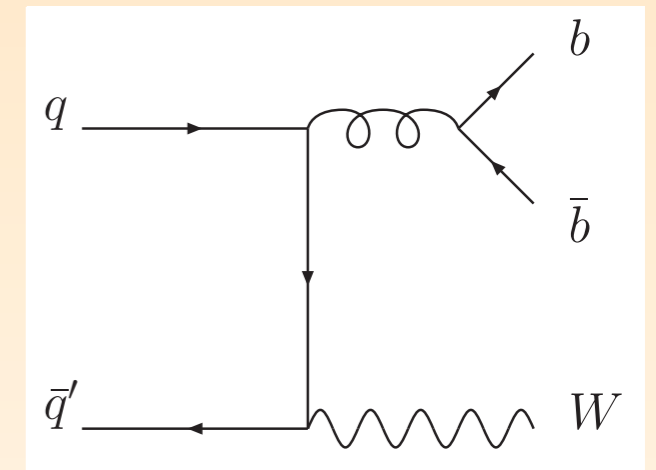
- reduced K factors in the lower jet-multiplicity bins,
- relative contribution from tW stays quite constant when going from LO to NLO.

W-BOSON + B JETS

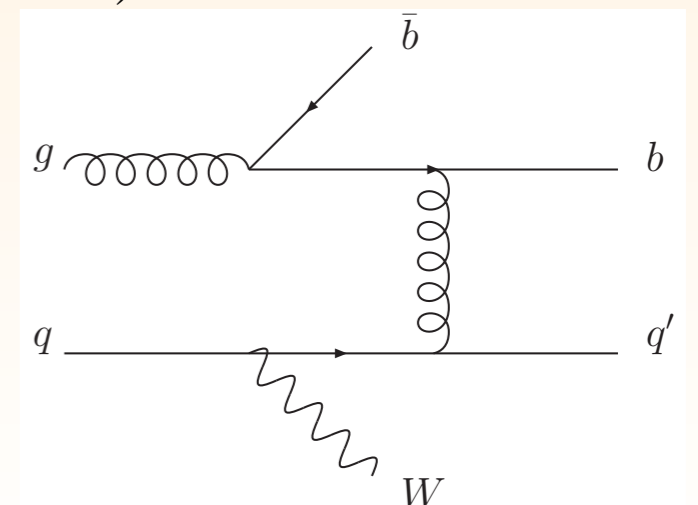
- ✱ For single top it is straight-forward to understand the two prescriptions
- ✱ It is more involved when trying to describe W-boson plus b-jets
- ✱ Each of the following need a separate description, e.g.
 - ✱ $W+1$ jets with 1 b tag (inclusive or exclusive)
 - ✱ $W+2$ jets with 1 b tag (inclusive or exclusive)
 - ✱ $W+2$ jets with 2 b tags (inclusive or exclusive)
 - ✱ $W+bb$ -jet (inclusive or exclusive)
[bb-jet is a jet containing two b-quarks]

4-FLAVOR SCHEME

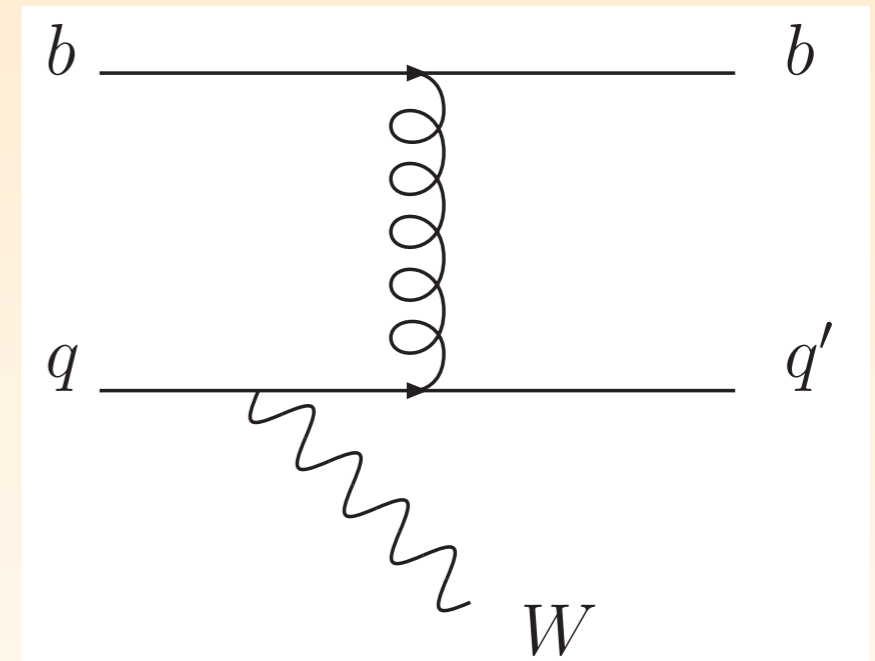
1. $W+1$ jets with 1 b tag (inclusive or exclusive)
2. $W+2$ jets with 1 b tag (inclusive or exclusive)
3. $W+2$ jets with 2 b tags (inclusive or exclusive)
4. $W+bb$ -jet (inclusive or exclusive)



- ✱ All of them are described by this process in the 4-flavor scheme
 - ✱ finite process (IR singularities regularized by the bottom mass)
 - ✱ known at NLO+PS (both in POWHEG and aMC@NLO)
- ✱ “ $W+2$ jets with 1 b tag (inclusive)” is also (better?) described by this diagram (and similar ones)
 - ✱ No careful study has been done to asses this. Easily possible with aMC@NLO

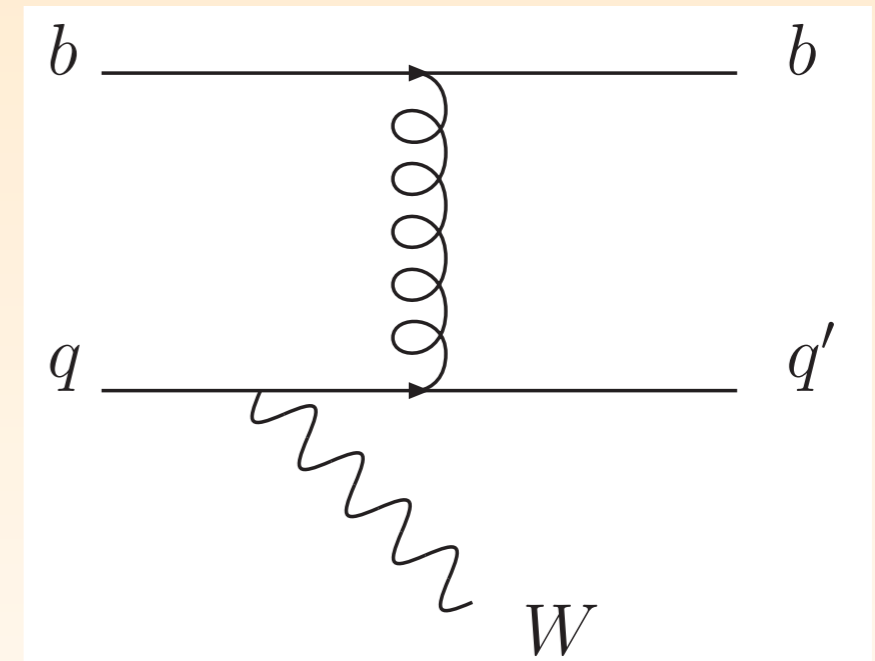


5-FLAVOR SCHEME



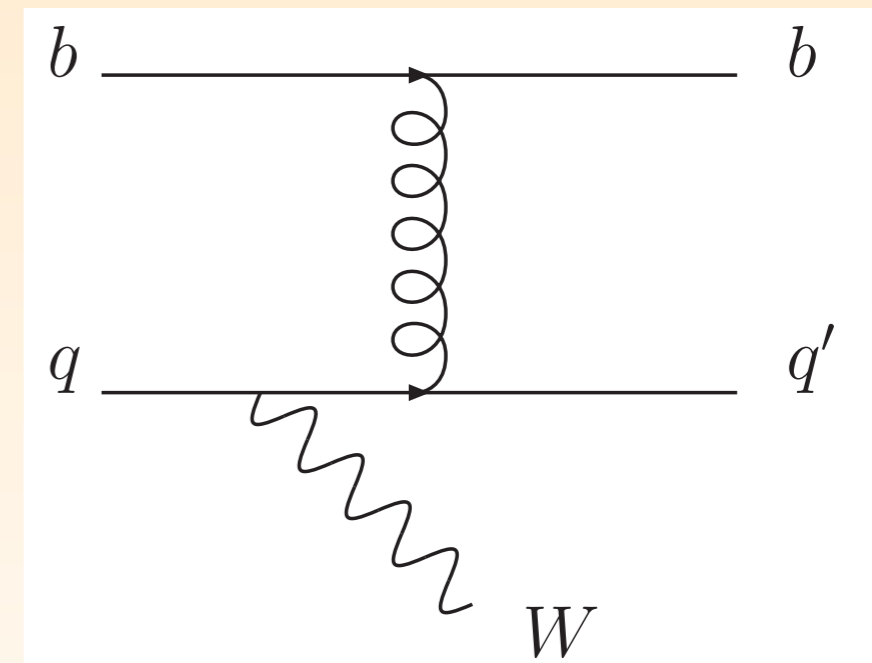
5-FLAVOR SCHEME

- ✱ When requiring 2 b tags, the same diagrams as in the 4-flavor scheme are appropriate here as well (in principle with a massless b -quark)



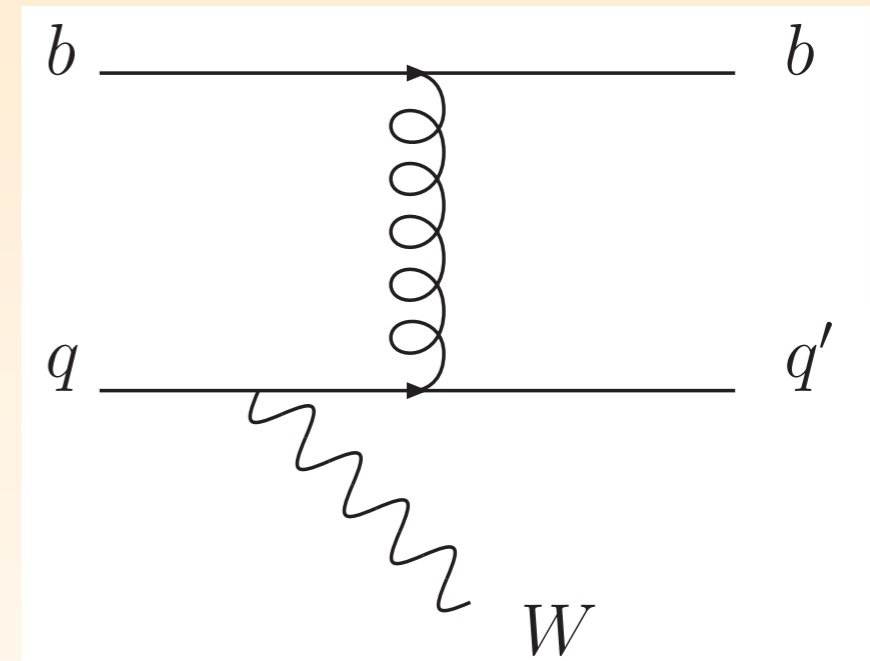
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- ✱ bb -jets can only be described by the parton shower



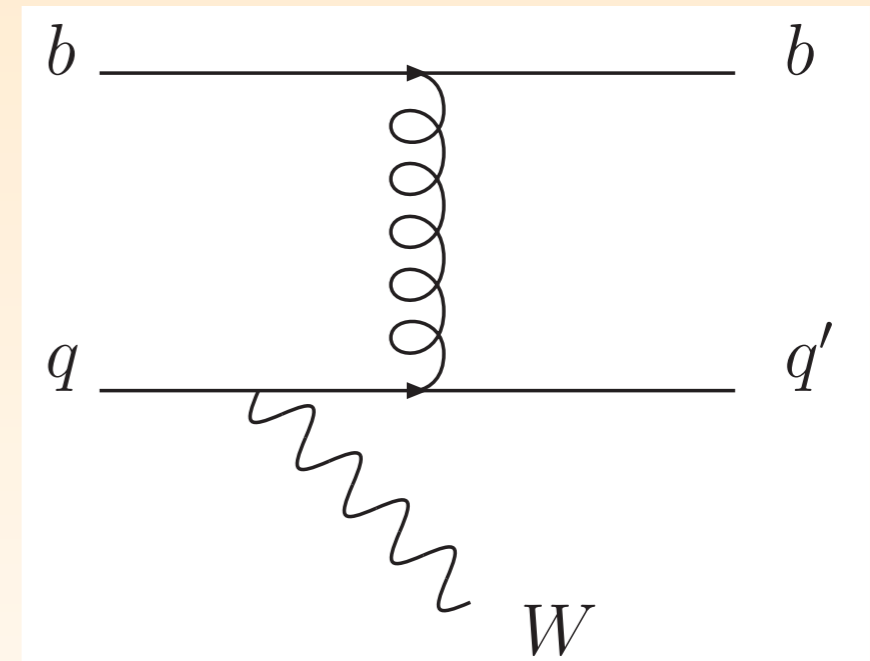
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- ✱ bb-jets can only be described by the parton shower
- ✱ When requiring only 1 b tag, there is a better description with initial state b-quarks
 - ✱ Smaller uncertainties compared to 4-flavor scheme for observables that are not sensitive to very soft/forward b quarks
 - ✱ Fixed order NLO study to combine the two approaches in one consistent description for $W+1,2$ jets with (at least) 1 b tag [[Caola et al. arXiv:1107.3714](#)], but non-trivial with parton shower



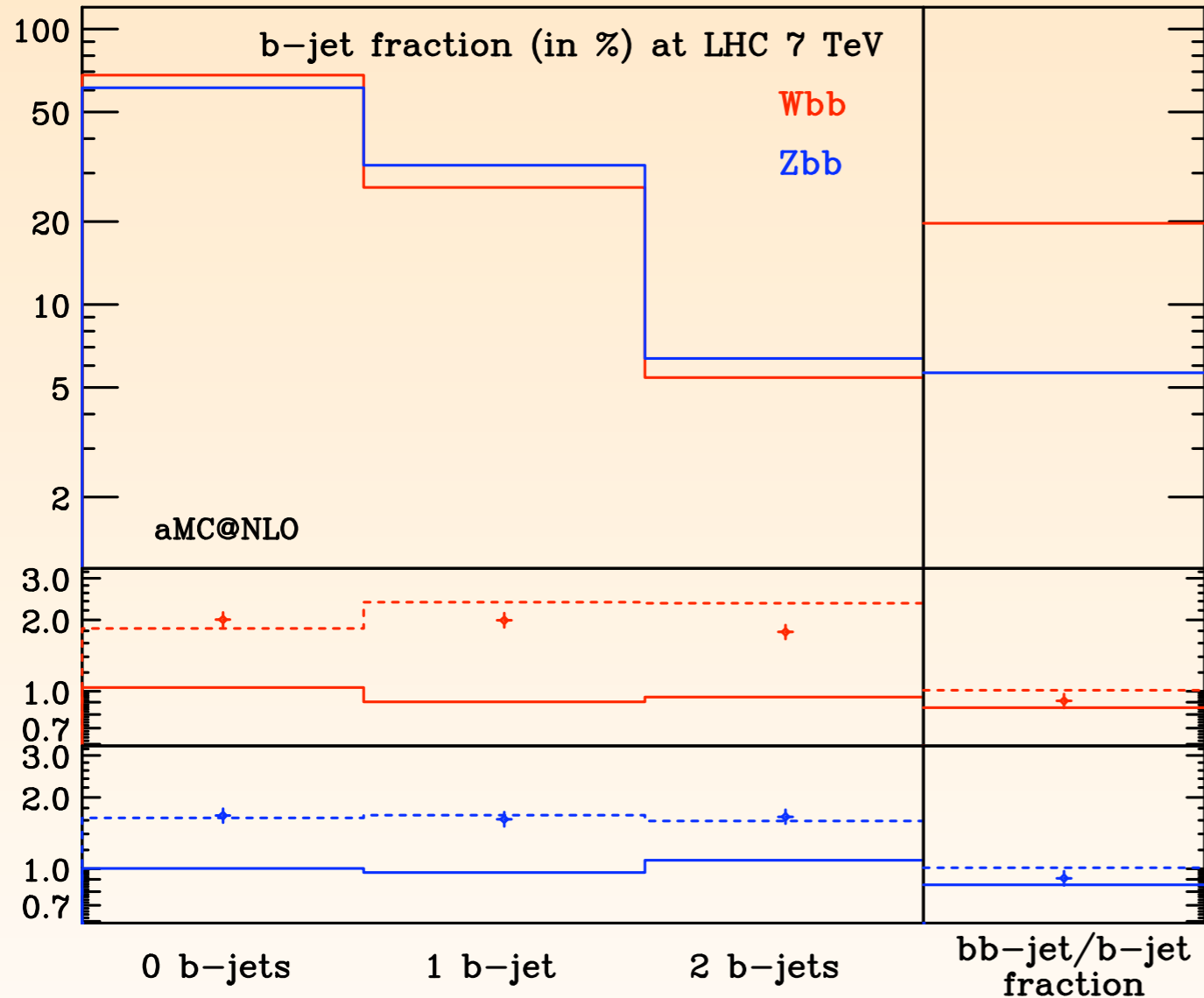
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 - ✱ 4-flavor scheme calculation is simpler in the sense that “one fits all”



PP \rightarrow WBB/ZBB

[RF, Frixione, Hirschi, Maltoni, Pittau & Torrielli, arXiv:1106.6019]



- ✱ Large NLO corrections
- ✱ Lower panels show the ratio of aMC@NLO with LO (crosses), NLO (solid) and LOwPS (dotted)
- ✱ NLO and aMC@NLO very similar and consistent
- ✱ Uncertainties can be estimated from (not shown):
 - ✱ Renormalisation & Factorisation scale dependence
 - ✱ Switching parton shower (Herwig vs. Pythia)

CONCLUSIONS

- ✱ Automation of NLO QCD corrections done
- ✱ But not always:
 - ✱ be careful when trying NLO QCD corrections to an EW process
- ✱ For single top, use NLO 4-flavour, but normalise to the total 5-flavour rate
- ✱ For $W+b$ background situation is not so obvious, because it depends much more on the precise observable (1 or 2 b-tags, inclusive/exclusive)
- ✱ Non-resonant contributions are being included at fixed order NLO (not yet with inclusion of the parton shower as well, but that should be soon)
- ✱ First results at NLO without separating Wt & $t\bar{t}$ have been presented. This amounts to computing the NLO corrections to the full $WWbb$