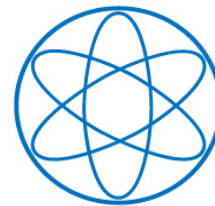


Dark Matter Hunt

Alejandro Ibarra

Technische Universität München



Cargese
July 2014

Direct detection

DM nucleus \rightarrow DM nucleus



Indirect detection

DM DM $\rightarrow \gamma X, e^+e^- \dots$ (annihilation)

DM $\rightarrow \gamma X, e^+X \dots$ (decay)

Collider searches

pp \rightarrow DM X

Direct detection

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Collider searches

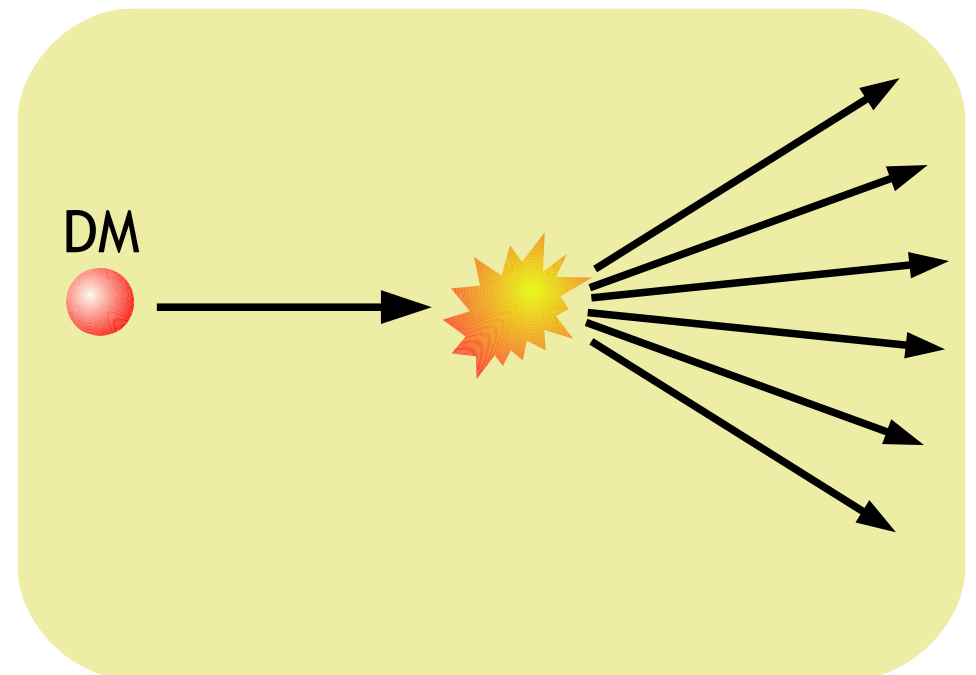
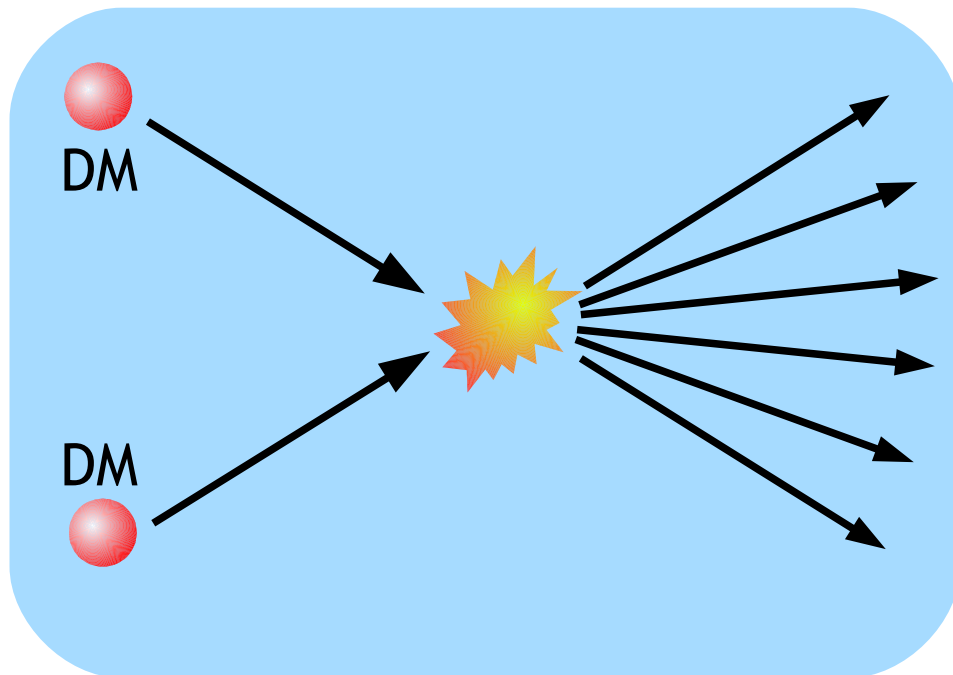
pp \rightarrow DM X

Indirect Dark Matter Searches

Indirect dark matter searches

General idea:

1) Dark matter particles annihilate or decay producing a flux of stable particles: photons, electrons, protons, positrons, antiprotons or (anti-)neutrinos.



Indirect dark matter searches

General idea:

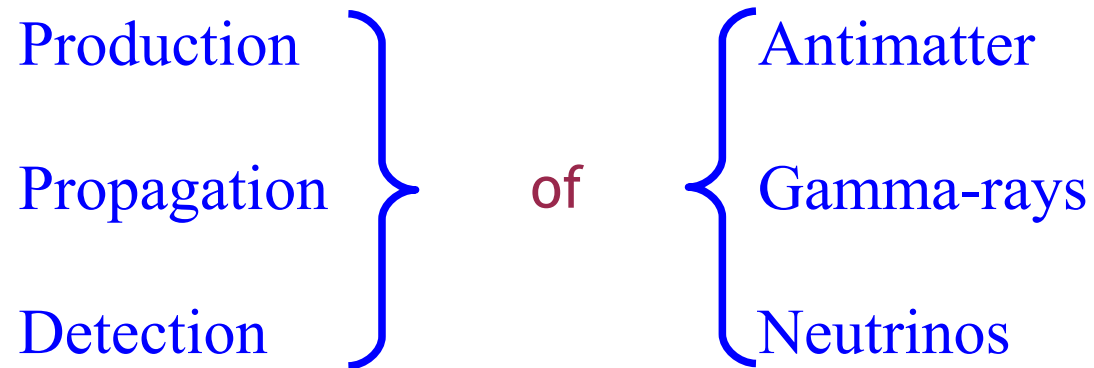
- 1) Dark matter particles annihilate or decay producing a flux of stable particles: photons, electrons, protons, positrons, antiprotons or (anti-)neutrinos.
- 2) These particles propagate through the galaxy and through the Solar System. Some of them will reach the Earth.

Indirect dark matter searches

General idea:

- 1) Dark matter particles annihilate or decay producing a flux of stable particles: photons, electrons, protons, positrons, antiprotons or (anti-)neutrinos.
- 2) These particles propagate through the galaxy and through the Solar System. Some of them will reach the Earth.
- 3) The products of the dark matter annihilations or decays are detected **together with other particles produced in astrophysical processes** (for example, cosmic ray collisions with nuclei in the interstellar medium). The existence of dark matter can then be inferred if there is a significant excess in the fluxes compared to the expected astrophysical backgrounds.

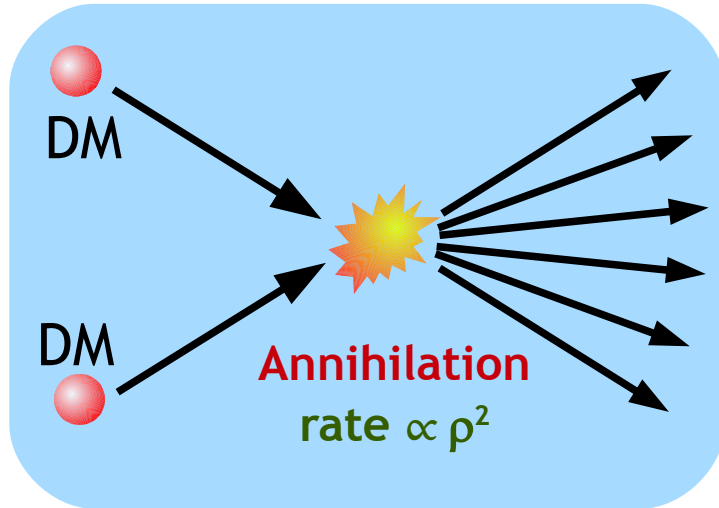
Indirect dark matter searches



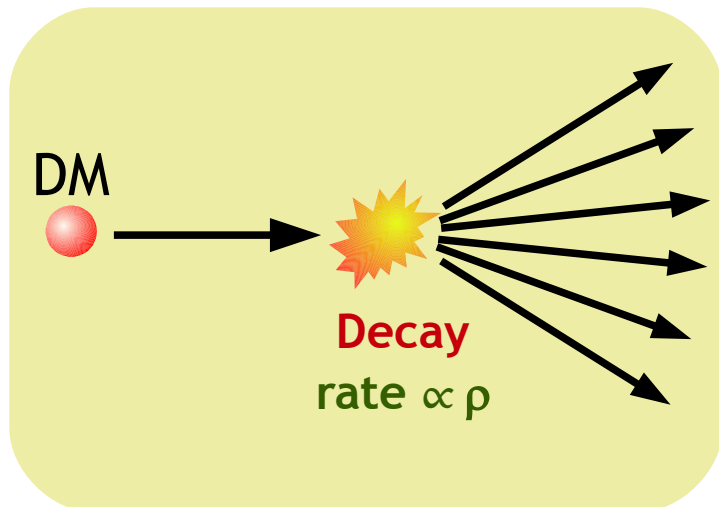
Antimatter

Production

The production is described by the **source function**: number of particles produced at a given position per unit volume, unit time and unit energy.

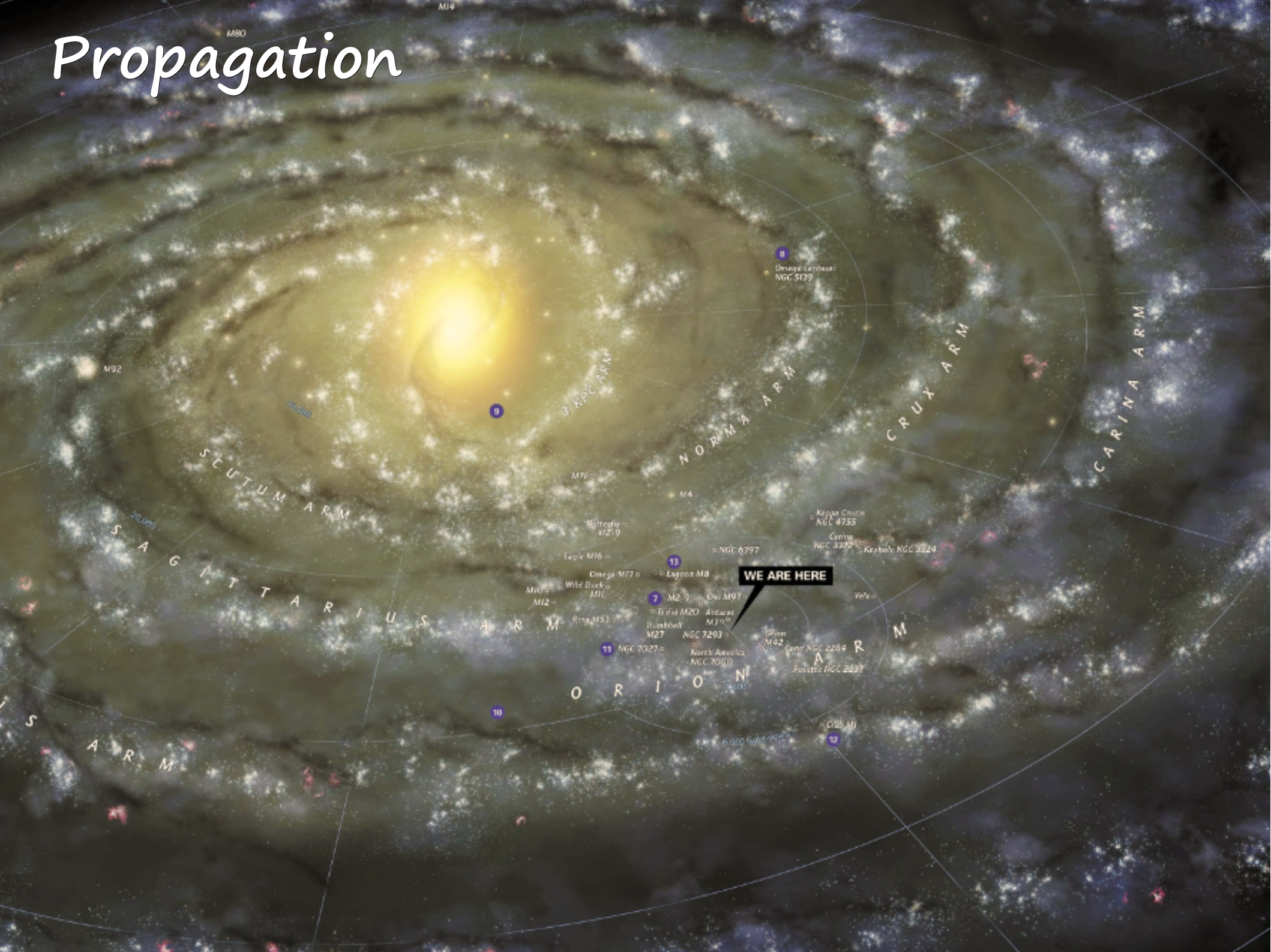


$$Q(E, \vec{r}) = \frac{1}{2} \frac{\rho^2(\vec{r})}{m_{\text{DM}}^2} \langle \sigma v \rangle \frac{dN}{dE}$$

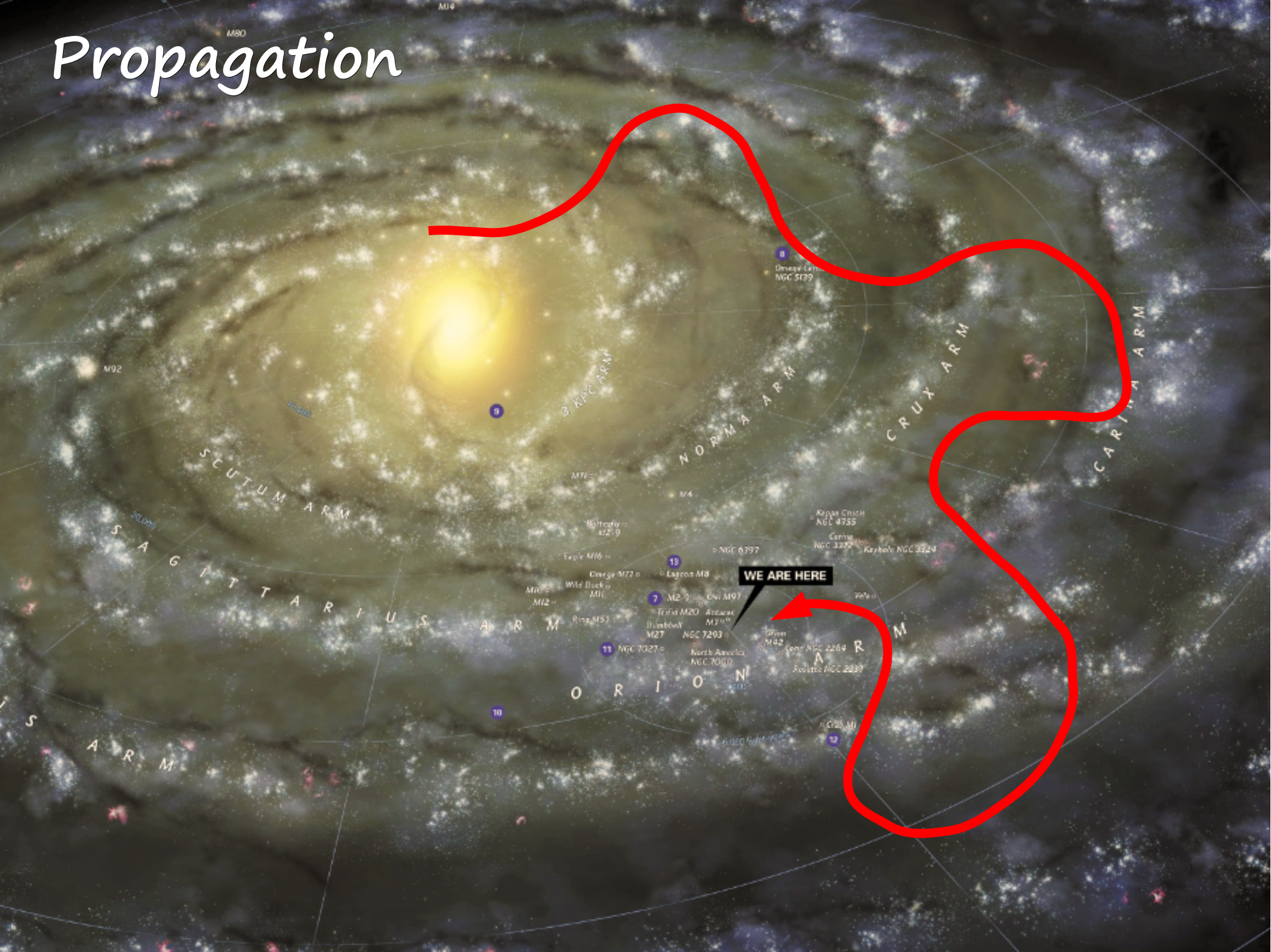


$$Q(E, \vec{r}) = \frac{\rho(\vec{r})}{m_{\text{DM}}} \frac{1}{\tau_{\text{DM}}} \frac{dN}{dE}$$

Propagation



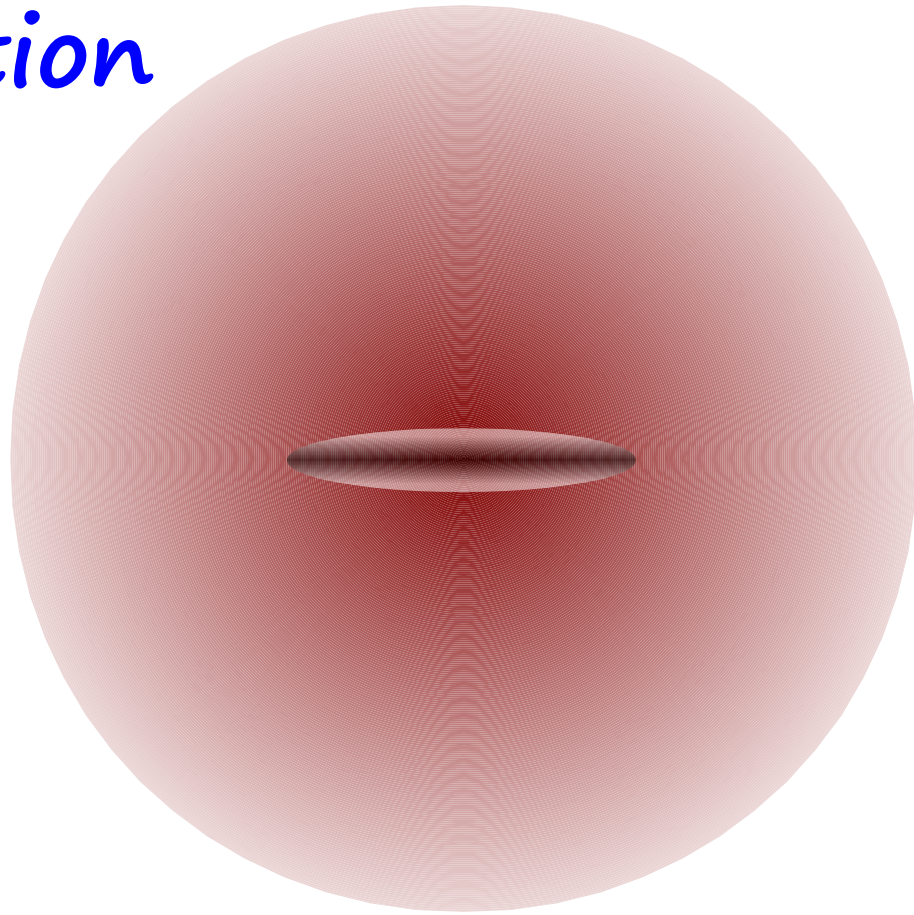
Propagation



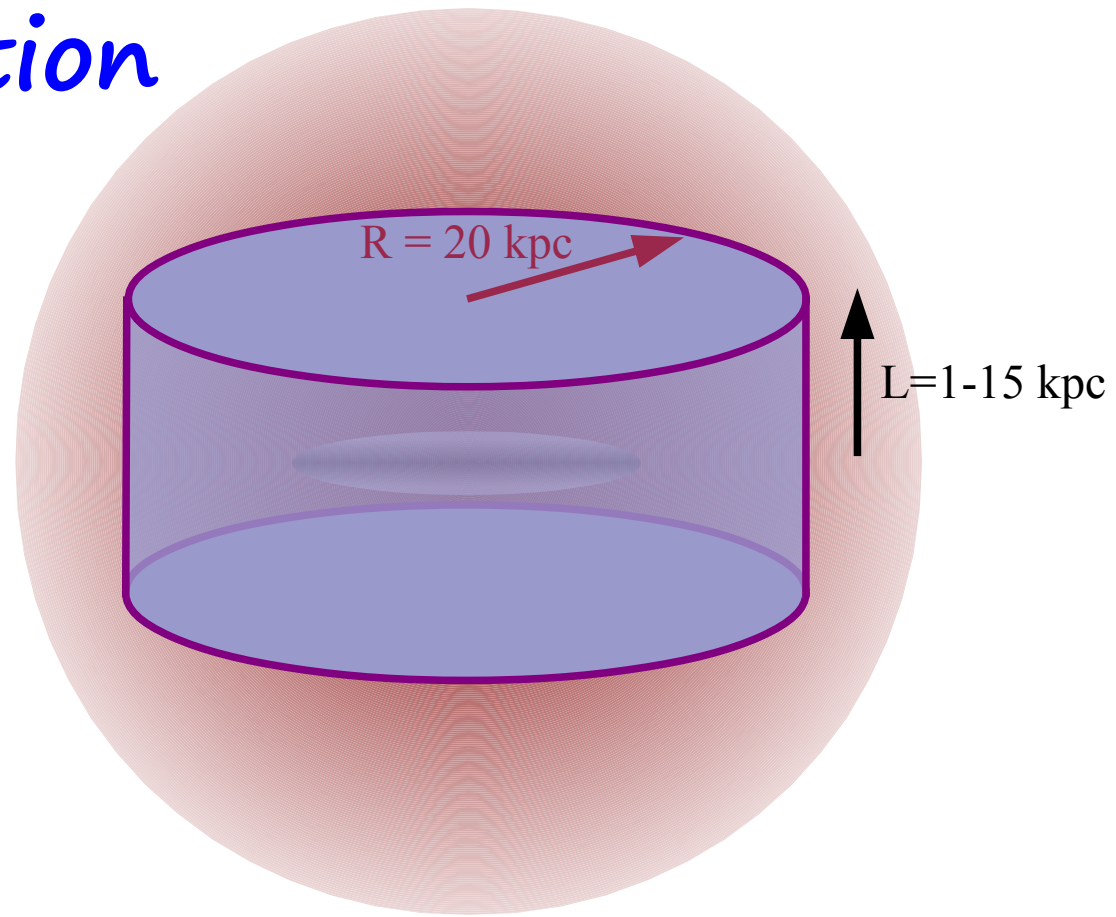
Propagation



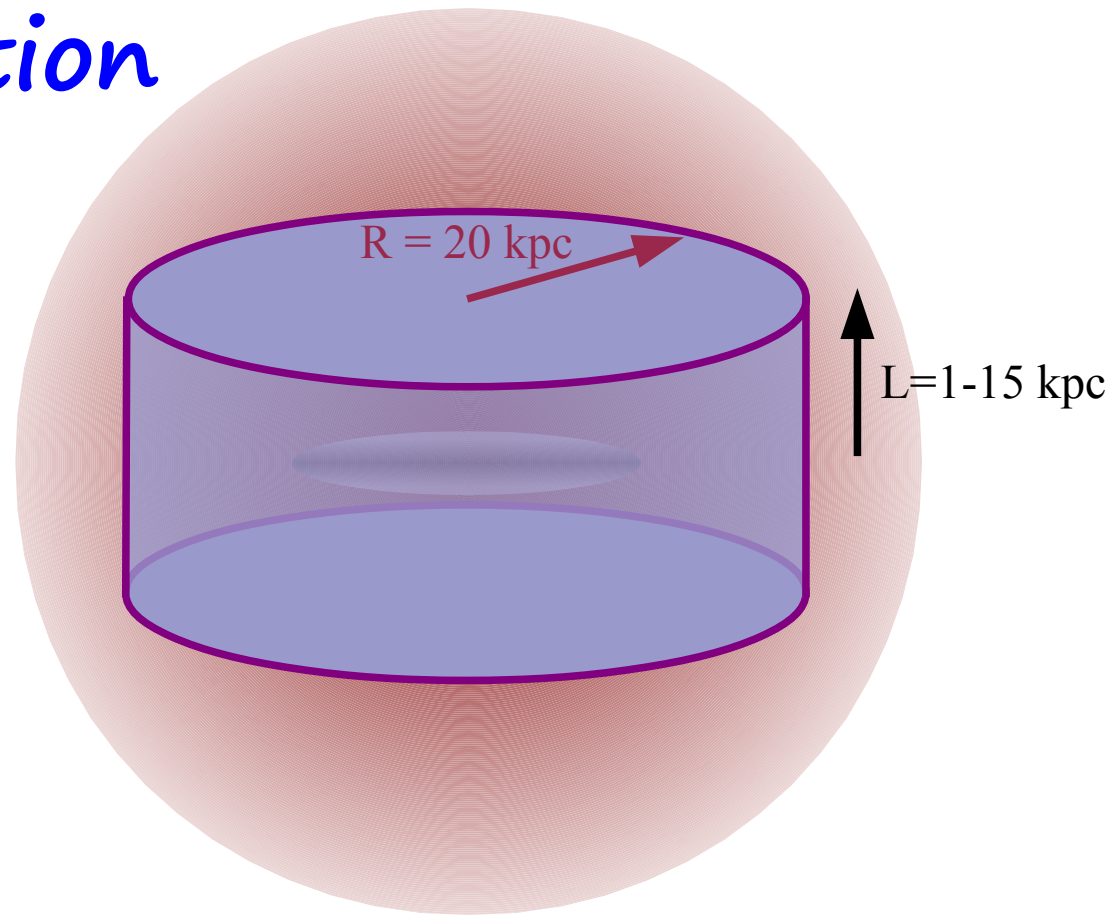
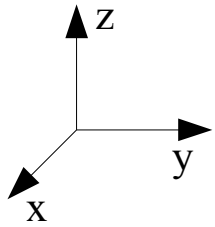
Propagation



Propagation



Propagation



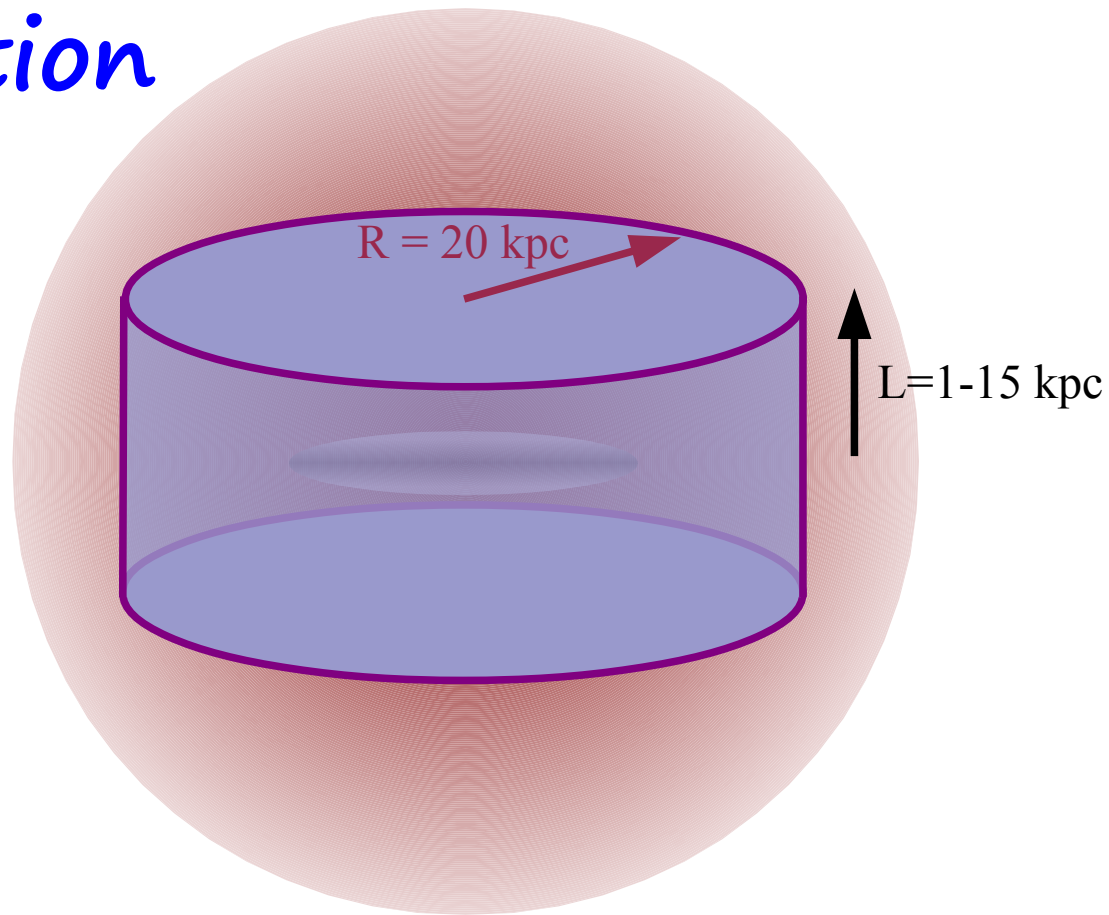
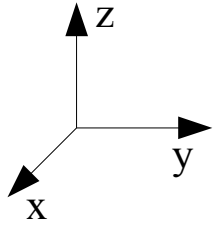
$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] + \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

f : number density of antiparticles per unit kinetic energy

interstellar antimatter flux:

$$\Phi^{\text{IS}}(T) = \frac{dN}{dt dS dT d\Omega} = \frac{v}{4\pi} f(T)$$

Propagation

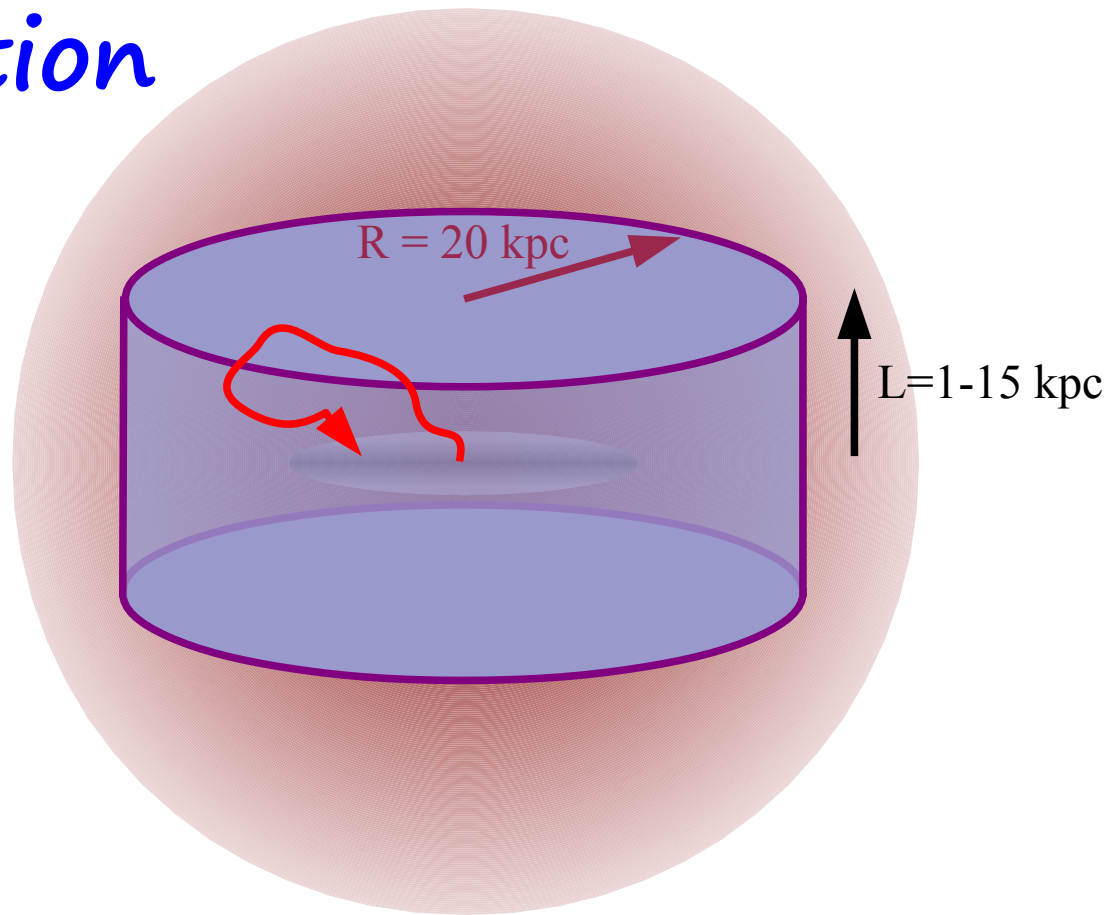
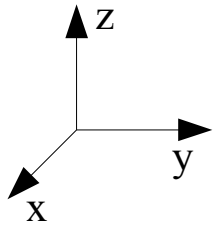


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Source term

$$Q(T, \vec{r}) = \begin{cases} \frac{1}{2} \frac{\rho^2(\vec{r})}{m_{\text{DM}}^2} \langle \sigma v \rangle \frac{dN}{dT} & \text{dark matter annihilation} \\ \frac{\rho(\vec{r})}{m_{\text{DM}}} \frac{1}{\tau_{\text{DM}}} \frac{dN}{dE} & \text{dark matter decay} \end{cases}$$

Propagation



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Source term

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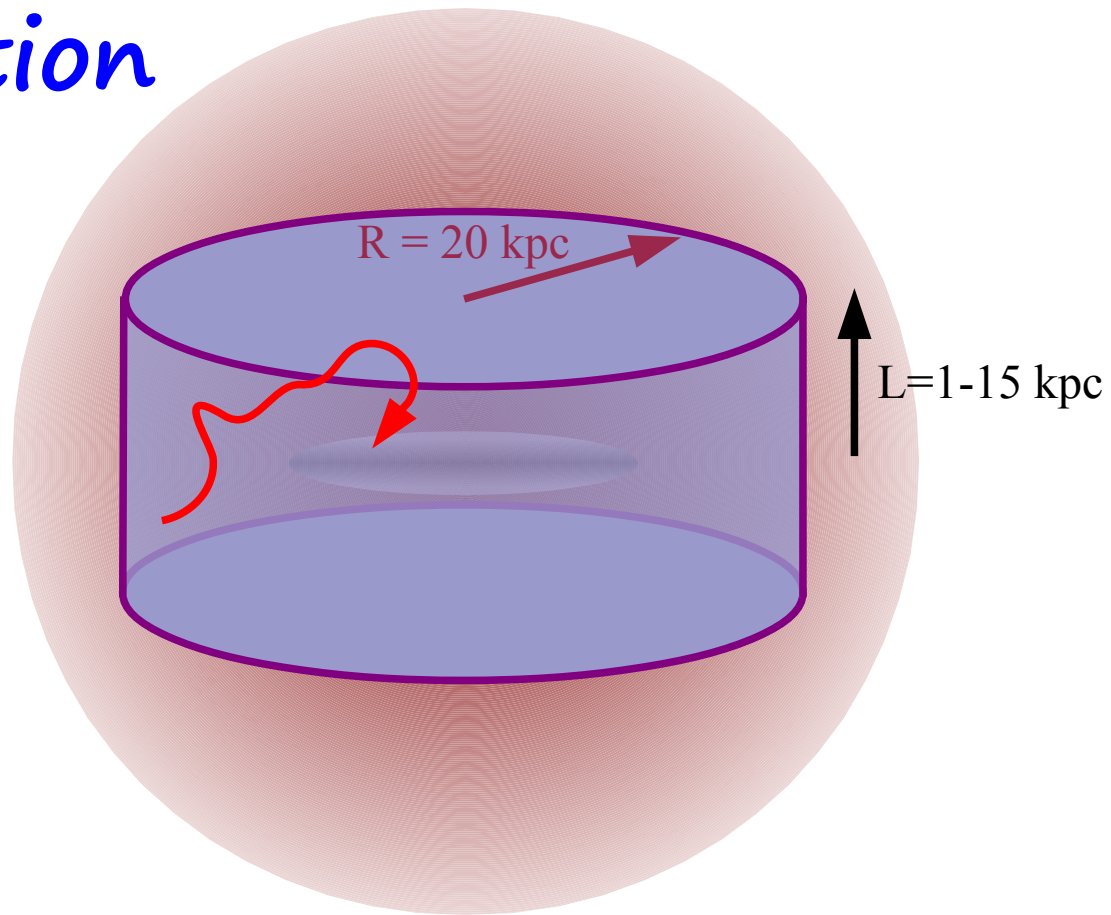
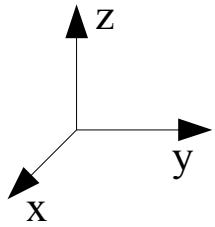
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dark matter annihilation

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dark matter decay

Propagation

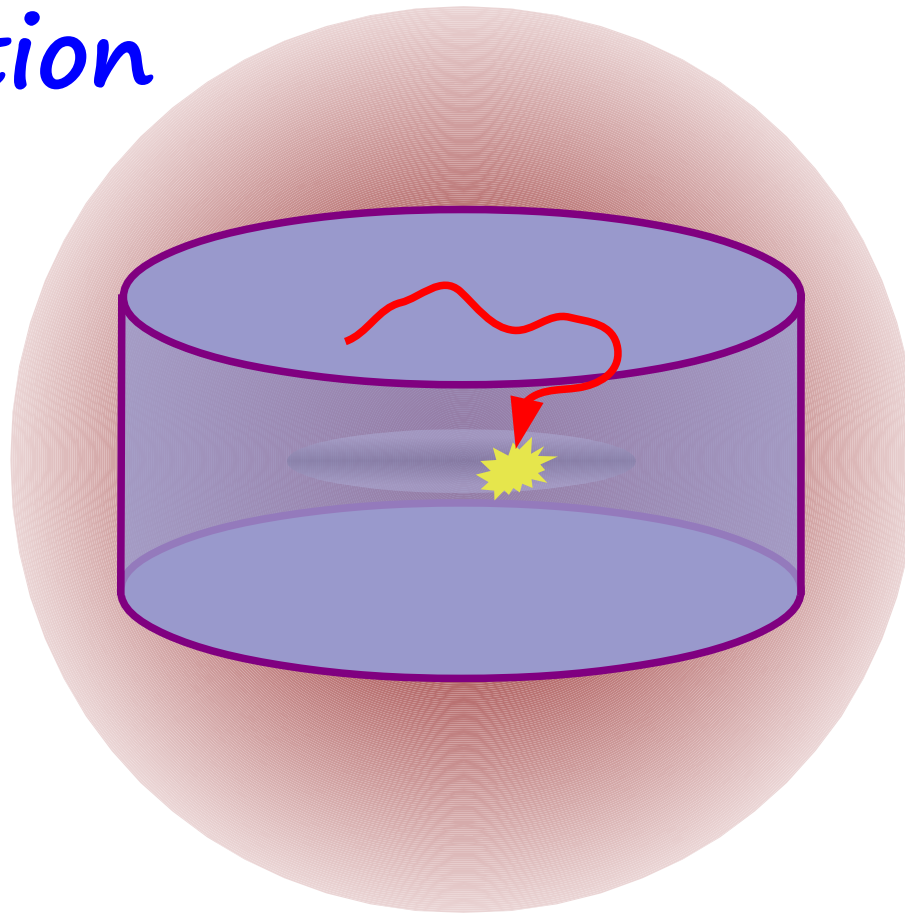
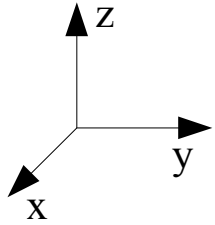


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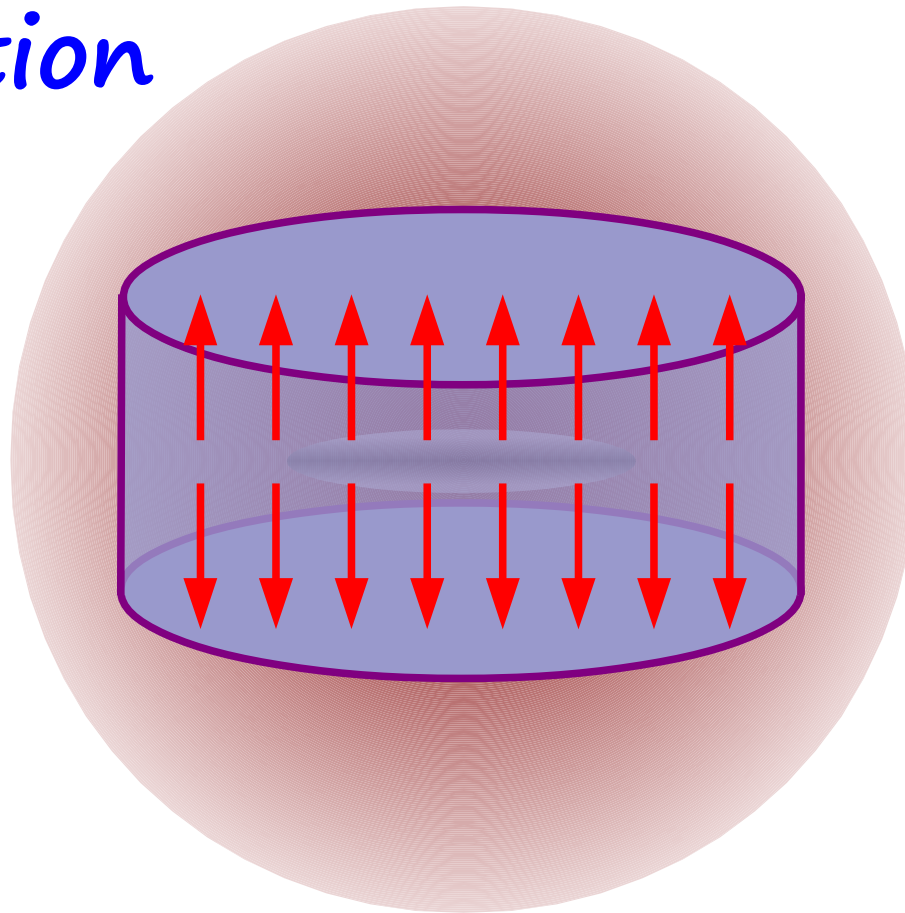
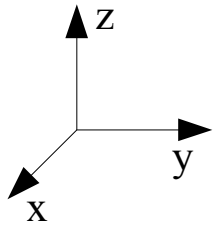
Annihilation term

Negligible for positrons.
For antiprotons,

$$\Gamma_{\text{ann}} = (n_{\text{H}} + 4^{2/3}n_{\text{He}})\sigma_{\bar{p}p}^{\text{ann}}v_{\bar{p}} .$$

$$\sigma_{\bar{p}p}^{\text{ann}}(T) = \begin{cases} 661 (1 + 0.0115 T^{-0.774} - 0.948 T^{0.0151}) \text{ mbarn} , & T < 15.5 \text{ GeV} , \\ 36 T^{-0.5} \text{ mbarn} , & T \geq 15.5 \text{ GeV} , \end{cases}$$

Propagation



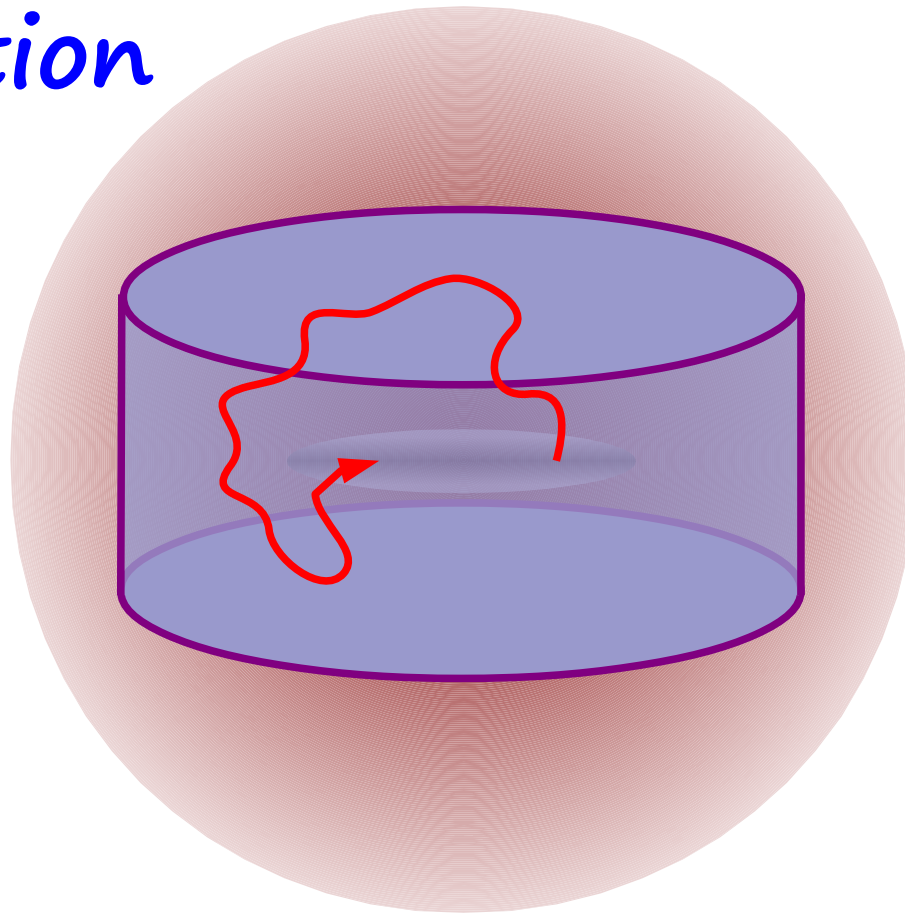
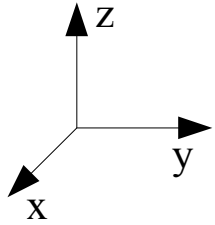
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Convection term

- Due to the Milky Way galactic wind.
- It drifts particles away from the Galactic disk.
- **Difficult to model.** Assume:

$$\vec{V}_c(\vec{r}) = V_e \text{sign}(z) \vec{k}$$

Propagation



$$0 = \frac{\partial f}{\partial t} = \nabla \cdot [K(T, \vec{r}) \nabla f] - \frac{\partial}{\partial T} [b(T, \vec{r}) f] - \nabla \cdot [\vec{V}_c(\vec{r}) f] - 2h\delta(z)\Gamma_{\text{ann}}f + Q(T, \vec{r}) .$$

Energy loss term

- Due to inverse Compton scattering on the interstellar radiation field (starlight, thermal radiation of dust, CMB) and synchrotron radiation.
- Negligible for antiprotons and antideuterons
- Can be modelled

- Energy loss due to Inverse Compton scattering: $e^+\gamma \rightarrow e^+\gamma$

$$b_{\text{ICS}}(E_e, \vec{r}) = \int_0^\infty d\epsilon \int_\epsilon^{E_\gamma^{\text{max}}} dE_\gamma (E_\gamma - \epsilon) \frac{d\sigma^{\text{IC}}(E_e, \epsilon)}{dE_\gamma} f_{\text{ISRF}}(\epsilon, \vec{r})$$

Number density
of photons in ISRF

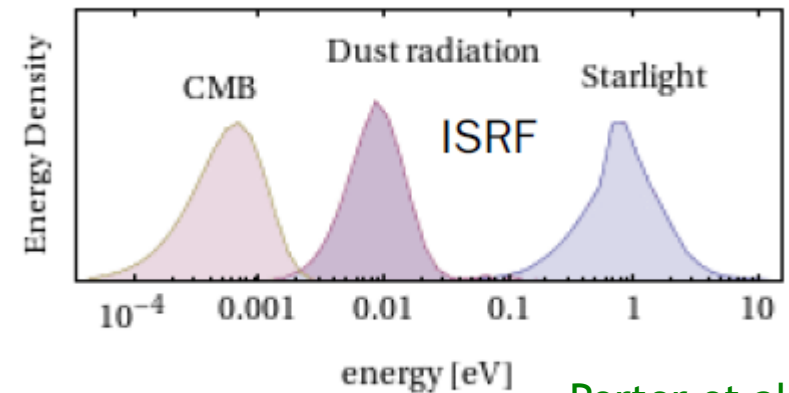
$$\frac{d\sigma^{\text{IC}}(E_e, \epsilon)}{dE_\gamma} = \frac{3}{4} \frac{\sigma_T}{\gamma_e^2 \epsilon} \times \left[2q \ln q + 1 + q - 2q^2 + \frac{1}{2} \frac{(q\Gamma)^2}{1 + q\Gamma} (1 - q) \right]$$

$\gamma_e = E_e/m_e \rightarrow$ Lorentz factor.

$\Gamma_e = 4 \gamma_e \epsilon/m_e$

$q = E_\gamma/\Gamma(E_e - E_\gamma)$

$\sigma_T = 0.67$ barn \rightarrow Compton scattering cross section
in the Thomson limit.



Porter et al.

- Energy loss due to synchrotron radiation:

$$b_{\text{sync}}(E_e, \vec{r}) = \frac{4}{3} \sigma_T \gamma_e^2 \frac{B^2}{2}$$

$$B = 6 \mu\text{G} \exp(-|z|/5\text{kpc} - r/20\text{kpc})$$

Approximately $b(E) = \frac{E^2}{E_0 \tau_E}$, with $E_0 = 1$ GeV and $\tau_E = 10^{16}$ s

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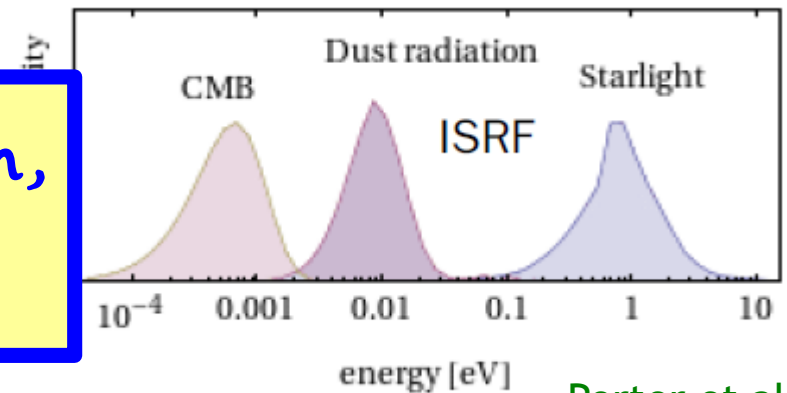
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Not very well known, though...



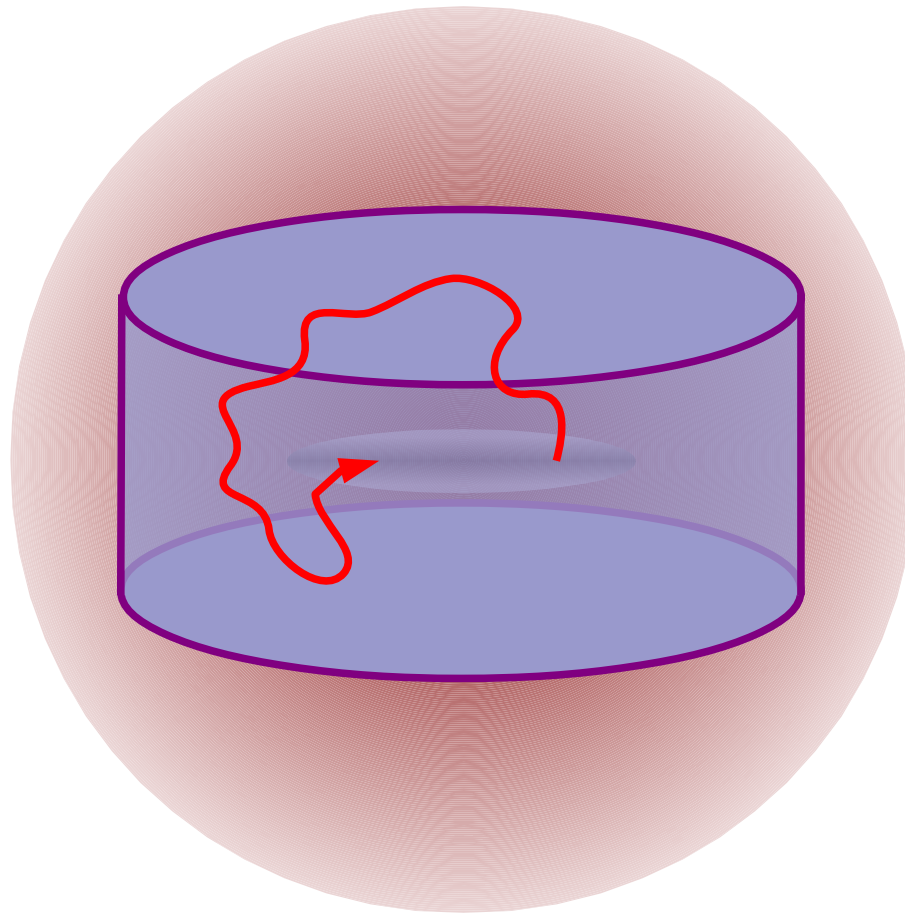
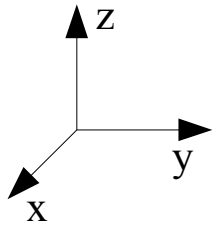
Porter et al.

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Approximately $b(E) = \frac{E^2}{E_0 \tau_E}$, with $E_0 = 1$ GeV and $\tau_E = 10^{16}$ s



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Diffusion term

- Due to the tangled magnetic field of the Galaxy.
- **Difficult to model.** Assume

$$K(T) = K_0 \beta \mathcal{R}^\delta$$

$$\left(\begin{array}{l} \beta = \text{velocity} \\ \mathcal{R} = \text{rigidity} \end{array} \right)$$

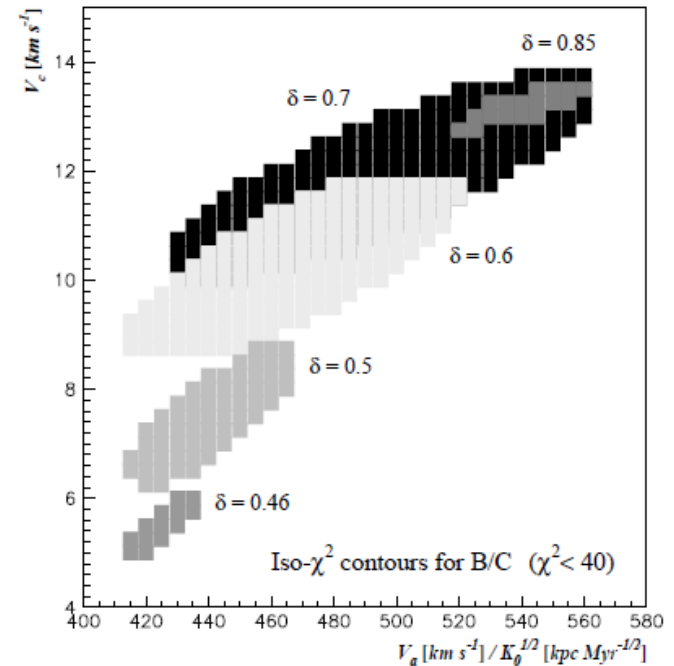
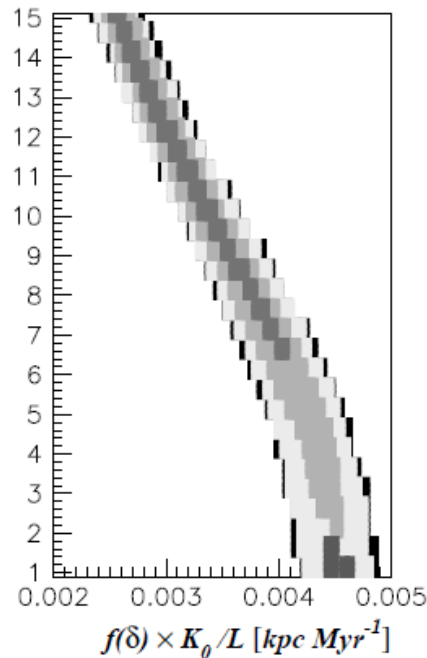
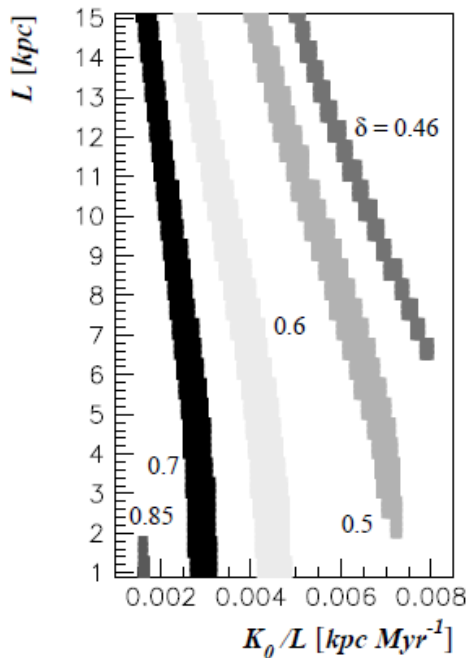
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$$K(T) = K_0 \beta \mathcal{R}^\delta$$

$$\vec{V}_c(\vec{r}) = V_c \text{sign}(z) \vec{k}$$

K_0 , δ , V_c (as well as L) must be determined with measurements of other cosmic ray species (mainly B/C ratio).

Iso- χ^2 contours for B/C ($\chi^2 < 40$)

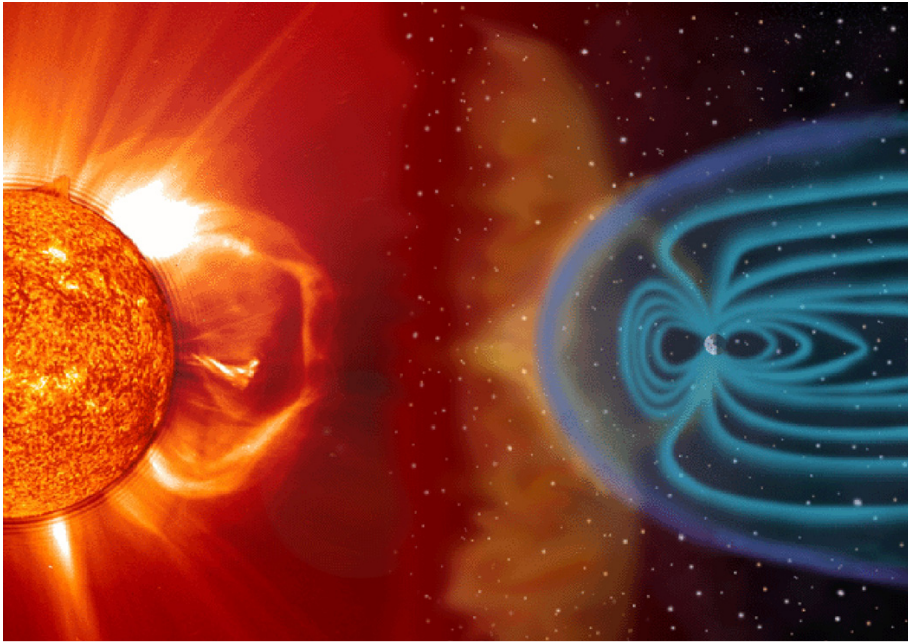


Model	δ	K_0 (kpc ² /Myr)	L (kpc)	V_c (km/s)
MIN	0.85	0.0016	1	13.5
MED	0.70	0.0112	4	12
MAX	0.46	0.0765	15	5

Maurin, Donato, Taillet, Salati '01



Propagation *inside* the Solar System



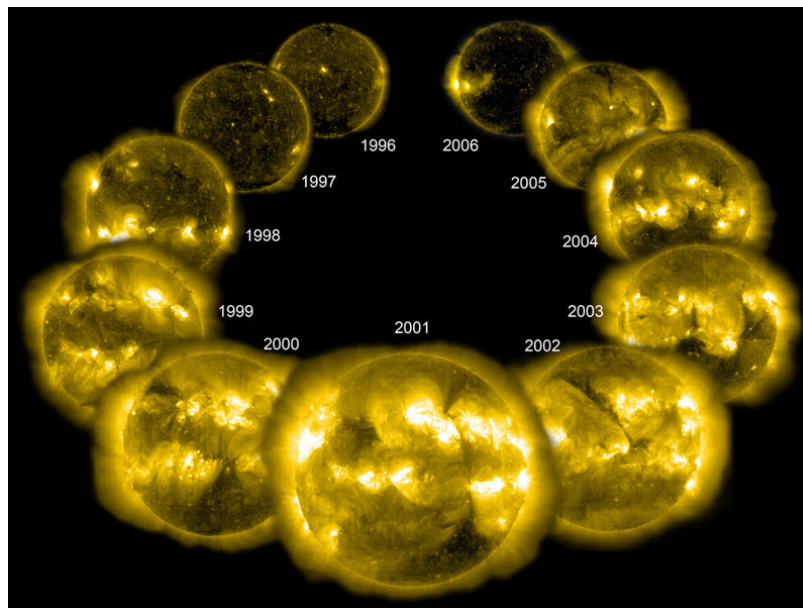
In the “force field approximation”, the flux at the top of the atmosphere (TOA) is related to the interstellar flux (IS) by

$$\Phi_{e^\pm}^{\text{TOA}}(E_{\text{TOA}}) = \frac{E_{\text{TOA}}^2}{E_{\text{IS}}^2} \Phi_{e^\pm}^{\text{IS}}(E_{\text{IS}})$$

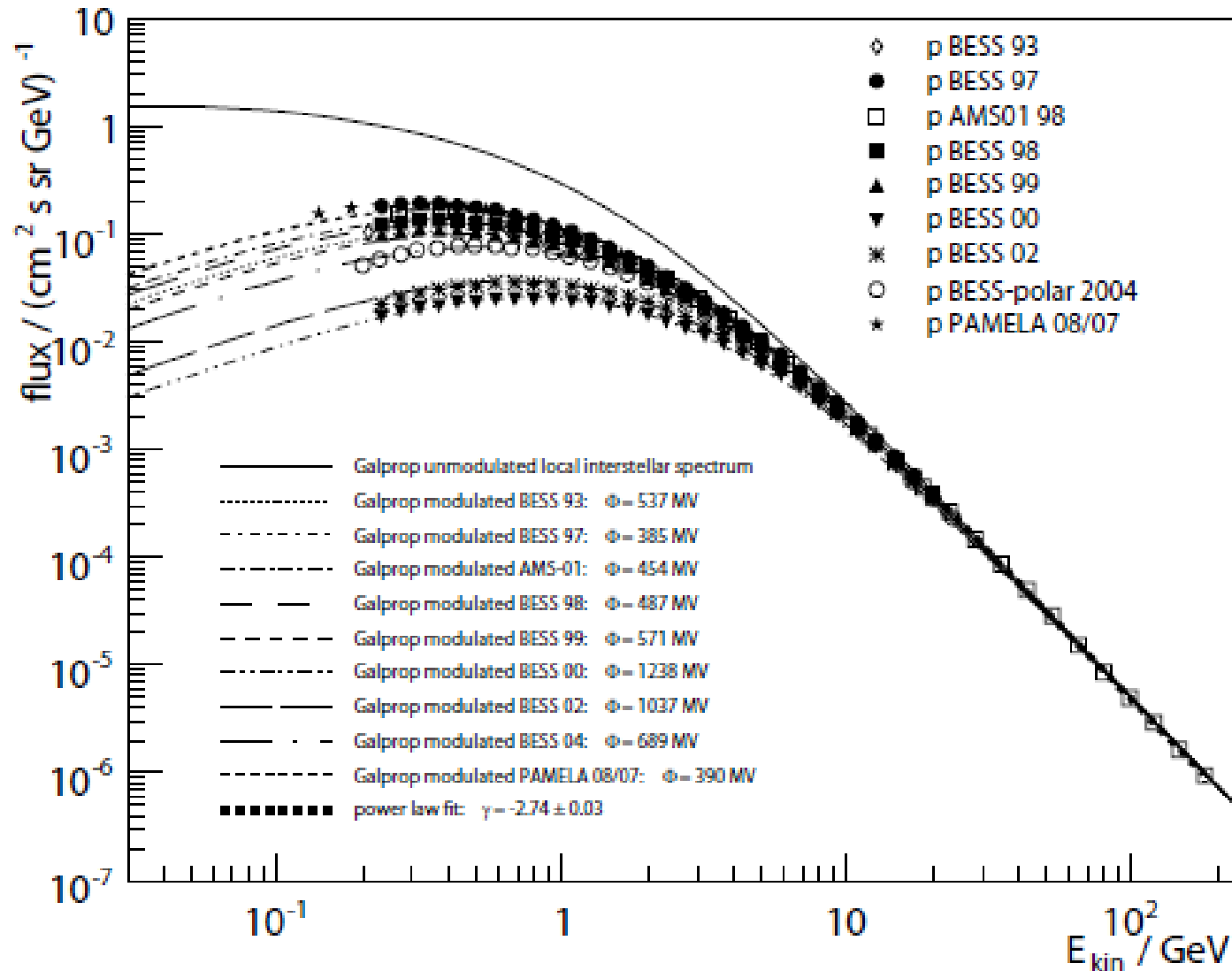
$$E_{\text{IS}} = E_{\text{TOA}} + \phi_F$$

↓
solar modulation parameter

$$\phi_F = 500 \text{ MV} - 1.3 \text{ GV}$$

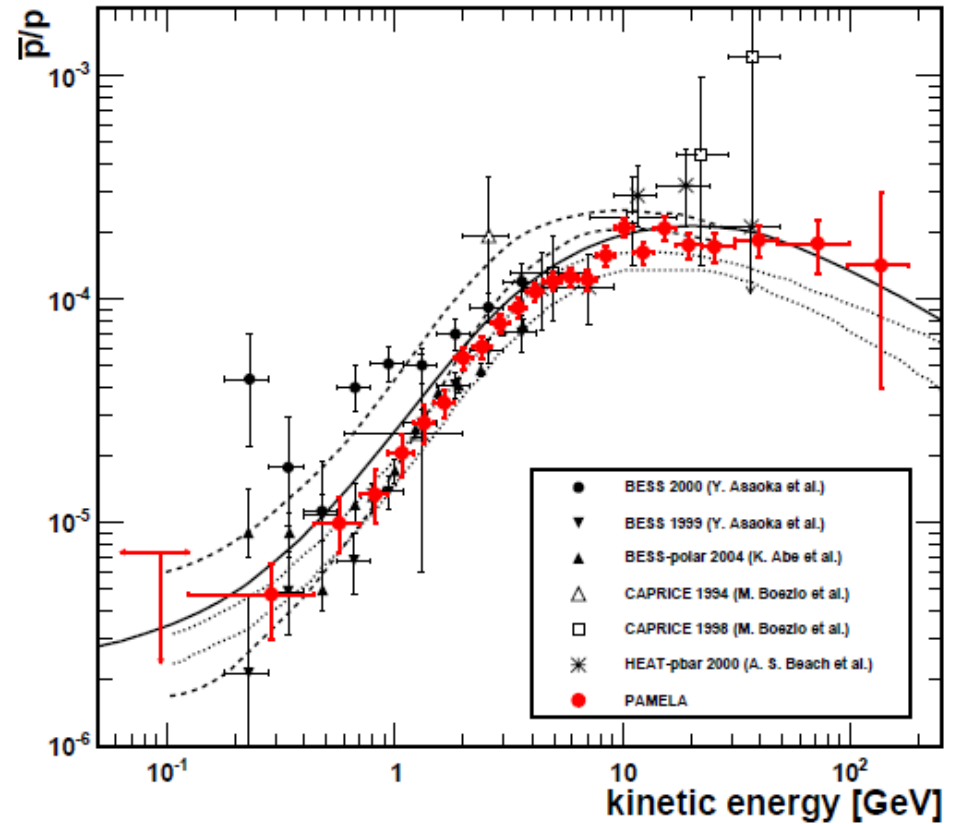
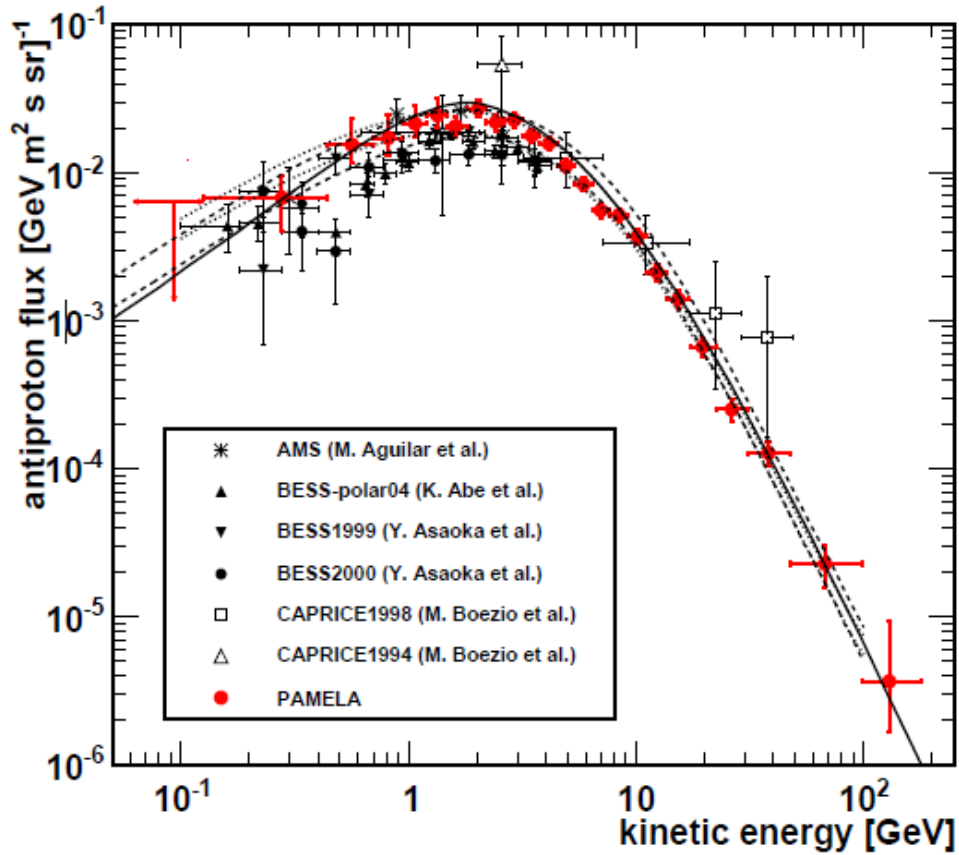


Cosmic ray **proton** spectrum as measured by BESS, AMS-01 and PAMELA



Gast, Schael '09

Experimental results: antiprotons



PAMELA collaboration
arXiv:1007.0821

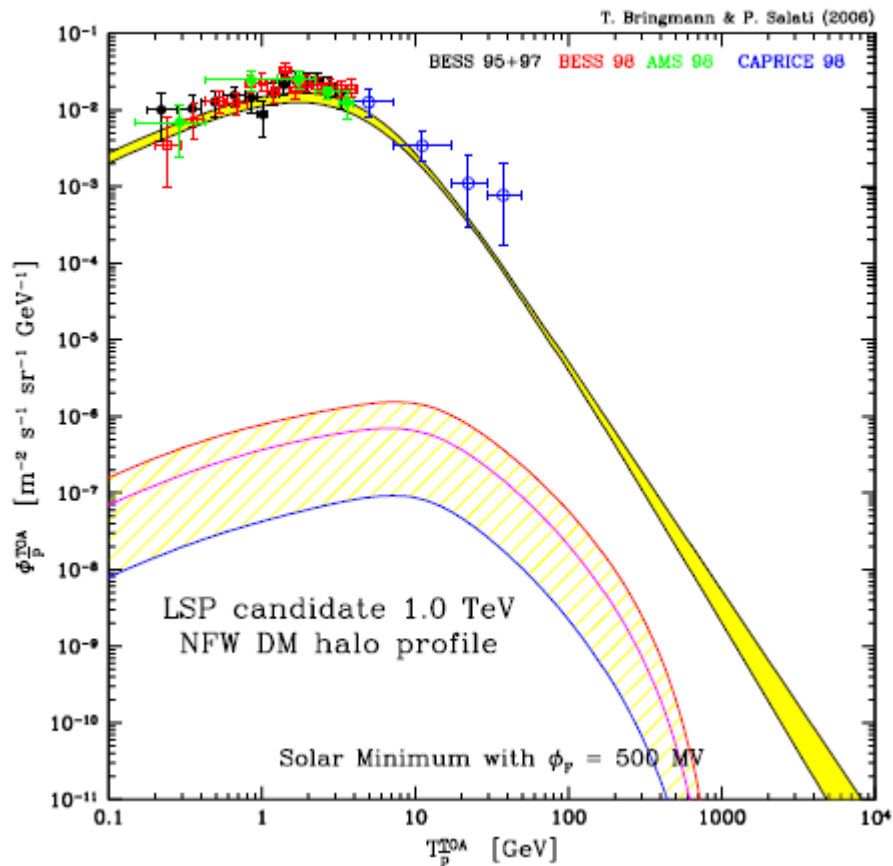
Fairly good agreement between the measurements and the theoretical predictions from collisions of cosmic rays on the interstellar medium $p p \rightarrow \bar{p} X$

Expectations from theory

A concrete example in the minimal supersymmetric standard model.

TeV $\times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$

DM model	m	$\langle \sigma_{\text{ann}} v \rangle$	$t\bar{t}$	$b\bar{b}$	$c\bar{c}$	$s\bar{s}$	$u\bar{u}$	$d\bar{d}$	ZZ	W^+W^-	HH	gg
LSP1.0	1.0	0.46	-	-	-	-	-	-	-	100	-	-

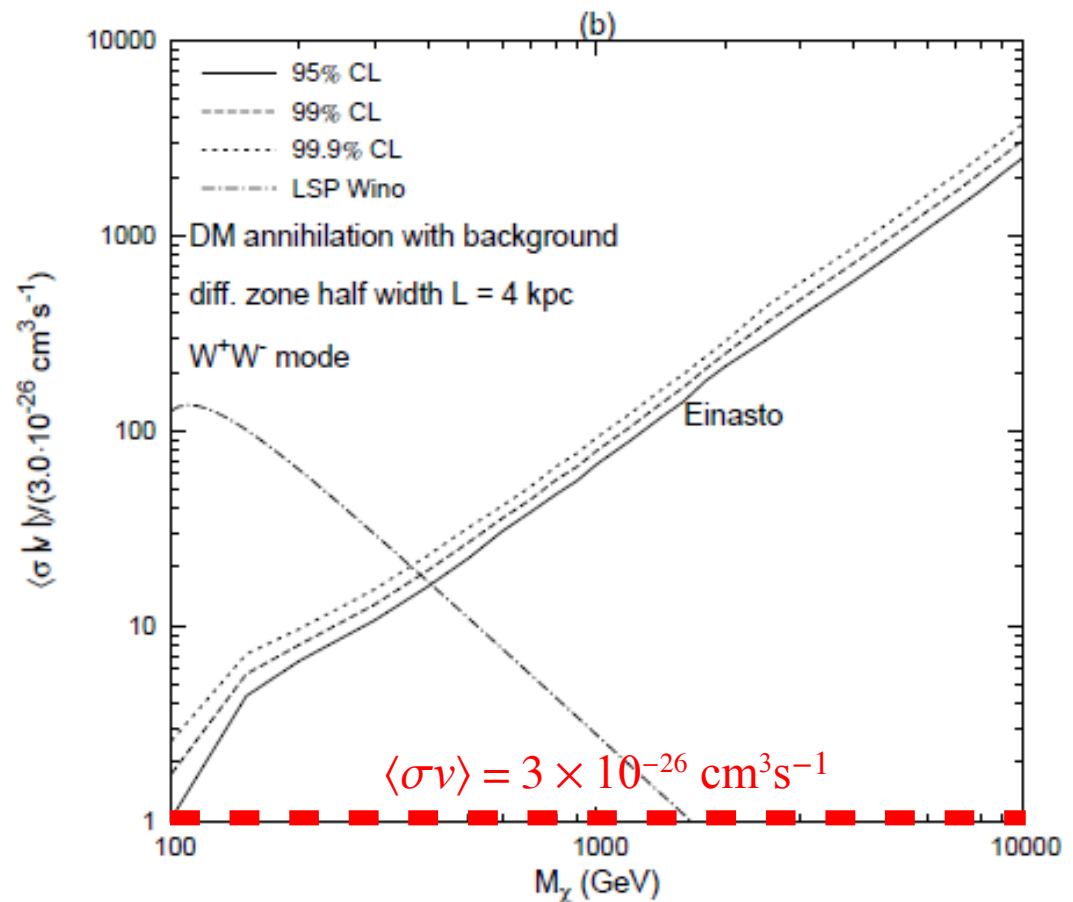
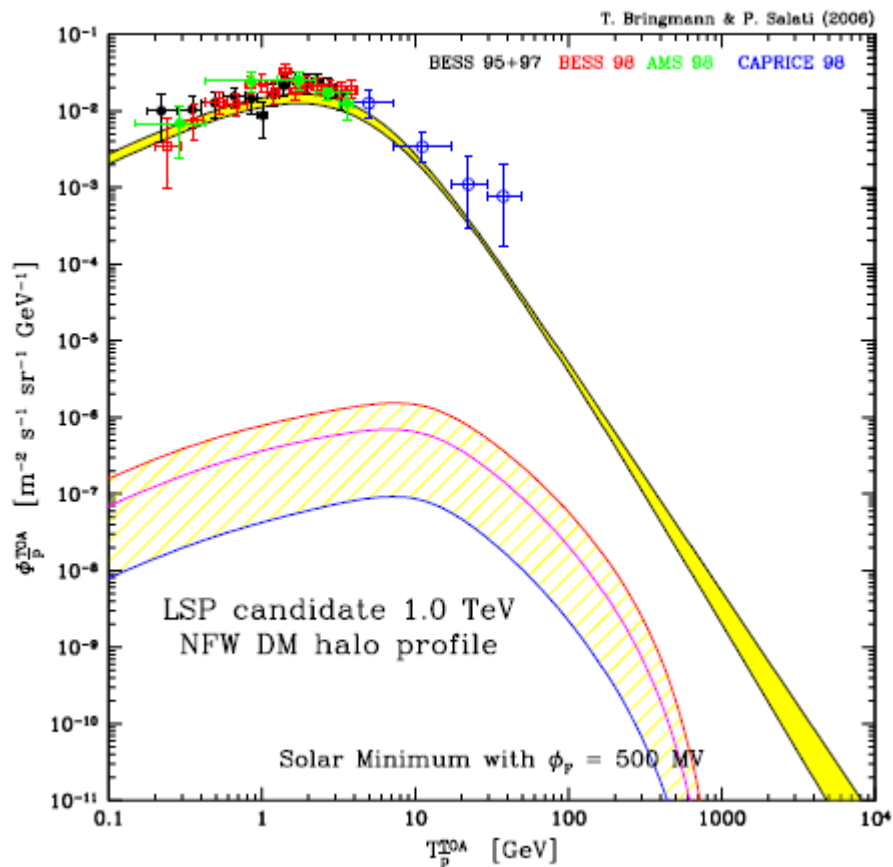


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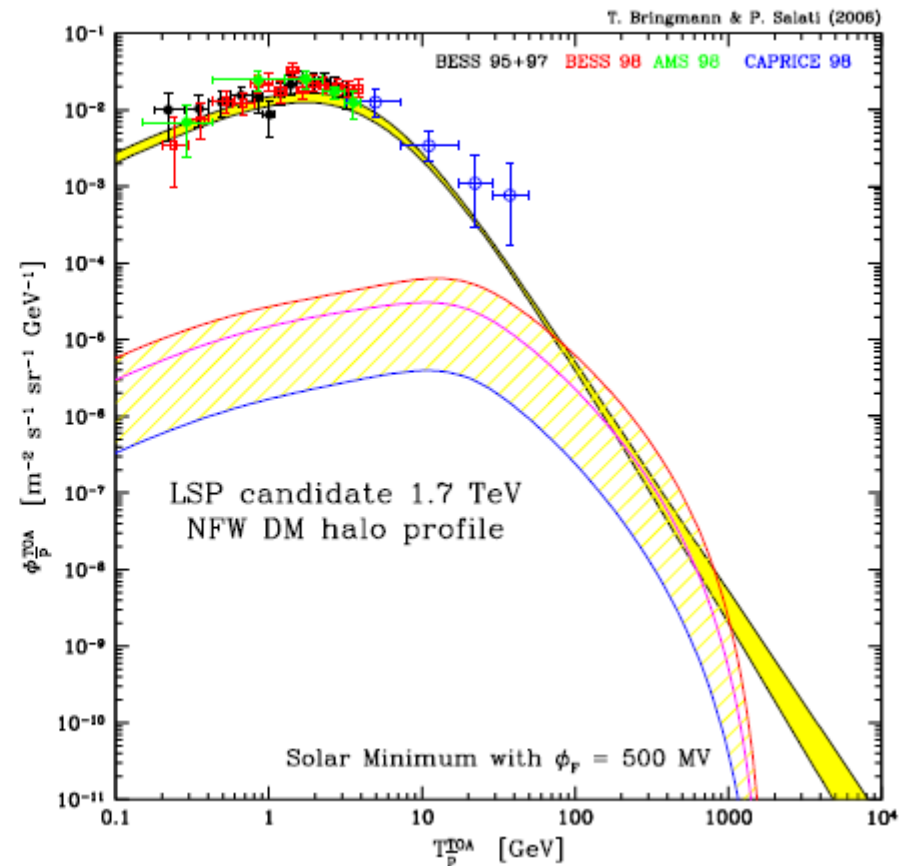
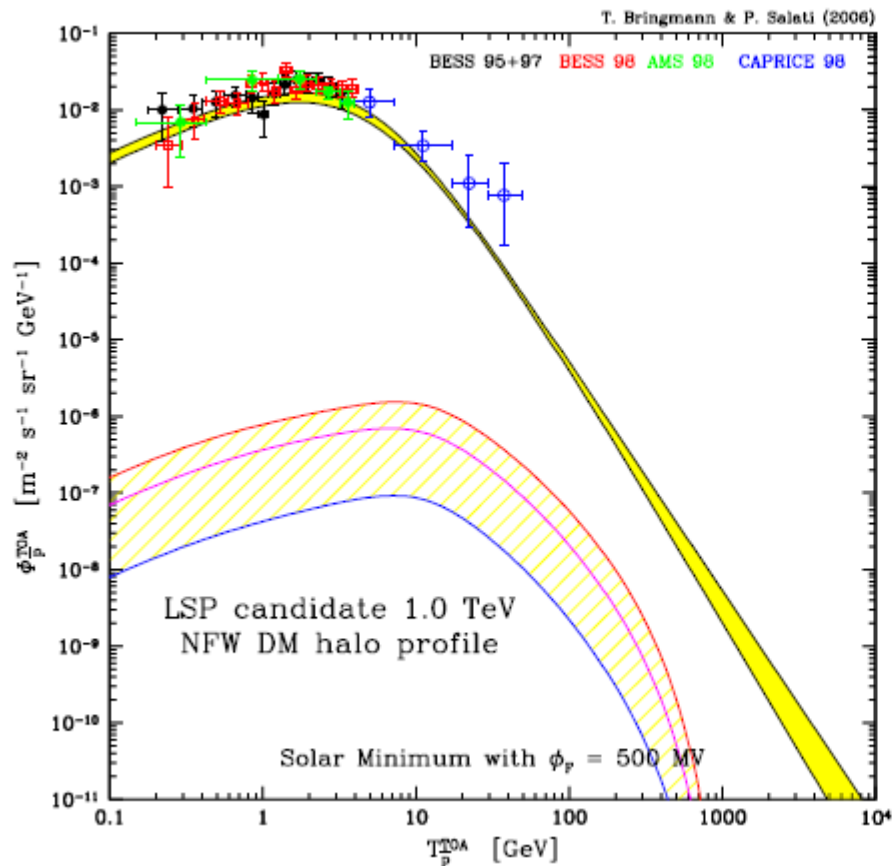


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LSP1.0	1.0	0.46	-	-	-	-	-	-	-	100	-	-
LSP1.7	1.7	102	-	-	-	-	-	-	20.1	79.9	-	-



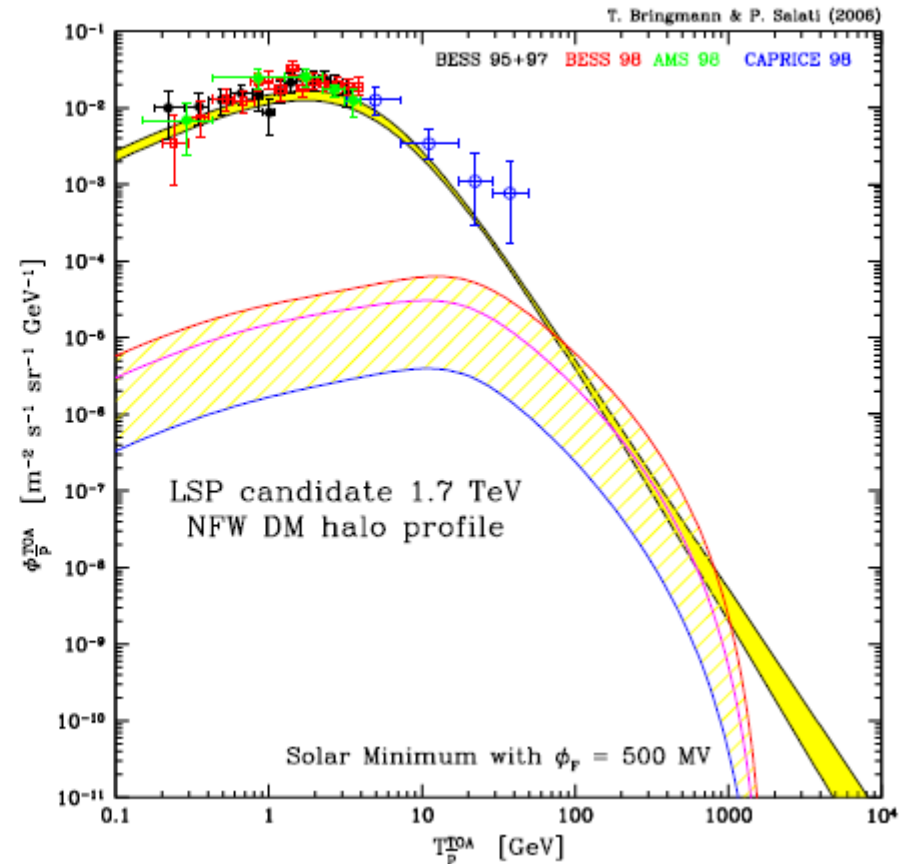
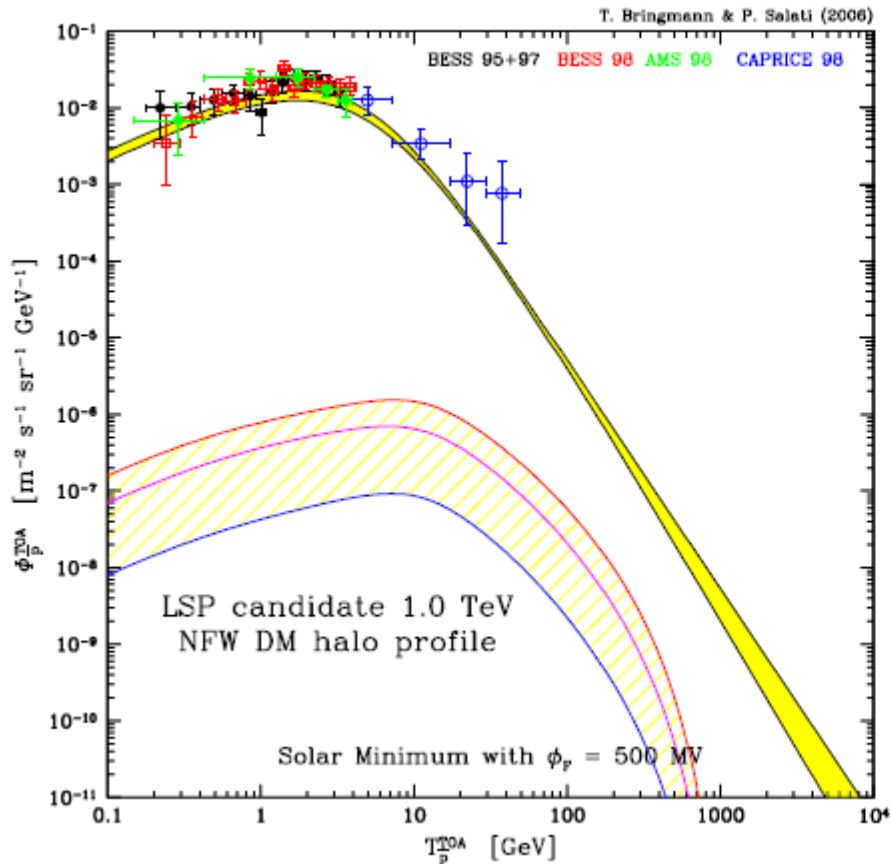
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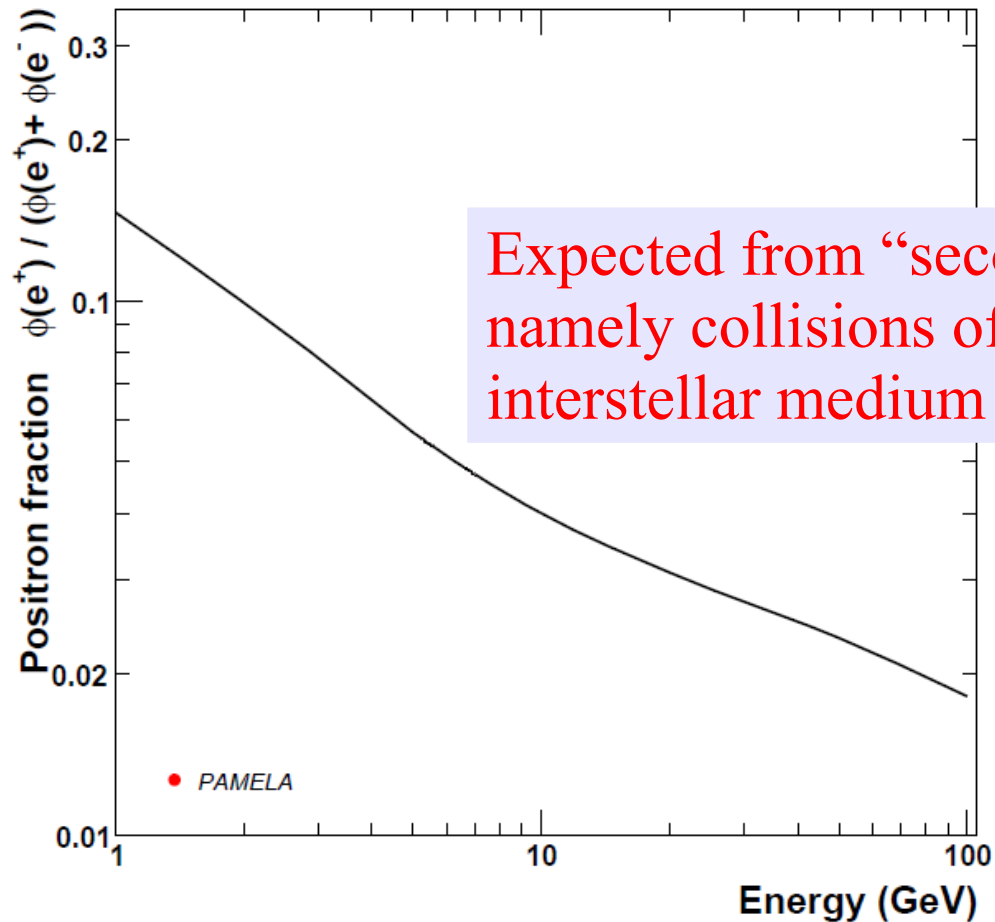
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DM model	m	$\langle \sigma_{\text{ann}} v \rangle$	$t\bar{t}$	$b\bar{b}$	$c\bar{c}$	$s\bar{s}$	$u\bar{u}$	$d\bar{d}$	ZZ	W^+W^-	HH	gg
LSP1.0	1.0	0.46	-	-	-	-	-	-	-	100	-	-
LSP1.7	1.7	102	-	-	-	-	-	-	100	500	-	-

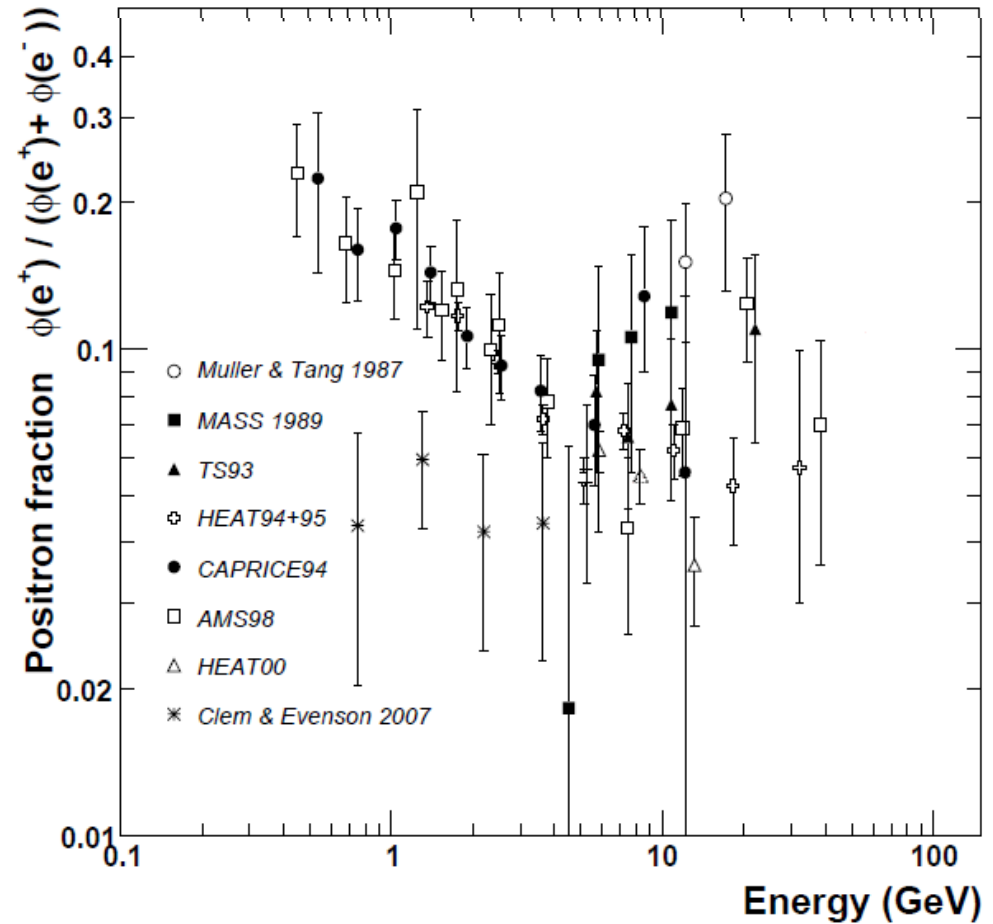
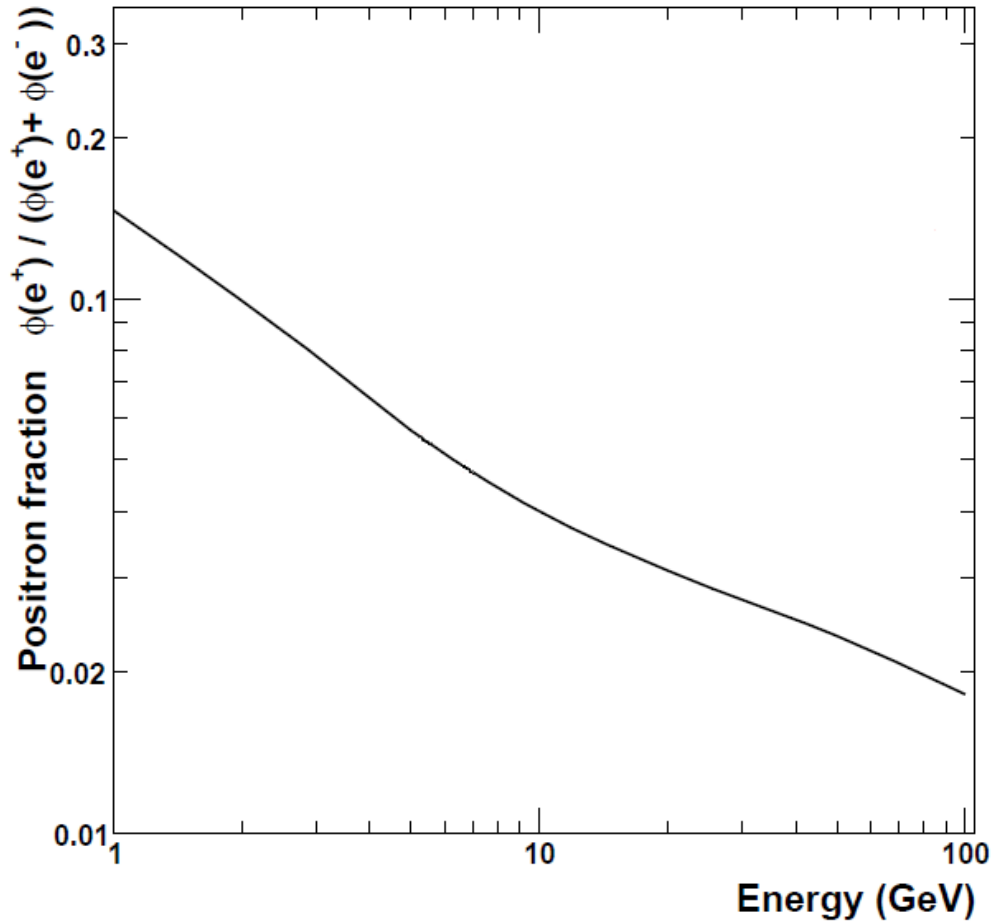
Annihilation rate "boosted"!



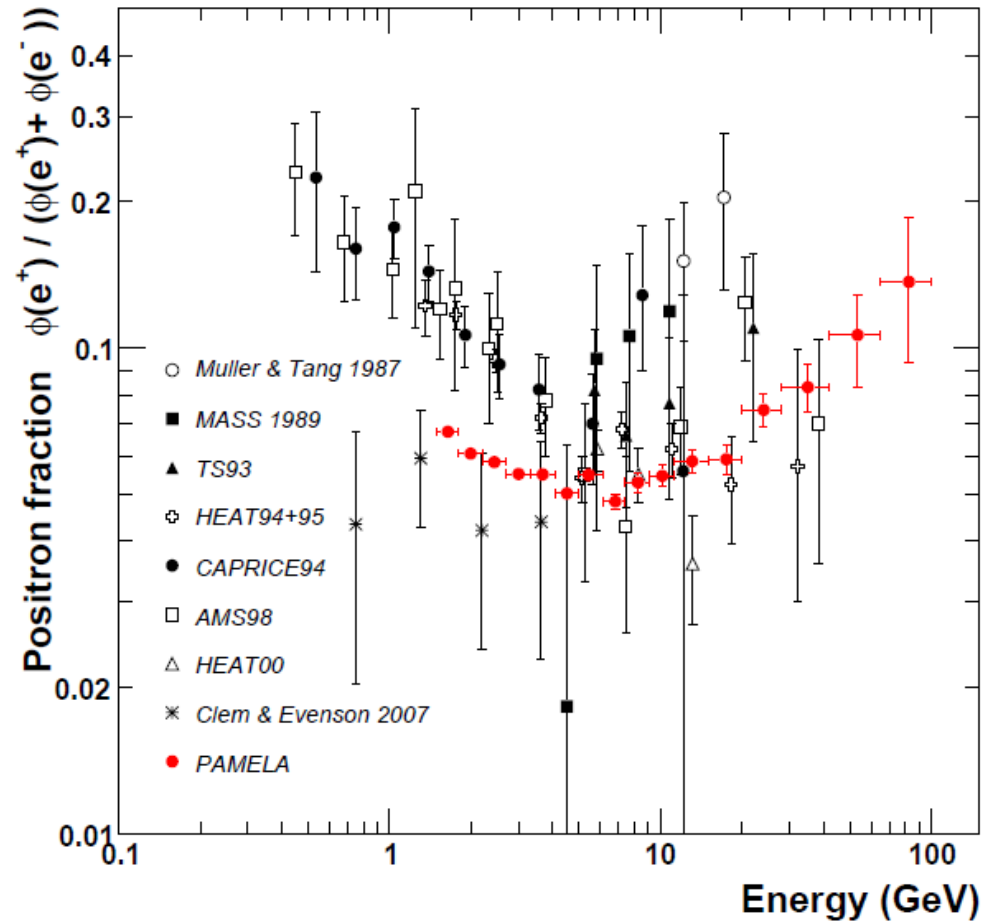
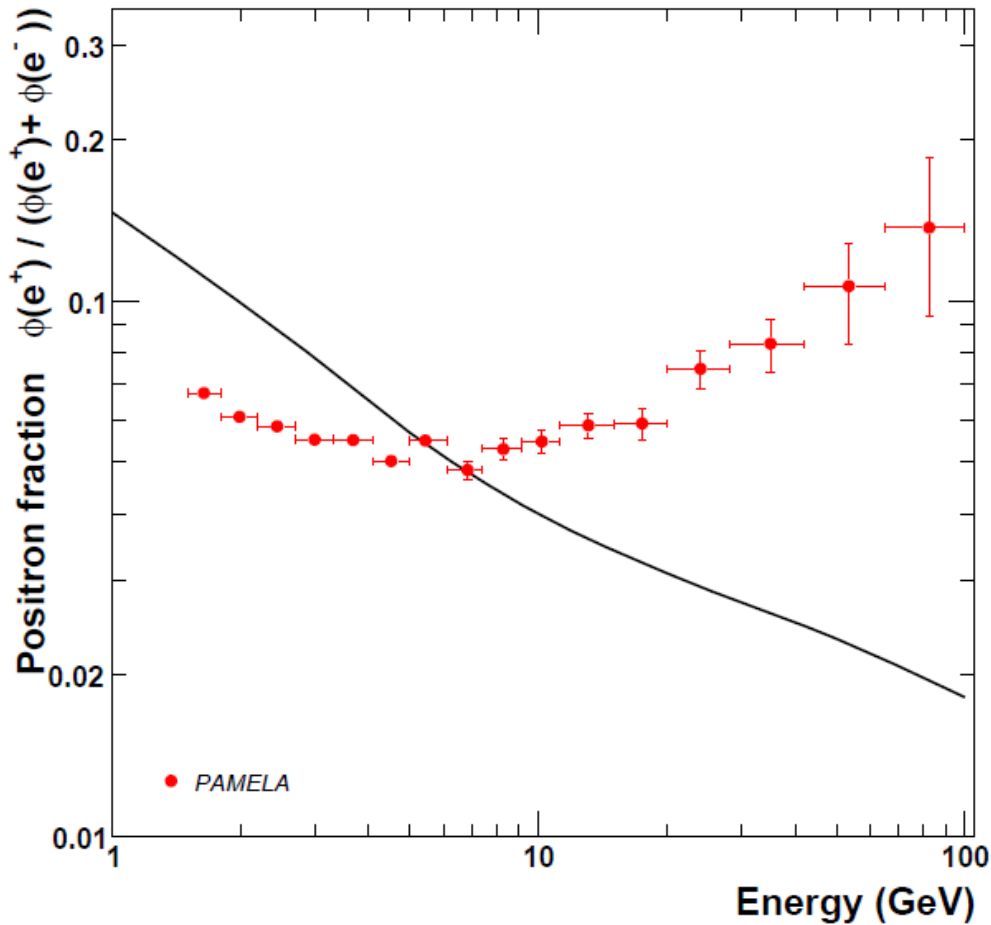
Experimental results: positrons



Experimental results: positrons

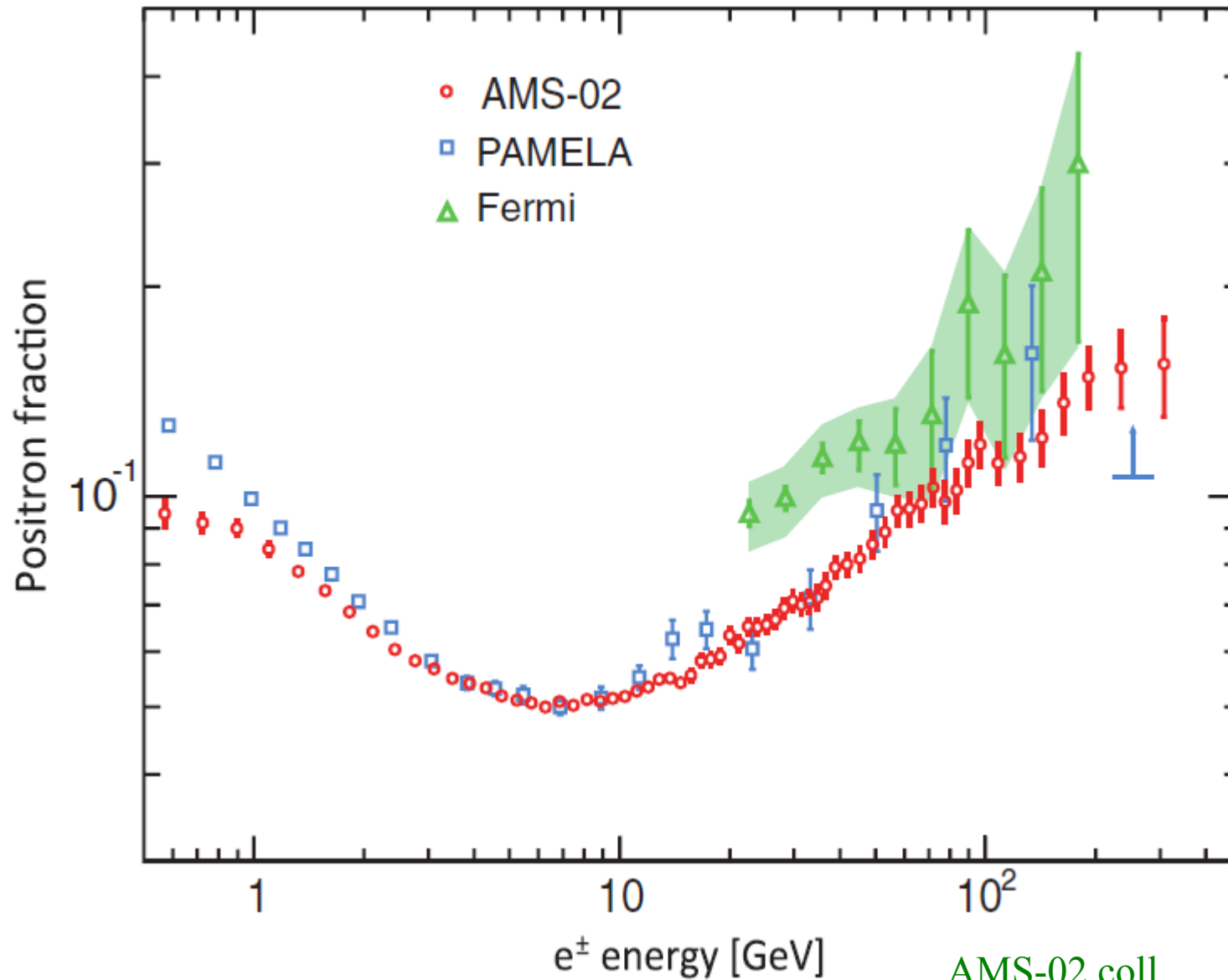


Experimental results: positrons



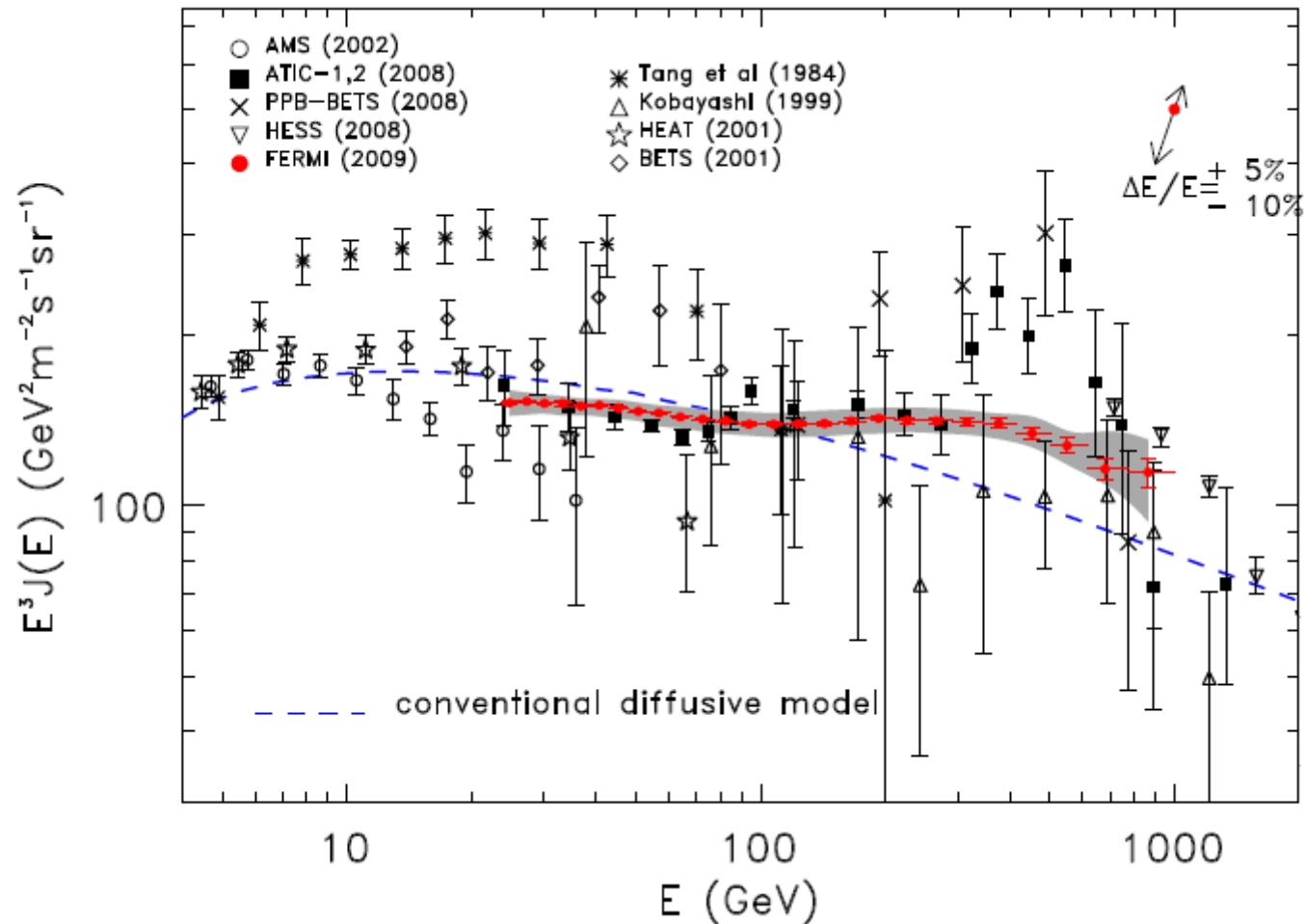
PAMELA coll.
arXiv:0810.4995

Experimental results: positrons



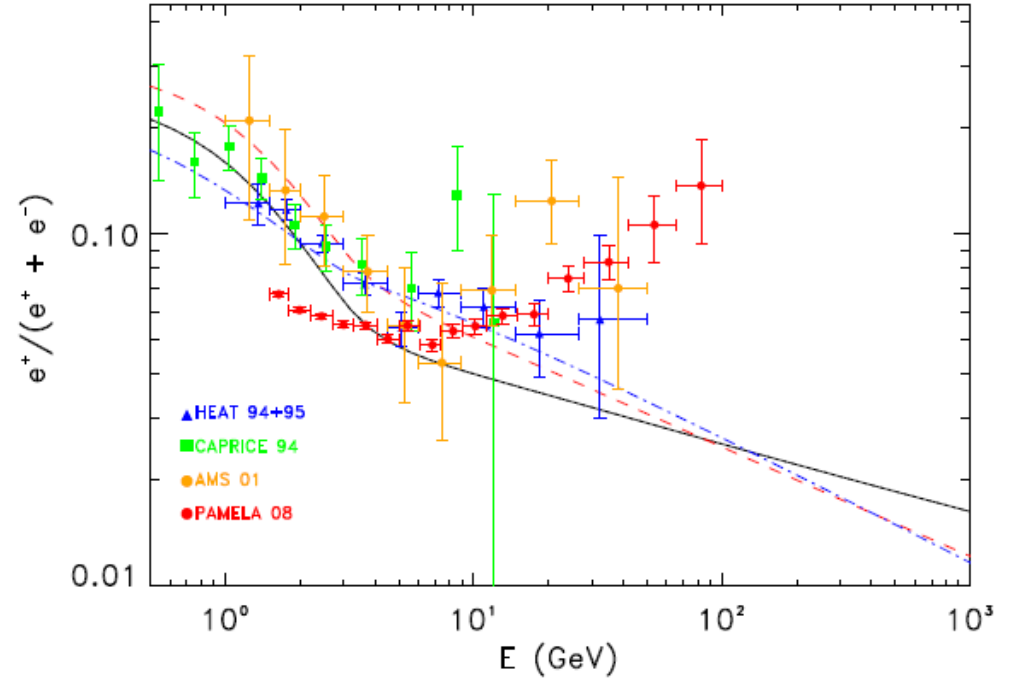
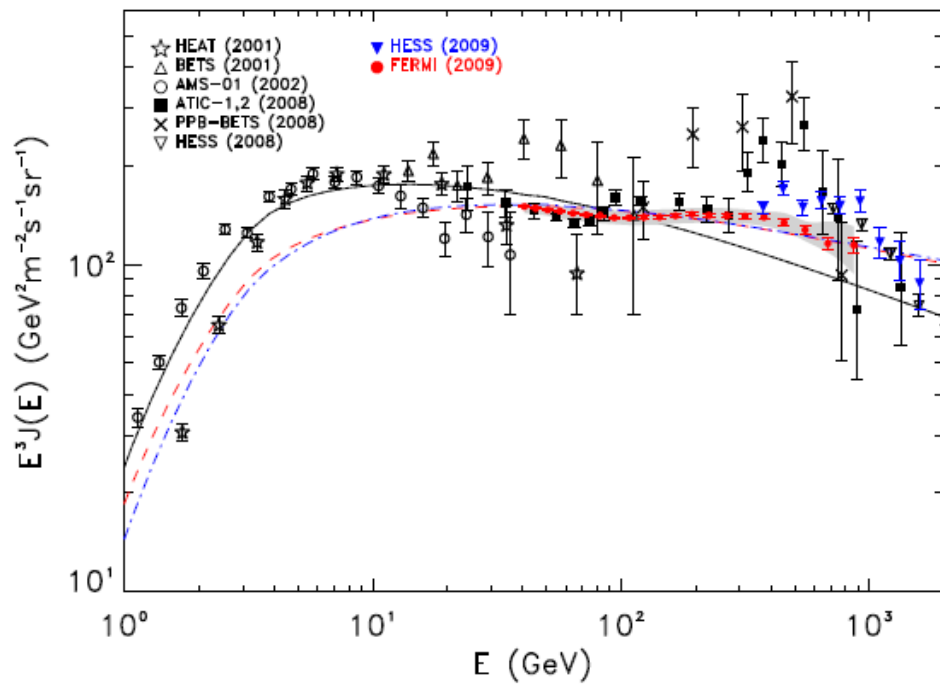
AMS-02 coll.
Phys.Rev.Lett. 110 (2013) 14, 141102

More puzzles: the electron+positron flux



Abdo et al.
ArXiv:0905.0025

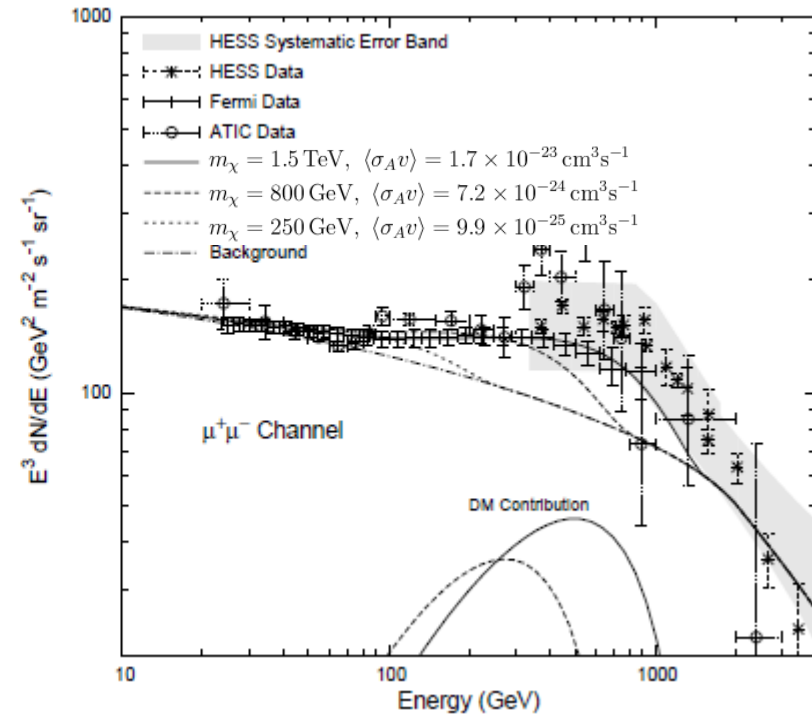
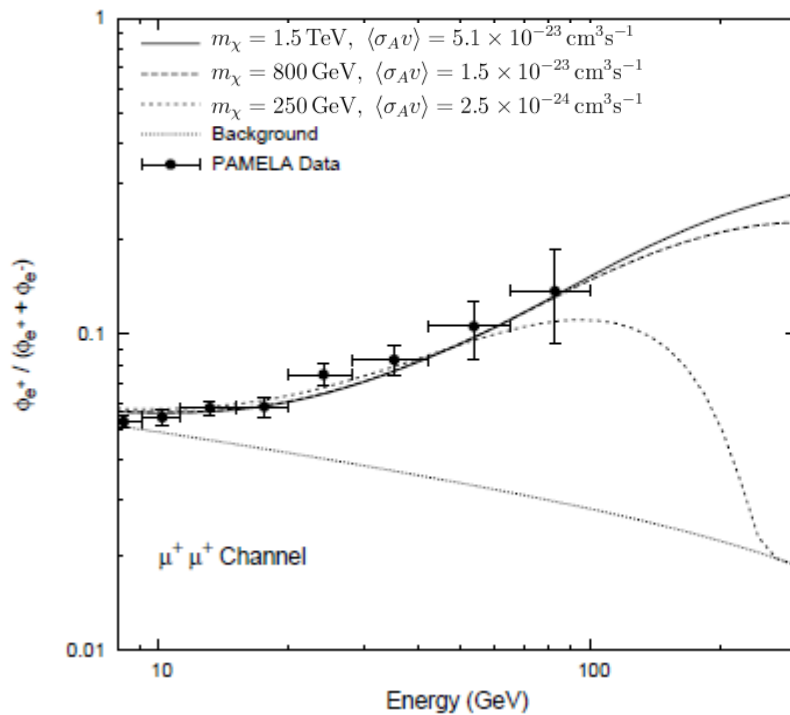
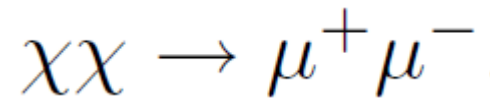
Present situation:



**Evidence for a primary component of positrons
(possibly accompanied by electrons)**

Dark matter interpretation

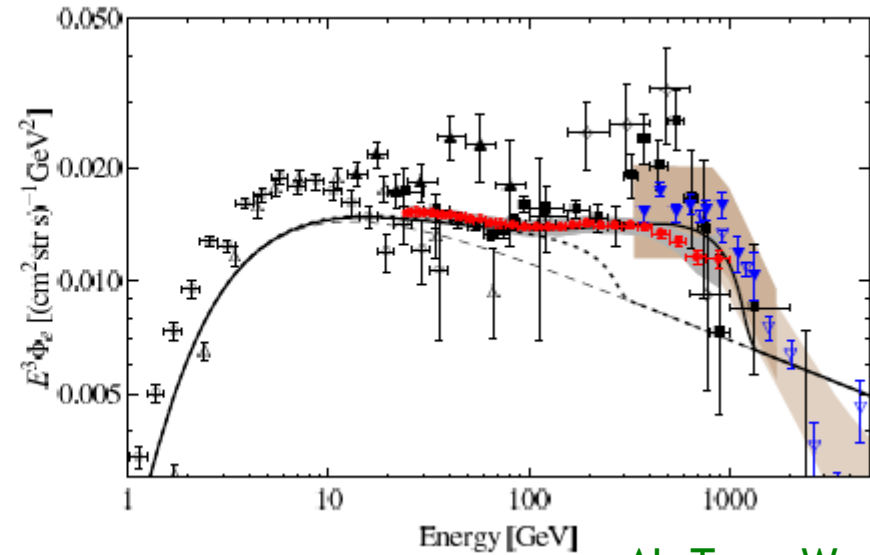
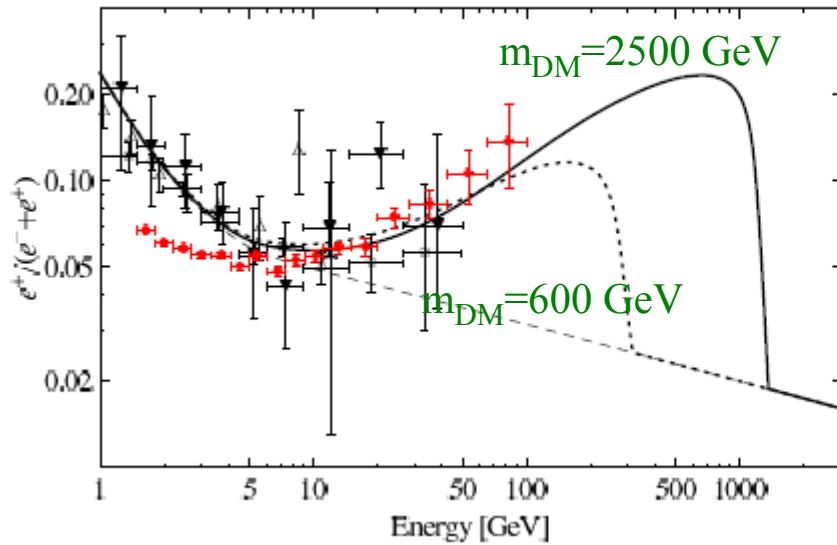
An electron/positron excess could arise from dark matter annihilations ...



Cholis et al.
arXiv:0811.3641

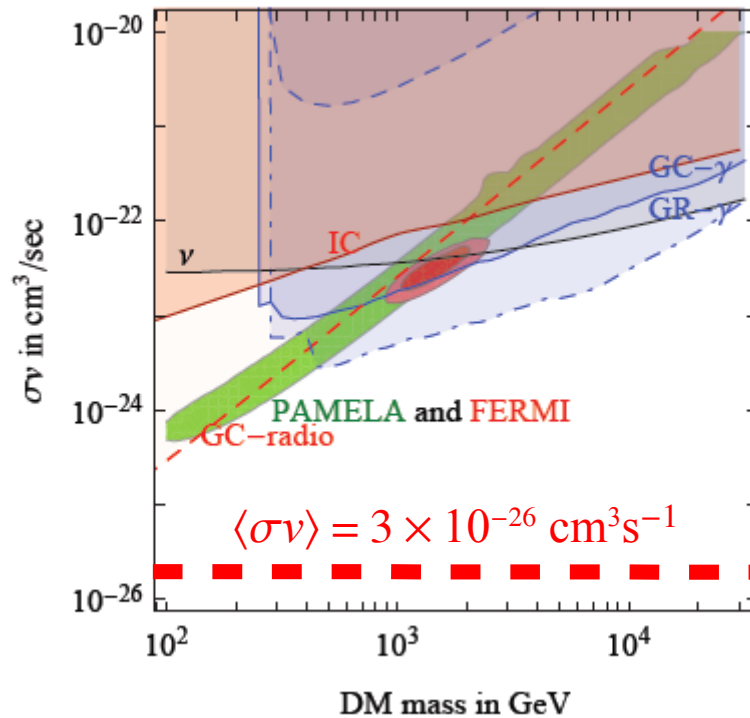
... or dark matter decays

“Democratic” decay $\psi \rightarrow l^+ l^- \nu$

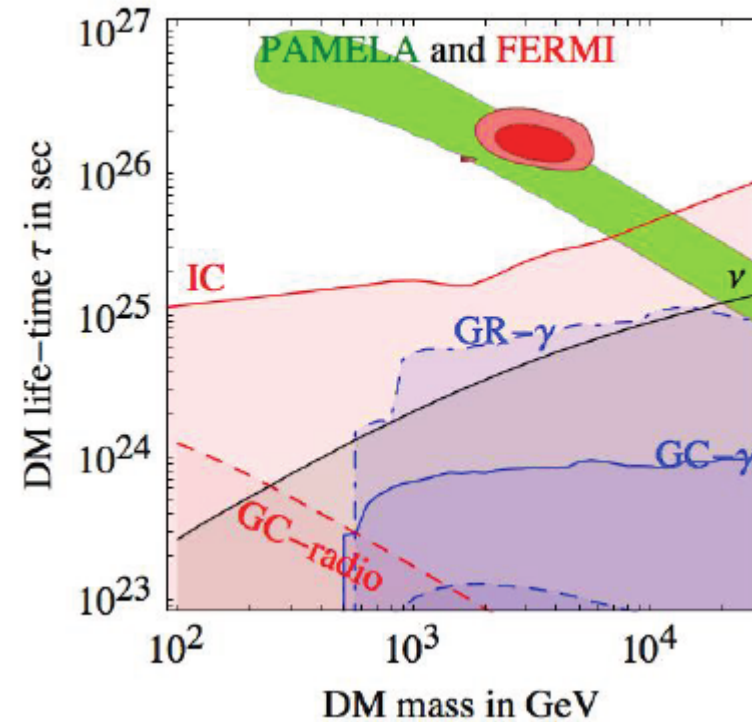


Al, Tran, Weniger
arXiv:0906.1571

DM DM $\rightarrow \mu^+ \mu^-$, Einasto profile



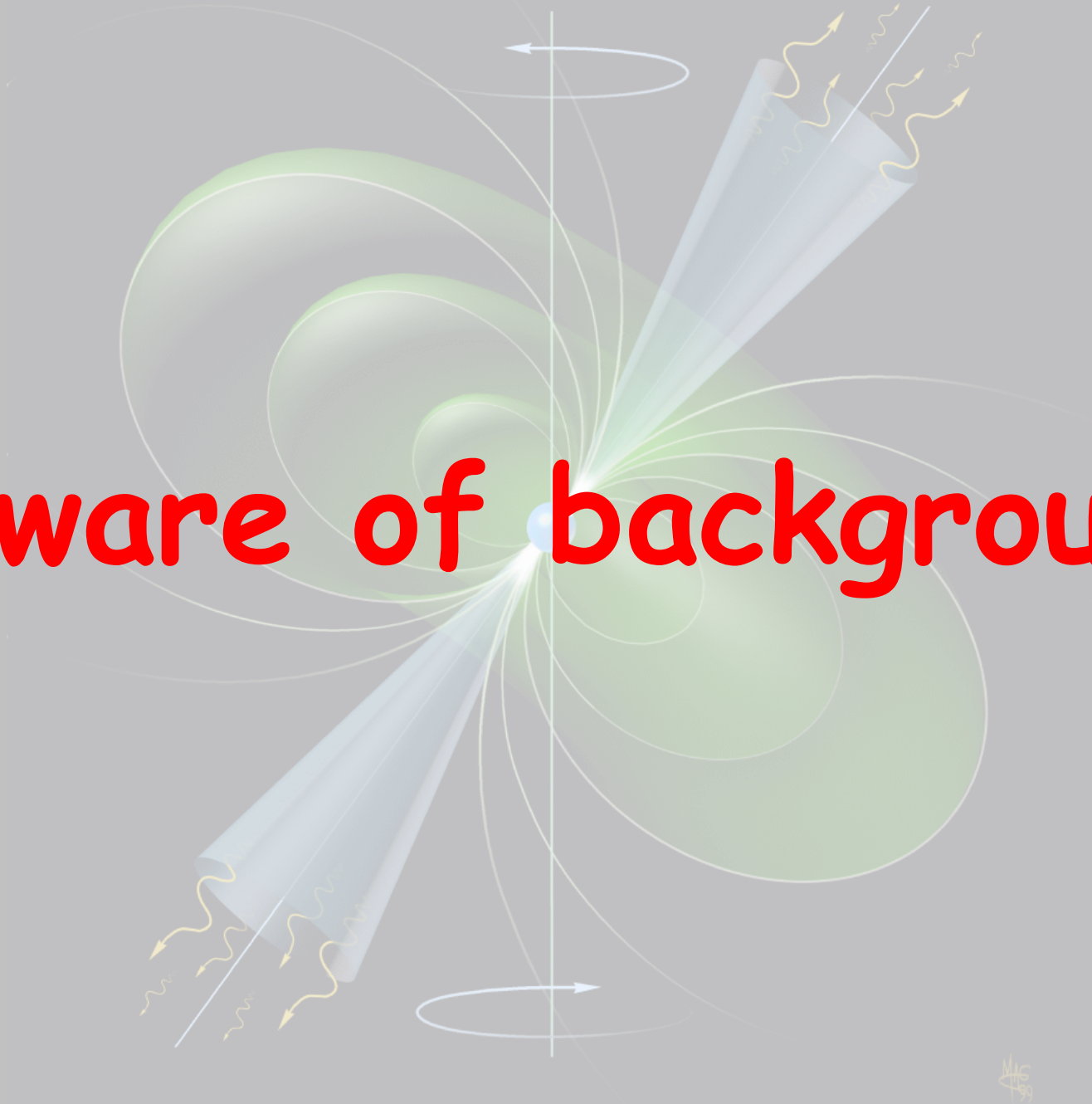
DM $\rightarrow \mu^+ \mu^-$, Einasto profile



Is this the first non-gravitational evidence of dark matter?

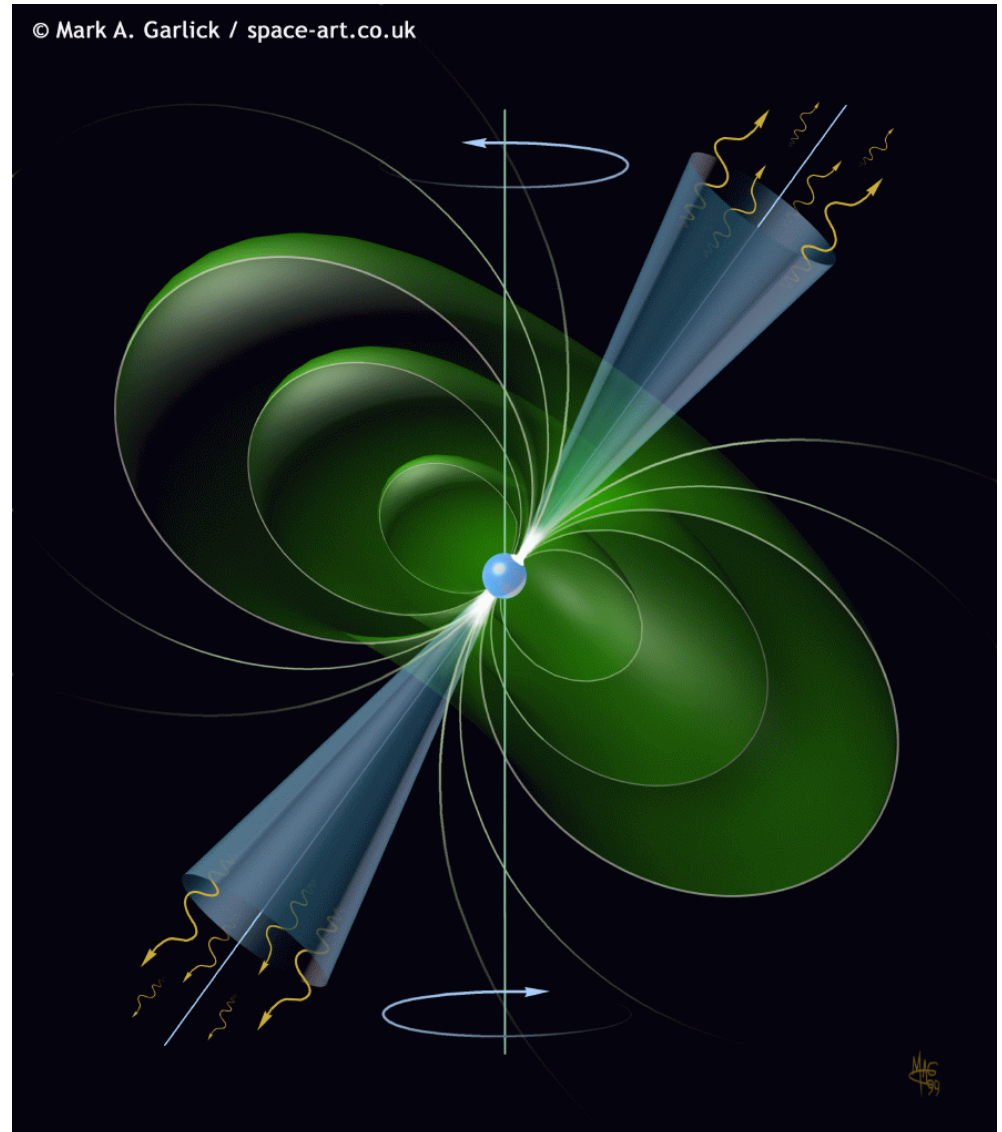
“Extraordinary claims require extraordinary evidence”
Carl Sagan

Beware of backgrounds!

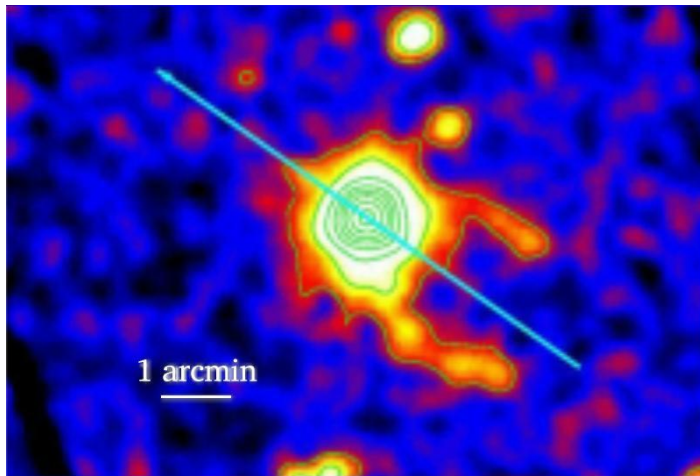


**Pulsars are sources
of high energy
electrons & positrons**

Atoyan, Aharonian, Völk '95
Chi, Cheng, Young '95
Grimani '04

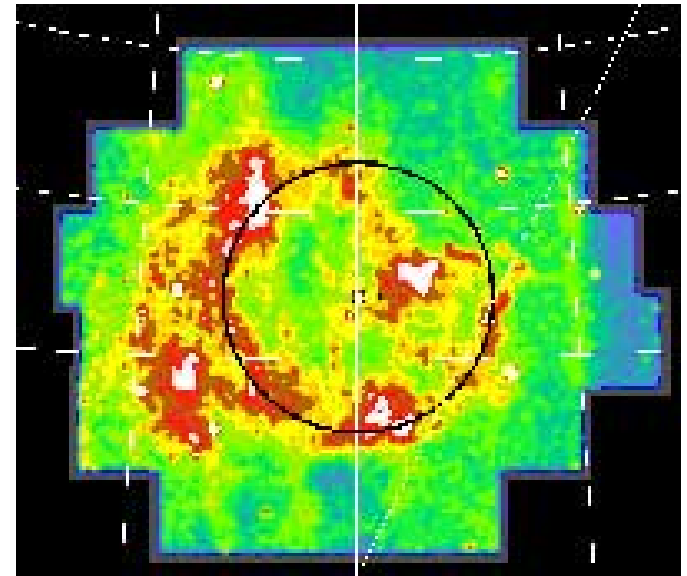


Pulsar explanation I: Geminga + Monogem



Geminga

T=370 000 years
D=157 pc

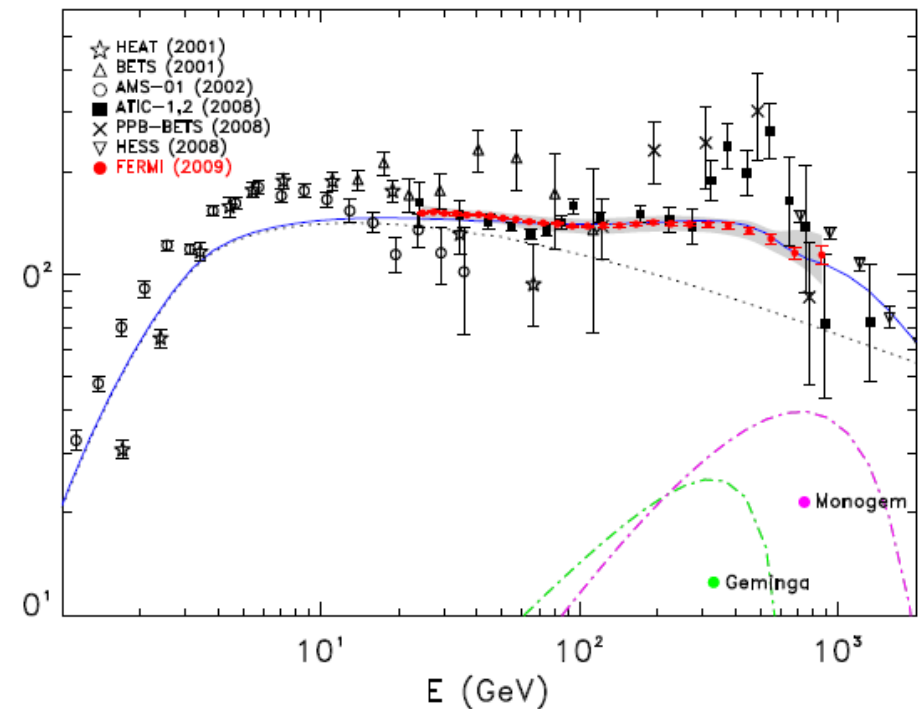
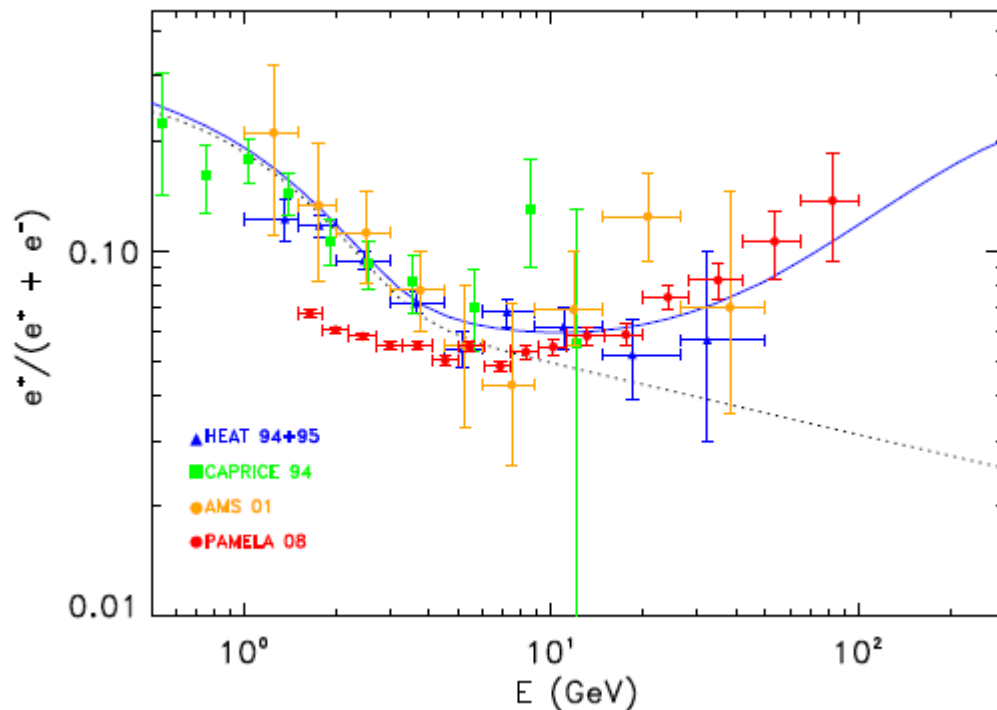


Monogem (B0656+14)

T=110 000 years
D=290 pc

Pulsar explanation I: Geminga + Monogem

Grasso et al.

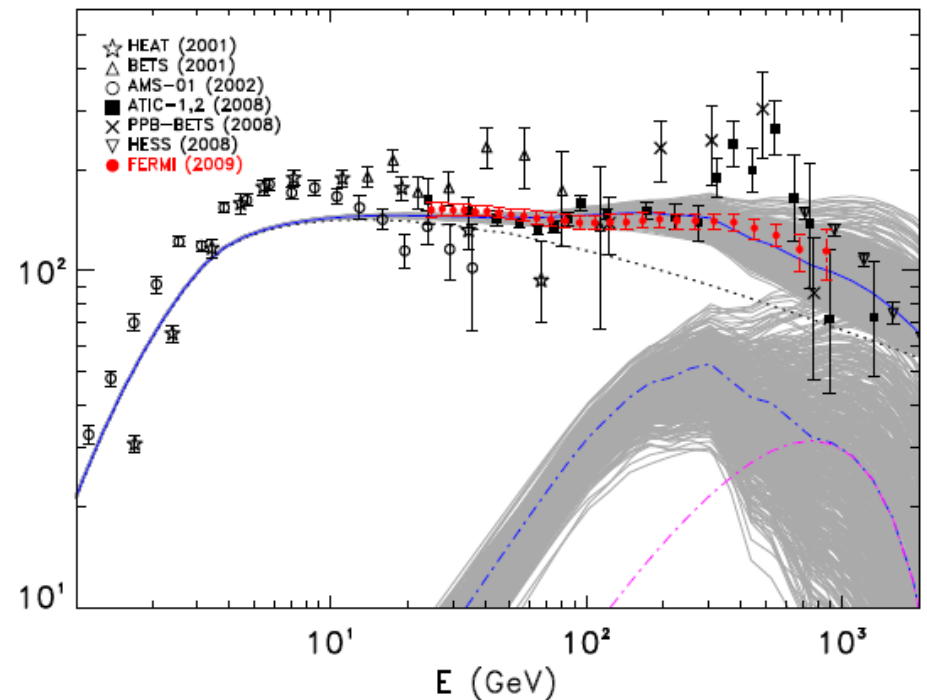
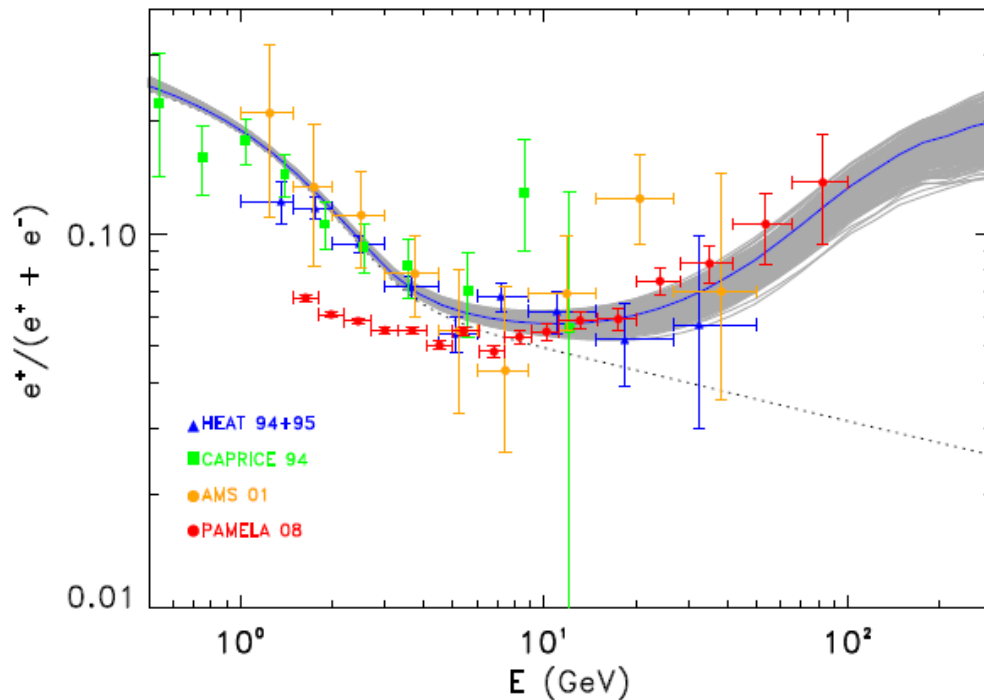


Nice agreement. However, it is not a prediction!

- $dN_e/dE_e \propto E_e^{-1.7} \exp(-E_e/1100 \text{ GeV})$
- Energy output in e^+e^- pairs: 40% of the spin-down rate

Pulsar explanation II: Multiple pulsars

Grasso et al.



- $dN_e/dE_e \propto E_e^{-\alpha} \exp(-E_e/E_0)$, $1.5 < \alpha < 1.9$, $800 \text{ GeV} < E_0 < 1400 \text{ GeV}$
- Energy output in e^+e^- pairs: between 10-30% of the spin-down rate

**The origin of the positron excess
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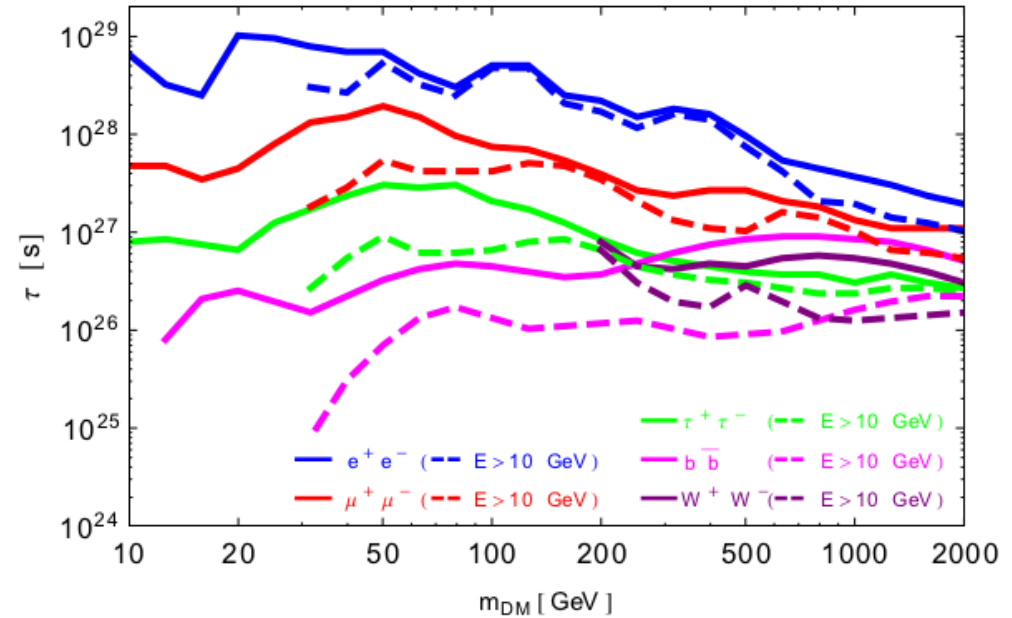
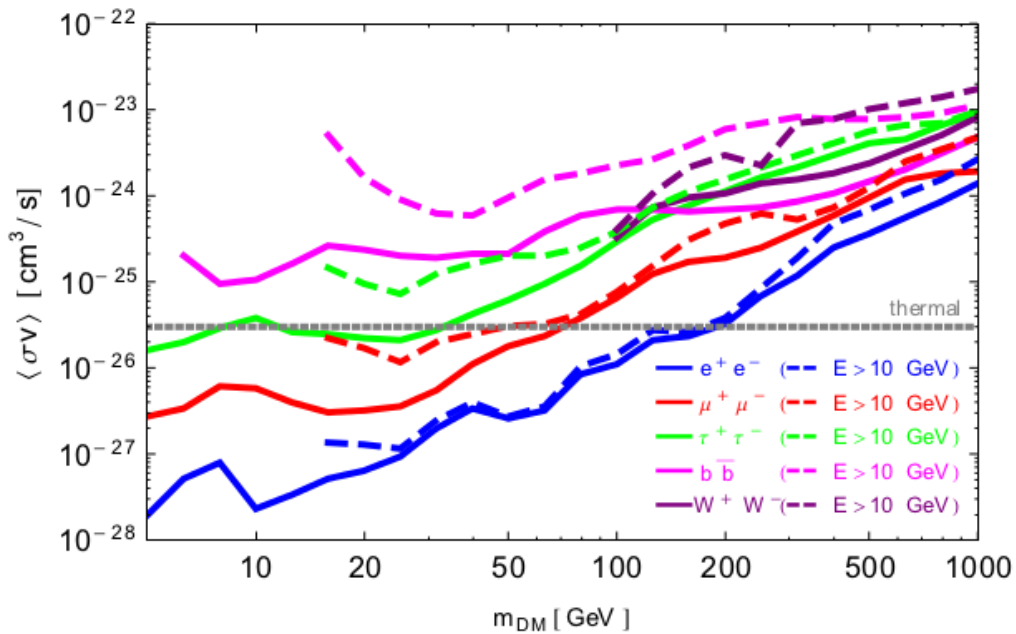
- Dark matter? Probably not.
- Pulsars? Perhaps yes.
- Something else? Perhaps yes.

The origin of the positron excess is still unclear:

- Dark matter? Probably not.
- Pulsars? Perhaps yes.
- Something else? Perhaps yes.
- Regardless of the origin of the positron excess, the positron data can be used to set limits on the dark matter parameters.

Latest limits from the positron fraction:

- Use AMS-02 data
- Make a fit of a model with secondary positrons + source + dark matter

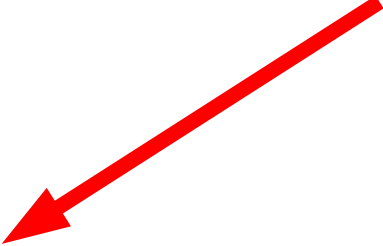


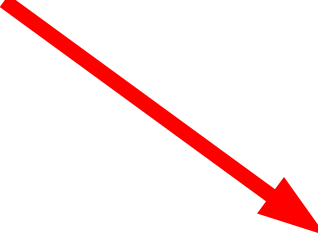
AI, Lamperstorfer, Silk '13
See also Bergström et al. '13

Gamma-rays

Production of gamma-rays

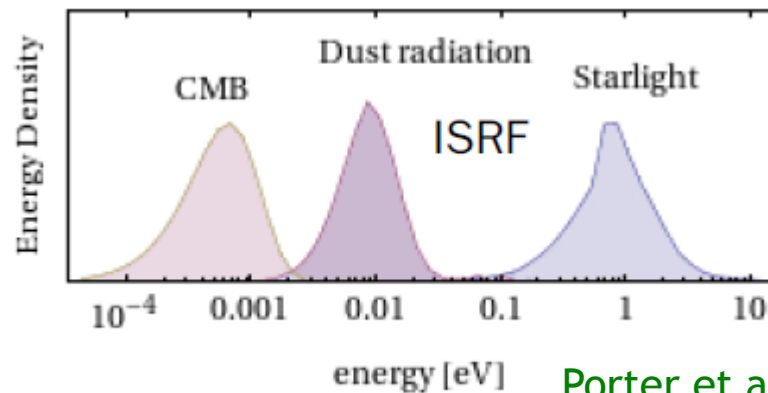
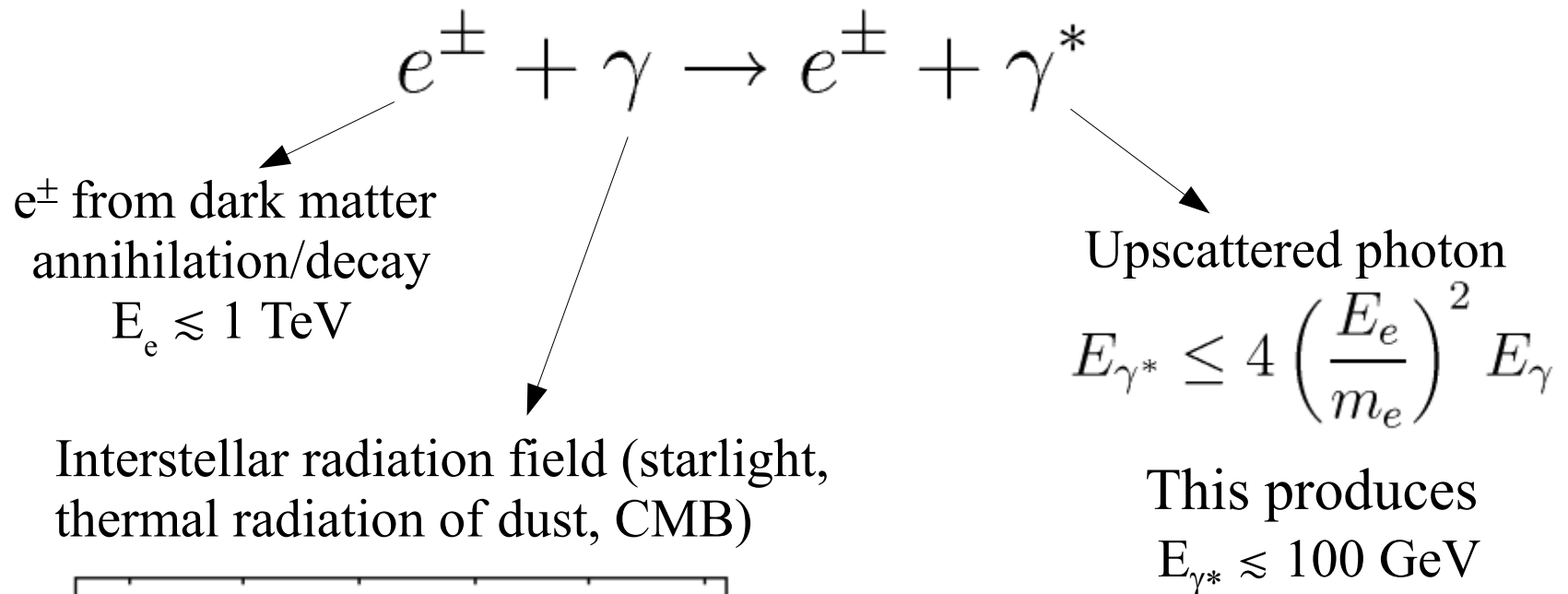
The gamma ray flux from dark matter annihilations/decays has two components:

- 
- Inverse Compton Scattering radiation of electrons/positrons produced in the annihilation/decay.
 - Always smooth spectrum.

- 
- Prompt radiation of gamma rays produced in the annihilation/decay (final state radiation, pion decay...)
 - May contain spectral features.

Inverse Compton Scattering radiation

The inverse Compton scattering of electrons/positrons from dark matter annihilation/decay with the interstellar and extragalactic radiation fields produces gamma rays.



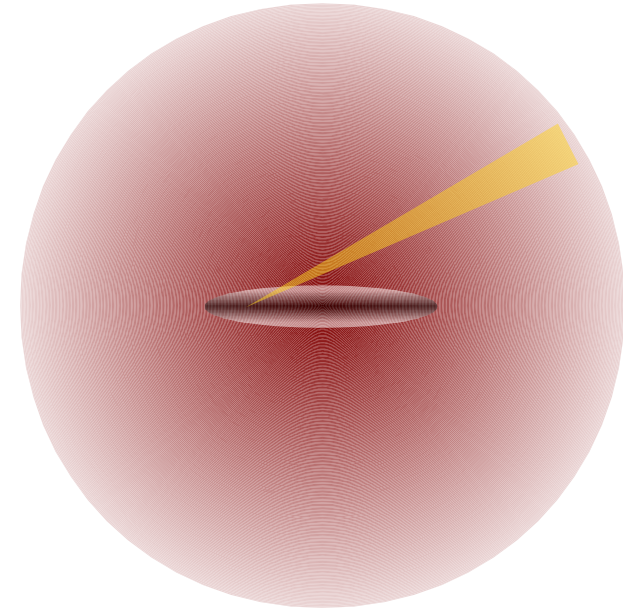
Prompt radiation

Annihilation

$$\frac{dJ}{dE_\gamma} = \frac{1}{4\pi} \underbrace{\left[\frac{\langle \sigma_{\text{ann}} v \rangle}{2m_{\text{DM}}^2} \sum_f \frac{dN_\gamma^f}{dE_\gamma} B_f \right]}_{\substack{\text{Source term} \\ \text{(particle physics)}}} \times \underbrace{\int_{\text{l.o.s.}} \rho^2(\vec{l}) d\vec{l}}_{\substack{\text{Line-of-sight integral} \\ \text{(astrophysics)}}$$

Decay

$$\frac{dJ}{dE_\gamma} = \frac{1}{4\pi} \underbrace{\left[\frac{1}{\tau_{\text{DM}} m_{\text{DM}}} \sum_f \frac{dN_\gamma^f}{dE_\gamma} B_f \right]}_{\substack{\text{Source term} \\ \text{(particle physics)}}} \times \underbrace{\int_{\text{l.o.s.}} \rho(\vec{l}) d\vec{l}}_{\substack{\text{Line-of-sight integral} \\ \text{(astrophysics)}}$$



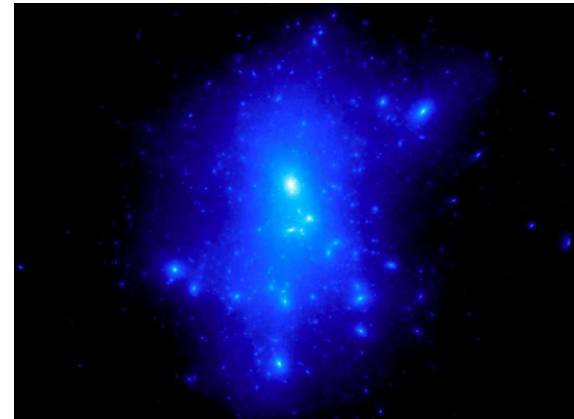
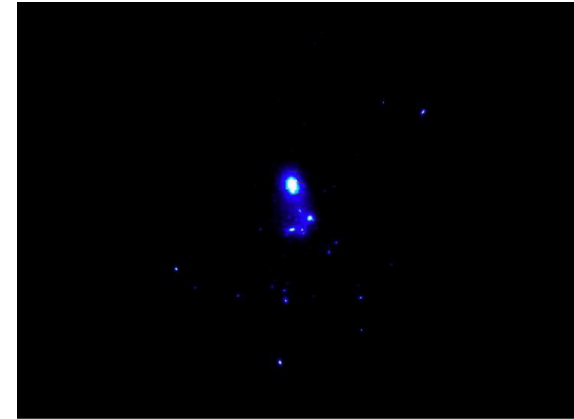
Prompt radiation

Annihilation

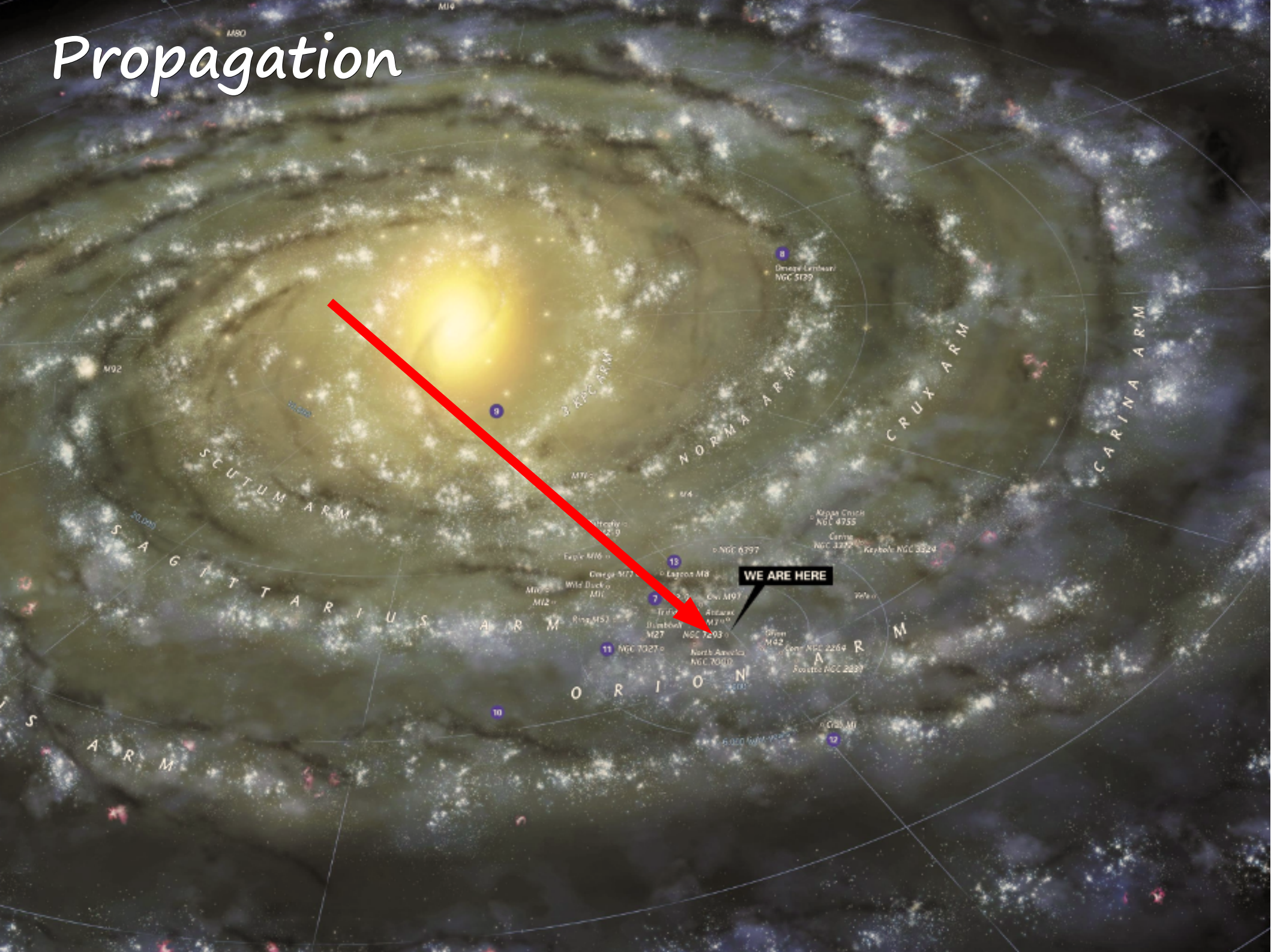
$$\frac{dJ}{dE_\gamma} = \frac{1}{4\pi} \underbrace{\left[\frac{\langle \sigma_{\text{ann}} v \rangle}{2m_{\text{DM}}^2} \sum_f \frac{dN_\gamma^f}{dE_\gamma} B_f \right]}_{\text{Source term (particle physics)}} \times \underbrace{\int_{\text{l.o.s.}} \rho^2(\vec{l}) d\vec{l}}_{\text{Line-of-sight integral (astrophysics)}}$$

Decay

$$\frac{dJ}{dE_\gamma} = \frac{1}{4\pi} \underbrace{\left[\frac{1}{\tau_{\text{DM}} m_{\text{DM}}} \sum_f \frac{dN_\gamma^f}{dE_\gamma} B_f \right]}_{\text{Source term (particle physics)}} \times \underbrace{\int_{\text{l.o.s.}} \rho(\vec{l}) d\vec{l}}_{\text{Line-of-sight integral (astrophysics)}}$$

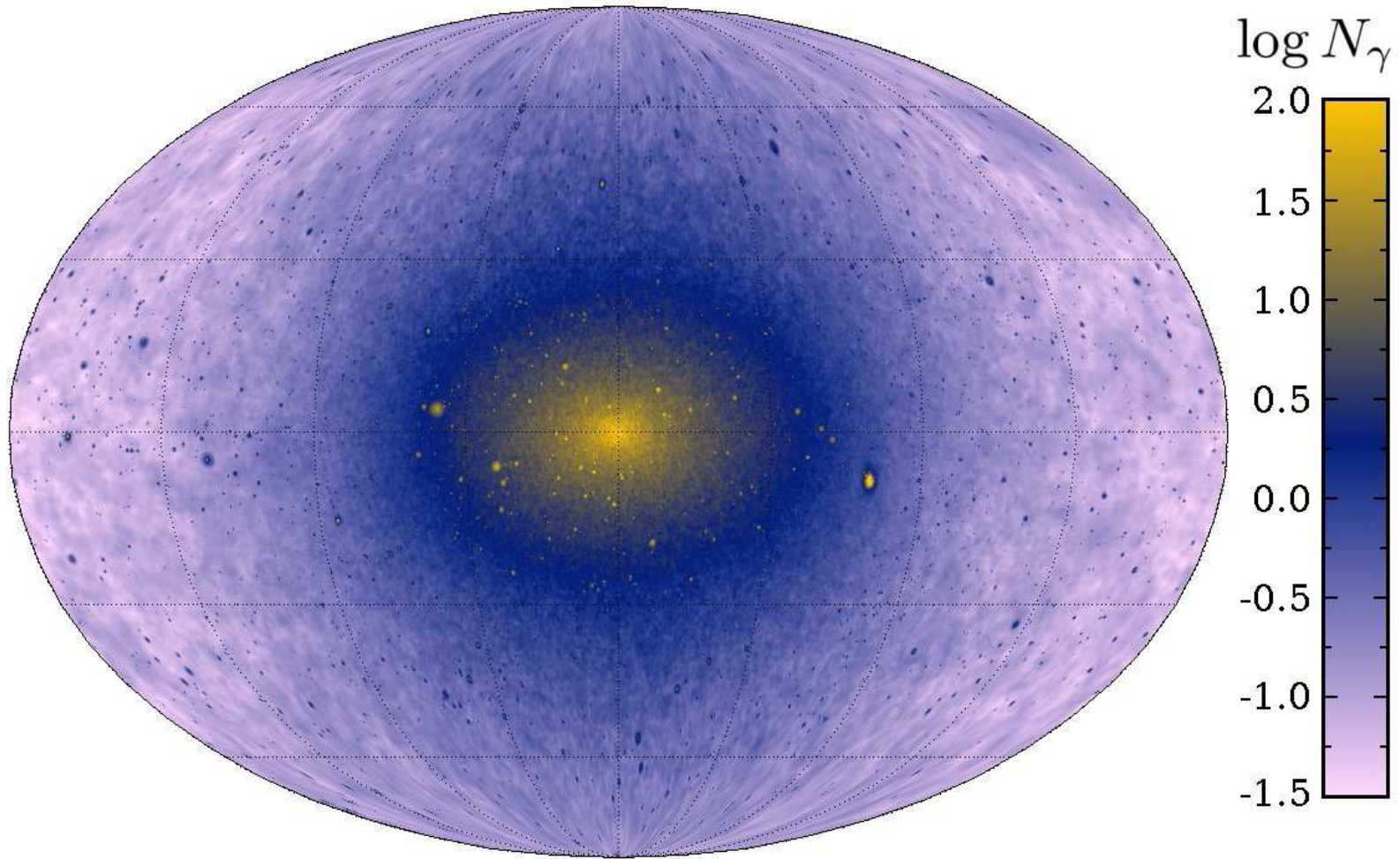


Propagation



Where to look for *annihilating* dark matter

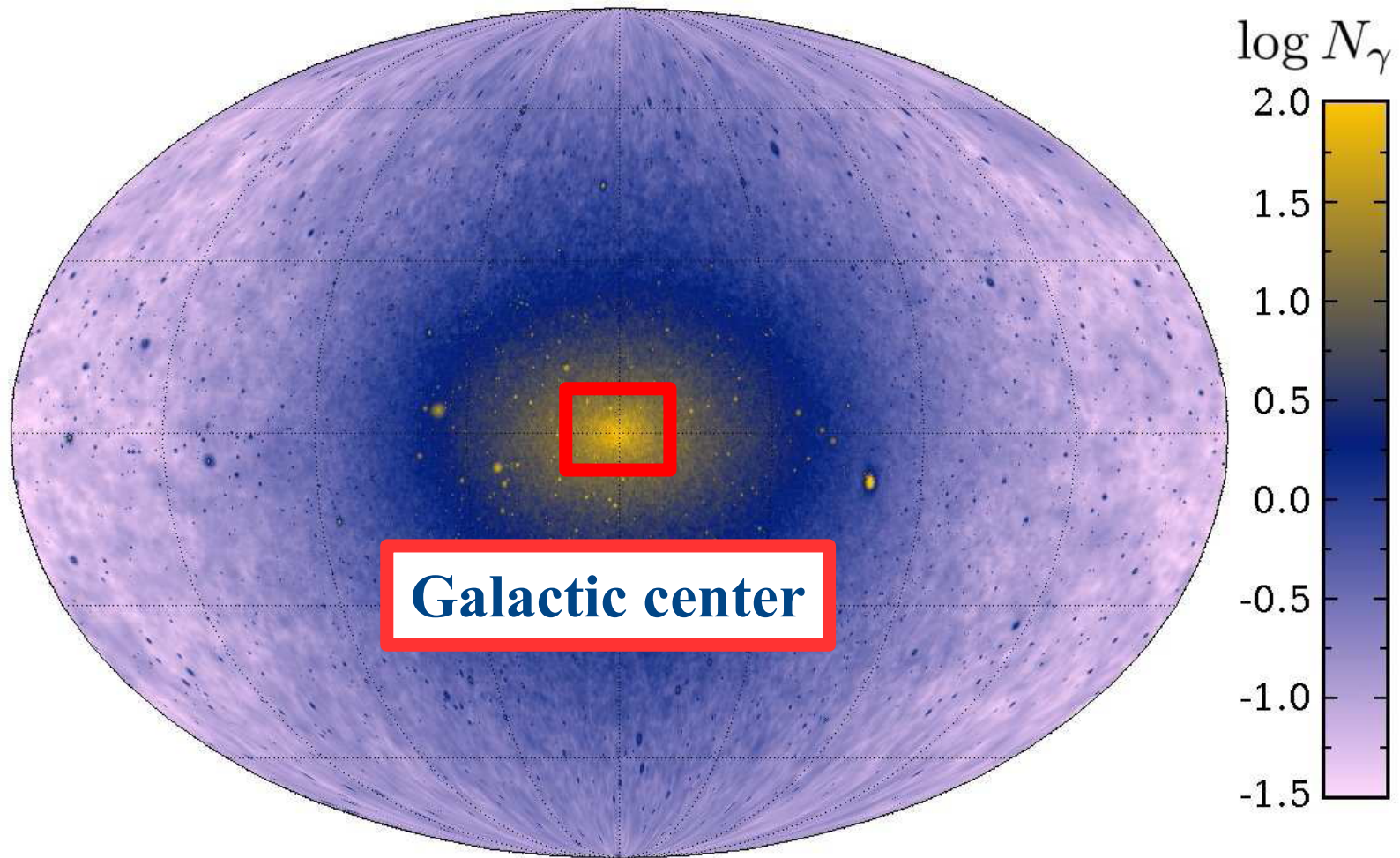
Baltz et al.
arXiv:0806.2911



Kuhlen, Diemand, Madau

Where to look for *annihilating* dark matter

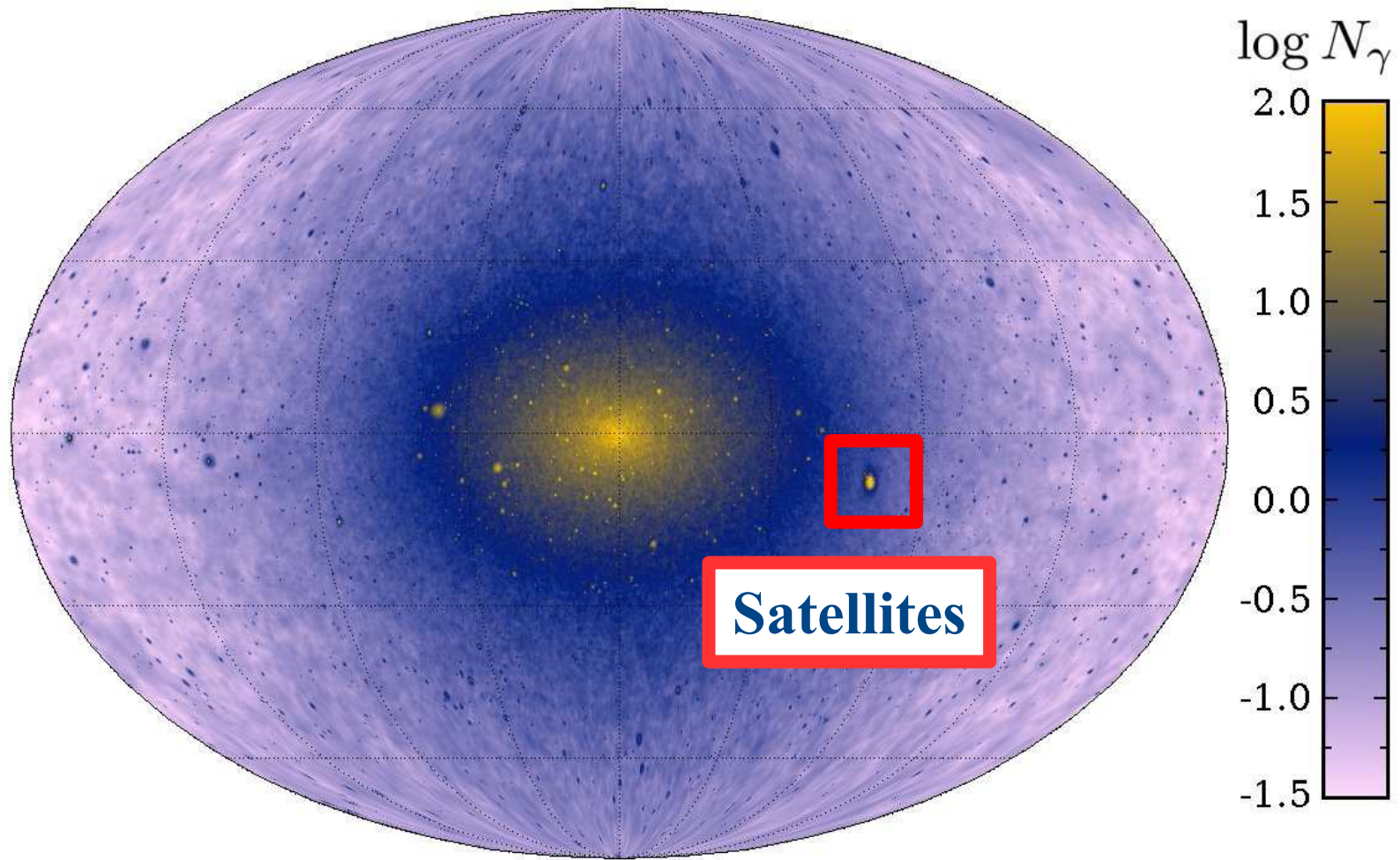
Baltz et al.
arXiv:0806.2911



Kuhlen, Diemand, Madau

Where to look for *annihilating* dark matter

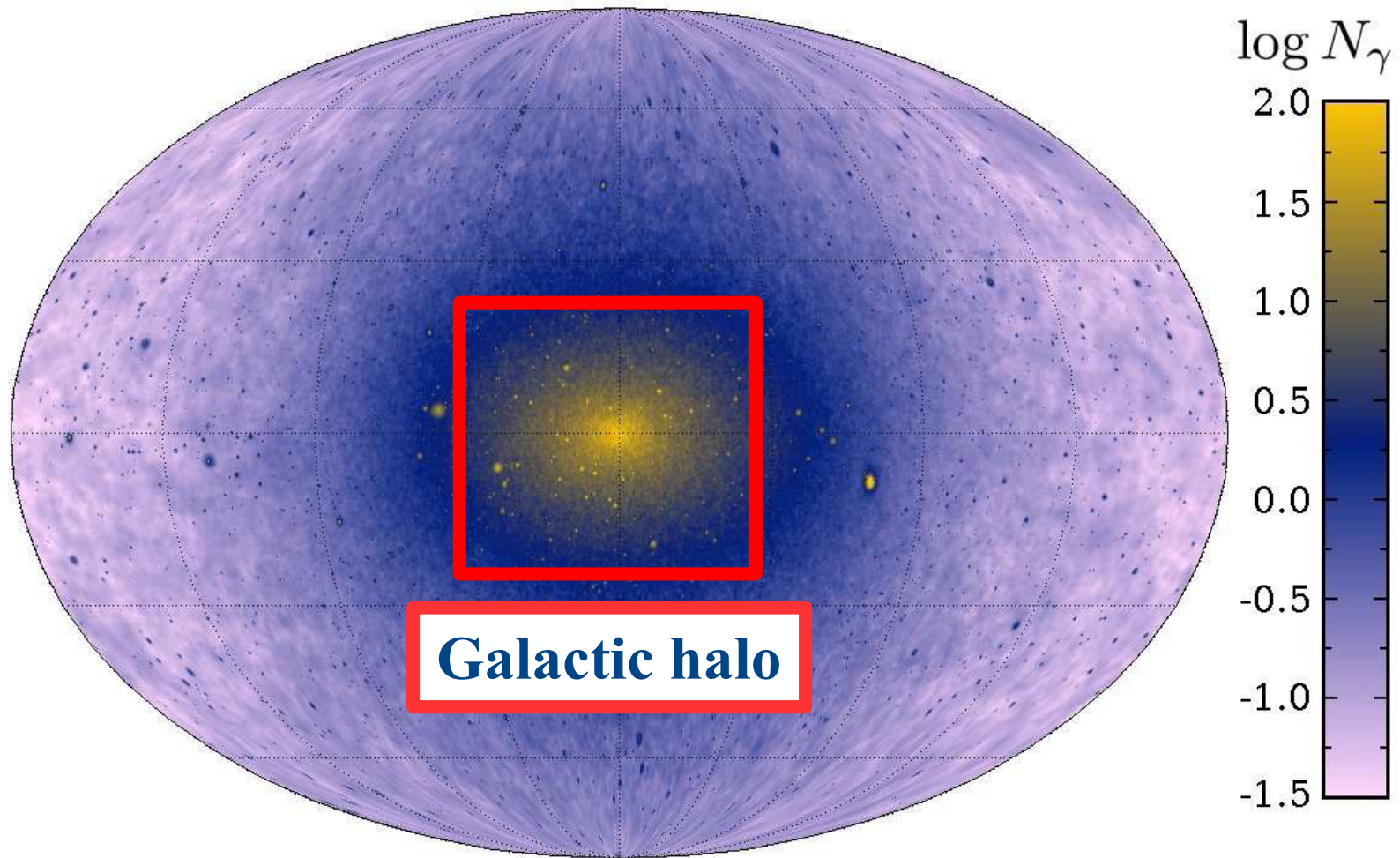
Baltz et al.
arXiv:0806.2911



Kuhlen, Diemand, Madau

Where to look for *annihilating* dark matter

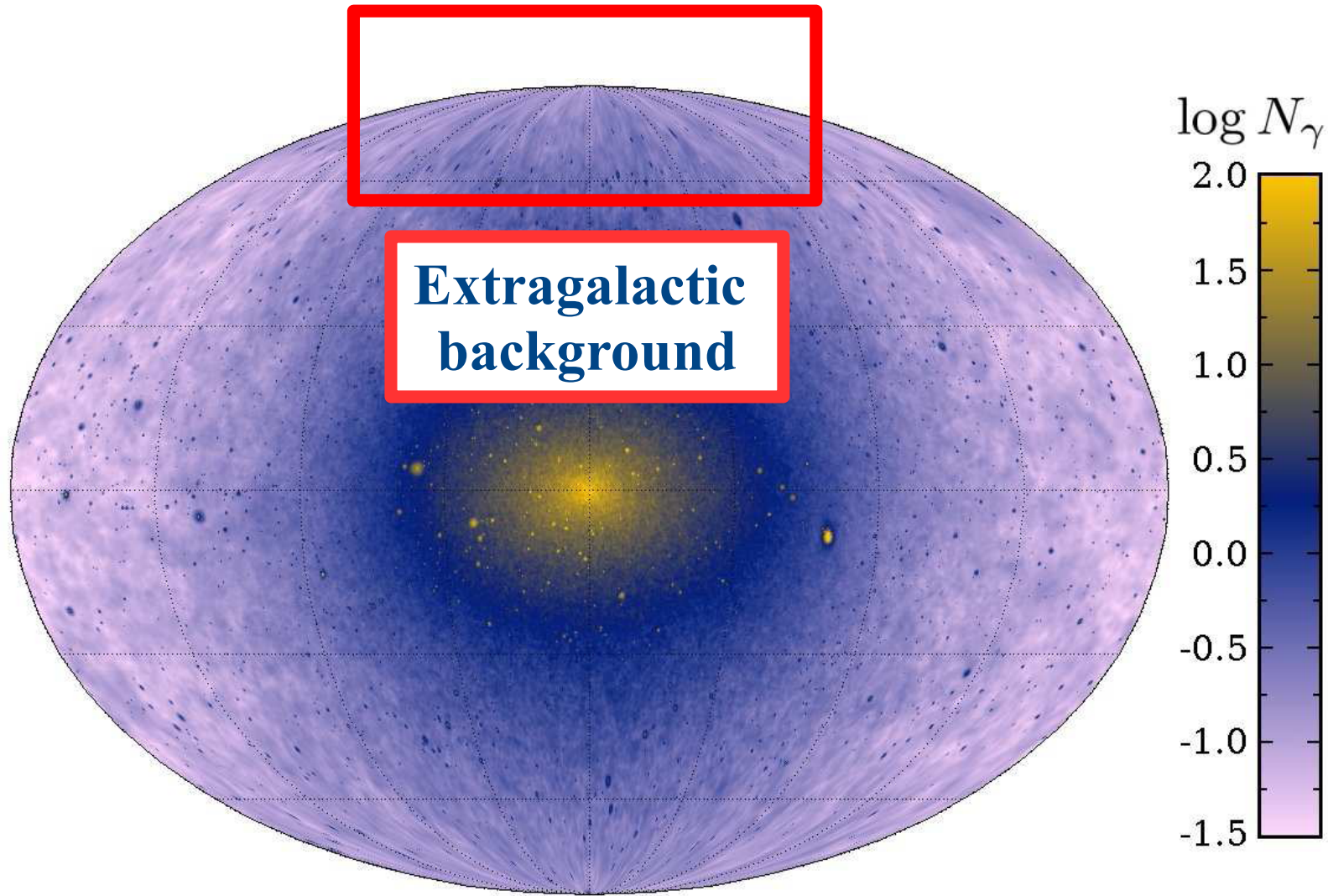
Baltz et al.
arXiv:0806.2911



Kuhlen, Diemand, Madau

Where to look for *annihilating* dark matter

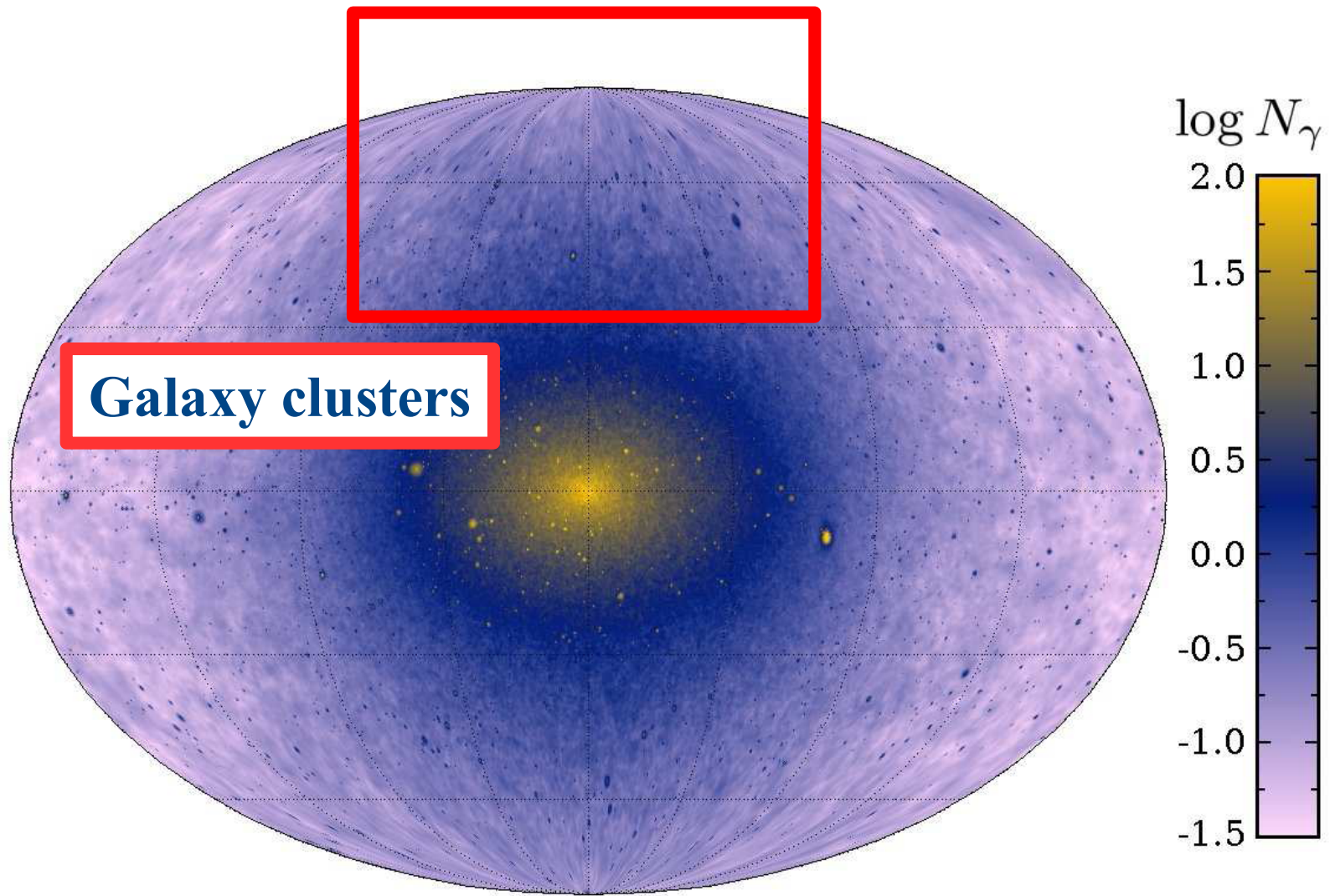
Baltz et al.
arXiv:0806.2911



Kuhlen, Diemand, Madau

Where to look for *annihilating* dark matter

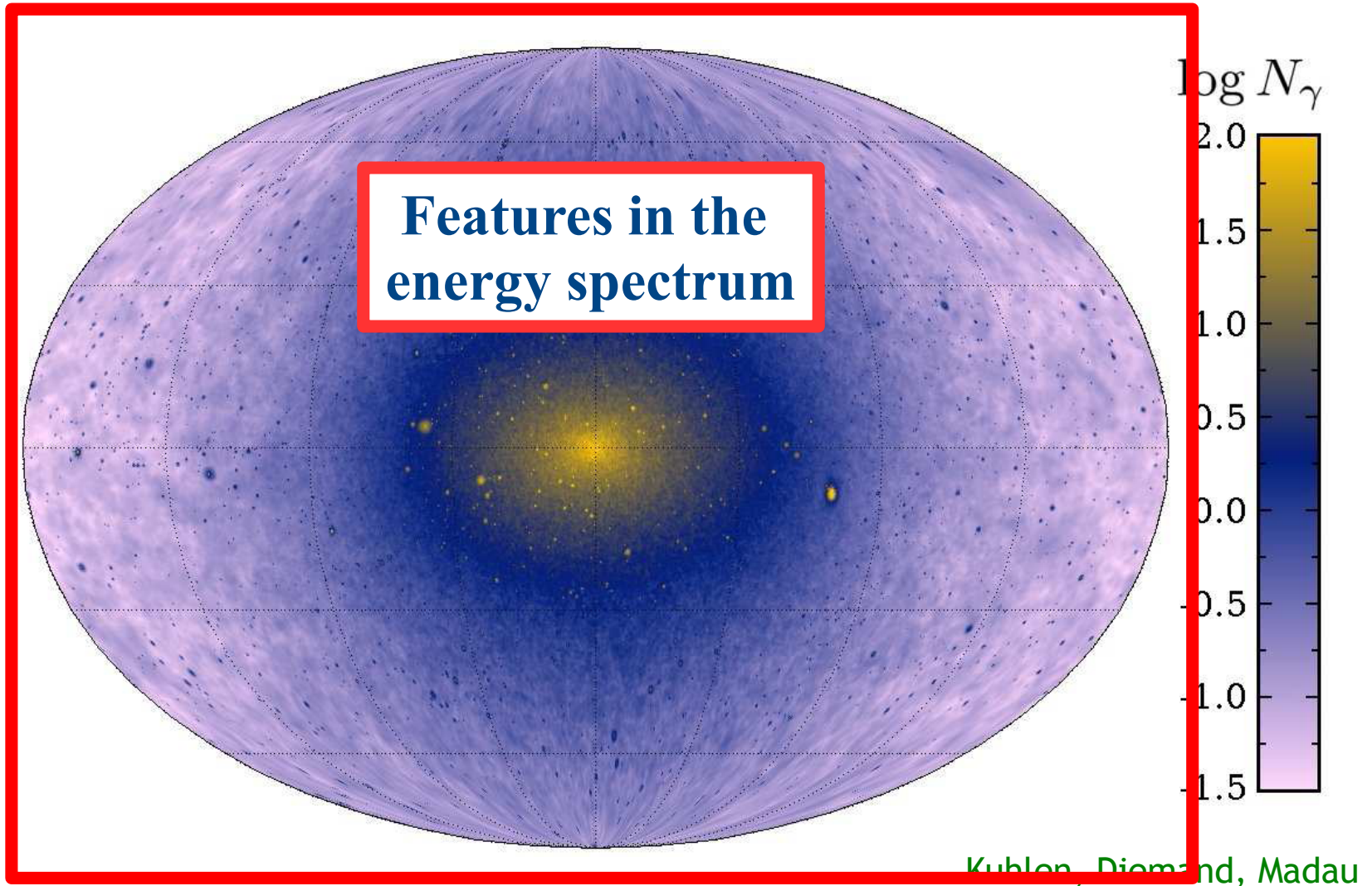
Baltz et al.
arXiv:0806.2911



Kuhlen, Diemand, Madau

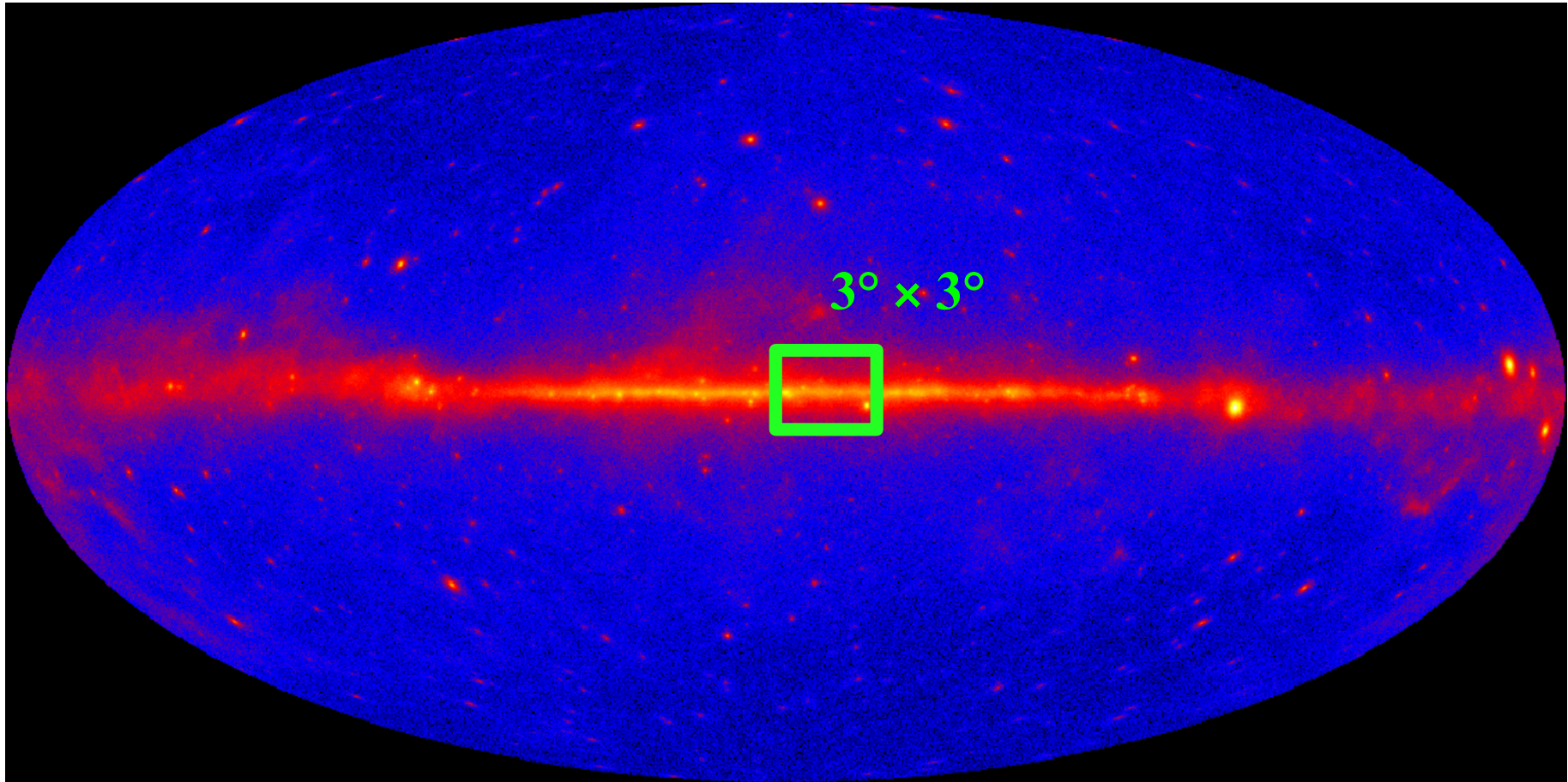
Where to look for *annihilating* dark matter

Baltz et al.
arXiv:0806.2911



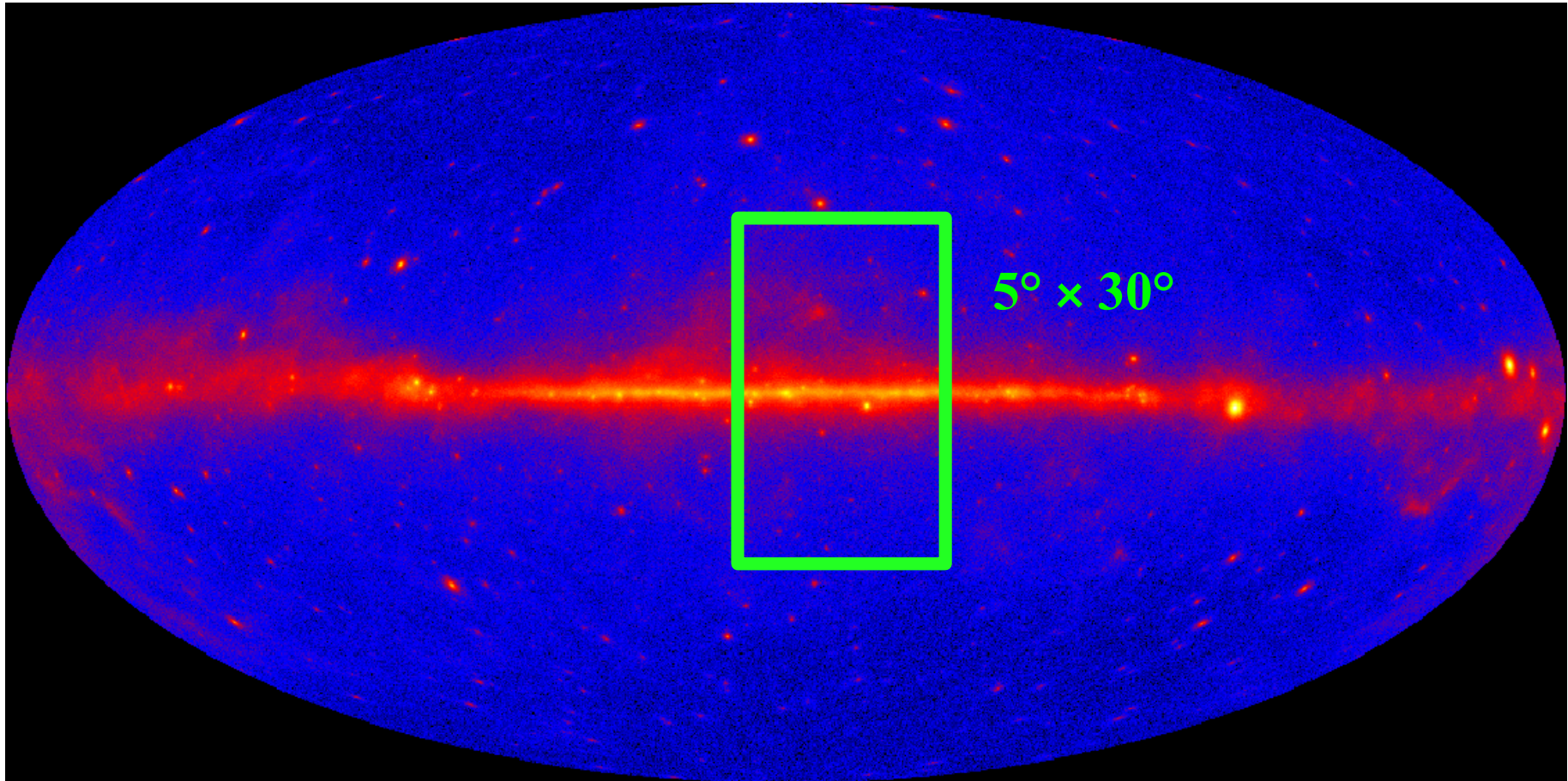
Diffuse Galactic emission

Divide the sky in different regions:



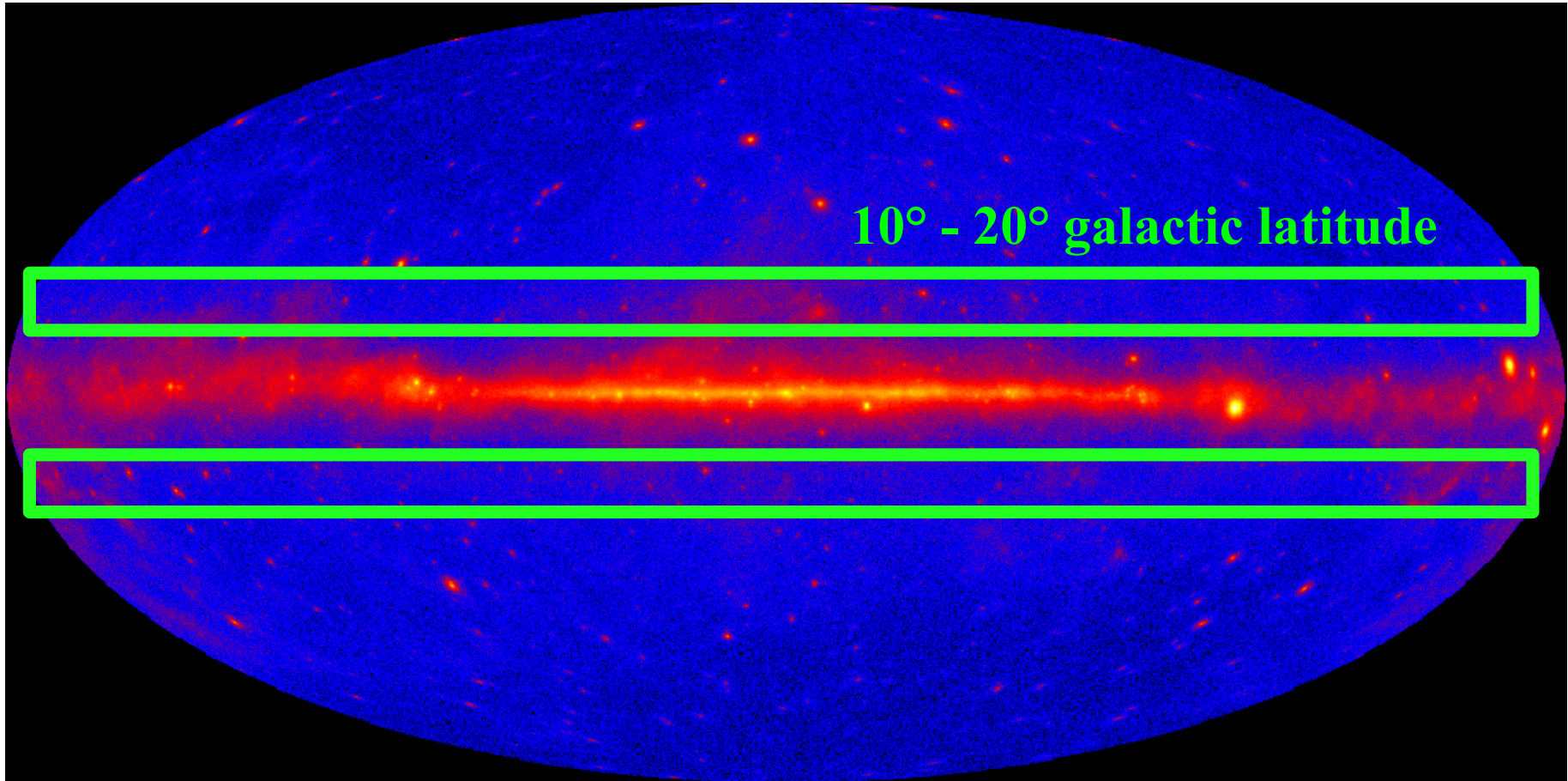
Diffuse Galactic emission

Divide the sky in different regions:



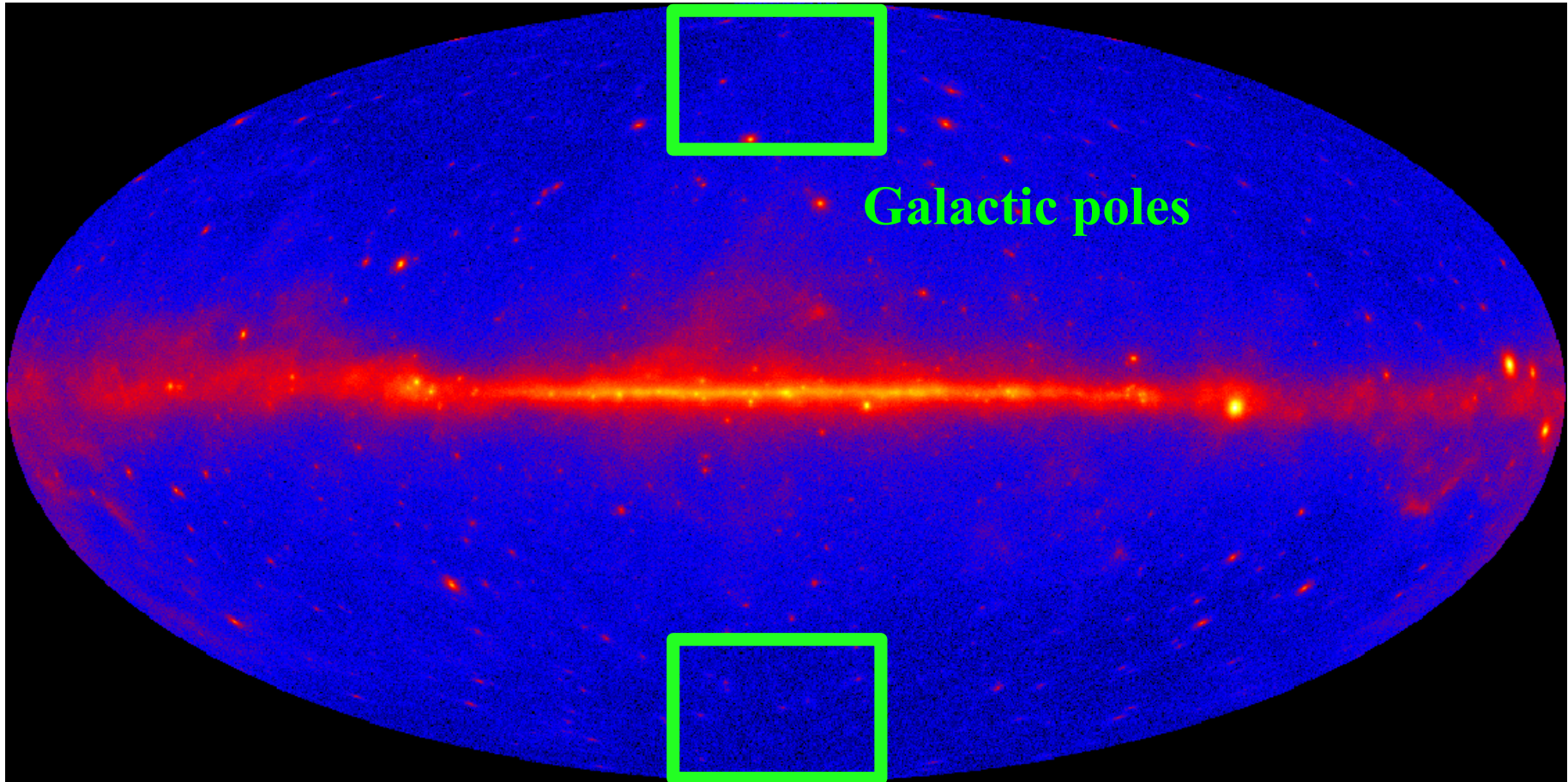
Diffuse Galactic emission

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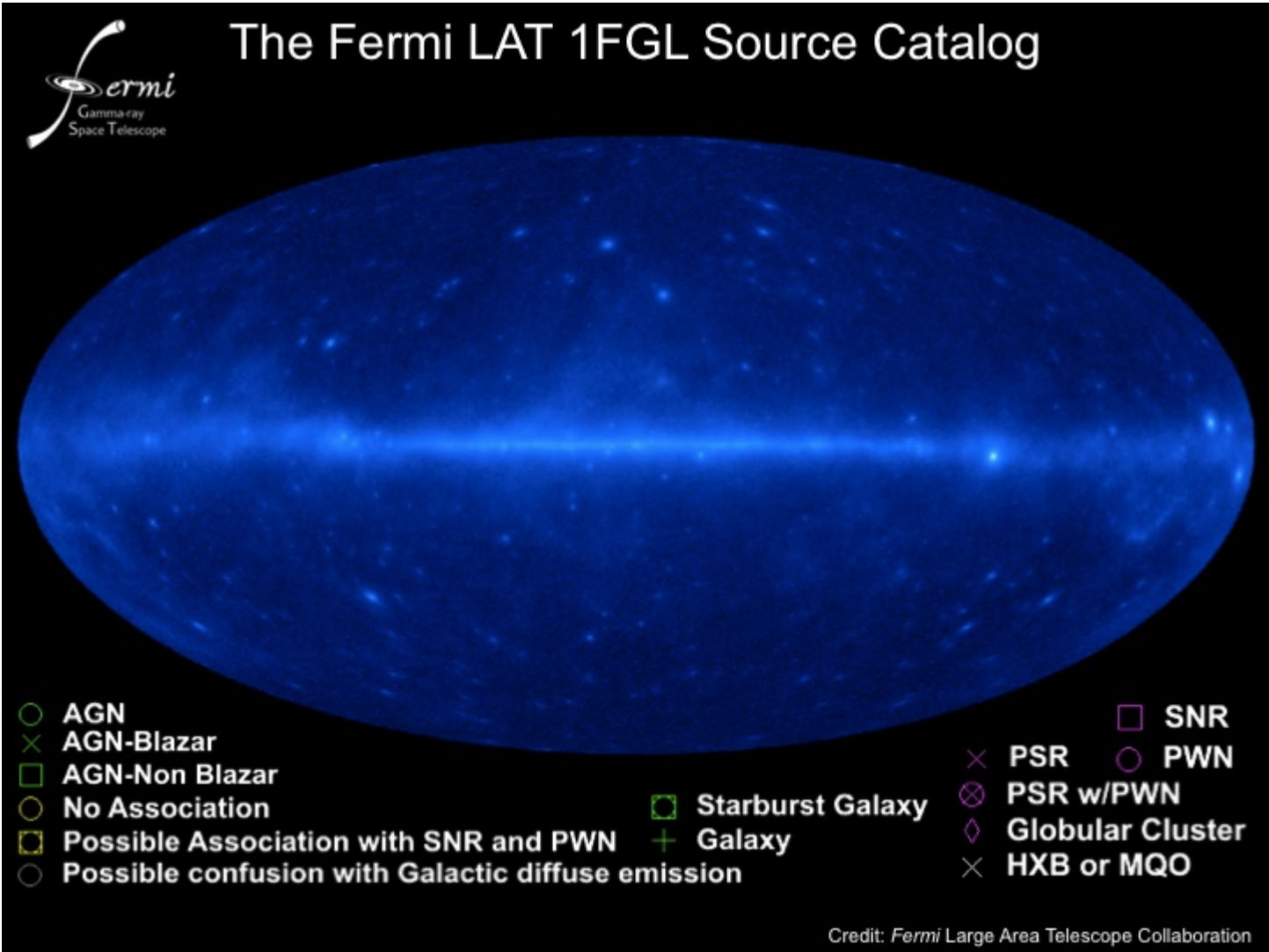
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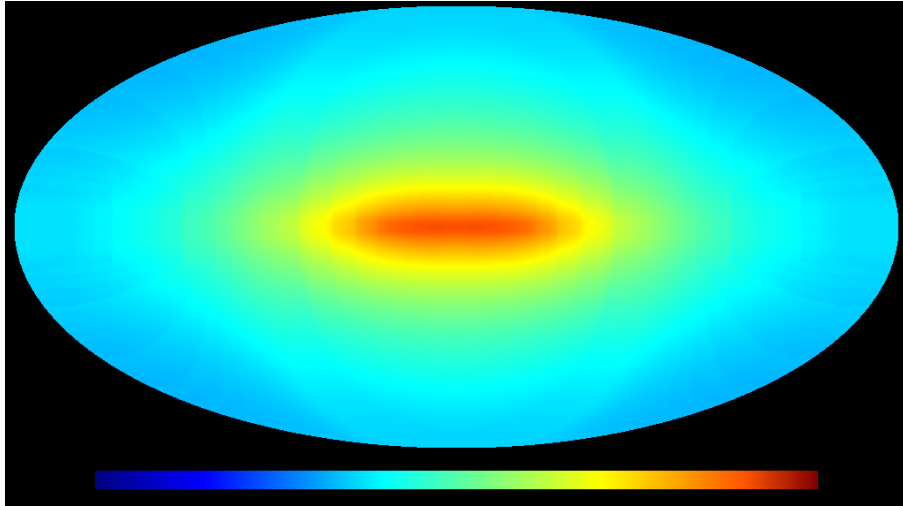


But beware of backgrounds when searching for dark matter...

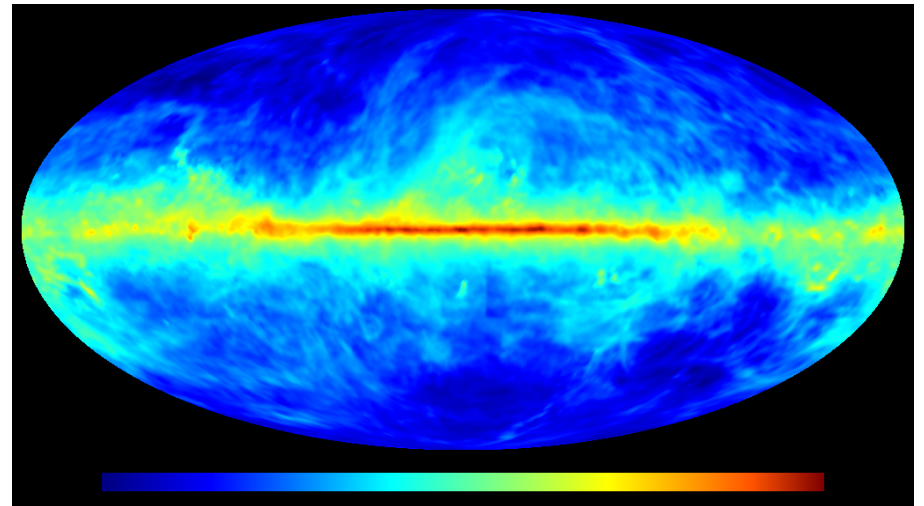
Background I: sources



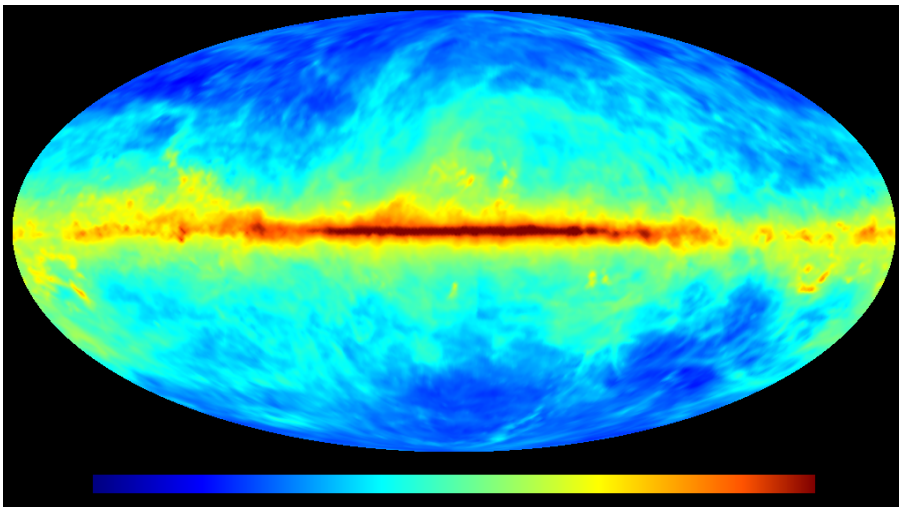
Background II: modelling of the diffuse emission



Inverse compton

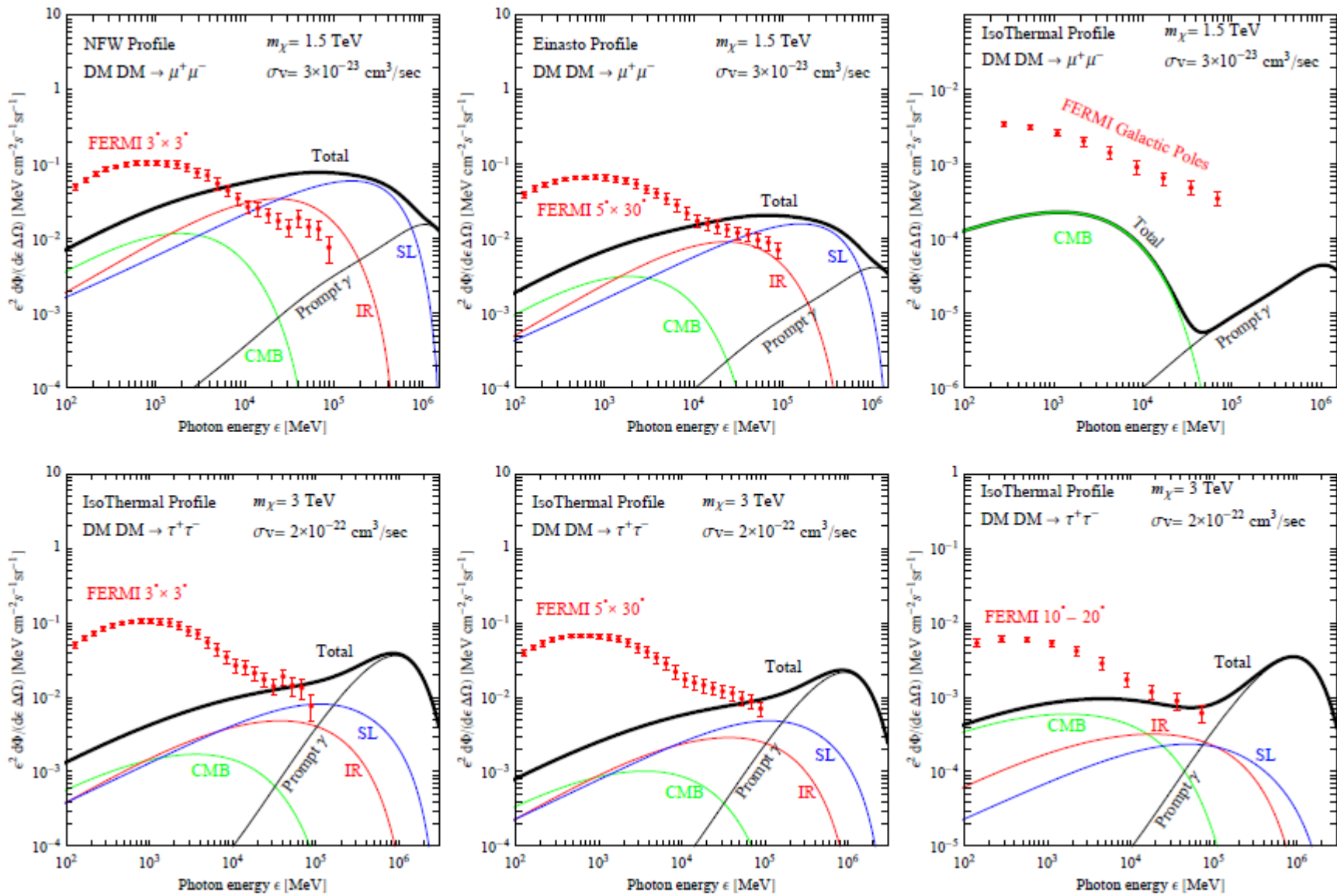


Bremmstrahlung

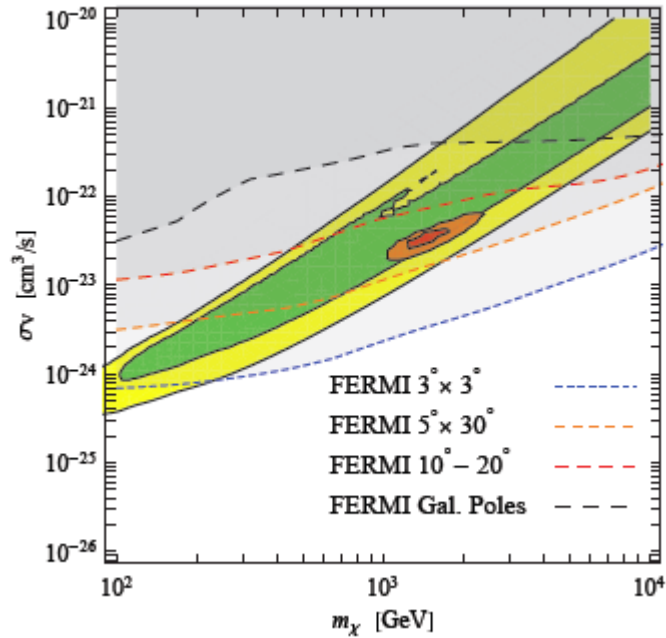


π^0 -decay

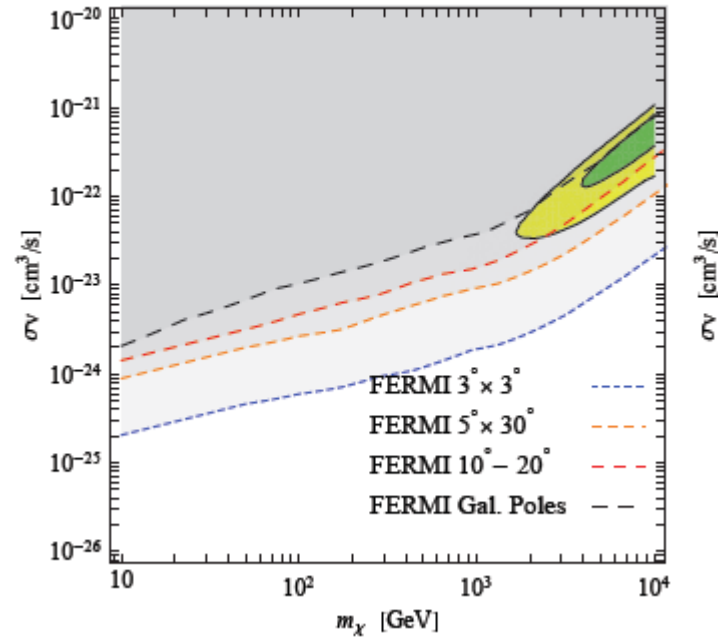
Conservative approach: demand that the flux from dark matter annihilation does not exceed the measured flux



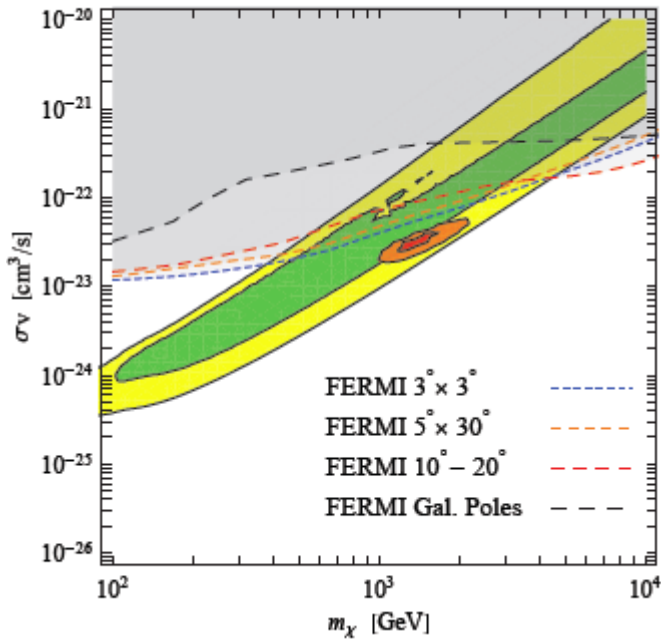
DM DM $\rightarrow \mu\mu$, Einasto profile



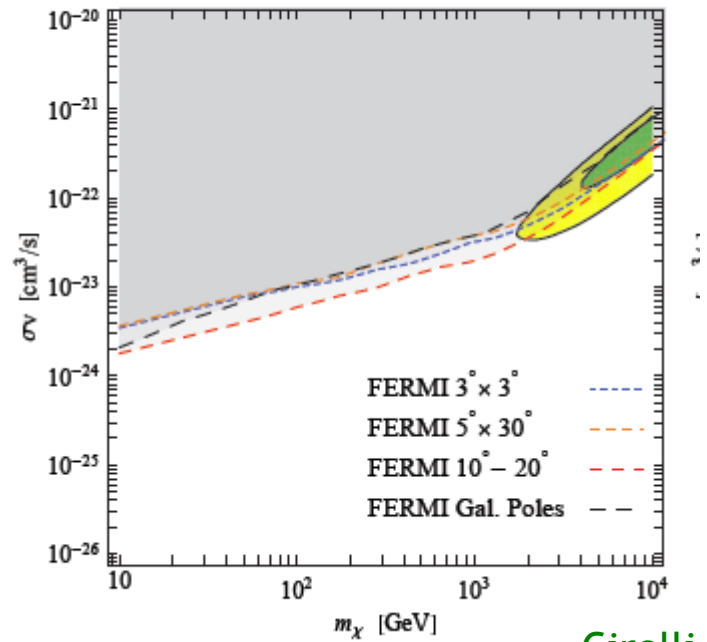
DM DM $\rightarrow bb$, Einasto profile



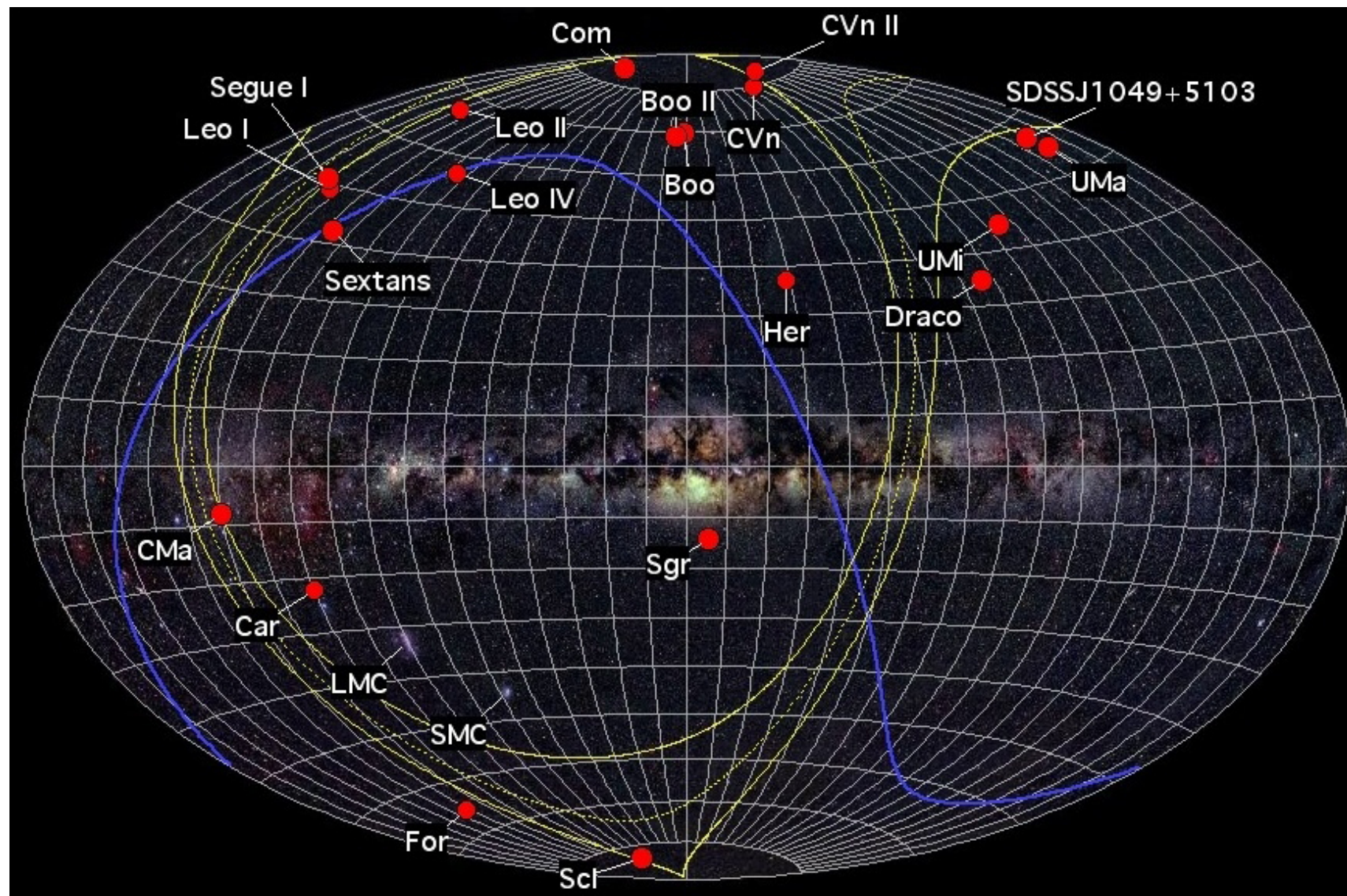
DM DM $\rightarrow \mu\mu$, Iso profile



DM DM $\rightarrow bb$, Iso profile



Dwarf spheroidal galaxies

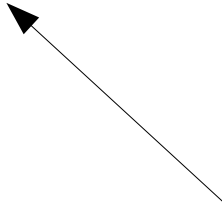


Name	Distance (kpc)	year of discovery	$M_{1/2}/L_{1/2}$ ref. 8	l	b	Ref.
Ursa Major II	30 ± 5	2006	4000^{+3700}_{-2100}	152.46	37.44	1,2
Segue 2	35	2009	650	149.4	-38.01	3
Willman 1	38 ± 7	2004	770^{+930}_{-440}	158.57	56.78	1
Coma Berenices	44 ± 4	2006	1100^{+800}_{-500}	241.9	83.6	1,2
Bootes II	46	2007	1800??	353.69	68.87	6,7
Bootes I	62 ± 3	2006	1700^{+1400}_{-700}	358.08	69.62	6
Ursa Minor	66 ± 3	1954	290^{+140}_{-90}	104.95	44.80	4,5
Sculptor	79 ± 4	1937	18^{+6}_{-5}	287.15	-83.16	4,5
Draco	76 ± 5	1954	200^{+80}_{-60}	86.37	34.72	4,5,9
Sextans	86 ± 4	1990	120^{+40}_{-35}	243.4	42.2	4,5
Ursa Major I	97 ± 4	2005	1800^{+1300}_{-700}	159.43	54.41	6
Hercules	132 ± 12	2006	1400^{+1200}_{-700}	28.73	36.87	6
Fornax	138 ± 8	1938	$8.7^{+2.8}_{-2.3}$	237.1	-65.7	4,5
Leo IV	160 ± 15	2006	260^{+1000}_{-200}	265.44	56.51	6

Relatively close




High mass-to-light ratio:
dwarf galaxies contain large
amounts of dark matter



Assume a Navarro-Frenk-White dark matter halo profile inside the tidal radius:

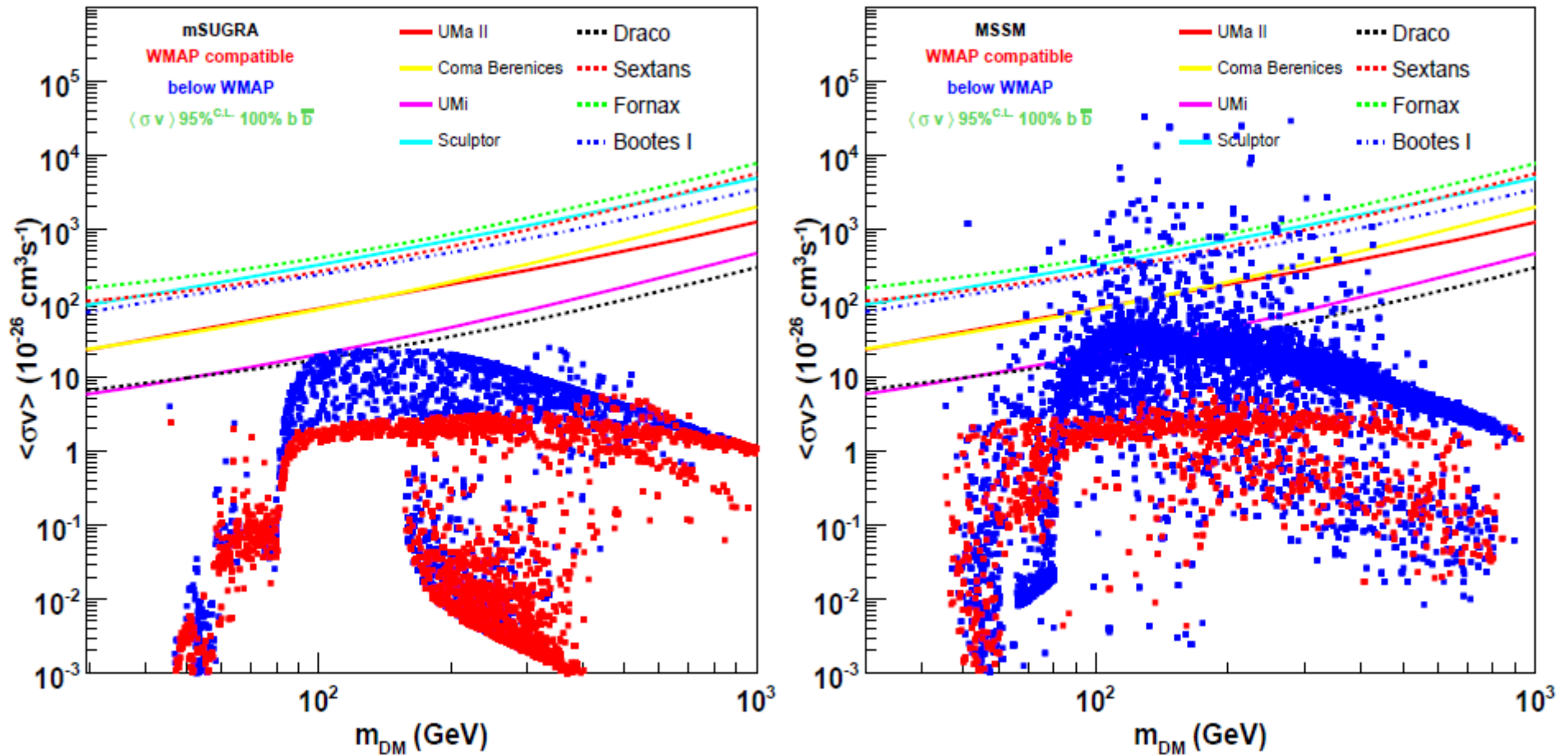
$$\rho(r) = \begin{cases} \frac{\rho_s r_s^3}{r(r_s+r)^2} & \text{for } r < r_t \\ 0 & \text{for } r \geq r_t \end{cases}$$

Name	ρ_s ($M_\odot \text{ pc}^{-3}$)	r_s (kpc)	J^{NFW} ($10^{19} \text{ GeV}^2 \text{ cm}^{-5}$)
Segue 1	1.65	0.05	0.97
Ursa Major II	0.17	0.25	0.57
Segue 2	0.61	0.06	0.1
Willman 1	0.417	0.17	0.84
Coma Berenices	0.232	0.22	0.42
Ursa Minor	0.04	0.97	0.35
Sculptor	0.063	0.52	0.12
Draco	0.13	0.50	0.43
Sextans	0.079	0.36	0.05
Fornax	0.04	1.00	0.11



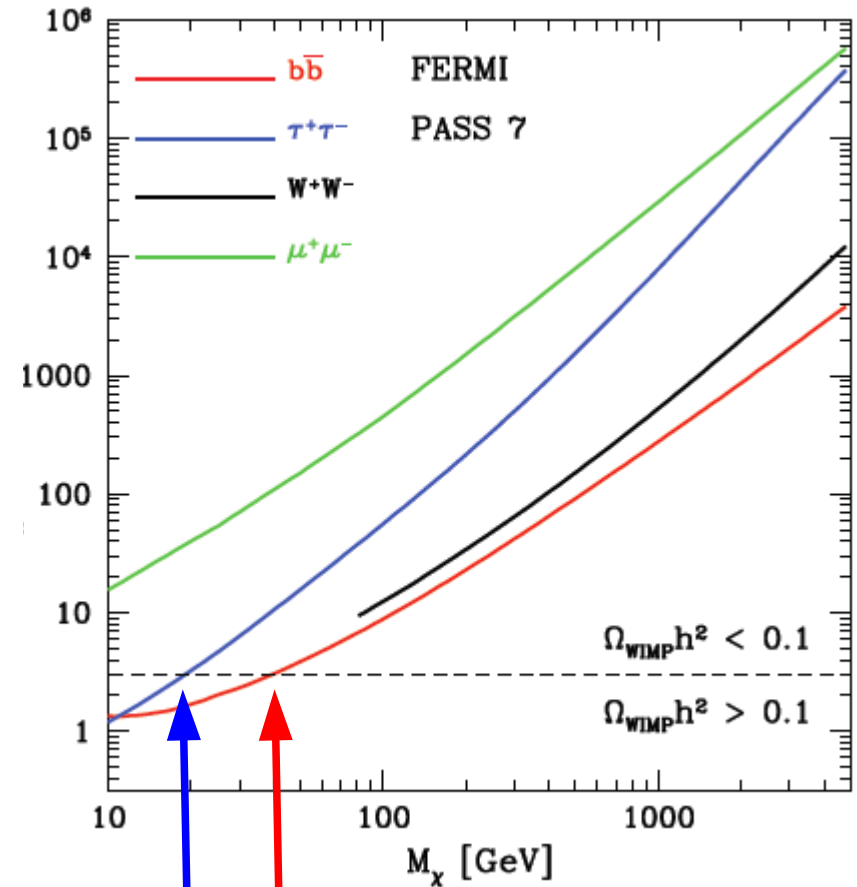
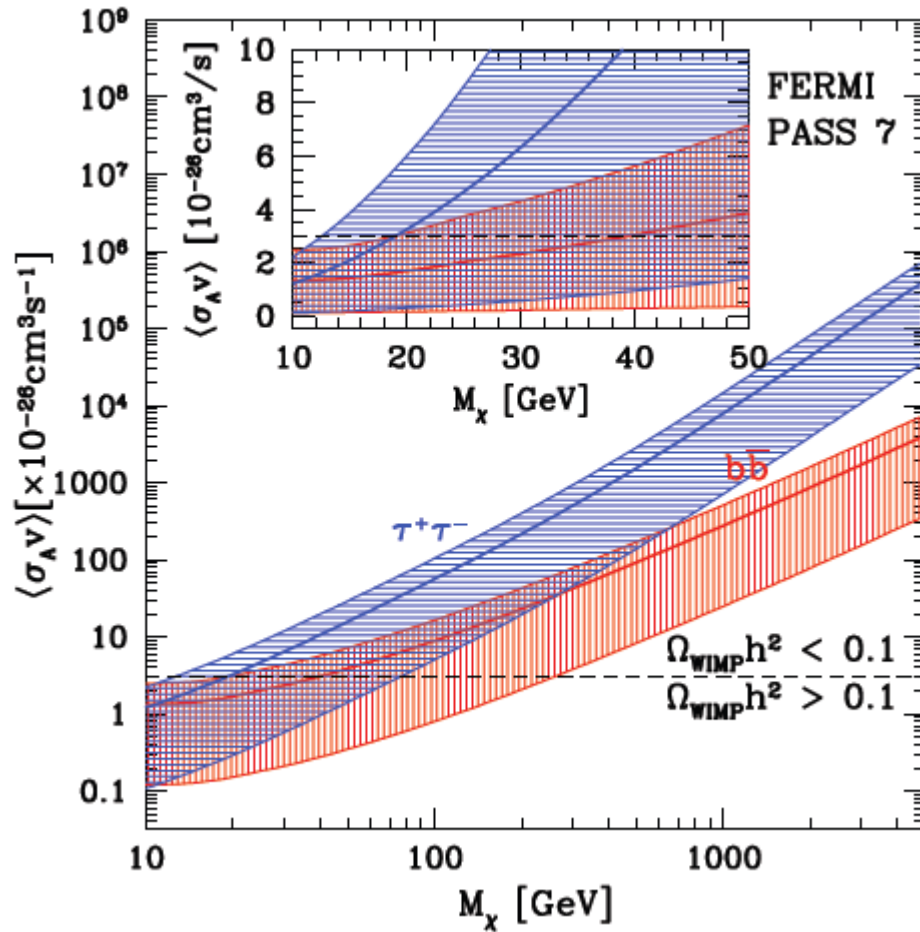
$$J(\psi) = \int_{\text{l.o.s}} dl(\psi) \rho^2(l(\psi))$$

Constraints on WIMP dark matter models



Closing in on light WIMP scenarios from dwarf galaxy observations

Geringer-Sameth, Koushiappas '11



$M_{\text{DM}} > 40 \text{ GeV}$ for $\text{DM DM} \rightarrow b\bar{b}$

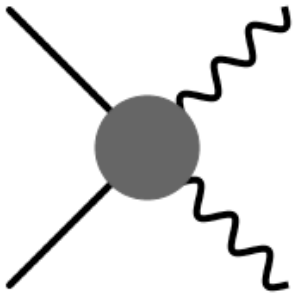
$M_{\text{DM}} > 19 \text{ GeV}$ for $\text{DM DM} \rightarrow \tau^+ \tau^-$

Gamma-ray features

“Smoking gun” for dark matter: no (known) astrophysical process can produce a sharp feature in the gamma-ray energy spectrum

Three gamma-ray spectral features have been identified:

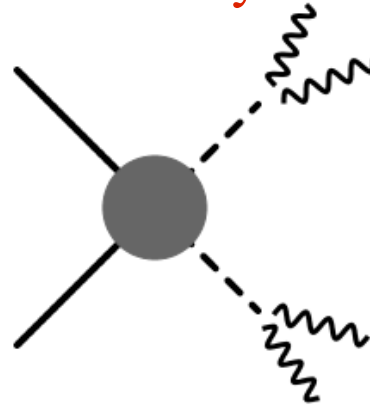
Gamma ray line



Srednicki, Theisen, Silk '86
Rudaz '86
Bergstrom, Snellman '88

$$\langle\sigma v\rangle^{\text{expected}} \lesssim 10^{-29} \text{ cm}^3 \text{ s}^{-1}$$

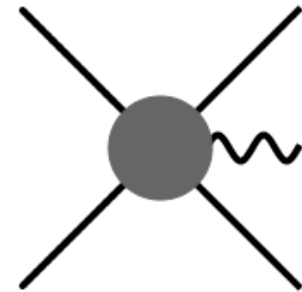
Gamma ray box



AI, Lopez Gehler, Pato '12

$$\langle\sigma v\rangle^{\text{expected}} \lesssim 10^{-26} \text{ cm}^3 \text{ s}^{-1}$$

Internal bremsstrahlung



Bergstrom '89
Flores, Olive, Rudaz '89
Bringmann, Bergstrom, Edsjo '08

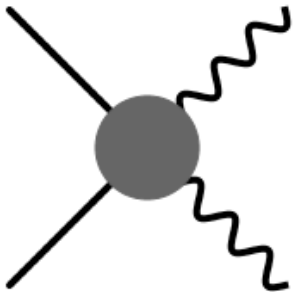
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Gamma-ray features

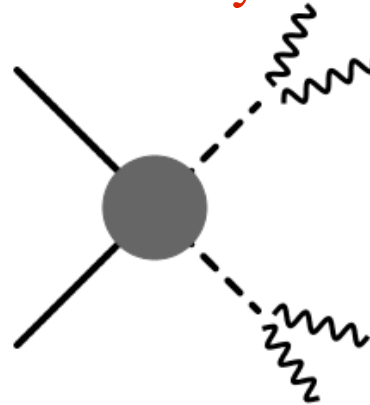
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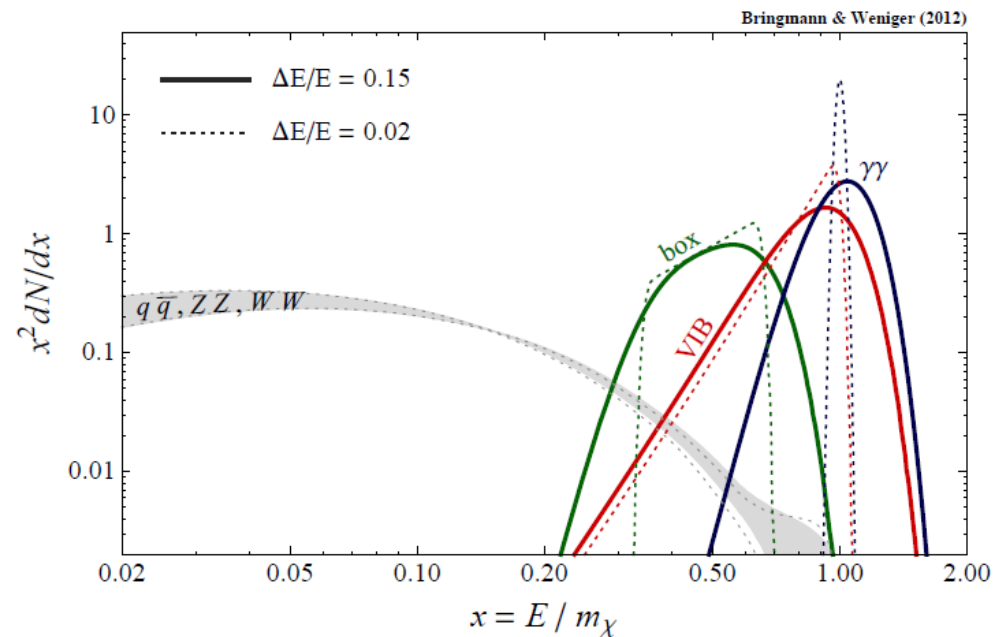
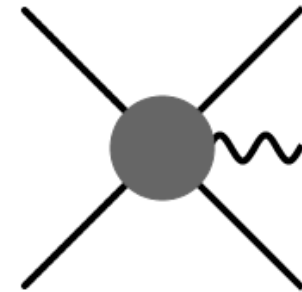
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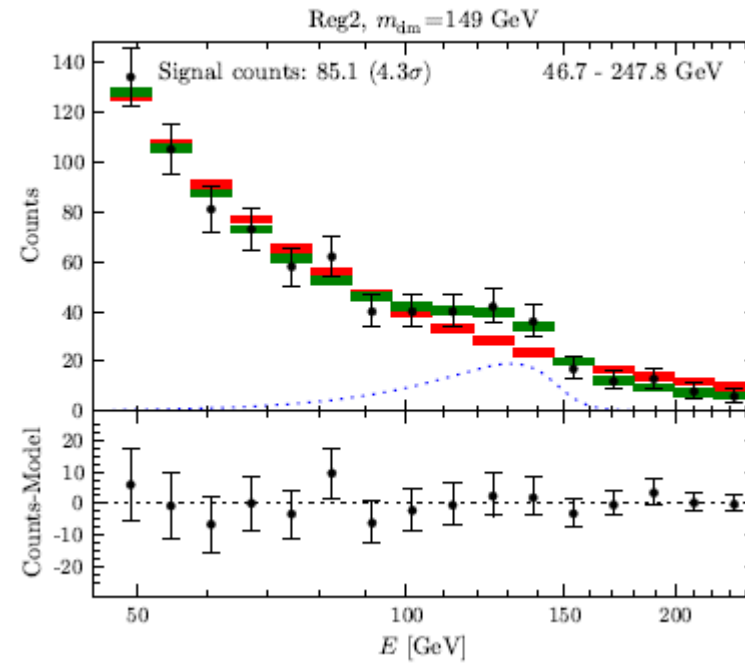
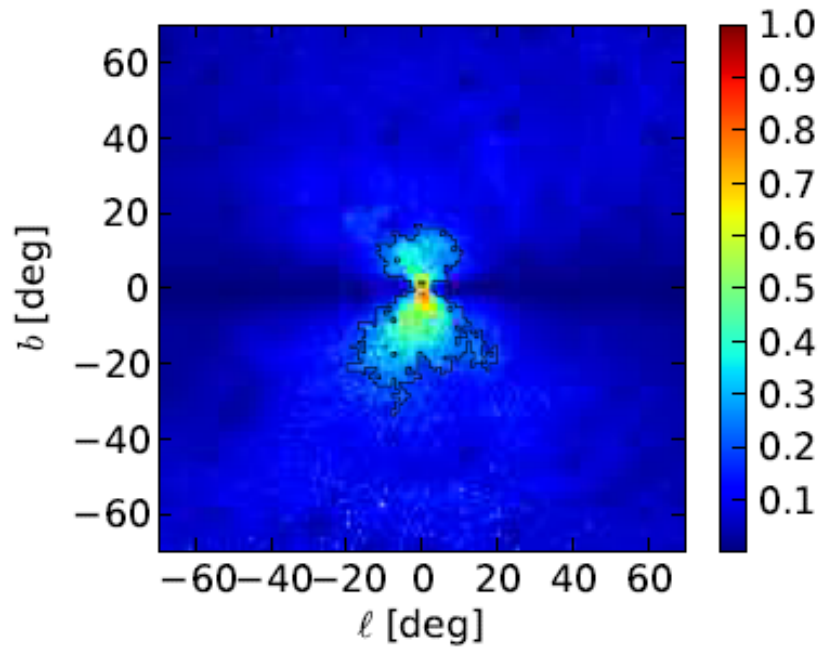
Gamma ray box



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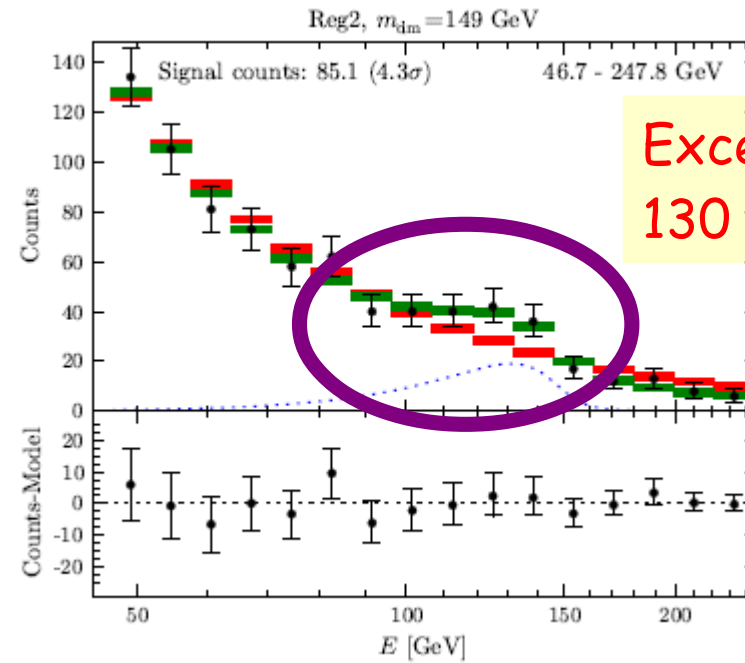
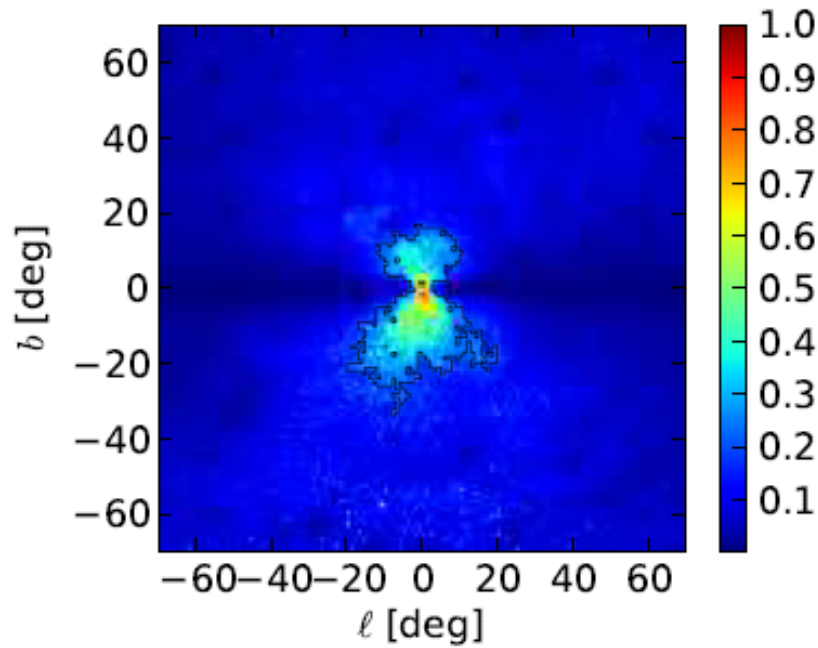


Searching for spectral features with the Fermi-LAT



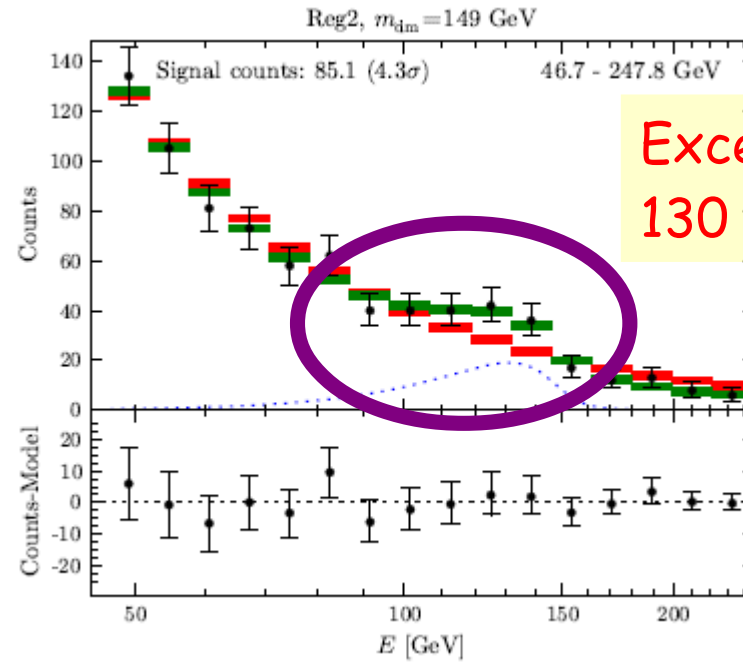
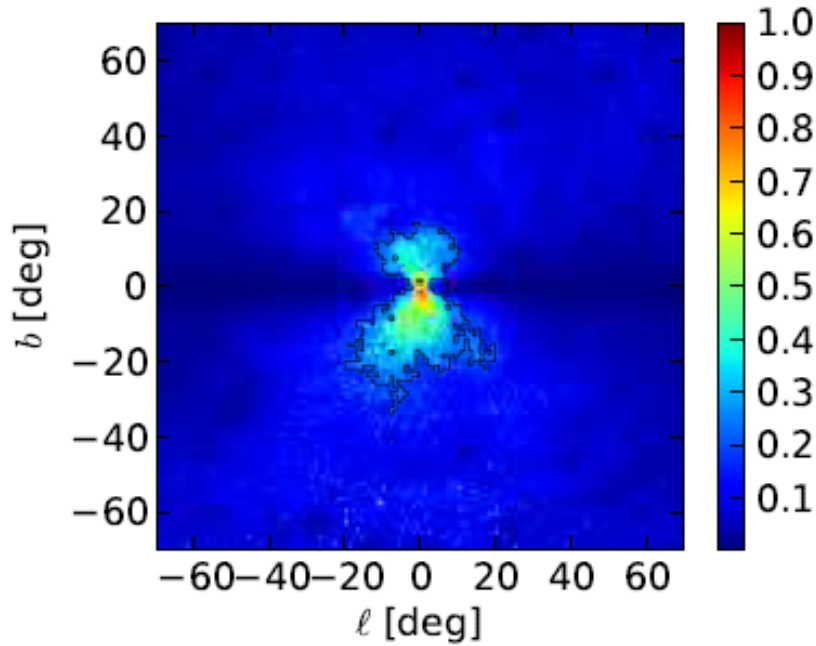
Bringmann, Huang,
AI, Vogl, Weniger
arXiv:1203.1312

Searching for spectral features with the Fermi-LAT

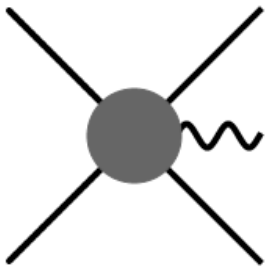


Bringmann, Huang,
AI, Vogl, Weniger
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Searching for spectral features with the Fermi-LAT



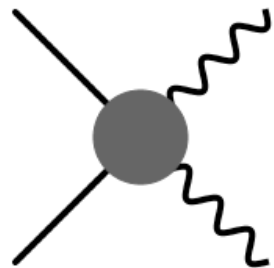
Bringmann, Huang,
AI, Vogl, Weniger
arXiv:1203.1312



$$m_\chi = (149 \pm 4) \text{ GeV}$$

$$\langle \sigma v \rangle = (5.7 \pm 1.4) \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}$$

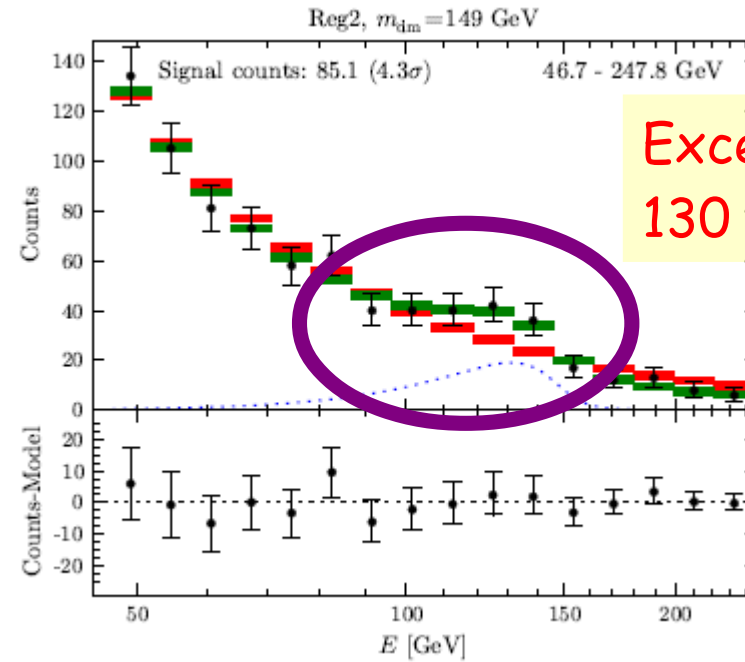
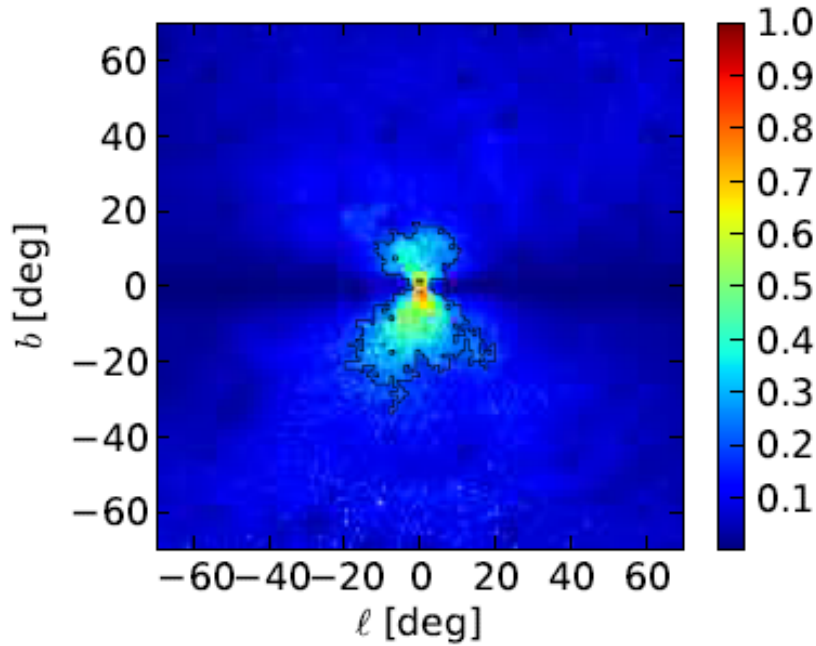
4.3 σ (3.1 σ with LEE)



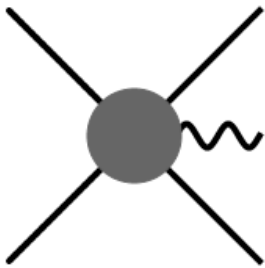
$$m_\chi \sim 130 \text{ GeV}$$

$$\langle \sigma v \rangle_{\chi\chi \rightarrow \gamma\gamma} \sim 10^{-27} \text{ cm}^3 \text{ s}^{-1}$$

Searching for spectral features with the Fermi-LAT



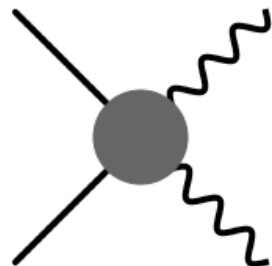
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$$m_\chi = 129.8 \pm 2.4^{+7}_{-13} \text{ GeV}$$

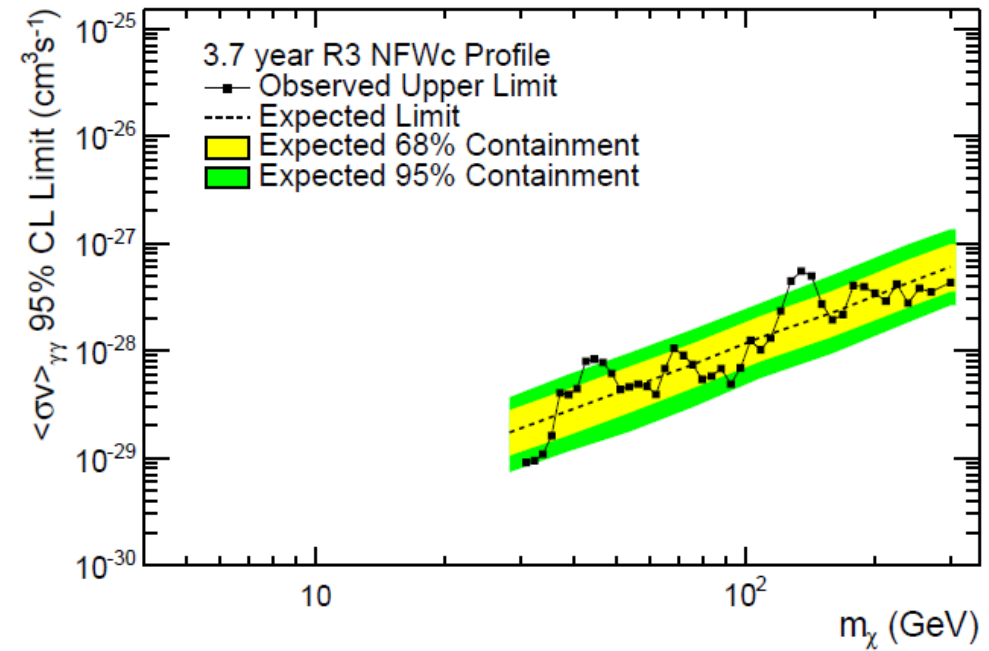
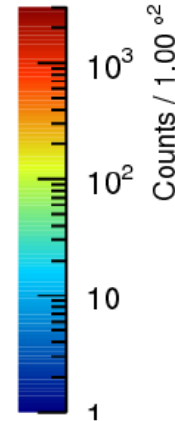
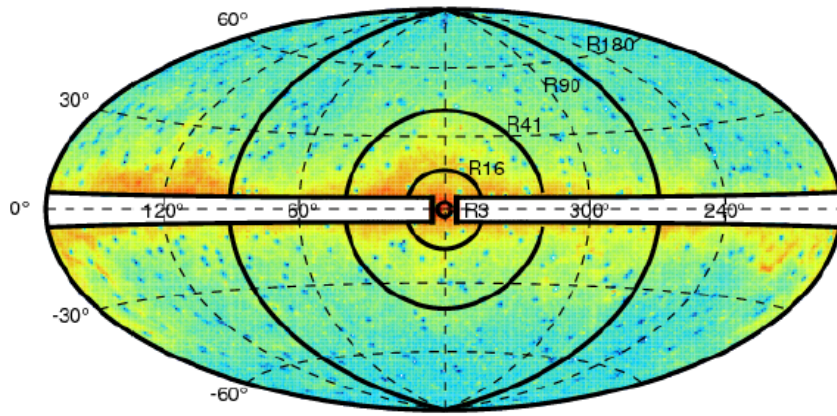
$$\langle \sigma v \rangle = (1.27 \pm 0.32^{+0.18}_{-0.28}) \times 10^{-27} \text{ cm}^3 \text{ s}^{-1}$$

4.6 σ (3.3 σ with LEE)

Weniger,
arXiv:1204.2797

Latest news on the 130 GeV excess

Fermi-LAT collaboration
arXiv:1305.5597

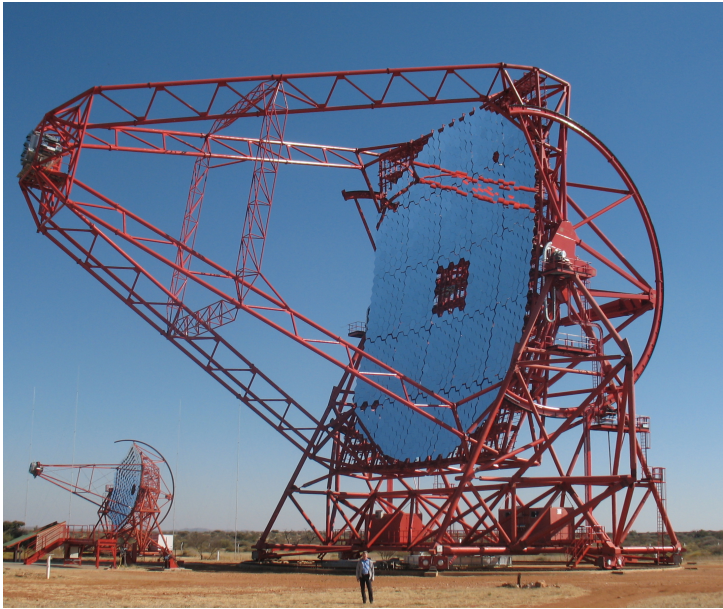


Significance reduced to 3.3σ (1.6σ with LEE)

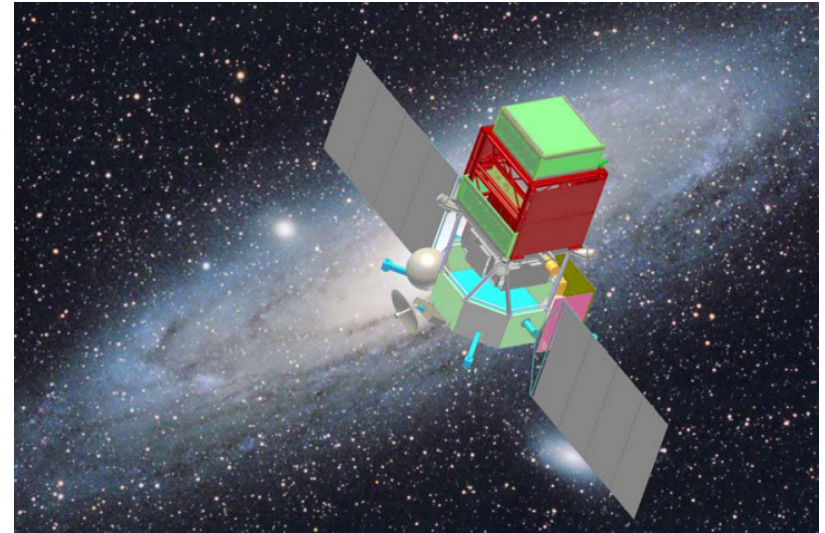
The 130 GeV excess could be just a statistical fluke

Bright future for dark matter searches using gamma-rays!

H.E.S.S. II – in operation



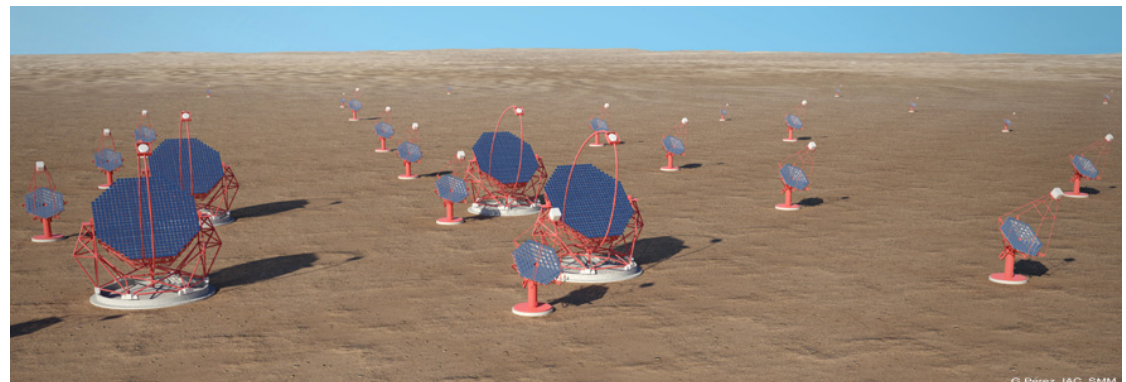
GAMMA 400 – Launch in 2018



DAMPE – Launch in 2015



CTA – Construction starting in 2017



Direct Dark Matter Searches

Direct dark matter searches

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General idea:

1) The Sun (and the Earth) is moving through a “gas” of dark matter particles. Or, from our point of view, there is a flux of dark matter particles going through the Earth.

Direct dark matter searches

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- 2) Once in a while a dark matter particle will interact with a nucleus.

Direct dark matter searches

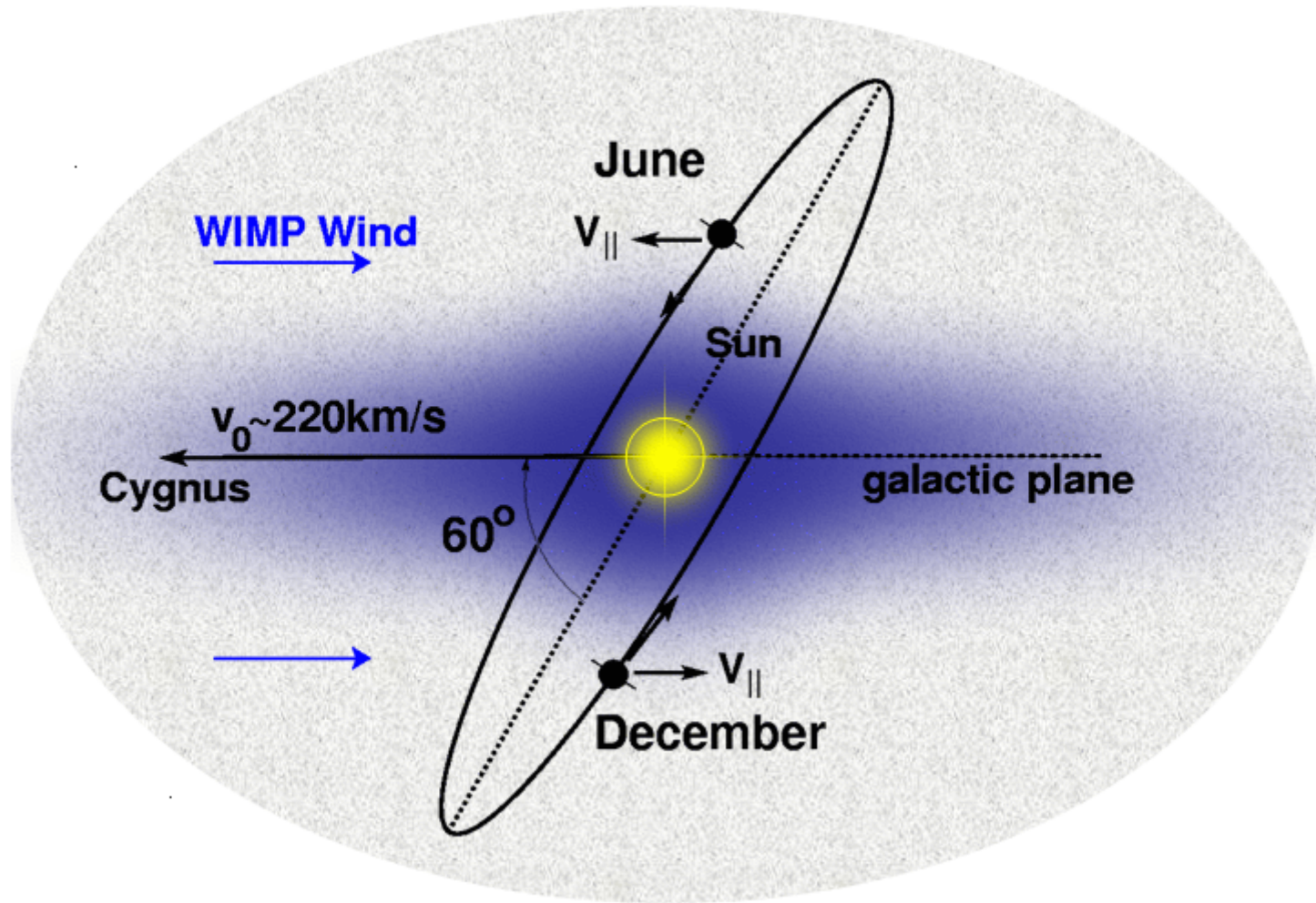
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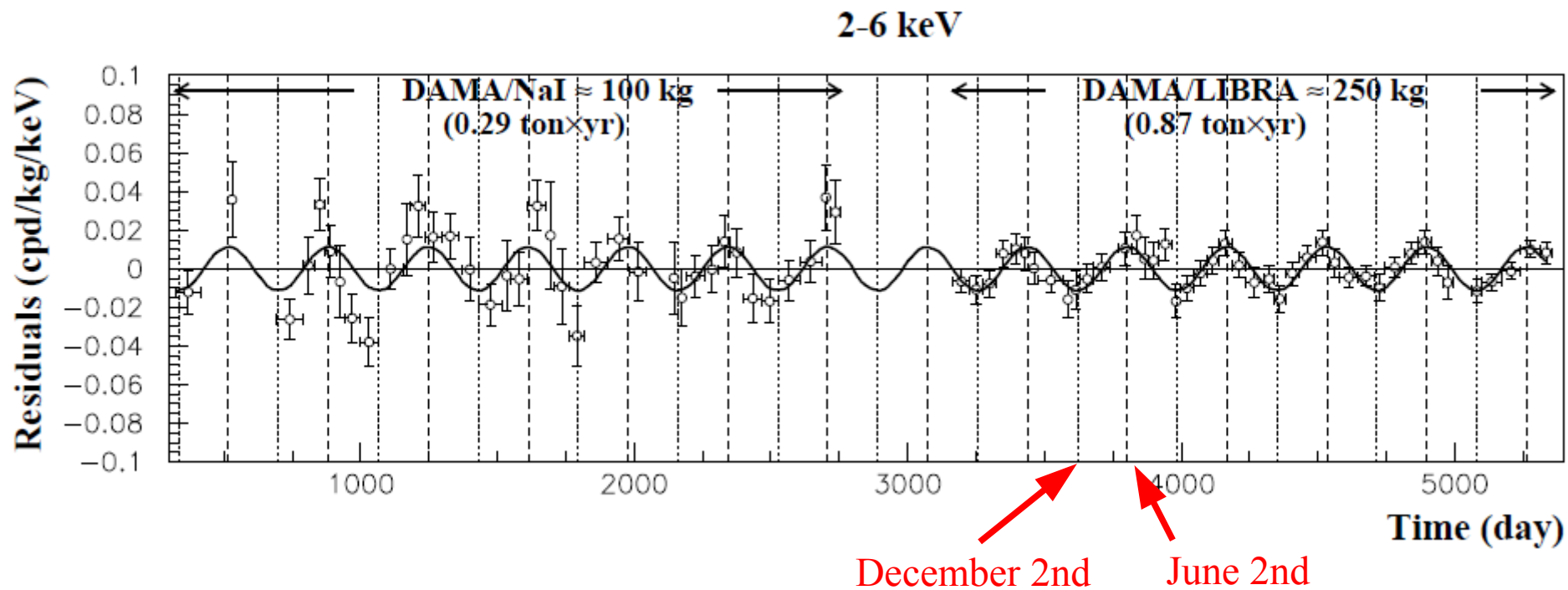
- 1) The Sun (and the Earth) is moving through a “gas” of dark matter particles. Or, from our point of view, there is a flux of dark matter particles going through the Earth.
- 2) Once in a while a dark matter particle will interact with a nucleus.
- 3) The nucleus gains momentum and recoils. The existence of dark matter can then be inferred if there is a significant excess in the number of recoils compared to the expected recoils induced by natural radioactivity in your lab or in your detector.

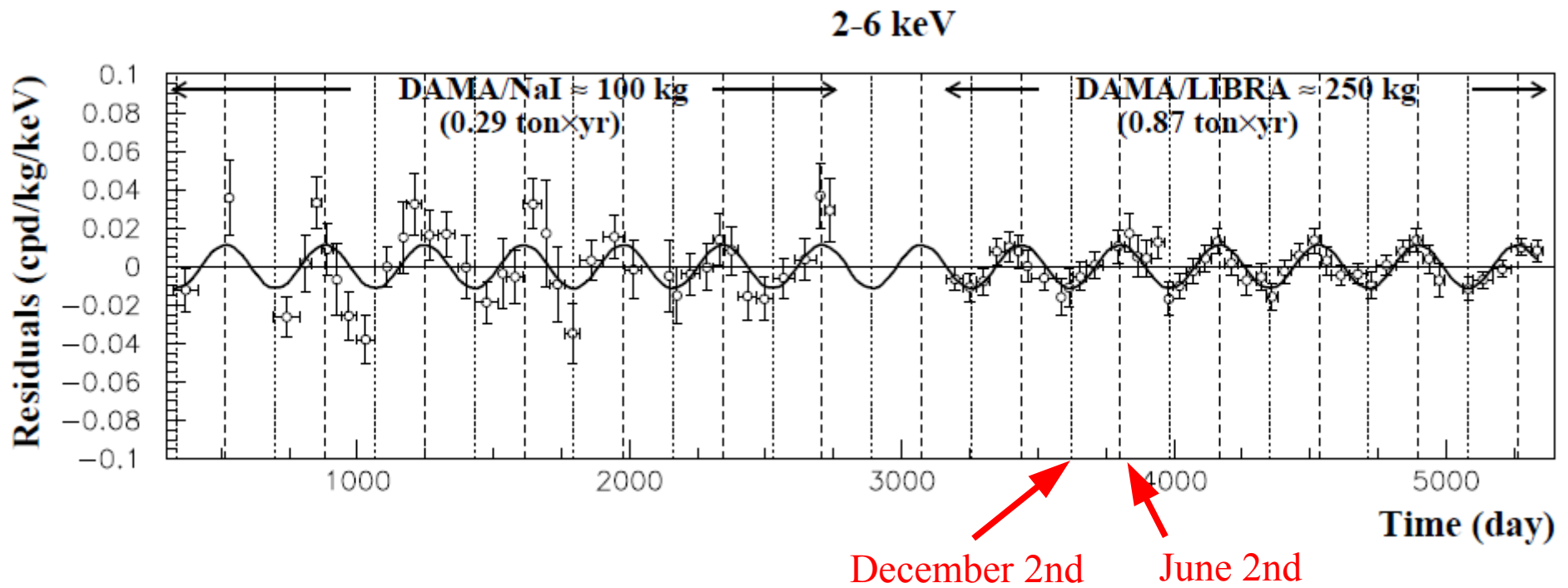
Simple idea ...

... but very challenging in practice!

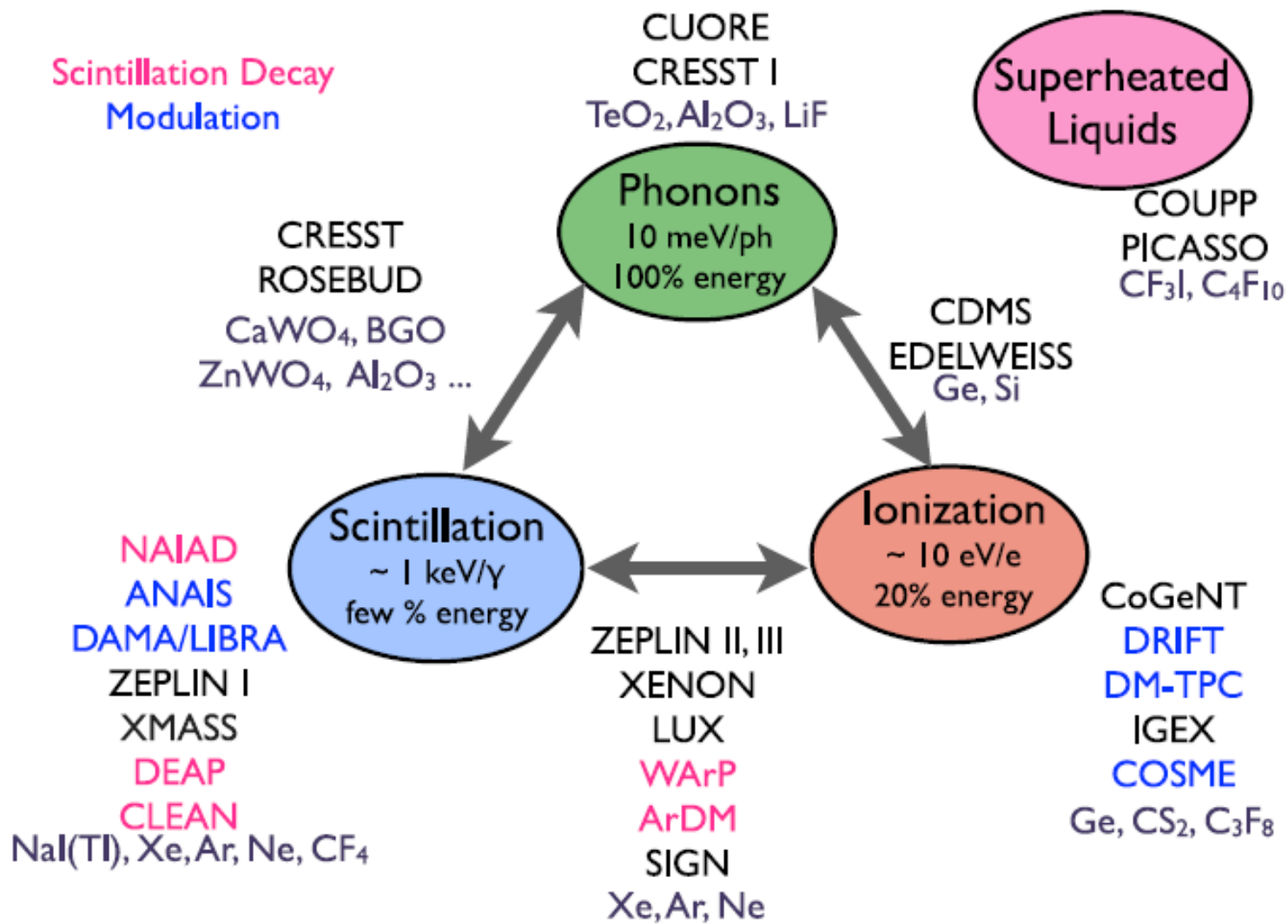
Annual modulation

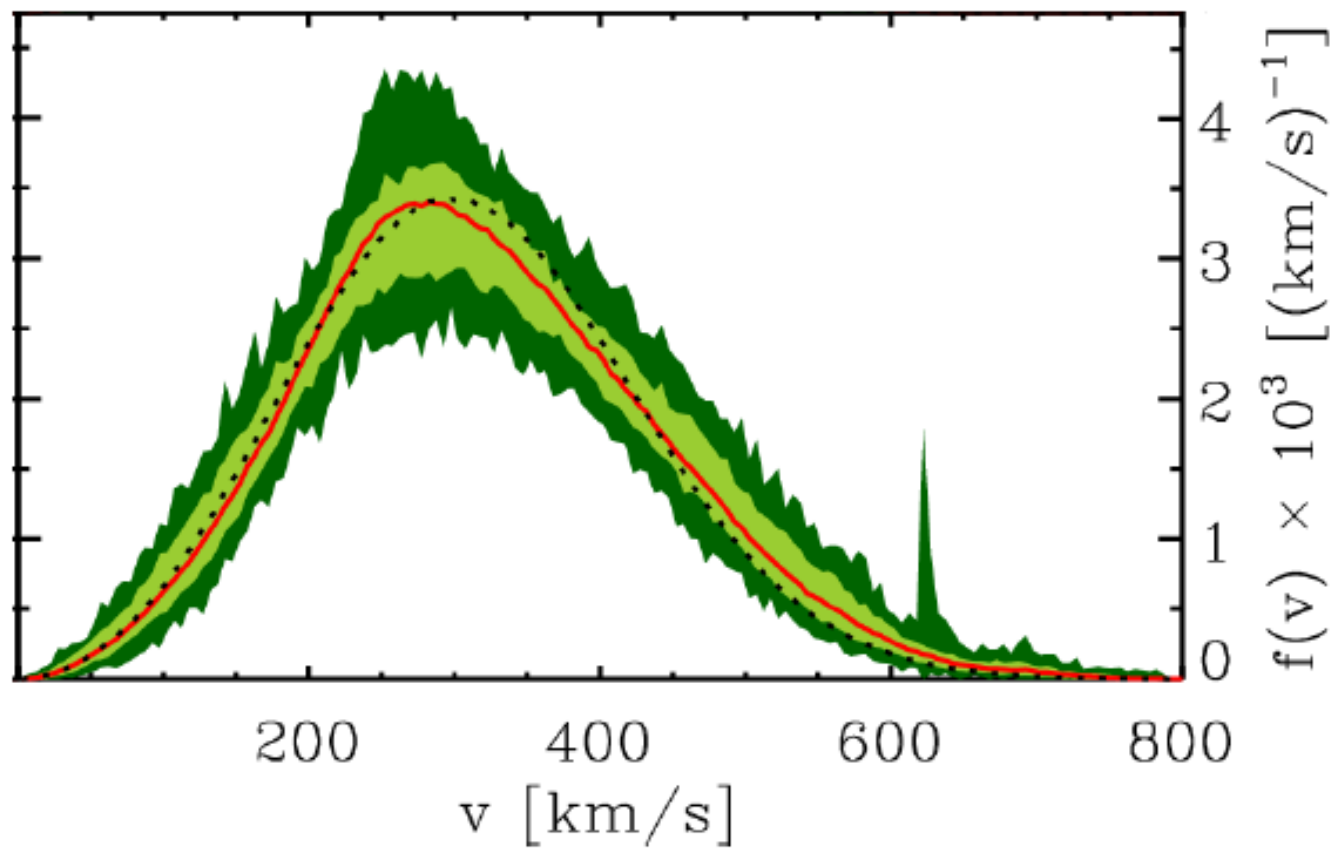




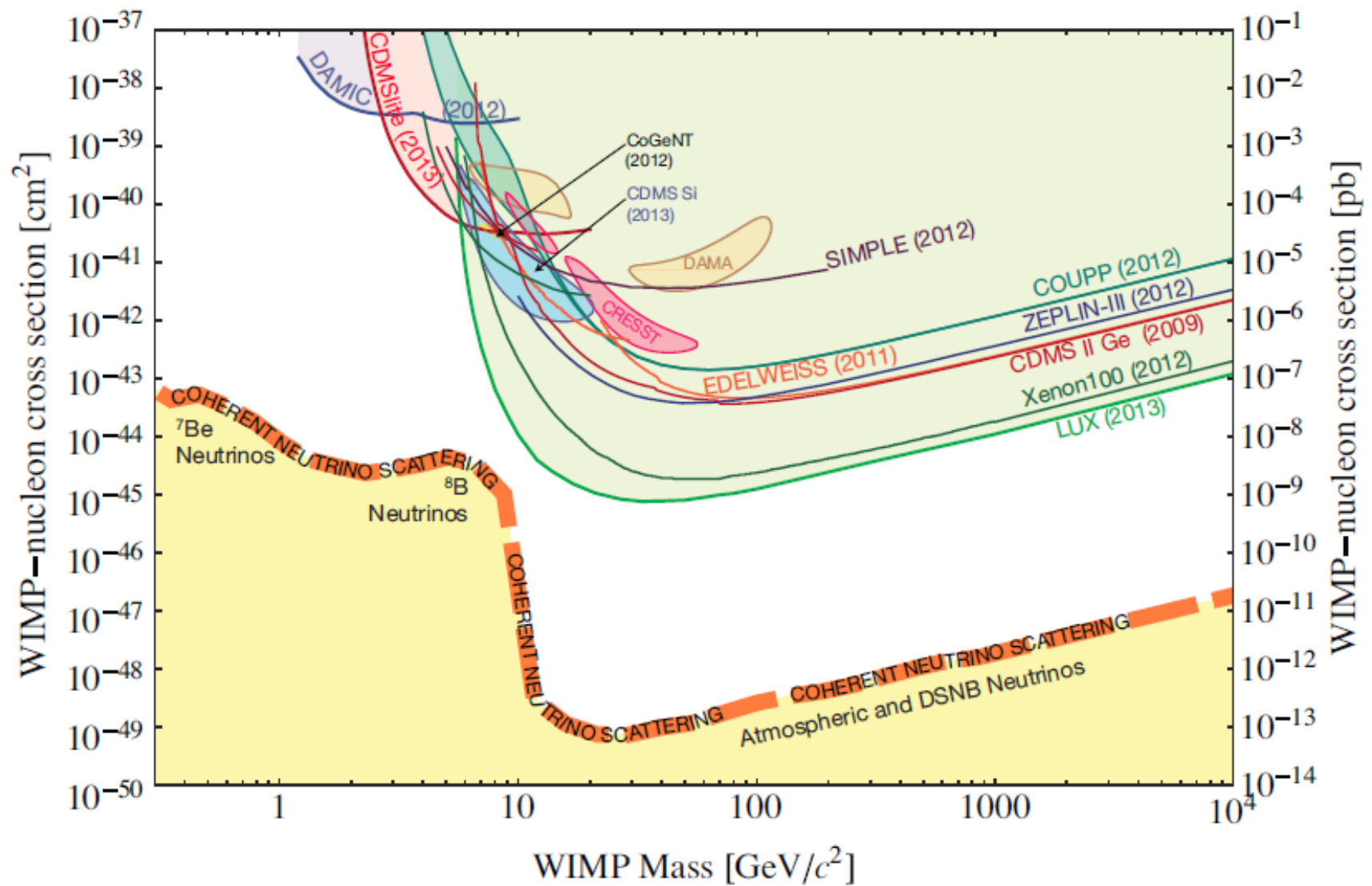


DM interpretation very controversial! More later...

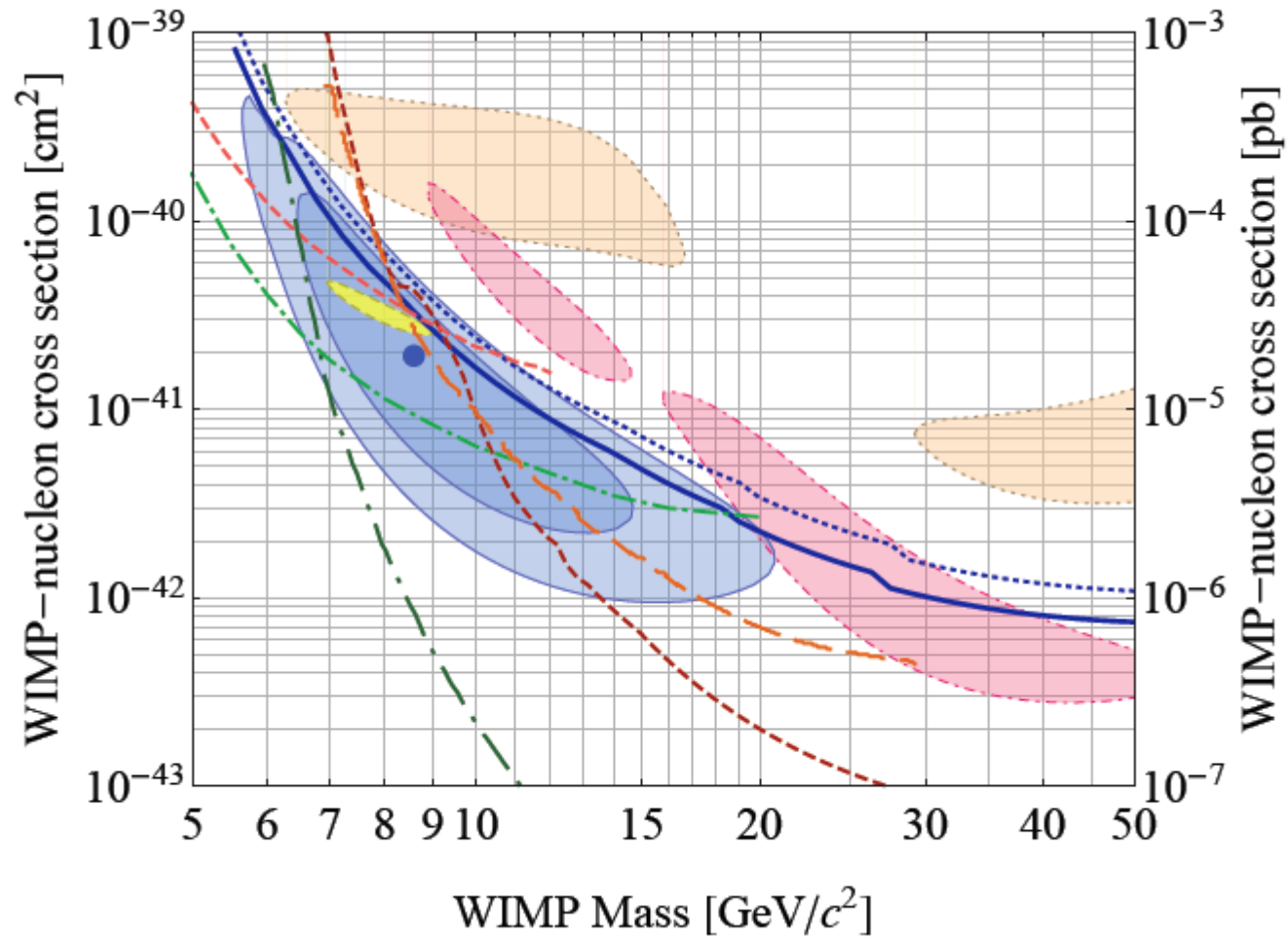




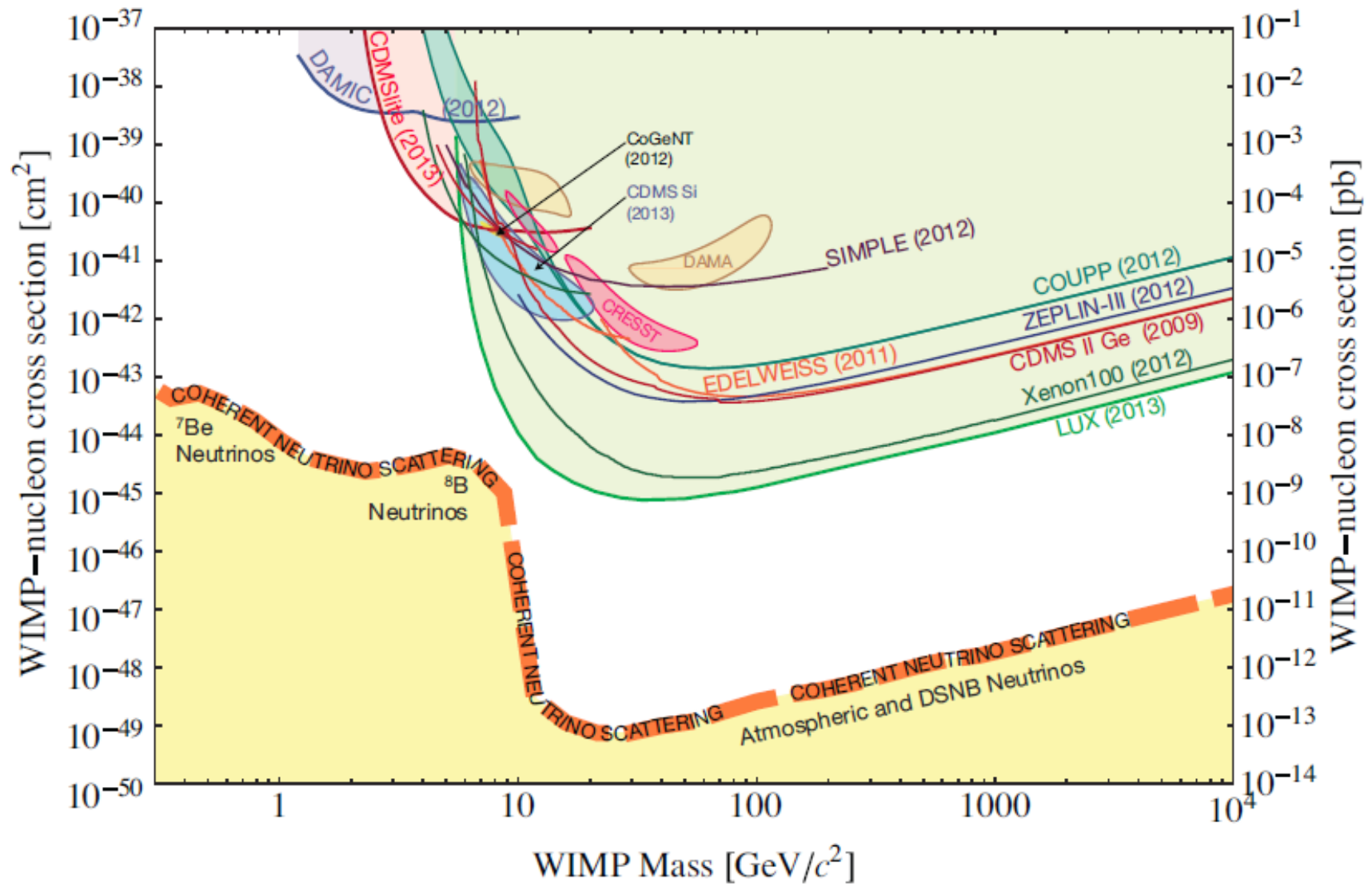
Kuhlen et al.'09



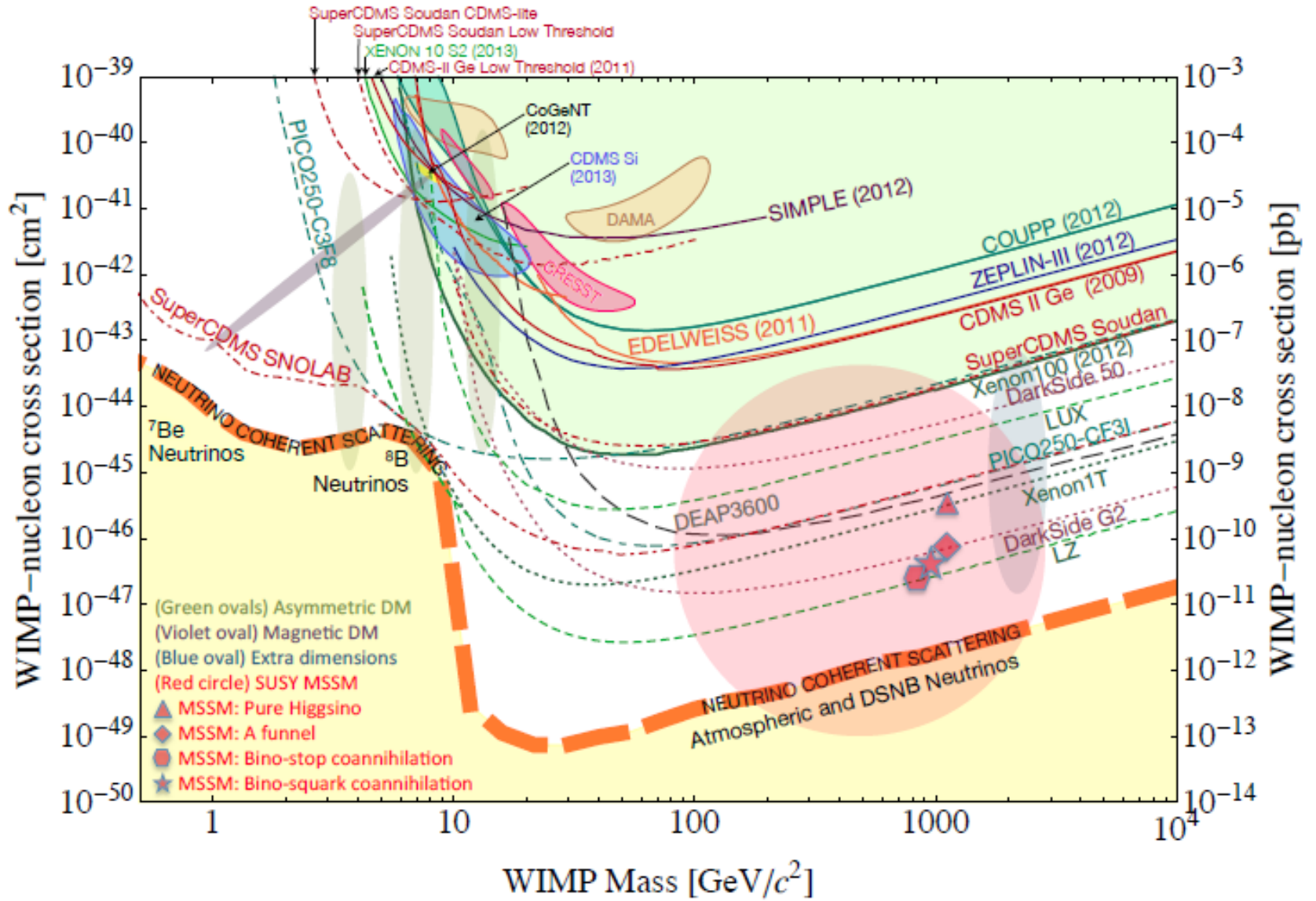
Billard,, Figueroa-Feliciano, Strigari '14



arXiv:1304.4279



Billard,, Figueroa-Feliciano, Strigari '14



List of conclusions

End of list

Concluding remarks

1- Zwicky's observations of 1933

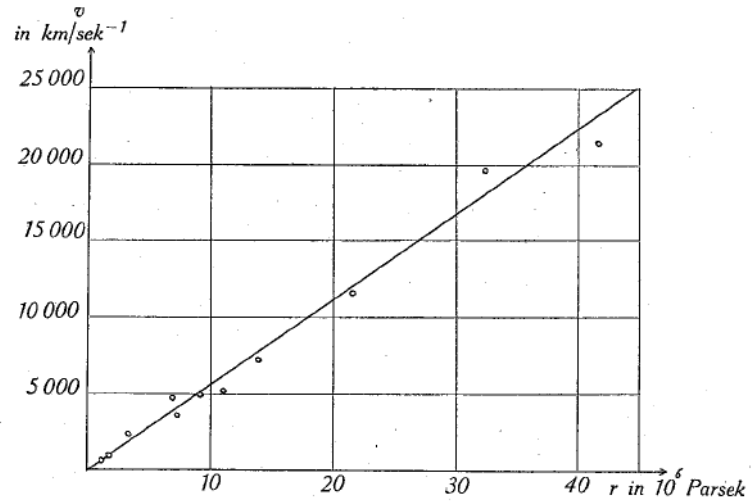


Fig. 2.

Concluding remarks

1- Zwicky's observations of 1933

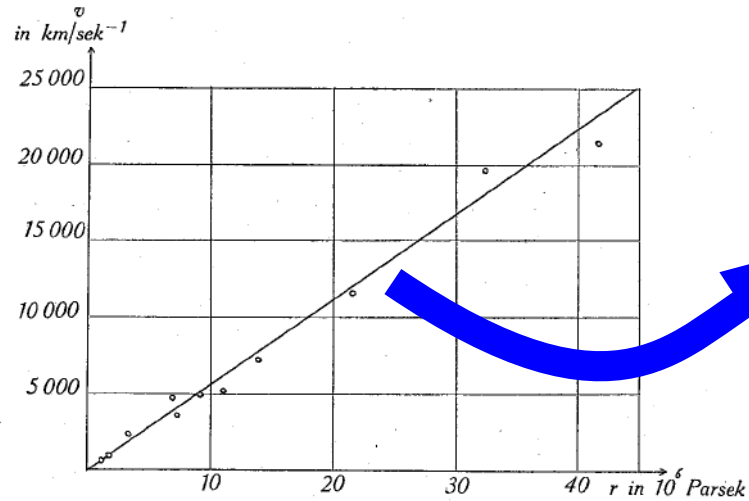
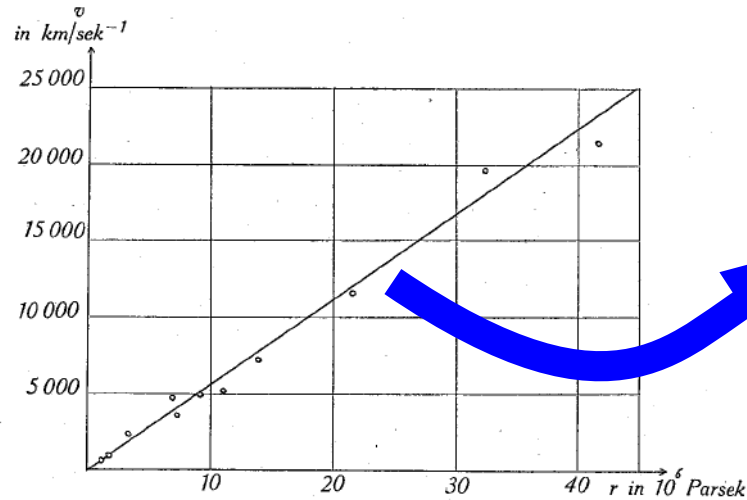


Fig. 2.

80 years later, we still don't know what is producing this.

Concluding remarks

1- Zwicky's observations of 1933

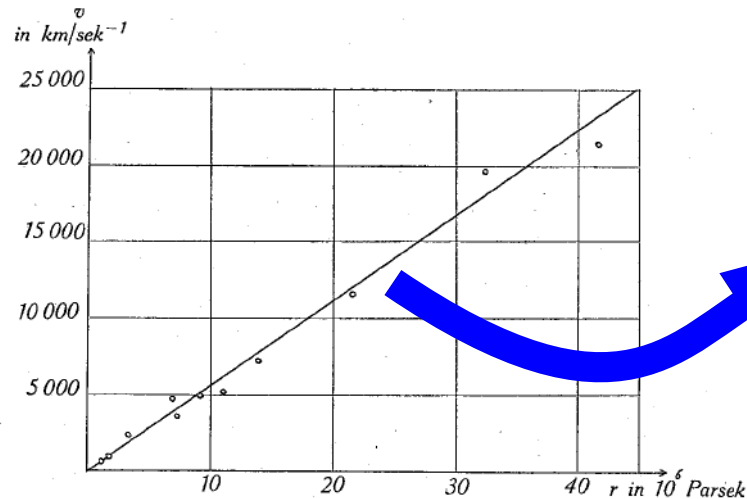


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2- If the dark matter is constituted by WIMPs, there are good chances to observe new signals **in this decade**. Exciting times ahead!

Concluding remarks

1- Zwicky's observations of 1933



80 years later, we still don't know what is producing this.

2- If the dark matter is constituted by WIMPs, there are good chances to observe new signals **in this decade**. Exciting times ahead!

3- BUT, the dark matter particle could not be a WIMP. Or perhaps the astronomical observations of galaxies, clusters of galaxies, etc. are explained by something completely different (not yet proposed).

Keep an open mind!

Thank you for your attention!