

The Higgs and EWPO in R-symmetric SUSY

[arXiv:1408.xxxx]

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Cargèse
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R-Symmetry

- Symmetry allowed by SUSY algebra
- For (N=1)-SUSY: $U(1)_R$ symmetry
- Superfield components have different charges
- SM fields have $Q_R = 0$ (R-parity is subgroup)
 - no rapid proton decay
 - forbids A-terms e.g. $A y_u \tilde{Q} H_u \tilde{U}$
 - forbids Majorana masses and μ term

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Already excluded?

One way to fix it: Dirac masses (MRSSM; *Kribs et.al. arXiv:0712.2039*)

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Soft breaking

$$M_i \tilde{\lambda}_i \tilde{\lambda}_i \rightarrow M_i^D \tilde{\lambda}_i^a \psi_j^a; \{i, j\} \in \{\{G, O\}, \{W, T\}, \{B, S\}\};$$

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Additional fields

		$SU(3)_C$	$SU(2)_L$	$U(1)_Y$	$U(1)_R$
Singlet	\hat{S}	1	1	0	0
Triplet	\hat{T}	1	3	0	0
Octet	\hat{O}	8	1	0	0
R-Higgses	\hat{R}_u	1	2	$-1/2$	2
	\hat{R}_d	1	2	$1/2$	2

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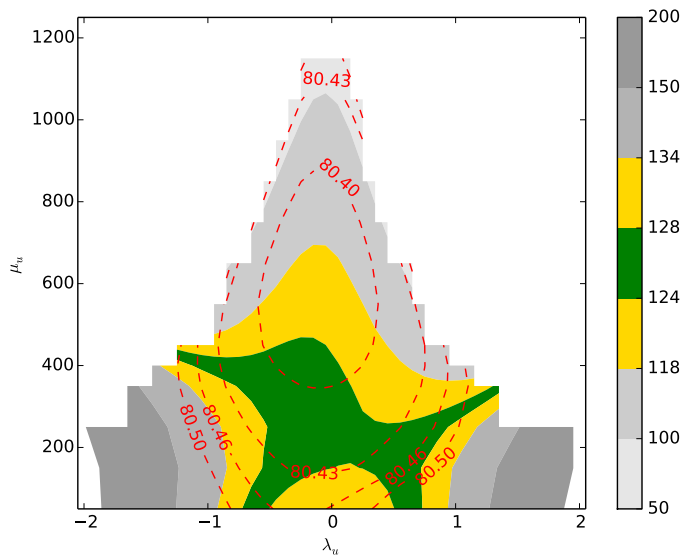
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Superpotential

$$\begin{aligned} \mathcal{W} = & y_e \hat{H}_d \hat{L} \hat{E} + y_d \hat{H}_d \hat{Q} \hat{d} - y_u \hat{H}_u \hat{Q} \hat{U} + \\ & \mu_d \hat{H}_d \hat{R}_d + \mu_u \hat{H}_u \hat{R}_u + \\ & \lambda_d \hat{H}_d \hat{R}_d \hat{S} + \lambda_u \hat{H}_u \hat{R}_u \hat{S} + \Lambda_d \hat{H}_d \hat{T} \hat{R}_d + \Lambda_u \hat{H}_u \hat{T} \hat{R}_u \end{aligned}$$

Higgs mass vs. W mass (at one-loop)



Conclusions

- Introduction to R-symmetry for model building:
→ global $U(1)$ symmetry allowed by SUSY algebra
- Different μ than in the MSSM, different λ 's than in the NMSSM
→ MRSSM with different phenomenology with complementary features
- lightest Higgs mass possible, as well as W mass

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Thanks for the attention!

References



MRSSM G. D. Kribs, E. Poppitz and N. Weiner, “Flavor in supersymmetry with an extended R-symmetry,” *Phys. Rev. D* **78** (2008) 055010 [arXiv:0712.2039 [hep-ph]].

Backup

SUSY and R-Symmetry

usual: $\{Q, Q\} = \{\bar{Q}, \bar{Q}\} = 0$

in "Haag-Łopuszański-Sohnius-Theorem": $\{Q_\alpha^L, Q_\beta^M\} = \epsilon_{\alpha\beta} \sum_I (a^I)^{LM} R^I$

$[R, Q] = -Q, \quad [R, \bar{Q}] = \bar{Q},$

\rightarrow in $\mathcal{N} = 1$ is $U(1)_R$

R-Symmetry

- additional symmetry allowed by SUSY algebra described in “Haag-Łopuszański-Sohnius-Theorem”
- For $N = 1$ SUSY it is a global $U(1)_R$ symmetry
→ charged Spinor coordinates:
 $Q_R(\theta) = 1, Q_R(\bar{\theta}) = -1; (\theta \rightarrow e^{i\alpha}\theta, \bar{\theta} \rightarrow e^{-i\alpha}\bar{\theta})$
- Lagrangian has to be invariant:
 - terms in superpotential have $Q_R = 2; \mathcal{L}_W = \int d^2\theta \mathcal{W}$
 - softbreaking terms have $Q_R = 0$
- SM fields have $Q_R = 0$

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R charges of component fields

	Q_R	scalar	vector	fermionic
vector superfield	0	-	0	1
chiral superfield	Q	Q	-	$Q - 1$

→ Higgs superfield: $Q_R = 0$; lepton and quark superfields: $Q_R = 1$

Dirac masses for soft breaking

$$\mathcal{L}_{\text{dirac}} = \int d\theta^2 \hat{S} \hat{W}'^\alpha \hat{W}_{B\alpha} = -M^D \tilde{\lambda}_B \psi_S^a + \sqrt{2} M^D D_B^a \phi_S^a$$

using Spurion: $\hat{W}'^\alpha = \sqrt{2} M^D \theta^\alpha$

Lightest Higgs mass - Tree level

EWSB

- Separating h_u, h_d, s, t in vevs, scalar and pseudo-scalar part
- Forming mass eigenstates with (4×4) mass matrices

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For $v_S, v_T \ll v$ and $(M^D)^2 \ll m_{\text{soft}}^2$
($\lambda_u = -\lambda_d = \lambda, \Lambda_u = \Lambda_d = \Lambda, \mu_u = \mu_d = \mu$):

$$m_{h, \text{tree}}^2 = \left(m_Z^2 - v^2 \left(\frac{(g_1 M_B^D + \sqrt{2} \lambda \mu)^2}{8(M_B^D)^2 + 2m_S^2} + \frac{(g_2 M_W^D + \Lambda \mu)^2}{8(M_W^D)^2 + 2m_T^2} \right) \right) \cos(2\beta)$$

No enhancement from singlet like in the NMSSM, additional states lower lightest mass on tree level \rightarrow one loop even more important

Lightest Higgs mass - One Loop I

Analytical approximation from the effective potential:

$$\begin{aligned}(\Delta m_h^2)_{1L} = & \frac{3v^2}{16\pi^2} \left[\frac{4\lambda^4 + 4\lambda^2\Lambda^2 + 5\Lambda^4}{8} \log \left(\frac{m_{R_u}^2}{Q^2} \right) \right. \\ & + \left(\frac{\lambda^4}{2} - \frac{\lambda^2\Lambda^2}{2} \frac{m_S^2}{m_T^2 - m_S^2} \right) \log \left(\frac{m_S^2}{Q^2} \right) \\ & + \left(\frac{5}{8}\Lambda^4 + \frac{\lambda^2\Lambda^2}{2} \frac{m_T^2}{m_T^2 - m_S^2} \right) \log \left(\frac{m_T^2}{Q^2} \right) \\ & - \left(\frac{5}{4}\Lambda^4 - \lambda^2\Lambda^2 \frac{(M_W^D)^2}{(M_B^D)^2 - (M_W^D)^2} \right) \log \left(\frac{(M_W^D)^2}{Q^2} \right) \\ & - \left(\lambda^4 + \lambda^2\Lambda^2 \frac{(M_B^D)^2}{(M_B^D)^2 - (M_W^D)^2} \right) \log \left(\frac{(M_B^D)^2}{Q^2} \right) \\ & \left. + \frac{\Lambda^2\lambda^2}{2} \right] + \text{stop contrb.}\end{aligned}$$

λ or/and Λ can give sizeable contribution, reducing the need for large stop