DECOUPLING OF HEAVY SNEUTRINOS IN LOW-SCALE SEESAW MODELS

BASED ON ARXIV:1312.5318

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MEASUREMENT OF NEUTRINO MIXING:

- $m_{\nu} \neq 0 \Rightarrow$ calls for an extension of the SM
- lepton flavor is violated in the neutral sector

WHAT ABOUT CHARGED LEPTON FLAVOR VIOLATION?

- can arise at 1-loop
- examples: $\mu \to e\gamma, \, \mu \to 3e, \, \mu e$ conversion in nuclei, ...
- can be especially large and low-scale seesaw models

SUPERSYMMETRY:

New sources of cLFV, especially by slepton mixing





usually in supersymmetric models:

- γ penguin contributions dominate
- $\mu \to e\gamma$ most likely to observe



BUT [Hirsch,Vicente,Staub '12; Abada,Das,Vicente,Weiland '12]: large enhancement of Z penguins in low-scale seesaw models: $Z - \text{peng}/\gamma - \text{peng} \sim (M_{SUSY}/M_Z)^4$



more explicitly: minimal supersymmetric inverse seesaw model

$$W = W_{\rm MSSM} + Y_{\nu} \hat{\nu}^c \hat{L} \hat{H}_u + M_R \, \hat{\nu}^c \hat{N}_S + \frac{\mu_N}{2} \hat{N}_S \hat{N}_S$$



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LOW-SCALE SEESAW MODELS



- constant nondecoupling term no $1/\Lambda_{NP}$ suppression? \Rightarrow critical look necessary
- recalculation: amplitude in [Arganda,Herrero '05] is wrong by a constant term

 $[{\rm MEK}, {\rm Porod}, {\rm Staub}, {\rm Abada}, {\rm Vicente}, {\rm Weiland~'13}]$



LOW-SCALE SEESAW

MODELS

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Outlook

- \blacksquare \exists differences in box form factors in the literature
- independent calculation of these operators + all still missing diagrams in a SUSY framework

[Abada,MEK,Porod,Staub,Vicente,Weiland arXiv:1407.XXXX]

 calculating flavour observables made easy with FlavorKit [Porod,Staub,Vicente '14]



LOW-SCALE SEESAW MODELS

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BACKUP SLIDES



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SUSY algebra

$$\begin{split} \{Q_{\alpha}, Q_{\beta}\} &= \{\bar{Q}_{\dot{\alpha}}, \bar{Q}_{\dot{\beta}}\} = 0\\ \{Q_{\alpha}, \bar{Q}_{\dot{\beta}}\} &= 2\sigma^{\mu}_{\alpha\dot{\beta}}P_{\mu}\\ [P^{\mu}, Q_{\alpha}] &= 0 \end{split}$$

Defining quantity of a SUSY model: Superpotential

$$\begin{split} W(\Phi) &= \frac{1}{2} m_{ij} \Phi_i \Phi_j + \frac{1}{3} \lambda_{ijk} \Phi_i \Phi_j \Phi_k \\ W^i &= \frac{\partial}{\partial \phi_i} W(\phi), \qquad W^{ij} = \frac{\partial^2}{\partial \phi_i \partial \phi_j} W(\phi) \\ \text{such that} \\ \mathscr{L}_{\text{chiral}} &= \partial_\mu \phi_i^* \partial^\mu \phi_i + i \bar{\psi}_i \bar{\sigma}^\mu \partial_\mu \psi_i \\ &- \underbrace{\frac{1}{2} (W^{ij} \psi_i \psi_j + W_{ij}^* \bar{\psi}^i \bar{\psi}^j)}_{\text{fermion mass terms and Yukawa interactions}} - \underbrace{W^i W_i^*}_{\text{scalar potential}} \end{split}$$