

Localization on round three-sphere revisited

Akinori Tanaka (Osaka Univ.)

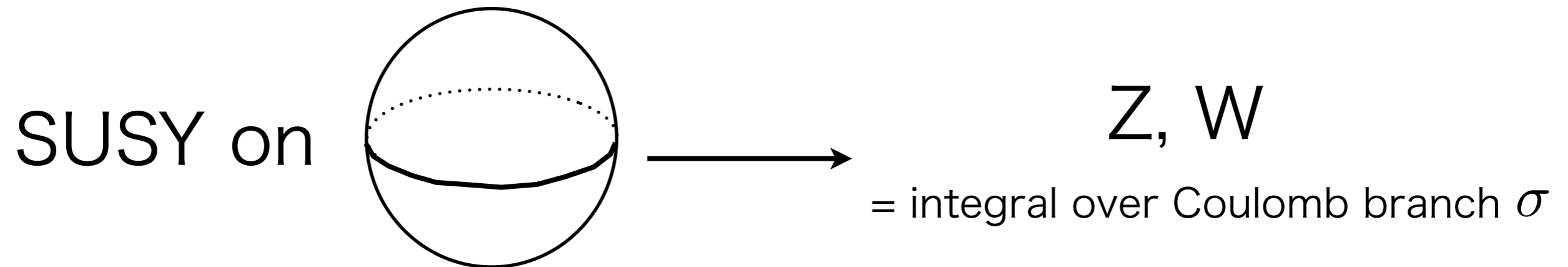
2013/11/30 PASCOS
@GIS conference center
Taipei, Taiwan

arXiv:1309.4992

History of Localization on 3-sphere

2009: Kapustin, Willett, Yaakov

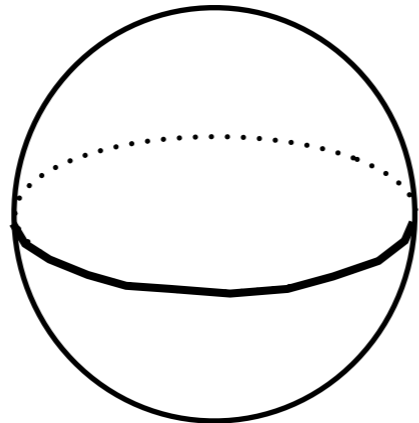
arXiv:0909.4559



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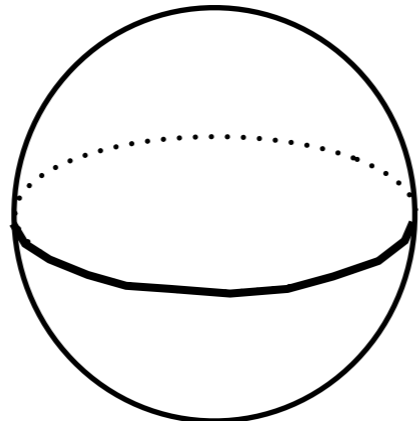
SUSY on  \longrightarrow Z, W
= integral over Coulomb branch σ

$$\mathcal{N} = 2 \text{ Vector} \longrightarrow \int d\sigma \prod_{\alpha > 0} \sinh^2(\pi \alpha \cdot \sigma)$$

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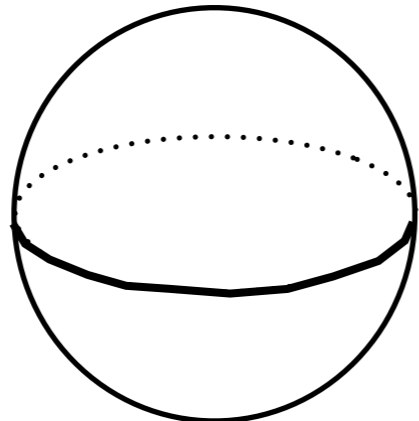
$$\mathcal{N} = 2 \text{ Vector} \longrightarrow \int d\sigma \prod_{\alpha > 0} \sinh^2(\pi \alpha \cdot \sigma)$$

$$\mathcal{N} = 2 \text{ Canonical Matter} \longrightarrow \prod_{\rho: \text{weight}} \left(\frac{1}{\cosh \pi \rho \cdot \sigma} \right)^{\frac{1}{2}}$$

History of Localization on 3-sphere

2010: Hama, Hosomichi, Lee arXiv:1012.3512

Jafferis arXiv:1012.3210

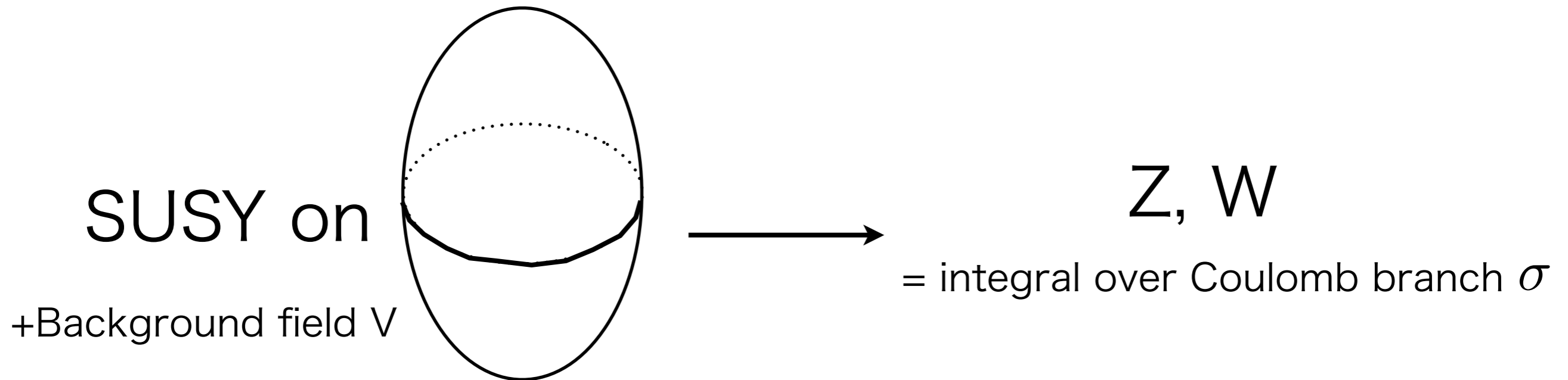
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$$\mathcal{N} = 2 \text{ Anomalous Matter} \longrightarrow \prod_{\rho: \text{weight}} s_{b=1}(i - i\Delta - \rho \cdot \sigma)$$

History of Localization on 3-sphere

2011: Hama, Hosomichi, Lee arXiv:1102.4716

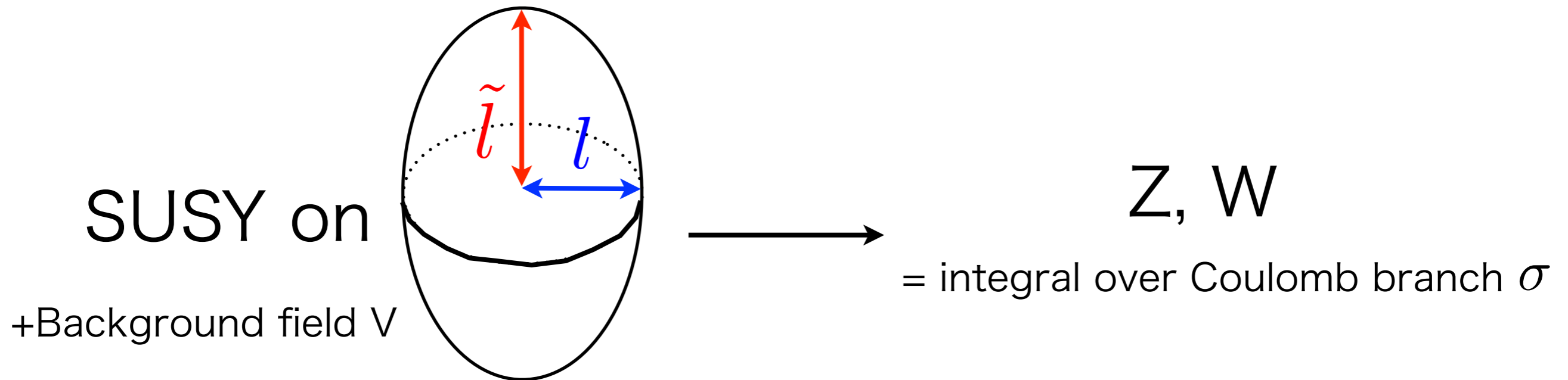


$$\mathcal{N} = 2 \text{ Vector} \longrightarrow \int d\sigma \prod_{\alpha > 0} \sinh \left(\sqrt{\tilde{l}/l} \pi \alpha \cdot \sigma \right) \sinh \left(\sqrt{l/\tilde{l}} \pi \alpha \cdot \sigma \right)$$

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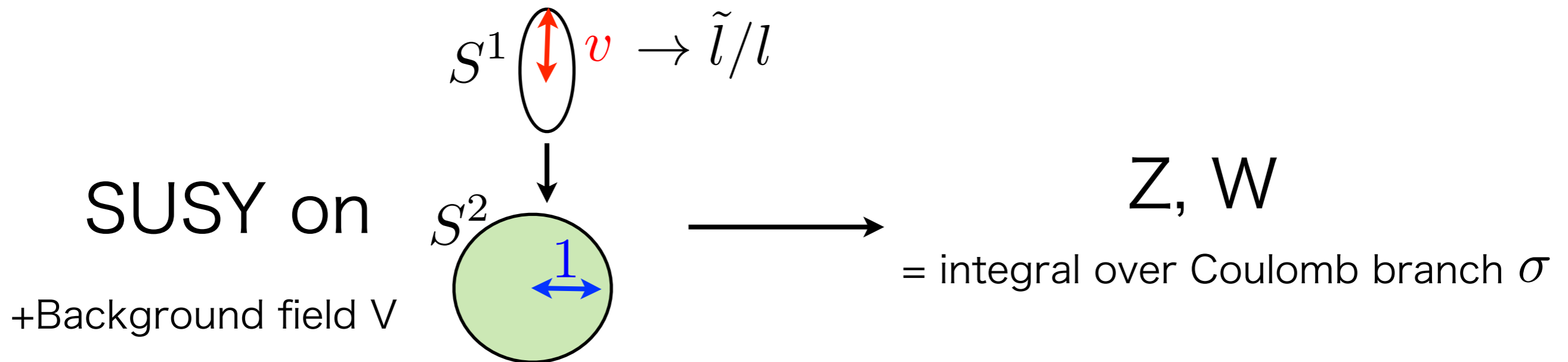


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History of Localization on 3-sphere

2011: Imamura, Yokoyama arXiv:1109.4734

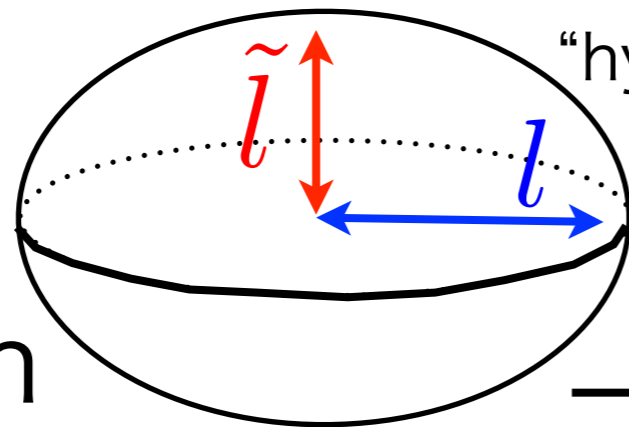


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History of Localization on 3-sphere

2011: Martelli, Passias, Sparks [arXiv:1100.6400](https://arxiv.org/abs/1100.6400)



“hyperbolic” squashed sphere

SUSY on

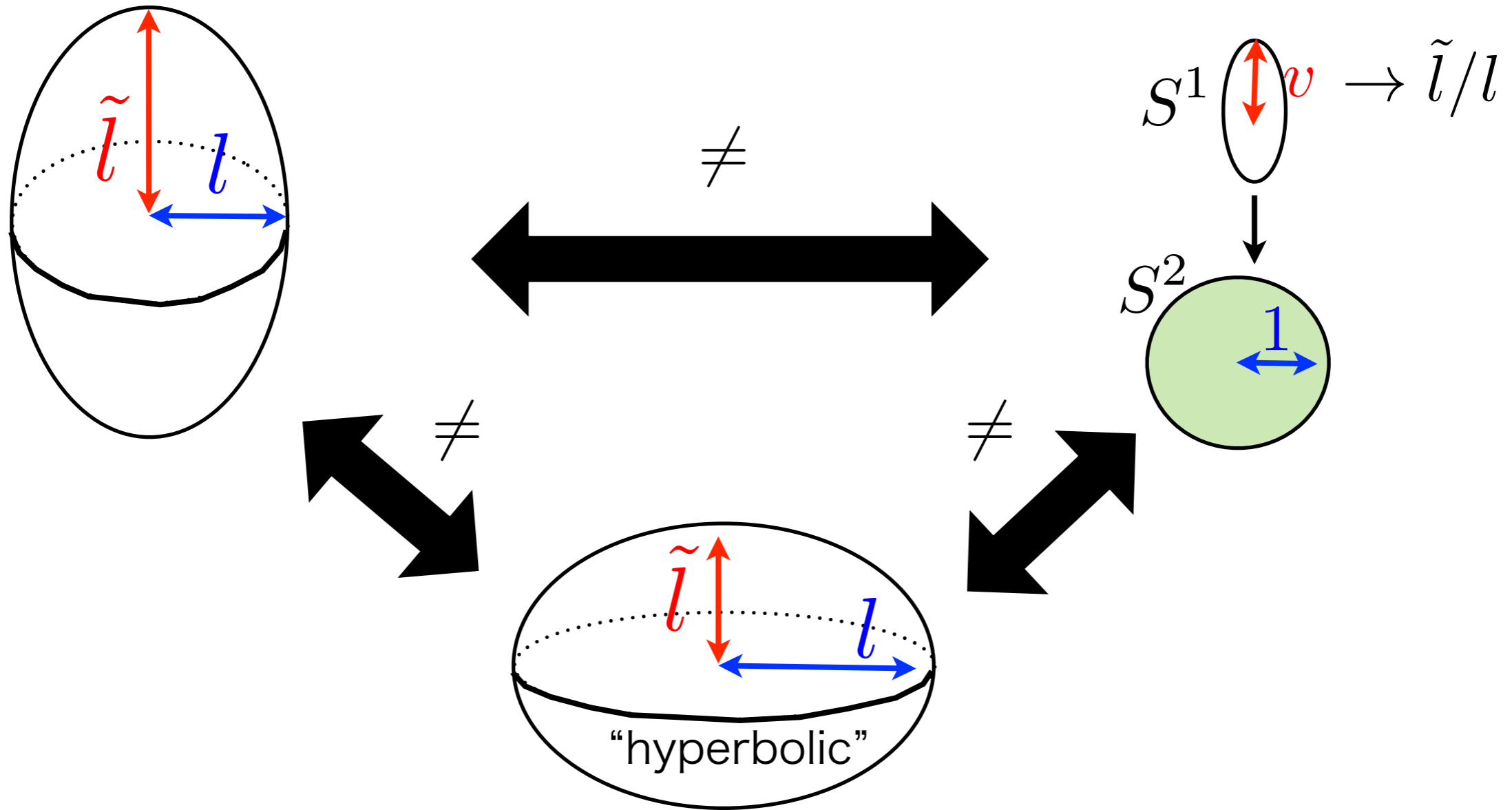
+Background field V

Z, W

= integral over Coulomb branch σ

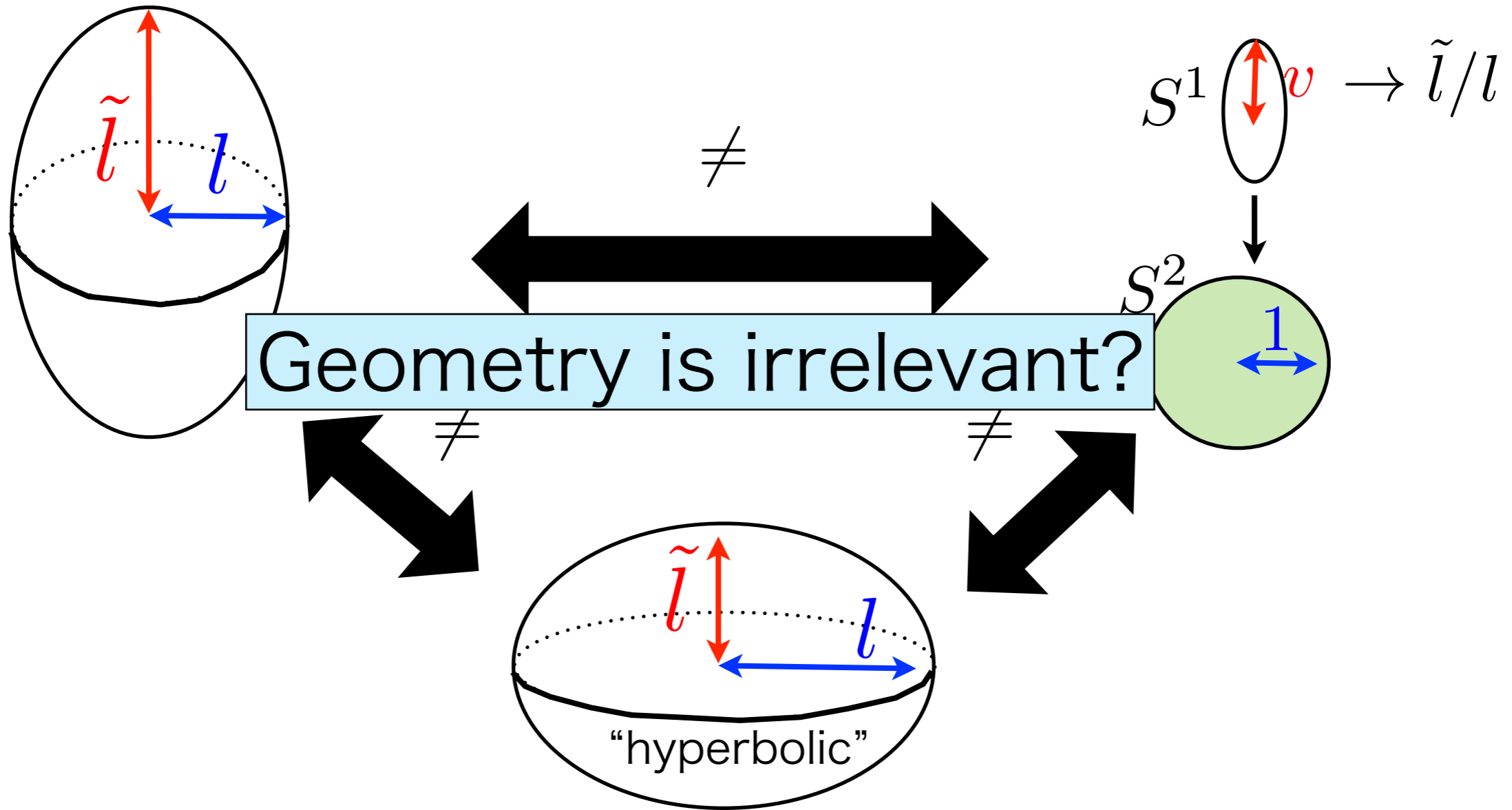
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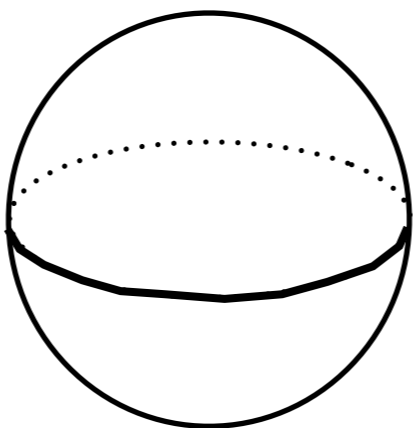
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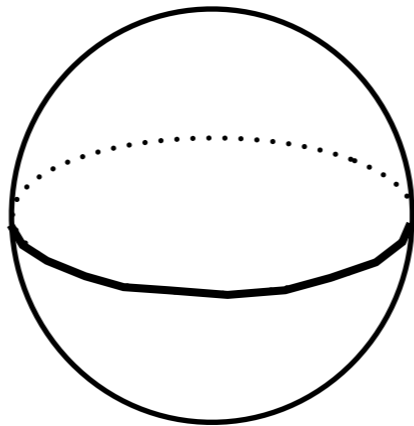
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+Background field V

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SUSY on round sphere

Applications

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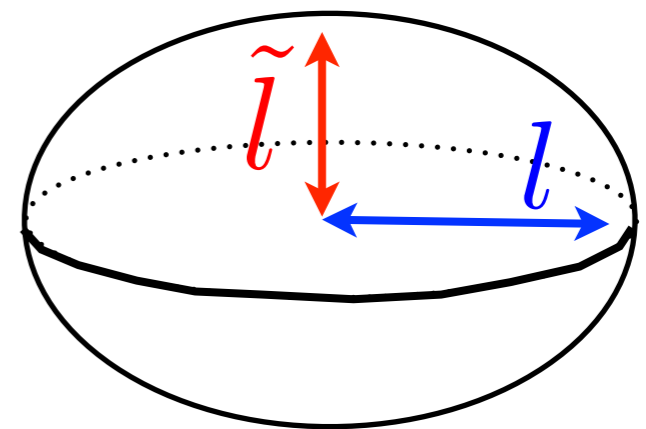
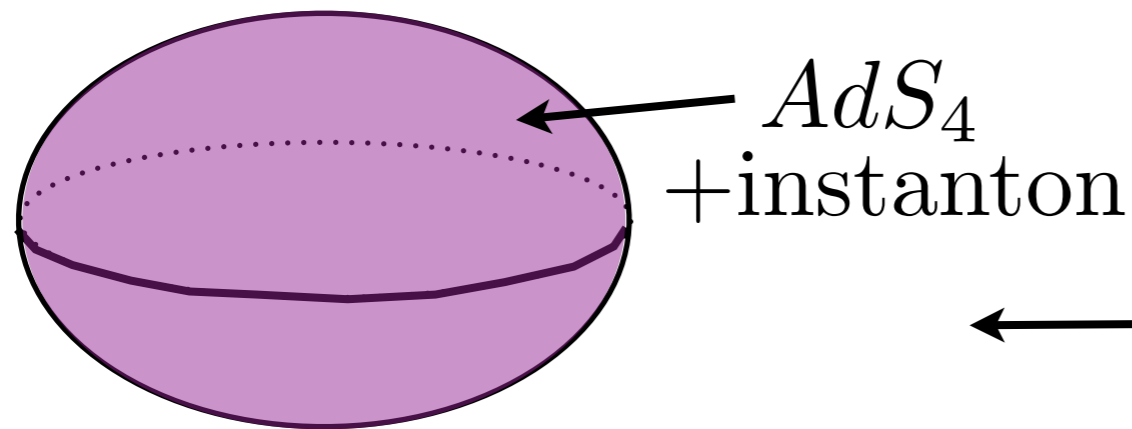
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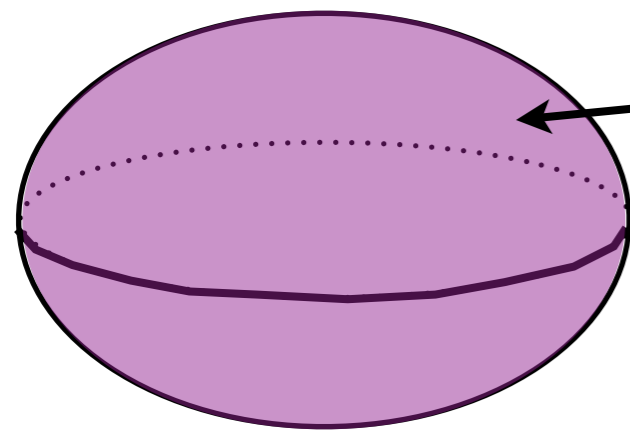
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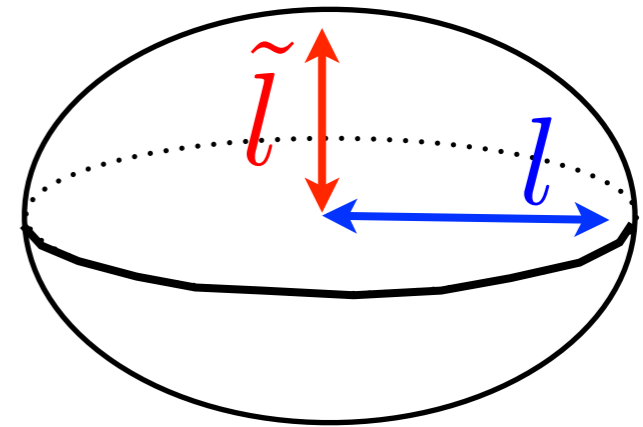
"Hyperbolic" squashed sphere
+Background U(1) vector
Killing spinors on 3D

4D Gravity

Field theory



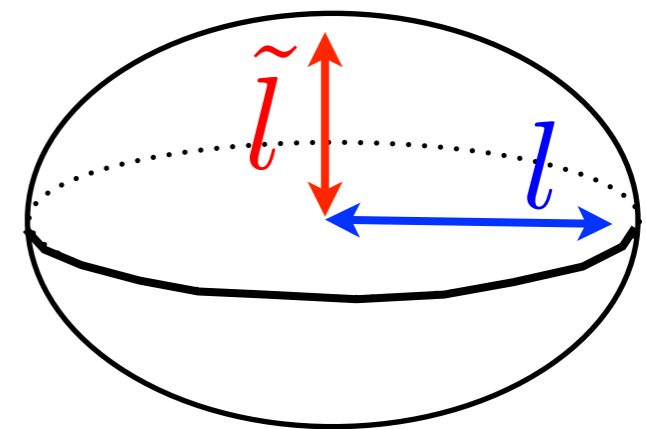
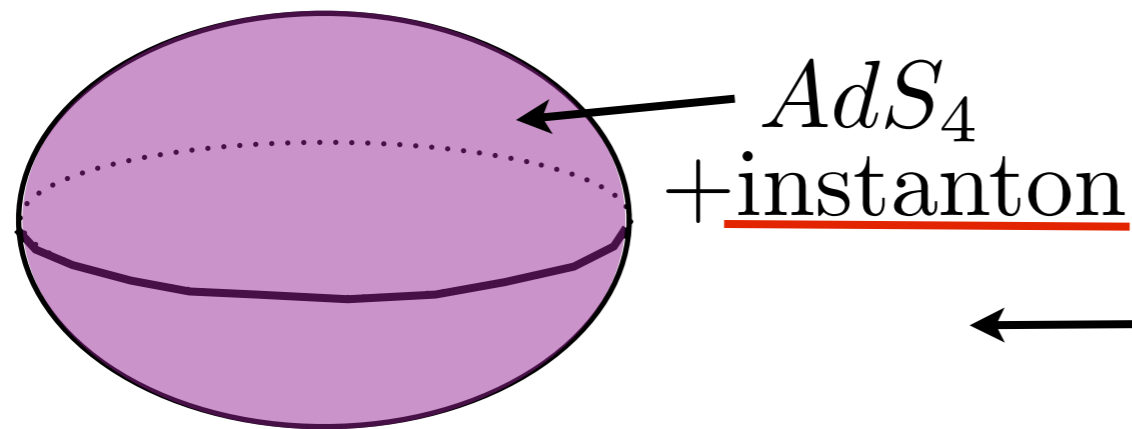
AdS_4
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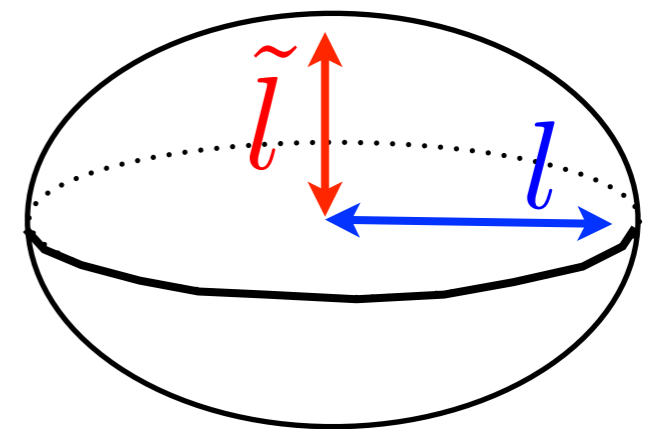
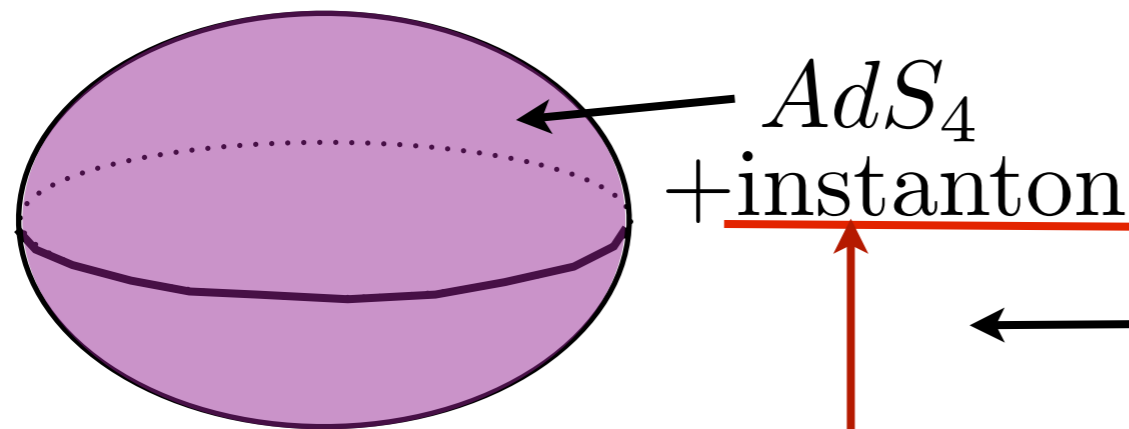
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Einstein-Maxwell type action

$$-\frac{1}{16\pi G_4} \int d^4x \sqrt{g} (R - 2\Lambda - F_{\mu\nu} F^{\mu\nu} + \text{gravitino})$$

4D Gravity

Field theory



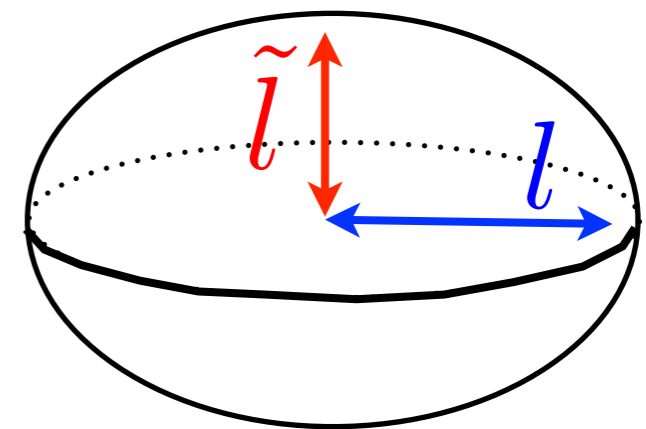
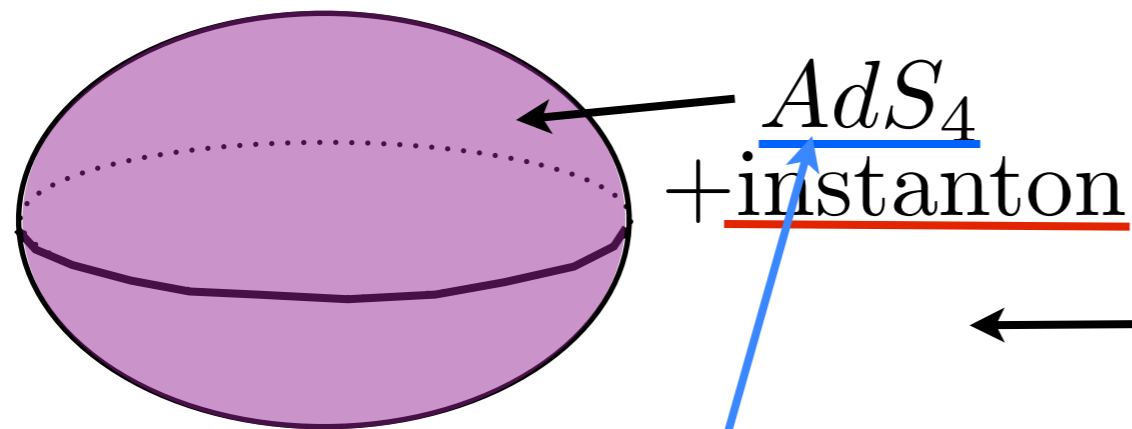
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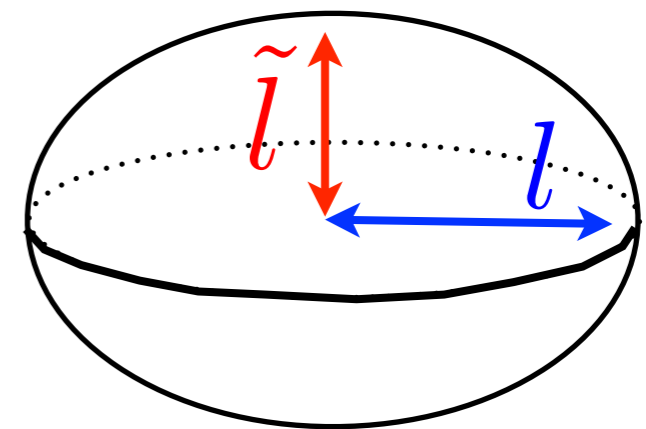
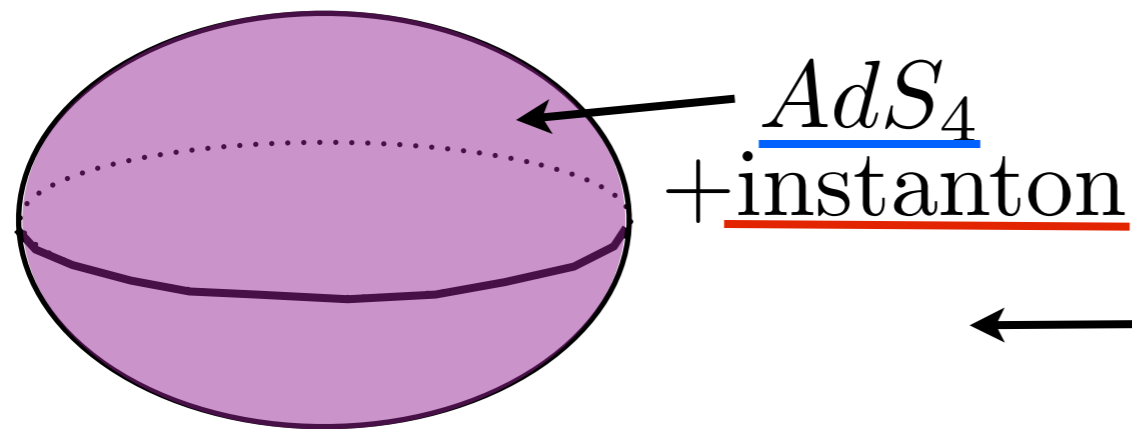
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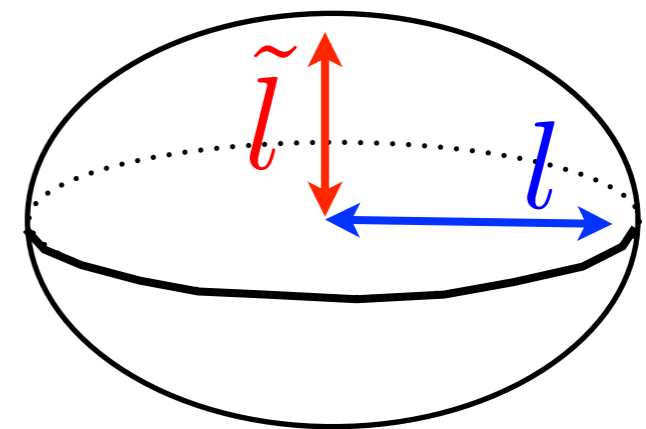
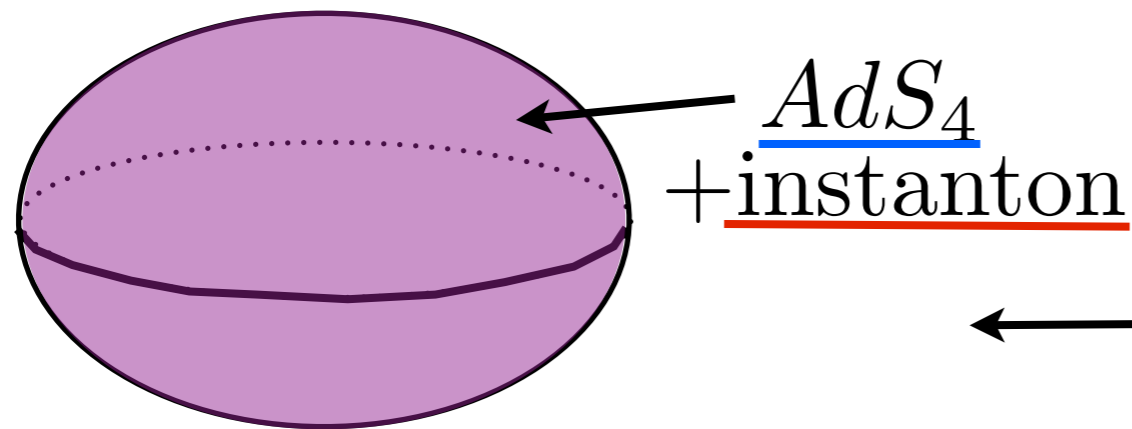
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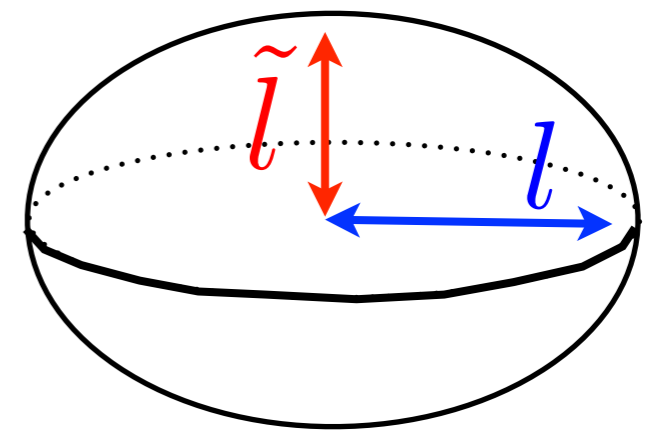
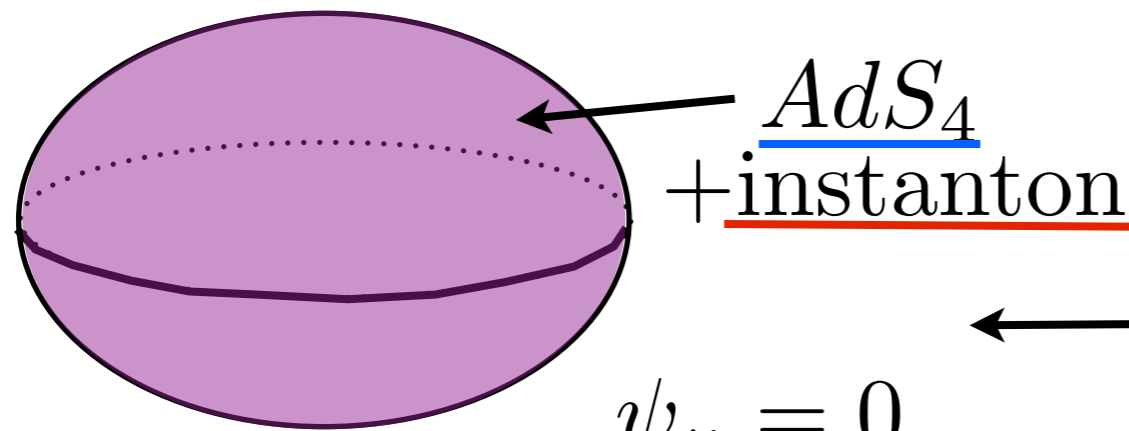
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4D Gravity

Field theory



$$\psi_\mu = 0$$

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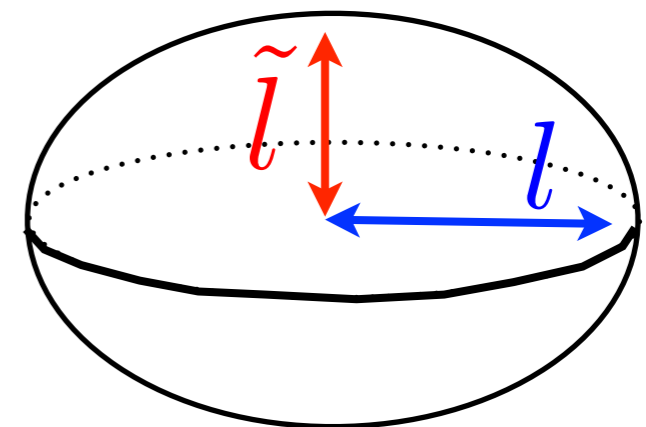
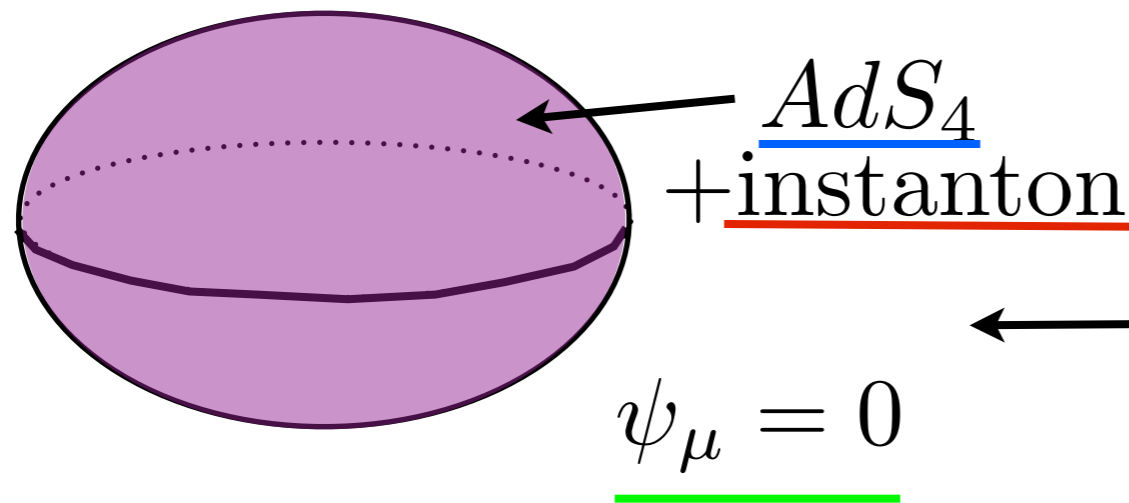


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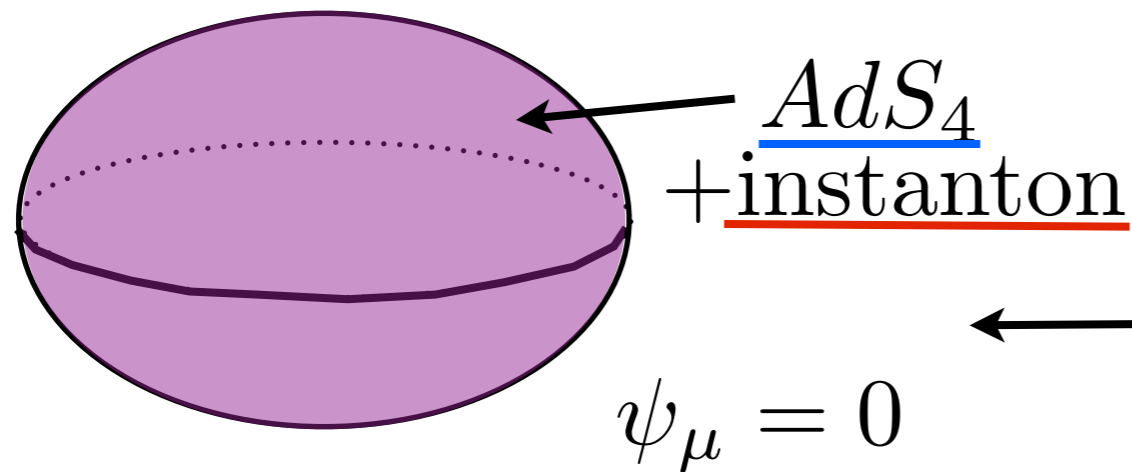


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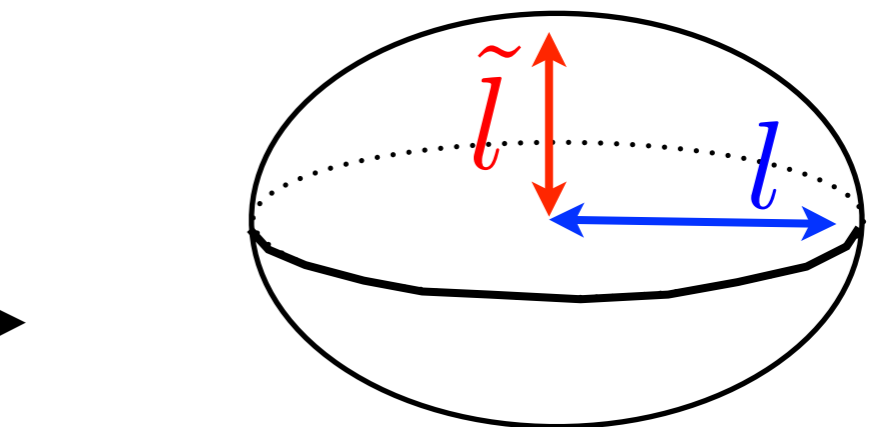
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4D Gravity

Field theory



BPS cond: $\delta\psi_\mu = 0$



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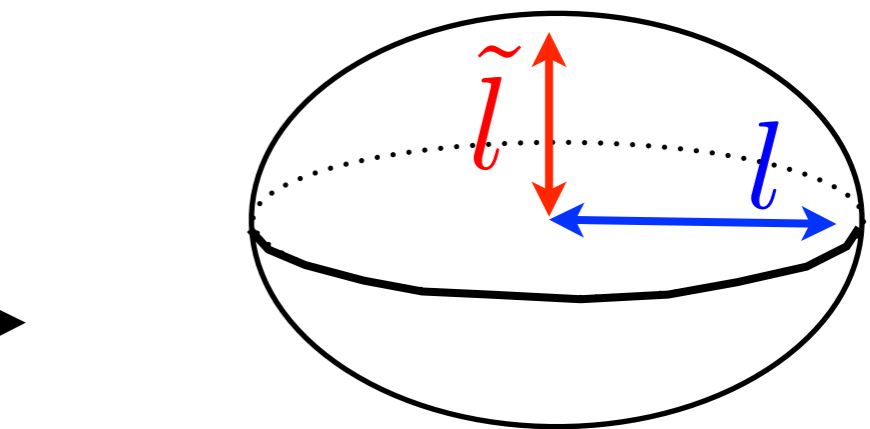
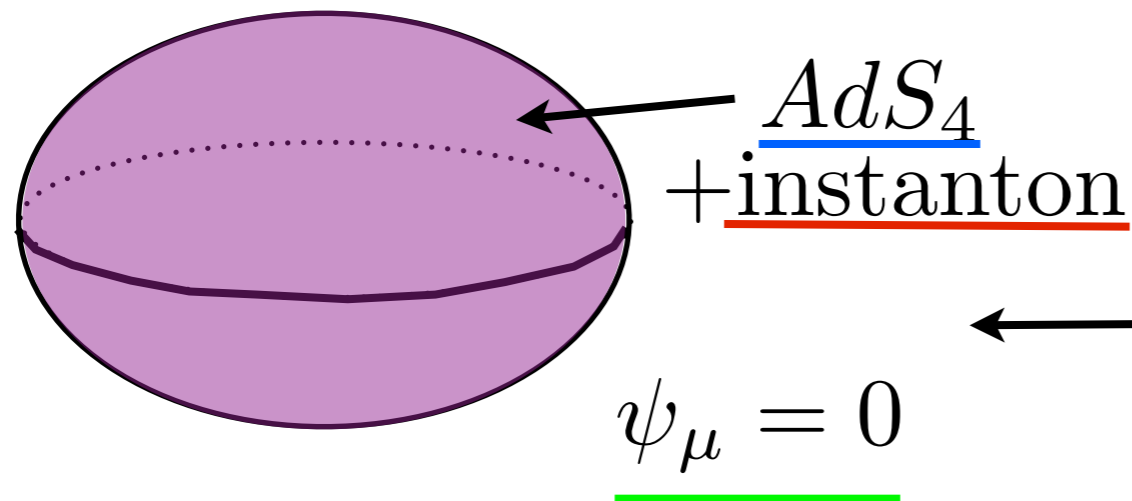


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4D Gravity

Field theory



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➔ Killing spinor on 4D

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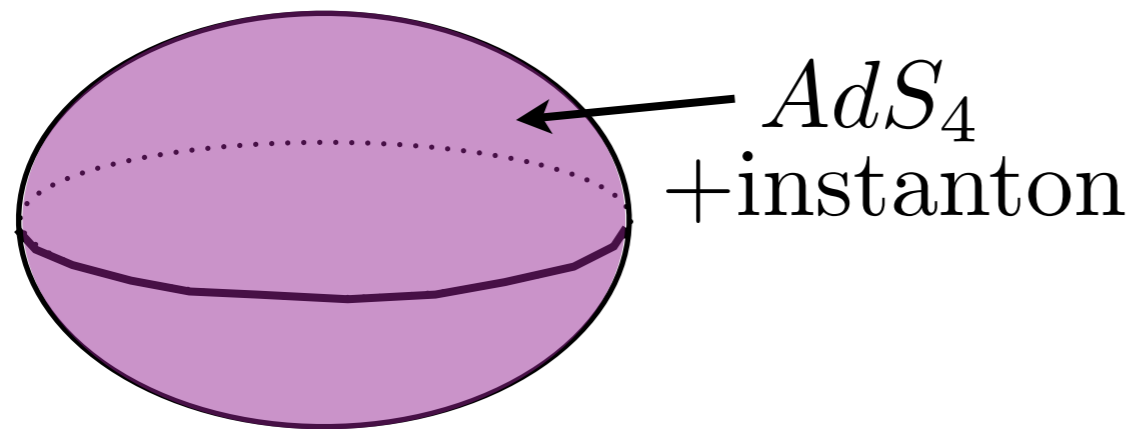
Killing spinors on 3D

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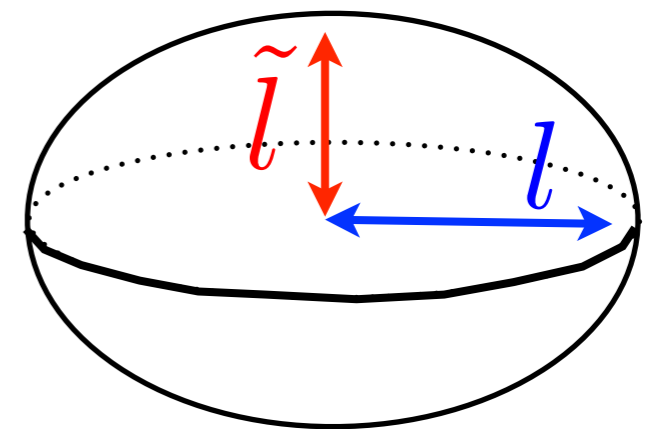
Field theory



BPS cond: $\delta\psi_\mu = 0$

➔ Killing spinor on 4D

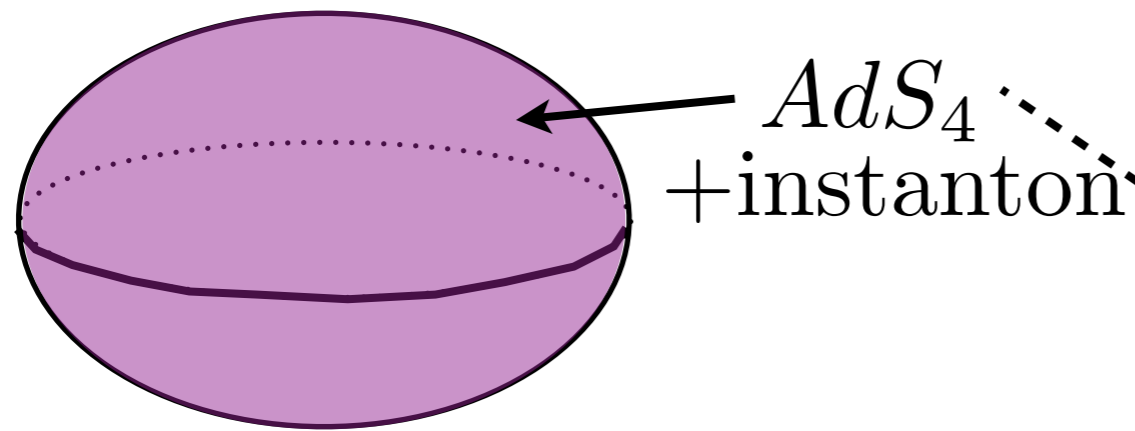
4D Gravity



“Hyperbolic” squashed sphere
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Killing spinors on 3D

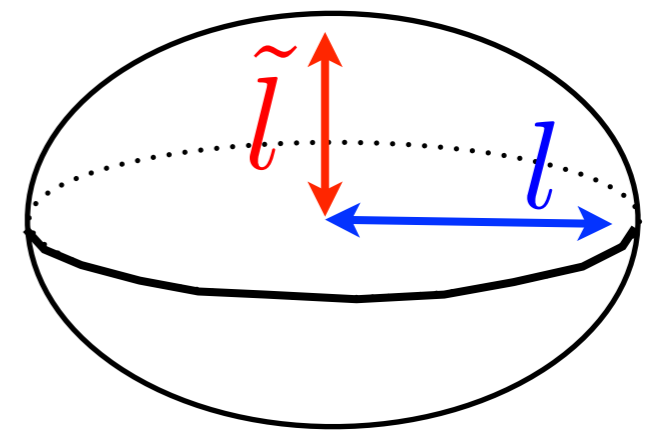
Field theory



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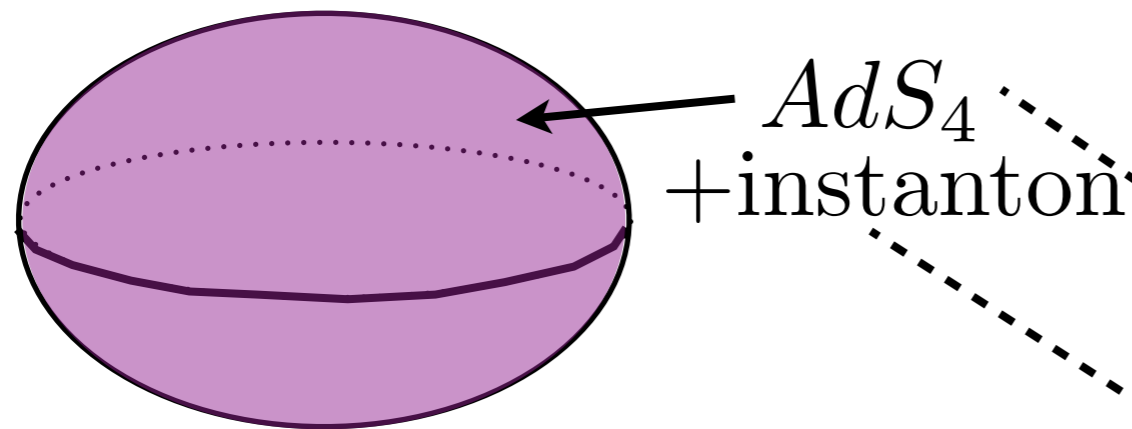
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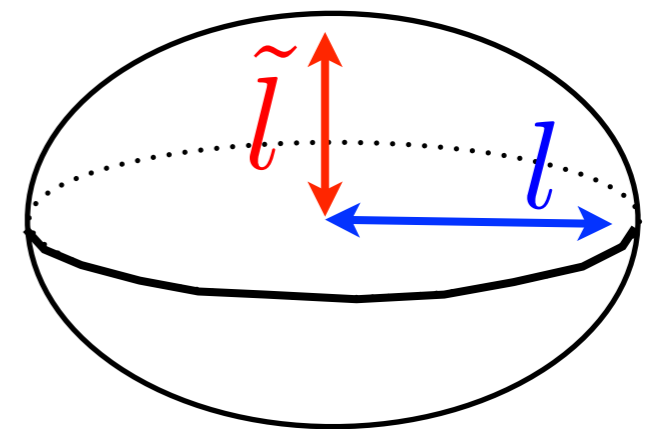
Killing spinors on 3D

Field theory



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➔ Killing spinor on 4D

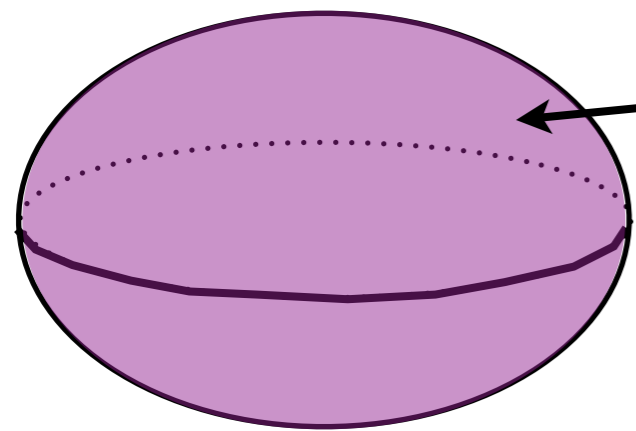


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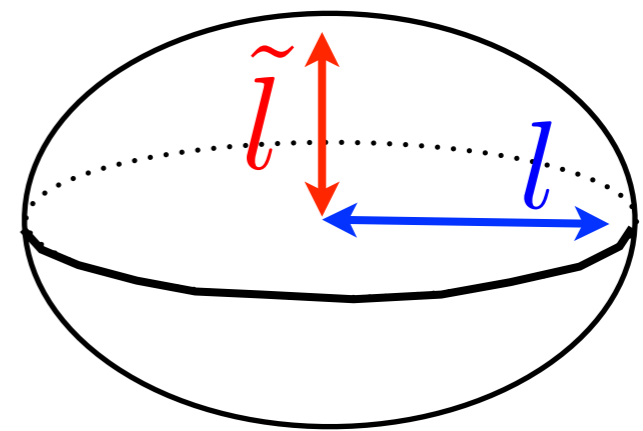
Killing spinors on 3D

4D Gravity

Field theory



AdS_4
+ instanton



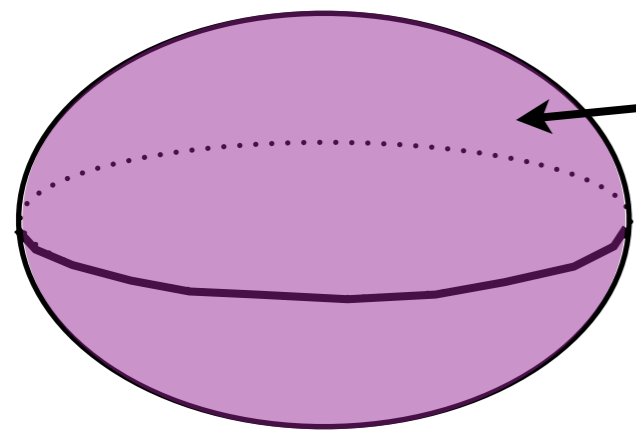
"Hyperbolic" squashed sphere
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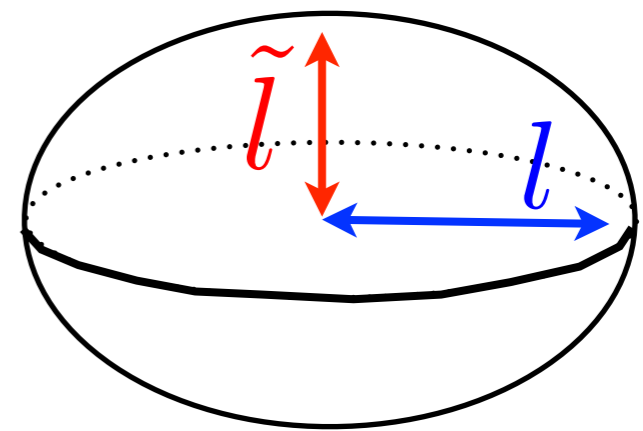
➔ Killing spinor on 4D \longleftrightarrow Killing spinors on 3D

4D Gravity

Field theory



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BPS cond: $\delta\psi_\mu = 0$

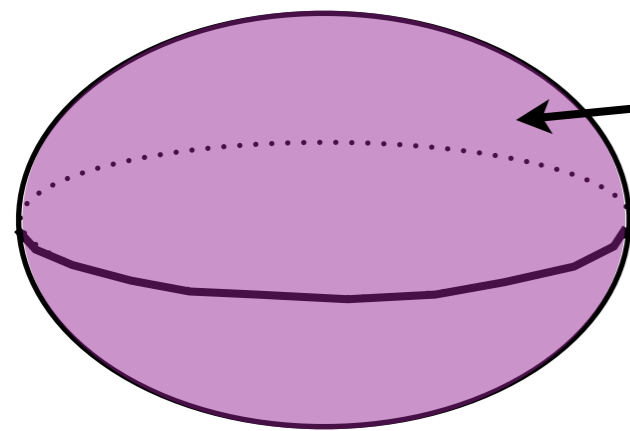
➔ Killing spinor on 4D ➔ Killing spinors on 3D

$\log Z_{grav}$

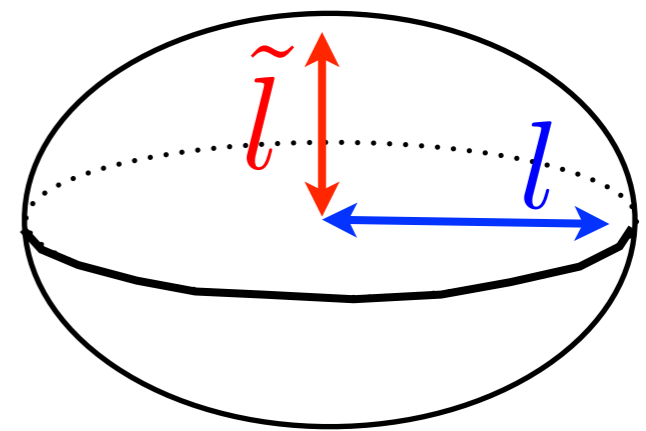
$\log Z_{field}$

4D Gravity

Field theory



AdS_4
+ instanton



“Hyperbolic” squashed sphere
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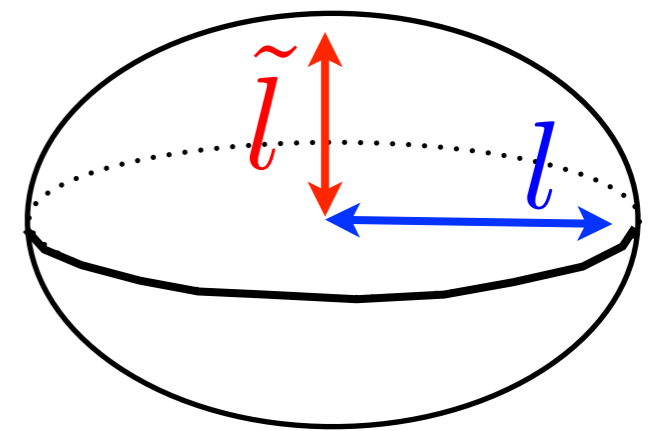
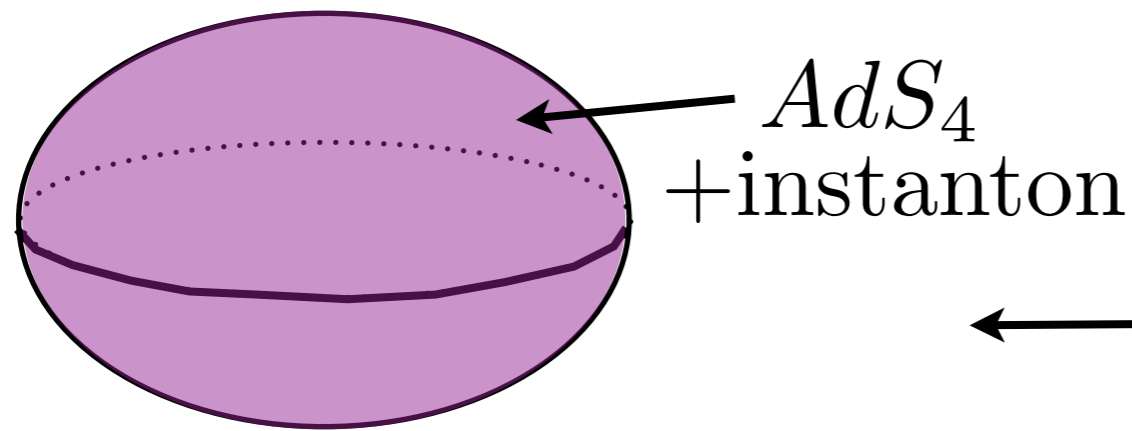
BPS cond: $\delta\psi_\mu = 0$

➔ Killing spinor on 4D \dots ➔ Killing spinors on 3D

$$\log Z_{grav} \xleftarrow{\text{Large N}} \log Z_{field}$$

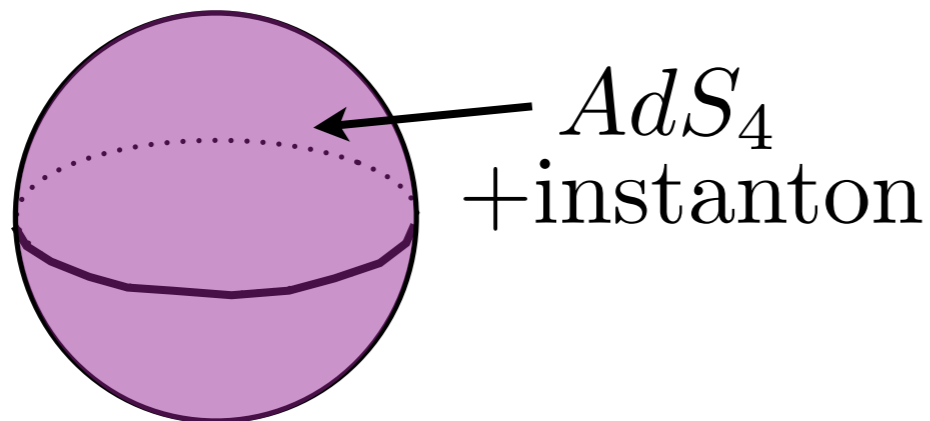
4D Gravity

Field theory



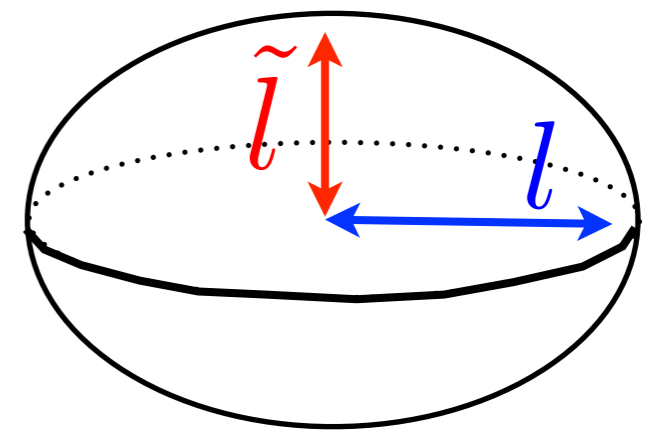
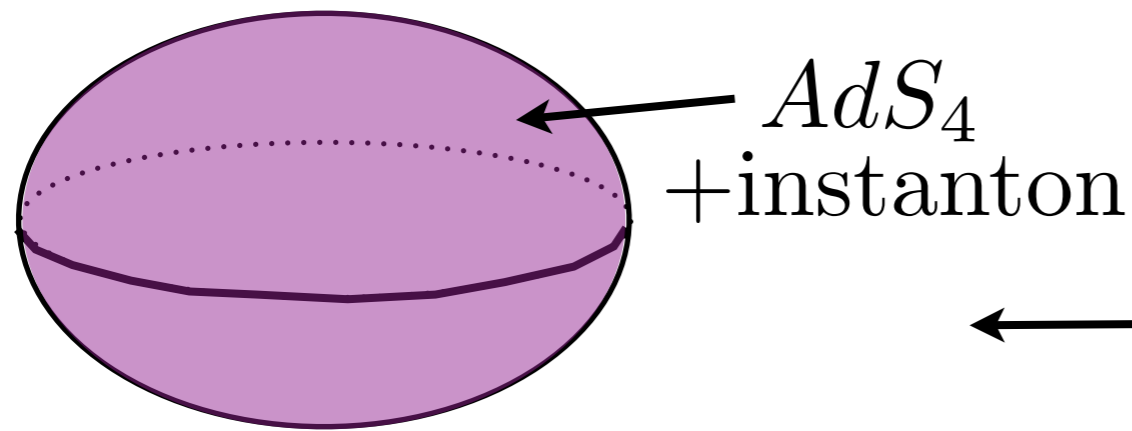
||
||
changing "slices"

"Hyperbolic" squashed sphere
+ Background U(1) vector
Killing spinors on 3D



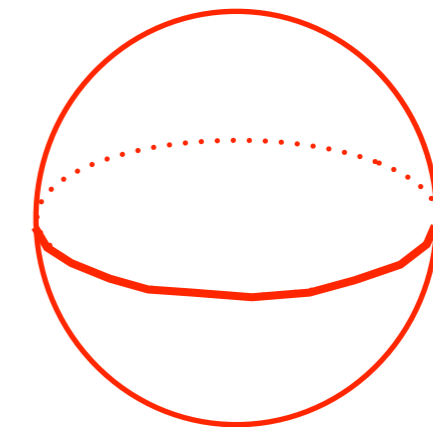
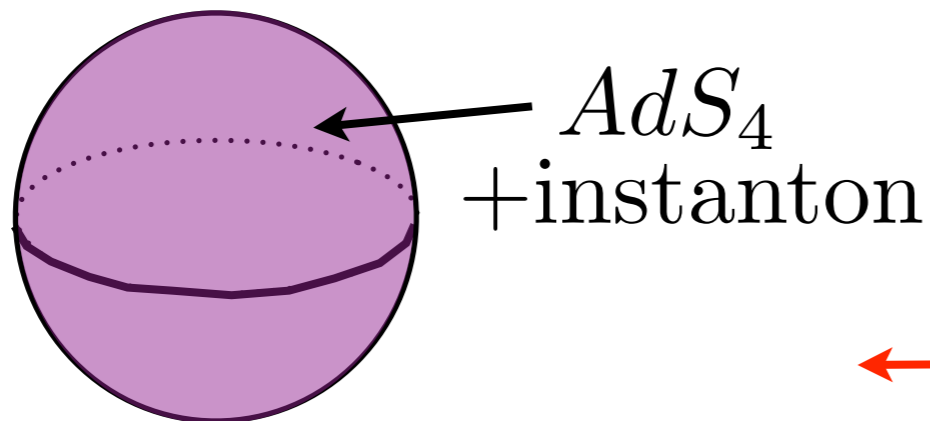
4D Gravity

Field theory



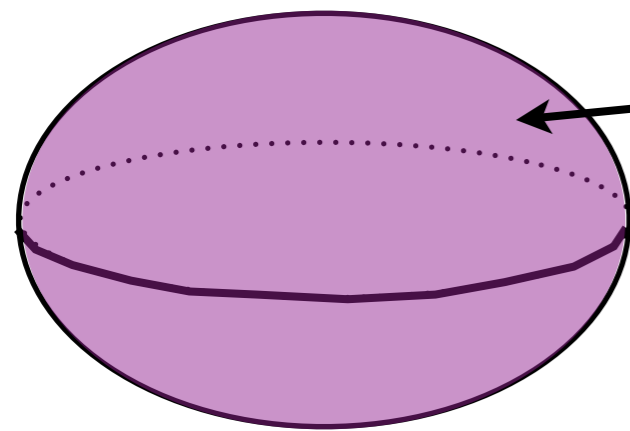
changing "slices"

"Hyperbolic" squashed sphere
+ Background U(1) vector
Killing spinors on 3D



4D Gravity

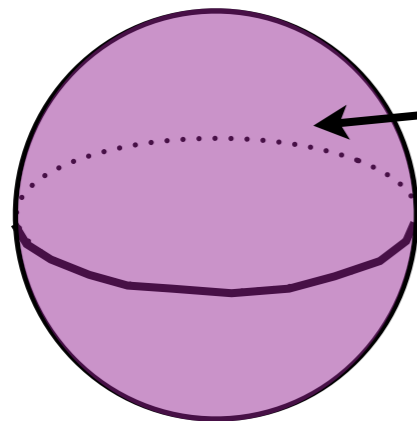
Field theory



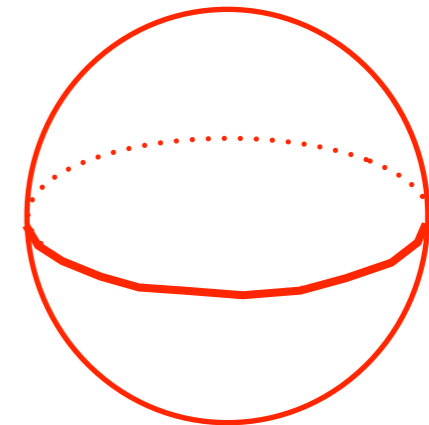
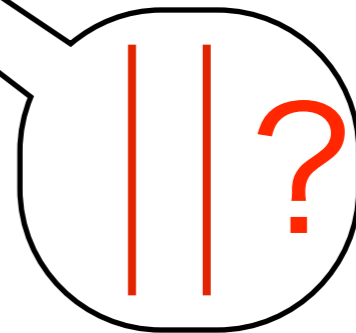
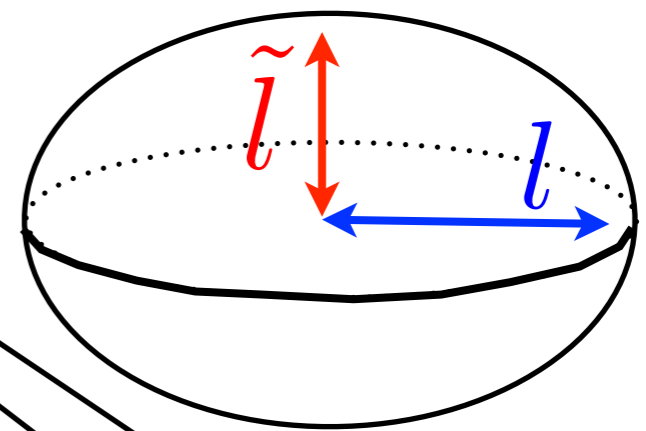
AdS_4
+ instanton



changing "slices"

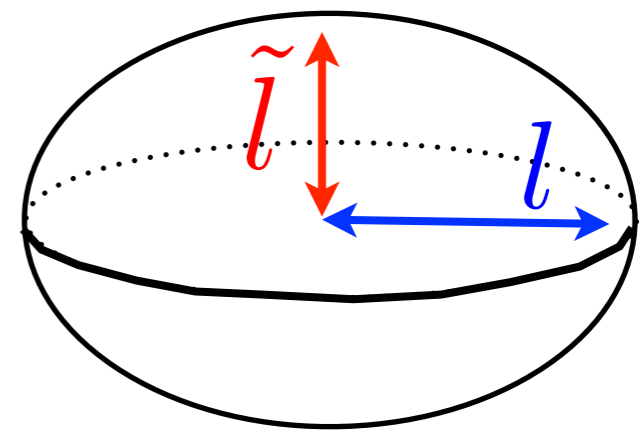
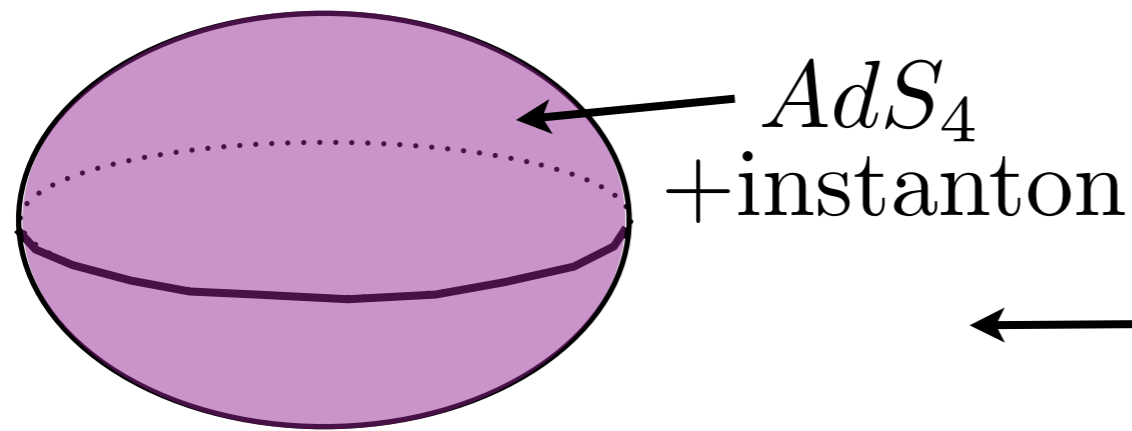


AdS_4
+ instanton



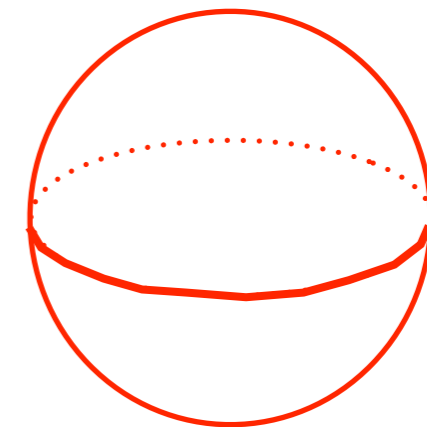
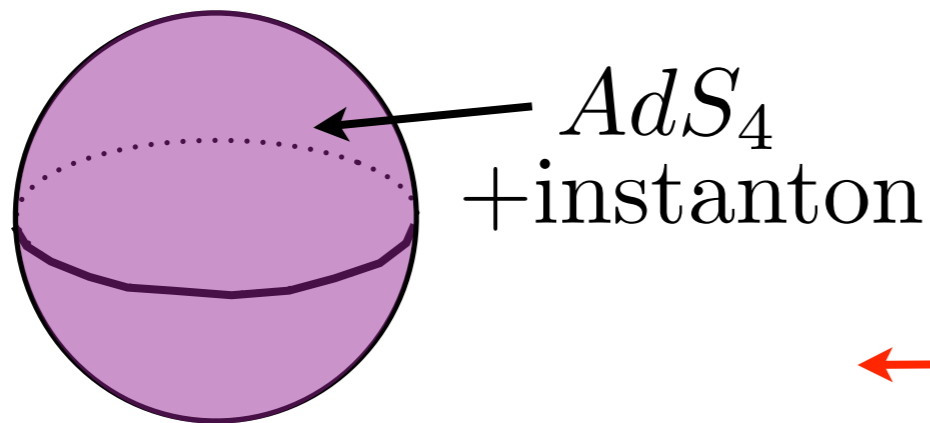
4D Gravity

Field theory



|| changing "slices" ||

Today's talk ||



4D Gravity

Field theory

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✓ History & Intro

✓ How to construct?

○ SUSY on round sphere

Applications

Metric

$$ds^2 = d\theta^2 + \cos^2 \theta d\phi^2 + \sin^2 \theta d\chi^2$$

Killing spinor

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

SUSY

$$\delta_\epsilon A_\mu = -\frac{i}{2} \bar{\lambda} \gamma_\mu \epsilon, \text{ etc}$$

$$\delta_\epsilon \phi = 0, \text{ etc}$$

Squashed sphere's case

Metric

$$ds^2 = d\theta^2 + \cos^2 \theta d\phi^2 + \sin^2 \theta d\chi^2$$

$$\leftrightarrow ds^2 = f^2(\theta) d\theta^2 + l^2 \cos^2 \theta d\phi^2 + \tilde{l}^2 \sin^2 \theta d\chi^2$$

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Action $\left(f^2(\theta) = \sin^2 \theta + \frac{1}{s^2} \cos^2 \theta\right)$

$$\mathcal{L}_{\text{CS}} = \text{Tr} \left[\frac{1}{\sqrt{g}} \varepsilon^{\mu\nu\lambda} (A_\mu \partial_\nu A_\lambda - \frac{2i}{3} A_\mu A_\nu A_\lambda) - \bar{\lambda} \lambda + 2D\sigma \right].$$

$$\begin{aligned} \mathcal{L}_{\text{YM}} = (2sf) \text{Tr} & \left(\frac{1}{4} F^{\mu\nu} F_{\mu\nu} + \frac{1}{2} \mathcal{D}^i \sigma \mathcal{D}_i \sigma + \frac{1}{2} (\mathcal{D}_\theta \sigma + \frac{f'}{f} \sigma)^2 + \frac{1}{2} (D + \frac{1}{sf^2} \sigma)^2 \right. \\ & \left. + \frac{i}{2} \bar{\lambda} \gamma^i \mathcal{D}_i \lambda + \frac{i}{2} \bar{\lambda} \gamma^\theta (\mathcal{D}_\theta \lambda + \frac{1}{2} \frac{f'}{f} \lambda) + \frac{i}{2} \bar{\lambda} [\sigma, \lambda] - \frac{1}{4} \frac{1}{sf^2} \bar{\lambda} \lambda \right) \end{aligned}$$

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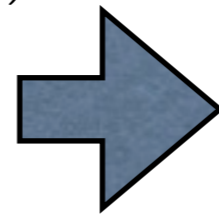
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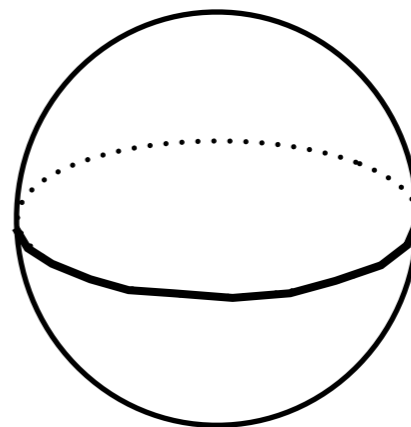
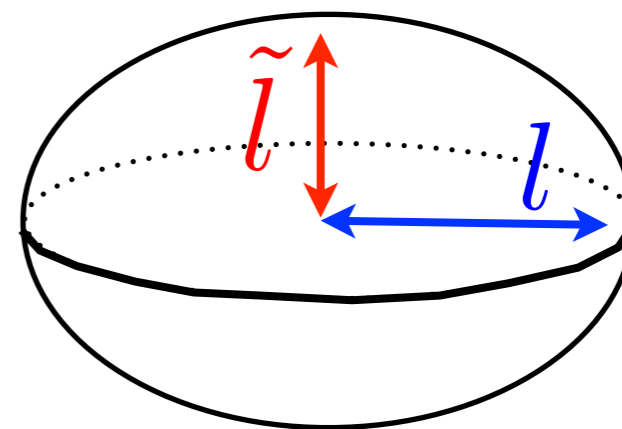
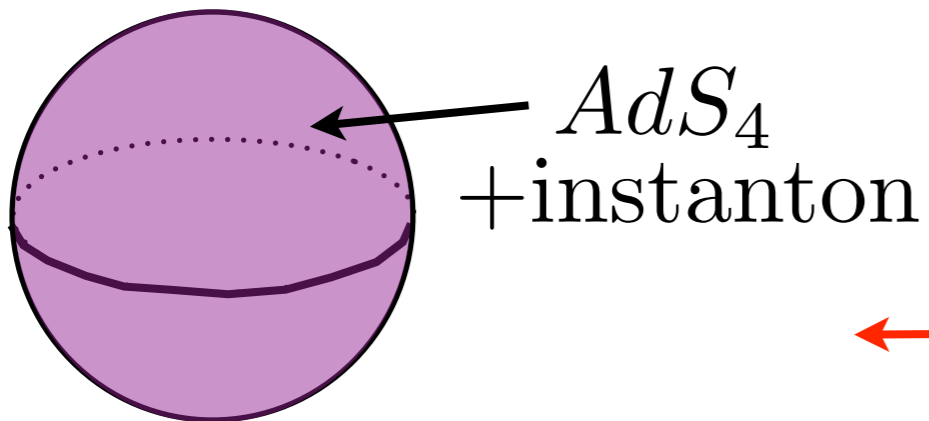
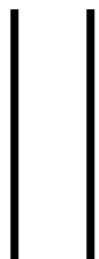
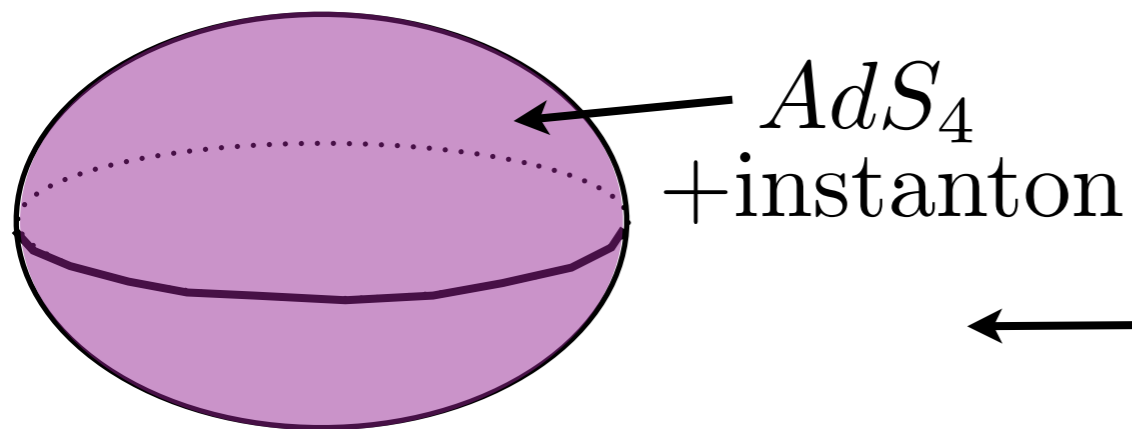
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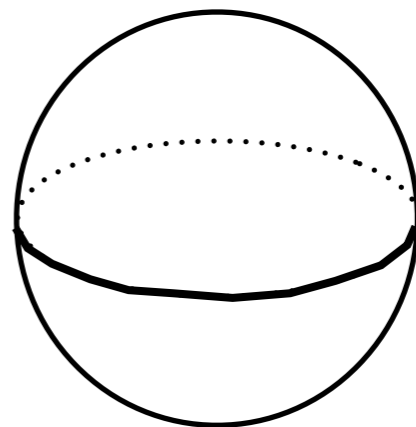
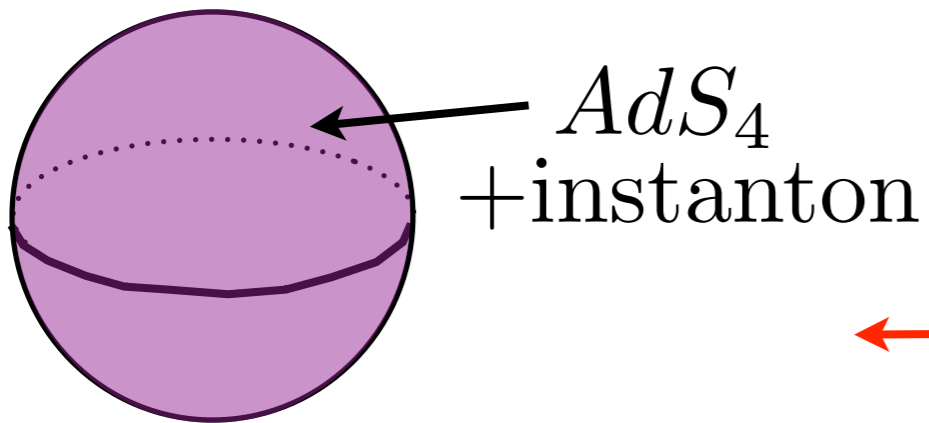
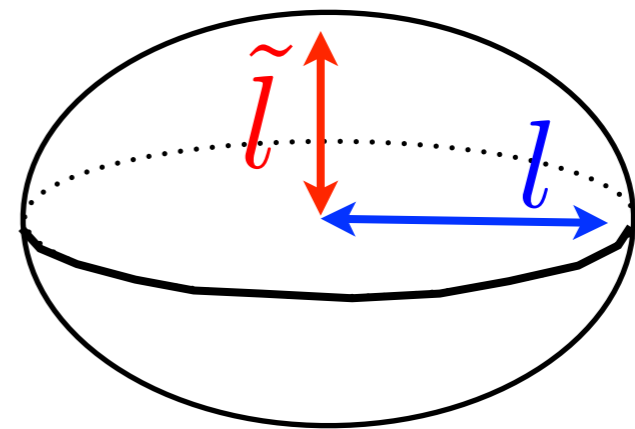
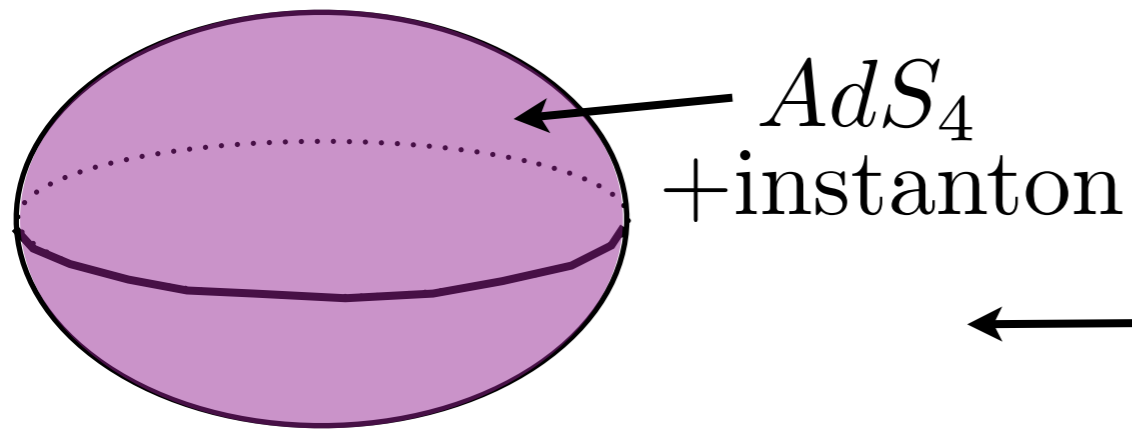


Double sine
(depends on s)



4D Gravity

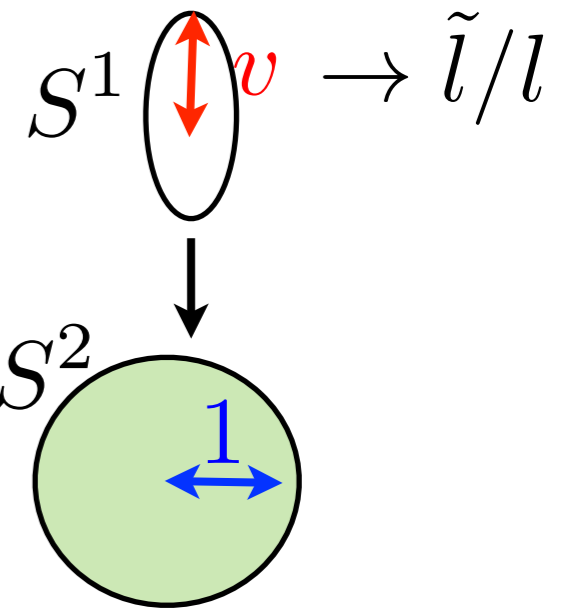
Field theory



4D Gravity

Field theory

Related works:



Nian arXiv:1309.3266

← solving a puzzle for S^2

Alday, Martelli, Richmond, Sparks arXiv:1307.6848

← emergence of double sine function
with a broader class of metrics

Closset, Dumitrescu, Festuccia, Komargodski

arXiv:1309.5876

← (geometrical) parameter dependence
for partition function in 4D and 3D

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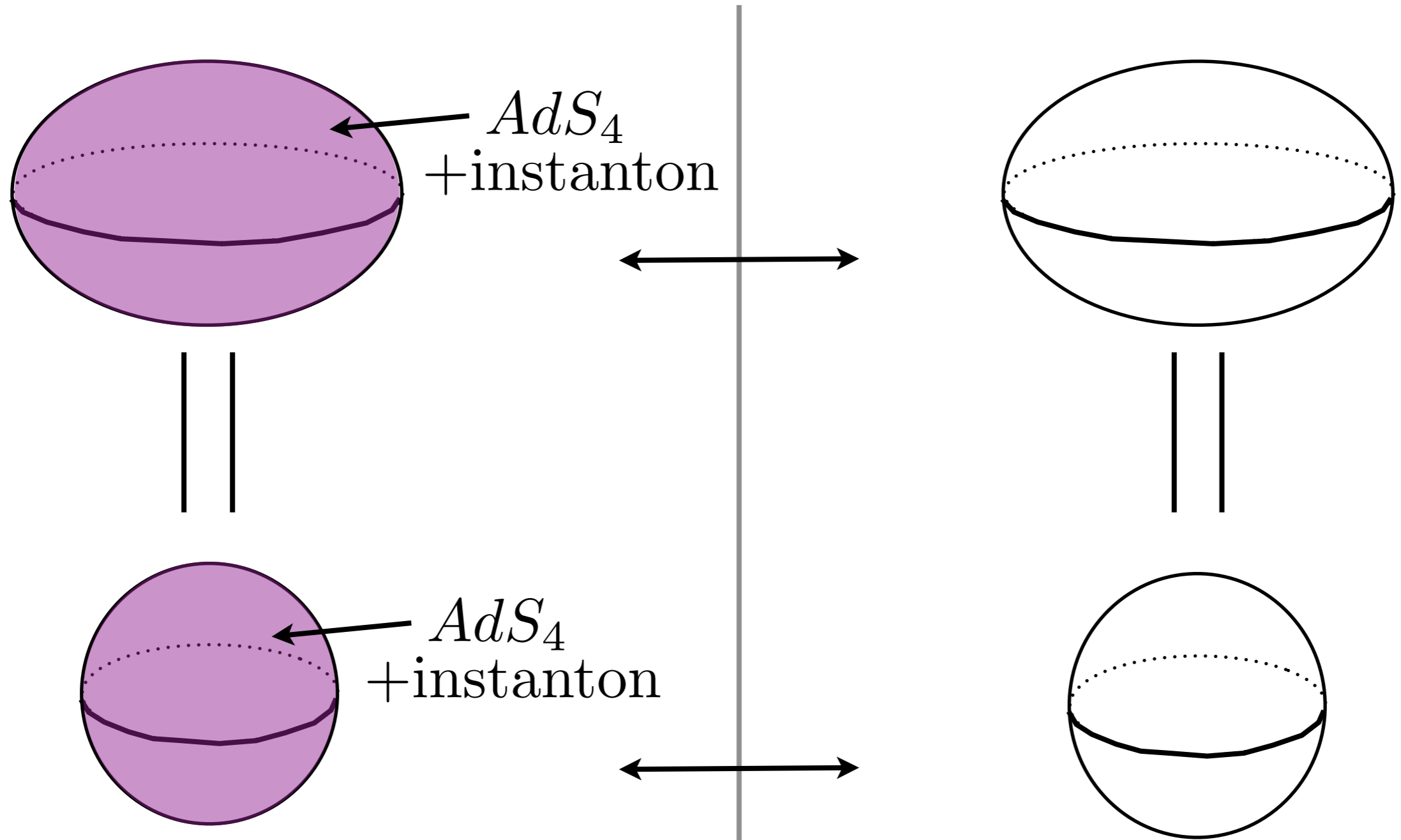
✓ History & Intro

✓ How to construct?

✓ SUSY on round sphere

○ Applications

1. Wilson loop

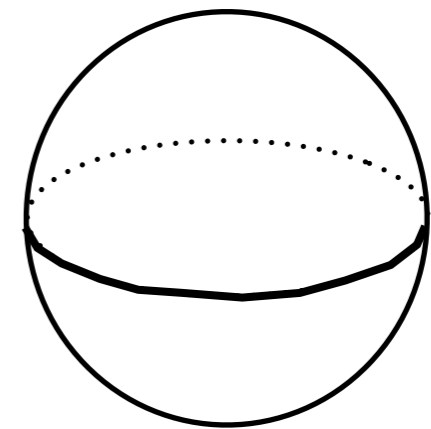
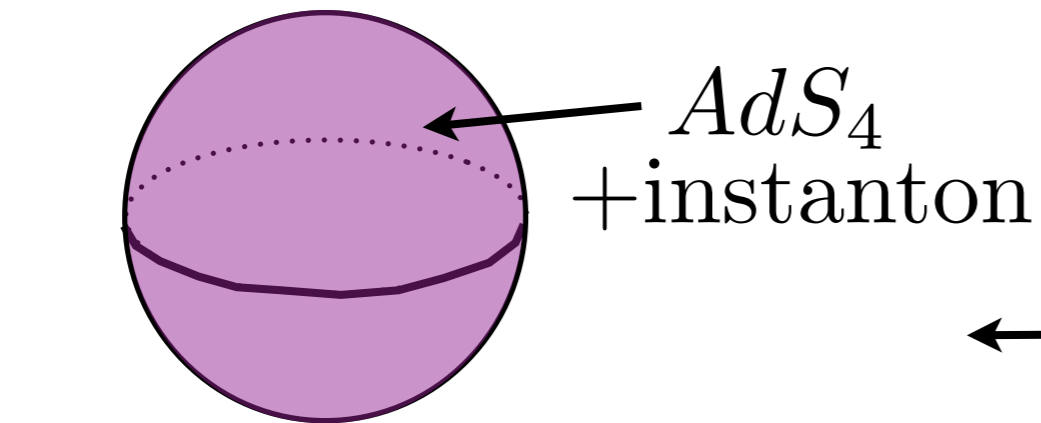
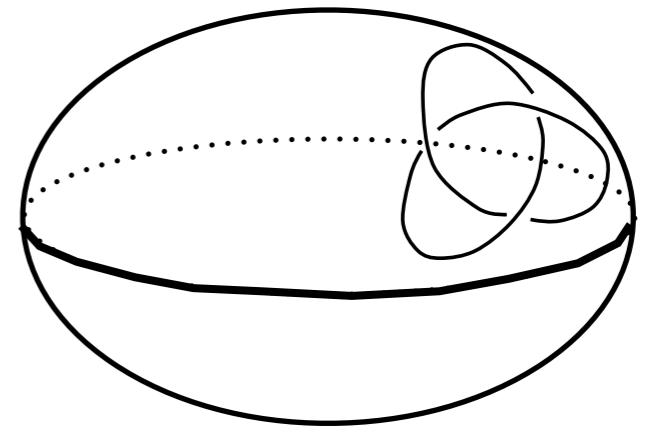
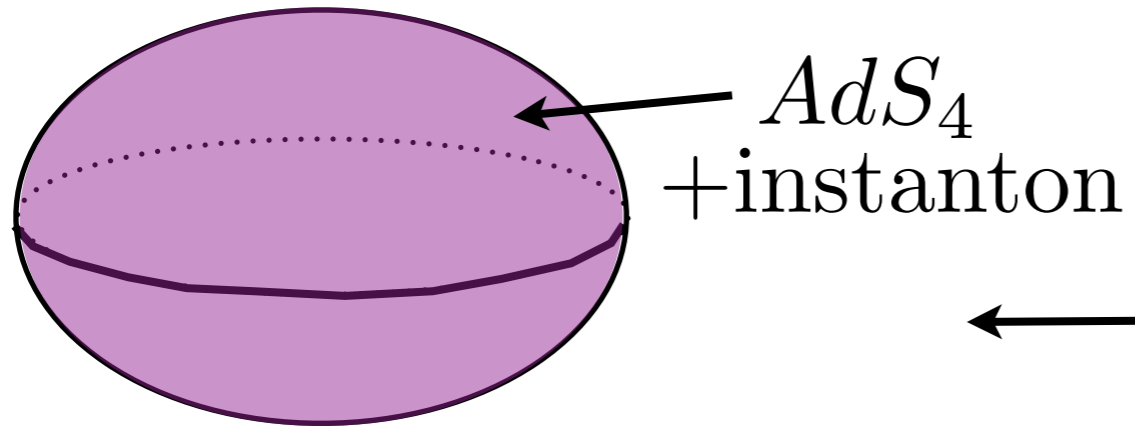


4D Gravity

Field theory

1. Wilson loop

1/2 BPS



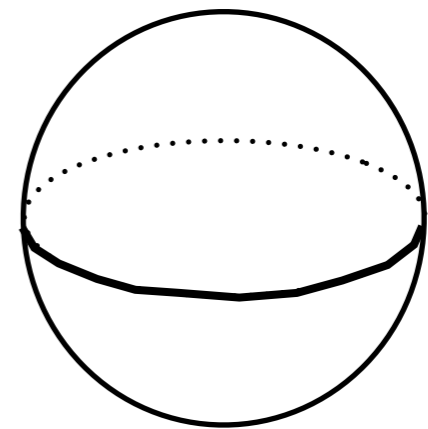
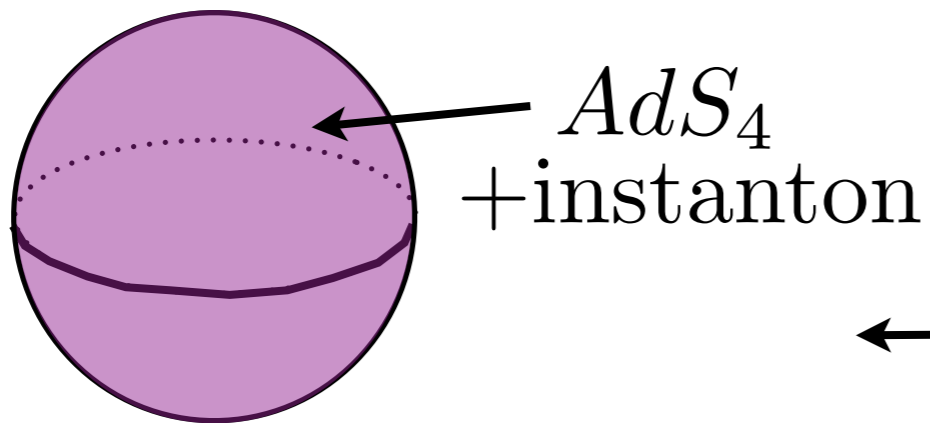
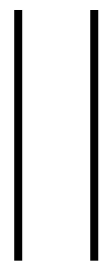
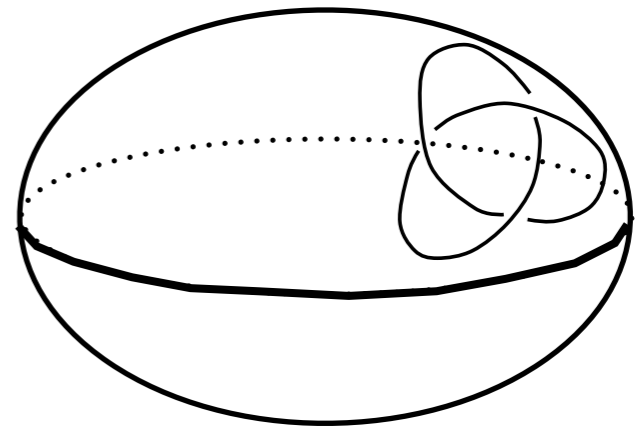
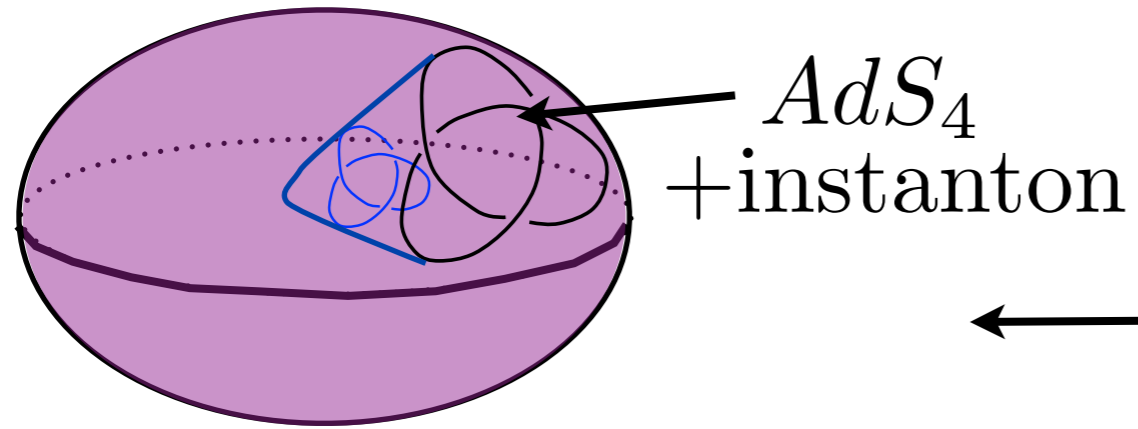
4D Gravity

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knotted surface?

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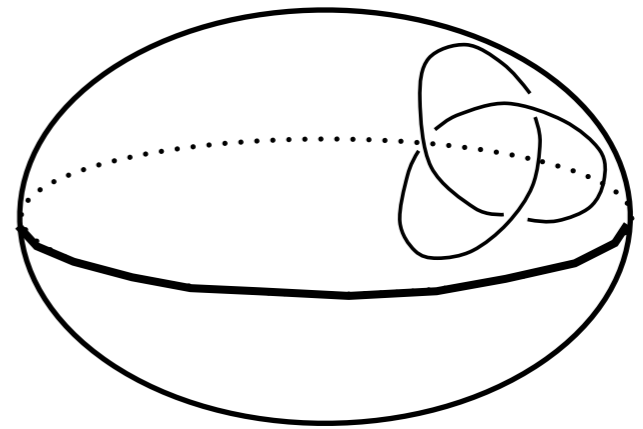
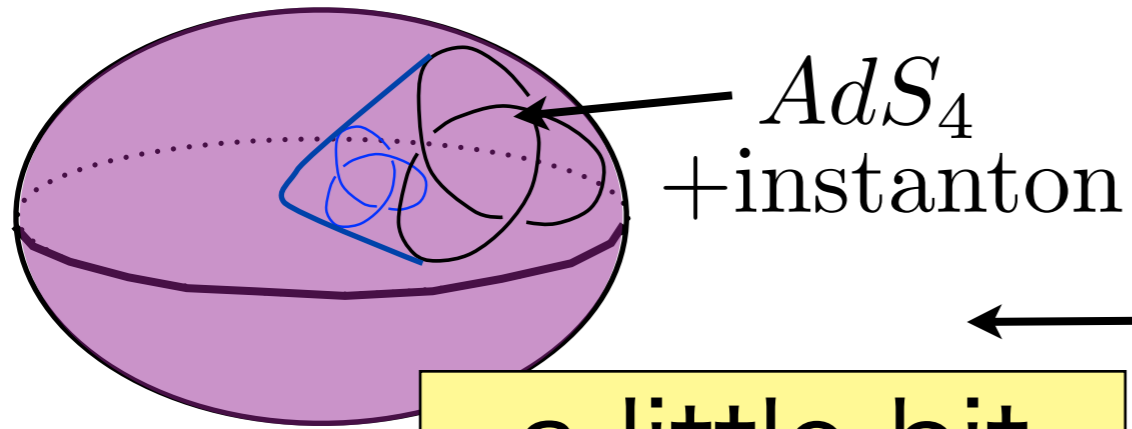
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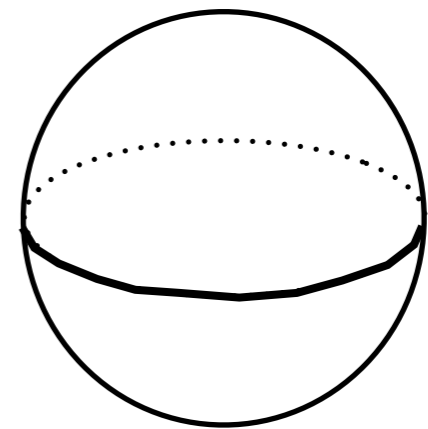
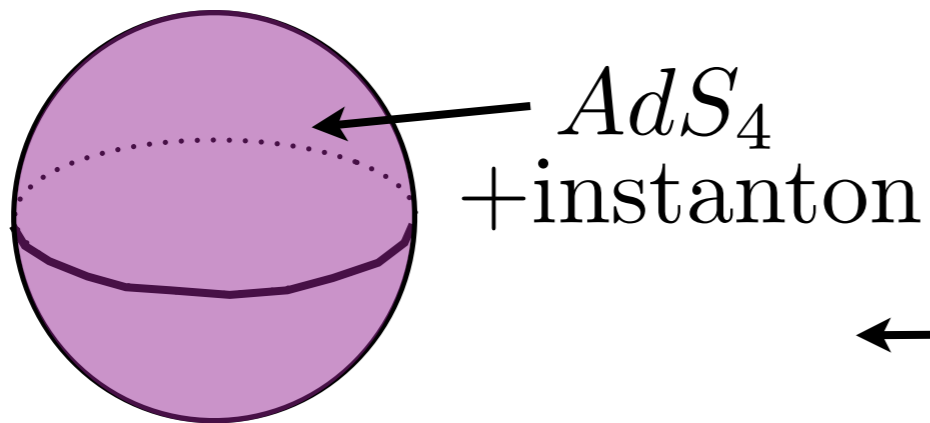
1. Wilson loop

knotted surface?

1/2 BPS



a little bit complicated

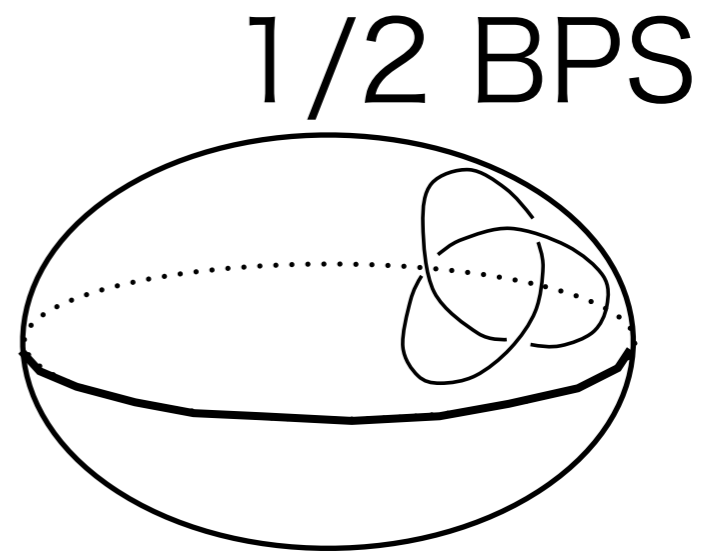
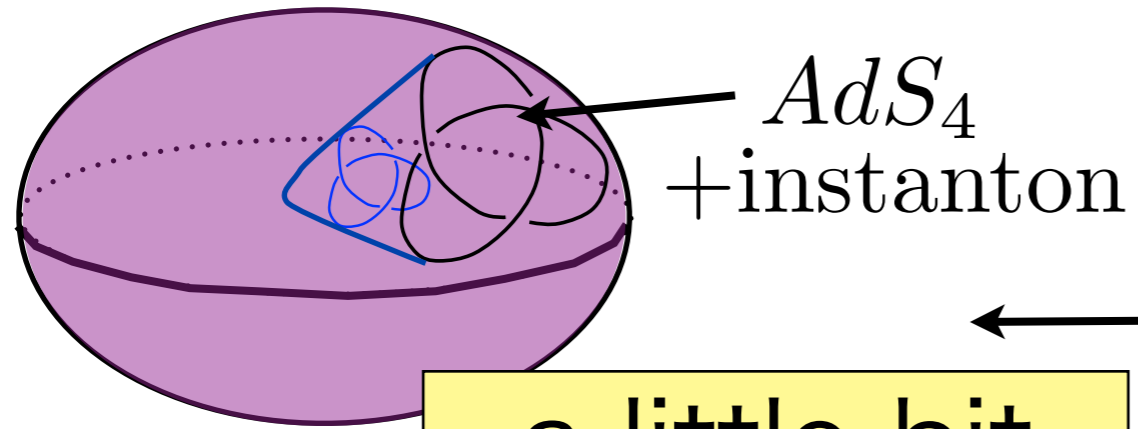


4D Gravity

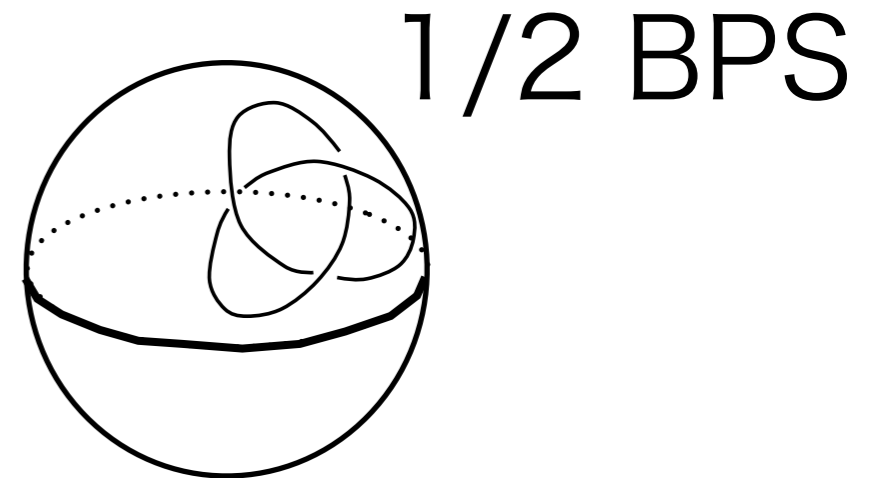
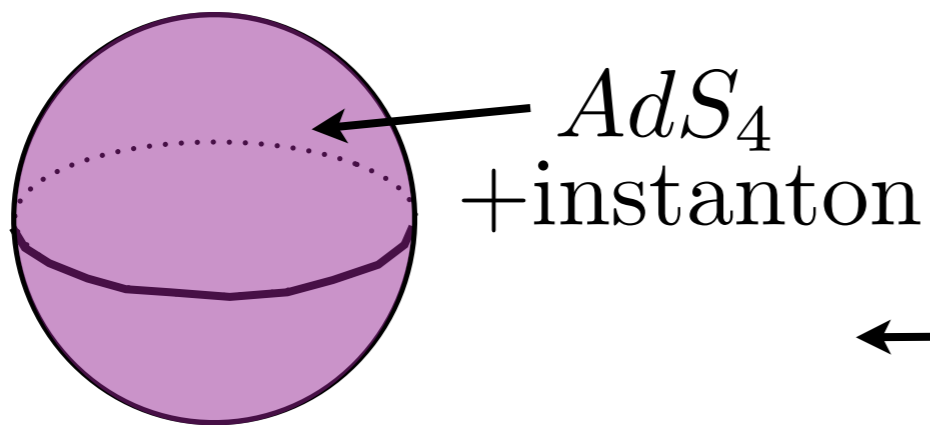
Field theory

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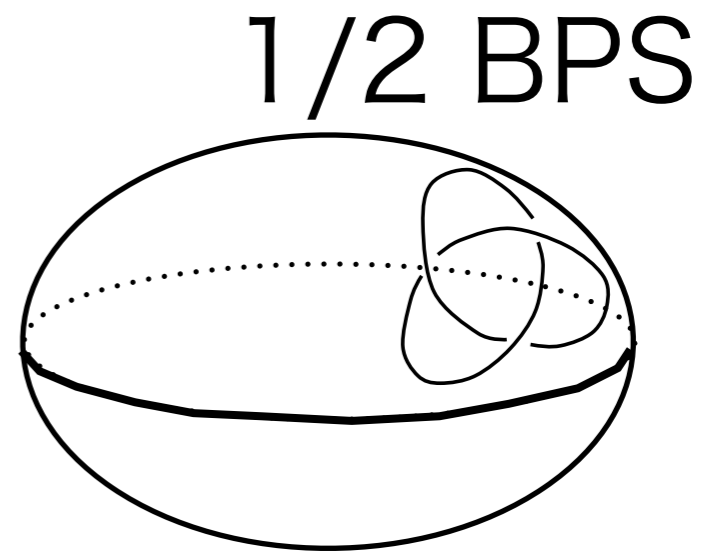
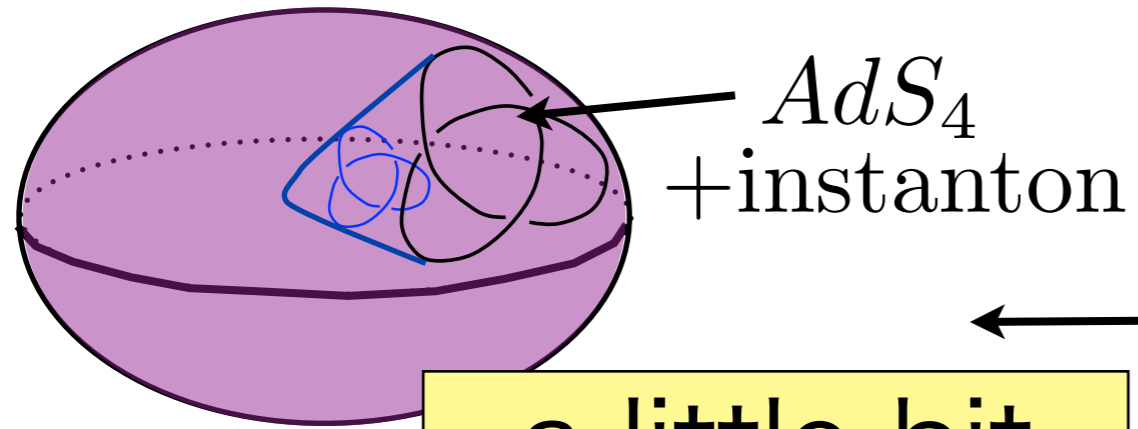


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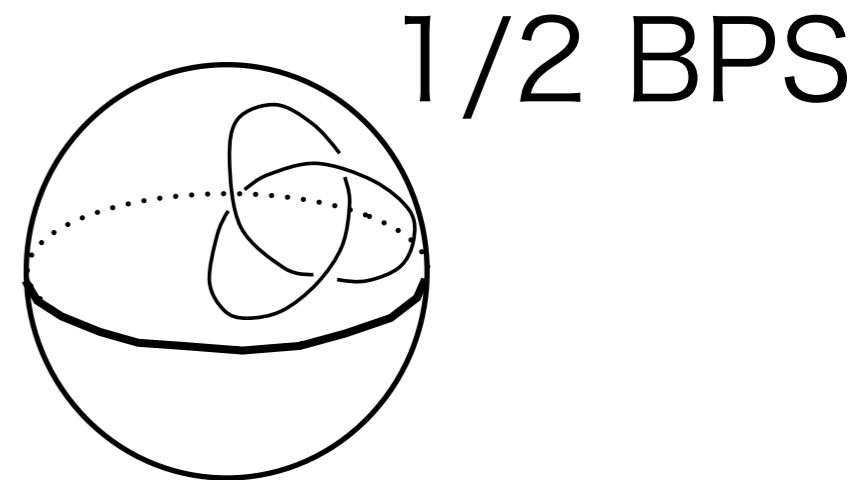
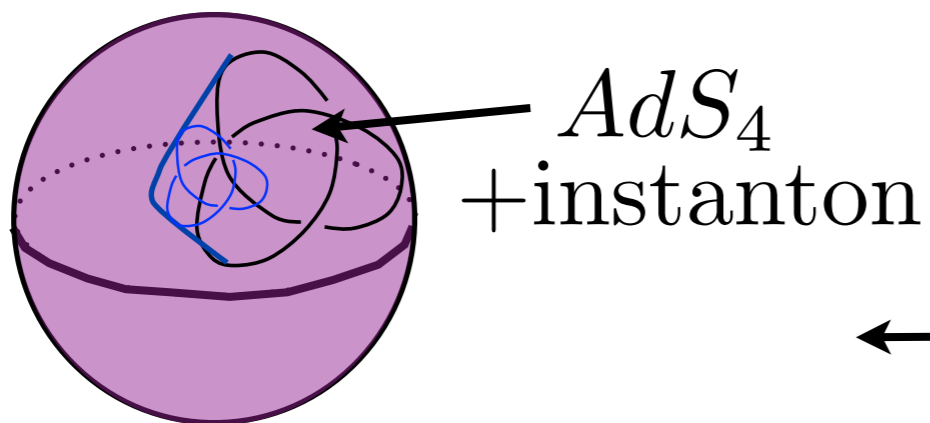
Field theory

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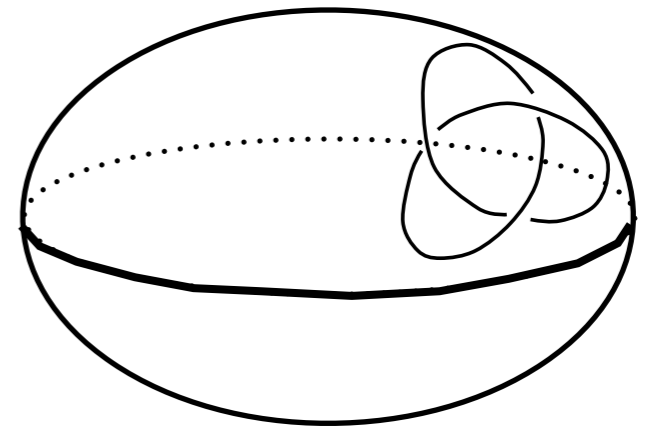
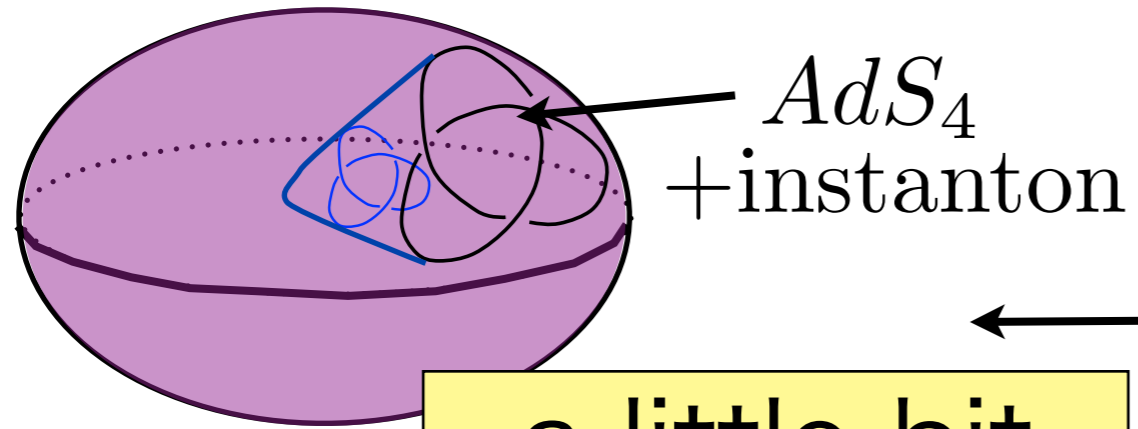
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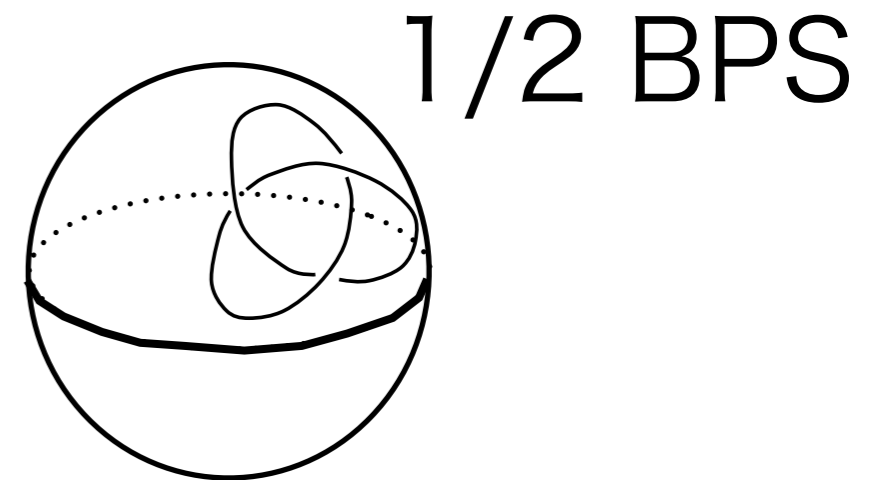
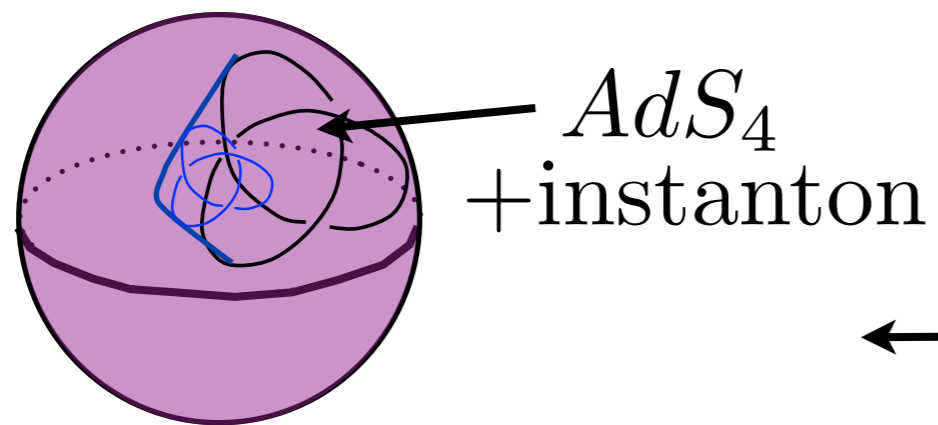
1. Wilson loop

knotted surface?

1/2 BPS



a little bit complicated



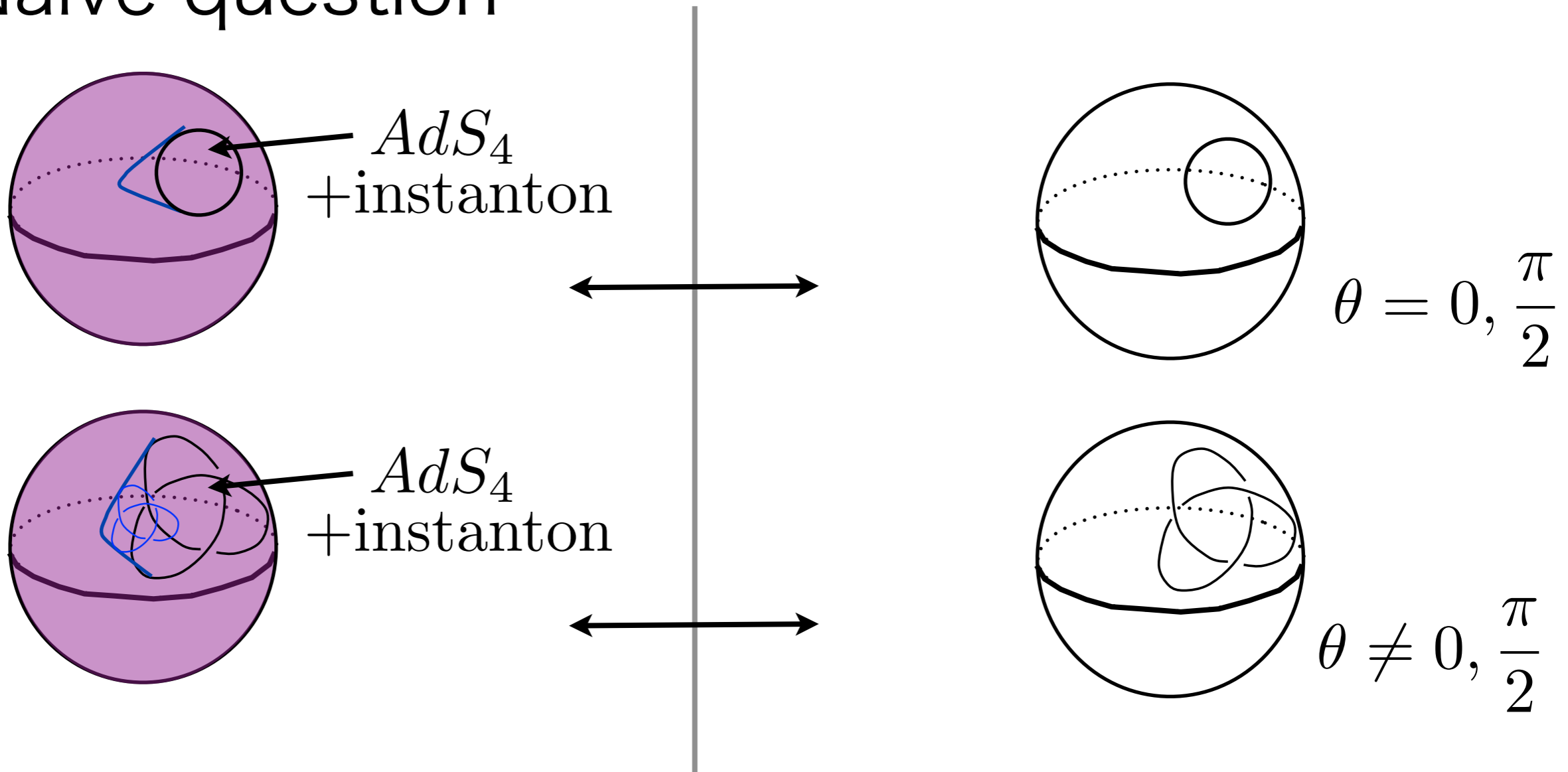
analysis is simplified?

4D Gravity

Field theory

1. Wilson loop

Naive question

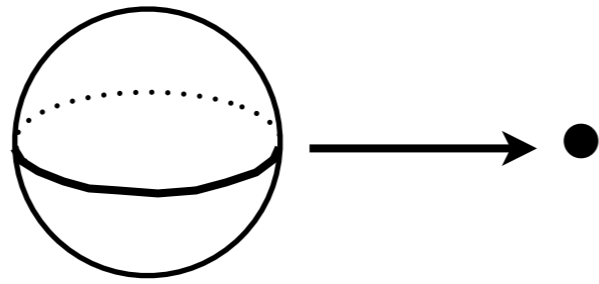


- Each surfaces are BPS? (expected so)
- Rotating in internal Sasaki-Einstein 7-mfd?

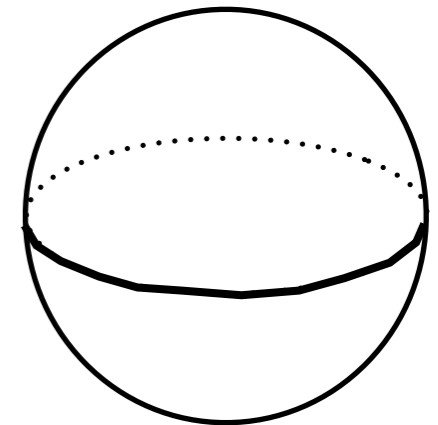
2013: Farquet, Sparks

arXiv:1304.0784

2. Large N reduction?



We can construct 1 para SUSY on



→ It guarantees the uses of
invariant 1-form.

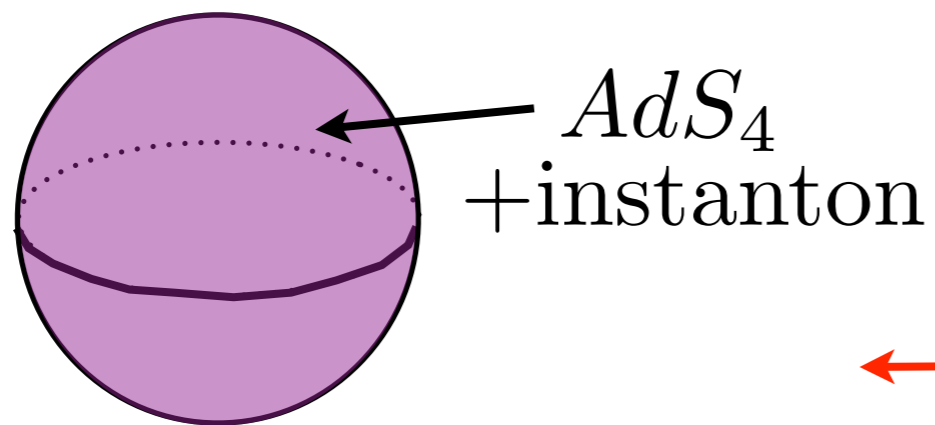
(needed field redefinition)

Kapustin et. al. case

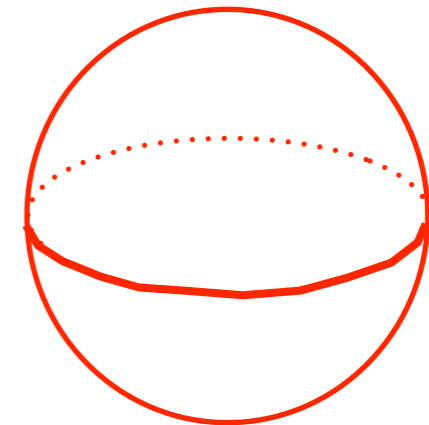
2012: Honda, Yoshida arXiv:1203.1019

2012: Asano, Ishiki, Okada, Shimasaki arXiv:1203.0559

3. Other dimension case via AdS/CFT



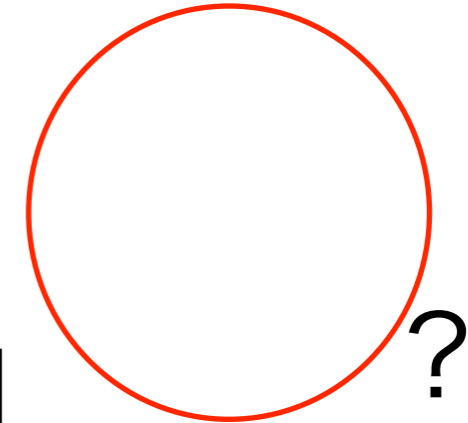
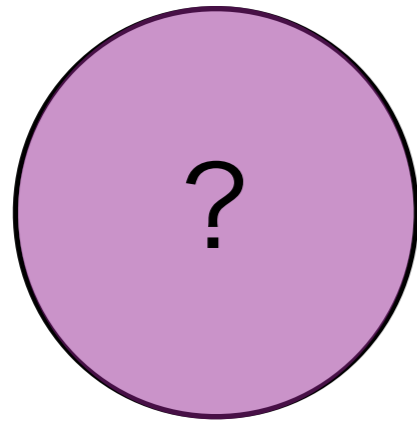
We construct it
inspired by AdS/CFT.



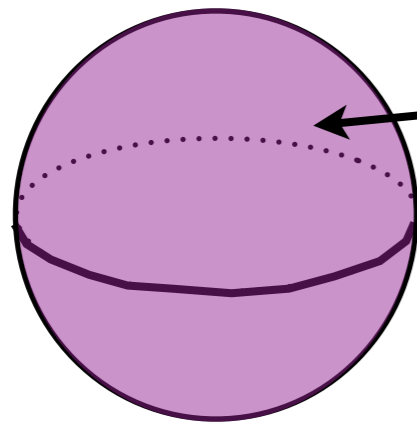
4D Gravity

Field theory

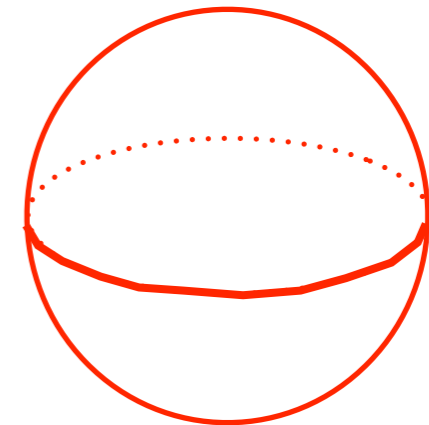
3. Other dimension case via AdS/CFT



It may be possible to construct new examples.



AdS_4
+ instanton



4D Gravity

Field theory

Thank you
for your attention.

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Vector multiplet

$\mathcal{N} = 2$ Matter multiplet

$$Z(t) = \int [\mathcal{D}A \dots \mathcal{D}\phi \dots] e^{\frac{ik}{4\pi} S_{\text{CS}} - t(\delta_{\bar{\epsilon}}\text{-exact})}$$

$$t \rightarrow \infty$$

steepest decent method around $(\delta_{\bar{\epsilon}}\text{-exact}) = 0$
turns to be exact.

$$Z(\infty) = \int d\sigma_0 e^{\frac{ik}{4\pi} S_{\text{CS}}^{cl}} \frac{\det \Delta_f}{\det \Delta_b}$$

$$A = 0, \bar{\lambda} = \lambda = 0, \sigma = \frac{\sigma_0}{f}, D = -\frac{\sigma_0}{sf^3}$$

$$\Delta_b \sim i\alpha(\sigma_0) - *df$$

$$\Delta_f = if\gamma^\mu \mathcal{D}_\mu + \frac{if'}{2}\gamma_3 + i\alpha(\sigma_0) - \frac{1}{2sf}$$

$$\mathcal{L}_{\text{YM}} = \delta_{\bar{\epsilon}}\text{-exact}$$

$$\phi = \bar{\phi} = \dots = 0$$

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1)ff' \mathcal{D}_\theta - \dots$$

$$\Delta_f = -if\gamma^\mu \mathcal{D}_\mu - if'(\Delta - \frac{1}{2})\gamma_3 - \dots$$

$$\mathcal{L}_{\text{mat}} = \delta_{\bar{\epsilon}}\text{-exact}$$

SUSY on round sphere

$\mathcal{N} = 2$ Vector multiplet

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$$\Delta_f = if\gamma^\mu \mathcal{D}_\mu + \frac{if'}{2}\gamma_3 + i\alpha(\sigma_0) - \frac{1}{2sf}$$

$$\frac{\det \Delta_f}{\det \Delta_b}$$

$$\Delta_b \mathcal{B} = M \mathcal{B}$$

$$\Delta_f \Lambda = M \Lambda$$

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

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cf. arXiv:1012.3512

Pairing structure

$$\Delta_b \mathcal{B} = M \mathcal{B}$$

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Pairing structure

$$\Delta_b \mathcal{B} = M \mathcal{B}$$

$$\downarrow \Lambda = \gamma^\mu \in \mathcal{B}_\mu$$

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$$\frac{\det \Delta_f}{\det \Delta_b}$$

cf. arXiv:1012.3512

Pairing structure

$$\Delta_b \mathcal{B} = M \mathcal{B}$$

$$\mathcal{B} = f^{-1} \left(d(f\bar{\epsilon}\Lambda) + [iM + \alpha(\sigma_0)]\bar{\epsilon}\gamma_\mu \Lambda dx^\mu \right)$$

$$\Delta_f \Lambda = M \Lambda$$

SUSY on round sphere

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Pairing structure

$$\Delta_b \mathcal{B}_{rel} = M_b \mathcal{B}_{rel}$$



0

$$\Delta_b \mathcal{B} = M \mathcal{B}$$



$$\Delta_f \Lambda = M \Lambda$$

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0

$$\Delta_b \mathcal{B} = M \mathcal{B}$$



$$\Delta_f \Lambda = M \Lambda$$



0



$$\Delta_f \Lambda_{rel} = M_f \Lambda_{rel}$$

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$$\Delta_b \mathcal{B} = M \mathcal{B}$$



$$\Delta_f \Lambda = M \Lambda$$



0



$$\Delta_f \Lambda_{rel} = M_f \Lambda_{rel}$$

$$\frac{M}{M} = 1$$

SUSY on round sphere

$\mathcal{N} = 2$ Vector multiplet

$$\frac{\det \Delta_f}{\det \Delta_b} = \prod_{m,n} \frac{M_f}{M_b}$$

$$\Delta_b \mathcal{B}_{rel} = M_b \mathcal{B}_{rel}$$



0

$$M_b = \frac{1}{s} m + n + i\alpha \cdot \sigma_0$$

$$m, n \leq -1$$

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

0



$$\Delta_f \Lambda_{rel} = M_f \Lambda_{rel}$$

$$M_f = \frac{1}{s} m + n + i\alpha \cdot \sigma_0$$

$$m, n \geq 0 \text{ except for } m = n = 0$$

SUSY on round sphere

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$$\frac{\det \Delta_f}{\det \Delta_b} = \prod_{m,n} \frac{M_f}{M_b}$$

$$= \prod_{\alpha \in \Delta_+} 4 \sinh(\pi \alpha(\sigma_0)) \sinh s(\pi \alpha(\sigma_0))$$

$$\Delta_b \mathcal{B}_{rel} = M_b \mathcal{B}_{rel}$$



0

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$$= \prod_{\alpha \in \Delta_+} 4 \sinh(\pi \alpha(\sigma_0)) \sinh s(\pi \alpha(\sigma_0))$$

We recover the result on squashed sphere!

$$\Delta_b \mathcal{B}_{rel} = M_b \mathcal{B}_{rel}$$



0

$$M_b = \frac{1}{s} m + n + i \alpha \cdot \sigma_0$$

$$m, n \leq -1$$

0



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$$\mathcal{L}_{\text{mat}} = \delta_{\bar{\epsilon}}\text{-exact}$$

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$$\frac{\det \Delta_f}{\det \Delta_b}$$

$$\Delta_b \Phi = M(M - 2i\sigma_0) \Phi$$

$$\Psi_1, \Psi_2 : \Delta_f = M, M - 2i\sigma_0$$

SUSY on round sphere

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Pairing structure

$$\Delta_b \Phi = M(M - 2i\sigma_0) \Phi$$

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Pairing structure

$$\Delta_b \Phi = M(M - 2i\sigma_0) \Phi$$

$$\Psi_1 = f^{-1} \epsilon \Phi$$

$$\Psi_2 = i \gamma^\mu \epsilon \mathcal{D}_\mu \Phi + i \frac{\sigma_0}{f} \epsilon \Phi - \frac{\Delta}{s f^2} \epsilon \Phi - i \Delta \frac{f'}{f} \bar{\epsilon} \Phi$$

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$$\frac{\det \Delta_f}{\det \Delta_b}$$

Pairing structure

$$\Delta_b \Phi = M(M - 2i\sigma_0) \Phi$$

$$\Phi = \bar{\epsilon} \Psi \uparrow$$

$$\Psi_1, \Psi_2 : \Delta_f = M, M - 2i\sigma_0$$

SUSY on round sphere

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$$\frac{\det \Delta_f}{\det \Delta_b}$$

Pairing structure

$$\Delta_b \Phi_{rel} = M_b (M_b - 2i\sigma_0) \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = M_b, M_b - 2i\sigma_0$$

$$\Delta_b \Phi = M (M - 2i\sigma_0) \Phi$$



$$\Psi_1, \Psi_2 : \Delta_f = M, M - 2i\sigma_0$$



SUSY on round sphere

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Pairing structure

$$\Delta_b \Phi_{rel} = M_b (M_b - 2i\sigma_0) \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, M_b - 2i\sigma_0$$

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Pairing structure

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, \cancel{M_b - 2i\sigma_0}$$

$$\Delta_b \Phi = M (M - 2i\sigma_0) \Phi$$



$$\Psi_1, \Psi_2 : \Delta_f = M, M - 2i\sigma_0$$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Matter multiplet

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1) f f' \mathcal{D}_\theta - \dots$$

$$\Delta_f = -i f \gamma^\mu \mathcal{D}_\mu - i f' \left(\Delta - \frac{1}{2}\right) \gamma_3 - \dots$$

$$\frac{\det \Delta_f}{\det \Delta_b}$$

Pairing structure

$$\Delta_b \Phi_{rel} = \underbrace{M_b}_{\text{circled}} \underbrace{(M_b - 2i\sigma_0)}_{\text{crossed out}} \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = \underbrace{M_b}_{\text{crossed out}}, \underbrace{(M_b - 2i\sigma_0)}_{\text{crossed out}}$$

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$$0 = \bar{\epsilon} \Psi_{rel}$$



$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

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$\mathcal{N} = 2$ Matter multiplet

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$$\frac{\det \Delta_f}{\det \Delta_b}$$

cf. arXiv:1012.3512

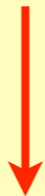
Pairing structure

$$\Delta_b \Phi_{rel} = \cancel{M_b} (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, \cancel{M_b - 2i\sigma_0}$$

$$\Delta_b \Phi = M(M - 2i\sigma_0) \Phi$$



$$\Psi_1, \Psi_2 : \Delta_f = M, M - 2i\sigma_0$$

$$\frac{M(M - 2i\sigma_0)}{M(M - 2i\sigma_0)} = 1$$

$$0 = \bar{\epsilon} \Psi_{rel}$$



$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

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$\mathcal{N} = 2$ Matter multiplet

$$\frac{\det \Delta_f}{\det \Delta_b} = \prod_{m,n} \frac{M_f}{M_b}$$

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b} \cancel{2i\sigma_0}) \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b} \cancel{M_b} \cancel{2i\sigma_0}$$

$$M_b = i\rho \cdot \sigma_0 + \frac{1}{s} m + n - \frac{\Delta - 2}{2} \left(\frac{1}{s} + 1 \right)$$

$$m, n \geq 0$$

$$0 = \bar{\epsilon} \Psi_{rel}$$



$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

$$M_f = i\rho \cdot \sigma_0 - \frac{1}{s} m - n - \frac{\Delta}{2} \left(\frac{1}{s} + 1 \right)$$

$$m, n \leq 0$$

SUSY on round sphere

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$\mathcal{N} = 2$ Matter multiplet

$$\frac{\det \Delta_f}{\det \Delta_b} = \prod_{m,n} \frac{M_f}{M_b} = \prod_{\rho \in \text{weight}} s_b \left(\frac{i(\sqrt{\frac{1}{s}} + \sqrt{s})}{2} (1 - \Delta) - \sqrt{s} \rho \cdot \sigma_0 \right)$$

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$\mathcal{N} = 2$ Matter multiplet

$$\frac{\det \Delta_f}{\det \Delta_b} = \prod_{m,n} \frac{M_f}{M_b} = \prod_{\rho \in \text{weight}} s_b \left(\frac{i(\sqrt{\frac{1}{s}} + \sqrt{s})}{2} (1 - \Delta) - \sqrt{s} \rho \cdot \sigma_0 \right)$$

We recover the result on squashed sphere!

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, \cancel{M_b - 2i\sigma_0}$$

$$M_b = i\rho \cdot \sigma_0 + \frac{1}{s} m + n - \frac{\Delta - 2}{2} \left(\frac{1}{s} + 1 \right)$$

$$m, n \geq 0$$

$$0 = \bar{\epsilon} \Psi_{rel}$$



$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

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SUSY on round sphere

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$\mathcal{N} = 2$ Vector multiplet

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$$Z(t) = \int [\mathcal{D}A \dots \mathcal{D}\phi \dots] e^{\frac{ik}{4\pi} S_{\text{CS}} - t(\delta_{\bar{\epsilon}}\text{-exact})}$$

$$t \rightarrow \infty$$

steepest decent method around $(\delta_{\bar{\epsilon}}\text{-exact}) = 0$
turns to be exact.

$$Z(\infty) = \int d\sigma_0 e^{\frac{ik}{4\pi} S_{\text{CS}}^{cl}} \frac{\det \Delta_f}{\det \Delta_b}$$

$$A = 0, \bar{\lambda} = \lambda = 0, \sigma = \frac{\sigma_0}{f}, D = -\frac{\sigma_0}{sf^3}$$

$$\frac{\det \Delta_f}{\det \Delta_b}$$

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$$= \prod_{\alpha \in \Delta_+} 4 \sinh(\pi \alpha(\sigma_0)) \sinh s(\pi \alpha(\sigma_0))$$

$$\mathcal{L}_{\text{YM}} = \delta_{\bar{\epsilon}}\text{-exact}$$

$$\phi = \bar{\phi} = \dots = 0$$

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$$\mathcal{L}_{\text{mat}} = \delta_{\bar{\epsilon}}\text{-exact}$$

How about S_{CS}^{cl} ?

Classical configuration : $A = 0, \bar{\lambda} = \lambda = 0, \sigma = \frac{\sigma_0}{f}, D = -\frac{\sigma_0}{sf^3}$

$$Z(\infty) = \int d\sigma_0 e^{\frac{ik}{4\pi} S_{\text{CS}}^{\text{cl}}} \frac{\det \Delta_f}{\det \Delta_b}$$

$$\mathcal{L}_{\text{CS}} = \text{Tr} \left[\frac{1}{\sqrt{g}} \varepsilon^{\mu\nu\lambda} (A_\mu \partial_\nu A_\lambda - \frac{2i}{3} A_\mu A_\nu A_\lambda) - \bar{\lambda} \lambda + 2D\sigma \right] = \text{Tr} \left[-2 \frac{\sigma_0^2}{sf^4} \right]$$

$$S_{\text{CS}}^{\text{cl}} = \int \sqrt{g} d^3x \mathcal{L}_{\text{CS}}$$

$$= (2\pi)^2 \int \sin \theta \cos \theta d\theta \text{Tr} \left[-2 \frac{\sigma_0^2}{sf^4} \right]$$

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$$= -4\pi^2 s \text{Tr} \sigma_0^2 \quad : \text{Equivalent to squashed one}$$

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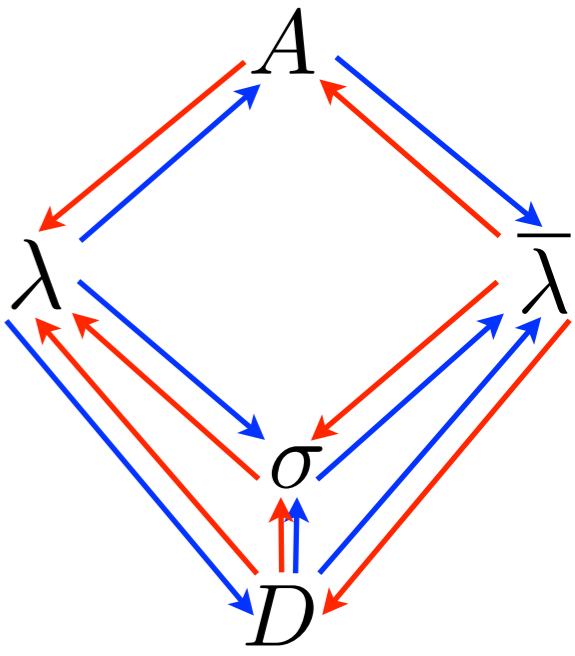
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How about Wilson loop?

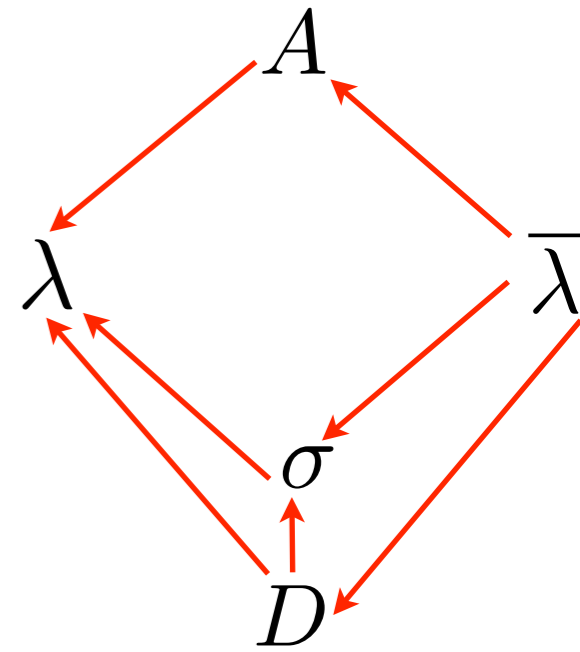
$$\delta_{\bar{\epsilon}}(\text{Wilson loop}) = 0$$

$$\text{Wilson loop} = \text{Tr}_R \mathcal{P} \exp \left(\oint_C dt (iA_\mu \dot{x}^\mu + \sigma |\dot{x}|) \right)$$



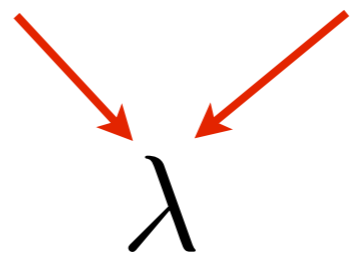
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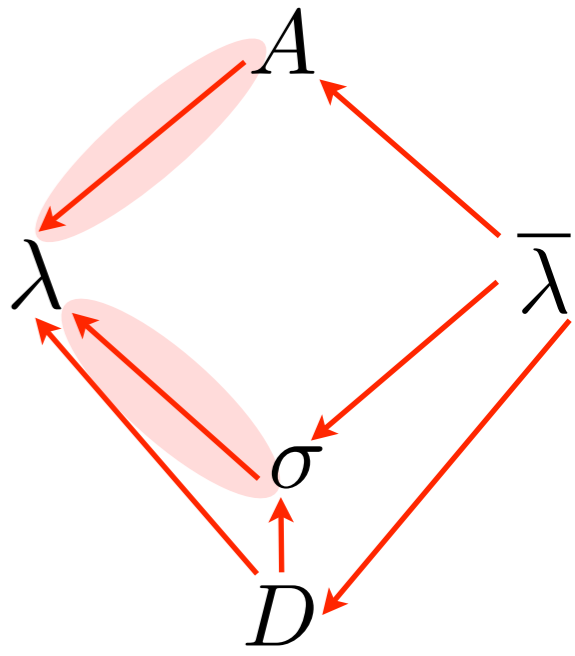
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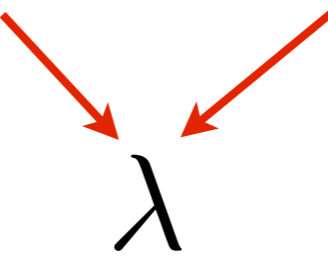
Wilson loop = $\text{Tr}_R \mathcal{P} \exp \left(\oint_C dt (i A_\mu \dot{x}^\mu + \sigma |\dot{x}|) \right)$



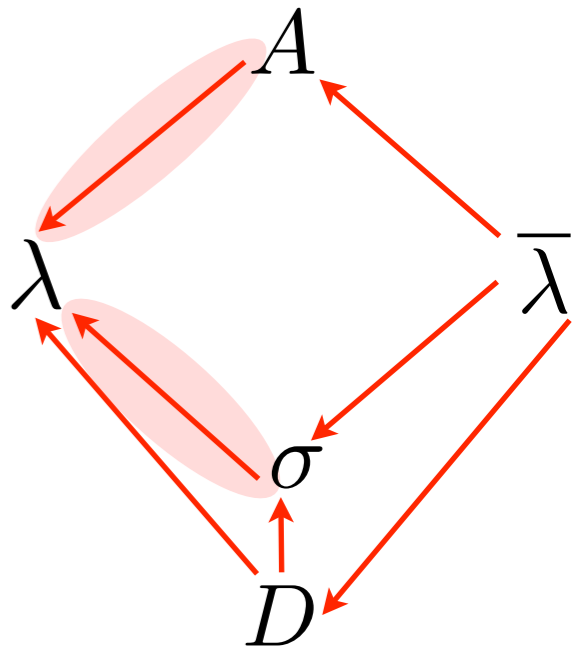


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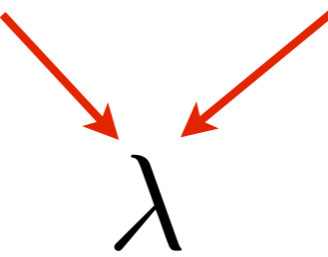
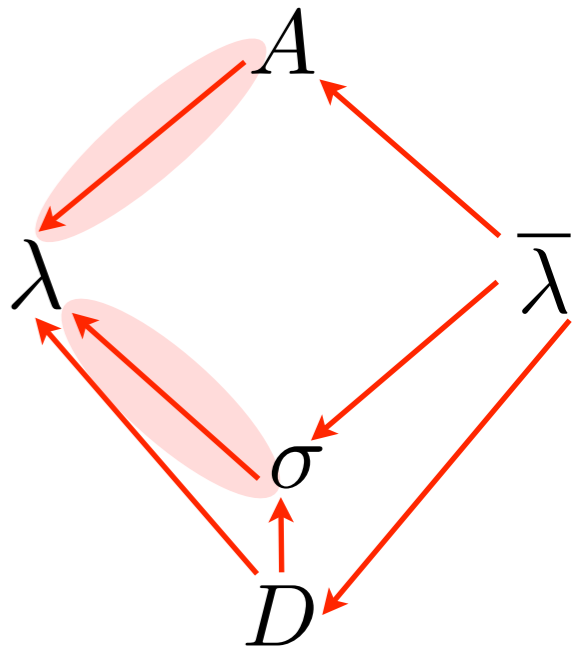


$$\delta_{\bar{\epsilon}}(\text{Wilson loop}) \propto \bar{\epsilon} (\gamma_\mu \dot{x}^\mu + |\dot{x}|) \lambda$$



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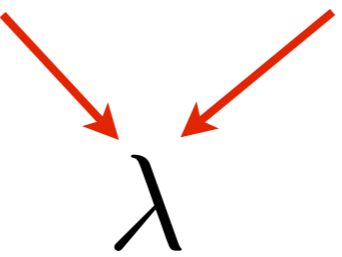
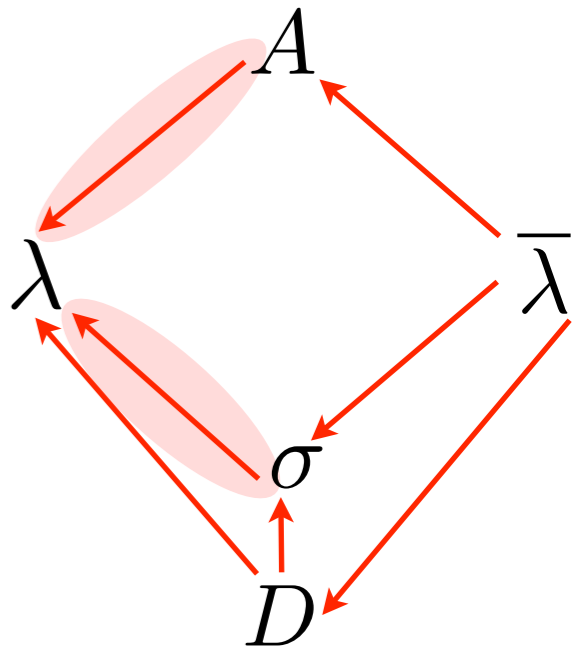


$$\delta_{\bar{\epsilon}}(\text{Wilson loop}) \propto \bar{\epsilon} (\gamma_\mu \dot{x}^\mu + |\dot{x}|) \lambda = 0$$

Equations for the Wilson loop's contour!

$$\delta_{\bar{\epsilon}}(\text{Wilson loop}) = 0$$

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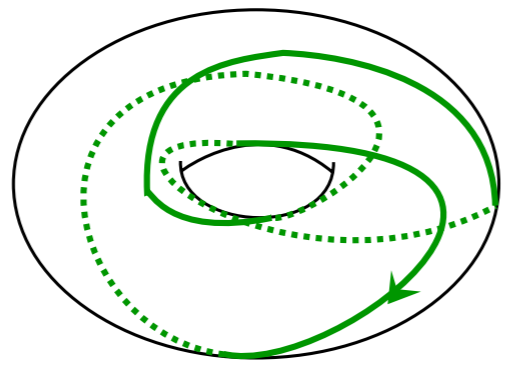


$$\propto \bar{\epsilon} (\gamma_\mu \dot{x}^\mu + |\dot{x}|) = 0$$

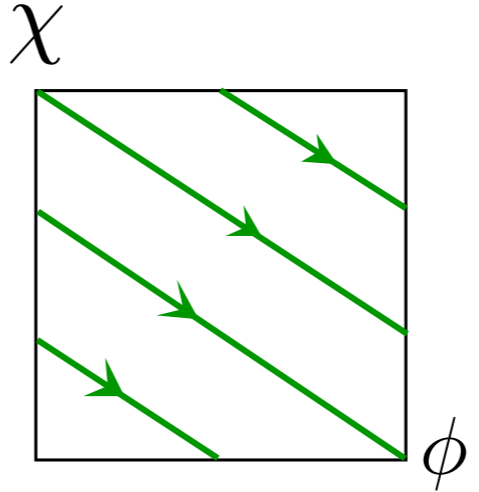


Equations for the Wilson loop's contour!

Torus knot



=



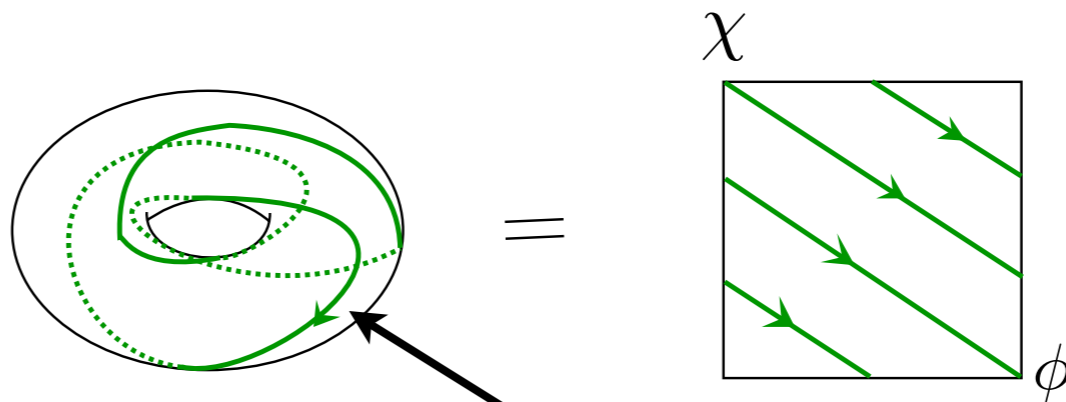
SUSY on round sphere

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$\mathcal{N} = 2$ Vector multiplet

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$$Z(\infty) = \int d\sigma_0 e^{\frac{ik}{4\pi} S_{\text{CS}}^{\text{cl}}} \frac{\det \Delta_f}{\det \Delta_b} \quad 1/2 \text{ BPS}$$

$$A = 0, \bar{\lambda} = \lambda = 0, \sigma = \frac{\sigma_0}{f}, D = -\frac{\sigma_0}{sf^3}$$

$$\frac{\det \Delta_f}{\det \Delta_b} = \prod_{\alpha \in \Delta_+} 4 \sinh(\pi \alpha(\sigma_0)) \sinh s(\pi \alpha(\sigma_0))$$

$\mathcal{L}_{\text{YM}} = \delta_{\bar{\epsilon}}$ -exact

$$\phi = \bar{\phi} = \dots = 0$$

$$\frac{\det \Delta_f}{\det \Delta_b} = \prod_{\rho \in \text{weight}} s_b \left(\frac{i(\sqrt{\frac{1}{s}} + \sqrt{s})}{2} (1 - \Delta) - \sqrt{s} \rho \cdot \sigma_0 \right)$$

$\mathcal{L}_{\text{mat}} = \delta_{\bar{\epsilon}}$ -exact

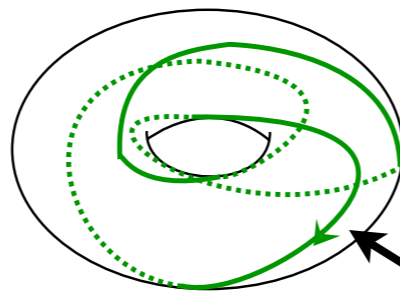
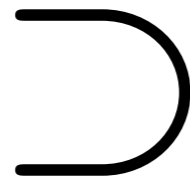
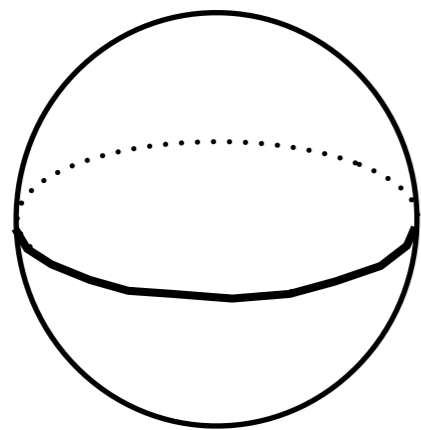
SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

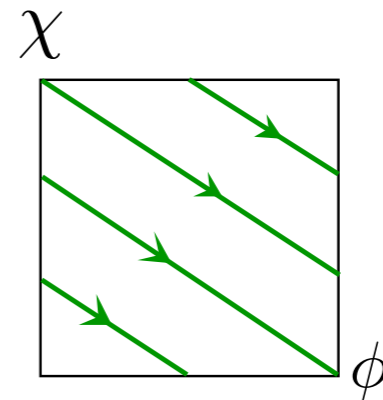
$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Vector multiplet

$\mathcal{N} = 2$ Matter multiplet



=



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1/2 BPS

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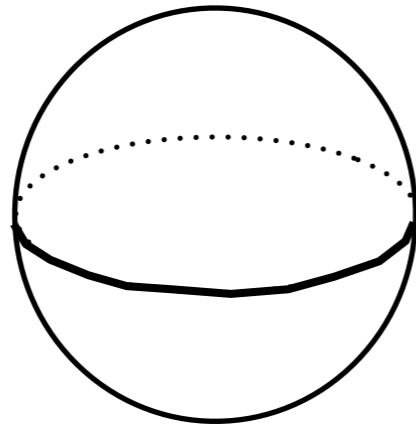
$$= \prod_{\rho \in \text{weight}} s_b \left(\frac{i(\sqrt{\frac{1}{s}} + \sqrt{s})}{2} (1 - \Delta) - \sqrt{s} \rho \cdot \sigma_0 \right)$$

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In introduction

Is it possible to get deformed results

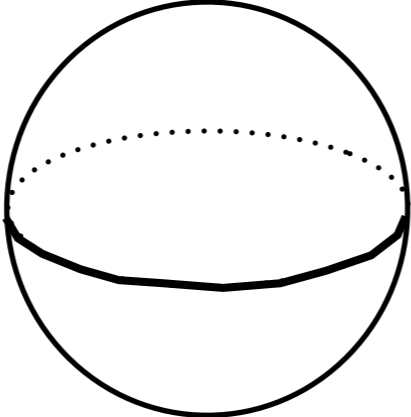
even on



?

In introduction

Is it possible to get deformed results

even on  ?

Now

The answer is **YES**.

(Based on **direct check**)

vector

SUSY on round sphere

$\mathcal{N} = 2$ Vector multiplet

• SUSY transformation

→ Just Replacing usual Killing spinors with our unusual Killing spinors.

$$\delta_\epsilon A_\mu = -\frac{i}{2} \bar{\lambda} \gamma_\mu \epsilon, \quad \delta_{\bar{\epsilon}} A_\mu = -\frac{i}{2} \bar{\epsilon} \gamma_\mu \lambda,$$

$$\delta_\epsilon \sigma = +\frac{1}{2} \bar{\lambda} \epsilon, \quad \delta_{\bar{\epsilon}} \sigma = +\frac{1}{2} \bar{\epsilon} \lambda,$$

$$\delta_\epsilon \lambda = \frac{1}{2} \gamma^{\mu\nu} \epsilon F_{\mu\nu} - D\epsilon + i\gamma^\mu \epsilon \mathcal{D}_\mu \sigma + \frac{2i}{3} \sigma \gamma^\mu \mathcal{D}_\mu \epsilon, \quad \delta_{\bar{\epsilon}} \lambda = 0,$$

$$\delta_\epsilon \bar{\lambda} = 0, \quad \delta_{\bar{\epsilon}} \bar{\lambda} = \frac{1}{2} \gamma^{\mu\nu} \bar{\epsilon} F_{\mu\nu} + D\bar{\epsilon} - i\gamma^\mu \bar{\epsilon} \mathcal{D}_\mu \sigma - \frac{2i}{3} \sigma \gamma^\mu \mathcal{D}_\mu \bar{\epsilon},$$

$$\delta_\epsilon D = +\frac{i}{2} \mathcal{D}_\mu \bar{\lambda} \gamma^\mu \epsilon - \frac{i}{2} [\bar{\lambda} \epsilon, \sigma] + \frac{i}{6} \bar{\lambda} \gamma^\mu \mathcal{D}_\mu \epsilon, \quad \delta_{\bar{\epsilon}} D = -\frac{i}{2} \bar{\epsilon} \gamma^\mu \mathcal{D}_\mu \lambda + \frac{i}{2} [\bar{\epsilon} \lambda, \sigma] - \frac{i}{6} \mathcal{D}_\mu \bar{\epsilon} \gamma^\mu \lambda.$$

$$f^2(\theta) = \sin^2 \theta + \frac{1}{s^2} \cos^2 \theta$$

$$\mathcal{D}_\mu \epsilon = \frac{i}{2s f^2} \gamma_\mu \epsilon - \frac{1}{2} \frac{f'}{f} \gamma_\mu \bar{\epsilon}$$

$$\mathcal{D}_\mu \bar{\epsilon} = \frac{i}{2s f^2} \gamma_\mu \bar{\epsilon} - \frac{1}{2} \frac{f'}{f} \gamma_\mu \epsilon$$

SUSY on round sphere

$\mathcal{N} = 2$ Vector multiplet

• SUSY transformation

→ Just Replacing usual Killing spinors with our unusual Killing spinors.

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\} A_\mu = i v^\nu \partial_\nu A_\mu + i \partial_\mu v^\nu A_\nu + \mathcal{D}_\mu \Lambda,$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\} \sigma = i v^\mu \partial_\mu \sigma + i[\Lambda, \sigma]$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\} \lambda = i v^\mu \partial_\mu \lambda + \frac{i}{4} \Theta_{\mu\nu} \gamma^{\mu\nu} \lambda + i[\Lambda, \lambda] + \alpha \lambda,$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\} \bar{\lambda} = i v^\mu \partial_\mu \bar{\lambda} + \frac{i}{4} \Theta_{\mu\nu} \gamma^{\mu\nu} \bar{\lambda} + i[\Lambda, \bar{\lambda}] - \alpha \bar{\lambda},$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\} D = i v^\mu \partial_\mu D + i[\Lambda, D],$$

$$f^2(\theta) = \sin^2 \theta + \frac{1}{s^2} \cos^2 \theta$$

$$\mathcal{D}_\mu \epsilon = \frac{i}{2s f^2} \gamma_\mu \epsilon - \frac{1}{2} \frac{f'}{f} \gamma_\mu \bar{\epsilon}$$

$$\mathcal{D}_\mu \bar{\epsilon} = \frac{i}{2s f^2} \gamma_\mu \bar{\epsilon} - \frac{1}{2} \frac{f'}{f} \gamma_\mu \epsilon$$

$\left(\begin{array}{l} v^\mu, \Lambda, \text{ same definition in} \\ \Theta_{\mu\nu}, \alpha \text{ Hama, Hosomichi, Lee} \end{array} \right)$

matter

SUSY on round sphere

$\mathcal{N} = 2$ Matter multiplet

- SUSY transformation

→ Just Replacing usual Killing spinors with our unusual Killing spinors.

$$\delta_\epsilon \phi = 0, \quad \delta_{\bar{\epsilon}} \phi = \bar{\epsilon} \psi,$$

$$\delta_\epsilon \bar{\phi} = \epsilon \bar{\psi}, \quad \delta_{\bar{\epsilon}} \bar{\phi} = 0,$$

$$\delta_\epsilon \psi = i\gamma^\mu \epsilon \mathcal{D}_\mu \phi + i\epsilon \sigma \phi + \frac{2\Delta i}{3} \gamma^\mu \mathcal{D}_\mu \epsilon \phi, \quad \delta_{\bar{\epsilon}} \psi = \bar{\epsilon} F,$$

$$\delta_\epsilon \bar{\psi} = \bar{F} \epsilon, \quad \delta_{\bar{\epsilon}} \bar{\psi} = i\gamma^\mu \bar{\epsilon} \mathcal{D}_\mu \bar{\phi} + i\bar{\phi} \sigma \bar{\epsilon} + \frac{2\Delta i}{3} \bar{\phi} \gamma^\mu \mathcal{D}_\mu \bar{\epsilon},$$

$$\delta_\epsilon F = \epsilon (i\gamma^\mu \mathcal{D}_\mu \psi - i\sigma \psi - i\lambda \phi) + \frac{i}{3} (2\Delta - 1) \mathcal{D}_\mu \epsilon \gamma^\mu \psi, \quad \delta_{\bar{\epsilon}} F = 0,$$

$$\delta_\epsilon \bar{F} = 0, \quad \delta_{\bar{\epsilon}} \bar{F} = \bar{\epsilon} (i\gamma^\mu \mathcal{D}_\mu \bar{\psi} - i\bar{\psi} \sigma + i\bar{\phi} \lambda) + \frac{i}{3} (2\Delta - 1) \mathcal{D}_\mu \bar{\epsilon} \gamma^\mu \bar{\psi}.$$

$$f^2(\theta) = \sin^2 \theta + \frac{1}{s^2} \cos^2 \theta$$

$$\mathcal{D}_\mu \epsilon = \frac{i}{2sf^2} \gamma_\mu \epsilon - \frac{1}{2} \frac{f'}{f} \gamma_\mu \bar{\epsilon}$$

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SUSY on round sphere

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- SUSY transformation

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$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\}\phi = i v^\mu \partial_\mu \phi + i \Lambda \phi - \Delta \alpha \phi,$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\}\bar{\phi} = i v^\mu \partial_\mu \bar{\phi} - i \bar{\phi} \Lambda + \Delta \alpha \bar{\phi},$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\}\psi = i v^\mu \partial_\mu \psi + \frac{i}{4} \Theta_{\mu\nu} \gamma^{\mu\nu} \psi + i \Lambda \psi + (1 - \Delta) \alpha \psi,$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\}\bar{\psi} = i v^\mu \partial_\mu \bar{\psi} + \frac{i}{4} \Theta_{\mu\nu} \gamma^{\mu\nu} \bar{\psi} - i \bar{\psi} \Lambda + (\Delta - 1) \alpha \bar{\psi},$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\}F = i v^\mu \partial_\mu F + i \Lambda F + (2 - \Delta) \alpha F,$$

$$\{\delta_\epsilon, \delta_{\bar{\epsilon}}\}\bar{F} = i v^\mu \partial_\mu \bar{F} - i \bar{F} \Lambda + (\Delta - 2) \alpha \bar{F},$$

$$f^2(\theta) = \sin^2 \theta + \frac{1}{s^2} \cos^2 \theta$$

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Poincare-Hopf

What is the spirit of localization?

Simple but interesting example $(\mu = 1, 2) x^\mu \in \mathcal{M}$ (2D manifold)

Bosonic d.o.f : (x^μ, p_μ)

Ferminic d.o.f : $(\psi^\mu, \bar{\psi}_\mu)$

SUSY: $x^\mu \xrightarrow{\delta} \psi^\mu \xrightarrow{\delta} 0$

$\bar{\psi}_\mu \xrightarrow{\delta} p_\mu \xrightarrow{\delta} 0$

Partition function :

$$Z(t) = \int \prod_{\mu} dp_{\mu} dx^{\mu} d\psi^{\mu} d\bar{\psi}_{\mu} e^{-\delta V_0 - t\delta V}$$

V_0 : a certain function

$$V = i\bar{\psi}_{\mu} v^{\mu}(x)$$

$v^{\mu}(x)$: a certain vector field on \mathcal{M}

t : “coupling constant”

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$\delta^2 = 0$ Nilpotent

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Nilpotent

$$t = 0$$

$$\frac{1}{4\pi} \int d^2x \sqrt{g} R$$

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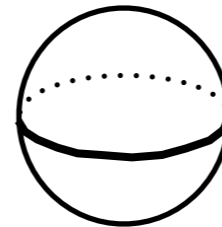
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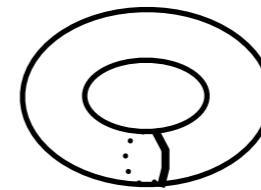
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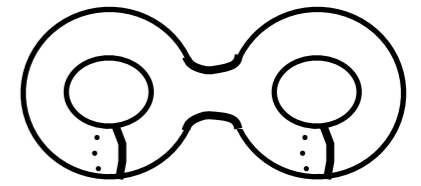
$$= \chi(\mathcal{M}) = 2(1 - g)$$



2



0



-2

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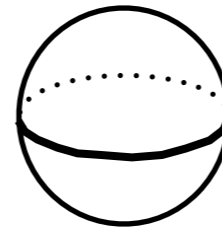
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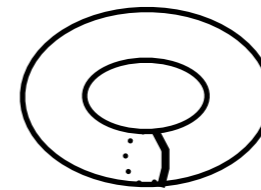
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$$\frac{1}{4\pi} \int d^2x \sqrt{g} R$$

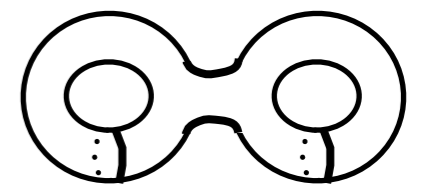
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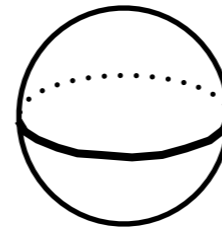
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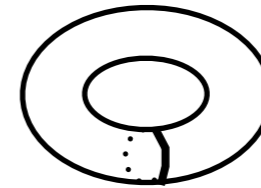
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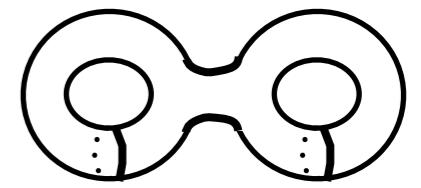
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Steepest decent method is exact

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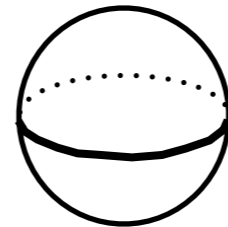
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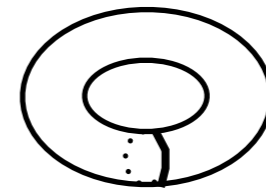
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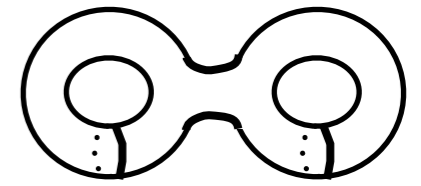
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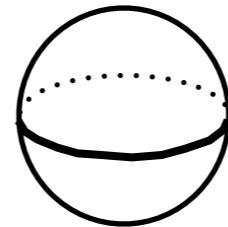
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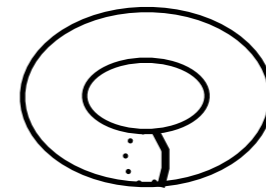
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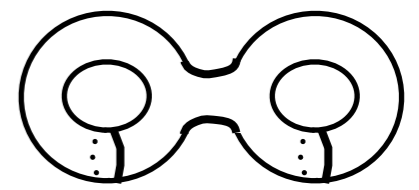
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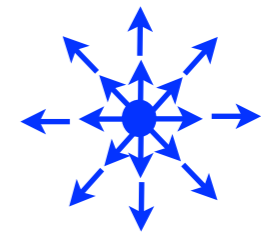
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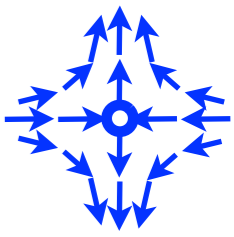
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+1



-1

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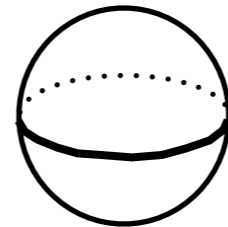
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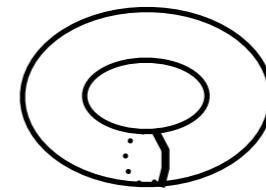
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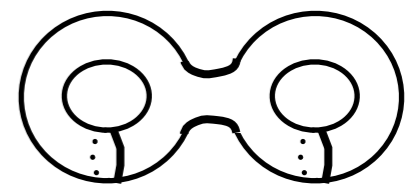
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2



0



-2

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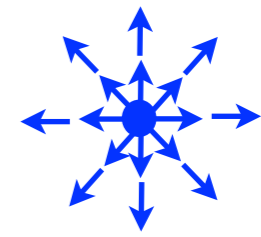
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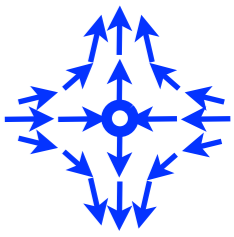
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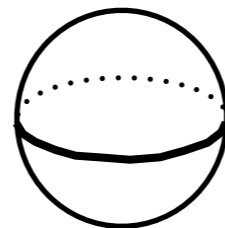
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+1



-1



What is the spirit of localization?

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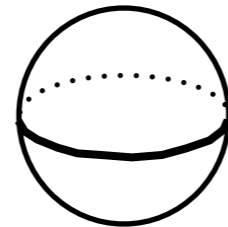
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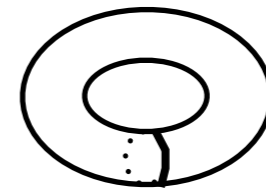
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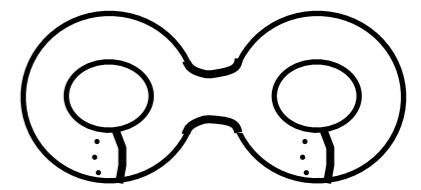
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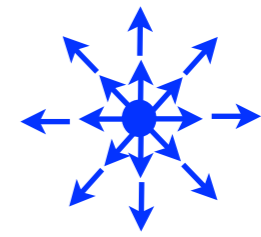
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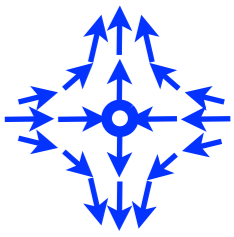
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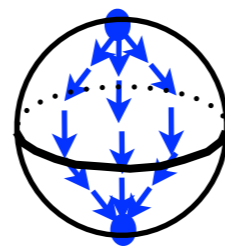
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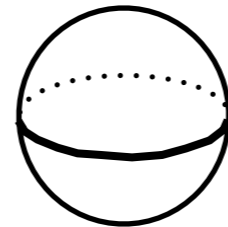
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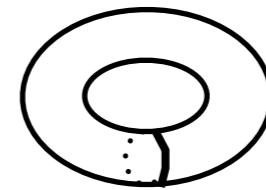
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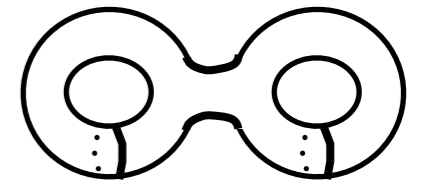
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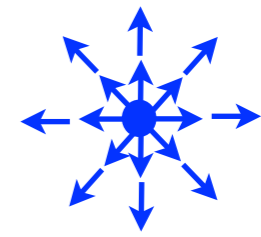
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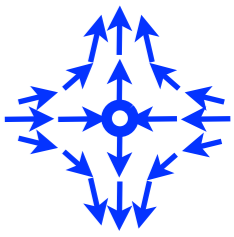
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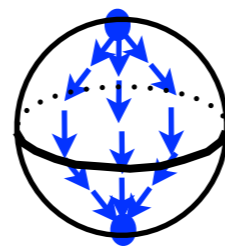
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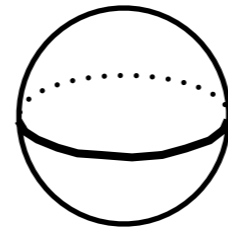
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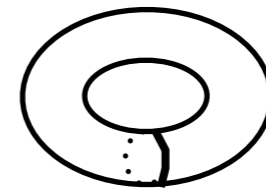
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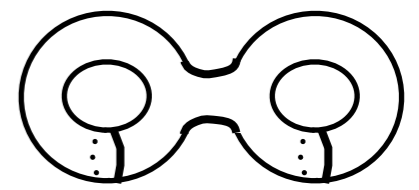
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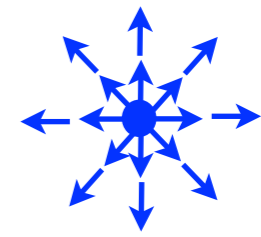
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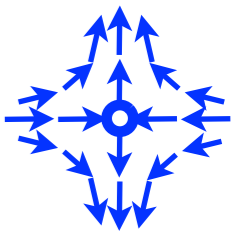
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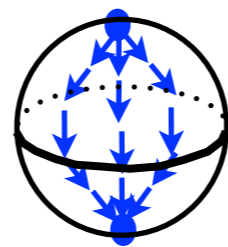
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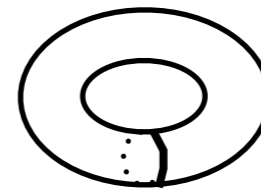
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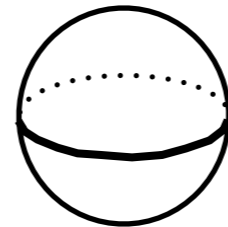
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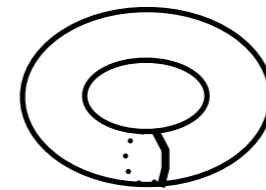
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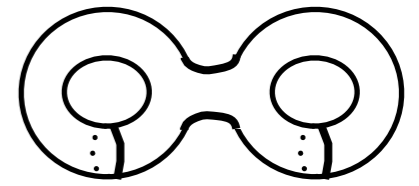
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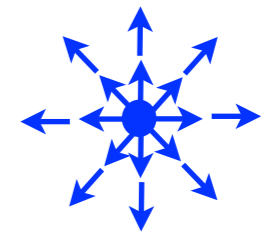
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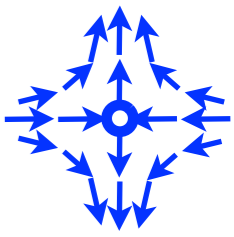
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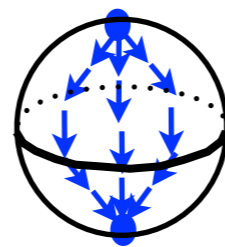
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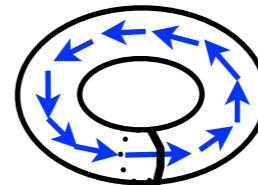
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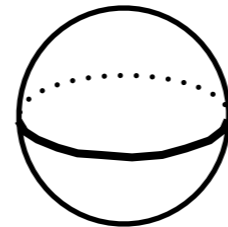
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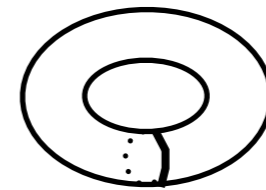
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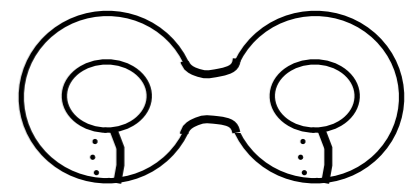
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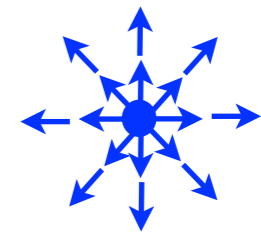
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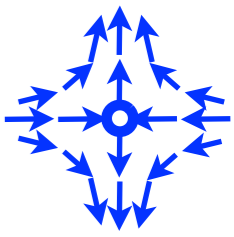
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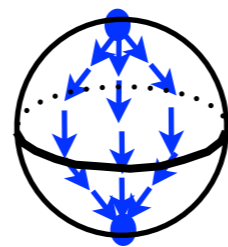
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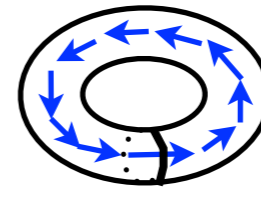
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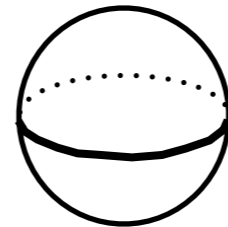
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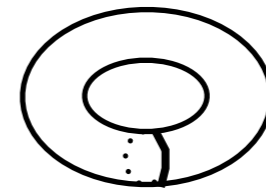
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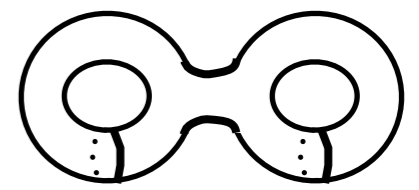
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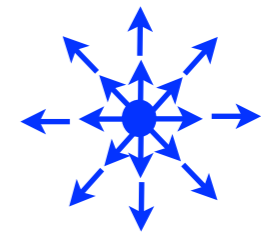
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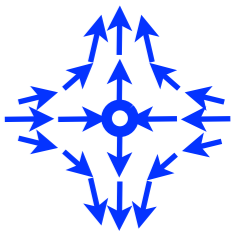
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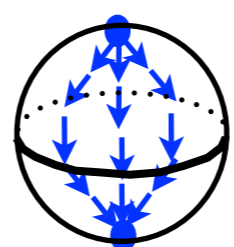
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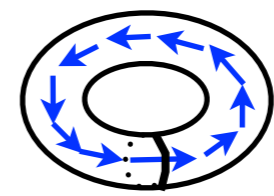
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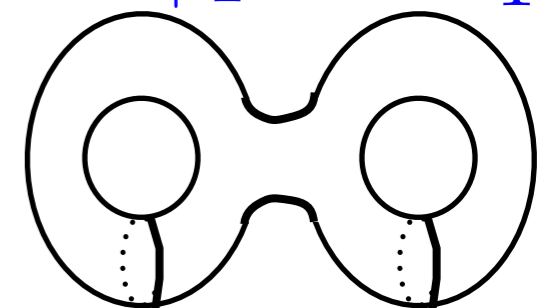
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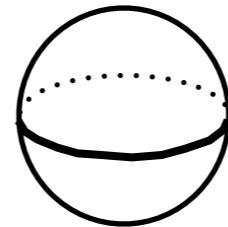
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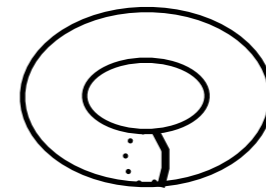
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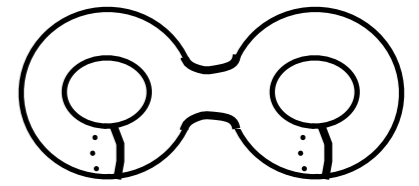
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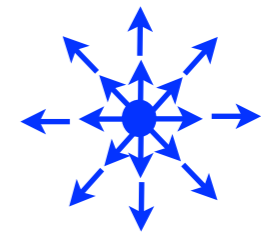
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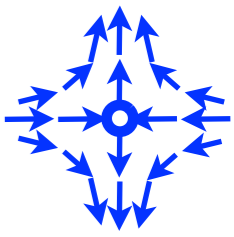
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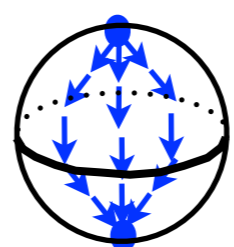
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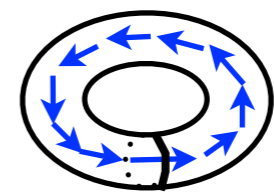
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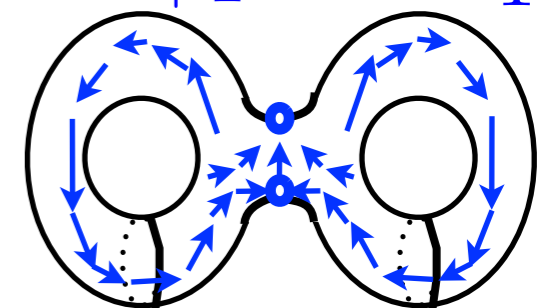
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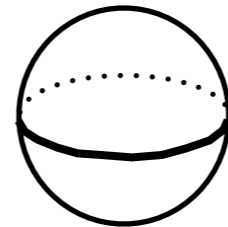
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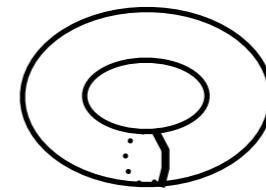
$t = 0$

$$\frac{1}{4\pi} \int d^2x \sqrt{g} R$$

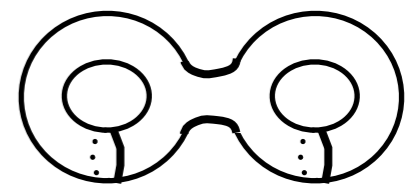
$$= \chi(\mathcal{M}) = 2(1 - g)$$



2



0



-2

Partition function :

$$Z(t) = \int \prod_{\mu} dp_{\mu} dx^{\mu} d\psi^{\mu} d\bar{\psi}_{\mu} e^{-\delta V_0 - t \delta V}$$

V_0 : a certain function

$$V = i\bar{\psi}_{\mu} v^{\mu}(x)$$

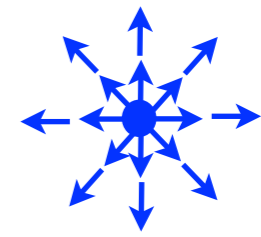
$v^{\mu}(x)$: a certain vector field on \mathcal{M}

t : “coupling constant”

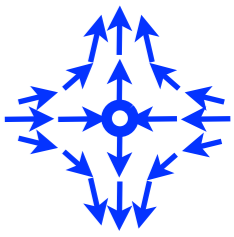
$t \rightarrow \infty$

$$\frac{1}{2\pi} \int d^2x d^2\psi d^2\bar{\psi} e^{-\frac{t^2}{2} v^2(x) + it \nabla_{\mu} v^{\nu} \bar{\psi}_{\nu} \psi^{\mu}}$$

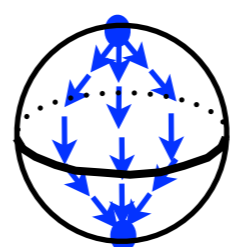
$$\sum_{v^{\mu}(x_p)=0} \frac{\det \partial_{\mu} v^{\nu}(x_p)}{|\det \partial_{\mu} v^{\nu}(x_p)|}$$



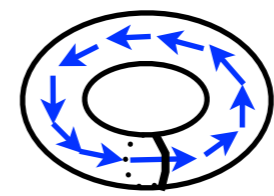
+1



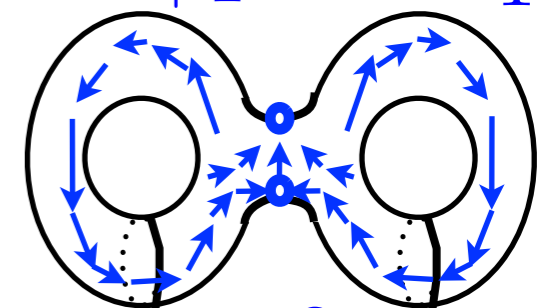
-1



2



0



-2

What is the spirit of localization?

Simple but interesting example $(\mu = 1, 2) x^\mu \in \mathcal{M}$ (2D manifold)

Bosonic d.o.f : (x^μ, p_μ)

Ferminic d.o.f : $(\psi^\mu, \bar{\psi}_\mu)$

SUSY: $x^\mu \xrightarrow{\delta} \psi^\mu \xrightarrow{\delta} 0$

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$$\delta^2 = 0$$

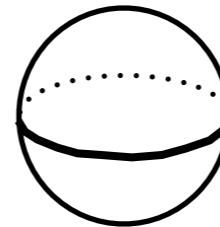
Nilpotent

does not depend on t !?

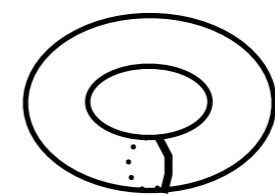
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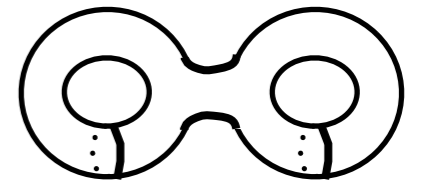
$t = 0$



2



0



-2

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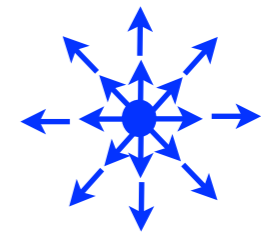
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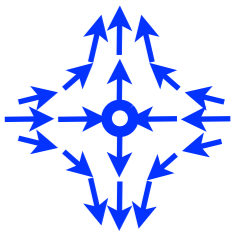
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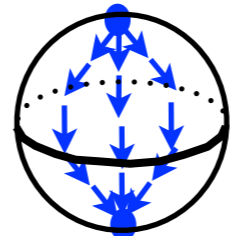
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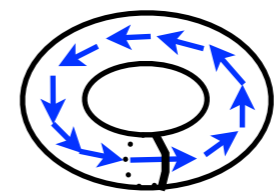
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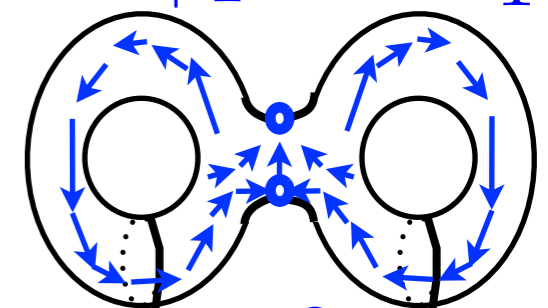
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$$\frac{d}{dt} Z(t) = \int \prod_{\mu} dp_{\mu} dx^{\mu} d\psi^{\mu} d\bar{\psi}_{\mu} \frac{d}{dt} e^{-\delta V_0 - t\delta V}$$

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$$= \int \prod_{\mu} dp_{\mu} dx^{\mu} d\psi^{\mu} d\bar{\psi}_{\mu} \delta(-V e^{-\delta V_0 - t\delta V})$$

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$\downarrow \because \delta e^{-\delta V_0 - t\delta V} = 0$

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$$= \int \prod_{\mu} \delta(dp_{\mu} dx^{\mu} d\psi^{\mu} d\bar{\psi}_{\mu}) (V e^{-\delta V_0 - t\delta V})$$

$$= 0$$

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$$= 0$$

$\because \delta e^{-\delta V_0 - t\delta V} = 0$
 \downarrow
 $= 0 \quad \because \text{SUSY}$

1-loop details

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$


$\mathcal{N} = 2$ Vector multiplet

$$\Delta_b \sim i\alpha(\sigma_0) - *df$$

$$\Delta_f = if\gamma^\mu \mathcal{D}_\mu + \frac{if'}{2}\gamma_3 + i\alpha(\sigma_0) - \frac{1}{2sf}$$


$$\frac{\det \Delta_f}{\det \Delta_b} = \frac{\prod_f M_f}{\prod_b M_b}$$

$$\Delta_b \mathcal{B}_{rel} = M_b \mathcal{B}_{rel}$$



$$0 = \gamma^\mu \epsilon \mathcal{B}_{rel \mu}$$

$$f^{-1} \left(d(f\bar{\epsilon}\Lambda_{rel}) + [iM_f + \alpha(\sigma_0)]\bar{\epsilon}\gamma_\mu \Lambda_{rel} dx^\mu \right) = 0$$



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these terms does not affect

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$$y \sim (\cos \theta)^{(-m-1)} (\sin \theta)^{(-n-1)}$$

$\rightarrow m, n \leq -1$

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We recover the result on squashed sphere!

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$$0 = \gamma^\mu \epsilon \mathcal{B}_{rel \mu}$$

$$\Delta_f \Lambda_{rel} = M_f \Lambda_{rel}$$

$$M_b = \frac{1}{s}m + n + i\alpha(\sigma_0),$$

$$0 = \left(\partial_\theta - \frac{\sin \theta}{\cos \theta} \left(\frac{1}{f}m + 1 \right) + \frac{\cos \theta}{\sin \theta} \left(\frac{1}{sf}n + 1 \right) \right) y.$$

$$y \sim (\cos \theta)^{(-m-1)} (\sin \theta)^{(-n-1)}$$

$\rightarrow m, n \leq -1$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Matter multiplet

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1) f f' \mathcal{D}_\theta - \dots$$

$$\Delta_f = -i f \gamma^\mu \mathcal{D}_\mu - i f' \left(\Delta - \frac{1}{2}\right) \gamma_3 - \dots$$

$$\frac{\det \Delta_f}{\det \Delta_b} = \frac{\prod_f M_f}{\prod_b M_b}$$

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$



$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, \cancel{M_b - 2i\sigma_0}$$

$$0 = \bar{\epsilon} \Psi_{rel}$$

$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

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$\mathcal{N} = 2$ Matter multiplet

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1) f f' \mathcal{D}_\theta - \dots$$

$$\Delta_f = -i f \gamma^\mu \mathcal{D}_\mu - i f' (\Delta - \frac{1}{2}) \gamma_3 - \dots$$

$$\frac{\det \Delta_f}{\det \Delta_b} = \frac{\prod_f M_f}{\prod_b M_b}$$

$$M_f = i\sigma_0 - \frac{m}{s} - n - \frac{\Delta}{2} \left(\frac{1}{s} + 1 \right)$$

$$(f \mathcal{D}_\theta + f' \Delta) \Phi = -i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi \Phi - \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi \Phi$$

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$

$$0 = \bar{\epsilon} \Psi_{rel}$$

$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, \cancel{M_b - 2i\sigma_0}$$

$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

$$M_b = i\sigma_0 + \frac{m}{s} + n - \frac{\Delta - 2}{2} \left(\frac{1}{s} + 1 \right)$$

$$(f \mathcal{D}_\theta + f' (\Delta + 1)) F = i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi F + \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi F$$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Matter multiplet

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1) f f' \mathcal{D}_\theta - \dots$$

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$$\frac{\det \Delta_f}{\det \Delta_b} = \frac{\prod_f M_f}{\prod_b M_b}$$

$$M_f = i\sigma_0 - \frac{m}{s} - n - \frac{\Delta}{2} \left(\frac{1}{s} + 1 \right)$$

$$(f \mathcal{D}_\theta + f' \Delta) \Phi = -i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi \Phi - \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi \Phi$$

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$

these terms does not affect

$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, \cancel{M_b - 2i\sigma_0}$$

$$M_b = i\sigma_0 + \frac{m}{s} + n - \frac{\Delta - 2}{2} \left(\frac{1}{s} + 1 \right)$$

$$(f \mathcal{D}_\theta + f' (\Delta + 1)) F = i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi F + \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi F$$

$$0 = \bar{\epsilon} \Psi_{rel}$$

$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Matter multiplet

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1) f f' \mathcal{D}_\theta - \dots$$

$$\Delta_f = -i f \gamma^\mu \mathcal{D}_\mu - i f' (\Delta - \frac{1}{2}) \gamma_3 - \dots$$

$$\frac{\det \Delta_f}{\det \Delta_b} = \frac{\prod_f M_f}{\prod_b M_b}$$

$$\Phi \sim \begin{cases} (\sin \theta)^{-n} & \theta \sim 0 \\ (\cos \theta)^{-m} & \theta \sim \frac{\pi}{2} \end{cases}$$

$$\rightarrow m \leq 0, n \leq 0$$

$$M_f = i\sigma_0 - \frac{m}{s} - n - \frac{\Delta}{2} \left(\frac{1}{s} + 1 \right)$$

$$(f \mathcal{D}_\theta + f' \Delta) \Phi = -i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi \Phi - \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi \Phi$$

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$

$$0 = \bar{\epsilon} \Psi_{rel}$$

$$M \Psi_1 = \Psi_2 : \Delta_f = \cancel{M_b}, \cancel{M_b - 2i\sigma_0}$$

$$\Delta_f \Psi_{rel} = M_f \Psi_{rel}$$

$$M_b = i\sigma_0 + \frac{m}{s} + n - \frac{\Delta - 2}{2} \left(\frac{1}{s} + 1 \right)$$

$$\left(f \mathcal{D}_\theta + f' (\Delta + 1) \right) F = i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi F + \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi F$$

$$F \sim \cos^m \theta \sin^n \theta$$

$$\rightarrow m, n \geq 0$$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Matter multiplet

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1) f f' \mathcal{D}_\theta - \dots$$

$$\Delta_f = -i f \gamma^\mu \mathcal{D}_\mu - i f' \left(\Delta - \frac{1}{2}\right) \gamma_3 - \dots$$

$$\frac{\det \Delta_f}{\det \Delta_b} = \frac{\prod_f M_f}{\prod_b M_b} = \prod_{\rho \in \text{weight}} s_b \left(\frac{i(\sqrt{\frac{1}{s}} + \sqrt{s})}{2} (1 - \Delta) - \sqrt{s\rho} \cdot \sigma_0 \right)$$

$$\Phi \sim \begin{cases} (\sin \theta)^{-n} & \theta \sim 0 \\ (\cos \theta)^{-m} & \theta \sim \frac{\pi}{2} \end{cases}$$

$$\rightarrow m \leq 0, n \leq 0$$

$$M_f = i\sigma_0 - \frac{m}{s} - n - \frac{\Delta}{2} \left(\frac{1}{s} + 1 \right)$$

$$(f\mathcal{D}_\theta + f'\Delta)\Phi = -i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi \Phi - \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi \Phi$$

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$

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$$M_b = i\sigma_0 + \frac{m}{s} + n - \frac{\Delta - 2}{2} \left(\frac{1}{s} + 1 \right)$$

$$\left(f\mathcal{D}_\theta + f'(\Delta + 1) \right) F = i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi F + \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi F$$

$$F \sim \cos^m \theta \sin^n \theta$$

$$\rightarrow m, n \geq 0$$

SUSY on round sphere

$$\epsilon = \begin{pmatrix} -(\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$$\bar{\epsilon} = \begin{pmatrix} (\cos \theta + i s \sin \theta)^{1/2} \\ (\cos \theta - i s \sin \theta)^{1/2} \end{pmatrix}$$

$\mathcal{N} = 2$ Matter multiplet

We recover the result on squashed sphere!

$$\Delta_b = -f^2 \mathcal{D}^2 - (2\Delta - 1) f f' \mathcal{D}_\theta - \dots$$

$$\Delta_f = -i f \gamma^\mu \mathcal{D}_\mu - i f' \left(\Delta - \frac{1}{2}\right) \gamma_3 - \dots$$

$$\frac{\det \Delta_f}{\det \Delta_b} = \frac{\prod_f M_f}{\prod_b M_b} = \prod_{\rho \in \text{weight}} s_b \left(\frac{i(\sqrt{\frac{1}{s}} + \sqrt{s})}{2} (1 - \Delta) - \sqrt{s\rho} \cdot \sigma_0 \right)$$

$$\Phi \sim \begin{cases} (\sin \theta)^{-n} & \theta \sim 0 \\ (\cos \theta)^{-m} & \theta \sim \frac{\pi}{2} \end{cases}$$

$$\rightarrow m \leq 0, n \leq 0$$

$$M_f = i\sigma_0 - \frac{m}{s} - n - \frac{\Delta}{2} \left(\frac{1}{s} + 1\right)$$

$$(f\mathcal{D}_\theta + f'\Delta)\Phi = -i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi \Phi - \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi \Phi$$

$$\Delta_b \Phi_{rel} = M_b (\cancel{M_b - 2i\sigma_0}) \Phi_{rel}$$

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$$M_b = i\sigma_0 + \frac{m}{s} + n - \frac{\Delta - 2}{2} \left(\frac{1}{s} + 1\right)$$

$$\left(f\mathcal{D}_\theta + f'(\Delta + 1)\right)F = i \frac{\sin \theta}{\cos \theta} \mathcal{D}_\phi F + \frac{i \cos \theta}{s \sin \theta} \mathcal{D}_\chi F$$

$$F \sim \cos^m \theta \sin^n \theta$$

$$\rightarrow m, n \geq 0$$