

Heterotic Line Bundle Models



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based on: [arXiv:1311.1941](https://arxiv.org/abs/1311.1941), [arXiv:1307.4787](https://arxiv.org/abs/1307.4787), [arXiv:1106.4804](https://arxiv.org/abs/1106.4804), [arXiv:1202.1757](https://arxiv.org/abs/1202.1757),
with Lara Anderson, Evgeny Buchbinder, Andrei Constantin, James Gray
and Eran Palti

Overview

- Introduction: Heterotic line bundle models
- Arena: Specific Calabi-Yau manifolds and line bundles
- An exhaustive scan over favourable Cicy's
- An example
- Continuation to non-Abelian bundles
- Conclusion and outlook

Introduction: Heterotic line bundle models

Data to define a heterotic line bundle model we need:

- A Calabi-Yau 3-fold X
- A line bundle sum $V = L_1 \oplus \cdots \oplus L_5$ on X ,
 $c_1(V) = 0$, so structure group is
 $S(U(1)^5) \subset SU(5) \subset E_8$
- vanishing slopes $\mu(L_a) \equiv c_1(L_a) \wedge J^2 \stackrel{!}{=} 0$
- Anomaly: $c_2(TX) - c_2(V) - c_2(\tilde{V}) = [C]$
in practice: $c_2(V) \leq c_2(TX)$

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standard-like model
(hopefully) with
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$\mathbf{1}_{\mathbf{e}_a - \mathbf{e}_b}$	$\mathbf{e}_a - \mathbf{e}_b$	$L_a \otimes L_b^*$	$V \otimes V^*$
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$\langle S^\alpha \rangle = 0$: line bundle model, $\langle S^\alpha \rangle \neq 0$: non-Abelian bundle

U(1) symmetries constrain 4d theory, e.g. superpotential:

$$W = \mu(S) H \bar{H} + Y_{pq}^{(d)}(S) H \bar{\mathbf{5}}^p \mathbf{10}^q + Y_{pq}^{(u)}(S) \mathbf{10}^p \mathbf{10}^q + \dots$$

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Two ways to explore non-Abelian bundles:

- VEVs $\langle S^\alpha \rangle \neq 0$, spontaneously breaks U(1)s
- Construct non-Abelian bundles which "split" to line bundle sum

Arena: complete intersection CY manifolds (CICYs)

CICYs defined as common zero locus $X = \{p_i = 0\} \subset \mathcal{A}$ of homogeneous polynomials p_i in ambient space $\mathcal{A} = \bigotimes_{r=1}^m \mathbb{P}^{n_r}$.

for example: quintic $X \sim [\mathbb{P}^4 | 5]$ or bi-cubic $X \sim \left[\begin{array}{c|c} \mathbb{P}^2 & 3 \\ \hline \mathbb{P}^2 & 3 \end{array} \right]$

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Complete classification of about 8000 spaces

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Classification of freely-acting discrete symmetries

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Line bundle cohomology can be computed.

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Line bundles on CY manifolds

Line bundles, L , are classified by their first Chern class:

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Write $L = \mathcal{O}_X(\mathbf{k})$ where $\mathbf{k} = (k^i)$ is an integer vector.

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Rank 5 line bundle sum:

$$V = \bigoplus_{a=1}^5 \mathcal{O}_X(\mathbf{k}_a) \qquad c_1(V) \sim \sum_{a=1}^5 \mathbf{k}_a \stackrel{!}{=} 0$$

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No a priori bounds on k_a^i , so for $-k_{\max} \leq k_a^i \leq k_{\max}$ we have

$$\sim (2k_{\max} + 1)^{4h^{1,1}(X)} \text{ line bundle sums } V$$

Last year:

Scan for favourable Cicys with $h^{1,1}(X) \leq 5$ (60 spaces) and

$$k_{\max} = 2, 3 \quad \longrightarrow \quad \sim 10^{12} \text{ bundles}$$

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*standard model: SM gauge group times (anomalous) U(1)s, exact MSSM matter spectrum, one or more pairs of Higgs doublets, no exotics charged under standard model group.

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Feasible because some constraints can be checked while line bundle sum is built up, e.g if

$$h^1(X, L) > 3|\Gamma|$$

we do not need to consider line bundle L (too many families).

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Scan over all (k_a^i) with $|k_a^i| \leq k_{\max}$ and find number of viable models as a function of k_{\max} :

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Table 6: *Number of models as a function of k_{\max} on CICYs with $h^{1,1}(X) = 6$. Total number of models: 41036*

$X, \Gamma $	$k_m = 1$	$k_m = 2$	$k_m = 3$	$k_m = 4$	$k_m = 5$	$k_m = 6$	$k_m = 7$	$k_m = 8$	$k_m = 9$	$k_m = 10, 11, 12, 13$
3413, 3	0	2278	2897	2906	2906	2906				
4190, 2	11	766	1175	1243	1246	1247	1249	1249	1249	
5273, 2	29	4895	7149	7738	7799	7810	7810	7810		
5302, 2	0	4314	5978	6360	6369	6369	6369			
5302, 4	0	11705	16988	17687	17793	17838	17868	17868	17868	
5425, 2	0	2381	3083	3305	3337	3337	3337			
5958, 2	0	148	224	240	253	253	253			
6655, 5	0	92	178	189	194	194	198	201	202	203
6738, 2	1	2733	4116	4346	4386	4393	4399	4399	4399	

Number of consistent SU(5) GUT models with correct indices:

$h^{1,1}(X)$	1	2	3	4	5	6	total
#models	0	0	6	552	21731	41036	63325

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Have started a similar programme on CY manifolds defined in toric 4-folds (Kreuzer-Skarke list) -> Chuang Sun's talk

An example

CY data: ■ Cicy 7862, Symmetry 3

$$X = \begin{pmatrix} 2 \\ 2 \\ 2 \\ 2 \end{pmatrix}$$

$$\eta(X) = -128 \quad h^{1,1}(X) = 4 \quad h^{2,1}(X) = 68 \quad c_2(\text{TX}) = \{24, 24, 24, 24\}$$

$$\kappa = 12 t_1 t_2 t_3 + 12 t_1 t_2 t_4 + 12 t_1 t_3 t_4 + 12 t_2 t_3 t_4$$

symmetry: 3 order: 4

Abelian: True block diagonal: True factors: {2, 2}

$$\text{Action on coordinates: } \left\{ \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & -1 \end{pmatrix}, \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \end{pmatrix} \right\}$$

Action on polynomials: {(1), (1)}

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$\mathbb{Z}_2 \times \mathbb{Z}_2$ generators

bundle data:

■ Basic properties

standard model? **True** massless U(1): **1** number of $5 \bar{5}$ pairs: **3** $c_2(V) = \{24, 8, 20, 12\}$

$$V: (k_a^i) = \begin{pmatrix} -1 & -1 & 0 & 1 & 1 \\ 0 & -3 & 1 & 1 & 1 \\ 0 & 2 & -1 & -1 & 0 \\ 1 & 2 & 0 & -1 & -2 \end{pmatrix}$$

Cohomology of V:

L_2	=	$\{-1, -3, 2, 2\}$	$h[L_2]$	=	$\{0, 8, 0, 0\}$	$h[L_2, R]$	=	$\{\{0, 0, 0, 0\}, \{2, 2, 2, 2\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
L_5	=	$\{1, 1, 0, -2\}$	$h[L_5]$	=	$\{0, 4, 0, 0\}$	$h[L_5, R]$	=	$\{\{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
$L_2 \times L_4$	=	$\{0, -2, 1, 1\}$	$h[L_2 \times L_4]$	=	$\{0, 4, 0, 0\}$	$h[L_2 \times L_4, R]$	=	$\{\{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
$L_2 \times L_5$	=	$\{0, -2, 2, 0\}$	$h[L_2 \times L_5]$	=	$\{0, 3, 3, 0\}$	$h[L_2 \times L_5, R]$	=	$\{\{0, 0, 0, 0\}, \{0, 1, 1, 1\}, \{0, 1, 1, 1\}, \{0, 0, 0, 0\}\}$
$L_4 \times L_5$	=	$\{2, 2, -1, -3\}$	$h[L_4 \times L_5]$	=	$\{0, 8, 0, 0\}$	$h[L_4 \times L_5, R]$	=	$\{\{0, 0, 0, 0\}, \{2, 2, 2, 2\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
$L_1 \times L_2^*$	=	$\{0, 3, -2, -1\}$	$h[L_1 \times L_2^*]$	=	$\{0, 0, 12, 0\}$	$h[L_1 \times L_2^*, R]$	=	$\{\{0, 0, 0, 0\}, \{0, 0, 0, 0\}, \{3, 3, 3, 3\}, \{0, 0, 0, 0\}\}$
$L_1 \times L_5^*$	=	$\{-2, -1, 0, 3\}$	$h[L_1 \times L_5^*]$	=	$\{0, 0, 12, 0\}$	$h[L_1 \times L_5^*, R]$	=	$\{\{0, 0, 0, 0\}, \{0, 0, 0, 0\}, \{3, 3, 3, 3\}, \{0, 0, 0, 0\}\}$
$L_2 \times L_3^*$	=	$\{-1, -4, 3, 2\}$	$h[L_2 \times L_3^*]$	=	$\{0, 20, 0, 0\}$	$h[L_2 \times L_3^*, R]$	=	$\{\{0, 0, 0, 0\}, \{5, 5, 5, 5\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
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$L_3 \times L_5^*$	=	$\{-1, 0, -1, 2\}$	$h[L_3 \times L_5^*]$	=	$\{0, 0, 4, 0\}$	$h[L_3 \times L_5^*, R]$	=	$\{\{0, 0, 0, 0\}, \{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}\}$

Wilson line: $\{\{0, 0\}, \{0, 1\}\}$ Equivariant structure: $\{\{0, 0\}, \{0, 0\}, \{0, 0\}, \{0, 0\}, \{0, 0\}\}$ Higgs pairs: **1**

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Transfer format: $\{\{6, 1, 1, 4, 6, 5, 9, 9, 8, 10, 1, 7, 17\}, \{6, 6, -1, -1, -1, -1\}\}$

$\text{rk}(Y^{(u)}) = \{2, 2\}$ $\text{rk}(Y^{(d)}) = \{0, 0\}$ dim. 4 operators absent: $\{\text{True}, \text{True}\}$ dim. 5 operators absent: $\{\text{True}, \text{True}\}$

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← integer matrix defining line bundle sum

Cohomology of V:

L_2	=	$\{-1, -3, 2, 2\}$	$h[L_2]$	=	$\{0, 8, 0, 0\}$	$h[L_2, R]$	=	$\{\{0, 0, 0, 0\}, \{2, 2, 2, 2\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
L_5	=	$\{1, 1, 0, -2\}$	$h[L_5]$	=	$\{0, 4, 0, 0\}$	$h[L_5, R]$	=	$\{\{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
$L_2 \times L_4$	=	$\{0, -2, 1, 1\}$	$h[L_2 \times L_4]$	=	$\{0, 4, 0, 0\}$	$h[L_2 \times L_4, R]$	=	$\{\{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
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$L_3 \times L_5^*$	=	$\{-1, 0, -1, 2\}$	$h[L_3 \times L_5^*]$	=	$\{0, 0, 4, 0\}$	$h[L_3 \times L_5^*, R]$	=	$\{\{0, 0, 0, 0\}, \{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}\}$

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$L_2 \times L_4$	=	$\{0, -2, 1, 1\}$	$h[L_2 \times L_4]$	=	$\{0, 4, 0, 0\}$	$h[L_2 \times L_4, R]$	=	$\{\{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}, \{0, 0, 0, 0\}\}$
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$L_1 \times L_5^*$	=	$\{-2, -1, 0, 3\}$	$h[L_1 \times L_5^*]$	=	$\{0, 0, 12, 0\}$	$h[L_1 \times L_5^*, R]$	=	$\{\{0, 0, 0, 0\}, \{0, 0, 0, 0\}, \{3, 3, 3, 3\}, \{0, 0, 0, 0\}\}$
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$L_3 \times L_5^*$	=	$\{-1, 0, -1, 2\}$	$h[L_3 \times L_5^*]$	=	$\{0, 0, 4, 0\}$	$h[L_3 \times L_5^*, R]$	=	$\{\{0, 0, 0, 0\}, \{0, 0, 0, 0\}, \{1, 1, 1, 1\}, \{0, 0, 0, 0\}\}$

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$3\ 1_{2,1}, 3\ 1_{5,1}, 5\ 1_{2,3}, 3\ 1_{2,4}, 1_{5,3}$

allowed operators:

■ Operators

basic superpotential terms:

$$\overline{H}10^p 10^q: Y^{(u)} = \begin{pmatrix} (0) & (0) & (1) \\ (0) & (0) & (1) \\ (1) & (1) & (0) \end{pmatrix}$$

$$H\overline{5}^p 10^q: Y^{(d)} = \begin{pmatrix} (0) & (0) & (0) \\ (0) & (0) & (0) \\ (0) & (0) & (0) \end{pmatrix}$$

$$H\overline{H}: \mu = \{1\}$$

$$W_{\text{sing}} = \{0\}$$

R-parity violating terms in superpotential:

$$\overline{H}L^p: \rho = \begin{pmatrix} 0 \\ S_{2,4} \\ S_{2,4} \end{pmatrix}$$

$$10^p \overline{5}^q \overline{5}^r: \lambda = \{ \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \} \}$$

Dimension 5 operators in superpotential:

$$\overline{5}^p 10^q 10^r 10^s: \lambda' = \{ \{ \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \} \}$$

D-terms:

$$\text{FI-terms: } k^i_{a\kappa_i} = \begin{pmatrix} 4 t_1 t_2 + 4 t_1 t_3 - 4 t_2 t_4 - 4 t_3 t_4 \\ 16 t_1 t_2 - 4 t_1 t_3 + 4 t_2 t_3 - 4 t_1 t_4 + 4 t_2 t_4 - 16 t_3 t_4 \\ -4 t_1 t_2 + 4 t_1 t_3 - 4 t_2 t_4 + 4 t_3 t_4 \\ -8 t_1 t_2 + 8 t_3 t_4 \\ -8 t_1 t_2 - 4 t_1 t_3 - 4 t_2 t_3 + 4 t_1 t_4 + 4 t_2 t_4 + 8 t_3 t_4 \end{pmatrix}$$

$$\text{singlet D-terms: } q_{\alpha a} S^\alpha \overline{S}^{\bar{\beta}} = \begin{pmatrix} -S_{2,1} S^\dagger_{2,1} - S_{5,1} S^\dagger_{5,1} \\ S_{2,1} S^\dagger_{2,1} + S_{2,3} S^\dagger_{2,3} + S_{2,4} S^\dagger_{2,4} \\ -S_{2,3} S^\dagger_{2,3} - S_{5,3} S^\dagger_{5,3} \\ -S_{2,4} S^\dagger_{2,4} \\ S_{5,1} S^\dagger_{5,1} + S_{5,3} S^\dagger_{5,3} \end{pmatrix}$$

allowed operators:

■ Operators

basic superpotential terms:

$$\overline{H}10^p 10^q: Y^{(u)} = \begin{pmatrix} (0) & (0) & (1) \\ (0) & (0) & (1) \\ (1) & (1) & (0) \end{pmatrix} \leftarrow \text{rank 2}$$

$$H\overline{5}^p 10^q: Y^{(d)} = \begin{pmatrix} (0) & (0) & (0) \\ (0) & (0) & (0) \\ (0) & (0) & (0) \end{pmatrix}$$

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Dimension 5 operators in superpotential:

$$\overline{5}^p 10^q 10^r 10^s: \lambda' = \{ \{ \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \}, \{ \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \}, \{ \{0\}, \{0\}, \{0\} \} \} \}$$

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$$\text{singlet D-terms: } q_{\alpha a} S^\alpha \overline{S}^\beta = \begin{pmatrix} -S_{2,1} S^\dagger_{2,1} - S_{5,1} S^\dagger_{5,1} \\ S_{2,1} S^\dagger_{2,1} + S_{2,3} S^\dagger_{2,3} + S_{2,4} S^\dagger_{2,4} \\ -S_{2,3} S^\dagger_{2,3} - S_{5,3} S^\dagger_{5,3} \\ -S_{2,4} S^\dagger_{2,4} \\ S_{5,1} S^\dagger_{5,1} + S_{5,3} S^\dagger_{5,3} \end{pmatrix}$$

allowed operators:

■ Operators

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$$H\overline{H}: \mu = \{1\}$$

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R-parity violating terms in superpotential:

$$\overline{H}L^p: \rho = \begin{pmatrix} 0 \\ S_{2,4} \\ S_{2,4} \end{pmatrix}$$

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Dimension 5 operators in superpotential:

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$$\text{FI-terms: } k^i_{a\kappa_i} = \begin{pmatrix} 4t_1 t_2 + 4t_1 t_3 - 4t_2 t_4 - 4t_3 t_4 \\ 16t_1 t_2 - 4t_1 t_3 + 4t_2 t_3 - 4t_1 t_4 + 4t_2 t_4 - 16t_3 t_4 \\ -4t_1 t_2 + 4t_1 t_3 - 4t_2 t_4 + 4t_3 t_4 \\ -8t_1 t_2 + 8t_3 t_4 \\ -8t_1 t_2 - 4t_1 t_3 - 4t_2 t_3 + 4t_1 t_4 + 4t_2 t_4 + 8t_3 t_4 \end{pmatrix}$$

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proton stable

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Continuation to non-Abelian bundles

Superpot. for example: $W = \lambda_i \bar{H}_{2,5} (Q_2^{(i)} u_5 + Q_5 u_2^{(i)}) + \rho_{\alpha i} \mathbf{1}_{2,4}^{(\alpha)} L_{4,5}^{(i)} \bar{H}_{2,5}$

Singlets for example: $3 \mathbf{1}_{2,1}, 3 \mathbf{1}_{5,1}, 5 \mathbf{1}_{2,3}, 3 \mathbf{1}_{2,4}, \mathbf{1}_{5,3}$

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Suggests: massless Higgs doublet pair throughout moduli space as long as $\langle \mathbf{1}_{2,4}^{\alpha} \rangle = 0$.

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What happens to bundle when singlets are switched on?

- All singlet VEVs non-zero: $V = \bigoplus_{a=1}^5 L_a \rightarrow \tilde{V}$
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Continuation to non-Abelian bundles

Superpot. for example: $W = \lambda_i \bar{H}_{2,5} (Q_2^{(i)} u_5 + Q_5 u_2^{(i)}) + \rho_{\alpha i} \mathbf{1}_{2,4}^{(\alpha)} L_{4,5}^{(i)} \bar{H}_{2,5}$

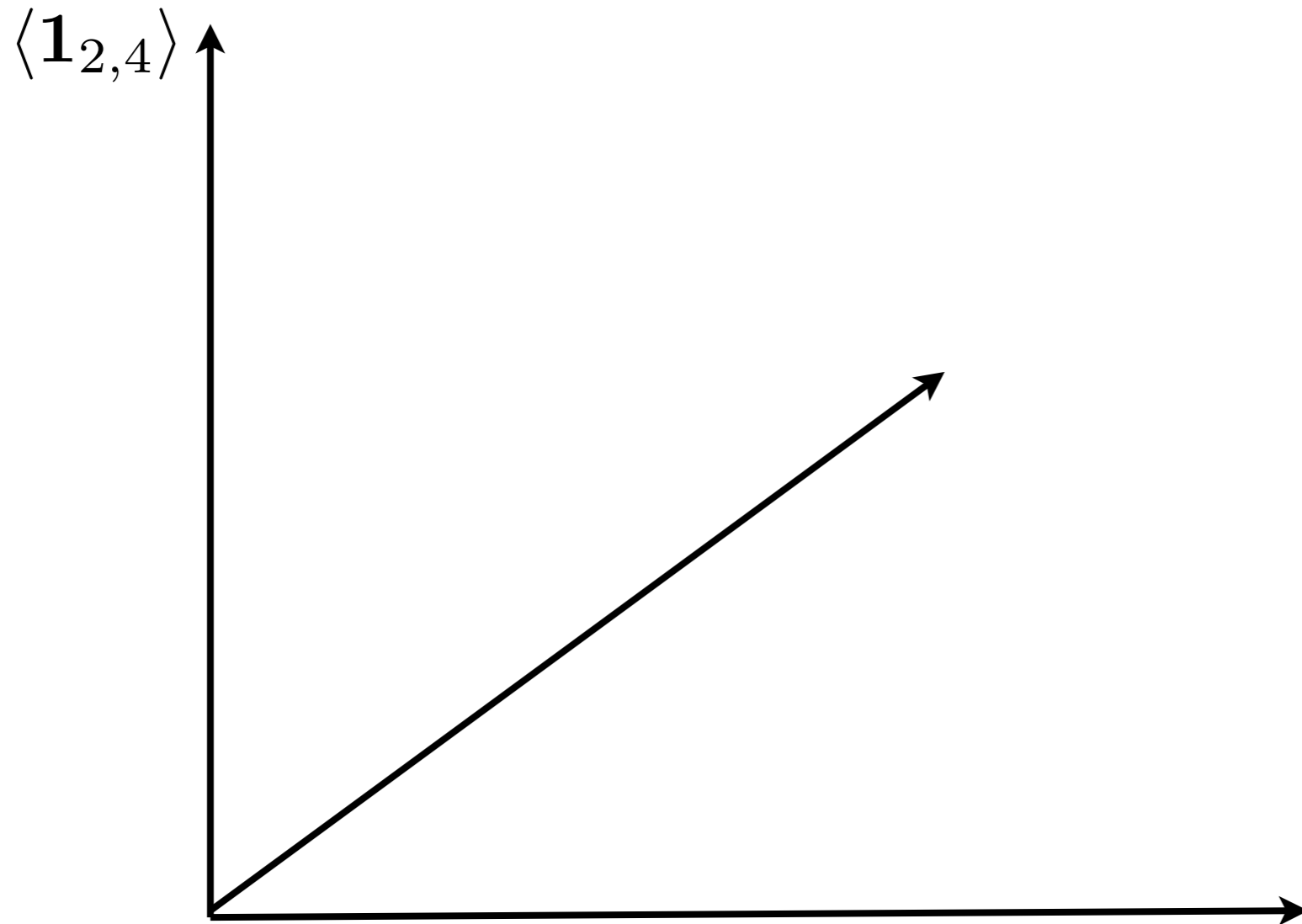
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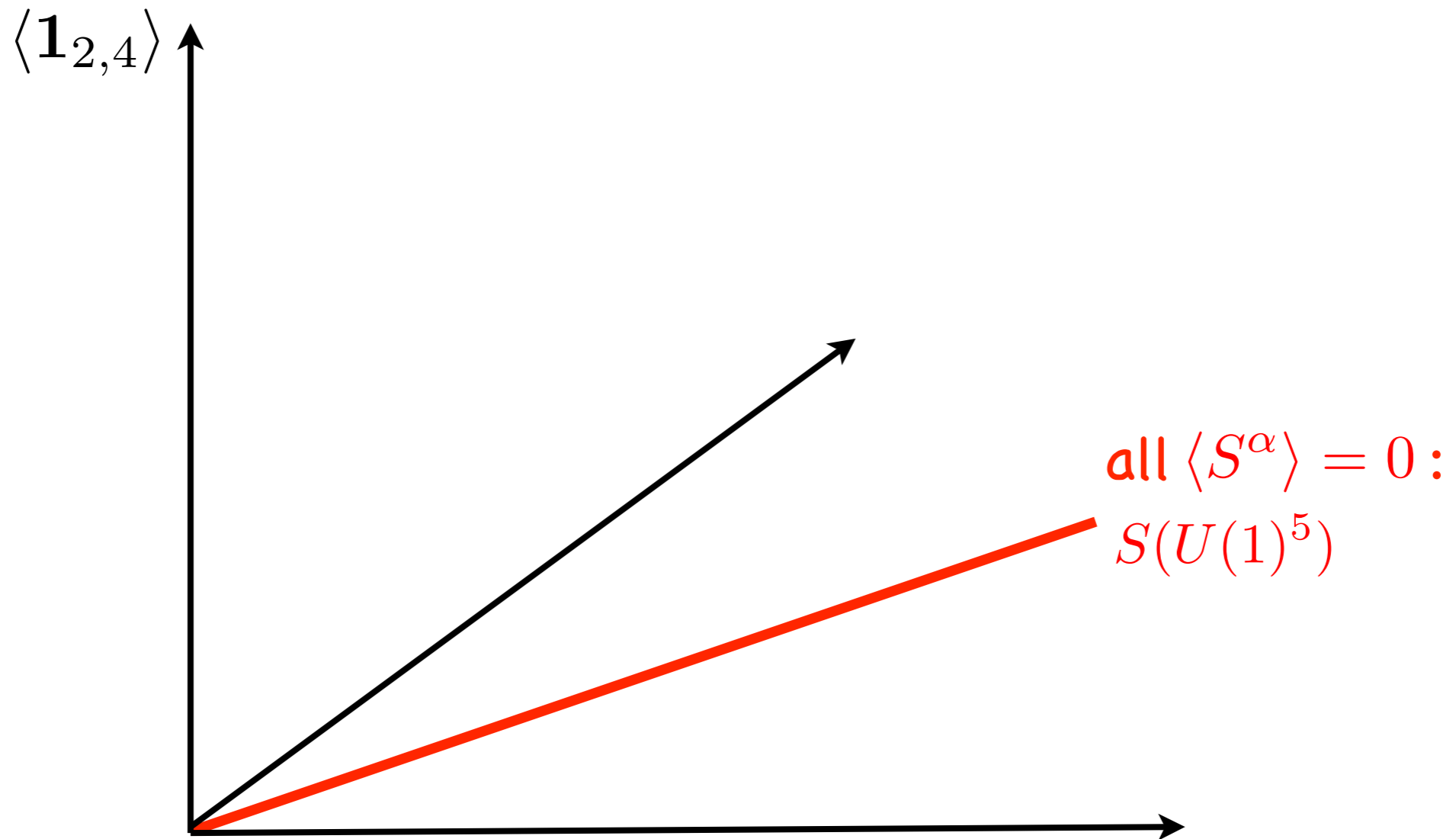
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 $U_{B-L}(1)$

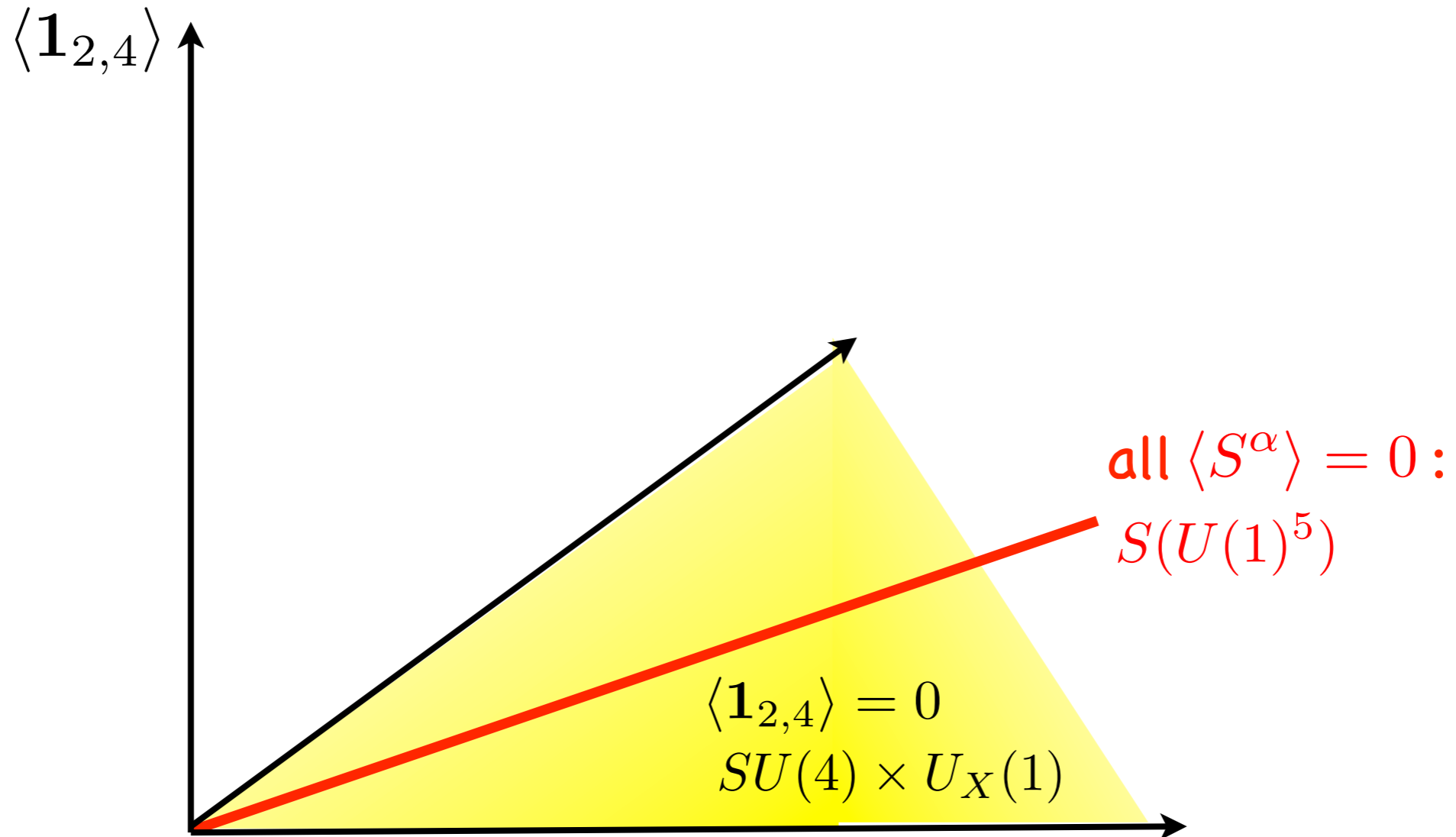
Schematic structure of moduli space:



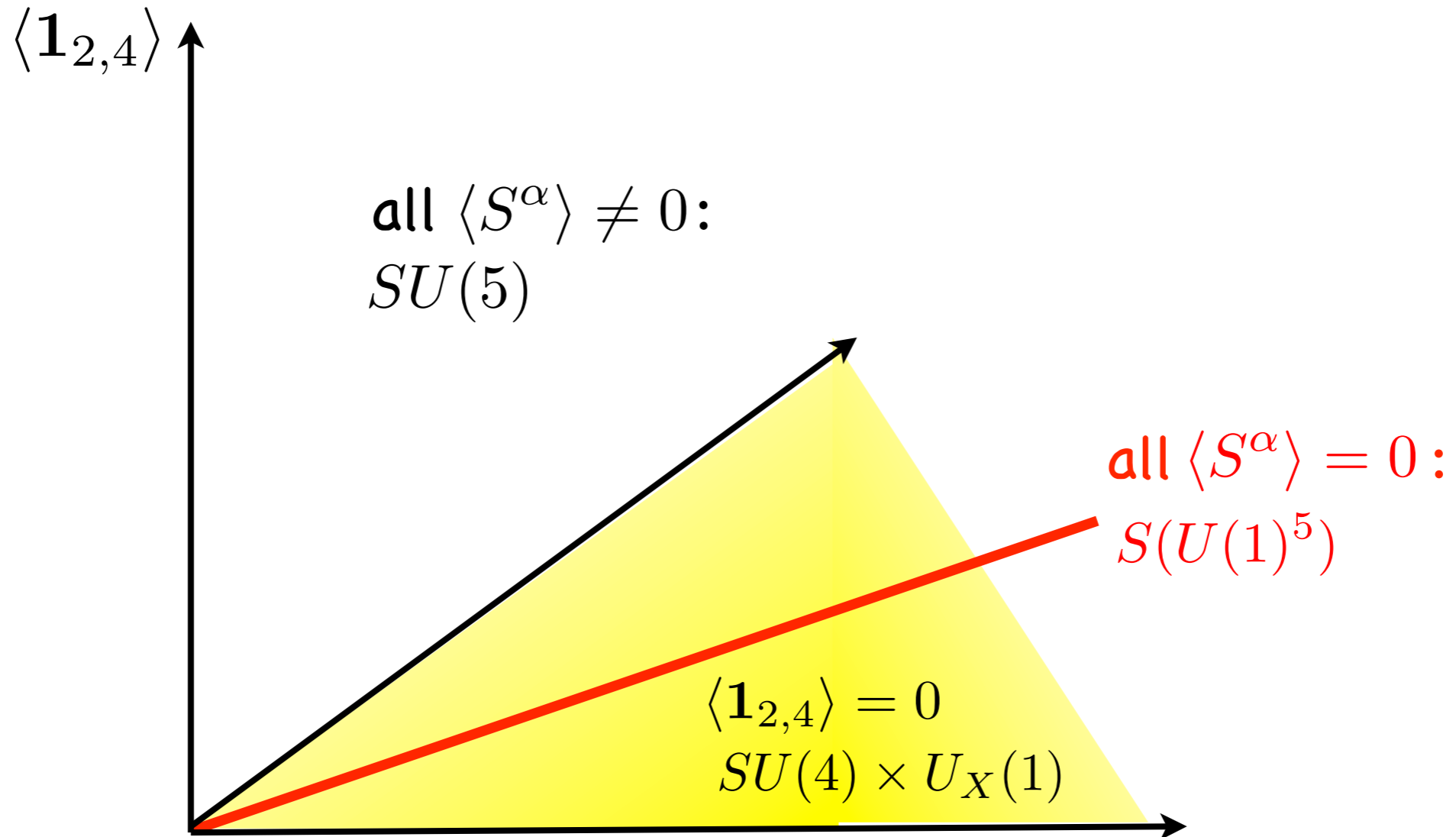
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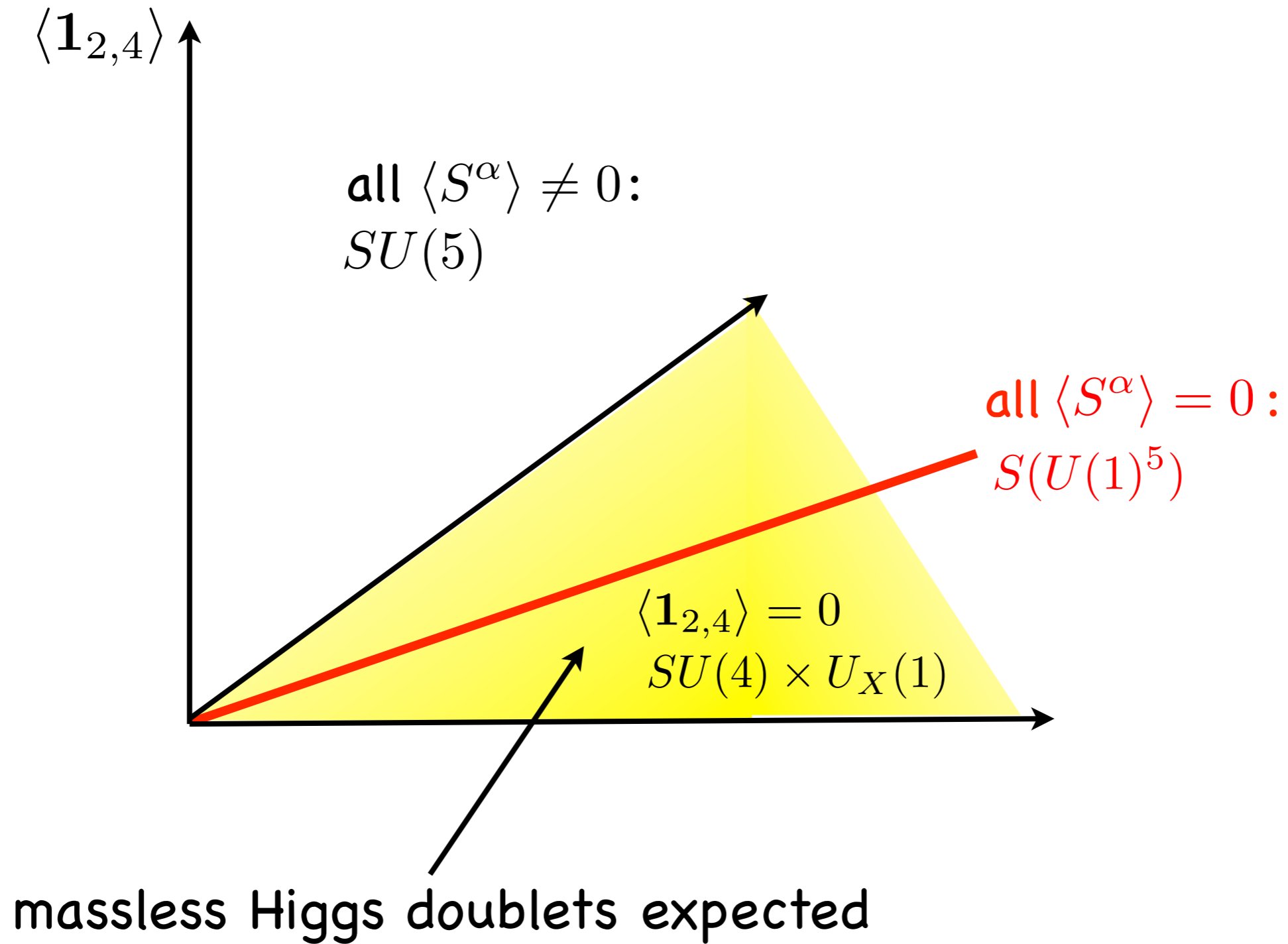
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Schematic structure of moduli space:



Check fate of Higgs by constructing non-Abelian bundle

Recall: $L_1 = \mathcal{O}_X(-1, 0, 0, 1)$, $L_2 = \mathcal{O}_X(-1, -3, 2, 2)$, $L_3 = \mathcal{O}_X(0, 1, -1, 0)$
 $L_4 = \mathcal{O}_X(1, 1, -1, -1)$, $L_5 = \mathcal{O}_X(1, 1, 0, -2)$

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1) Extension bundles

For $V_1 = L_2 \oplus L_5$, $V_2 = L_1 \oplus L_3 \oplus L_4$ define extension

$$0 \longrightarrow V_1 \longrightarrow \tilde{V} \longrightarrow V_2 \longrightarrow 0$$

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1) Extension bundles

For $V_1 = L_2 \oplus L_5$, $V_2 = L_1 \oplus L_3 \oplus L_4$ define extension

$$0 \longrightarrow V_1 \longrightarrow \tilde{V} \longrightarrow V_2 \longrightarrow 0$$

Compute $\#5 = h^2(X, \wedge^2 \tilde{V}) = \begin{cases} 3 & \text{for } \langle \mathbf{1}_{2,4} \rangle = 0 \\ 0 & \text{for } \langle \mathbf{1}_{2,4} \rangle \neq 0 \end{cases}$

2) Monads

$$0 \longrightarrow \tilde{V} \longrightarrow B \xrightarrow{f} C \longrightarrow 0$$

$$B \sim \begin{bmatrix} -1 & -1 & -1 & 0 & 1 & 1 & 1 \\ 0 & -1 & -1 & 1 & 1 & 1 & 1 \\ 0 & 2 & 2 & -1 & -1 & 0 & 0 \\ 1 & 2 & 2 & 0 & -1 & 0 & 0 \end{bmatrix}$$

$$C \sim \begin{bmatrix} -1 & 1 \\ 1 & 1 \\ 2 & 0 \\ 2 & 2 \end{bmatrix}$$

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L_1 L_2 L_3 L_4 L_5

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$$B \sim \begin{bmatrix} -1 & -1 & -1 & 0 & 1 & 1 & 1 \\ 0 & -1 & -1 & 1 & 1 & 1 & 1 \\ 0 & 2 & 2 & -1 & -1 & 0 & 0 \\ 1 & 2 & 2 & 0 & -1 & 0 & 0 \end{bmatrix}$$

$L_1 \quad L_2 \quad L_3 \quad L_4 \quad L_5$

$$C \sim \begin{bmatrix} -1 & 1 \\ 1 & 1 \\ 2 & 0 \\ 2 & 2 \end{bmatrix}$$

$$f \sim \begin{pmatrix} f_{(0,1,2,1)} & f_{(0,2,0,0)} & f'_{(0,2,0,0)} & 0 & 0 & 0 & 0 \\ f_{(2,1,0,1)} & 0 & 0 & f_{(1,0,1,2)} & f_{(0,0,1,3)} & f_{(0,0,0,2)} & f'_{(0,0,0,2)} \end{pmatrix}$$

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bundle splits if zero

We can show for $\langle \mathbf{1}_{2,4} \rangle = 0$, $\tilde{V} = U \oplus L_4$:

- bundle \tilde{V} is supersymmetric
- $\#5 = h^2(X, \wedge^2 \tilde{V}) = 3$

Features of the $SU(4) \times U_X(1)$ model:

- $U_X(1) \longrightarrow U_{B-L}(1)$
- μ -term forbidden
- dangerous dim. 4 terms forbidden by $U_{B-L}(1)$
- $\bar{5} 10 10 10$ operators still absent, due to existence of line bundle locus

Conclusions and outlook

- We can “mass-produce” heterotic CY standard models from line bundles. 2000 models have been found from 200 GUT models.
- We have found all viable line bundle GUT models on favourable Cicyes with freely-acting symmetries: **35000 models**
- These models will lead to a large number of standard models which will form the starting point for a detailed phenom. analysis.
- Higgs can be kept light away from the line bundle locus.
- We can study the non-Abelian continuation of line bundle models: “Unexpected” absences of operators due to line bundle locus.
- What is the total number of string standard models?

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Thanks