

Studies of  $b \rightarrow (s, d)(\mu^+ \mu^-, \gamma)$  transitions at LHCb

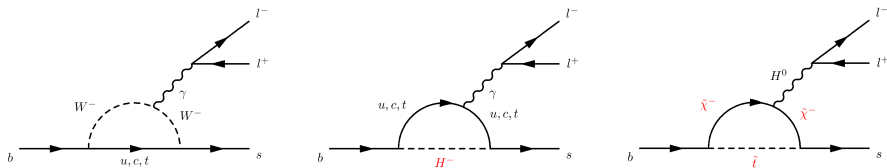
PASCOS 2013

Monday, November 25, 2013

Simon Wright, University of Cambridge  
on behalf of the LHCb collaboration

- $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ 
  - Angular analysis.
  - Results using form-factor independent observables.
- $B^+ \rightarrow K^+ \mu^+ \mu^-$ 
  - Observation of a low-recoil resonance.
  - Measurement of direct  $CP$  asymmetry.
- $B^+ \rightarrow K^+ \pi^- \pi^+ \gamma$ 
  - $CP$  and up-down asymmetries.

- Study of flavour changing neutral current decays that have no tree-level Feynman diagrams.
- Hence proceed via loop and box diagrams, and New Physics can enter through the loops.



- Theoretical framework via an effective Hamiltonian:
  - Wilson coefficients ( $C_i$ ), describing short-distance interactions
  - Operators, ( $\mathcal{O}_i$ ), describing long-distance interactions

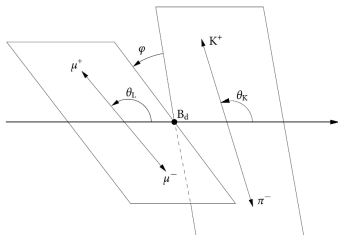
$$\mathcal{H}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_{i=1}^{10} (C_i^{\text{SM}} + \Delta C_i^{\text{NP}}) \mathcal{O}_i$$

$$B^0 \rightarrow K^{*0}(\rightarrow K^+\pi^-)\mu^+\mu^-$$

Decay distribution summing over  $B^0$  and  $\bar{B}^0$  mesons:

$$\begin{aligned} \frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{d\theta_I d\theta_K d\varphi dq^2} = & \frac{9}{32\pi} \left[ \frac{3}{4}(1 - F_L) \sin^2 \theta_K + F_L \cos^2 \theta_K + \frac{1}{4}(1 - F_L) \sin^2 \theta_K \cos 2\theta_I \right. \\ & - F_L \cos^2 \theta_K \cos 2\theta_I + S_3 \sin^2 \theta_K \sin^2 \theta_I \cos 2\varphi \\ & + S_4 \sin 2\theta_K \sin 2\theta_I \cos \varphi + S_5 \sin 2\theta_K \sin \theta_I \cos \varphi \\ & + S_6^s \sin^2 \theta_K \cos \theta_I + S_7 \sin 2\theta_K \sin \theta_I \sin \varphi \\ & \left. + S_8 \sin 2\theta_K \sin 2\theta_I \sin \varphi + S_9 \sin^2 \theta_K \sin^2 \theta_I \sin 2\varphi \right] \end{aligned}$$

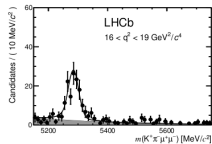
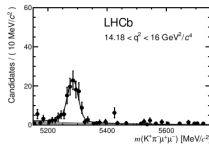
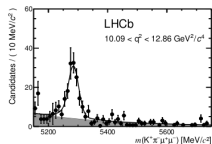
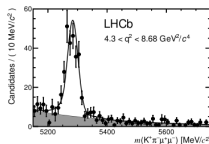
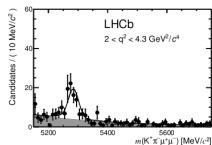
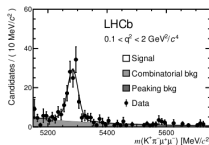
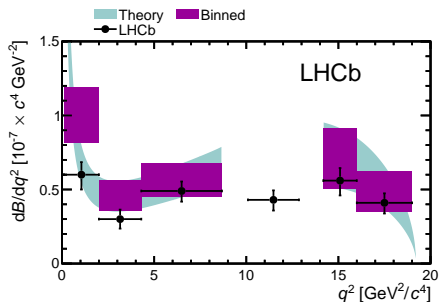
(from Altmannshofer et al JHEP 01 (2009) 019)



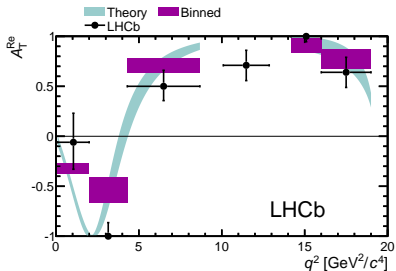
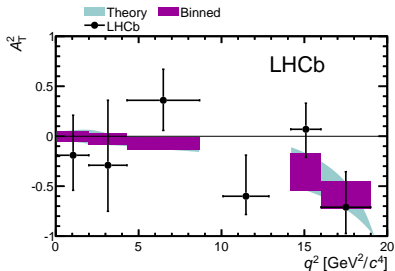
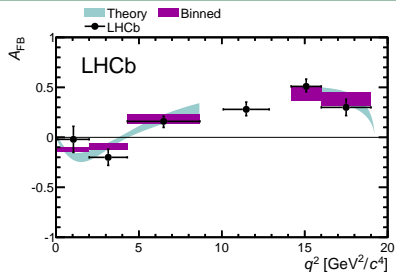
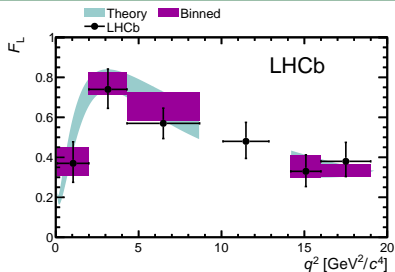
Observables include:

- $F_L$ , the  $K^{*0}$  longitudinal polarisation fraction.
- $A_{\text{FB}} = \frac{4}{3} S_6^s$ , the  $\mu^+\mu^-$  forward-backward asymmetry.
- $A_{\text{T}}^2 = 2S_3/(1 - F_L)$  and  $A_{\text{T}}^{\text{Re}} = \frac{4}{3} A_{\text{FB}}/(1 - F_L)$ , a pair of  $K^{*0}$  transverse asymmetries.

- Differential branching fraction obtained by performing an extended unbinned maximum likelihood fit to the  $K^+ \pi^- \mu^+ \mu^-$  invariant mass distribution in each of six  $q^2$  bins.

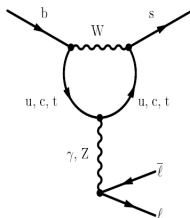
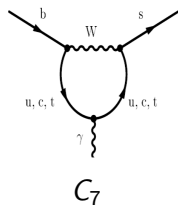


# Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (JHEP 1308 (2013) 131)



■ Good agreement with theory predictions.

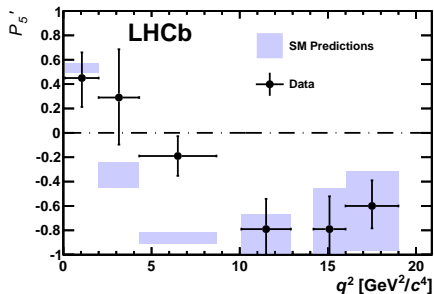
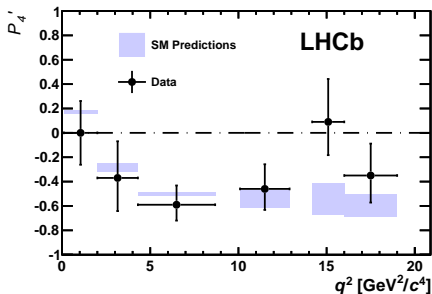
- The large theoretical uncertainties on these observables are, in part, due to large contributions the hadronic form factors.
- Combinations of  $F_L$  and  $S_i$  can have reduced form factor uncertainties.
- At large recoil (low  $q^2$ ), the combination  $P'_{i=4,5,6,8} = \frac{S_{i=4,5,7,8}}{\sqrt{F_L(1-F_L)}}$  is largely free of these uncertainties (arXiv:1303.5794).
- These observables are sensitive to New Physics in the Wilson coefficients  $C_7$ ,  $C_9$  and  $C_{10}$ :



$C_9$  = vector  
component

$C_{10}$  = axial-vector  
component

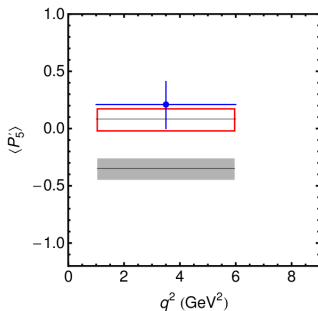
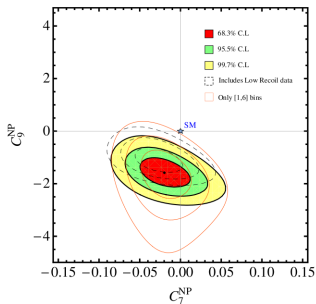
- $P'_6$  and  $P'_8$  found to be close to their SM predictions, which are close to zero, over the full  $q^2$  range.



- In the  $4.30 < q^2 < 8.68 \text{ GeV}^2/c^4$  bin for  $P'_5$ , there is a  $3.7\sigma$  discrepancy between the measurement and the prediction.
- Considering 24 independent measurements, the significance drops to  $2.8\sigma$ .
- In the range  $1.0 < q^2 < 6.0 \text{ GeV}^2/c^4$ , the deviation from SM is  $2.5\sigma$ .



# New observables in $B^0 \rightarrow K^{*0} \mu^+ \mu^-$



Data

SM prediction

$$C_9^{\text{NP}} = -1.5$$

(Descotes-Genon et al, PRD 88, 074002 (2013))

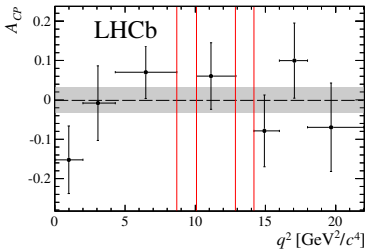
- Can potentially explain this discrepancy by a negative New Physics contribution to the Wilson coefficient  $C_9$ .
- However, this cannot occur in models such as MSSM or partial-compositeness (Altmannshofer and Straub, arXiv:1308.1501)...
- ...but it could in models with a flavour-changing neutral gauge boson, a  $Z'$ , with  $m_{Z'} \sim 7 \text{ TeV}$ . (Gauld et al, arXiv:1308.1959; Buras and Girschbach, arXiv:1309.2466)

$$\mathcal{A}_{CP}(q^2) = \frac{\Gamma(B^- \rightarrow K^- \mu^+ \mu^-) - \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}{\Gamma(B^- \rightarrow K^- \mu^+ \mu^-) + \Gamma(B^+ \rightarrow K^+ \mu^+ \mu^-)}$$

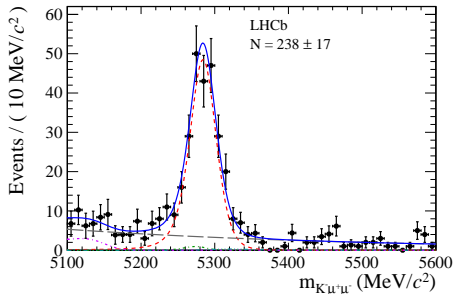
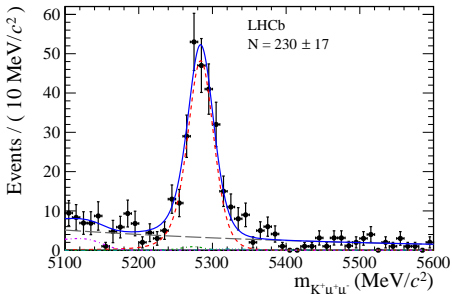
- $\mathcal{A}_{CP} \sim 10^{-4}$  in the Standard Model, and should be similar to  $\mathcal{A}_{CP}(B^0 \rightarrow K^{*0} \mu^+ \mu^-) = -0.072 \pm 0.040$  (PRL 110 031801).
- Analysis performed using 2011 LHCb data set ( $1.0 \text{ fb}^{-1}$ ).
- Use  $B^+ \rightarrow J/\psi K^+$  as a control channel to account for production and detection asymmetries:

$$\mathcal{A}_{CP} = \mathcal{A}_{RAW}(B^+ \rightarrow K^+ \mu^+ \mu^-) - \mathcal{A}_{RAW}(B^+ \rightarrow J/\psi K^+) + \mathcal{A}_{CP}(B^+ \rightarrow J/\psi K^+).$$

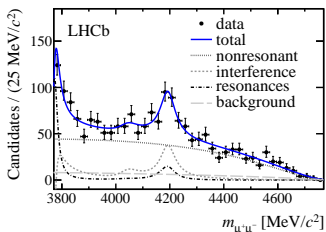
- $CP$  asymmetry extracted from simultaneous unbinned likelihood fit of  $B^+ \rightarrow J/\psi K^+$  and  $B^+ \rightarrow K^+ \mu^+ \mu^-$   $m_{K\mu\mu}$  distributions in bins of  $q^2$ .
- Average of results for both magnet polarities taken to reduce detector effects.



- $\mathcal{A}_{CP}$  over the full  $q^2$  range is the average of each  $q^2$  bin weighted by signal yield and efficiency.
- $\mathcal{A}_{CP} = 0.000 \pm 0.033(\text{stat.}) \pm 0.005(\text{syst.}) \pm 0.007(J/\psi K)$ .
- World's best measurement by a factor of 4, and consistent with both SM and  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  measurement.
- Mass fits for one polarity below:



- Using both the 2011 and 2012 data sets ( $3 \text{ fb}^{-1}$ ), look at low-recoil region (high  $q^2$ ).
- For high  $q^2$ , able to investigate the structure of the resonances coming from above the open charm threshold.

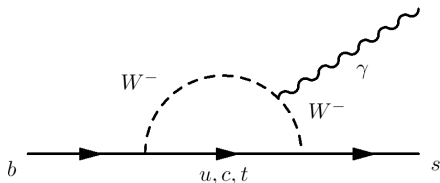


	Unconstrained	$\psi(4160)$
$\mathcal{B}[\times 10^{-9}]$	$3.9^{+0.7}_{-0.6}$	$3.5^{+0.9}_{-0.8}$
Mass [MeV/c <sup>2</sup> ]	$4191^{+9}_{-8}$	$4190 \pm 5$
Width [MeV/c <sup>2</sup> ]	$65^{+22}_{-16}$	$66 \pm 12$
Phase [rad]	$-1.7 \pm 0.3$	$-1.8 \pm 0.3$

Resonance at low-recoil observed with significance  $> 6\sigma$ .

Consistent with  $\psi(4160)$  resonance seen by the BES collaboration.

- Radiative decays are also mediated by penguin diagrams in the SM.
- However, their signature is a high- $E_T$  photon in the final state.



- For the above  $b \rightarrow s\gamma$  transition, the SM photon is predominantly left-handed, with weak amplitudes satisfying  $|c_L|^2 \gg |c_R|^2$ .
- Define photon polarisation  $\lambda_\gamma$ :

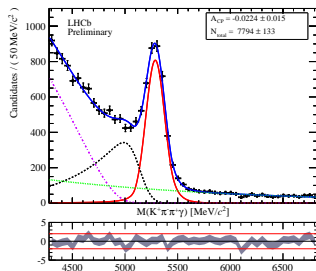
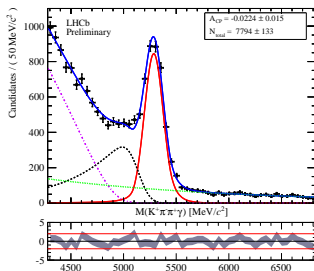
$$\lambda_\gamma = \frac{|c_R|^2 - |c_L|^2}{|c_R|^2 + |c_L|^2}$$

so that  $\lambda_\gamma \simeq -1$  ( $+1$ ) for  $B^-$  ( $B^+$ ) decays.

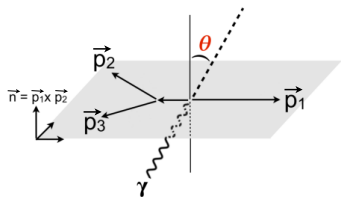
- However, several BSM models predict that the photon acquires a significant right-handed component.

- Consider  $B \rightarrow K_{\text{res}}(\rightarrow K^+ \pi^- \pi^+) \gamma$  decays corresponding to  $\mathcal{L} = 2 \text{ fb}^{-1}$ .
- $\mathcal{A}_{CP}$  is calculated similarly to  $B^+ \rightarrow K^+ \mu^+ \mu^-$ , using  $B^+ \rightarrow J/\psi K^+$  as a control mode again to account for production and detection asymmetries.

Partially  
reconstructed  
background  
Missing pion  
background  
Combinatorial  
background  
Signal



$$\mathcal{A}_{CP} = -0.007 \pm 0.015(\text{stat.}) \pm 0.008(\text{syst.})$$

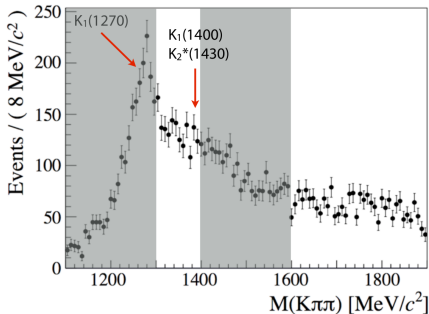


- Also measure the up-down asymmetry:

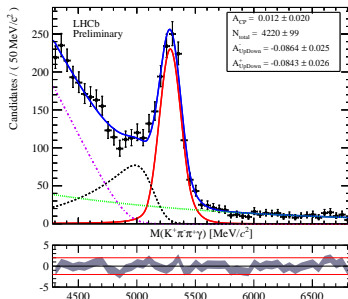
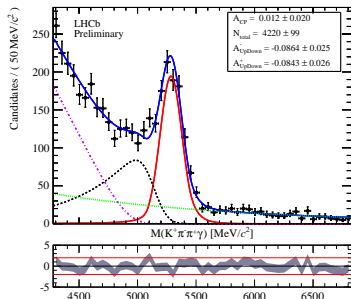
$$\mathcal{A}_{ud} = \frac{\int_0^1 d \cos \theta \frac{d\Gamma}{d \cos \theta} - \int_{-1}^0 d \cos \theta \frac{d\Gamma}{d \cos \theta}}{\int_{-1}^1 d \cos \theta \frac{d\Gamma}{d \cos \theta}}$$

$$= \frac{3}{4} \lambda_\gamma \frac{\int ds ds_{13} ds_{23} \text{Im}[n \cdot (\mathcal{J} \times \mathcal{J}^*)]}{\int ds ds_{13} ds_{23} |\mathcal{J}^2|}$$

- So  $\mathcal{A}_{ud} \propto \lambda_\gamma$ , and we extract the significance compared to a no polarisation hypothesis.
- Perform mass fits for  $\mathcal{A}_{ud}$  simultaneously on separated  $B^+$  and  $B^-$  data sets.
- Choose specific region in  $m_{K^+ \pi^- \pi^+}$  to conduct analysis (PRD 66 (2002) 054008).



## Mass fits for $B^+$ mesons:



$$\mathcal{A}_{ud}^+ = -0.084 \pm 0.026(\text{stat.}) \pm \begin{matrix} 0.003 \\ 0.004 \end{matrix}(\text{syst.})$$

$$\mathcal{A}_{ud}^- = -0.086 \pm 0.025(\text{stat.}) \pm 0.002(\text{syst.})$$

- These are significances of  $3.2\sigma$  and  $3.4\sigma$  from the no polarisation hypothesis.
- When combined, we obtain the first evidence of photon polarisation in  $b \rightarrow s\gamma$  with a significance of  $4.6\sigma$

$$\mathcal{A}_{ud} = -0.085 \pm 0.019(\text{stat.}) \pm 0.003(\text{syst.})$$



- With the 2011 data set ( $1 \text{ fb}^{-1}$  at  $\sqrt{s} = 7 \text{ TeV}$ ), LHCb has produced multiple interesting results:
  - Angular analyses of  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$  and  $B_s^0 \rightarrow \phi \mu^+ \mu^-$ , highlighting a potential discrepancy in the observable  $P'_5$ .
  - World's best measurement of  $CP$  asymmetry in  $B^+ \rightarrow K^+ \mu^+ \mu^-$ .
  - The most accurate measurement of  $\mathcal{B}(\Lambda_b \rightarrow \Lambda \mu^+ \mu^-)$ .
- Using the  $2 \text{ fb}^{-1}$  of  $\sqrt{s} = 8 \text{ TeV}$  data taken in 2012, evidence has been found for photon polarisation in  $b \rightarrow s \gamma$  transitions.
- And combining both data sets, a resonance at low-recoil has been observed in  $B^+ \rightarrow K^+ \mu^+ \mu^-$  decays.
- This is just the start of the improvements and new discoveries that could arise from the full analysis of the 2012 data!

# Happy St. Catharine's Day!

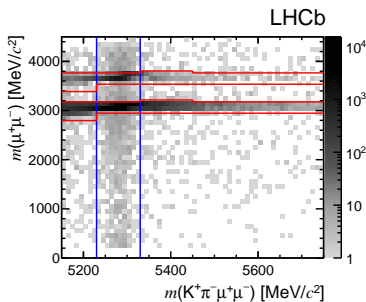


[www.caths.cam.ac.uk](http://www.caths.cam.ac.uk)

Catherine of Alexandria



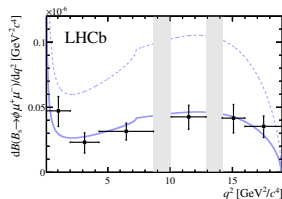
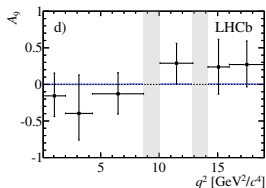
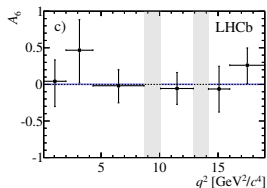
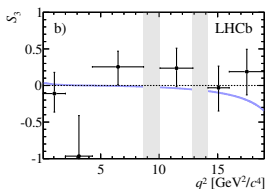
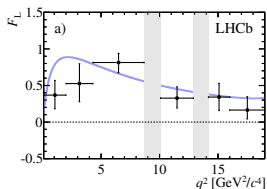
BACKUP



- Analysis performed in six bins of  $q^2$ , as well as the region  $1 < q^2 < 6 \text{ GeV}^2/c^4$ .
- The charmonium resonance regions corresponding to  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow \psi(2S) K^{*0}$  are vetoed.

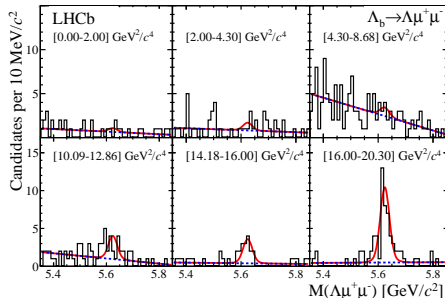
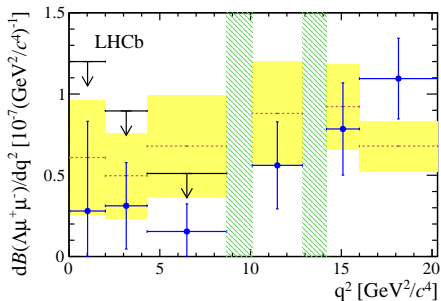
- The signal is selected using a boosted decision tree, and several additional vetoes are implemented to remove peaking backgrounds.
- An acceptance correction, compensating for discrepancies in the angular distributions caused by the selection, is performed using simulated Monte Carlo samples.
- The mode  $B^0 \rightarrow J/\psi K^{*0}$  is used to improve the agreement between data and simulation.

- Physics similar to  $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ , but final state does not self-tag.
- Perform a very similar angular analysis.



- $B_s^0 \rightarrow \phi \mu^+ \mu^-$  branching fraction lower than the Standard Model (dashed line).
- Angular observables agree with SM predictions.

- Analysis on baryonic modes can probe helicity structure of Hamiltonian, and also different hadronic physics than the  $B$  meson decays.
- Using  $1 \text{ fb}^{-1}$ , start by measuring the differential branching fraction, using the control mode  $\Lambda_b \rightarrow J/\psi \Lambda$ .



$78 \pm 12$   $\Lambda_b \rightarrow \Lambda \mu^+ \mu^-$  decays found, predominantly at high  $q^2$ .

$$\mathcal{B}(\Lambda_b \rightarrow \Lambda \mu^+ \mu^-) = 0.96 \pm 0.16(\text{stat.}) \pm 0.13(\text{syst.}) \pm 0.21(\text{norm.})$$