

Supersymmetric Dark Matter post run I at the LHC

Which Supersymmetric Model?

- ✦ MSSM with R-Parity (still more than 100 parameters)

Which Supersymmetric Model?

- ✦ MSSM with R-Parity (still more than 100 parameters)
- ✦ Gaugino mass Unification : $m_{1/2}$

Which Supersymmetric Model?

- ✦ MSSM with R-Parity (still more than 100 parameters)
- ✦ Gaugino mass Unification : $m_{1/2}$
- ✦ A-term Unification : A_0

Which Supersymmetric Model?

- ✦ MSSM with R-Parity (still more than 100 parameters)
- ✦ Gaugino mass Unification : $m_{1/2}$
- ✦ A-term Unification : A_0
- ✦ Scalar mass unification : m_0

Which Supersymmetric Model?

- ✦ MSSM with R-Parity (still more than 100 parameters)
- ✦ Gaugino mass Unification : $m_{1/2}$
- ✦ A-term Unification : A_0
- ✦ Scalar mass unification : m_0

CMSSM

Which Supersymmetric Model?

- ✦ MSSM with R-Parity (still more than 100 parameters)
- ✦ Gaugino mass Unification : $m_{1/2}$
- ✦ A-term Unification : A_0
- ✦ Scalar mass unification : m_0

CMSSM

add

parameter for
ratio of Higgs
vevs: $\tan \beta$

Which Supersymmetric Model?

- ✦ CMSSM (4+ parameters)
- ✦ mSUGRA (3+ parameters)
- ✦ NUHM (5,6+ parameters)
- ✦ (mini) Split SUSY (2+ parameters)

The CMSSM

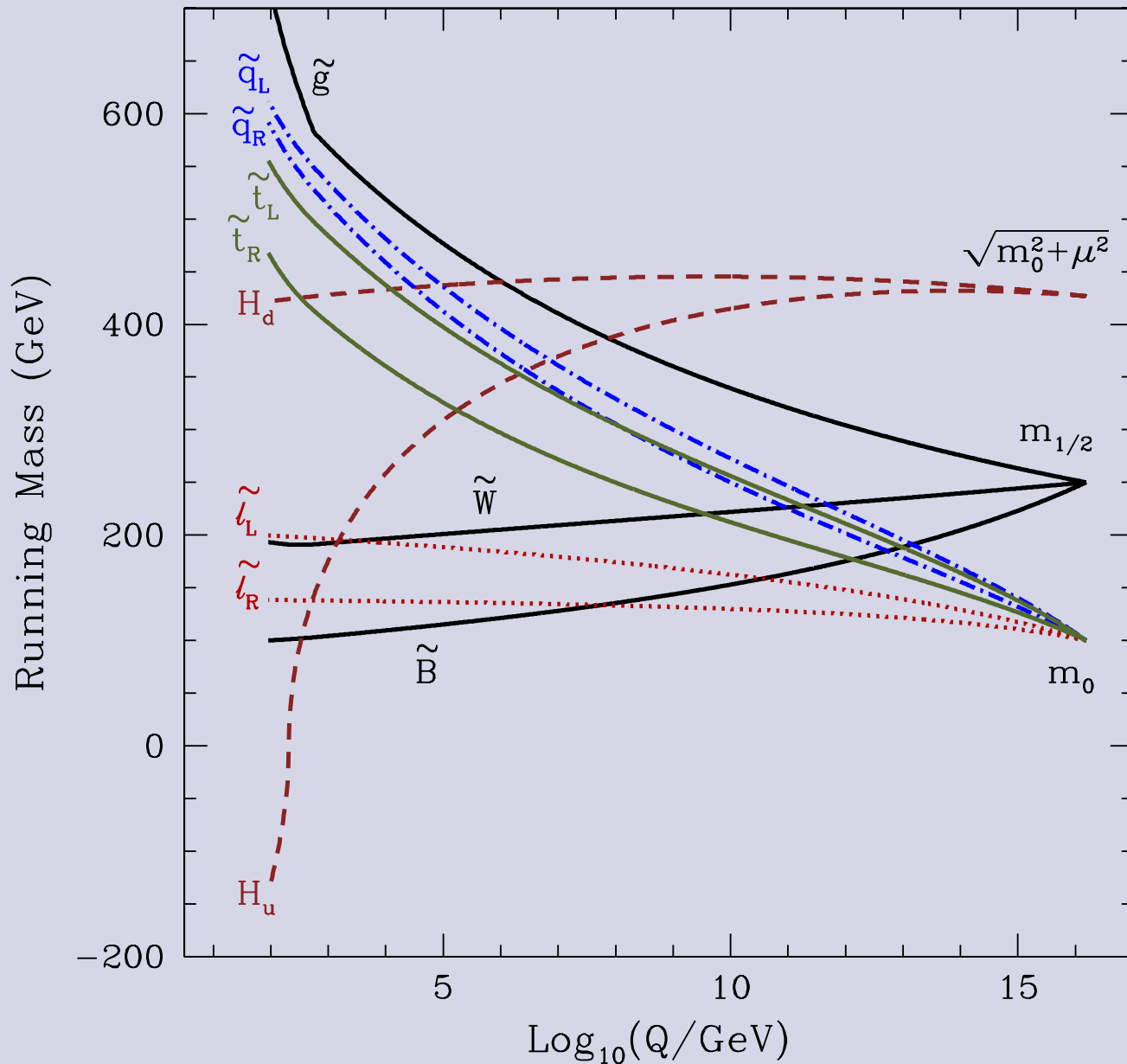
Parameters: $m_{1/2}$, m_0 , A_0 , $\tan \beta$, $\text{sgn}(\mu)$ $\{m_{3/2}\}$

Electroweak Symmetry Breaking conditions:

$$\mu^2 = \frac{m_1^2 - m_2^2 \tan^2 \beta + \frac{1}{2} M_Z^2 (1 - \tan^2 \beta) + \Delta_\mu^{(1)}}{\tan^2 \beta - 1 + \Delta_\mu^{(2)}}$$

$$B\mu = -\frac{1}{2}(m_1^2 + m_2^2 + 2\mu^2) \sin 2\beta + \Delta_B$$

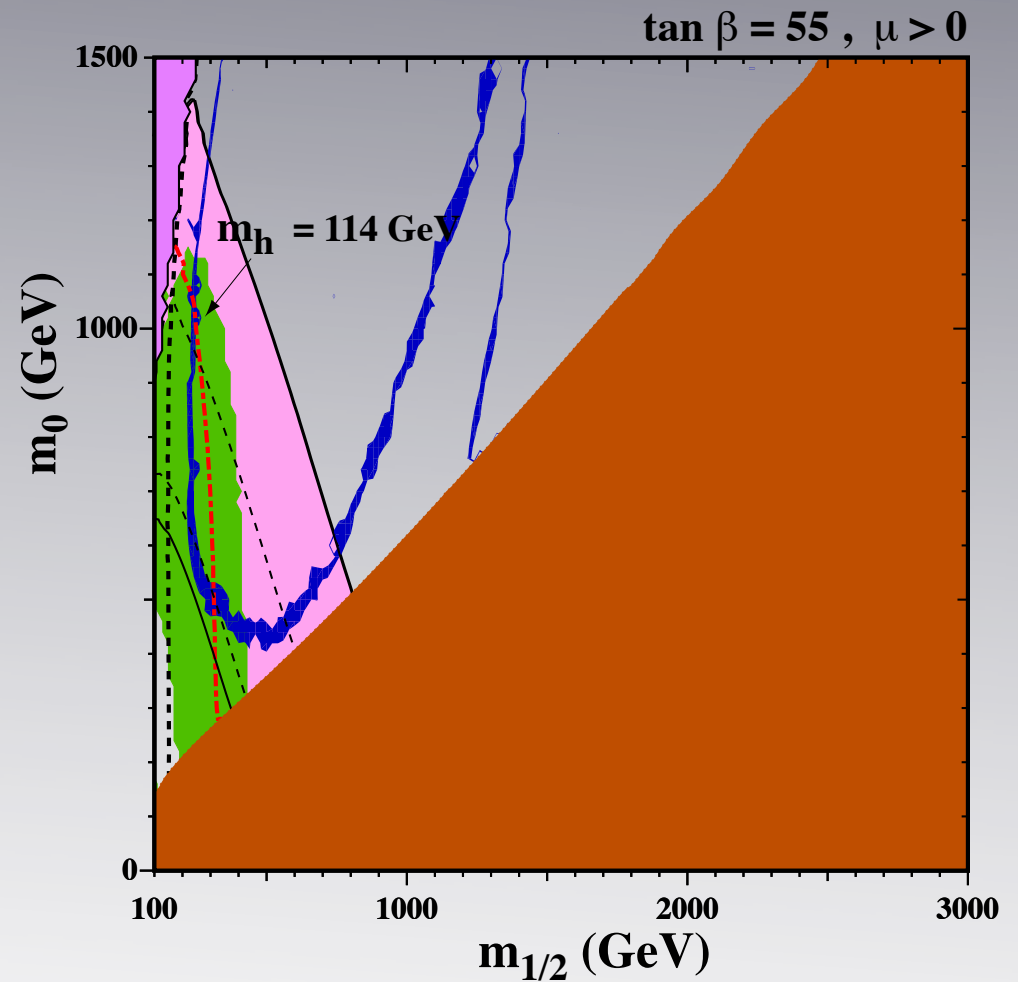
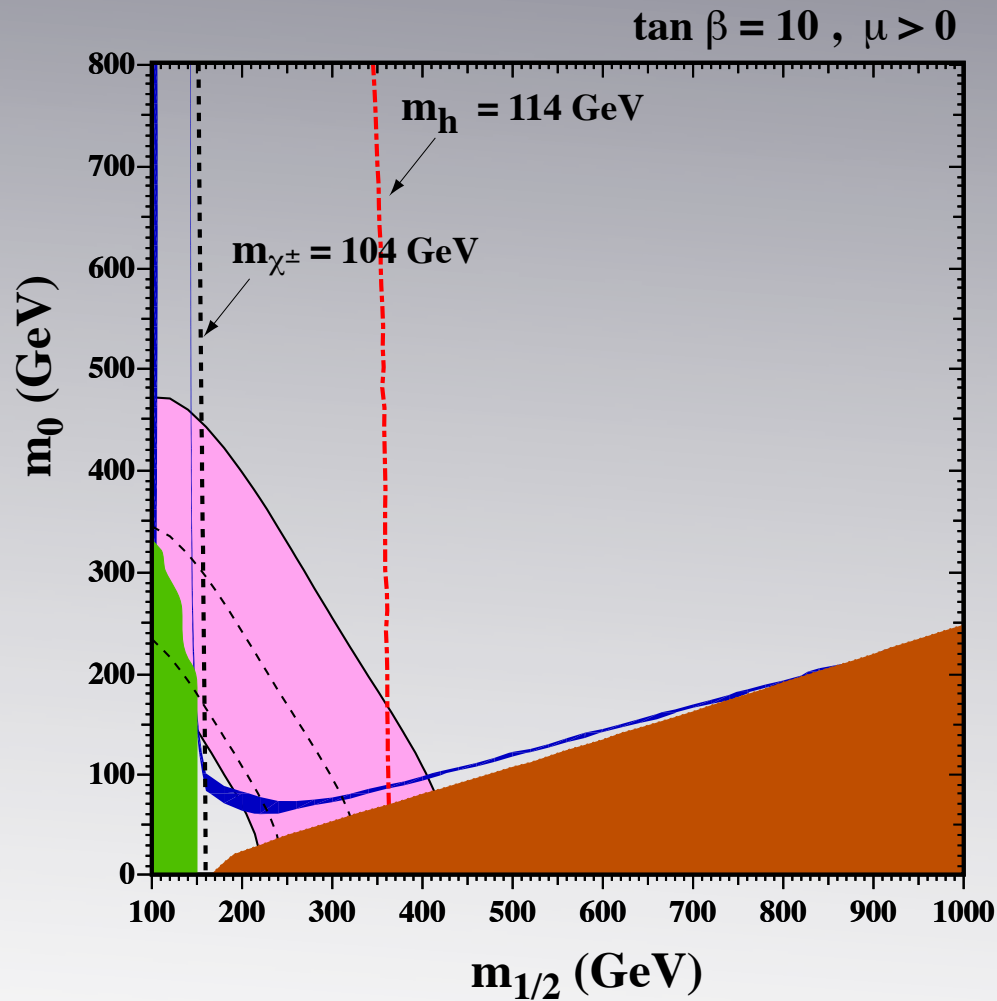
CMSSM Spectra



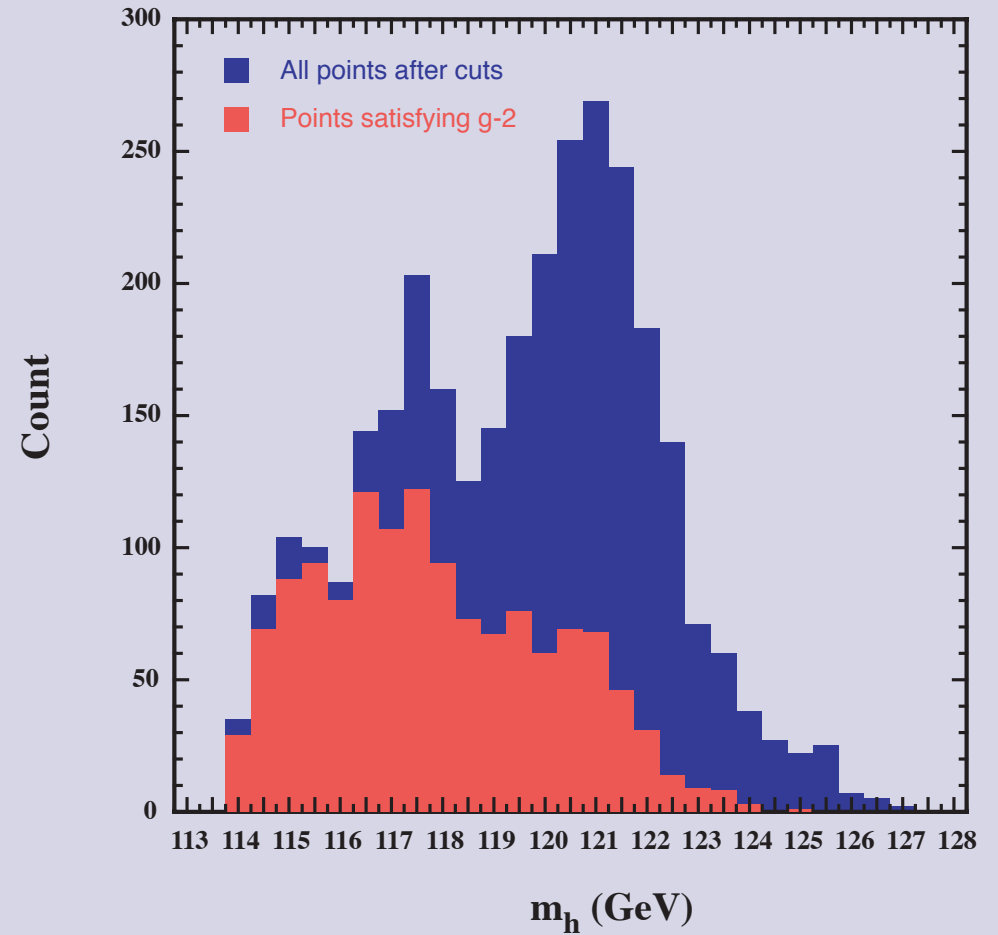
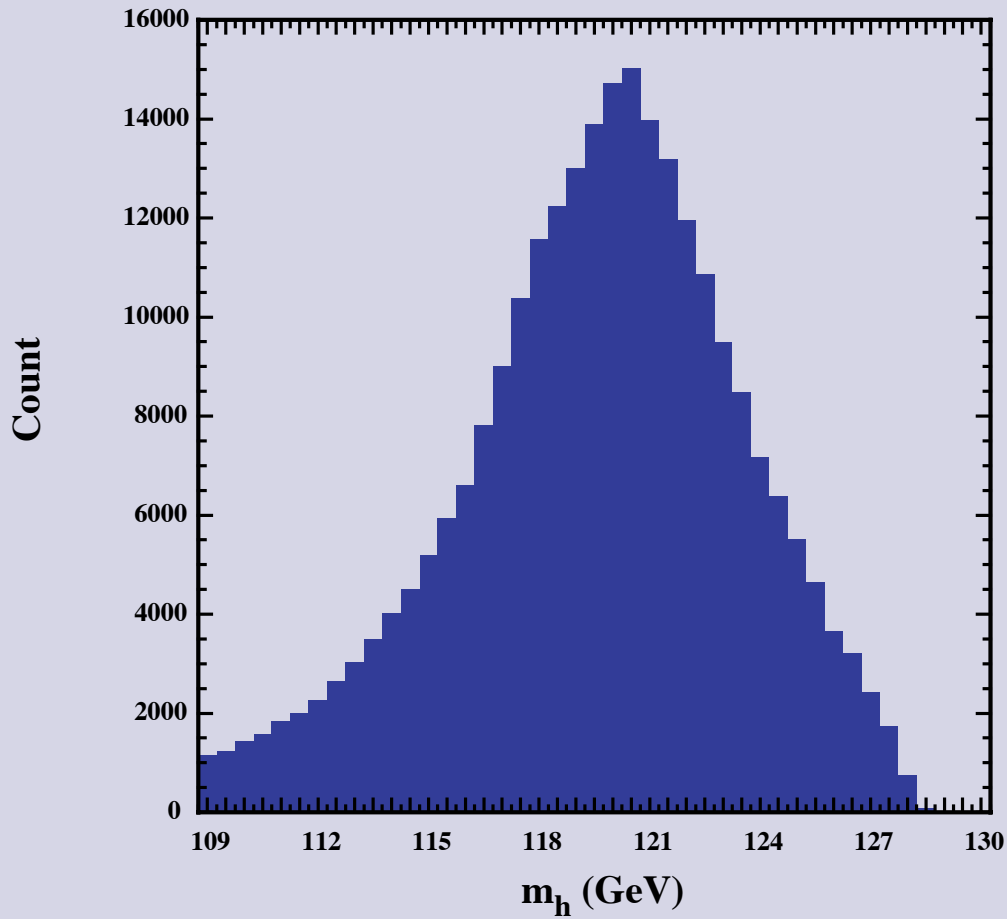
Unification to
rich spectrum
+
EWSB

Falk

$m_{1/2} - m_0$ planes



The Higgs mass in the CMSSM



mSUGRA models

e.g. Barbieri, Ferrara, Savoy

$$G = \varphi \varphi^* + z z^* + \ln |W|^2; \quad W = f(z) + g(\varphi)$$

Scalar Potential (N=1):

$$V = e^{(|z|^2 + |\varphi|^2)} \left[\left| \frac{\partial f}{\partial z} + z^* (f(z) + g(\varphi)) \right|^2 + \left| \frac{\partial g}{\partial \varphi} + \varphi^* (f(z) + g(\varphi)) \right|^2 - 3|f(z) + g(\varphi)|^2 \right]$$

In the low energy limit ($M_P \rightarrow \infty$),

$$V = \left| \frac{\partial g}{\partial \phi^i} \right|^2 + \left(A_0 g^{(3)} + B_0 g^{(2)} + h.c. \right) + m_{3/2}^2 \phi^i \phi_i^*$$

where

$$A_0 g^{(3)} = \left(\phi^i \frac{\partial g^{(3)}}{\partial \phi^i} - 3g^{(3)} \right) m_{3/2} + z^* \left(z f^* + \frac{\partial f^*}{\partial z^*} \right) g^{(3)}$$

For example,

Polonyi: $f(z) = m_0 (z + \beta)$;

With $\langle z \rangle = \sqrt{3} - 1$ for $\beta = 2 - \sqrt{3}$

$$m_0 = m_{3/2} ; \quad A_0 = (3 - \sqrt{3}) m_0 ; \quad B_0 = A_0 - m_0$$

mSUGRA

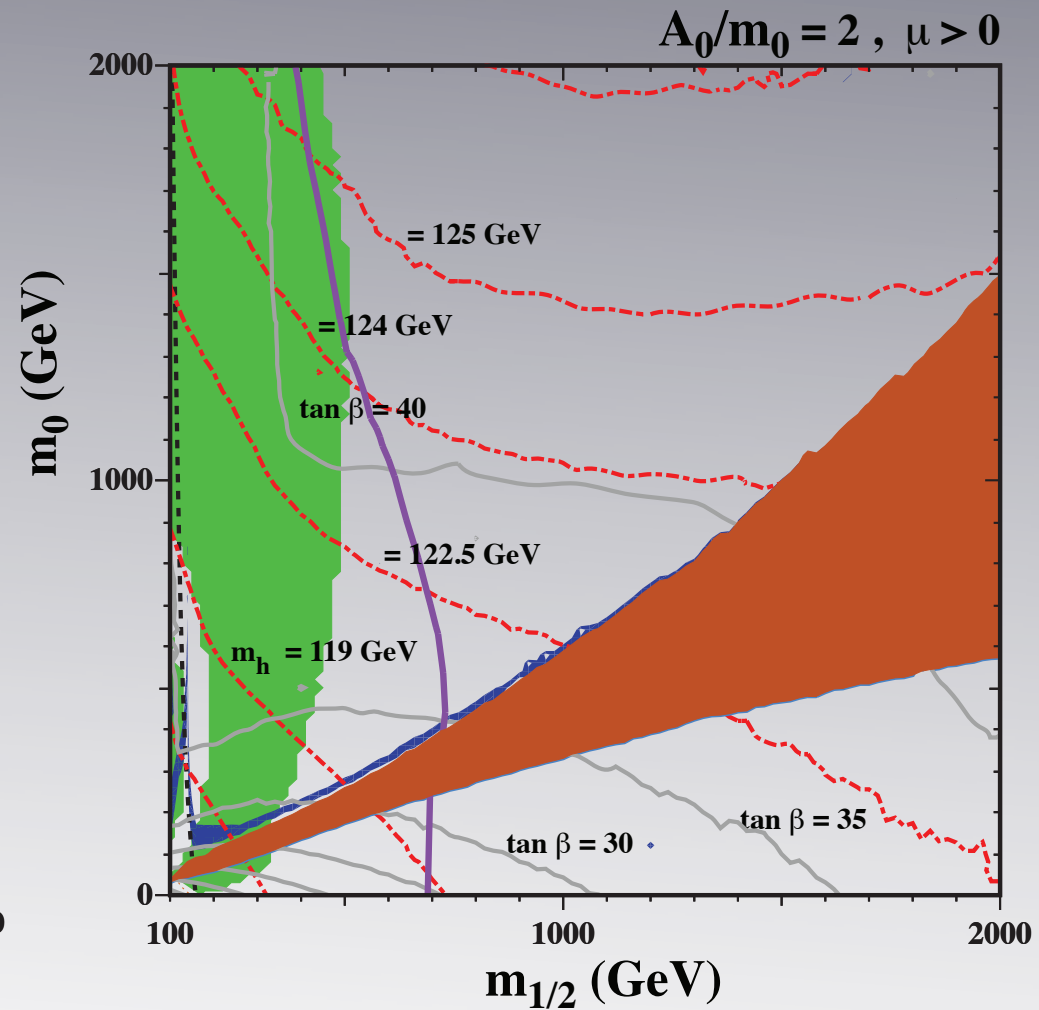
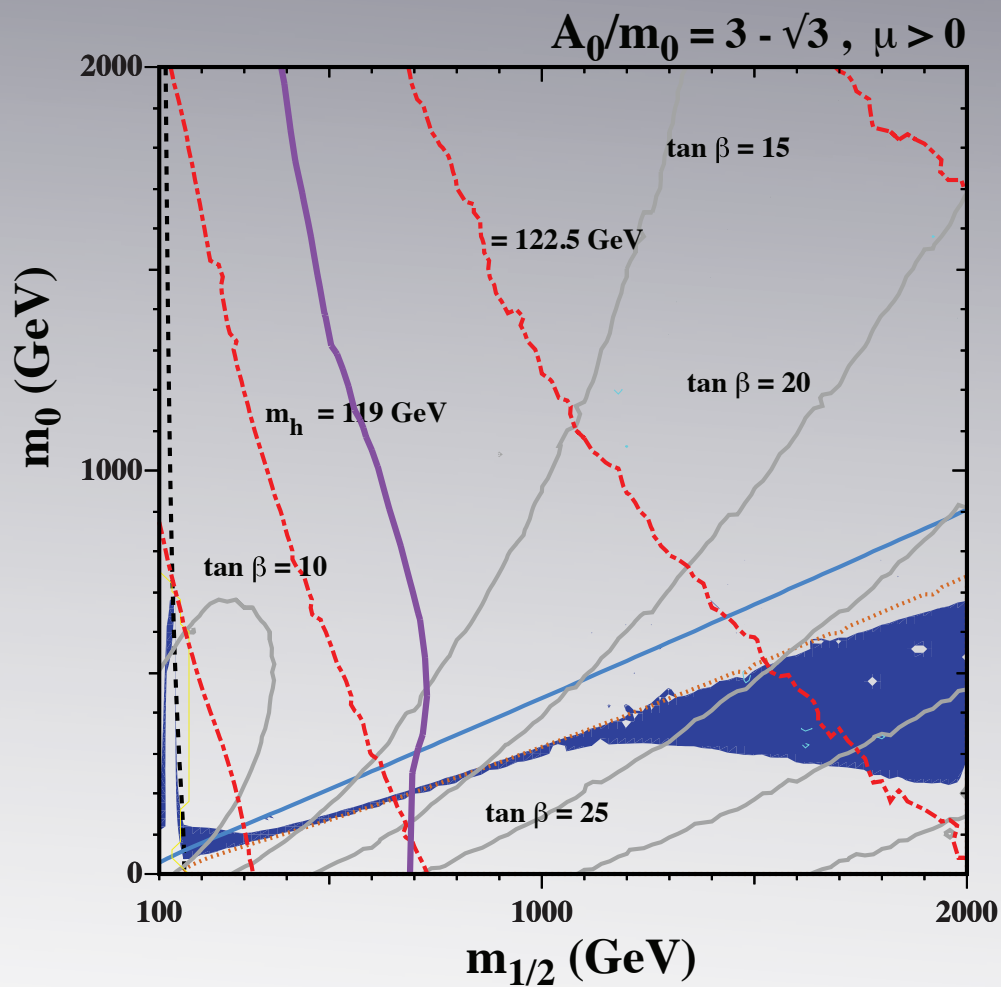
Parameters: $m_{1/2}$, $m_{3/2}$, A_0 , $\text{sgn}(\mu)$

Electroweak Symmetry Breaking conditions used to solve for $\tan\beta$:

$$\mu^2 = \frac{m_1^2 - m_2^2 \tan^2 \beta + \frac{1}{2} M_Z^2 (1 - \tan^2 \beta) + \Delta_\mu^{(1)}}{\tan^2 \beta - 1 + \Delta_\mu^{(2)}}$$

$$B\mu = -\frac{1}{2} (m_1^2 + m_2^2 + 2\mu^2) \sin 2\beta + \Delta_B$$

mSUGRA planes



Mastercode - MCMC

Long list of observables to
constrain CMSSM parameter space

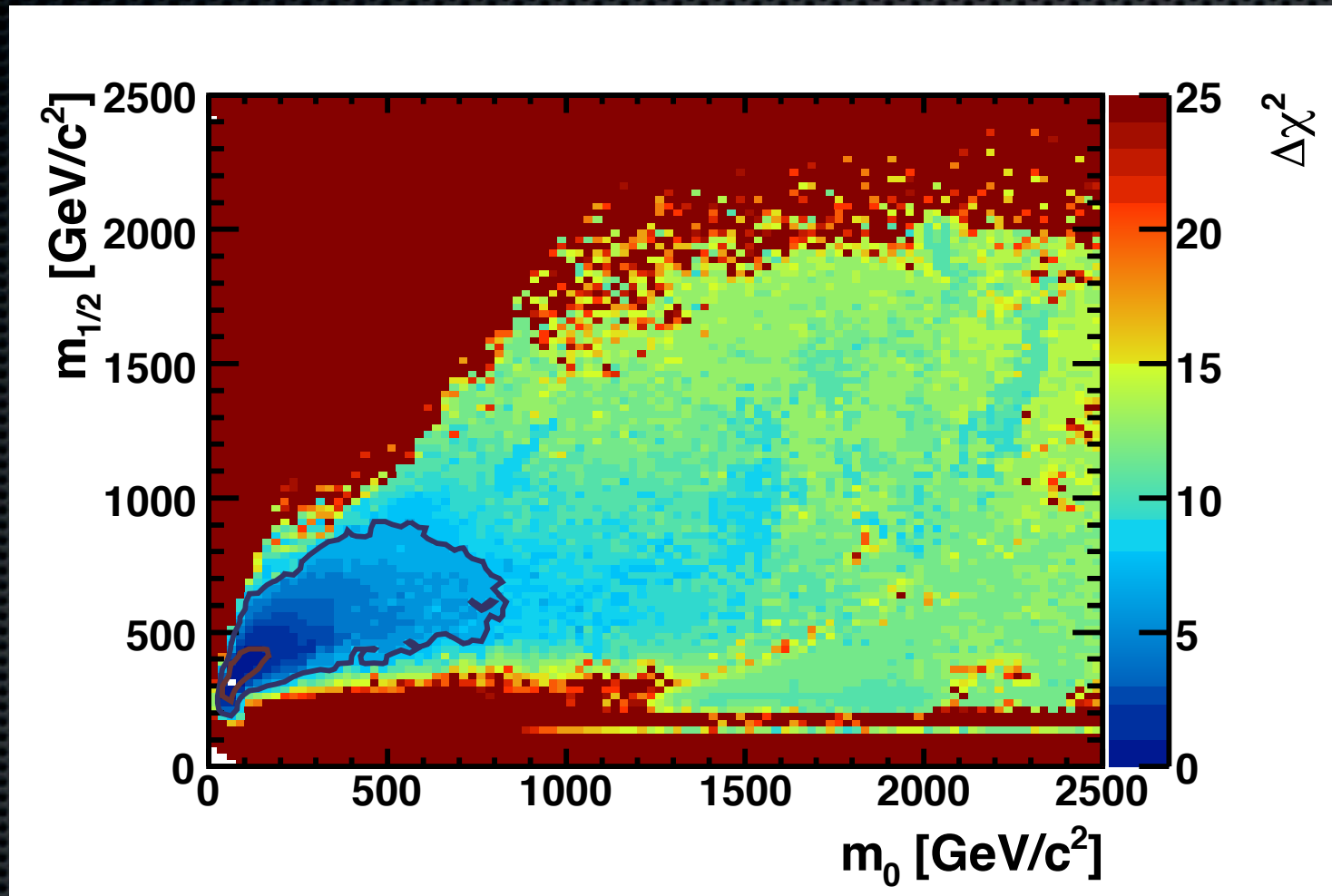
- ✦ MCMC technique to sample efficiently the SUSY parameter space, and thereby construct the χ^2 probability function
- ✦ Combines SoftSusy, FeynHiggs, SuperFla, SuperIso, MicrOmegas, and SSARD
- ✦ Purely frequentist approach (no priors) and relies only on the value of χ^2 at the point sampled and not on the distribution of sampled points.
- ✦ 70 million points sampled (CMSSM)

$$\begin{aligned}\chi^2 = & \sum_i^N \frac{(C_i - P_i)^2}{\sigma(C_i)^2 + \sigma(P_i)^2} \\ & + \chi^2(M_h) + \chi^2(\text{BR}(B_s \rightarrow \mu\mu)) \\ & + \chi^2(\text{SUSY search limits}) \\ & + \sum_i^M \frac{(f_{\text{SM}_i}^{\text{obs}} - f_{\text{SM}_i}^{\text{fit}})^2}{\sigma(f_{\text{SM}_i})^2}\end{aligned}$$

$\Delta\chi^2$ map of $m_0 - m_{1/2}$ plane

Mastercode

2009

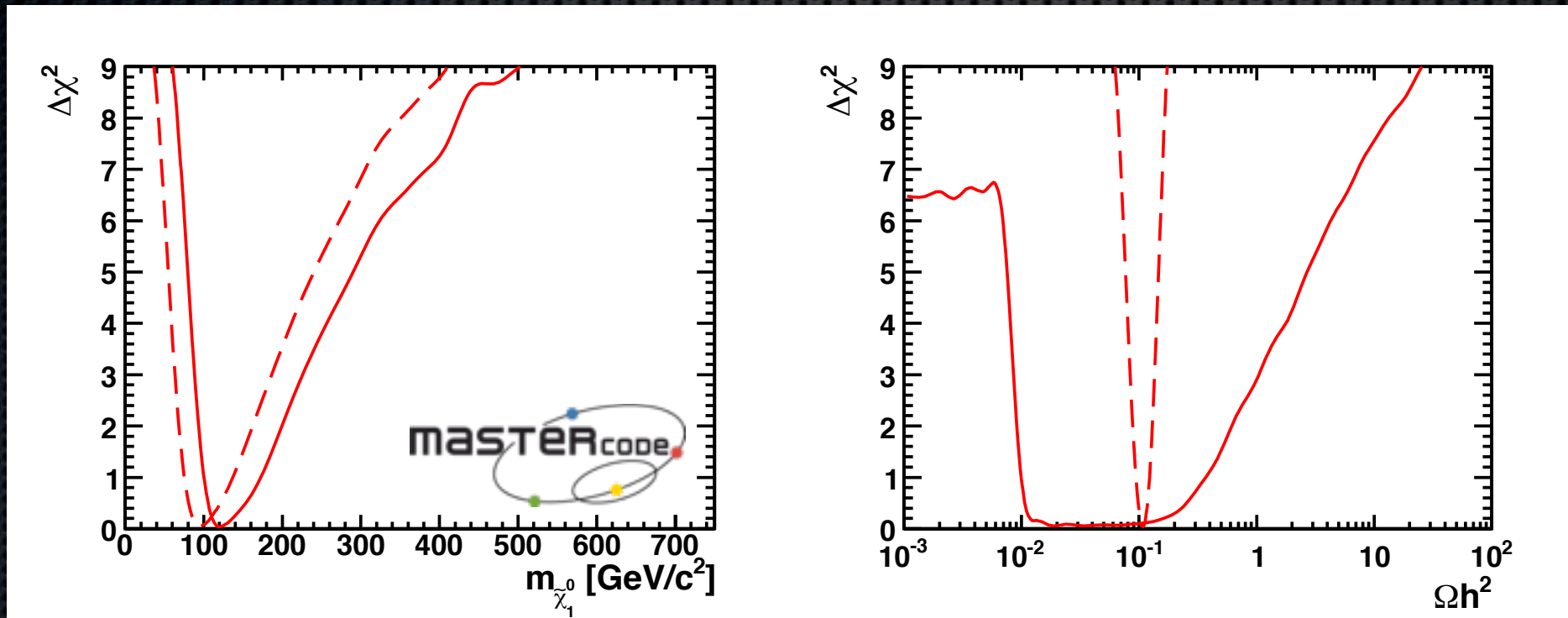


✦ CMSSM

Buchmueller, Cavanaugh, De Roeck, Ellis, Flacher, Heinemeyer,
Isidori, Olive, Ronga, Weiglein

Neutralino mass and Relic Density from MCMC analysis

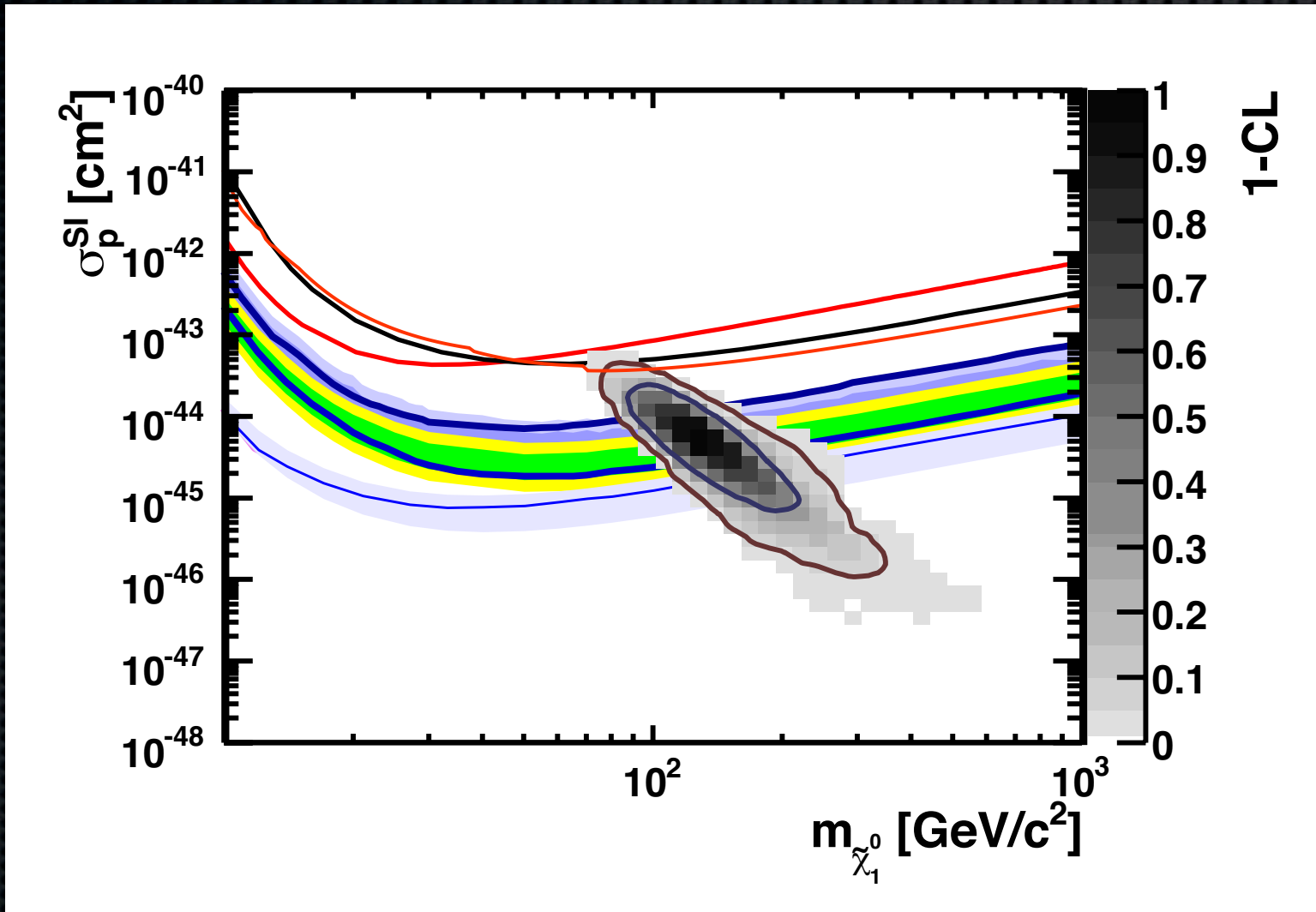
Mastercode 2009



Elastic scattering cross-section

Mastercode

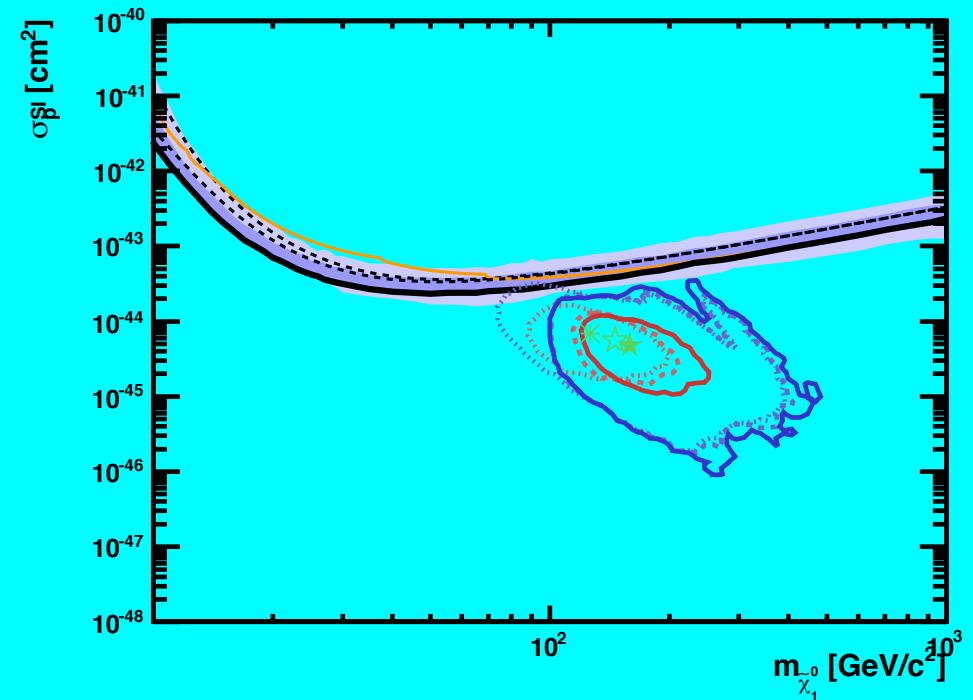
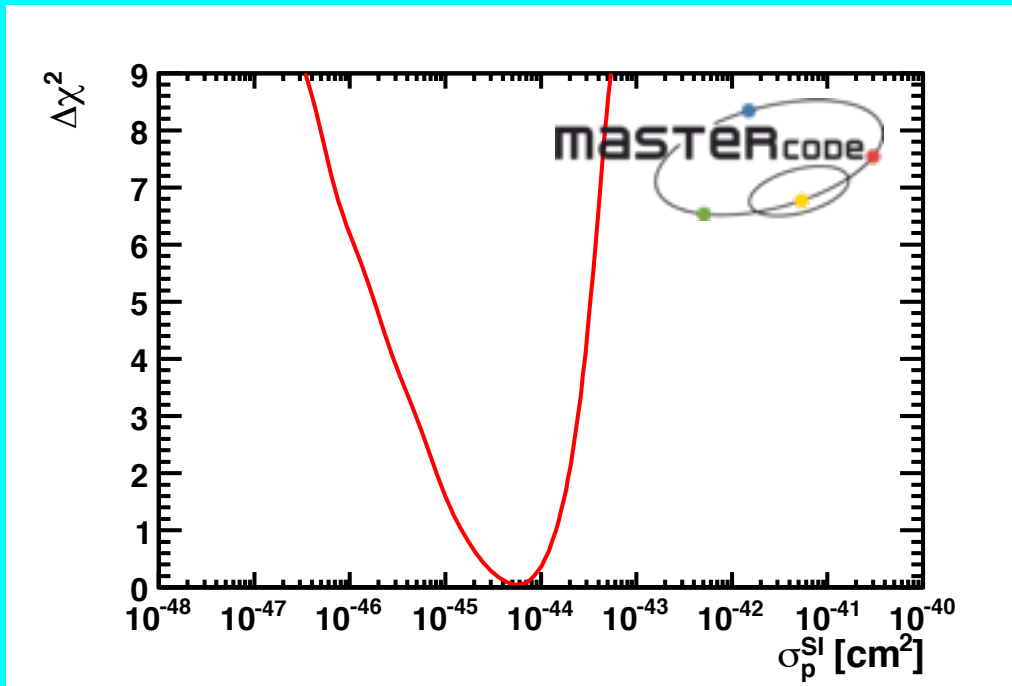
2009



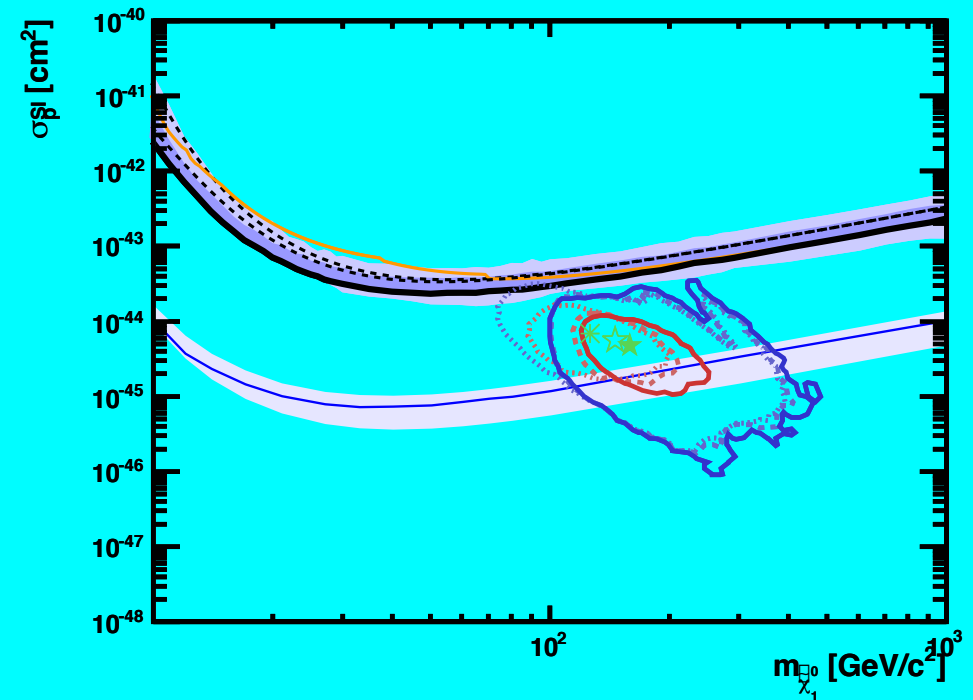
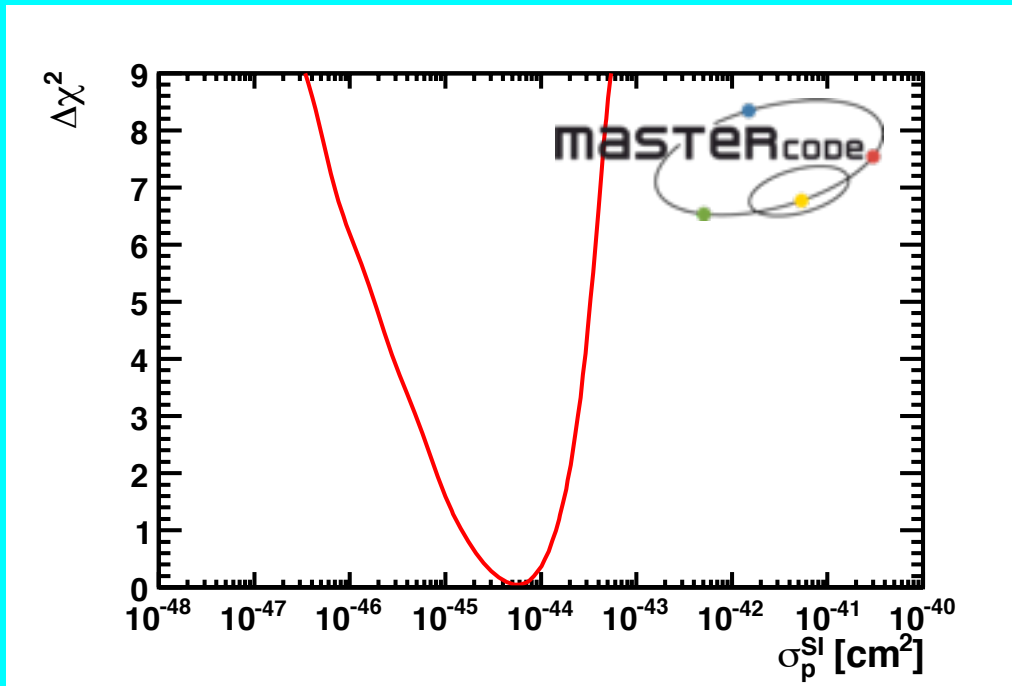
CMSSM

Buchmueller, Cavanaugh, De Roeck, Ellis, Flacher, Heinemeyer,
Isidori, Olive, Ronga, Weiglein

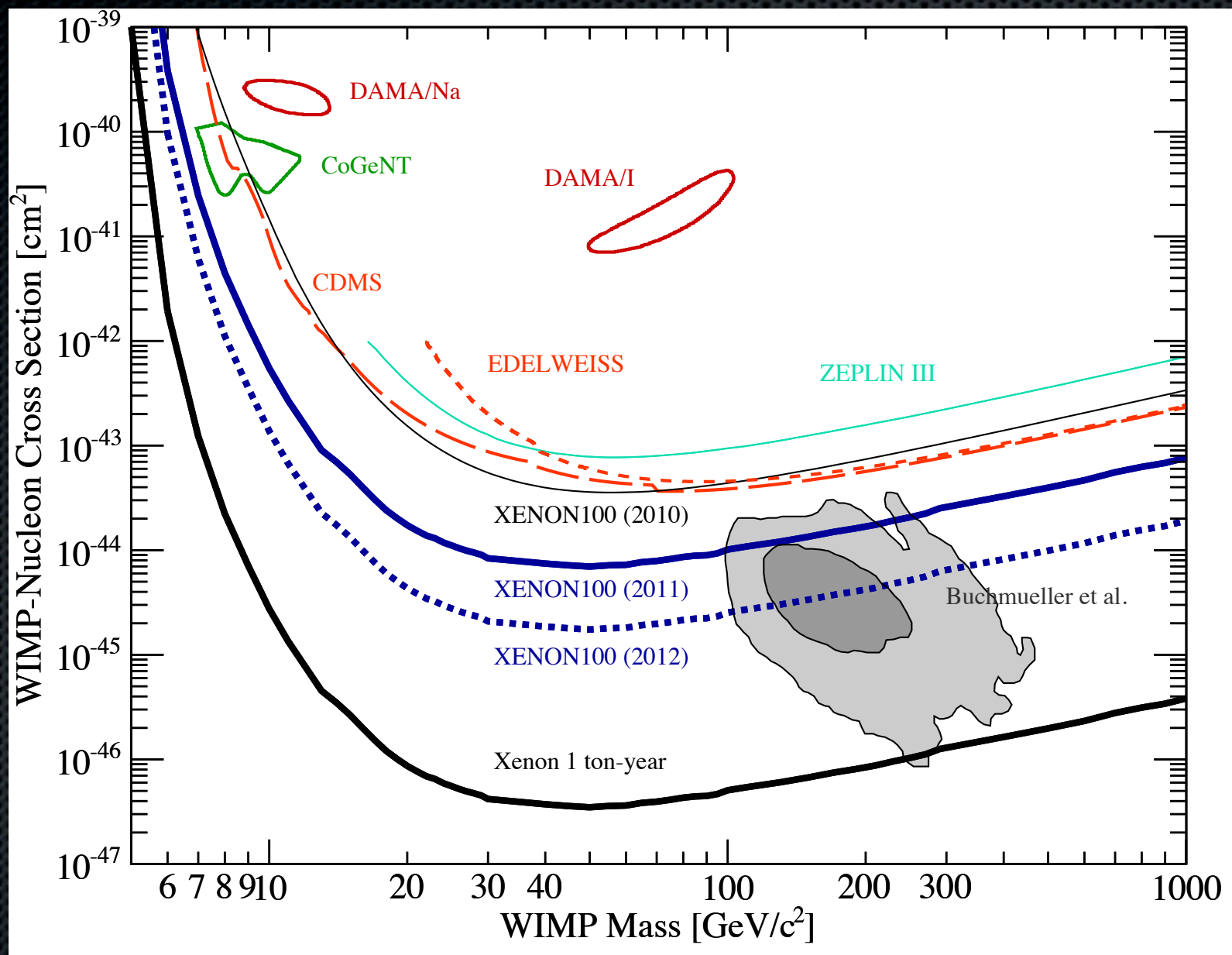
Elastic cross section from MCMC analysis



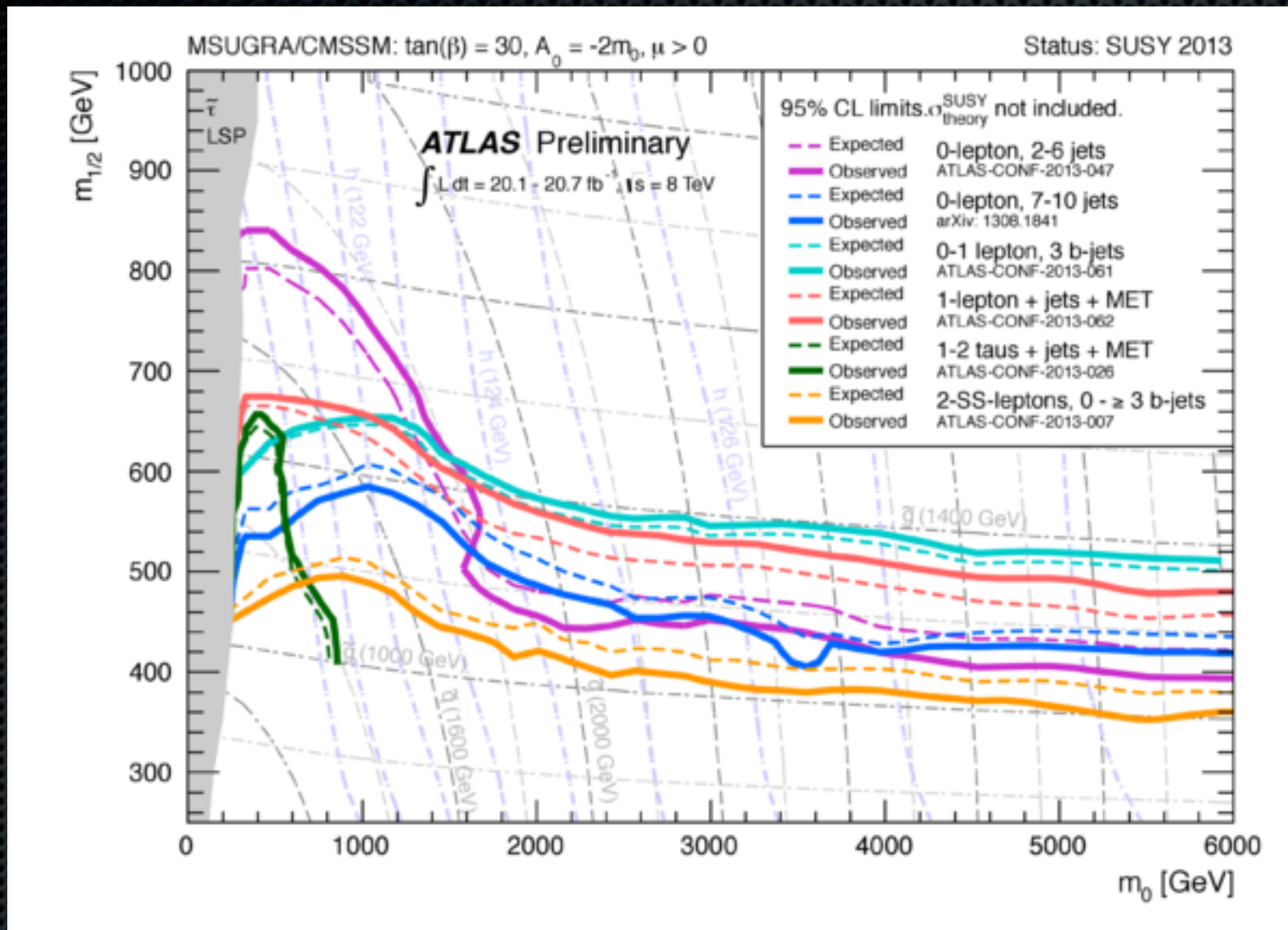
Elastic cross section from MCMC analysis



Most recent result from XENON100

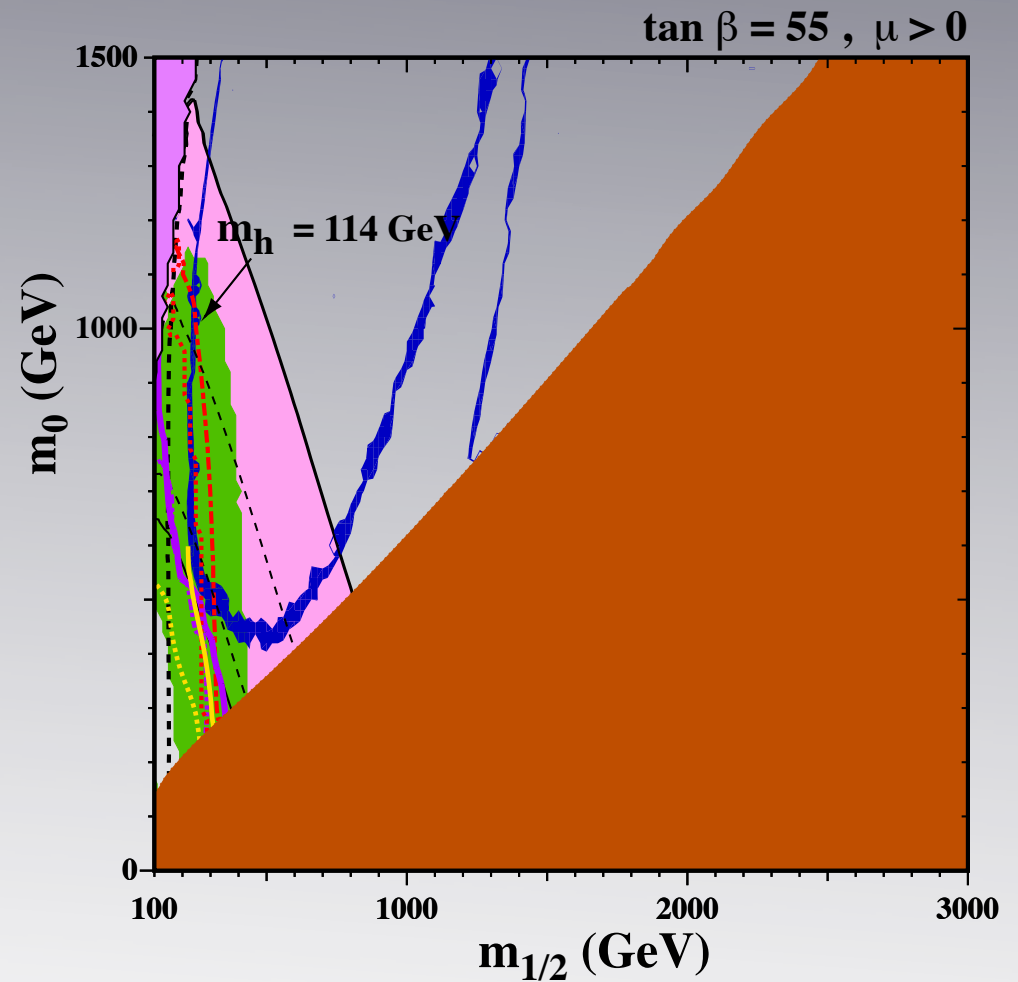
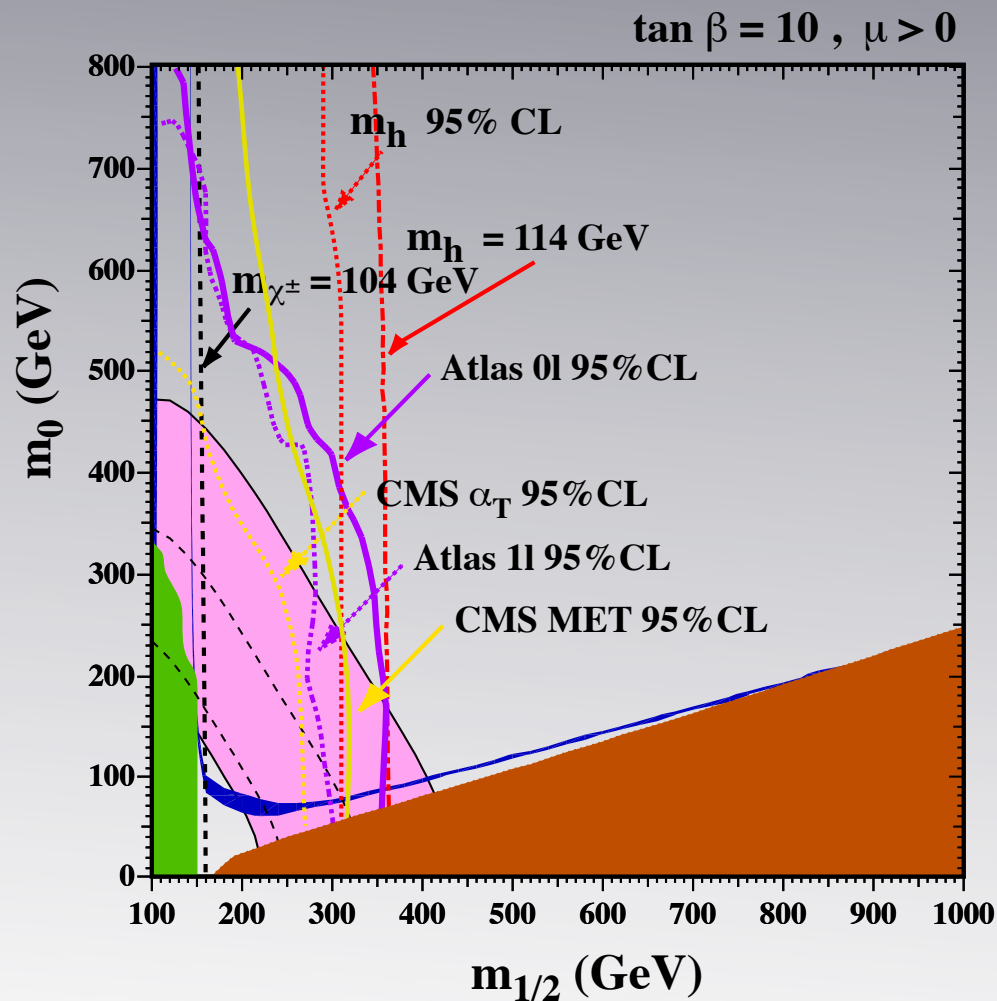


ATLAS Results from run I

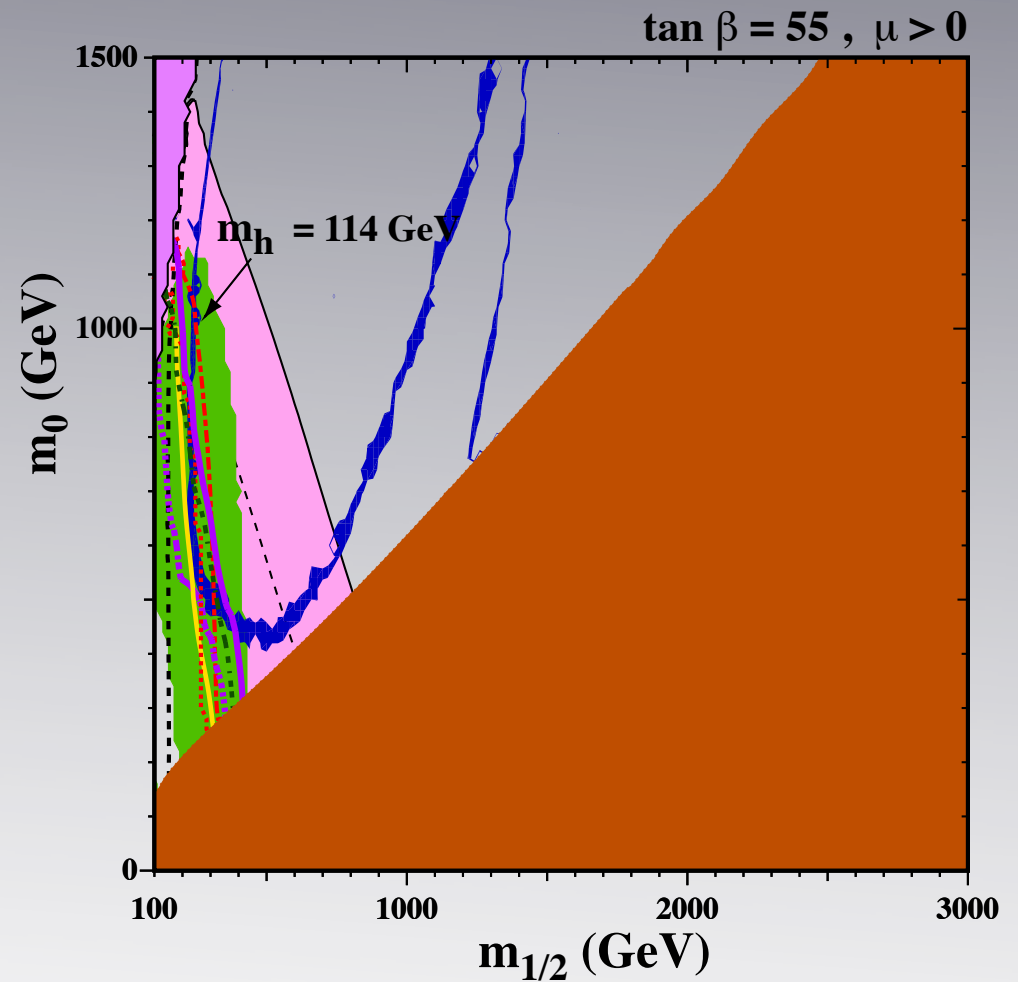
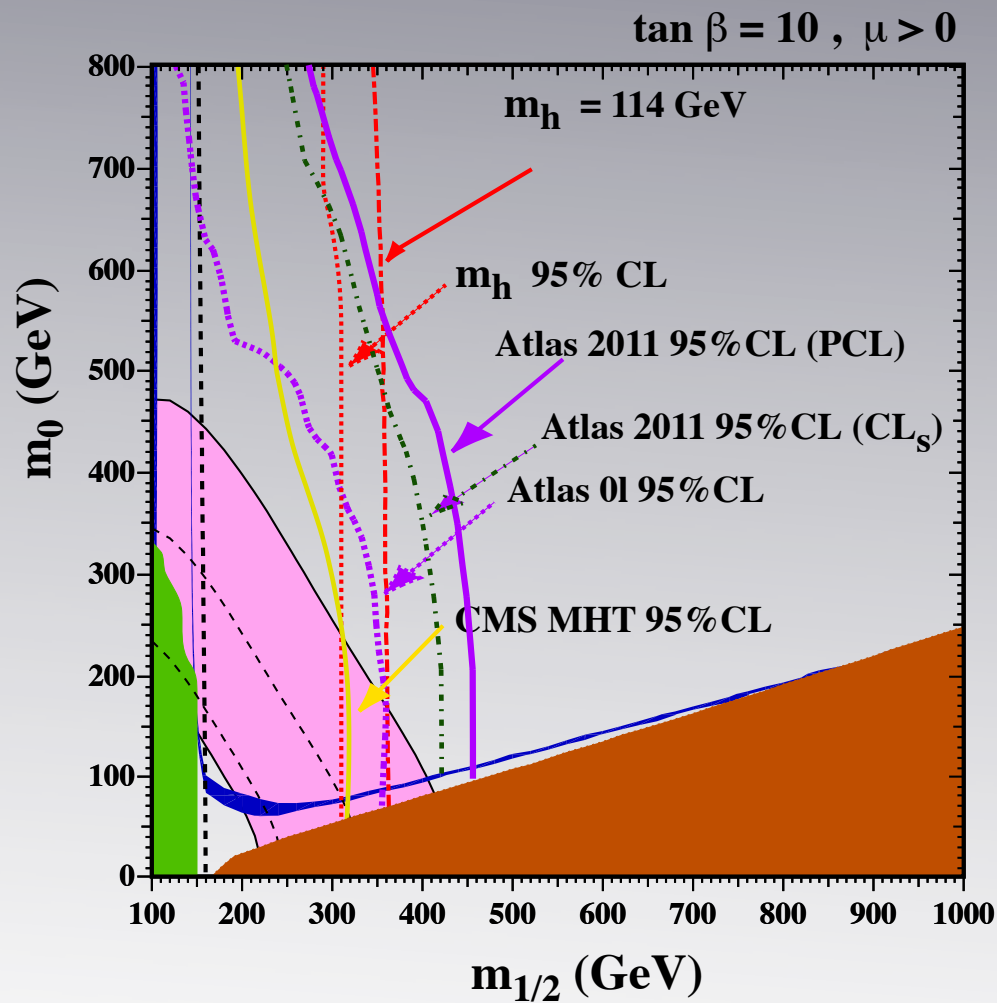


$\sim 20.7 \text{ fb}^{-1}$ @ 8 TeV

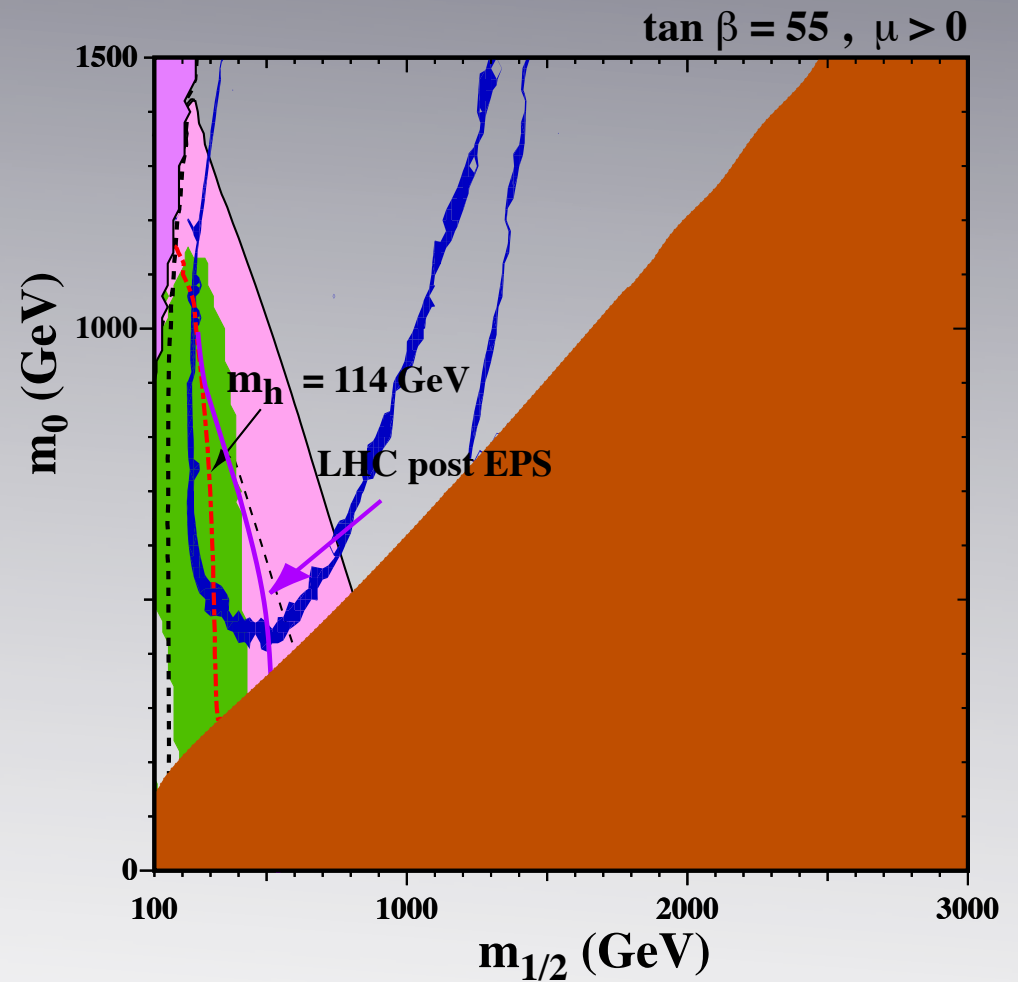
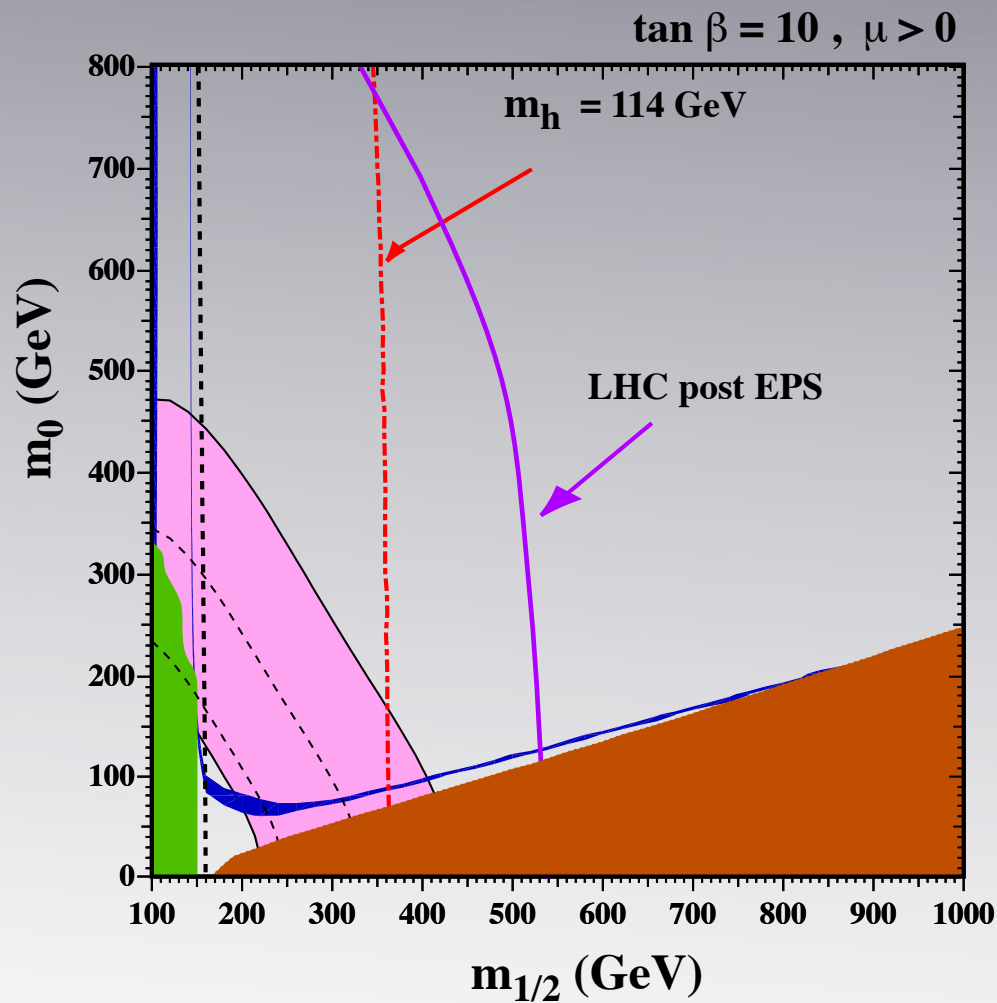
$m_{1/2} - m_0$ planes incl. LHC



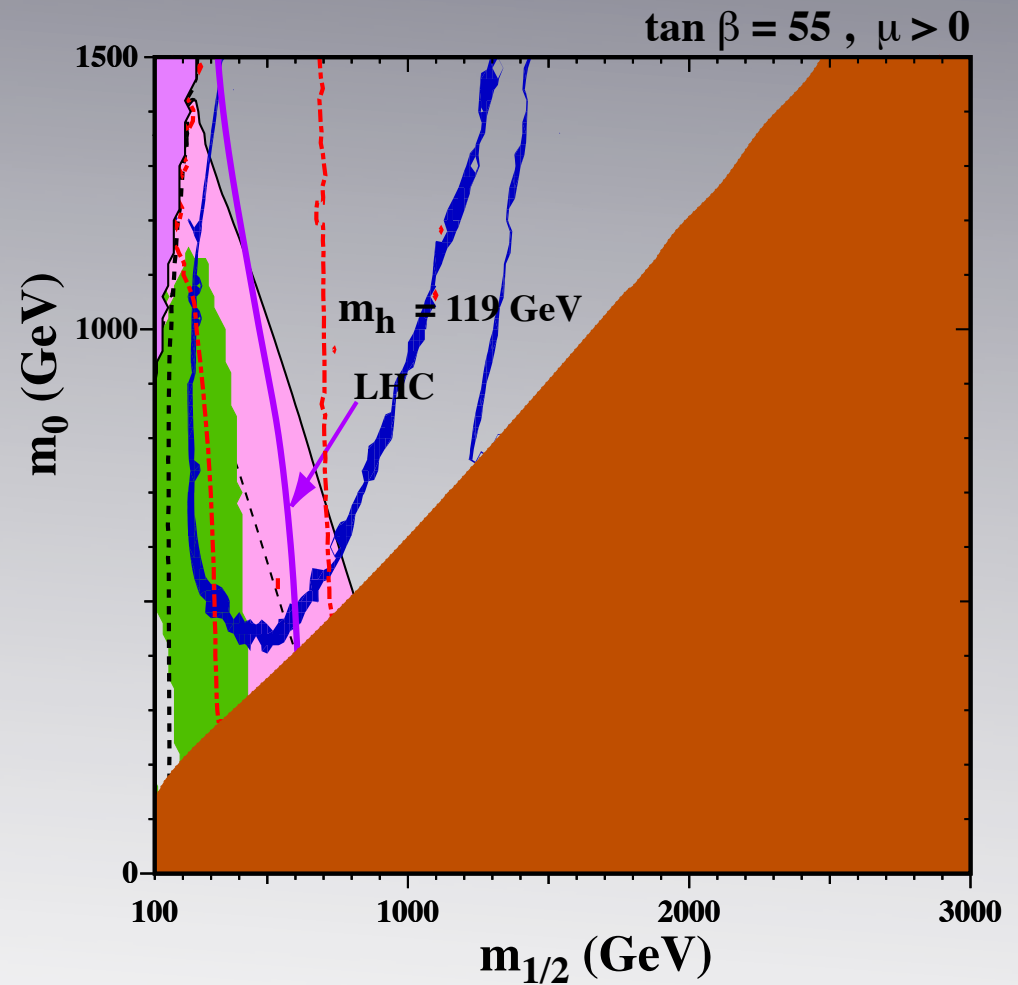
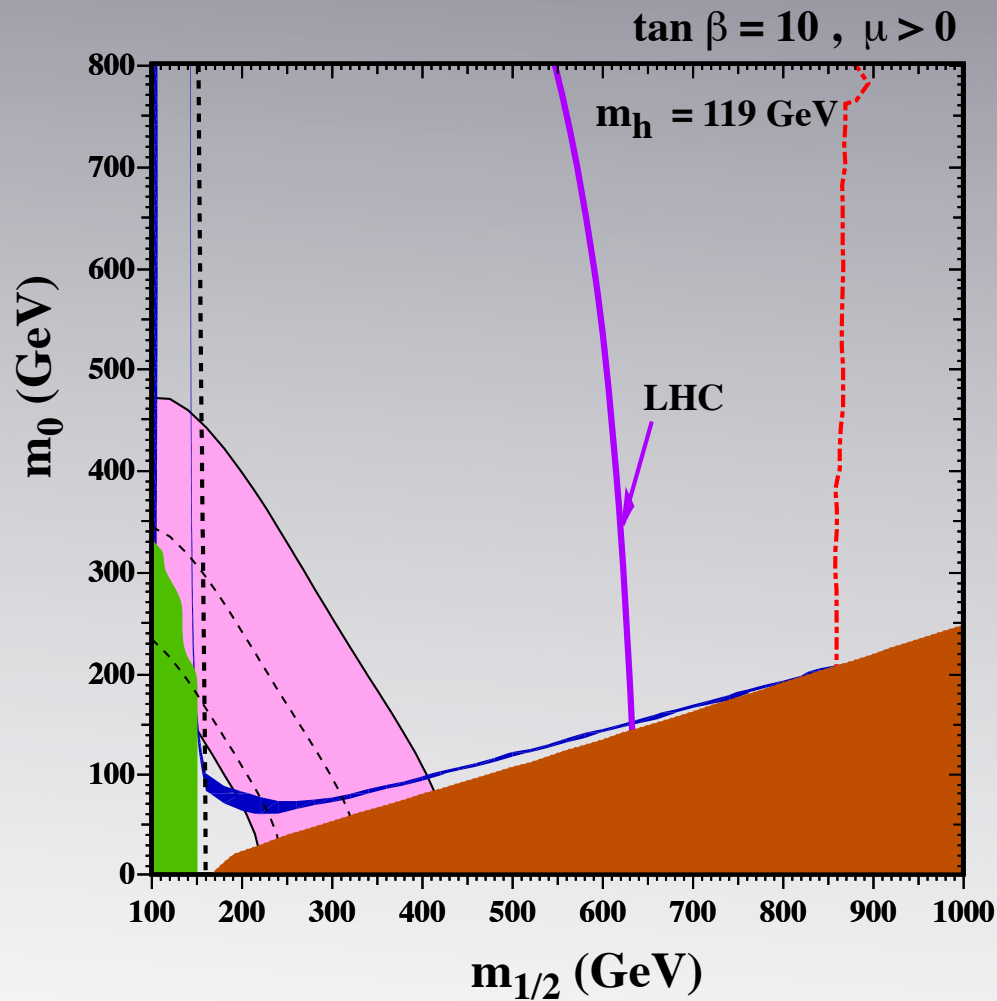
$m_{1/2} - m_0$ planes incl. LHC



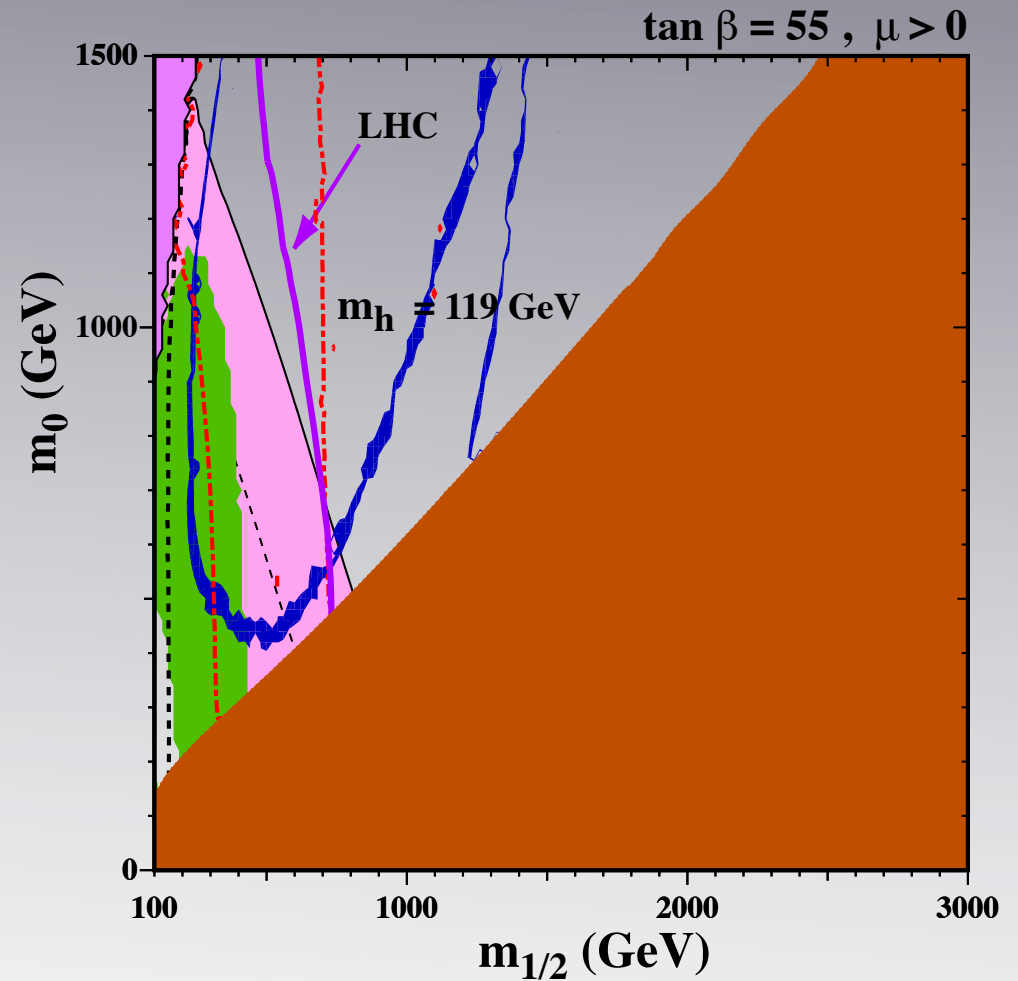
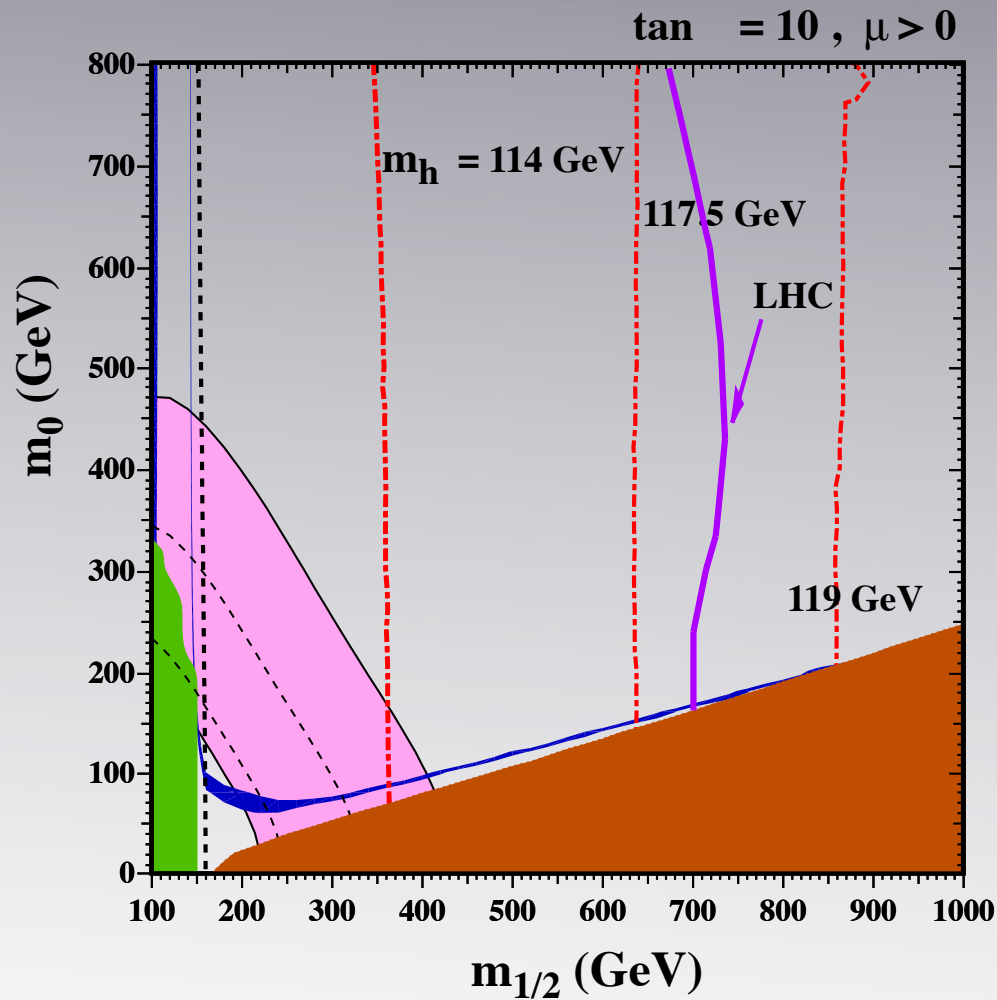
$m_{1/2} - m_0$ planes incl. LHC



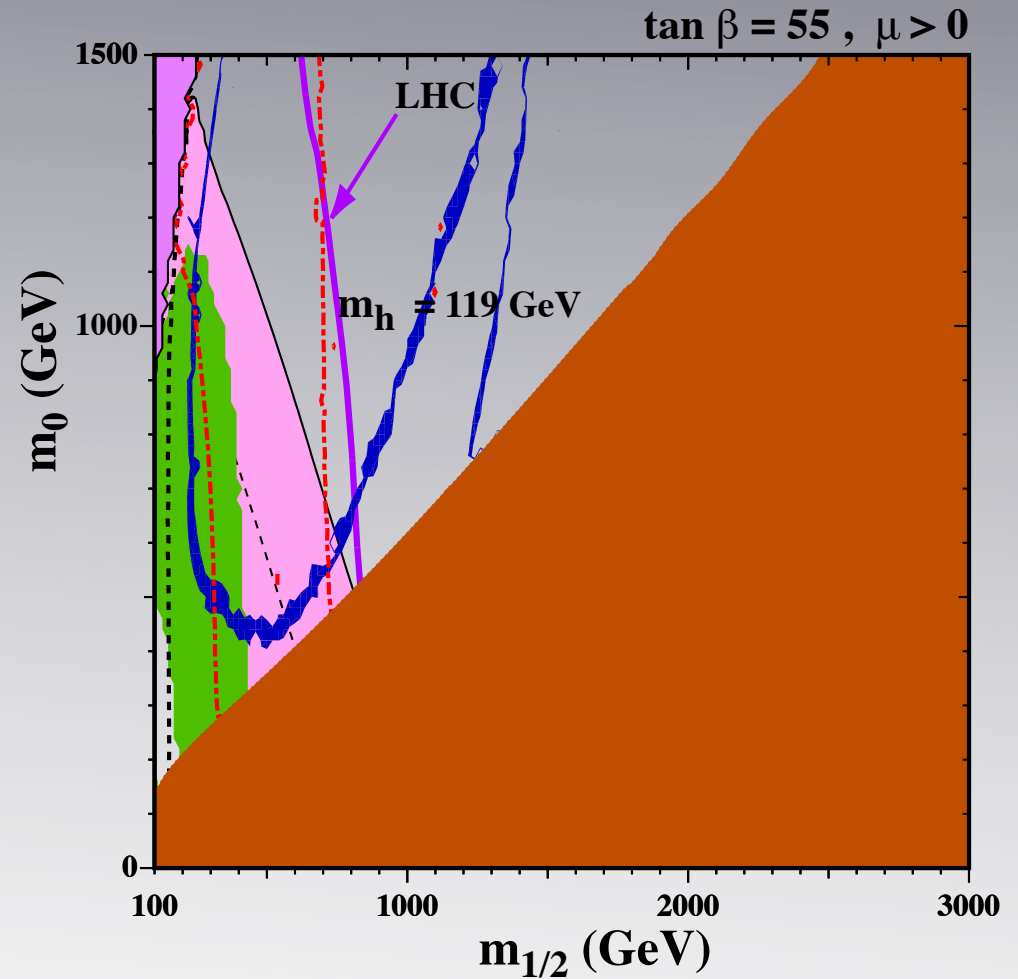
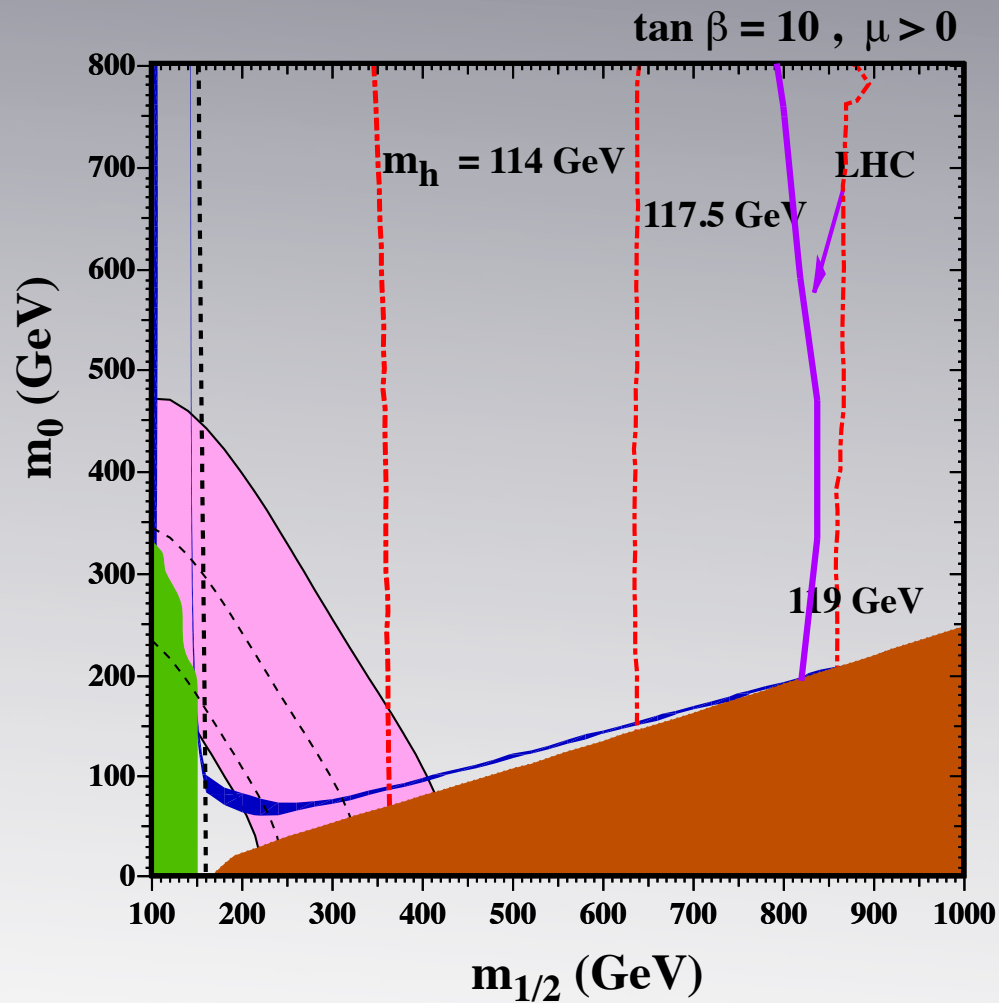
$m_{1/2} - m_0$ planes incl. LHC



$m_{1/2} - m_0$ planes incl. LHC

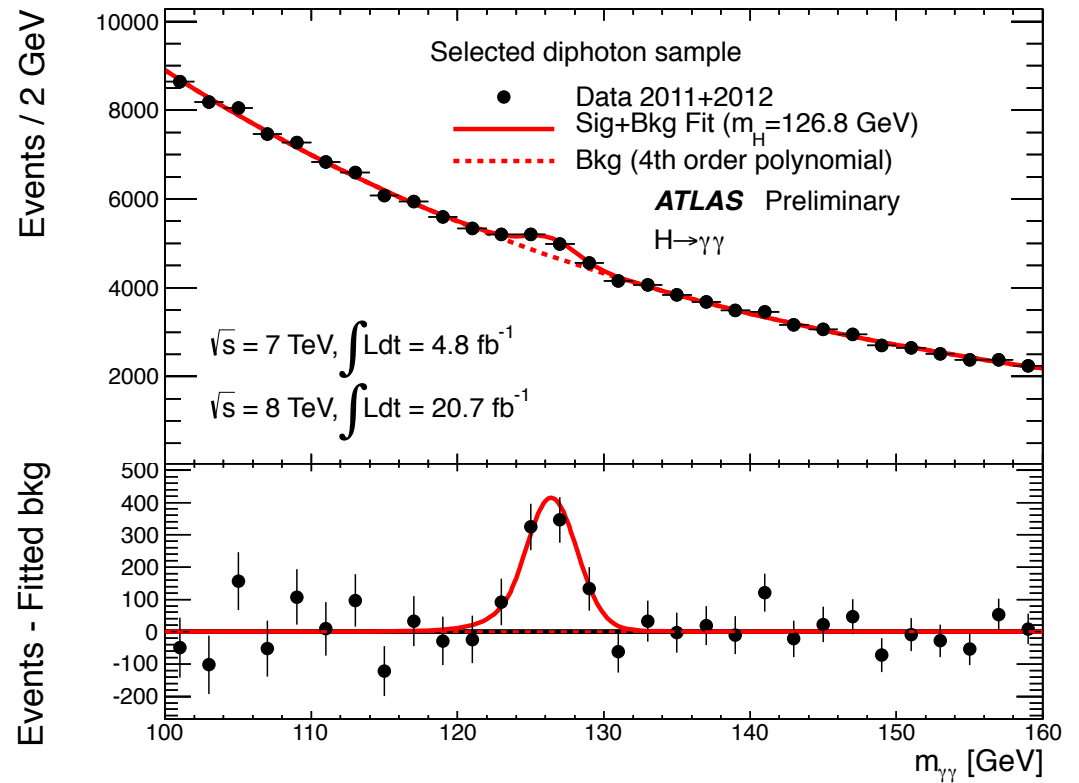
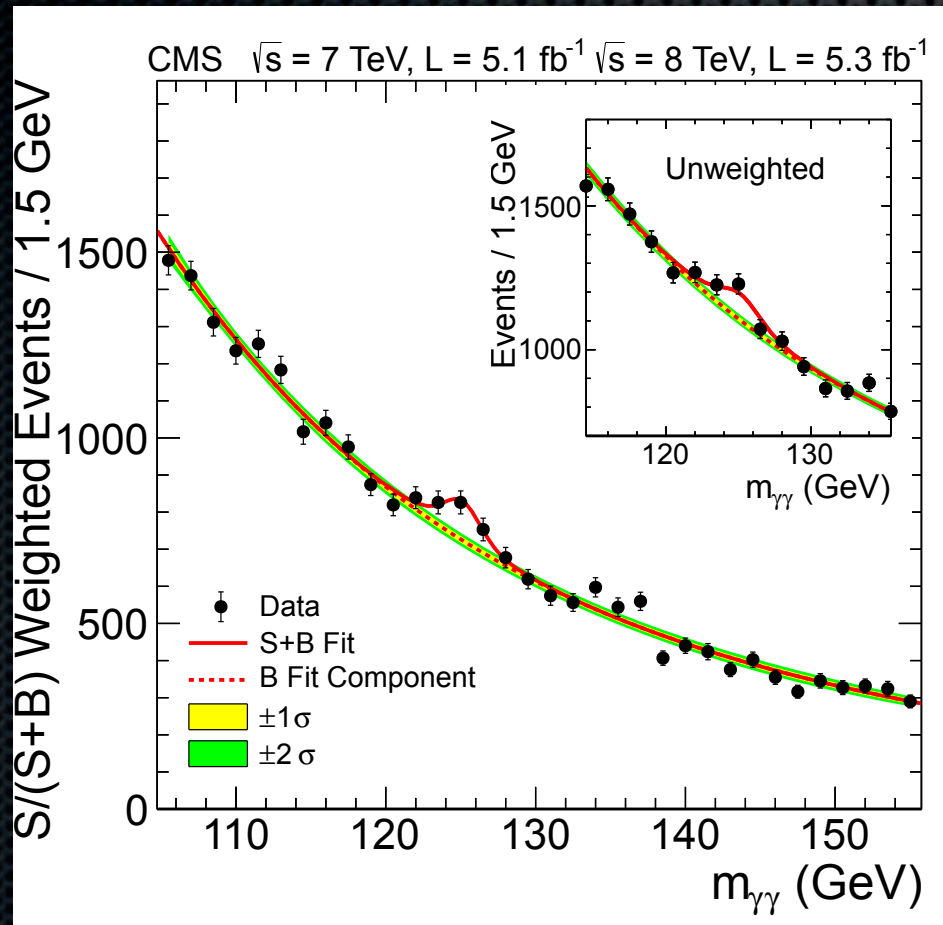


$m_{1/2} - m_0$ planes incl. LHC



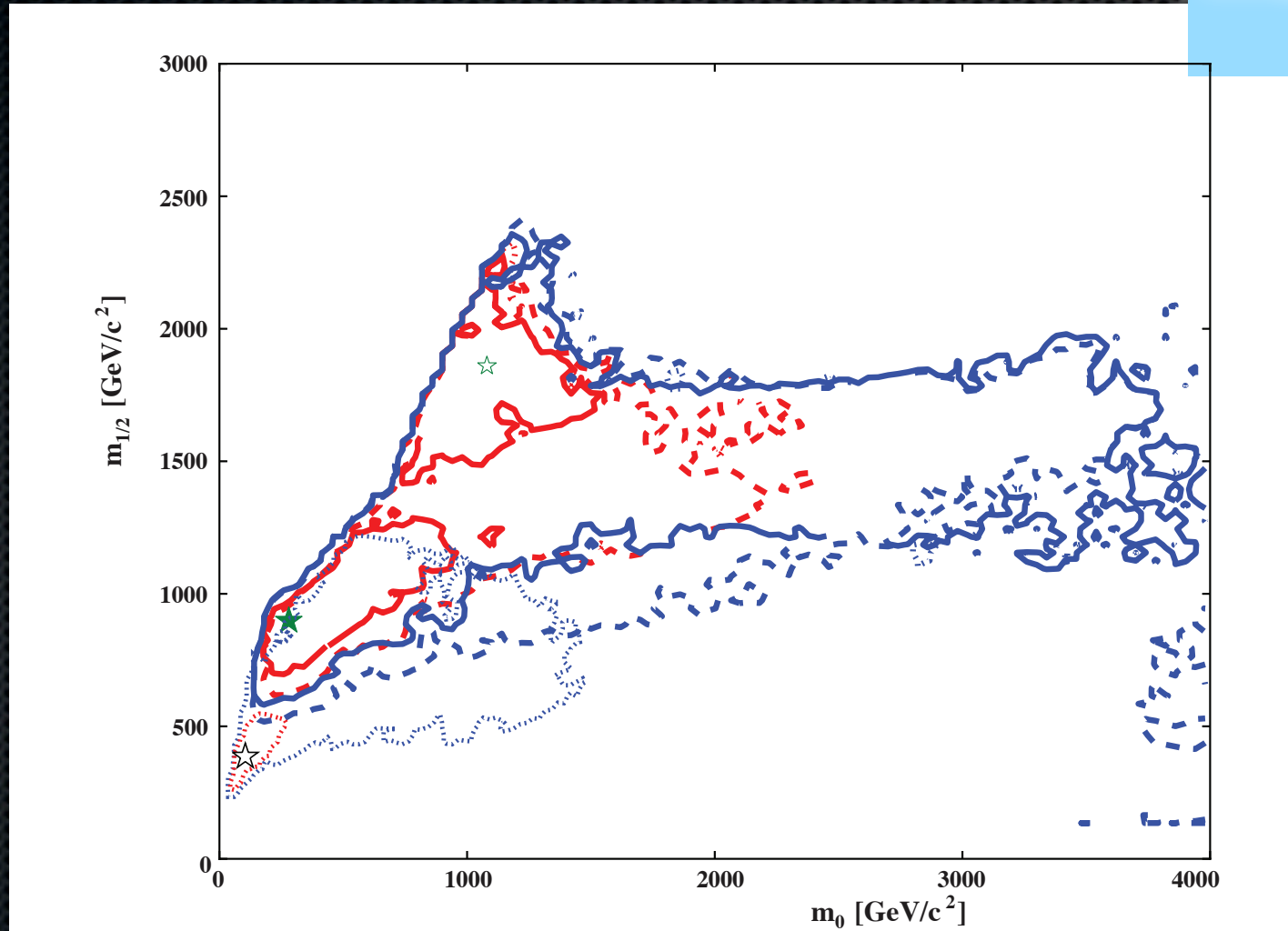
The Higgs Search

The LHC @ $\sim 20.7/\text{fb}$



$\Delta\chi^2$ map of $m_0 - m_{1/2}$ plane

Limits at $\sim 5 \text{ fb}^{-1}$



Buchmueller, Cavanaugh, Citron, De Roeck, Dolan, Ellis, Flacher, Heinemeyer, Isidori, Marrouche, Martinez Santos, Nakach, Olive, Rogerson, Ronga, de Vries, Weiglein

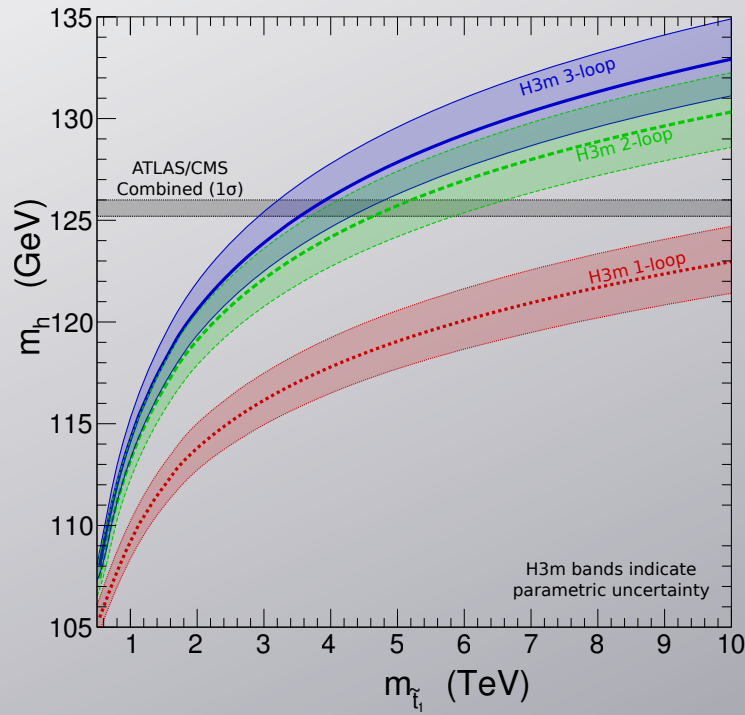
COMPARISON OF BEST FIT POINTS PRE AND POST LHC

Model	Data set	Minimum χ^2 /d.o.f.	Probability	m_0 (GeV)	$m_{1/2}$ (GeV)	A_0 (GeV)	$\tan \beta$
CMSSM	pre-LHC	21.5/20	37 %	90	360	-400	15
	LHC _{1/fb}	31.0/23	12%	1120	1870	1220	46
	ATLAS _{5/fb} (low)	32.8/23	8.5%	300	910	1320	16
	ATLAS _{5/fb} (high)	33.0/23	8.0%	1070	1890	1020	45
NUHM1	pre-LHC	20.8/18	29 %	110	340	520	13
	LHC _{1/fb}	28.9/22	15%	270	920	1730	27
	ATLAS _{5/fb} (low)	31.3/22	9.1%	240	970	1860	16
	ATLAS _{5/fb} (high)	31.8/22	8.1%	1010	2810	2080	39

p-value of SM = 9% (32.7/23) - but note: does not include dark matter

Buchmueller, Cavanaugh, Citron, De Roeck, Dolan, Ellis, Flacher, Heinemeyer, Isidori, Marrouche, Martinez Santos, Nakach, Olive, Rogerson, Ronga, de Vries, Weiglein

New Higgs Mass Calculations



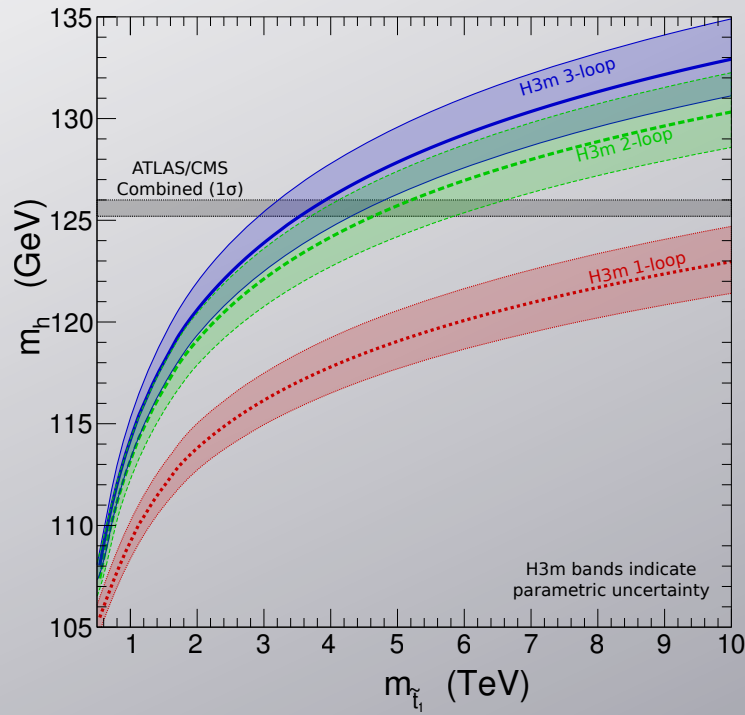
Feng, Kant, Profumo, Sanford

Includes dominant $O(\alpha_t \alpha_s^2)$ corrections

FeynHiggs 2.10.0

to include next-to-leading logs $\text{Log}(m_{\tilde{t}}/m_t)$ to all orders

New Higgs Mass Calculations



Feng, Kant, Profumo, Sanford

Includes dominant $O(\alpha_t \alpha_s^2)$ corrections

FeynHiggs 2.10.0

to include next-to-leading logs $\text{Log}(m_{\tilde{t}}/m_t)$ to all orders

$B \rightarrow \mu\mu$

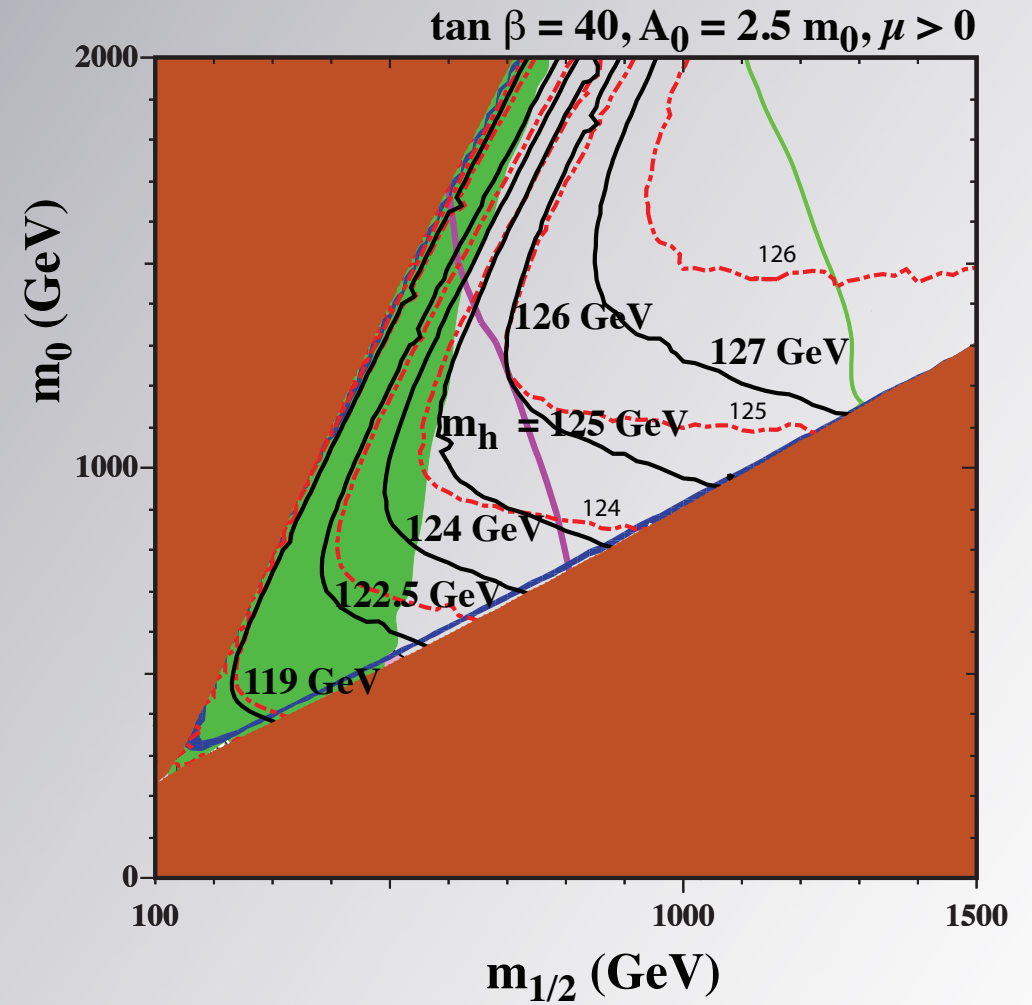
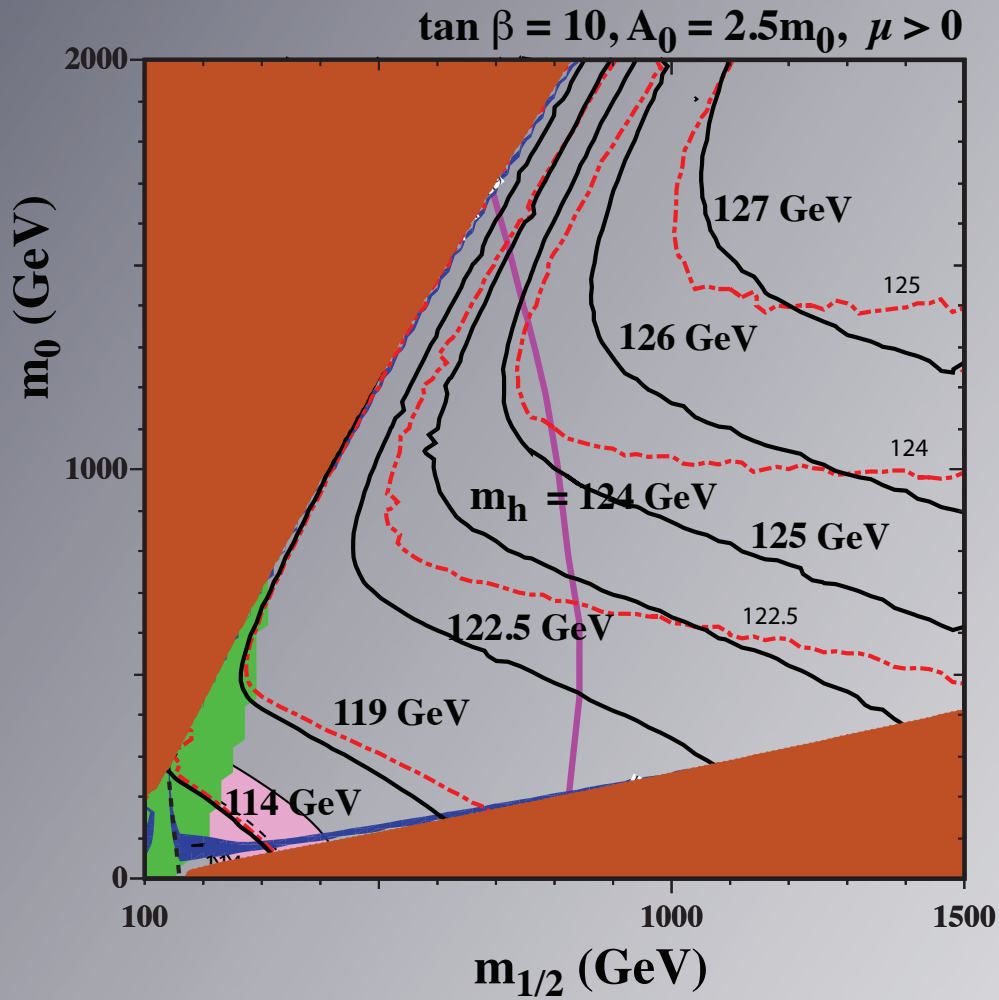
CMS: $\text{BR}(B_s \rightarrow \mu^+ \mu^-) = (3.0_{-0.9}^{+1.0}) \times 10^{-9}$,

LHCb: $\text{BR}(B_s \rightarrow \mu^+ \mu^-) = (2.9_{-1.0}^{+1.1}) \times 10^{-9}$,

Combined: $\left(\frac{\text{BR}(B_{s,d} \rightarrow \mu^+ \mu^-)_{EXP}}{\text{BR}(B_{s,d} \rightarrow \mu^+ \mu^-)_{SM}} \right)_{TA} = 0.94_{-0.21}^{+0.22}$.

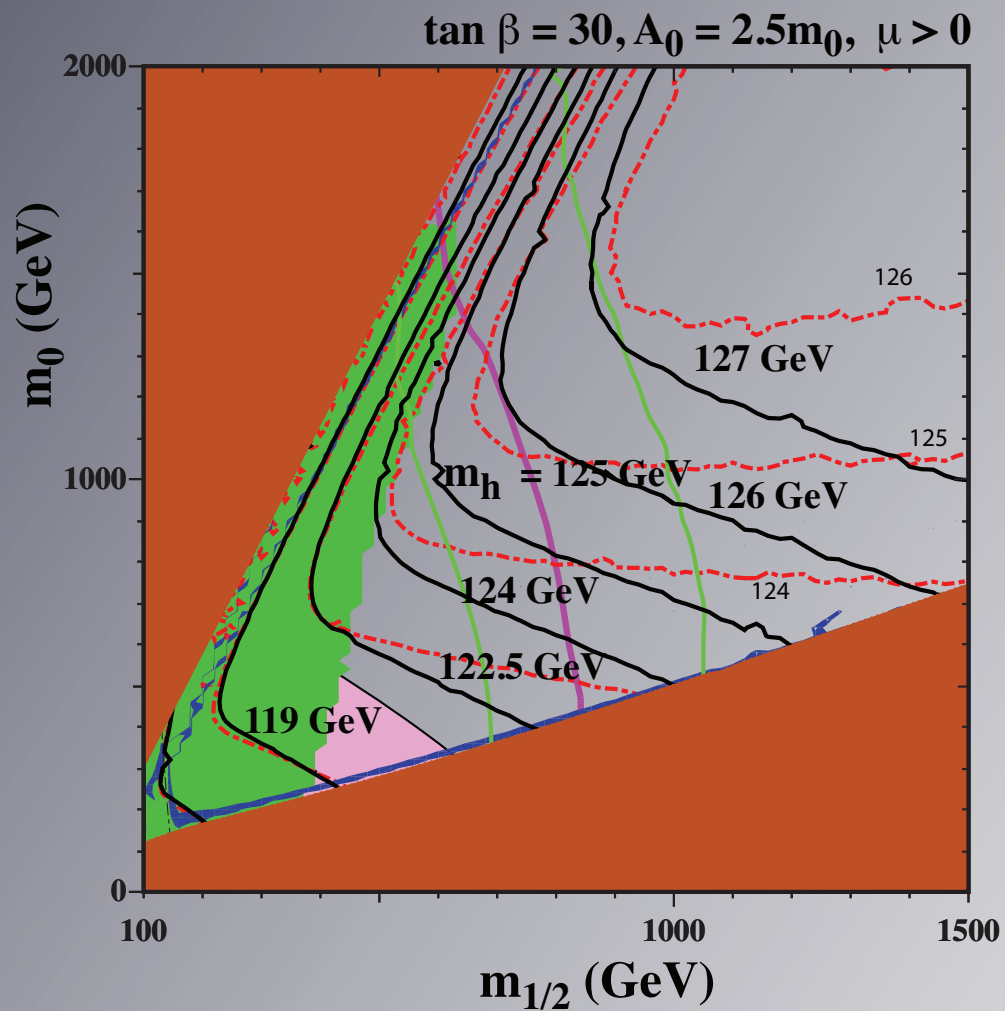
Buchmuller et al.

High and low $\tan \beta$ gone!

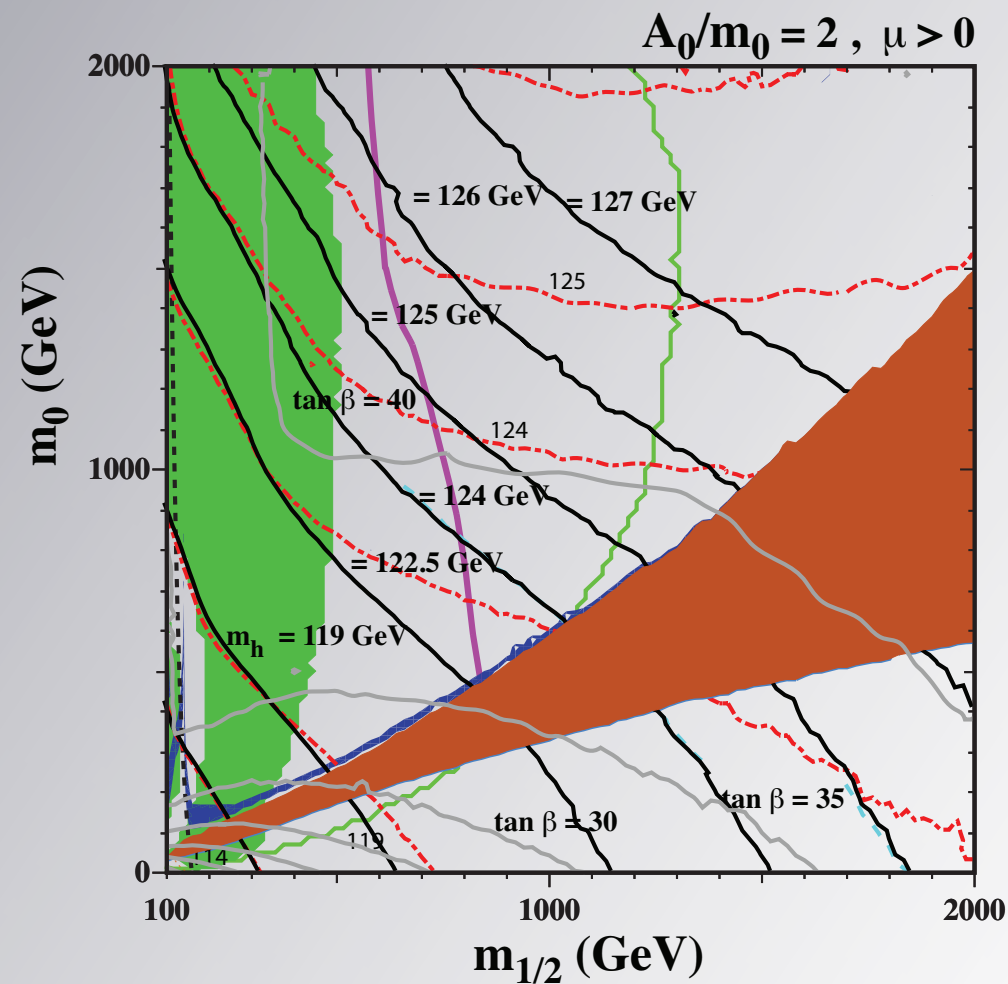


Buchmueller, Dolan, Ellis, Hahn, Heinemeyer, Hollik, Marrouche, Olive, Rzehak, de Vries, Weiglein

Something left?



CMSSM



mSUGRA

Buchmueller, Dolan, Ellis, Hahn, Heinemeyer, Hollik, Marrouche, Olive, Rzehak, de Vries, Weiglein

Mastercode - MCMC

Long list of observables to
constrain CMSSM parameter space

- ✦ MCMC technique to sample efficiently the SUSY parameter space, and thereby construct the χ^2 probability function
- ✦ Combines SoftSusy, FeynHiggs, SuperFla, SuperIso, MicrOmegas, and SSARD
- ✦ Purely frequentist approach (no priors) and relies only on the value of χ^2 at the point sampled and not on the distribution of sampled points.
- ✦ 70 million points sampled (CMSSM)

$$\begin{aligned}\chi^2 = & \sum_i^N \frac{(C_i - P_i)^2}{\sigma(C_i)^2 + \sigma(P_i)^2} \\ & + \chi^2(M_h) + \chi^2(\text{BR}(B_s \rightarrow \mu\mu)) \\ & + \chi^2(\text{SUSY search limits}) \\ & + \sum_i^M \frac{(f_{\text{SM}_i}^{\text{obs}} - f_{\text{SM}_i}^{\text{fit}})^2}{\sigma(f_{\text{SM}_i})^2}\end{aligned}$$

Mastercode - MultiNest

Long list of observables to
constrain CMSSM parameter space

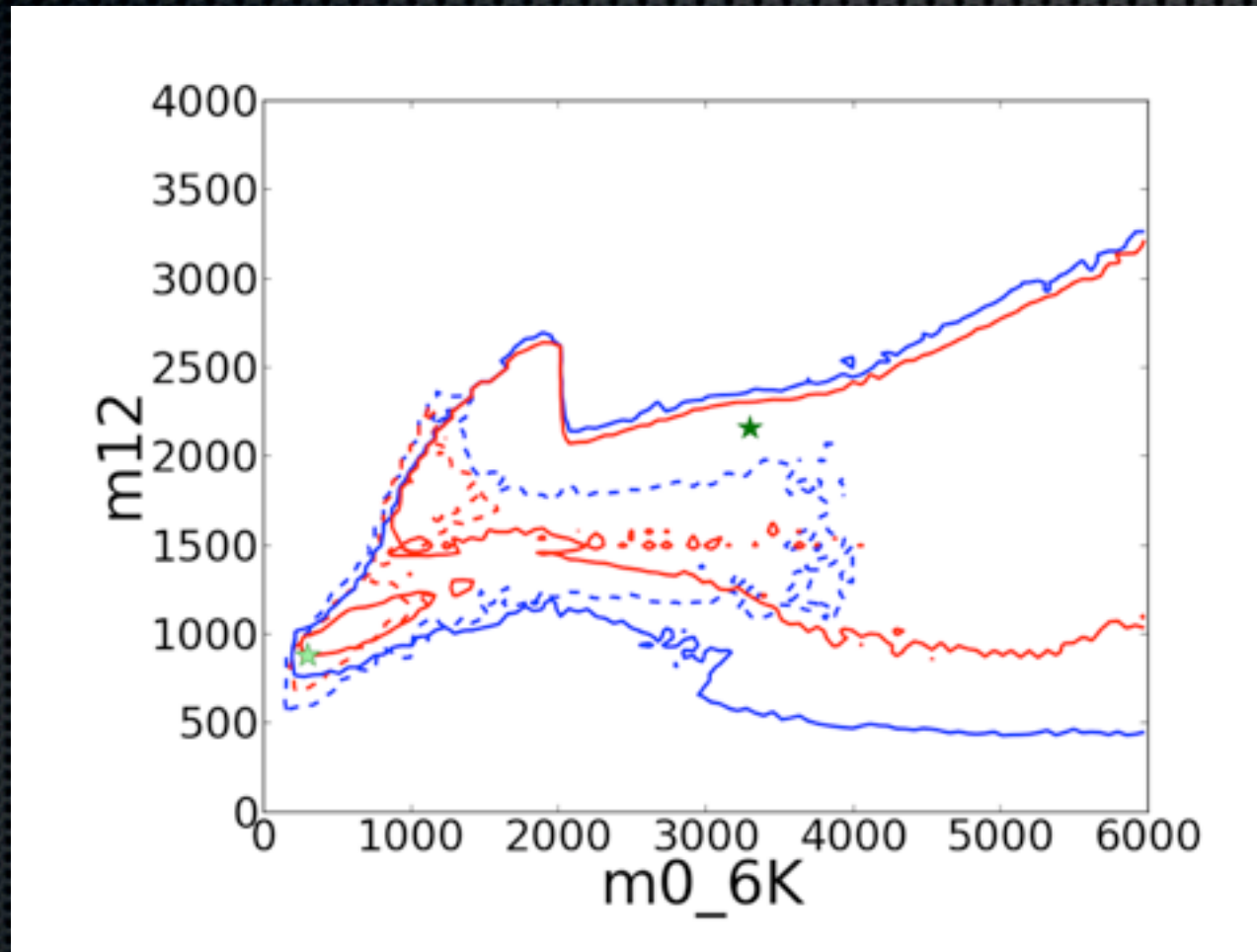
- ✦ MCMC technique to sample efficiently the SUSY parameter space, and thereby construct the χ^2 probability function
- ✦ Combines SoftSusy, FeynHiggs, SuperFla, SuperIso, MicrOmegas, and SSARD
- ✦ Purely frequentist approach (no priors) and relies only on the value of χ^2 at the point sampled and not on the distribution of sampled points.
- ✦ 12 million points sampled (CMSSM)

$$\begin{aligned}\chi^2 = & \sum_i^N \frac{(C_i - P_i)^2}{\sigma(C_i)^2 + \sigma(P_i)^2} \\ & + \chi^2(M_h) + \chi^2(\text{BR}(B_s \rightarrow \mu\mu)) \\ & + \chi^2(\text{SUSY search limits}) \\ & + \sum_i^M \frac{(f_{\text{SM}_i}^{\text{obs}} - f_{\text{SM}_i}^{\text{fit}})^2}{\sigma(f_{\text{SM}_i})^2}\end{aligned}$$

Buchmueller, Cavanaugh, De Roeck, Dolan, Ellis, Flacher, Heinemeyer, Isidori, Marrouche, Martinez Santos, Olive, Rogerson, Ronga, de Vries, Weiglein

$\Delta\chi^2$ map of $m_0 - m_{1/2}$ plane

Final run I



Preliminary

Buchmueller, Cavanaugh, De Roeck, Dolan, Ellis, Flacher,
Heinemeyer, Isidori, Marrouche, Martinez Santos, Olive, Rogerson,
Ronga, de Vries, Weiglein

Preliminary

COMPARISON OF BEST FIT POINTS PRE AND POST LHC

Model	Data set	Minimum $\chi^2/\text{d.o.f.}$	Prob- ability	m_0 (GeV)	$m_{1/2}$ (GeV)	A_0 (GeV)	$\tan \beta$
CMSSM $\mu > 0$	ATLAS 7 TeV	32.6/23	8.8%	340	910	2670	12
	ATLAS _{20/fb} (low)	35.6/23	4.5%	710	1070	3580	21
	ATLAS _{20/fb} (high)	34.9/23	5.3%	3310	2180	-1490	51
CMSSM $\mu < 0$	ATLAS _{20/fb} (low)	37.8/23	2.7%	2100	660	4930	11
	ATLAS _{20/fb} (high)	36.9/23	3.3%	6490	2430	-3300	36
NUHM1 $\mu > 0$	ATLAS 7 TeV	30.4/22	10.9%	360	1080	4990	9
	ATLAS _{20/fb} (low)	33.1/22	6.0%	470	1270	5700	11
	ATLAS _{20/fb} (high)	32.7/22	6.6%	1380	3420	-3140	39
“SM”	ATLAS _{20/fb} (high)	36.5/24	5.0%	-	-	-	-

Buchmueller, Cavanaugh, De Roeck, Dolan, Ellis, Flacher,
Heinemeyer, Isidori, Marrouche, Martinez Santos, Olive, Rogerson,
Ronga, de Vries, Weiglein

Effective four-fermion Lagrangian

$$\mathcal{L} = \bar{\chi} \gamma^\mu \gamma^5 \chi \bar{q}_i \gamma_\mu (\alpha_{1i} + \alpha_{2i} \gamma^5) q_i + \alpha_{3i} \bar{\chi} \chi \bar{q}_i q_i \\ + \alpha_{4i} \bar{\chi} \gamma^5 \chi \bar{q}_i \gamma^5 q_i + \alpha_{5i} \bar{\chi} \chi \bar{q}_i \gamma^5 q_i + \alpha_{6i} \bar{\chi} \gamma^5 \chi \bar{q}_i q_i$$

The terms proportional to $\alpha_1, \alpha_4, \alpha_5, \alpha_6$,
lead to velocity-dependent cross sections

Remaining terms:

α_2 : Spin-dependent cross section

α_3 : Spin-independent cross section

Uncertainties from hadronic matrix elements

The scalar cross section

$$\sigma_3 = \frac{4m_r^2}{\pi} [Z f_p + (A - Z) f_n]^2$$

where

$$\frac{f_p}{m_p} = \sum_{q=u,d,s} f_{Tq}^{(p)} \frac{\alpha_{3q}}{m_q} + \frac{2}{27} f_{TG}^{(p)} \sum_{c,b,t} \frac{\alpha_{3q}}{m_q}$$

and

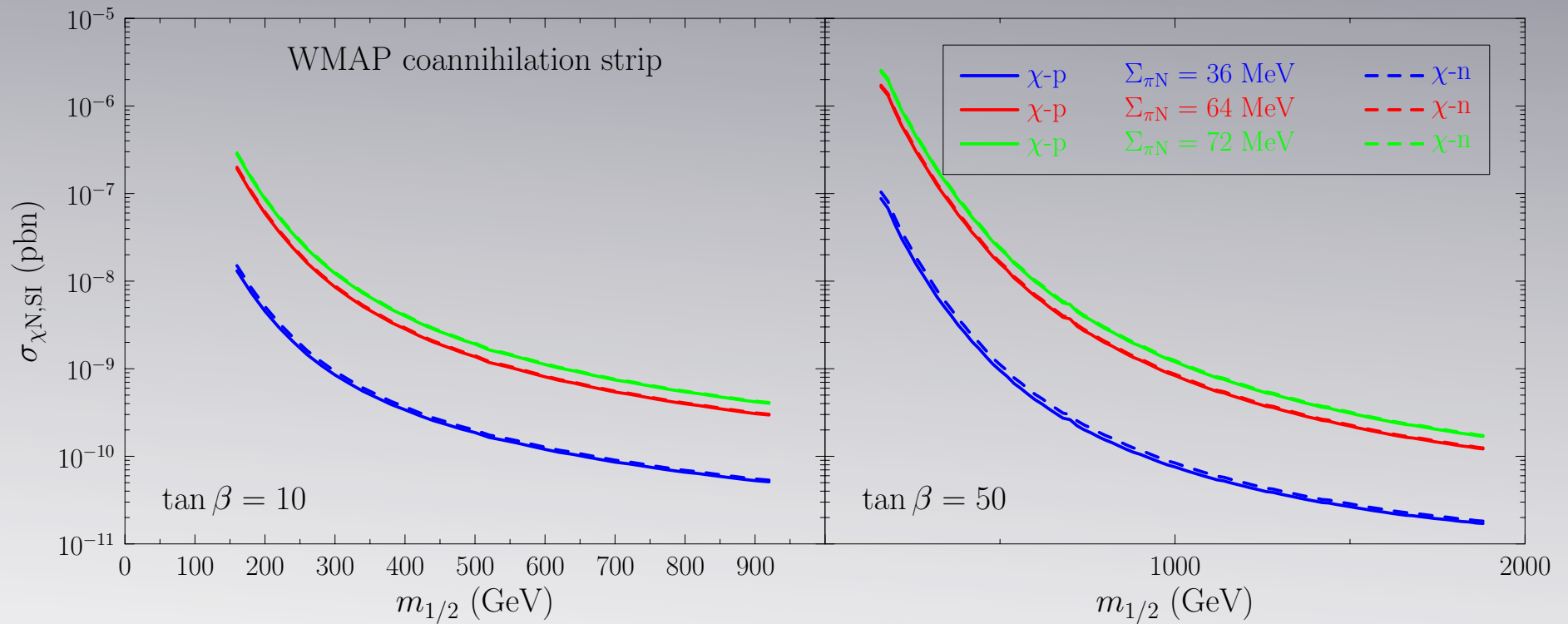
$$m_N f_{Tq}^{(N)} \equiv \langle N | m_q \bar{q}q | N \rangle ,$$

determined by

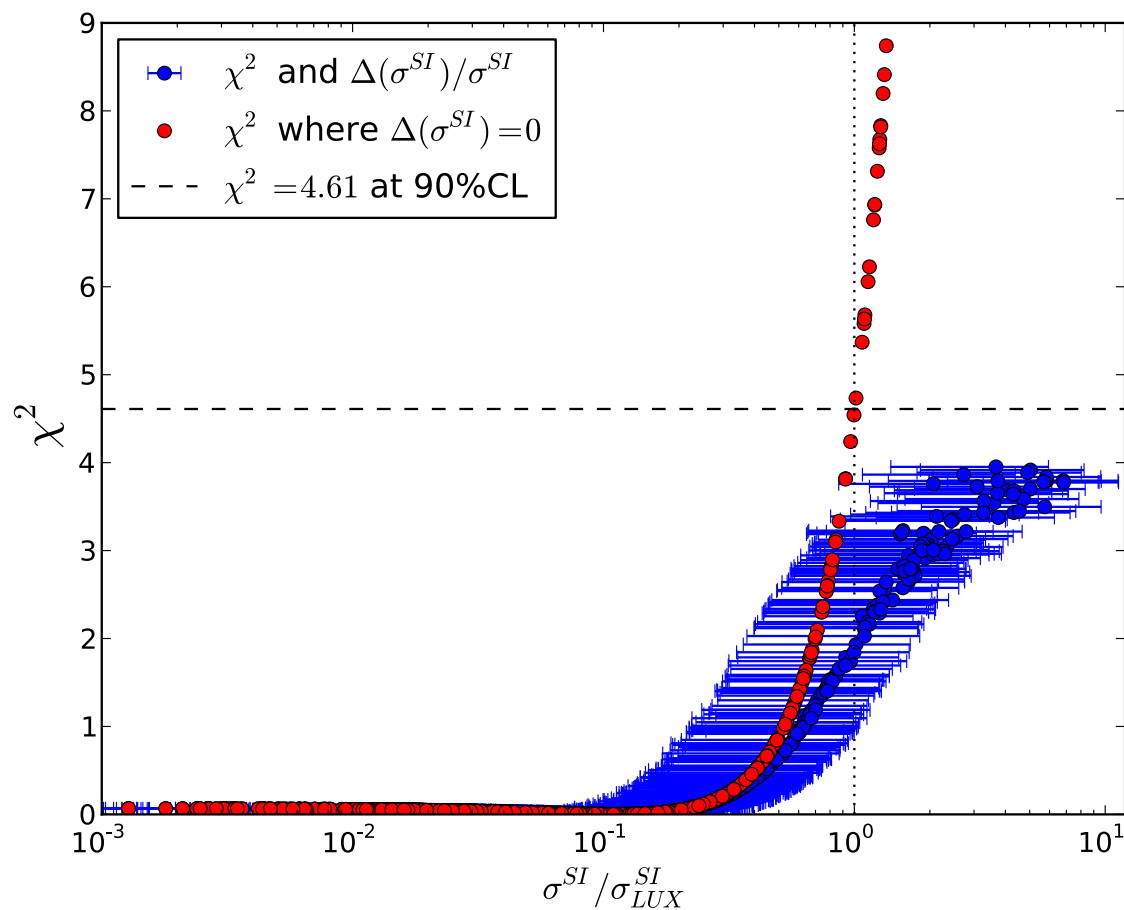
$$f_{T_{u,d}} \propto \Sigma_{\pi N} \quad f_{T_s} \propto \Sigma_{\pi N} y \quad y = 1 - \sigma_0 / \Sigma_{\pi N}$$

$$\Sigma_{\pi N} = 50 \pm 7 \text{ MeV} \quad \sigma_0 = 36 \pm 7 \text{ MeV}$$

Uncertainties due to $\Sigma_{\pi N}$



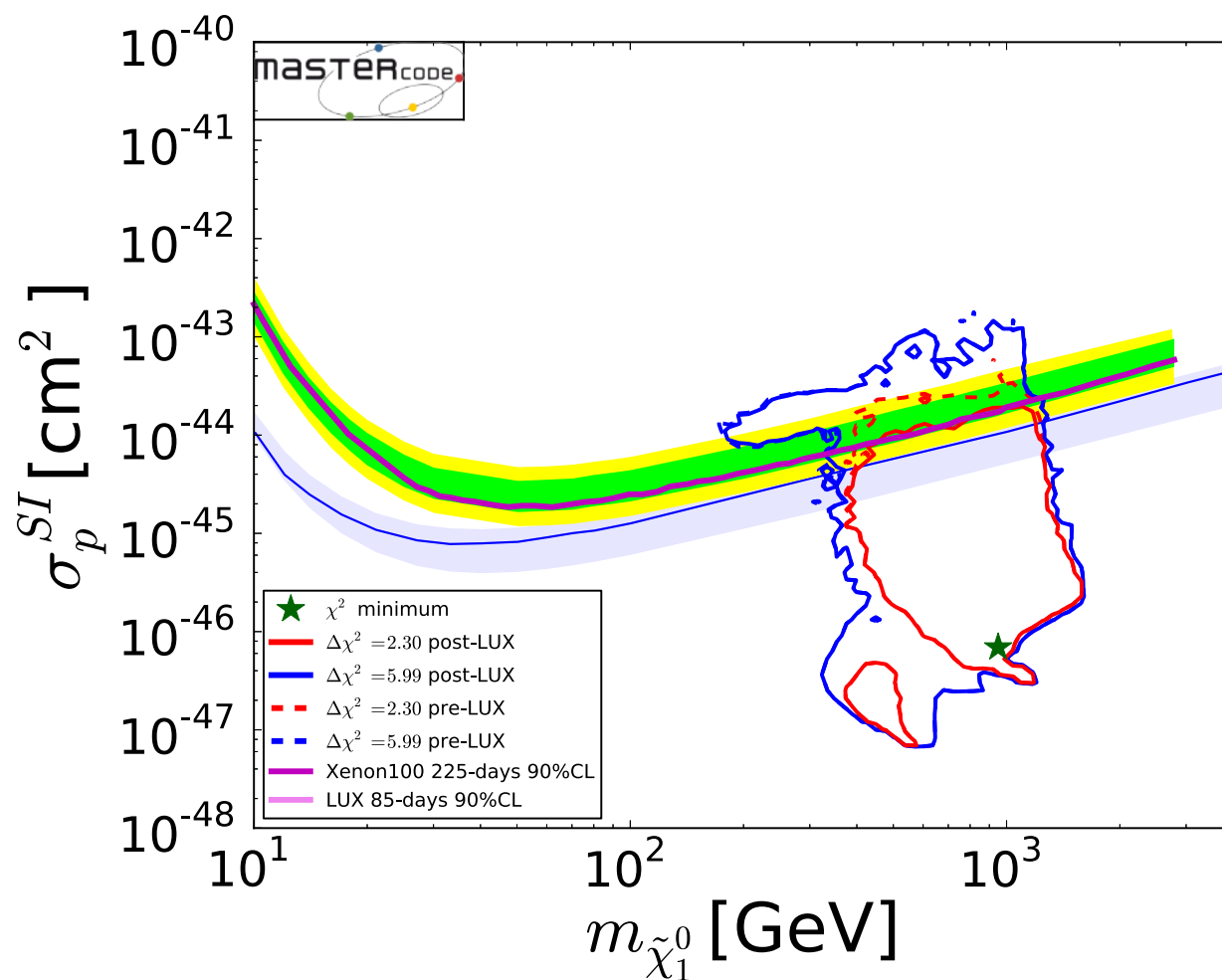
Elastic cross sections



Preliminary

Buchmueller, Cavanaugh, De Roeck, Dolan, Ellis, Flacher, Heinemeyer, Isidori, Marrouche, Martinez Santos, Olive, Rogerson, Ronga, de Vries, Weiglein

Elastic cross sections



Preliminary

Buchmueller, Cavanaugh, De Roeck, Dolan, Ellis, Flacher, Heinemeyer, Isidori, Marrouche, Martinez Santos, Olive, Rogerson, Ronga, de Vries, Weiglein

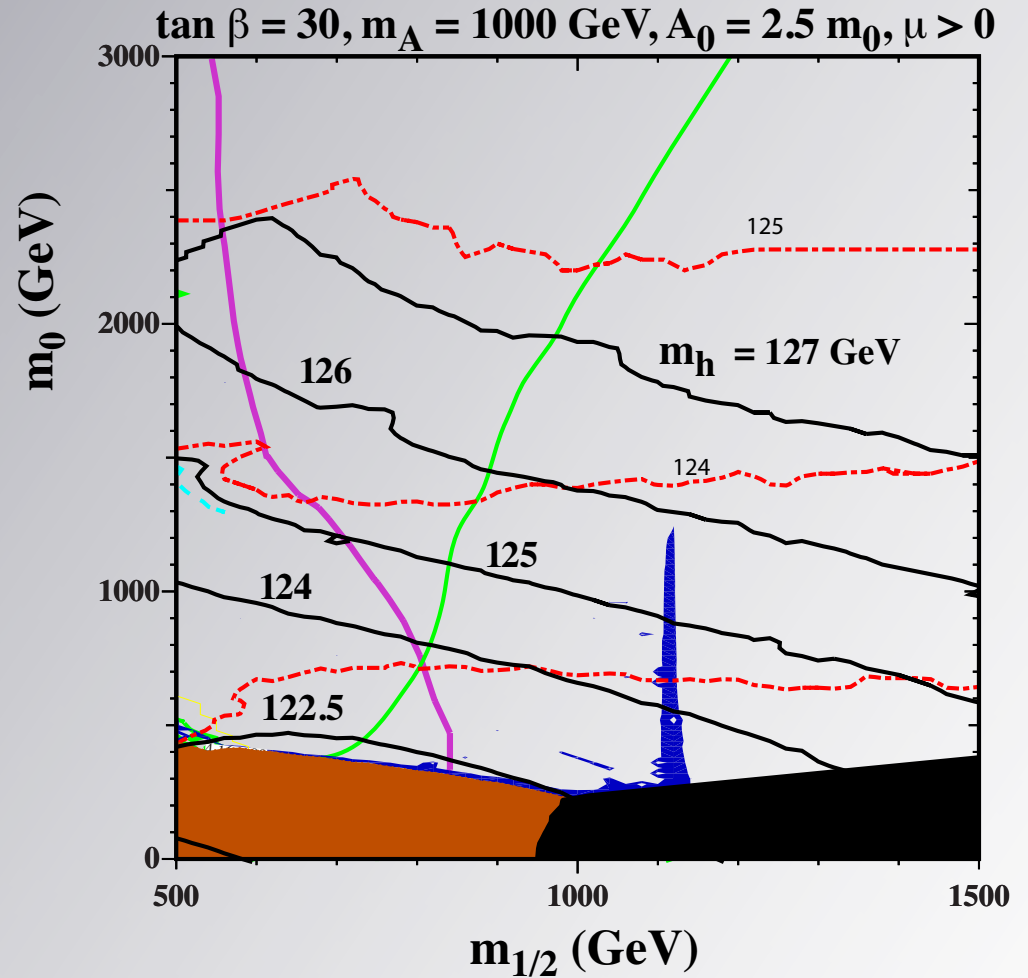
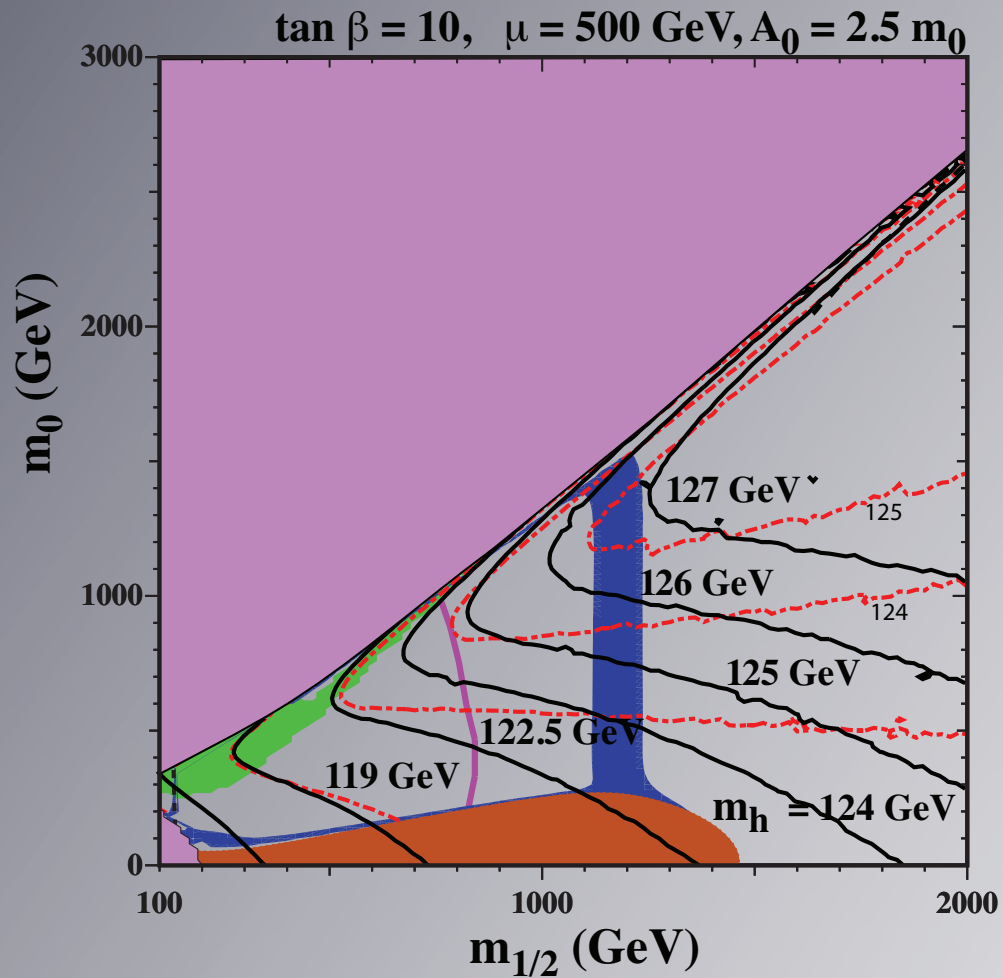
May require more general models
which are concordant with LHC MET;
Higgs; and $B_s \rightarrow \mu^+\mu^-$; and Dark Matter

May require more general models
which are concordant with LHC MET;
Higgs; and $B_s \rightarrow \mu^+\mu^-$; and Dark Matter

Other Possibilities

- ✦ NUHM1,2: $m_1^2 = m_2^2 \neq m_0^2$, $m_1^2 \neq m_2^2 \neq m_0^2$
 - ✦ μ and/or m_A free
- ✦ subGUT models: $M_{in} < M_{GUT}$
 - ✦ with or without mSUGRA

NUHM1 models with μ or m_A free



Ellis, Luo, Olive, Sandick

Buchmueller, Dolan, Ellis, Hahn, Heinemeyer, Hollik, Marrouche, Olive, Rzehak, de Vries, Weiglein

Moving beyond the CMSSM-like models

Moving beyond the CMSSM-like models

Models with
Strongly Stabilized Moduli;
Pure Gravity Mediation (PGM)

Moving beyond the CMSSM-like models

Models with Strongly Stabilized Moduli; Pure Gravity Mediation (PGM)

- ✦ Usually ignored in phenomenological studies of the MSSM
- ✦ In general, many moduli:
- ✦ Volume Modulus: destabilization
- ✦ Polonyi-like fields: cosmological entropy production; gravitino production; LSP production....

Consider a Polonyi-like modulus
but with a non-minimal kinetic term

$$K = Z\bar{Z} - \frac{(Z\bar{Z})^2}{\Lambda^2}$$

Dine et al,
Kitano

and Polonyi superpotential

$$W = \mu^2(Z + \nu)$$

$$\langle z \rangle_{\text{Min}} \simeq \frac{\Lambda^2}{\sqrt{6}}, \quad \langle \chi \rangle = 0, \quad \nu \simeq \frac{1}{\sqrt{3}}$$

where $Z = \frac{1}{\sqrt{2}}(z + i\chi)$

Impact on Phenomenology

$$m_{3/2} = \langle e^{K/2} W \rangle \simeq \mu^2 / \sqrt{3}$$

$$m_{z,\chi}^2 \simeq \frac{12 m_{3/2}^2}{\Lambda^2} \gg m_{3/2}^2$$

Soft scalar masses $m_0^2 = m_{3/2}^2$

A terms $A_0 \simeq \frac{1}{2} m_{3/2} \Lambda^2 + \text{anomalies}$

gaugino masses anomalies

Impact on Phenomenology

$$m_{3/2} = \langle e^{K/2} W \rangle \simeq \mu^2 / \sqrt{3}$$

$$m_{z,\chi}^2 \simeq \frac{12 m_{3/2}^2}{\Lambda^2} \gg m_{3/2}^2$$

Soft scalar masses $m_0^2 = m_{3/2}^2$

A terms $A_0 \simeq \frac{1}{2} m_{3/2} \Lambda^2$ + anomalies

gaugino masses anomalies

Massive scalar sector as in split susy, with anomaly mediation for A-terms and gaugino masses

Pure Gravity Mediation

- ✦ Two parameter model!

- ✦ $m_0 = m_{3/2}; \tan \beta$

- ✦ gaugino masses (and A-terms) generated through loops

$$M_1 = \frac{33}{5} \frac{g_1^2}{16\pi^2} m_{3/2} ,$$

$$M_2 = \frac{g_2^2}{16\pi^2} m_{3/2} ,$$

$$M_3 = -3 \frac{g_3^2}{16\pi^2} m_{3/2} .$$

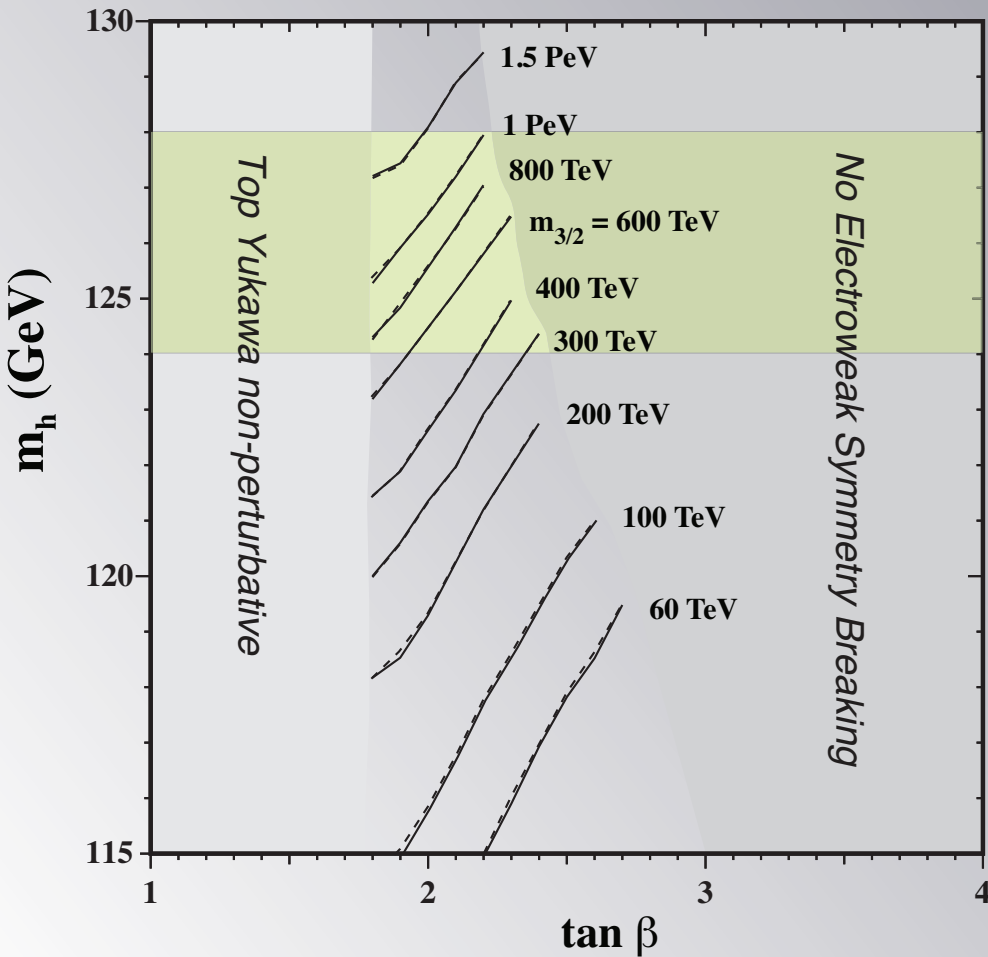
- ✦ \Rightarrow Push towards very large masses

Pure Gravity Mediation

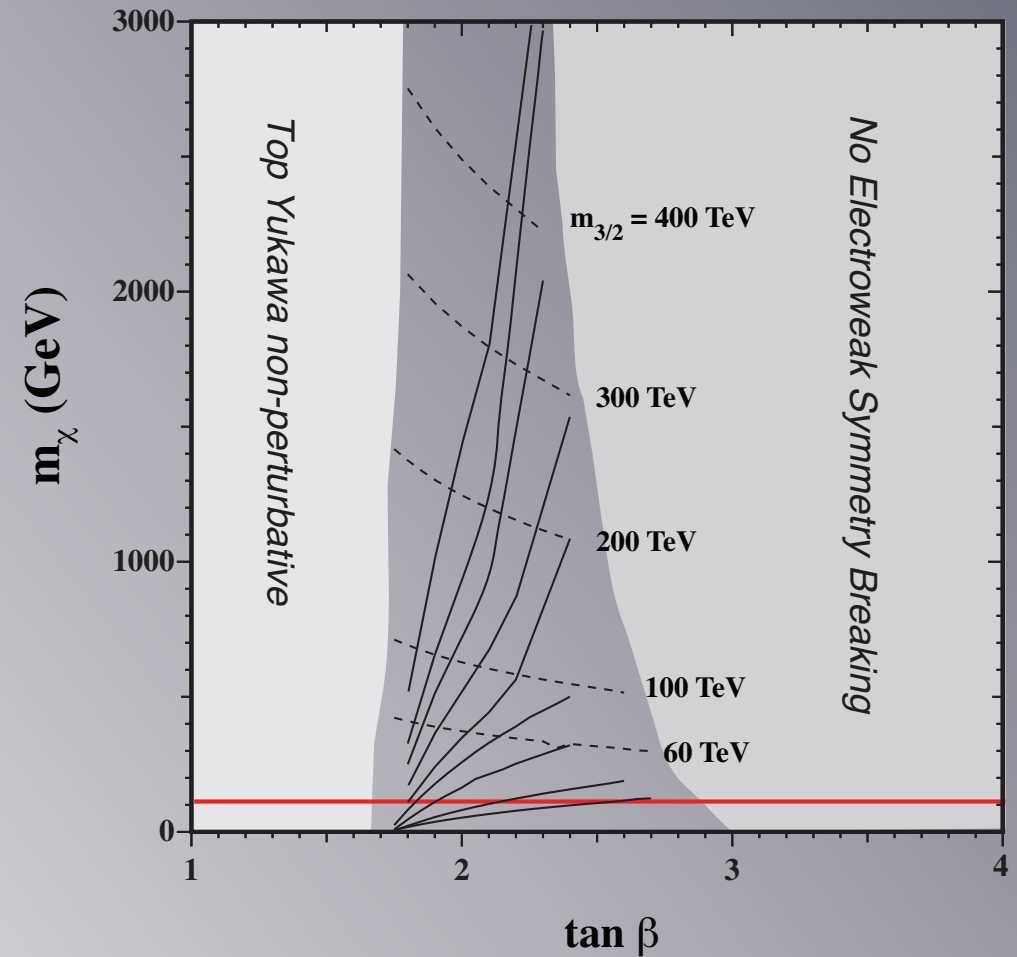
- ✦ The sfermion and gravitino have masses $O(100)$ TeV.
- ✦ The higgsino and the heavier Higgs boson also have masses $O(100)$ TeV.
- ✦ The gaugino masses are in the range of hundreds to thousands of GeV.
- ✦ The LSP is the neutral wino which is nearly degenerate with the charged wino.
- ✦ The lightest Higgs boson mass is consistent with the observed Higgs-like boson, i.e. $m_h \sim 125 - 126$ GeV.

Phenomenological Aspects

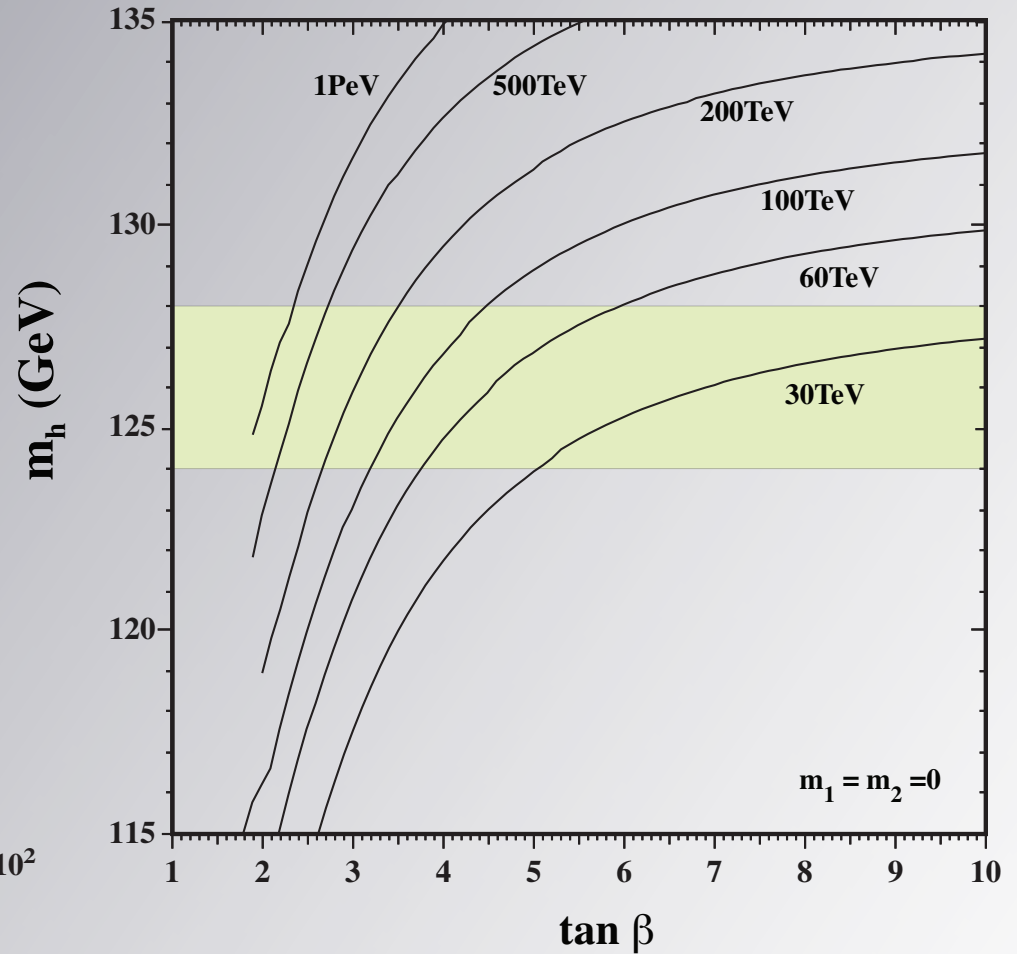
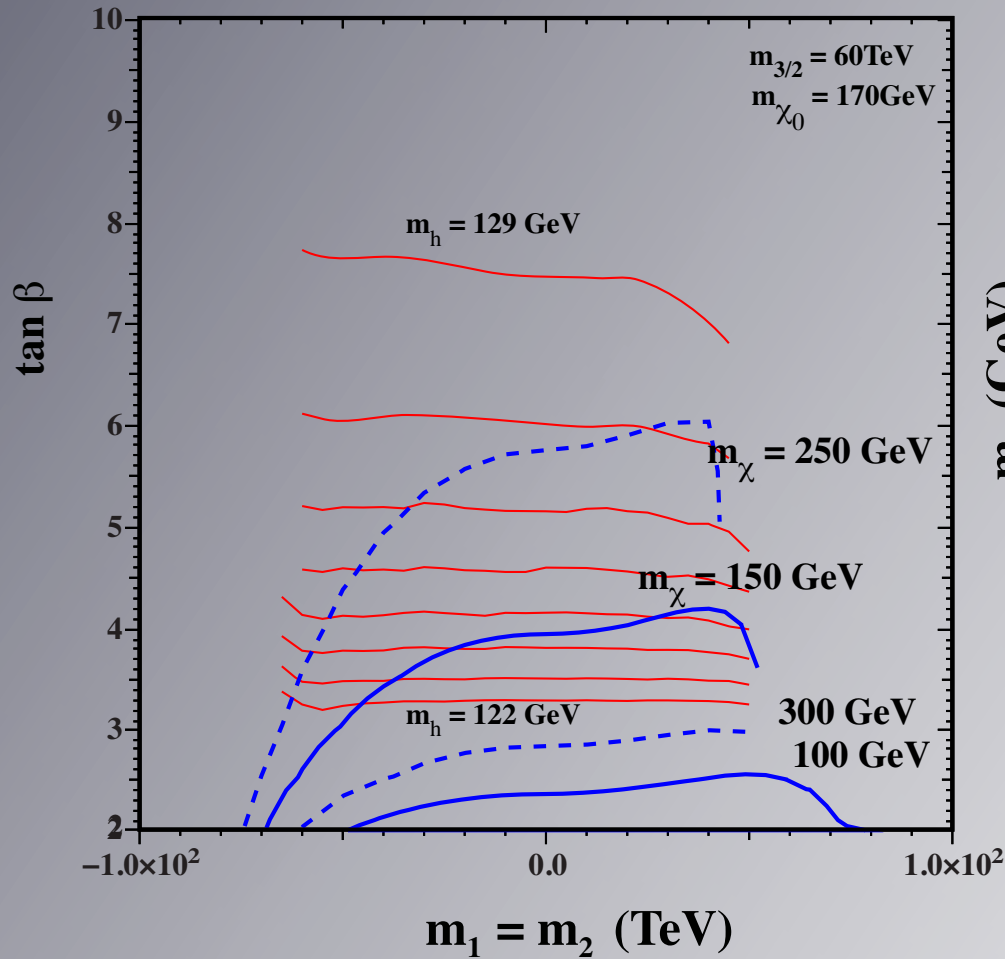
Higgs Mass



Neutralino mass

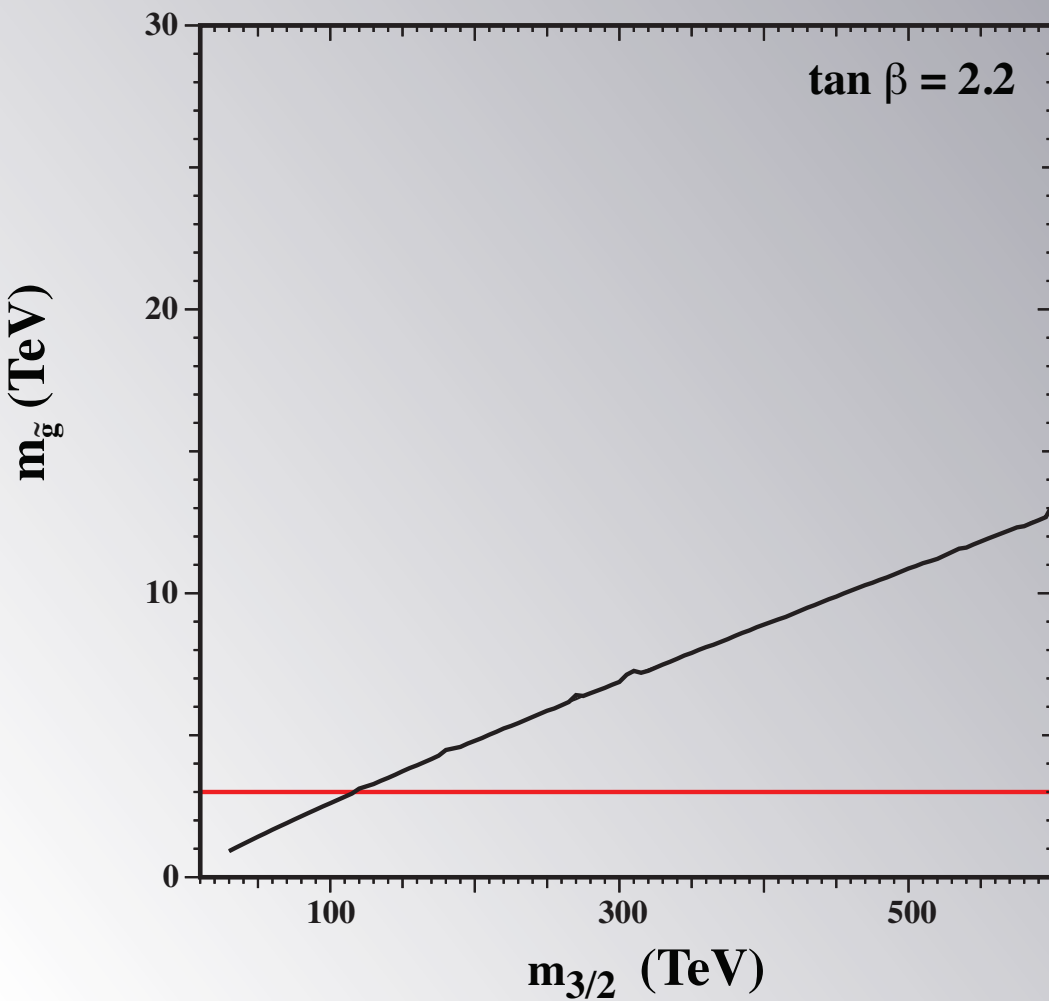


Somewhat more freedom with non-universal Higgs masses

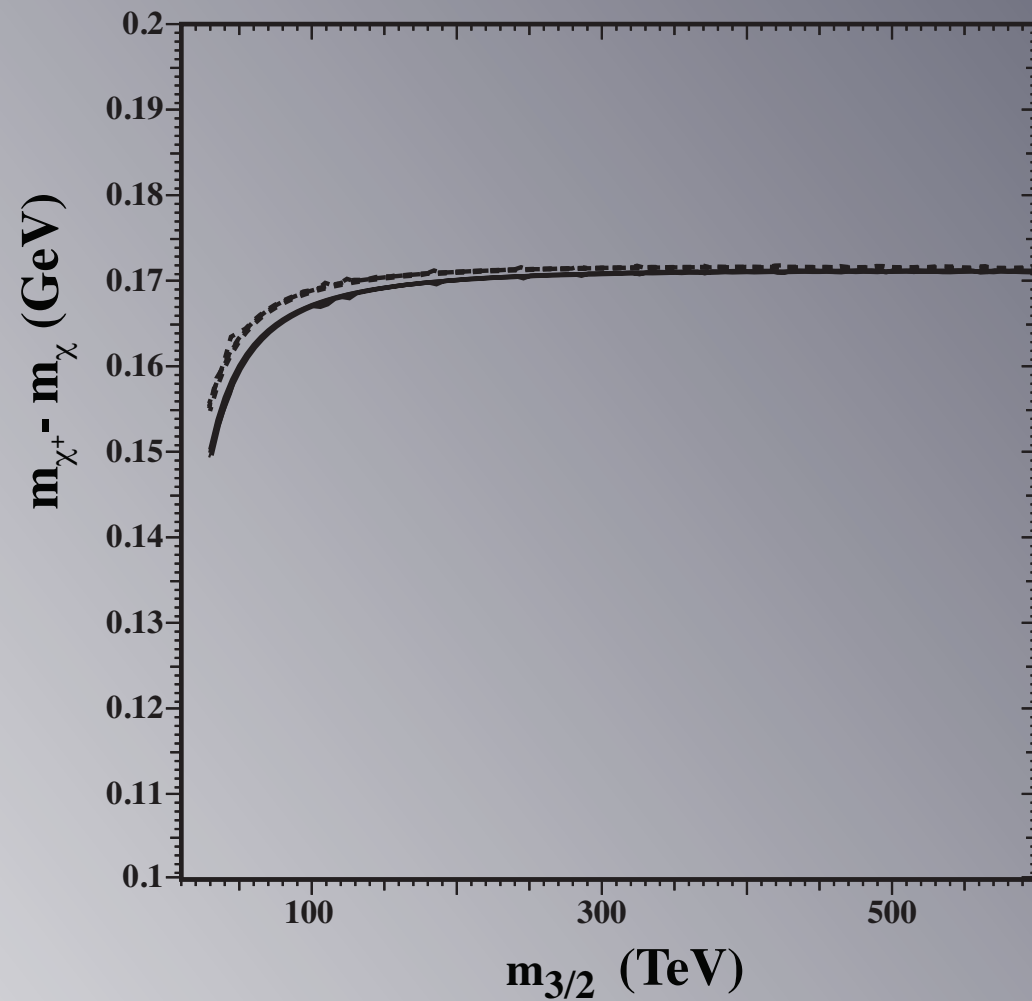


Phenomenological Aspects

gluino Mass



chargino mass



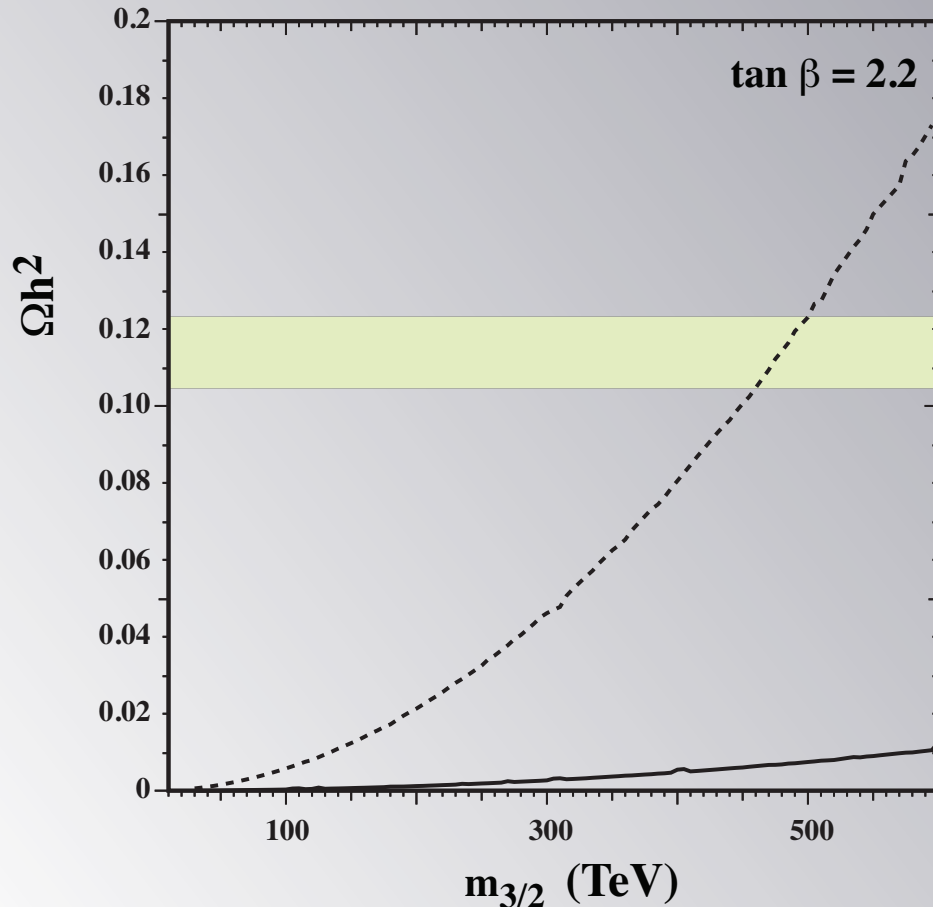
Dark Matter

- Dark matter is something else (axion)
- LSPs from gravitino or moduli (Z) decay
- $m_{3/2} \sim 650 \text{ TeV}$, and $\Omega h^2 \sim 0.11$

$$\Omega_\chi h^2 = \frac{m_\chi}{m_{3/2}} \Omega_{3/2} h^2 = 0.4 \left(\frac{m_\chi}{\text{TeV}} \right) \left(\frac{T_R}{10^{10} \text{ GeV}} \right)$$

Other Phenomenological Aspects

Dark Matter:
LSP is a wino



Potential problem for wino
dark matter from Fermi/HESS
(Fan + Reese;
Cohen, Lianti, Pierce, Slatyer)

Summary

- ✦ LHC susy and Higgs searches have pushed CMSSM-like models to “corners”
- ✦ Though many phenomenological solutions are still viable
- ✦ Models with strong moduli stabilization:
 - ✦ easier for inflation,
 - ✦ no cosmological problems
 - ✦ interesting phenomenology
- ✦ Heavy scalar spectrum with anomaly mediated gaugino masses
- ✦ Challenge lies in detection strategies