

Photons signaling a QCD phase transition in neutron star mergers

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arXiv:1305.7397, PRD 88, 083006 (2013)



Motivation from Cosmology/Astrophysics

short GRBs thought to signal **neutron star-neutron star mergers**

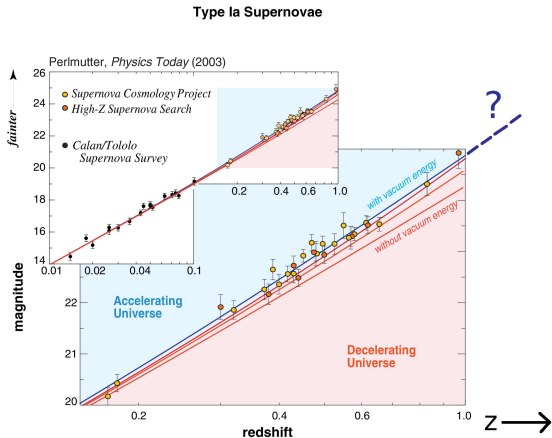
Total energy of burst:

$$E_{\text{tot}} \simeq 10^{51} - 10^{53} \text{ erg}$$

Visible at cosmological distances $z > 1$

1. Use microphysics to predict common spectral feature \rightarrow standard candle for larger z

2. Solidify short GRB - neutron star merger relationship



Neutron star binaries

- binary merger rate $2 - 60(10^6 \text{ yr})^{-1}$ per galaxy

C. Kim, V. Kalogera and D. R. Lorimer, *Ap. J.* **584** (2003) 985

- expected to be origin of “short” Gamma Ray Bursts (SGRBs) (along with neutron star-black hole mergers)

- primary motivation for gravitational wave detectors

LIGO/Virgo collaboration, *Class. Quantum Grav.* **27**, 173001, (2010)

Author	NS-NS		BH-NS		Method
	LIGO	AdLIGO	LIGO	AdLIGO	
Kim et al. [142]	5e-3	27			Empirical
Nakar et al. [194]		~2		~20.0	SGRBs
Guetta & Stella [126]	7.0e-3	22	7.0e-2	220	SGRBs
Voss & Tauris [318]	6.0e-4	2.0	1.2e-3	4.0	Pop. Synth. - SFR
de Freitas Pacheco et al. [79]	8.0e-4	6.0			Pop. Synth. - SFR
Kalogera et al. [139]	1.0e-2	35	4.0e-3	20	Pop. Synth. - NS-NS
O’Shaughnessy et al. [214]	1.0e-2	10	1.0e-2	10	Pop. Synth. - NS-NS

Event rate estimates from SFR=star formation rate, NS-NS=observed population of binary NSs

Faber & Rasio, *Living Rev. Relativity*, **15** (2012)

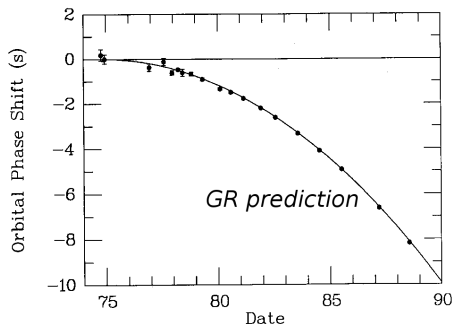
Why and how neutron star binaries merge

1. quasi-equilibrium decay
due to Gravitational radiation

$$\frac{dE_{\text{GW}}}{dt} \simeq -\frac{32}{5} \frac{\mu^2 M}{M^2 r} \left(\frac{GM}{r}\right)^4$$

$$M = m_1 + m_2 \quad \mu = \frac{m_1 m_2}{M}$$

Radiates angular momentum
making orbit more circular



Pulsar B1913+16

Other binary pulsars now identified —

Hailed as one of most precise tests of General Relativity

milliseconds before the collision

Kinematics of collision follow from radiation loss
contact at $r_{\text{coll}} \simeq 2R_*$

Kinetic+Potential energy of binary system with GR corrections:

$$K + V = -\frac{\mu x}{2} \left(1 - \frac{3}{4} \left(1 + \frac{\mu}{9M} \right) x - \left(27 - 19 \frac{\mu}{M} + \frac{2\mu^2}{3M^2} \right) \frac{x^2}{8} + \dots \right)$$

GR gauge invariant variable:

$$x = (GM\omega)^{2/3} \simeq \frac{GM}{r} \simeq v^2$$

Subtracting Newtonian potential $V \simeq -G\mu Mx/r$

$$v_r^2 = \left(19 - \frac{5\mu}{3M} \right) \frac{x^2}{4} \simeq 0.20c^2$$

- estimated correction $\sim +5\%$ dissipation of angular momentum

Conditions of the matter in the collision

From radial velocity at collision, $v_r^2 \simeq 0.2c^2$

lower estimate of kinetic energy per baryon

$$\frac{E}{N} \simeq \frac{1}{2} \frac{m_N}{1 + \delta M/M} v_r^2 \simeq 85 \text{ MeV} \simeq T$$

(mass defect $\delta M/M$ corrects gravitational binding in counting baryons)

Near surface of initial stars : $n/n_{\text{sat}} \simeq 0.15 - 0.6$

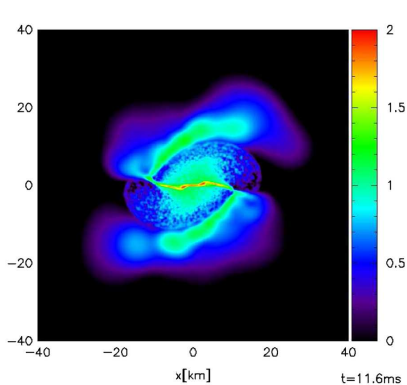
Conservation of baryon number where matter overlaps in collision

\Rightarrow amplify density by factor 2 – 4

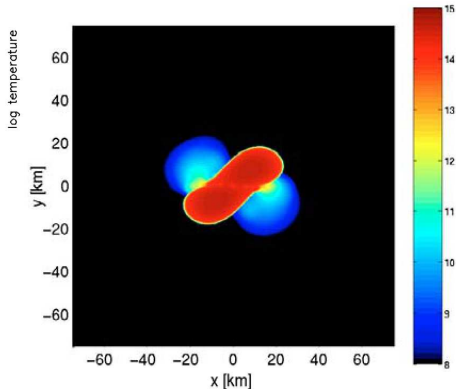
Expected density $n/n_{\text{sat}} \simeq 0.3 - 2.4$

Baryonic chemical potential $\mu_B \simeq 944 - 1240 \text{ MeV} \simeq 3\mu_q$

Numerical simulation



left, temperature in units MeV

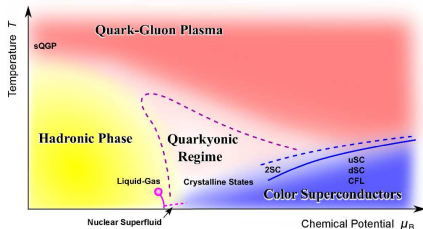
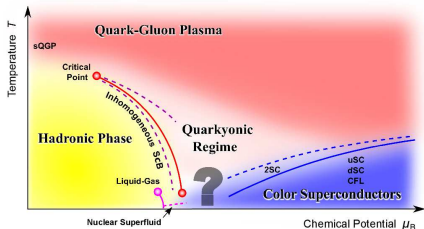


right: log plot of density in g/cm^3
 $n_{\text{sat}} \rightarrow 10^{14} \text{g}/\text{cm}^3$

Bauswein, Janka & Oechslin, PRD 82 (2010) 084043;

Oechslin, Bauswein & Janka, A & A, 467, 395 (2007).

Sensitive region of the QCD phase diagram

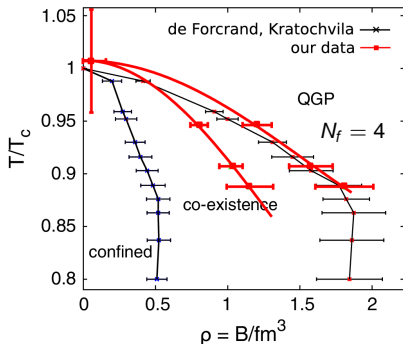


Fukushima & Sasaki, arXiv:1301.6377

- quarkyonic phase?
- order of transition?

Li, Alexandru, Liu and Meng,

PRD 82, 054502 (2010) →



Comparison of timescales

QCD timescales in general many orders of magnitude shorter:

$$\tau_{\text{QCD}} \lesssim 10^{-20} \text{ s} \ll \tau_{\text{weak}} \sim 10^{-7} - 10^{-6} \text{ s} \ll \tau_{\text{merger}} \sim 10^{-3} \text{ s}$$

- QCD timescales:

kinetic equilibrium: $\tau \simeq (n\sigma v)^{-1} \sim 10^{-23} \text{ s}$

bulk transport: thermal conductivity ($\sim 10^{-20} \text{ s}$), energy loss (we show: $\sim 10^{-14} \text{ s}$)

- Weak interactions: “fast” Urca processes $\tau_{\text{weak}}^{-1} \simeq \frac{1}{T^4} \left. \frac{dE}{dt} \right|_{\text{Urca}}$
e.g. in quark phases $u + e^- \rightarrow d + \nu_e$, $d \rightarrow u + e^- + \bar{\nu}_e$

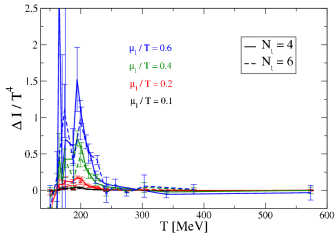
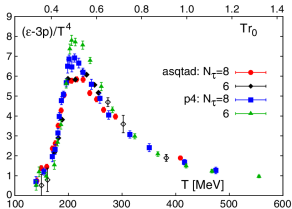
- Merger timescale: final orbital period $\tau_{\text{merger}} \sim 1 \text{ ms}$
numerical simulations see system relax (quasi-)equilibrium state
 $\tau \sim 5 - 10 \text{ ms}$

★ Matter close to equilibrium with respect to QCD processes

We will find that all timescales remain shorter than τ_{weak}

→ self-consistent to neglect weak flavor changing reactions

Energy-momentum trace peaks near phase change

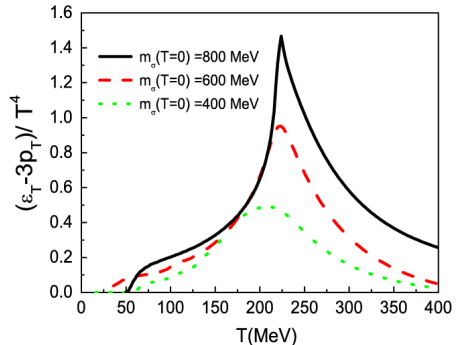


Lattice results:

MILC, PRD 80, 014504 (2009);

PoS LATTICE, 2008, 171 (2008)

$$g_{\mu\nu}\Theta^{\mu\nu} = \Theta^{\mu}_{\mu} = \epsilon - 3p$$



$O(4)$ model (same universality class as QCD)
same qualitative behavior across models and parameters

Li & Huang, PRD80, 034023 (2009)

Energy-momentum Trace as collective mode

Collective modes described in finite temperature field theory, allows consistent inclusion of low-energy effective couplings

Energy-momentum two-point function (=propagator of collective mode)

$$\langle \Theta_{\mu\mu}(t, \vec{x}) \Theta_{\mu\mu}(t', \vec{x}') \rangle = G_{\mu\mu, \nu\nu}(t, \vec{x}; t', \vec{x}')$$

Spectral representation : $\frac{1}{\omega} \text{Im} \left[G_{\mu\mu, \nu\nu}(\omega, \vec{k}) \right] = \frac{\rho_{\sigma}(\omega)}{\omega} \simeq \frac{9}{\pi} \zeta$

ζ = bulk viscosity, the dissipation that damps collective mode

see Kapusta & Gale, *Finite-T Field Theory*; Meyer JHEP **1004**, 099 (2010)

decay of collective mode seen in imaginary part of polarization

$$\Gamma = \text{Im}\Pi = \text{Im} \left[\text{Diagram} \right]$$

hadron level theory — no explicit quark or gluon degrees of freedom

How photons couple: Triangle anomaly

Effective theory for collective mode

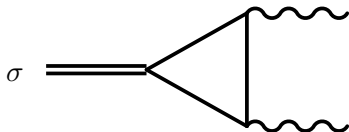
$$\langle 0 | \Theta_{\mu}^{\mu} | \sigma \rangle = m_{\sigma}^2 f_{\sigma}$$

(like pion low-energy parameters)

Quarks carry electric charge

⇒ photons couple through loop

$$\mathcal{L}_{\text{eff}} = g_{\sigma\gamma\gamma} \sigma F^{\mu\nu} F_{\mu\nu}$$



Ellis & Lanik, PLB **175**, 83 (1986)

Matching onto 2-photon width of $f_0(600)$ meson: [PDG]

(lowest state with right quantum numbers)

$$g_{\sigma\gamma\gamma} \simeq (50 \text{ GeV})^{-1} \quad m_{\sigma} \simeq 550 \text{ MeV} \quad f_{\sigma} \simeq 100 \text{ MeV}$$

[Basar, Kharzeev, Skokov, Phys. Rev. Lett. **109**, 202303 (2012)]

★ Effective theory valid for momenta $k \ll 50 \text{ GeV}$

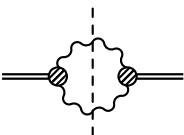
parameters from $\sigma - f_0$ meson in vacuum

** assume in-medium properties contained in spectral function $\rho_{\sigma}(\omega, \vec{k})$

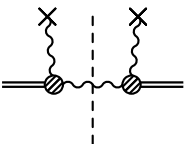
Photon production: two cases

Calculate imaginary part of photon polarization (at $k \ll 50$ GeV)

a) final state \rightarrow two real photons

$$\Gamma = \text{Im}\Pi = \text{---} \text{---} \text{---} \text{---}$$
A Feynman diagram representing the imaginary part of the photon polarization function, $\Gamma = \text{Im}\Pi$. It shows two external fermion lines (represented by double lines) entering from the left and exiting to the right. These lines are connected by a loop of a fermion (represented by a shaded circle). The loop is connected to a wavy line representing a photon, which then splits into two real photons (represented by wavy lines) exiting to the right. A vertical dashed line is drawn through the center of the loop.

b) recoils off external magnetic field \rightarrow one photon

$$\Gamma = \text{Im}\Pi = \text{---} \text{---} \text{---} \text{---}$$
A Feynman diagram representing the imaginary part of the photon polarization function, $\Gamma = \text{Im}\Pi$. It shows two external fermion lines (represented by double lines) entering from the left and exiting to the right. These lines are connected by a loop of a fermion (represented by a shaded circle). The loop is connected to a wavy line representing a photon, which then recoils off an external magnetic field (represented by a vertical dashed line). The photon then splits into one real photon (represented by a wavy line) exiting to the right and another photon (represented by a wavy line) exiting upwards, marked with an 'X'.

subdominant to 2-photon unless stellar $|\vec{B}| \sim 100 B_{\text{QED}}$

$B_{\text{QED}} = 4.41 \times 10^{13}$ Gauss – (in fact, magnetic fields of this scale and higher are inferred from observations)

Energy loss rates

Obtain energy release by integrating weighted with k^μ

$$\Gamma_{2\gamma}^\mu = \int \frac{d^3k_1}{(2\pi)^3 2\omega_1} \frac{d^3k_2}{(2\pi)^3 2\omega_2} (k_1^\mu + k_2^\mu) \omega_1 \omega_2 \frac{d\Gamma_{2\gamma}}{d^3k_1 d^3k_2}$$

Dense plasma \rightarrow Refractive medium: finite mean free path

$$l_f \sim (em_p)^{-1} \sim 100 \text{ fm}$$

\Rightarrow effective photon mass: fix photon energy from medium to vacuum:

$$\vec{k}_{\text{vac}}^2 = \omega^2$$

$$\text{Energy emitted from plasma: } \frac{dE}{dVdt} = \Gamma^0$$

$$\left. \frac{dE}{d^4x} \right|_{2\gamma} = 7.13 \frac{\text{TeV}}{\text{fm}^3 \text{s}} \frac{\zeta}{s} \frac{s}{s_0} \left(\frac{m_p}{15 \text{ MeV}} \right)^{10} I_{2\gamma}(m_p/T)$$

$$\left. \frac{dE}{d^4x} \right|_{B\gamma} = 5.51 \frac{\text{GeV}}{\text{fm}^3 \text{s}} \frac{\zeta}{s} \frac{s}{s_0} \left(\frac{m_p}{15 \text{ MeV}} \right)^6 \frac{\vec{B}^2}{B_{\text{QED}}^2} I_{B\gamma}(m_p/T)$$

$s_0 =$ entropy of (u, d, e) plasma at $T = 50 \text{ MeV}$ and $2n_{\text{nuc}}$

Energy loss and surface cooling

Timescale of energy loss:

$$\frac{1}{\tau_E} = \frac{1}{\varepsilon} \frac{dE}{dVdt} \simeq (10^{-6} \text{ s})^{-1} \left(\frac{m_p}{15 \text{ MeV}} \right)^{10} \zeta \frac{s/\varepsilon}{s (s/\varepsilon)_0} I_{2\gamma}$$

entropy-to-energy ratio s/ε normalized to $T = 50 \text{ MeV}$ and $2n_{\text{sat}}$ Much less than weak reaction timescale!

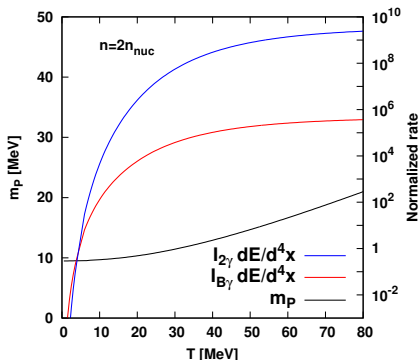
Compare to heat conduction

$$\tau_\kappa = \frac{c_V R^2}{\kappa} = 2 \cdot 10^{-21} \text{ s} \frac{R^2}{l_f^2}$$

$$c_V \simeq \pi^2 \sum_f n_f T / \mu_f$$

$$\text{thermal conductivity } \kappa \simeq \sum_f \mu_f^2 / \alpha_S$$

Heiselberg & Pethick, PRD **48**, 2916 (1993)



Other sources of photons

- Thermal emission: phenomenologically simple background

$$\sigma_{SB} T^4 \simeq 4 \times 10^{21} \frac{\text{MeV}}{\text{fm}^2 \text{s}} \quad (\text{Stefan-Boltzmann estimate})$$

- Bulk viscosity–anomaly

$$\frac{dE}{dVdt} l_f \simeq 4.4 \times 10^{12} (100 \text{ fm}) \frac{\text{MeV}}{\text{fm}^3 \text{s}} \quad \text{at } T \sim 30 \text{ MeV}$$

Emitted from surface layer: $\sim l_f = 100 \text{ fm}$

estimate assumes $\zeta/T^3 \sim O(1)$ – can be enhanced near transition

- Electromagnetic bremsstrahlung, $p + p \rightarrow p + p + \gamma$ (also e^-)

$$\left. \frac{dE}{dVdt} \right|_{\text{brem}} l_f = n_p \frac{2}{3} \alpha \dot{v}^2 \gamma^6 \simeq 7.2 \times 10^9 \frac{\text{MeV}}{\text{fm}^3 \text{s}} (100 \text{ fm})$$

Suppressed by small QED coupling, smaller number of charged particles $n_p \simeq 0.01 n_B$

Conclusions

- Neutron star mergers are violent, involving high densities and moderate temperatures $T \sim 20 - 100$ MeV
- QCD processes very fast \Rightarrow nuclear matter equilibrates, probe new region of phase diagram
- Photon emission due to dissipation of collective modes depends on bulk properties ζ, κ of matter during collision
- Knowing QCD physics even qualitatively provides opportunity as a source of cosmic rays and signatures may be able to identify NS mergers

Extra Slides

Photons in-medium and energy emitted

In hot dense matter of charged particles, photon has self-energy due to fluctuations into fermion pairs. For $|\vec{k}|, \omega \gg T, |\mu_f|$, plasma mass

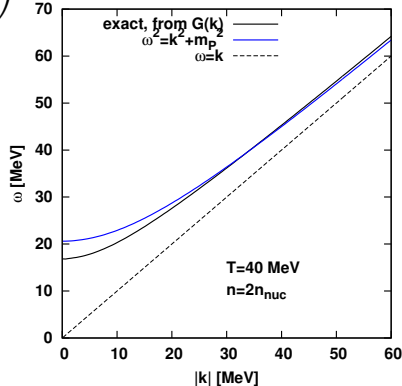
$$m_P^2 = \frac{1}{2} \sum_f (Q_f e)^2 \left(\frac{T^2}{3} + \frac{\mu_f^2}{\pi^2} \right) \simeq (10 - 20 \text{ MeV})^2$$

Mean free path

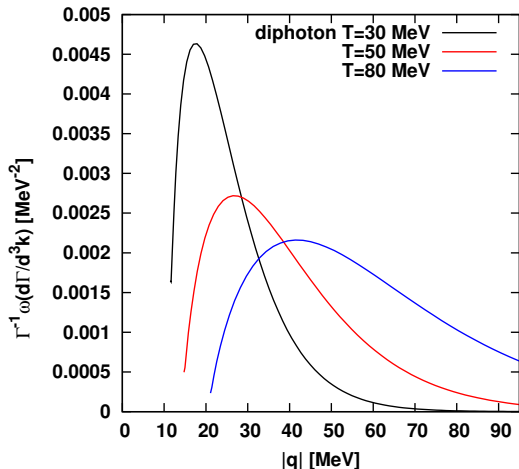
$$l_f \sim (em_P)^{-1} \sim 100 \text{ fm}$$

from $\text{Im} [\text{photon self-energy}]$

\Rightarrow only boundary layer releases energy



Direct photon spectrum



“direct”=no secondary scattering

Differential spectrum
Function of T and density

★ candidate feature(s) to identify NS merger as source

★ source of cosmic rays

to do:

calculate e^+e^- conversion,
plasma scattering

differs from massive particle spectrum due to nontrivial dispersion relation

Opportunities with pions → neutrinos

Nucleon-nucleon scattering into pions

1. Further source of photons:

$N + N \rightarrow N + N + \pi^0$ for $N = n, p$, then decay $\pi \rightarrow 2\gamma$

$$\left. \frac{dE}{dVdt} \right|_{\pi} \uparrow_{f,\pi} = m_{\pi} \sigma n_B^2 v F \simeq 5.3 \times 10^{18} (n_B^{-1/3}) \frac{\text{MeV}}{\text{fm}^3 \text{ s}}$$

$F \sim e^{-m_{\pi}/T}$ = statistical factor for Pauli blocking in degenerate gas

★ large rate due to cross section ($\sigma = \text{a few} \times 10 \mu\text{b}$)

★ also expected to be distinct spectral component due to pion mass

2. Source of prompt neutrinos

$n + n \rightarrow n + p + \pi^-$ $n + p \rightarrow n + n + \pi^+$ $n + p \rightarrow p + p + \pi^-$

charged pions subsequently decay into charged lepton plus neutrino

⇒ if muon, then up to 3 neutrinos per NN collision:

$$3 \left. \frac{dN}{dVdt} \right|_{\pi} I_{f,\pi} = \sigma_{NN\pi} n_B^2 v F \simeq 3.9 \times 10^{16} (n_B^{-1/3}) \frac{1}{\text{fm}^3 \text{ s}}$$

Fuzzy picture of neutron stars

Nuclear “saturation” density

$$n_{\text{sat}} \simeq 0.17 \text{ fm}^{-3}$$

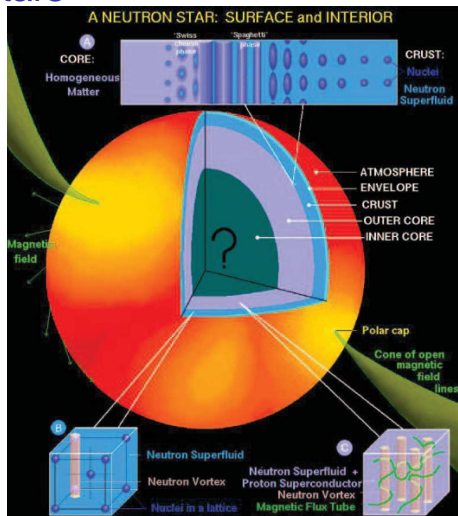
1. crust: mostly neutrons

0.2% – 2% protons believed
in a Wigner-Seitz lattice of
neutron-rich nuclei

neutron superfluid suggested
by glitch phenomenon

2. core: ??

Lattimer & Prakash, Science 304, 536 (2004); PhysRept 442, 109 (2007)



How photons couple, 1: Conformal anomaly

Energy-momentum Trace Θ_{μ}^{μ} breaks Scale invariance
(quantum effect, running of the coupling constant)
 \Rightarrow **anomalous** source of dilatational current

$$\partial_{\mu} S^{\mu} = \Theta_{\mu}^{\mu} = \frac{\beta(\alpha)}{4\alpha} G_{\mu\nu}^a G^{a,\mu\nu} + \sum_f (1 + \gamma_f) m_f \bar{q}_f q_f$$

β =QCD renormalization group function, γ_f =quark anomalous dimension
** also small, **explicit** breaking by quark masses

J. Ellis, Nucl. Phys. B **22**, 478 (1970); M. Chanowitz and J. Ellis, Phys. Lett. B **40**, 397 (1972)

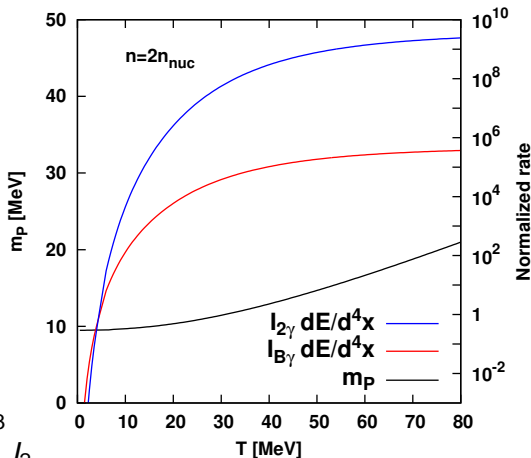
Low-energy effective theory for broken symmetry:
acting on vacuum Θ_{μ}^{μ} operator creates scalar, color singlet σ meson

$$\langle 0 | \Theta_{\mu}^{\mu} | \sigma \rangle = m_{\sigma}^2 f_{\sigma}$$

Ellis & Lanik, Phys. Lett. B **150**, 289 (1985), **175**, 83 (1986)

Energy loss rates

s_0 = entropy of (u, d, e)
 plasma at $T = 50$ MeV and
 $2n_{\text{nuc}}$

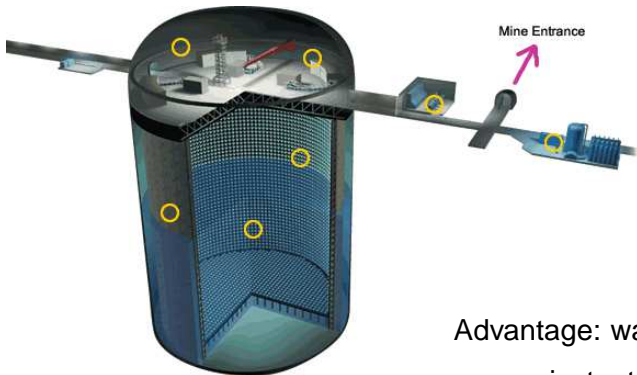


$$\left. \frac{dE}{d^4x} \right|_{2\gamma} \propto 7.13 \frac{\text{TeV}}{\text{fm}^3 \text{s}} \zeta \frac{\text{s}}{s_0} \left(\frac{T}{20 \text{ MeV}} \right)^3 l_{2\gamma}$$

$$\left. \frac{dE}{d^4x} \right|_{B\gamma} \propto 5.51 \frac{\text{GeV}}{\text{fm}^3 \text{s}} \zeta \frac{\text{s}}{s_0} \left(\frac{T}{20 \text{ MeV}} \right)^3 \frac{\vec{B}^2}{B_{\text{crit}}^2} l_{B\gamma}$$

Experimental Input, Now...

Super Kamiokande, ARA (LeCosPA/NTU project),
... IceCube and other neutrino
experiments



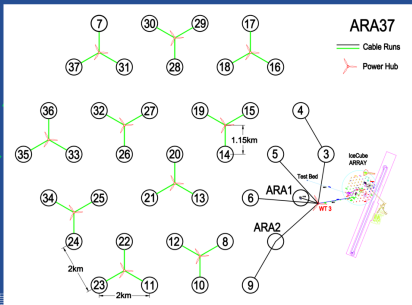
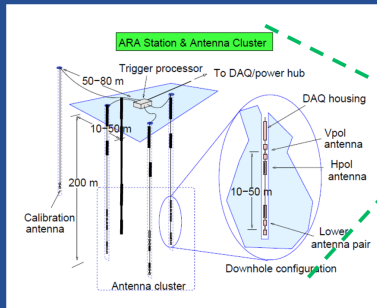
from SuperK's website

Advantage: watching whole sky
means instantaneous detection
and coverage of events

ARA37 (Askaryan Radio Array)

37 4-string, 16-antenna stations covering 200km² with 3-5 v/yr

Taiwan team will contribute 10 stations, or 1/4 of ARA.



Experimental Input,... Coming

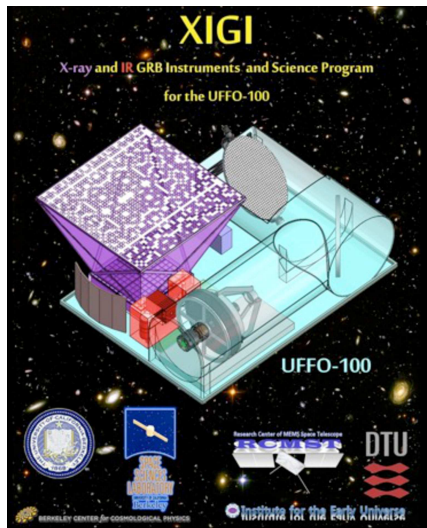
UFFO = Ultra Fast Flash Observatory

- slewing mirror:
millisecond response
to GRB

≥ 50 events/year

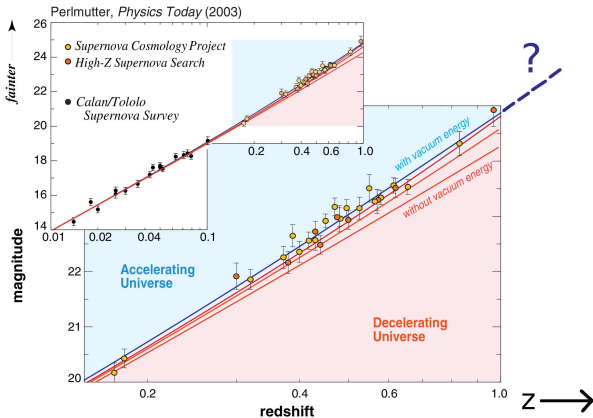
Correlate events with
neutrinos and high
energy cosmic rays

LeCosPA/NTU project!



Brighter standard candles needed!

Type Ia Supernovae



Extend Hubble measurement to $z > 1$, **Probe dark energy**

e.g. possible time dependent Equation of state $w = w_0 + w_1(1 - a) + \dots$