

# Avalanche simulation in noble gas + isobutane mixtures

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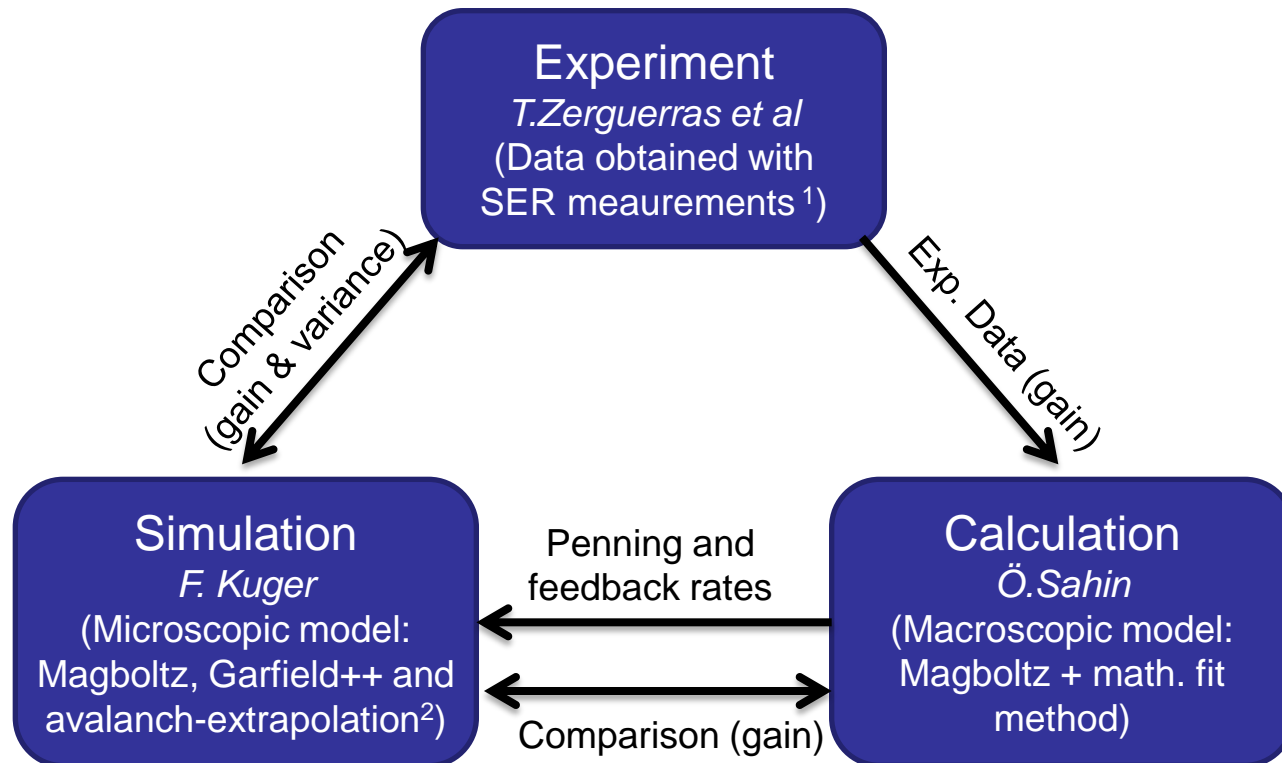
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Studies on penning transfer and feedback in electron avalanche formation, gas gain behavior and gain variances using a comparison between

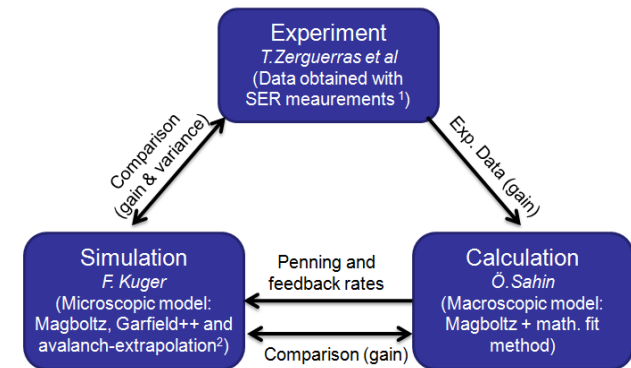


<sup>1</sup>T. Zerguerras et al: *New results on gas gain fluctuations in a Micromegas detector*, MPGD Conference- July 2013 – Zaragoza

<sup>2</sup> F. Kuger et al: *Simulation of gas gain variance in noblegas + isobutane mixtures*, 12<sup>th</sup> RD51 Collaboration Meeting – Oct 2013 – CERN.



As reported during the 12<sup>th</sup> RD51 Collaboration Meeting (Oct 2013) huge differences in gas gain have been observed between the results yielded by the macroscopic and the microscopic model.



- Further investigation on the used methods (e.g. FEM field maps vs. analytic fields, avalanche extrapolation)
- Enhanced understanding of the fundamental differences between microscopic treatment in Garfield++ and macroscopic model calculation
- Better agreement between gain data from simulation and experiment, necessary premise to trust in variance comparison which is not accessible in the calculation method.



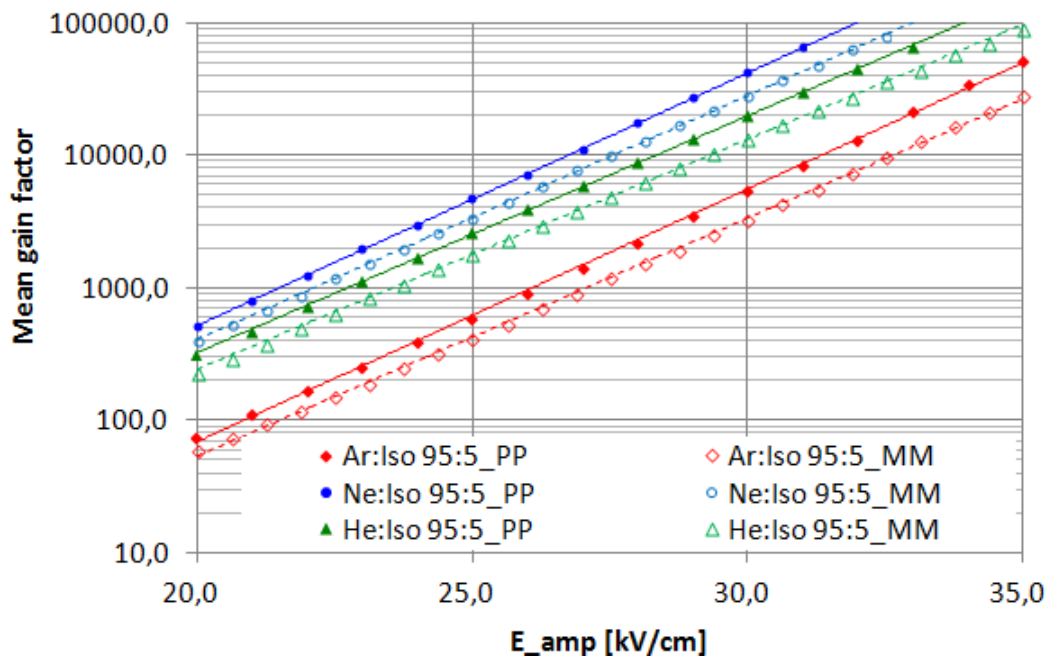
## Simulation

using FEM field maps modeling a Micromegas (MM)

## Calculation

based on a homogeneous field assumption (like a parallel plate setup - PP)

→ Switching to analytic field setup in Garfield, modeling the homogeneous field in a PP



→ PP leads to significant higher gain values than the MM layout.

→ Using PP for further investigation  
+ much faster / less CPU intensive  
+ avoiding introduction of new errors  
- less realistic model

→ For further studies both methods should be adopted to a MM approach



Coping with large avalanches using a mathematical extrapolation.

**Simulation**

The problem:

- ⚡ full avalanche simulation is limited towards high gain due to CPU limits
- ⚡ the larger the avalanche grows, the more statistic becomes dominant
- Full simulation for the first (physical significant) part of the avalanche + mathematical extrapolation yields valid results using less CPU

The model:

- The avalanche development can be divided into independent steps.
- During the first step, a single starting electron is amplified into n electrons, with a probability  $p_1(n)$ , which depends (mostly) on the spatial step length and the electric field.
- Each following step is identical as long as the step length and the electric field are equal.

Mathematical derivation yields:

Mean

RMS

$$\bar{p}_n = \bar{p}_1^n, \quad \text{RMS}(p_n) = \text{RMS}(p_1) \sqrt{\bar{p}_1^{n-1} \sum_{i=0}^{n-1} \bar{p}_1^i}$$



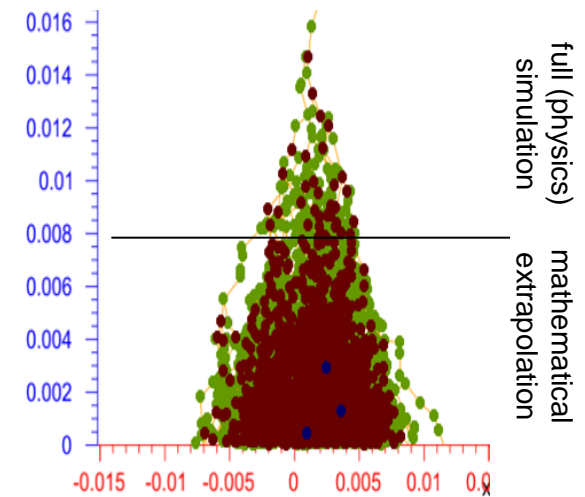
## Steps of the simulation:

- Full simulation of the avalanches (gap = 160μm) in control region (20-28 kV/cm) →  $\bar{g}_{full}$
- Simulation of the first step avalanches (gap = 80μm) in the full region (20-35 kV/cm) →  $\bar{g}_{step}$
- 'Verification' of the step size and the extrapolation exponent  $x$ :

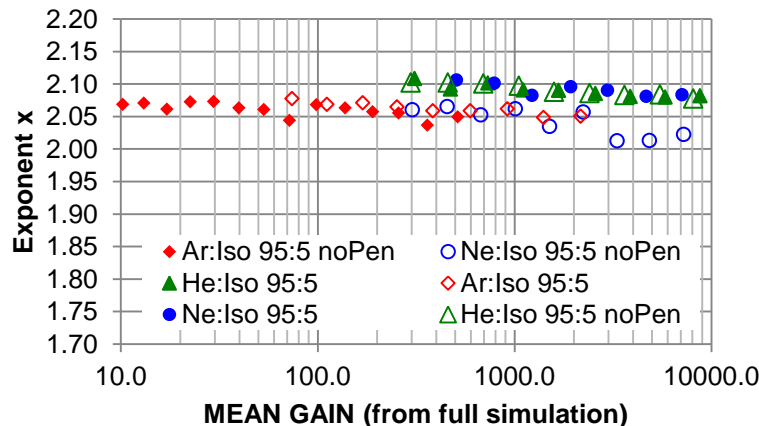
$$x = \frac{\lg \bar{g}_{full}}{\lg \bar{g}_{step}}$$

- Calculation of mean gain and RMS, using  $x$ ,  $\bar{g}_{step}$  and the formula shown above

## Simulation



Calculated Exponents for Extrapolation over MEAN GAIN- half Step



- ⚡ Exponents do not fit the expectation of  $x = 160\mu\text{m} / 80\mu\text{m} = 2$
- ⚡ Systematic effect visible: first step seems to be underestimated /second overestimated.
- ⚡ Effect seems to depend on the gas mixture



Simulation

What is the difference between one full step and two half steps?

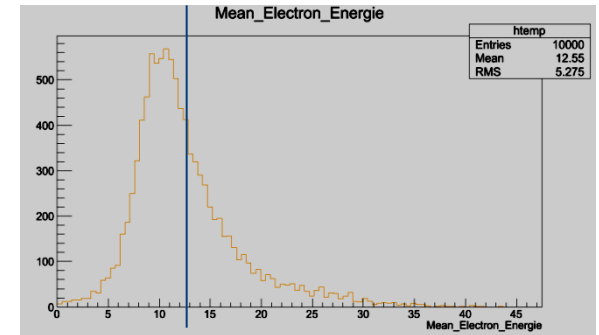
➤ Constraints on the 'half-way-plain', like the starting energy of the electron  $E_0$

*Full simulation:*

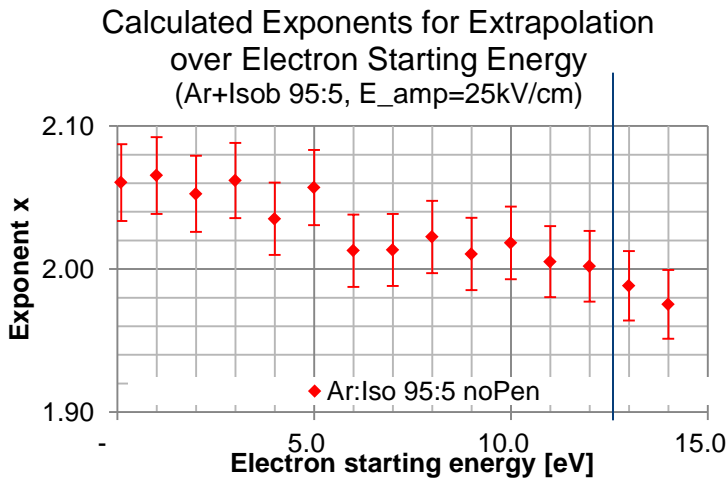
No restrictions, 'random'  
energy after previous steps  
→  $E_0$  is distributed

*Step simulation:*

Reset of each electron  
to starting conditions  
→  $E_0$  is fixed



→ Variation of electron starting energy verifies this assumption:



→ Reason for  $x \neq 2$  is perfectly understood

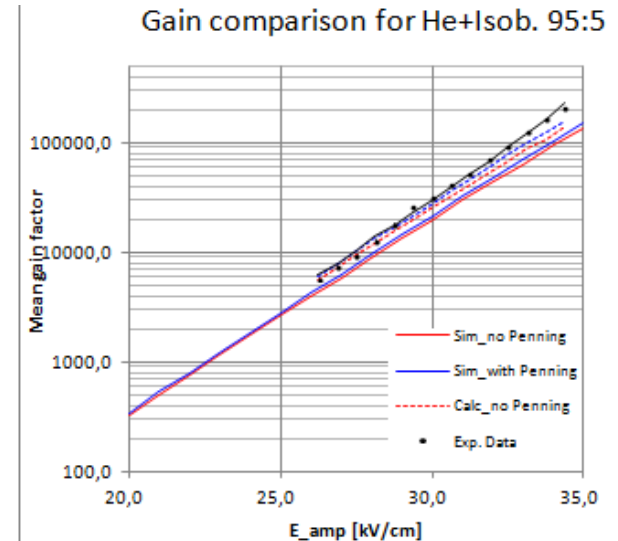
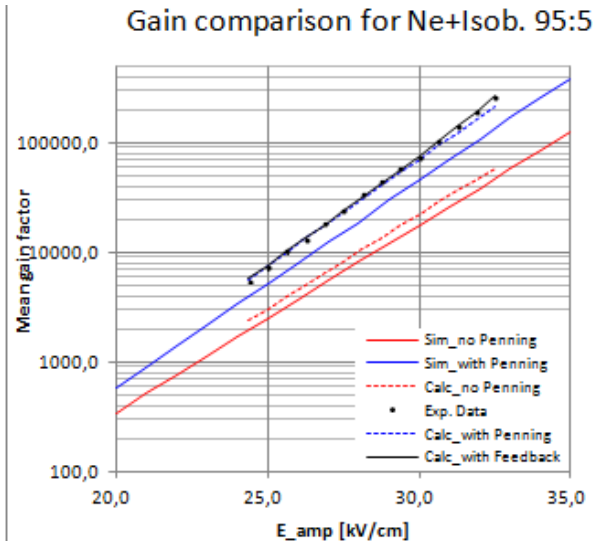
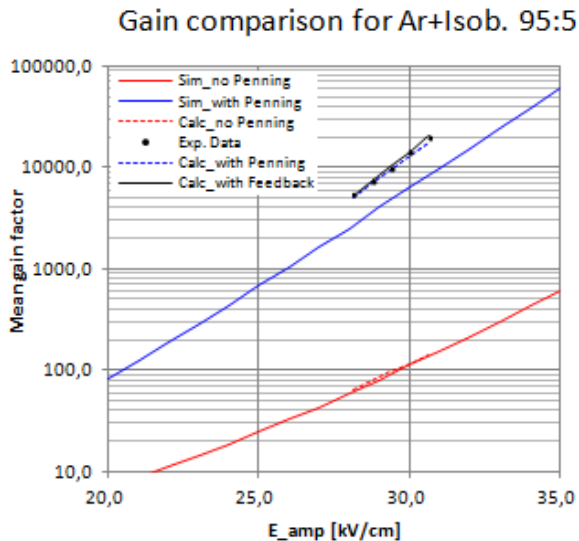
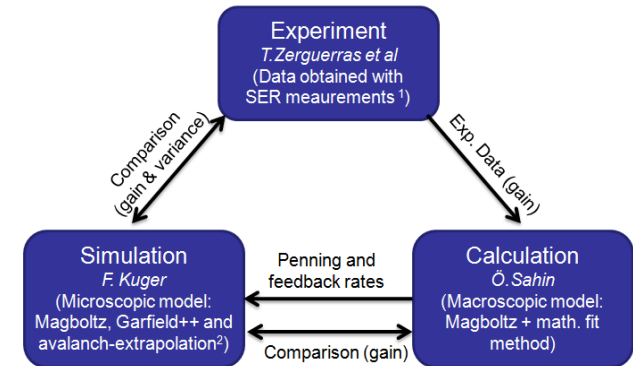
→ Forcing  $E_0$  to obtain  $x = 2$  is not necessary  
and would introduce an error in gain values

→  $E_0$  is set accordingly to simulated energy-  
distribution at the end of a drift process  
(simulation fitting to exp. drift circumstances)



Results comparison after parameter alignment (pressure, temperature, gas input parameter like excitation levels...):

- Still discrepancies in the 'without penning effect – data' (Ar: 2%, Ne: ~20%, He: ~25%)
- Differences in the effect of penning transfer using the same penning transfer rates! (most pronounced in Ar)

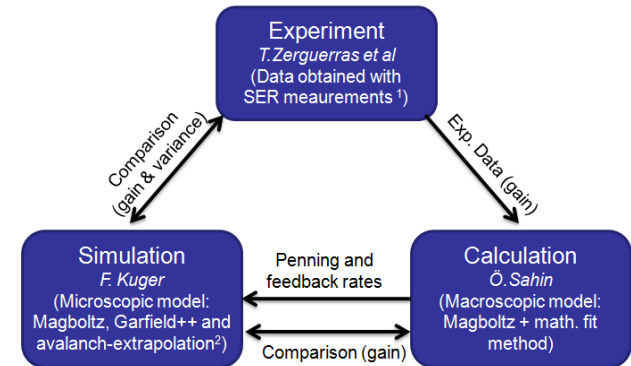




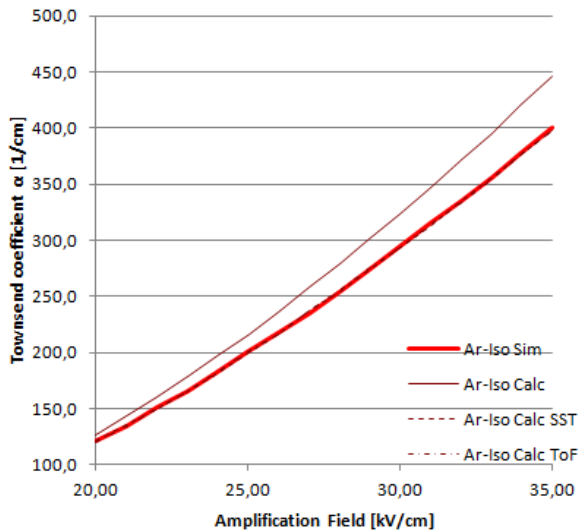


Both methods (Sim and Calc) use the same Magboltz input!  
→ Nevertheless the 'raw' data differs

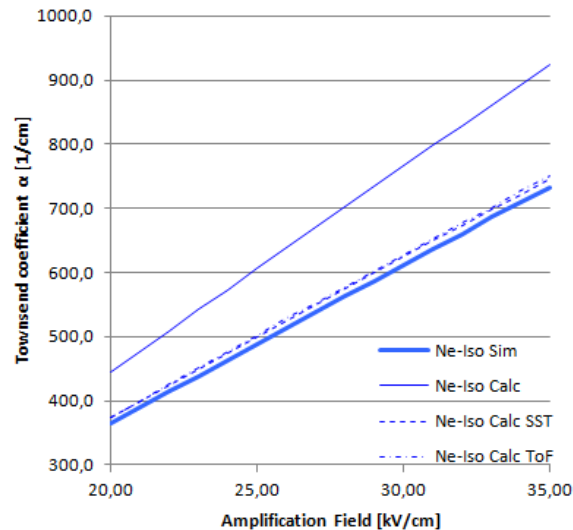
→ Comparison of the Townsend coefficients yield  
(Input parameters in calculation, extracted by magboltz;  
Results derived from microscopic simulation in Garfield++)



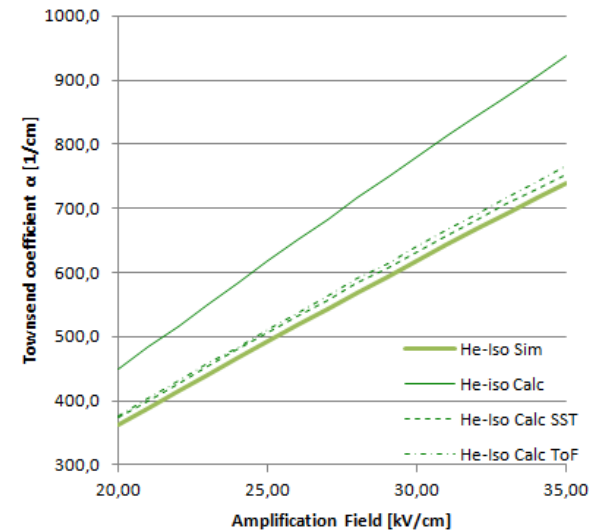
Comparison Townsend coefficients  $\alpha$   
Ar - Iso 95:5



Comparison Townsend coefficients  $\alpha$   
Ne - Isob 95:5



Comparison Townsend coefficients  $\alpha$   
He - Isob 95:5



\* SST = Steady State Townsend, ToF = Time of Flight



**Calculation**

The mean gain depends on the spatial length of the avalanche development

→ Gain is a steady function in d.

$$G_0 = e^{\alpha d}$$

**Simulation**

Gain depends on the number of fully finished amplification processes

→ Gain has the characteristic of a step function in d, with steps every mean free path length  $\lambda$ .

→ Steps are less pronounced if  $d/\lambda \gg 1$ .

Short example assuming fixed interaction length between ionisation and  $d = 3.5 \lambda$

**Calc**

$$G_0 = 2^{3.5} = 11,3$$

**Sim**

$$G_0 = 2^3 = 8 + \text{energy distribution}$$

When calculation Townsend coefficients with Magboltz:  $d/\lambda \gg 1$  is fulfilled.

Running Garfield++ simulations in a 'half step method' (80 $\mu$ m):

$$d/\lambda_{Ar} \sim 19 > d/\lambda_{Ne} > d/\lambda_{He}$$

Which fits to the observed discrepancies. Numerical validation is in progress.



Although a lot of aspects have been under study, no satisfying explanation has been found.

## Calculation

## Simulation

Gain value  
generation

$$G_1 = e^{\alpha_{Penning} d}$$

$$\alpha_{Penning} = \alpha \left( 1 + r_p \frac{f_{exc}}{f_{ion}} \right)$$

Garfield ++ algorithm:

If excited state with energy > lowest ionisation threshold occurs, ionisation takes place with probability  $r_p$ .

Excitation levels &  
Cross sections

Checked for all noble gases! One isobutane state @ 14.0eV not included in Garfield++ Sim – to check if included in Calc. (min. influence)

Electrons energy  
after ionisation

?

Energy difference transferred to electron starting energy

→ Investigation work on this topic is in progress.

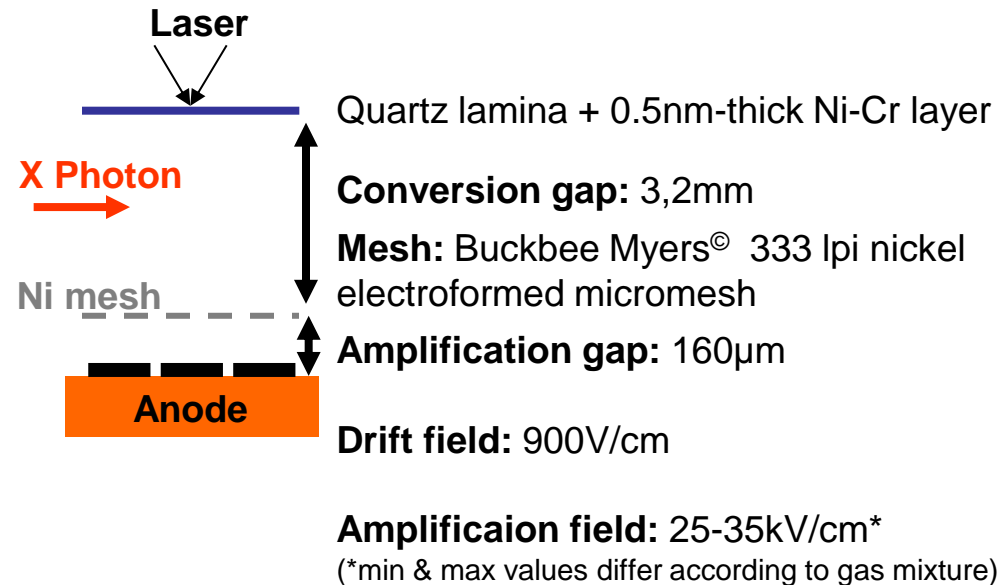
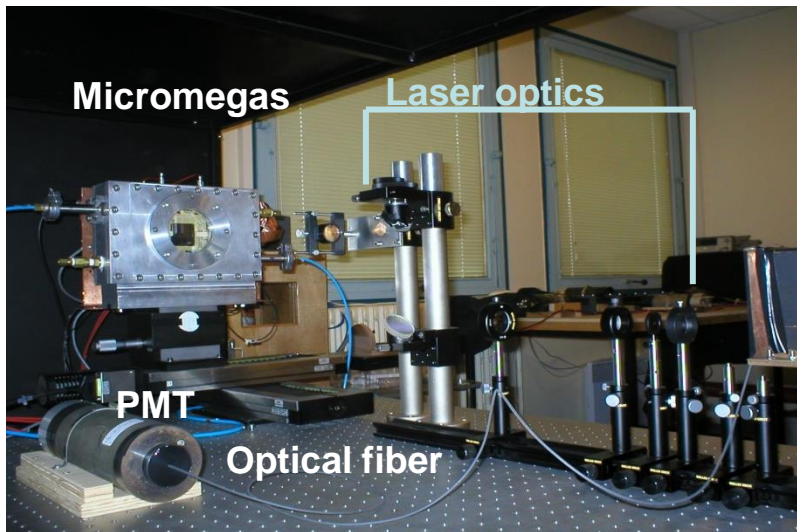
→ Ideas, remarks and contributions are highly welcome!



- The detailed inquiry in the simulation method provided a satisfying level of understanding in the avalanche extrapolation method and revealed some error sources (MM vs. PP, starting energy, measurement conditions...)
- Simulation and Calculation still differ significantly:
  - Difference on 'without-penning-data' is understood but not yet proven
  - Different impact of penning transfer is not yet understood
- Comparison to experimental data is on a satisfying level. (results to be presented soon). Work will go on relying on the microscopic approach.

*Thank you for your attention!*

MPGD characterization with a point-like electron source ( $<100 \mu\text{m}$ ) of variable intensity produced by a 337nm UV laser. Single electrons can be produced at the cathode and their spectra measured with low noise background.



**Gas mixtures:** 95% Ar/He/Ne + 5%  $i\text{C}_4\text{H}_{10}$ , 1atm

« Comparisons with calculations are needed [...] and could help quantifying Penning effect in the three tested mixtures. »<sup>1</sup>

T. Zerguerras – MPGD 2013

<sup>1</sup>T. Zerguerras et al: *New results on gas gain fluctuations in a Micromegas detector*, MPGD Conference- July 2013 – Zaragoza.

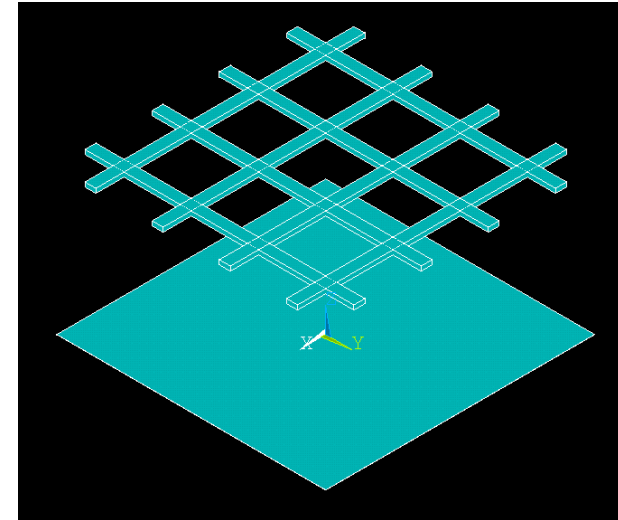
OLD

Calculation of the **electrical field** with *FEM* in **ANSYS 14** using the *smart meshing* with finest granularity.

Layout describes a micromegas with an electroformed mesh (thickness:  $5\mu\text{m}$ , pitch:  $76\mu\text{m}$ , window-size:  $64\mu\text{m}$ ), in fixed  $160\mu\text{m}$  distance to the anode.

Voltages:

$U_{\text{anode}}=0\text{V}$ ,  $U_{\text{mesh}}:-320\text{V to }-560\text{ V}$ ,  $U_{\text{cathode}}=U_{\text{mesh}}-288\text{V}$ ,  
 → Amplification fields: 20-35 kV/cm \*



NEW

Calculation of the electrical field of a **parallel plate** setup.

Using *FEM* in **ANSYS 14** in the first approach, including drift of the electron along  $300\mu\text{m}$  in a fixed driftfield.

The *analytic component* - class in **Garfield++** has been used in the second approach, neglecting the drift region, starting the electron at the upper border of the amplification region.



\* Calculation assumes a homogeneous field ( $E=\Delta U/d$ ).

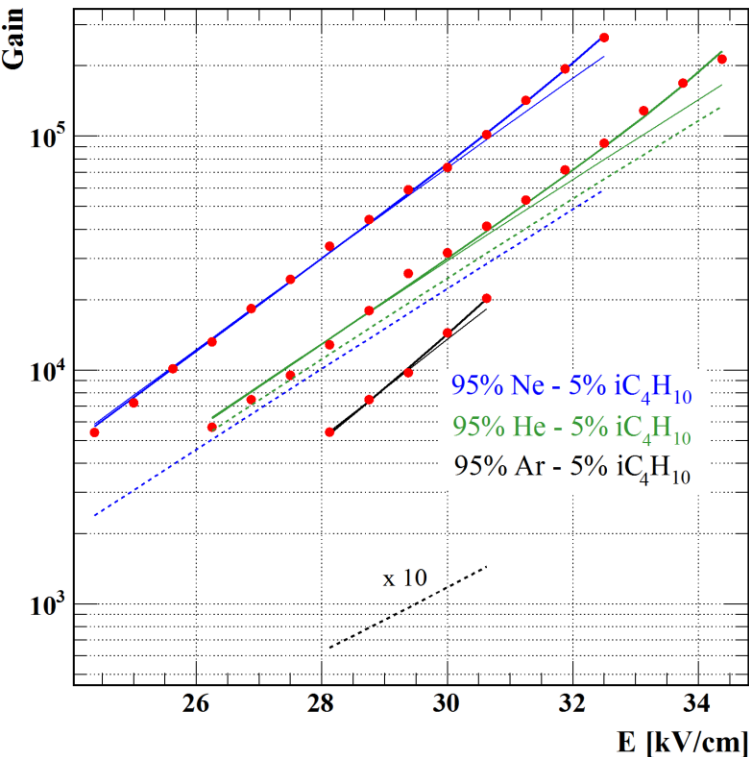
# Macroscopic calculation model (Formulas and fit results)



Gain without penning effect:  $G_0 = e^{\alpha d}$

Gain with penning effect:  $G_1 = e^{\alpha_{Penning} d}$        $\alpha_{Penning} = \alpha \left( 1 + r_p \frac{f_{exc}}{f_{ion}} \right)$

Gain with penning effect and feedback term:  $G_2 = \frac{G_1}{1 - \beta G_1}$



Mixture	$r_p$	$\beta$
He <i>i</i> -C <sub>4</sub> H <sub>10</sub>	0.044 ± 0.012	2.01 ± 0.36 10 <sup>-6</sup>
Ne <i>i</i> -C <sub>4</sub> H <sub>10</sub>	0.3223 ± 0.0054	0.99 ± 0.18 10 <sup>-6</sup>
Ar <i>i</i> -C <sub>4</sub> H <sub>10</sub>	0.2659 ± 0.0022	7.6 ± 3.4 10 <sup>-6</sup>