



Tools for Higgs Physics

Cross Section

ggF

- HIGLU** (NNLO QCD+NLO EW)
- iHixs** (NNLO QCD+NLO EW)
- FeHiPro** (NNLO QCD+NLO EW)
- HNNLO, HRes** (NNLO+NNLL QCD)
- SusHi** (NNLO QCD)
- RGHiggs** (NNLO+NNLL QCD)
- ggHiggs** (approx. NNNLO QCD)

VBF

- VV2H** (NLO QCD)
- VBFNLO** (NLO QCD)
- HAWK** (NLO QCD+EW)
- VBF@NNLO** (NNLO QCD)

WH/ZH

- V2HV** (NLO QCD)
- HAWK** (NLO QCD+EW)
- VH@NNLO** (NNLO)

ttH

- HQQ** (LO QCD)

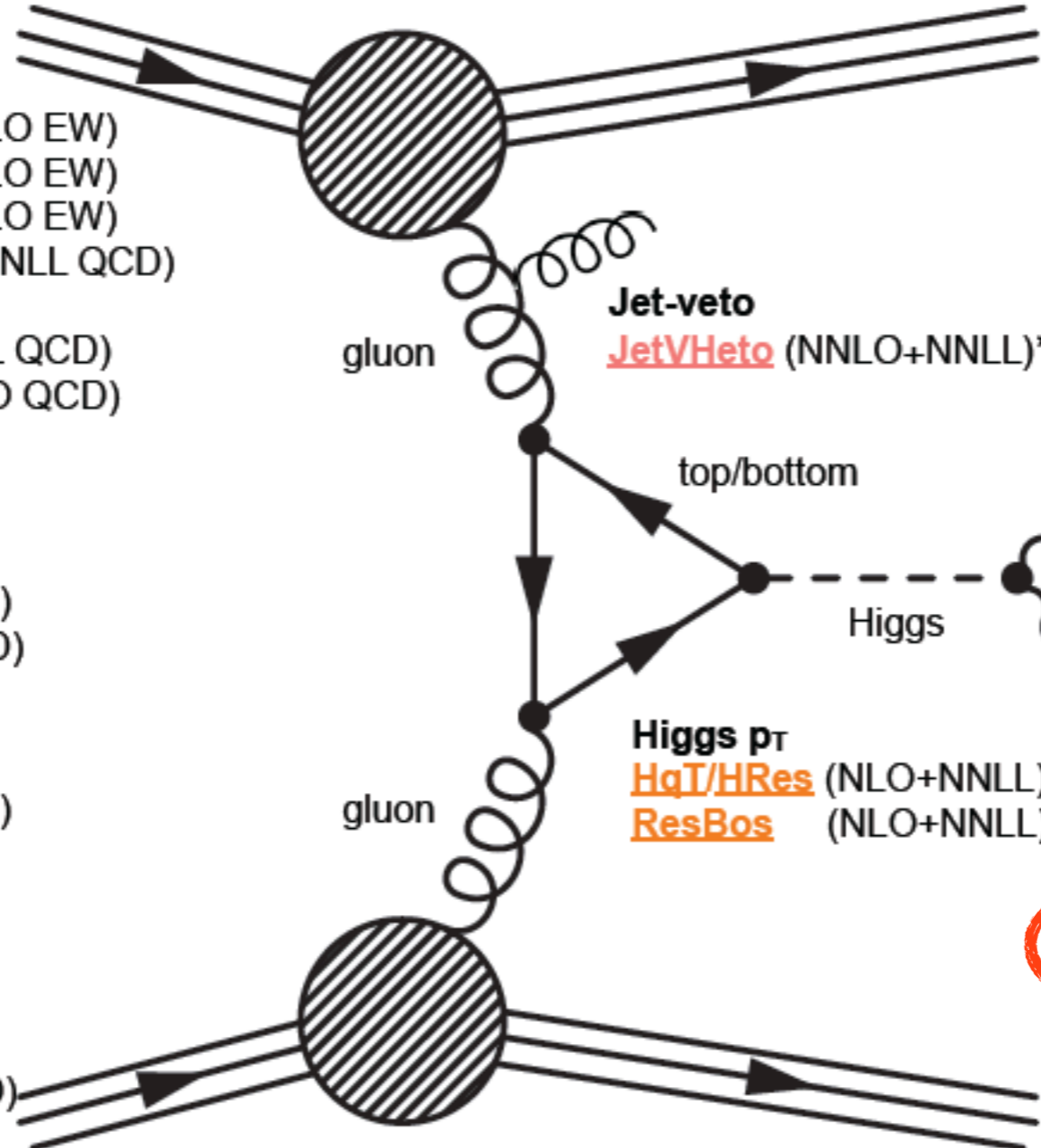
bbH

- bbh@NNLO** (NNLO QCD)

HH

- HPAIR** (NLO QCD)

+ private codes.



PDF: **MSTW, CTEQ, NNPDF, etc.**
LHAPDF, HOPPET, APFEL

NLO MC
POWHEG MiNLO
MadGrapn5 aMC@NLO
SHERPA MEPS@NLO

LO MC
gg2VV

NLO ME
MC2M, MG5 aMC@NLO

W/Z
Higgs Decay
HDECAY (NLO++)
Prophecy4f (NLO)

W/Z

Higgs Properties
MELA/JHU, MEKD
MG5 aMC@NLO (HC)

in this talk

MSSM/2HDM
FeynHiggs, CPSuperH
SusHi+2HDMC
HIGLU+HDECAY

* NLO+NNLL in differential

Higgs Characterisation

via the FeynRules and MadGraph5_aMC@NLO frameworks

Kentarou Mawatari

(Vrije Universities Brussel and International Solvay Institutes)

▶ Sec. II in YR3 of the LHC Higgs Cross Section Working Group [arXiv:1307.1347]

▶ Artoisenet, de Aquino, Demartin, Frederix, Frixione, Maltoni, Mandal, Mathews, KM, Ravindran, Seth, Torrielli, Zaro

“A framework for Higgs characterisation” JHEP11(2013)043 [arXiv:1306.6464]

▶ Maltoni, KM, Zaro

“Higgs characterisation via VBF/VH” EPJC74(2014)2710 [arXiv:1311.1829]

▶ Alloul, Fuks, Sanz

“Phenomenology of the Higgs Effective Lagrangian via FR” JHEP04(2014)110 [arXiv:1310.5150]

Contents

- Introduction
 - Effective field theory
 - Gauge/mass basis implementations
- Higgs characterisation framework (mass basis)
 - Effective Lagrangians -- $X(J=0,1,2)$
 - NLO QCD effects
- Summary

Is this the Standard Model scalar boson?

- ➡ Higgs boson precision measurement
- ➡ determination of **the Higgs boson Lagrangian**
 - **the structure of the operators**, linked to the spin/
parity of a Higgs boson
 - ▶ distributions
 - **the coupling strength**
 - ▶ rate
- ➡ How do we approach to get them?

Effective field theory approach

- Given the fact that only a 125 GeV SM-like boson and nothing else so far, the effective field theory approach is one of the best way to explore BSM effects.
- ▶ All new particles and phenomena are assumed to appear at some scale Λ .
- ▶ Not predictive at scales larger than $\Lambda \rightarrow$ **loss of unitarity**
- ▶ Below Λ , all new physics effects are parametrized by higher dimensional gauge invariant operators made of SM fields. \rightarrow **many parameters 59!**
- ▶ No assumption on the form of new physics \rightarrow **model independent**
- ▶ Renormalisable order by order in the scale $\Lambda \rightarrow$ **systematically improvable**

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots \quad \mathcal{L}_6 = \sum_i C_i Q_i$$

Buchmuller&Wyler 1986 ...
Grzadkowski et al. 2010

SM Lagrangian up to D6

$$\begin{aligned} \mathcal{L}_{\text{SILH}} = & \frac{\tilde{c}_R}{2v^2} \partial^\mu [\Phi^\dagger \Phi] \partial_\mu [\Phi^\dagger \Phi] + \frac{\tilde{c}_T}{2v^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] - \frac{\tilde{c}_\lambda}{v^2} [H^\dagger H]^3 \\ & - \left[\frac{\tilde{c}_u}{v^2} y_u \Phi^\dagger \Phi \Phi^\dagger \cdot Q_L u_R + \frac{\tilde{c}_d}{v^2} y_d \Phi^\dagger \Phi \Phi^\dagger Q_L d_R + \frac{\tilde{c}_e}{v^2} y_e \Phi^\dagger \Phi \Phi^\dagger L_L e_R + \text{h.c.} \right] \\ & + \frac{ig}{m_W^2} \tilde{c}_{TW} [\Phi^\dagger T_{2k} \overleftrightarrow{D}^\mu \Phi] D^\nu W_{\mu\nu}^k + \frac{ig'}{2m_W^2} \tilde{c}_B [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] \partial^\nu B_{\mu\nu} \\ & + \frac{2ig}{m_W^2} \tilde{c}_{HW} [D^\mu \Phi^\dagger T_{2k} D^\nu \Phi] W_{\mu\nu}^k + \frac{ig'}{m_W^2} \tilde{c}_{HB} [D^\mu \Phi^\dagger D^\nu \Phi] B_{\mu\nu} \\ & + \frac{g^2}{m_W^2} \tilde{c}_1 \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g_s^2}{m_W^2} \tilde{c}_2 \Phi^\dagger \Phi G_{\mu\nu}^a G_a^{\mu\nu}, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{F_1} = & \frac{i\tilde{c}_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu Q_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\tilde{c}'_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu T_{2k} Q_L] [\Phi^\dagger T_2^k \overleftrightarrow{D}_\mu \Phi] \\ & + \frac{i\tilde{c}_{Hu}}{v^2} [\bar{u}_R \gamma^\mu u_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{i\tilde{c}_{Hd}}{v^2} [\bar{d}_R \gamma^\mu d_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] \\ & - \left[\frac{i\tilde{c}_{He}}{v^2} [\bar{e}_R \gamma^\mu e_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \text{h.c.} \right] \\ & + \frac{i\tilde{c}_{HL}}{v^2} [\bar{L}_L \gamma^\mu L_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\tilde{c}'_{HL}}{v^2} [\bar{L}_L \gamma^\mu T_{2k} L_L] [\Phi^\dagger T_2^k \overleftrightarrow{D}_\mu \Phi] \\ & + \frac{i\tilde{c}_{He}}{v^2} [\bar{e}_R \gamma^\mu e_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi], \end{aligned}$$

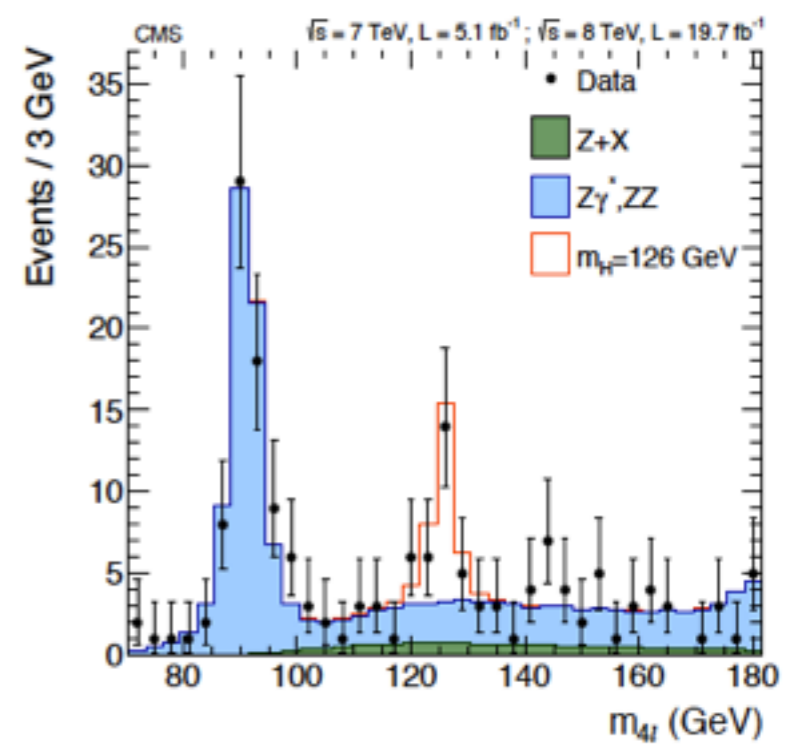
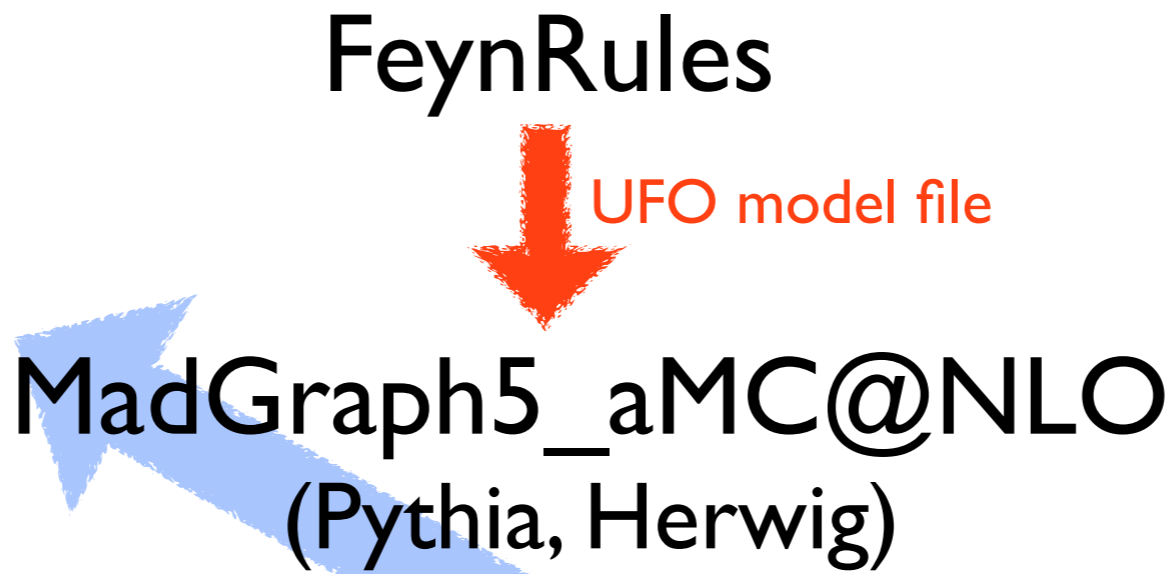
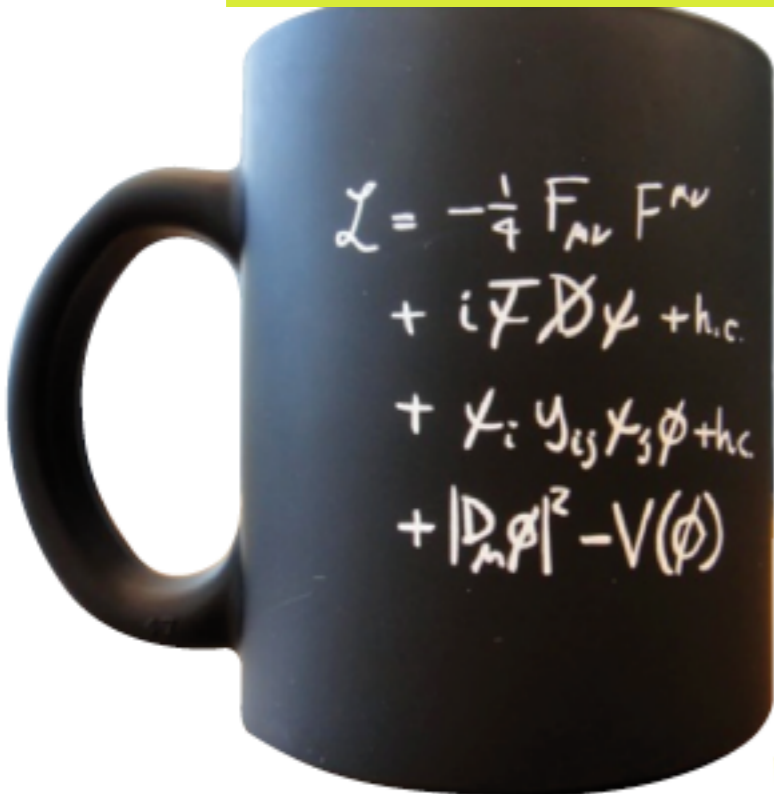
$$\begin{aligned} \mathcal{L}_{CP} = & \frac{ig}{m_W^2} \tilde{c}_{HW} D^\mu \Phi^\dagger T_{2k} D^\nu \Phi \tilde{W}_{\mu\nu}^k + \frac{ig'}{m_W^2} \tilde{c}_{HB} D^\mu \Phi^\dagger D^\nu \Phi \tilde{B}_{\mu\nu} + \frac{g^2}{m_W^2} \tilde{c}_1 \Phi^\dagger \Phi B_{\mu\nu} \tilde{B}^{\mu\nu} \\ & + \frac{g_s^2}{m_W^2} \tilde{c}_2 \Phi^\dagger \Phi G_{\mu\nu}^a \tilde{G}_a^{\mu\nu} + \frac{g^3}{m_W^2} \tilde{c}_{3W} \epsilon_{ijk} W_{\mu\nu}^i W_{\rho\sigma}^j \tilde{W}^{\rho\sigma k} + \frac{g_s^3}{m_W^2} \tilde{c}_{3G} f_{abc} G_{\mu\nu}^a G^{\nu\rho b} \tilde{G}^{\rho\sigma c} \end{aligned}$$

$$\begin{aligned} \mathcal{L}_G = & \frac{g^3}{m_W^2} \tilde{c}_{3W} \epsilon_{ijk} W_{\mu\nu}^i W^{\nu\rho j} W^{\rho\sigma k} + \frac{g_s^3}{m_W^2} \tilde{c}_{3G} f_{abc} G_{\mu\nu}^a G^{\nu\rho b} G^{\rho\sigma c} + \frac{\tilde{c}_{2W}}{m_W^2} D^\mu W_{\mu\nu}^k D_\rho W_k^{\rho\nu} \\ & + \frac{\tilde{c}_{2B}}{m_W^2} \partial^\mu B_{\mu\nu} \partial_\rho B^{\rho\nu} + \frac{\tilde{c}_{2G}}{m_W^2} D^\mu G_{\mu\nu}^a D_\rho G_a^{\rho\nu}, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{F_2} = & \left[-\frac{2g'}{m_W^2} \tilde{c}_{uB} y_u \Phi^\dagger \cdot \bar{Q}_L \gamma^{\mu\nu} u_R B_{\mu\nu} - \frac{4g}{m_W^2} \tilde{c}_{uW} y_u \Phi^\dagger \cdot (\bar{Q}_L T_{2k}) \gamma^{\mu\nu} u_R W_{\mu\nu}^k \right. \\ & - \frac{4g_s}{m_W^2} \tilde{c}_{uG} y_u \Phi^\dagger \cdot \bar{Q}_L \gamma^{\mu\nu} T_a u_R G_{\mu\nu}^a + \frac{2g'}{m_W^2} \tilde{c}_{dB} y_d \Phi^\dagger \bar{Q}_L \gamma^{\mu\nu} d_R B_{\mu\nu} \\ & + \frac{4g}{m_W^2} \tilde{c}_{dW} y_d \Phi^\dagger (\bar{Q}_L T_{2k}) \gamma^{\mu\nu} d_R W_{\mu\nu}^k + \frac{4g_s}{m_W^2} \tilde{c}_{dG} y_d \Phi^\dagger \bar{Q}_L \gamma^{\mu\nu} T_a d_R G_{\mu\nu}^a \\ & \left. + \frac{2g'}{m_W^2} \tilde{c}_{eB} y_e \Phi^\dagger \bar{L}_L \gamma^{\mu\nu} e_R B_{\mu\nu} + \frac{4g}{m_W^2} \tilde{c}_{eW} y_e \Phi^\dagger (\bar{L}_L T_{2k}) \gamma^{\mu\nu} e_R W_{\mu\nu}^k + \text{h.c.} \right] \end{aligned}$$

from Contino, Ghezzi, Grojean, Muhlleitner, Spira [arXiv: 1303.3876]

Lagrangian (TH) \Leftrightarrow Data (EXP)



Higgs effective Lagrangian before vs. after EW symmetry breaking

- D6 (the gauge basis): HEL [Alloul, Fuks, Sanz, arXiv:1310.5150]
 - ▶ Only using Standard Model gauge-eigenstates
 - ▶ Several operators may be associated with a single coupling (in the mass basis)
 - ▶ One operator associated with several couplings (in the mass basis)
 - ▶ <https://feynrules.irmp.ucl.ac.be/wiki/HEL>
- D5 (the mass basis): HC [Artoisenet et al., arXiv:1306.6464]
 - ▶ Couplings of the physical Higgs boson to the Standard Model (physical) states
 - ▶ One operator associated with a single coupling (and Lorentz structure)
 - ▶ No assumption on the Higgs boson spin
 - ▶ <https://feynrules.irmp.ucl.ac.be/wiki/HiggsCharacterisation>

D6 Higgs Effective Lagrangian

[from Contino, Ghezzi, Grojean, Muhlleitner, Spira (JHEP '13)]

[Alloul, Fuks, Sanz (1310.5150)]

$$\begin{aligned} \mathcal{L}_{\text{SILH}} = & \frac{\bar{c}_H}{2v^2} \partial^\mu [\Phi^\dagger \Phi] \partial_\mu [\Phi^\dagger \Phi] + \frac{\bar{c}_T}{2v^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] - \frac{\bar{c}_\lambda}{v^2} [H^\dagger H]^3 \\ & - \left[\frac{\bar{c}_u}{v^2} y_u \Phi^\dagger \Phi \Phi^\dagger \cdot \bar{Q}_L u_R + \frac{\bar{c}_d}{v^2} y_d \Phi^\dagger \Phi \Phi^\dagger \bar{Q}_L d_R + \frac{\bar{c}_l}{v^2} y_l \Phi^\dagger \Phi \Phi^\dagger \bar{L}_L e_R + \text{h.c.} \right] \\ & + \frac{ig}{m_W^2} \bar{c}_W [\Phi^\dagger T_{2k} \overleftrightarrow{D}^\mu \Phi] D^\nu W_{\mu\nu}^k + \frac{ig'}{2m_W^2} \bar{c}_B [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] \partial^\nu B_{\mu\nu} \\ & + \frac{2ig}{m_W^2} \bar{c}_{HW} [D^\mu \Phi^\dagger T_{2k} D^\nu \Phi] W_{\mu\nu}^k + \frac{ig'}{m_W^2} \bar{c}_{HB} [D^\mu \Phi^\dagger D^\nu \Phi] B_{\mu\nu} \\ & + \frac{\bar{g}'^2}{m_W^2} c_\gamma \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{\bar{g}_s^2}{m_W^2} c_g \Phi^\dagger \Phi G_{\mu\nu}^a G_a^{\mu\nu}, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{CP} = & \frac{ig}{m_W^2} \bar{c}_{HW} D^\mu \Phi^\dagger T_{2k} D^\nu \Phi \tilde{W}_{\mu\nu}^k + \frac{ig'}{m_W^2} \bar{c}_{HB} D^\mu \Phi^\dagger D^\nu \Phi \tilde{B}_{\mu\nu} + \frac{g'^2}{m_W^2} \bar{c}_\gamma \Phi^\dagger \Phi B_{\mu\nu} \tilde{B}^{\mu\nu} \\ & + \frac{g_s^2}{m_W^2} \bar{c}_g \Phi^\dagger \Phi G_{\mu\nu}^a \tilde{G}_a^{\mu\nu} + \frac{g^3}{m_W^2} \bar{c}_{3W} \epsilon_{ijk} W_{\mu\nu}^i W_{\nu\rho}^j \tilde{W}^{\rho\mu k} + \frac{g_s^3}{m_W^2} \bar{c}_{3G} f_{abc} G_{\mu\nu}^a G_{\nu\rho}^b \tilde{G}^{\rho\mu c} \end{aligned}$$

$$\begin{aligned} \mathcal{L}_G = & \frac{g^3}{m_W^2} \bar{c}_{3W} \epsilon_{ijk} W_{\mu\nu}^i W_{\nu\rho}^j W^{\rho\mu k} + \frac{g_s^3}{m_W^2} \bar{c}_{3G} f_{abc} G_{\mu\nu}^a G_{\nu\rho}^b G^{\rho\mu c} + \frac{\bar{c}_{2W}}{m_W^2} D^\mu W_{\mu\nu}^k D_\rho W_k^{\rho\nu} \\ & + \frac{\bar{c}_{2B}}{m_W^2} \partial^\mu B_{\mu\nu} \partial_\rho B^{\rho\nu} + \frac{\bar{c}_{2G}}{m_W^2} D^\mu G_{\mu\nu}^a D_\rho G_a^{\rho\nu}, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{F_1} = & \frac{i\bar{c}_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu Q_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\bar{c}'_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu T_{2k} Q_L] [\Phi^\dagger T_2^k \overleftrightarrow{D}_\mu \Phi] \\ & + \frac{i\bar{c}_{Hu}}{v^2} [\bar{u}_R \gamma^\mu u_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{i\bar{c}_{Hd}}{v^2} [\bar{d}_R \gamma^\mu d_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] \\ & - \left[\frac{i\bar{c}_{Hud}}{v^2} [\bar{u}_R \gamma^\mu d_R] [\Phi \cdot \overleftrightarrow{D}_\mu \Phi] + \text{h.c.} \right] \\ & + \frac{i\bar{c}_{HL}}{v^2} [\bar{L}_L \gamma^\mu L_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\bar{c}'_{HL}}{v^2} [\bar{L}_L \gamma^\mu T_{2k} L_L] [\Phi^\dagger T_2^k \overleftrightarrow{D}_\mu \Phi] \\ & + \frac{i\bar{c}_{He}}{v^2} [\bar{e}_R \gamma^\mu e_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi], \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{F_2} = & \left[-\frac{2g'}{m_W^2} \bar{c}_{uB} y_u \Phi^\dagger \cdot \bar{Q}_L \gamma^{\mu\nu} u_R B_{\mu\nu} - \frac{4g}{m_W^2} \bar{c}_{uW} y_u \Phi^\dagger \cdot (\bar{Q}_L T_{2k}) \gamma^{\mu\nu} u_R W_{\mu\nu}^k \right. \\ & - \frac{4g_s}{m_W^2} \bar{c}_{uG} y_u \Phi^\dagger \cdot \bar{Q}_L \gamma^{\mu\nu} T_a u_R G_{\mu\nu}^a + \frac{2g'}{m_W^2} \bar{c}_{dB} y_d \Phi \bar{Q}_L \gamma^{\mu\nu} d_R B_{\mu\nu} \\ & + \frac{4g}{m_W^2} \bar{c}_{dW} y_d \Phi (\bar{Q}_L T_{2k}) \gamma^{\mu\nu} d_R W_{\mu\nu}^k + \frac{4g_s}{m_W^2} \bar{c}_{dG} y_d \Phi \bar{Q}_L \gamma^{\mu\nu} T_a d_R G_{\mu\nu}^a \\ & \left. + \frac{2g'}{m_W^2} \bar{c}_{eB} y_l \Phi \bar{L}_L \gamma^{\mu\nu} e_R B_{\mu\nu} + \frac{4g}{m_W^2} \bar{c}_{eW} y_l \Phi (\bar{L}_L T_{2k}) \gamma^{\mu\nu} e_R W_{\mu\nu}^k + \text{h.c.} \right] \end{aligned}$$

◆ The model file is publicly available. (<https://feynrules.irmp.ucl.ac.be/wiki/HEL>)

Mapping between the D6 and D5 operators

HC [arXiv: 1306.6464]

HEL [arXiv: 1310.5150]

$$\mathcal{L}_0^f = - \sum_{f=t,b,\tau} \bar{\psi}_f (c_\alpha \kappa_{Hff} g_{Hff} + i s_\alpha \kappa_{Aff} g_{Aff} \gamma_5) \psi_f X_0$$

$$\mathcal{L}_0^V = \left\{ c_\alpha \kappa_{SM} \left[\frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \right.$$

$$- \frac{1}{4} \left[c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma} A_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma} A_{\mu\nu} \tilde{A}^{\mu\nu} \right]$$

$$- \frac{1}{2} \left[c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma} Z_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{AZ\gamma} g_{AZ\gamma} Z_{\mu\nu} \tilde{A}^{\mu\nu} \right]$$

$$- \frac{1}{4} \left[c_\alpha \kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + s_\alpha \kappa_{Agg} g_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right]$$

$$- \frac{1}{4\Lambda} \left[c_\alpha \kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + s_\alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right]$$

$$- \frac{1}{2\Lambda} \left[c_\alpha \kappa_{HWW} W_{\mu\nu}^+ W^{-\mu\nu} + s_\alpha \kappa_{AWW} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} \right]$$

$$- \frac{1}{\Lambda} c_\alpha \left[\kappa_{H\theta\gamma} Z_\nu \partial_\mu \right.$$

$$\left. + (\kappa_{H\theta W} W_\nu \partial_\mu - \kappa_{H\theta W} W_\mu \partial_\nu) \right]$$

Eq. (2.25)	Ref. [46]	Section 2.1
g_{hgg}	$c_\alpha \kappa_{Hgg} g_{Hgg}$	$g_H - \frac{4\bar{c}_g g_s^2 v}{m_W^2}$
\tilde{g}_{hgg}	$s_\alpha \kappa_{Agg} g_{Agg}$	$-\frac{4\bar{c}_g g_s^2 v}{m_W^2}$
$g_{h\gamma\gamma}$	$c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma}$	$a_H - \frac{8g\bar{c}_\gamma s_W^2}{m_W}$
$\tilde{g}_{h\gamma\gamma}$	$s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma}$	$-\frac{8g\bar{c}_\gamma s_W^2}{m_W}$
$g_{hzz}^{(1)}$	$\frac{1}{\Lambda} c_\alpha \kappa_{HZZ}$	$\frac{2g}{c_W^2 m_W} [\bar{c}_{HB} s_W^2 - 4\bar{c}_\gamma s_W^4 + c_W^2 \bar{c}_{HW}]$
\tilde{g}_{hzz}	$\frac{1}{\Lambda} s_\alpha \kappa_{AZZ}$	$\frac{2g}{c_W^2 m_W} [\bar{c}_{HB} s_W^2 - 4\bar{c}_\gamma s_W^4 + c_W^2 \bar{c}_{HW}]$
$g_{hzz}^{(2)}$	$\frac{1}{\Lambda} c_\alpha \kappa_{H\theta Z}$	$\frac{g}{c_W^2 m_W} [(\bar{c}_{HW} + \bar{c}_W) c_W^2 + (\bar{c}_B + \bar{c}_{HB}) s_W^2]$
$g_{hzz}^{(3)}$	$c_\alpha \kappa_{SM} g_{HZZ}$	$\frac{gm_W}{c_W^2} \left[1 - \frac{1}{2}\bar{c}_H - 2\bar{c}_T + 8\bar{c}_\gamma \frac{s_W^4}{c_W^2} \right]$
$g_{haz}^{(1)}$	$c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma}$	$\frac{gs_W}{c_W m_W} [\bar{c}_{HW} - \bar{c}_{HB} + 8\bar{c}_\gamma s_W^2]$
\tilde{g}_{haz}	$s_\alpha \kappa_{AZ\gamma} g_{AZ\gamma}$	$\frac{gs_W}{c_W m_W} [\bar{c}_{HW} - \bar{c}_{HB} + 8\bar{c}_\gamma s_W^2]$
g_{hzw}	$c_\alpha \kappa_{H\theta W}$	$\frac{gs_W}{m_W} [\bar{c}_{HW} - \bar{c}_{HB} - \bar{c}_B + \bar{c}_W]$
$g_{hww}^{(2)}$	$\frac{1}{\Lambda} c_\alpha \kappa_{H\theta W}$	$\frac{g}{m_W} [\bar{c}_W + \bar{c}_{HW}]$

Two approaches equivalent as they can be mapped into one another.

$$V_{\mu\nu} = \partial_\mu V_\nu - \partial_\nu V_\mu \quad (V = A, Z, W^\pm), \quad V_{\mu\nu} = \frac{1}{2} \epsilon_{\mu\nu\rho\sigma} V^{\rho\sigma}$$

$$G_{\mu\nu}^a = \partial_\mu G_\nu^a - \partial_\nu G_\mu^a + g_s f^{abc} G_\mu^b G_\nu^c,$$

Higgs Characterisation model

- We implemented an effective Lagrangian featuring bosons $X(J^P=0^+,0^-,1^+,1^-,2^+)$ in FeynRules.
 - ▶ **Effective field theory** approach, valid up to a cutoff scale Λ
 - ▶ Only **one new bosonic state** $X(J^P)$ at the EW scale (No other state below the cutoff Λ)
 - ▶ Any new physics is described by the lowest dimensional operators.

The parametrization is based on the recent work [Englert, Goncalves-Netto, KM, Plehn (2013)].

Effective Lagrangian -- spin0

$$\mathcal{L}_0^f = - \sum_{f=t,b,\tau} \bar{\psi}_f (c_\alpha \kappa_{Hff} g_{Hff} + i s_\alpha \kappa_{Aff} g_{Aff} \gamma_5) \psi_f X_0$$

$$\mathcal{L}_0^V = \left\{ c_\alpha \kappa_{SM} \left[\frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \right. \\ - \frac{1}{4} \left[c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma} A_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma} A_{\mu\nu} \tilde{A}^{\mu\nu} \right] \\ - \frac{1}{2} \left[c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma} Z_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{AZ\gamma} g_{AZ\gamma} Z_{\mu\nu} \tilde{A}^{\mu\nu} \right] \\ - \frac{1}{4} \left[c_\alpha \kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + s_\alpha \kappa_{Agg} g_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right] \\ - \frac{1}{4} \frac{1}{\Lambda} \left[c_\alpha \kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + s_\alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] \\ - \frac{1}{2} \frac{1}{\Lambda} \left[c_\alpha \kappa_{HWW} W_{\mu\nu}^+ W^{-\mu\nu} + s_\alpha \kappa_{AWW} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} \right] \\ - \frac{1}{\Lambda} c_\alpha \left[\kappa_{H\partial\gamma} Z_\nu \partial_\mu A^{\mu\nu} + \kappa_{H\partial Z} Z_\nu \partial_\mu Z^{\mu\nu} \right. \\ \left. + (\kappa_{H\partial W} W_\nu^+ \partial_\mu W^{-\mu\nu} + h.c.) \right] \Big\} X_0$$

parameter	description
Λ [GeV]	cutoff scale
c_α ($\equiv \cos \alpha$)	mixing between 0^+ and 0^-
κ_i	dimensionless coupling parameter

```
#####
## INFORMATION FOR FRBLOCK
#####
Block frblock
  1 1.000000e+03 # Lambda
  2 1.000000e+00 # ca
  3 1.000000e+00 # kSM
  4 1.000000e+00 # kHtt
  5 1.000000e+00 # kAtt
  6 1.000000e+00 # kHbb
  7 1.000000e+00 # kAbb
  8 1.000000e+00 # kHll
  9 1.000000e+00 # kAll
 10 1.000000e+00 # kHaa
 11 1.000000e+00 # kAaa
 12 1.000000e+00 # kHza
 13 1.000000e+00 # kAza
 14 1.000000e+00 # kHgg
 15 1.000000e+00 # kAgg
 16 0.000000e+00 # kHzz
 17 0.000000e+00 # kAzz
 18 0.000000e+00 # kHww
 19 0.000000e+00 # kAww
 20 0.000000e+00 # kHda
 21 0.000000e+00 # kHdz
 22 0.000000e+00 # kHdwR
 23 0.000000e+00 # kHdwI
```

Effective Lagrangian -- spin0

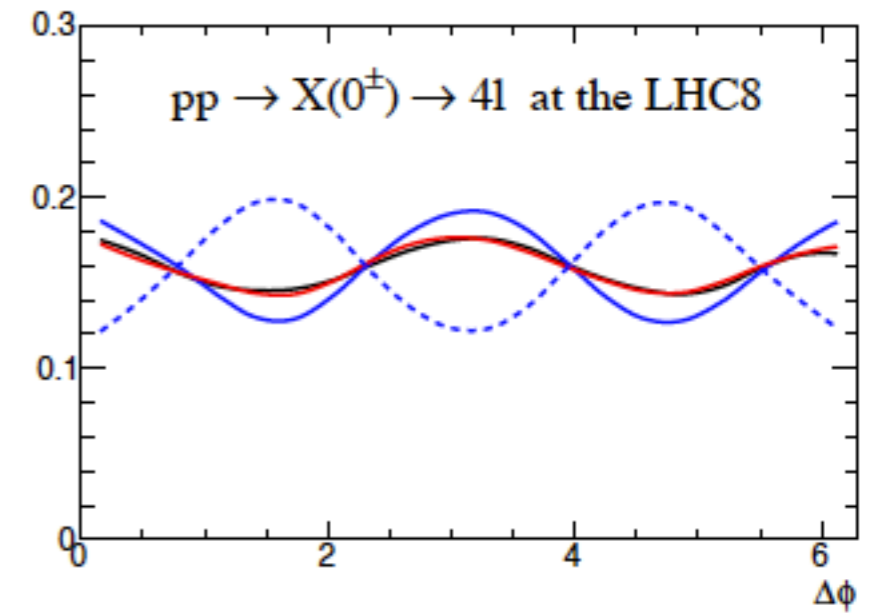
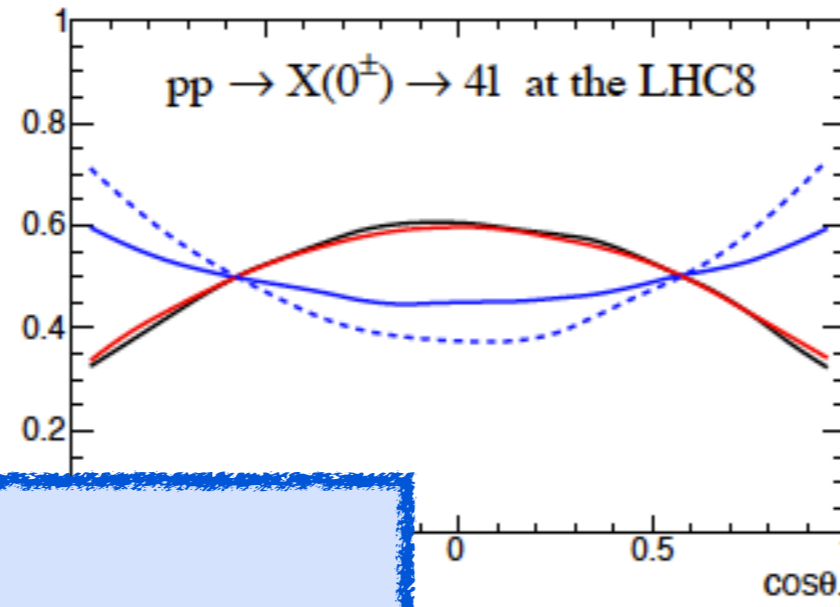
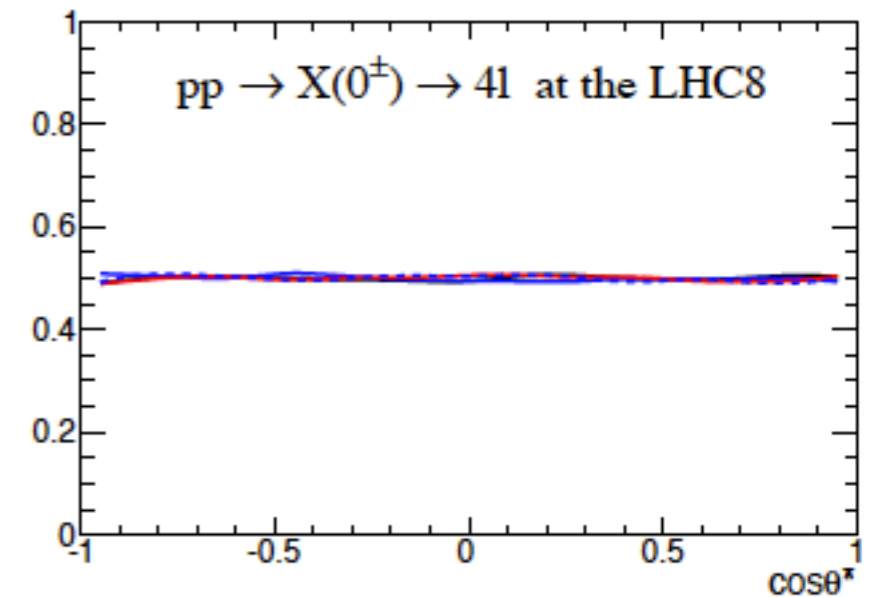
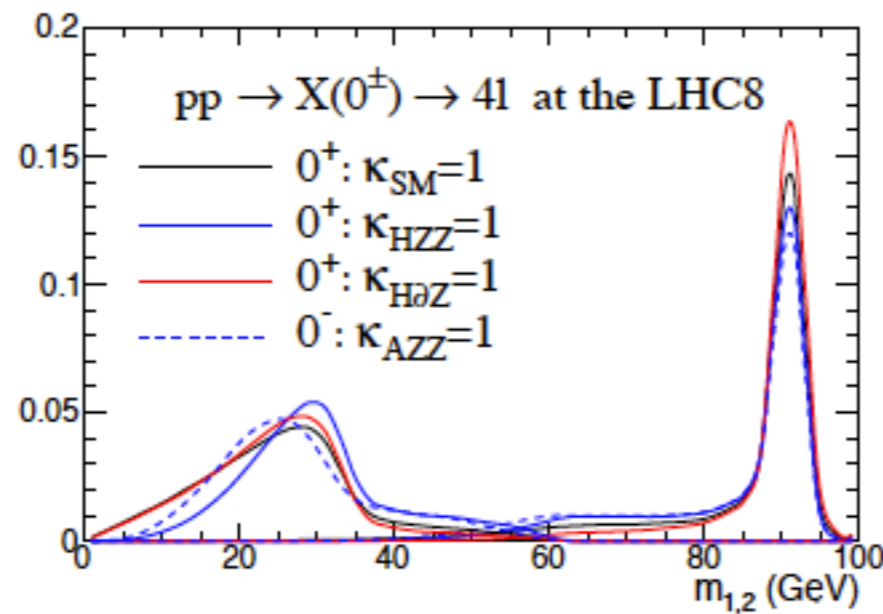
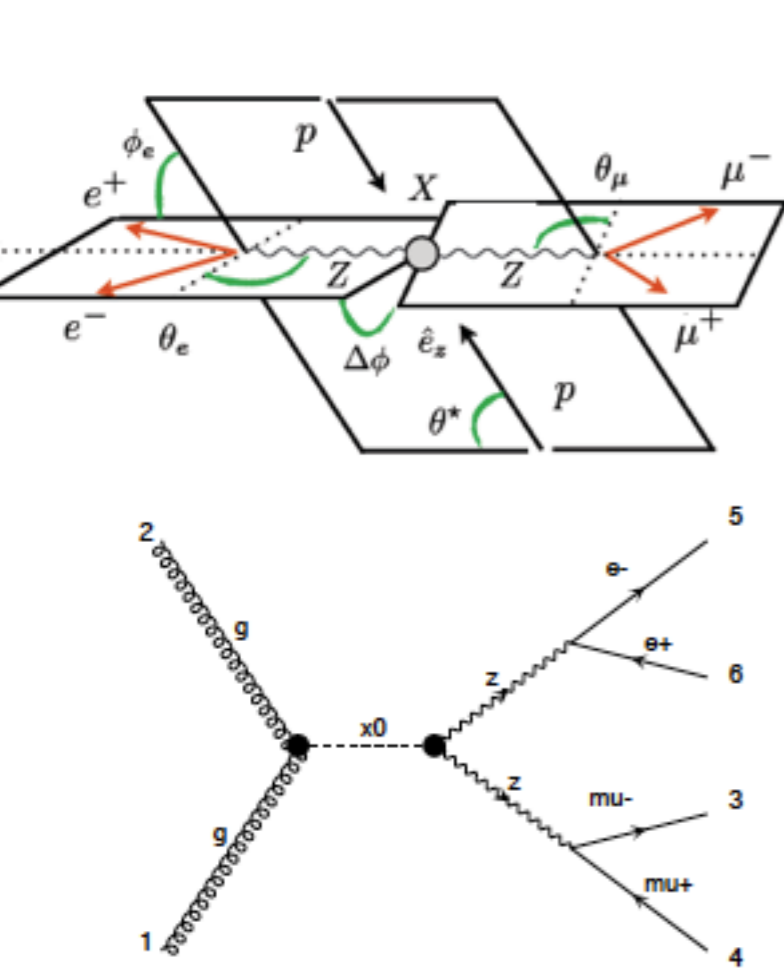
$$\mathcal{L}_0^f = - \sum_{f=t,b,\tau} \bar{\psi}_f (c_\alpha \kappa_{Hff} g_{Hff} + i s_\alpha \kappa_{Aff} g_{Aff} \gamma_5) \psi_f X_0$$

$$\mathcal{L}_0^V = \left\{ c_\alpha \kappa_{SM} \left[\frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \right. \\ - \frac{1}{4} \left[c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma} A_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma} A_{\mu\nu} \tilde{A}^{\mu\nu} \right] \\ - \frac{1}{2} \left[c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma} Z_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{AZ\gamma} g_{AZ\gamma} Z_{\mu\nu} \tilde{A}^{\mu\nu} \right] \\ - \frac{1}{4} \left[c_\alpha \kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + s_\alpha \kappa_{Agg} g_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right] \\ - \frac{1}{4} \frac{1}{\Lambda} \left[c_\alpha \kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + s_\alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] \\ - \frac{1}{2} \frac{1}{\Lambda} \left[c_\alpha \kappa_{HWW} W_{\mu\nu}^+ W^{-\mu\nu} + s_\alpha \kappa_{AWW} W_{\mu\nu}^+ \tilde{W}^{-\mu\nu} \right] \\ - \frac{1}{\Lambda} c_\alpha \left[\kappa_{H\theta\gamma} Z_\nu \partial_\mu A^{\mu\nu} + \kappa_{H\theta Z} Z_\nu \partial_\mu Z^{\mu\nu} \right. \\ \left. + (\kappa_{H\theta W} W_\nu^+ \partial_\mu W^{-\mu\nu} + h.c.) \right] \Big\} X_0$$

Dimensionful **couplings g** are set as internal parameters so as to reproduce a **SM Higgs** for $\kappa=1$.

$g_{X_{yy'}}$ $\times v$	ff	ZZ/WW	$\gamma\gamma$	$Z\gamma$	gg
H	m_f	$2m_{Z/W}^2$	$47\alpha_{EM}/18\pi$	$C(94 \cos^2 \theta_W - 13)/9\pi$	$-\alpha_s/3\pi$
A	m_f	0	$4\alpha_{EM}/3\pi$	$2C(8 \cos^2 \theta_W - 5)/3\pi$	$\alpha_s/2\pi$

Mass and angular distributions -- spin0



```

./bin/mg5_aMC
>import model HC
>generate p p > x0, x0 > mu- mu+ e- e+
>launch
    
```

Effective Lagrangian -- spin 1

- The most general interactions at the lowest canonical dimension:

$$\mathcal{L}_1^f = \sum_{f=q,\ell} \bar{\psi}_f \gamma_\mu (\kappa_{f_a} a_f - \kappa_{f_b} b_f \gamma_5) \psi_f X_1^\mu$$

$$\begin{aligned} \mathcal{L}_1^W = & i\kappa_{W_1} g_{WWZ} (W_{\mu\nu}^+ W^{-\mu} - W_{\mu\nu}^- W^{+\mu}) X_1^\nu + i\kappa_{W_2} g_{WWZ} W_\mu^+ W_\nu^- X_1^{\mu\nu} \\ & - \kappa_{W_3} W_\mu^+ W_\nu^- (\partial^\mu X_1^\nu + \partial^\nu X_1^\mu) \\ & + i\kappa_{W_4} W_\mu^+ W_\nu^- \tilde{X}_1^{\mu\nu} - \kappa_{W_5} \epsilon_{\mu\nu\rho\sigma} [W^{+\mu} (\partial^\rho W^{-\nu}) - (\partial^\rho W^{+\mu}) W^{-\nu}] X_1^\sigma \end{aligned}$$

$$\mathcal{L}_1^Z = -\kappa_{Z_1} Z_{\mu\nu} Z^\mu X_1^\nu - \kappa_{Z_3} X_1^\mu (\partial^\nu Z_\mu) Z_\nu - \kappa_{Z_5} \epsilon_{\mu\nu\rho\sigma} X_1^\mu Z^\nu (\partial^\rho Z^\sigma)$$

- Parity conservation implies that

▶ for X_{1-} $\kappa_{f_b} = \kappa_{V_4} = \kappa_{V_5} = 0$

▶ for X_{1+} $\kappa_{f_a} = \kappa_{V_1} = \kappa_{V_2} = \kappa_{V_3} = 0$

Effective Lagrangian -- spin2

- via the energy-momentum tensor of the SM fields, starting from D5:

$$\mathcal{L}_2^f = -\frac{1}{\Lambda} \sum_{f=q,\ell} \kappa_f T_{\mu\nu}^f X_2^{\mu\nu}$$

$$\mathcal{L}_2^V = -\frac{1}{\Lambda} \sum_{V=Z,W,\gamma,g} \kappa_V T_{\mu\nu}^V X_2^{\mu\nu}$$

► The E-M tensor for QED:

$$T_{\mu\nu}^f = -g_{\mu\nu} \left[\bar{\psi}_f (i\gamma^\rho D_\rho - m_f) \psi_f - \frac{1}{2} \partial^\rho (\bar{\psi}_f i\gamma_\rho \psi_f) \right]$$

$$+ \left[\frac{1}{2} \bar{\psi}_f i\gamma_\mu D_\nu \psi_f - \frac{1}{4} \partial_\mu (\bar{\psi}_f i\gamma_\nu \psi_f) + (\mu \leftrightarrow \nu) \right],$$

$$T_{\mu\nu}^\gamma = -g_{\mu\nu} \left[-\frac{1}{4} A^{\rho\sigma} A_{\rho\sigma} + \partial^\rho \partial^\sigma A_\sigma A_\rho + \frac{1}{2} (\partial^\rho A_\rho)^2 \right]$$

$$- A_\mu^\rho A_{\nu\rho} + \partial_\mu \partial^\rho A_\rho A_\nu + \partial_\nu \partial^\rho A_\rho A_\mu,$$



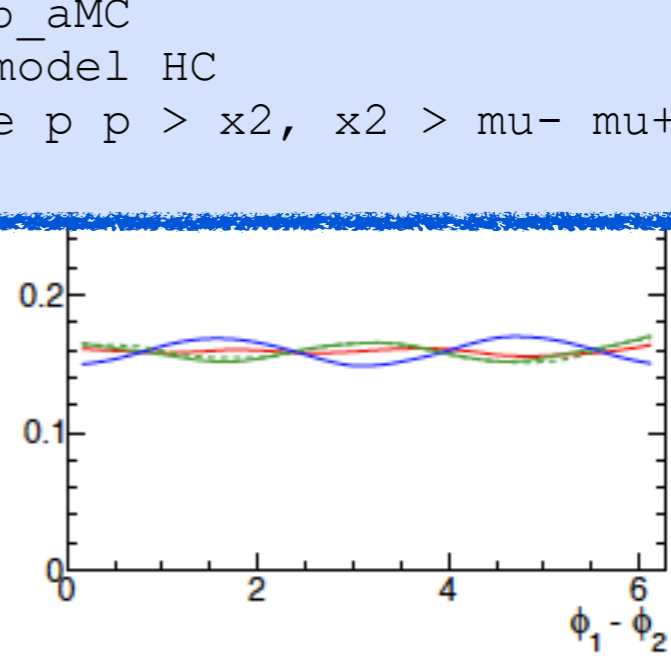
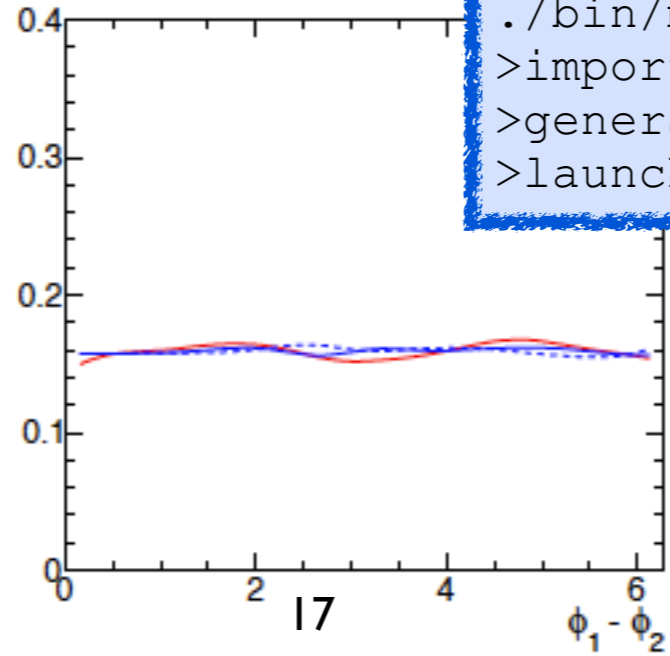
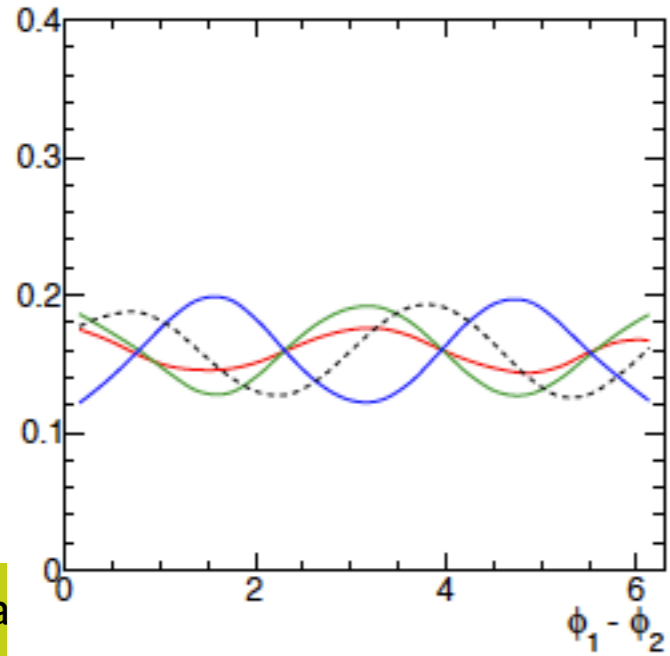
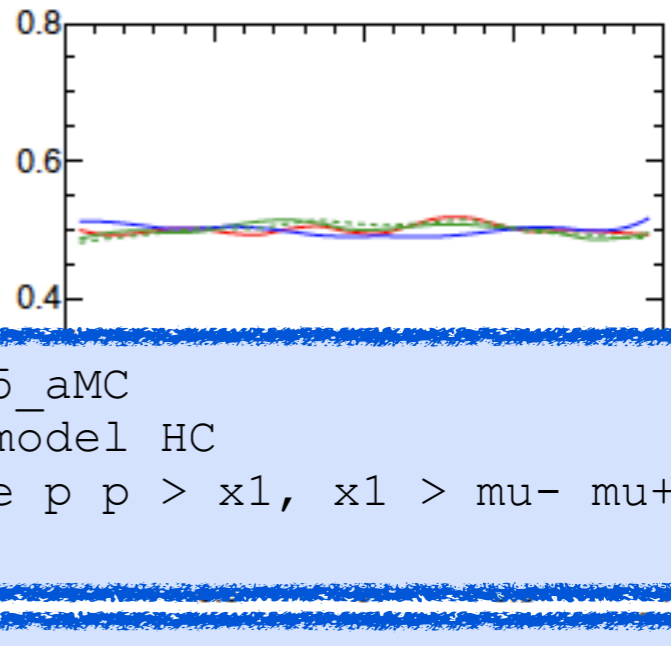
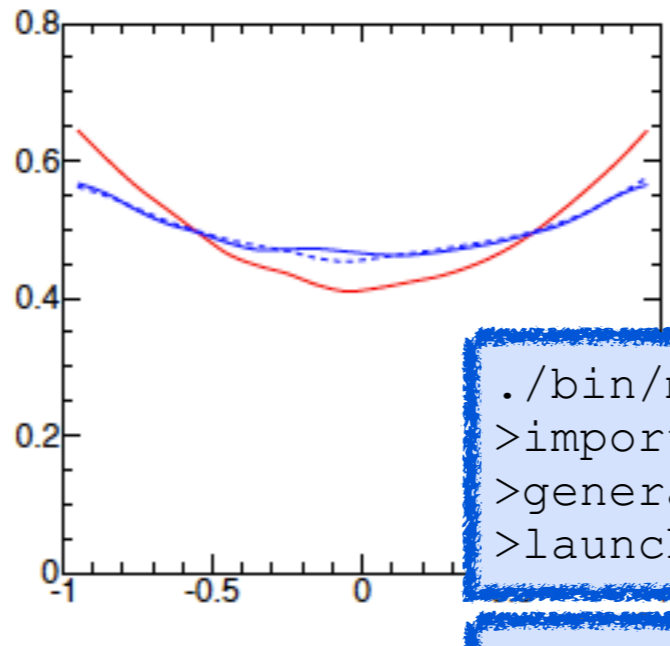
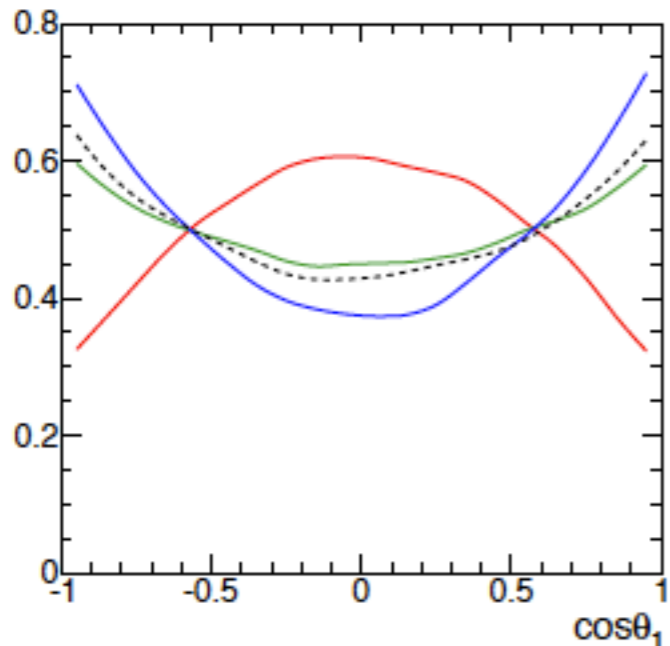
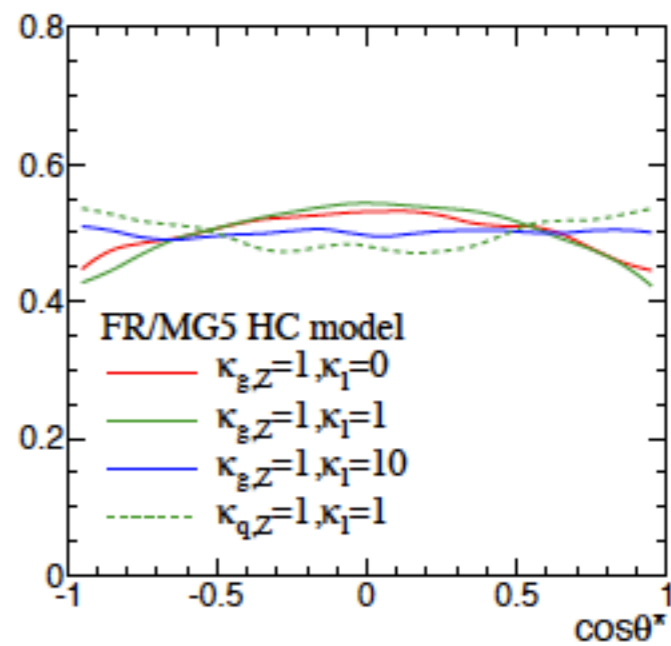
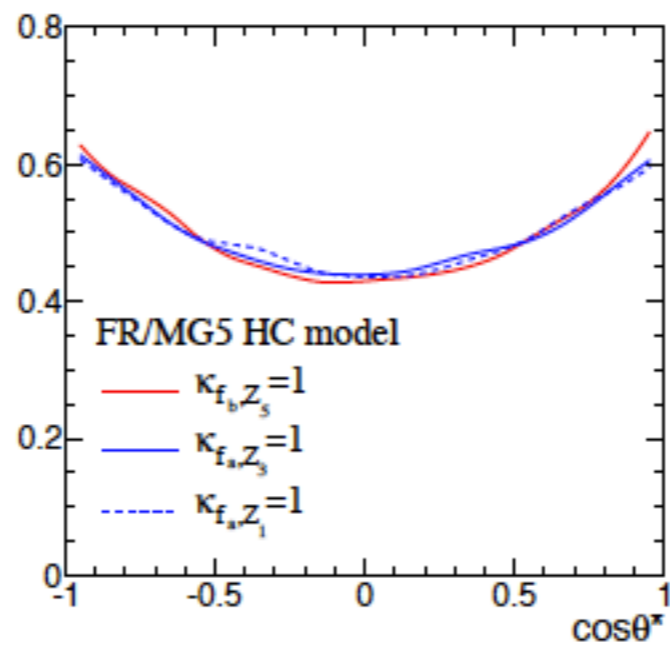
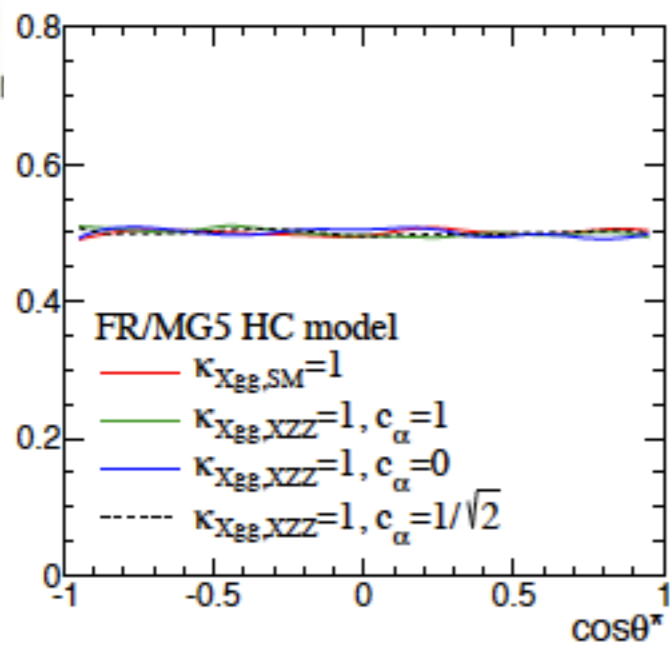
V
U
B



spin-0

spin-1

spin-2

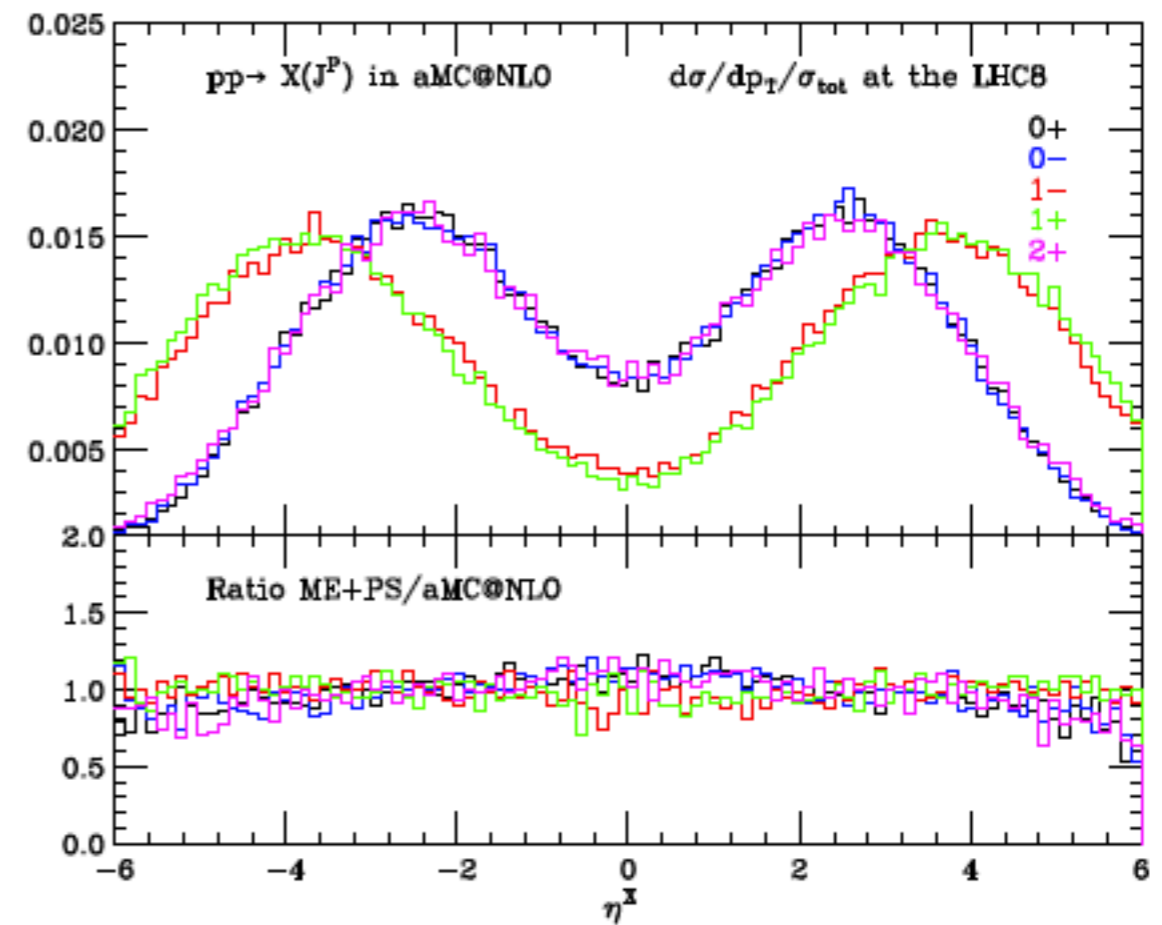
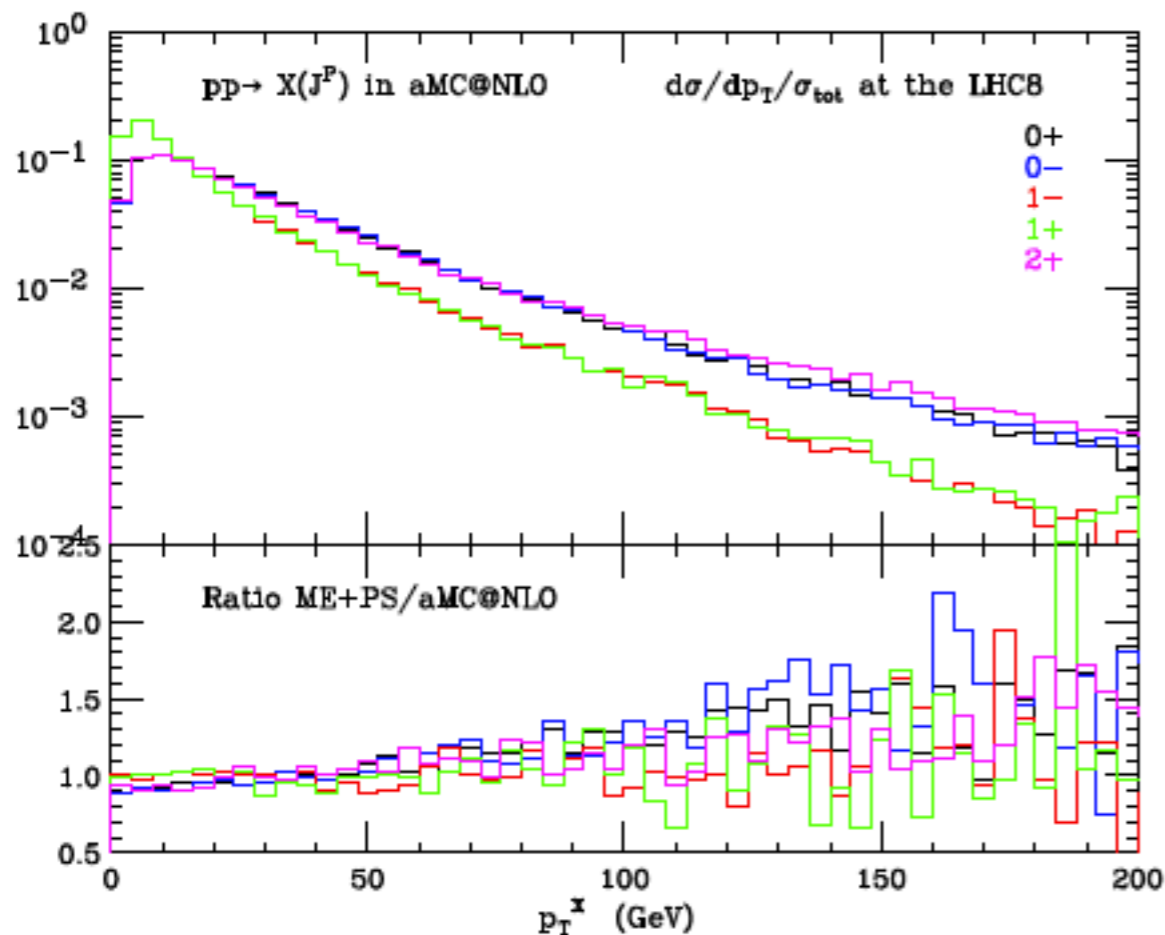


```
./bin/mg5_aMC
>import model HC
>generate p p > x1, x1 > mu- mu+ e- e+
>launch
```

```
./bin/mg5_aMC
>import model HC
>generate p p > x2, x2 > mu- mu+ e- e+
>launch
```

Higher order effects in QCD

- The LO predictions can be systematically improved by including the effects due to the emission of QCD partons.
 - ▶ LO Matrix-Element/Parton-Shower merging [[ME+PS](#)]
 - ▶ full-NLO matrix element with parton-shower [[aMC@NLO+Pythia/Herwig](#)]



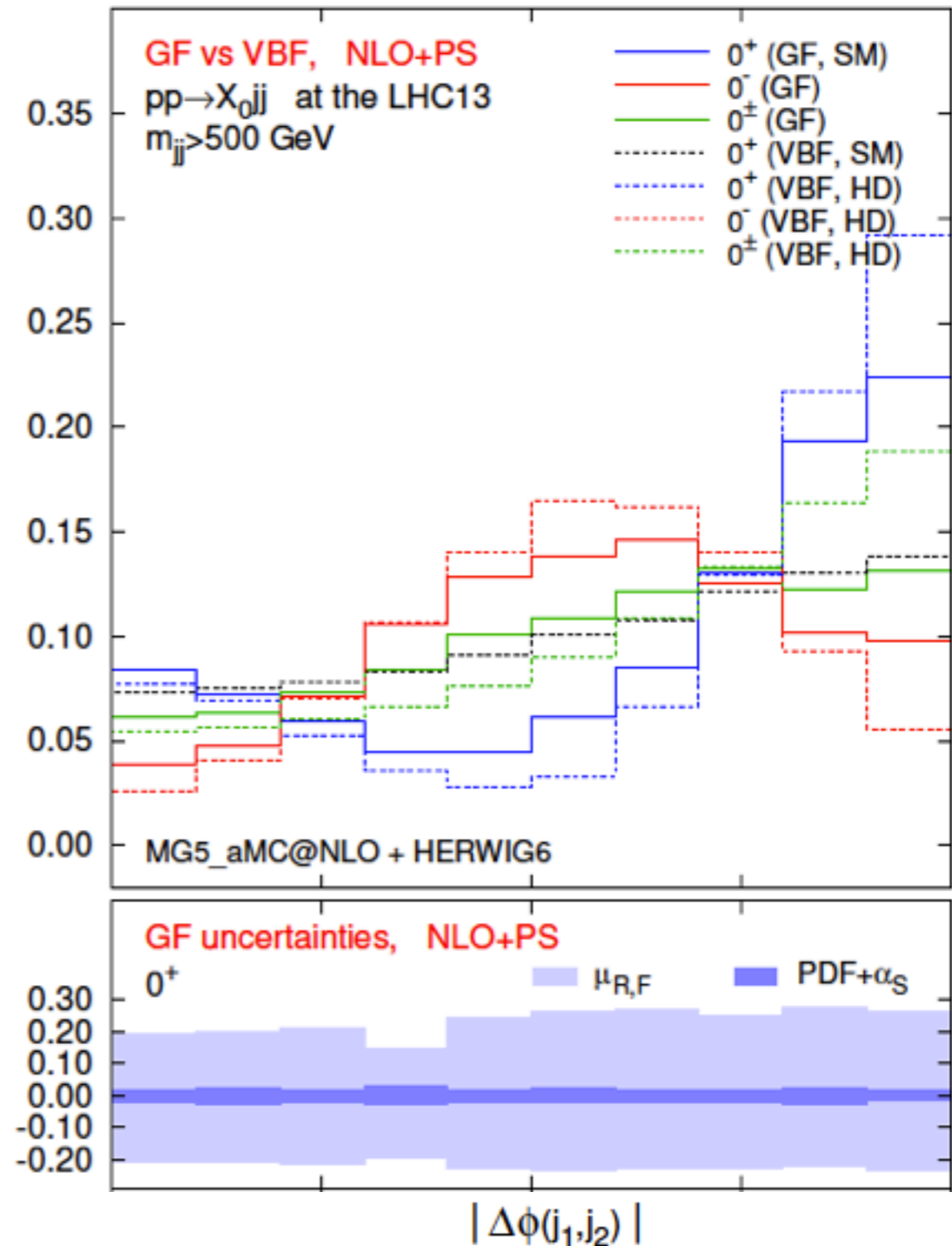
VBF

Maltoni, KM, Zaro [arXiv:1311.1829]

```
./bin/mg5_aMC
> import model HC_NLO
> generate p p > x0 j j QCD=0 [QCD]
> launch
```

Demartin, Maltoni, KM, Page, Pittau, Zaro [in progress]

```
./bin/mg5_aMC
> import model HC_NLO-heft
> generate p p > x0 j j / t [QCD]
> launch
```



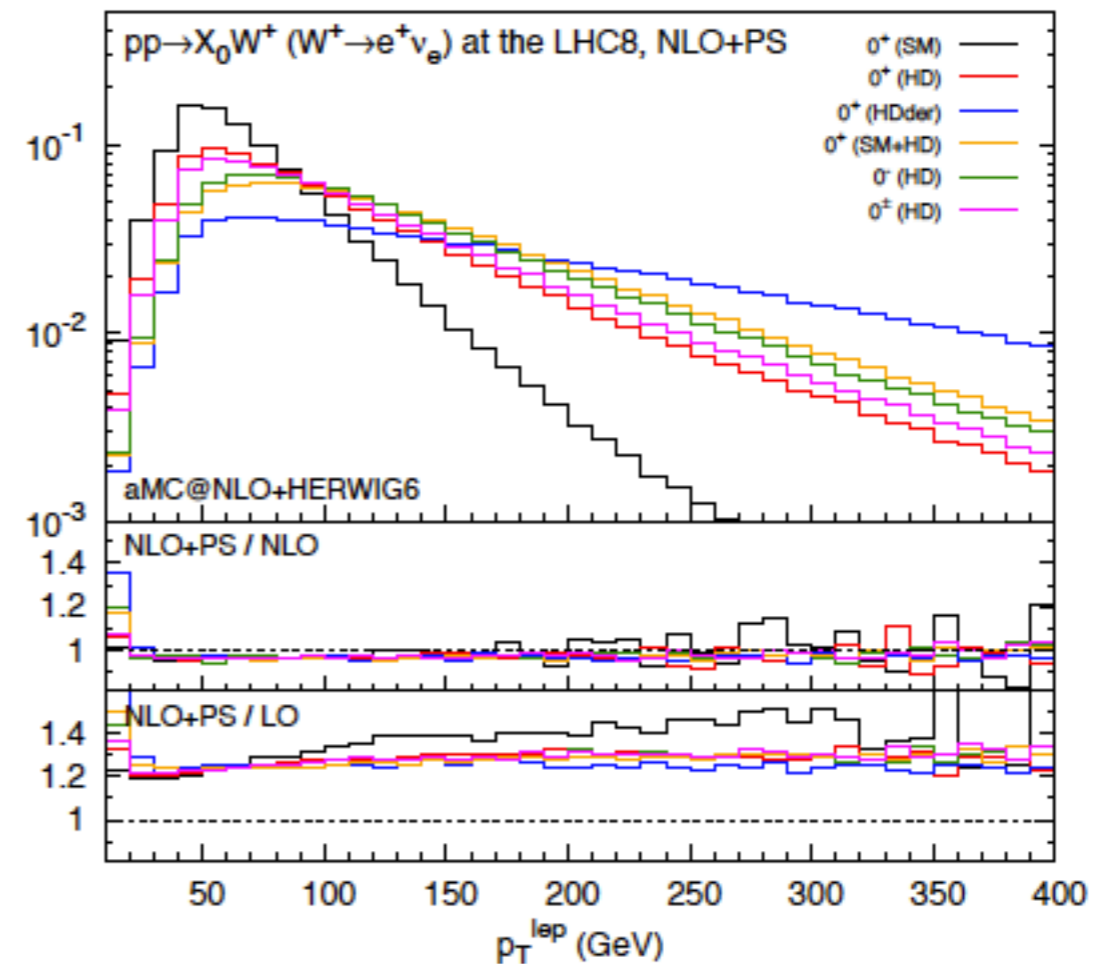
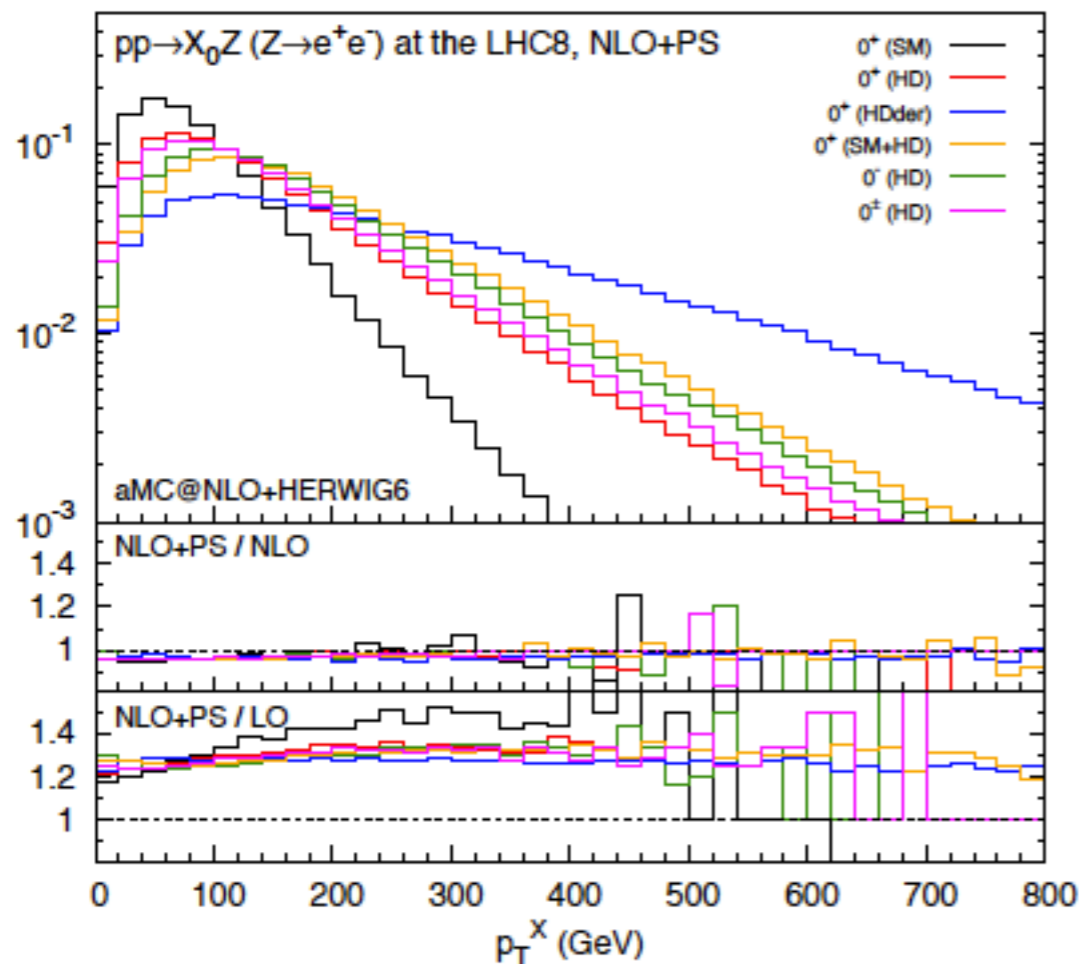
scenario		σ_{LO} (pb)		σ_{NLO} (pb)	
LHC 8 TeV	0^+	1.351(1)	+67.1 +4.3% -36.8 -4.3%	1.702(6)	+19.7 +1.7% -20.8 -1.7%
	0^-	2.951(3)	+67.2 +4.4% -36.8 -4.4%	3.660(15)	+19.1 +1.7% -20.6 -1.7%
	0^\pm	2.142(2)	+67.1 +4.4% -36.8 -4.4%	2.687(10)	+19.6 +1.7% -20.8 -1.7%
LHC 13 TeV	0^+	4.265(4)	+61.5 +3.3% -34.9 -3.3%	5.092(23)	+15.4 +1.2% -17.9 -1.2%
	0^-	9.304(9)	+61.6 +3.4% -34.9 -3.4%	11.29(4)	+16.0 +1.2% -18.2 -1.2%
	0^\pm	6.775(6)	+61.5 +3.3% -34.9 -3.3%	8.055(35)	+15.8 +1.2% -18.2 -1.2%

VH

Maltoni, KM, Zaro [arXiv:1311.1829]

```
./bin/mg5_aMC
> import model HC_NLO
> generate p p > x0 e+ e- [QCD]
> launch
```

scenario	σ_{LO} (fb)	σ_{NLO} (fb)	K
0^+ (SM)	39.58(3) $+0.1\%$ -0.6%	51.22(5) $+2.2\%$ -1.8%	1.29
0^+ (HD)	13.51(1) $+1.5\%$ -1.7%	17.51(1) $+1.9\%$ -1.3%	1.30
0^+ (HDder)	324.2(2) $+4.7\%$ -4.3%	416.1(4) $+2.3\%$ -2.1%	1.28
0^+ (SM+HD)	118.8(1) $+3.0\%$ -2.9%	154.2(1) $+1.8\%$ -1.6%	1.30
0^- (HD)	8.386(7) $+2.6\%$ -2.6%	10.89(1) $+1.8\%$ -1.5%	1.30
0^\pm (HD)	10.96(1) $+1.9\%$ -2.1%	14.22(1) $+1.8\%$ -1.3%	1.30



Summary

- After the discovery of a Higgs-like resonance at the LHC, the main focus of the analyses now is **the determination of the Higgs Lagrangian**.
- This includes
 - **the structure of the operators**, linked to the spin/parity of a ‘Higgs’ boson.
 - **the coupling strength**.
- MC tools to study the property of the SM-like boson are publicly available, e.g. JHUGen, MEKD, HC, HEL.
- Event generation at NLO is possible for (several) spin 0,1,2 hypothesis and can be used to validate merged samples.

HiggsCharacterisation - FeynRules

http://feynrules.irmp.ucl.ac.be/wiki/HiggsCharacterisation

HiggsCharacterisation - FeynRules

The Higgs Characterisation model

Authors

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Description of the model

This is a model file for the characterisation of the boson recently discovered at the LHC. Our effective lagrangian consists of the SM (except for the Higgs itself), expressed through the physical degrees of freedom present below the EWSB scale, plus a new bosonic state $X(J^P)$ with spin/parity assignments $J^P = 0^+, 0^-, 1^+, 1^-,$ or 2^+ . The new state can couple to SM particles via interactions of the lowest possible dimensions. In addition, the state 0^+ is allowed to mix with the 0^- one, and can interact with SM particles with higher-dimensional operators beyond those of the SM. See more details in

- [1306.6464](#) : P. Artoisenet, P. de Aquino, F. Demartin, R. Frederix, S. Frixione, F. Maltoni, M. K. Mandal, P. Mathews, K. Mawatari, V. Ravindran, S. Seth, P. Torrielli, M. Zaro, "A framework for Higgs characterisation" (JHEP11(2013)043).
- [1307.5607](#) : P. de Aquino, K. Mawatari, "Characterising a Higgs-like resonance at the LHC" (Proceedings for HPNP2013).
- [1311.1829](#) : F. Maltoni, K. Mawatari, M. Zaro, "Higgs characterisation via vector-boson fusion and associated production: NLO and parton-shower effects" (EPJC74(2014)2710).

Model files for LO

- [HC.fr](#) : the main model file.
- [SM_HC.fr](#) : This model requires the modified Standard Model Implementation of [FeynRules](#).
- [Massless.rst](#), [Cabibbo.rst](#) : SM restriction files.
- [HC.nb](#) : this is an example Mathematica notebook that loads the model, calculates the Feynman rules and extract the model files within the UFO format.
- [HC_UFO.zip](#) : The model files in UFO format (for MadGraph5).

Model files for NLO (only for the J=0 case)

A few remarks before use; see the README file in the model.

- [HC_NLO_X0_UFO.zip](#) : The model files in UFO format (for MadGraph5).