



What does a Higgs Tell us about SUSY?

Based on work with
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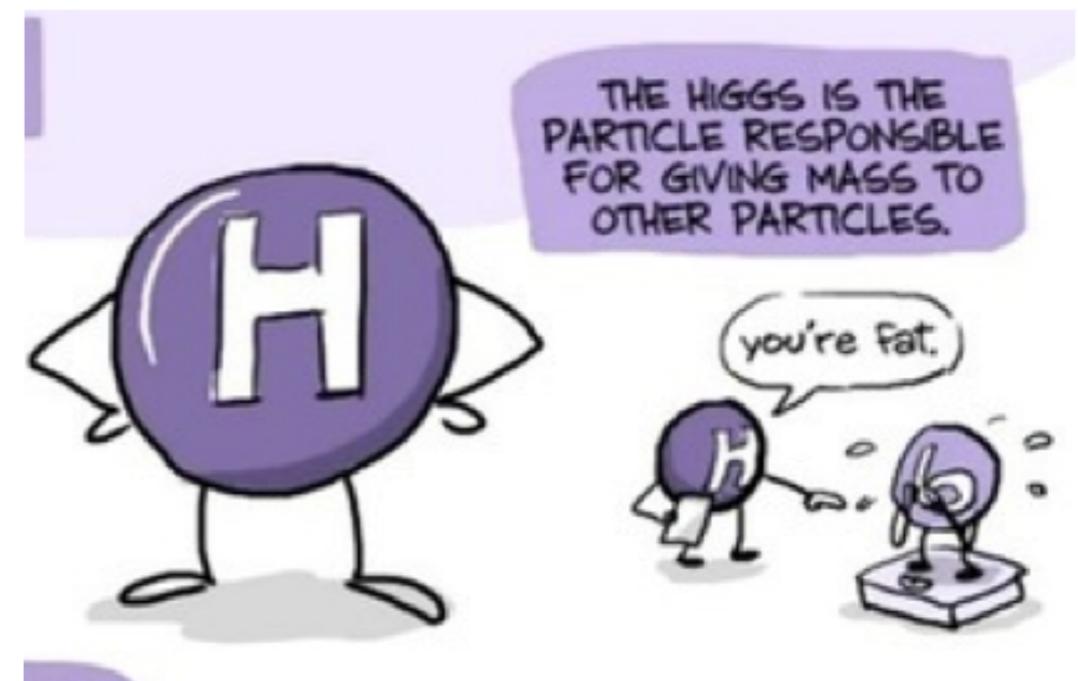
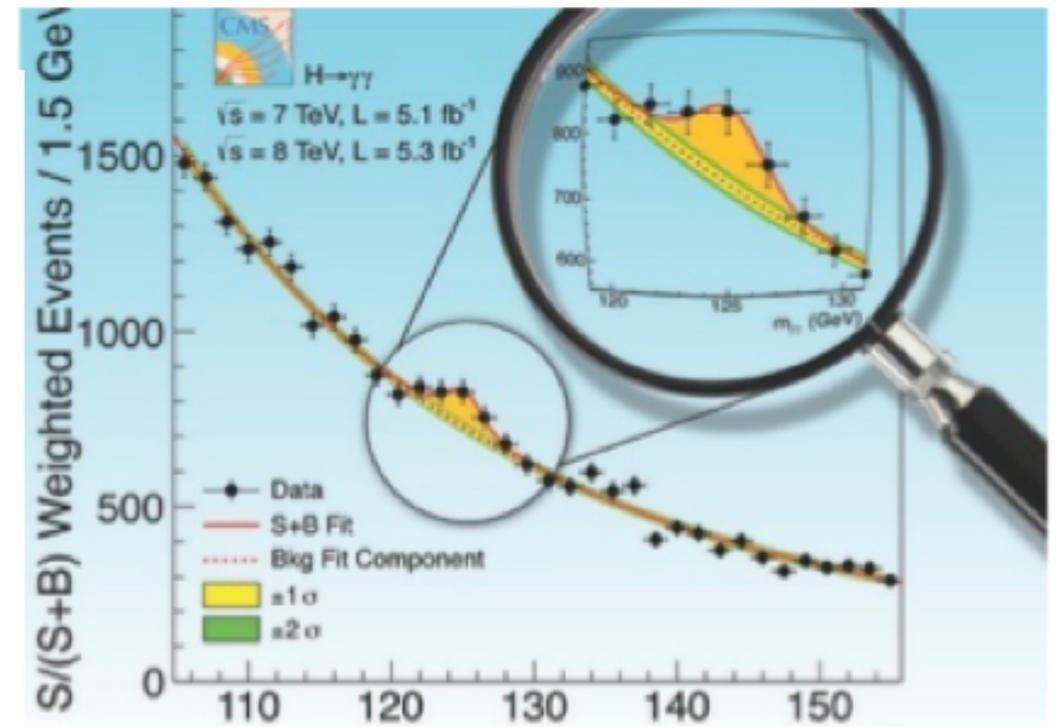
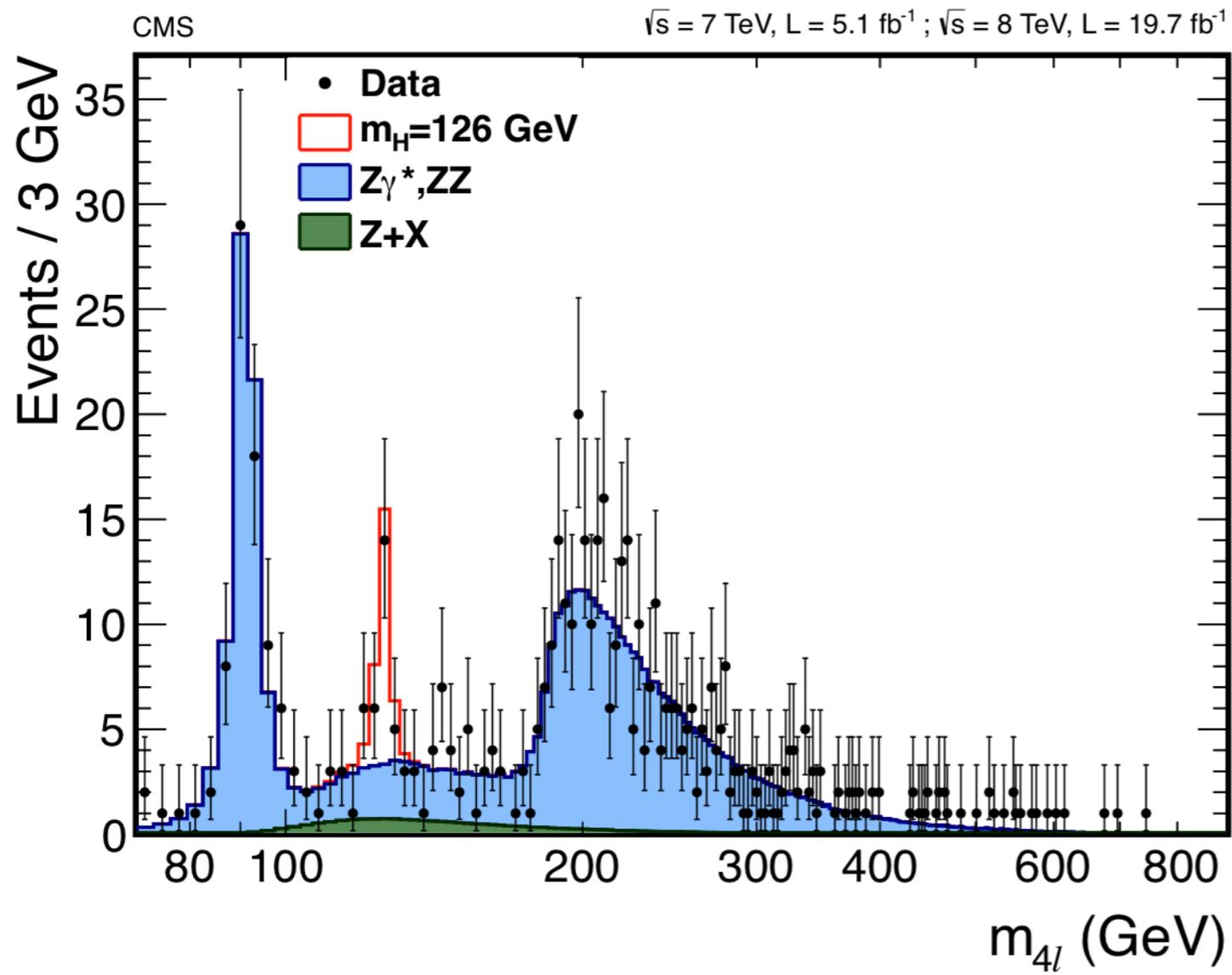
Mitchell Workshop on Collider and Dark Matter Physics

Peisi Huang

ANL & University of Chicago



What do We Know about the Higgs?



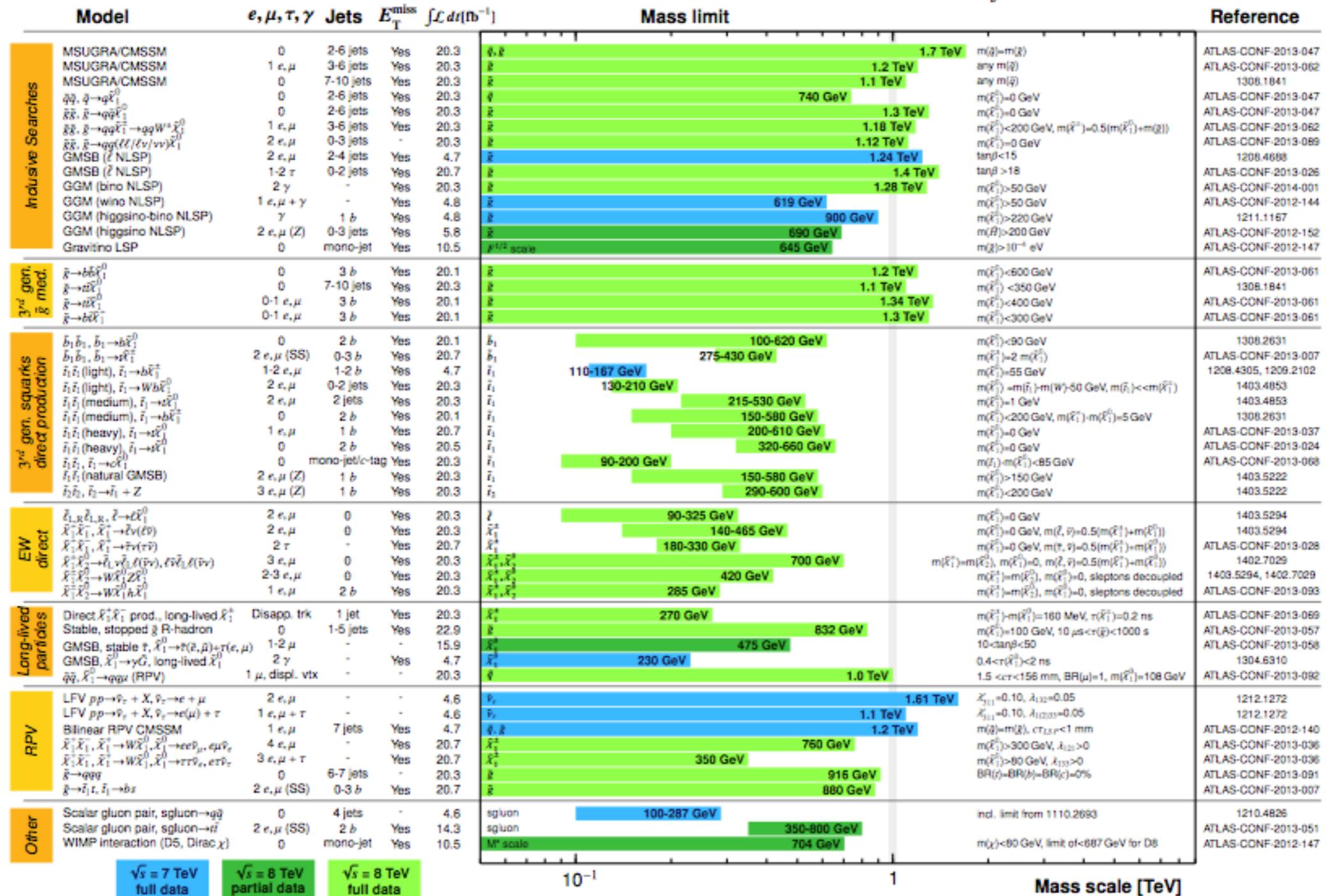
What do We Know about SUSY?

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: Moriond 2014

ATLAS Preliminary

$\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$ $\sqrt{s} = 7, 8 \text{ TeV}$



*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

What can We Learn about SUSY from Higgs?

- Calculate higgs mass in MSSM
- Consider the inverse problem, higgs to SUSY
- Search for SUSY with a higgs in the final states

Higgs Mass, MSSM

- Two higgs doublets H_u and H_d
- D-terms $\frac{1}{8}(g^2 + g'^2)(|H_u^0|^2 - |H_d^0|^2)^2$
- Tree level bound $m_h \leq m_Z \cos 2\beta$
- Important quantum corrections to the higgs mass from stops and sbottoms/staus for large $\tan\beta$

Loop Corrections to the MSSM Higgs Mass

- 1 loop • CP-odd higgs mass m_A , $\tan\beta$, stop mass and mixing.

$$M_{\tilde{t}} = \begin{pmatrix} m_{\tilde{t}_L}^2 + m_t^2 + D_L & -m_t X_t \\ -m_t X_t & m_{\tilde{t}_R}^2 + m_t^2 + D_R \end{pmatrix}$$

- For moderate to large value of $\tan\beta$

Polynomial dependence on the stop mixing

$$m_h^2 \simeq M_Z^2 \cos^2 2\beta + \frac{3g^2}{8\pi^2} \frac{m_t^4}{m_W^2} \left[\ln \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{X_t^2}{12m_{\tilde{t}}^2} \right) \right]$$

Logarithmic growth with stop mass

2 loop

A lot of efforts: Carena, Casas, Degrassi, Ellis, Espinosa, Haber, Heinemeyer, Hempfling, Hoang, Hollik, Martin, Quiros, Ridolfi, Slavich, Wagner, Weiglein, Zvriner, Zhang...

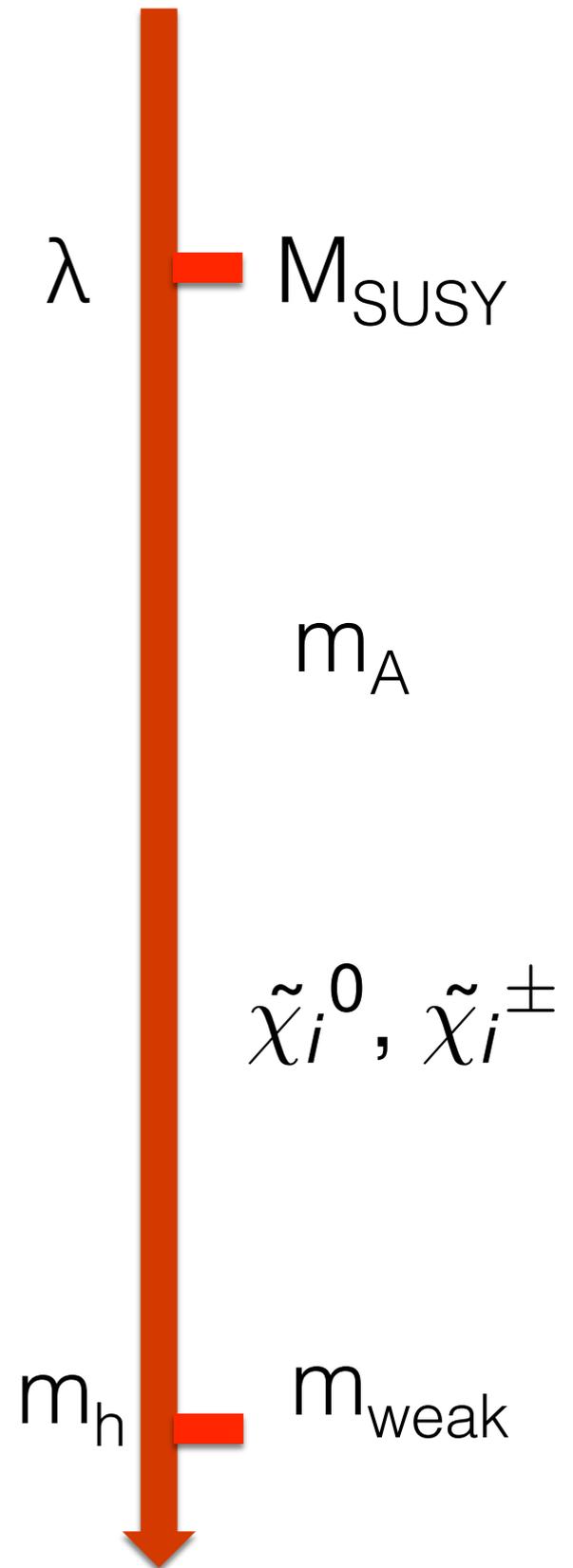
In public codes: FeynHiggs, Isajet, SuSpect, Softsusy, Spheno

Higher Order Corrections to Higgs Mass

- Diagrammatic calculations:
 - ❖ Up to partial three loop order
 - Harlander, Kant, Mihaila and Steinhauser.
 - ❖ In a mathematica code H3M
 - ❖ 6 particular mass hierarchies
- Effective potential method:
 - ❖ Derivatives of the MSSM potential $V(h_1, h_2)$ evaluated at the vev
 - ❖ Zero external momentum approximation of the diagrammatic calculation
 - ❖ Up to partial three loop order, Martin

Higher Order Corrections to Higgs Mass

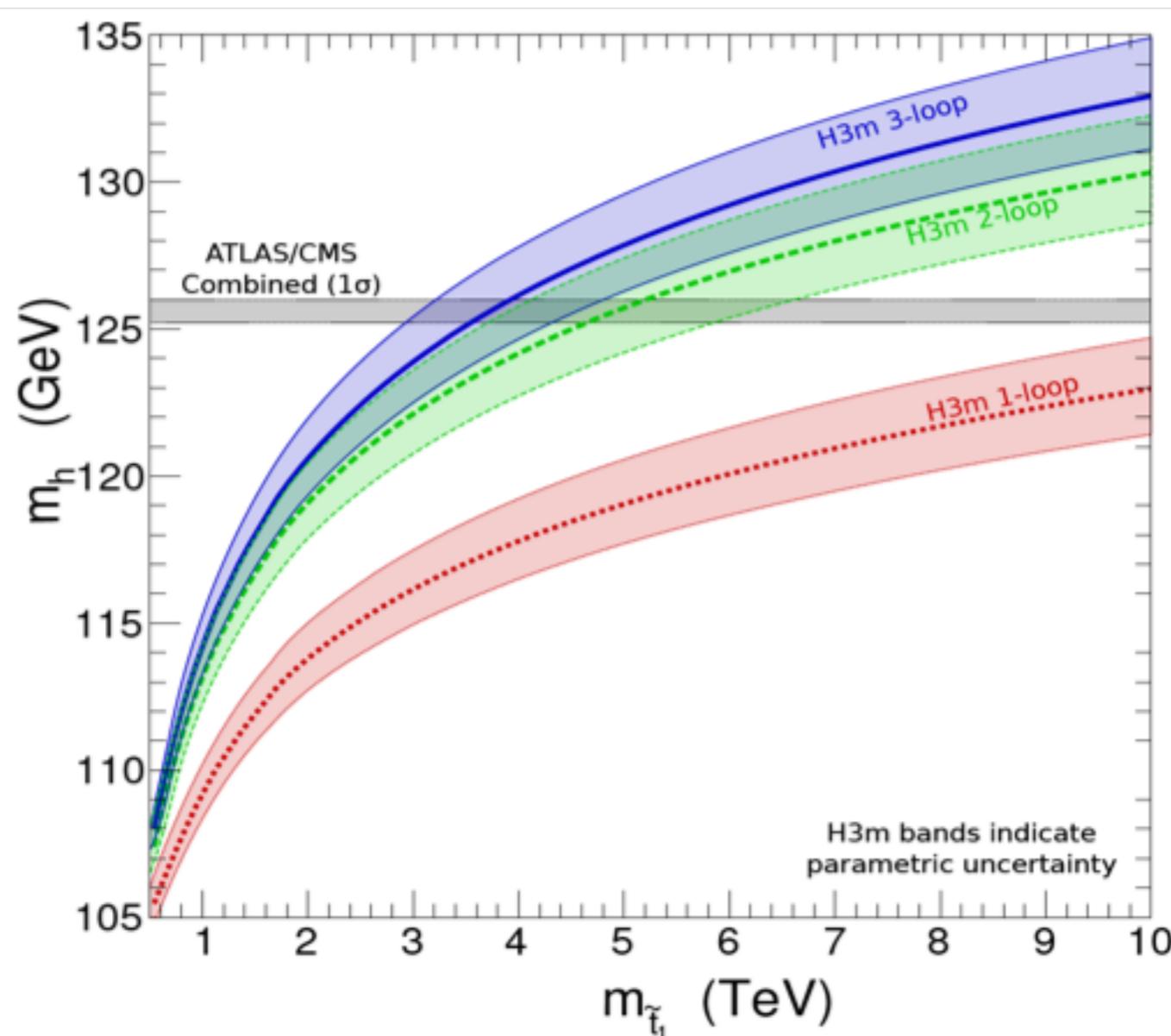
- Effective field theory analysis
 - ❖ Running the effective theory couplings (quartic couplings) down to the electroweak scale
 - ❖ Integrate out the MSSM particles at their thresholds
 - ❖ Evaluate the higgs pole mass at the electroweak scale
 - ❖ Partial three loop order, Degrandi, Di Vita, Elia-Miro, Espinosa, Giudice, Isidori and Strumia; Draper, Lee and Wagner



Higher Order Corrections to Higgs Mass

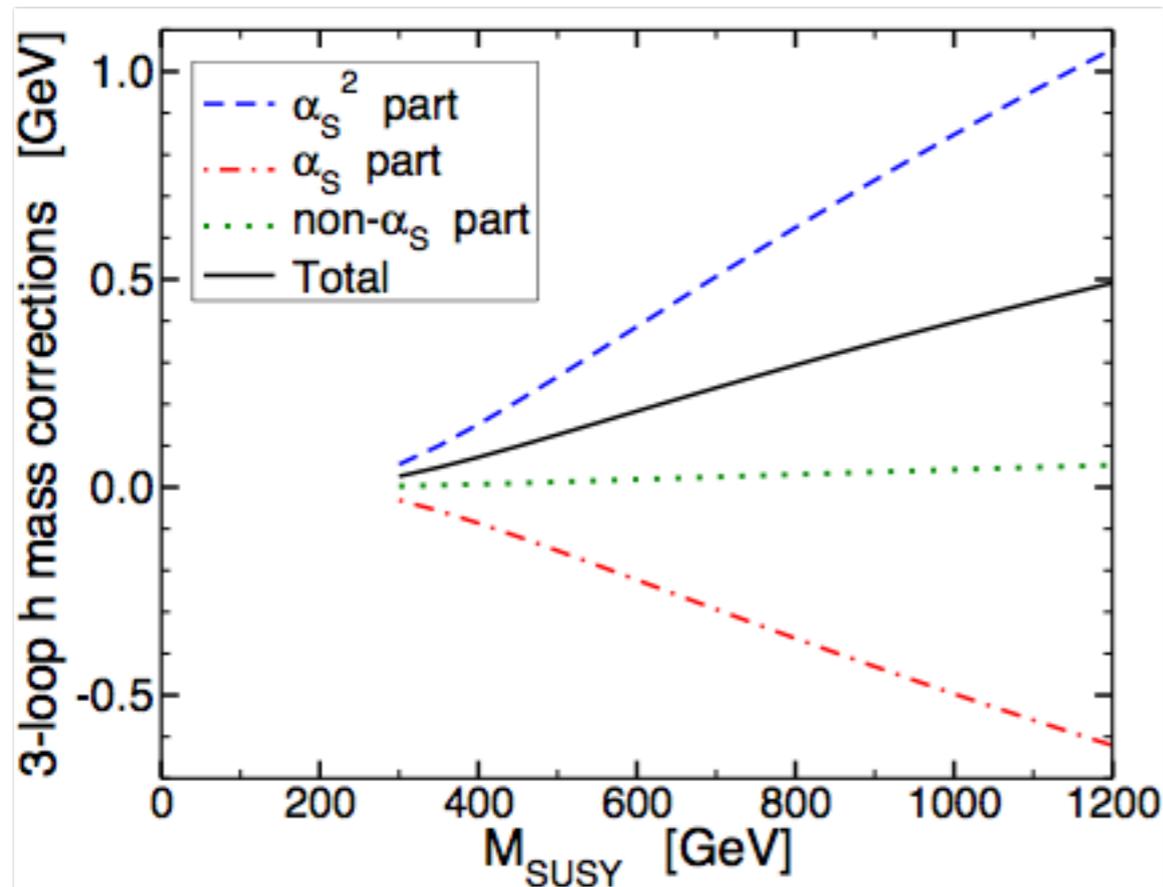
- Low SUSY scales, fixed order calculations (diagrammatic calculations and effective potential method) are most accurate.
 - ❖ Keep a large set of non-logarithmic terms
- High SUSY scales, especially in the simplified models, EFT analysis is more accurate.
 - ❖ Resume large log terms

Higher Order Corrections to Higgs Mass, Diagrammatic Calculation



- Degenerate, non-mixed stops
- Gaugino mass unification and gluino mass is 1.5 TeV
- $\tan\beta = 20$, $\mu = 200$ GeV
- The uncertainties are from top mass and $\alpha_s(m_Z)$
- Stop mass ~ 3 TeV

Higher Order Corrections to Higgs Mass, Effective Potential

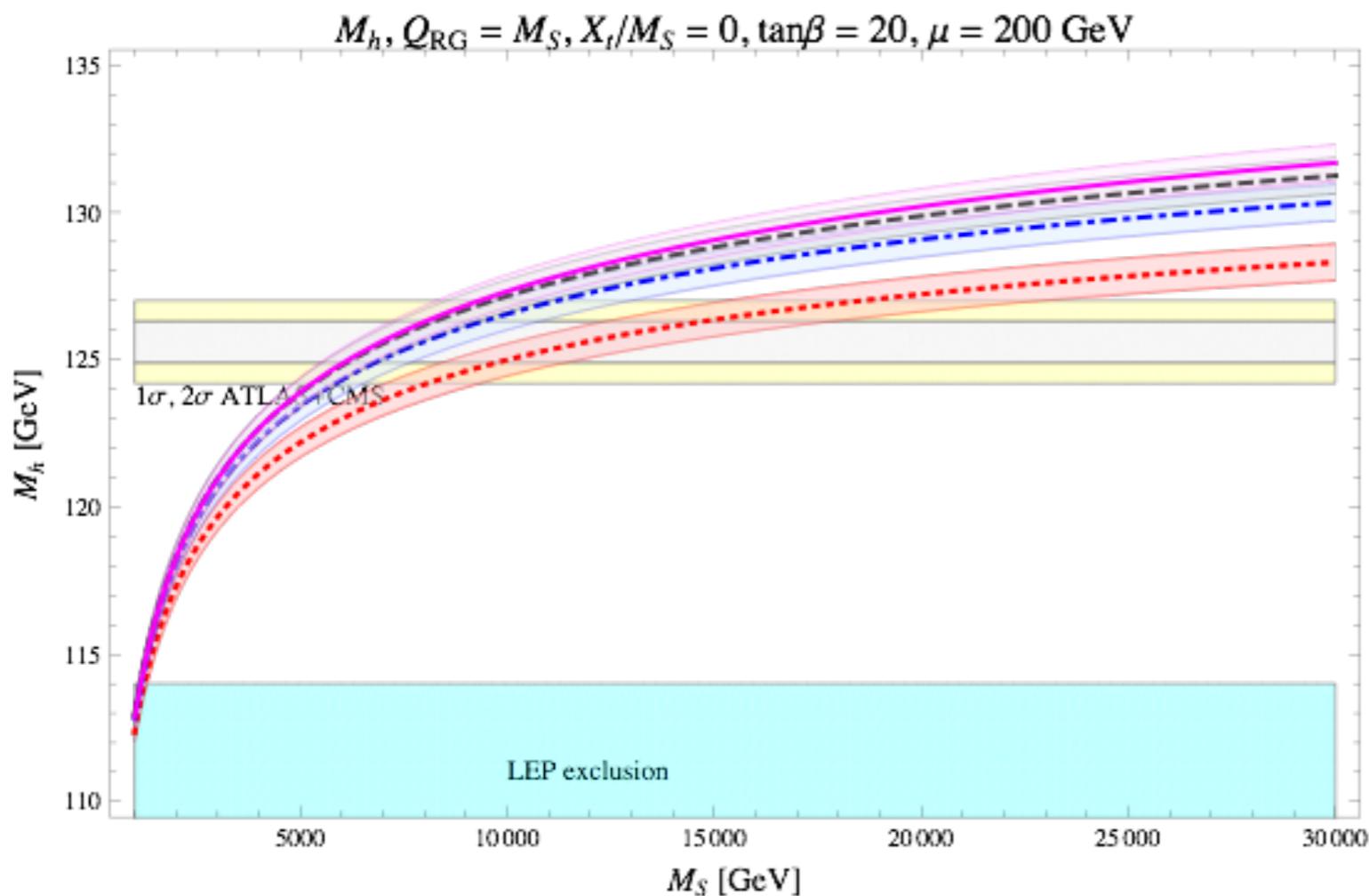


- Degenerate and non-mixed stops
- Large $\tan\beta$
- Leading logs
- Large stop mass

Martin, 2007

Higher Order Corrections to Higgs Mass, EFT Analysis

1loop 2loop 3loop



- Degenerate, non-mixed stops
- All other SUSY particles are decoupled
- $\tan\beta = 20, \mu = 200 \text{ GeV}$
- Stop mass $\sim 7 \text{ TeV}$

Draper, Lee and Wagner

What about the Inverse Problem?

$$\lambda_{SM} = \lambda_{MSSM} \quad \text{at } M_{SUSY}$$

- Start from the electroweak scale, calculate λ
- run λ , coupled with g_1, g_2, g_3, y_t up: $\lambda_{SM}(Q)$
- integrate out the MSSM particles at their threshold
- Calculate $\lambda_{MSSM}(Q, m_{\tilde{g}}, m_A, m_{\tilde{q}} \dots)$ from V_{eff}
- Estimate the SUSY scale:
where λ_{SM} and λ_{MSSM} match


$$\lambda_{SM} \quad \text{at } m_{\text{weak}}$$

$$\tilde{\chi}_i^0, \tilde{\chi}_i^\pm$$

$$m_A$$

What about the Inverse Problem?

Ingredients

- Higgs mass is known from the experiment
- Quartic coupling λ is known to 2 loop order

Bezrukov, Kalmykov, Kniehl, and Shaposhnikov 2012

- β -function for higgs quartic coupling is known to 3 loop order

Chetyrkin and Zoller, 2013

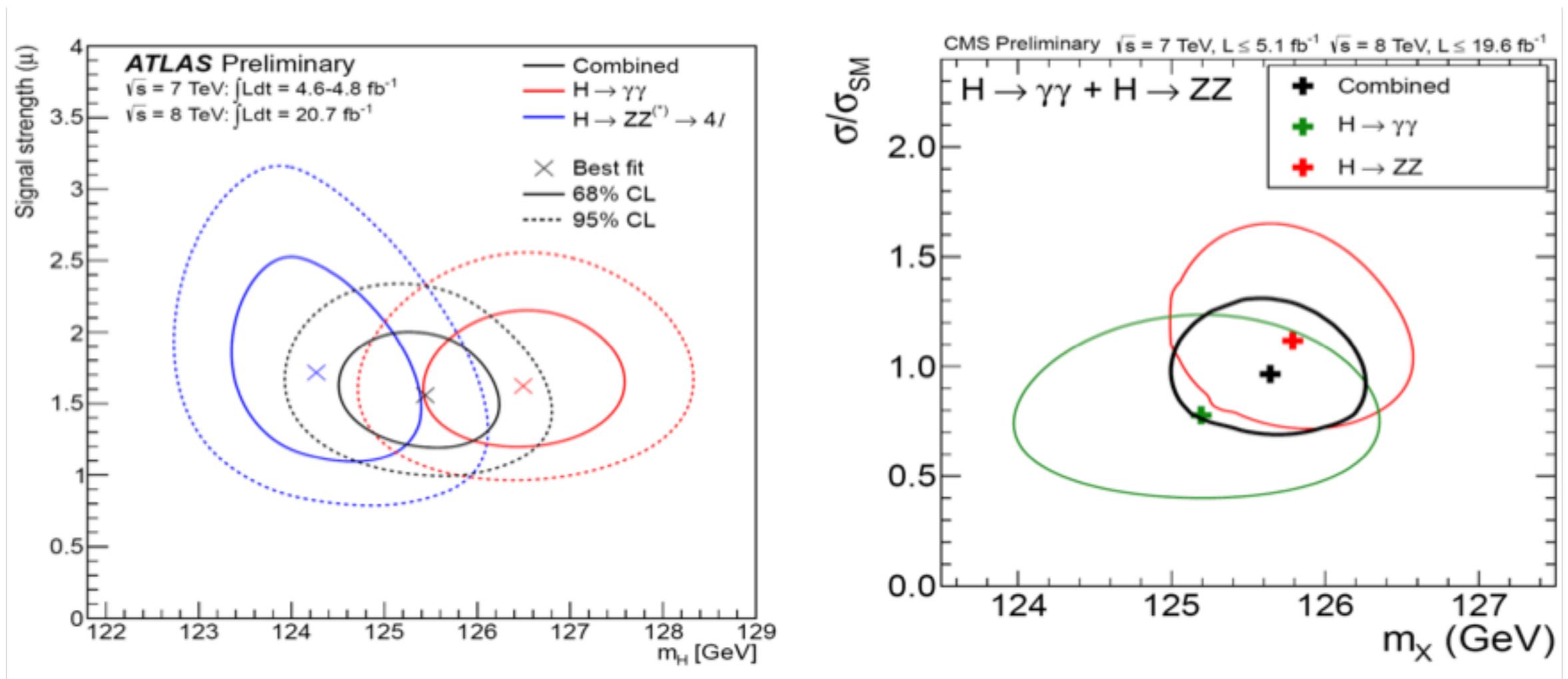
- SM RGE is known to 3 loop order

Chetyrkin and Zoller, 2012

- MSSM effective potential is known to 2 loop order

Espinosa and Zhang, 2001

Boundary Conditions at Weak Scale, Higgs Mass



$$\left. \begin{array}{l} \text{ATLAS (combined) : } 125.5 \pm 0.2^{+0.5}_{-0.6} \text{ GeV} \\ \text{CMS (combined) : } 125.7 \pm 0.3 \pm 0.3 \text{ GeV} \end{array} \right\} \Rightarrow 125.6 \pm 0.4 \text{ GeV}$$

Boundary Conditions at Weak Scale

$$\alpha_s(m_Z) = 0.1184 \pm 0.0007$$

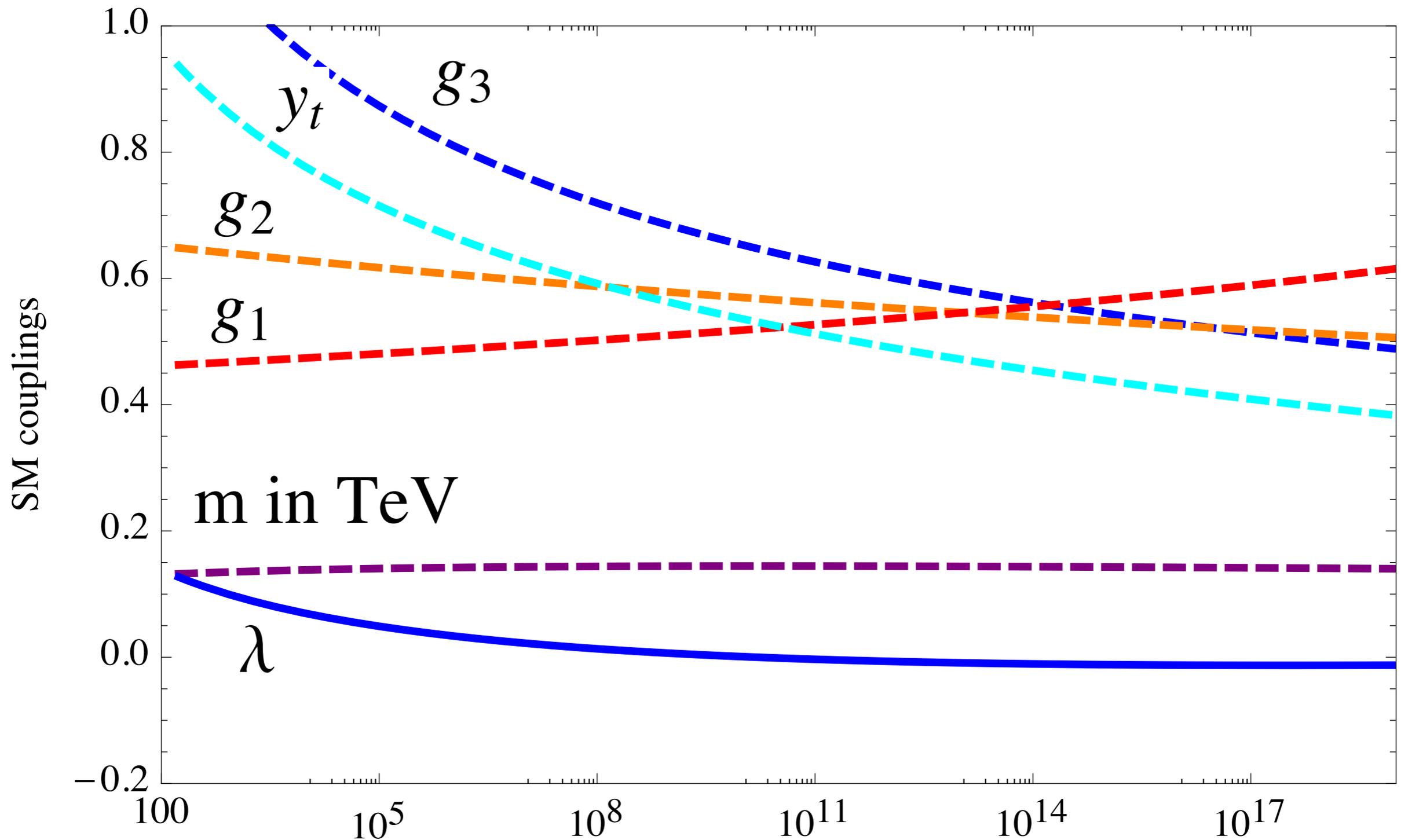
$$m_t = 173.2 \pm 0.87 \text{ (Tevatron)}$$

	LO 0 loop	NLO 1 loop	NNLO 2 loop
g_2	$2m_W/V$	full	full
g	$2\sqrt{m_Z^2 - m_W^2}/V$	full	full
y_t	$\sqrt{2}m_t/V$	$O(\alpha_s), O(\alpha)$	$O(\alpha_s^2, \alpha_s\alpha_{1,2})$
λ	$m_h^2/2V^2$	full	full

$$V \equiv (\sqrt{2}G_\mu)^{-1/2}$$

Sirlin; Tarrach; Hempfling and Kneihl; Buttazzo, Degrassi, Giardino, Giudice, Sala, Salvio and Strumia

Running from Weak Scale to SUSY Scale



Matching at the SUSY Scale

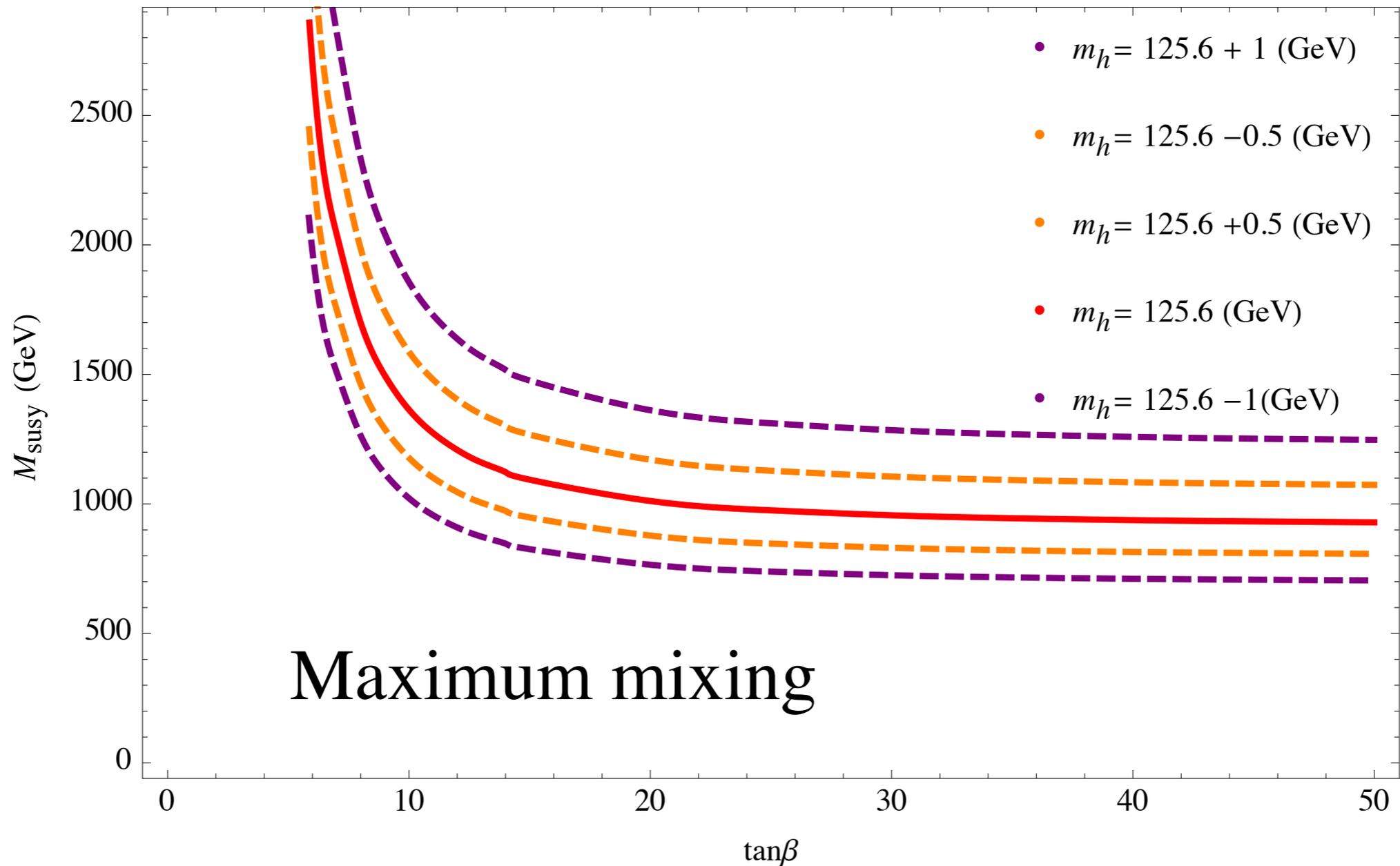
$$\lambda_{SUSY}^0 = \frac{(g^2 + g'^2) \cos 2\beta}{8}$$

$$V_{eff}(h) \equiv V_0(h) + V_{loop}(h) = V_0(h) + V_1(h) + V_2(h) + \dots$$

$$\lambda_{SUSY} = \lambda_{SUSY}^0 + d\lambda^1 + d\lambda^2$$

$$\begin{aligned}
 d\lambda &= g_Z^2 c_{2\beta}^2 \Delta Z_h^{SUSY} + 6\kappa \left[\zeta_{m1}^2 \ln \frac{m_{\tilde{t}_1}^2}{Q^2} + \zeta_{m2}^2 \ln \frac{m_{\tilde{t}_2}^2}{Q^2} + \eta_m B(\tilde{t}_1, \tilde{t}_2) \right] + 6\kappa \left(\frac{g^2}{4} + \frac{g'^2}{12} \right)^2 c_{2\beta}^2 \ln \frac{m_Q^2}{Q^2} \\
 &+ \kappa \frac{9g^4 + g'^4}{6} c_{2\beta}^2 \ln \frac{m_{Q_{12}}^2}{Q^2} + \kappa \frac{g'^4}{6} c_{2\beta}^2 \left(2 \ln \frac{m_{D_{1,2}}^2}{Q^2} + \ln \frac{m_D^2}{Q^2} + 8 \ln \frac{m_{U_{1,2}}^2}{Q^2} \right) \\
 &+ \kappa \frac{3(g^4 + g'^4)}{4} c_{2\beta}^2 \ln \frac{m_L^2}{Q^2} + \kappa \frac{3g'^4}{2} c_{2\beta}^2 \ln \frac{m_E^2}{Q^2} \\
 &+ \frac{\kappa}{8} \left[\left\{ 4g^4 - 4g^2 g_Z^2 c_{2\beta}^2 + g_Z^4 \left(2 + \frac{3}{2} s_{4\beta}^2 \right) \right\} \ln \frac{m_A^2}{Q^2} - 3g_Z^4 s_{4\beta}^2 \right] + d\lambda_{\tilde{C}_i}^{(1)} + d\lambda_{\tilde{N}_i}^{(1)} - \kappa \left(g^4 + \frac{g_Z^4}{2} \right) \\
 &+ \kappa^2 36 h_t^6 s_\beta^6 X_t^2 B^{(1)}(\tilde{t}_1, \tilde{t}_2) \left[2 \ln \frac{m_t^{T,2}}{Q^2} + \frac{16}{3} - \left(1 - \frac{X_t^2}{2s} \right)^2 \ln \frac{m_{\tilde{t}_1}^2}{Q^2} - \left(1 + \frac{X_t^2}{2s} \right)^2 \ln \frac{m_{\tilde{t}_2}^2}{Q^2} - \frac{X_t^4}{2s^2} B(\tilde{t}_1, \tilde{t}_2) \right] \\
 &+ d\lambda_{Yukawa}^{(2)} + d\lambda_{QCD}^{(2)} + \kappa^2 24 \hat{y}_t^4 \left[(g_3^2 c_{y_t} + \hat{y}_t^2 c'_{y_t}) \left(2 \ln \frac{m_t^{T,2}}{Q^2} + \frac{16}{3} \right) + g_3^2 c_{y_t} + \hat{y}_t^2 (c'_{y_t} + c_v) \right] \quad (65)
 \end{aligned}$$

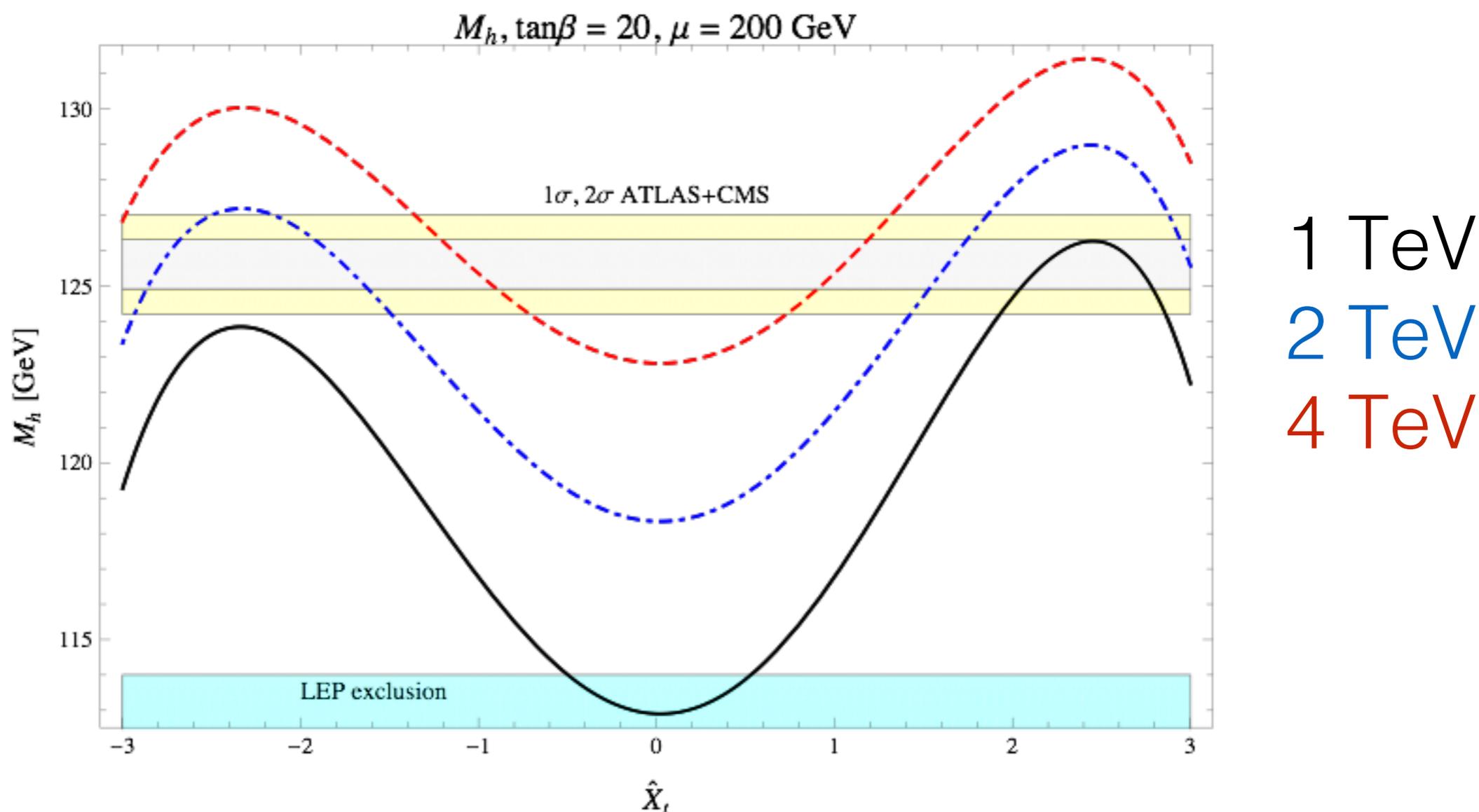
Results



Maximum mixing, large $\tan\beta$, stop mass ~ 1 TeV

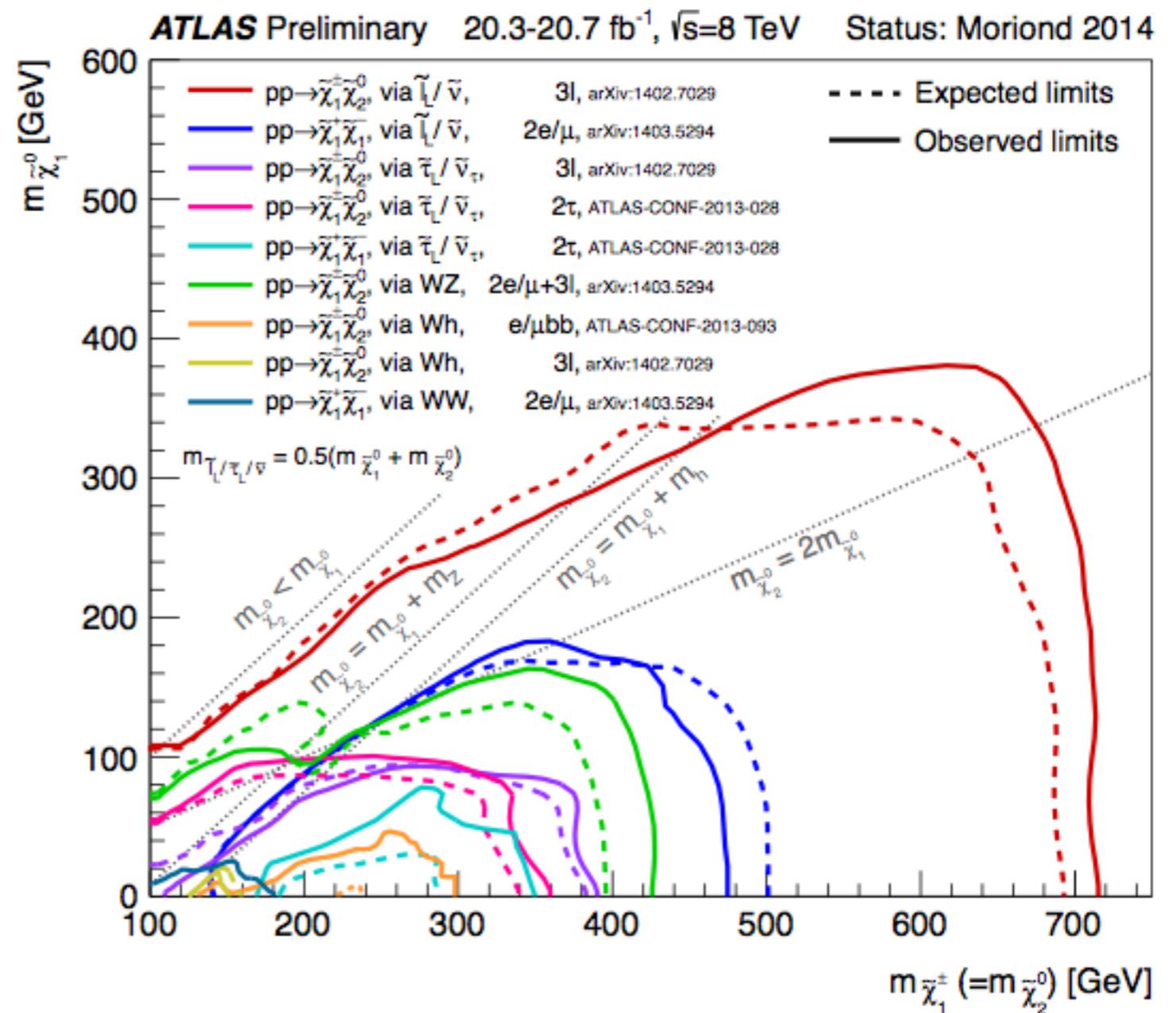
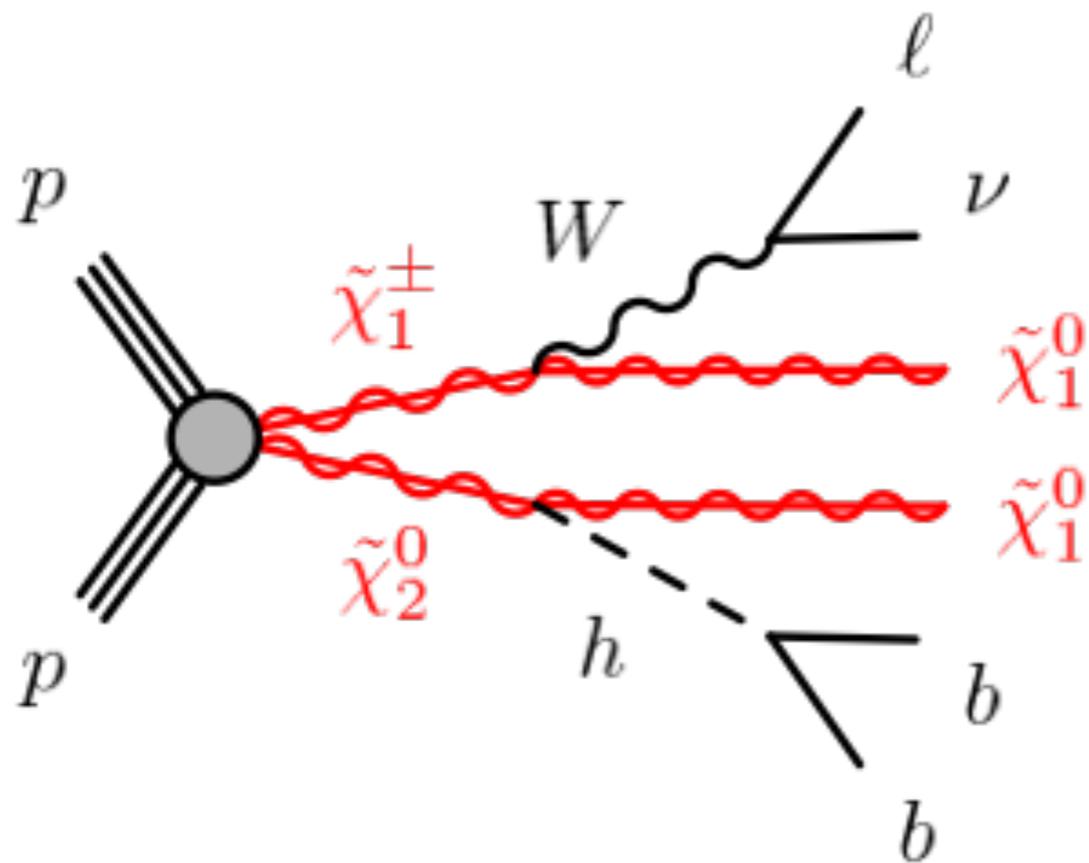
Compare with EFT Analysis

Draper, Lee and Wagner



for large $\tan\beta$, stop mass around 1TeV

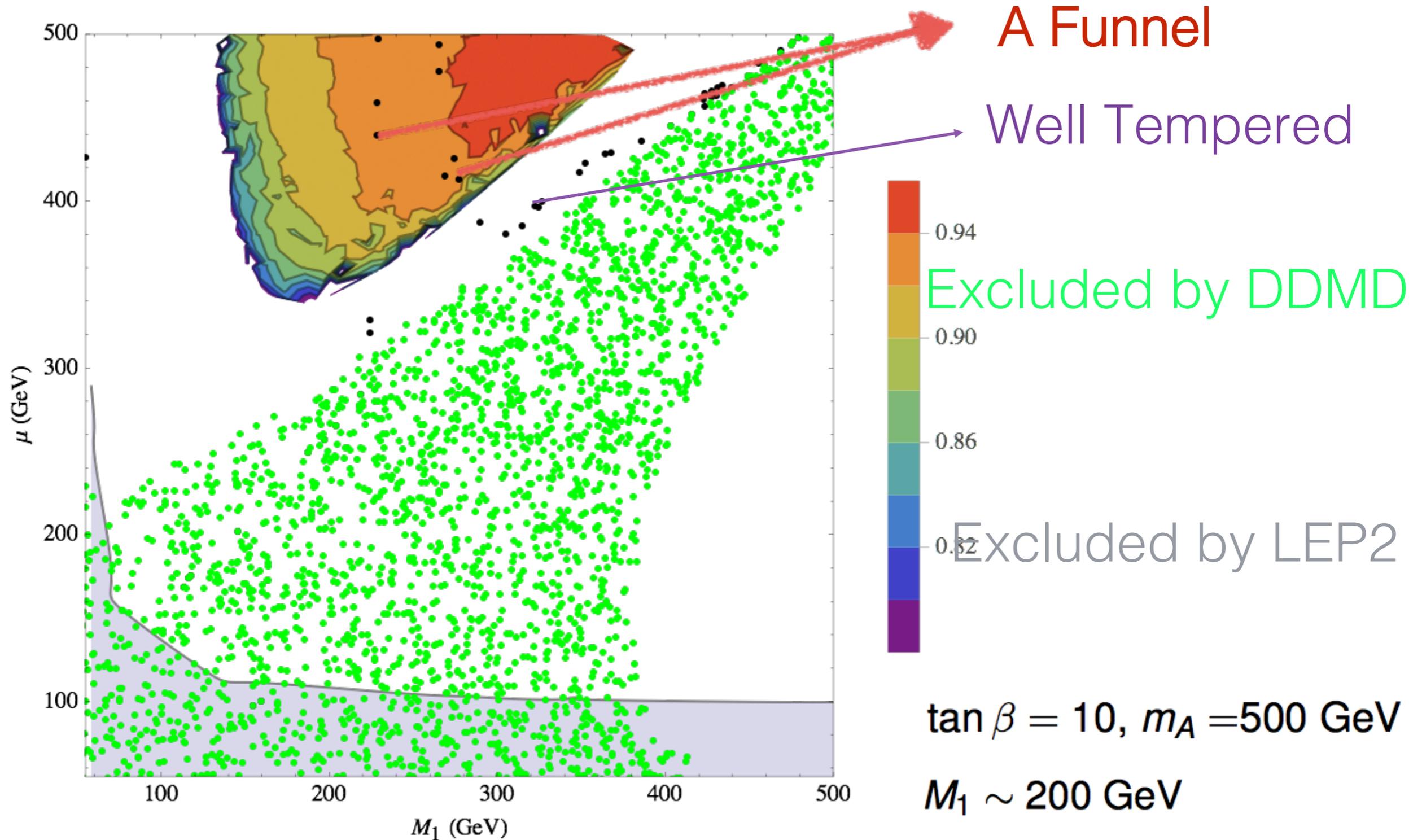
Search for SUSY with a Higgs in the Final State



Parameter Space

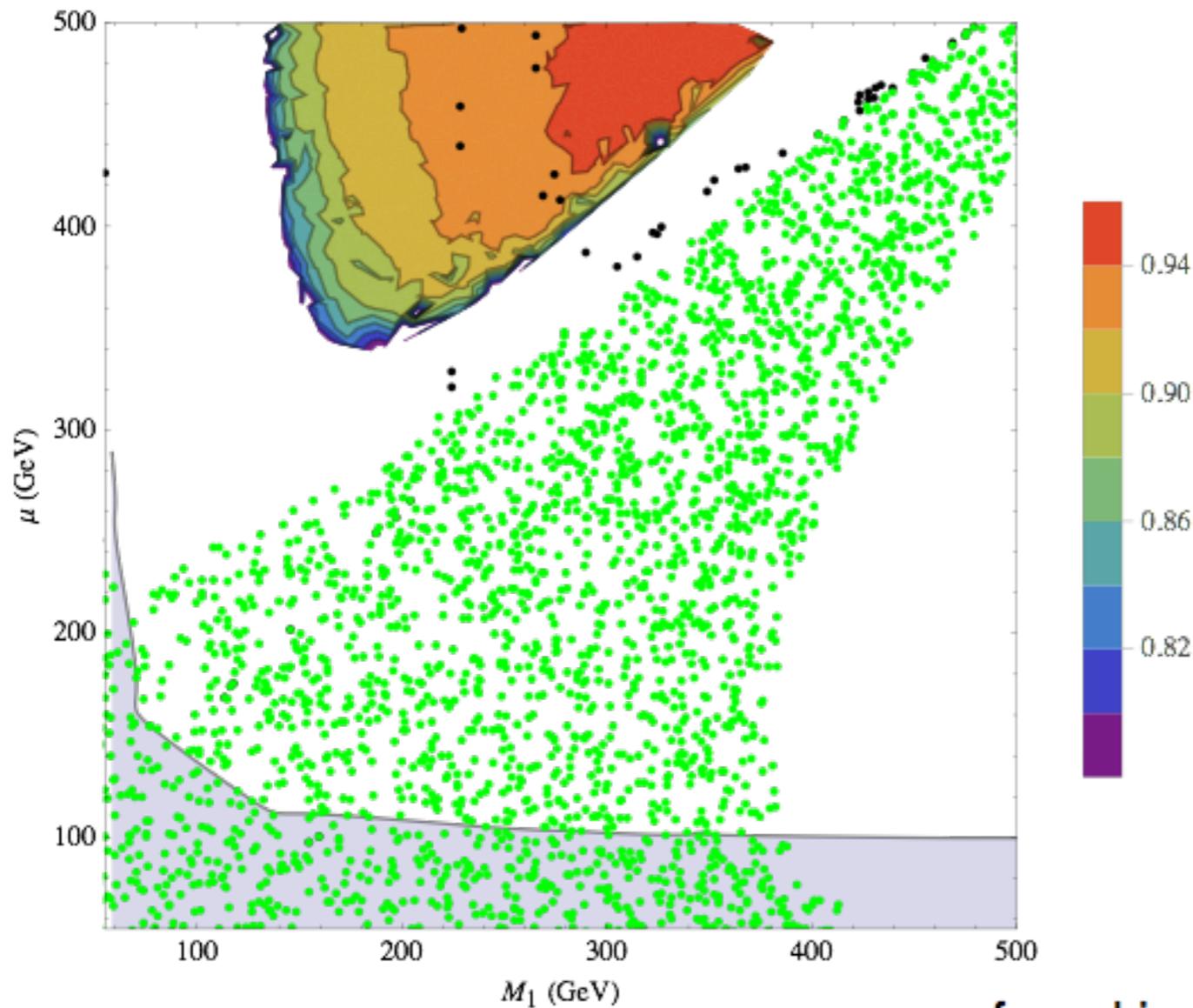
- right relic density : $\Omega h^2 = 0.1138 \pm 0.0045$
- allowed by the direct detection experiments(LUX)
- allowed by CMS $H, A \rightarrow \tau\tau$ searches
- higgs is the dominant decay mode for $\tilde{\chi}_i^0$ and/or $\tilde{\chi}_i^\pm$
- $m_{\tilde{\chi}_1^\pm} > 103.5$ GeV, from LEP

Right Relic Density, Black DDMD, Green



Higgs in Neutralino and Chargino Decays, **Contours**

$$L = -D_\mu H_u^+ D^\mu H_u - i\bar{\tilde{H}}_u \not{D}\tilde{H}_u - \sqrt{2}g'Y_{H_u}\tilde{B}\tilde{H}_u H_u^* - \sqrt{2}g\tilde{W}^a\tilde{H}_u t^a H_u^* + (u \leftrightarrow d)$$



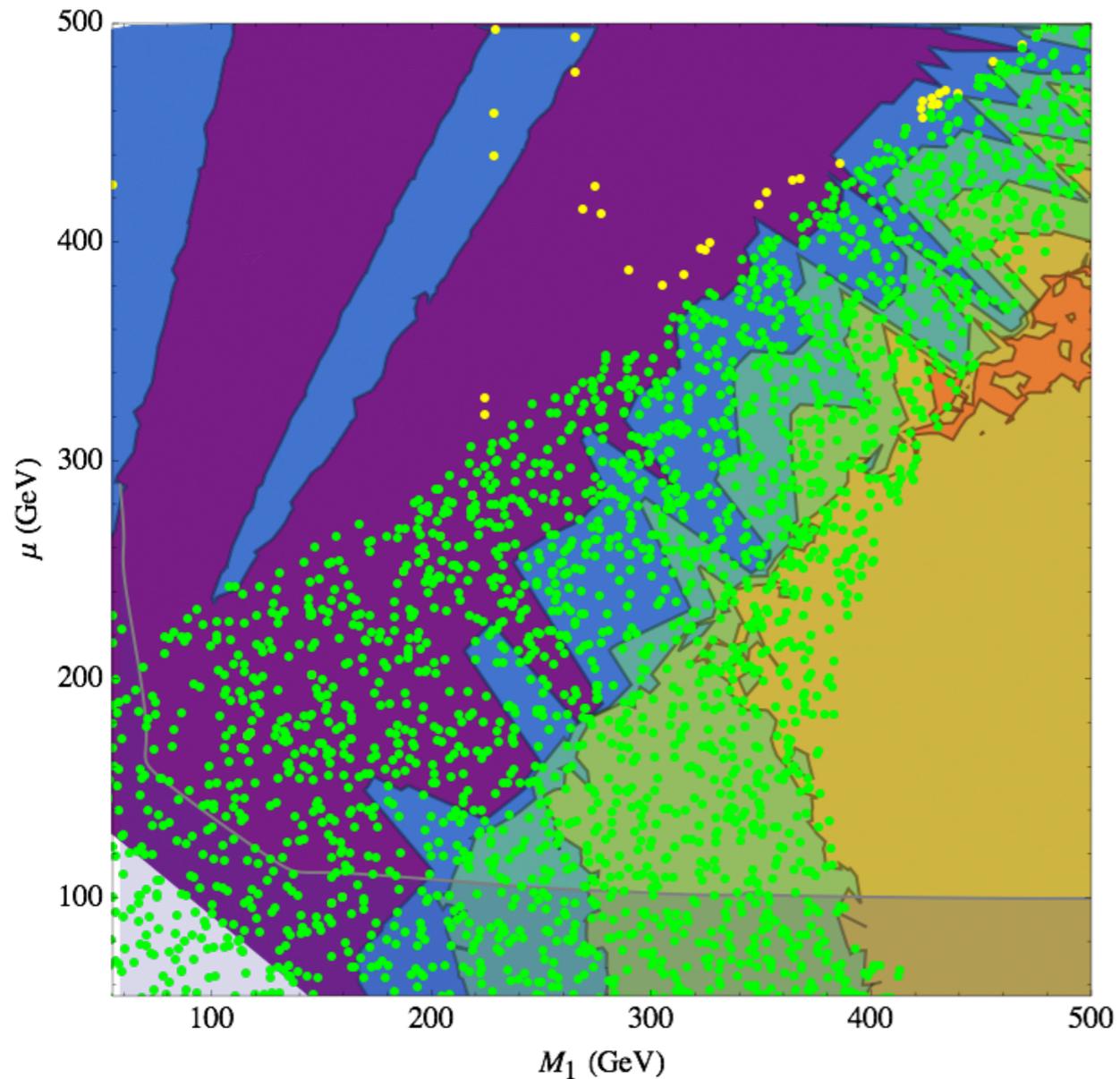
$$\tilde{\chi} = N_{i1} \tilde{B} + N_{i2} \tilde{W} + N_{i3} \tilde{H}_d + N_{i4} \tilde{H}_u$$

$$\tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 + h$$

from kinematics $M_2 - M_1 \approx M_1 > m_h$ and $\mu - M_1 > m_h$

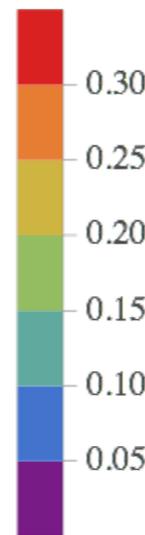
Higgs in Neutralino and Chargino Decays, **C**ontours

$$L = -D_\mu H_u^+ D^\mu H_u - i\bar{\tilde{H}}_u \not{D}\tilde{H}_u - \sqrt{2}g'Y_{H_u}\tilde{B}\tilde{H}_u H_u^* - \sqrt{2}g\tilde{W}^a\tilde{H}_u t^a H_u^* + (u \leftrightarrow d)$$



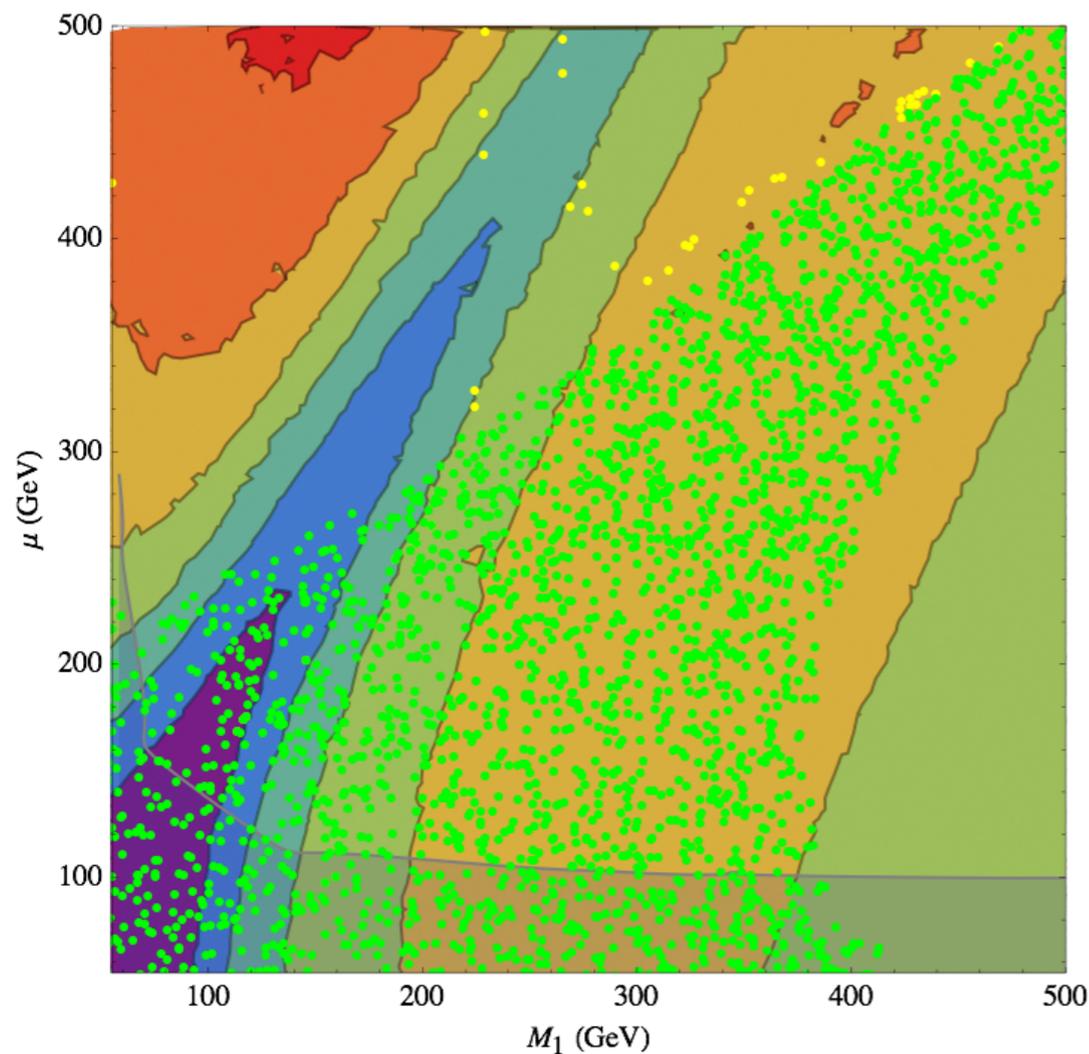
$$\tilde{\chi} = N_{i1} \tilde{B} + N_{i2} \tilde{W} + N_{i3} \tilde{H}_d + N_{i4} \tilde{H}_u$$

$$\tilde{\chi}_3^0 \rightarrow \tilde{\chi}_1^0 + h$$

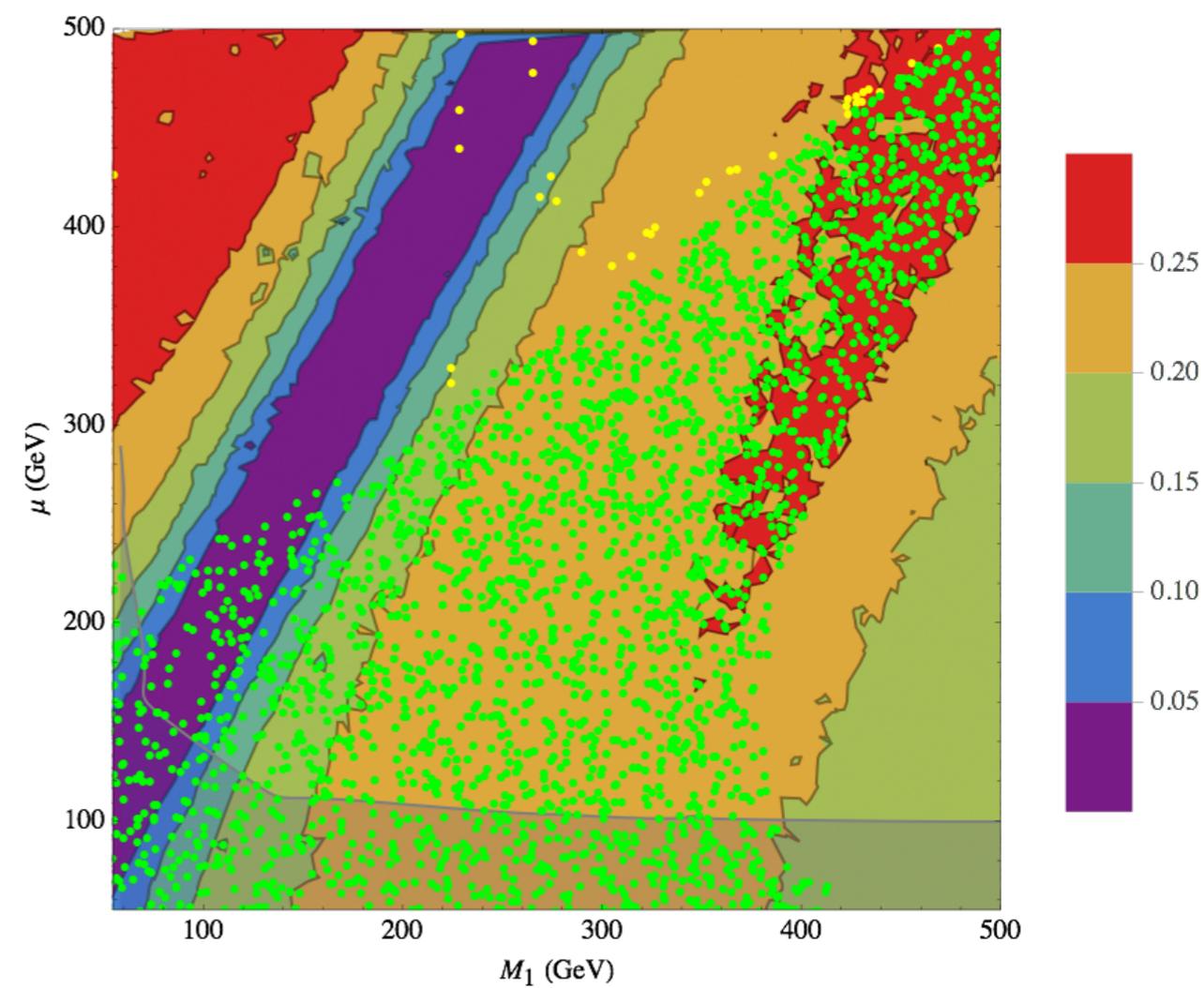


Higgs in Neutralino and Chargino Decays, Winos

$$L = -D_\mu H_u^+ D^\mu H_u - i\bar{\tilde{H}}_u \not{D} \tilde{H}_u - \sqrt{2}g' Y_{H_u} \tilde{B} \tilde{H}_u H_u^* - \sqrt{2}g \tilde{W}^a \tilde{H}_u t^a H_u^* + (u \leftrightarrow d)$$

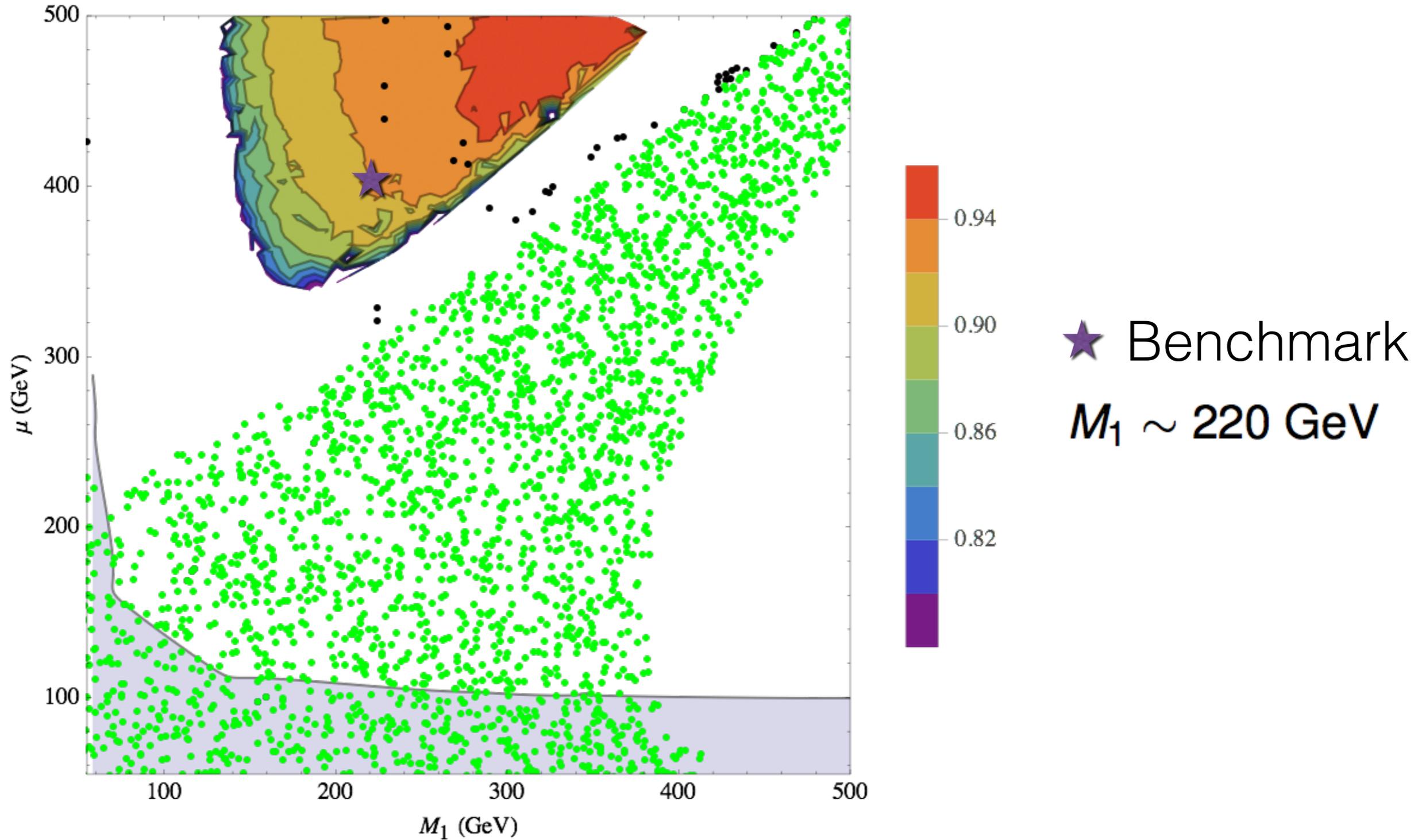


$$\tilde{\chi}_4^0 \rightarrow \tilde{\chi}_i^0 + h$$

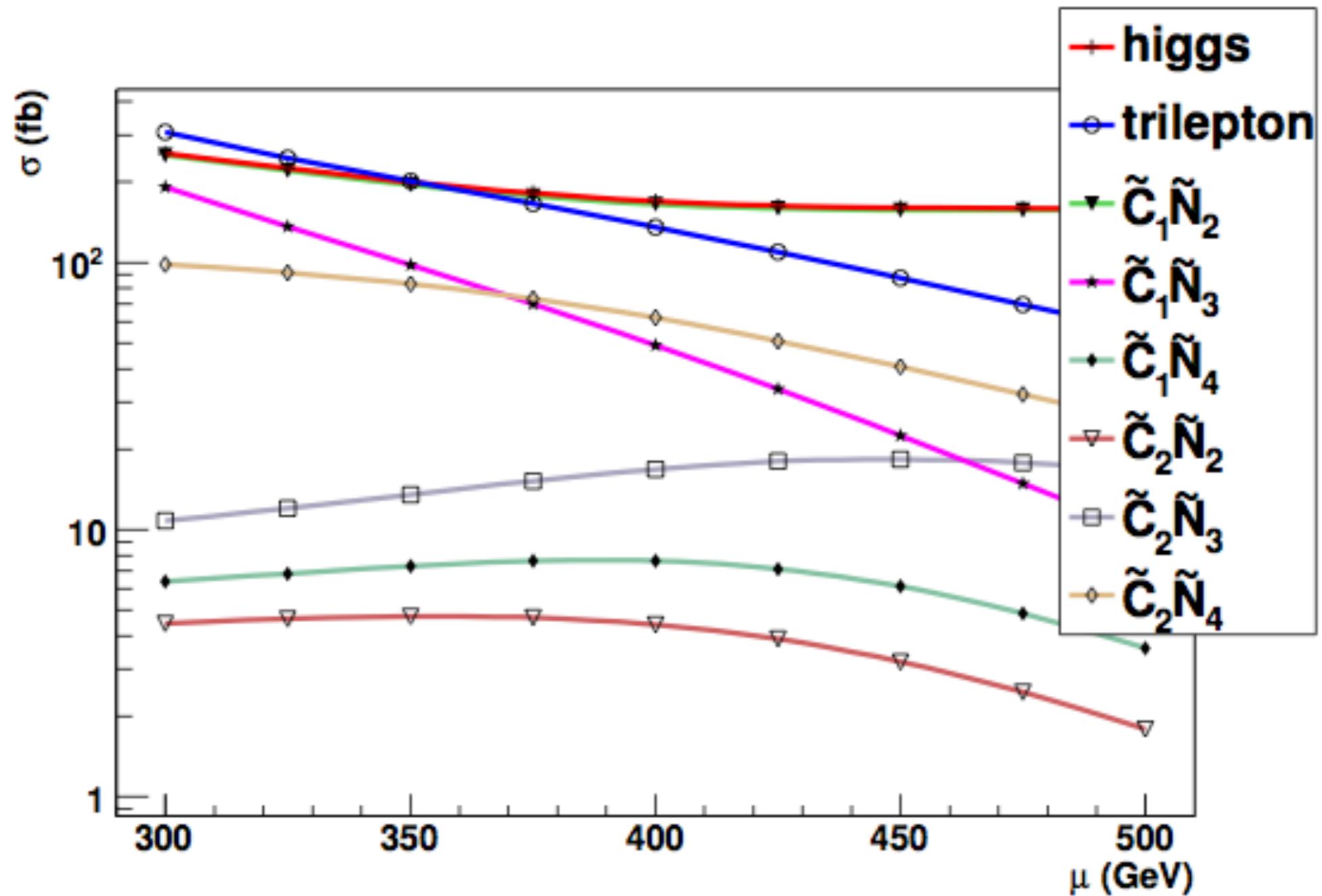


$$\tilde{\chi}_2^\pm \rightarrow \tilde{\chi}_1^\pm + h$$

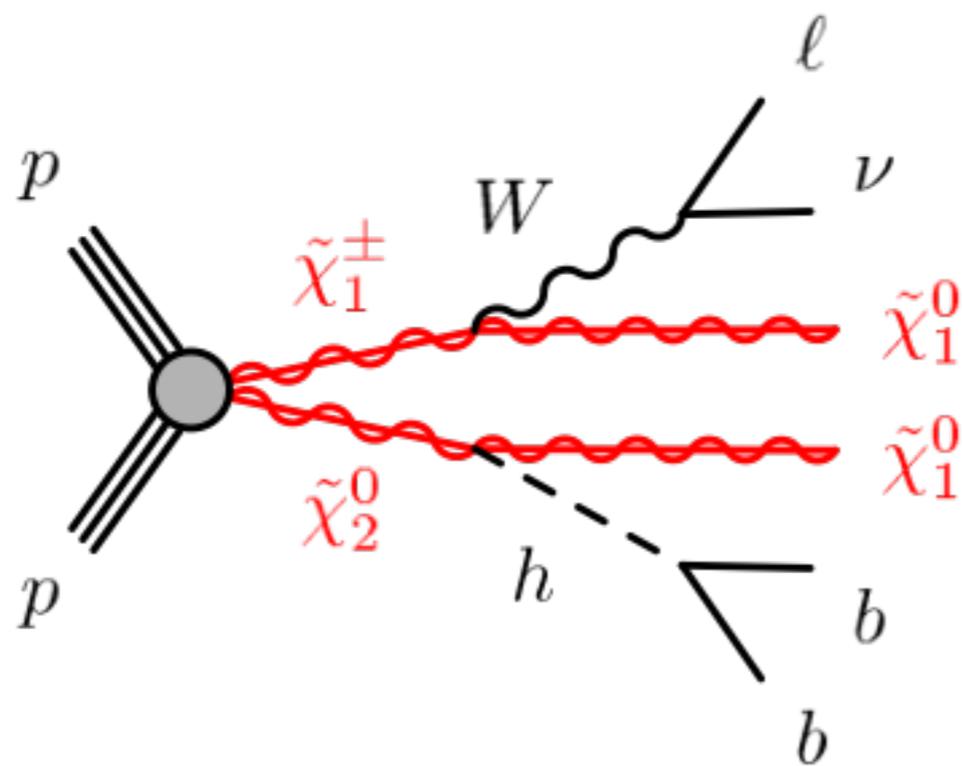
Collider Searches



Collider Searches



Collider Searches



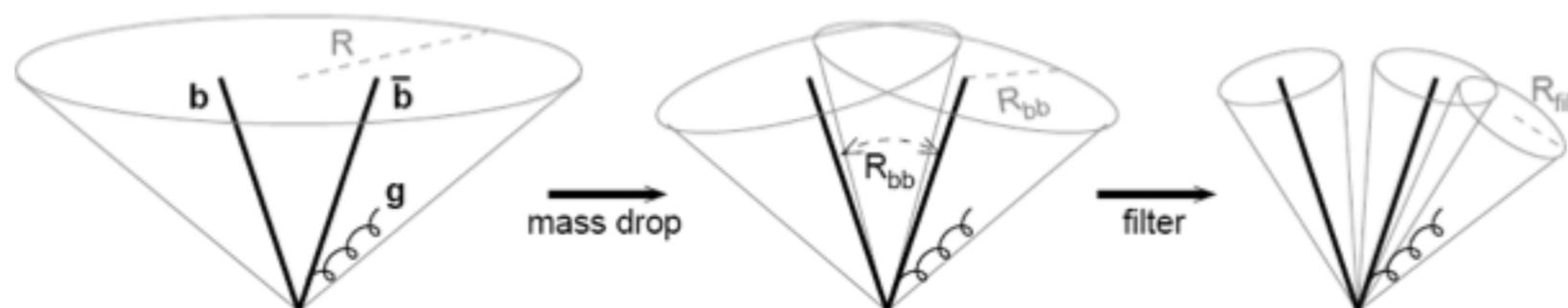
- Higgs can be boosted
- Look for 1 lepton, 1 fat jet that can be reconstructed as a higgs and missing energy

Reconstruct the Higgs

- Use the Cambridge-Aachen jet algorithm
 - ❖ Recombines the closest pair of objects in the events up to R
- For a boosted fat jet
 - ❖ Undo the last clustering step
 - ❖ Check “mass drop” and “subjettiness”

$$\mu = \frac{m_{j1}}{m_j} \quad y = \frac{\min(p_{T1}, p_{T2}) \Delta R_{12}^2}{m_j^2}$$

- ❖ If there are significant mass drop and symmetric subjet, recluster the remaining constituents with a smaller R (jet filter)



Collider Searches

- 1 isolated lepton, 1 fat jet reconstructed as a higgs($m_j = 125 \pm 4$ GeV) and $\cancel{E}_T > 100$ GeV
- main background: Wjj , dijet, $t\bar{t}$

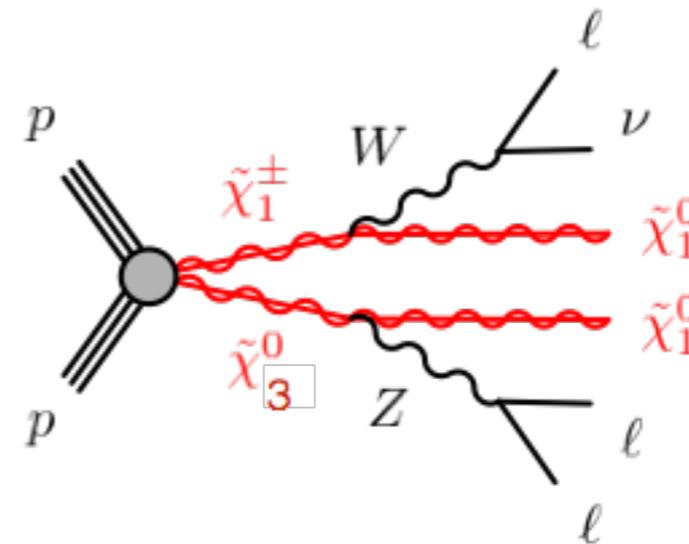
S	$Wjj+\text{dijet}$	$t\bar{t}$	B	S/\sqrt{B}
27.8	10	88.5	99	2.8

100 fb⁻¹, 14 TeV

Tri-lepton Channel



ATLAS NOTE
ATLAS-CONF-2013-035
March 24, 2013



Search for direct production of charginos and neutralinos in events with three leptons and missing transverse momentum in 21 fb^{-1} of pp collisions at $\sqrt{s} = 8 \text{ TeV}$ with the ATLAS detector

Selection	SRnoZa	SRnoZb	SRnoZc	SRZa	SRZb	SRZc
m_{SFOS} [GeV]	<60	60–81.2	<81.2 or >101.2	81.2–101.2	81.2–101.2	81.2–101.2
$E_{\text{T}}^{\text{miss}}$ [GeV]	>50	>75	>75	75–120	75–120	>120
m_{T} [GeV]	–	–	>110	<110	>110	>110
$p_{\text{T}}^{3^{\text{rd}} \ell}$ [GeV]	>10	>10	>30	>10	>10	>10
SR veto	SRnoZc	SRnoZc	–	–	–	–

main background: WZ, ZZ, $t\bar{t}V$, tri-boson

safe at 8 TeV, 20.7 fb^{-1} , $S/\sqrt{B} \sim 1.7$ for 14 TeV 100 fb^{-1}

Conclusion

- Higgs mass is known to ~ 0.4 GeV level, which motivates higgs mass calculation to a higher order
- Run the quartic coupling up, and match to MSSM at SUSY scale, to predict the SUSY scale for given parameters
- New search channel for LHC: neutralino/chargino production with a boosted higgs in the decay chain