

Weak interactions of kaons from lattice QCD

BEACH 2014

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RBC and UKQCD Collaborations

Outline

- Lattice QCD in 2014
- First order electroweak:
 - f_K/f_π , $Kl3$
 - $K \rightarrow \pi \pi$ decay
- Second order electroweak
 - $K_L - K_S$ mass difference
 - Long distance parts of ε_K .
 - Rare kaon decays

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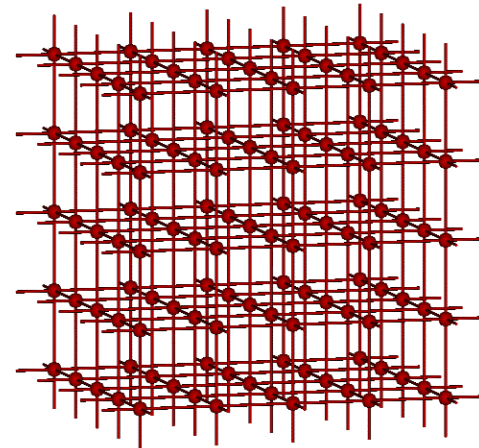
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Lattice QCD

2014

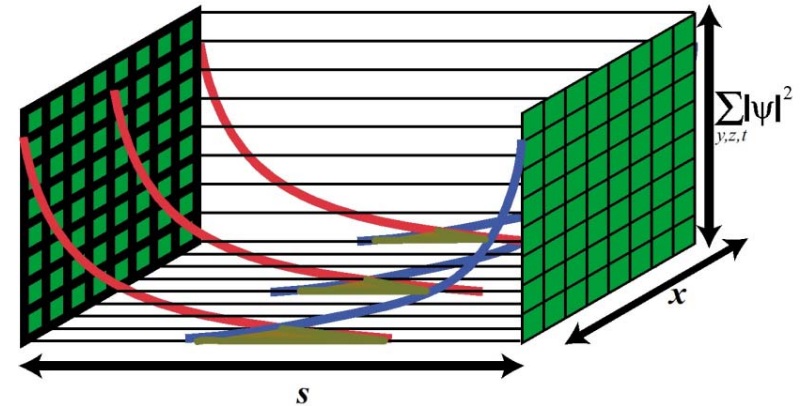
Lattice QCD

- First-principles treatment of low-energy, non-perturbative QCD.
- All approximations understood and controlled:
 - Non-zero lattice spacing: $a \rightarrow 0$.
 - Finite volume: $L \rightarrow \infty$
 - Typically neglect E&M and $m_u \neq m_d$,
 $\alpha_{\text{EM}} \ll 1$
- Supports not only rough phenomenology but also accurate theoretical physics (where it can be applied).



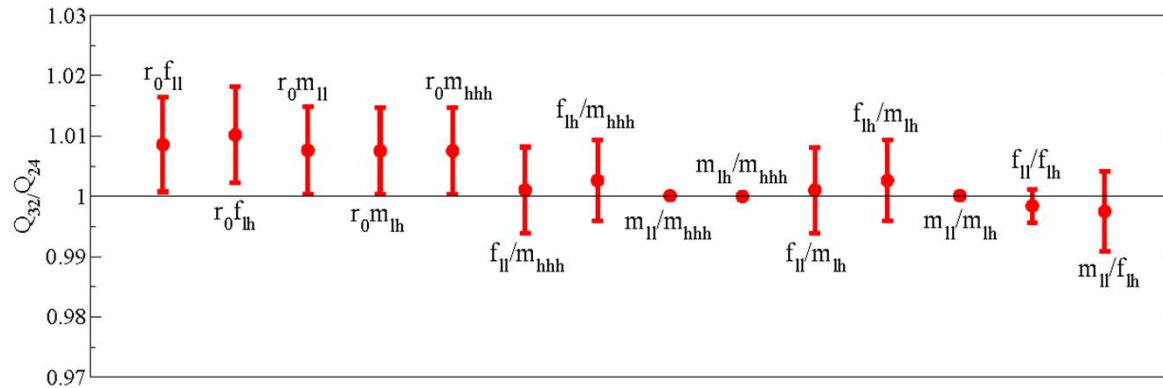
Domain wall fermions

- Use 5-dimensional Wilson fermion action with open BC on $s = 0$ and $s = L_s - 1$ slice.
- 4-D chiral bound states form on the $s = 0$ and $L_s - 1$ walls.
- 5-D propagating states are large-action, lattice artifacts.
- 4-D states disappear as $p \rightarrow \pi/a$, solving the doubling problem.
- Accurate chiral symmetry **at all energies**, broken by left-right mixing: residual chiral symmetry breaking.



DWF at low energy

- At low energy $E \ll 1/a$, 5-D DWF theory looks like a chiral 4-D theory (QCD) with small chiral asymmetry:
 - Leading, dim-3 operator: $m_{\text{res}} \bar{q} q$ (mass term)
 - Next leading dim-5 operator: $m_{\text{res}} \bar{q} \sigma^{\nu\nu} F^{\nu\nu} q$ (clover term)
- Very small discretization errors:

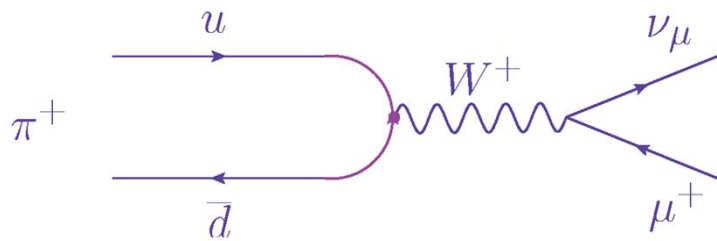


Ratios of dimensionless combinations of physical quantities
computed using $1/a = 1.73$ and 2.28 GeV.

Current state-of-the-art

- Physical $m_\pi=135$ MeV and $L = 4 - 6$ fm.
- Generate $48^3 \times 96$ and $64^3 \times 128$ ensembles.
- Complete set of measurements takes 5.3 hours on a 32-rack BG/Q machine (**sustains 1 Pflops**)
- Large collaboration essential:
 - Highly optimized code (64 threads, SPI comms., wide-vector FP)
 - Sophisticated algorithms (deflation, FG $(\Delta t)^3$ integrator)
 - Complex measurement strategies (NPR, G-parity BC, 4-pt functions, all-mode-averaging, all-to-all propagators)

Simple example: f_π



$$\langle 0 | \bar{d} \gamma^5 \gamma^\mu u | \pi^+(\vec{p}) \rangle = f_\pi \frac{p^\mu}{\sqrt{4E_\pi(\vec{p})}}$$

$$f_\pi = N \sum_{\vec{r}} \frac{\langle A^0(\vec{r}, t) O_\pi(t=0) \rangle}{\langle O_\pi^\dagger(t) O_\pi(t=0) \rangle^{\frac{1}{2}}} e^{m_\pi t/2}$$

- 2012 (elaborate chiral fit): $f_\pi = 127(3)_{\text{stat}}(3)_{\text{sys}} \text{ MeV}$
- 2013 ($m_\pi=135 \text{ MeV}$): $f_\pi = 130.0(0.3)_{\text{stat}} \text{ MeV}$ (40 configs.)
- Experiment: $f_\pi = 130.4(0.04)(0.2) \text{ MeV}$

Semi-leptonic decays

- f_K / f_π

- RBC/UKQCD (prelim.)

$$f_K / f_\pi = 1.199(4)_{\text{stat}}(4)_{\text{chi}}(2)_{\text{FV}}$$

- FNAL/MILC (arXiv:1407.3772)

$$f_K / f_\pi = 1.1956(10)_{\text{stat}}(+23/-14)_{\text{cont}}$$

$$(10)_{\text{FV}}(5)_{\text{EM}}$$

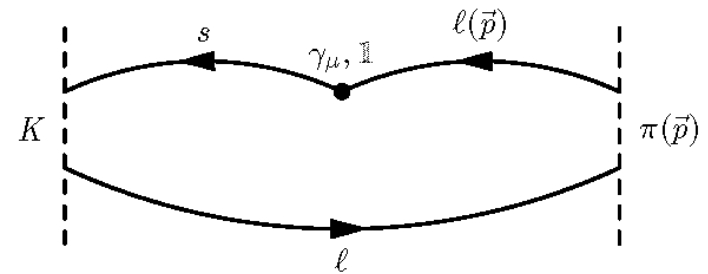
- $Kl3: f_+$

- FNAL/MILC (PRL 112 (2014) 112001)

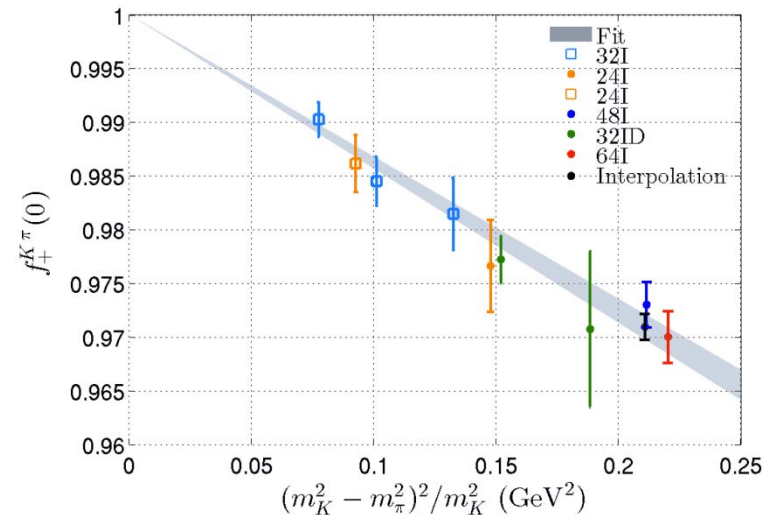
$$f_+(0) = 0.9704(24)_{\text{chi-cont}}(22)_{\text{sys}}$$

- ETMc, (preliminary)

$$f_+(0) = 0.9683(50)_{\text{stat+fit}}(42)_{\text{chiral}}$$



RBC/UKQCD

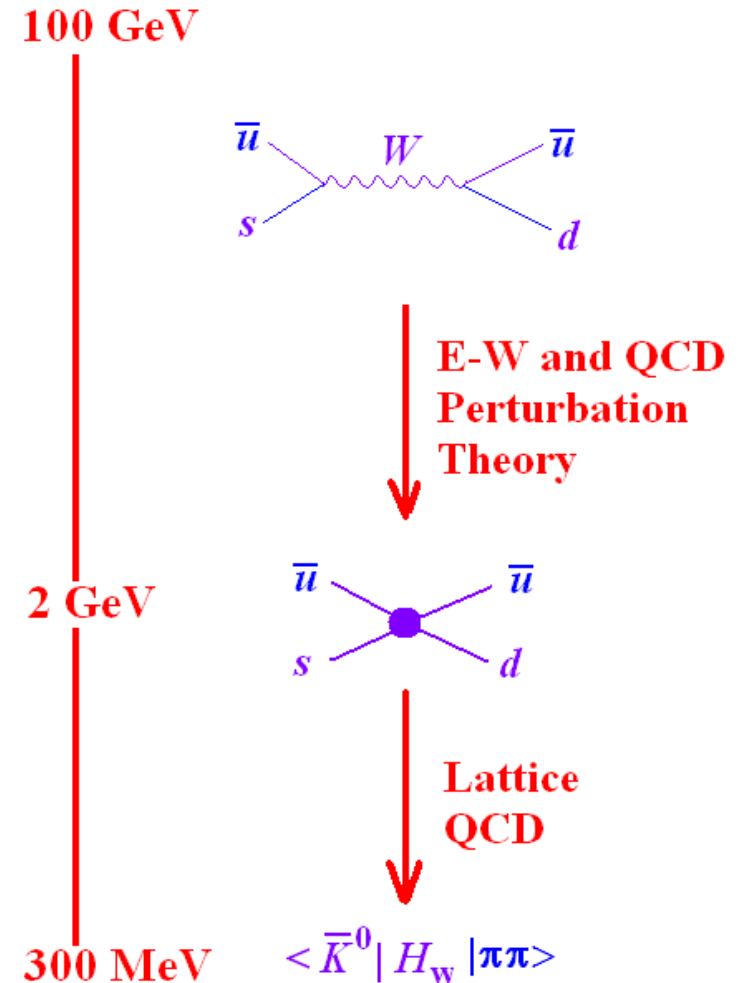


$\Delta S=1$ Weak Interactions

- Represent weak interactions by local four-quark Lagrangian

$$\mathcal{H}^{(\Delta S=1)} = \frac{G_F}{\sqrt{2}} V_{ud} V_{us}^* \left\{ \sum_{i=1}^{10} \left[z_i(\mu) - \frac{V_{td} V_{ts}^*}{V_{ud} V_{us}^*} y_i(\mu) \right] Q_i \right\}$$

- $V_{qq'}$ – CKM matrix elements
- z_i and y_i – Wilson Coefficients
- Q_i – four-quark operators



Four quark operators

- **Current-current operators**

$$Q_1 \equiv (\bar{s}_\alpha d_\alpha)_{V-A} (\bar{u}_\beta u_\beta)_{V-A}$$

$$Q_2 \equiv (\bar{s}_\alpha d_\beta)_{V-A} (\bar{u}_\beta u_\alpha)_{V-A}$$

- **QCD Penguins**

$$Q_3 \equiv (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\beta)_{V-A}$$

$$Q_4 \equiv (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\alpha)_{V-A}$$

$$Q_5 \equiv (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\beta)_{V+A}$$

$$Q_6 \equiv (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} (\bar{q}_\beta q_\alpha)_{V+A}$$

- **Electro-Weak Penguins**

$$Q_7 \equiv \frac{3}{2} (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\beta)_{V+A}$$

$$Q_8 \equiv \frac{3}{2} (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\alpha)_{V+A}$$

$$Q_9 \equiv \frac{3}{2} (\bar{s}_\alpha d_\alpha)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\beta)_{V-A}$$

$$Q_{10} \equiv \frac{3}{2} (\bar{s}_\alpha d_\beta)_{V-A} \sum_{q=u,d,s} e_q (\bar{q}_\beta q_\alpha)_{V-A}$$

$K \rightarrow \pi \pi$ decay

$K \rightarrow \pi \pi$ phenomenology

- Final $\pi\pi$ states can have $I = 0$ or 2.

$$\langle \pi\pi(I = 2) | H_w | K^0 \rangle = A_2 e^{i\delta_2} \quad \Delta I = 3/2$$

$$\langle \pi\pi(I = 0) | H_w | K^0 \rangle = A_0 e^{i\delta_0} \quad \Delta I = 1/2$$

- Direct CP violation in this decay is characterized by:

$$\epsilon' = \frac{i e^{\delta_2 - \delta_0}}{\sqrt{2}} \left| \frac{A_2}{A_0} \right| \left(\frac{\text{Im} A_2}{\text{Re} A_2} - \frac{\text{Im} A_0}{\text{Re} A_0} \right)$$

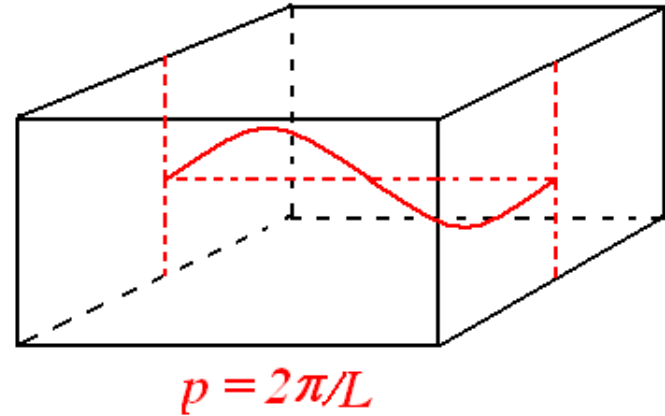
- $K^0 - \bar{K}^0$ mixing gives indirect CP violation:

$$\epsilon_K = \frac{i}{2} \left\{ \frac{\text{Im} M_{0\bar{0}} - \frac{i}{2} \text{Im} \Gamma_{0\bar{0}}}{\text{Re} M_{0\bar{0}} - \frac{i}{2} \text{Re} \Gamma_{0\bar{0}}} \right\} + i \frac{\text{Im} A_0}{\text{Re} A_0}$$

Lattice Aspects

Physical $\pi\pi$ states – Lellouch-Lüscher

- Euclidean e^{-Ht} projects onto $|\pi\pi(\vec{p}=0)\rangle$
- Exploit finite-volume quantization.
- Adjust volume so 1st or 2nd excited state has correct p .
- Impose boundary conditions so ground state has physical p
 - $\Delta I = 3/2$: impose anti-periodic BC on d quark
 - $\Delta I = 1/2$: impose G-parity BC
- Correctly include $\pi - \pi$ interactions, including normalization.



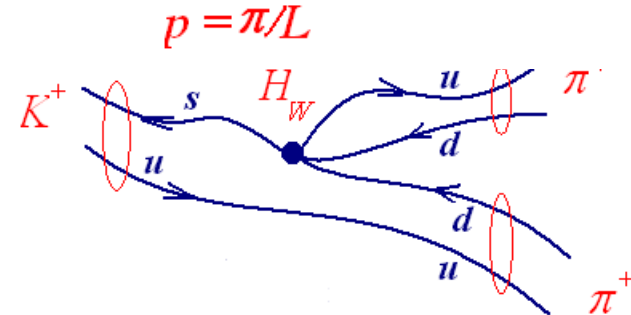
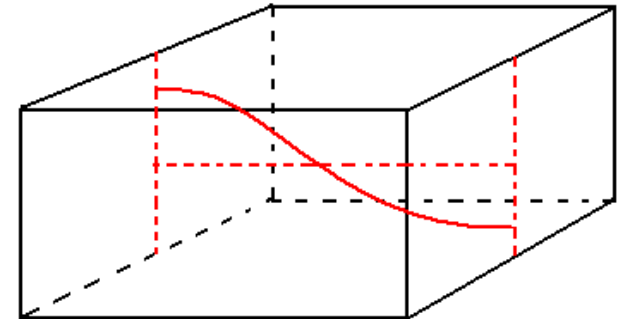
$$\Delta I = 3/2$$

$\Delta I = 3/2 \quad K \rightarrow \pi \pi$

- Three operators contribute
 $O^{(27,1)}$, $O^{(8,8)}$ and $O^{(8,8)_m}$.
- **Achieve essentially physical kinematics!**

(146 configurations)

- $m_\pi = 142.9(1.1)$ MeV
- $m_K = 511.3(3.9)$ MeV
- $E_{\pi\pi} = 492(5.5)$ MeV



- $\text{Re}(A_2) = (1.436 \pm 0.063_{\text{stat}} \pm 0.258_{\text{sys}}) 10^{-8} \text{ GeV}$

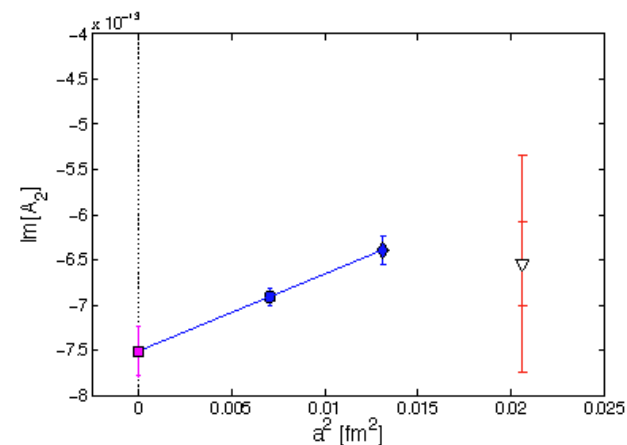
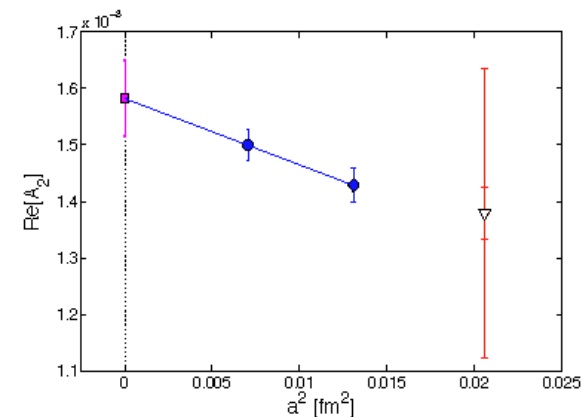
Experiment: $1.479(4) 10^{-8} \text{ GeV}$

- $\text{Im}(A_2) = -(6.29 \pm 0.46_{\text{stat}} \pm 1.20_{\text{sys}}) 10^{-13} \text{ GeV}$

$\Delta I = 3/2$: Next results

(Tadeusz Janowski)

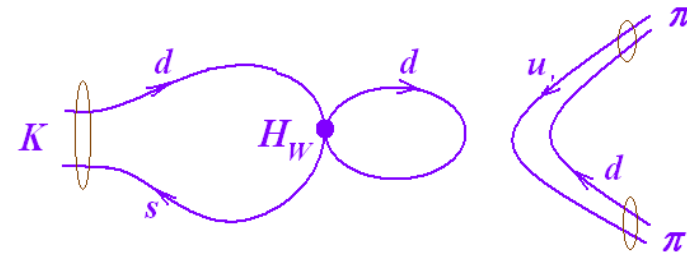
- Use two new large ensembles to remove a^2 error ($m_\pi=135$ MeV, $L=5.4$ fm)
 - $48^3 \times 96$, $1/a=1.73$ GeV
 - $64^3 \times 128$, $1/a=2.28$ GeV
- First continuum results, (**preliminary**):
 - $\text{Re}(A_2) = (1.583 \pm 0.067_{\text{stat}}) \times 10^{-8}$ GeV
 - $\text{Im}(A_2) = - (7.51 \pm 27_{\text{stat}}) \times 10^{-13}$ GeV
- Experiment: $\text{Re}(A_2) = 1.479(4) 10^{-8}$ GeV



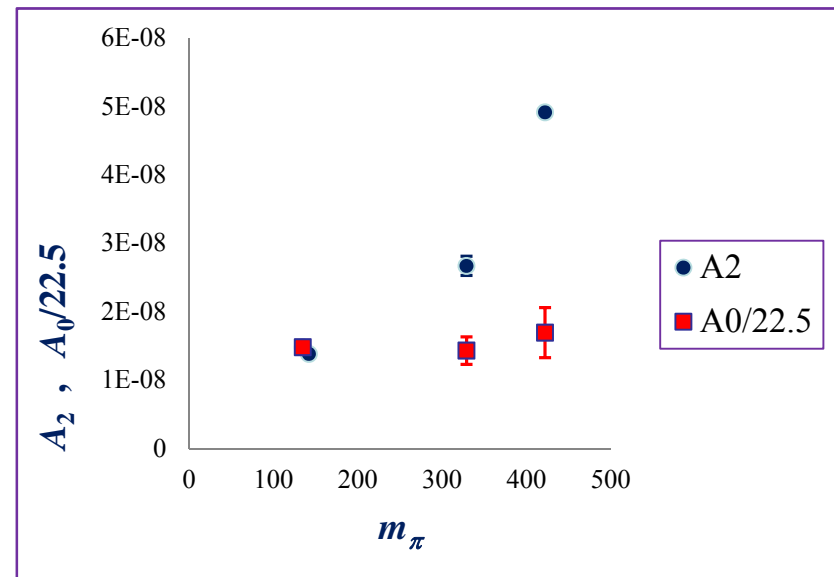
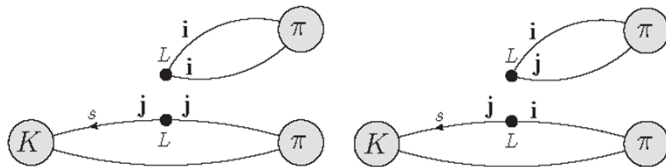
$$\Delta I = 1/2$$

$$\Delta I = 1/2 \quad K \rightarrow \pi \pi$$

- Made much more difficult by disconnected diagrams:

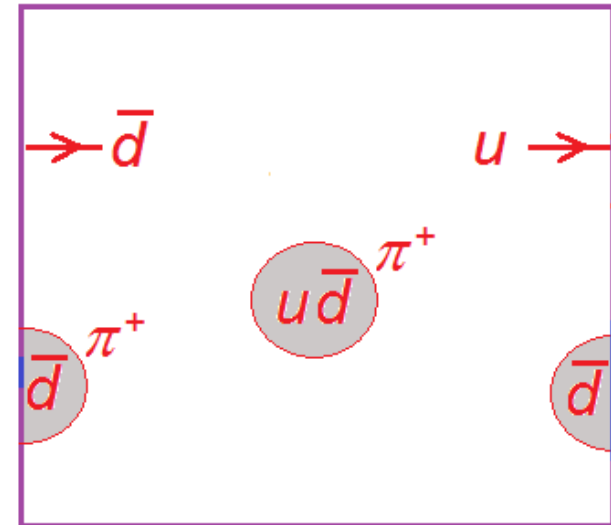


- Many more diagrams (48) than $\Delta I = 3/2$.
- Initial threshold decay calculation successful (Qi Liu)
 - $\text{Re}(A_0)$: 25% stat errors
 - $\text{Im}(A_0)$: 50% stat errors
- $\Delta I = 1/2$ rule “explained”

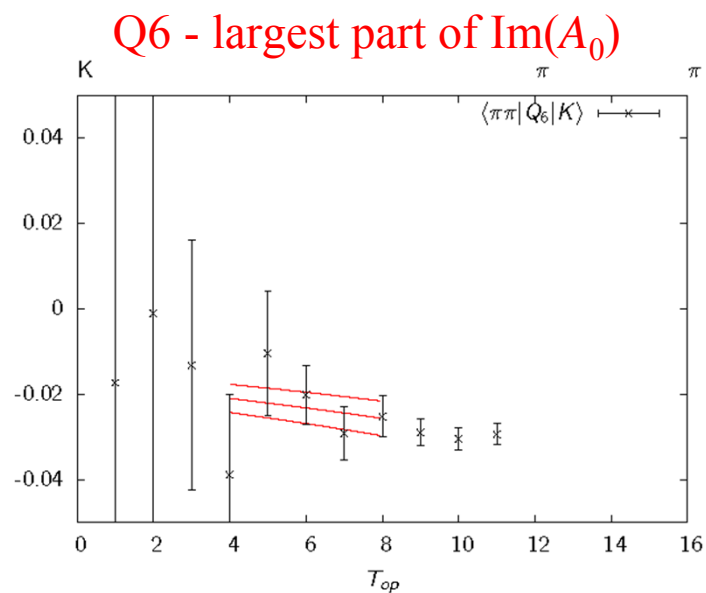
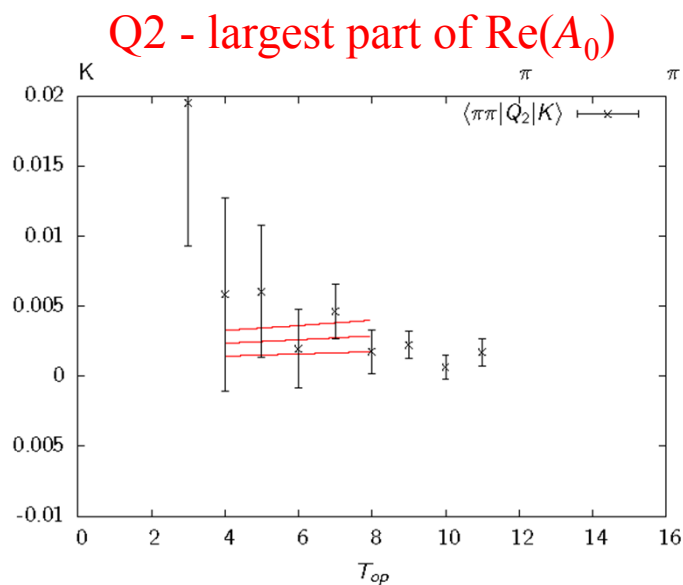


$\Delta I = 1/2$ $K \rightarrow \pi \pi$: **Physical kinematics**

- Goal is a 20% calculation of ε'/ε with all errors controlled
- Use $32^3 \times 64$ volume with $1/a = 1.37$ GeV
- Achieve $p = 205$ MeV from **G-parity** boundary conditions in 3 directions
- Requires new **G-parity** ensembles



$\Delta I = 1/2$ $K \rightarrow \pi\pi$: Current status



$$\langle \pi\pi_{I=0} | Q_2 | K \rangle = (1.92 \pm 0.75) \times 10^{-3}$$

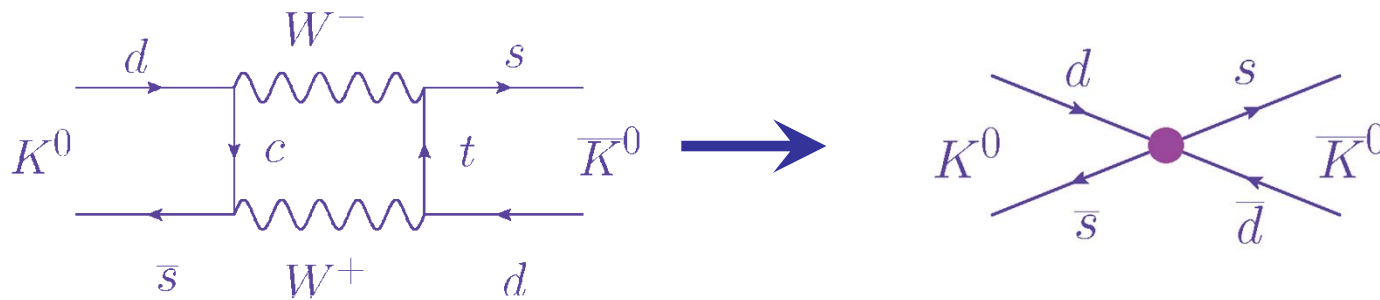
$$\langle \pi\pi_{I=0} | Q_6 | K \rangle = (-1.71 \pm 0.27) \times 10^{-3}$$

- 70 configurations
- Result expected in 1 year

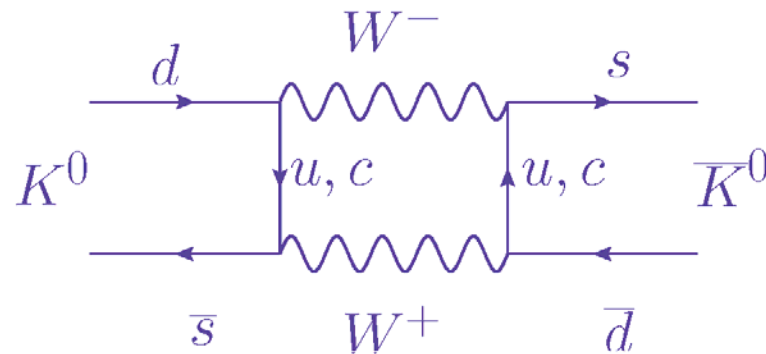
$K_L - K_S$ mass difference

$K^0 - \bar{K}^0$ Mixing

- CP violating: $p \sim m_t$ $\epsilon_K = \frac{i}{2} \left\{ \frac{\text{Im} M_{00} - \frac{i}{2} \text{Im} \Gamma_{00}}{\text{Re} M_{00} - \frac{i}{2} \text{Re} \Gamma_{00}} \right\} + i \frac{\text{Im} A_0}{\text{Re} A_0}$



- CP conserving: $p \leq m_c$ $m_{K_S} - m_{K_L} = 2\text{Re}\{M_{00}\}$

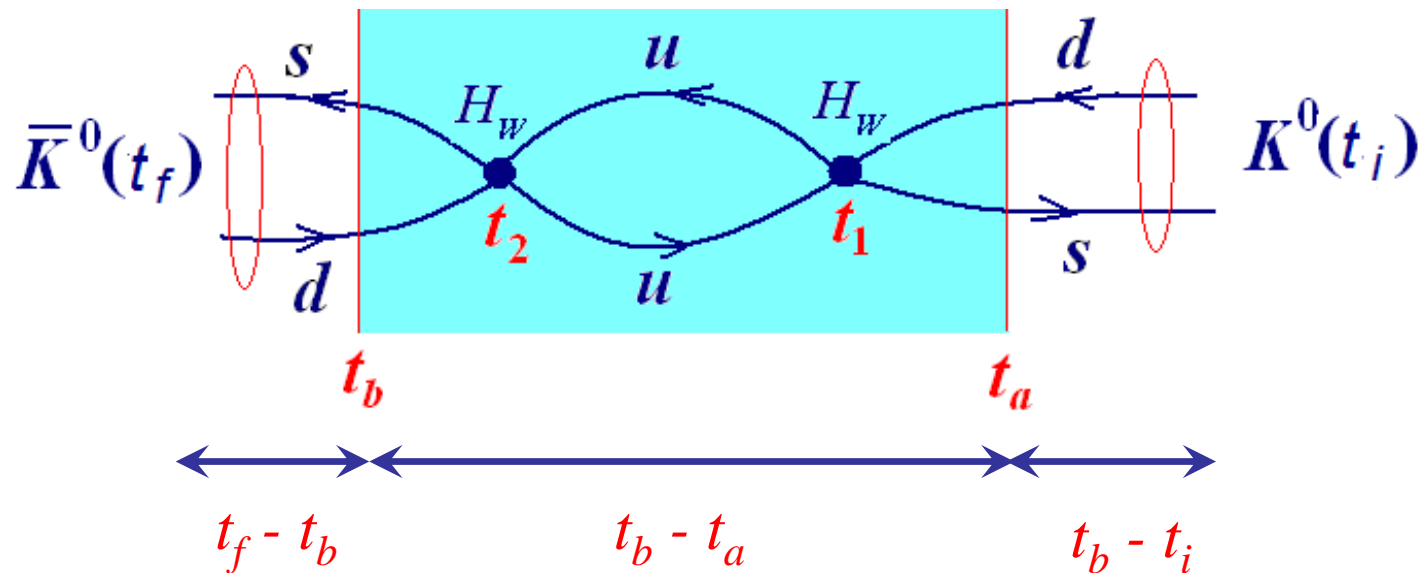


Lattice Version

(Jianglei Yu)

- Evaluate standard, Euclidean, 2nd order $K^0 - \bar{K}^0$ amplitude:

$$\mathcal{A} = \langle 0 | T \left(K^0(t_f) \frac{1}{2} \int_{t_a}^{t_b} dt_2 \int_{t_a}^{t_b} dt_1 H_W(t_2) H_W(t_1) K^0(t_i) \right) | 0 \rangle$$



Interpret Lattice Result

$$\mathcal{A} = N_K^2 e^{-M_K(t_f - t_i)} \sum_n \frac{\langle \bar{K}^0 | H_W | n \rangle \langle n | H_W | K^0 \rangle}{M_K - E_n} \left(\overset{\textcircled{1.}}{- (t_b - t_a)} - \overset{\textcircled{2.}}{\frac{1}{M_K - E_n}} + \frac{e^{(M_K - E_n)(t_b - t_a)}}{M_K - E_n} \right)$$

1. Δm_K^{FV}

2. Uninteresting constant

3. Growing or decreasing exponential:

states with $E_n < m_K$ must be removed!

- Finite volume correction:

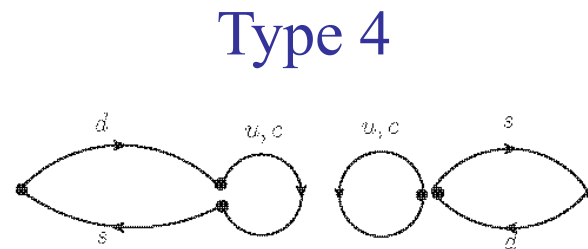
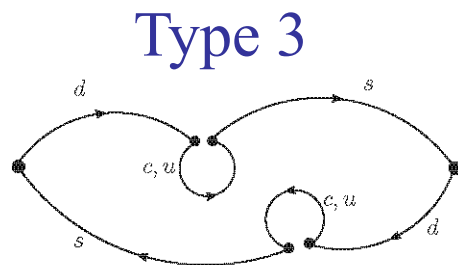
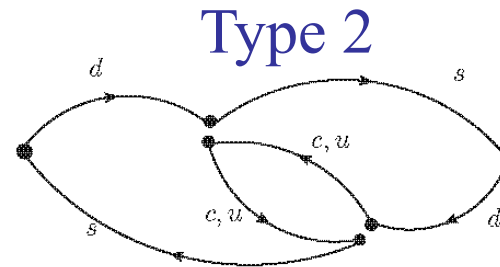
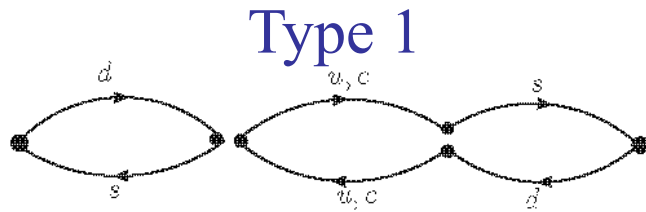
$$M_{K_L} - M_{K_S} = 2 \sum_n \frac{\langle \bar{K}^0 | H_W | n \rangle \langle n | H_W | K^0 \rangle}{M_K - E_{n_0}} - \frac{E_{n_0}^2}{2k_n M_K} \frac{d(\phi + \delta_0)}{dk} \Big|_{m_K} V |\langle n_0 | H_W | K^0 \rangle|^2 \cot(\phi + \delta_0) \Big|_{m_K}$$

(N.H. Christ, X. Feng, G. Martinelli, C.T. Sachrajda)

Lattice setup

(Jianglei Yu)

- Must include charm quark (GIM $u-d$ cancellation)
- Two calculations performed
 - $16^3 \times 32$, $m_p = 420$ MeV, types 1 & 2 (arXiv:1212.5931)
 - $24^3 \times 64$, $m_p = 330$ MeV, all graphs (arXiv:1406.0916)



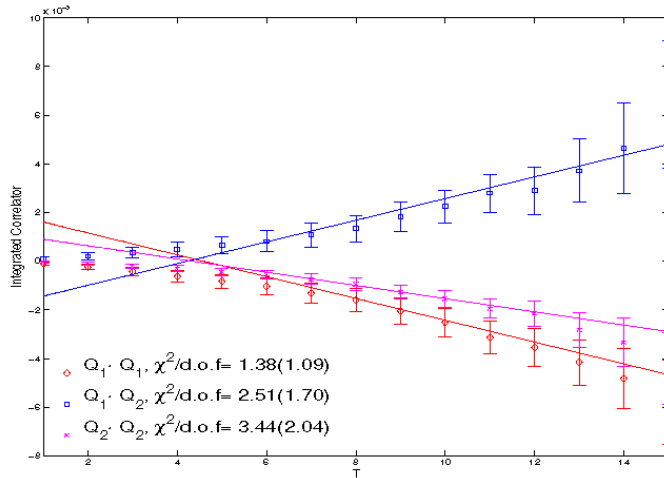
ΔM_K – results

Diagrams	$Q_1 \cdot Q_1$	$Q_1 \cdot Q_2$	$Q_2 \cdot Q_2$	ΔM_K	
Type 1,2	1.485(80)	1.567(38)	3.678(56)	6.730(96)	$\times 10^{-12}$ MeV
All	0.754(42)	-0.16(15)	2.70(18)	3.30(34)	

- Result:
 $\Delta M_K = 3.30(34) 10^{-12}$ MeV
- $\Delta M_K^{\text{expt}} = 3.483(6) 10^{-12}$ MeV
- Agreement fortuitous!
- Unphysical, $m_\pi = 330$ MeV
- Active charm but $m_c a = 0.55$

ΔM_K – Next Steps

(Ziyuan Bai)



- Begin $32^3 \times 64$ calculation
- Case study for $2M_\pi < M_K$
- Prepare for $80^2 \times 96 \times 192$, $1/a=3.0$ GeV.
- Long dist. contribution to ε_K being coded.

- $M_\pi = 170$ MeV
- $1/a = 1.35$ MeV
- 250 configurations

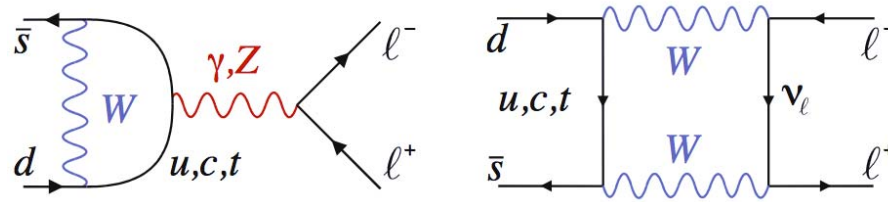
Charm mass	ΔM_K
750 MeV	4.6(13)
592 MeV	3.8(17)

I	$E_{2\pi}$ (MeV)	ΔM_K (10^{-12} MeV)
0	337.7(30)	-0.133(99)
2	343.5(25)	$-6.54(25) \times 10^{-4}$

Rare Kaon Decays

Rare Kaon Decays

(Xu Feng, Antonin Portelli)



- Can lattice methods be of use for rare K decays?
 - $K \rightarrow \pi + l + \bar{l}$ or $K \rightarrow \pi + \nu + \bar{\nu}$
 - Semi-leptonic decays with $\Delta S=1$ are second order in the standard model and hence sensitive to new physics.
- $K_L \rightarrow \pi^0 + l + \bar{l}$: determine the sign of the indirect CP violating amplitude.
- $K^+ \rightarrow \pi^+ + \nu + \bar{\nu}$: calculate the long distance ($l \geq 1/m_c$) part of charm contribution. Small ($\approx 4\%$) but leading theoretical uncertainty.
- **First calculations underway.**

Outlook

- DWF with physical quark masses reproduce QCD at the $\leq 3\%$ level on a $48^3 \times 64$ lattice.
- NPR gives continuum-like control of operator normalization and mixing.
- Theoretical advances allow rescattering effects to be correctly computed in Euclidean space (so far only for low energy π - π states).
- Many critical quantities can now be computed:
 - $K \rightarrow \pi \pi$, $\Delta I=3/2$ and $1/2$, ε'/ε
 - $m_{K_L} - m_{K_S}$ long dist. contribution to ε
 - $K \rightarrow \pi l \bar{l}$, $K \rightarrow \pi \nu \bar{\nu}$
 - Quark effects in $g_\mu - 2$ from HVP and HLbL at $O(\alpha^3)$