

# Anisotropic flow generated by hard partons in medium

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# Azimuthal anisotropy of hadronic momentum distributions

- parametrized by Fourier expansion

$$\frac{dN}{p_t dp_t dy d\phi} = \frac{1}{2\pi} \frac{dN}{p_t dp_t dy} \left( 1 + 2 \sum_{n=1}^{\infty} v_n(p_t, y) \cos(n(\phi - \phi_n)) \right)$$

- summation over many events in symmetric collision at midrapidity  
 $\Rightarrow$  symmetry constraints:  $\phi_n = 0, n = 2, 4, 6, \dots$
- all  $v_n$ 's non-vanishing in individual events

# Anisotropic expansion

- generic effect: blue-shift  
⇒ more particles and higher  $p_t$  in direction of stronger transverse flow
- link between the observable spectrum and the expansion of the fireball
- expansion results from the pressure gradients
- anisotropic expansion  $\Leftarrow$  anisotropic pressure gradients in initial conditions

# Hydrodynamics – state of the art

- Conservation laws

$$\partial_\mu T^{\mu\nu} = 0$$

- energy momentum tensor

$$T^{\mu\nu} = (\epsilon + p)u^\mu u^\nu + p g^{\mu\nu} + \Pi^{\mu\nu}$$

with stress tensor  $\Pi^{\mu\nu} = \pi^{\mu\nu} + \Delta^{\mu\nu}\Pi$

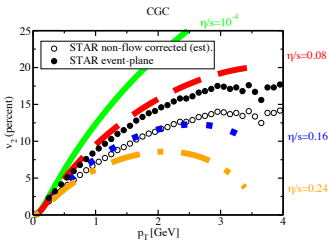
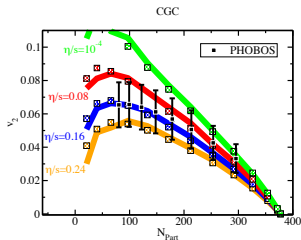
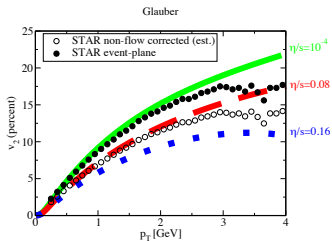
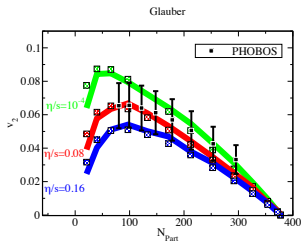
(split into traceless shear and non-traceless bulk contribution)

- viscous corrections

$$\begin{aligned}\pi^{\mu\nu} &= \eta(\epsilon) \left( \nabla^\mu u^\nu + \nabla^\nu u^\mu - \frac{2}{3} \Delta^{\mu\nu} \nabla_\alpha u^\alpha \right) \\ \Pi &= \zeta(\epsilon) \nabla_\alpha u^\alpha\end{aligned}$$

- Equation of State  $p = p(\epsilon)$

# Initial conditions – an ambiguity (illustration)

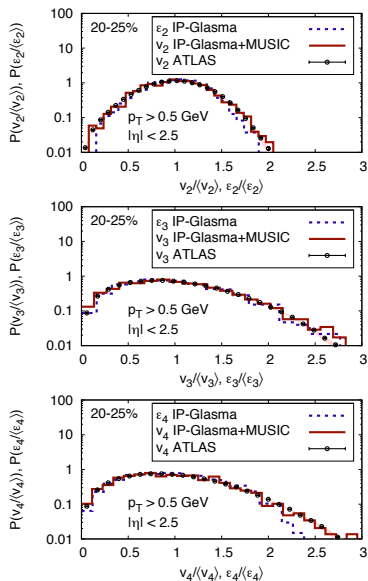


[M. Luzum, P. Romatschke: Phys. Rev. C **78** (2008) 034915]

# Fluctuating initial conditions

- Use the **fluctuations** of  $v_n$ 's to get the access to initial conditions.
- fluctuations of  $v_n$ 's seem to follow those of spatial anisotropies  $\varepsilon_n$ 's

[Ch. Gale et al.:  
Phys. Rev. Lett. **110** (2013) 012302]



# Motivation

## We want

- Equation of State
- transport properties (viscosities)

## Then we must

- disentangle the influence of (fluctuating) initial conditions
- get under control all other effects influencing the anisotropies of hadronic distributions

## Here we propose

a novel mechanism which contributes to anisotropies of hadronic distributions.

# The idea

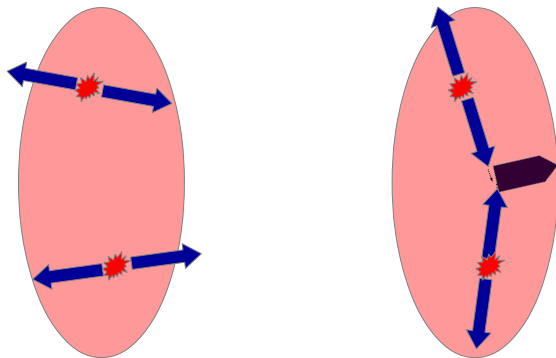
- At the LHC there is copious production of hard partons – may have more than one pair in single event.
- Their momentum is deposited into medium over some time span  
⇒ collective flow, wakes, **streams**
- Anisotropic flow – event by event
- Elliptic flow after summation over all events.



## Anisotropic flow from isotropic jets

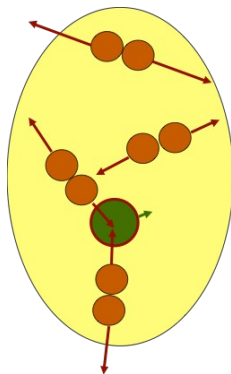
Streams are more likely to merge if they are directed out of reaction plane

- ⇒ less contribution to flow out of plane
- ⇒ enhance  $v_2$  correlated with the reaction plane
- ⇒ also contribute to  $v_3$



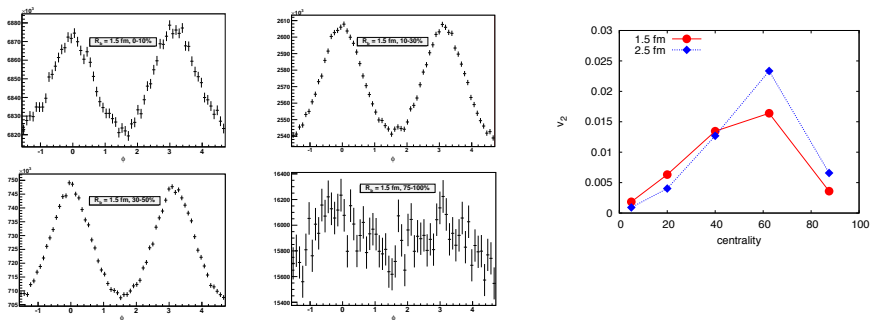
## Check the idea with a toy model

- Streams represented by drops
- Pairs of drops back-to-back (with some  $k_t$  smearing)
- Drops merge after they meet
- Size of the drop represents the radius of the stream
- Pions evaporate from droplets ( $T = 175$  MeV)



# Toy model – results

## Azimuthal distribution of hadrons



[B. Tomášik, P. Lévai: J.Phys.G **38** (2011) 095101]

# Hydrodynamic implementation

[B. Betz et al.: Phys. Rev. C **79** (2009) 034902]

Ideal hydrodynamics with source term

$$\partial_\mu T^{\mu\nu} = J^\nu$$

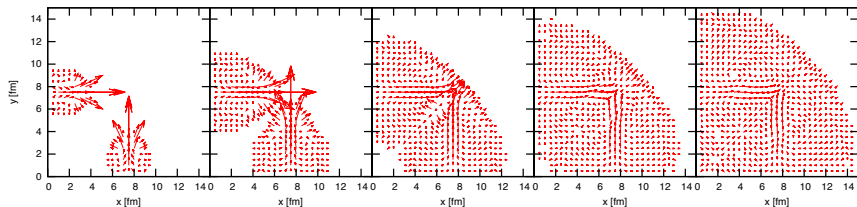
$$J^\nu = \sum_i \frac{1}{(2\pi\sigma_i^2)^{3/2}} \exp\left(-\frac{(\vec{x} - \vec{x}_{\text{jet},i})^2}{2\sigma_i^2}\right) \left(\frac{dE_i}{dt}, \frac{d\vec{P}_i}{dt}\right)$$

with  $\sigma = 0.3$  fm

# Test of the concept: static medium

Two streams meet perpendicularly

Plot momentum density



[M. Schulc, B. Tomášik: J. Phys. G **40** (2013) 125104]

# Hydrodynamic simulations of nuclear collisions

- 3+1D ideal hydrodynamics
- EoS from P. Petreczky, P. Huovinen: Nucl. Phys. A **897** (2010) 26
- **smooth** initial energy density scaled with

$$W(x, y; b) = (1 - \alpha)n_w(x, y; b) + \alpha n_{\text{bin}}(x, y; b)$$

with  $\alpha = 0.16$ ,  $\varepsilon(0, 0, 0) = 60 \text{ GeV}/\text{fm}^3$  at  $\tau_0 = 0.55 \text{ fm}/c$   
rapidity plateau over 10 units of rapidity

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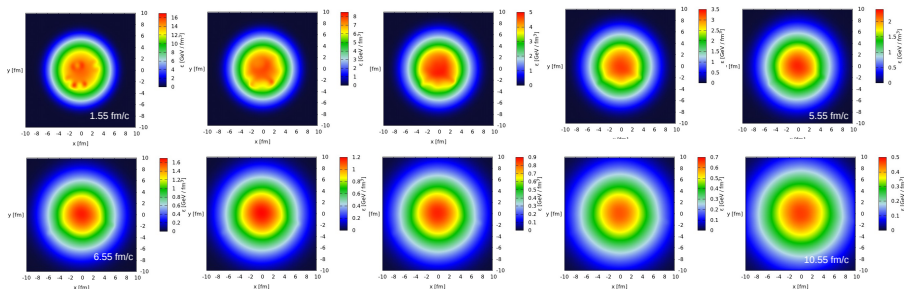
$$\frac{dE}{dx} = \frac{dE}{dx} \Big|_0 \frac{s}{s_0}$$

- fluctuating number of jet pairs

# Illustration: evolution of energy density

Evolution of an event with four pairs of jets at the beginning.

frames follow with time delay  $1\text{fm}/c$



# Results from ultra-central collisions

Anisotropy coefficients

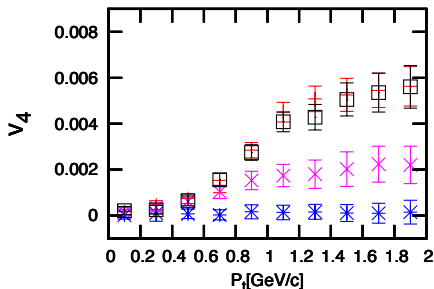
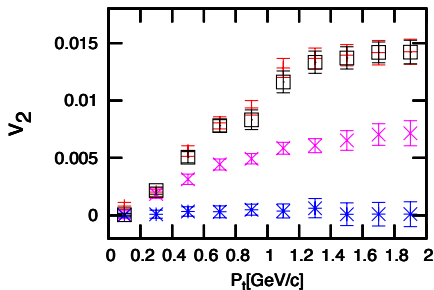
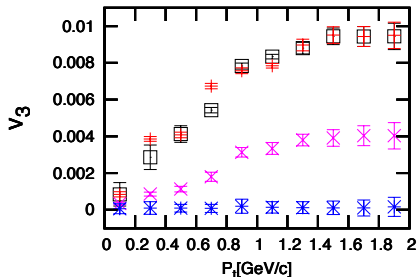
compare:

$dE/dx = 7$  GeV/fm

$dE/dx = 4$  GeV/fm

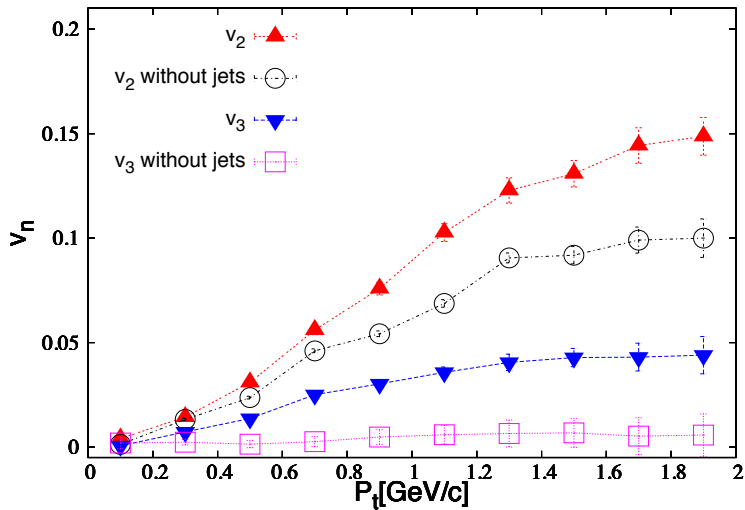
hot spots

smooth initial conditions





# Results from 30–40% centrality



## Similar approaches

- Y. Tachibana, T. Hirano: Phys. Rev. C **90** (2014) 021902  
reponse of medium to only one dijet
- R.P.G. Andrade, J. Noronha, G. Denicol: arxiv:1403.1789  
one dijet, 2+1D hydrodynamics
- S. Floerchinger and K. Zapp: arxiv:1407.1782  
1+1D hydrodynamics

# Conclusions and Outlook

- Momentum deposition from hard partons gives large contribution to anisotropic flow  
⇒ must be included in simulations
- The interplay of many induced streams is important
- Outlook: simulations with viscous hydrodynamics and fluctuating initial conditions