

Fermionic Symmetry Protected Topological Phase Induced by Interaction

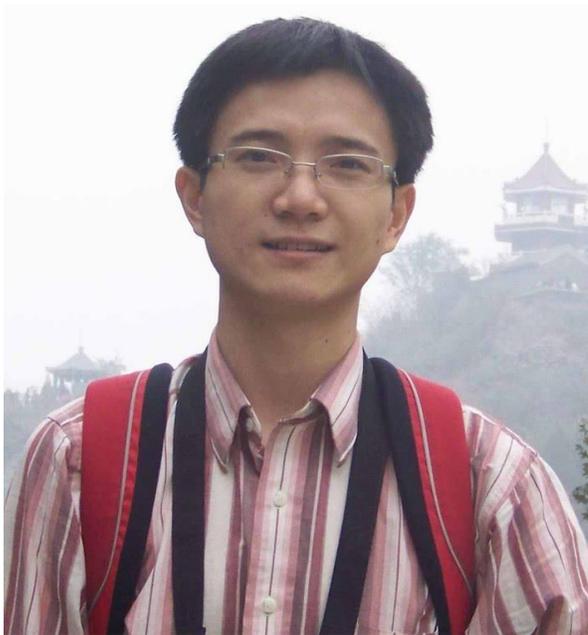
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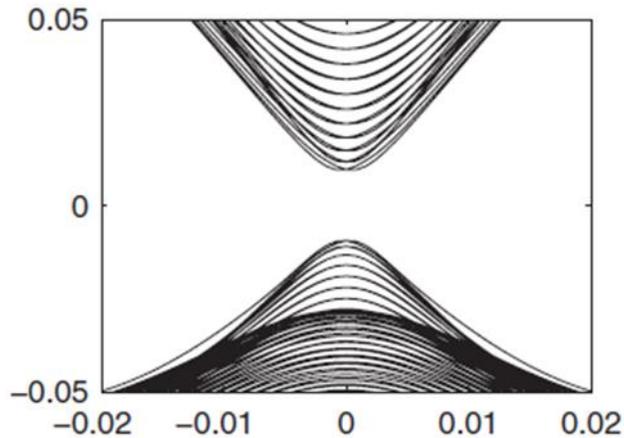


Hongchen Jiang
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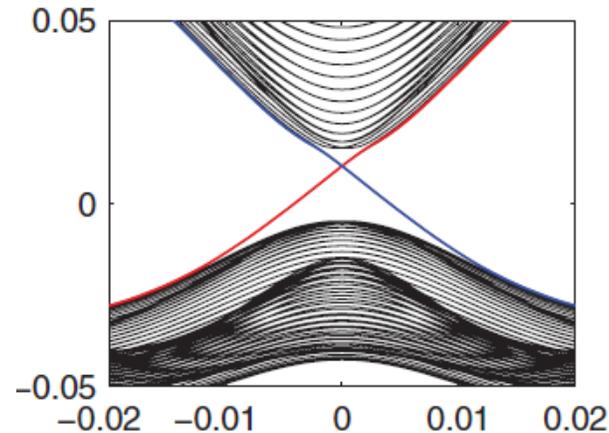
Outline

1. Background and motivation
2. Study of a specific model
3. Conclusion

Background and Motivation



Band insulator



2D topological insulator

Key properties of TI:

- Bulk gap
- Gapless edge states protected by time reversal symmetry and U(1)
- No symmetry breaking

Mass terms

$$B(\psi_{\uparrow}^{\dagger}\psi_{\downarrow} + \psi_{\downarrow}^{\dagger}\psi_{\uparrow})$$

breaks T

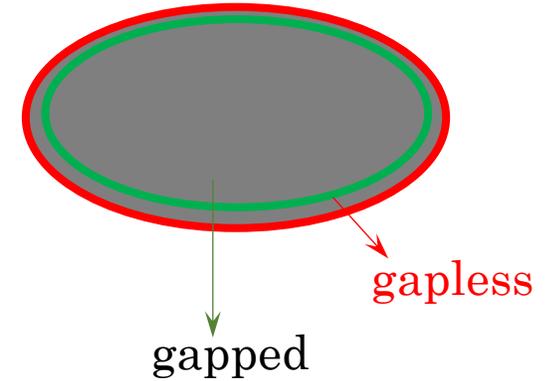
$$\Delta(\psi_{\uparrow}^{\dagger}\psi_{\downarrow}^{\dagger} + \psi_{\downarrow}\psi_{\uparrow})$$

breaks $U(1)$

Symmetry Protected Topological phase (SPT phase)

Key points of SPT

- No symmetry breaking
- bulk is gapped
- boundary excitation is gapless
- protected by symmetry



Classification of SPT

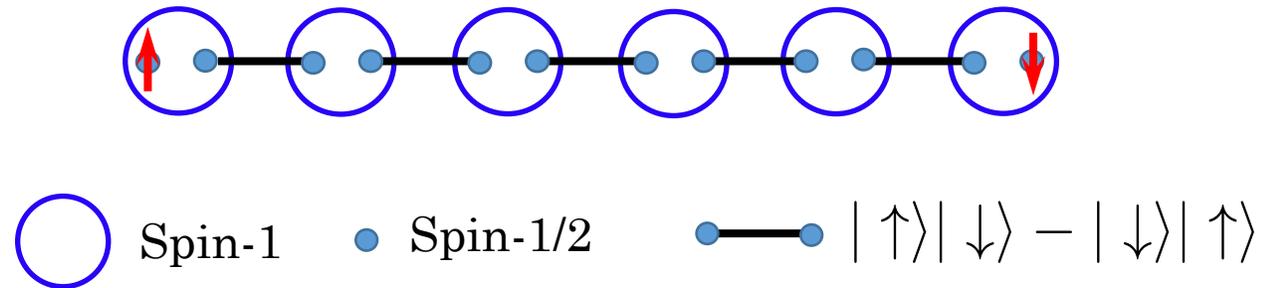
- Free fermion systems are classified by K-theory

q	$\pi_0(R_q)$	$d = 1$	$d = 2$	$d = 3$
0	\mathbb{Z}		no symmetry ($p_x + ip_y$, e.g., SrRu)	T only ($^3\text{He-B}$)
1	\mathbb{Z}_2	no symmetry (Majorana chain)	T only ($(p_x + ip_y)\uparrow + (p_x - ip_y)\downarrow$)	T and Q (BiSb)
2	\mathbb{Z}_2	T only ($(\text{TMTSF})_2\text{X}$)	T and Q (HgTe)	
3	0	T and Q		
4	\mathbb{Z}			
5	0			
6	0			
7	0			no symmetry

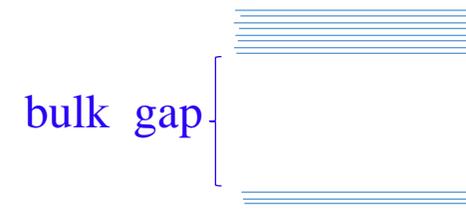
■ SPT phases also exist in Bosonic systems

- $S=1$ Haldane phase (Bosonic) $SO(3)$ symmetry

$$H = \sum_i J_1 \mathbf{S}_i \cdot \mathbf{S}_{i+1}$$



- Finite excitation gap
- No symmetry breaking
- $SO(3)$ symmetry protected spin-1/2 **edge states**
- Degeneracy of entanglement spectrum



■ entanglement spectrum for 1D



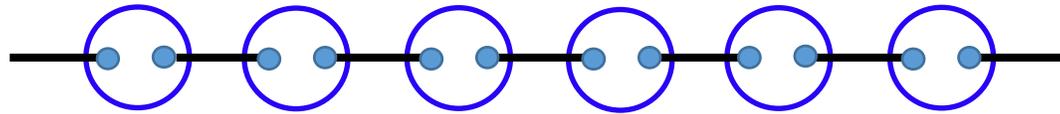
$\rho_A = \text{Tr}_B (\rho)$, ρ density matrix of the ground state

Entanglement spectrum: eigenvalues of ρ_A

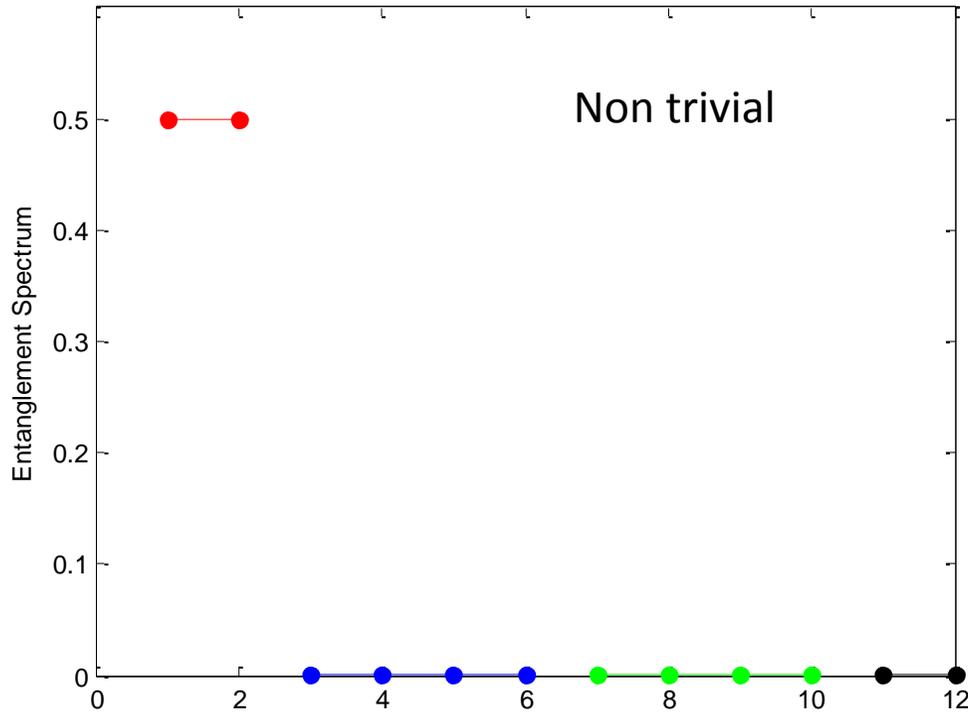
Entanglement spectrum

$$H = \sum_i J_1 \mathbf{S}_i \cdot \mathbf{S}_{i+1} + J_2 \mathbf{S}_i \cdot \mathbf{S}_{i+2}$$

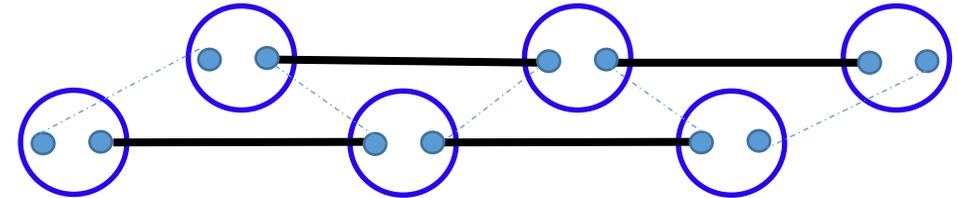
$SO(3)$ symmetry



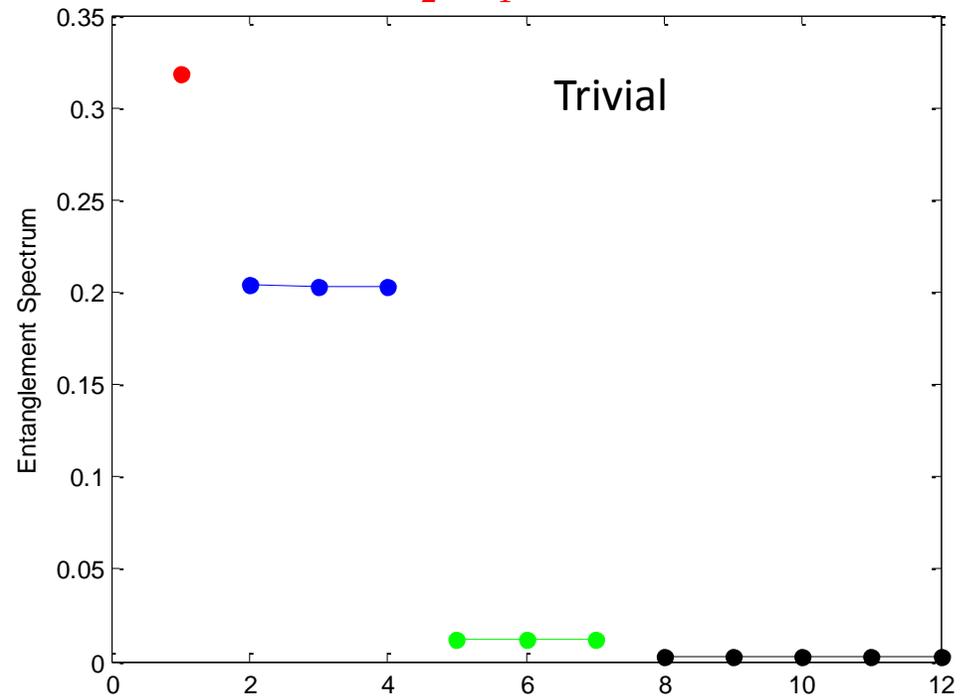
$J_2/J_1=0.3$



Non trivial



$J_2/J_1=0.9$



Trivial

Entanglement spectrum

Entanglement spectrum

Classification of Bosonic SPT

■ 1D by the projective representation

$$H = \sum_i J_1 \mathbf{S}_i \cdot \mathbf{S}_{i+1} + J_2 \mathbf{S}_i \cdot \mathbf{S}_{i+2} \quad : \text{2 phases}$$

$SO(3)$ symmetry

$SO(3)$: 2 projective representation

Xie Chen, Zheng-Cheng Gu, and Xiao-Gang Wen, 2011

■ 2D or higher by group cohomology

Chen, Gu, **ZXL**, Wen, Science, 338, 1604 (2012);
Chen, Gu, **ZXL**, Wen, PRB, 87, 155114 (2013)

Classification of Interacting Fermionic SPT

- 1D by projective representation
- 2D or higher partially by super-cohomology
- effect of interaction

1D superconductor

Symmetry	Free classification	With interactions
$U(1) \times Z_2^T$	\mathbb{Z}	\mathbb{Z}_4
$Z_n \times Z_2^T$ (n even)	\mathbb{Z}	\mathbb{Z}_4
$Z_n \times Z_2^T$ (n odd and $n > 1$)	\mathbb{Z}	\mathbb{Z}_2

2D superconductor

Time reversal and mirror reflect	\mathbb{Z}	\mathbb{Z}_8
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It seems that interactions reduce the classification. However, it maybe not true!

Model

$$H = H_0 + H_U + H_J$$

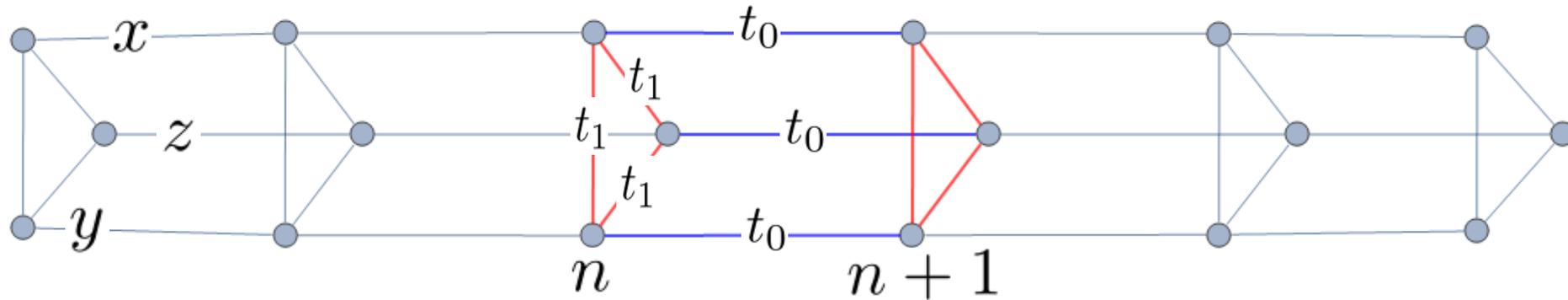
$$H_0 = -\left\{ \sum_{i,\alpha} t_0 (c_{i,\alpha}^\dagger c_{i+1,\alpha}) + \sum_{i,\{\alpha,\hat{\alpha}\}} t_1 c_{i,\alpha}^\dagger c_{i,\hat{\alpha}} + h.c. \right\} + \lambda N_i$$

$$H_U = U \sum_i (N_i - 1)^2 \quad H_J = J \sum_i \mathbf{S}_i \cdot \mathbf{S}_{i+1}$$

$$N_i = n_{i,x} + n_{i,y} + n_{i,z}$$

$$S^\alpha = \sum_{\beta,\gamma} i \varepsilon^{\alpha\beta\gamma} c_\beta^\dagger c_\gamma$$

$$\alpha = x, y, z, \{\alpha, \hat{\alpha}\} = \{x, y\}, \{y, z\}, \{z, x\}$$



Symmetry and classification

■ The model : U(1) and time reversal symmetry T

$$U_\theta c_\alpha U_\theta^{-1} = c_\alpha e^{i\theta}$$
$$T c_\alpha T^{-1} = -c_\alpha. \quad T^2 = 1$$

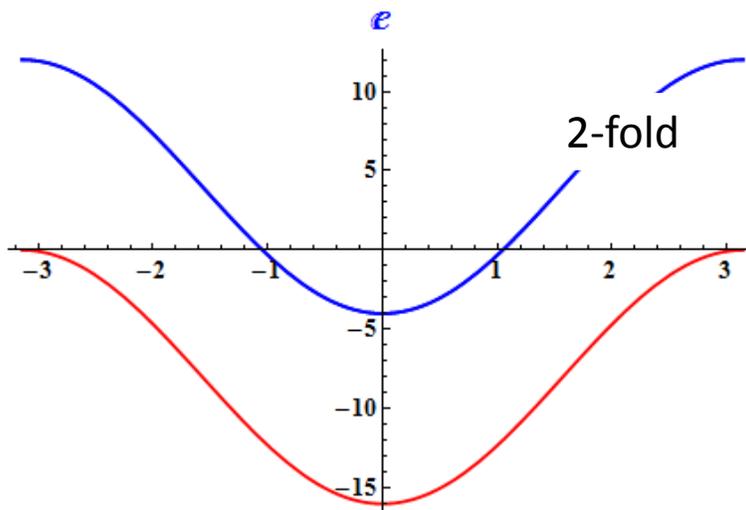
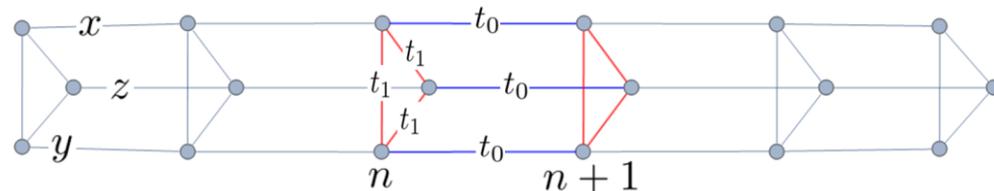
■ Classification

free case, only the trivial phase.

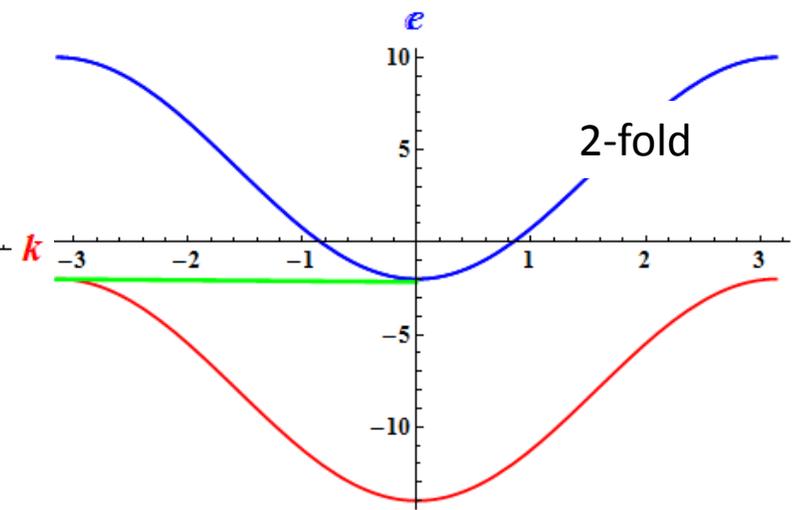
2 projective representation → two SPT phases.

q	$\pi_0(\mathcal{R}_q)$	$d = 1$
0	\mathbb{Z}	
1	\mathbb{Z}_2	no symmetry (Majorana chain)
2	\mathbb{Z}_2	T only ((TMTSF) ₂ X)
3	0	T and Q
4	\mathbb{Z}	
5	0	
6	0	
7	0	

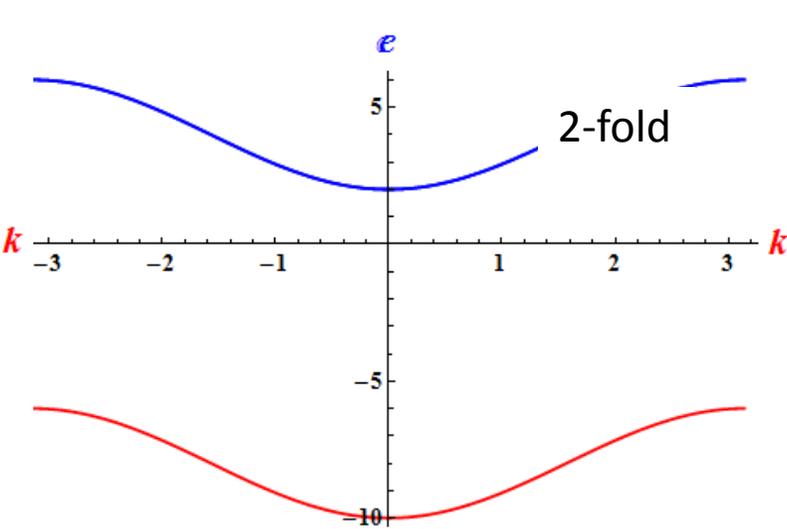
Non-interacting limit $U=0, J=0$



$$t_1 < \frac{8}{3}t_0$$



$$t_1 = \frac{8}{3}t_0$$

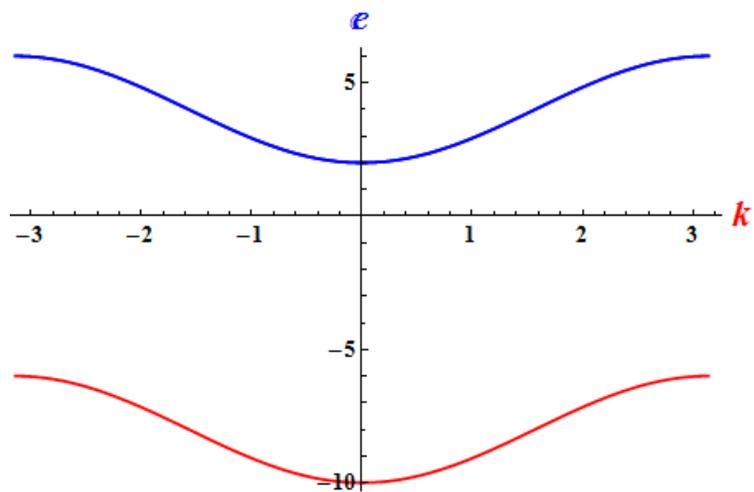
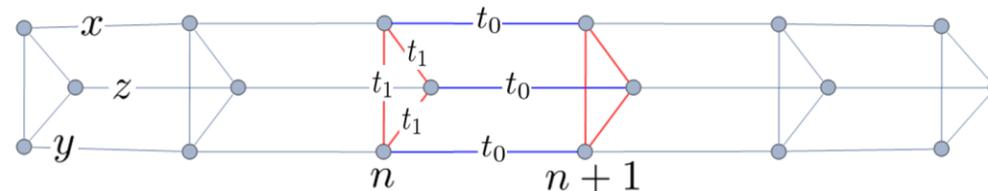


$$t_1 > \frac{8}{3}t_0$$

Non-interacting limit $U=0, J=0$

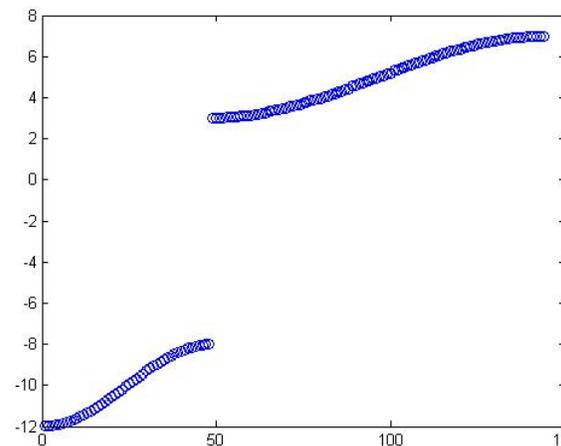
SPT only focus on gapped phases.

In $\frac{1}{3}$ filling case



$$t_1 > \frac{8}{3}t_0$$

In region of $t_1 > \frac{8}{3}t_0$, *only trivial phase*, i.e band insulator

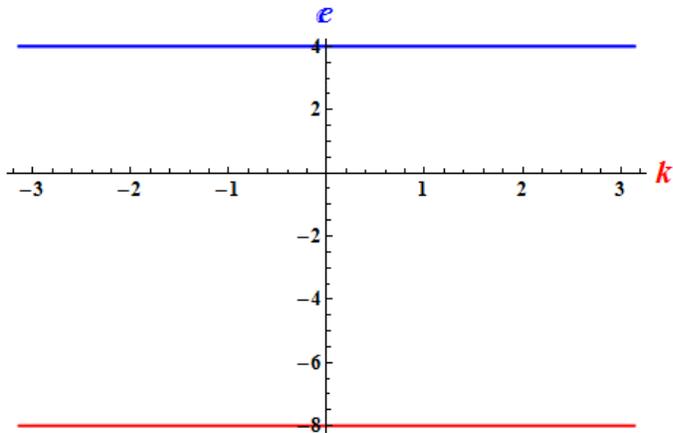


Open boundary, *no zero mode*, indicating the trivialness of phase

Interacting case

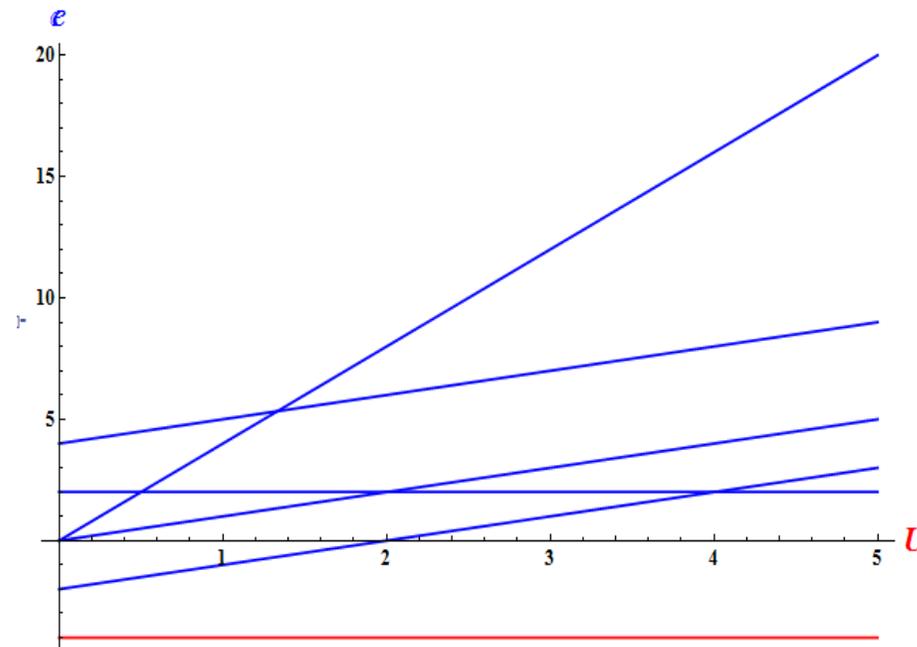
$$t_1 \gg t_0$$

Add Hubbard interaction $H_u = U \sum_i (N_i - 1)^2$



$$\frac{t_1}{t_0} \rightarrow \infty, U = 0$$

the charge are *localized* at each rung in ground state.



the band insulator and the Mott insulator are *adiabatically connected*

Interacting case

- Add Heisenberg-like interaction $H_J = J \sum_i \mathbf{S}_i \cdot \mathbf{S}_{i+1}$,
- This term competes to exhibit the Haldane phase, so in the limit $J \rightarrow \infty$, the system would fall into Haldane phase.

Interacting case

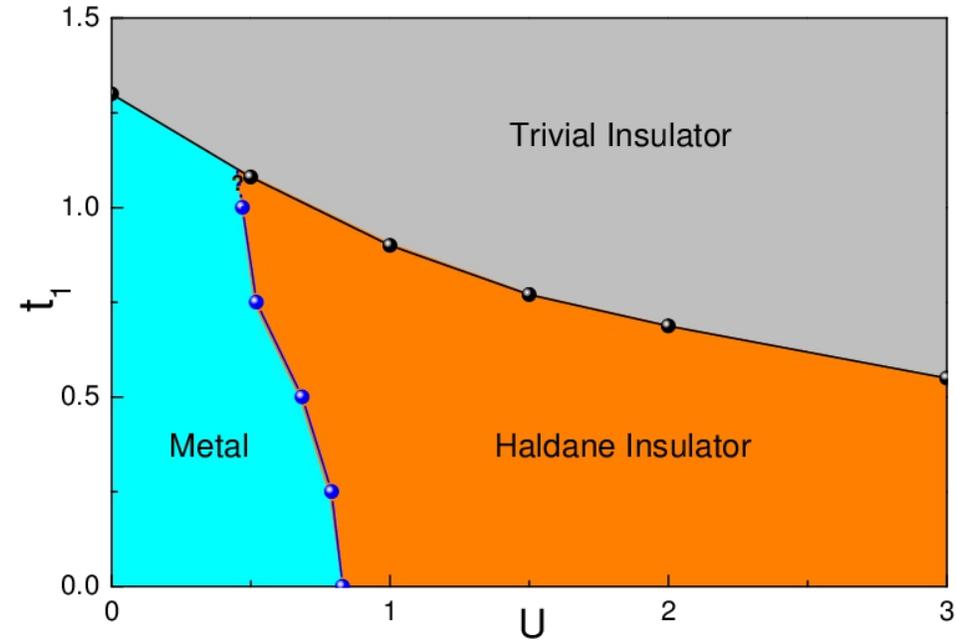
t_0 terms \rightarrow metal phase,
 t_1 and U \rightarrow trivial insulator phase,
 J terms \rightarrow nontrivial Haldane phase.

Competition

For **finite-parameters region**, the system maybe exist the **trivial and nontrivial** gapped phases.

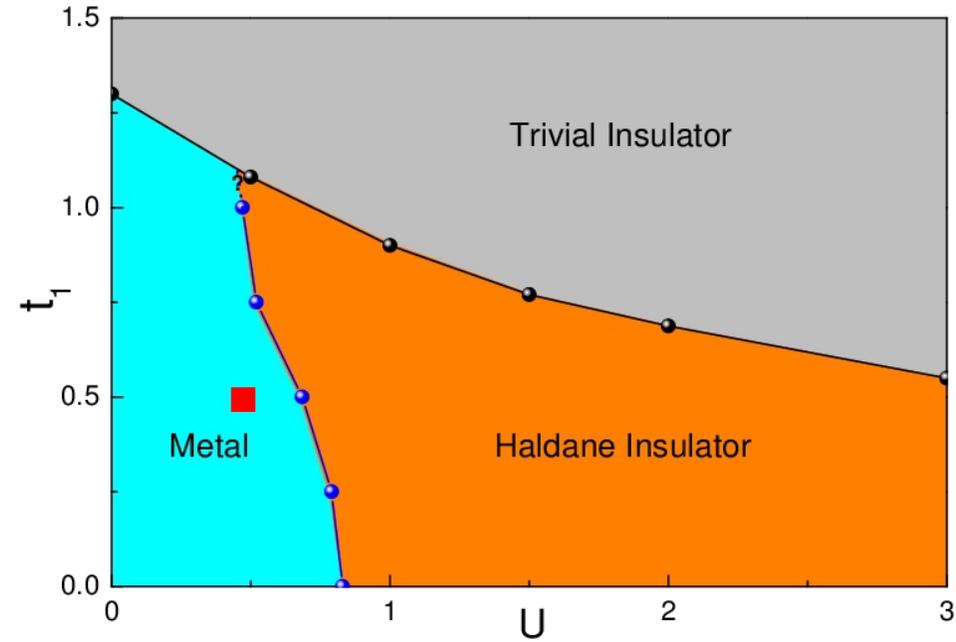
Interacting case

Set $t_0 = 1, J = 0.5$ in DMRG simulation



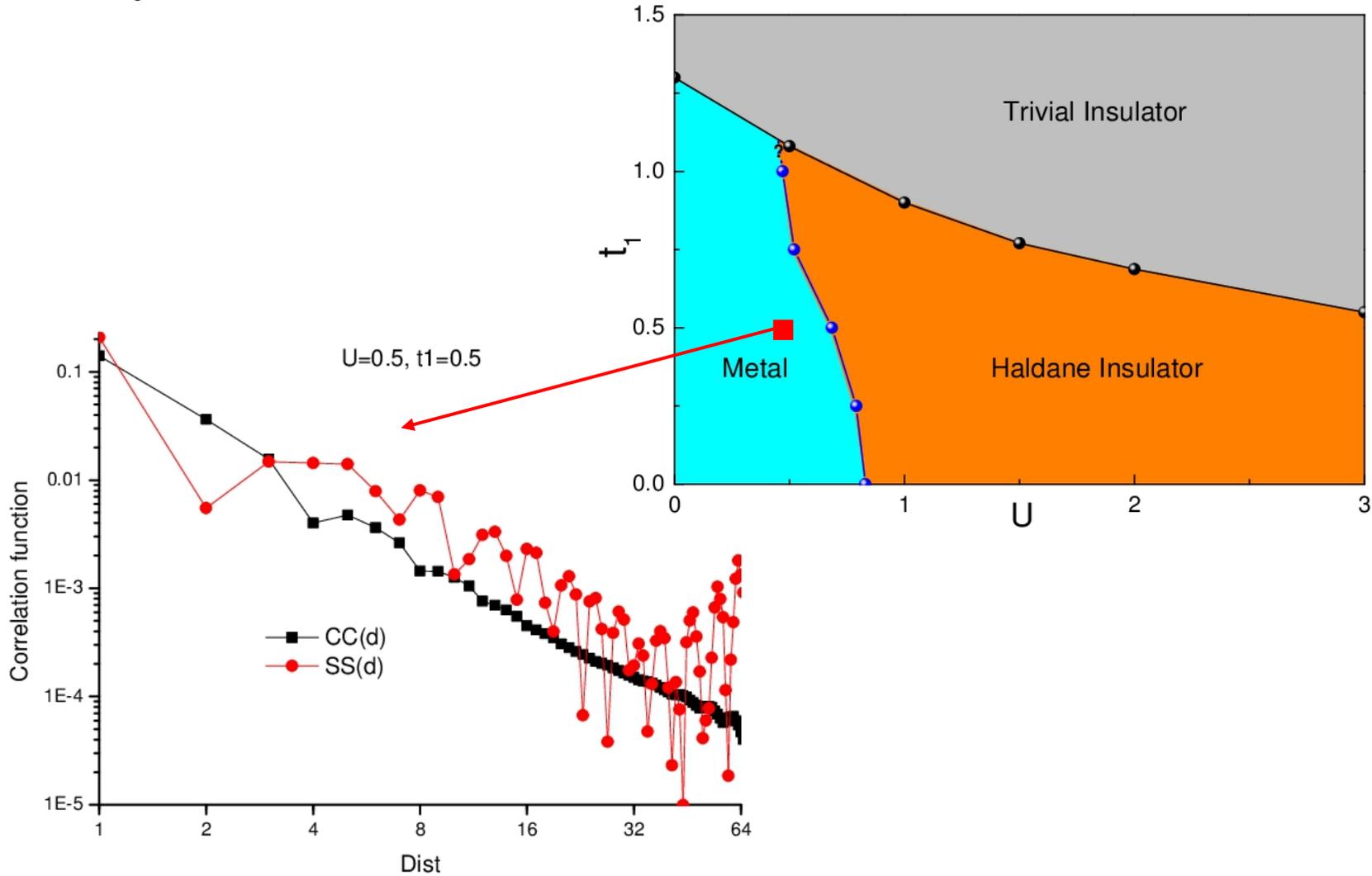
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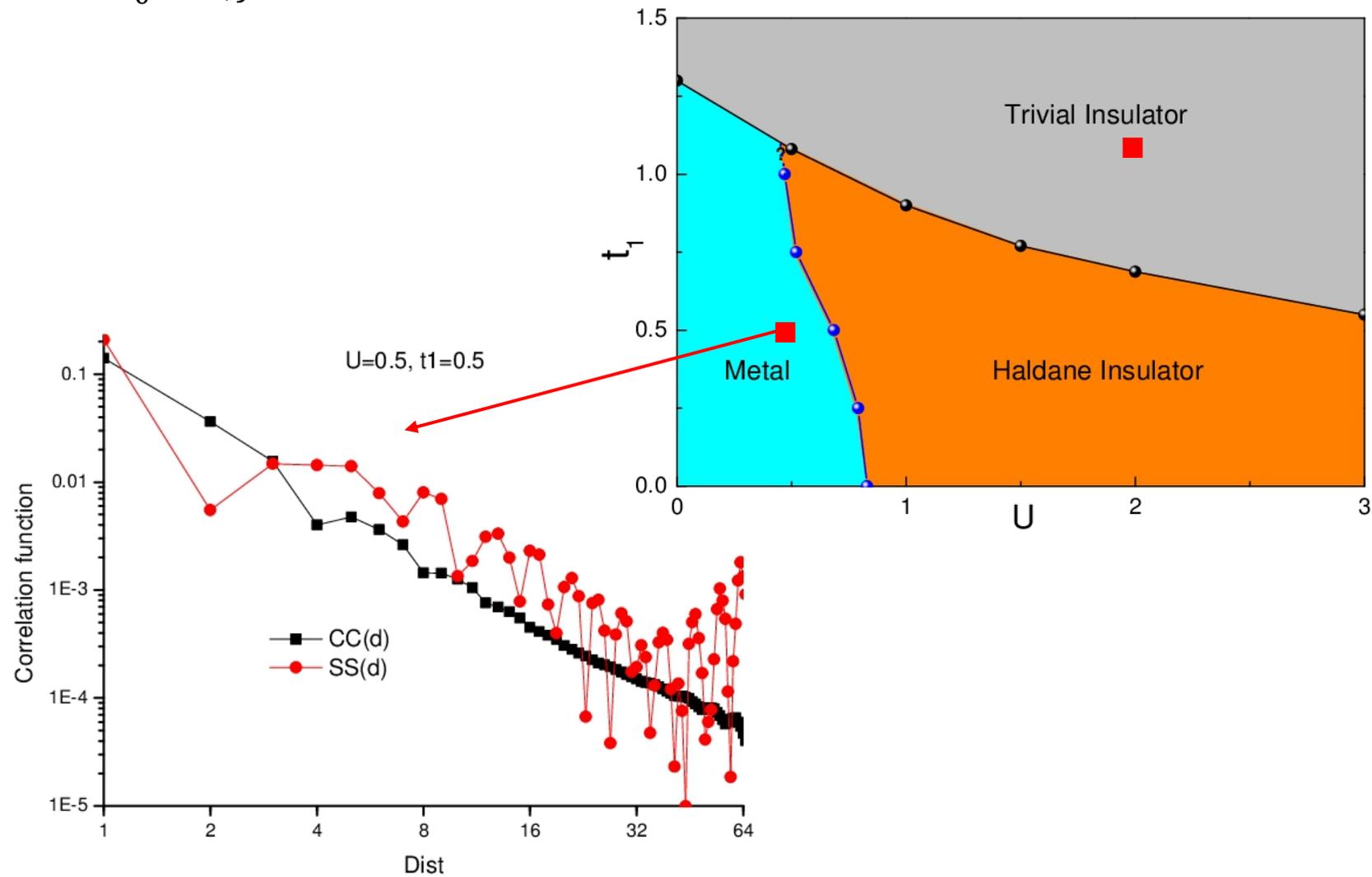
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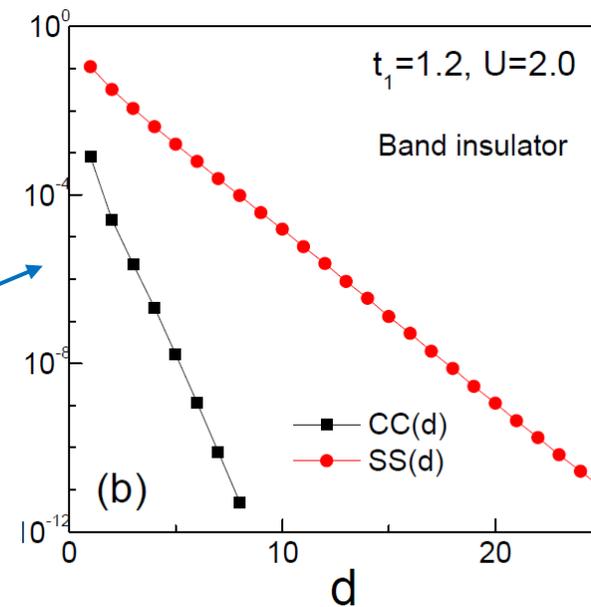
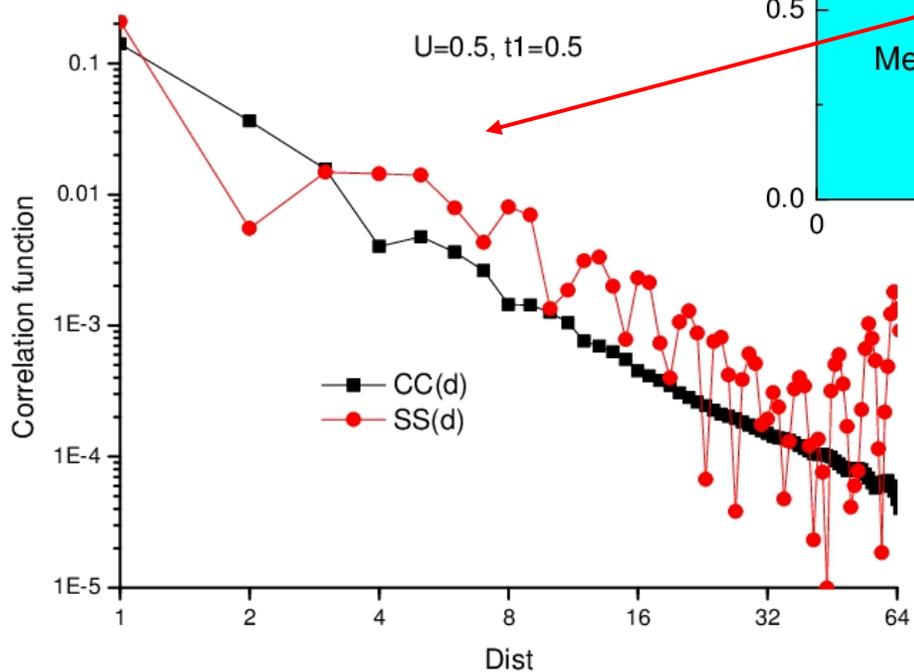
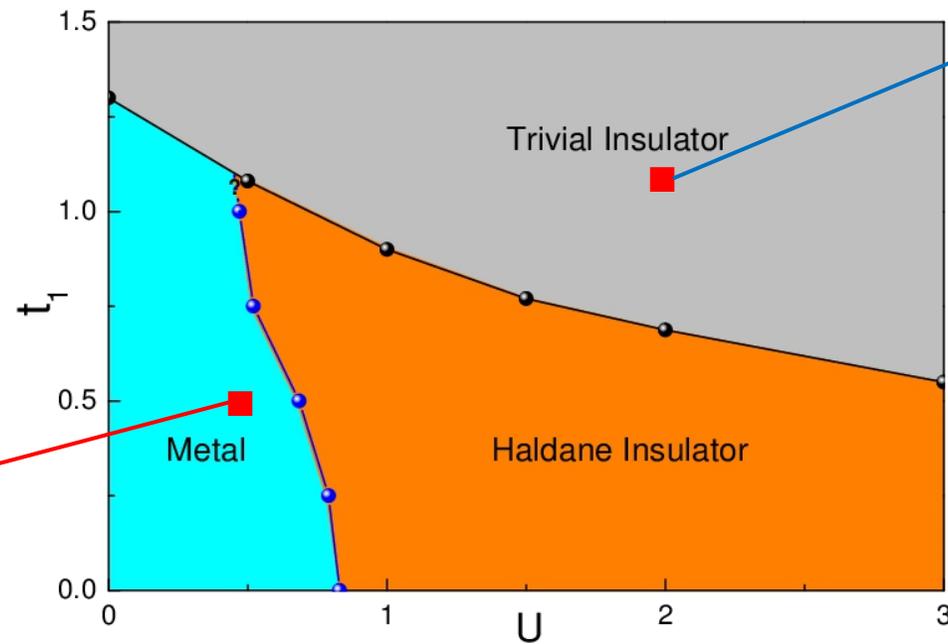
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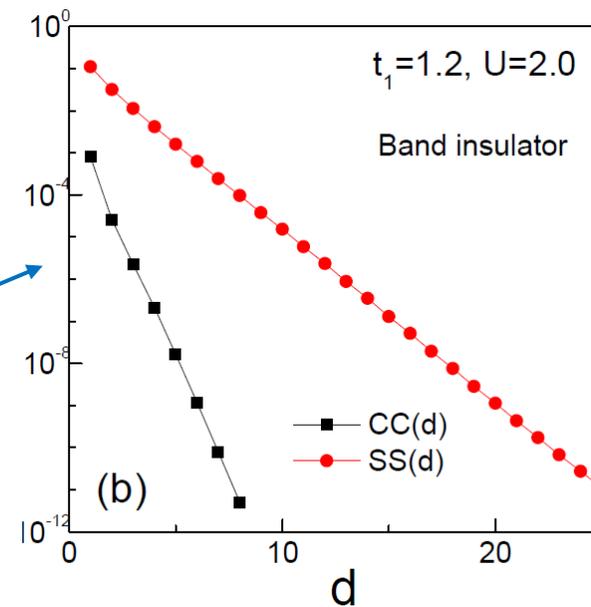
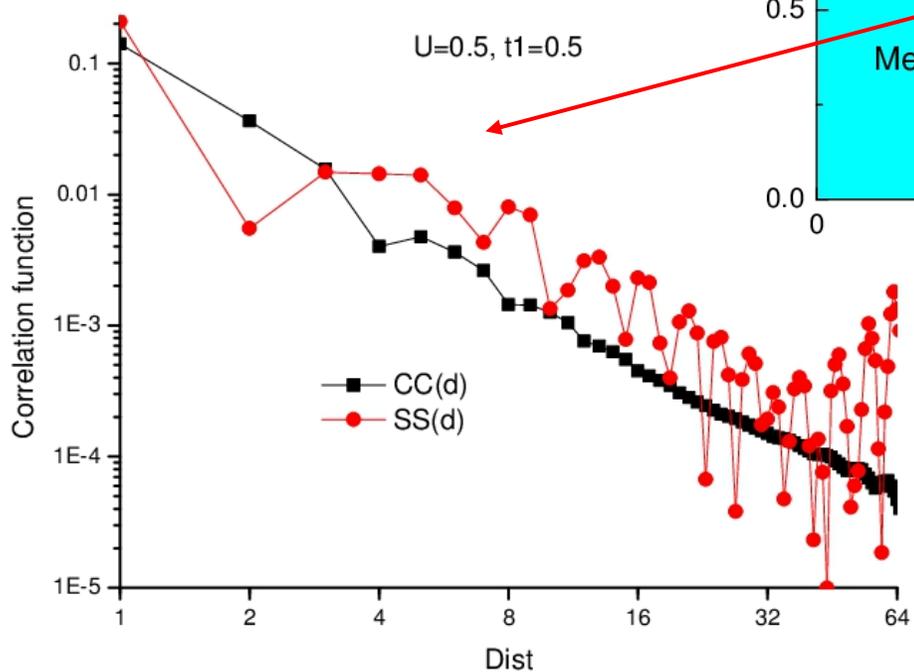
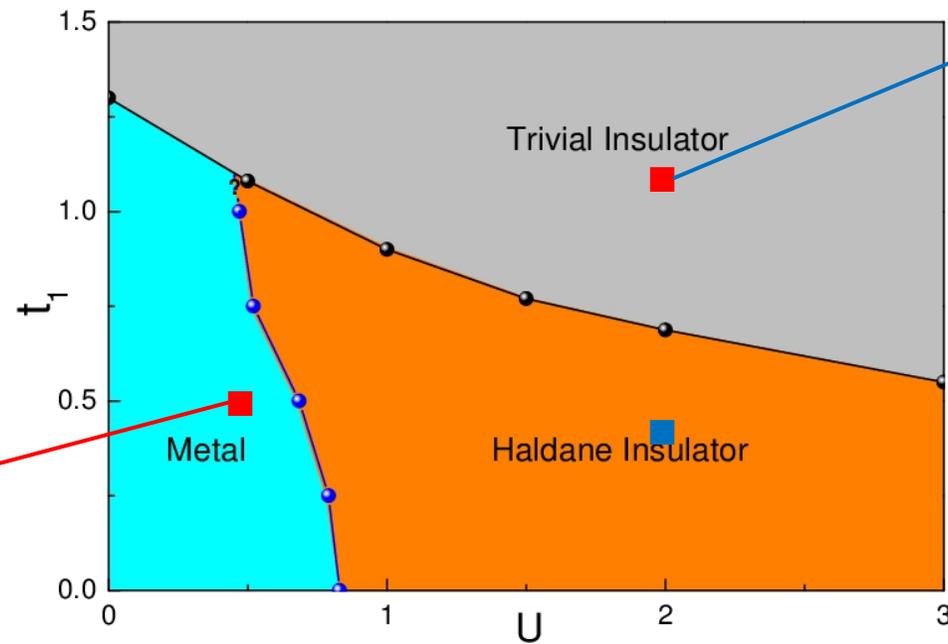
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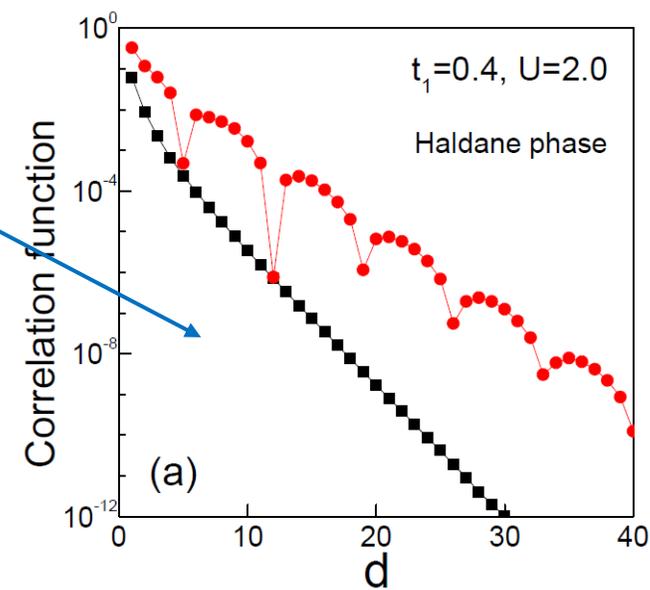
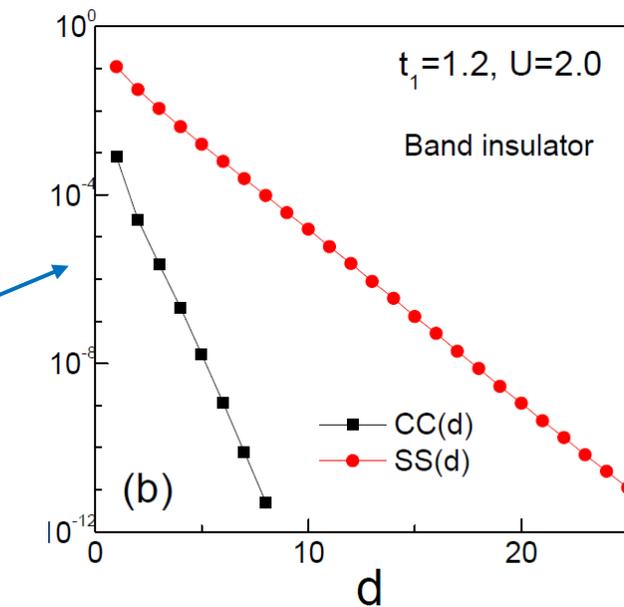
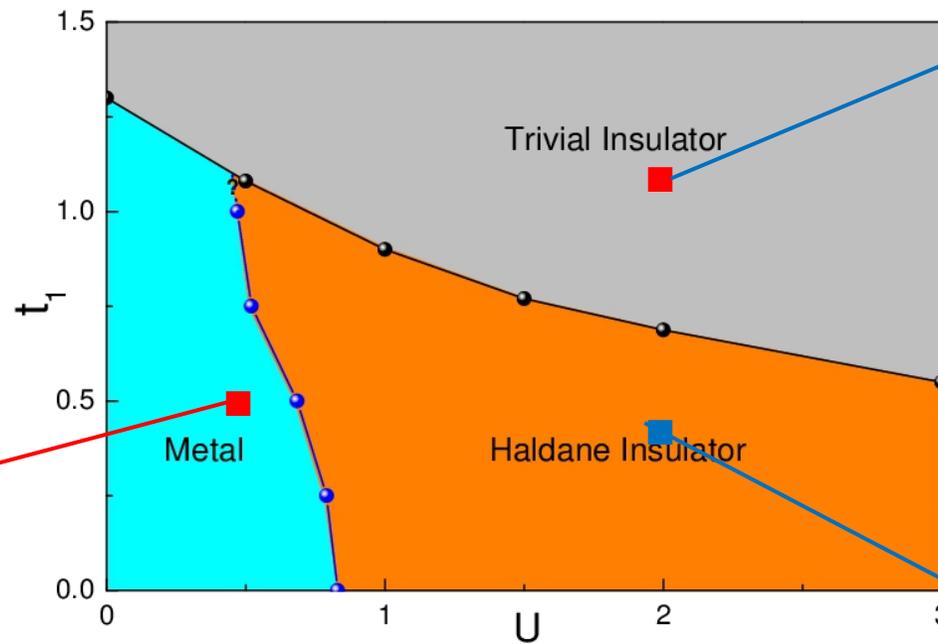
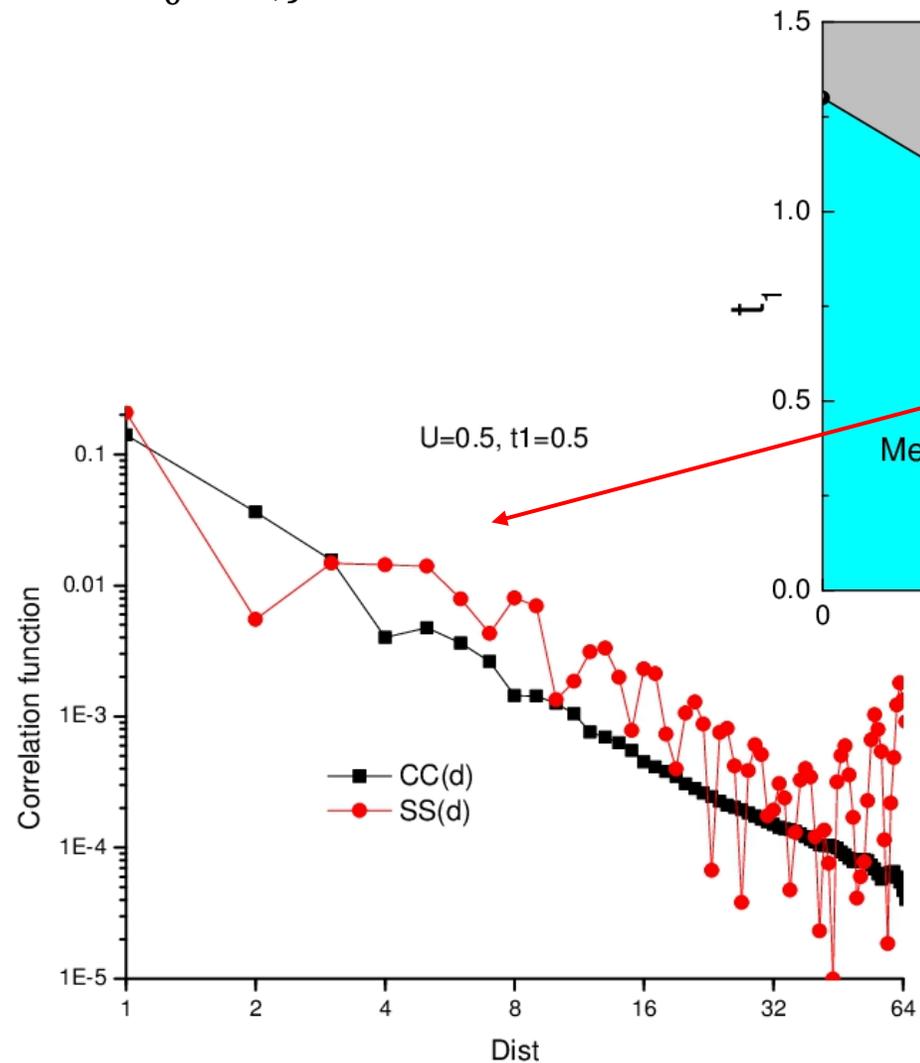
Interacting case

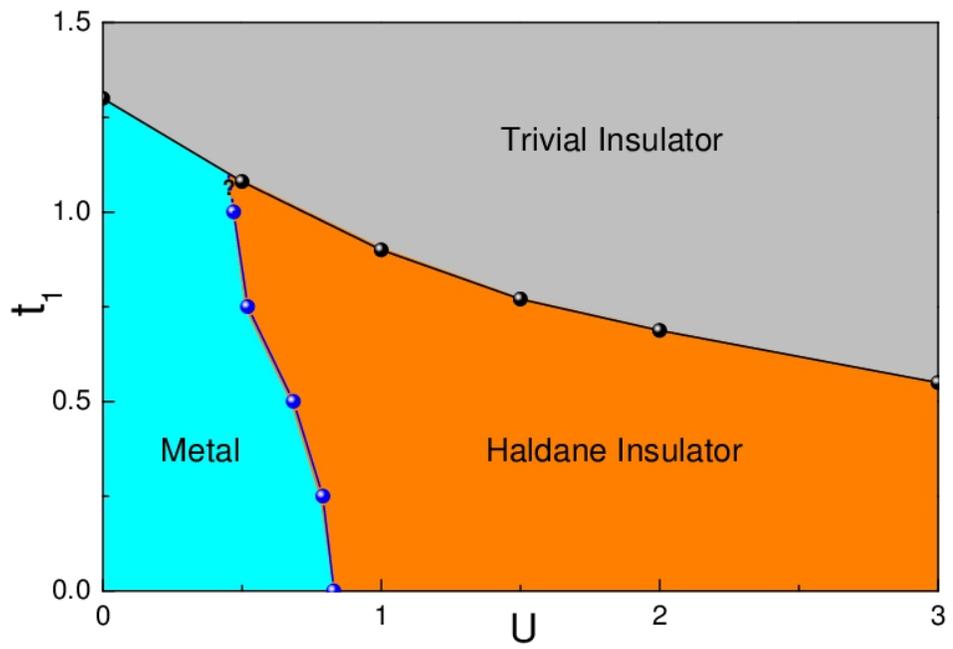
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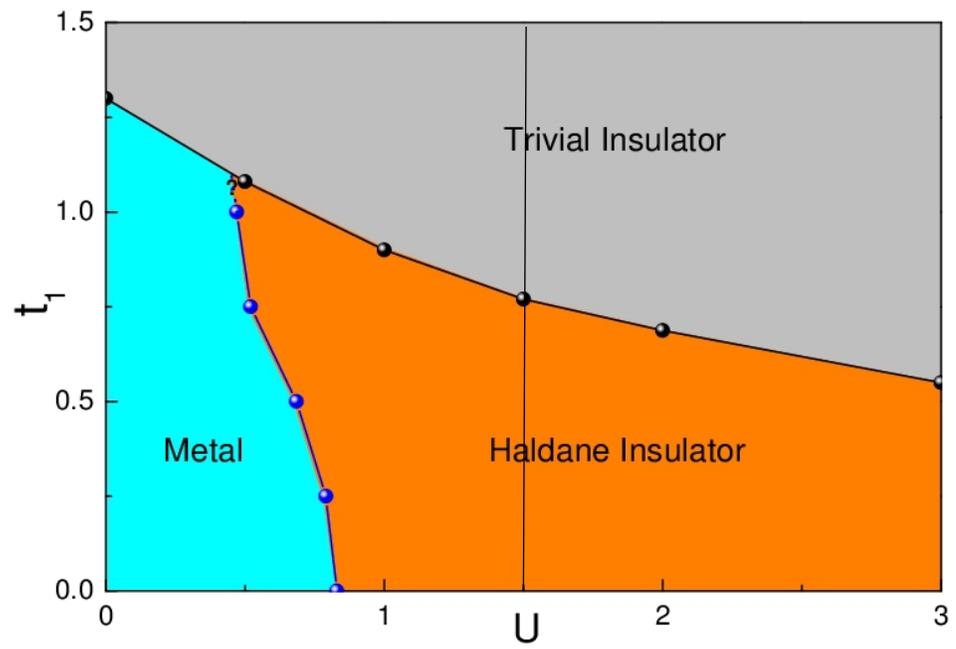


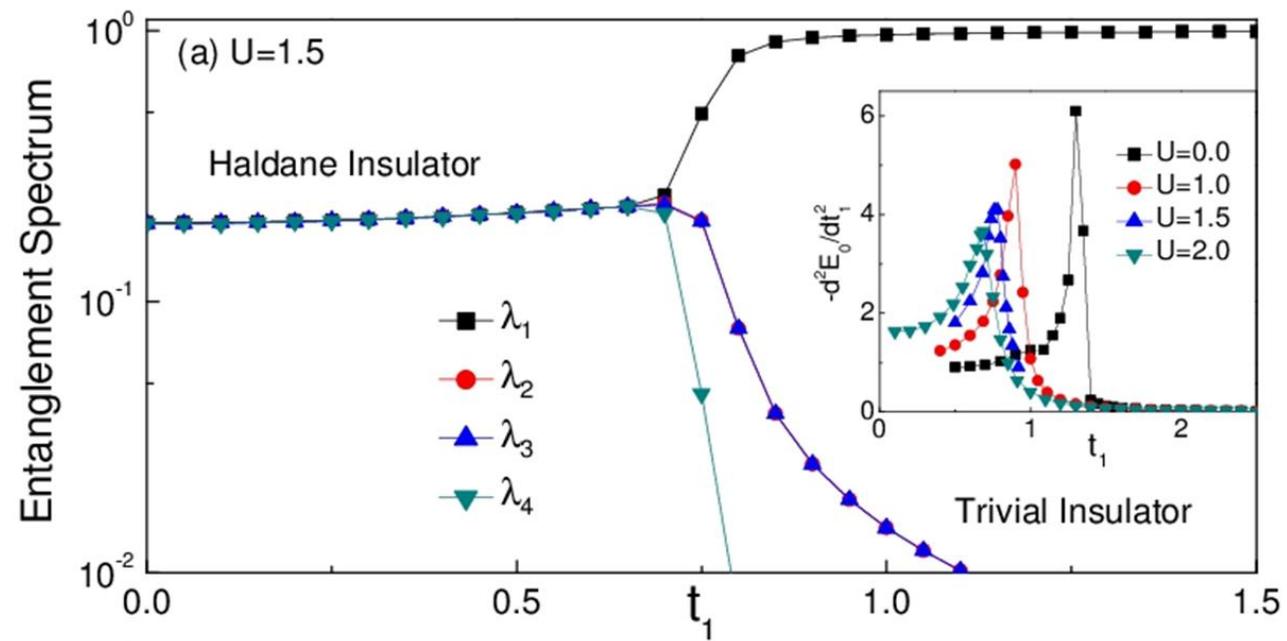
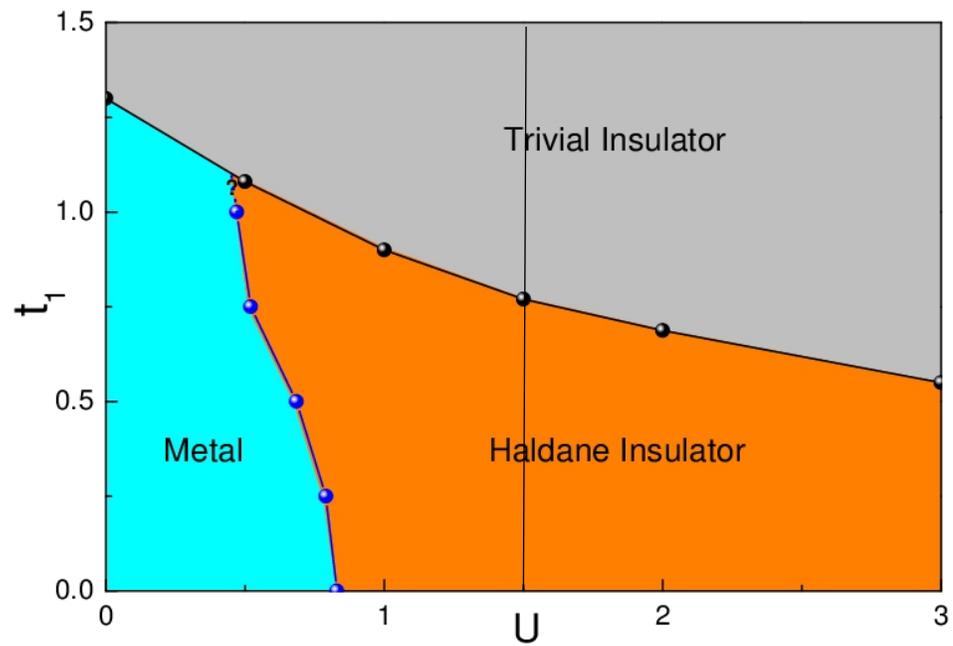
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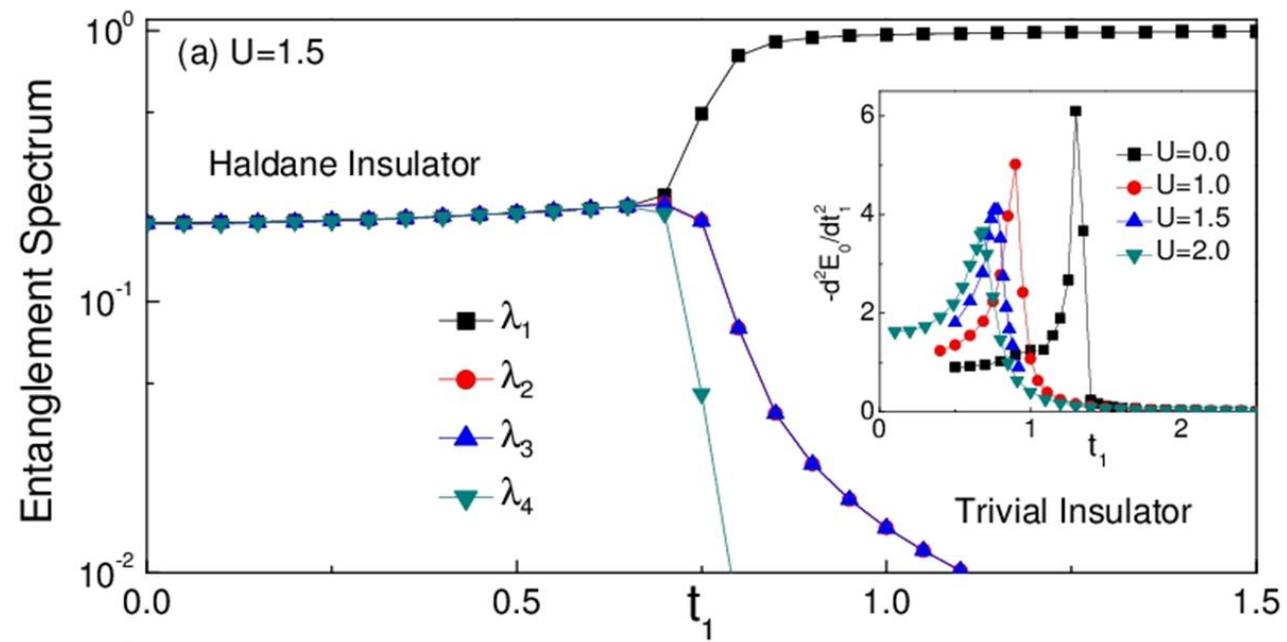
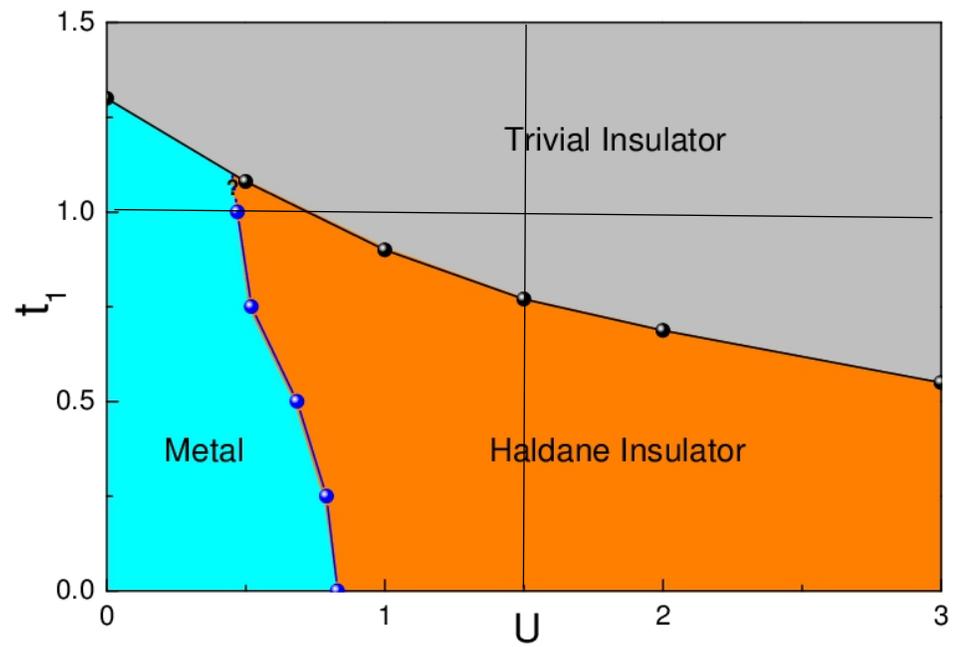
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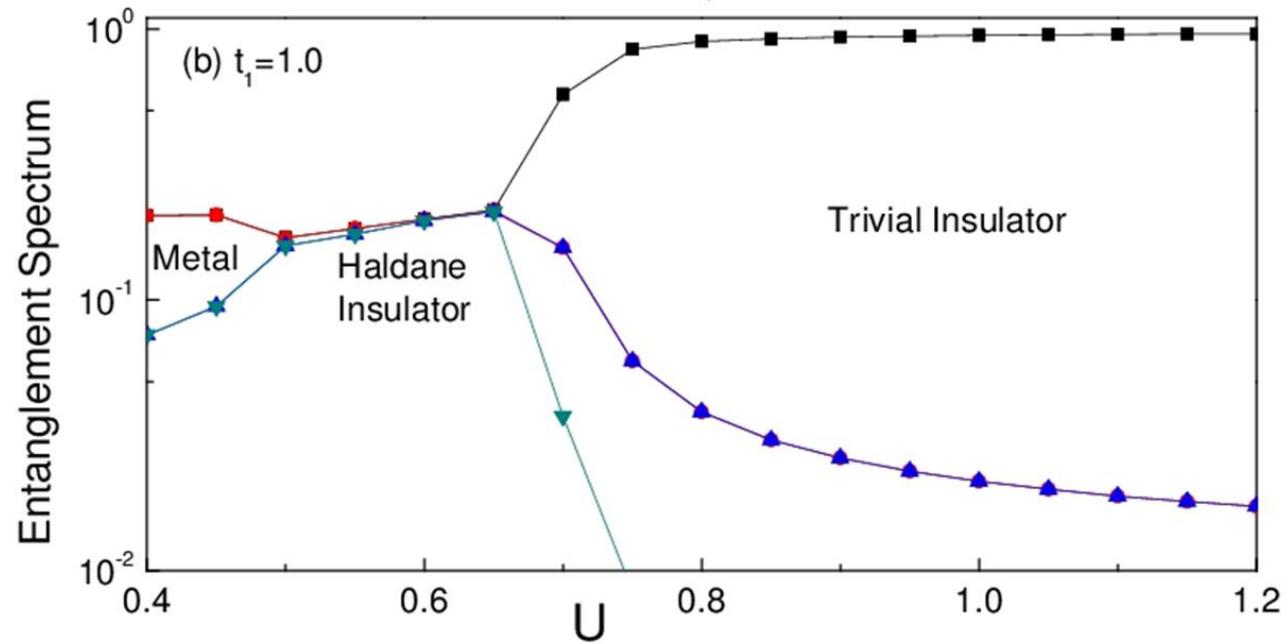
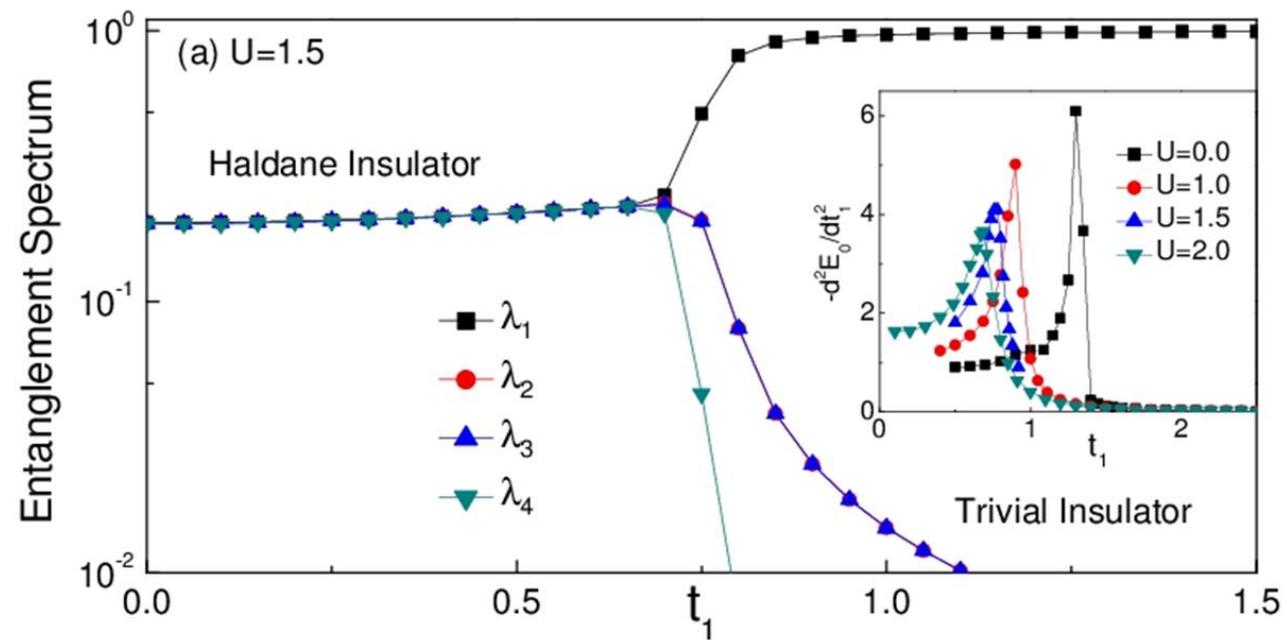
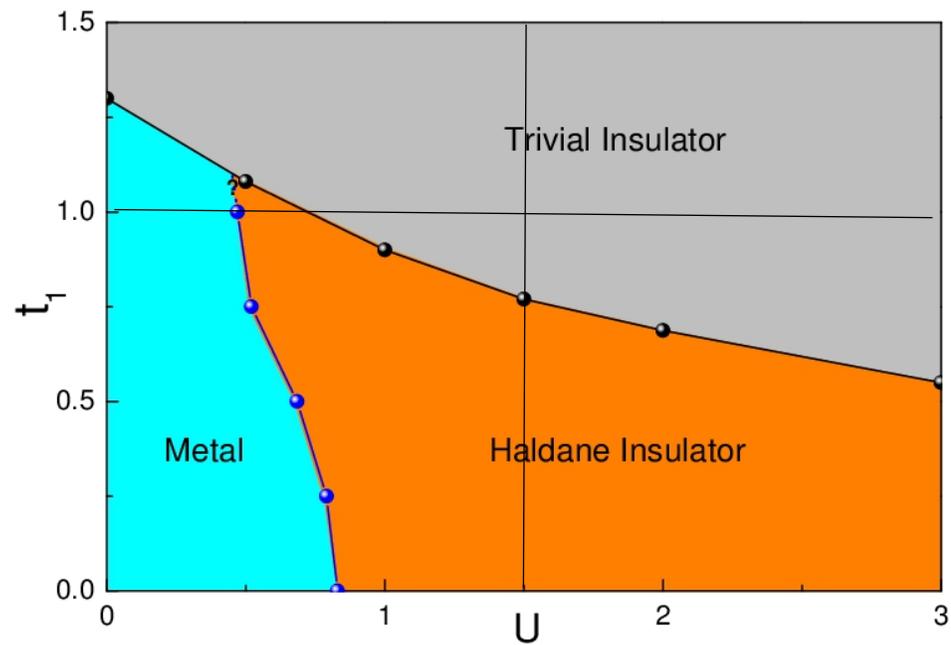


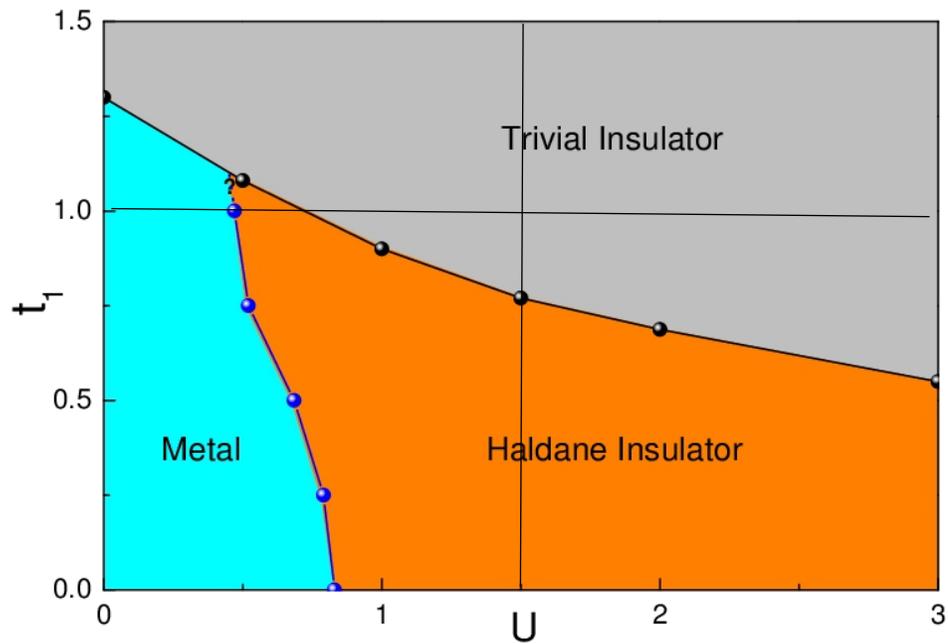




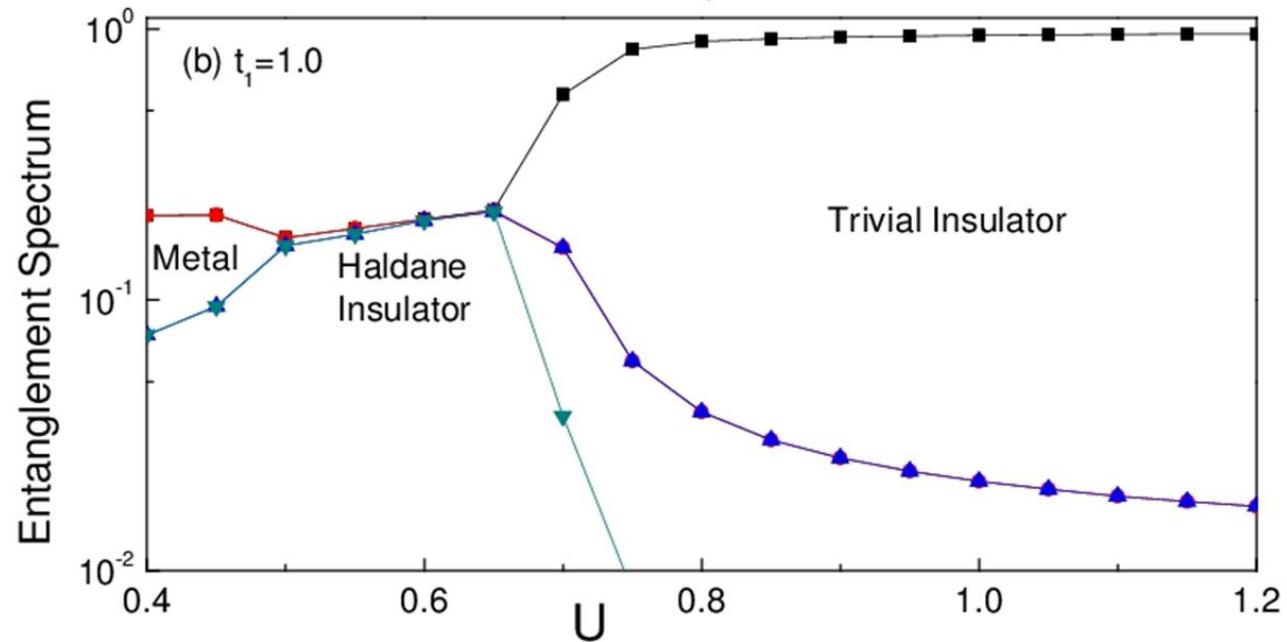
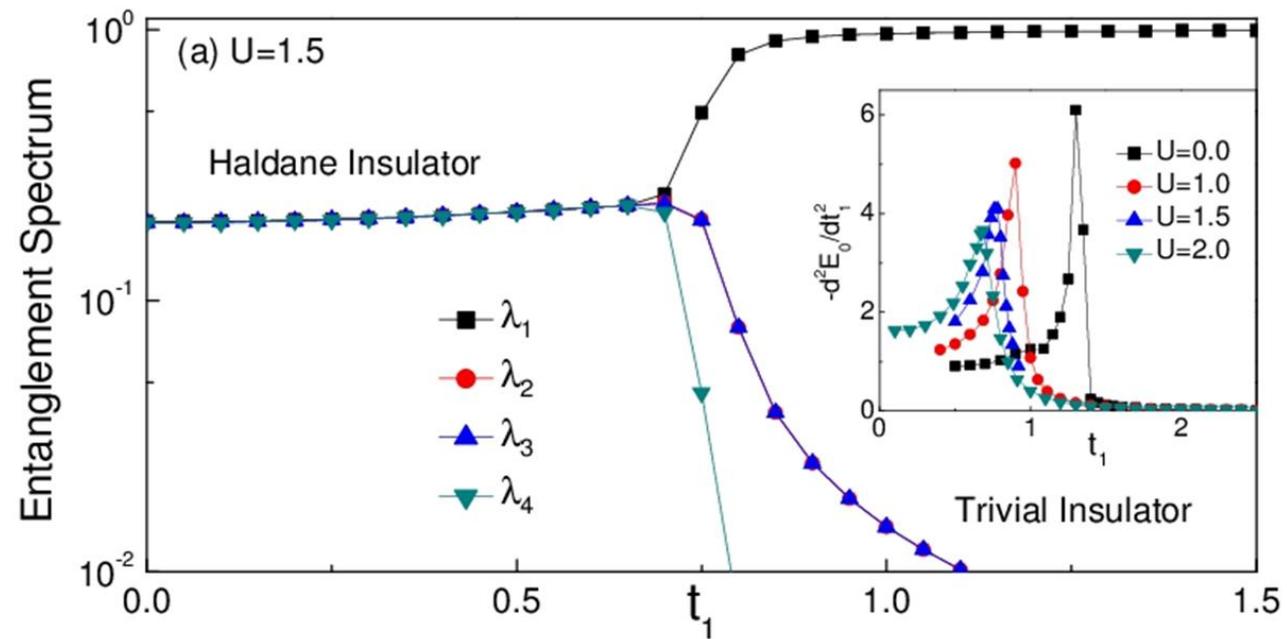








We find two phases in the interacting case:
trivial insulator and Haldane insulator



Conclusion

- introduction of SPT
- Interaction can induce new SPT phases for fermionic systems.

Thanks for your attention!