

TMDs STUDIES WITH A FIXED TARGET EXPERIMENT USING THE LHC BEAMS



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OUTLINE

- ❑ Why a fixed-target experiment at the LHC ?
- ❑ LHC beam extraction using a bent crystal
- ❑ Luminosities
- ❑ Physics highlights for AFTER@LHC
- ❑ TMDs studies with AFTER@LHC
- ❑ First simulations

WHY A FIXED TARGET EXPERIMENT AT THE LHC ?

❑ **Several advantages of the fixed-target mode wrt to the collider mode**

- Accessing the high Feynman x_F domain ($x_F = p_z/p_{zmax}$)
- Achieving **high luminosities** thanks to dense targets
- Easier to **change the target** type (\neq atomic mass)
- Possibility to **polarize the target**

→ Open the possibility for a spin physics program!

❑ **Without affecting the LHC performances**

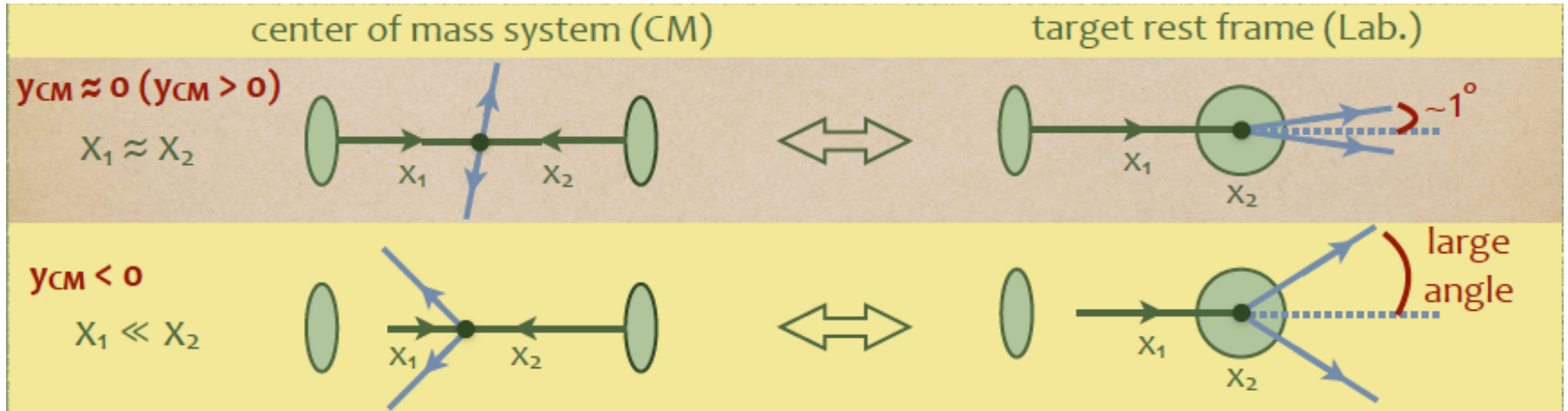
→ Recycling of the beam loss by inserting a bent crystal in the halo (7σ) of the LHC beam

- ❑ With an outstanding luminosity, yet without pile-up
- ❑ With modern detection techniques
- ❑ Virtually no limit on particle-species studies (except top quark)

AFTER@LHC would definitely be a unique experiment

WHY A FIXED TARGET EXPERIMENT AT THE LHC ?

- Provide a novel testing ground for QCD in the high x frontier: $x = [0.3-1]$



- Entire CM forward hemisphere ($y_{CM} > 0$) within $0^\circ < \theta_{lab} < 1^\circ$ (small detector and high multiplicities \rightarrow large occupancies)
- Backward physics** ($y_{CM} < 0$) : larger angle in the laboratory frame (low occupancies, no constraint from beam pipe). **Access to parton with momentum fraction $x_2 \rightarrow 1$ in the target**

7 TeV proton beam on a fixed target

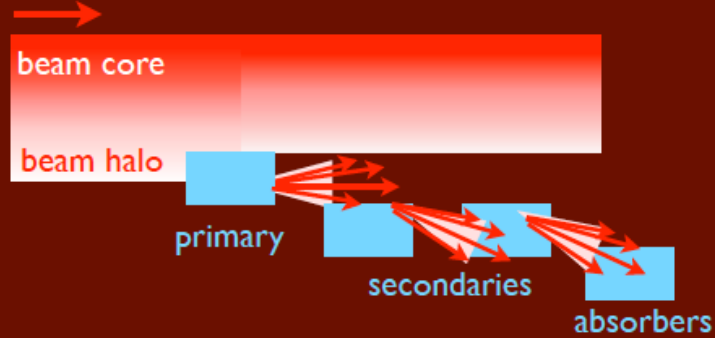
CMS energy: $\sqrt{s} = \sqrt{2m_N E_p} \approx 115 \text{ GeV}$	Rapidity shift: $y_{CM} = 0 \rightarrow y_{lab} = 4.8$
Boost: $\gamma = \sqrt{s} / (2m_p) \approx 60$	

2.76 TeV Pb beam on a fixed target

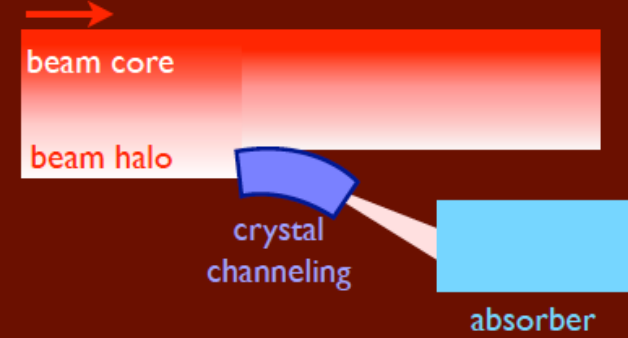
CMS energy: $\sqrt{s_{NN}} = \sqrt{2m_N E_{Pb}} \approx 72 \text{ GeV}$	Rapidity shift: $y_{CM} = 0 \rightarrow y_{lab} = 4.3$
Boost: $\gamma \approx 40$	

LHC BEAM EXTRACTION USING A BENT CRYSTAL

standard collimation



crystal-based collimation (ideally)



Standard collimation today

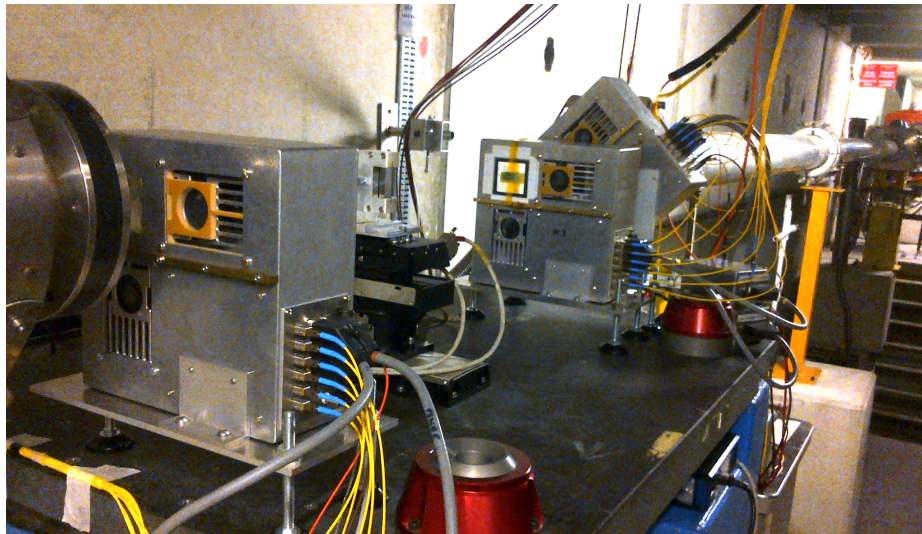


Crystal-based collimation
- UA9 (@SPS)
- LUA9 (@LHC)

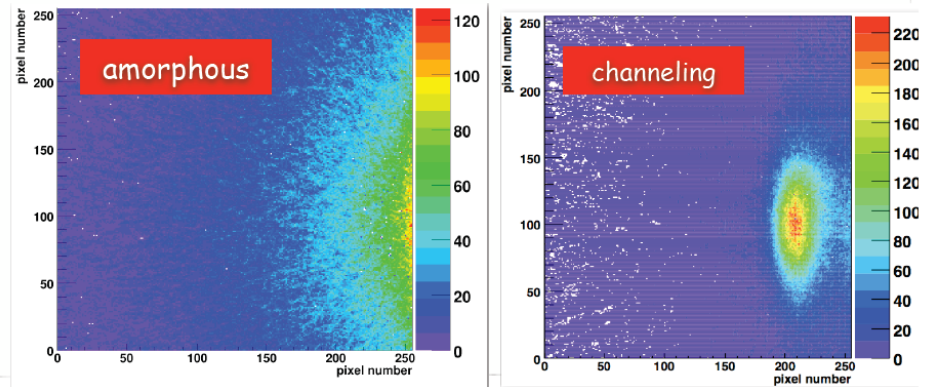


To beam extraction
- CRYSBEAM (@SPS then LHC)
- AFTER@LHC

UA9 experiment @ SPS, 15/10/2014



Direct view of the channeled beam



S. Montesano, W. Scandale, Joint LUA9-AFTER meeting, Nov. 2013

LHC BEAM EXTRACTION USING A BENT CRYSTAL

Beam collimation @LHC: amorphous collimator, inefficiency of 0.2% (3.5TeV p beam)
→ Expected bent crystal inefficiency : 0.02%

UA9: test @SPS on the crystal with proton and ion beams

LUA9 (beam bending experiment using crystal): approved by LHCC

- 2 bent crystals installed in IR7 during LS1
- 2015/2016 first tests with beams

Proton beam extraction:

- Single or multi-pass extraction efficiency of 50%
- LHC beam loss $\sim 10^9$ p/s → Extracted beam : 5×10^8 p/s
- Extremely small emittance:
→ beam size (in the extraction direction) 950m after the extraction: 0.3mm

Ion beam extraction:

- Successfully tested at the SPS
- Should also work at the LHC

LUMINOSITIES IN pH, pA @ $\sqrt{s_{NN}} = 115 \text{ GeV}$ AND Pb-A $\sqrt{s_{NN}} = 72 \text{ GeV}$

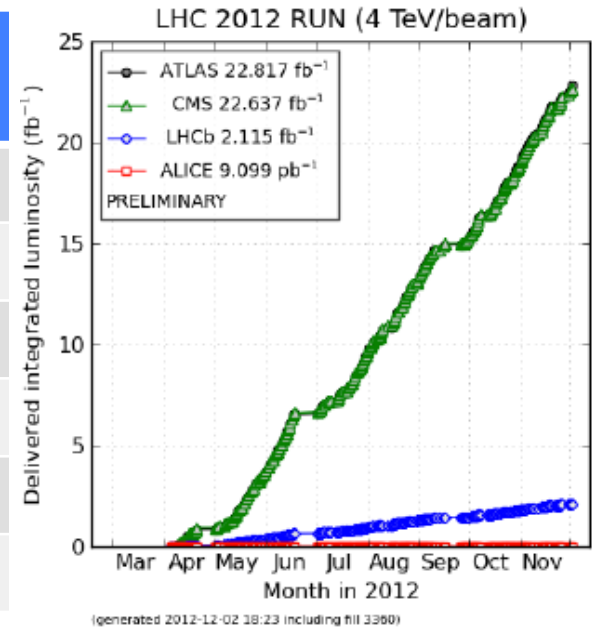
Instantaneous luminosity: $L = \phi_{\text{beam}} \times N_{\text{target}} = \phi_{\text{beam}} \times (\rho \times l \times N_A) / A$

With l target thickness, $\phi_{\text{beam}} = 5 \times 10^8 p^+ s^{-1}$ (1/2 of the beam loss)

Integrated luminosity: assuming 10^7 s of p beam and 10^6 s of Pb beam (LHC year)

In pH and pA (115 GeV/c)

Target	ρ (g.cm ⁻³)	A	L ($\mu\text{b}^{-1} \cdot \text{s}^{-1}$)	$\int L$ ($\text{pb}^{-1} \cdot \text{yr}^{-1}$)
Liq H ₂ (1m)	0.07	1	2000	20000
Liq D ₂ (1m)	0.16	2	2400	24000
Be (1cm)	1.85	9	62	620
Cu (1cm)	8.96	64	42	420
W (1cm)	19.1	185	31	310
Pb (1cm)	11.35	207	16	160



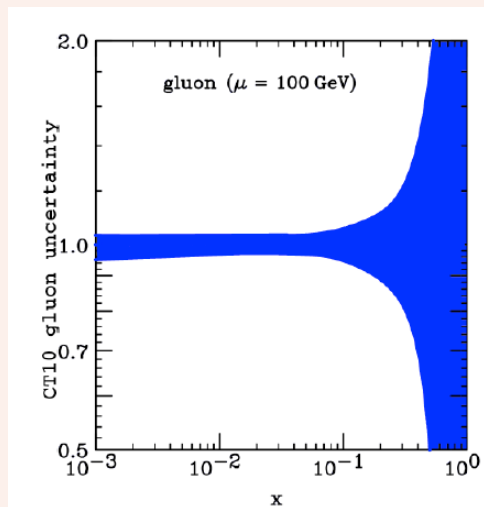
Luminosity comparable to the LHC itself (with 1m long H₂ (D₂) target)
 3 orders of magnitude larger than PHENIX@RHIC
 (run14pp@200GeV: 12pb⁻¹, run14dAu@200GeV: 0.15pb⁻¹)

PHYSICS HIGHLIGHTS FOR AFTER@LHC

p-p and p-A @ $\sqrt{s_{NN}} = 115$ GeV

Nucleon partonic structure

- Gluon PDF in the proton
→ large uncertainty at large x. DIS not ideal
- $g_n(x) = g_p(x)$?
- Experimental probes: quarkonia, isolated photons, high p_T jets
- Multiple probes essential to check factorization



Heavy quark distribution at large x in the proton

- Intrinsic charm?
- Experimental probes: open heavy flavors

Spin Physics

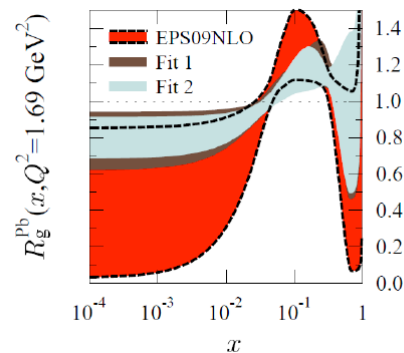
Gluon Sivers effect
Linearly polarized gluons: $h_1^{\perp g}$
Single Spin Asymetry in HF and DY studies

W and Z bosons production near threshold?

Pb-A @ $\sqrt{s_{NN}} = 72$ GeV

Gluon distribution in nucleus at large x

- Large uncertainty at high x
- EIC, LHeC experiments do not help much



Quark Gluon Plasma

- Υ sequential suppression
- Quarkonium excited state suppression
 - Jet-HF quenching
 - Direct photons

Ultra-peripheral collisions

More details in [Physics Reports 522 \(2013\) 239](#)

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Physics opportunities of a fixed-target experiment using LHC beams

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TMDs STUDIES IN AFTER@LHC (WITH A POLARIZED TARGET)

See also presentation by J.P. Lansberg, Gluon TMDs and Quarkonium production in unpolarised and polarised Proton-Proton collisions, on Monday 20th October, s1, 16:30

(Gluon) Sivers effects with a transversely polarized target

Gluon Sivers effect: correlation between the gluon transverse momentum k_T and the proton spin

- ❑ **The target rapidity region ($x_F < 0$) corresponds to high x^\uparrow ($x_F \rightarrow -1$) where the k_T - spin correlation is the largest**
- ❑ Transverse single spin asymmetries studied using **gluon sensitives probes**:
 - quarkonia (J/ψ , Y , χ_c)
 - B & D mesons production
 - γ , γ -jet, γ - γ also J/ψ - γ

TMDs STUDIES WITH AFTER@LHC (WITH A POLARIZED TARGET)

(Quark) Sivers effects with a transversely polarized target

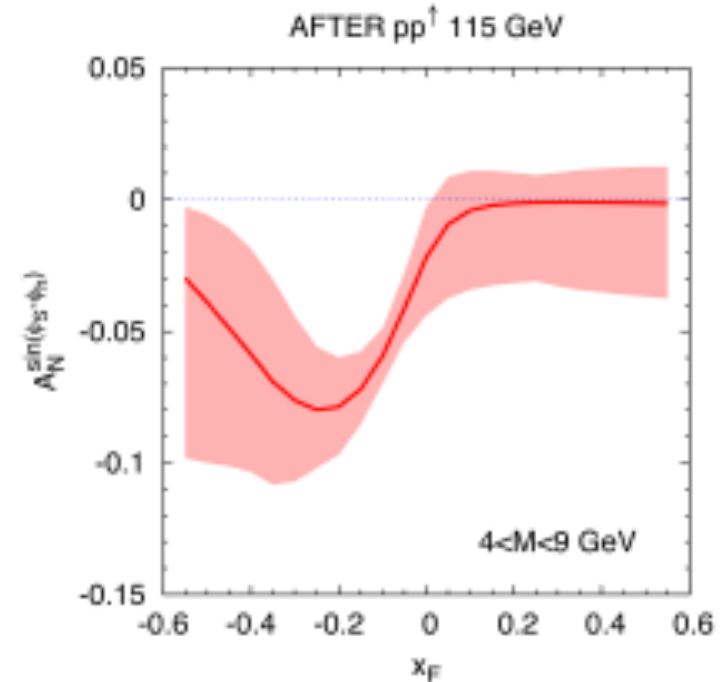
□ Can be probed with the Drell-Yan process

Experiment	particles	energy (GeV)	\sqrt{s} (GeV)	x_p^\uparrow	\mathcal{L} ($\text{nb}^{-1}\text{s}^{-1}$)
AFTER	$p + p^\uparrow$	7000	115	0.01 \div 0.9	1
COMPASS	$\pi^\pm + p^\uparrow$	160	17.4	0.2 \div 0.3	2
COMPASS (low mass)	$\pi^\pm + p^\uparrow$	160	17.4	\sim 0.05	2
RHIC	$p^\uparrow + p$	collider	500	0.05 \div 0.1	0.2
J-PARC	$p^\uparrow + p$	50	10	0.5 \div 0.9	1000
PANDA (low mass)	$\bar{p} + p^\uparrow$	15	5.5	0.2 \div 0.4	0.2
PAX	$p^\uparrow + \bar{p}$	collider	14	0.1 \div 0.9	0.002
NICA	$p^\uparrow + p$	collider	20	0.1 \div 0.8	0.001
RHIC	$p^\uparrow + p$	250	22	0.2 \div 0.5	2
Int.Target 1					
RHIC	$p^\uparrow + p$	250	22	0.2 \div 0.5	60
Int.Target 2					
P1027	$p^\uparrow + p$	120	15	0.35 \div 0.85	400-1000
P1039	$p + p^\uparrow$	120	15	0.1 \div 0.3	400-1000

Relevant parameters for the future proposed polarized DY experiments

S. J. Brodsky et al., Phys. Rep. 522 (2013) 239

V. Barone et al., Prog. Part. Nucl. Phys. 65 (2010) 267



Prediction for AFTER

M. Anselmo, ECT*, Feb. 2013
(Courtesy U. d'Alessio)

Asymmetry up to 10% predicted in DY for the target rapidity region ($x_F < 0$)

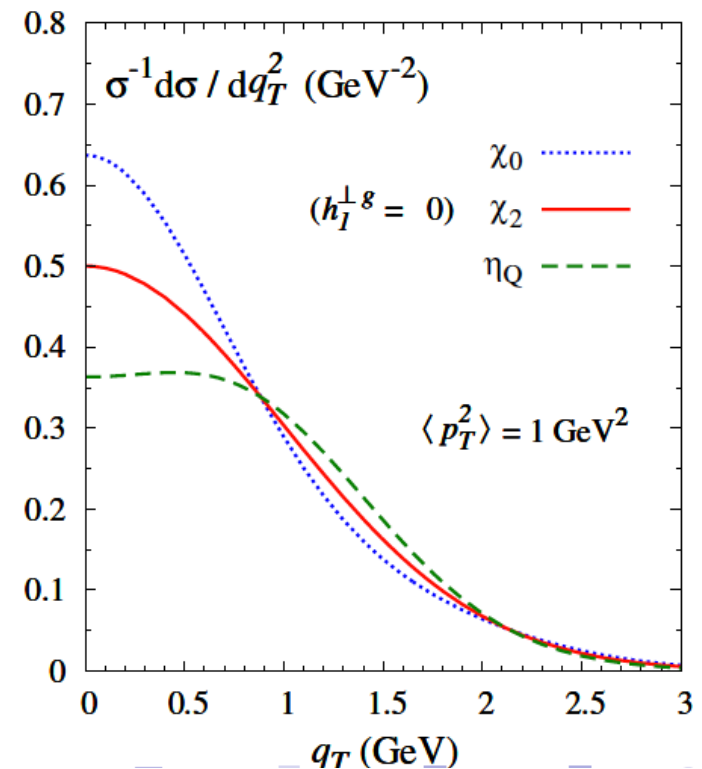
TMDs STUDIES WITH AFTER@LHC (WITHOUT A POLARIZED TARGET)

Access to the distribution of linearly polarized gluons ($h_1^{\perp g}$)

« Boers-Mulder » effect: correlation between the parton k_T and its spin
 For gluons, it is encoded in $h_1^{\perp g}$



- Low- p_T C-even quarkonium production is a good probe of the gluon TMDs.
- The low- p_T spectra of scalar and pseudo-scalar quarkonium (χ_{c0} , χ_{b0} , η_c , η_b) are affected differently by the linearly polarized gluons in unpolarized nucleons
 - Boost: better access to low- p_T C-even quarkonia
 - Still challenging experimentally (first study of η_c in collider by LHCb for $p_T > 6$ GeV/c)
 - [LHCb 1409.3612](#)
 - If possible somewhere, it is at AFTER@LHC
- Back-to-back $J/\psi + \gamma$ is also a good probe of gluon TMDs



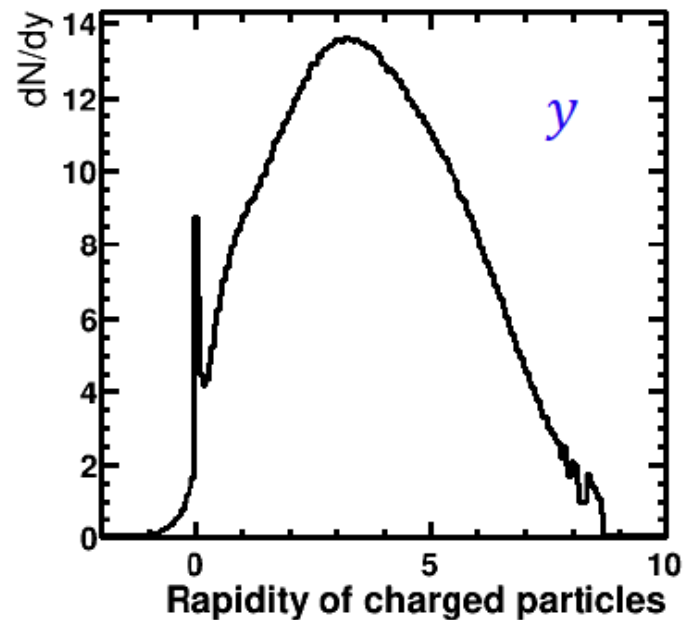
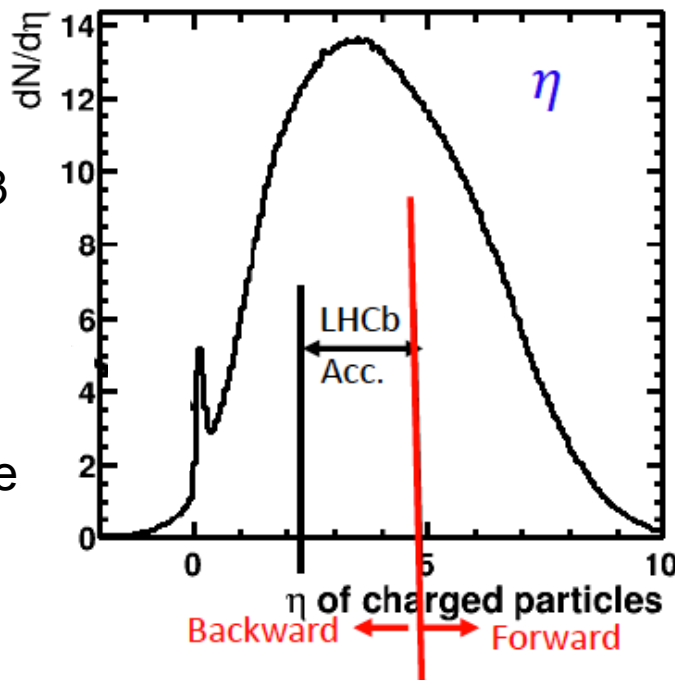
Boer, Pisano, PRD 86 (2012) 094007

Simulations of a 7 TeV proton beam on a Pb target ($\sqrt{s_{NN}} = 115 \text{ GeV}$)

- ❑ Full LHCb simulation and standard reconstruction
- ❑ Study the resolution at vertex, the occupancy in the pixels...
- ❑ Compare multiplicities in AFTER with LHCb pA run

- ❑ Simulations with HIJING version 1.383bs.2
- ❑ 10 000 events generated, no pile-up

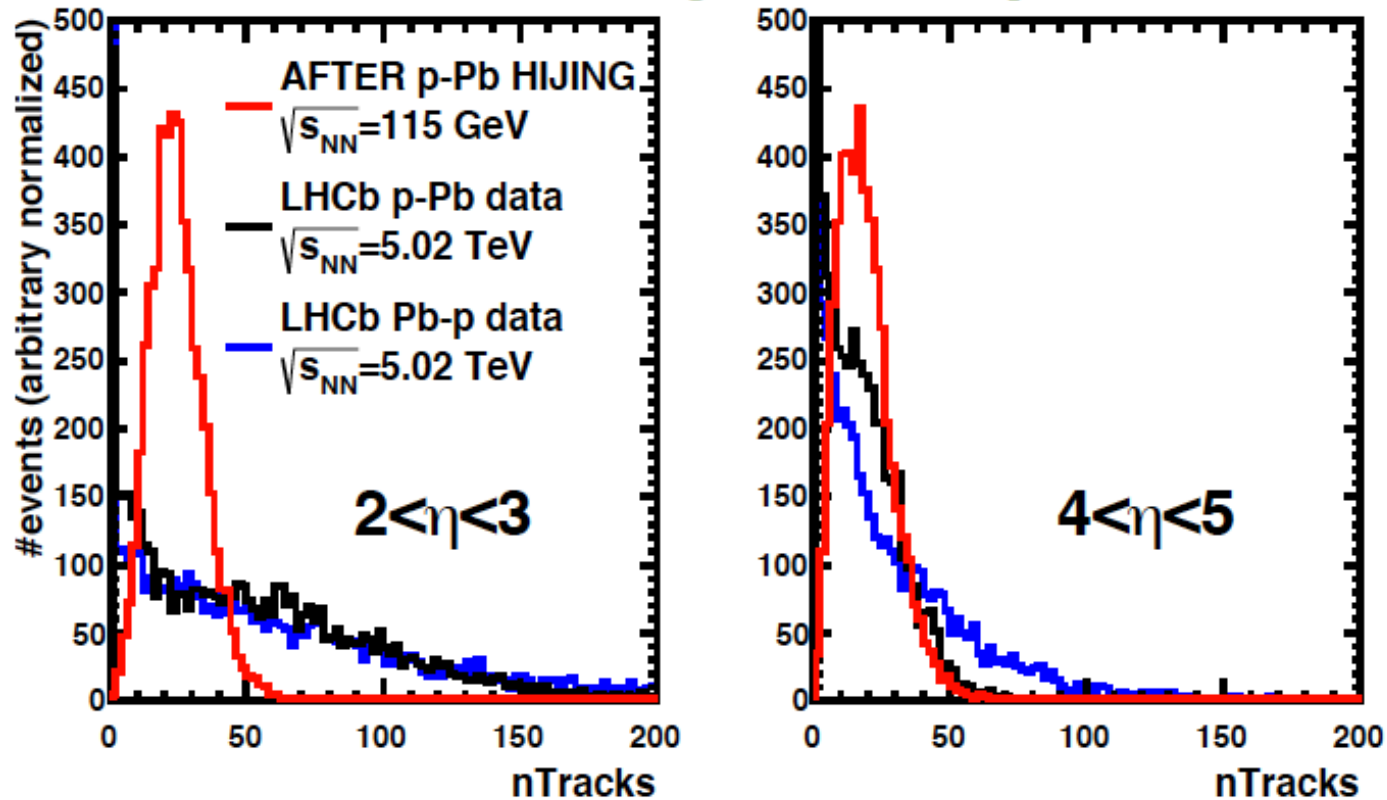
Z. Yang, AFTER workshop les Houches, January 2014



$dN/d\eta \sim 13$
Generated
charged
particles in
LHCb
acceptance

Simulations of a 7 TeV proton beam on a Pb target ($\sqrt{s_{NN}} = 115$ GeV)

Z. Yang, AFTER workshop les Houches, January 2014



- ❑ Probability for high track multiplicity in AFTER (pPb @ 115 GeV) is lower than the one measured by LHCb (pPb/Pbp @ 5.02 TeV)
- ❑ No problem for a LHCb-like detector to cope with the multiplicity of pPb collisions at $\sqrt{s_{NN}} = 115$ GeV in $2 < \eta < 5$

QUARKONIUM CASE

Expected quarkonium yields

In pA and pA (115 GeV)

Target	$\int \mathcal{L} \text{ (fb}^{-1}\text{,yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ = $A\mathcal{L}B\sigma_\Psi$	$N(\Upsilon) \text{ yr}^{-1}$ = $A\mathcal{L}B\sigma_\Upsilon$
1 m Liq. H ₂	20	4.0 10 ⁸	8.0 10 ⁵
1 m Liq. D ₂	24	9.6 10 ⁸	1.9 10 ⁶
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 10 ⁷ 1.4 10 ⁹	1.8 10 ⁵ 7.2 10 ⁶
RHIC pp 200GeV	1.2 10 ⁻²	4.8 10 ⁵	1.2 10 ³

Target	$\int \mathcal{L} \text{ (fb}^{-1}\text{,yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ = $A\mathcal{L}B\sigma_\Psi$	$N(\Upsilon) \text{ yr}^{-1}$ = $A\mathcal{L}B\sigma_\Upsilon$
1cm Be	0.62	1.1 10 ⁸	2.2 10 ⁵
1cm Cu	0.42	5.3 10 ⁸	1.1 10 ⁶
1cm W	0.31	1.1 10 ⁹	2.3 10 ⁶
1cm Pb	0.16	6.7 10 ⁸	1.3 10 ⁶
LHC pPb 8.8 TeV	10 ⁻⁴	1.0 10 ⁷	7.5 10 ⁴
RHIC dAu 200GeV	1.5 10 ⁻⁴	2.4 10 ⁶	5.9 10 ³
RHIC dAu 62GeV	3.8 10 ⁻⁶	1.2 10 ⁴	18

pp : 1000 times more statistics than at RHIC ($\sqrt{s} = 200 \text{ GeV}$) and comparable statistics to LHCb with a 1m H₂ target
 pA: 100 times more statistics than at RHIC (dAu $\sqrt{s} = 200 \text{ GeV}$) with a 1cm Pb target

In PbA (72 GeV)

Target	$\int \mathcal{L} \text{ (nb}^{-1}\text{,yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_\Psi$	$N(\Upsilon) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_\Upsilon$
1 m Liq. H ₂	800	3.4 10 ⁶	6.9 10 ³
1cm Be	25	9.1 10 ⁵	1.9 10 ³
1cm Cu	17	4.3 10 ⁶	0.9 10 ³
1cm W	13	9.7 10 ⁶	1.9 10 ⁴
1cm Pb	7	5.7 10 ⁶	1.1 10 ⁴
LHC PbPb 5.5 TeV	0.5	7.3 10 ⁶	3.6 10 ⁴
RHIC AuAu 200GeV	2.8	4.4 10 ⁶	1.1 10 ⁴
RHIC AuAu 62GeV	0.13	4.0 10 ⁴	61

PbA: similar statistics as at RHIC (Au-Au $\sqrt{s_{NN}} = 200 \text{ GeV}$) and 2 orders of magnitude larger than at RHIC (Au-Au $\sqrt{s_{NN}} = 62 \text{ GeV}$) with a 1cm thick Pb target

Detailed study of quarkonium production and nuclear effects

Detailed study of quarkonium states

FAST SIMULATIONS FOR QUARKONIA ($pp \sqrt{s} = 115 \text{ GeV}$) USING LHCb RECONSTRUCTION PARAMETERS

- ❑ Simulations with Pythia 8.185
- ❑ LHCb detector is NOT simulated but LHCb reconstruction parameters are introduced in the fast simulation (resolution, analysis cuts, efficiencies...)

Requirements

Momentum resolution : $\Delta p/p = 0.5\%$

Muon identification efficiency: 98%

Cuts at the single muon level

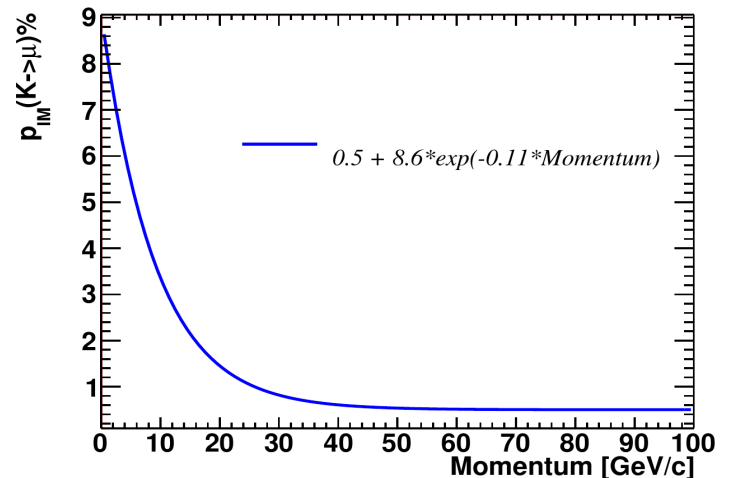
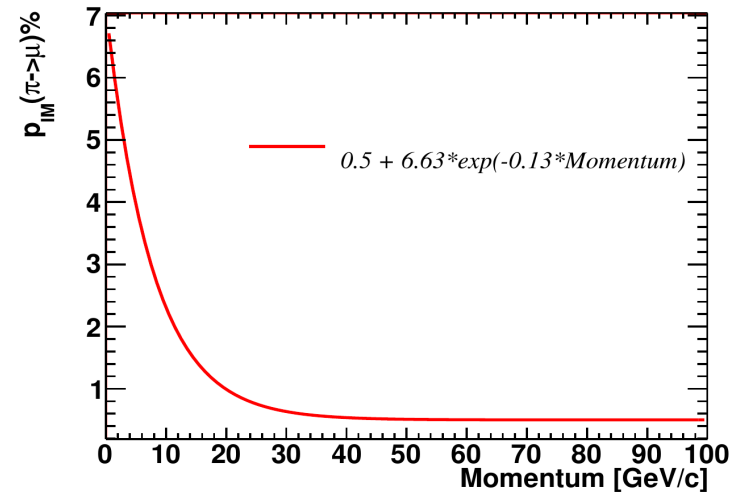
$2 < \eta_{\mu} < 5$

$p_{T}^{\mu} > 0.7 \text{ GeV}/c$

Muon misidentification

If π and K decay before the calorimeters (12m), they are rejected by the tracking
Otherwise a misidentification probability is applied

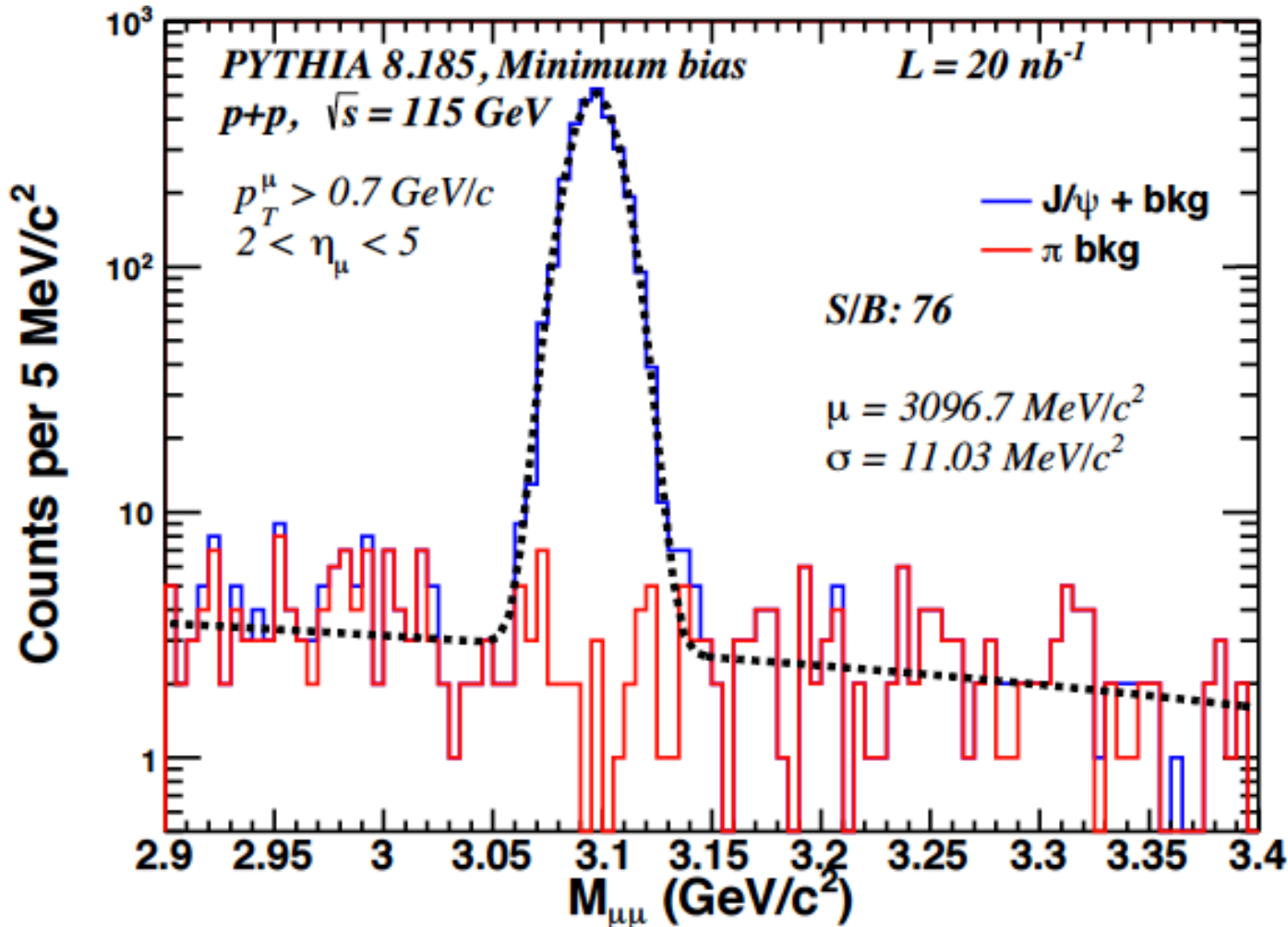
Performance of the muon identification at LHCb,
F. Achilli et al, arXiv:1306.0249



J/ψ → μ⁺μ⁻ IN MINIMUM BIAS pp COLLISIONS @ 115 GeV

□ For 1m of H target and 10 seconds of data taking!

B. Trzeciak, July 2014, Orsay

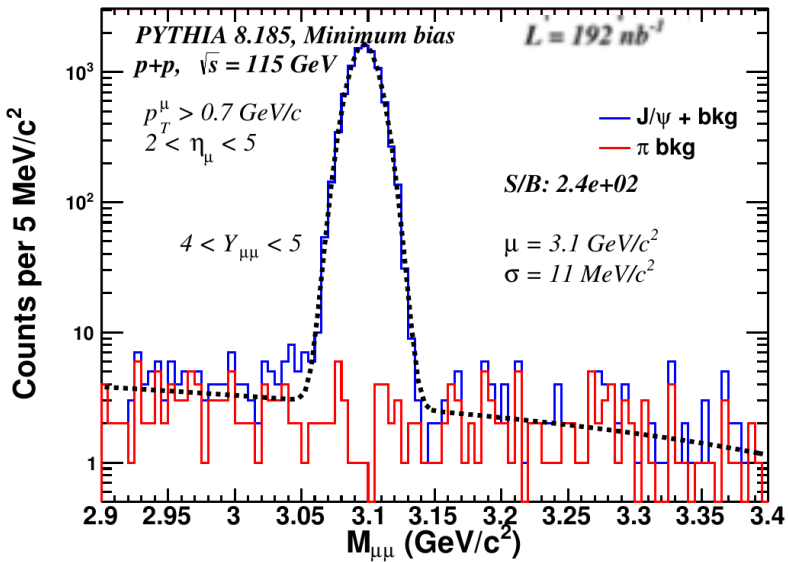
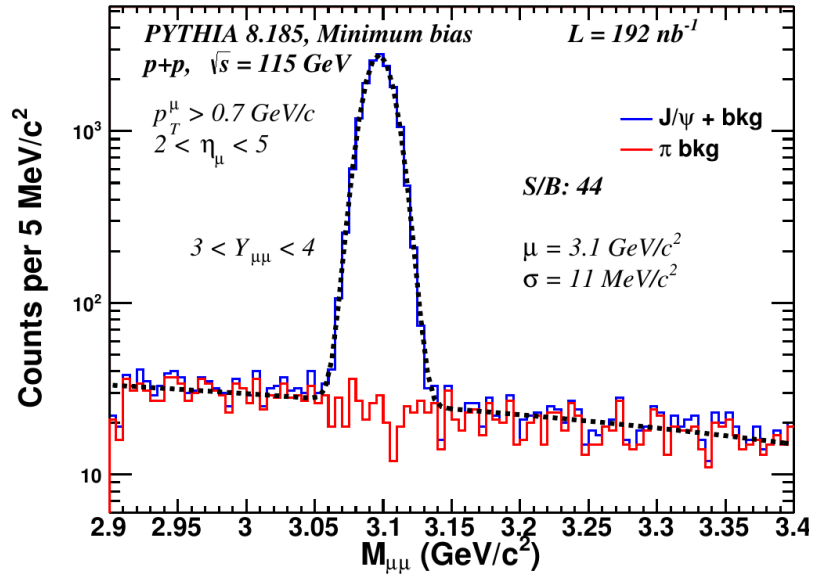
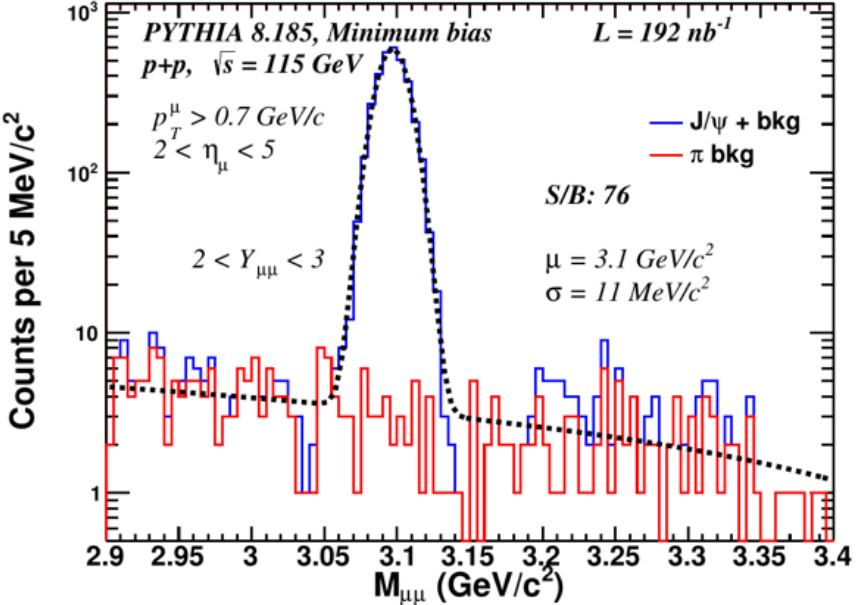


Misidentified pions is the dominant source of background

J/ψ → μ⁺μ⁻ IN MINIMUM BIAS pp COLLISIONS @ 115 GEV (BINS IN RAPIDITY)

□ For 1m of H target and 1.5 minute of data taking

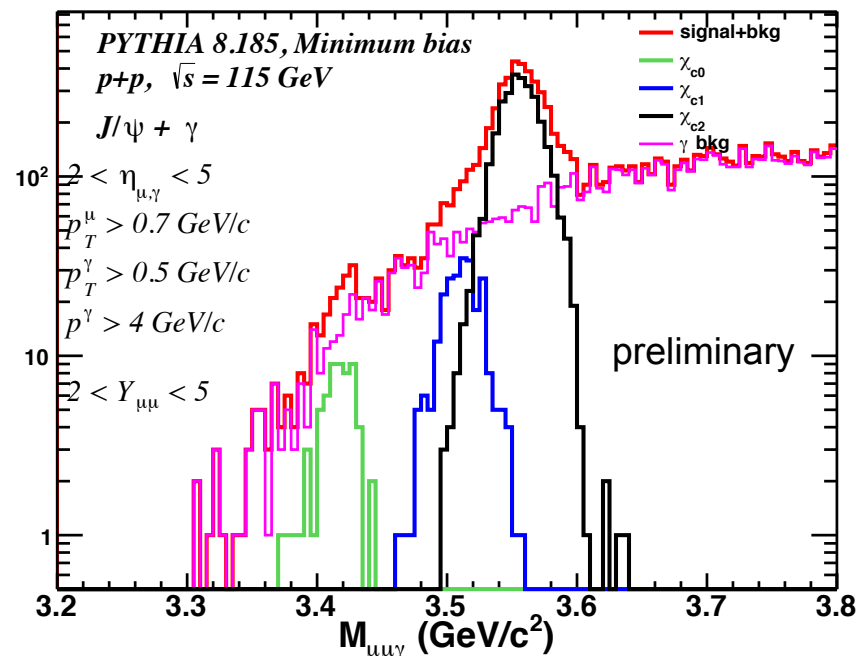
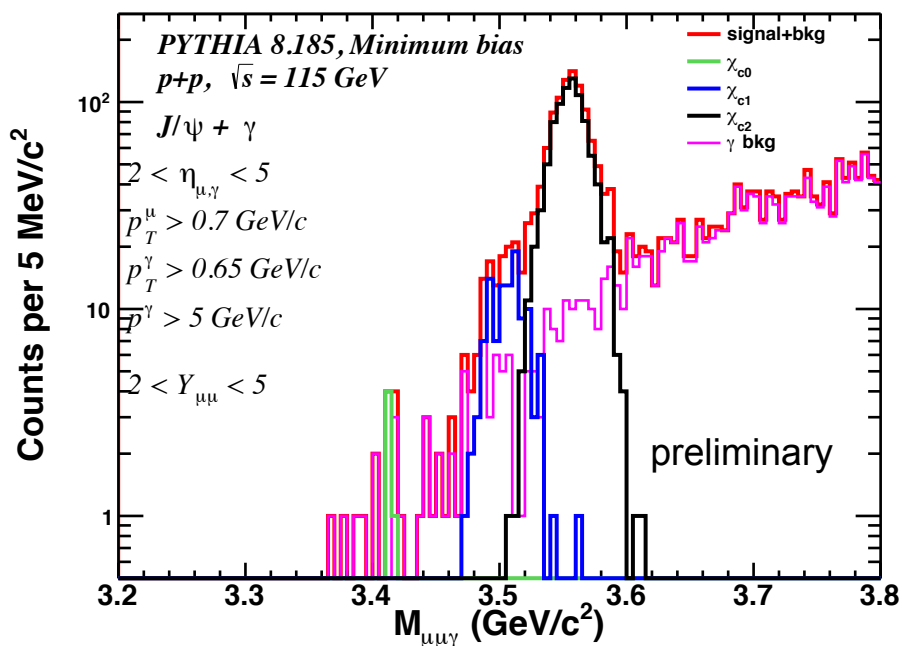
B. Trzeciak, July 2014, Orsay



$\chi_c \rightarrow \mu^+ \mu^- \gamma$ IN MINIMUM BIAS pp COLLISIONS @ 115 GEV

- Preliminary studies of the χ_c also started
- $L_{\text{int}} = 192 \text{ nb}^{-1}$, for 1m of H target and 1.5 minute of data taking

B. Trzeciak, July 2014, Orsay



- Hope to be able to reach p_T of the χ_c down to 0
- χ_{c1} and χ_{c2} separation

SUMMARY

- ❑ AFTER@LHC provides a novel testing ground for QCD in the high x frontier
- ❑ High luminosities are achievable in pp, pA @ 115 GeV and PbA @ 72 GeV using dense targets and without affecting the LHC beam
- ❑ TMDs studies can be performed with/without polarizing the target thanks to e.g. low- p_T quarkonia
- ❑ First fast simulations with LHCb like setup are promising

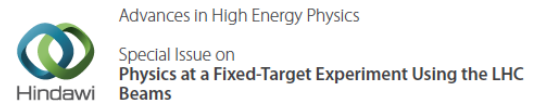
- ❑ What's next :

- Special issue in Advances in High Energy Physics (submission deadline in March 2015)

Everybody is welcome to contribute

- Expression of interest expected in 2015
- AFTER week @ CERN, 17-21 November 2014

- ❑ The AFTER web page : after.in2p3.fr



CALL FOR PAPERS

Fixed-target experiments (FTE) have brought essential contributions to particle and nuclear physics. They have led to particle discoveries (Ω , Λ , ...) and evidence for the novel dynamics of quarks and gluons in heavy-ion collisions. In accessing high x_F and in offering options for (un-) polarized proton and nuclear targets, they have also led to the observation of surprising QCD phenomena. They offer specific advantages compared to collider experiments: access to high x_F , high luminosities, target versatility, and polarisation.

The LHC 7 TeV protons on targets release a c.m.s. energy close to 115 GeV (72 GeV with Pb), in a range never explored so far, significantly higher than that at SPS and far from RHIC. The production of quarkonia, DY, heavy flavours, jets, and γ in pA collisions can be studied with statistics previously unheard of and in the backward region, $x_F < 0$, which is uncharted. High precision QCD measurements can also obviously be carried out in pp and pd collisions with H_1 and D_1 targets. With the 50 TeV protons of the future circular collider (FCC), the c.m.s. energy could reach 300 GeV for original studies of W and Z bosons, and perhaps H^0 , production in pp and pA collisions.

With the LHC Pb beam, one can study the quark-gluon plasma (QGP) from the viewpoint of the nucleus rest frame after its formation. Thanks to modern technologies, studies of, for instance, direct γ and quarkonium P-waves production in heavy-ion collisions can be envisioned.

Polarising the target allows one to study single-spin correlations including the Sivers effect, hence, the correlation between the parton k_T and the nucleon spin.

We intend to publish a special issue on the physics at such a FTE using the LHC or FCC beams. The editors welcome original research articles and review articles from both theorists and experimentalists.

Potential topics include, but are not limited to:

- Heavy-quark and gluon content at large x
- TMDs and single-spin asymmetries
- Heavy-flavour studies in pA and AA collisions at FTEs
- W, Z, and H^0 production near threshold
- Target polarisation
- Secondary beams
- Simulation tools for high-energy physics
- Beam collimation and extraction with bent crystals
- Machine feasibility and radiological aspects
- Connection between UHEICR studies and FTEs

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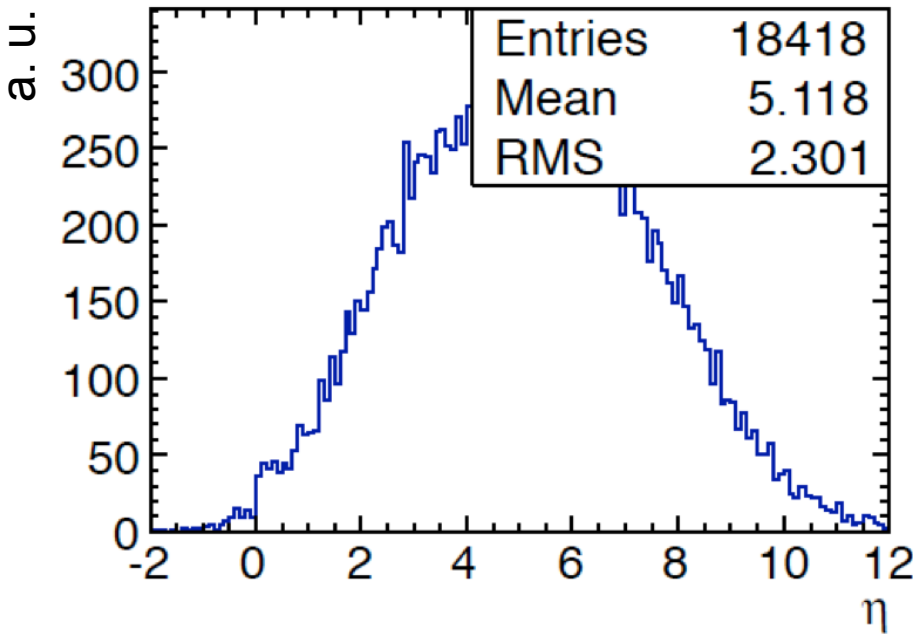
First Round of Reviews
Friday, 12 June 2015

Publication Date
Friday, 7 August 2015

BACK UP

NUMBER OF CHARGED PARTICLES IN MB pp @ $\sqrt{s} = 115$ GEV

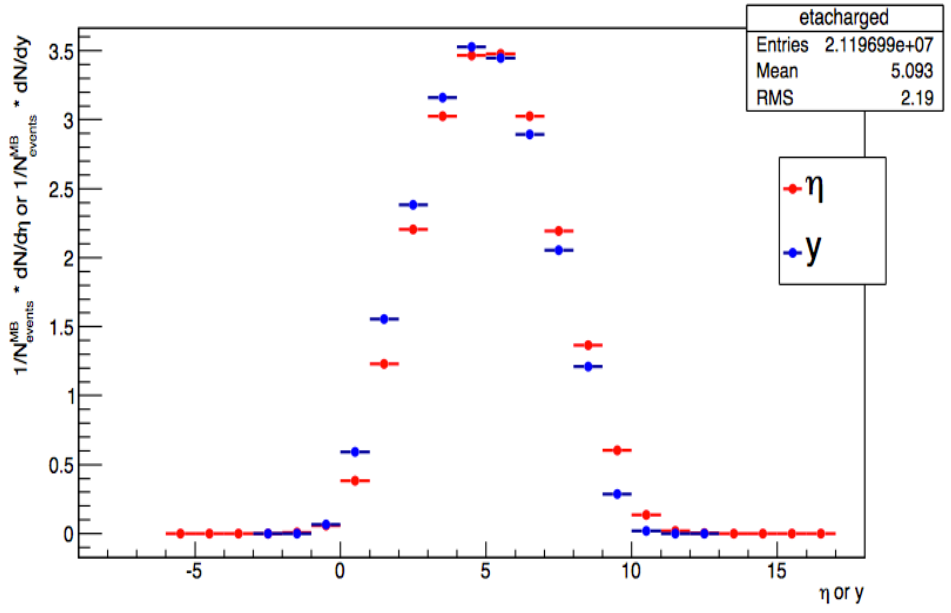
AFTER workshop les Houches, January 2014
 AFTER simulation group



EPOS 1.6.5

Number of generated events: 1000

$$dN_{ch}/d\eta |_{\eta=0} \sim 3$$



PYTHIA 8.170

Number of generated events: 10^6

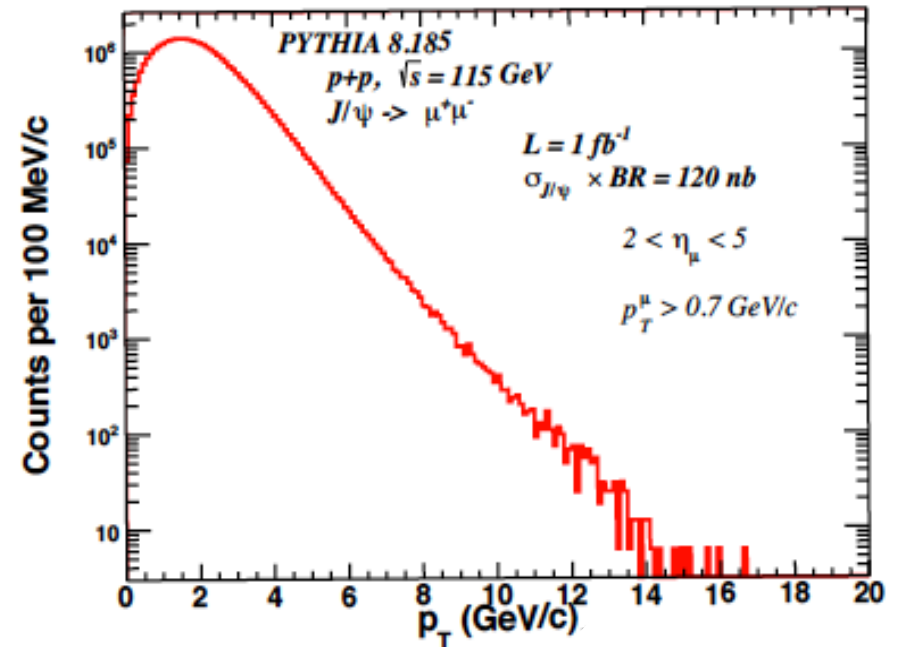
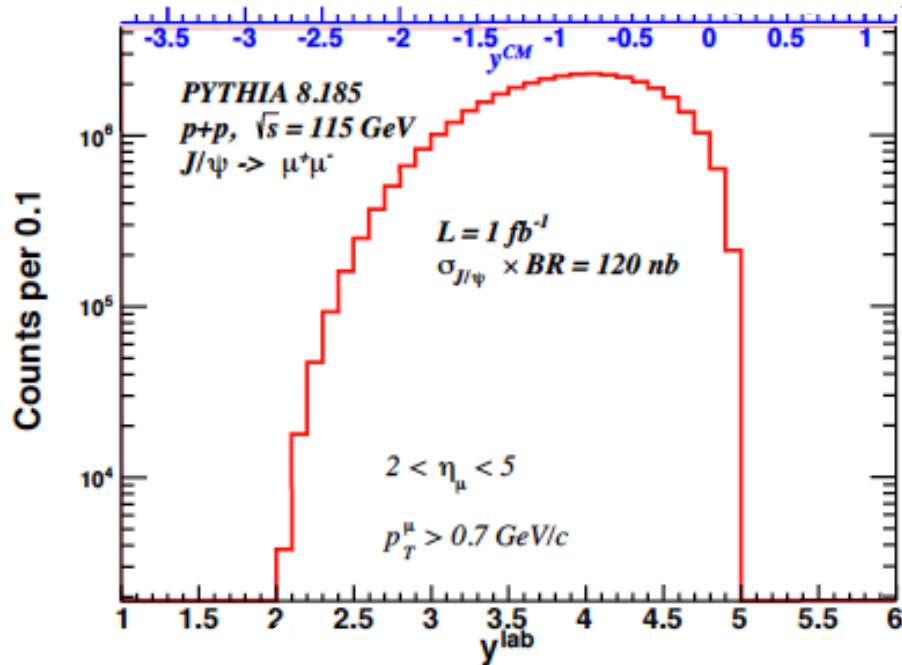
$$dN_{ch}/d\eta |_{\eta=0} \sim 3.5$$

Rapidity shift: $\Delta y = \tan^{-1}\beta \approx 4.8$
 $y_{CM} = 0 \rightarrow y_{lab} \approx 4.8$

$J/\psi \rightarrow \mu^+\mu^-$ IN MB pp @ 115 GEV (Y_{LAB} AND P_{T} REACH)

□ For 1m of H target and 2 weeks of data taking

B. Trzeciak, July 2014, Orsay



Large statistics allow one:

- To reach large p_{T}
- Large y_{lab} acceptance $2 < y_{\text{lab}} < 5$

ACCESSING THE LARGE x GLUON PDF

PYTHIA simulation
 $\sigma(y) / \sigma(y=0.4)$
 statistics for one month
 5% acceptance considered

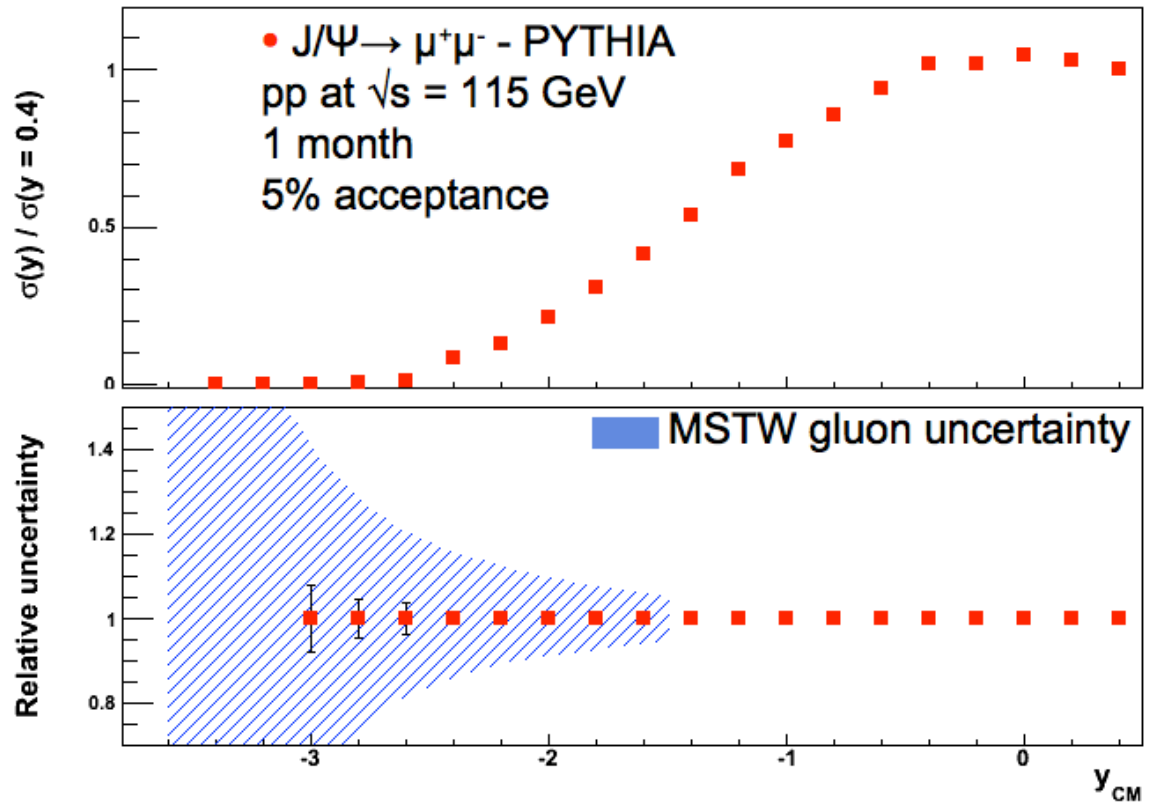
Statistical relative uncertainty
 Large statistics allow to access
 very backward region

Gluon uncertainty from
 MSTWPDF
 - only for the gluon content of
 the target
 - assuming

$$x_g = M_{J/\psi} / \sqrt{s} e^{-y_{CM}}$$

J/ψ
 $y_{CM} \sim 0 \rightarrow x_g = 0.03$
 $y_{CM} \sim -3.6 \rightarrow x_g = 1$

Y: larger x_g for same y_{CM}
 $y_{CM} \sim 0 \rightarrow x_g = 0.08$
 $y_{CM} \sim -2.4 \rightarrow x_g = 1$

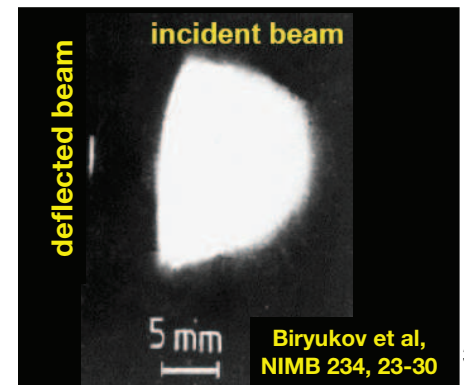
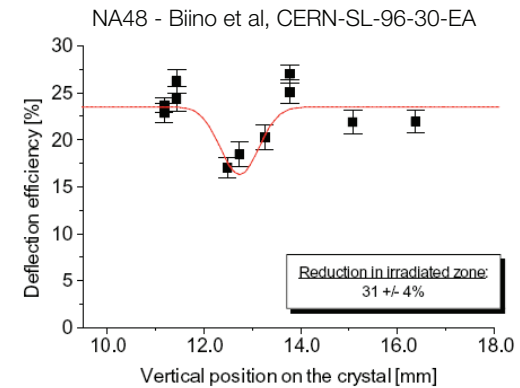


⇒ Backward measurements allow to access large x gluon pdf

Crystal resistance to irradiation



- **IHEP U-70** (Biryukov et al, NIMB 234, 23-30):
 - 70 GeV protons, 50 ms spills of **10^{14} protons every 9.6 s**, several minutes irradiation
 - equivalent to 2 nominal LHC bunches for 500 turns every 10 s
 - 5 mm silicon crystal, **channeling efficiency unchanged**
- **SPS North Area - NA48** (Biino et al, CERN-SL-96-30-EA):
 - 450 GeV protons, 2.4 s spill of 5×10^{12} protons every 14.4 s, one year irradiation, **2.4×10^{20} protons/cm²** in total,
 - equivalent to several year of operation for a primary collimator in LHC
 - $10 \times 50 \times 0.9$ mm³ silicon crystal, 0.8×0.3 mm² area irradiated, **channeling efficiency reduced by 30%**.
- **HRMT16-UA9CRY** (HiRadMat facility, November 2012):
 - 440 GeV protons, up to 288 bunches **in 7.2 μ s**, 1.1×10^{11} protons per bunch (**3×10^{13} protons** in total)
 - energy deposition comparable to an asynchronous beam dump in LHC
 - 3 mm long silicon crystal, **no damage to the crystal after accurate visual inspection**, more tests planned to assess possible crystal lattice damage
 - **accurate FLUKA simulation of energy deposition** and residual dose



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