

# The Twin Higgs: Can Naturalness Hide?

Twin Higgs: Zackaria Chacko, Hock-Seng Goh, RH hep-ph/0506256  
Folded SUSY: Brudman, Chacko, Goh, RH hep-ph/0609152

Liberal use of Nathaniel's work.

Discussions N. Arkani-Hamed, N. Craig, Z. Chacko, G. Perez.



# Where is everybody?

- \* LHC run I:  
We found a Higgs. Nothing else so far.
- \* We know how EW symmetry is broken.
- \* The burning question:  
**Is the EW scale natural or tuned?**

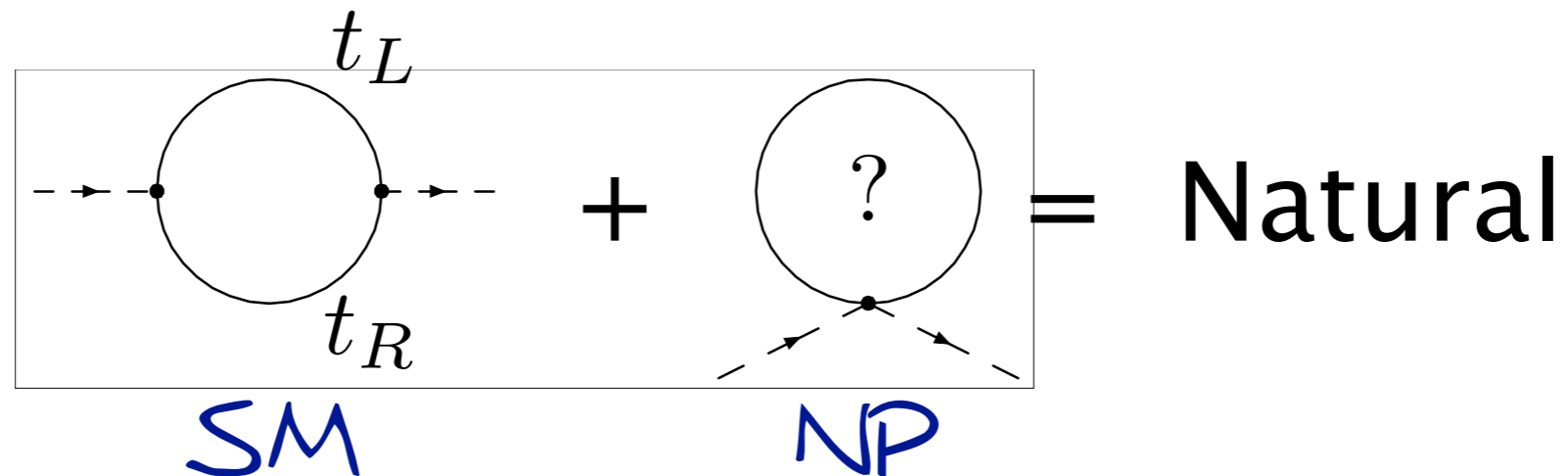
Our hope:

LHC will address this question.  
(Evidence of naturalness will be found).

This may still happen in run II.  
But what if it doesn't? Is the world tuned?

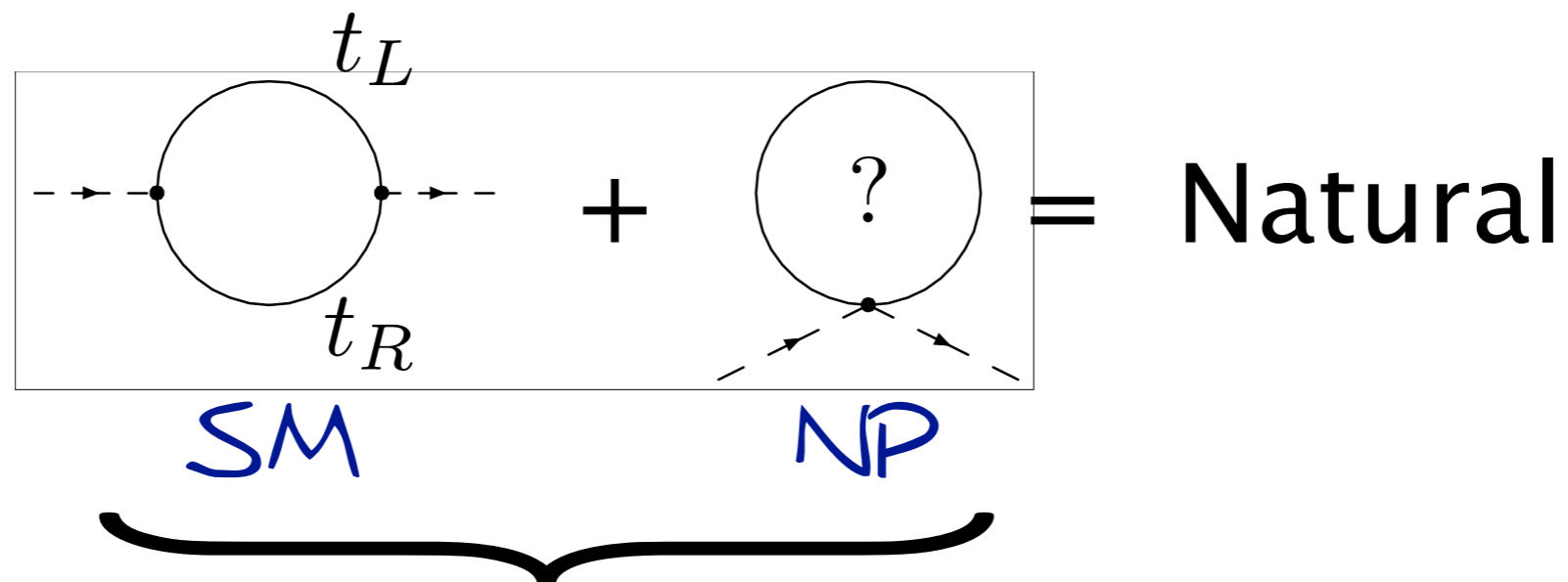
# Naturalness and LHC

- \* Why did we expect LHC to find the evidence for naturalness?



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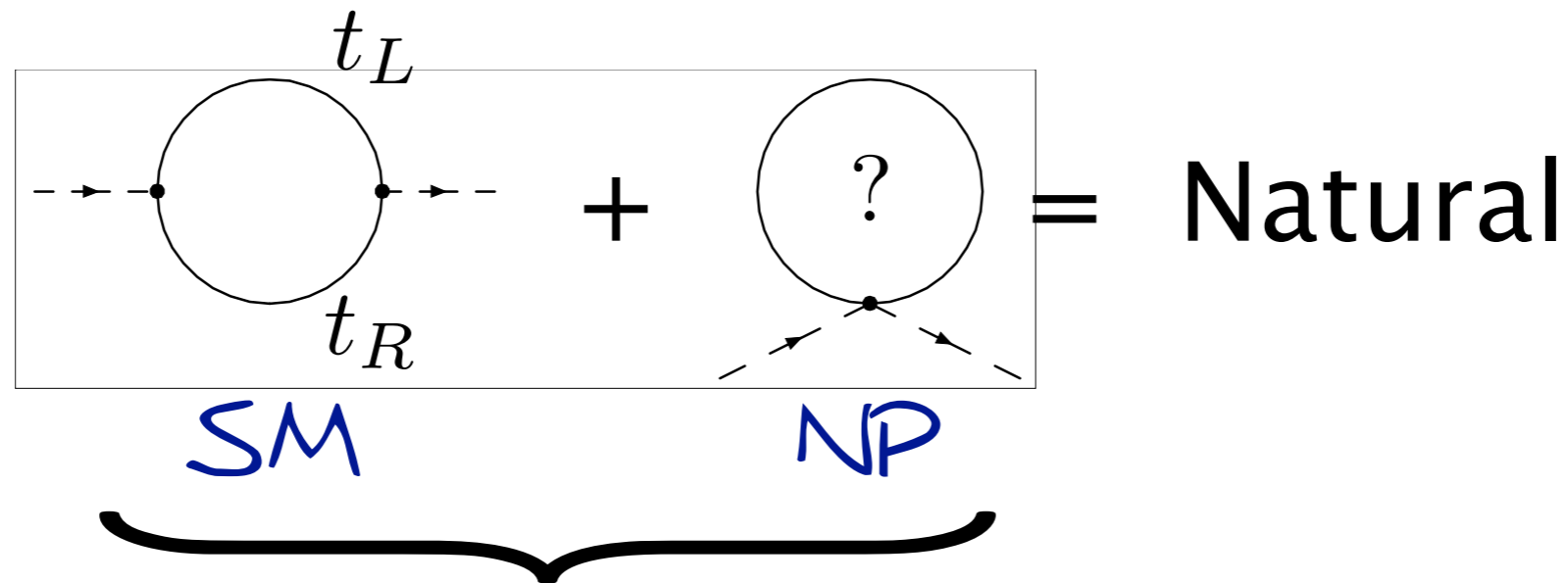


NP is related to the SM top by a symmetry.  
NP is around a TeV.

Colored NP at a TeV!  
Will be produced abundantly at LHC!

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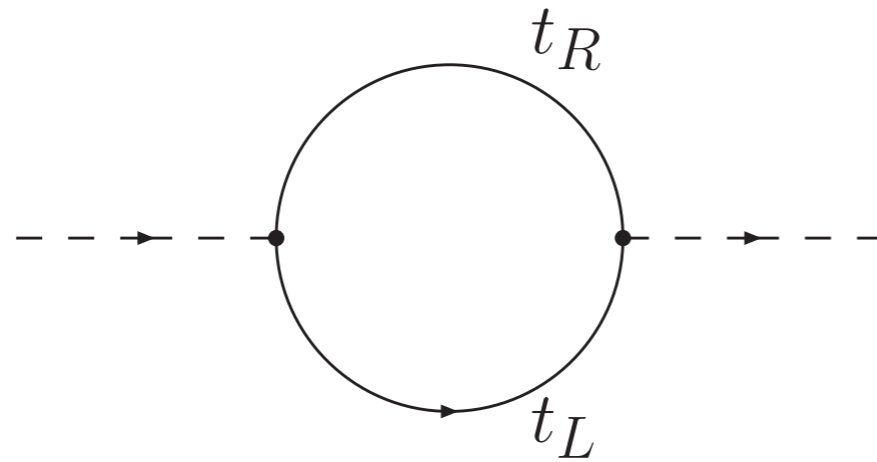


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**Not Necessarily True.**

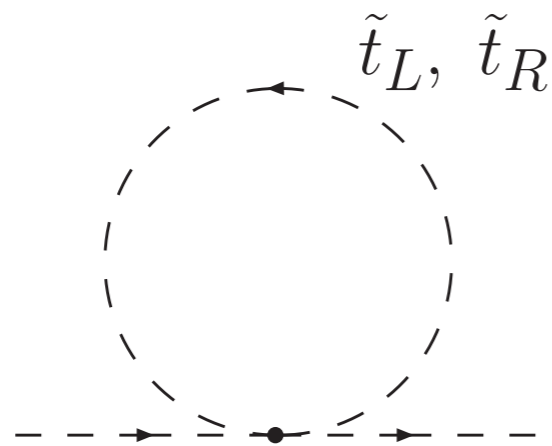
# Just a Factor of 3



Standard Model

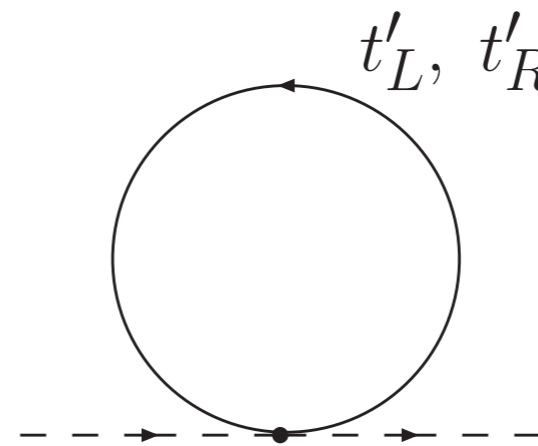
color factor:

$$\times 3$$



Supersymmetry

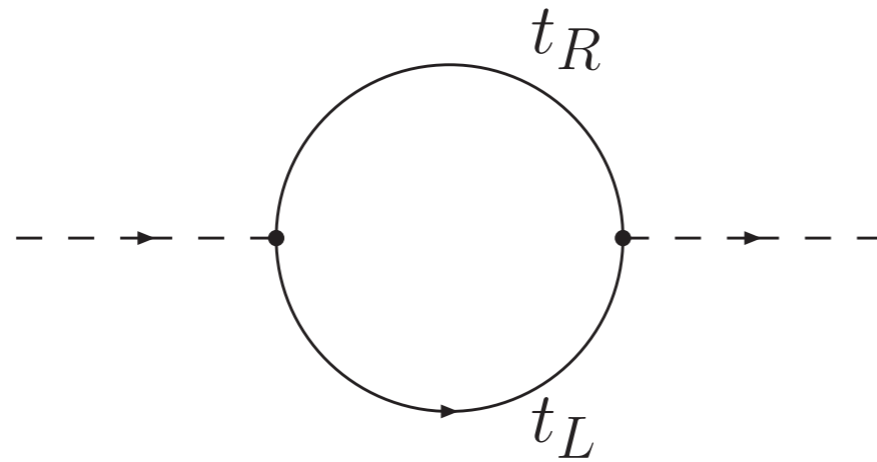
or



Little Higgs

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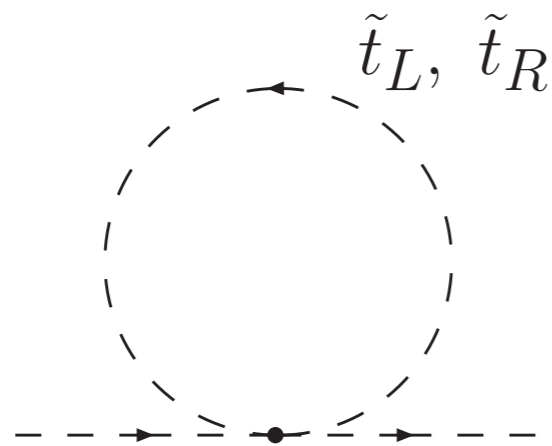
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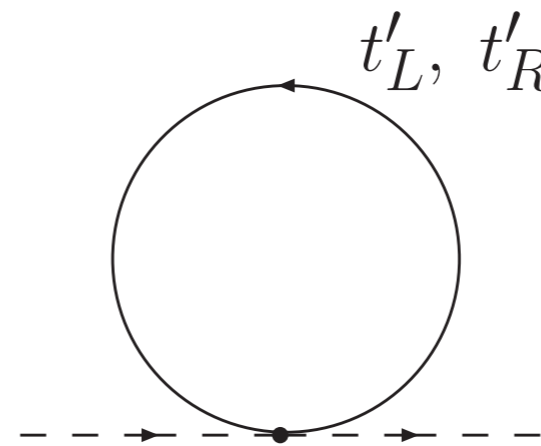
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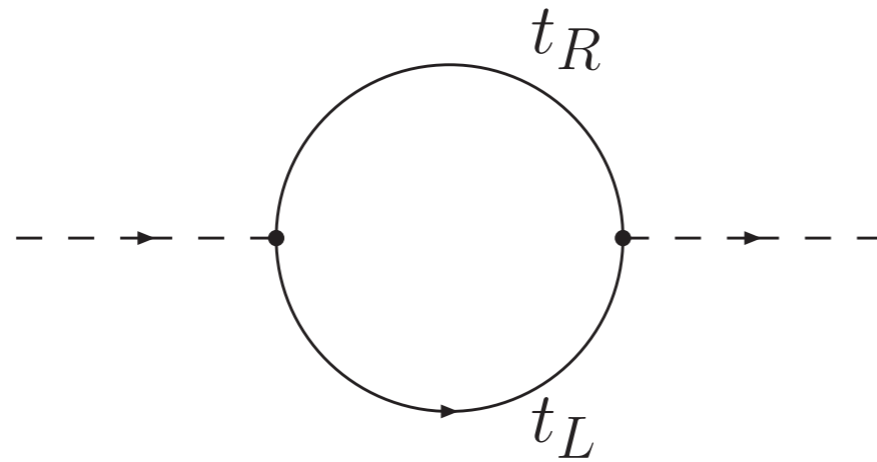
or



Little Higgs

$$\times \cancel{3} \quad 3'$$

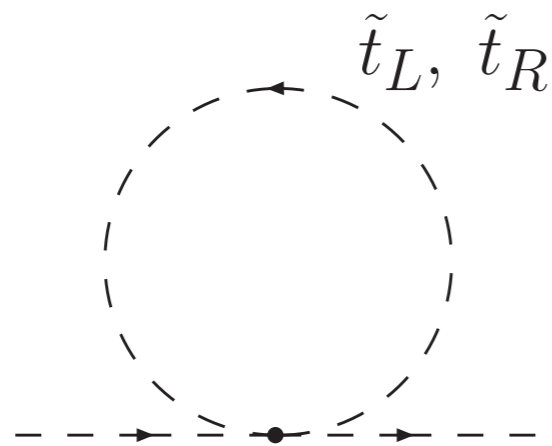
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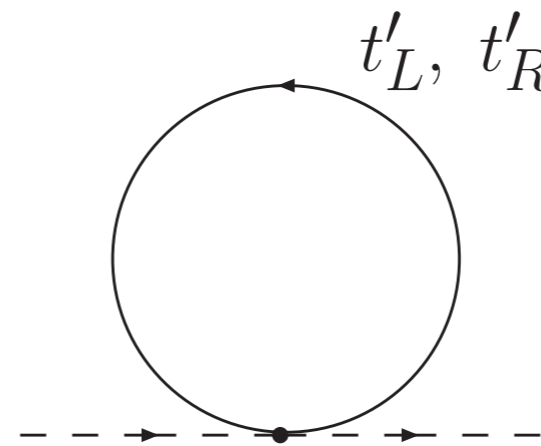
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Supersymmetry

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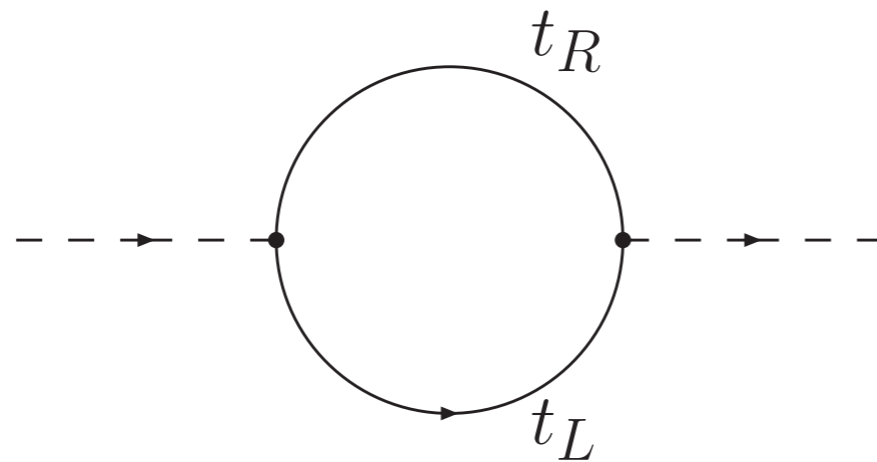
$$\times \cancel{3}$$

$$3'$$

*symmetry does not commute with color.*



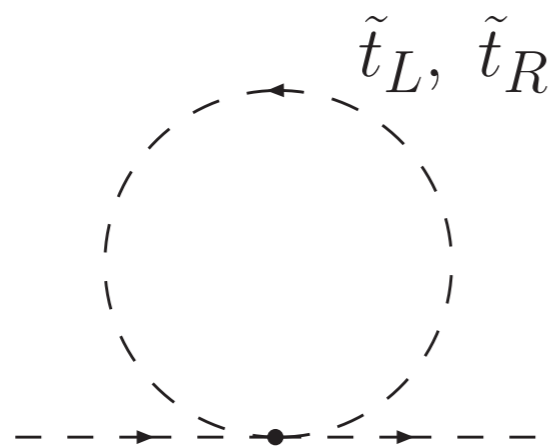
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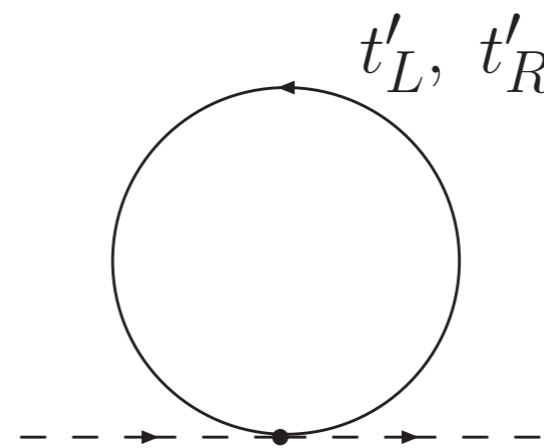
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Supersymmetry

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$$\times 3$$

~~3~~  
3'  
symmetry does not commute with color.

## Folded SUSY

Burdman, Chacko, Goh, RH (06')

## Twin Higgs

Chacko, Goh, RH (05')

# Outline

- \* Review: Twin Higgs

- o The Mechanism.
- o Varieties of twin models.

- \* What can LHC discover?

- \* What can Future Colliders Discover?

- \* A no lose theorem?

To "prove" a theorem we should pick the difficult path at every step.


# The Mechanism.

The Higgs is a PNCB of an approximate  $SU(4)$ .

# A Toy Example

- \* A global  $SU(4)$  symmetry w/ one fundamental:

$$V(H) = -m^2 |H|^2 + \lambda |H|^4$$


$$\langle |H|^2 \rangle = \frac{M^2}{2\lambda} \equiv f^2$$


$$SU(4) \longrightarrow SU(3)$$

7 Goldstones

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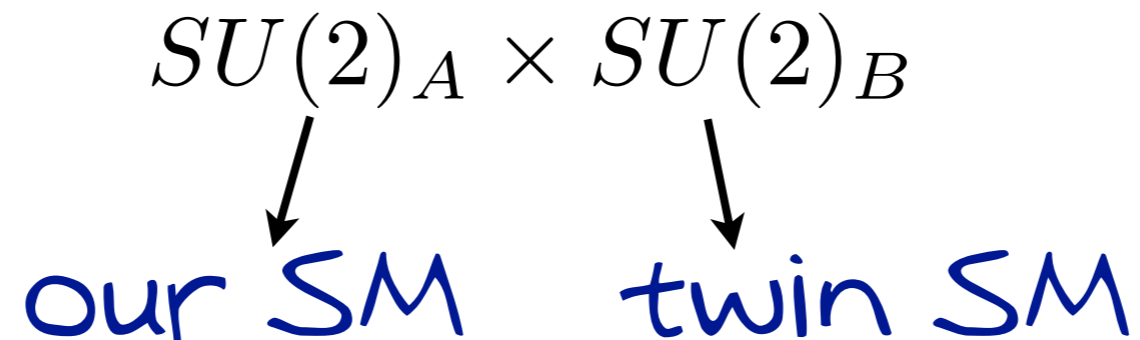
taxonomy:  
"weakly coupled":  $\lambda \sim O(1\text{-few})$   
(linear sigma model)  
"strongly coupled":  $\lambda \sim 16\pi^2$   
(non-linear sigma model)

$$SU(4) \longrightarrow SU(3)$$

7 Goldstones

$$SU(2)_A \times SU(2)_B$$

- \* Gauge a subgroup (a.k.a  $Z_2$  orbifold of  $SU(4)$ ):



- \* In some basis,  $H$  transforms as

$$H = \begin{pmatrix} H_A \\ H_B \end{pmatrix} \begin{array}{l} 6 \text{ eaten.} \\ 1 \text{ Goldstone left.} \end{array}$$

- \* Gauging  $SU(2)_A \times SU(2)_B$  breaks global  $SU(4)$ .

# Radiative Corrections

\* At 1-loop:

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\* Impose a  $Z_2$  “twin” symmetry:

$$A \longleftrightarrow B$$

$$g_A = g_B$$



$$\Delta V = \frac{9g^2 \Lambda^2}{64\pi^2} \left( H_A^\dagger H_A + H_B^\dagger H_B \right) \quad \text{SU(4) invariant!}$$

**Does not give a Goldstone mass.**

# Twin Mechanism

~~(Global Symmetry)~~ + (Discrete Symmetry)



Quadratic terms are globally symmetric.  
No quadratic divergences.

- \* Quartic terms can violate global symmetry.  
Goldstone mass only log divergent.

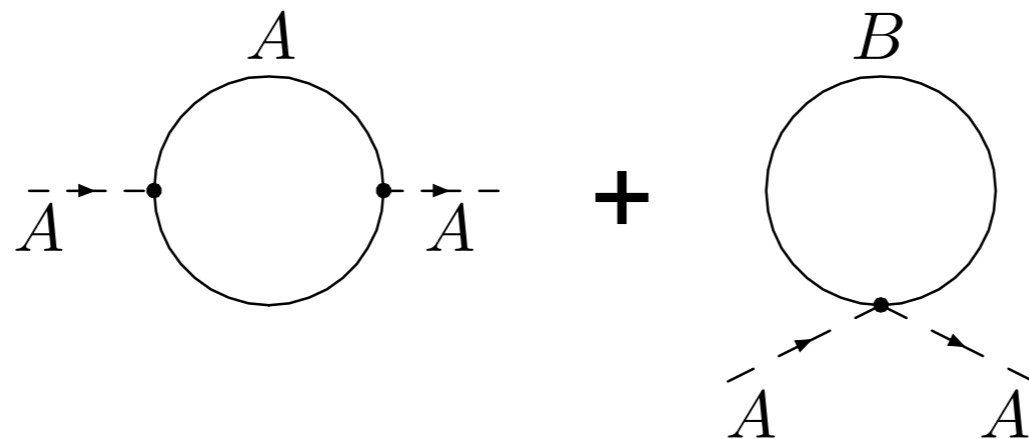
$$SM_A \times SM_B$$

\* Double all of the SM. Impose  $Z_2$ .  
 (a.k.a orbifold of  $SU(6) \times SU(4)$  by a  $Z_2$ ).

\* In particular  $\mathcal{L} \supset y_t H_A \bar{t}_A t_A + y_t H_B \bar{t}_B t_B$

$Z_2$ : quadratic divergence has the form

\*  $c\Lambda^2 (|H_A|^2 + |H_B|^2)$   
*SU(4) invariant.*



# SU(4) breaking

- \* Radiative corrections induce

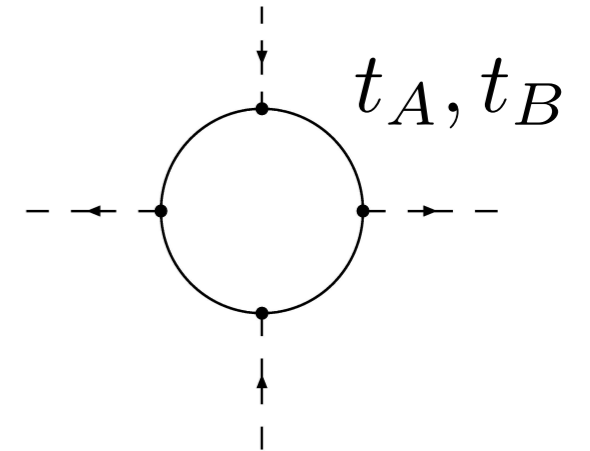
$$\Delta V = \kappa (|H_A|^4 + |H_B|^4)$$

with  $\kappa \sim \frac{y_t^4}{16\pi^2} \log \frac{\Lambda}{f}$

- \* Goldstone mass is  $m_h \sim \frac{y_t^2}{4\pi} f$ .

- \* Adding mixed “top partners” at 5-6 TeV keeps this quartic finite with a correct Higgs mass.

$$\begin{aligned} Q_L &= (\mathbf{6}, \bar{\mathbf{4}}) \\ &= (\mathbf{3}, \mathbf{2}; \mathbf{1}, \mathbf{1}) + (\mathbf{1}, \mathbf{1}; \mathbf{3}, \mathbf{2}) + (\mathbf{3}, \mathbf{1}; \mathbf{1}, \mathbf{2}) + (\mathbf{1}, \mathbf{2}; \mathbf{3}, \mathbf{1}) \end{aligned}$$



# Soft Breaking

- \* The potential as is gives  $v_A = v_B \sim \frac{f}{\sqrt{2}}$
- \* But then  $\Lambda \sim 4\pi f$  is too low.

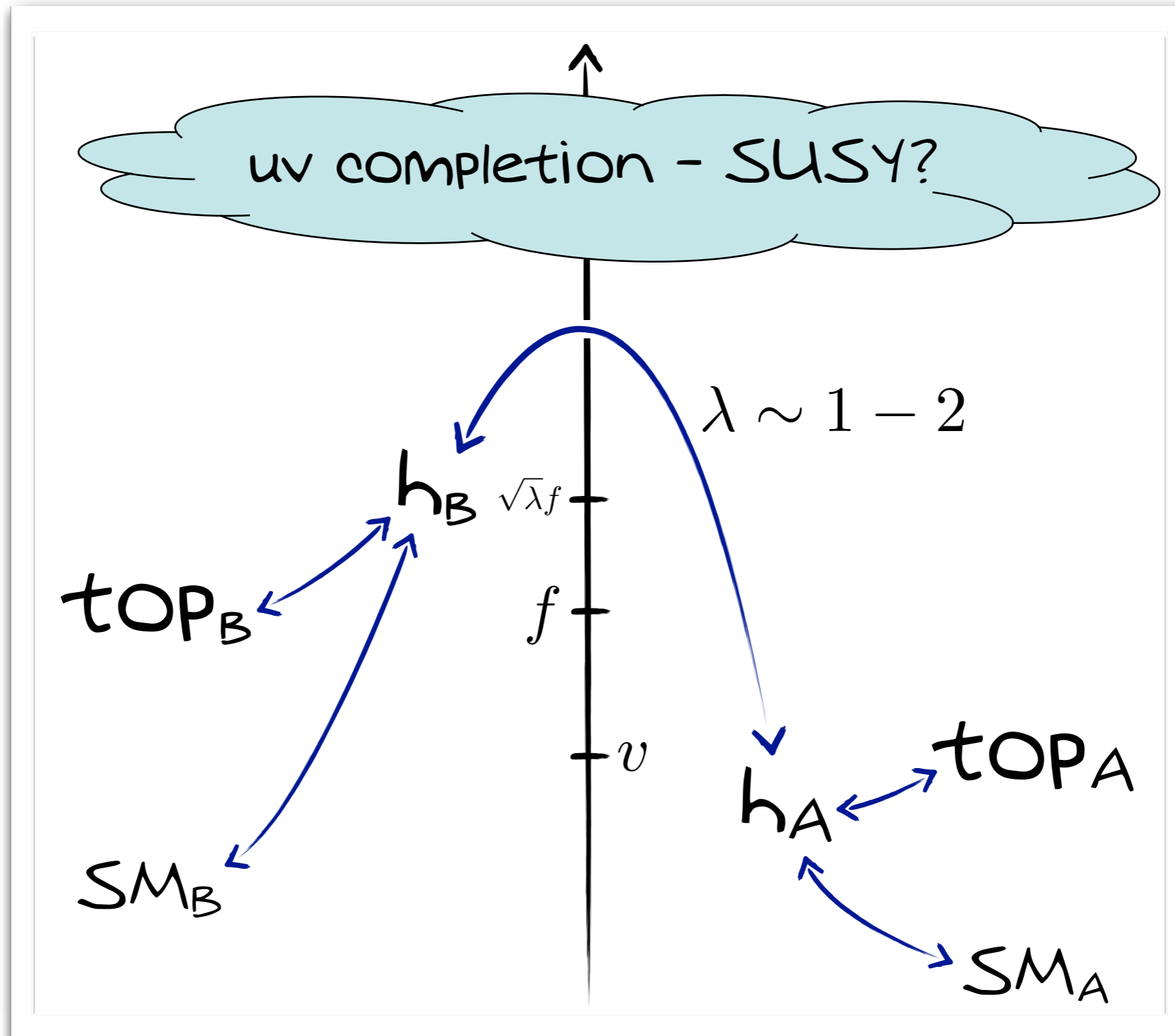
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- \* Add  $V_{soft} = \mu^2 |H_A|^2$  to get  $v < f$ .
- \* A mild cancelation needed to get EW scale

$\Lambda_{(\text{TeV})}$	$f_{(\text{GeV})}$	$M_{(\text{TeV})}$	$M_B_{(\text{TeV})}$	$\mu_{(\text{GeV})}$	$m_h_{(\text{GeV})}$	Tuning
10	800	6	1	239	122	0.134
6	500	5.5	1	145	121	0.378
10	800	—	0	355	166	0.112
6	500	—	0	203	153	0.307

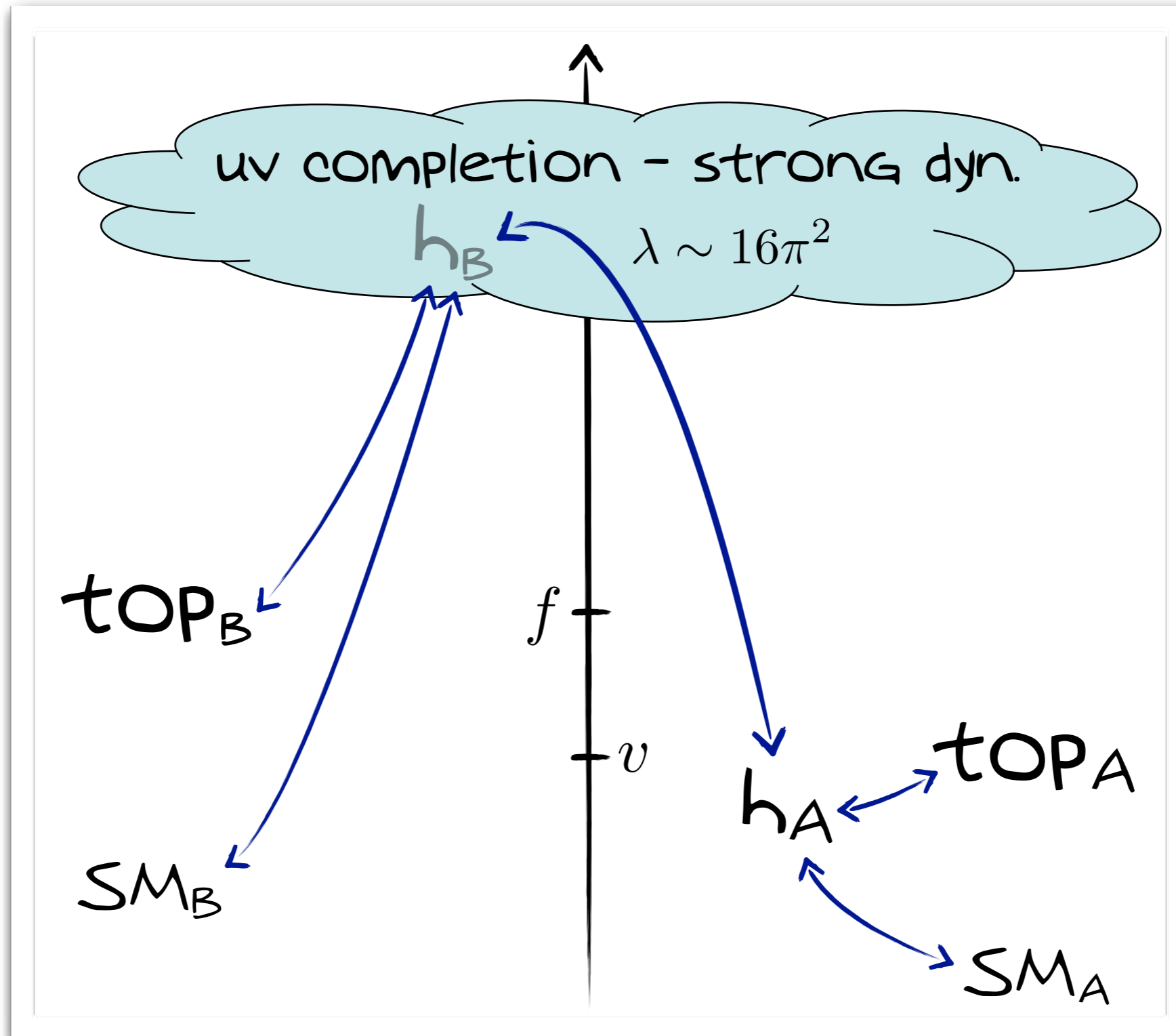
So let's summarize the moving parts....

# Sketch - Weak Coupling





# Sketch - Strong Coupling

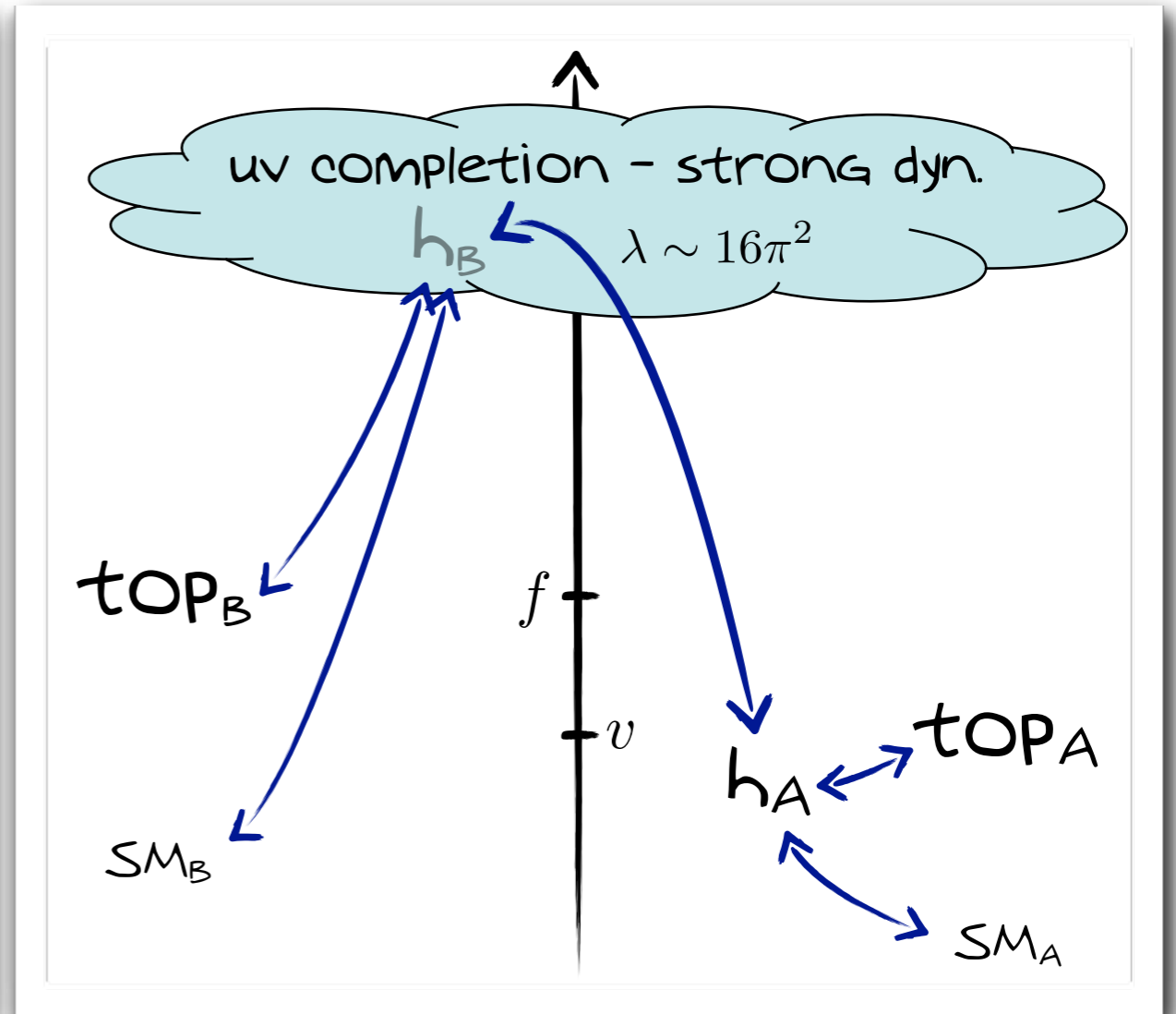
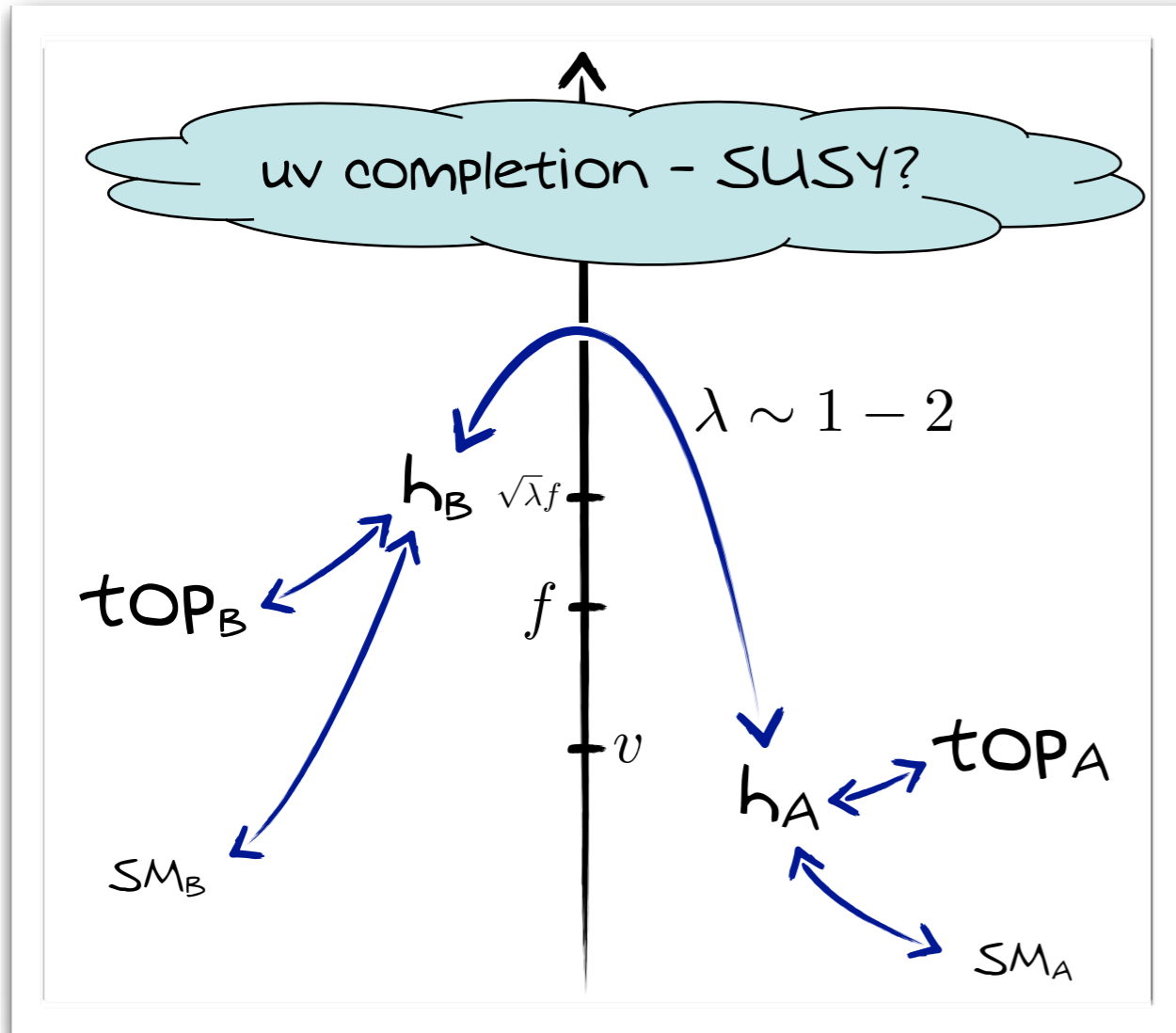


**With these moving parts  
we have a natural model.  
Yet *all* new states below a few TeV  
are complete SM singlets.**

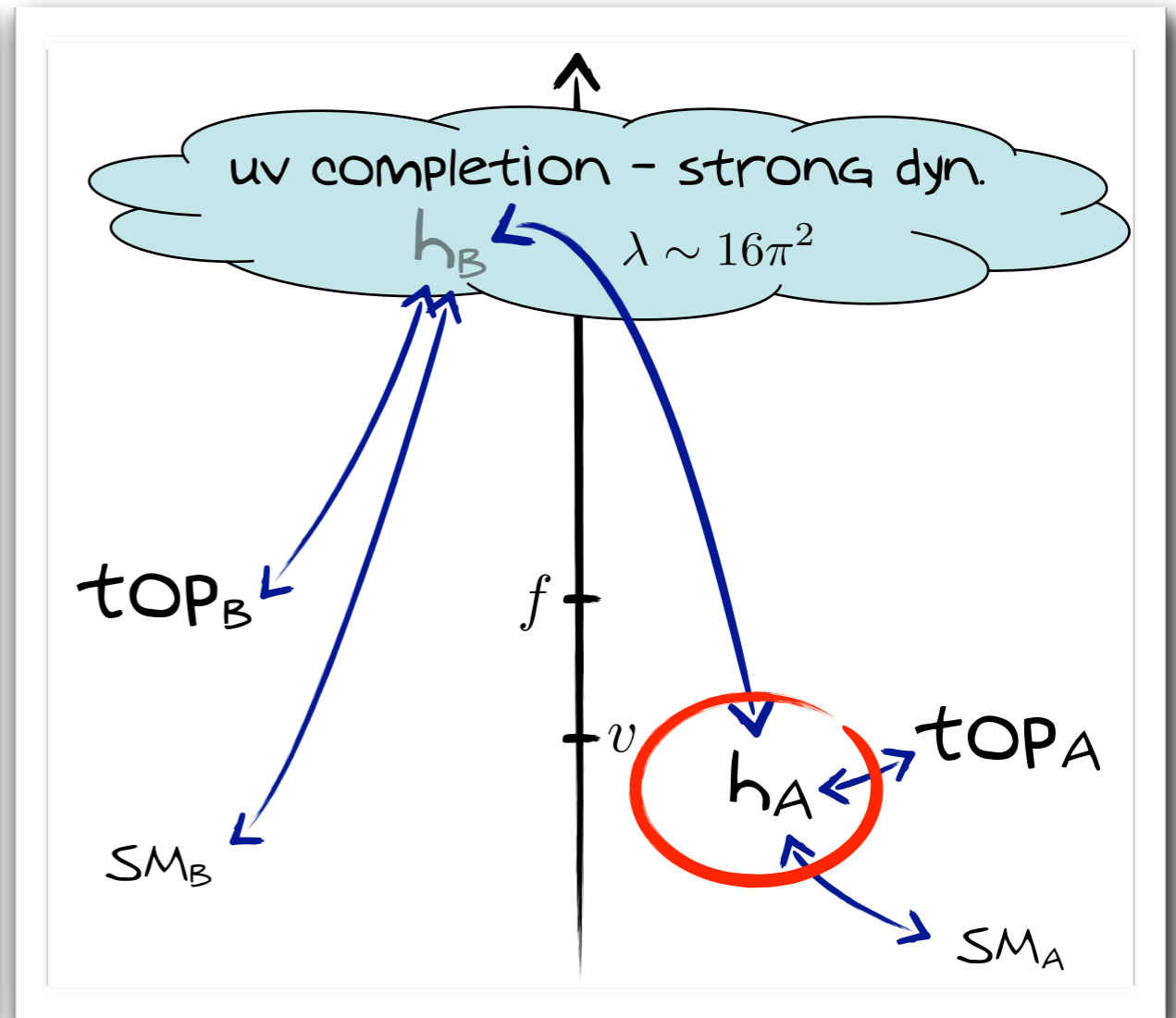
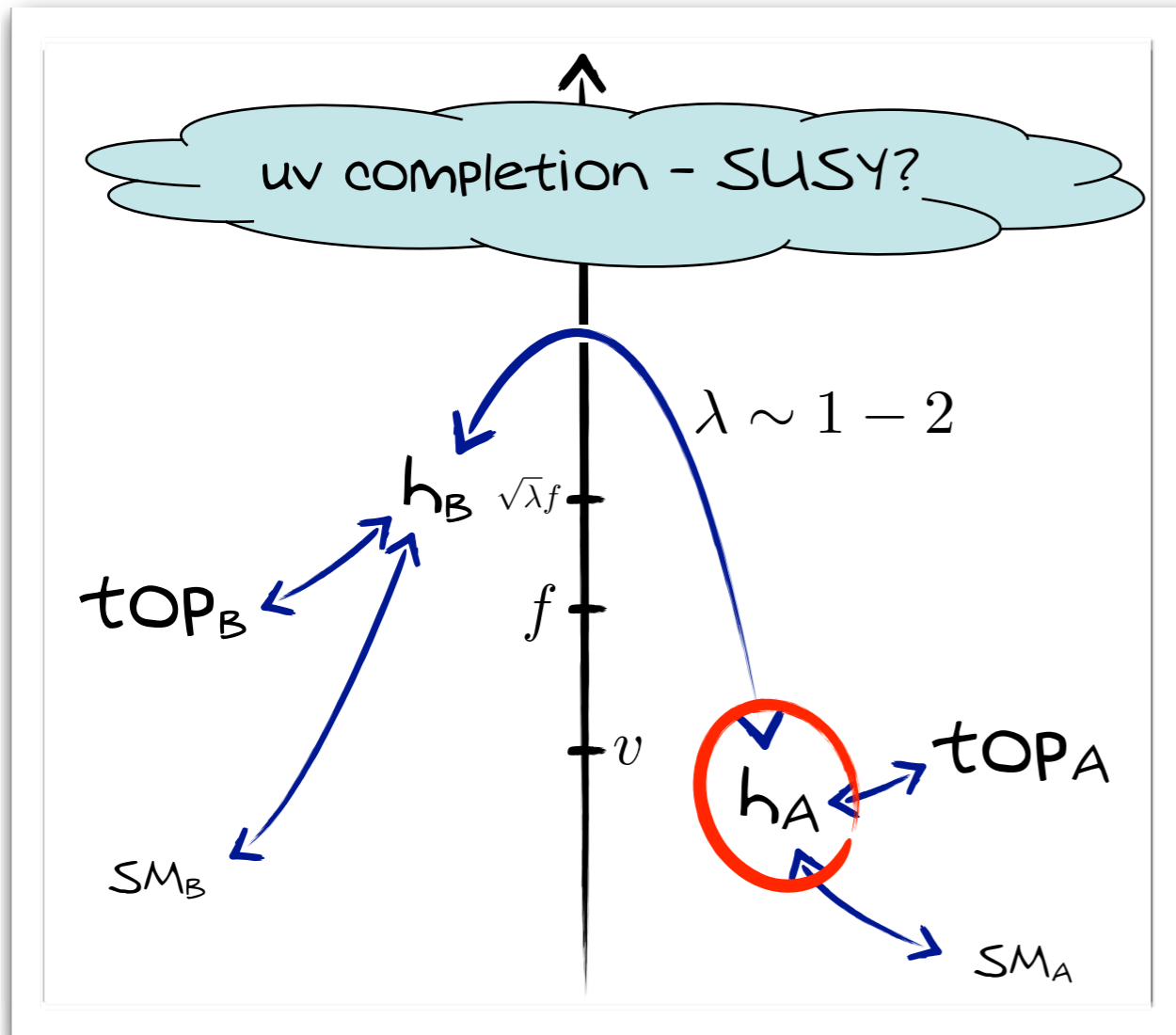
**LHC signatures will be subtle at best.**

Still, what are its signatures?  
LHC? Higgs Factory?  
100 TeV?

# Where can we look for signs of NP?



Where can we look for signs of NP?



## Precision Higgs tests:

Modified Higgs couplings and invisible decays.

# Precision Higgs

- \* All SM Higgs couplings are universally suppressed:

$$\Gamma_{h \rightarrow X} = \cos^2 \left( \frac{v}{f} \right) \Gamma_{h \rightarrow X}^{\text{SM}} \approx \left( 1 - \frac{v^2}{f^2} \right) \Gamma_{h \rightarrow X}^{\text{SM}}$$

- \* Ratios of Higgs rates are SM-like.

- \* A correlated invisible decay

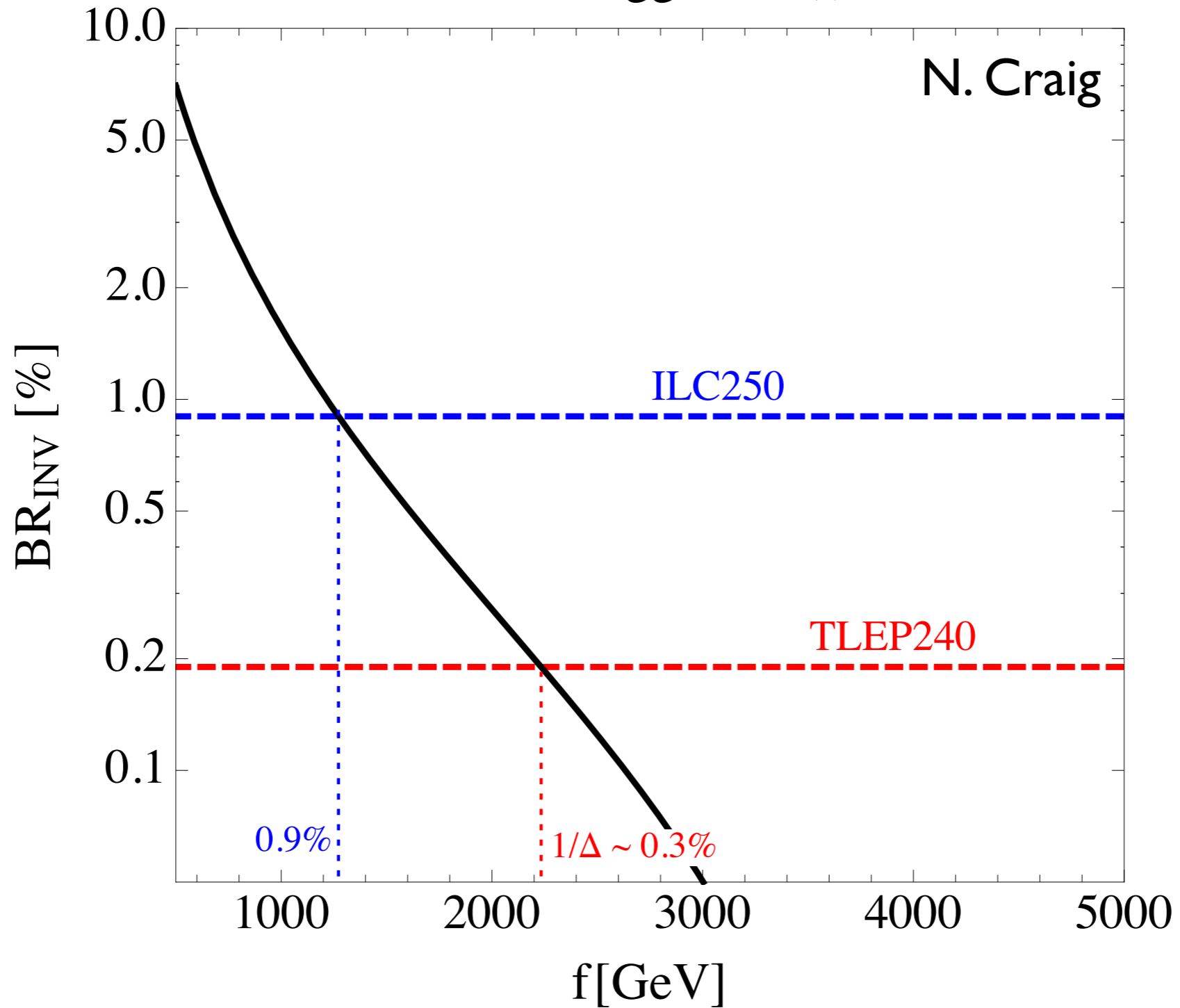
$$\Gamma_{h \rightarrow \text{inv}} = \sin^2 \left( \frac{v}{f} \right) \Gamma_{h \rightarrow b\bar{b}}^{\text{SM}} \approx \frac{v^2}{f^2} \Gamma_{h \rightarrow b\bar{b}}^{\text{SM}}$$

(assuming the  $b$  Yukawa respects the  $Z_2$  !!!).

- \* Both are set by the tuning in the model,  $\frac{2v^2}{f^2}$ .

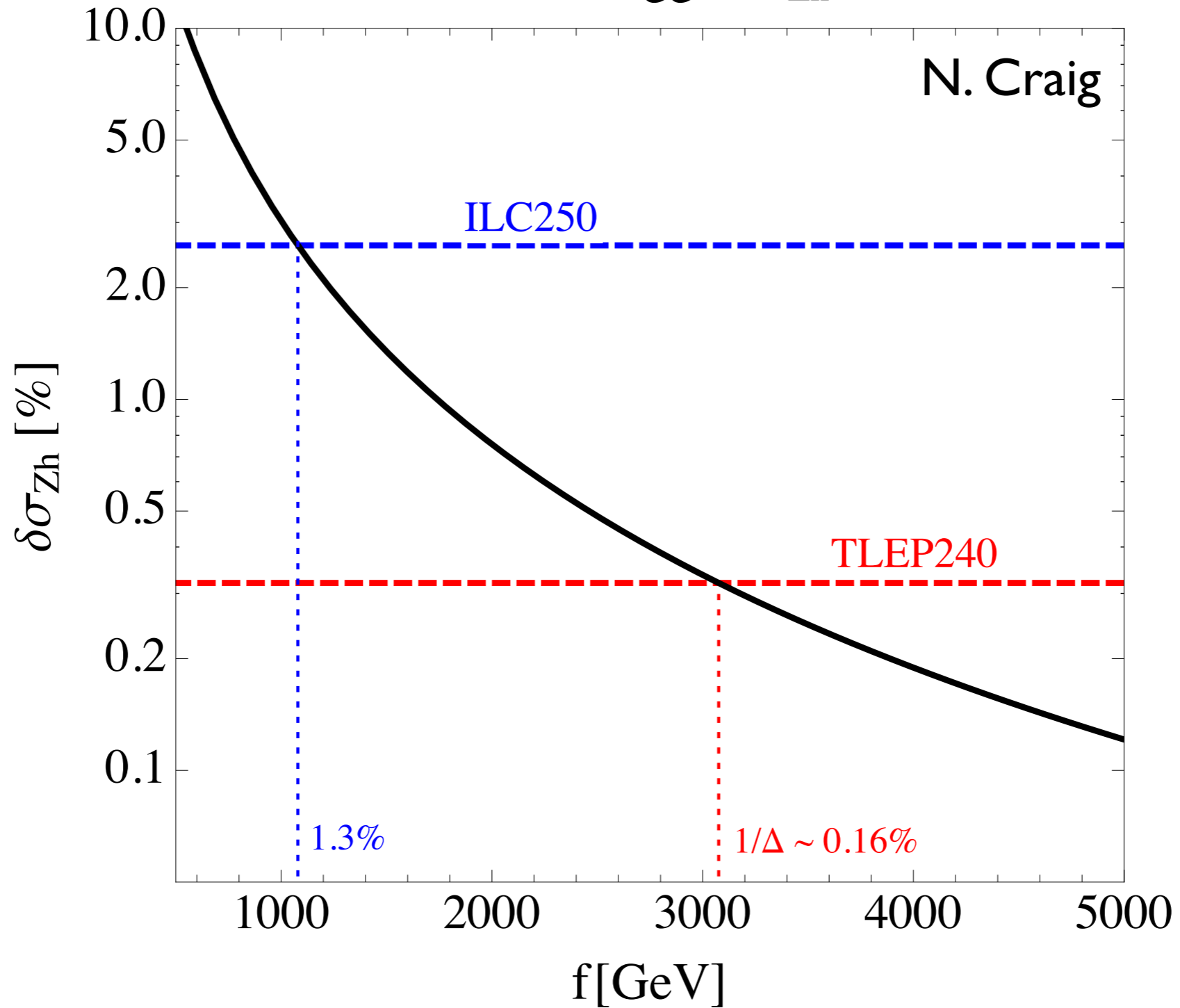
# Precision Higgs

Twin Higgs  $BR_{INV}$



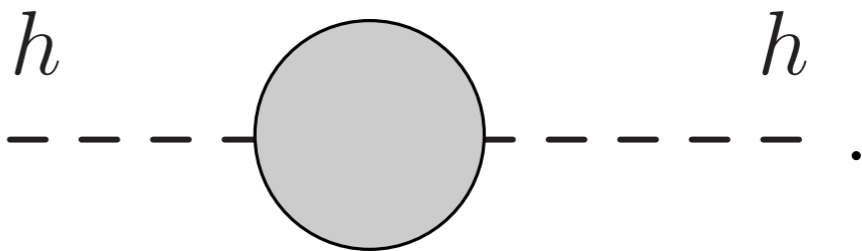
# Precision Higgs

Twin Higgs  $\delta\sigma_{Zh}$



# Non-Colored Partners

- \* But what if we came up with a model in which there is *no* tree level mixing?

$$\delta Z_h, \delta m_h^2 \sim \begin{array}{c} h \\ \text{---} \end{array} \text{---} \text{---} \text{---} \begin{array}{c} \text{---} \\ h \end{array} .$$


$$\mathcal{L} = \mathcal{L}_{SM} + \frac{1}{2} \delta Z_h (\partial_\mu h)^2 + \dots$$

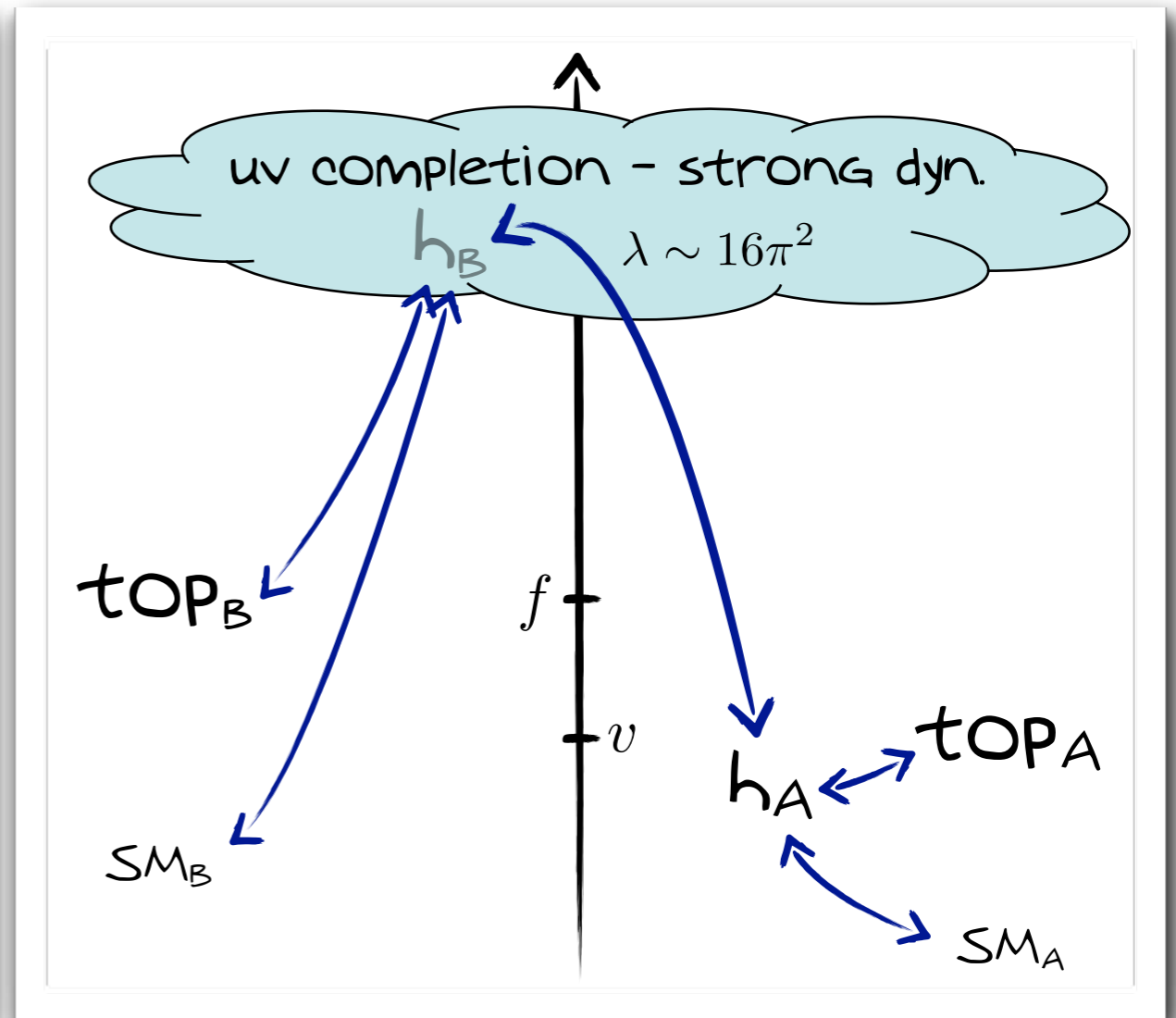
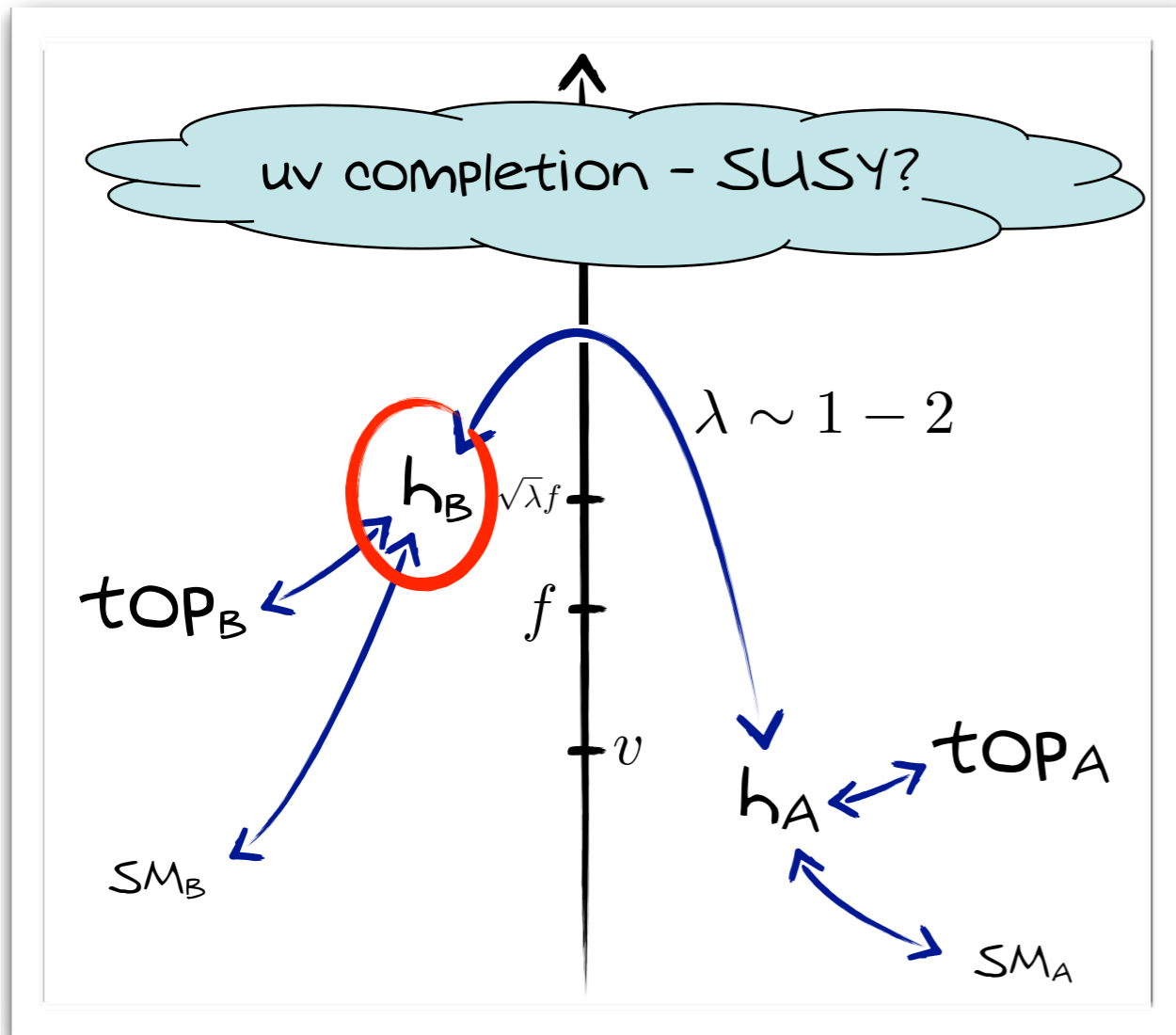
All Higgs couplings are modified by  $\delta Z_h$ .

of order 1% modifications.

torn from Craig, Englert, McCullough 1305.5252.



Where can we look for signs of NP?



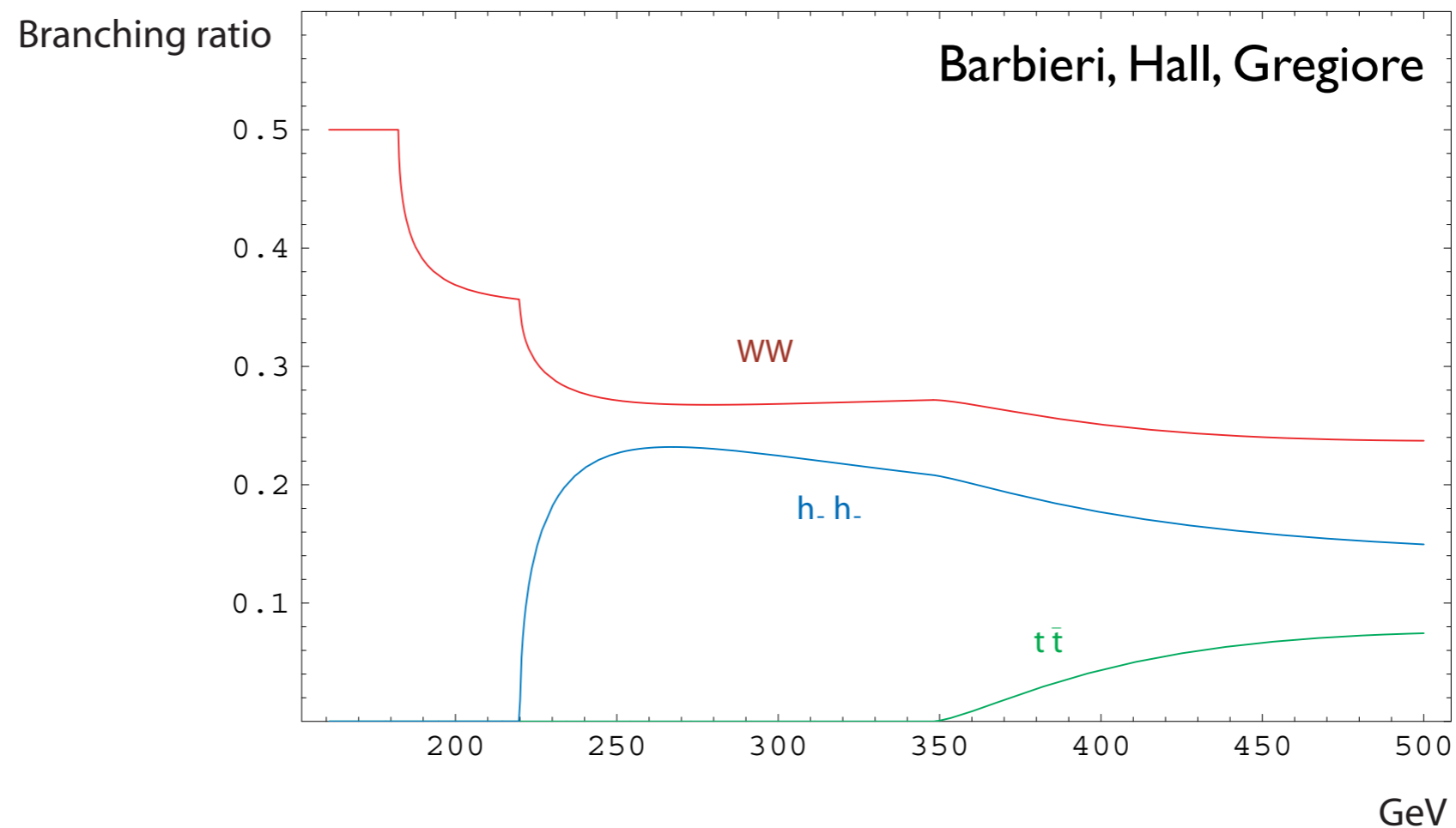
## UV Completion - linear model:

A heavy HIGGS.

Decays to invisibly or to  $W$ 's, HIGGSes, tops.

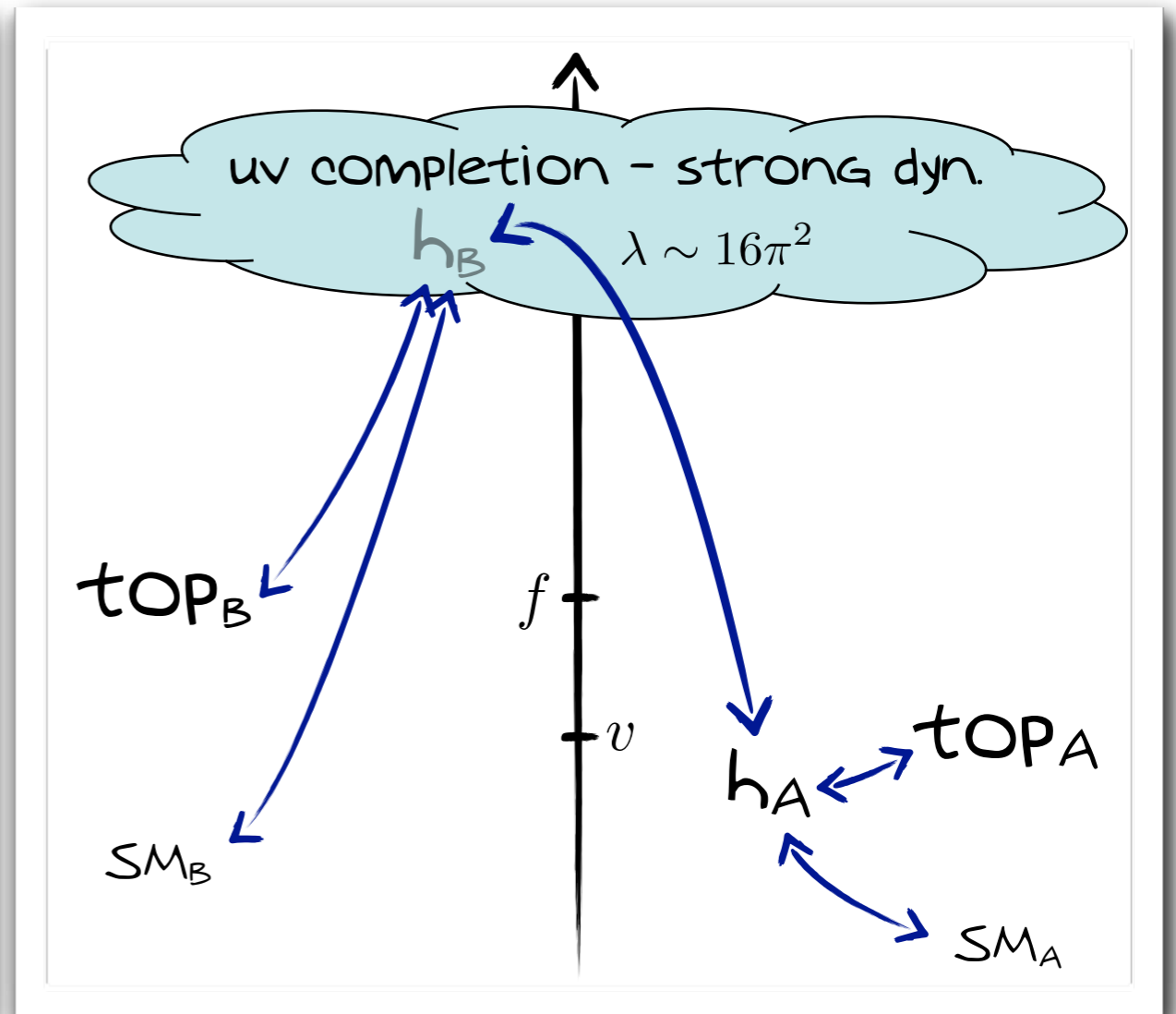
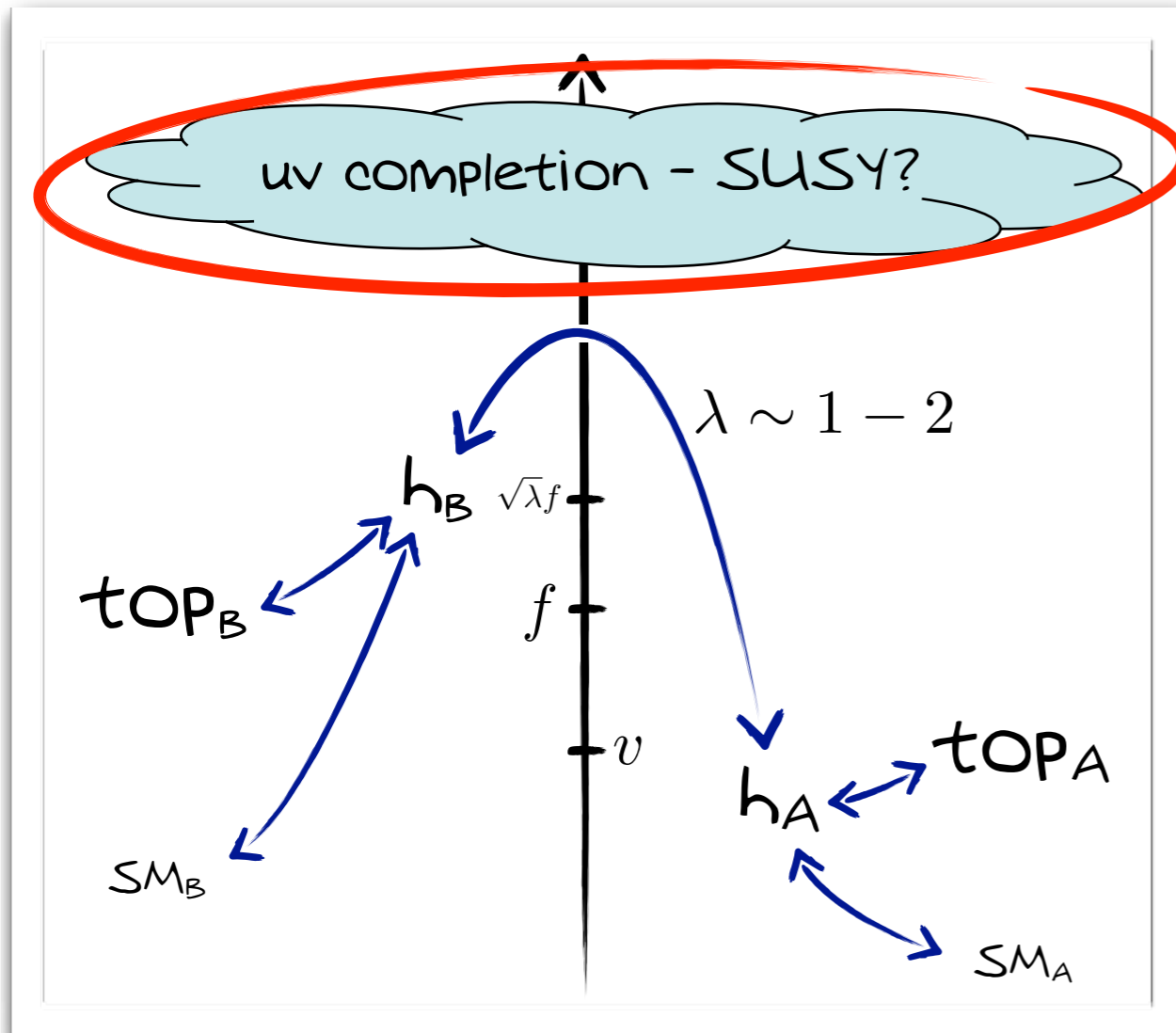
# Heavy Higgs

- \* The heavy Higgs decays  $\sim 50\%$  invisibly.
- \* The rest of the BR is dominated by  $WW$ ,  $hh$ , and  $t\bar{t}$ .



may be LHC accessible. (see Nate's talk)

Where can we look for signs of NP?

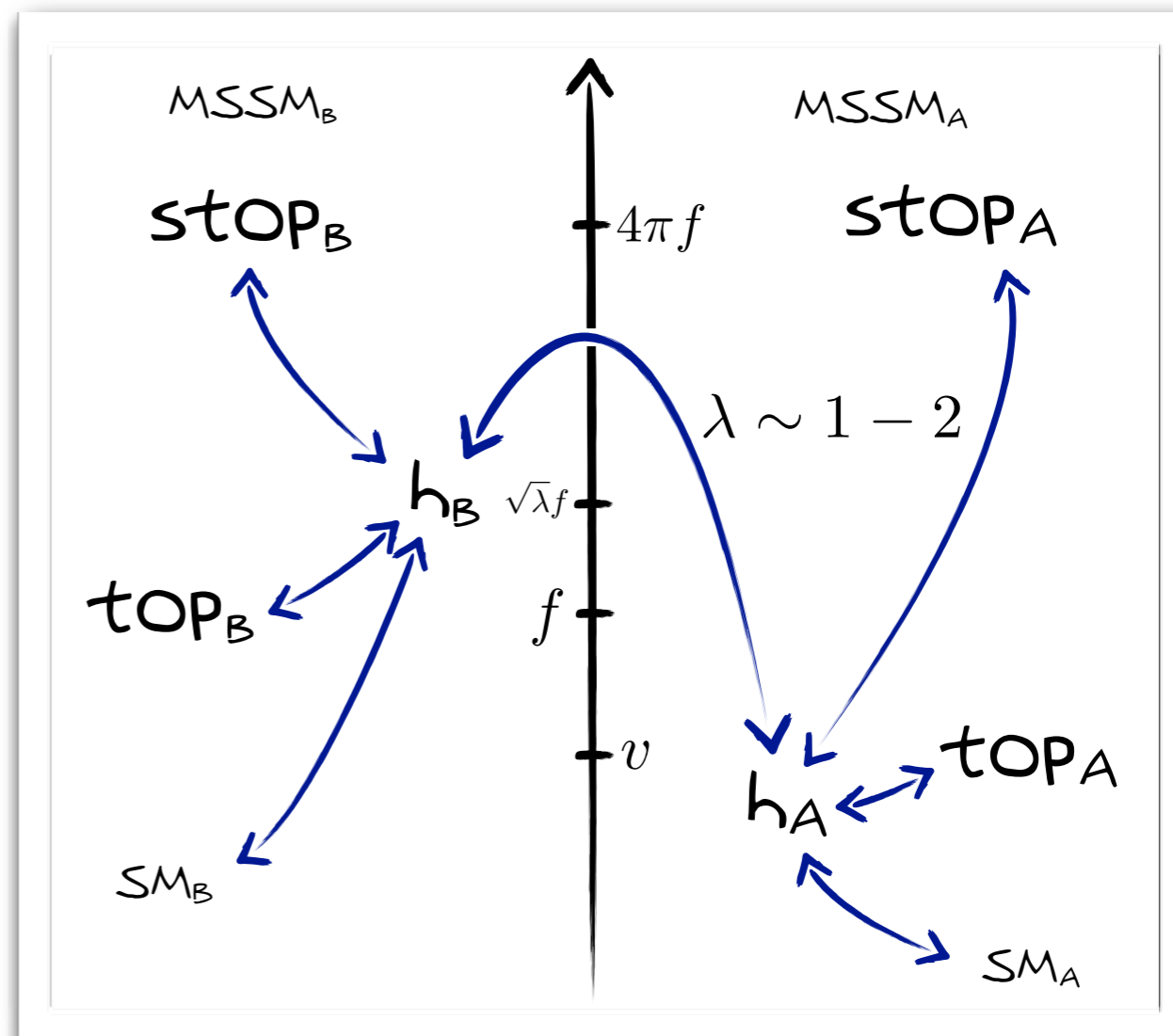


## UV Completion - linear model:

Re-do of the search for naturalness,  
But now at a few TeV.

# Twin SUSY

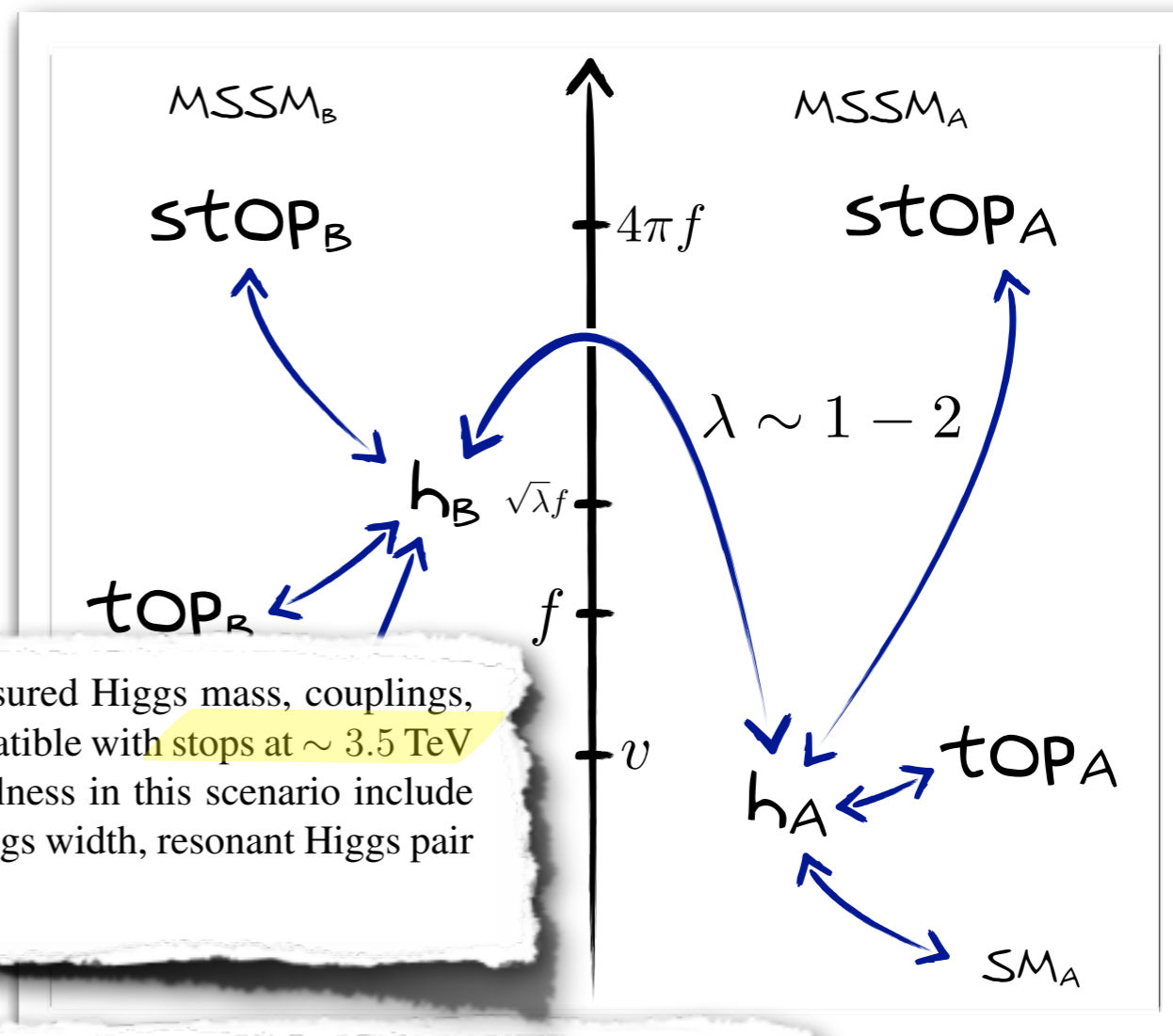
- \* Introduce SUSY at  $\sim$ few TeV.
- \* MSSM solves the big hierarchy. The scale  $f$  is natural.
- \* Doubling of MSSM Does the last bit, from  $f$  to  $v$ .



Stops accessible only at 100 TeV.

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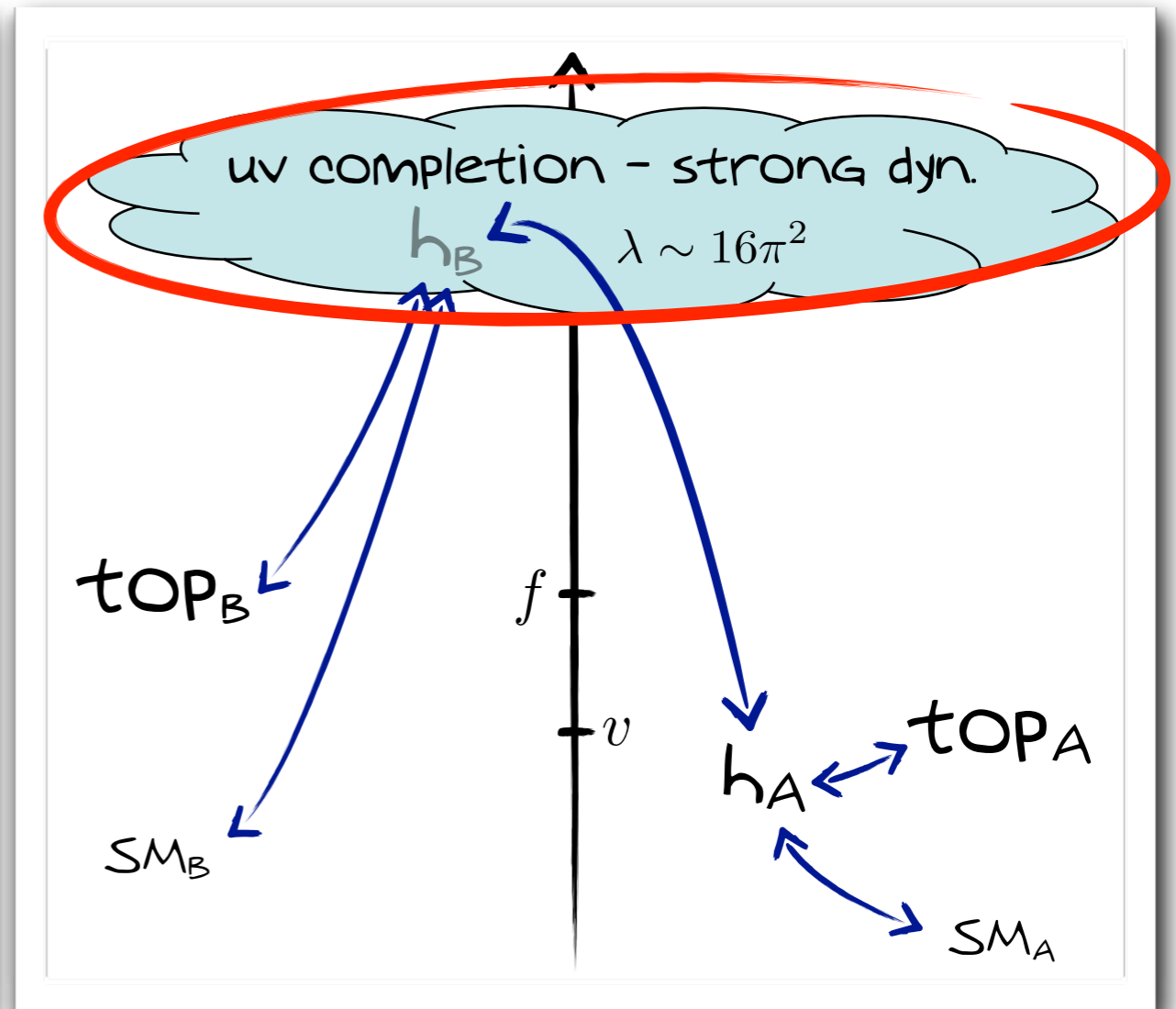
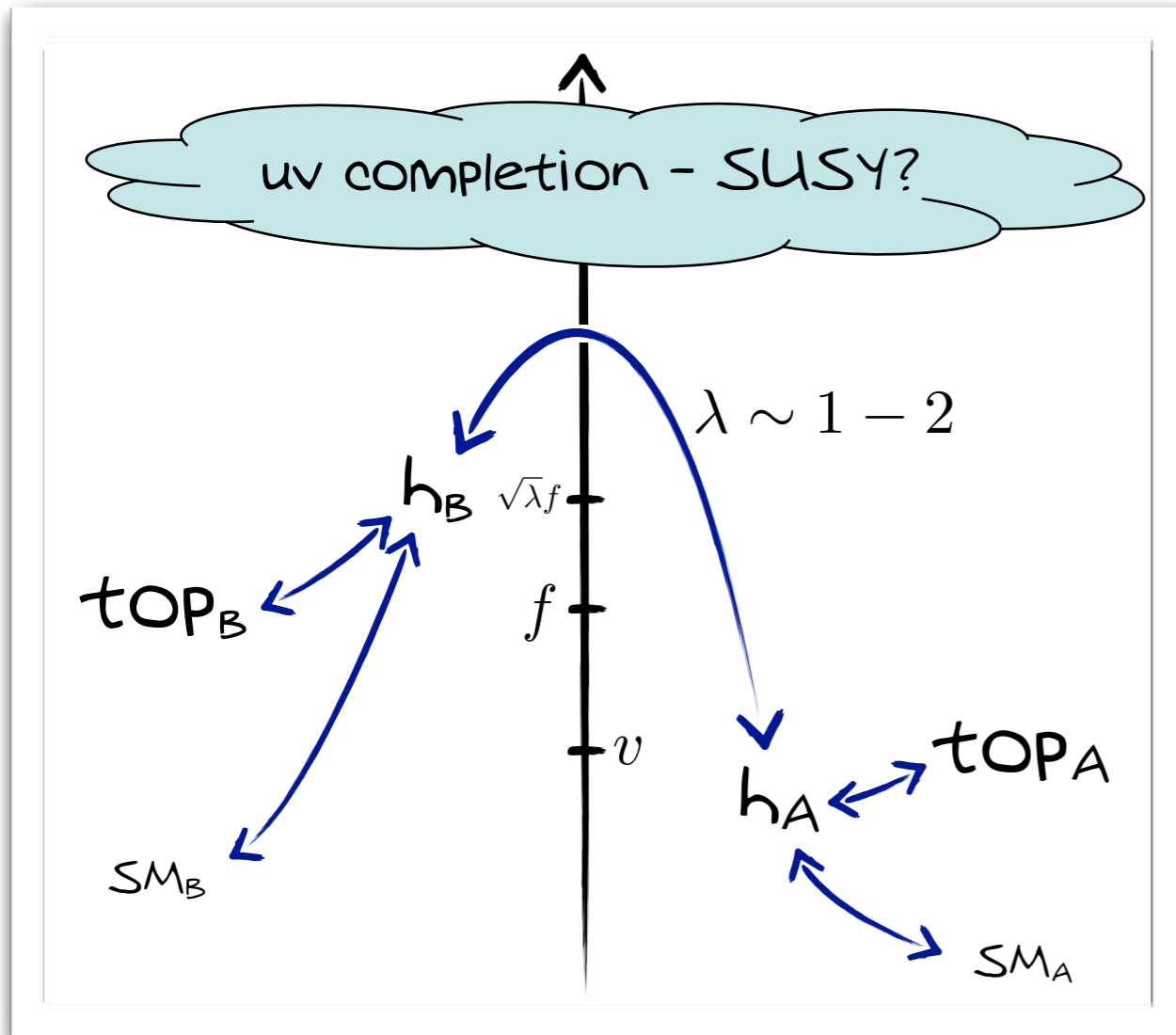


We find the measured Higgs mass, couplings, and percent-level naturalness of the weak scale are compatible with stops at  $\sim 3.5$  TeV and higgsinos at  $\sim 1$  TeV. The primary signs of naturalness in this scenario include modifications of Higgs couplings, a modest invisible Higgs width, resonant Higgs pair production, and an invisibly-decaying heavy Higgs.

Furthermore, if we discard the requirement of perturbative MSSM-like gauge coupling unification, a Higgs compositeness scale of  $\sim 50$  TeV allows 10%-level tuning with superpartners entirely out of reach of the LHC.

Stops accessible only at 100 TeV.

Where can we look for signs of NP?



**UV Completion - non-linear model:**  
Strong dynamics at 5-10 TeV. Resonances...

# New Strong Dynamics

- \* We imagine new QCD-like dynamics at 5-10 TeV.
- \* The Higgs doublet is a pion of this “techni-QCD”.

The constituents *must* be charged under EW.

The top *could* be a composite as well.



Colored states.  
Piece of cake at 100 TeV.

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$t_{\text{OP}_A}$                        $t_{\text{OP}_B}$                       *exotics*

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Can we make a generic statement?

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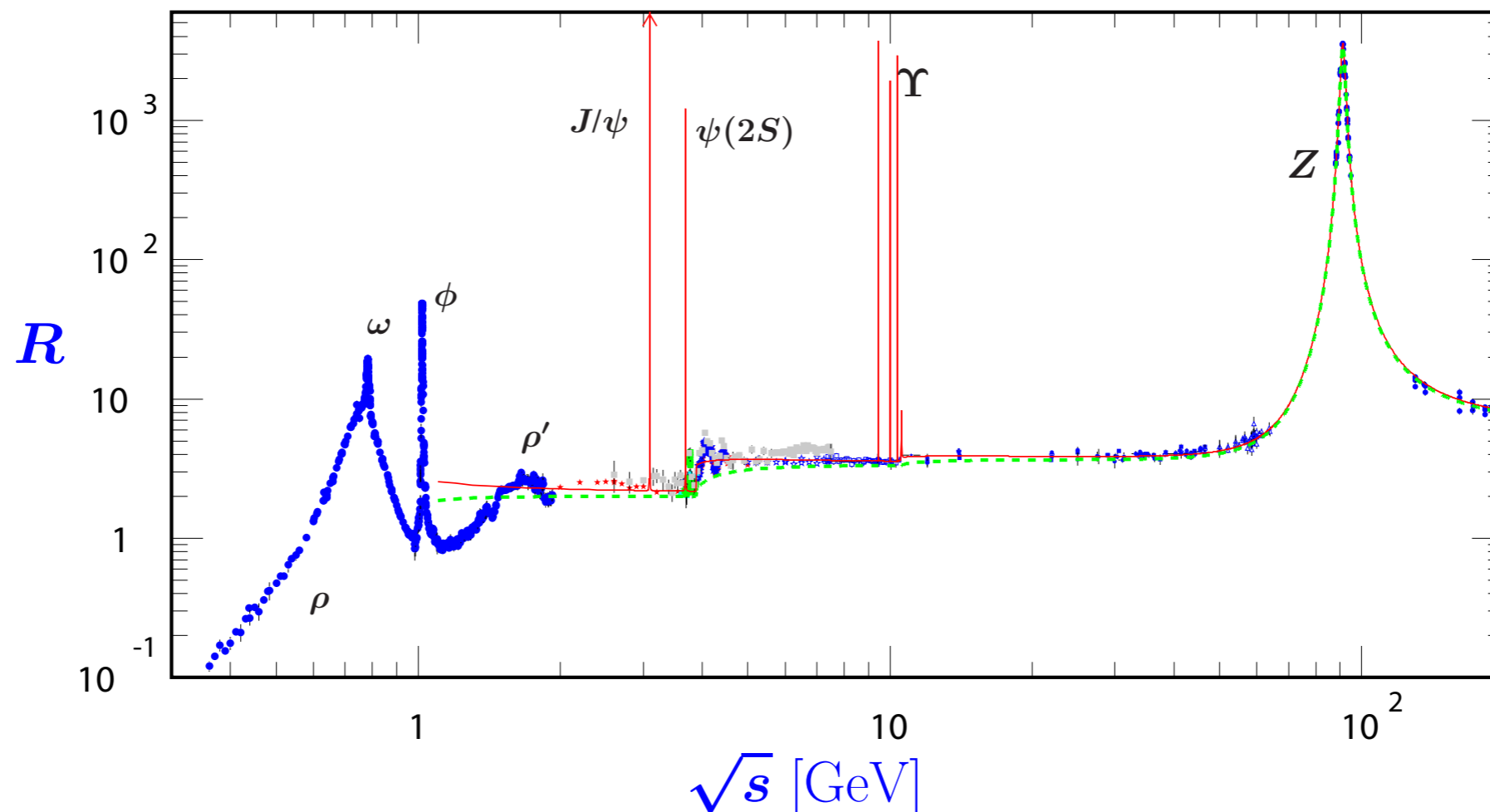


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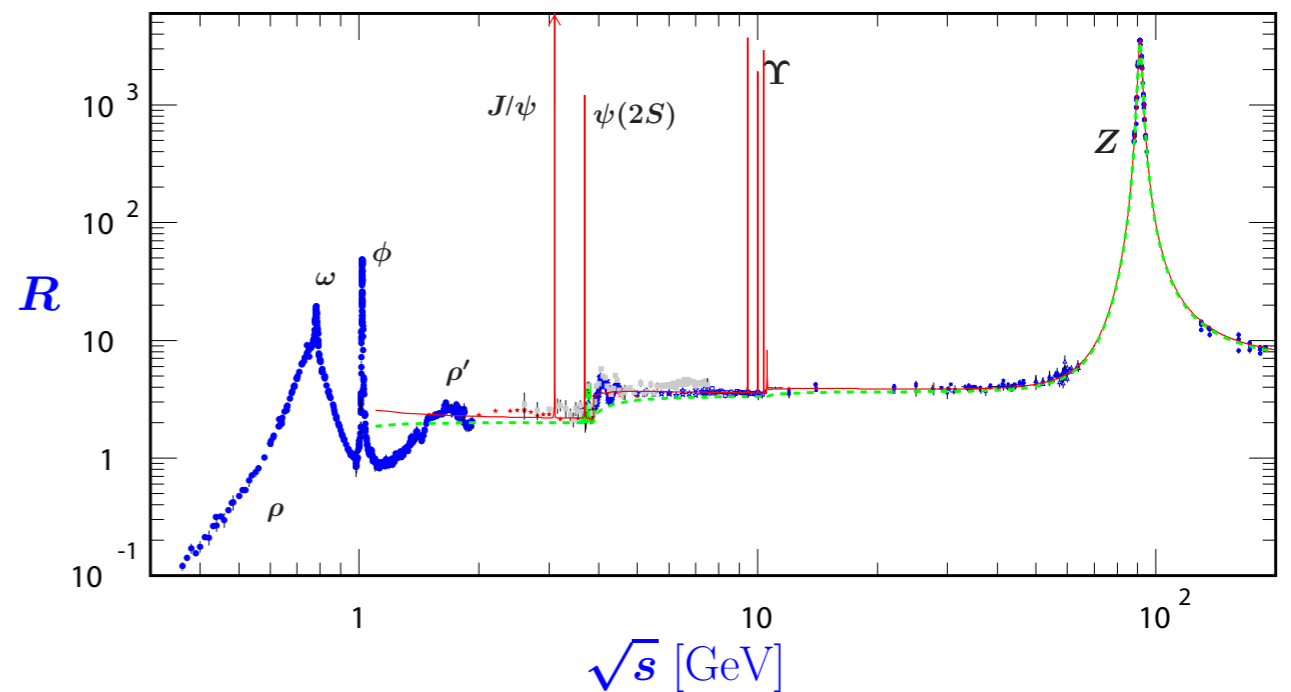
# Weak Resonances

- \* What is the 100 TeV sensitivity to weakly charged techi-sectors? (for any composite Higgs setup).
- \* We have one precedent for probing a strong sector with external weak probes:



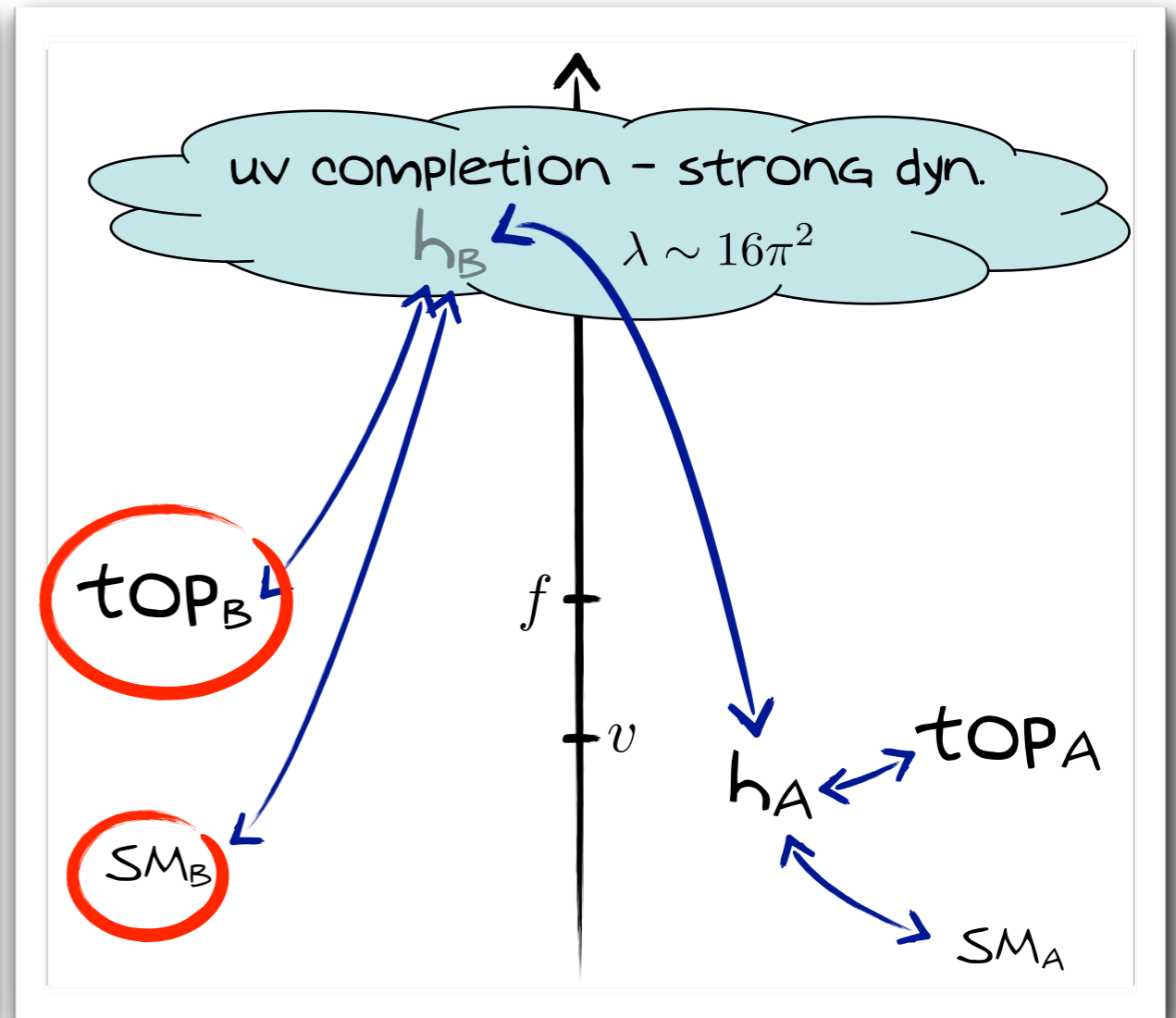
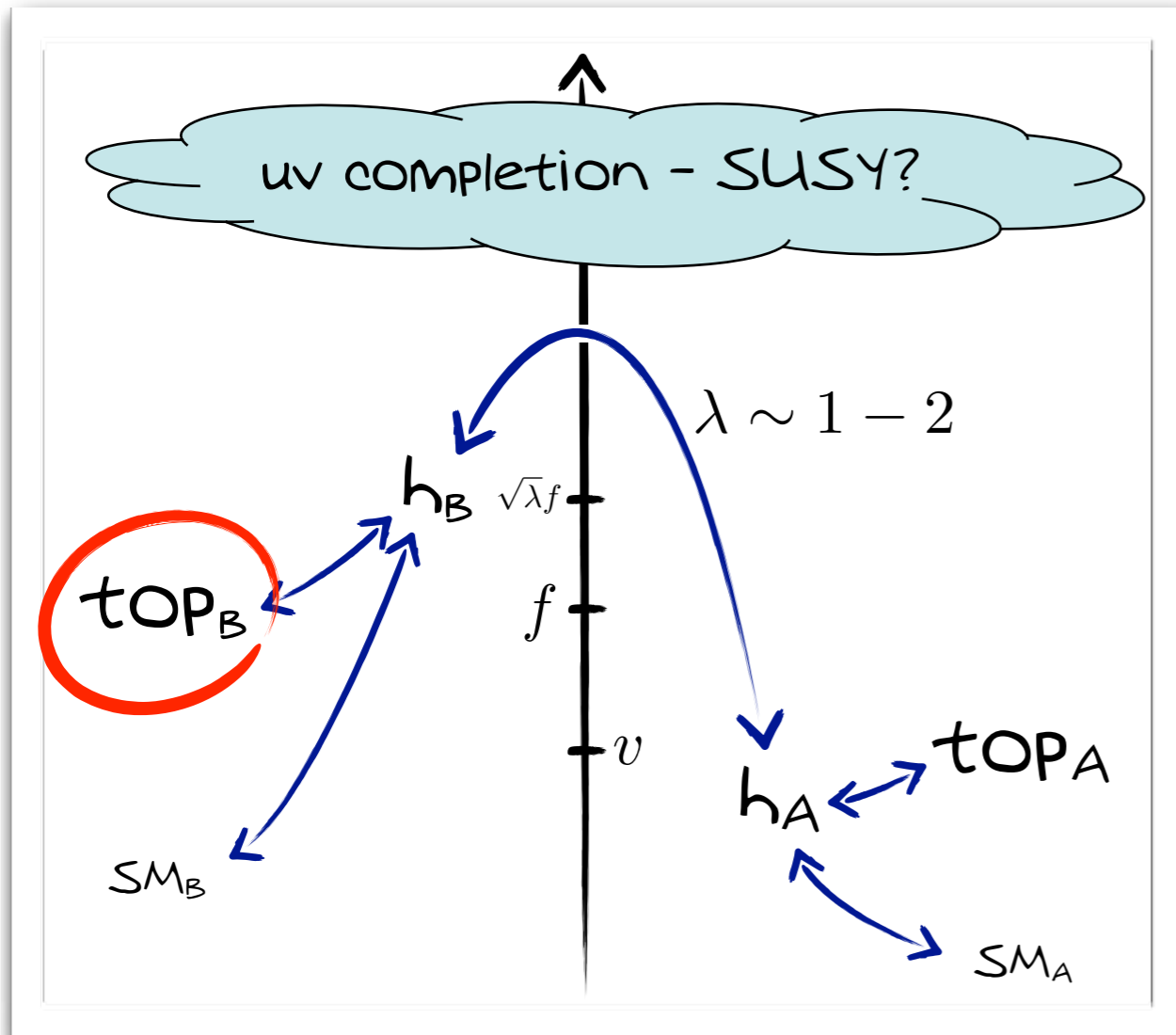
# Weak Resonances

- \* It seems reasonable that there will be
  - Wide  $W'$ -like states (rho-like) going to Higgses (pions).
  - Narrower states that must undergo weak decay.
  - Pair produced states.



Calls for a taxonomical study Based on quantum numbers and spins of resonances.

Where can we look for signs of NP?

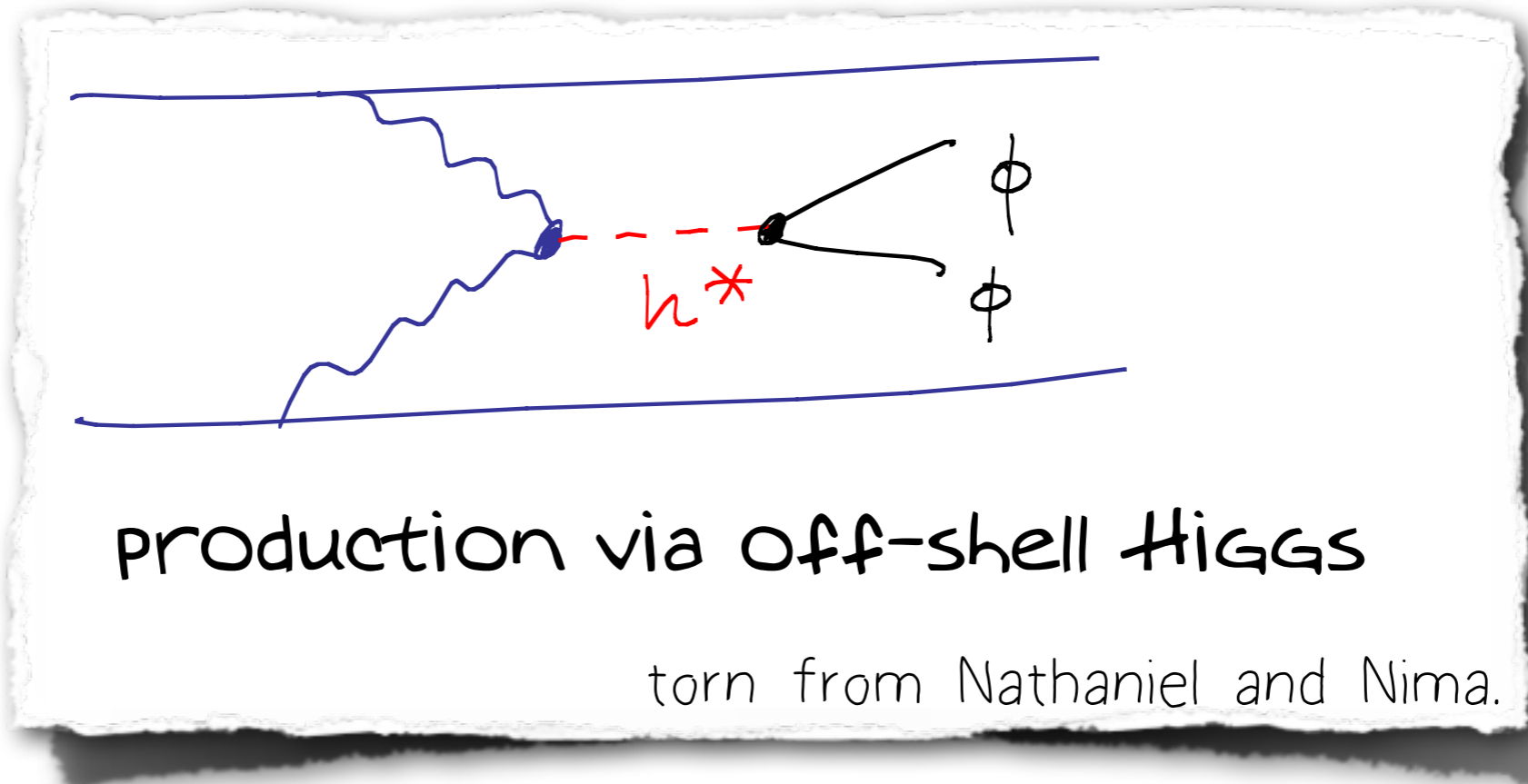


## The Top Partners:

Direct production of top partners  
(or other  $SM_B$  states).

# Colorless Tops

- \* Minimally we have a coupling of the partner to the Higgs:



- \* Partner decays invisibly - background study needed.
- \* Partner could decay visibly, Hidden-valley style! (high multiplicity from hidden showers, displaced vertices, etc.)

# Cosmology

- \* Twin baryons could be dark matter.
- \* Light twin matter is dangerous during BBN.
  - Small  $Z_2$  breaking in light yukawa sector. Make twin matter Heavier than a few GeV. (Barbieri, Gregoire, Hall)
  - More entropy production in our QCD transition?
  - Late inflation?
- \* LHC implications are drastic enough to overshadow these issues.

# Conclusion and Outlook

- \* We have natural models with colorless partners (including some models I did not discuss).
- \* Twin Higgs can be remarkably elusive at LHC.
- \* Seems within reach of a 100 TeV machine & TLEP:
  - o Precision Higgs studies and invisible decay (or hidden valley decays).
  - o Producing UV completion states.
  - o Producing top partners.

*A no-loose theorem is within reach.*

**Deleted scenes.**

# LHC Signals

- \* A standard model Higgs.



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- \* If we are lucky we can measure-
  - Higgs decay to invisibles,  $\text{BR} \sim O(v^2/f^2)$ .
  - Modification of  $ZZh, WWh, tth, h^3, \dots$   
also of  $O(v/f)$ . (correlations).

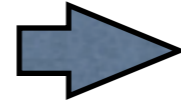
# LHC Signals

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  - Modification of  $ZZh, WW h, tth, h^3, \dots$   
also of  $O(v/f)$ . (correlations).
- \* If we are *really* lucky (work in progress):  
Twin hadrons decay back to SM with displaced vertices. A “Hidden Valley” signal (Strassler and Zurek).

# Folded SUSY (Disney version)

Usually:

Supersymmetry commutes  
with gauge transformations.



Superpartners *always*  
have the same quantum  
numbers as SM counterparts.

How can we get non-colored partners?

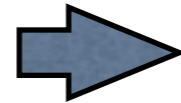
Inspiration: **The Large-N Orbifold Correspondence**

Kachru & Silverstein (98)  
Bershadsky & Johansen (98)  
Schmaltz (99)

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$t$

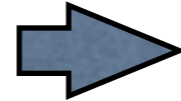
Inspiration: **The Large-N Orbifold Correspondence**

Kachru & Silverstein (98)  
Bershadsky & Johansen (98)  
Schmaltz (99)

# Folded SUSY (Disney version)

Usually:

Supersymmetry commutes  
with gauge transformations.



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How can we get non-colored partners?

$$\begin{array}{c} t \\ \updownarrow Q^\alpha \\ \tilde{t} \end{array}$$

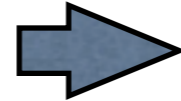
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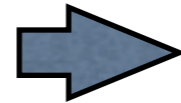
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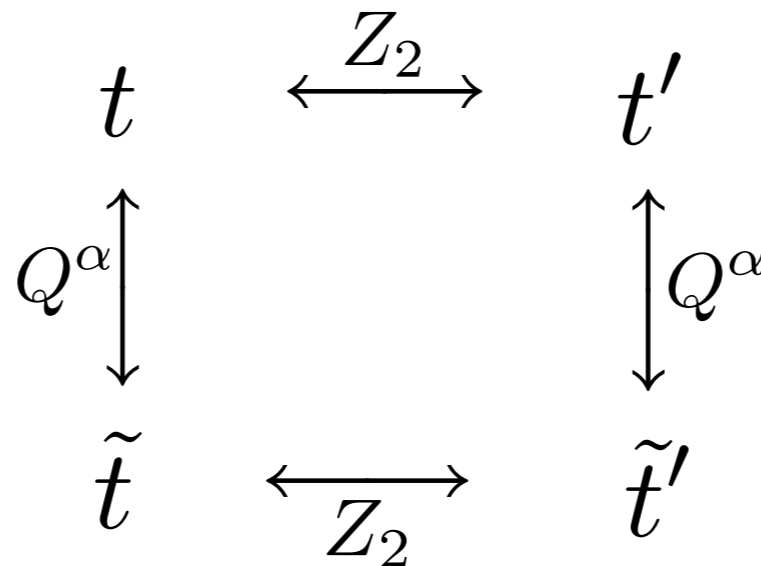
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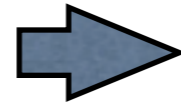
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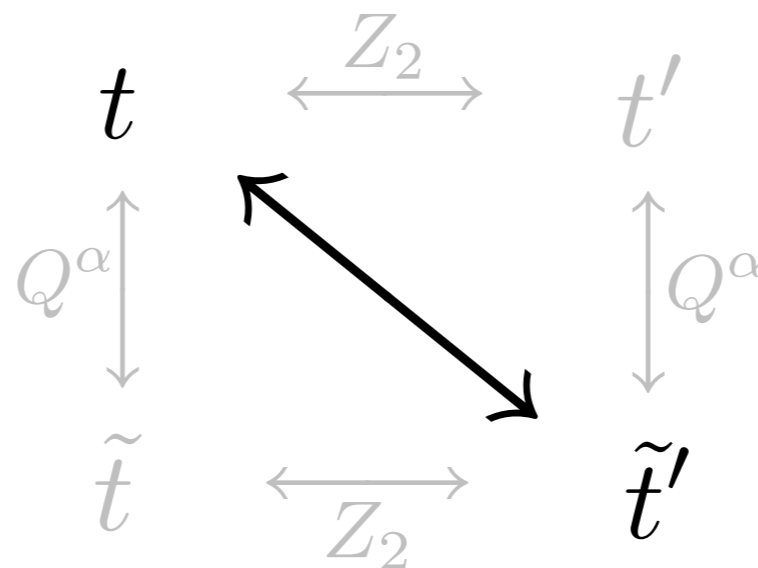
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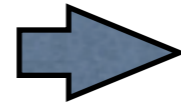
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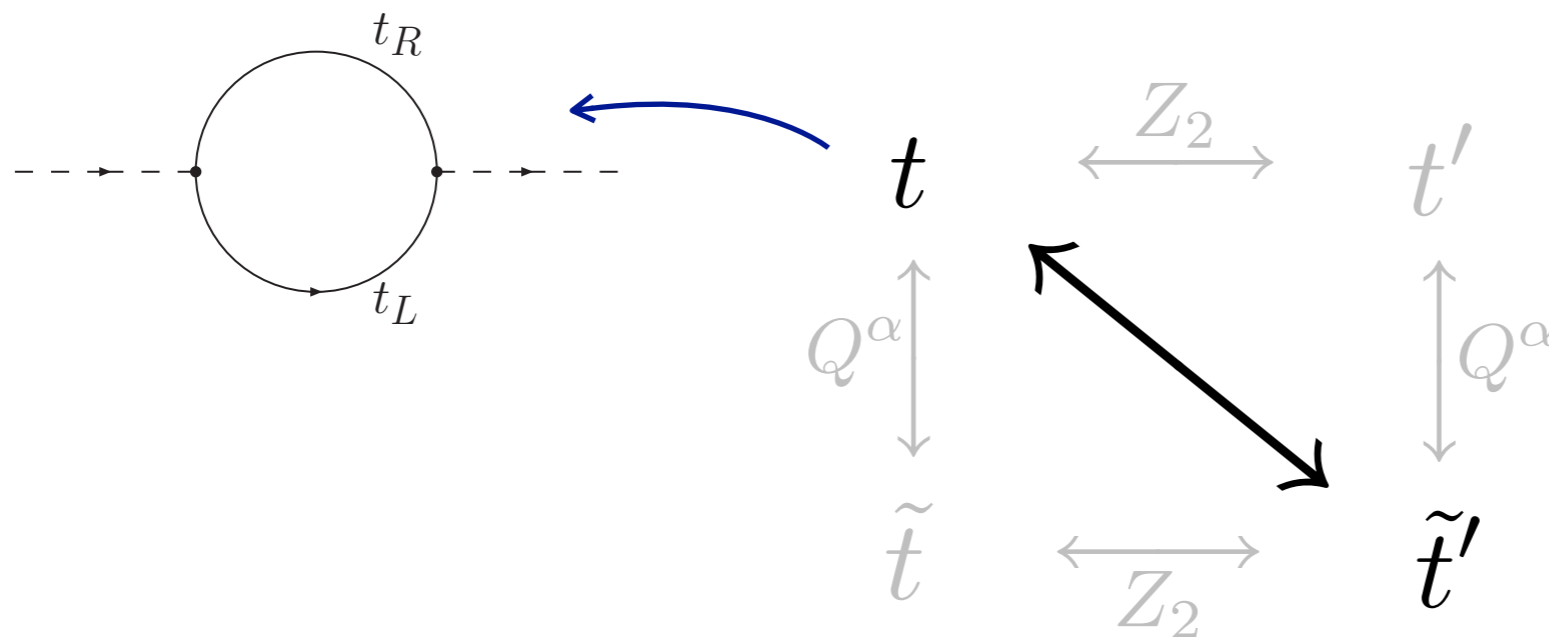
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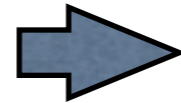


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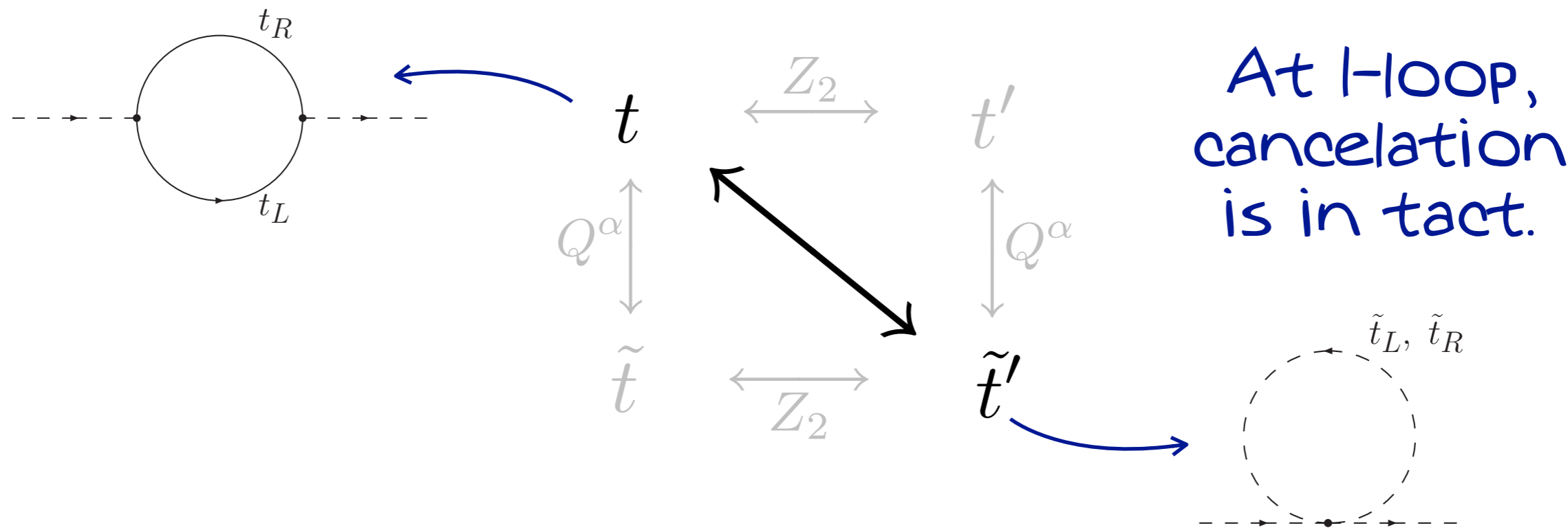
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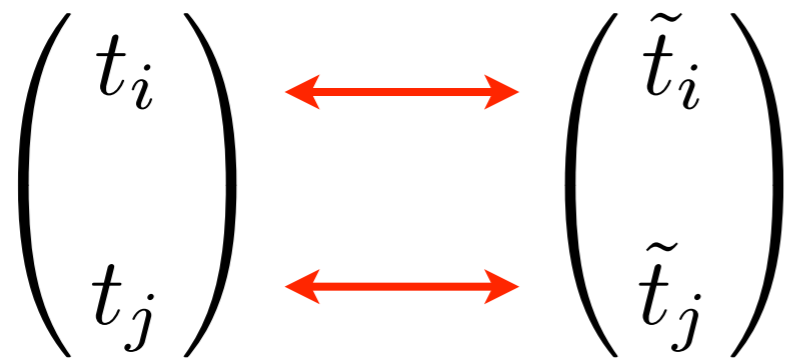


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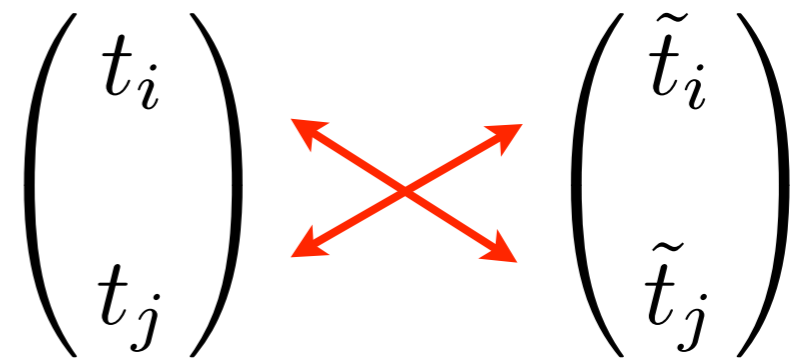
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# Folded SUSY

\* The Higgs is protected twice:



Supersymmetric



Folded-Supersymmetric

We get to choose which states to keep at low energies.