Electroweak B+L Violation at High Energies.

Part 2: Rate Estimates and Signatures

Andreas Ringwald

BSM Opportunities at 100 TeV CERN, Geneva, February 10-11, 2014





Recap

Cross-sections for exclusive B+L violation rapidly growing below

$$4\pi \; \frac{M_W}{\alpha_W} \simeq 30 ~ {\rm TeV}$$

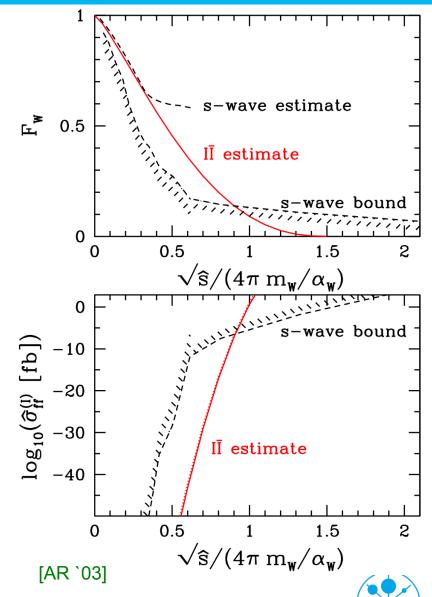
Total cross-section grows exponentially; dominated by multiple emission of EW gauge bosons

$$\hat{\sigma}_{\mathrm{ff}}^{(I_W)} \approx \frac{1}{m_W^2} \left(\frac{2\pi}{\alpha_W}\right)^{7/2} \mathrm{e}^{-\frac{4\pi}{\alpha_W} F_W(\epsilon)}$$

$$\epsilon \equiv \sqrt{\hat{s}}/(4\pi m_W/\alpha_W) \simeq \sqrt{\hat{s}}/(30 \text{ TeV})$$

$$F_W(\epsilon) ~~=~~ 1 - rac{3^{4/3}}{2} \, \epsilon^{4/3} + rac{3}{2} \, \epsilon^2 + \mathcal{O}(\epsilon^{8/3})$$

- For \(\epsilon > 0.3 \dot 0.75\) only estimates, educated guesses and bounds
- Need future hadron collider or search for analogous QCD processes



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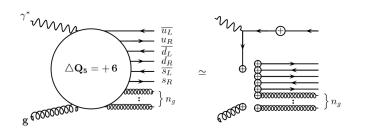
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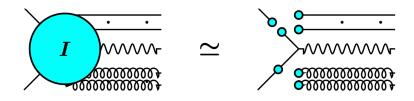
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Hard QCD-instanton induced processes in

DIS



virtual vector boson production



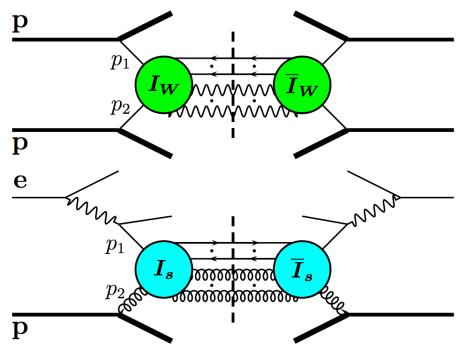


Instanton-Antiinstanton Valley Cross-Section Estimate

> Total cross-section via optical theorem [V.V. Khoze, AR <u>`91; AR, F. Schrempp</u> `98] $\hat{\sigma}_{p_1p_2}^{(I)} \sim \int d^4R \int_0^{\infty} d\rho \int_0^{\infty} d\overline{\rho} D(\rho) D(\overline{\rho}) \int dU e^{-\frac{4\pi}{\alpha g} \Omega\left(U, \frac{R^2}{\rho \overline{\rho}}, ...\right)} e^{i(p_1 + p_2) \cdot R - \sum_{i=1}^2 \sqrt{-p_i^2} (\rho + \overline{\rho})}$

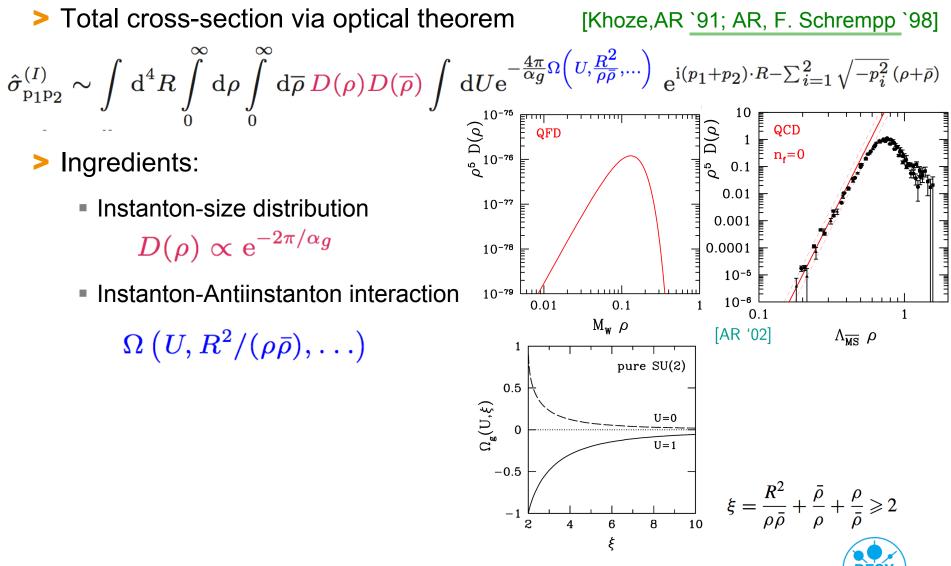
- > Ingredients:
 - Instanton-size distribution $D(
 ho) \propto {
 m e}^{-2\pi/lpha g}$
 - Instanton-Antiinstanton interaction

 $\Omega\left(U,R^2/(
hoar
ho),\ldots
ight)$





Instanton-Antiinstanton Valley Cross-Section Estimate



Instanton-Antiinstanton Valley Cross-Section Estimate

• Saddle point evaluation:

$$\hat{\sigma}^{(I)} \propto e^{-\Gamma_*} \equiv e^{-\frac{4\pi}{\alpha_g} F_g(\epsilon)},$$

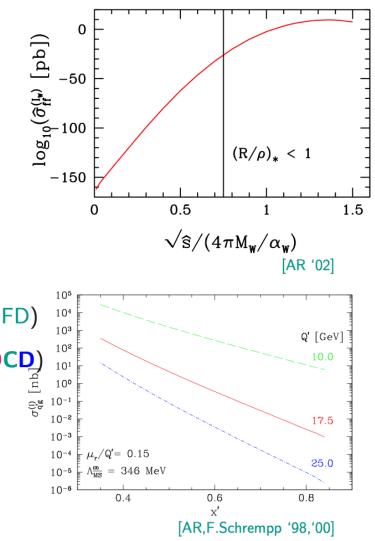
where

$$\epsilon \equiv \begin{cases} \sqrt{\hat{s}}/(4\pi M_W/\alpha_W) & (\text{QFD}) \\ \sqrt{\hat{s}}/Q' \equiv \sqrt{1/x'-1} & (\text{QCD}) \end{cases}$$

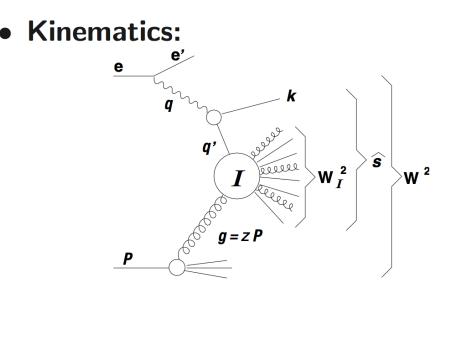
is a scaled cm energy and

$$\begin{split} \mathbf{F}_{g} &= 1 + \Omega_{g}(1,\xi_{*}) + \\ & \begin{cases} -(\xi_{*}-2) \frac{\partial}{\partial\xi_{*}} \Omega_{g}(1,\xi_{*}) \\ 0 \end{cases}_{|\xi_{*}=2 + \left(\frac{R}{\rho}\right)_{*}^{2}} & (\mathbf{Q}) \end{cases} \end{split}$$

• Increasing $\epsilon \Rightarrow$ smaller $(R/\rho)_*$ probed \Rightarrow cross-section grow due to attractive nature of Ω_g in perturbative semiclassical regime







Deep-inelastic scattering variables: $S = (e + P)^2$ $Q^2 = -q^2 = -(e - e')^2$ $x_{Bj} = Q^2 / (2P \cdot q)$ $y_{Bj} = Q^2 / (S x_{Bj})$ $W^2 = (q + P)^2 = Q^2(1/x_{Bj} - 1)$ $\hat{s} = (q + g)^2$ $z = x_{Bj} (1 + \hat{s}/Q^2)$

Variables of instanton-subprocess:

 $Q'^{2} = -q'^{2} = -(q-k)^{2}$ $x' = Q'^{2} / (2 g \cdot q')$ $W_{I}^{2} = (q'+g)^{2} = Q'^{2}(1/x'-1)$

• "Fiducial" kinematical region from lattice constraints: [AR,F.Schrempp '99;'01]

$$\left(\rho^* \Lambda_{\overline{\mathrm{MS}}}^{(0)} \!\lesssim\! 0.4, \frac{R^*}{\rho^*} \!\gtrsim\! 1.0\right) \Rightarrow \left(Q' / \Lambda_{\overline{\mathrm{MS}}}^{(n_f)} \!\gtrsim\! 30.8, x' \!\gtrsim\! 0.35\right)$$



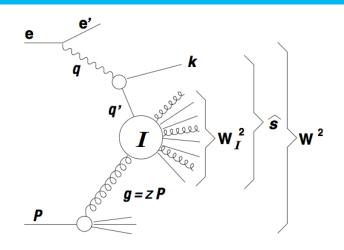
Event generator QCDINS 2.0:

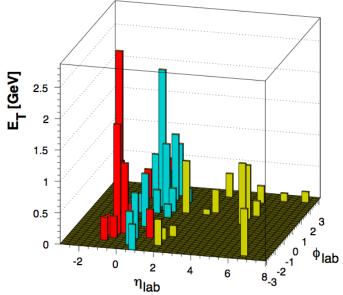
[Gibbs,AR,F.Schrempp '95; AR,F.Schrempp '00]

- Hard subprocess:
 - isotropic in $q^\prime g~{\rm CM}$
 - flavour democratic
 - large parton multiplicity

$$\langle n_q + n_g \rangle = 2 n_f - 1 + \mathcal{O}(1) / \alpha_s \gtrsim 8,$$

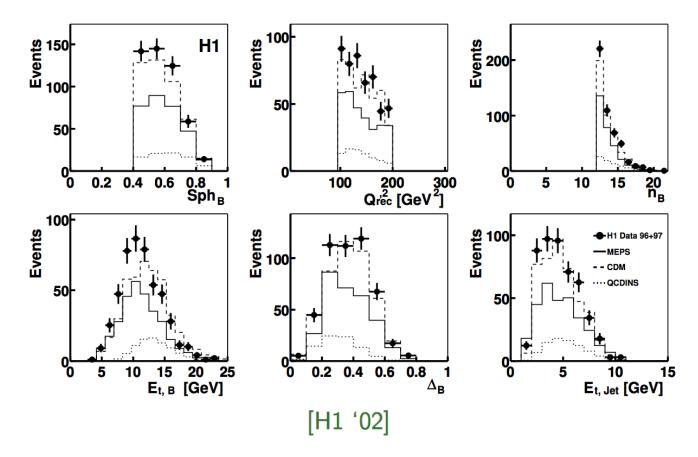
- Parton shower (HERWIG)
- Hadronization (HERWIG or JET-SET)





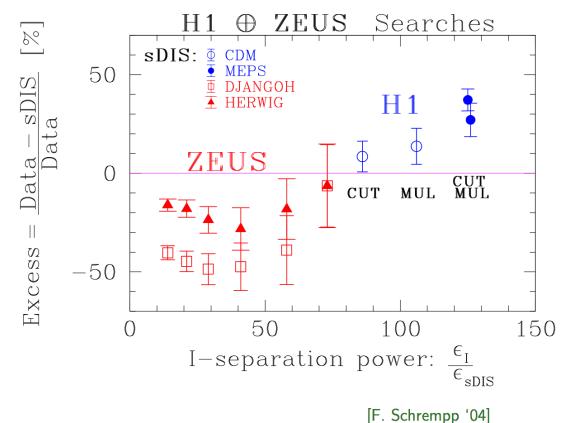


- Instanton-enriched samples by cuts on discriminating observables
- Large uncertainties in predictions of standard DIS processes





- Instanton-enriched samples by cuts on discriminating observables
- Large uncertainties in predictions of standard DIS processes
- H1/ZEUS "excess" increases with separation power (ratio of efficiencies)

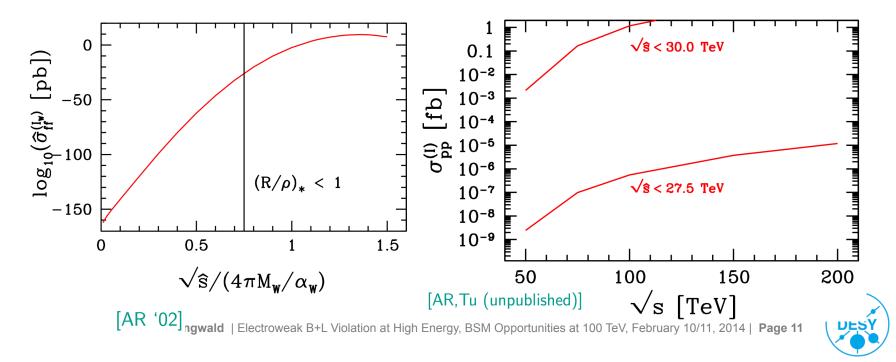




QFD-Instanton Induced Processes at Future Hadron Coll.

> H1/ZEUS limits on QCD-instanton induced processes suggest:

- Instanton-antiinstanton estimate reliable, as long as $(R/\rho)_* \geq 1$
- For $(R/\rho)_* < 0.5 \div 1$, rapid growth, as implied by Ω , stops.
- Implications for QFD-instantons:
 - $(R/\rho)_* < 0.5 \div 1$ corresponds to $\epsilon < 0.75 \div 1.15$, $\sqrt{\hat{s}} < 22 \div 35$ TeV
 - At these energies, cross-section estimate reaches observable values



Phenomenology of QFD-instantons

[AR,F.Schrempp,Wetterich '91; Gibbs,AR,Webber,Zadrozny '94]

- No background from perturbative Standard Model processes by requiring
 - ≥ 4 identified charged $e\mbox{'s}$ or $\mu\mbox{'s}$
 - $E_T \ge$ several TeV
- Event generator **HERBVI**:

[Gibbs,Webber '95]

- B-violation cannot be established
- L-violation verifiable: measure

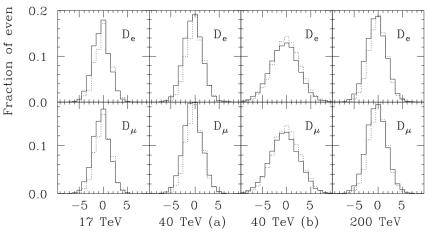
 $D_{\ell} = N_{\ell^{-}} - N_{\ell^{+}};$

need $\sim 10^3$ events

[Gibbs,AR,Webber,Zadrozny '94]

Simulations performed			
Energy (TeV)		n_B estimate	$\sqrt{\hat{s}_0}$ (TeV)
17		$1/lpha_W$	5
40	(a)	$1/lpha_W$	18
40	(b)	LOME	18
200		$1/lpha_W$	18

[Gibbs,AR,Webber,Zadrozny '94]





Summary

- Electroweak B+L violation central building block of our understanding of baryogenesis in big bang cosmology
- > Characteristic scale of B+L violation

$$4\pi \frac{M_W}{\alpha_W} \simeq 30 \text{ TeV}$$

- Cross-sections for B+L violating processes
 - exponentiall small below this scale
 - exponentially growing below this scale
 - may reach observable values near this scale
- New computational methods needed
- > Unique opportunity for a 100 TeV collider!
- In the meantime, search for small-size QCD-instanton induced processes



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- [V.V. Khoze,AR `91] "Nonperturbative contribution to total crosssections in nonAbelian gauge theories", Physics Letters B 259 (1991) 106
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