

- **Compact Compton sources**

Alessandro Variola
In2p3 - CNRS

LA3NET
Salamanca

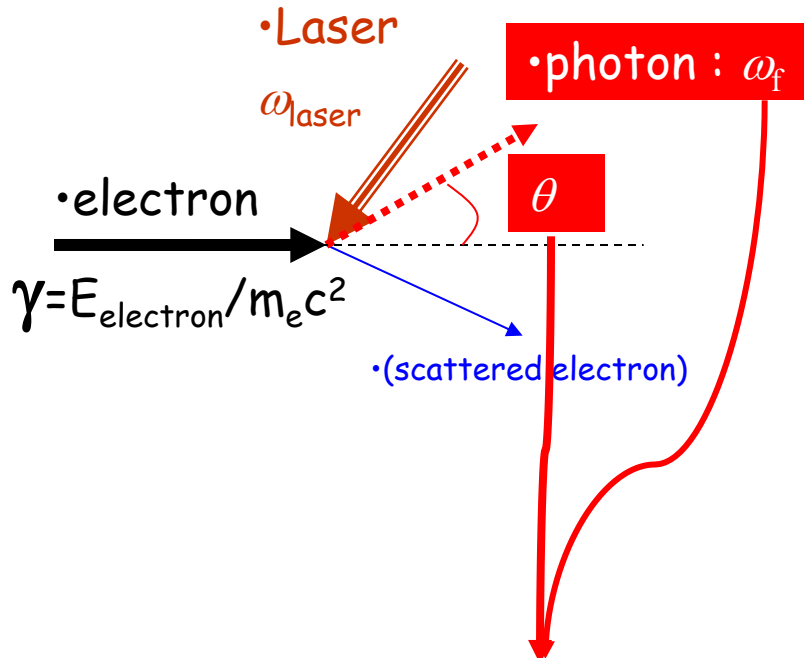
- Introduction : Compton Effect for compact sources
- Compact Compton Sources Why
- Compact Compton Sources How?
- Projects

Compact Sources : Compton backscattering

- **Compact sources -> Compton Backscattering effect (CBS)**
- **Why CBS for Compact Sources?**
 - CBS is by far the most efficient photon energy amplifier : $\omega_{\text{diff}} = 4\gamma^2 \omega_{\text{laser}}$, for example if $\gamma \sim 100 \Rightarrow$ it is possible to have at one's disposal hard X rays (50 keV) with a relatively low energy electron machine.
 - **But for a light source:** $\sigma \sim 6.6524 \cdot 10^{-25} \text{ cm}^2$, it is low!!!! Need to increase lasers and accelerators performances
 - Two kind of compact source. Strong impact on the accelerator and on the optical system choice :
 - High average flux. Mainly based on high repetition frequency. Storage rings / SRF linacs / High average power lasers coupled with FP cavities.
 - High peak brilliance. Mainly based on the laser pulse energy. Warm linacs/ High power lasers/ Optical re-circulators to increase the flux

- CBS attractiveness :
- 1) relativistic boost and directivity $\Rightarrow f = 1/\gamma$ around the electron direction
- 2) Energy angle dependence \Rightarrow monochromatic by diaphragm
- 3) Polarized if needed
- 4) Backscattered spectrum cut off \Rightarrow Energy dependence on collision angle

Compton backscattering



•Dynamics of the process

- Compton Scattering γ (laser) + e \rightarrow γ' + e'
- Photon(laser) + electron scattering
 - Fano, JOSA39(1949)859;
- Simulation programme: CAIN from Yokoya-san
<http://lcdev.kek.jp/~yokoya/CAIN/cain235/>

•We are interested by using the scattered photon

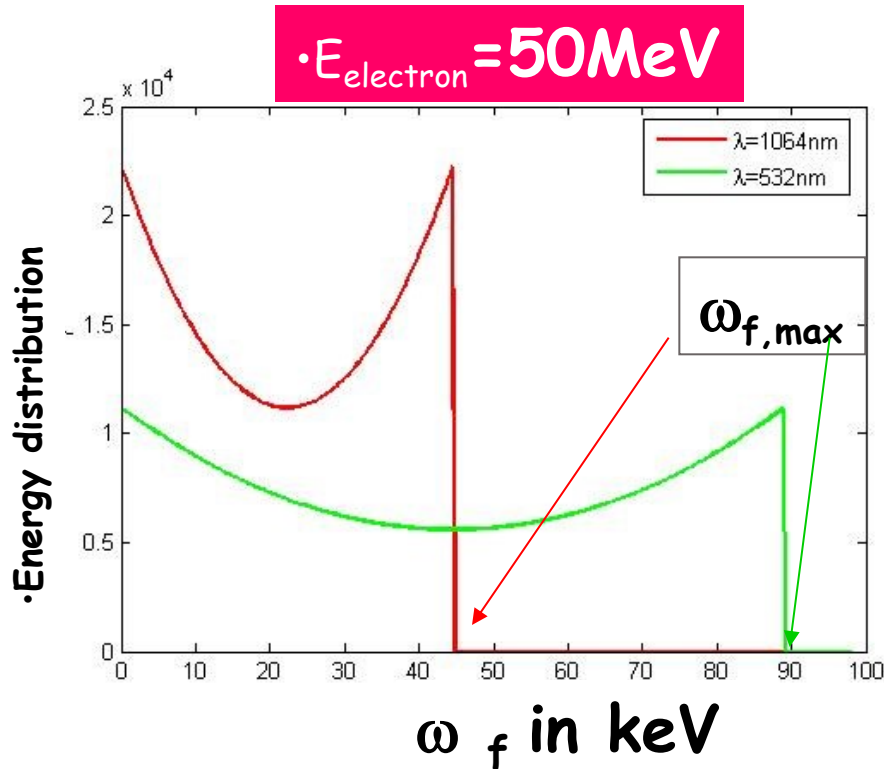
•Scattered photon properties given by the Compton differential cross-section:

$$\frac{d\sigma}{d\Omega^*} = \sigma_0 + \sigma_1 + \sigma_2 + \sigma_3 + \sigma_4$$

•Independent from polarisations

•Polarisation of the 4 particles are observed

• 1st interest (and 4th): the energy boost



• Sprangle et al. JAP72(1992)5032

Energy distribution ~flat with

$$\omega_{f,max} = 4\gamma^2 \omega_{\text{laser}}$$

with $\gamma \sim 100$ ($E_{\text{electron}} = 50\text{MeV}$)



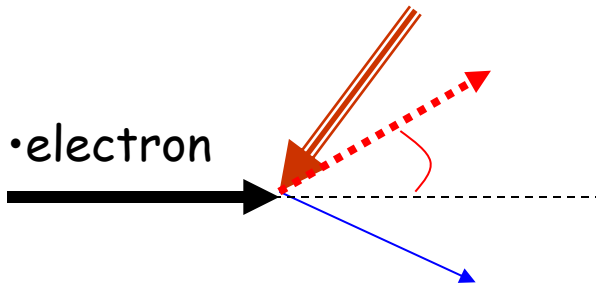
$$\omega_{f,max} = 45000\text{eV if } \omega_{\text{laser}} \sim 1\text{eV}$$

Angular emission $\Rightarrow 10$ mrad

Compton scattering is the most powerful mechanism to boost photon energies

The cut off is dependent on the incidence angle !! (factor two up to $\pi/2$)...we can play to make short pulses

- 2nd interest: the angular energy correlation,
- Monochromatic sources

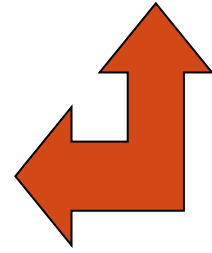
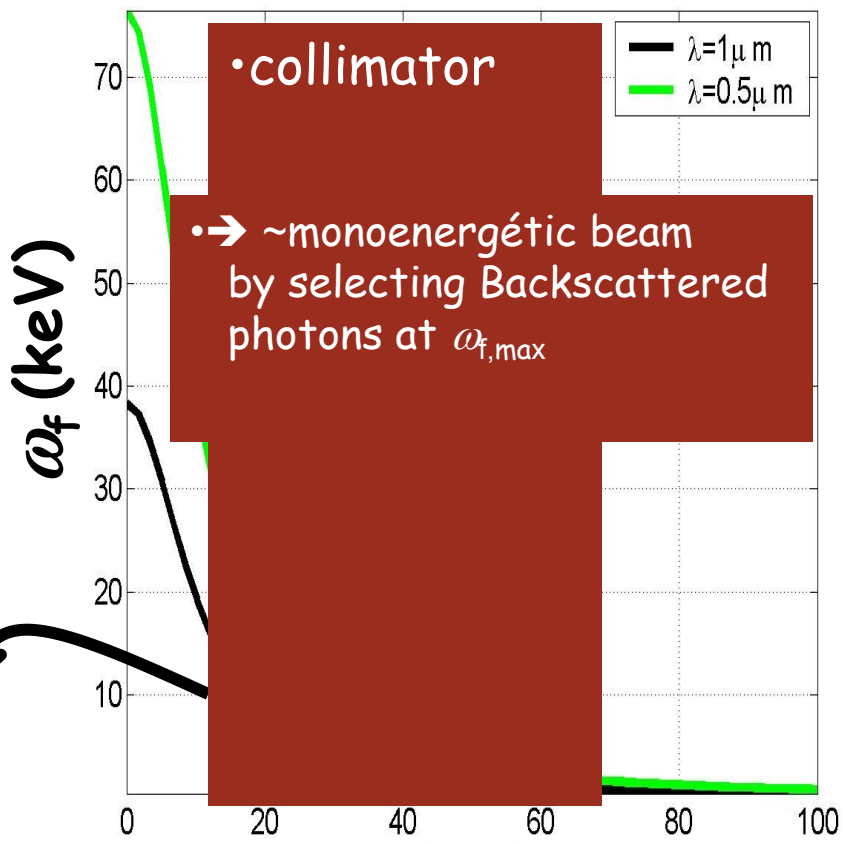
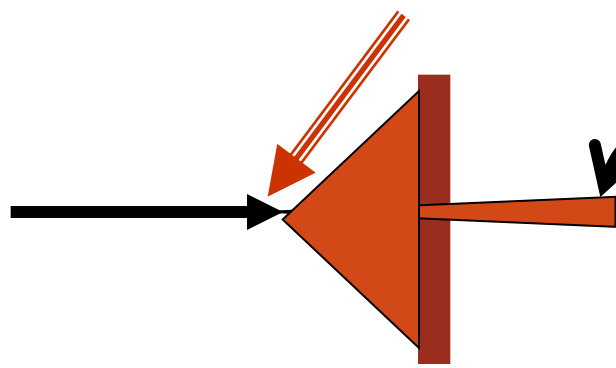


• Compton scattering

• Photon_{laser}+e → photon+e'

2 body process → $\omega_f = f(\theta)$

• Sprangle et al. JAP72(1992)5032



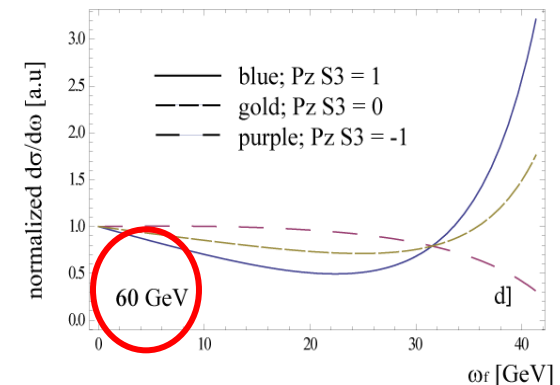
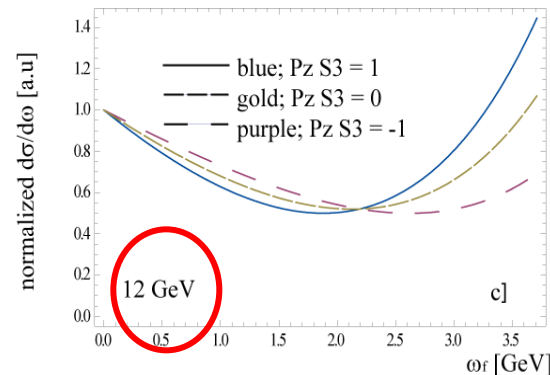
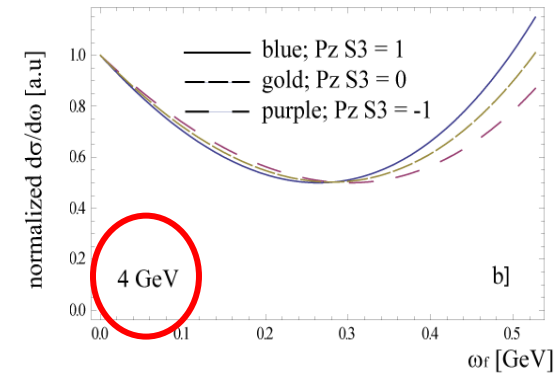
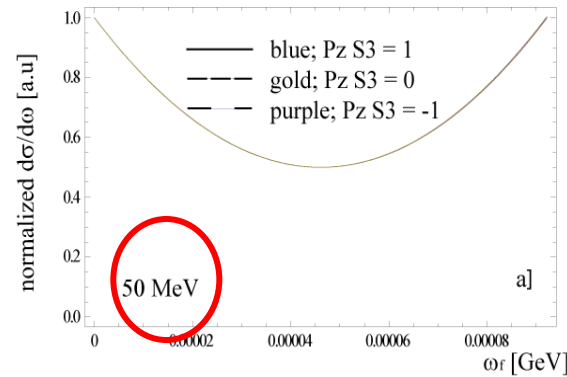
• $E_{\text{electron}} = 50 \text{ MeV}$

θ (mrad)

•3rd interest: incident electron and laser polarisation effects

•Differential Compton cross-section with 2 polarisations observed (energy distribution):

- $d\sigma/d\omega_f = A_0(\omega_f) - P_e S_3 A_1(\omega_f)$
- A_0, A_1 : known (QED)
- S_3 : laser degree of circular polarisation
- P_e : e^- longitudinal polarisation



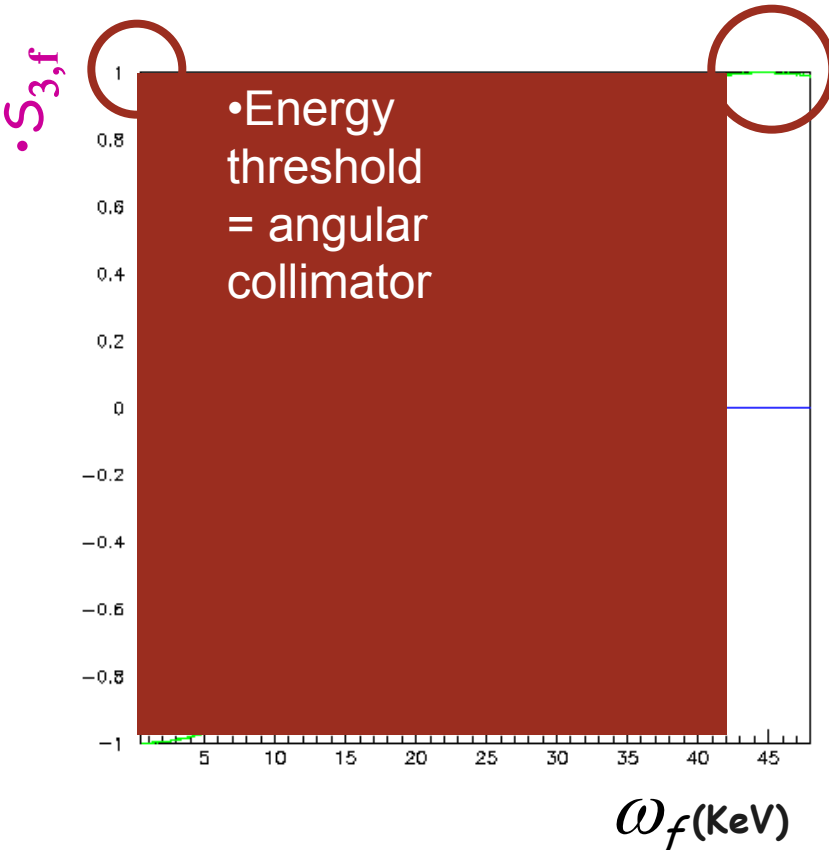
→ Knowing S_3 one can determine the polarisation of electrons above $\sim 4\text{GeV}$

→ electron/positron Compton polarimeters used in accelerators

e.g. Barber et al. Nucl.Instrum.Meth.A329(1993)79

•4th interest: polarisation effects in the final state

$$\text{---} \cdot (P_e, S_{3,\text{laser}}) = (0, 1)$$



$$\cdot S_{3f}=1 \text{ for } \omega_f = \omega_{f,max} \text{ \& } S_{3,laser}=1$$

- Compton scattering acts as a mirror for circular polarisation at low energy **if** highest values of ω_f are selected
- (i.e. backscattered photons are selected)

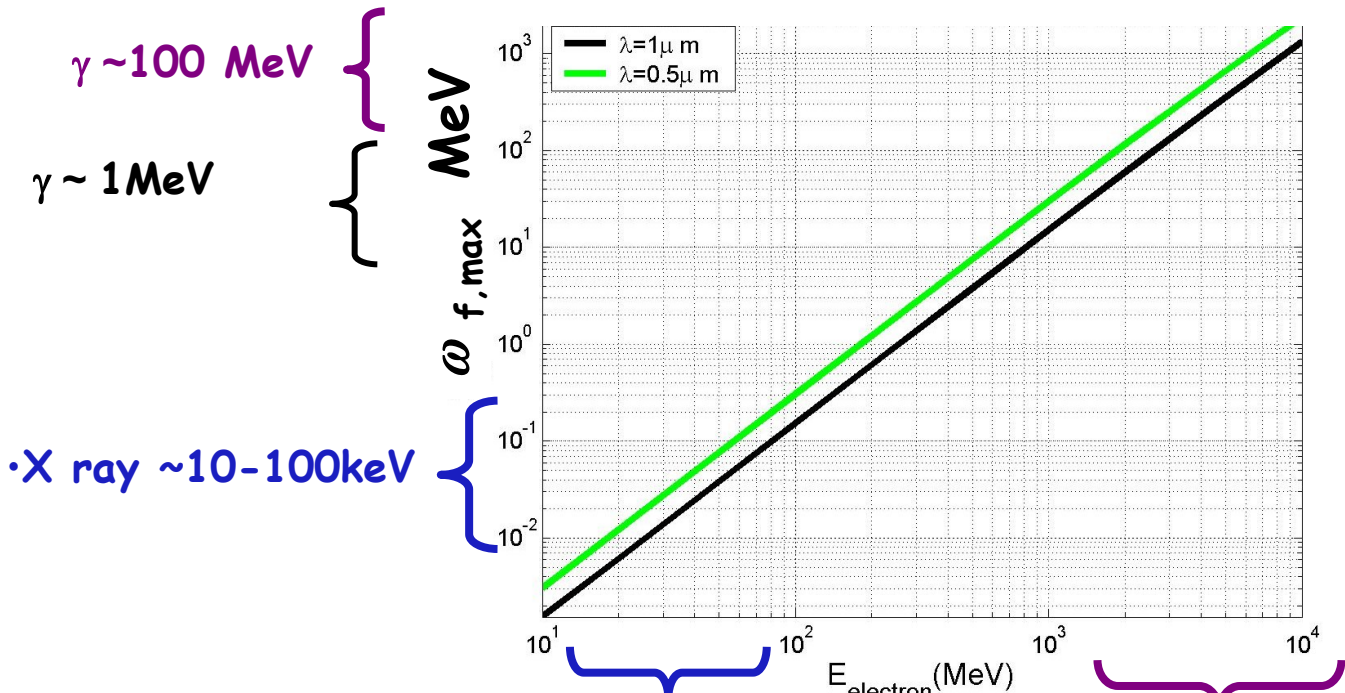
→ $\gamma\gamma$ collider ($E_{\text{electron}} \sim 250-500 \text{ GeV}$)

Ginzburg et al. NIM219(1984)5

→ Polarised positron source ($E_{\text{electron}} \sim 1 \text{ GeV}$)

→ Polarised Light source

• **SUMMARY** : Applications of Compton scattering.
 Quasi monochromatic X/ γ ray beam



$\gamma \sim 100$ MeV }
 $\gamma \sim 1$ MeV }

• X ray $\sim 10-100$ keV }

• Elec. $\sim 20-100$ MeV

• Elec. $\sim 100-500$ MeV

• Elec. 1 GeV

- Low energy applications
- Medical: radiography & radiotherapy
 - Museology
 - Material science
 - crystallography

- Nuclear fluorescence applications
- Nuclear survey
 - Nuclear waste management

- High energy applications
- Compton polarimeter \rightarrow LEP energy measurement
 - Laser wire
 - gg collider
 - Polarised positron source

A more classical point of view...intense laser as an helical undulator

- Helical undulator and laser wave ($v=c=1 \rightarrow$ ultrarelativistic case)

- Fields : Undulator
$$H_{LabSys} = B(e_1 \cos(kz) + \rho e_2 \sin(kz)), \quad k = 2\pi/\lambda_u$$

$$\rho = \pm 1 = \text{Helicity (p}_\perp \text{ sense of rotation)}$$

- In the Electron rest system an E field appears:

$$H_{RSys} = \gamma B(e_1 \cos(\theta) + \rho e_2 \sin(\theta)),$$

$$E_{RSys} = \gamma B(-\rho e_1 \sin(\theta) + e_2 \cos(\theta)),$$

- The laser field is :
$$A_{LabSys} = a(e_1 \cos(kz) + \rho e_2 \sin(kz)), \quad k = 2\pi/\lambda_l$$

And in rest system

$$H_{RSys} = -2\rho\gamma k_l a(e_1 \cos(\mathcal{G}) + \rho e_2 \sin(\mathcal{G}))$$

$$E_{RSys} = -2\rho\gamma k_l a(-\rho e_1 \sin(\mathcal{G}) + e_2 \cos(\mathcal{G}))$$

- Same ' kind ' of trajectory solution :
$$\bar{p}_\perp = \langle \bar{p}_\perp \rangle + p_\perp (e_1 \cos(\theta) + \rho e_2 \sin(\theta)),$$

Field intensity : important factor

- Radiation depends on the laser intensity parameter

$$\xi^2 = \frac{e^2 \langle A^2 \rangle}{m^2} \quad \text{that in the undulator case} = K^2 = \frac{e^2 B^2}{(k m)^2}$$

For a laser $\xi^2 = 3.66 \cdot 10^{-19} \lambda^2 (\text{mm}) P (\text{W}/\text{cm}^2)$

For an undulator = $(0.935 \text{ B(T)} \lambda_u (\text{cm}))^2$

- Meaning (going back to equations) $p_{\perp} = \xi m$

So in the ultrarelativistic case and $v_{\perp} \ll 1$

$$\xi^2 = \frac{\langle p_{\perp}^2 \rangle - \langle p_{\perp} \rangle^2}{m^2} \quad \text{deviation from a straight line (average angular deviation)}$$

$$\xi^2 \approx \alpha N_{\text{int}} \quad N_{\text{int}} = \text{number of laser photons in the interacting volume } V_{\text{int}}$$

$$V_{\text{int}} = \lambda_c^2 \lambda_{\text{laser}} \quad \text{with } \lambda_c = 3.86 \cdot 10^{-11} \text{ cm}$$

- Why a "compact" Compton $X(\gamma)$ -ray source ?

▶ **In many scientific domains**

synchrotron sources are currently the only machines in term of brightness to perform and carry out **the most ambitious analyses and searches** requiring **~ 10-100 KeV X-rays**.

▶ **Synchrotron sources :**

- very powerful, but,
- not very "practical" for some applications,
- limited access time.

▶ **With Compact sources :**

Methods currently used at synchrotrons (diffraction, absorption, diffusion, imaging, spectroscopy...) could be largely developed in **a laboratory size environment (hospitals, labs, museums)**.

What about 'compact' source for nuclear physics photons (MeV range)?

- Nuclear safety
- Nuclear waste management
- Radioisotopes detection

The compact X-sources today

15

▶ X-ray tubes

- The most efficient are rotating anodes
- Rigaku $\sim 10^{10}$ ph/sec , polychromatic

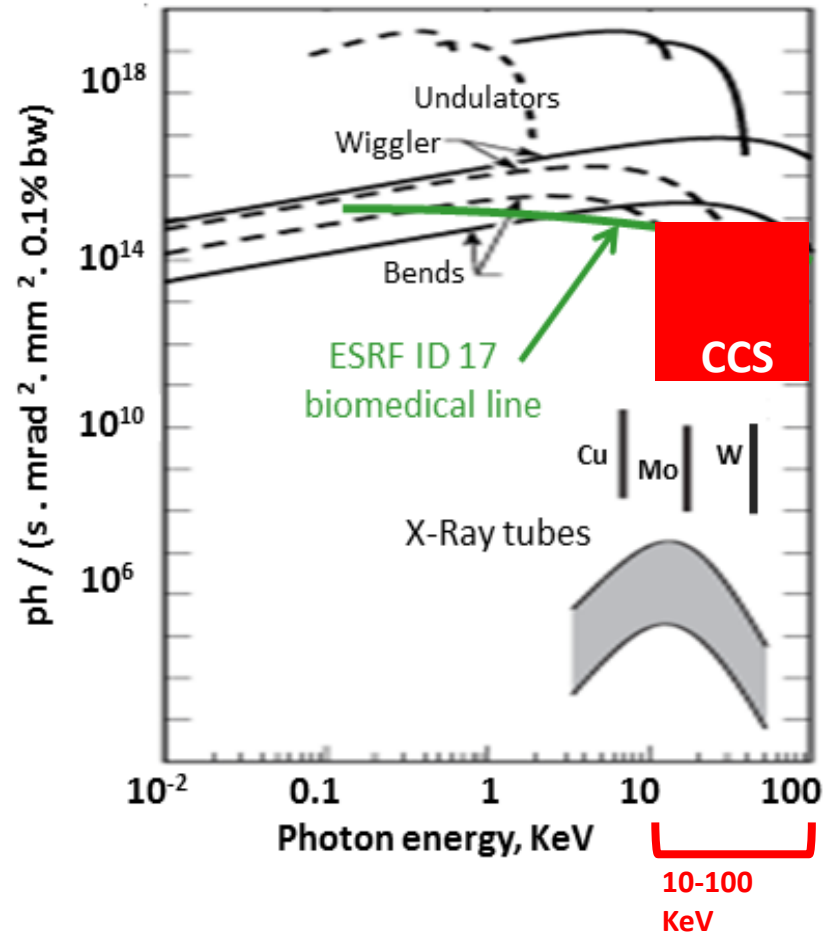
▶ Plasma sources

Ultra-short pulses \sim fs ,
but very low fluxes.

→ These sources does not allow to carry out many of the techniques used at synchrotrons

▶ Compact Compton Sources (CCS)

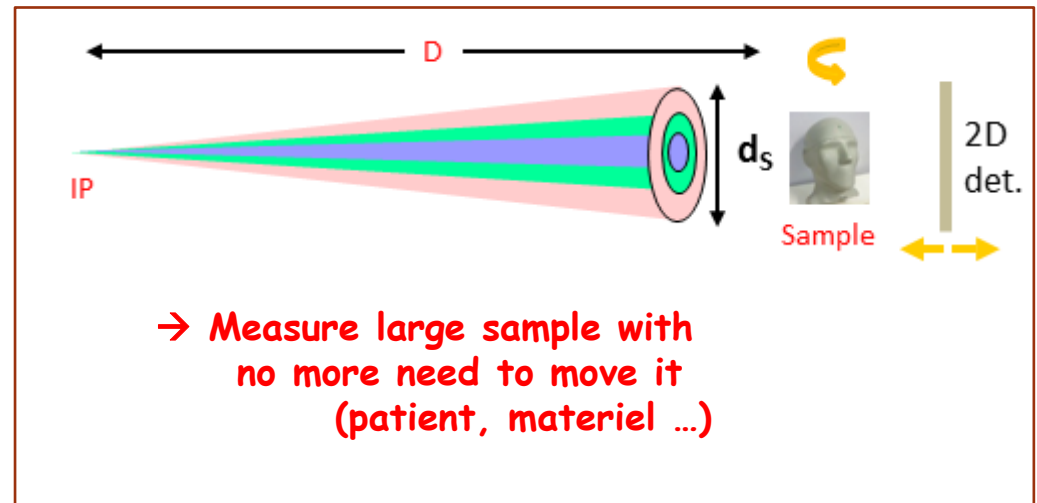
- **Compactness** (surface ~ 100 m²)
(Integration in hospitals, labs, museums)
- **Relative high intensity** (10^{12} - 10^{14} ph/sec)
- **Tunable beam** (Linac configuration)
- **High quality beam** (brightness 10^{11} - 10^{15} ph/sec/ mm² / 0.1% bw / mrad²)



1. Using the 2D divergent beam (biomedical and cultural heritage applications)

- Conventional radiography
- K-edge subtraction imaging
- Phase contrast imaging
- Magnification
- Radiotherapy

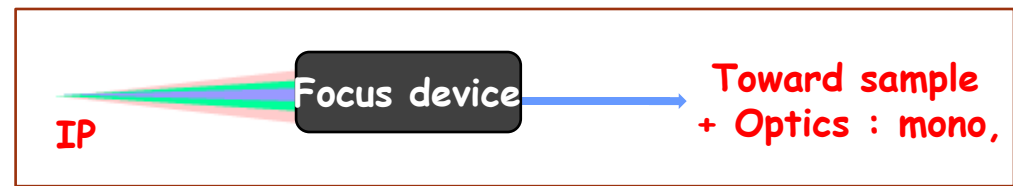
IMAGING



Pink beam (3-30% bw)

2. Using the central part of the beam (cultural heritage / material science applications)

- Fluorescence Spectroscopy
- XANES Spectroscopy
- Diffraction
 - Structural analyses
 - Pump-probe experiments

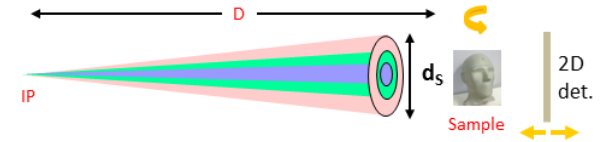


Quasi-monochromatic beam ($\sim 0.1\% - 0.01\%$ bw)

- Example : medical application

Potential of applications of X-ray CCS: Examples

1. Using the 2D divergent beam (biomedical and cultural heritage applications)



- **Conventional radiography** →
 - K-edge subtraction imaging
 - Phase contrast imaging
 - Magnification
 - Radiotherapy
- **High energy** ($\sim 80\text{KeV}$) to test high-Z element drug
 - **No need of monochromaticity** (pink beam, bw $\sim 30\%$)

Ex. : Human head phantom



- 5 mrad opening angle
- $d_s = 12\text{ cm}$ at $D \sim 15\text{ m}$
- $6 \cdot 10^{12}\text{ ph/s}$
- bw 60-90 KeV

CCS and Synchrotron (ID17/ESRF) → comparable

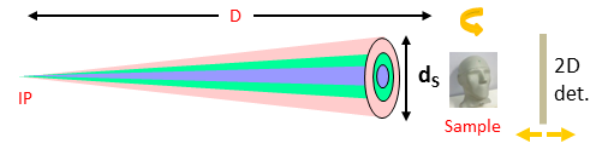
- reduction of the dose
- better image quality

compared to **hospital sources**
(broad spectrum, low flux)

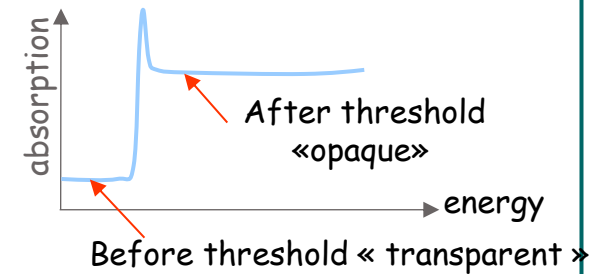
1. Using the 2D divergent beam

(biomedical and cultural heritage applications)

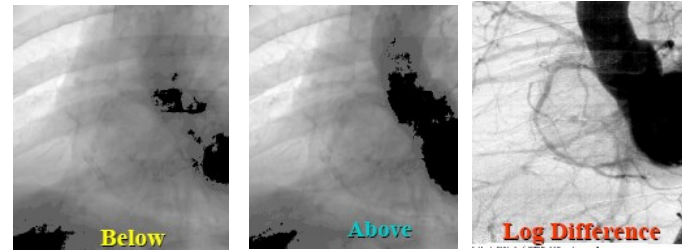
- Conventional radiography
- **K-edge subtraction imaging**
- Phase contrast imaging
- Magnification
- Radiotherapy



- **Tunable energy**
- **bw 2-3%**



K-edge at ESRF (using a contrast agent)



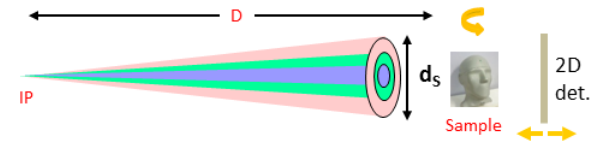
→ The difference of both increase the contrast



- 45 KeV, 2-3% bw
- $d_s = 4 \text{ cm}$ at $D \sim 15 \text{ m}$
- 10^{12} ph/s

Potential of applications of X-ray CCS: Examples

1. Using the 2D divergent beam (biomedical and cultural heritage applications)

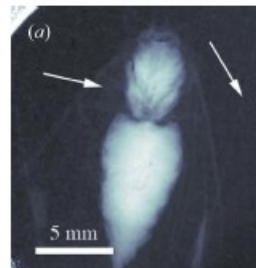


- Conventional radiography
- K-edge subtraction imaging
- **Phase contrast imaging**
- Magnification
- Radiotherapy

- **bw 2-3%**
- **Small source size** (to have transv. coherence)

[Synch. Rad. 16, 2009, 43-47]

CS Lyncean Tech. (only CCS in operation in the world)



13.5 KeV , 3% bw
 10^9 ph/sec
 $\sigma = 165 \mu\text{m}$

Proof of principle

standard absorption phase-contrast

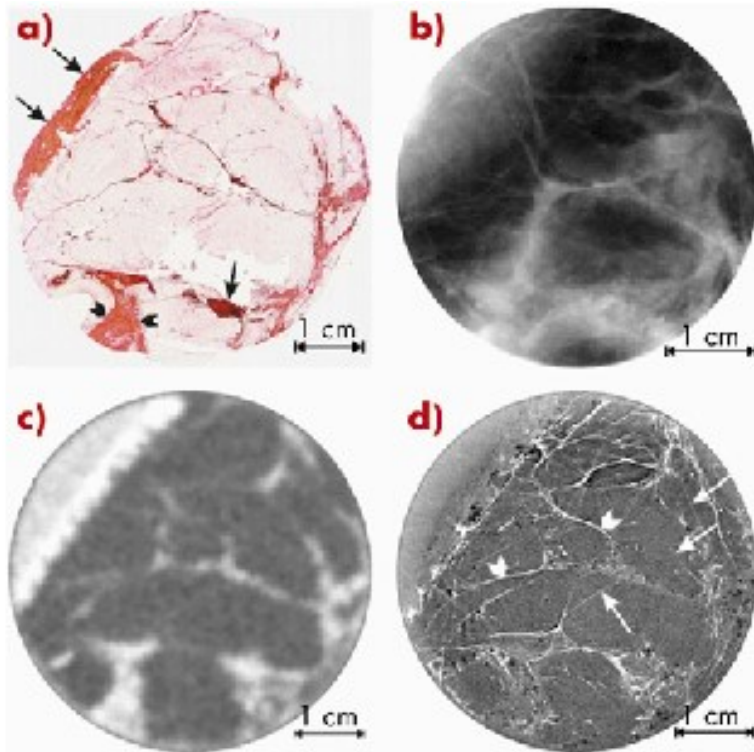


- 70 KeV, 2-3% bw, $\sigma \sim 70 \mu\text{m}$
- $d_s = 4 \text{ cm}$ at $D \sim 15 \text{ m}$
- 10^{12} ph/s

Hospital sources

(large focal spot size, broad spectrum, low flux)

Mapping of a breast tissue sample



a) Histological section
(used as a standard for interpretation)

b) Clinical planar screen-film
mammogram taken at the hospital

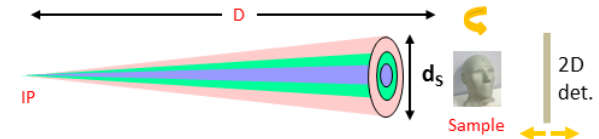
c) Clinical scanner

d) ID17 ESRF (Phase contrast imaging)
Same dose as c)

Stronger contrast

→ Improvement in the visualisation of
the morphology and of the overall
architecture of the breast tissues

1. Using the 2D divergent beam (biomedical and cultural heritage applications)



- Conventional radiography
- K-edge subtraction imaging
- Phase contrast imaging
- Magnification
- **Radiotherapy**

- **High energy** ($\sim 80\text{KeV}$)
- **bw** $\sim 10\%$



- $80\text{ KeV} \pm 10\text{ KeV}$
- $d_s = 5\text{ cm}$ at $D \sim 10\text{ m}$
- $3 \cdot 10^{12}\text{ ph/s}$

Ex. : Human head tumor
(tumor deliver dose $\sim 10\text{-}20\text{ Gy}$)

- **ThomX** $\rightarrow 9\text{ mGy/sec} \rightarrow 20\text{-}30\text{ min}$ of irradiation
- **ESRF/ID17** ($\sim 6\text{ mGy/sec}$)
- **Hospital sources** \rightarrow broad spectrum,
and continuously operation not possible

- Example : cultural heritage application

Examples : Cultural Heritage

Conventional X Radiography & IR Reflectography



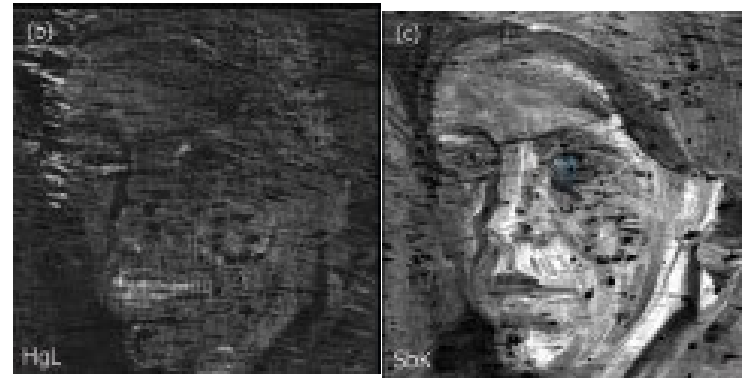
X Transmission IR Reflectography

(Anal. Chem. 80, 2008, 6436-6442)

Vincent van Gogh "Un coin d'herbe" (1887) at synchrotron DESY

Analyses at synchrotron DESY (non destructive)

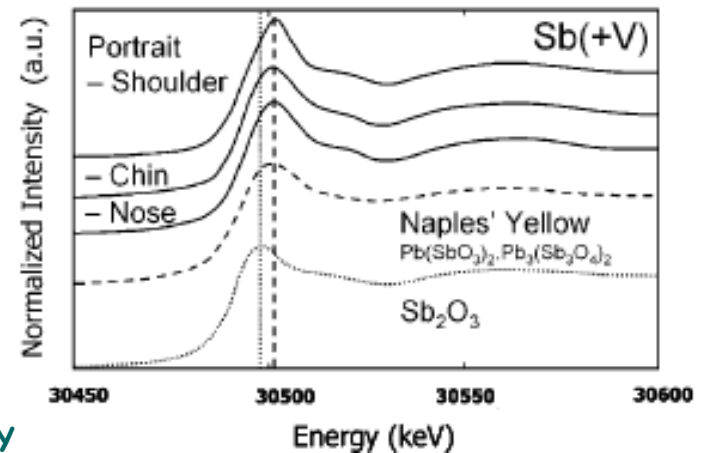
Fluorescence X (38.5 KeV) → Pigment identification



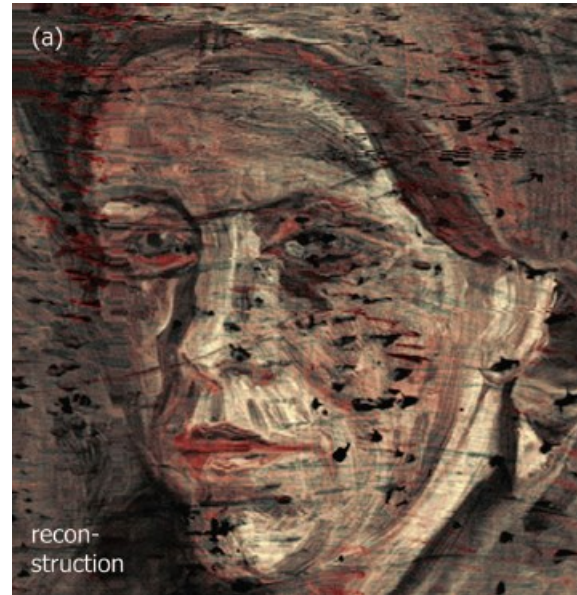
Vermillon

Antimoine

XANES : absorption at the K-edge of an element to identify precisely the chemical components (colors)



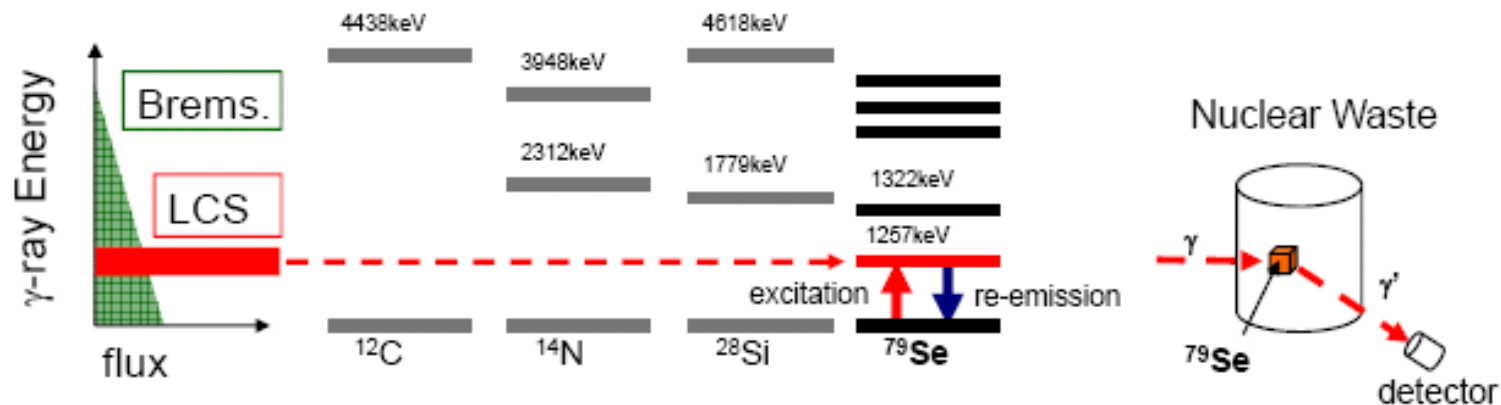
Colored reconstruction



Work painted
some years before

- Example : nuclear science application

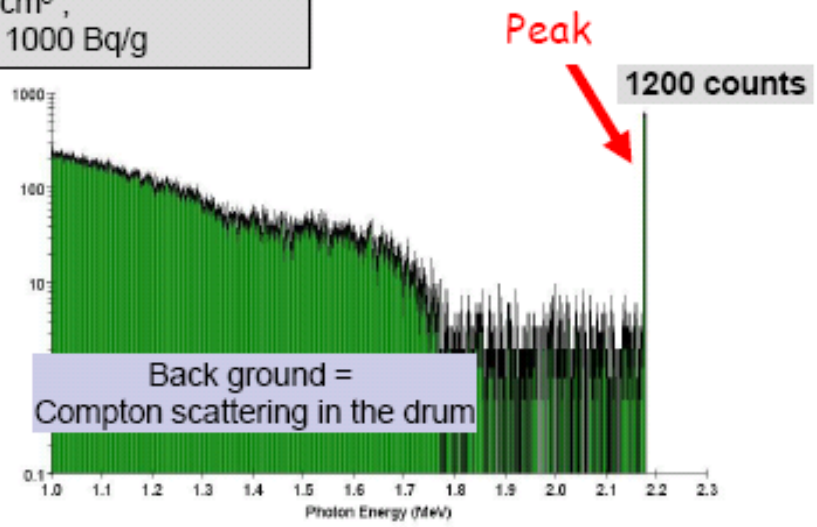
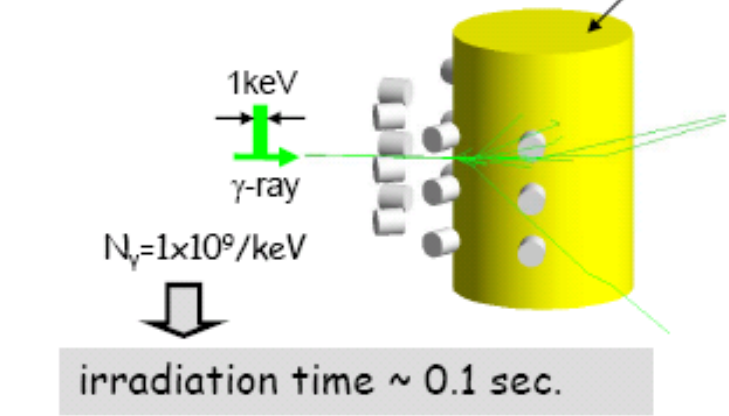
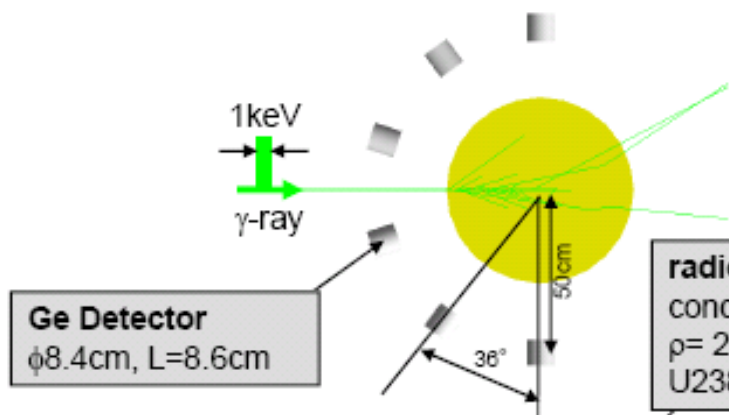
■ Nondestructive Assay by Nuclear Resonant Fluorescence



- Irradiation of γ -rays tuned at a NRF energy of nuclide to detect
- Detection of scattered γ -rays by energy-resolved detectors
- NRF is a unique fingerprint of nuclides \rightarrow radioactive and stable nuclides can be detected
- Using MeV γ -rays \rightarrow applicable to thick objects

Nondestructive Detection of Isotopes

tuning the γ -ray energy to the NRF
 ↓
 separation of the NRF signal from the background



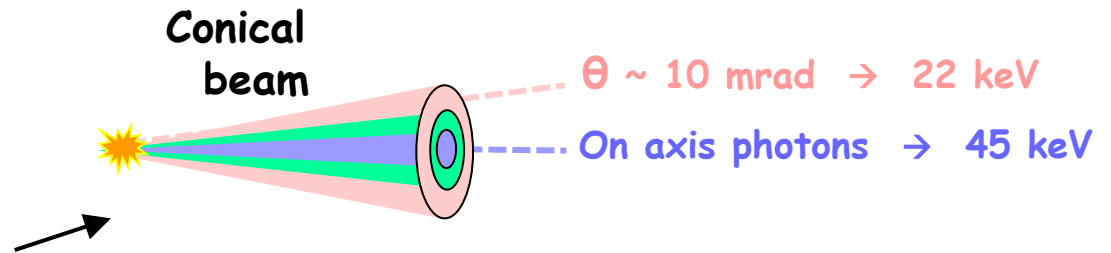
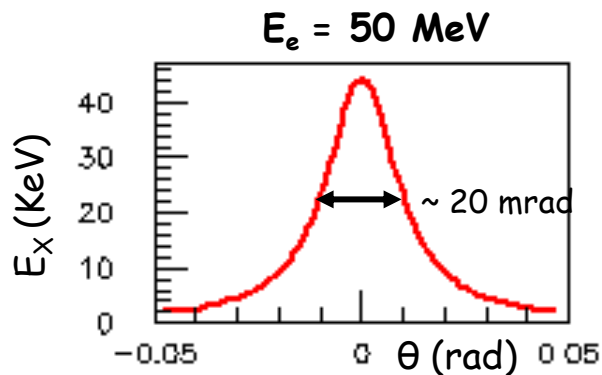
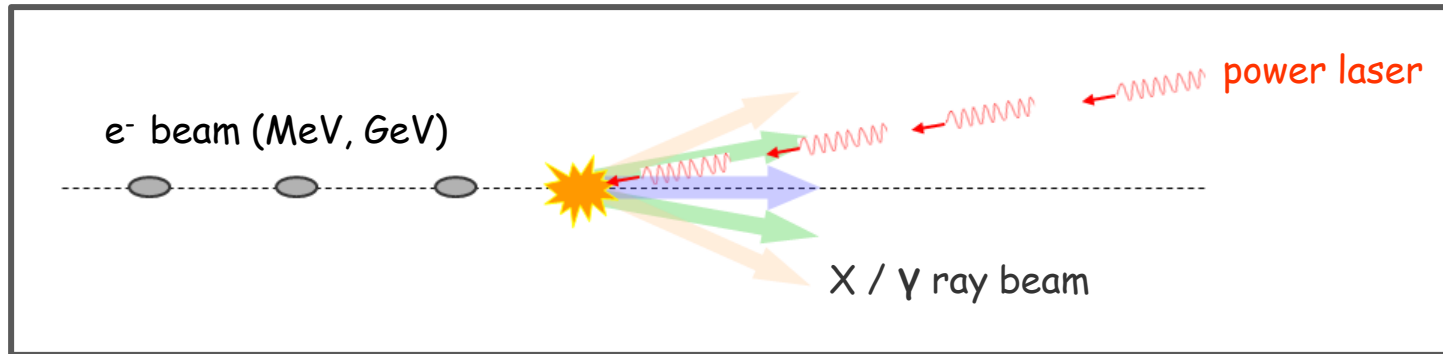
$\sigma_{\text{NRF}} = 28 \text{ mbarn} \cdot \text{keV} @ 2.176 \text{ MeV}$

Simulation by GEANT4 (with NRF routine)

Compact Compton Sources (CCS) → How ?

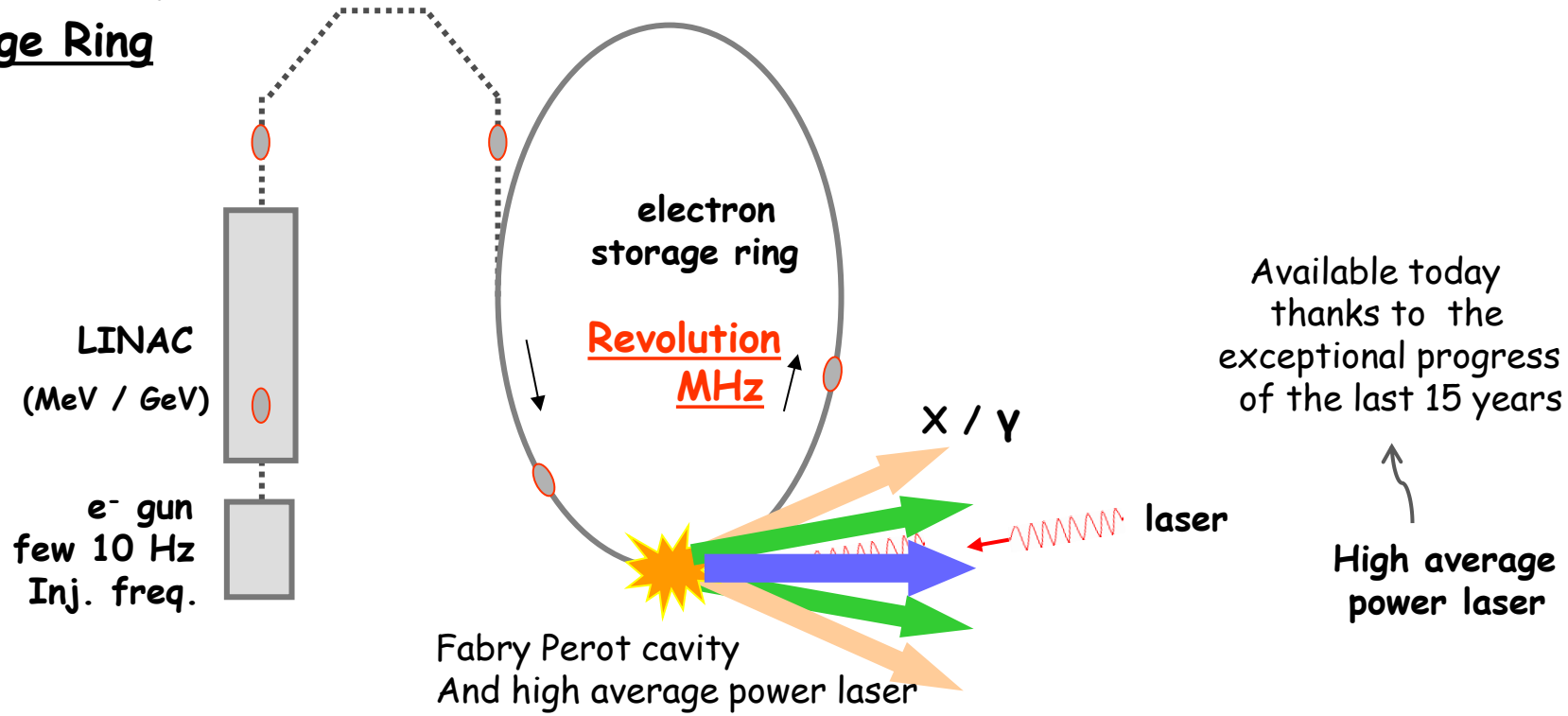
→ Principle

→ Characteristics / Specifications

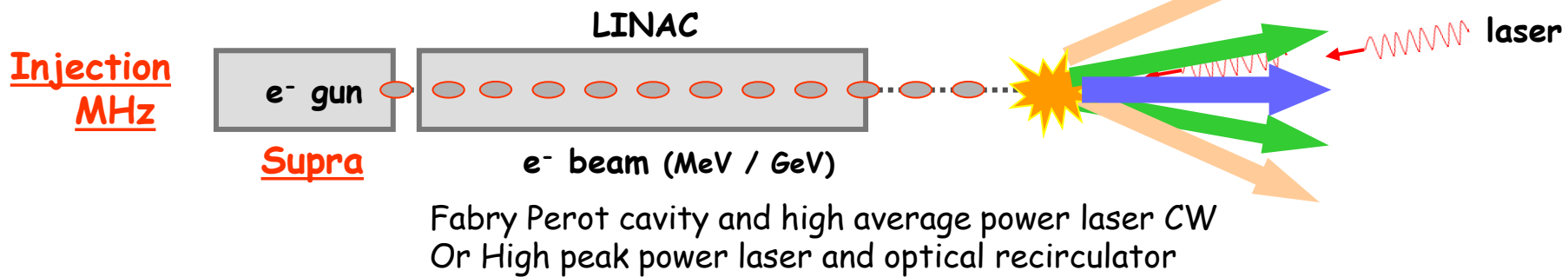


High flux ($10^{12} - 10^{14}$ ph/sec) \rightarrow
Increase f_{rep} e^- / laser ($\sim 10-100$ MHz) \rightarrow 2 main schemes

Storage Ring



Linac (or ERL)



- Source Parameters and Characteristics

- **Compactness** (surface $\sim 100 \text{ m}^2$) \rightarrow Integration in an hospital, a lab, a museum

- X-ray energy is tunable (LINAC is better)

$$E_x \sim E_e^2 E_{ph} \rightarrow E_x \text{ can "easily" be tuned by changing the electron beam energy}$$

- X-ray polarisation is controlled

The polarisation of the X-ray beam is governed by the laser polarisation.

- Relative High intensity beam

$$\text{CCS flux} \sim 10^{12} - 10^{14} \text{ ph / sec}$$

- High quality beam

- Transverse coherence
- High brightness

$$\text{Flux} \sim \frac{\sigma_{\text{compt}} N_e N_\gamma f_{\text{rep}}}{2\pi (\sigma_e^2 + \sigma_\gamma^2)} \quad (= \text{nb ph/sec})$$

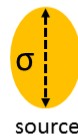
High flux \rightarrow lots of photons
lots of electrons
small beam sizes
large f_{rep}

Quality of the X-ray beam \approx Quality of the e- beam

• High quality beam

L : distance source-sample

σ : transverse source size



L



d_{tc}

$$d_{tc} \sim \frac{L}{\sigma E_x (\text{KeV})} (10^{-9})$$

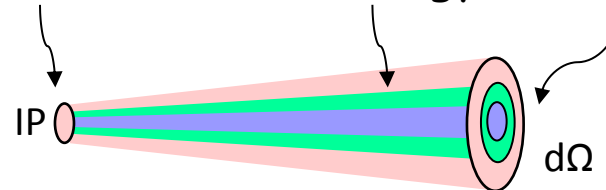
A large transverse coherence requires a small source size \rightarrow a small e- beam size

- Transverse coherence

- High brightness

Flux

$$Br_x = \frac{\text{Flux}}{(\text{mm}^2 \text{ source}) (\text{energy bin}) (\text{mrad})^2}$$



CCS $Br_x \sim 10^{11} - 10^{15}$ ph / sec / mm^2 / 0.1% bw / mrad^2

Brightness of a high flux CCS

$$Br_x \sim \frac{\text{Flux}}{(\text{mm}^2 \text{ source}) (dE_x/E_x) (\text{mrad})^2}$$

X-ray spectral bandwidth

laser energy spread
 $dE_x/E_x \sim dE_{ph}/E_{ph}$

laser divergence
 $dE_x/E_x \sim \frac{1}{4} (d\theta_{ph})^2$

e- energy spread
 $dE_x/E_x \sim 2 (d\gamma/\gamma)$

$dE_x/E_x \sim 10^{-3} \otimes 10^{-5} \otimes 10^{-3} \otimes 10^{-3} - \text{few } 10^{-2}$

(dominant effect)

e- beam divergence $\rightarrow dE_x/E_x \sim (\gamma \sigma_e)^2$

Typical e- divergence values:

$\sigma_e' \sim 0.5 \text{ mrad} \rightarrow dE_x/E_x \sim 0.25 \%$

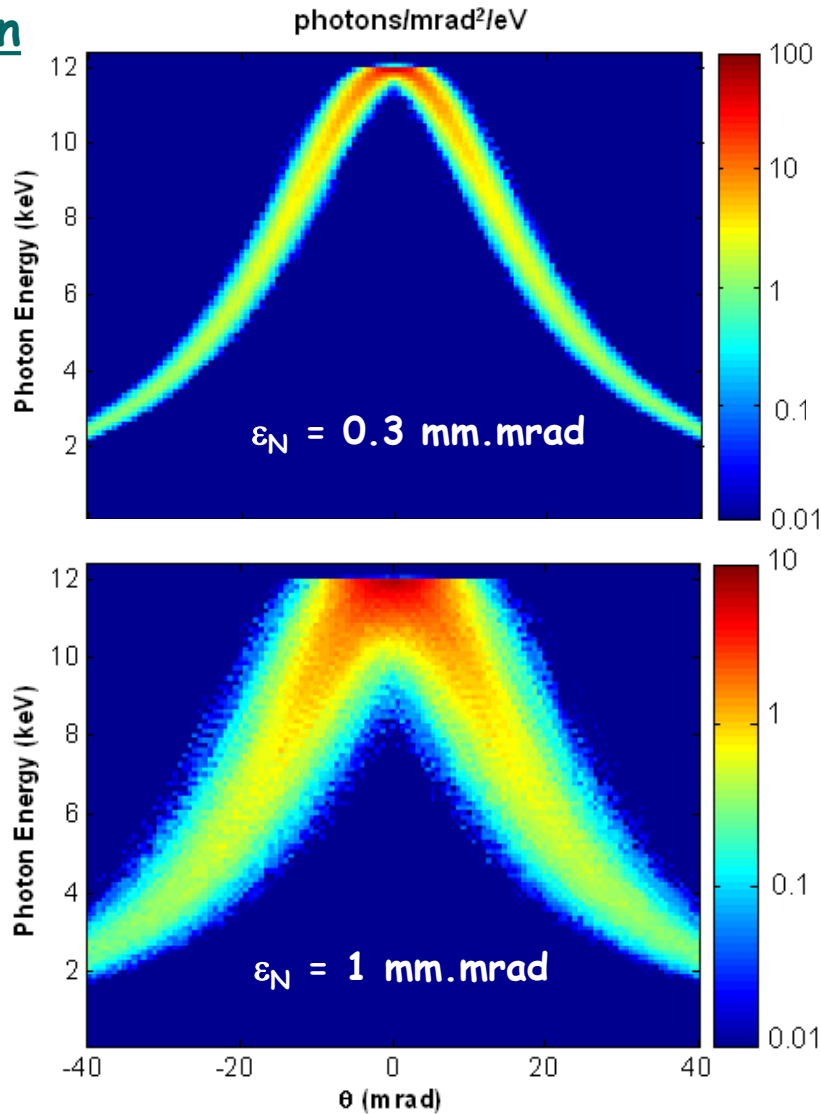
$\sigma_e' \sim 2 \text{ mrad} \rightarrow dE_x/E_x \sim 4 \%$

$$Br_x \sim \frac{\text{Flux}}{(\text{mm}^2 \text{ source}) (dE_x/E_x) (\text{mrad})^2} \sim \frac{\text{Flux} \cdot \gamma^2}{\epsilon_N^2}$$

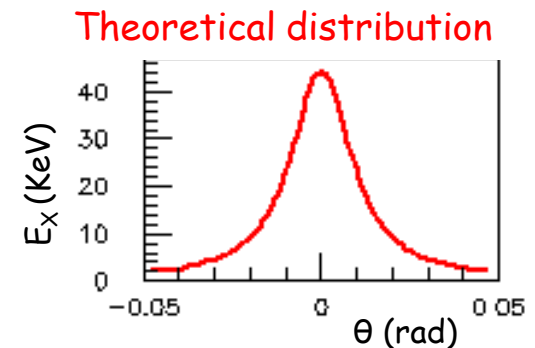
$\sigma_e^2 \sigma_e'^2 = \epsilon^2$

→ High brightness = small e- beam emittance

→ ... Illustration



[W.S. Graves, Alghero Workshop 2008]



→ The e^- beam emittance is a crucial parameter for the brightness

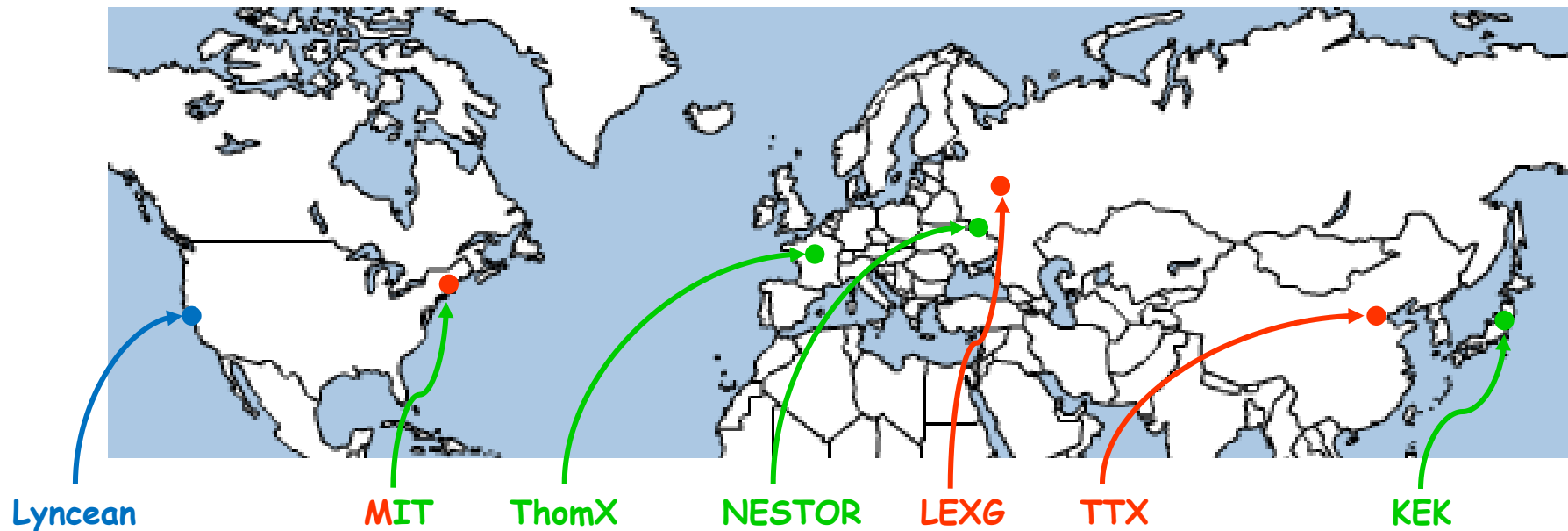
High flux X-ray CCS projects

→ Illustration through projects

Compact Compton projects (X-ray flux > 10^{12} ph/sec)

Project	Type	f_{rep} (MHz)	E_X (KeV)	ph/sec	σ_S (μm)	τ_S (ps)
Lyncean	SR	80	10-20	10^{11}	50	
TTX	SR	25	20-80	10^{12}	50	10
LEXG	SR (sc)	100	33	10^{13}	20	10
NESTOR	SR	20	30-500	10^{13}	70	50
ThomX	SR	20	20-90	10^{13}	70	10
KEK QB	Linac (sc)	160	35	10^{13}	10	
KEK ERL	Linac (ERL/sc)	130	67	10^{13}	30	3
KEK CSR	Linac (ERL/sc)	1300	0.04-4	10^{13-14}		0.1-1
MIT	Linac (sc)	100	3-30	10^{14}	2	0.1

- In operation
- Funded
- Project

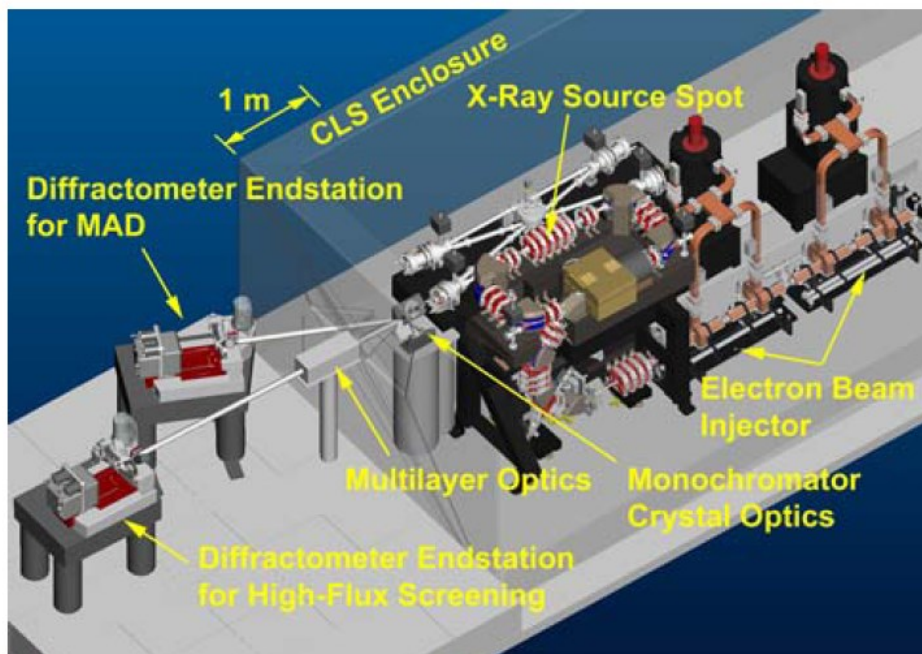


LYNCEANTECH : Diffraction : 3D structural analyses

3D structure of macro-molecules (proteins)

Knowledge of the structure of a protein → acced to its function in the cell

First determination of the 3D structure of a protein at CS Lyncean Tech.

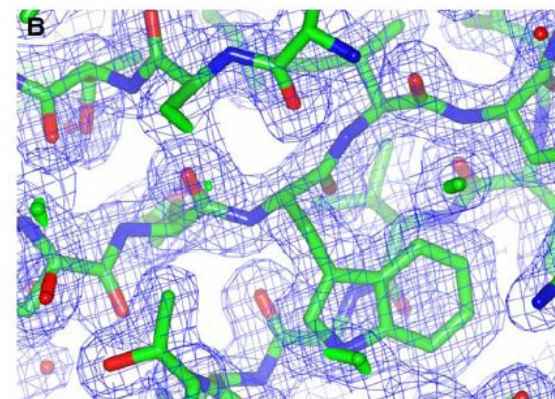
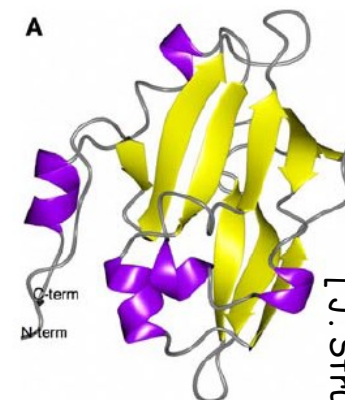


$5 \cdot 10^6$ ph/sec
X beam : 120 μ m on crystal

$E=15$ KeV
 $\Delta E/E = 1.4\%$

Protein MytuGCSPH (crystal size : 250 X 250 X 100 μ m)

Fig. 6 Structure and electron density of MytuGCSPH. **a** Overview of the structure for MytuGCSPH in ribbon representation. **b** Electron density from the MytuGCSPH CLS data set at 2.0 Å resolution centered around Trp 14



[J. Struct. Funct. Gen. 11, 2010, 91-100]

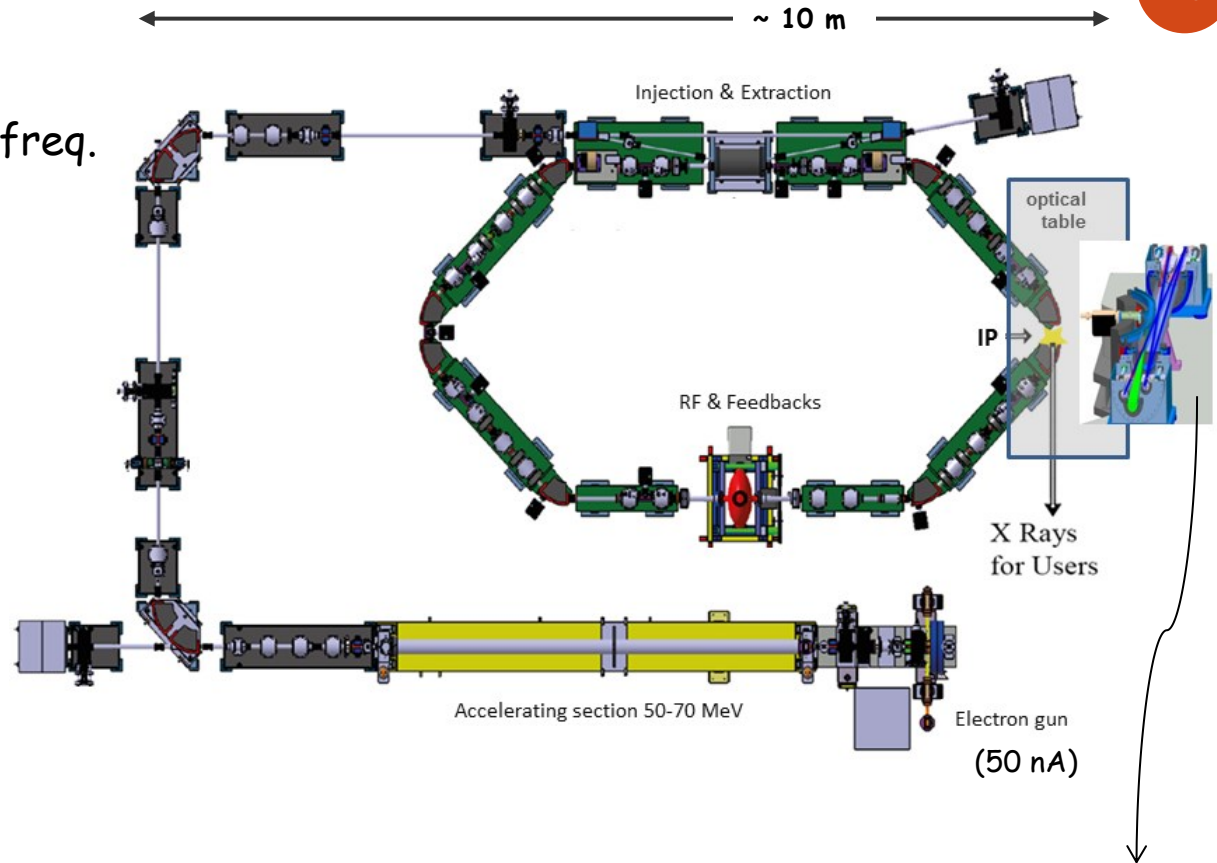
Storage ring : ThomX design

40

Electron machine

- 1 nc / bunch , 50 Hz inj. freq.
- Ring, 20 MHz freq.
- $\sigma_e \sim 70 \mu\text{m}$
- $\epsilon_N \sim 4 \text{ mm.mrad}$
- $\tau_e \sim 10\text{-}20 \text{ ps}$

Machine funded
In construction

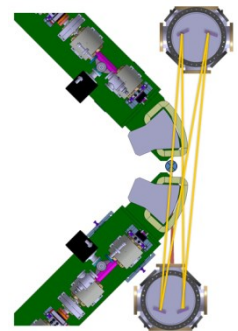


X-ray beam

Flux	10^{13}
Brightness	10^{11}
Size	$70 \mu\text{m}$

Laser /Cavity system

- Laser $\sim 1\text{W}$
- Optical fiber amplification (100 W) 2-3 $\mu\text{J/pulse}$
- Optical FP cavity amplification (gain 10000)
- 1 MW stored inside the cavity (20-30 mJ/pulse)



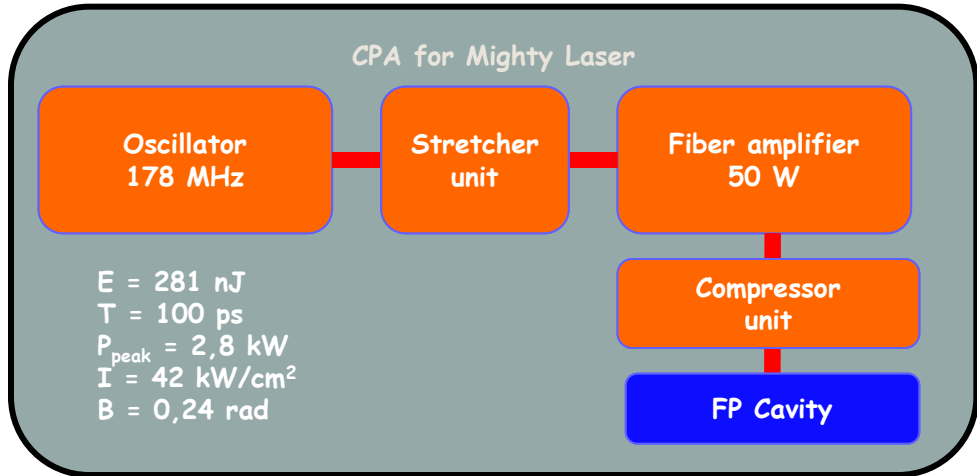
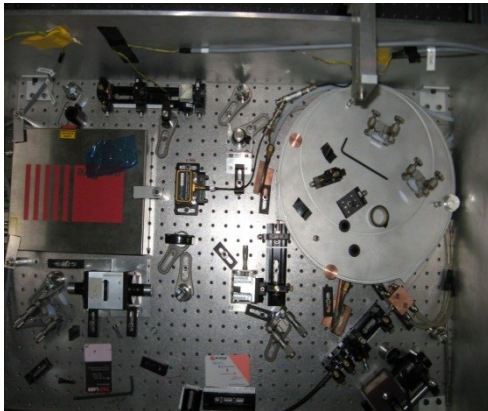
Expected beams characteristics

41

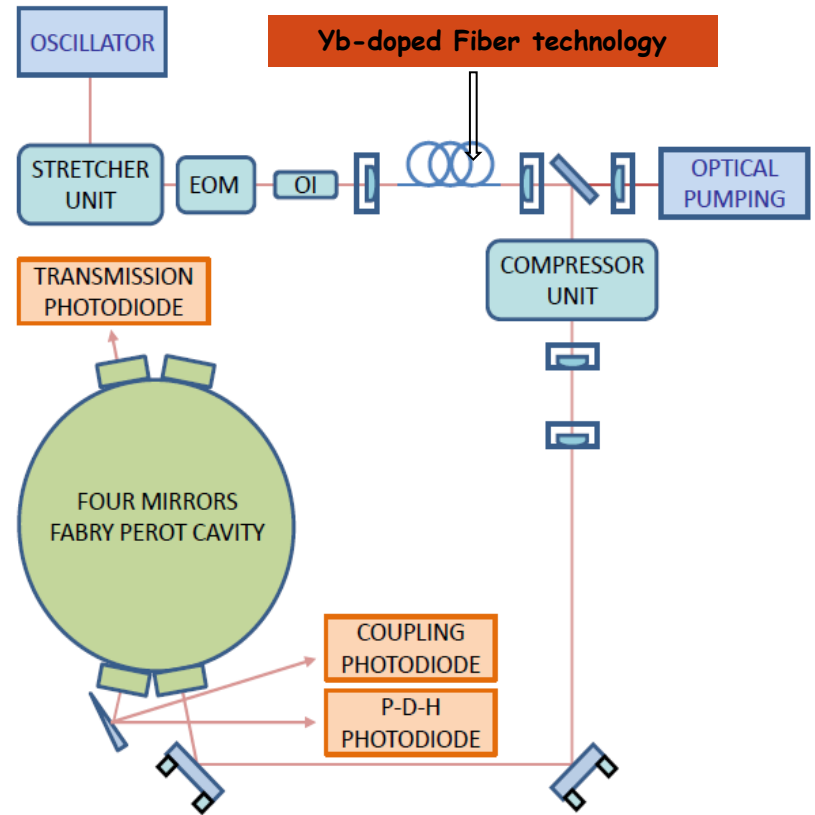
Injector		Ring	
Charge	1 nC	Energy	50 MeV (70 MeV possible)
Laser wavelength and pulse power	266 nm, 100 μ J	Circumference	16.8 m
Gun Q and Rs	14400, 49 MW/m	Crossing-Angle (full)	2 degrees
Gun accelerating gradient	100 MV/m @ 9.4 MW	$B_{x,y}$ @ IP	0.2 m
Normalized r.m.s emittance	8π mm mrad	Emittance x,y (without IBS and Compton)	$3 \cdot 10^{-8}$ m
Energy spread	0.36%	Bunch length (@ 20 ms)	30 ps
Bunch length	3.7 ps	Beam current	17.84 mA
Laser and FP cavity		RF frequency	500 MHz
Laser wavelength	1030 nm	Transverse / longitudinal damping time	1 s / 0.5 s
Laser and FP cavity Frep	36 MHz	RF Voltage	300 kV
Laser Power	50 - 100 W	Revolution frequency	17.8 MHz
FP cavity finesse / gain	30000 / 10000	σ_x @ IP (injection)	78 mm
FP waist	70 μ m	Tune x / y	3.4 / 1.74
		Momentum compaction factor α_c	0.013
		Final Energy spread	0.6 %
Source			
Photon energy cut off	46 keV (@50 MeV), 90 keV (@ 70 MeV)		
Total Flux	10^{11} - 10^{13} ph/sec		
Bandwidth (with diaphragm)	1 % - 10%		
Divergence	$1/\gamma \sim 10$ mrad without diaphragm @ 50 MeV		

Injector, ring, laser, Fabry-Perot resonator and the source

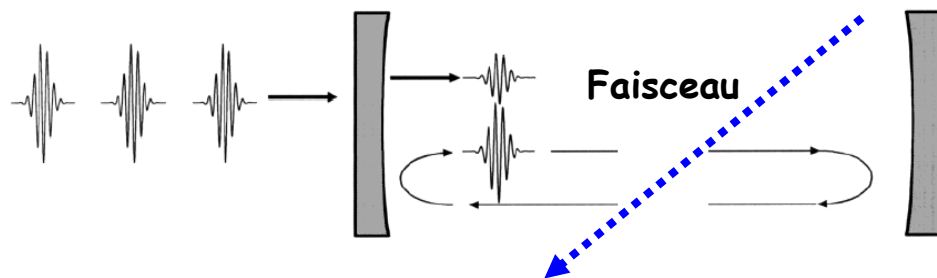
Laser



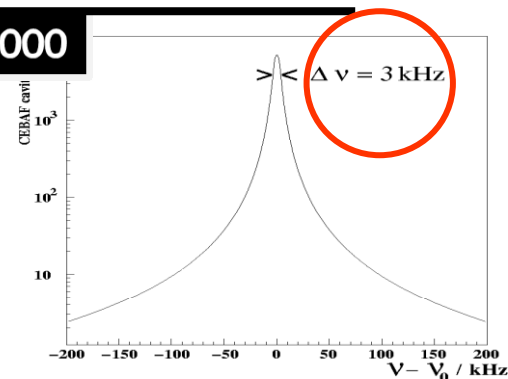
Compatible with low noise amplification



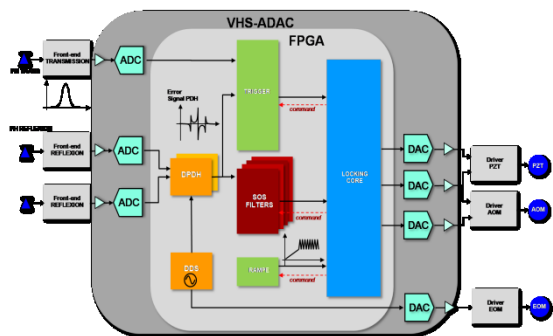
Fabry-Perot cavity



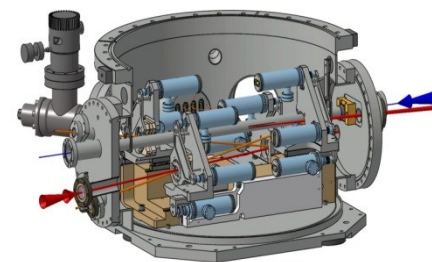
$Gain = 1/(1-R) \sim 10000$



Digital Pound-Drever-Hall feedback



Vacuum and mechanics : MightyLaser experience



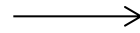
PLIC and MightyLaser : record in stable finesse locking (30000).

Beam Dynamics

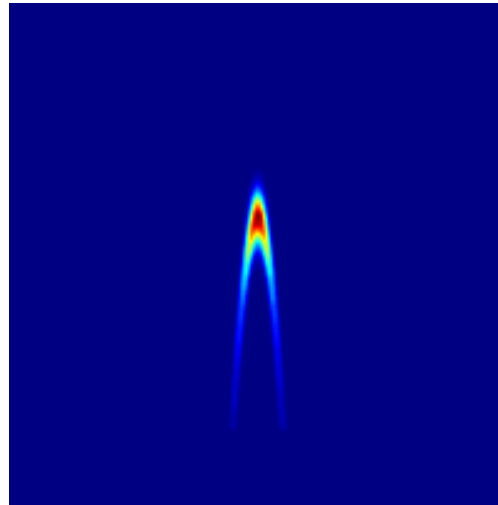
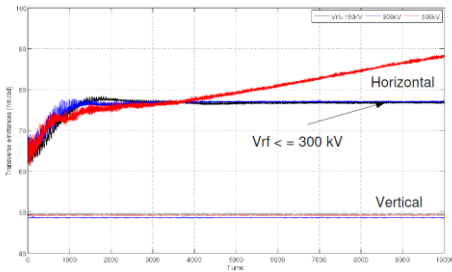
- Transient dynamics (no equilibrium)
- Compton recoil
- Collective instabilities
- CSR, Ions, Vacuum scattering, IBS.....
- Injection mismatching



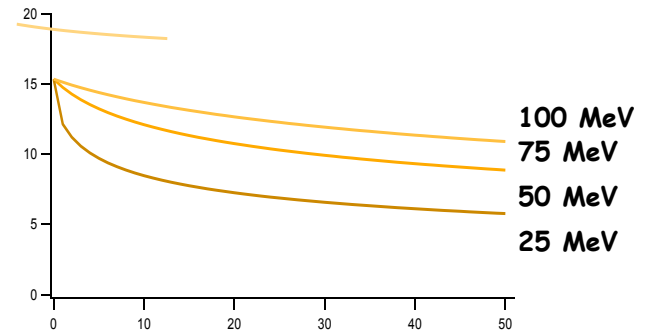
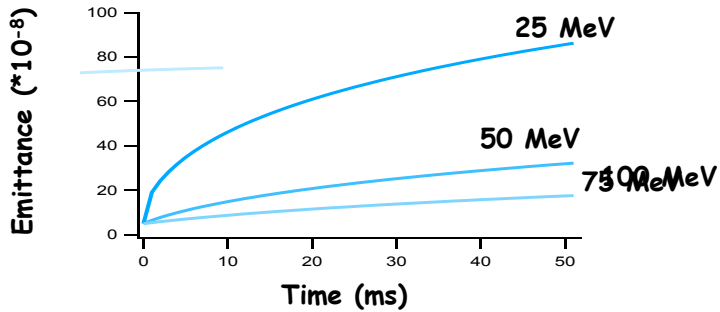
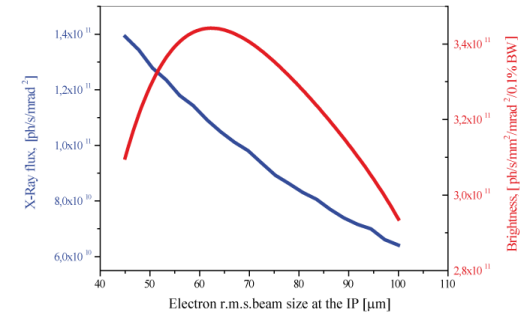
Feedbacks!!!
Ion clearing
Simulations



3 Phases
Injection
Turbulent regime
Stabilization (thousand turns)



Careful to the brilliance !!!



100 MeV
75 MeV
50 MeV
25 MeV

Electron machine

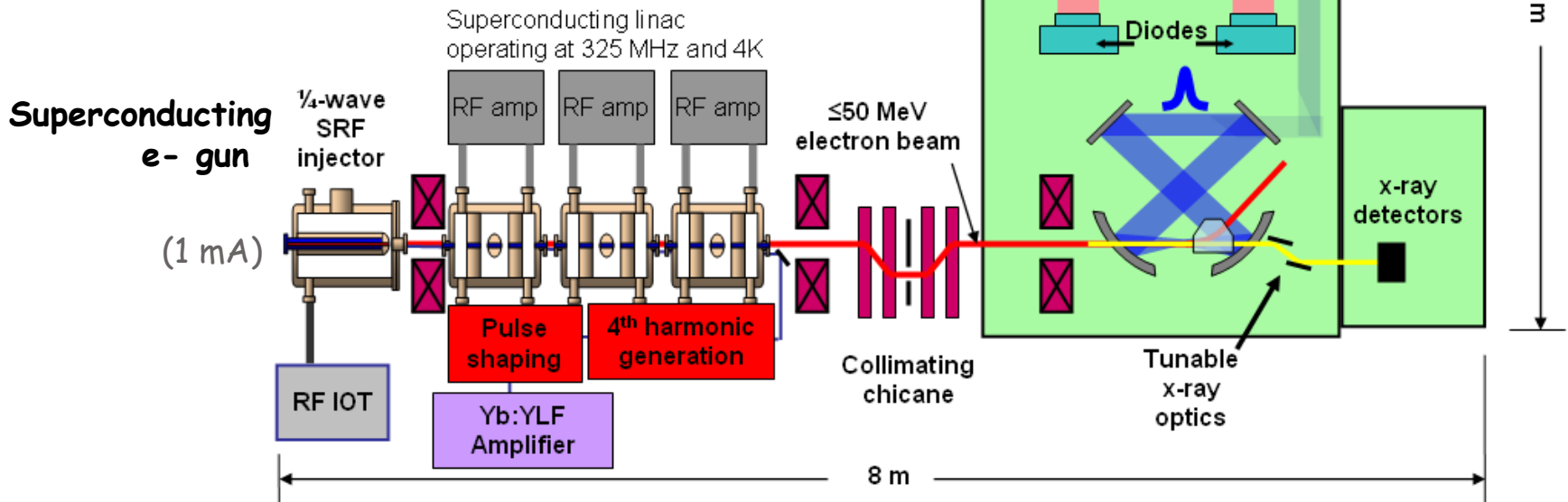
- 0.01 nc / bunch , 100 MHz inj. freq.
- $\sigma_e \sim 2 \mu\text{m}$
- $\epsilon_N \sim 0.1 \text{ mm.mrad}$
- $\tau_e \sim 0.1 \text{ ps}$

Laser /Cavity system

- Laser $\sim 10\text{W}$
- Cryogenic fiber amplification (1 KW) $10 \mu\text{J/pulse}$
- Optical cavity amplification (gain 1000)
- 1 MW stored inside the cavity (10 mJ/pulse)

X-ray beam

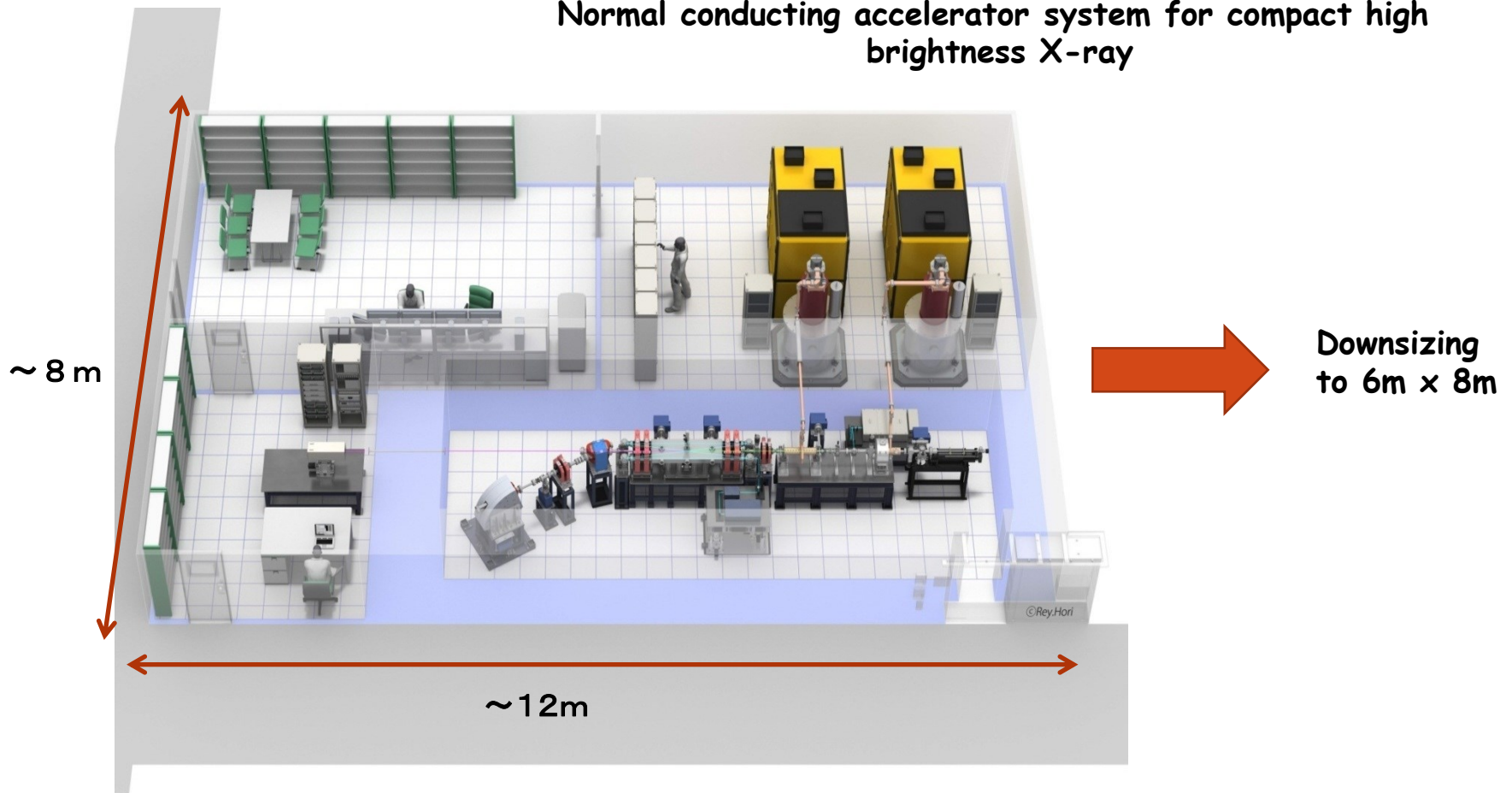
Flux	10^{14}
Brightness	10^{15}
Size	$2 \mu\text{m}$



New Quantum Beam Technology Program(NQBTP) is supported by MEXT from 2013.8 to 2018.3 (~5 years project).

Approved project included two Japanese Companies at least and the development for CW super conducting acceleration technologies. Normal conducting accelerator system and super conducting accelerator system for compact high brightness X-ray source should be realized by joint research with companies.

Normal conducting accelerator system for compact high brightness X-ray

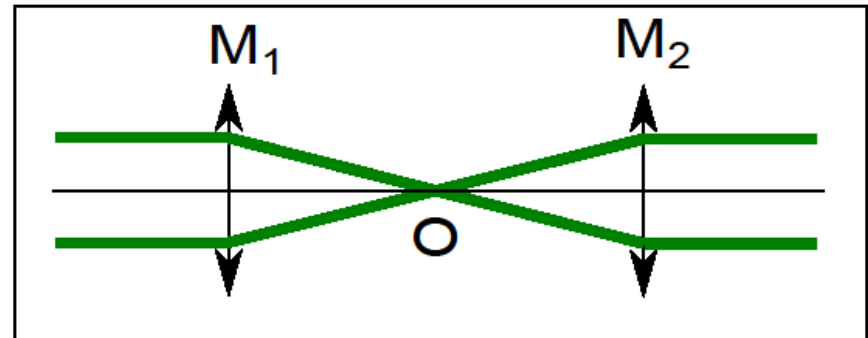
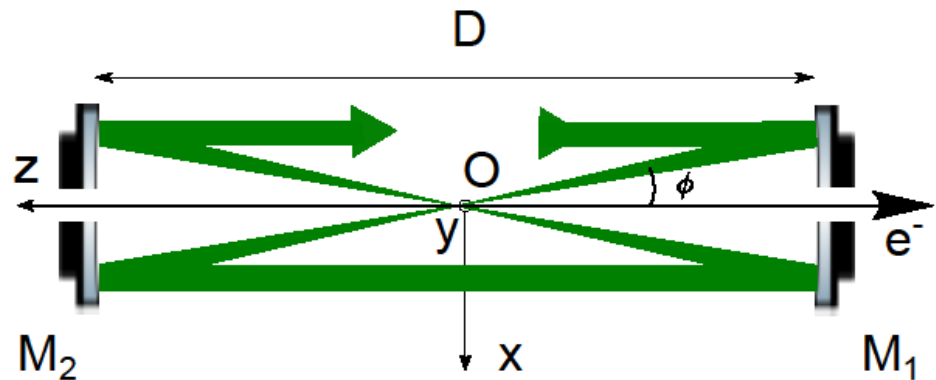


ELI-NP gamma source specifications

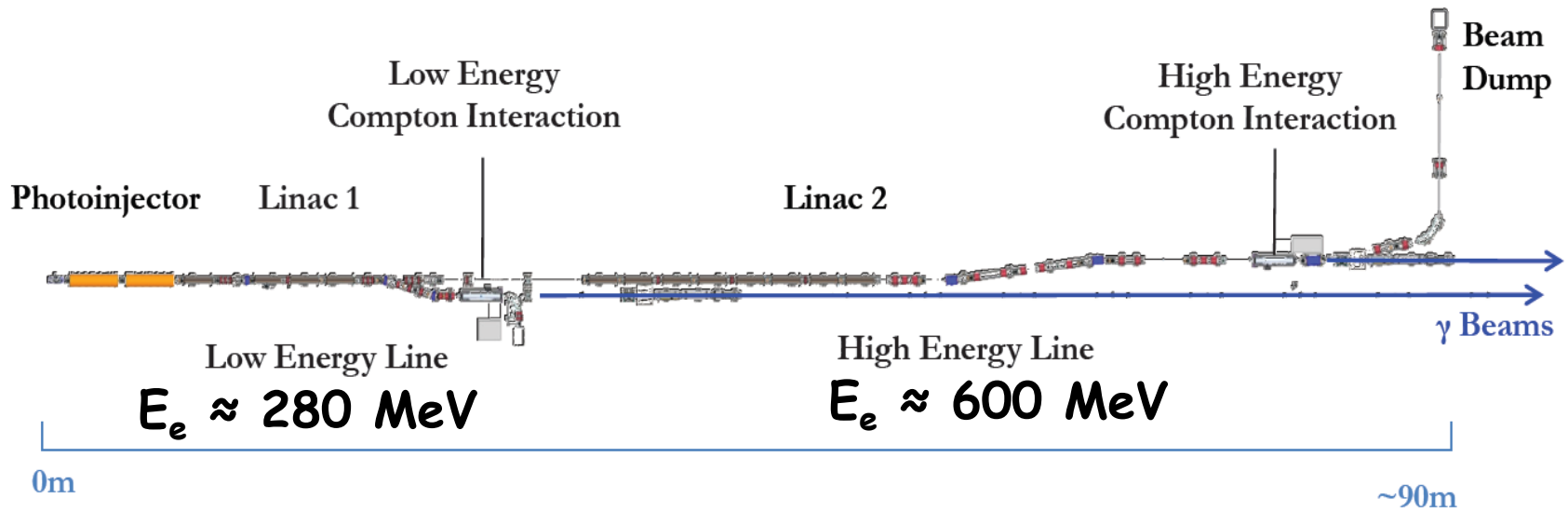
A future application to X rays?

- Energies γ (E_γ) : 0.2 - 19.5 MeV
- Bandwidth ($\Delta E/E$) : 0.5%
- Spectral density (flux) : 5000 $\gamma/(s.eV)$
- Linear polarization : 95%

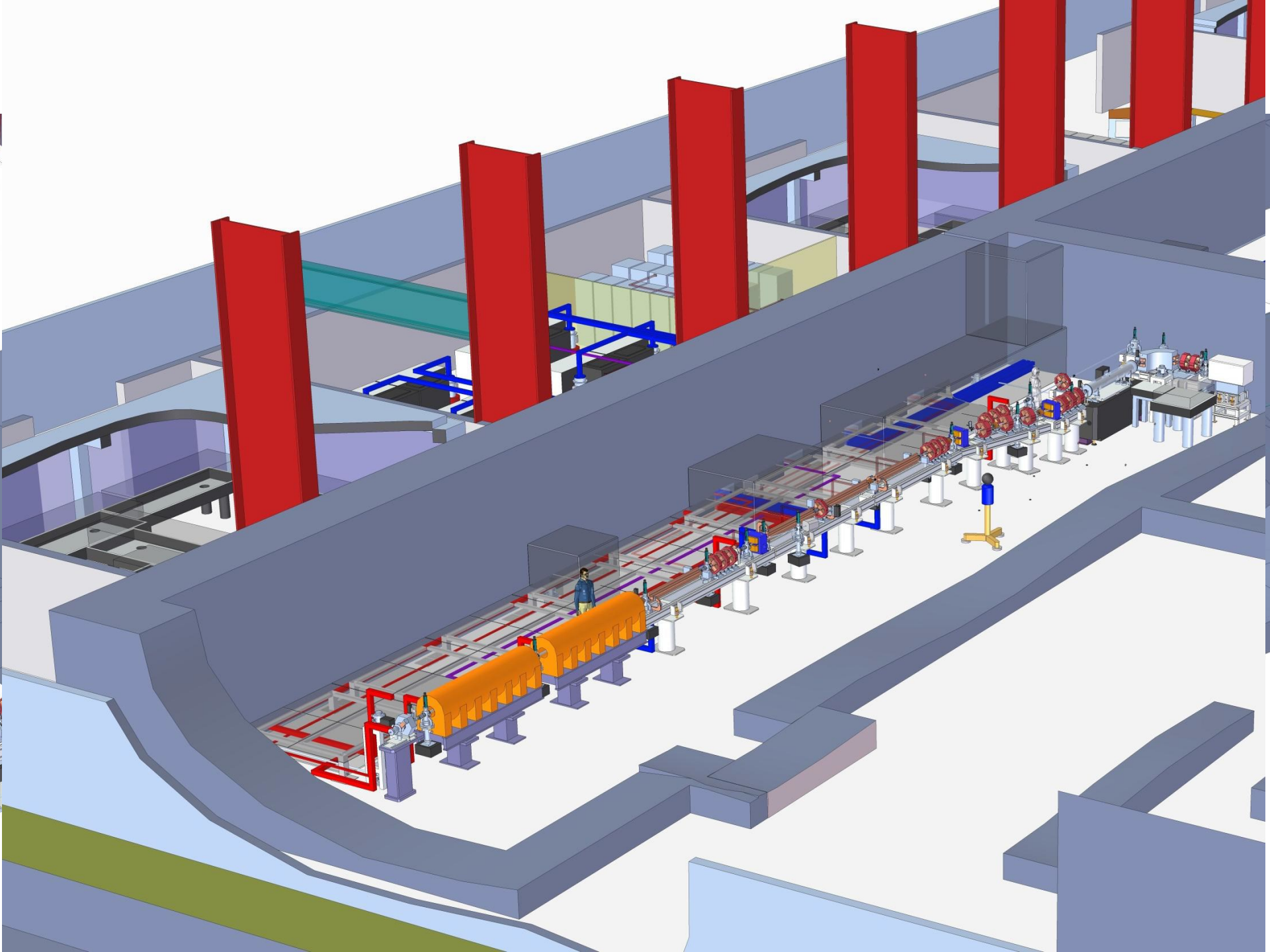
⇒ LINAC multi-bunch 100Hz
+
optical circulator



Source layout



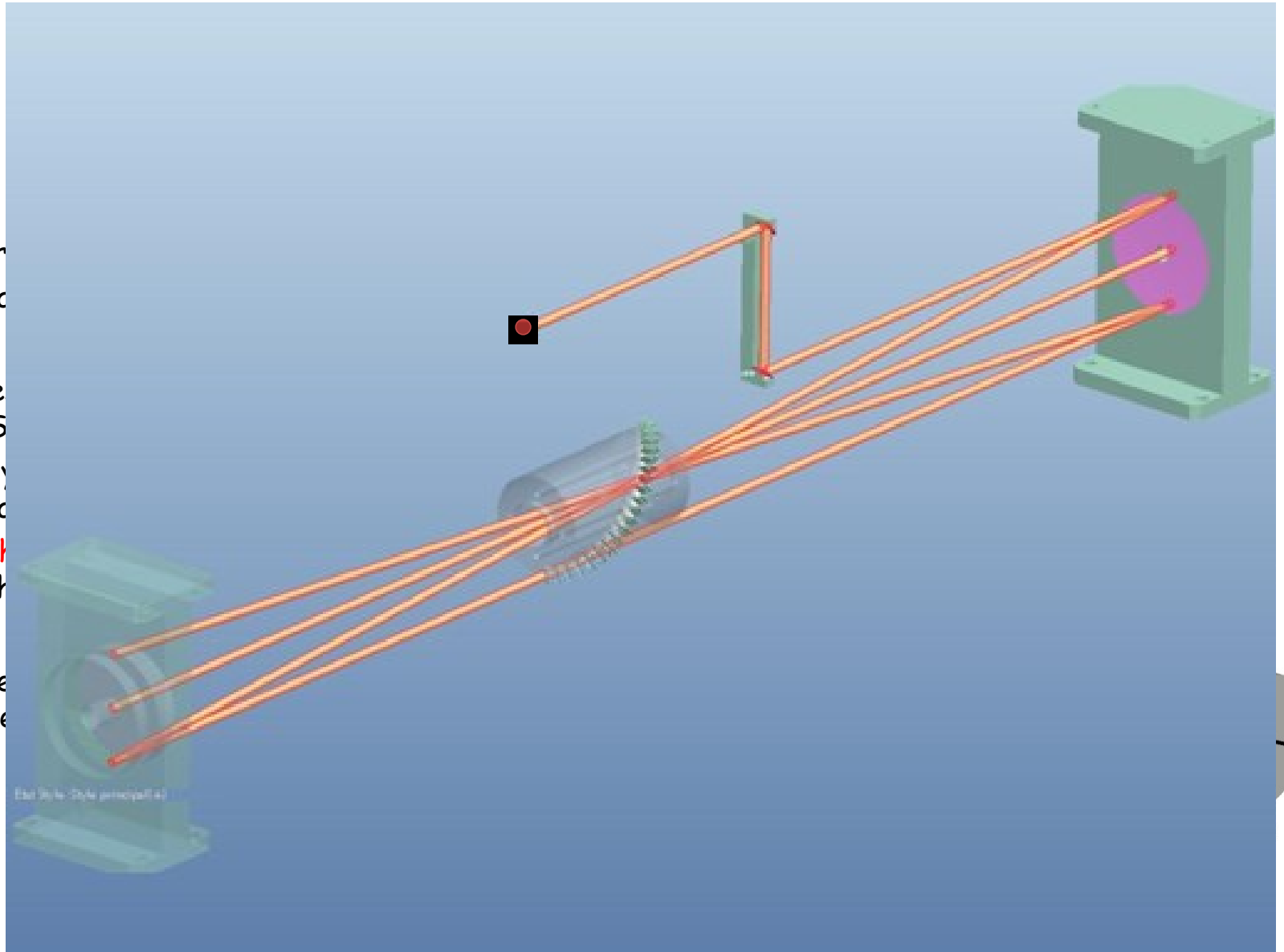
- **2 points d'interactions :**
 - 1 lasers 200mJ Yb@515nm (3.5ps) par point d'interaction
 - LINAC hybride bandes S et C (~100 - 720 MeV)



Système optique : circulateur

- 2 mir
- Pa
- Paire (MPS
- Sy
- pa
- Ch
- ch

⇒ Angle faible



Etat 3D - Style principal #1

Outlook

- Compton effect has various advantages as far as compact sources are concerned. ENERGY BOOST
- Complementarity to SR
- New generation of 'compact' X rays and nuclear gamma sources
- Different fields of application. Each requires a custom design and choice of accelerator and laser systems. High average power or high peak brilliance
- Worldwide increasing interest. Many projects ongoing
- Fascinating projects, they join the top of th accelerator and laser science with the laser-beam interaction physics

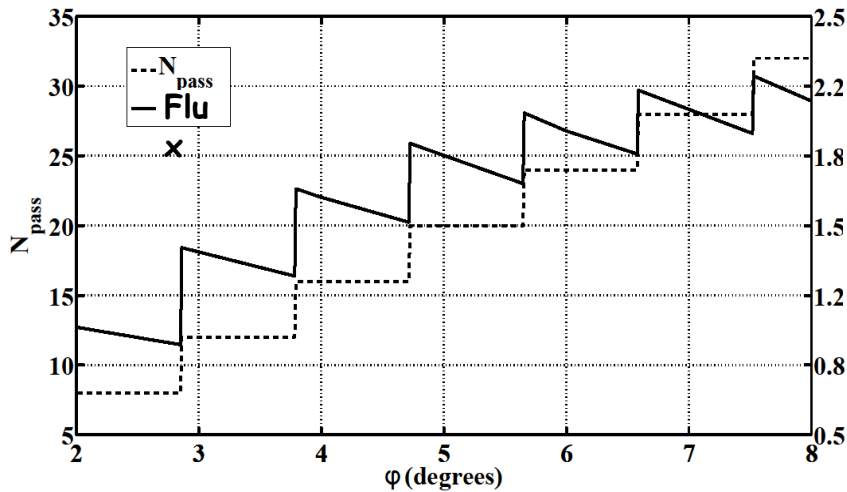
Thank you

Design

Design du circulateur

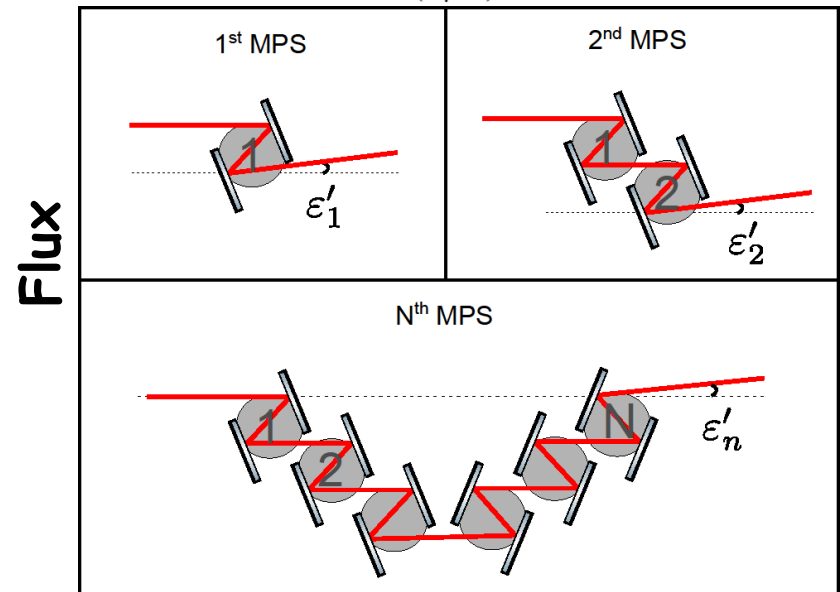
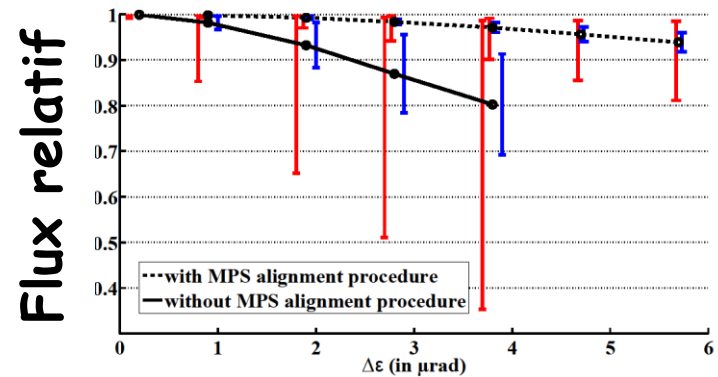
- $\omega_0 \uparrow \Rightarrow \text{flux} \downarrow$
- $\phi \uparrow \Rightarrow$
 - flux \downarrow
 - Nb passages $\uparrow \Rightarrow \text{flux} \uparrow$

\Rightarrow Optimisation



$\Rightarrow 32$ passages, $\phi = 7.54^\circ$

Parallélisme des paires de miroirs à face parallèles



Potential of applications of X-ray CCS: Examples

2. Using the central part of the beam (cultural heritage / material science applications)

- Fluorescence Spectroscopy
- XANES Spectroscopy

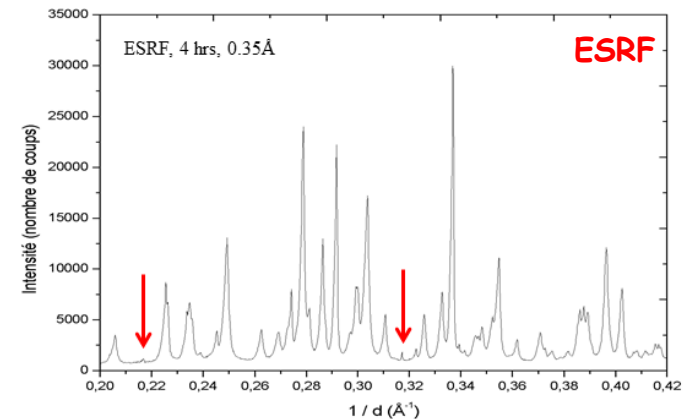
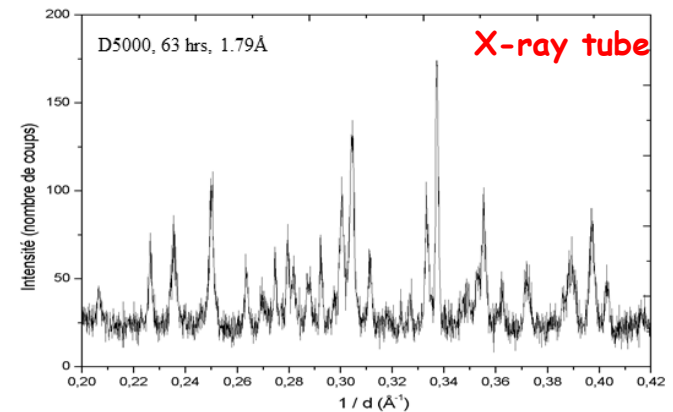
- Diffraction

→ Structural analyses

→ Pump-probe experiments

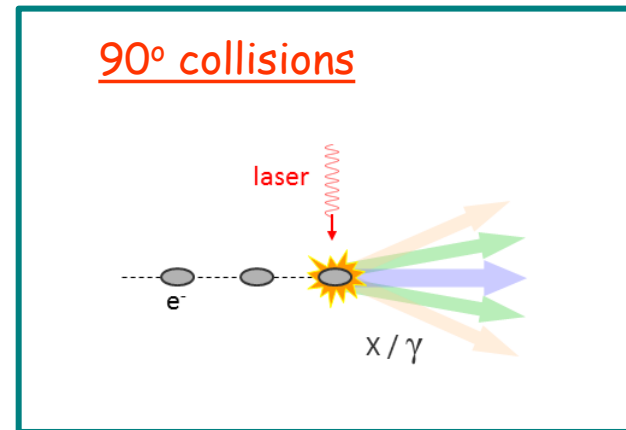
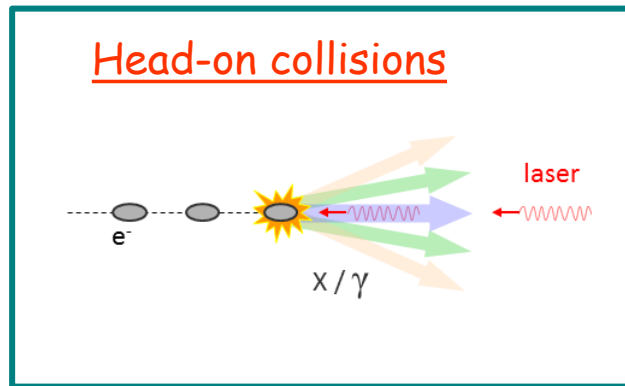
- Quasi-monochromatic beam
- Tunable energy

Diffraction patterns



→ ESRF : minor phases present in the powder visible

Pulse duration

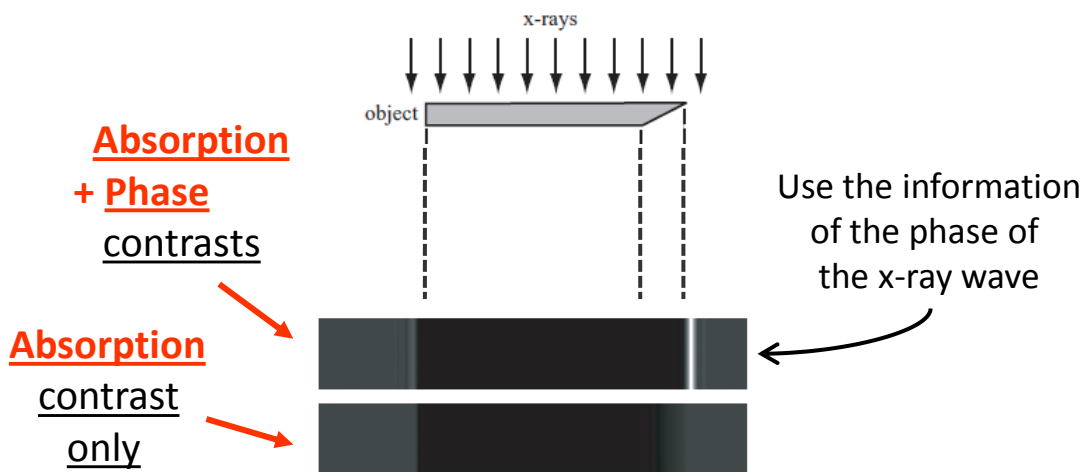


→ **0.1 - 10 ps**

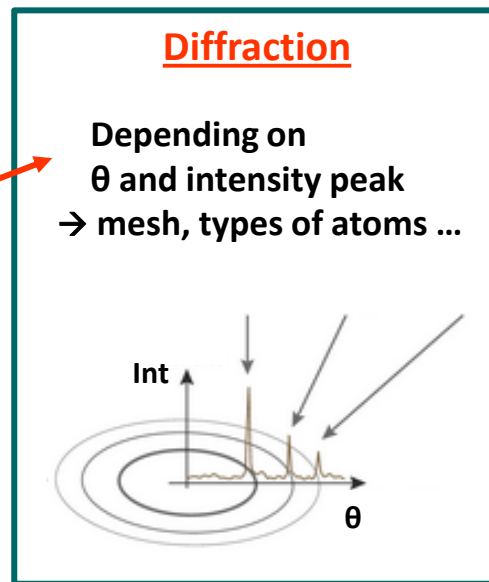
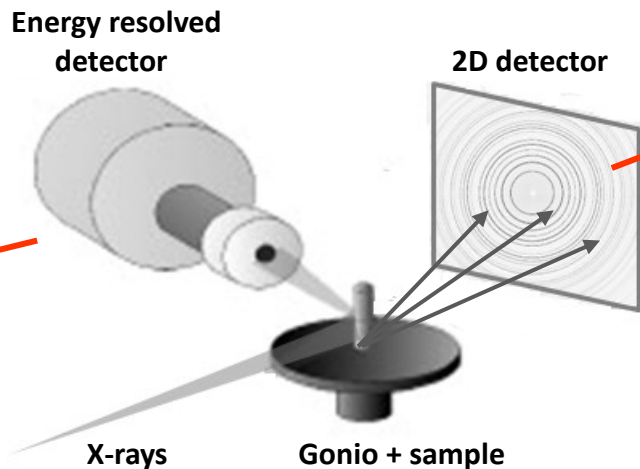
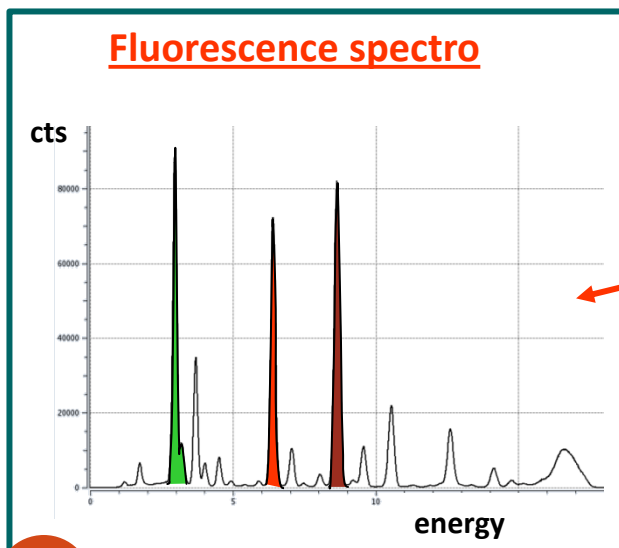
(depending on the e^- and laser pulses length, and the collision angle)

Imaging

→ Study of **absorption contrast and phase contrasts** of the **transmitted photons**



Diffraction / Spectroscopy → Study of the reemitted photon spectra



Interactions wave/matter described by the refraction index

$$N = n_R + i n_A$$

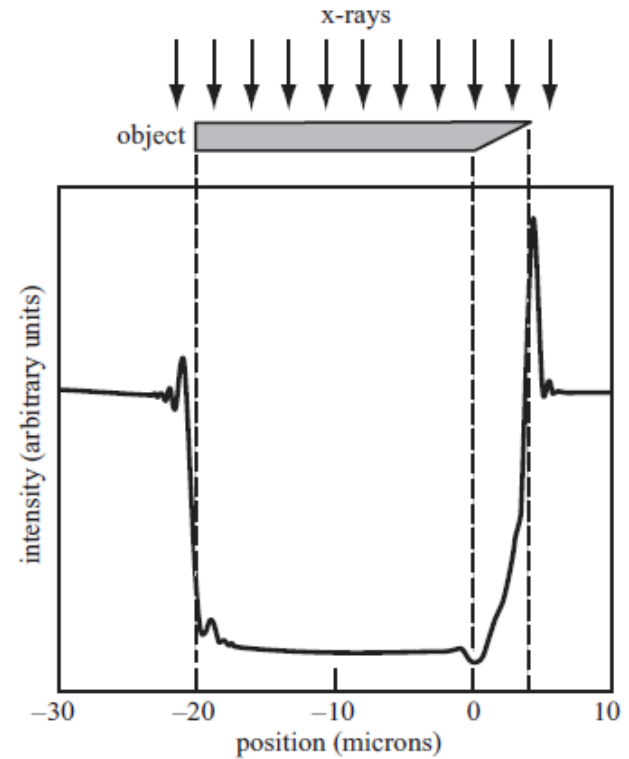
Phenomena of
→ refraction
→ diffraction
→ interferences

Absorption Contrast

(effects related to the phase of the wave)

Wave superposition
→ classical phenomena
refraction/diffraction
(Fresnel integrals)

→ **Phase Contrast**

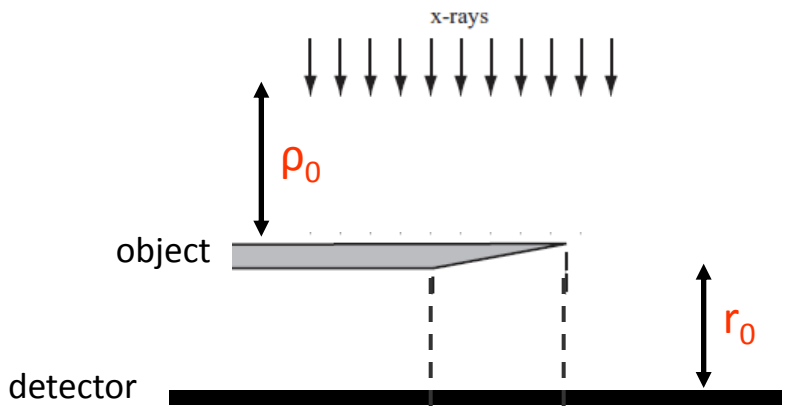


Absorption + Refraction

Absorption only

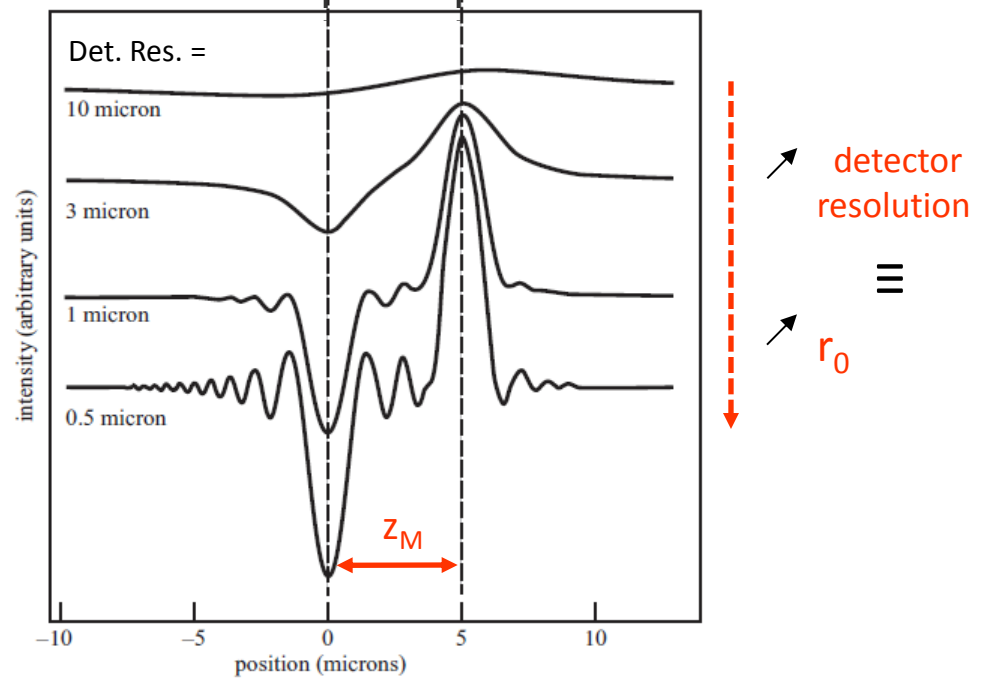


source : $\lambda, \Delta\lambda, \sigma$



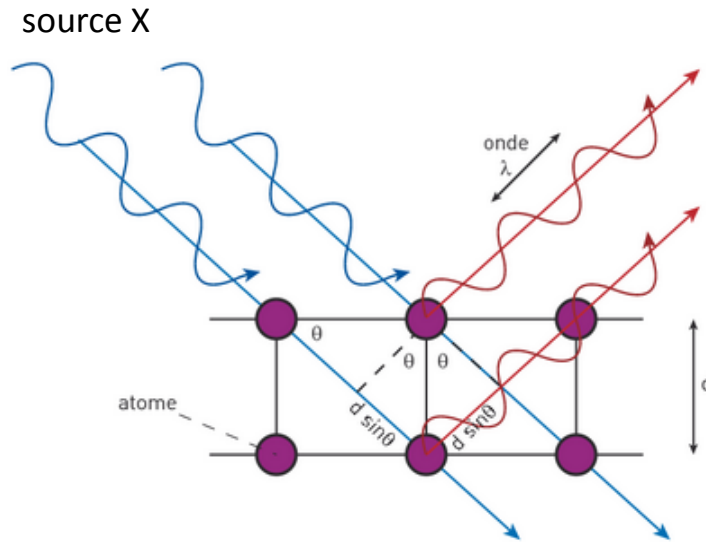
Conditions to detect the main maximum (frange)

- $\Delta z_M / z_M \sim \Delta\lambda / \lambda \ll 1$
- $(r_0 / \rho_0) \sigma \ll z_M$



Diffraction

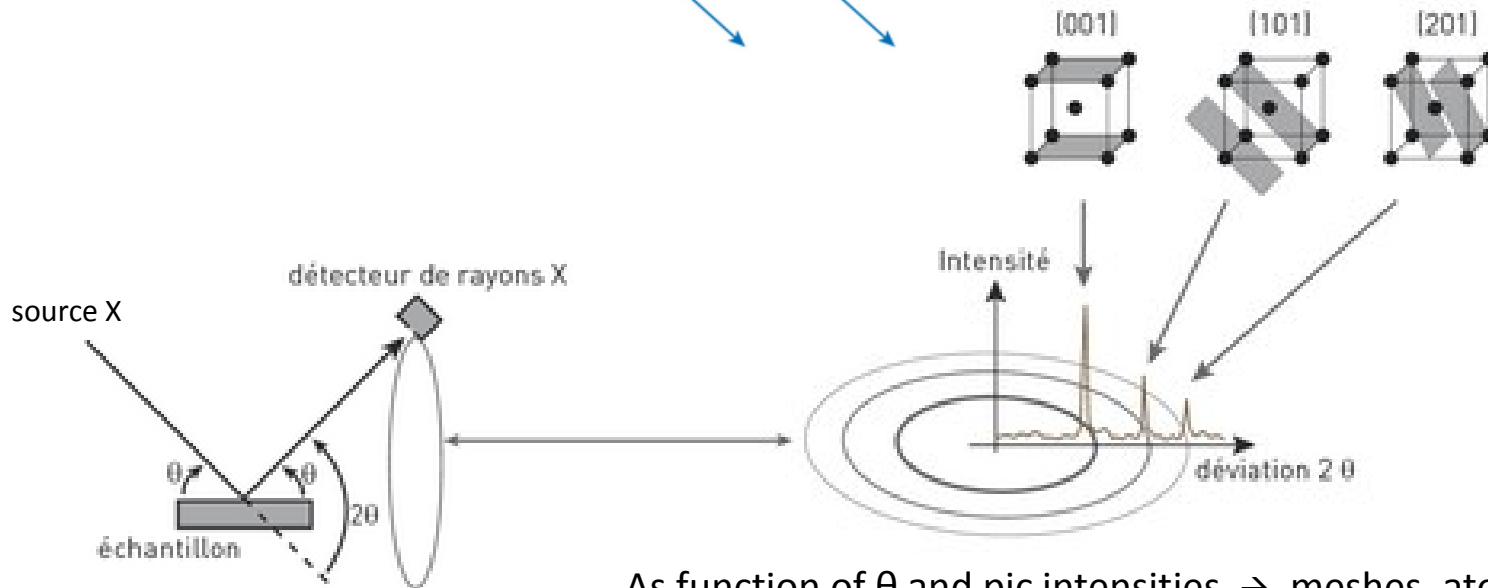
Diffraction



Crystallised material, powders ...

→ Diffused waves

→ Interferences

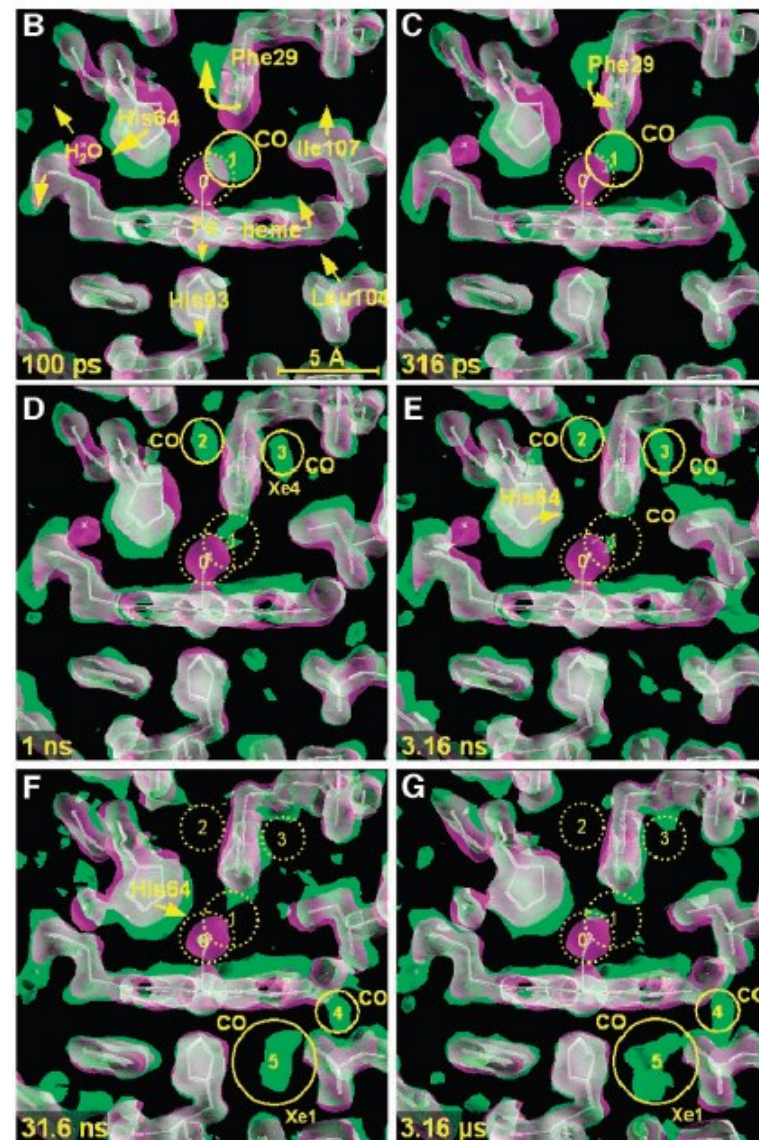


As function of θ and pic intensities → meshes, atoms types ...

- Visualisation in real time of atomic/molecular processes
- “Pump-probe” experiments
Currently: resolution ~ 100 ps)

“Live” view of the interaction
between a protein and a ligand
(ESRF, Resolutions 2Å, 100 ps)

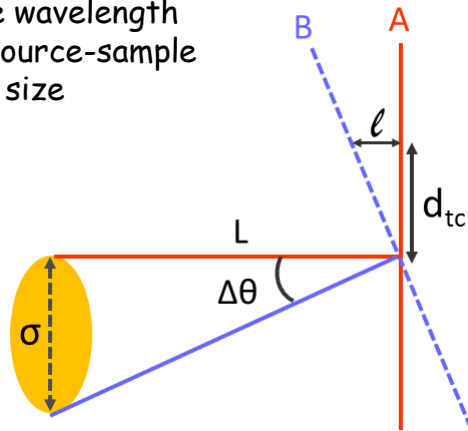
- Time scale of molecular movements ~ 10 -100 fs
- **Compact ICS with ultra-short pulses**



[Science 300, 1944, 2003]

Transverse coherence

λ : source wavelength
 L : dist. source-sample
 σ : source size



The wave fronts A et B arrive in phase if $l \sim \lambda$

$$\rightarrow d_{tc} \sim \frac{L}{\sigma E_x (\text{KeV})} (10^{-9})$$

A large transverse coherence
 requires a **small source size**
 → a **small e⁻ beam size**