

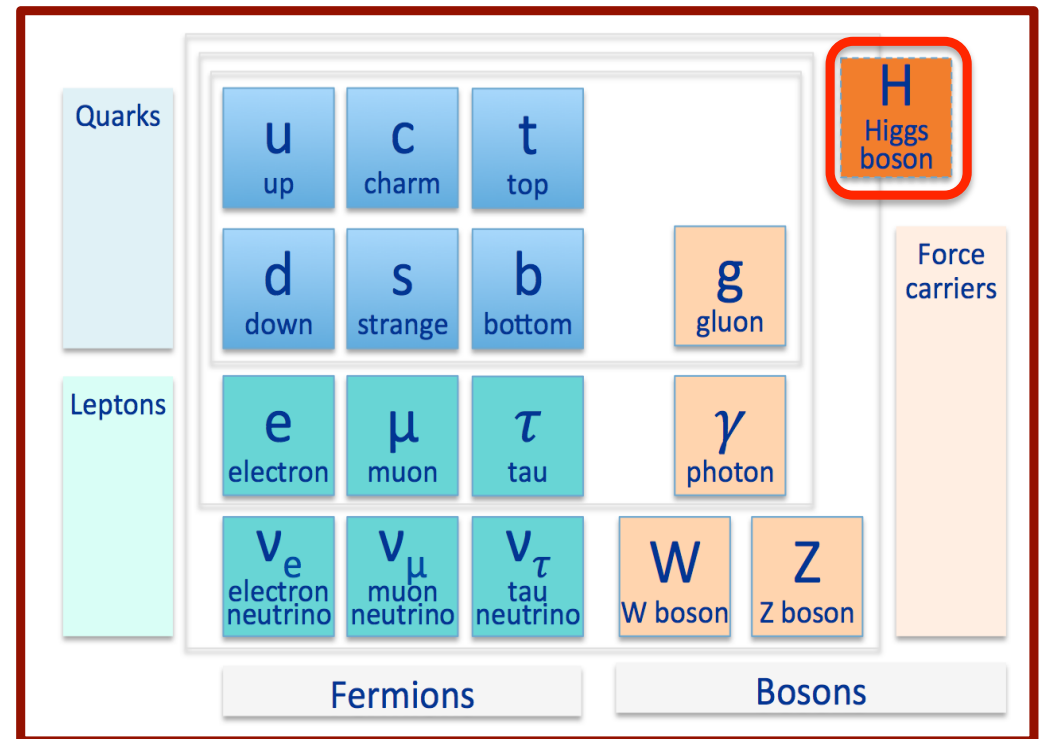
Beyond the (already observed) Higgs

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Zimányi School, Budapest, 6 December 2013

Status quo: the Standard Model after LHC Run1

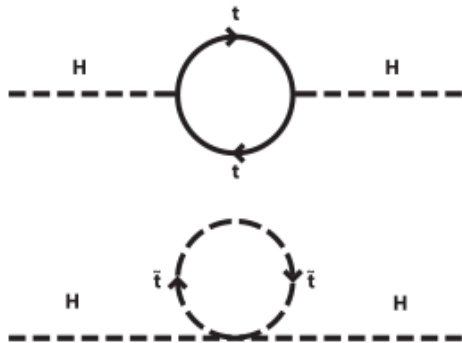
- ATLAS and CMS discovered a Higgs boson with $m = 125.5$ GeV, so far compatible with SM
- No other new particle seen at LHC
- SM complete and describes with huge success phenomena observed in collider experiments
- **SM can not be the final theory!**
- Deviations due to new physics (in the Higgs sector and elsewhere) might be subtle
- Precision measurements might give the key to discoveries
- ... parallel to direct searches for new physics!



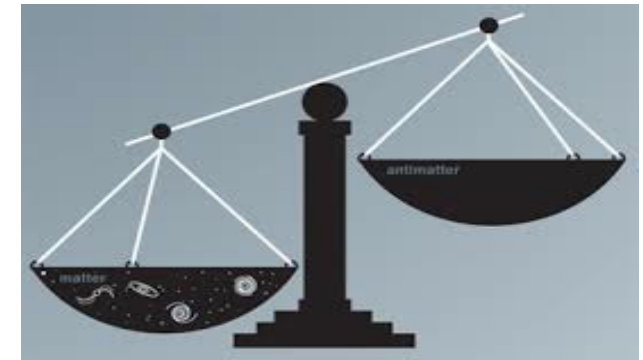
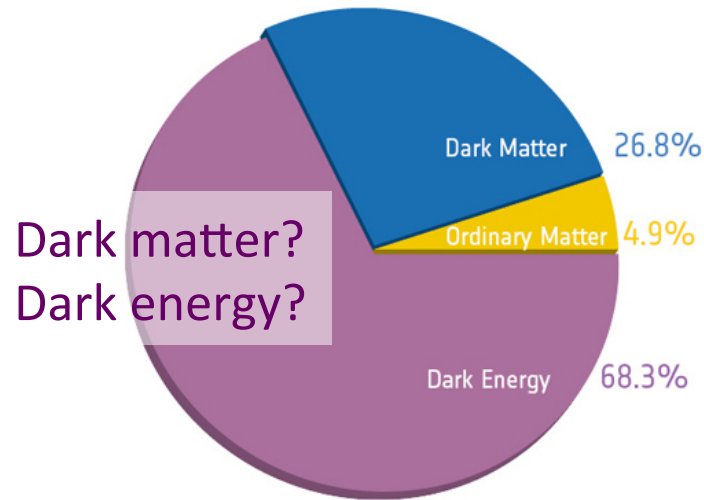
- BEH mechanism provides mass to the W and Z bosons
- Higgs boson regularizes weak vector-boson scattering and makes the theory consistent

Why go beyond the Standard Model?

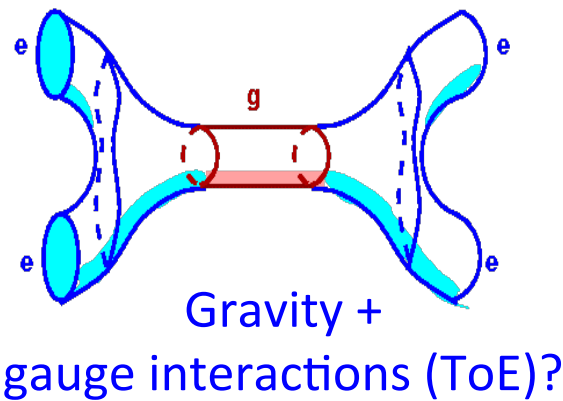
Conceptual weaknesses, observational hints!



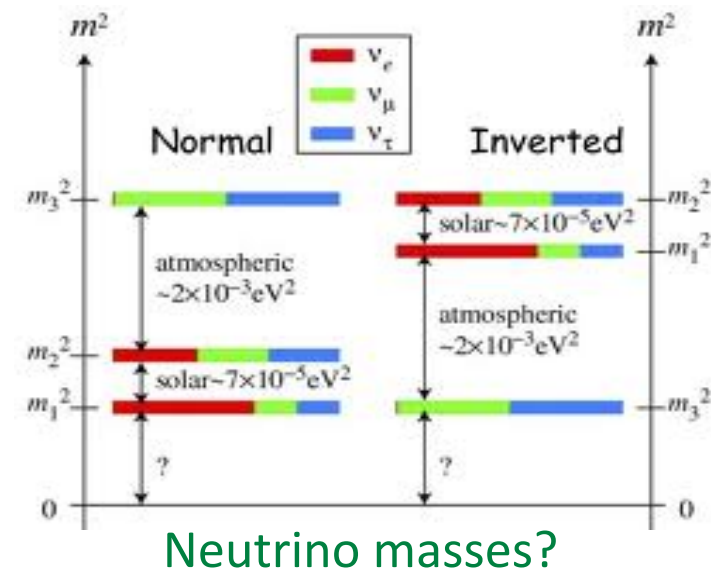
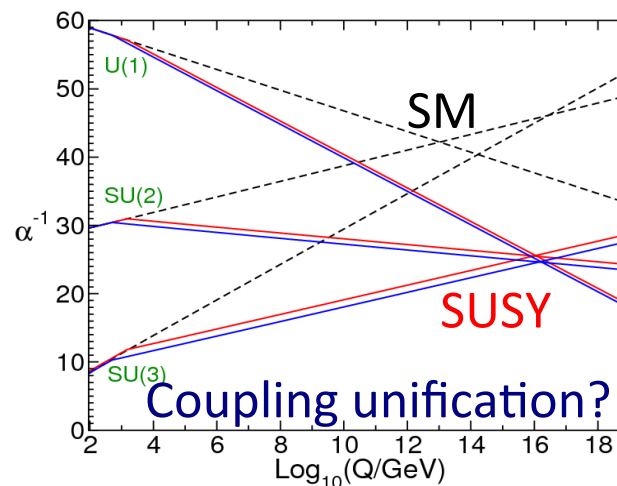
Higgs mass stability?
(Hierarchy problem)



Matter – anti-matter
asymmetry? (CP violation)



Gravity +
gauge interactions (ToE)?



The way ahead

to understand EWSB and discover the more fundamental theory

Some questions after the Higgs discovery:

- Properties as predicted by SM? Spin, parity, CP violation, couplings? With what precision can we extract them?
- Can we measure couplings to 2nd generation fermions (muon, c quark)?
- Can we reconstruct the Higgs potential? Can we measure Higgs self coupling, i.e. di-Higgs production?
- Does the Higgs boson regularize the VV scattering fully? Do other processes beyond the SM contribute?
- Is the new boson alone? Can we discover its (neutral and/or charged) partners?
- Is it a fundamental scalar or a composite particle?
- What stabilizes the Higgs mass against radiative corrections? SUSY? ED? Something else?
- Does the Higgs open a window to new physics?

The program:

- 1) Measure precisely the properties of the Higgs boson
- 2) Look for additional Higgs bosons
- 3) Study massive EW vector boson scattering and triple gauge boson production
 - measure Quartic Gauge Couplings / test the SM gauge structure
 - sensitive to the scalar sector / EWSB model
- 4) Look directly for new phenomena
(SUSY, new gauge bosons W'/Z' , extra dimensions... and the unexpected!)

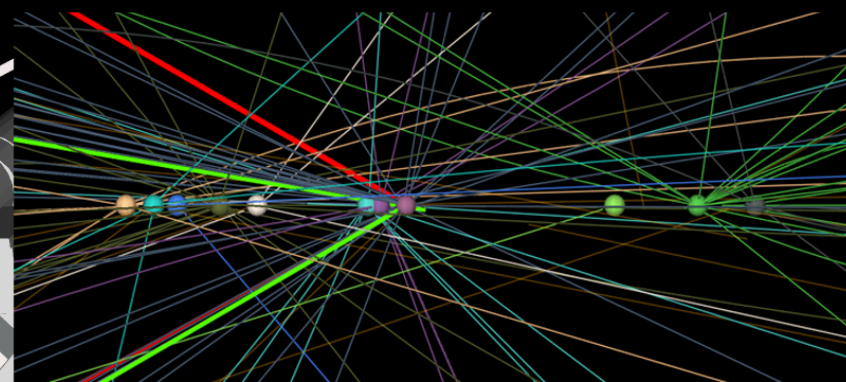
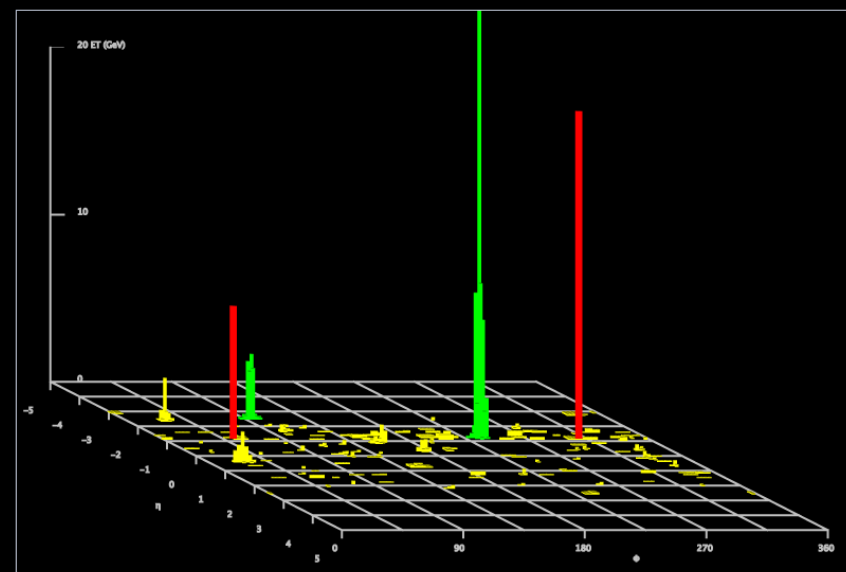
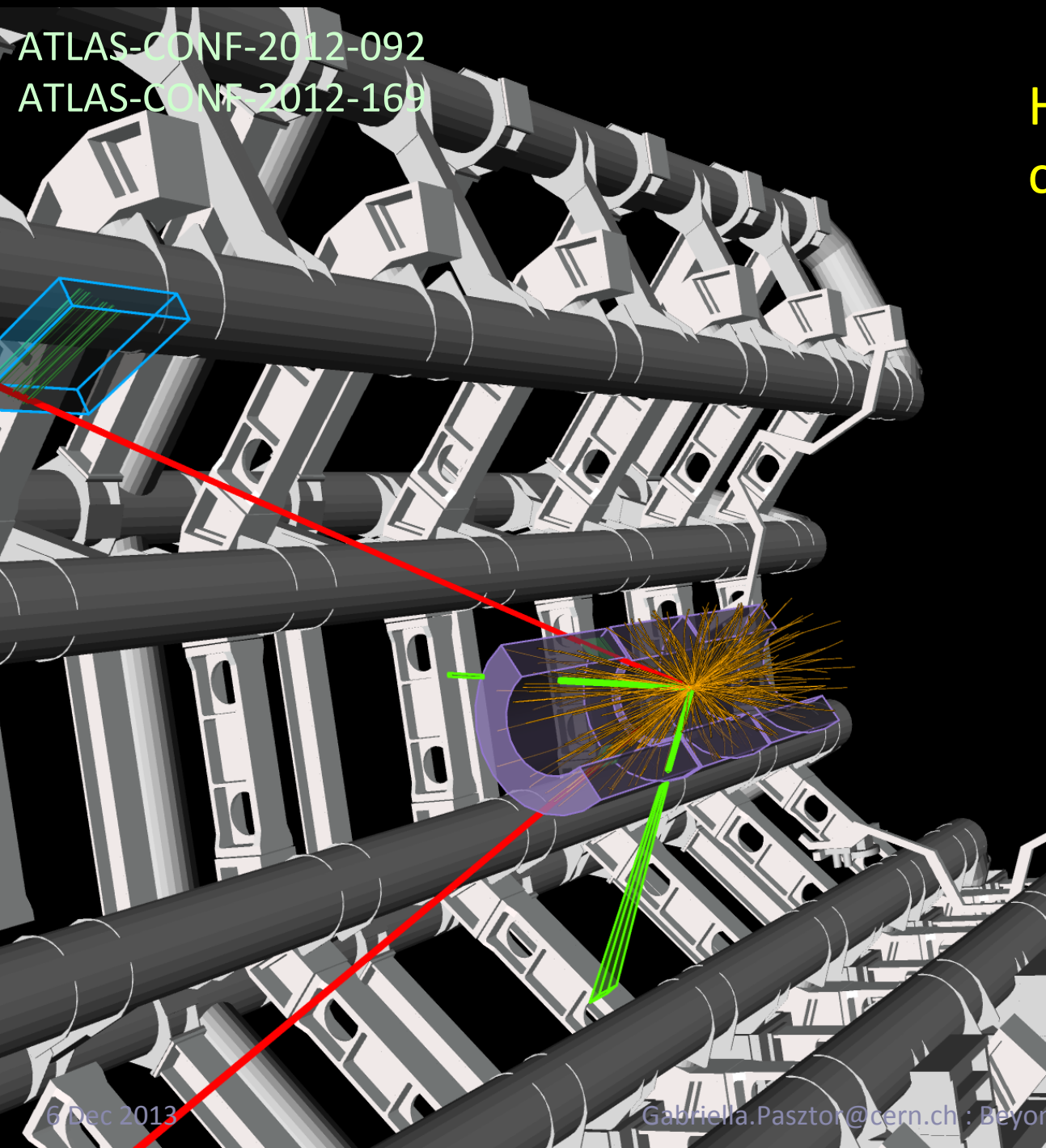
Higgs property measurements

ATLAS-CONF-2012-092
ATLAS-CONF-2012-169

$H \rightarrow ZZ^* \rightarrow ee \mu\mu$
candidate


ATLAS
EXPERIMENT
<http://atlas.ch>

Run: 205113
Event: 12611816
Date: 2012-06-18
Time: 11:07:47 CEST

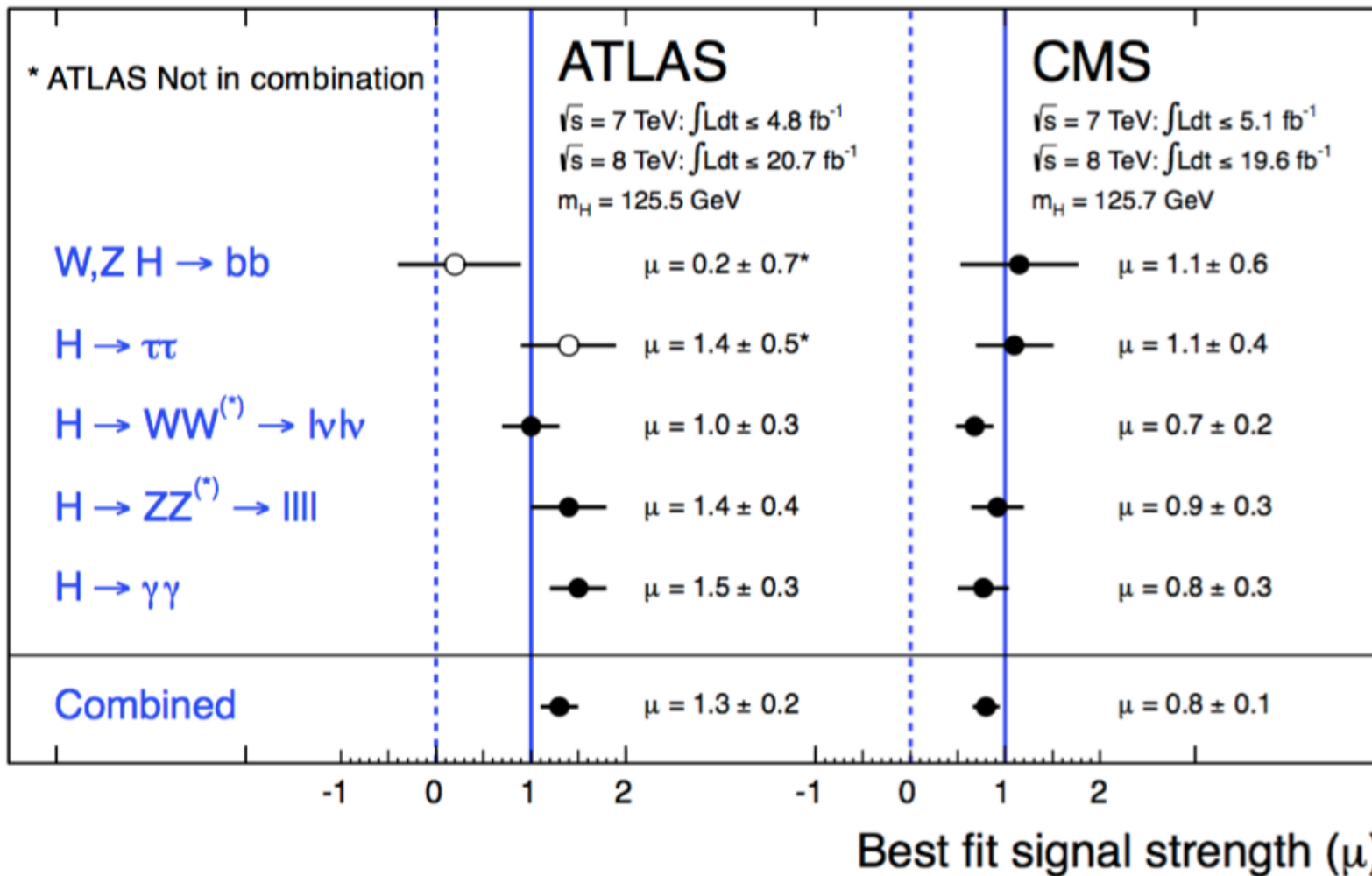


Higgs production cross-section

$$\mu = \frac{\sigma \cdot BR}{(\sigma \cdot BR)_{SM}}$$

PLB 726 (2013) 88-119

CMS-PAS-HIG-13-005



Tevatron combined
 (bb, $\tau\tau$, $\gamma\gamma$, WW^*):
 $\mu = 1.44 \pm 0.6$

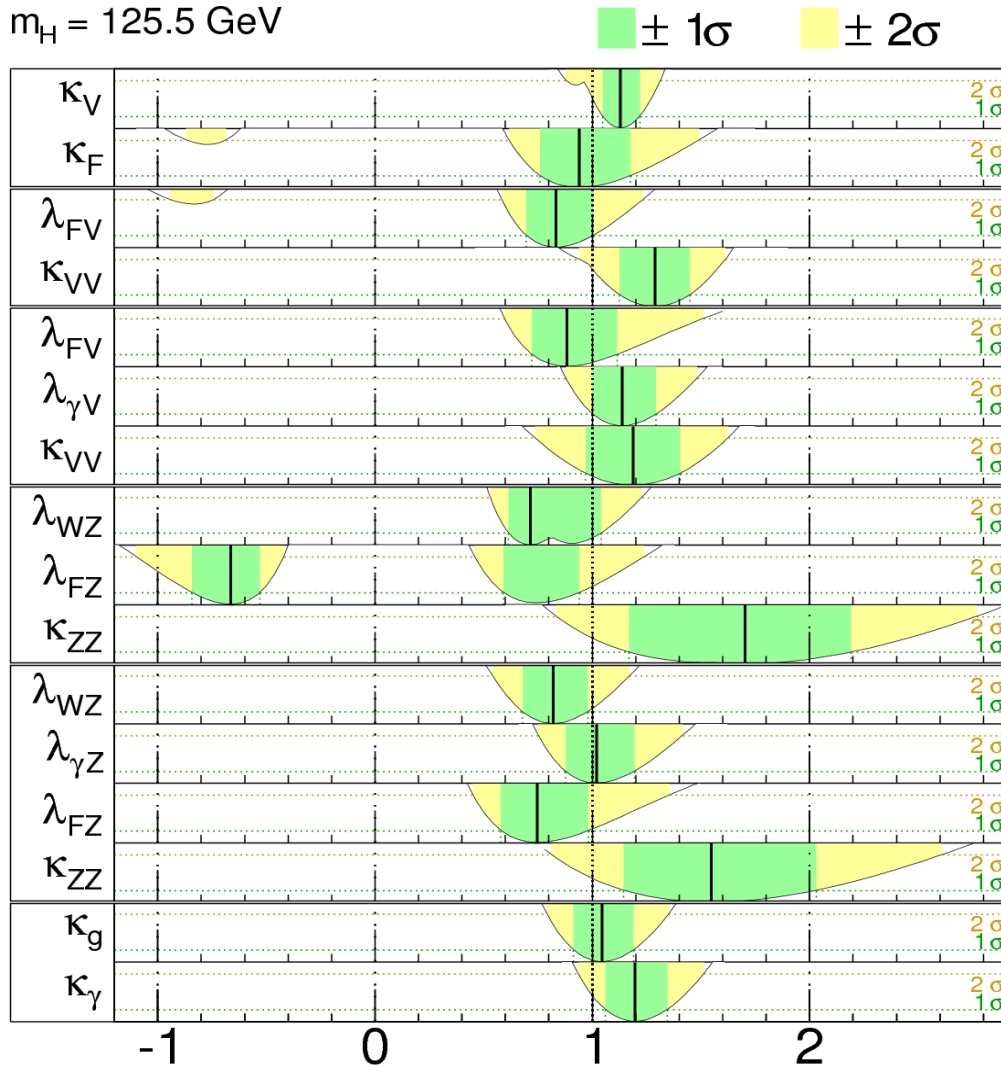
Excellent agreement with SM within current precision of 10-20%

Coupling constant scale factors

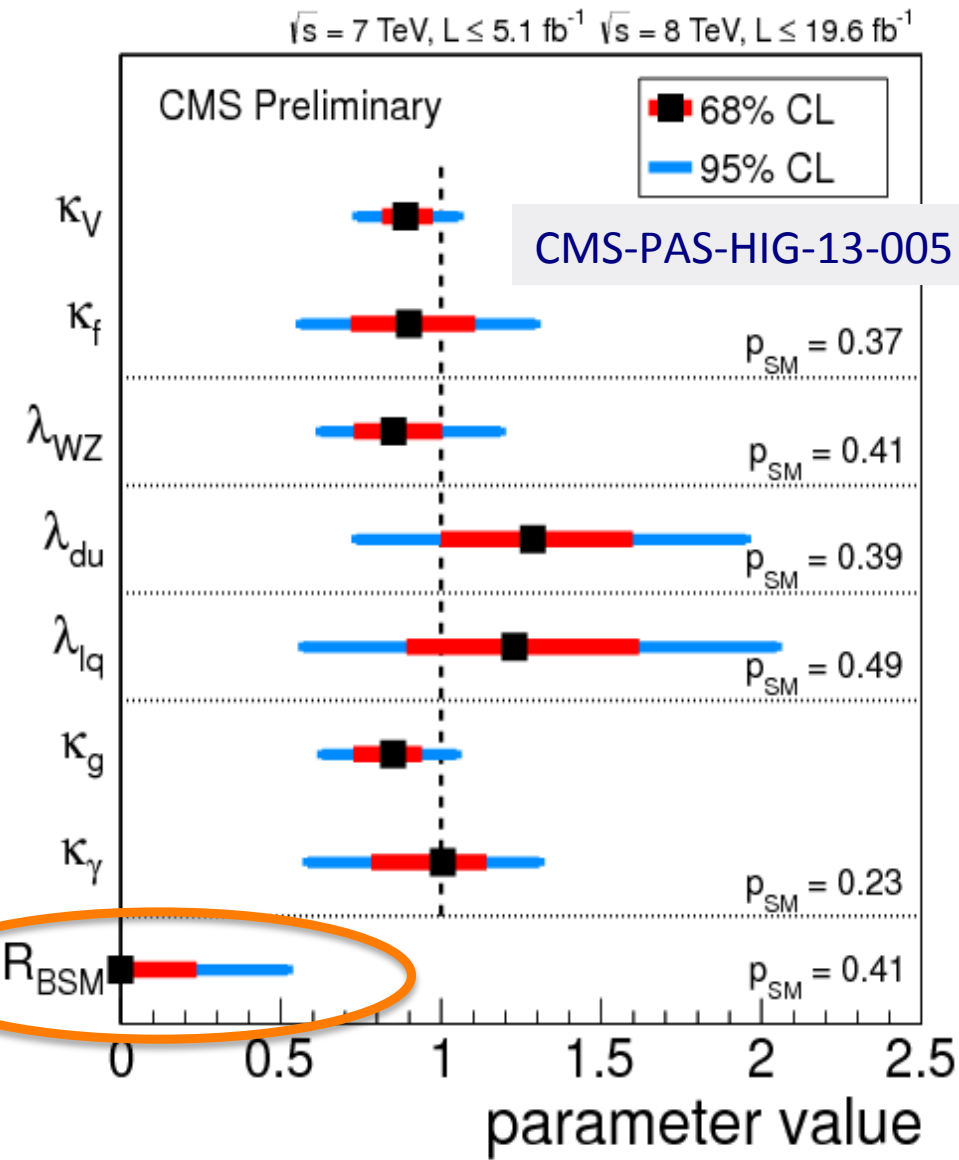
$$\kappa_i \equiv g_i / g_i^{\text{SM}}$$

Measuring different production and decay channels....

$m_H = 125.5 \text{ GeV}$



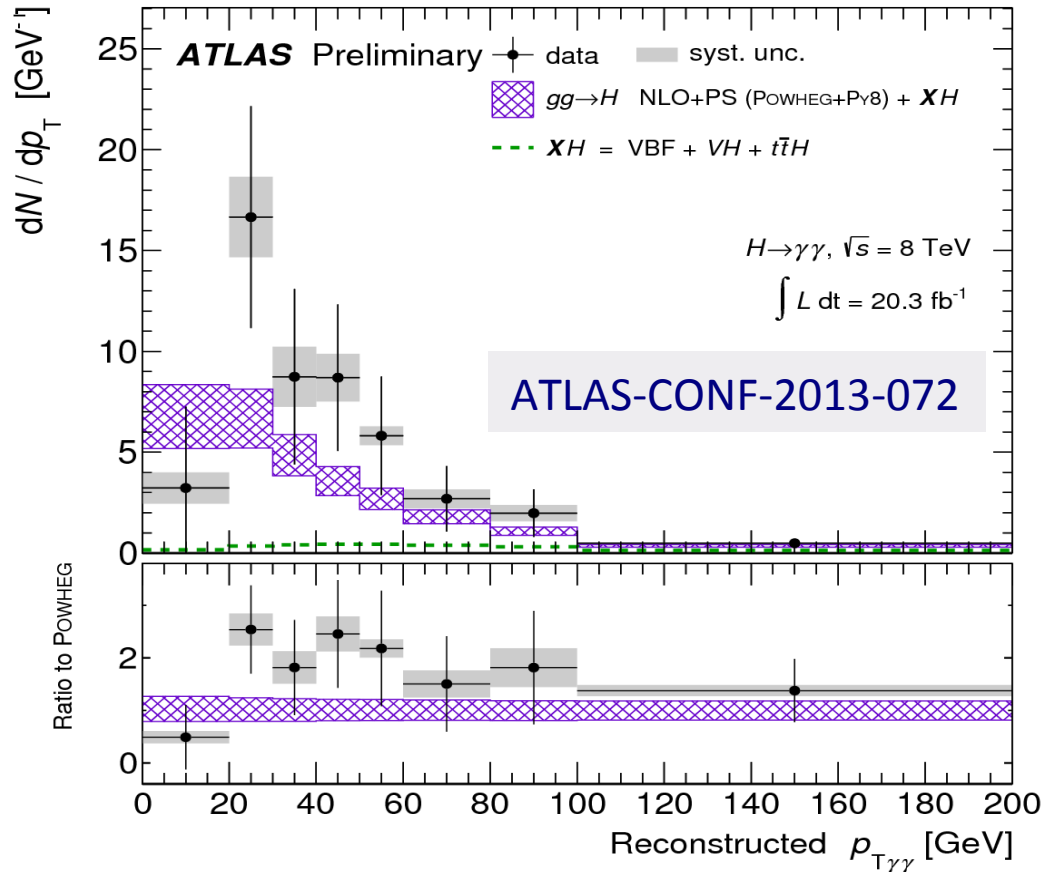
PLB 726 (2013) 88-119



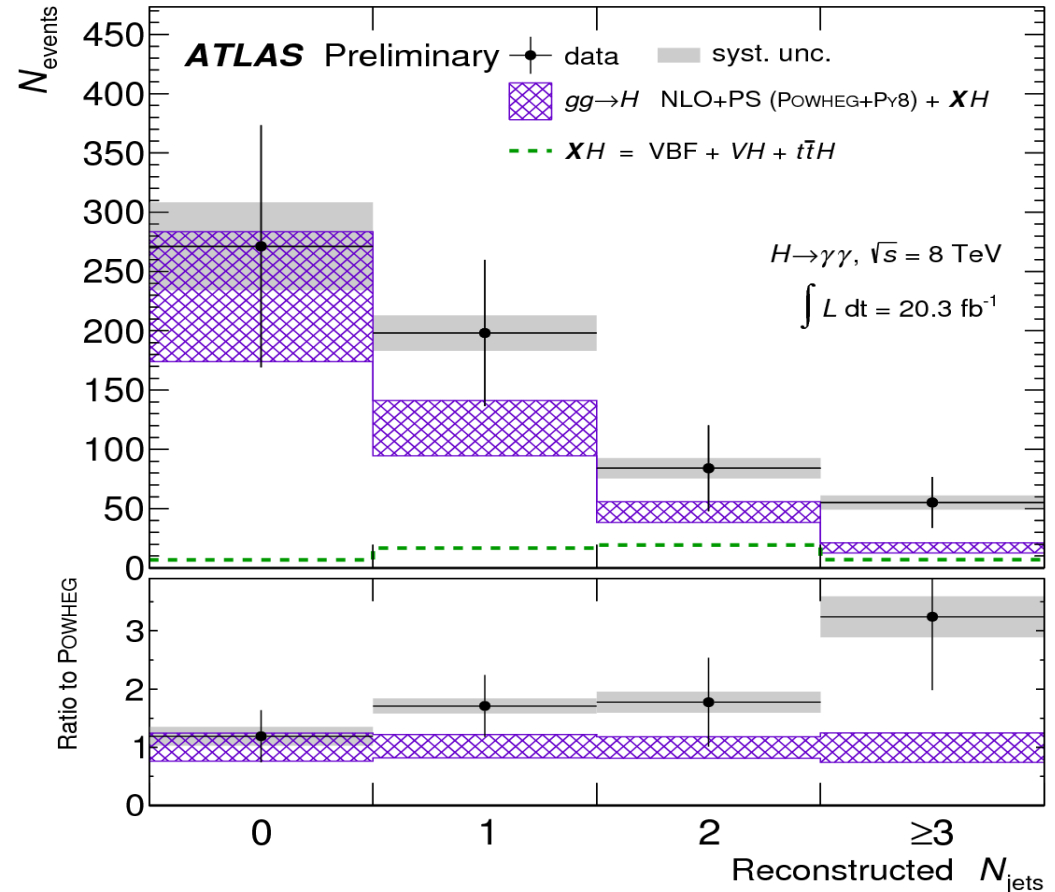
There is still place for new physics, exotic decays!

Entering the Higgs measurement era: $H \rightarrow \gamma\gamma$ differential cross-sections

Test perturbative QCD calculations for the dominant ggH process



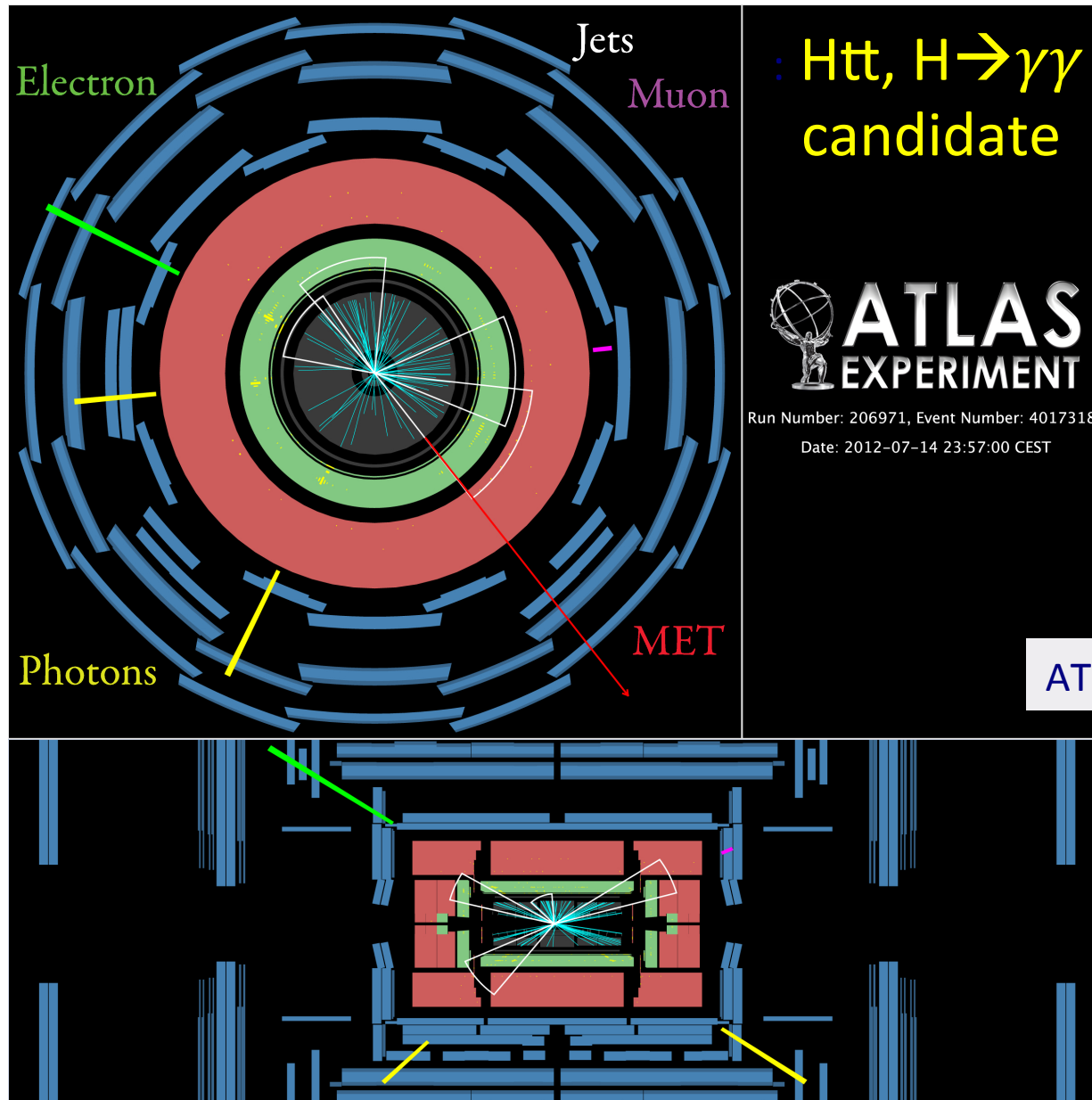
Sensitive to relative rate of production via different mechanisms



Reduced model dependence!

If the BSM model has non-SM like tensor structure, i.e. predicts different detector acceptances, selection efficiencies, fiducial cross-sections unfolded to particle level provide a better source of information

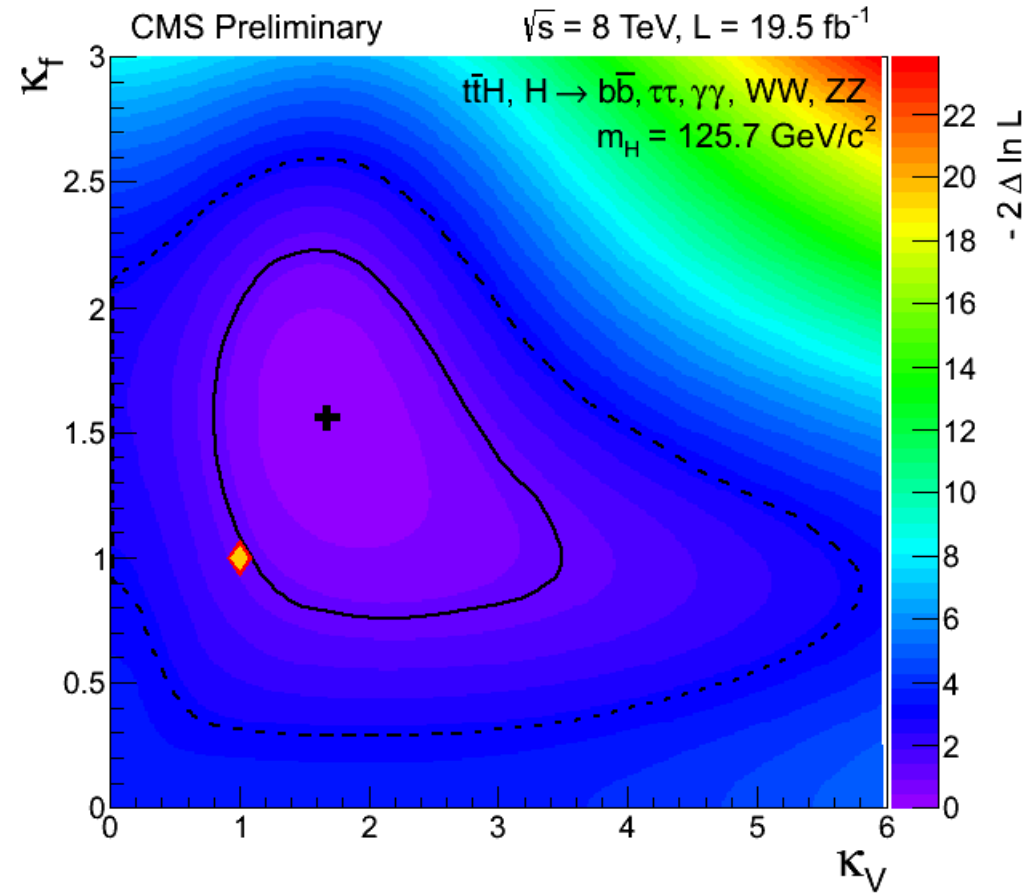
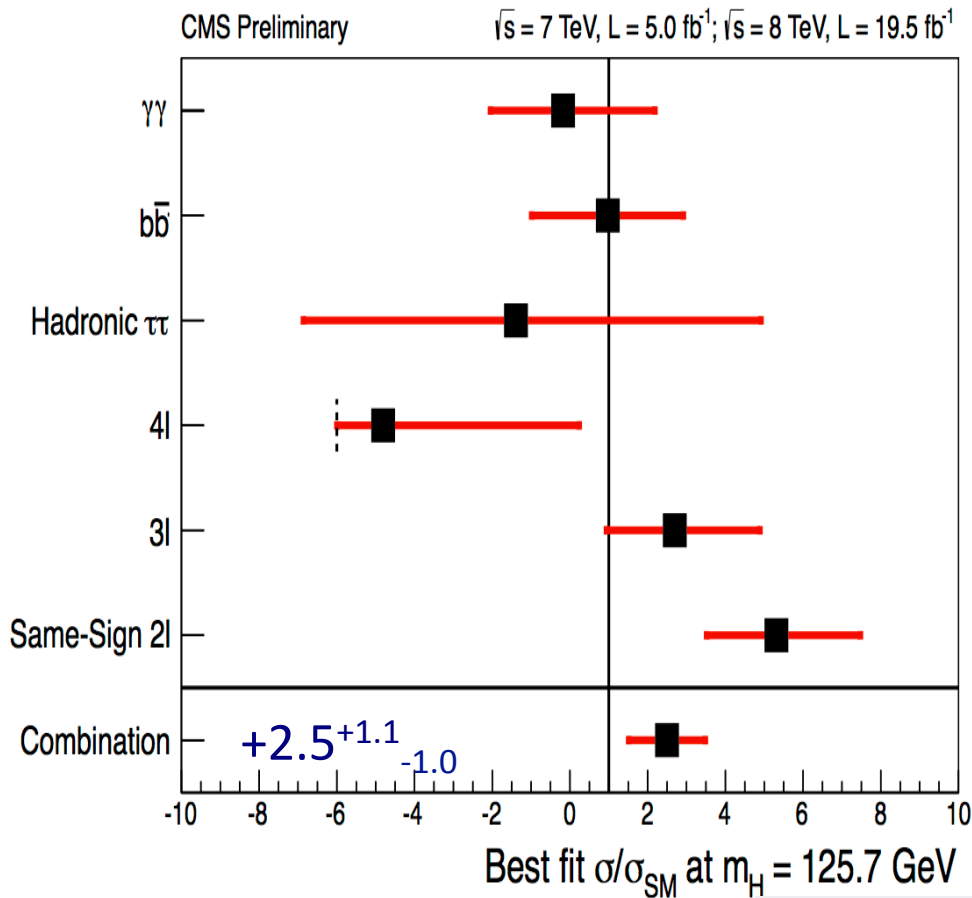
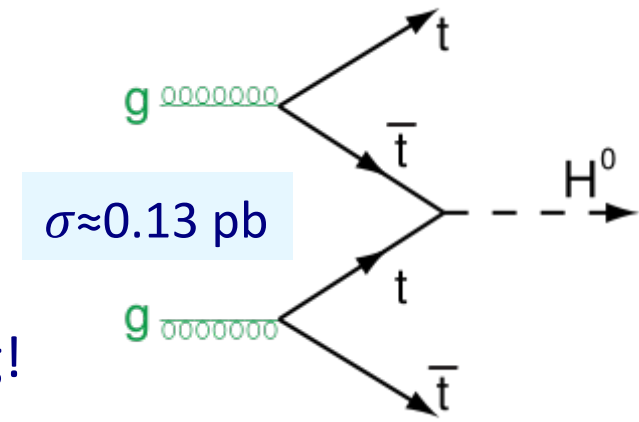
Measuring rare processes



Search for ttH production

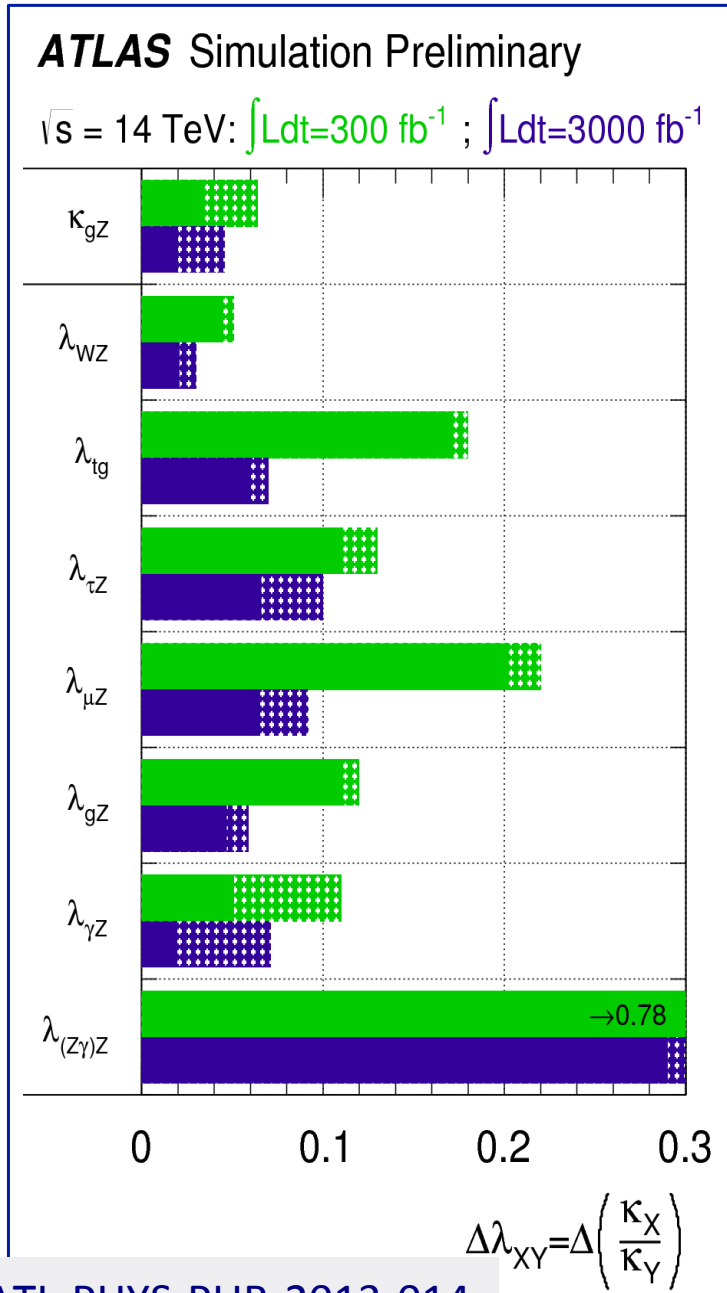
Direct access to Htt coupling

Large background, small cross-section → very challenging!

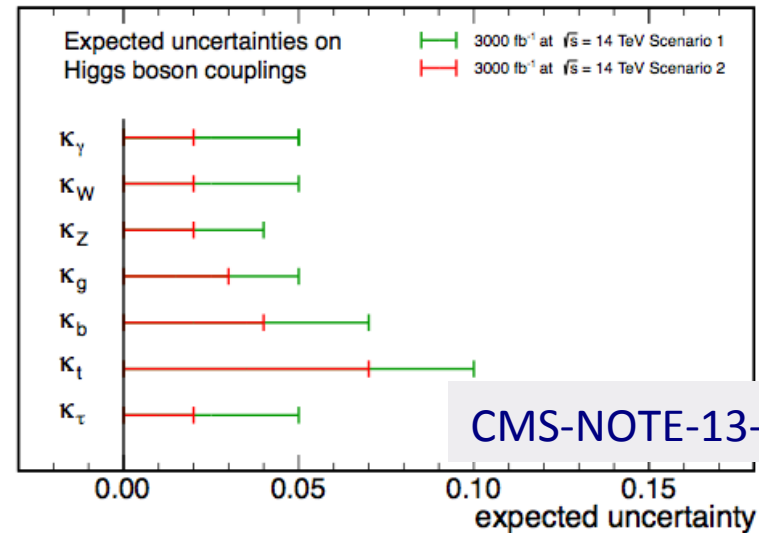
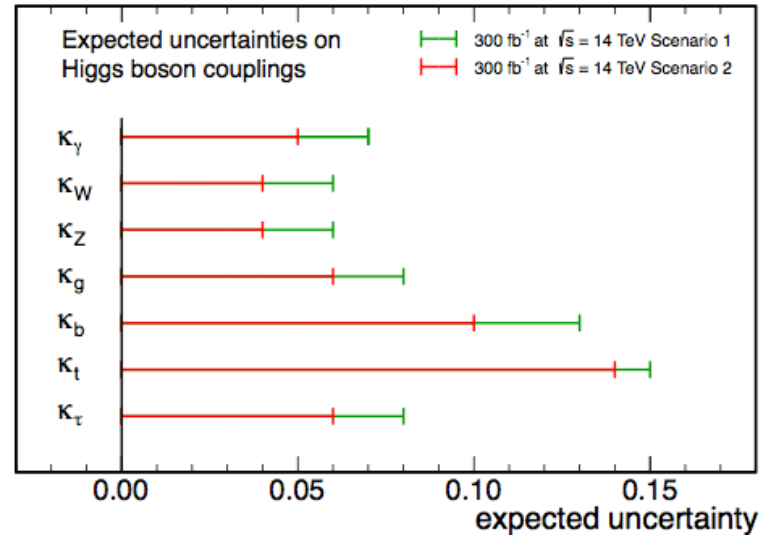


<https://twiki.cern.ch/twiki/bin/view/CMSPublic/ttHCombinationTWiki>

Projections for Run 2 and HL-LHC ($E_{cm}=14$ TeV)



CMS Projection



CMS-NOTE-13-002

Most difficult task:
 Measurement of
 Higgs self-interaction
 CMS+ATLAS
 combined: ~30%

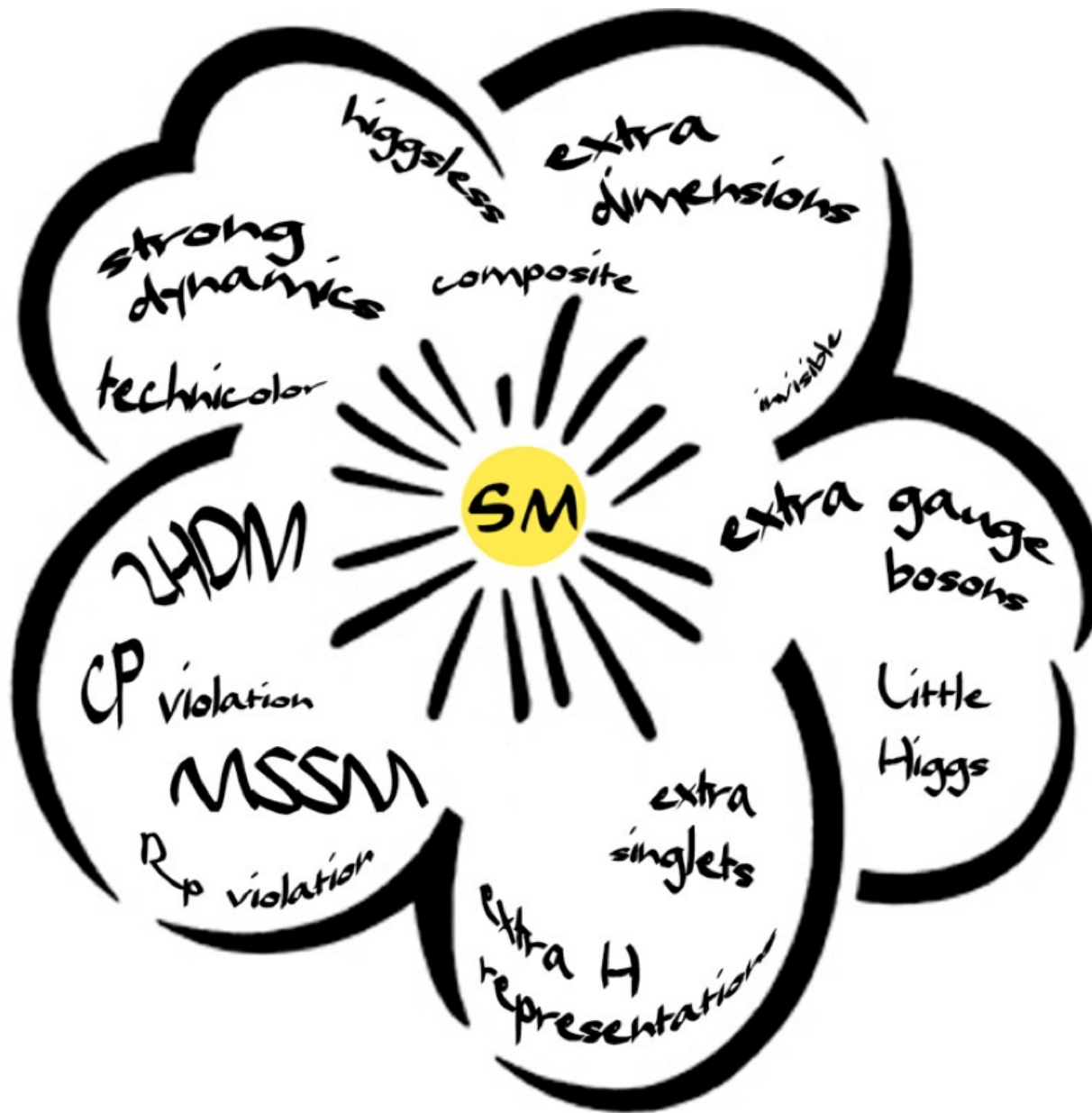
- Theory „requires”
 <20% typically
 to see SUSY
 contributions

arXiv:1305.6397v2 [hep-ph]

CMS scenario 1: syst uncertainties as in Run1

scenario 2: etheory syst x 0,5; experimental $\sim 1/\sqrt{L}$

Going beyond...



Extended Higgs sector

Standard Model

1 complex scalar doublet
1 physical Higgs boson: H

2HDM

2 complex scalar doublets
5 physical Higgs bosons: h, H, A, H^+, H^-

MSSM (Type-II 2HDM)

2 complex scalar doublets
5 physical Higgs bosons: h, H, A, H^+, H^-

NMSSM (~~μ -problem of MSSM~~)

2 complex scalar doublets + 1 singlet
7 physical Higgs bosons:
 $h_1, h_2, h_3, a_1, a_2, h^+, h^-$

Additional SM-like Higgs

(\rightarrow high-mass searches)

Fermiophobic Higgs

Invisible Higgs

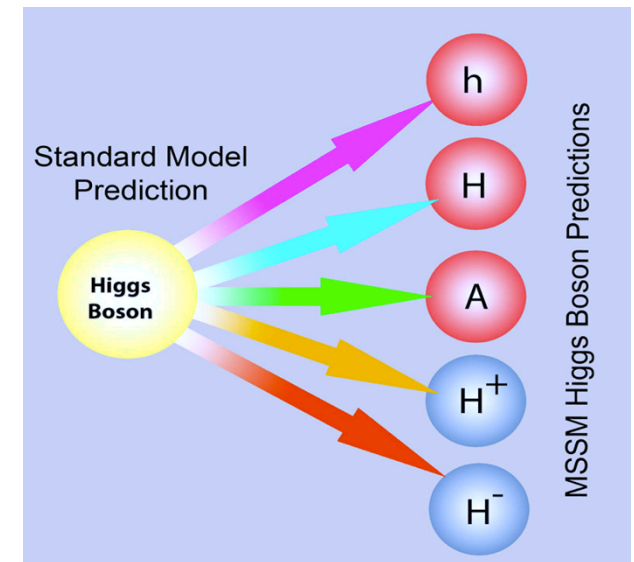
e.g. decaying to neutral LSP

„Exotic” Higgs

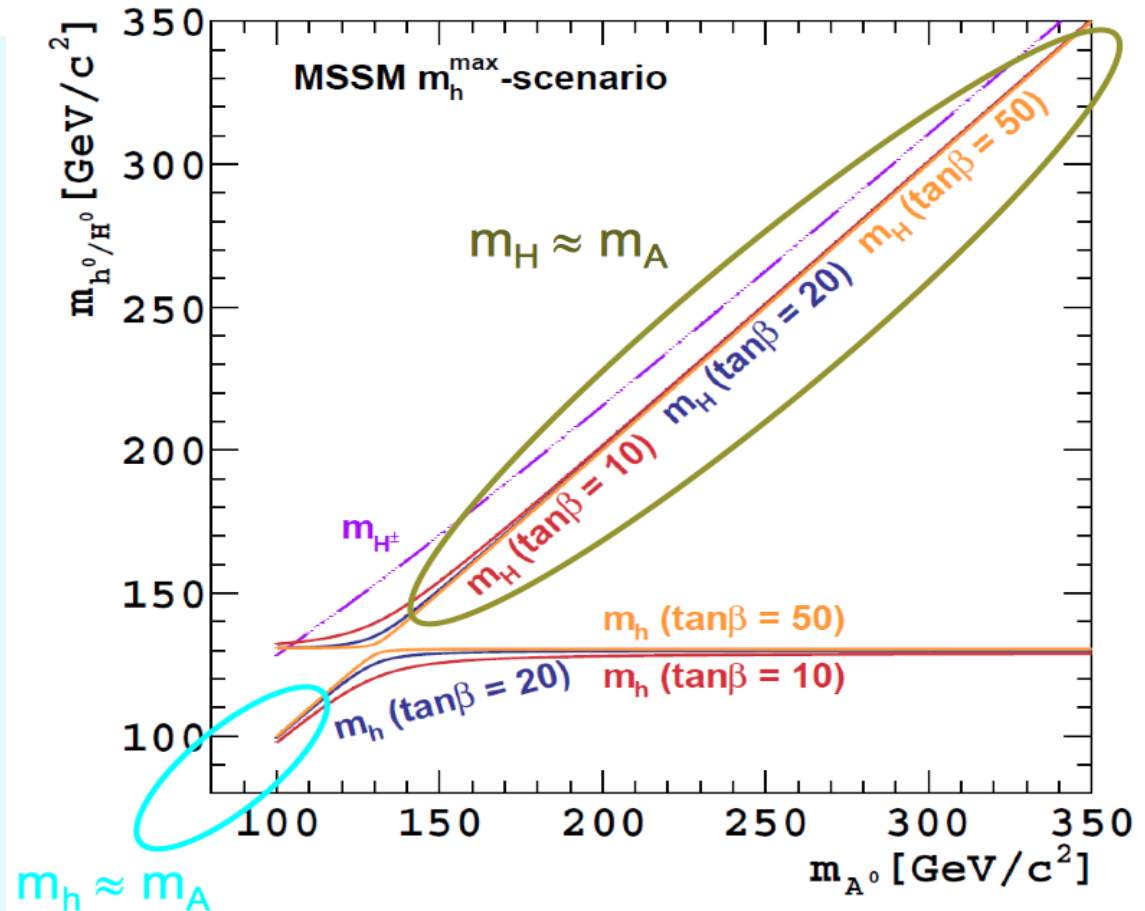
e.g. decaying to lepton-jets in
hidden-valley SUSY

MSSM Higgs sector

- 2 complex scalar doublets \rightarrow 5 physical Higgs bosons
- Tree-level parameters: m_A , $\tan\beta$
- Many more parameters after radiative corrections
 \rightarrow need benchmark scenarios
- Show only m_h^{\max} in this talk

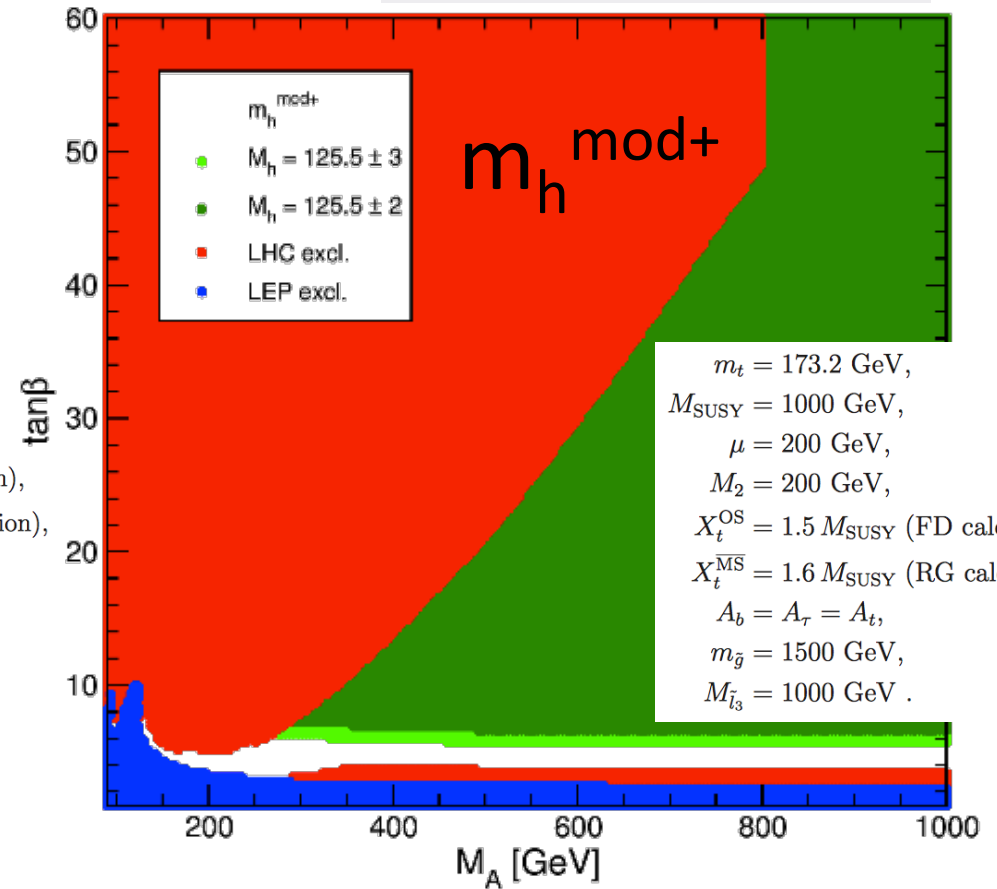
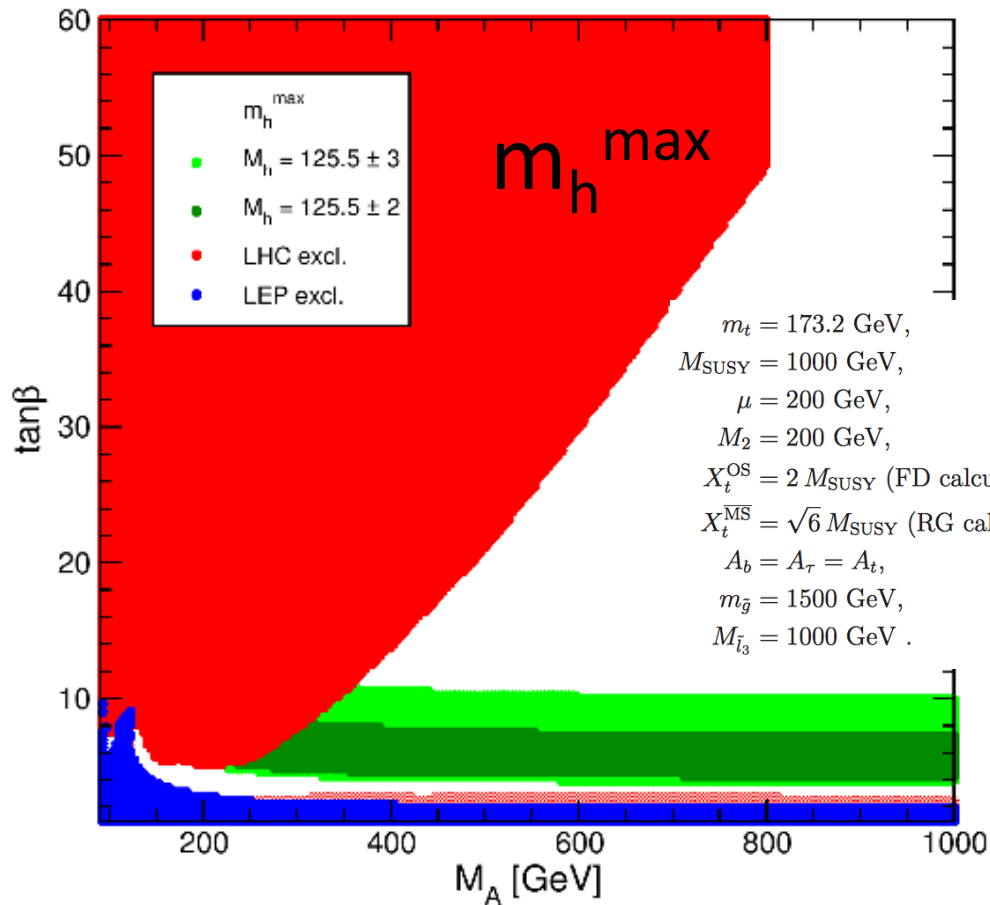


- In MSSM, at large m_A ($\gg m_Z$), h usually becomes SM-like (decoupling limit)
 \rightarrow **Direct search for additional Higgs states important!**
- Rise of phenomenological-MSSM (pMSSM- n): general scan of the n most important parameters (while fixing the less important ones to fit precision data)
- Remember: while general MSSM has 100+ parameters from the soft susy breaking terms of the Lagrangian, relations are expected among them depending on how SUSY is broken)



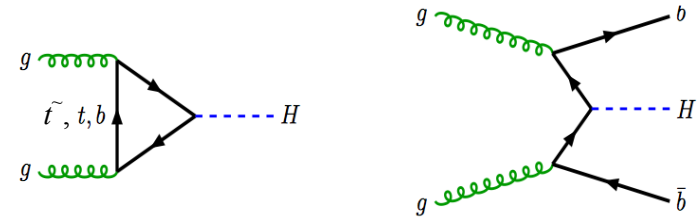
Consistency with $m_h = 125$ GeV

arXiv:1302.7033 [hep-ph]

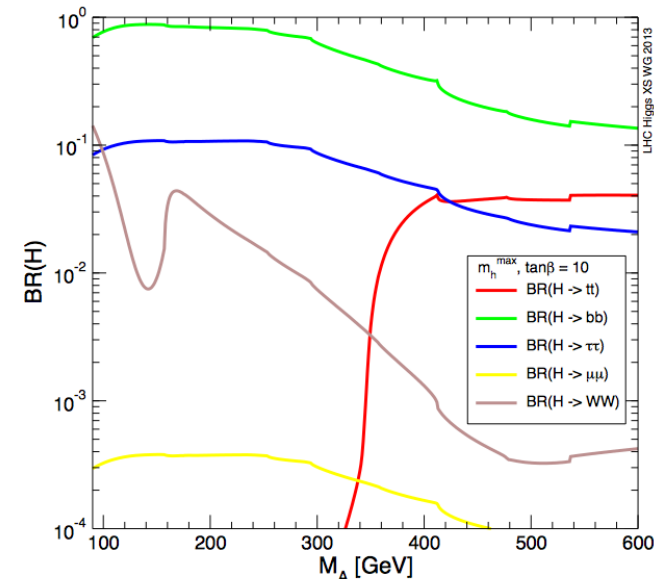
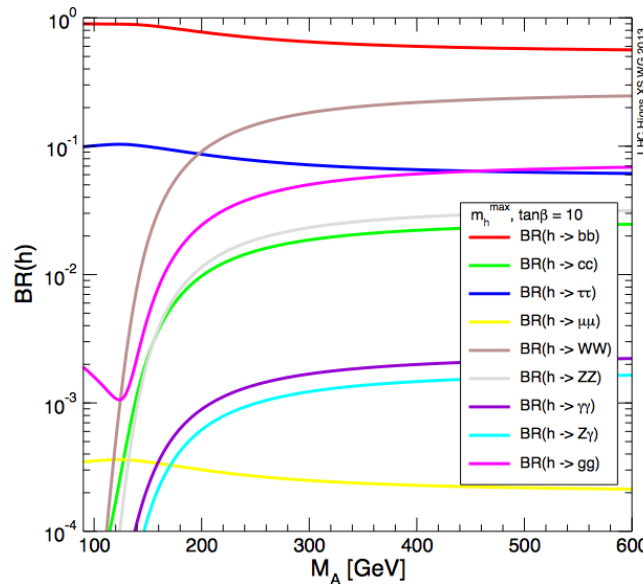
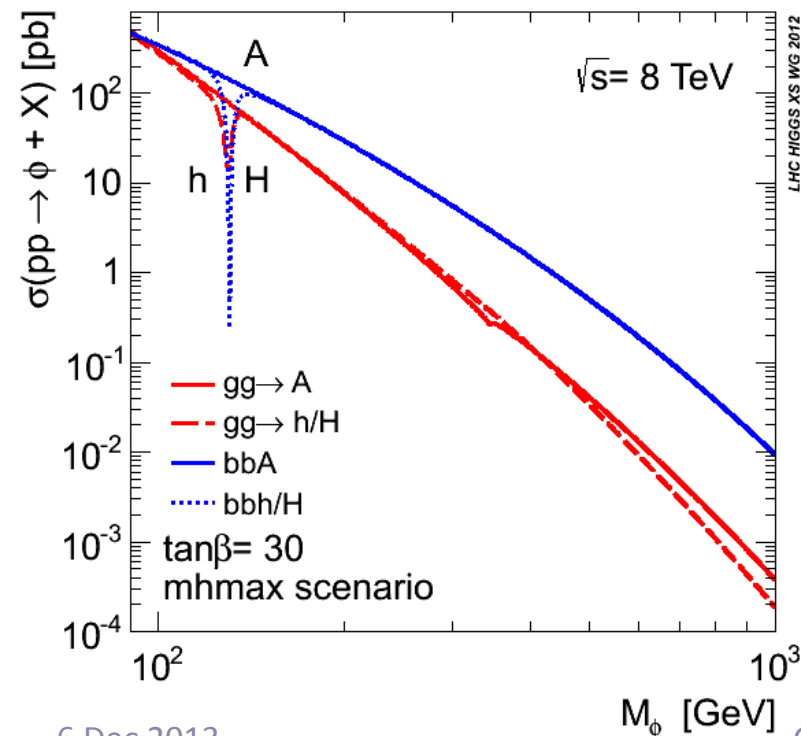
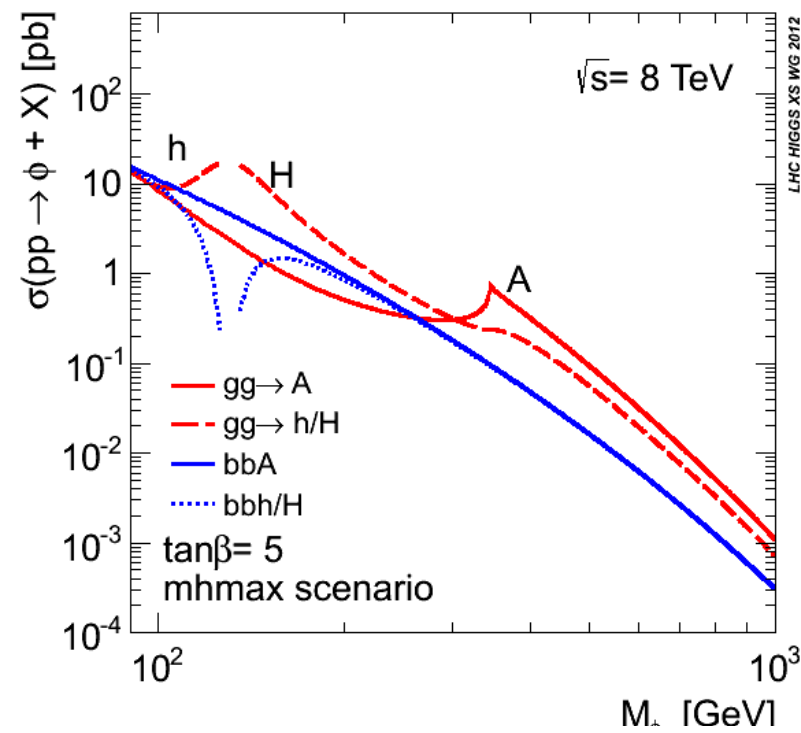


Production and decay

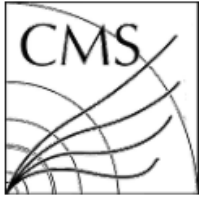
- Associated $bb\phi$ production plays an important role



- Strongly enhanced cross-section at high $\tan\beta$
- Decays $\phi \rightarrow bb, \tau\tau$ are important also at high mass
- $\phi \rightarrow bb$ very challenging (huge background)
- $\phi \rightarrow \mu\mu$ very low BR but excellent resolution (could separate H / A when degenerate)

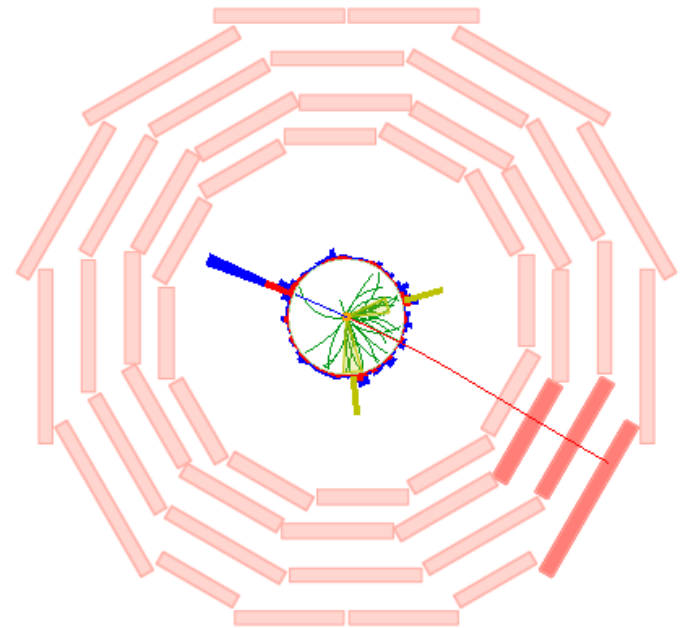
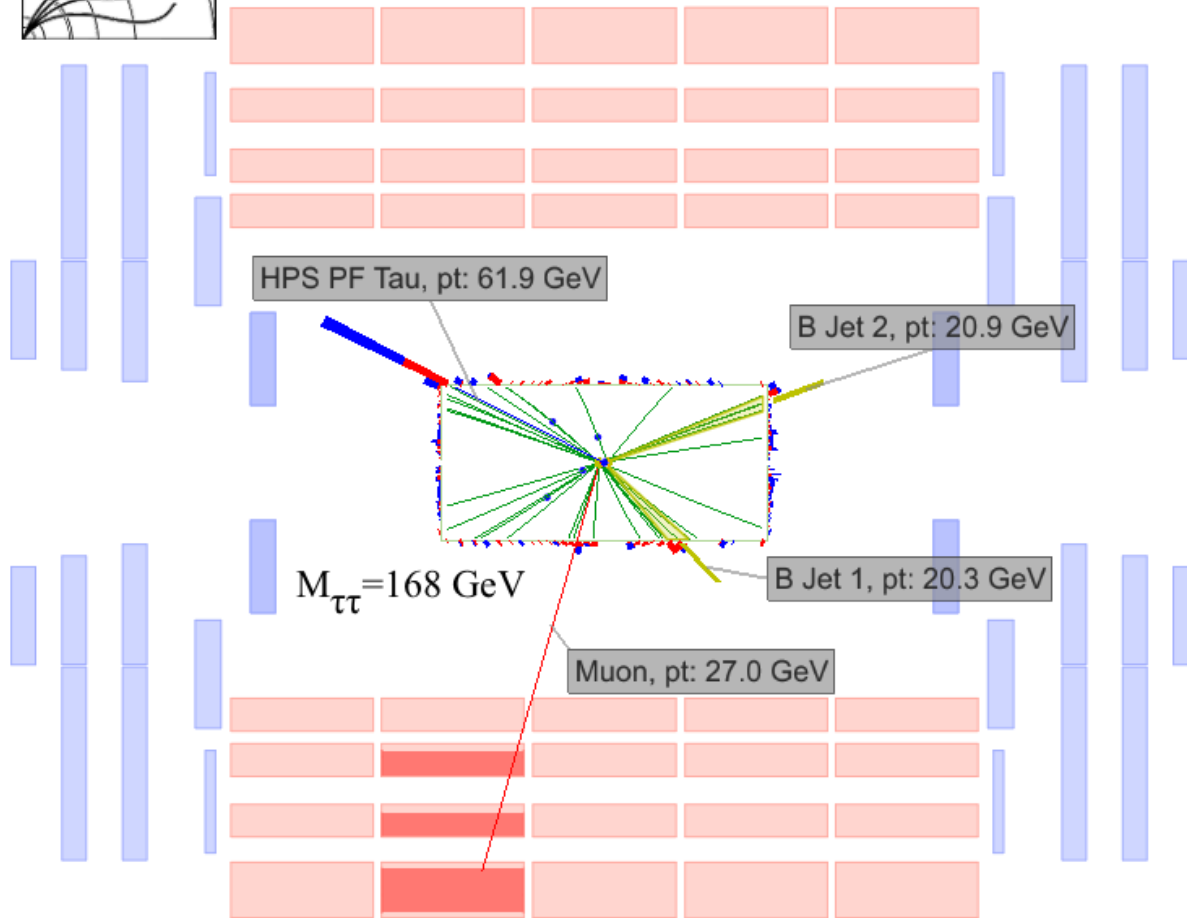


Search for $bb\phi \rightarrow bb(\tau\tau)$



CMS-PAS-HIG-13-021

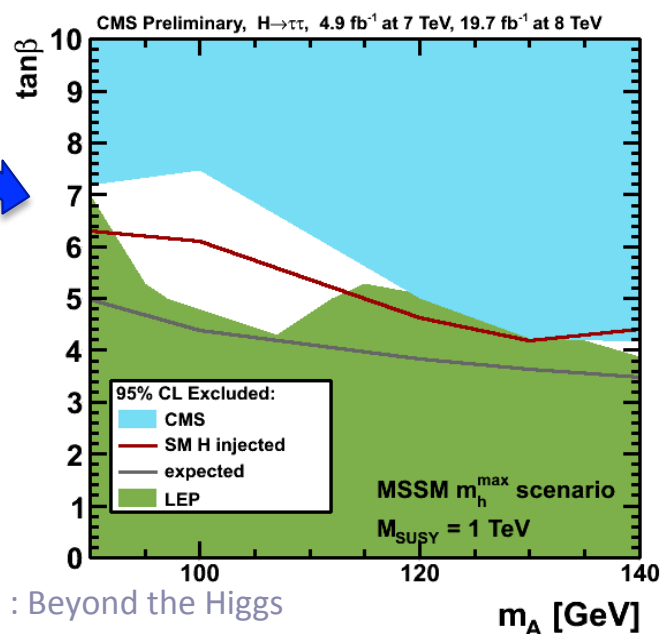
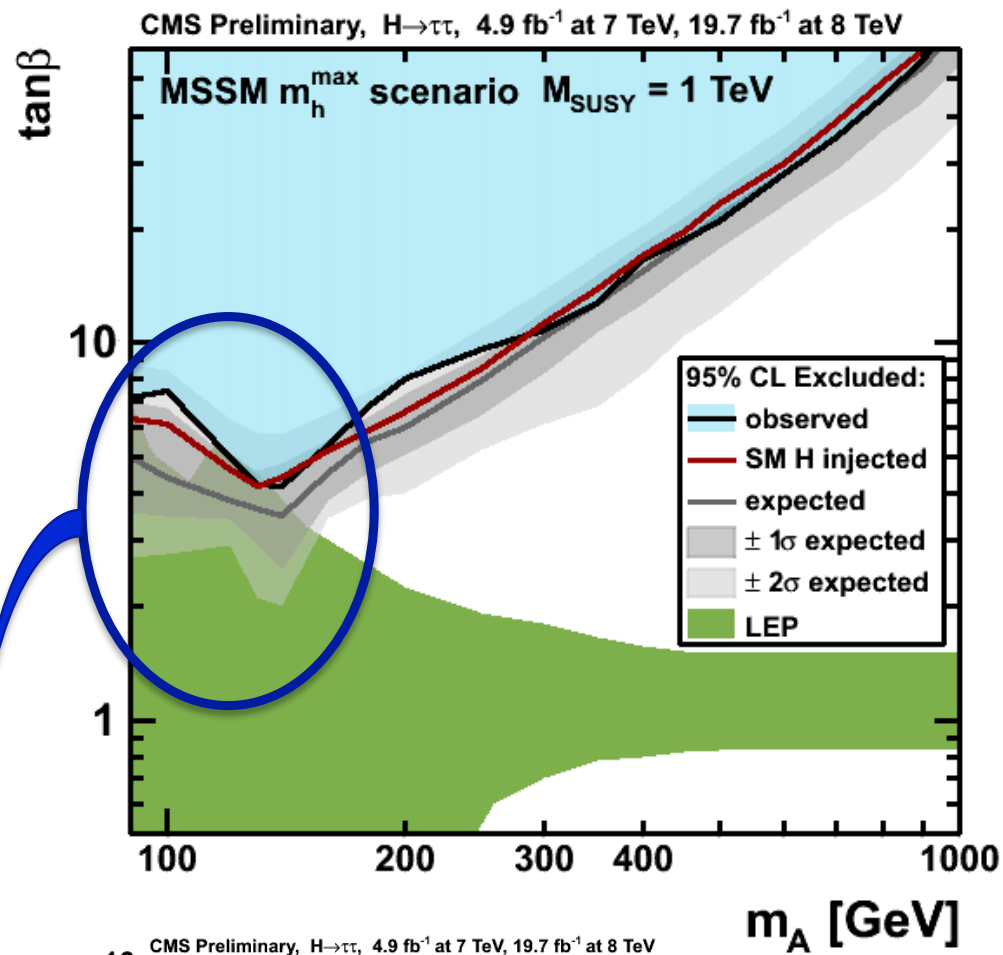
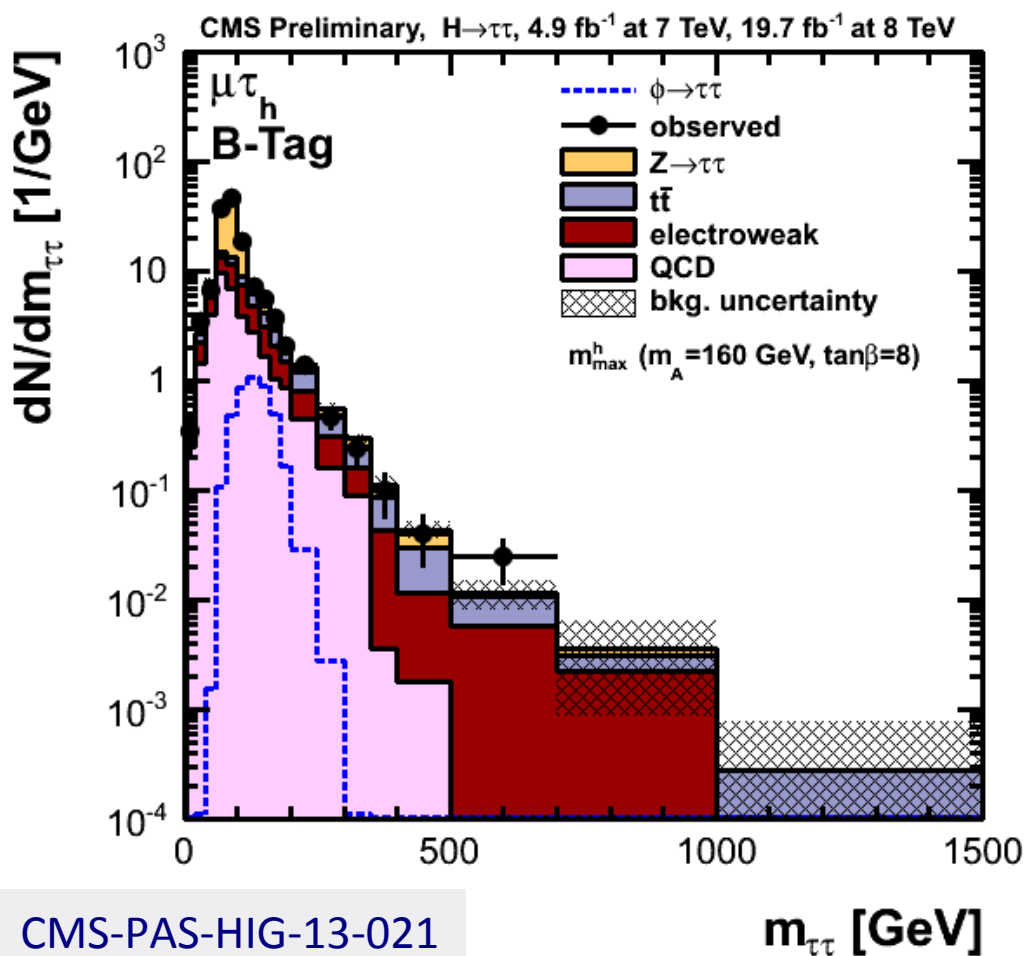
CMS Experiment at LHC, CERN
Data recorded: Mon Oct 3 03:07:23 2011 CEST
Run/Event: 177730 / 2113660794



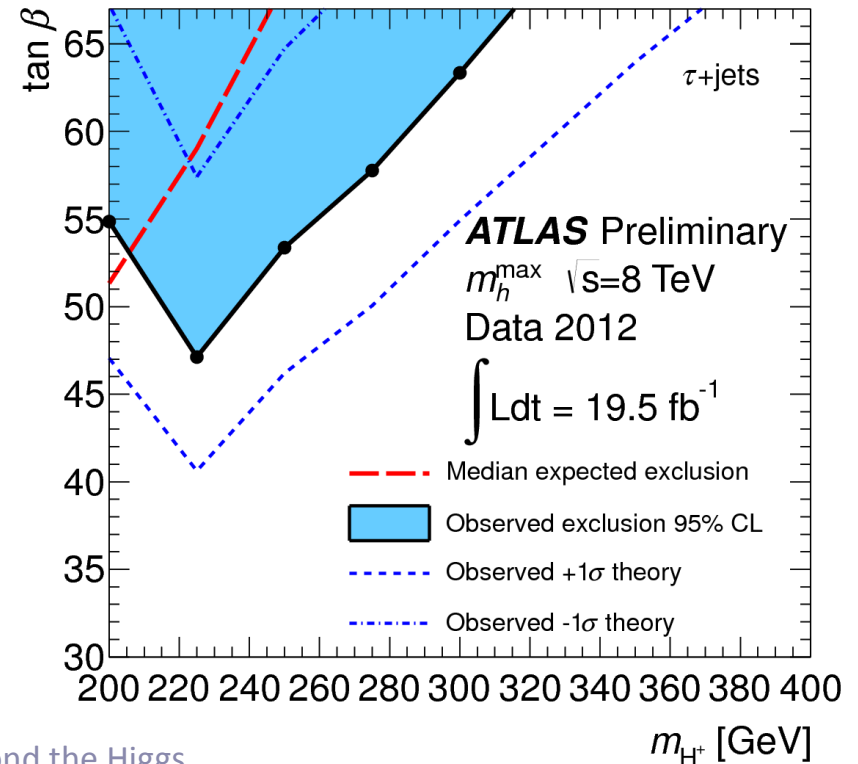
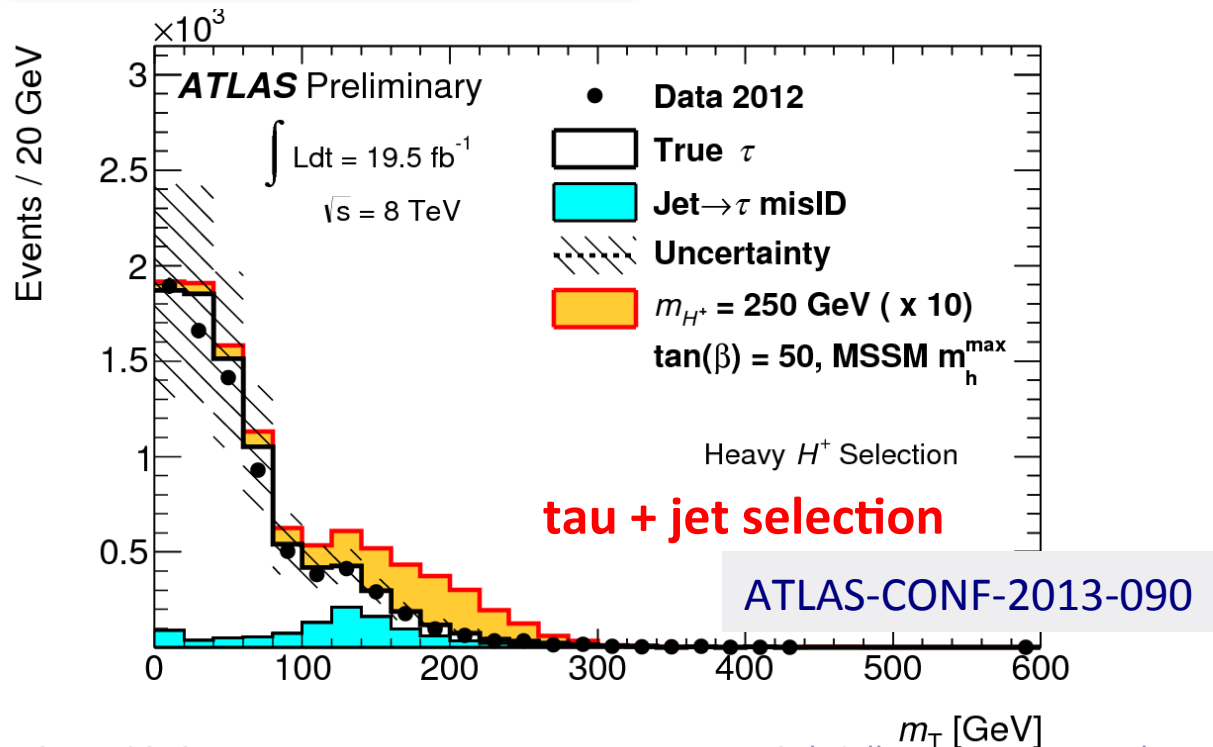
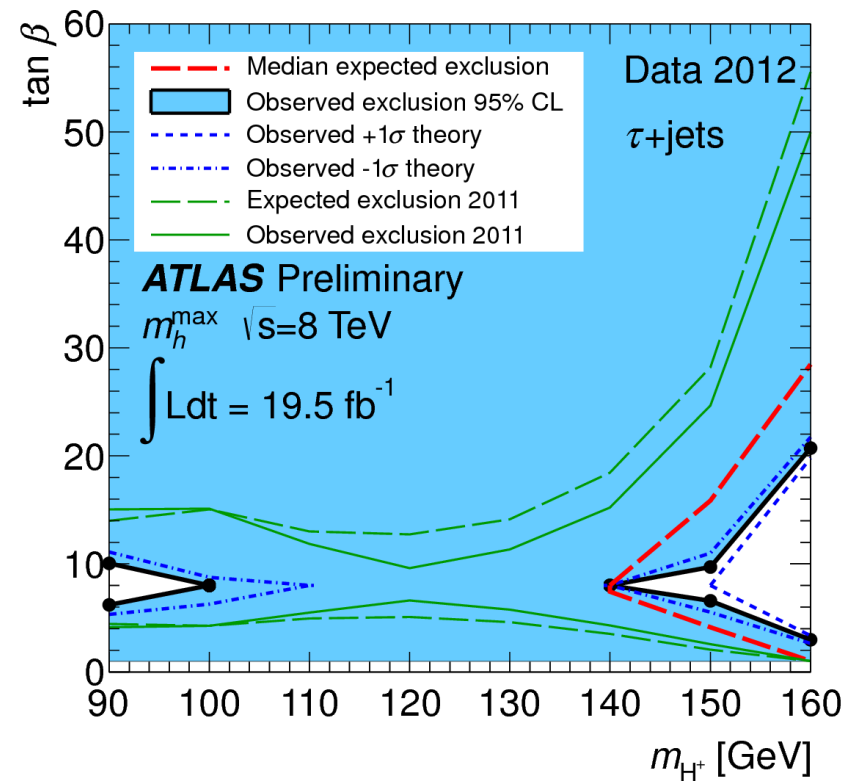
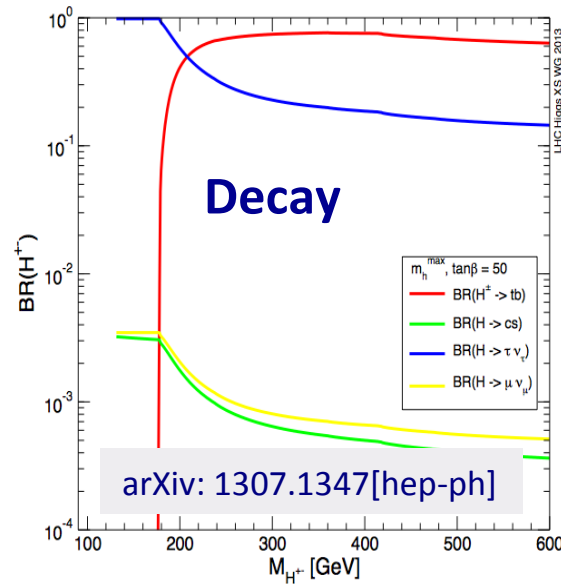
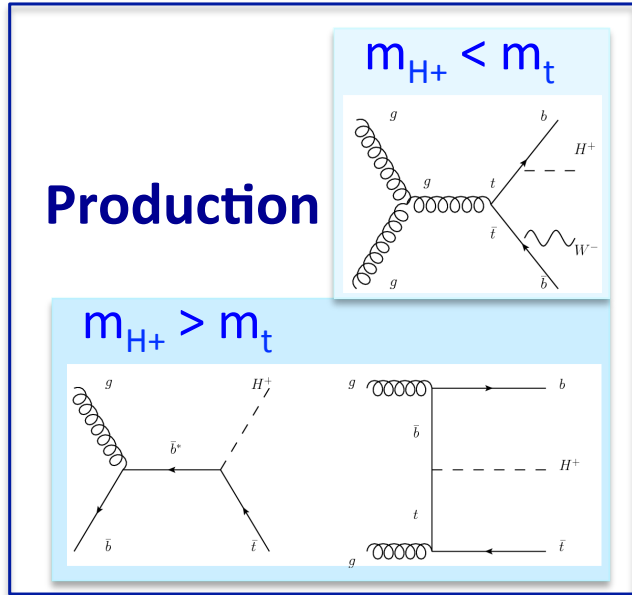
$bb\phi \rightarrow bb(\tau\tau) \rightarrow bb(\text{had } \mu)$ candidate

Search for $\phi \rightarrow \tau\tau$

Background composition differs significantly from tau final state to tau final state, but all distributions are well described by background only hypothesis



Search for Charged Higgs



Searching for exotic Higgs decays

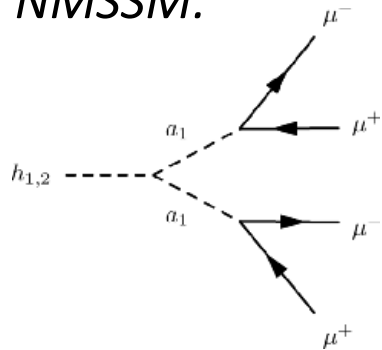
Many „beyond MSSM” extensions, e.g.

- NMSSM (to solve the μ -problem of MSSM): 7 physical Higgs bosons: $h_1, h_2, h_3, a_1, a_2, h^+, h^-$
- Hidden-valley theories with a dark sector (dark-particles can have low mass!)

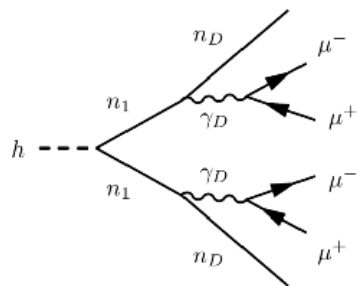
Exotic decay modes open, e.g.

$$h \rightarrow 2a + X \rightarrow 4\mu + X$$

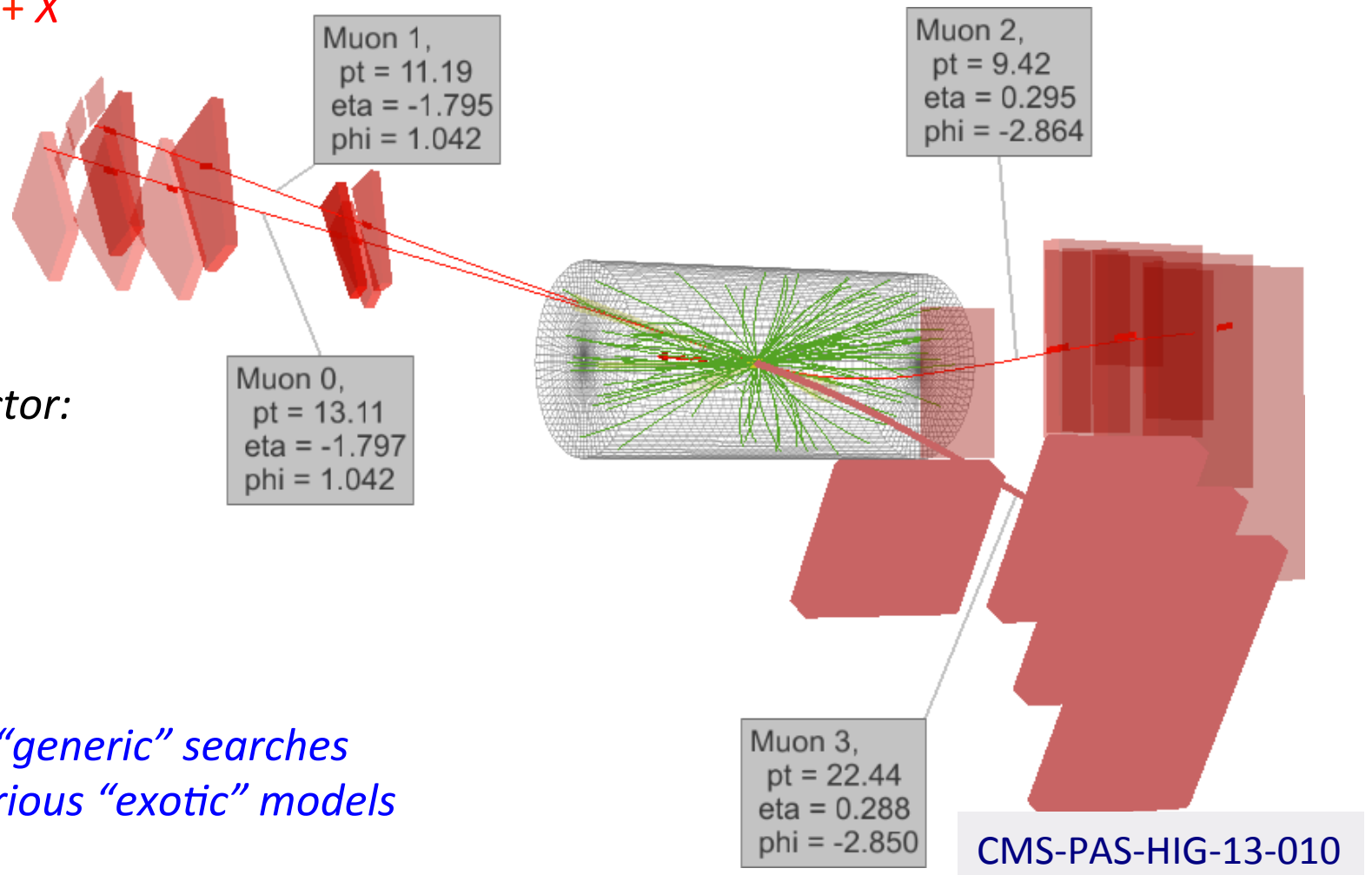
NMSSM:



SUSY with a dark sector:



→ Topology-based, “generic” searches interpreted in various “exotic” models



Search for $h \rightarrow 2a + X \rightarrow 4\mu + X$ at low m_a

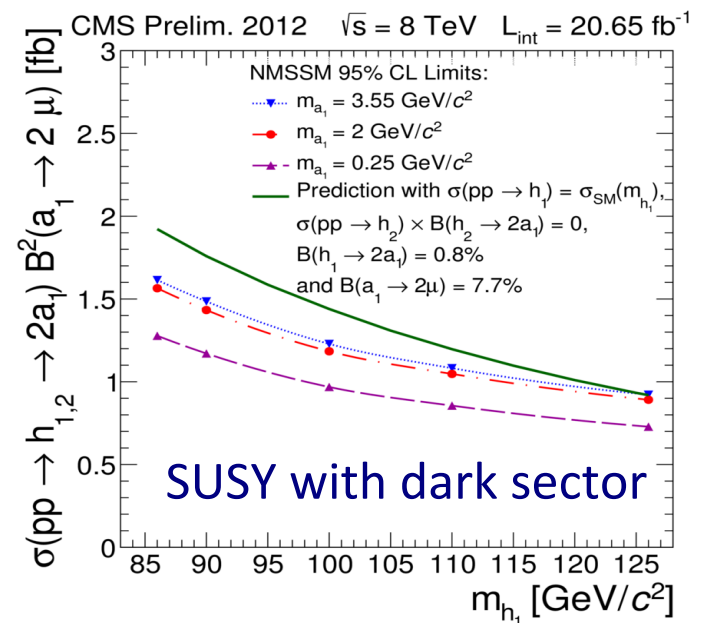
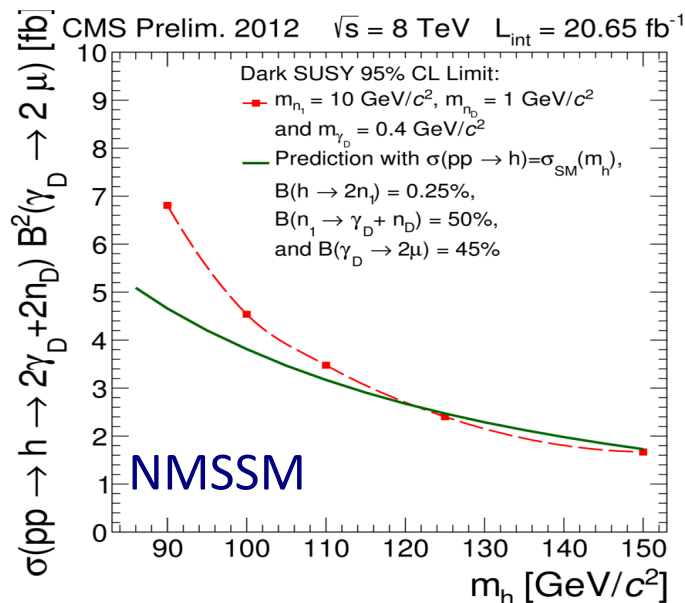
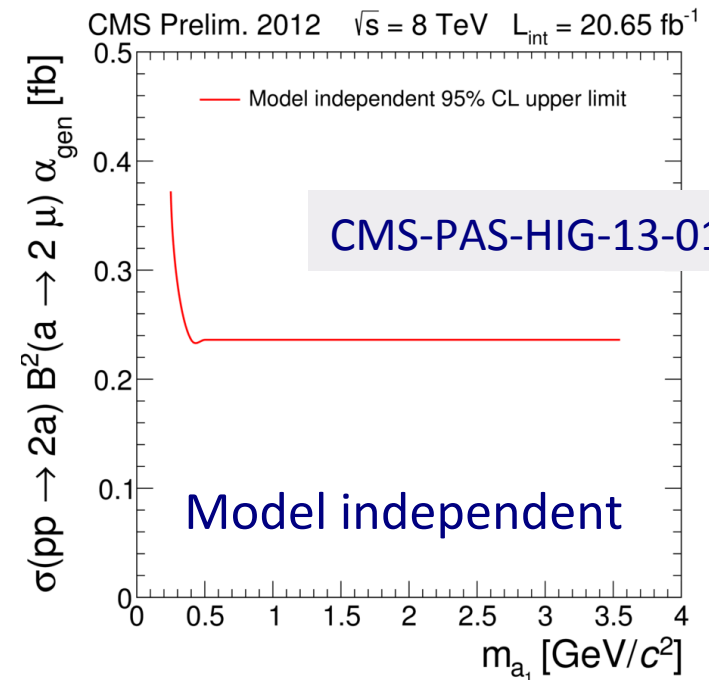
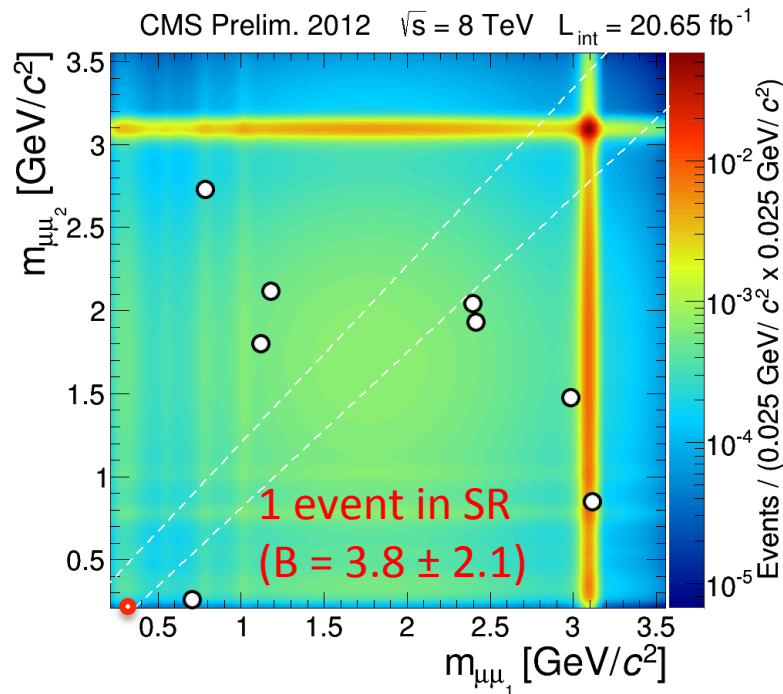
Search for decays into two very light bosons ($m < 2 m_\tau$) decaying to OS, boosted muon pairs

NMSSM

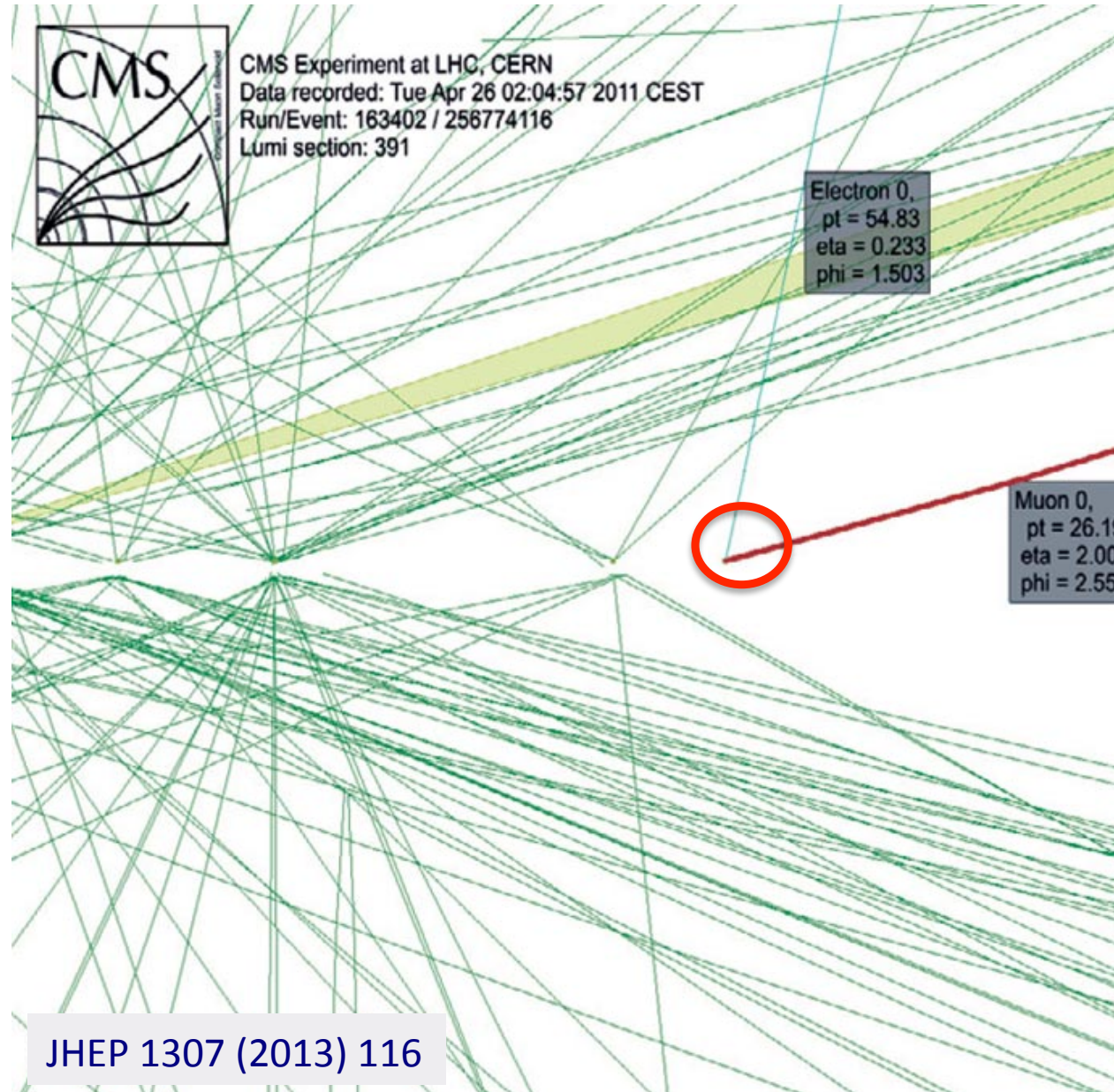
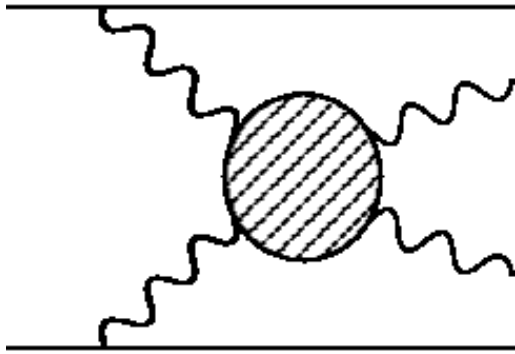
- h_1 or h_2 corresponds to 125 GeV Higgs
- a new CP-odd Higgs

Hidden-valley

- h decays to lightest neutralino, which decays to a dark fermion + a (low-mass) dark photon that weakly couples to SM fermions
- cascades of dark particle decays might lead to larger multiplicity lepton-jets



What if no additional state is seen?



Exclusive $\gamma\gamma \rightarrow WW$ production candidate

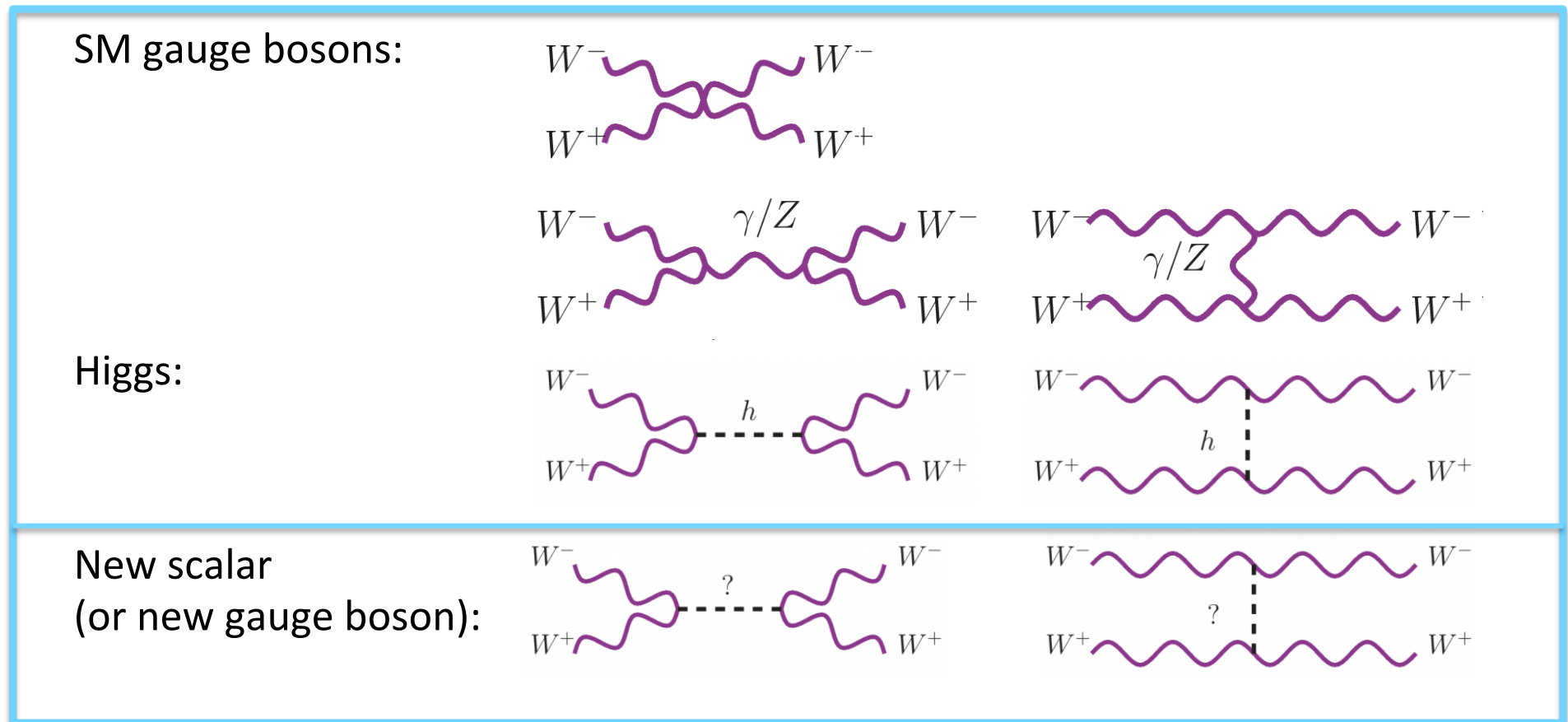
Gabriella.Pasztor@cern.ch : Beyond the Higgs

VV scattering: a probe of EWSB

Vector boson scattering is „intimately” connected to EWSB and new physics

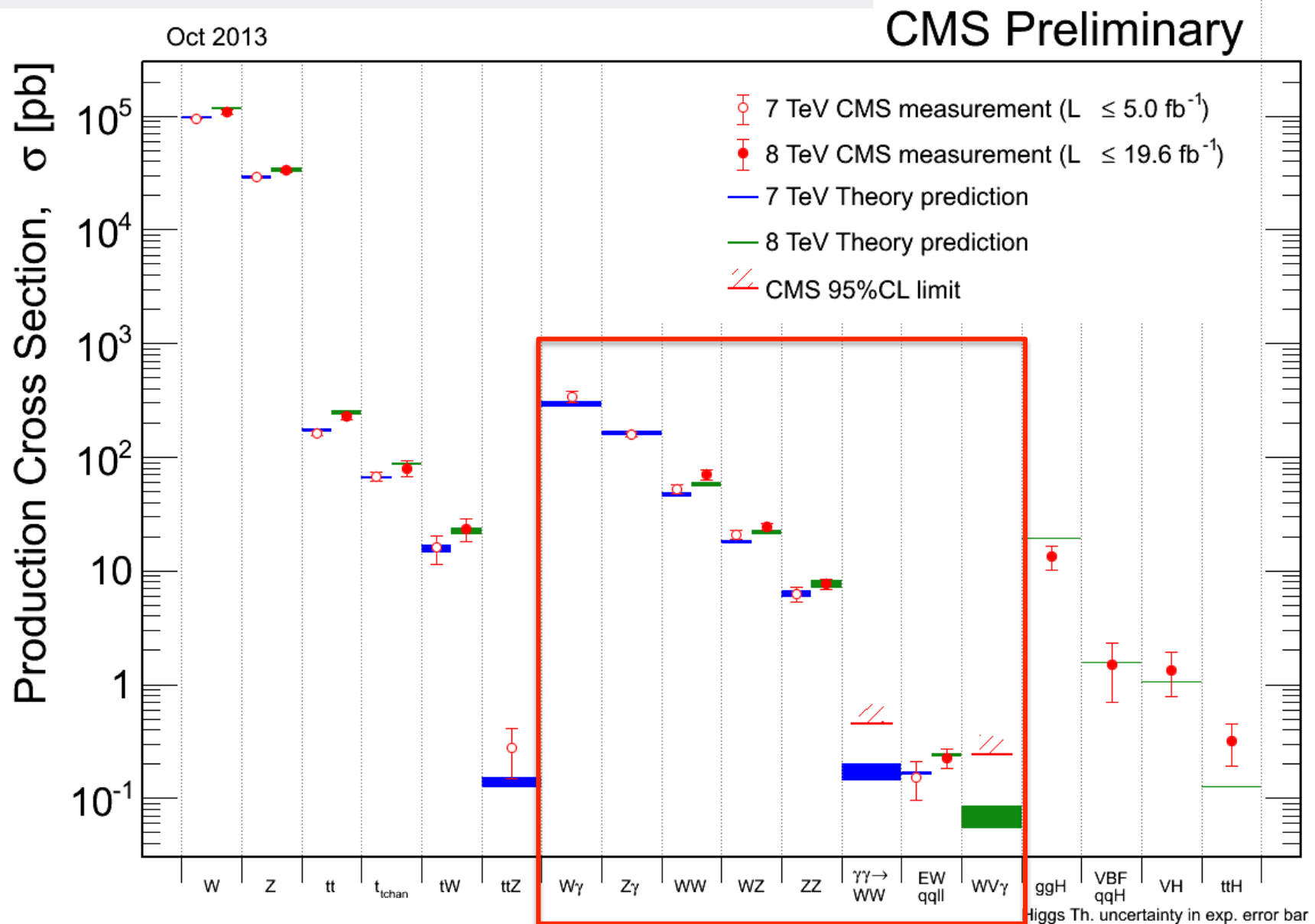
- In SM, unitarity in VV scattering is restored by Higgs exchange: $\sigma \sim O(E^2) - O(E^2) \rightarrow O(E^0)$
- If HVV coupling is not exactly the SM value, unitarity is not realized [$\sigma \sim O(E^2)$] or „delayed” until a new high-mass state enters

Even if no new physics is observed directly (finite energy reach, large backgrounds), VV scattering can reveal its existence

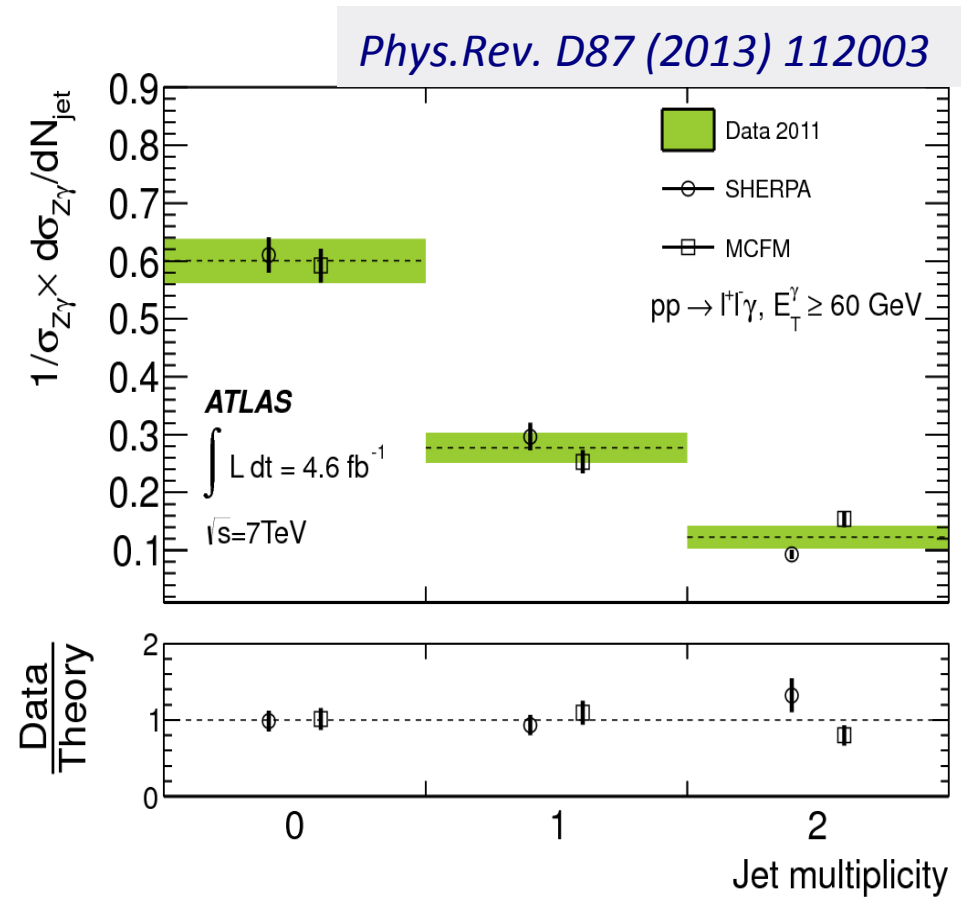
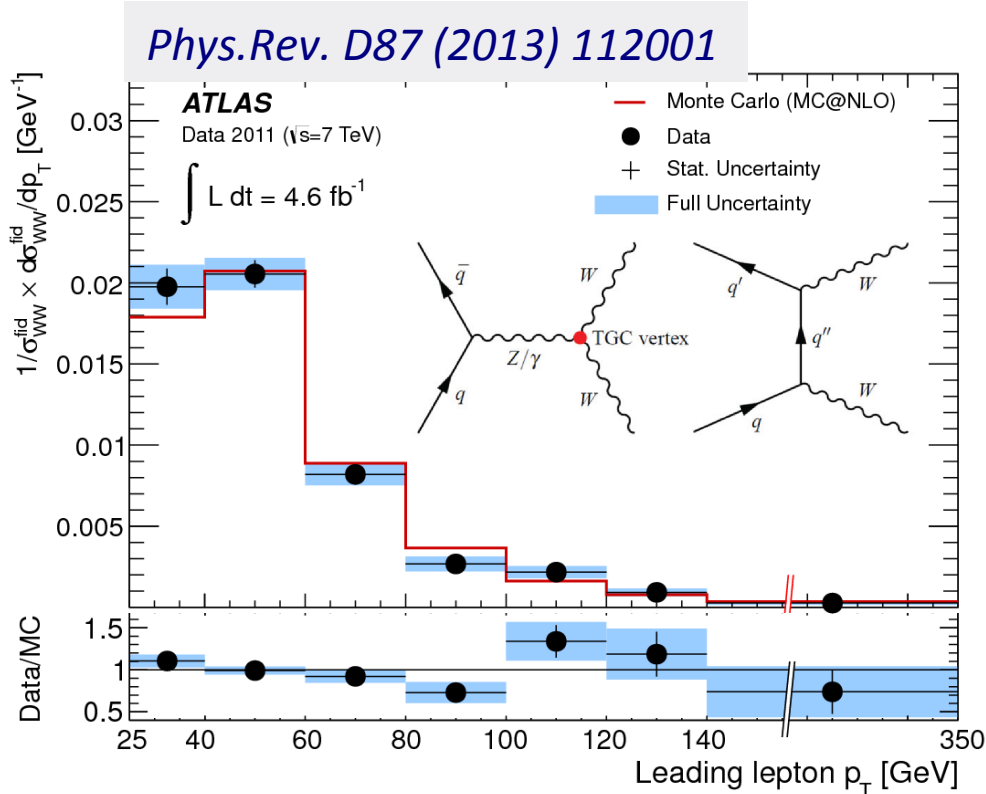


State-of-the-art in multi-boson studies

<https://twiki.cern.ch/twiki/bin/view/CMSPublic/PhysicsResultsSMP>

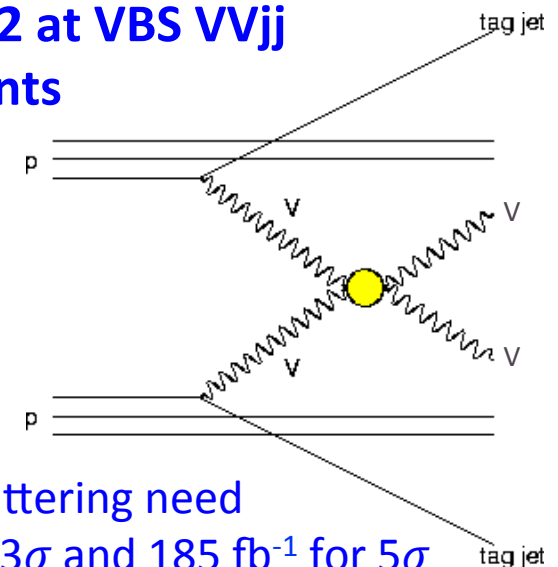


Diboson measurements going differential...



- Jet multiplicity distribution only available for $l\bar{l}\gamma$ and $l\nu\gamma$ measurements
- Especial interesting for inclusive $VVjj$ as a step towards vector boson scattering measurements, where forward jets tag the event

Aim in Run 2 at VBS $VVjj$ measurements

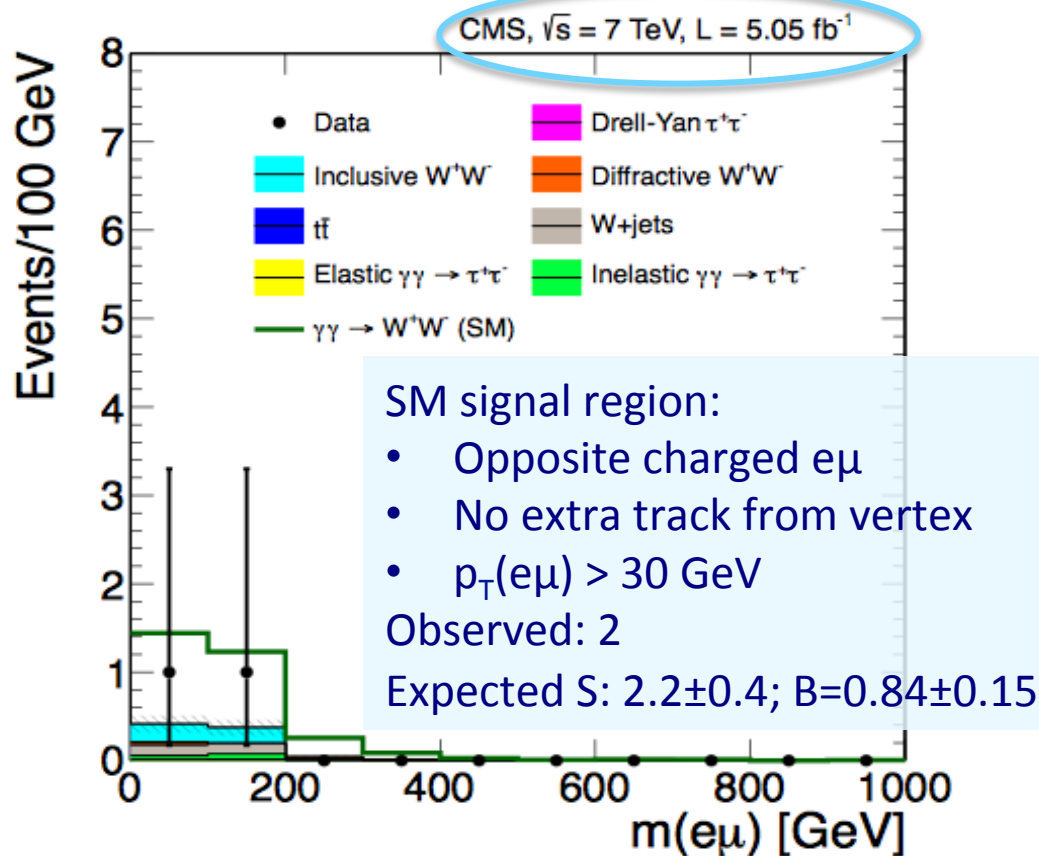
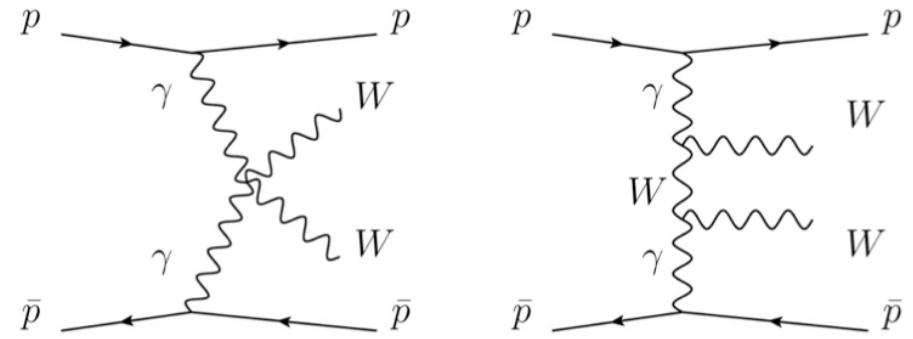


SM WZ scattering need

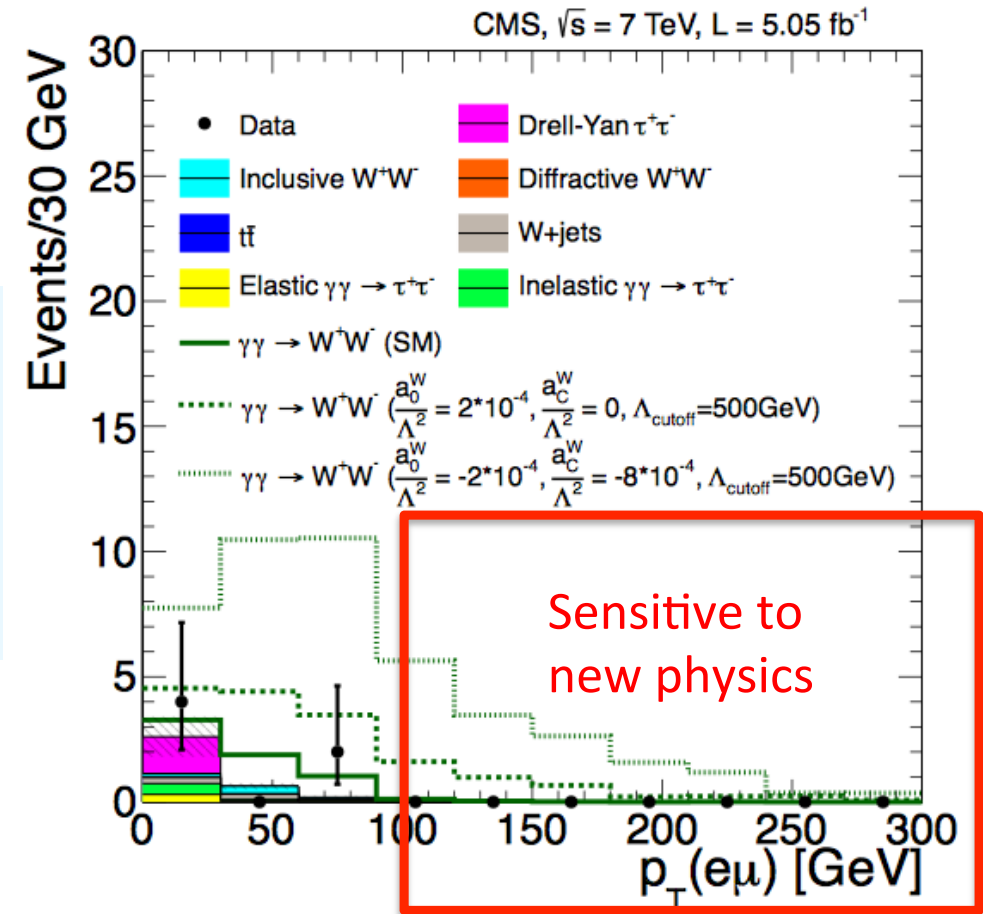
75 fb^{-1} for 3σ and 185 fb^{-1} for 5σ

Exclusive $\gamma\gamma \rightarrow WW \rightarrow e\nu\mu\nu$

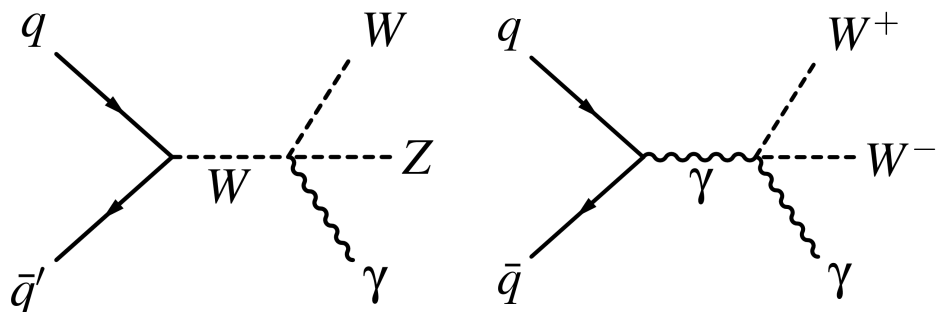
- Access $\gamma\gamma WW$ quartic coupling
- Statistics limited measurement
- High $p_T(e\mu)$ sensitive to anomalous couplings from new physics



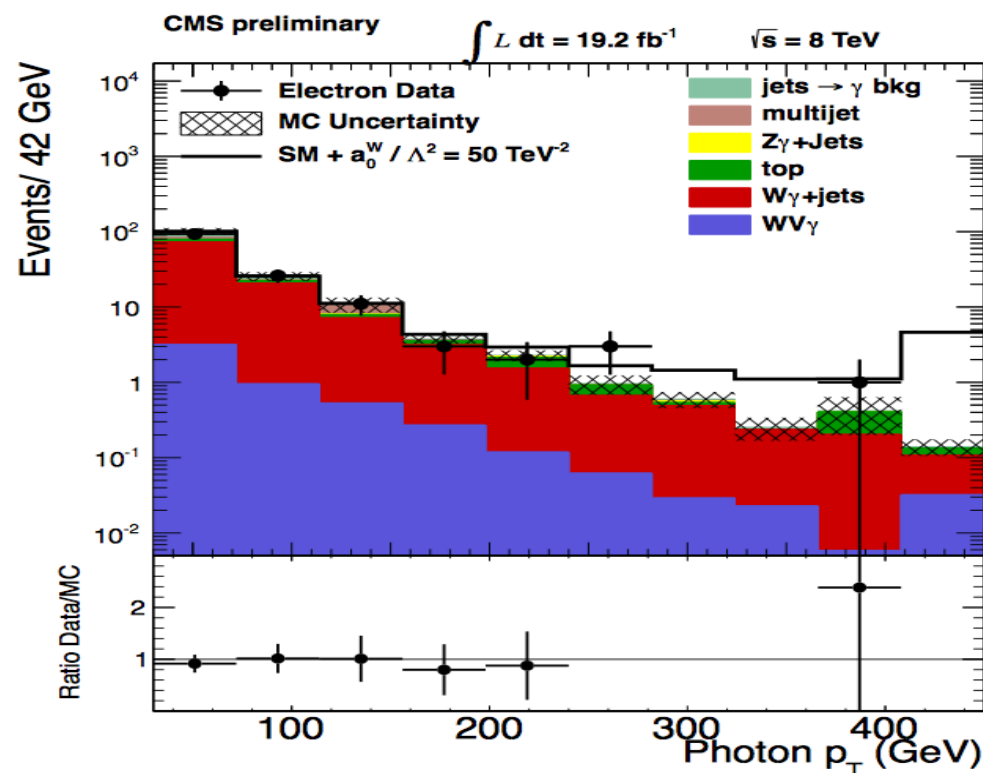
$\sigma = 2.2^{+3.3}_{-2.0}$ fb (< 10.6 fb @ 95%CL)
 $\sigma_{SM} = 4.0 \pm 0.7$ fb



Tri-boson production



- Leptonic W decay $\rightarrow e, \mu$
- Hadronic W/Z $\rightarrow 2$ jets
- Photon with $p_T > 10$ GeV



Process	muon channel number of events	electron channel number of events
W γ +jets	$136.9 \pm 3.5 \pm 9.2 \pm 0.0$	$101.6 \pm 2.9 \pm 8.0 \pm 0.0$
WV+jet, jet $\rightarrow \gamma$	$33.1 \pm 1.3 \pm 4.6 \pm 0.0$	$21.3 \pm 1.0 \pm 3.1 \pm 0.0$
MC $t\bar{t}\gamma$	$12.5 \pm 0.8 \pm 2.9 \pm 0.5$	$9.1 \pm 0.7 \pm 2.1 \pm 0.4$
MC single top	$2.8 \pm 0.8 \pm 0.2 \pm 0.1$	$1.7 \pm 0.6 \pm 0.1 \pm 0.1$
MC Z γ +jets	$1.7 \pm 0.1 \pm 0.1 \pm 0.1$	$1.5 \pm 0.1 \pm 0.1 \pm 0.1$
multijets	$<0.2 \pm 0.0 \pm 0.1 \pm 0.0$	$7.2 \pm 3.6 \pm 3.6 \pm 0.0$
SM WW γ	$6.3 \pm 0.1 \pm 1.5 \pm 0.3$	$4.7 \pm 0.1 \pm 1.1 \pm 0.2$
SM WZ γ	$0.6 \pm 0.0 \pm 0.1 \pm 0.0$	$0.5 \pm 0.0 \pm 0.1 \pm 0.0$
Total predicted	$193.9 \pm 3.9 \pm 10.8 \pm 1.0$	$147.6 \pm 4.8 \pm 9.6 \pm 0.7$
Data	183	139

Cross-section not yet accessible

Upper limit: 241 fb
(3.4 x SM)

CMS-PAS-SMP-13-009

Anomalous quartic gauge coupling limits

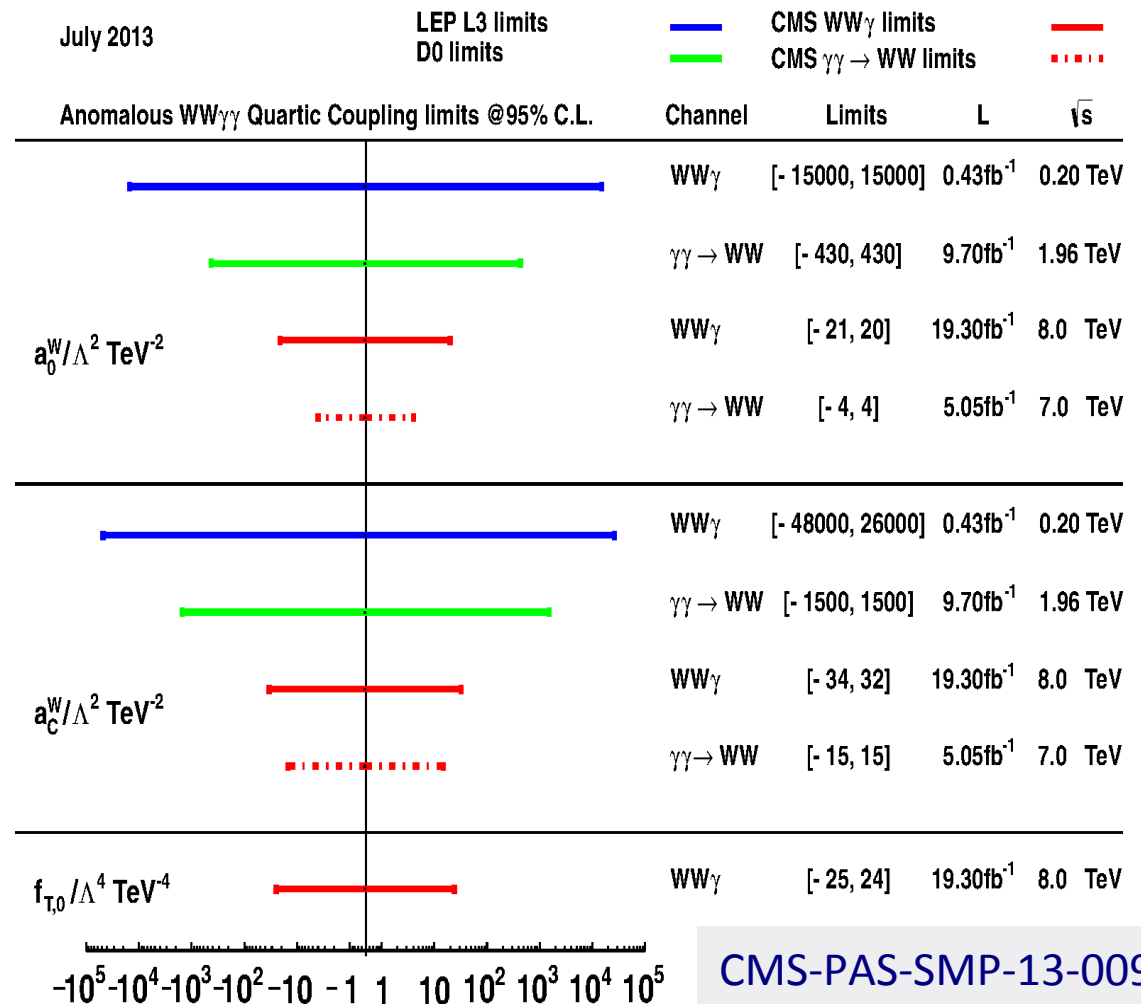
- EFT Lagrangian parametrizing possible new physics

$$L = L_{SM} + \sum_d \sum_i \frac{C_i^{(d)}}{\Lambda^{d-4}} O_i^{(d)}$$

- Dim-6 and dim-8 operators by integrating out the new degrees of freedom
- Coefficients to be calculated from more complete high-energy theory
- Unitarity violation
→ form-factors, cut-off
- Large theoretical uncertainties

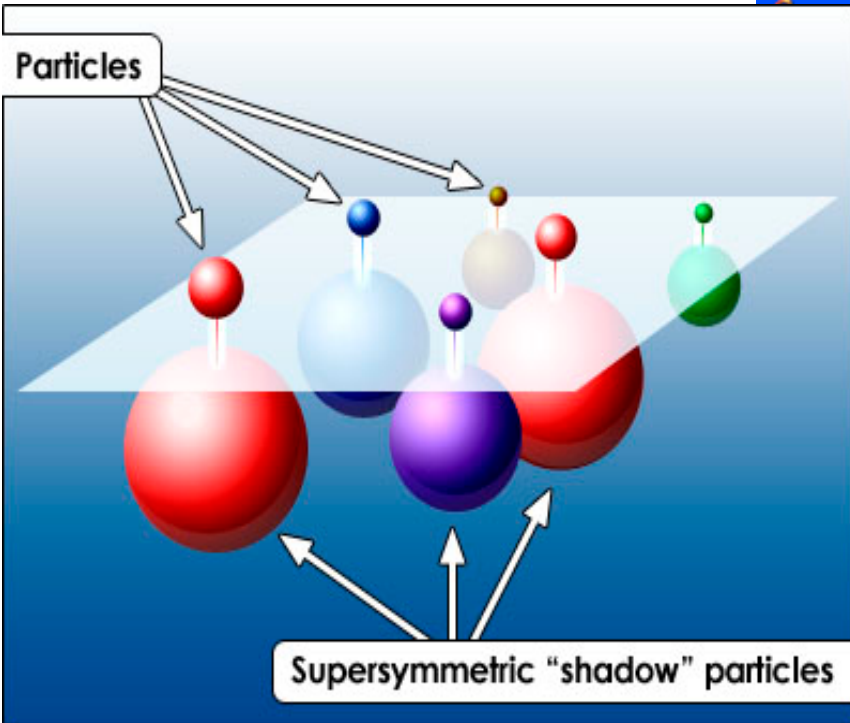
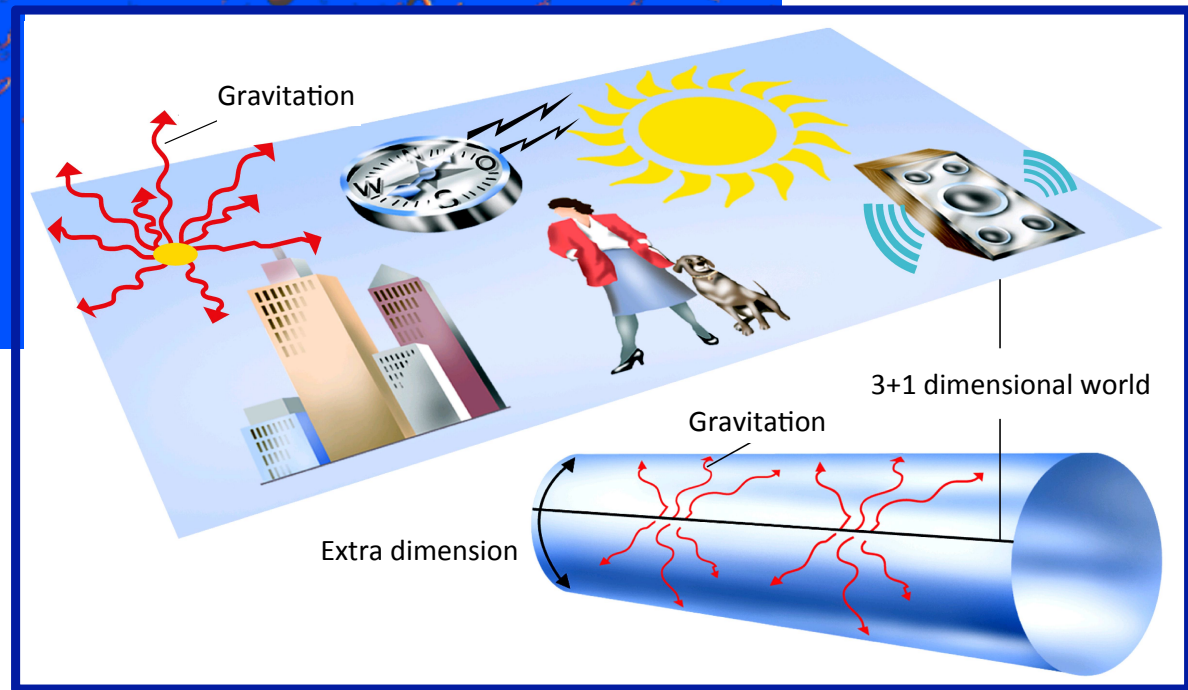
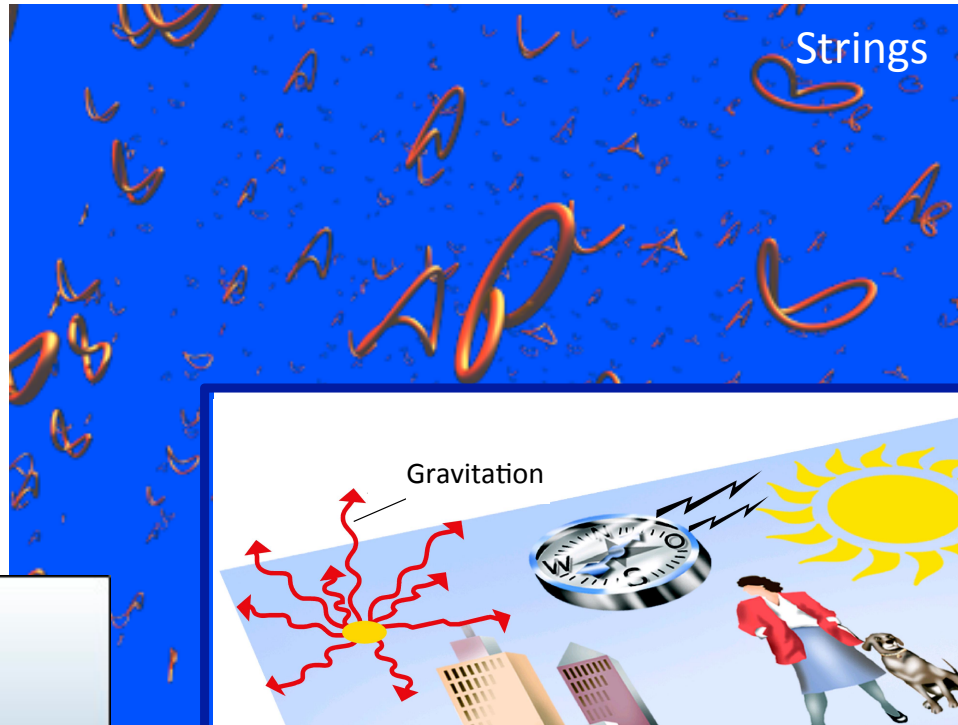
$$L_6^0 = \frac{e^2 a_0^W}{8 \Lambda^2} F_{\mu\nu} F^{\mu\nu} W^{+\alpha} W_{\alpha}^{-} - \frac{e^2}{16 \cos^2 \Theta_W} \frac{a_0^Z}{\Lambda^2} F_{\mu\nu} F^{\mu\nu} Z^{\alpha} Z_{\alpha}$$

$$L_6^C = \frac{-e^2 a_C^W}{16 \Lambda^2} F_{\mu\alpha} F^{\mu\beta} (W^{+\alpha} W_{\beta}^{-} + W^{-\alpha} W_{\beta}^{+}) - \frac{e^2}{16 \cos^2 \Theta_W} \frac{a_C^Z}{\Lambda^2} F_{\mu\alpha} F^{\mu\beta} Z^{\alpha} Z_{\beta}$$



CMS-PAS-SMP-13-009

Beyond the Higgs sector



ATLAS SUSY Searches* - 95% CL Lower Limits

Status: SUSY 2013

ATLAS Preliminary

$\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$ $\sqrt{s} = 7, 8 \text{ TeV}$

Model	e, μ, τ, γ	Jets	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Mass limit	Reference		
Inclusive Searches	MSUGRA/CMSSM	0	2-6 jets	Yes	20.3	\tilde{q}, \tilde{g} 1.7 TeV	$m(\tilde{q})=m(\tilde{g})$	ATLAS-CONF-2013-047
	MSUGRA/CMSSM	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.2 TeV	any $m(\tilde{q})$	ATLAS-CONF-2013-062
	MSUGRA/CMSSM	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	any $m(\tilde{q})$	1308.1841
	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{q} 740 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-047
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{\chi}_1^0$	0	2-6 jets	Yes	20.3	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-047
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{\chi}_1^0 \rightarrow qqW^\pm \tilde{\chi}_1^0$	1 e, μ	3-6 jets	Yes	20.3	\tilde{g} 1.18 TeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$	ATLAS-CONF-2013-062
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$	2 e, μ	0-3 jets	-	20.3	\tilde{g} 1.12 TeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-089
	GMSB ($\tilde{\ell}$ NLSP)	2 e, μ	2-4 jets	Yes	4.7	\tilde{g} 1.24 TeV	$\tan\beta < 15$	1208.4688
	GMSB ($\tilde{\ell}$ NLSP)	1-2 τ	0-2 jets	Yes	20.7	\tilde{g} 1.4 TeV	$\tan\beta > 18$	ATLAS-CONF-2013-026
	GGM (bino NLSP)	2 γ	-	Yes	4.8	\tilde{g} 1.07 TeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$	1209.0753
GGM (wino NLSP)	1 $e, \mu + \gamma$	-	Yes	4.8	\tilde{g} 619 GeV	$m(\tilde{\chi}_1^0) > 50 \text{ GeV}$	ATLAS-CONF-2012-144	
GGM (higgsino-bino NLSP)	γ	1 b	Yes	4.8	\tilde{g} 900 GeV	$m(\tilde{\chi}_1^0) > 220 \text{ GeV}$	1211.1167	
GGM (higgsino NLSP)	2 $e, \mu (Z)$	0-3 jets	Yes	5.8	\tilde{g} 690 GeV	$m(\tilde{H}) > 200 \text{ GeV}$	ATLAS-CONF-2012-152	
Gravitino LSP	0	mono-jet	Yes	10.5	$F^{1/2}$ scale 645 GeV	$m(\tilde{g}) > 10^{-4} \text{ eV}$	ATLAS-CONF-2012-147	
3 rd gen. \tilde{g} med.	$\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$	0	3 b	Yes	20.1	\tilde{g} 1.2 TeV	$m(\tilde{\chi}_1^0) < 600 \text{ GeV}$	ATLAS-CONF-2013-061
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$	0	7-10 jets	Yes	20.3	\tilde{g} 1.1 TeV	$m(\tilde{\chi}_1^0) < 350 \text{ GeV}$	1308.1841
	$\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^\pm$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.34 TeV	$m(\tilde{\chi}_1^0) < 400 \text{ GeV}$	ATLAS-CONF-2013-061
	$\tilde{g} \rightarrow b\tilde{t}\tilde{\chi}_1^\pm$	0-1 e, μ	3 b	Yes	20.1	\tilde{g} 1.3 TeV	$m(\tilde{\chi}_1^0) < 300 \text{ GeV}$	ATLAS-CONF-2013-061
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$	0	2 b	Yes	20.1	\tilde{b}_1 100-620 GeV	$m(\tilde{\chi}_1^0) < 90 \text{ GeV}$	1308.2631
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^0$	2 e, μ (SS)	0-3 b	Yes	20.7	\tilde{b}_1 275-430 GeV	$m(\tilde{\chi}_1^0)=2 m(\tilde{\chi}_2^0)$	ATLAS-CONF-2013-007
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	1-2 e, μ	1-2 b	Yes	4.7	\tilde{t}_1 110-167 GeV	$m(\tilde{\chi}_1^0)=55 \text{ GeV}$	1208.4305, 1209.2102
	$\tilde{t}_1\tilde{t}_1$ (light), $\tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$	2 e, μ	0-2 jets	Yes	20.3	\tilde{t}_1 130-220 GeV	$m(\tilde{\chi}_1^0)=m(\tilde{t}_1)-m(W)-50 \text{ GeV}, m(\tilde{t}_1) < m(\tilde{\chi}_1^\pm)$	ATLAS-CONF-2013-048
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	2 e, μ	2 jets	Yes	20.3	\tilde{t}_1 225-525 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-065
	$\tilde{t}_1\tilde{t}_1$ (medium), $\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$	0	2 b	Yes	20.1	\tilde{t}_1 150-580 GeV	$m(\tilde{\chi}_1^0) < 200 \text{ GeV}, m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=5 \text{ GeV}$	1308.2631
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	1 e, μ	1 b	Yes	20.7	\tilde{t}_1 200-610 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-037
	$\tilde{t}_1\tilde{t}_1$ (heavy), $\tilde{t}_1 \rightarrow t\tilde{\chi}_1^\pm$	0	2 b	Yes	20.5	\tilde{t}_1 320-660 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-024
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	0	mono-jet/c-tag	Yes	20.3	\tilde{t}_1 90-200 GeV	$m(\tilde{t}_1)-m(\tilde{\chi}_1^0) < 85 \text{ GeV}$	ATLAS-CONF-2013-068
	$\tilde{t}_1\tilde{t}_1$ (natural GMSB)	2 $e, \mu (Z)$	1 b	Yes	20.7	\tilde{t}_1 500 GeV	$m(\tilde{\chi}_1^0) > 150 \text{ GeV}$	ATLAS-CONF-2013-025
$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 $e, \mu (Z)$	1 b	Yes	20.7	\tilde{t}_2 271-520 GeV	$m(\tilde{t}_1)=m(\tilde{\chi}_1^\pm)+180 \text{ GeV}$	ATLAS-CONF-2013-025	
EW direct	$\tilde{\ell}_L, \tilde{R}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ	0	Yes	20.3	$\tilde{\ell}$ 85-315 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}$	ATLAS-CONF-2013-049
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\ell}\nu(\ell\bar{\nu})$	2 e, μ	0	Yes	20.3	$\tilde{\chi}_1^\pm$ 125-450 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	ATLAS-CONF-2013-049
	$\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\tau}\nu(\tau\bar{\nu})$	2 τ	0	Yes	20.7	$\tilde{\chi}_1^\pm$ 180-330 GeV	$m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\tau}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	ATLAS-CONF-2013-028
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow \tilde{\ell}\nu\tilde{\ell}\ell(\tilde{\nu}\bar{\nu}), \tilde{\ell}\bar{\nu}\tilde{\ell}\ell(\tilde{\nu}\bar{\nu})$	3 e, μ	0	Yes	20.7	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 600 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \bar{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$	ATLAS-CONF-2013-035
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 Z\tilde{\chi}_1^0$	3 e, μ	0	Yes	20.7	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 315 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0$, sleptons decoupled	ATLAS-CONF-2013-035
	$\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0$	1 e, μ	2 b	Yes	20.3	$\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 285 GeV	$m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0$, sleptons decoupled	ATLAS-CONF-2013-093
Long-lived particles	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	Yes	20.3	$\tilde{\chi}_1^\pm$ 270 GeV	$m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)=0.2 \text{ ns}$	ATLAS-CONF-2013-069
	Stable, stopped \tilde{g} R-hadron	0	1-5 jets	Yes	22.9	\tilde{g} 832 GeV	$m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s} < \tau(\tilde{g}) < 1000 \text{ s}$	ATLAS-CONF-2013-057
	GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e, \mu)$	1-2 μ	-	-	15.9	$\tilde{\chi}_1^0$ 475 GeV	$10 < \tan\beta < 50$	ATLAS-CONF-2013-058
	GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$	2 γ	-	Yes	4.7	$\tilde{\chi}_1^0$ 230 GeV	$0.4 < \tau(\tilde{\chi}_1^0) < 2 \text{ ns}$	1304.6310
$\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow qq\mu$ (RPV)	1 μ , displ. vtx	-	-	20.3	\tilde{q} 1.0 TeV	$1.5 < c\tau < 156 \text{ mm}, BR(\mu)=1, m(\tilde{\chi}_1^0)=108 \text{ GeV}$	ATLAS-CONF-2013-092	
RPV	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$	2 e, μ	-	-	4.6	$\tilde{\nu}_\tau$ 1.61 TeV	$\lambda'_{311}=0.10, \lambda_{132}=0.05$	1212.1272
	LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$	1 $e, \mu + \tau$	-	-	4.6	$\tilde{\nu}_\tau$ 1.1 TeV	$\lambda'_{311}=0.10, \lambda_{1(2)33}=0.05$	1212.1272
	Billinear RPV CMSSM	1 e, μ	7 jets	Yes	4.7	\tilde{q}, \tilde{g} 1.2 TeV	$m(\tilde{q})=m(\tilde{g}), c\tau_{LSP} < 1 \text{ mm}$	ATLAS-CONF-2012-140
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$	4 e, μ	-	Yes	20.7	$\tilde{\chi}_1^\pm$ 760 GeV	$m(\tilde{\chi}_1^0) > 300 \text{ GeV}, \lambda_{121} > 0$	ATLAS-CONF-2013-036
	$\tilde{\chi}_1^+\tilde{\chi}_1^-, \tilde{\chi}_1^0 \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$	3 $e, \mu + \tau$	-	Yes	20.7	$\tilde{\chi}_1^\pm$ 350 GeV	$m(\tilde{\chi}_1^0) > 80 \text{ GeV}, \lambda_{133} > 0$	ATLAS-CONF-2013-036
	$\tilde{g} \rightarrow qq\tilde{q}$	0	6-7 jets	-	20.3	\tilde{g} 916 GeV	$BR(t)=BR(b)=BR(c)=0\%$	ATLAS-CONF-2013-091
$\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$	2 e, μ (SS)	0-3 b	Yes	20.7	\tilde{g} 880 GeV		ATLAS-CONF-2013-007	
Other	Scalar gluon pair, $sgluon \rightarrow q\bar{q}$	0	4 jets	-	4.6	sgluon 100-287 GeV	incl. limit from 1110.2693	1210.4826
	Scalar gluon pair, $sgluon \rightarrow t\bar{t}$	2 e, μ (SS)	1 b	Yes	14.3	sgluon 800 GeV		ATLAS-CONF-2013-051
	WIMP interaction (D5, Dirac χ)	0	mono-jet	Yes	10.5	M^* scale 704 GeV	$m(\chi) < 80 \text{ GeV}$, limit of $< 687 \text{ GeV}$ for D8	ATLAS-CONF-2012-147

$\sqrt{s} = 7 \text{ TeV}$
full data

$\sqrt{s} = 8 \text{ TeV}$
partial data

$\sqrt{s} = 8 \text{ TeV}$
full data

10⁻¹

1

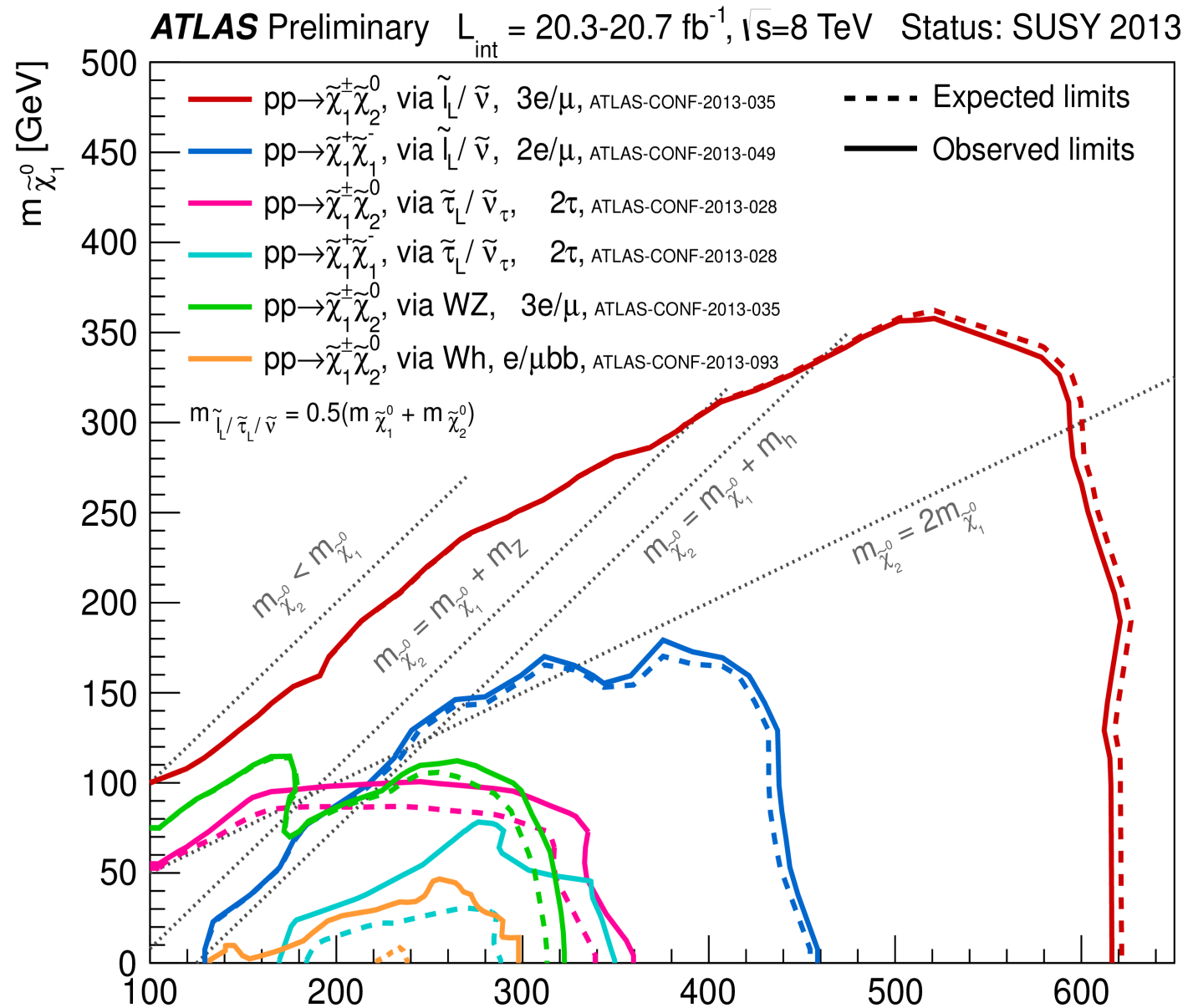
Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Constraints on gaugino production

No hint anywhere for new SUSY phenomena

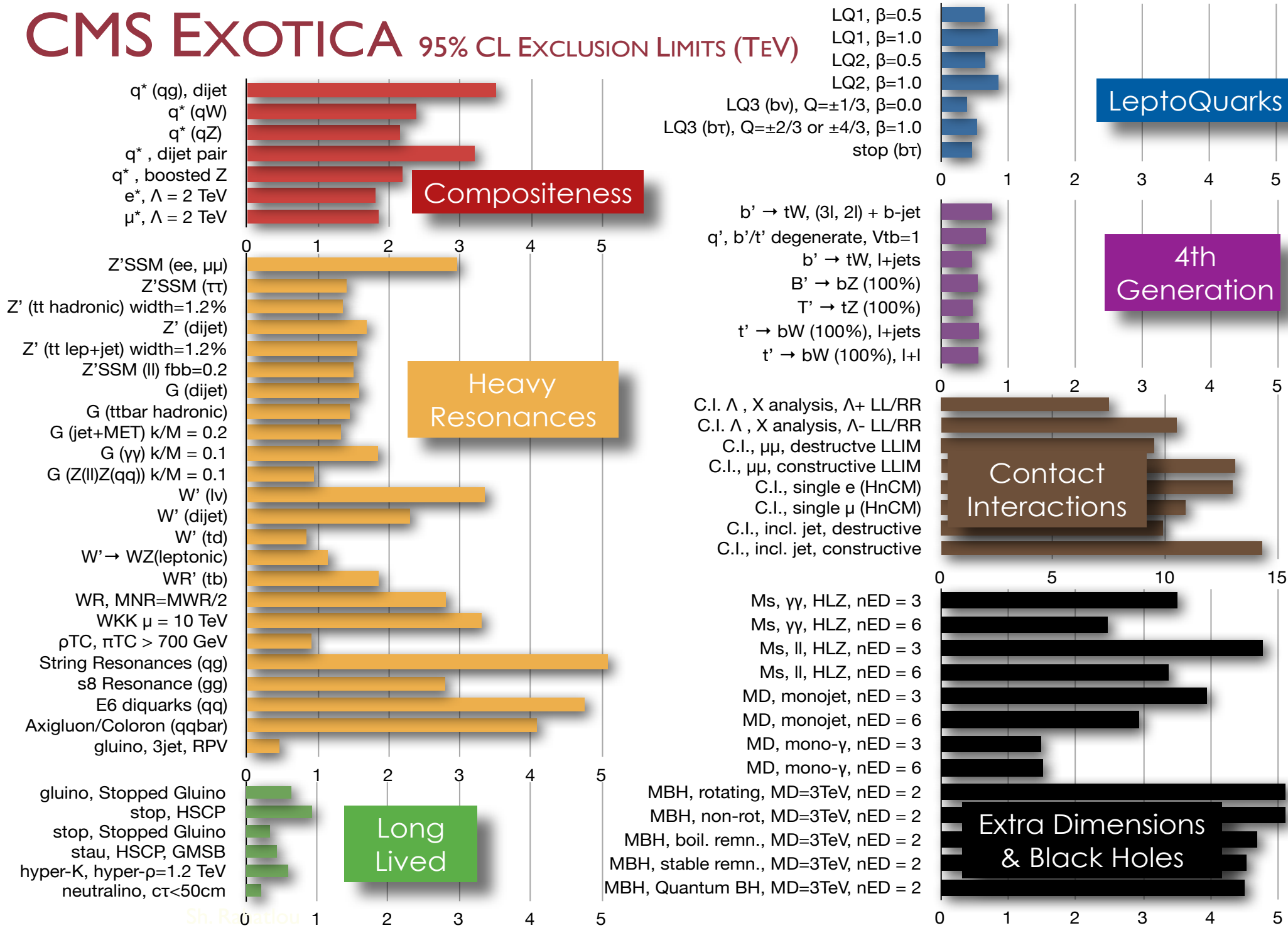
Natural-SUSY seems less and less probable (Higgs mass & direct search results)



<https://twiki.cern.ch/twiki/bin/view/AtlasPublic/CombinedSummaryPlots#SusySummary>

$m_{\tilde{\chi}_1^\pm} (=m_{\tilde{\chi}_2^0})$ [GeV]

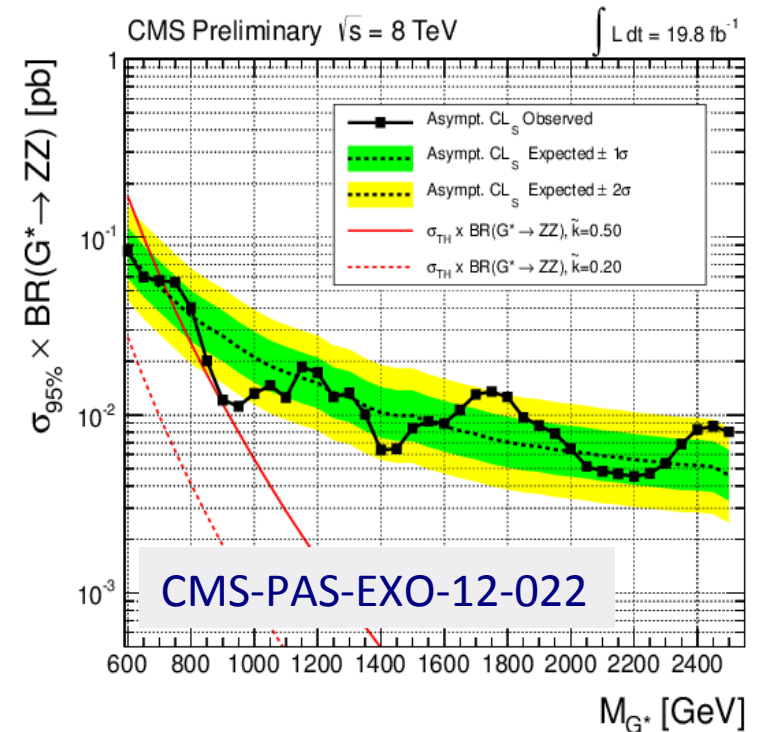
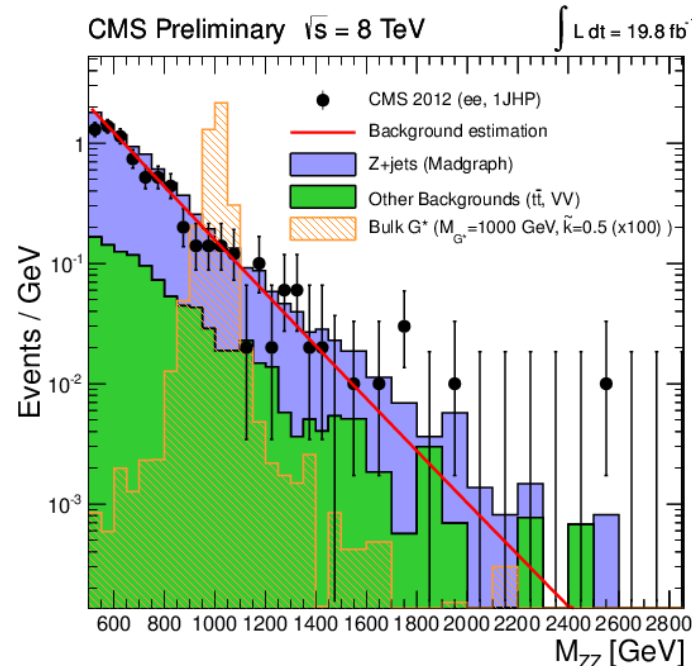
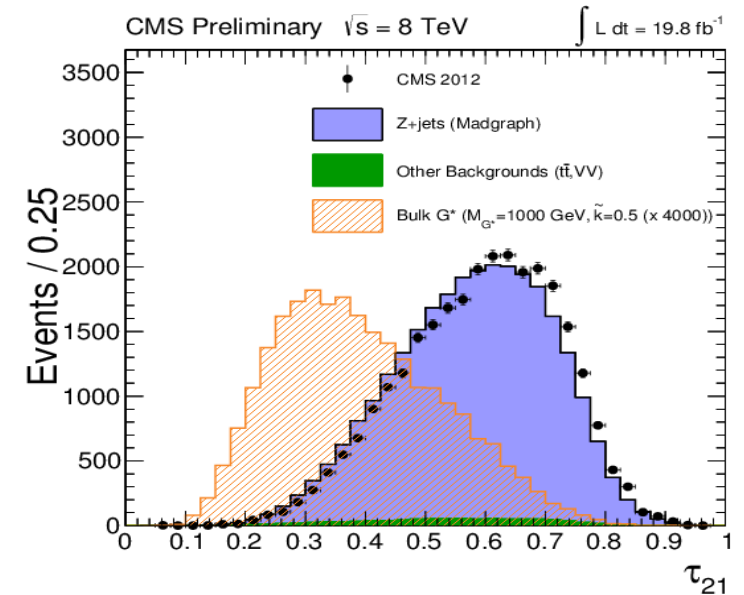
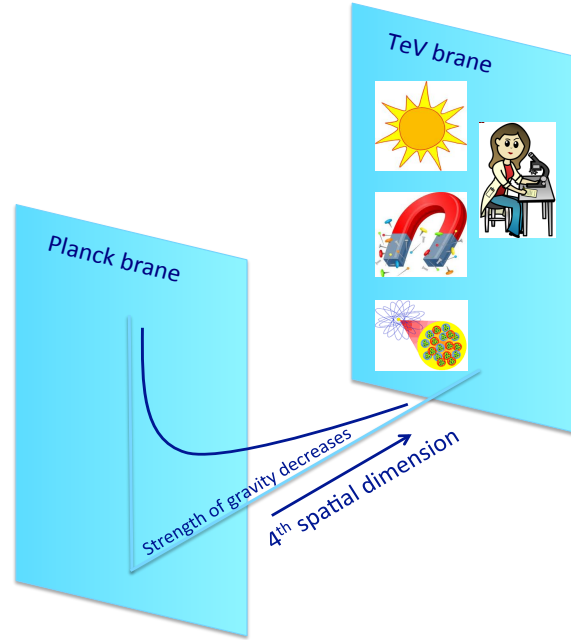
CMS EXOTICA 95% CL EXCLUSION LIMITS (TeV)



<https://twiki.cern.ch/twiki/pub/CMSPublic/PhysicsResultsEXO/CMS-EXO-Moriond2013.pdf>

Search for extra dimensions: RS graviton

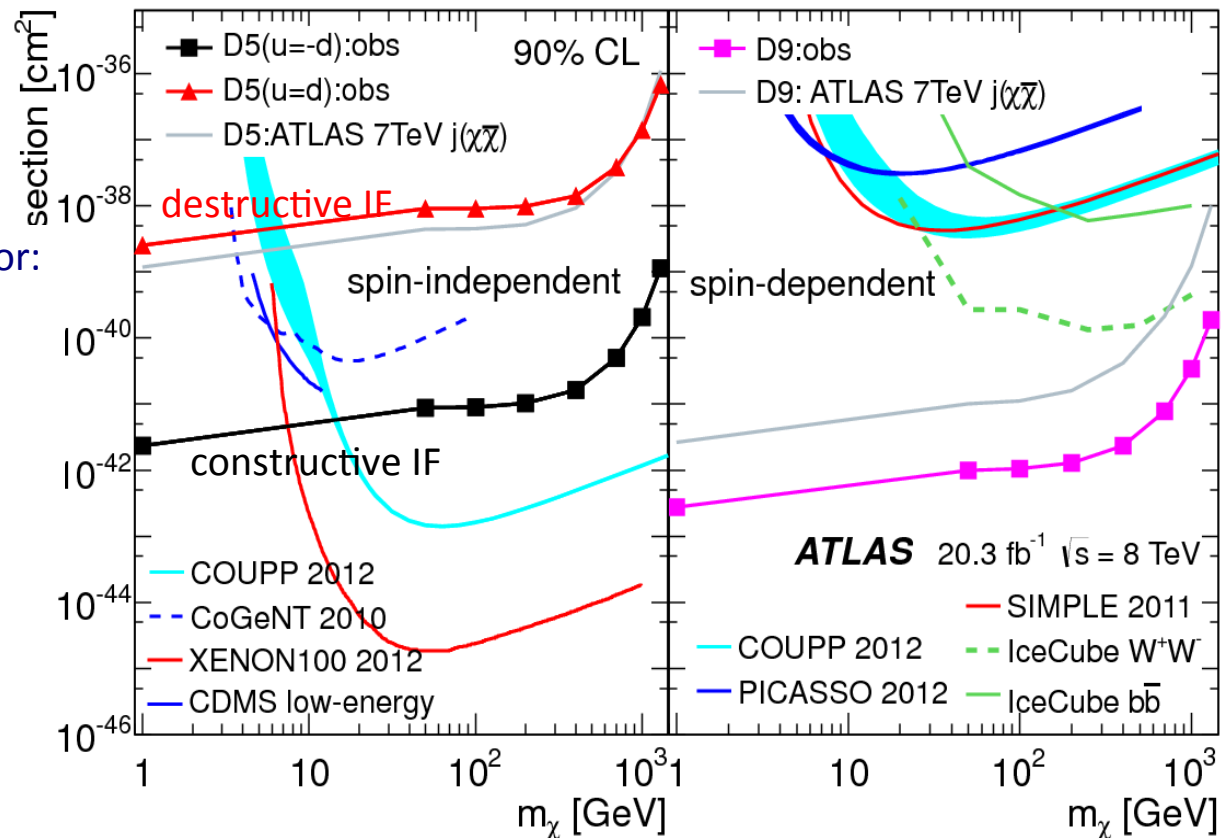
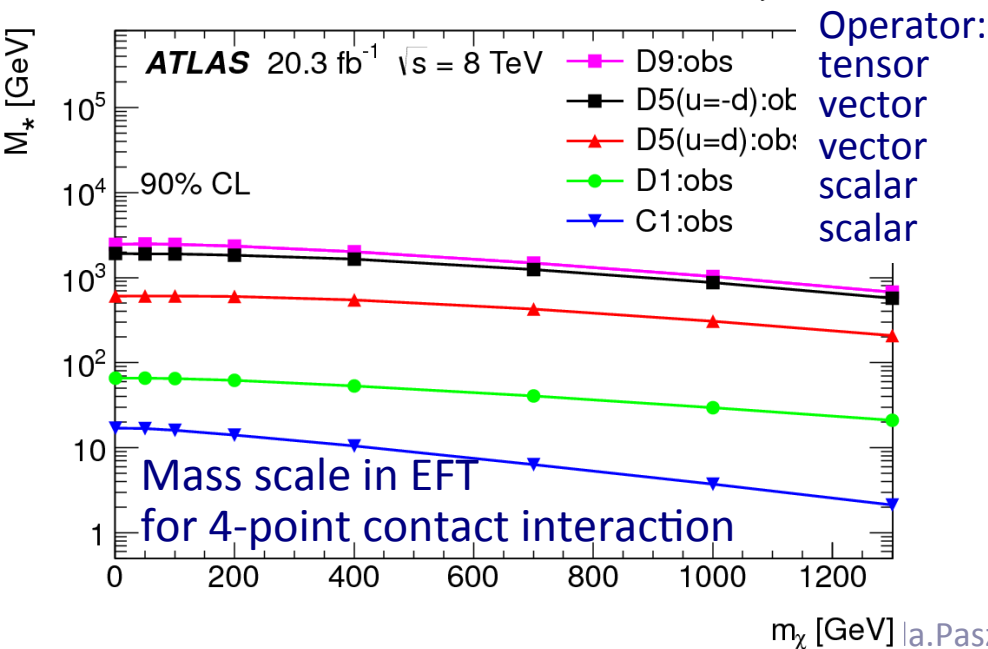
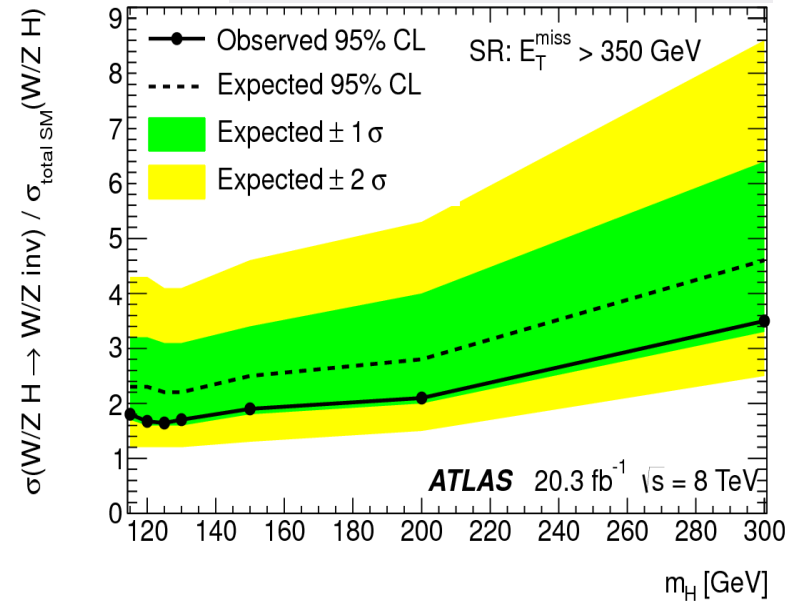
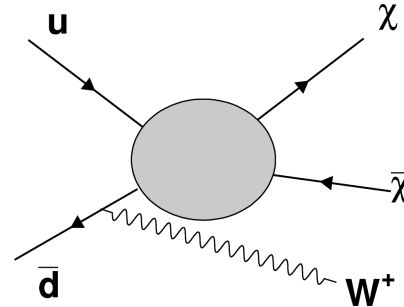
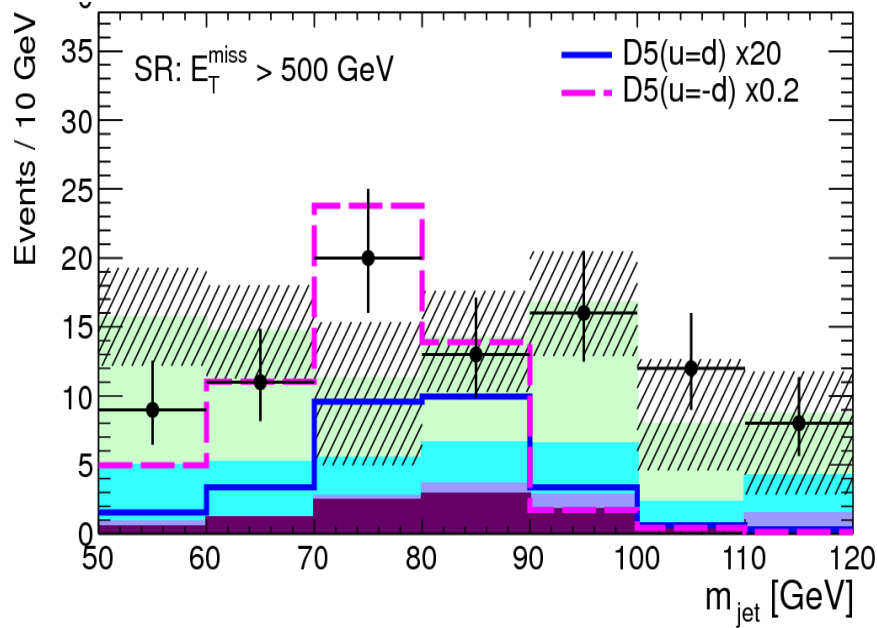
- Randall-Sundrum model proposed to solve the hierarchy problem assuming a 5th dimension with a warped geometry
- Only gravity can propagate in the ED
- In 3+1 D, KK graviton resonances appear
- Search for $G_{KK} \rightarrow ZZ$ in $qqll$ final state
- Hadronic Z boosted: jet substructure techniques



Search for dark matter

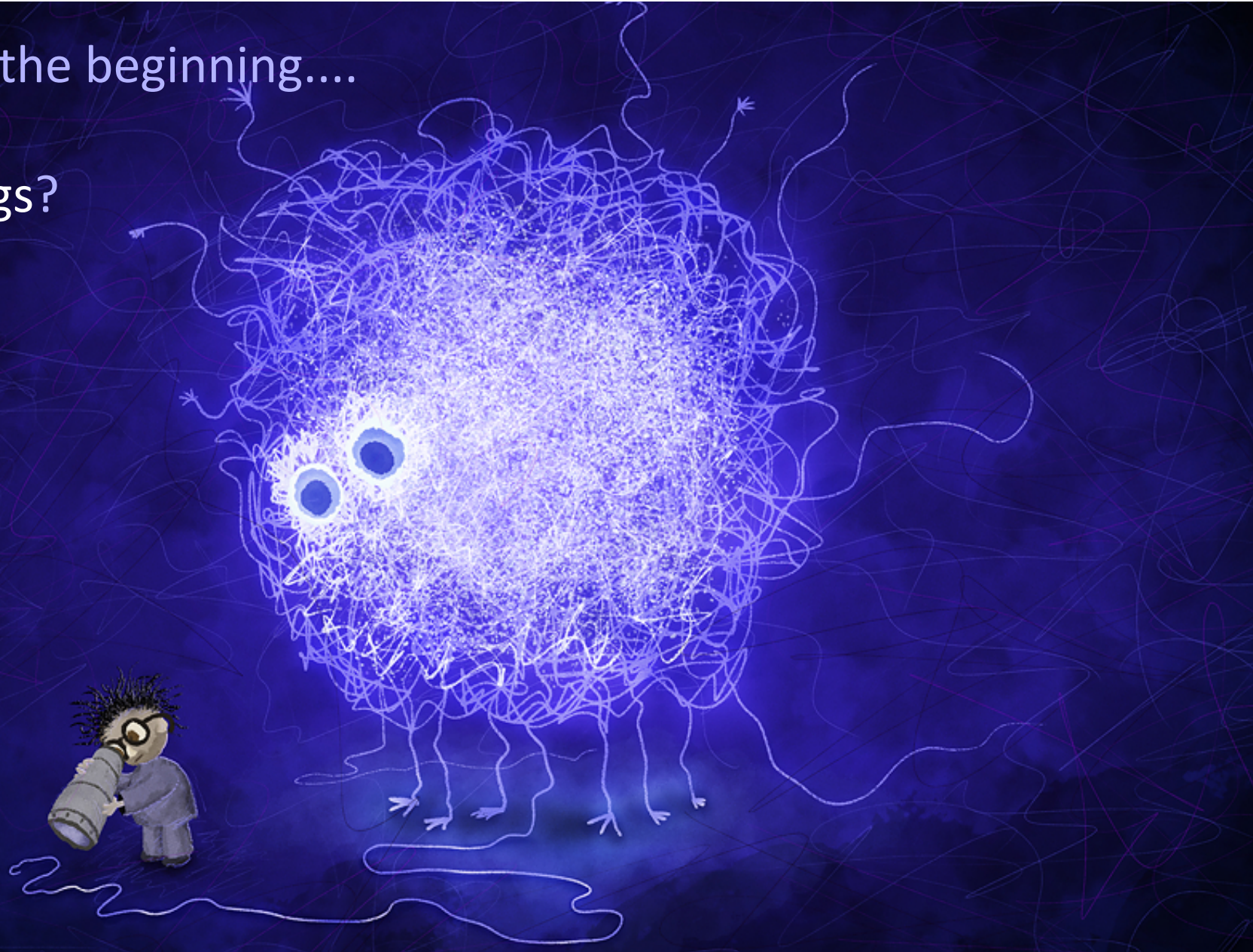
arXiv:1309.4017 [hep-ex]

Pair-produced dark matter tagged by initial-state W/Z emission in boosted hadronic final state



As exiting as ever...

We are only at the beginning....
What waits
behind the Higgs?



Extra

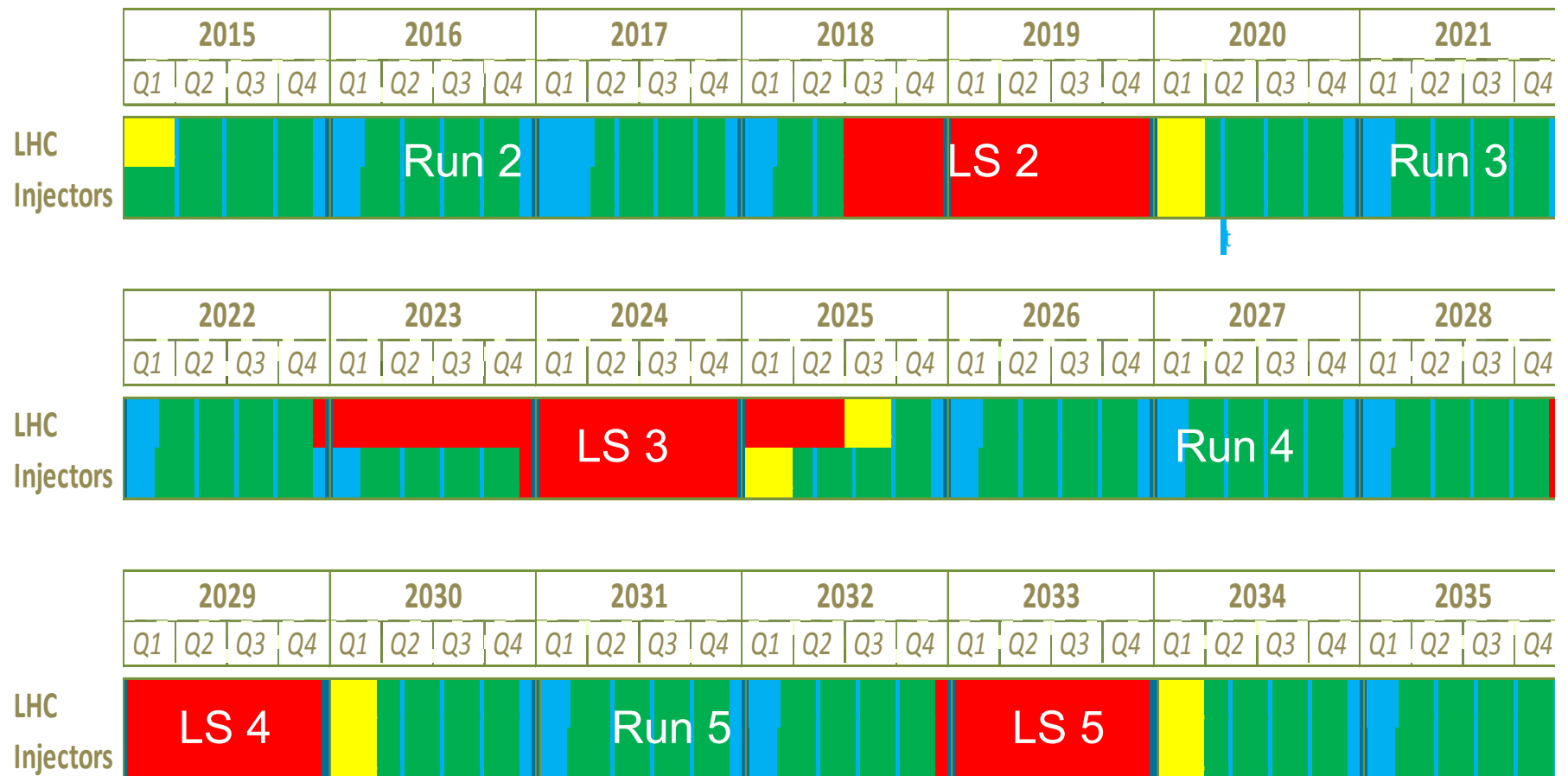
LHC schedule beyond LS1

Only EYETS (19 weeks) (no Linac4 connection during Run2)

LS2 starting in 2018 (July) 18 months + 3months BC (Beam Commissioning)

LS3 LHC: starting in 2023 => 30 months + 3 BC

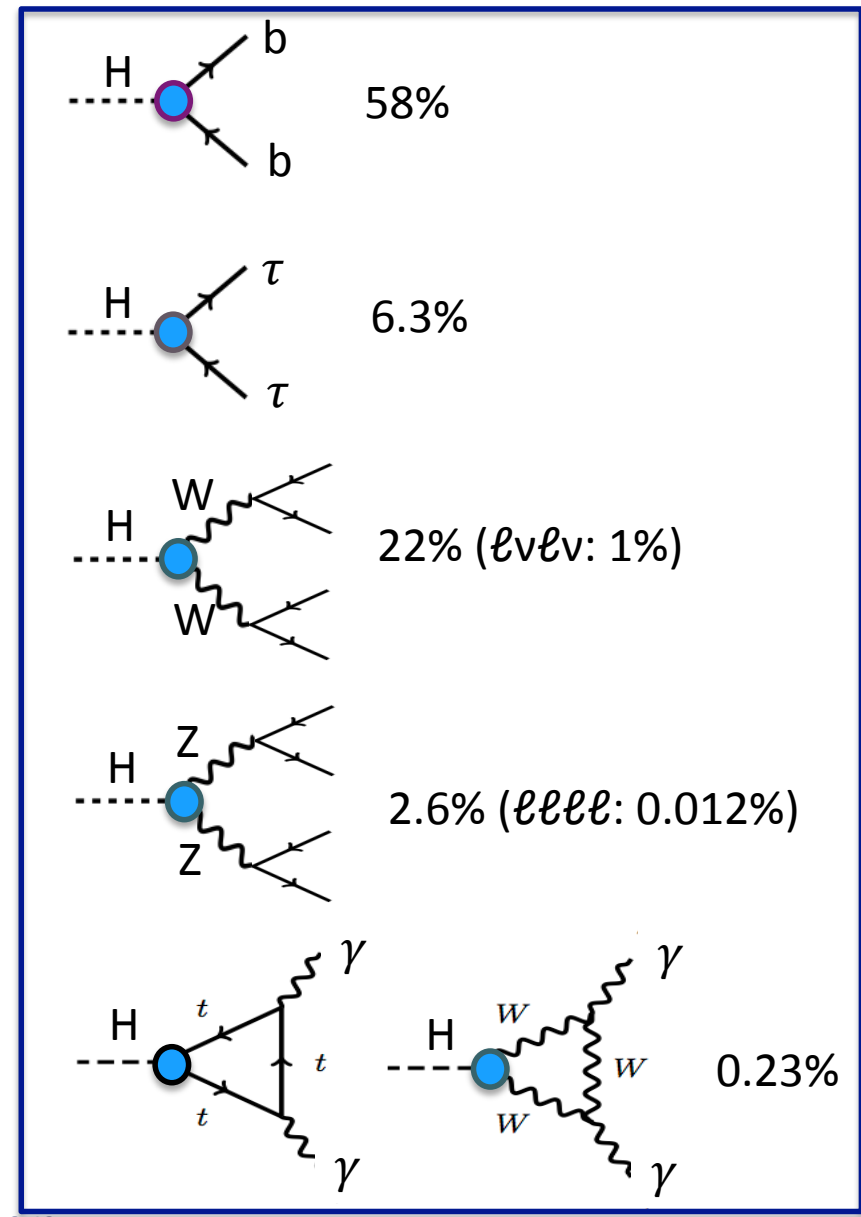
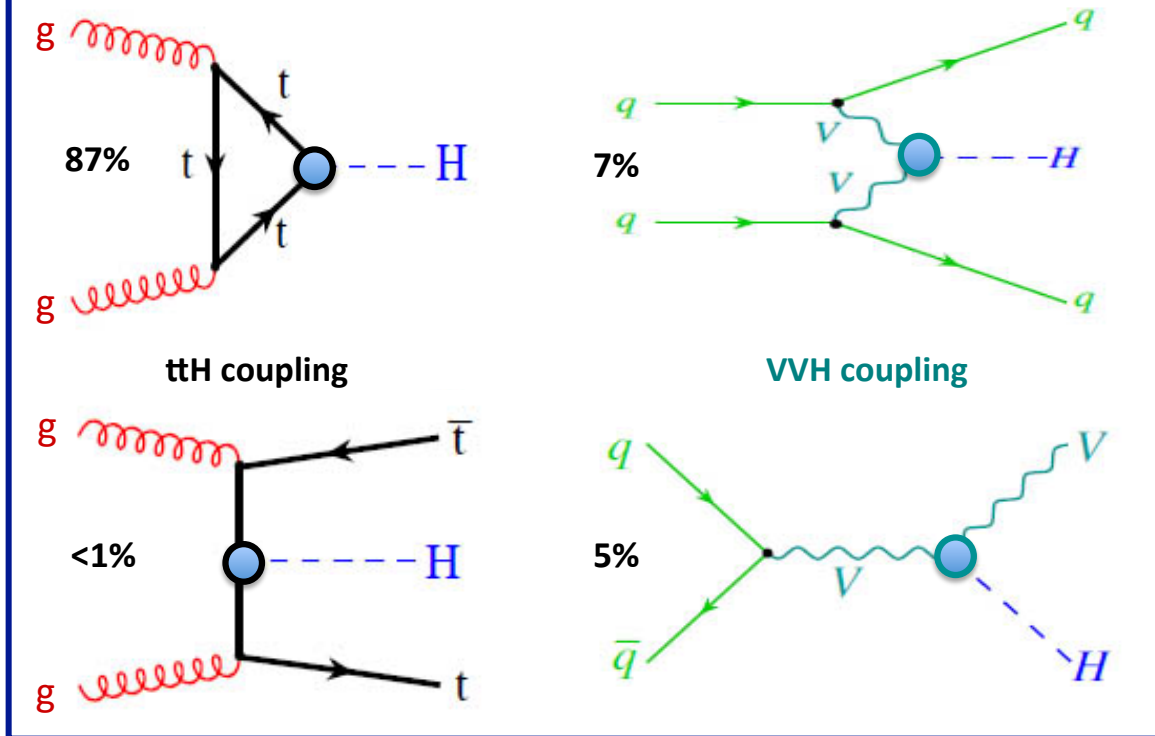
injectors: in 2024 => 13 months + 3 BC



LHC schedule approved by CERN management and LHC experiments spokespersons and technical coordinators
Monday 2nd December 2013

Higgs production and decay

$$\sigma_{\text{total}} = 22 \text{ pb @ } \sqrt{s} = 8 \text{ TeV, } m_H = 125 \text{ GeV}$$

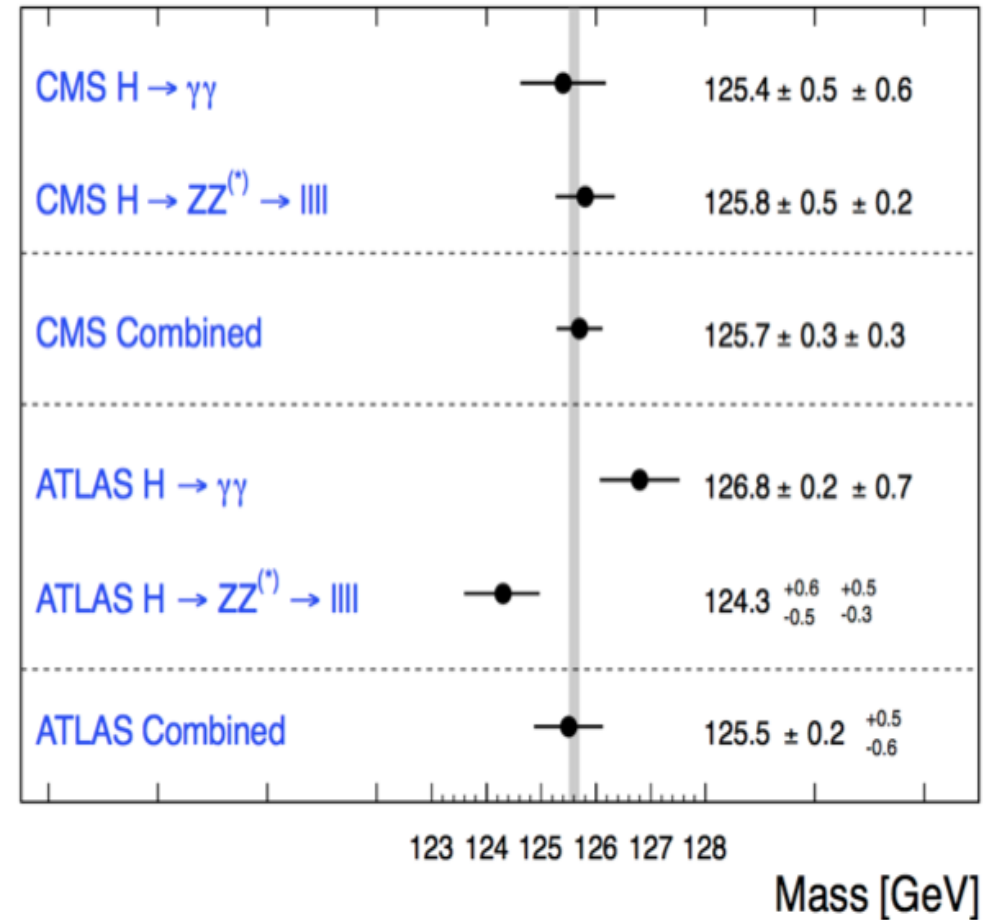
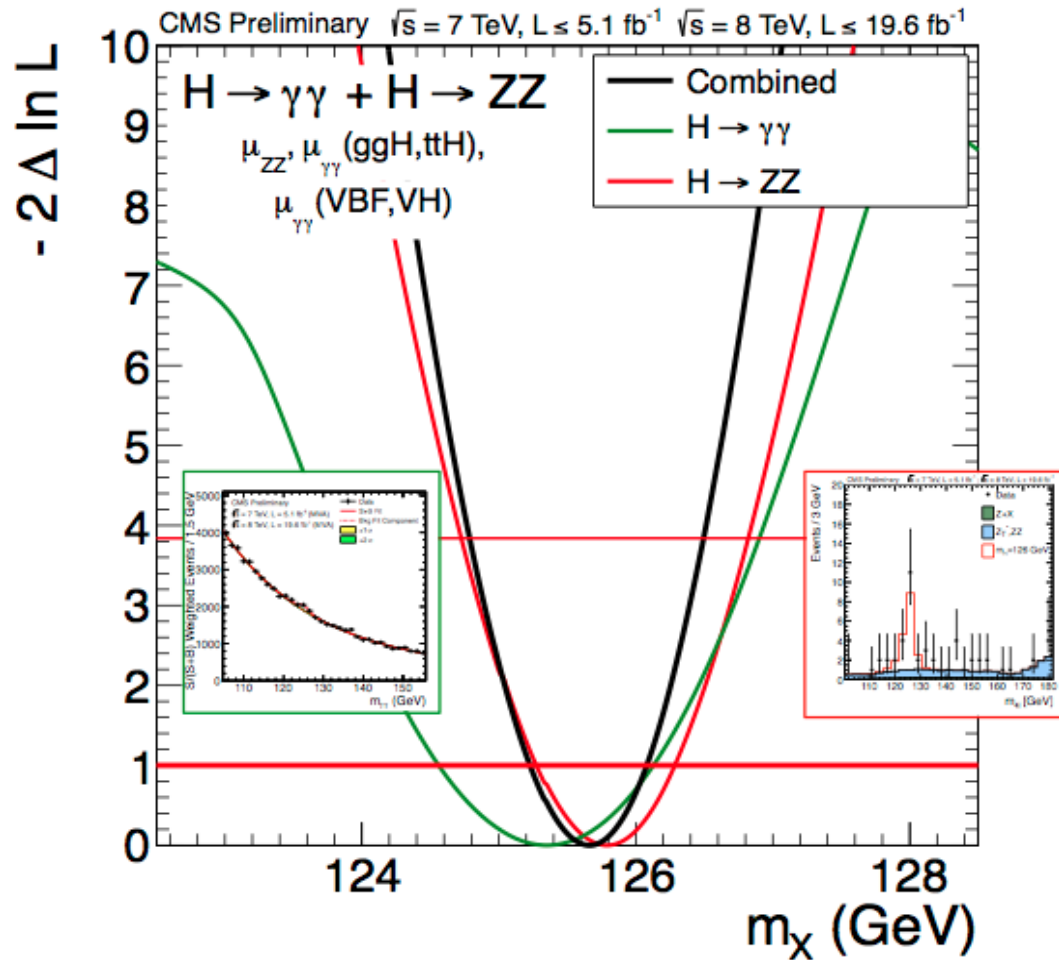


Measuring different production and decay modes, the different coupling constants can be determined

$$g(Hff) \propto m_{\text{fermion}}$$

$$g(HVV) \propto m_{\text{bozon}}^2$$

Higgs mass



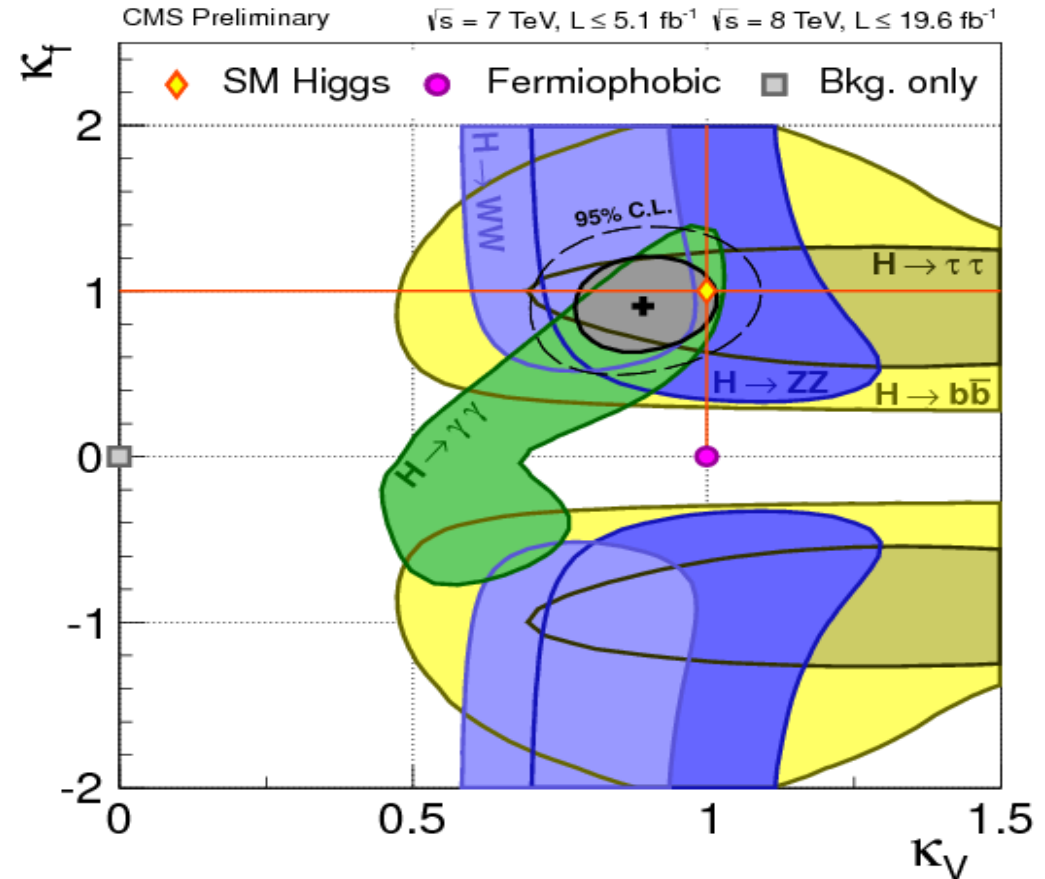
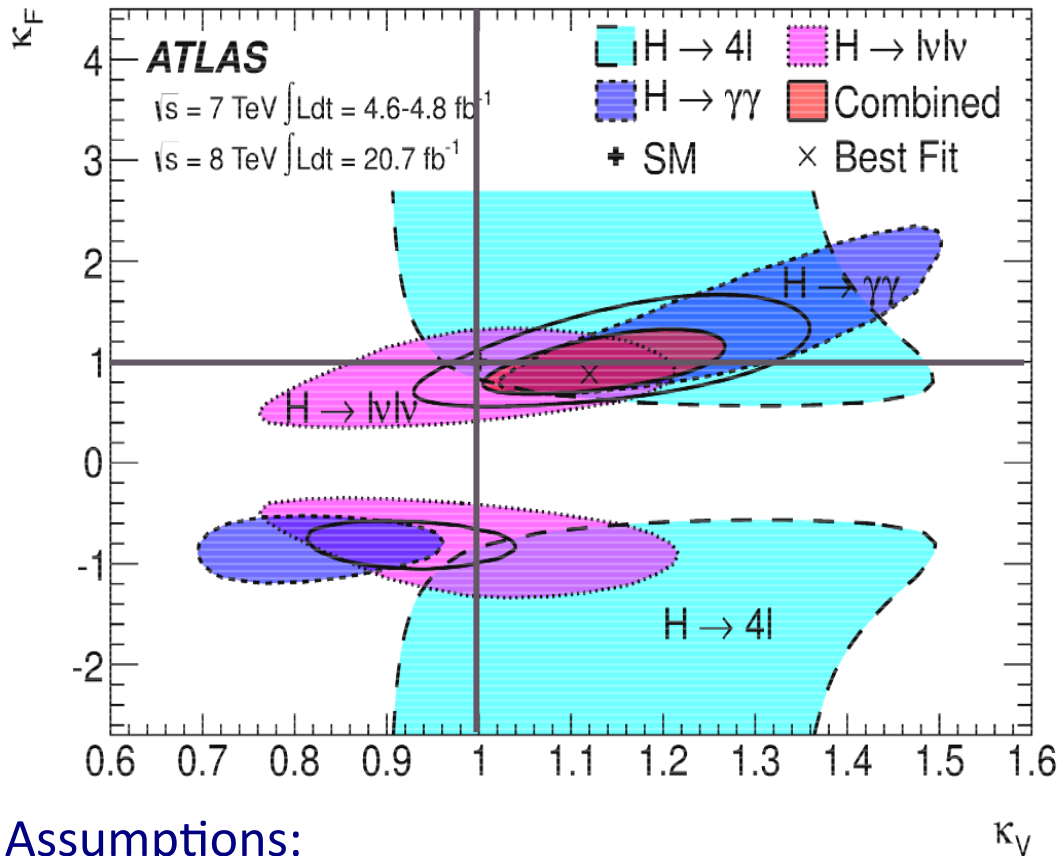
Final results require final Run1 calibration

Constrains both SUSY and composite-Higgs models! (See an example later...)

Coupling constant scale factors

$$\kappa_i \equiv g_i / g_i^{\text{SM}}$$

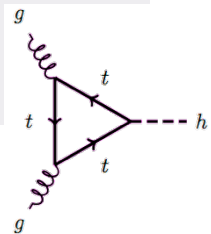
κ_V VS. κ_F



Assumptions:

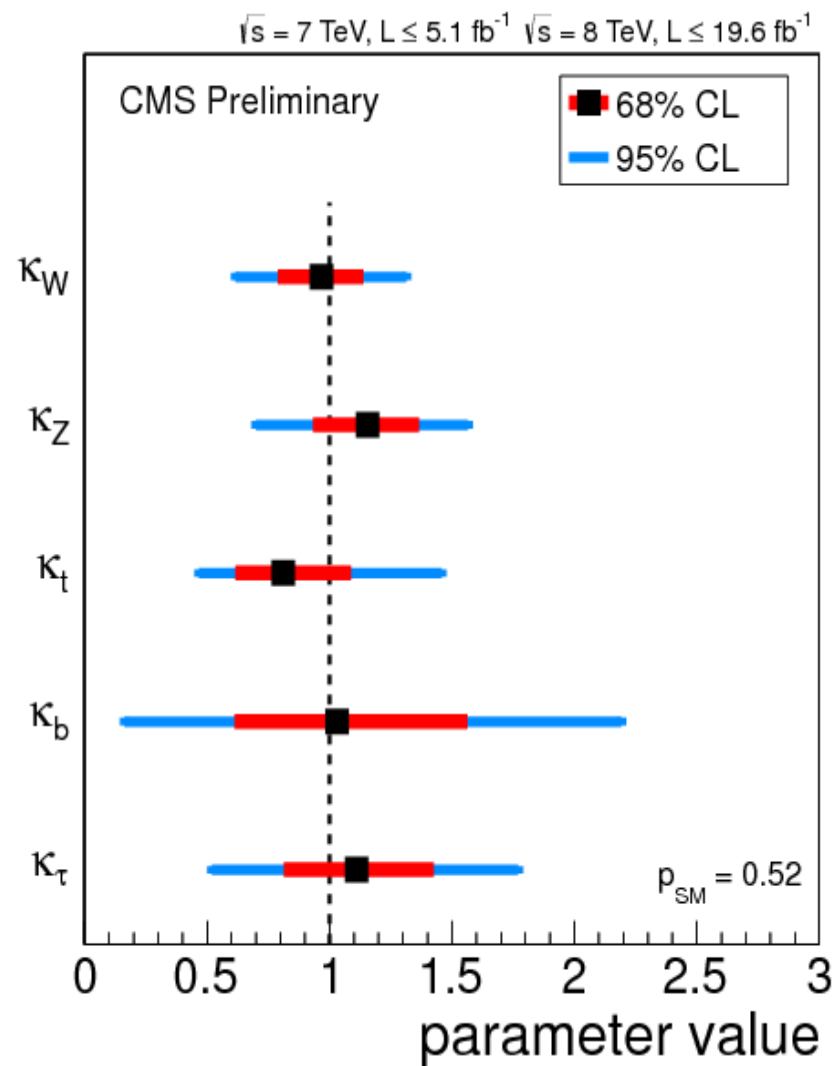
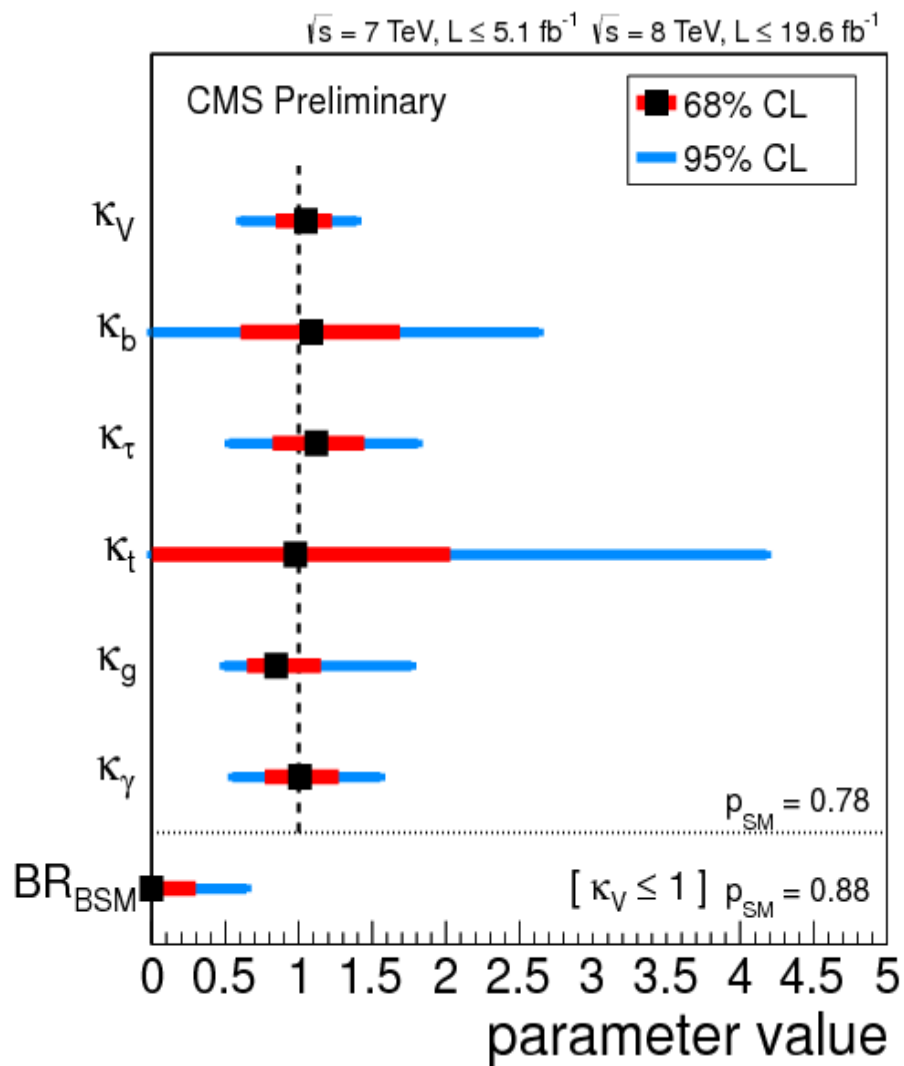
- All fermion (t, b, τ , ..) couplings scale with κ_F
- All heavy vector boson (W, Z) couplings scale with κ_V
- No other new physics contributes to the total width:
 $\kappa_g(\kappa_F \kappa_V), \kappa_\gamma(\kappa_F \kappa_V)$
- Total width Γ_H scales with $\kappa_H^2 \approx 0.7 \kappa_F^2 + 0.3 \kappa_V^2$

- **Agreement with SM within 10-20% precision**
- **$\kappa_F=0$ excluded with $>5\sigma$ (from indirect ggH loop)**

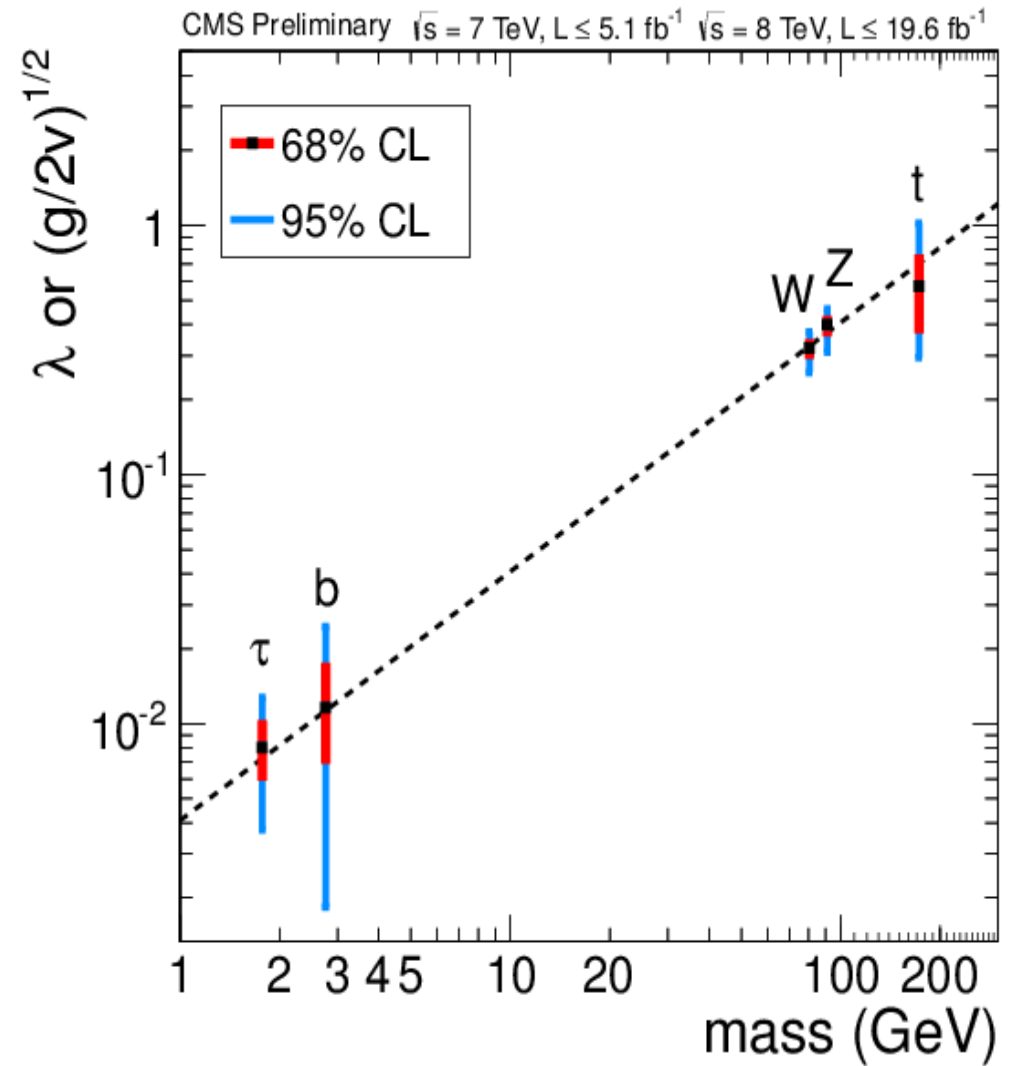
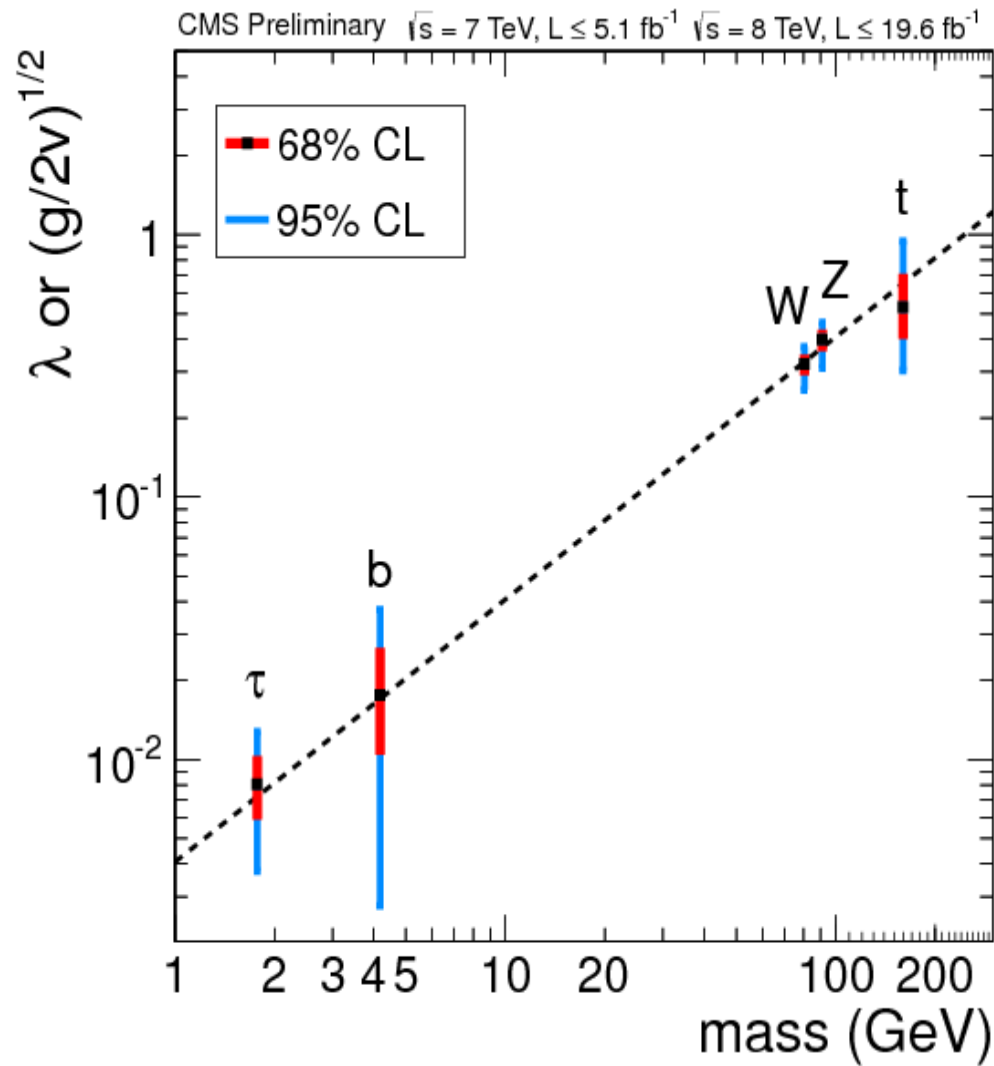


Coupling constant scale factors

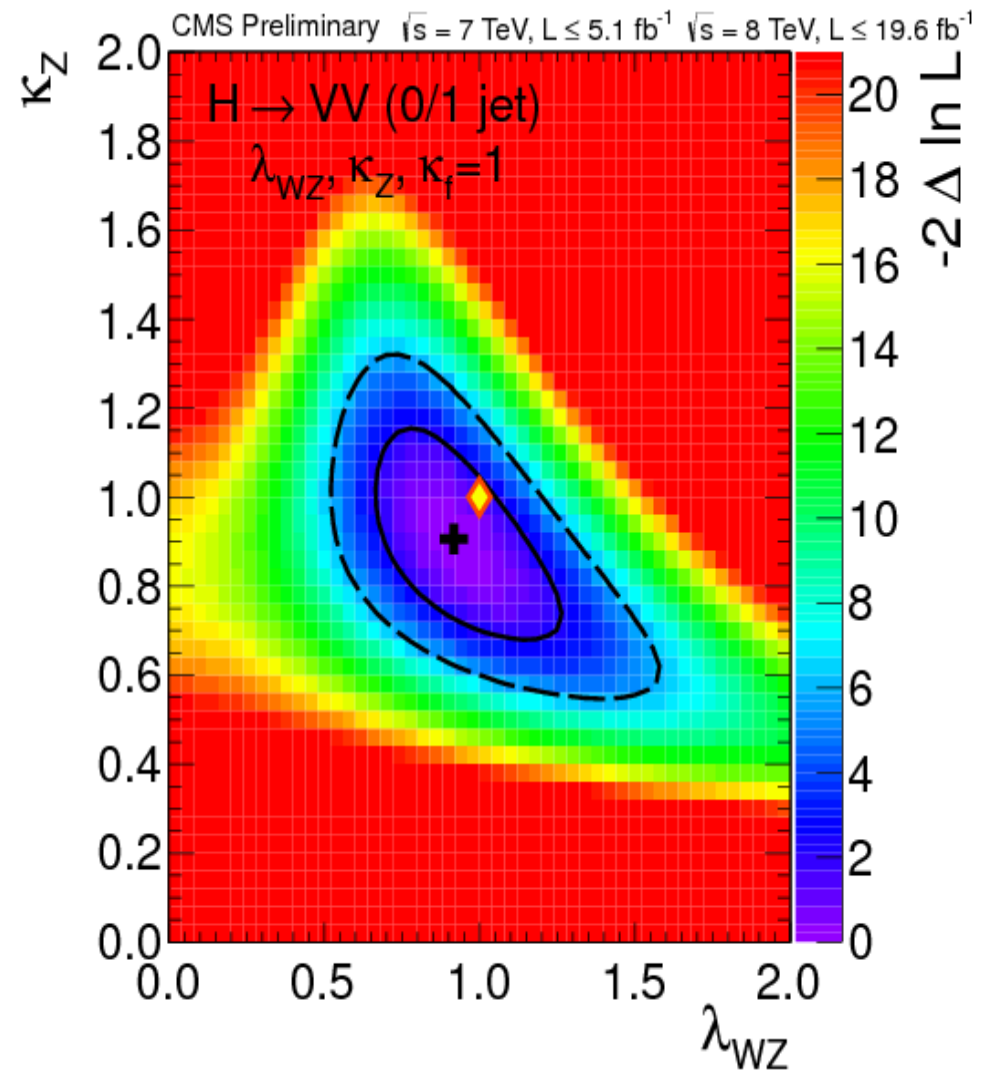
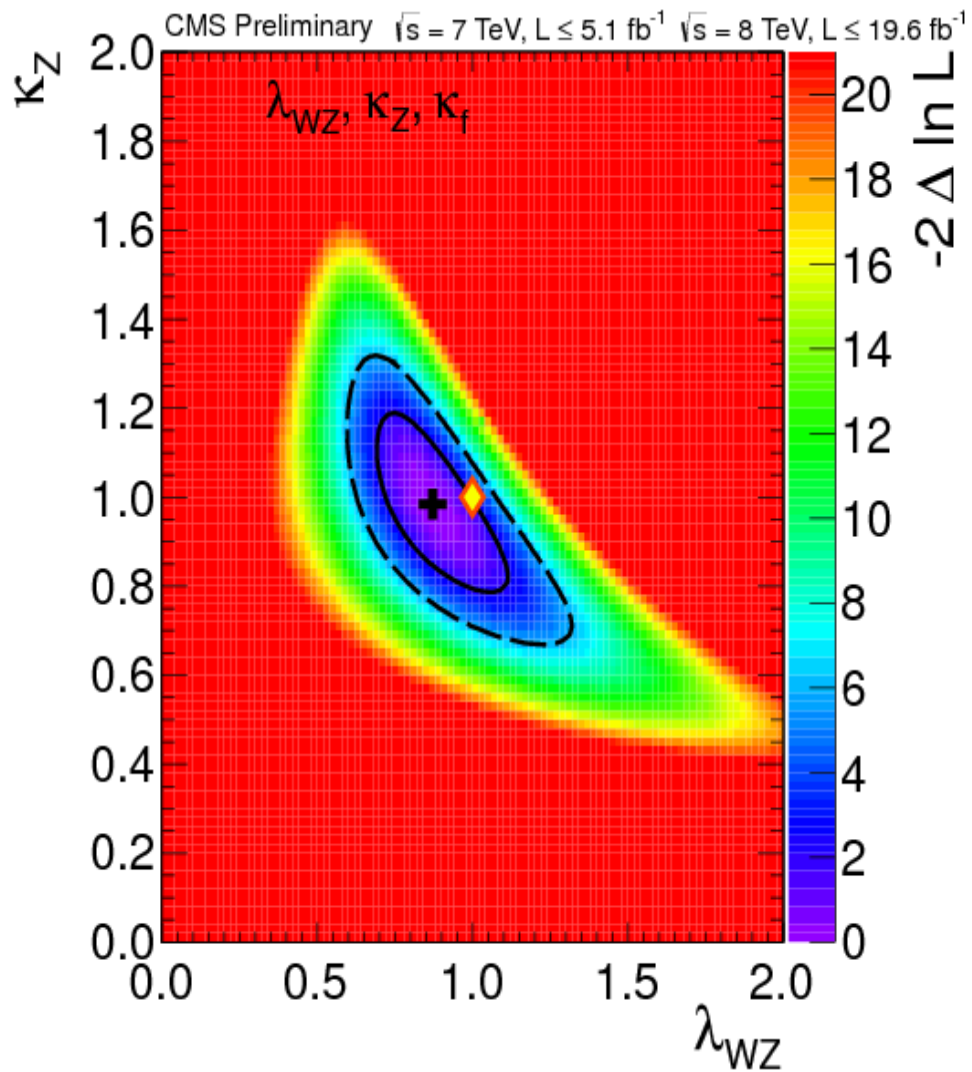
$$\kappa_i \equiv g_i / g_i^{\text{SM}}$$



Higgs coupling vs. mass

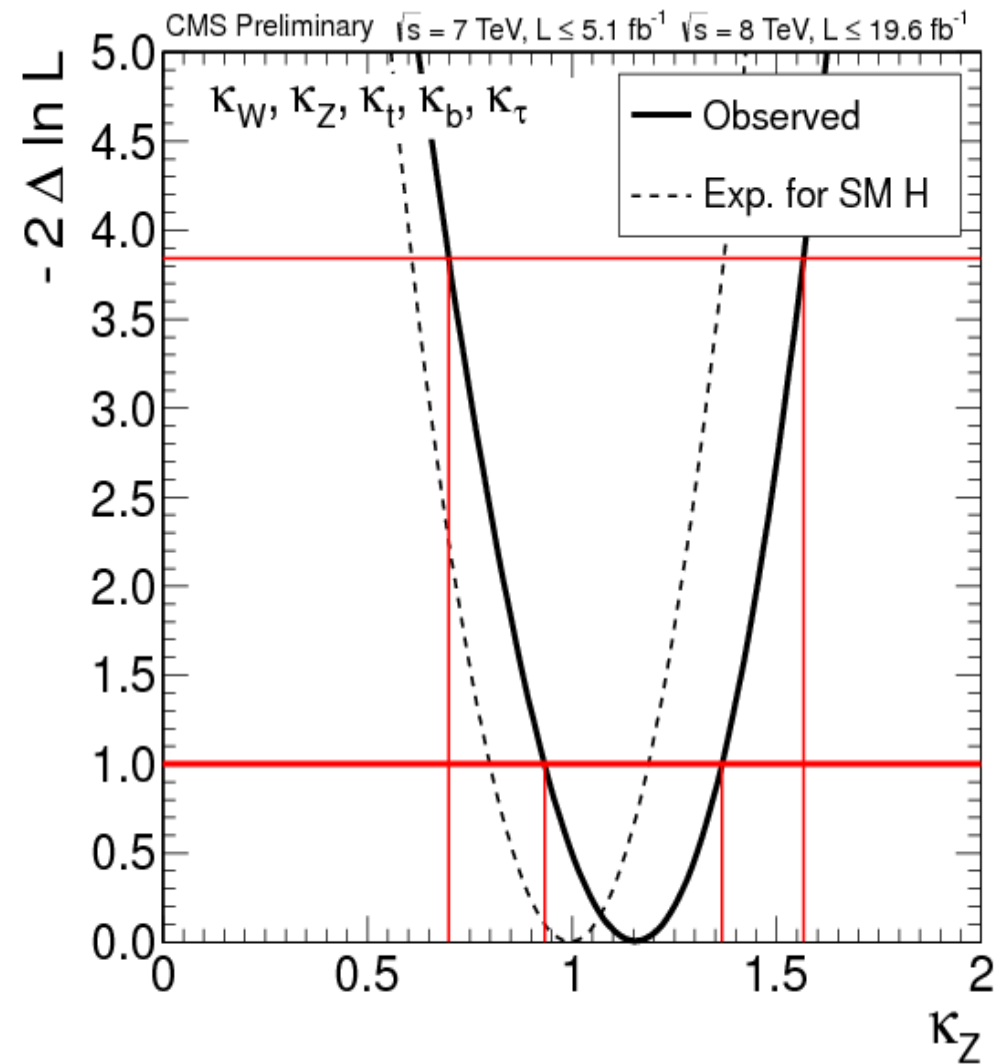
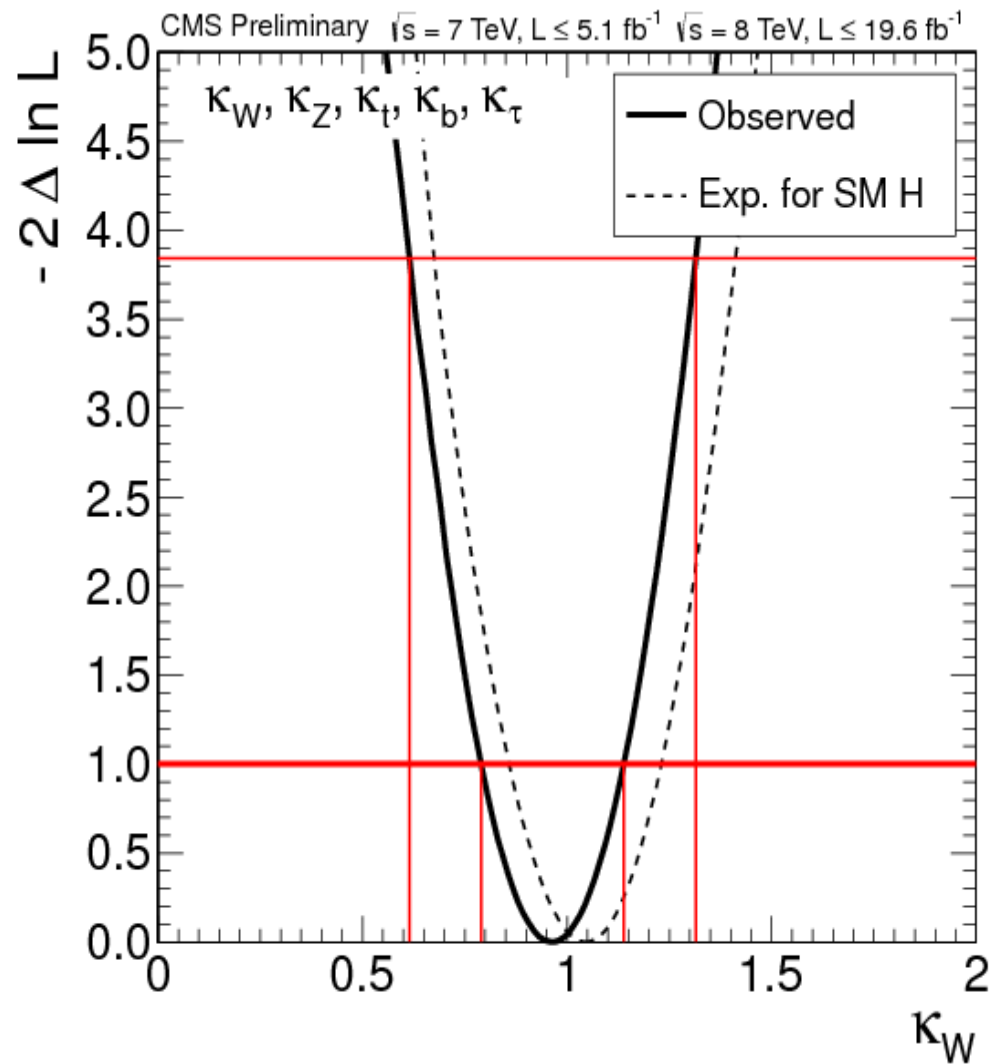


Test of custodial symmetry

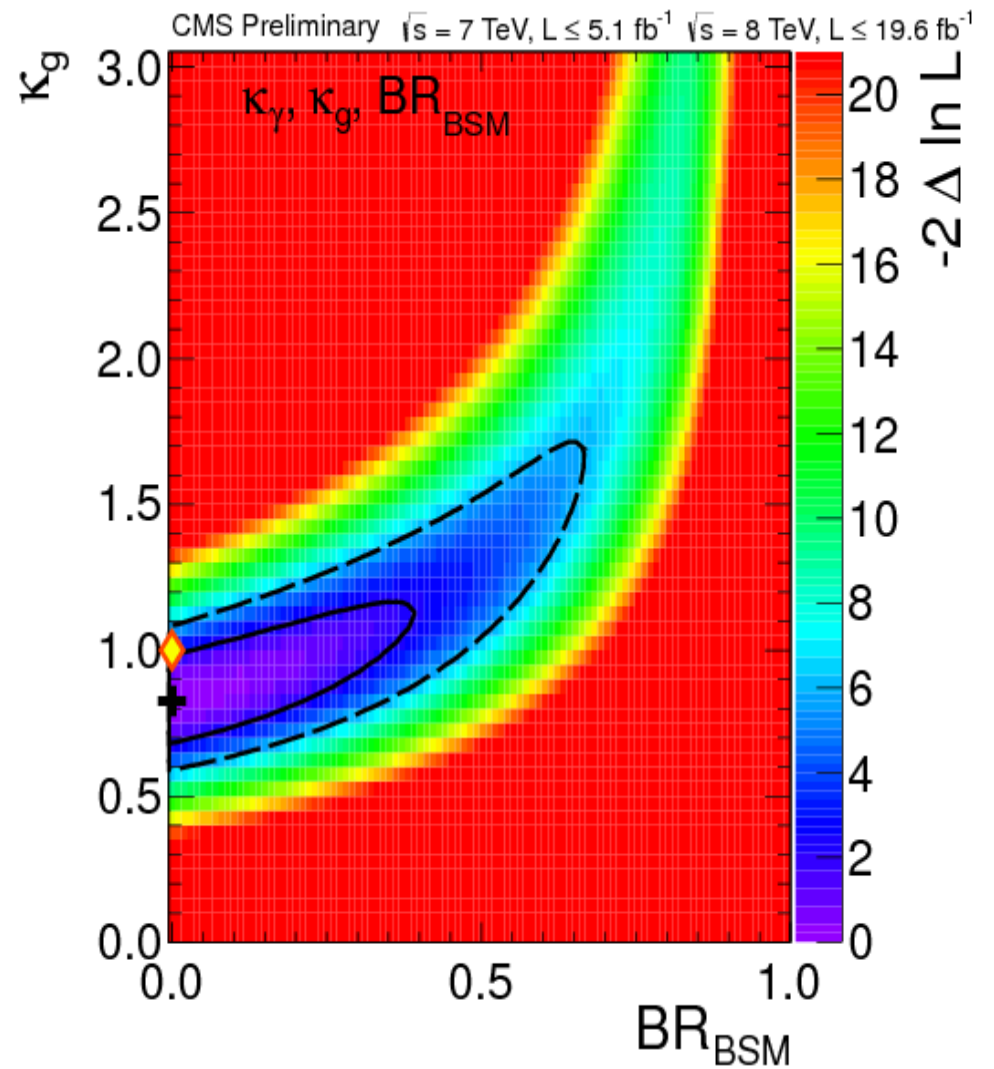
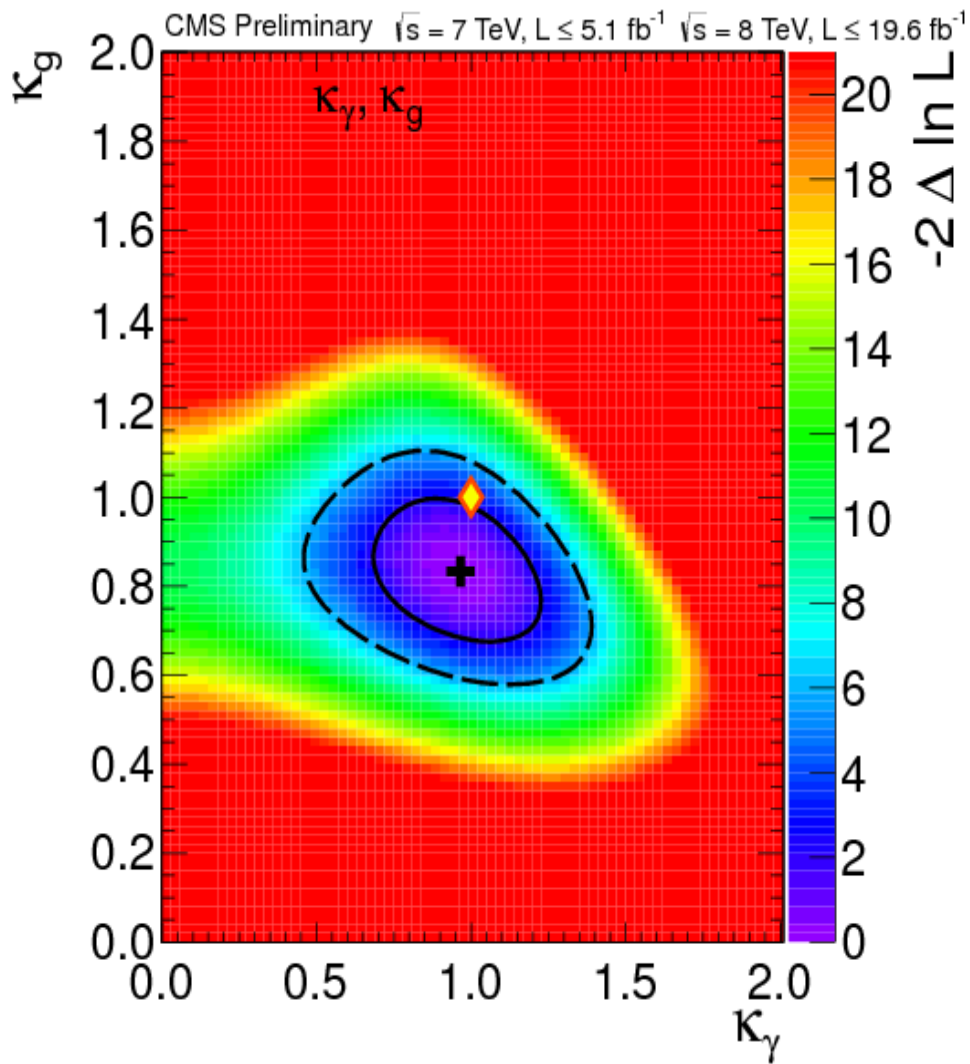


Generic search for deviation from SM

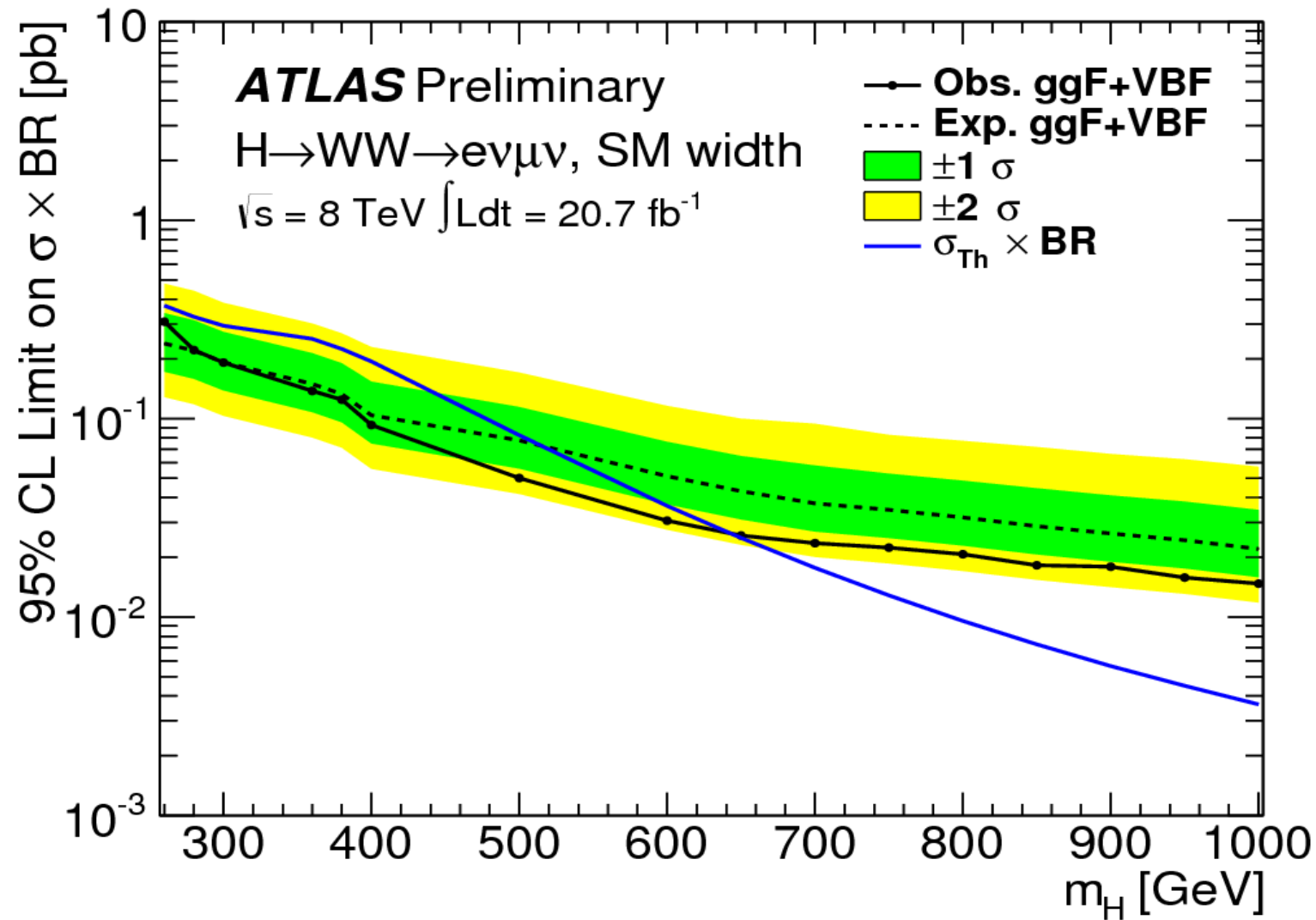
by profiling the other coupling modifiers



Search for new physics in loops



High-mass searches: an example

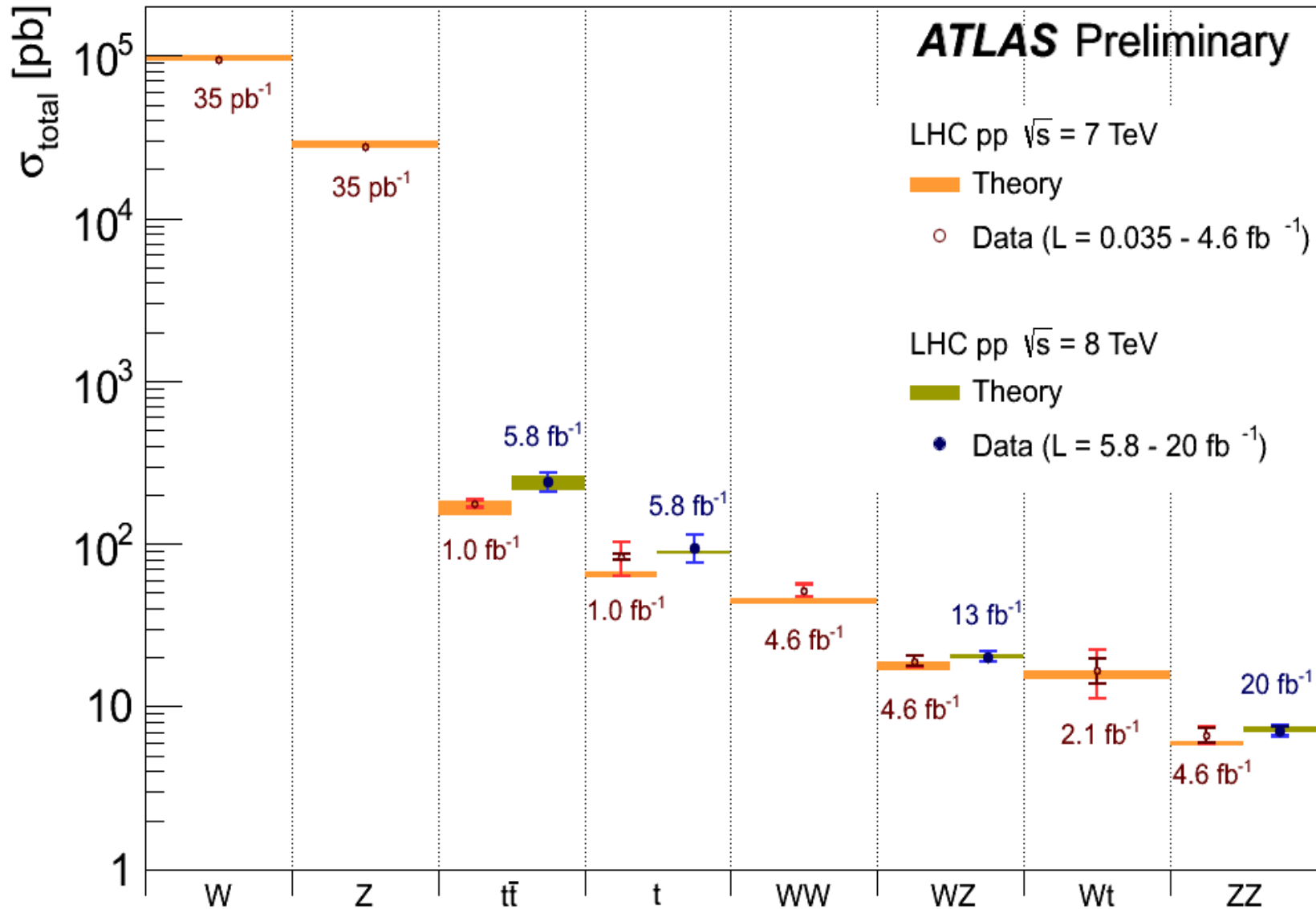


The μ -problem of SUSY

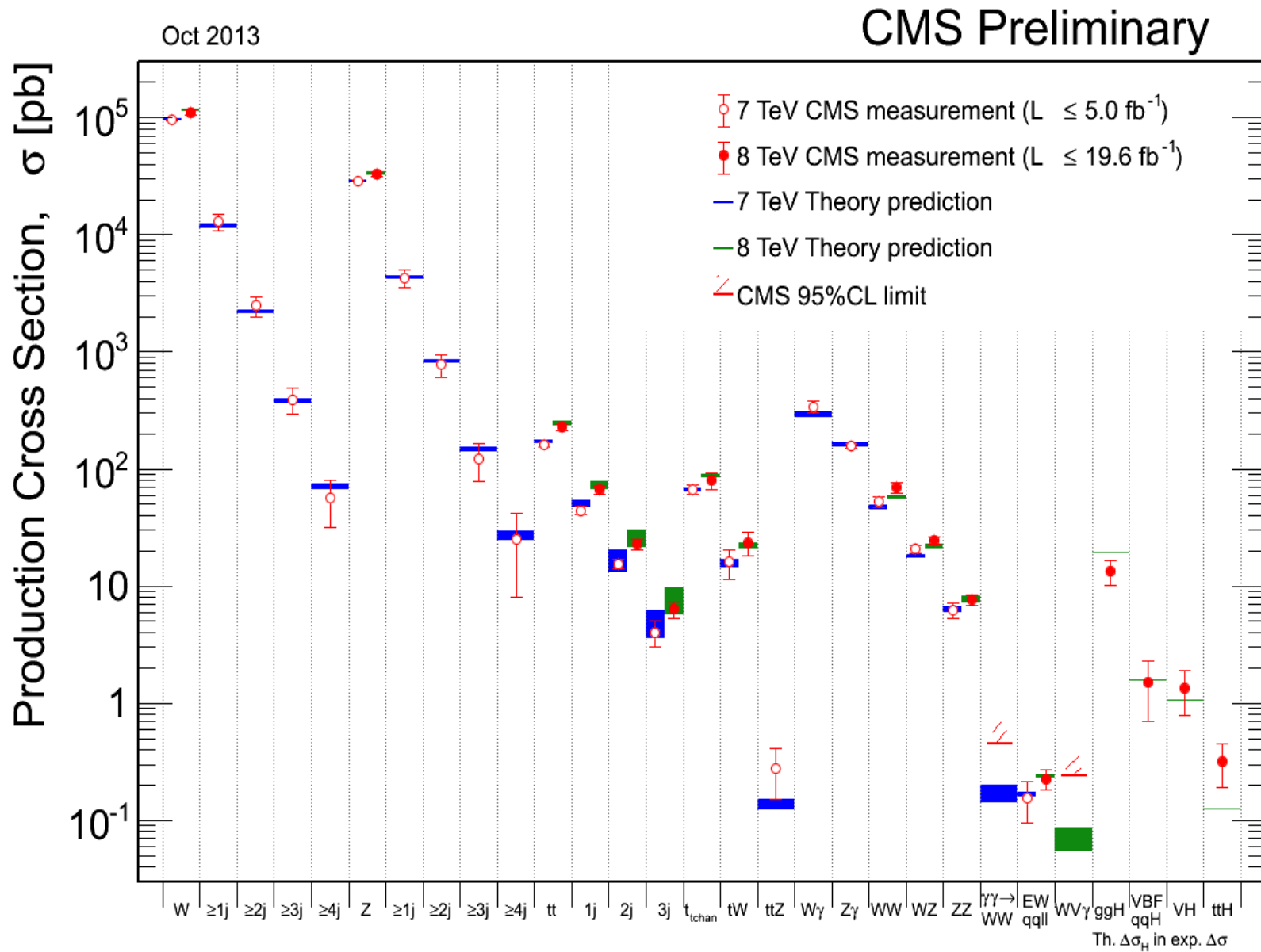
$$W_{\text{MSSM}} = \hat{Q}\hat{H}_u\mathbf{h}_u\hat{U}^C + \hat{H}_d\hat{Q}\mathbf{h}_d\hat{D}^C + \hat{H}_d\hat{L}\mathbf{h}_e\hat{E}^C + \mu\hat{H}_u\hat{H}_d$$

- μ is in the superpotential, so it is present before supersymmetry breaking (via soft terms). μ should know nothing about the electroweak scale.
- If $\mu = 0$ then there is no mixing between the two Higgs doublets. Any breaking of electroweak symmetry generated in the up-quark sector (by $M_H^2 < 0$) could not be communicated to the down-quark sector. The down-type quarks and leptons would remain massless.
- If $\mu = M_{\text{Planck}}$ then the Higgs bosons and their higgsino partners would gain Planck scale masses, in contradiction with upper bounds from triviality and precision electroweak data.
- For phenomenologically acceptable supersymmetry, the μ -parameter must be of the order of the electroweak scale.
- One way to link the μ -parameter with the electroweak scale is to make it a vacuum expectation value. Introduce a new iso-singlet neutral colorless chiral superfield \hat{S} coupling together the usual two Higgs doublet superfields. The scalar part of this is $\lambda S H_u H_d$. If S gains a vacuum expectation value we generate an effective μ -term. \rightarrow NMSSM

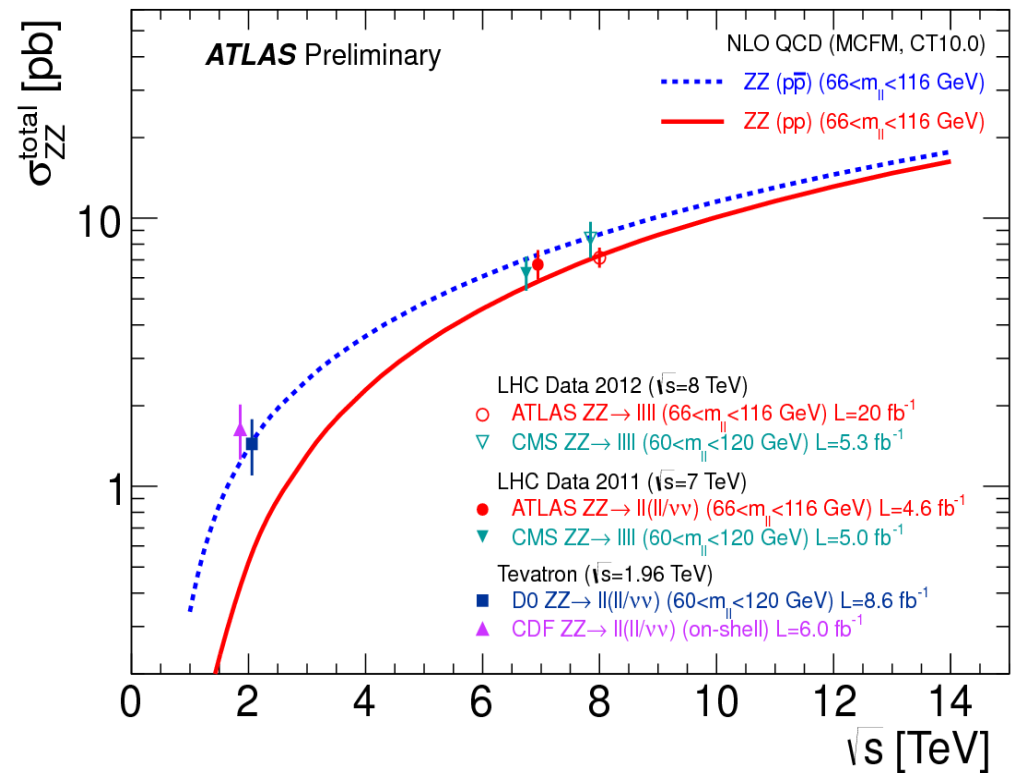
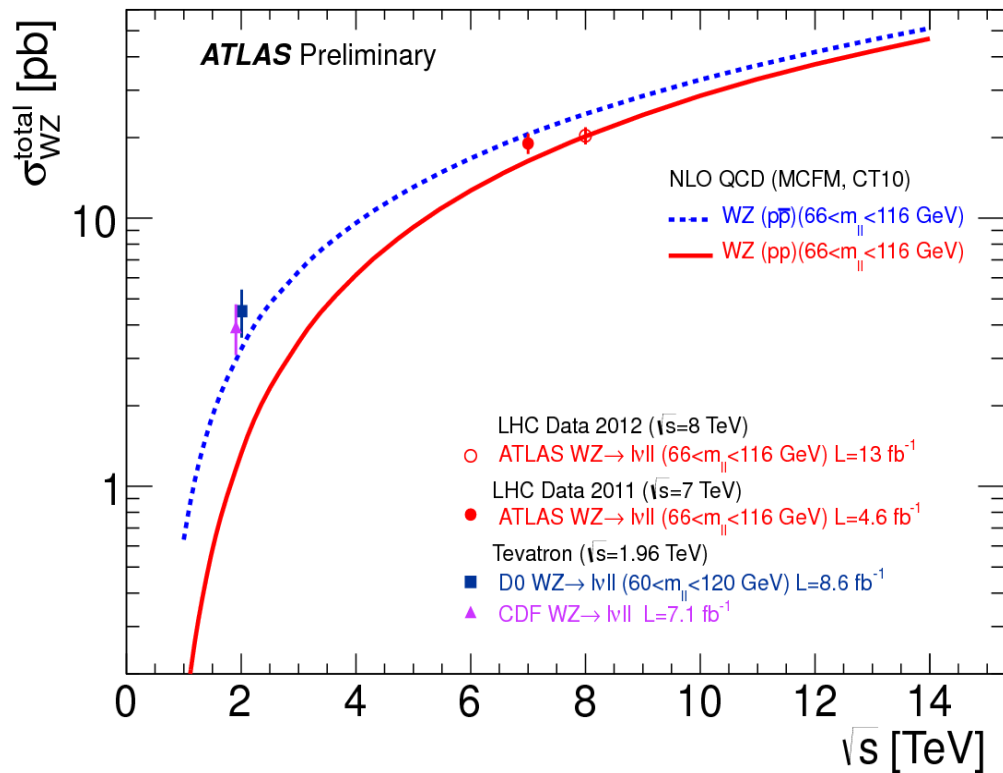
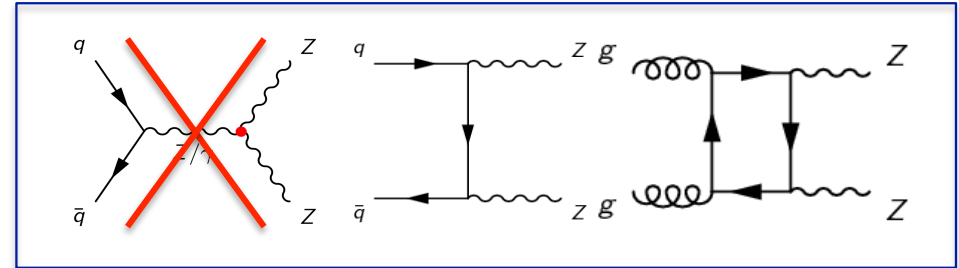
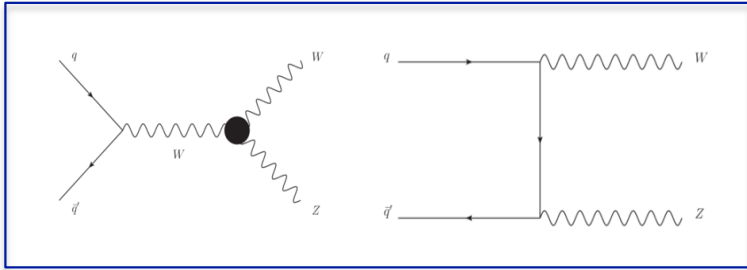
The SM strairway



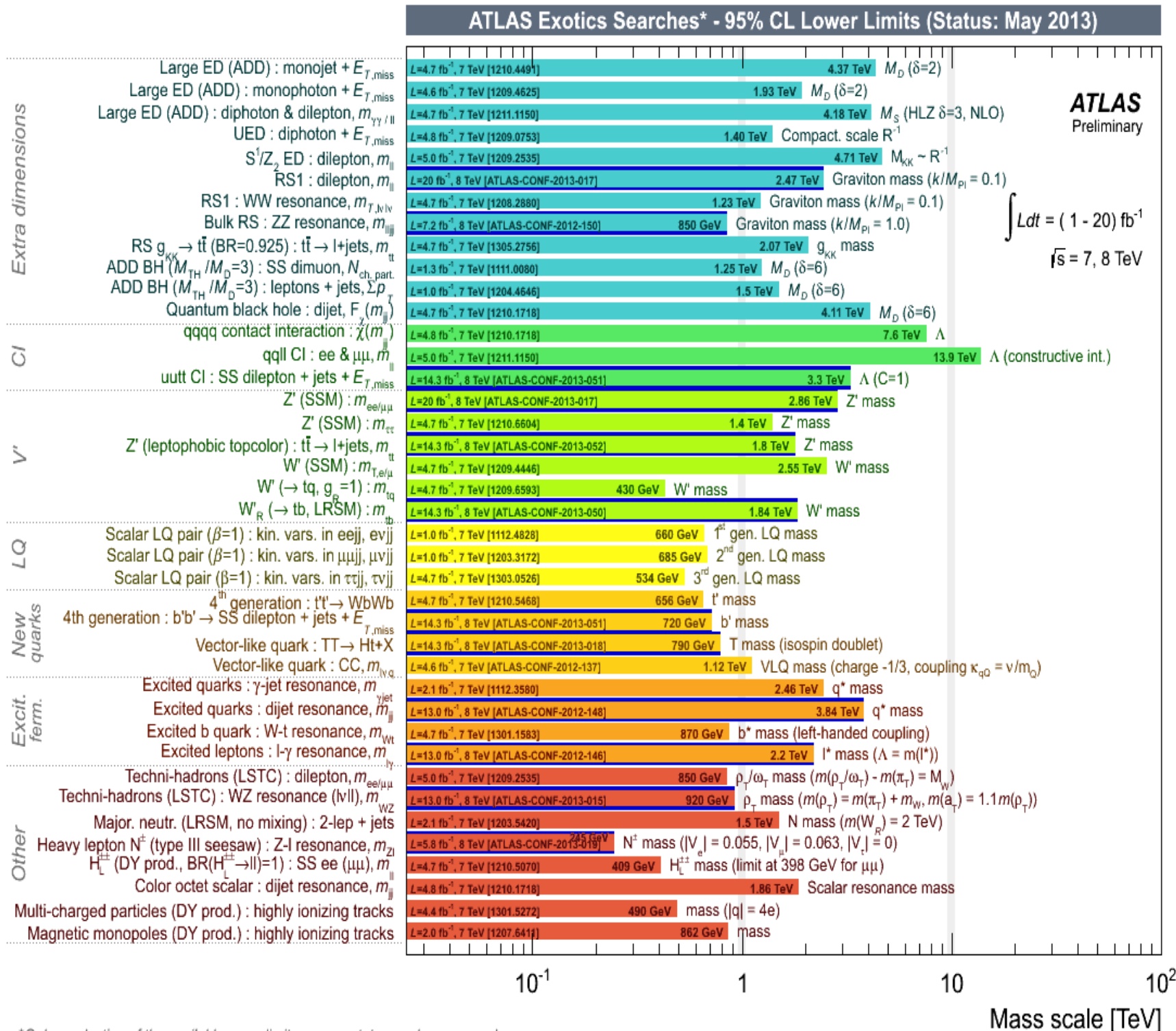
The SM strairway



Diboson measurements

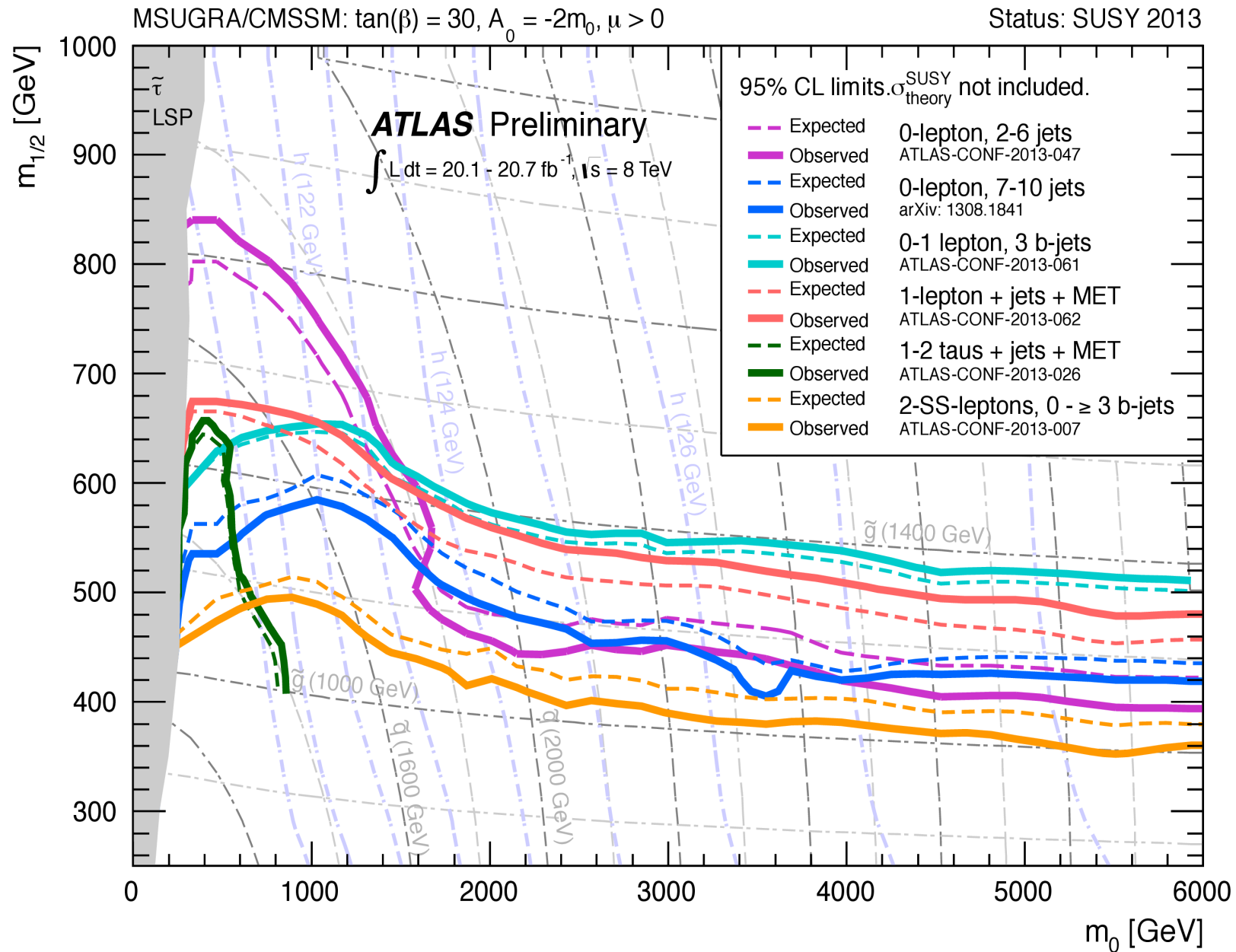


Testing „exotic” models

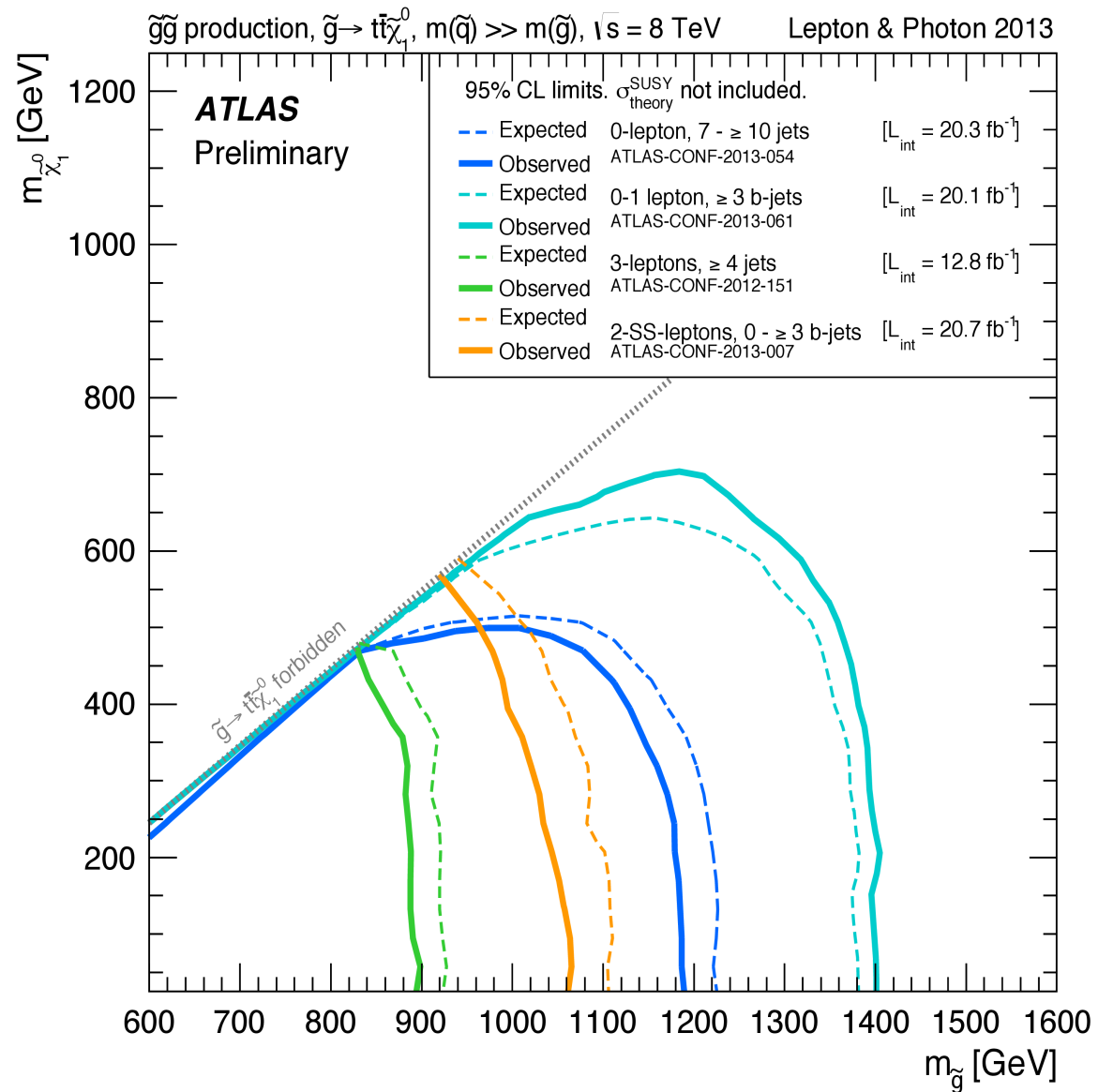


*Only a selection of the available mass limits on new states or phenomena shown

MSUGRA / CMSSM constraints

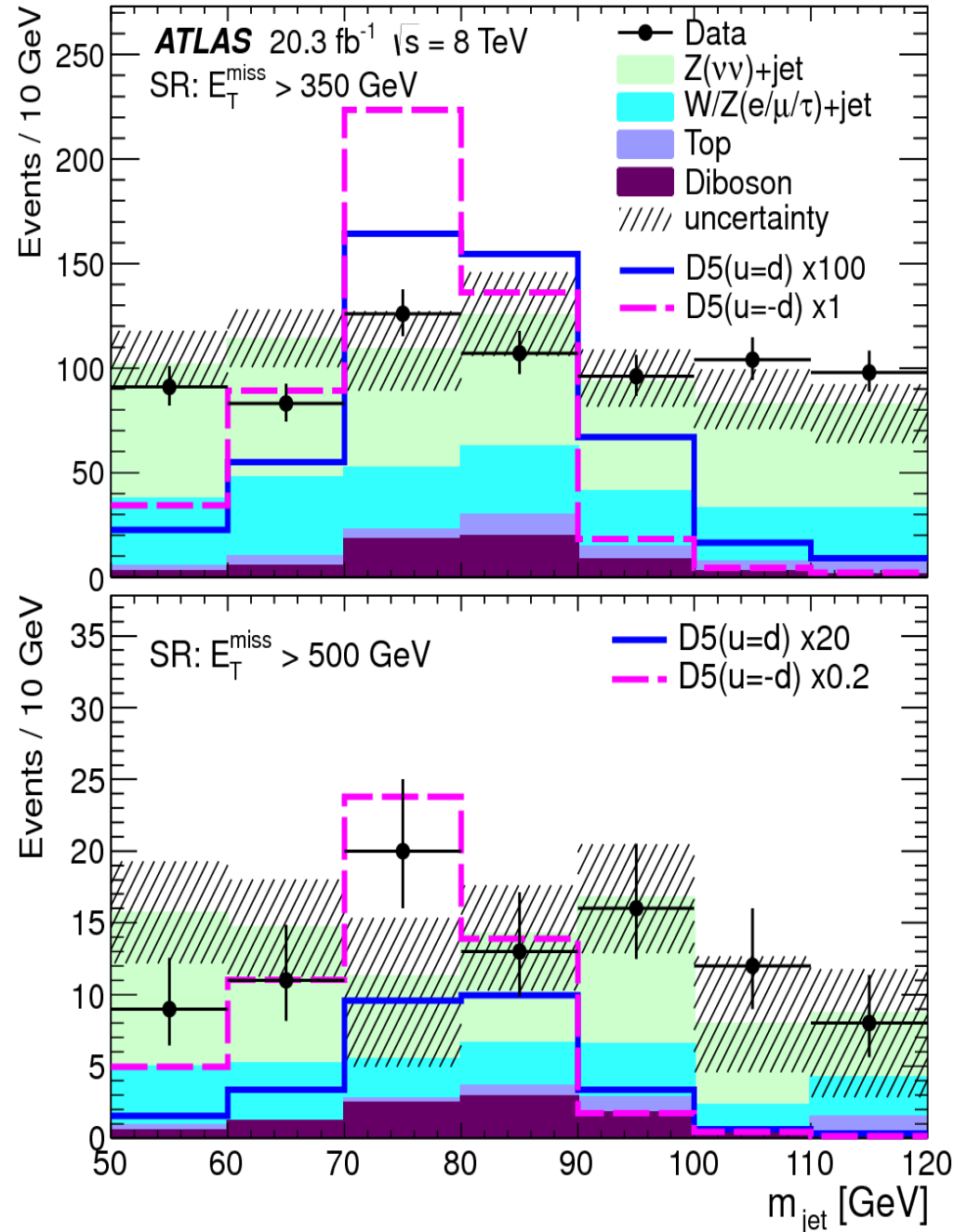
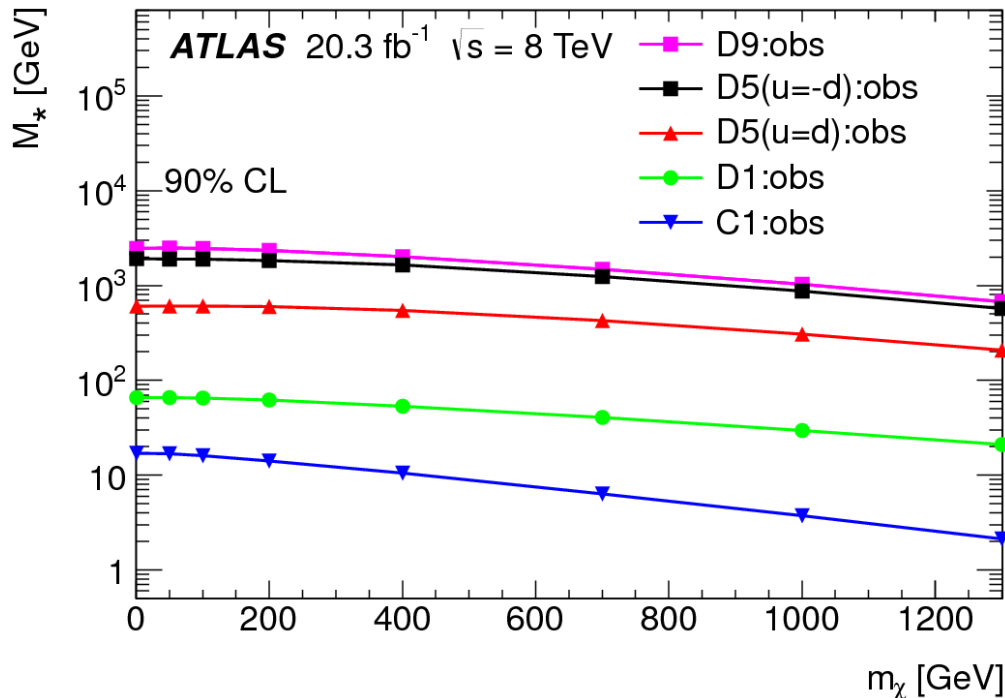
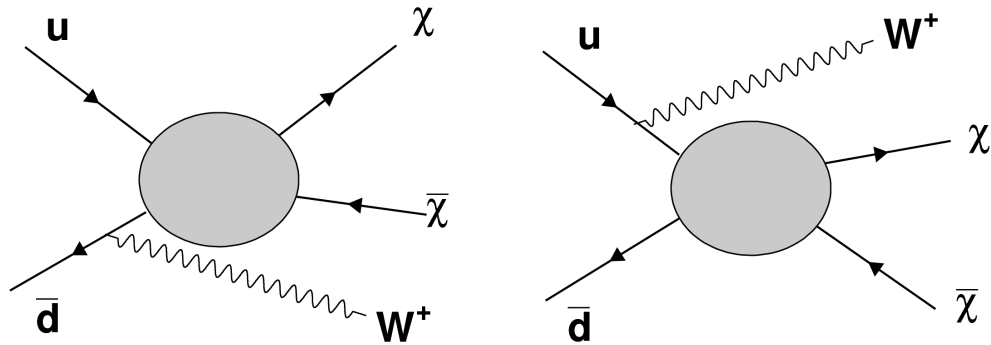


Stop searches



Search for dark matter

Search for pair-produced dark matter using initial state W/Z emission in hadronic final state



Search for dark matter

- Compare to direct dark matter detection experiments: best sensitivity at low masses
- Can be interpreted in models with invisibly decaying Higgs boson: complementary to dedicated searches

