## Measurement of longitudinal Parameters



## Measurement of longitudinal parameter:

- > Definition of longitudinal phase space
- > Proton LINAC: Determination of mean energy
- > Determination of longitudinal emittance
- > Bunch length measurement for non-relativistic beams
- > Bunch length measurement for relativistic beams
- > Summary

#### **Longitudinal** ↔ **transverse correspondences:**

- ➤ position relative to rf
   ⇔ transverse center-of-mass
- ➤ momentum or energy spread ↔ transverse divergence

## Measurement of longitudinal Parameters



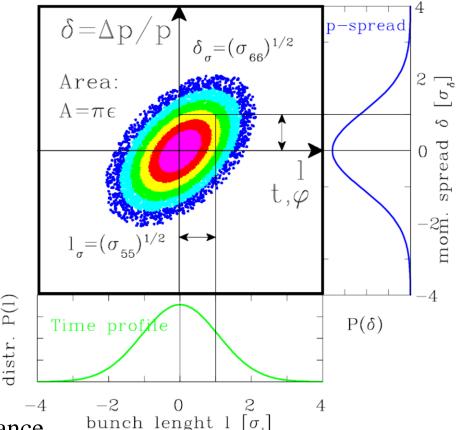
The longitudinal dynamics is described by the longitudinal emittance as given by:

- Spread of the bunches *l* in time, length *or* rf-phase.
- Momentum spread  $\delta = \Delta p/p$ , or energy spread  $\Delta W/W$

$$\Rightarrow \varepsilon_{long} = \frac{1}{\pi} \int_{A} dl \cdot d\delta$$

The normalized value is preserved:

$$\varepsilon_{long}^{norm} = \beta \gamma \cdot \varepsilon_{long}$$



#### **Discussed devices:**

- ➤ Pick-ups for bunch length and emittance.
- $\triangleright$  Special detectors (low  $E_{kin}$  protons), streak cameras & ele.-optical modulation (e<sup>-</sup>)

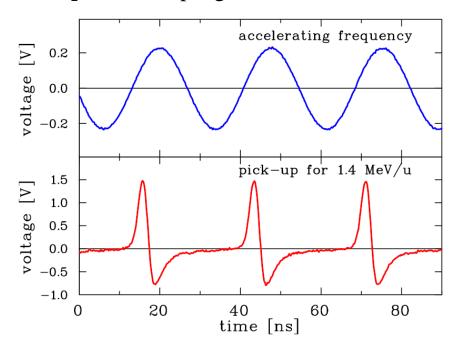
## The Bunch Position measured by a Pick-Up

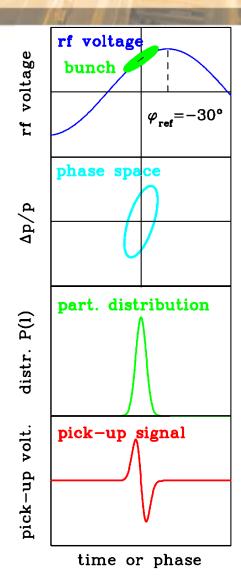


#### The **bunch position** is given relative to the accelerating rf.

e.g.  $\varphi_{ref}$ =-30° inside a rf cavity must be well aligned for optimal acceleration Transverse correspondence: Beam position

#### **Example:** Pick-up signal and 36 MHz rf at GSI-LINAC:







### **Outline:**

- > Definition of longitudinal phase space
- ➤ Proton LINAC: Determination of mean energy used for alignment of cavities phase and amplitude
- > Determination of longitudinal emittance
- > Bunch length measurement for non-relativistic beams
- > Bunch length measurement for relativistic beams
- > Summary

## Determination of non-relativistic mean Energy using Pick-Ups



The energy delivered by a LINAC is sensitive to the mechanics, rf-phase and amplitude.

For non-relativistic energies at proton LINACs time-of-flight (TOF) with two pick-ups is used:

$$\beta c = \frac{L}{NT + t_{\text{scope}}}$$

 $\rightarrow$  the velocity  $\beta$  is measured.

*Example*: Time-of-flight signal from two pick-ups at 1.4 MeV/u:

The reading is  $t_{scope} = 15.82(5)$ ns with

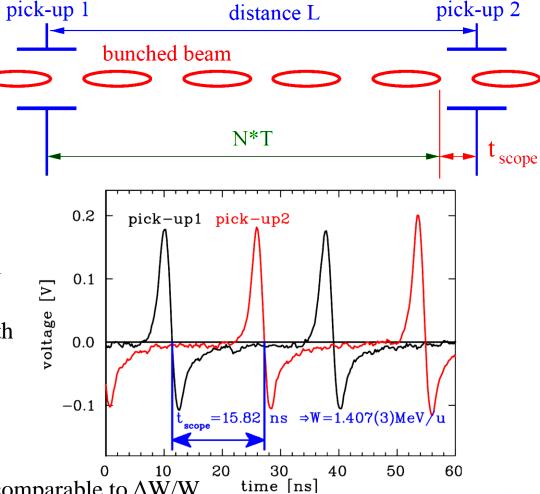
$$f_{rf} = 36.136 \text{MHz} \Leftrightarrow T = 27.673 \text{ns}$$

$$L = 1.629 \text{ m and } N = 3$$

$$\Rightarrow \beta = 0.05497(7)$$

$$\Rightarrow$$
 *W*=1.407(3) MeV/u

The accuracy is typically 0.1 % i.e. comparable to  $\Delta W/W$ 



## Precision of TOF Measurement for non-relativistic Energy



The precision of TOF is given by the accuracy in time and distance reading:

$$\frac{\Delta \beta}{\beta} = \sqrt{\left(\frac{\Delta L}{L}\right)^2 + \left(\frac{\Delta t}{NT + t_{\text{scope}}}\right)^2}$$

Accuracy of scope reading  $\Delta t \approx 100$  ps, uncertainty in distance  $\Delta L \approx 1$  mm.

**Example:** GSI-LINAC: L = 3.25 m and  $f_{rf} = 36$  MHz:

location (GSI-slang)		RFQ	IH1	IH2	AL4
energy $W$	[MeV/u]	0.12	0.75	1.4	11.4
velocity $\beta$	%	1.6	4.0	5.5	15.5
total TOF	[ns]	677	271	197	70
bunch spacing $\beta c/f_{rf}$	[cm]	13	33	45	129
Number of bunches $N$		25	9	7	2
resolution $\Delta W/W$	%	0.07	0.10	0.12	0.22

- $\triangleright$  The accuracy is typically 0.1 % (same order of magnitude as  $\Delta W/W$ )
- > The length has to be matched to the velocity
- $\triangleright$  Due to the distance of  $\approx$  3 m, different solutions for the # of bunches N are possible
- → A third pick-up has to be installed closed by, to get an unique solution.

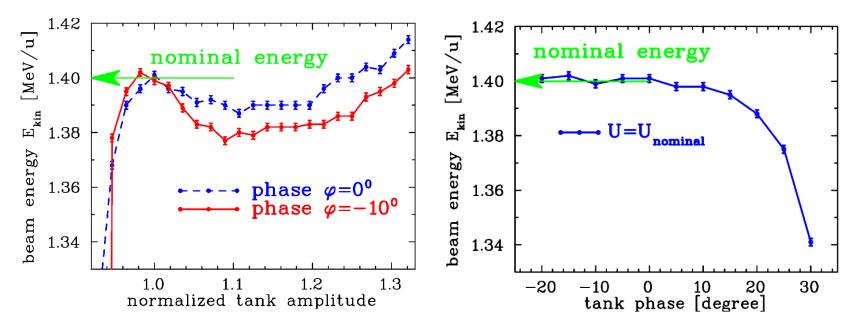
## Cavity Alignment using a TOF Measurement



The mean energy is important for the matching between LINAC module. It depends on phase and amplitude of the rf wave inside the cavities.

*Example:* Energy at GSI LINAC (nominal energy 1.400 MeV/u):

(distance between pick-ups:  $L = 1.97 \text{ m} \Rightarrow N = 4 \text{ bunches}$ )



>Proton LINACs: Amplitude and phase should be carefully aligned by precise TOF

**Electron LINACs**: Due to relativistic velocity, TOF is not applicable.



#### **Outline:**

- > Definition of longitudinal phase space
- > Proton LINAC: Determination of mean energy used for alignment of cavities phase and amplitude
- > Determination of longitudinal emittance
  - LINAC: variation of bunch length
  - **Synchrotron: Topographic reconstruction**
- **Bunch length measurement for non-relativistic beams**
- > Bunch length measurement for relativistic beams
- > Summary

# Longitudinal Emittance by linear Transformation using a Buncher

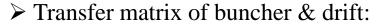


#### **Longitudinal focusing:**

Variation of the bunch shape by a rf-buncher

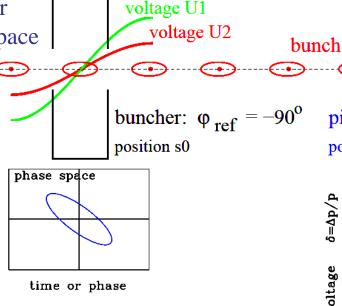
→ components 5 and 6 from 6-dim phase-space

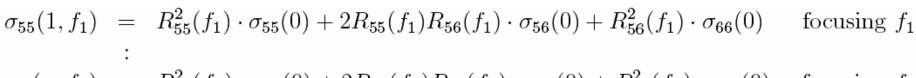
Transversal corres.: Quadrupole variation —

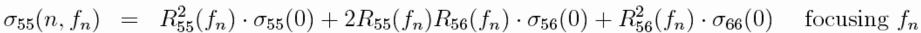


$$R_{buncher} = \begin{pmatrix} 1 & 0 \\ -1/f & 1 \end{pmatrix}, R_{drift} = \begin{pmatrix} 1 & L/\gamma^2 \\ 0 & 1 \end{pmatrix}$$
with focal length:  $1/f = \frac{2\pi f_{rf}}{Anv^2} \cdot U$ 

- > Variation of buncher amplitude *U*  $\Rightarrow$  different bunch width at  $s_1$ : beam matrix  $\Delta t^2_{max} = \sigma_{55}(1, f)$
- $\triangleright$  System of redundant linear equations for  $\sigma_{ii}(0)$ :







pick-up

position s1

phase space

pick-up signal

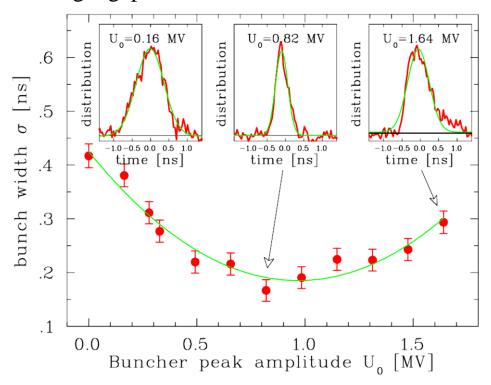
time or phase

## Result of a longitudinal Emittance Measurement



#### **Example** GSI LINAC:

The voltage at the single gap -resonator is varied for 11.4 MeV/u Ni<sup>14+</sup> beam, 31 m drift:

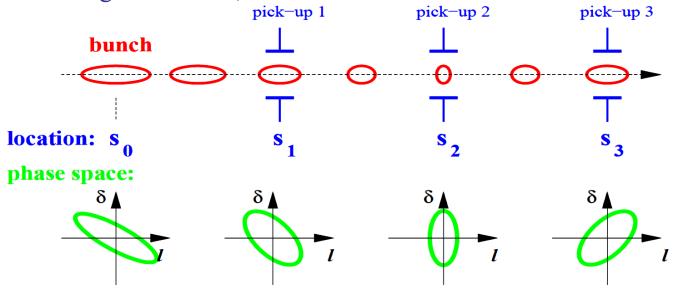


- > The structure of short bunches can be determined with special monitor
- ➤ This example: The resolution is better than 50 ps or 2° for 108 MHz
- > Typical bunch length at proton LINACs: 30 to 200 ps

## Longitudinal Emittance within a Transfer-Line



As for the 'three grid' method, the emittance can be determined in a transfer line.



The system of redundant linear equations with the transfer matrix  $\mathbf{R}(\mathbf{i})$  to location  $s_i$ :

$$\begin{array}{lll} \sigma_{55}(1) & = & R_{55}^2(1) \cdot \sigma_{55}(0) + 2R_{55}(1)R_{56}(1) \cdot \sigma_{56}(0) + R_{56}^2(1) \cdot \sigma_{66}(0) & \mathbf{R}(1) : s_0 \to s_1 \\ & : \\ \sigma_{55}(n) & = & R_{55}^2(n) \cdot \sigma_{55}(0) + 2R_{55}(n)R_{56}(n) \cdot \sigma_{56}(0) + R_{56}^2(n) \cdot \sigma_{66}(0) & \mathbf{R}(n) : s_0 \to s_n \end{array}$$

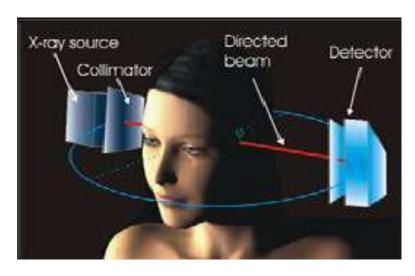
**Assumptions:**  $\triangleright$  Bunches much longer than pick-up or relativistic E -field:  $E_{\perp} >> E_{\parallel}$   $\triangleright$  Gaussian distribution without space-charge effects.

# Longitudinal Emittance using tomographic Reconstruction



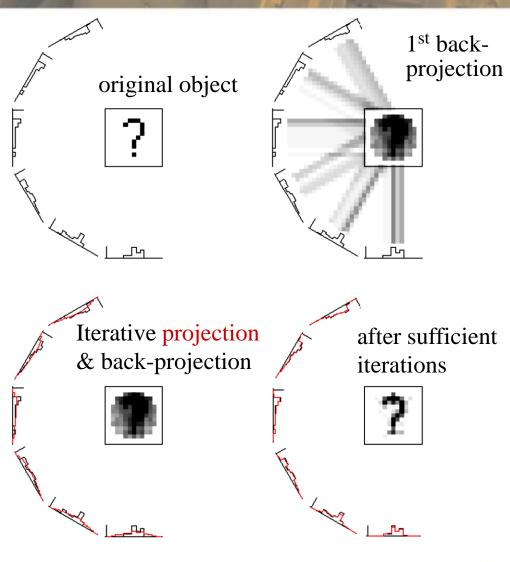
# Tomography is medical image method Tomography:

2-dim reconstruction of sufficient 1-dim projections



### Algebraic back projection:

Iterative process by redistributing the 2-dim image and considering the differences to the previous iteration step.



## Longitudinal Emittance using tomographic Reconstruction



# Tomography is medical image method Tomography:

2-dim reconstruction of sufficient 1-dim projections

### Application at accelerators:

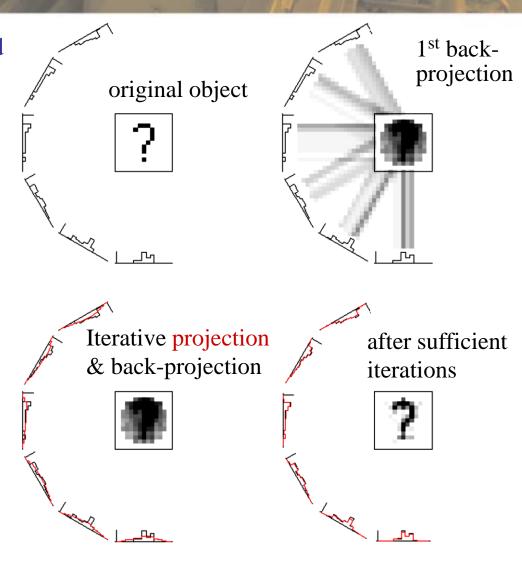
Longitudinal emittance evolution in synchrotrons.

#### Bunch observation:

Each revolution, the bunch shape changes a bit due to synchrotron oscillations. Fulfilled condition:  $f_{synch} << f_{ref}$ .

### Algebraic back projection:

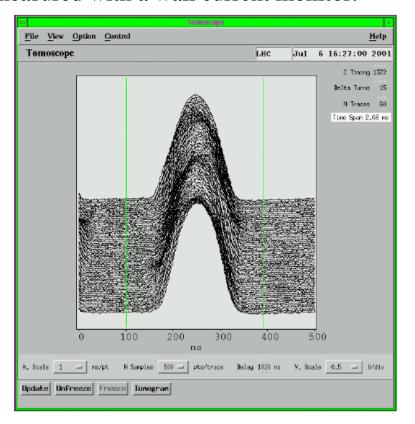
Iterative process by redistributing the 2-dim image and considering the differences to the previous iteration step.

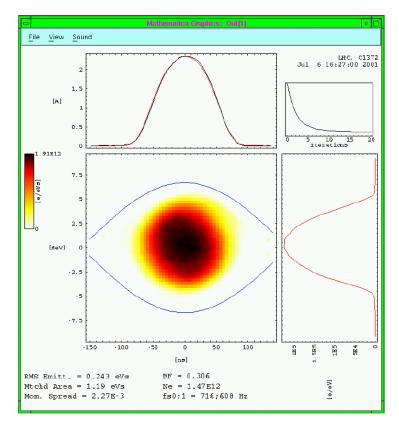


# Results of tomographic Reconstruction at a Synchrotron I



Bunches from 500 turns at the CERN PS and the phase space for the first time slice, measured with a wall current monitor:



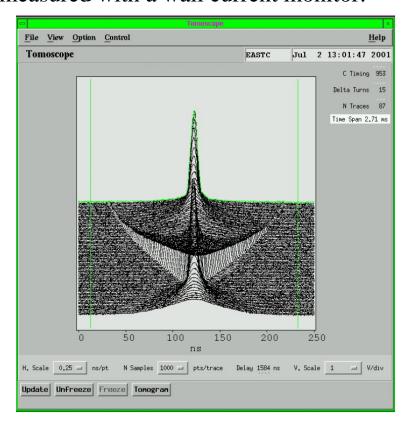


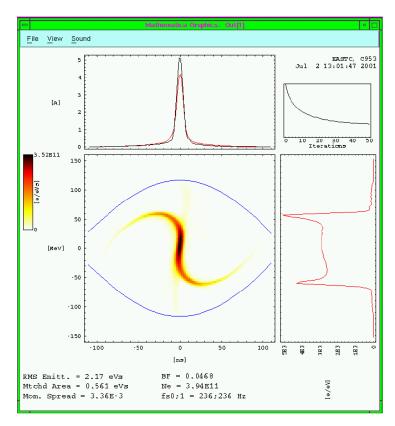
Typical bucket filling. Important knowledge for bunch 'gymnastics'.

# Results of tomographic Reconstruction at a Synchrotron II



Bunches from 500 turns at the CERN PS and the phase space for the first time slice, measured with a wall current monitor:





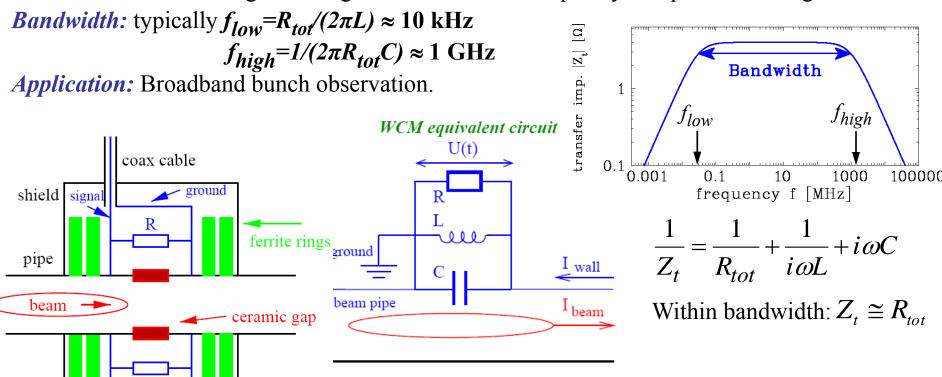
Mismatched bunch shown oscillations and filamentation due to 'bunch-rotation'.

#### Resistive Wall Current Monitor



#### Broadband observation of bunches can be performed with a resistive Wall Current Monitor

- **Principle:**  $\triangleright$  Ceramic gap bridged with n = 10...100 resistors of  $R = 10...100 \Omega$ 
  - $\triangleright$  Measurement of voltage drop for  $R_{tot} = R/n = 1...10 \Omega$
  - $\triangleright$  Ferrite rings with high  $L \rightarrow$  forces low frequency components through resistors



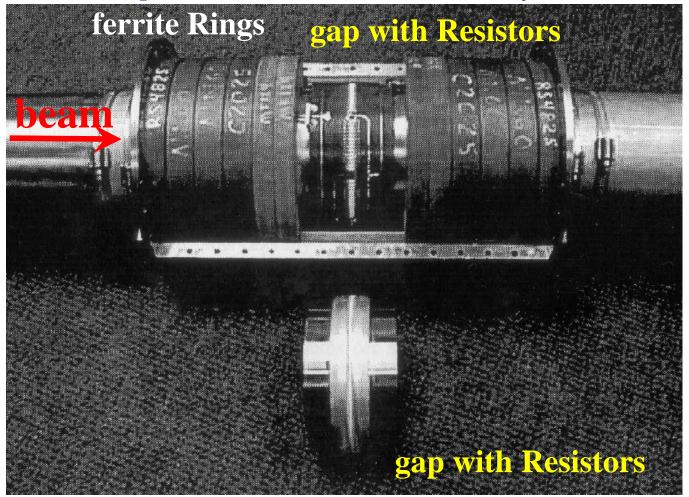
to ground

to signal

#### Resistive Wall Current Monitor



## **Example**: Realization at Fermi-Laboratory





### **Outline:**

- > Definition of longitudinal phase space
- ➤ Proton LINAC: Determination of mean energy used for alignment of cavities phase and amplitude
- ➤ Determination of longitudinal emittance LINAC: variation of bunch length Synchrotron: Topographic reconstruction
- ➤ Bunch length measurement for non-relativistic beams

  Determination of particle arrival
- > Bunch length measurement for relativistic beams
- > Summary

# Bunch Structure at low $E_{kin}$ : Not possible with Pick-Ups



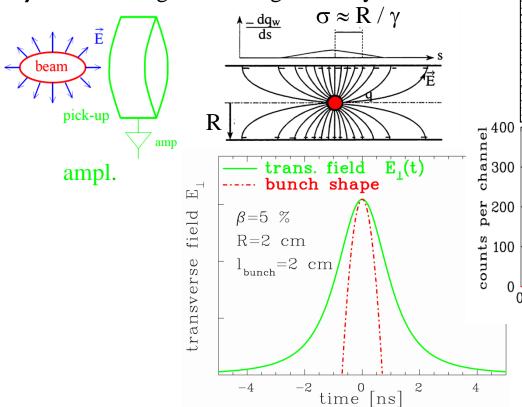
#### Pick-ups are used for:

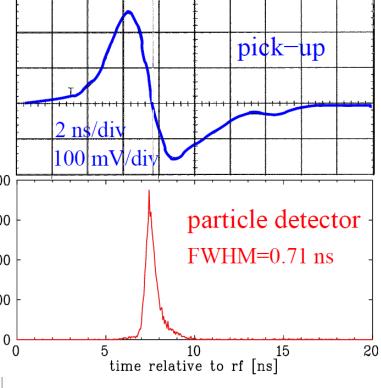
- > precise for bunch-center relative to rf
- > course image of bunch shape

*Example*: Comparison pick-up – particle counter:  $Ar^{1+}$  with 1.4 MeV/u ( $\beta = 5.5\%$ )

#### **But:**

For  $\beta \ll 1 \rightarrow \text{long. } E\text{-field significantly modified:}$ 





⇒ the pick-up signal is insensitive to bunch 'fine-structure'

## Low Velocity Effect: General Consideration

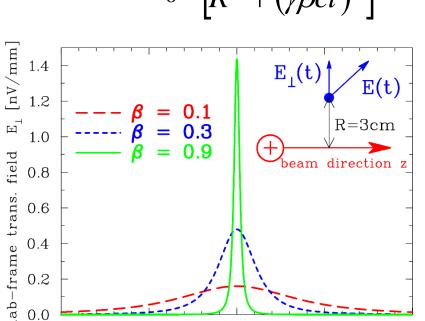


### Lorentz transformation of single point-like charge:

Lorentz boost and transformation of time:  $E_{\perp}(t) = \gamma \cdot E'_{\perp}(t')$  and  $t \rightarrow t'$ 

Trans.  $E_{\perp}$  lab.-frame of a point charge:

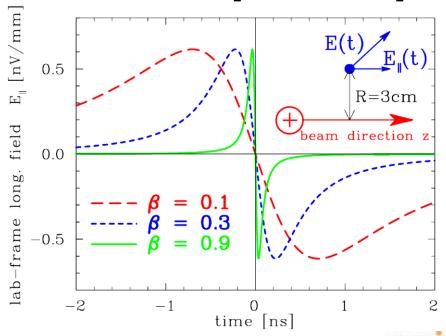
$$E_{\perp}(t) = \frac{e}{4\pi\varepsilon_0} \cdot \frac{\gamma R}{\left[R^2 + (\gamma \beta ct)^2\right]^{3/2}}$$



time [ns]

Long.  $E_{//}$  lab.-frame of a point charge:

$$E_{||}(t) = -\frac{e}{4\pi\varepsilon_0} \cdot \frac{\gamma\beta ct}{\left[R^2 + (\gamma\beta ct)^2\right]^{3/2}}$$



0.4

0.2

## Broadband coaxial Faraday Cups for Bunch Structure



The bunch structure can be observed with cups, having a bandwidth up to several GHz.

Bandwidth and rise time: BW [GHz] =  $0.3/t_{rise}$ [ns]

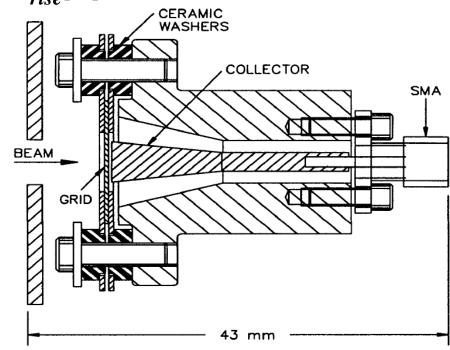
Impedance of a coaxial transmission line:

$$Z_0 = \frac{Z_c}{2\pi} \cdot \ln \frac{r_{\text{outer}}}{r_{\text{inner}}}$$

with 
$$Z_c = \sqrt{\frac{\mu_0 \mu_r}{\varepsilon_0 \varepsilon_r}}$$

→ impedance matching to prevent for reflections

Voltage reflection:  $\rho_V = \frac{Z - Z_0}{Z + Z_0}$ Voltage Standing Wave Ratio: VSWR  $= \frac{Z}{Z_0} = \frac{1 + \rho_V}{1 - \rho_V}$ 



$$=\frac{Z}{Z_0} = \frac{1+\rho_V}{1-\rho_V}$$

 $Z = Z_0$ : no reflection.  $Z = 0 \Rightarrow \rho_V = -1$ : short circuit.  $Z = \infty \Rightarrow \rho_V = 1$ : open circuit.

# Realization of a Broadband coaxial Faraday Cup





# Bunch Structure using secondary Electrons for low Ekin Protons

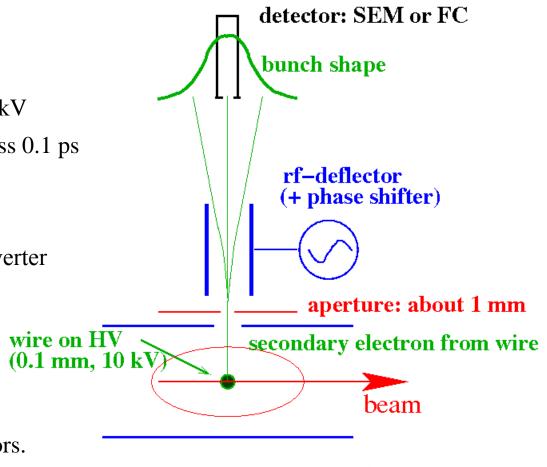


Secondary e<sup>-</sup> liberated from a wire carrying the time information.

→ Bunch Shape Monitor (BSM)

#### **Working principle:**

- $\triangleright$  insertion of a 0.1 mm wire at  $\approx$  10 kV
- > emission of secondary e within less 0.1 ps
- ➤ secondary e<sup>-</sup> are accelerated
- > toward an rf-deflector
- > rf-deflector as 'time-to-space' converter
- > detector with a thin slit
- > slow shift of the phase
- $\triangleright$  resolution  $\approx 1^{\circ} < 10 \text{ ps}$
- ➤ Measurements are comparable to that obtained with particle detectors.

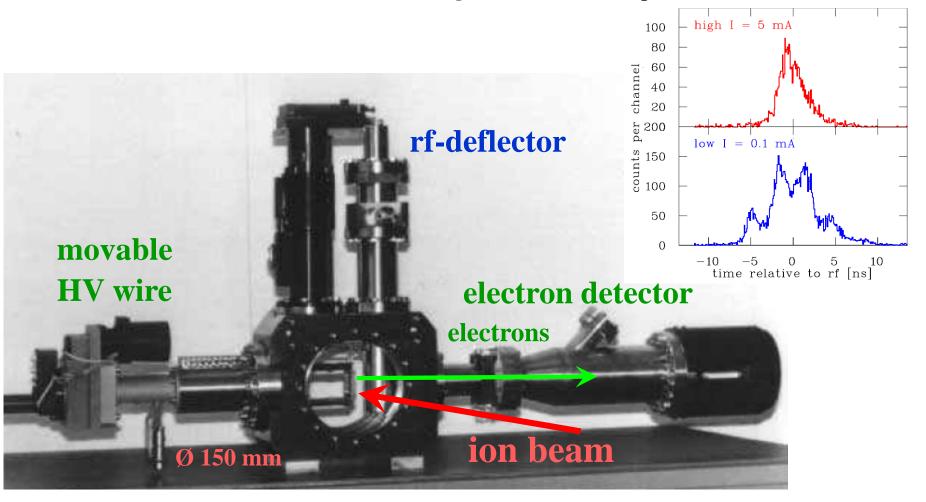


SEM: secondary electron multiplier

# Realization of Bunch Shape Monitor at CERN LINAC2



Example: The bunch shape at 120 keV/u for 120 keV/u:





### **Outline:**

- > Definition of longitudinal phase space
- > Proton LINAC: Determination of mean energy used for alignment of cavities phase and amplitude
- ➤ Determination of longitudinal emittance LINAC: variation of bunch length Synchrotron: Topographic reconstruction
- ➤ Bunch length measurement for non-relativistic beams Determination of particle arrival
- ➤ Bunch length measurement for relativistic beams

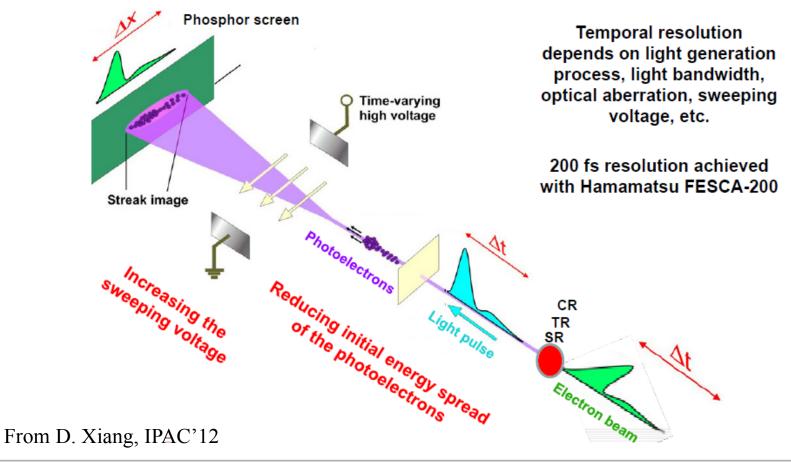
  Synchrotron light monitor and electro-optical modulation of a laser beam
- > Summary

## Bunch Length Measurement for relativistic e<sup>-</sup>



Electron bunches are too short ( $\sigma_t$  < 300 ps) to be covered by the bandwidth of pick-ups (f < 1 GHz  $\Leftrightarrow t_{rise}$  >300 ps) for structure determination.

 $\rightarrow$  Time resolved observation of synchr. light with a streak camera: Resolution  $\approx 1$  ps.

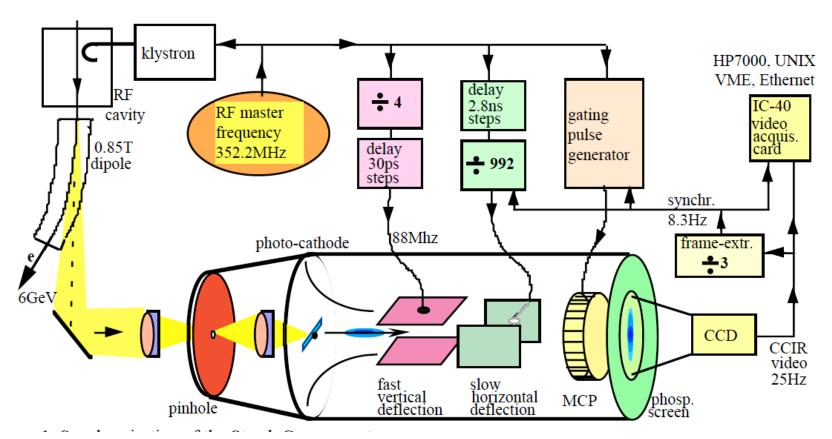


## Bunch Length Measurement for relativistic e<sup>-</sup>



Electron bunches are too short ( $\sigma_t$  < 300 ps) to be covered by the bandwidth of pick-ups (f < 1 GHz  $\Leftrightarrow t_{rise}$  >300 ps) for structure determination.

 $\rightarrow$  Time resolved observation of synchr. light with a streak camera: Resolution  $\approx 1$  ps.



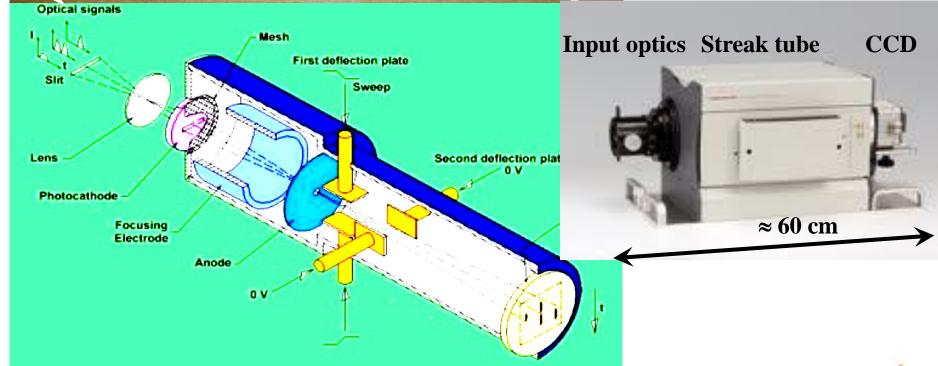
### Technical Realization of Streak Camera





Hardware of a streak camera

Time resolution down to 0.5 ps:



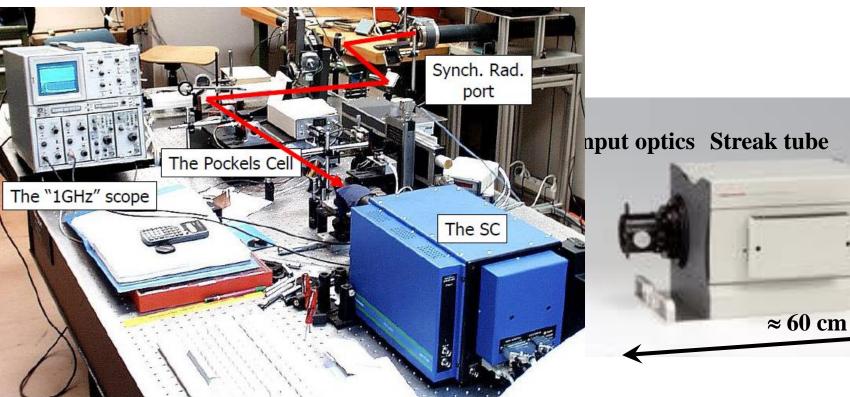
### Technical Realization of Streak Camera



**CCD** 

Hardware of a streak camera

Time resolution down to 0.5 ps:



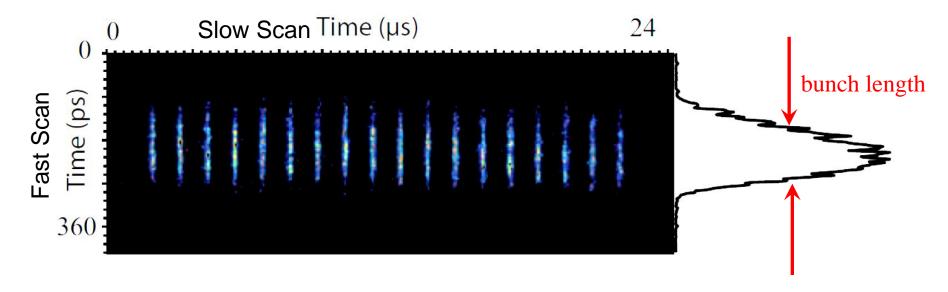
The Streak Camera setup at ELETTRA, Trieste, Italy

## Results of Bunch Length Measurement by a Streak Camera



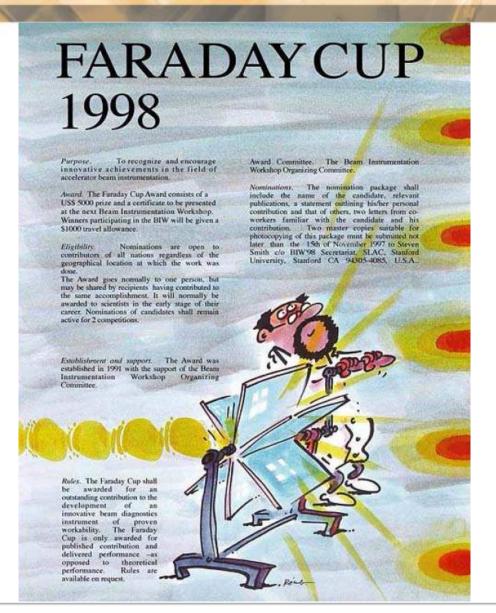
The streak camera delivers a fast scan in vertical direction (here 360 ps full scale) and a slower scan in horizontal direction (24  $\mu$ s).

*Example:* Bunch length at the synch. Light source SOLEIL for  $U_{rf} = 2$  MV for slow direction 24  $\mu$ s and scaling for fast scan 360 ps: measure  $\sigma_t = 35$  ps.



### The Artist View of a Streak Camera



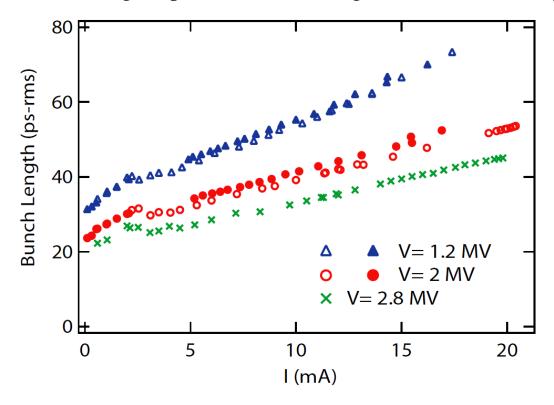


## The Importance of Bunch Length by Streak Camera



Short bunches are desired by the synchrotron light users for time resolved spectroscopy. The bunch focusing is changed by the rf-amplitude.

*Example:* Bunch length  $\sigma_t$  as a function of stored current (space-charge de-focusing, impedance broadening) for different rf-amplitudes at SOLEIL:



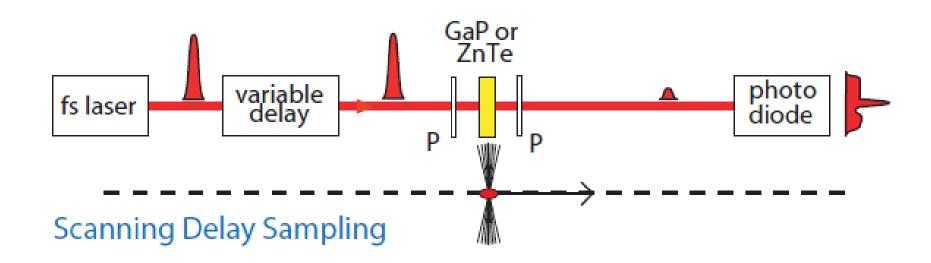
## Bunch length measurement by electro-optical methods



#### For Free Electron Lasers → bunch length below 1 ps is achieved

- → below resolution of streak camera
- $\rightarrow$  short laser pulses with  $t \approx 10$  fs and electro-optical modulator

Electro optical modulator: birefringent, rotation angle depends on external electric field Relativistic electron bunches: transverse field  $E_{\perp, lab} = \gamma E_{\perp, rest}$  carries the time information Scanning of delay between bunch and laser  $\rightarrow$  time profile after several pulses.



From S.P.Jamison et al., EPAC 2006

## Pockels Effect



Crystal with different index of refractivity on orthogonal planes

- $\Rightarrow$  Bire-fringent
- ⇒ external electric field changes the phase between orthogonal wave:

'Pockels effect' with polarization  $P = \varepsilon_0 \chi E$ 

## **Diagnostics method:**

Probing bunch's electric field by polarized laser beam

Polarized Input

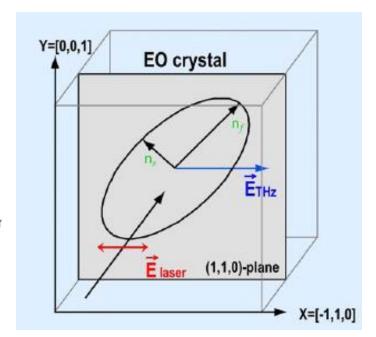
Polarized Input

Notitage

Crystal Ring Electrodes

Unpolarized Light

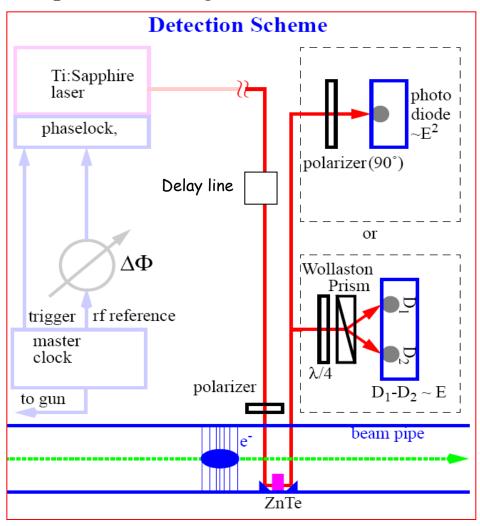
Optical Aperture in Electrode

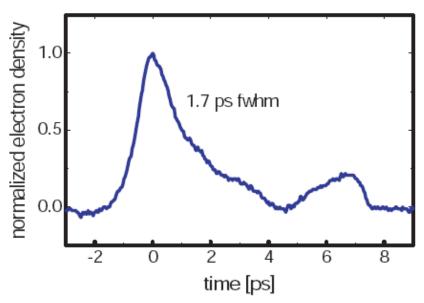


## Realization of EOS Scanning



## Setup of a scanning EOS method



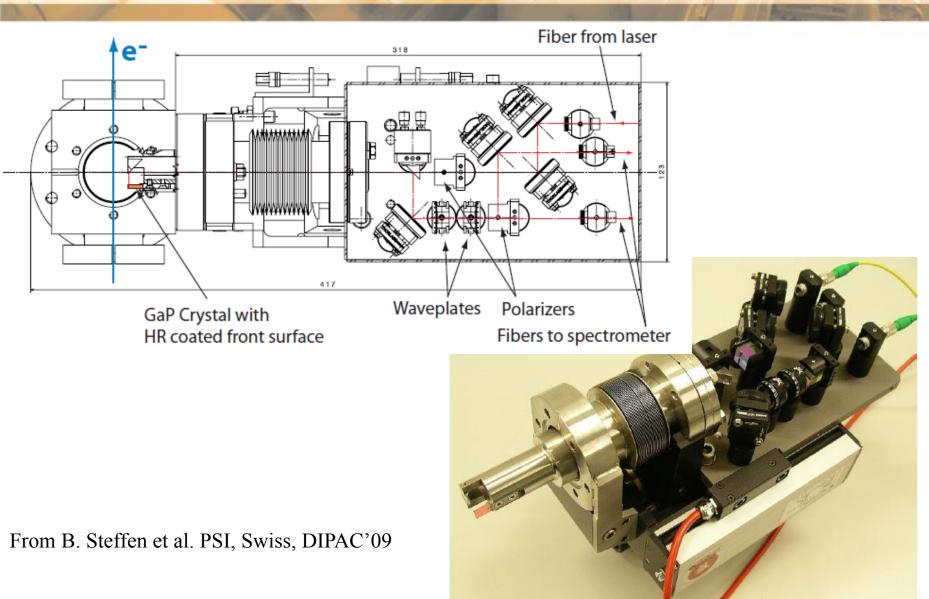


Using 12fs Ti:Al2O3 laser at 800nm and ZnTe crystal 0.5mm thick and a beam of 46MeV, 200pC, 2ps.

X. Yan et al, PRL 85, 3404 (2000)

## Hardware of a compact EOS Scanning Setup

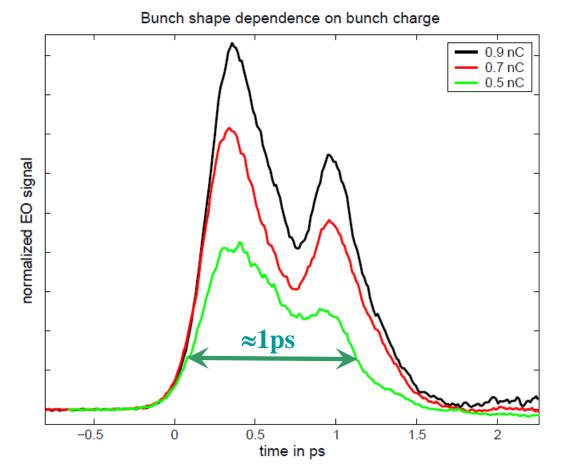




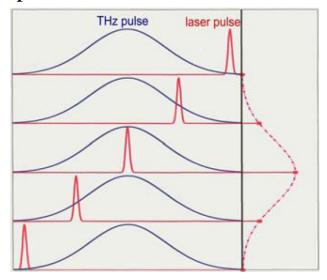
## Measurement of Bunch Shape at FEL-Facility



#### **Example**: Bunch length at FEL test facility FLASH



# Scanning of the short laser pulse relative to bunch:

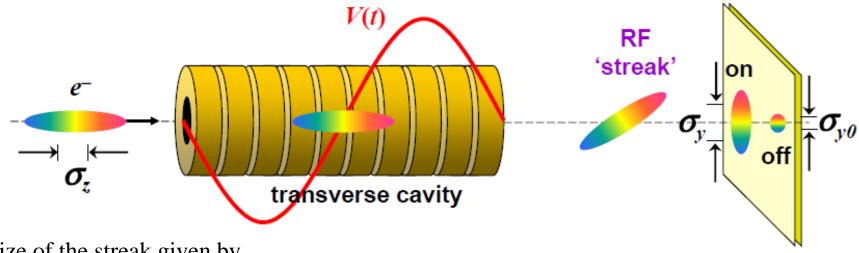


Results at FLASH, Hamburg, see B. Steffen et al., FEL Conf. Stanford, p. 549, 2005.

## Bunch Length by rf-Deflection: Principle



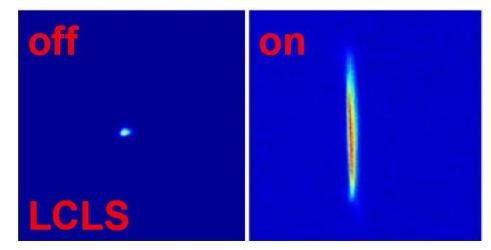
## Transversal deflection of the bunch i.e. time-to-space conversion



Size of the streak given by

$$\sigma_{y} = \sqrt{\sigma_{y0}^{2} + R_{35} \cdot k \cdot \sigma_{z}^{2}}$$

k is determined by the rf-power 
$$k = \frac{2\pi e \cdot U_{rf}}{\lambda_{rf} E}$$



From D. Xiang, IPAC'12

## Bunch Length by rf-Deflection: Hardware



Transversal deflection of the bunch

i.e. time-to-space conversion

Example: Cavity at FERMI, Trieste, Italy

Typical beam size
Length
Frequency
Max. rf power
Total trans. volt.
Time resolution

Beam energy

320 MeV 0.2 mm 0.5 m 2.998 GHz 5 MW 4.9 MV

70 fs

Bunch Length from Injector (6 ps fwhm)

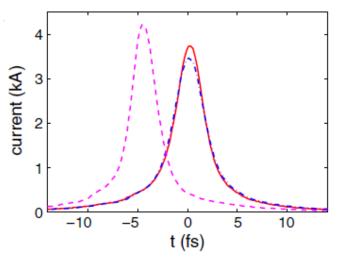
Bunch I

Bunch length compression (1ps fwhm)

From M. Veronese, BIW'12

## Bunch Length by rf-Deflection: Results



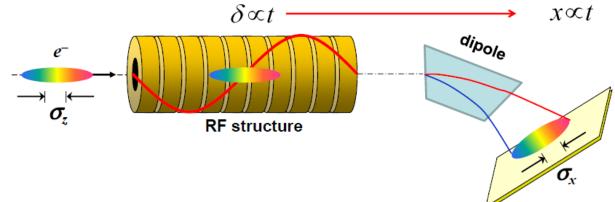


0.3 0.2 (%) 0.1 -0.1 -0.2 -0.3 -10 -5 0 5 10 t (fs) Example: Short bunch generation at Free Electron Laser Facility LCLS, Stanford, USA

Time resolution down to (true) fs range!

**Vertical deflection** with cavity → time spectum **Horizontal bending** 

- →energy (momentum) distribution via dispersion
- ⇒ Longitudinal phase space determined!



From D. Xiang, PRST-AB, 13, 094001 (2010)

## Summary of longitudinal Measurements



### **Longitudinal** ↔ **transverse correspondences**:

- $\triangleright$  position relative to rf  $\leftrightarrow$  transverse center-of-mass
- ➤ bunch structure in time ↔ transverse profile in space
- $\triangleright$  momentum or energy spread  $\leftrightarrow$  transverse divergence.

#### **Determination uses:**

- *Broadband pick-ups:* ➤ position relative to rf, mean energy
  - right emittance at transfer lines or synchrotron via tomography assumption: bunches longer than pick-up.
- Particle detectors:
- ➤ TOF or secondary e<sup>-</sup> from wire
  - → for non-relativistic proton beams reason: *E*-field does not reflect bunch shape.

- Streak cameras:
- time resolved monitoring of synchrotron radiation
  - $\rightarrow$  for relativistic e<sup>-</sup>-beams,  $t_{bunch} < 1$ ns reason: too short bunches for rf electronics.

- Laser scanning:
- > Electro-optical modulation of short laser pulse
  - $\rightarrow$  highest time resolution.