

Testing fundamental physical principles with entangled neutral K mesons



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Fundamental tests possible with entangled neutral kaons

- a) Test of Quantum coherence** Bertlmann, Grimus, Hiesmayr PR D60 (1999) 114032
KLOE coll. PLB 642(2006) 315, FP40 (2010) 852
CPLEAR PLB 422 (1999) 339
- b) Test of CPT symmetry + Quantum coherence**
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- c) Test of Lorentz and CPT symmetry**
Kostelecky PRD61 (1999) 016002, PRD64 (2001) 076001
KLOE coll. PLB 730 (2014) 89
- d) Direct Test of T (time-reversal) symmetry**
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- e) Bell's inequality test**
Hiesmayr, A.D.D. et al. EPJC (2012) 72:1856
- f) Kaonic quantum eraser (Bohr's complementarity)**
Bramon, Garbarino, Hiesmayr PRL (2004) 020405
- g) Test of collapse models**
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T, CP and CPT violation parameters for neutral kaons

$$|\Psi\rangle = a(t)|K^0\rangle + b(t)|\bar{K}^0\rangle \quad i\frac{\partial}{\partial t}|\Psi\rangle = \mathbf{H}|\Psi\rangle \quad \mathbf{H} = \mathbf{M} - \frac{i}{2}\mathbf{\Gamma} = \begin{pmatrix} m_{11} & m_{12} \\ m_{21} & m_{22} \end{pmatrix} - \frac{i}{2}\begin{pmatrix} \Gamma_{11} & \Gamma_{12} \\ \Gamma_{21} & \Gamma_{22} \end{pmatrix}$$

$$|K_{S,L}\rangle = \frac{1}{\sqrt{2(1+|\varepsilon_{S,L}|^2)}} \left[(1 + \varepsilon_{S,L})|K^0\rangle \pm (1 - \varepsilon_{S,L})|\bar{K}^0\rangle \right] \quad \lambda_S = m_S - \frac{i}{2}\Gamma_S \quad , \quad \lambda_L = m_L - \frac{i}{2}\Gamma_L$$

$$|K_{S,L}(t)\rangle = e^{-i\lambda_{S,L}t} |K_{S,L}(0)\rangle$$

T, CP and CPT violation parameters:

T viol. $\varepsilon = \frac{H_{12} - H_{21}}{2(\lambda_S - \lambda_L)} = \frac{-i\Im m_{12} - \Im \Gamma_{12}/2}{\Delta m + i\Delta\Gamma/2}$

CP viol. $\varepsilon_{S,L} = \varepsilon \pm \delta$

CPT viol. $\delta = \frac{H_{11} - H_{22}}{2(\lambda_S - \lambda_L)} = \frac{1}{2} \frac{(m_{22} - m_{11}) - (i/2)(\Gamma_{22} - \Gamma_{11})}{\Delta m + i\Delta\Gamma/2}$

$|\varepsilon| \cong 2.232 \times 10^{-3}$

CPT violation: $|\delta| < \sim 10^{-4} \Rightarrow \left| \frac{m_{K^0} - m_{\bar{K}^0}}{m_K} \right| < 10^{-18}$

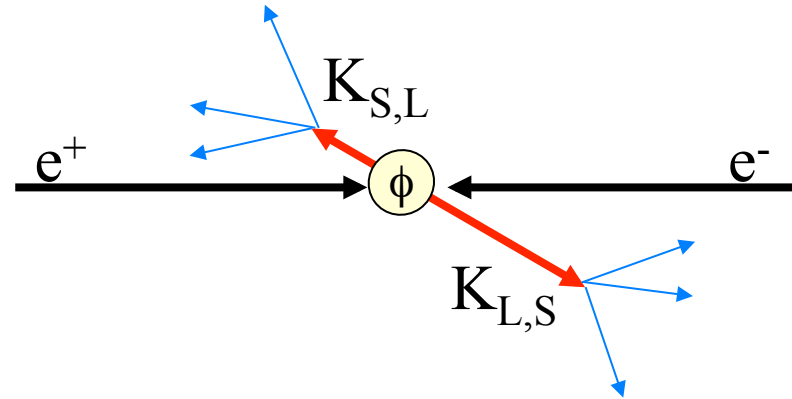
Neutral kaons at a ϕ -factory

Production of the vector meson ϕ in e^+e^- annihilations:

- $e^+e^- \rightarrow \phi$ $\sigma_\phi \sim 3 \mu\text{b}$
 $W = m_\phi = 1019.4 \text{ MeV}$
- $\text{BR}(\phi \rightarrow K^0\bar{K}^0) \sim 34\%$
- $\sim 10^6$ neutral kaon pairs per pb^{-1} produced in an antisymmetric quantum state with $J^{PC} = 1^{--}$:

$$\mathbf{p}_K = 110 \text{ MeV}/c$$

$$\lambda_S = 6 \text{ mm} \quad \lambda_L = 3.5 \text{ m}$$



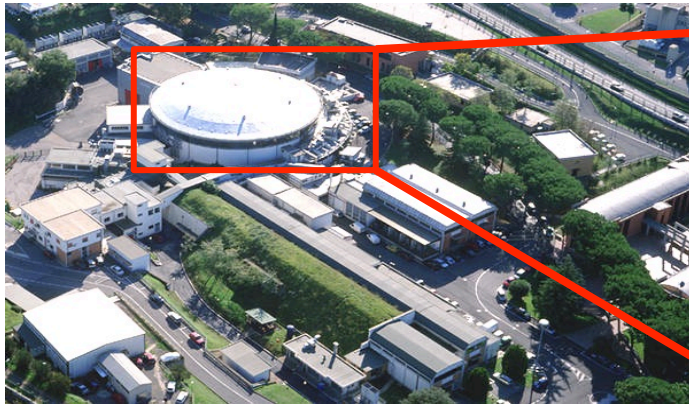
$$|i\rangle = \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right]$$

$$= \frac{N}{\sqrt{2}} \left[|K_S(\vec{p})\rangle |K_L(-\vec{p})\rangle - |K_L(\vec{p})\rangle |K_S(-\vec{p})\rangle \right]$$

$$N = \sqrt{(1 + |\varepsilon_S|^2)(1 + |\varepsilon_L|^2)} / (1 - \varepsilon_S \varepsilon_L) \cong 1$$

The KLOE detector at the Frascati ϕ -factory DAFNE

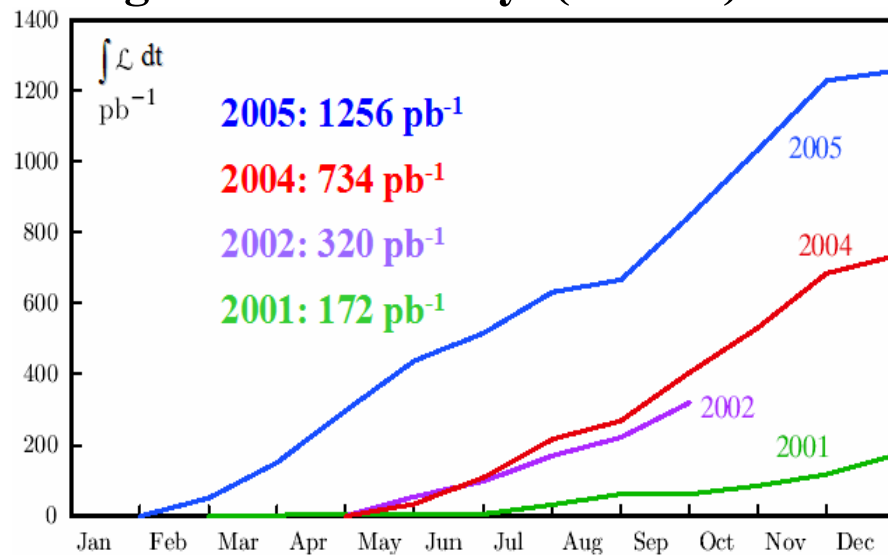
DAFNE
collider



KLOE detector



Integrated luminosity (KLOE)



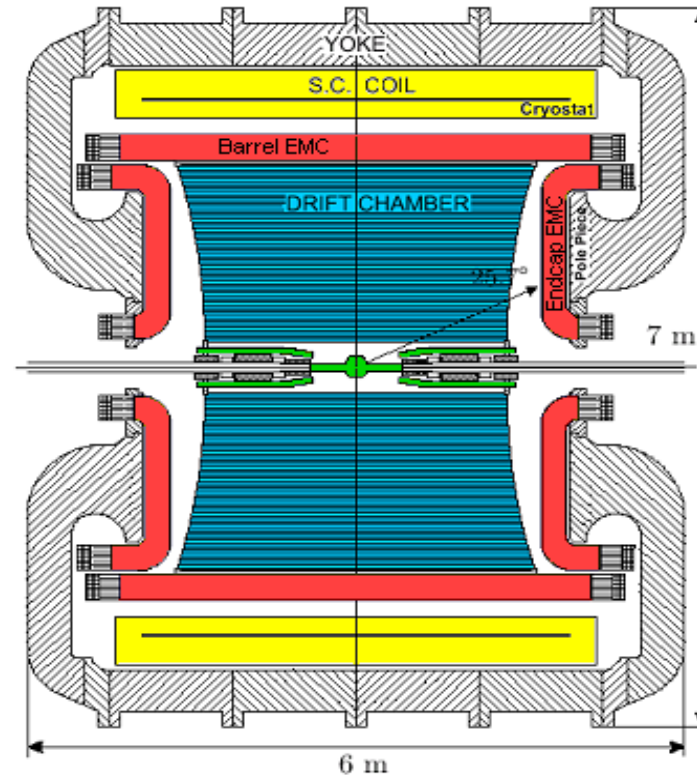
Total KLOE $\int \mathcal{L} dt \sim 2.5 \text{ fb}^{-1}$
(2001 - 05) $\rightarrow \sim 2.5 \times 10^9 \text{ K}_S \text{K}_L \text{ pairs}$

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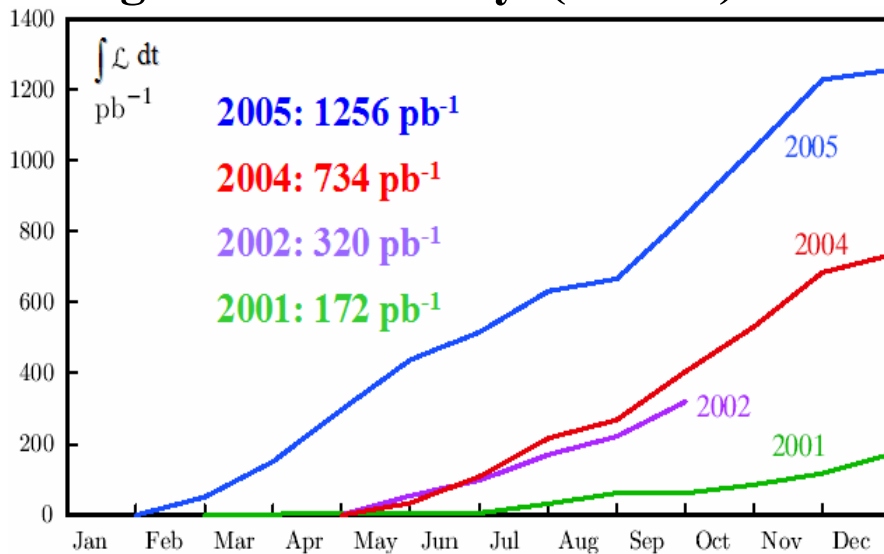
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KLOE detector



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Lead/scintillating fiber calorimeter
 drift chamber
 4 m diameter \times 3.3 m length
 helium based gas mixture

Neutral kaon interferometry

$$|i\rangle = \frac{N}{\sqrt{2}} \left[|K_S(\vec{p})\rangle |K_L(-\vec{p})\rangle - |K_L(\vec{p})\rangle |K_S(-\vec{p})\rangle \right]$$

Double differential time distribution:

$$I(f_1, t_1; f_2, t_2) = C_{12} \left\{ |\eta_1|^2 e^{-\Gamma_L t_1 - \Gamma_S t_2} + |\eta_2|^2 e^{-\Gamma_S t_1 - \Gamma_L t_2} \right.$$

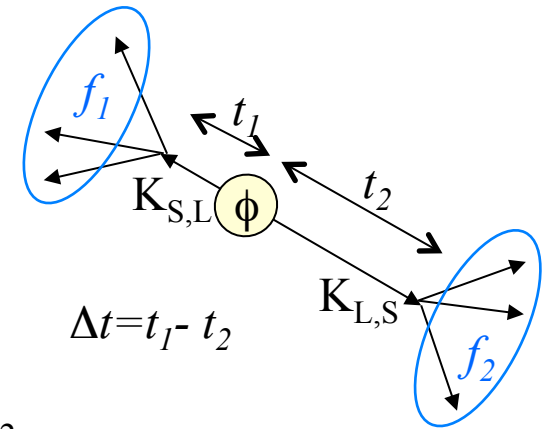
$$\left. - 2|\eta_1||\eta_2| e^{-(\Gamma_S + \Gamma_L)(t_1 + t_2)/2} \cos \left[\Delta m(t_2 - t_1) + \phi_1 - \phi_2 \right] \right\}$$

where $t_1(t_2)$ is the proper time of one (the other) kaon decay into f_1 (f_2) final state and:

$$\eta_i = |\eta_i| e^{i\phi_i} = \langle f_i | T | K_L \rangle / \langle f_i | T | K_S \rangle$$

$$C_{12} = \frac{|N|^2}{2} \left| \langle f_1 | T | K_S \rangle \langle f_2 | T | K_S \rangle \right|^2$$

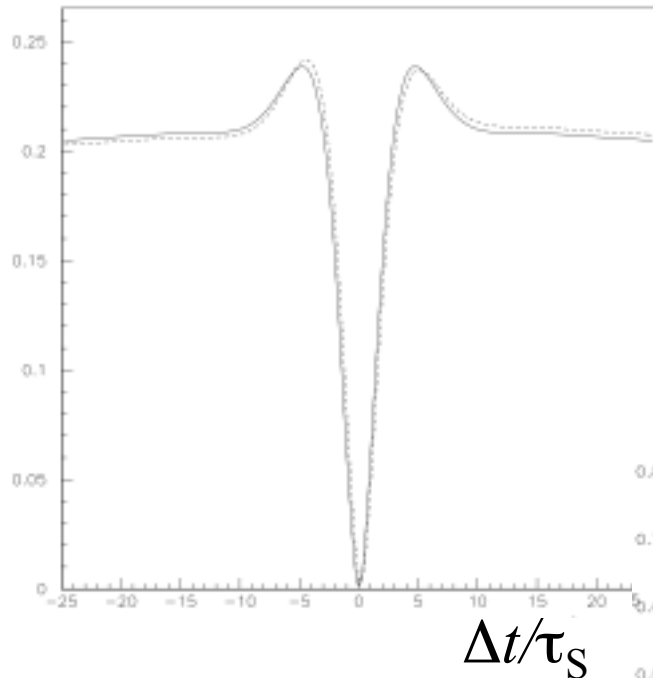
From these distributions for various final states f_i one can measure the following quantities: Γ_S , Γ_L , Δm , $|\eta_i|$, $\phi_i \equiv \arg(\eta_i)$



**characteristic interference term
at a ϕ -factory \Rightarrow interferometry**

Neutral kaon interferometry: main observables

$I(\Delta t)$ (a.u)

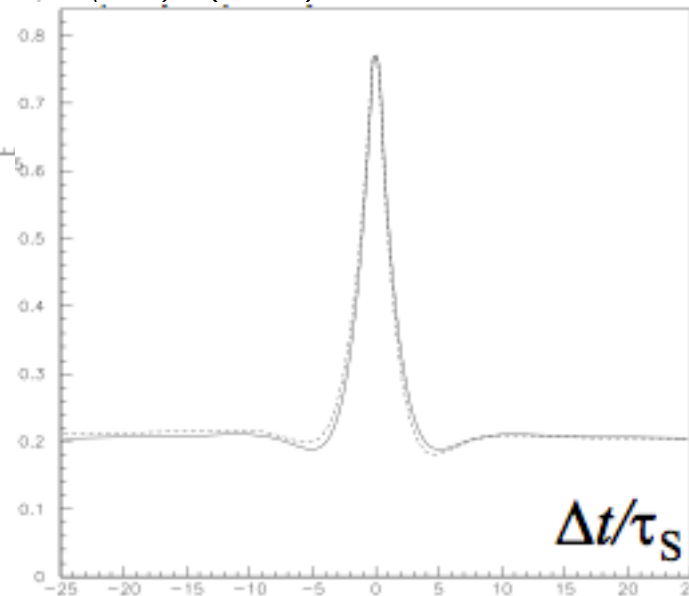


$$\Re\delta + \Re x_-$$

$$\Im\delta + \Im x_+$$

$$\phi \rightarrow K_S K_L \rightarrow \pi^+ \ell^- \bar{\nu} \quad \pi^- \ell^+ \nu$$

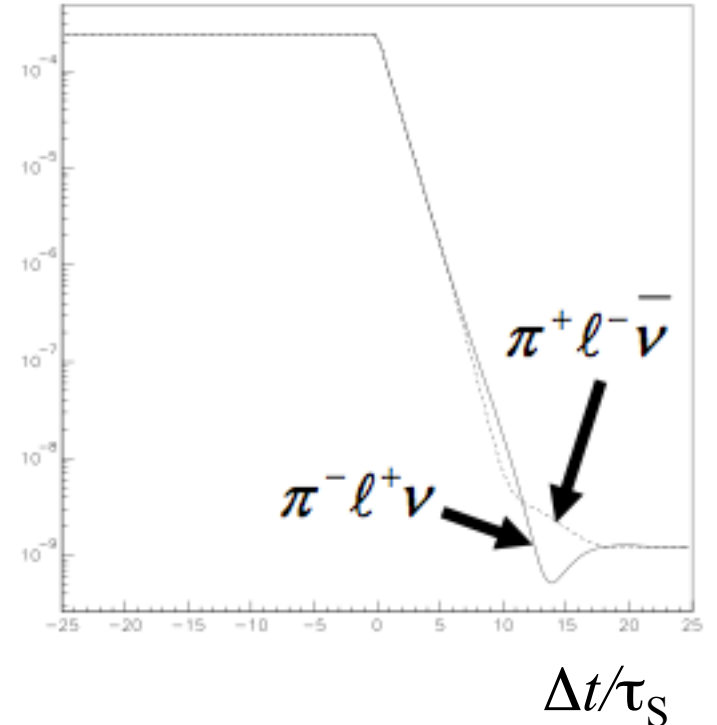
$I(\Delta t)$ (a.u)



$$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^0 \pi^0$$

$$\Re\left(\frac{\varepsilon'}{\varepsilon}\right) \quad \Im\left(\frac{\varepsilon'}{\varepsilon}\right)$$

$I(\Delta t)$ (a.u)



$$\phi \rightarrow K_S K_L \rightarrow \pi\pi \quad \pi\ell\nu$$

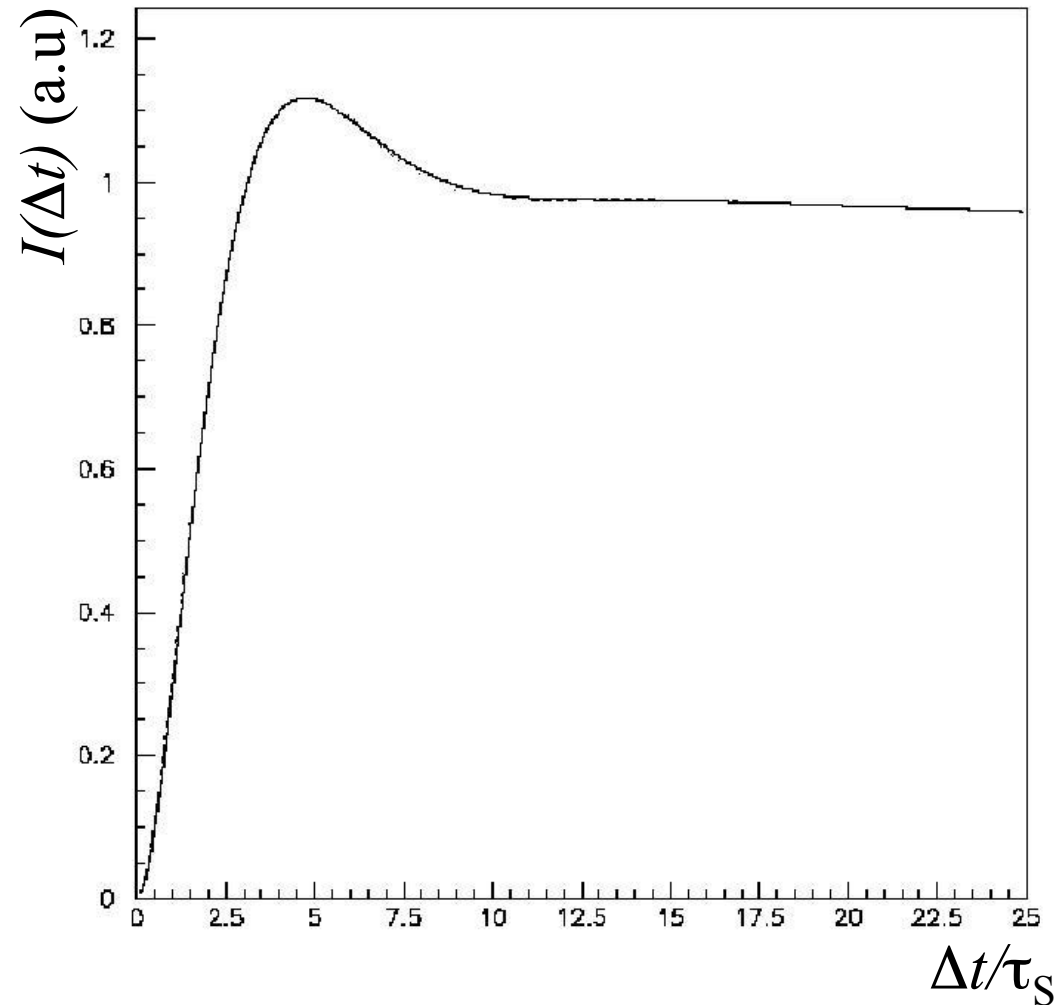
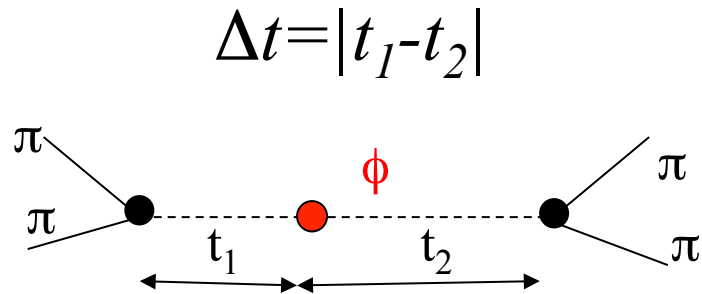
$$A_L = 2\Re\varepsilon - \Re\delta - \Re y - \Re x_-$$

$$\phi_{\pi\pi}$$

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$

$$|i\rangle = \frac{1}{\sqrt{2}} \left[|K^0\rangle |\bar{K}^0\rangle - |\bar{K}^0\rangle |K^0\rangle \right]$$

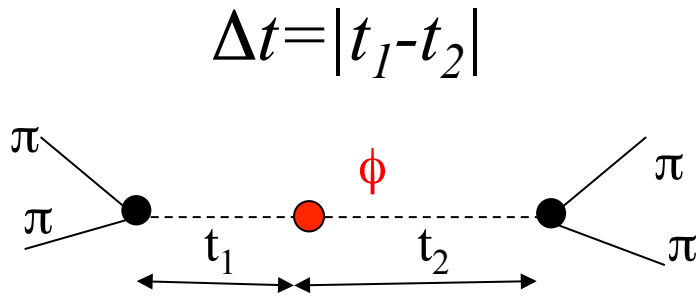
Same final state for both kaons: $f_1 = f_2 = \pi^+ \pi^-$



$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$

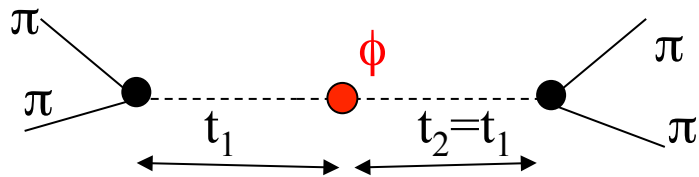
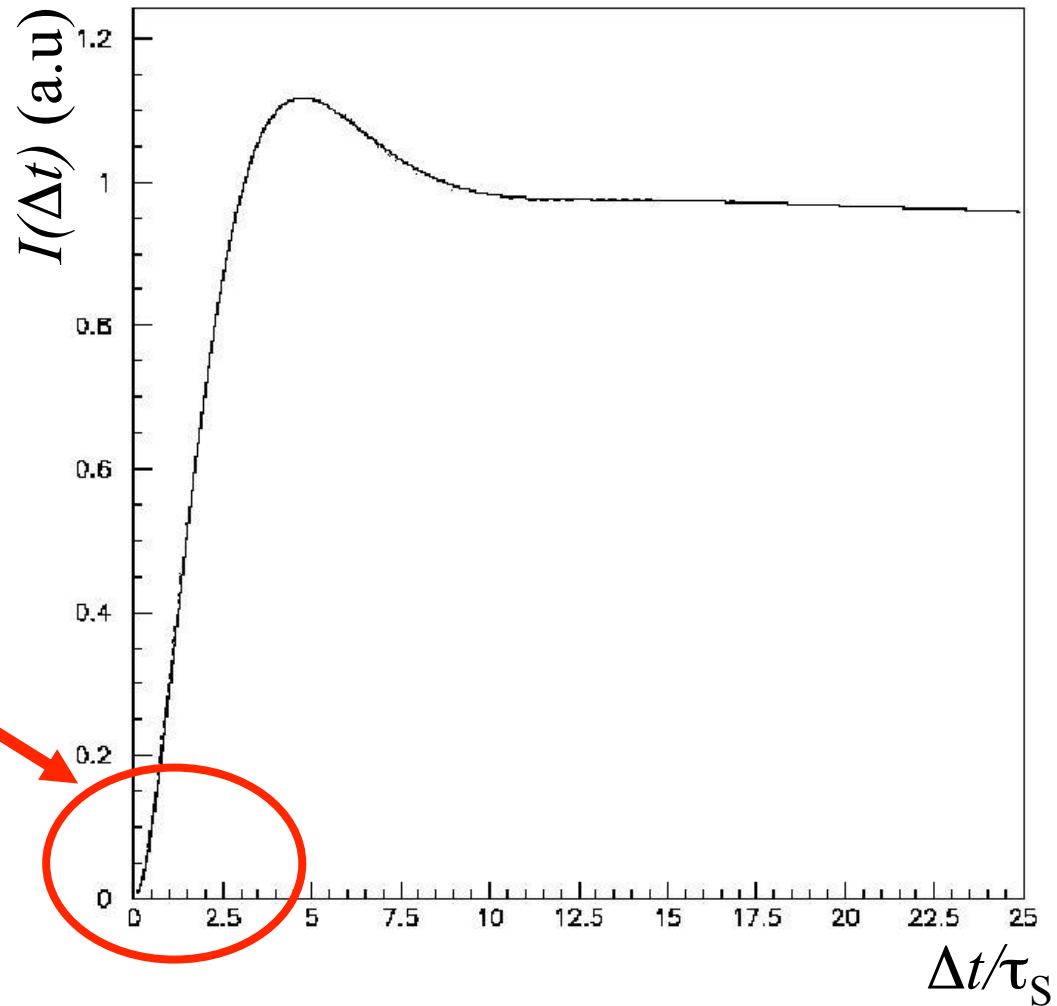
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EPR correlation:

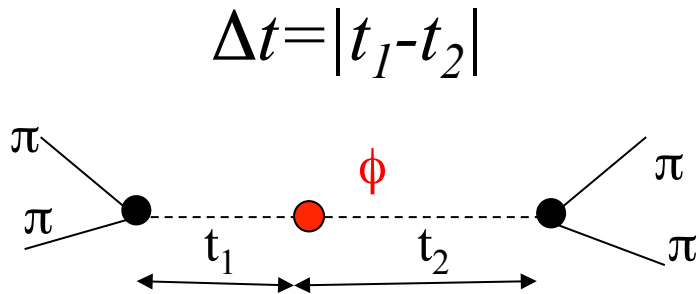
no simultaneous decays
($\Delta t=0$) in the same
final state due to the
destructive
quantum interference



$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$

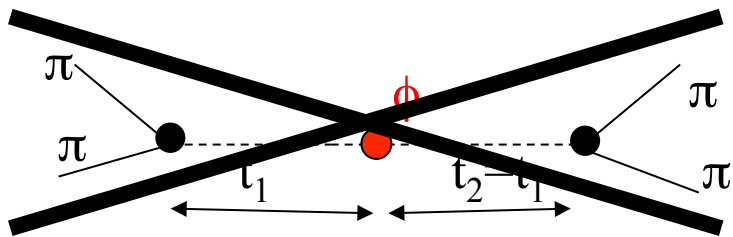
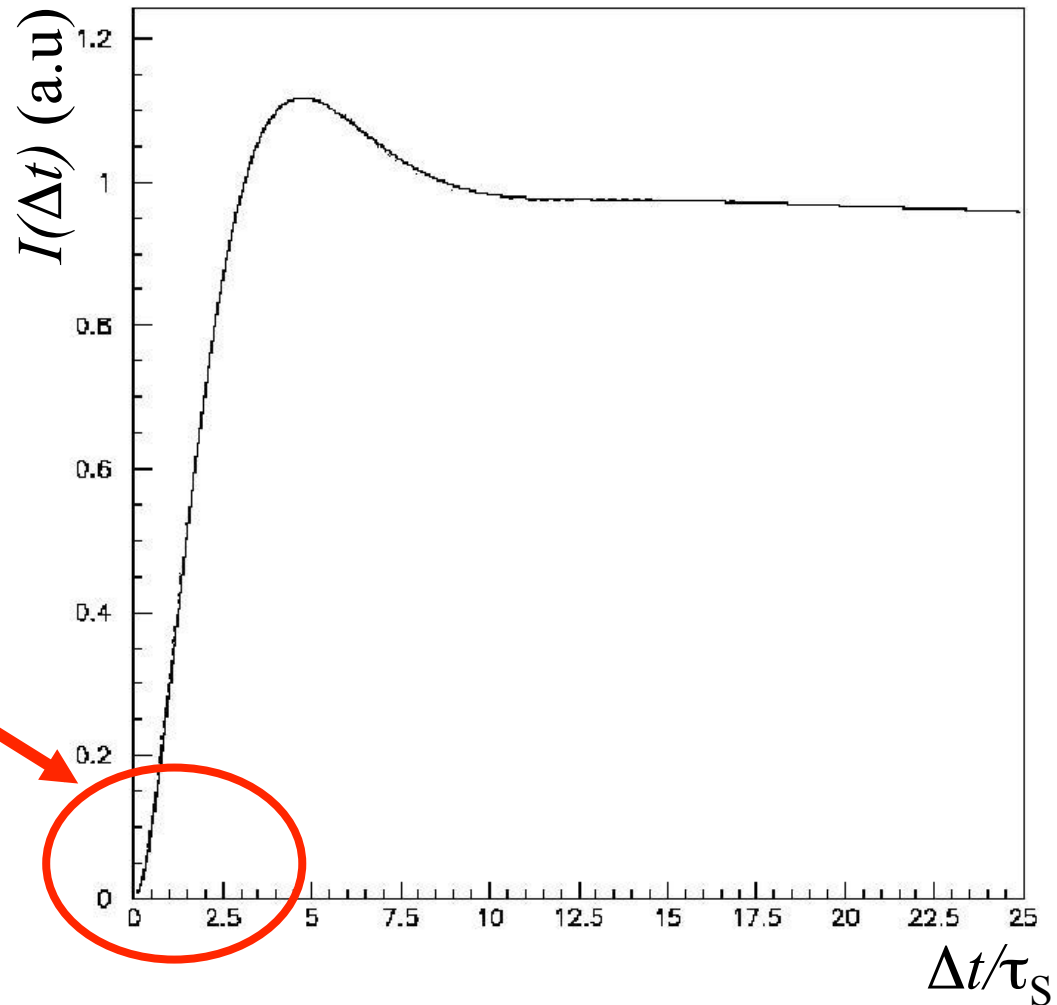
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**test of quantum coherence
(or search for decoherence effects)**

$\phi \rightarrow \mathbf{K}_S \mathbf{K}_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: test of quantum coherence

$$|i\rangle = \frac{1}{\sqrt{2}} \left[|K^0\rangle |\bar{K}^0\rangle - |\bar{K}^0\rangle |K^0\rangle \right]$$

$$I(\pi^+ \pi^-, \pi^+ \pi^-; \Delta t) = \frac{N}{2} \left[\left| \langle \pi^+ \pi^-, \pi^+ \pi^- | K^0 \bar{K}^0(\Delta t) \rangle \right|^2 + \left| \langle \pi^+ \pi^-, \pi^+ \pi^- | \bar{K}^0 K^0(\Delta t) \rangle \right|^2 - 2\Re \left(\langle \pi^+ \pi^-, \pi^+ \pi^- | K^0 \bar{K}^0(\Delta t) \rangle \langle \pi^+ \pi^-, \pi^+ \pi^- | \bar{K}^0 K^0(\Delta t) \rangle^* \right) \right]$$

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$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: test of quantum coherence

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Decoherence parameter:

$$\xi_{00} = 0 \quad \rightarrow \quad \text{QM}$$

$$\xi_{00} = 1 \quad \rightarrow \quad \text{total decoherence}$$

(also known as Furry's hypothesis
or spontaneous factorization)

[W.Furry, PR 49 (1936) 393]

Bertlmann, Grimus, Hiesmayr PR D60 (1999) 114032

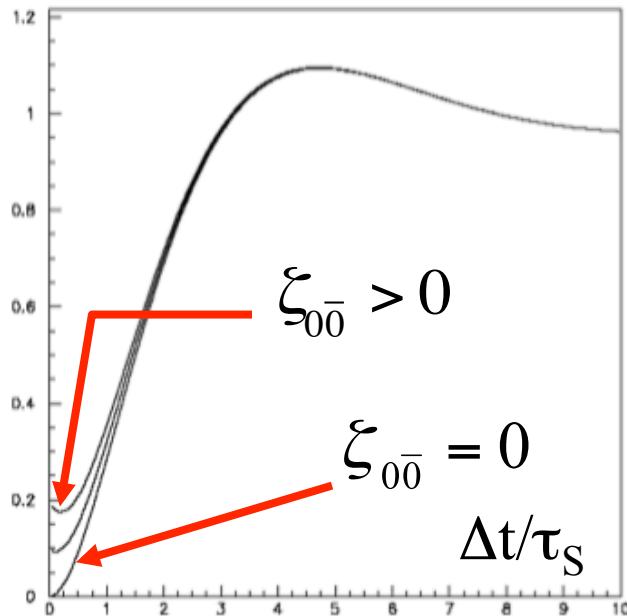
Bertlmann, Durstberger, Hiesmayr PRA 68 012111 (2003)

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: test of quantum coherence

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$I(\Delta t)$ (a.u.)



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$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: test of quantum coherence

- Analysed data: $L=1.5 \text{ fb}^{-1}$
- Fit including Δt resolution and efficiency effects + regeneration

KLOE result: [PLB 642\(2006\) 315](#)
[Found. Phys. 40 \(2010\) 852](#)

$$\xi_{0\bar{0}} = (1.4 \pm 9.5_{\text{STAT}} \pm 3.8_{\text{SYST}}) \times 10^{-7}$$

Observable suppressed by CP

violation: $|\eta_{+-}|^2 \sim |\epsilon|^2 \sim 10^{-6}$

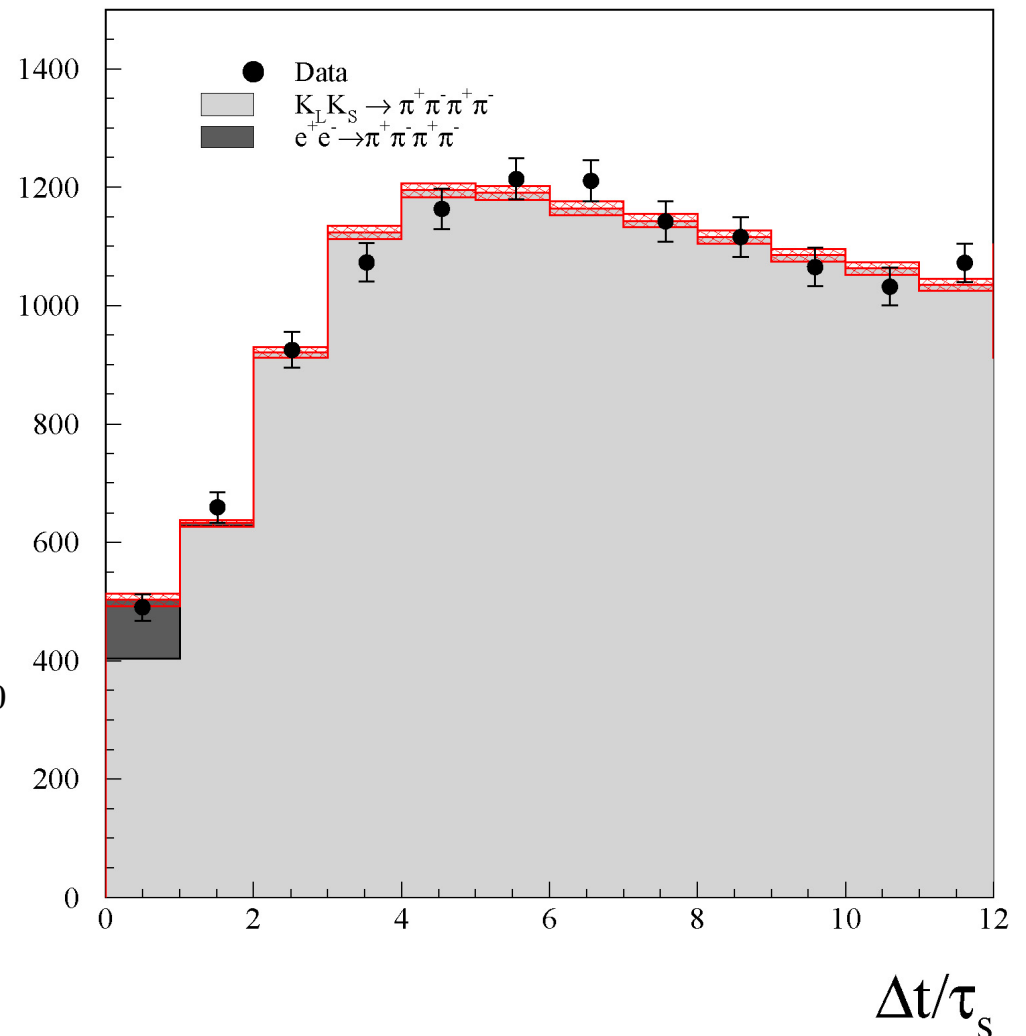
\Rightarrow terms $\xi_{00}/|\eta_{+-}|^2 \Rightarrow$ high sensitivity to ξ_{00}

From CPLEAR data, Bertlmann et al.
 (PR D60 (1999) 114032) obtain:

$$\xi_{0\bar{0}} = 0.4 \pm 0.7$$

In the B-meson system, BELLE coll.
 (PRL 99 (2007) 131802) obtains:

$$\xi_{0\bar{0}}^B = 0.029 \pm 0.057$$



$\Delta t/\tau_s$

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: test of quantum coherence

- Analysed data: $L=1.5 \text{ fb}^{-1}$
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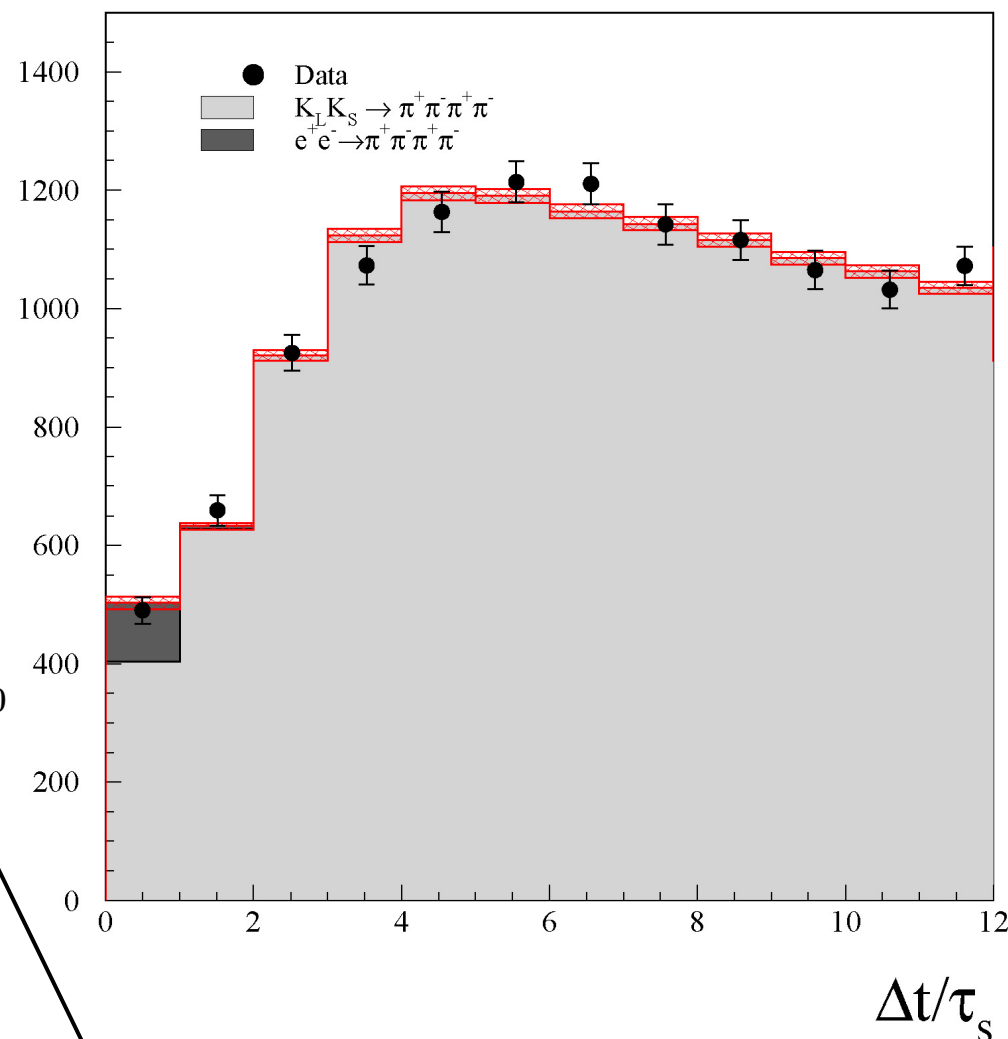
Observable suppressed by CP violation: $|\eta_{+-}|^2 \sim |\epsilon|^2 \sim 10^{-6}$
 \Rightarrow terms $\xi_{00}/|\eta_{+-}|^2 \Rightarrow$ high sensitivity to ξ_{00}

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Best precision achievable in an entangled system

Search for decoherence and CPT violation effects

Decoherence and CPT violation

Modified Liouville – von Neumann equation for the density matrix of the kaon system:

$$\dot{\rho}(t) = \underbrace{-iH\rho + i\rho H^+}_{\text{QM}} + L(\rho)$$

← extra term inducing decoherence:
pure state => mixed state

Possible decoherence due quantum gravity effects:

Black hole information loss paradox => Possible decoherence near a black hole.

Hawking [1] suggested that at a microscopic level, in a quantum gravity picture, non-trivial space-time fluctuations (generically space-time foam) could give rise to decoherence effects, which would necessarily entail a violation of CPT [2].

J. Ellis et al.[3-6] => model of decoherence for neutral kaons => 3 new CPTV param. α, β, γ :

$$L(\rho) = L(\rho; \alpha, \beta, \gamma)$$

$$\alpha, \gamma > 0 \quad , \quad \alpha\gamma > \beta^2$$

$$\text{At most: } \alpha, \beta, \gamma = O\left(\frac{M_K^2}{M_{\text{PLANCK}}}\right) \approx 2 \times 10^{-20} \text{ GeV}$$

[1] Hawking, Comm.Math.Phys.87 (1982) 395; [2] Wald, PR D21 (1980) 2742; [3] Ellis et. al, NP B241 (1984) 381; PRD53 (1996)3846 [4] Huet, Peskin, NP B434 (1995) 3; [5] Benatti, Floreanini, NPB511 (1998) 550 [6] Bernabeu, Ellis, Mavromatos, Nanopoulos, Papavassiliou: Handbook on kaon interferometry [hep-ph/0607322]

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: decoherence and CPT violation

Study of time evolution of **single kaons** decaying in $\pi^+ \pi^-$ and semileptonic final state

CPLEAR **PLB 364, 239 (1999)**

$$\alpha = (-0.5 \pm 2.8) \times 10^{-17} \text{ GeV}$$

$$\beta = (2.5 \pm 2.3) \times 10^{-19} \text{ GeV}$$

$$\gamma = (1.1 \pm 2.5) \times 10^{-21} \text{ GeV}$$

single kaons

In the complete positivity hypothesis

$$\alpha = \gamma, \quad \beta = 0$$

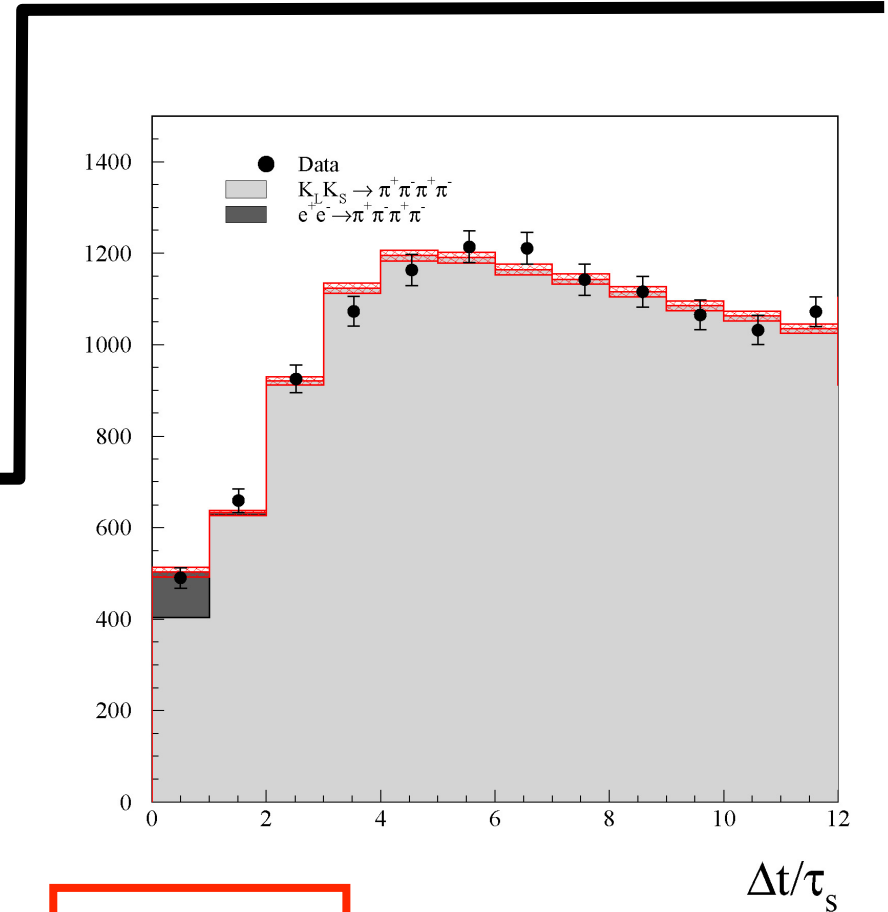
=> only one independent parameter: γ

The fit with $I(\pi^+ \pi^-, \pi^+ \pi^-; \Delta t, \gamma)$ gives:

KLOE result $L = 1.5 \text{ fb}^{-1}$

$$\gamma = (0.7 \pm 1.2_{STAT} \pm 0.3_{SYST}) \times 10^{-21} \text{ GeV}$$

PLB 642(2006) 315
Found. Phys. 40 (2010) 852



entangled kaons

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: CPT violation in entangled K states

In presence of decoherence and CPT violation induced by quantum gravity (CPT operator “ill-defined”) the definition of the particle-antiparticle states could be modified. This in turn could induce a breakdown of the correlations imposed by Bose statistics (EPR correlations) to the kaon state:

[Bernabeu, et al. PRL 92 (2004) 131601, NPB744 (2006) 180].

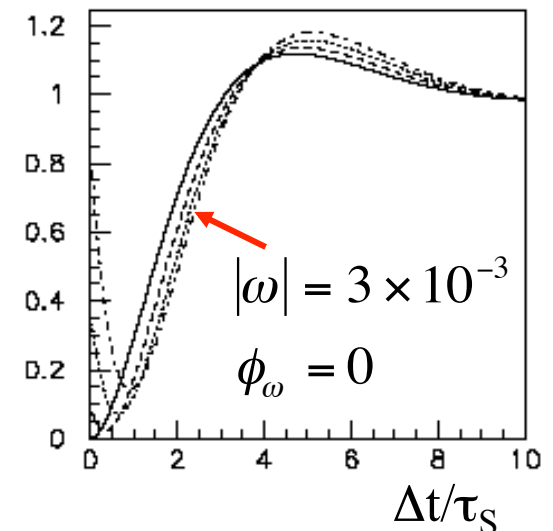
$$|i\rangle \propto (|K^0\rangle|\bar{K}^0\rangle - |\bar{K}^0\rangle|K^0\rangle) + \omega(|K^0\rangle|\bar{K}^0\rangle + |\bar{K}^0\rangle|K^0\rangle)$$

$$\propto (|K_S\rangle|K_L\rangle - |K_L\rangle|K_S\rangle) + \omega(|K_S\rangle|K_S\rangle - |K_L\rangle|K_L\rangle)$$

at most one expects:

$$|\omega|^2 = O\left(\frac{E^2/M_{PLANCK}}{\Delta\Gamma}\right) \approx 10^{-5} \Rightarrow |\omega| \sim 10^{-3}$$

$I(\pi^+\pi^-, \pi^+\pi^-; \Delta t)$ (a.u.)



In some microscopic models of space-time foam arising from non-critical string theory:

[Bernabeu, Mavromatos, Sarkar PRD 74 (2006) 045014]

$$|\omega| \sim 10^{-4} \div 10^{-5}$$

The maximum sensitivity to ω is expected for $f_1=f_2=\pi^+\pi^-$

All CPTV effects induced by QG ($\alpha, \beta, \gamma, \omega$) could be simultaneously disentangled.

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$: CPT violation in entangled K states

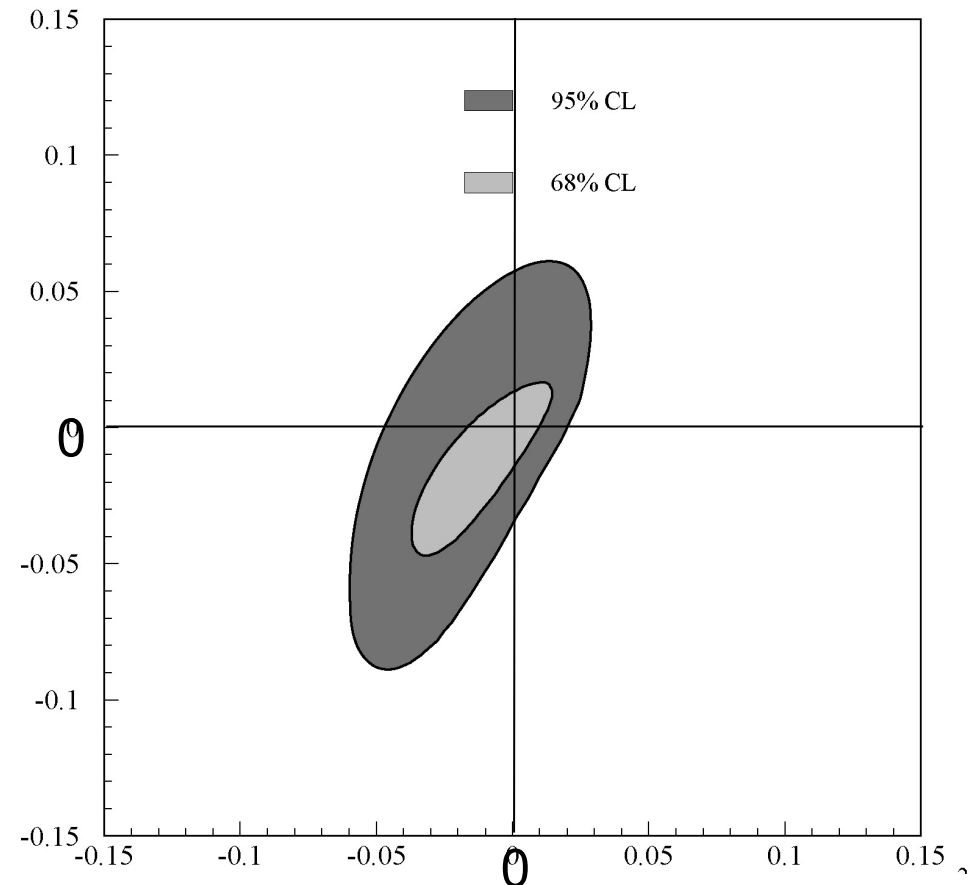
Fit of $I(\pi^+ \pi^-, \pi^+ \pi^-; \Delta t, \omega)$:

- Analysed data: 1.5 fb^{-1}

KLOE result: [PLB 642\(2006\) 315](#)
[Found. Phys. 40 \(2010\) 852](#)

$$\Re \omega = \left(-1.6_{-2.1}^{+3.0} \text{STAT} \pm 0.4_{\text{SYST}} \right) \times 10^{-4}$$
$$\Im \omega = \left(-1.7_{-3.0}^{+3.3} \text{STAT} \pm 1.2_{\text{SYST}} \right) \times 10^{-4}$$
$$|\omega| < 1.0 \times 10^{-3} \quad \text{at } 95\% \text{ C.L.}$$

$\text{Im } \omega \times 10^{-2}$



In the B system [Alvarez, Bernabeu, Nebot JHEP 0611, 087]:

$$-0.0084 \leq \Re \omega \leq 0.0100 \quad \text{at } 95\% \text{ C.L.}$$

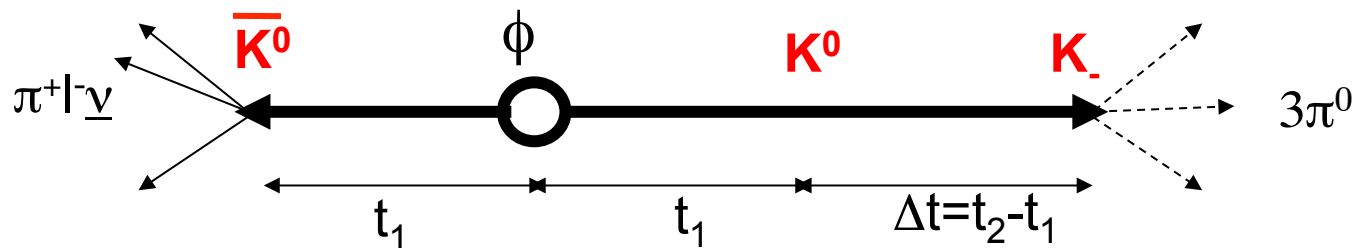
Direct test of time reversal symmetry

Direct test of Time Reversal symmetry with neutral kaons

- EPR correlations at a ϕ -factory (or B-factory) can be exploited to study other transitions involving also “CP states” K_+ and K_- (K_1, K_2)

$$\begin{aligned}
 |i\rangle &= \frac{1}{\sqrt{2}} \left[|K^0(\vec{p})\rangle |\bar{K}^0(-\vec{p})\rangle - |\bar{K}^0(\vec{p})\rangle |K^0(-\vec{p})\rangle \right] \\
 &= \frac{1}{\sqrt{2}} \left[|K_+(\vec{p})\rangle |K_-(-\vec{p})\rangle - |K_-(-\vec{p})\rangle |K_+(\vec{p})\rangle \right]
 \end{aligned}$$

- decay as filtering measurement
- entanglement \rightarrow preparation of state

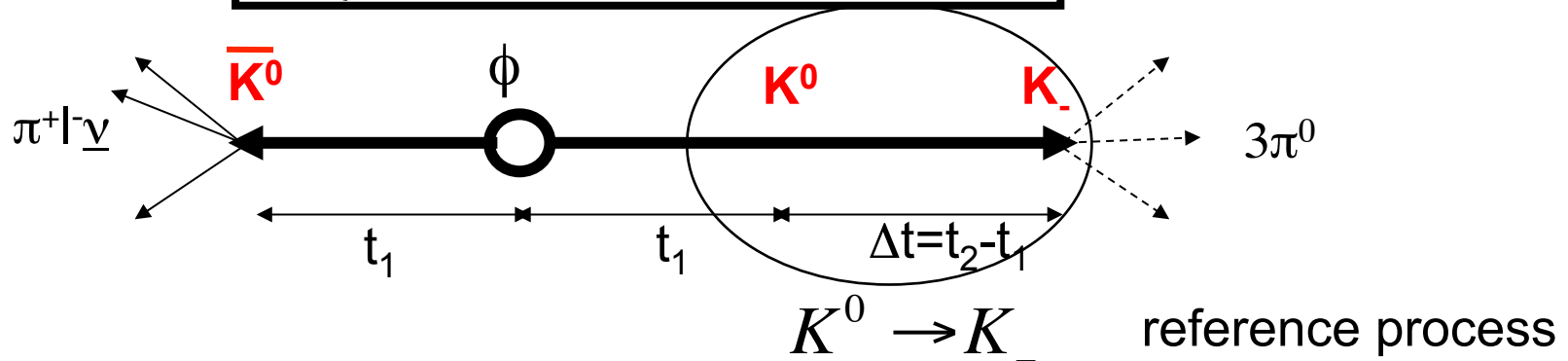


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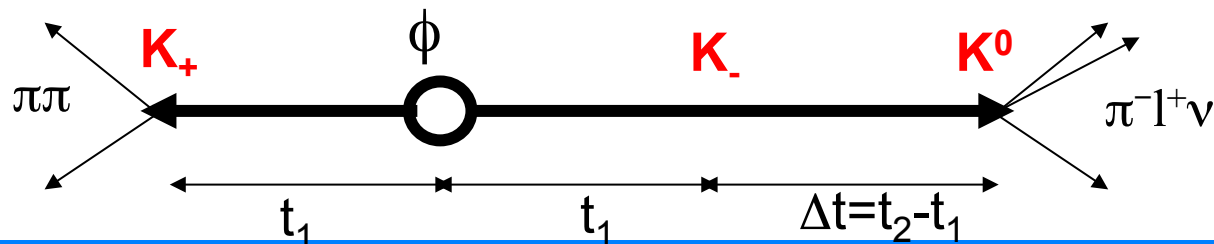
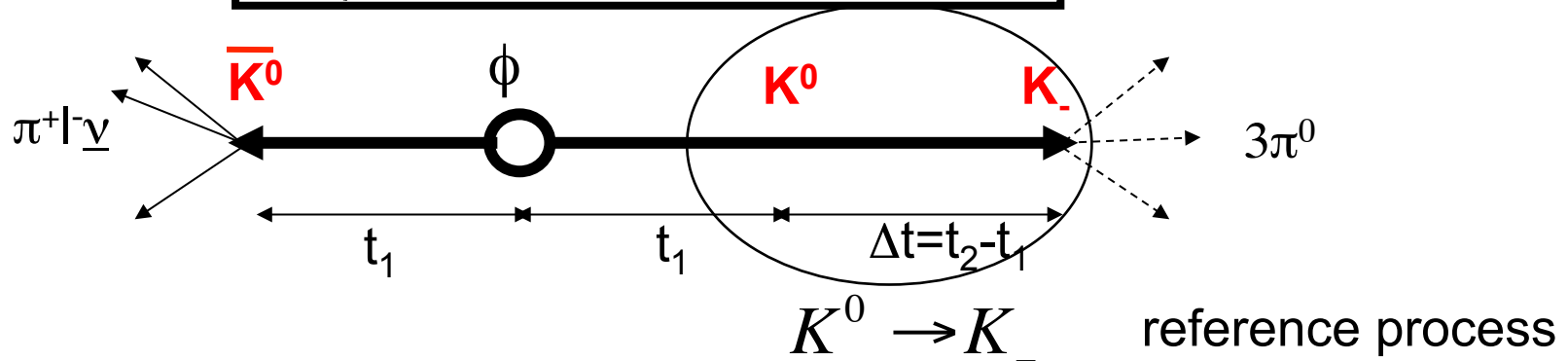


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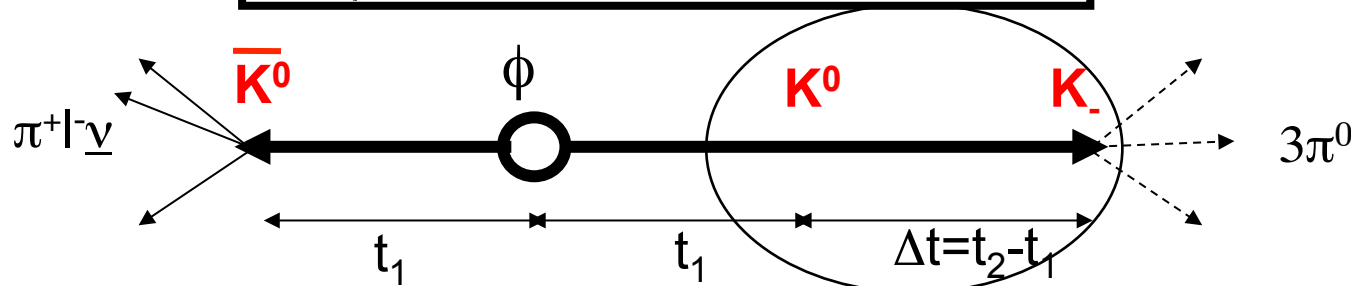


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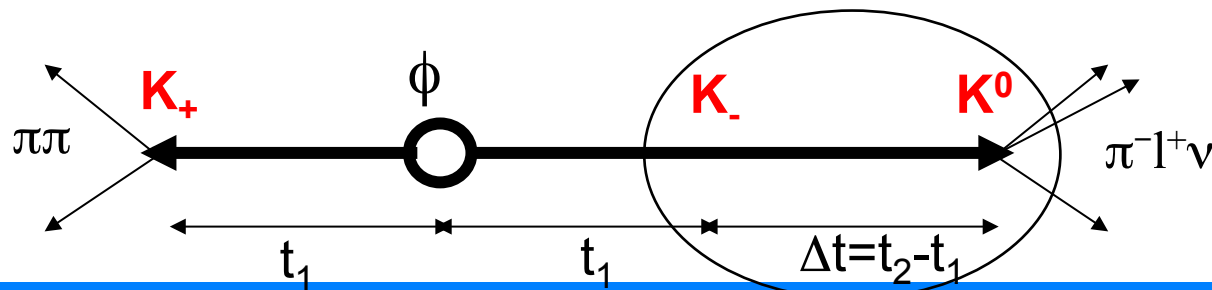
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$K^0 \rightarrow K_-$ reference process

$K_- \rightarrow K^0$ T-conjugated process

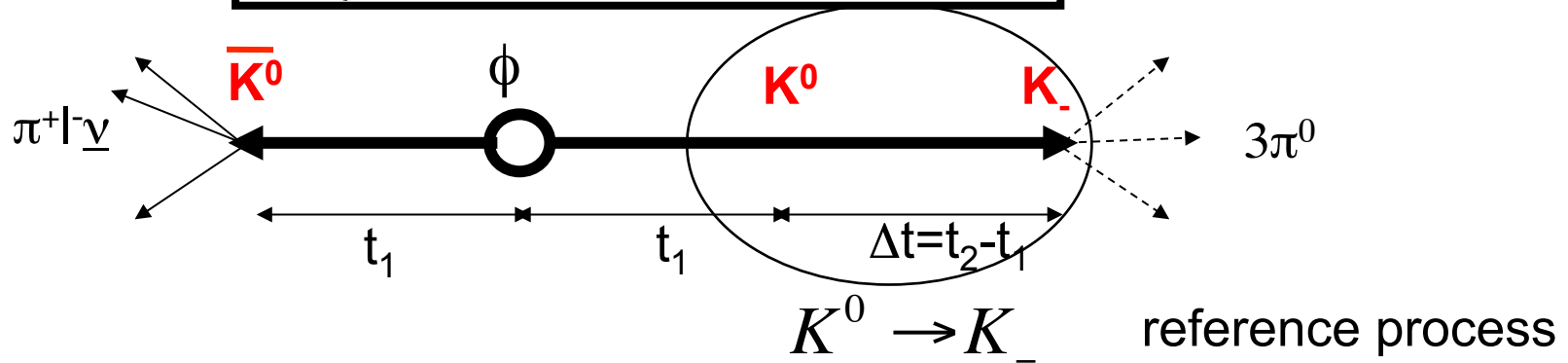


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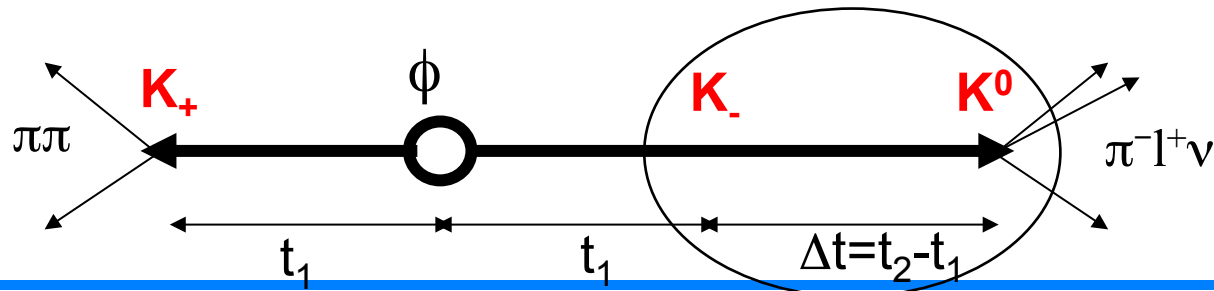
- decay as filtering measurement
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Note: CP-conjugated process

$$\bar{K}^0 \rightarrow K_-$$

$$K_- \rightarrow K^0 \quad \text{T-conjugated process}$$

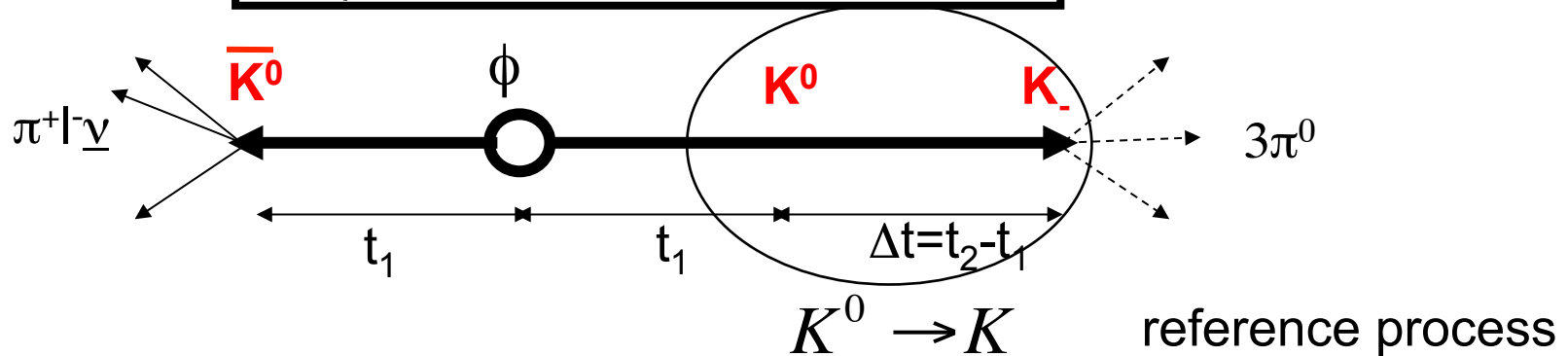


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 \end{aligned}$$

- decay as filtering measurement
- entanglement \rightarrow preparation of state



$$I(\pi\pi, l^+; \Delta t) = C(\pi\pi, l^+) \times P[K_-(0) \rightarrow K^0(\Delta t)]$$

In general with $f_{\bar{X}}$ decaying before f_Y , i.e. $\Delta t > 0$:

$$I(f_{\bar{X}}, f_Y; \Delta t) = C(f_{\bar{X}}, f_Y) \times P[K_X(0) \rightarrow K_Y(\Delta t)]$$

with $C(f_{\bar{X}}, f_Y) = \frac{1}{2(\Gamma_S + \Gamma_L)} |\langle f_{\bar{X}} | T | \bar{K}_X \rangle \langle f_Y | T | K_Y \rangle|^2$

Direct test of Time Reversal symmetry with neutral kaons

T symmetry test

Reference		T -conjugate	
Transition	Final state	Transition	Final state
$\bar{K}^0 \rightarrow K_-$	$(\ell^+, \pi^0 \pi^0 \pi^0)$	$K_- \rightarrow \bar{K}^0$	$(\pi^0 \pi^0 \pi^0, \ell^-)$
$K_+ \rightarrow K^0$	$(\pi^0 \pi^0 \pi^0, \ell^+)$	$K^0 \rightarrow K_+$	$(\ell^-, \pi \pi)$
$\bar{K}^0 \rightarrow K_+$	$(\ell^+, \pi \pi)$	$K_+ \rightarrow \bar{K}^0$	$(\pi^0 \pi^0 \pi^0, \ell^-)$
$K_- \rightarrow K^0$	$(\pi \pi, \ell^+)$	$K^0 \rightarrow K_-$	$(\ell^-, \pi \pi)$

One can define the following ratios of probabilities:

$$\begin{aligned}
 R_1(\Delta t) &= P [K^0(0) \rightarrow K_+(\Delta t)] / P [K_+(0) \rightarrow K^0(\Delta t)] \\
 R_2(\Delta t) &= P [K^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow K^0(\Delta t)] \\
 R_3(\Delta t) &= P [\bar{K}^0(0) \rightarrow K_+(\Delta t)] / P [K_+(0) \rightarrow \bar{K}^0(\Delta t)] \\
 R_4(\Delta t) &= P [\bar{K}^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow \bar{K}^0(\Delta t)] .
 \end{aligned}$$

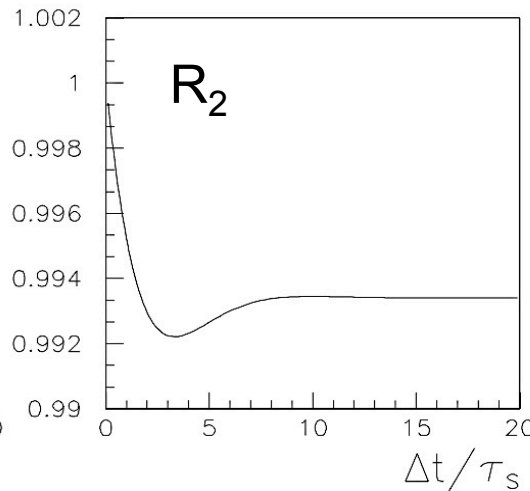
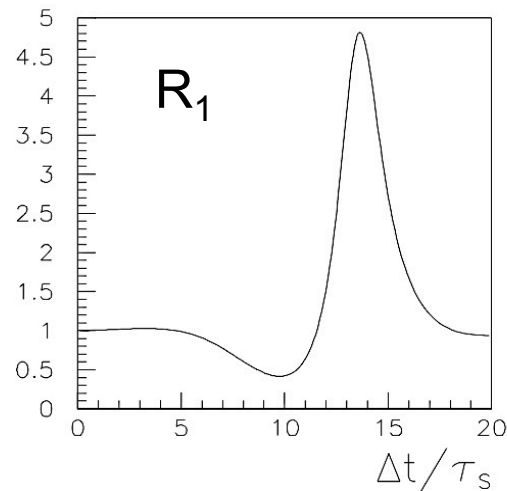
Any deviation from $R_i=1$ constitutes a violation of T-symmetry

J. Bernabeu, A.D.D., P. Villanueva: NPB 868 (2013) 102

Test feasible at KLOE-2 with $L=O(10 \text{ fb}^{-1})$

Direct test of Time Reversal symmetry with neutral kaons

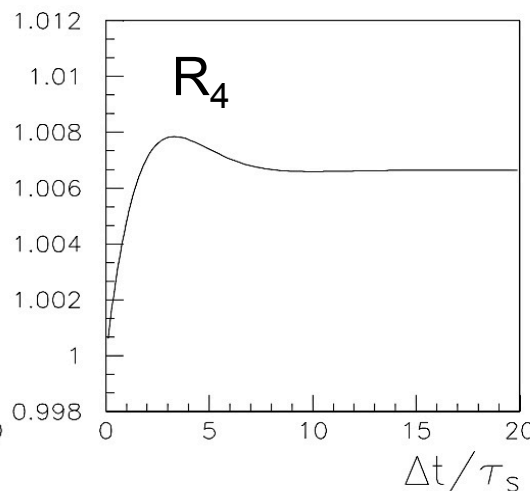
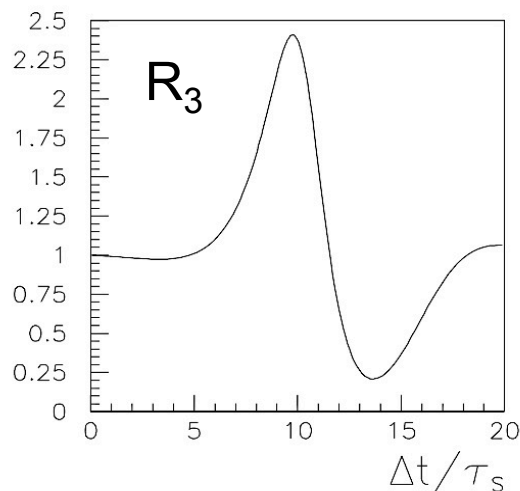
Any deviation from $R_i=1$ constitutes a direct evidence of T-symmetry violation



$$R_i(\Delta t=0)=1$$

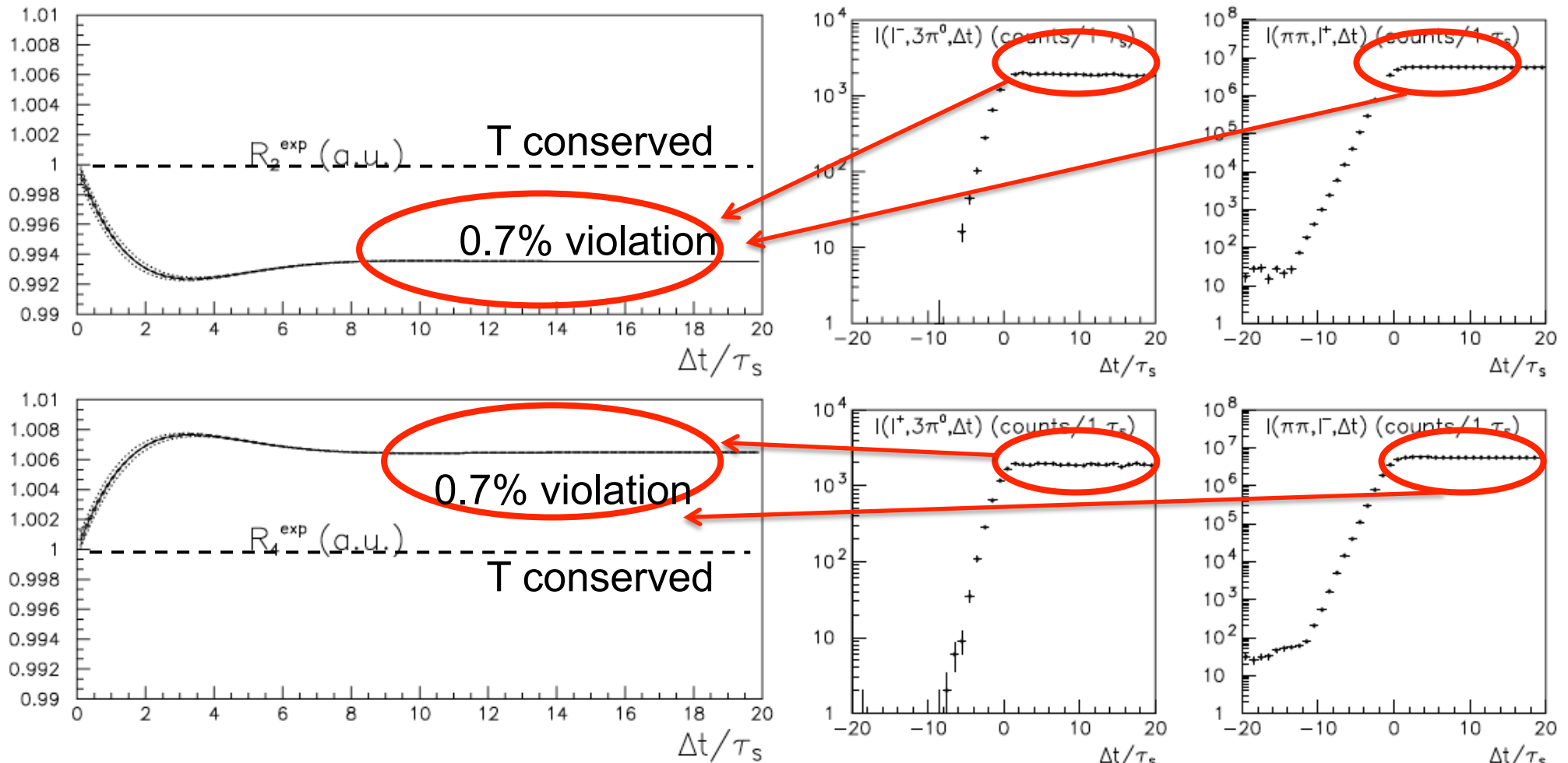
$$R_2(\Delta t \gg \tau_S)=1-4\text{Re}(\varepsilon)$$

$$R_4(\Delta t \gg \tau_S)=1+4\text{Re}(\varepsilon)$$



Direct test of Time Reversal symmetry with neutral kaons

toy MC with $L=10 \text{ fb}^{-1}$



$$R_2(\Delta t \gg \tau_S) = 1 - 4\text{Re}(\epsilon) \sim 0.993$$

$$R_4(\Delta t \gg \tau_S) = 1 + 4\text{Re}(\epsilon) \sim 1.007$$

Direct test of CPT symmetry with neutral kaons

CPT symmetry test

Reference		<i>CPT</i> -conjugate	
Transition	Decay products	Transition	Decay products
$K^0 \rightarrow K_+$	$(\ell^-, \pi\pi)$	$K_+ \rightarrow \bar{K}^0$	$(3\pi^0, \ell^-)$
$K^0 \rightarrow K_-$	$(\ell^-, 3\pi^0)$	$K_- \rightarrow \bar{K}^0$	$(\pi\pi, \ell^-)$
$\bar{K}^0 \rightarrow K_+$	$(\ell^+, \pi\pi)$	$K_+ \rightarrow K^0$	$(3\pi^0, \ell^+)$
$\bar{K}^0 \rightarrow K_-$	$(\ell^+, 3\pi^0)$	$K_- \rightarrow K^0$	$(\pi\pi, \ell^+)$

One can define the following ratios of probabilities:

$$R_{1,CPT}(\Delta t) = P [K^0(0) \rightarrow K_+(\Delta t)] / P [K_+(0) \rightarrow \bar{K}^0(\Delta t)]$$

$$R_{2,CPT}(\Delta t) = P [K^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow \bar{K}^0(\Delta t)]$$

$$R_{3,CPT}(\Delta t) = P [\bar{K}^0(0) \rightarrow K_+(\Delta t)] / P [K_+(0) \rightarrow K^0(\Delta t)]$$

$$R_{4,CPT}(\Delta t) = P [\bar{K}^0(0) \rightarrow K_-(\Delta t)] / P [K_-(0) \rightarrow K^0(\Delta t)]$$

Any deviation from $R_{i,CPT}=1$ constitutes a violation of CPT-symmetry

J. Bernabeu, A.D.D., P. Villanueva: NPB 868 (2013) 102

Significant test at KLOE-2 with $L=O(10 \text{ fb}^{-1})$, feasibility study ongoing

Future perspectives

KLOE-2 at upgraded DAΦNE

DAΦNE upgraded in luminosity:

- new scheme of the interaction region (crabbed waist scheme) at DAΦNE (proposal by P. Raimondi)
- increase L by a factor $\times 3$ demonstrated by an experimental test

KLOE-2 experiment:

- extend the KLOE physics program at DAΦNE upgraded in luminosity
- Collect $O(10) \text{ fb}^{-1}$ of integrated luminosity in the next 2-3 years

Physics program

(see [EPJC 68 \(2010\) 619-681](#))

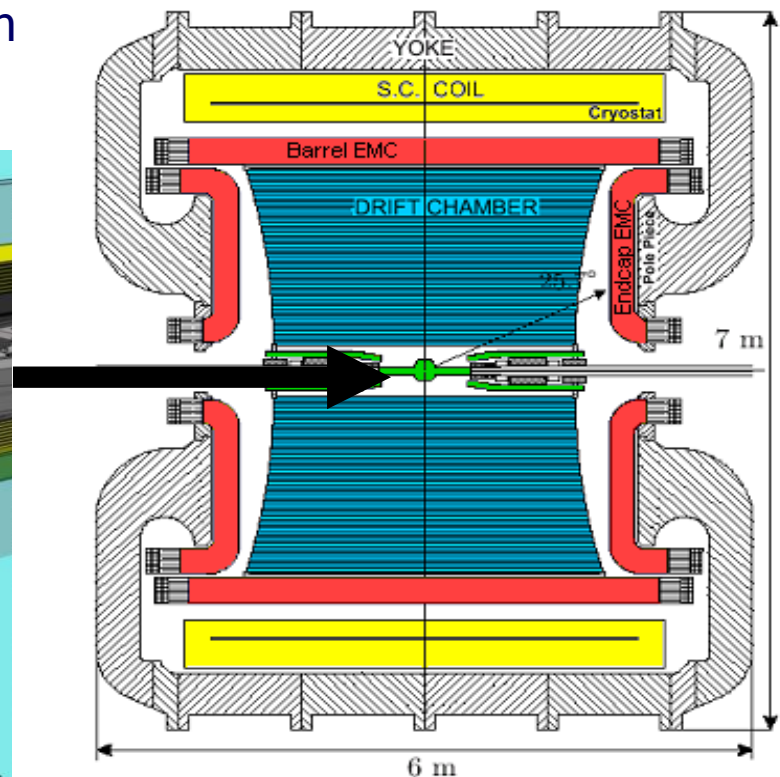
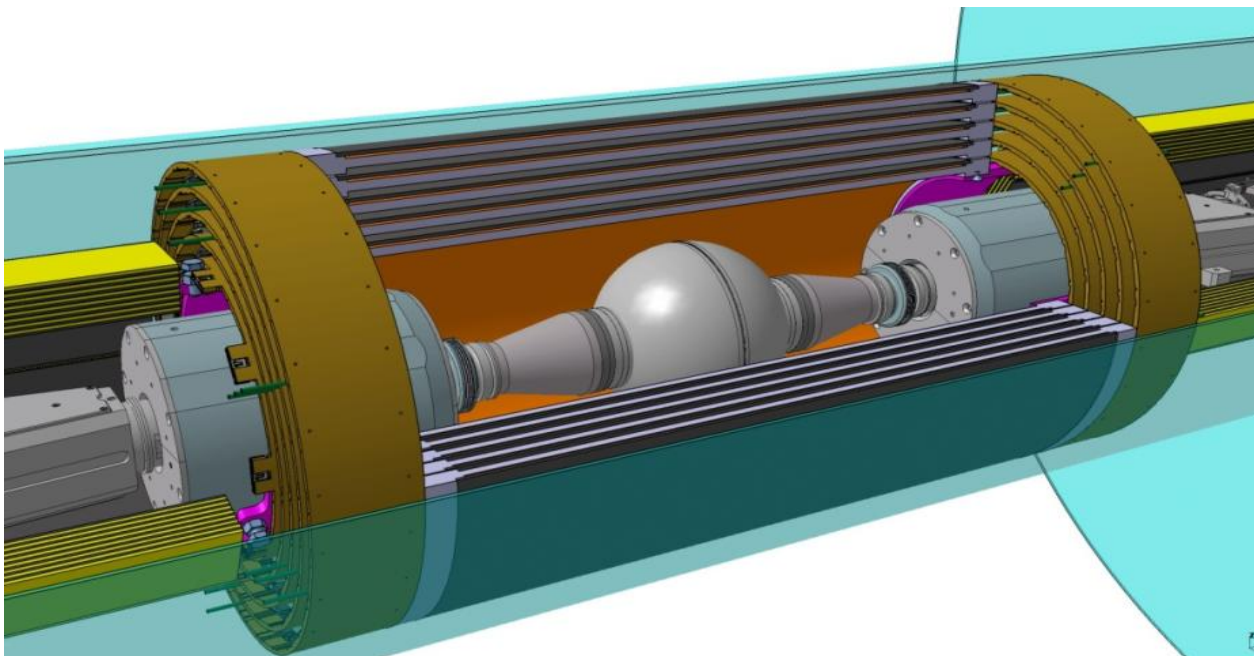
- Neutral kaon interferometry, CPT symmetry & QM tests
- Kaon physics, CKM, LFV, rare K_S decays
- η, η' physics
- Light scalars, $\gamma\gamma$ physics
- Hadron cross section at low energy, a_μ
- Dark forces: search for light U boson

Detector upgrade:

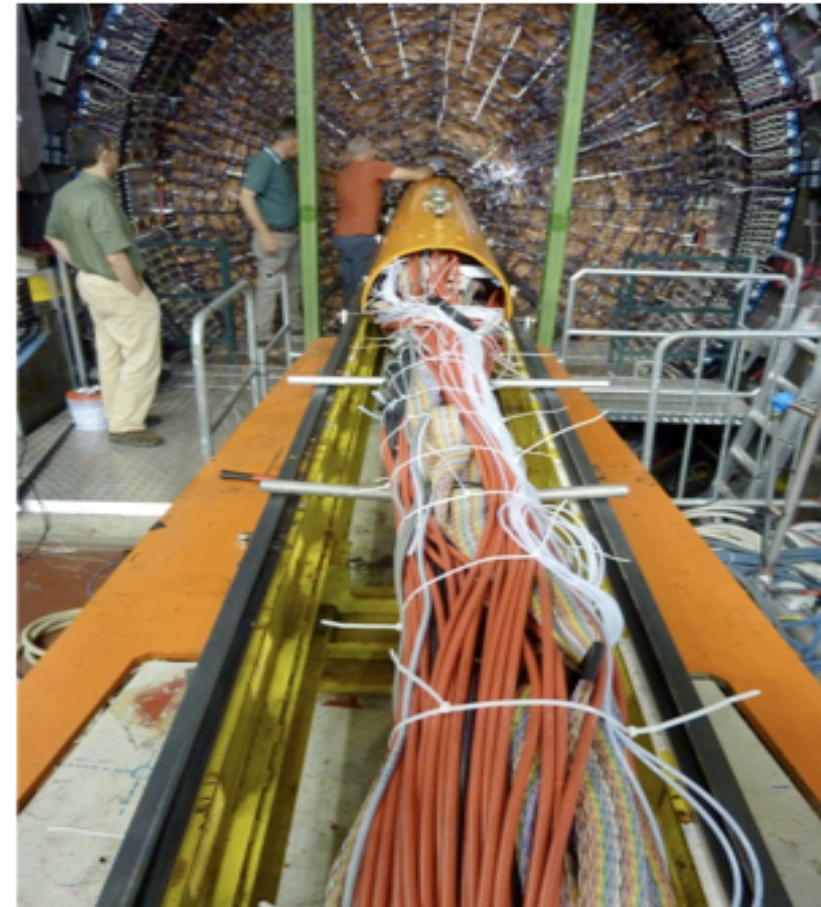
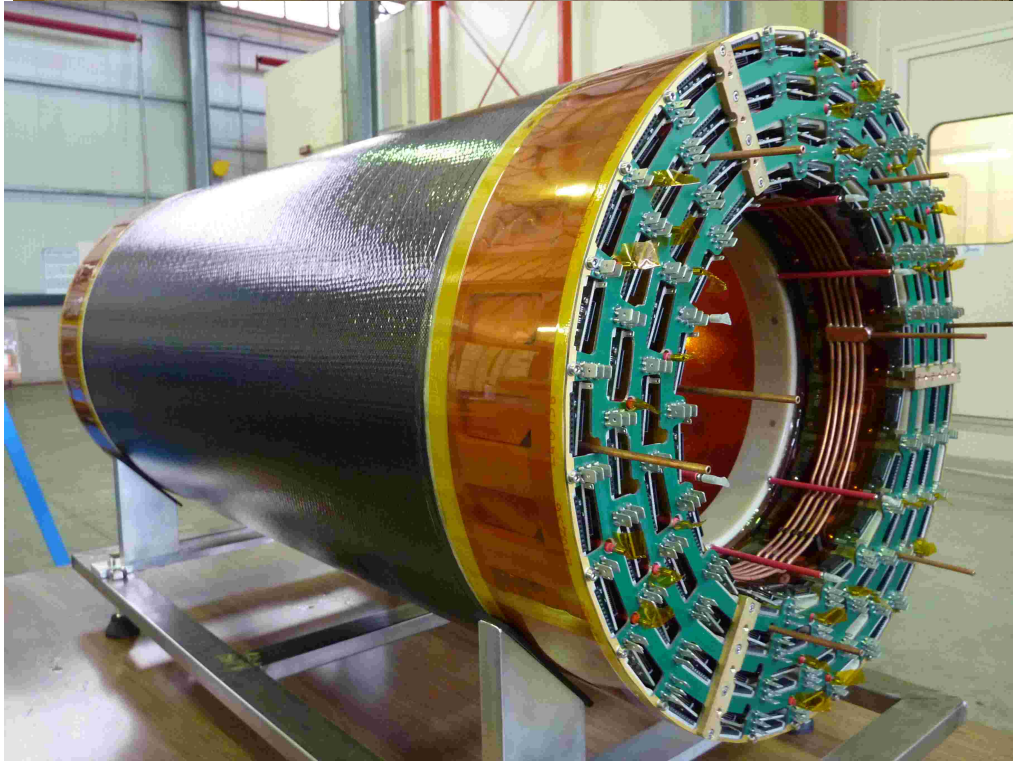
- $\gamma\gamma$ tagging system
- inner tracker
- small angle and quad calorimeters
- FEE maintenance and upgrade
- Computing and networking update
- etc.. (Trigger, software, ...)

Inner tracker at KLOE-2

- 4 independent tracking layers for a fine vertex reconstruction of K_S and η
- $200 \mu\text{m}$ $\sigma_{r\phi}$ and $500 \mu\text{m}$ σ_z spatial resolutions with XV readout
- 700 mm active length
- from 150 to 250 mm radii
- 1.8% X_0 total radiation length in the active region
- Realized with **Cylindrical-GEM** detectors



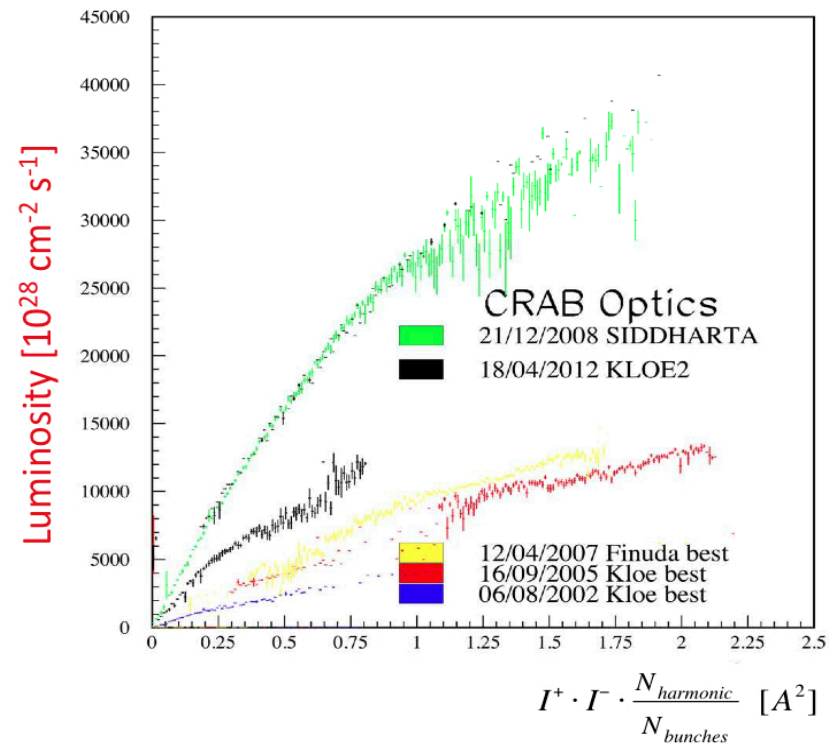
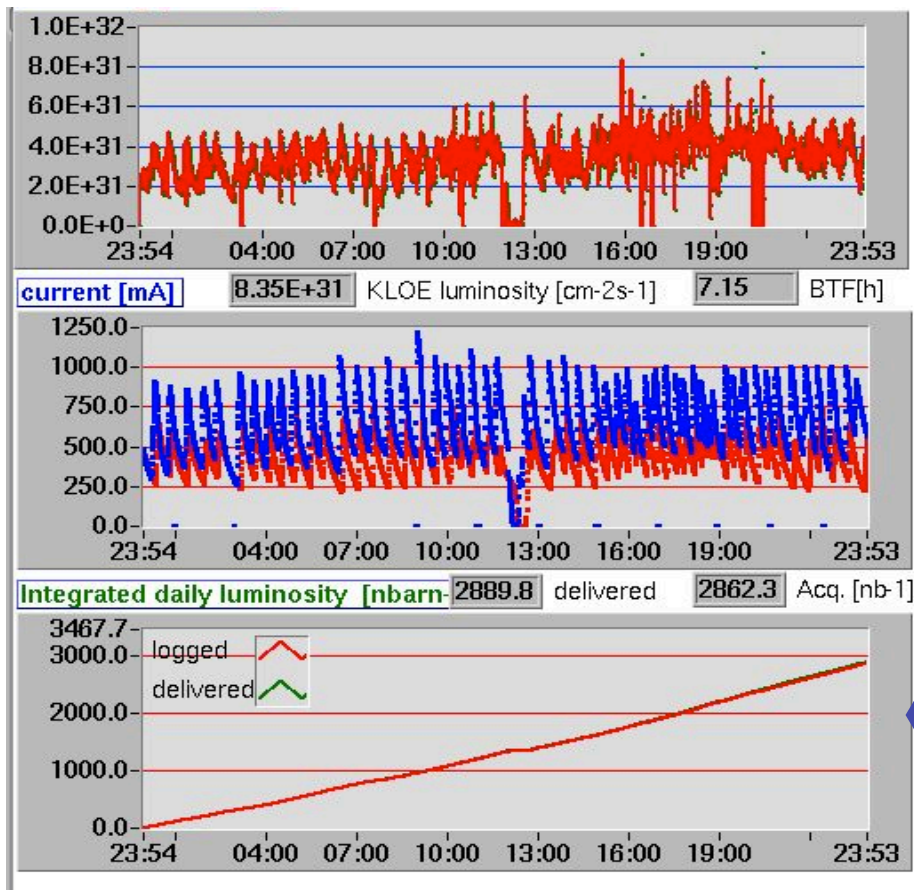
Inner tracker at KLOE



- Construction and installation inside KLOE completed (July 2013)
- Commissioning in progress

DAΦNE luminosity upgrade

- DAΦNE upgrade (2008): new interaction scheme; large beam crossing angle + crabbed waist sextupoles



- Installation of the new DAΦNE IR + KLOE-2 upgrades completed in July 2013

- Feb. 2014: one week after restarting collisions
- Crab sextupoles off
- Uptime ~ 100%

Prospects for KLOE-2

Param.	Present best published measurement	KLOE-2 (IT) L=5 fb ⁻¹	KLOE-2 (IT) L=10 fb ⁻¹
ξ_{00}	$(0.1 \pm 1.0) \times 10^{-6}$	$\pm 0.26 \times 10^{-6}$	$\pm 0.18 \times 10^{-6}$
ξ_{SL}	$(0.3 \pm 1.9) \times 10^{-2}$	$\pm 0.49 \times 10^{-2}$	$\pm 0.35 \times 10^{-2}$
α	$(-0.5 \pm 2.8) \times 10^{-17} \text{ GeV}$	$\pm 5.0 \times 10^{-17} \text{ GeV}$	$\pm 3.5 \times 10^{-17} \text{ GeV}$
β	$(2.5 \pm 2.3) \times 10^{-19} \text{ GeV}$	$\pm 0.50 \times 10^{-19} \text{ GeV}$	$\pm 0.35 \times 10^{-19} \text{ GeV}$
γ	$(1.1 \pm 2.5) \times 10^{-21} \text{ GeV}$ compl. pos. hyp. $(0.7 \pm 1.2) \times 10^{-21} \text{ GeV}$	$\pm 0.75 \times 10^{-21} \text{ GeV}$ compl. pos. hyp. $\pm 0.33 \times 10^{-21} \text{ GeV}$	$\pm 0.53 \times 10^{-21} \text{ GeV}$ compl. pos. hyp. $\pm 0.23 \times 10^{-21} \text{ GeV}$
Re(ω)	$(-1.6 \pm 2.6) \times 10^{-4}$	$\pm 0.70 \times 10^{-4}$	$\pm 0.49 \times 10^{-4}$
Im(ω)	$(-1.7 \pm 3.4) \times 10^{-4}$	$\pm 0.86 \times 10^{-4}$	$\pm 0.61 \times 10^{-4}$
Δa_0	$(-6.2 \pm 8.8) \times 10^{-18} \text{ GeV}$	$\pm 4.8 \times 10^{-18} \text{ GeV}$	$\pm 3.4 \times 10^{-18} \text{ GeV}$
Δa_Z	$(-0.7 \pm 1.0) \times 10^{-18} \text{ GeV}$	$\pm 0.6 \times 10^{-18} \text{ GeV}$	$\pm 0.4 \times 10^{-18} \text{ GeV}$
Δa_X	$(3.3 \pm 2.2) \times 10^{-18} \text{ GeV}$	$\pm 0.76 \times 10^{-18} \text{ GeV}$	$\pm 0.54 \times 10^{-18} \text{ GeV}$
Δa_Y	$(-0.7 \pm 2.0) \times 10^{-18} \text{ GeV}$	$\pm 0.76 \times 10^{-18} \text{ GeV}$	$\pm 0.54 \times 10^{-18} \text{ GeV}$

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ elsewhere??

Is it possible to perform these tests of Quantum Mechanics and CPT symmetry with entangled neutral kaons elsewhere than a ϕ -factory?

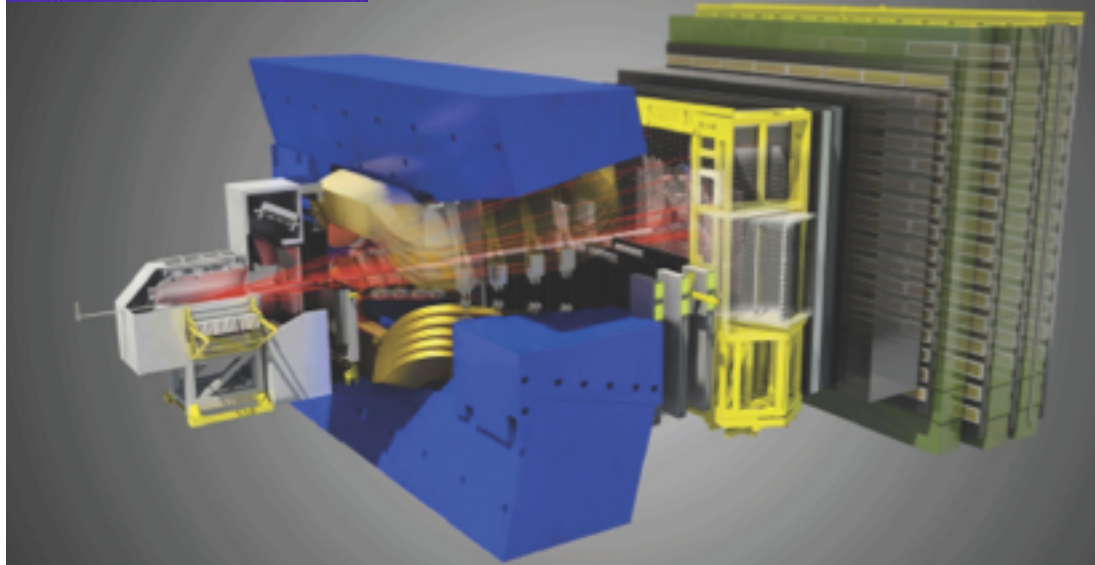
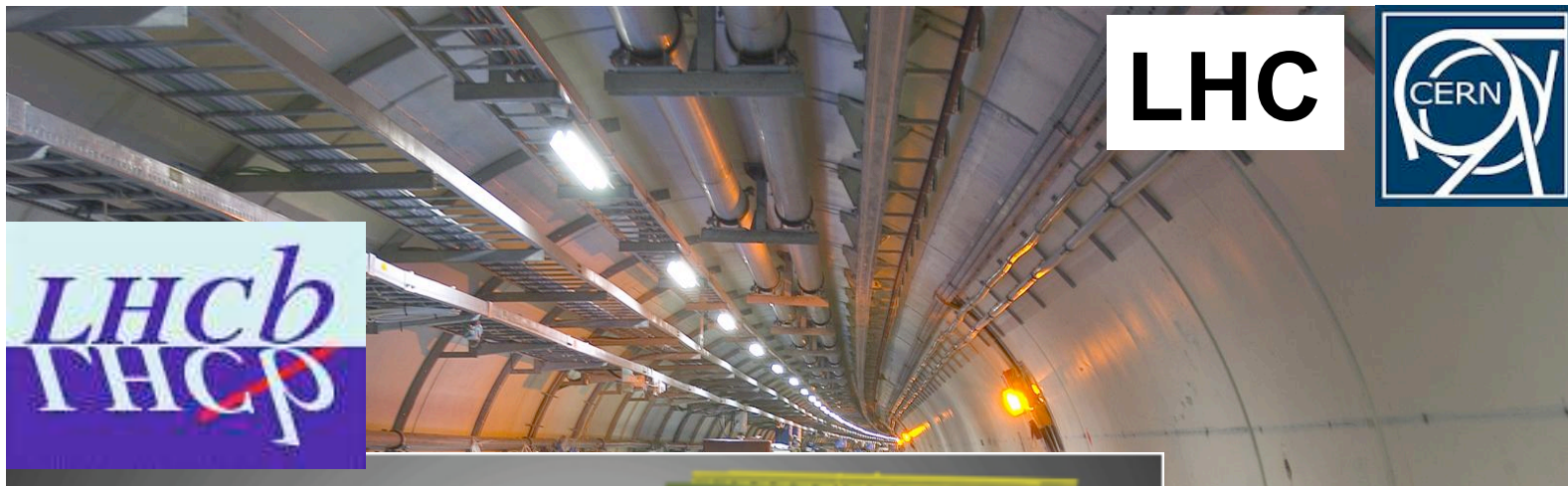
As we have seen the process $\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ is very rich in Physics and plays a key role in this respect.

From the experimental point of view the main advantage of this process is the “closed” kinematics which makes it “easily” observable – in principle – also in a “hadron machine” environment.

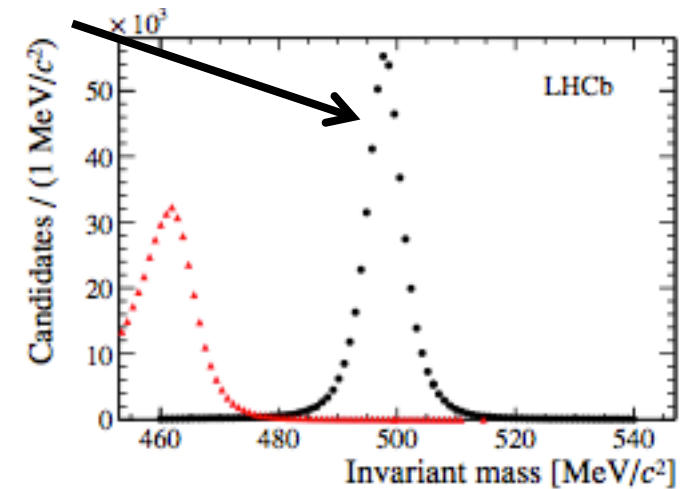
$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ at LHC??



$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ at LHC??



(from JHEP 01(2013) 090
search for $K_S \rightarrow \mu^+ \mu^-$)
 $K_S \rightarrow \pi^+ \pi^-$ $\sigma \sim 4 \text{ MeV}/c^2$



=> feasibility study, A.D.D. + T. Ruf (CERN), ongoing

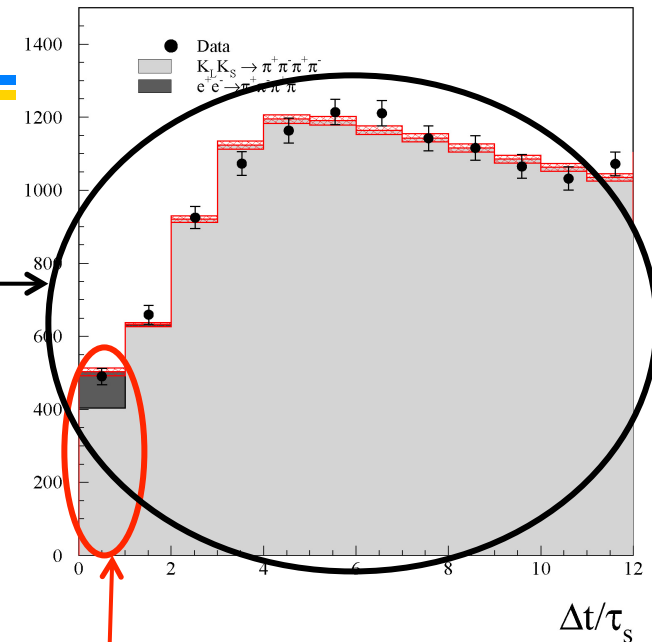
$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ at LHCb: statistics

at KLOE

$\lambda_S \sim 6$ mm ($P=110$ MeV/c)

$L=1.5$ fb⁻¹

4.4×10^9 ϕ 's



At LHCb:

$\sim 0.2 < \lambda_S \sim 0.8$ m ($\sim 3 < P < \sim 20$ GeV/c)

$\sigma(pp \rightarrow \phi X) \sim 1800$ $\mu\text{b} \rightarrow 1.8 \times 10^{12}$ ϕ 's /fb⁻¹

(in LHCb acceptance [PLB 703\(2011\) 267](#))

NOT including the efficiency one expects:

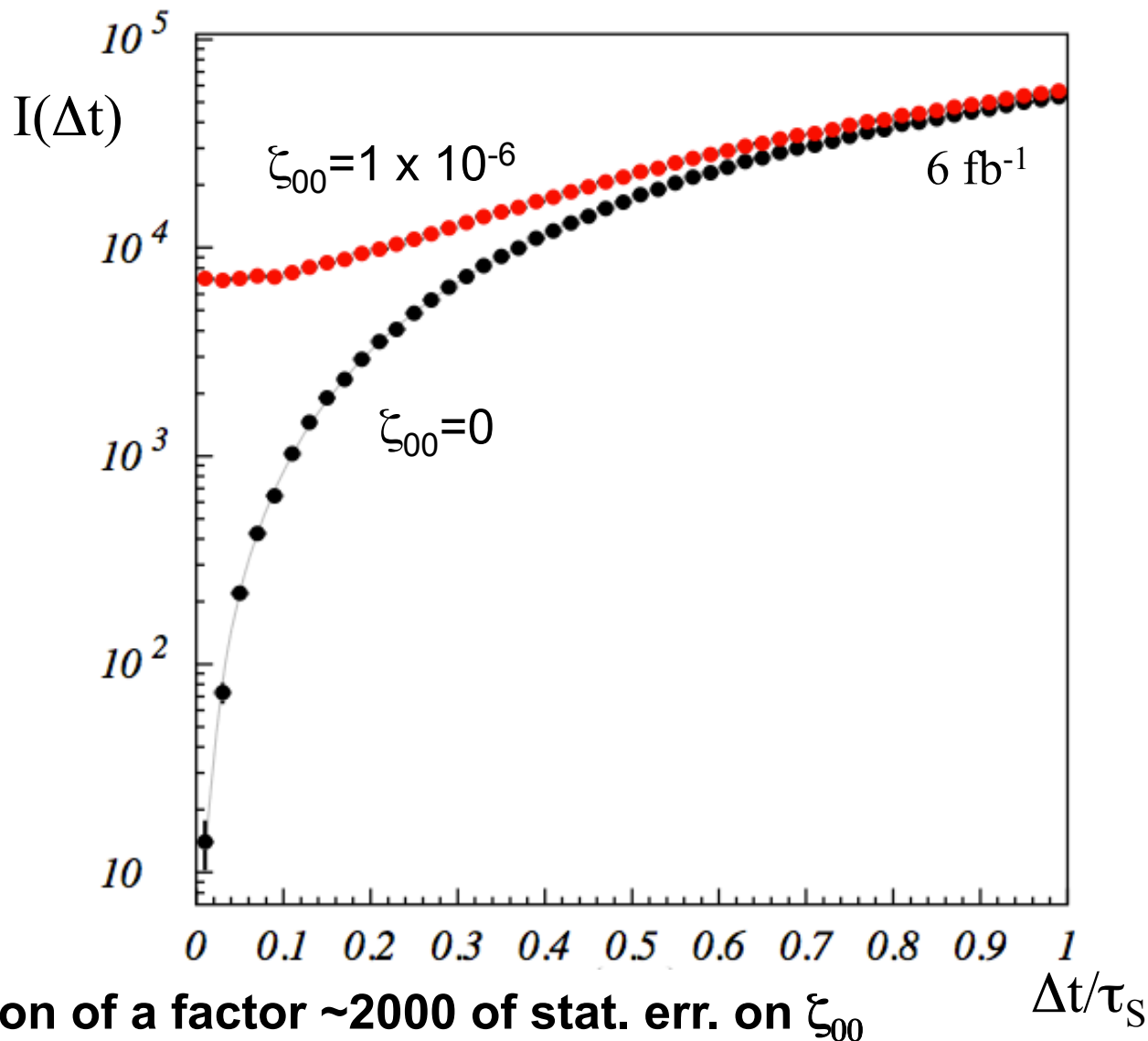
$N(\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-) = 1.5 \times 10^5$ evts/fb⁻¹ in 0-1 τ_S range

x 6 fb⁻¹ $\sim 9.0 \times 10^5$ evts

$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ at LHCb: toy MC simulation

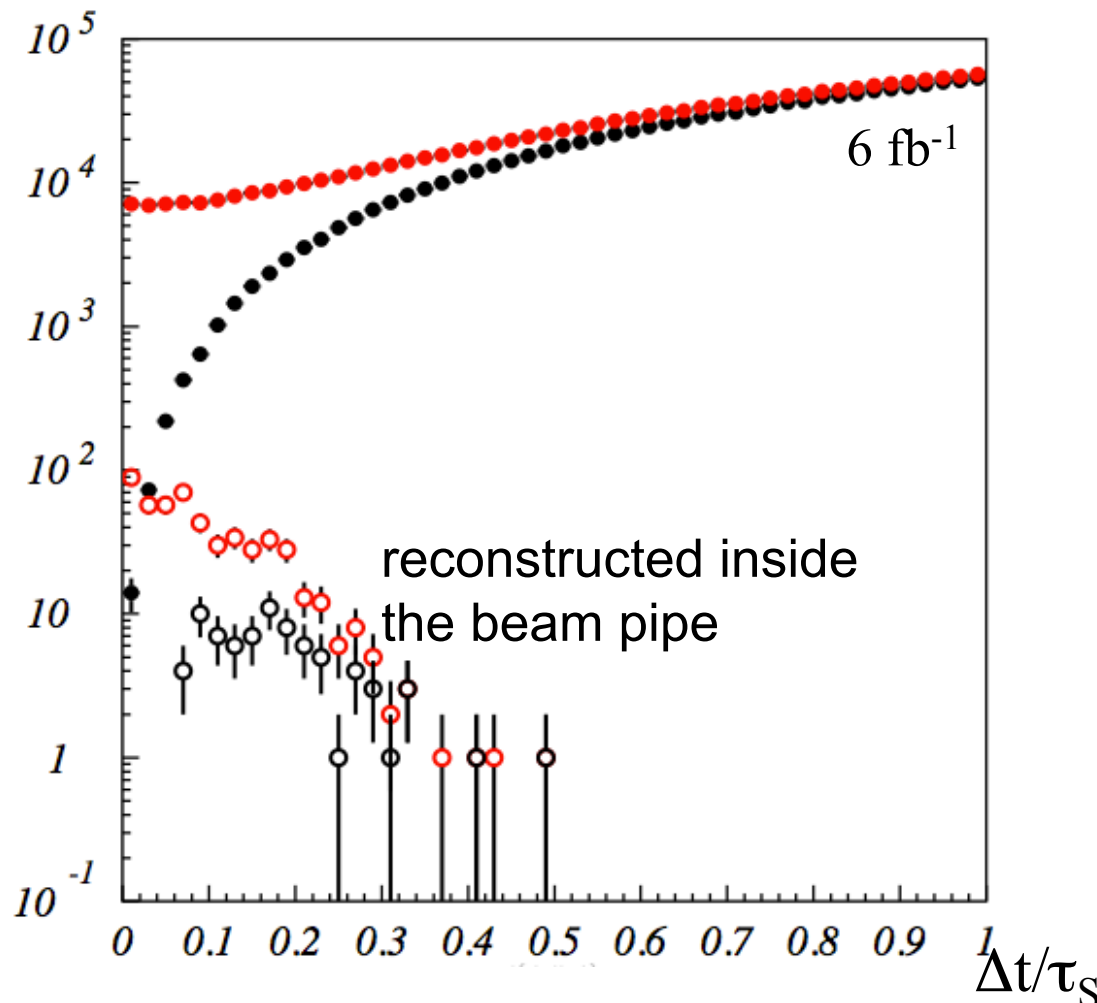
FIT in 0 -1 τ_S range
 N events = 10^6 ($\sim 6\text{fb}^{-1}$)

	$\sigma(\xi_{00})$
KLOE 1.5 fb^{-1}	$\pm 1.0 \times 10^{-6}$
LHCb 6 fb^{-1}	$\pm 0.4 \times 10^{-9}$



$\phi \rightarrow K_S K_L \rightarrow \pi^+ \pi^- \pi^+ \pi^-$ at LHCb: toy MC simulation

-In order to avoid K_S regeneration on the beam pipe, both kaon vertices are required to be reconstructed inside the beam pipe: $r < 5$ mm



- Δt resolution effects are totally negligible
- only $0 < \Delta t < 0.3 \tau_S$ range can be reconstructed (both kaons decaying at short times), which is the most relevant for CPT and QM tests!!
- large background to be rejected/subtracted
- trigger efficiency to be increased

Prospects for KLOE-2 and LHCb

PRELIMINARY

Para m.	Present best published measurement	KLOE-2 (IT) L=5 fb ⁻¹	KLOE-2 (IT) L=10 fb ⁻¹	LHCb L=6 fb ⁻¹
ξ_{00}	$(0.1 \pm 1.0) \times 10^{-6}$	$\pm 0.26 \times 10^{-6}$	$\pm 0.18 \times 10^{-6}$	$\pm 0.04 \times 10^{-6}$
ξ_{SL}	$(0.3 \pm 1.9) \times 10^{-2}$	$\pm 0.49 \times 10^{-2}$	$\pm 0.35 \times 10^{-2}$	$\pm 0.11 \times 10^{-2}$
α	$(-0.5 \pm 2.8) \times 10^{-17}$ GeV	$\pm 5.0 \times 10^{-17}$ GeV	$\pm 3.5 \times 10^{-17}$ GeV	
β	$(2.5 \pm 2.3) \times 10^{-19}$ GeV	$\pm 0.50 \times 10^{-19}$ GeV	$\pm 0.35 \times 10^{-19}$ GeV	
γ	$(1.1 \pm 2.5) \times 10^{-21}$ GeV compl. pos. hyp. $(0.7 \pm 1.2) \times 10^{-21}$ GeV	$\pm 0.75 \times 10^{-21}$ GeV compl. pos. hyp. $\pm 0.33 \times 10^{-21}$ GeV	$\pm 0.53 \times 10^{-21}$ GeV compl. pos. hyp. $\pm 0.23 \times 10^{-21}$ GeV	compl. pos. hyp. $\pm 0.07 \times 10^{-21}$ GeV
Re(ω)	$(-1.6 \pm 2.6) \times 10^{-4}$	$\pm 0.70 \times 10^{-4}$	$\pm 0.49 \times 10^{-4}$	$\pm 0.70 \times 10^{-4}$
Im(ω)	$(-1.7 \pm 3.4) \times 10^{-4}$	$\pm 0.86 \times 10^{-4}$	$\pm 0.61 \times 10^{-4}$	$\pm 1.1 \times 10^{-4}$
Δa_0	$(-6.2 \pm 8.8) \times 10^{-18}$ GeV	$\pm 4.8 \times 10^{-18}$ GeV	$\pm 3.4 \times 10^{-18}$ GeV	
Δa_Z	$(-0.7 \pm 1.0) \times 10^{-18}$ GeV	$\pm 0.6 \times 10^{-18}$ GeV	$\pm 0.4 \times 10^{-18}$ GeV	
Δa_X	$(3.3 \pm 2.2) \times 10^{-18}$ GeV	$\pm 0.76 \times 10^{-18}$ GeV	$\pm 0.54 \times 10^{-18}$ GeV	
Δa_Y	$(-0.7 \pm 2.0) \times 10^{-18}$ GeV	$\pm 0.76 \times 10^{-18}$ GeV	$\pm 0.54 \times 10^{-18}$ GeV	

Conclusions

- The entangled neutral kaon system at a ϕ -factory is an excellent laboratory for studies of fundamental physical principles, e.g. the study of discrete symmetries and of the basic principles of Quantum Mechanics;
- **Several parameters related to possible decoherence and/or CPT violation (CPT violation and Lorentz symmetry breaking) have been recently measured at KLOE, in some cases with a precision reaching the interesting Planck's scale region;**
- All results are consistent with no symmetry violation and no decoherence

- Neutral kaon interferometry, CPT symmetry and QM tests are one of the main issues of the **KLOE-2** physics program. (EPJC 68 (2010) 619-681)
- The precision of several tests could be improved by about one order of magnitude.

- **LHCb** might further improve these limits.
The main advantages are: (i) copious ϕ production rate
(ii) excellent reconstruction capabilities in the interesting $\Delta t \sim 0$ region
(iii) accurate Δt reconstruction \rightarrow negligible Δt resolution
- The main challenging issues are: (iv) background rejection and subtraction
(v) trigger efficiency; \rightarrow feasibility study ongoing.