Compactifications on Generalized Geometries

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Introduction

Phenomenological models in string theory:

space-time background

$$M_{10} = M_4 \times Y_6$$

• N=1 (spontaneously broken) supersymmetry

realized as

- ⇒ Heterotic string with Y₆: Calabi-Yau threefold/orbifold
- □ Type II/I with space-time filling D-branes and orientifold-planes

Y₆: generalized Calabi-Yau orientifold (with background fluxes)

Problems:

- mechanism for spontaneous supersymmetry breaking
- stabilization of moduli
- "realistic" models of particle physics
- model for inflation
- cosmological constant

they might all be (un-) related

Purpose of this talk:

discuss string compactifications on generalized geometries

Compactification:

Space-time background: $\mathbf{M}_{10} = \mathbf{M_4} \times \mathbf{Y_6}$

Lorentz-group: $SO(1,9) \rightarrow SO(1,3) \times SO(6)$

10D Supercharge: $\mathbf{16} \rightarrow (\mathbf{2}, \mathbf{4}) \oplus (\mathbf{\bar{2}}, \mathbf{\bar{4}})$

Impose:

1. existence of 4D supercharge \Rightarrow existence of globally def. spinor η

 \Rightarrow Y₆ has reduced structure group SU(3)

$$SO(6) \rightarrow SU(3)$$
 s.t. $4 \rightarrow 3 + 1$

2. background preserves supersymmetry

$$\delta \Psi = \nabla \eta + (\gamma \cdot \mathbf{F}) \eta + \dots = 0$$
, $\mathbf{F} = \text{background fluxes}$

- $\nabla \eta = 0 \Rightarrow \mathbf{Y_6}$ is Calabi-Yau manifold
- ullet here: study manifolds with SU(3)-structure, i.e. $\nabla \eta \neq \mathbf{0}$

Manifolds with SU(3) structure: [Gray, Hervella, Salamon, Chiossi, Hitchin,...]

characterized by two tensors J, Ω (follows from existence of η)

$$(1,1)$$
-form

$$\mathbf{J_{mn}} = \eta^\dagger \gamma_{[\mathbf{m}} \gamma_{\mathbf{n}]} \eta \; , \qquad \mathbf{dJ}
eq \mathbf{0}$$

 $\Rightarrow (3,0)$ -form

$$\mathbf{\Omega_{mnp}} = \eta^{\dagger} \gamma_{[\mathbf{m}} \gamma_{\mathbf{n}} \gamma_{\mathbf{p}]} \eta , \qquad \mathbf{d} \Omega \neq \mathbf{0}$$

Remarks:

- $dJ, d\Omega \sim \text{(intrinsic) torsion of } Y_6$
- Calabi-Yau: $\nabla \eta = 0 \Rightarrow dJ = d\Omega = 0$
 - ⇒ torsion parameterizes the deviation from Calabi-Yau

Manifolds with $SU(3) \times SU(3)$ structure:

[Jescheck, Witt; Grana, Minasian, Petrini, Tomasiello; Grana, Waldram, JL]

In type II string theory one can be slightly more general:

choose different spinors η^1, η^2 for the two gravitini $\Psi^{1,2}$

each η def. SU(3)-structure \Rightarrow together: $\mathbf{SU(3)} \times \mathbf{SU(3)}$ -structure (characterized by pair $\mathbf{J^{1,2}}, \mathbf{\Omega^{1,2}}$)

<u>Hitchin:</u> embed in $SU(3) \times SU(3)$ in O(6,6) acting on $T \oplus T^*$

 \Rightarrow structure characterized by two pure spinors Φ^+, Φ^- of O(6,6)

$$\mathbf{\Phi}^+ = \mathbf{e}^{\mathbf{B}} \eta_+^{\mathbf{1}} \otimes \bar{\eta}_+^{\mathbf{2}} \simeq \sum \mathbf{\Phi}_{\mathrm{even}}^+ , \qquad \mathbf{\Phi}^- = \mathbf{e}^{\mathbf{B}} \eta_+^{\mathbf{1}} \otimes \bar{\eta}_-^{\mathbf{2}} \simeq \sum \mathbf{\Phi}_{\mathrm{odd}}^+ ,$$

$$SU(3)$$
 structure $(\eta^1 = \eta^2)$: $\Phi^+ = e^{B+iJ}$, $\Phi^- = e^B\Omega$,

Low energy effective action:

$$S = \int_{\mathbf{M_4}} \frac{1}{2} \mathbf{R} - \mathbf{g_{ab}}(\mathbf{z}) \, \mathbf{D}_{\mu} \mathbf{z^a} \mathbf{D}^{\mu} \mathbf{z^b} - \mathbf{V}(\mathbf{z}) + \dots$$

 \Rightarrow Type II string theory: S is N=2 gauged supergravity (before orientifolding)

- $\mathbf{z}^{\mathbf{a}}$: coordinates of scalar manifold \mathcal{M}
 - correspond to deformations of $\mathbf{B}, \mathbf{J}, \mathbf{\Omega}$ or $\mathbf{\Phi}^+, \mathbf{\Phi}^-$
 - scalars from RR-sector

$$N = 2$$
 constraint: $\mathcal{M} = \mathcal{M}_{SK} \times \mathcal{M}_{QK}$

type IIA:
$$\mathcal{M}_{SK} = \mathcal{M}_{\Phi^+}$$
, $\mathcal{M}_{QK} \supset \mathcal{M}_{\Phi^-}$
type IIB: $\mathcal{M}_{SK} = \mathcal{M}_{\Phi^-}$, $\mathcal{M}_{QK} \supset \mathcal{M}_{\Phi^+}$

type IIB:
$$\mathcal{M}_{SK} = \mathcal{M}_{\Phi^-}$$
, $\mathcal{M}_{QK} \supset \mathcal{M}_{\Phi^+}$

Metric g_{ab} : special Kähler metric on $\mathcal{M} = \mathcal{M}_{\Phi^+} \times \mathcal{M}_{\Phi^-}$ [Hitchin, Graña, Gurrieri, Micu, Waldram, JL,...]

$$\mathbf{e}^{-\mathbf{K}_{\Phi^{+}}} = \int_{\mathbf{Y}} \langle \mathbf{\Phi}^{+}, \overline{\mathbf{\Phi}}^{+} \rangle = \mathbf{X}^{\mathbf{A}} \overline{\mathbf{F}}_{\mathbf{A}} - \overline{\mathbf{X}}^{\mathbf{A}} \mathbf{F}_{\mathbf{A}}$$

$$= \int_{\mathbf{Y}} \mathbf{J} \wedge \mathbf{J} \wedge \mathbf{J} , \quad \text{for} \quad \mathbf{\Phi}^{+} = \mathbf{e}^{\mathbf{B} + \mathbf{i} \mathbf{J}} ,$$

$$egin{array}{lll} \mathbf{e^{-K_{\Phi^{-}}}} &=& \int_{\mathbf{Y}} \langle \mathbf{\Phi^{-}}, \overline{\mathbf{\Phi}^{-}}
angle &=& \mathbf{Z^{I}} ar{\mathcal{F}_{I}} - \mathbf{ar{Z}^{I}} \mathcal{F}_{I} \ &=& \int_{\mathbf{Y}} \mathbf{\Omega_{3}} \wedge \overline{\mathbf{\Omega}_{3}} \;, \qquad \mathrm{for} \qquad \mathbf{\Phi^{-}} = \mathbf{\Omega_{3}} \end{array}$$

where $\langle \Phi^+, \overline{\Phi}^+ \rangle = \Phi_0^+ \wedge \overline{\Phi}_6^+ - \Phi_2^+ \wedge \overline{\Phi}_4^+ + \Phi_4^+ \wedge \overline{\Phi}_2^+ - \Phi_6^+ \wedge \overline{\Phi}_0^+$, etc. and $F_A(X)$, $\mathcal{F}_I(Z)$ are N=2 prepotentials

- $e^{-\mathbf{K}}$ is quartic invariant of O(6,6) (Hitchin functional)
- $rac{1}{2}$ for SU(3) same expression as in Calabi-Yau compactifications

Potential: is derived from Killing prepotential (or superpotential) \vec{P}

戊 ⅡA

$$\mathbf{P}^{1} + \mathbf{i}\mathbf{P}^{2} = e^{\frac{1}{2}\mathbf{K}_{\Phi}^{-} + \phi^{(4)}} \int_{Y_{6}} \langle \mathbf{\Phi}^{+}, d\mathbf{\Phi}^{-} \rangle , \qquad \mathbf{P}^{3} = -e^{2\phi^{(4)}} \int_{Y_{6}} \langle \mathbf{\Phi}^{+}, F_{\mathsf{A}} \rangle$$
$$F \equiv \sum_{\mathsf{RR-forms}} F^{\mathsf{RR}}$$

□ IIB

$$\mathbf{P^1} + i\mathbf{P^2} = e^{\frac{1}{2}\mathbf{K_{\Phi}^+} + \phi^{(4)}} \int_{Y_6} \langle \mathbf{\Phi}^-, d\mathbf{\Phi}^+ \rangle , \qquad \mathbf{P^3} = e^{2\phi^{(4)}} \int_{Y_6} \langle \mathbf{\Phi}^-, F_{\mathsf{B}} \rangle$$

Note:

- NS 3-form flux H included
- ullet gauged N=2 supergravity with charged fields in hypermultiplets
- as expected **P** depends on torsion and flux

'Formal aspects'

Mirror symmetry for generalized geometries corresponds to

$$\mathbf{\Phi}^+(\mathbf{Y}) = \mathbf{\Phi}^-(\mathbf{\tilde{Y}}) \qquad \mathbf{\Phi}^-(\mathbf{Y}) = \mathbf{\Phi}^+(\mathbf{\tilde{Y}})$$

manifest in the large volume (supergravity) limit.

Non-perturbative dualities with flux/torsion

Het. on
$$\mathbf{K3} \times \mathbf{T^2}$$
 with flux \leftrightarrow IIA on $\mathbf{SU(3)} \times \mathbf{SU(3)}$ can be embedded for M-Theory on $\mathbf{7d}\text{-}SU(3)$ -manifold [Aharony,Berkooz,Micu,JL, see Micu's talk]

non-geometric backgrounds [Hull; Shelton, Wecht, Taylor,...] can be included into $SU(3) \times SU(3)$ -structure formalism

'Formal aspects'

- Heterotic string compactified on SU(3)-structure manifolds [Becker, Becker, . . . ; Gurrieri, Lukas, Micu; Benmachiche, Martinez, JL]
- type II compactifications on SU(2)-structure manifolds $\Rightarrow \mbox{gauged } N=4 \mbox{ supergravity in the same formalism}$ [Spanjaard, Triendl, JL]
- \Rightarrow heterotic compactifications on SU(2)-structure manifolds [Martinez,Micu,JL]
- RR-scalars can be incorporated into geometrical description $\Rightarrow E_7$ -covariant formulation of N=2 backgrounds

[Graña, Waldram, JL]

Orientifolds

[Benmachiche, Grimm, JL, Shelton, Taylor, Wecht, Micu, Palti, Tasinato, ...]

Orientifold projection: $N=2 \rightarrow N=1$

- truncates spectrum
- ullet chooses complex structure on field space and selects a Kähler subspace in \mathcal{M}_{QK}
- repackages P into W&D
 - large volume/large complex structure limit:

$$\mathbf{W} \sim \kappa_{\mathbf{abc}} \phi^{\mathbf{a}} \phi^{\mathbf{b}} \phi^{\mathbf{c}}$$

IIA: W cubic in Kähler moduli, linear in c.s. moduli IIB: W cubic in c.s. moduli, linear in Kähler moduli

 $-\ W$ depends on both type of moduli

N=1 de Sitter backgrounds [Covi, Gomez-Reino, Gross, Palma, Scrucca, JL]

motivation: relevant for early and late universe

problem: sGoldstino is potentially tachyonic

mass is entirely determined by Kähler geometry:

$$egin{array}{cccc} rac{\Lambda}{\mathbf{m_{3/2}^2 M_{PL}^2}} \gg \mathbf{1} & \Rightarrow & \mathbf{Sec} \leq 0 \ & & & & & \\ rac{\Lambda}{\mathbf{m_{3/2}^2 M_{PL}^2}}
ightarrow \mathbf{0} & \Rightarrow & \mathbf{Sec} \leq rac{2}{3} \end{array}$$

sectional curvature: $\mathbf{Sec} := |\mathbf{G}|^{-4} \, \mathbf{R}_{\mathbf{i}\bar{\mathbf{i}}\mathbf{m}\bar{\mathbf{n}}} \mathbf{G}^{\mathbf{i}} \mathbf{G}^{\mathbf{j}} \mathbf{G}^{\mathbf{m}} \mathbf{G}^{\bar{\mathbf{n}}}$

Summary

- ightharpoonup discussed backgrounds with $SU(3) \times SU(3)$ structure
 - NS scalar manifold is product of special geometries $\mathbf{K} = \mathbf{K}_{\Phi^+} + \mathbf{K}_{\Phi^-} \text{ is given by Hitchin functionals}$ and independent of torsion
 - (super)potential depends on torsion and background fluxes
 - generalized mirror symmetry intact
 - non-geometric backgrounds can be included into formalism

⇒ phenomenological questions:

- stabilization of moduli
- inclusion of D-branes
- supersymmetry breaking
- inflationary scenarios