Status of D6-brane Flux Models and their Effective Field Theories

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Outline

- Short introduction
- Mini-review on known type-IIA vacua
- Consistency constraints in compactifications
- Constructing the EFT
 - bulk sector
 - vector fields
 - chiral sector
- Leftover and conclusions

Renewed interests in type-IIA compactifications:

- New 4d classical vacua (more generic than in type-IIB and Heterotic cases)
- Improvement in understanding of generalized geometries (i.e. non CY compactifications with fluxes and torsion...)
- AdS₄/CFT₃ in IIA/M-theory

(may provide non-perturbative description of IIA string in AdS backgrounds)

Intersecting brane model building – phenomenology

(easy geometrical construction of chiral gauge theories)

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Final goal:

- Find a "reliable" string construction (aka vacuum) compatible with experiments -

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(start from the strings, find vacuum, stab. moduli...)

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– bottom-up

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key – "try to find a good compromise":

derive an EFT from string theory *such that* it allows to implement both the constraints and the corrections from string theory *and* to model a successful low energy description

Mini-review on type-IIA 'classical' known vacua I

• Susy AdS₄:

- Many types:
 - Toroidal Orbifolds ($T^6/Z_2 \times Z_2$, $T^6/Z_3 \times Z_3$, ...);
 - Group manifolds (S3xS3, twisted tori...);
 - Coset manifolds (CP₃, SU(3)/U(1)xU(1), ...);
 - ...
- All moduli stabilized;
- infinite number, with small g_s large V_6 ;
- good for AdS/CFT;
- not good for phenomenology;

GV, Zwirner, DeWolfe et al, Camara, Font, Ibanez...

> Behrndt, Cvetic, Lust, Tsimpis...

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Non susy AdS₄:

(twisted)-Toroidal Orbifolds (+fluxes+sources);

Camara, Font, Ibanez...

- All moduli stabilized;
- good for phenomenology (uplifting with D-terms?, other F-terms?...);

GV, Zwirner, DeWolfe et al, Camara, Font, Ibanez...

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Mini-review on type-IIA 'classical' known vacua II

...continue

• Susy Mink₄:

- flat directions with geometrical fluxes;
- All moduli fixed with <u>non-geometrical fluxes?</u>
- large V_6 and small g_s ? (Micu, Palti, Tasinato '07);
- not good for phenomenology:

Aldazabal et al.

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No-scale models:

• (twisted)-Toroidal Orbifolds (+fluxes+sources)

see e.g. Derendinger et al.,... Camara, Grana...

- good for phenomenogy (after no-scale moduli stabilized radiatively...)
- Ω-dominated? (*J*-dominated? non geom.)

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dS4 vacua?

• Silverstein '07 (EFT?)

From the compactification to the EFT

closed strings \rightarrow metric [$g(J,\Omega)$], dilaton [Φ], $[B_{\rm NS}]$ -field, RR-fields [$C^{(p)}$], ... open strings \rightarrow D6-D6 gauge sector [A_{μ} , Z^k],

D6-D6' chiral sector [ϕ^i]

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N=1 4d Effective SUGRA determined by:

- Kahler function $G(q,q^*) = K(q,q^*) + \log|W(q)|^2$
- Gauge kinetic function $f_{AB}(q)$
- Killing vectors X(q)

Most of the 4d EFT can be derived from dim.red. from 10d SUGRA EFT only partially known for generic background Explicit formulae known for particular orbifolds

RR-tadpole ⇔ RR Bianchi Id.

$$d_H G = Q_{RR} \quad \Rightarrow \quad \omega G^{(2)} + H \wedge G^{(0)} = \Sigma [\pi_6]$$

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$$d_H d_H = 0, \Rightarrow \omega \omega = 0, \omega H = 0$$

with sources
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Localized tadpoles ⇔ loc. Bianchi Id.

$$d_H Q_{RR} = 0 \Rightarrow \omega[\pi_6] = 0, H \wedge [\pi_6] = 0$$
 $Q_{NSNS} G = 0 \Rightarrow [v_5] \wedge G = 0, [\kappa_5] G = 0$
(+ branes on branes)

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Crucial for gauge (+susy) inv. of the EFT:

- classical: gauge inv. of W
- perturbative: cancellation of 1-loop anomalies
- non-perturbative: gauge inv. of instanton corrections

1. EFT for bulk fields

Kahler potential:
$$K = -\log\left(\frac{1}{3!}\int J \wedge J \wedge J\right) - 2 \log\left(\frac{1}{3!}\int \Omega' \wedge \overline{\Omega}'\right)$$

From dim. red. on $T^6/Z_2 \times Z_2 \Rightarrow$

$$K = -\log(su_1u_2u_3) - \log(t_1t_2t_3) + \dots$$
 (twisted sector)

see Lust, Reffert, Stieberger...

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Superpotential:

$$W = \frac{1}{4} \int_{\mathcal{M}_6} e^{iJ^c} \left(\overline{\mathbf{G}} - id_H \Omega'^c \right)$$

$$J^c \equiv J + i B$$
, $\Omega'^c \equiv \operatorname{Re} \Omega' + i C^{(3)}$,

$$d_H \equiv d + \omega + \overline{H}$$
 $\overline{\mathbf{G}} = \sum_{p=even} \overline{G}^{(p)}$

D2 domain wall argument: $\Delta W \sim \int \Delta G^{(6)}$

a la' Gukov Vafa Witten

$$W = \int G^{(6)}$$

with $G^{(6)}$ complexification of the solution of $d_HG=0$

Gauge kinetic function: (from DBI+WZ)

$$f_{ab} = \delta_{ab} N_a \left(\operatorname{Re} \Omega'_{\pi_a} + i \int_{\pi_a} C^{(3)} \right) = \delta_{ab} N_a \Omega'^c_{\pi_a}$$

From dim. red. on $T^6/Z_2 \times Z_2 \Rightarrow$

$$f = N T_6 \left(m_1 m_2 m_3 S - m_1 n_2 n_3 U_1 - n_1 m_2 n_3 U_2 - n_1 n_2 m_3 U_3 \right)$$

Cremades, Ibanez, Marchesano '02

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Killing vectors: (of the gauging)

$$i X^{S} = -2 N \mu_{6} n_{1} n_{2} n_{3} = -2 N \mu_{6} q^{0},$$

$$i X^{U_{A}} = 2 N \mu_{6} n_{A} m_{B} m_{C} = 2 N \mu_{6} q^{A},$$

D-terms:
$$(= i K_a X^a)$$

$$\operatorname{Im}\widetilde{\Omega}_{\pi} = \frac{\sqrt{su_1u_2u_3} \ D}{N T_6}$$

$$D = N \mu_6 \left(\frac{n_1 n_2 n_3}{s} - \sum_{A=1}^3 \frac{n_A m_B m_C}{u_A} \right) = N \mu_6 \left(\frac{q^0}{s} - \sum_{A=1}^3 \frac{q^A}{u_A} \right)$$

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...but from DBI reduction on generic configurations

Re
$$f = N T_6 \sqrt{(\operatorname{Re} \widetilde{\Omega}_{\pi})^2 + (\operatorname{Im} \widetilde{\Omega}_{\pi})^2}$$

$$V_D = \frac{1}{2} \frac{1}{N T_6 \operatorname{Re} \widetilde{\Omega}_{\pi}} D^2 \frac{2}{\sqrt{1 + \left(\frac{\operatorname{Im} \widetilde{\Omega}_{\pi}}{\operatorname{Re} \widetilde{\Omega}_{\pi}}\right)^2 + 1}}$$

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Right SUGRA results only in the 2-der. limit

i.e.
$$\left| \frac{\operatorname{Im} \widetilde{\Omega}_{\pi}}{\operatorname{Re} \widetilde{\Omega}_{\pi}} \right| \ll 1$$

GV, Zwirner '06

lesson for anti-D-branes....

2. EFT for the D6-D6 sector

Kahler potential: for arbitrary configuration of D6-branes

From dim.red. of DBI action (*involving highly non-trivial field-redefinition*) (assuming $S_{DBI} = \Sigma_k S_{DBI}(D6_k)$ from dim. red. on T^6/Z_2xZ_2)

$$K = -\log\left[t_1t_2t_3\left(s + \sum_{A=1}^{3} p_A^a \frac{z_A^{i_a 2}}{t_A}\right) \prod_{A=1}^{3} \left(u_A - p_0^a \frac{z_A^{i_a 2}}{t_A} + \sum_{B,C=1}^{3} d_{ABC} p_B^a \frac{z_C^{i_a 2}}{t_C}\right)\right]$$

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up to quadratic fluctuations it gives:

$$K = -\log(t_1 t_2 t_3 s u_1 u_2 u_3) + p_0^a \sum_{A=1}^3 \frac{z_A^{i_a 2}}{t_A u_A} - \sum_{A=1}^3 \left(p_A^a \frac{z_A^{i_a 2}}{s t_A} + \sum_{B,C=1}^3 d_{ABC} p_B^a \frac{z_A^{i_a 2}}{t_A u_C} \right) + \mathcal{O}(z^4)$$

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Notice that $z_A = Re Z_A$

i.e. there are Giudice-Masiero-like terms (like in Heterotic) good for μ -terms

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$$(m_A^2 + n_A^2)\sqrt{\frac{\Upsilon_B \Upsilon_C}{\Upsilon_A}} \qquad \qquad \Upsilon_A = \frac{m_A^{a\,2}}{u_B' u_C'} + \frac{n_A^{a\,2}}{s' u_A'}$$

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Notice also that mixed $O(z^2z^2)$ terms do not agree in the N=2 limit with the prepotential proposed by Ferrara, Minasian, Sagnotti,... Antoniadis et al... connected with ρ -problem, see below...

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from dim.red. + DBI fluctuations

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$$\omega_{YM}^c = Z^{i_a} \wedge \omega Z^{i_a} + \frac{2}{3} i f_{i_a i_b i_c} Z^{i_a} \wedge Z^{i_b} \wedge Z^{i_c}$$

$$Z^{i_a} = \left[Z_1^{i_a} \frac{n_1^a dx^5 + m_1^a dx^6}{\sqrt{(m_1^a)^2 + (n_1^a)^2}} + Z_2^{i_a} \frac{n_2^a dx^7 + m_2^a dx^8}{\sqrt{(m_2^a)^2 + (n_2^a)^2}} + Z_3^{i_a} \frac{n_3^a dx^9 + m_3^a dx^{10}}{\sqrt{(m_3^a)^2 + (n_3^a)^2}} \right]$$

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in agreement with *Martucci, Marchesano, Camara, Grana* but valid also for susy-breaking and/or non-Mink vacua, automatically holomorphic

D-terms:

• D6-D6 sector uncharged under U(1) part

$$D_a = \frac{q_a^0}{s'} - \sum_{A=1}^3 \frac{q_a^A}{u_A'} = \frac{q_a^0}{s + \sum_{A=1}^3 p_A^a \frac{z_A^{i_a \, 2}}{t_A}} - \sum_{A=1}^3 \frac{q_a^A}{u_A - p_0^a \frac{z_A^{i_a \, 2}}{t_A} + \sum_{B,C=1}^3 d_{ABC} p_B^a \frac{z_C^{i_a \, 2}}{t_C}}$$

Only non-abelian D-term

$$D_{i_a} = iK_{Z_A^{i_b}}(X_{i_a})_A^{i_b} = Q_A^a f_{i_a j_a k_a} \frac{z_A^{j_a} \zeta_A^{k_c}}{t_A} \qquad Q_A^a = \frac{p_0^a}{u_A'} - \frac{p_A^a}{s'} - \frac{p_B^a}{u_C'} - \frac{p_C^a}{u_B'}$$

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Gauge kinetic function:

$$f_{a} \propto \Omega'_{\pi_{a}}^{c} = \Omega_{\pi_{a}}^{c} + \sum_{b} \int_{\pi_{a}} Z^{i_{b}} \phi^{i_{b}}[\pi_{b}] = \Omega_{\pi_{a}}^{c} + \sum_{b} (non \ hol.) \times I_{ab}$$

dual to the type-IIB " ρ -problem":

- solution 1: (Berg, Haack, Kors '04) 1-loop from open strings (see Akerblom et al. '07 for IIA)
- solution 2: (Baumann et al. '06) tree level from closed-strings inter-brane backreaction

D-terms:

D6-D6 sector uncharged under U(1) part

$$D_a = \frac{q_a^0}{s'} - \sum_{A=1}^3 \frac{q_a^A}{u_A'} = \frac{q_a^0}{s + \sum_{A=1}^3 p_A^a \frac{z_A^{i_a \, 2}}{t_A}} - \sum_{A=1}^3 \frac{q_a^A}{u_A - p_0^a \frac{z_A^{i_a \, 2}}{t_A} + \sum_{B,C=1}^3 d_{ABC} p_B^a \frac{z_C^{i_a \, 2}}{t_C}}$$

Only non-abelian D-term

$$D_{i_a} = iK_{Z_A^{i_b}}(X_{i_a})_A^{i_b} = Q_A^a f_{i_a j_a k_a} \frac{z_A^{j_a} \zeta_A^{k_c}}{t_A} \qquad Q_A^a = \frac{p_0^a}{u_A'} - \frac{p_A^a}{s'} - \frac{p_B^a}{u_C'} - \frac{p_C^a}{u_B'}$$

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Gauge kinetic function:

$$f_{a} \propto \Omega'_{\pi_{a}}^{c} = \Omega_{\pi_{a}}^{c} + \sum_{b} \int_{\pi_{a}} Z^{i_{b}} \phi^{i_{b}}[\pi_{b}] = \Omega_{\pi_{a}}^{c} + \sum_{b} (non \ hol.) \times I_{ab}$$

dual to the type-IIB " ρ -problem":

- solution 1: (Berg, Haack, Kors '04) 1-loop from open strings (see Akerblom et al. '07 for IIA)
- solution 2: (Baumann et al. '06) tree level from closed-strings inter-brane backreaction

Conclusions:

EFT is known up to quadratic fluctuations, higher-power terms only partially known, not under control...

(not from supergravity)

Kahler potential: from string amplitudes (Lust et al... Bertolini et al...)

$$K_{ij}^{ab} = \delta_{ij} S^{-\alpha} \prod_{I=1}^{3} U_{I}^{-(\beta+\xi \theta_{ab}^{I})} T_{I}^{-(\gamma+\zeta \theta_{ab}^{I})} \sqrt{\frac{\Gamma(\theta_{ab}^{1})\Gamma(\theta_{ab}^{2})\Gamma(1+\theta_{ab}^{3})}{\Gamma(1-\theta_{ab}^{1})\Gamma(1-\theta_{ab}^{2})\Gamma(-\theta_{ab}^{3})}},$$

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$$K_{ij}^{ab} = \delta_{ij} S^{-\frac{1}{4}} \prod_{J=1}^{3} U_{J}^{-\frac{1}{4}} T_{J}^{-\left(\frac{1}{2} \pm \frac{1}{2} \operatorname{sign}(I_{ab}) \theta_{ab}^{J}\right)} \sqrt{\frac{\Gamma(\theta_{ab}^{1}) \Gamma(\theta_{ab}^{2}) \Gamma(1 + \theta_{ab}^{3})}{\Gamma(1 - \theta_{ab}^{1}) \Gamma(1 - \theta_{ab}^{2}) \Gamma(-\theta_{ab}^{3})}},$$

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Superpotential:

$$W = Y_{ijk} \phi^{i} \phi^{j} \phi^{k} + \mu_{1ijk} \phi^{i} \phi^{j} Z^{k} + \mu_{2ij} \phi^{i} \phi^{j} + \dots$$

$$Y_{ijk} \text{ from w.s. instantons,} \\ \text{(~IIB perturbative calculus)} \qquad Y_{ijk} = h_{qu} \sigma_{abc} \prod_{r=1}^n \vartheta \begin{bmatrix} \delta^{(r)} \\ \phi^{(r)} \end{bmatrix} (\kappa^{(r)}) \\ \text{(~IIB perturbative calculus)} \qquad Cremades, Ibanez, Marchesano, Cvetic, Papadimitriou (~IIB) (~I$$

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Killing vectors: known (gauge transformation properties of chiral fields)

Gauge kinetic functions: (corrections are higher order, from threshold....)

- KK5-monopoles:
 - $\omega \omega \neq 0$ (alternative description of twisted sector, stabilization, ...)
 - Phenomenology (N=1 configurations)
 - dual to D6-branes
 - D6-KK5 intersecting models → KK6 M-theory configuration
 - May help with dS vacua? (Silverstein '07)

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- Other important subjects not discussed here:
 - α' , g_s corrections to EFT (threshold corrections...);
 - Non-perturbative corrections (E2-instantons...);
 - Inflation...

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Next step – Start putting everything together....