



# XI<sup>th</sup> Quark Confinement and the Hadron Spectrum

Modeling the influence of string collective phenomena on the long range rapidity correlations between the transverse momentum and the multiplicities

Evgeny Andronov, Vladimir Vechernin

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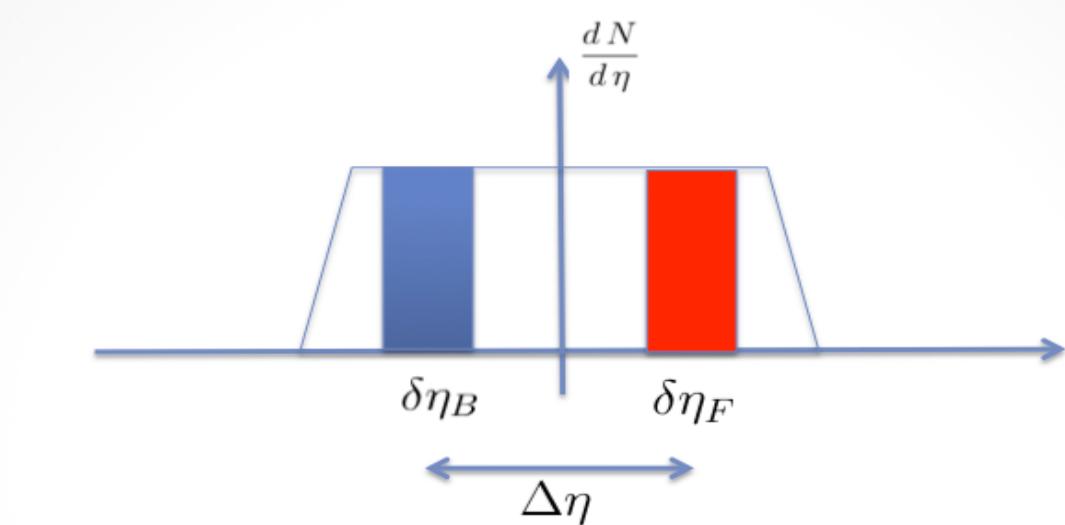
SPbSU, Department of High Energy and Elementary Particles Physics,  
Laboratory of Ultra-High Energy Physics

St. Petersburg, 12/09/14

# Outline

- Introduction
- Model with two types of strings
- Results for the n-n and pT-n correlation parameters
- Summary

# Long-range pseudorapidity correlations



Two definitions of the correlation parameter for two observables B and F

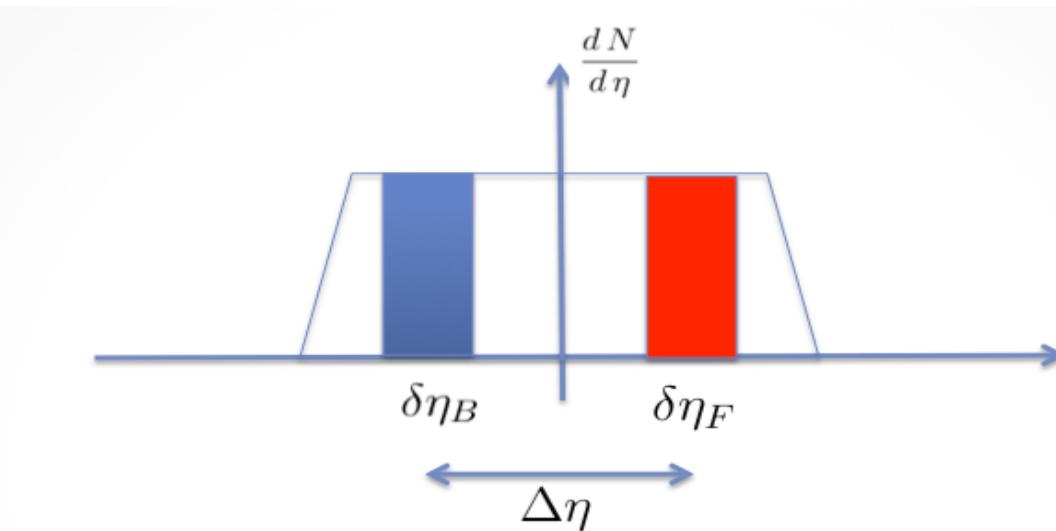
1) Linear regression

$$\langle B \rangle_F = a + b \cdot F$$

2) Correlator formula

$$b = \frac{\langle B \cdot F \rangle - \langle B \rangle \cdot \langle F \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

# Long-range pseudorapidity correlations



Two definitions of the correlation parameter for two observables B and F

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$$b = \frac{\langle B \cdot F \rangle - \langle B \rangle \cdot \langle F \rangle}{\langle F^2 \rangle - \langle F \rangle^2}$$

# Long-range pseudorapidity correlations

## Observable types

n – charged particles multiplicity

$$p_t = \frac{1}{n} \sum_{i=1}^n p_t^{(i)}$$

- event mean value of transverse momentum

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## Correlation types

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- $pT$ - $n$
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# Long-range pseudorapidity correlations

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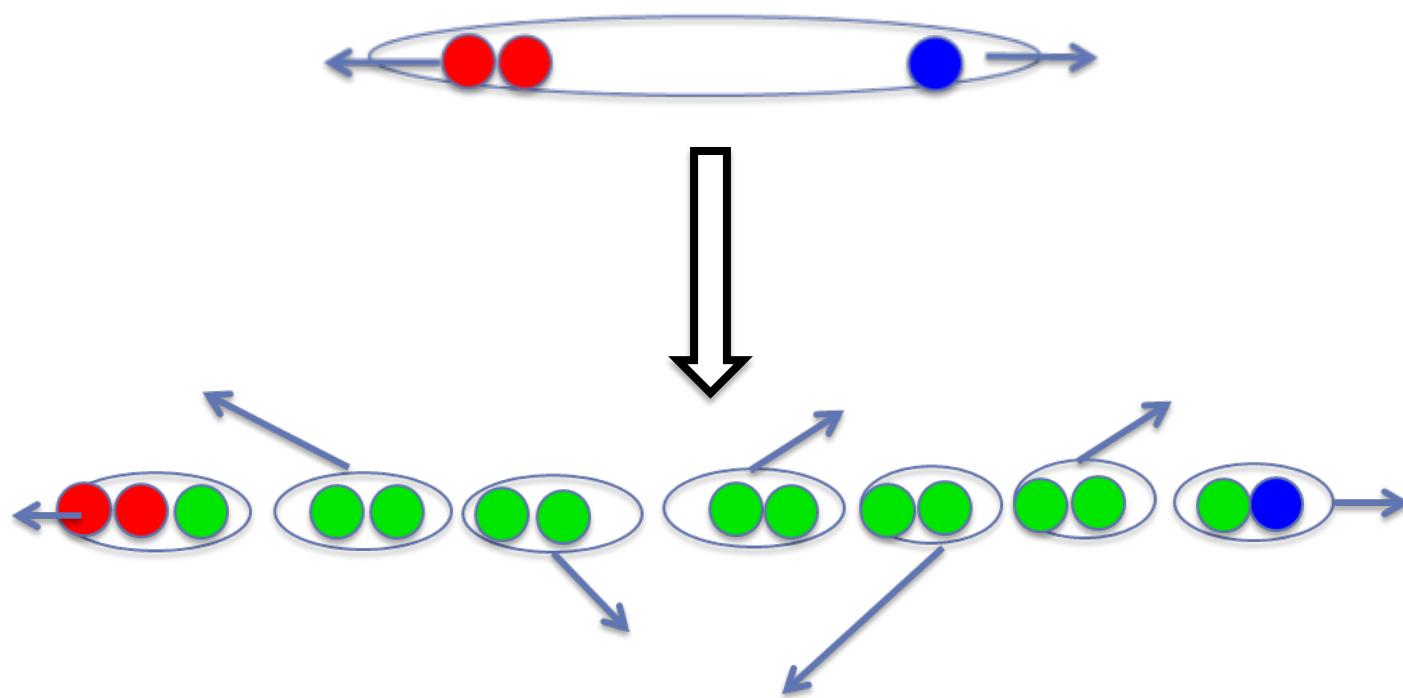
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## Correlation types

- $n$ - $n$
- $pT$ - $n$
- ~~$pT$ - $pT$~~

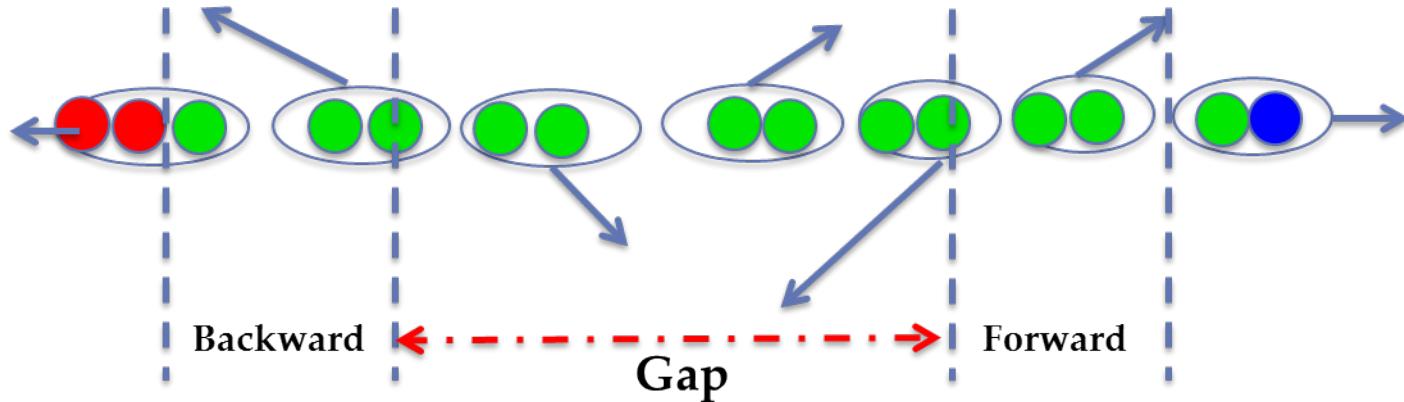
# Mechanism of particle production in the model with independent strings



A.Capella, U.P.Sukhatme, C.I.Tan and J.Tran Thanh Van, Phys. Lett. B81 (1979) 68;  
Phys. Rep. 236 (1994) 225.

A.B.Kaidalov, Phys. Lett., 116B(1982)459

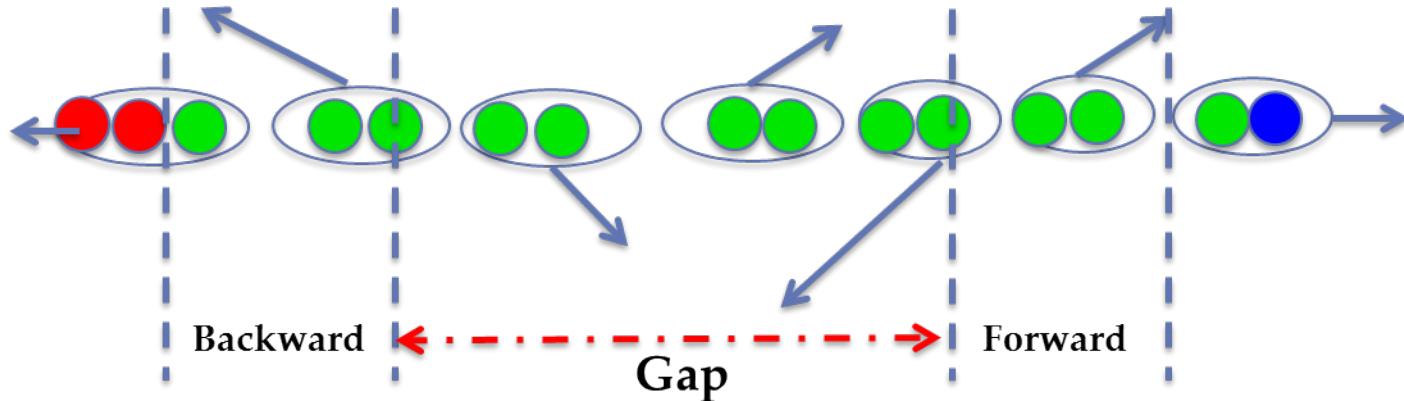
# Single string case



$$\langle n_B \rangle = \sum_B B \cdot P(B) = \langle n_B \rangle_{one}$$

$$\langle n_F \rangle = \sum_F F \cdot P(F) = \langle n_F \rangle_{one}$$

# Single string case



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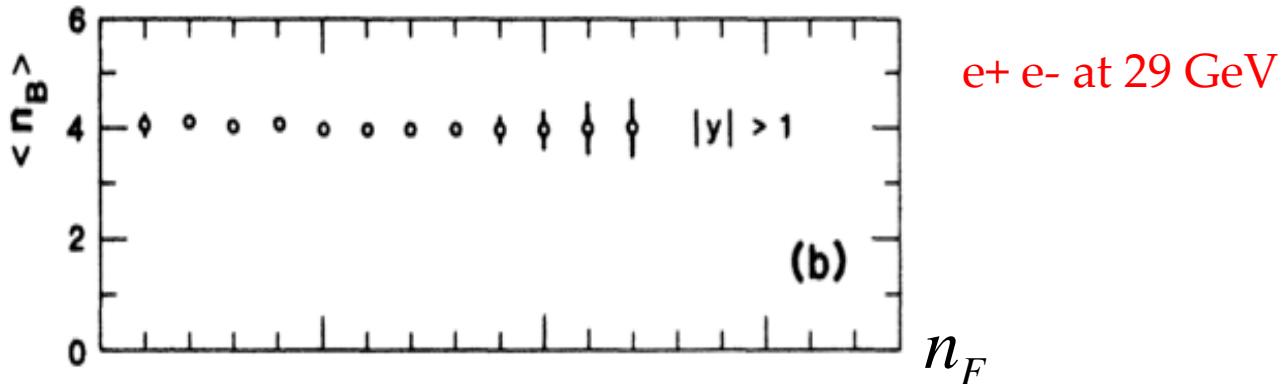
$$\langle n_F \rangle = \sum_F F \cdot P(F) = \langle n_F \rangle_{one}$$

In case of sufficiently large gap between windows one  
string produces particles in both windows  
**independently!**

$$(P(F, B) = P(B) \cdot P(F))$$

$$\langle n_B n_F \rangle = \sum_{F, B} F \cdot B \cdot P(F, B) = \langle n_B \rangle \cdot \langle n_F \rangle \rightarrow b_{n-n} = 0$$

# Single string case



$$\langle n_B \rangle = \sum_B B \cdot P(B) = \langle n_B \rangle_{one}$$

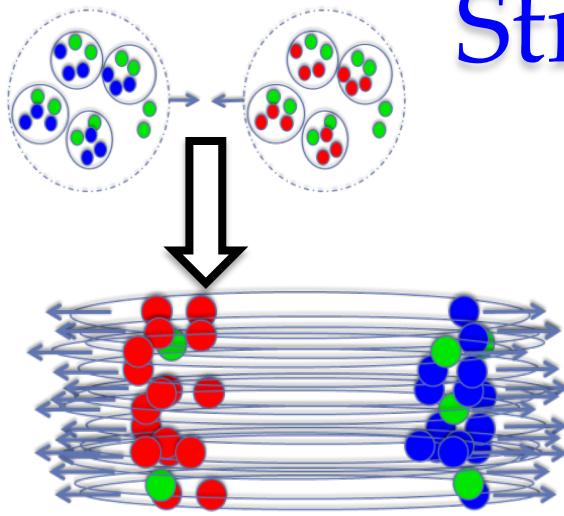
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Experiment:  $b = 0.002 \pm 0.006$  Phys.Rev.D vol.34, num.11(1986)

# String fusion model



Single string

String  
( $r \sim 0.2\text{fm}$ )

Overlapping strings



Multiplicity

$$\langle n \rangle_{one}$$

$$\langle n \rangle_{new} = \sqrt{N_{str}} \langle n \rangle_{one}$$

Transverse  
momentum

$$\langle p_t \rangle_{one}$$

$$\langle p_t^2 \rangle_{new} = \sqrt{N_{str}} \langle p_t^2 \rangle_{one}$$

M.A.Braun and C.Pajares, Phys. Rev. Lett. **85** (2000) 4864;

M.A.Braun and C.Pajares, Phys. Lett. **B287** (1992) 154; Nucl. Phys. **B390** (1993) 542, 549;

N.S.Amelin, M.A.Braun and C.Pajares, Phys. Lett. **B306** (1993) 312;

M.A.Braun, C.Pajares and V.V.Vechernin, Internal Note/FMD ALICE---INT---2001---16

# Long-range correlations. General remarks.

- LRC are governed by the fluctuations in number of strings and by the string fusion effects
- n-n correlation coefficient is zero without these fluctuations and fusion effects
- pT-n correlation coefficient is zero without fusion effects.

See also at indico:  
1) 09/09/14 talk by V.Kovalenko  
2) 09/09/14 poster by I. Altsybeev  
3) 09/09/14 poster by D. Neverov



# Model with two types of strings

N primary strings (N – even)

$$P_N(N_2) = C_{N/2}^{N_2} r^{N_2} (1 - r)^{N/2 - N_2}$$

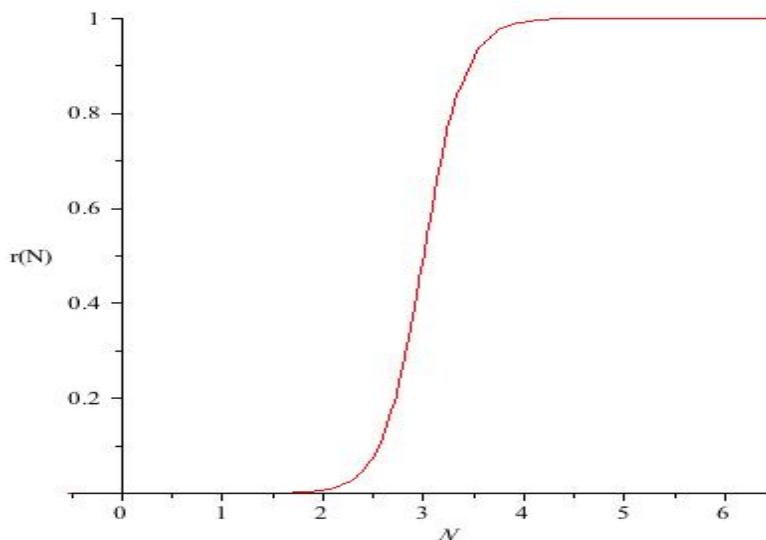
$$N_2 ; N_1 = N - 2N_2$$

Analytical results for n-n correlation coefficient  
Only negative pT-n correlations!

E. Andronov, V. Vechernin, PoS(QFTHEP2013), 054 (2014).

$$r(N) = \frac{1}{1 + e^{-\frac{N - shift}{slope}}}$$

Only MC simulations



## Long-range n-n correlation parameter. Monte-Carlo simulations.

$$\omega [N] = \frac{\bar{N}}{D_N}$$

- mean value of the number of primary strings

- scaled variance of the number of primary strings

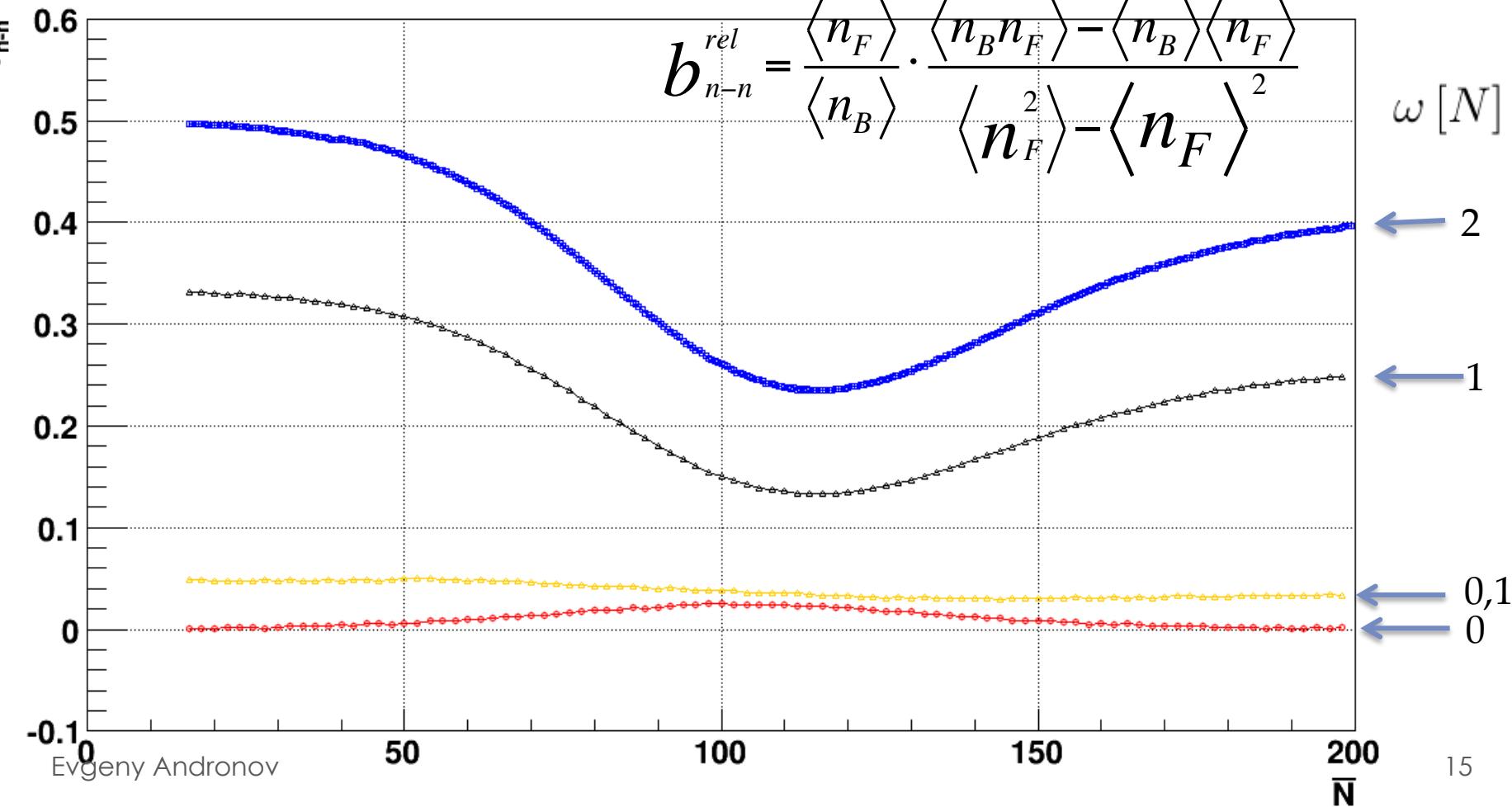
Shift=100

Slope=20

$$\bar{\mu} = D_\mu = 0.5$$

- parameters of string decay

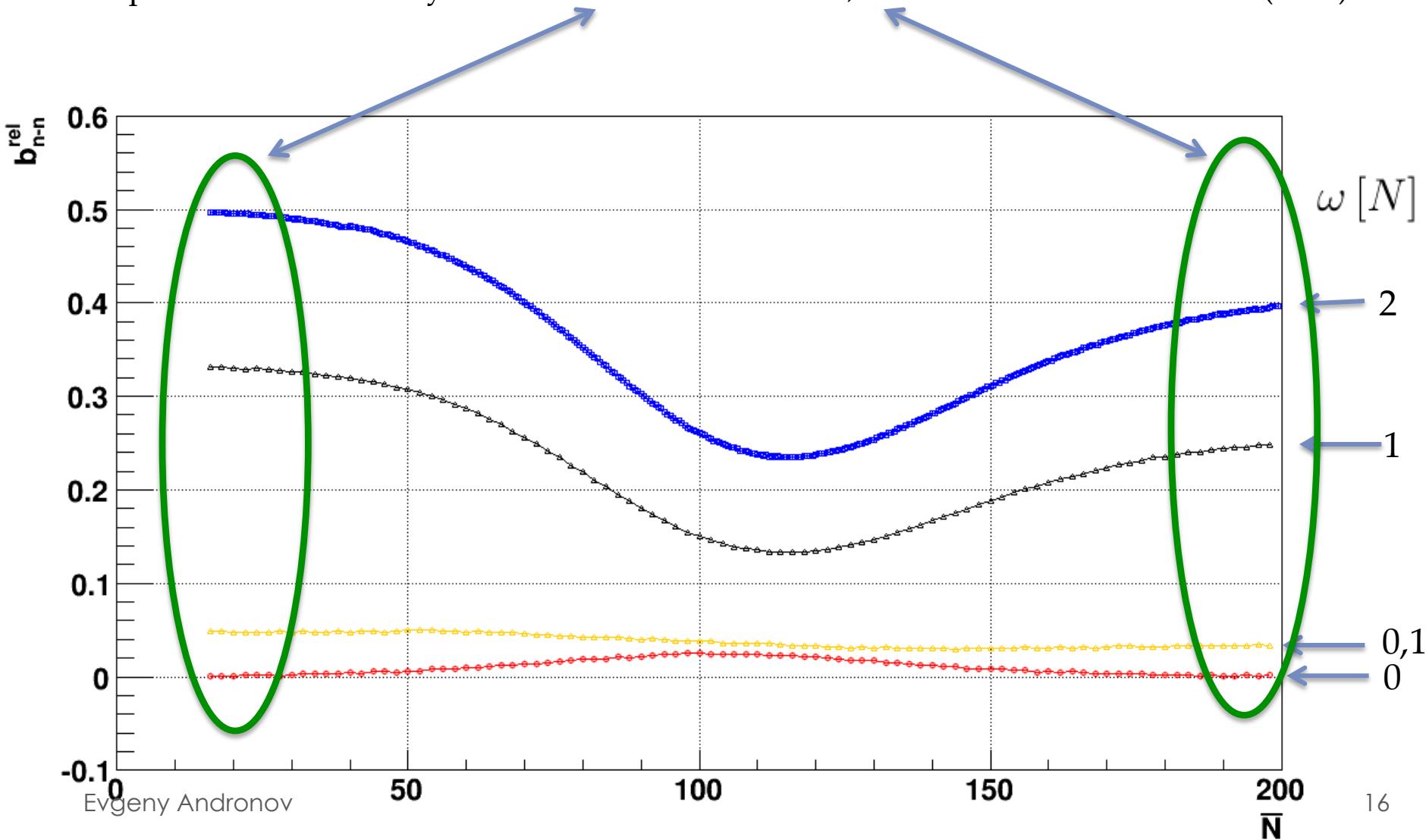
$$b_{n-n}^{rel} = \frac{\langle n_F \rangle}{\langle n_B \rangle} \cdot \frac{\langle n_B n_F \rangle - \langle n_B \rangle \langle n_F \rangle}{\langle n_F^2 \rangle - \langle n_F \rangle^2}$$



# Long-range n-n correlation parameter. Monte-Carlo simulations.

Single type of the string limits ( $r=0$  or  $r=1$ )

Correspondence to the analytical results – V.V. Vechernin, Proc. Of XX Baldin ISHEPP (2011)

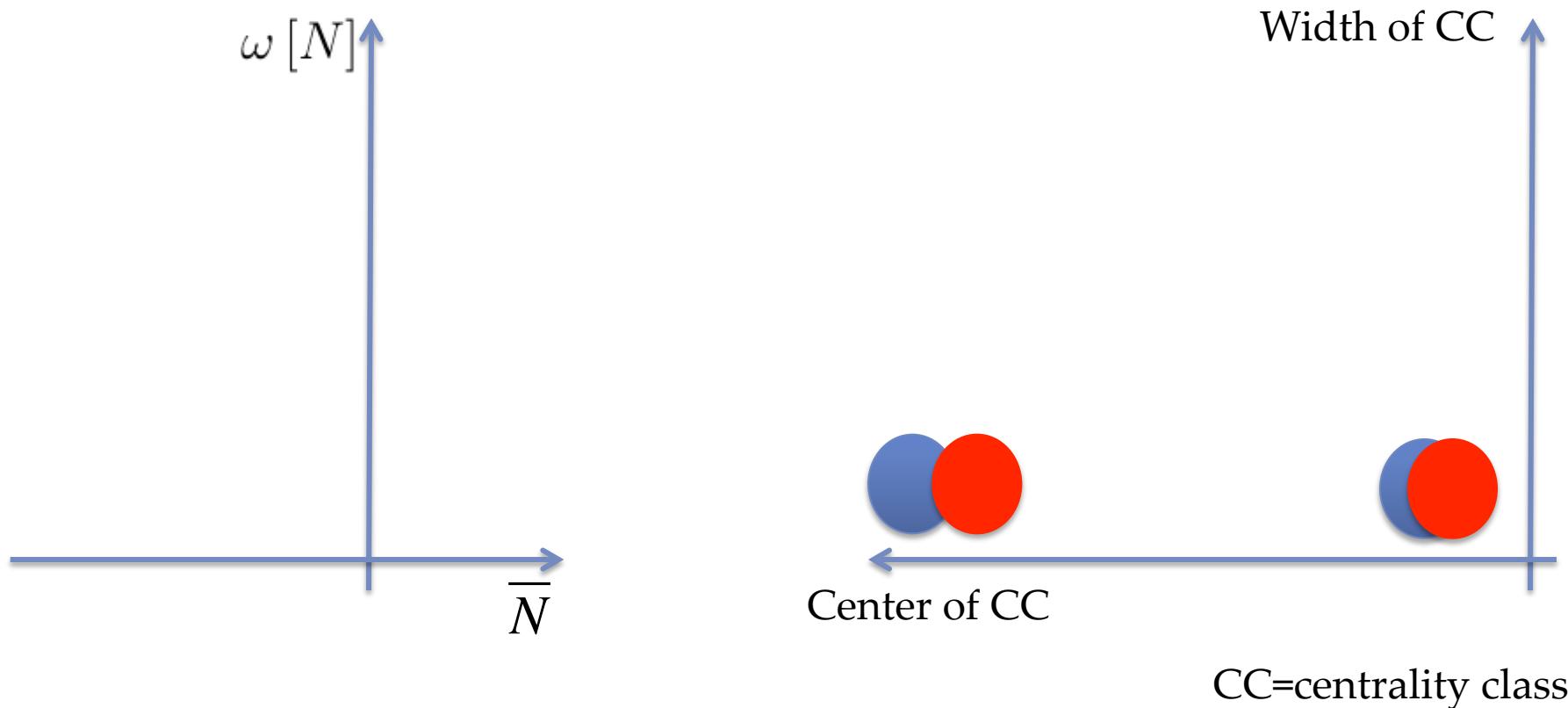


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## Long-range n-n correlation parameter. Connection with experiment.

By varying the width and the position of the centrality class one can scan our plot in two directions and search for the predicted effects



## Long-range pT-n correlation parameter. Monte-Carlo simulations.

$$\omega [N] = \frac{D_N}{\bar{N}}$$

- mean value of the number of primary strings

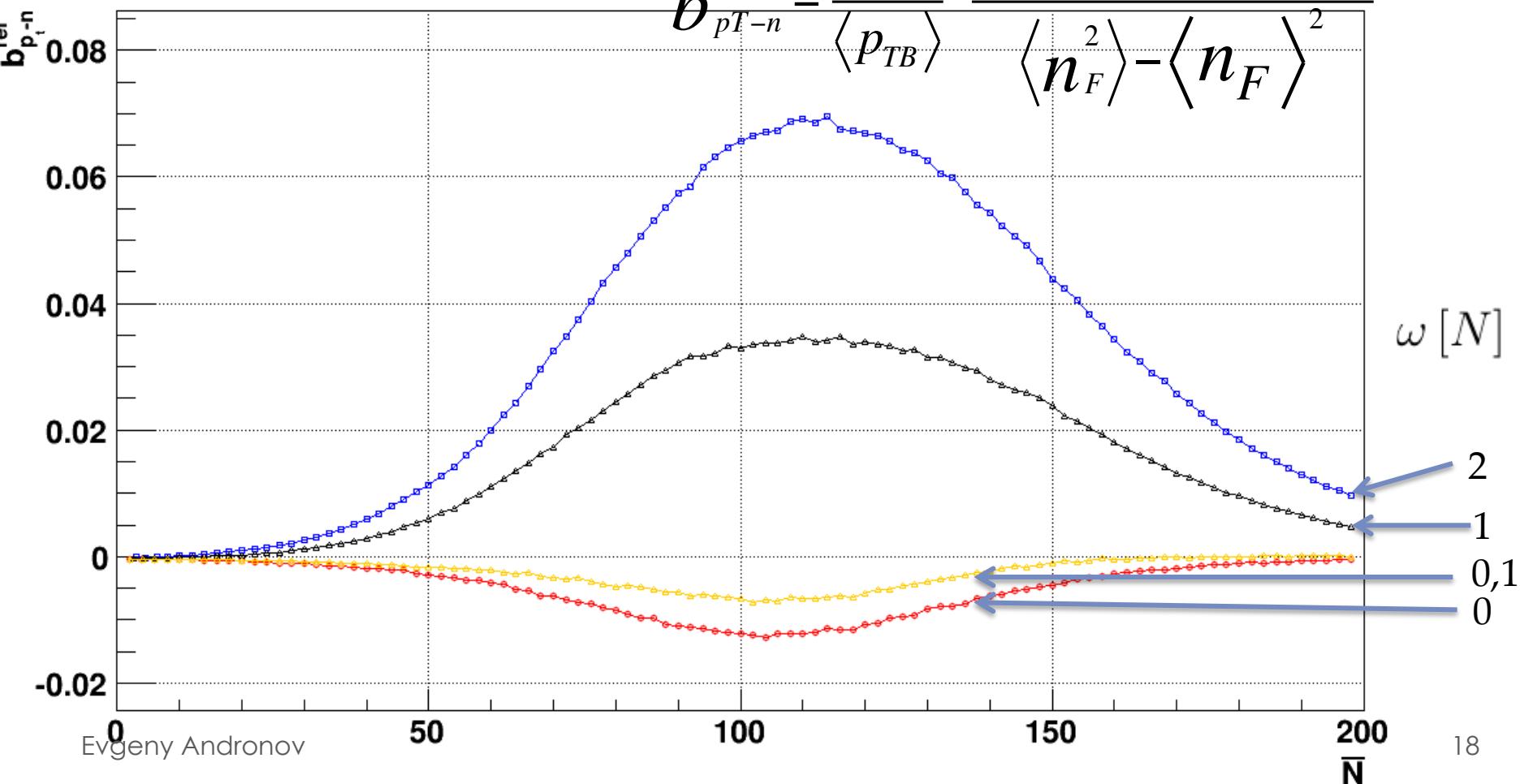
Shift=100

- the scaled variance of the number of primary strings

Slope=20

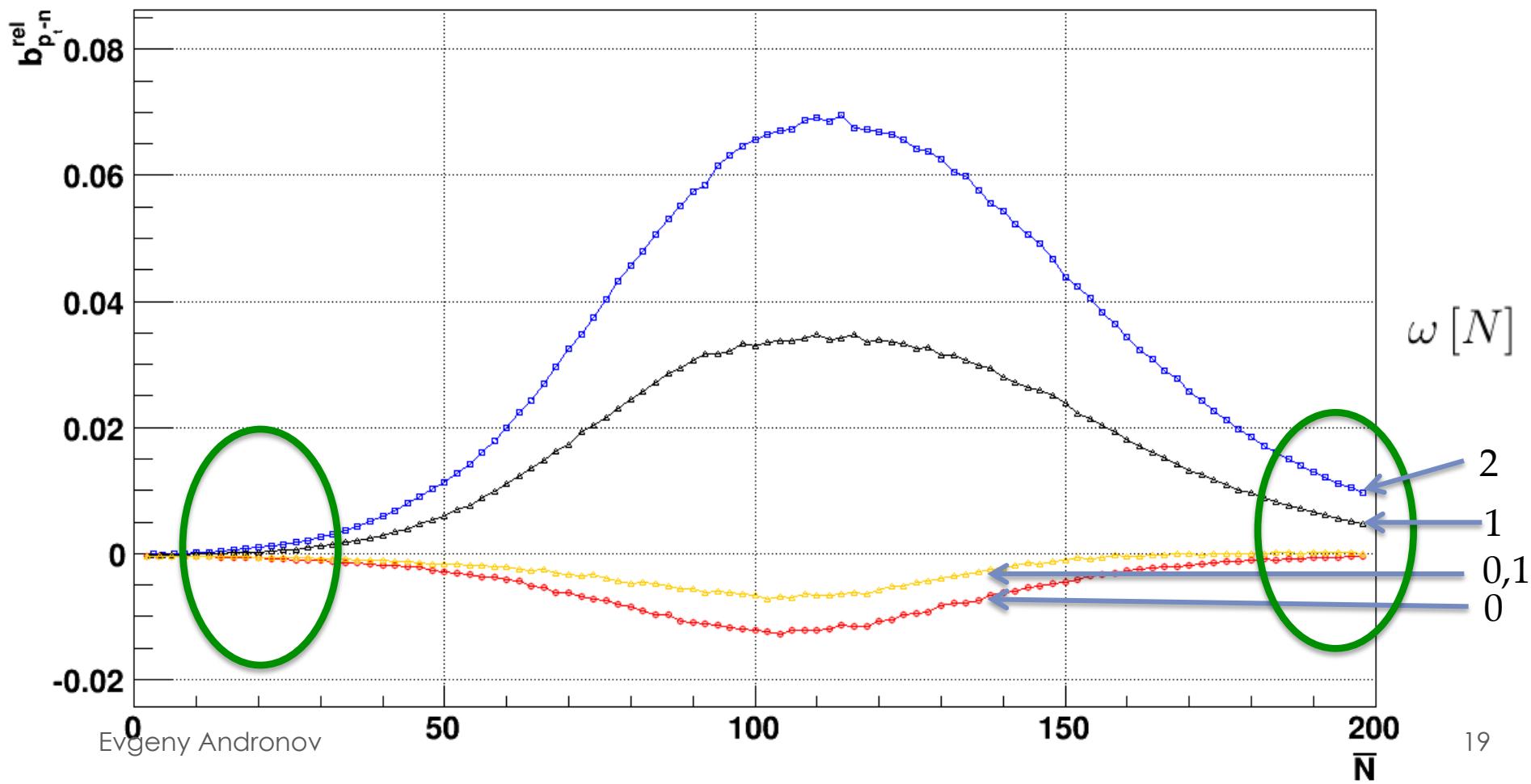
$$\bar{\mu} = D_\mu = 0.5 \quad \text{- parameters of string decay}$$

$$b_{pT-n}^{rel} = \frac{\langle n_F \rangle}{\langle p_{TB} \rangle} \cdot \frac{\langle p_{TB} n_F \rangle - \langle p_{TB} \rangle \langle n_F \rangle}{\langle n_F^2 \rangle - \langle n_F \rangle^2}$$



# Long-range pT-n correlation parameter. Monte-Carlo simulations.

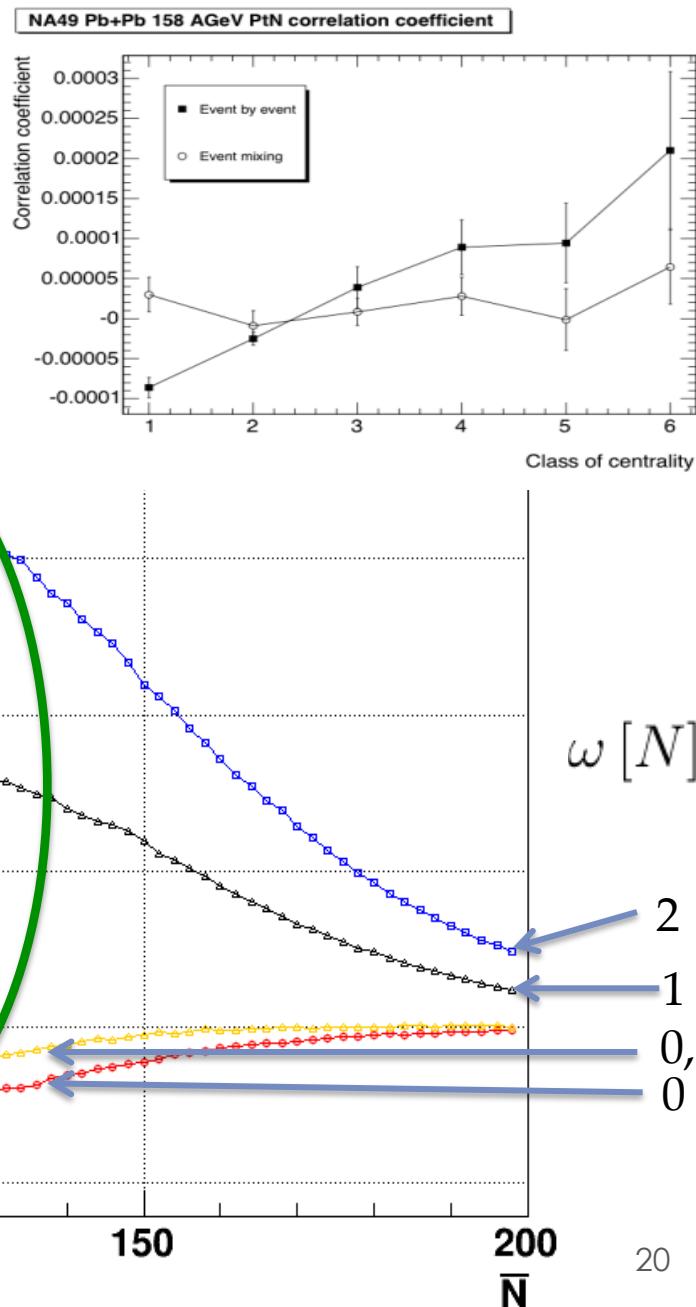
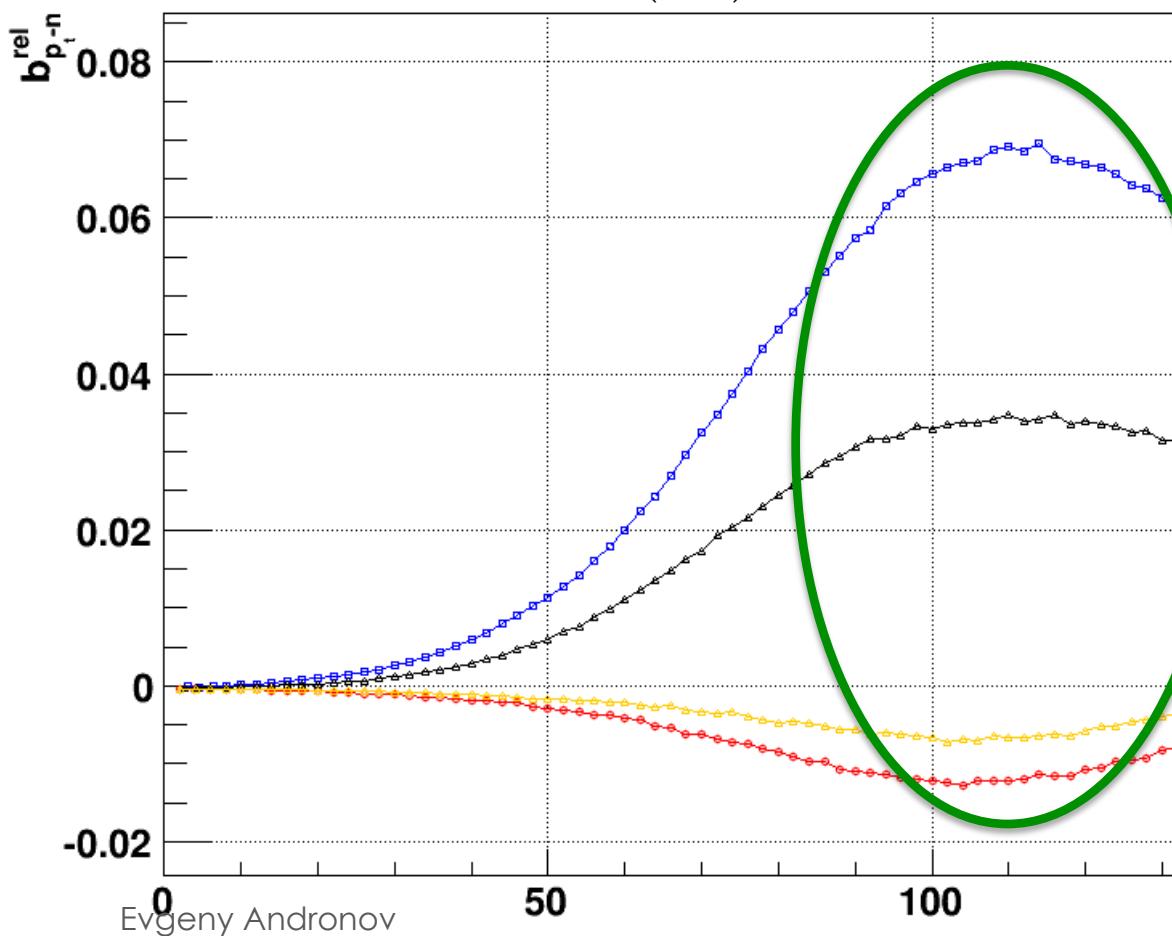
Absence of correlations without fusion effects!



# Long-range pT-n correlation parameter. Monte-Carlo simulations.

Transition from the negative values to positive.

G. Feofilov et al. (NA49 Collaboration), Proc. of Baldin ISHEPP XVII, Volume I, 222 (2005)



# Summary

- Development of the model with 2 types of strings. Predictions for the  $n$ - $n$  and  $pT$ - $n$  correlation coefficients with Monte-Carlo simulations.
- The calculation results predict the non-monotonic behavior of the correlation coefficients with the growth of the mean number of primary strings, i.e. with the increase of the collision centrality.
- Taking into account that fusion parameter depends on the number of primary emitters in the event enables to describe transition from the negative values of the  $pT$ - $n$  correlation coefficients to the positive ones.

Thank you for your  
attention!

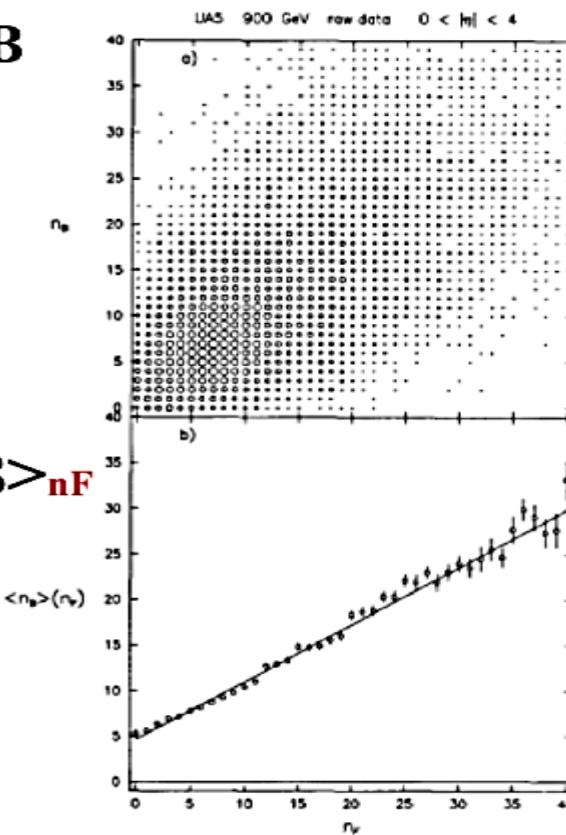
# Back-up

# Experimental studies on LRC

p+(anti-)p, 900 GeV (1988)

Linear regression

**nB**



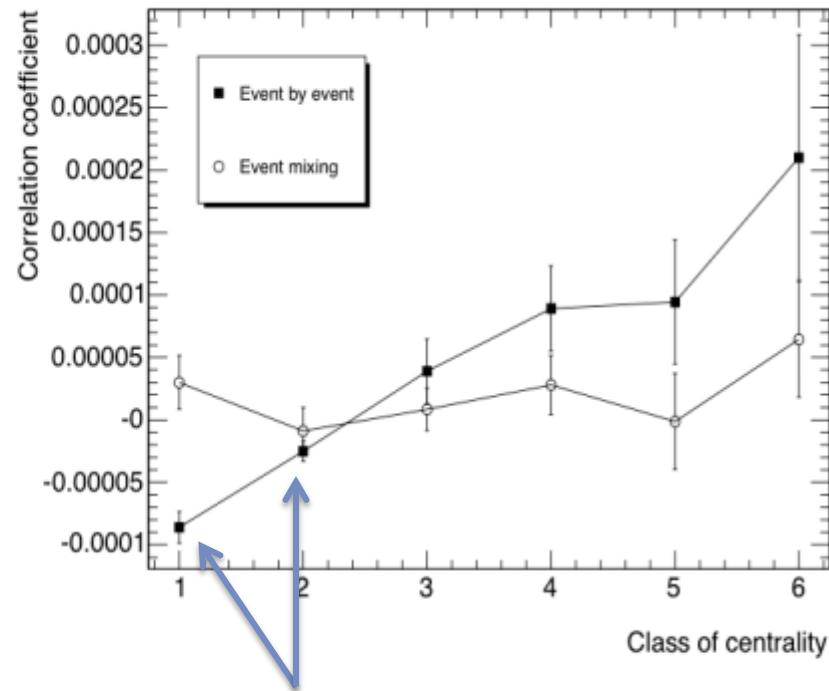
**nF**

**nF**

Pb+Pb, 158 AGeV/c (2005)

**pT-n**

NA49 Pb+Pb 158 AGeV PtN correlation coefficient



Negative pT-n correlations

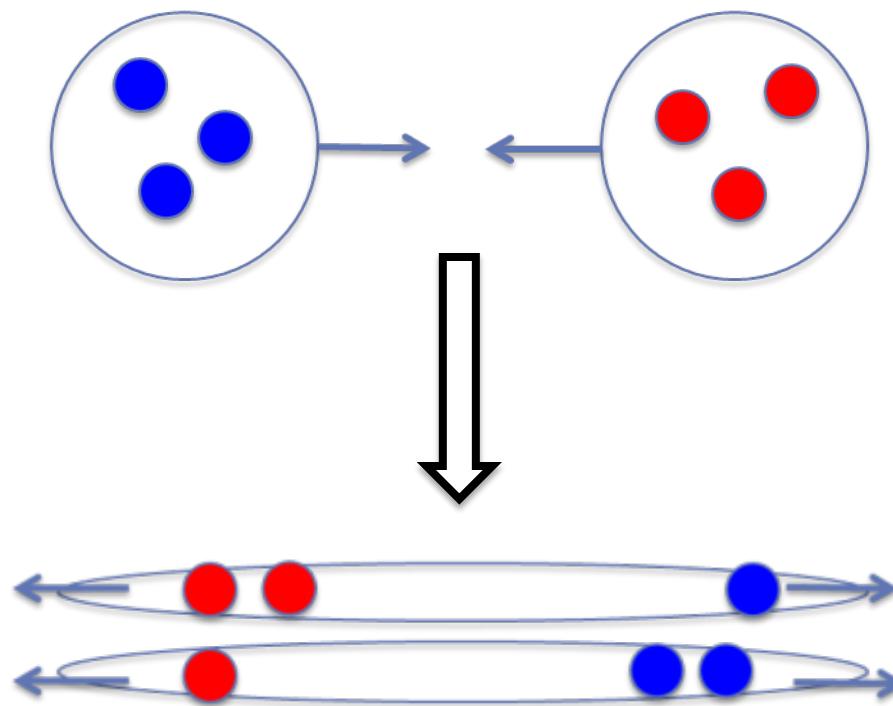
[1] R.E. Ansorge et al. (UA5 Collaboration), Z. Phys., C37-191, (1988).

[2] G. Feofilov et al. (NA49 Collaboration), Proc. of Baldin ISHEPP XVII, Volume I, 222 (2005)

# Two-stage scenario of particles production.

## I stage: strings creation.

p-p, low energies



[1] A.Capella, U.P.Sukhatme, C.I.Tan and J.Tran Thanh Van, Phys. Lett. B81 (1979) 68; Phys. Rep. 236 (1994) 225.

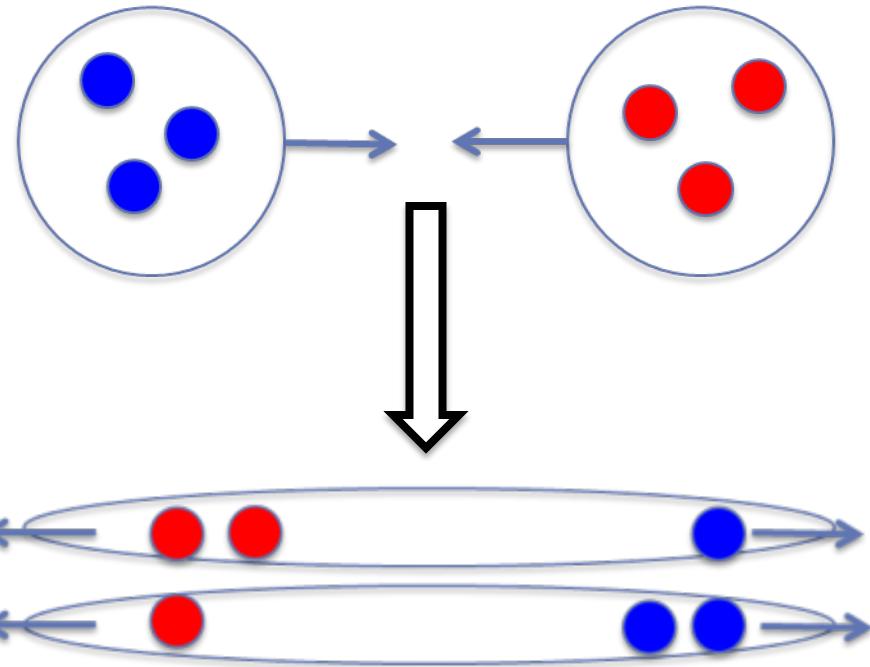
● Evgeny Kaidalov, Phys. Lett., 116B(1982)459

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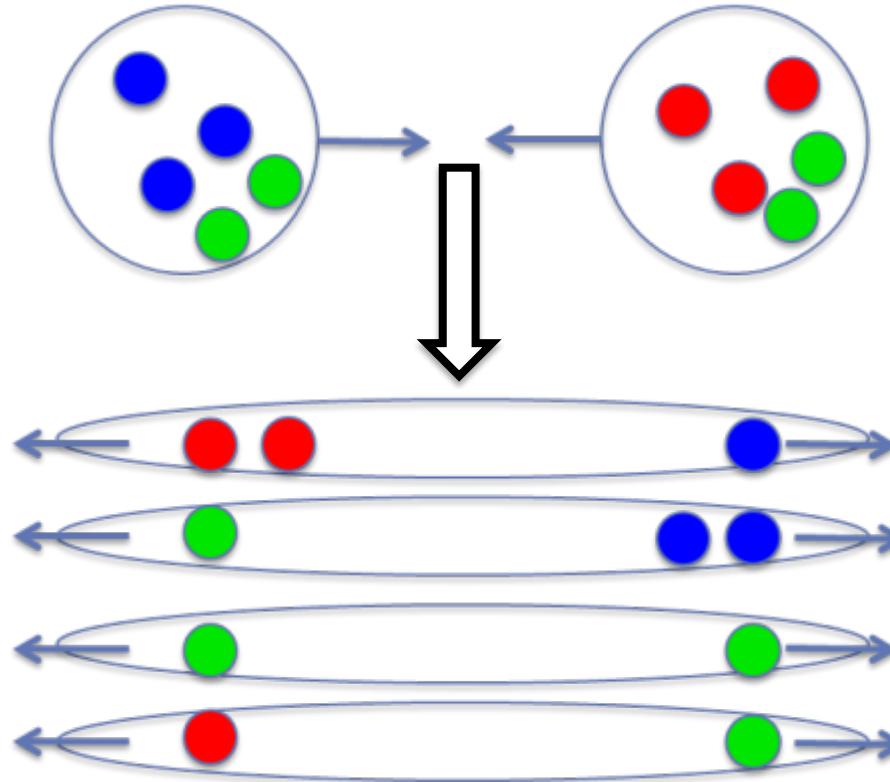
## I stage: strings creation.

●  $q_{sea}$

p-p, **low** energies



p-p, **high** energies



[1] A.Capella, U.P.Sukhatme, C.I.Tan and J.Tran Thanh Van, Phys. Lett. B81 (1979) 68; Phys. Rep. 236 (1994) 225.

● Evgeny [2] A.B.Kaidalov, Phys. Lett., 116B(1982)459

# Model with two types of emitters.

Non-fused

$$\overline{N_1} = \overline{\mu_{F1}} = \overline{\mu_{B1}} = \overline{\mu}$$

$$D_{\mu F1} = D_{\mu B1} = D_\mu$$

$$\langle p_{tB1} \rangle_{1-string} = \overline{k}_1$$

Fused

$$\overline{N_2} = \overline{\mu_{F2}} = \overline{\mu_{B2}} = \sqrt{2} \overline{\mu}$$

$$D_{\mu F2} = D_{\mu B2} = \sqrt{2} D_\mu$$

$$\langle p_{tB2} \rangle_{1-string} = \overline{k}_2 = 2^{1/4} \overline{k}_1$$

## Model with two types of emitters.

$$b_{n-n} = \frac{D_{N_1} \bar{\mu} + 2D_{N_2} \bar{\mu} + 2\sqrt{2} \operatorname{cov}(N_1, N_2)}{\overline{N_1} \omega[\mu] + \overline{N_2} \omega[\mu] + D_{N_1} \bar{\mu} + 2D_{N_2} \bar{\mu} + 2\sqrt{2} \operatorname{cov}(N_1, N_2)}$$



Transition to one-type case

$$b_{n-n} = \frac{D_N \bar{\mu}}{\overline{N} \omega[\mu] + D_N \bar{\mu}}$$

## Model with two types of emitters.

$$\langle p_{tB} \rangle = \bar{k} \sum_{B^{(1)}, B^{(2)}} \frac{B^{(1)} + 2^{1/4} B^{(2)}}{B^{(1)} + B^{(2)}} \sum_{N_1, N_2} \omega(N_1, N_2) P_{N_1}(B^{(1)}) P_{N_2}(B^{(2)})$$

$$\langle p_{tB} n_F \rangle = \bar{\mu k} \sum_{B^{(1)}, B^{(2)}} \frac{B^{(1)} + 2^{1/4} B^{(2)}}{B^{(1)} + B^{(2)}} \sum_{N_1, N_2} (N_1 + \sqrt{2} N_2) \omega(N_1, N_2) P_{N_1}(B^{(1)}) P_{N_2}(B^{(2)})$$

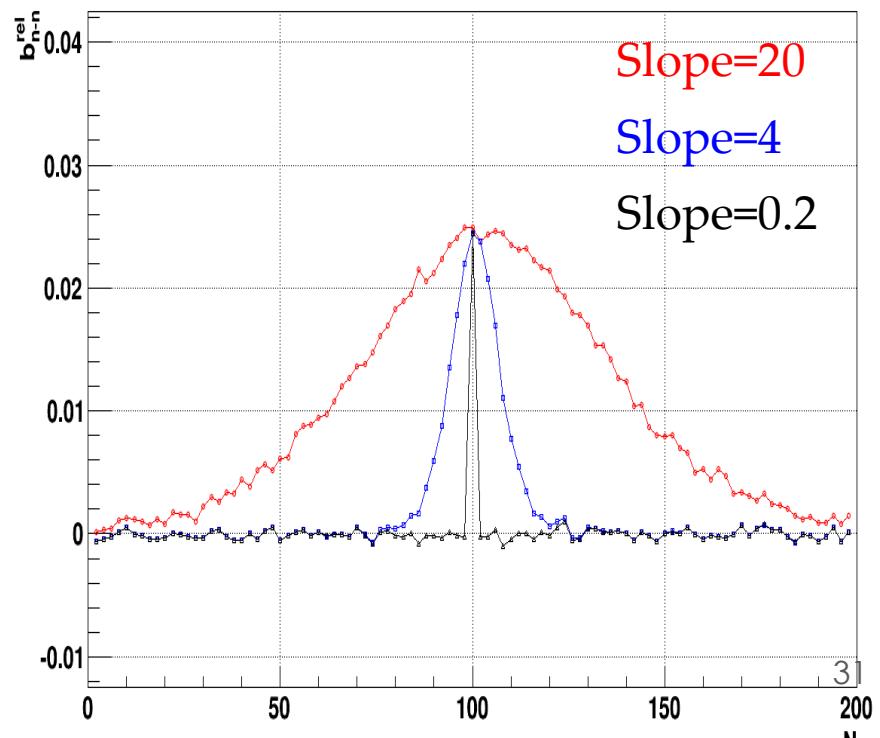
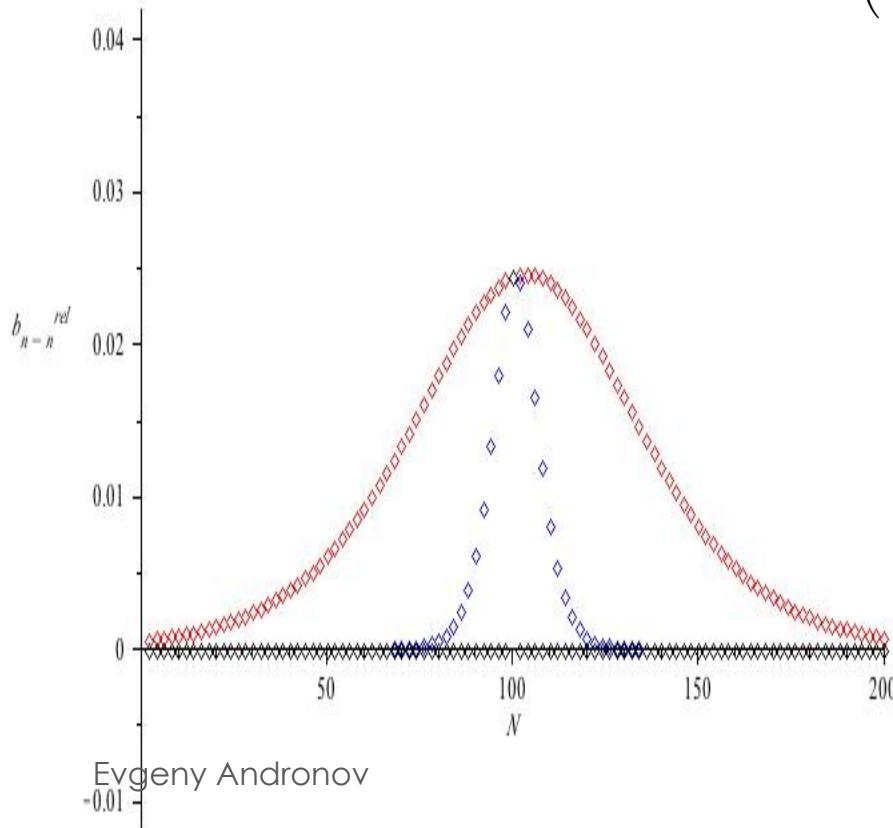
# Model with two types of emitters.

N primary strings in all events

$$b_{n-n} = \frac{\bar{\mu} * N * r(N) * (1 - r(N)) * (3 - 2\sqrt{2})}{D_\mu * N * (1 - r(N) + \sqrt{2}/2 * r(N)) + \bar{\mu} * N * r(N) * (1 - r(N)) * (3 - 2\sqrt{2})}$$

# Comparison of the n-n correlation coefficients without fluctuations in the number of primary strings.

$$b_{n-n}^{rel} = \frac{0.25 \cdot N \cdot r(N) \cdot (1 - r(N)) \cdot (3 - 2\sqrt{2})}{0.25 \cdot N \cdot r(N) \cdot (1 - r(N)) \cdot (3 - 2\sqrt{2}) + 0.5 \cdot N \cdot \left(1 - r(N) + \frac{\sqrt{2}}{2}r(N)\right)}.$$



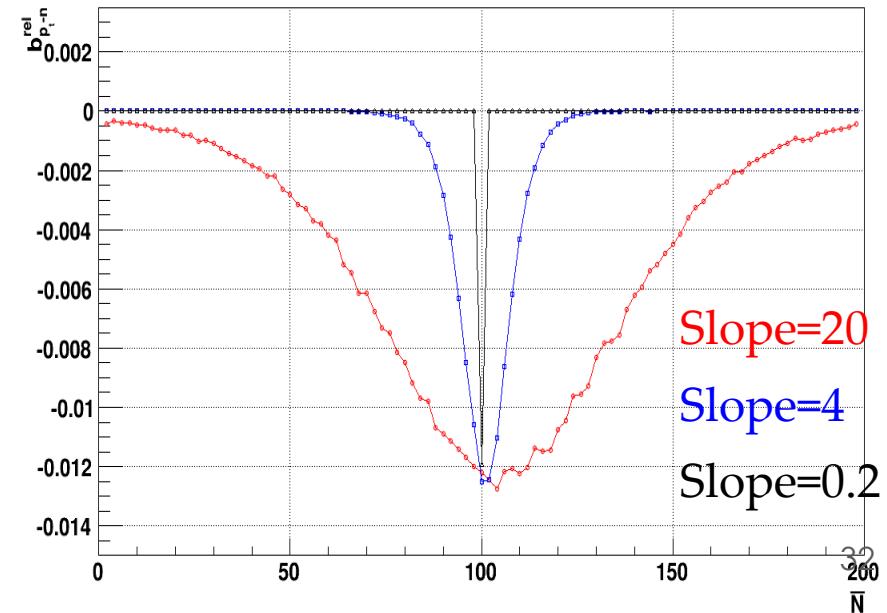
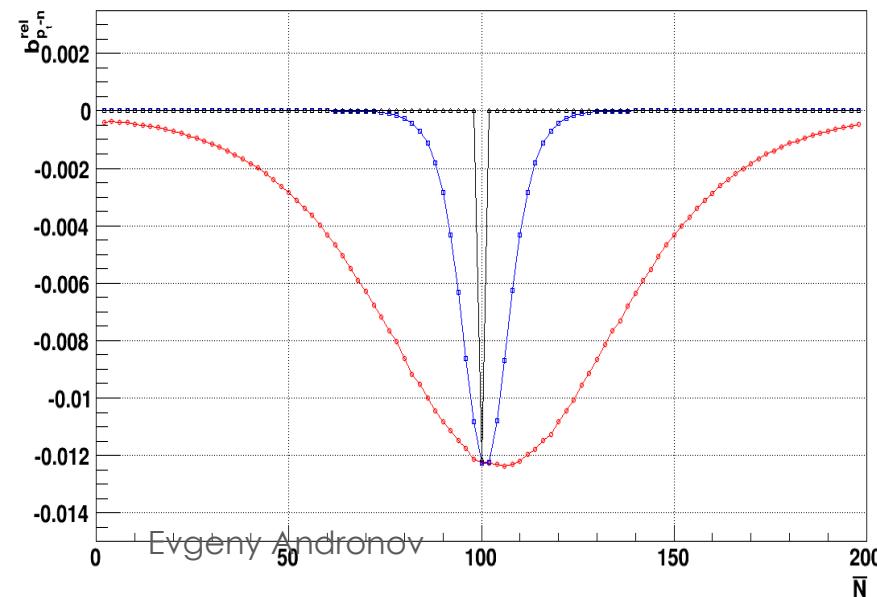
# Comparison of the pT-n correlation coefficients.

$$\langle p_{TB} \rangle = \sum_{B^{(1)}, B^{(2)}} \frac{\bar{k}_1 B^{(1)} + \bar{k}_2 B^{(2)}}{B^{(1)} + B^{(2)}} \sum_{N_1, N_2} q(N_1, N_2) P_{N_1}(B^{(1)}) P_{N_2}(B^{(2)});$$

$$\langle p_{TB} n_F \rangle = \sum_{B^{(1)}, B^{(2)}} \frac{\bar{k}_1 B^{(1)} + \bar{k}_2 B^{(2)}}{B^{(1)} + B^{(2)}} \sum_{N_1, N_2} q(N_1, N_2) (N_1 \bar{\mu}_{F^{(1)}} + N_2 \bar{\mu}_{F^{(2)}}) P_{N_1}(B^{(1)}) P_{N_2}(B^{(2)}).$$

Approximation:  $B^{(1)} + B^{(2)} \approx N_1 \bar{\mu}_{B1} + N_2 \bar{\mu}_{B2}$

Without approximation



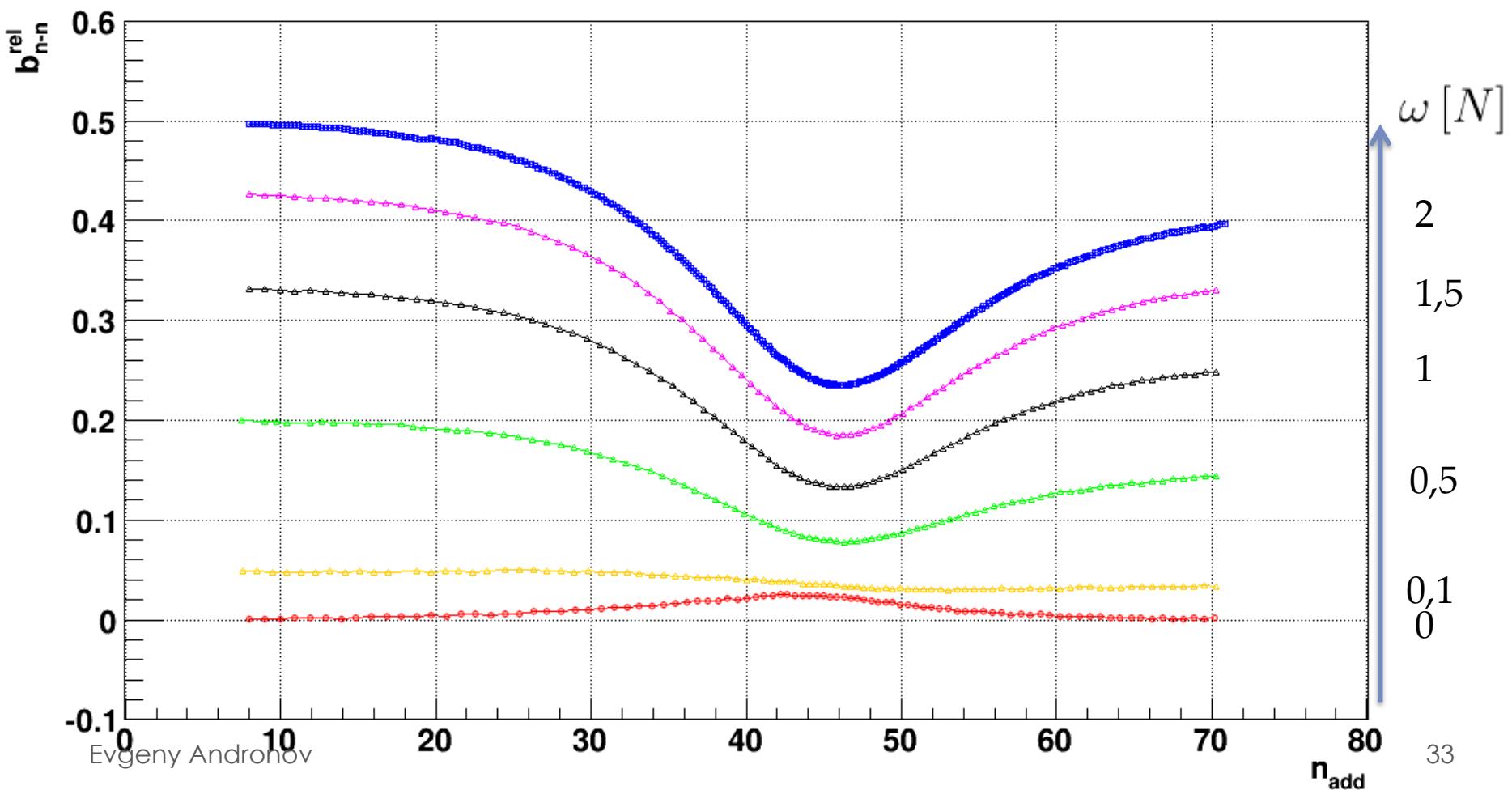
$$\bar{\mu} = D_\mu = 0.5$$

## Long-range n-n correlations

$$r(N) = \frac{1}{\frac{N-shift}{slope}} = \frac{1}{1+e^{-\frac{N-shift}{slope}}}$$

Shift=100  
Slope=20

$\omega [N] = \frac{D_N}{\bar{N}}$  - the scaled variance of the number of primary strings



$$\bar{\mu} = D_\mu = 0.5$$

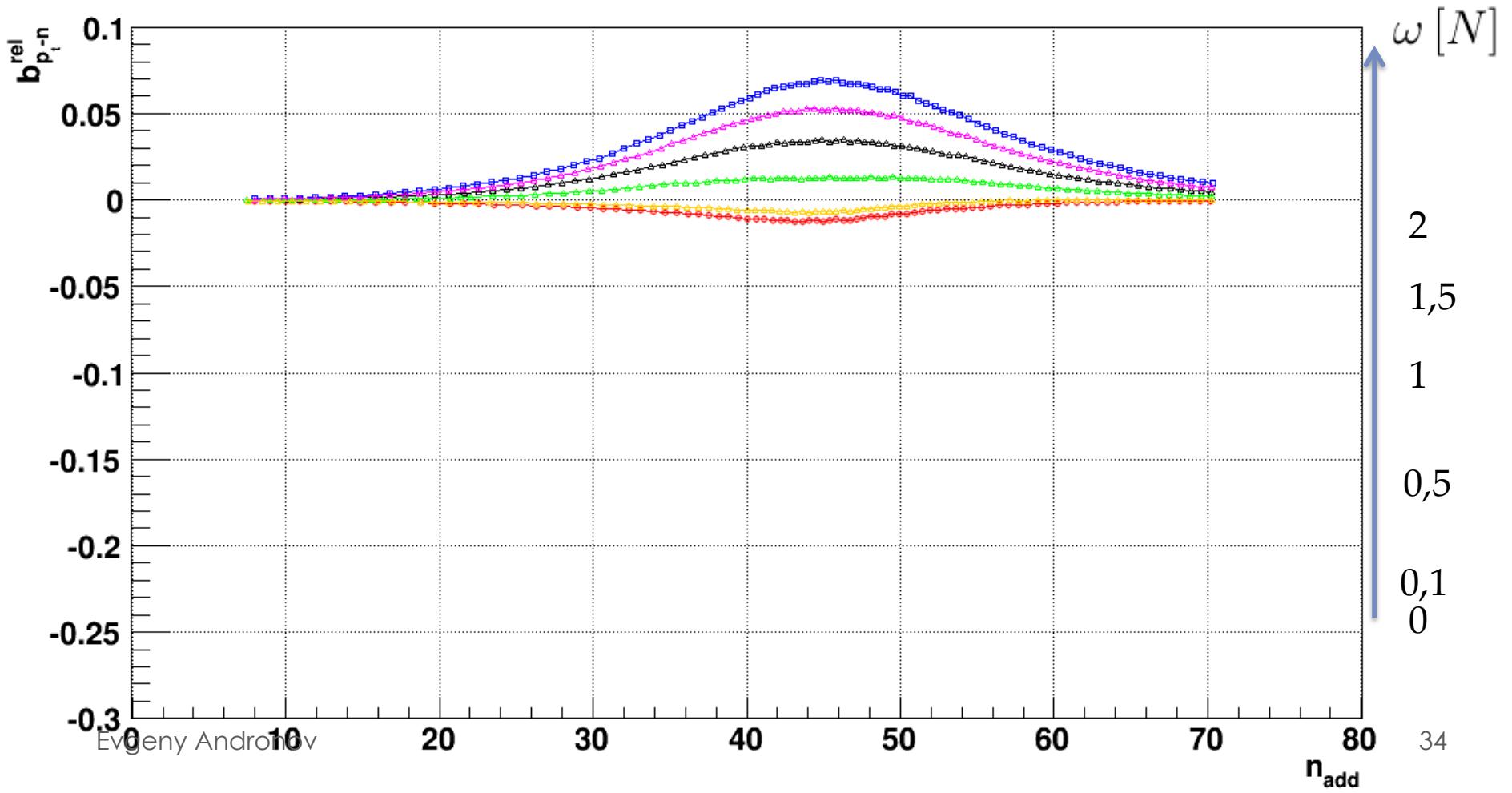
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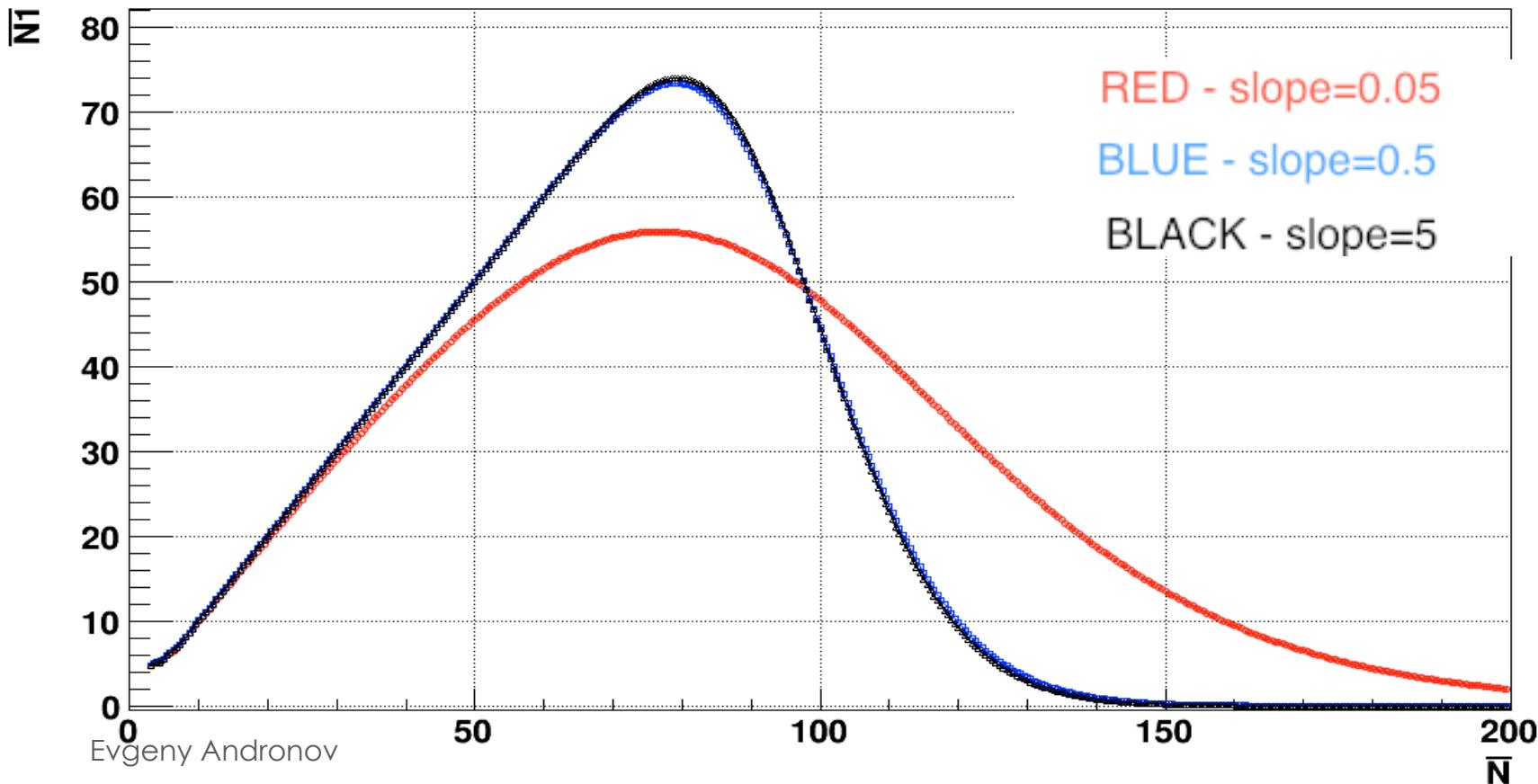
$$\bar{\mu} = D_\mu = 0.5$$

# Monte-Carlo generator

Fluctuating number of strings  $N$ ,  $w[N]=2$

Shift=100

MC

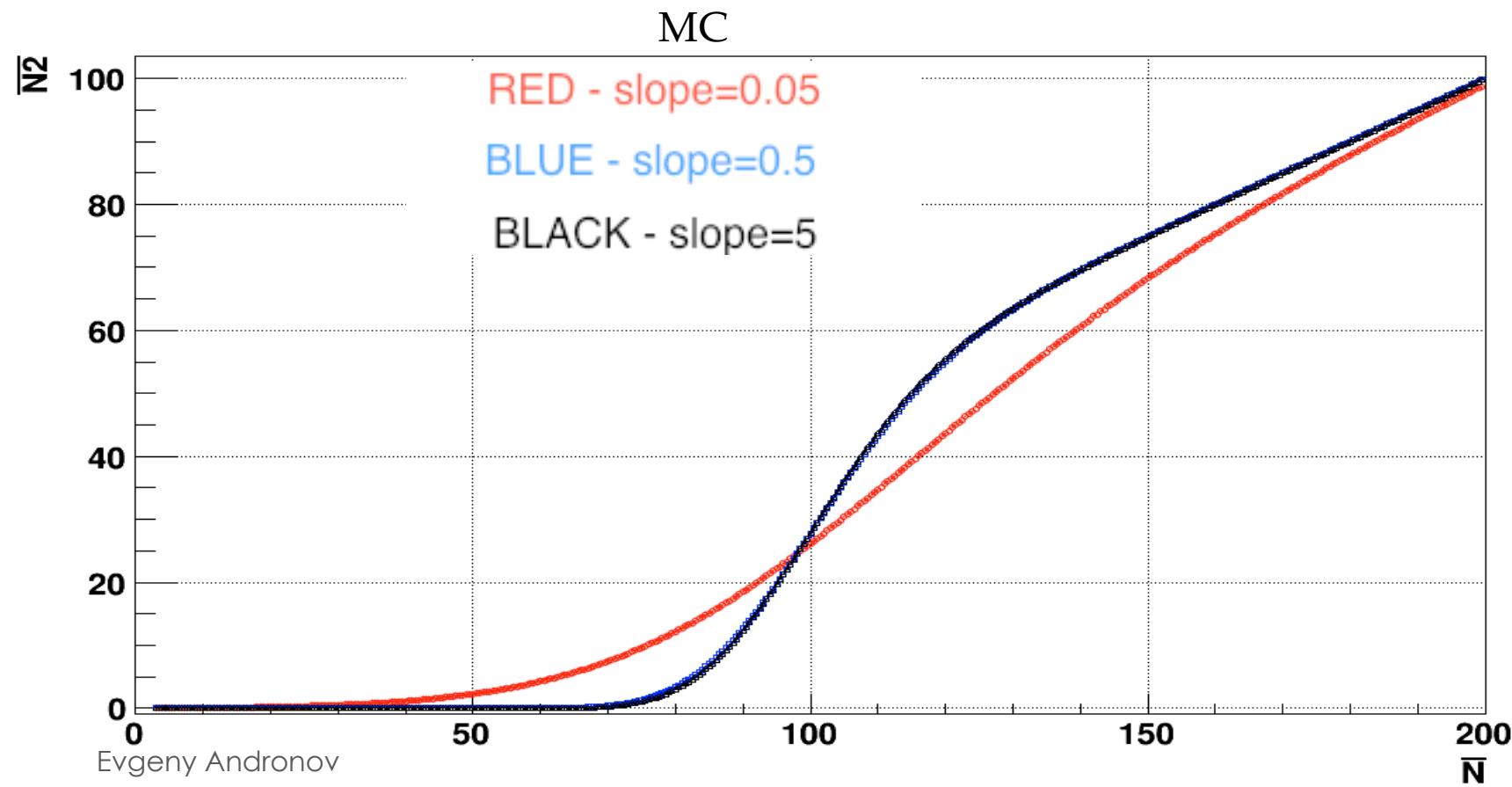


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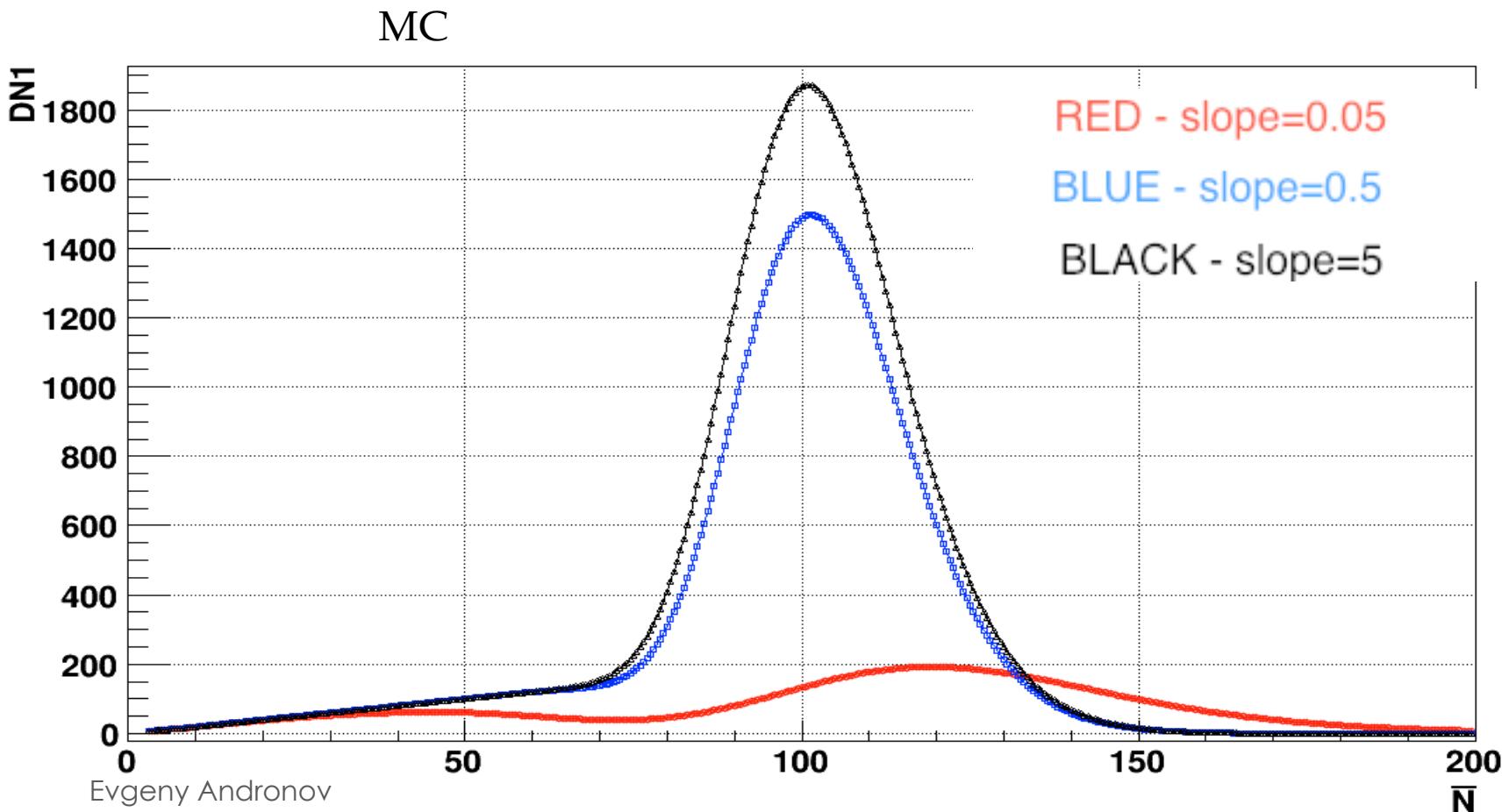


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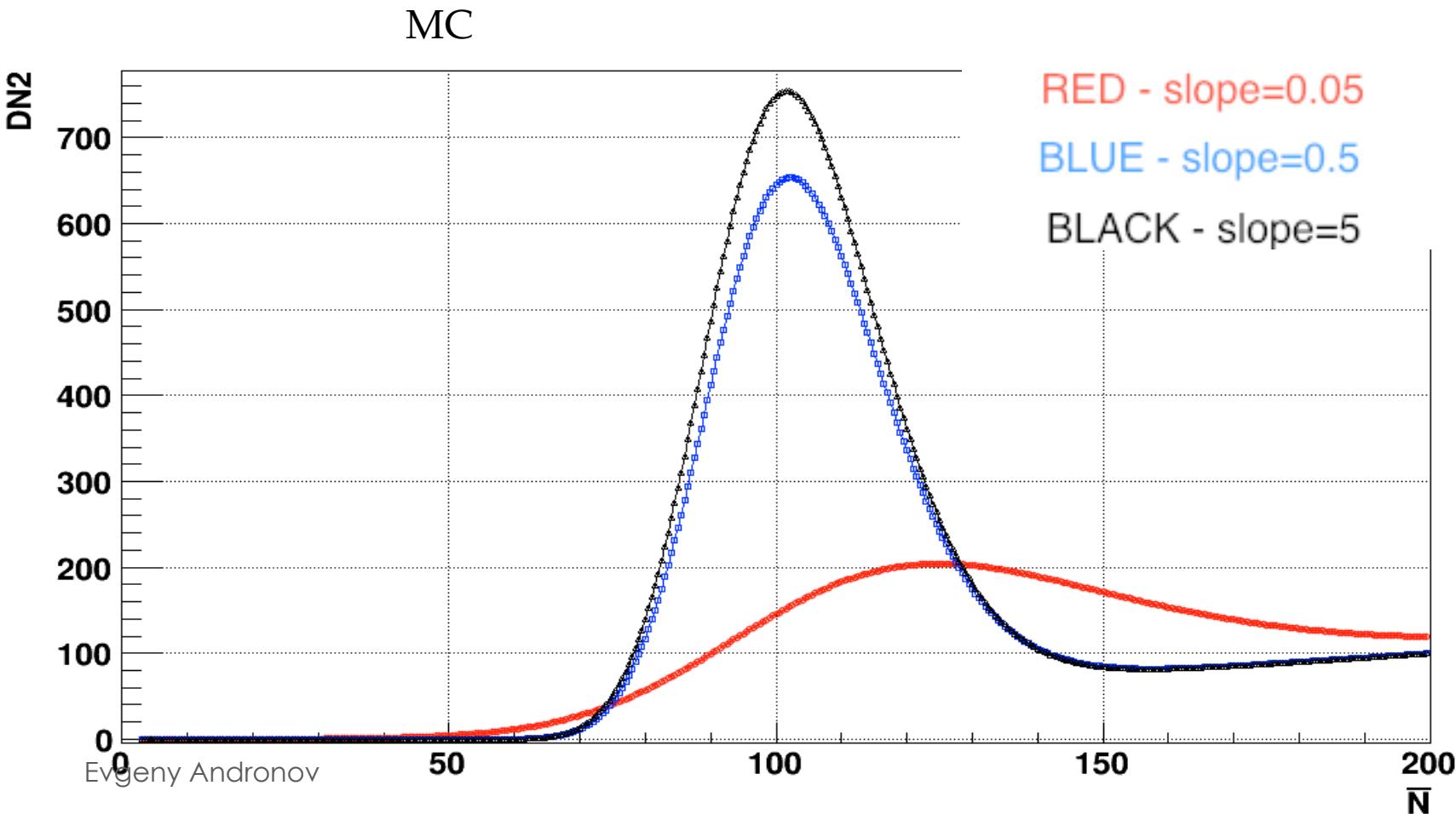


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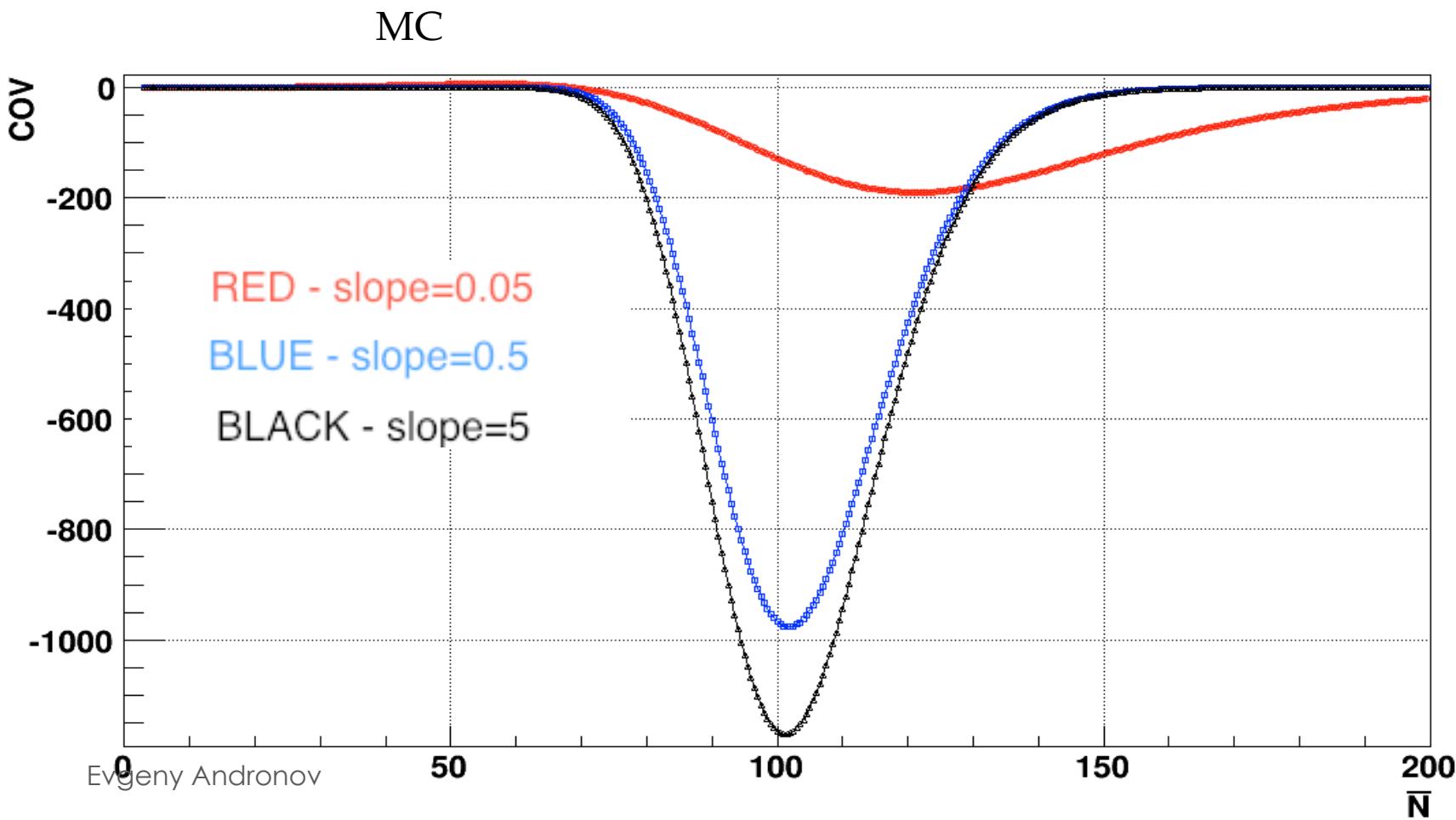


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Shift=100

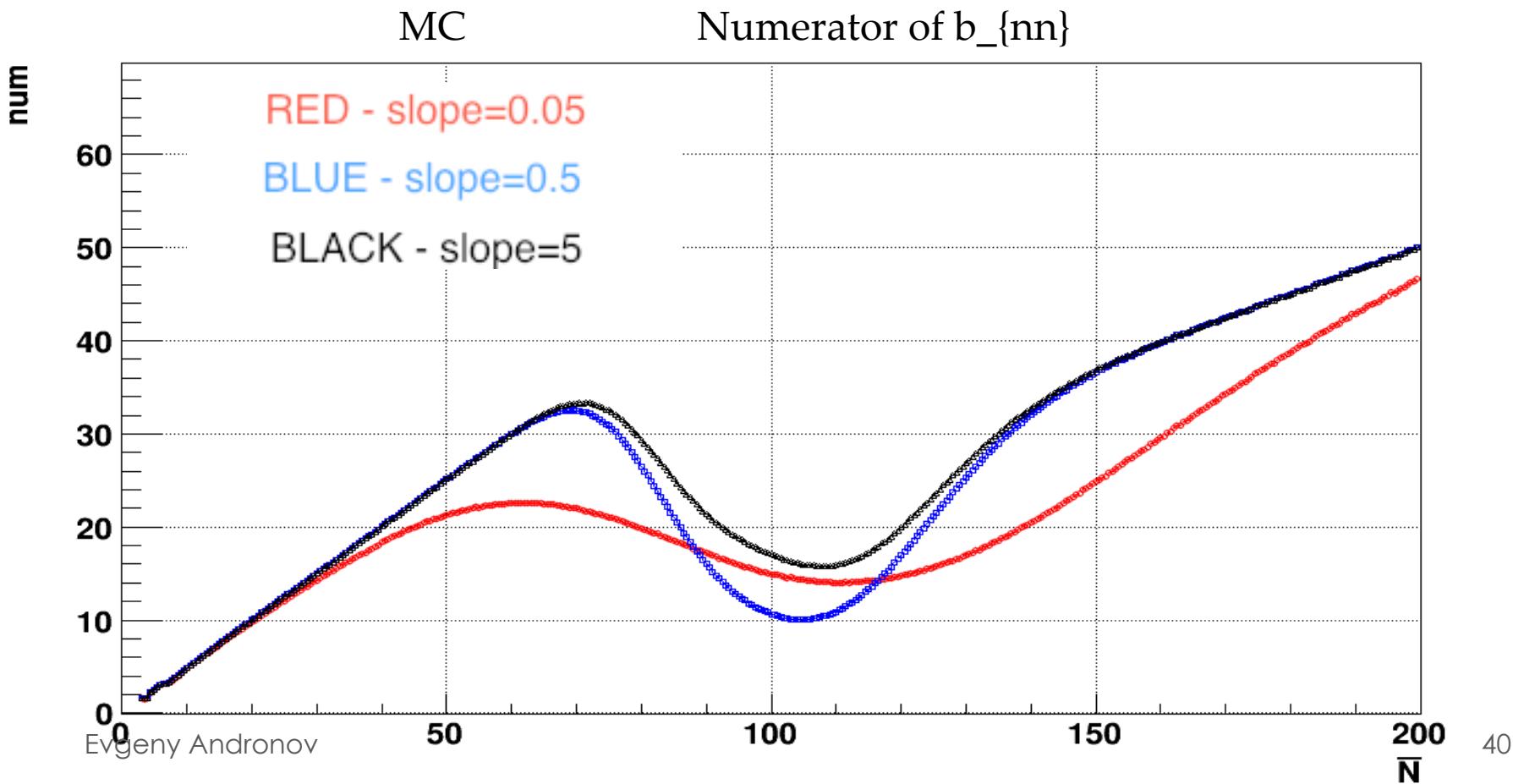


$$\bar{\mu} = D_\mu = 0.5$$

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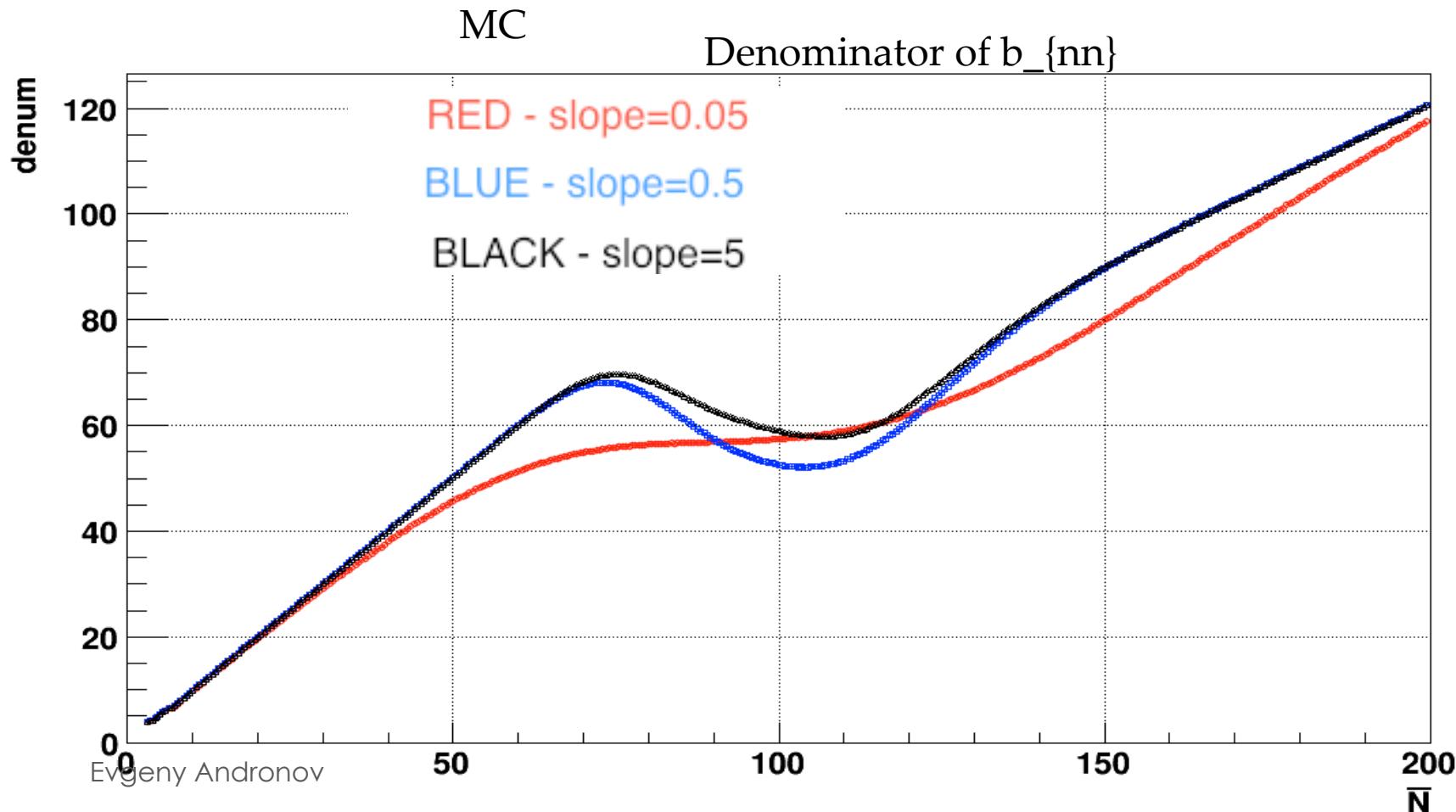


$$\bar{\mu} = D_\mu = 0.5$$

# Monte-Carlo generator

Fluctuating number of strings  $N$ ,  $w[N]=2$

Shift=100



RED - slope=0.05

BLUE - slope=0.5

BLACK - slope=5

MC

$\bar{\mu} = D_\mu = 0.5$

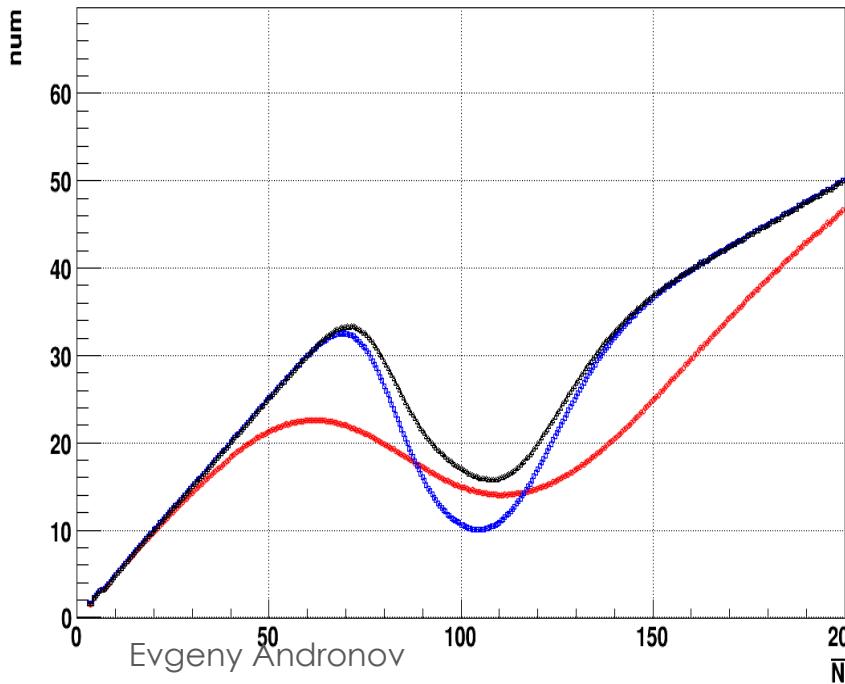
# Monte-Carlo generator

Fluctuating number of strings N, w[N]=2

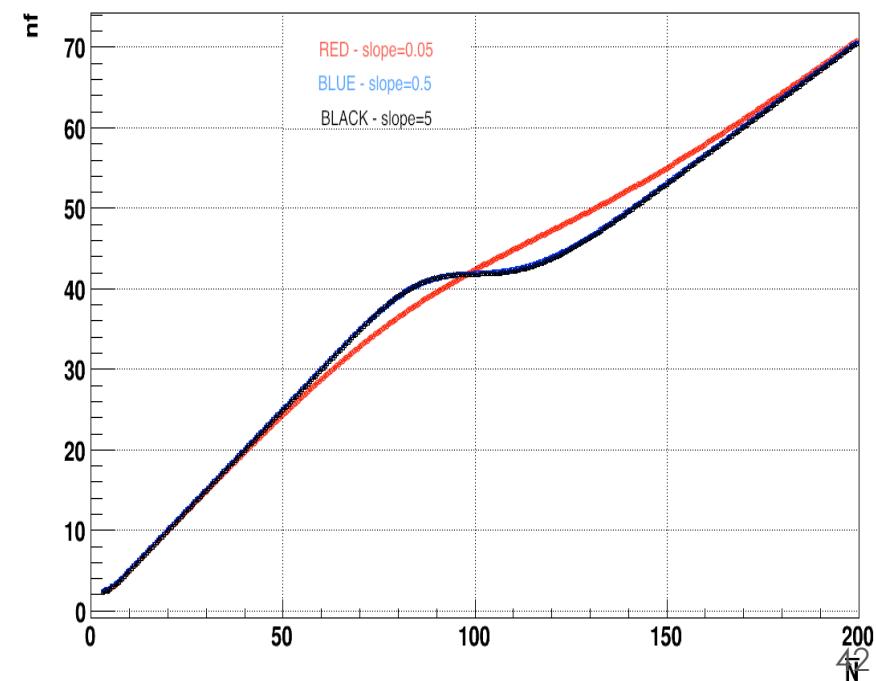
Shift=100

$$b_{n-n} = \frac{num}{num + \langle n_F \rangle}$$

Numerator of  $b_{\{nn\}}$



$\langle nF \rangle$



MC

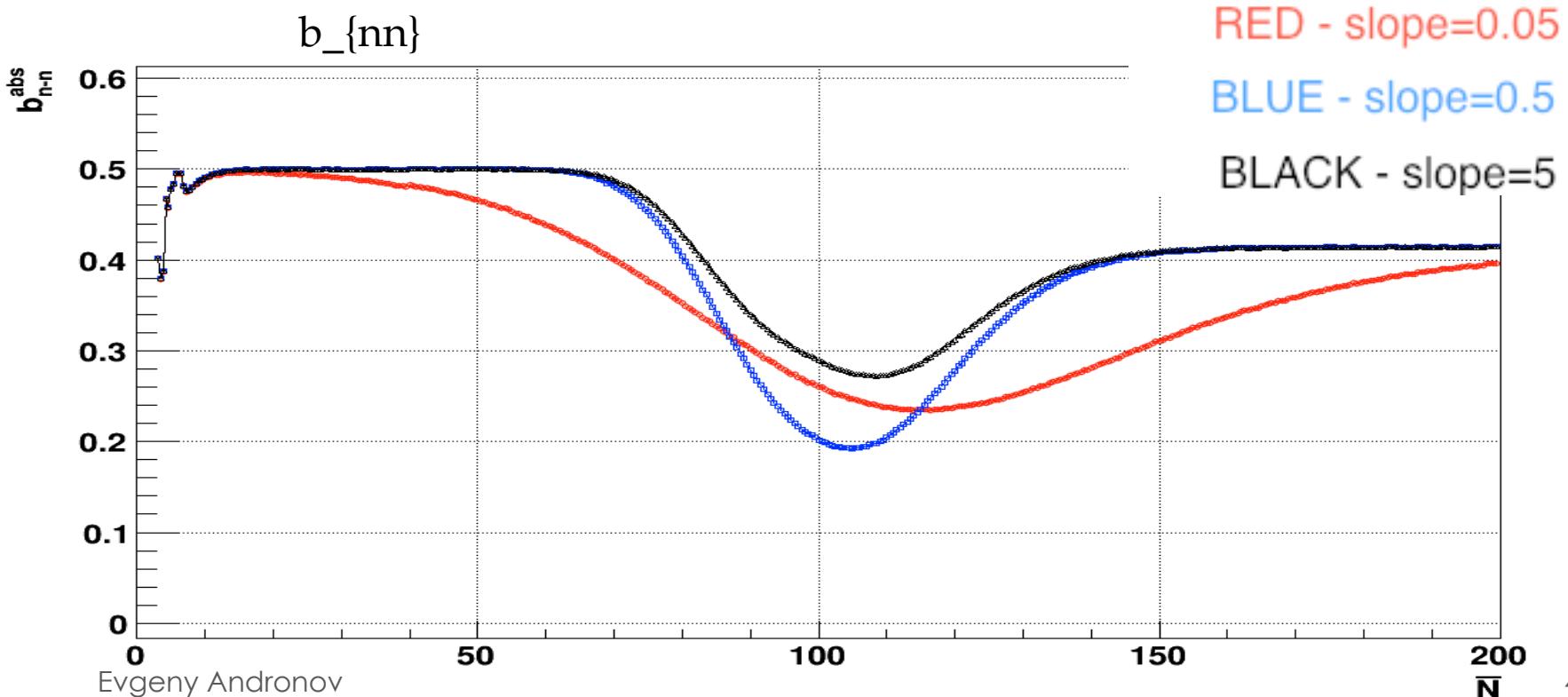
$$\bar{\mu} = D_\mu = 0.5$$

# Monte-Carlo generator

Fluctuating number of strings N, w[N]=2

$$b_{n-n} = \frac{num}{num + \langle n_F \rangle}$$

Shift=100



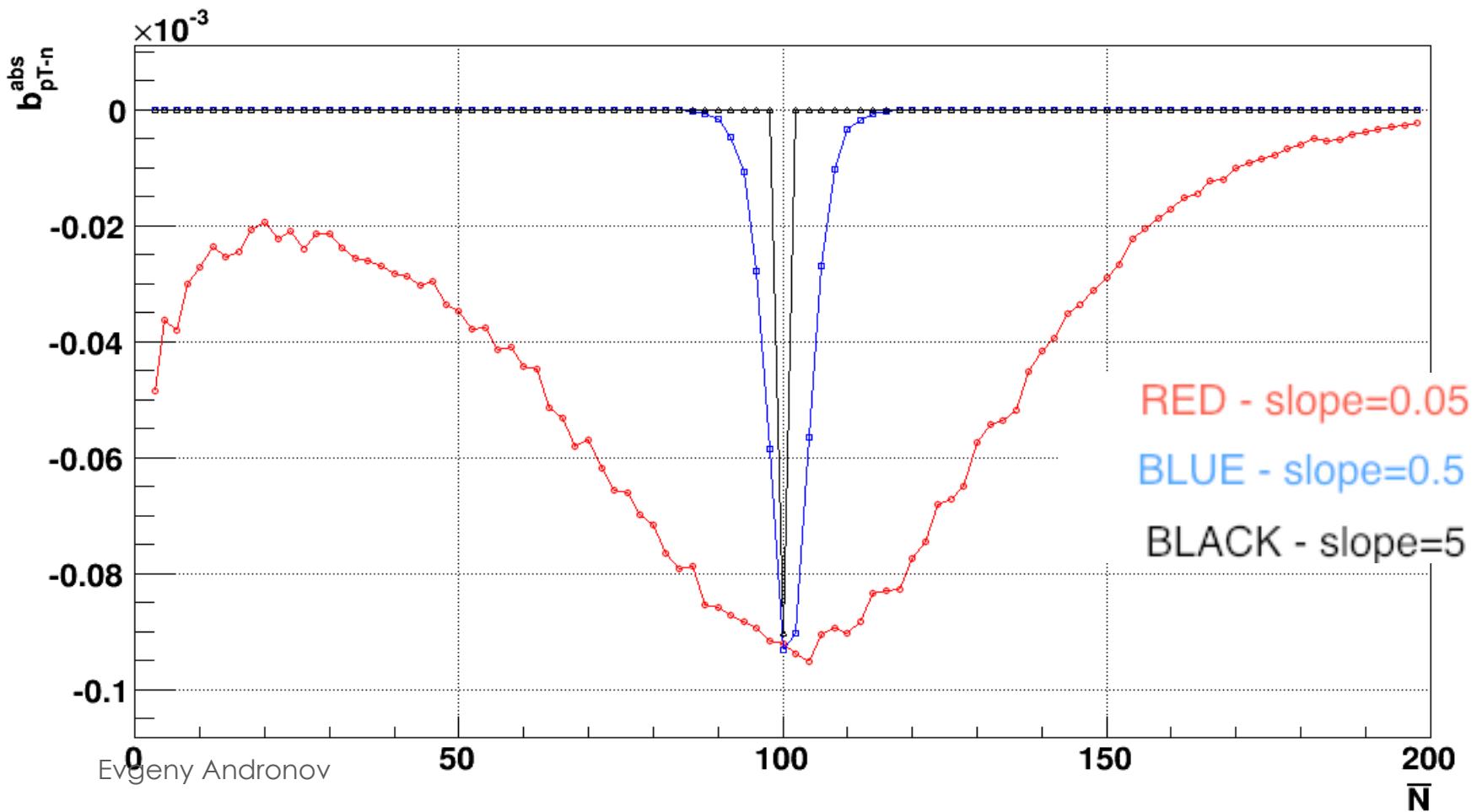
MC

$$\bar{\mu} = D_\mu = 0.5$$

# Monte-Carlo generator

Nonfluctuating number of strings N

Shift=100



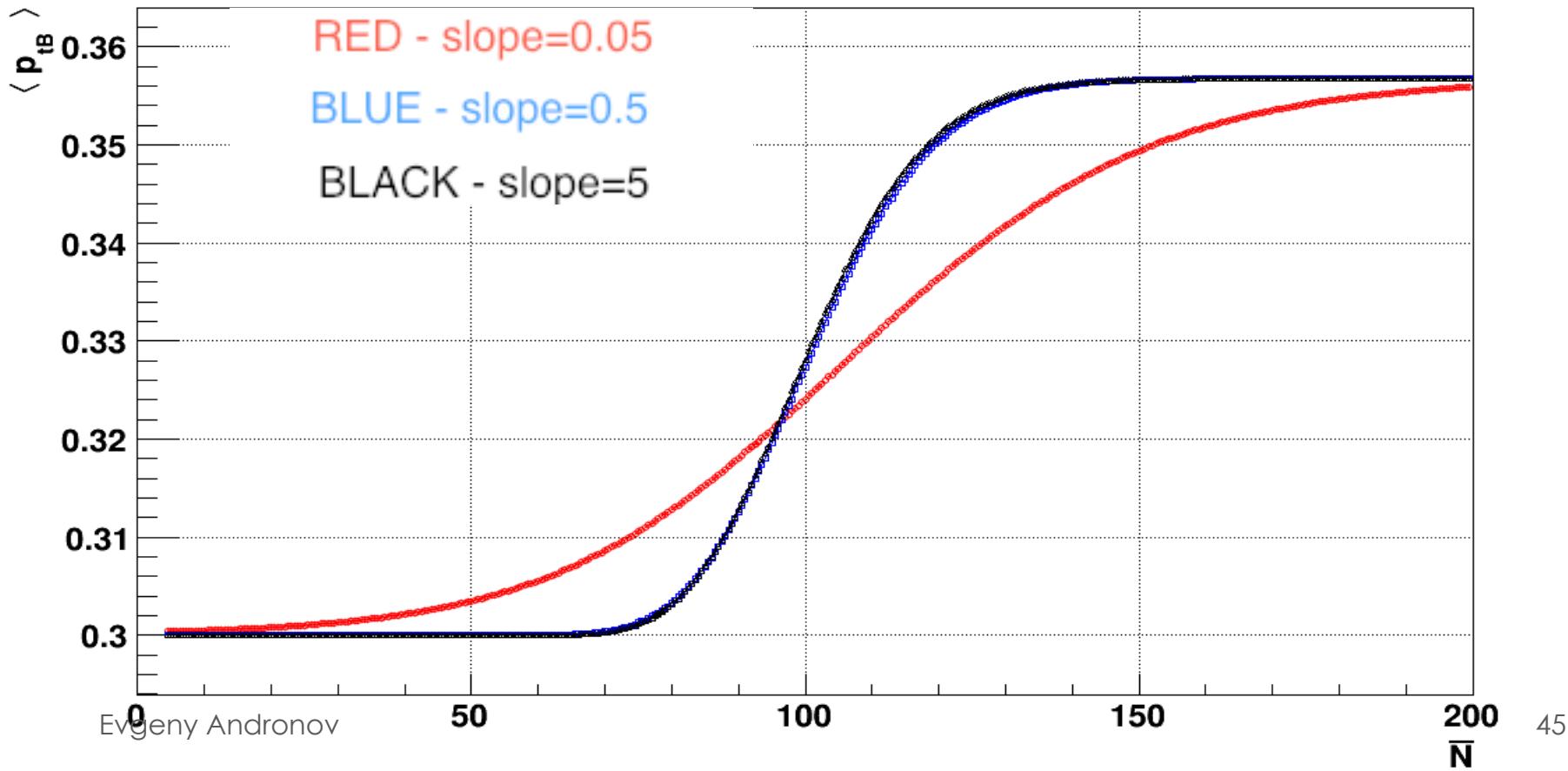
MC

$\bar{\mu} = D_\mu = 0.5$

# Monte-Carlo generator

Fluctuating number of strings N,  $w[N]=2$

Shift=100



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# Monte-Carlo generator

Fluctuating number of strings N, w[N]=2

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