

Quark Confinement and the Hadron Spectrum XI

Transverse Momentum Dependence of Spectra of Cumulative Particles Produced from Droplets of Dense Nuclear Matter

Vladimir Vechernin
Saint-Petersburg State University
v-vechernin@ya.ru

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Cumulative Particle Production

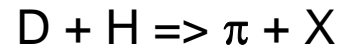
Production of particles from nuclei in a region, kinematically forbidden for reactions with free nucleons.

Cumulative Pion Production

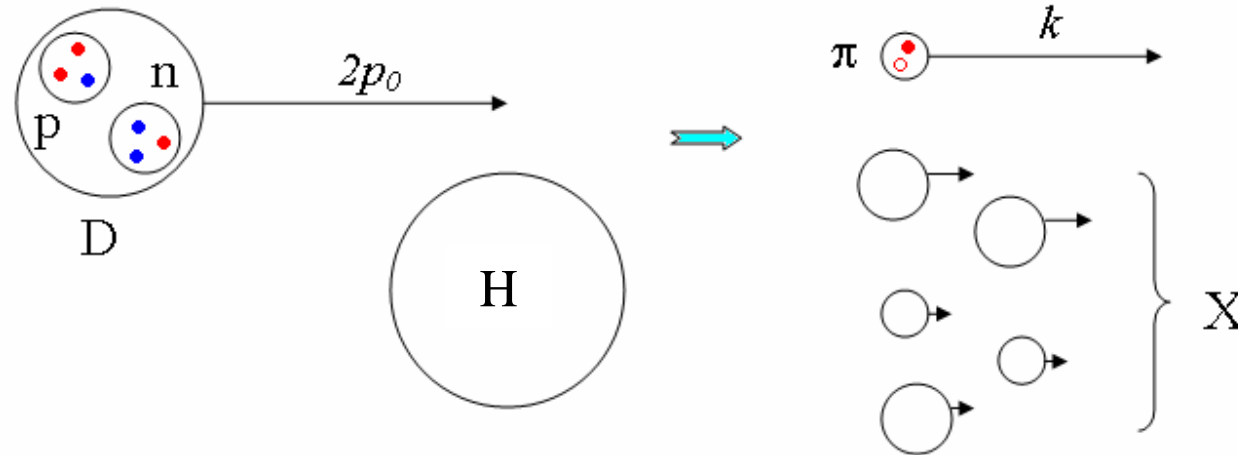
Fragmentation of deuterons, D, on some target, H.

Baldin A.M. et al., Yad.Fiz.18 (1973) 79

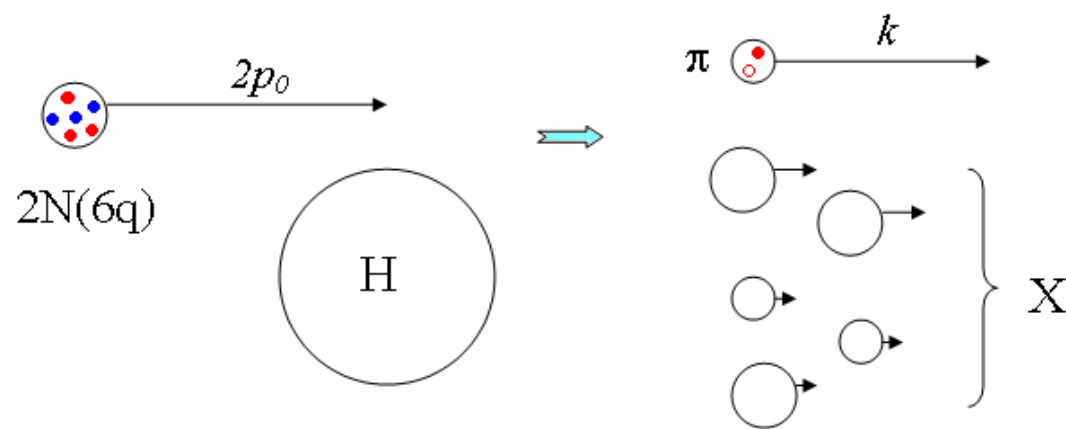
Nuclotron@Dubna ($p_0=5 \text{ GeV}/c$)



$p_0 \gg m_N$: $p_0 < k < 2p_0$ - cumulative pions



Flucton - droplet of dense cold nuclear matter (2N fluctuation – 6 quark state)
Blokhintsev D.I., JETP 33 (1957) 1295



Theoretical description near threshold: $k \rightarrow 2p_0$, $x = k/p_0 \rightarrow 2$

$1 < x < 2$ - the cumulative region ($1 < x < f$ - for the fN flucton)

Brodsky S.J., Chertok B.T. Phys.Rev. D14 (1976) 3003; Phys.Rev.Lett. 37 (1976) 269
Schmidt I.A., Blankenbecler R. Phys.Rev. D15 (1977) 3321

Quark counting rules: $\sim \Delta^{2n-3}$

n - the number of constituents, $n = 6$

Δ - the deviation from the threshold, $\Delta = 2 - x$, $\Delta \ll 1$

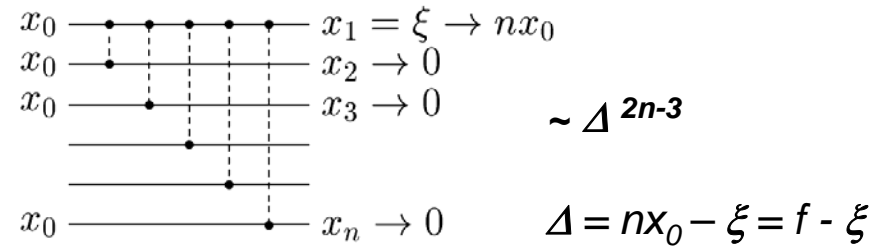
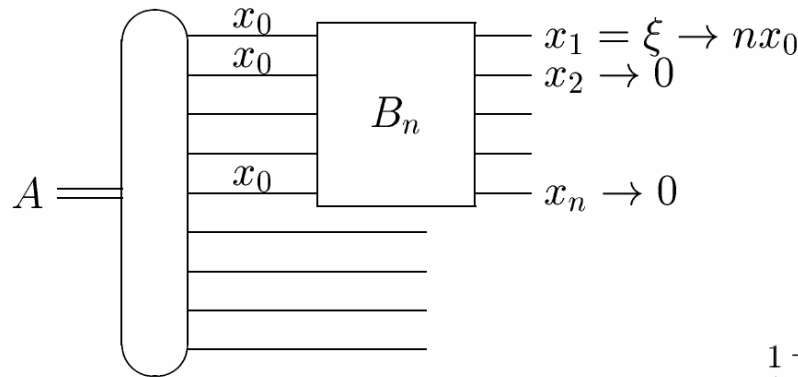
Description of the hadron asymptotics at $x \rightarrow 1$ by the intrinsic diagrams in QCD in light-cone gauge

Brodsky S.J., Hoyer P., Mueller A., Tang W.-K., Nucl.Phys. B369 (1992) 519

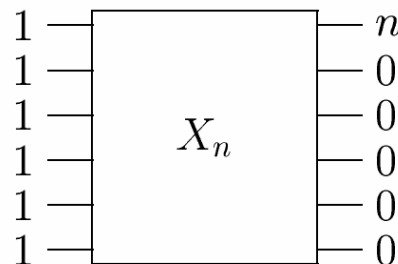
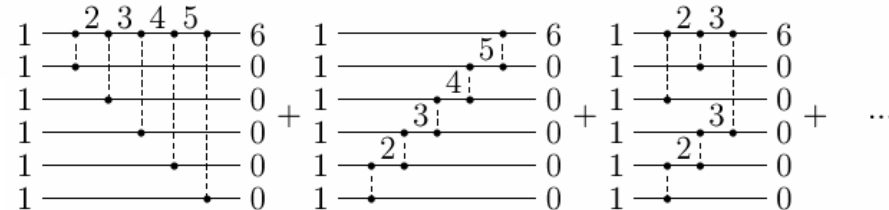
Description of the flucton asymptotic at $x \rightarrow f$,

f - the number of nucleons in flucton, n - the number of quarks in flucton, $x_0 = f/n$.

M.A. Braun, V.V. Vechernin, Nucl.Phys. B427 (1994) 614.



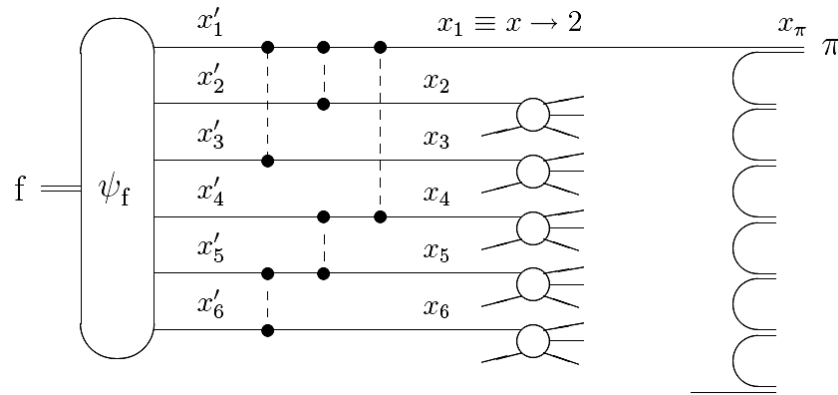
$p=n-1: \Delta^{2n-3} = \Delta^{2p-1}$



$$= \sum_{k=1}^{n-1} C_{n-2}^{k-1} \left[\begin{array}{c} \text{Box } X_k \text{ with } k \text{ lines on the right} \\ \text{Box } X_{n-k} \text{ with } n-k \text{ lines on the right} \end{array} \right]$$

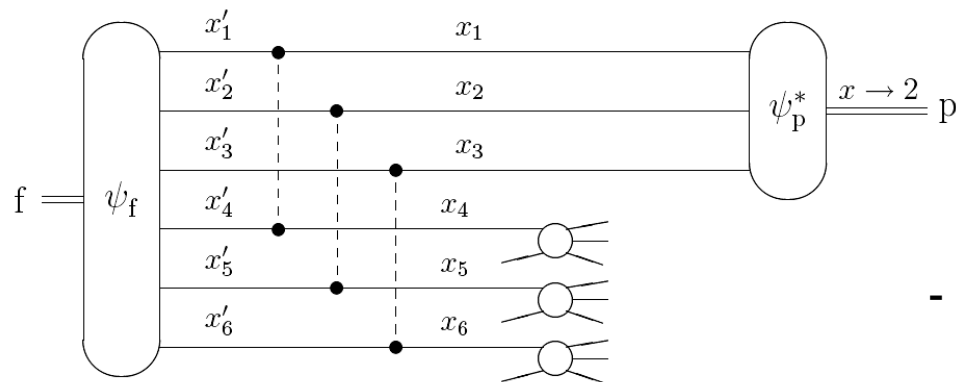
- the recurrence relation

Coherent Quark Coalescence and Production of Cumulative Protons



- the cumulative pion production

k_T – dependence:
M.A. Braun, V.V. Vechernin,
*Phys.Atom.Nucl. **63**, 1831 (2000)*



- the cumulative proton production

coherent quark coalescence mechanism:

M.A. Braun, V.V. Vechernin,
*Nucl.Phys. **B92**, 156 (2001);*
*Theor.Math.Phys **139**, 766 (2004)*

$$p_1=p_2=p_3=1$$

$$W_j: j=1,2,3.$$

$$n=n_1+n_2+n_3$$

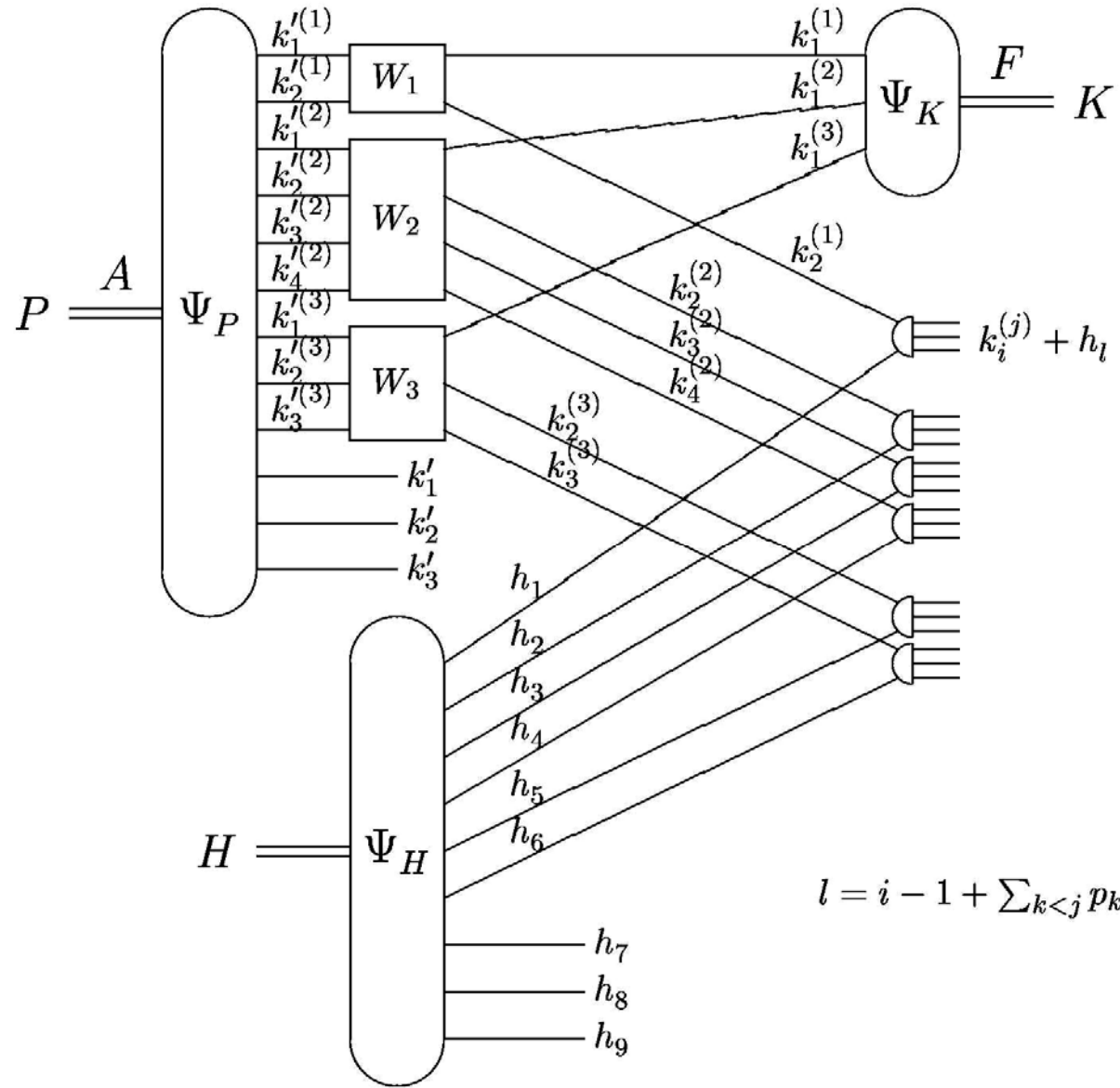
$$p_1=n_1-1$$

$$p_2=n_2-1$$

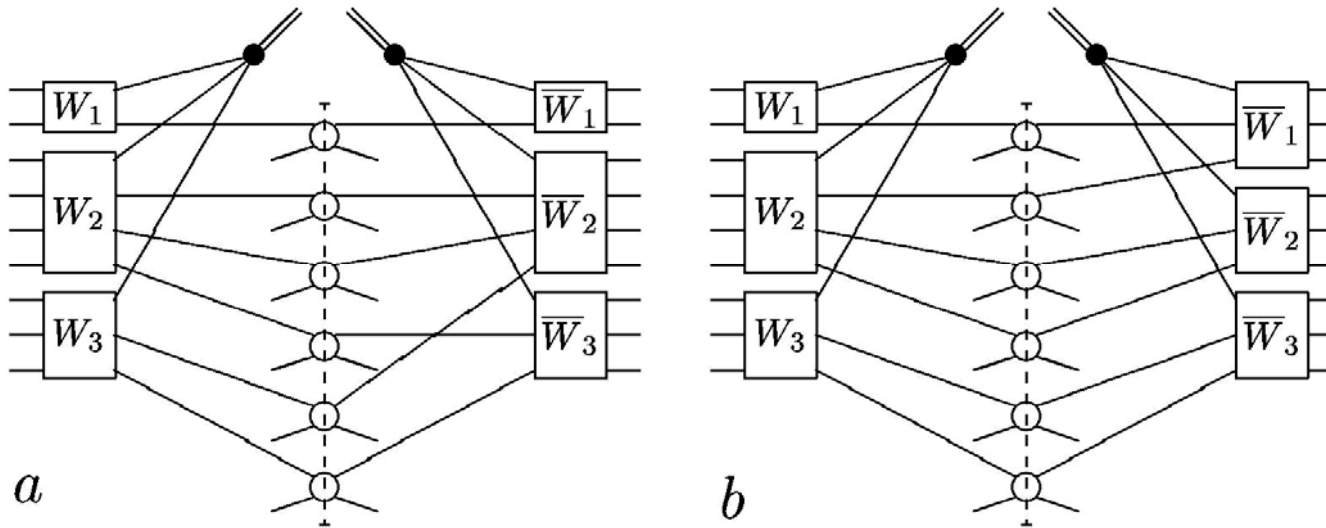
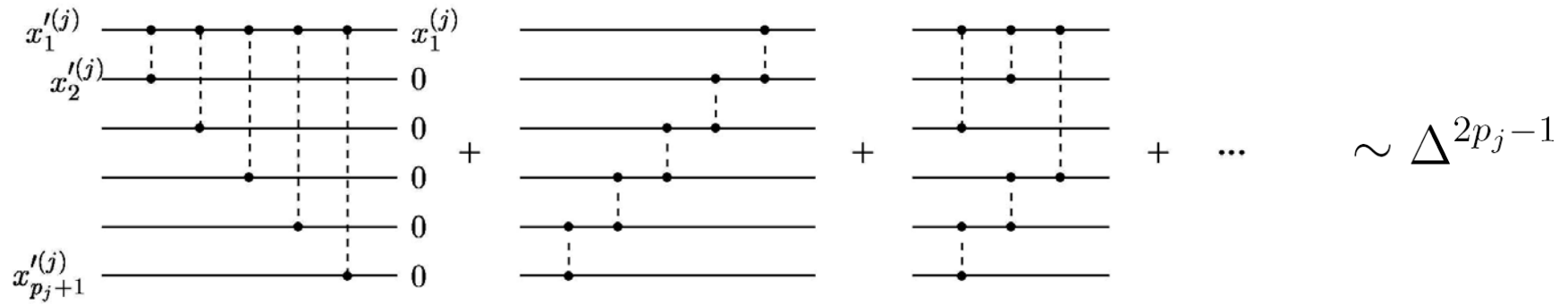
$$p_3=n_3-1$$

$$p=p_1+p_2+p_3=1+3+2=6$$

$$n=p+3=9$$

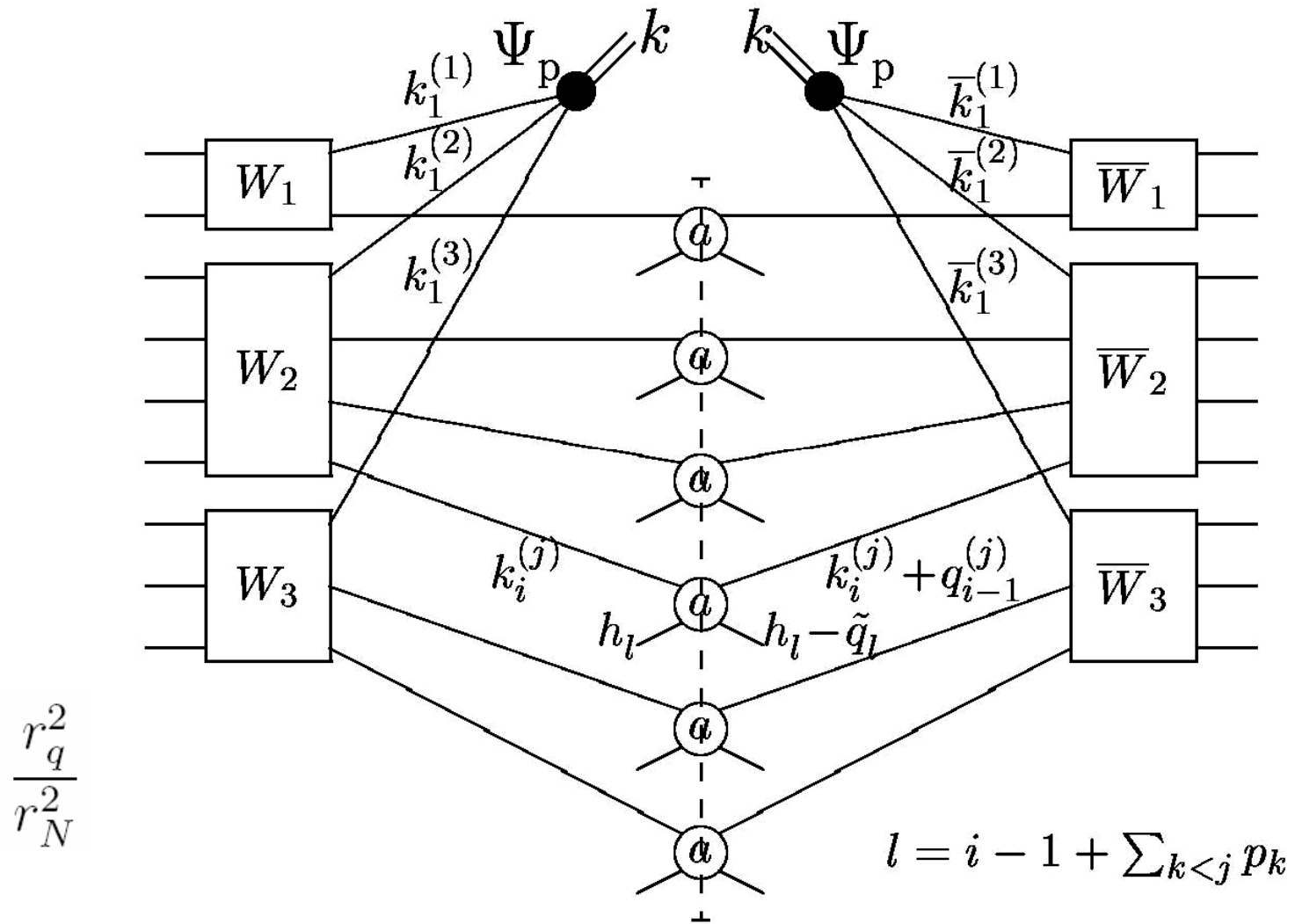


Contributions to the blobs W_j :



The examples of two types of non-diagonal contributions to the cross section of cumulative proton production:

a – all $p_j = \bar{p}_j$, b – some $p_j \neq \bar{p}_j$



The diagonal contribution to the cross section of cumulative proton production.
 Note the presence of the interference effects also in this case!

$$\sigma_{pion}(x, k_{\perp}; p) = C(p) (x_{frag} - x)^{2p-1} f_p \left(\frac{k_{\perp}}{m} \right)$$

$$x < x_{frag}(p) = 1/3 + p/3$$

$\rho=n-1$

M.A. Braun, V.V. Vechernin, Phys.Atom.Nucl. 63, 1831 (2000)

$$\sigma_{prot}(x, k_{\perp}; p_1, p_2, p_3) = C(p_1, p_2, p_3) (x_{coal} - x)^{2p-1} f_{p_1} \left(\frac{k_{\perp}}{3m} \right) f_{p_2} \left(\frac{k_{\perp}}{3m} \right) f_{p_3} \left(\frac{k_{\perp}}{3m} \right)$$

$$x < x_{coal}(p) = 1 + p/3, \quad p = p_1 + p_2 + p_3$$

$\rho=n-3$

M.A. Braun, V.V. Vechernin, Theor.Math.Phys. 139, 766 (2004)

$$f_p(t) = 2\pi \int_0^{\infty} dz z J_0(tz) [z K_1(z)]^p$$

$J_0(z)$ - the Bessel function, $K_1(z)$ - the modified Bessel function.

$$(2\pi)^{-2} \int f_p(|\mathbf{b}|) d^2\mathbf{b} = (2\pi)^{-1} \int_0^{\infty} f_p(t) t dt = 1$$

Note that for $p=1$ it can be simplified to $f_1(t) = 4\pi/(t^2 + 1)^2$

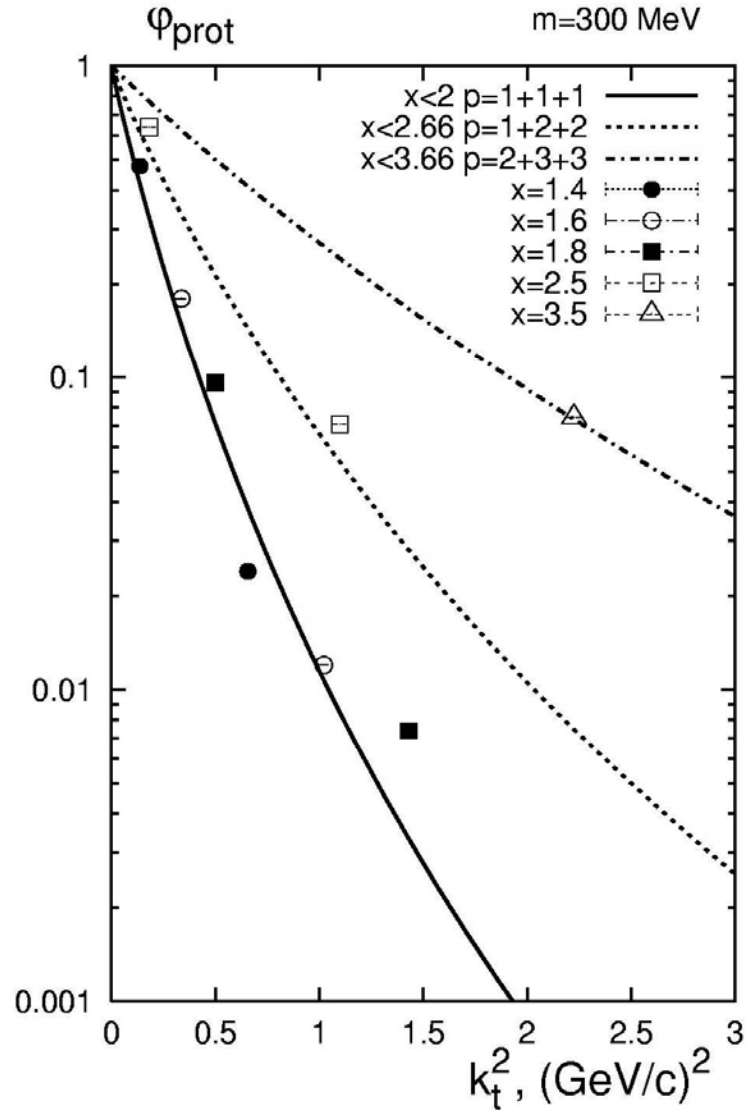
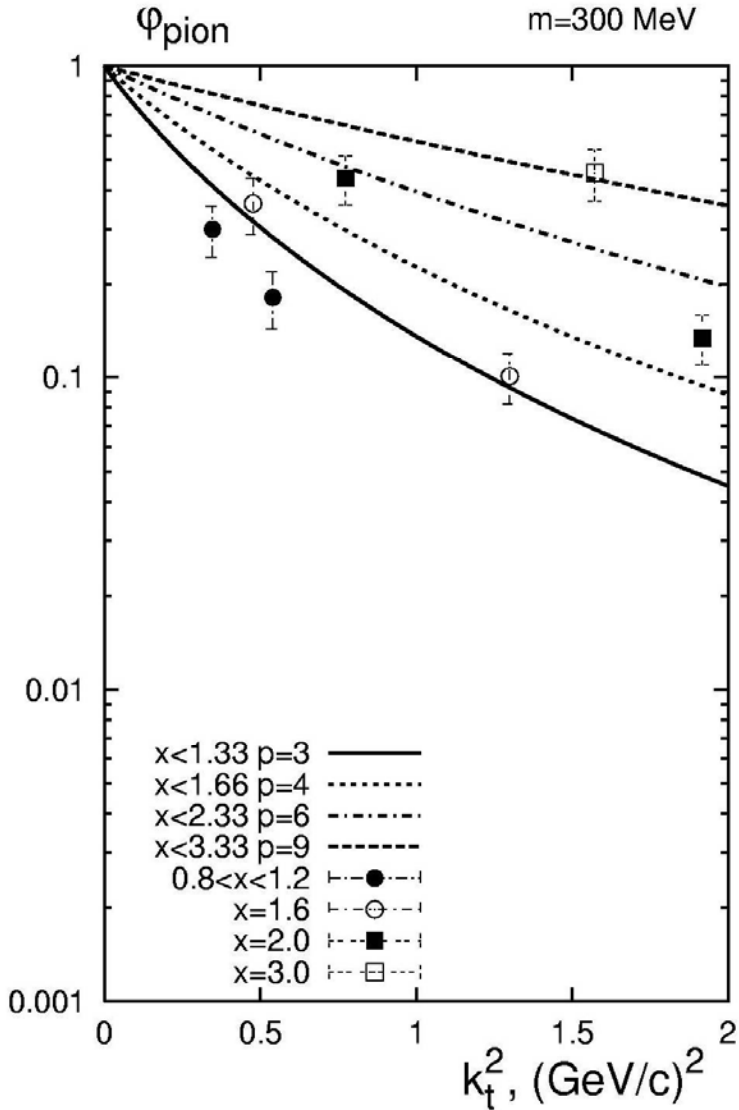
$$\varphi_{pion}(k_{\perp}, p) \equiv \sigma_{pion}(x, k_{\perp}; p) / \sigma_{pion}(x, 0; p) = f_p \left(\frac{k_{\perp}}{m} \right) / f_p(0)$$

$$\varphi_{prot}(k_{\perp}, p) \equiv \sigma_{prot}(x, k_{\perp}; p) / \sigma_{prot}(x, 0; p)$$

$$\varphi_{prot}(k_{\perp}, p) = \frac{\sum_{p_1, p_2, p_3} \delta_{p, p_1+p_2+p_3} C(p_1, p_2, p_3) f_{p_1} \left(\frac{k_{\perp}}{3m} \right) f_{p_2} \left(\frac{k_{\perp}}{3m} \right) f_{p_3} \left(\frac{k_{\perp}}{3m} \right)}{\sum_{p_1, p_2, p_3} \delta_{p, p_1+p_2+p_3} C(p_1, p_2, p_3) f_{p_1}(0) f_{p_2}(0) f_{p_3}(0) \dots}$$

$$\varphi_{prot}(k_{\perp}, p_1, p_2, p_3) \equiv \frac{\sigma_{prot}(x, k_{\perp}; p_1, p_2, p_3)}{\sigma_{prot}(x, 0; p_1, p_2, p_3)} = \frac{f_{p_1} \left(\frac{k_{\perp}}{3m} \right) f_{p_2} \left(\frac{k_{\perp}}{3m} \right) f_{p_3} \left(\frac{k_{\perp}}{3m} \right)}{f_{p_1}(0) f_{p_2}(0) f_{p_3}(0)}$$

No free parameters (!) only m – the constituent quark mass: $m = 300 \text{ MeV}$.



S.V. Boyarinov et al., Sov.J.Nucl.Phys. **46**, 871 (1987)

S.V. Boyarinov et al., Physics of Atomic Nuclei **57**, 1379 (1994)

S.V. Boyarinov et al., Sov.J.Nucl.Phys. **55**, 917 (1992)