

# Magnetohydrodynamics, charged currents and directed flow in heavy ion collisions

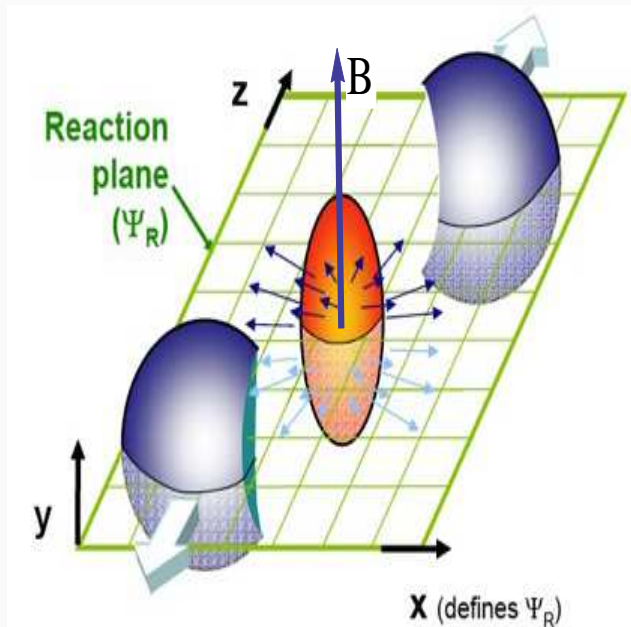
**Umut Gürsoy**

**Utrecht University**

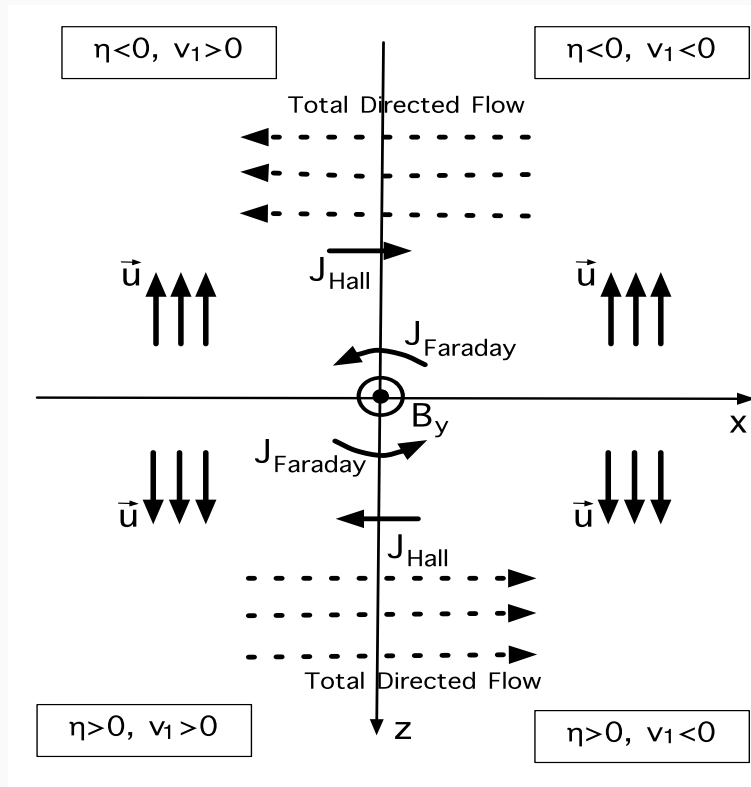
**Quark Confinement XI, St. Petersburg 12.09.2014**

**with D. Kharzeev and K. Rajagopal**  
**Phys. Rev. C, 089 (2014), arXiv:1401.3805**

# Heavy ion collisions and magnetic fields



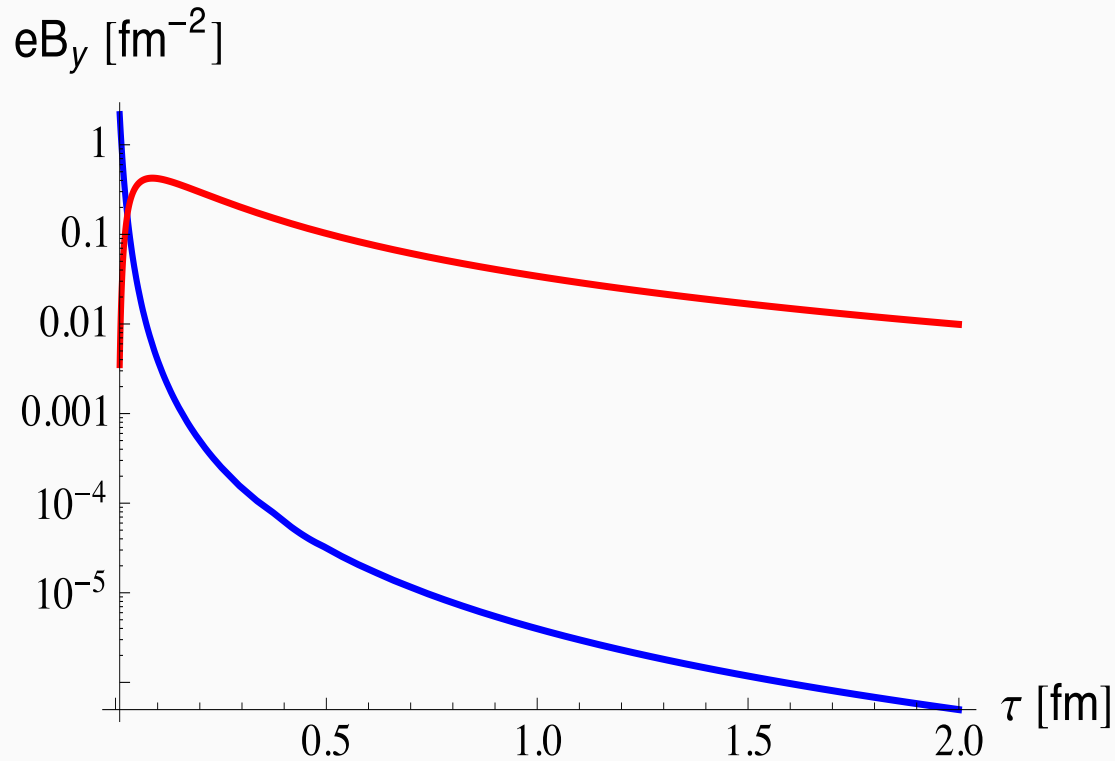
- Initial magnitude of  $B$
- Bio-Savart:  $B_0 \sim \gamma Z e \frac{b}{R^3} \Rightarrow eB \approx 5 - 15 \times m_\pi^2$  at RHIC (LHC).
- In this talk  $b = 7\text{fm}$  and  $R = 7\text{fm}$ .
- Motivation: find observables that are directly tied to the presence of  $B$



“Classical” currents in charged and expanding medium:

- Faraday currents  $\vec{J}_F \sim \sigma \vec{E}_F$  with  $\nabla \times \vec{E}_F = -\frac{\partial \vec{B}}{\partial t}$
- Hall currents  $\vec{J}_H \sim \sigma \vec{E}_H$  with  $\vec{E}_H = \vec{u} \times \vec{B}$

# Time profile of B at LHC



with  $\sigma = 0.023\text{fm}^{-1}$  and with  $\sigma = 0$

- Simplifying assumption **hard-sphere distribution** for **spectators** and **participants**
- For participants **empirical distribution** over  $Y$ : [Kharzeev et al. 2007](#)

$$f(Y_b) = (4 \sinh(Y_0/2))^{-1} e^{Y_b/2}, \quad -Y_0 \leq Y_b \leq Y_0$$

# Perturbative Magnetohydrodynamics

- Suppose  $u^\mu(x)$  with no back reaction of electromagnetic fields and  $\vec{E}, \vec{B}$  are known
- Go to the **comoving frame** by  $\Lambda(-\vec{u})$  e.g.  $F'_{\mu\nu} = (\Lambda \cdot F \cdot F)_{\mu\nu}$
- Compute the **stationary velocity**:

$$m \frac{d\langle v_B^\vec{} \rangle}{dt} = q \langle v_B^\vec{} \rangle \times \vec{B}' + q \vec{E}' - \mu m \langle v_B^\vec{} \rangle = 0,$$

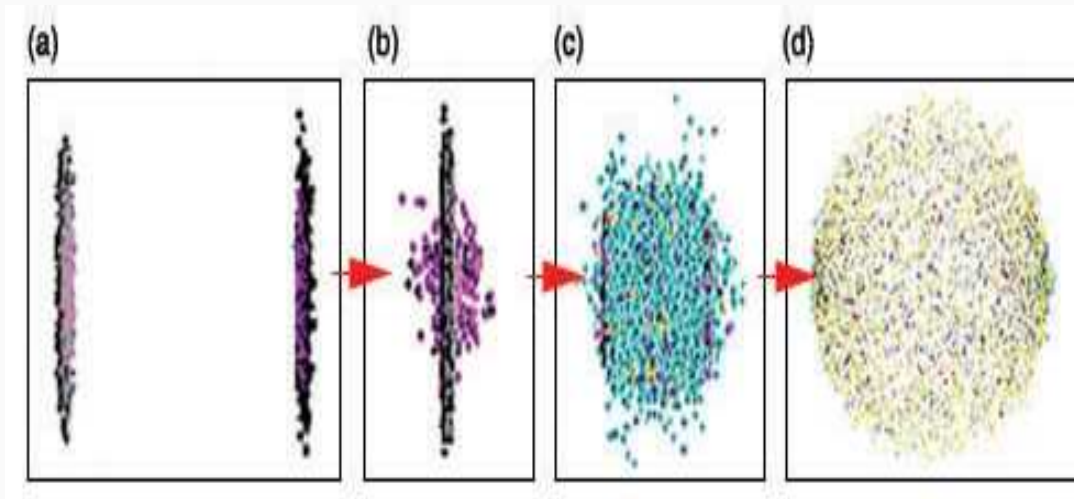
$\mu m$  the drag coefficient; e.g. from AdS/CFT:  $\frac{\pi\sqrt{\lambda}}{2} T^2$

- Go back to the center of mass frame: by  $V^\mu = \Lambda(\vec{u})^\mu_\nu v_B^\nu$
- $V^\mu$  contains both  $u^\mu$  and  $v_B^\mu$   
 $\Rightarrow$  construct observables from  $V$

# Assumptions

- Validity of **Perturbative magnetohydro**  $\Rightarrow |\vec{v}_B| \ll |\vec{u}|$
- Validity of **classical** force equation  $\Rightarrow$  magnetic energy  $E_B \ll \frac{2\pi h}{\lambda}$  with  $\lambda \approx R = 7$  fm.
- **Both are checked in our set-up**
- **First may be violated in reality**

# Constructing $u^\mu$ for the expanding fluid



- Start from the **Bjorken flow**: Bjorken '83
  1. **Boost invariance** along  $z$ :  $\xi = z\partial_t + t\partial_x$
  2. **Rotation around  $z$** :  $\xi = x\partial_y - y\partial_x$
  3. **Translations in transverse plane**:  $\xi = \partial_x$  and  $\xi = \partial_y$
- Solution to  $[\xi, u] = 0$  is  $u = \partial_\tau$  ( $ds^2 = -d\tau^2 + \tau^2 d\eta^2 + dx_\perp^2 + x_\perp^2 d\phi^2$ )
- Fine except transverse translations

# Gubser's flow solution

Gubser '10

- Begin by **Bjorken's flow**
- **Replace**  $\xi_i = \partial_x, \partial_y$  with  $\xi_i = \partial_i + q^2 [2x^i x^\mu \partial_\mu - x^\mu x_\mu \partial_i]$
- Solution to  $[\xi, u] = 0$  is  $u = \cosh \kappa \partial_\tau + \sinh \kappa \partial_\perp$  with
$$\kappa = \frac{2q^2 \tau x_\perp}{1 + q^2 \tau^2 + q^2 x_\perp^2}$$
- Solution to Hydrodynamics:  $\nabla_\mu T^{\mu\nu} = 0$  with
$$\epsilon = \frac{\hat{\epsilon}_0}{\tau^{4/3}} \frac{(2q)^{8/3}}{[1 + 2q^2(\tau^2 + x_\perp^2) + q^4(\tau^2 - x_\perp^2)^2]^{4/3}}$$
- Also analytic dissipative correction with  $\eta/S$ .
- Two parameters to fix: **Initial energy**  $\hat{\epsilon}_0$  and “system size”  $1/q$



- Gubser's solution has **axial symmetry** around the beam axis
- How is this consistent with **non-zero  $B$**  then?

- Gubser's solution has **axial symmetry** around the beam axis
- How is this consistent with **non-zero B** then?
- This is actually **precisely** what we want:
- Adding electromagnetism on an axial plasma  $\Rightarrow$  extract effects on only **charge identified hadrons!**
- e.g.  $v_1, v_2$ , etc. flow coefficients will be only for charged constituents, neutral background subtracted.
- We will see clearly:  $v_1 = v_2 = \dots = 0$  for **Gubser flow**  $u^\mu$  but non-zero for **Gubser + electrodynamics**  $V^\mu = u^\mu + v_B^\mu$

# How to test Gubser's flow?

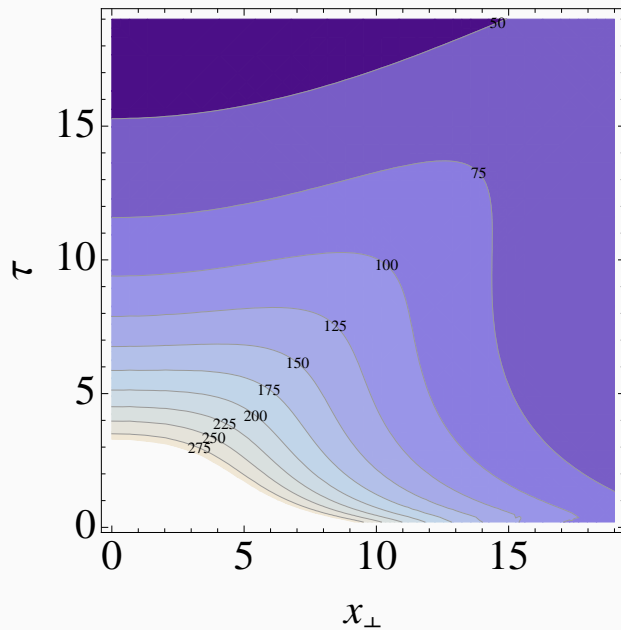
- Hadron spectrum from hydrodynamic flow: **Cooper-Frye**:

$$S_i = p^0 \frac{dN_i}{dp^3} = -\frac{g_i}{(2\pi)^3} \int d\Sigma_\mu p^\mu F\left(\frac{p^\mu V_\mu}{T_f}\right)$$

# How to test Gubser's flow?

- Hadron spectrum from hydrodynamic flow: **Cooper-Frye:**

$$S_i = p^0 \frac{dN_i}{dp^3} = -\frac{g_i}{(2\pi)^3} \int d\Sigma_\mu p^\mu F\left(\frac{p^\mu V_\mu}{T_f}\right)$$

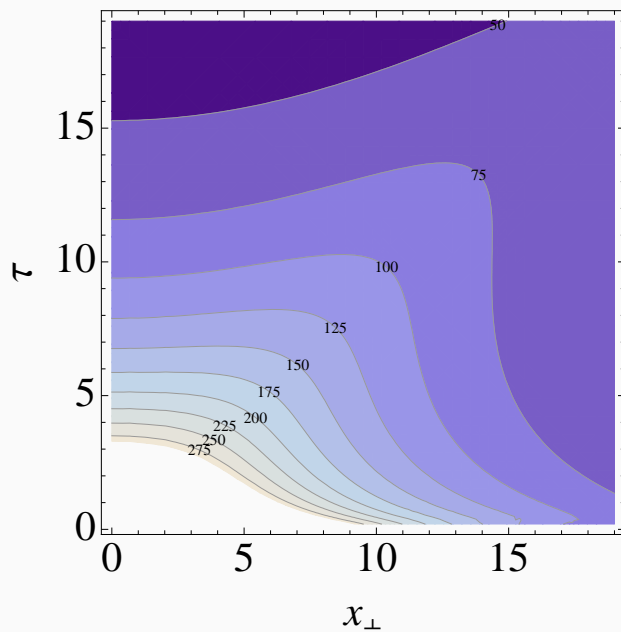


- Isothermal “freezeout surface”
- $T_f$  is the freezeout temperature,  
 $T_f \approx 130$  MeV
- Assume Boltzmann distribution:  
 $F(x) = e^x$

# How to test Gubser's flow?

- Hadron spectrum from hydrodynamic flow: **Cooper-Frye**:

$$S_i = p^0 \frac{dN_i}{dp^3} = -\frac{g_i}{(2\pi)^3} \int d\Sigma_\mu p^\mu F\left(\frac{p^\mu V_\mu}{T_f}\right)$$



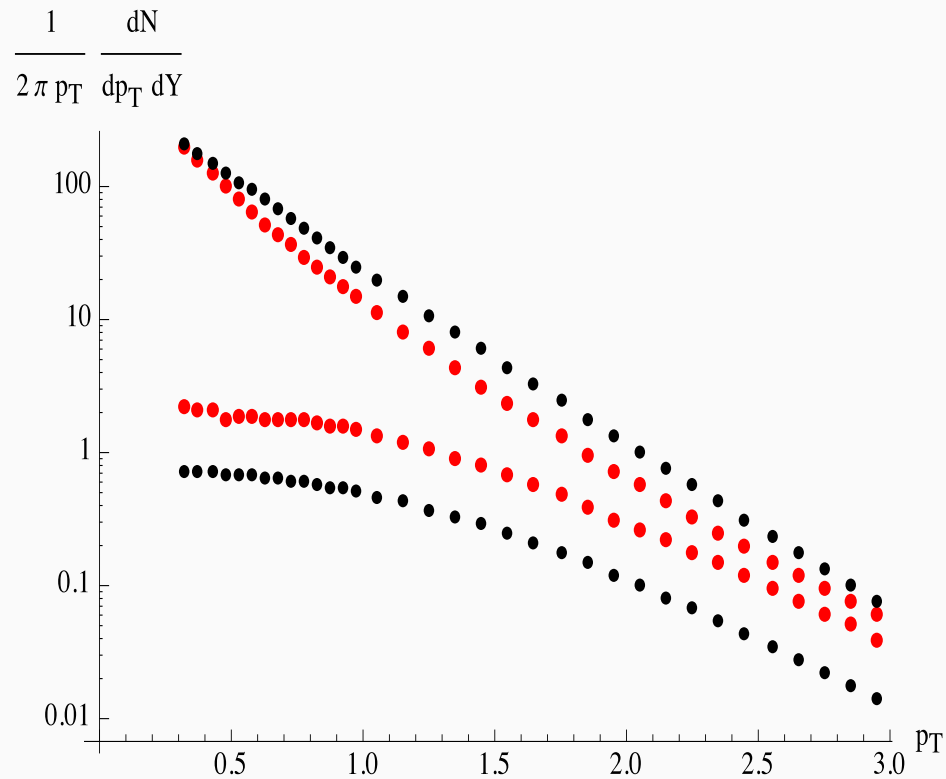
- Isothermal “freezeout surface”
- $T_f$  is the freezeout temperature,  
 $T_f \approx 130$  MeV
- Assume Boltzmann distribution:  
 $F(x) = e^x$

- $S_i(p_T) =$

$$\frac{g_i}{2\pi^2} \int dx_\perp x_\perp \tau_f \left\{ K_1\left(\frac{m_T u^\tau}{T_f}\right) I_0\left(\frac{p_T u^\perp}{T_f}\right) - \tau'_f p_T K_0\left(\frac{m_T u^\tau}{T_f}\right) I_1\left(\frac{p_T u^\perp}{T_f}\right) \right\}$$

- **Gubser's flow is independent of  $\Phi_p$  and  $Y \Rightarrow v_n = 0$**

# Cooper-Frye and parameter fixing



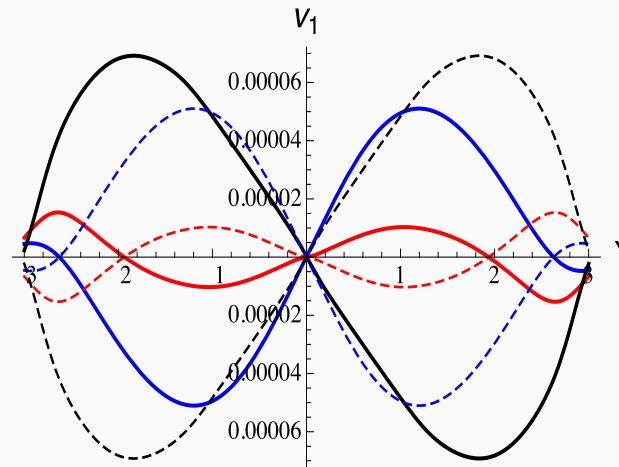
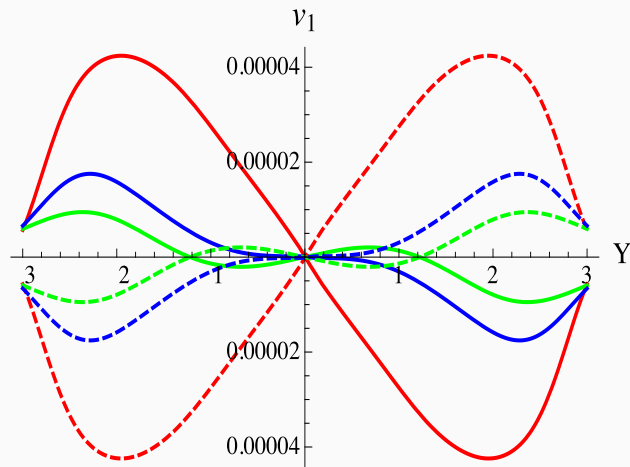
Upper  $\Rightarrow$  pions, lower  $\Rightarrow$  protons; red  $\Rightarrow$  ALICE, black  $\Rightarrow$  hydro

- Demand realistic comparison with **ALICE data** for **pions and protons** and reasonable **hadronization temperature**  
 $T_h \approx 400 - 550 \text{ MeV}$
- Optimal solution  $q^{-1} = 6.5 \text{ fm}$  and  $\hat{\epsilon}_0 = (8.7)^4$ .

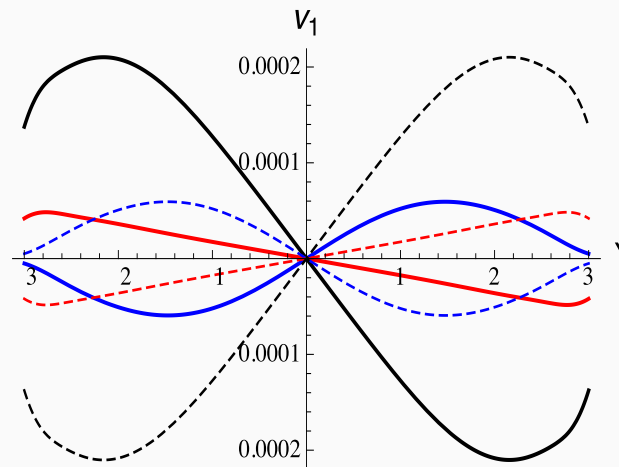
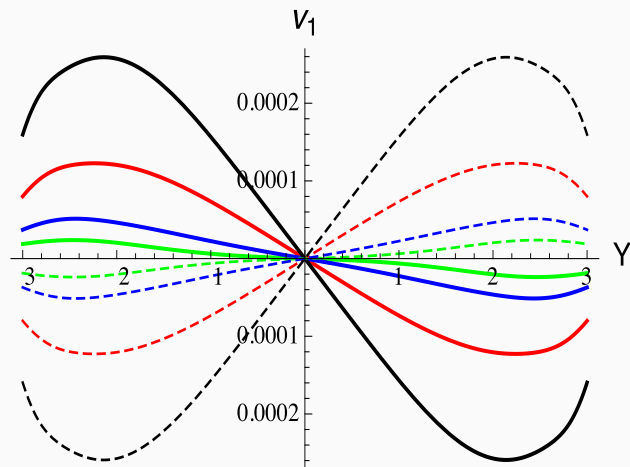
# Results

- Calculated  $B$
- Fixed Gubser's flow parameters  $\Rightarrow u^\mu$
- Solve classical force equation **electromagnetic force = drag**
- Do Cooper-Frye to calculate  $v_n$
- The simplest and most direct effect: **directed flow**  $v_1$ :

- Pions and protons at LHC

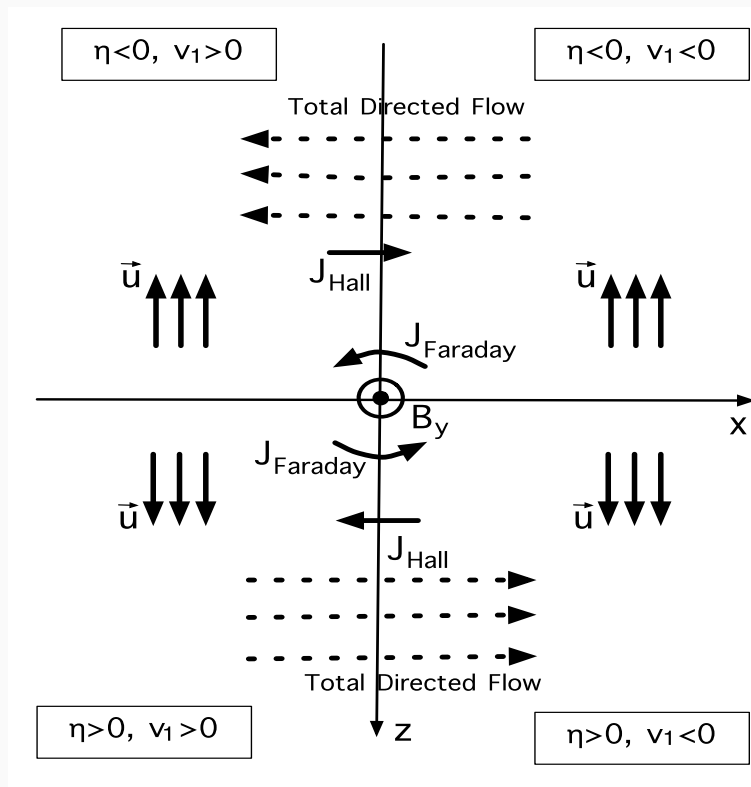


- Pions and protons at RHIC



Color coding:  $p_T = 0.25$  (green),  $0.5$  (blue),  $1$  (red),  $2$  (black) GeV





- Effect is small:  $\sim 5 \times 10^{-5}$  at LHC,  $\sim 2 \times 10^{-4}$  at RHIC.
- Specific features for detection:  $v_1(-Y) = -v_1(Y)$ ,  $v_1^+ = -v_1^-$

# Proposal for observables

- Define  $A_1^{+-}(Y_1, Y_2) = v_1^+(Y_1) - v_1^-(Y_2)$ ,  
 $A_1^{++}(Y_1, Y_2) = v_1^+(Y_1) - v_1^+(Y_2)$ , etc.  
to eliminate **charge independent contributions** to  $v_1$  produced in event-by-event fluctuations
- Look at **quadratic observables**  
 $C_1^{+-,+-}(Y, Y) = \langle A_1^{+-}(Y, Y) A_1^{+-}(Y, Y) \rangle = 4 \langle v_1^+(Y) v_1^+(Y) \rangle$   
to eliminate event-by-event fluctuations in **direction of B**.
- To be compared with data ...

## Summary:

- Calculated the contribution of the **time-varying B** in an **expanding plasma**, using a **perturbative approach to magnetohydrodynamics**.
- Effect **odd under charge and rapidity**.
- Competition between **Faraday** and **“Hall”** effects.
- However **the magnitude is small**.

## Summary:

- Calculated the contribution of the **time-varying B** in an **expanding plasma**, using a **perturbative approach to magnetohydrodynamics**.
- Effect **odd under charge and rapidity**.
- Competition between **Faraday** and “**Hall**” effects.
- However **the magnitude is small**.

## Outlook:

- Time dependence of  $\sigma, \mu, T$  etc.
- More realistic hydrodynamics.
- Backreaction of EM on hydro  $\Rightarrow$  full magnetohydrodynamics
- More realistic distributions for the sources
- Compute **charge identified**  $v_n$  for  $n \geq 2$ .

THANK YOU !