# The muon g-2: DSE status on light-by-light 

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Quark Confinement and the Hadron Spectrum XI
St. Petersburg, Russia
Sep 9, 2014

## Muon g-2

- Magnetic moment of a fermion due to its spin:

$$
\vec{\mu}=g \frac{e}{2 m} \vec{S} \quad \text { Pointlike fermion: } g=2
$$

Fermion with structure has anomalous magnetic moment: $\quad a=\frac{g-2}{2}$

- Electron \& muon anomalous magnetic moments among the most precisely measured \& theoretically calculated quantities:

- measured with precision $10^{-12}$ for electron and $10^{-10}$ for muon


## Muon g-2

- QED corrections: overwhelming part, calculated up to $\mathrm{O}\left(\alpha^{5}\right)$ :

- Electroweak and QCD corrections very small: $10^{-12}$ for electron, $10^{-8}$ for muon
- Electroweak corrections up to 2-loop:

- QCD:


Hadronic vacuum polarization


Hadronic light-by-light scattering
$a_{\mu}\left[10^{-10}\right]$
Jegerlehner, Nyffeler,

| Exp: | 11659208.9 | $(6.3)$ |
| :--- | ---: | ---: |
| QED: | 11658471.9 | $(0.0)$ |
| EW: | 15.3 | $(0.2)$ |

Hadronic:

| • VP (LO+HO) | 685.1 | $(4.3)$ |
| :--- | ---: | ---: |
| •LBL | 10.5 | $(2.6)$ |
| SM: | 11659182.8 | $(4.9)$ |
| Diff: | 26.1 | $(8.0)$ |

- Total SM prediction deviates from measured $a_{\mu}$ by $\sim 3 \sigma$ : new physics?
- Theory uncertainty dominated by QCD! Is QCD contribution under control?


## Light-by-light scattering



## Light-by-light scattering


 $+$



$+$

$\pi, K$ loop

$$
\begin{equation*}
-2 \tag{-10}
\end{equation*}
$$

Model results:
ENJL \& MD models
Bijnens 1995
Hakayawa 1995,
Knecht 2002,
Melnikov 2004,
Prades 2009,
Jegerlehner 2009,
Pauk 2014


Quark loop

- Constituent quark loop known analytically: 6 ... 8

- ENJL: VM poles by summing up quark bubbles Bijnens 1995
$\gamma^{\mu}-\gamma_{T}^{\mu} \frac{Q^{2}}{Q^{2}+m_{V}^{2}}$
Large reduction: 2
How to improve on this?


## Quark loop



- Quark mass is not a constant: dressed quark propagator has nonperturbatively enhanced quark mass function (DSE, Lattice, ...)

$$
S_{0}(p)=\frac{-i \not p+m}{p^{2}+m^{2}} \rightarrow S(p)=\frac{1}{A\left(p^{2}\right)} \frac{-i \not p+M\left(p^{2}\right)}{p^{2}+M^{2}\left(p^{2}\right)}
$$

- Quark-photon vertex is not bare!

$$
\begin{array}{ll}
\Gamma^{\mu}(k, Q)=\left[i \gamma^{\mu} \Sigma_{A}+2 k^{\mu}\left(i k \Delta_{A}+\Delta_{B}\right)\right]
\end{array}+\left[i \sum_{j=1}^{8} f_{j} \tau_{j}^{\mu}(k, Q)\right]
$$



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\Gamma^{\mu}(k, Q)=\left[i \gamma^{\mu} \Sigma_{A}+2 k^{\mu}\left(i k_{k} \Delta_{A}+\Delta_{B}\right)\right]+\left[i \sum_{j=1}^{8} f_{j} \tau_{j}^{\mu}(k, Q)\right]
$$



Ball-Chiu vertex, depends only on quark propagator Ball, Chiu, PRD 22 (1980)
necessary for electromagnetic gauge invariance!

$$
Q^{\mu} \Gamma^{\mu}(k, Q)=S^{-1}\left(k+\begin{array}{c}
Q \\
2
\end{array}\right)-S^{-1}\left(k-\begin{array}{c}
Q \\
2
\end{array}\right)
$$

| $A\left(p^{2}\right)$ | $M\left(p^{2}\right)$ | $\gamma^{\mu}$ | $\Gamma_{T}^{\mu}$ | $a_{\mu}\left[10^{-10}\right]$ |
| :---: | :---: | :---: | :---: | :---: |
| 1 | 0.2 GeV | 1 | 0 | 10 |
| 1 | $M\left(p^{2}\right)$ | 1 | 0 | 10 |
| $A\left(p^{2}\right)$ | $M\left(p^{2}\right)$ | 1 | 0 | 5 |
| $A\left(p^{2}\right)$ | $M\left(p^{2}\right)$ | $\Sigma_{A}$ | 0 | 10 |
| $A\left(p^{2}\right)$ | $M\left(p^{2}\right)$ | $\Sigma_{A}$ | $k=0$ | 4 |
| $A\left(p^{2}\right)$ | $M\left(p^{2}\right)$ | $\Sigma_{A}$ | Full | 10 |

- DSE result for quark loop (including strange \& charm):

$$
a_{\mu}=10.7 \times 10^{-10}
$$

- full Ball-Chiu vertex problematic


## Quark-photon vertex

- Quark-photon vertex: (vector current: $J^{\mu}=\bar{\psi} \gamma^{\mu} \psi$ )
$\langle 0| \mathrm{T} J^{\mu}(x) \psi\left(x_{1}\right) \bar{\psi}\left(x_{2}\right)|0\rangle$


Quark-photon vertex has $\rho$-meson poles: 'vector-meson dominance'

- Hadronic vacuum polarization = vector current correlator
$\langle 0| \mathbf{T} J^{\mu}(x) J^{\nu}(y)|0\rangle$

- BSE for quark four-point function \& quark photon vertex:

determine vertex dynamically from a given $q \bar{q}$ kernel, e.g. rainbow-ladder (= gluon exchange)

Maris \& Tandy, PRC 61 (2000)

## Context matters

## Hadron physics from QCD's Dyson-Schwinger \& bound-state equations

- Mesons from Bethe-Salpeter equation: meson spectra, form factors, PDFs, GPDs, ... Chang et al., Commun. Theor. Phys. 58 (2012), ...
- Baryons from covariant Faddeev equation: octet \& decuplet masses, nucleon \& $\Delta$ form factors GE, Alkofer, Krassnigg, Nicmorus, PRL 104 (2010), GE, PRD 84 (2011), Sanchis-Alepuz, Fischer, 1408.5577, ...


Compton scattering, pion electroproduction, ... GE, Fischer, PRD 85 (2012), PRD 87 (2013)

- Tetraquarks:
see talk by Walter Heupel
Friday,
- Beyond rainbow-ladder:
see talk by Richard Williams



## Quark-photon vertex

Structure of quark-photon vertex is reflected in hadron form factors GE, PRD 84(2011)
Experimentally (sketch):


Calculated:
(Sketch)



- Deriving hadronic currents: Kvinikhidze, Blankleider, PRC 60 (1999)
- Ball-Chiu part is dominant (em. gauge invariance): charge, magnetic moments
- Transverse part changes slope and charge radii. No pion cloud in RL $\Rightarrow$ timelike $\rho$-meson poles


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## Pion form factor



Spacelike and timelike region:
A. Krassnigg (Schladming 2010) extension of Maris \& Tandy, Nucl. Phys. Proc. Suppl. 161 (2006)

Include pion cloud:
Kubrak et al., in preparation

## Structure of the $\gamma Y Y \gamma$ amplitude



3 independent momenta:

$$
\begin{aligned}
p & =p_{2}+p_{3} \\
q & =p_{3}+p_{1} \\
k & =p_{1}+p_{2}
\end{aligned}
$$

6 Lorentz invariants:

$$
p^{2}, \quad q^{2}, \quad k^{2}, \quad p \cdot q, \quad p \cdot k, \quad q \cdot k
$$

Bose symmetry:

$$
\begin{aligned}
& \Gamma^{\mu \nu \rho \sigma}(p, q, k)=\sum_{i=1}^{136} f_{i}(\ldots) \tau_{i}^{\mu \nu \rho \sigma}(p, q, k) \\
& \stackrel{!}{=} \text { symmetric } \\
& \text { S4 multiplets }
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- Arrange the 24 permutations of $\psi_{1234}$ into multiplets:

| Singlet | Triplets | Doublets | Antitriplets | Antisinglet |
| :---: | :---: | :---: | :---: | :---: |
| $\mathcal{S}$ | $\mathcal{T}_{i}^{+}=\left[\begin{array}{l}\bullet \\ \bullet \\ \bullet\end{array}\right]$ | $\mathcal{D}_{j}=\left[\begin{array}{l}\bullet \\ \bullet\end{array}\right]$ | $\mathcal{T}_{i}^{-}=\left[\begin{array}{l}\bullet \\ \bullet \\ \bullet\end{array}\right]$ | $\mathcal{A}$ |

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- Arrange the 24 permutations of $\psi_{1234}$ into multiplets:

- 6 Lorentz invariants form singlet $\mathcal{S}_{0}$, doublet $\mathcal{D}$, triplet $\mathcal{T}^{+}$


## Phase space

- Singlet: symmetric variable, carries overall scale:
$\mathcal{S}_{0}=\frac{p^{2}+q^{2}+k^{2}}{4}=\frac{p_{1}^{2}+p_{2}^{2}+p_{3}^{2}+p_{4}^{2}}{4}$
- Doublet: $\mathcal{D}=\left[\begin{array}{l}a \\ s\end{array}\right]$

Mandelstam triangle,
2-photon poles (pion, scalar, axialvector, ...)


- Triplet: $\mathcal{T}=\left[\begin{array}{l}u \\ v \\ w\end{array}\right]$
tetrahedron bounded by $p_{i}^{2}=0$, vector-meson poles



## Phase space

- fixed doublet variables $\Rightarrow$ complicated geometric object inside tetrahedron:

$r=0$

$r=1$
$\varphi=0$

$r=0.9$
$\varphi=\pi$

relevant for g-2


## Tensor basis I

- construct all possible multiplets from generic seed elements:
138 elements, but only 136 independent
- removing the "wrong ones" leads to kinematic singularities!
- Dressing functions form multiplets too $\Rightarrow$ expand them into singlets:

$$
\left[\begin{array}{l}
f_{1}\left(\mathcal{S}_{0}, \mathcal{D}_{0}, \mathcal{T}_{0}\right) \\
f_{2}\left(\mathcal{S}_{0}, \mathcal{D}_{0}, \mathcal{T}_{0}\right)
\end{array}\right]=c_{1}\left(\mathcal{S}_{0}\right)[\bullet]_{1}+c_{2}\left(\mathcal{S}_{0}\right)[\bullet]_{2}+\ldots
$$

Singlets depend (almost) only on $\mathcal{S}_{0}$, dependence on $\nabla, \Delta$ absorbed in basis

This works extremely well!

| $n$ | Seed | $\#$ | Multiplet type |
| :--- | :--- | :--- | :--- |
| 0 | $\delta^{\mu \nu} \delta^{\rho \sigma}$ | 3 | $\mathcal{S}, \mathcal{D}_{1}$ |
| 2 | $\delta^{\mu \nu} k^{\rho} k^{\sigma}$ | 6 | $\mathcal{S}, \mathcal{D}_{1}, \mathcal{T}_{1}^{+}$ |
|  | $\delta^{\mu \nu} p^{\rho} p^{\sigma}$ | 12 | $\mathcal{S}, \mathcal{D}_{1}, \mathcal{D}_{2}, \mathcal{T}_{1}^{ \pm}, \mathcal{A}$ |
|  | $\delta^{\mu \nu} p^{\rho} q^{\sigma}$ | 12 | $\mathcal{S}, \mathcal{D}_{1}, \mathcal{T}_{1}^{+}, \mathcal{T}_{2}^{ \pm}$ |
|  | $\delta^{\mu \nu} p^{\rho} k^{\sigma}$ | 24 | $\mathcal{S}, \mathcal{D}_{1}, \mathcal{D}_{2}, \mathcal{T}_{1}^{ \pm}, \mathcal{T}_{2}^{ \pm}, \mathcal{T}_{3}^{ \pm}, \mathcal{A}$ |
| 4 | $p^{\mu} p^{\nu} p^{\rho} p^{\sigma}$ | 3 | $\mathcal{S}, \mathcal{D}_{1}$ |
|  | $p^{\mu} p^{\nu} q^{\rho} q^{\sigma}$ | 6 | $\mathcal{S}, \mathcal{D}_{1}, \mathcal{T}_{1}^{-}$ |
|  | $p^{\mu} p^{\nu} k^{\rho} k^{\sigma}$ | 10 | $\mathcal{S},\left(\mathcal{D}_{1}, \mathcal{D}_{2}, \mathcal{T}_{1}^{ \pm}, \mathcal{A}\right.$ |
|  | $p^{\mu} q^{\nu} k^{\rho} k^{\sigma}$ | 12 | $\mathcal{S}, \mathcal{D}_{1}, \mathcal{T}_{1}^{+}, \mathcal{T}_{2}^{ \pm}$ |
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Singlets depend (almost) only on $\mathcal{S}_{0}$, dependence on $\nabla, \Delta$ absorbed in basis


This works extremely well!
$\rightarrow$ YYYY amplitude calculated in full kinematics!

- reproduces previous DSE results for g-2, but complete Ball-Chiu vertex still erratic...?


## Tensor basis II

To make gauge invariance explicit, split YYyץ amplitude into
$\Gamma=\Gamma_{\text {Gauge }}+\Gamma_{\text {Transverse }}$

- Transverse part: 41 tensors, at least quartic in photon momenta, dominant tensors $\sim$ pion + scalar exchange

- 'Gauge part': 95 tensors, must vanish if pYYץ amplitude gauge invariant!
... but is it?



## Results: quark loop with $m_{q}=$ const

Photon four-point function: $S_{o}$ dependence for fixed doublet \& triplet variables


## Results: quark loop from DSE

Photon four-point function: $S_{0}$ dependence for fixed doublet \& triplet variables





goal:

- find transverse basis w/o kin. singularities
- isolate dominant tensor structures


## Gauge invariance

Full $\mathrm{Y} Y Y \gamma$ amplitude at quark level, derived from gauge invariance:
GE, Fischer, PRD 85 (2012), Goecke, Fischer, Williams, PRD 87 (2013)



Quark loop + all 2-photon poles from T-matrix (pion, scalar, axialvector, ...)

- no double-counting!
- gauge artifacts in quark loop must be cancelled by offshell structure of T-matrix!


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- no double-counting!
- gauge artifacts in quark loop must be cancelled by offshell structure of T-matrix!
- Quark Compton vertex already determined from nucleon Compton scattering:


GE, Fischer, PRD 87 (2013),
PoS Conf. X (2012)

## Summary

Muon g-2: theory uncertainty dominated by QCD


LBL: need to get QCD contribution under control!

- yypy amplitude = quark loop + T-matrix, no double counting, gauge invariant!
- need to understand structure of $\mathrm{p} \gamma \mathrm{Y} \mathrm{p}$ amplitude
- dressed quarks \& vertices have impact, QCD prediction for LBL may change!

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## Hadronic:

- VP (LO+HO) 685.1
- LBL 10.5 (2.6) ?
SM
11659182.8 (4.9)
Diff:
26.1 (8.0)


## Electron vs. muon g-2

| Exp: | 11596521.81 |  |
| :---: | :---: | :---: |
| QED: | $\begin{array}{r} 11596521.71 \\ .81 \end{array}$ | $\begin{aligned} & (0.09) \\ & (0.08) \end{aligned}$ |
| EW: | 0.00 |  |
| Hadronic: | c: 0.02 |  |
| SM: | 11596521.73 .83 | $\begin{aligned} & (0.09) \\ & (0.08) \end{aligned}$ |

$$
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$$

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Bijnens, Prades, Mod. Phys. Lett. A22 (2007) Jegerlehner, Nyffeler, Phys. Rept. 477 (2009) Hagiwara et al., J. Phys. G 38 (2011)

