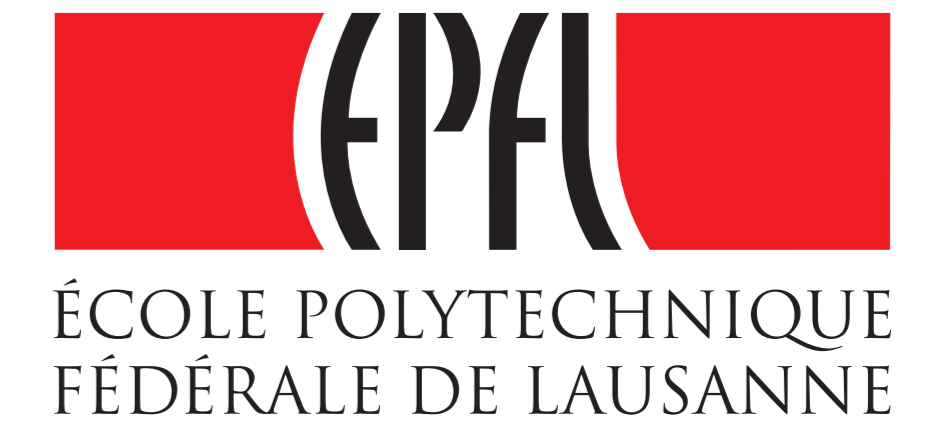


The in-medium heavy-quark potential from quenched and dynamical lattice QCD



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Physics motivation

■ **Heavy quarkonium:** the bound states of a heavy quark and antiquark are intensely studied probes in relativistic heavy-ion collisions. Bottomonium ($b\bar{b}$) is of particular interest as its yields have recently been obtained by the CMS collaboration with unprecedented precision. They show a suppression of excited states in the presence of a hot medium. Once we understand the real-time in-medium evolution of quarkonium it is possible to extract from the measured yields the properties of the surrounding quark-gluon plasma. [CMS collaboration] PRL 109 (2012) 222301

■ **In-medium heavy quark potential:** The separation of scales $m_Q \gg T$ and $m_Q \gg \Lambda_{\text{QCD}}$ allows describing the interaction between heavy QQ and its environment in the language of a Schrödinger equation with in-medium potential. Initially the potential was modeled using e.g. the color singlet free energies measured on the lattice. Modern effective field theories such as NRQCD and pNRQCD define V^{QQ} directly from a QCD observable, the real-time Wilson loop. Brambilla et al. Rev.Mod.Phys. 77 (2005)

$$V^{\text{QQ}}(r) = \lim_{t \rightarrow \infty} \frac{i \partial_t W(r, t)}{W(r, t)} \quad W(r, t) = \left\langle \text{Tr} \left[\exp \left(-i g \int_{\square} dz^\mu A_\mu(z) \right) \right] \right\rangle_{\text{QCD}@T>0}$$

Brambilla, Ghiglieri, Vairo and Petreczky PRD 78 (2008) 014017, for a brief review see A.R. MPLA 28 (2013) 133000

■ **Perturbative evaluation of the potential at high T:** The hard thermal loop approximation enabled the evaluation of the real-time Wilson loop and revealed, V^{QQ} is in general complex. Laine et al. JHEP03 (2007) 054

■ $\text{Re}[V]$ shows Debye screening (consistent with the classic idea of Matsui and Satz)

■ $\text{Im}[V]$ from medium gluon scattering (Landau damping) and absorption (Singlet to Octet transition) Beraudo et al. NPA 806:312,2008 Brambilla et al. PRD 78 (2008) 014017

■ **Goal:** Extract the values of the complex in-medium heavy quark potential non-perturbatively from lattice QCD simulations at phenomenologically relevant $T \sim T_c$

■ **Challenge 1:** Potential is a real-time object but lattice QCD simulations are performed in Euclidean time. **Strategy:** Express the potential in terms of spectral functions

■ **Challenge 2:** Spectral functions are not directly accessible on the lattice, they require an analytic continuation. **Strategy:** Bayesian inference of spectra from Euclidean data

In-medium potential from lattice QCD

■ V^{QQ} defined in real-time: connection to Euclidean lattice QCD via spectral functions

■ Spectral representation of the real-time Wilson loop is a simple Fourier transform and ρ positive definite A.R., T.Hatsuda & S.Sasaki PoS LAT2009 (2009) 162

$$W(r, t) = \int_{-\infty}^{\infty} d\omega e^{-i\omega t} \rho(r, \omega) \quad \longleftrightarrow \quad W(r, \tau) = \int_{-\infty}^{\infty} d\omega e^{-\omega \tau} \rho(r, \omega)$$

$$V^{\text{QQ}}(r) = \lim_{t \rightarrow \infty} \frac{\int_{-\infty}^{\infty} d\omega \omega e^{i\omega t} \rho(r, \omega)}{\int_{-\infty}^{\infty} d\omega e^{i\omega t} \rho(r, \omega)} \quad \longleftarrow \quad \text{Bayesian spectral reconstruction}$$

A.R., T.Hatsuda & S.Sasaki PRL 108 (2012) 162001

■ Extracting $\rho(r, \omega)$ from $W(r, \tau)$ amounts to an inverse Laplace transform: ill-posed problem (Bayesian inference)

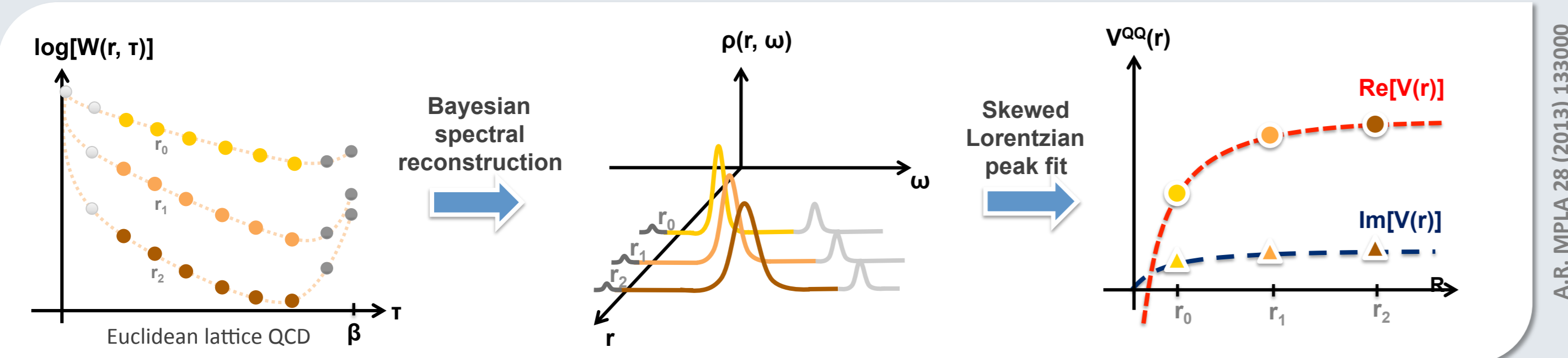
■ The most general form of the Wilson loop spectrum: disentangling late and early times

$$\rho(r, \omega) \propto \frac{|\text{Im}V(r)| \cos[\text{Re}\sigma_\infty(r)] - (\text{Re}V(r) - \omega) \sin[\text{Re}\sigma_\infty(r)]}{\text{Im}V(r)^2 + (\text{Re}V(r) - \omega)^2} + c_0(r) + c_1(r)(\text{Re}V(r) - \omega) + \dots$$

Y.Burnier, A.R. Phys.Rev. D86 (2012) 051503

■ Potential represents late times physics: encoded in position and width of a low lying skewed Lorentzian peak.

■ **Summary:** From Euclidean correlators to the in-medium heavy quark potential



A.R. MPLA 28 (2013) 133000

Improved Bayesian spectral reconstruction

■ An **ill-posed problem:** extract a spectral function $\rho(\omega) = \rho_l$ along N_ω frequencies from $N_\tau \ll N_\omega$ noisy datapoints $D(\tau_i) = D_i$ (likelihood L fit alone is underdetermined)

$$D_i^\rho = \sum_{l=1}^{N_\omega} \Delta\omega_l e^{-\omega_l \tau_i} \rho_l \quad L = \frac{1}{2} \sum_{i,j} (D_i - D_i^\rho) C_{ij}^{-1} (D_j - D_j^\rho) \quad C_{ij} \text{ covariance matrix}$$

■ Bayes Theorem: Incorporation of **prior information (I)** regularizes the χ^2 fit

$$P[\rho|D, I] \propto P[D|\rho, I] P[\rho|I] \quad P[D|\rho, I] = \exp[-L - \gamma(L - N_\tau)^2] \quad \text{Since the correct } \rho \text{ leads to } L \sim N_\tau$$

■ **Improved prior functional:** enforces (1) positive definiteness of ρ (2) independence of the result from dimension of ρ (3) smoothness of ρ , where data does not imprint peaks

$$P[\rho|I] = e^S \quad S = \alpha \sum_{l=1}^{N_\omega} \Delta\omega_l \left(1 - \frac{\rho_l}{m_l} + \log \left[\frac{\rho_l}{m_l} \right] \right) \quad m_l \text{ default model: correct spectrum in the absence of data}$$

Y.Burnier, A. Rothkopf Phys.Rev.Lett. 111 (2013) 18, 182003

■ Different from Maximum Entropy Method: **S is not entropy and does not contain flat directions**

■ α is integrated out analytically $P[\alpha] = 1$, Bayesian solution as maximum of the posterior

$$P[\rho|D, I] \propto P[D|\rho, I] \int_0^\infty d\alpha P[\rho|I, \alpha] \quad \left. \frac{\delta}{\delta \rho} P[\rho|D, I] \right|_{\rho = \rho^{\text{BR}}} = 0$$

■ Implementation via LBFGS algorithm: **no restriction of the search space** allows higher than MEM resolution

■ Due to absence of flat direction in S , true convergence to a global extremum (no artificial cutoff in step size)

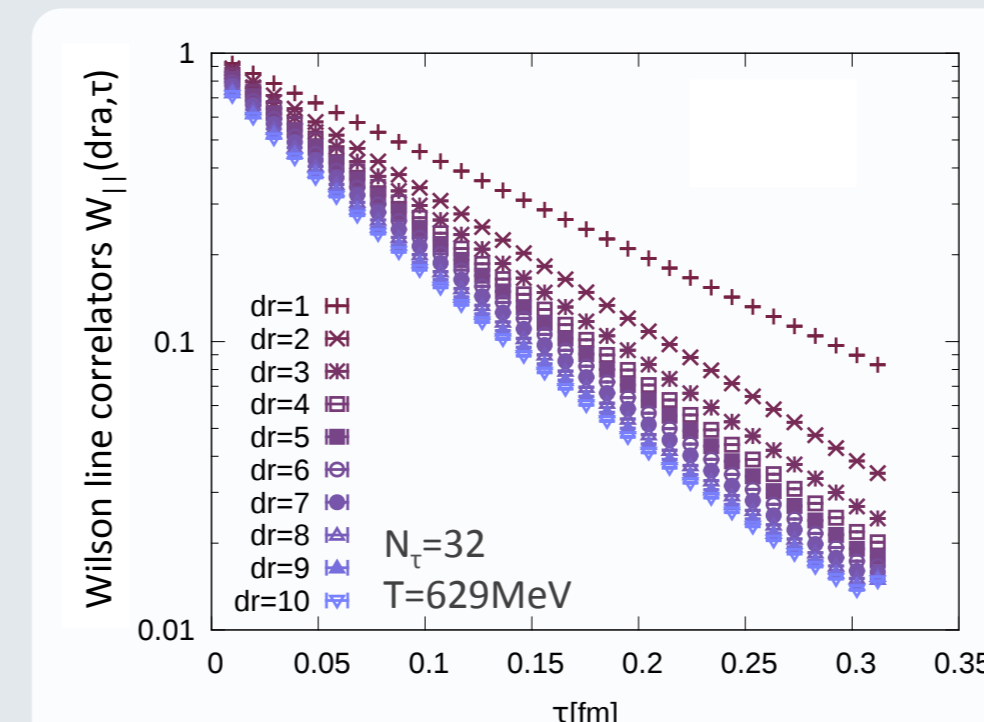
Quenched Lattice QCD

■ Fixed scale approach: anisotropic naïve Wilson action $\beta=7.0$ $\xi=a_s/a_t=4$ $a_s=0.039\text{fm}$

N_t	24	32	40	48	56	64	72	80	96
T/T_c	3.11	2.33	1.86	1.55	1.33	1.17	1.04	0.93	0.78
N_{meas}	3270	2030	1940	1110	1410	910	860	1190	1800

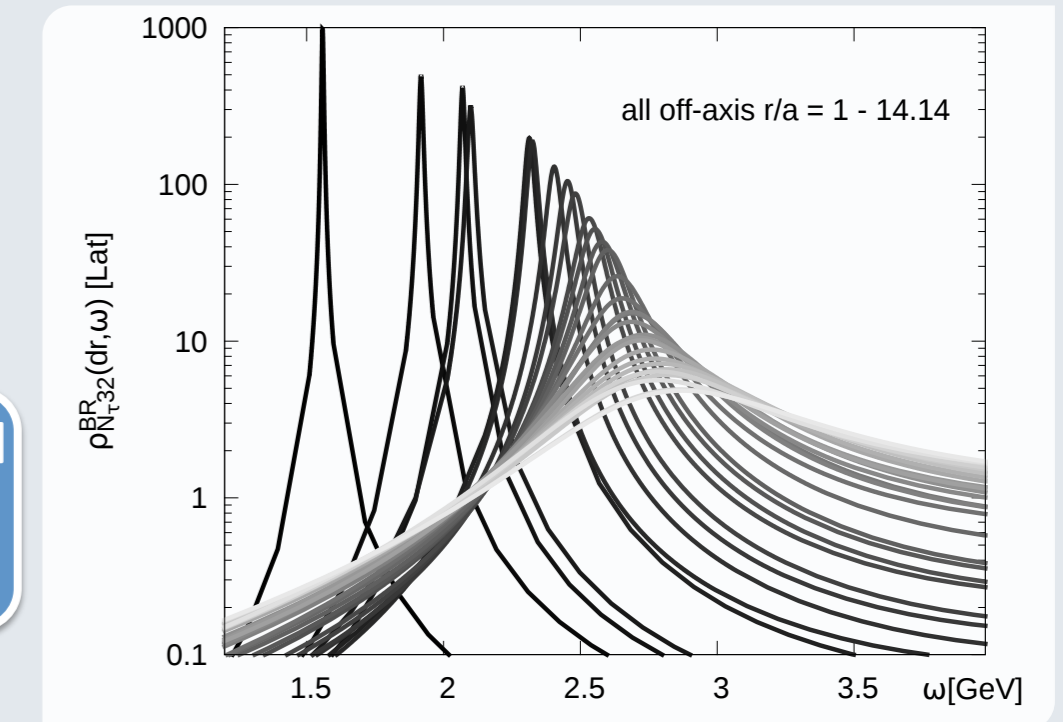
$T_c=271\text{MeV}$
close to continuum

■ Spectra from Wilson Line correlators in Coulomb gauge $W_{11}(r, \tau)$ to avoid cusp divergences in the Wilson loop



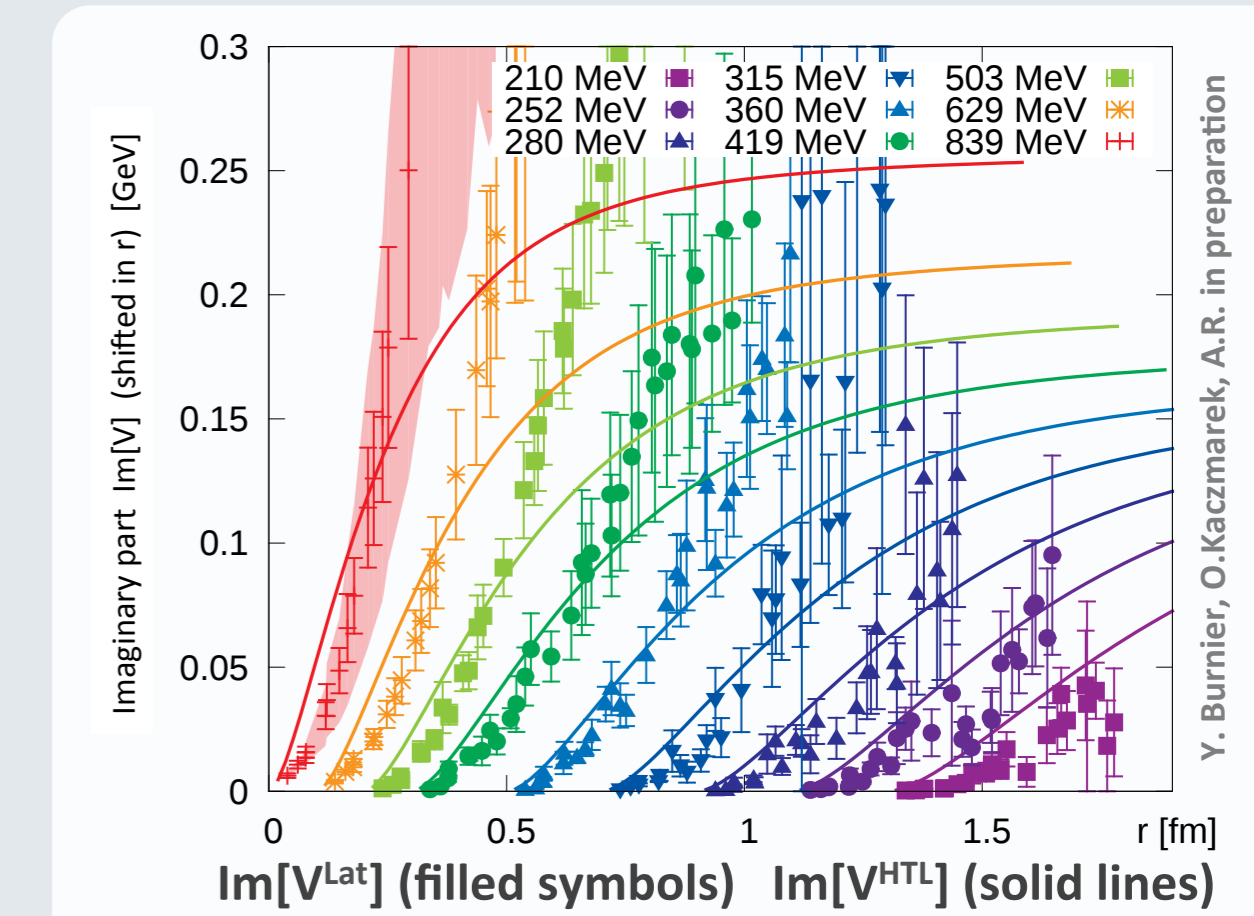
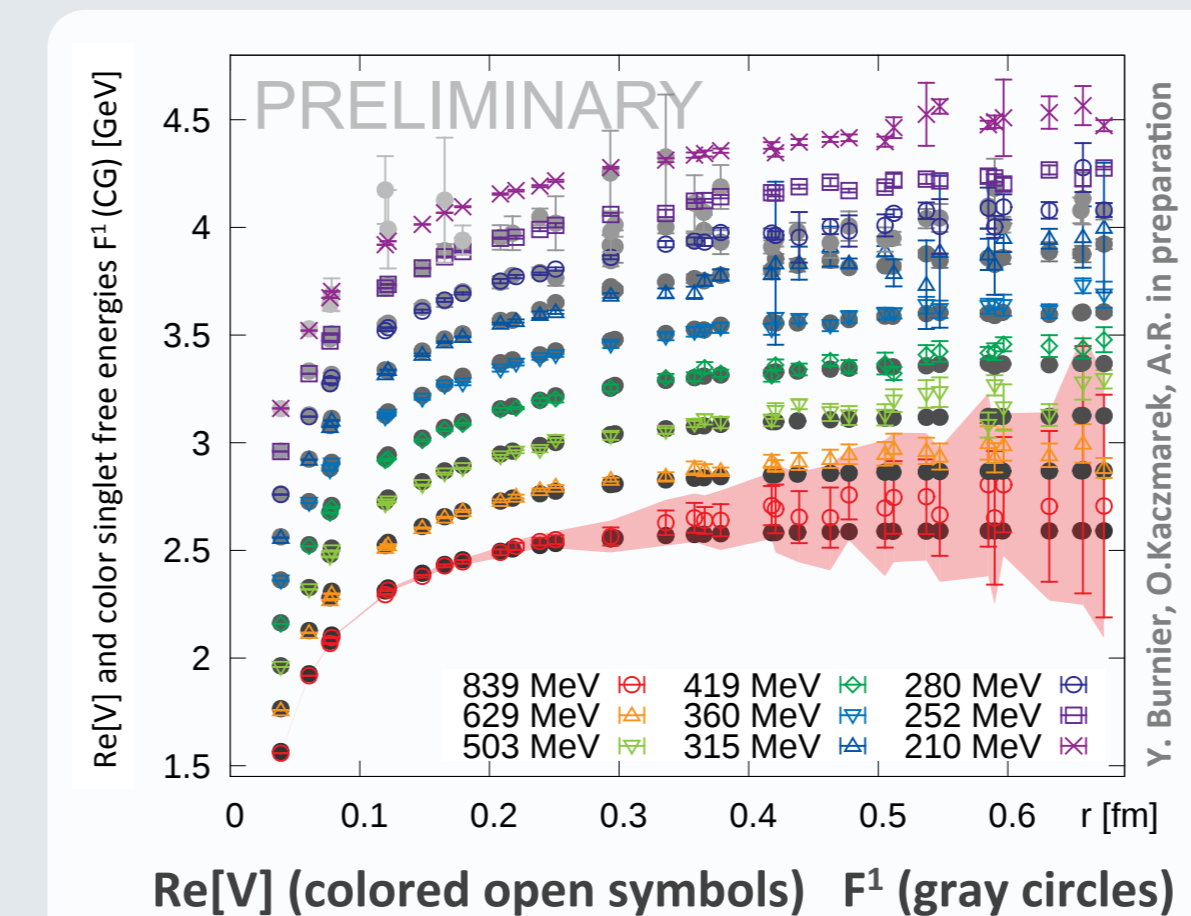
Bayesian spectral reconstruction

$N_\omega=4000, I_\omega^{\text{num}}=[-13, 14.5]$
 $\tau^{\text{max}}_{\text{num}}=20, m(\omega)=\text{const.}$
512bits prec., $\Delta^{\text{min}}=10^{-60}$



■ Wilson line correlators indicate $\text{Im}[V] \neq 0$ above T_c : curvature along τ corresponds to a finite spectral width

■ Long standing question: $\text{Re}[V]$ and Coulomb gauge color singlet free energies $F^1(r) = T \log[W_{11}(r, \beta)]$ similar?



■ Error bars denote statistical uncertainty (Jackknife), error band adds systematics (default model dependence)

■ Spatial distances are corrected for lattice discretization artefacts by comparing free correlators to continuum Necco, Sommer Nucl.Phys. B622 (2002) 328-346

■ **Re[V] transitions from confining to Debye screened behavior and lies close to the color singlet free energies** compare with: Y. Burnier et al., JHEP 01 (2010) 054

■ **Im[V]** is reconstructed robustly at small r : **similar magnitude** as predictions from **HTL perturbation theory**

■ Testing in progress: changing T by N_t modifies reconstruction accuracy, need to quantify this N_t dependence

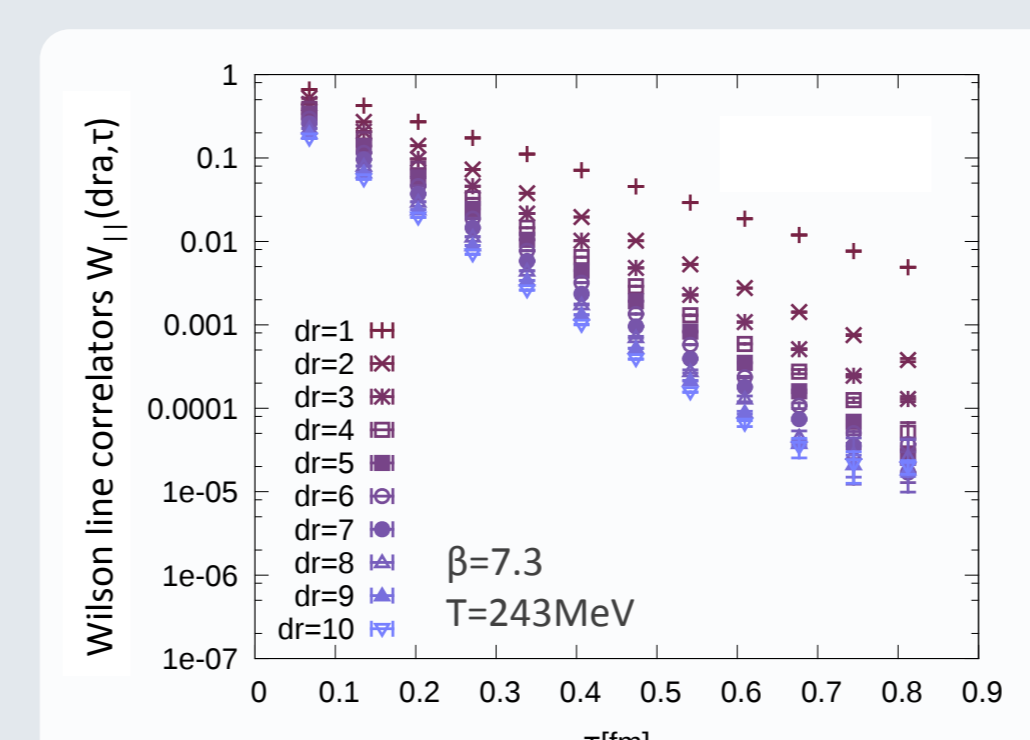
Dynamical $N_f=2+1$ Lattice QCD (HotQCD)

■ Fixed box approach: isotropic lattices with asqtad action $48^3 \times 12$ ($m_l/m_s=0.05$)

β	6.80	6.90	7.00	7.125	7.25	7.30	7.48
T/T_c	0.85	0.94	1.04	1.18	1.33	1.39	1.64
a [fm]	0.111	0.100	0.090	0.080	0.071	0.068	0.057
N_{meas}	1295	1340	1015	840	1620	1150	1130

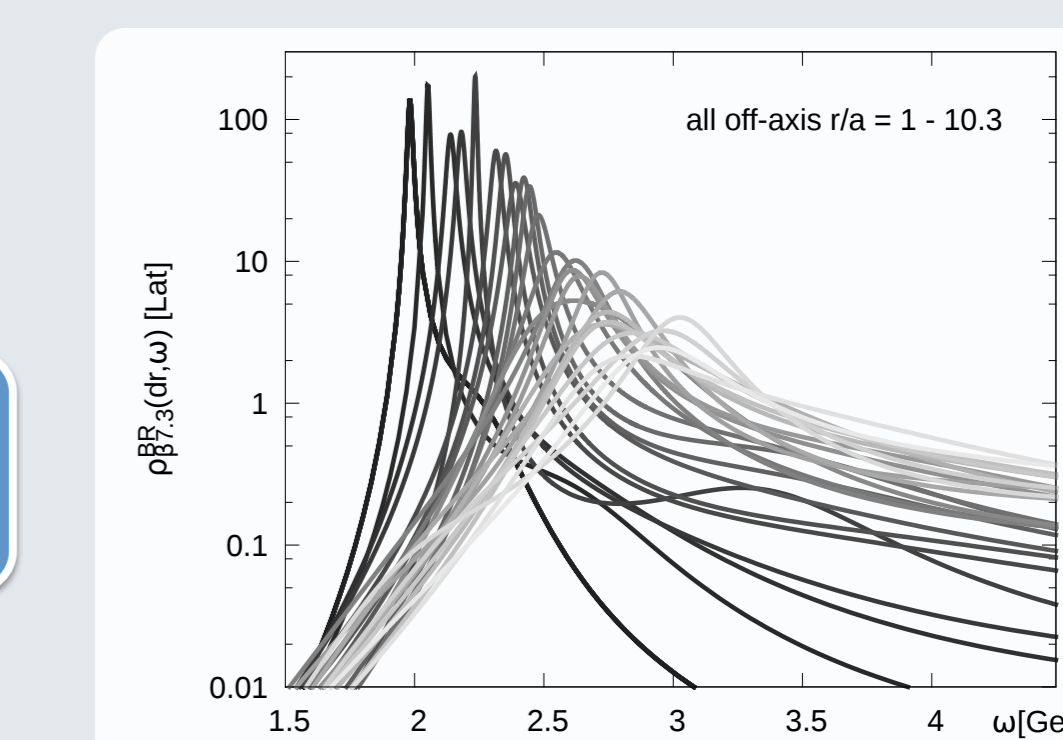
$T_c=174(3)\text{MeV}$

$m_\tau \sim 300\text{MeV}$

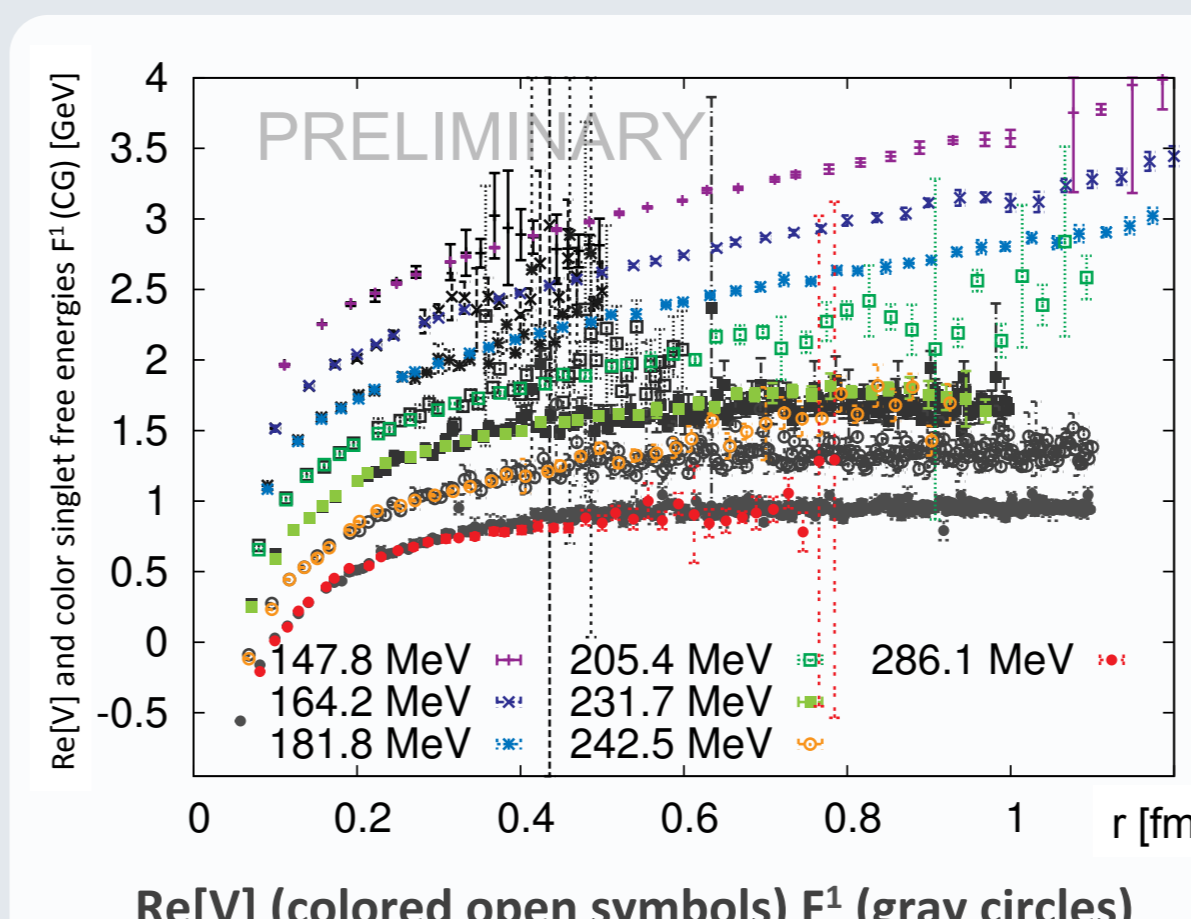


Bayesian spectral reconstruction

$N_\omega=4600, I_\omega^{\text{num}}=[-11, 12]$
 $\tau^{\text{max}}_{\text{num}}=20, m(\omega)=\text{const.}$
512bits prec., $\Delta^{\text{min}}=10^{-60}$



■ Wilson line correlators above T_c show increased curvature with T and R : spectral width present, thus $\text{Im}[V] \neq 0$



■ Investigate $\text{Re}[V]$ only, since $N_t=12$ does not allow quantitative determination of spectral widths

■ Since we use all available datapoints: stable reconstruction of linear rising part of $\text{Re}[V]$ at low T possible (F^1 too noisy)

■ Clear change from confinement to Debye screening visible, **Re[V] lies close to the color singlet free energies**

■ Quantitative testing in progress: Limitations of reconstruction accuracy at different T

Conclusion

■ We showed how the effective field theory based definition of the real-time in-medium heavy quark potential V^{QQ} can be used to express its values in terms of spectral functions.

■ A novel Bayesian reconstruction method for spectral functions enables the controlled extraction of $\text{Re}[V^{\text{QQ}}]$ and $\text{Im}[V^{\text{QQ}}]$ from non-perturbative and first-principles lattice QCD simulations in Euclidean time.

■ Our study on quenched and fully dynamical lattices shows that the **real-part** of the in-medium potential lies **close to the color singlet free energies** in Coulomb gauge. The extraction of the **imaginary part** reveals that its values are of the **same order of magnitude as the predictions from resummed HTL perturbation theory** down to close above the phase transition.

■ Y.B. is supported by SNF grant PZ00P2-142524 and A.R. acknowledges funding from the DFG