

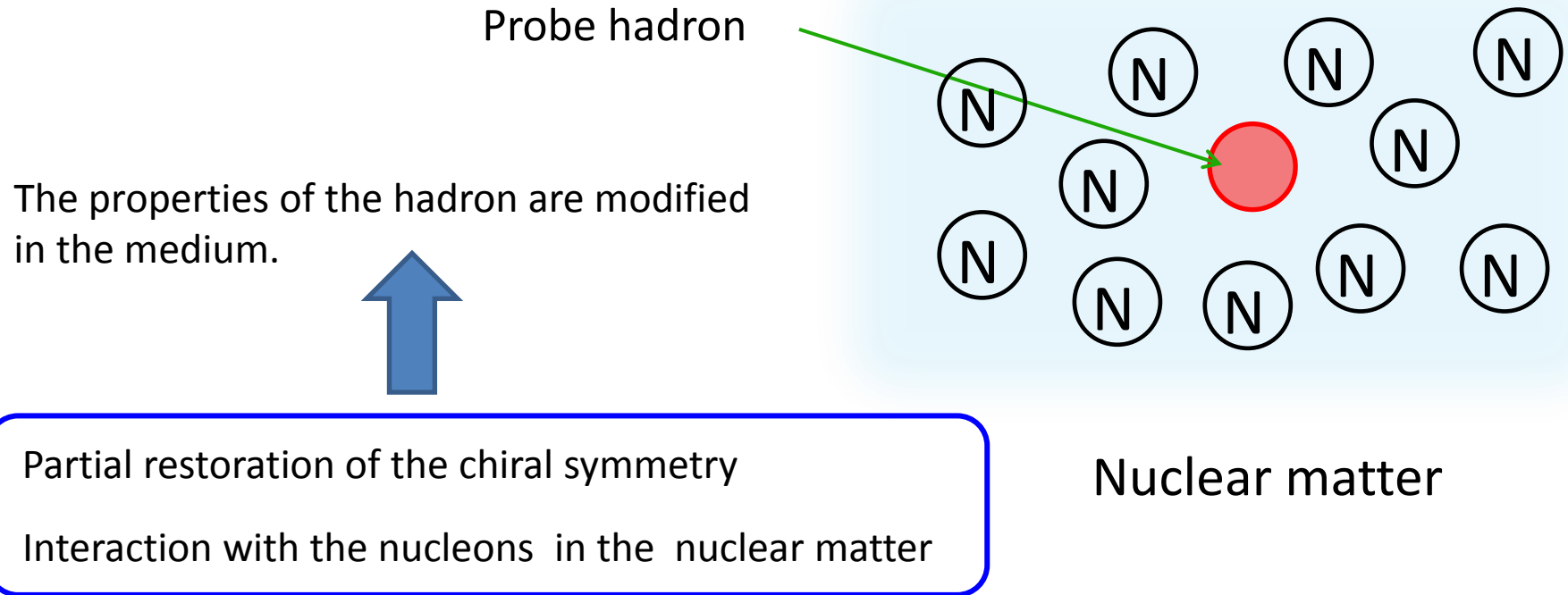
# An analysis of the nucleon spectral function in the nuclear medium from QCD sum rule

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# Introduction

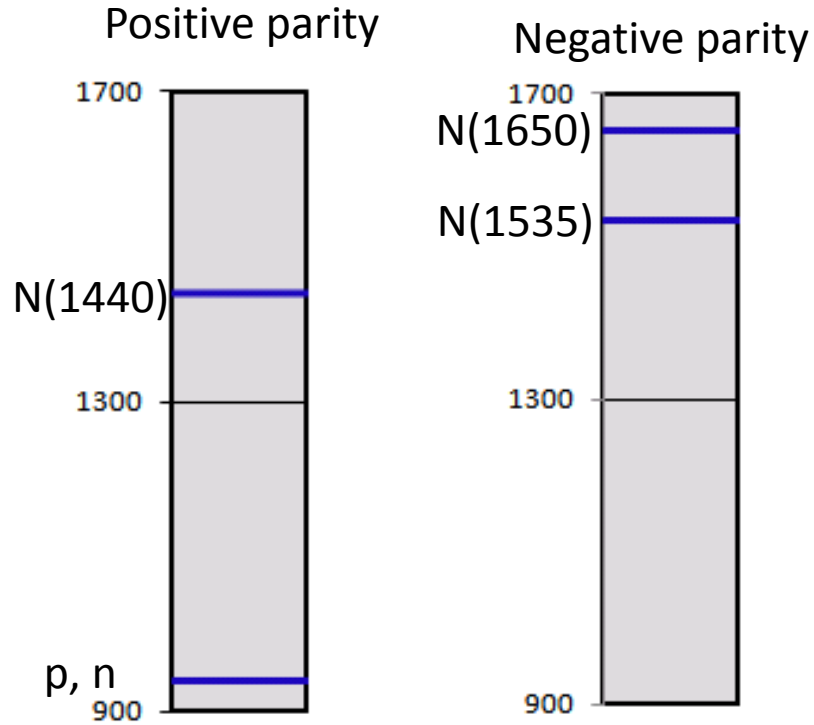
- Hadron properties in the nuclear medium



We focus on the nucleon ground state and excited state.

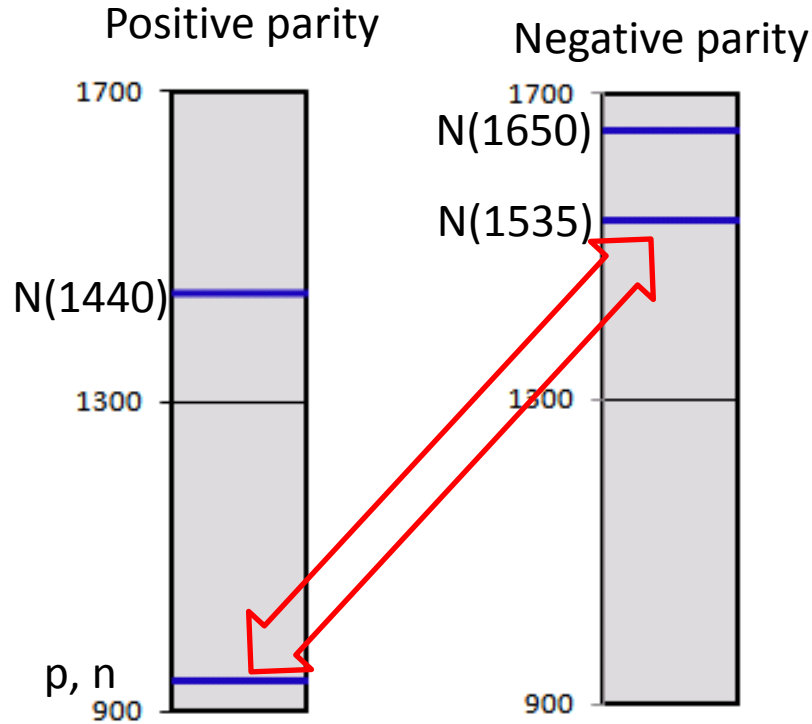
# Introduction

Mass spectrum of the nucleons



# Introduction

## Mass spectrum of the nucleons

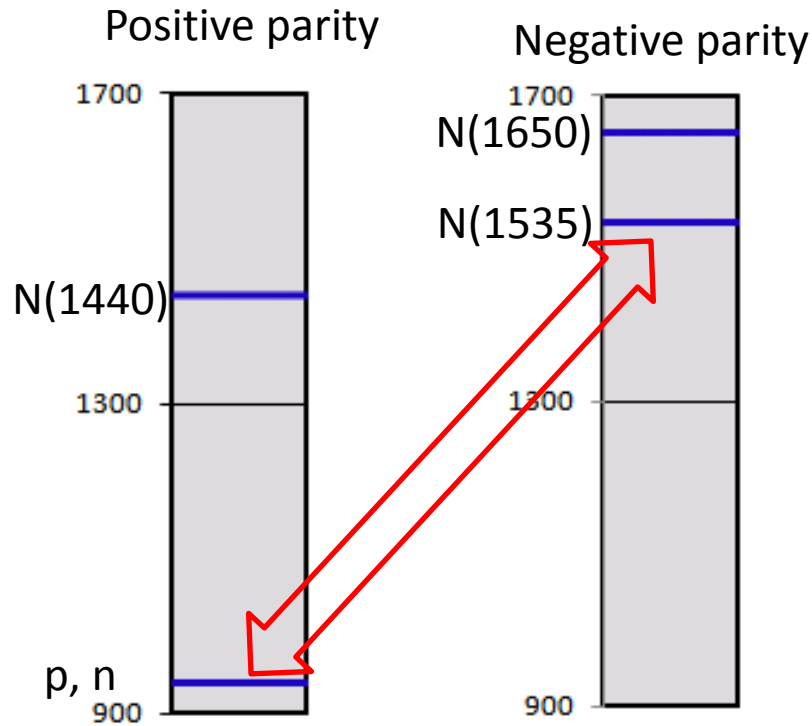


- The mass difference between nucleon ground state and N(1535) is about 600 MeV.
- It is predicted that Chiral symmetry breaking cause these difference.

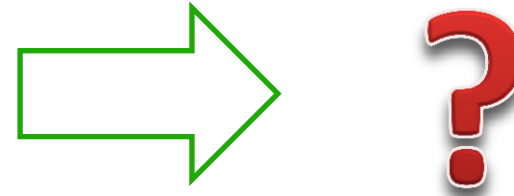
When chiral symmetry is restored, the mass spectrum will change.

# Introduction

Mass spectrum of the nucleons

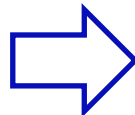


In the nuclear matter



When chiral symmetry is restored, the mass spectrum will change.

To investigate these properties from QCD, non perturbative method is needed.

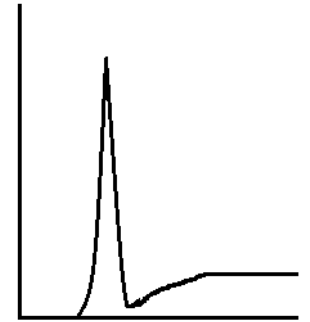


Analysis of QCD sum rule in nuclear matter

# Nucleon QCD sum rule

$$\begin{aligned}\Pi(q) &\equiv i \int e^{iqx} \langle 0 | T[\eta(x) \bar{\eta}(0)] | 0 \rangle d^4x \\ &= \int_0^\infty \frac{1}{\pi} \frac{\text{Im}\Pi(t)}{t - q^2} dt = \int_0^\infty \frac{\rho(t)}{t - q^2} dt\end{aligned}$$

Hadronic spectral function



is calculated by the operator product expansion (OPE)

Non perturbative contributions are expressed by some Condensates.

➡  $\langle \bar{q}q \rangle, \langle \frac{\alpha_s}{\pi} G^2 \rangle \dots$

An order parameter of chiral symmetry

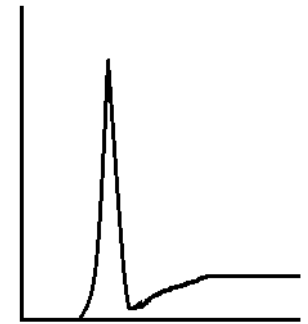
We apply this method to the analyses in the nuclear matter.

# Nucleon QCD sum rule

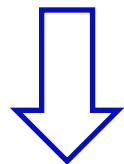
$$\Pi(q) \equiv i \int e^{iqx} \langle 0 | T[\eta(x) \bar{\eta}(0)] | 0 \rangle d^4x$$

Hadronic spectral function

$$= \int_0^\infty \frac{1}{\pi} \frac{\text{Im}\Pi(t)}{t - q^2} dt = \int_0^\infty \frac{\rho(t)}{t - q^2} dt$$



is calculated by the operator product expansion (OPE)



Application for the analyses in nuclear matter

$$i \int e^{iqx} \langle 0 | T[\eta(x) \bar{\eta}(0)] | 0 \rangle d^4x \quad \Rightarrow \quad i \int d^4x e^{iqx} \langle \underline{\Psi}_0 | T[\eta(x) \bar{\eta}(0)] | \underline{\Psi}_0 \rangle$$

Modification:  $\langle 0 | O_i | 0 \rangle \quad \Rightarrow \quad \langle \Psi_0 | O_i | \Psi_0 \rangle \quad \Psi_0 : \text{Ground state of nuclear matter}$

Chiral condensate:  $\langle \bar{q}q \rangle_0 \quad \Rightarrow \quad \langle \bar{q}q \rangle_\rho = \langle \bar{q}q \rangle_0 + \frac{\sigma_N}{2m_q} \rho + \dots$

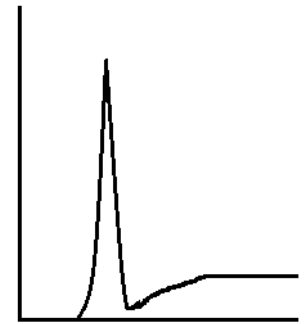
New condensate:  $\quad \Rightarrow \quad \langle q^\dagger q \rangle_\rho = \frac{3}{2} \rho$

# Nucleon QCD sum rule

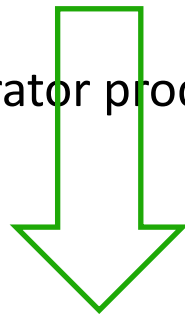
$$\Pi(q) \equiv i \int e^{iqx} \langle 0 | T[\eta(x) \bar{\eta}(0)] | 0 \rangle d^4x$$

Hadronic spectral function

$$= \int_0^\infty \frac{1}{\pi} \frac{\text{Im}\Pi(t)}{t - q^2} dt = \int_0^\infty \frac{\rho(t)}{t - q^2} dt$$



is calculated by the operator product expansion (OPE)



Gaussian sum rule

$$G(s, \tau) = \int_0^\infty \rho(\omega) \frac{1}{\sqrt{4\pi\tau}} \exp\left(-\frac{(\omega^2 - s)^2}{4\tau}\right) d\omega$$

is calculated by OPE

$\tau, s$ : parameter



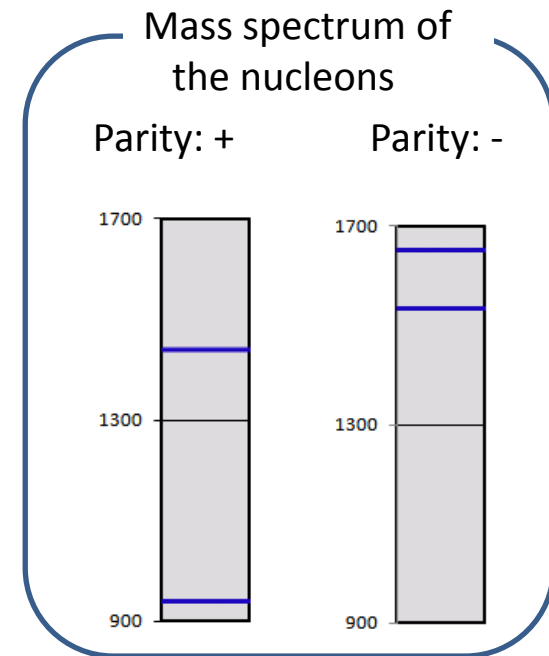
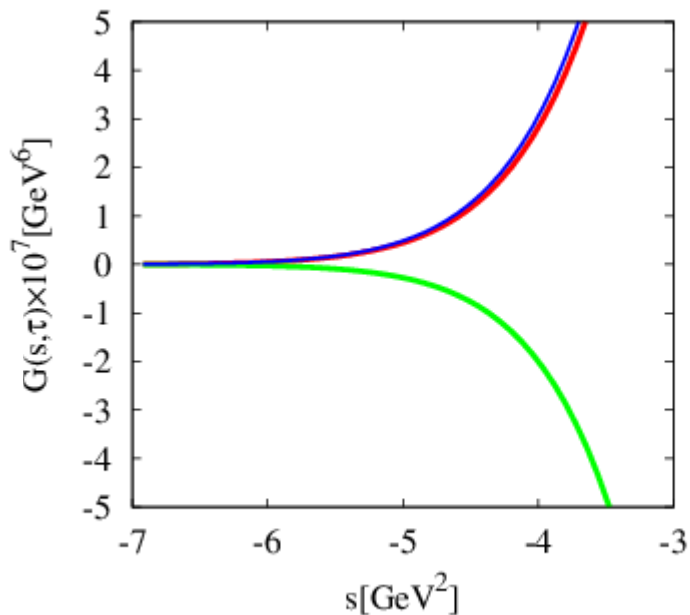
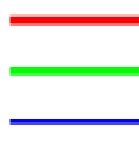
# Nucleon QCD sum rule

The behavior of the OPE data in the vacuum

$$G^{\oplus}(s, \tau) = \left[ C_0(s, \tau) + C_4(s, \tau) \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle + C_6(s, \tau) \langle \bar{q}q \rangle^2 \right]$$

$$\pm \left[ C_3(s, \tau) \langle \bar{q}q \rangle + C_5(s, \tau) \langle \bar{q}g\sigma \cdot Gq \rangle + C_7(s, \tau) \langle \bar{q}q \rangle \left\langle \frac{\alpha_s}{\pi} G^2 \right\rangle \right]$$

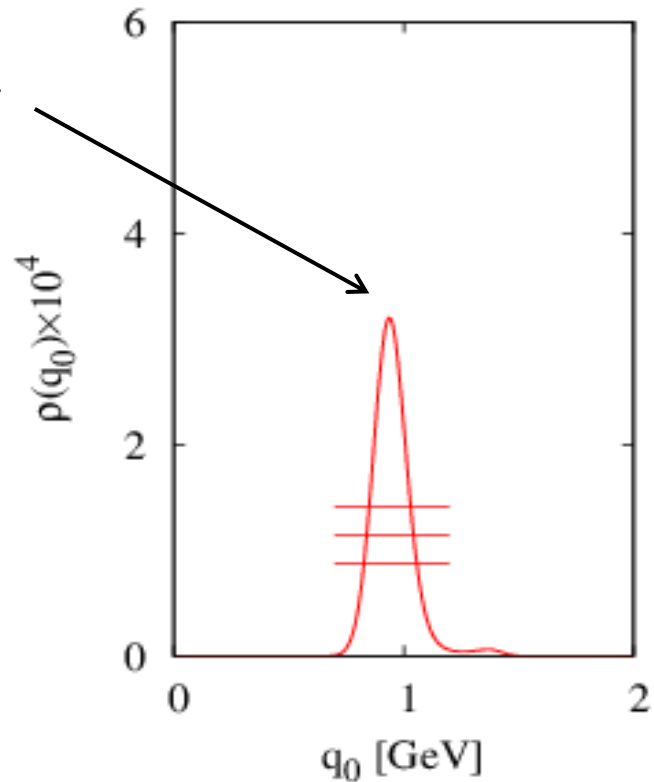
Positive parity OPE  
 Negative parity OPE  
 $\langle \bar{q}q \rangle$  term



The difference between positive parity and negative parity is mainly caused by chiral condensate term.

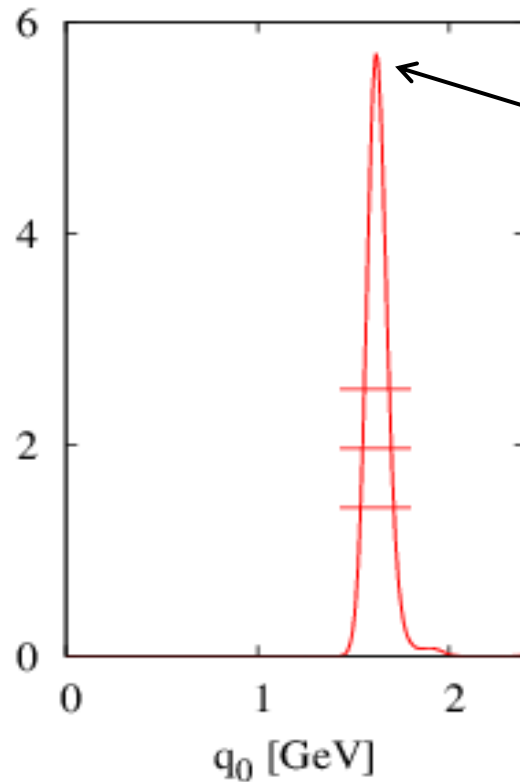
# Nucleon QCD sum rule

Positive parity



930MeV

Negative parity



1620 MeV

In both positive and negative parity, the peaks are found.

In the negative parity analysis, the peak correspond to the N(1535) or (and) N(1650).

# Nucleon QCD sum rule

The behavior of the OPE data in the nuclear matter

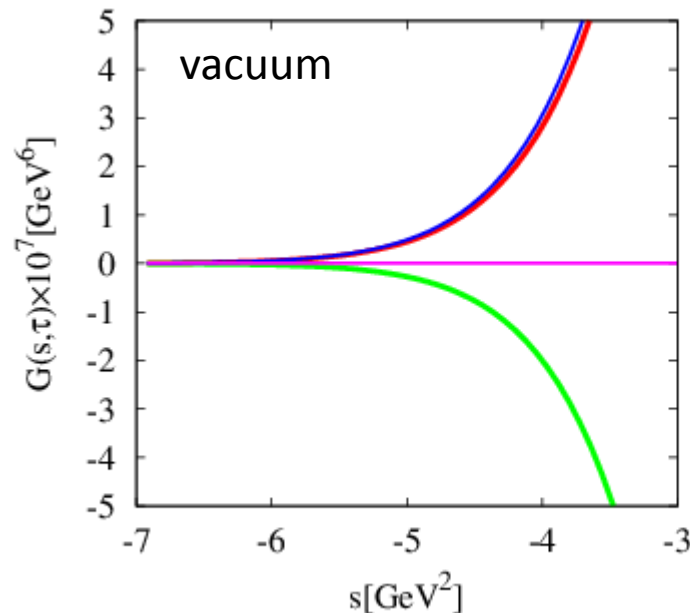
Vacuum

$$\Pi(q) \equiv i \int e^{iqx} \langle 0 | T[\eta(x) \bar{\eta}(0)] | 0 \rangle d^4x$$

Nuclear matter

$$i \int d^4x e^{iqx} \langle \Psi_0 | T[\eta(x) \bar{\eta}(0)] | \Psi_0 \rangle$$

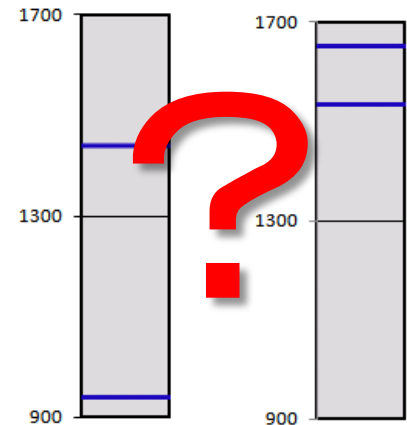
- Positive parity OPE
  - Negative parity OPE
  - - -  $\langle \bar{q}q \rangle_\rho$  term
  - - -  $\langle q^\dagger q \rangle_\rho$  term
- $\rho_N$ : nuclear matter density



Mass spectrum of the nucleons

Parity: +

Parity: -



Positive parity: OPE data decreases.

Positive parity:  $+ C_1 \langle \bar{q}q \rangle_\rho + C_2 \langle q^\dagger q \rangle_\rho$

Negative parity: OPE data does not decrease

Negative parity:  $- C_1 \langle \bar{q}q \rangle_\rho + C_2 \langle q^\dagger q \rangle_\rho$

# Nucleon QCD sum rule

The behavior of the OPE data in the nuclear matter

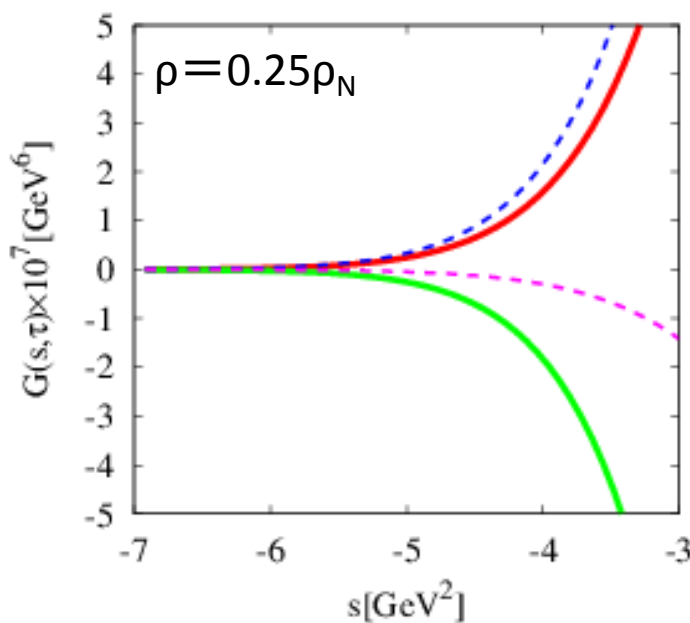
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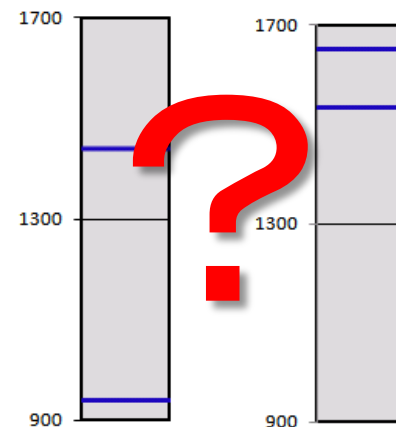
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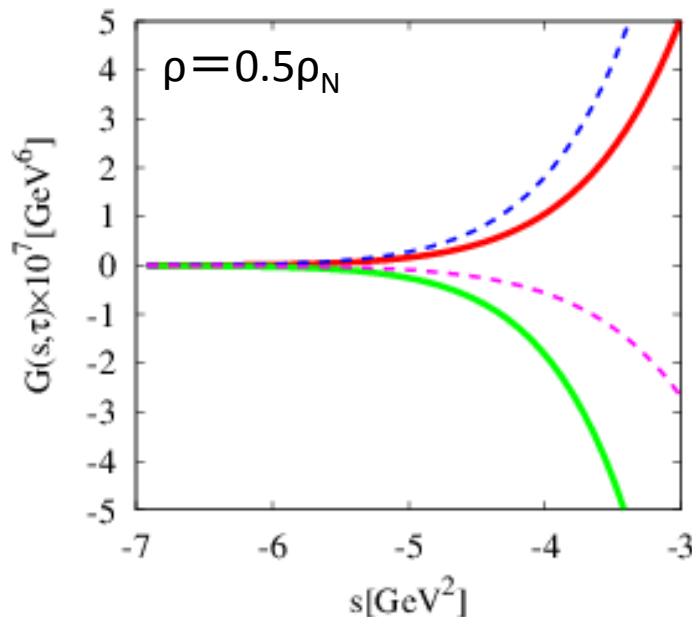
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Nuclear matter

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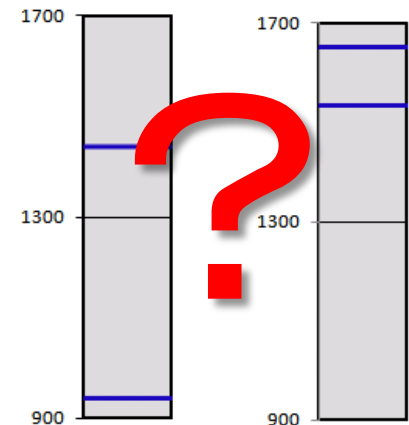
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# Nucleon QCD sum rule

The behavior of the OPE data in the nuclear matter

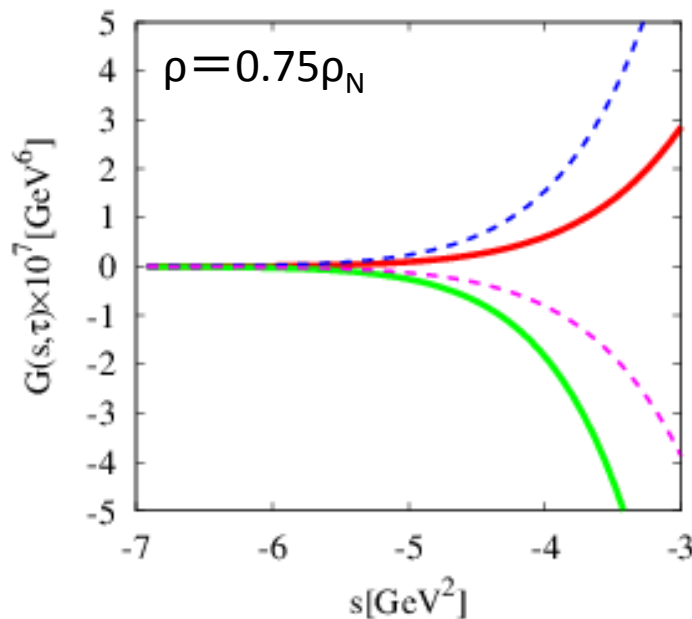
Vacuum

$$\Pi(q) \equiv i \int e^{iqx} \langle 0 | T[\eta(x) \bar{\eta}(0)] | 0 \rangle d^4x$$

Nuclear matter

$$i \int d^4x e^{iqx} \langle \Psi_0 | T[\eta(x) \bar{\eta}(0)] | \Psi_0 \rangle$$

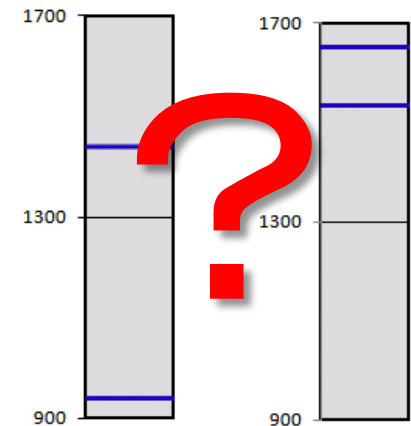
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Mass spectrum of the nucleons

Parity: +

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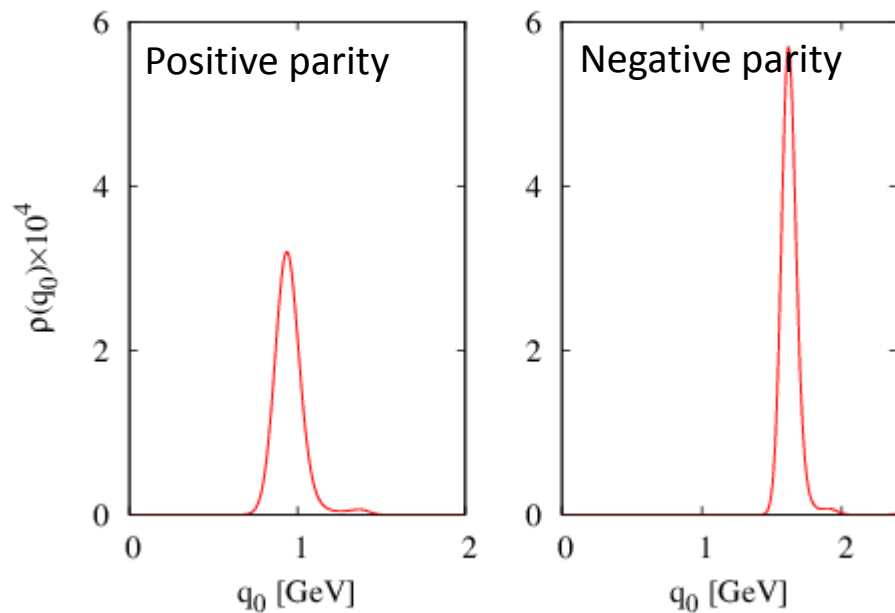
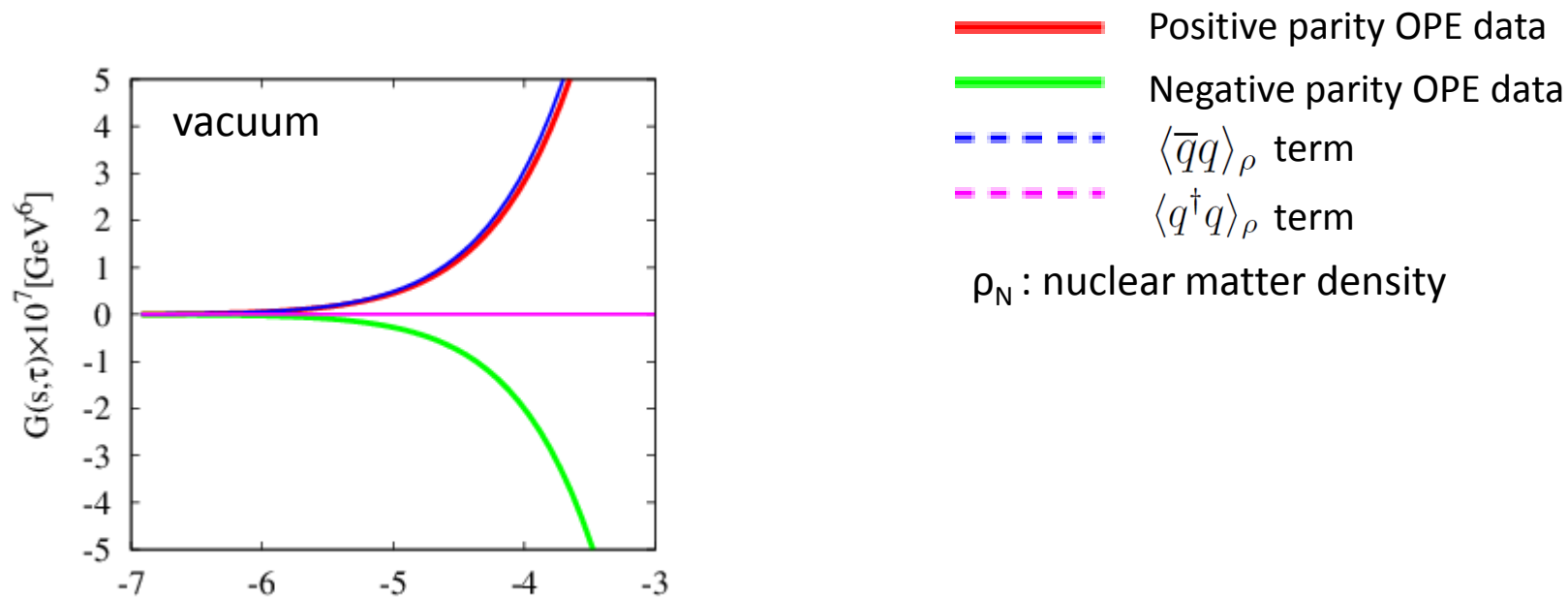
Positive parity: OPE data decreases.

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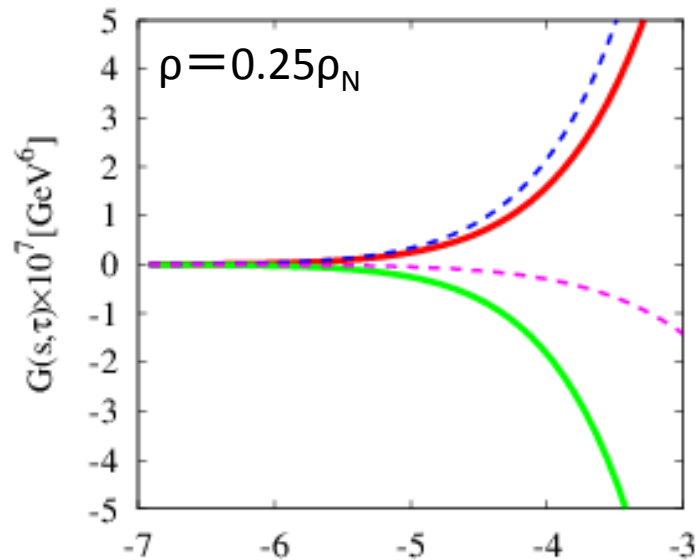
Negative parity:  $- C_1 \langle \bar{q}q \rangle_\rho + C_2 \langle q^\dagger q \rangle_\rho$

# Nucleon QCD sum rule

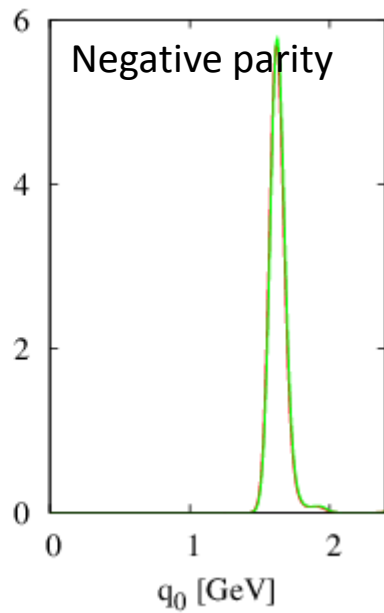
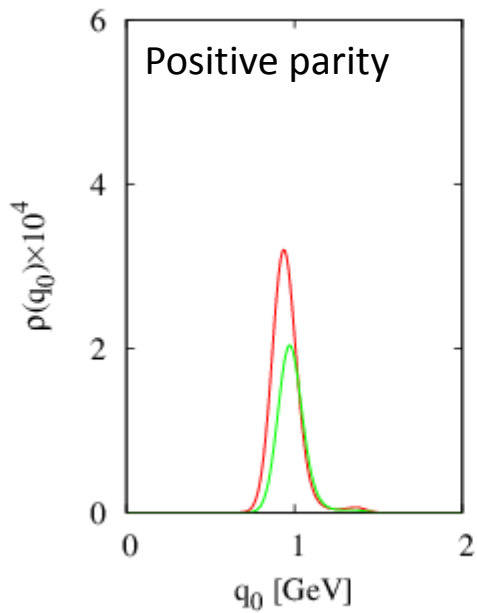


— : Vacuum

# Nucleon QCD sum rule



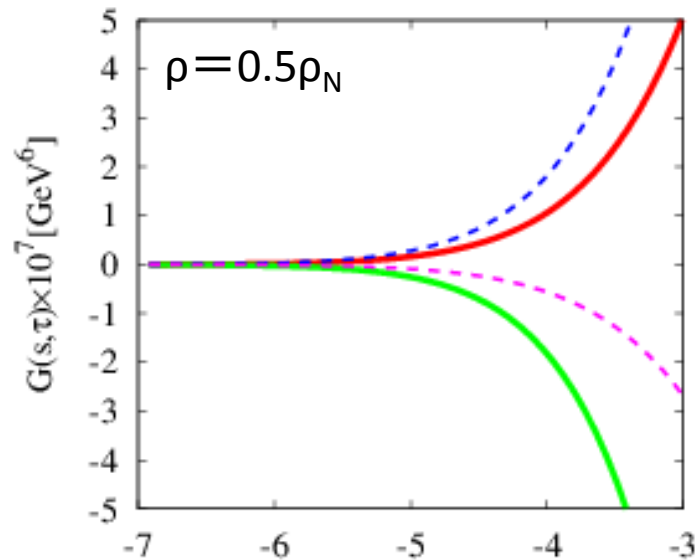
- Positive parity OPE data
  - Negative parity OPE data
  - - -  $\langle \bar{q}q \rangle_\rho$  term
  - - -  $\langle q^\dagger q \rangle_\rho$  term
- $\rho_N$  : nuclear matter density



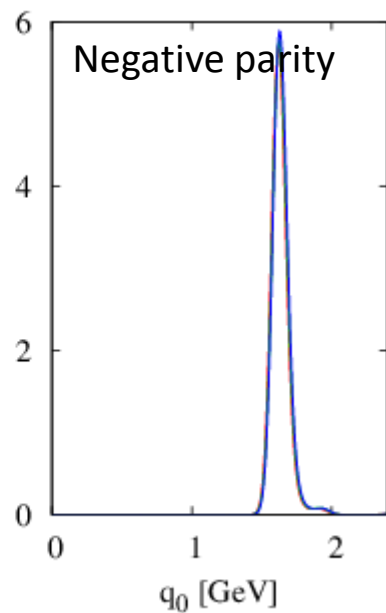
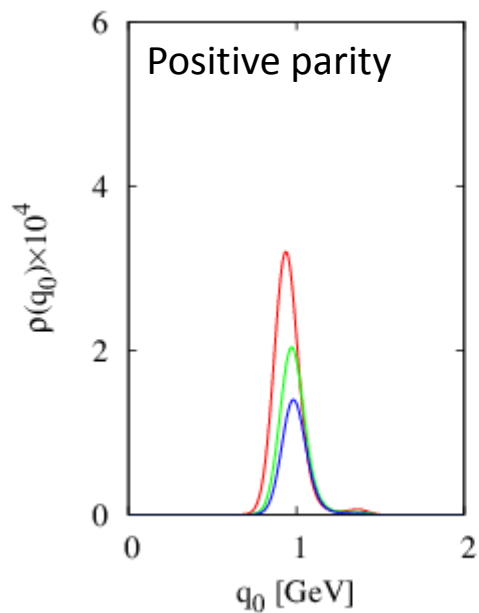
- : Vacuum
- :  $\rho = 0.25\rho_0$



# Nucleon QCD sum rule

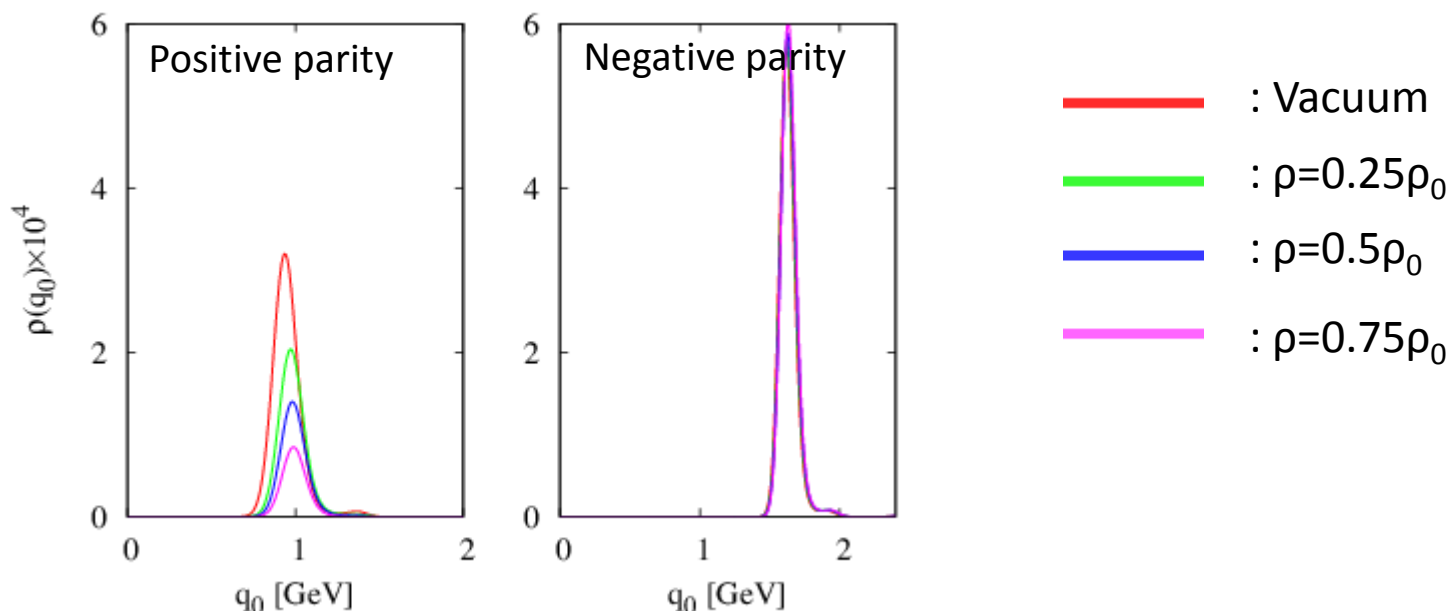
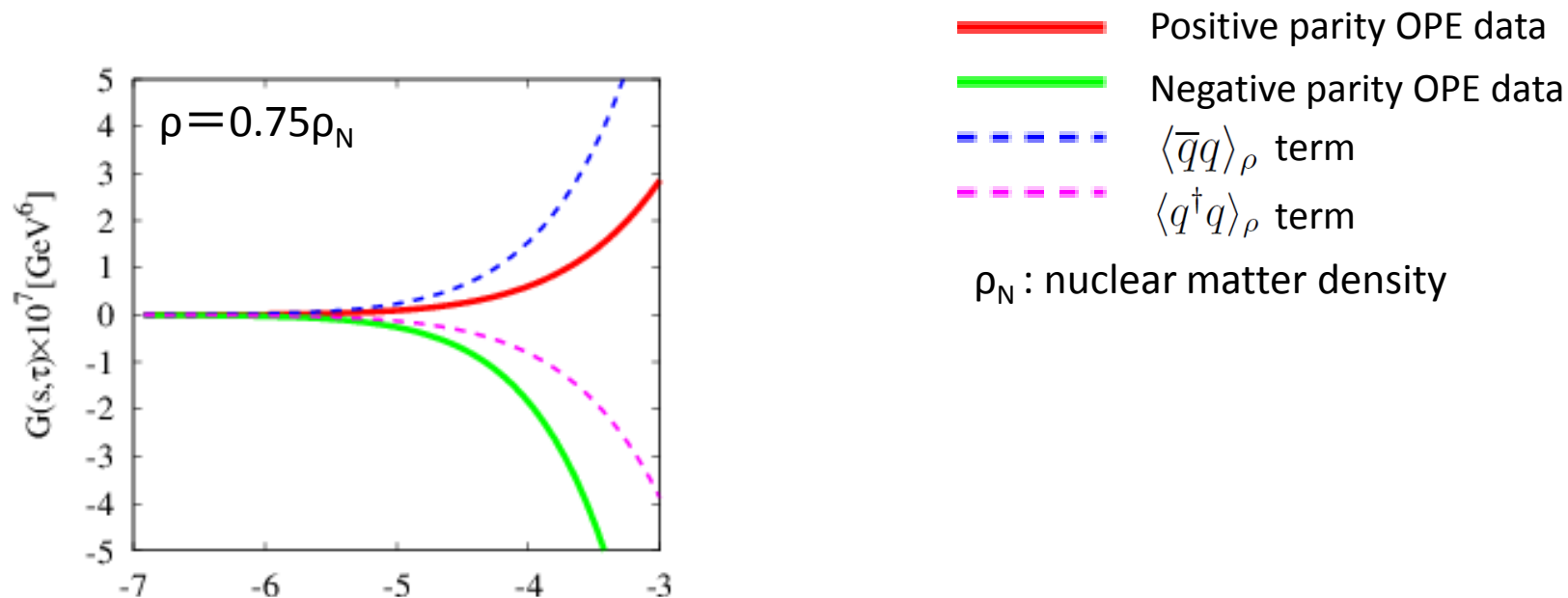


- Positive parity OPE data
  - Negative parity OPE data
  - - -  $\langle \bar{q}q \rangle_\rho$  term
  - - -  $\langle q^\dagger q \rangle_\rho$  term
- $\rho_N$  : nuclear matter density



- : Vacuum
- :  $\rho = 0.25\rho_0$
- :  $\rho = 0.5\rho_0$

# Nucleon QCD sum rule



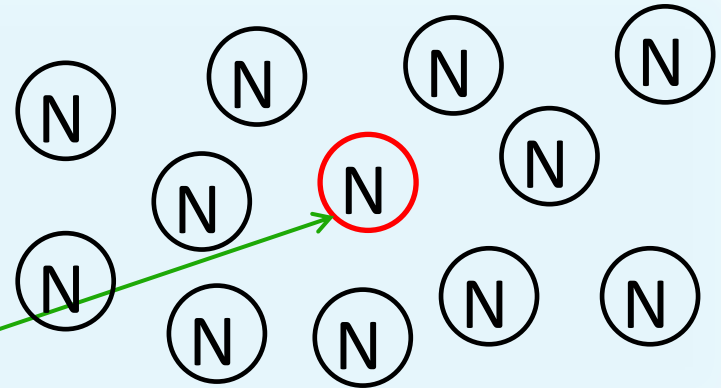
# Nucleon QCD sum rule

Vacuum

$$\Pi(q) = i \int d^4x e^{iqx} \langle \underline{0} | T[\eta(x) \bar{\eta}(0)] | \underline{0} \rangle$$

Nuclear matter

$$\Pi(q) = i \int d^4x e^{iqx} \langle \underline{\Psi_0} | T[\eta(x) \bar{\eta}(0)] | \underline{\Psi_0} \rangle$$



Probe nucleon

Propagator:  $\frac{1}{\not{q} - M + i\epsilon}$

Propagator:  $\frac{1}{\not{q} - M - \Sigma(q) + i\epsilon}$

$$\frac{\not{q} - \not{v}\Sigma_v + M^*}{(q_0 - E + i\epsilon)(q_0 + \bar{E} - i\epsilon)}$$

Pole of positive energy state:  $E = \sqrt{M^{*2} + \vec{q}^2} + \Sigma_v$

Effective mass:  $M^* = M + \Sigma_s$

Pole of negative energy state:  $-\bar{E} = -\sqrt{M^{*2} + \vec{q}^2} + \Sigma_v$

# Nucleon QCD sum rule

Investigation of  $M_{0\pm}^*, \Sigma_{0\pm}^v$

$$\Pi(q) = i \int d^4x e^{iqx} \theta(x_0) \langle \Psi_0 | T[\eta(x) \bar{\eta}(0)] | \Psi_0 \rangle \quad E = \sqrt{q^2 + M^{*2}} + \Sigma^v$$

$$= \not{q} \Pi_1(q^2, q \cdot u) + \Pi_2(q^2, q \cdot u) + \not{u} \Pi_u(q^2, q \cdot u)$$

$$= \sum_n |\lambda_{n+}|^2 \left( \frac{(\sqrt{M^{*2} + \bar{q}^2} \gamma_0 + M^*)}{2\sqrt{M^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^+ + i\epsilon} + \frac{\gamma_i q^i}{2\sqrt{M^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^+ + i\epsilon} \right)$$

+ (contribution of negative parity states)

$$q_0 \Pi_1 + \Pi_2 \quad \Rightarrow \quad \sum |\lambda_+|^2 \frac{E^+ + M_+^*}{2\sqrt{M_+^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^+ + i\epsilon} + |\lambda_-|^2 \frac{E^- - M_-^*}{2\sqrt{M_-^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^- + i\epsilon},$$

$$q_0 \Pi_1 - \Pi_2 \quad \Rightarrow \quad \sum |\lambda_+|^2 \frac{E^+ - M_+^*}{2\sqrt{M_+^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^+ + i\epsilon} + |\lambda_-|^2 \frac{E^- + M_-^*}{2\sqrt{M_-^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^- + i\epsilon},$$

$$\Pi_u \quad \Rightarrow \quad \sum |\lambda_+|^2 \frac{-\Sigma_+^v}{2\sqrt{M_+^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^+ + i\epsilon} + |\lambda_-|^2 \frac{-\Sigma_-^v}{2\sqrt{M_-^{*2} + \bar{q}^2}} \frac{1}{q_0 - E^- + i\epsilon}.$$

By fitting the phenomenological side and OPE side, we can investigate  $M_{0\pm}^*, \Sigma_{0\pm}^v$ .

# Nucleon QCD sum rule

$\rho_N$ : nuclear matter density

$n_0$ : nuclear matter density

	Vacuum	$n=0.25n_0$	$n=0.5n_0$	$n=0.75n_0$	
Positive parity	$M_{0+}^*$	930	850	710	470
	$\Sigma_{0+}^v$	0	120	270	500
Negative parity	$M_{0-}^*$	1620	1630	1650	1680
	$\Sigma_{0-}^v$	0	0	-20	-50

# Summary

- We analyze the nucleon spectral function by using QCD sum rules with MEM
- We find that the difference between the positive and negative parity spectral function is mainly caused by the chiral condensate.
- The information of not only the ground state but also the negative parity excited state is extracted
- We apply this method to the analyses in nuclear medium and investigate the effective masses and the vector self-energies.