pQCD predictions for the LHC

Tools for multijet phenomenology

Steffen Schumann



II. Physikalisches Institut, Universität Göttingen

LHC Days in Split 02/10/14



pQCD challenges at the LHC

ATLAS $Z(
ightarrow e^+e^-/\mu^+\mu^-)$ +jets analysis [JHEP 1307 (2013) 032]



inclusive jet rates

- large jet multiplicities
- wide range of kinematics/scales
 - hard Z associated with jets [requires FO calculations]
 - high-p_T jets associated by Z [expect large logarithmic corrections]
- \rightsquigarrow high-multiplicity FO calculations
- $\rightsquigarrow\,$ parton-shower resummation important
- $\rightsquigarrow\,$ combining different jet-multi final states

A B > A B >

pQCD challenges at the LHC

ATLAS $Z(
ightarrow e^+e^-/\mu^+\mu^-)$ +jets analysis [JHEP 1307 (2013) 032]



inclusive xsec ratios

- large jet multiplicities
- wide range of kinematics/scales
 - hard Z associated with jets [requires FO calculations]
 - high-p_T jets associated by Z [expect large logarithmic corrections]
- \rightsquigarrow high-multiplicity FO calculations
- $\rightsquigarrow\,$ parton-shower resummation important
- $\rightsquigarrow\,$ combining different jet-multi final states

同 ト イ ヨ ト イ ヨ ト

 \rightsquigarrow interesting patterns in xsec ratios

- NLO QCD predictions the new standard
- ME+PS matching towards automation at NLO
- Applications scaling patterns, jet-rate extrapolation

(同) (ヨ) (ヨ)

э

How to address an LHC measurement with pQCD?

the partonic cross section

$$rac{d\sigma}{dX} \propto \sum_{ab,n} \int \,\mathrm{d} x_a \mathrm{d} x_b \mathrm{d} \Phi_n(p_1,\ldots,p_n) \; f_a(x_a,\mu_F^2) f_b(x_b,\mu_F^2) \left|\mathcal{M}_{ab
ightarrow n}
ight|^2 \;
ho_n(p_1,\ldots,p_n)$$

- Φ_n is the *n*-particle phase space, $f_{a,b}$ the proton PDFs
- $\mathcal{M}_{2 \rightarrow n}$ invariant matrix element [approx. for large n]
- measurement function ρ_n projects out observable

$$\rho_n = \begin{cases} 1 & : \text{ total cross section} \\ \delta(X - \chi_n(p_1, \dots, p_n)) & : \text{ differential cross section} \end{cases}$$

 \sim implements acceptance cuts, observable definition \sim evaluated in fixed-order perturbation theory [hard process] \sim semi-classical approximations for large *n* [parton shower, resummation] \sim model parton-hadron transition, remnant interactions [hadronization, underlying event]

Hard Processes at Next-to-Leading Order QCD

Anatomy of NLO QCD calculations [in dim. regularization $d = 4 - 2\epsilon$]

• (UV renormalized) virtual-corrections $\sigma^V \sim \text{IR divergent}$

• real-emission $\sigma^R \rightarrow \mathsf{IR}$ divergent

 \rightsquigarrow for IR safe observables sum is finite

3

Hard Processes at Next-to-Leading Order QCD

Anatomy of NLO QCD calculations [in dim. regularization $d = 4 - 2\epsilon$]

• (UV renormalized) virtual-corrections $\sigma^V \sim \text{IR divergent}$

• real-emission $\sigma^R \sim \text{IR divergent}$

 \rightsquigarrow for IR safe observables sum is finite

Dipole subtraction method [Catani, Seymour Nucl. Phys. B 485 (1997) 291]

$$\sigma_{2 \to n}^{NLO} = \int_{n} \left[\mathbf{d}^{(4)} \sigma^{B} + \int_{\text{loop}} \mathbf{d}^{(d)} \sigma^{V} + \int_{\mathbf{1}} \mathbf{d}^{(d)} \sigma^{A} \right]_{\epsilon=0} + \int_{n+1} \left[\mathbf{d}^{(4)} \sigma^{R} - \mathbf{d}^{(4)} \sigma^{A} \right]$$

• subtraction terms yield local approximation for the real emission process

• describe the amplitude in the soft & collinear limits $[1/\epsilon \text{ and } 1/\epsilon^2 \text{ poles}]$

$$\int_{n+1} \mathrm{d}^{(d)} \sigma^{\mathrm{A}} = \sum_{\substack{\mathrm{dipoles} \\ \mathrm{spin-\& color correlations}}} \int_{n} \mathrm{d}^{(d)} \sigma^{\mathrm{B}} \otimes \int_{1} \mathrm{d}^{(d)} V_{\mathrm{dipole}} \xrightarrow{} \mathrm{universal dipole terms}$$

Hard Processes at Next-to-Leading Order QCD

The emerging picture: a fully differential NLO calculation

$$\sigma_{2 \to n}^{NLO} = \int_{n+1} \left[\mathrm{d}^{(4)} \sigma^{\mathrm{R}} - \mathrm{d}^{(4)} \sigma^{\mathrm{A}} \right] + \int_{n} \left[\mathrm{d}^{(4)} \sigma^{B} + \int_{\mathrm{loop}} \mathrm{d}^{(d)} \sigma^{\mathrm{V}} + \int_{1} \mathrm{d}^{(d)} \sigma^{\mathrm{A}} \right]_{\epsilon=0}$$

Monte-Carlo codes

- all the tree-level bits
- subtraction of singularities
- efficient phase-space integration

One-Loop codes

- Loop amplitudes, *i.e.* $2\Re(\mathcal{A}_V \mathcal{A}_B^{\dagger})$
- Loop integration

 $\rightsquigarrow~1/\epsilon,~1/\epsilon^2$ coefficients & finite terms

some recent NLO calculations:

2009
$$W$$
 + 3jets, $t\overline{t}$ + 1jet

2010 W + 4jets, Z + 3jets

- 2011 Z + 4jets, $t\overline{t}$ + 2jets, 4jets
- 2012 γ + 3jets
- 2013 W + 5jets, 5jets

2014 $\gamma\gamma$ + 3jets

OL tools & names:

- BlackHat: Bern et al.
- HelacNLO: Bevilacqua et al.
- OpenLoops: Pozzorini et al.
- GoSam: Cullen et al.
- NJET: Biedermann et al.
- MadLoops: Hirschi et al.

NLO QCD predictions: W + 5 jets



同 ト イ ヨ ト イ ヨ ト

э

NLO QCD predictions: W + 5 jets



first calculation completed recently

- BLACKHAT+SHERPA: Bern et al. [Phys. Rev. D 88 (2013) 1, 014025]
 - BLACKHAT: on-shell methods for one-loop amplitudes [arXiv:0808.0941]
 - SHERPA: dipole subtraction, real-emission, phase space, stearing
 - $\,\hookrightarrow\,$ fully differential partonic event generator with NLO accuracy

NLO QCD calculations: W + 5 jets

BLACKHAT+SHERPA: 7 TeV LHC predictions [Phys. Rev. D 88 (2013) 1, 014025]

• consider anti- k_t jets with $p_T^{\rm jet} > 25$ GeV & R=0.5

process	W ⁻ – LO	$W^ NLO$	W ⁺ – LO	$W^+ - NLO$
xsec [pb]	$1.076(0.003)^{+0.985}_{-0.480}$	$0.77(0.02)^{+0.07}_{-0.19}$	$2.005(0.006)^{+1.815}_{-0.888}$	$1.45(0.04)^{+0.12}_{-0.34}$



 $\begin{array}{l} \underline{\text{central scale}} \\ \mu_R = \mu_F = \mu = \hat{H}'_T/2 \\ \hat{H}'_T \equiv \sum_i \rho^{\rm jet}_{T,i} + \sqrt{M^2_W + \rho^2_{T,W}} \end{array}$

 $\frac{\text{scale variations}}{\mu/2, \mu/\sqrt{2}, \mu, \sqrt{2}\mu, 2\mu}$



Approximating multi-parton production

The QCD Parton Shower picture



- construct explicitely the initial- & final-state partons history/fate
- successive branching of incoming and outgoing legs

 \sim exclusive partonic final states with $\mathcal{P}_{\rm tot}=\mathcal{P}_{\rm hard}\cdot\mathcal{P}_{\rm IS}\cdot\mathcal{P}_{\rm FS}$

• evolve parton ensemble from high- to low scale $\mathcal{O}(1 {
m GeV}^2)$

 \rightsquigarrow link the hard process to universal hadronization models

• model intra-jet energy flows: jets become multi-parton objects

Matching exact matrix elements with parton showers

The art of combining matrix elements with parton showers

- model few hardest emissions by exact matrix elements
- account for subsequent splittings, i.e. soft and/or collinear emissions
- avoid any double counting or dead regions of emission phase space
- preserve fixed-order & logarithmic precision of the calculation
- ground-breaking work:
 - multileg tree-level matching: Catani et al. JHEP 0111 (2001) 063
 - NLO + Parton Shower: Frixione, Webber JHEP 0206 (2002) 029

伺 と く き と く き と

Matching exact matrix elements with parton showers

The art of combining matrix elements with parton showers

- model few hardest emissions by exact matrix elements
- account for subsequent splittings, i.e. soft and/or collinear emissions
- avoid any double counting or dead regions of emission phase space
- preserve fixed-order & logarithmic precision of the calculation
- ground-breaking work:
 - multileg tree-level matching: Catani et al. JHEP 0111 (2001) 063
 - NLO + Parton Shower: Frixione, Webber JHEP 0206 (2002) 029



Matching exact matrix elements with parton showers

The art of combining matrix elements with parton showers

- model few hardest emissions by exact matrix elements
- account for subsequent splittings, i.e. soft and/or collinear emissions
- avoid any double counting or dead regions of emission phase space
- preserve fixed-order & logarithmic precision of the calculation
- ground-breaking work:
 - multileg tree-level matching: Catani et al. JHEP 0111 (2001) 063
 - NLO + Parton Shower: Frixione, Webber JHEP 0206 (2002) 029



ME+PS: pull the rabbit – merging NLO processes

a new standard emerging - merging at NLO

- combination of NLO processes of varying jet multi, matched with shower
- option to add in higher-multiplicity tree-level matrix elements
- \hookrightarrow needs notion of exclusive NLO cross sections, e.g. vetoed showers

MEPS@NLO



available approaches

• merging of MC@NLO-type processes in SHERPA [Höche et al. arXiv:1207.5030]

W + 0, 1, 2 at NLO & $W + \ge 3$ at LO

 aMC@NLO merging approach [Frederix, Frixione arXiv:1209.6215]

adding up $\textit{W}+0,1~\mathrm{MC@NLO}$ processes

 CKKW-L at NLO [Lönnblad, Prestel arXiv:1211.7278]

merging W + 0, 1j at NLO & PYTHIA8

 Merging of POWHEG processes [Hamilton et al. arXiv:1212.4504]

combine W/Z + 0, 1 POWHEG generators

ロト ・ 同ト ・ ヨト ・ ヨト

3

ME+PS: merging NLO processes in SHERPA

calculational setup MEPS@NLO [Höche et al. JHEP 1304 (2013) 027]

- W + 0, 1, 2j at NLO matched MC@NLO style \int on the flight merged
- W + 3,4j at LO matched using truncated shower \int into inclusive sample
- CKKW-style scale setting: $\alpha_s(\mu_R^2)^n = \prod \alpha_s(\mu_i^2)$
- factorization & renormalization scale variations $\mu/2$ & 2μ



 \hookrightarrow largely reduced systematics for merged NLO samples

 \hookrightarrow leading-order systematics where dominated by tree-level, i.e. high H_T

consider ratios of successive exclusive jet rates [Gerwick et al. JHEP 1210 (2012) 162]

$$R_{(n+1)/n} = \frac{\sigma_{n+1}^{\text{excl}}}{\sigma_n^{\text{excl}}}$$

- \sim *n* counting (radiated) jets in addition to the core process
- $\rightsquigarrow\,$ ratios quite stable against QCD corrections, reduced systematics
- \sim analytically accessible via resummed jet rates [Gerwick et al. JHEP 1304 (2013) 089]

同 ト イ ヨ ト イ ヨ ト

consider ratios of successive exclusive jet rates [Gerwick et al. JHEP 1210 (2012) 162]

$$R_{(n+1)/n} = \frac{\sigma_{n+1}^{\text{excl}}}{\sigma_n^{\text{excl}}}$$

- \sim *n* counting (radiated) jets in addition to the core process
- $\rightsquigarrow\,$ ratios quite stable against QCD corrections, reduced systematics
- ightarrow analytically accessible via resummed jet rates [Gerwick et al. JHEP 1304 (2013) 089]

Poisson scaling

defined via
$$R_{(n+1)/n} \equiv \frac{\overline{n}}{n+1}$$

- jet rates given by a Poisson dist. $\sigma_n \sim \frac{\bar{n}^n e^{-\bar{n}}}{n!}$
- independent emission picture
- inclusive ratios more cumbersome
- like soft-photon emission in QED
- \hookrightarrow driven by large emission prob.





consider ratios of successive exclusive jet rates [Gerwick et al. JHEP 1210 (2012) 162]

$$R_{(n+1)/n} = \frac{\sigma_{n+1}^{\text{excl}}}{\sigma_n^{\text{excl}}}$$

- \sim *n* counting (radiated) jets in addition to the core process
- $\rightsquigarrow\,$ ratios quite stable against QCD corrections, reduced systematics
- \sim analytically accessible via resummed jet rates [Gerwick et al. JHEP 1304 (2013) 089]

staircase scaling

defined via
$$R_{(n+1)/n} \equiv R = \text{const}$$

 $\sim \sigma_n = \sigma_0 R^n$

- excl. & incl. scale the same • $\frac{\sigma_{n+1}^{\text{incl}}}{\sigma_{n-1}^{\text{incl}}} = \frac{\sigma_{n+1}^{\text{excl}}}{\sigma_{n-1}^{\text{excl}}} = R$
- first mentioned for W/Z + jets

[Ellis, Kleiss, Stirling '85 & Berends, Giele, Kuijf '89]

 \hookrightarrow induced by democratic jet cuts



ATLAS $Z(
ightarrow e^+e^-/\mu^+\mu^-)$ +jets analysis [JHEP 1307 (2013) 032]



 \sim confirms analytic findings based on resummed jet rates \sim PDF/energetic suppression for low-multi bins in democratic selection

[Gerwick et al. JHEP 1210 (2012) 162 & JHEP 1304 (2013) 089]

3

э

Applications: jet-rate extrapolation

Predicting W + 6j with BLACKHAT+SHERPA

[Bern et al. Phys. Rev. D 88 (2013) 1, 014025 & arXiv:1407.6564]



phenomenological fits motivated by scaling arguments:

$$R_{n/(n-1)}^{\text{NLO, }W^{-}} = 0.248 \pm 0.008 - (0.009 \pm 0.002) n$$
$$R_{n/(n-1)}^{\text{NLO, }W^{+}} = 0.263 \pm 0.009 - (0.009 \pm 0.003) n$$

NLO rate estimates

 $W^- + 6 ext{ jets}: ext{ 0.15} \pm 0.01 ext{ pb} ext{ \& } W^+ + 6 ext{ jets}: ext{ 0.30} \pm 0.03 ext{ pb}$

Summary: pQCD predictions for the LHC

Prepare(d) for the LHC challenge

- precise predictions for the Standard Model
 - multileg tree-level & one-loop matrix elements
 - improved parton-shower & matching algorithms
- invoke models for non-perturbative physics
 - capture main features of low-energy dynamics
 - testable/tunable against experimental data

Current activities

- further improving our Standard Model predictions
 - improving the logarithmic accuracy of shower algorithms
 - matching NLO calculations of varying multiplicity with showers
 - attaching parton showers to NNLO matrix elements
- BSM searches in a QCD environment
 - new analysis techniques, e.g. subjets, event shapes
 - NLO matrix-element and shower matching for signals



同 ト イ ヨ ト イ ヨ ト