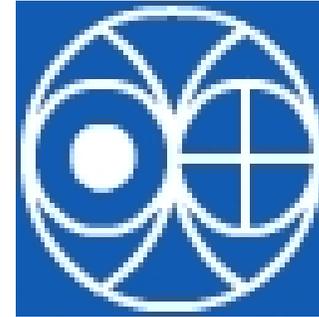


An Overview of Neutrino Physics

Srubabati Goswami

Physical Research Laboratory, Ahmedabad, India

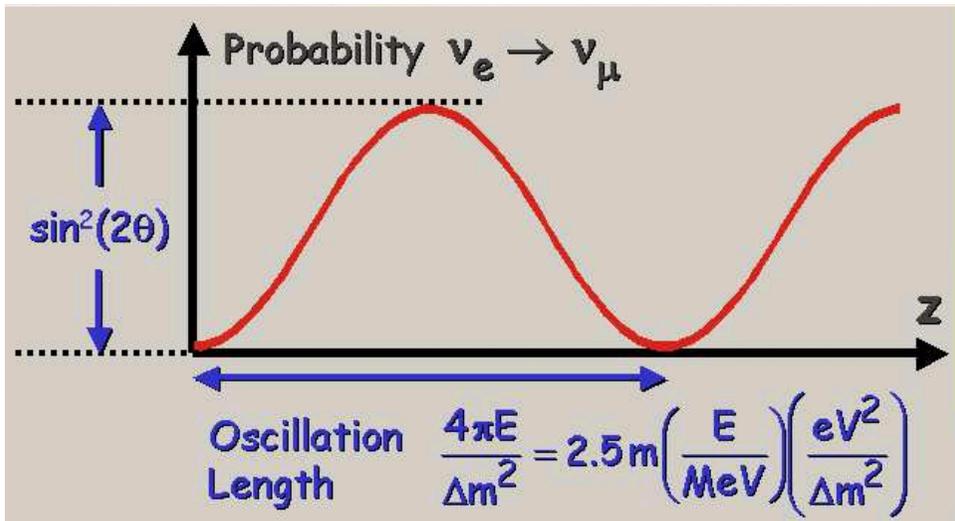


The Outline

- Neutrino Oscillation in Vacuum
- Matter Effect
- Solar Neutrino problem and its resolution
- Atmospheric neutrino problem and its resolution
- Current Status of Oscillation Parameters
- Future Perspective

Oscillation Probability in Vacuum

$$P_{\nu_e \nu_\mu} = \sin^2 2\theta \sin^2(\Delta m^2 L / 4E) = \sin^2 2\theta \sin^2(\pi L / \lambda)$$



$$\Delta m^2 = m_2^2 - m_1^2$$

$\theta \rightarrow$ mixing angle

$L \rightarrow$ Distance travelled (in m/Km)

$E \rightarrow \nu$ Energy (in MeV/GeV)

- Neutrino Oscillation requires
 - Non-zero neutrino mass
 - Non-zero mixing angles
 - Oscillation effect $\Delta m^2 \sim E/L$

● Oscillation Wavelength

$$\lambda = 4\pi E / \Delta m^2$$

$$\lambda = 2.5m(E/MeV)(eV^2/\Delta m^2)$$

- $\lambda \gg L, \sin^2(\pi L/\lambda) \rightarrow 0$
- $\lambda \ll L, \sin^2(\pi L/\lambda) \rightarrow 1/2$
- $\lambda \sim 2L, \sin^2(\pi L/\lambda) \sim 1 \rightarrow \Delta m^2 \sim E/L$

Neutrino Oscillation

- Sensitivity to Δm^2 in various oscillation experiments

Experiment	Energy(E)	Distance(L)	Δm^2
Reactor	1 MeV	100m	10^{-2} eV^2
Low E Accelerator	10 MeV	100m	10^{-1} eV^2
High E Accelerator	1 GeV	1 km	1 eV^2
Atmospheric ν	1 GeV	10,000 km	10^{-4} eV^2
Solar Neutrinos	10 MeV	10^{11} m	10^{-10} eV^2
Long-Baseline Reactor	1 MeV	1 km	10^{-3} eV^2
Long-Baseline Accelerator	1 GeV	1000km	10^{-3} eV^2

Matter effects: Two Flavours

- Mass eigenvalues in matter :

$$\frac{m_{1m,2m}^2}{2E} = \frac{1}{2E} \left(\frac{A}{2} \mp \frac{1}{2} \sqrt{(\Delta m^2 \cos 2\theta - A)^2 + (\Delta m^2 \sin 2\theta)^2} \right)$$

- The mass squared difference in matter:

$$\Delta m_m^2 = \sqrt{(A - \Delta m^2 \cos 2\theta)^2 + \sin^2 2\theta}$$

- Effective mixing angle θ_M in matter :

$$\tan 2\theta_M = \frac{\Delta m_{21}^2 \sin 2\theta}{\Delta m_{21}^2 \cos 2\theta - A}$$

$$\Delta m^2 \cos 2\theta = A = 2\sqrt{2}G_F n_e,$$

$\theta_M \rightarrow \pi/4$ MSW Resonance

- Mixing angle **Maximal**

L. Wolfenstein, PRD 17, 1978

S.P. Mikheyev, A.Yu. Smirnov, SJNP 42, 1985

- Mass squared difference in Matter **Minimal**

Intrinsic Neutrino Properties $\Delta m^2, \theta$

- Survival probability in constant density matter is

$$P_{ee}^m = 1 - \sin^2 2\theta_m \sin^2 \left(\frac{\Delta m_m^2 L}{4E} \right)$$

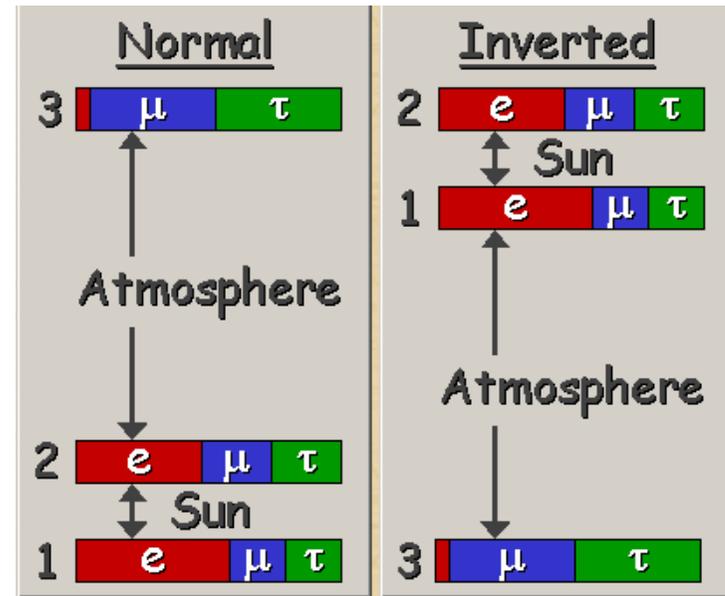
Three Neutrino oscillation parameters

- 9 unknown parameters for three neutrino flavours
 - 3 masses, m_1 , m_2 and m_3
 - 3 mixing angles
 - 3 phases

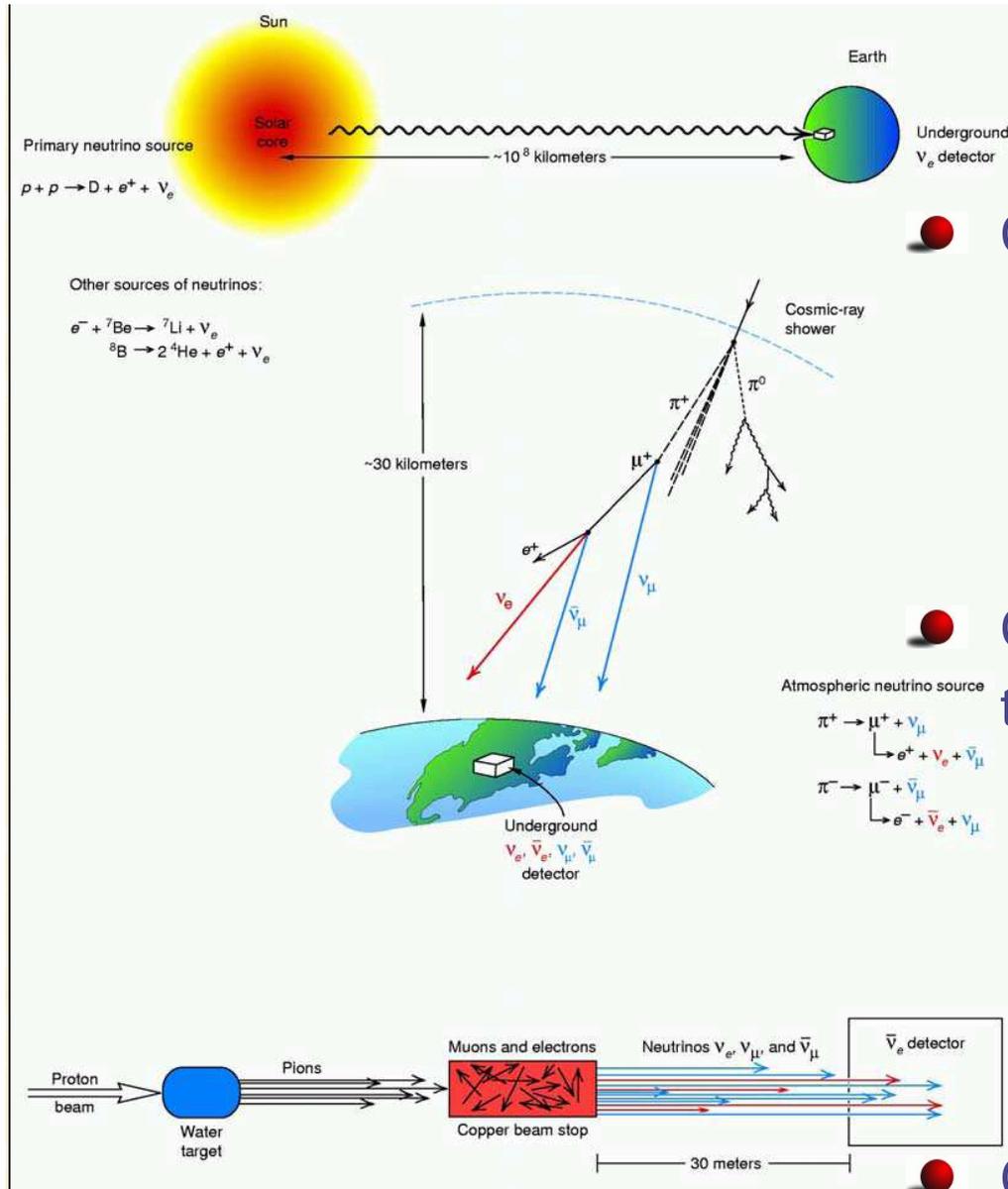
$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} 1 & & \\ & c_{23} & s_{23} \\ & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & & e^{-i\delta} s_{13} \\ & 1 & \\ -e^{i\delta} s_{13} & & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} \\ -s_{12} & c_{12} \\ & & 1 \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

$c_{12} = \cos\theta_{12}$ etc., δ CP-violating phase

- Oscillation experiments – sensitive to mass squared differences
- Solar : $\Delta m_{21}^2 = m_2^2 - m_1^2$, θ_{12}
- Atmospheric : $\Delta m_{31}^2 = m_3^2 - m_1^2$, θ_{23}
- Reactor Neutrinos : θ_{13}
- For small θ_{13} approximate 2 generation works



Experimental Evidences for Neutrino Oscillation



Solar Neutrinos

Confirmation from reactor neutrinos

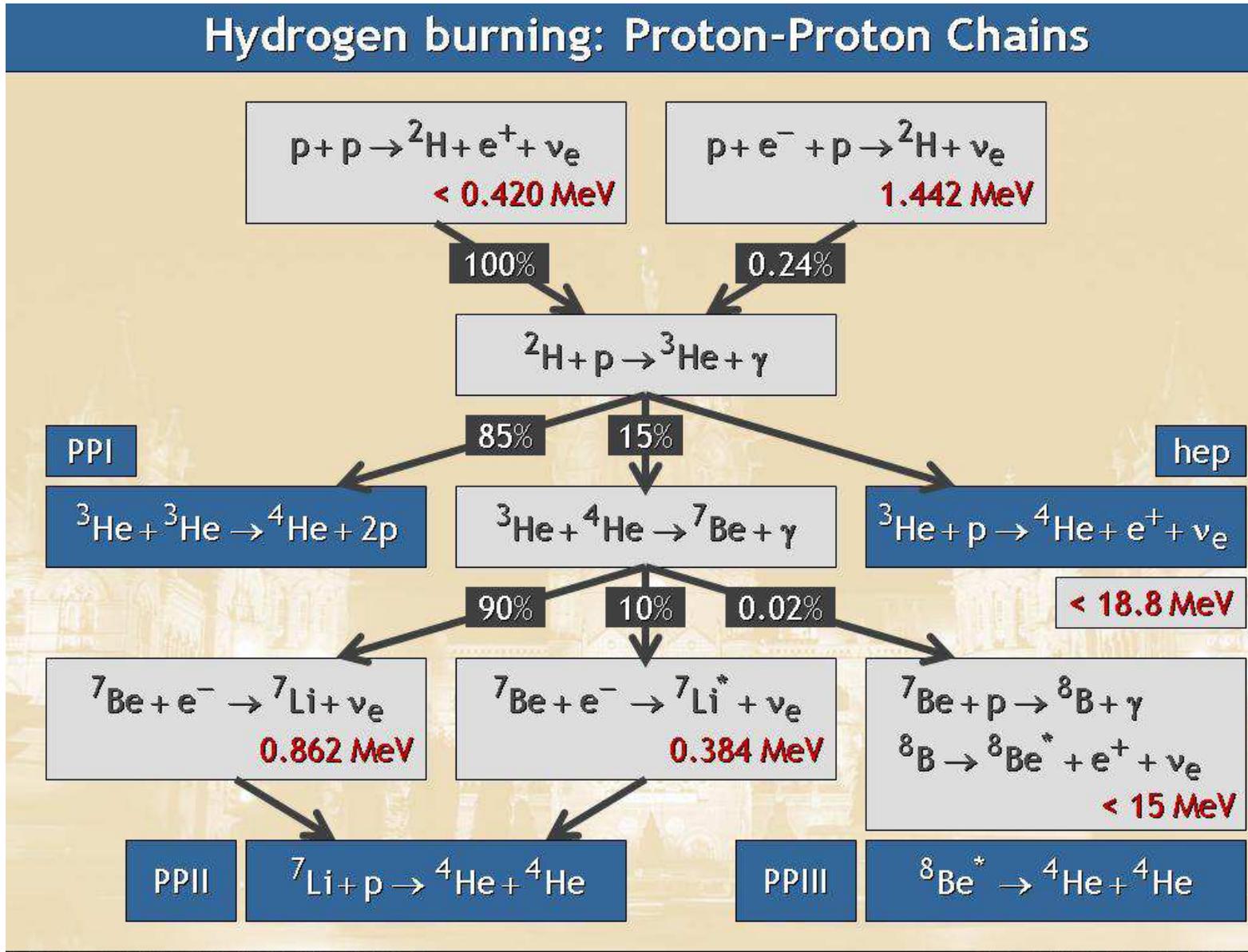
Atmospheric Neutrinos

Confirmation from accelerator neutrinos

LSND Experiment

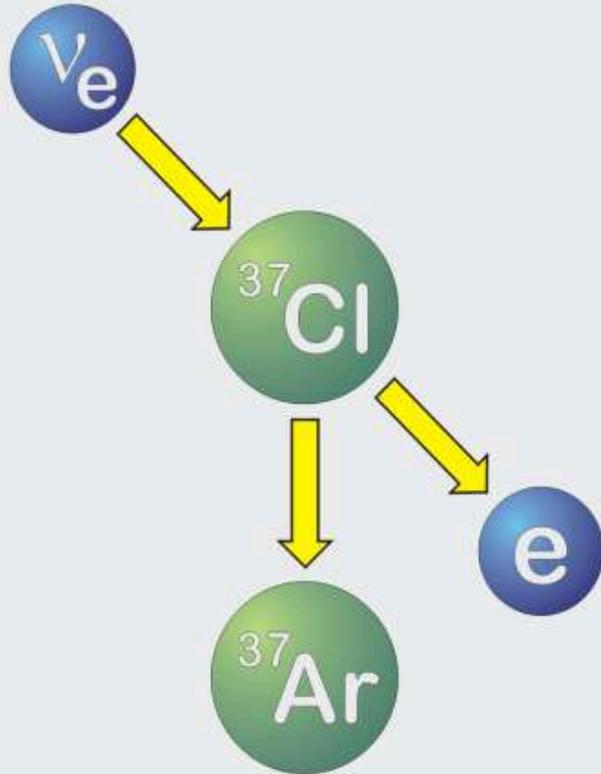
Confirmation from MiniBoone (?)

Nuclear Reaction inside the Sun

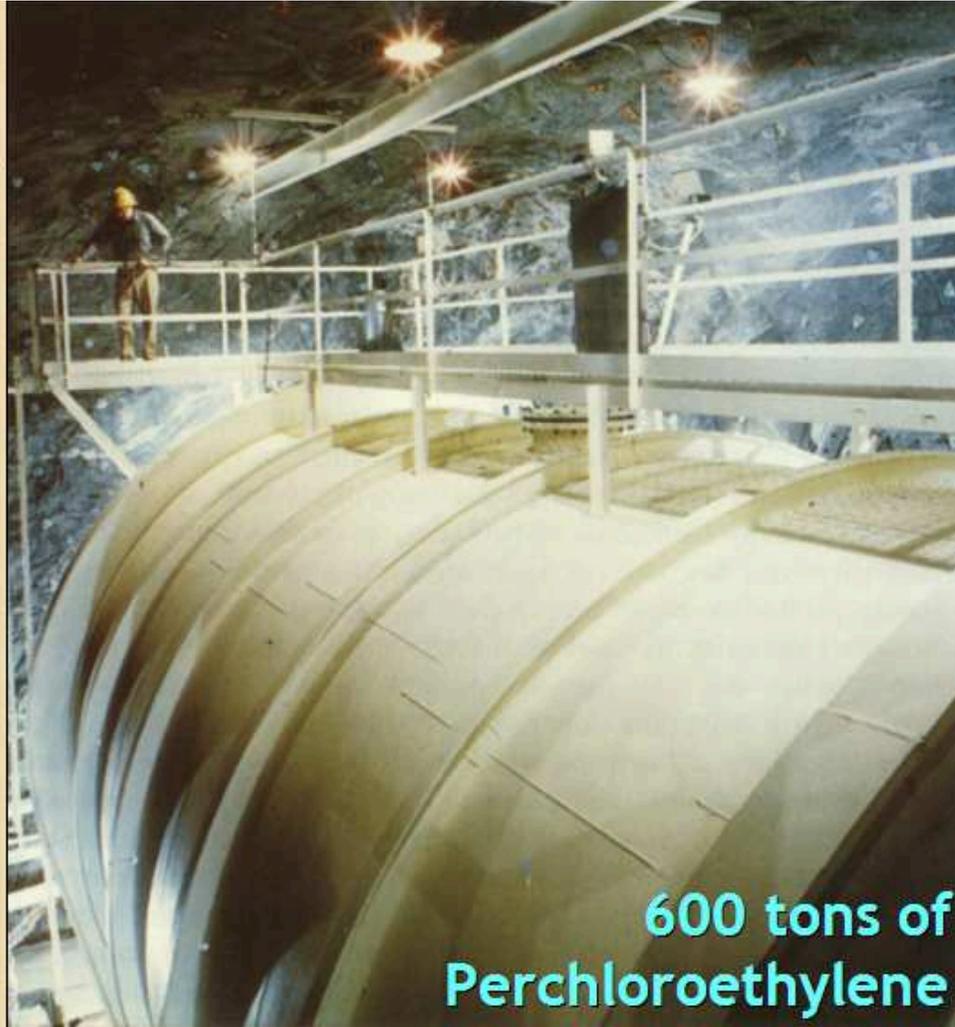


First Detection of Solar Neutrinos

Inverse beta decay
of chlorine



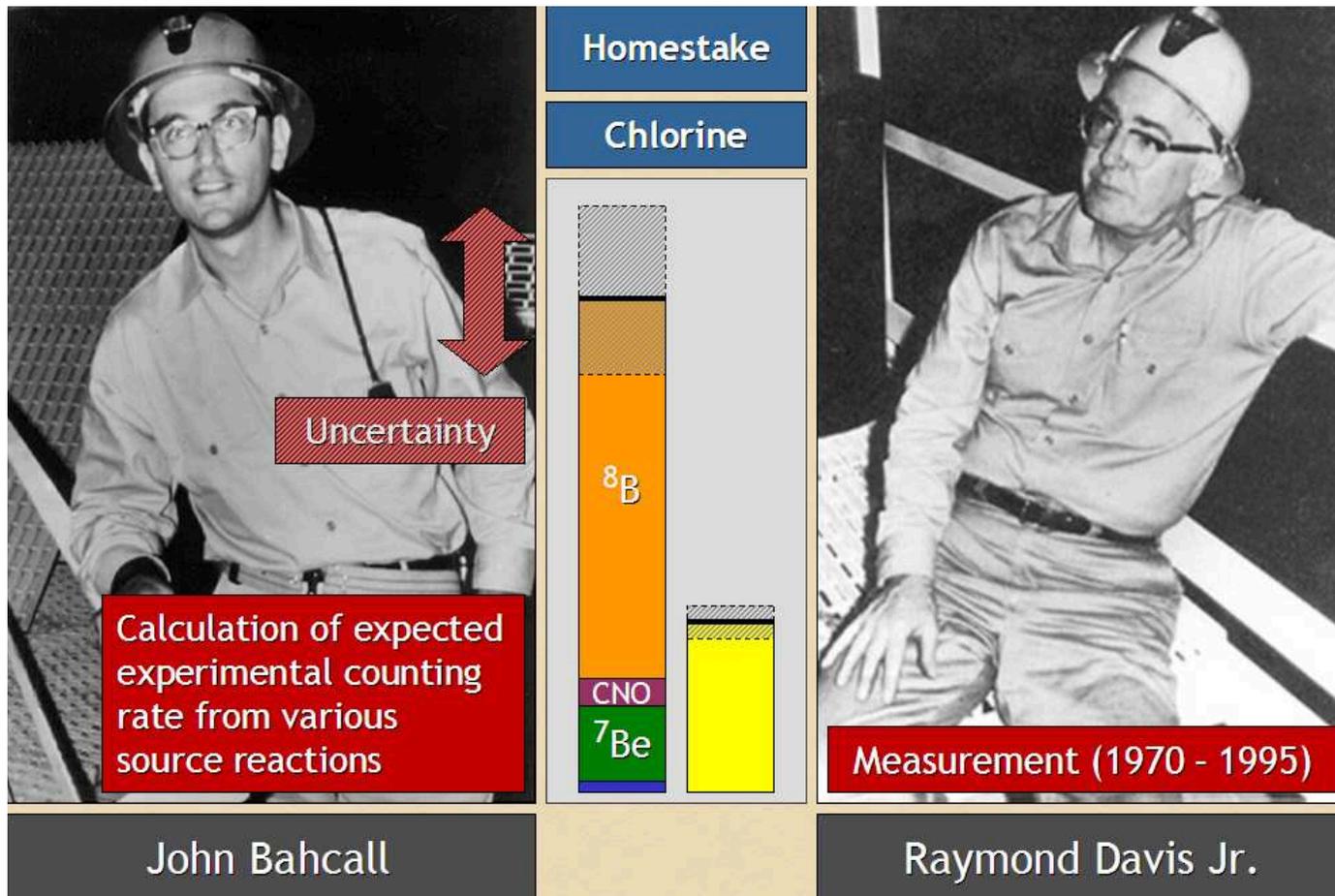
Radiochemical
Method



Homestake solar neutrino
observatory (1967–2002)

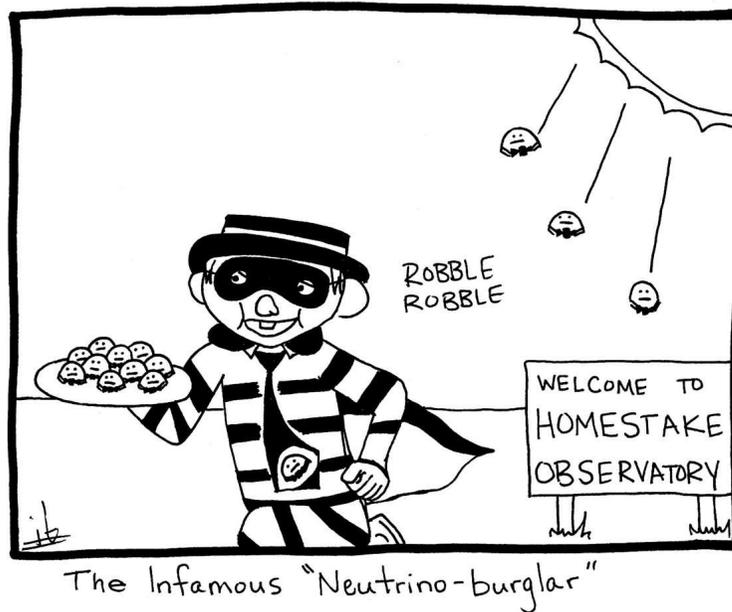
First Detection of Solar Neutrinos

- Pioneered by John Bahcall and Raymond Davies
- To and verify the hypothesis of energy generation in stars



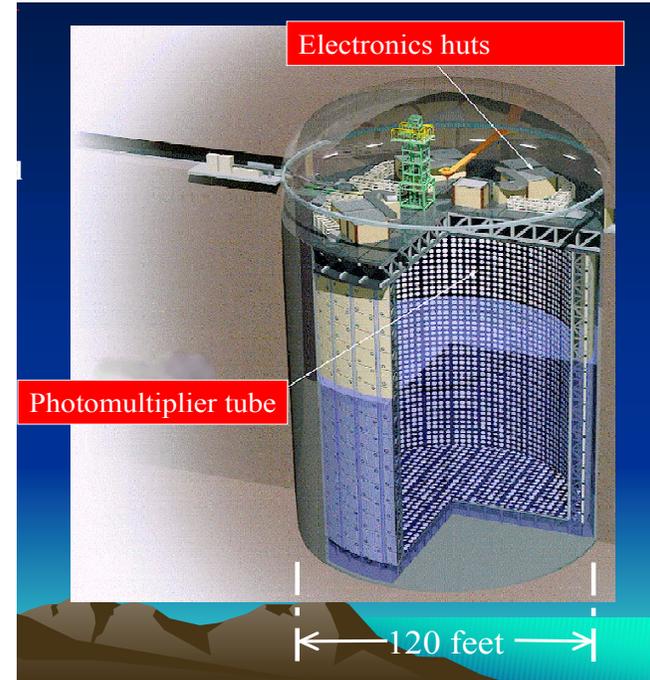
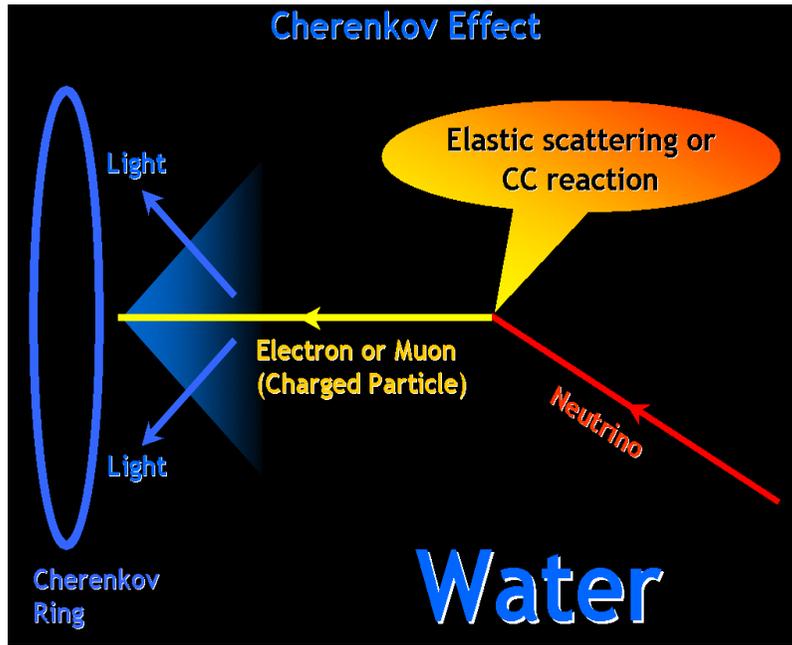
- Only $\frac{1}{3}$ rd of the predicted solar electron neutrinos observed

Where are the missing solar neutrinos going ?



- The experiment is wrong – difficult to detect a handful of Ar atoms
- Solar Model calculations are wrong
- Solar electron neutrinos getting converted to muon or tau neutrinos
- New Experiments were planned to check the above results

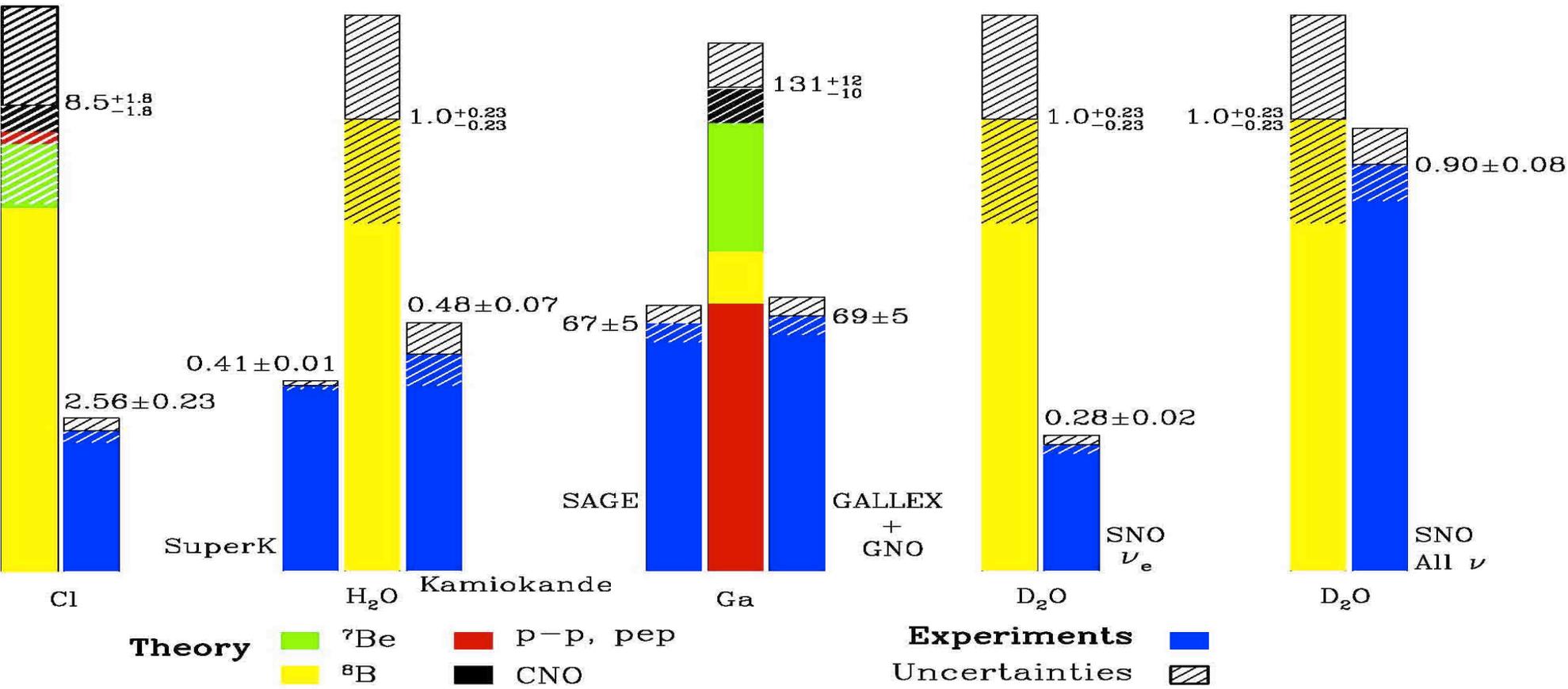
Solar Neutrino Experiments



- $e^- + \nu_x \rightarrow e^- + \nu_x$
- Detection of Solar neutrinos in real time
- KamioKande, Japan (1991)

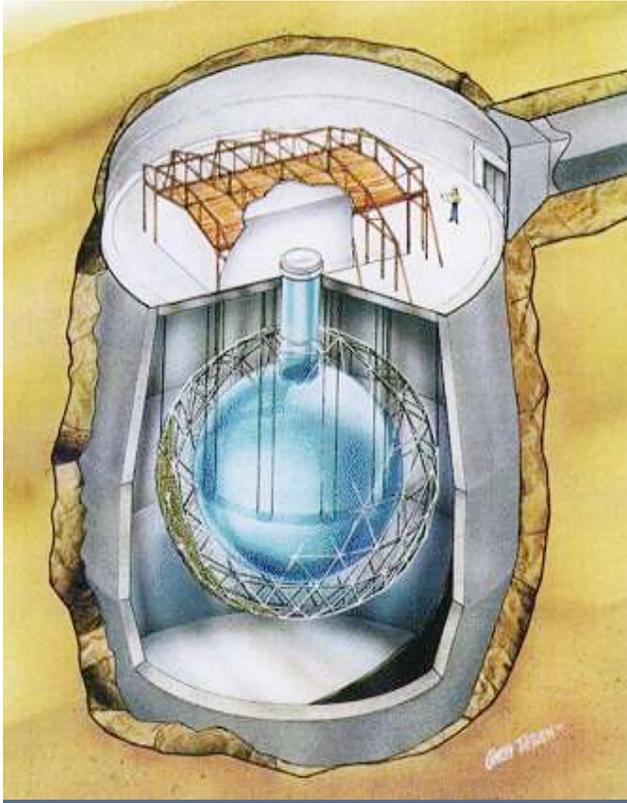
- SuperKamiokande (1991-)
 - 50,000 ton of purified water
 - 3000 feet underground Japanese Alps
 - 13,000 photomultiplier tubes

Solar Neutrinos: Experimental Results



- Observed ν_e flux < Theoretical Prediction → solar Neutrino Problem
- Remained unresolved for 30 years

Sudbury Neutrino Observatory



- 1000 Tonnes Heavy Water (D_2O)
- 7000 feet underground in Sudbury Mine Canada
- 10,000 photomultiplier tubes



$$\frac{CC}{NC} = \frac{\nu_e}{\nu_e + \nu_\mu + \nu_\tau} < 1$$

→ indicates conversion of ν_e to ν_μ, ν_τ

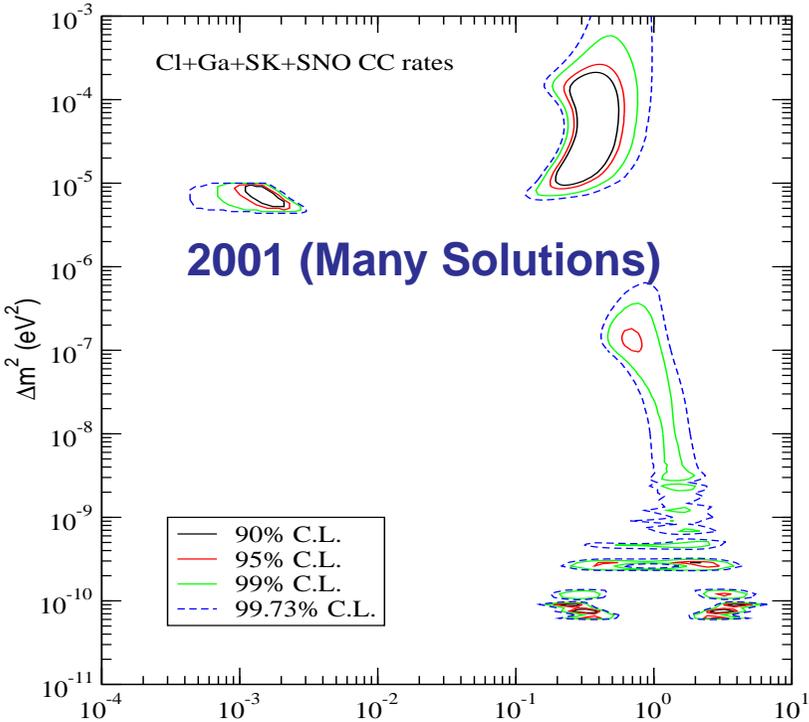
7σ evidence for a non-zero $\nu_{\mu,\tau}$ flux from SUN

Solution to the Solar Neutrino Problem

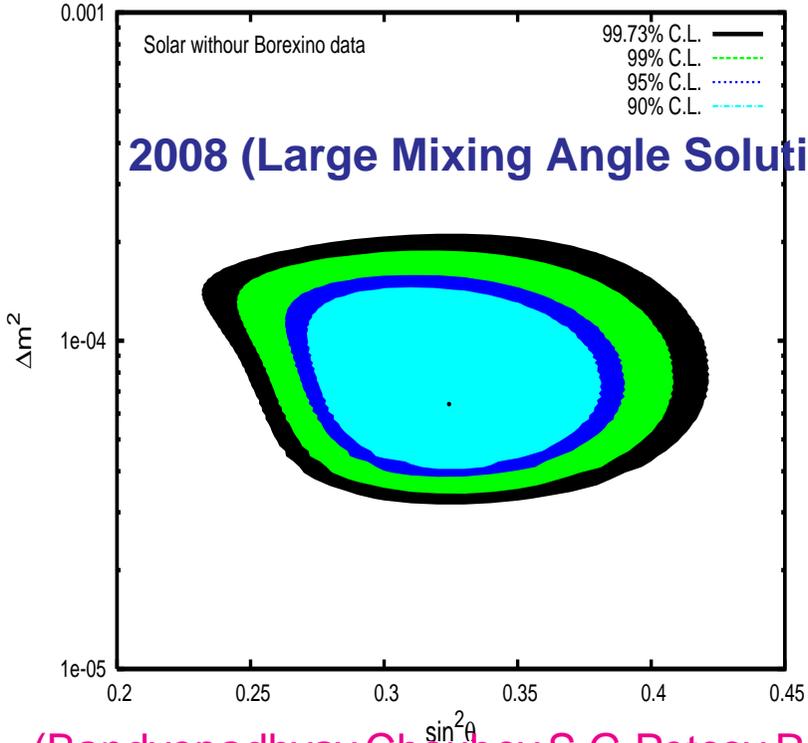
- Neutrinos undergo **MSW resonant flavour conversion** of ν_e inside the sun

$$\Delta m_{21}^2 \cos 2\theta_{12} - 2\sqrt{2}G_F n_e E = 0$$

- Sensitive to the sign of Δm_{21}^2



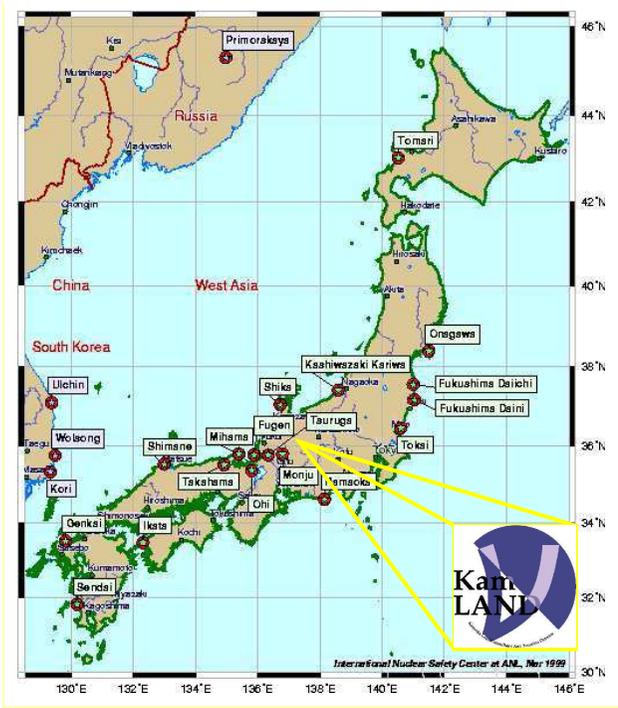
(Bandyopadhyay, Choubey, S.G, Kar, PLB, 2001)



(Bandyopadhyay, Choubey, S.G, Petcov, Roy, 2008)

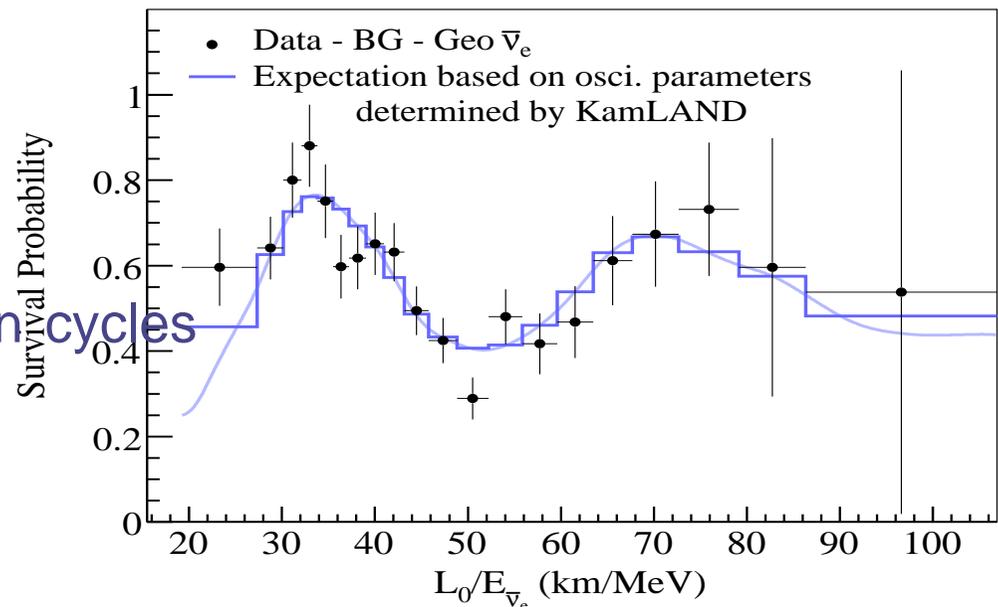
Confirmed by KamLAND experiment using reactor neutrinos

KamLAND



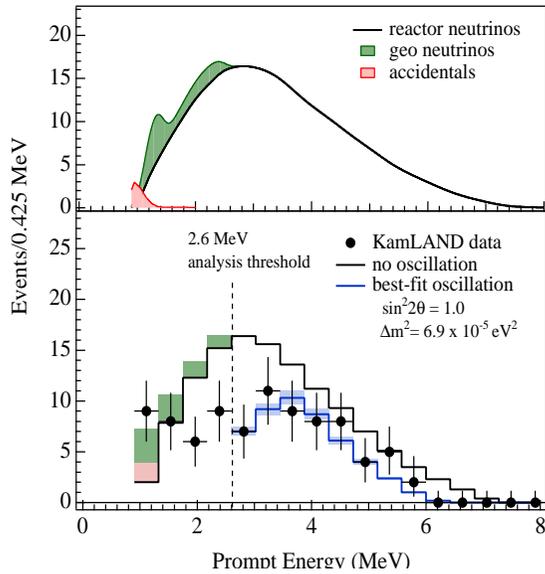
- 1 kton liquid scintillator neutrino detector at Kamioka, Japan
- detects antineutrinos coming from Japanese nuclear reactors through:
$$\bar{\nu}_e + p \rightarrow n + e^+$$
- $E_\nu \sim 3 \text{ MeV}$, $L \sim 1.8 \times 10^5 \text{ m}$, $\Delta m^2 \sim 1.6 \times 10^{-5} \text{ eV}^2$

- Data: 2002 - 2007
- Almost two complete oscillation cycles are observed

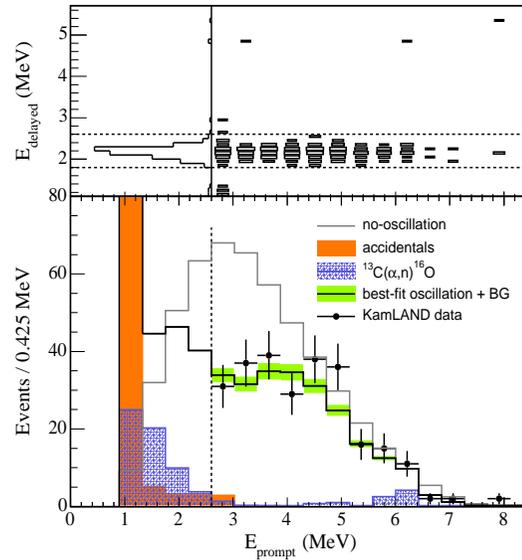


Impact of KamLAND on Oscillation Parameters

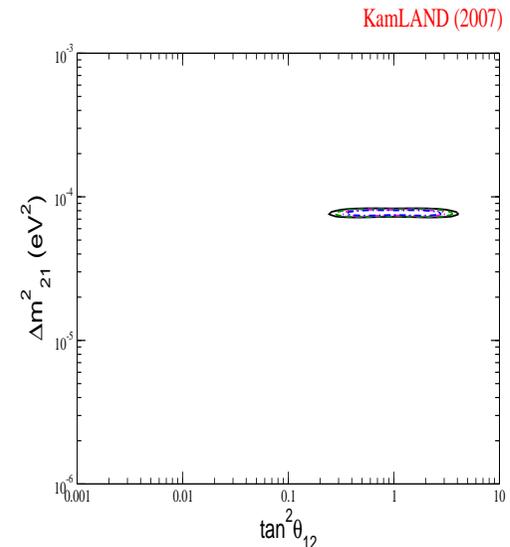
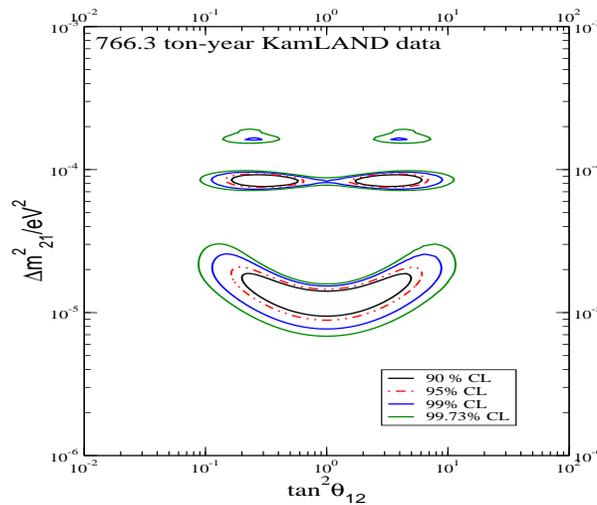
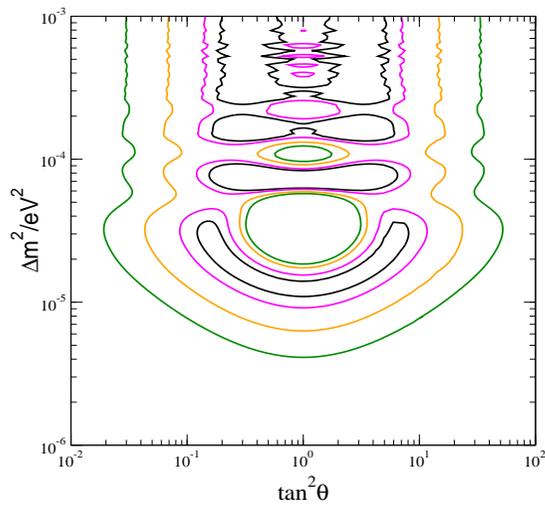
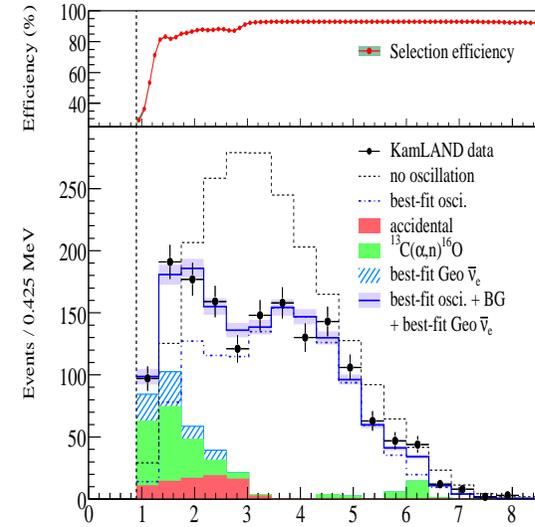
2003



2004

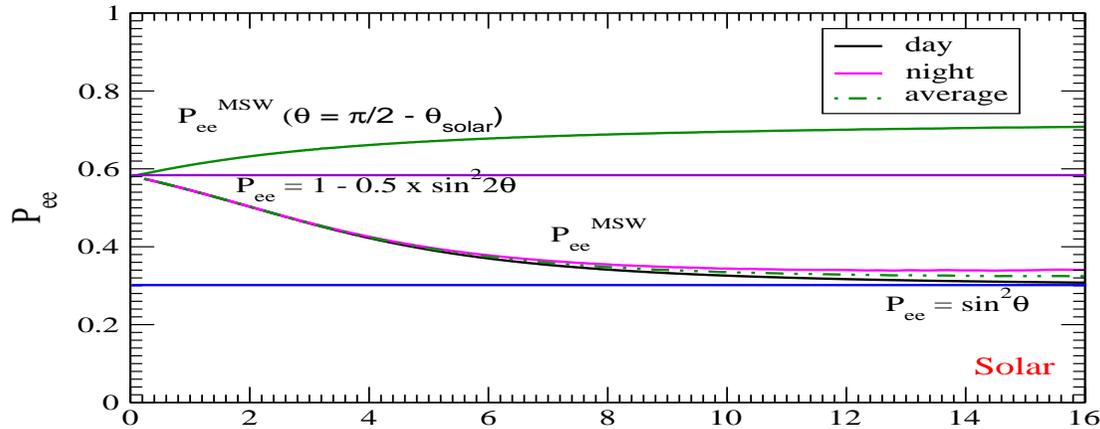


2007



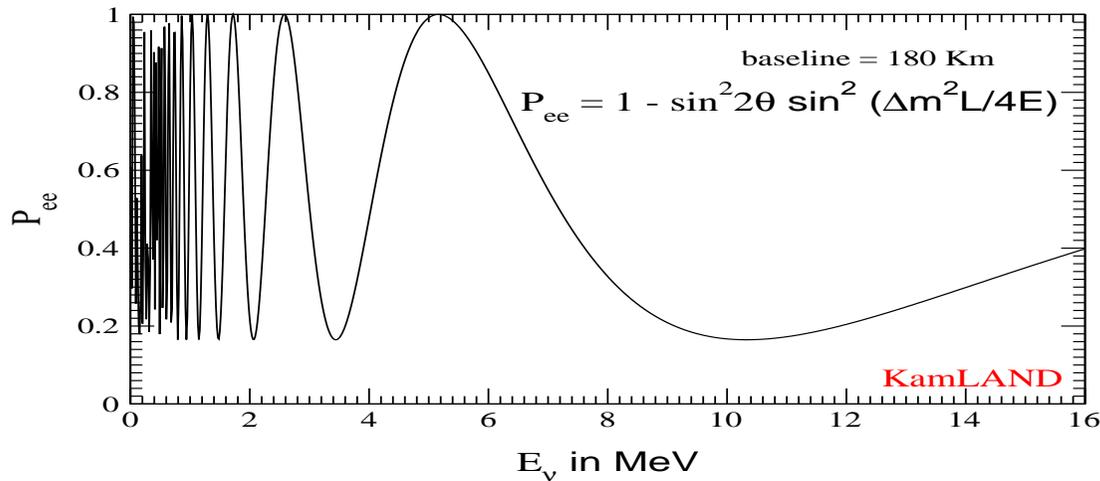
Bandyopadhyay, Choubey, Goswami, Petcov, Roy 2003, 2005, 2008

Survival Probabilities for solar and KamLAND



Solar Neutrinos

Δm_{21}^2 effects are averaged out



KamLAND

can probe the L/E dependence of the oscillations in the LMA region \rightarrow unprecedented sensitivity to Δm^2

Atmospheric neutrinos . . .

Cosmic Ray + $A_{air} \rightarrow \pi^+ + \dots$

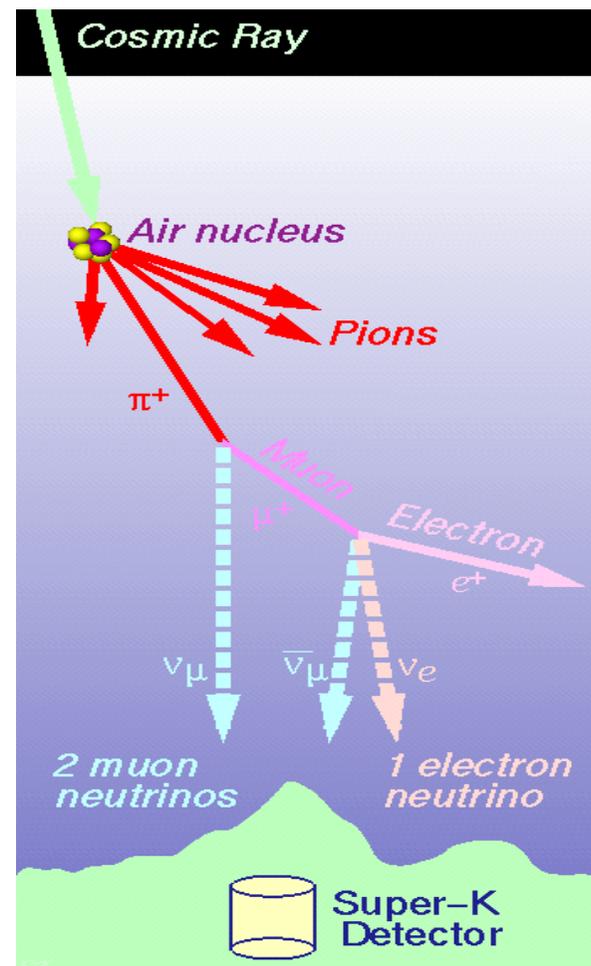
$\pi^+ \rightarrow \mu^+ + \nu_\mu$

$\mu^+ \rightarrow e^+ + \bar{\nu}_\mu + \nu_e$

- Energy: 100 MeV - TeV
- Pathlength: 15 -13,000 km
- Provides broad L/E band

$\nu_\mu : \nu_e = 2: 1$ (expected)

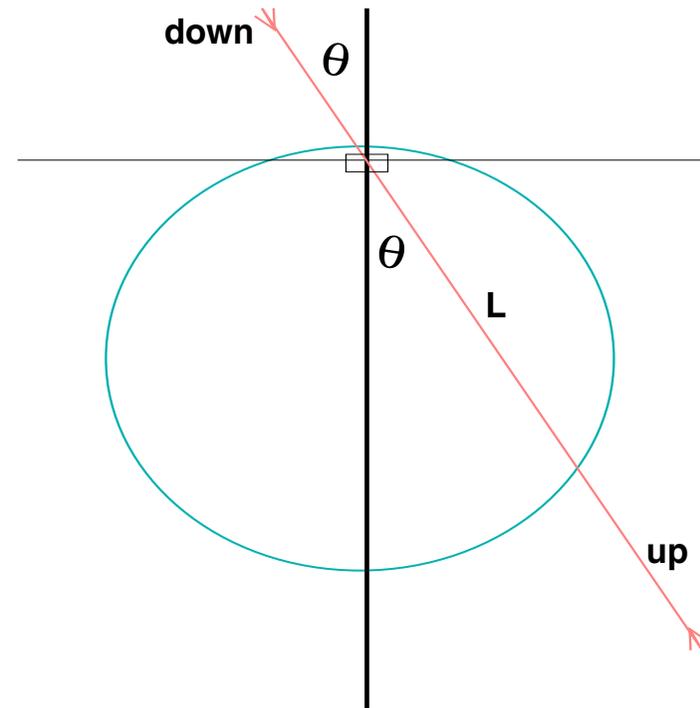
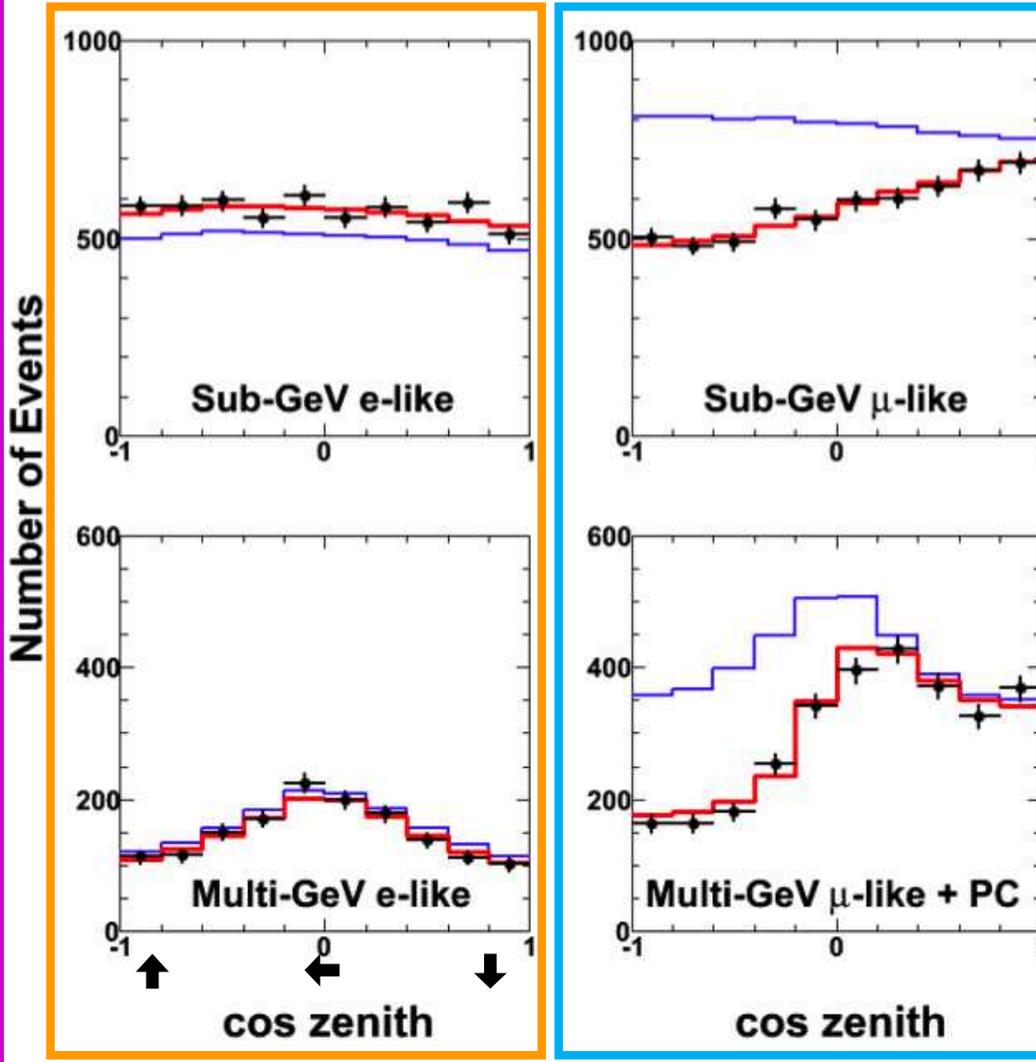
$\nu_\mu / \nu_e \sim 0.9 - 1$ (observed)



$\Rightarrow \nu_\mu$ conversion

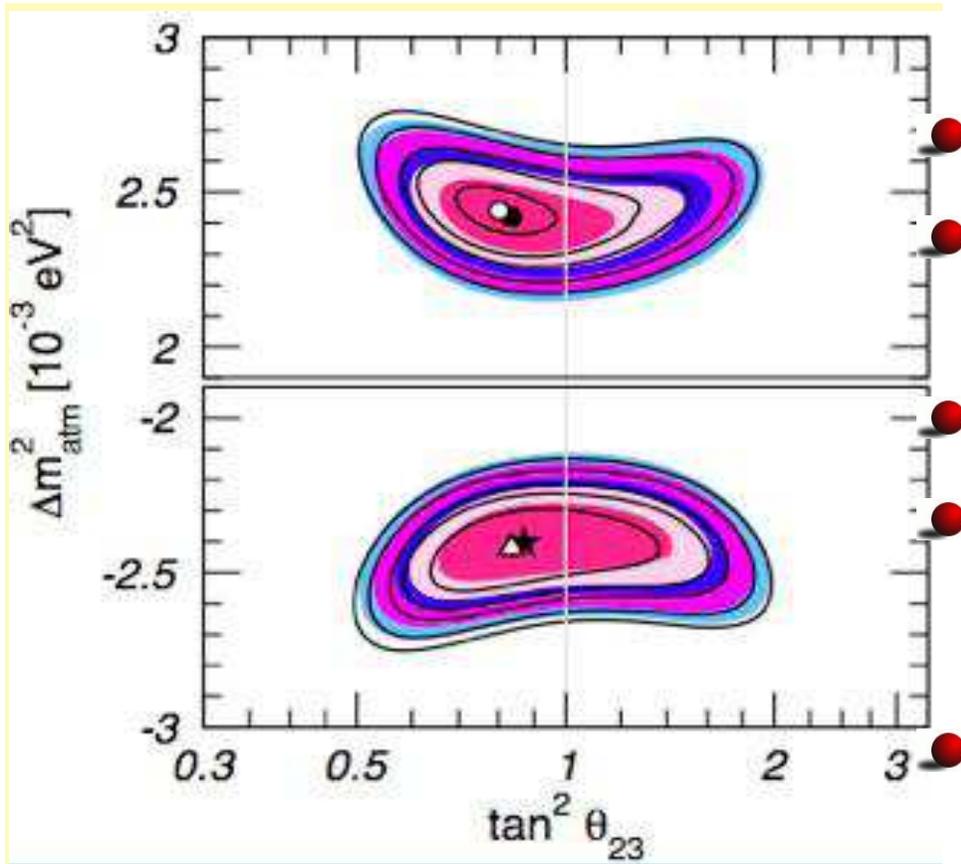
Detection of Atmospheric neutrinos at SK

SK-I+II+III, 2806 Days



- μ -events less than theoretical prediction
- The flux is up-down symmetric for $E > \text{a few GeV}$
- Up-down asymmetry in multi-GeV μ -events \implies neutrino oscillations

Solution to the Atmospheric Neutrino Problem



● Neutrinos have mass and mixing

● Dominant solution $\nu_{\mu} - \nu_{\tau}$ oscillation in vacuum

● Not very sensitive to matter effects

● Relevant probability

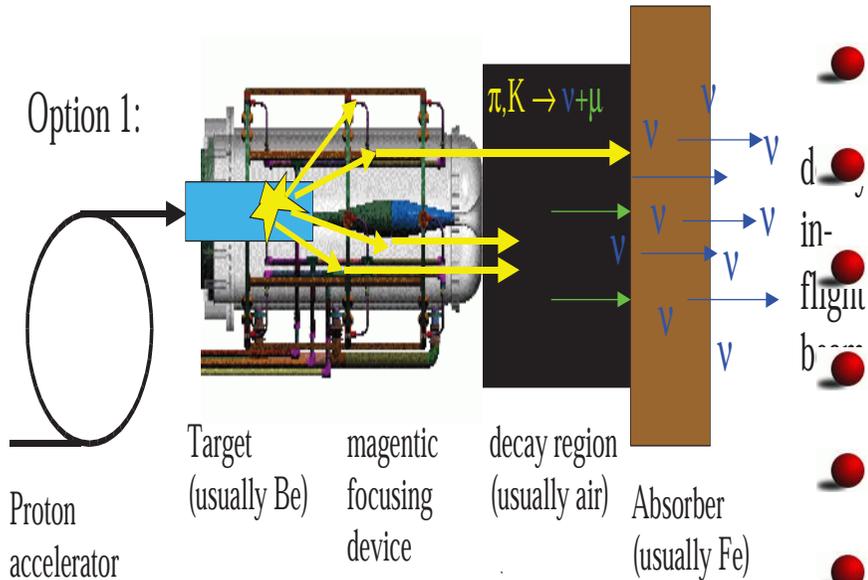
$$P_{\mu\mu} = 1 - \sin^2 2\theta_{23} \sin^2 \left(\frac{\Delta m_{31}^2 L}{4E} \right)$$

● No information on $|(\Delta m_{31}^2)|$.

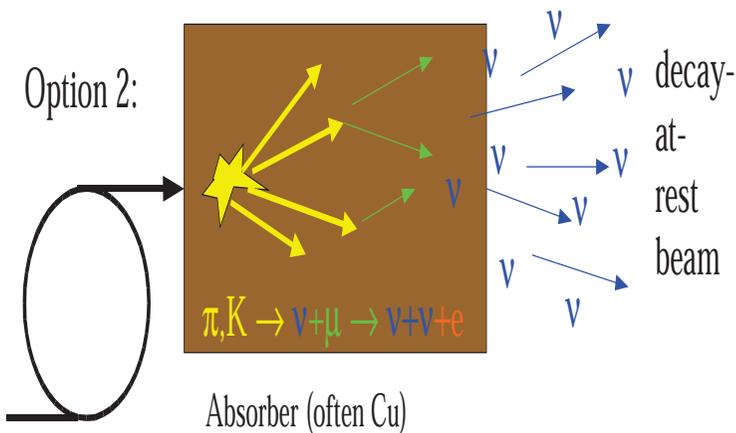
Fogli. et al., 2012

Long Baseline Accelerator Experiments

Making Neutrino Beams at Proton Accelerators



- Beam with well known properties
- Tune L/E to study interesting Δm^2
- Distance > 100 km
- L is a particular choice
- E is always a distribution
- Decay at rest: $E_{max} \sim 50 GeV$
- Decay in flight $E_{max} \sim 100 GeV$



Long Baseline Accelerator Experiments: T2K



- Beam with well known properties from J-PARC
- Long Baseline Distance ~ 295 km
- $L/E \sim 500$ GeV to study Δm_{atm}^2



Physics Goals

- ν_e appearance $\implies \theta_{13}$
- ν_μ disappearance $\implies \Delta m_{31}^2, \theta_{23}$

Super Kamiokande
“far” detector (FD)

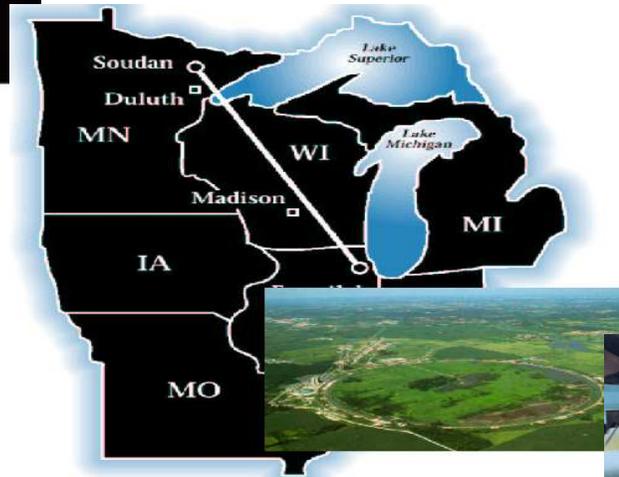


Confirmed the atmospheric neutrino oscillations using man-made sources

Long Baseline Accelerator Experiments: MINOS



- Beam with well known properties
- Long Baseline Distance > 100 km
- L/E ~ 500 GeV to study Δm_{atm}^2

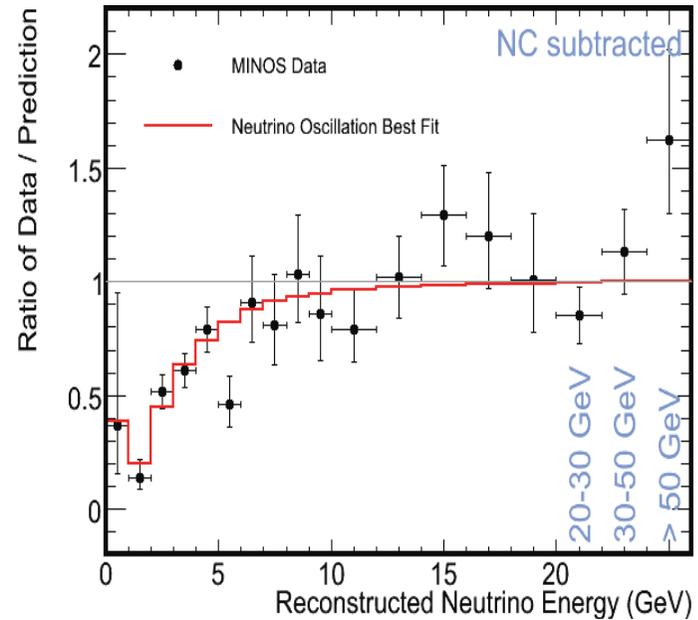
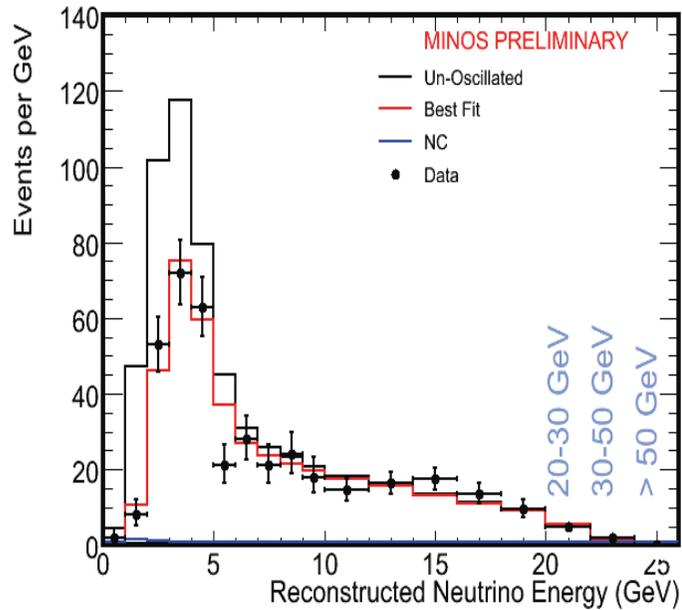


Physics Goals

- ν_{μ} disappearance $\implies \Delta m_{31}^2, \theta_{23}$
- $\bar{\nu}_{\mu}$ disappearance $\implies \overline{\Delta m_{31}^2}, \overline{\theta_{23}}$
- ν_e appearance $\implies \theta_{13}$



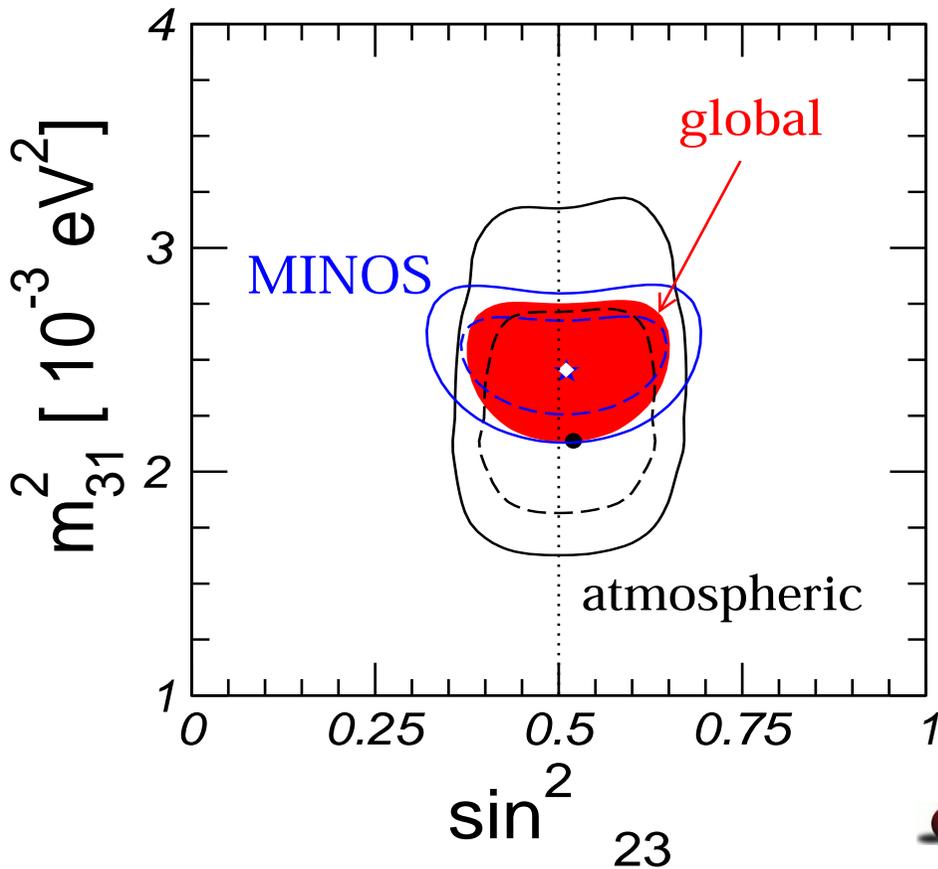
First MINOS Result



MINOS and K2K confirmed the atmospheric neutrino oscillations using man-made sources

MINOS Collaboration: [arXiv:0806.2237](https://arxiv.org/abs/0806.2237)

Atmospheric neutrino Oscillation Parameter



$$|\Delta m_{atm}^2| = 2.4 \times 10^{-3} \text{eV}^2$$

$$\sin^2 \theta_{23} = 0.5$$

- Two generation $\nu_{\mu} - \nu_{\tau}$ oscillation

- Matter Effect does not play role

- Relevant probability

$$P_{\mu\mu} = 1 - \sin^2 2\theta_{atm} \sin^2 \left(\frac{\Delta m_{atm}^2 L}{4E} \right)$$

$$(\theta_{atm} \equiv \theta_{23}, \Delta m_{atm}^2 \equiv \Delta m_{31}^2)$$

- $\theta_{23} - (\pi/2 - \theta_{23})$ symmetry

- No information on $\text{sgn}(\Delta m_{atm}^2)$.

- 3 σ Precision:

$|\Delta m_{atm}^2| \sim 12\%$ mainly by MINOS

$\sin^2 \theta_{23} \sim 24\%$ mainly by Atmospheric

The Last of the Mixing angles



- Till 2011 : $\theta_{13} < 13^\circ$
- New results from three reactor experiments
- Long Baseline: distance ~ 1 km
- $L/E \sim \text{Km/MeV}$
- $P_{ee} = 1 - \sin^2 2\theta_{13} \sin^2 4\pi E / \Delta m_{31}^2$



● $\theta_{13} \sim 9^\circ$



Summary of Three Neutrino Oscillation Parameters

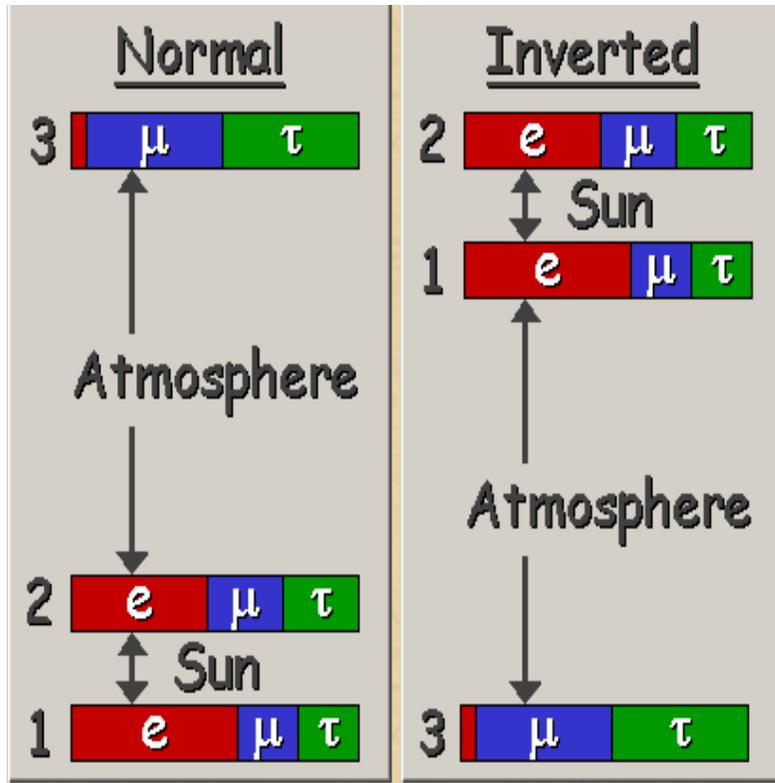
- Significant progress in measurement of neutrino oscillation parameters in the last 15 years
- The last of the unknown mixing angles θ_{13} is measured (2012)

current status of oscillation parameters

Parameter	Best fit	3σ range
$\Delta m_{21}^2 (eV^2)$	7.5×10^{-5}	$(7.0 - 8.2) \times 10^{-5}$
$ \Delta m_{31}^2 (eV^2)$	2.4×10^{-3}	$(2.2 - 2.6) \times 10^{-3}$
$\theta_{12} (degree)$	33.6	30.6 - 36.8
$\theta_{23} (degree)$	38.4	35.1 - 53.0
$\theta_{13} (degree)$	8.9	7.5 - 10.2

- No constraint on the CP-phases

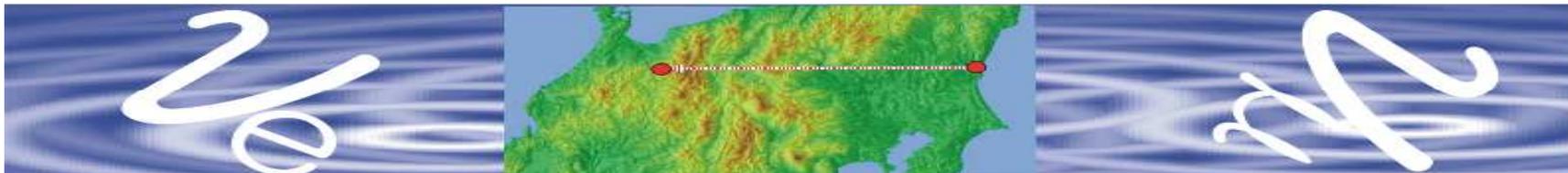
The outstanding issues



- The Neutrino Mass hierarchy
- Determination of CP phases
- The absolute neutrino mass scale
- The nature of Neutrinos : Dirac or Majorana
- Mass hierarchy can be determined by exploiting the earth matter effect if θ_{13} is appreciable.
- Determination of CP phase is possible only if θ_{13} is non-zero.

Future Longbaseline Experiments

- **Superbeam Experiments** (conventional accelerator beam but more power)
 - T2K: – $E \sim 0.76$ GeV, $L = 295$ km
 - NO ν A: – $E \sim 1 - 3$ GeV, $L = 810$ km
 - T2K taking data, NO ν A to start soon
 - LBNO in Europe ~ 2200 km ,LBNE in US ~ 1300 km
 - Utilizes **flavour conversion** of neutrinos passing through **earths matter**
 - Sensitive to **hierarchy** and δ_{cp}
 - Most useful is the conversion probability ν_e to ν_μ
- **Beta-Beams, Neutrino-Factory** (low and high Energy)
 - Powerful source, adequate flux at large distance ($\sim 3000 - 7000$ Km)
 - Important for small θ_{13}



Matter effects: Three Flavours

- The effective Hamiltonian is

$$\tilde{H} = \frac{1}{2E} \left[U \begin{pmatrix} m_1^2 & 0 & 0 \\ 0 & m_2^2 & 0 \\ 0 & 0 & m_3^2 \end{pmatrix} U^\dagger + \begin{pmatrix} A & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right]$$

excluding terms \propto identity matrix. $U = R_{23}R_{13}R_{12}$

Matter effects: Three Flavours

- The effective Hamiltonian is

$$\tilde{H} = \frac{1}{2E} \left[U \begin{pmatrix} m_1^2 & 0 & 0 \\ 0 & m_2^2 & 0 \\ 0 & 0 & m_3^2 \end{pmatrix} U^\dagger + \begin{pmatrix} A & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right]$$

excluding terms \propto identity matrix. $U = R_{23}R_{13}R_{12}$

- Subtracting m_1^2 from the first part,

$$\tilde{H} = \frac{1}{2E} \left[U \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + \begin{pmatrix} A & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right]$$

Matter effects: Three Flavours

- The effective Hamiltonian is

$$\tilde{H} = \frac{1}{2E} \left[U \begin{pmatrix} m_1^2 & 0 & 0 \\ 0 & m_2^2 & 0 \\ 0 & 0 & m_3^2 \end{pmatrix} U^\dagger + \begin{pmatrix} A & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right]$$

excluding terms \propto identity matrix. $U = R_{23}R_{13}R_{12}$

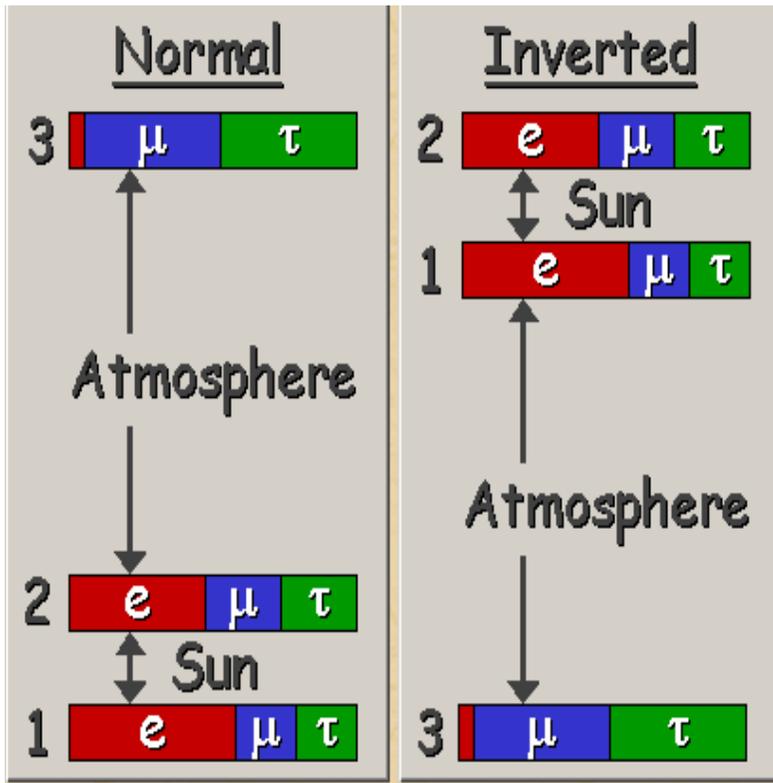
- 1 MSD ($\Delta m_{21}^2 = 0$) limit

$$\tilde{H} = \frac{1}{2E} \left[U \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} U^\dagger + \begin{pmatrix} A & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right]$$

→ Resonance in the 1-3 sector

Matter Effect, θ_{13} and hierarchy

$$\bullet \quad \tan 2\theta_{13}^m = \frac{\Delta m_{31}^2 \sin 2\theta_{13}}{\Delta m_{31}^2 \cos 2\theta_{13} \pm 2\sqrt{2}G_F n_e E}$$



- For $\Delta m_{31}^2 > 0$ matter resonance in neutrinos
- For $\Delta m_{31}^2 < 0$ matter resonance in anti neutrinos
- Experiments that can differentiate between neutrinos and antineutrinos can probe **matter effects**
- Experiments sensitive to **matter effects** can probe the mass hierarchy
- Larger θ_{13} , larger will be matter effects

The Golden Channel

$$P_{e\mu} = |\cos \theta_{23} A_S e^{i\delta} + \sin \theta_{23} A_A|^2$$

- $A_S \rightarrow$ **Solar** amplitude depends on Δm_{21}^2 and θ_{12}
- $A_A \rightarrow$ **atmospheric** amplitude depends on Δm_{31}^2 and θ_{13}
- **CP violation** arises from the **interference** term
- Absence of **CP violation** requires either $A_S = 0$ or $A_A = 0$.

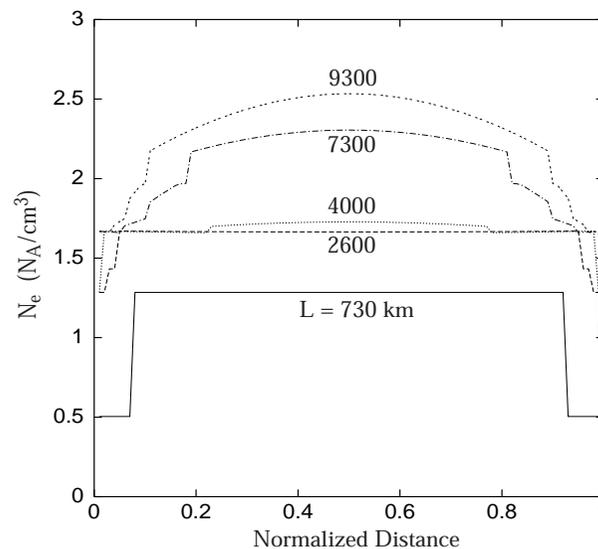
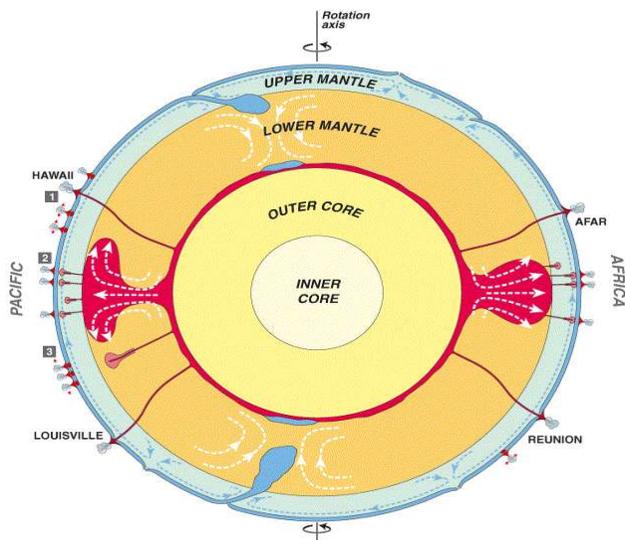
A.Yu. Smirnov, hep/ph 0610198

- **GOLDEN** because it can probe all three unknowns θ_{13} , $\text{sgn}\Delta m_{31}^2$, δ_{CP}

The Golden Channel

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2} \\
 &+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos(\Delta - \delta_{\text{CP}}) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin(1-\hat{A})\Delta}{(1-\hat{A})} \\
 &+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

- $\alpha = \Delta m_{31}^2 / \Delta m_{31}^2 \approx 0.04$ $\sin^2 \theta_{13} \sim 0.01$
- $\hat{A} \equiv 2\sqrt{2}G_F n_e E_\nu / \Delta m_{31}^2$, $\Delta \equiv \Delta m_{31}^2 L / (4E_\nu)$,
- Expanded in small parameters α and $\sin^2 \theta_{13}$ (constant matter density)



Attack of the Clones

- δ_{CP} can vary from (0 to 2π) and creates the problem of **Parameter Degeneracies**

● $(\theta_{13}, \delta_{CP})$ intrinsic degeneracy

Burguet-Castell, Gavela, Gomez-Cadenas, Hernandez, Mena, hep-ph/0103258

● $(\text{sgn}(\Delta m_{31}^2), \delta_{CP})$ degeneracy

Minakata, Nunokawa, hep-ph/0108085

● $(\theta_{23}, \pi/2 - \theta_{23})$ degeneracy

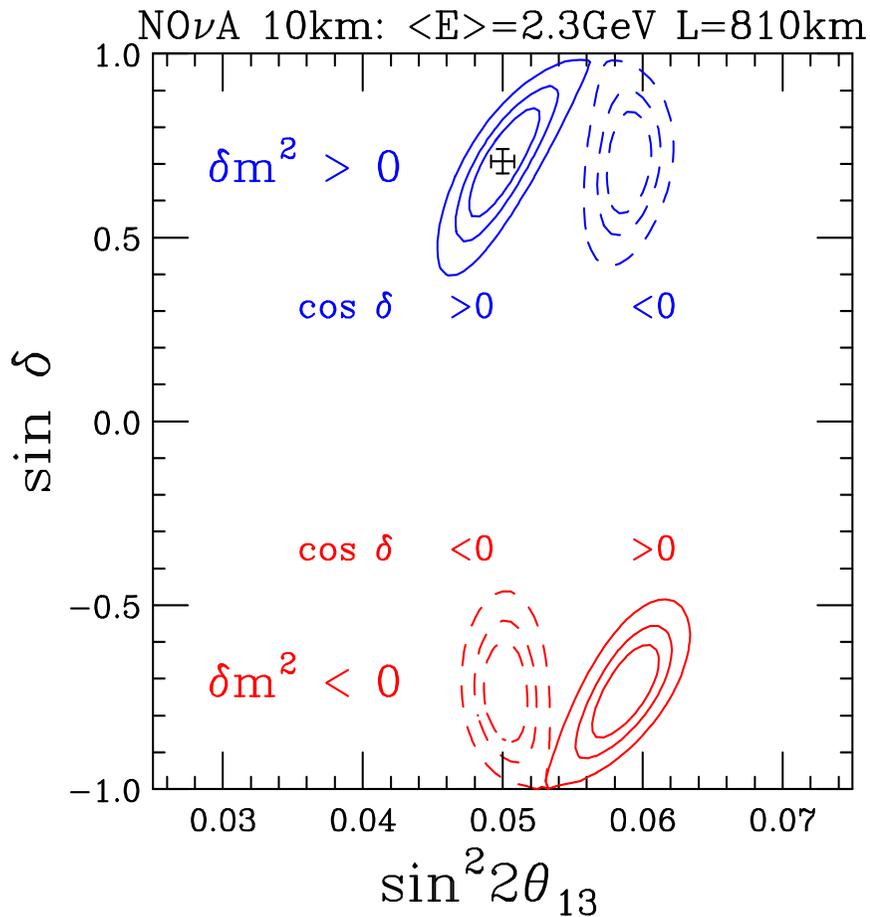
Fogli, Lisi, hep-ph/9604415



- Different set of parameters producing equally good fit to the data
- Give rise to multiple solutions → **Eightfold degeneracy**

The Phantom Menace

- Ghost (Degenerate) Solutions in $(\delta - \theta_{13})$ plane
- Unambiguous determination of parameters difficult



Mena and Parke, 2005

A New Hope: Magic Baseline

$$P_{e\mu} \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2} \\ + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos(\Delta - \delta_{CP}) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin(1-\hat{A})\Delta}{(1-\hat{A})} \\ + \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}$$

● If $\frac{\sin(\hat{A}\Delta)}{\hat{A}} = 0$

$\implies P_{e\mu}$ independent of δ_{CP}

A New Hope: Magic Baseline

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2} \\
 &+ \cancel{\alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos(\Delta - \delta_{CP}) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin(1-\hat{A})\Delta}{(1-\hat{A})}} \\
 &+ \cancel{\alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}}
 \end{aligned}$$

A New Hope: Magic Baseline

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2} \\
 &+ \cancel{\alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos(\Delta - \delta_{CP}) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin(1-\hat{A})\Delta}{(1-\hat{A})}} \\
 &+ \cancel{\alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}}
 \end{aligned}$$

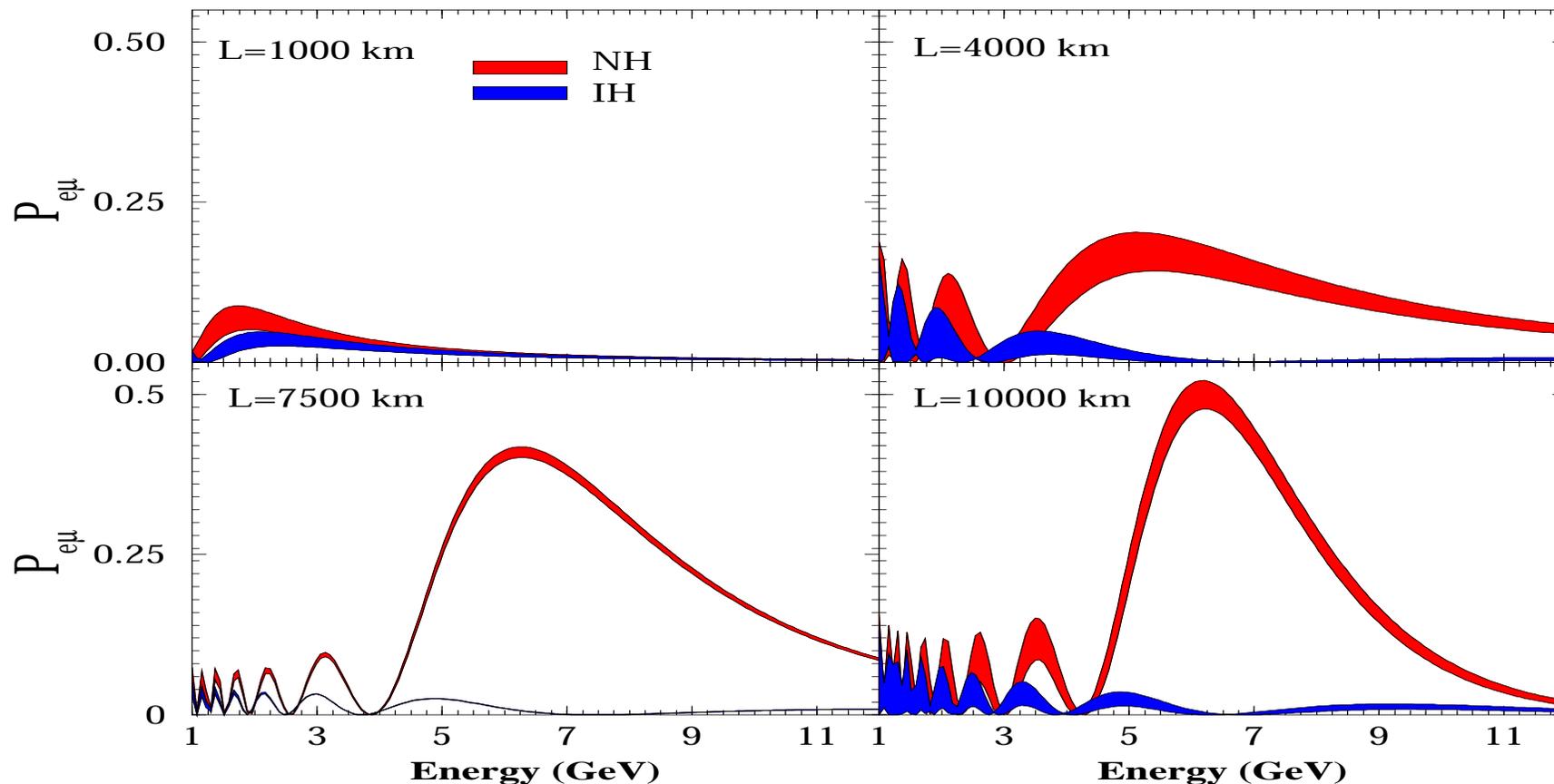
- $\sin(\hat{A}\Delta) \simeq 0 \Rightarrow \frac{1}{\sqrt{2}} G_F n_e L = \pi \Rightarrow L_{magic} \simeq 7690 \text{ km}$
- Independent of neutrino parameters and energy
- True for both NH and IH

Barger, Marfatia, Whisnant, hep-ph/0112119

Huber, Winter, hep-ph/0301257

Smirnov, hep-ph/0610198

$P_{e\mu}$ for NH and IH at long baselines

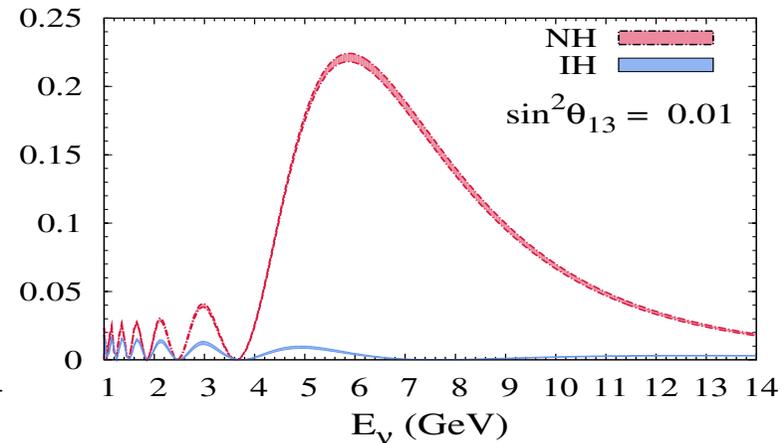
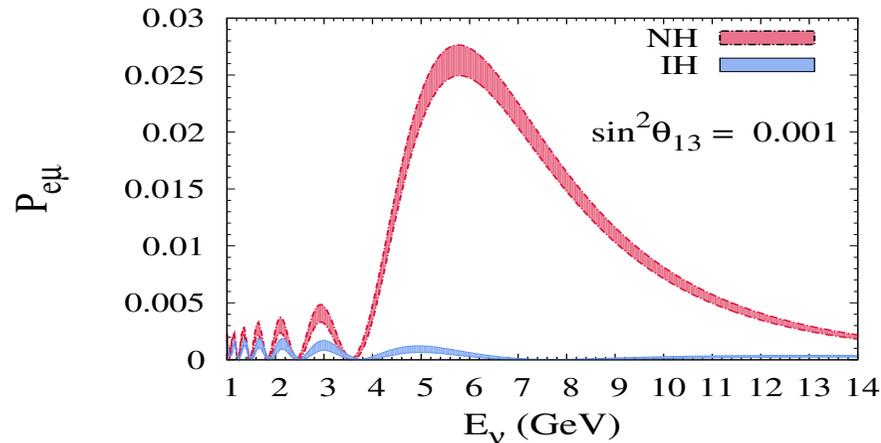


- At ~ 7500 km δ_{CP} dependence negligible
- $(\delta_{CP}, \text{sgn}(\Delta m_{\text{atm}}^2))$ degeneracies vanish
- Clean measurement of $\text{sgn}(\Delta m_{\text{atm}}^2)$

Agarwalla, Choubey, Raychaudhuri, hep-ph/0610333

The Magic baseline: Absence of CP sensitivity

$$P_{e\mu} \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2}$$



⇒ No CP sensitivity

- Recommendation of International Design Study of Neutrino Factory Group: another experiment at 4000 km for δ_{CP} with $E_\mu = 25$ GeV
- Requires high acceleration of the muons, also $1/r^2$ fall in flux
- Can there be a single experiment at a shorter baseline and lower energy that can determine all the three parameters ?

Enters The Bi-Magic baseline

$$\begin{aligned}
 P_{e\mu} &\simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2} \\
 &+ \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos(\Delta - \delta_{CP}) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin(1-\hat{A})\Delta}{(1-\hat{A})} \\
 &+ \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

- The condition $\sin(\hat{A}\Delta) \simeq 0$ implies vanishing of δ_{CP} dependent terms
- If we instead make $\sin[(1 - \hat{A})\Delta] = 0$ the δ_{CP} dependent term can vanish.



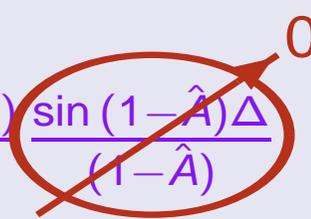
Enters The Bi-Magic baseline

$$P_{e\mu} \simeq \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2} + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos(\Delta - \delta_{CP}) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin(1-\hat{A})\Delta}{(1-\hat{A})} + \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}$$

- In that case $P_{e\mu} \approx \mathcal{O}(\alpha^2) \rightarrow$ small
- The condition $\sin(\hat{A}\Delta) \simeq 0$ is valid for both NH and IH
- But this condition depends on hierarchy



Enters The Bi-Magic baseline

$$\begin{aligned}
 P_{e\mu} \simeq & \sin^2 \theta_{23} \sin^2 2\theta_{13} \frac{\sin^2 (1-\hat{A})\Delta}{(1-\hat{A})^2} \\
 & + \alpha \sin 2\theta_{13} \sin 2\theta_{12} \sin 2\theta_{23} \cos(\Delta - \delta_{\text{CP}}) \frac{\sin(\hat{A}\Delta)}{\hat{A}} \frac{\sin(1-\hat{A})\Delta}{(1-\hat{A})} \\
 & + \alpha^2 \cos^2 \theta_{23} \sin^2 2\theta_{12} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$


• For IH $\hat{A} = -\hat{A}, \Delta = -\Delta$

• Magic condition depends on hierarchy

IH-NoCP

$$(1 + |\hat{A}|) \cdot |\Delta| = n\pi, n > 0$$

NH-NoCP

$$(1 - |\hat{A}|) \cdot |\Delta| = n\pi, n \neq 0$$

• Demand: Maximum hierarchy sensitivity

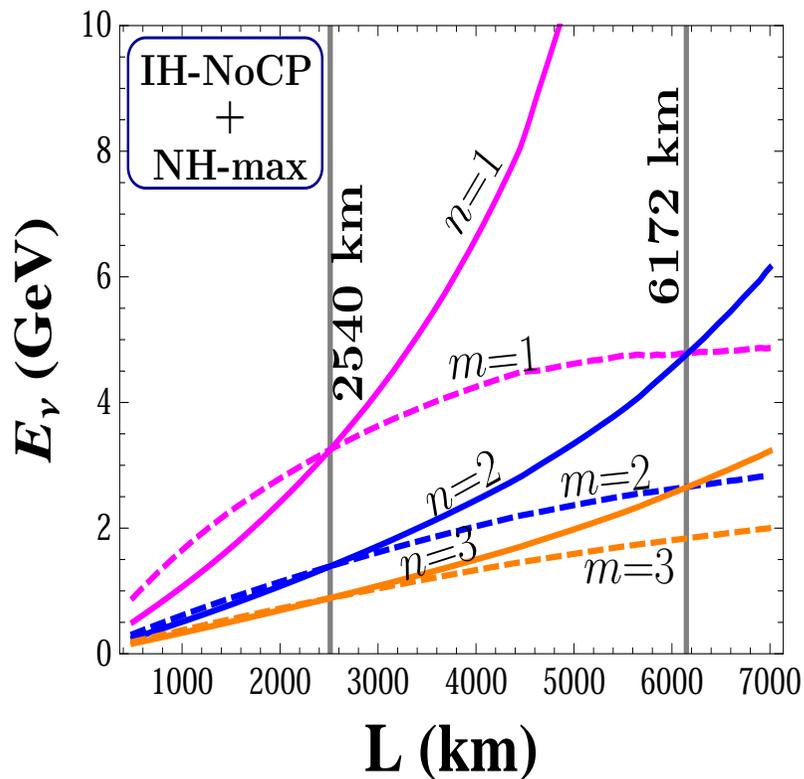
NH-max

$$(1 - |\hat{A}|) \cdot |\Delta| = (m - \frac{1}{2})\pi$$

IH-max

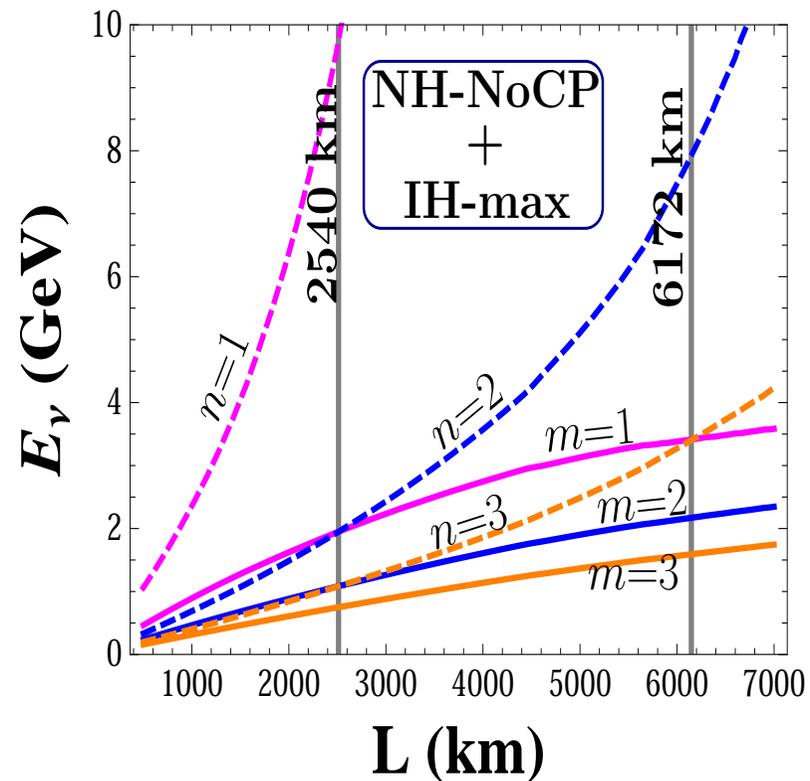
$$(1 + |\hat{A}|) \cdot |\Delta| = (m - \frac{1}{2})\pi$$

The Bi-Magic baseline



- $n = 1$ and $m = 1$
 $L \approx 2540$ km
 $E_\nu \equiv E_{IH} \approx 3.3$ GeV,

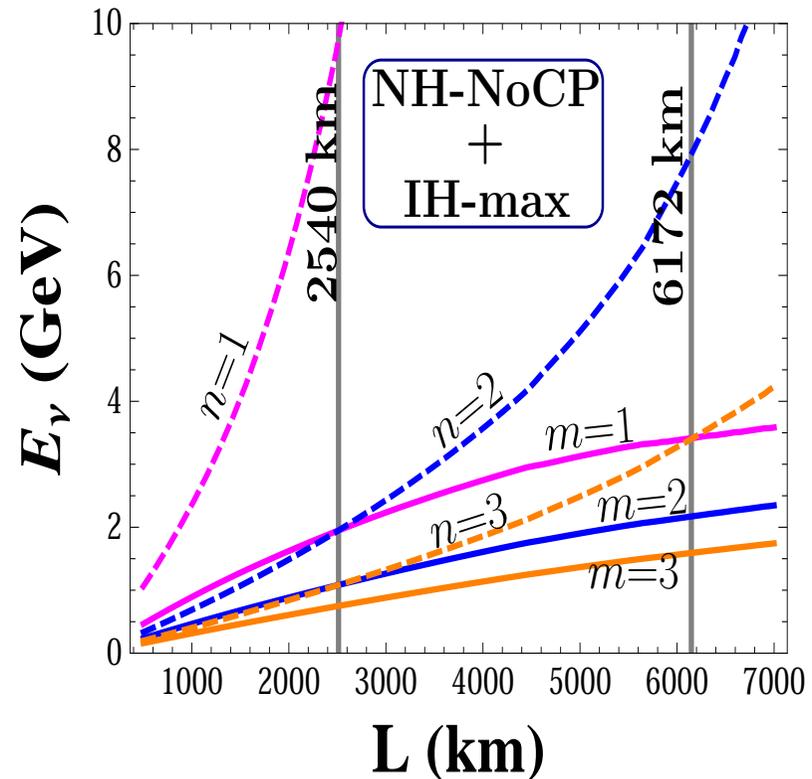
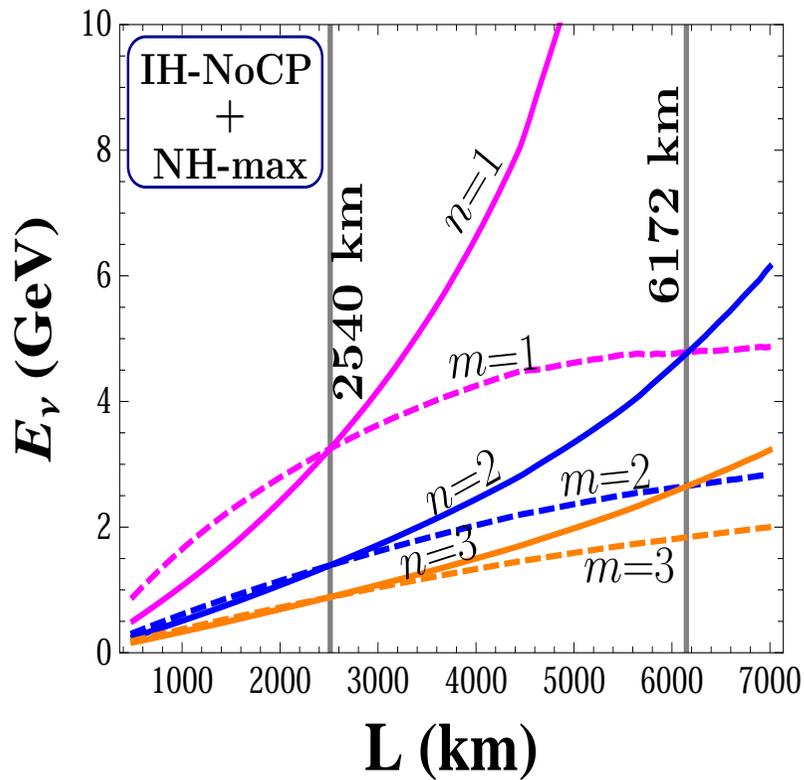
Raut, Singh, UmaShankar, 2009



- $n = 1$ and $m = 2$,
 $L \approx 2540$ km
 $E_\nu \equiv E_{NH} \approx 1.9$ GeV.

Dighe, Goswami, Ray, 2010

The Bi-Magic baseline



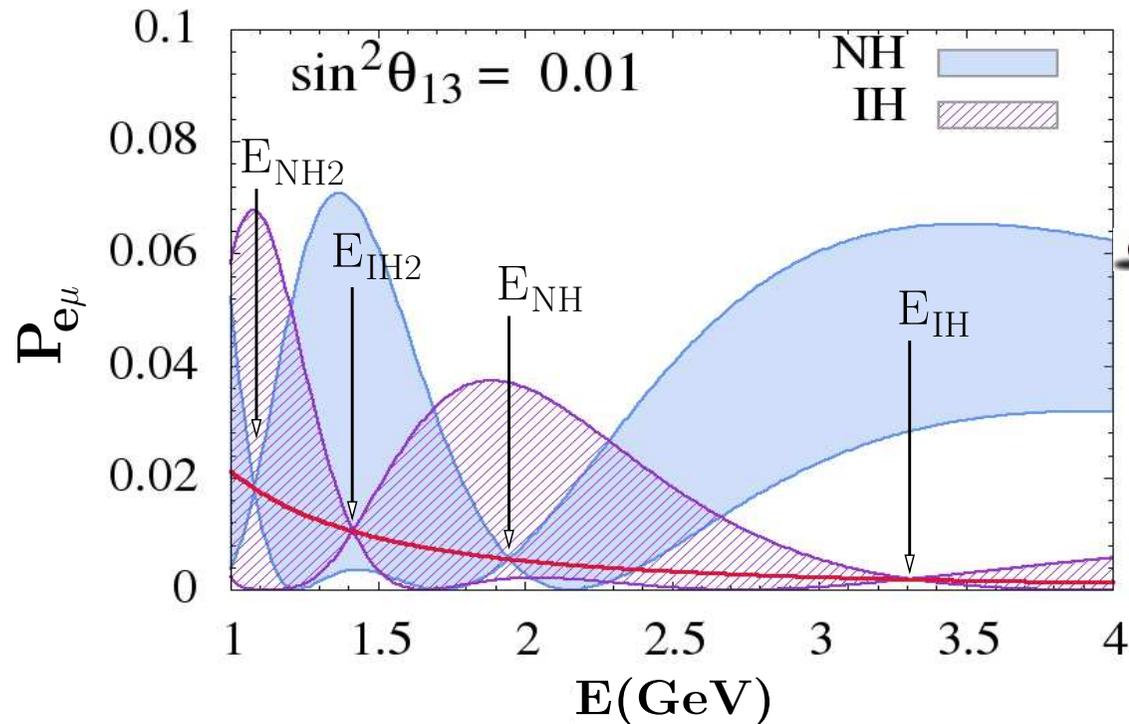
- Lowest **Bi-magic** Baseline \sim **2540 km**
- Higher $n, m \Rightarrow$ **Lower E_ν** , to satisfy no CP condition **low flux, low efficiency**
- More **bimagics** 6172 km, 8950 km, 106900 km

Baselines close Bi-Magic baseline

- CERN - PhyÅsalmi : 2288 km (LAGUNA)
- CERN - GranCanaria: 2780 km
- BNL- Homestake : 2540 km
- Fermilab-Icicle Creek : 2610 km
- Fermilab - SanJacinto: 2610 km

For a compilation of baselines from different accelerator facilities and underground laboratories see [Agarwalla et. al. arXiv:1012.1872 \[hep-ph\]](#)

$P_{e\mu}$ at the Bi-Magic baseline



• NH, IH interchanged for antineutrinos

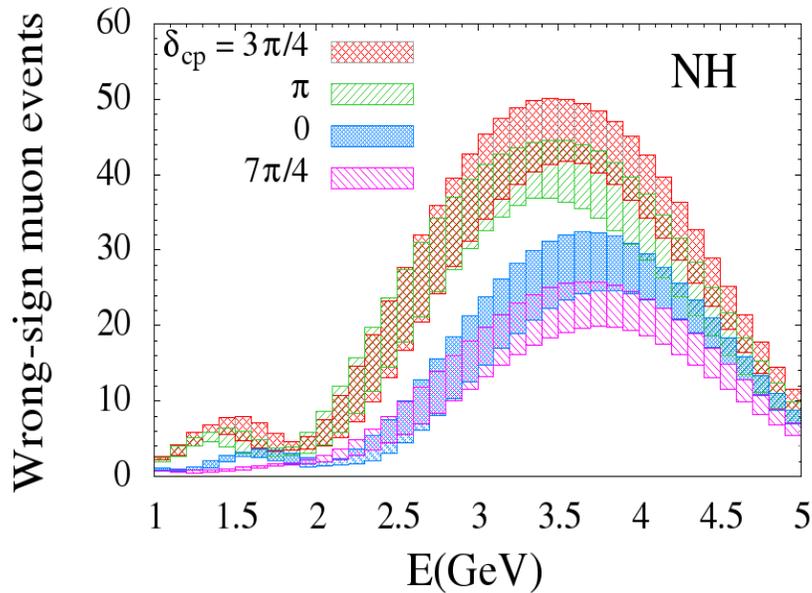
$$E_{IH} = 1.9 \text{ GeV}$$

- NH probability independent of δ_{CP} but δ_{CP} band in IH large
- δ_{CP} sensitivity for IH

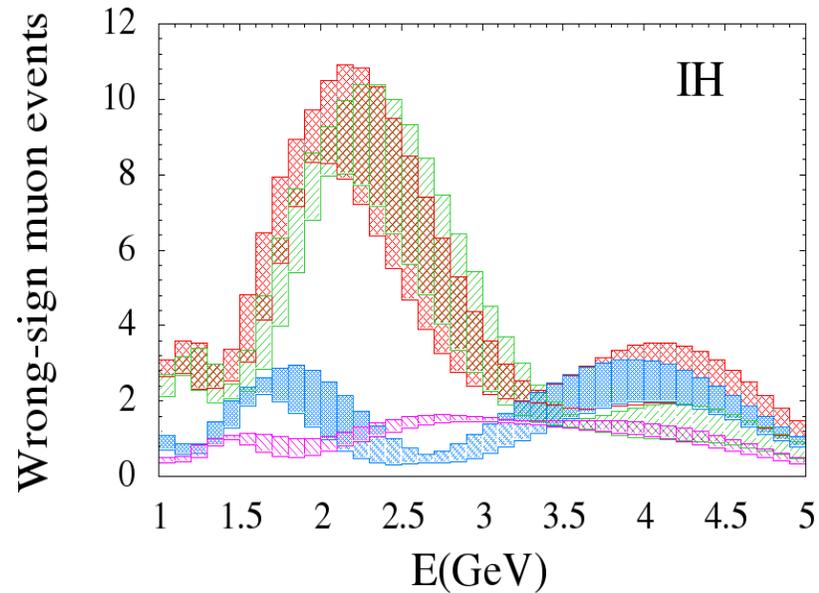
$$E_{IH} = 3.3 \text{ GeV}$$

- IH probability independent of δ_{CP} and non-overlapping with NH
- δ_{CP} sensitivity for NH

The Return of the CP sensitivity



$$E_{IH} = 3.3 \text{ GeV}$$



$$E_{NH} = 1.9 \text{ GeV}$$

$$P_{e\mu}(\text{IH}) \approx \mathcal{O}(\alpha^2)$$

$$P_{e\mu}(\text{NH}) \approx 18\alpha^2 s_{12}^2 c_{12}^2 c_{23}^2 + 9s_{13}^2 s_{23}^2 - 18\sqrt{2}\alpha s_{12} c_{12} s_{23} c_{23} s_{13} \cos(\delta_{CP} + \pi/4)$$

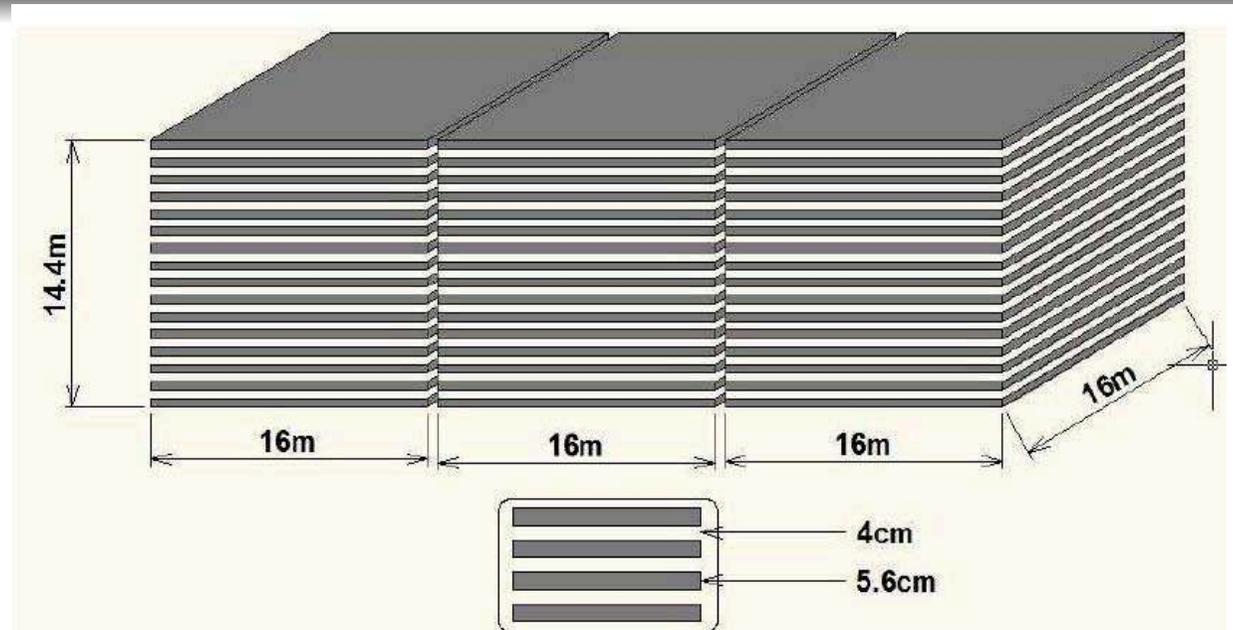
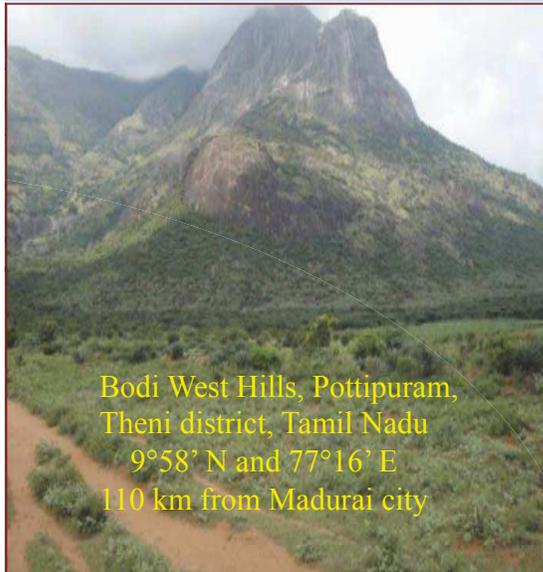
$$P_{e\mu}(\text{NH}) \approx \mathcal{O}(\alpha^2)$$

$$P_{e\mu}(\text{IH}) \approx 50\alpha^2 s_{12}^2 c_{12}^2 c_{23}^2 + \frac{25}{9} s_{13}^2 s_{23}^2 - \frac{50\sqrt{2}}{3} \alpha s_{12} c_{12} s_{23} c_{23} s_{13} \cos(\delta_{CP} + \pi/4)$$

$P_{e\mu}$ largest at $\delta_{CP} = 3\pi/4$, lowest at $\delta_{CP} = 7\pi/4$.

Can Atmospheric Neutrinos help ?

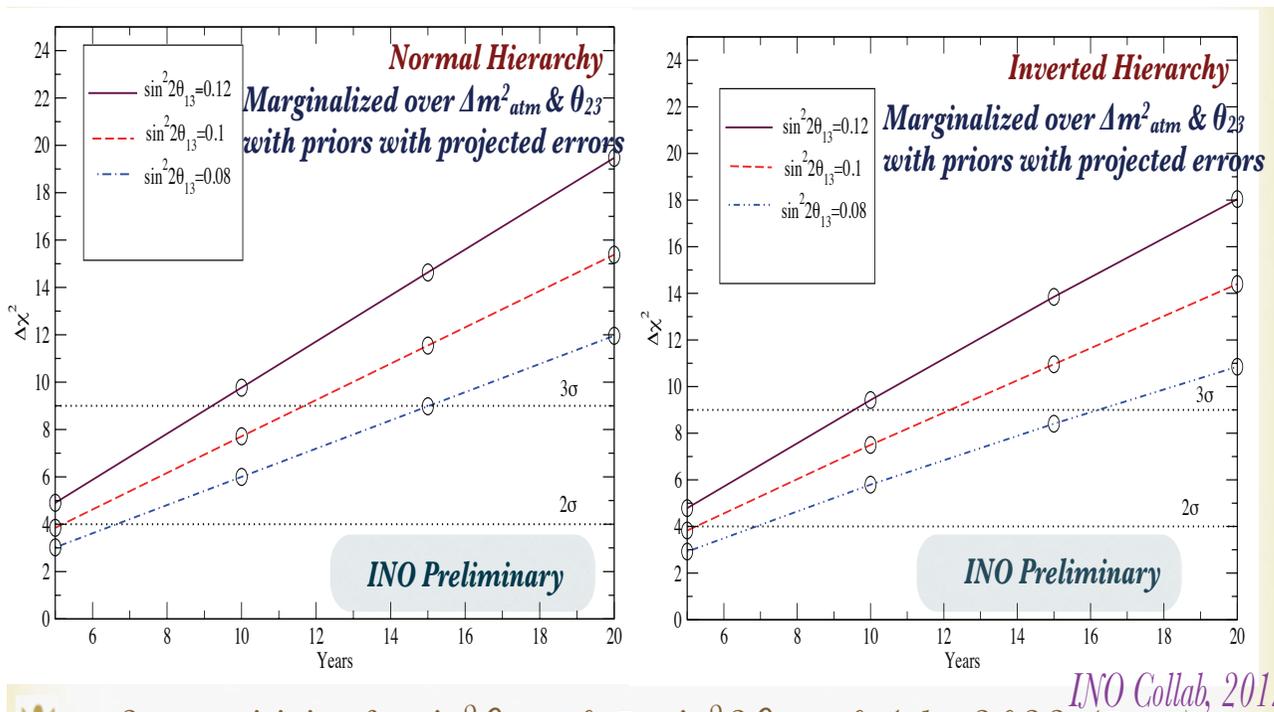
- Atmospheric neutrinos travel through a large path length in earth's matter
- Large measured $\theta_{13} \implies$ appreciable matter effect
- Three proposals
 - (i) HyperKamiokande – upgraded version of SuperK, sensitive to both electron and muons
 - (ii) India-Based Neutrino Observatory (INO) – sensitive to muons, have charge sensitivity
 - (iii) PINGU – Huge statistics



- 50 kton magnetized iron detector Sensitive to **muons**
- Good **Energy determination** from
 - Track **length**
 - Track **curvature** in a magnetic field
- **Directionality** from tracking and ns timing resolution
- **Charge identification** from track curvature in magnetic field
- **Hadron Shower reconstruction**
- Coming Soon in Bodi Hills, Tamilnadu

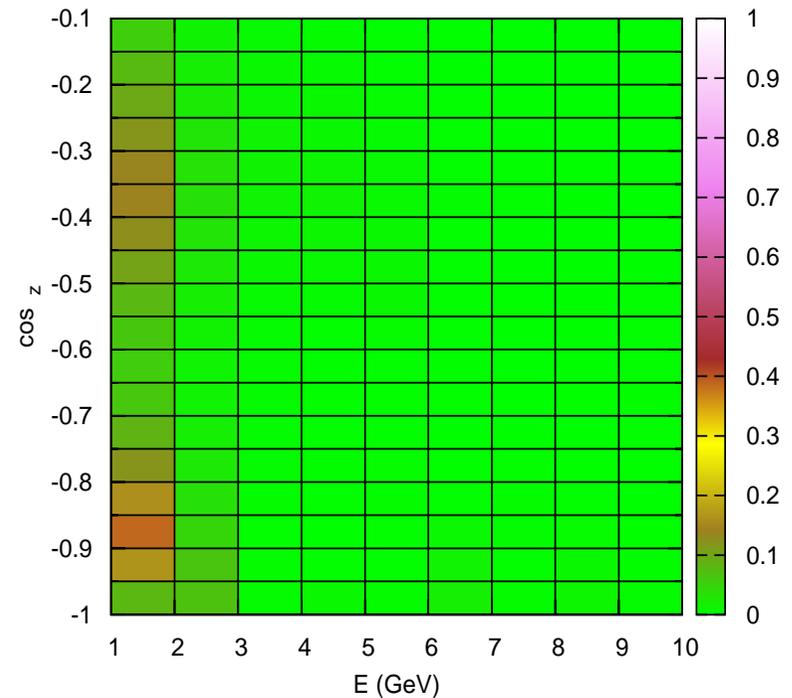
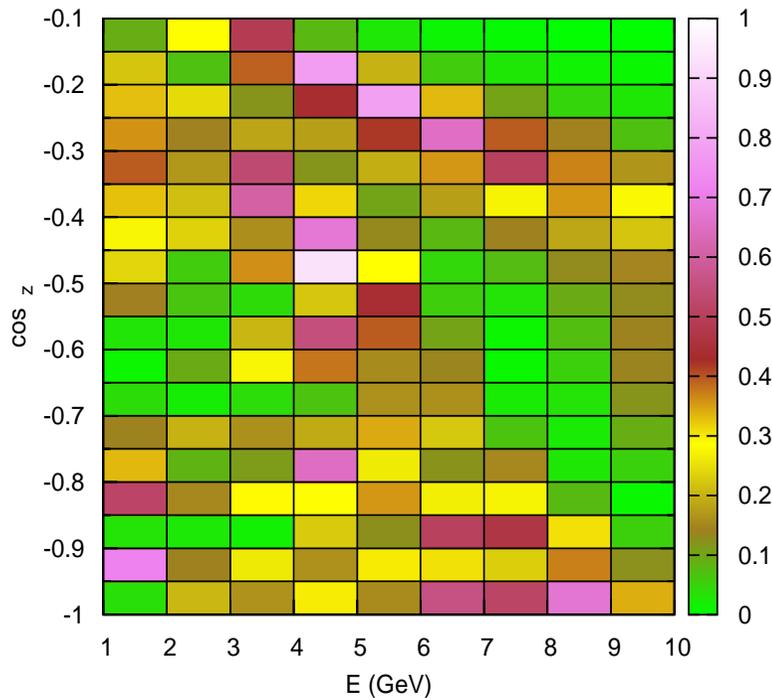
Hierarchy Sensitivity from INO simulation

- Using 50 Kton detector and events generated from NUANCE
- Using ICAL resolutions in Muon energy and Zenith Angle
- Using efficiency and Charge-ID from ICAL simulations



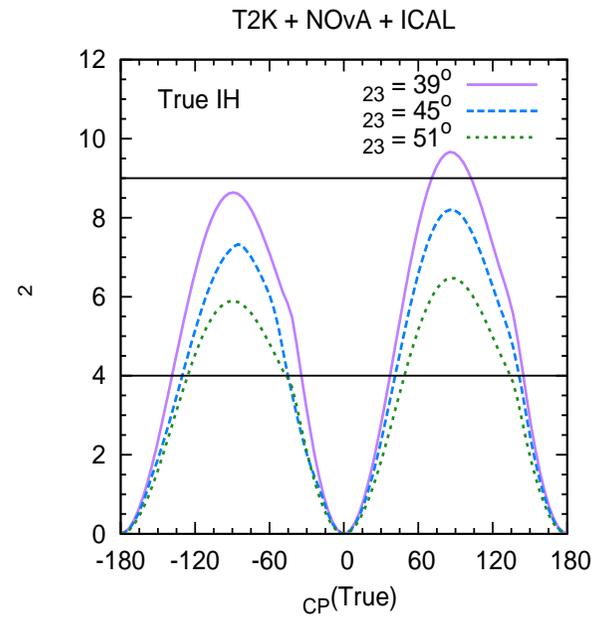
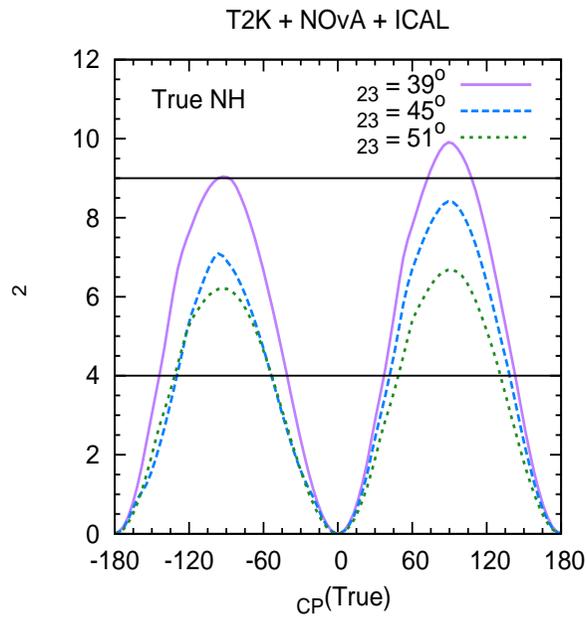
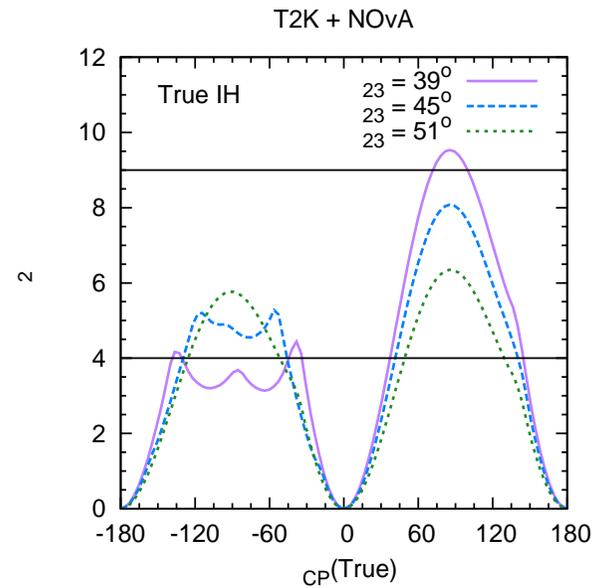
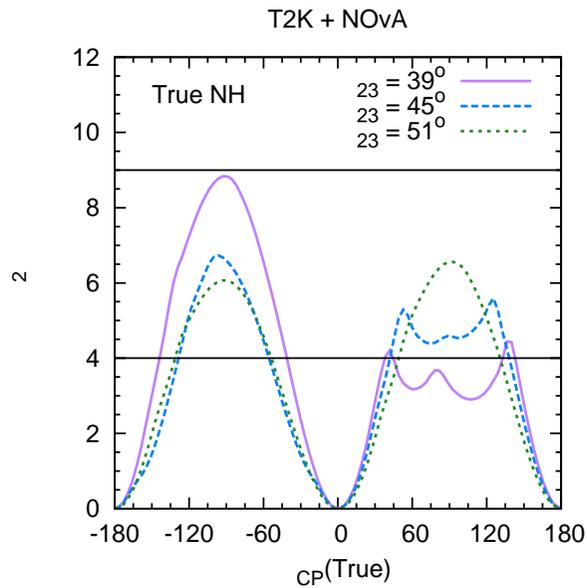
- For $\sin^2 2\theta_{13} = 0.1 \sim 2.7 \sigma$ sensitivity in 10 years

δ_{CP} sensitivity of Atmospheric Neutrinos

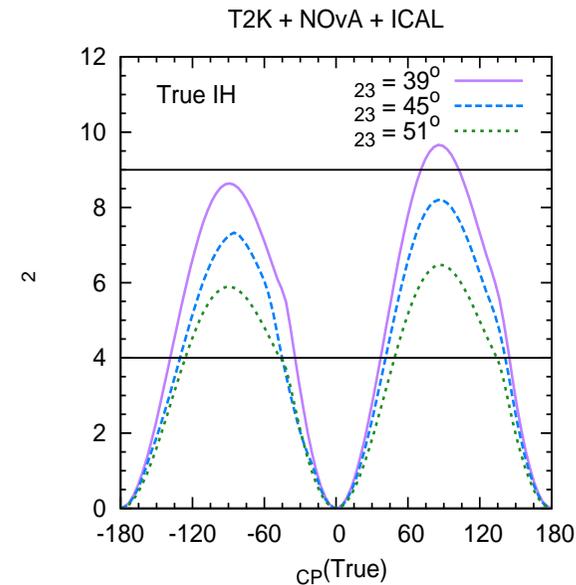
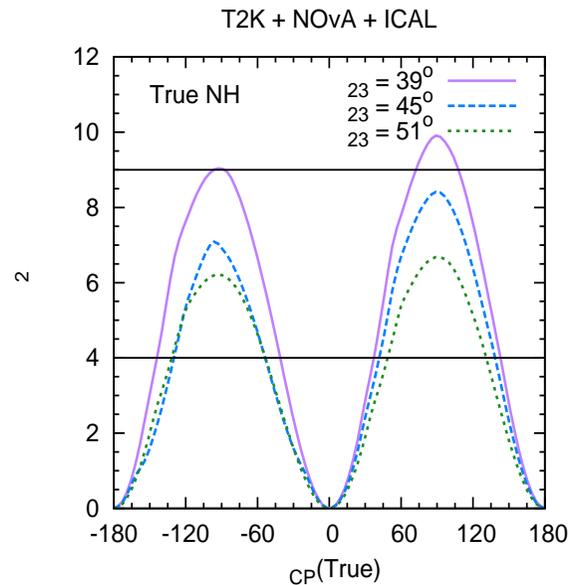


- $S_\mu + S_{\bar{\mu}} = (\Delta N_\mu)^2 / N_\mu(\text{avg})$,
- ΔN_μ is the maximum difference in events by varying δ_{CP} , $N_\mu(\text{avg})$ is the average number of events over all values of δ_{CP} .
- A measure of δ_{CP} -sensitivity in the $E - \cos\theta_z$ plane; the grid represents bins in energy and $\cos\theta_z$.
- $\sin^2 2\theta_{13} = 0.1$, $\sin^2 \theta_{23} = 0.5$ and NH.

$\text{NO}\nu\text{A}+\text{T2K}+\text{ICAL}$ CP violation discovery



NO ν A+T2K+ICAL CP violation discovery



- $\sin^2 2\theta_{13} = 0.1$, 500 kt yr exposure
- 3σ sensitivity possible in the wrong hierarchy region also.
- For unfavourable parameter values the first signature of CPV can come after including **ICAL data to T2K+NO ν A**
- Unfavourable for T2K+NO ν A :(\rightarrow favourable for INO :)