Inert 2HDM with local $U(1)_H$ gauge symmetry and related issues

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and arXiv:1405.2138

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Two Higgs doublet model

- Many high-energy models predict extra Higgs doublets.

  - SUSY, GUT, flavor symmetric models, etc.

- Two Higgs doublet model could be an effective theory of a high-energy theory.

- Two (or multi) Higgs doublet model itself is interesting.

  - Higgs physics (heavy Higgs, pseudoscalar, charged Higgs physics)

  - dark matter physics (one of Higgs scalar or extra fermions could be CDM.) Ma,PRD73;Barbieri,Hall,Rychkov,PRD74

  - baryon asymmetry of the Universe Shu,Zhang,PRL111

  - neutrino mass generation Kanemura,Matsui,Sugiyama,PLB727

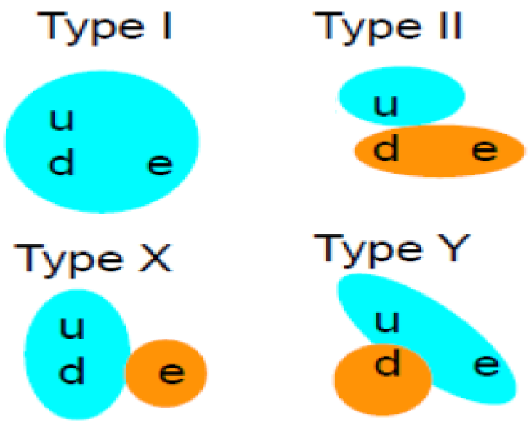
  - can resolve experimental anomalies (top A_{FB} at Tevatron, $B \to D(*) \tau \nu$ at BABAR) Ko,Omura,Yu,EPJC73;JHEP1303

2HDM with Z_2 symmetry (2HDMwZ_2)

- One of the simplest models to extend the SM Higgs sector.
- In general, flavor changing neutral currents (FCNCs) appear.
- A simple way to avoid the FCNC problem is to assign ad hoc Z_2 symmetry.

NFC criterion by Glashow and Weinberg (1977)

Type	H_1	H_2	U_R	D_R	E_R	N_R	Q_L, L
I	$+$	$-$	$+$	$+$	$+$	$+$	$+$
II	$+$	$-$	$+$	$-$	$-$	$+$	$+$
X	$+$	$-$	$+$	$+$	$-$	$-$	$+$
Y	$+$	$-$	$+$	$-$	$+$	$-$	$+$



Fermions of same electric charges get their masses from one Higgs VEV.

$$\mathcal{L} = \bar{L}_i (y^E_{1ij} H_1 + \cancel{y^E_{2ij} H_2}) E_{Rj} + \text{H.c.}$$

or vice versa

NO FCNC at tree level.

Generic problems of 2HDM

• It is well known that discrete symmetry could generate a domain wall problem when it is spontaneously broken.

• Usually the Z_2 symmetry is assumed to be broken softly by a dim-2 operator, $H_1^\dagger H_2$ term.

The softly broken Z_2 symmetric 2HDM potential

$$V = m_1^2 H_1^\dagger H_1 + m_2^2 H_2^\dagger H_2 - (m_{12}^2 H_1^\dagger H_2 + h.c.) + \frac{1}{2}\lambda_1 (H_1^\dagger H_1)^2 + \frac{1}{2}\lambda_2 (H_2^\dagger H_2)^2$$
$$+ \lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 (H_1^\dagger H_2)(H_2^\dagger H_1) + \frac{1}{2}\lambda_5 [(H_1^\dagger H_2)^2 + h.c.]$$

• the origin of the Z_2 symmetry and the softly breaking term?

Z_2 symmetry in 2HDM can be replaced by new $U(1)_H$ symmetry associated with Higgs flavors.

Type-I 2HDM

- Only one Higgs couples with fermions.

$$V_y = y_{ij}^U \overline{Q}_{Li} \tilde{H}_1 U_{Rj} + y_{ij}^D \overline{Q}_{Li} H_1 D_{Rj} + y_{ij}^E \overline{L}_i H_1 E_{Rj} + y_{ij}^N \overline{L}_i \tilde{H}_1 N_{Rj}$$

- anomaly free $U(1)_H$ without extra fermions except RH neutrinos.

U_R	D_R	Q_L	L	E_R	N_R	H_1
u	d	$\frac{(u+d)}{2}$	$\frac{-3(u+d)}{2}$	$-(2u+d)$	$-(u+2d)$	$\frac{(u-d)}{2}$

2 parameters

- In other Types of 2HDMs, extra fermions are required in order to cancel gauge anomaly.

→ one of extra fermions can be a candidate for the cold dark matter.

Type-I 2HDM

- Only one Higgs couples with fermions.

$$V_y = y_{ij}^U \overline{Q}_{Li} \tilde{H}_1 U_{Rj} + y_{ij}^D \overline{Q}_{Li} H_1 D_{Rj} + y_{ij}^E \overline{L}_i H_1 E_{Rj} + y_{ij}^N \overline{L}_i \tilde{H}_1 N_{Rj}$$

- anomaly free $U(1)_H$ without no extra fermions except RH neutrinos.

U_R	D_R	Q_R	L	E_R	N_R	H_1	Type
u	d	$\frac{(u+d)}{2}$	$\frac{-3(u+d)}{2}$	$-(2u+d)$	$-(u+2d)$	$\frac{(u-d)}{2}$	
0	0	0	0	0	0	0	$h_2 \neq 0$
$1/3$	$1/3$	$1/3$	-1	-1	-1	0	$U(1)_{B-L}$
1	-1	0	0	-1	1	1	$U(1)_R$
$2/3$	$-1/3$	$1/6$	$-1/2$	-1	0	$1/2$	$U(1)_Y$

Ko,Omura,Yu, PLB717,202(2013)

- SM fermions are $U(1)_H$ singlets.
- Z_H is fermiophobic and Higgphilic.

$U(1)_H$ symmetry beomes Dark Gauge Symmetry

Type-II 2HDM

• H_1 couples to the up-type fermions, while H_2 couples to the down-type fermions.

$$V_y = y_{ij}^U \bar{Q}_{Li} \tilde{H}_1 U_{Rj} + y_{ij}^D \bar{Q}_{Li} H_2 D_{Rj} + y_{ij}^E \bar{L}_i H_2 E_{Rj} + y_{ij}^N \bar{L}_i \tilde{H}_1 N_{Rj}$$

U_R	D_R	Q_L	L	E_R	N_R	H_1	H_2
u	0	0	0	0	u	u	0

• Requires extra chiral fermions for cancellation of gauge anomaly.

$U(1)_H^3$

$U(1)_Y U(1)_H^2$

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)_H$
q_{Li}	3	1	2/3	$\hat{Q}_L = u + \hat{Q}_R$
q_{Ri}	3	1	2/3	$\hat{Q}_R$
n_{Li}	1	1	0	$\hat{n}_L = u + \hat{n}_R$
n_{Ri}	1	1	0	$\hat{n}_R$

+ Two vector-like pairs of $SU(2)_L$

Mixing between new chiral fermions and SM fermions is prohibited by $U(1)_H$ charge assignment.

One of extra fermions could be a candidate for CDM.

Type-II 2HDM

• H_1 couples to the up-type fermions, while H_2 couples to the down-type fermions.

$$V_y = y_{ij}^U \bar{Q}_{Li} \tilde{H}_1 U_{Rj} + y_{ij}^D \bar{Q}_{Li} H_2 D_{Rj} + y_{ij}^E \bar{L}_i H_2 E_{Rj} + y_{ij}^N \bar{L}_i \tilde{H}_1 N_{Rj}$$

U_R	D_R	Q_L	L	E_R	N_R	H_1	H_2
u	0	0	0	0	u	u	0

• Requires extra chiral fermions for cancellation of gauge anomaly.

for example, $E_6 \to SO(10) \times U(1)_\psi \to SU(5) \times U(1)_\chi \times U(1)_\psi$.

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)_H$	$U(1)_\psi$	$U(1)_\chi$	$U(1)_\eta$
Q^i	3	2	1/6	$-1/3$	1	-1	-2
U_R^i	3	1	2/3	2/3	-1	1	2
D_R^i	3	1	$-1/3$	$-1/3$	-1	-3	-1
L_i	1	2	$-1/2$	0	1	3	1
E_R^i	1	1	-1	0	-1	1	2
N_R^i	1	1	0	1	-1	5	5
H_1	1	2	1/2	0	2	2	-1
H_2	1	2	1/2	1	-2	2	4

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)_H$	$U(1)_\psi$	$U(1)_\chi$	$U(1)_\eta$
q_L^i	3	1	$-1/3$	2/3	-2	2	4
q_R^i	3	1	$-1/3$	$-1/3$	2	2	-1
l_L^i	1	2	$-1/2$	0	-2	-2	1
l_R^i	1	2	$-1/2$	-1	2	-2	-4
n_L^i	1	1	0	-1	4	0	-5

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)_H$	$U(1)_\psi$	$U(1)_\chi$	$U(1)_\eta$
Φ	1	1	0	1	-4	0	5

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U_R	D_R	Q_L	L	E_R	N_R	H_1	H_2
u	0	0	0	0	u	u	0

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Q^i	3	2	1/6	−1/3	1	−1	−2
U_R^i	3	1	2/3	2/3	−1	1	2
D_R^i	3	1	−1/3	−1/3	−1	−3	−1
L_i	1	2	−1/2	0	1	3	1
E_R^i	1	1	−1	0	−1	1	2
N_R^i	1	1	0	1	−1	5	5
H_1	1	2	1/2	0	2	2	−1
H_2	1	2	1/2	1	−2	2	4

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)_H$	$U(1)_\psi$	$U(1)_\chi$	$U(1)_\eta$
q_L^i	3	1	−1/3	2/3	−2	2	4
q_R^i	3	1	−1/3	−1/3	2	2	−1
l_L^i	1	2	−1/2	0	−2	−2	1
l_R^i	1	2	−1/2	−1	2	−2	−4
n_L^i	1	1	0	−1	4	0	−5

	$SU(3)$	$SU(2)$	$U(1)_Y$	$U(1)_H$	$U(1)_\psi$	$U(1)_\chi$	$U(1)_\eta$
Φ	1	1	0	1	−4	0	5

Leptophoobic E6 by JLRosner

Constraints

- experimental and theoretical constraints

Perturbativity
Unitarity
Vacuum stability

$m_h \sim 126$ GeV

$|m_{H^+} - m_A|$

$|m_{H^+} - m_H|$

$\sin(\beta - \alpha)$

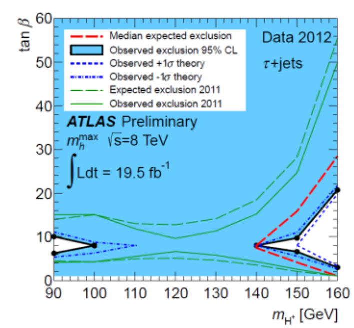
$\tan\beta$

m_{H^+}

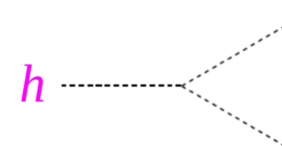
SM-like Higgs

m_H

EWPOs
small mass differences
requied

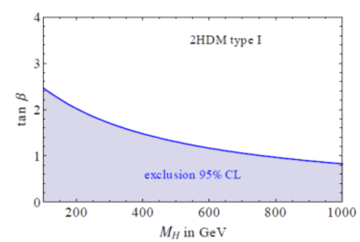


Invisible Higgs decay



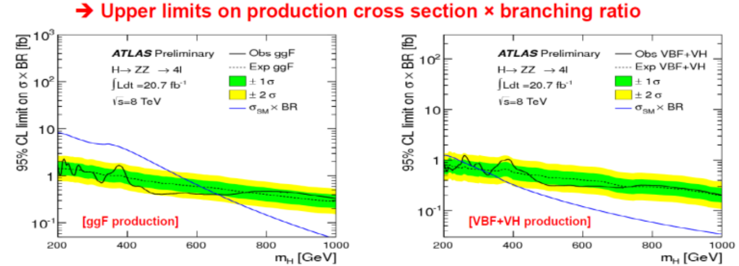
Exotic top decay

$b \to s\gamma$



Hermann,Misiak,Steinhauser,
JHEP1211 (2012) 036

→ Upper limits on production cross section × branching ratio



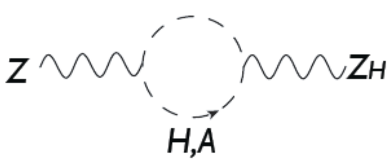
Heavy Higgs search at LHC

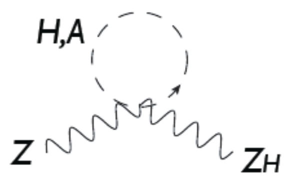
Z-Z_H mixing

- tree-level mixing ($v_i \neq 0$)

$$\Delta M^2_{ZZ_H} = -\frac{\hat{M}_Z}{v} g_H \sum_{i=1}^{2} q_{H_i} v_i^2 . \qquad \tan 2\xi = \frac{2\Delta M^2_{ZZ_H}}{\hat{M}^2_{Z_H} - \hat{M}^2_Z}$$

- loop-level mixing ($v_1=0$, $v_2 \neq 0$)





$$\kappa_Z = \frac{q_H g_H e c_W}{16\pi^2 s_W}\left\{\frac{1}{3}\ln\left(\frac{m_A^2}{m_{H^+}^2}\right) - \frac{1}{6}\frac{m_A^2 - m_H^2}{m_A^2}\right\},$$

$$\kappa_\gamma = \frac{q_H g_H e}{16\pi^2}\left\{\frac{1}{3}\ln\left(\frac{m_A^2}{m_{H^+}^2}\right) - \frac{1}{6}\frac{m_A^2 - m_H^2}{m_A^2}\right\},$$

$$\Delta M^2_{Z_H Z} = -\frac{q_H g_H e}{32\pi^2 s_W c_W}(m_A^2 - m_H^2).$$

$$-\frac{\kappa_Z}{2} F_Z^{\mu\nu} F_{H\mu\nu} - \frac{\kappa_\gamma}{2} F_\gamma^{\mu\nu} F_{H\mu\nu} + \Delta M^2_{Z_H Z} \hat{Z}^\mu \hat{Z}_{H\mu}$$

The mixing can appear because of $SU(2)_L$ $U(1)_Y$ breaking effects.

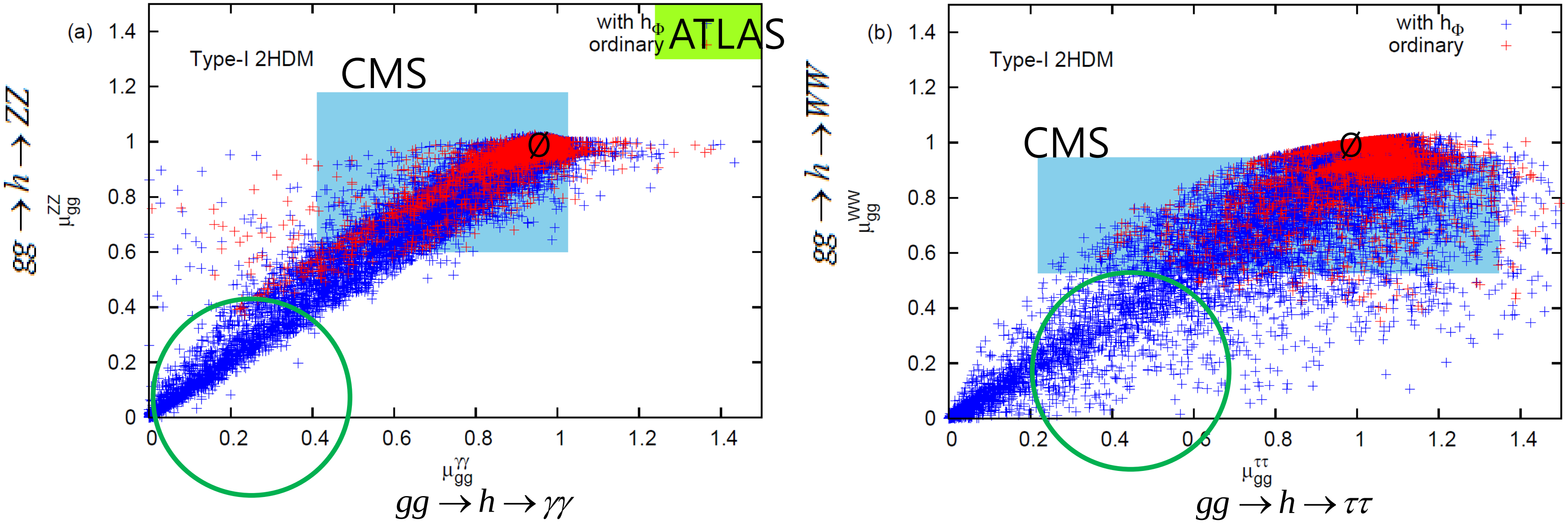
- collider bound depends on the $U(1)_H$ charge assignment.

- In the fermiophobic Z_H case, the Z_H boson can be produced through the Z-Z_H mixing and the bound for the mixing angle is

$$\sin\xi \text{ d } O(10^{-2}) \sim O(10^{-3})$$

Type-I 2HDM with h_Φ

- Assume the Z_H boson is heavy.



- consistent with CMS in the 1σ level while consistent with ATLAS in the 2σ.
- difficult to distinguish because the current experimental values are consistent with the SM prediction.
- essential to discover the extra scalar bosons and the new gauge boson.

Ko,Omura,Yu, JHEP1401 (2014) 016

Inert Doublet Model (IDMwZ_2)

- a 2HDM ~ one of the simplest extension

- One of Higgs doublets does not develop VEV and exact Z_2 symmetry is imposed. (Still there is a stability issue with Z_2)

- The new Higgs doublet does not participate in the EW symmetry breaking.

- Under the Z_2 symmetry, SM particles are even, but the new Higgs doublet is odd.

- Viable DM candidate

$$H_1 = \begin{pmatrix} H^+ \\ \frac{1}{\sqrt{2}}(H + iA) \end{pmatrix}, \quad H_2 = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}}(v + h + iG^0) \end{pmatrix}$$

DM candidates

SM-like Higgs

Inert Doublet Model (IDMwZ_2)

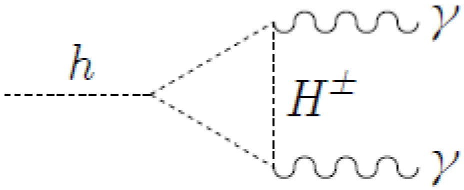
• CP-conserving potential

forbidden by the Z_2 symmetry

$$V = \mu_1 (H_1^\dagger H_1) + \mu_2 (H_2^\dagger H_2) - \mu_{12} (H_1^\dagger H_2 + \text{h.c.}) + \frac{\lambda_1}{2}(H_1^\dagger H_1)^2 + \frac{\lambda_2}{2}(H_2^\dagger H_2)^2$$
$$+\lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 |H_1^\dagger H_2|^2 + \frac{\lambda_5}{2}\{(H_1^\dagger H_2)^2 + h.c.\}.$$

• Type-I Yukawa interactions ~ only H_2 couples to the SM fermions.

• The h decay to two photons receives additional contribution through charged Higgs loop.



• H,A,$H^\pm$ ~ do not couple to SM fermions at tree level.

Inert Double Model (IDMwU$(1)_H$)

- We replace the Z_2 symmetry by U(1) gauge symmetry.
- A SM-singlet Phi has to be added.
- Without Phi, Z_H boson becomes massless.

$$V = (m_1^2 + \tilde{\lambda}_1 |\Phi|^2)(H_1^\dagger H_1) + (m_2^2 + \tilde{\lambda}_2 |\Phi|^2)(H_2^\dagger H_2) - (m_{12}^2 H_1^\dagger H_2 + \mathrm{h.c.})$$
$$+ \frac{\lambda_1}{2}(H_1^\dagger H_1)^2 + \frac{\lambda_2}{2}(H_2^\dagger H_2)^2 + \lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 |H_1^\dagger H_2|^2$$
$$+ \frac{\lambda_5}{2}\{(H_1^\dagger H_2)^2 + h.c.\} + m_\Phi^2 |\Phi|^2 + \lambda_\Phi |\Phi|^4$$

- Phi breaks the $U(1)_H$ symmetry while H_2 breaks the EW symmetry.
- The remnant symmetry of $U(1)_H$ is the origin of the exact Z_2 symmetry.

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forbidden
by the Z_2 symmetry

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$$+ \frac{\lambda_1}{2}(H_1^\dagger H_1)^2 + \frac{\lambda_2}{2}(H_2^\dagger H_2)^2 + \lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 |H_1^\dagger H_2|^2$$

$$+ \frac{\lambda_5}{2}\{(H_1^\dagger H_2)^2 + h.c.\} + m_\Phi^2 |\Phi|^2 + \lambda_\Phi |\Phi|^4$$

forbidden by the $U(1)_H$ symmetry ($q_{H_2}=0, q_{H_1} \neq 0$)

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- The remnant symmetry of $U(1)_H$ is the origin of the exact Z_2 symmetry.

Inert Double Model (IDMwU(1)$_H$)

- IDM + SM-singlet Phi.

forbidden by the Z_2 symmetry

$$V = (m_1^2 + \tilde{\lambda}_1 |\Phi|^2)(H_1^\dagger H_1) + (m_2^2 + \tilde{\lambda}_2 |\Phi|^2)(H_2^\dagger H_2) - (m_{12}^2 H_1^\dagger H_2 + \mathrm{h.c.})$$

$$+ \frac{\lambda_1}{2}(H_1^\dagger H_1)^2 + \frac{\lambda_2}{2}(H_2^\dagger H_2)^2 + \lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 |H_1^\dagger H_2|^2$$

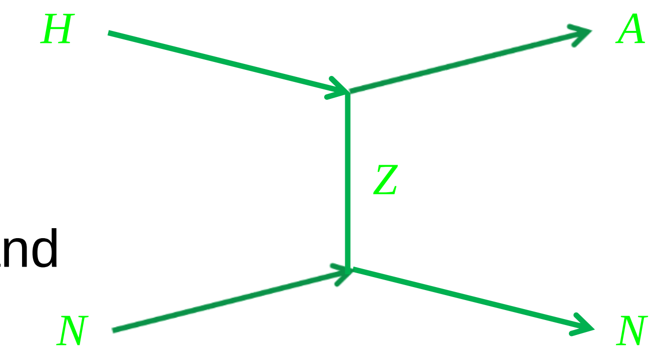
$$+ \frac{\lambda_5}{2}\{(H_1^\dagger H_2)^2 + h.c.\} + m_\Phi^2 |\Phi|^2 + \lambda_\Phi |\Phi|^4$$

forbidden by the U(1)$_H$ symmetry ($q_{H_2}=0, q_{H_1} \neq 0$)

- Without λ_5, H and A are degenerate.

$$m_A = \sqrt{m_H^2 - \lambda_5 v^2}$$

- Direct searches for DM at XENON100 and LUX exclude this degenerate case.



Inert Double Model (IDMwU(1)$_H$)

- IDM + SM-singlet Phi.

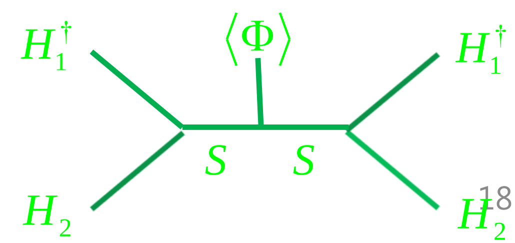
forbidden by the Z_2 symmetry

$$V = (m_1^2 + \tilde{\lambda}_1 |\Phi|^2)(H_1^\dagger H_1) + (m_2^2 + \tilde{\lambda}_2 |\Phi|^2)(H_2^\dagger H_2) - (m_{12}^2 H_1^\dagger H_2 + \mathrm{h.c.})$$
$$+ \frac{\lambda_1}{2}(H_1^\dagger H_1)^2 + \frac{\lambda_2}{2}(H_2^\dagger H_2)^2 + \lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 |H_1^\dagger H_2|^2$$
$$+ \left\{ c_l \left(\frac{\Phi}{\Lambda}\right)^l (H_1^\dagger H_2)^2 + h.c. \right\} + m_\Phi^2 |\Phi|^2 + \lambda_\Phi |\Phi|^4$$

- The λ_5 term can effectively be generated by a higher-dimensional operator.
- It could be realized by introducing a singlet S charged under U(1)$_H$ with $q_S = q_{H_1}$.

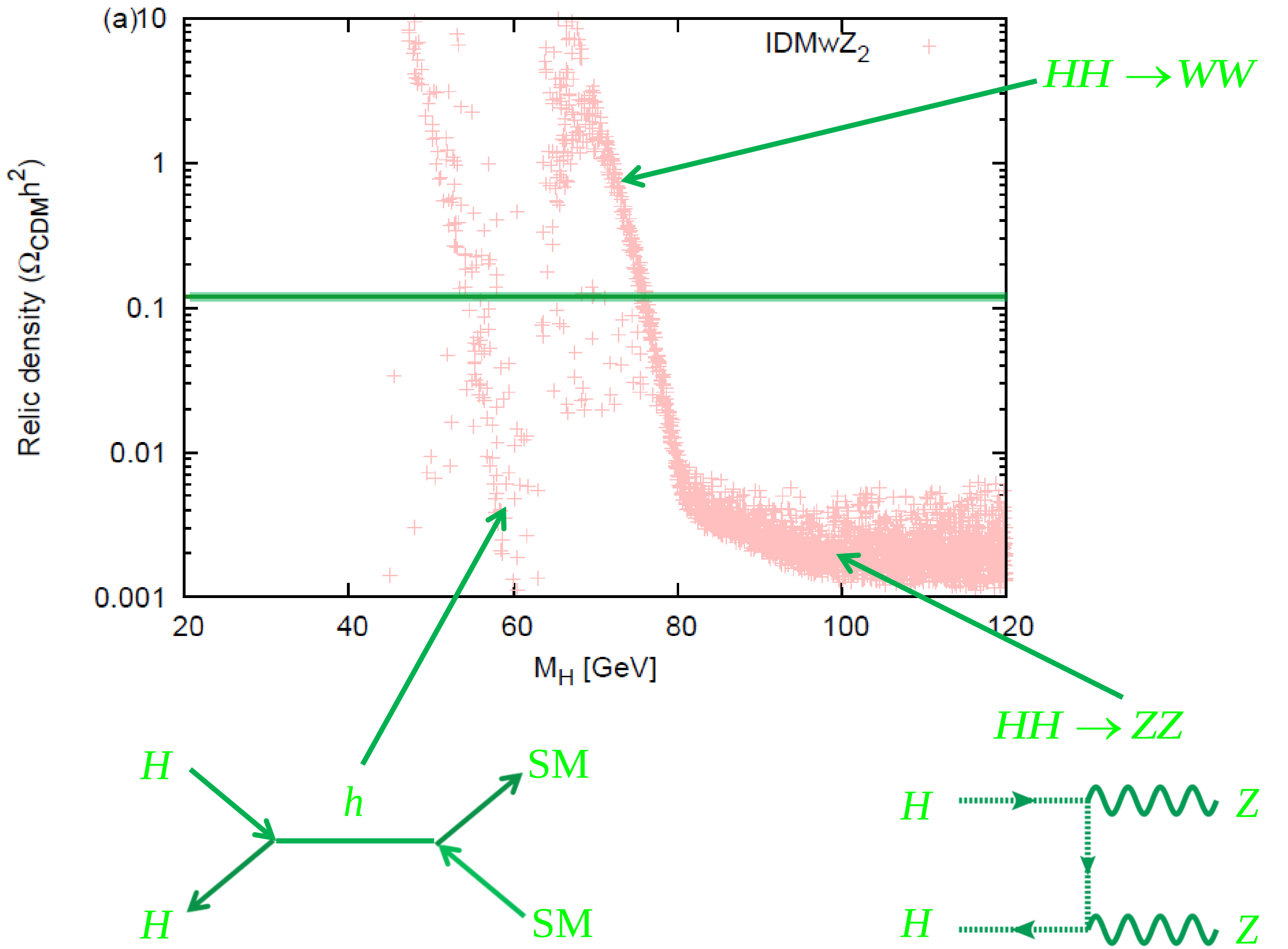
$$V_\Phi(|\Phi|^2, |S|^2) + V_H(H_i, H_i^\dagger) + \lambda_S(\Phi) S^2 + \lambda_H(S) H_1^\dagger H_2 + h.c..$$

$$\lambda_H = \lambda_H^0 S \qquad \lambda_5 \sim \frac{(\lambda_H^0)^2}{2} \frac{\Delta m^2}{m_{Re(S)}^2 m_{Im(S)}^2},$$



Relic density (low mass)

$$\Omega_{\text{CDM}}h^2 = 0.1199 \pm 0.0027$$



+ IDMwZ$_2$

LUX bound is satisfied.

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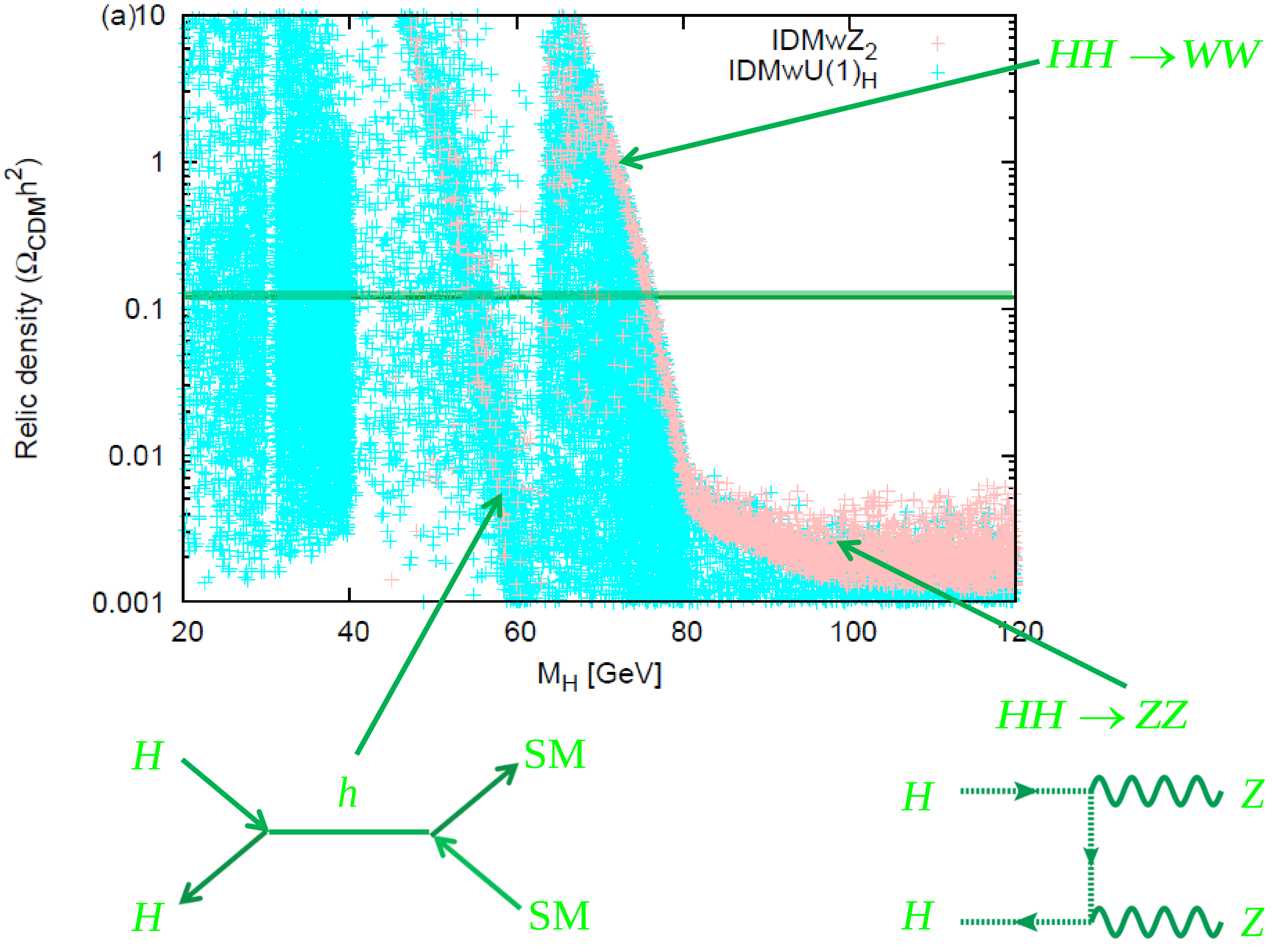

+ $IDMwZ_2$
+ $IDMwU(1)_H$

LUX bound is satisfied.

Relic density (low mass)

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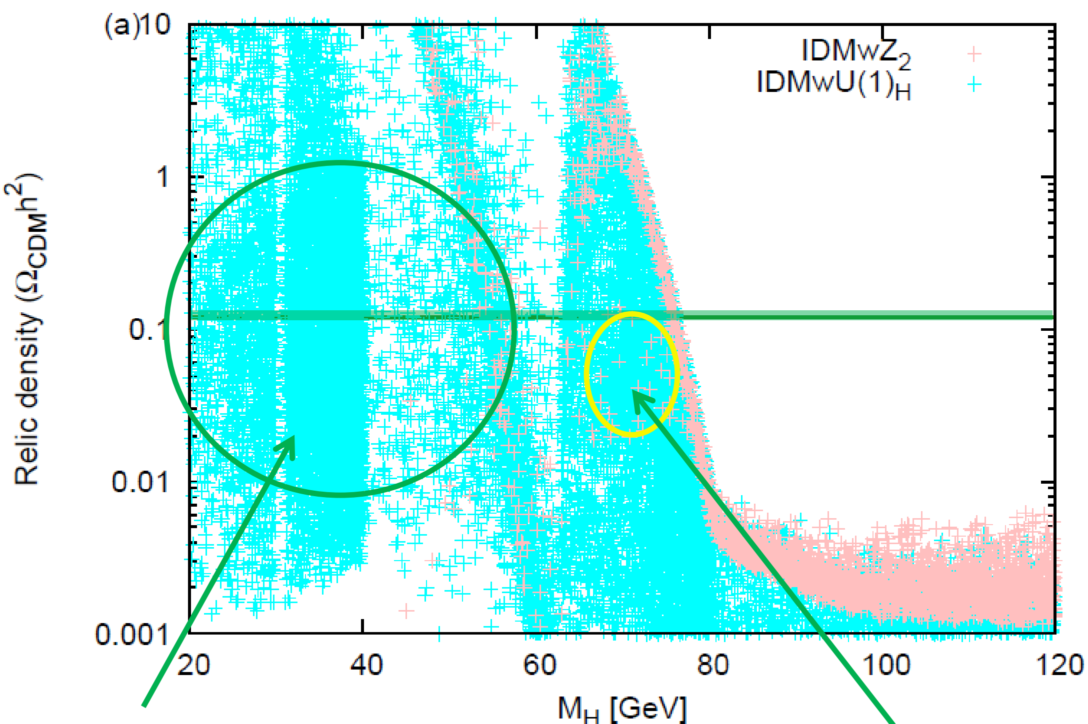
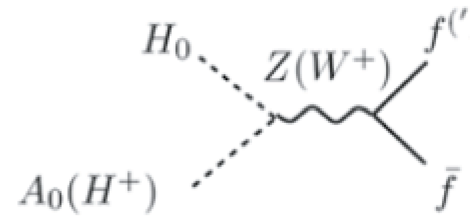


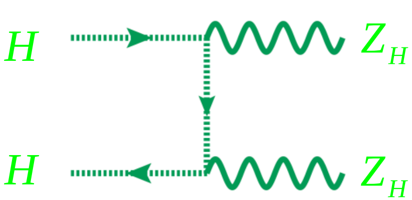
+ IDMwZ_2
+ IDMwU(1)$_H$

LUX bound is satisfied.

Co-annihilation



$HH \to Z_H Z_H, ZZ_H$

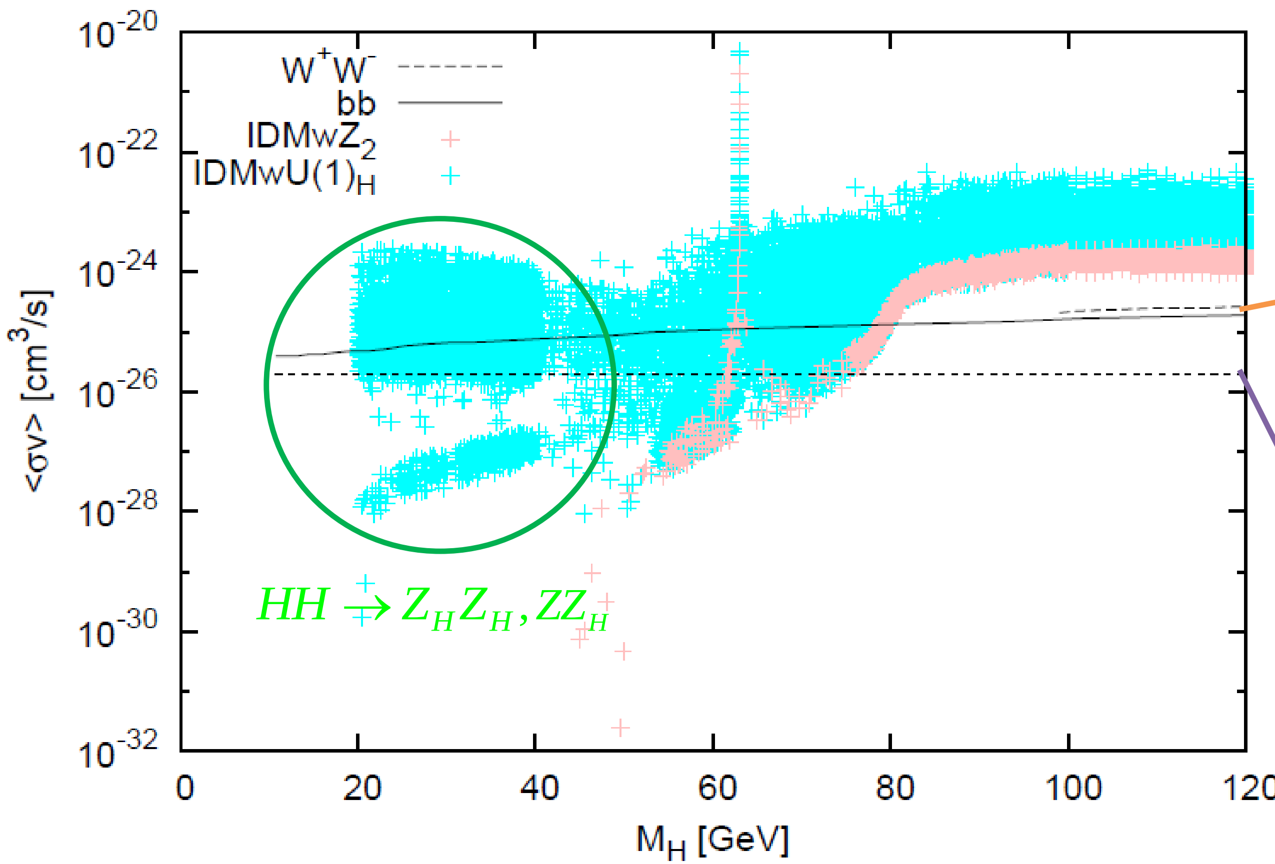


$HA, HH^{\pm} \to \mathrm{SM} + \mathrm{SM}^{(\prime)}$

$H^+H^- \to A + Z_H, Z + Z_H, \ldots$

Could accommodate the gamma ray excess from GC (work in preparation)

Indirect searches (low mass)



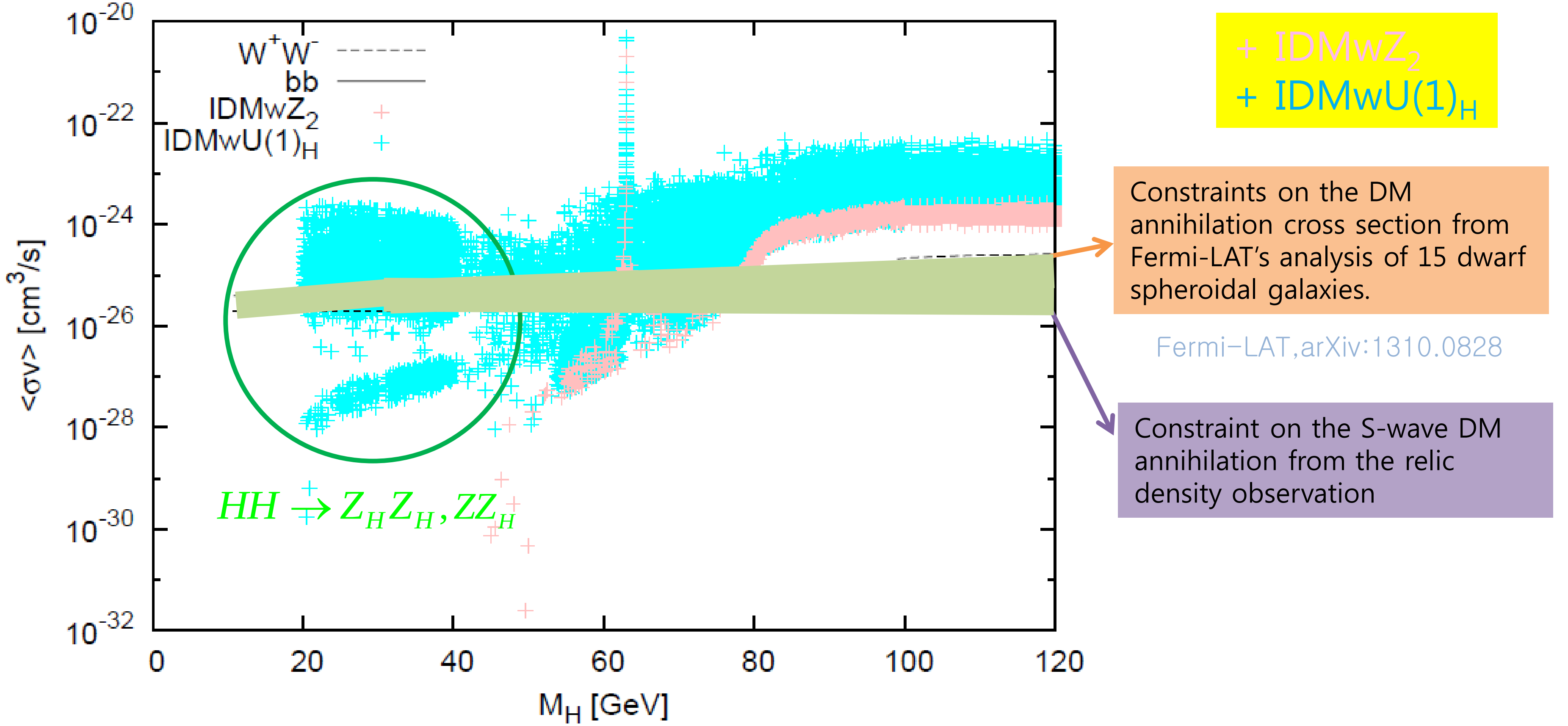
Constraints on the DM annihilation cross section from Fermi-LAT's analysis of 15 dwarf spheroidal galaxies.

Fermi−LAT,arXiv:1310.0828

Constraint on the S-wave DM annihilation from the relic density observation

- All points satisfy constraints from the relic density observation and LUX experiments.

Indirect searches (low mass)



- But, indirect DM signals depend on the decay patterns of produced particles from annihilation or decay of DMs.

Gamma ray flux from DM annihilation

• Dwarf spheroidal galaxies are excellent targets to search for annihilating DM signatures because of DM-dominant nature without astrophysical backgrounds like hot gas.

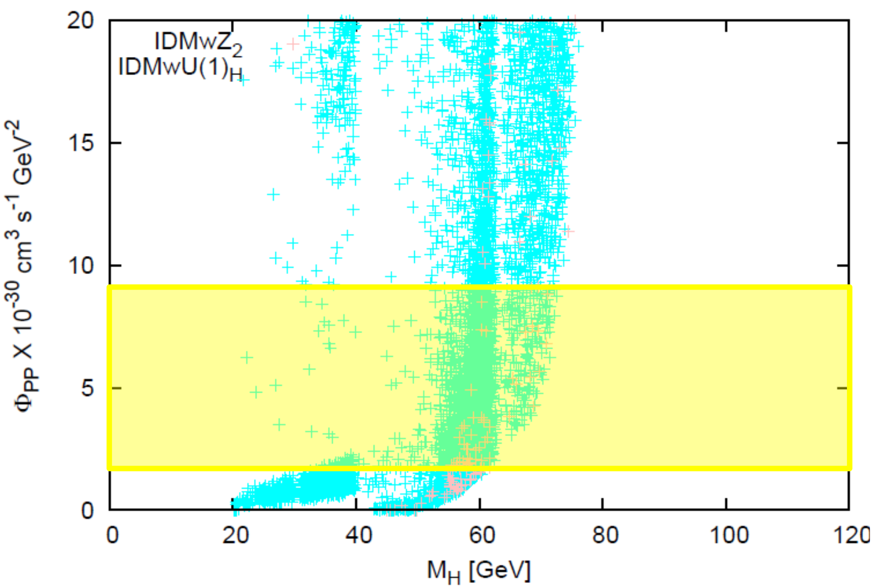
$$\phi_s(\Delta\Omega) = \underbrace{\frac{1}{4\pi}\frac{\langle\sigma v\rangle}{2m_{\rm DM}^2}\int_{E_{\rm min}}^{E_{\rm max}}\frac{{\rm d}N_\gamma}{{\rm d}E_\gamma}{\rm d}E_\gamma}_{\Phi_{\rm PP}} \cdot \underbrace{\int_{\Delta\Omega}\left\{\int_{\rm l.o.s.}\rho^2(\boldsymbol{r}){\rm d}l\right\}{\rm d}\Omega'}_{\text{J-factor}} .$$

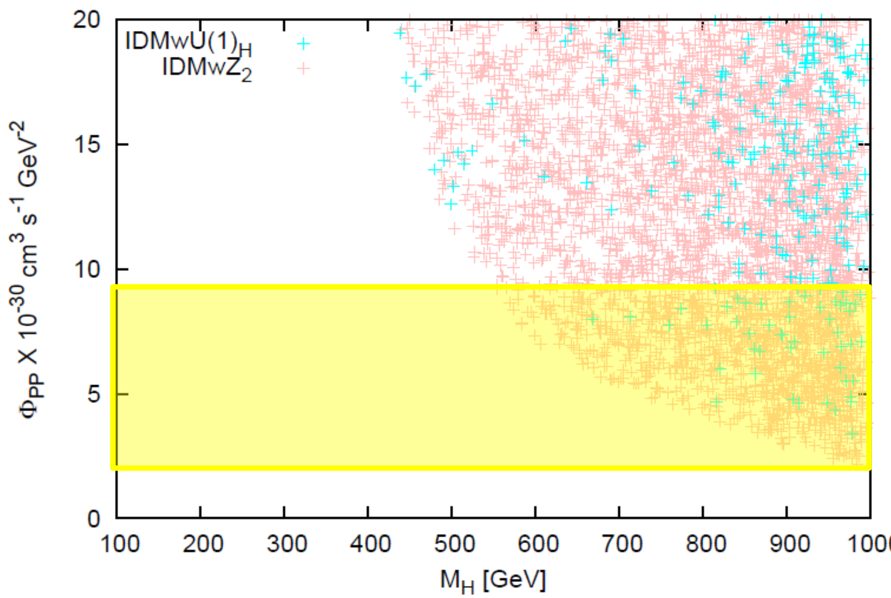
The final γ-ray spectrum.

contains information about the distribution of DM.

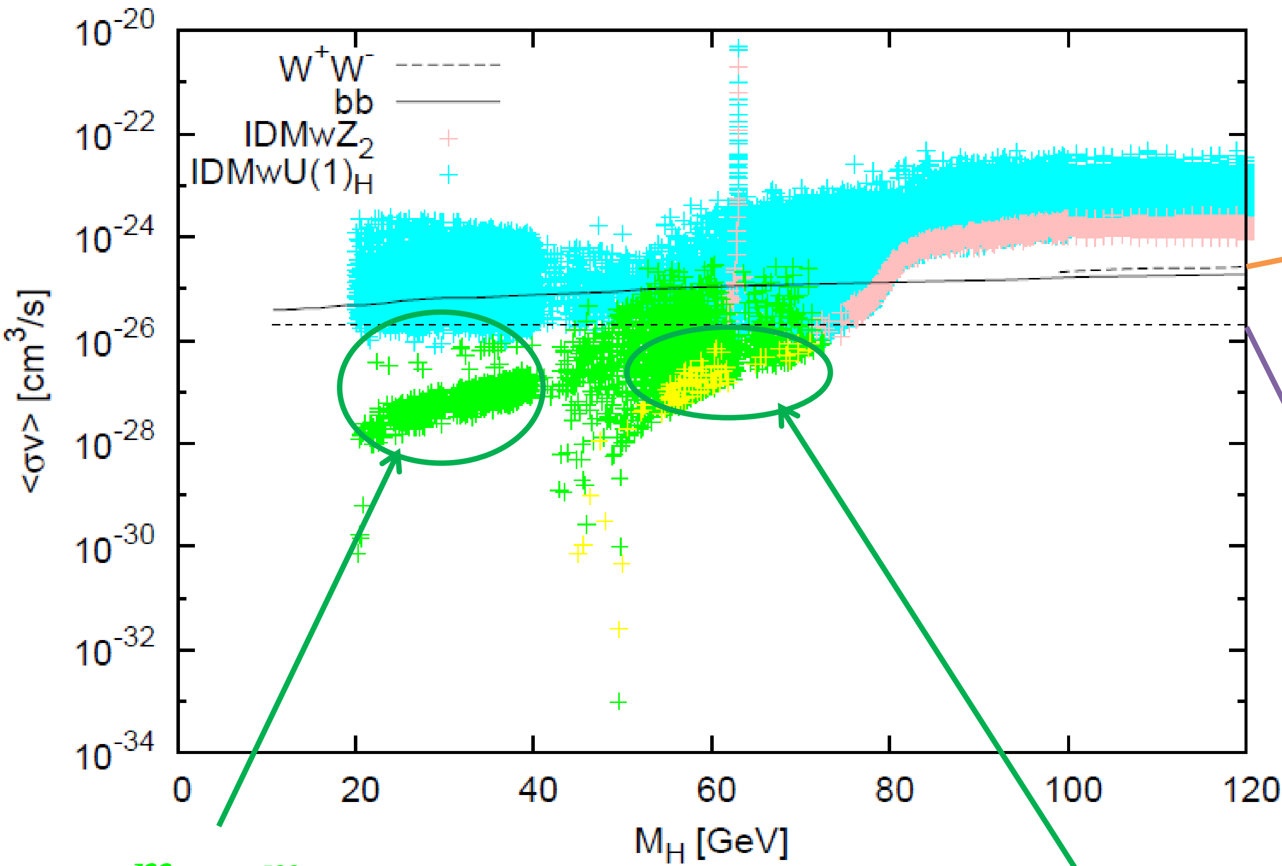
A 95% upper bound is $\Phi_{\rm PP} = 5.0^{+4.3}_{-4.5}\times 10^{-30}\,{\rm cm^3 s^{-1} GeV^{-2}}$

Geringer-Sameth,Koushiappas, PRL107





Indirect searches (low mass)



$m_H \sim m_{Z_H}$

+ IDMwZ$_2$
+ IDMwU(1)$_H$

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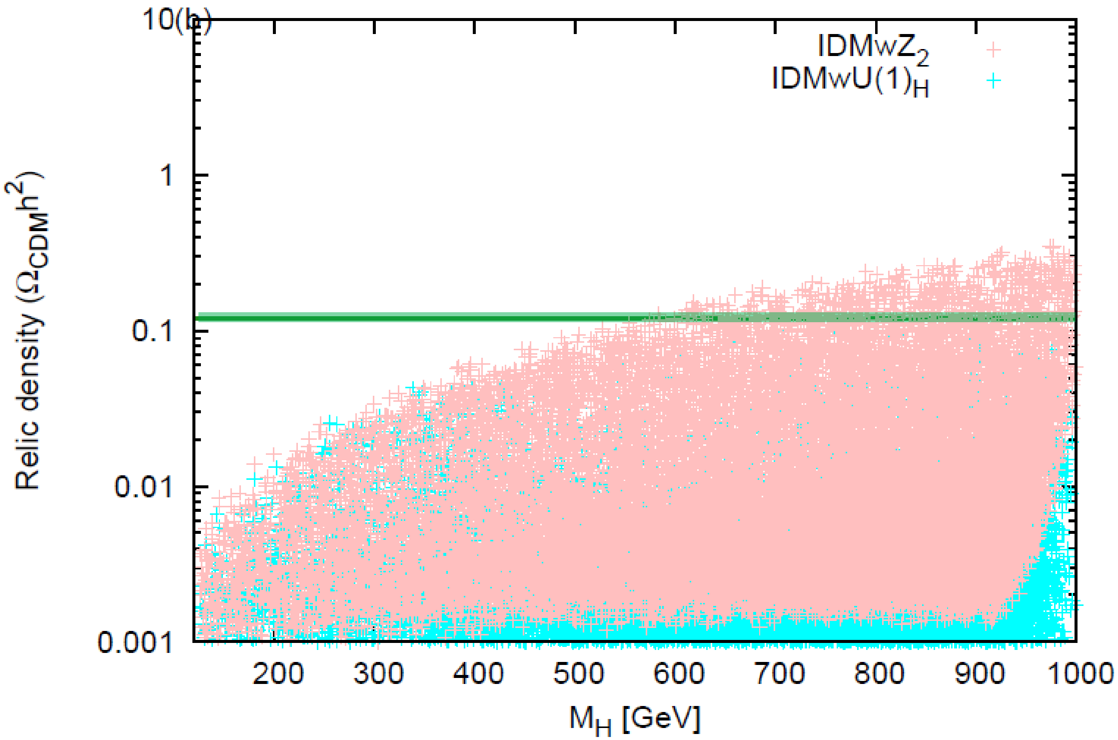
Fermi-LAT,arXiv:1310.0828

Constraint on the S-wave DM annihilation from the relic density observation

Co-annihilation

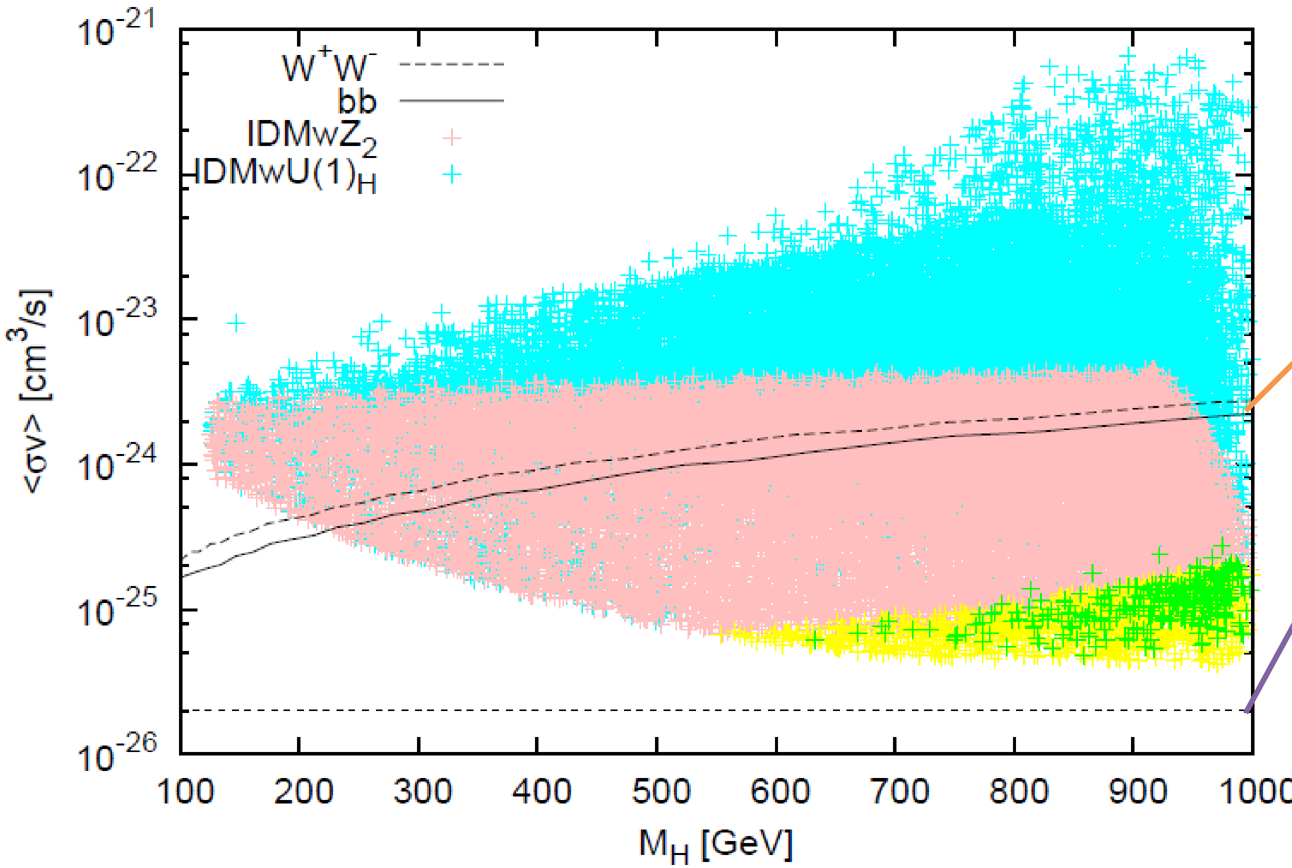
Relic density (high mass)

$$\Omega_{\mathrm{CDM}}h^2 = 0.1199 \pm 0.0027$$



+ $IDMwZ_2$
+ $IDMwU(1)_H$

Indirect searches (high mass)



+ IDMwZ_2
+ IDMwU$(1)_H$

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Constraint on the S-wave DM annihilation from the relic density observation

Conclusions

• 2HDM may be an effective theory of a high-energy theory and useful to test the underlying theory.

• 2HDM can easily be extended to a gauged model and the U(1) gauge symmetry could be the origin of Z_2 symmetry.

• The U(1) extension to inert doublet model could introduce dark matter candidates whose stability are guaranteed by the remnant discrete gauge symmetry of $U(1)_H$.

• In type-I, a light CDM scenario is possible in the IDMw$U(1)_H$.

• Type-II with local $U(1)_H$ is under study (stay tuned)

General aspects for Local Dark Gauge Symmetry

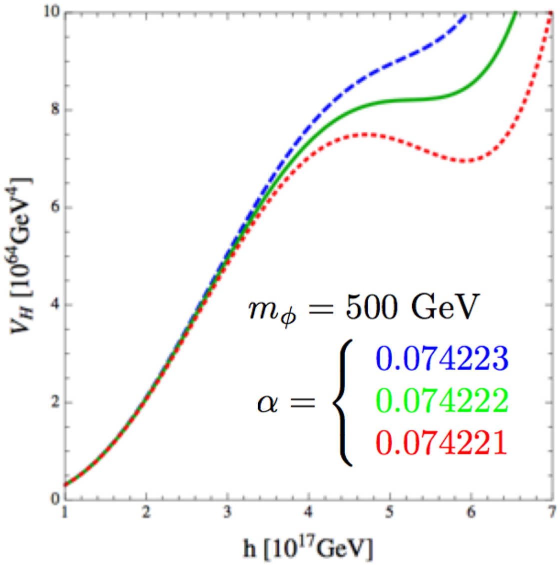
- Stability of EW scale CDM guaranteed dynamically like QED + electron

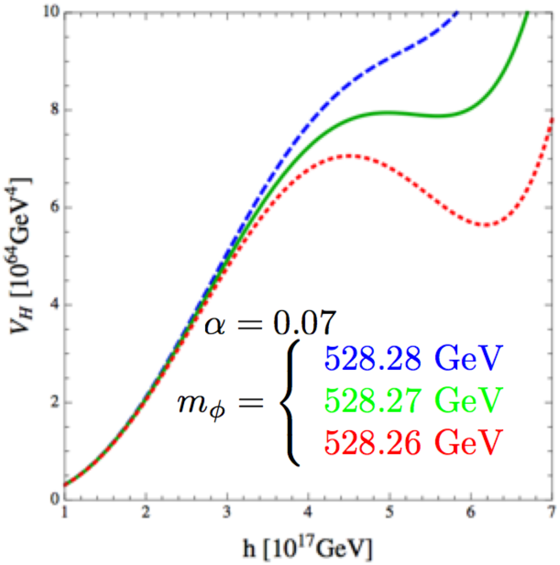
- Massive dark Higgs and dark gauge boson appear in addition to the DM and can modify the DM phenomenology completely
- Allowed DM mass region can be lower than DM models without fark gauge symmetry
- Sizable self interaction among CDM possible by light dark Higgs or dark gauge bosons
- Higgs can have nonstandard decay modes into dark Higgs and/or dark gauge boson, and the signal strengths universally reduced from 1

- Dark Higgs stabilizes the EW vacuum up to Planck scale, and also opens a new dim for the Higgs inflation (r~0.1 possible)

- These points are discussed in detail in recent papers with my collaboratators

Higgs portal Higgs inflation



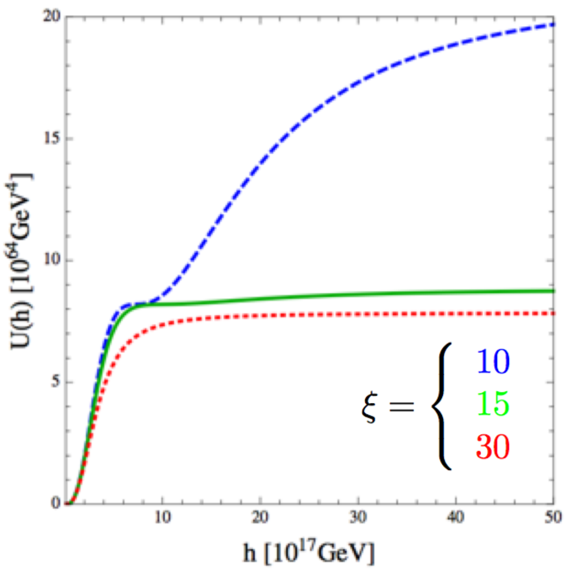


$m_t = 173.2$ GeV
$M_h = 125.5$ GeV

Ko, Park arXiv: 1405.1635

* Inflection point control

(α, m_ϕ) & $\lambda_{\Phi H}$



Result of numerical analysis

$k_* \times$ Mpc	N_e	$h_*/M_{\rm Pl}$	ϵ_*	η_*	$10^9 P_S$	n_s	r
0.002	59	0.83	0.00448	−0.02465	2.2639	0.9238	0.0717
0.05	56	0.72	0.00525	−0.0019	2.1777	0.9647	0.084

- Result depends very sensitively on α, m_Φ and $\lambda_{\Phi H}$ -

H.P.H.I allows Higgs inflation matching to BICEP2 result without resorting to m_t and M_h.

(with S.Baek, Suyong Choi, P. Gondolo, T. Hur, D.W.Jung, Sunghoon Jung, J.Y.Lee, W.I.Park, E.Senaha, Yong Tang in various combinations)

- Strongly interacting hidden sector (0709.1218 PLB, 1103.2571 PRL)
- Light DM in leptophobic Z' model (1106.0885 PRD)
- Singlet fermion dark matter (1112.1847 JHEP)
- Higgs portal vector dark matter (1212.2131 JHEP)
- Vacuum structure and stability issues (1209.4163 JHEP)
- Singlet portal extensions of the standard seesaw models with unbroken dark symmetry (1303.4280 JHEP)
- Hidden sector Monopole, VDM and DR (1311.1035)
- Self-interacting scalar DM with local Z3 symmetry (1402.6449)
- And a few more, including Higgs-portal assisted Higgs inflation, Higgs portal VDM for gamma ray excess from GC, and DM-sterile nu's etc.

Back up

Higgs Potential

- in the ordinary 2HDM with Z_2 symmetry

$$V = m_1^2 H_1^\dagger H_1 + m_2^2 H_2^\dagger H_2 - (m_{12}^2 H_1^\dagger H_2 + h.c.) + \frac{1}{2}\lambda_1 (H_1^\dagger H_1)^2 + \frac{1}{2}\lambda_2 (H_2^\dagger H_2)^2$$
$$+\lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 (H_1^\dagger H_2)(H_2^\dagger H_1) + \frac{1}{2}\lambda_5 [(H_1^\dagger H_2)^2 + h.c.].$$

not invariant under $U(1)_H$

- in the 2HDM with $U(1)_H$, we include an extra singlet scalar Φ, which makes Z_H heavy.

$$V = \hat{m}_1^2(|\Phi|^2) H_1^\dagger H_1 + \hat{m}_2^2(|\Phi|^2) H_2^\dagger H_2 - \left(m_3^2(\Phi) H_1^\dagger H_2 + h.c.\right)$$
$$+\frac{\lambda_1}{2}(H_1^\dagger H_1)^2 + \frac{\lambda_2}{2}(H_2^\dagger H_2)^2 + \lambda_3 (H_1^\dagger H_1)(H_2^\dagger H_2) + \lambda_4 |H_1^\dagger H_2|^2$$
$$+m_\Phi^2 |\Phi|^2 + \lambda_\Phi |\Phi|^4.$$

$H_1^\dagger H_2 \Phi$

invariant under $U(1)_H$

no $\lambda 5$ terms!

- neutral Higgs

$$\begin{pmatrix} h_\Phi \\ h_1 \\ h_2 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\alpha & -\sin\alpha \\ 0 & \sin\alpha & \cos\alpha \end{pmatrix} \begin{pmatrix} \cos\alpha_1 & 0 & -\sin\alpha_1 \\ 0 & 1 & 0 \\ \sin\alpha_1 & 0 & \cos\alpha_1 \end{pmatrix} \begin{pmatrix} \cos\alpha_2 & -\sin\alpha_2 & 0 \\ \sin\alpha_2 & \cos\alpha_2 & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \tilde{h} \\ H \\ h \end{pmatrix}$$

- a pair of charged Higgs + 1 pseudoscalar Higgs + 3 neutral Higgs bosons