



Differential Transverse Momentum Correlations in p-Pb and Pb-Pb in ALICE

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ALICE

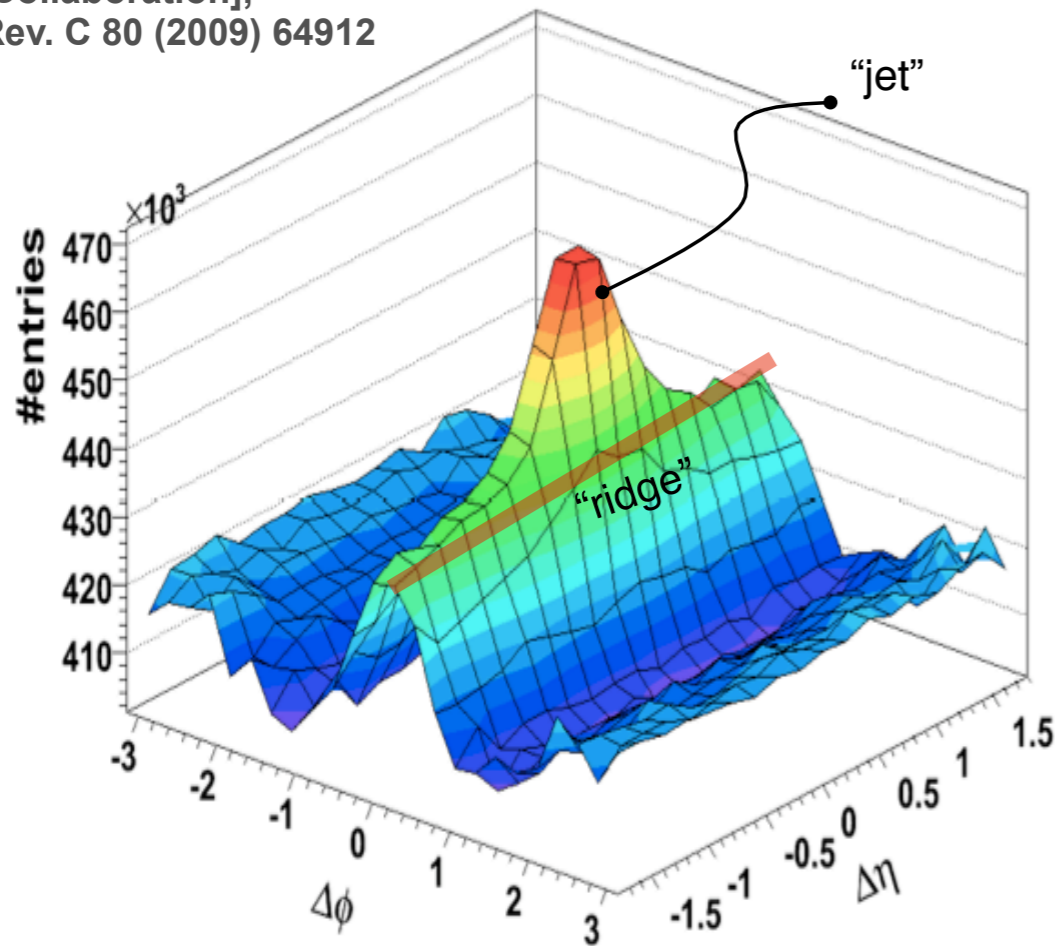
A JOURNEY OF DISCOVERY

Outline

- Introduction to 2-particle correlation features
Near and away-side ($\Delta\eta, \Delta\varphi$) correlation
- Observables
Number (“old”) and transverse momentum correlation (“new”)
- Physics motivation
Gain additional information on underlying particle production mechanisms
- Results (p-Pb vs. Pb-Pb systems)
Correlation functions and Fourier decompositions
- Summary

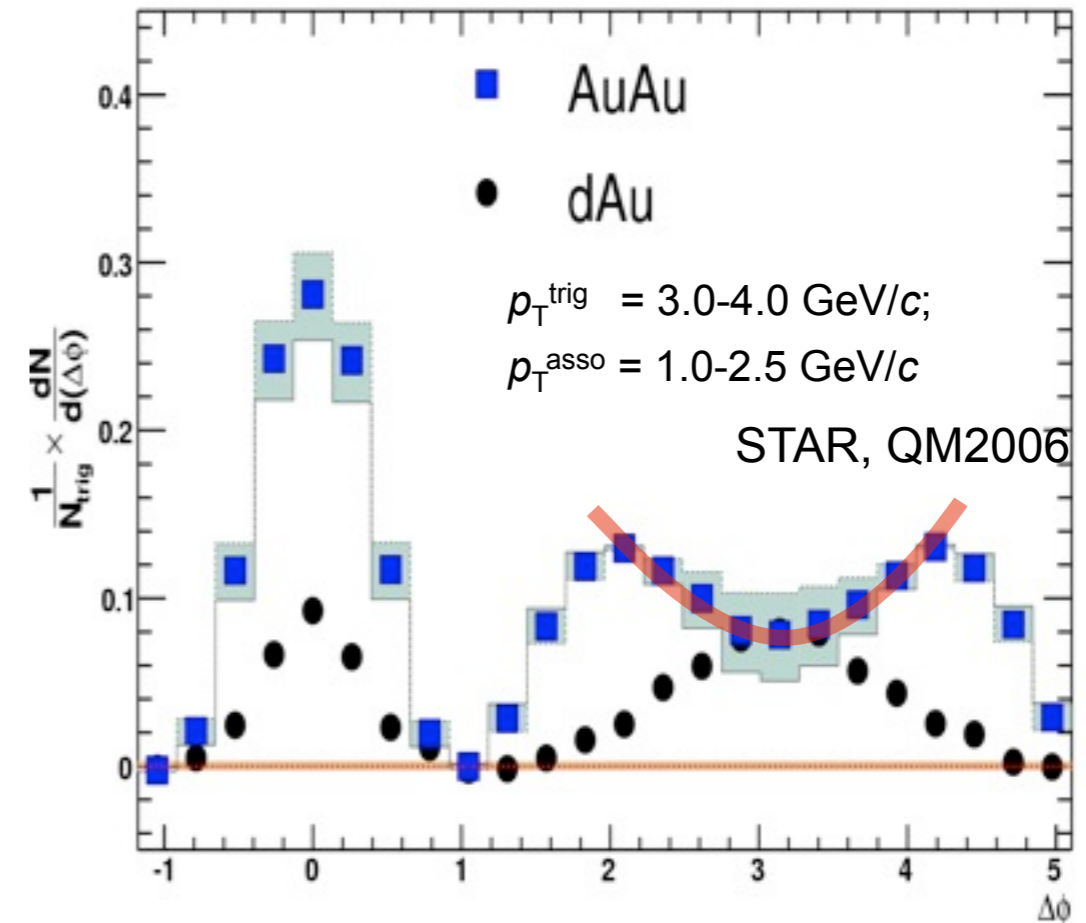
“The Ridge”

[STAR Collaboration],
Phys. Rev. C 80 (2009) 64912



“The Valley”

Phys. Rev. Lett. 102, 052302



“With elliptic flow (v_2) only subtracted”

Au-Au collisions exhibit a long range rapidity correlation called the ridge - not seen in pp (at the same collision energy)

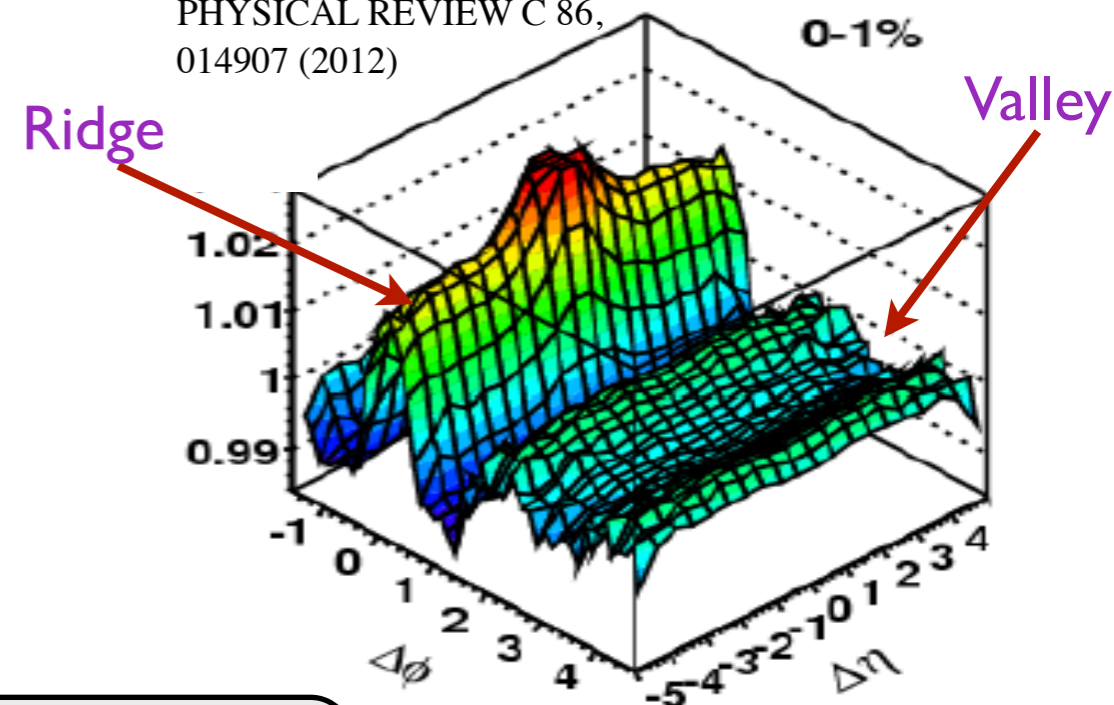
“The valley”, double bump peak, observed in central Au-Au
 (2-particle correlation with background subtracted)

Physics Processes:

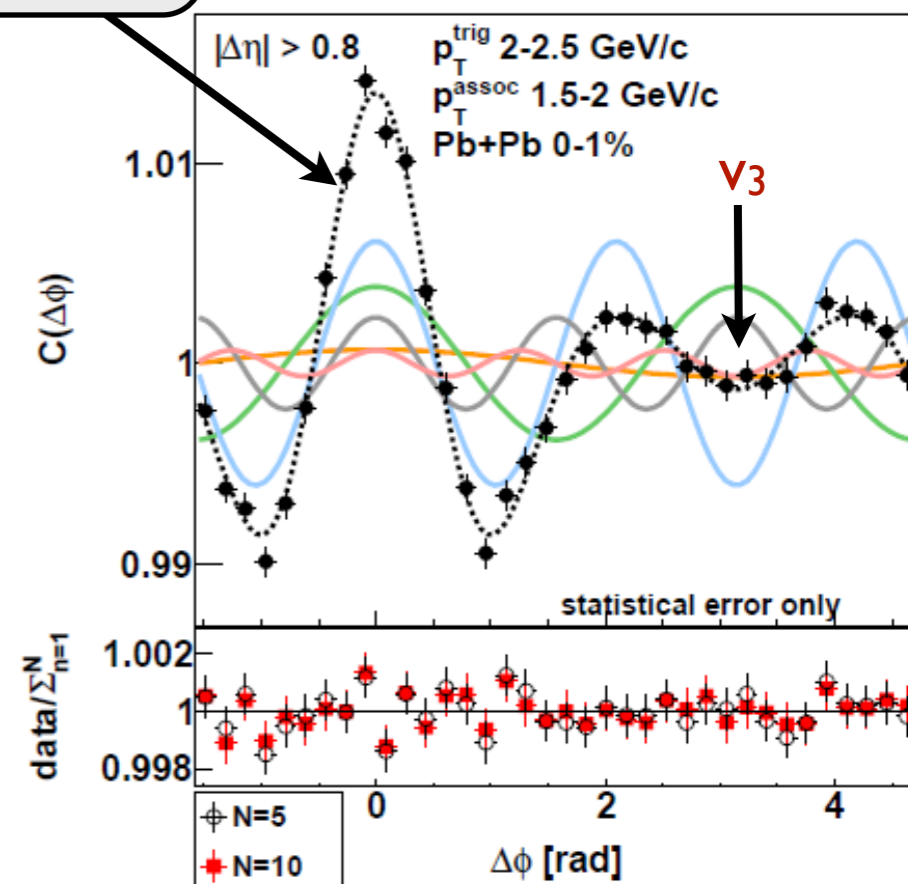
- Lumpy initial conditions also affect in long range rapidity correlations
 - Initial geometry fluctuation leads to even and odd harmonics in particle azimuthal distributions
- Hydrodynamic evolution
spatial anisotropy -> momentum anisotropy

Two Particle Correlation (ATLAS)

PHYSICAL REVIEW C 86,
014907 (2012)



$v_1 + v_2 + v_3 + v_4 + v_5$



And many open questions ...

- How large are the higher harmonics?
- Are all harmonics driven by initial geometry?
- Are the amplitudes of the harmonics correlated?
- At what p_T do jets matter?
- Relative contributions of flow and “non-flow”?
- Is the mach cone (really) dead?
- etc, etc

Scope for new correlation observables ...

2-particle correlation functions - definitions

Distribution of “correlated pairs”

$$C_2(x_1, x_2) = \rho_2(x_1, x_2) - \rho_1(x_1)\rho_2(x_2) \quad x \equiv \{y, p_T\} \quad \rho(x) = \frac{1}{\sigma} \frac{d\sigma}{dx}$$

Number correlation function:

$$R_2(x_1, x_2) = \frac{C_2(x_1, x_2)}{\rho_1(x_1)\rho_1(x_2)}$$

R_2 is a robust observable;
single track efficiencies cancel
out of this ratio

Transverse momentum correlation function:

$$\langle \Delta p_{T,1} \Delta p_{T,2} \rangle = \frac{\int \rho_2 \Delta p_{T,1} \Delta p_{T,2} dp_{T,1} dp_{T,2}}{\int \rho_2 dp_{T,1} dp_{T,2}}$$

measures deviation from the
inclusive average momentum

$$\Delta p_T = p_{T,i} - \langle p_T \rangle$$
$$\langle p_{T,i} \rangle = \frac{\int \rho_1(\mathbf{p}_i) p_{T,i} dp_{T,i}}{\int \rho_1(\mathbf{p}_i) dp_{T,i}}$$

Why $\Delta p_T \Delta p_T (\Delta \eta, \Delta \varphi)$ is interesting ?

- Additional sensitivity to hardness of particles and their correlations (jet vs. non-jet)
- **Positive correlation:** both particles have p_T either higher or lower than the average p_T
- **Negative correlation:** one particle has a higher p_T than average whereas the associated particle has p_T lower than average
- **Integral part of this observable:** sensitive to temperature fluctuation, average momentum fluctuation
- **Flow factorization (see next slide...)**

Physics motivation II: factorization

- Factorization expected when correlations driven by geometry only

$$V_{n\Delta}(p_1, p_2) = v_n(p_1)v_n(p_2)$$

- Factorization expected to break down when jet and geometry have similar contributions

Geometrical flow model

Sharma and Pruneau, Phys Rev C,
79, 024905, 2009

Factorization for $\Delta p_T \Delta p_T$

$$\Delta p_T \Delta p_T(\Delta\varphi) = \frac{2 \sum_n (v_n^{p_1} - \langle p_T \rangle v_n^1) (v_n^{p_2} - \langle p_T \rangle v_n^2)}{1 + 2 \sum_n v_n^1 v_n^2 \cos(n\Delta\varphi)}$$

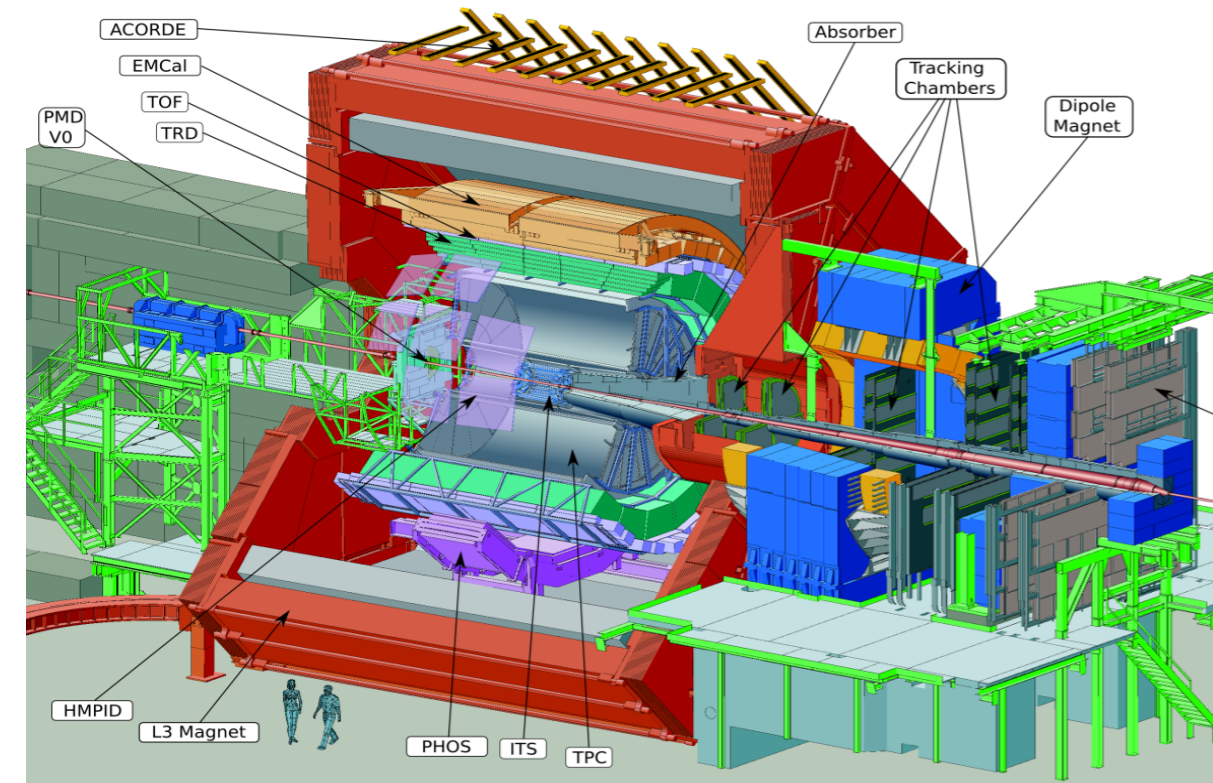
used later as:

Eq. 2

Consequence:

FDC of $\Delta p_T \Delta p_T$ at large $\Delta\eta$ should scale with flow coefficients

Experimental setup and methodology



- 4 different charge combinations: (+ -), (- +), (+ +), and (- -)
- Results for different charge combinations combined into **charge independent (CI)**

$$CI = \frac{1}{2} \{LS + US\}$$

TPC: gas drift detector \longrightarrow *charged particles*
 ITS: Silicon detector \longrightarrow *charged particles*

The centrality is selected using the magnitude of the VZERO (V0) signal (\sim multiplicity)

- kinematic selection

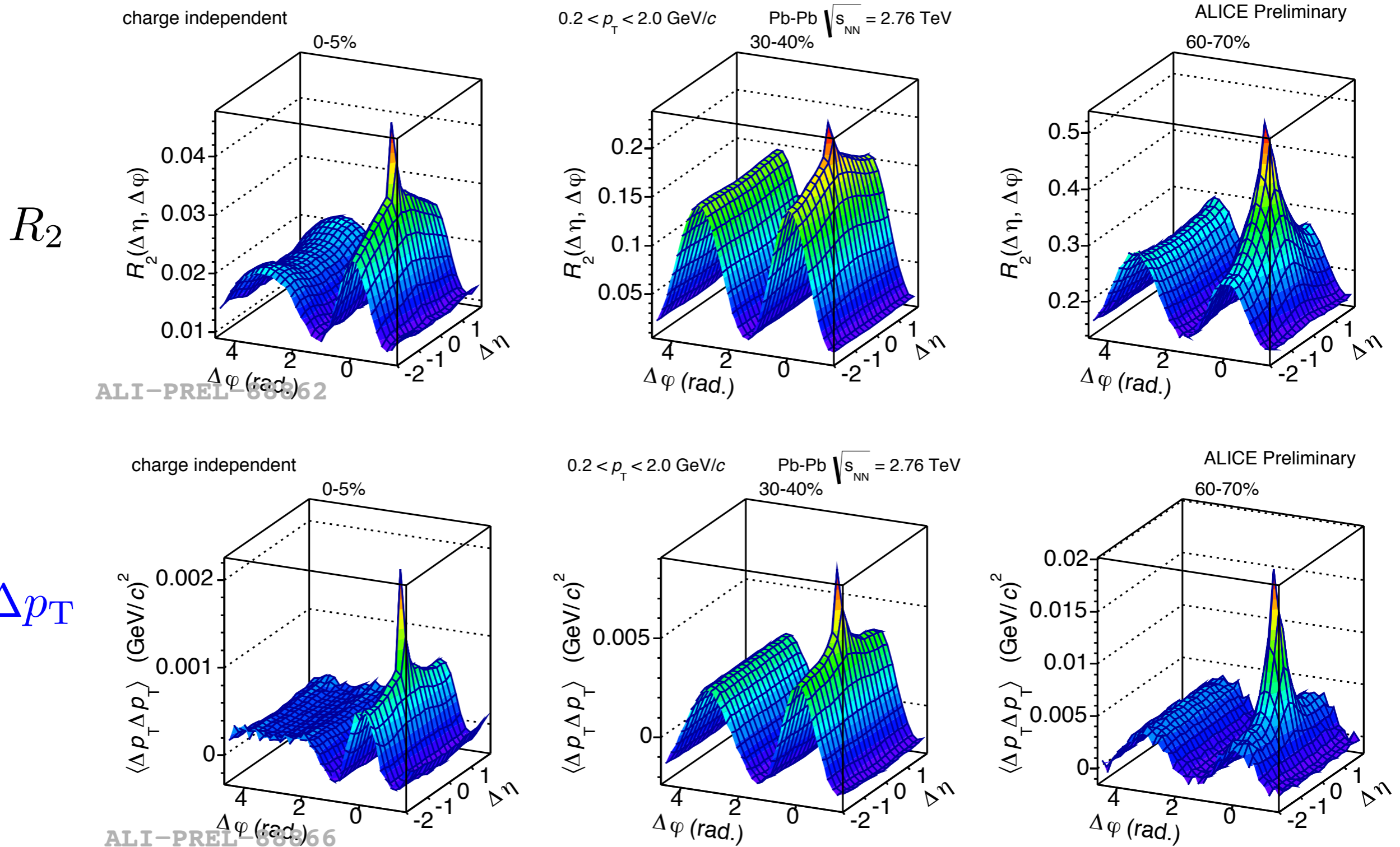
$$|\eta| < 1.0, \quad 0 < \varphi < 2\pi$$

$$0.2 < p_T < 2.0 \text{ GeV}/c$$

$$LS = \frac{1}{2} \{((++) + (---))\}$$

$$US = \frac{1}{2} \{((+-) + (-+))\}$$

Correlation functions in Pb-Pb

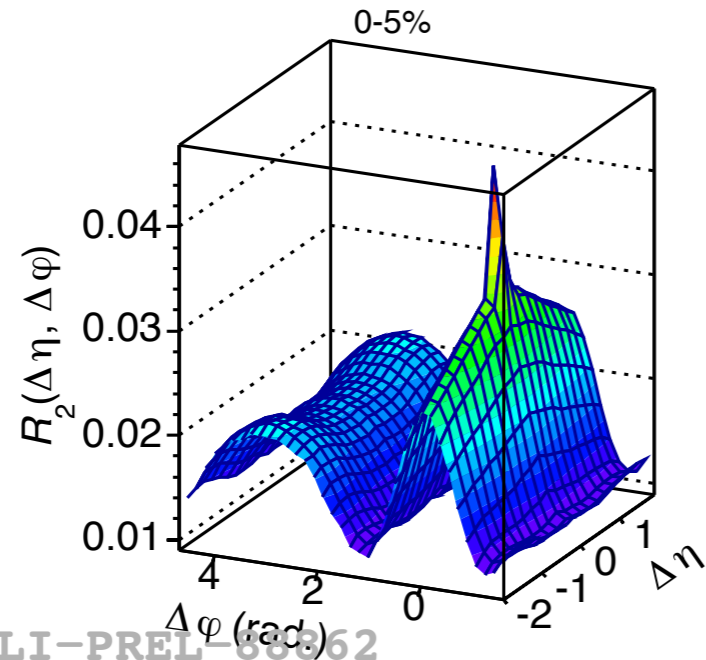


- Away-side 'double bump' in most central collision for $\Delta p_T \Delta p_T$

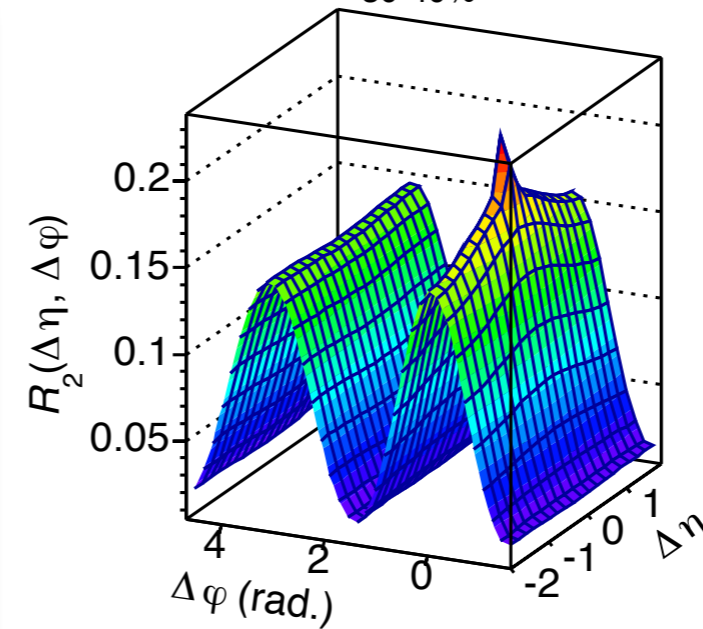
Correlation functions in Pb-Pb

R_2

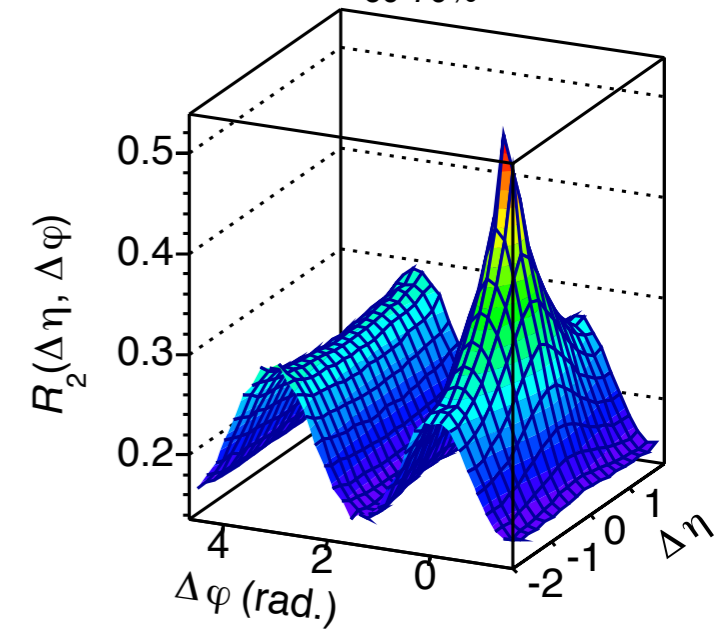
charge independent



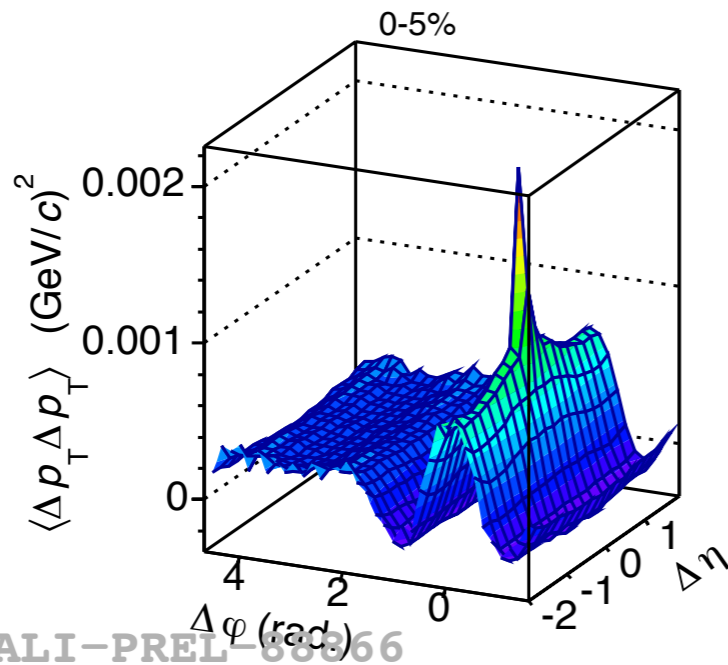
$0.2 < p_T < 2.0 \text{ GeV}/c$ Pb-Pb $\sqrt{s_{NN}} = 2.76 \text{ TeV}$



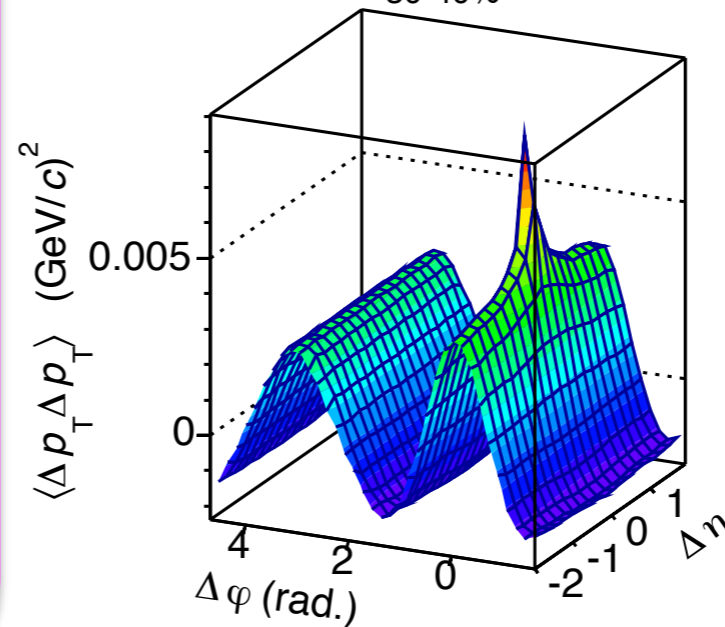
ALICE Preliminary
60-70%



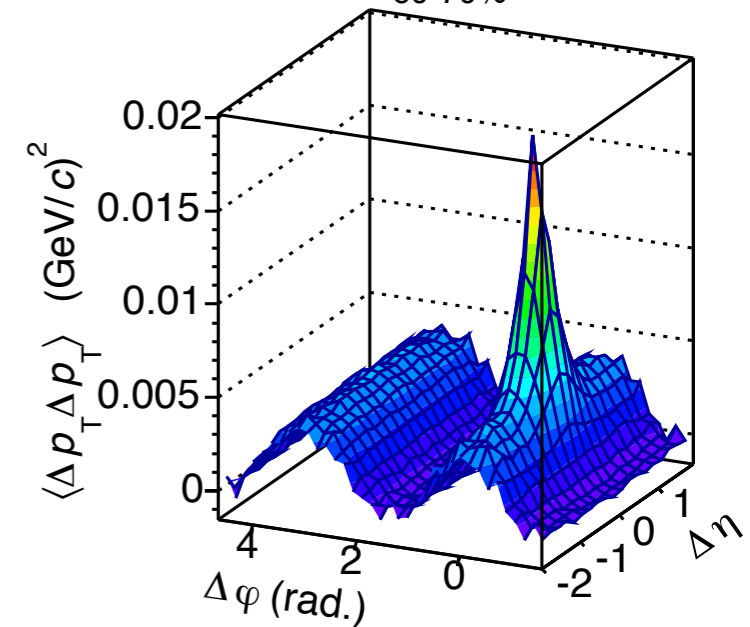
charge independent



$0.2 < p_T < 2.0 \text{ GeV}/c$ Pb-Pb $\sqrt{s_{NN}} = 2.76 \text{ TeV}$



ALICE Preliminary
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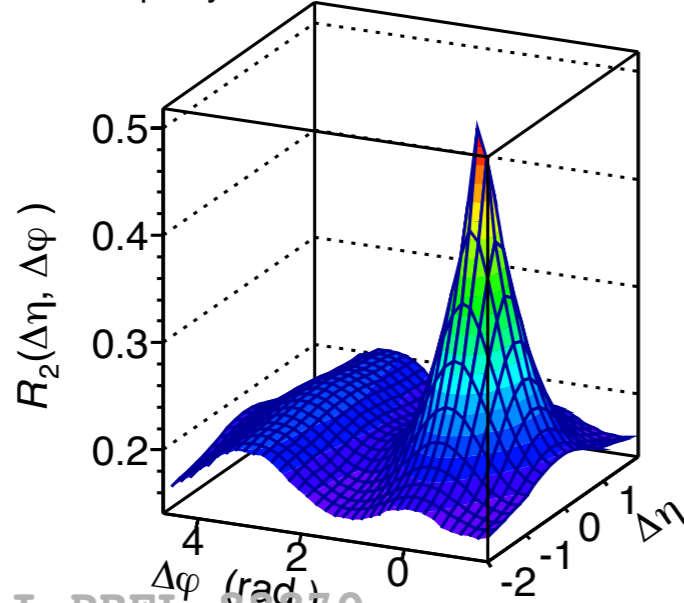
$\Delta p_T \Delta p_T$

- Away-side 'double bump' in most central collision for $\Delta p_T \Delta p_T$

correlation functions in p-Pb

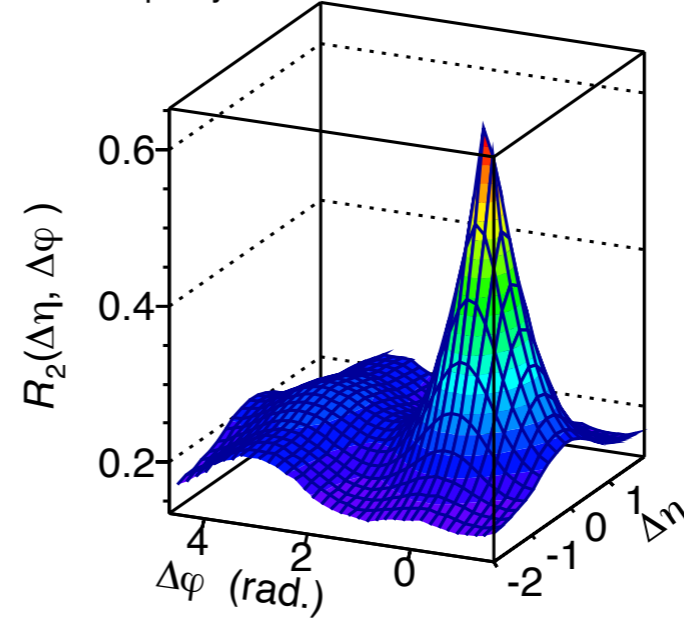
R_2

charge independent
VOA Multiplicity Class 0-20%

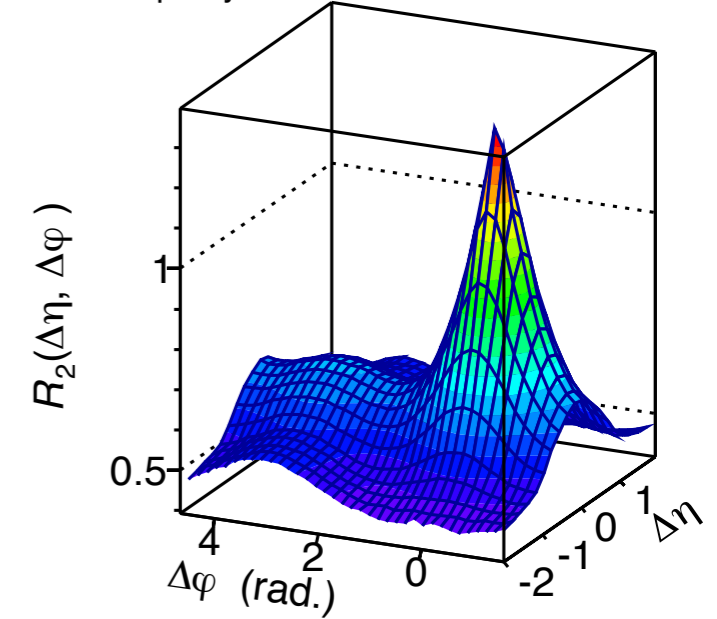


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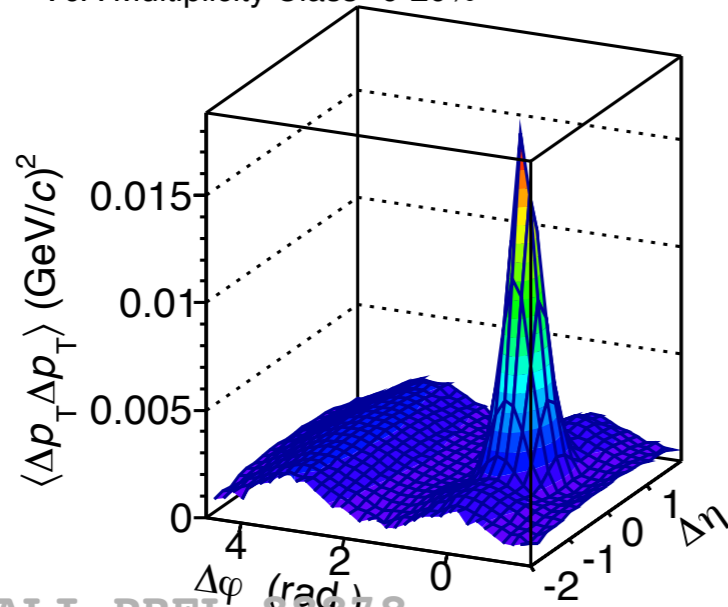
$0.2 < p_T < 2.0$ GeV/c p-Pb $\sqrt{s_{NN}} = 5.02$ TeV
VOA Multiplicity Class 20-40%



ALICE Preliminary
VOA Multiplicity Class 60-80%

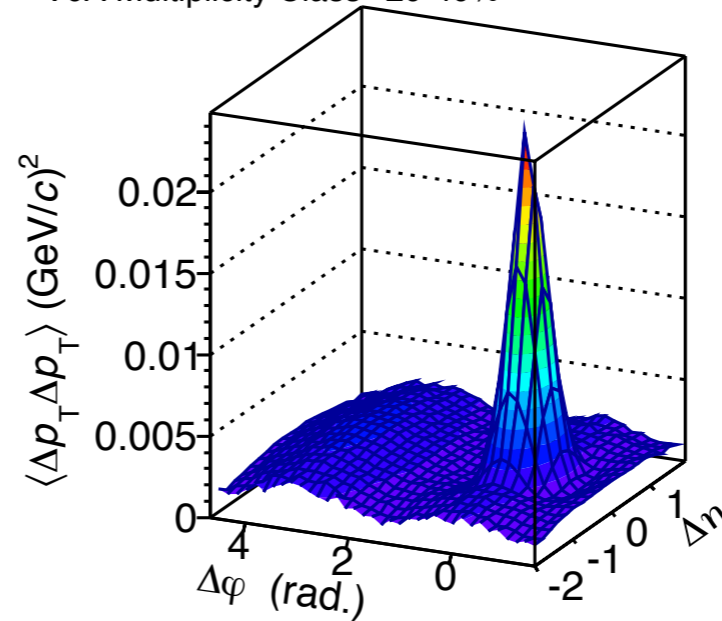


charge independent
VOA Multiplicity Class 0-20%

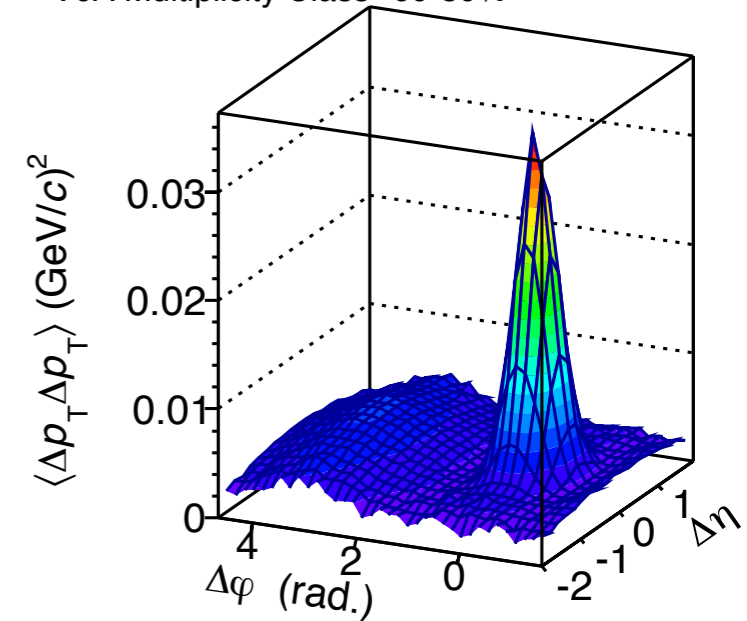


ALI-PREL-88878

$0.2 < p_T < 2.0$ GeV/c p-Pb $\sqrt{s_{NN}} = 5.02$ TeV
VOA Multiplicity Class 20-40%



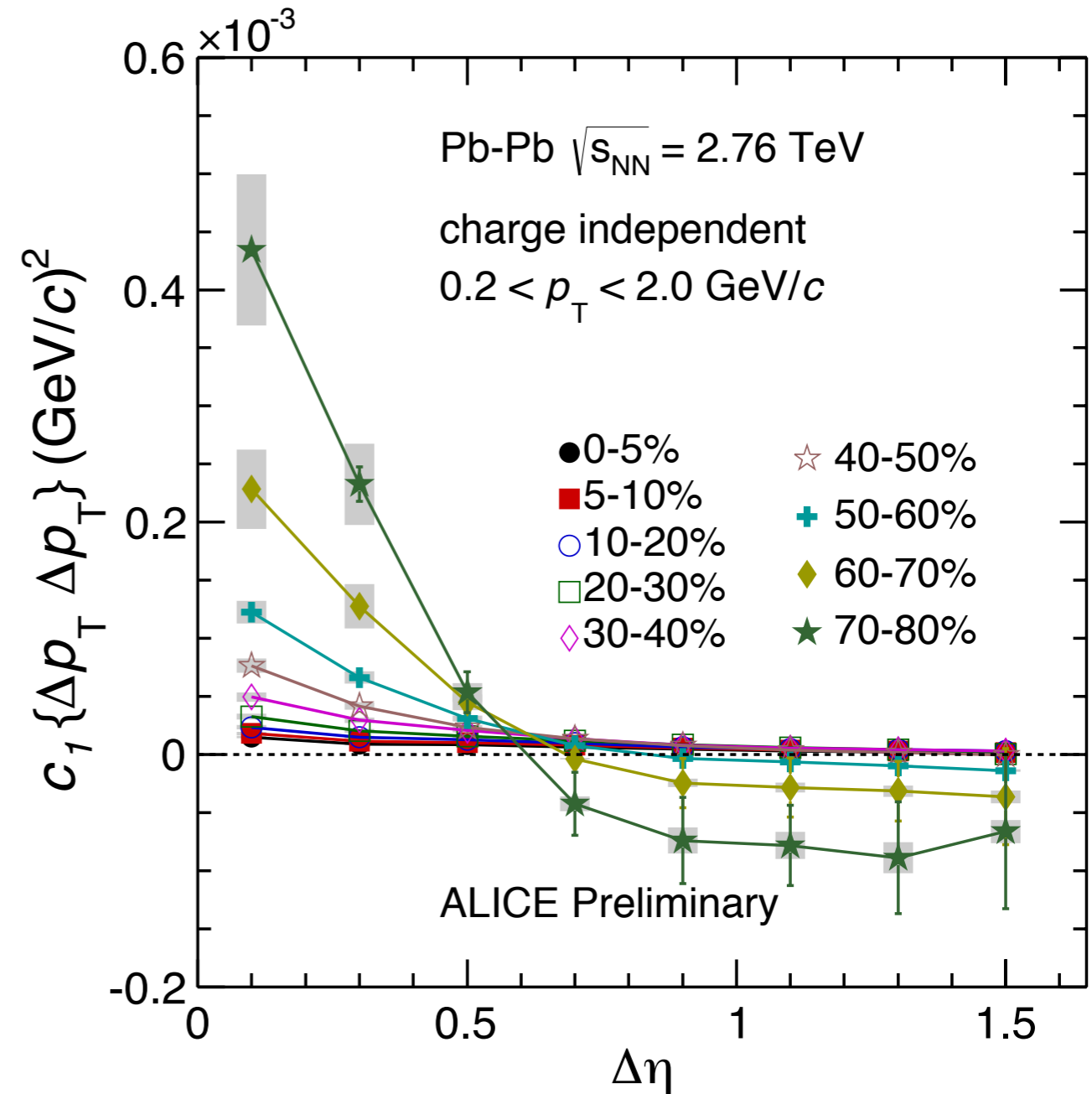
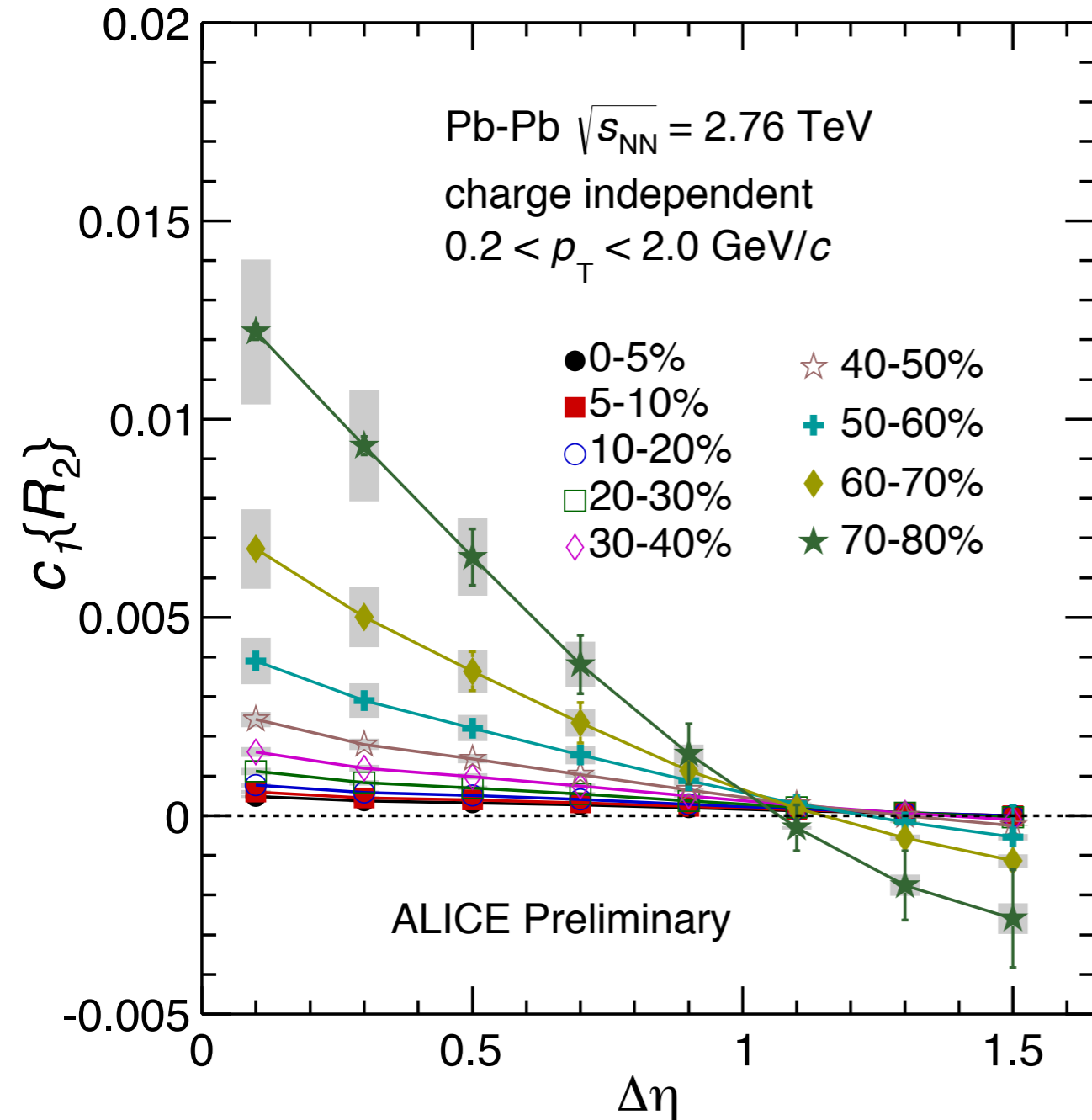
ALICE Preliminary
VOA Multiplicity Class 60-80%



$\Delta p_T \Delta p_T$

- Diminishing $\cos(\Delta\varphi)$ modulations
- Narrower near side peak in $\Delta p_T \Delta p_T$

Harmonic coefficient (c_1) in Pb-Pb

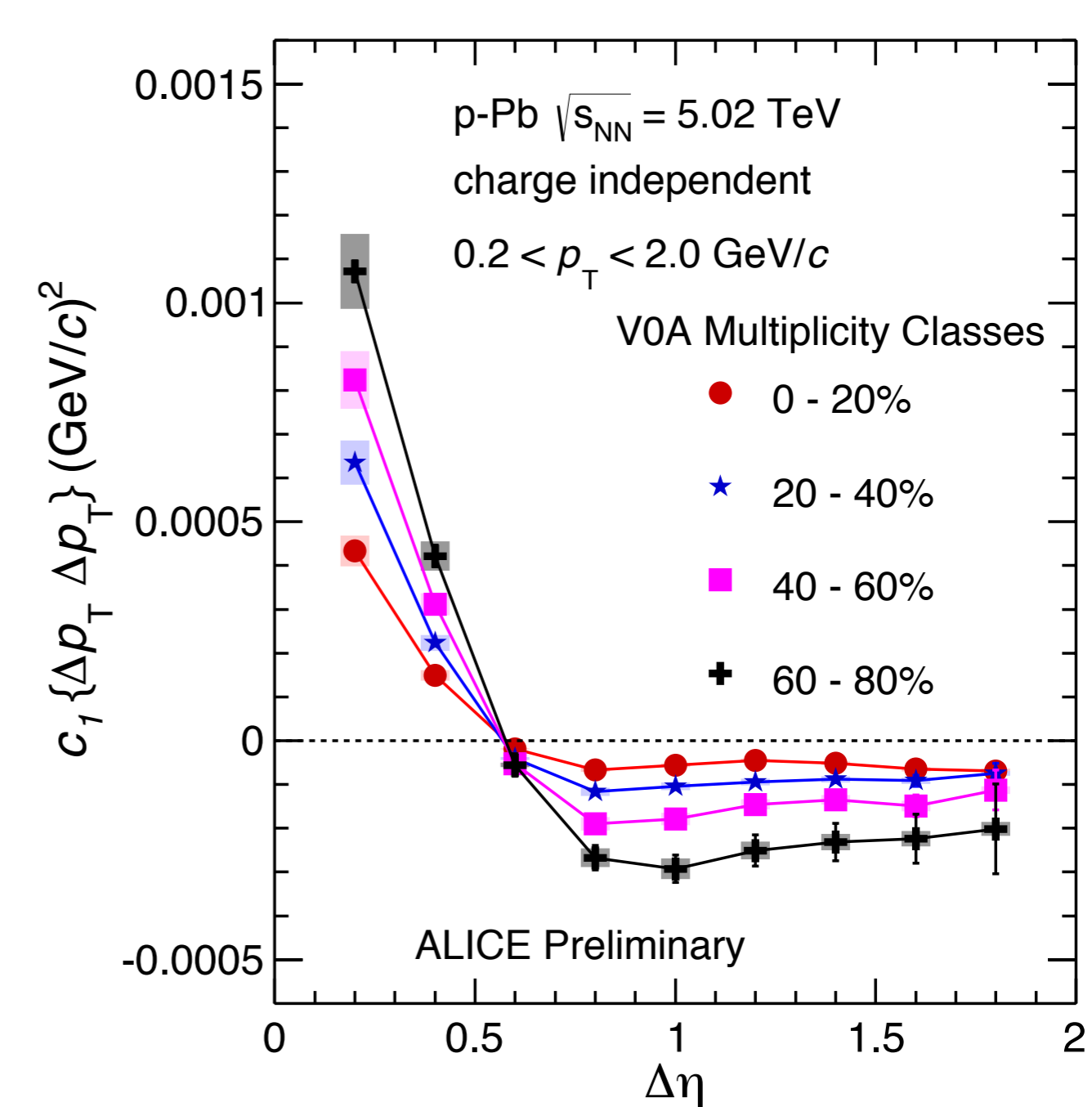
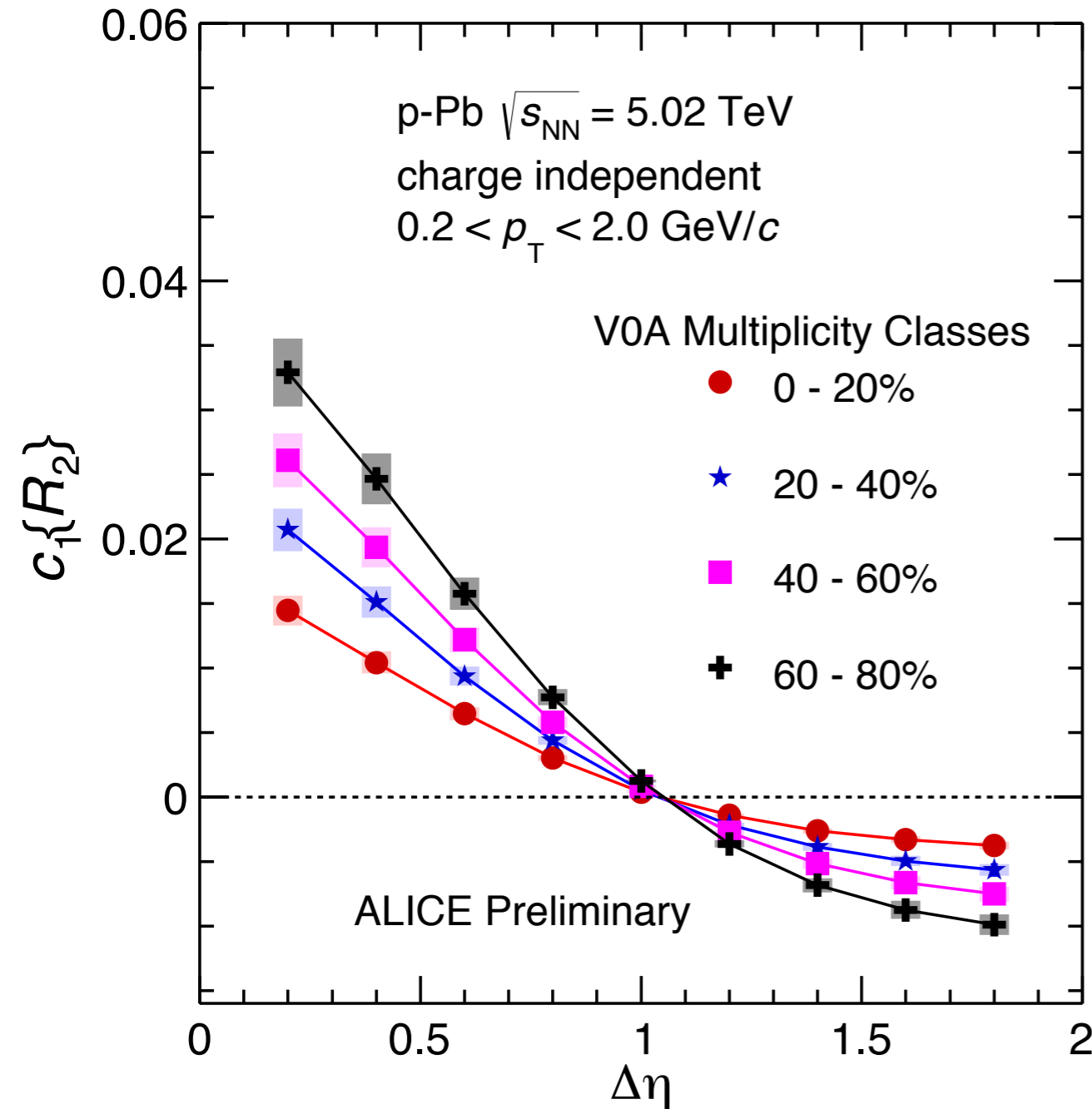


- $\Delta p_T \Delta p_T$ changes sign at lower $\Delta \eta$ than R_2

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Harmonic coefficient (c_1) in p-Pb

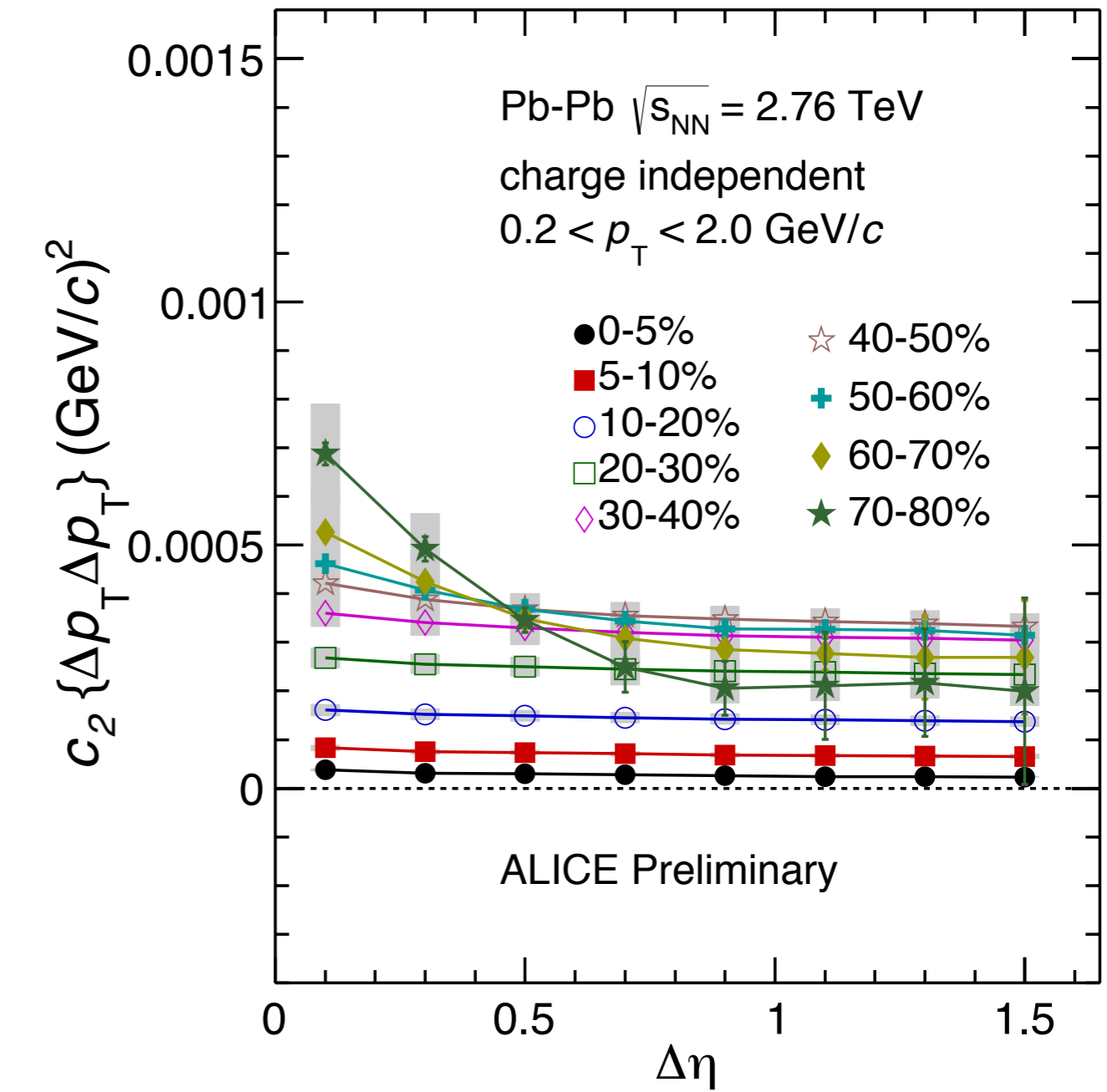
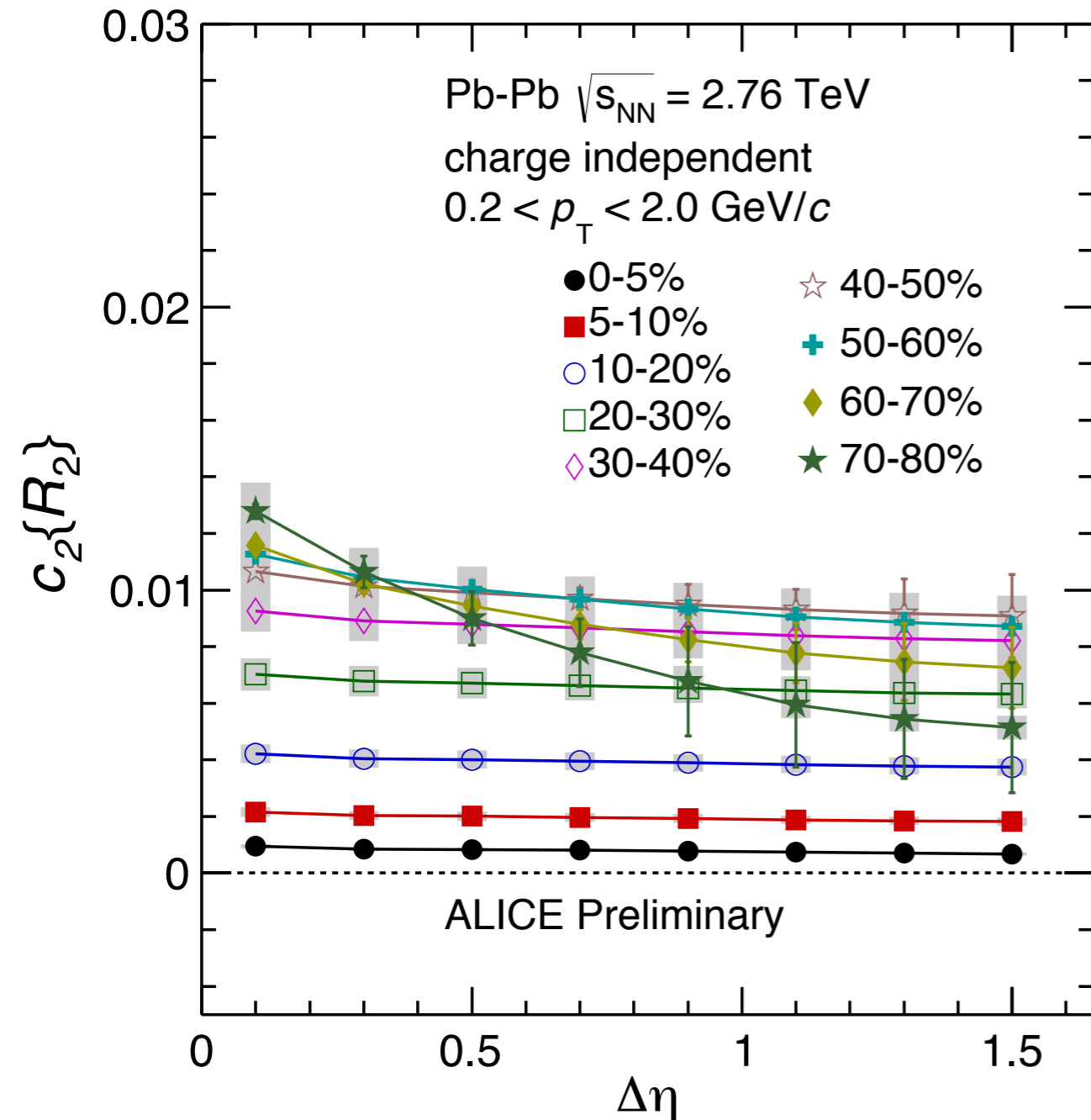


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R_2 : 'longer' range correlation
 $\Delta p_T \Delta p_T$: 'shorter' range correlation

Harmonic coefficient (c_2) in Pb-Pb

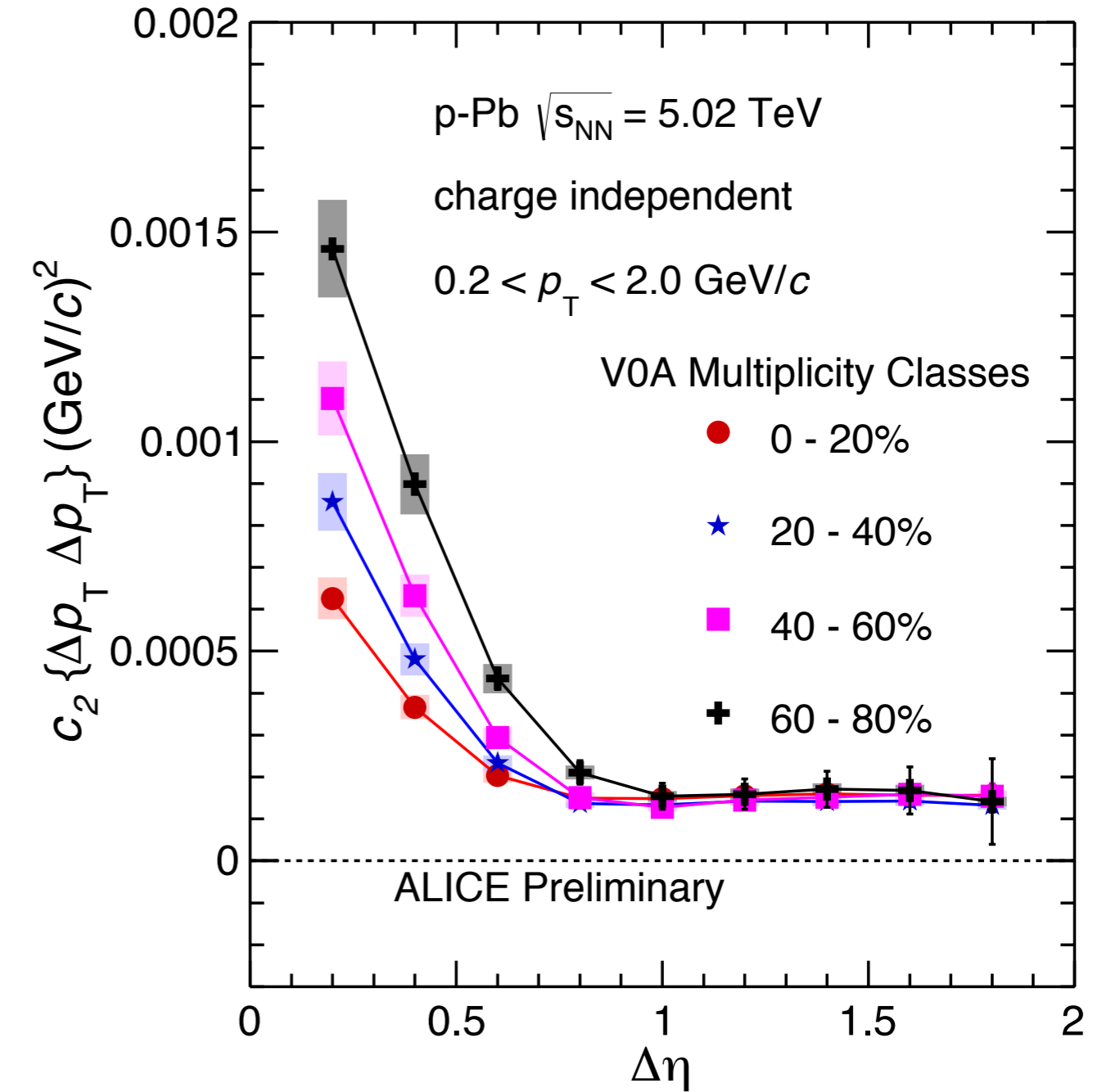
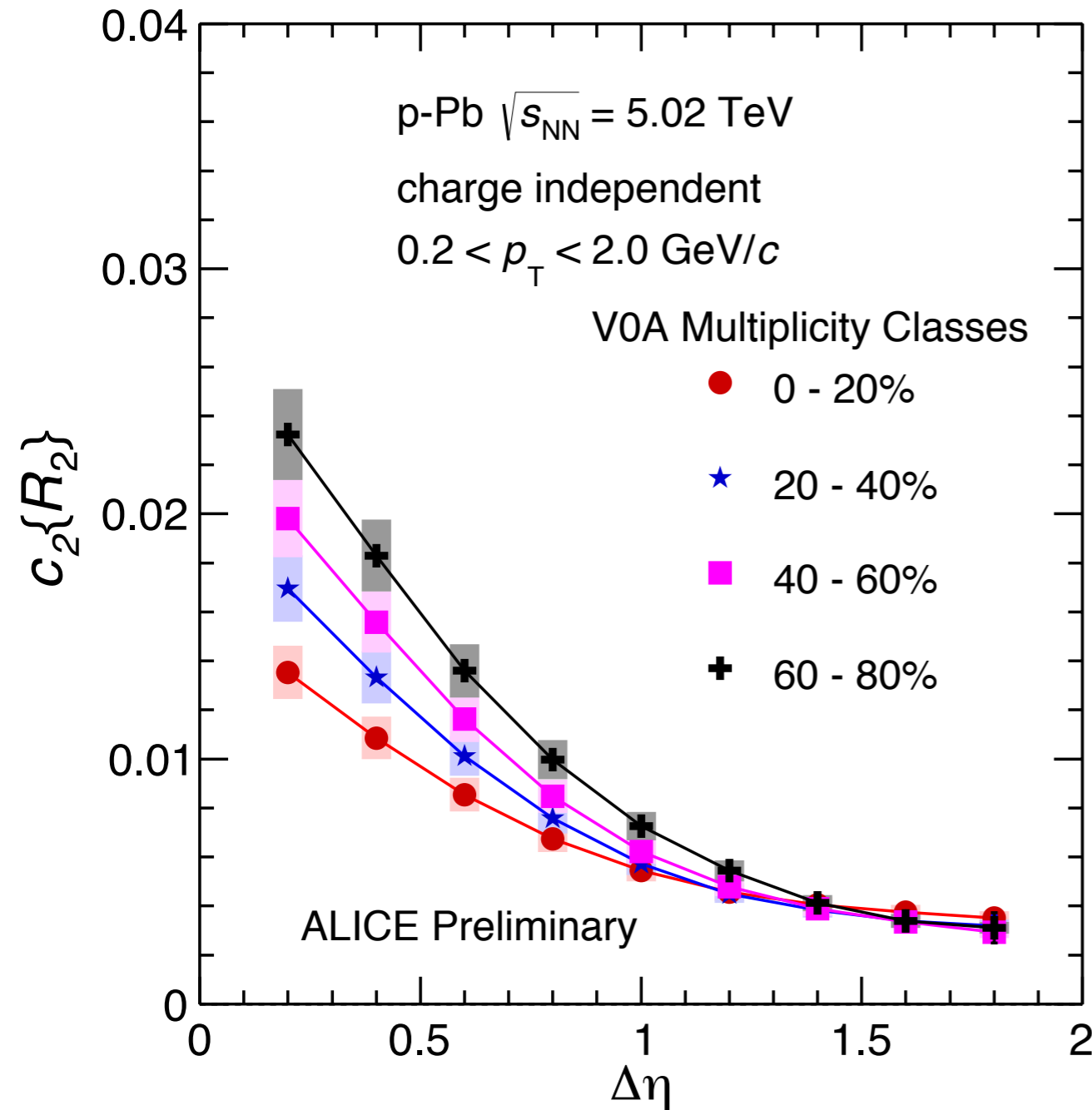


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- Independent of $\Delta\eta$ for 0-60% centrality for R_2 and $\Delta p_T \Delta p_T$
- Centrality dependent of c_2
(increases with decreasing centrality for 0-60%)

Harmonic coefficient (c_2) in p-Pb



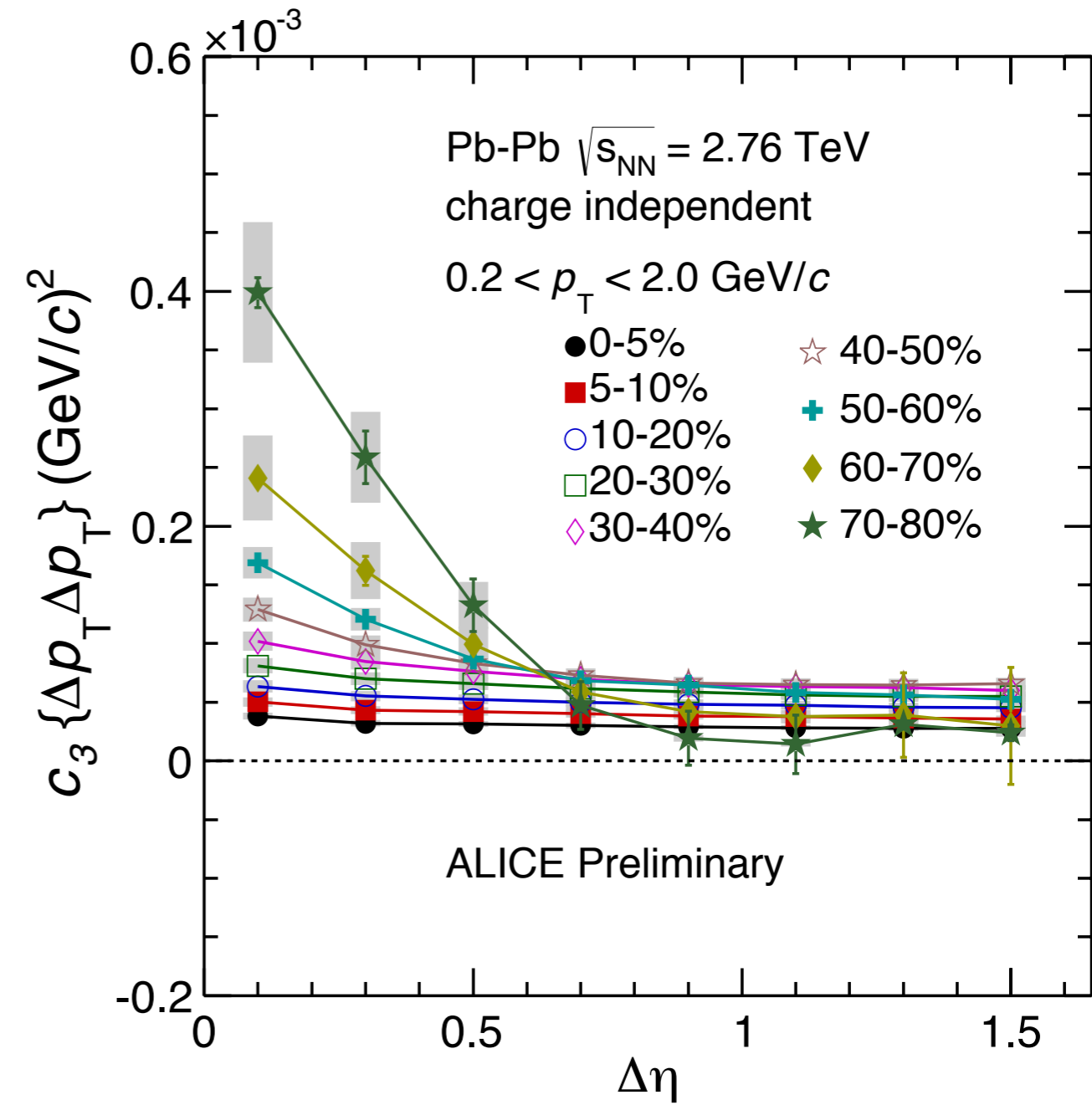
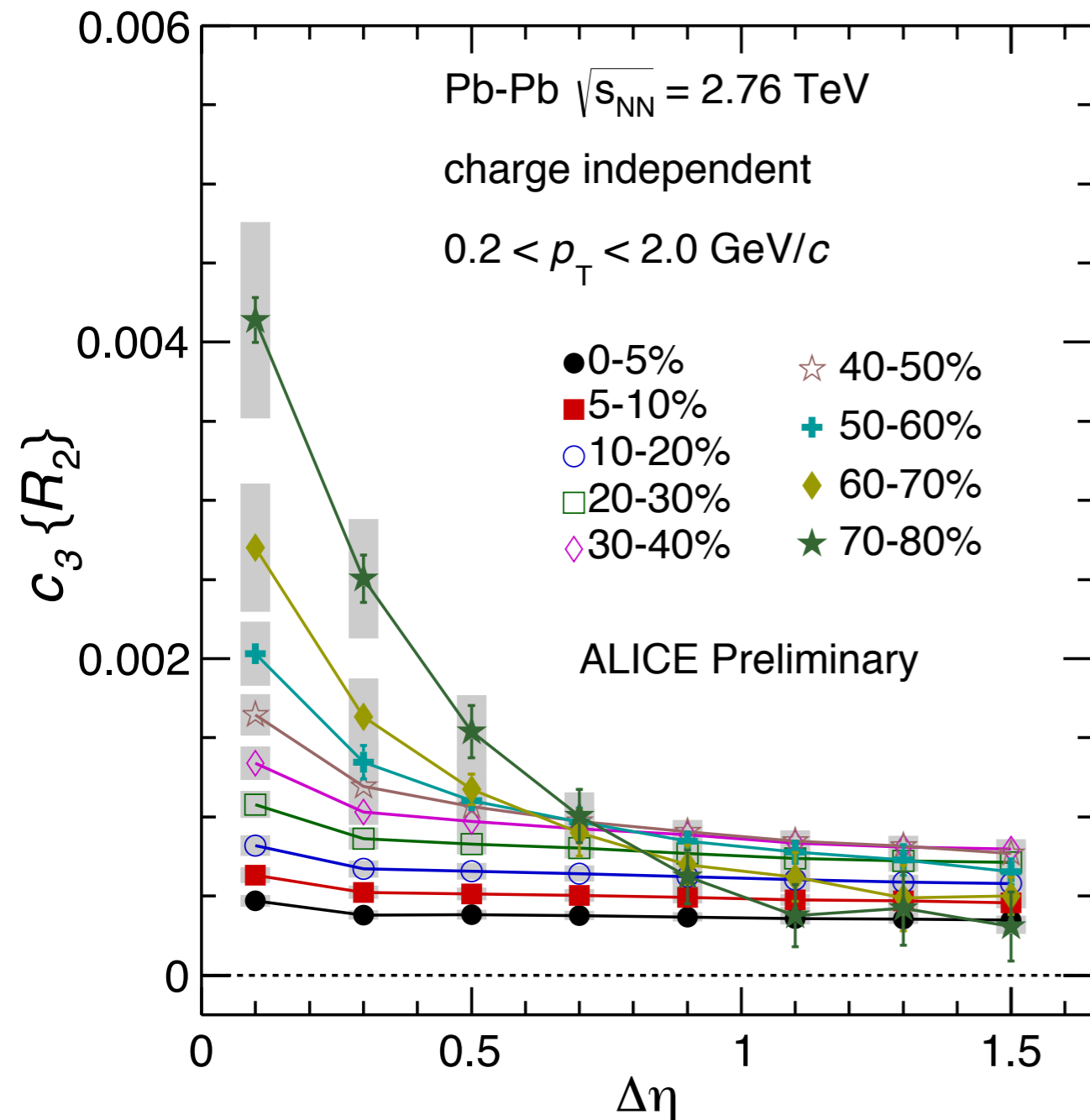
$c_2\{\Delta p_T \Delta p_T\}$ saturates at smaller $\Delta\eta$ than $c_2\{R_2\}$

For $c_2\{\Delta p_T \Delta p_T\}$ all 4 centralities merge at $\Delta\eta \geq 1.0$

ALI-PREL-88902

ALI-PREL-88926

Harmonic coefficient (c_3) in Pb-Pb

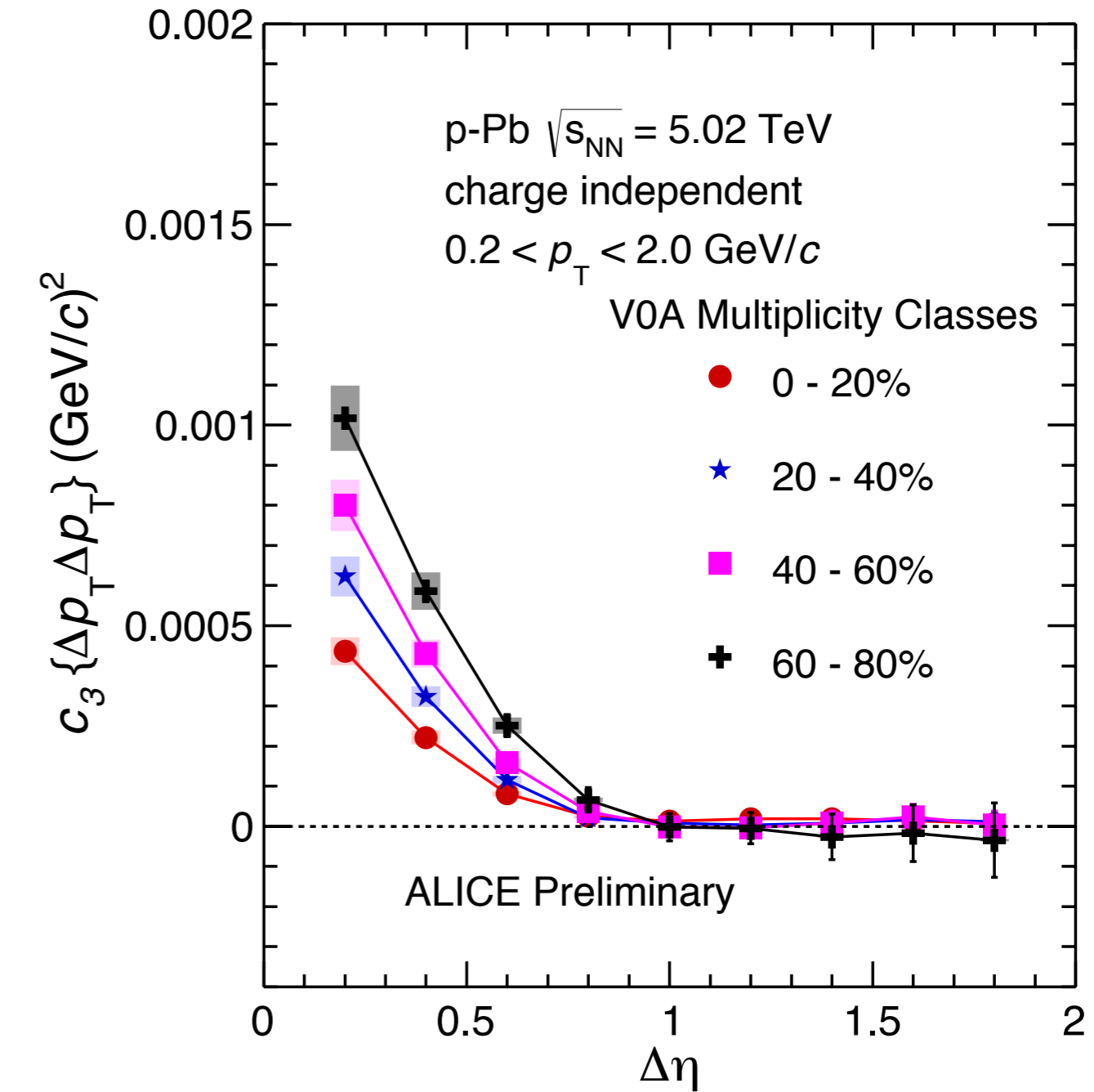
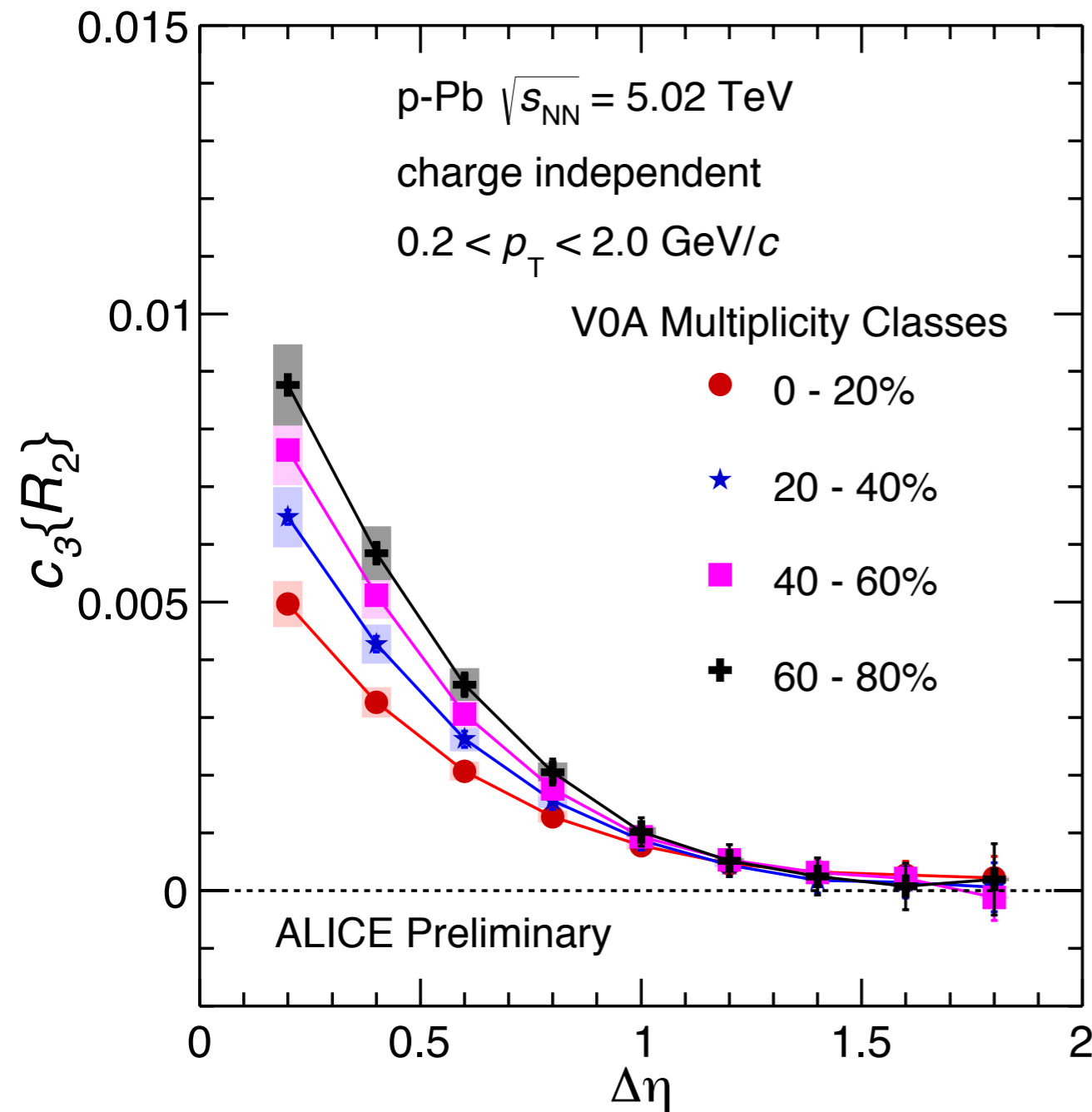


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- Independent of $\Delta\eta$ for 0-60% centrality at $\Delta\eta > 0.5$

Harmonic coefficient (c_3) in p-Pb

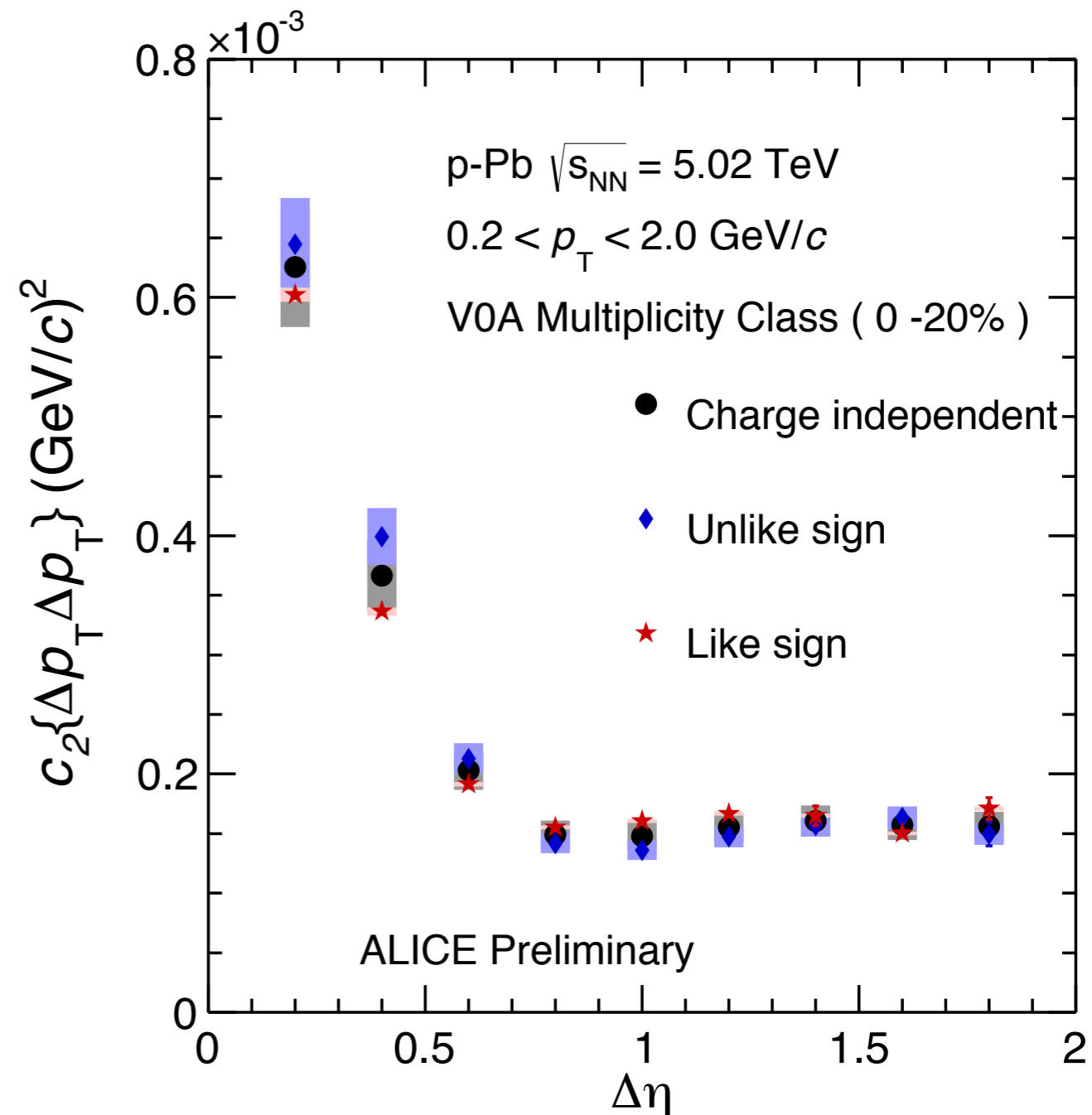
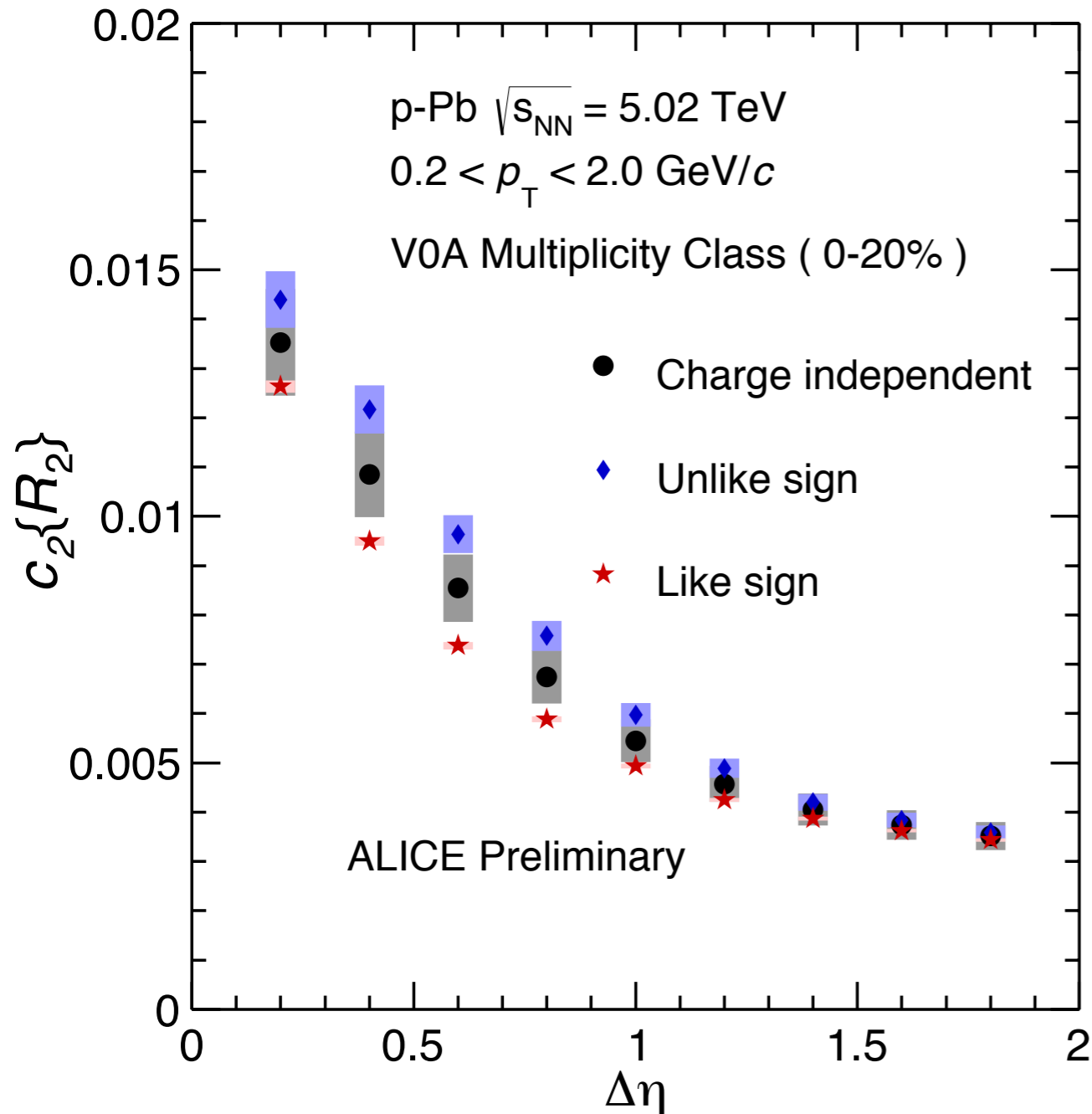


ALI-PREL-88906

ALI-PREL-88930

- $c_3\{\Delta p_T \Delta p_T\}$ saturates and consistent with 'zero' at $\Delta\eta \geq 1.0$

Like sign vs. unlike sign (c_2) in p-Pb



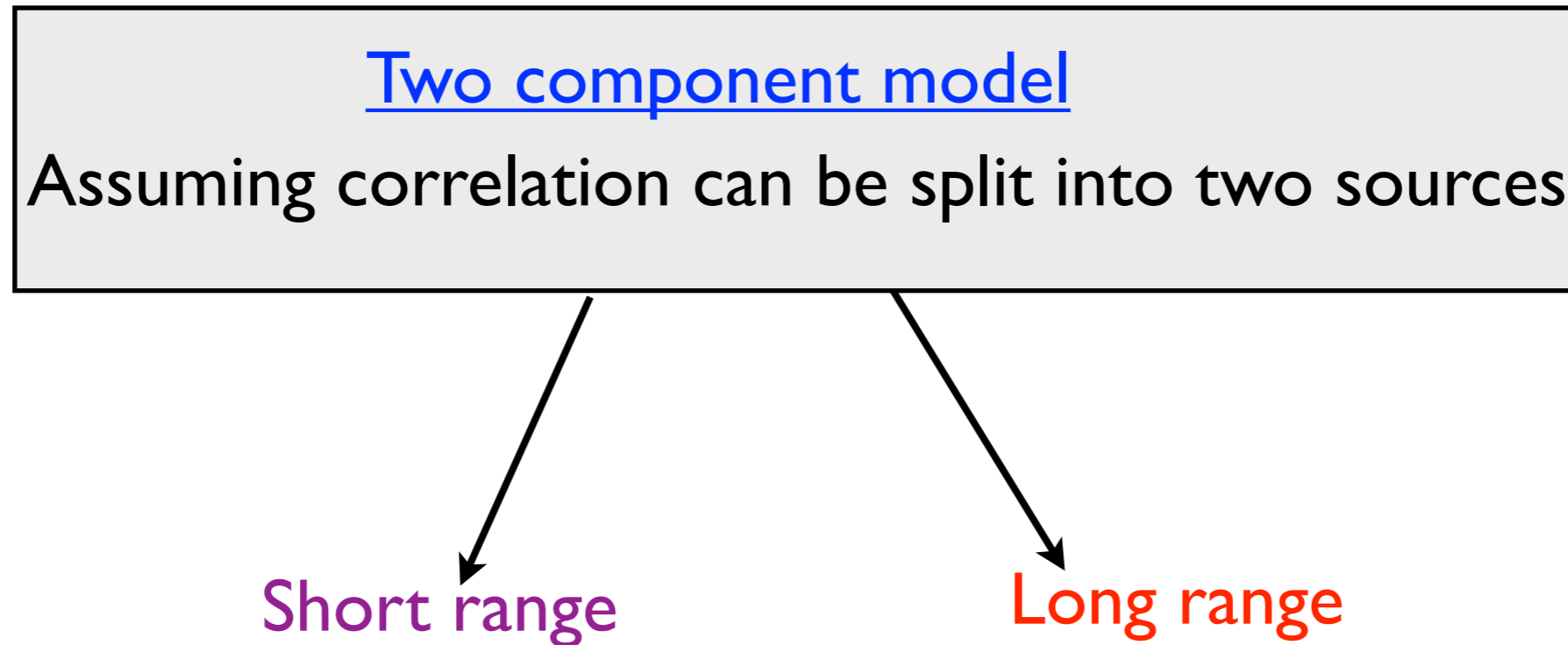
ALI-PREL-88934

- US > LS at low $\Delta\eta$
- Contributions from charge conservation and resonances

ALI-PREL-88938

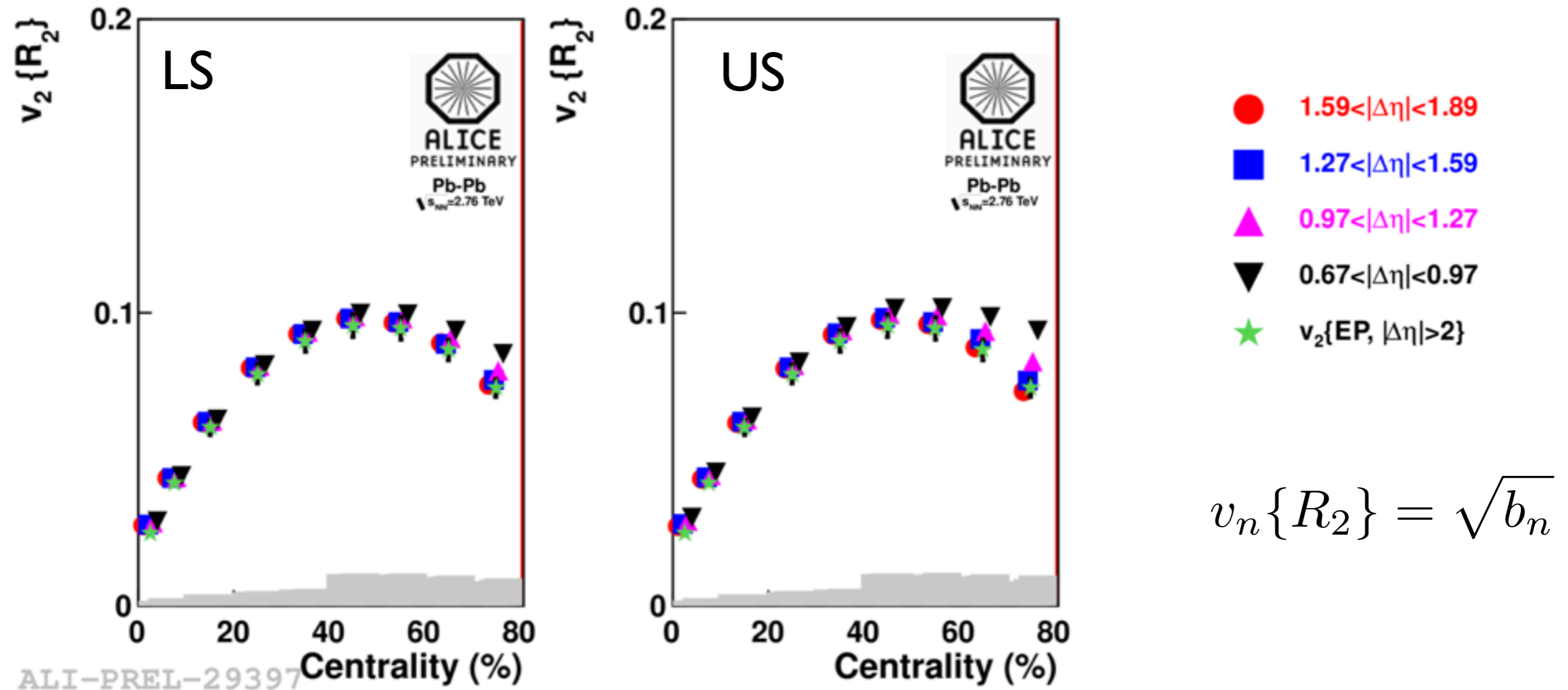
- US and LS merge at large $\Delta\eta$
- Earlier in $c_2\{\Delta p_T \Delta p_T\}$

What do we learn from the new observable $\Delta p_T \Delta p_T$?



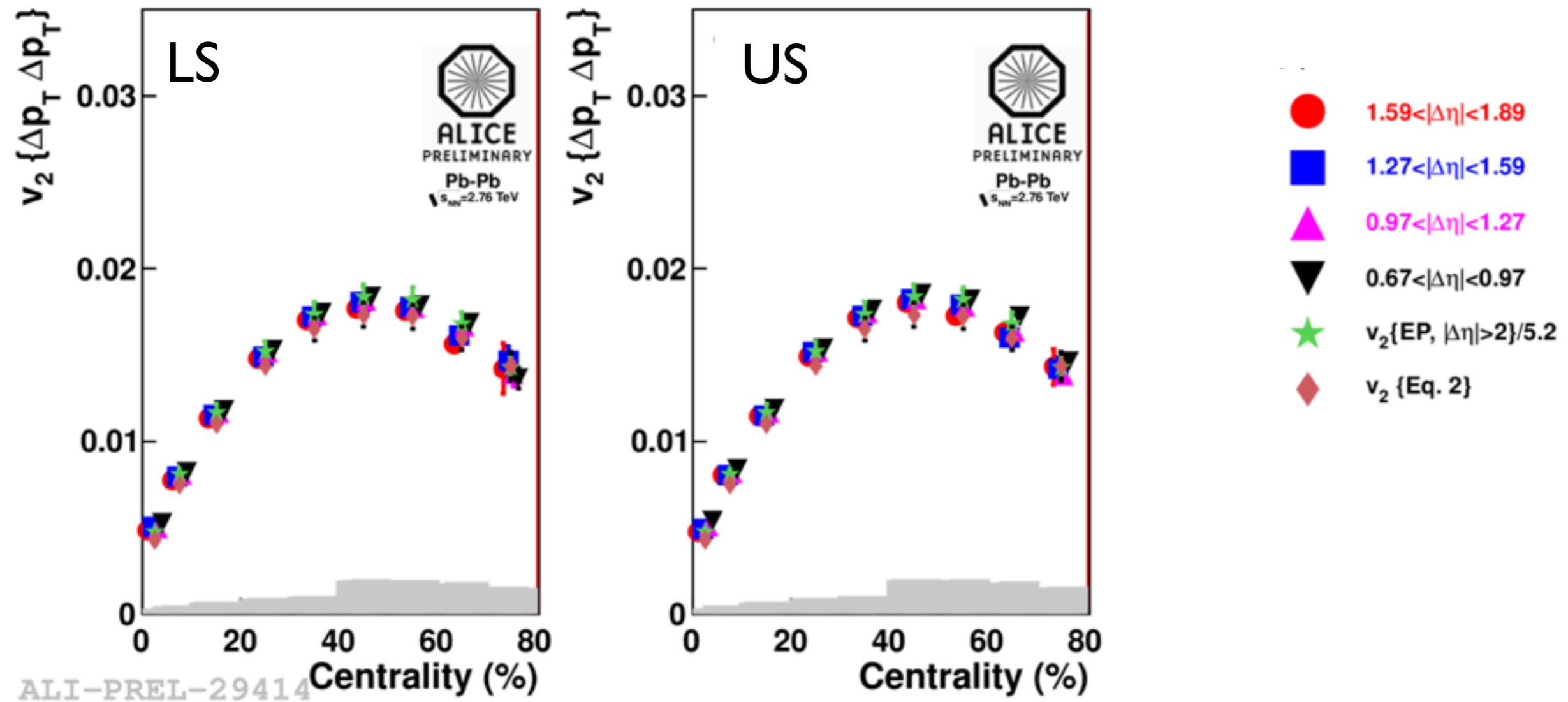
- The saturation of $\Delta p_T \Delta p_T$ at smaller $\Delta \eta$ than R_2 might be an indication that this observable has a “shorter” range correlation
- Is $\Delta p_T \Delta p_T = 0$ for long range correlation ?

$v_2\{R_2\}$ vs. centrality in Pb-Pb



- Measurements compared to $v_2\{EP\}$ by ALICE
 - Good agreement for large eta ranges $\Delta\eta > 1.0$

$v_2\{\Delta p_T \Delta p_T\}$ vs. centrality in Pb-Pb

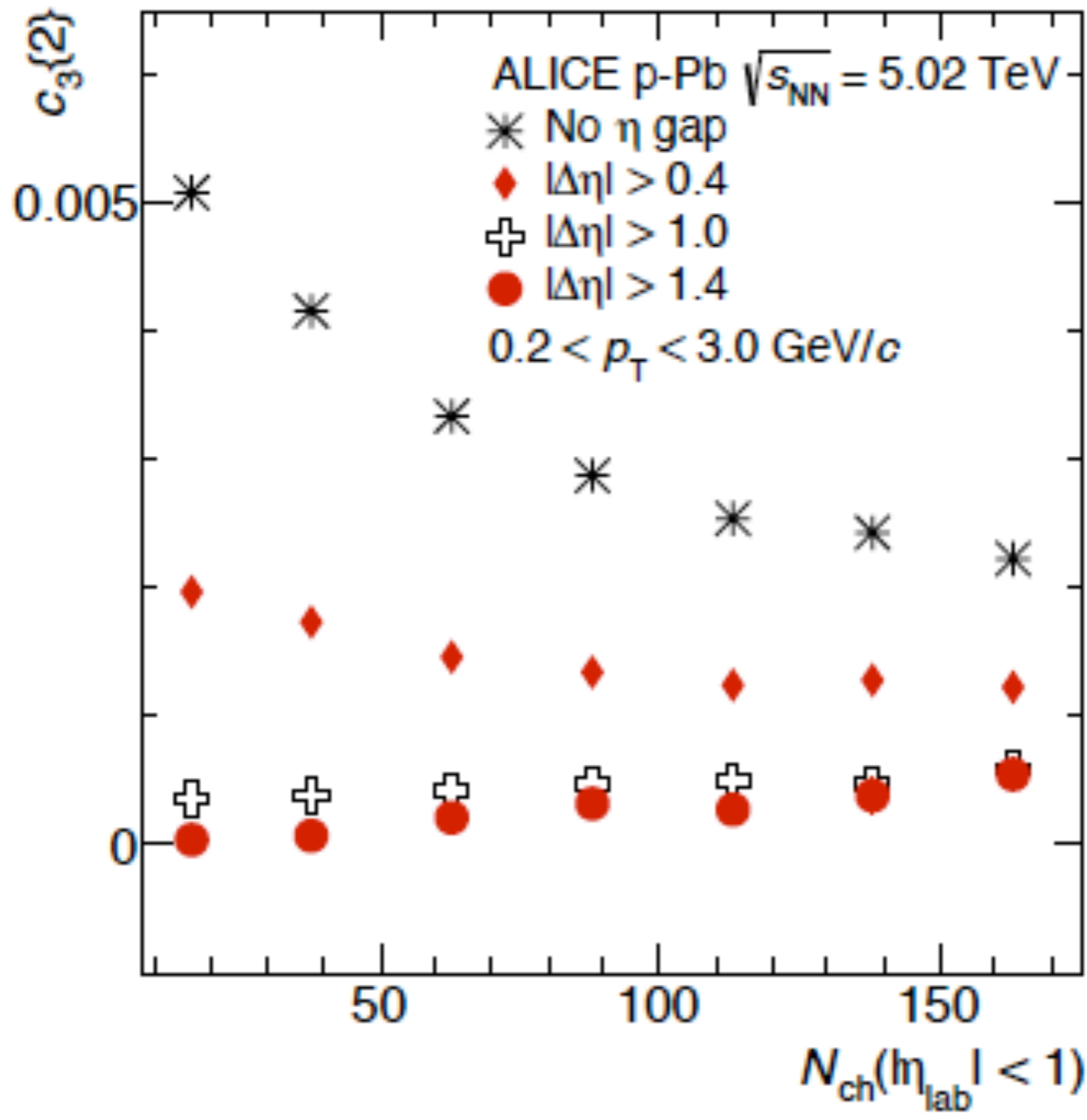


- Centrality dependence in excellent agreement w/scaled $v_2\{\text{EP}\}$ coefficients
- Excellent agreement w/coefficients “predicted” based on **geometric flow model** (Eq. 2) coefficients

Summary

- First measurement of the differential observable $\Delta p_T \Delta p_T$
- $\Delta p_T \Delta p_T$ shows an away-side double bump in most central (0-5%) Pb-Pb
 - Not seen in R_2 (0-5% central) correlation function
- Harmonic coefficient $c_1 \{ \Delta p_T \Delta p_T \}$ changes sign at lower $\Delta \eta$ than $c_1 \{ R_2 \}$
- Harmonic coefficient c_2 for $\Delta p_T \Delta p_T$ saturates earlier than R_2
 - $\Delta p_T \Delta p_T$ behaves more like a “shorter” range correlation
- Fourier coefficient of $\Delta p_T \Delta p_T$ scale according to geometrical model --- Yet another form of “factorization” !
 - Additional support for models describing correlations based on initial geometry and fluctuations in initial conditions

Back up Slides ...



[arXiv:1406.2474](https://arxiv.org/abs/1406.2474)

Two particle Number correlations and Flow Factorization

- Two-Particle Correlation -- Fourier Decomposition

$$\frac{dN_{pairs}}{d\Delta\phi} \propto 1 + 2 \sum_{n=1}^{\infty} V_{n\Delta}(p_1, p_2) \cos(n(\phi_1 - \phi_2))$$

- Flow Factorization

$$V_{n\Delta}(p_1, p_2) = \left\langle \left\langle e^{in(\phi_1 - \phi_2)} \right\rangle \right\rangle = \left\langle \left\langle e^{in(\phi_1 - \Psi_n)} \right\rangle \left\langle e^{in(\phi_2 - \Psi_n)} \right\rangle \right\rangle = \langle v_n \{2\}(p_1) v_n \{2\}(p_2) \rangle$$

- Factorization Expected When Correlations Driven by ...
 - Geometry Only
 - Jet Only(?)
- Factorization Expected to Break Down
 - When Jet & Geometry have 'similar' contributions

Transverse momentum correlation

$$\langle \Delta p_{t,1} \Delta p_{t,2} \rangle (\eta_1, \phi_1, \eta_2, \phi_2) = \frac{\int \rho_2(\eta_1, \phi_1, p_{t1}, \eta_2, \phi_2, p_{t2}) \Delta p_{t1} \Delta p_{t2} dp_{t1} dp_{t2}}{\int \rho_2(\eta_1, \phi_1, p_{t1}, \eta_2, \phi_2, p_{t2}) dp_{t1} dp_{t2}}$$

$$\langle \Delta p_t \Delta p_t \rangle (\Delta \eta, \Delta \varphi) = \frac{\sum_{\text{events accept}} \sum \Delta p_{t,1} \Delta p_{t,2}}{\sum_{\text{events}} n_{\text{pairs}}(\Delta \eta, \delta \varphi)}$$

where, $\Delta p_t = p_{t,i} - \langle p_t \rangle$

$$\langle p_{t,i} \rangle = \frac{\int \rho_1(\mathbf{p}_i) p_{t,i} dp_{t,i}}{\int \rho_1(\mathbf{p}_i) dp_{t,i}}$$

$\Delta\eta$ Slices - Fourier Decomposition Fits

$$F(\Delta\varphi) = b_0 + 2 \sum_{n=1}^6 b_n \cos(n\Delta\varphi)$$

Fourier harmonics:

$$c_n \{R_2\} = \frac{b_n}{b_0 + 1}$$

$$c_n \{\Delta p_T \Delta p_T\} = b_n$$

Can we test factorization for $\Delta p_T \Delta p_T$?

Hypothesis: the correlations in particles production are only due to flow

$$P(\eta, \phi, p_{\perp} | \{\Psi_n\}) = P(\eta, p_{\perp}) \left(1 + 2 \sum_n v_n(\eta, p_{\perp}) \cos(n(\phi - \Psi_n)) \right)$$

Sharma and Pruneau, Phys Rev C,
79, 024905, 2009, arXiv:0810.0716

Prediction for $\langle \Delta p_{t,1} \Delta p_{t,2} \rangle$

$$\langle \Delta p_t \Delta p_t \rangle(\eta_1, \eta_2, \Delta\phi) = \frac{2 \sum_n \left(v_n^p(\eta_1) - \langle p_t \rangle(\eta_1) v_n(\eta_1) \right) \left(v_n^p(\eta_2) - \langle p_t \rangle(\eta_2) v_n(\eta_1) \right) \cos(n\Delta\phi)}{1 + 2 \sum_n v_n(\eta_1) v_n(\eta_1) \cos(n\Delta\phi)}$$

$$v_n(\eta) = \frac{1}{P(\eta)} \int P(\eta, p_t) v_n(\eta) dp_t \quad \text{Regular Flow Coefficients}$$

$$v_n^{p_t}(\eta) = \frac{1}{P(\eta)} \int P(\eta, p_t) v_n(\eta) p_t dp_t \quad \text{p}_T \text{ weighted Flow Coefficients}$$

Consequence: FDC of $\langle \Delta p_{t,1} \Delta p_{t,2} \rangle$ at large $\Delta\eta$ should scale with flow coefficients

Factorization

- 2-particle correlation -- Fourier decomposition (FDC)

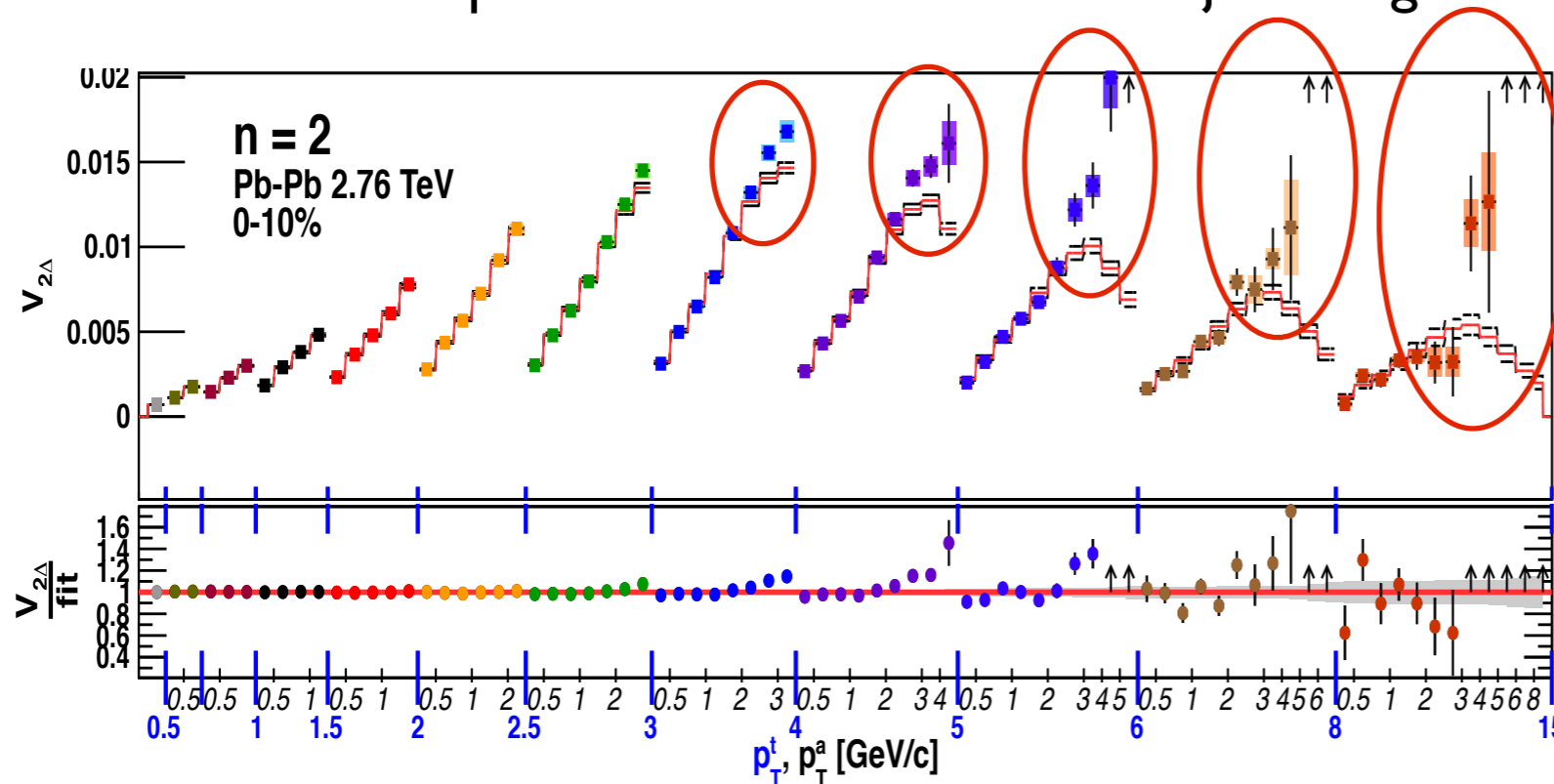
$$\frac{dN_{pairs}}{d\Delta\phi} \propto 1 + 2 \sum_{n=1}^{\infty} V_{n\Delta}(p_1, p_2) \cos(n(\phi_1 - \phi_2))$$

ALICE $v_n\{2\}$
PRL 107 032301 (2011)

- Factorization expected when correlations driven by geometry only

$$V_{n\Delta} = v_n(p_1) v_n(p_2)$$

- Factorization expected to break down when jet and geometry have similar contributions



Physics Letters B 708 (2012) 249–264

- Fit supports factorization at low p_T^a
⇒ suggests flow correlations.

- Fit deviates in jet-dominated high p_T^a region
⇒ collective description less appropriate

‘Geometrical flow model’

Sharma and Pruneau, Phys Rev C, 79, 024905, 2009, arXiv: 0810.0716

Can we test factorization for $\Delta p_T \Delta p_T$?

Consequence: FDC of $\langle \Delta p_{t,1} \Delta p_{t,2} \rangle$ at large $\Delta\eta$ should scale with flow coefficients