Pb-Pb results

Conclusions

 ↑ production in hadron collisions at forward rapidity with ALICE at the LHC

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# Bottomonium in hadron collisions

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The study of quarkonium ( $c\bar{c}$  and  $b\bar{b}$  mesons) production in hadron collisions plays an important role for different reasons:

- the  $Q\bar{Q}$  binding into quarkonium states is a non-perturbative process still not well understood: new data in pp collisions help to validate theoretical models;
- colour-screening model predicts the quarkonium suppression in AA collisions and less tightly bound states melt at a lower temperature: quarkonia give important information about the properties of the deconfined medium;
- cold nuclear matter effects are competing mechanisms: pA collisions allow to disentangle these effects from the hot ones.



- Theoretical calculations for bottomonium production are more robust due to the higher mass of the *b* quark.
- The probability of  $\Upsilon$  regeneration by  $b\bar{b}$  recombination is much smaller than that for the the J/ $\psi$ .
- The measurement of the  $\Upsilon$  in pA collisions allows a study of CNM effects in a different kinematic regime, complementing the J/ $\psi$  studies.

For these reasons this talk is focused only on  $\Upsilon$  production: results on charmonia will be shown by Javier Martin Blanco  $(J/\psi)$  and Marco Leoncino  $(\psi(2S))$ .

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## ALICE

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Bottomonium in hadron collisions ALICE pp results Pb-Pb results Conclusions A Large Ion Collider Experiment

- ALICE is the LHC experiment dedicated to the study of ultrarelativistic heavy-ion collisions.
- It participates also to the LHC pp and p-Pb program.
- At forward rapidity (2.5 < y < 4) quarkonium states are reconstructed via the dimuon decay down to transverse momentum (p<sub>T</sub>) equal to 0 with the Muon Spectrometer.



• V0 and T0 detectors are also used in the analyses for triggering purposes, while the SPD is used for primary vertex reconstruction.

Results in pp collisions at 
$$\sqrt{s} = 7$$
 TeV

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- The analysis is based on a data sample corresponding to an integrated luminosity of 1.35 pb<sup>-1</sup>.
- The inclusive production cross sections of Υ(1S) and Υ(2S) are measured as a function of p<sub>T</sub> and rapidity y.



- Results are published on EPJC 74 (2014) 2974.
- Data at 8 TeV are being analysed.

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#### Total and differential production cross sections



- The integrated values measured by ALICE in 2.5 < y < 4 and  $0 < p_T < 12 \text{ GeV}/c$  are:  $\sigma_{\Upsilon(1S)} = 54.2 \pm 5.0 \pm 6.7 \text{ nb}$   $\sigma_{\Upsilon(2S)} = 18.4 \pm 3.7 \pm 2.9 \text{ nb}$
- The  $p_T$  and y-differential cross sections compared to the values reported by LHCb (EPJC 72 (2012) 2025) show a good agreement for both resonances. They complement the measurements performed by CMS at midrapidity.

### Model comparison: *p*<sub>T</sub> dependence



- CSM predictions are scaled by a factor 1/0.6 to account for the feed-down from  $\Upsilon(2S)$ ,  $\Upsilon(3S)$  and  $\chi_b$ :
  - LO calculation underestimates the data for  $p_{\rm T}>4~{\rm GeV}/c;$
  - NLO calculation reproduces the data at low p<sub>T</sub>, but it still underestimates the cross section over the full range;
  - a good agreement is achieved at NNLO\*, over a limited p<sub>T</sub> range and with large uncertainties.
- NRQCD (with feed-down) overestimates the data, but the disagreement becomes smaller at higher p<sub>T</sub>.

#### Model comparison: y dep<u>endence</u>



- LO CSM calculations integrated over p<sub>T</sub> down to 0 are evaluated as a function of the rapidity with a large theoretical uncertainty.
- These calculations are scaled by the factor 1/0.6 and have no free parameters.
- The magnitude of the calculations is in agreement with the measurements.

# Results in Pb–Pb collisions at $\sqrt{s_{\rm NN}} = 2.76$ TeV

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- The analysed data sample corresponds to an integrated luminosity of 69 μb<sup>-1</sup>.
- The in-medium modification is evaluated through the nuclear modification factor:

$$R_{\rm AA} = \frac{Y_{\rm AA}}{\langle T_{\rm AA} \rangle \cdot \sigma_{\rm pp}}$$

( $Y_{AA}$  is the yield,  $T_{AA}$  is the nuclear overlap function,  $\sigma_{pp}$  is the pp reference cross section).

• Results are available on arXiv:1405.4493 (submitted to PLB).



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#### Comparison with CMS results



- ALICE and CMS (PRL, 109 (2012) 222301) measure inclusively the  $\Upsilon(1S)$  state down to  $p_{\rm T}=0$  in two complementary rapidity regions.
- In central collisions the suppression is stronger at forward rapidity than at midrapidity.
- The value of the  $\Upsilon(1S)$   $R_{AA}$  in 2.5 < y < 4 is significantly lower than in |y| < 2.4.

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#### Comparison with a dynamical model



- The evolving QGP is described by means of a dynamical model which includes the suppression of the different bottomonium states, but not CNM effects nor recombination.
- Two different initial temperature rapidity profiles: boost-invariant plateau and Gaussian. For each of them 3 values of  $4\pi\eta/s = \{1, 2, 3\}$ .
- None of the calculations reproduce the ALICE data. The rapidity trend measured by ALICE and CMS is opposite to what foreseen by the model.

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#### Comparison with a transport model



- The model accounts for both suppression and regeneration mechanisms.
- Cold nuclear matter effects are considered by means of an effective absorption cross section from 0 and 2 mb, including shadowing, nuclear absorption and Cronin effect.
- The measured  $R_{AA}$  is overestimated by the calculation which, however, reproduces the decreasing trend. The model predicts a  $R_{AA}$  almost constant as a function of rapidity, not supported by the data.

# Results in p–Pb collisions at $\sqrt{s_{\rm NN}} = 5.02$ TeV

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Pb 2.03<VCMS<3.53 -4.46<y\_ms<-2.96 The cold nuclear matter effects are quantified by the  $R_{\rm nPb}$ :

collisions.

$$R_{\rm pPb} = \frac{\sigma_{\rm pPb}}{A_{\rm Pb} \cdot \sigma_{\rm pp}}$$

 $(\sigma_{nPb} \text{ and } \sigma_{PP} \text{ are the } \Upsilon \text{ cross sections, } A_{Pb} \text{ is}$ the Pb mass number).

 $\chi^2/ndf = 1.07$ m..... (GeV/c2 ALI-PERF-51284 Events/(100 MeV/c<sup>2</sup>) Pb-p  $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ -4.46<y\_\_\_\_<-2.96 EKHORMAN:  $N_{T(1S)} = 152 \pm 19$  $m_{11(1S)} = 9.45 \pm 0.02 \text{ GeV}/c^2$  $\sigma_{T(1S)} = 0.159 \pm 0.022 \text{ GeV}/c^2$  $\chi^2$ /ndf = 0.80 m,,, (GeV/c2

m<sub>r(1S)</sub> = 9.45±0.02 GeV/c<sup>2</sup>  $\sigma_{1'(1S)} = 0.16 \pm 0.022 \text{ GeV}/c^2$  ALICE

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## $\Upsilon(1S) R_{pPb}$ measurements



- The  $\Upsilon(1S)$  is suppressed at forward rapidity, while at backward rapidity the  $R_{\rm pPb}$  is compatible with unity within uncertainties, disfavouring a strong gluon anti-shadowing.
- At positive  $y_{cms}$  the  $\Upsilon(1S)$  and  $J/\psi R_{pPb}$  are rather similar. At negative rapidity, the  $J/\psi R_{pPb}$  is systematically above that of  $\Upsilon(1S)$ , even if they are consistent within uncertainties.
- EPS09 shadowing at LO predicts within uncertainties the measured  $R_{pPb}$ at forward rapidity, while the NLO calculation underestimate the suppression of the  $\Upsilon(1S)$ .

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## $\Upsilon(1S) R_{pPb}$ compared to theoretical models



- The parton energy loss with EPS09 calculation reproduces well the data at forward rapidity, while data at backward rapidity are in a better agreement with parton energy loss only calculation.
- The calculation based on the CGC slightly underestimates the  $R_{pPb}$ , but it is not able to reproduce the  $J/\psi$  measurements in the same rapidity range.

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### Conclusions

#### pp collisions:

- the  $p_T$  and y-differential production cross sections are in good agreement with measurements by LHCb and complement the results at midrapidity from CMS;
- CSM calculations underestimate the data at large p<sub>T</sub>. The leading-p<sub>T</sub> NNLO helps to reduce the disagreement but with larger uncertainties.

Pb-Pb collisions:

- all models underestimate the suppression at forward rapidity.

#### p-Pb collisions:

- the  $\Upsilon(1S) R_{pPb}$  is consistent with unity at backward rapidity suggesting a small gluon anti-shadowing;
- models tend to overestimate the measurements and cannot describe the full rapidity dependence;