

# Heavy flavours in nucleus-nucleus collisions: quenching, flow and correlations

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# Heavy Flavor in the QGP: the conceptual setup

- Description of **soft observables** based on **hydrodynamics**, assuming to deal with **a system close to local thermal equilibrium** (no matter why);
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NB At high- $p_T$  the interest in heavy flavor is no longer related to thermalization, but to the study of the **mass** and **color charge dependence** of **jet-quenching** (not addressed in this talk)

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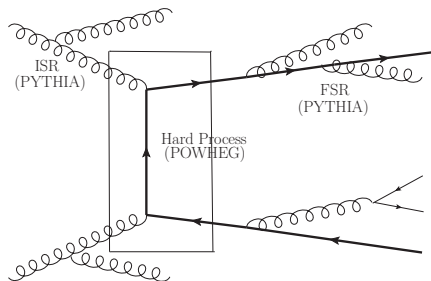
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NB for realistic temperatures  $g \sim 2$ , so that one can wonder *whether a charm is really “heavy”*. We will always assume also the last condition to be satisfied, and we will follow HQ propagation through a **Langevin equation** (*soft momentum-exchange limit of the Boltzmann equation*)



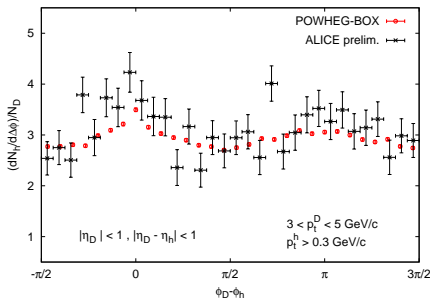
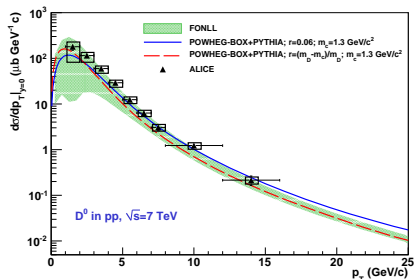
# Simulating the initial hard production



- Powerful pQCD tools<sup>1</sup> are available to simulate the initial  $Q\bar{Q}$  production, interfacing the output of a **NLO event-generator** (POWHEG, MC@NLO) for the **hard process** with a **parton-shower** (PYTHIA, HERWIG) describing **Initial** and **Final State Radiation**.
- This provides a *fully exclusive information on the final state*, also thanks to the simulation of the **Underlying Event** and **hadronization**.

<sup>1</sup>For a **systematic comparison** (POWHEG vs MC@NLO vs FONLL): M. Cacciari *et al.*, JHEP 1210 (2012) 137.

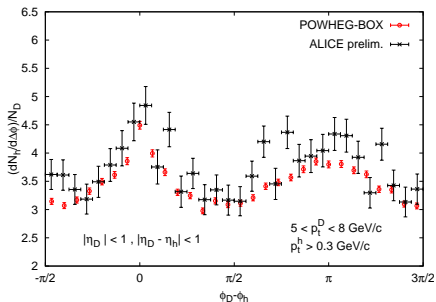
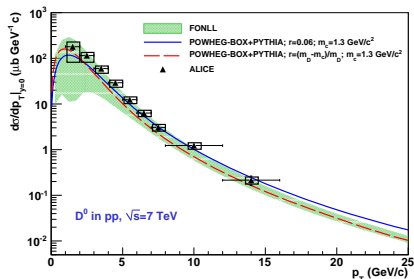
# HF production in $pp$ collisions: results



- Besides reproducing the inclusive  $p_T$ -spectra...<sup>2</sup>
- ...the POWHEG+PYTHIA setup allows also the comparison with  $D-h$  correlation data which start getting available
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<sup>2</sup>W.M. Alberico *et al*, Eur.Phys.J. C73 (2013) 2481

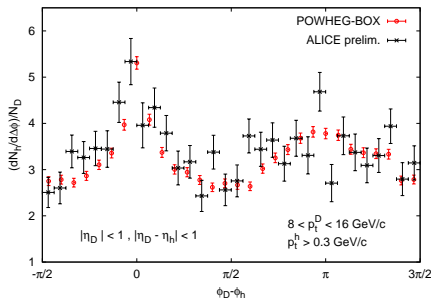
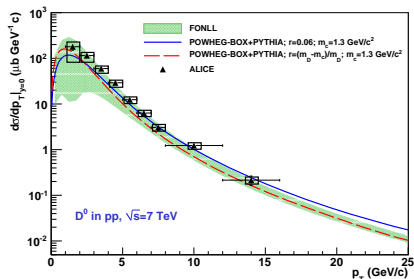
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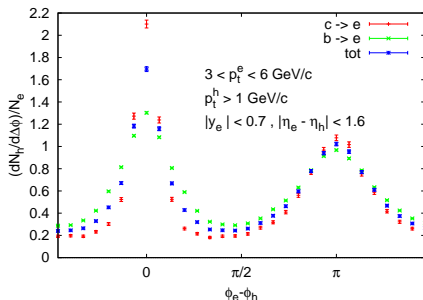
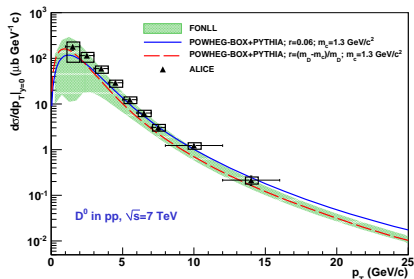
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# HF in nucleus-nucleus collisions

- The **theoretical tools**: **transport calculations** (for this talk: the relativistic **Langevin equation**)
- The main **questions** we wish to answer: **how close/far** are heavy quarks go **to/from thermalization**? Are final (hadronic) observables able to answer this question?

# The relativistic Langevin equation

The **Langevin equation** provides an algorithm to follow the dynamics of each heavy quark in the medium, updating its momentum and position

$$\frac{\Delta p^i}{\Delta t} = - \underbrace{\eta_D(\mathbf{p}) p^i}_{\text{determ.}} + \underbrace{\xi^i(t)}_{\text{stochastic}},$$

with the properties of the noise encoded in

$$\langle \xi^i(\mathbf{p}_t) \xi^j(\mathbf{p}_{t'}) \rangle = b^{ij}(\mathbf{p}_t) \frac{\delta_{tt'}}{\Delta t} \quad b^{ij}(\mathbf{p}) \equiv \kappa_{\parallel}(\mathbf{p}) \hat{p}^i \hat{p}^j + \kappa_{\perp}(\mathbf{p}) (\delta^{ij} - \hat{p}^i \hat{p}^j)$$

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**Transport coefficients** to calculate:

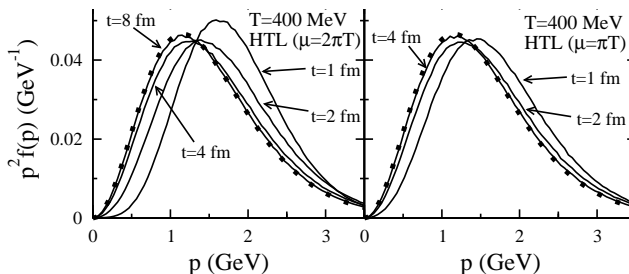
- **Momentum diffusion**  $\kappa_{\perp} \equiv \frac{1}{2} \frac{\langle \Delta p_{\perp}^2 \rangle}{\Delta t}$  and  $\kappa_{\parallel} \equiv \frac{\langle \Delta p_{\parallel}^2 \rangle}{\Delta t}$ ;
- **Friction** term (dependent on the **discretization scheme!**)

$$\eta_D^{\text{Ito}}(p) = \frac{\kappa_{\parallel}(p)}{2TE_p} - \frac{1}{E_p^2} \left[ (1 - v^2) \frac{\partial \kappa_{\parallel}(p)}{\partial v^2} + \frac{d-1}{2} \frac{\kappa_{\parallel}(p) - \kappa_{\perp}(p)}{v^2} \right]$$

fixed in order to insure approach to equilibrium (**Einstein relation**, **fluctuation-dissipation** theorem):



# A first check: thermalization in a static medium



For  $t \gg 1/\eta_D$  one approaches a relativistic Maxwell-Jüttner distribution<sup>3</sup>

$$f_{MJ}(p) \equiv \frac{e^{-E_p/T}}{4\pi M^2 T K_2(M/T)}, \quad \text{with } \int d^3p f_{MJ}(p) = 1$$

(Test with a sample of  $c$  quarks with  $p_0 = 2$  GeV/ $c$  and weak-coupling pQCD+HTL transport coefficients)

<sup>3</sup>A.B., A. De Pace, W.M. Alberico and A. Molinari, NPA 831, 59 (2009)

# The realistic case: expanding fireball

- $Q\bar{Q}$  pairs initially **produced with** the **POWHEG-BOX** package (with nPDFs) and **distributed** in the transverse plane **according to**  $n_{\text{coll}}(\mathbf{x}_{\perp})$  from (optical) Glauber model;

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- **update** of the HQ momentum and position **to be done** at each step **in the local fluid rest-frame**
  - $u^{\mu}(x)$  used to perform the boost to the **fluid rest-frame**;
  - $T(x)$  used to set the value of the **transport coefficients**with  $u^{\mu}(x)$  and  $T(x)$  fields taken from the output of **hydro codes**<sup>4</sup>;
- Procedure iterated **until hadronization**

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# Heavy quark thermalization?

Wondering **whether heavy quarks thermalize** entails a number of related questions...

- Are **theoretical tools** able to describe their approach to *thermal* equilibrium in a evolving medium?
- What are the **indications** coming **from experiment**? Are final hadronic/leptonic observables able to provide an unambiguous answer on *what happens in the partonic stage*?
- What could be the **role of experiments at larger  $\sqrt{s_{NN}}$** ? Higher temperature and radial flow, but also much harder initial  $Q\bar{Q}$  spectrum...

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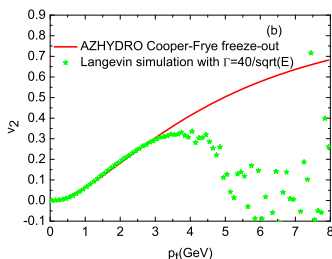
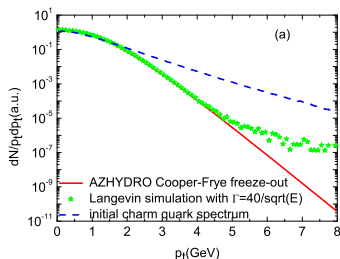
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NB thermal equilibrium of HQ's at the end of the QGP phase is assumed in the description of **hidden and open charm production** within the **Statistical Hadronization Model**: answering this question may support or rule out such an hypothesis

# Validation of the theoretical tools

In the limit of **large transport coefficients** heavy quarks should reach **local thermal equilibrium** and decouple from the medium as the other light particles, according to the Cooper-Frye formula:

$$E(dN/d^3p) = \int_{\Sigma_{fo}} \frac{p^\mu \cdot d\Sigma_\mu}{(2\pi)^3} \exp[-p \cdot u / T_{fo}]$$



This was verified to be actually the case (M. He, R.J. Fries and R. Rapp, PRC 86, 014903).

# Experimental indications

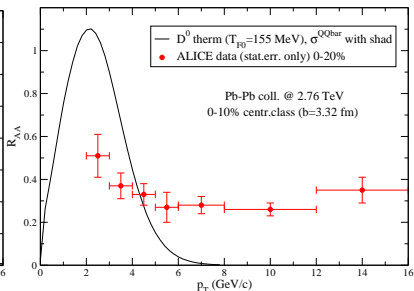
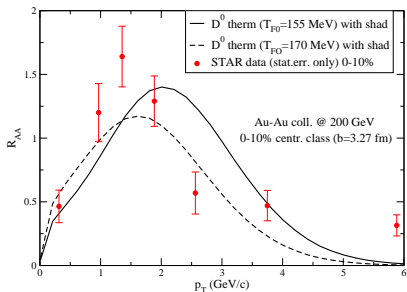
It is possible to compare the *experimental D-meson  $R_{AA}$*  with the *theoretical expectation in the case of kinetic equilibrium*

- Spectrum in  $pp$  given by POWHEG+PYTHIA setup
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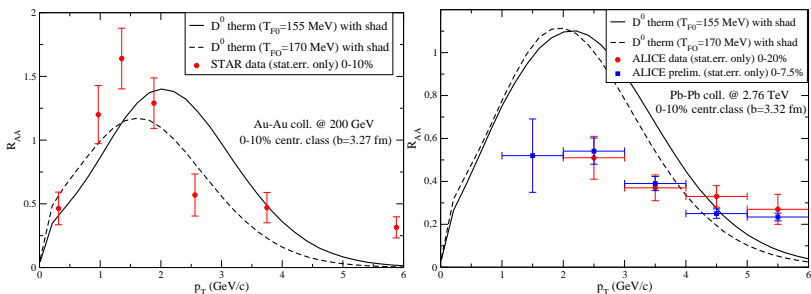
Evidence of a **bump from radial flow at RHIC**, while more data at low- $p_T$  (waiting for ALICE ITS upgrade) necessary at LHC; in any case **charm partially out of kinetic equilibrium**, at least at LHC



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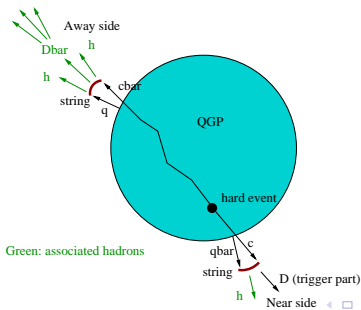
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# From quarks to hadrons

In-medium hadronization may affect the  $R_{AA}$  and  $v_2$  of final D-mesons due to the *collective flow of light quarks*. We tried to estimate the effect through this *model* interfaced to our POWLANG transport code:

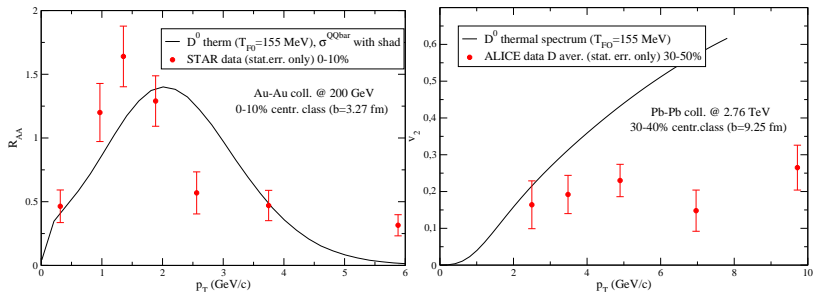
- At  $T_{dec}$  c-quarks coupled to light  $\bar{q}$ 's from a local *thermal distribution*, eventually boosted ( $u_{fluid}^\mu \neq 0$ ) to the lab frame;
- *Strings are formed* and given to PYTHIA 6.4 to simulate their fragmentation and produce the final hadrons ( $D + \pi + \dots$ )

One can address the study of  $D-h$  and  $e-h$  correlations in AA collisions



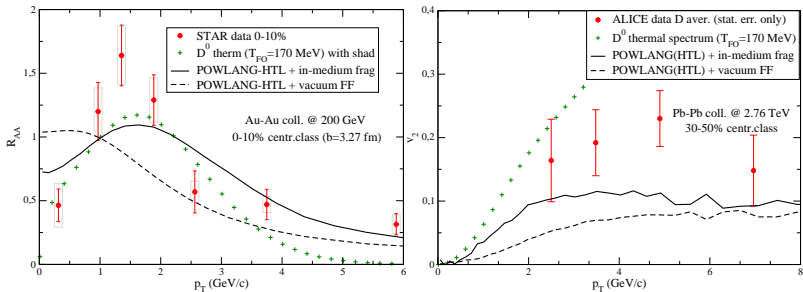
# From quarks to hadrons: effect on $R_{AA}$ and $v_2$

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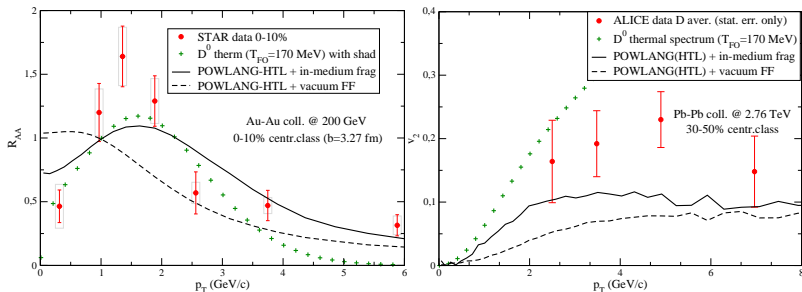
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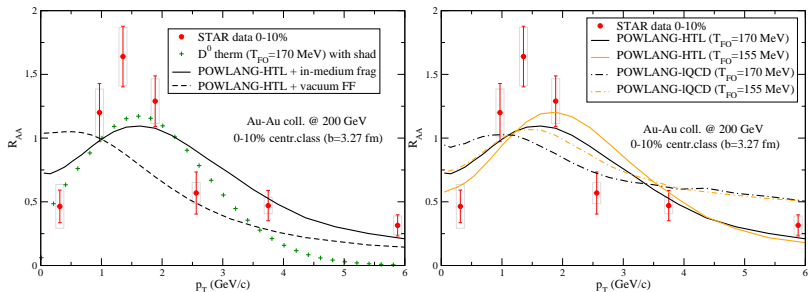
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However, comparing *transport results with/without the boost* due to  $u_{\text{fluid}}^\mu$ , at least part of the effect might be due to the **radial and elliptic flow of the light partons** from the medium picked-up at hadronization. **Rescattering in the hadronic phase** and its effect on  $v_2$  should be also investigated (in progress)!

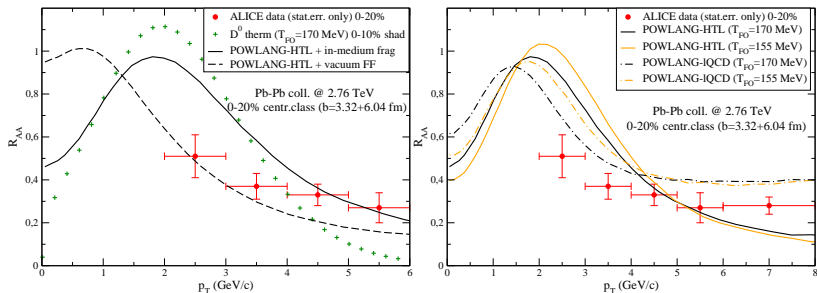
# D-meson $R_{AA}$ at RHIC



It is possible to perform a systematic study of different choices of

- **Hadronization** scheme (left panel)
- **Transport coefficients** (weak-coupling pQCD+HTL vs non-perturbative I-QCD) and **decoupling temperature** (right panel)

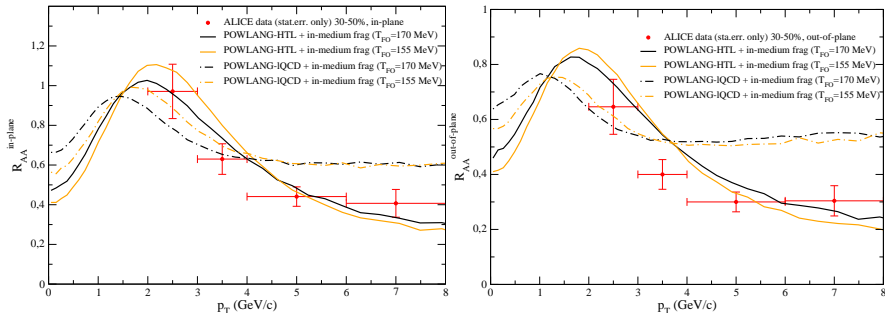
# D-meson $R_{AA}$ at LHC



Experimental data for central (0–20%) Pb-Pb collisions at LHC display a strong quenching, but – at least with the present bins and  $p_T$  range – don't show strong signatures of the bump from radial flow predicted by “thermal” and “transport +  $Q\bar{q}_{\text{therm}}$ -string fragmentation” curves.

# D meson $R_{AA}$ : in-plane vs out-of-plane

One can study di  $R_{AA}$  in- and out-of-plane in non-central (30–50%) Pb-Pb collisions at LHC:

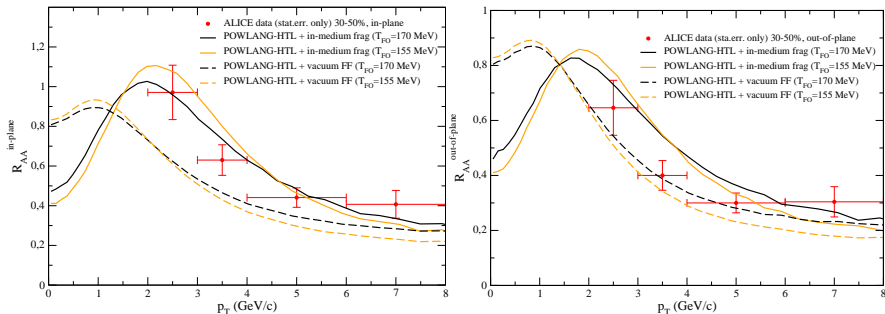


- Data better described by weak-coupling (pQCD+HTL) transport coefficients;



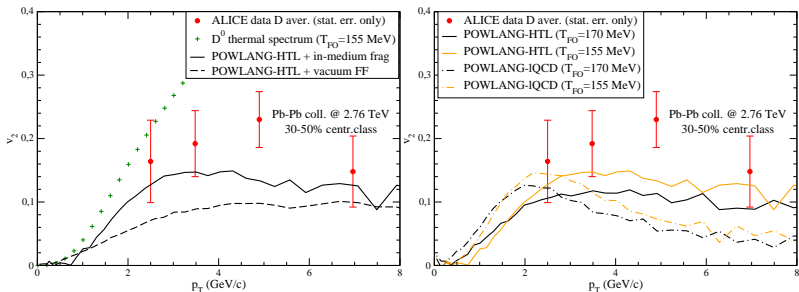
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- $Q\bar{q}_{\text{therm}}$ -string fragmentation describes data slightly better than in-vacuum independent Fragmentation Functions.

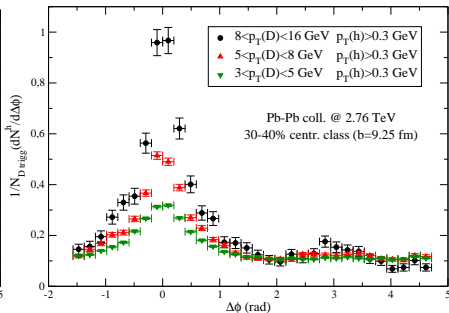
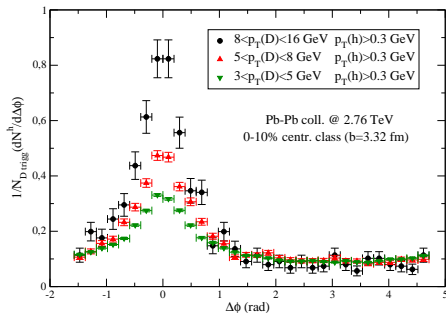
# D-meson $v_2$ at LHC



Concerning  $D$ -meson  $v_2$  in non-central (30–50%) Pb-Pb collisions:

- $Q\bar{q}_{\text{therm}}$ -string fragmentation routine significantly improves our transport model predictions compared to the data;
- HTL curves with a lower decoupling temperature display the best agreement with ALICE data

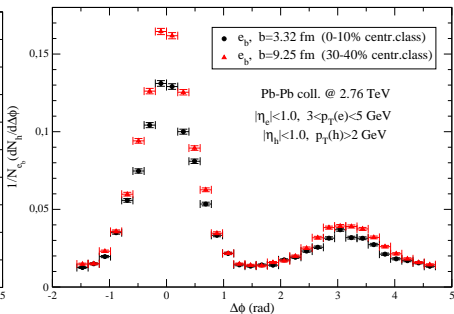
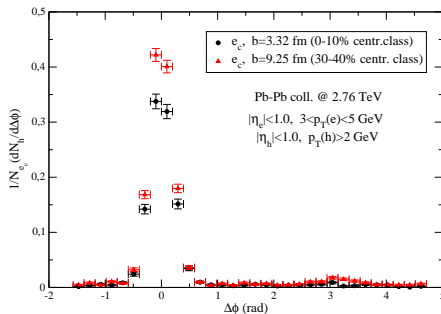
# Azimuthal correlations: $D-h$



Away-side peak strongly suppressed  
both in central and semi-central collisions

# Azimuthal correlations: $e-h$

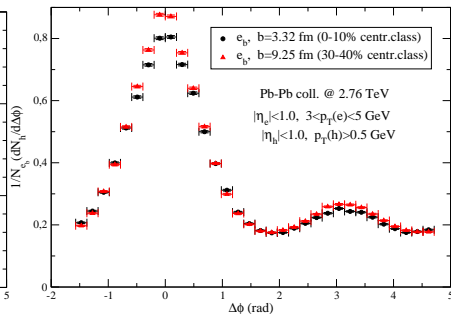
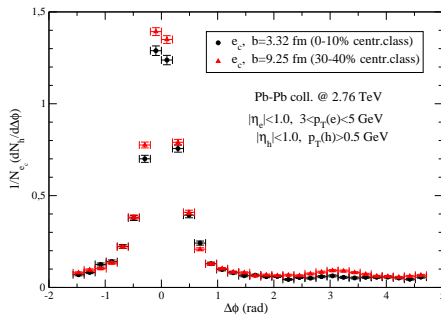
We plot the separate  $e_c$  (left) and  $e_b$  (right) contributions from charm and beauty decays



- charm away-side peak always strongly suppressed for any centrality and  $p_T^{\text{ass}}$  cut;
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# Summary and perspectives

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- We have presented a **new hadronization routine** (**recombination followed by string fragmentation**) interfaced to our partonic transport code, which **improves the agreement** of the latter **with experimental data**:
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- We have applied our setup to the study of **azimuthal heavy-flavour** ( $D-h$  and  $e-h$ ) **correlations**. **Which information can be obtained on the HQ-medium interaction?**