# ${\mathrm J}/\psi$  and  $\psi(2S)$  measurement in  $p{+}p$  collisions at  $\sqrt{s}=$ 200 and 500 GeV in the STAR experiment

Barbara Trzeciak<sup>1</sup> for the STAR Collaboration

<sup>1</sup>Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Brehova 7, 115 19 Praha 1, Czech Republic

E-mail: trzecbar@fjfi.cvut.cz

Abstract. In this paper, results on the  $J/\psi$  cross section and polarization measured via the **Abstract.** In this paper, results on the  $J/\psi$  cross section and potarization measured via the dielectron decay channel at mid-rapidity in  $p + p$  collisions at  $\sqrt{s} = 200$  and 500 GeV in the STAR experiment are discussed. The first measurement of  $\psi(2S)$  to  $J/\psi$  ratio at  $\sqrt{s} = 500$ GeV is also reported.

### 1. Introduction

 $J/\psi$  and  $\psi(2S)$  are bound states of charm (c) and anti-charm (c) quarks. Charmonium physical states have to be colorless, however they can be formed via a color-singlet (CS) or color-octet (CO) intermediate  $c\bar{c}$  state. One of the first models of the charmonium production, the Color Singlet Model (CSM) [1], assumed that  $J/\psi$  is created through the color-singlet state only. This early prediction failed to describe the measured charmonium cross section which has led to the development of new models. For example, Non-Relativistic QCD (NRQCD) [1] calculations were proposed in which a  $c\bar{c}$  color-octet intermediate states, in addition to a color-singlet states, can bind to form charmonia.

However, the charmonium production mechanism in elementary particle collisions is not yet exactly known. For many years measurements of the  $J/\psi$  cross section have been used to test different  $J/\psi$  production models. While many models can describe relatively well the experimental data on the  $J/\psi$  cross section in  $p + p$  collisions [2–9], they have different predictions for the J/ $\psi$  polarization. Therefore, measurements of the J/ $\psi$  polarization may allow to discriminate among different models and provide new insight into the  $J/\psi$  production mechanism.

### 2. Charmonium measurements in STAR

In STAR, charmonia have been measured so far via the dielectron decay channel. The STAR detector [10] is a multi-purpose detector that has large acceptance at mid-rapidity,  $|\eta|$  < 1 with a full azimuthal coverage. Electrons can be identified using the Time Projection Chamber (TPC) [11] through ionization energy loss  $(dE/dx)$  measurement. The Time Of Flight (TOF) detector [12] greatly enhances the electron identification capability at low momenta where the  $dE/dx$  bands for electrons and hadrons cross each other. At high  $p_T$ , electron identification can be improved by the Barrel Electromagnetic Calorimeter (BEMC) [13] which measures electron energy and shower shape. The BEMC is also used to trigger on high- $p_T$  electrons (HT trigger). Minimum Bias (MB) events are triggered by the Vertex Position Detectors (VPD) [14].

# 3. J/ $\psi$  measurements in  $p{+}p$  at  $\sqrt{s} = 200 \,\, \mathrm{GeV}$

STAR has measurements in  $p+p$  at  $\sqrt{s} = 200$  GeV<br>STAR has measured inclusive  $J/\psi$   $p_T$  spectra and polarization in  $p+p$  collisions at  $\sqrt{s} = 200$ GeV via the dielectron decay channel  $(B_{ee} = 5.9\%)$  at mid-rapidity  $(|y| < 1)$ . These results are compared to different model predictions to understand  $J/\psi$  production mechanism in elementary collisions.

Left panel of Fig. 1 shows STAR low and high- $p_T$  measurements of  $J/\psi$   $p_T$  spectra [3, 15] compared to model predictions. The Color Evaporation Model (CEM) [16] for prompt  $J/\psi$  can describe the  $p_T$  spectrum reasonably well, except the region around  $p_T \approx 3 \text{ GeV}/c$  where it overpredicts the data. NLO NRQCD calculations with color-singlet and color-octet transitions [17] for prompt  $J/\psi$  match the data for  $p_T > 4$  GeV/c. NNLO\* CS model [18] for direct  $J/\psi$ production under-predicts the STAR data, but the prediction does not include contributions from  $\psi(2S)$ ,  $\chi_C$  and B-meson decays to  $J/\psi$ .



**Figure 1.** Left:  $J/\psi$  invariant cross section vs  $p_T$  in  $p+p$  collisions at  $\sqrt{s} = 200$  GeV at midrapidity at low [15] and high  $p_T$  [3] shown as blue squares and red circles, respectively, compared to different model predictions [16–18]. Right: Polarization parameter  $\lambda_{\theta}$  vs  $J/\psi$   $p_T$  for  $|y|$  < 1 [19] compared to the PHENIX measurement [20] and two model predictions [21, 22].

In  $p+p$  collisions at  $\sqrt{s} = 200 \text{ GeV}$  STAR has also measured  $\text{J}/\psi$  polarization parameter  $\lambda_{\theta}$  in the helicity frame at mid-rapidity and  $2 < p_T < 6$  GeV/c [19].  $J/\psi$  polarization is analyzed via the angular distribution of the decay electrons that is described by:  $\frac{d^2N}{d(\cos\theta)d\phi} \propto 1 + \lambda_{\theta}\cos^2\theta +$  $\lambda_{\phi} \sin^2 \theta \cos 2\phi + \lambda_{\theta\phi} \sin 2\theta \cos \phi$ , where  $\theta$  and  $\phi$  are polar and azimuthal angles, respectively;  $\lambda_{\theta}$ ,  $\lambda_{\phi}$  and  $\lambda_{\theta\phi}$  are the angular decay coefficients. The  $p_T$  dependence of  $\lambda_{\theta}$  is shown on the right panel of Fig. 1 with low- $p_T$  PHENIX results [20] and compared to NRQCD calculations [21] and the  $NLO^+$  CSM prediction [22]. A trend observed in the RHIC data is towards longitudinal polarization as  $p_T$  increases and, within experimental and theoretical uncertainties, the result is consistent with the  $NLO^+$  CSM model.

The inclusive  $J/\psi$  production is a combination of prompt and non-prompt  $J/\psi$ . The prompt J/ $\psi$  production consists of the direct one (∼60%) and feed-down from excited states  $\psi(2S)(\sim10\%)$  and  $\chi_C(\sim30\%)$ , while non-prompt J/ $\psi$  originate from B-hadron decays. STAR has estimated the contribution from B-meson decays using a measurement of azimuthal angular correlation between high- $p_T J/\psi$  and charged hadrons [2, 3]. The relative contribution of Bhadron decays to inclusive J/ $\psi$  yield is strongly  $p_T$  dependent and it is 10-25% for  $4 < p_T <$ 12  $GeV/c$ , as it is shown on the left panel of Fig. 2. The measurement is consistent with the FONLL+CEM prediction [23, 24].



**Figure 2.** Left: relative contribution from B-meson decays to inclusive  $J/\psi$  production in  $p+p$ **rigure 2.** Left: relative contribution from B-meson decays to inclusive  $J/\psi$  production in  $p+p$  at  $\sqrt{s} = 200$  GeV [3] compared to FONLL+CEM calculations [23, 24]. Right: ratio of  $\psi(2S)$  to at  $\sqrt{s} = 200$  GeV [5] compared to FONLE+CEM calculations [25, 24]. Kight: Tatio of  $\psi(25)$  to  $J/\psi$  in  $p + p$  collisions at  $\sqrt{s} = 500$  GeV from STAR (red circle) compared to results from other experiments at different energies.

# 4. J/ $\psi$  and  $\psi(2S)$  measurements in  $p+p$  at  $\sqrt{s} =$  500 GeV

In order to further test the charmonium production mechanism and constrain the feed-down contribution from the excited states to the inclusive  $J/\psi$  production, the  $J/\psi$  and  $\psi(2S)$  signals contribution from the excred states to the inclusive  $J/\psi$  production, the  $J/\psi$  and  $\psi$ (25) signals were extracted in  $p + p$  collisions at  $\sqrt{s} = 500$  GeV at mid-rapidity. The  $J/\psi$   $p_T$  spectrum were extracted in  $p + p$  comsions at  $\sqrt{s} = 300$  GeV at mid-rapidity. The  $J/\psi$   $p_T$  spectrum is shown on the left panel of Fig. 3. The STAR results at  $\sqrt{s} = 500$  GeV (full circles) are is shown on the left panel of Fig. 5. The STAK results at  $\sqrt{s} = 300$  GeV (fun circles) are compared to those at  $\sqrt{s} = 200$  GeV (open circles) and with measurements of other experiments in  $p+\bar{p}$  collisions at different energies. The STAR measurements cover  $p_T$  range of 4 - 20 GeV/c with a good precision. It was also observed that  $J/\psi$  cross section follows the  $x_T$  scaling:  $\frac{d^2\sigma}{2\pi p_T dp_T dy} = g(x_T)/(\sqrt{s})^n$ , where  $x_T = 2p_T/\sqrt{s}$ , with  $n = 5.6 \pm 0.2$  at mid-rapidity and  $p_T > 5$  GeV/c for a wide range of colliding energies [2]. At  $\sqrt{s} = 500$  GeV the same  $x_T$  scaling of high- $p_T$  J/ $\psi$  production is seen, as shown on the right panel of Fig. 3.

Right panel of Fig. 2 shows  $\psi(2S)/J/\psi$  ratio from STAR (red full circle) compared to measurements of other experiments at different colliding energies, in  $p + p$  and  $p + A$  collisions. The STAR data point is consistent with the observed trend, and no collision energy dependence of the  $\psi(2S)$  to  $J/\psi$  ratio is seen with the current precision.

The  $\psi(z)$  to  $J/\psi$  ratio is seen with the current precision.<br>The statistics available at  $\sqrt{s}$  = 500 GeV will allow us to extract the frame invariant polarization parameter, also in different reference frames, providing model independent information about the  $J/\psi$  polarization [25]. It will be possible to measure the azimuthal polarization parameter,  $\lambda_{\phi}$ , and improve precision of the  $\lambda_{\theta}$  measurement. Analysis of  $J/\psi$ polarization parameter,  $\lambda_{\phi}$ , and improve<br>polarization at  $\sqrt{s} = 500$  GeV is ongoing.

### 5. Summary

In summary, STAR has measured the inclusive  $J/\psi$  cross section and polarization in  $p+p$ in summary, 5.1AK has measured the inclusive  $J/\psi$  cross section and polarization in  $p+p$  collisions at  $\sqrt{s} = 200 \text{ GeV}$  as a function of  $p_T$ . The measurements are compared to different model predictions of the J/ $\psi$  production. The  $p_T$  spectrum is described well by the NRQCD calculations while the measured polarization parameter  $\lambda_{\theta}$  is consistent with the NLO<sup>+</sup> CSM executations while the measured polarization parameter  $\lambda_{\theta}$  is consistent with the NEO. GeW prediction. STAR new result for  $J/\psi$  at  $\sqrt{s} = 500$  GeV extends  $p_T$  reach up to 20 GeV/c. prediction. STAK hew result for  $J/\psi$  at  $\sqrt{s} = 500$  GeV extends  $p_T$  reach up to 20 GeV/c.<br>The first measurement of  $\psi(2S)/J/\psi$  ratio in  $p+p$  collisions at  $\sqrt{s} = 500$  GeV is reported and compared with results from other experiments. No collision energy dependence is observed.



Figure 3. J/ $\psi$  invariant cross section vs  $p<sub>T</sub>$ , left panel, and invariant cross section multiplied by  $\sqrt{s}^{5.6}$  vs  $x_T$ , right panel, in  $p+p$  collisions at  $\sqrt{s} = 500$  GeV at mid-rapidity shown as full circles compared to measurements at different energies.

### Acknowledgements

This publication was supported by the European social fund within the framework of realizing the project ,,Support of inter-sectoral mobility and quality enhancement of research teams at Czech Technical University in Prague", CZ.1.07/2.3.00/30.0034.

#### References

- [1] Braaten E, Fleming S and Yuan T C 1996 Ann. Rev. Nucl. Part. Sci. 46 197–235 (Preprint hep-ph/9602374)
- [2] Abelev B et al. (STAR Collaboration) 2009 Phys. Rev. C 80 041902 (Preprint 0904.0439)
- [3] Adamczyk L et al. (STAR Collaboration) 2013 Phys. Lett. **B722** 55–62 (Preprint 1208.2736)
- [4] Adare A et al. (PHENIX Collaboration) 2012 Phys.Rev. D85 092004 (Preprint 1105.1966)
- [5] Abe F et al. (CDF Collaboration) 1997 Phys. Rev. Lett. 79(4) 572–577
- [6] Acosta D et al. (CDF Collaboration) 2005 Phys.Rev. D71 032001 (Preprint hep-ex/0412071)
- [7] Aad G et al. (ATLAS Collaboration) 2011 Nucl. Phys. **B850** 387-444 (Preprint 1104.3038)
- [8] Khachatryan V et al. (CMS Collaboration) 2011 Eur.Phys.J. C71 1575 (Preprint 1011.4193)
- [9] Aaij R et al. (LHCb Collaboration) 2011 Eur.Phys.J. C71 1645 (Preprint 1103.0423)
- [10] Ackermann K et al. (STAR Collaboration) 2003 Nucl. Instrum. Meth. A 499 624–632
- [11] Anderson M et al. 2003 Nucl. Instrum. Meth. A 499 659–678 (Preprint nucl-ex/0301015)
- [12] Llope W J et al. 2012 Nucl. Instrum. Meth. A 661 110–113
- [13] Beddo M et al. (STAR Collaboration) 2003 Nucl. Instrum. Meth. A 499 725–739
- [14] Llope W J et al. 2004 Nucl. Instrum. Meth. A 522 252–273 (Preprint nucl-ex/0308022)
- [15] Kosarzewski L (STAR Collaboration) 2012 Acta Phys.Polon.Supp. 5 543–548
- [16] Frawley A D, Ullrich T and Vogt R 2008 Phys.Rept. 462 125–175 (Preprint 0806.1013)
- [17] Ma Y Q, Wang K and Chao K T 2011 Phys. Rev. D84 114001 (Preprint 1012.1030)
- [18] Artoisenet P et al. 2008 Phys.Rev.Lett. 101 152001 (Preprint 0806.3282)
- [19] Adamczyk L et al. (STAR Collaboration) 2013 Phys.Lett. B739 180 (Preprint 1311.1621)
- [20] Adare A et al. (PHENIX Collaboration) 2010 Phys. Rev. D 82(1) 012001
- [21] Chung H S, Yu C, Kim S and Lee J 2010 Phys. Rev. D 81(1) 014020
- [22] Lansberg J 2011 Phys. Lett. B 695 149–156 (Preprint 1003.4319)
- [23] Bedjidian M, Blaschke D, Bodwin G T, Carrer N, Cole B et al. 2004 (Preprint hep-ph/0311048)
- [24] Cacciari M, Nason P and Vogt R 2005 Phys.Rev.Lett. 95 122001 (Preprint hep-ph/0502203)
- [25] Faccioli P, Lourenco C, Seixas J and Wohri H K 2010 Eur. Phys. J. C 69 657–673 (Preprint 1006.2738)