

Effective Theory approach to Higgs analyses

Veronica Sanz (Sussex)

Benasque

Shouldn't we focus on discovery searches?

New Physics could be **heavy**
as compared with the channel we look at

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New Physics could be **heavy**
as compared with the channel we look at
EFT: expansion in higher-dimensional
operators (HDOs)

Buchmuller and Wyler. NPB (86)

model independent

Systematic studies

One operator, corrs.

Translation to thy

Advantages of EFT

Boring and necessary details

Bottom-up approach:
operators w / SM particles and symmetries,
plus the **newcomer**, the **Higgs**

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Realization of EWSB

Linear or non-linear

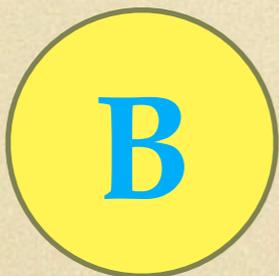
Boring and necessary details

Bottom-up approach:
operators w/ SM particles and symmetries,
plus the **newcomer**, the **Higgs**



Realization of EWSB

Linear or non-linear



And the Higgs could be

Weak doublet or singlet

Once this choice is made, expand...

$$\frac{1}{\Lambda^2}$$

Integrating out new physics

$$\frac{v^2}{f^2}$$

Non-linearity

$$U = e^{i\Pi(h)/f}$$

...order-by-order

For example, some operators
Higgs-massive vector bosons

ex.

$$\mathcal{L}_{eff} = \sum_i \frac{f_i}{\Lambda^2} \mathcal{O}_i$$

$$\mathcal{O}_W = (D_\mu \Phi)^\dagger \widehat{W}^{\mu\nu} (D_\nu \Phi)$$

$$\mathcal{O}_B = (D_\mu \Phi)^\dagger (D_\nu \Phi) \widehat{B}^{\mu\nu}$$

$$\mathcal{O}_{WW} = \Phi^\dagger \widehat{W}^{\mu\nu} \widehat{W}_{\mu\nu} \Phi$$

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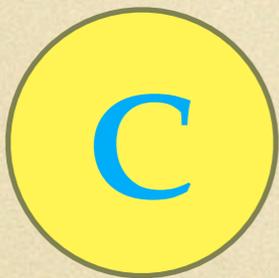
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UV theory: tree-level or loop

may need a model bias

ex. SILH

$$\frac{2igc_{HW}}{m_W^2} (D^\mu \Phi^\dagger) \widehat{W}_{\mu\nu} (D^\nu \Phi)$$

redundancies trade off operators using EOM

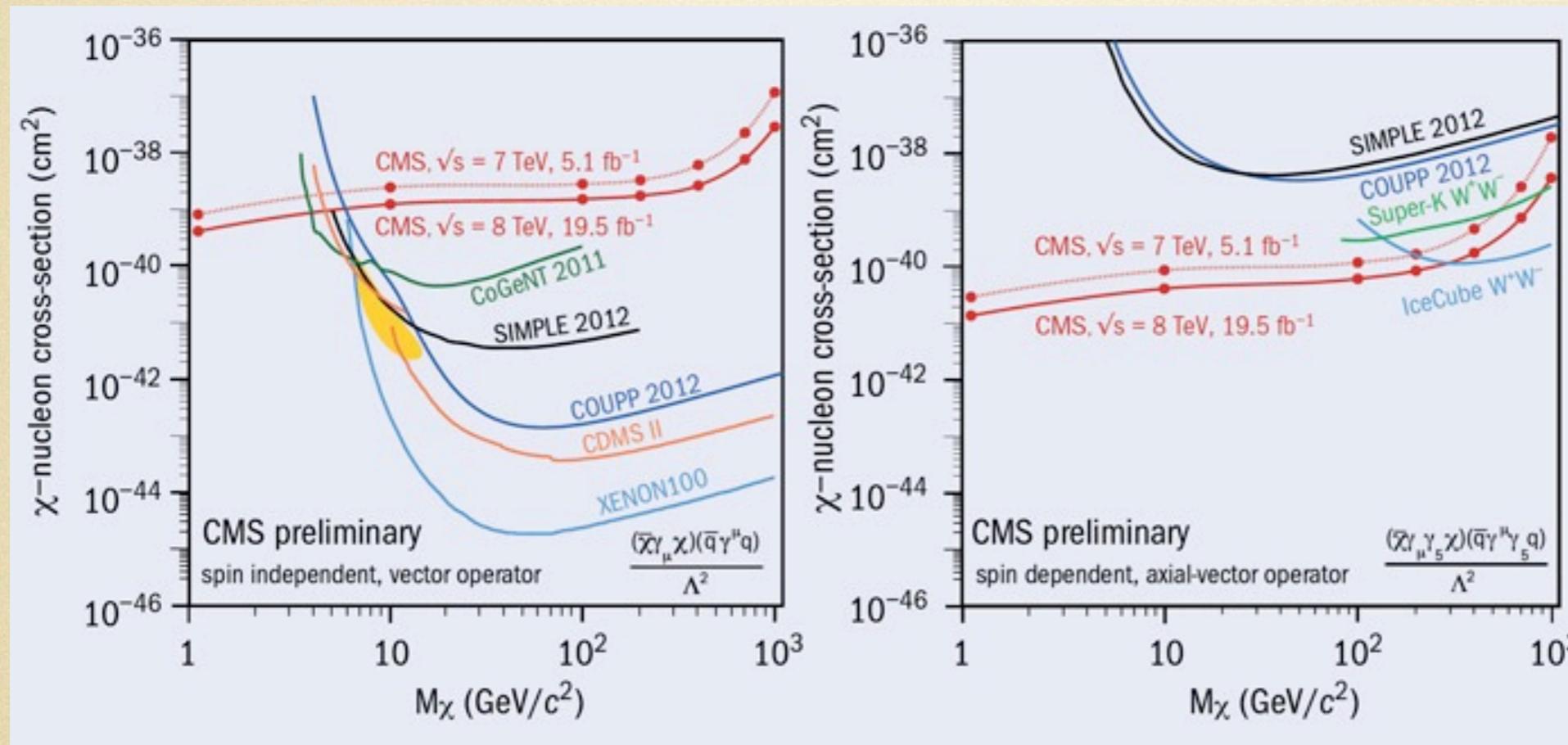
D Choice of basis

And, finally

Observables as a function
of HDOs coefficients

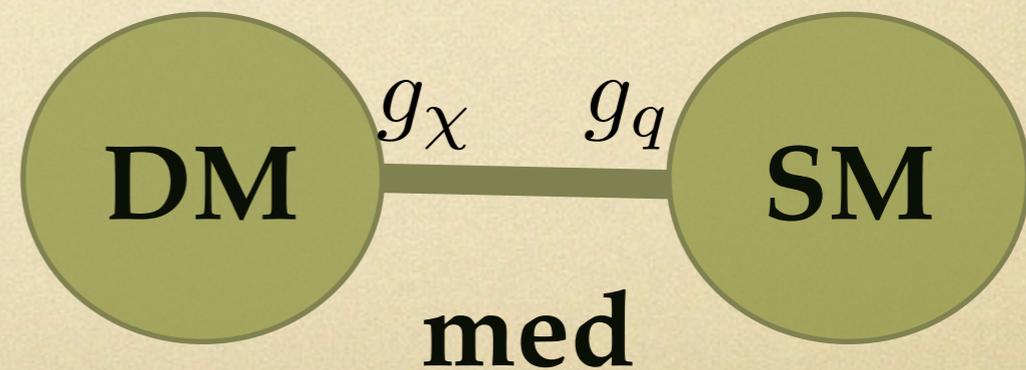
Beware of abusing the HDO framework

ex. Monojet searches as DM searches



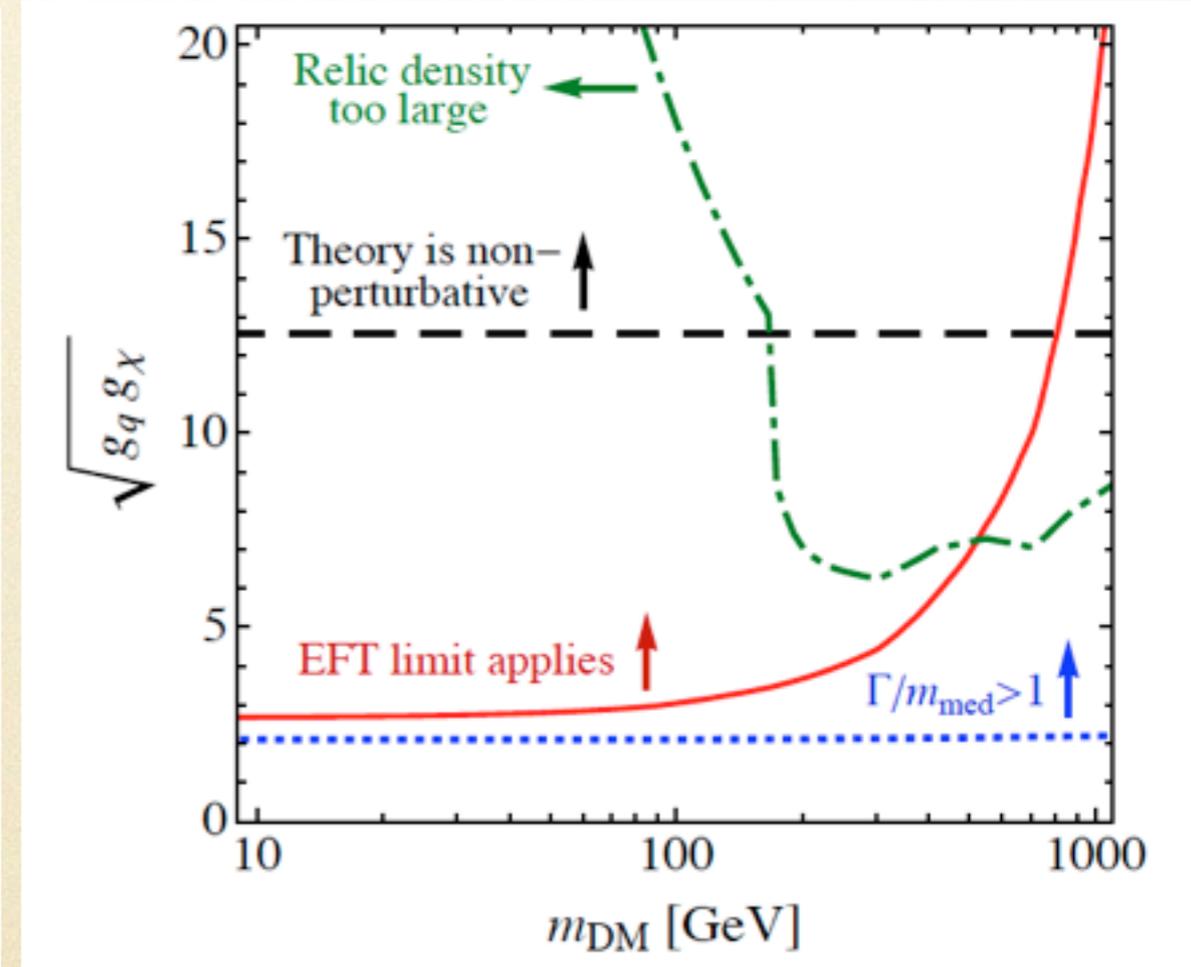
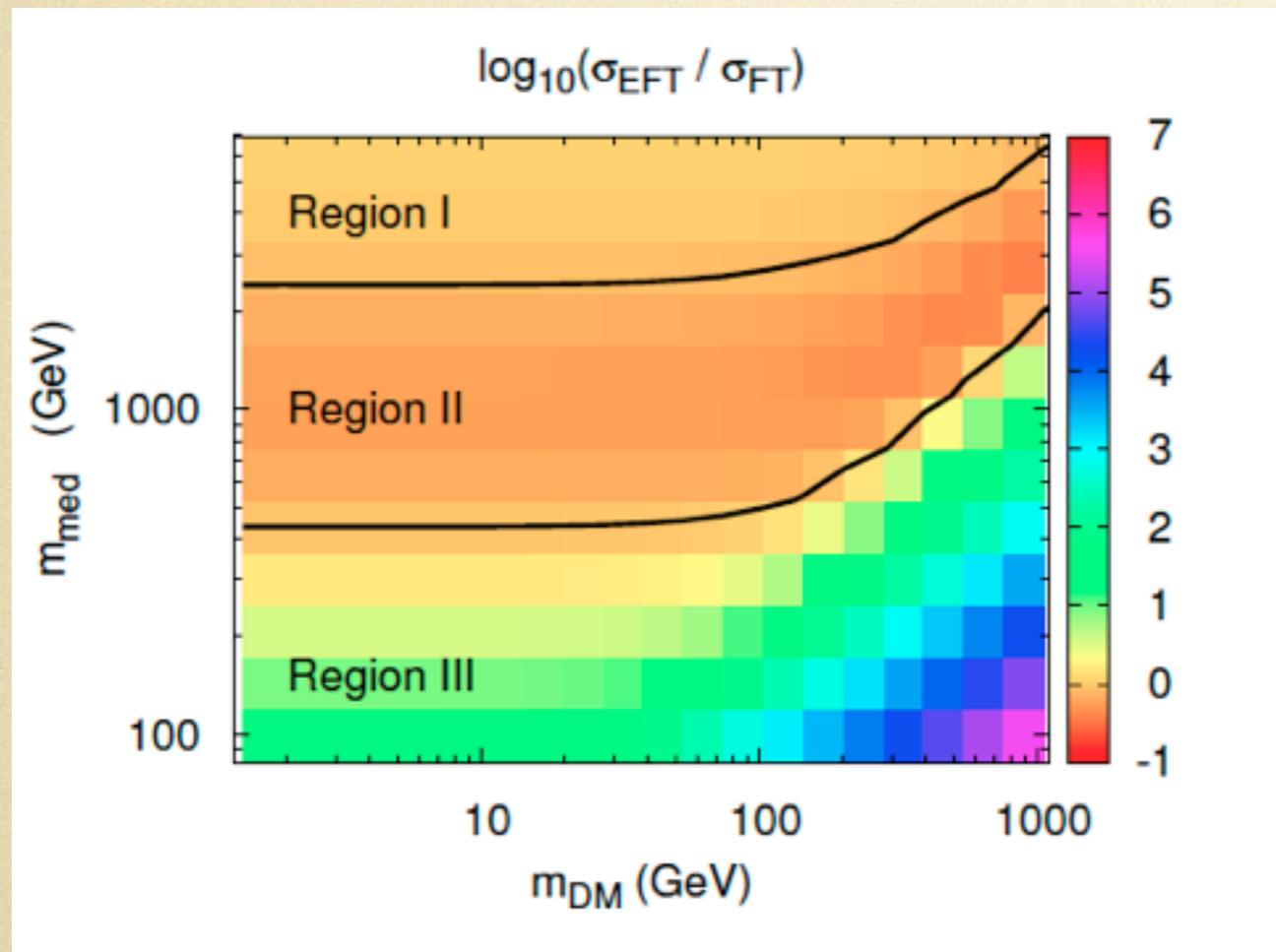
HDO for DM

$$\frac{1}{\Lambda^2} \bar{\chi} \gamma_\mu \chi \bar{q} \gamma^\mu q$$



Beware of abusing the HDO framework

ex. Monojet searches as DM searches



Buchmueller, Dolan and McCabe. 1308.6799

Limits valid to 20% if

$$m_{\text{med}} > 2.5 \text{ TeV}$$

$$\sqrt{g_\chi g_q} > 3$$

$$\Gamma/m_{\text{med}} > 1$$

Expansion may (non-trivially) *fail*

Some operators could have large coefficients

ex. $\frac{c_6}{\Lambda^2} (D^\mu \Phi^\dagger) \hat{W}_{\mu\nu} (D^\nu \Phi)$ could be loop-ind

$\frac{c_8}{\Lambda^4} (D^\rho \Phi^\dagger) (D_\rho \Phi) \hat{W}_{\mu\nu} \hat{W}^{\mu\nu}$ could dominate

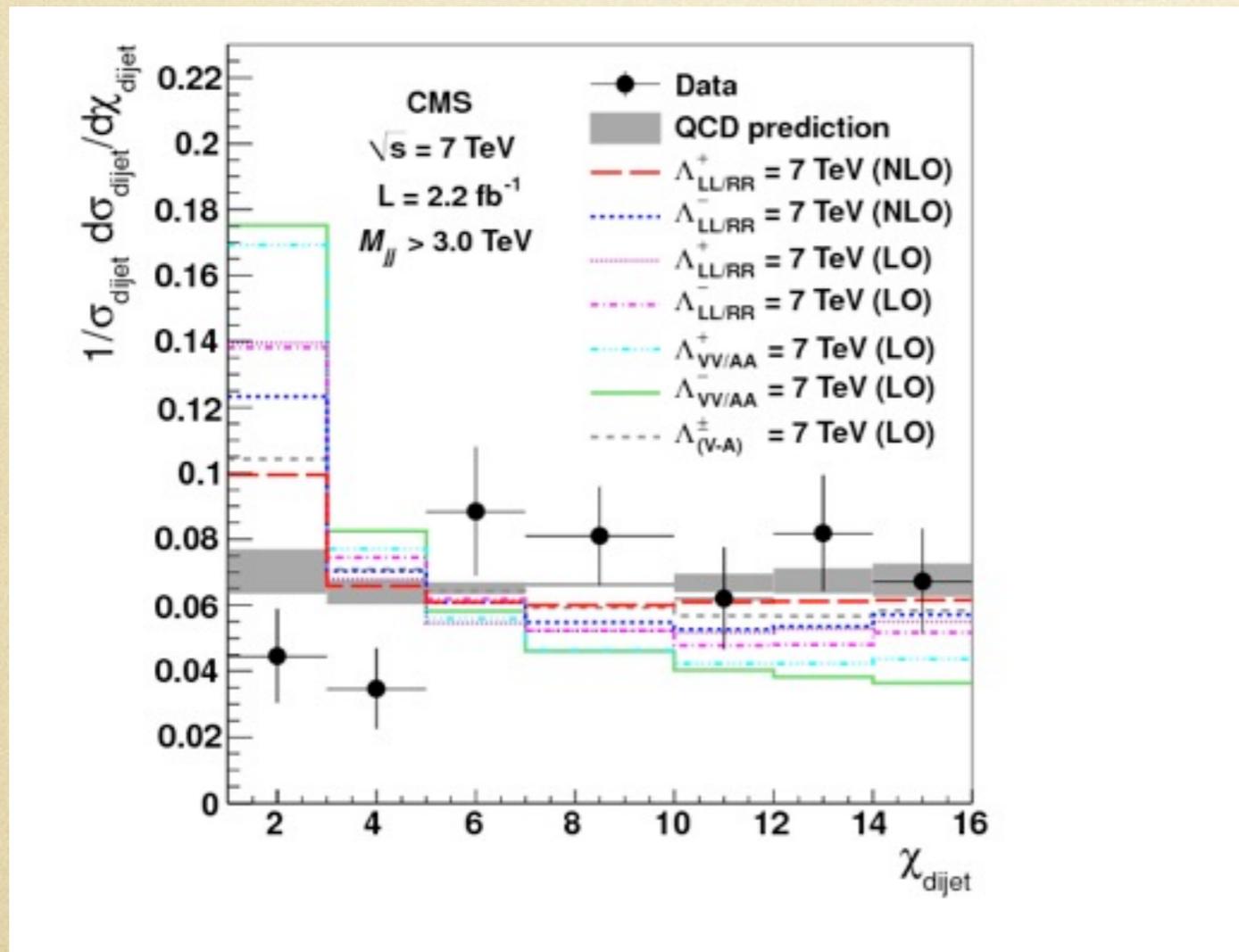
specially dangerous when

$\Lambda \sim m_W$ and high- p_T

How do we look for HDOs?

HDOs affect momentum dependence:
angular, p_T and inv mass distributions

Usual searches,

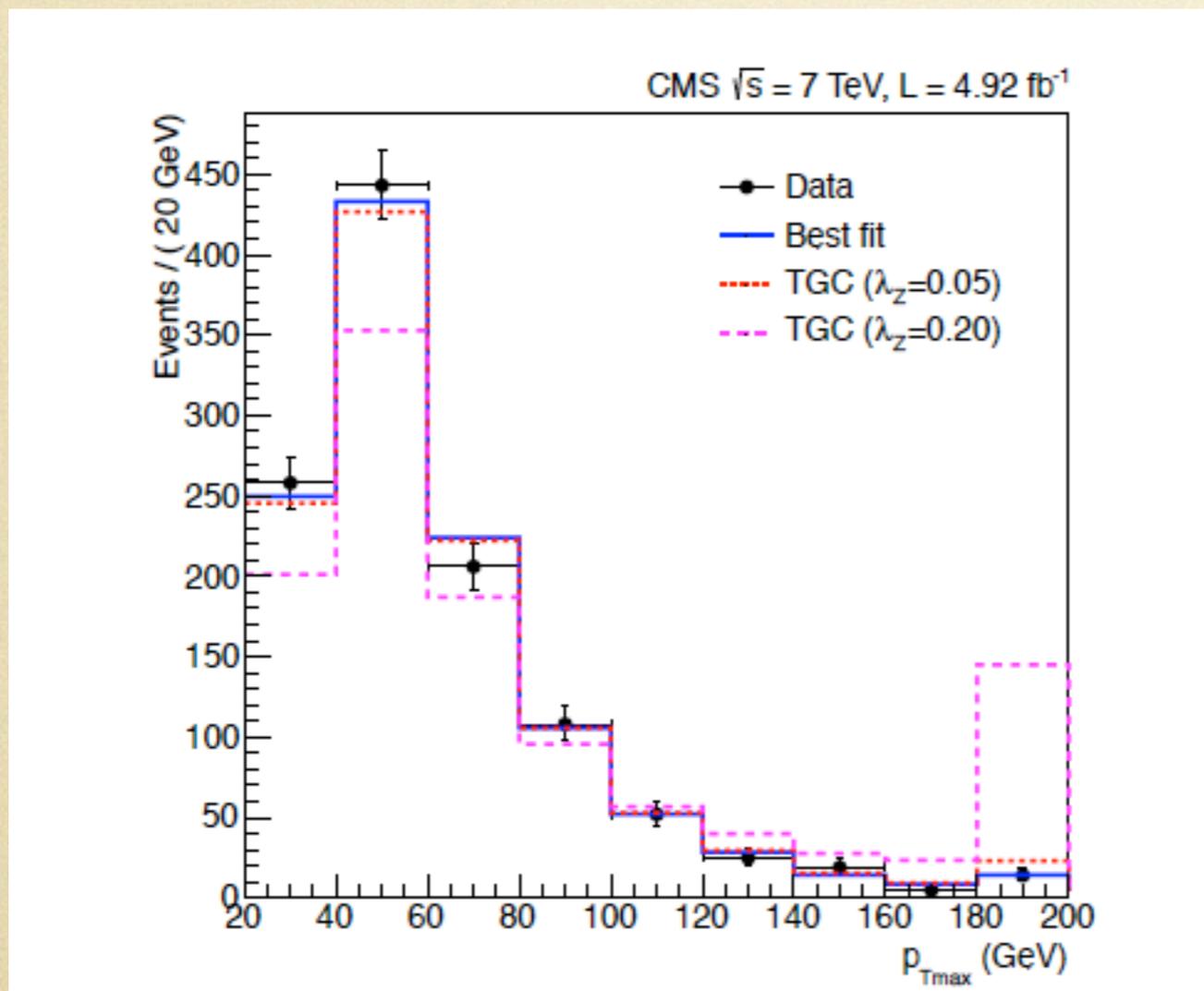


ex. dijet searches

Dijet angular distribution

HDOs affect momentum dependence:
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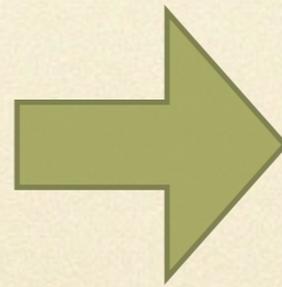
ex. TGCs

leading lepton p_T

And what about the Higgs?

Higgs anomalous couplings

HDOs generate
HVV interactions
with more
derivatives



$$-\frac{1}{4}h \underline{g_{hVV}^{(1)}} V_{\mu\nu} V^{\mu\nu}$$

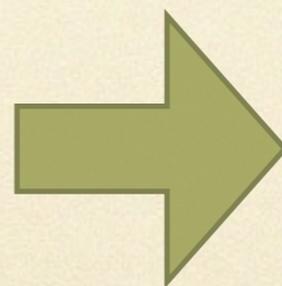
$$-h \underline{g_{hVV}^{(2)}} V_\nu \partial_\mu V^{\mu\nu}$$

$$-\frac{1}{4}h \underline{\tilde{g}_{hVV}} V_{\mu\nu} \tilde{V}^{\mu\nu}$$

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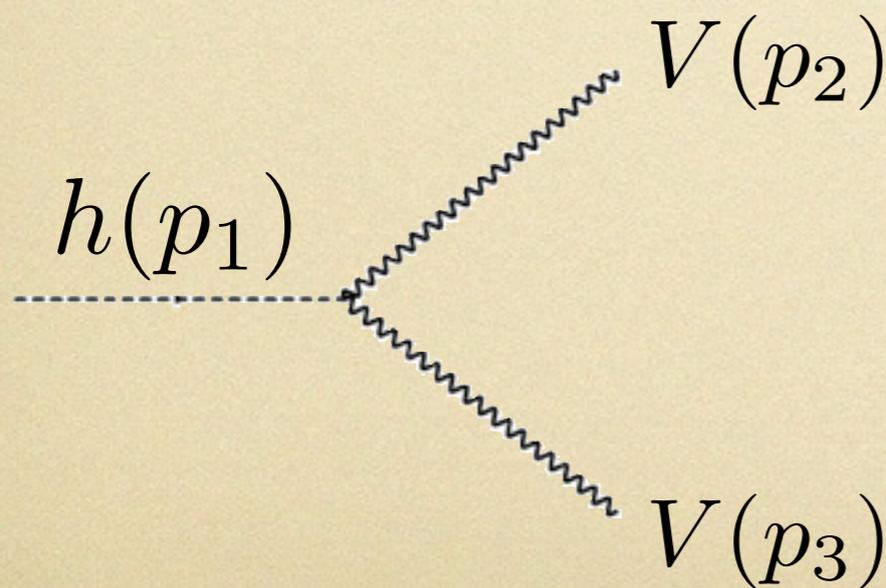
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 & -h \underline{g_{hVV}^{(2)}} V_\nu \partial_\mu V^{\mu\nu} \\
 & -\frac{1}{4} h \underline{\tilde{g}_{hVV}} V_{\mu\nu} \tilde{V}^{\mu\nu}
 \end{aligned}$$

ex. Feynman rule if $m_h > 2m_V$



$$i\eta_{\mu\nu} \left(\underline{g_{hVV}^{(1)}} \left(\frac{\hat{s}}{2} - m_V^2 \right) + 2 \underline{g_{hVV}^{(2)}} m_V^2 \right)$$

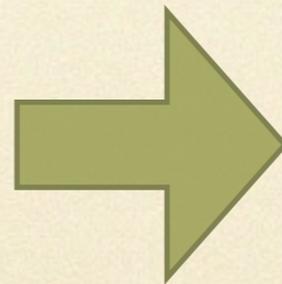
$$-i \underline{g_{hVV}^{(1)}} p_3^\mu p_2^\nu$$

$$-i \underline{\tilde{g}_{hVV}} \epsilon^{\mu\nu\alpha\beta} p_{2,\alpha} p_{3,\beta}$$

And what about the Higgs?

Higgs anomalous couplings

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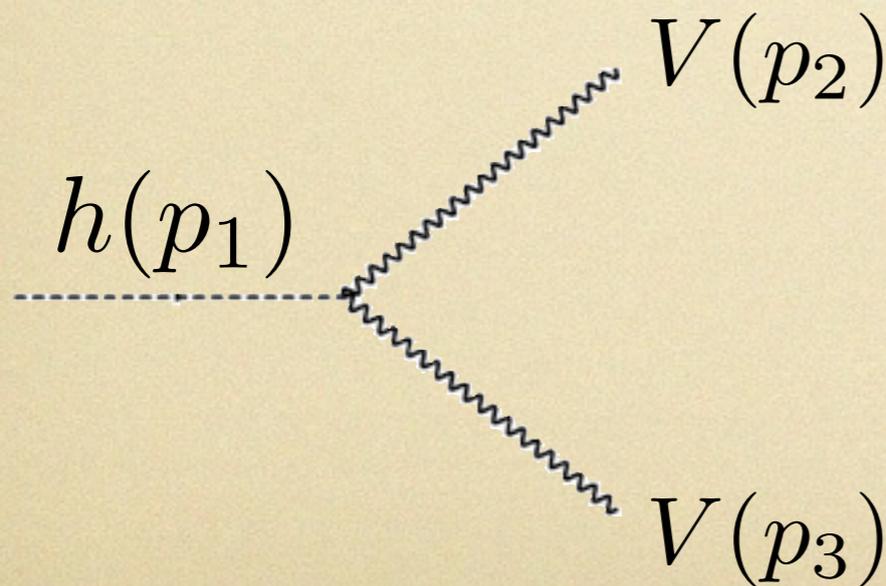


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total rates, COM, angular, inv mass and PT distributions

$$-i \underline{g_{hVV}^{(1)}} p_3^\mu p_2^\nu$$

$$-i \underline{\tilde{g}_{hVV}} \epsilon^{\mu\nu\alpha\beta} p_{2,\alpha} p_{3,\beta}$$

Framework for EFT studies

Higgs BRs

eHDECAY

Contino, Ghezzi, Grojean, Muhlleitner and Spira.

1303.3876

Production rates and kinematic distributions

depend on cuts

need radiation and detector effects

Simulation tools

We need a framework

$$\mathcal{L}_{eff} = \sum_i \frac{f_i}{\Lambda^2} \mathcal{O}_i$$

coefficients

**Collider
simulation**

observables

Limit coefficients
= new physics

A couple of tools

1. Feynrules HDOs involving Higgs and TGCs

Alloul, Fuks, VS. 1310.5150

links to CalcHEP, LoopTools, Madgraph...

HDOs -> Madgraph -> Pythia... -> FastSim / FullSim

2. QCD NLO HDOs involving Higgs and TGCs

VS and Williams. In prep.

MCFM and POWHEG

Pythia, Herwig... -> FastSim / FullSim

Before we move onto Higgs physics

Reviewing LEP constraints

Three CP-conserving operators affect TGCs

$$\begin{aligned}
 \mathcal{L}_{\text{TGC}}^{D=6} &= \frac{c_{WB} g_L g_Y}{m_W^2} B_{\mu\nu} W_{\mu\nu}^i H^\dagger \sigma^i H + \frac{i c_W g_L}{2m_W^2} \left(H^\dagger \sigma^i \overleftrightarrow{D}^\mu H \right) (D^\nu W_{\mu\nu})^i + \frac{c_{3W} g_L^3}{m_W^2} \epsilon^{ijk} W_{\mu\nu}^i W_{\nu\rho}^j W_{\rho\mu}^k \\
 &+ \tilde{c}_{WB} \frac{g_L g_Y}{m_W^2} \tilde{B}_{\mu\nu} W_{\mu\nu}^i H^\dagger \sigma^i H + \frac{\tilde{c}_{3W} g_L^3}{m_W^2} \epsilon^{ijk} W_{\mu\nu}^i W_{\nu\rho}^j \tilde{W}_{\rho\mu}^k.
 \end{aligned} \tag{8}$$

$$\begin{aligned}
 \mathcal{L}_{\text{TGC}}^+ &= i(1 + \delta g_1^V) (W_{\mu\nu}^+ W_\mu^- - W_{\mu\nu}^- W_\mu^+) V_\nu + i(1 + \delta \kappa_V) V_{\mu\nu} W_\mu^+ W_\nu^- \\
 &+ i \frac{\lambda_V}{m_W^2} W_{\mu\nu}^+ W_{\nu\rho}^- V_{\rho\mu} - g_5^V \epsilon_{\mu\nu\rho\sigma} (W_\mu^+ \partial_\rho W_\nu^- - \partial_\rho W_\mu^+ W_\nu^-) V_\sigma,
 \end{aligned}$$

c_W and c_{WB}
 affect Higgs physics
 and S-parameter,
 but more independent
 operators involved

dim-6 and TGCs

$$\begin{aligned}
 \delta \kappa_\gamma &= 4c_{WB}, \\
 \delta \kappa_Z &= -4 \frac{g_Y^2}{g_L^2} c_{WB} - \frac{g_L^2 + g_Y^2}{2g_L^2} \\
 \delta g_Z &= -\frac{g_L^2 + g_Y^2}{2g_L^2} c_W, \\
 \lambda_\gamma = \lambda_Z &= -6g_L^2 c_{3W}, \\
 \tilde{\kappa}_\gamma &= 4\tilde{c}_{WB}, \\
 \tilde{\kappa}_Z &= -4 \frac{g_Y^2}{g_L^2} \tilde{c}_{WB}, \\
 \tilde{\lambda}_\gamma = \tilde{\lambda}_Z &= -6g_L^2 \tilde{c}_{3W},
 \end{aligned}$$

LEP fit to WW observables

Falkowski, Fichet, Mohan, Riva, VS.
Les Houches proceedings. To appear.

Total cross section and diff distributions from
ALEPH, DELPHI, L3, OPAL, LEP Electroweak Collaboration.

Phys.Rept. 532

SM NLO from RACOONWW

$$c_{WB} = -0.01 \pm 0.03, \quad c_W = 1.18 \pm 0.56, \quad c_{3W} = -0.30 \pm 0.16,$$

partial cancellation in the direction $c_{3W} + 0.3 c_W$

due to polarization of main contribution to WW

the fit leads to no meaningful limits, strong sensitivity to
quadratic terms (= higher-orders)

to compare with setting by
hand $c_{3W}=0$

$$c_{WB} = 0.03 \pm 0.02, \quad c_W = 0.18 \pm 0.07,$$

Consequences:

1. TGCs at the LHC are terribly important (no cancellation there)
2. LEP TGCs do not constrain Higgs operators to the extent it was thought before by, e.g.

Pomarol and Riva. 1308.2803

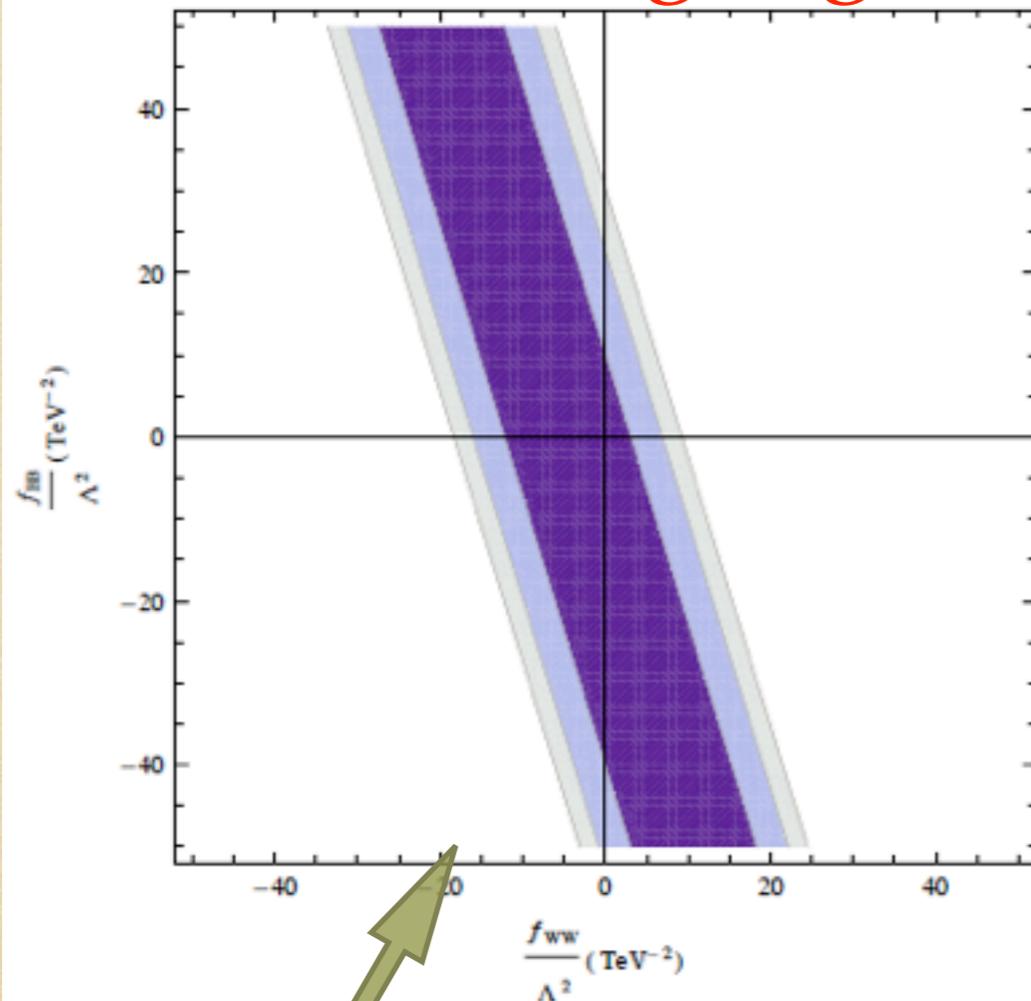
Corbett et al. 1211.4580

Higgs data truly complementary to TGCs

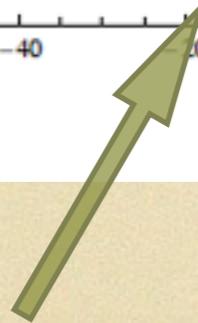
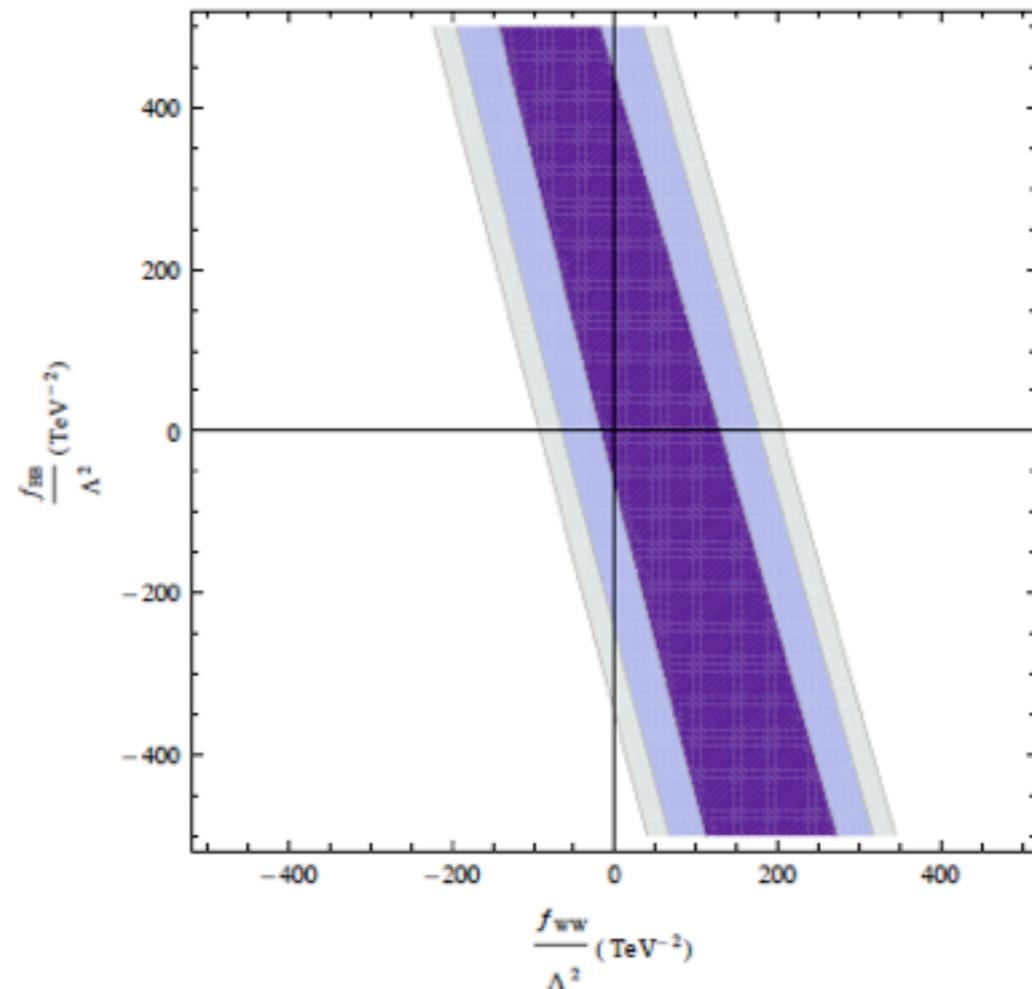
Same goes for EWPTs

renormalization **is** important

leading log



ren. scheme dep.



results of
Masso, VS. 1211.1320

Cheng, Dawson, Zhang. 1311.3107

Higgs data is complementary
to EWPT and TGCs

Complete analysis of dim-6 Higgs

Ellis, VS and You.
1404.3667

Number of independent operators in Higgs physics

CP-conserving

See Rohini's talk for CP-violating

In the SILH basis

$$\bar{c}_i \equiv \{ \bar{c}_H, \bar{c}_{t,b,\tau}, \bar{c}_W, \bar{c}_{HW}, \bar{c}_{HB}, \bar{c}_\gamma, \bar{c}_g \}.$$

(cB has been eliminated using S)

Other operators affect TGCs and EWPTs but not Higgs

$$\begin{aligned} \mathcal{L} \supset & \frac{\bar{c}_H}{2v^2} \partial^\mu [\Phi^\dagger \Phi] \partial_\mu [\Phi^\dagger \Phi] + \frac{g'^2 \bar{c}_\gamma}{m_W^2} \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{g_s^2 \bar{c}_g}{m_W^2} \Phi^\dagger \Phi G_{\mu\nu}^a G_a^{\mu\nu} \\ & + \frac{2ig \bar{c}_{HW}}{m_W^2} [D^\mu \Phi^\dagger T_{2k} D^\nu \Phi] W_{\mu\nu}^k + \frac{ig' \bar{c}_{HB}}{m_W^2} [D^\mu \Phi^\dagger D^\nu \Phi] B_{\mu\nu} \\ & + \frac{ig \bar{c}_W}{m_W^2} [\Phi^\dagger T_{2k} \overleftrightarrow{D}^\mu \Phi] D^\nu W_{\mu\nu}^k + \frac{ig' \bar{c}_B}{2m_W^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] \partial^\nu B_{\mu\nu} \\ & + \frac{\bar{c}_t}{v^2} y_t \Phi^\dagger \Phi \Phi^\dagger \cdot \bar{Q}_L t_R + \frac{\bar{c}_b}{v^2} y_b \Phi^\dagger \Phi \Phi \cdot \bar{Q}_L b_R + \frac{\bar{c}_\tau}{v^2} y_\tau \Phi^\dagger \Phi \Phi \cdot \bar{L}_L \tau_R. \end{aligned}$$

Global fit to **signal strengths**
and **kinematic distributions**

Conclusions of the analysis

1. Breaking of blind directions requires information on associated production (AP)
2. Kinematic distributions in AP is as sensitive (or more) than total rates

Global fit to **signal strengths**
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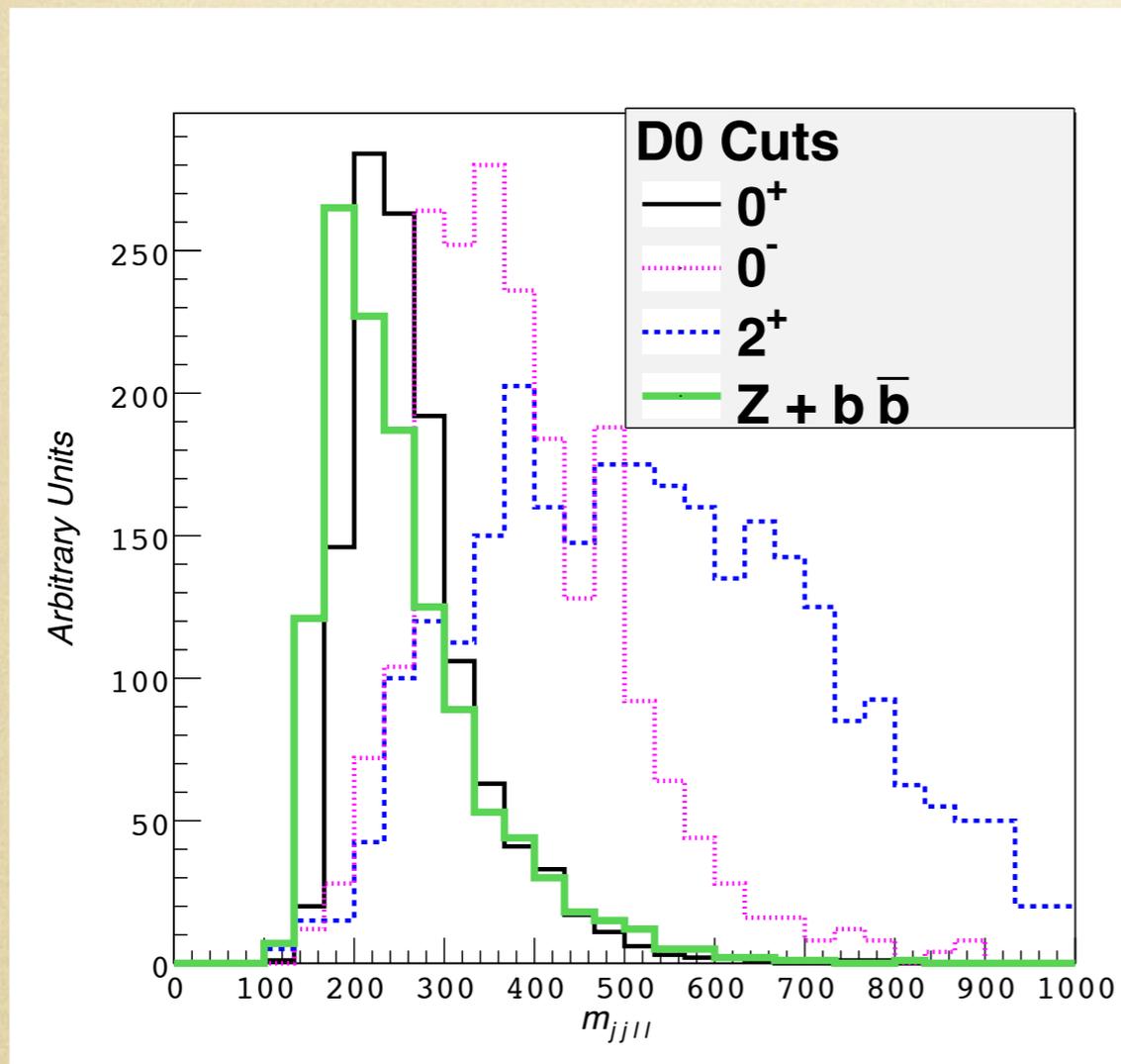
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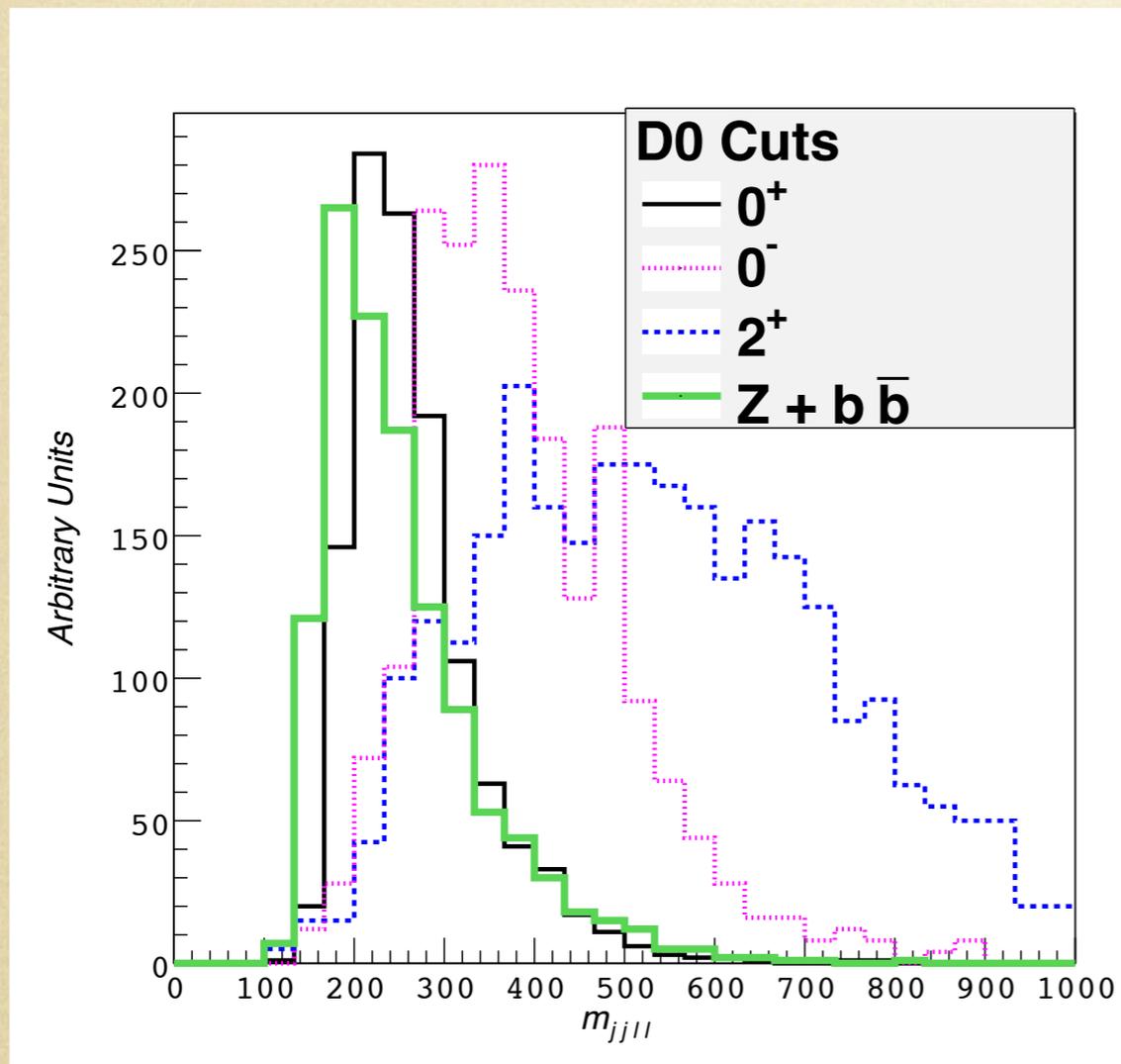
Kinematics of associated production

AP is very sensitive to the Lorentz structure of the vertex



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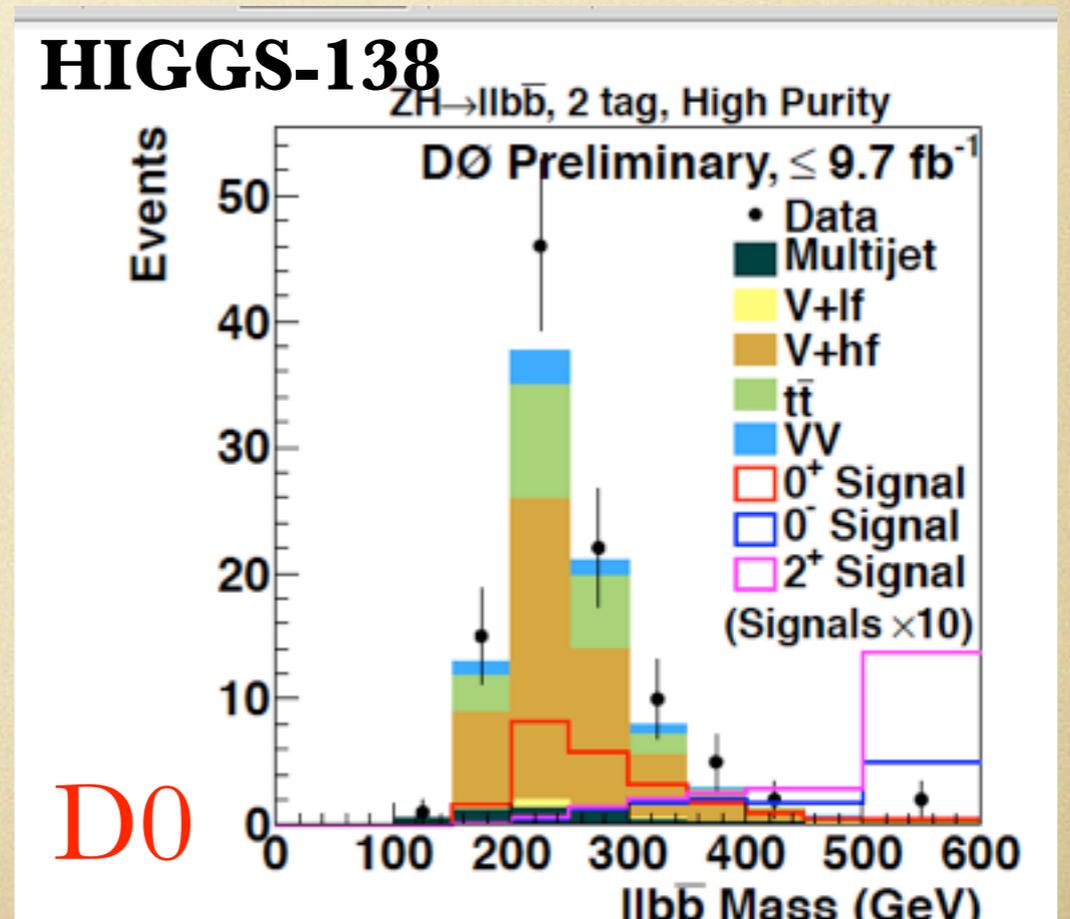


Ellis, Hwang, VS, You. 1208.6002

Test JCP of the Higgs

$$m_V h$$

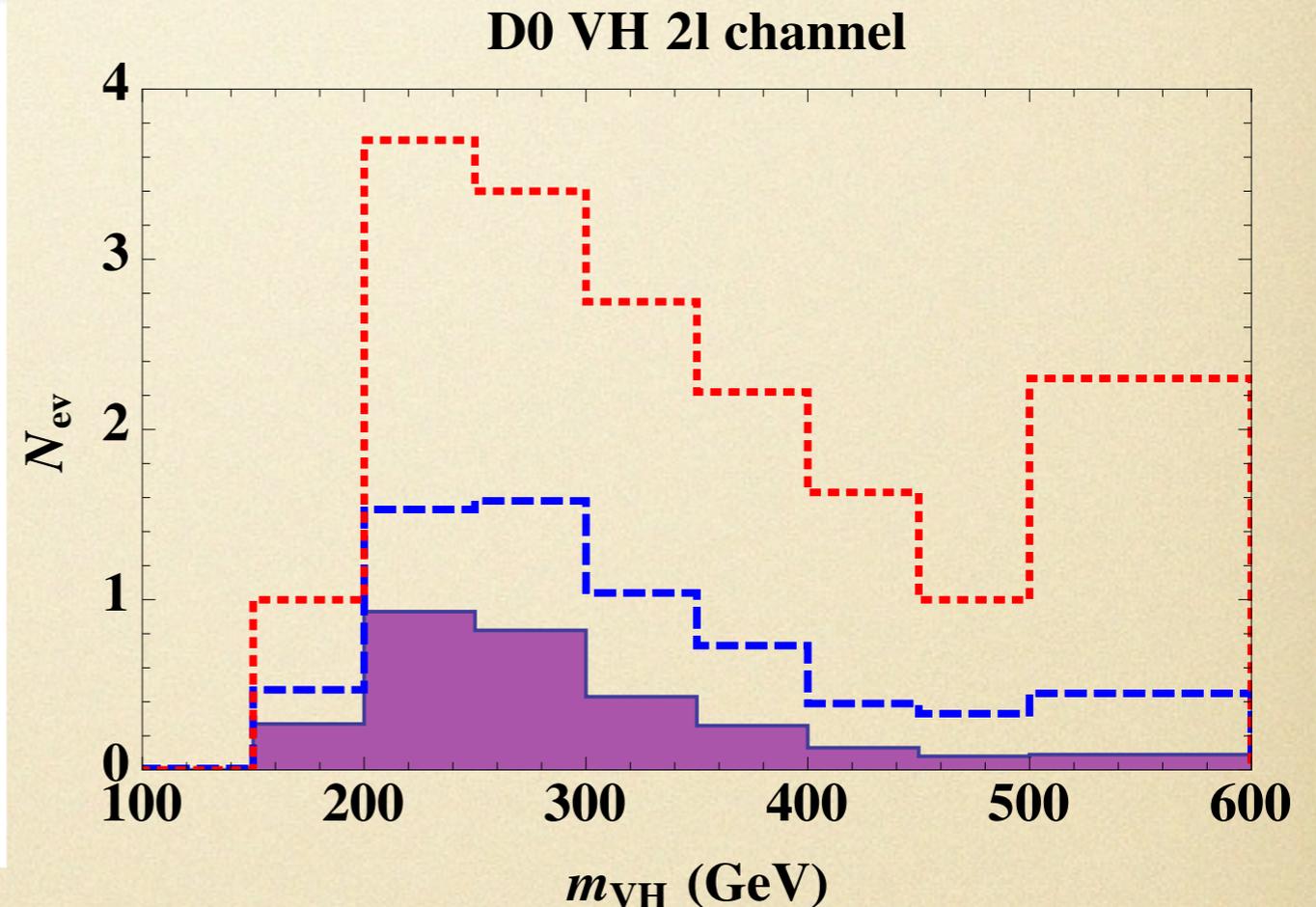
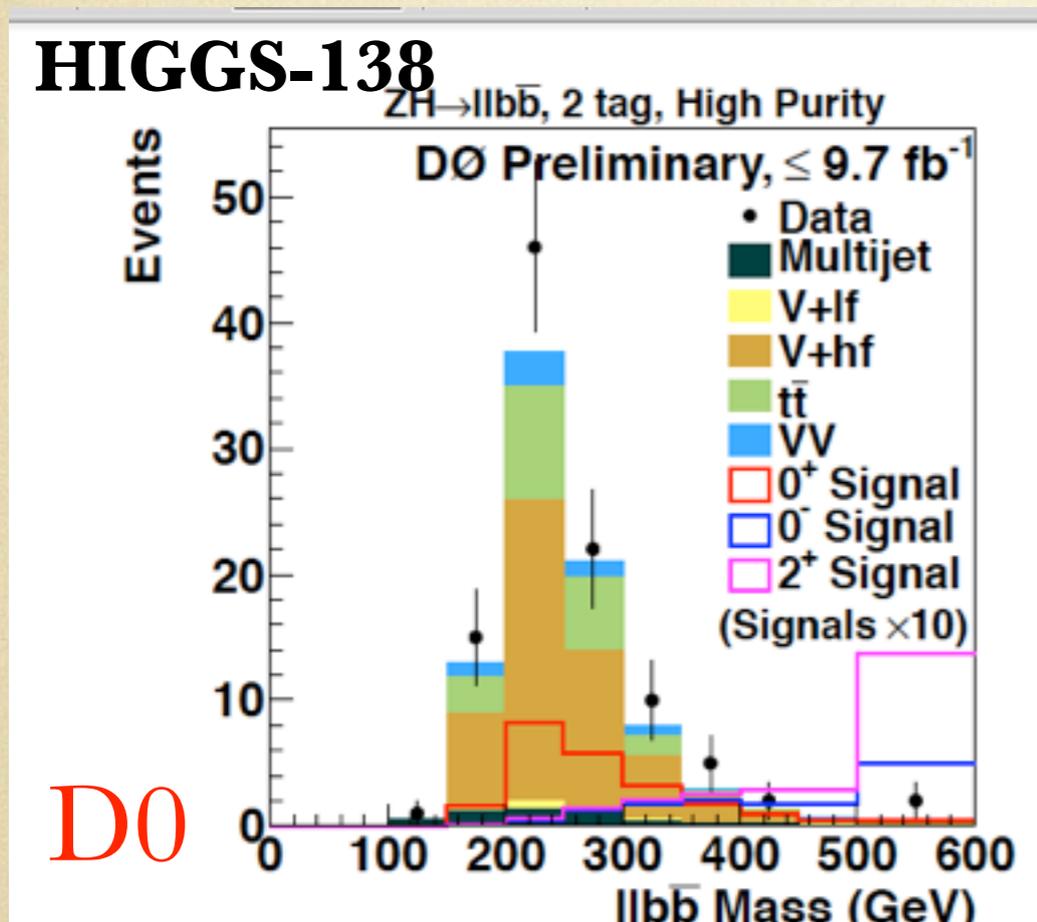
HIGGS-138



Kinematics of associated production

Higgs' HDOs also have a different Lorentz structure

HIGGS-138

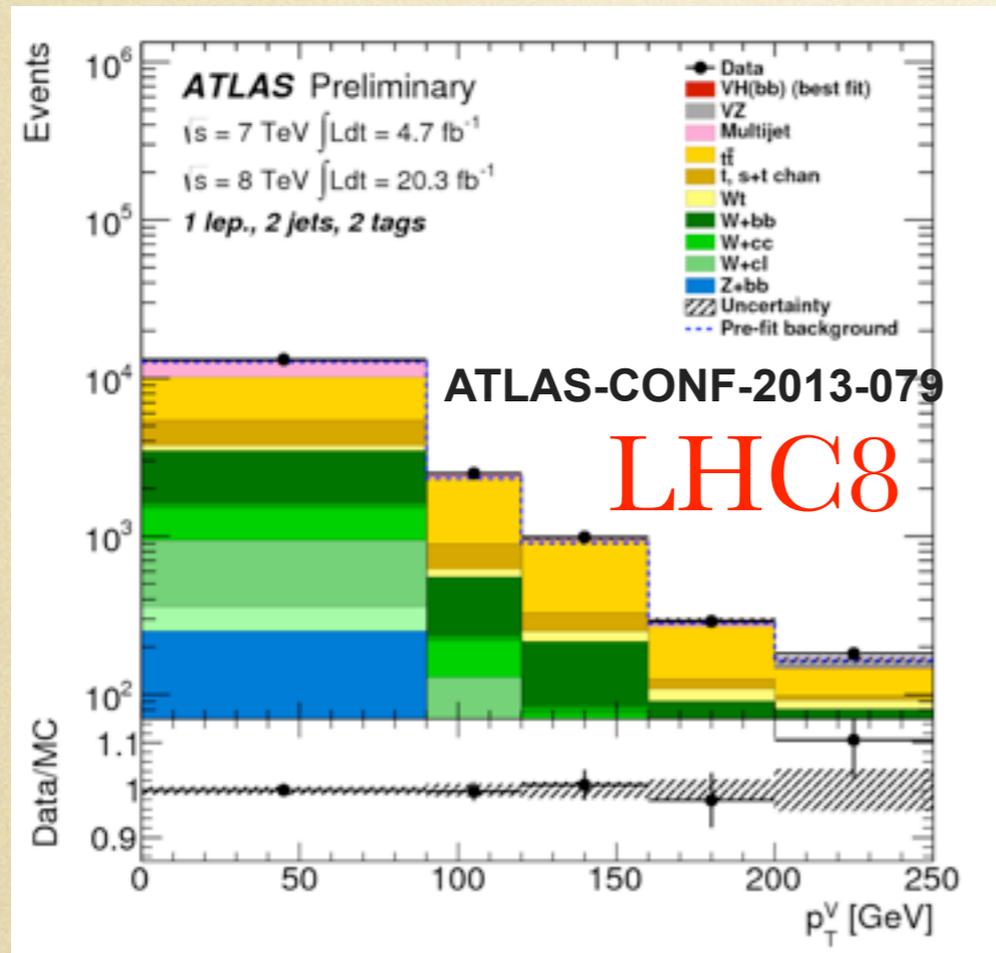


Feynrules \rightarrow MG5 \rightarrow pythia \rightarrow Delphes3
validate SM/BGs \Rightarrow expectation for HDOs

Last bin constraint better than D0 signal strength

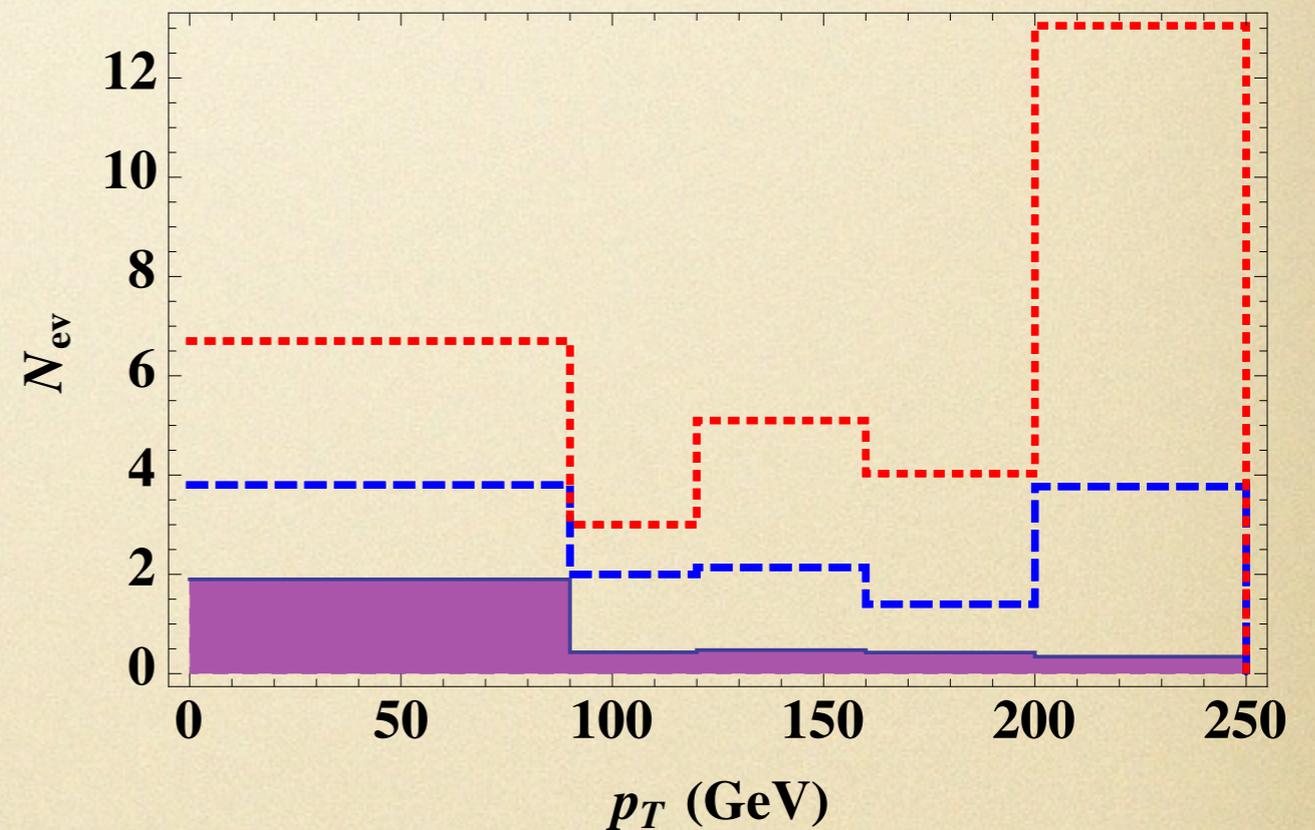
Kinematics of associated production

LHC data should be even better



only p_T^V available

LHC8 ATLAS VH



Feynrules -> MG5-> pythia->Delphes3
verified for SM/BGs => expectation for HDOs

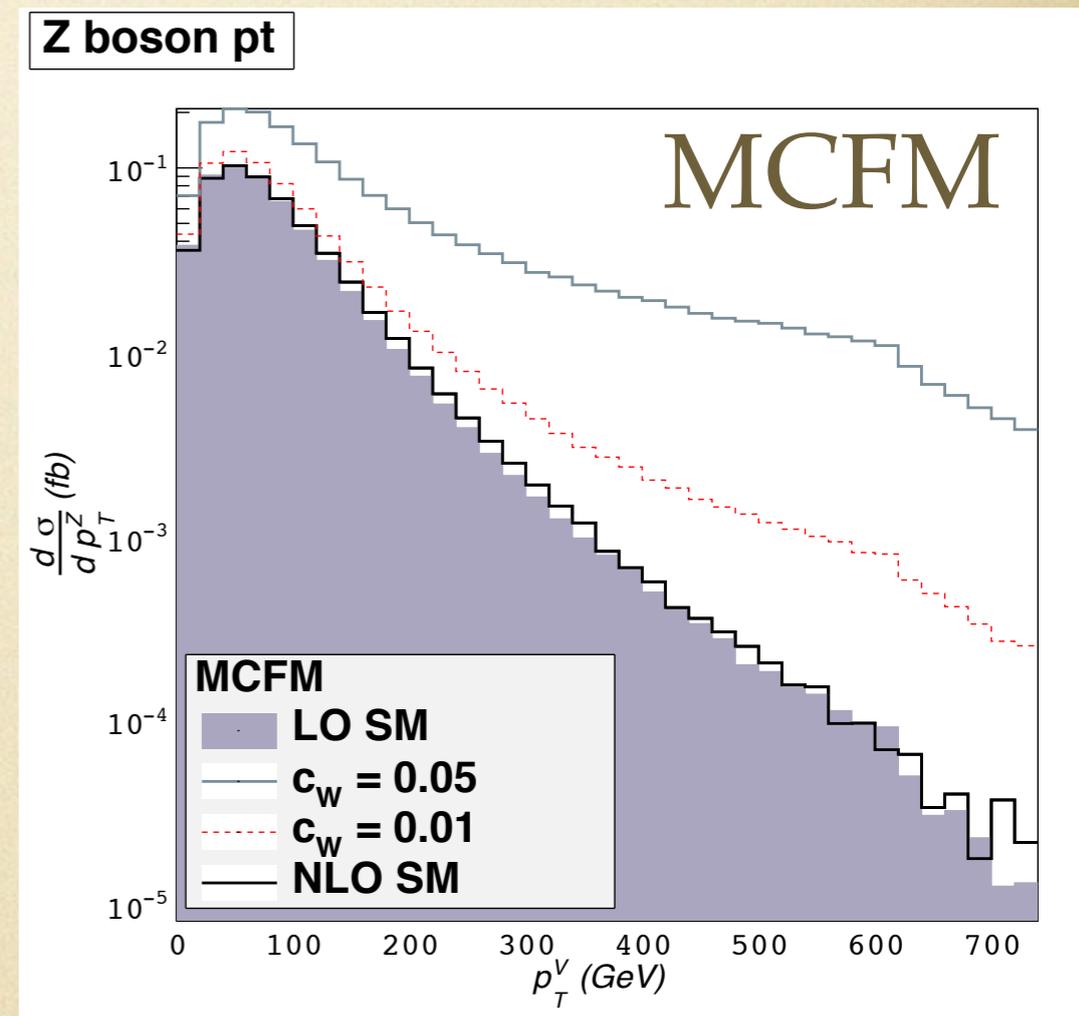
Kinematics of associated production

comment 1:

pTV is more sensitive than mVH to QCD NLO
but effect not yet at the level of operator values we can
bound

comment 2:

Sensitivity to quadratic orders
in c 's (dim-8) is less than current
errors.



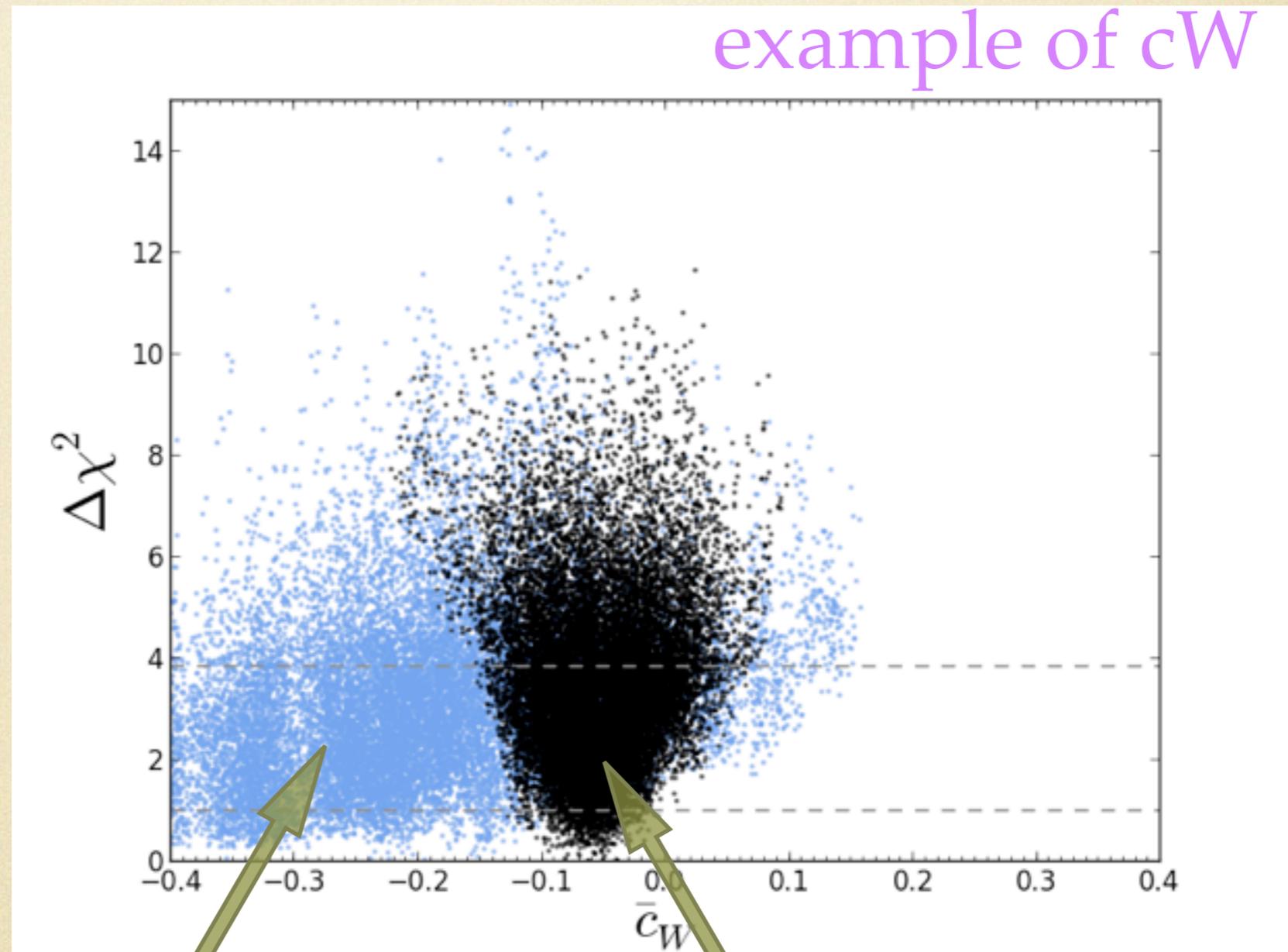
VS and Williams. In prep.

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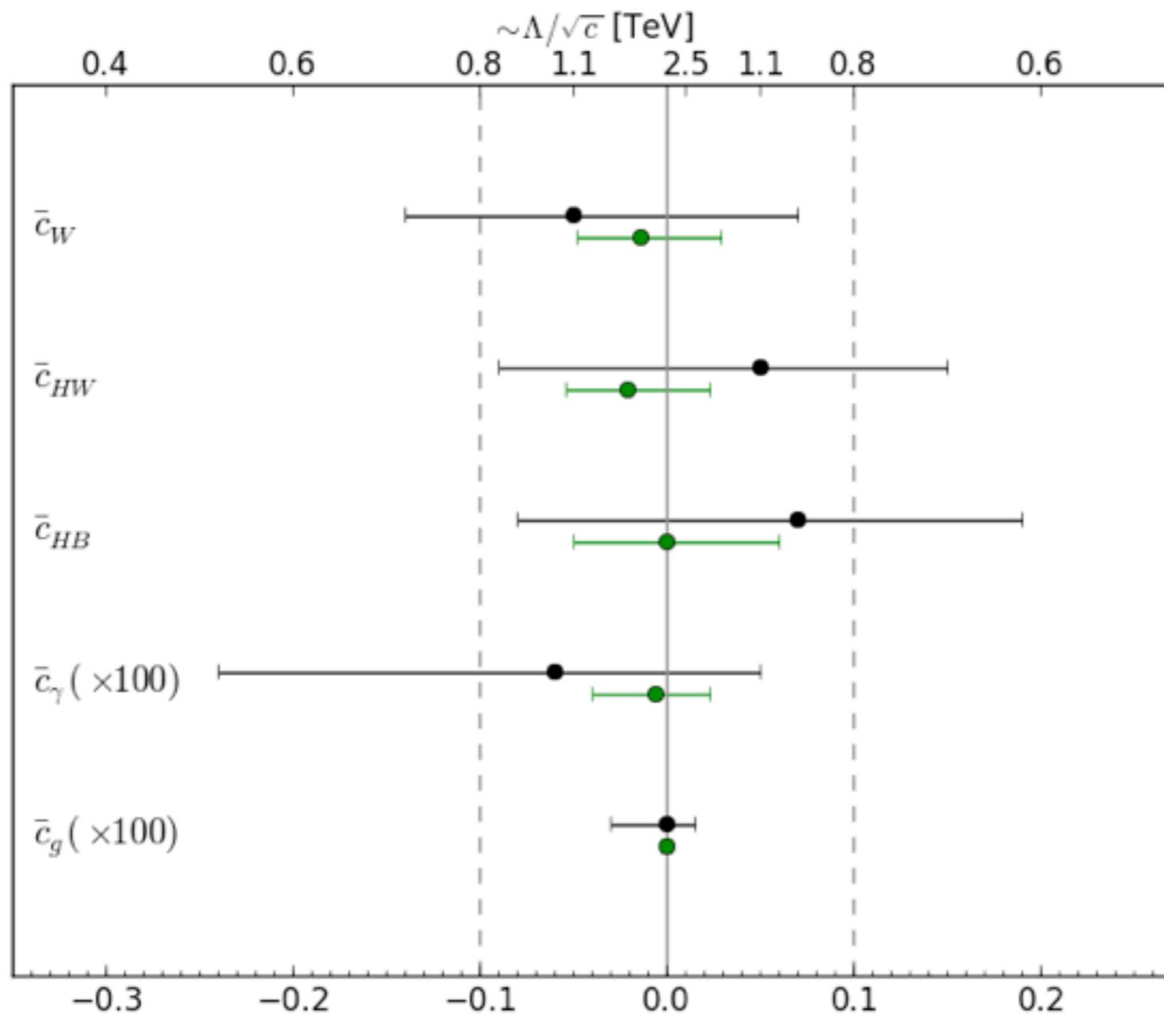
Global fit to 8 parameters



without AP

with AP

In summary



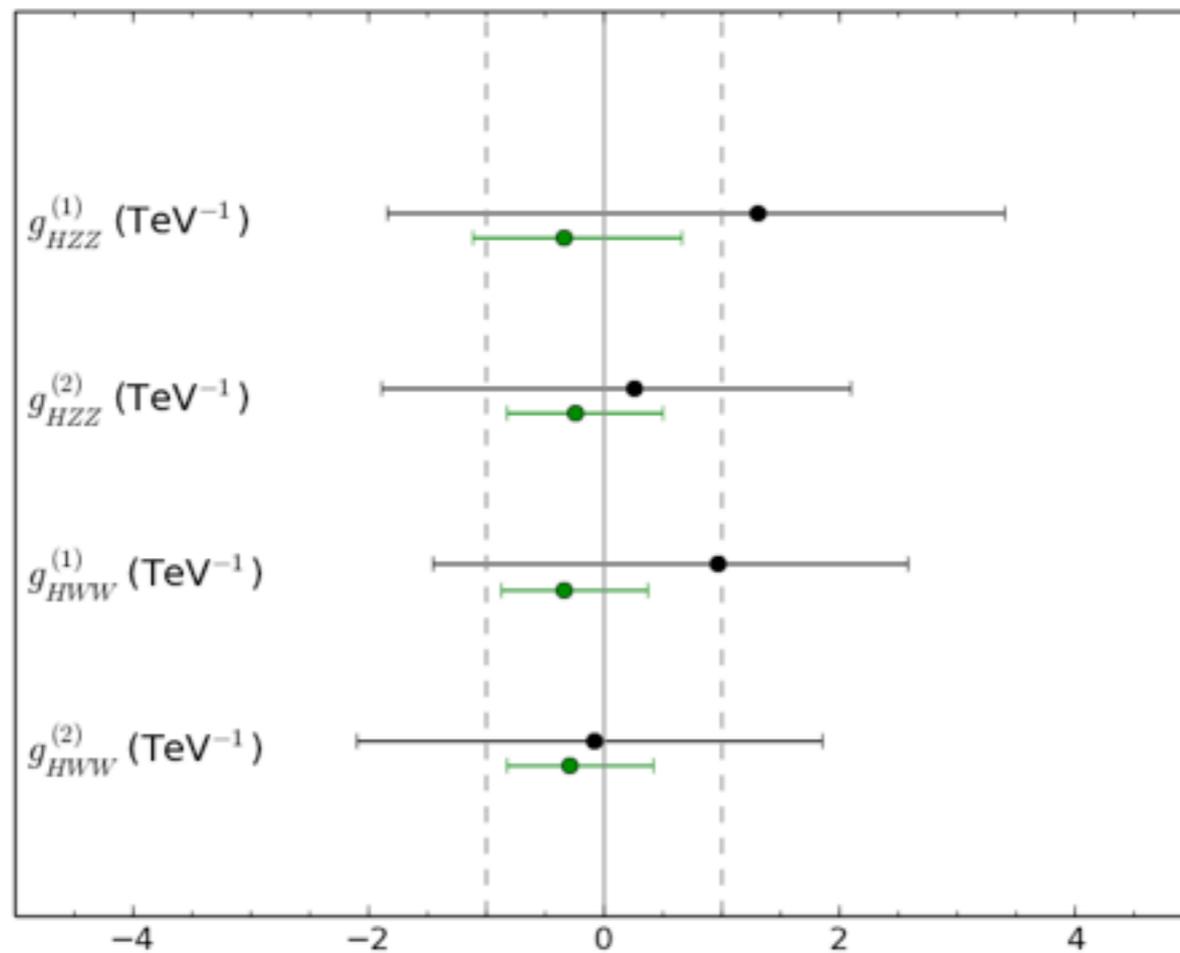
black global fit
green one-by-one fit

ct, cd, cH: no meaningful constraints

In summary

In terms of Higgs' anomalous couplings

$$\begin{aligned} \mathcal{L} \supset & -\frac{1}{4}g_{HZZ}^{(1)}Z_{\mu\nu}Z^{\mu\nu}h - g_{HZZ}^{(2)}Z_\nu\partial_\mu Z^{\mu\nu}h \\ & -\frac{1}{2}g_{HWW}^{(1)}W^{\mu\nu}W_{\mu\nu}^\dagger h - \left[g_{HWW}^{(2)}W^\nu\partial^\mu W_{\mu\nu}^\dagger h + \text{h.c.} \right], \end{aligned}$$



black global fit
green one-by-one fit

Conclusions

Absence of hints in direct searches

EFT approach to Higgs physics

Systematic, model-independent

Handles at LHC and Tevatron

beyond signal strengths

COM dependence, angular, p_T and inv mass

SM precision crucial: excess as genuine new physics

LEP constraints not as powerful as expected

Higgs physics is complementary

Complete global fit to Higgs physics

also, beyond signal strengths

Use of AP breaks degeneracy in fit, plus provides stronger constraints through kinematics

Framework for HDO studies

Feynrules HDOs involving Higgs and TGCs

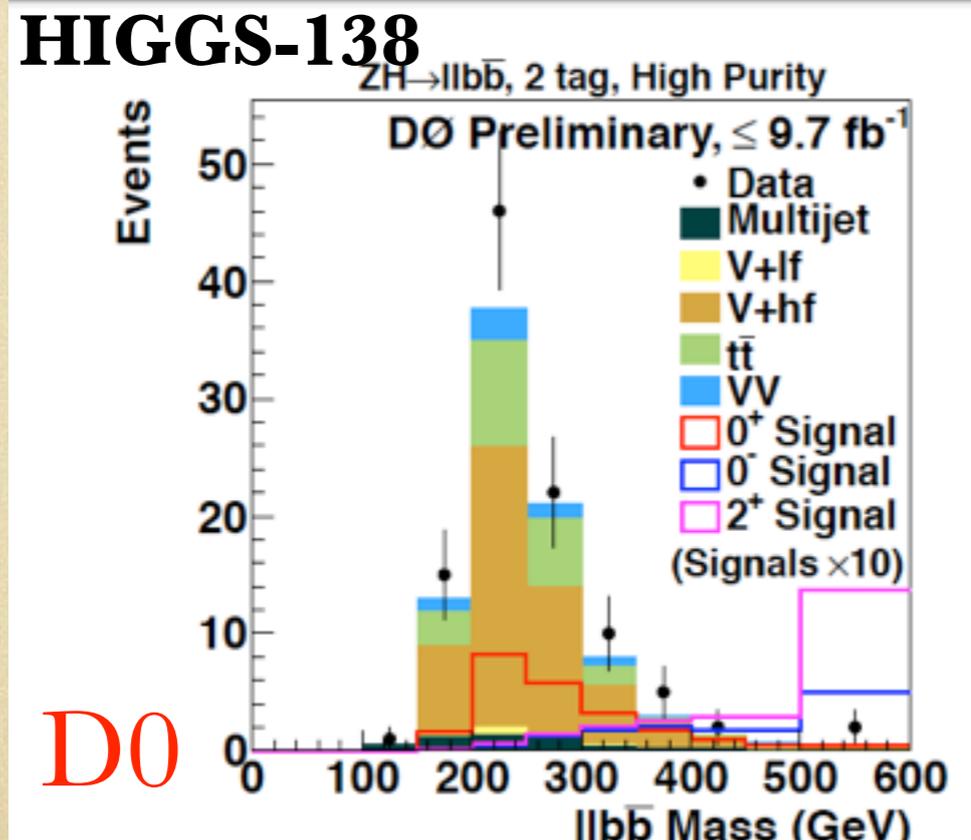
Alloul, Fuks, VS. 1310.5150

links to CalcHEP, LoopTools, Madgraph...

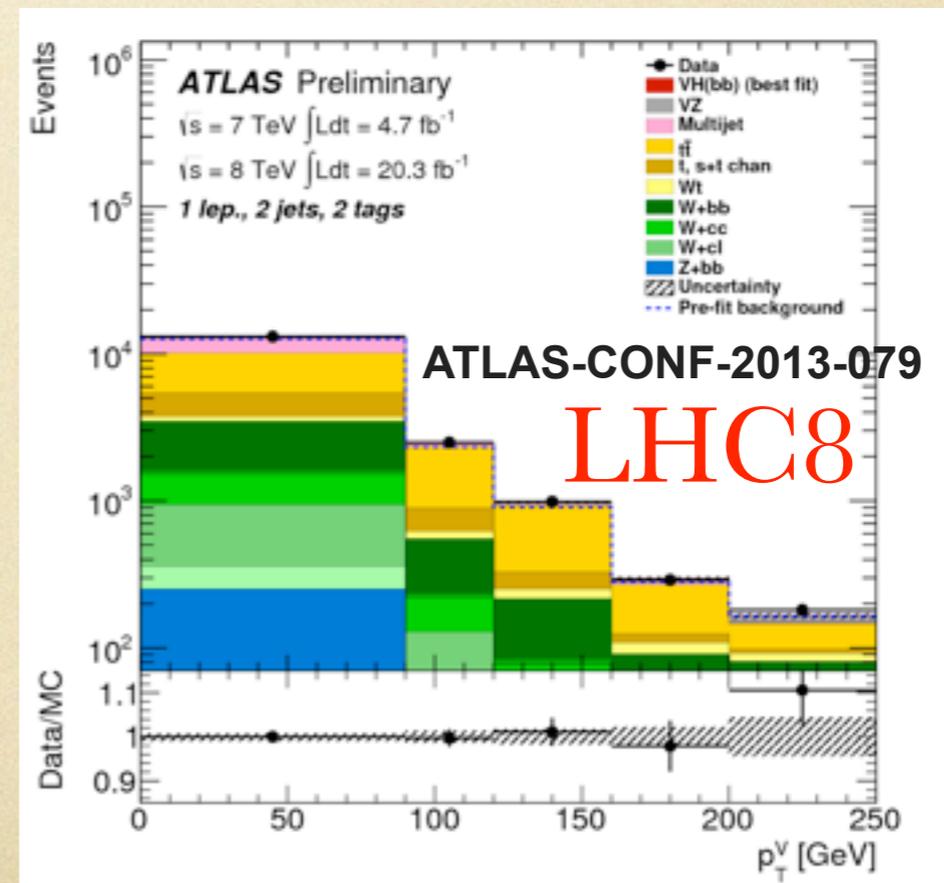
simulations: HDOs->Madgraph->Pythia... -> FastSim / FullSim

ex. Higgs in associated production

HIGGS-138



m_{Vh}



p_T^V

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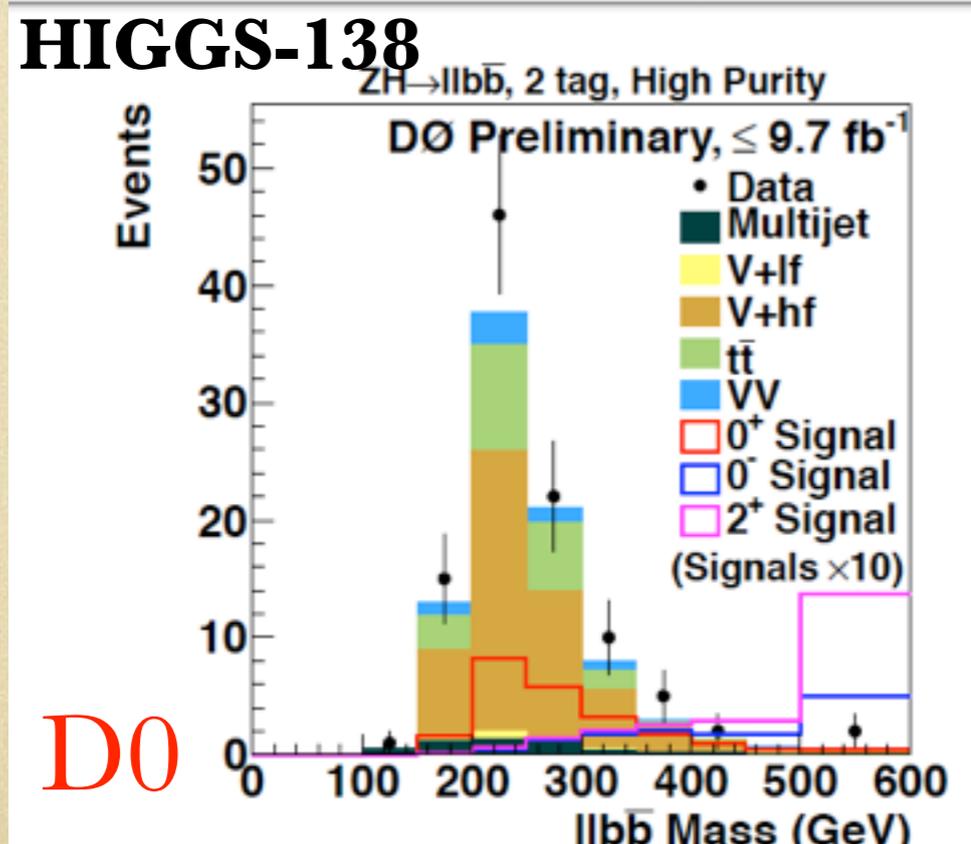
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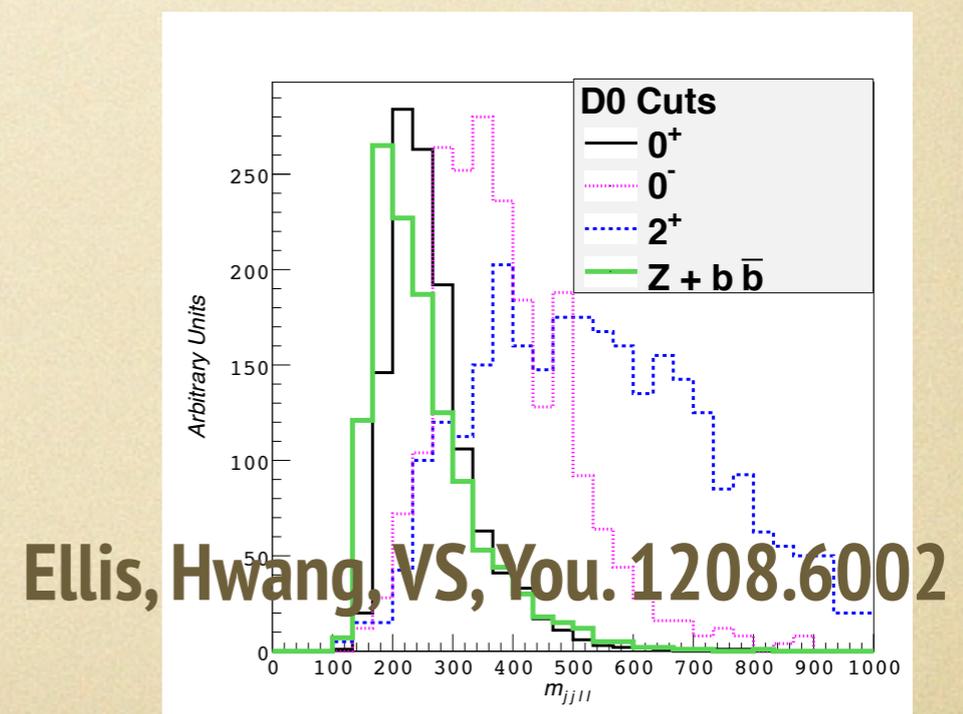
HIGGS-138



D0

m_{Vh}

Test other JCP



Ellis, Hwang, VS, You. 1208.6002