Higgs portal dark matter

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and invisible width

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Higgs Portal

- Second simplest model of WIMP dark matter
- Rich phenomenology: direct and indirect detection as well as collider signals possible
- Motivated: |H|^2 is lowest dimension SM singlet, so SM singlet dark matter may naturally couple to SM via this operator
- Oark matter could be scalar, fermion, or vector
- Simplest scalar case:

2 parameter model:

$$\Delta {\cal L}_{
m Scalar} = - \, {\lambda \over 4} S^2 H^\dagger H^{
m Burgess, \, Pospelov, \, ter \, Veldhuis, hep-ph/0011335}$$

phenomenology fully determined by mass m_{s} and coupling λ

Higgs Portal: scalar case

Dark matter annihilation proceeds via s-channel Higgs

Solution For given (m_s , λ) annihilation cross section is fixed

 Insisting on thermal abundance of dark matter fixes λ(ms)

For weak scale dark matter and order 0.1-1 coupling dark matter relic abundance in right ballpark



Scalar DM coupled to the Higgs



de Simone, Giudice, Strumia 1402.6287

Higgs Portal: scalar case



For thermal abundance low mass dark matter region, ms < mh/2 excluded

Actually, doubly excluded:
 by direct detection constraints
 and by Higgs invisible width





de Simone, Giudice, Strumia 1402.6287 Higgs Portal: constraints from direct detection $c_1^N S^2 ar{N} N$

 Higgs exchange induces effective coupling of dark matter to nucleons

For given (ms, λ)
 scattering cross section
 on nucleons fixed

For λ(ms) fixed by thermal relic abundance, direct detection safely excluded for ms >7 GeV

Jessie Shelton, private com







Higgs Portal: constraints from direct detection

For ms < mh/2 Higgs can decay directly to dark matter, leading to invisible width

Branching fraction is huge for thermal region of parameter space

$$egin{split} \Gamma(h o ext{invisible}) &= rac{\lambda^2 v^2}{64\pi m_h} \sqrt{1 - 4m_S^2/m_h^2} \ rac{\Gamma(h o ext{invisible})}{\Gamma_h^{ ext{SM}}} pprox \left(rac{\lambda}{0.04}
ight)^2 \end{split}$$

Jessie Shelton, private com





Limits on exotic Higgs branching fraction

Assuming Higgs couplings to SM fixed



Br(h→exotic) $\lesssim 18\%$ at 95% CL

Limits on exotic Higgs branching fraction

Allowing some Higgs couplings to SM to float



Br(h→exotic) \lesssim 30% at 95% CL

Limits on exotic Higgs branching fraction

Compare direct and indirect width constraints



Limits on invisible Higgs branching fraction

Assuming Higgs couplings to SM fixed



Limits on invisible Higgs branching fraction Assuming Higgs coupling to gluons floating

Effect of monojet constraints recast from CMS and ATLAS monojets searches by Djouadi et al. 1205.3169



Scalar Higgs portal summary

- Barring biggest conspiracy since Roswell
 Br(h→invisible) ≤ 30% (and most likely ≤ 20%)
- Direct limits from CMS and ATLAS yield
 Br(h→invisible) ≤ 58% for SM Higgs coupling to V
- So For ms < mh/2 this translates to λ ≥ 0.03 (0.02) and excludes couplings corresponding to thermal WIMP
- ${\ensuremath{ \circ }}$ For 10 GeV < ms comparable constraints on λ from direct detection



Light Higgs portal DM should be non-thermal. Then probed via (small) invisible Higgs width at LHC or in direct detection experiments

Higgs portal DM exactly at resonance (ms=mh/2) or heavier than 100 GeV can be thermal. Then probed in direct detection experiments. What about LHC?

Fermion and vector Higgs portal

- Fermionic or vector dark matter can also couple via Higgs portal, though in this case 0 portal interaction is non-renormalizable
- For light dark matter conclusions similar as in scalar case... 0
- For heavy dark matter thermal cross section $\Delta \mathcal{L}_{ ext{Fermion}}$ 0 excluded by direct detection except for purely imaginary y where direct detection cross section is spin dependent and velocity suppressed



Lopez-Honorez, Schwetz, Zupan, arXiv:1203.2064. Djouadi, Lebedev, Mambrini, Quevillon, arXiv:1112.3299.

 $\psi \psi H^{\dagger} H$

+ h.c.

$$\Delta \mathcal{L}_{\rm Vector} = -\,\lambda_V V_\mu V_\mu H^\dagger H$$

thermal

 10^{3}

2v

Fermion and vector Higgs portal

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 Lopez-Honorez, Schwetz, Zup
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- For heavy dark matter thermal cross section excluded by direct detection except for purely imaginary y where direct detection cross section is spin dependent and velocity CMS, 1404.1344 ATLAS, 1402.3244 $\Delta \mathcal{L}$ Vector



Lopez-Honorez, Schwetz, Zupan, arXiv:1203.2064. Djouadi, Lebedev, Mambrini, Quevillon, arXiv:1112.3299.

 $-\psi\psi H^{\dagger}H$

$$\Delta \mathcal{L}_{\text{Vector}} = -\lambda_V V_\mu V_\mu H^\dagger H$$

2v



DM-nucleon cross section a SI [pb]

+ h.c.

Take-away

Higgs portal DM is a very motivated and very predictive scenario for dark matter.

- Predicts direct and indirect detection signals as well as collider signals
- For ms < 60 GeV region of parameter space corresponding to thermal relic abundance is safely excluded by both LHC and direct detection experiments
- Heavier region will be probed by direct detection in near future
- Can we probe heavier Higgs portal DM at LHC? Or at ILC/TLEP? Or at 100 TeV