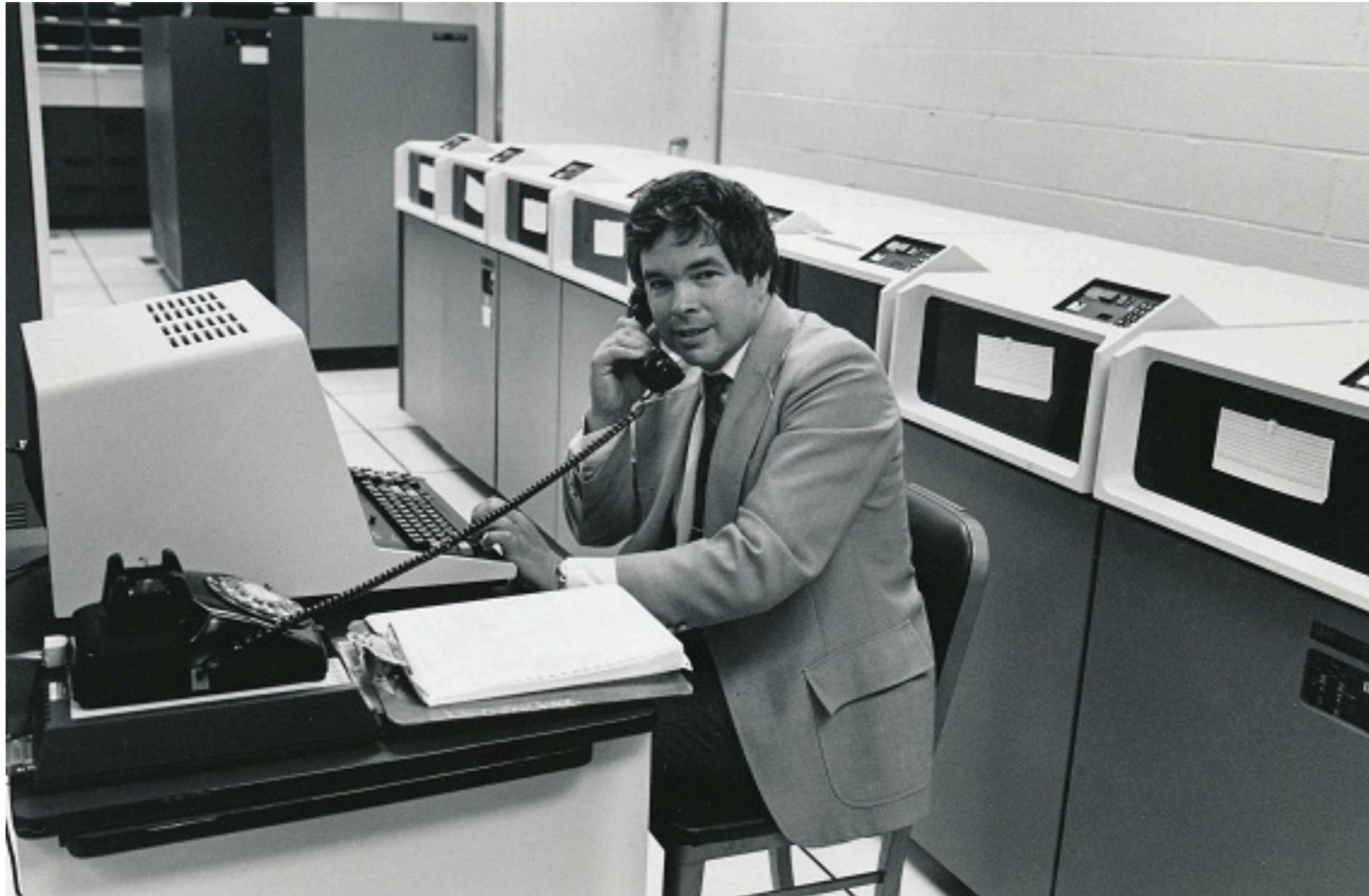


Recent Progress in Strongly Coupled CFT

Slava Rychkov

Ken Wilson



Unsolved problems - I

- where Laws of Nature are not known/hypothetical
 - ▶ high T_c superconductivity
 - ▶ quantum gravity
 - ▶ Beyond SM...

Unsolved problems - 2

- where Laws of Nature are relatively certain, but extracting predictions remains hard*
 - ▶ turbulence
 - ▶ strongly coupled QFT (QCD + many condensed matter examples)

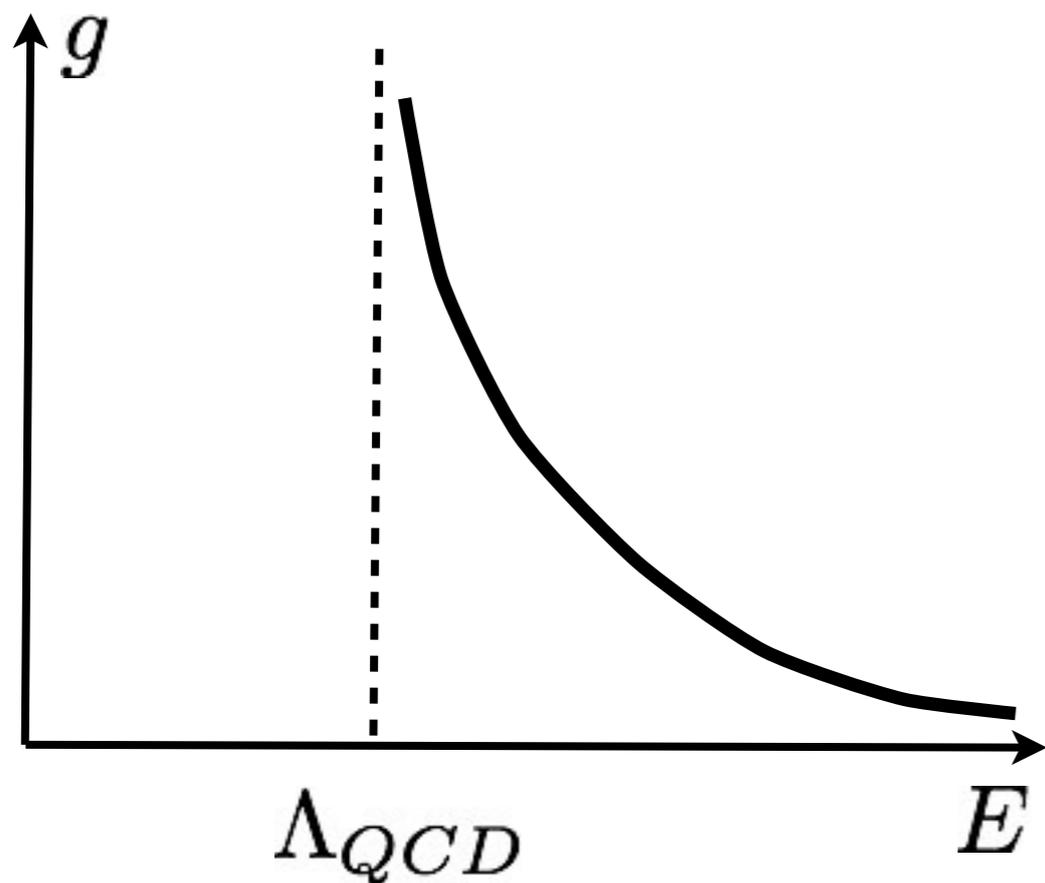
Computational complexity:
can we find better algorithms?

*Hard = nontrivial even with supercomputers

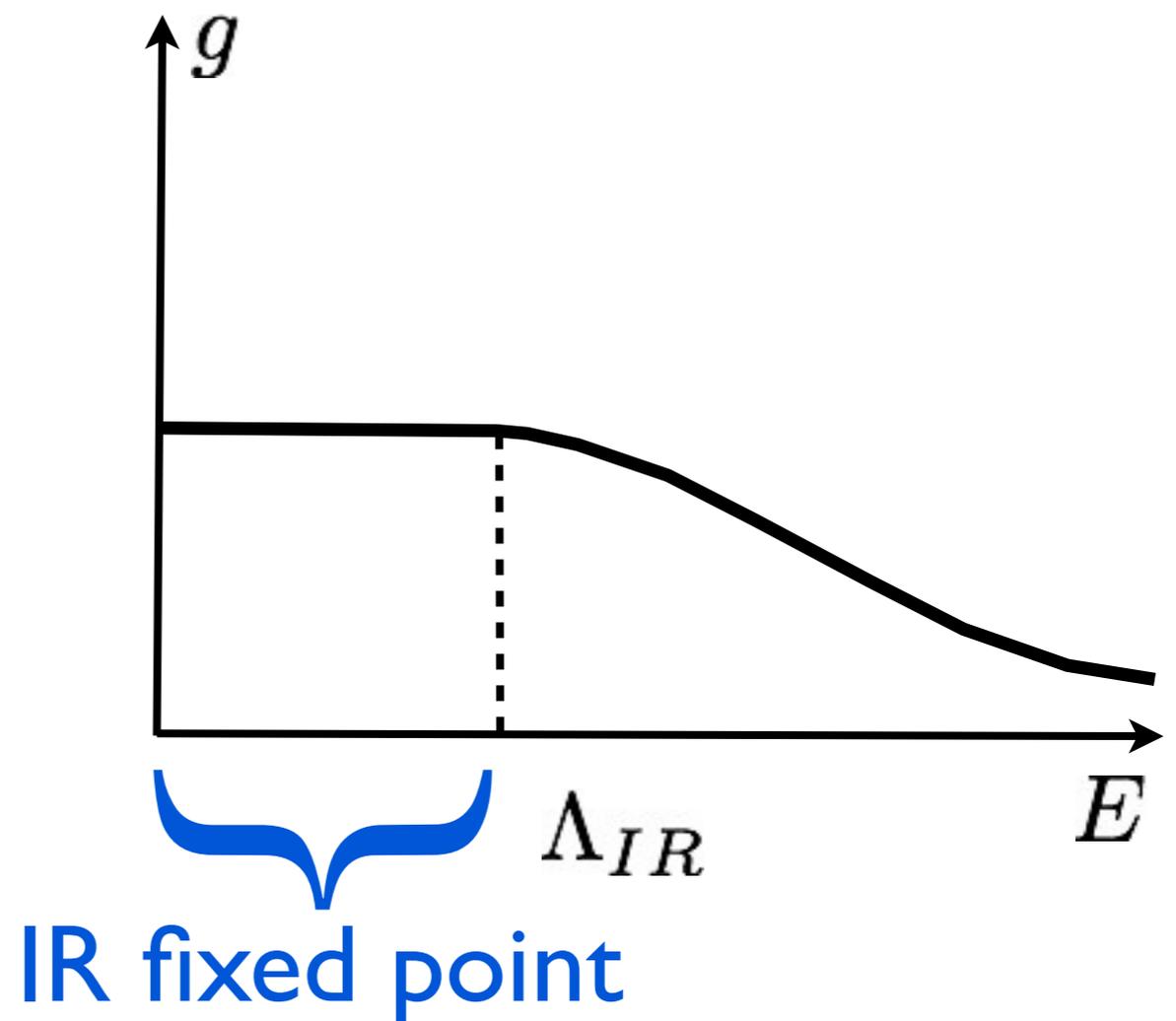
Conformal Field Theories

Running coupling in QCD:

$$N_c = 3, \quad N_f = 3$$



$$N_c = 3, \quad 12 \lesssim N_f \leq 16$$



IR fixed points are

- **scale invariant** (by definition)
- **conformally invariant** (generically, and perhaps always in unitary theories)

Condition for conformal invariance: $T_{\mu}^{\mu} = 0$ (1)

Condition for scale invariance: $T_{\mu}^{\mu} = \partial_{\lambda} V^{\lambda}$ (2)

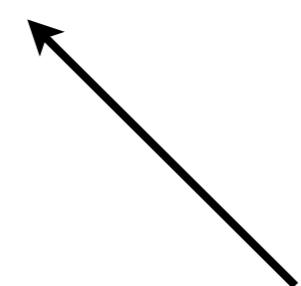
Generically $(2) \Rightarrow (1)$

Scale invariance

$$x \rightarrow \lambda x$$

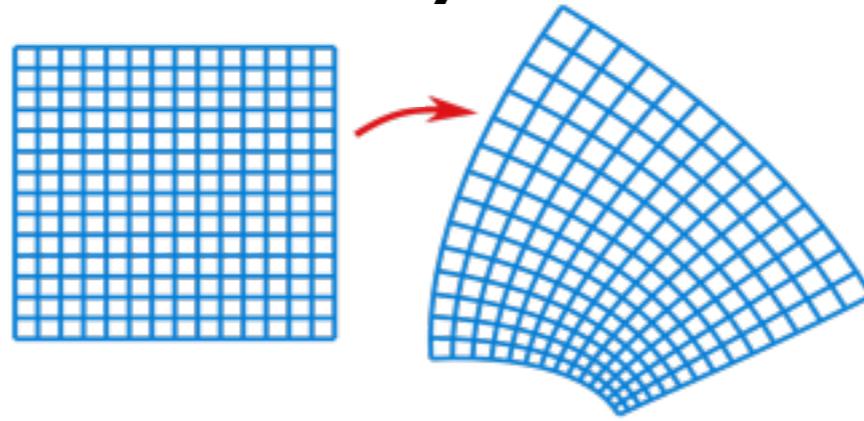
$$\langle \mathcal{O}(x)\mathcal{O}(y) \rangle \propto \frac{1}{|x-y|^{2\Delta_{\mathcal{O}}}}$$

scaling dimension



Conformal invariance

$x \rightarrow x' = f(x)$ s.t. locally rotation + dilatation



Physically...

Theorem [Liouville, 1850]

In $d \geq 3$ the group of conformal transformations is

- finite-dimensional
- isomorphic to $SO(d+1, 1)$
- generated by $M_{\mu\nu}$, P_μ , D , and K_μ

$$K_\mu = I \circ P_\mu \circ I = 2x_\mu(x\partial) - x^2\partial_\mu$$

special conformal transformation

What does it buy you?

For starters, also 3-point functions are fixed:

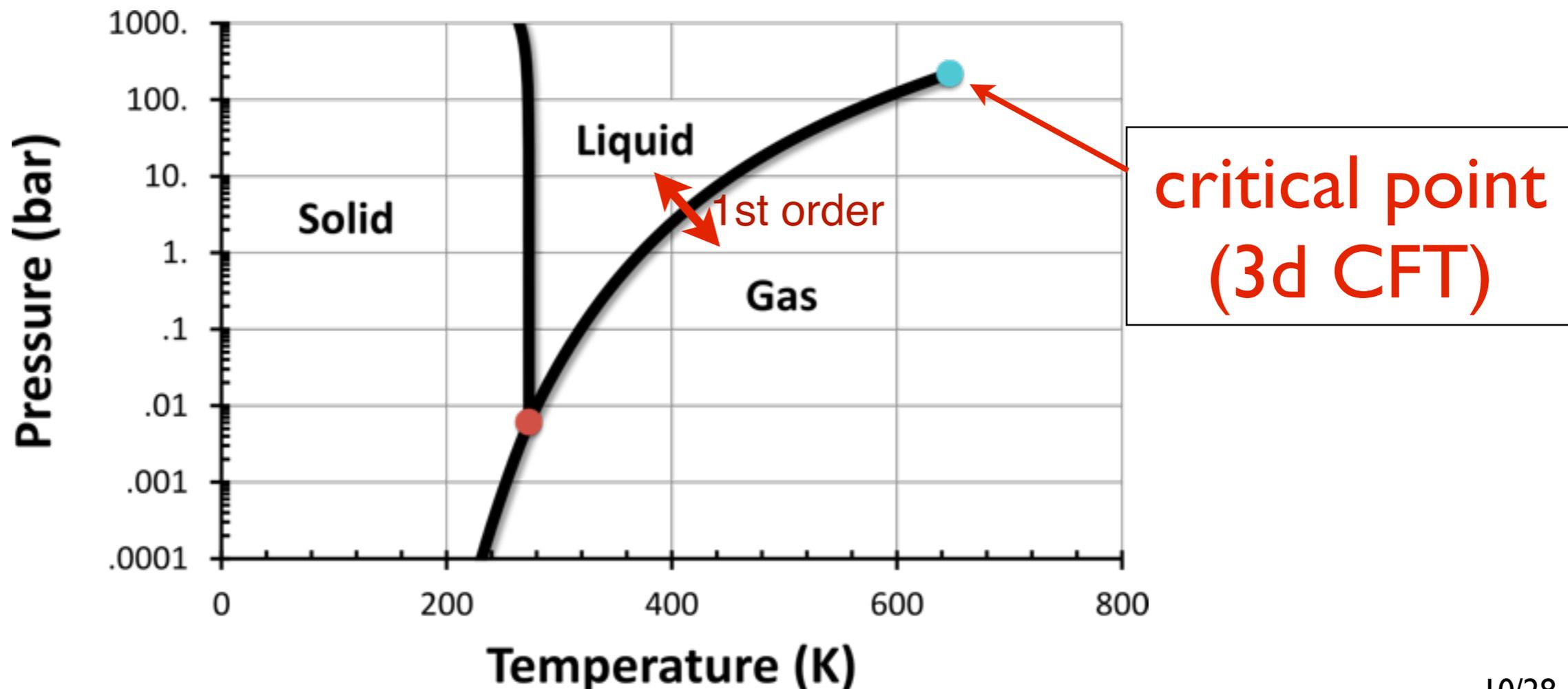
$$\langle \mathcal{O}(x)\mathcal{O}(y)\mathcal{O}(z) \rangle \propto \frac{1}{|x-y|^{\Delta_{\mathcal{O}}}|x-z|^{\Delta_{\mathcal{O}}}|y-z|^{\Delta_{\mathcal{O}}}}$$

[Polyakov 1971]

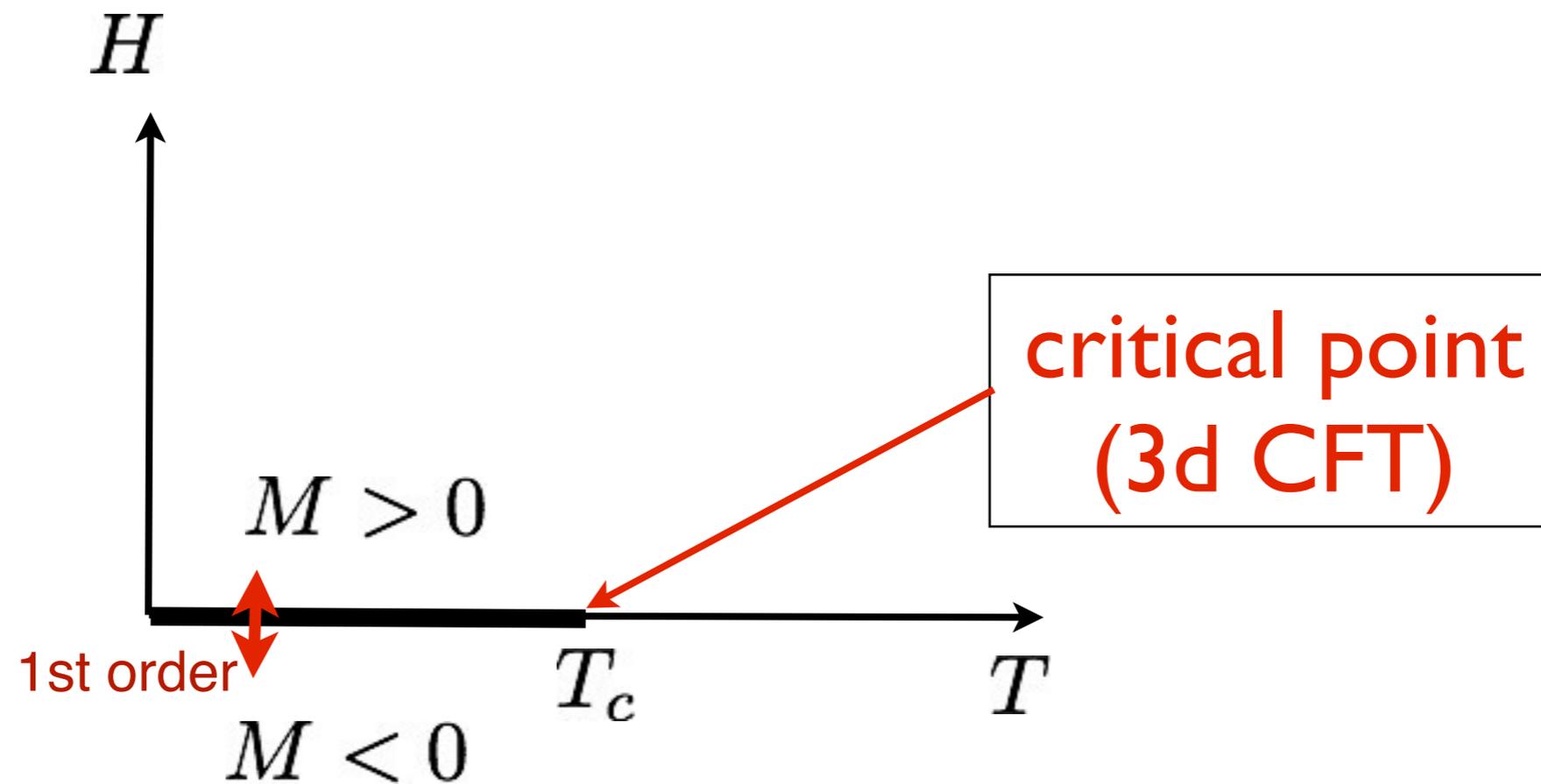
CFTs in the real world

Rich source: condensed matter and statistical physics
(2nd order phase transitions)

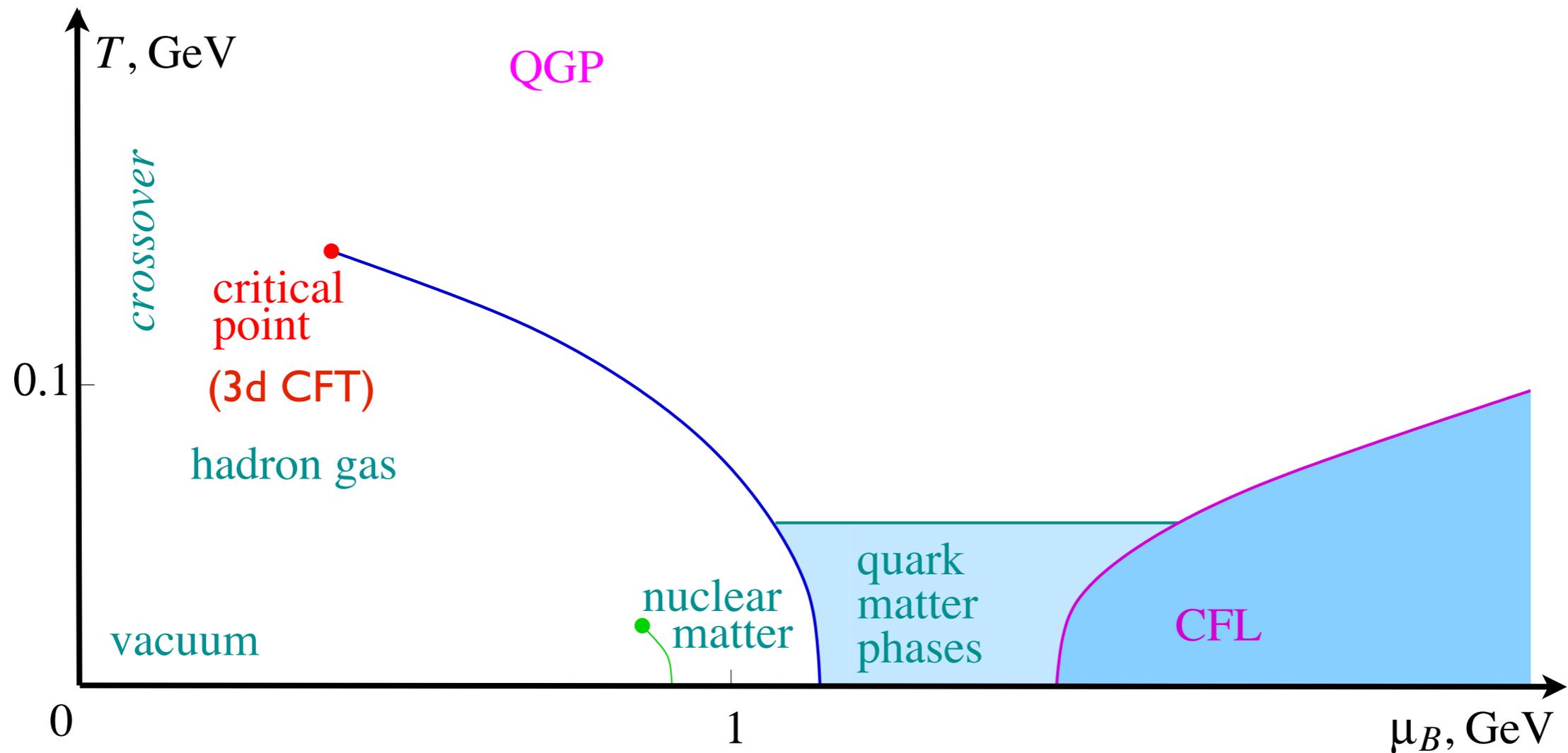
Example A: **liquid-vapor transitions**



Example B: uniaxial ferromagnets



Example C: QCD matter



from a talk by Stephanov'2004

Universality

Critical points fall into large
“universality classes”, described by the same CFT.

[Justification: RG theory of phase transitions]

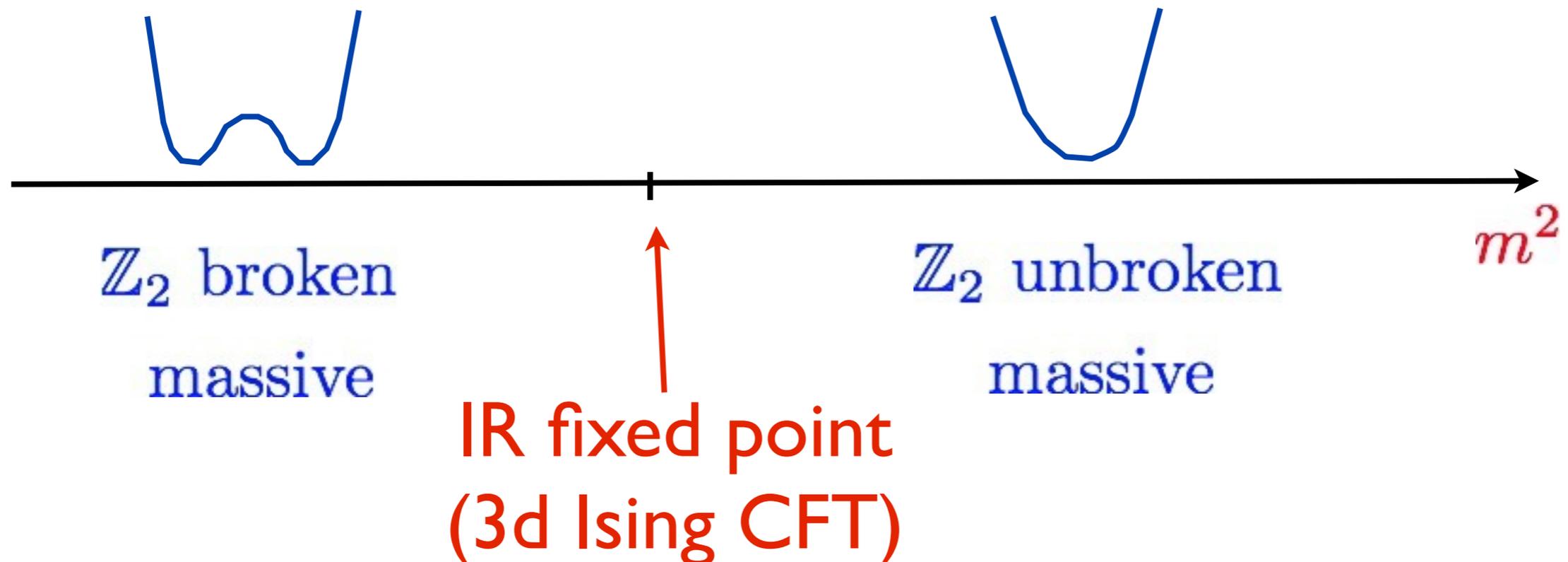
Critical points A,B,C are all described by the same CFT
(have the same operator spectrum and correlation functions)

“3d Ising model CFT”

3d Ising CFT: RG definition

$$S = \int d^d x \left\{ \frac{1}{2} (\partial\phi)^2 + \frac{m^2}{2} \phi^2 + \lambda \phi^4 \right\}$$

3d IR physics is different from 4d,
because $[\phi]=1/2$ and λ is relevant, $[\lambda]=1$



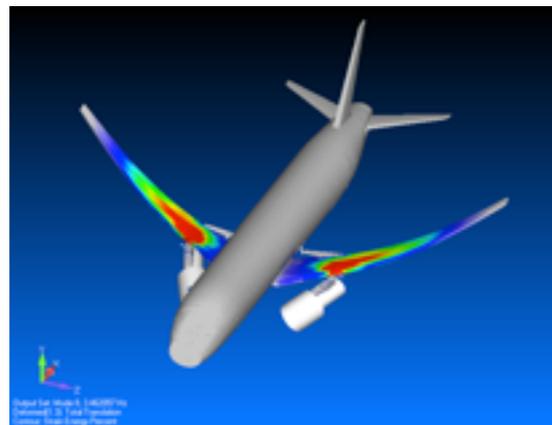
What do we need to find-1

- Scaling dimensions of operators Δ_i normalization

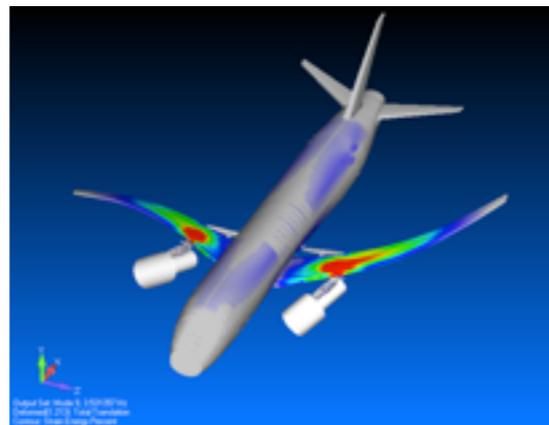
$$\langle \mathcal{O}_i(x) \mathcal{O}_j(y) \rangle = \frac{\delta_{ij}}{|x - y|^{2\Delta_i}}$$

there are ∞ many operators of each spin

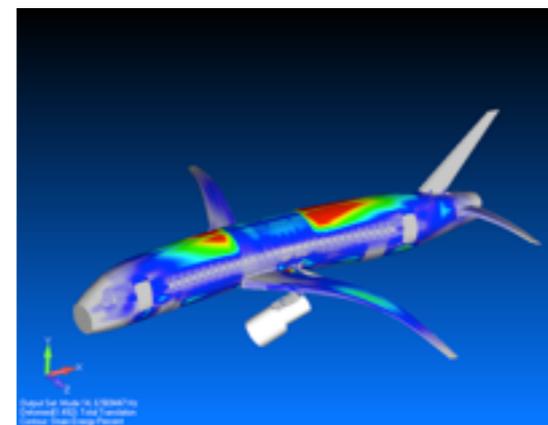
analogy: eigenmodes in mechanics:



ω_1



ω_2



ω_3

What do we need to find-2

- 3-point function couplings f_{ijk} :

$$\langle \mathcal{O}_1(x) \mathcal{O}_2(y) \mathcal{O}_3(z) \rangle = \frac{f_{123}}{|x-y|^{h_{123}} |x-z|^{h_{132}} |y-z|^{h_{231}}}$$

$$h_{ijk} = \Delta_i + \Delta_j - \Delta_k$$

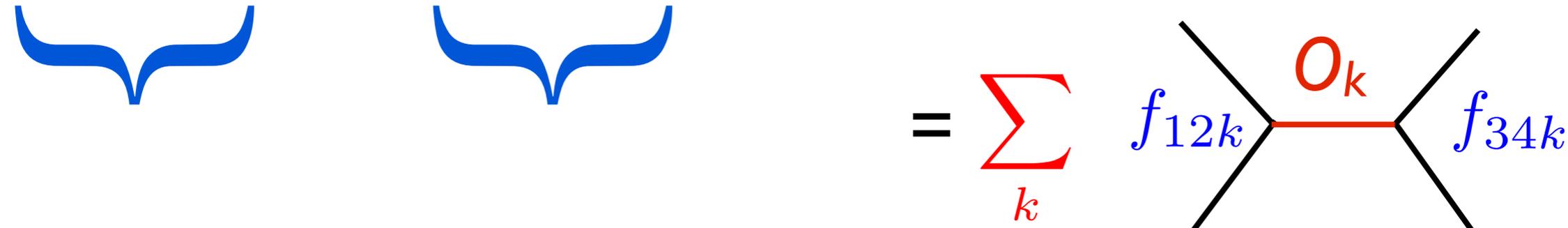
Unlike in engineering, eigenmodes interact nonlinearly

Δ_i 's and f_{ijk} 's determine everything

This follows from Operator Product Expansion:

$$O_i(x) \times O_j(0) = \sum \frac{f_{ijk}}{|x|^{\Delta_i + \Delta_j - \Delta_k}} O_k(0) + \text{derivatives}$$

In CFT's, OPE has finite radius of convergence and can be used to reconstruct higher point functions, e.g.:

$$\langle O_1(x_1) O_2(x_2) O_3(x_3) O_4(x_4) \rangle = \sum f_{12k} f_{34k} \times (\text{2-point fns})$$


The diagram shows a sum over k of a diagram with four external legs. The top-left and bottom-right legs are labeled f_{12k} and f_{34k} respectively. The top-right and bottom-left legs are labeled O_k . The internal line connecting the two vertices is a red horizontal line labeled O_k .

Bootstrap idea

Ferrara, Gatto, Grillo 1973; Polyakov 1974

To fix Δ_i 's and f_{ijk} 's impose that 4pt functions be unique:

$$\langle O_1(x_1)O_2(x_2)O_3(x_3)O_4(x_4) \rangle$$

The diagrammatic equation shows two ways to sum over an internal operator O_k to form a 4-point function. On the left, a red summation symbol \sum_k is followed by a diagram with a central red horizontal line labeled O_k . Two black lines enter from the left, labeled f_{12k} , and two black lines exit to the right, labeled f_{34k} . On the right, another red summation symbol \sum_k is followed by a diagram with a central red vertical line labeled O_k . Two black lines enter from the top, labeled f_{14k} , and two black lines exit from the bottom, labeled f_{23k} . An equals sign is placed between the two diagrams.

Was shown to work in $d=2$ (Belavin, Polyakov, Zamolodchikov 1984)

Work in $d \geq 3$ restarted only recently

(Rattazzi, S.R., Tonni, Vichi 2008)

$d=2$

vs

$d \geq 3$

- * conformal algebra is infinite dimensional (Virasoro)

- * there are interesting CFTs (e.g. 2d Ising model) with a **finite** # of primary operators of **known** dimension

- * bootstrap needs to fix f_{ijk} 's only: a problem of **fin-dim** linear algebra

- * conformal algebra is finite dimensional, $SO(d+1,1)$

- * any CFT contains **infinitely** many primaries; their dimensions are **unknown**

- * bootstrap is a problem of **infinite-dimensional** analysis

Exponential decoupling thm

(Pappadopulo, S.R, Espin, Rattazzi 2012)

$$\left| \sum_{\substack{k \\ \Delta_k \geq \Delta_*}} f_{12k} \begin{array}{c} \diagup \\ \diagdown \end{array} \begin{array}{c} O_k \\ \text{---} \end{array} \begin{array}{c} \diagdown \\ \diagup \end{array} f_{34k} \right| \lesssim \exp(-C\Delta_*)$$

(valid for any unitary CFT, and any d)

\Rightarrow makes bootstrap numerically tractable

$O(100)$ papers since 2008

Red = to do

$2 < d \leq 4$ and $d=2$ $SL(2, \mathbb{C})$

- bounds (dims, ope coeffs, central charges)
- extremal spectrum studies
- numerical techniques (simplex method, SDPA, dual/direct) **ellipsoid method** minor method by Gliozzi
- global syms
- impact of SUSY
- large $N \leftrightarrow$ AdS
- lightcone results: large spin, small twist **numerical impact?**
- **several correlators**
- **basis optimization**
- **external states with spin (T, J)**

CFT_d with bdry

Conformal defects \leftrightarrow $d=1$ bootstrap **$d \rightarrow 1$ limit**

[Study of 2d CFT torus partition functions]

$d=2$ non-rational Virasoro bootstrap

Bootstrap on other geometries ($\mathbb{R}^{d-1} \times S^1$)

Conformal blocks

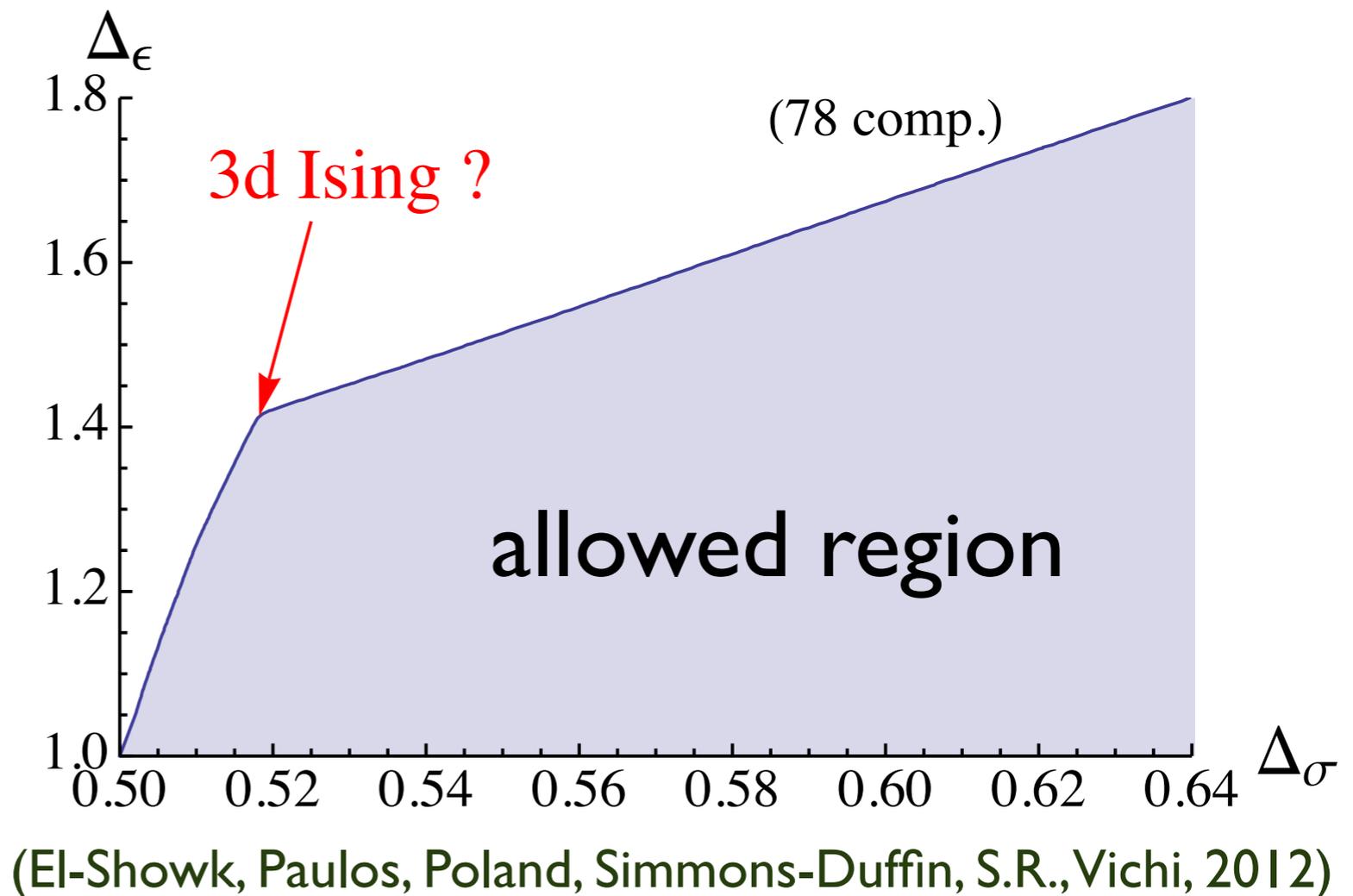
- exact expressions
- power series expansions
- recursions
- for ops with spin
- large d limit

Bootstrap analysis for 3d Ising

Two most important operators:

σ (Z_2 odd), ε (Z_2 even)

$$\sigma \times \sigma = 1 + \varepsilon + \dots$$



Ingredients which go into this plot:

- conformal symmetry & unitarity
- bootstrap equation for $\langle \sigma \sigma \sigma \sigma \rangle$

Conjecture: 3d Ising spectrum maximizes scalar gap

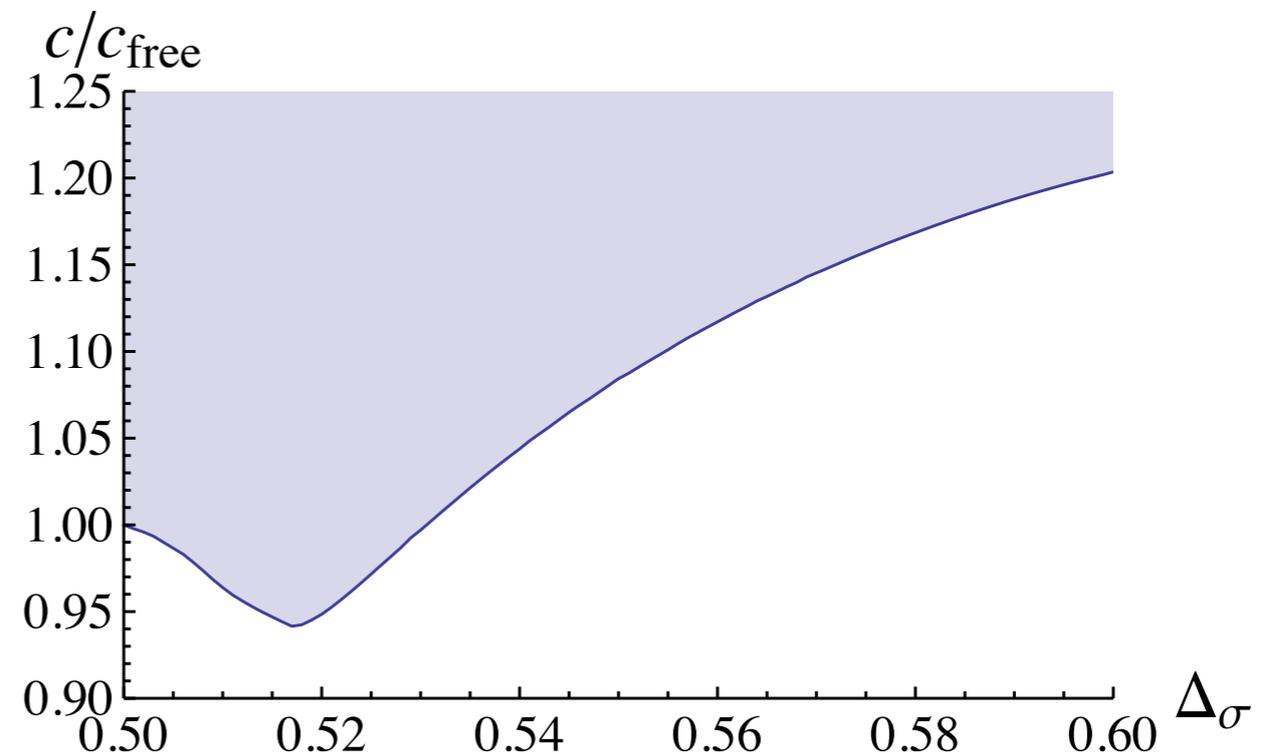
c-Minimization

Another interesting quantity is central charge c :

$$\langle T_{\mu\nu}(x)T_{\lambda\sigma}(y) \rangle \propto c$$

$$\sigma \times \sigma = 1 + \varepsilon + T_{\mu\nu} + \dots$$

contribution goes as c^{-1}

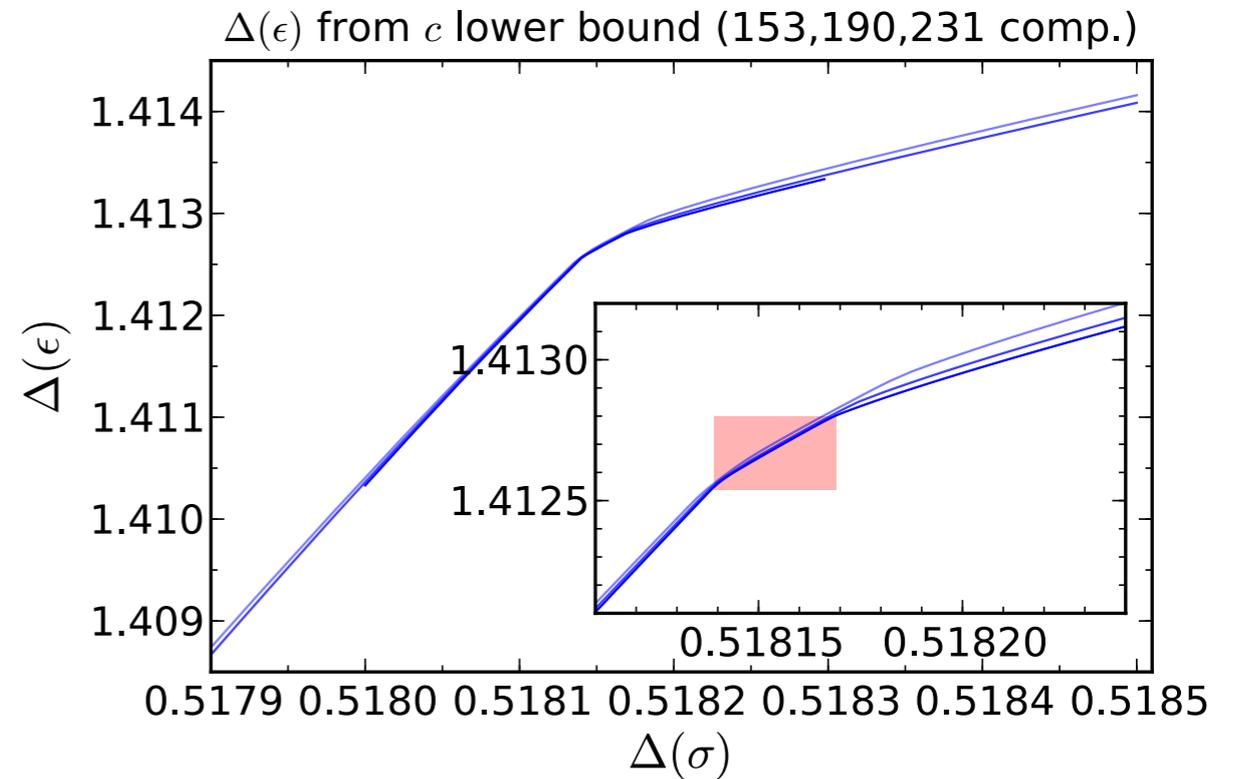
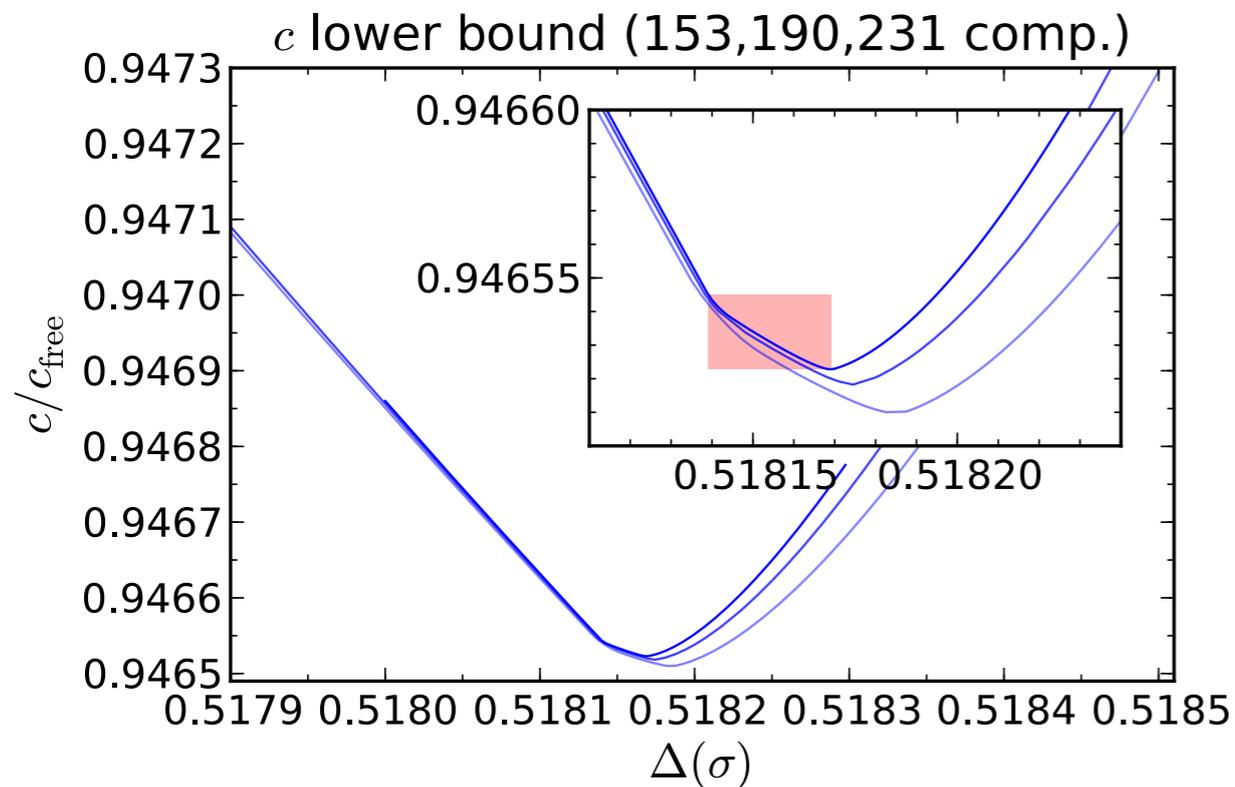


(El-Showk, Paulos, Poland, Simmons-Duffin, S.R., Vichi, 2012)

Conjecture: 3d Ising CFT has minimal central charge

Precision bootstrap

(El-Showk, Paulos, Poland, Simmons-Duffin, S.R., Vichi, 2014)



spin & \mathbb{Z}_2	name	Δ	OPE coefficient
$\ell = 0, \mathbb{Z}_2 = -$	σ	0.518154(15)	
$\ell = 0, \mathbb{Z}_2 = +$	ϵ	1.41267(13)	$f_{\sigma\sigma\epsilon}^2 = 1.10636(9)$
	ϵ'	3.8303(18)	$f_{\sigma\sigma\epsilon'}^2 = 0.002810(6)$
$\ell = 2, \mathbb{Z}_2 = +$	T	3	$c/c_{\text{free}} = 0.946534(11)$
	T'	5.500(15)	$f_{\sigma\sigma T'}^2 = 2.97(2) \times 10^{-4}$

Results by traditional techniques

$$\Delta_\sigma, \Delta_\epsilon, \Delta_{\epsilon'} \leftrightarrow \eta, \nu, \omega$$

	ref	year	Method	ν	η	ω
Zinn-Justin {	[22]	1998	ϵ -exp	0.63050(250)	0.03650(500)	0.814(18)
	[22]	1998	3D exp	0.63040(130)	0.03350(250)	0.799(11)
Vicari	[23]	2002	HT	0.63012(16)	0.03639(15)	0.825(50)
Blöte	[24]	2003	MC	0.63020(12)	0.03680(20)	0.821(5)
Hasenbusch	[25]	2010	MC	0.63002(10)	0.03627(10)	0.832(6)
		this work		0.62999(5)	0.03631(3)	0.8303(18)

Numerical resources:

2-3 CPU years

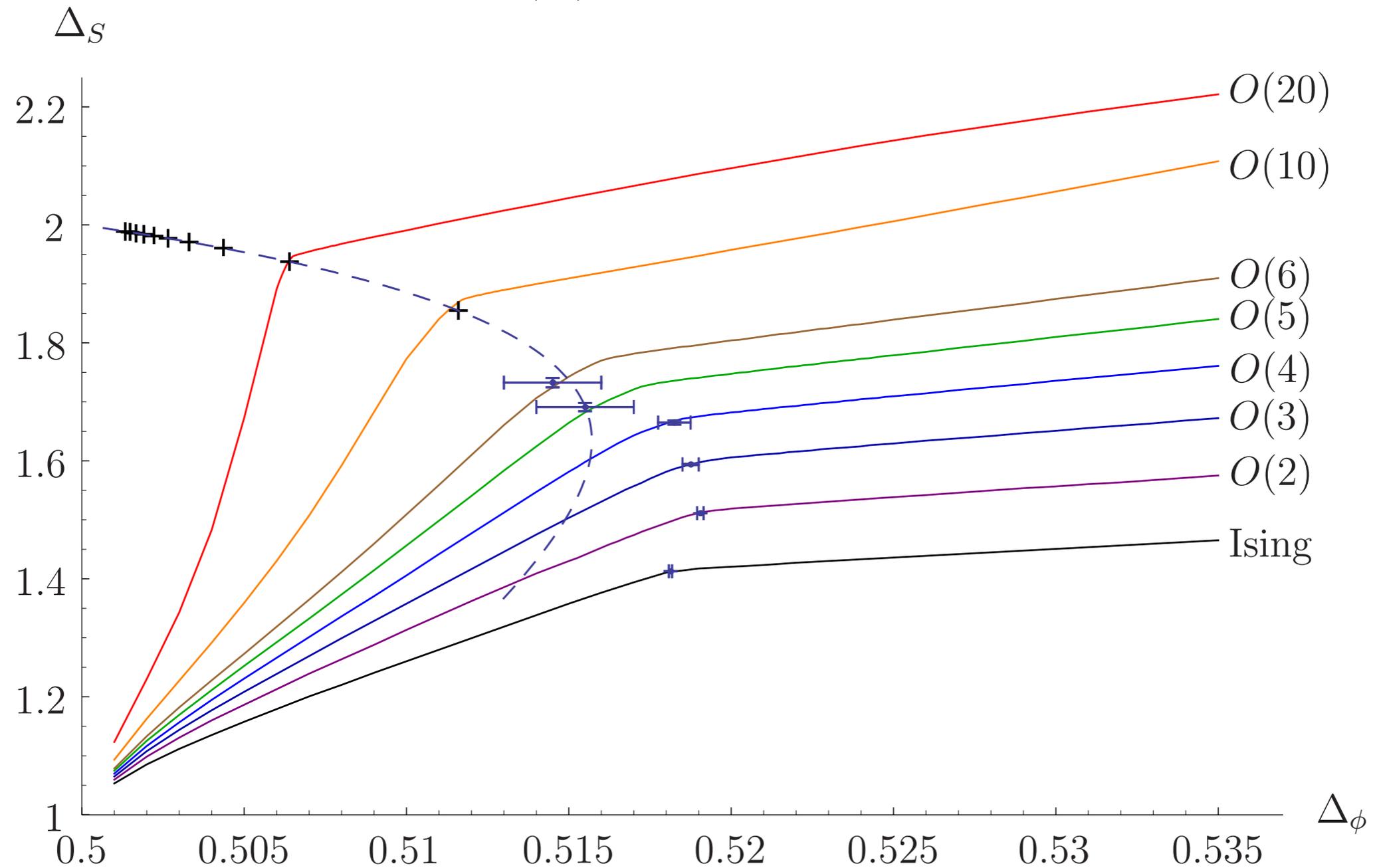
vs 30 CPU years for Hasenbusch

(1000 CPU years would be needed to reach our precision)

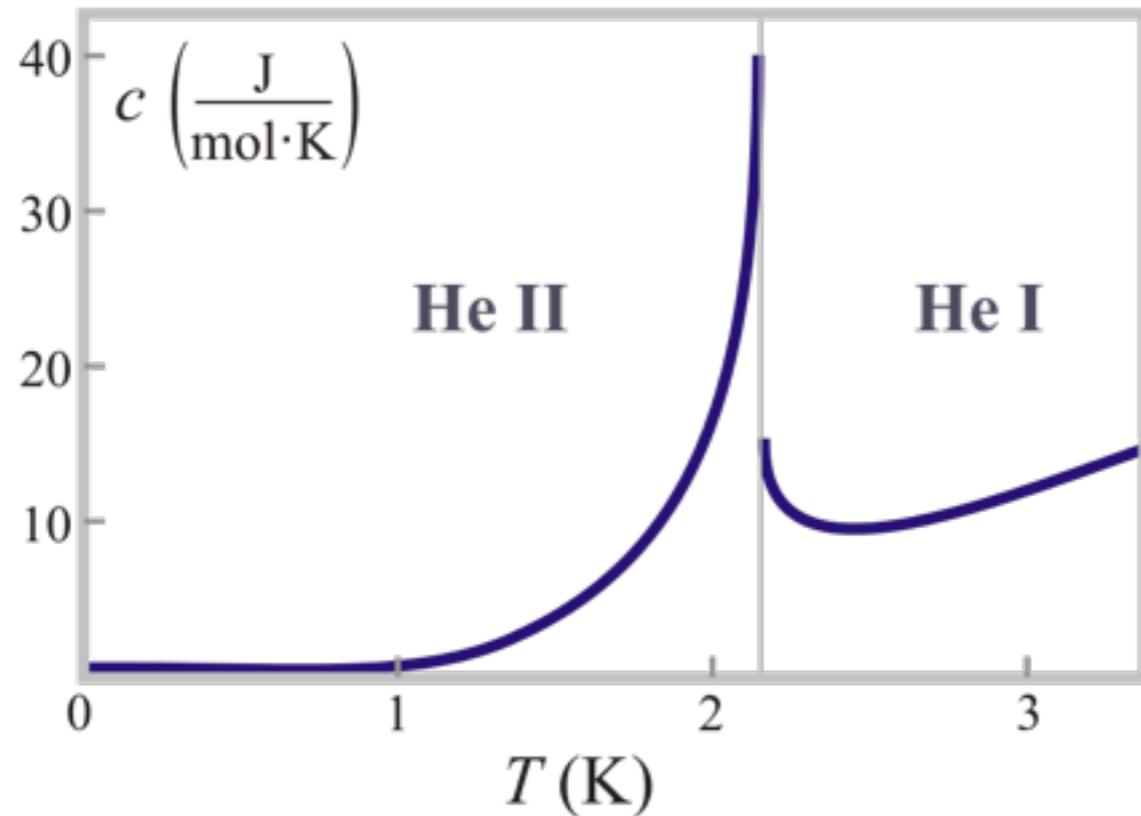
$O(N)$ models

(Kos, Poland, Simmons-Duffin, 2013)

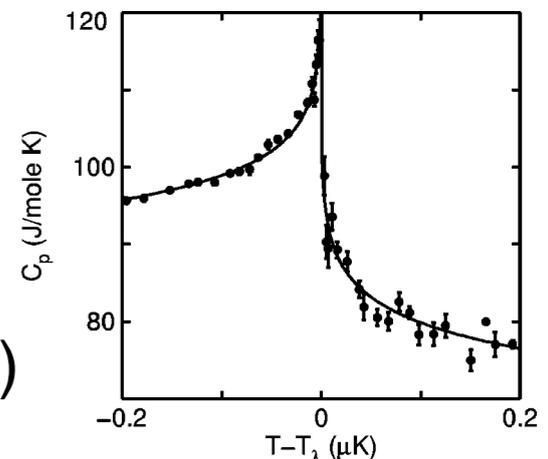
$O(N)$ Singlet Bounds



O(2) controversy



λ -point of liquid Helium-4



Space shuttle measurement (1992): $\Delta_\epsilon = 1.5094(2)$

Best lattice prediction(2006): $\Delta_\epsilon = 1.5112(2)$

Best RG prediction (Guida,Zinn-Justin 1998): $\Delta_\epsilon = 1.5081(33)$ **inconclusive**

Bootstrap should be able to settle this

Conclusions

General lesson:

much to learn from conformal bootstrap,
which in $d \geq 3$ was unjustly neglected for 40 years

Particular lesson:

some models have extremal properties \Rightarrow more tractable

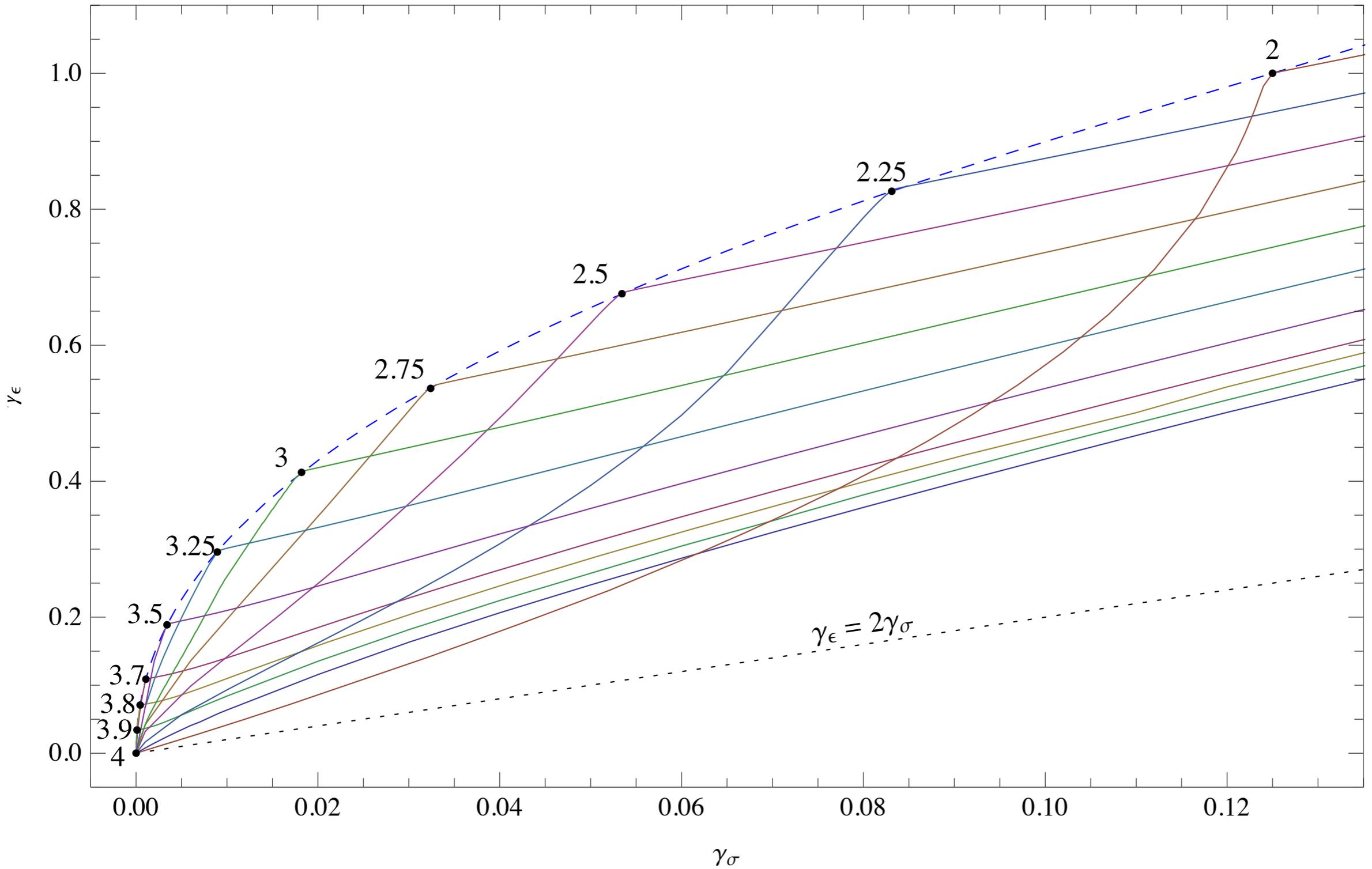
3d Ising CFT is one of them. **Does it mean it's exactly solvable?**

Big picture:

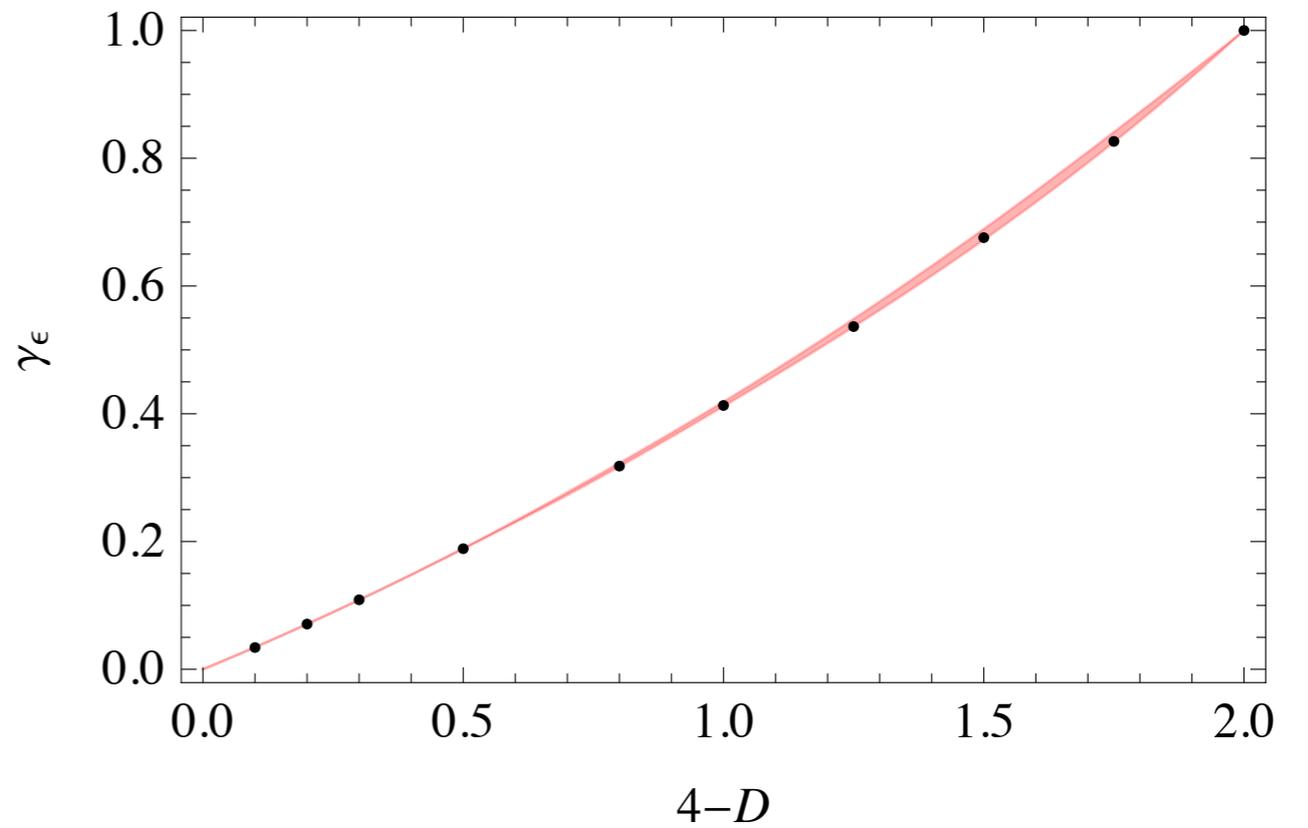
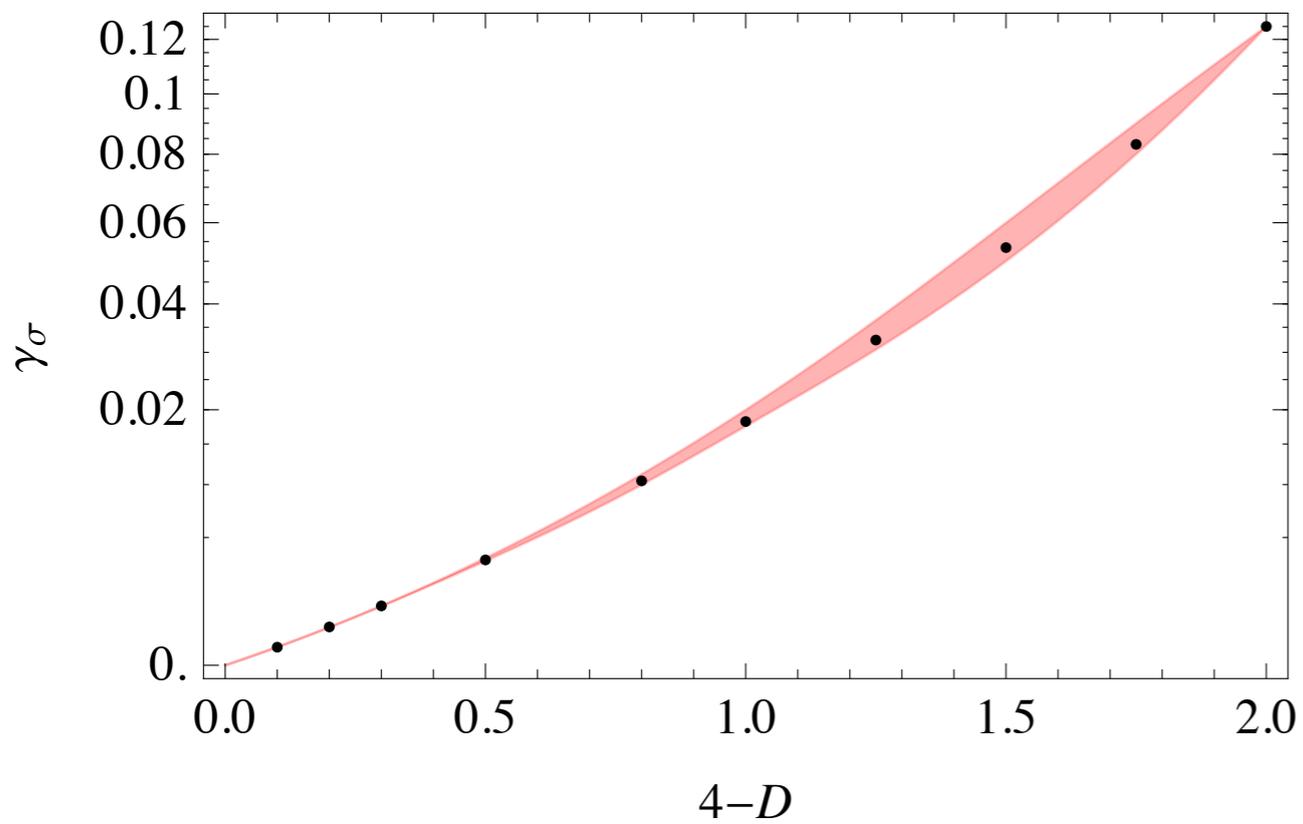
CFTs are important, but RG flows are even more important
They have to be solved as well

Backup

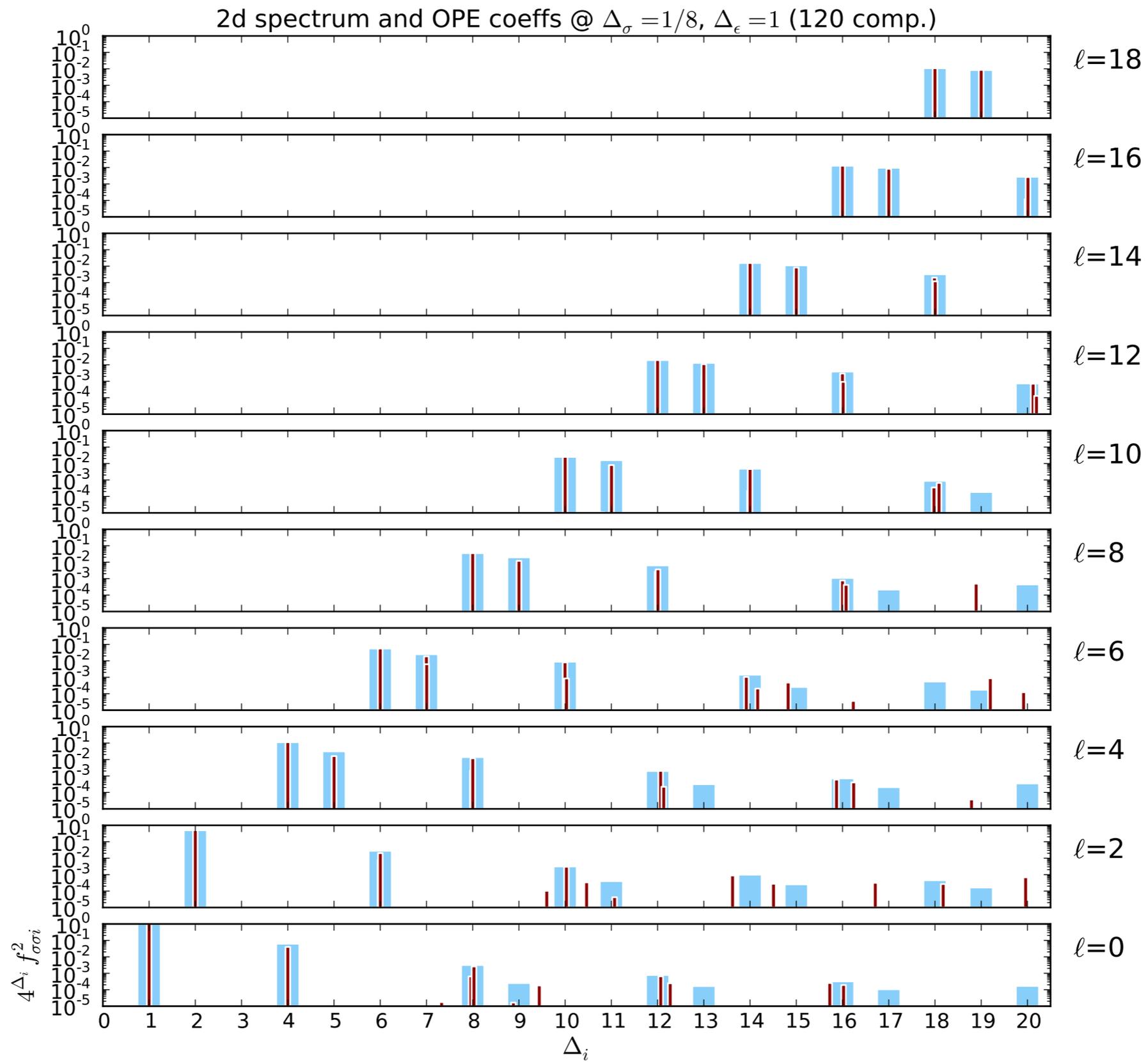
Fractional d



Comparing with $4-\epsilon$ expansion



2d



2d

