# Exact Solutions of 2d Supersymmetric gauge theories

Abhijit Gadde, Caltech

w. Sergei Gukov and Pavel Putrov

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- IR fixed points
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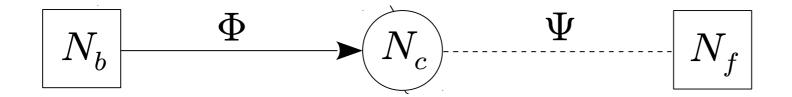
$$\Psi = \psi_{-} - \sqrt{2}\theta^{+}G - i\theta^{+}\bar{\theta}^{+}\partial_{+}\psi_{-}$$

- Complex left-moving fermion
- Vector multiplet:
  - ullet Gauge invariant d.o.f.: Fermi multiplet  $\Lambda$

• Similar to 4d N=1 SQCD, but 2 types of matter

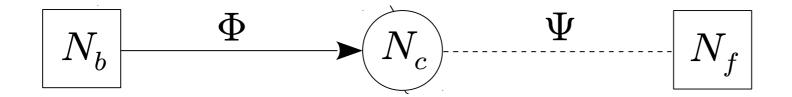
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$$\int d\theta \, \Gamma P \Phi + \int d\theta \, \Lambda$$

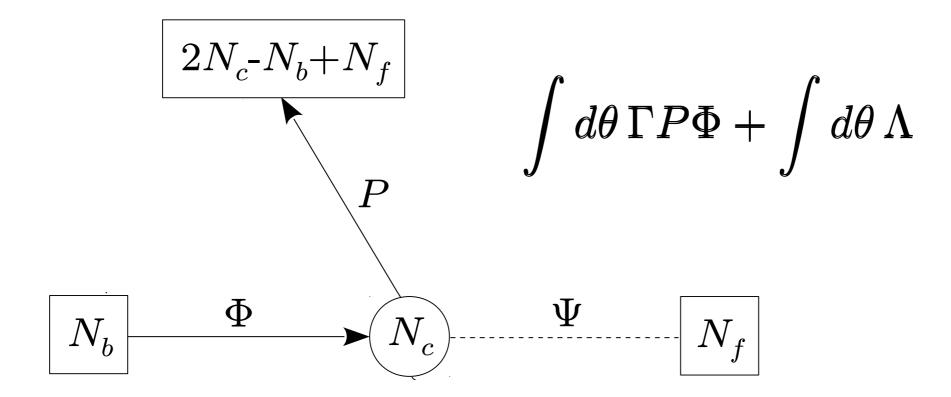


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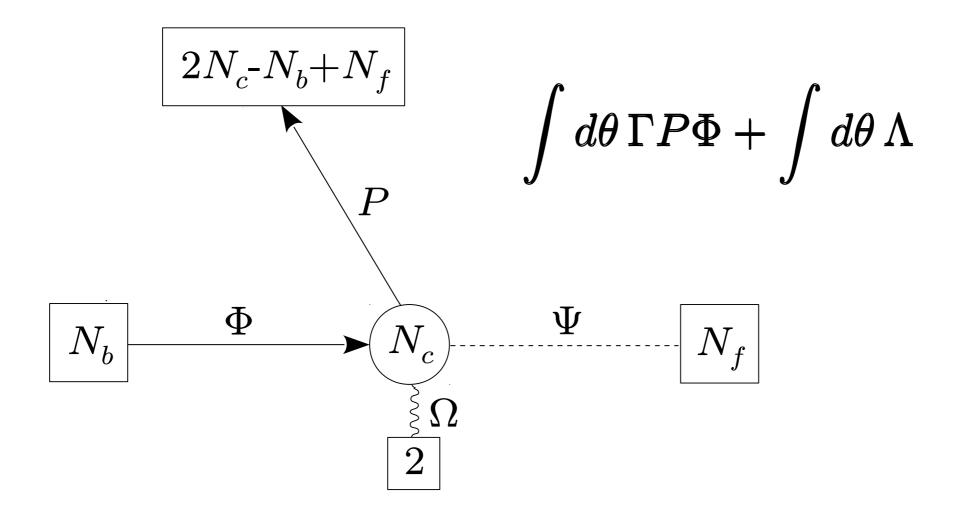
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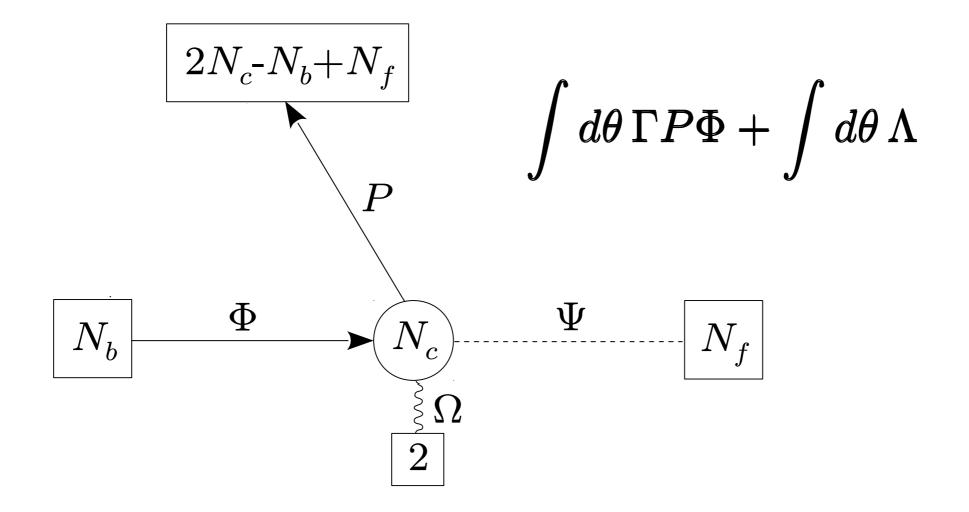
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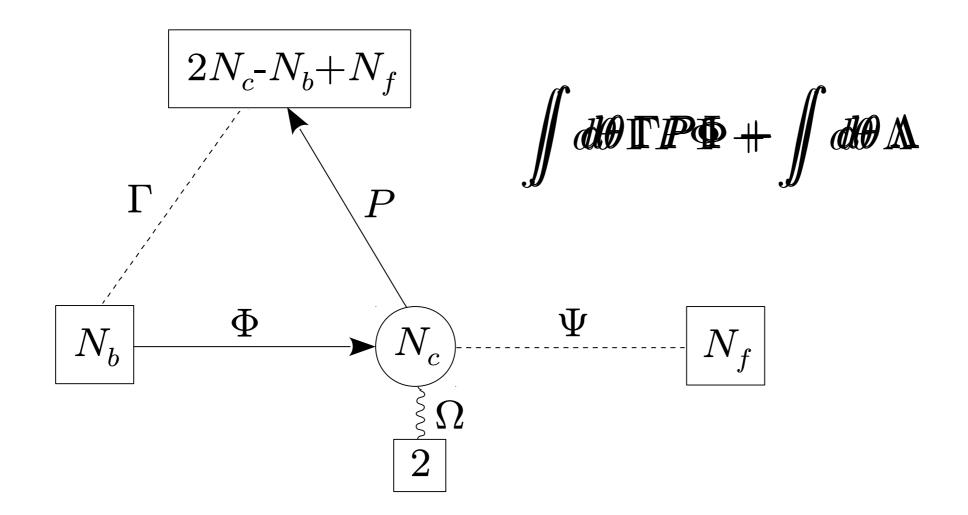
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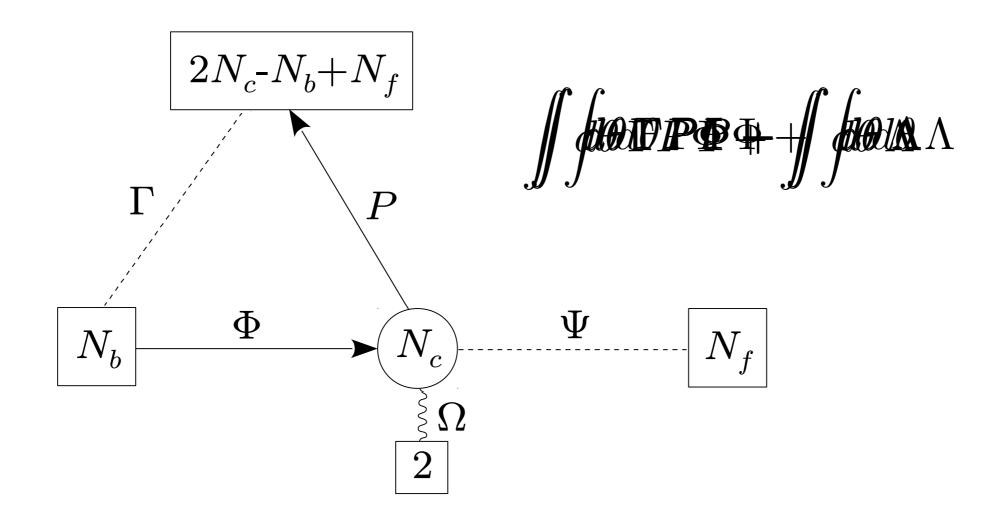
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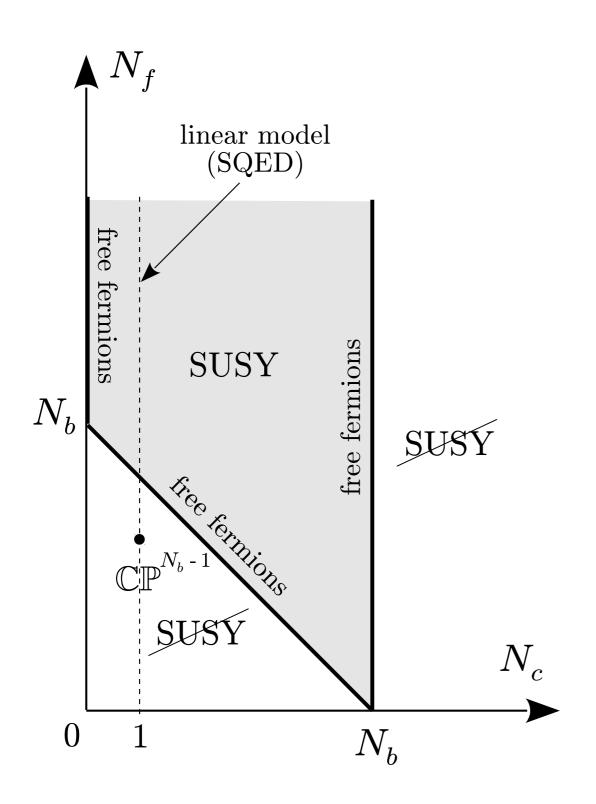


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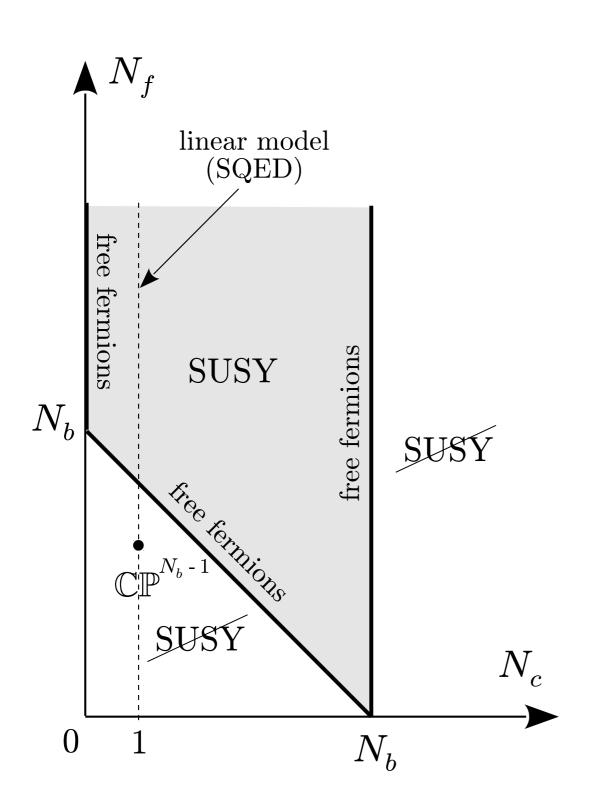


## Phase diagram

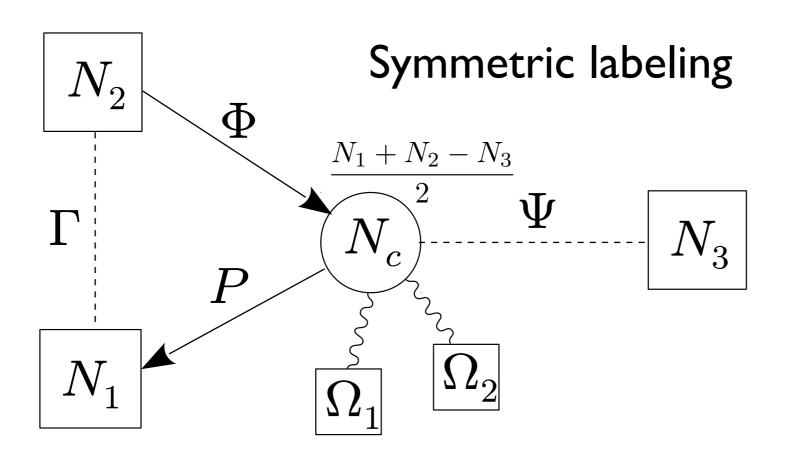
#### Phase diagram

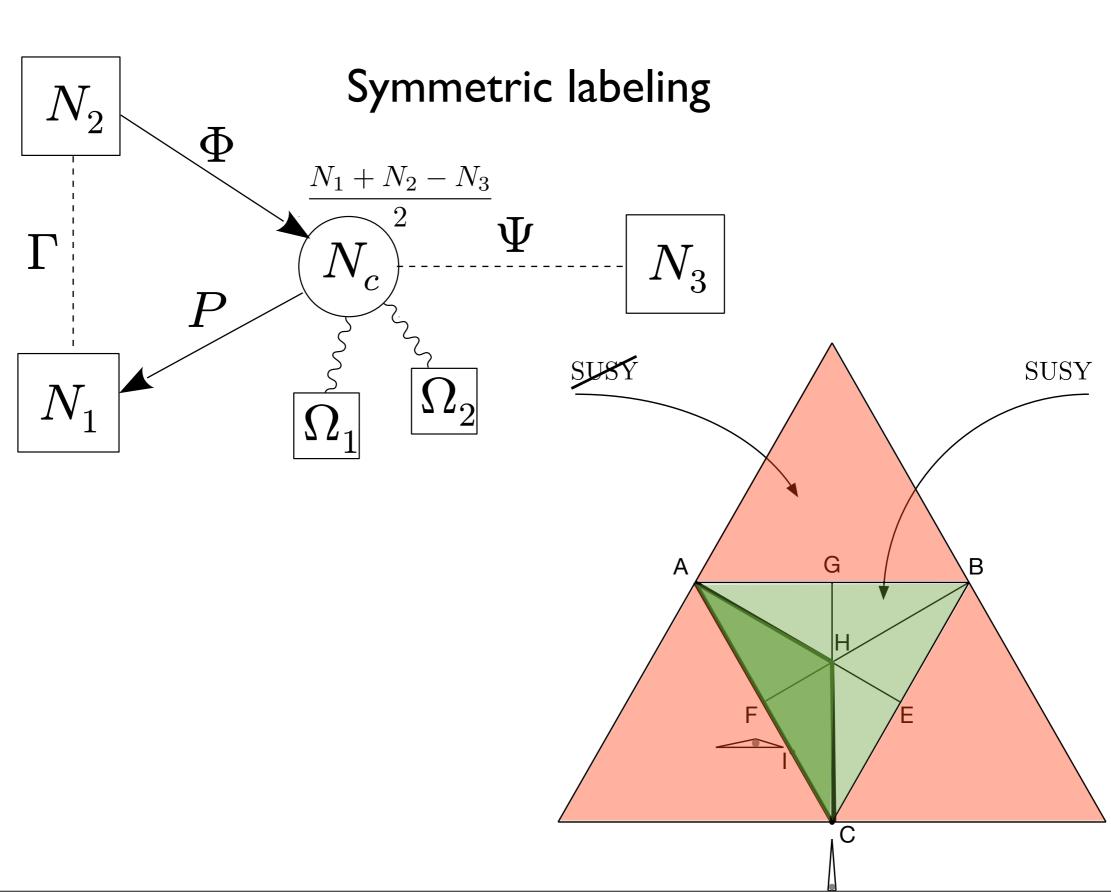


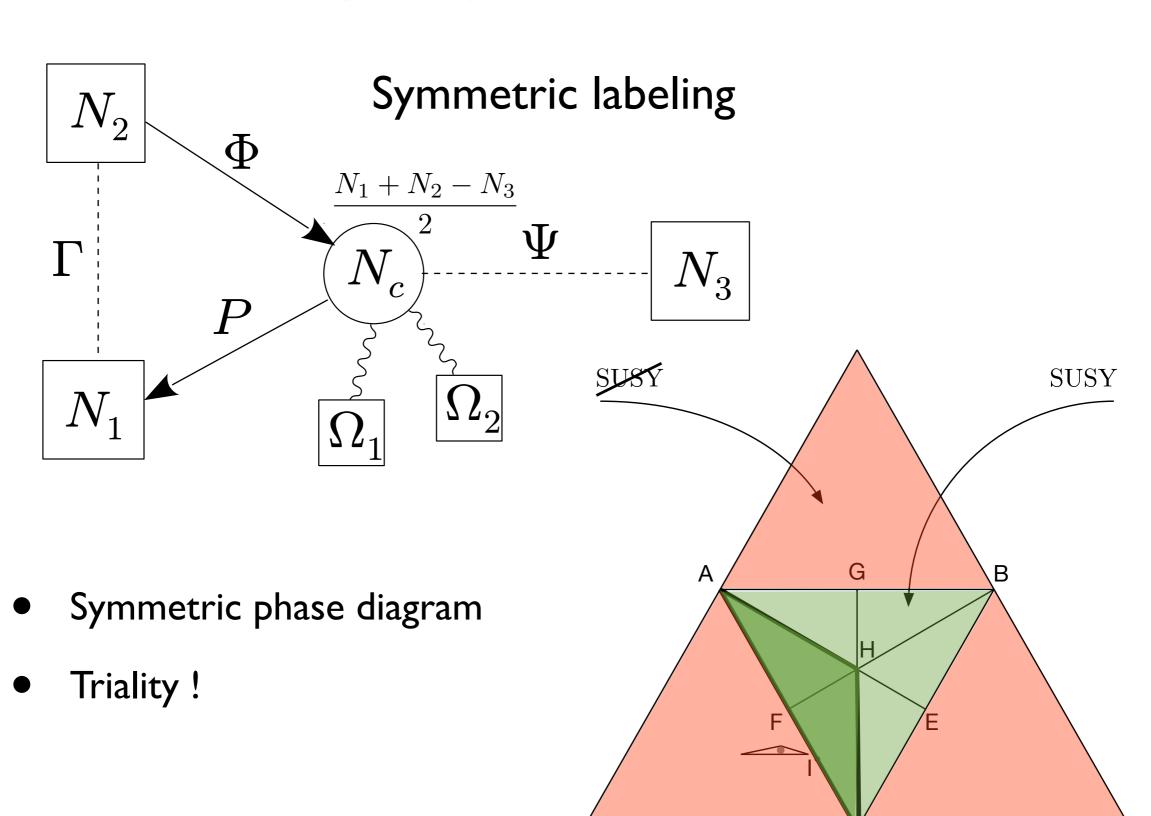
#### Phase diagram



- Produced by superconformal index
- SUSY preserved in the shaded region
- SUSY broken outside
- Gapless theory at all points
- We will focus on the shaded region







$$\operatorname{Tr} \gamma^{3} J_{\mathrm{SU}(\mathrm{N}_{1})} J_{\mathrm{SU}(\mathrm{N}_{1})} = \frac{N_{1} + N_{2} - N_{3}}{2} T_{P}(\Box) - N_{2} T_{\Gamma}(\overline{\Box}) = -\frac{1}{4} (-N_{1} + N_{2} + N_{3})$$

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#### Abelian symmetries

Pick a basis so that cross anomalies are zero

	$\Phi$	_	P	$\Gamma$	I	$\Omega_2$
$U(1)_{(1)}$	0	0	1	-1	$-N_1$	0
$U(1)_{(2)}$	-1	_	0		$-N_2$	_
$U(1)_{(3)}$	0	1	0	0	0	$N_3$

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$$Tr\gamma^{3}J_{U(1)_{1}}J_{U(1)_{1}} = -NN_{1}/2$$

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Poincare symmetry

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Virasoro

Poincare symmetry

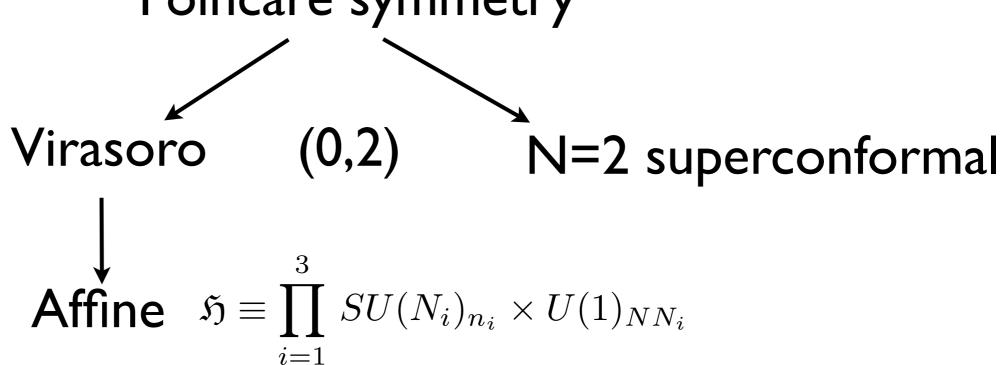
Virasoro

N=2 superconformal

Poincare symmetry

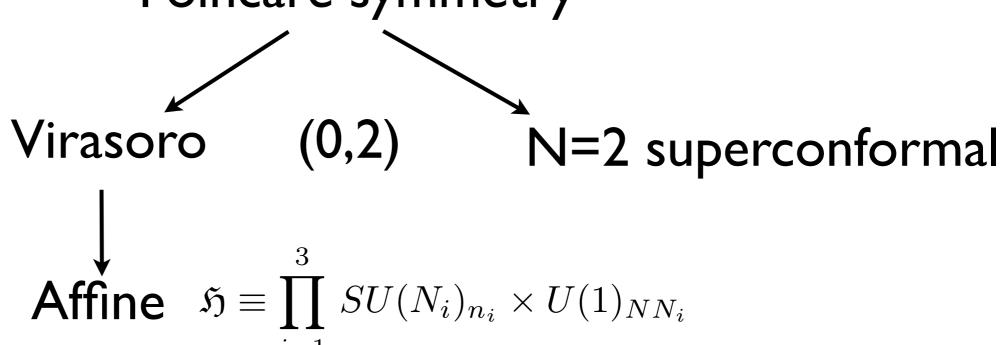
Virasoro (0,2) N=2 superconformal

Poincare symmetry



### Low energy physics

Poincare symmetry

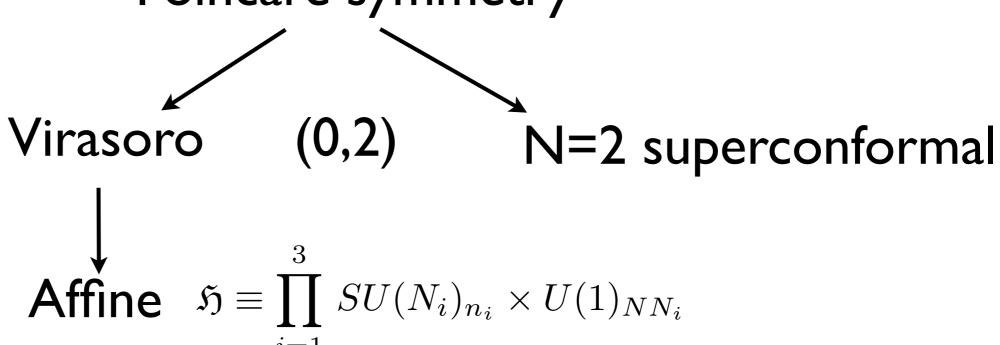


 The central charges can be determined from c-extremization and gravitational anomaly

$$c_{\rm R} = 3 {\rm Tr} \gamma^3 R R \,, \qquad c_{\rm R} - c_{\rm L} = {\rm Tr} \gamma^3$$

### Low energy physics

Poincare symmetry



 The central charges can be determined from c-extremization and gravitational anomaly

$$\begin{split} c_{\rm R} &= 3 {\rm Tr} \gamma^3 R R \;, \qquad c_{\rm R} - c_{\rm L} = {\rm Tr} \gamma^3 \\ c_{\rm R} &= \frac{3}{4} \frac{(-N_1 + N_2 + N_3)(N_1 - N_2 + N_3)(N_1 + N_2 - N_3)}{N_1 + N_2 + N_3} \\ c_{\rm L} &= c_{\rm R} - \frac{1}{4} (N_1^2 + N_2^2 + N_3^2 - 2 N_1 N_2 - 2 N_2 N_3 - 2 N_3 N_1) + 2 \end{split}$$

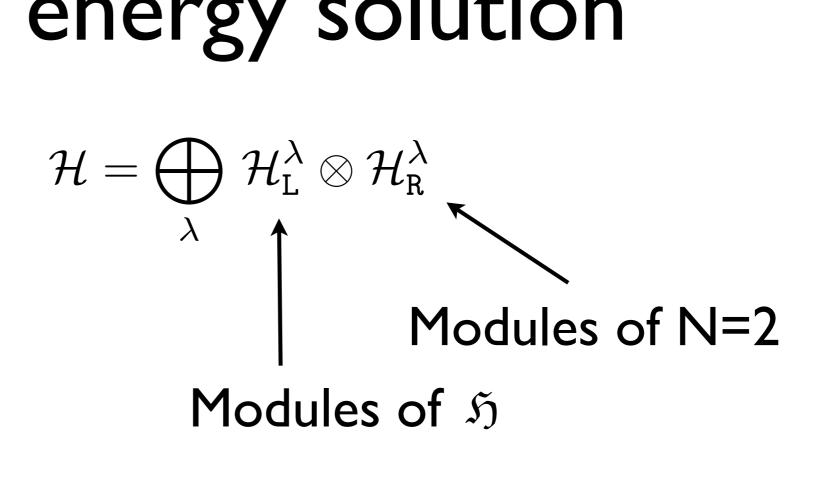
$$\mathcal{H} = igoplus_{\lambda}^{\lambda} \mathcal{H}_{\mathtt{L}}^{\lambda} \otimes \mathcal{H}_{\mathtt{R}}^{\lambda}$$

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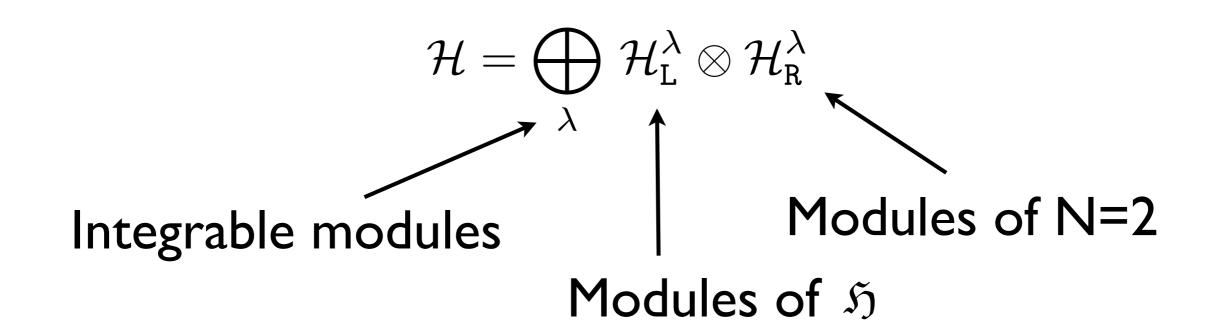
$$\downarrow^{\lambda} \qquad \uparrow$$

$$\downarrow^{\lambda} \qquad \qquad \mathsf{Modules of } \mathfrak{H}$$

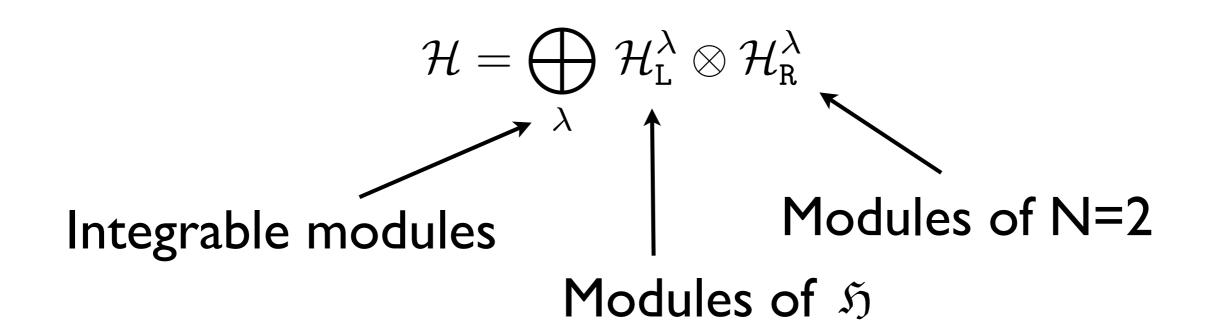
energy solution 
$$\mathcal{H} = \bigoplus_{\lambda} \mathcal{H}_{L}^{\lambda} \otimes \mathcal{H}_{R}^{\lambda}$$
 Modules of N=2 Modules of  $\mathfrak{H}$ 



Sugawara central charge =  $C_{\rm L}$ 



• Sugawara central charge =  $C_{\rm L}$ 



- Sugawara central charge =  $c_{\rm L}$
- Immense simplification: rational CFT
- ullet Modular invariance of the partition function helps fix  $\, {\cal H}^{\lambda}_{R} \,$

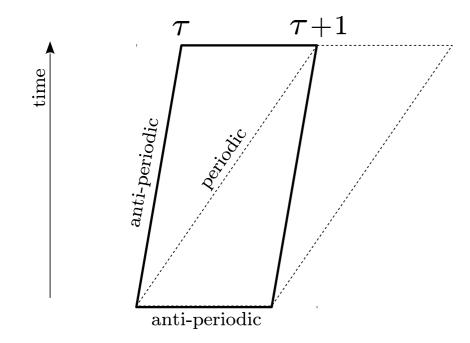
$$Z(\tau, \xi_i; \bar{\tau}, \bar{\eta}) := \operatorname{Tr}_{\mathcal{H}} e^{2\pi i (\tau L_0 + \sum_i \xi_i H_0^i - \bar{\tau} \bar{L}_0 - \bar{\eta} \bar{J}_0)}$$

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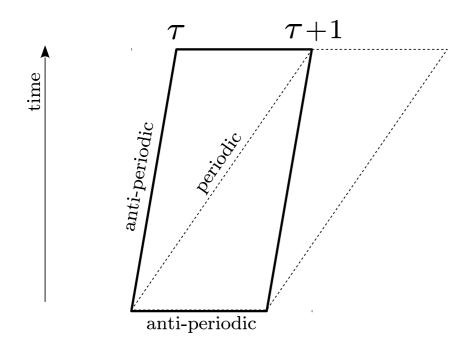
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55 symmetry Cartan R symmetry Cartan



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ullet Invariant under S and  $T^2$ 

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 $\tau \qquad \tau+1$ 

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anti-periodic

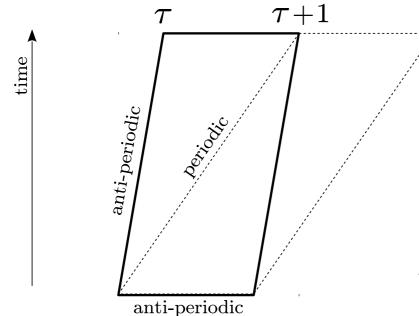
Affine character

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$$Z(\tau, \xi; \bar{\tau}, \bar{\eta}) = Z(-\frac{1}{\tau}, \frac{\xi}{\tau}; -\frac{1}{\bar{\tau}}, -\frac{\bar{\eta}}{\bar{\tau}})$$
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- ullet K is NOT the anti-holomorphic affine character of  ${\mathfrak H}$
- ullet Note that characters of level-rank dual  $\, {\mathfrak H}^t \,$  transform with  $\bar{S}$

$$\mathfrak{H}^t = \prod SU(n_i)_{N_i} \times U(1)_{N_{n_i}}$$

#### To summarize

- K is an N=2 character with central charge cR
- It transforms as a character of holomorphic  $\mathfrak{H}^t$  under modular S-transformation
- Singlet under all affine symmetries

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• K is a character of the Kazama- Suzuki coset  $[\mathfrak{G}]/[\mathfrak{H}^t]$ (For appropriate  $\mathfrak{G}$ )

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- $[\mathfrak{g}]_k$  is SUSY extension of WZW model  $\mathfrak{g}$  at level k
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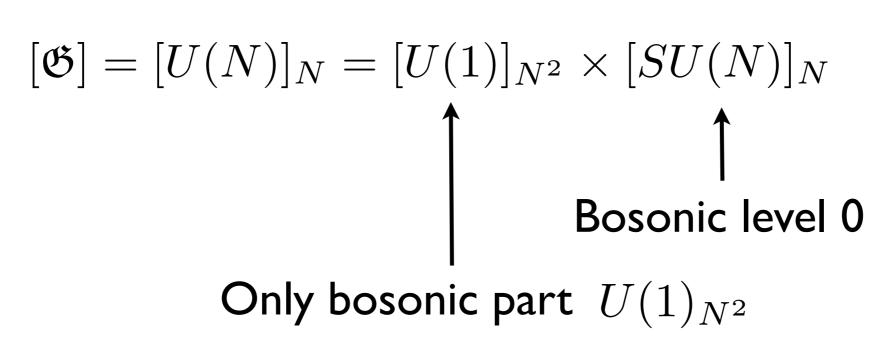
 Matching the right-moving central charge with that of the coset:

$$c_{[\mathfrak{G}]} = N^2$$

$$[\mathfrak{G}] = [U(N)]_N = [U(1)]_{N^2} \times [SU(N)]_N$$

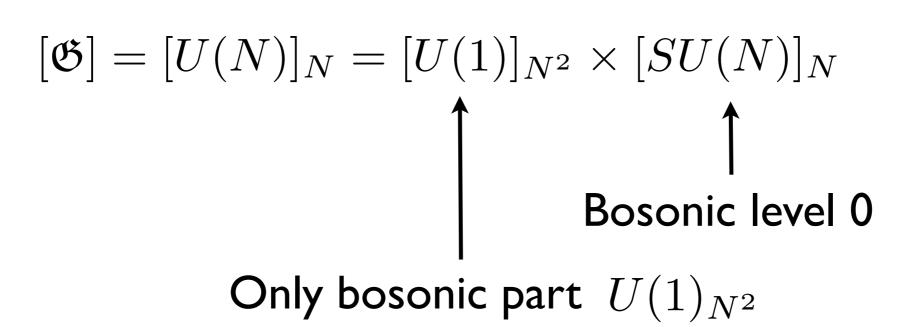
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$$\uparrow$$
 Bosonic level 0

$$[\mathfrak{G}]=[U(N)]_N=[U(1)]_{N^2}\times[SU(N)]_N$$
 
$$\uparrow$$
 Bosonic level 0 Only bosonic part  $U(1)_{N^2}$ 

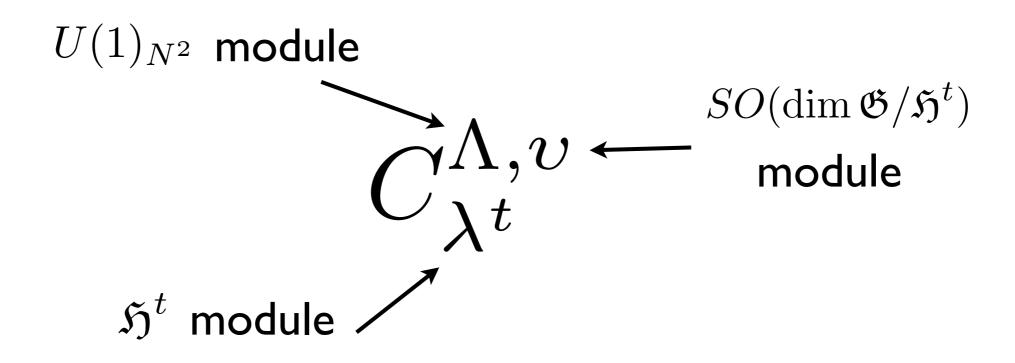


Coset character C is a branching function

$$C_{\lambda^t}^{\Lambda,\upsilon}$$



Coset character C is a branching function



### Solution

• We pick modular invariant combinations  $(\Lambda_0, \upsilon_0)$ 

Then 
$$K_{\lambda} = \sum_{\lambda^t} L_{\lambda,\lambda^t} \, C_{\lambda^t}^{\Lambda_0,\upsilon_0}$$

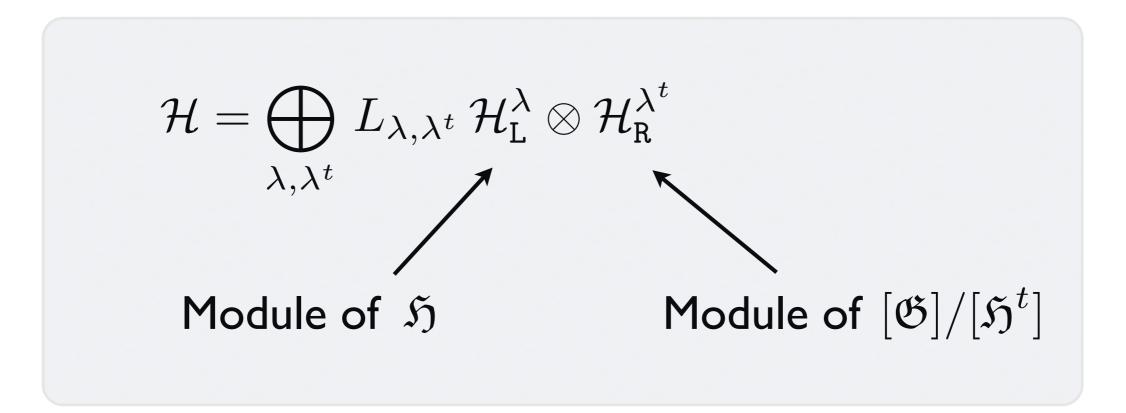
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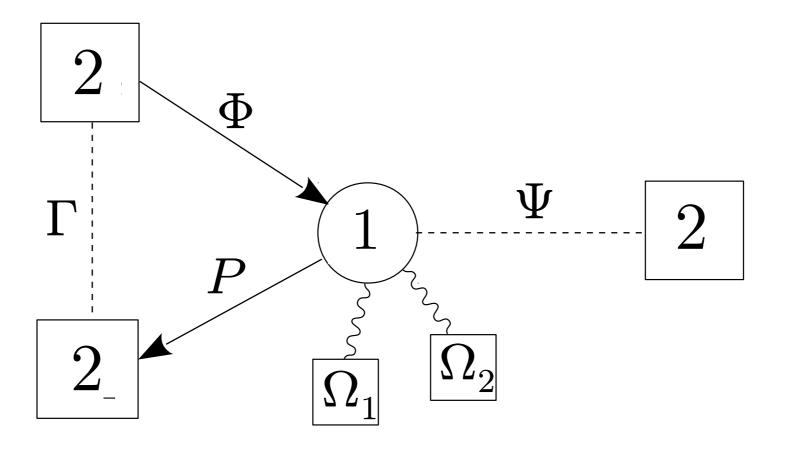
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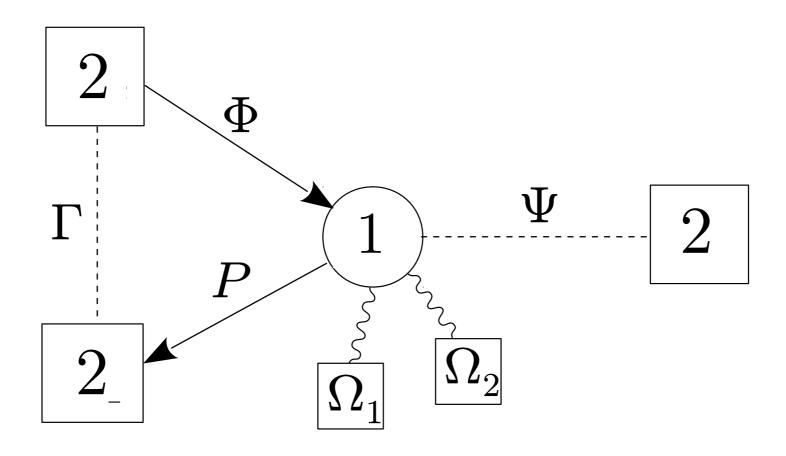
$$\mathcal{H} = \bigoplus_{\lambda,\lambda^t} L_{\lambda,\lambda^t} \, \mathcal{H}_{\mathtt{L}}^{\lambda} \otimes \mathcal{H}_{\mathtt{R}}^{\lambda^t}$$
 Module of  $\mathfrak{H}$  Module of  $[\mathfrak{G}]/[\mathfrak{H}^t]$ 

Matches with the UV computation of the index

## Example

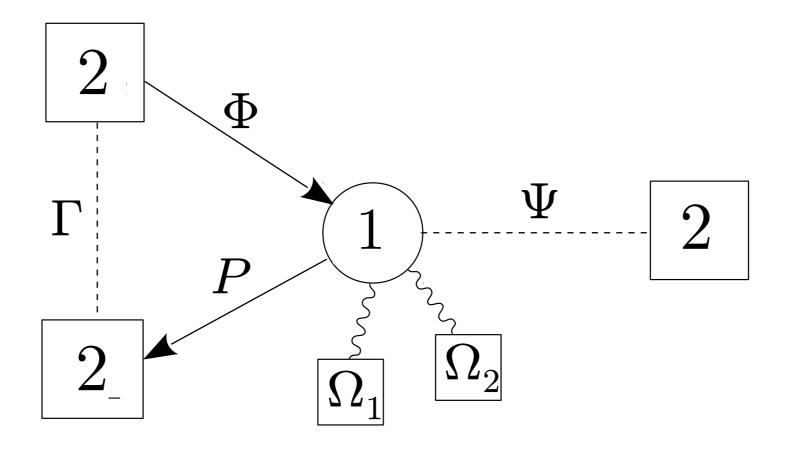


### Example



$$\mathfrak{H} = \left( SU(2)_1 \times U(1)_6 \right)^3 \qquad [\mathfrak{G}]/[\mathfrak{H}^t] = [U(3)]_3/[U(1)_3]^3$$

### Example



$$\mathfrak{H} = \left(SU(2)_1 \times U(1)_6\right)^3 \qquad [\mathfrak{G}]/[\mathfrak{H}^t] = [U(3)]_3/[U(1)_3]^3$$
 
$$\uparrow \qquad \qquad \qquad \qquad \mathsf{c=I \ minimal \ model}$$

$$Z_{\mathcal{T}_{222}} = \chi_{(0,0)}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \Big( \Xi_{0,0,0}(\tau) + \Xi_{1,1,1}(\tau) + \Xi_{-1,-1,-1}(\tau) \Big)$$

$$+ \chi_{(\frac{1}{6}, \frac{1}{3})}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \Big( \Xi_{1,0,-1}(\tau) + \Xi_{-1,1,0}(\tau) + \Xi_{0,-1,1}(\tau) \Big)$$

$$+ \chi_{(\frac{1}{6}, -\frac{1}{3})}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \Big( \Xi_{-1,0,1}(\tau) + \Xi_{1,-1,0}(\tau) + \Xi_{0,1,-1}(\tau) \Big)$$

#### where

$$\Xi_{a,b,c}(\tau,\xi_{1},\xi_{2},\xi_{3}) := \Xi_{a}(\tau,\xi_{1})\Xi_{b}(\tau,\xi_{2})\Xi_{c}(\tau,\xi_{3})$$

$$\Xi_{-1}(\tau,\xi) := \chi_{(\square,-1)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi) + \chi_{(\cdot,2)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi)$$

$$\Xi_{0}(\tau,\xi) := \chi_{(\cdot,0)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi) + \chi_{(\square,3)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi)$$

$$\Xi_{1}(\tau,\xi) := \chi_{(\square,1)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi) + \chi_{(\cdot,-2)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi).$$

$$Z_{\mathcal{T}_{222}} = \chi_{(0,0)}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \Big( \Xi_{0,0,0}(\tau) + \Xi_{1,1,1}(\tau) + \Xi_{-1,-1,-1}(\tau) \Big)$$

$$+ \chi_{(\frac{1}{6}, \frac{1}{3})}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \Big( \Xi_{1,0,-1}(\tau) + \Xi_{-1,1,0}(\tau) + \Xi_{0,-1,1}(\tau) \Big)$$

$$+ \chi_{(\frac{1}{6}, -\frac{1}{3})}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \Big( \Xi_{-1,0,1}(\tau) + \Xi_{1,-1,0}(\tau) + \Xi_{0,1,-1}(\tau) \Big)$$

#### where

$$\Xi_{a,b,c}(\tau,\xi_{1},\xi_{2},\xi_{3}) := \Xi_{a}(\tau,\xi_{1})\Xi_{b}(\tau,\xi_{2})\Xi_{c}(\tau,\xi_{3})$$

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$$\Xi_{0}(\tau,\xi) := \chi_{(\cdot,0)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi) + \chi_{(\square,3)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi)$$

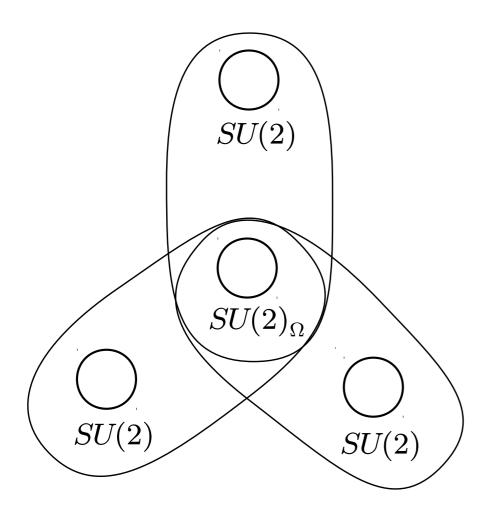
$$\Xi_{1}(\tau,\xi) := \chi_{(\square,1)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi) + \chi_{(\cdot,-2)}^{SU(2)_{1}\times U(1)_{6}}(\tau,\xi).$$

#### Remarkably

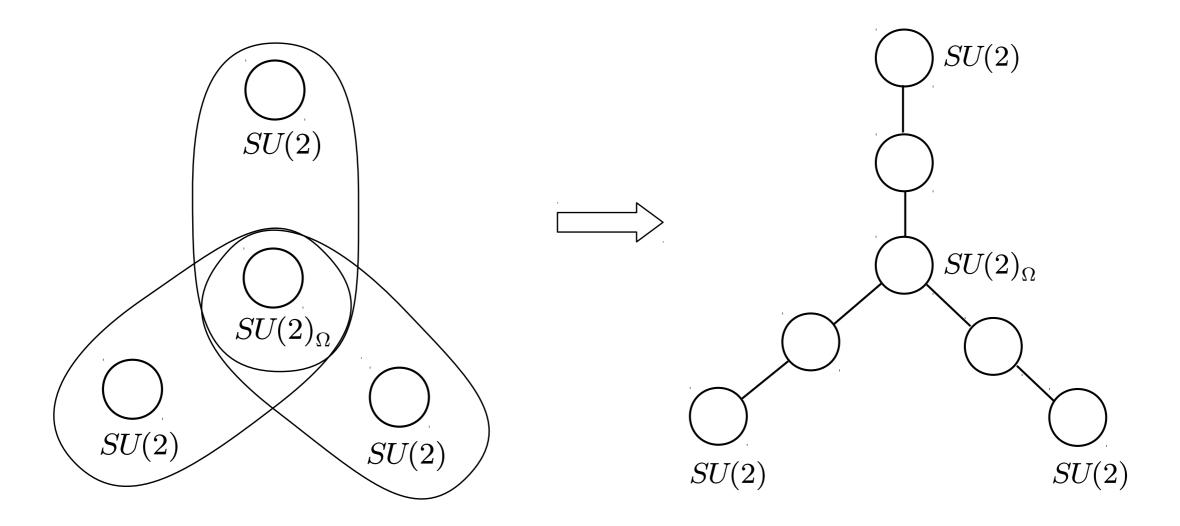
$$Z_{\mathcal{T}_{222}} = \chi_{(0,0)}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \chi_{\bullet}^{(E_6)_1}(\tau, \xi_i) + \chi_{(\frac{1}{6}, \frac{1}{3})}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \chi_{\square}^{(E_6)_1}(\tau, \xi_i) + \chi_{(\frac{1}{6}, -\frac{1}{3})}^{\mathcal{N}=2}(\overline{\tau}, \overline{\eta}) \chi_{\overline{\square}}^{(E_6)_1}(\tau, \xi_i)$$

### Triality and enhancement

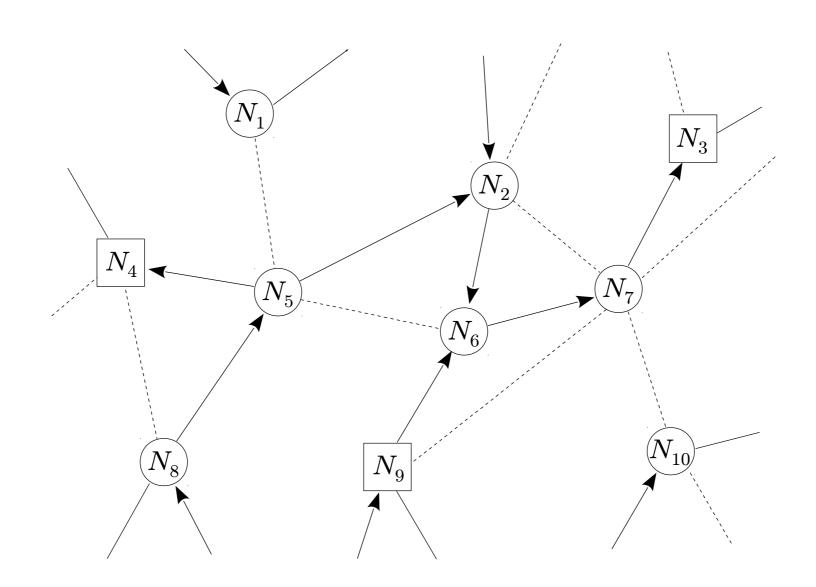
### Triality and enhancement



### Triality and enhancement



### Solution to a general quiver



# Thank you!