

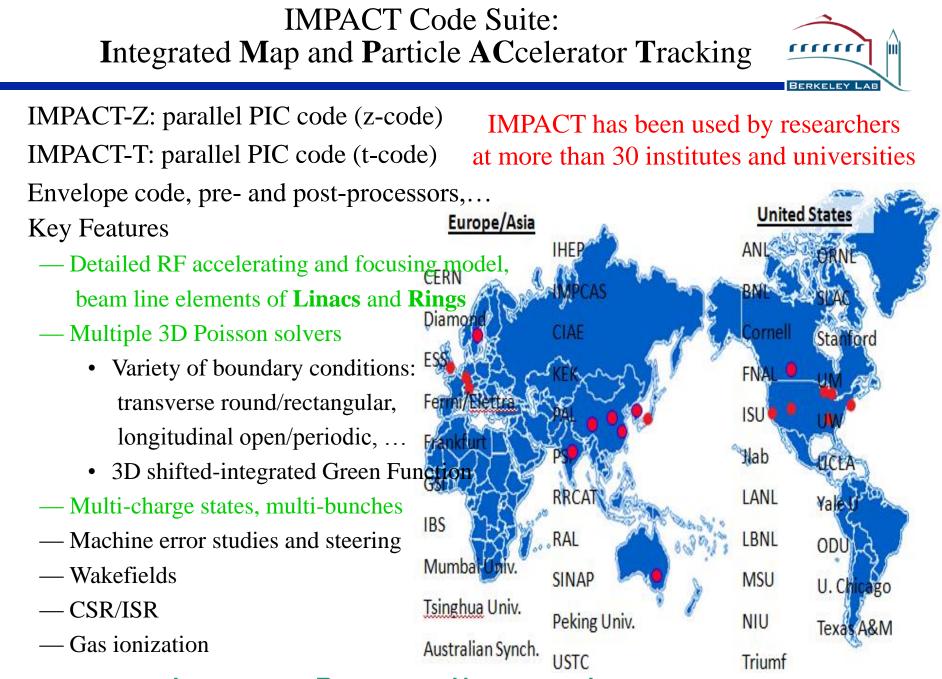
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IMPACT: Benchmarking

Ji Qiang

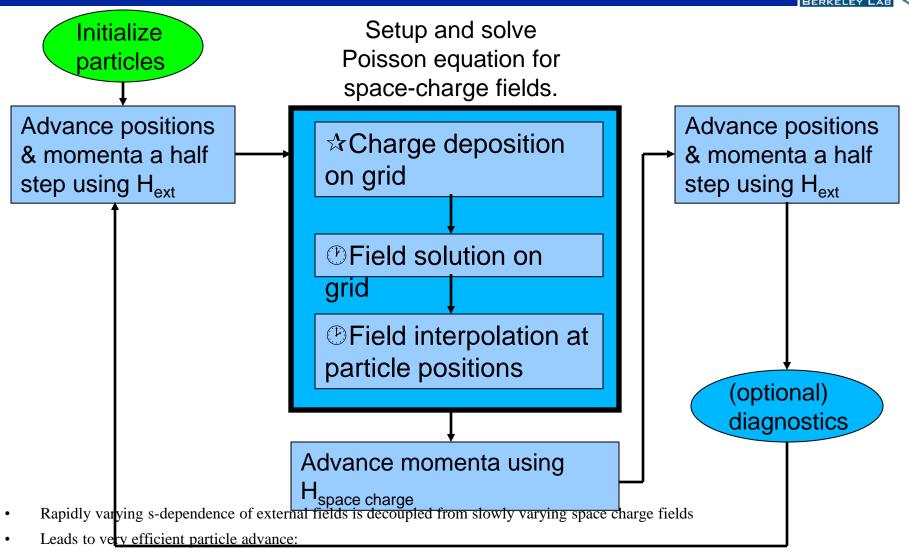
Lawrence Berkeley National Laboratory

CERN Space-Charge Collaboration Meeting, May 20-21, 2014



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Particle-In-Cell Simulation with Split-Operator Method



- Do not take tiny steps to push millions billions of particles
- Do take tiny steps to compute maps; then push particles w/ maps LAWRENCE BERKELEY NATIONAL LABORATORY

IMPACT-Z for Space-Charge Study in Ring (fully 6D symplectic tracking)

Drift – exact solution

Combined function bend – direct symplectic integrator Quadrupole – direct symplectic integrator RF cavity – thin lens kick Nonlinear Multipoles – thin lens kick Space-charge forces– lumped thin lens kick

An Example: Symlectic Tracking through a Combined Function Bend

$$\begin{split} H &= H_1 + H_2 \\ H_1 &= -\left(1 + \frac{x}{\rho_c}\right) \sqrt{(1+\delta)^2 - p_x^2 - p_y^2} + b_0 x + b_0 \frac{x^2}{2\rho_c} \\ H_2 &= V(x, y, \rho_c) \\ \end{split}$$

Exact solution for $H_I \quad x^f = \frac{\rho_c}{b_0} \left(\frac{1}{\rho_c} \sqrt{(1+\delta)^2 - p_x^{f^2} - p_y^2} - \frac{dp_x^f}{ds} - b_0\right) \\ p_x^f &= p_x \cos\left(\frac{s}{\rho_c}\right) + \left(\sqrt{(1+\delta)^2 - p_x^2 - p_y^2} - b_0(\rho_c + x)\right) \sin\left(\frac{s}{\rho_c}\right) \\ y^f &= y + \frac{p_y s}{b_0 \rho_c} + \frac{p_y}{b_0} \left\{ \arcsin\left(\frac{p_x}{\sqrt{(1+\delta)^2 - p_y^2}}\right) - \arcsin\left(\frac{p_x^f}{\sqrt{(1+\delta)^2 - p_y^2}}\right) \right\} \\ p_y^f &= p_y \\ \delta^f &= \delta \\ \tau^f &= \tau + \frac{(1+\delta)s}{b_0 \rho_c} + \frac{(1+\delta)}{b_0} \left\{ \arcsin\left(\frac{p_x}{\sqrt{(1+\delta)^2 - p_y^2}}\right) - \arcsin\left(\frac{p_x^f}{\sqrt{(1+\delta)^2 - p_y^2}}\right) \right\} \end{split}$

Transfer map for one step: $M(\tau) = M_2(\tau/2)M_1(\tau)M_2(\tau/2) + O(\tau^3)$

Ref: E. Forest, Beam Dynamics: A New Attitude and Framework, 1998.

Different Boundary/Beam Conditions Need Different Efficient Numerical Algorithms O(Nlog(N)) or O(N)

FFT based Green function method:

- Standard Green function: low aspect ratio beam
- Shifted Green function: separated particle and field domain
- Integrated Green function: large aspect ratio beam
- Non-uniform grid Green function: 2D radial non-uniform beam

Fully open boundary conditions

Spectral-finite difference method:

Partially open boundary Transverse regular pipe with longitudinal open

Multigrid spectral-finite difference method:

Transverse irregular pipe

J. Qiang, S. Paret, "Poisson solvers for self-consistent multiparticle simulations," ICFA Mini-Workshop on Beam-Beam Effects in Hadron Colliders, March 18-22, 2013.

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6

Green Function Solution of Poisson's Equation (I) (open boundary conditions)

$$\phi(r) = \int G(r, r') \rho(r') dr' \quad ; \ r = (x, y, z)$$

$$f(r_i) = h \overset{N}{\underset{i'=1}{\circ}} G(r_i - r_{i'}) \varGamma(r_{i'})$$

$$G(x, y, z) = 1/\sqrt{(x^2 + y^2 + z^2)}$$

Direct summation of the convolution scales as N² !!!! N – total number of grid points

FFT based Hockney's Algorithm /zero padding:- scales as (2N)log(2N)

- Ref: Hockney and Easwood, Computer Simulation using Particles, McGraw-Hill Book Company, New York, 1985.

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$$f_c(r_i) = h \overset{\text{\tiny 2N}}{=} G_c(r_i - r_{i'}) \Gamma_c(r_{i'})$$

$$f(r_i) = f_c(r_i) \text{ for } i = 1, N$$

This is NOT a spectral solver! AWRENCE BERKELEY NATIONAL LABORATORY Integrated Green Function Method (II)

(large aspect ratio beam with open boundary conditions)

$$f_c(r_i) = \mathop{\stackrel{2N}{\stackrel{}_{a'=1}}}_{i'=1}^{2N} G_i(r_i - r_i') \Gamma_c(r_i')$$

$$r') = \mathop{\check{}}_{a'} G_s(r, r') dr'$$

$$G_s(x, y, z)$$

 $G_s(x, y, z) = 1/\sqrt{(x^2 + y^2 + z^2)}$

rrrr

integrated Green function

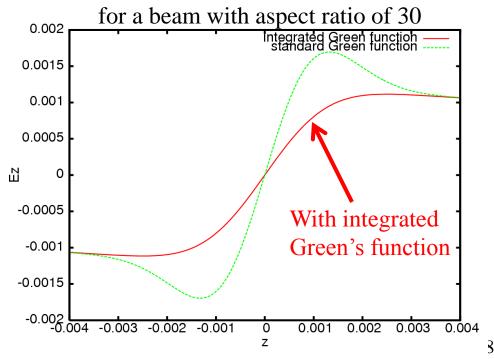
standard Green function

Integrated Green's function is needed for modeling large aspect ratio beams!

 $G_i(r)$

$$(O(N \log N))$$

R. D. Ryne, ICFA Beam DynamicsMini
Workshop on Space Charge Simulation,
Trinity College, Oxford, 2003
J. Qiang, S. Lidia, R. D. Ryne, and C. Limborg-Deprey,
Phys. Rev. ST Accel. Beams, vol 9, 044204 (2006).



Comparison between the IG and SG

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3D Poisson Solver with Transverse Rectangular Pipe (f) (Spectral-Finite Difference Method)

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial z^2} = -\frac{\rho}{\epsilon_0}$$

with boundary conditions

$$\phi(x = 0, y, z) = 0,$$

$$\phi(x = a, y, z) = 0,$$

$$\phi(x, y = 0, z) = 0,$$

$$\phi(x, y = b, z) = 0,$$

$$\phi(x, y, z = \pm \infty) = 0,$$

$$\rho(x, y, z) = \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \rho^{lm}(z) \sin(\alpha_l x) \sin(\beta_m y),$$

$$\phi(x, y, z) = \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \phi^{lm}(z) \sin(\alpha_l x) \sin(\beta_m y),$$

where

$$\rho^{lm}(z) = \frac{4}{ab} \int_{0}^{a} \int_{0}^{b} \rho(x, y, z) \sin(\alpha_l x) \sin(\beta_m y),$$

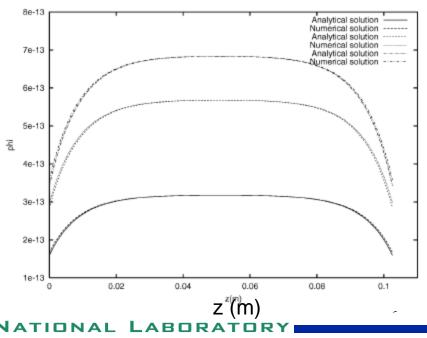
$$\phi^{lm}(z) = \frac{4}{ab} \int_{0}^{a} \int_{0}^{b} \phi(x, y, z) \sin(\alpha_l x) \sin(\beta_m y),$$

J. Qiang, and R. Ryne, Comput. Phys. Comm. 138, p. 18 (2001).

BFRKFIF

 $\frac{\partial^2 \phi^{lm}(z)}{\partial z^2} - \gamma_{lm}^2 \phi^{lm}(z) = -\frac{\rho^{lm}(z)}{\epsilon_0},$ $\frac{\phi_{n+1}^{lm} - 2\phi_n^{lm} + \phi_{n-1}^{lm}}{h_z^2} - \gamma_{lm}^2 \phi_n^{lm} = -\frac{\rho_n^{lm}}{\epsilon_0},$ $\phi_{-1}^{lm} = \exp(-\gamma_{lm}h_z)\phi_0^{lm}, \quad n = 0,$ $\phi_{N+1}^{lm} = \exp(-\gamma_{lm}h_z)\phi_N^{lm}, \quad n = N.$

Numerical Solutions vs. Analytical Solutions



3D Poisson Solver with Transverse Rectangular Pipe (ff) (Spectral-Green Function Method and 3D Spectral Method)

$$\frac{d\phi^{lm}(z)}{\partial z^{2}} - \gamma_{lm}^{2} \phi^{lm}(z) = -\frac{\rho^{lm}(z)}{\epsilon_{0}},$$

$$\begin{pmatrix} \phi^{lm}(z) = \frac{-1}{2\gamma_{lm}\varepsilon_{0}} \int \exp(-\gamma_{lm} |z-z'|)\rho^{lm}(z')dz' \\ \phi^{lm}(z) = \frac{-1}{2\gamma_{lm}\varepsilon_{0}} \sum_{j} \rho^{lm}(z'_{j})G(z-z'_{j}) \\ G(z-z') = \int_{z'-h/2}^{z'+h/2} \exp(-\gamma_{lm} |z-t|)dt \\ \phi^{lm}(z) = \sum_{n} \phi_{n}^{lm}H_{n}(z); \ \rho^{lm}(z) = \sum_{n} \rho_{n}^{lm}H_{n}(z) \\ \text{Hermite polynomial } H_{n}: \int_{-\infty}^{\infty} H_{n}(z)H_{m}(z)dz = 2^{n}n!\sqrt{\pi}\delta_{mn} \\ \frac{1}{4}\phi_{n-2}^{lm} - (\frac{1}{2}(2n+1) + \gamma_{lm}^{2})\phi_{n}^{lm} + (n+2)(n+1)\phi_{n+2}^{lm} = A^{2}\rho_{n}^{lm} \end{pmatrix}$$

10

Green function method: - efficiently handle long bunch

 ∂^2

spectral method: -reduce numerical noise with filtering in frequency domain

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3D Poisson Solver with Transverse Rectangular Pipe (Green Function Method)

$$\begin{split} G(x, x', y, y', z, z') &= \frac{1}{2\pi ab} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \frac{1}{\kappa_{mn}} \sin \frac{m\pi x}{a} \sin \frac{m\pi x'}{a} \sin \frac{n\pi y}{b} \sin \frac{n\pi y'}{b} e^{-\kappa_{mn}|z-z'|}, \\ G &= R(x - x', y - y', z - z') - R(x - x', y + y', z - z') \\ -R(x + x', y - y', z - z') + R(x + x', y + y', z - z'), \\ R(u, v, w) &= \frac{1}{2\pi ab} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \frac{1}{\kappa_{mn}} \cos \frac{m\pi u}{a} \cos \frac{n\pi v}{b} e^{-\kappa_{mn}|w|}. \\ \phi_{i,j,k}/(h_x h_y h_z) &= \mathcal{F}^{bbb} \{(\mathcal{F}^{fff} \rho_{i,j,k})(\mathcal{F}^{fff} R_{i,j,k})\} - \mathcal{F}^{bfb} \{(\mathcal{F}^{fff} \rho_{i,j,k})(\mathcal{F}^{bf} R_{i,j,k})\} \\ -\mathcal{F}^{fbb} \{(\mathcal{F}^{fff} \rho_{i,j,k})(\mathcal{F}^{bff} R_{i,j,k})\} + \mathcal{F}^{ffb} \{(\mathcal{F}^{fff} \rho_{i,j,k})(\mathcal{F}^{bbf} R_{i,j,k})\} \\ g_z &= \frac{1}{h_z^2} \int_{z_k}^{z_{k+1}} dz' \left[\rho_k(h_z - (z' - z_k)) + \rho_{k+1} (z' - z_k)\right] e^{-\kappa_{mn}|z-z'|}. \\ R_{int}(u, v, w) &= \frac{1}{2\pi ab} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \frac{1}{\kappa_{mn}} \cos \frac{m\pi u}{a} \cos \frac{n\pi v}{b} g_z(w). \\ g_z(w) &= \frac{1}{h_z^2 \kappa_{mn}^2} \left[2h_z \kappa_{mn} \delta_{w,0} + \left(e^{-\kappa_{mn}|w+h_z|} - 2e^{-\kappa_{mn}|w|} + e^{-\kappa_{mn}|w-h_z|} \right) \right], \\ Pro: computational domain needs to contain only the beam itself \\ \end{array}$$

Con: more numerical operations for Green function evaluation. R. Ryne, arXiv:1111, 4971, 2011 LAWRENCE BERKELEY NATIONAL LABORATORY



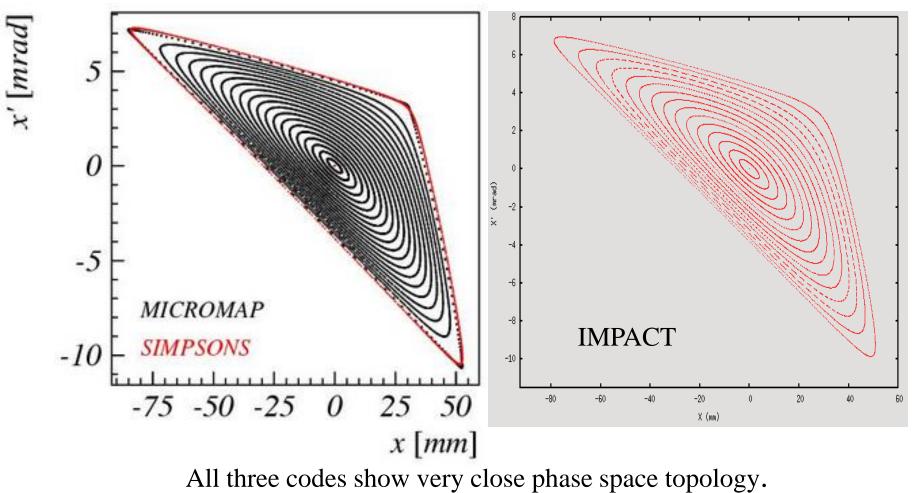
PARAMETERS for setps 1-6

| Lattice SIS18 | |
|-------------------------------------|--------------------------------------|
| Strength of the sextupole (when use | ed) $K_2 = 0.2 \text{ m}^{-2}$ |
| Maximum tuneshift | $\Delta Q_x = 0.1$ |
| Horizontal transverse size (rms) | $X_{rms} = 5 mm$ |
| Vertical transverse size (rms) | $Y_{rms} = 5 mm$ |
| Longitudinal size (rms) | $Z_{rms} = 40.35 \text{ m}$ |
| Horizontal emittance (2 o) | $\epsilon_x = 12.57 \text{ mm mrad}$ |
| Vertical emittance (2 o) | $\epsilon_y = 9.30 \text{ mm mrad}$ |
| Turns for 1 synchrotron oscillation | N _{synch} =15000 |
| Bunch length (4 σ_z) | $\tau = 3472.7 \text{ ns}$ |
| Kinetik energy | $E_k = 11.4 \text{ MeV/u}$ |
| Gamma transition | $\gamma_t = 5$ |
| momentum spread at $3\sigma_z$ | $\Delta p/p = 2.5 \times 10^{-4}$ |

http://web-docs.gsi.de/~giuliano/



Qx = 4.338, Qy = 3.2



13



PARAMETERS for setps 7-8

| Lattice SIS18 | |
|-------------------------------------|--------------------------------------|
| Strength of the sextupole (when use | ed) $K_2 = 0.2 \text{ m}^{-2}$ |
| Maximum tuneshift | $\Delta Q_x = 0.1$ |
| Horizontal transverse size (rms) | $X_{rms} = 5 mm$ |
| Vertical transverse size (rms) | $Y_{rms} = 5 mm$ |
| Longitudinal size (rms) | $Z_{ms} = 2.69 \text{ m}$ |
| Horizontal emittance (2 o) | $\epsilon_x = 12.57 \text{ mm mrad}$ |
| Vertical emittance (2 o) | $\epsilon_y = 9.30 \text{ mm mrad}$ |
| Turns for 1 synchrotron oscillation | N _{synch} =1000 |
| Bunch length $(4 \sigma_z)$ | $\tau = 231.51 \text{ns}$ |
| Kinetik energy | $E_k = 11.4 \text{ MeV/u}$ |
| Gamma transition | $\gamma_t = 5$ |
| momentum spread at $3\sigma_z$ | $\Delta p/p = 2.5 \times 10^{-4}$ |

Parameters for step 9:

- -Qx = 4.3604, Qy = 3.2
- RF cavity voltage: 152 V
- RF frequency: 0.214456 MHz
- X aperture : 40 cm
- Y aperture: 30 cm

Generate Initial Matched Beam Distribution

(Normal Form Transformation)



For a given particle phase space coordinate: $\xi = (x, p_x, y, p_y, z, p_z)^T$

after one turn through the ring $\xi^{n+1} = M\xi^n$ where *M* is the one-turn transfer matrix

Define a second order moments matrix of the distribution:

 $\Sigma = \begin{pmatrix} \langle xx \rangle & \langle xp_x \rangle & \langle xy \rangle & \langle xp_y \rangle & \langle xz \rangle & \langle xp_z \rangle \\ \langle p_x x \rangle & \langle p_x p_x \rangle & \langle p_x y \rangle & \langle p_x p_y \rangle & \langle p_x z \rangle & \langle p_x p_z \rangle \\ \langle yx \rangle & \langle yp_x \rangle & \langle yy \rangle & \langle yp_y \rangle & \langle yz \rangle & \langle yp_z \rangle \\ \langle p_y x \rangle & \langle p_y p_x \rangle & \langle p_y y \rangle & \langle p_y p_y \rangle & \langle p_y z \rangle & \langle p_y p_z \rangle \\ \langle zx \rangle & \langle zp_x \rangle & \langle zy \rangle & \langle zp_y \rangle & \langle zz \rangle & \langle zp_z \rangle \\ \langle p_z x \rangle & \langle p_z p_x \rangle & \langle p_z y \rangle & \langle p_z p_y \rangle & \langle p_z z \rangle & \langle p_z p_z \rangle \end{pmatrix},$ $\Sigma^{n+1} = \mathcal{M}\Sigma^n \mathcal{M}^T$

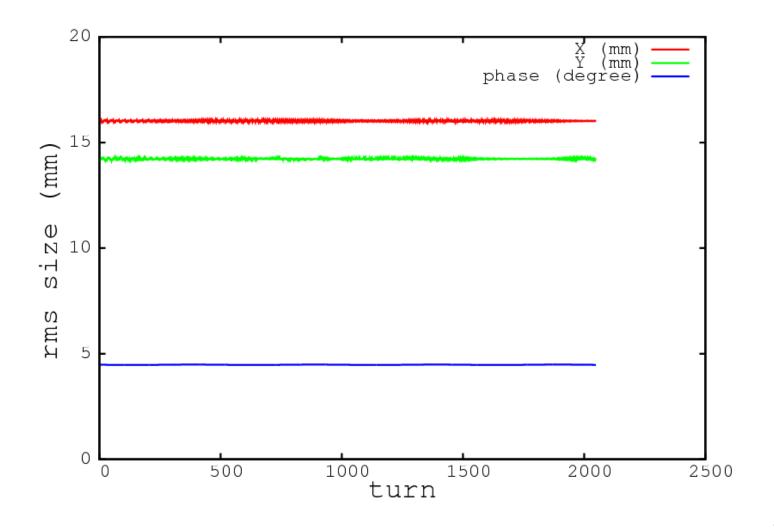
after one turn:

For an envelope matched beam: $\Sigma^n = M \Sigma^n M^T$ Let: $M = ANA^{-1}$ where N is the normal form matrix $(A^{-1}\Sigma^n A^{-T}) = N(A^{-1}\Sigma^n A^{-T})N^T$

For the given emittances in the phase space coordinate, one can find the corresponding emittances in the normal coordinates to generate a matched beam distribution.

15

Evolution of RMS Beam Sizes with the Initial Matched Distribution (no space-charge)

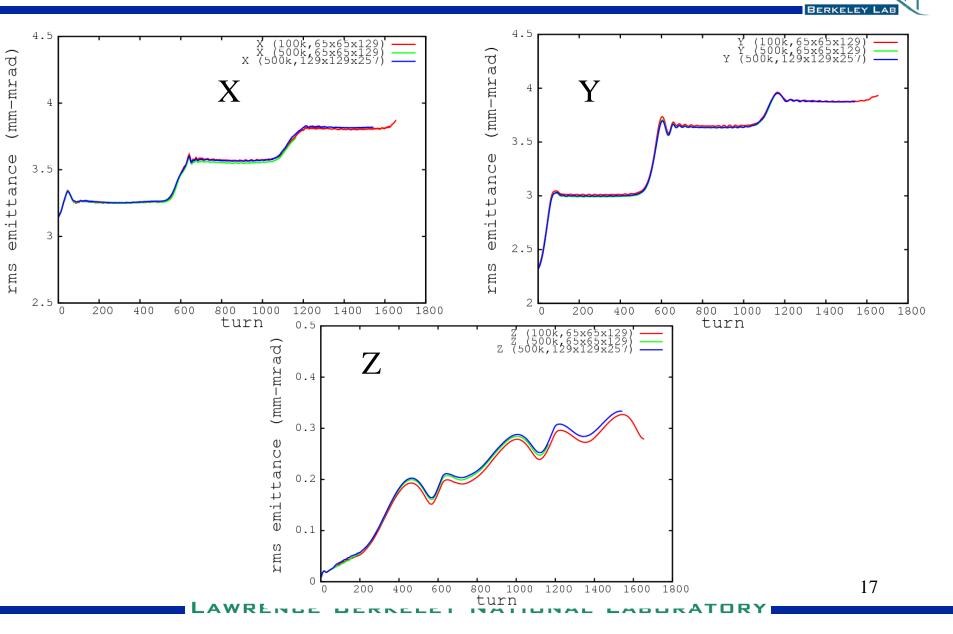


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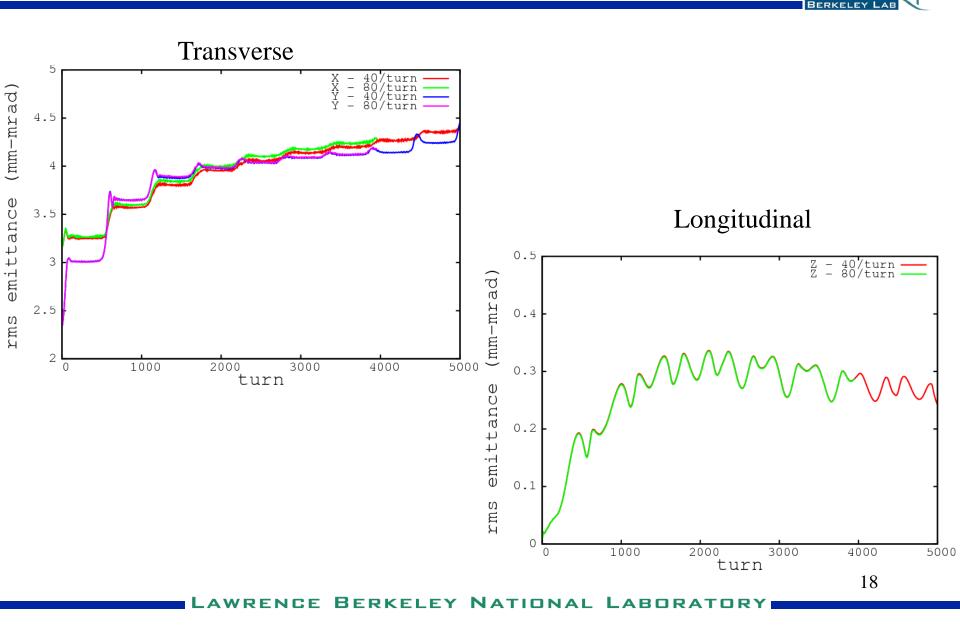
16

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Numerical Test of Convergence: Evolution of RMS Beam Emittances (1) (with space-charge)

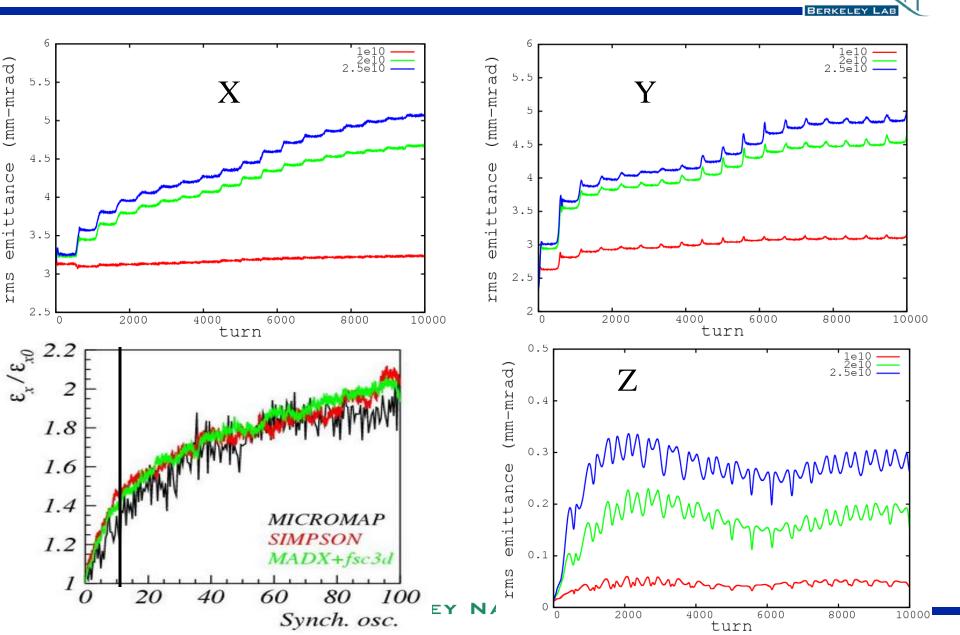


Numerical Test of Convergence: Evolution of RMS Beam Emittances (II) (with space-charge)



Evolution of RMS Beam Emittances: Different Beam Intensity

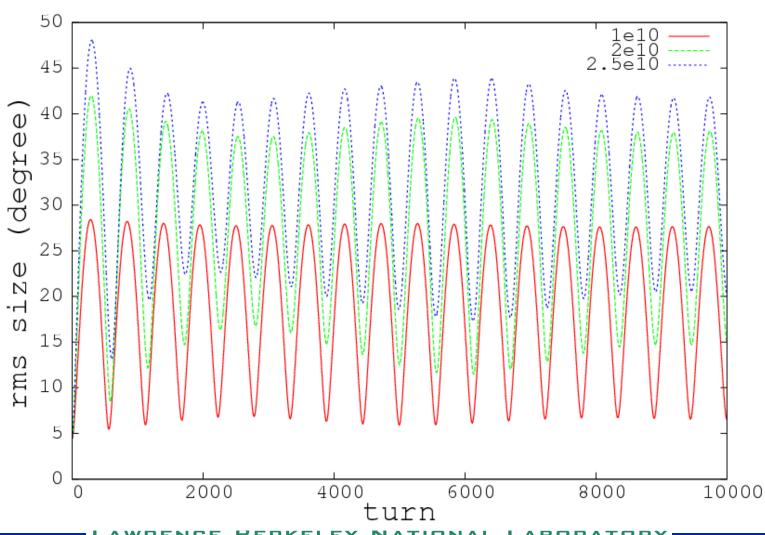
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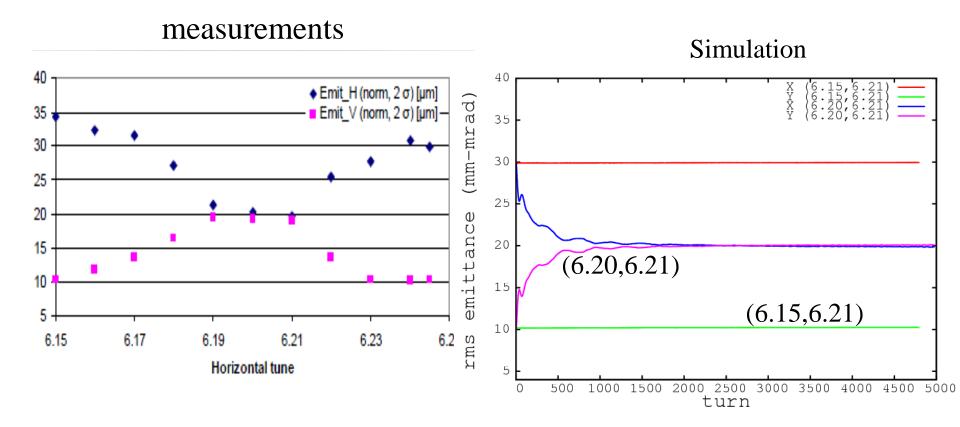
Physical parameters:

RF frequency = 3.5 MHzRF voltage = 27 kVEk = 1.4 GeVEmit_x = 7.5 mm-mradEmit_y = 2.5 mm-mradRms bunch length = 45 nsRms dp/p = 1.7×10^{-3}

Horizontal tune: 6.15 - 6.245Vertical tune: 6.21Synchrotron period: 1.5 ms Half Aperture = 7cm x 3.5cm $I = 1.0x10^{12}$

Refs: B. W. Montague, CERN-Report No. 68-38, CERN, 1968.
E. Metral et al., Proc. of EPAC 2004, p. 1894.
I. Hofmann et al., Proc. of EPAC 2004, p. 1960.

Evolution of Transverse RMS Emittances: Different Tunes



22

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Thank You for Your Attention!