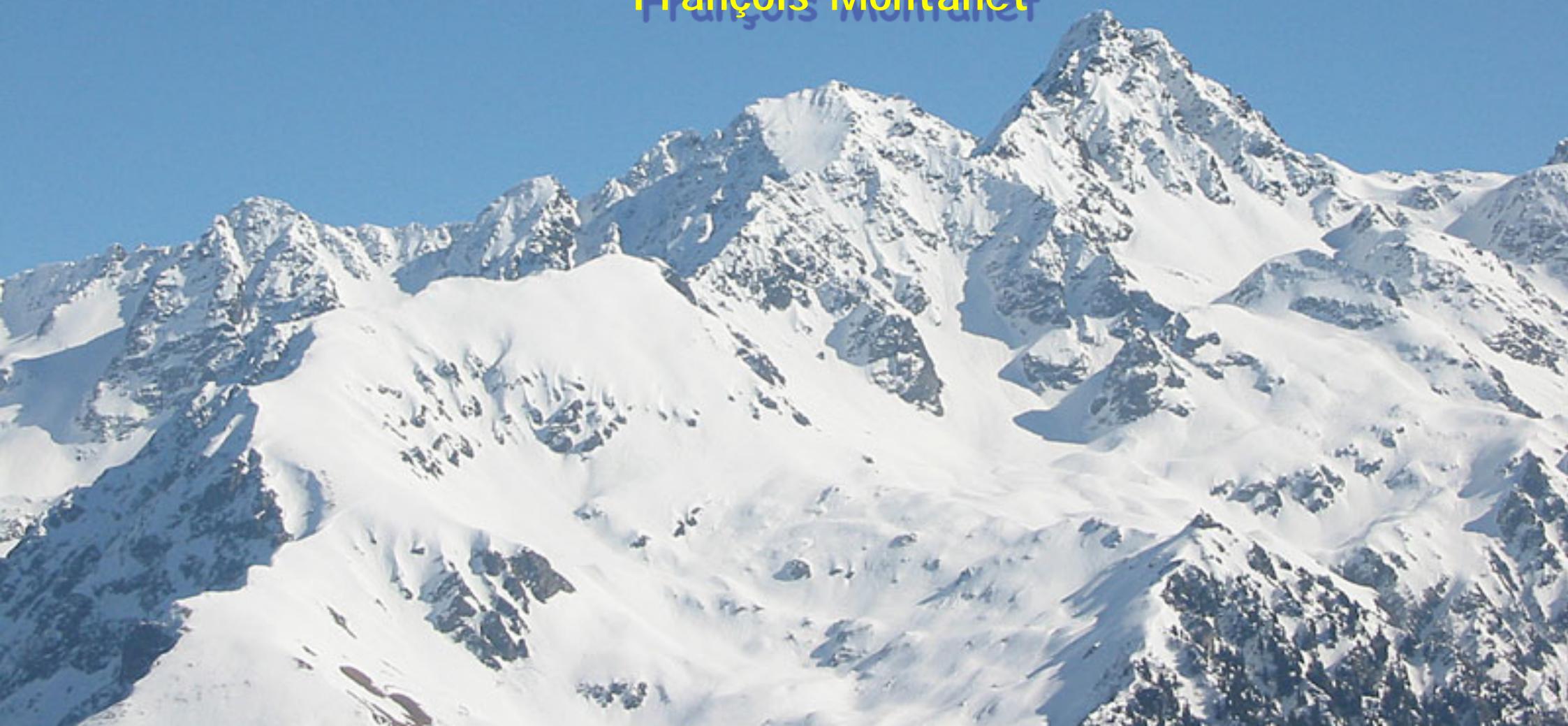


ASTROPARTICLES

ESIPAP – 2014
François Montanet



Plan of the course

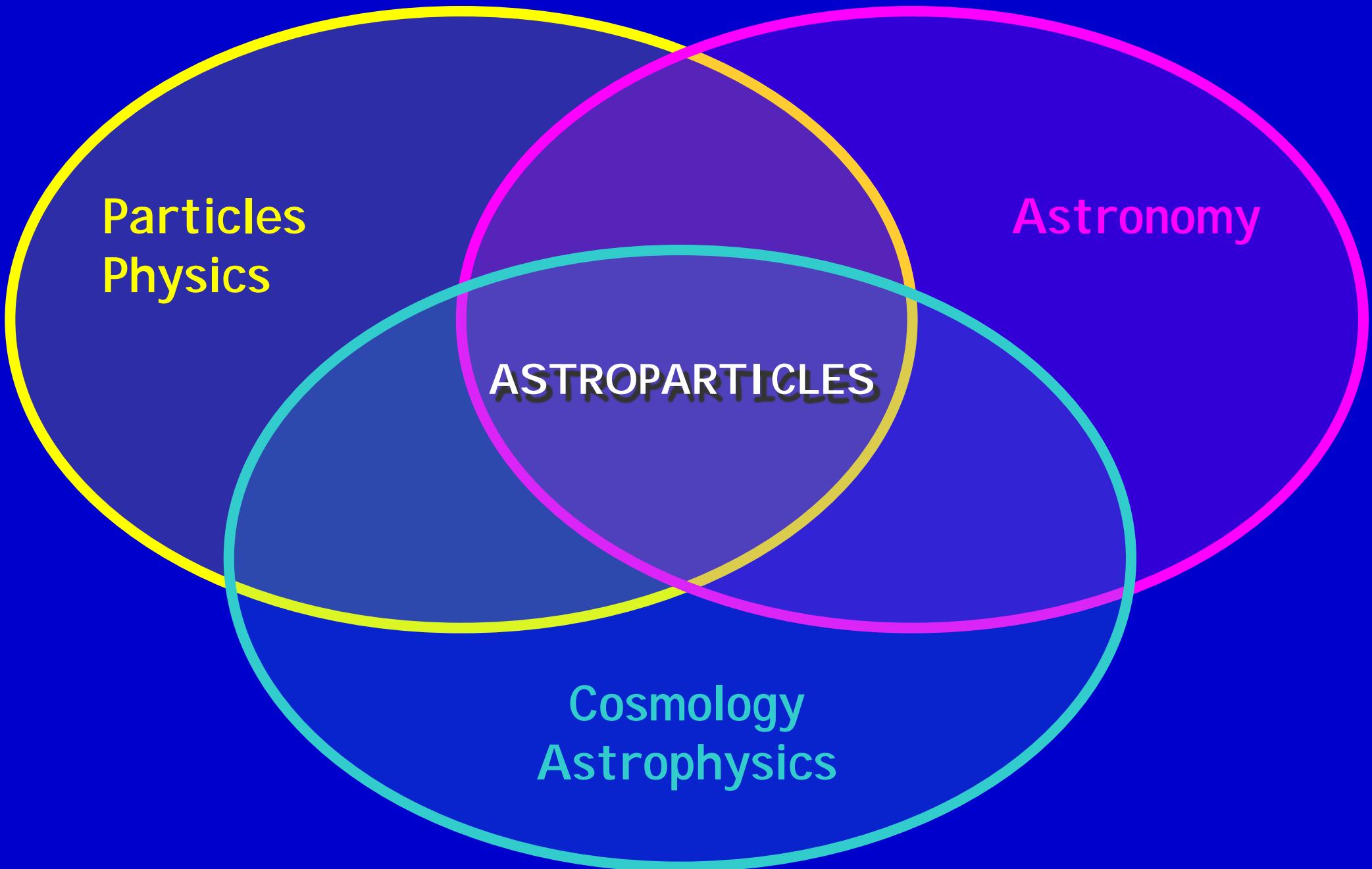
- Introduction & a very brief History
- Nature and properties of Astroparticles
- Propagation medium, IGM, ISM and atmosphere
- Astrophysical Sources
 - Astrophysical shocks
 - Fermi acceleration
 - Standard Model for the production of galactic CR, SNR
 - Gamma-ray sources, pulsars
 - AGN and other extragalactic sources
 - Neutrinos sources
- Cosmological Sources
 - Cosmological implications
 - Dark Matter as a source of CR
 - "top-down" type of sources at UHE
- Propagation
 - CR propagation in the Galaxy:
The Leaky box model
 - VHE g-rays propagation
 - UHECR propagation
- Observables & Observations
 - Primary CR detection
(on top of atmosphere)
 - Air Showers
Development Models
 - Gamma-ray (EM) induced showers
detection
 - Hadronic Showers
Models and Detection
 - UHECR detection
- Dark matter search
- Neutrino Physics with astroparticules

Why studying
Astroparticles

Open questions

Understand our Universe at extreme scales

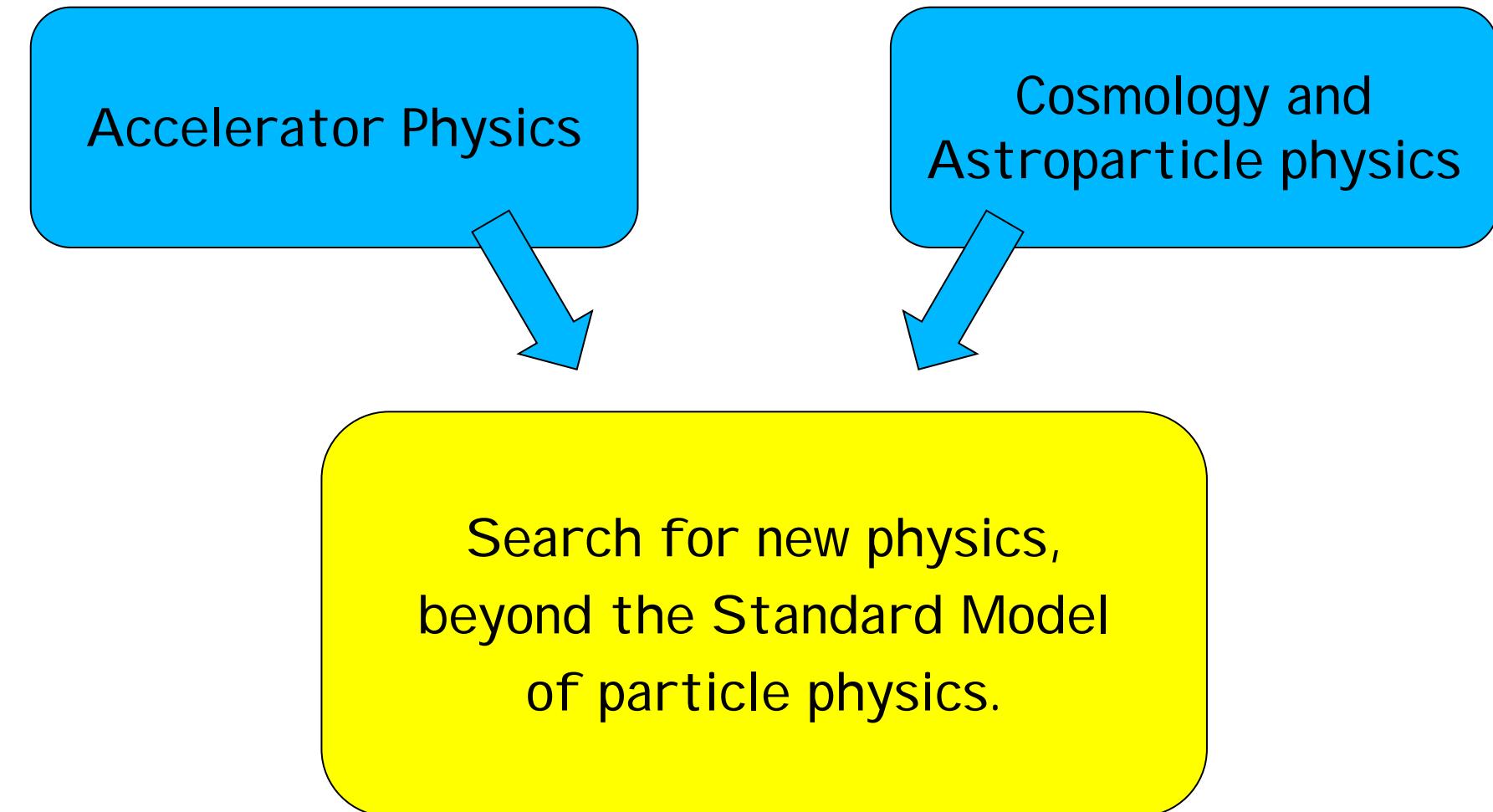
- The Higgs boson or the origin of mass
- Nature and mass of Neutrinos
- Fundamental symmetries: CP , supersymmetry
- New dimensions in physical space ?
- What is our Universe made of ?
- Sources and propagation of cosmic rays ?
- New Physics at $E \gg E(\text{LHC})$?



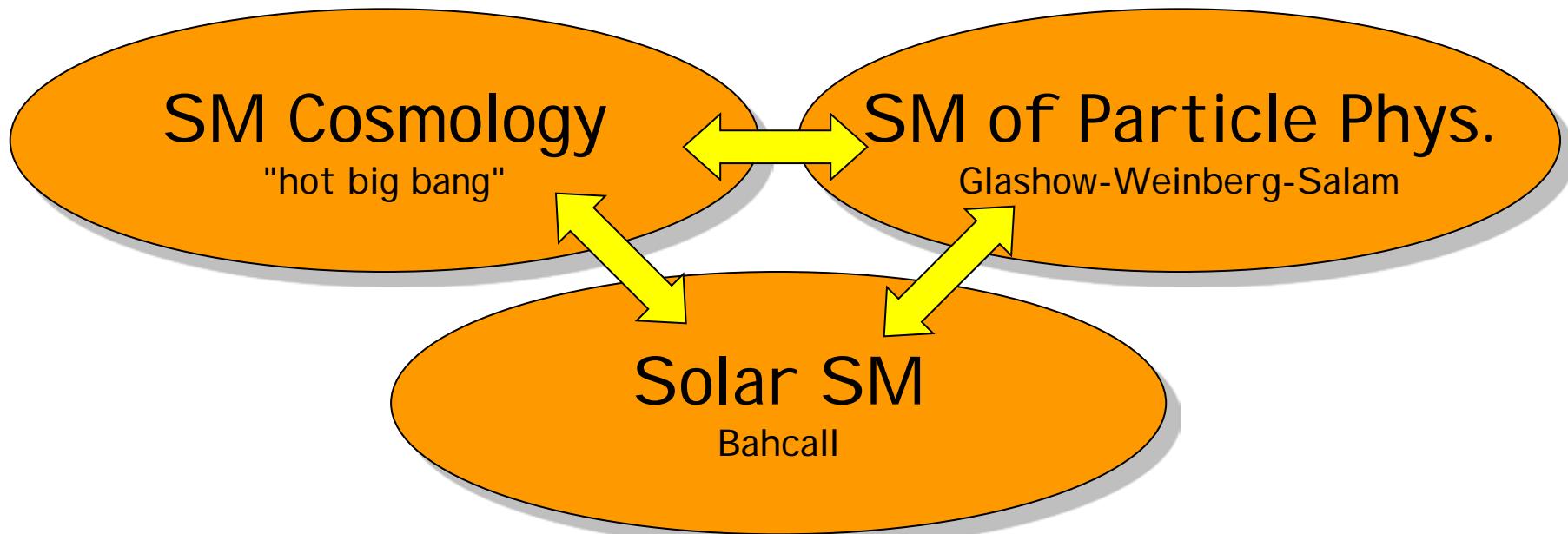
Astroparticles & HEP

2010-2011

F.Montanet Cours Astroparticules M2R PSA Grenoble



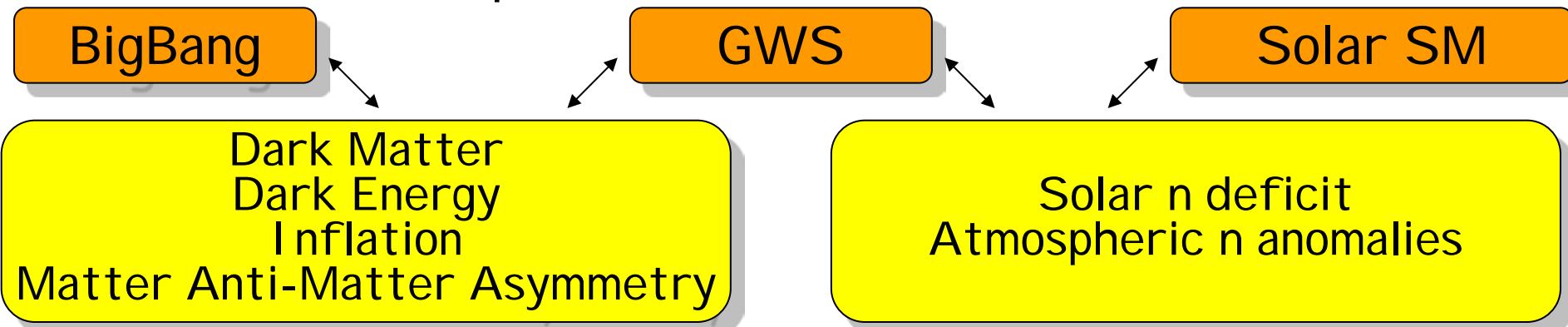
Matching "standard models"



Examples of happy unions...

Nucleosynthesis $\Rightarrow N_\nu$
 $\Omega h^2, LSS \dots \Rightarrow \sum m_\nu < \dots$

... as well as some disputes...



Direct searches



new particles production-observations
(Tevatron, LHC)

Indirect searches

$FCNC$, CP

$FV \rightarrow \mu \rightarrow e\gamma$

d_n^e

B physics

New Physics

This decade's grail:
solve the puzzle of
electroweak symmetry
breaking

Progress in Theory

Supergravity

→ Superstrings, M-Theory

Cosmology

Measure the parameters
of the Univers
and their evolution

Astroparticle Physics

Neutrino Physics

Cosmic Rays

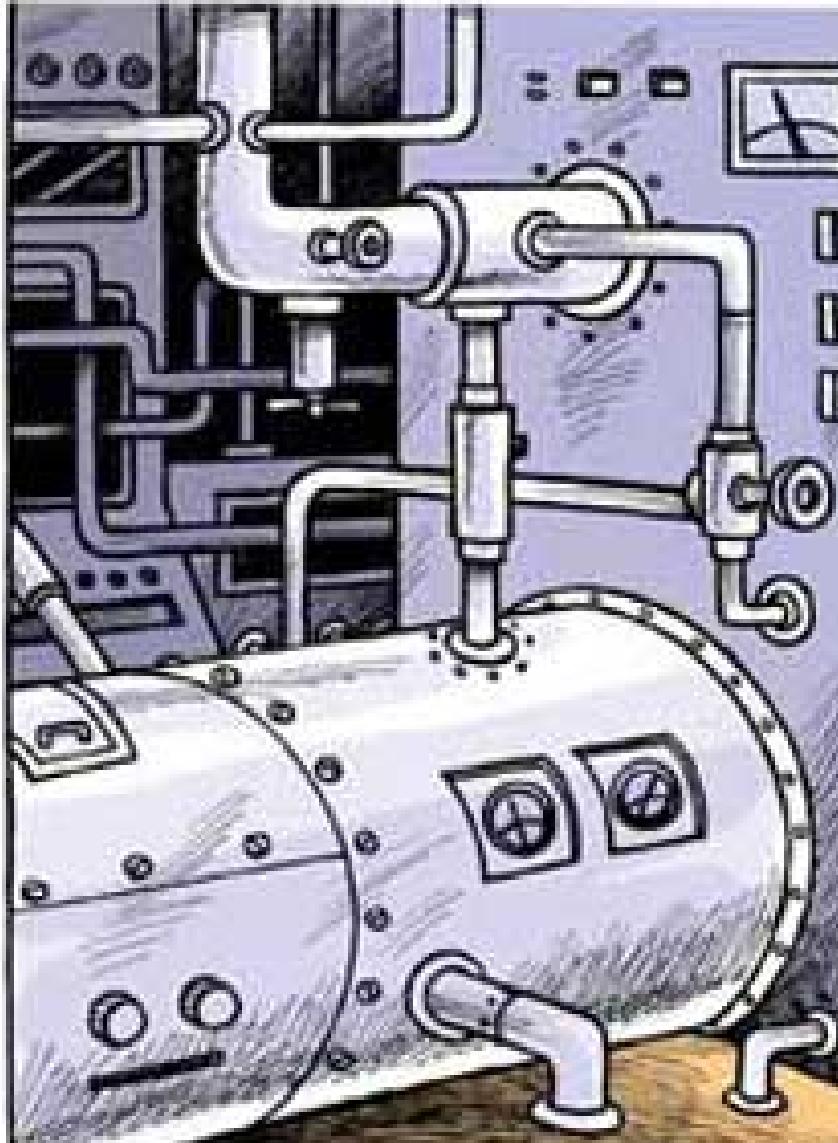
γ Astronomy

Gravitationnal Waves



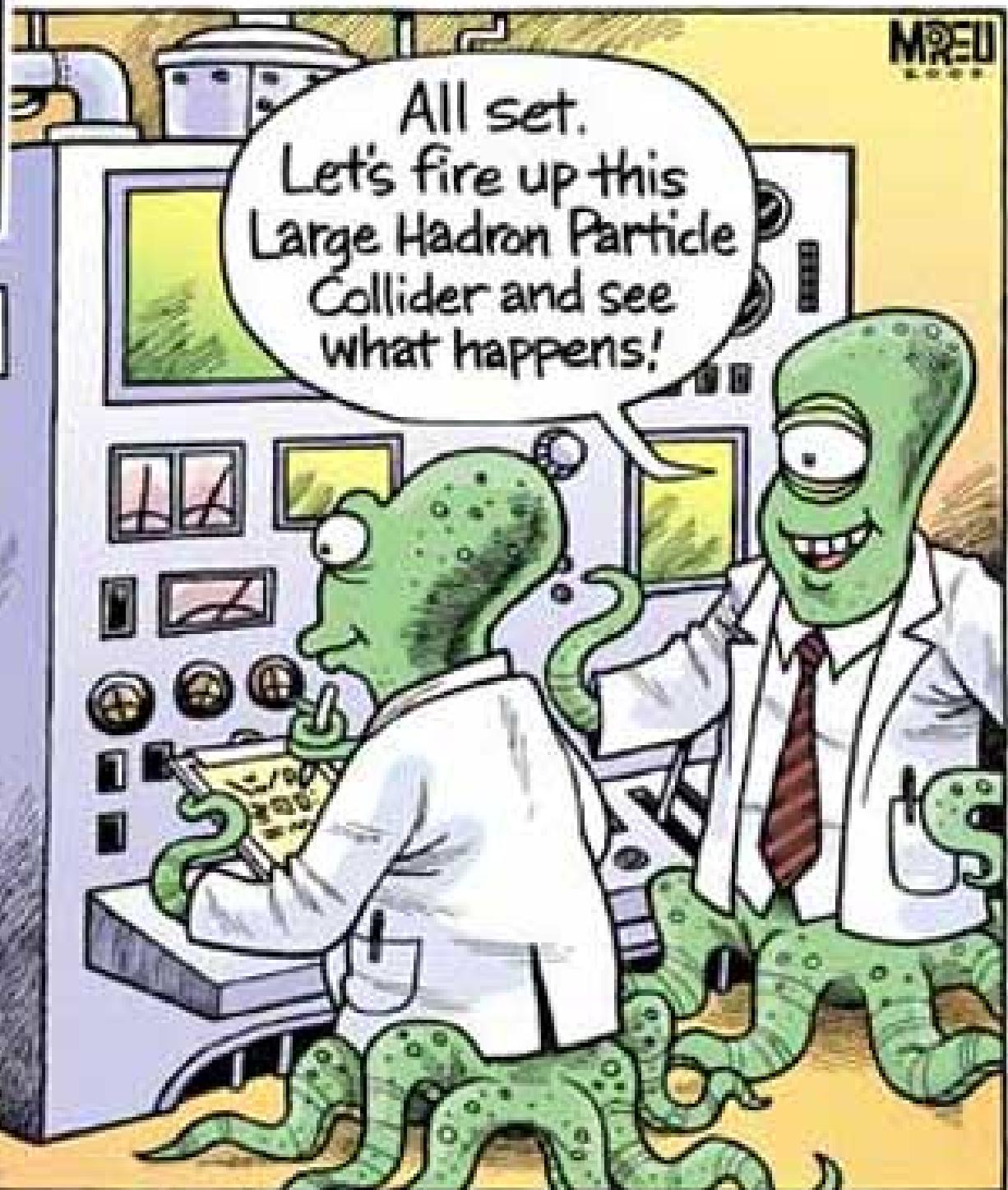
New Physics probes

13.8 BILLION YEARS AGO,
A FEW SECONDS BEFORE THE
CREATION OF OUR UNIVERSE...



MOREU

All set.
Let's fire up this
Large Hadron Particle
Collider and see
what happens!



What are "Astroparticles"

What we know

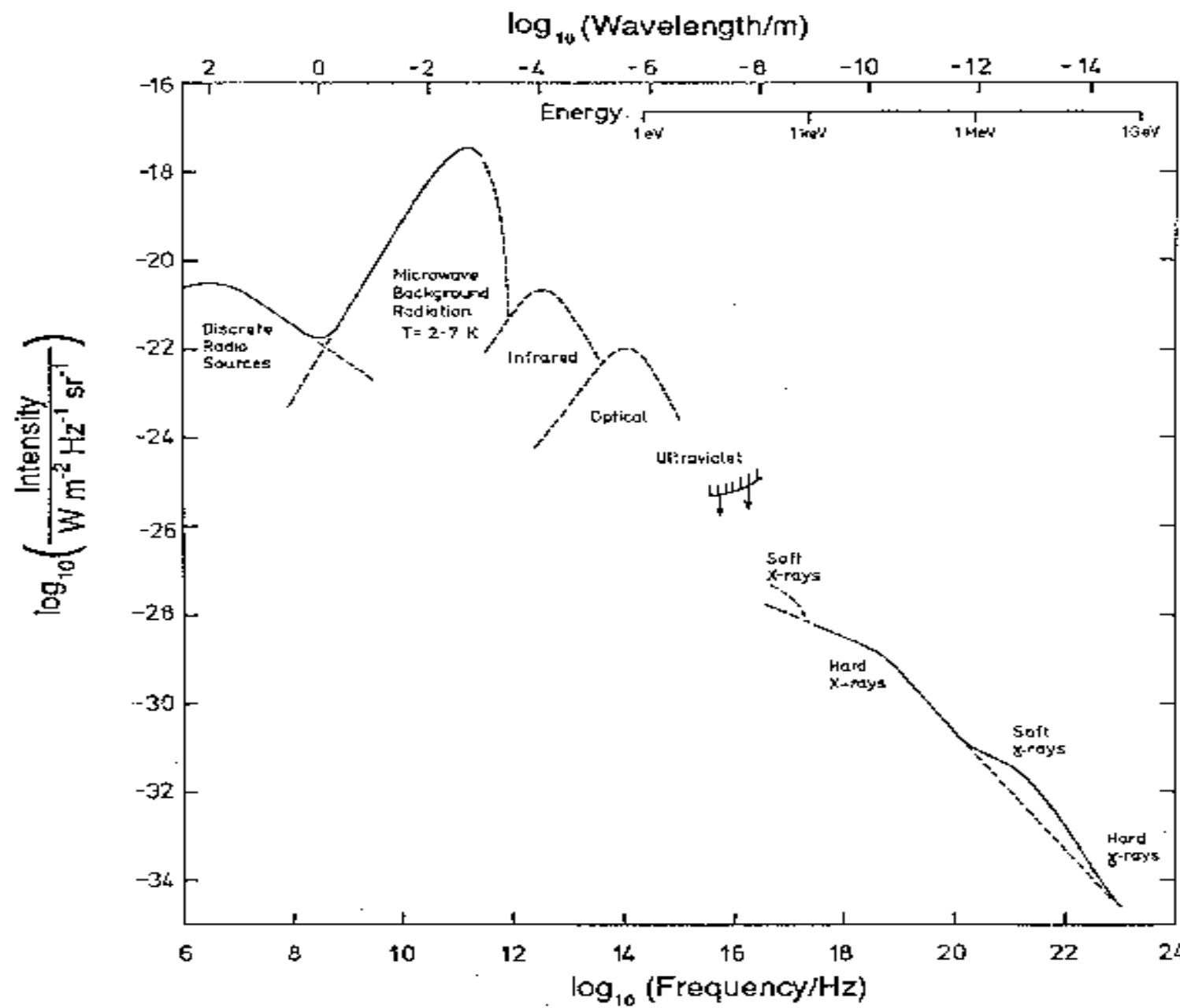
Let there be light !

All what we know in astrophysics is thanks to the light !

- Temperatures, stars masses, galaxies, magnetic fields, chemical composition, age of stars and structures...
- Nuclear reactions, galactic and extragalactic hydrodynamics, MHD, explosions, nucleosynthesis, past, future... EVERYTHING !

Well, almost everything...

A multi-wavelength sky



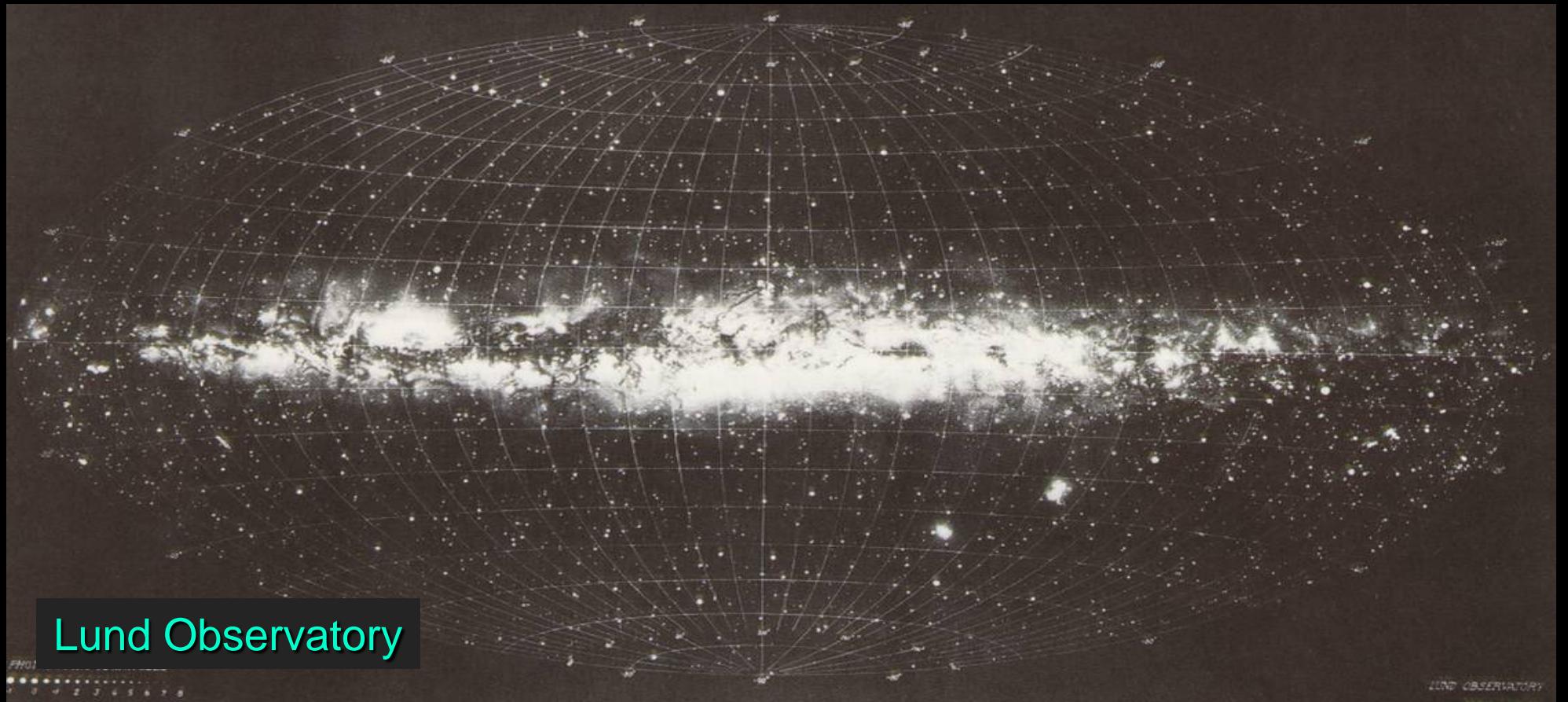
Our Galaxy

The optical Milky Way



Our Galaxy

The optical Milky Way



Lund Observatory

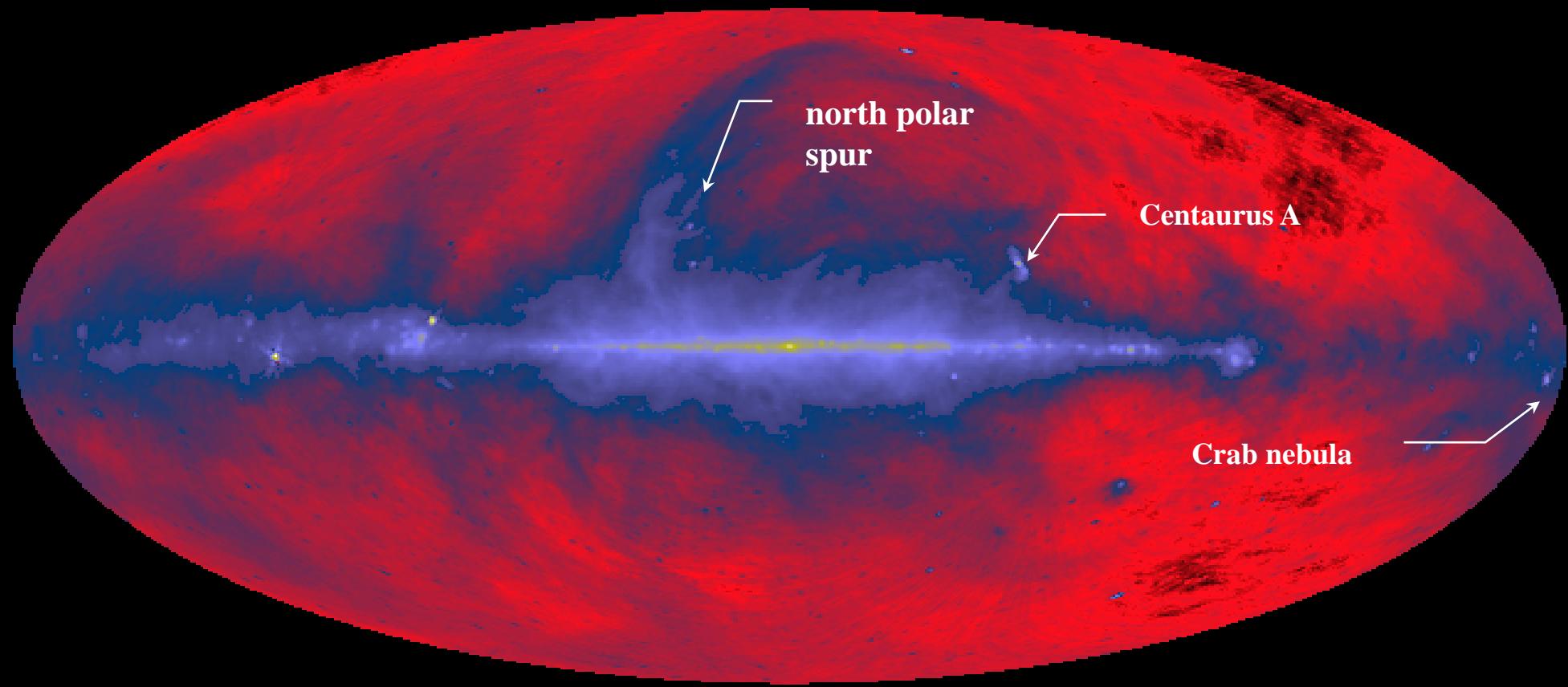
FNU1
0 1 2 3 4 5 6 7 8

LUND OBSERVATORY

Our Galaxy

The Milky Way : Radio at 73cm

$408 \text{ MHz} / 73.5 \text{ cm} / 1.6 \cdot 10^{-6} \text{ eV}$)

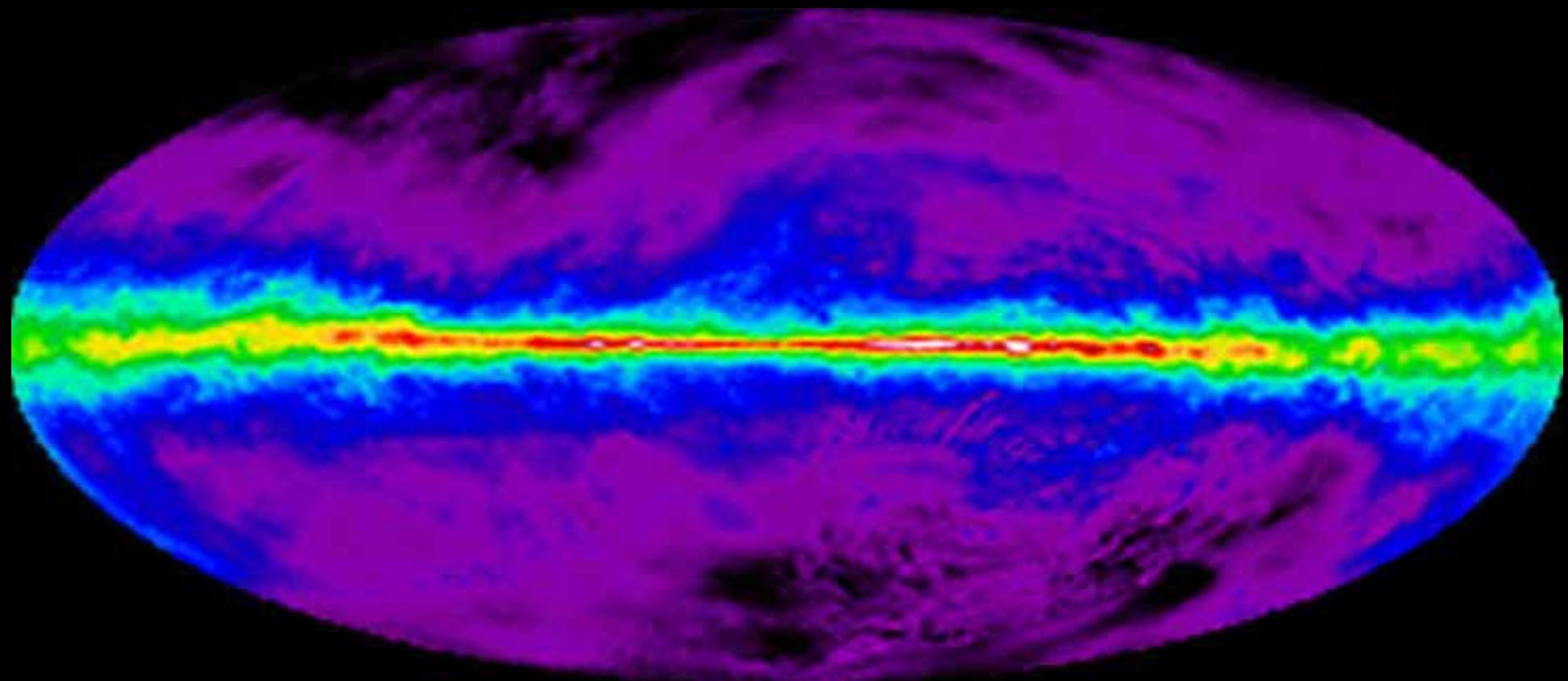


Essentially from the movement of ultra relativistic electrons probably issue from supernovae remnants in the galactic magnetic field.

Our Galaxy

The Milky Way : Radio at 21 cm

($\sim 1.42 \text{ GHz} / 21.1 \text{ cm} / 5.9 \cdot 10^{-6} \text{ eV}$)

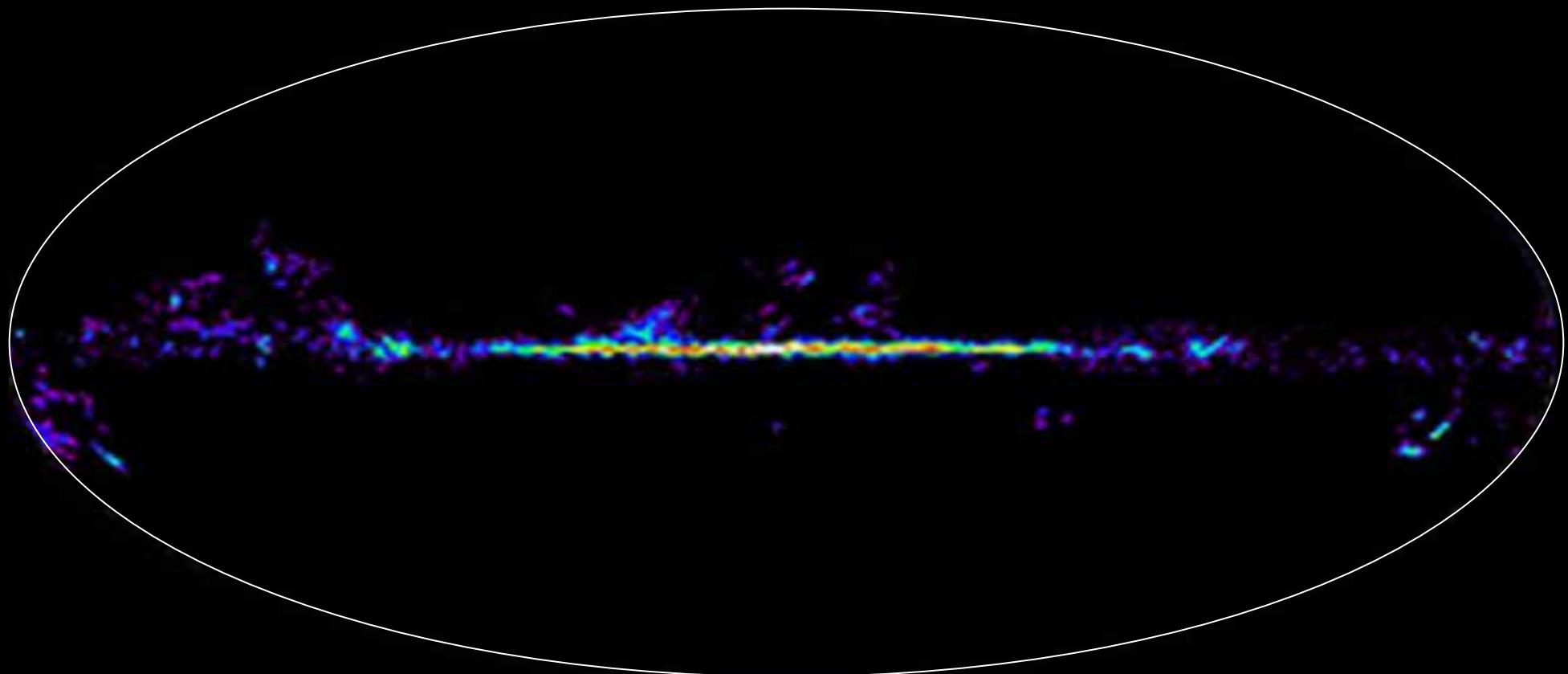


Hyperfine transition of hydrogen. Structures are due to the column density
of atomic hydrogen clouds along the line of sight.

Our Galaxy

The Milky Way : Radio at 2,6mm

Millimetric waves (115 GHz / 2.6 mm / $4.7 \cdot 10^{-4}$ eV)

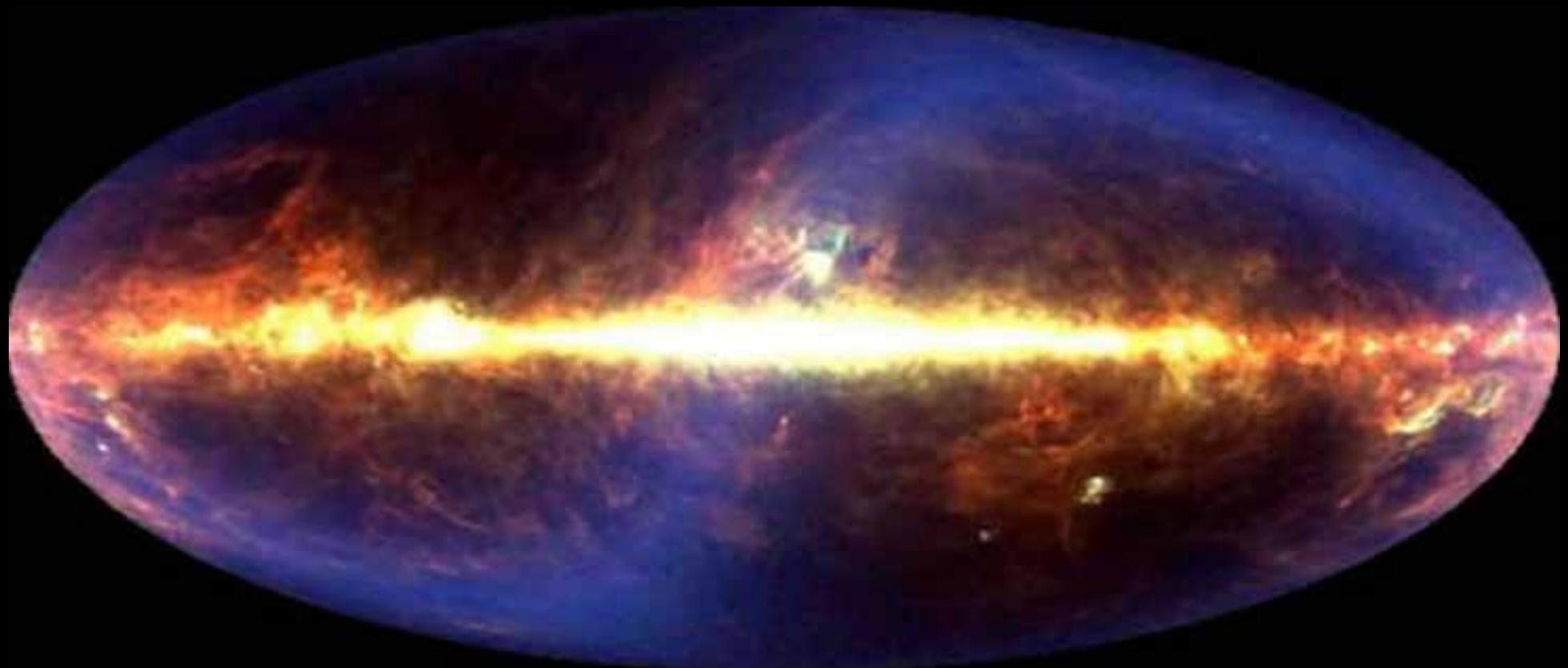


Rotation mode ray of carbon monoxide. One assumes that CO abundance is proportional to that of cold molecular hydrogen (directly undetectable).

Our Galaxy

Infra red

Infrared (3 10^3 to 25 10^3 GHz / 100 to 12 μm / 0.01 to 0.1 eV)



Thermal emission, due to interstellar dust heated by starlight.

Our Galaxy

Its structure is clearly visible in IR (COBE satellite).

Near Infrared (86 10³ à 240 10³ GHz / 1.25 à 3.5 μm / 0.35 à 1 eV).



Giant stars emission in the disk and in the bulb

Our Galaxy

Optical

Visible ($460 \cdot 10^3$ GHz / $0.65 \mu\text{m}$ / $2 \text{ eV} - \text{red}$)

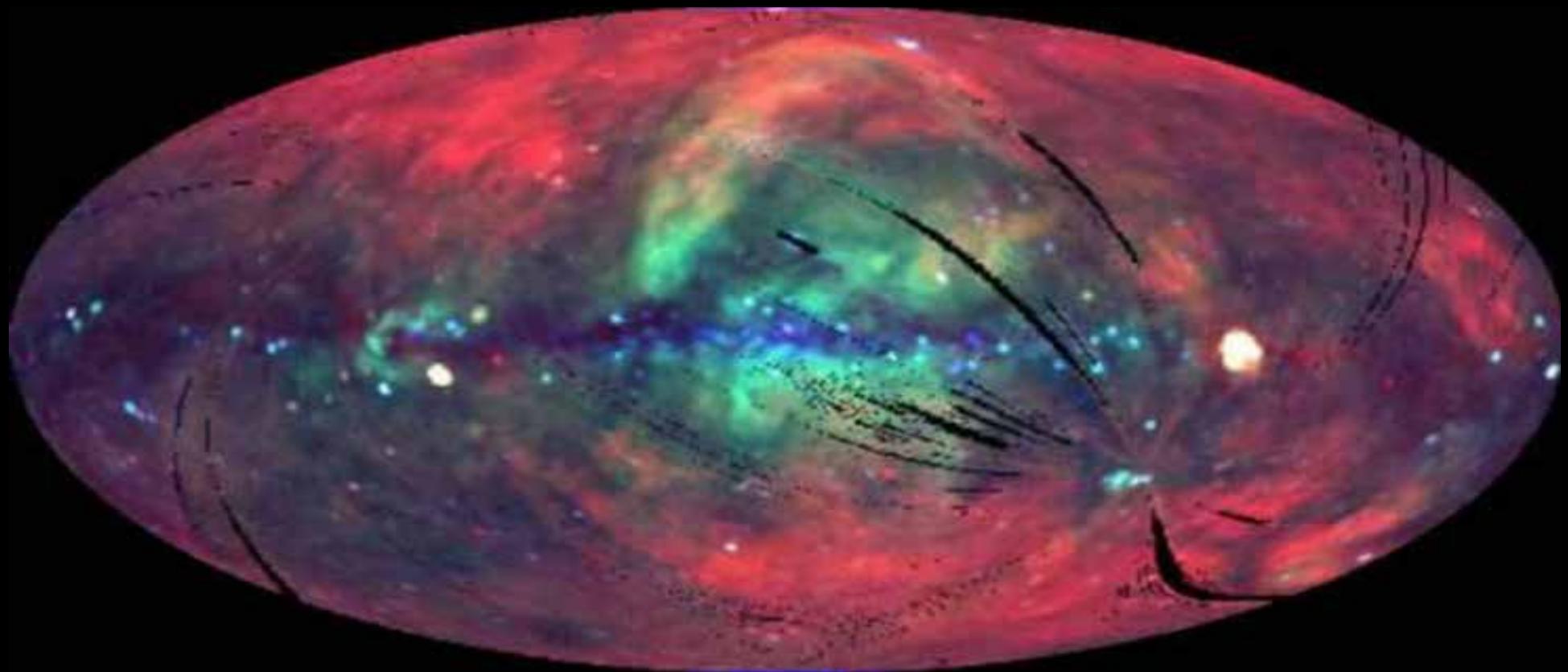


Visible light is absorbed by interstellar dust clouds.
Only stars close enough to the solar system (few parsec) are seen.

Our Galaxy

X-rays

X-rays (60.10⁶ to 360.10⁶ GHz / 5 to 8.3 nm / 0.25 to 1.5 keV).

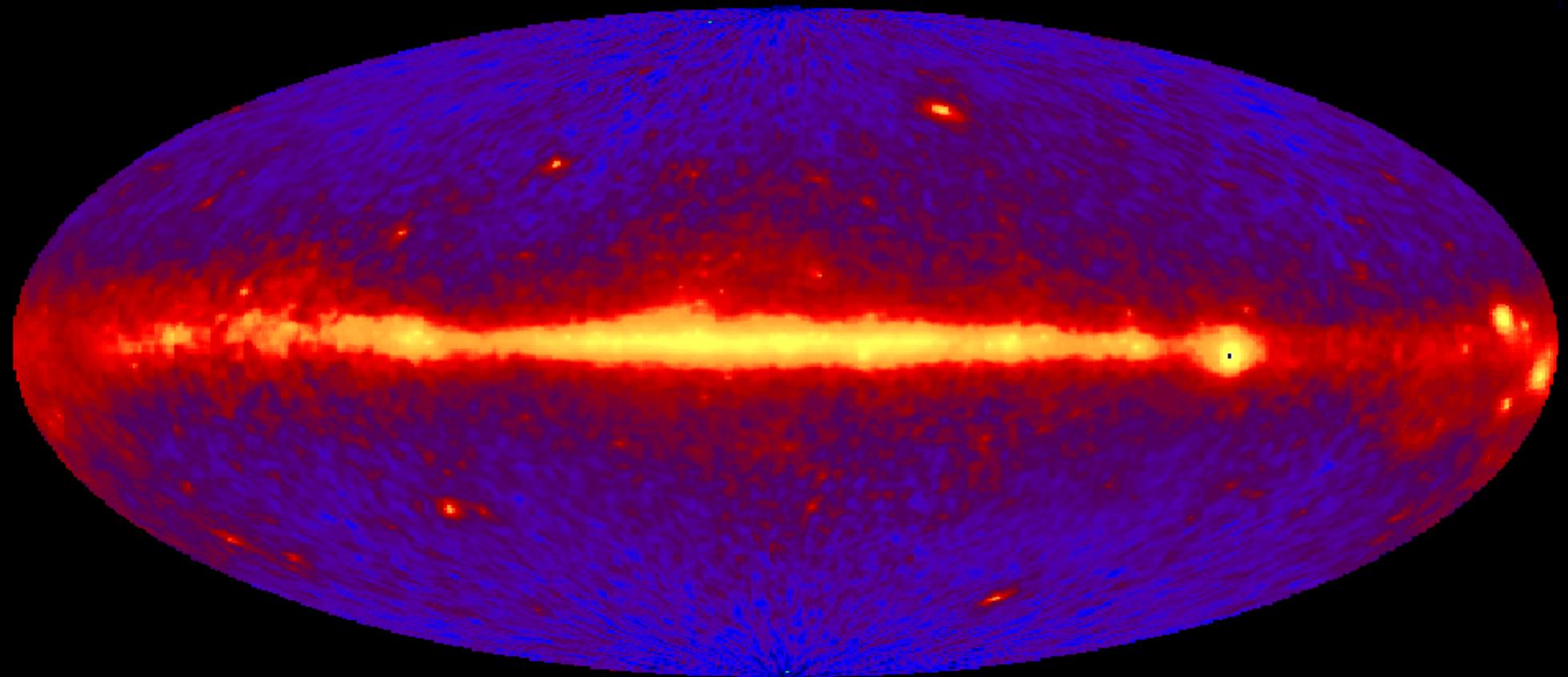


Diffuse X-ray emission from overheated and shocked gas.

Our Galaxy

Gamma-rays

Gamma-rays ($> 2.4 \cdot 10^{13} \text{ GHz}$ / $< 12.5 \text{ fm}$ / $> 100 \text{ MeV}$).

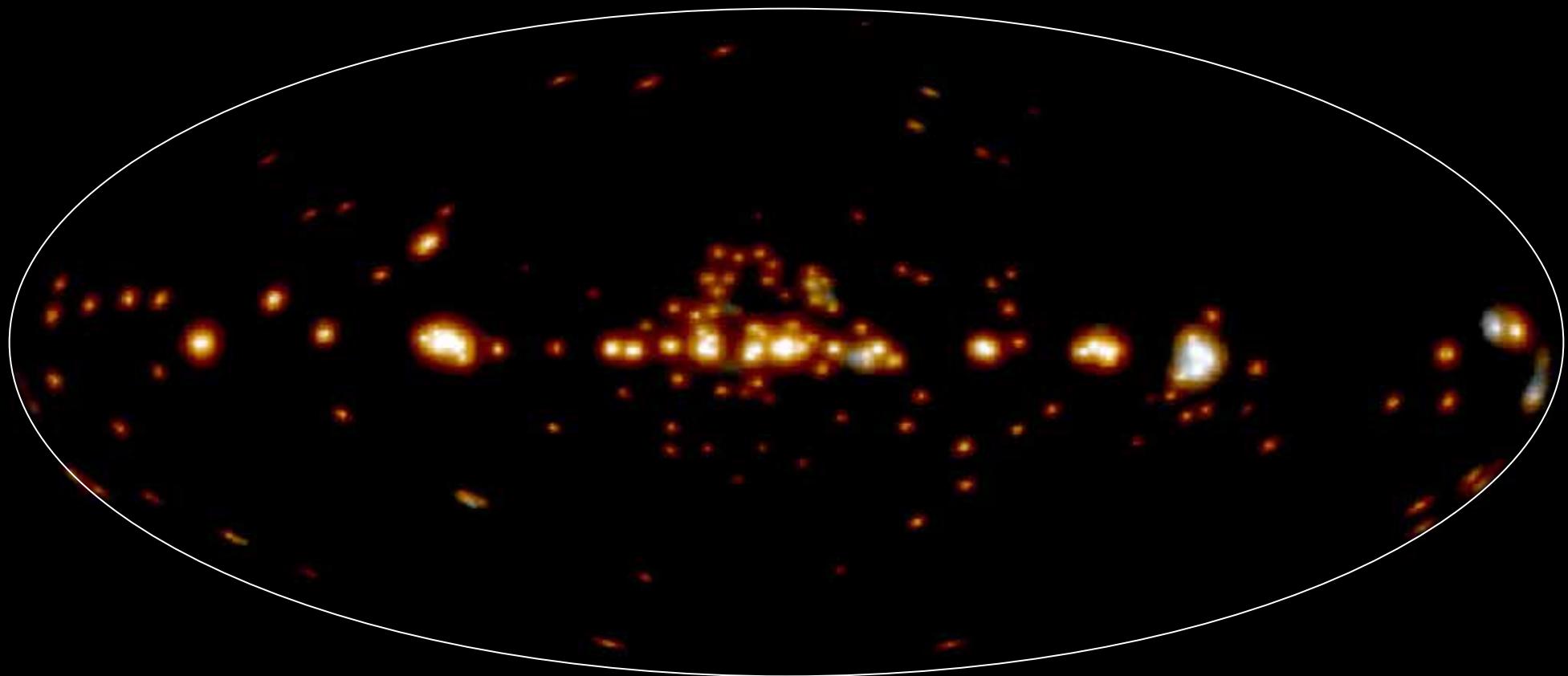


Photons (gammas) from the decay of neutral pions produced in the interaction of CR with interstellar matter, from the Bremsstrahlung of CR and from the inverse Compton of relativistic electrons with ambient photons.

Our Galaxy

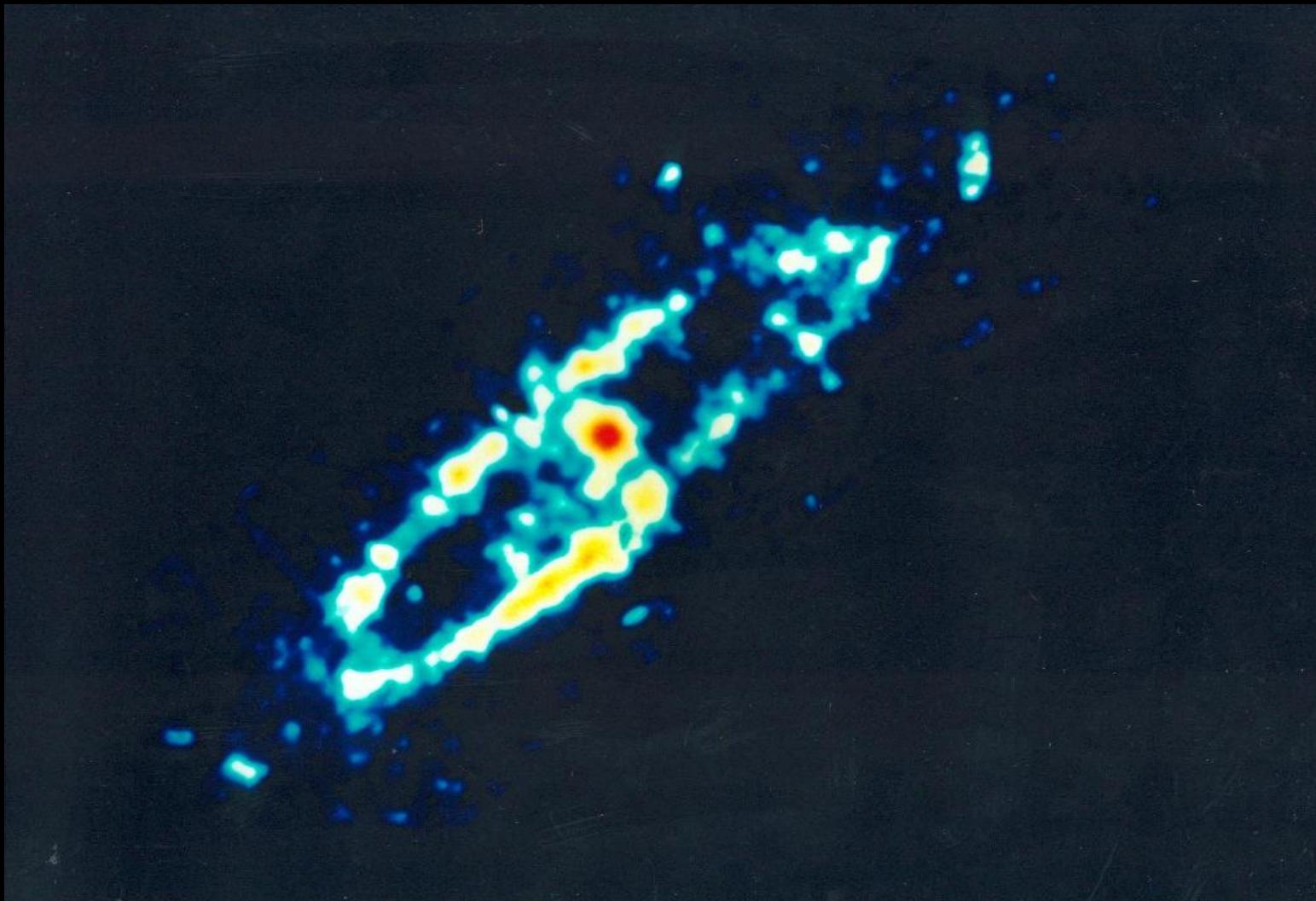
HE gamma-rays (>100MeV EGRET satellite)

Resolved point-like sources:
Binary systems, pulsars, SN remnants...





Andromeda (M31): IR

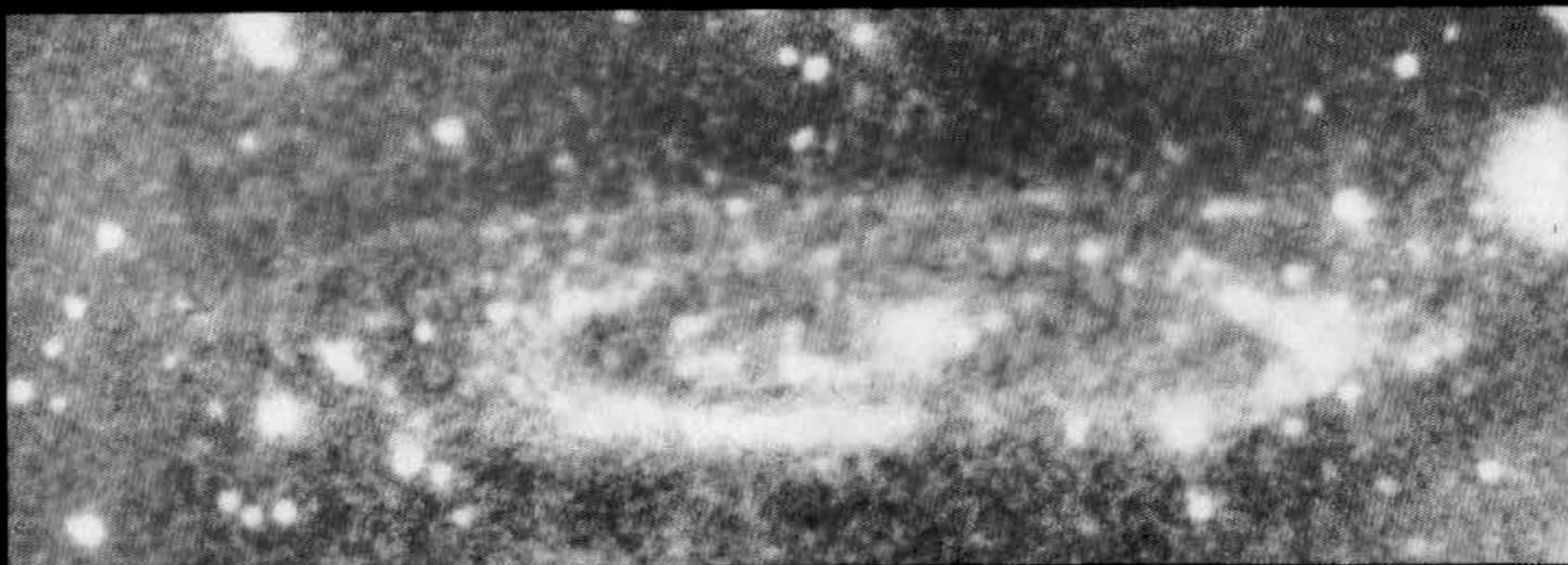


Star forming
regions in
spiral arms

Andromeda (M31): UV

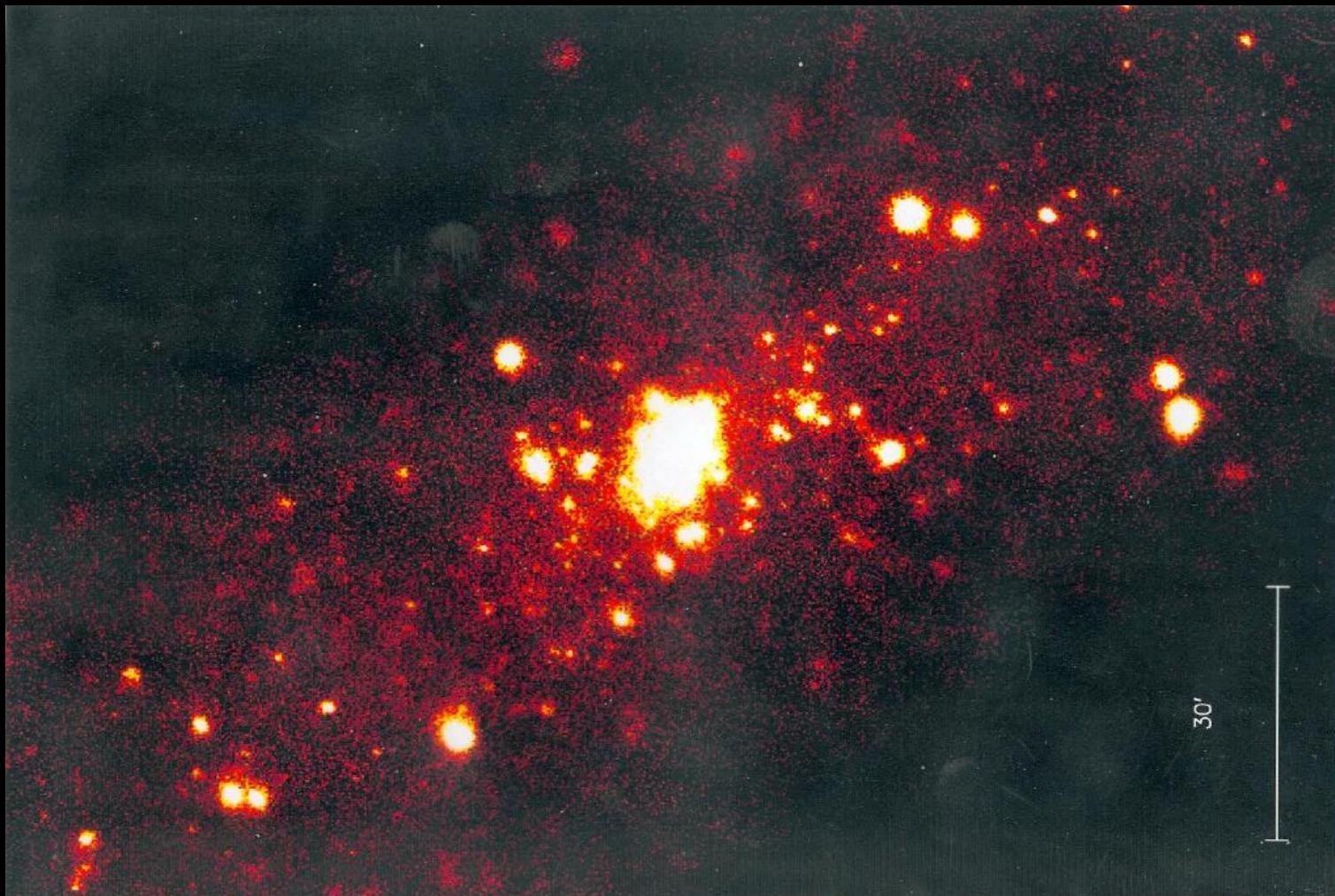
2014

*This photograph of the galaxy M31 reveals the prominence at ultraviolet wavelengths (2000 Å) of young stars in the spiral arms over the older population in the central bulge.
(B. Milliard/Laboratoire d'Astronomie Spatiale).*



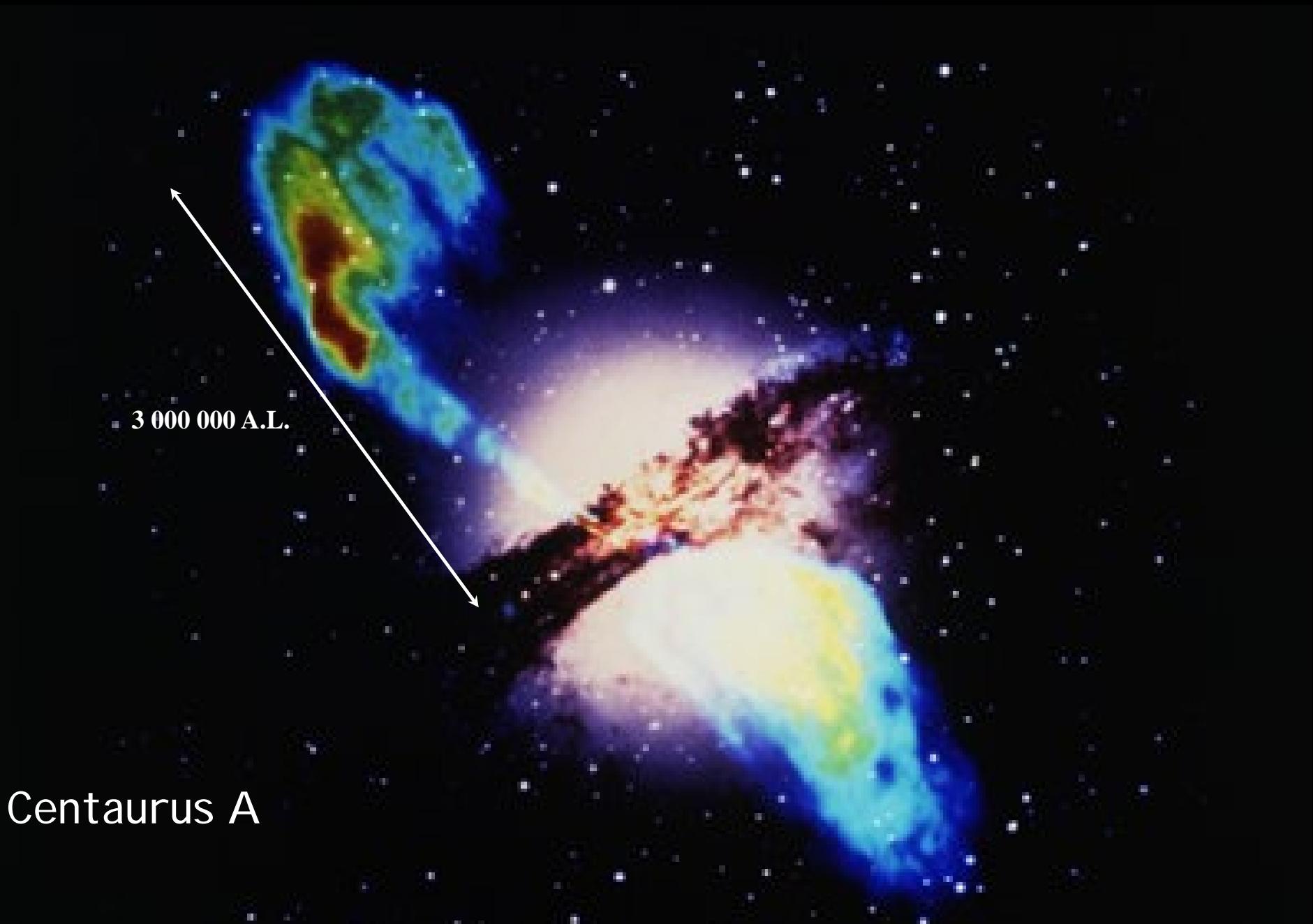
Young,
hot stars
in spiral
arms

Andromeda (M31): Xray



Xray binaries,
supernova
remnants, hot gas

Radio Galaxy

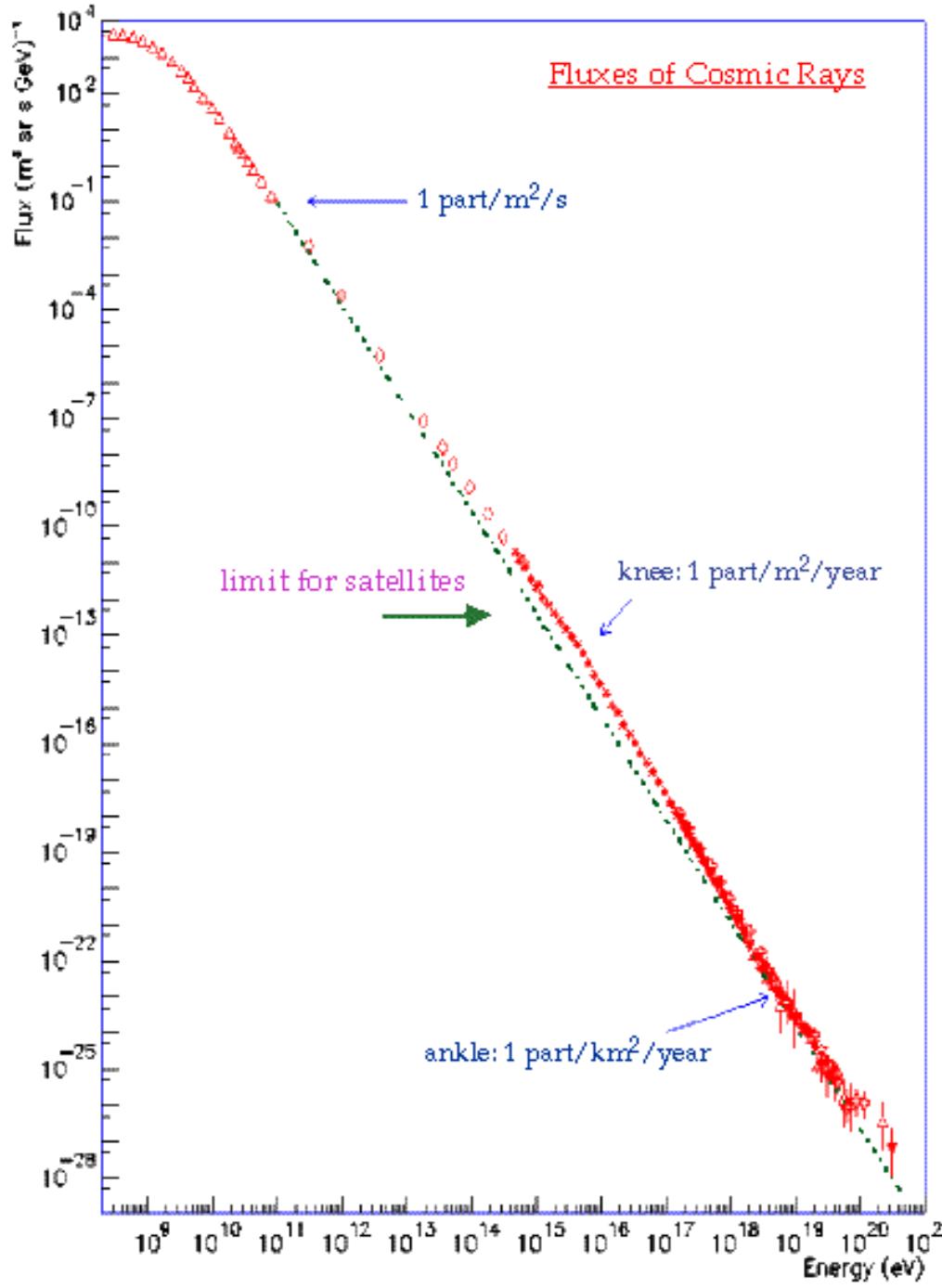


Let there be light !

All what we know in astrophysics is thanks to the light !

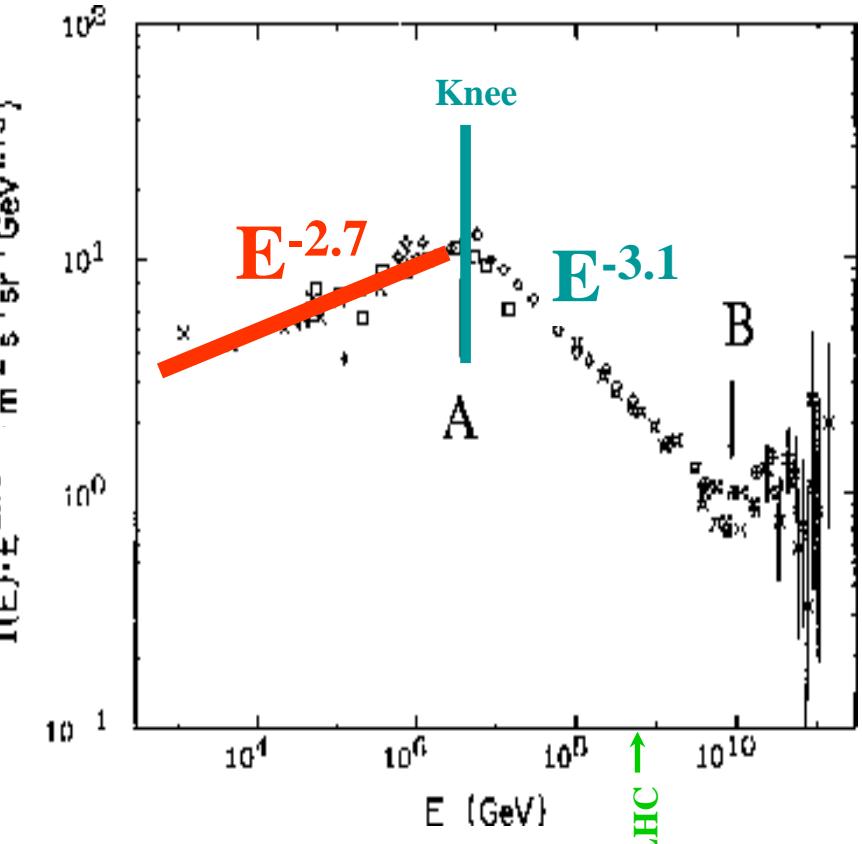
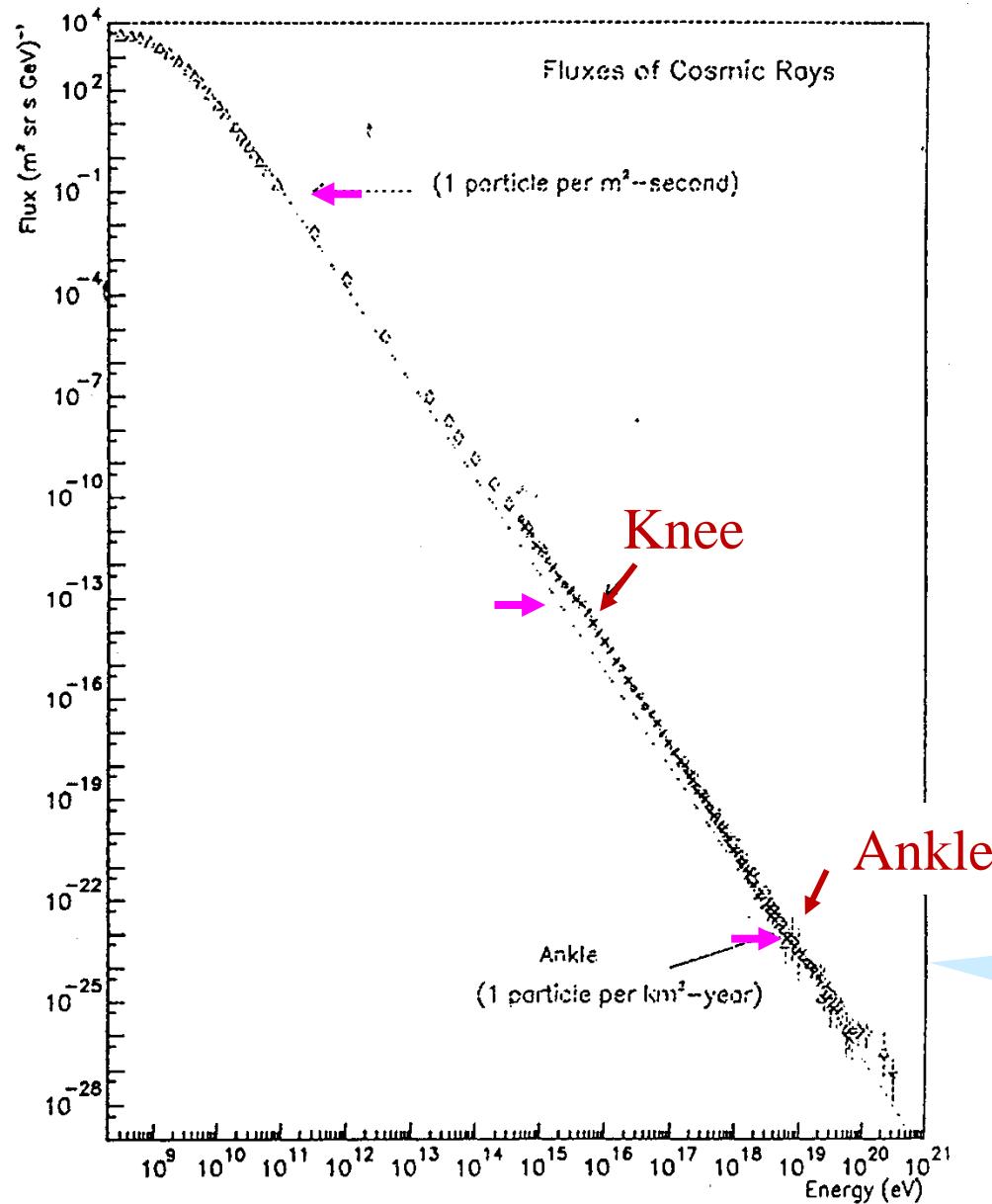
- Temperatures, stars masses, galaxies, magnetic fields, chemical composition, age of stars and structures...
- Nuclear reactions, galactic and extragalactic hydrodynamics, MHD, explosions, nucleosynthesis, past, future... EVERYTHING !
- Well, almost everything...
 - Non-luminous messengers : CR !
 - Rare but precious : $\sim 4 \text{ CR/cm}^2/\text{s}$
 $\sim 30 \mu\text{g/s}$ on entire earth (1kg per year !)
- CR astronomy is impossible...
 - Directions randomized by magnetic fields
 - What we would know if it was the same for photons !
- ...but not astrophysics !
 - Energy spectra and chemical composition tells us a lot...

The "all particles" spectrum



- Regular spectrum over 12 decades in energy, and **32 decades in flux !!!**
- Small break near $3 \times 10^{15}\text{eV}$: the "knee"
- An other one near 10^{18}eV : the "ankle"
- Spectrum badly known at the two extremities
 - Geomagnetic "shield"
 - + Solar modulation
 - Extreme rareness...

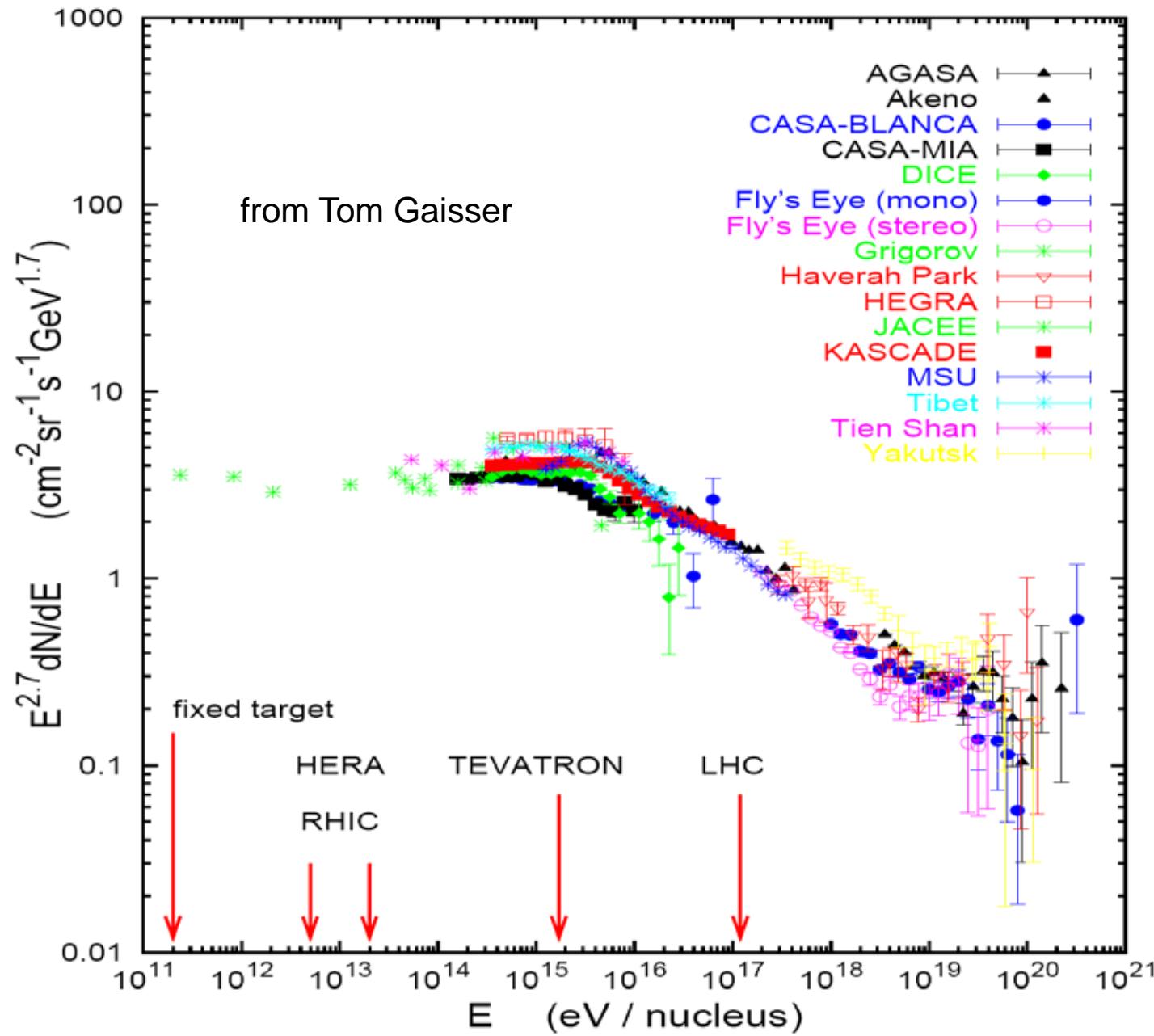
Cosmic-rays spectrum



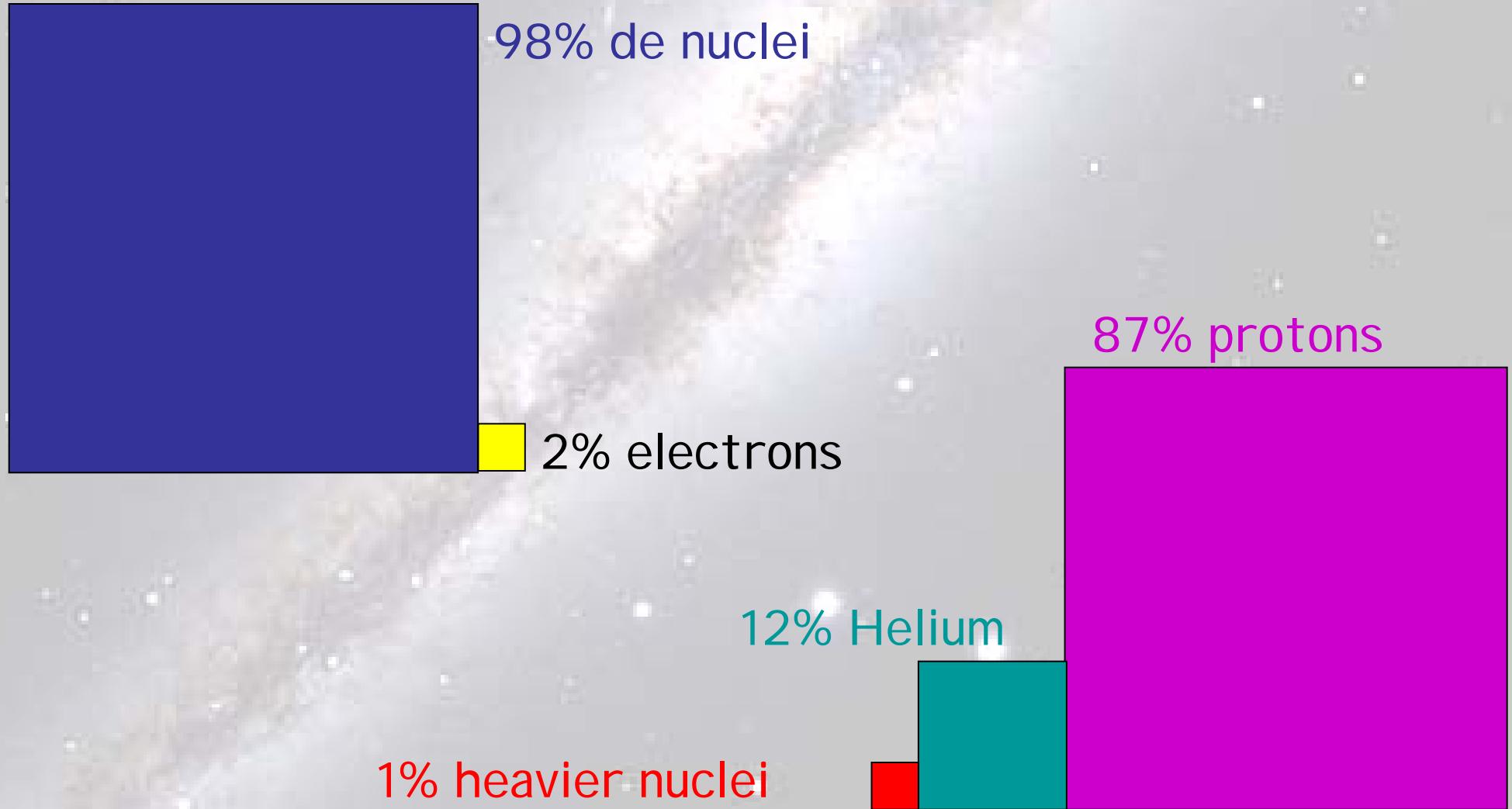
Origin ?
Galactic ?
Cosmologic ?
Astrophysics ?
New physics ?

CR Spectrum above a TeV

2014

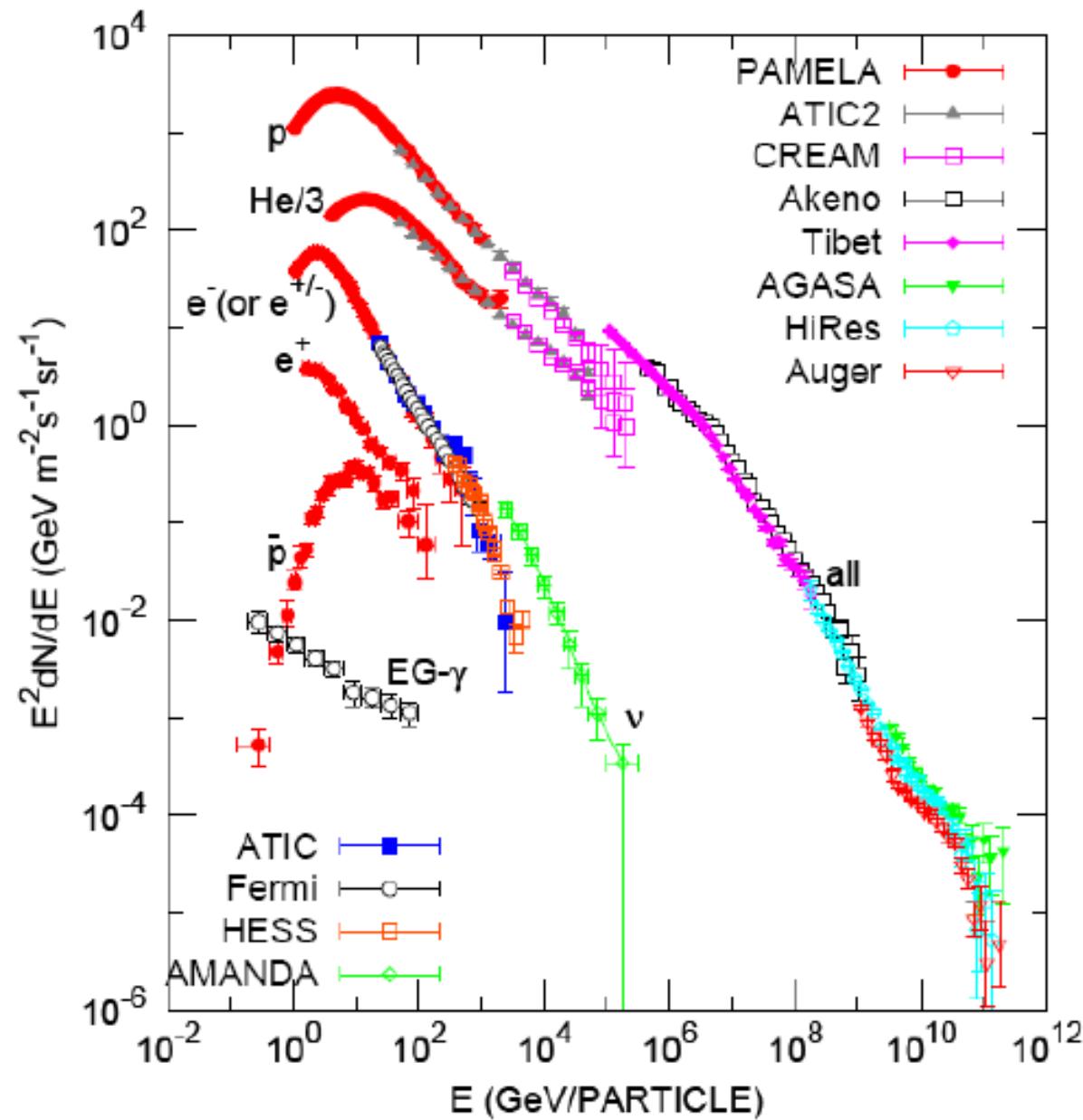


Charge cosmic rays composition



Flux : 4 RC/cm²/s ↳ 1 kg/year << 40 000 ton/year (meteorites)

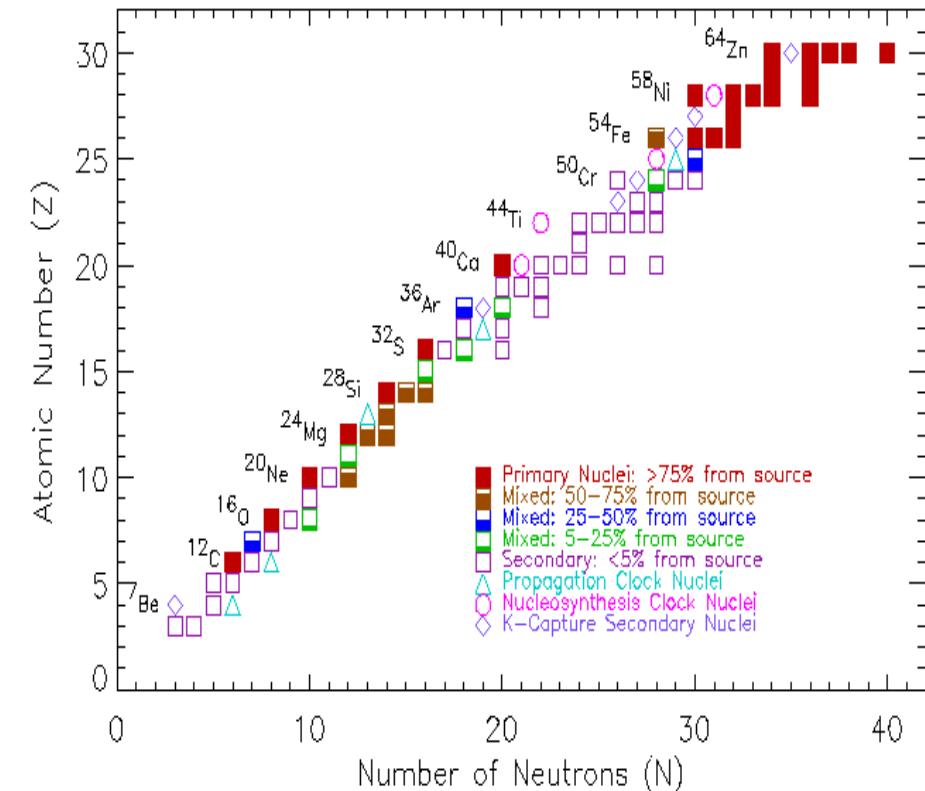
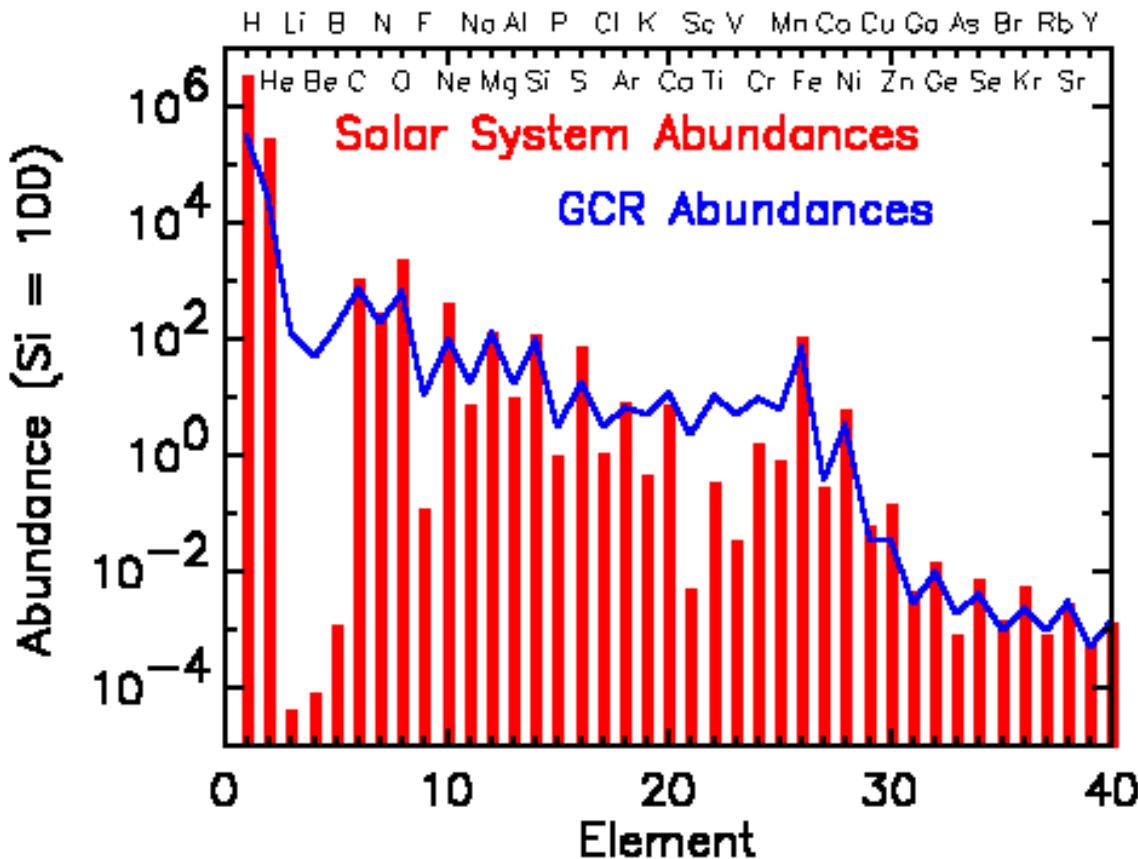
I dentified spectra



Overview of CR data on composition

- Chemical composition
 - Nuclei = 98% (H = 87%, He = 12%, "metals" = 1%)
 - Electrons = 2%
 - More or less standard composition (i.e. solar system) except for fewer H and He, presence of secondary nuclei, and a few "anomalies"...
- Secondary atoms
 - Li, Be, B : spallation of C, N, O (+ nuclei below the Fe peak)
 - Nuclear thicknesses traversed by CR : $X_{CR} = 6$ to 10 g/cm²
- Isotopic anomalies
 - $^{22}Ne \rightarrow$ link with massive stars
- Cosmic clocks
 - $^{10}Be \rightarrow ^{10}B$, $\tau \approx 4 \times 10^6$ years (as well as ^{26}Al , ^{36}Cl , ^{53}Mn , ^{54}Mn , ^{59}Ni)
 - $\tau_{RC} \approx 2 \times 10^7$ years
 - $\frac{X_{RC}}{c\tau_{RC}} \approx 0.2 \text{ part/cm}^3 \Rightarrow$ CR halo extention ($\approx 3\text{-}7$ kpc)

Nature of cosmic rays

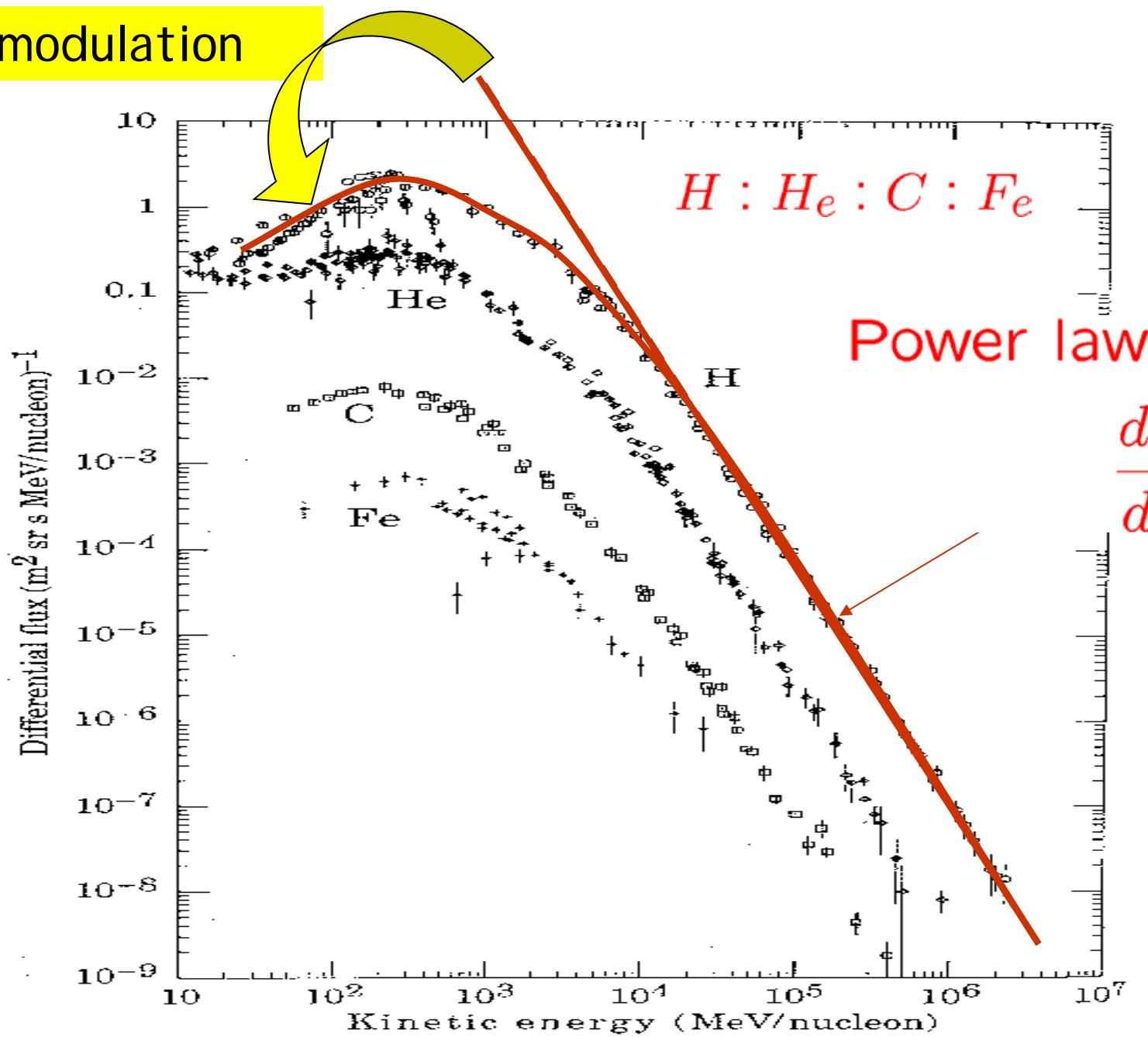


Abundances different in CR and local measurements
(Li Be B and Sub-Fe)

CR undergo spallations and produce secondary CR

Nature of primary cosmic rays

Solar modulation

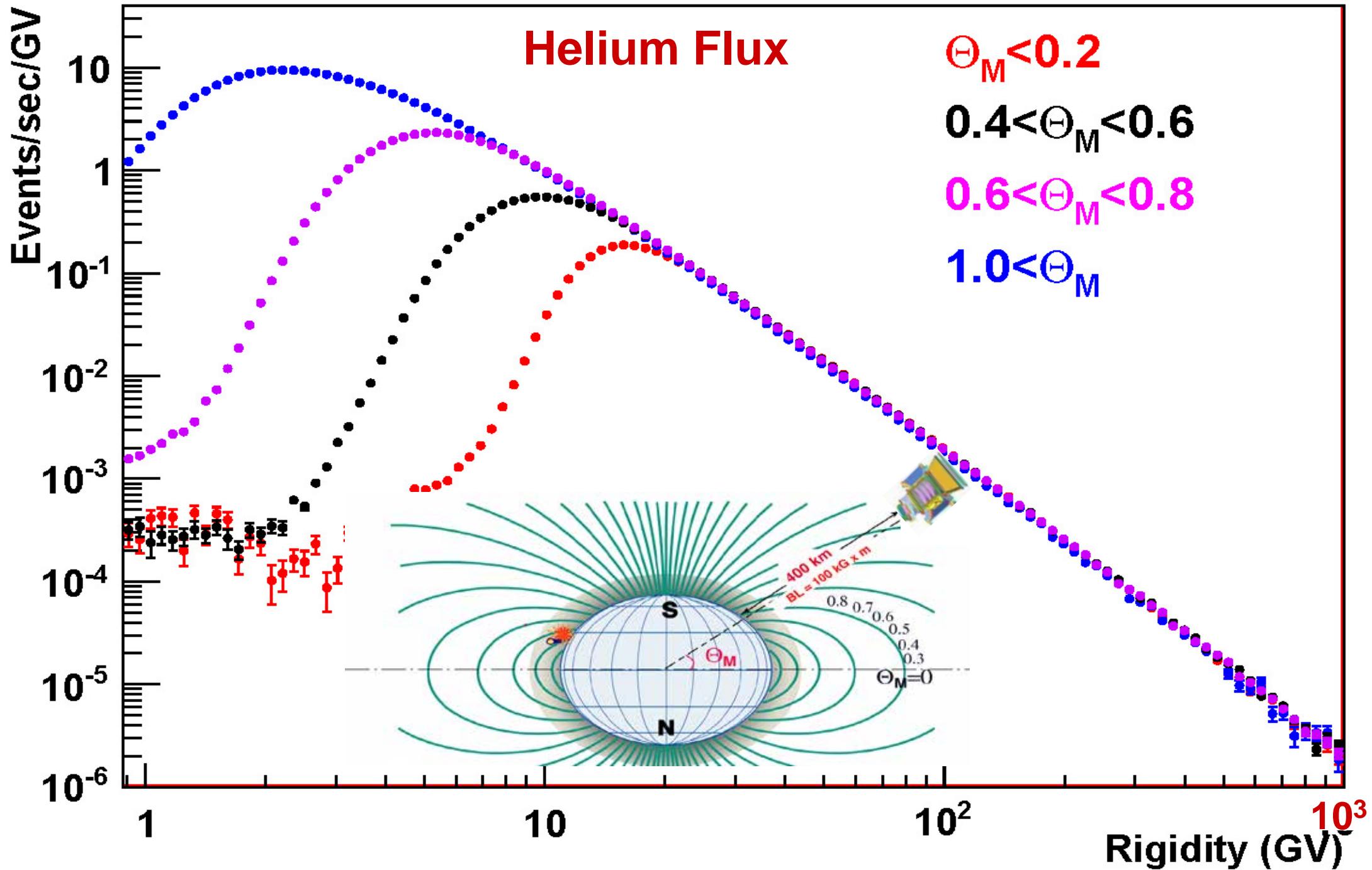


$H : H_e : C : F_e$

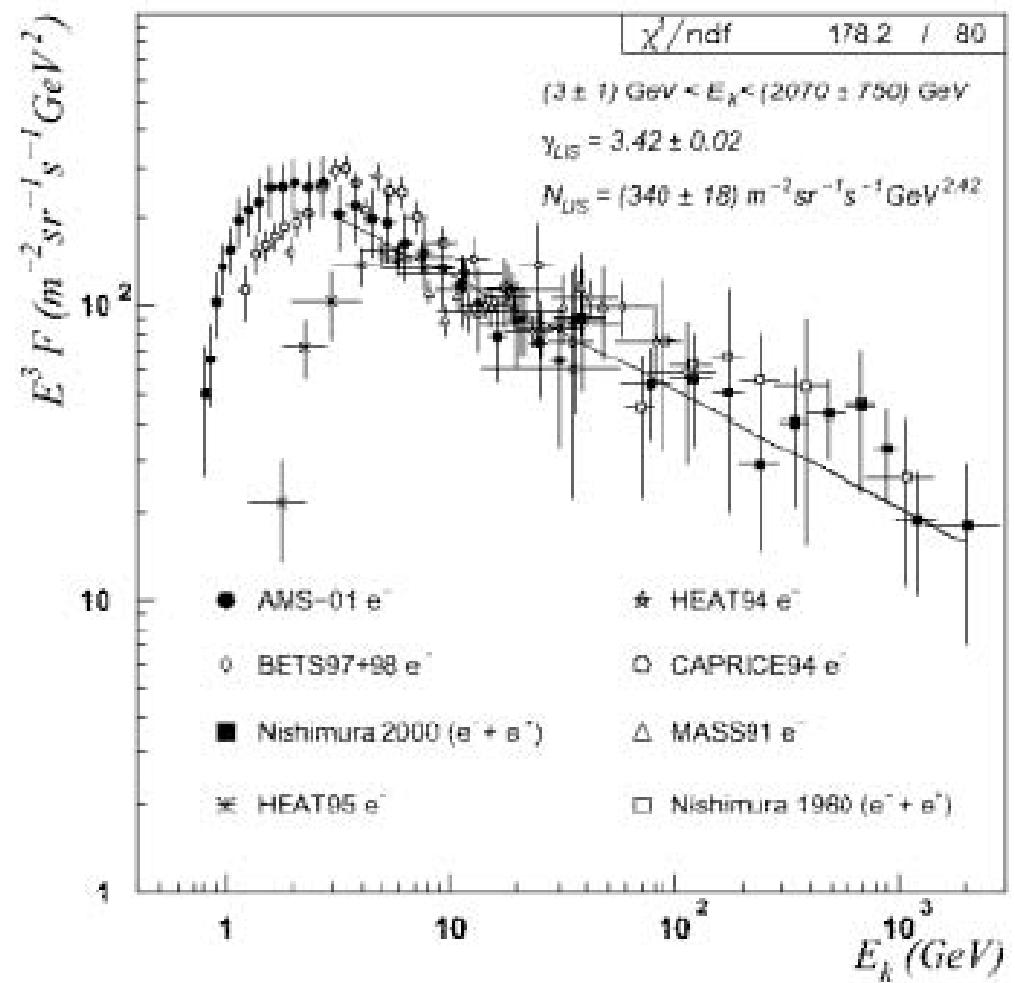
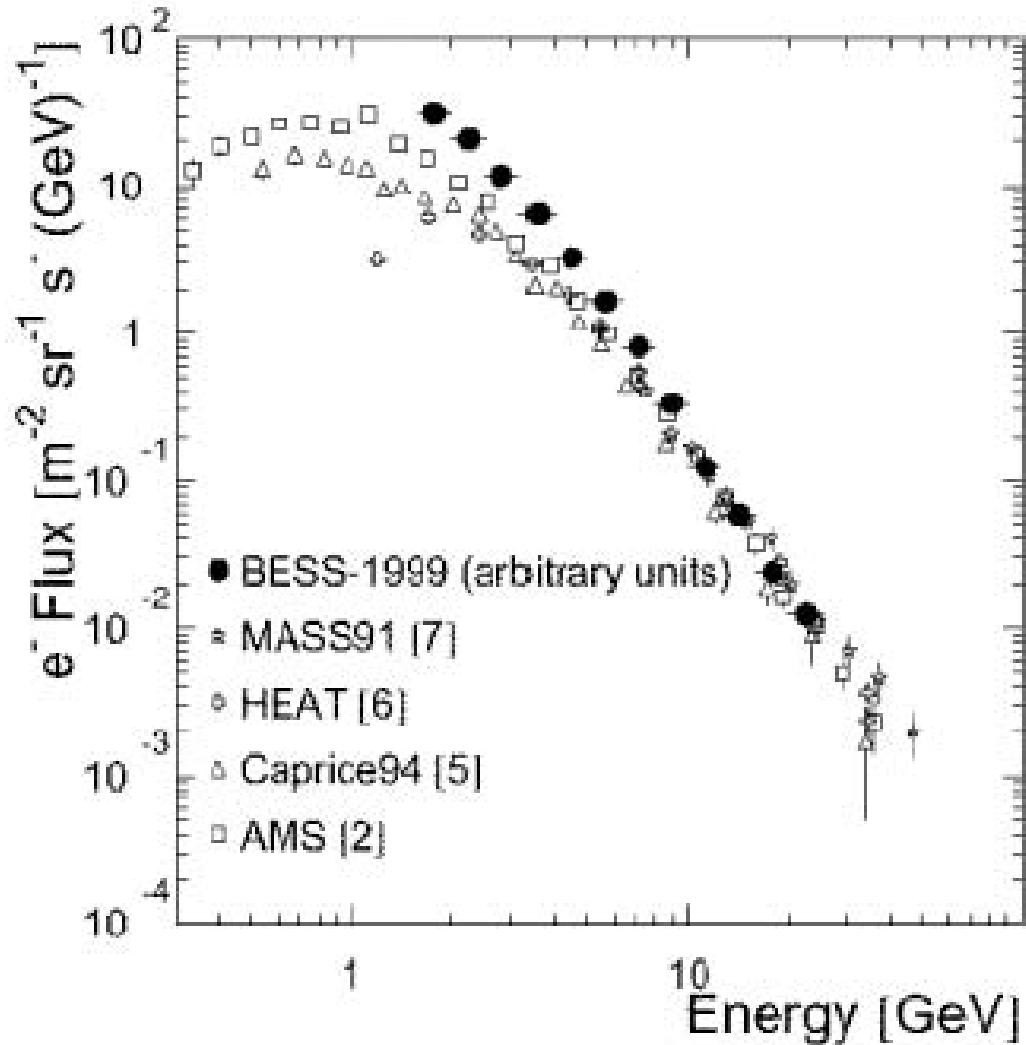
Power law

$$\frac{dN}{dE} \propto E^{-\alpha}$$

Data from AMS on ISS



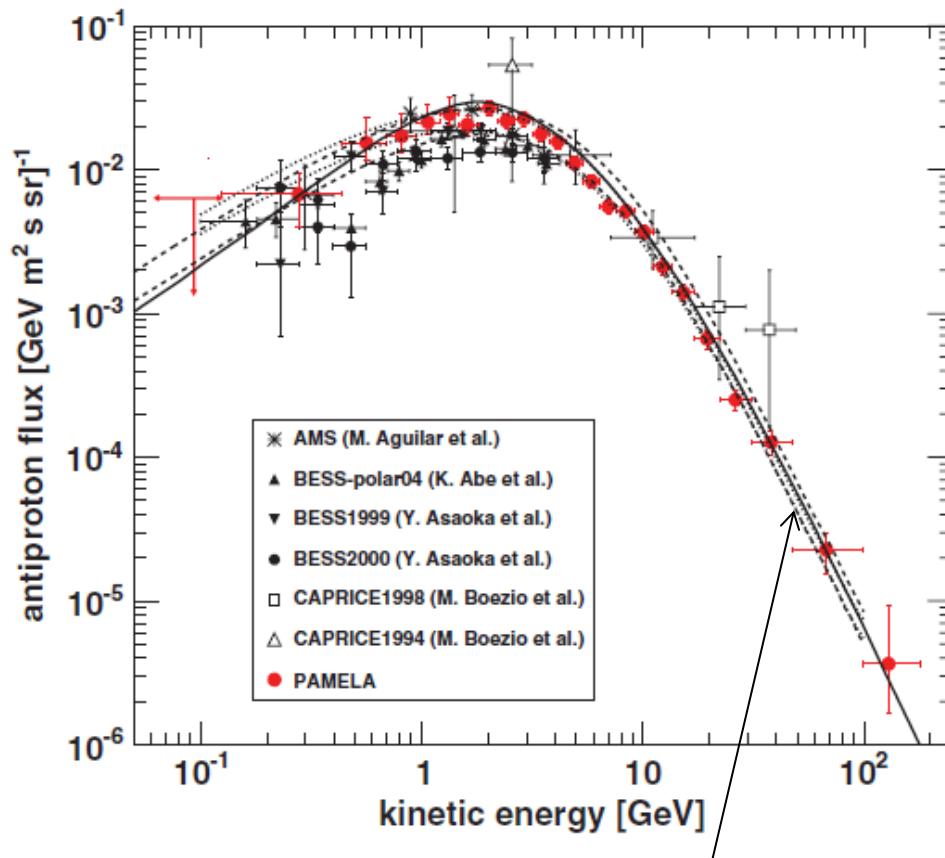
Electron primary flux



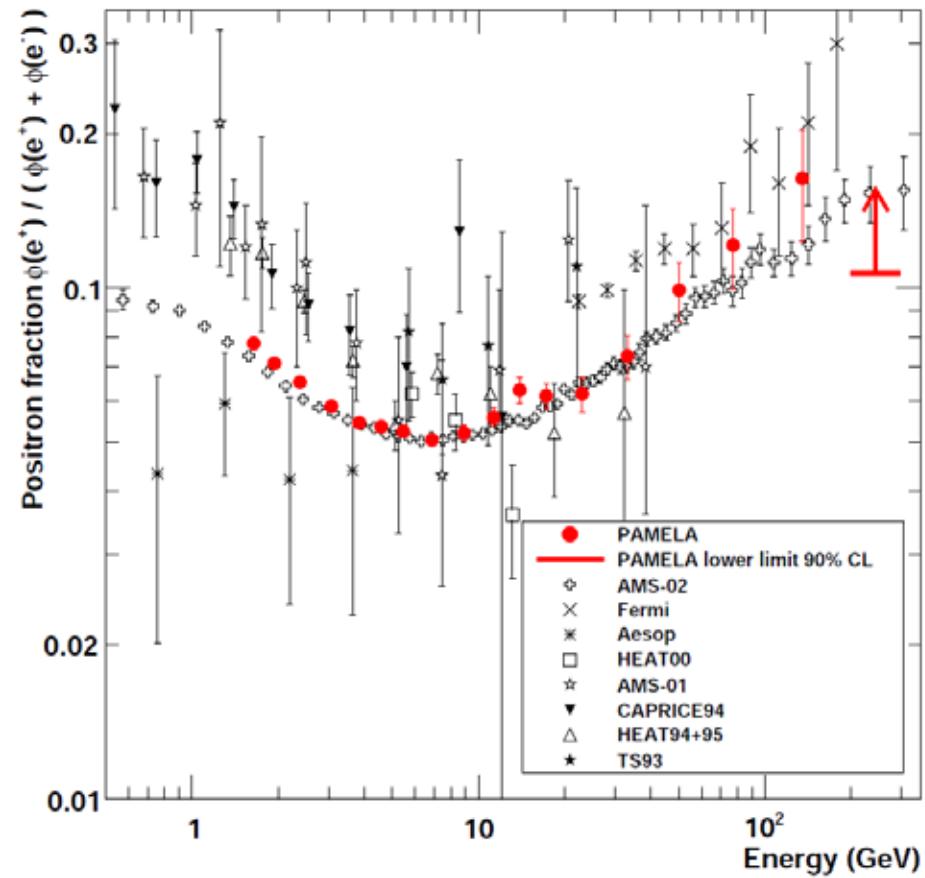
Nature of primary cosmic rays

Antimatter ?

\bar{p}



$$\frac{e^+}{e^+ + e^-}$$

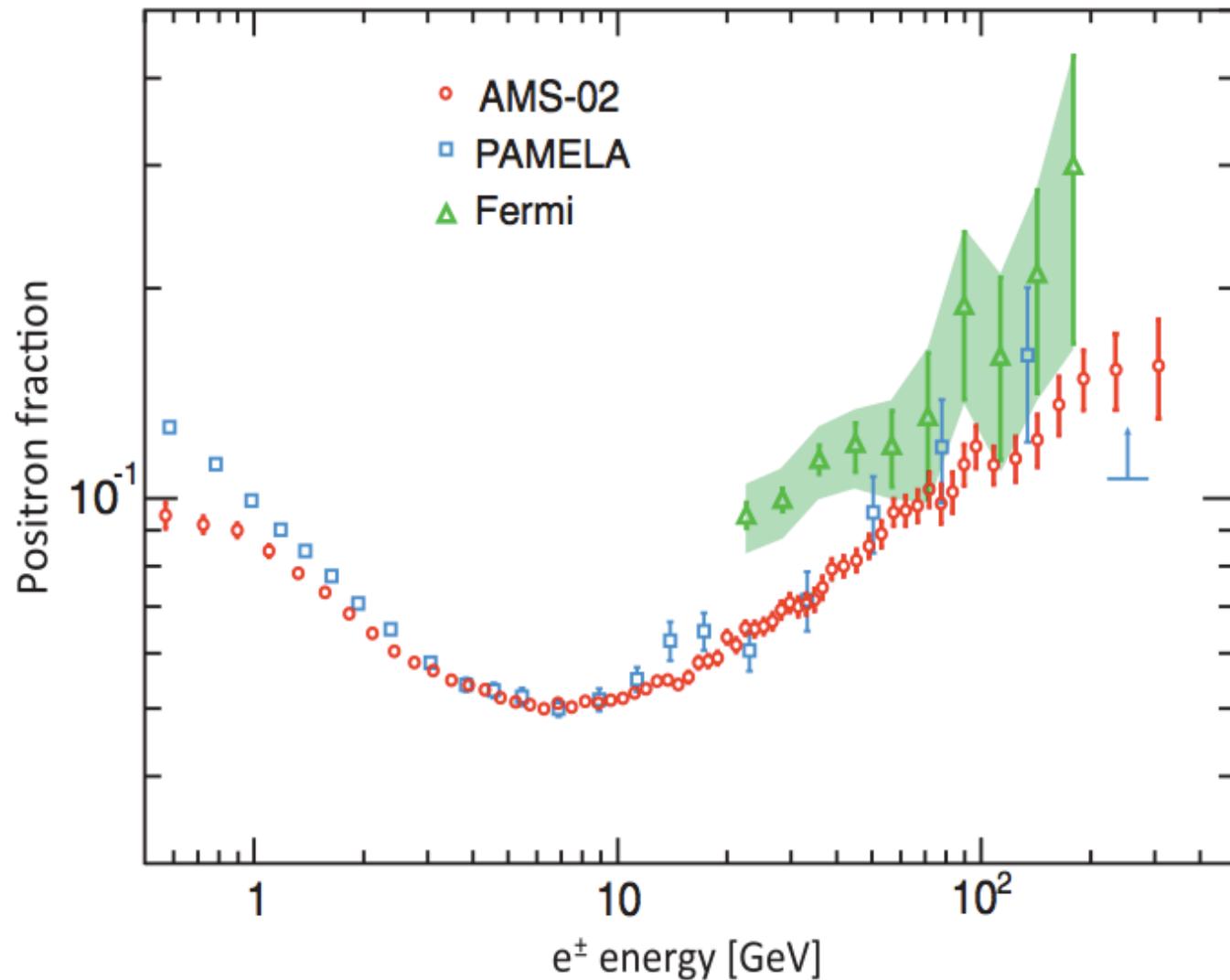


Pure secondary (standard) propagation

Positron fraction in primary flux

2014

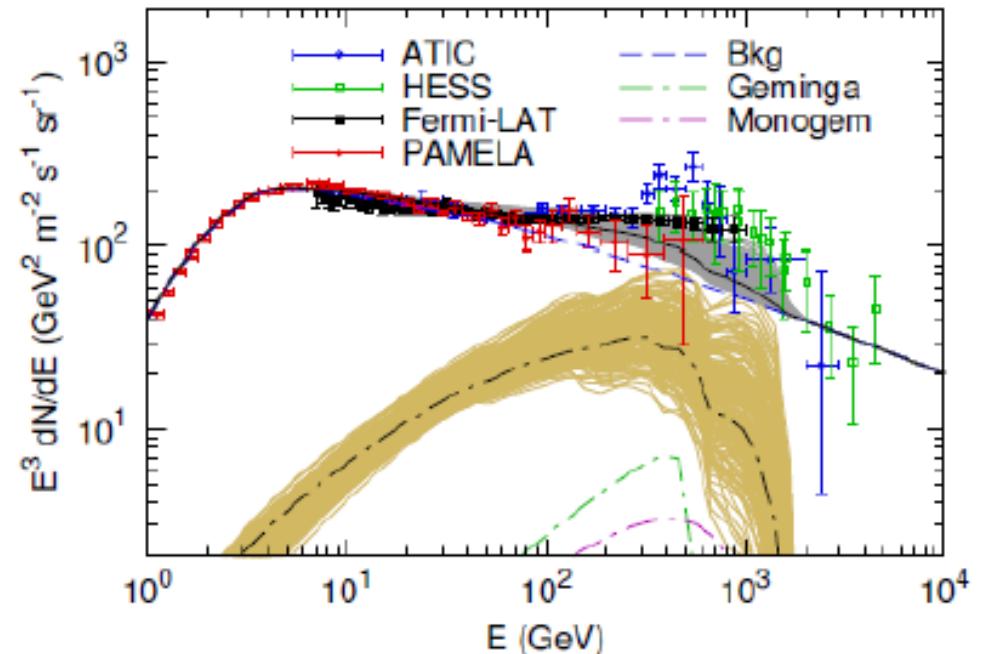
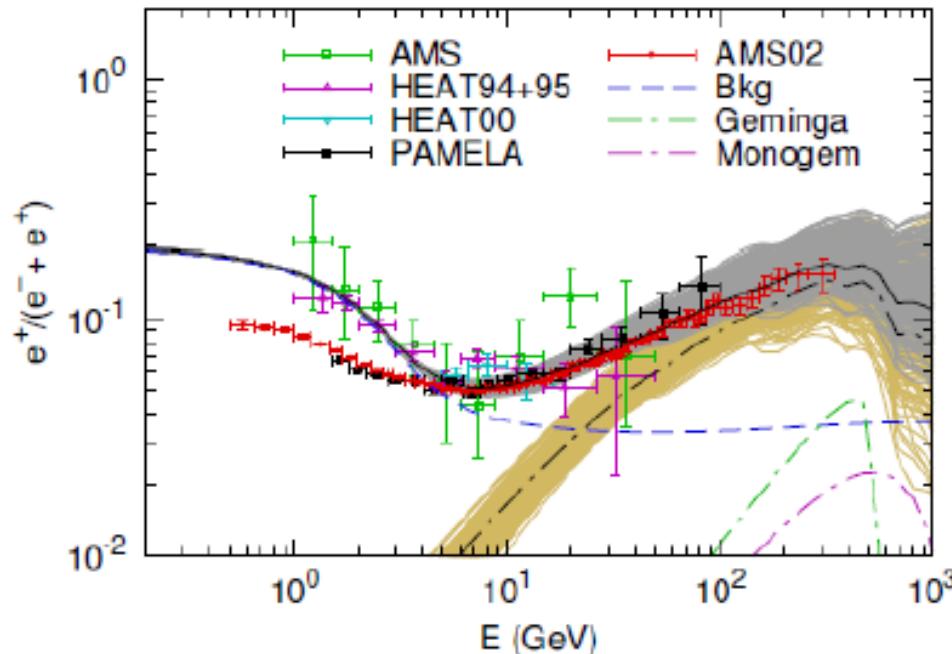
F.Montanet Experimental Astroparticle Physics ESIPAP



Antimatter search (Dark Matter ?)

2014

F.Montanet Experimental Astroparticle Physics ESIPAP



... or rather a boring "local" pulsar spoiling physicists hopes !

Gamma rays

2014

- Gamma-rays observed \gtrsim TeV
- Spectrum \pm understood up to MeV.
- Above, the diffuse spectrum and that of sources are very "hard", in $1/E^2$ and of still unclear origin...

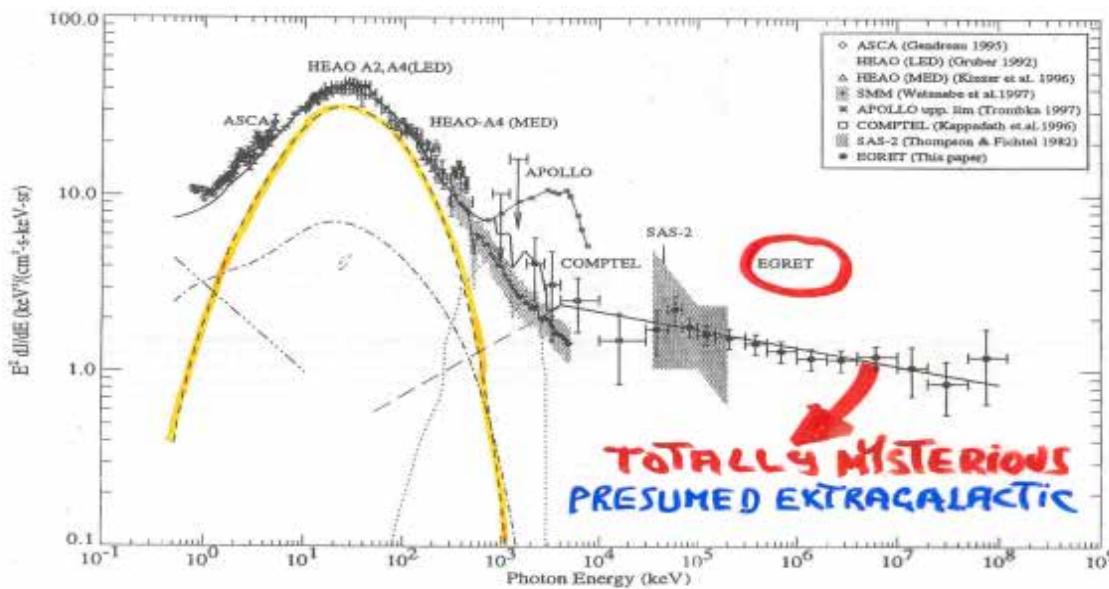
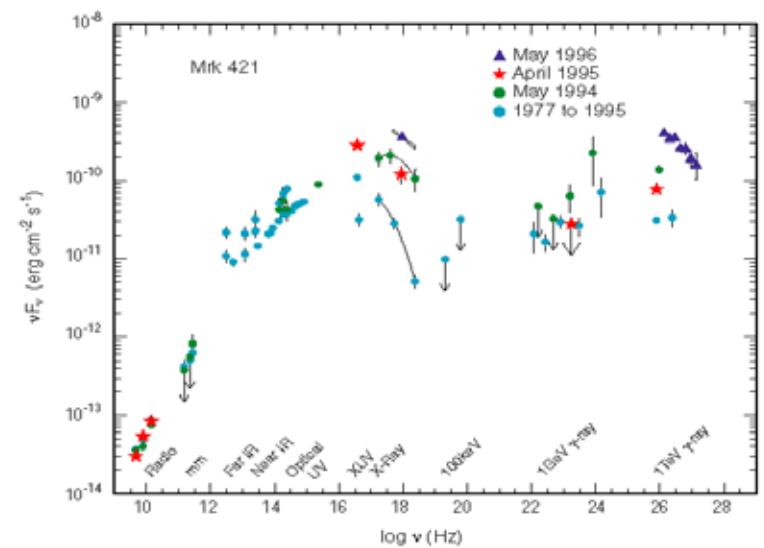


FIG. 10.—Mult>wavelength spectrum of the extragalactic gamma-ray spectrum from X-rays to high-energy gamma rays. The estimated contribution from Seyfert I (dot-dashed line), and Seyfert II (dashed) are from the model of Zdziarski (1996); steep-spectrum quasar contribution (triple-dot-dashed line) is taken from Chen, Fabian, & Gendreau (1997); Type Ia supernovae (dotted line) is from The et al. (1993). The blazar contribution below 4 MeV (long-dashed line) is derived assuming the average blazar spectrum breaks around 4 MeV (McNaron-Brown et al. 1995) to a power law with an index of ~ -1.7 . The thick solid line indicates the sum of all the components.



Why all this non thermal equilibrium radiation?

Gamma, diffuse emission

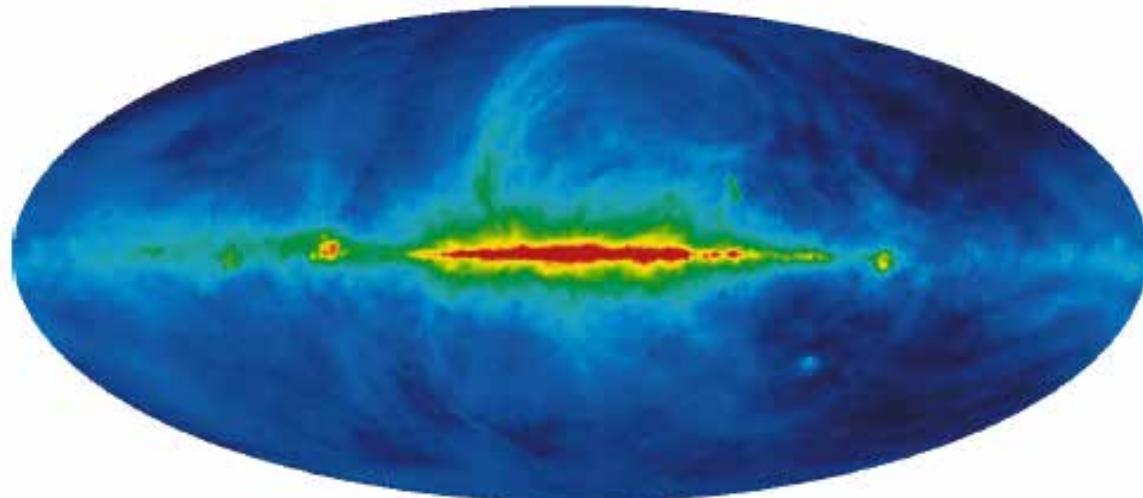
2014

Emission due to:

- the interactions of cosmic electrons with:
 - the magnetic fields (**synchrotron radiation** dominates the radio emission of the Galaxy up to a few GHz)
 - interstellar Matter (ISM); **bremsstrahlung** important below 100 MeV
 - Interstellar photon: **Inverse Compton** above GeV
- the **decay of p^0** produced when CR interact with protons and nuclei
 - $p^0 \rightarrow \gamma\gamma$ above 100 MeV
 - Concomitant emission of n in the decay of p^\pm

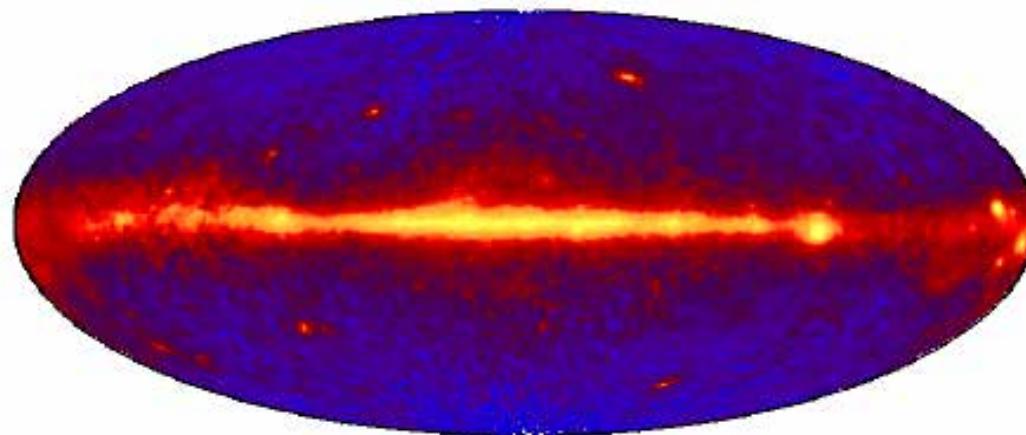
Gamma, diffuse emission

408 MHz



100 MeV

EGRET All-Sky Gamma-Ray Survey Above 100 MeV



Galactic or Extragalactic CR ?

Definite answer from EGRET in 1993 !

Hypothesis: if CR are extra or metagalactic, the density of CR is identical in our Galaxy and in its satellites

- Radio observations radio give the mass of gas M_H in the SMC
- M_H implies a measurable flux for SMC of: $2.5 \times 10^{-7} \text{ cm}^{-2} \cdot \text{s}^{-1}$
$$F_\gamma \propto M_H N_{CR} R_q$$
- EGRET gives an upper limit (at 95%CL): $< 0.5 \times 10^{-7} \text{ cm}^{-2} \cdot \text{s}^{-1}$
- The CR density is 5 times smaller within SMC

Cosmic rays are indeed Galactic !

The general problematic

- Thermal speeds ! RCUHE (few 10^{20} eV)



- From top to bottom (decay...)
- From bottom to top (acceleration)



- Energy losses (Synch., IC, p, pairs...)
- Destruction (photo-dissociation...)
- Escape probabilities
- Propagation in ISM and IGM (mag fields: deflection, confinement...)
- Re-acceleration



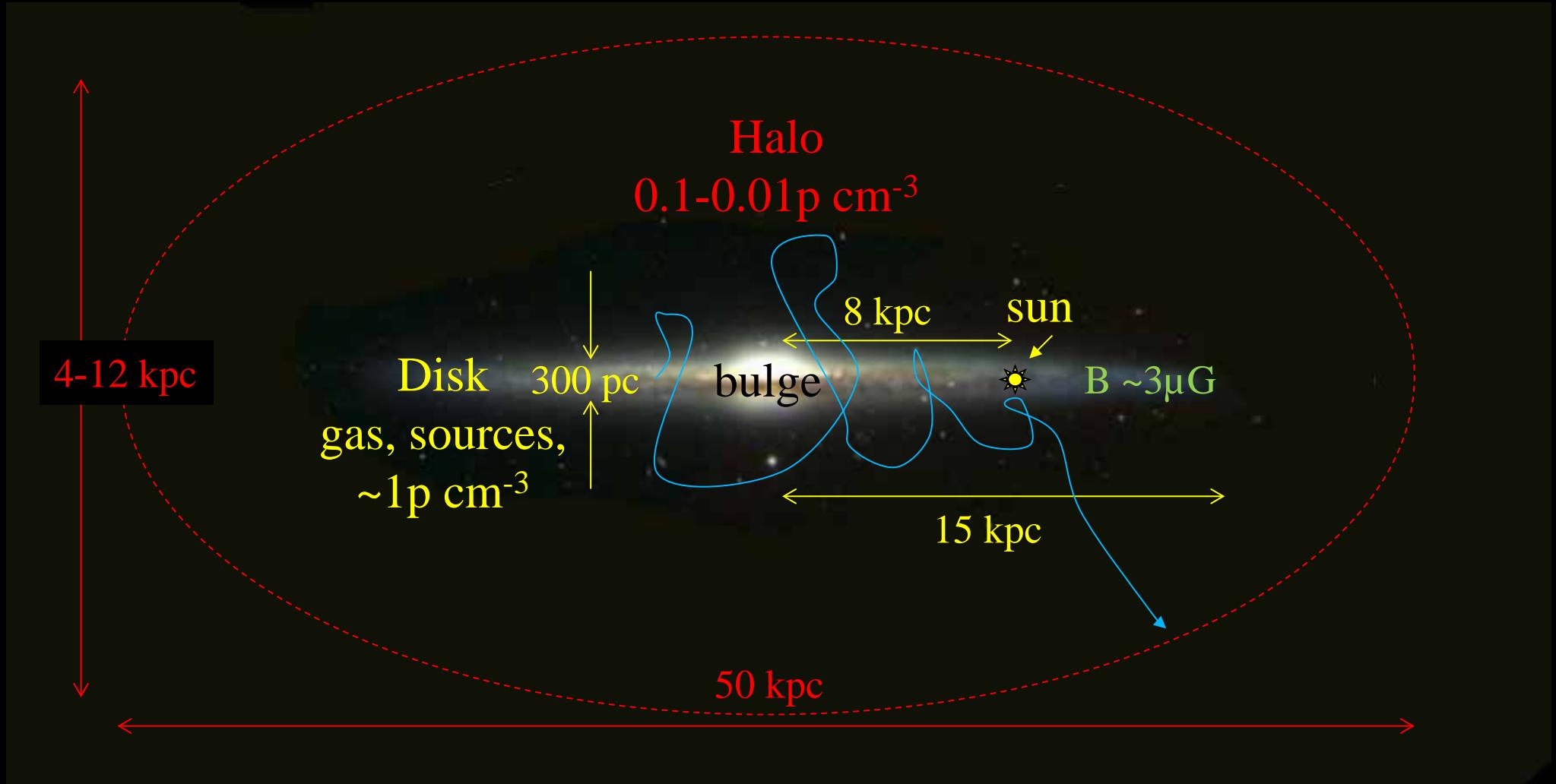
- Balloons, satellites...
- Air showers...
 - Cherenkov telescopes
 - Surface & Fluorescence Detectors



Propagation medium,
IGM, ISM and
atmosphere

Dimensions of the Milky Way

$$1 \text{ pc} \approx 3 \text{ l.y.} \approx 3 \times 10^{16} \text{ m}$$



Milky Way, a spiral galaxy

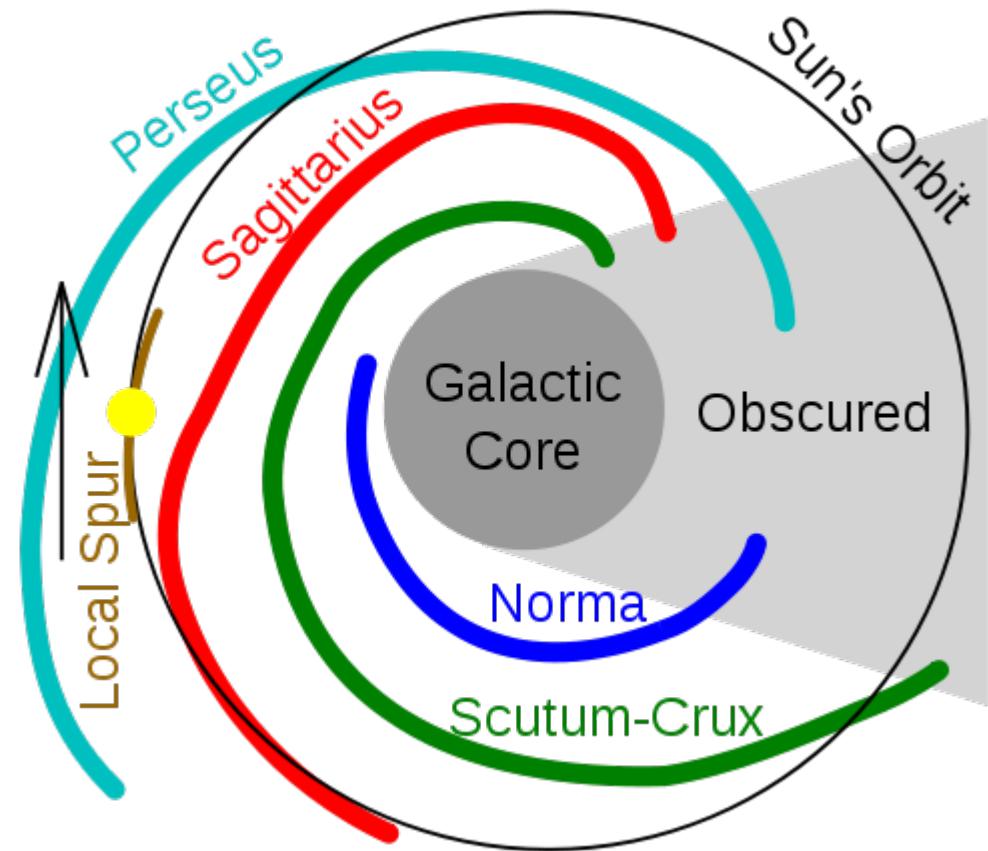


Milky Way, a spiral galaxy

2014

Local spur and neighboring arms
) local matter and B field inhomogeneity.

Mean "regular" B field ~ $3\mu\text{G}$ roughly parallel to spiral arms, more intense in between arms.

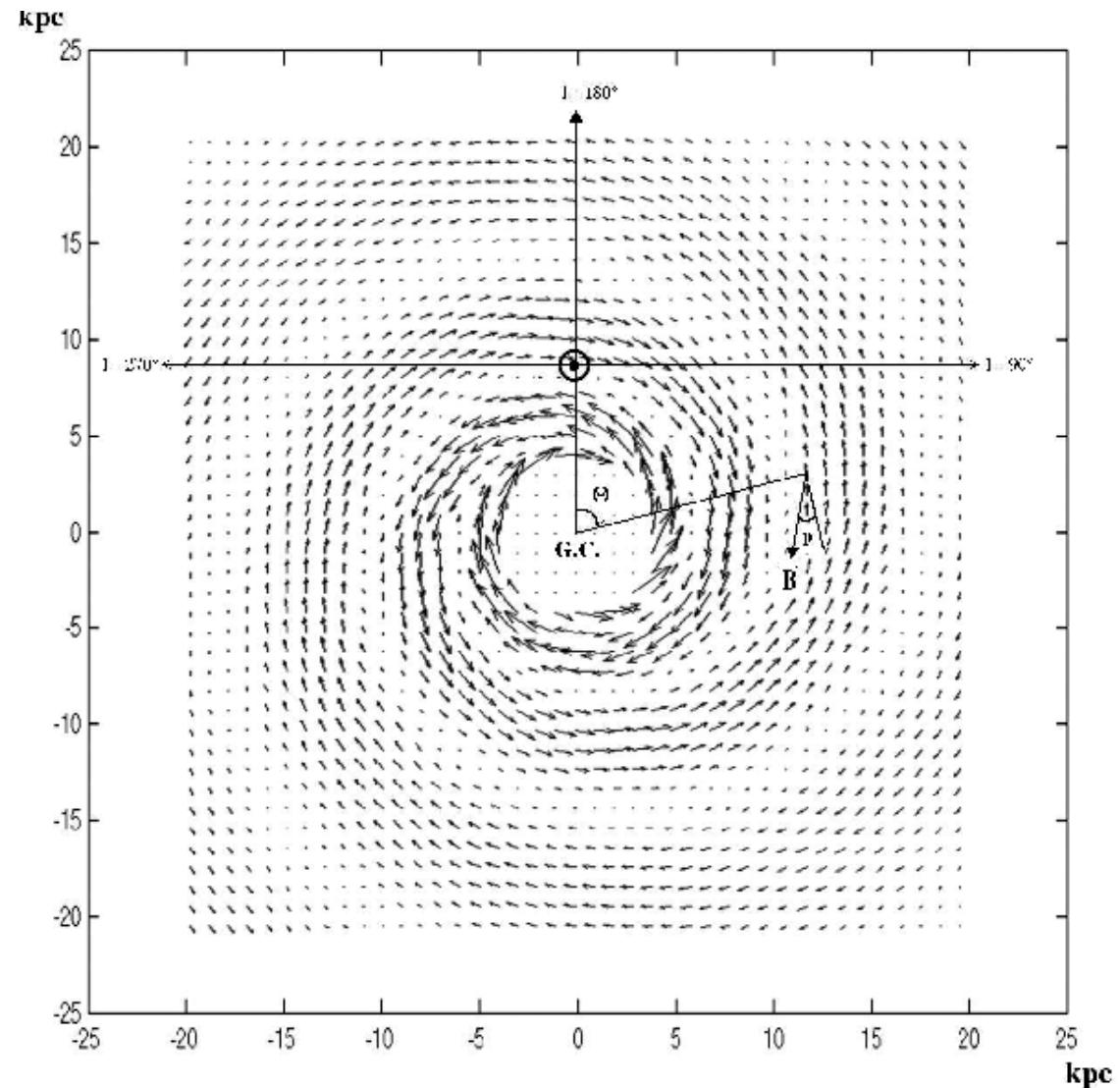


Milky Way, a spiral galaxy

2014

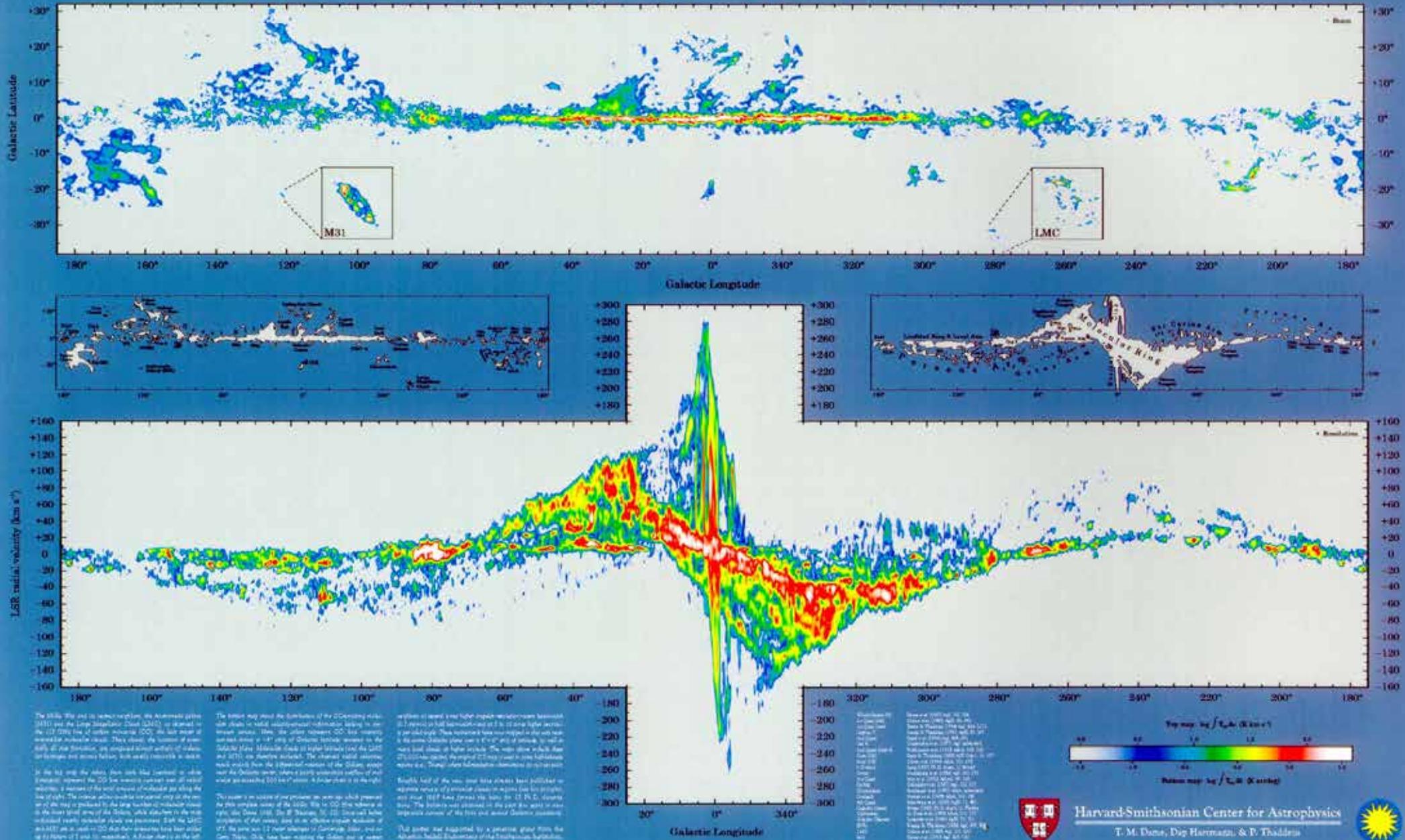
Local spur and
neighboring arms
) local matter and
B field inhomogeneity.

Mean "regular" B field
 $\sim 3\mu\text{G}$ roughly parallel
to spiral arms, more
intense in between arms.



A thick target...

The Milky Way in Molecular Clouds

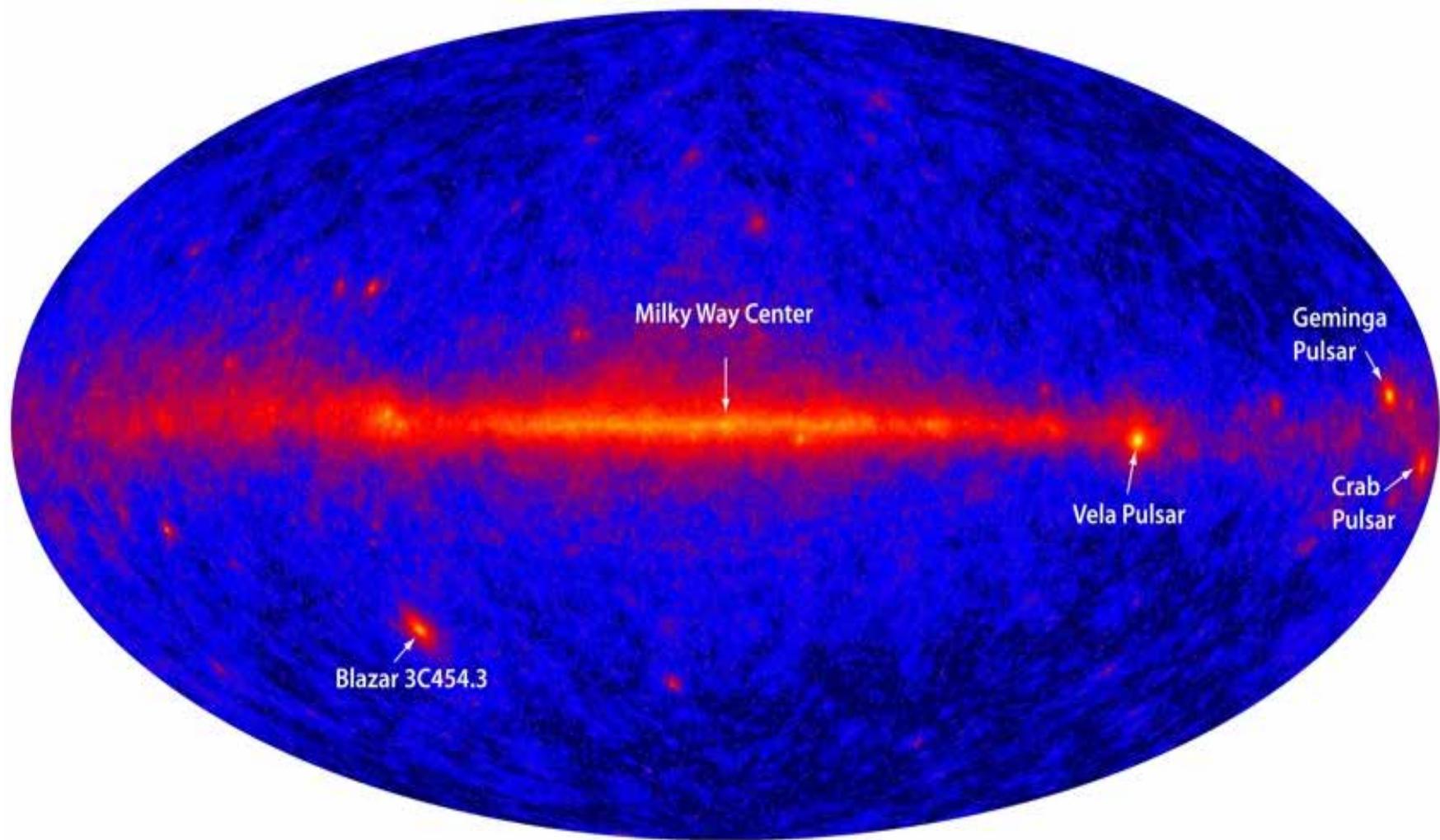


Harvard-Smithsonian Center for Astrophysics
T. M. Dame, Dag Hartmann, & P. Thaddeus



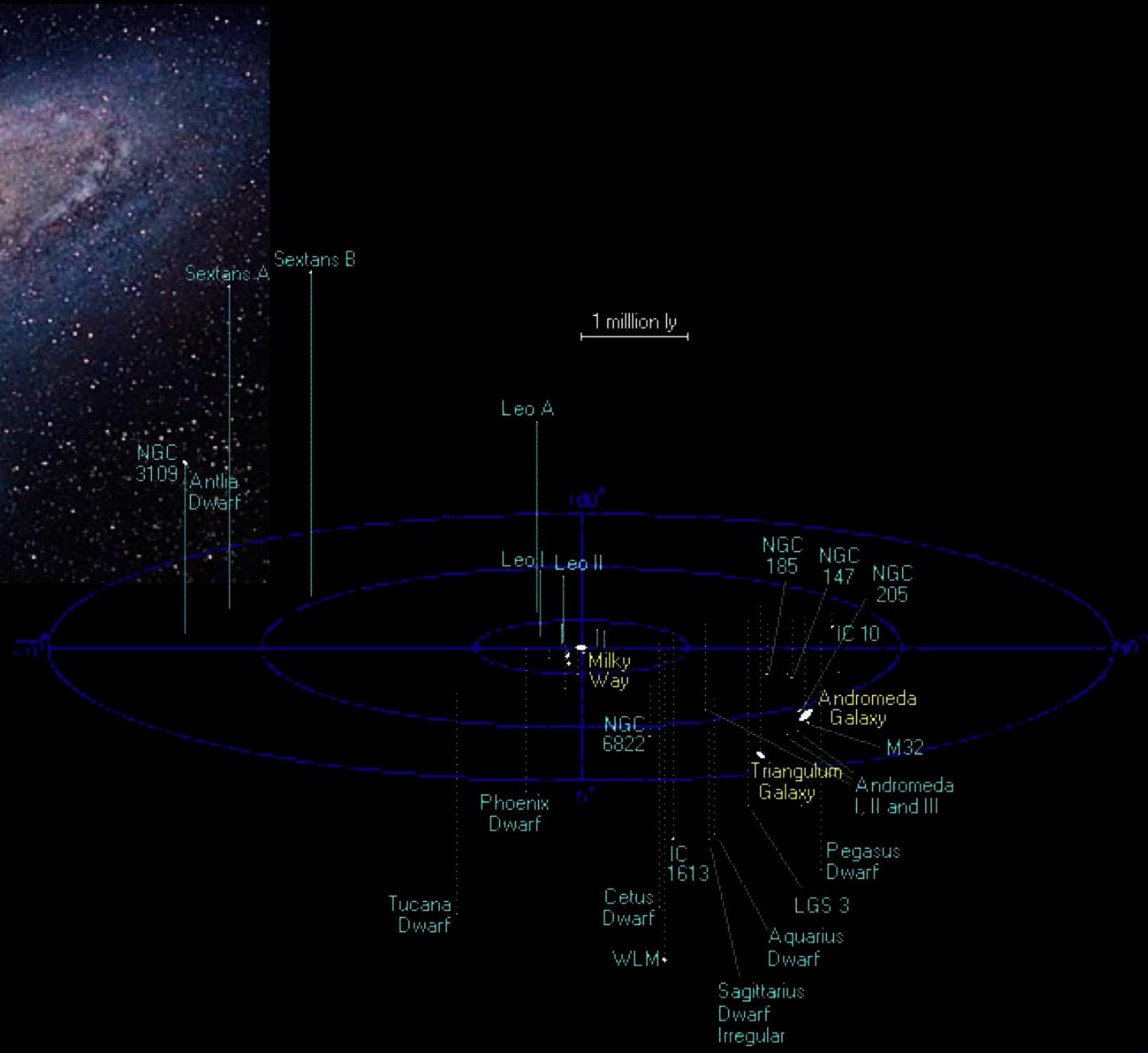
A thick target

- Diffuse gamma-ray emission from galactic CR interaction with matter (mostly molecular H clouds).



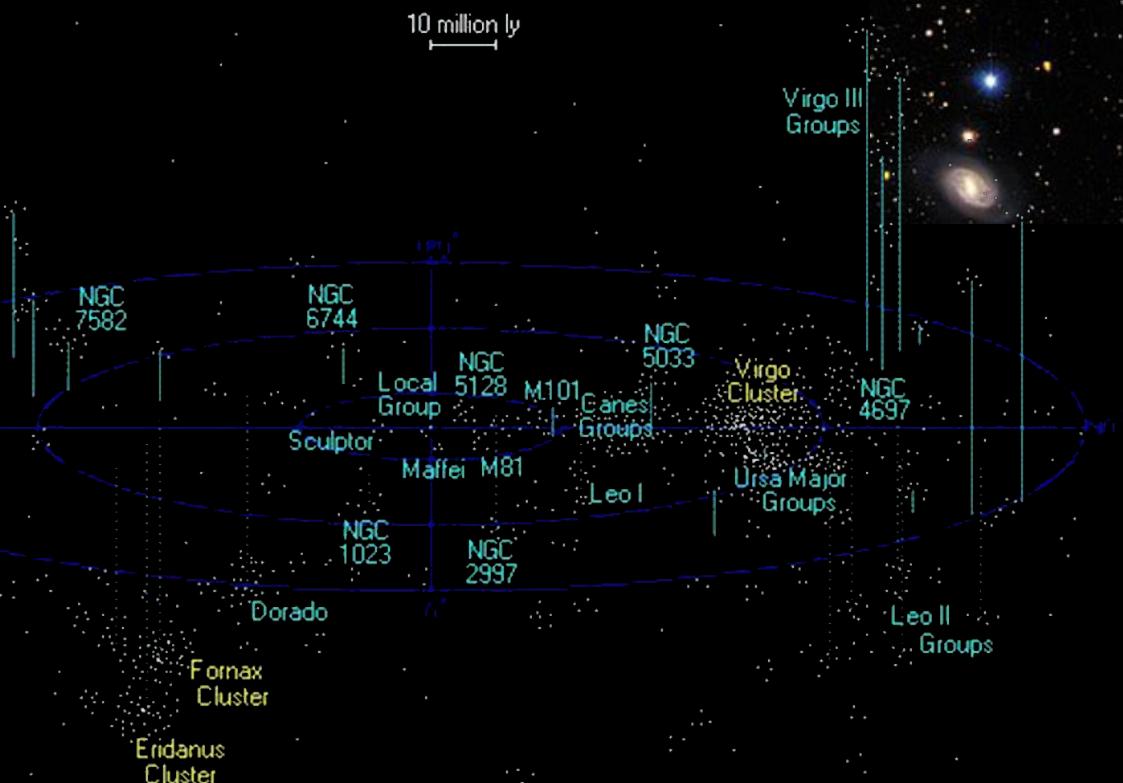
The nearby islands...

Andromeda (M31)
A twin of our Milky Way
slightly larger and (only)
distant by 780kpc.
Many small (dwarf)
galaxies are orbiting
around these twins.



The local group and the Virgo cluster

Our local group is at the periphery of the large Virgo supercluster (~2000 galaxies) at ~20Mpc

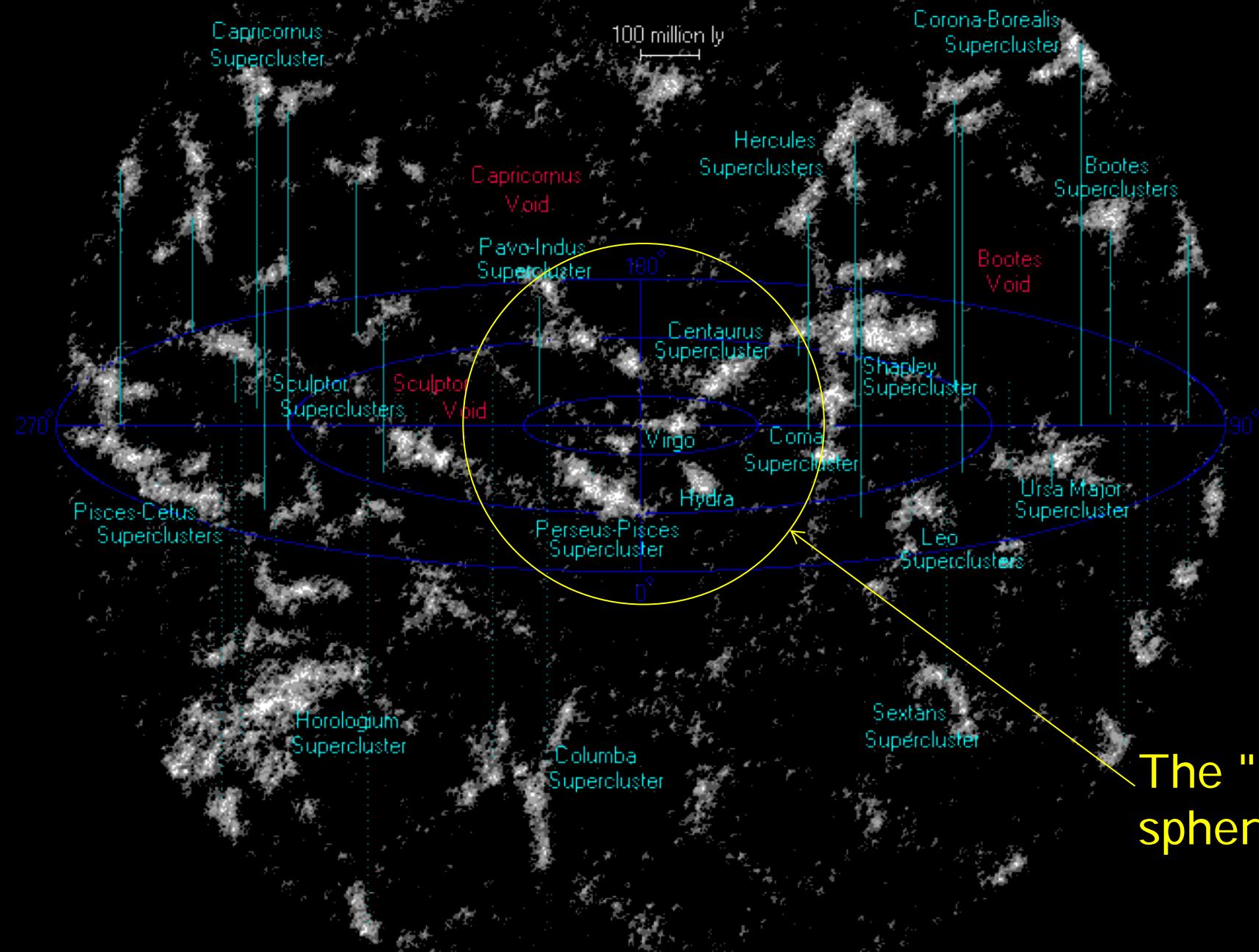


Another super cluster: Abel 1689

2014

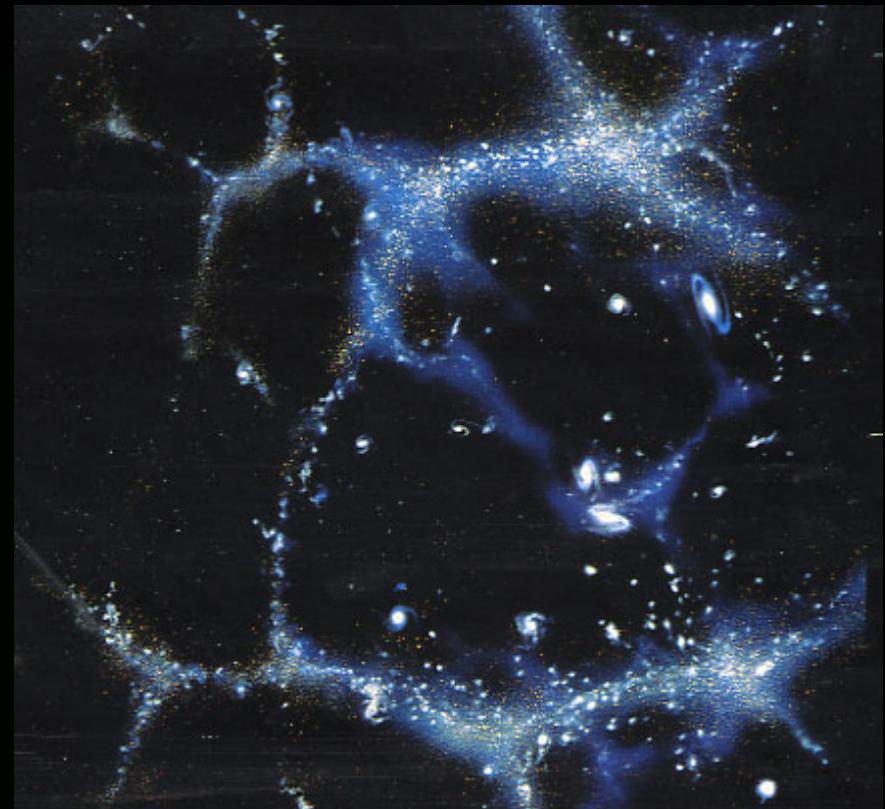
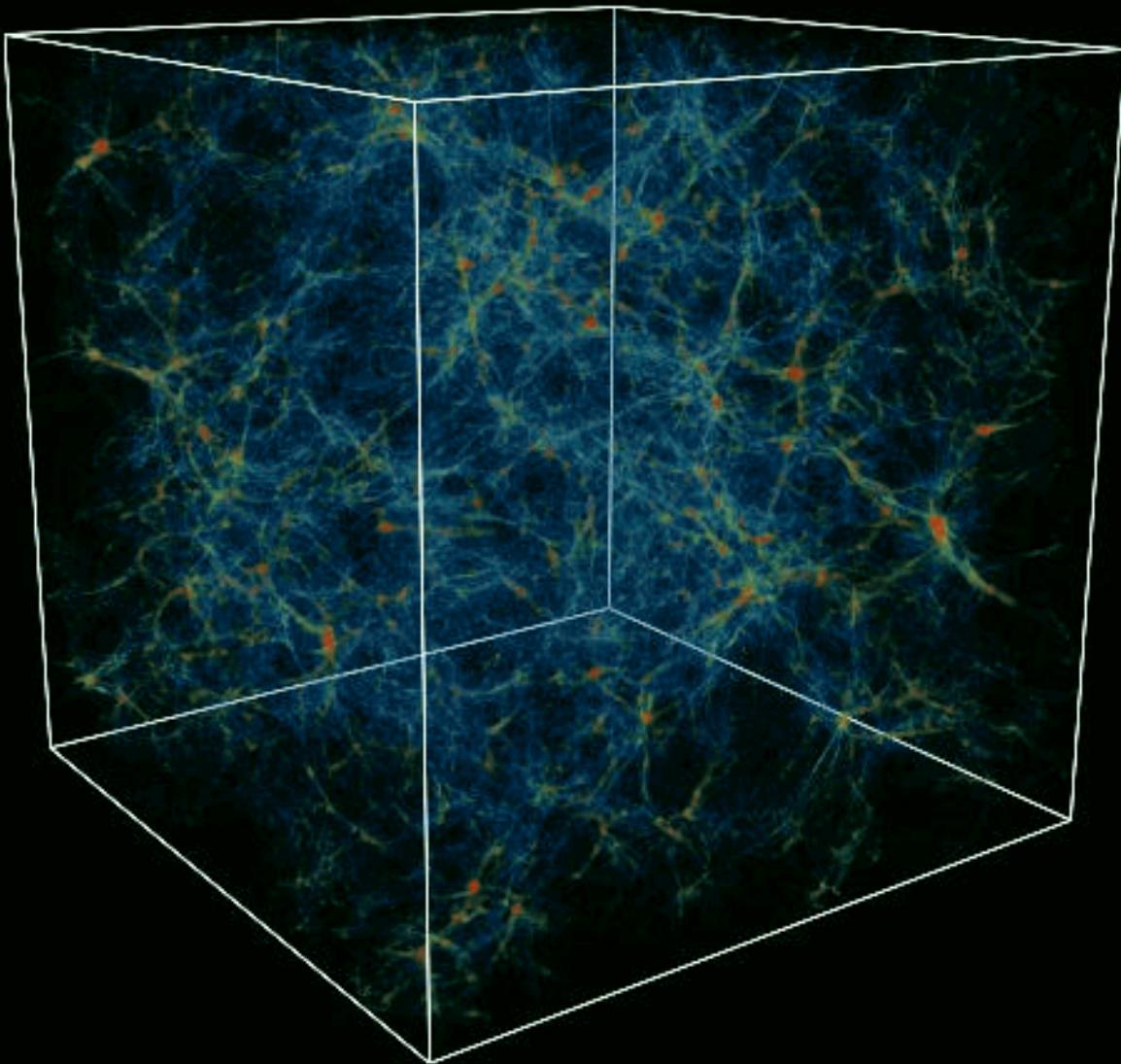
F. Montanet Experimental Astroparticle Physics ESIDAP

A 300Mpc horizon



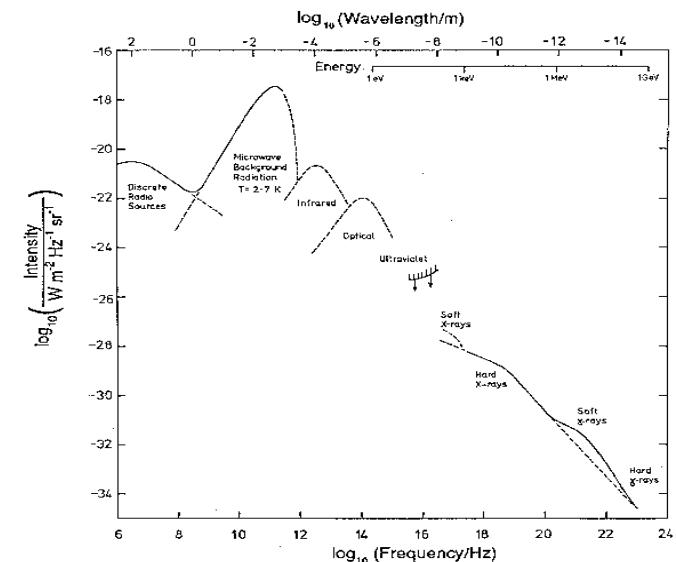
The "GZK"
sphere

Large scale filamentary structures



Vacuum is not emptiness !

- Inter Galactic Medium contains:
 - Magnetic fields (regular + random) are highly speculative and range from 2×10^{-6} nT (20pG) to 10^{-4} nT (1nG).
 - Very little matter (p, He, and a few electrons):
 - 1 proton / m³
 - Electromagnetic radiations:
 - 413 CMB photons per cm³
 - Also IR, radio photons...
 - Neutrinos:
 - Mostly C°B neutrinos (decoupled when universe was only 2" old!)
 - Today 1.95 K i.e. 1.7×10^{-3} eV
 - 336 ° (all species) per cm³
 - + Many mysterious dark matter WIMPs ...



The earth atmosphere

- An evident characteristic of the atmospheric medium is that of being inhomogeneous.
 - Its density, diminishes six orders of magnitude when the altitude above sea level passes from zero to 100km, and another additional six orders for the range 100km to 300km.
 - Although up to ~100km, the composition is nearly constant: 78.47% N, 21.05% O, 0.47% Ar and 0.03% other elements.
 - It follows a quasi exponential profile ("quasi" because T is not quite constant!)

$$\rho(h) = \rho_0 e^{-gMh/RT}$$

The earth atmosphere

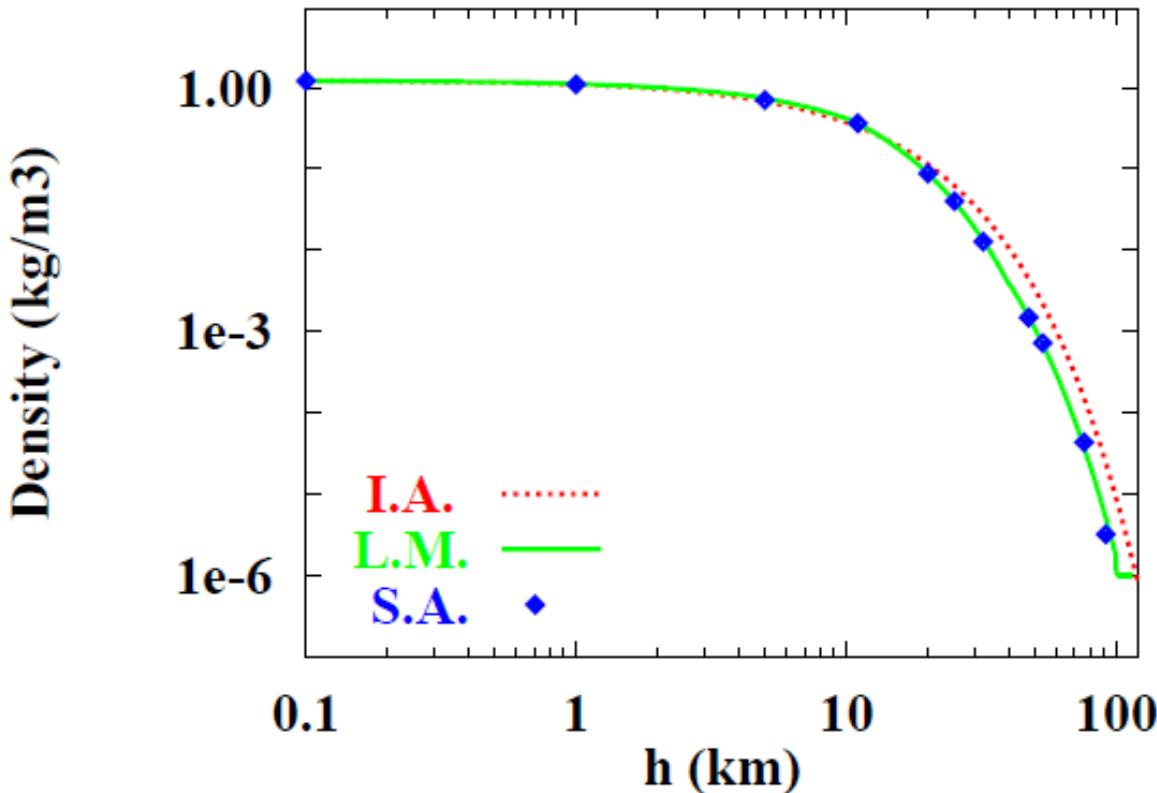


Figure 2.3. Density of the air as a function of the vertical altitude. The dots represent the US standard atmosphere data [14], while the full green line corresponds to Linsley's model [16] and the dashed red one to the isothermal atmosphere

$\rho(h) = \rho_0 e^{-gMh/RT}$ with
 $\rho_0 = 1.225 \text{ kg/m}^3$,
 $M = 28.966$ and $T = 288 \text{ K}$.

$$\rho(h) = \rho_0 e^{-gMh/RT}$$

Scale height $gM/RT \approx 9 \text{ km}$

The earth atmosphere

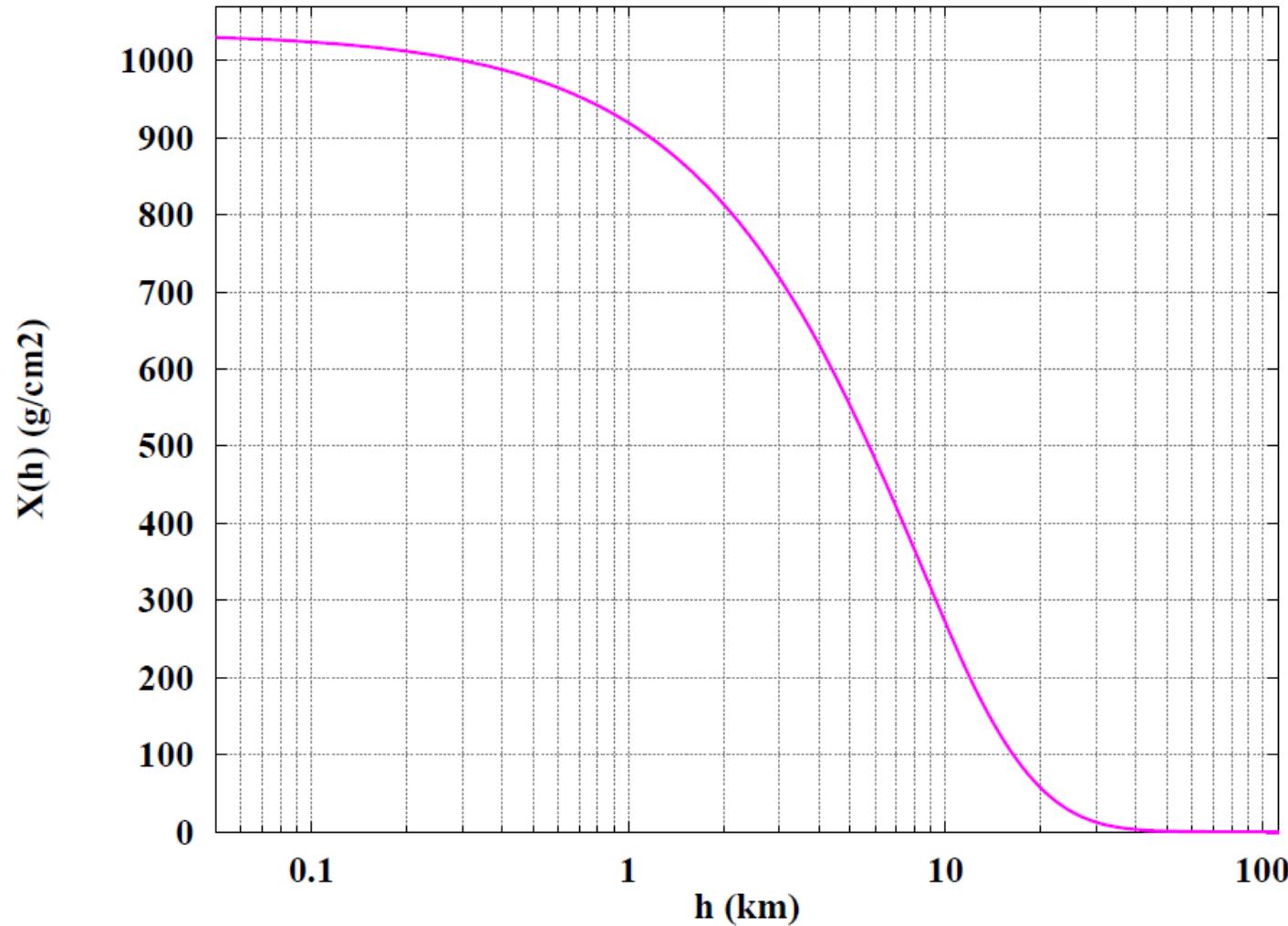
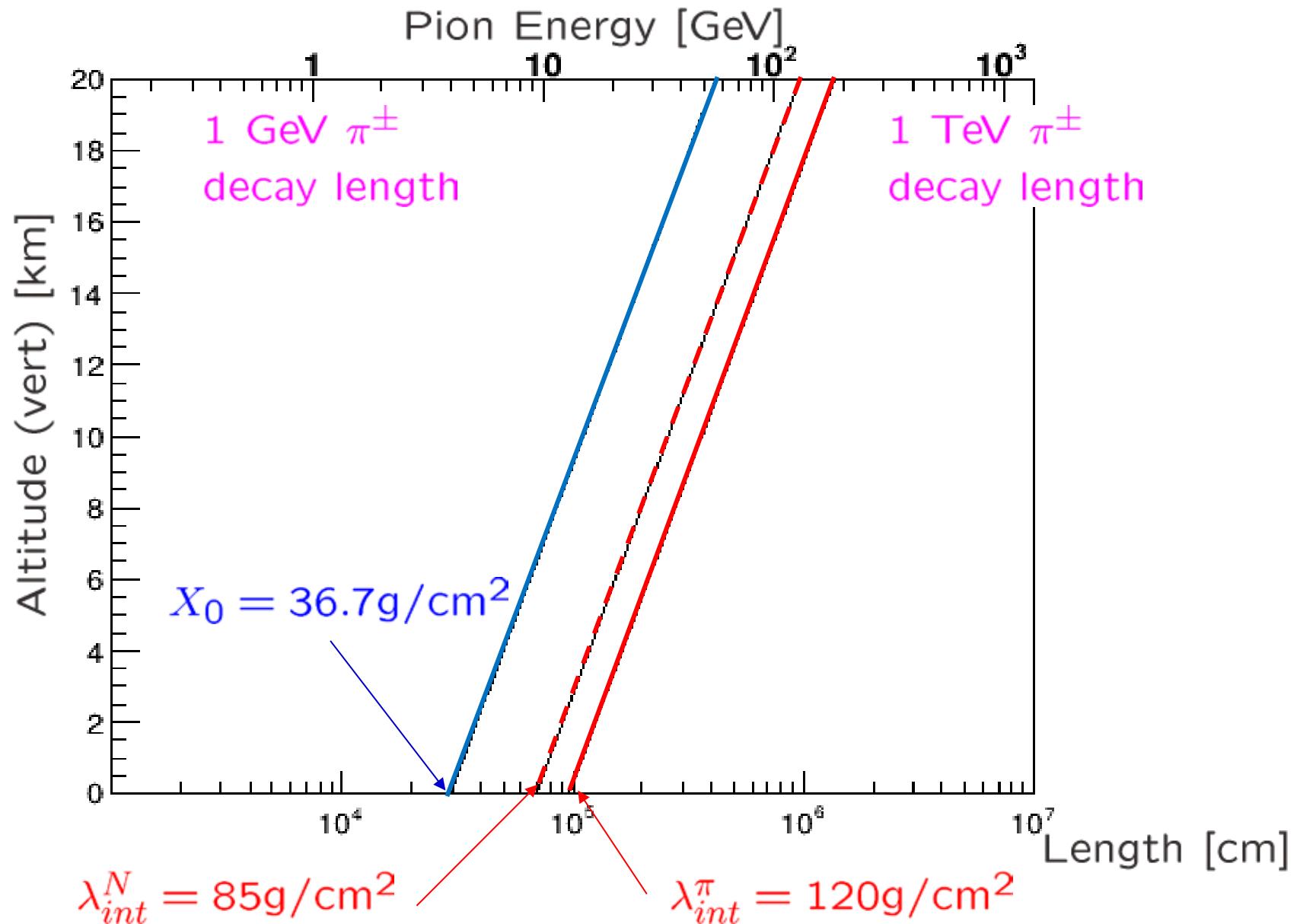


Figure 2.4. Vertical atmospheric depth, X_v , versus vertical altitude over sea level, h , accordingly with Linsley's model [16].

The earth atmosphere

- In terms of particle/radiation interaction with matter, the atmosphere is:
 - A total of $\frac{1}{4} 1000 \text{ g/cm}^2$ at sea level
 - So 1 atm ~ 12 interaction lengths ($\rho_N \frac{1}{4} 85 \text{ g/cm}^2$)
 - A vertical proton first interacts at $h \sim 15 \text{ km}$
 - One radiation length (at 1 atm) $X_0 = 36.6 \text{ g/cm}^2 \frac{1}{4} 300 \text{ m}$
 - One Moliere radius (at 1 atm) is $\frac{1}{2} R_M \frac{1}{4} 78 \text{ m}$
 - The Lorentz factor for a muon produced at $h = 10 \text{ km}$ to reach ground before decaying is $\gamma > 15 (> 1.6 \text{ GeV})$
 - Critical energy (EM) $E_c = 84.2 \text{ MeV}$

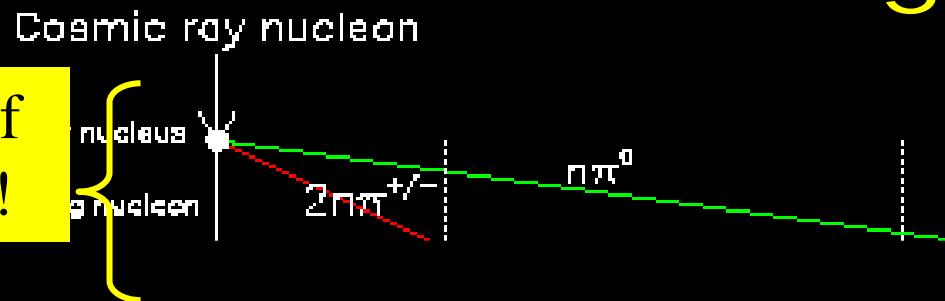
Interaction and radiation lengths in atmosphere



Shower development

er 2012

Physics well out of
reach of colliders !



At each step energy
is shared by more
numerous particles

"Hadronique"
shower

"Electro-Magnétique"
shower

$$\lambda_{interaction} < \beta\gamma c\tau_{decay}$$

$$\lambda_{radiation} < \lambda_{ionisation}$$

Maximum of
développement

Critical Energy E_c

No more
multiplication,
decrease by decay
and energy loss

$$\lambda_{interaction} > \beta\gamma c\tau_{decay}$$

$$\lambda_{radiation} > \lambda_{ionisation}$$

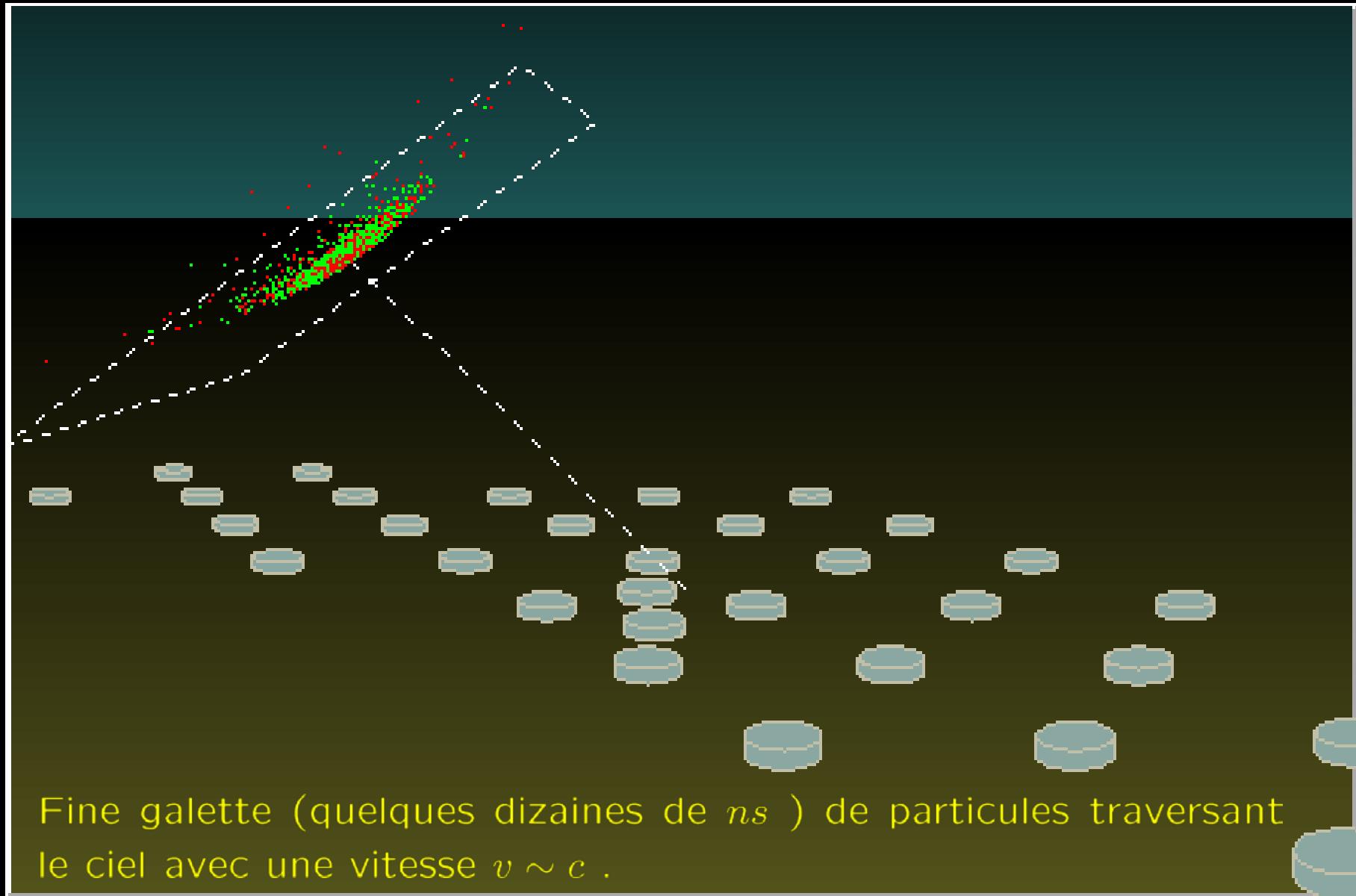
At ground, essentially
 $m^{\frac{1}{2}} g e^{\frac{1}{2}}$



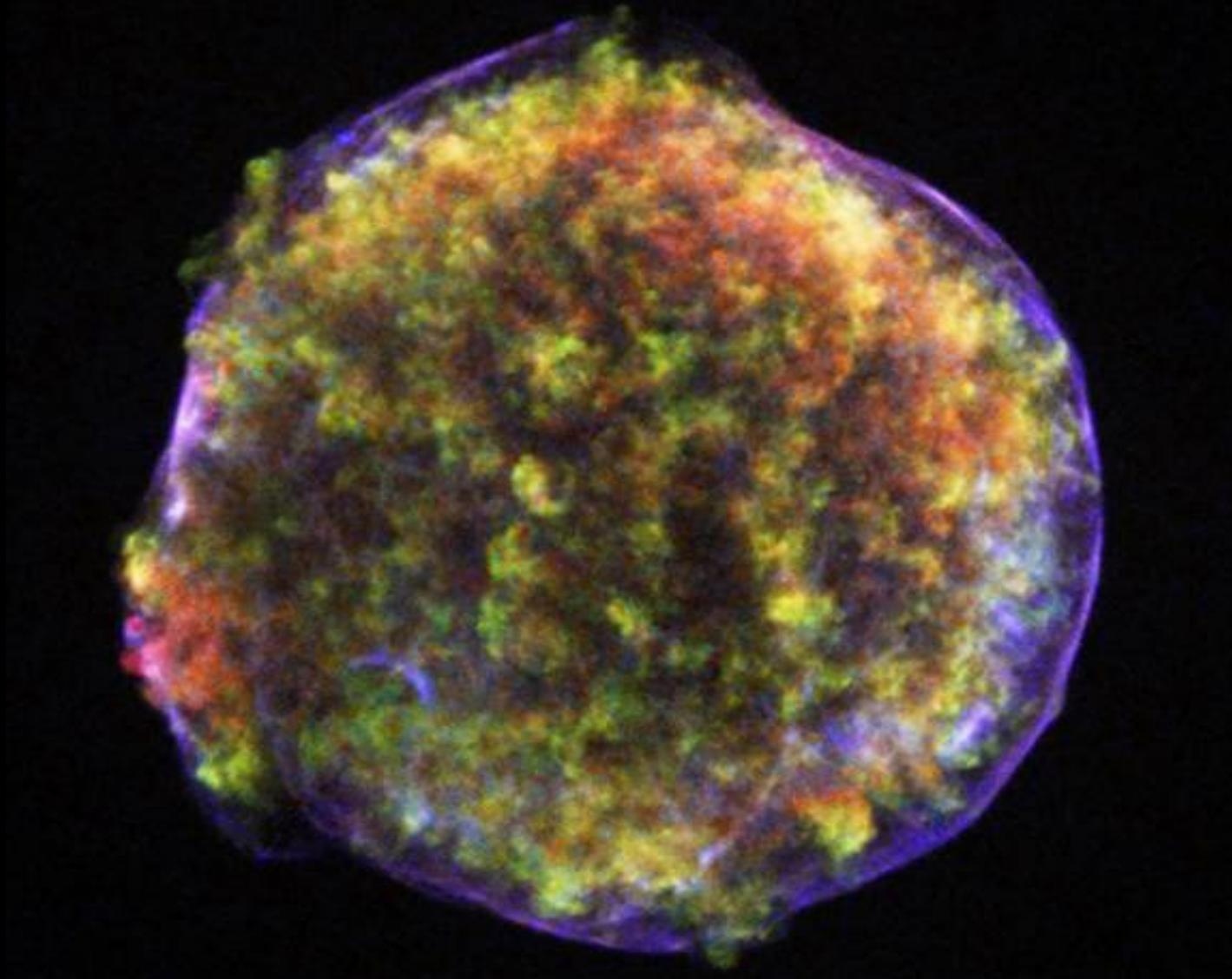
des RC"

F. Montanet "L2"

Structure en temps



Astrophysical Sources Cosmic Accelerators

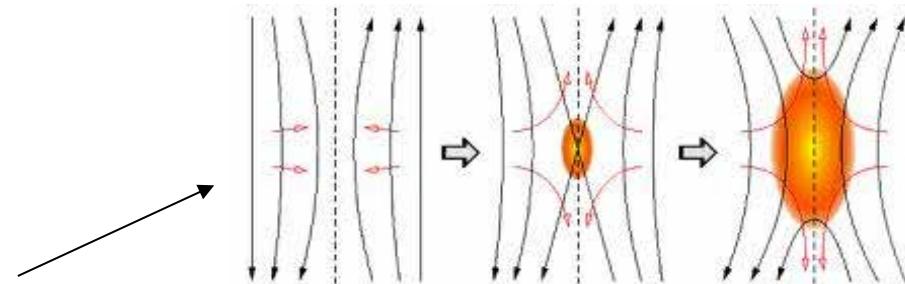


General ideas on acceleration

- Take the necessary energy somewhere...
 - Kinetic energy:
 - Translation: shock waves, moving clouds...
® Fermi acceleration
 - Rotation : pulsars, black holes, neutron stars
 - Gravitational energy
 - via accretion (® jets...): accretion disks (divers)
 - Electromagnetic (EM)
 - From turbulence, from compression, or from rotating magnets...
- In fine, charged particles interact with EM fields: $f = q(E + v' B)$
- Remember:
Astrophysical shocks are collisionless.
® Energy transfert through EM fields !

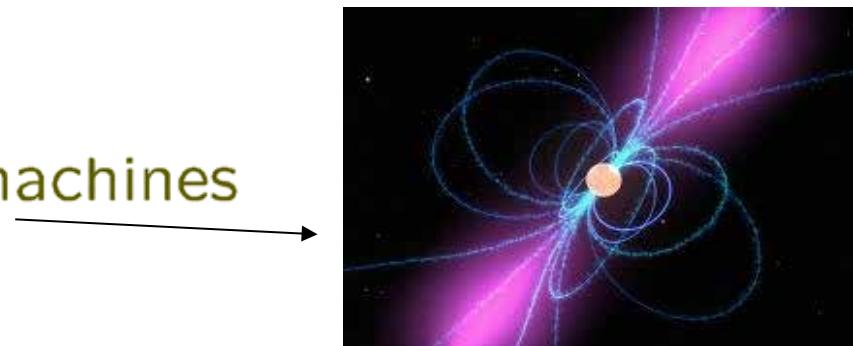
E and B fields in Universe

- * In ISM as on earth, $\langle E \rangle \approx 0$
 - ISM is neutral or conducting



- * Transient electric fields:
 - Magnetic re-connections (e.g. solar flares...)
 - EM waves

- * Producing E fields in EM "engines"
 - Astrophysical dynamos, induction machines



- * Magnetic fields :
 - $\epsilon_B \approx 1\text{eV/cm}^3 \approx \epsilon_{optique} \approx \epsilon_{CMB} \approx \epsilon_{CR}$!!!
 - Astrophysical plasmas :
ISM, stars, accretion disks, IGM, jets, etc...

Magnetic field production

- Large scale movements of ionized media
 - ® generating magnetic fields, magnetized clouds...
- Turbulence in interstellar medium
 - ® Magnetic turbulence, inhomogeneous B fields, plasma waves...
- Hydro and MHD instabilities
 - e.g. Rayleigh-Taylor in supernova remnants
- "Streaming" instabilities
 - CR generates waves in a magneto-actif plasma
 - ® creating the conditions for their own diffusion

Magnetic field production

In many cases, **equipartition** can be reached

- for ex: behind a shock wave :
thermal en. ~ kinetic en. ~ magnetic en.

- Energy exchange between macroscopic structures and individual particles
- individual particles may reach very high energies!

Magnetic fields and acceleration !

- **How is it possible at all?**

magnetic fields don't work ! ($\vec{F} \perp \vec{B}$)

- Well, variable $\vec{B}(t)$ fields do ! (example: *Betatron*)

$$\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

- In a different reference frame,

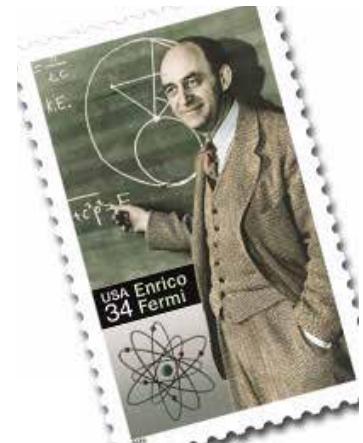
a pur \vec{B} field is feeled as a \vec{E} field...

$$\vec{E}' = \vec{v} \times \vec{B} \text{ (for } v \ll c\text{)}$$

- In principle, one can always identified an effective \vec{E} field that works, but the description in terms of \vec{B} fields is often simpler (and more physical !)

® Acceleration by "change of reference frame"

Principle of Fermi acceleration



The Ingredients :

- **A magnetic field \vec{B}**
= with a regular component \vec{B}_0
+ and irregular component $\delta\vec{B}$
- **A plasma** i.e. a good electrical conductor :
 $\vec{E} + \vec{u} \times \vec{B} = 0$ and $|E| \approx 0$
⇒ the magnetic field is "frozen" and moves with the plasma (Alfvén).
- **A CR population** coupled to the medium via the magnetic field \vec{B} . They scatter on the field irregularities. This diffusion processes are **collisionless** i.e. they conserve the particle energy. The MHD or Alfvén waves act as massive scattering centers (recoilless).

Fermi 1949 :

- first hypothesis of converging movements of MHD perturbations
⇒ "first order" acceleration, but where ?
- second more realistic hypothesis at that time: random movement of interstellar gas clouds (observed) or MHD perturbations
⇒ "second order acceleration."

Where to accelerate

- At creation :
 - For example: e^- extracted from the surface of a neutron star by an intense E field.
- Within the source neighborhood :
 - For example: Fermi acceleration in plasma shocks in a SNR.
- During transport:
 - "reacceleration" by shock waves and excitation of Alfvén waves during diffusive transport in the Galaxy.

Power laws and stochastic processes

- The power laws observed in differential energy spectra follow naturally from cyclic acceleration mechanisms with constant energy gain and constant escape probabilities:
 - Initial energy: E_0
 - Energy gain at each cycle: $\Delta E = \varepsilon E$
 - Particle energy after n iterations: $E_n = E_0(1 + \varepsilon)^n$
 - Escape probability from the acceleration zone: P_{esc}
 - Probability to remain in the acceleration zone: $(1 - P_{esc})^n$

Power laws and stochastic processes

- Particle energy after n iterations: $E_n = E_0(1 + \varepsilon)^n$
- Probability to remain in the acceleration zone: $(1 - P_{esc})^n$

Number of iterations to reach an energy E :

$$n = \frac{\ln(E/E_0)}{\ln(1 + \varepsilon)}$$

Proportion of particles accelerated up to an energy equal or greater than E :

$$N(\geq E) = N_0 \sum_{m=n}^{\infty} (1 - P_{esc})^m = N_0 \frac{(1 - P_{esc})^n}{P_{esc}}$$

thus :

$$\frac{\ln(P_{esc}N/N_0)}{\ln(1 - P_{esc})} = n = \frac{\ln(E/E_0)}{\ln(1 + \varepsilon)}$$

eliminating n :

$$N(\geq E) \propto \left(\frac{E}{E_0}\right)^{-\gamma}$$

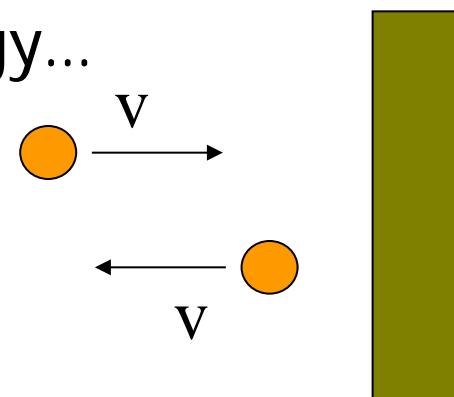
with $\gamma = \frac{-\ln(1 - P_{esc})}{\ln(1 + \varepsilon)} \approx \frac{P_{esc}}{\varepsilon} = \frac{1}{\varepsilon} \frac{T_{cycle}}{T_{esc}}$



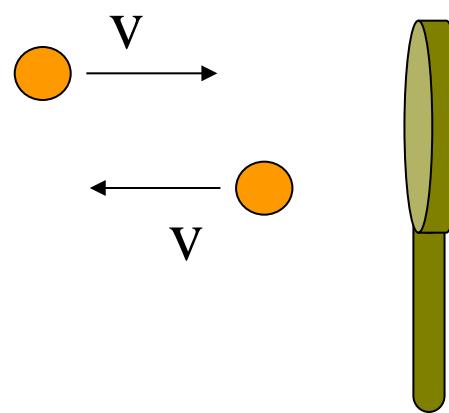
Power laws
are natural !

A small analogy...

- A tennis ball bouncing on a wall
 - neither gain nor loss of energy...



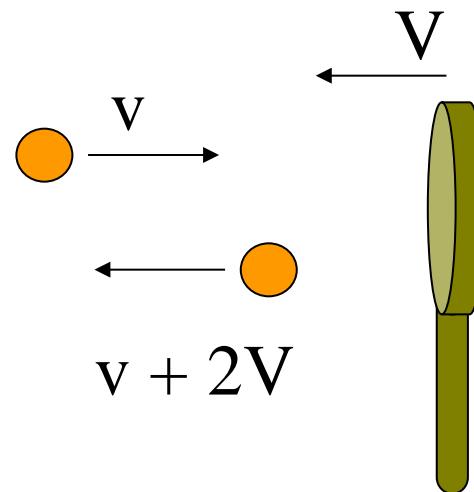
bounce = speed unchanged



Same thing with
a motionless racquet...

Then how does one accelerate a tennis ball ?!

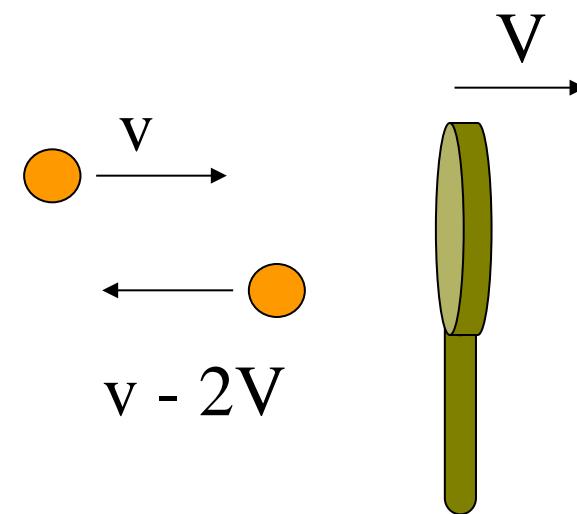
- Moving racquet
 - Neither gain nor loss of energy... in the racquet reference frame !



Speed unchanged with
respect to the racquet

® acceleration through a change of reference frame

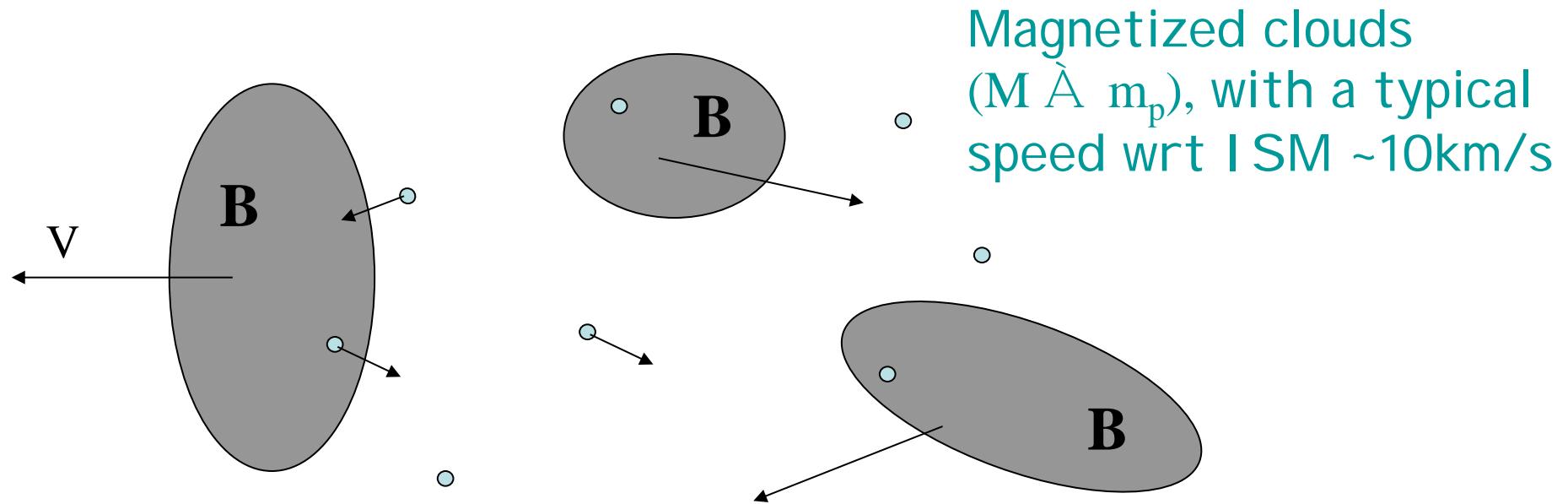
- A drop shot:



Particle deceleration !

Fermi Acceleration

- Ball ® charged particle
- Racquet ® "magnetic mirrors"



- Magnetic inhomogeneities or plasma waves also work...

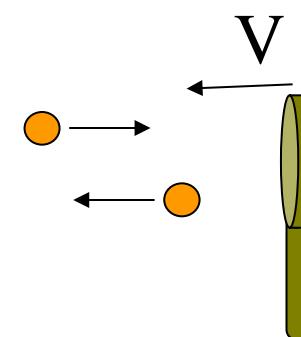
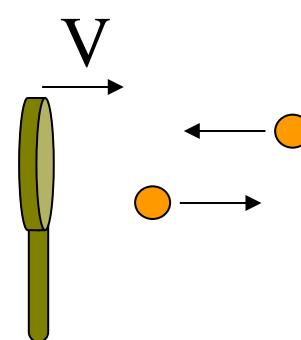
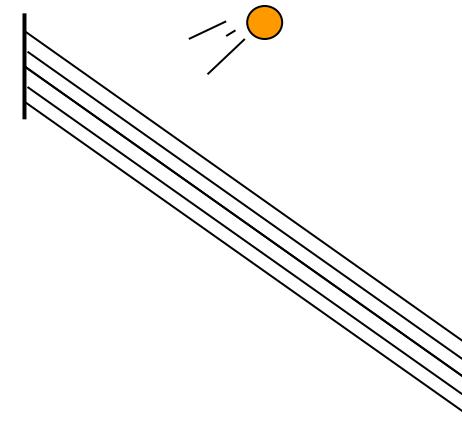
The essence of stochastic de Fermi acceleration

- 1 When a particle bounces on an **incoming** magnetic mirror, in a **head-on** collision, it **gains** energy.
- 2 When a particle bounces on a **receding** magnetic mirror that it catches back, it **loses** energy.
- 3 Head-on collisions are **more frequent** than receding collisions.

Þ Net energy gain in average (stochastic process)

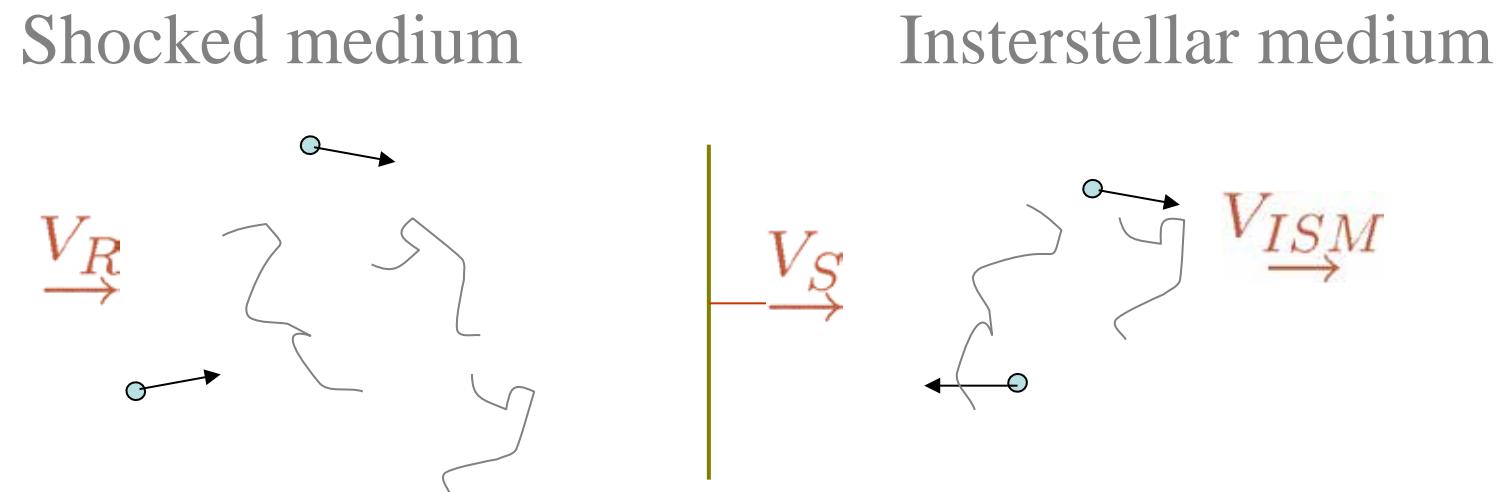
Add a second player...

- Converging flows...



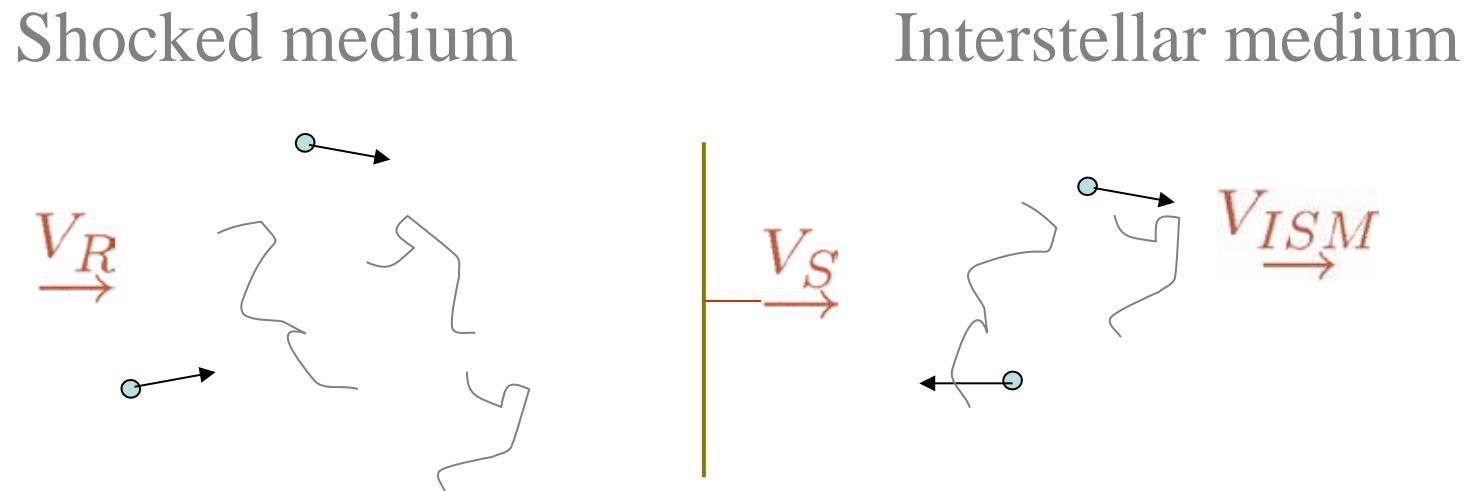
Shocks hydrodynamics

- Shock waves (e.g. supernova explosion) : expending plasma flow with a speed V_R much larger than the sound speed in the interstellar medium (ISM).



Shocks hydrodynamics

- Shock wave:

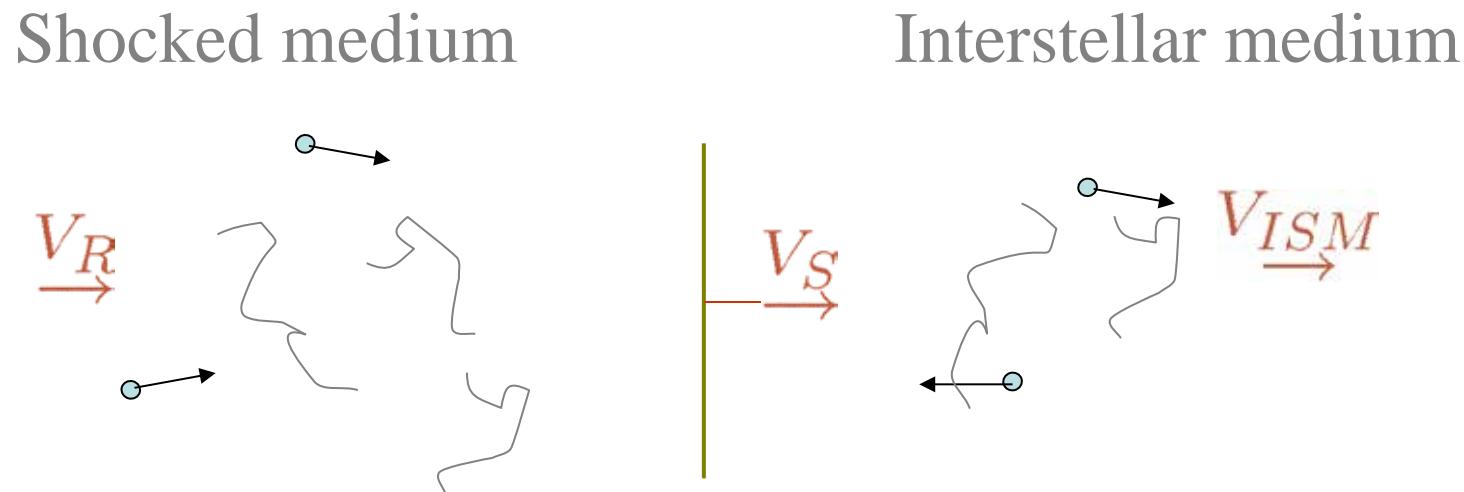


- The shock moves at a speed V_S which depends on V_R and the specific heat of both media.
- For an ionized ISM:

$$V_S \approx \frac{4}{3} V_R$$

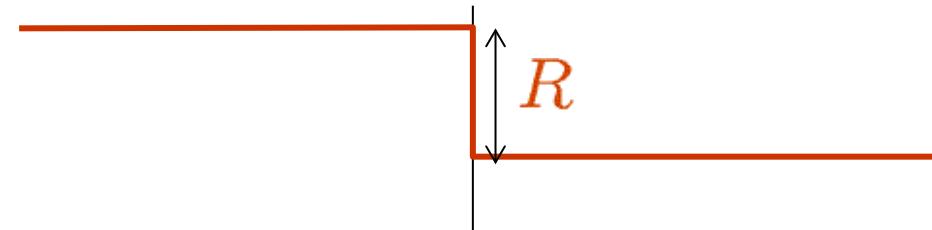
Shocks hydrodynamics

- Onde de choc :



- The shock intensity is characterized by the compression factor:

$$R = \frac{V_S/V_R}{V_S/V_R - 1} \approx 4$$

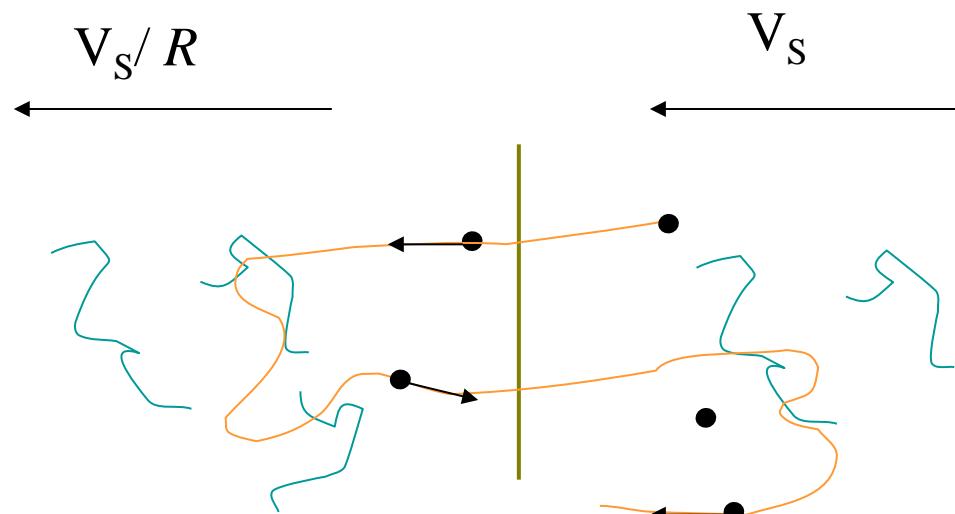


Shocks hydrodynamics

In the shock frame

Shocked medium

Interstellar medium

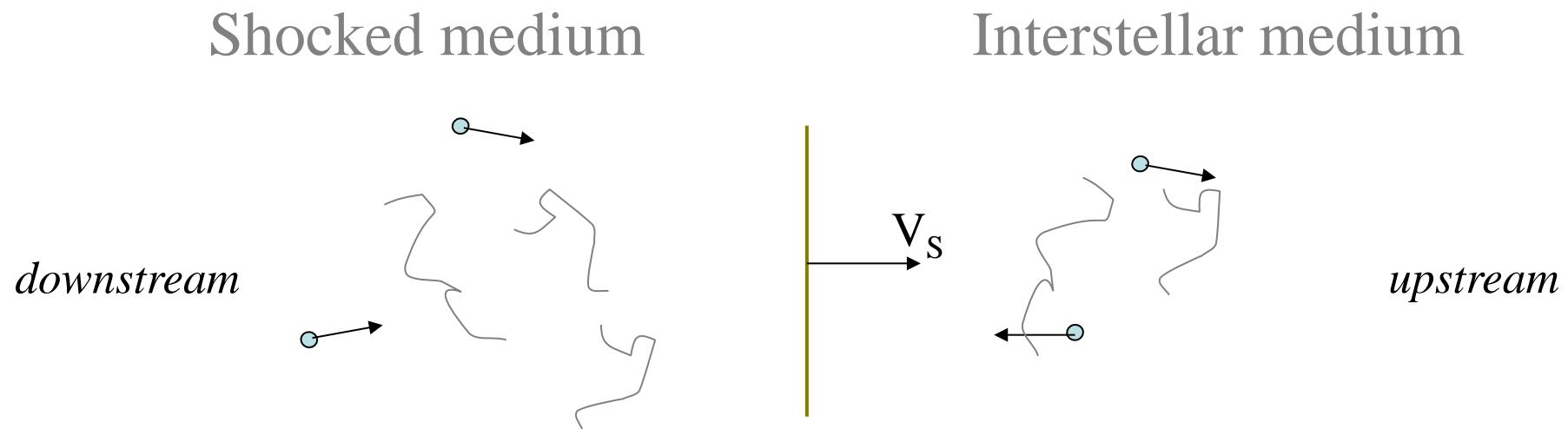


- In the shock frame, the upstream (non-shocked) medium flows toward the shock at a speed V_s and the downstream (shocked) medium flows away with a speed reduced by the compression factor (mass flow conservation) :

$$V_s / V_d = R \approx 4$$

Acceleration diffusive par onde de choc

- Onde de choc (e.g. explosion de supernova)

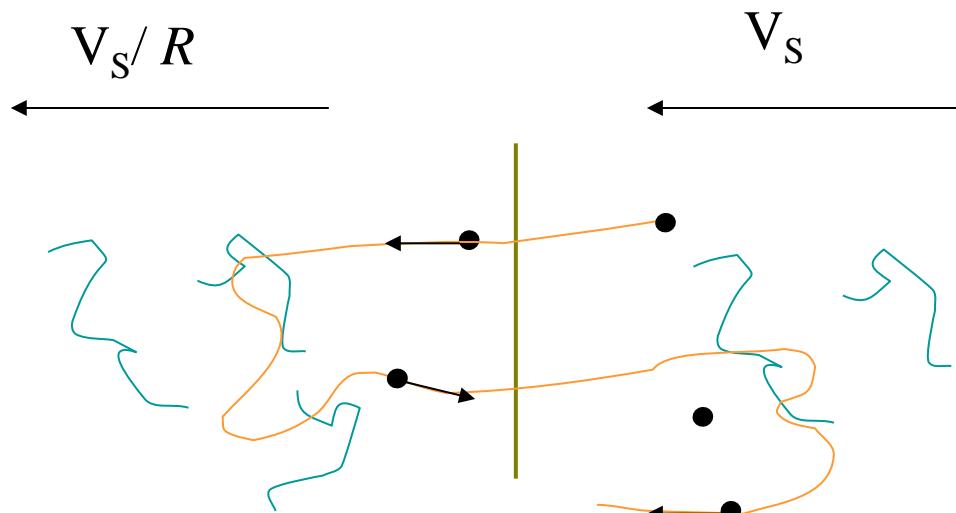


- Magnetic wave generation:
 - **Downstream** : by the shock (compression, turbulence, hydro and MHD instabilities, shear, etc.)
 - **Upstream** : by the accelerated cosmic rays themselves !
- ® 'isotropization' of the distribution
(in the local frame)

A win win process !

Shocked medium

Interstellar medium



In the shock frame

- At each shock crossing, one way or the other, the particle hits a "magnetic wall" with a relative speed:

$$V = (1 - 1/R)V_S$$

⑧ only head-on collisions...

Summary on acceleration

- Acceleration from interaction with fields
 - **E** field: e.g. induced by spinning magnets such as neutron stars (pulsars) or black holes...
 - **B** field: inhomogeneous moving fields
 - MHD waves
- Acceleration by reference frame transformation
 - Fermi stochastic acceleration (2nd order)
 - Diffusive shock acceleration diffusive (1st order)
- Power law are natural
 - Fermi type process ($\Delta E \propto E$, P_{ech})
 - Universal power law for non relativistic shocks ($N(E) \propto E^{-2}$)
- Cosmic rays up to the knee
 - CR power = power of SNe, $E_{\text{max}} \approx 10^{14} \text{ eV}$ hardly 10^{15} eV

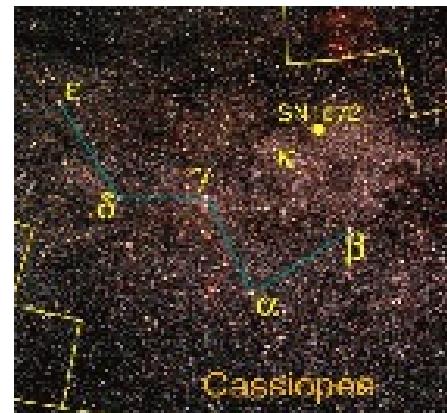
The CR standard model

- Analytic calculations, simulations and observations show that diffusive shock acceleration works !
- Supernovae and GCRs
 - Estimated efficiency of shock acceleration : ... **10 – 50%**
 - Power required to sustain CR energy density: $\frac{\varepsilon_{CR} \times V_{conf}}{\tau_{conf}}$
 $\sim 10^{41} \text{erg/s}$!
 - Power injected by SN power in the Galaxy: **10^{42}erg/s !**

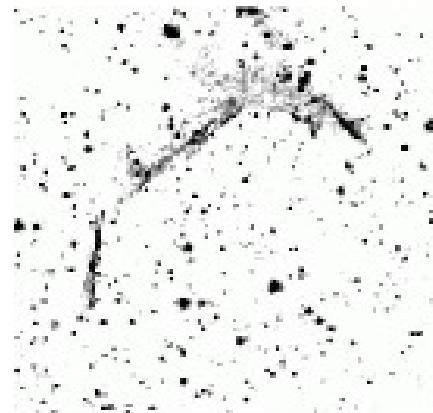
® Enough power for Galactic CR

Tycho, 11 November 1572...

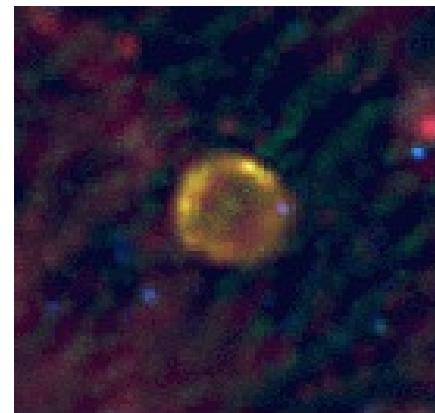
2009-2010



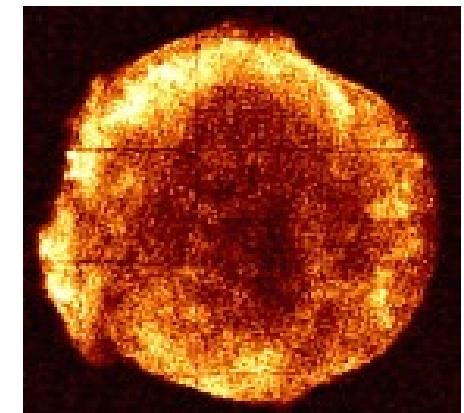
position



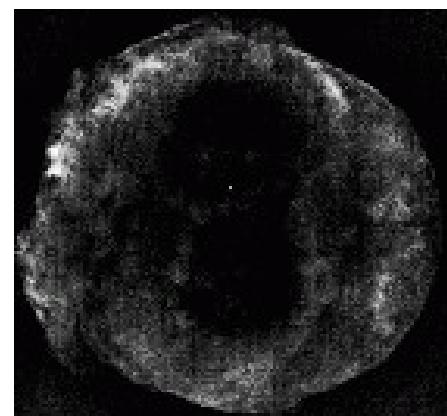
visible



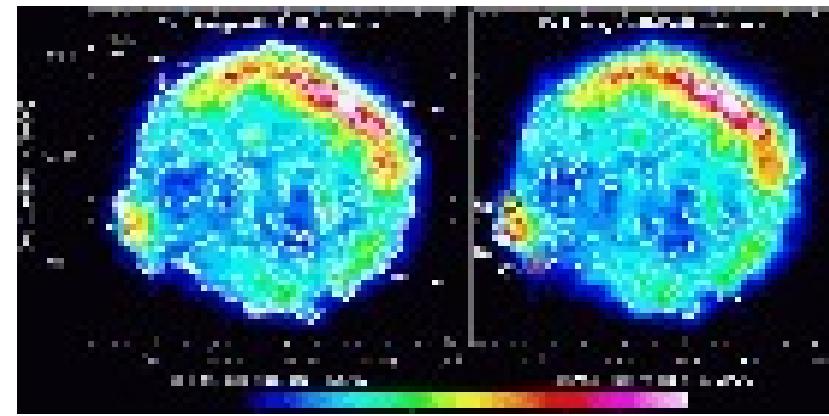
IRAS



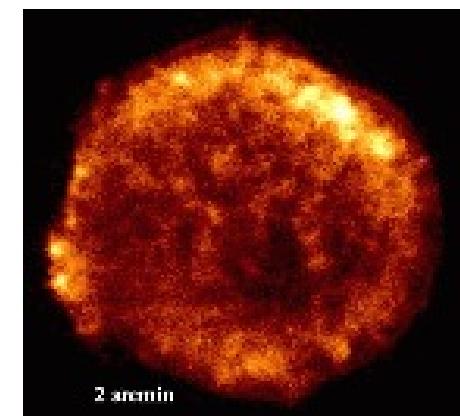
Km (VLA)



6 cm (VLA)



Si K (XMM) Fe K



X (ROSAT)

The CR standard model

- Proposed acceleration site, **isolated SNR**
 - Supernovæ : ejection of many solar masses of nuclear matter at supersonic speeds ($\sim 10\,000$ km/s) following massive star explosion.
 - Formation of a quasi spherical expanding shock wave that wipes out the interstellar medium (ionized beforehand by the progenitor's radiation).
 - Total kinetic energy injected by the explosion: 10^{51} erg ($= 10^{44}$ J).
 - Roughly 3 SNe per century within our Galaxy, which corresponds to an averaged power of 10^{42} erg/s (10^{35} W)
 - SNR are observed at all wavelength.
 - SNe explosion is essential to the Galaxy chemical content: heavy elements enrichment.

The CR standard model

- Shock waves in isolated SNe (SNR)
 - Source composition source ~ interstellar medium + modifications (ionizability, volatility, Z/A effects, ^{22}Ne ...)
 - Source spectrum: E^{-2} power law
 - Maximal energy reached: $E_{\max} \sim 10^{14}$ eV
- Energetics :
 - Measured flux / speed = CR density
 - CR density ' mean energy = energy density
 - Energy density ' confinement volume = total energy
 - Total energy / confinement time = necessary injected power
 - $P_{\text{CR}} \sim 1.5 \cdot 10^{41}$ erg/s
- Required efficiency $\sim 10\text{-}30\%$...

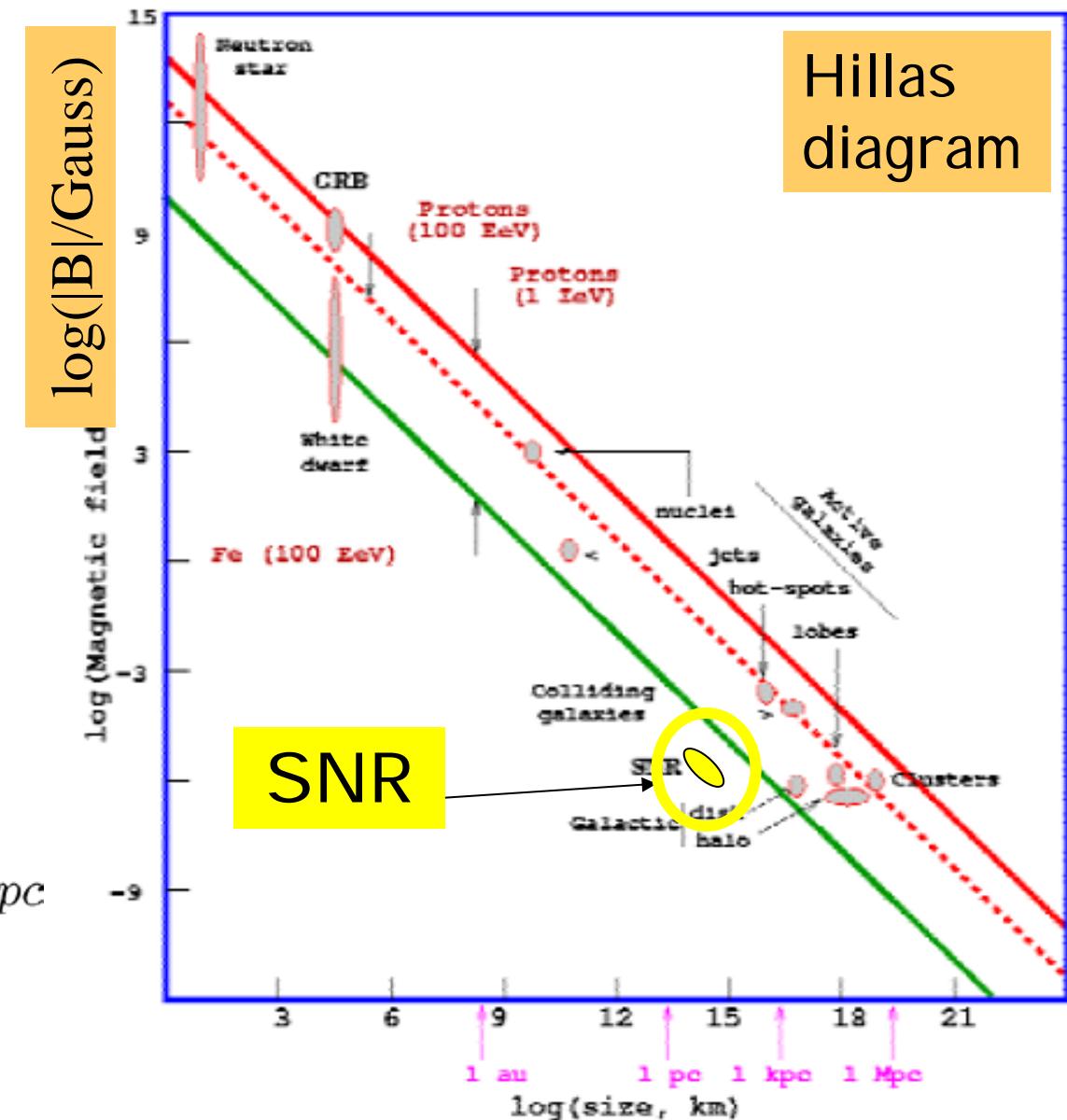
Finite size of confinement magnetic field



Larmor radius

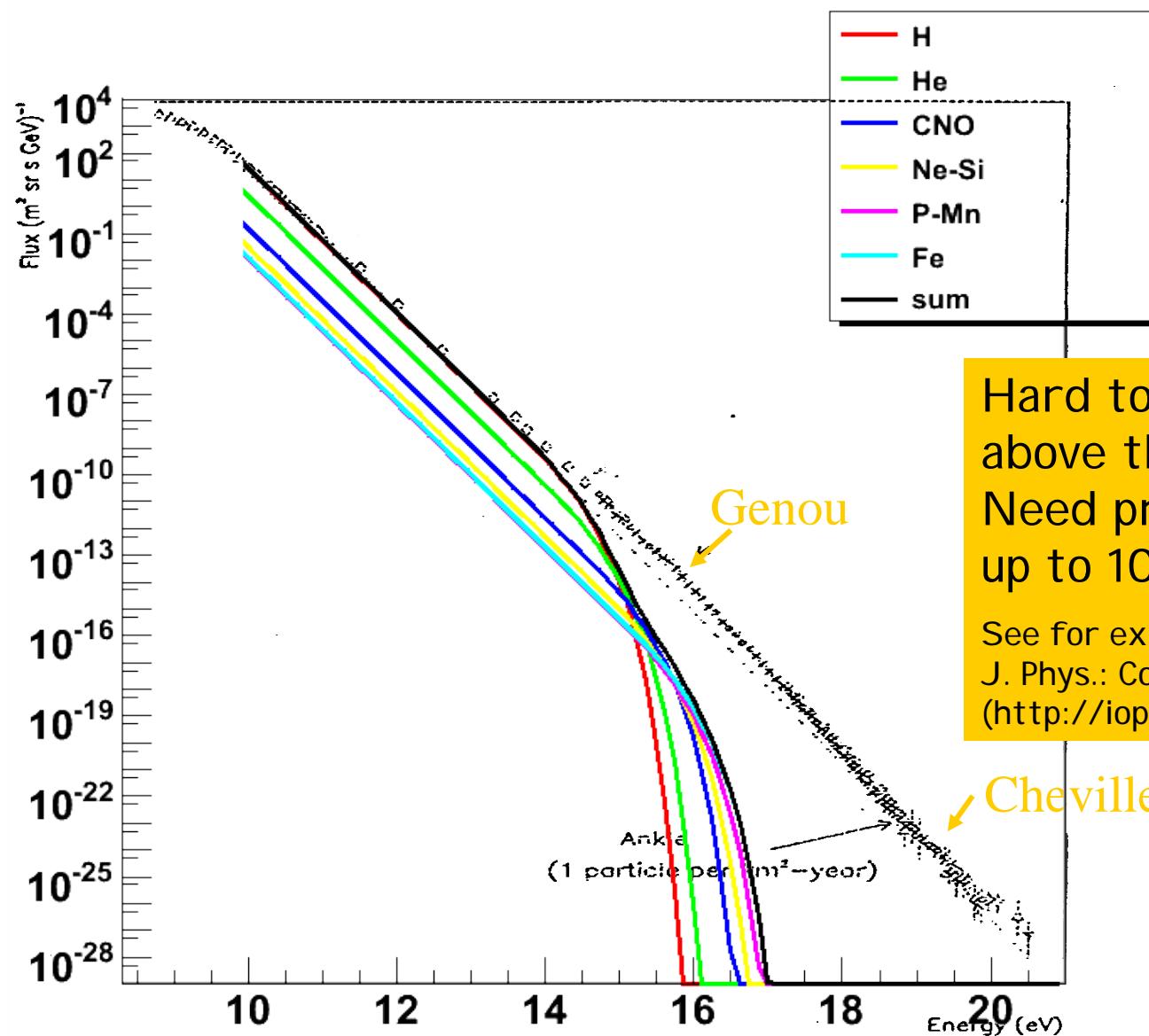
$$\begin{aligned} r_g \leq R \\ \Leftrightarrow \frac{p}{ZeB} \leq R \\ \Leftrightarrow E \leq Ze \times B \times R \\ \Leftrightarrow E \leq (10^{17} eV) Z B \mu G R_{pc} \end{aligned}$$

Finite size



$\log(R/\text{km})$

CR SM with Emax / Z

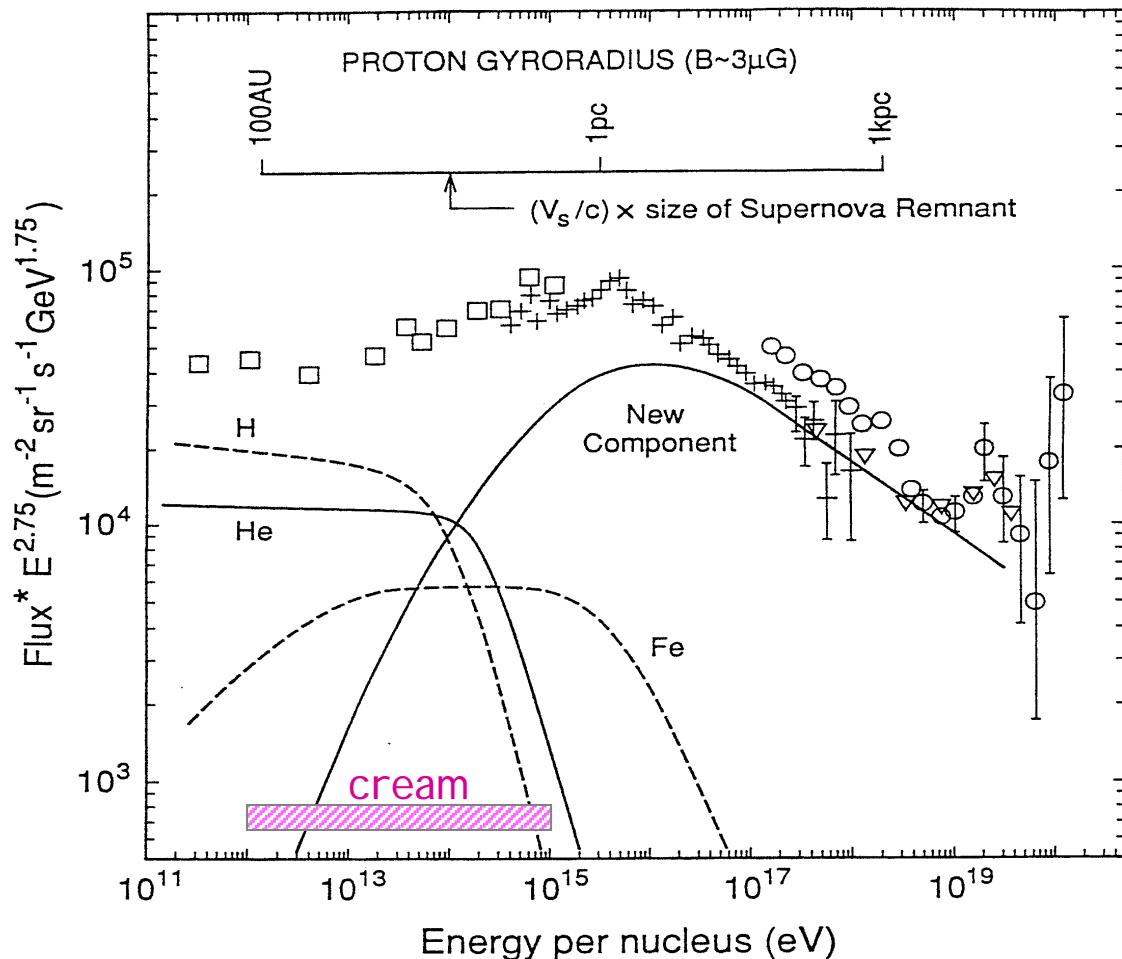


Hard to explain a single power law above the knee.
Need probably another component up to 10^{17}eV

See for ex : M. Hillas
J. Phys.: Conf. Ser. 47 (2006) 168
(<http://iopscience.iop.org/1742-6596/47/1/021>)

The knee

→ 2010



- Is the knee simply the consequence of an energy cut-off of accelerators (SNR) or is it due to :
 - a propagation effect ?
 - different CR sources ?
 - a physics threshold at $\sim E_{LHC}$?

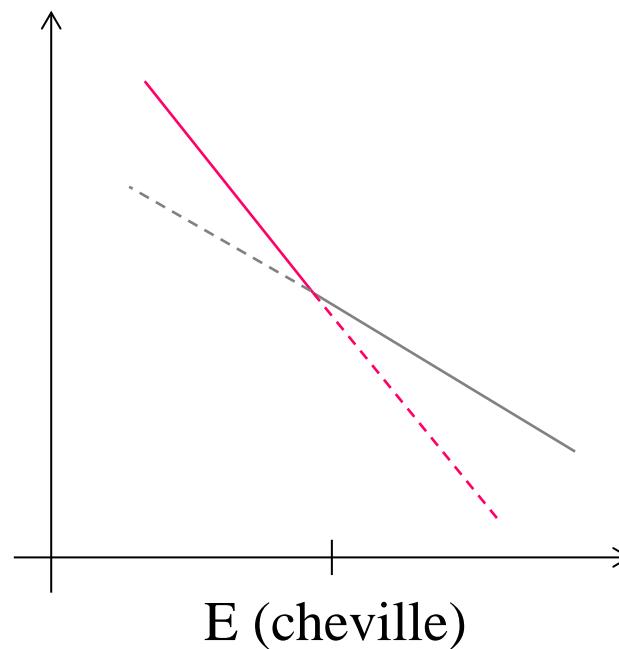
— The CR SM implies a change in composition at the knee energy.

SNR energy limit : $E_{max} \sim Z \cdot 10^{15}$ eV

- Data:
- Proton4 sat
 - Akeno
 - Yakutsk
 - Haverah Pk

Explaining the ankle is much easier...

- Two components with two different slopes...



- For exemple, galactic and extragalactic...

DIFFUSE GAMMA-RAY SOURCE

VHE gamma-rays sources

° -rays production processes:

- Electro-Magnetic processes :

- Bremsstrahlung

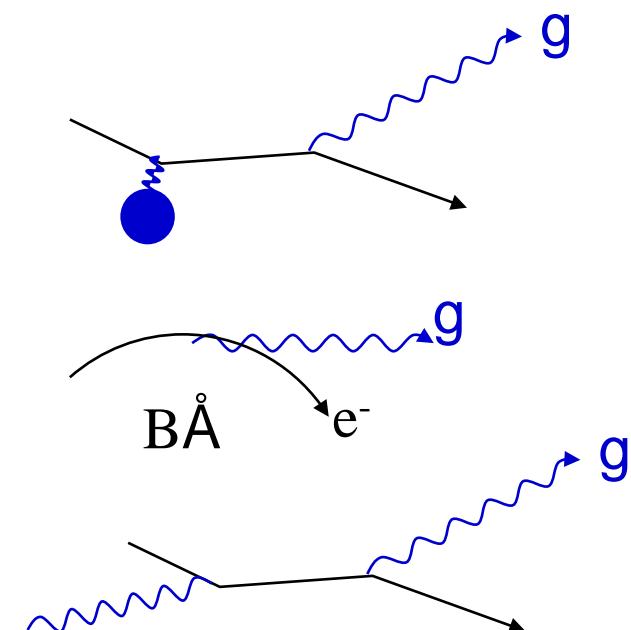
- Synchrotron radiation

- Inverse Compton

- Hadroniques processes:

$$(p, N) + (N\gamma) \rightarrow \pi_0 + \dots$$

$\hookrightarrow \gamma\gamma$



Hadronic production of gamma-rays

Source function [TK Gaisser]

$$q_k(E_k, r) = \int \frac{d\sigma_{i \rightarrow k}(E_k, E_i)}{dE_k} \left(\frac{c\rho(r)}{m} \right) \left(\frac{4\pi}{c} \phi(E_i) \right) dE_i$$

Fluxes on earth (neutrals)

$$\phi_{k=\gamma} = \frac{dN_k}{dAdE_k d\Omega} = \int \frac{q_k(E_k, r)}{4\pi r^2} d^3r = \int_0^{r_{max}} \frac{q_k(E_k, r)}{4\pi} dr d\Omega$$

Scale invariance : $E_\gamma/E_i = Z$ (Z indep. of E)

(for $E_k \gg m_i/2 = m_\pi/2 \approx 70\text{MeV}$)
(system mass scale).

$$\left. \right\} \Rightarrow \phi_k \propto \phi_i$$

Parent spectral density power law $\phi_i \propto E_i^{-\alpha}$
[Gaisser,Hansen,Berezinsky,Stanev]

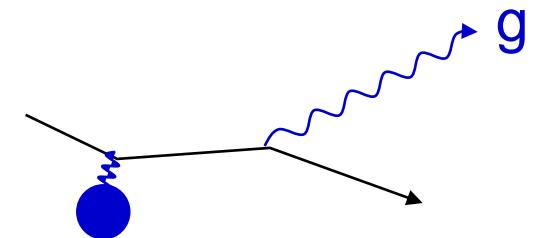
$$\phi_\gamma = 4\pi\rho \left[\frac{\sigma_{inel}}{m_N} \right] \left[\frac{2Z_{N \rightarrow \pi^0 \rightarrow \gamma}}{\alpha} \right] \phi_{CR}(E)$$

$$\frac{\phi_\gamma}{\phi_{CR}} \approx 6 \times 10^{-4} \times \left(\frac{\rho R}{[g.cm^{-2}]} \right) \quad \rho R = \text{target column density}$$

Bremsstrahlung

$$\frac{d\sigma_{e \rightarrow \gamma}(E_\gamma, E_e)}{dE_\gamma} = \frac{1}{E_e N_A X_0} \phi(z)$$

where $\phi(z)$ is a function of $z = E_\gamma/E_e$



For energies $> 70 MeV/c^2$, there is a scaling relation between the γ energy and that of the parent electron. (the only mass scale in the problem is the electron mass, much lower than 70 MeV).

⇒ the differential spectrum of the progenitor electrons ($\phi_e(E_e) \propto E^{-\alpha}$) is transmitted to the γ differential spectrum:

$$\phi_\gamma(E_\gamma) \propto E_\gamma^{-\alpha}$$

For $dN/dE_e = a_e E_e^{-\alpha}$ and $\alpha = 2.7$

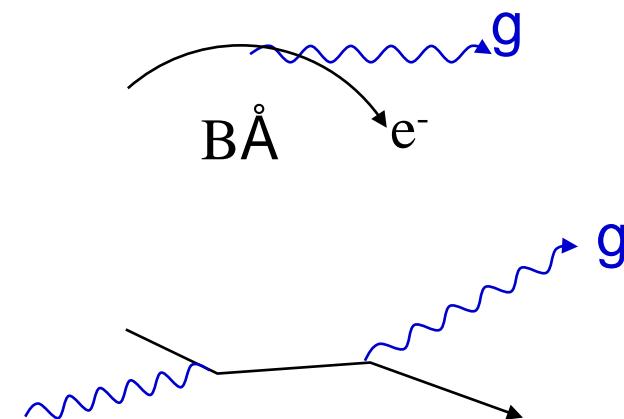
$q_{br} \approx 1.2 \times 10^{-25} a_e n E_\gamma^{-2.7}$ photons $\text{GeV}^{-1} \text{s}^{-1} \text{cm}^{-3}$

VHE Gamma sources

- Productions of $g > 10 \text{ MeV}$
 - Synchrotron radiation negligible.
 - Inverse Compton

For $\phi_e \propto E^{-\alpha}$

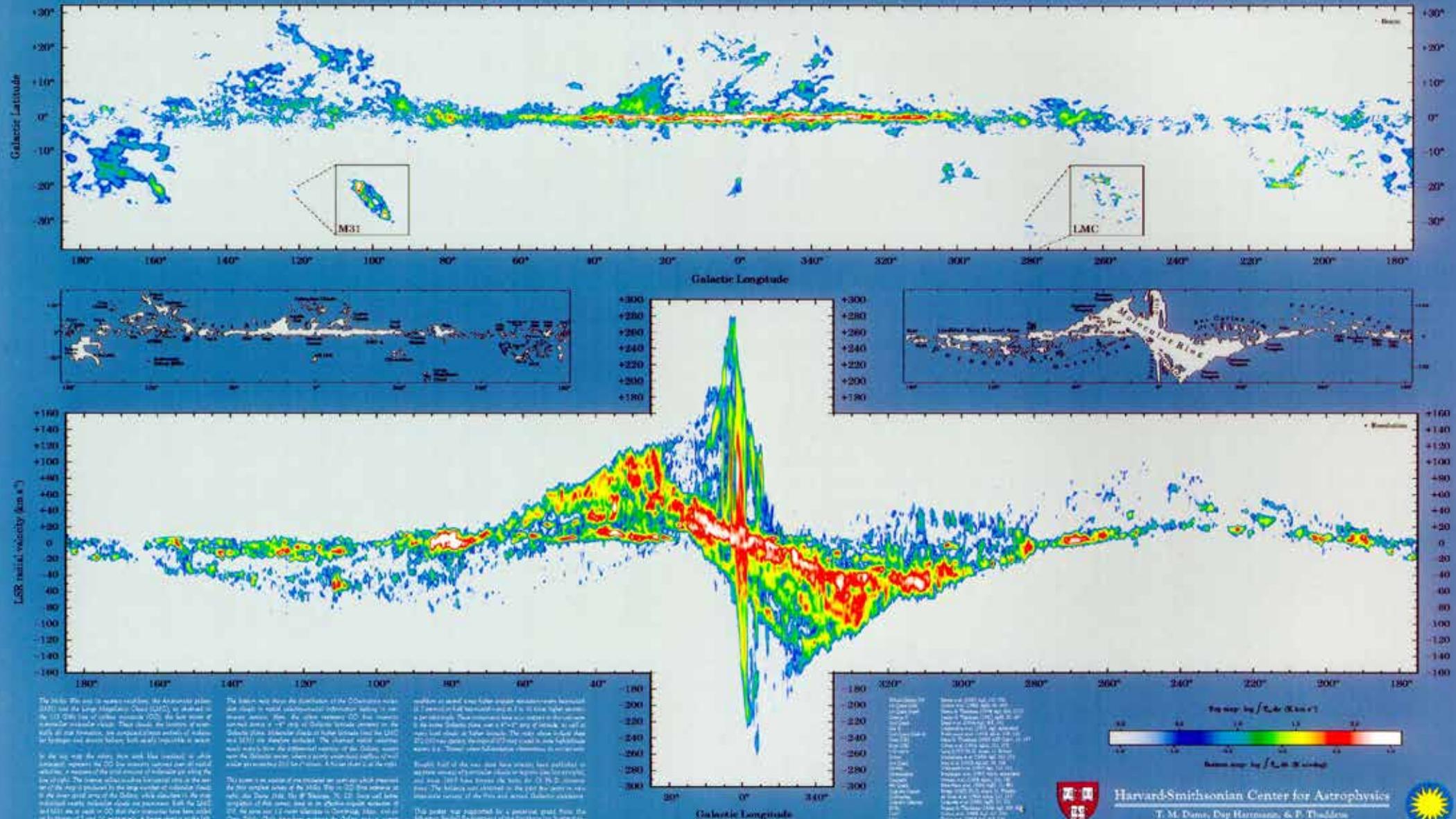
the IC spectrum is much flatter
 $\propto E^{-(\alpha+1)/2}$



but $N_e \ll N_p$ and the electron spectrum drops \downarrow above a few GeV
(synchrotron radiation energy losses)

The target...

The Milky Way in Molecular Clouds



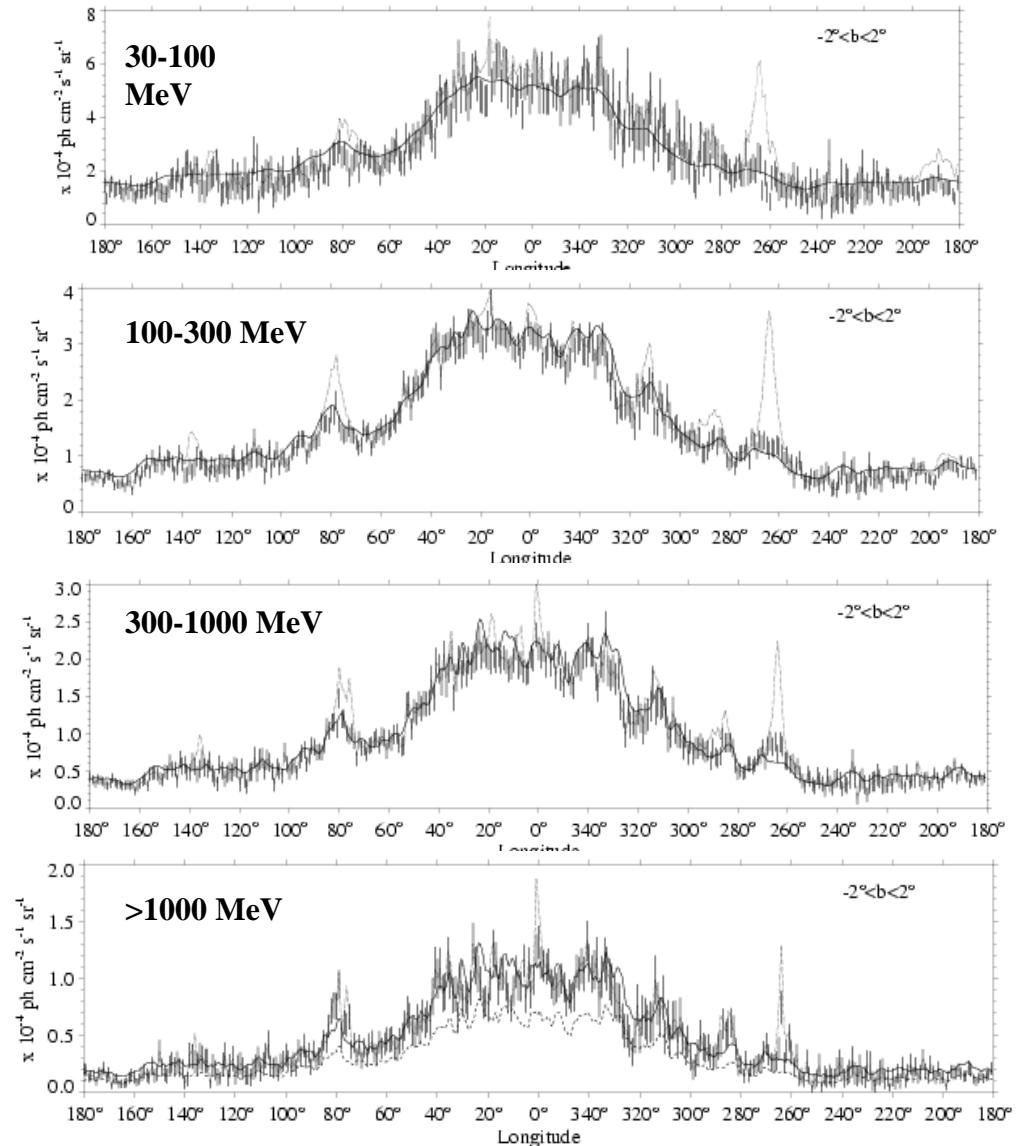
Harvard-Smithsonian Center for Astrophysics
T. M. Dame, Dag Hammarskjöld, & P. Thaddeus



Tracking back the CR flux elsewhere in the Galaxy

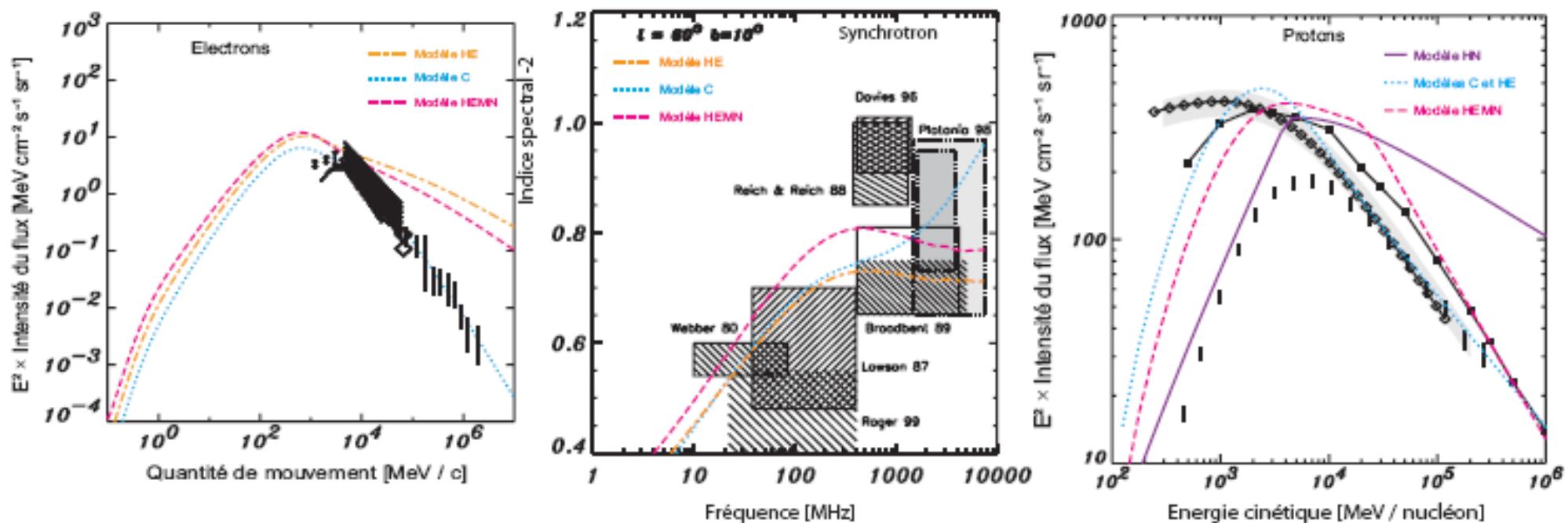
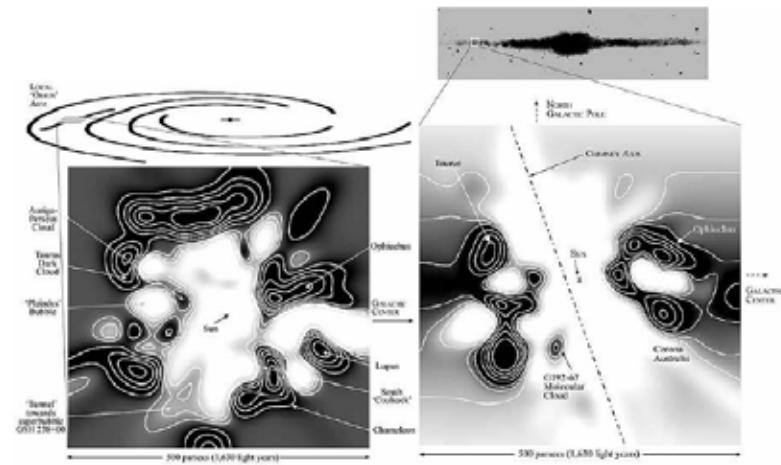
- Knowing the target column density from radio 21cm measurements and fitting the diffuse g flux, one can map the density, spectrum and composition of CR elsewhere in the galaxy.

- S.D.Hunter et al,
ApJ 481 (1997) 205
- M.Pohl and J.A.Esposito
ApJ, 507 (1998) 327
- S.LeBohec et al,
astro-ph/0003265
- Strong, A.W., Moskalenko, I.V.,
ApJ 509:213-228, 1998



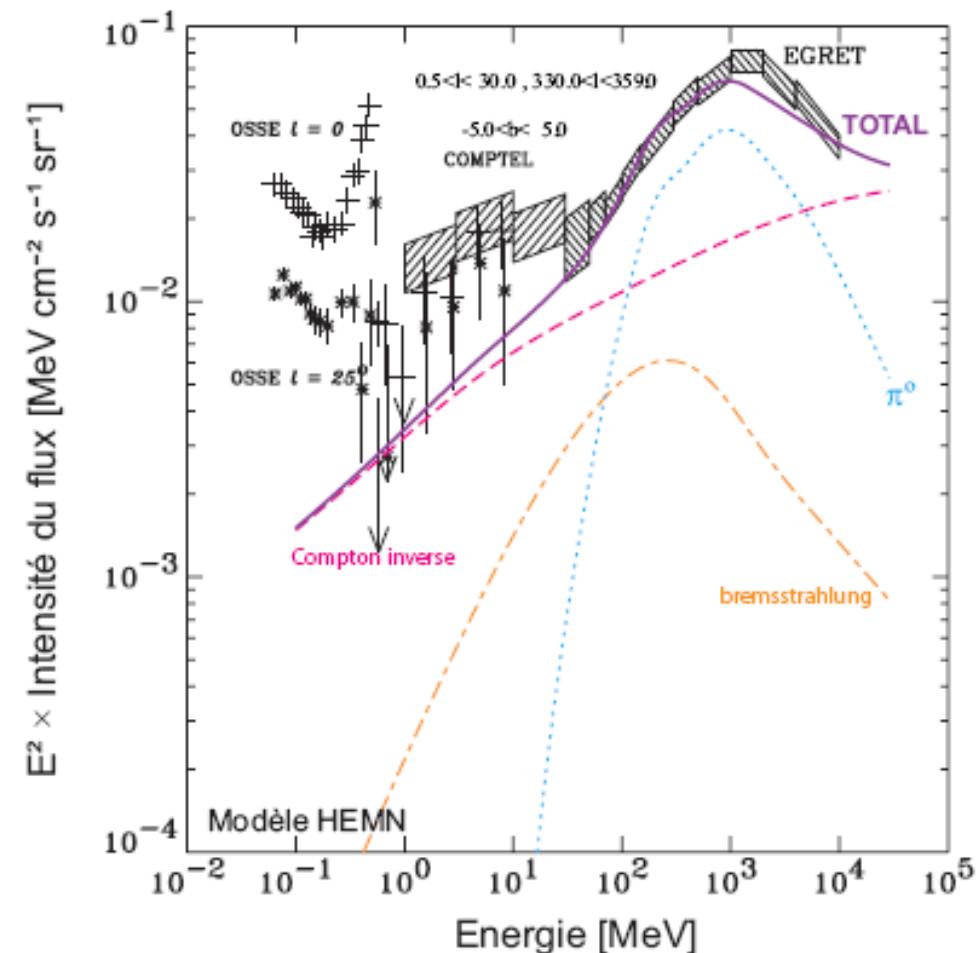
Tracking back the CR flux elsewhere in the Galaxy

- Other constrains
 - Protons local spectrum and flux
(altered <10GeV solar magnetic field)
 - Electrons local spectrum and flux
(influenced by the local bubble of matter under-density)
 - Radio measurements of the synchrotron emission



Tracking back the CR flux elsewhere in the Galaxy

- Refined predictions based on detailed simulations
(magnetic model of the galaxy, CR diffusion equation, matter density maps...):
for ex: GALPROP program

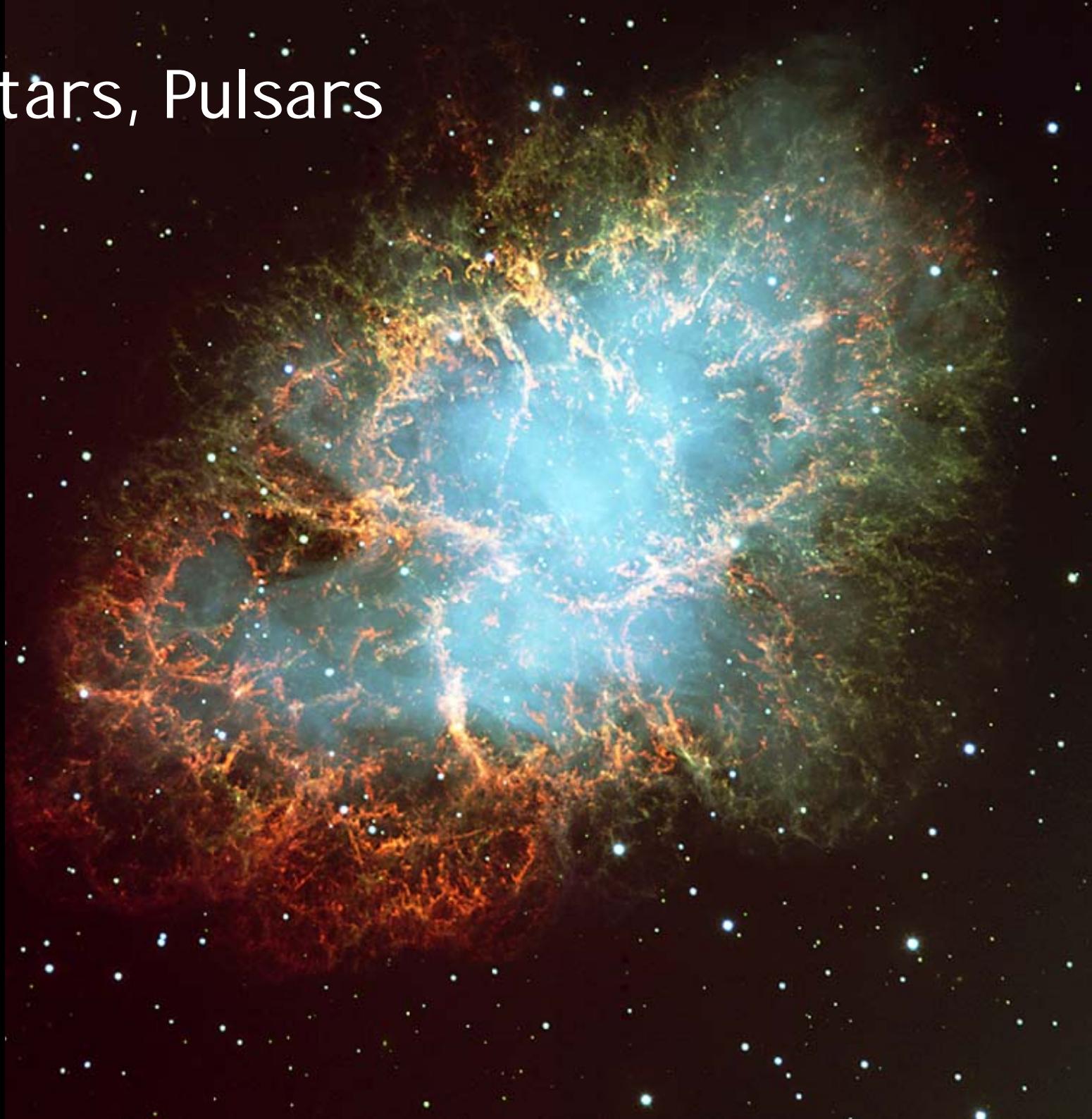


<http://www.gamma.mpe-garching.mpg.de/~aws/aws.html>

COMPACT OBJECT ENVIRONMENT : NEUTRON STARS AND PULSARS BLACK HOLES

Neutron stars, Pulsars

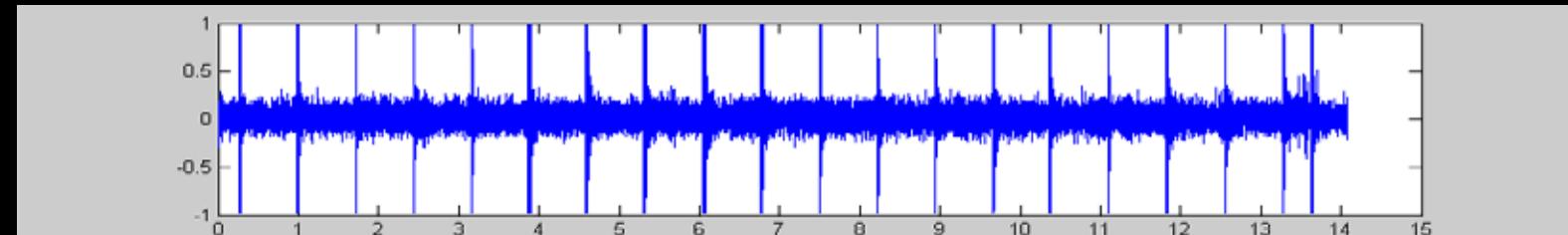
Journey into the
Crab nebula.



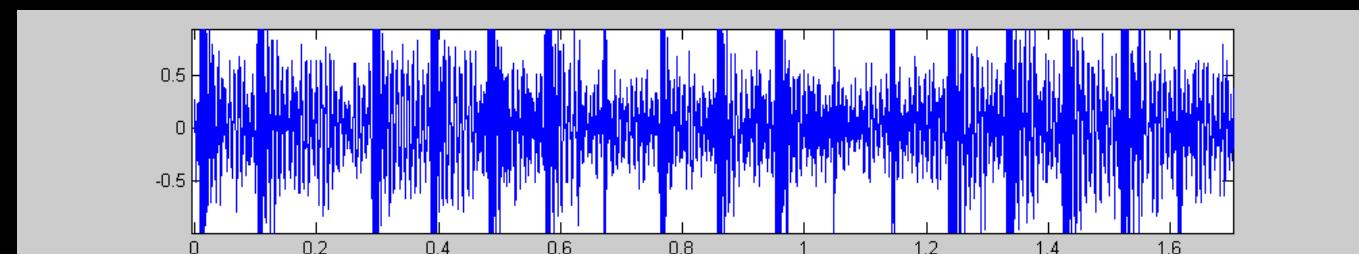
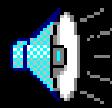
Here is radio pulsar...

2009-2010

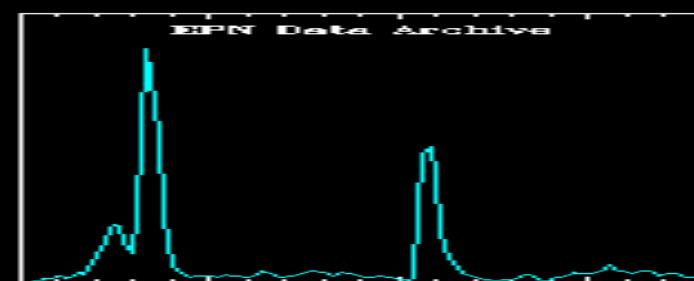
PSR0329+54
~1,40 Hz



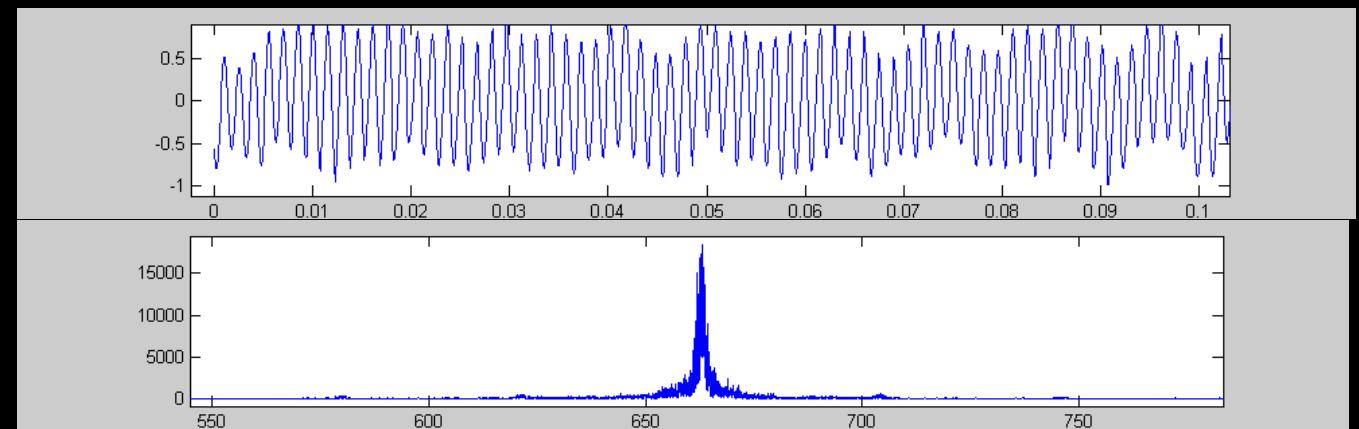
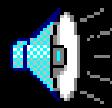
Vela
PSR0833-45
~11Hz



Crabe
PSR0531+21
~30 Hz



PSR1937+21
~641 Hz



Neutron stars

2009-2010

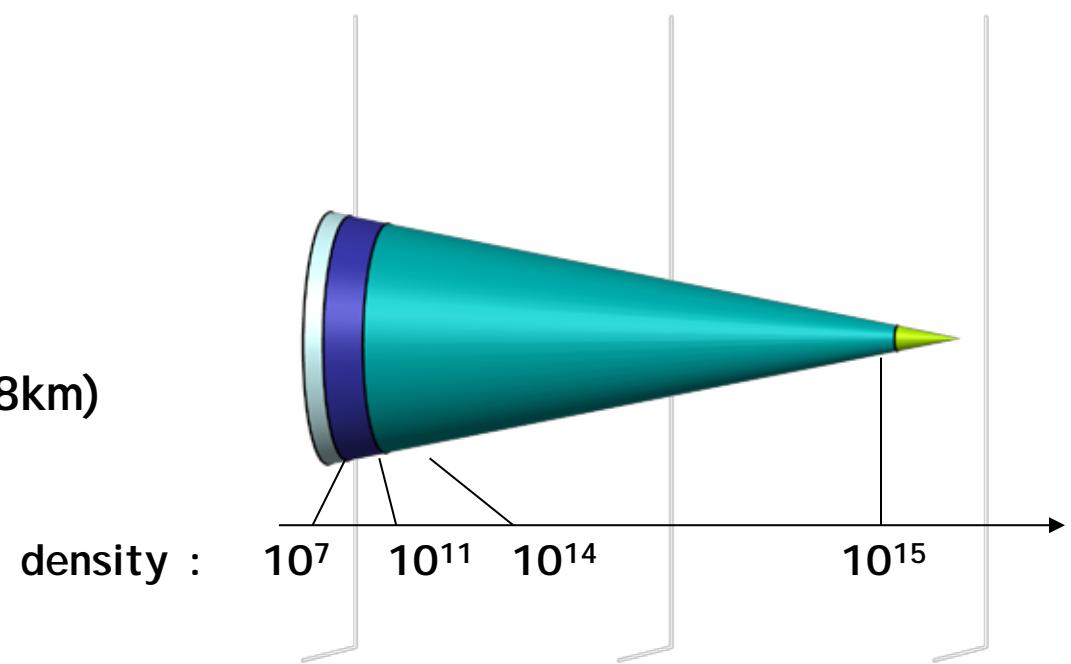
- Ultra compact objects:
 - Density ~ nuclear matter (10^{13} ' water !)
 - sphere ~ 10 km radius
 - gravitation at surface ~ $5 \cdot 10^{10}$ ' gravitation on earth
- Density and structure varies with :

■ Iron core (0,3km)

■ Neutrons and Nuclei (0,6km)

■ Neutrons forming an superfluid ocean (8km)

■ Unknown (1km)



Magnetic field of a neutron star

- Most celestial bodies bear a magnetic field (i.e. the sun: 10^{-3} Tesla)
- At the surface of a neutron star, it is ≈ 10 million Tesla
- This value is intrinsically linked to the rapid rotation
(conservation of rotation kinetic energy and magnetic energy \rightarrow concentration)
- A kind of giant **COSMIC DYNAMO** : **rotation of $\vec{B} \Rightarrow \vec{E} \Rightarrow 10^{18}V !!$**
- The electric force at the surface is \gg gravitationnal force !
 \rightarrow charged particles are **expelled** and **accelerated**
 $\Rightarrow 10^{38} e^-$ per second radiating synchrotron light.
- The strong anisotropy of the radiation is badly understood but probably due to the intense \vec{E} and \vec{B} fields near the "polar caps".

Unipolar Induction

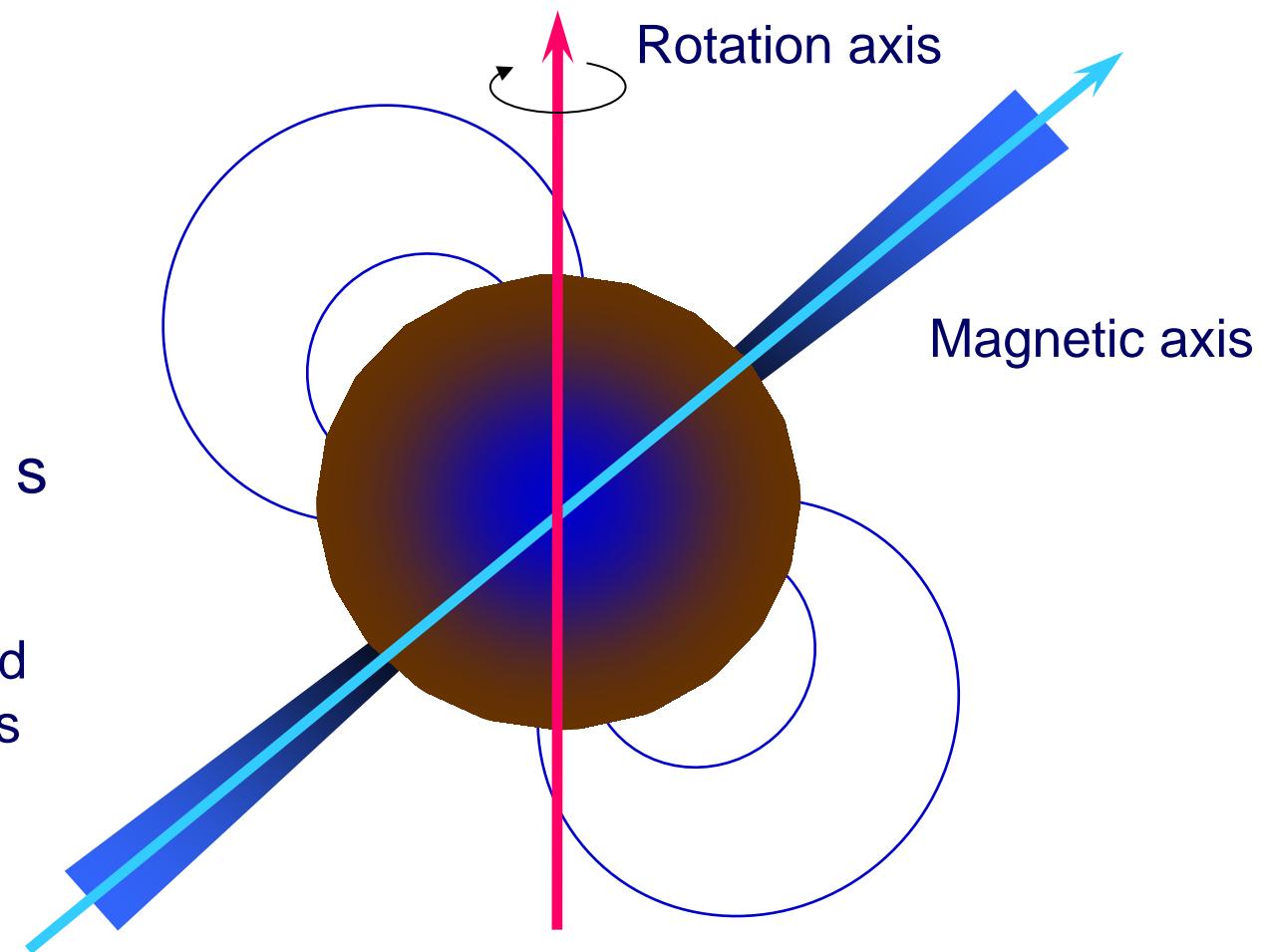
- **Neutron star :**
a rotating magnet with a magnetosphere

Neutron star

Mass = $1.4 M_{\odot}$

Radius = 10 km

Rotation period = 1 s



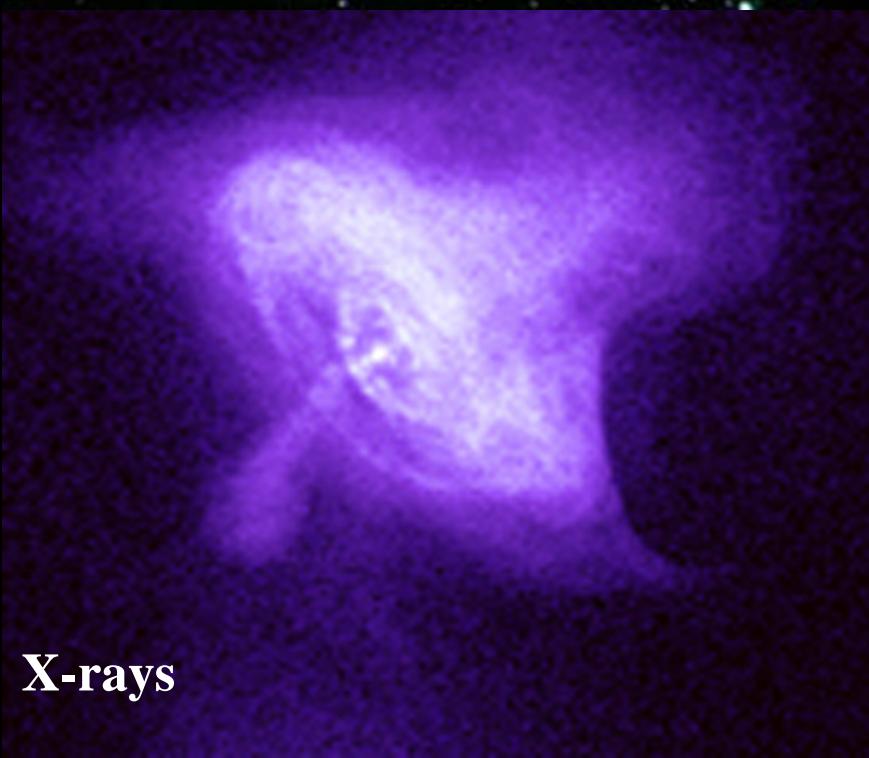
Induction unipolaire

- **Neutron star :**
a rotating magnet with a magnetosphere

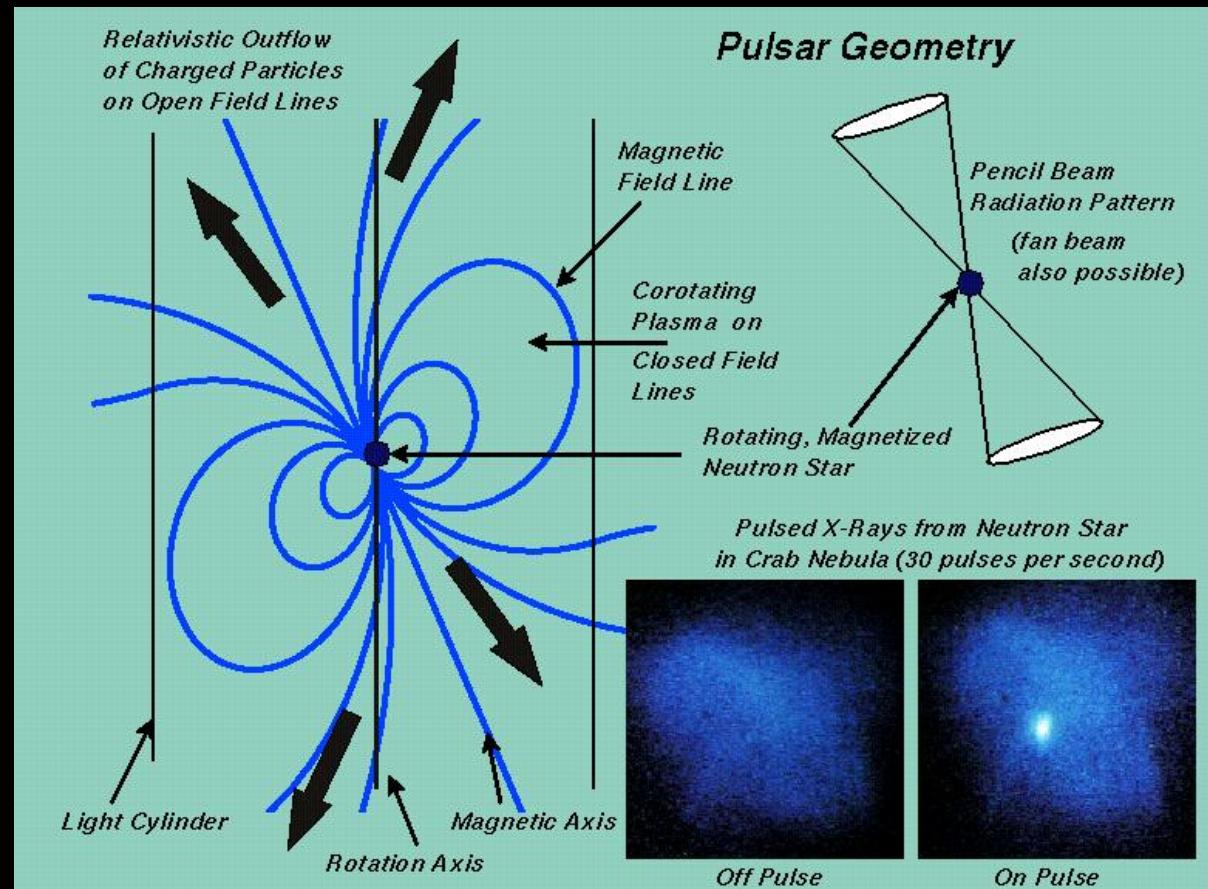
$$\rho_0 = \vec{\nabla} \cdot \left(\frac{(\vec{\Omega} \times \vec{r}) \times \vec{B}}{4\pi c} \right) = \frac{\vec{\Omega} \cdot \vec{B}}{2\pi c \left(1 - |\vec{\Omega} \times \vec{r}/c|^2 \right)}$$

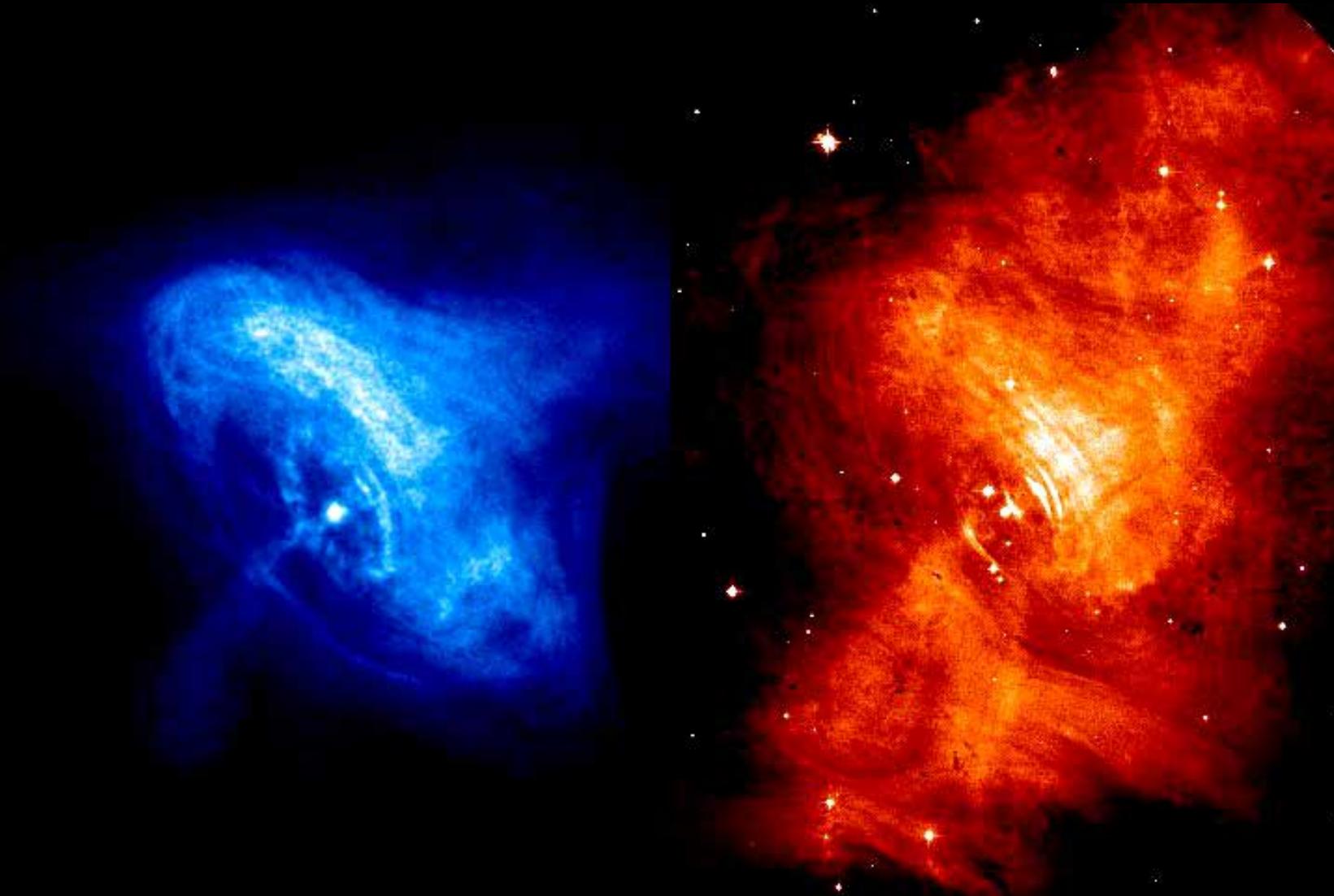
$$\Delta V \approx \frac{\Omega^2 B_s R^3}{c^2} = 3 \times 10^{16} \Omega_2^2 B_{13} R_6^3 \text{ Volts}$$

Optical



X-rays

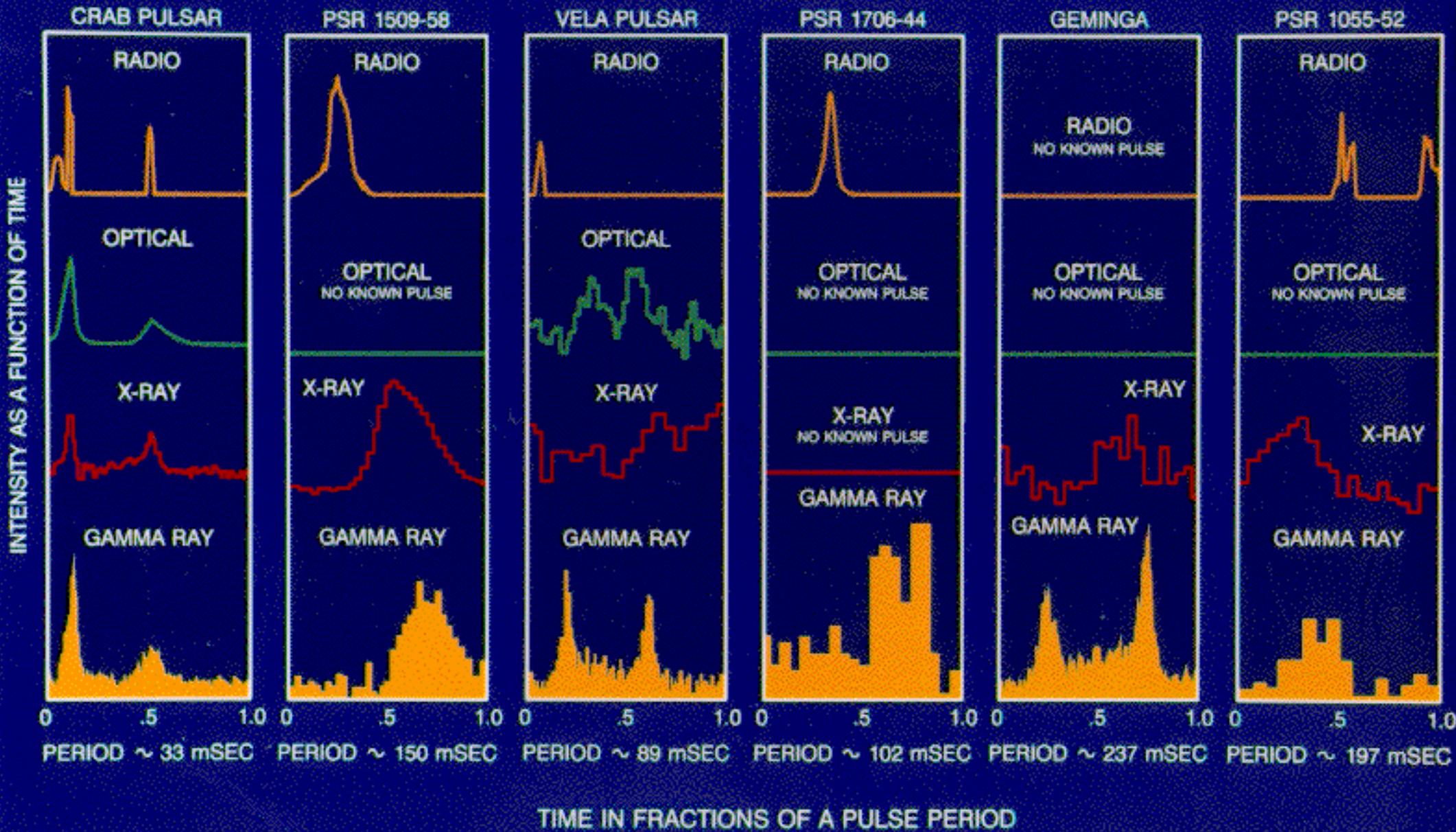




The Crab X-source, the central engine, X (Chandra) et optical (HST)

7 images from november 2000 to april 2001, showing the time and space coincidence between the shocks observed at the very center of the system (close to the pulsar) and the X-ray radiation produced by the accelerated VHE electrons

GAMMA-RAY PULSARS

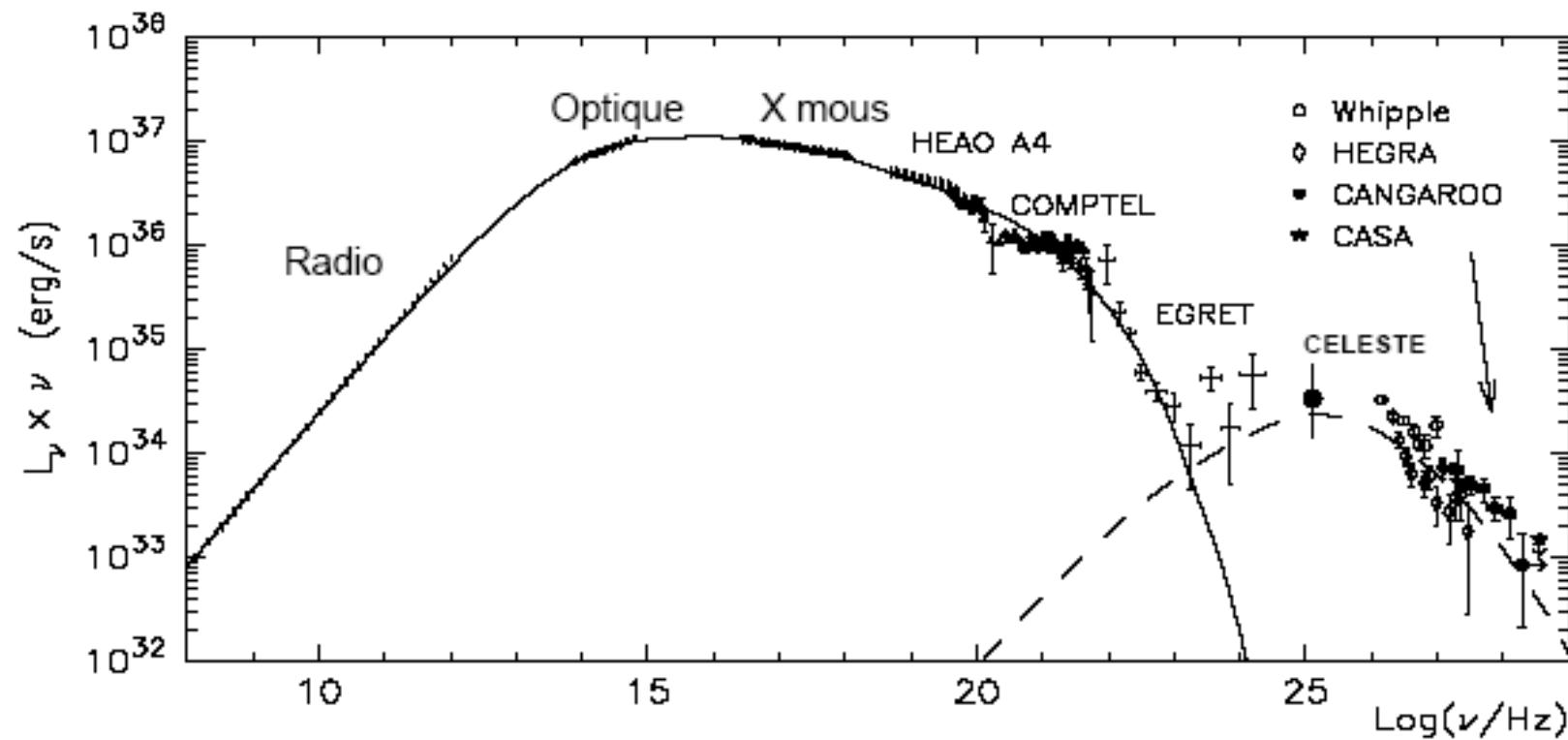


D665.001

The Crab Nebula a case study for VHE gamma-rays VHE

2009-2010

- Steady emission observed at all wavelengths.
- First point like gsource identified
- Intense flux: a "standard candle" for γ -ray observatories

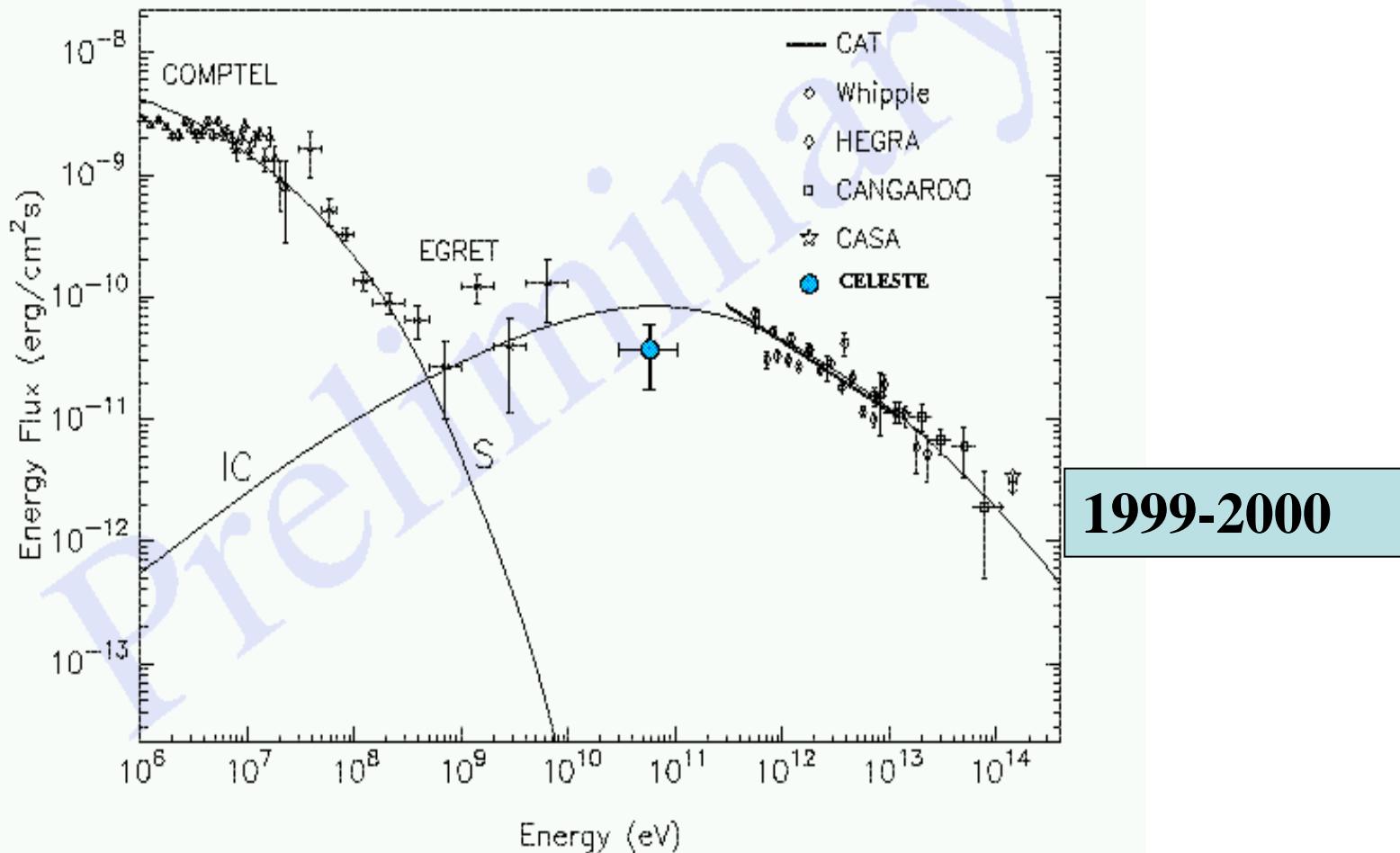


The Crab Nebula

a case study for VHE gamma-rays VHE

- Two distinct components

Preliminary evaluation from CELESTE data
of the Crab Nebula around 60 GeV



VHE photons sources

- γ -ray production:
 - Electro-magnetic processes :

- Bremsstrahlung

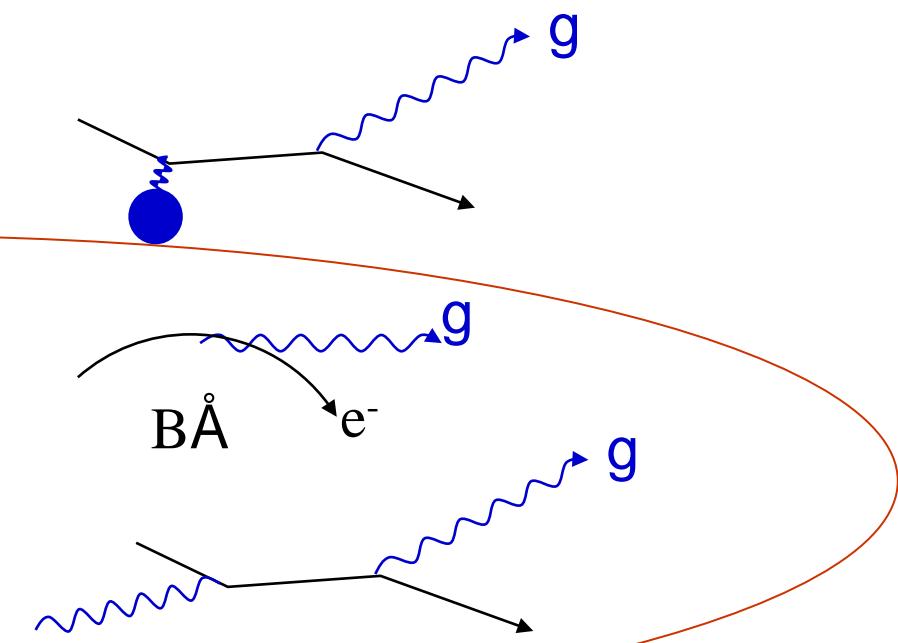
- Synchrotron radiation

- Inverse Compton scattering

- Hadronic processes :

$$(p, N) + (N\gamma) \rightarrow \pi_0 + \dots$$

$\hookrightarrow \gamma\gamma$



The Crab Nebula

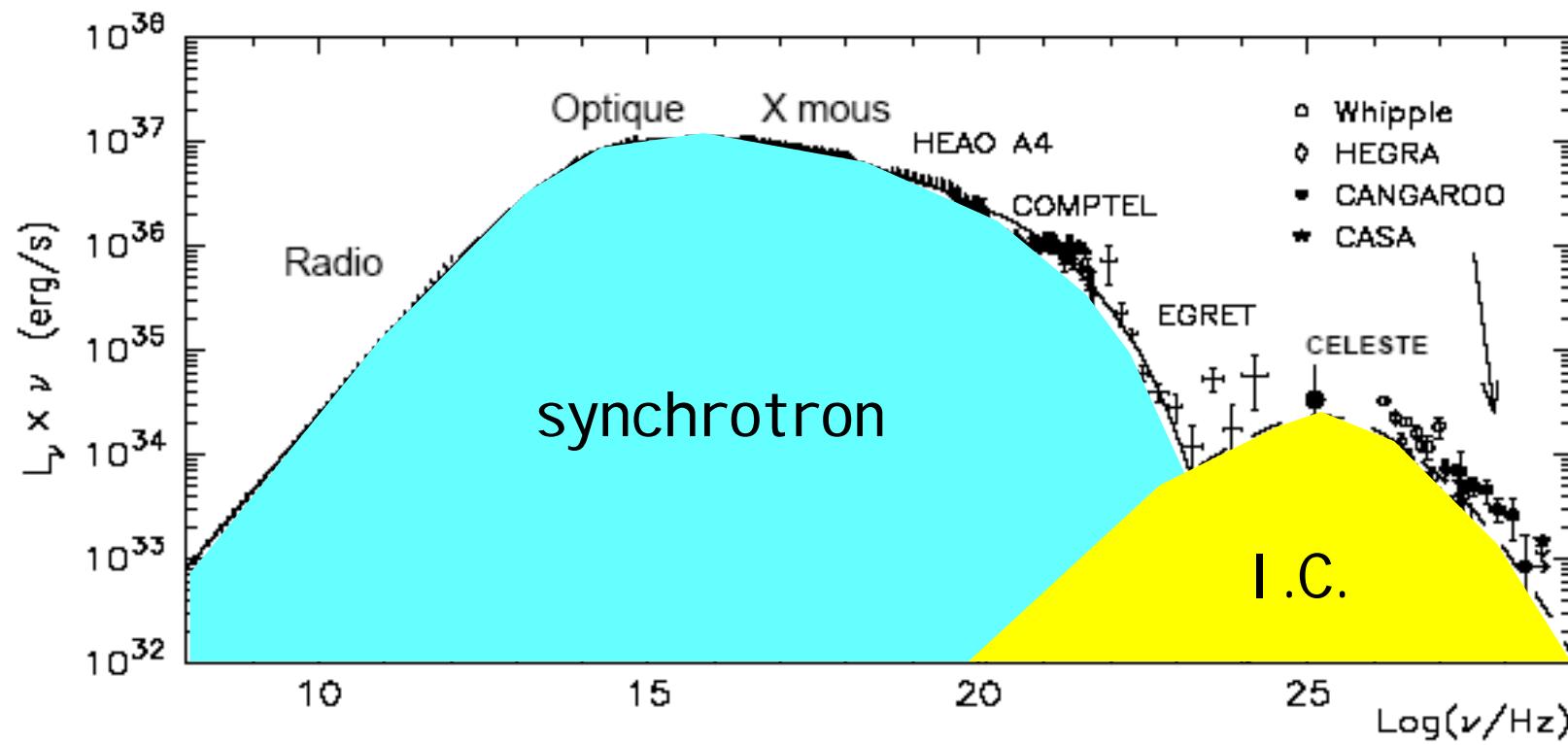
a case study for VHE gamma-rays VHE

- TeV emission is based on Inverse Compton scattering
- The e^- producing synchrotron radiation can interact with these synchrotron photons.
- If $\lambda_{\text{photon cible}} \ll \lambda_{\text{compton}} = \frac{h}{m_e c}$ it is possible to transfer most of the incident e^- energy to the photon.
ex: A $10^{13} eV$ e^- can boost an I.R. photon I.R. into a VHE gamma-ray.
- These "Self Synchrotron Compton" models reproduce both GeV and TeV spectra.

La Nébuleuse du Crabe et les rayons gamma VHE (~TeV)

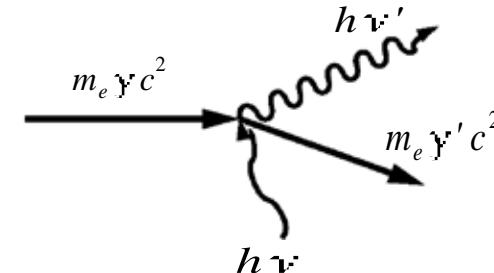
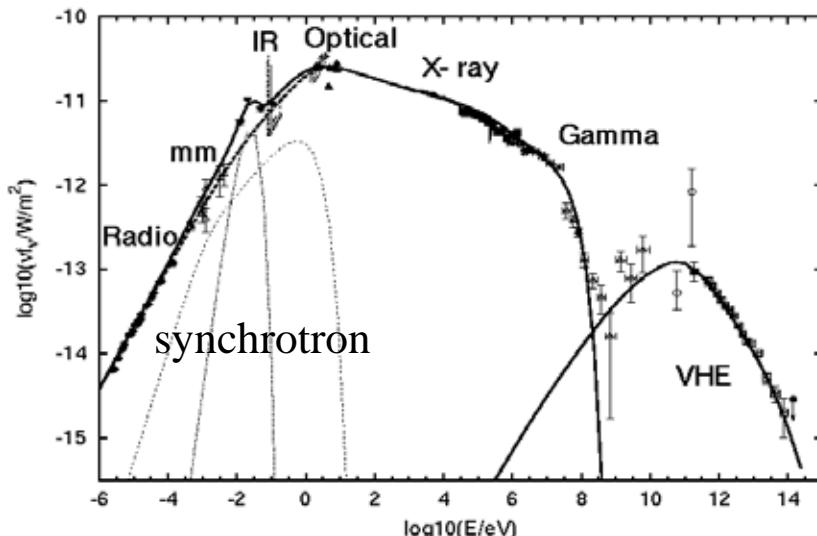
2009-2010

Model SSC : Synchrotron Self Compton



PWN emission mechanisms: the Crab Nebula

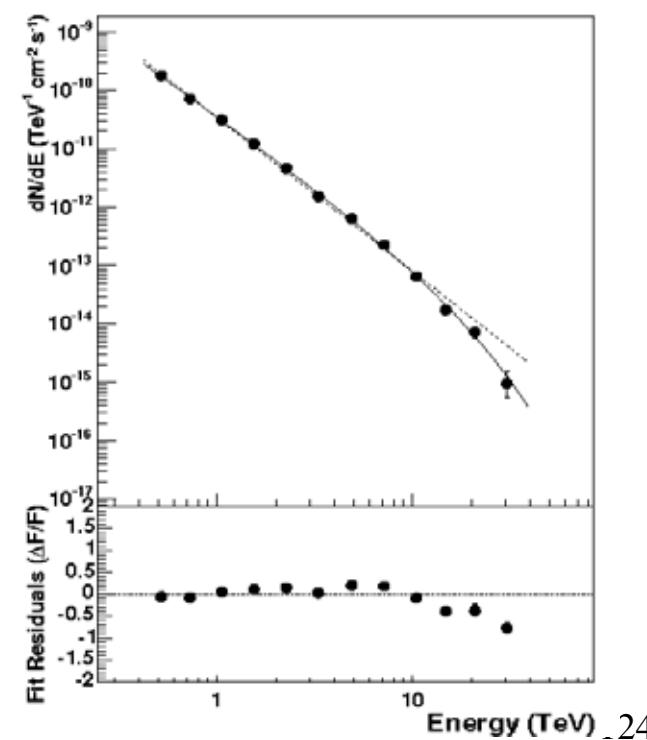
- Assume leptonic model: synchrotron and Inverse Compton emission
- Relativistic electrons and positrons created and accelerated by the pulsar



Target photons : CMB, interstellar IR, stellar photons, synchrotron (SSC)...

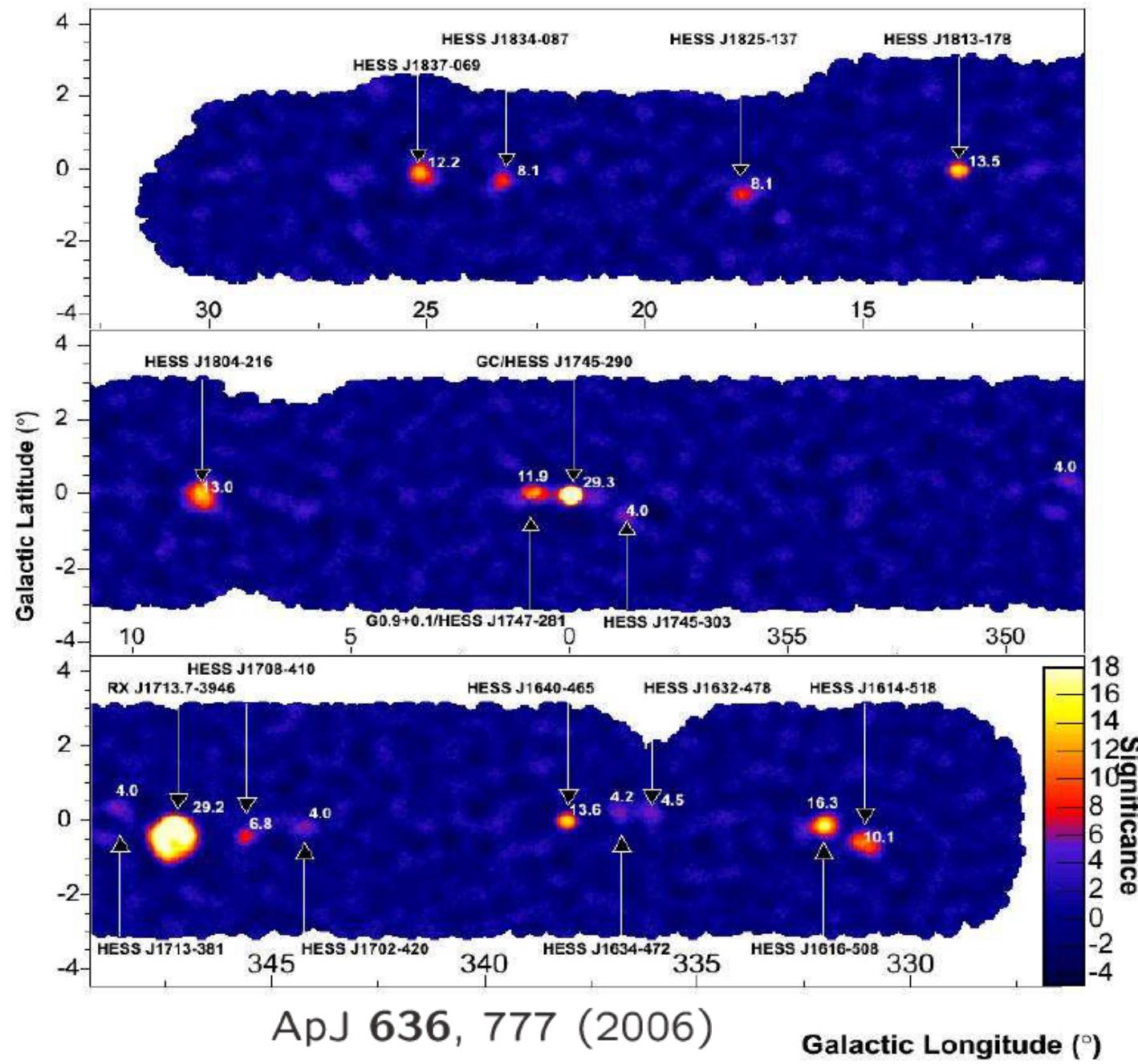


H.E.S.S. spectrum
(*A&A* 2006 in press,
astro-ph/0607333):
Spectral curvature,
Consistent with IC
expectations



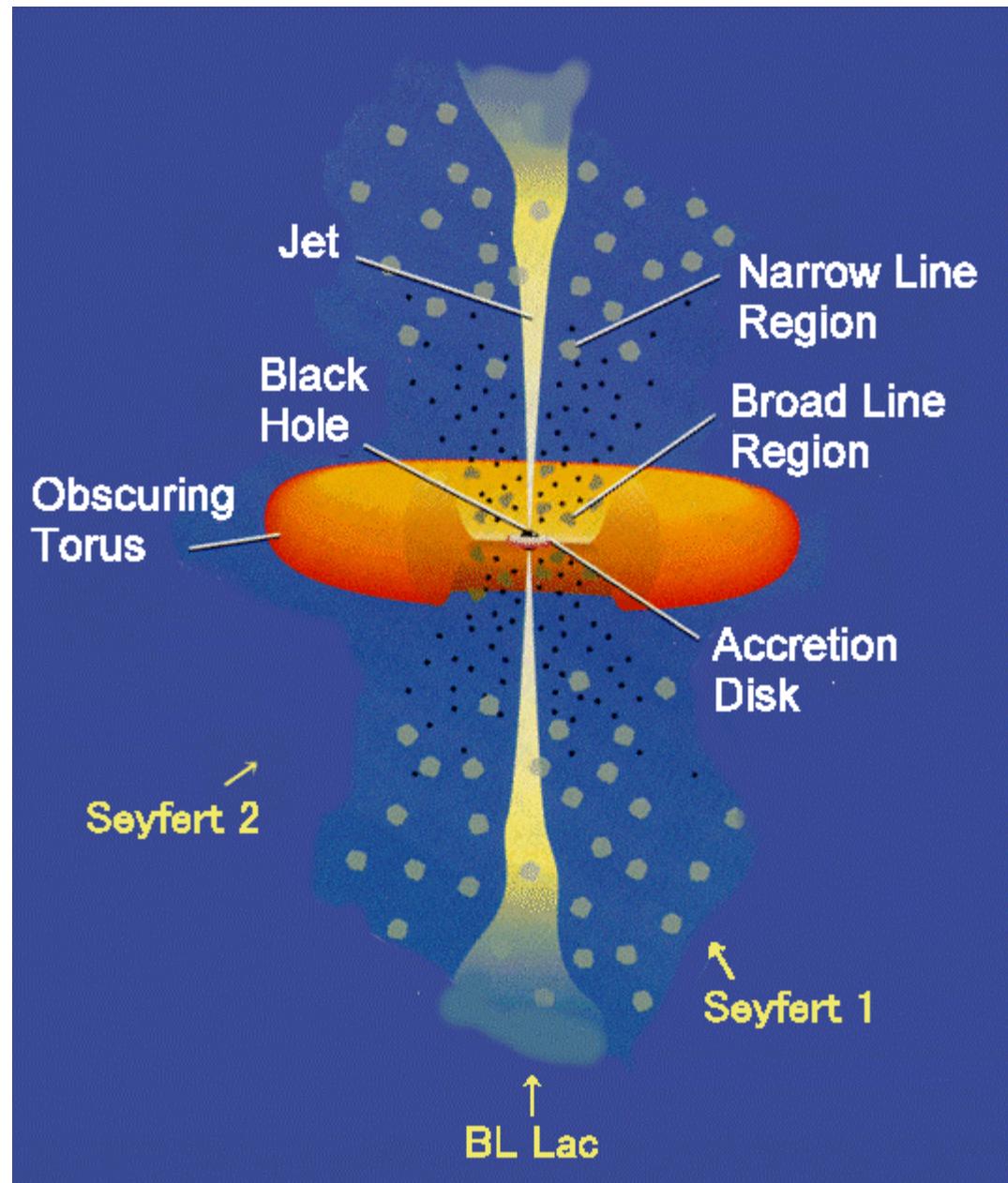
Radio, optical, X-rays

Many recently discovered sources in the Galactic plane by HESS, MAGIC and VERITAS large angle surveys



THE ENVIRONMENT OF SUPERMASSIVE BLACK HOLES: ACTIVE GALATIC NUCLEI

AGN: a unified scheme



Active Galactic Nuclei

- Galaxies with an important activity of its nucleus
- Numerous classes of objects:
 - Radio-galaxies
 - Radio-quasars
 - BL Lac objects
 - Violently Variables objects (VVO)
 - Radio Quiet Quasars
 - Seyfert Galaxies of type 1
 - Seyfert Galaxies of type 2
 - Low Ionization Nuclear Emission-Line Regions (LINERs)
 - Nuclear HII Regions
 - "Star Burst" Galaxies
- Arbitrary categories et badly mal defined.

What is an AGN

- Elementary AGN model:

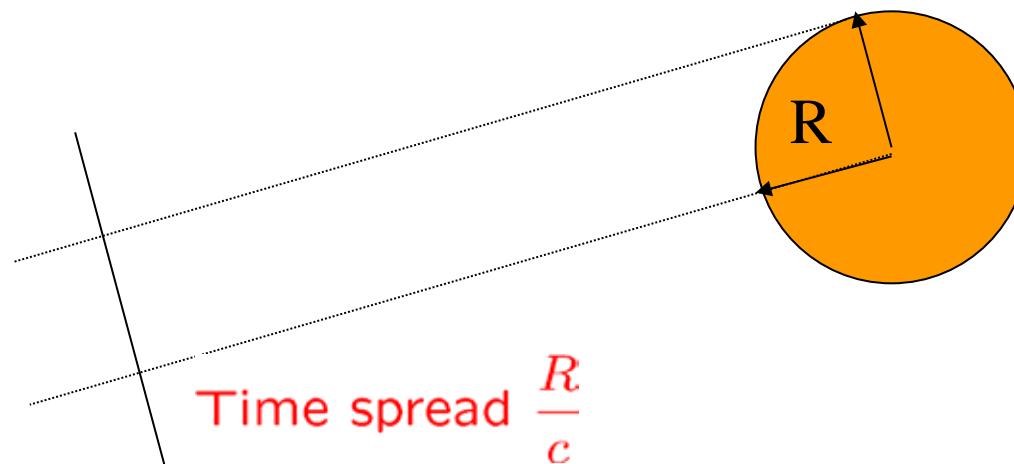
Discovered in the early 60ties when trying to identify the sources in the 3rd Cambridge Catalogue.

- Superluminous objects

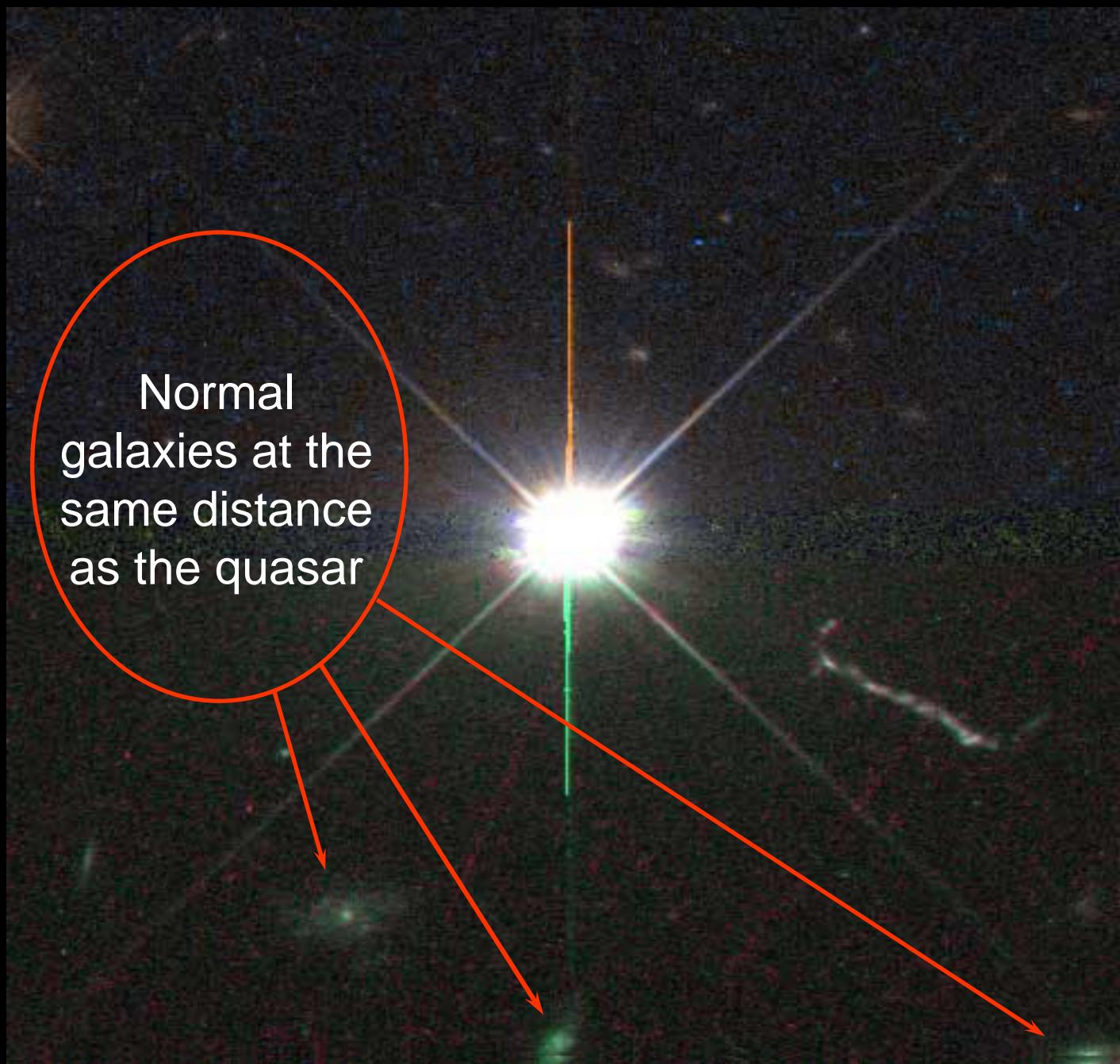
- Spectral redshift \propto cosmological distances
 - For 3C273 $\propto P^{1/4} 10^{40}$ Watts $\propto 10000$ the Milky Way $\propto 1000$ SN/year

- Very small objects

- Variability criteria $\propto d^{1/4}$ solar system size !!



The Quasar 3C 273



In 1963, 3C 273, the first quasar, was discovered. They can be up to **1000 more luminous** than the massive host galaxy. The emission originates in the **active nucleus** of the galaxy.

What is an AGN

- Elementary AGN model:

Discovered in the early 60ties when trying to identify the sources in the 3rd Cambridge Catalogue.

- Objects from the past

- The most distance quasars observed ($z \approx 5$) are at distances of the order of the radius of the **observable universe!**
 - No such objects in our close environment \circledR must be linked to a phase of the galaxy evolution.

What is an AGN

- Necessarily a gravitational engine

Thermonuclear falls short by a factor 10 to 100 !

Form of energy release	Efficiency of energy production (wrt mc^2)
• Chemical energy	10^{-9}
• Nuclear energy	10^{-2}
• Accretion of mass onto non-rotating black holes	6×10^{-2}
• Accretion onto maximally rotating Kerr black holes	0.42
• Rotational energy of maximally rotating Kerr black holes	0.29

What is an AGN

- Necessarily a gravitational engine

$$E_{tot} = 0 \Rightarrow E_{cin} = \frac{GmM}{R} \Rightarrow R \text{ must be small}$$

Huge mass, small volume \Rightarrow BLACK HOLE !!!

$$\text{With } R_S = \frac{2GM}{c^2} \Rightarrow E_{rad.max} \approx \frac{GM}{R_S} \frac{R_S}{R} m = \frac{1}{2} \frac{R_S}{R} mc^2$$

Accounting for relativistic corrections and ergosphere rotation
 \rightarrow conversion efficiency $m \rightarrow E$ few 10%

What is an AGN

- Necessarily a gravitational engine

One gets 10^{40} Watts with $15M_{\odot}/\text{year}$.

Over 10^8 years (age measured from radio lobes) $\Rightarrow 10^9 M_{\odot}$

$\Rightarrow R_S \approx 10^{13} m$ compatible with variability.

More over, one is below Eddington luminosity

(radiation pressure = gravitational forces)

$$\left(L \approx \frac{4\pi c}{\sigma_T} GMm_H \right) \text{ de } \left(\frac{\sigma_T L}{4\pi R^2} = \frac{GMm_H}{R^2} \right)$$

\Rightarrow PLAUSIBLE.

What is an AGN

- Accretion of matter:
 - In absence of angular momentum:
 \vec{v}_∞ directed toward the BH: $\frac{GMm_H}{R} = kT$
⇒ almost all is radiated in X or γ rays
INCOMPATIBLE with MEASUREMENTS
 - With angular momentum:
Intense radiation in the UV : **OBSERVED**
⇒ **Anisotropic Accretion**
in the shape of an ACCRETION DISK

What is an AGN

- **Jets:**

Relativistic particles constituting the jets are accelerated when reflecting on the boundaries of a magnetic "tunnel" induced by the rotating plasmas near the BH.

- **Hot spots**

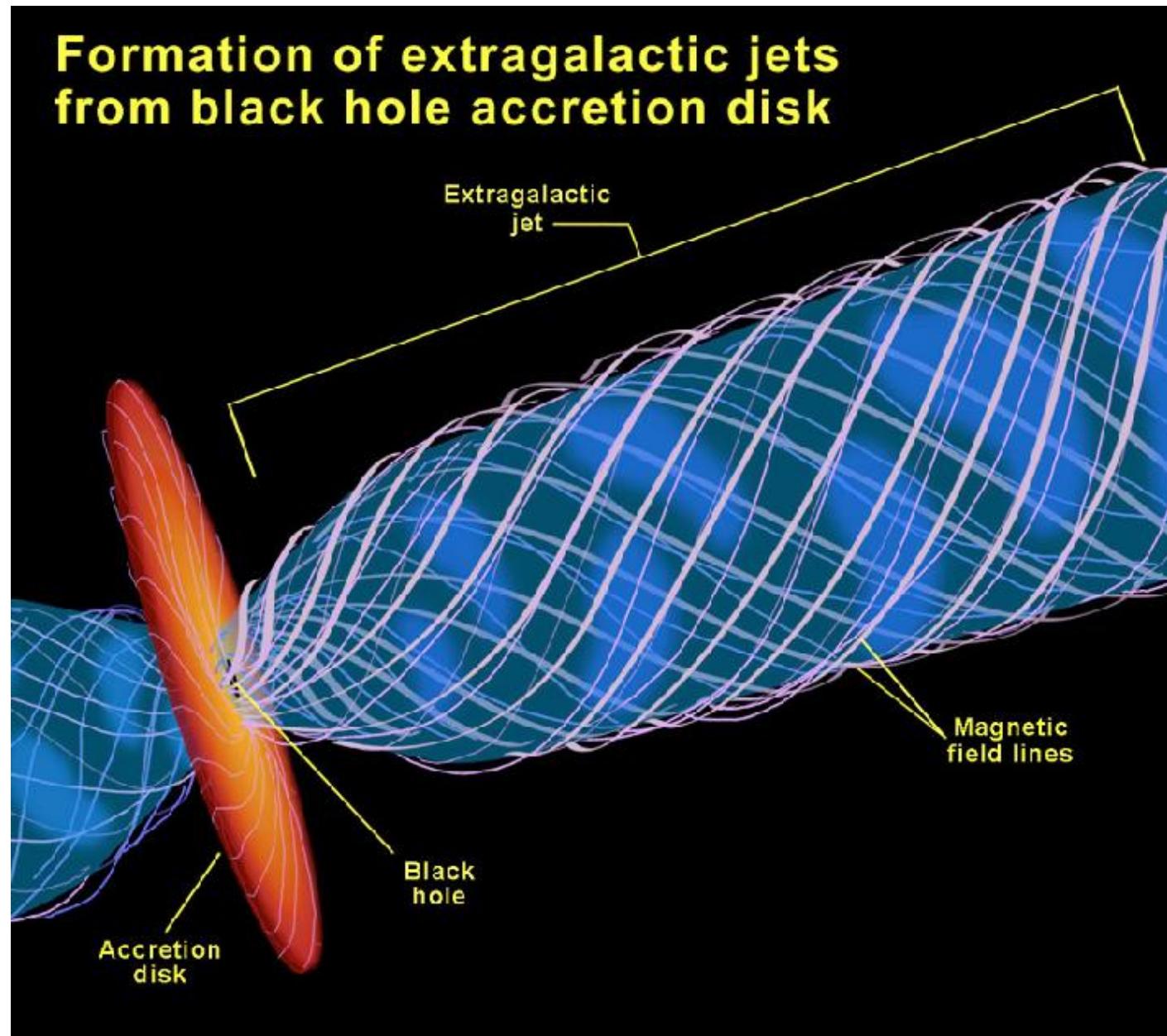
Terminal shocks (jet-IGM): large size moderate field regions
⇒ plausible place for UHECR production.

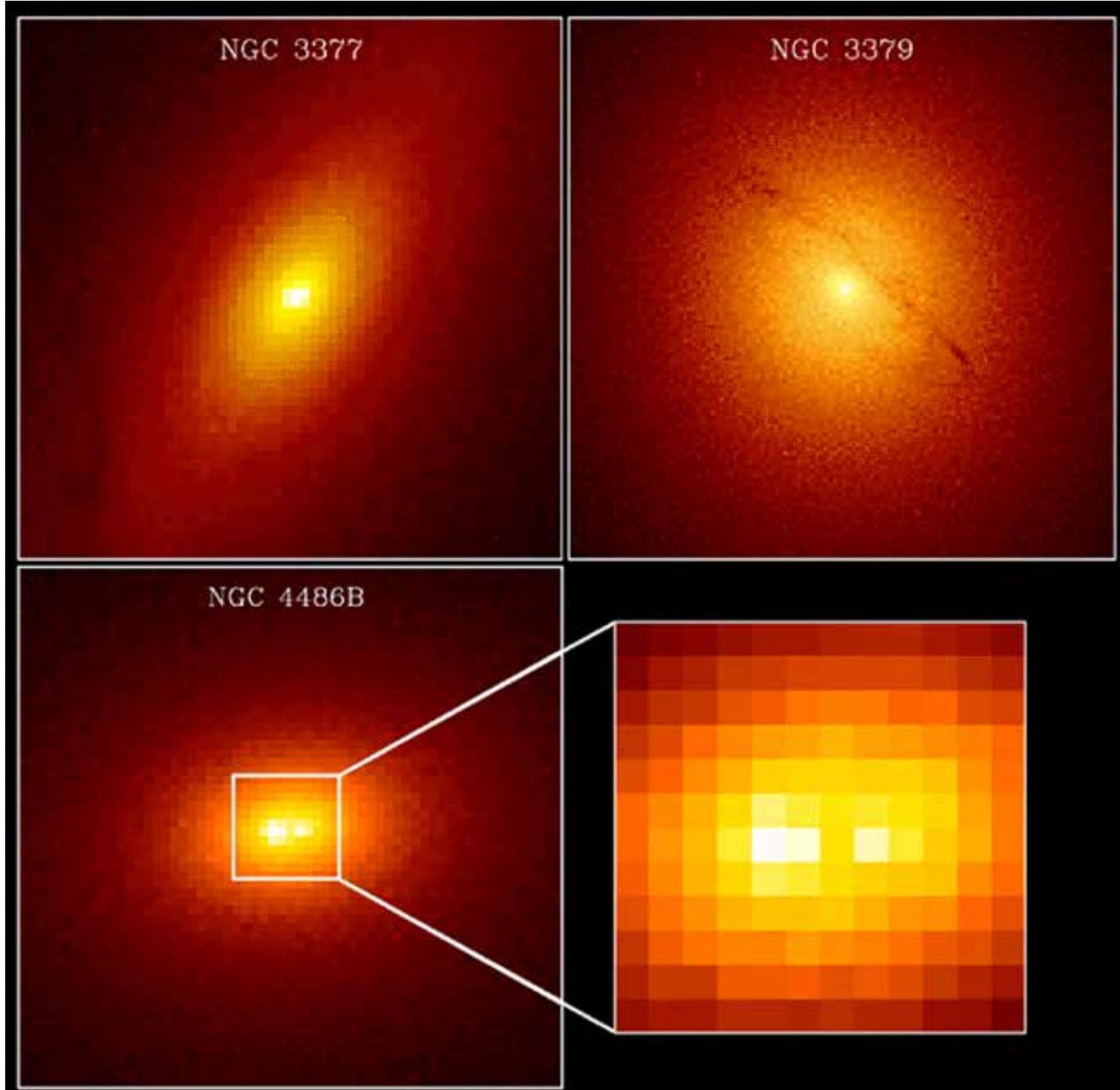
→ explains many observed phenomena with both leptonic and hadronic models.

Qu'est-ce qu'un AGN

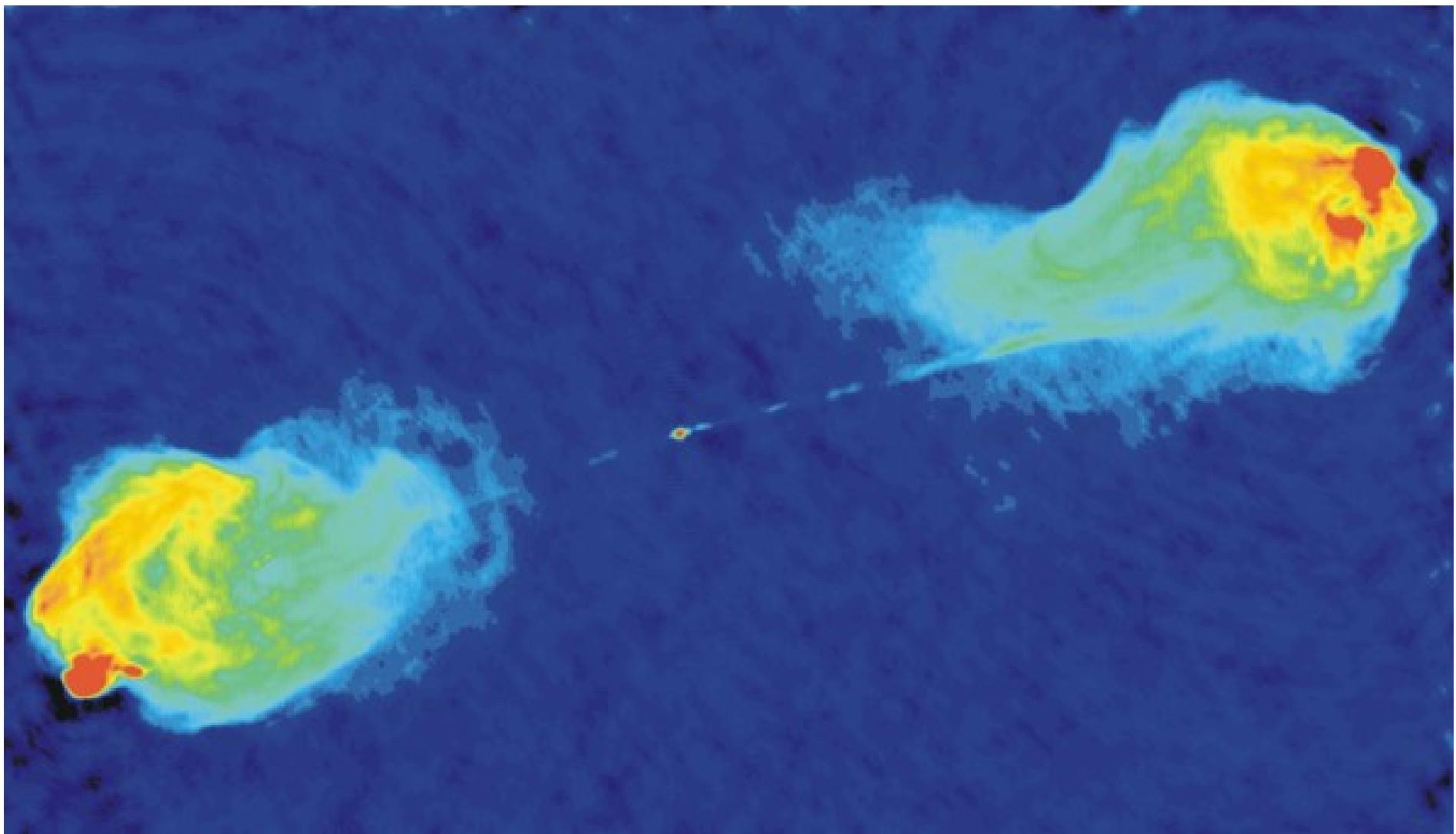
2009-2010

- Jets :

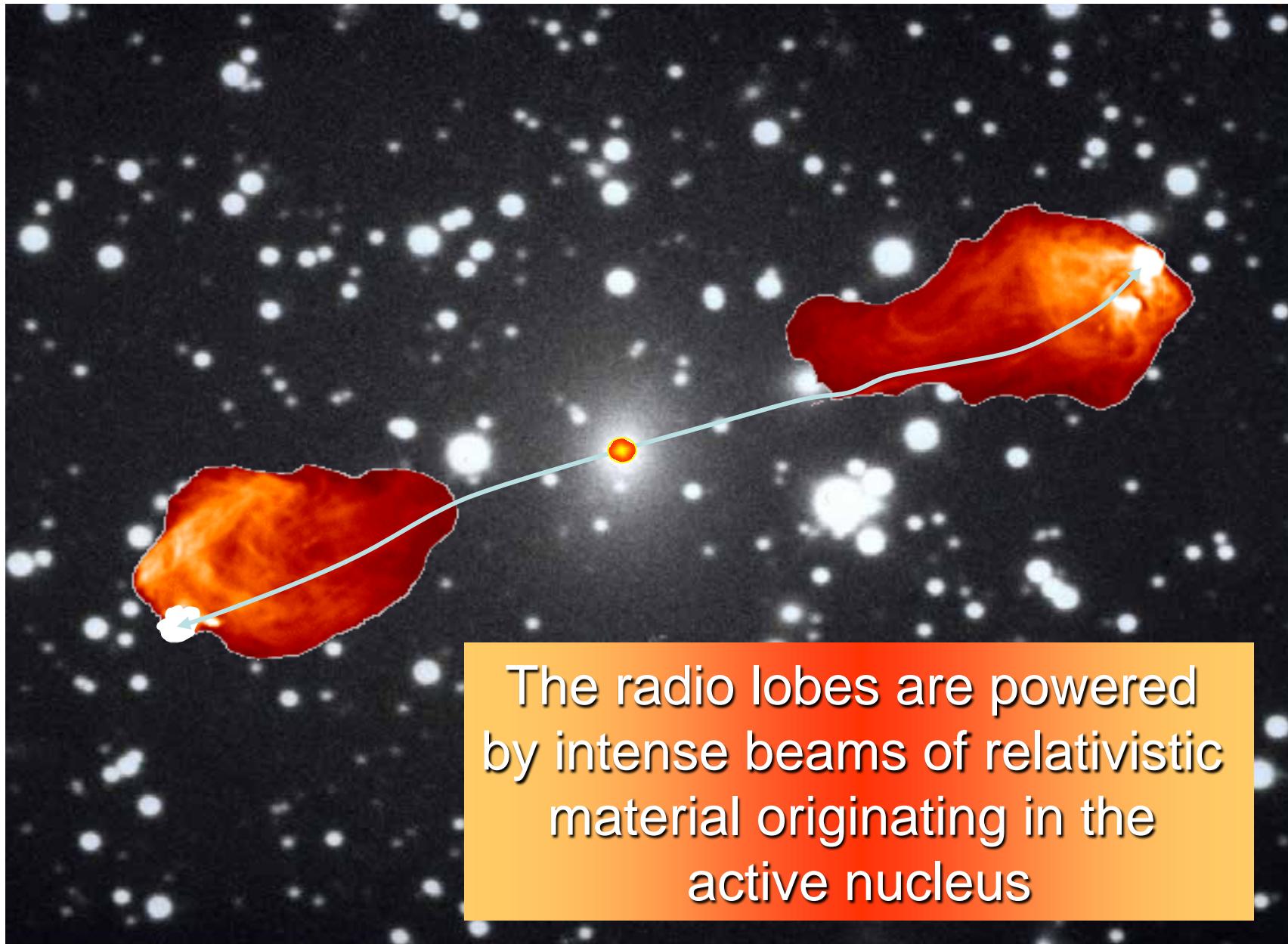




Cygnus A

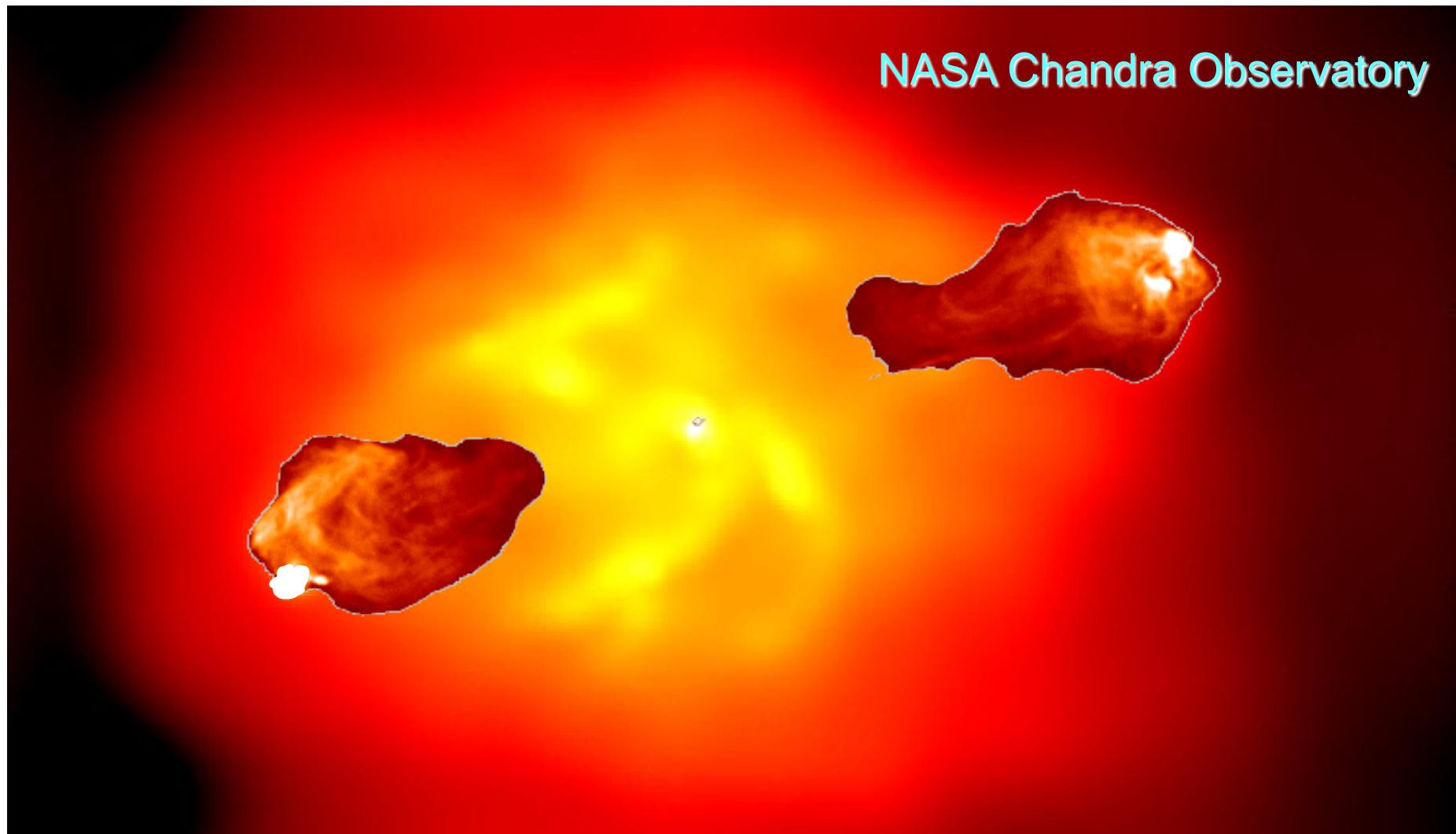


The Radio Galaxy Cygnus A



Cygnus A – X-ray image

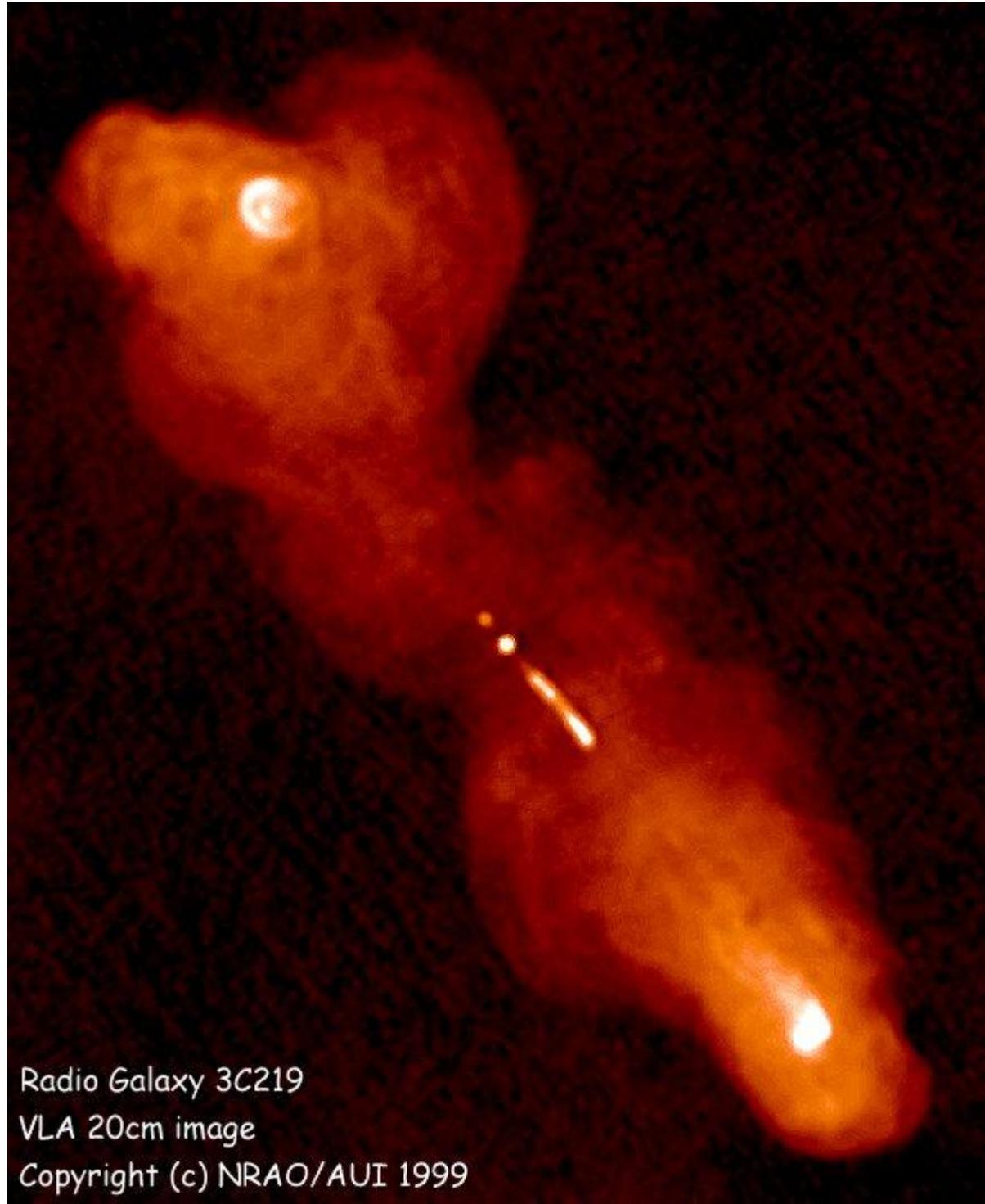
The image shows the distribution of hot intergalactic gas surrounding the radio source.



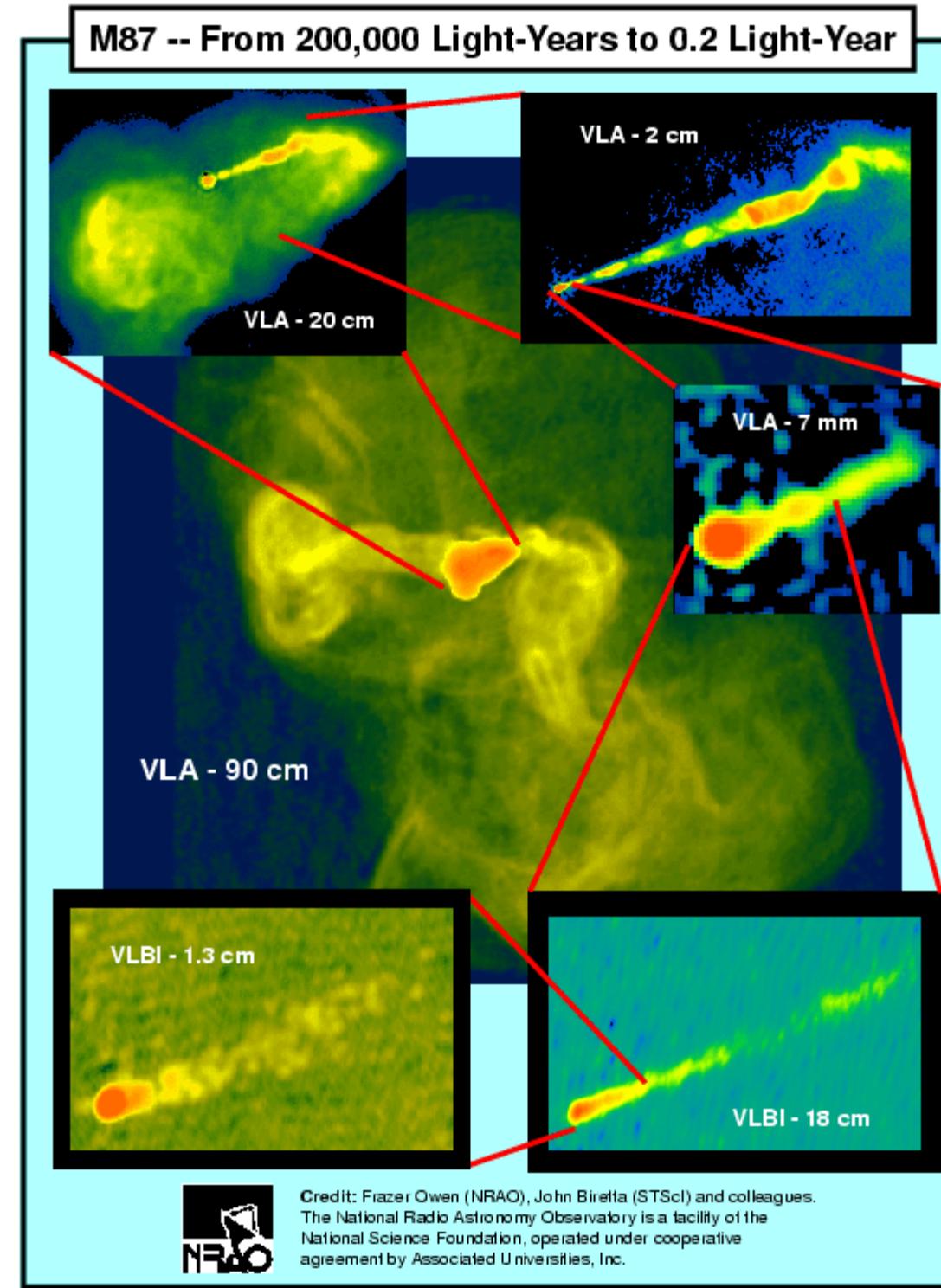
3C219



**BLUE = V BAND
RED = 22 CM**



Radio Galaxy 3C219
VLA 20cm image
Copyright (c) NRAO/AUI 1999

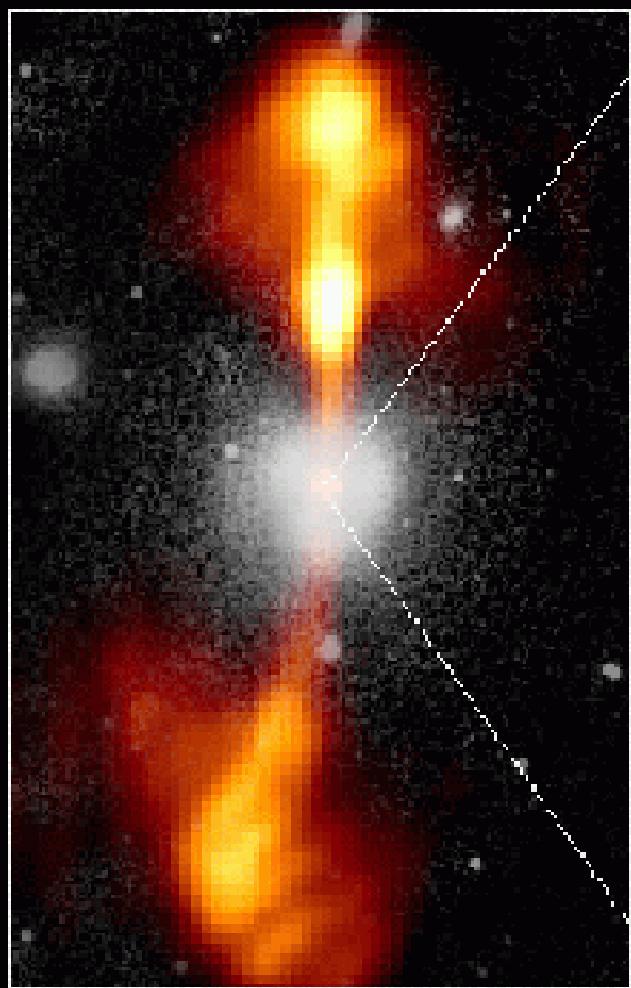


Core of Galaxy NGC 4261

Hubble Space Telescope

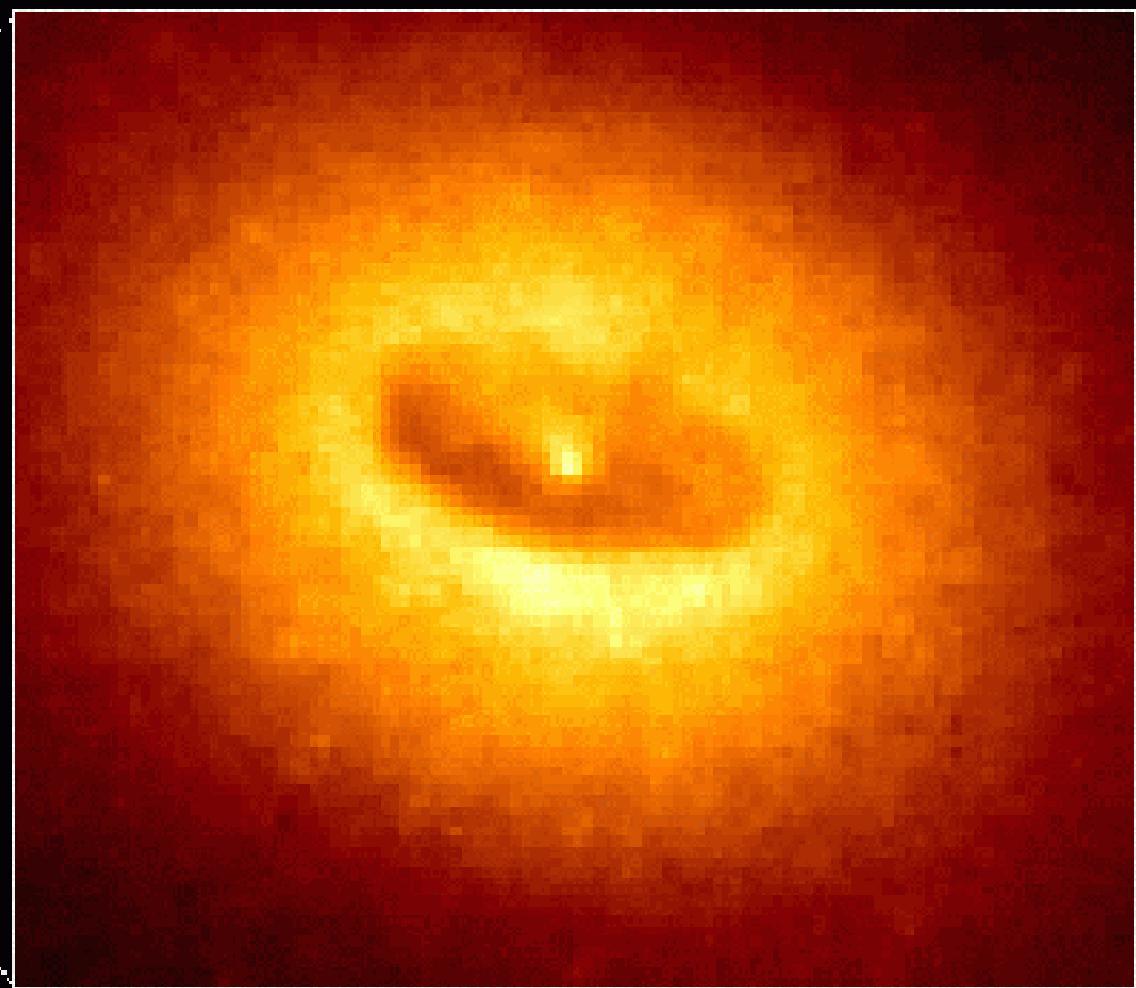
Wide Field / Planetary Camera

Ground-Based Optical/Radio Image



390 Arc Seconds
88,000 LIGHTYEARS

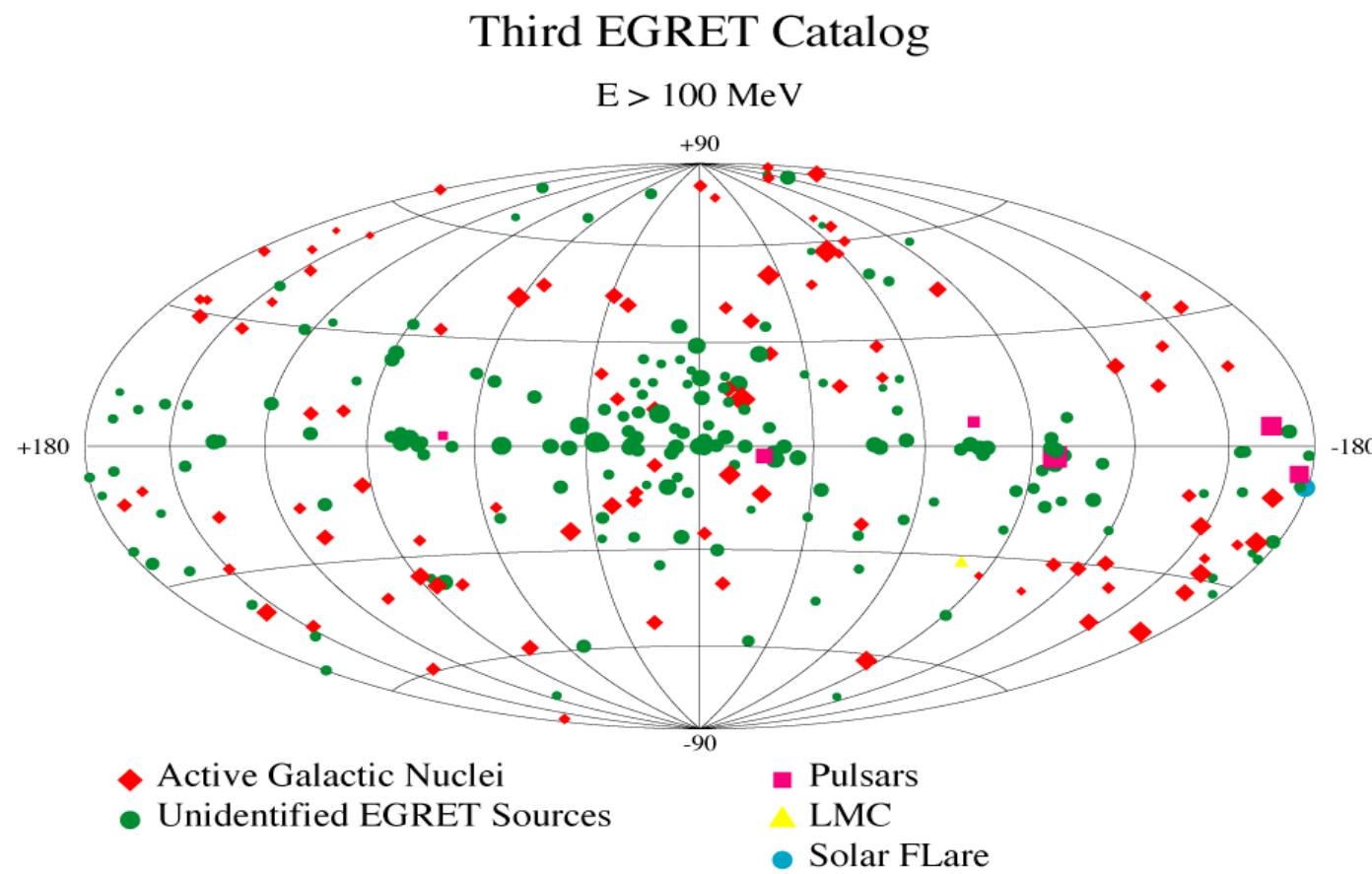
HST Image of a Gas and Dust Disk



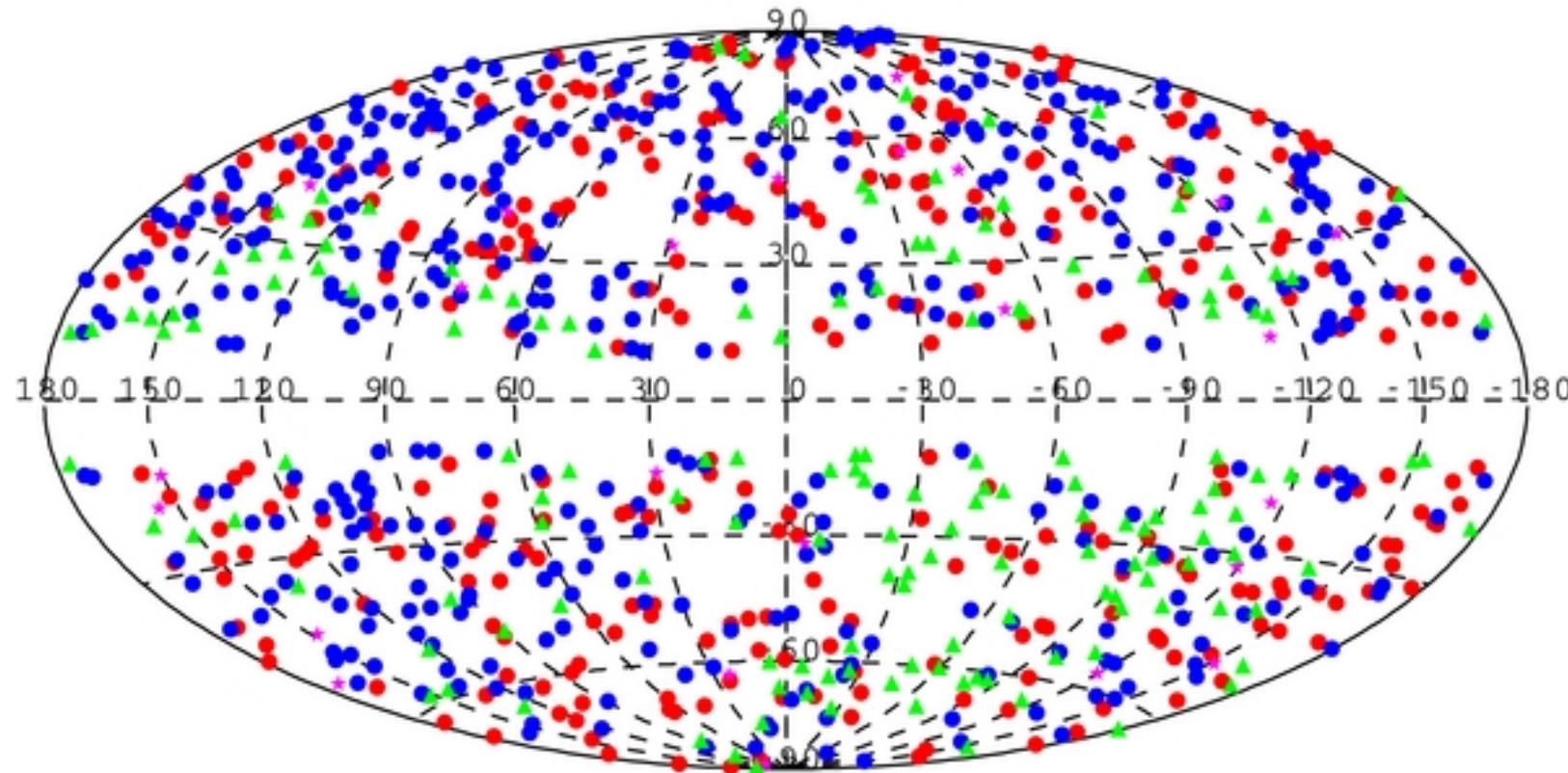
17 Arc Seconds
400 LIGHTYEARS

EGRET

- Before EGRET : 3C273
- After EGRET : 33 AGN a more than $5 \frac{3}{4}$
all Blazars



FERMI 2nd LAC catalogue



886 AGN with a clear GeV gamma-ray signal, identified with comprising 395 BL Lacertae objects (BL Lac objects), 310 flat-spectrum radio quasars (FSRQs), 157 candidate blazars of unknown type (i.e., with broadband blazar characteristics but with no optical spectral measurement yet), 8 misaligned AGNs, 4 narrow-line Seyfert 1 (NLS1s), 10 AGNs of other types, and 2 starburst galaxies

EGRET

- AGN basic results :
 - The γ -ray energy flux is often dominant
 - A naive estimate $(4\pi d^2 L) \approx 10^{47} \rightarrow 10^{49} \text{ergs.s}^{-1}$
 - $0,03 \leq z \leq 2,28$
 - $(30MeV - 30GeV)$ spectrum \sim a power law:
$$\frac{dN}{dE} \propto e^{-\alpha}$$
 with $1,3 \leq \alpha \leq 3,0$
 - Many sources are strongly and rapidly variables on time scales ~ 1 month.
 - Many known AGN are not detected in γ -rays.

EGRET

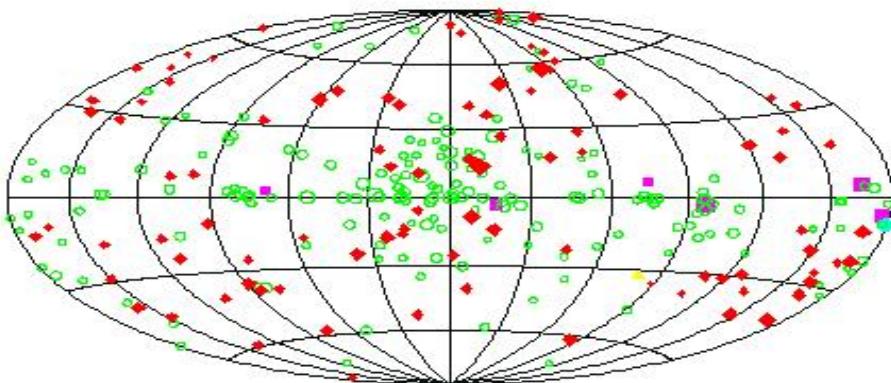
- Variabilité:
 - La plupart des Blazars EGRET présentent une variabilité sur des échelles de temps comprises entre **le jours et un mois**.

Source	T_{var} (jours)	$T_{var AGN}$ (jours)
0208-512	8	4
0235+164	<90	<46
0446+112	15	...
0528+134	2	0,65
1253-55	2	1,3
1406-076	6	2,4
1633+382	2	0,7
2251+158	4	2,2

- Les temps très courts **constraining la géométrie de la source :**
 1 jour \triangleright Rayon de Schwarzschild d'un trou noir de $10^{10} M_{\odot}$

Third EGRET Catalog

E > 100 MeV

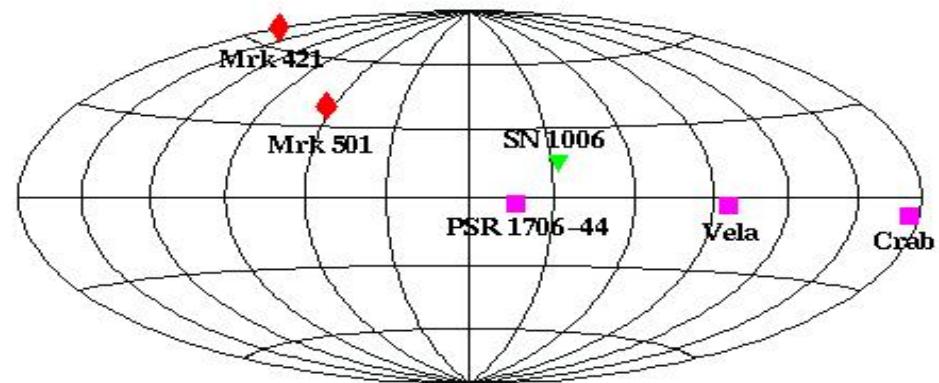


Non identified sources
 Blazars (AGN)
 Pulsars
 Supernova remnants
 LMC

After COS-B, the great success of EGRET has been
 the rather unexpected discovery that » 100 blazars
 are brilliant in the GeV energy range
with no turn-over at maximum energies

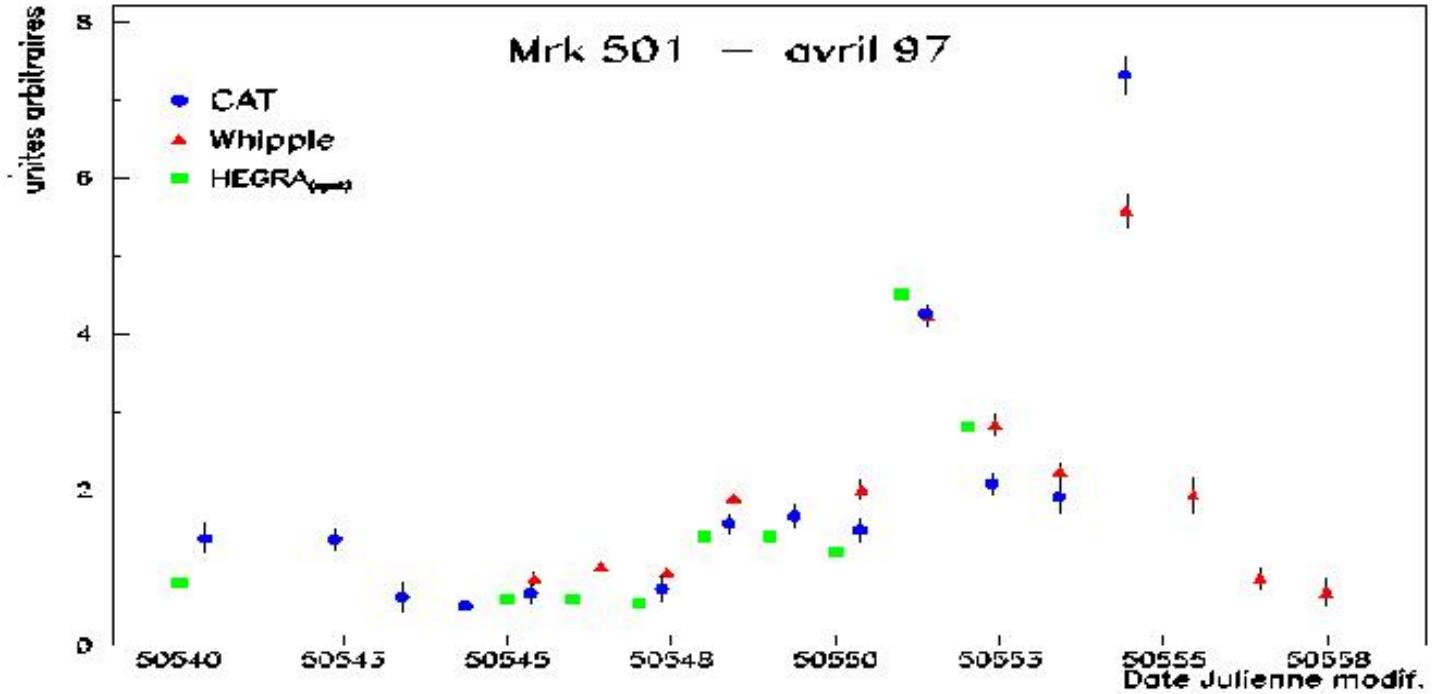
Sources seen from the ground

E > 250 GeV



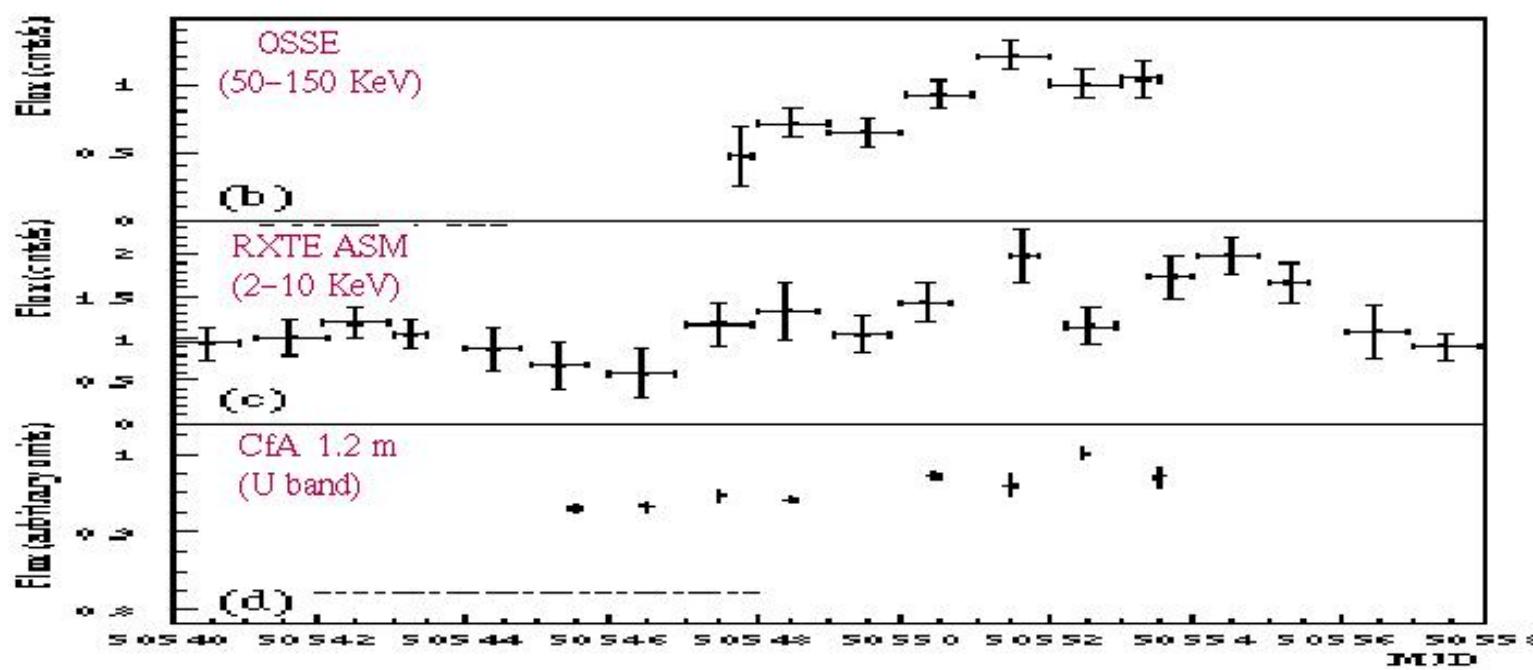
During about the same decade
 the TeV observations broke through,
 with no pre-notice.
 The a posteriori surprise is the scarcity of sources

- the intermediate region, from 30 to 300 GeV, remains unexplored
- the (few) TeV blazars are weak EGRET sources
- the absorption ($gg \rightarrow e^+e^-$) obscures the far Universe for TeV g's

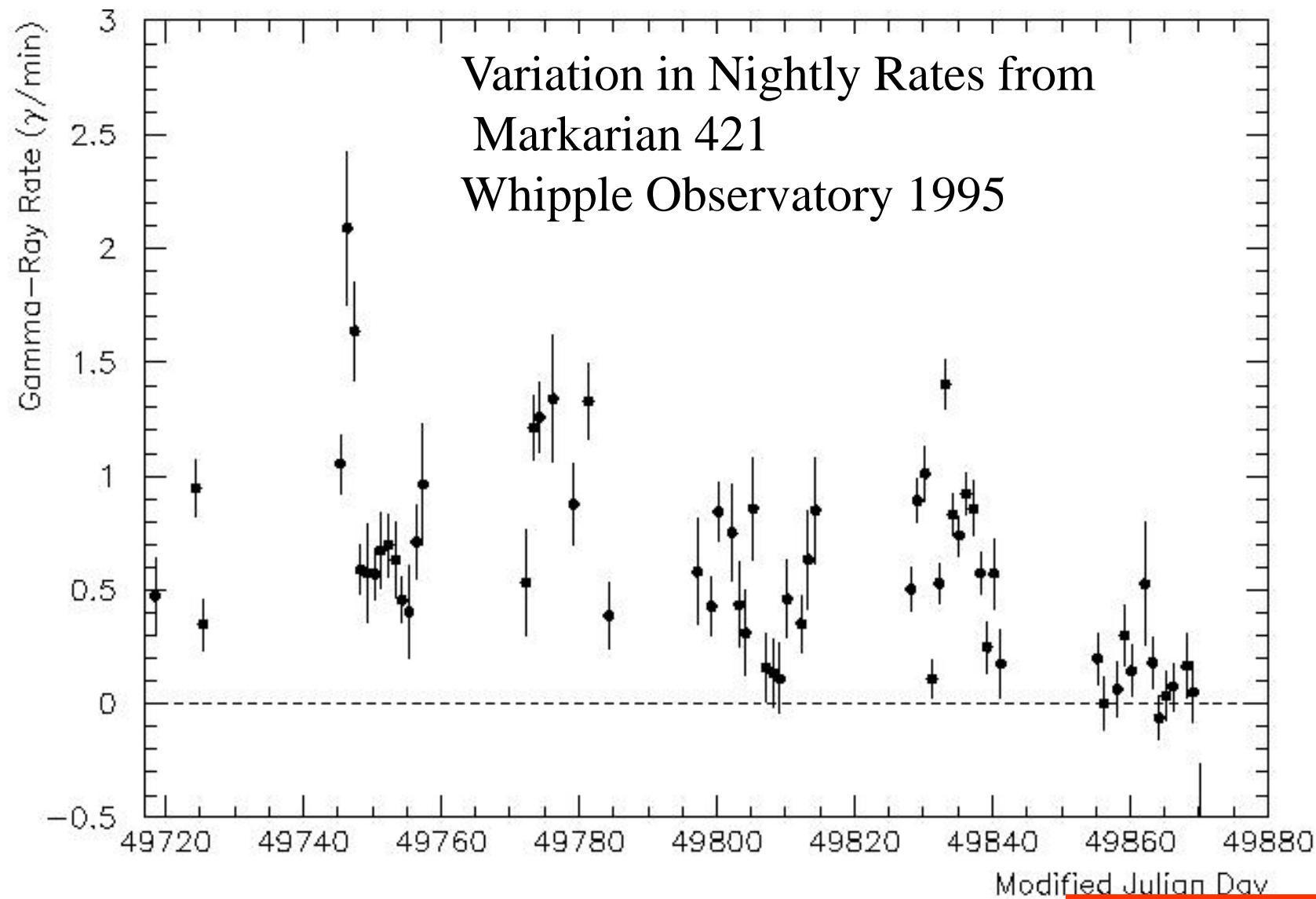


Mrk 501 flares of April 1997 showing variability on a day scale

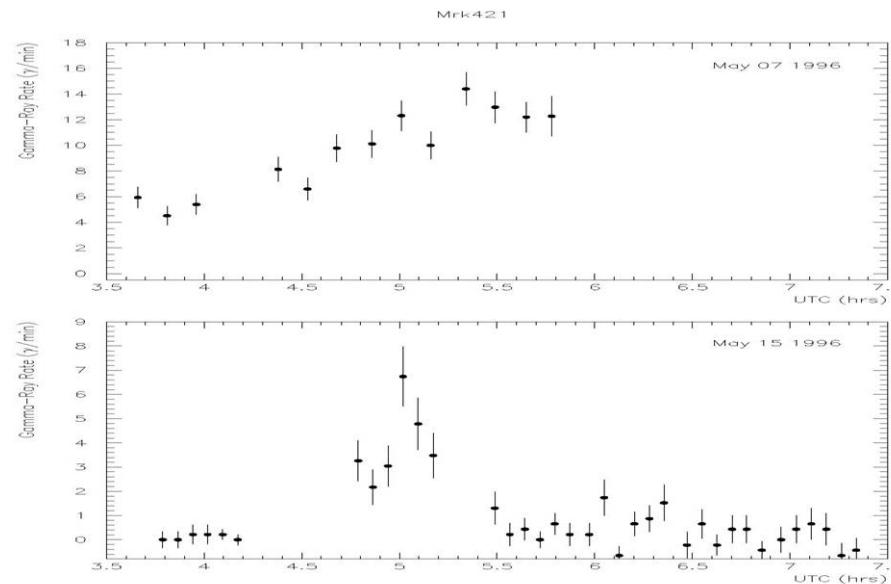
Even faster (<hr) variations have been evidenced by the Whipple on Mrk 421



Flux sensitivity
on the scale of
the hour is needed



Buckley et al. 1996



Markarian 421

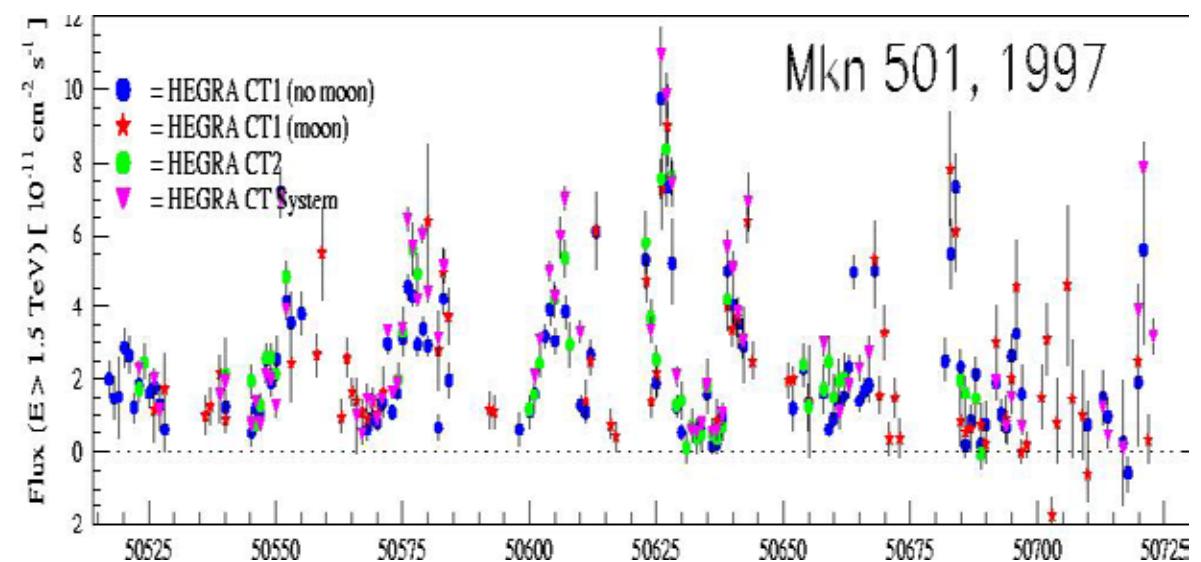
Whipple Observatory

Hours - minutes

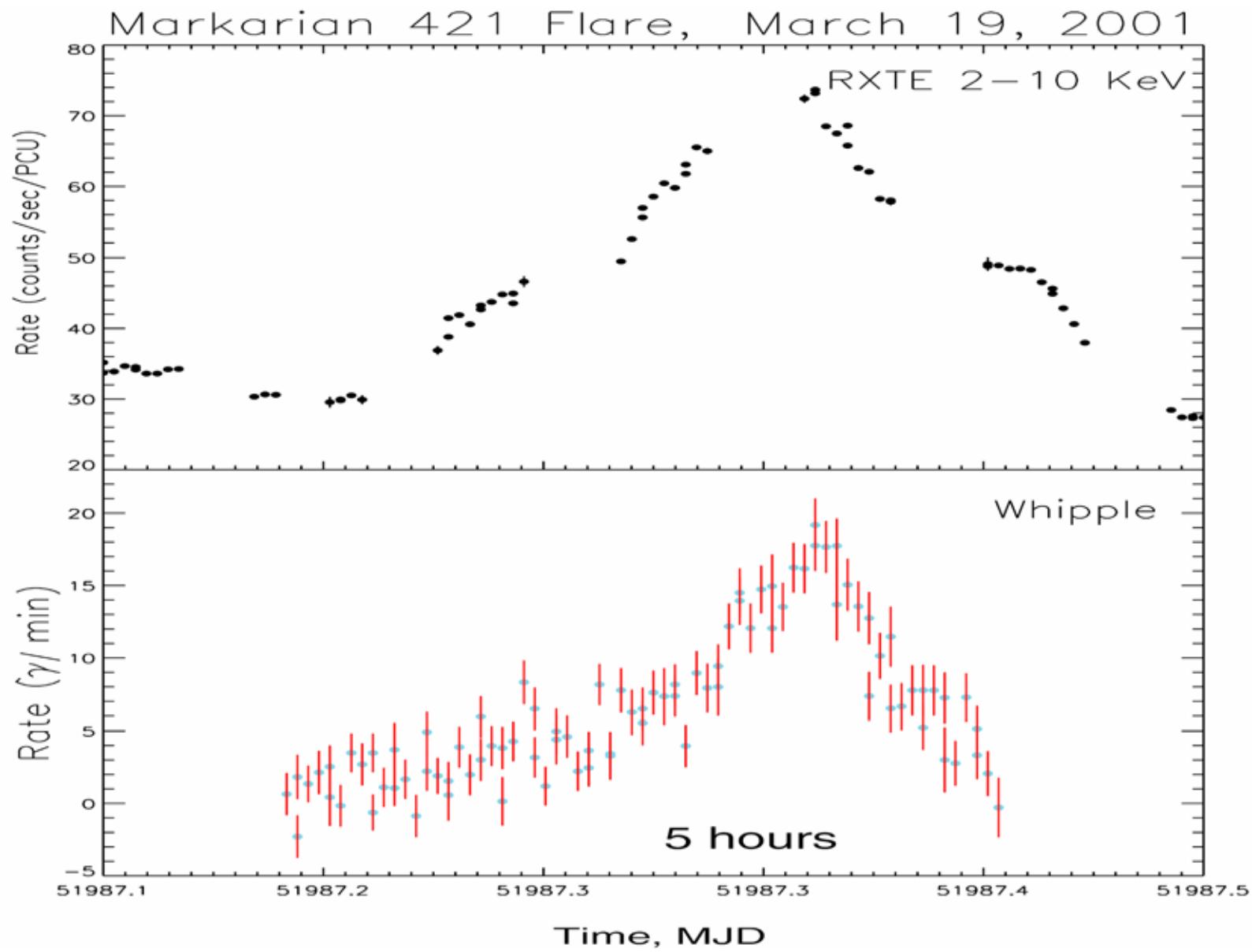
Markarian 501

HEGRA

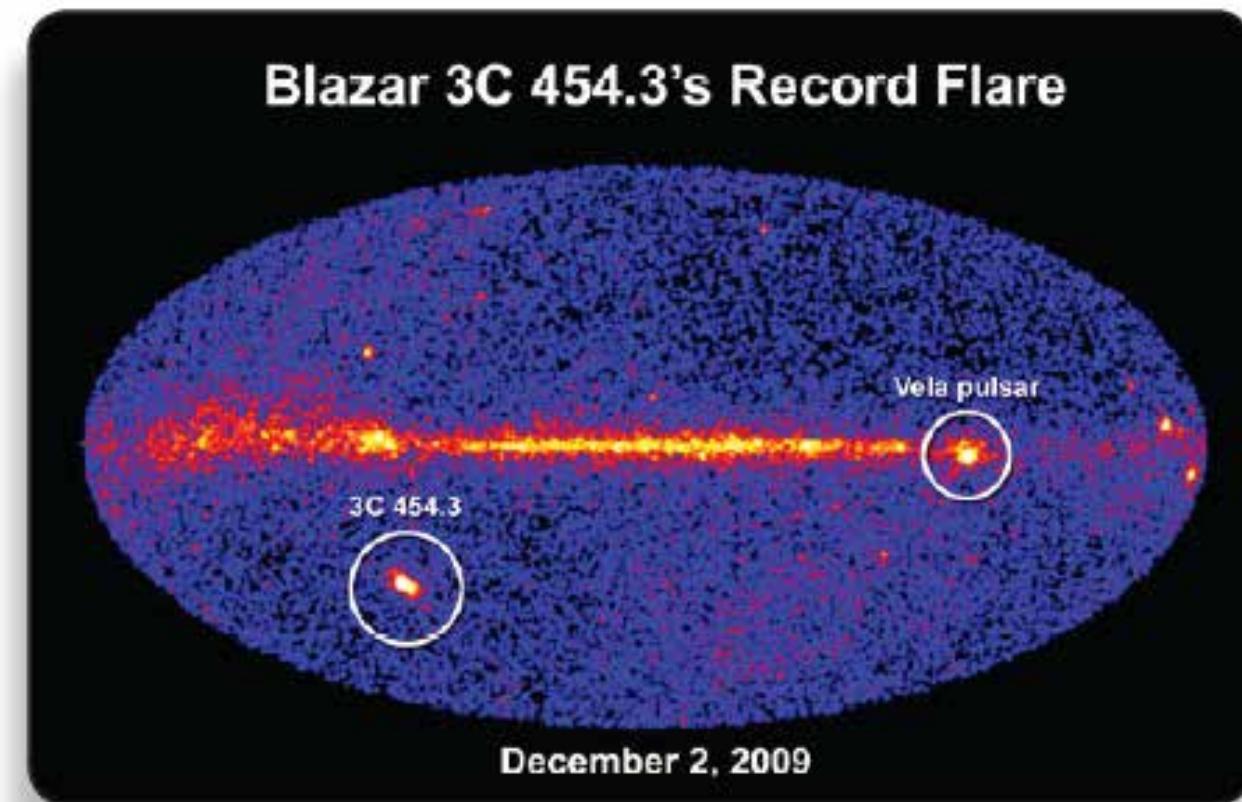
Days-months



Time coincidence X and TeV g

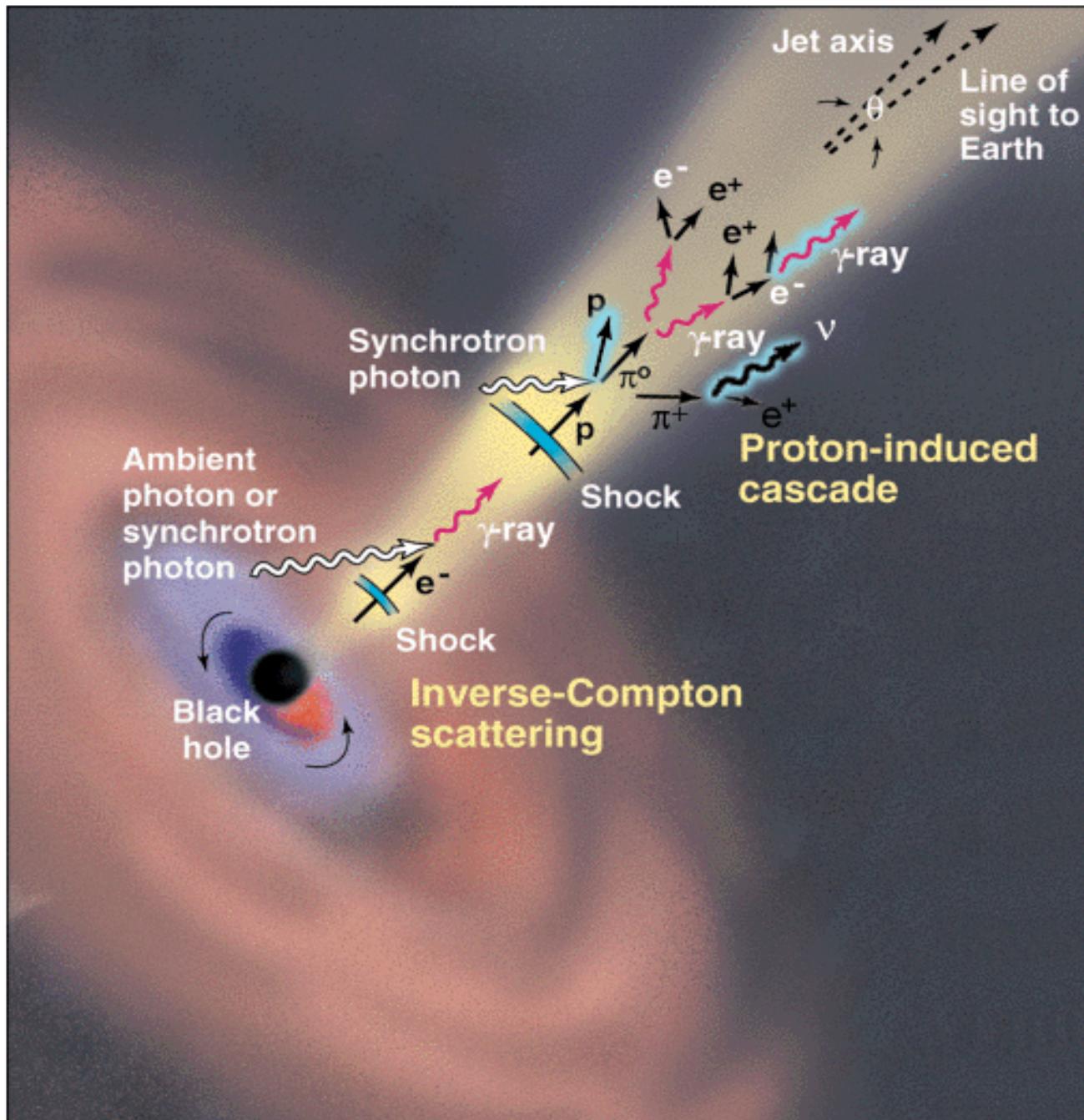


Flare de 3C454.3 : 2 x luminous than Vela although 10^6 x distant !



Unprecedented flares from the blazar 3C 454.3 in the constellation Pegasus now make it the brightest persistent gamma-ray source in the sky. That title usually goes to the Vela pulsar in our galaxy, which is millions of times closer (Credit: NASA/DOE/Fermi LAT Collaboration).

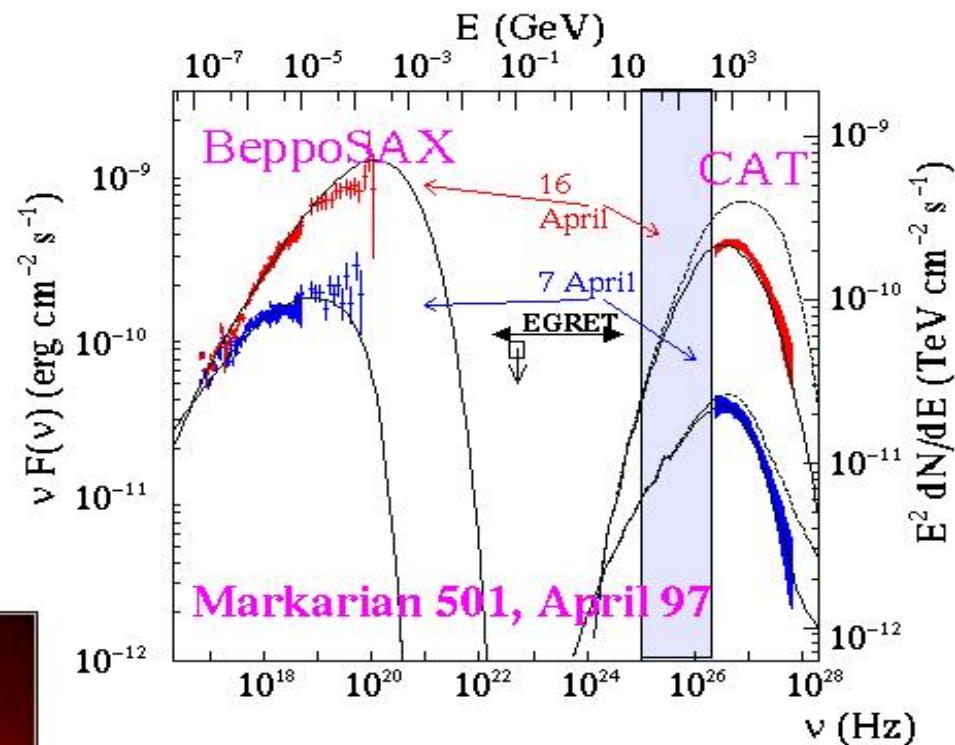
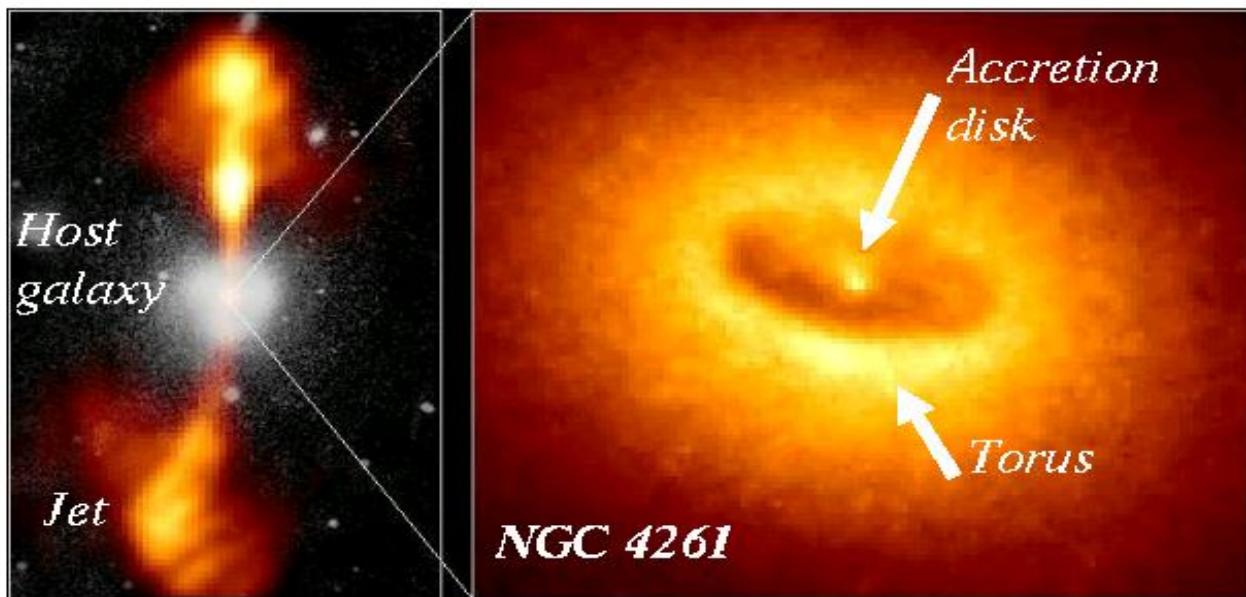
AGN Jet Emission Mechanisms



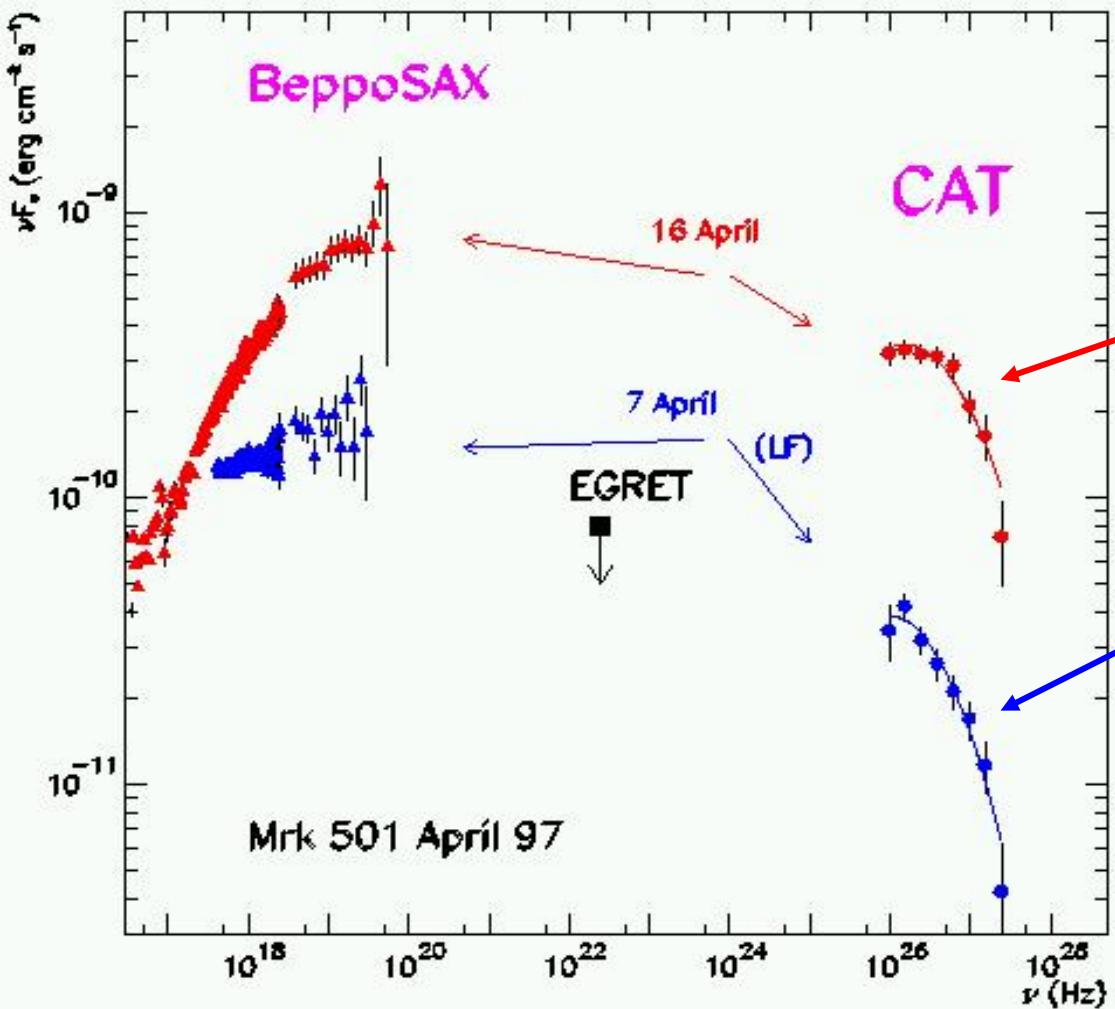
Cosmic Cannon:
Looking down the
Barrel of the Cannon

Electron Progenitors:
Synchrotron Self Compton
External Compton
Proton Progenitors:
Proton Cascades
Proton Synchrotron

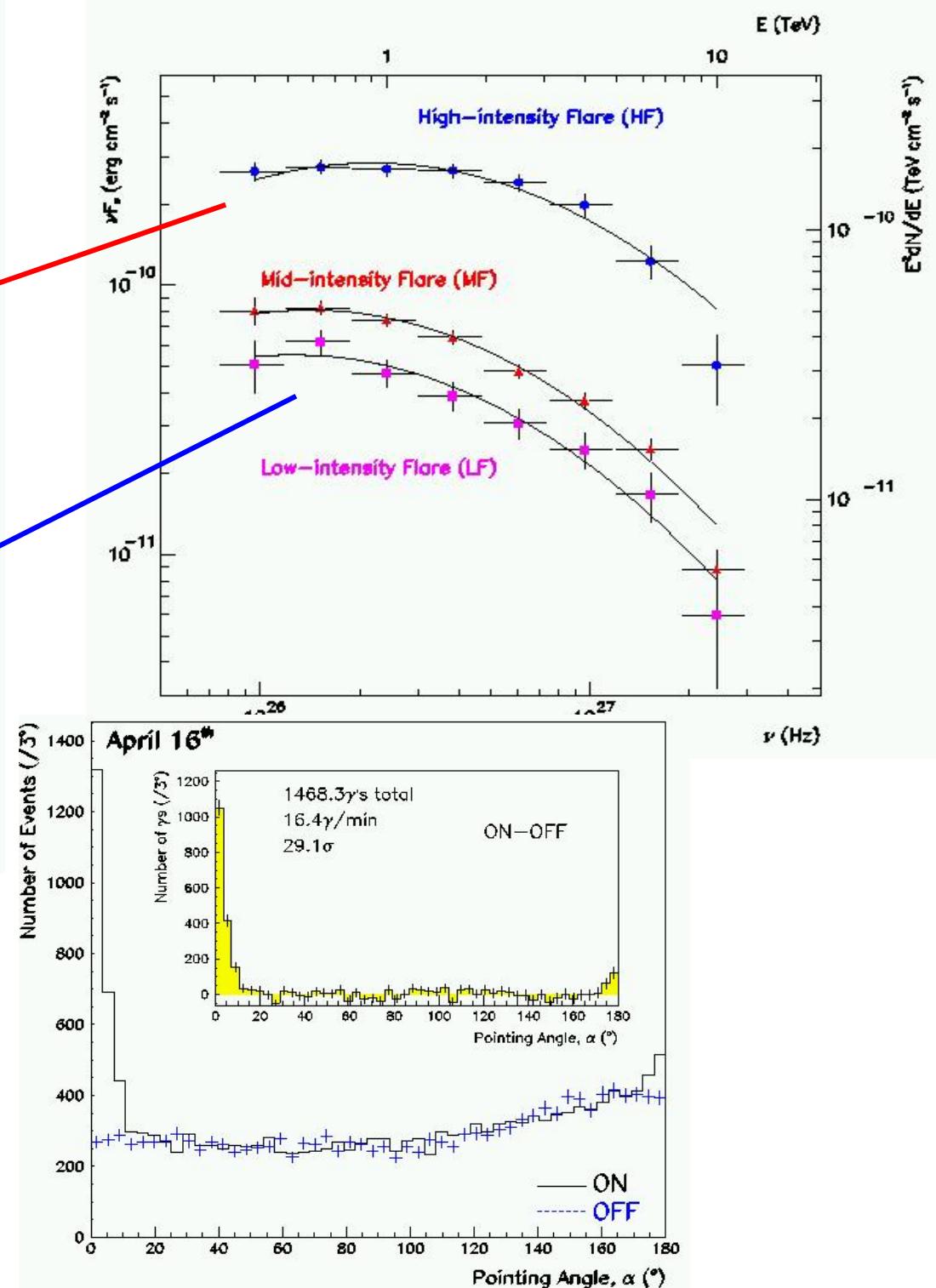
- ❑ AGN / Blazar : Markarian 501
 - ❑ Supermassive Black Hole ($M > 10^6 M_\odot$)
 - ❑ With an accretion disk (1 pc)
 - ❑ Surrounded by a torus of dust (100 pc)
- ❑ Radio loud:
Two jets (< 100 kpc)



- ❑ *Blazars*
- ❑ Near jet axis : $\theta \leq 1/\gamma$
→ High energy emission
- ❑ Strong variability on time scale of a *day* or even less

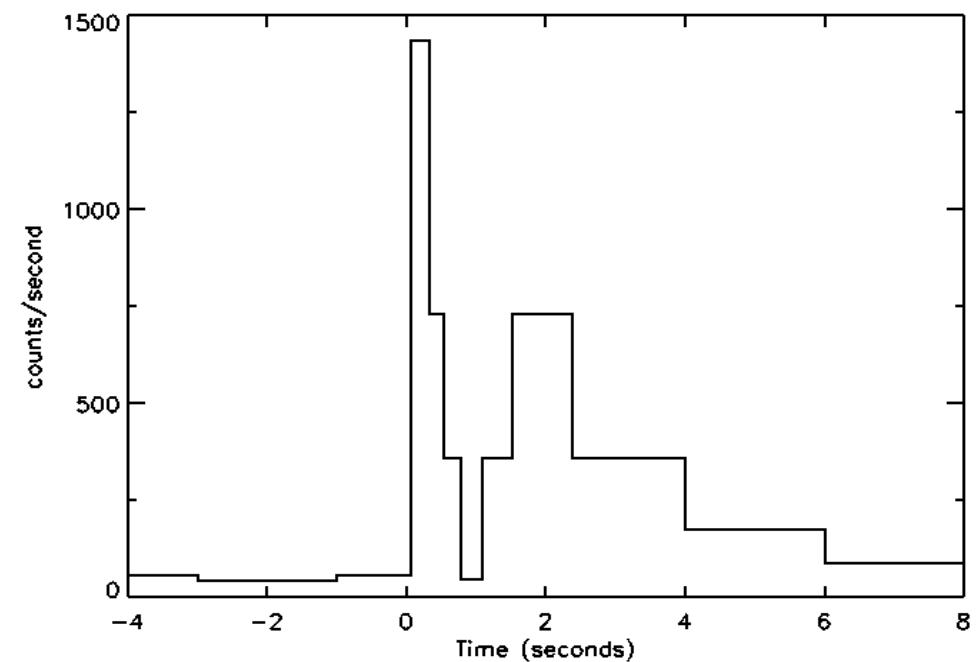
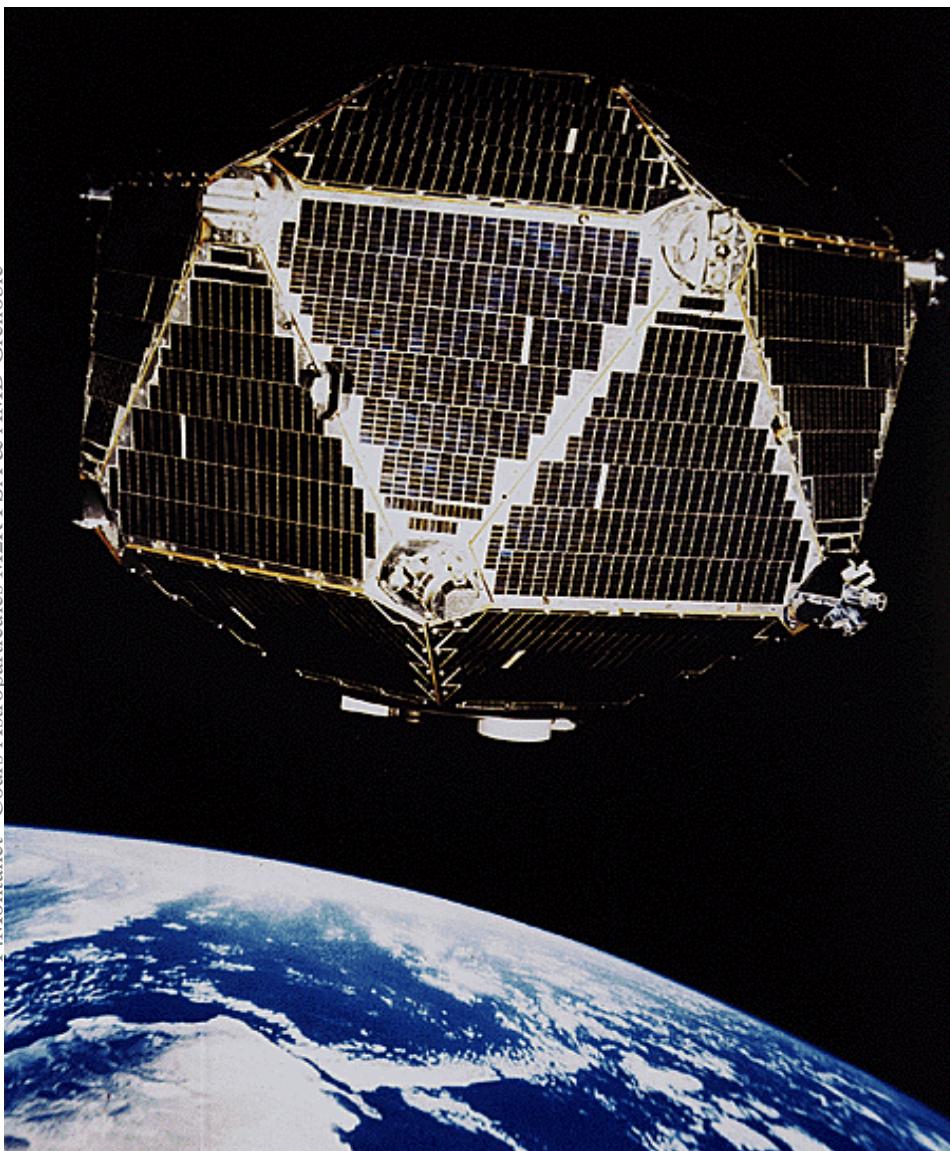


Markarian 501 Flare of April 1997



GAMMA RAY BURSTS

Vela Program (1969-1979)

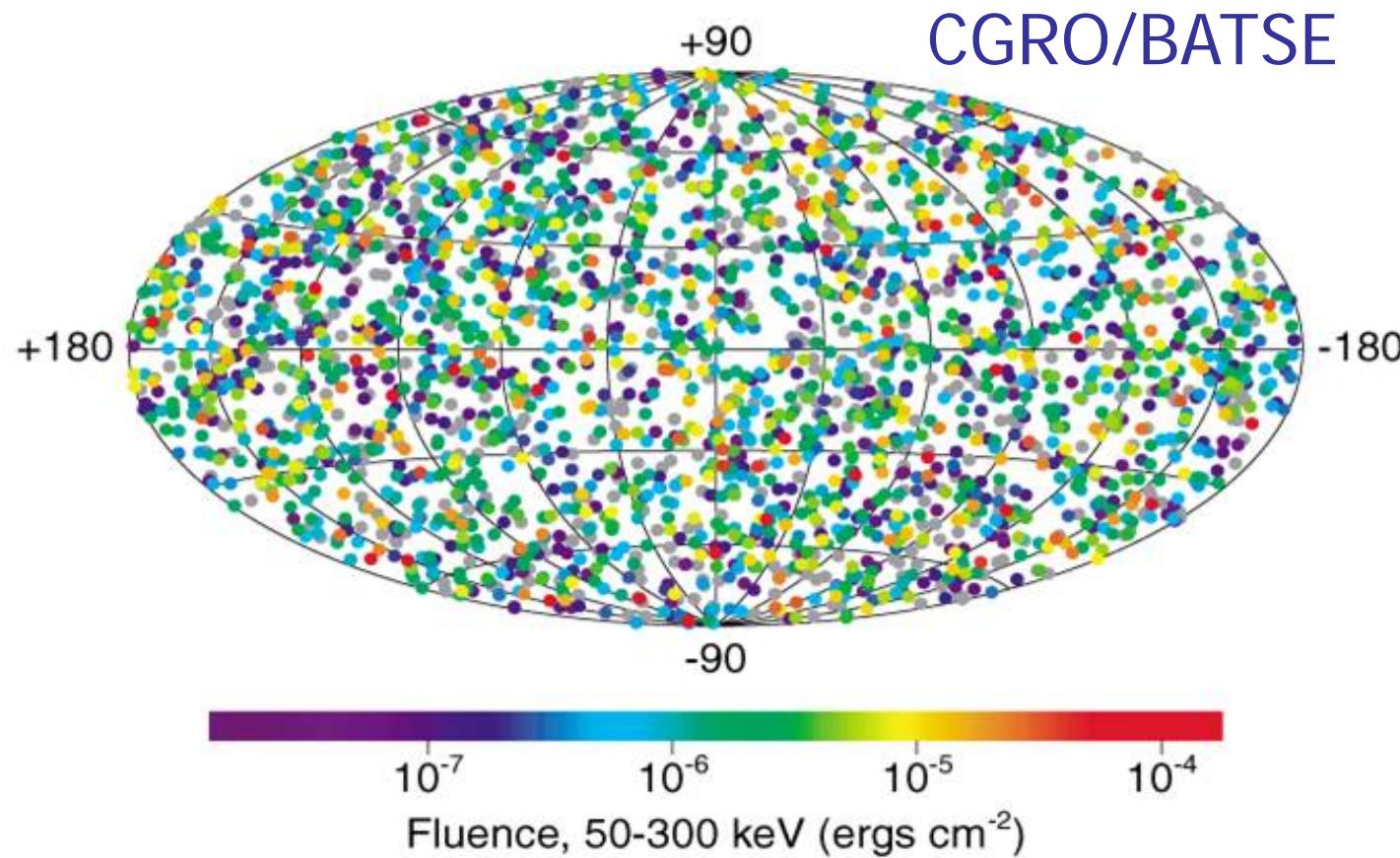


- Discovered in 1967 when spying thermonuclear bomb tests.
- A 30+ years old mystery unraveled in the 90ties !

Gamma-ray Burst Sky

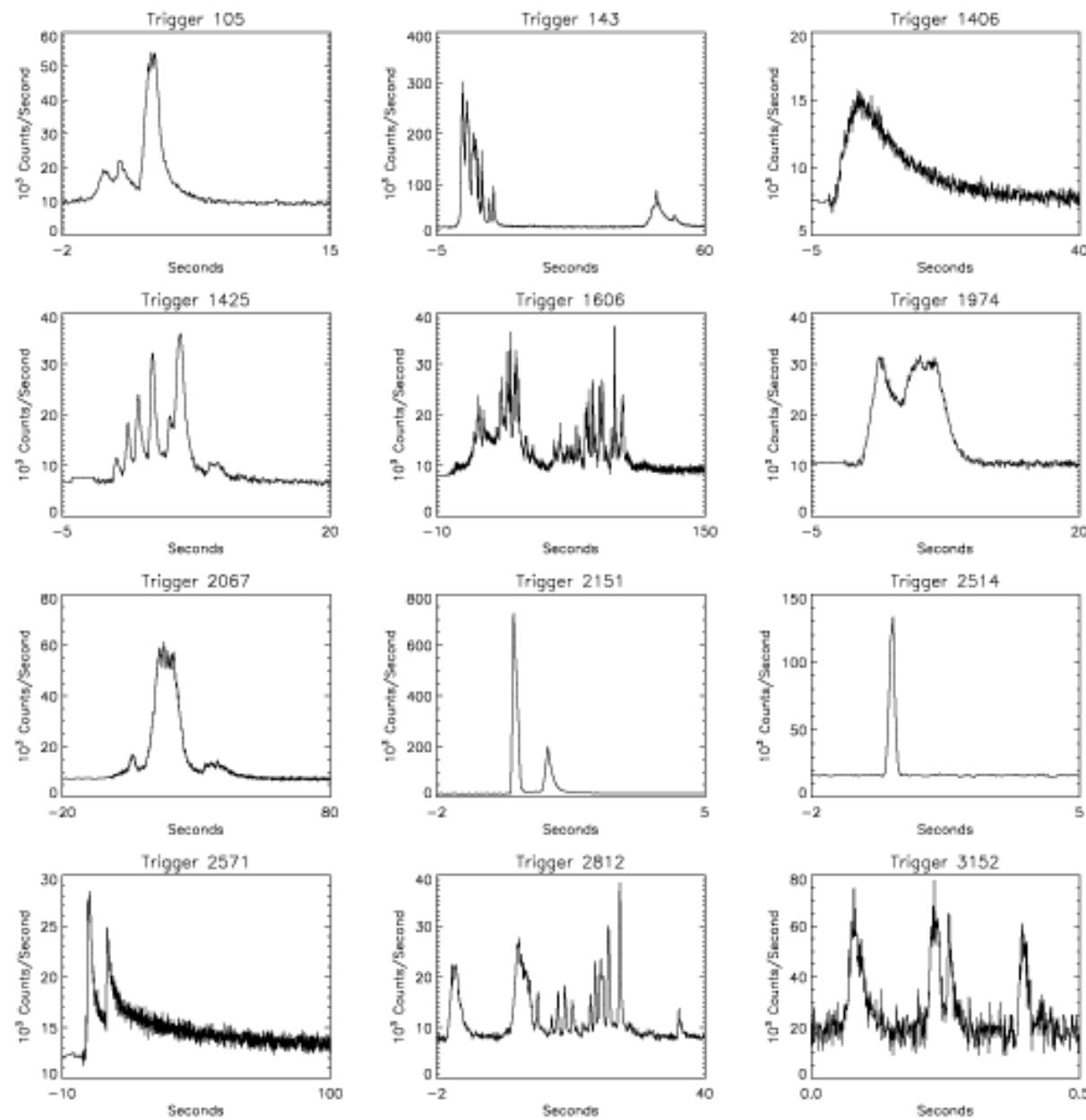
2009-2010

F.Montanet Cours Astroparticules M2R PSA & AMD Grenoble



~ Once a day, anywhere in the universe !

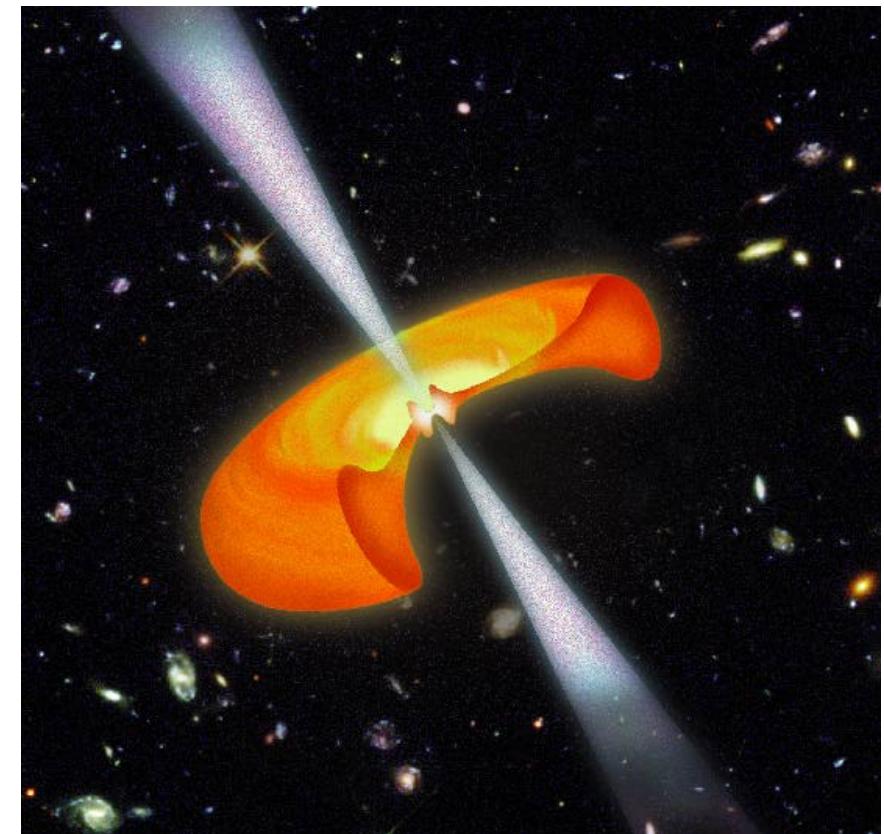
GRB



Hypernova

2009-2010

Death of a massive star ?

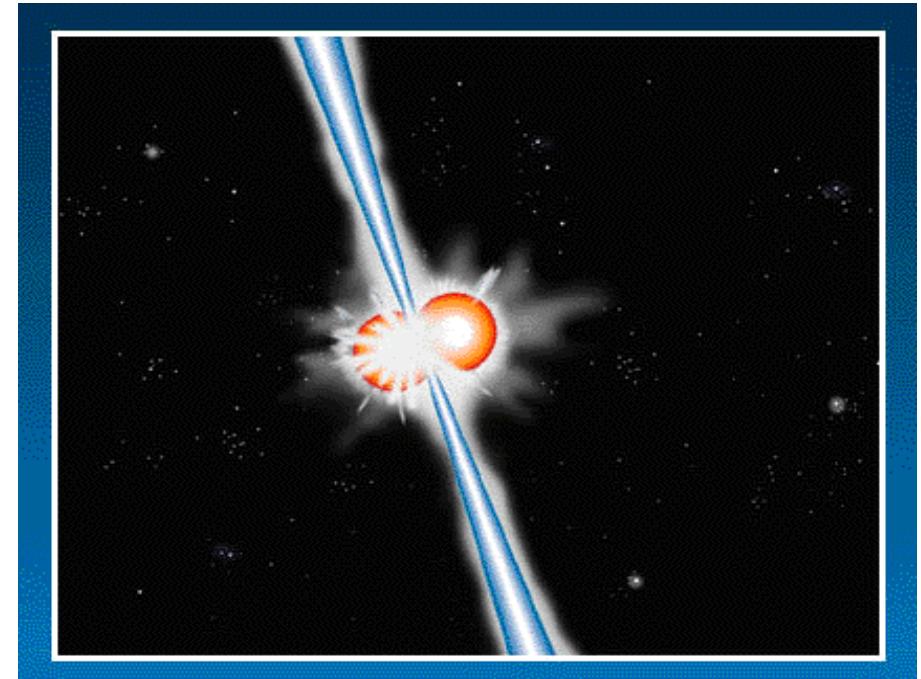
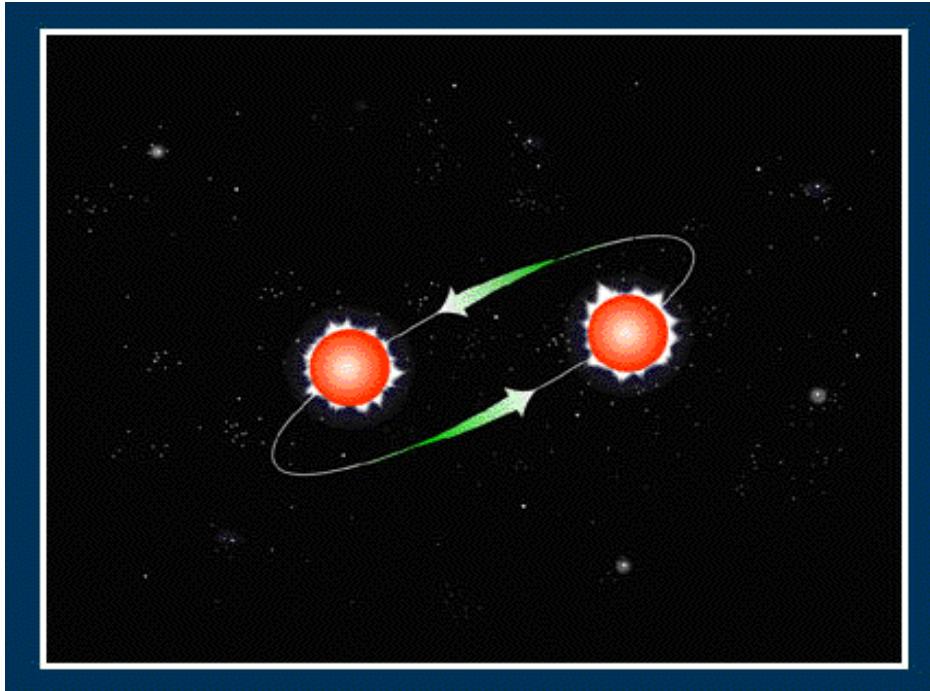


- A billion trillion times the power from the Sun

Catastrophic Mergers

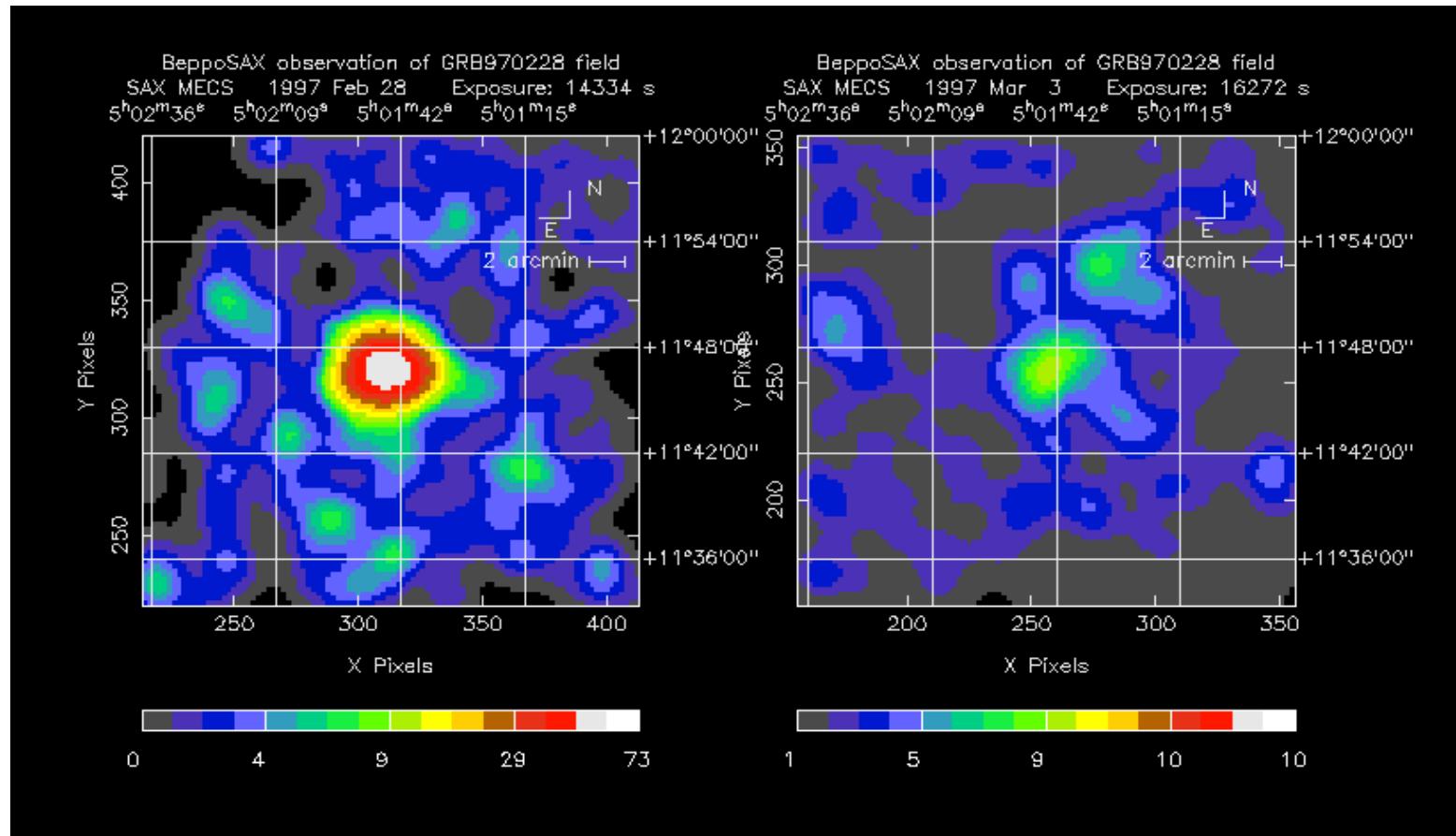
2009-2010

F.Montanet Cours Astroparticules M2R PSA & AMD Grenoble



- Spiral death of 2 neutron stars or black holes

Afterglow

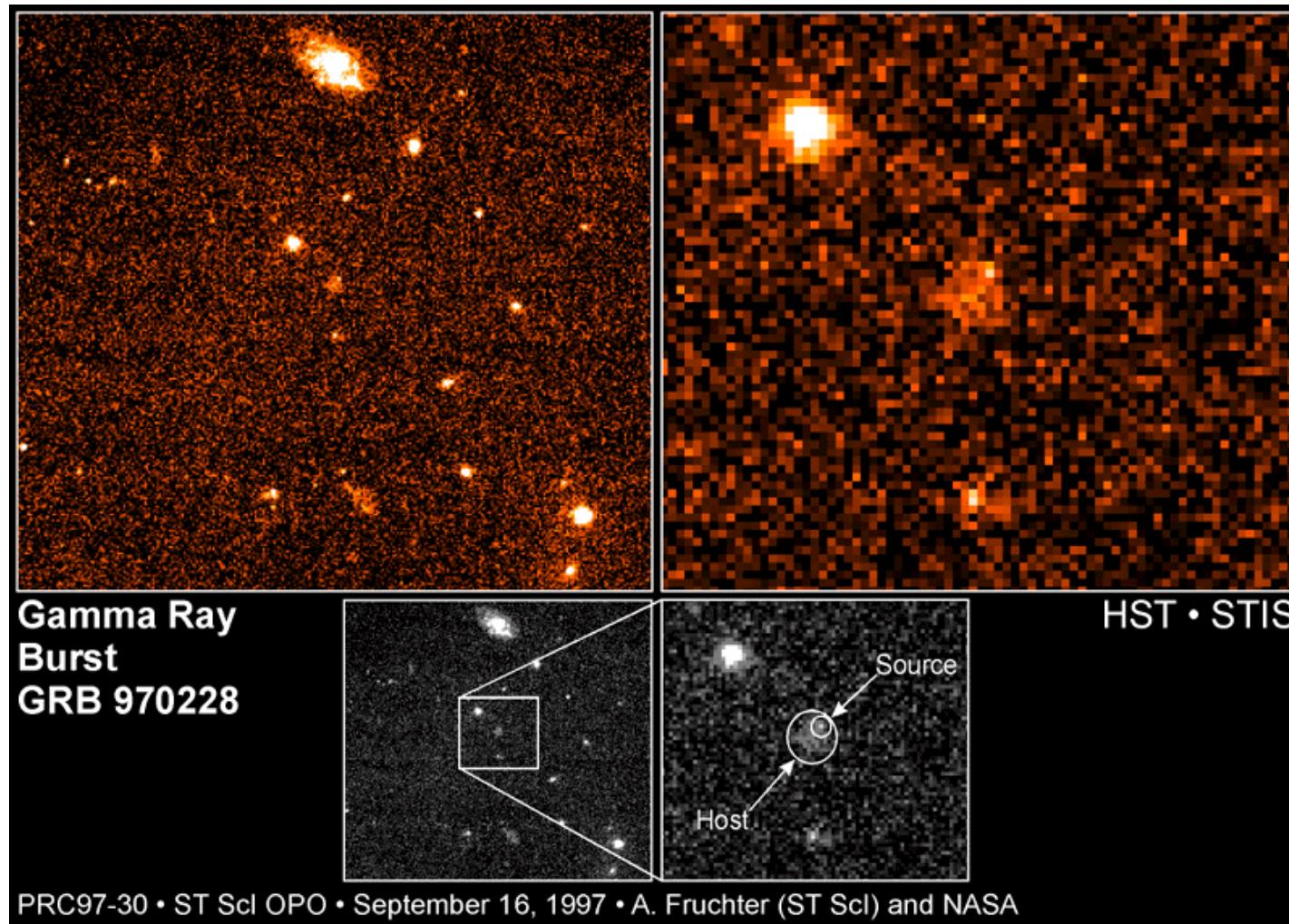


- Discovered in 1997 by BeppoSAX satellite

GRB afterglow en optique

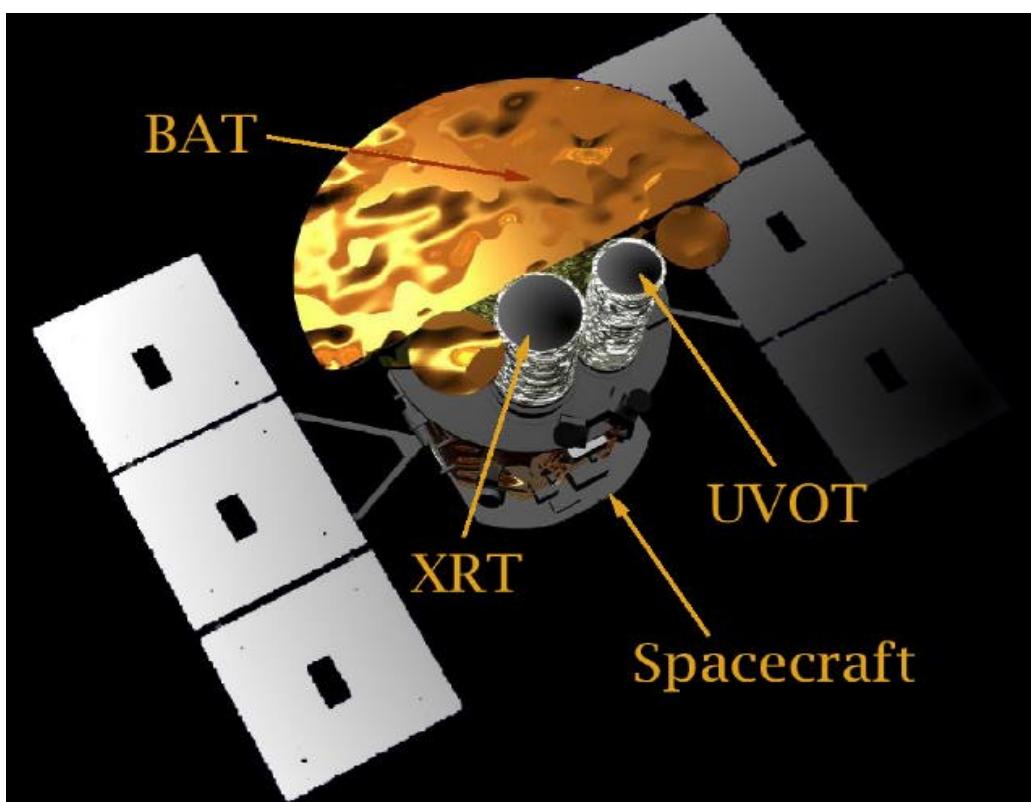
2009-2010

F.Montanet Cours Astroparticules M2R PSA & AMD Grenoble



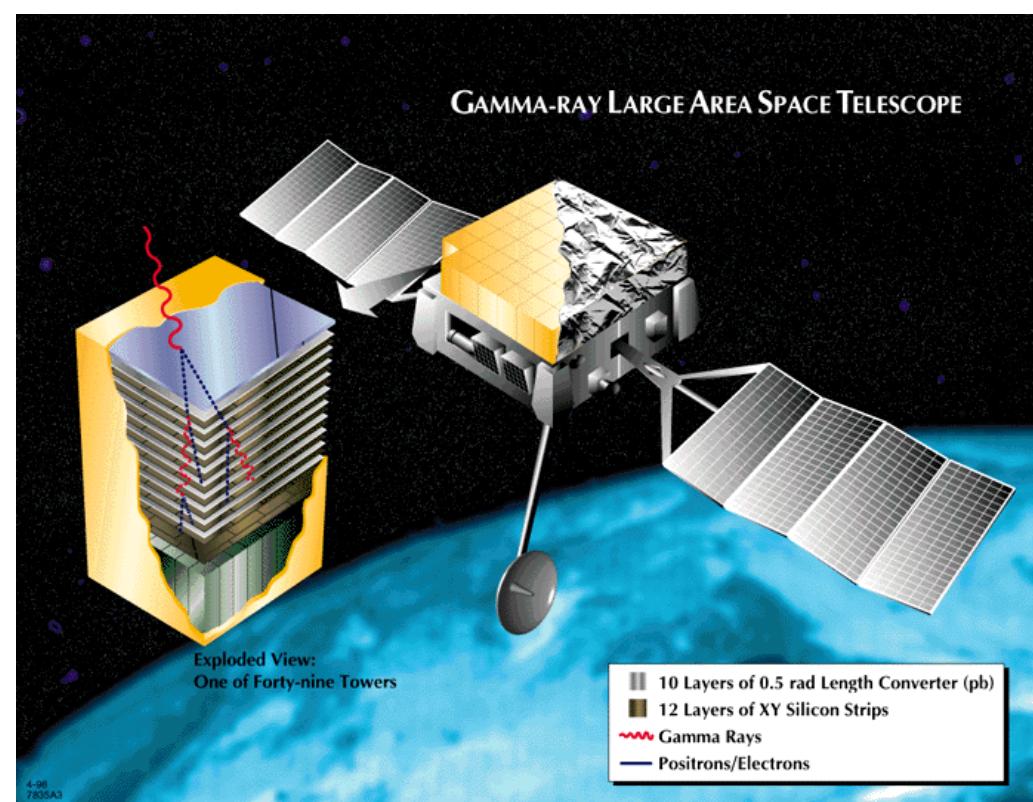
New Missions = Better Data

HETE II (launched 10/9/00)



INTEGRAL (launched 17/10/2002)

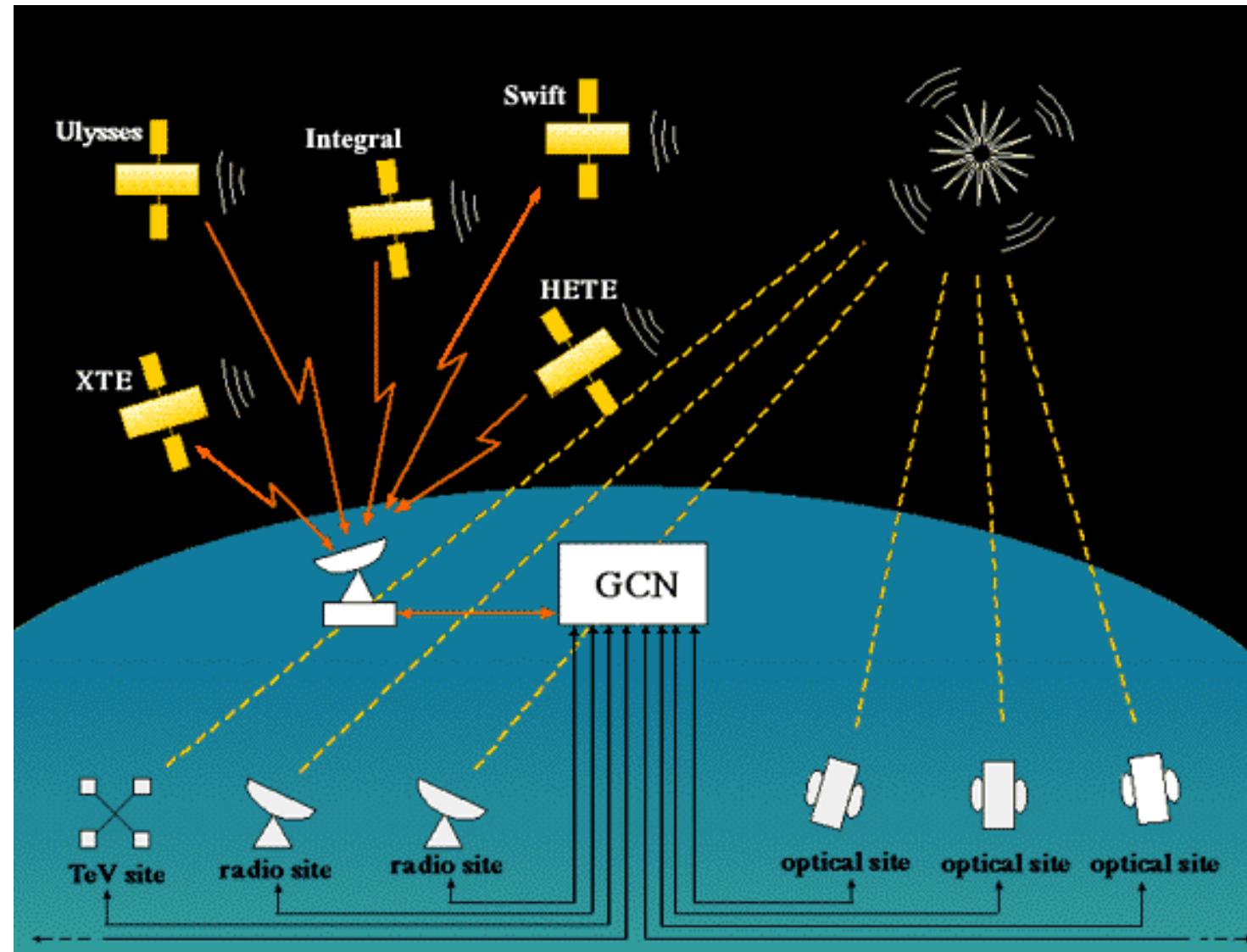
FERMI (launched 16/5/2008)



The Gamma ray bursts Coordinates Network

2009-2010

F.Montanet Cours Astroparticules M2R PSA & AMD Grenoble



GRB's

- Gamma-Ray Bursts : intense gamma-ray flux
 - 0.1 to 1 MeV, and up to 100 MeV
 - Emitted on a short time scale (~ 1 seconde !)
 - Observed ~once per day, with an isotropic distribution
 - Source at cosmological distances (most distant is $z = 9.4$!!)
- γ -ray luminosity: $\sim 10^{52}$ erg/s (10 \leq SNe !)
- Extreme variability in intensity and spectrum
 - Time scales from **10 ms** to **10 s**
 - Some very short **1 ms** variabilities \circledR internal shocks
- Clear bimodality suggesting the existence of two separate populations:
 - a "short" population with an average duration of about 0.3 seconds
 - a "long" population with an average duration of about 30 seconds.

UHE SOURCES

General limits on models

Shock waves

- Acceleration site confinement: $r_g = E/ZeBc < L$
 - Depends in fact on $V_{\text{shock}} = bc$:

$$E_{\max} \approx b' Ze' Bc' L$$

Unipolar induction

- Accelerate in one step (E field, $f_{\text{Lorentz}} \approx V_{\text{rot}}' B' R$)
 - ® No confinement necessary

Top-Down Models

- No acceleration at all !
 - Decay products of exotic physics states, supermassive particles at E_{GUT} , Topological Defects...

Bottom-Up

ZeVatrons

Astrophysical Accelerators reaching ZeV

Acceleration = Fermi-like diffuse acceleration.

Frist challenge, E_{\max} : Reach $\geq 10^{20}$ eV

(if 1 TeV is hard, guess for 10^9 TeV !!)

Second challenge, Propagation : B_{igm}
(determine the spectrum and the arrival)

Hillas criterion:

Magnetic confinement in the shock zone i.e.

$1 R_{\text{gyr}} < \text{accelerator size}$

Not many candidates survive!

Neutron stars (pulsars)

AGNs

Radio lobes

Clusters

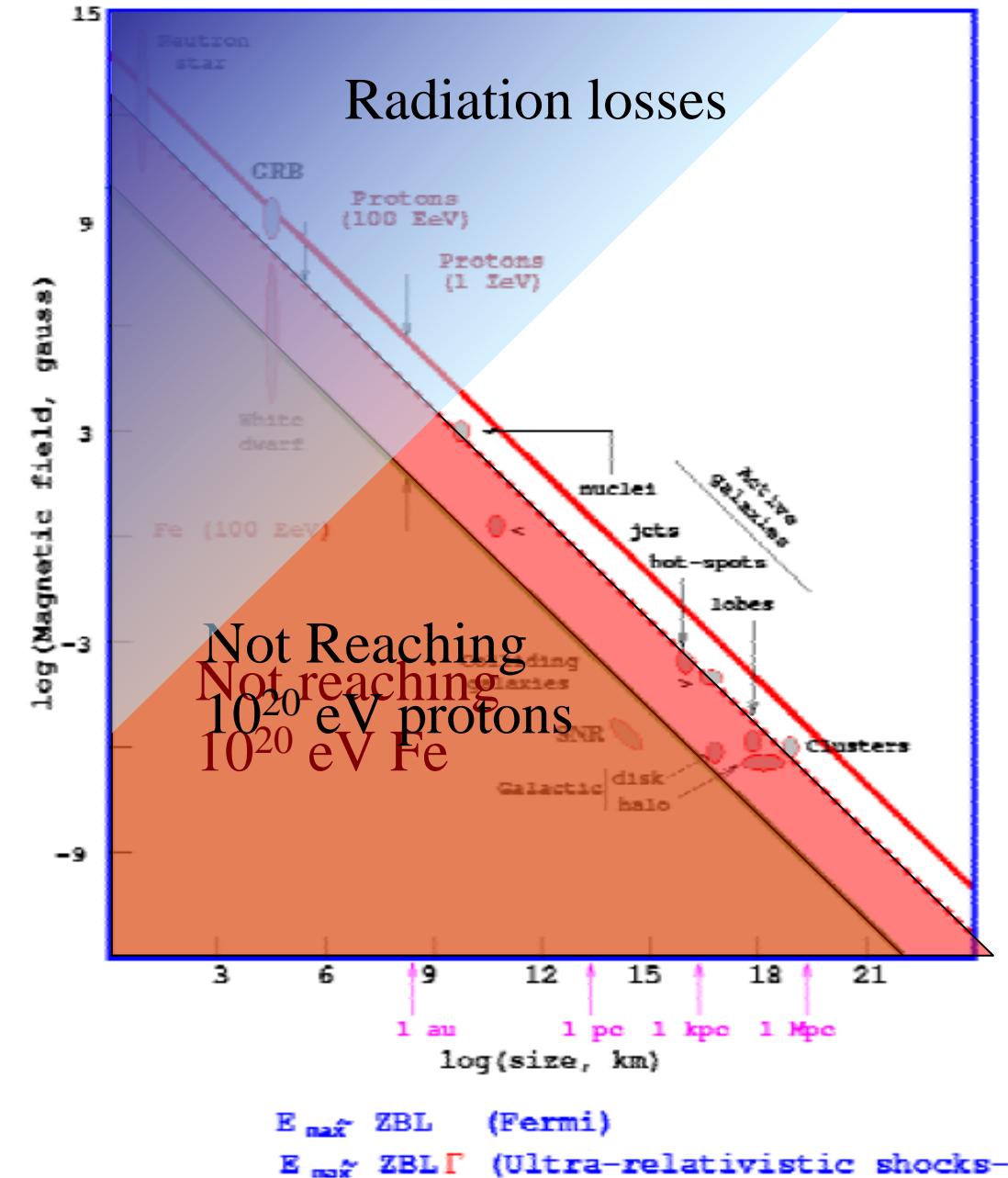
Colliding Galaxies/Clusters

Gamma Ray Bursts

Hillas diagram

Standard estimates for E_{\max} :

- Confinement :
 $r_g = E/(ZeB) < R \Rightarrow E < ZeBR$
- Unipolar inductor (pulsar)
 $E < ZeBR(\Omega R/c) \approx \beta_s ZeBR$
- Diffusive acceleration
 by non relativistic shocks:
 $\tau_{acc} \approx 10\kappa/u_s^2 < R/u_s$ avec $\kappa > r_g c$
 $\Rightarrow E < \beta_s ZeBR$
- Diffusive acceleration
 by relativistic shocks:
 $E < \Gamma_s ZeBR$
- General Hillas condition:
 $E < 0.9\beta\Gamma ZeB_{Gauss}R_{pc} \text{ ZeV}$



Relativistic shocks

- Acceleration / G^2 works fine for a couple of cycle
- After that it fails for mere kinematical reasons
- But this is still very efficient (\gg standard shocks)
- Confinement is easier
 - A weak deflexion is enough : $\delta\theta \approx 1/\Gamma_s \Rightarrow r_g < R_s/\Gamma_s$
 - $\Rightarrow E_{max} \approx \Gamma_s \times$ larger
 - \Rightarrow one can reach the limits induced by energy losses

BOTTOM -UP

Galactic pulsars

Extragalactiques radio galaxy lobes

Gamma Ray Bursts

Protons, Iron, Nuclei?

Spectral index

Explaining isotropy is not trivial

Angular coincidences to be confirmed...

Top-Down

Topological defects, superheavy relics
with $M \sim$ GUT scale that is $\sim 10^{16}$ GeV

- Energy $\gtrsim 10^{20}$ eV easy!
(QCD fragmentation spectrum QCD with $M \sim 10^{24}$ eV!!)
- Explaining the flux is not trivial !!
(natural density scale is $\sim H_0^{-1}$)
- Composition of UHECR is the clue (photons + neutrinos) !!

Low energy gamma-rays constraint

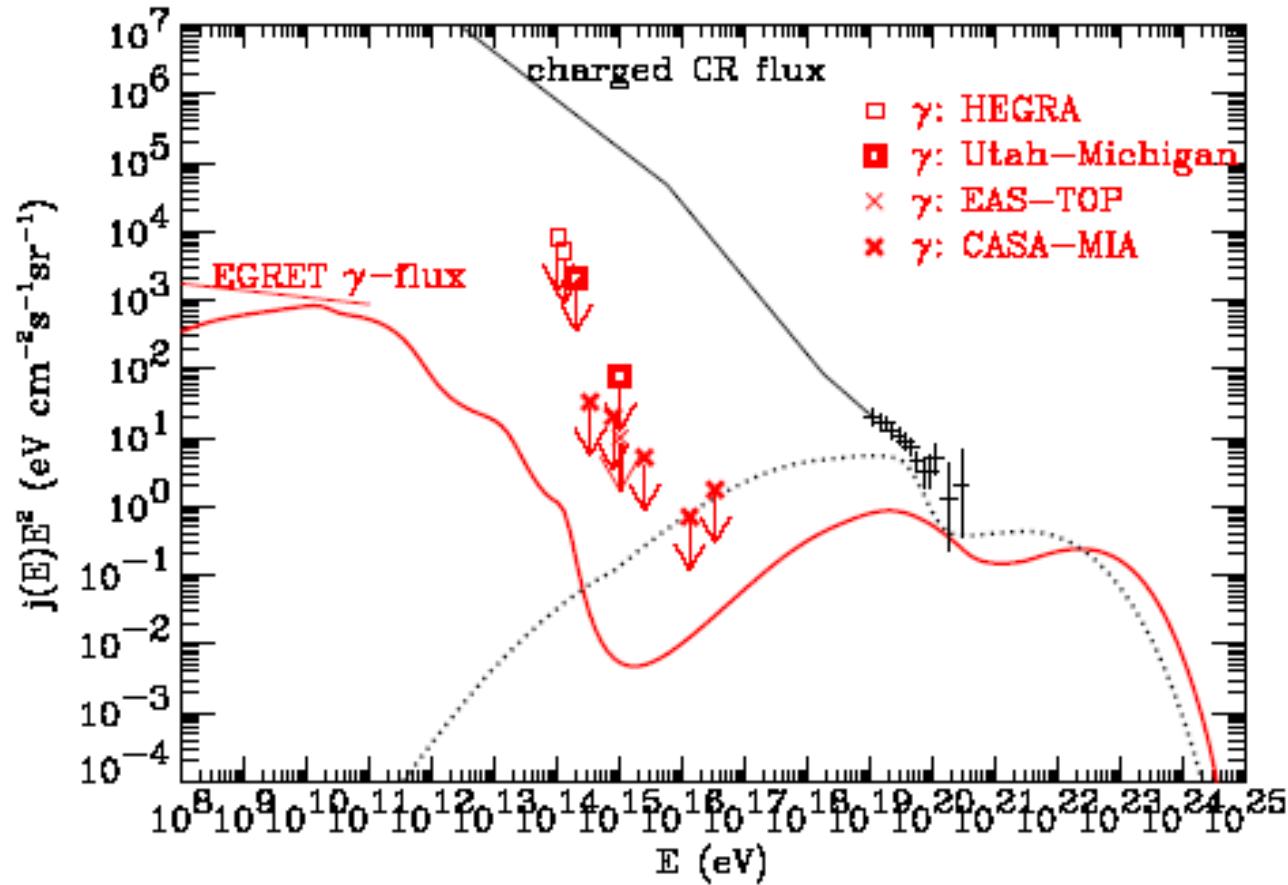


Figure 28: Predictions for the differential fluxes of γ -rays (solid line) and protons and neutrons (dotted line) in a TD model characterized by $p = 1$, $m_X = 10^{16}$ GeV, and the decay mode $X \rightarrow q + q$, assuming the supersymmetric modification of the fragmentation function, Eq. (57), with a fraction of about 10% nucleons. The calculation used the code described in Ref. [206] and assumed the strongest URB version shown in Fig. 10 and an $EGMF \ll 10^{-11}$ G. 1 sigma error bars are the combined data from the Haverah Park [3], the Fly's Eye [7], and the AGASA [8] experiments above 10^{19} eV. Also shown are piecewise power law fits to the observed charged CR flux (thick solid line) and the EGRET measurement of the diffuse γ -ray flux between 30 MeV and 100 GeV [185] (solid line on left margin). Points with arrows represent upper limits on the γ -ray flux from the HEGRA [257], the Utah-Michigan [510], the EAS-TOP [511], and the CASA-MIA [258] experiments, as indicated.

Top-Down Signatures

Composition:

Flux de Photons, Neutrinos \wedge Protons

The current (AUGER) limits on UHE neutrino and photon flux
already kill most Top-Down models !!

Spectrum:

QCD-like fragmentation spectrum quite "hard"

Cosmography:

Halo distribution!! (SHRs & TDs locales)
or ~ Homogeneous

and even more exotic stuff...

Strongly interacting neutrinos

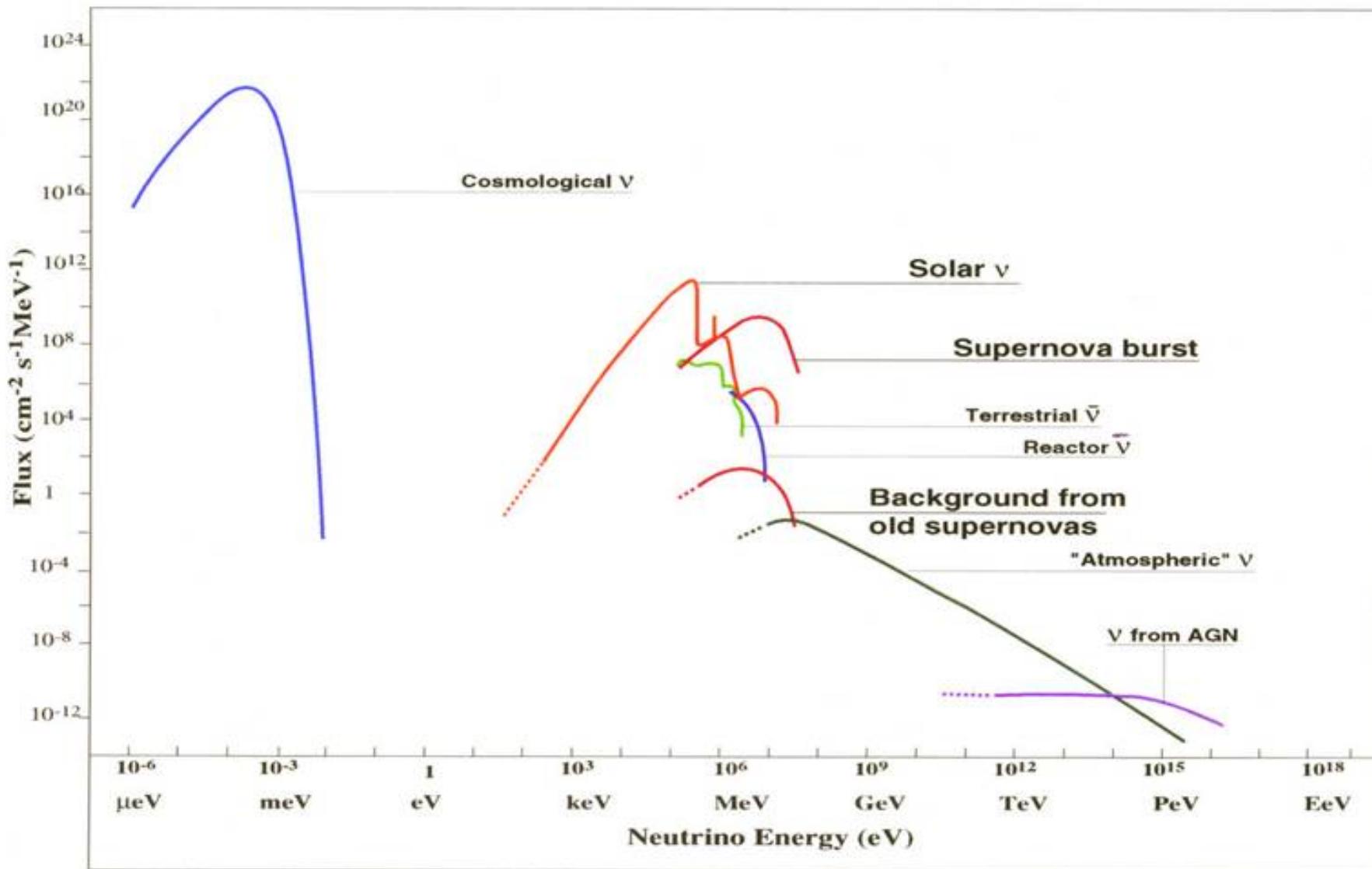
Lorentz Invariance Violation

Special Relativity Violation

etc...

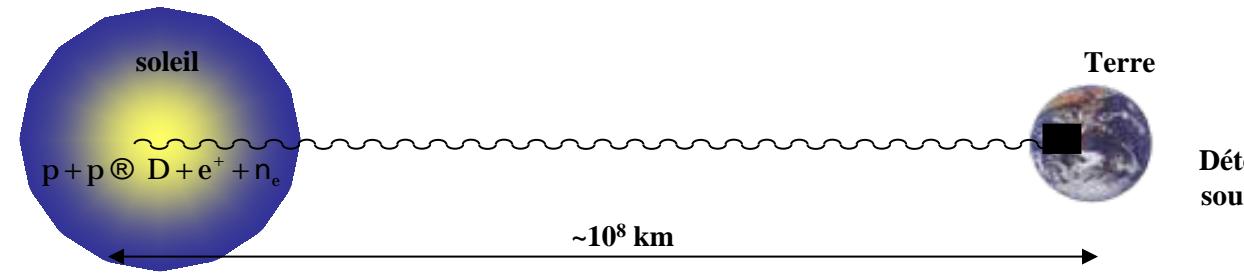
NEUTRINOS SOURCES

The overall neutrino spectrum

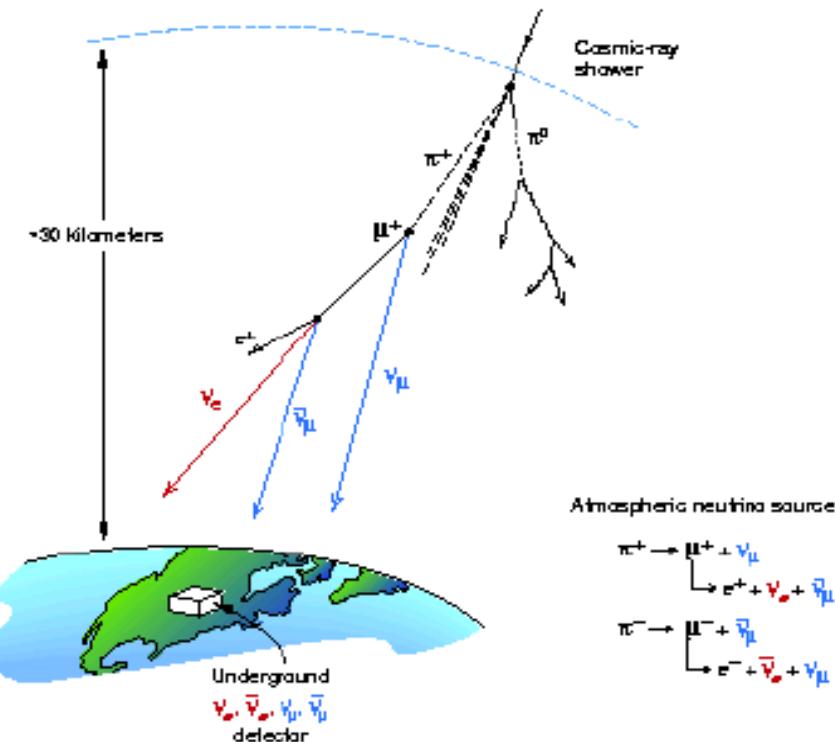


Natural Neutrinos Sources

Solar
Neutrinos

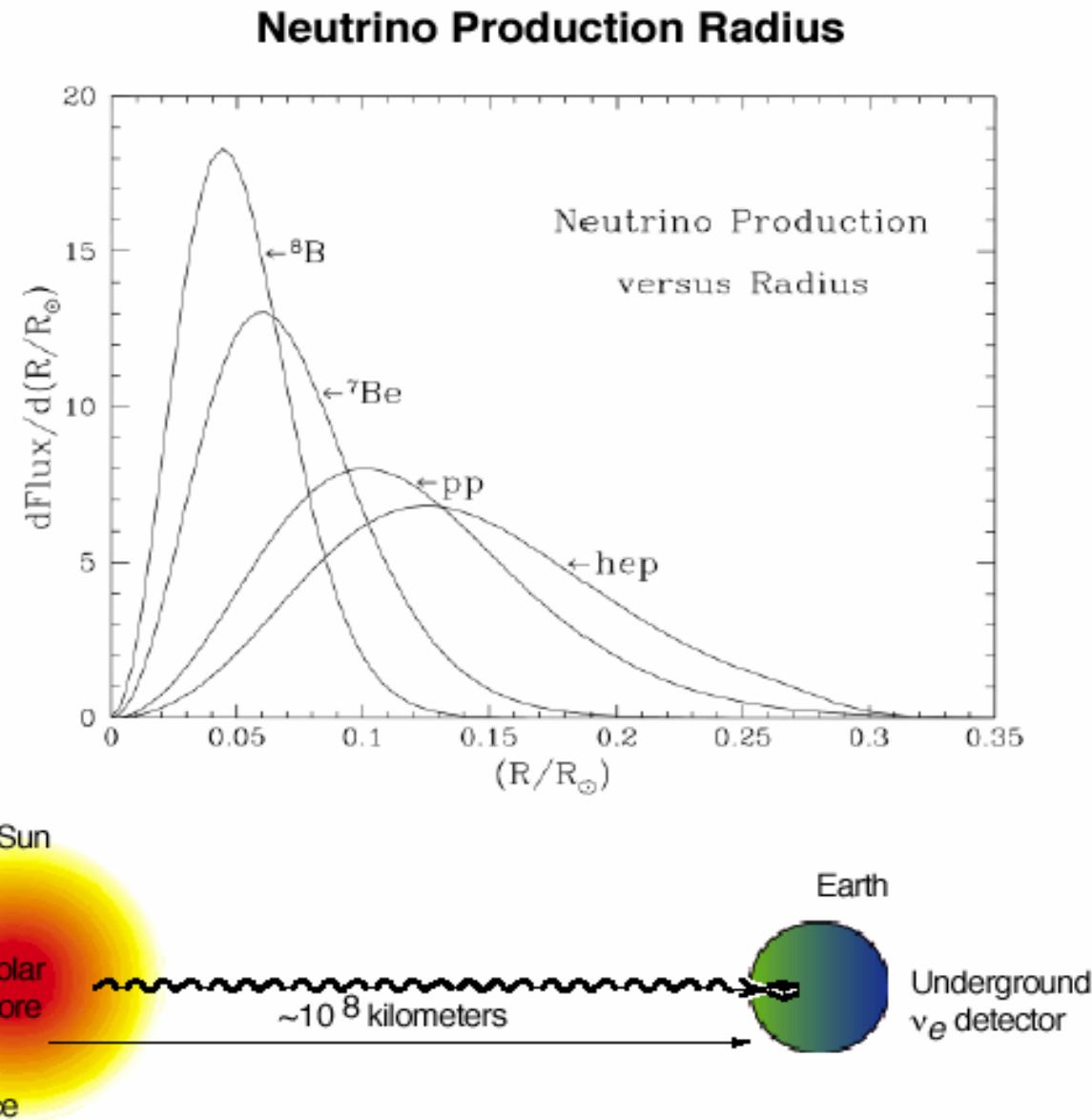
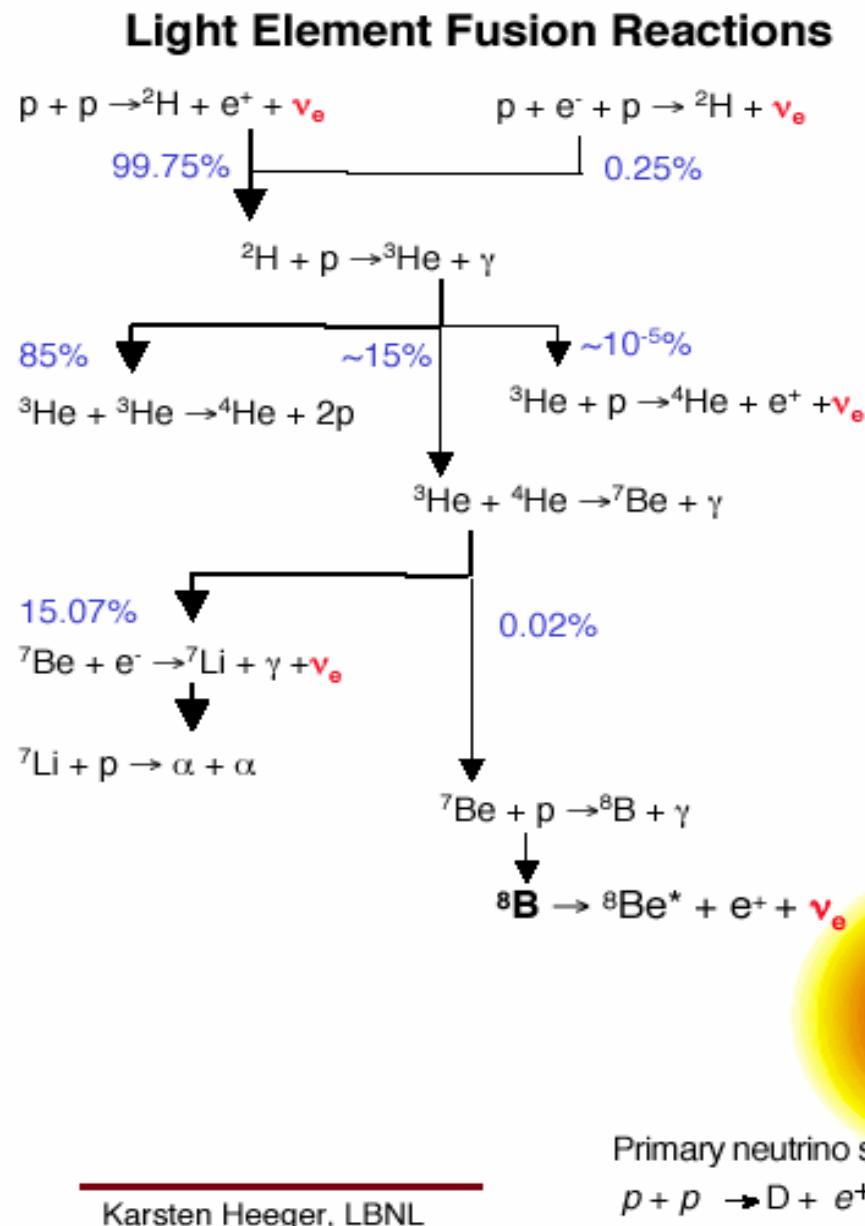


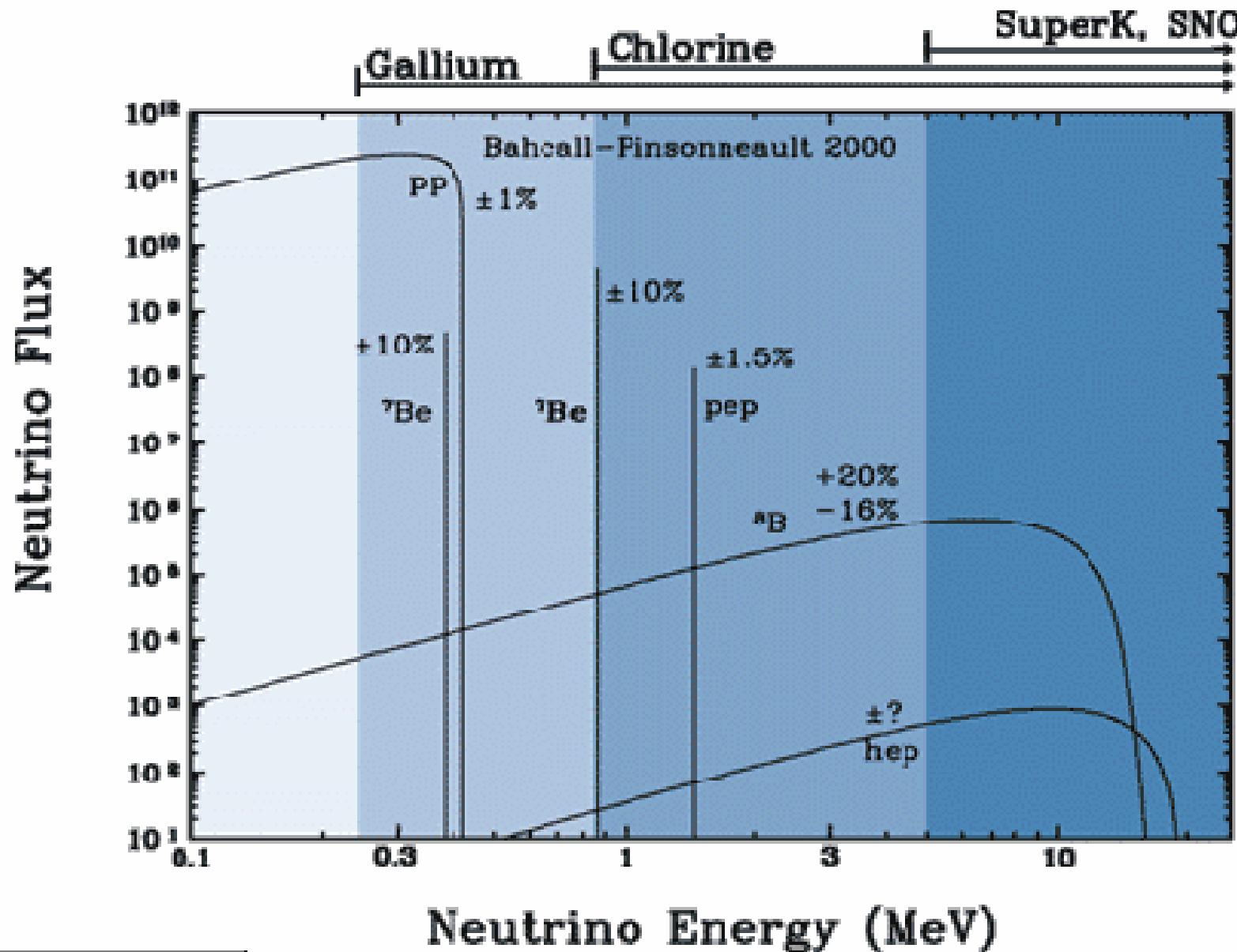
Atmospheric
Neutrinos



Also,
Super Novae (SN1987A),
Neutrinos de beam-dump cosmiques (AGN, GRB...),
Neutrinos cosmogéniques, Neutrinos reliques du Big Bang...

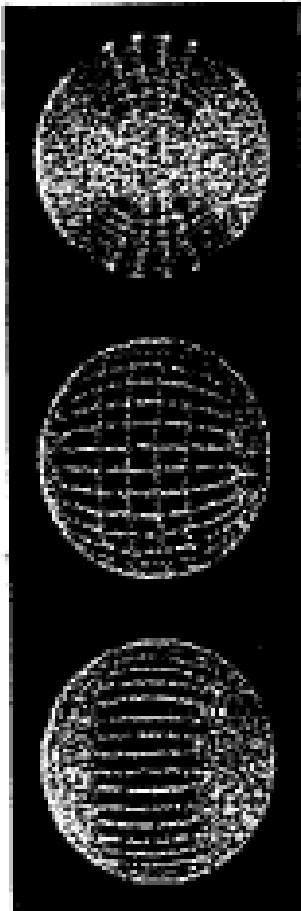
Neutrino Production in the Sun





Solar Standard Model

- *NEW* 39 experts agree on cross secs & systematics!
Rev. Mod. Physics, Oct 1998, astro-ph/9805121



- How well do we know T ?
 $\phi(pp) \sim T^{-1.2}$, $\phi(Be) \sim T^8$, $\phi(B) \sim T^{18}$
→ Helioseismology as a test...
 - The sun is a resonant cavity frequencies depend on $U = P/\rho$.
 - The SSM predicts U
Measured/SSM $U(r)$ agree $< 1\%$!
 - Temperature Uncertainty:
SSM alone: $\Delta T/T = \pm 2.7\%$
SSM + Helio: $\Delta T/T = \pm 1.4\%$

SNe neutrinos

Normal stellar situation: the fusion thermal and radiation pressure compensate the gravitation pressure.

During the collapse, the gravitation binding energy ($\frac{1}{4} 3 \cdot 10^{53}$ erg) cannot escape in an other form than neutrino-antineutrino pairs.

99% of the energy in the form of neutrinos

1% in the form of kinetic energy

0.01% of the energy in optical photons.

The neutron star is opaque to neutrinos. The diffusion and escape time is of the order of 1 second.

$$\langle E n_e \rangle \gg 11 \text{ MeV}$$

$$\langle E \bar{n}_e \rangle \gg 16 \text{ MeV}$$

SN1987a

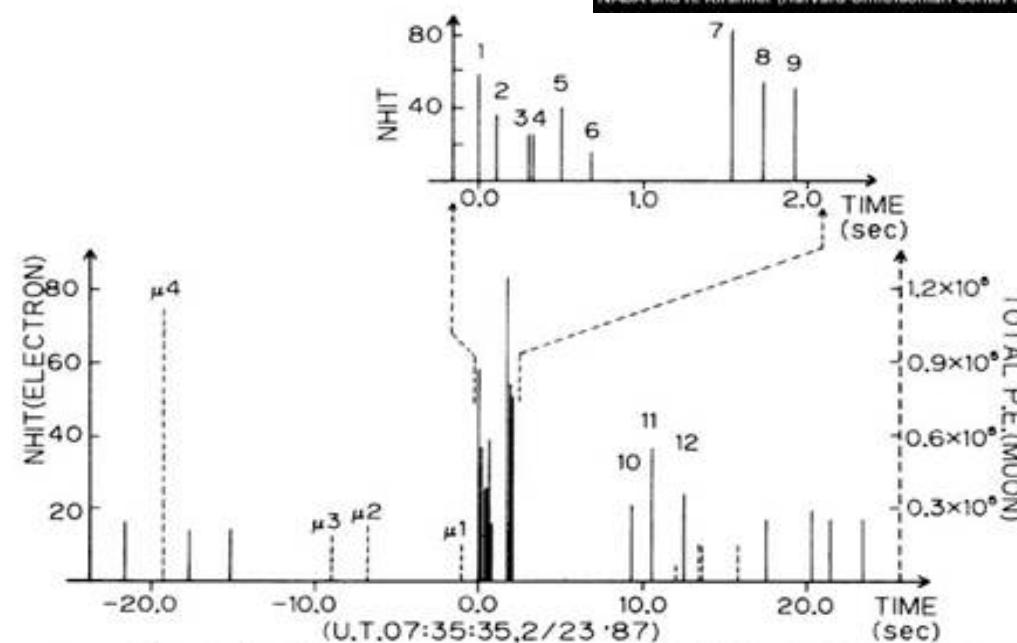
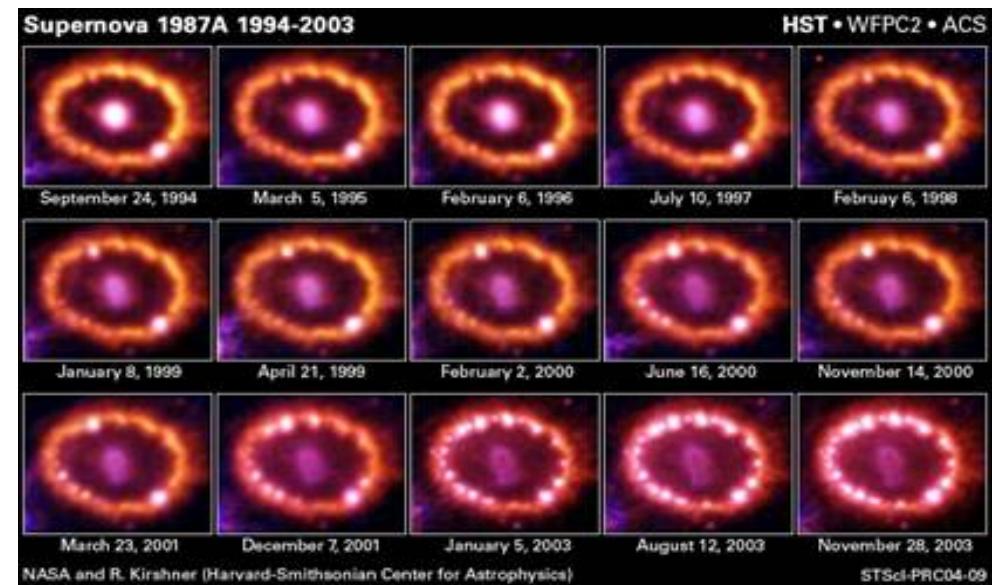
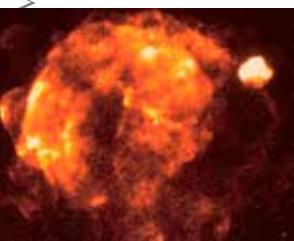


FIG. 2. The time sequence of events in a 45-sec interval centered on 07:35:35 UT, 23 February 1987. The vertical height of each line represents the relative energy of the event. Solid lines represent low-energy electron events in units of the number of hit PMT's, N_{hit} (left-hand scale). Dashed lines represent muon events in units of the number of photoelectrons (right-hand scale). Events μ_1 - μ_4 are muon events which precede the electron burst at time zero. The upper right figure is the 0-2-sec time interval on an expanded scale.

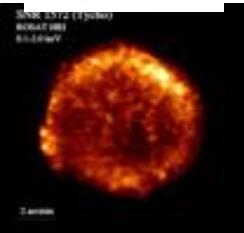
SuperNovae Remnants

M2R PSA -- 2008-2009

Vela

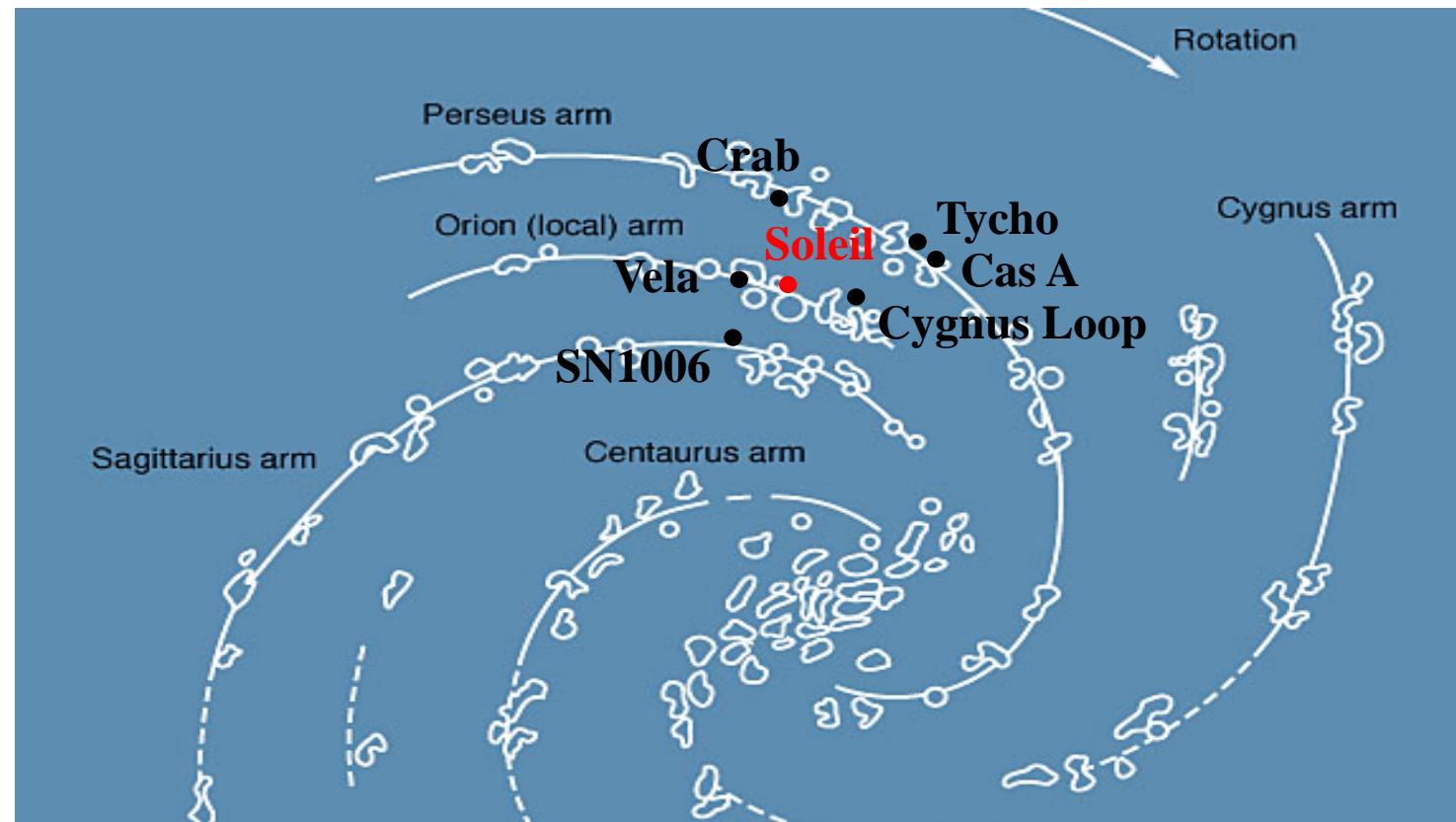


Tycho



François Mon

Cygnus



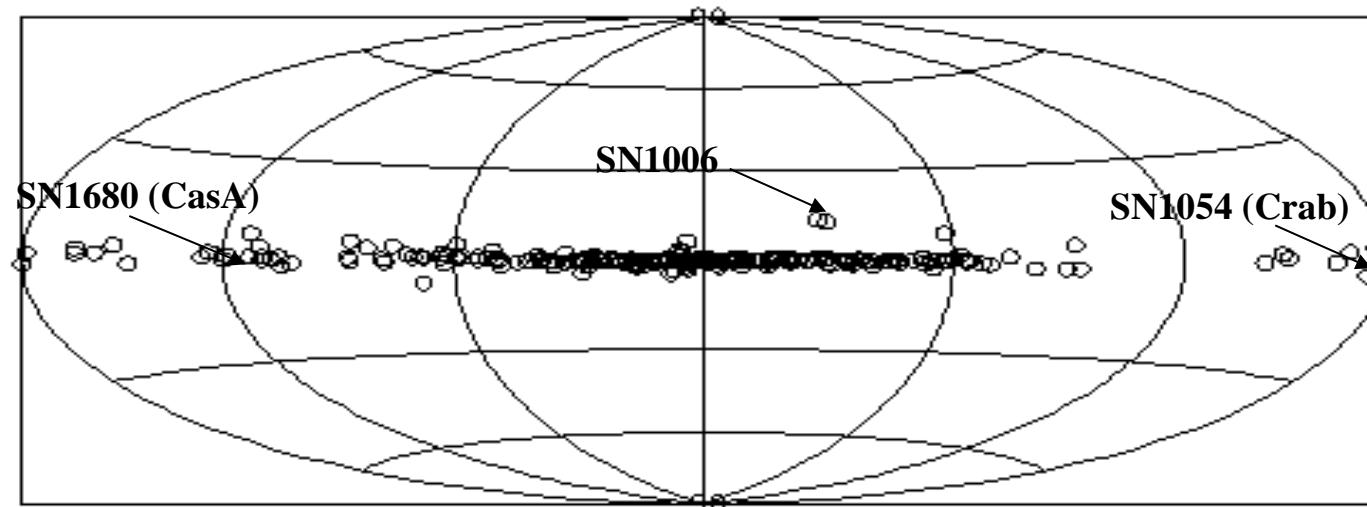
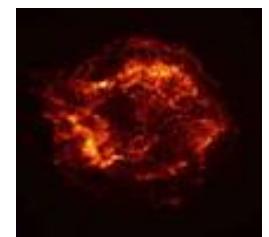
Crab



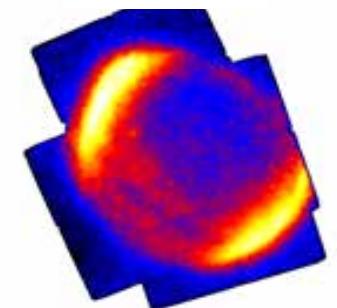
Kepler



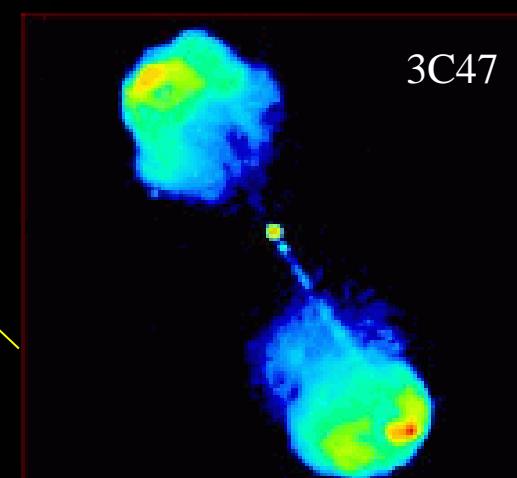
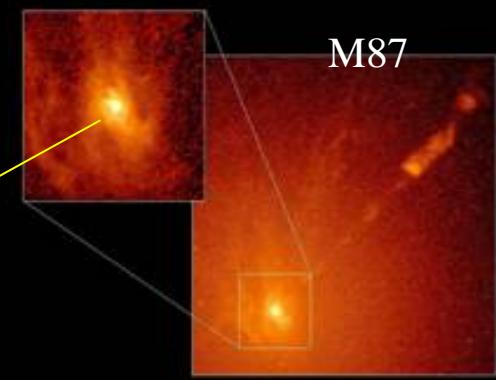
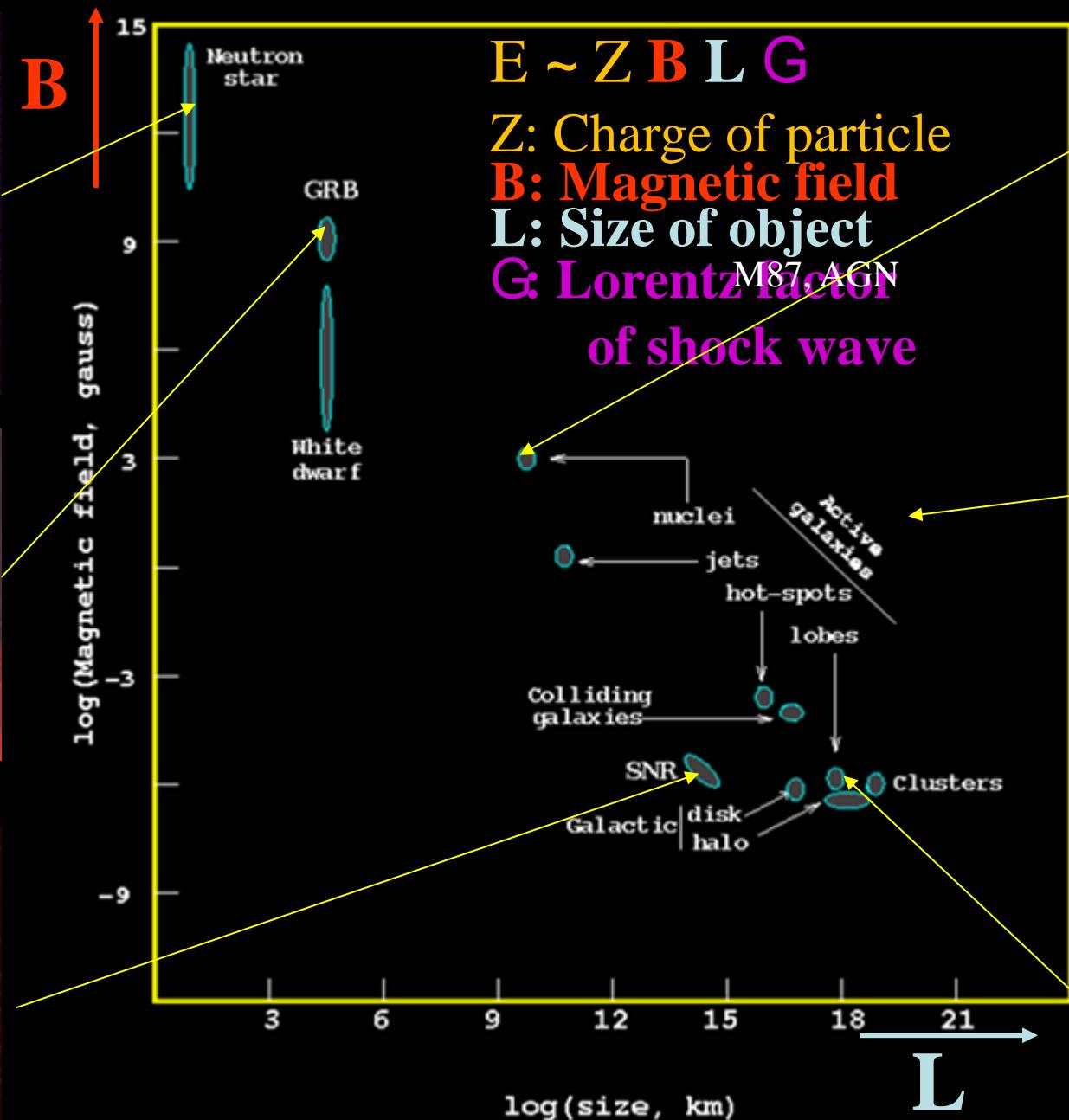
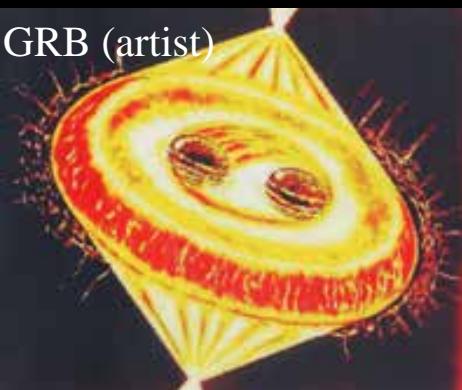
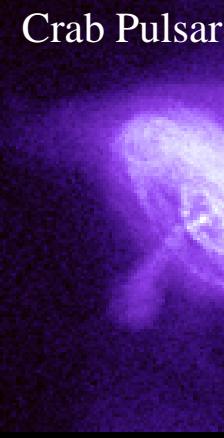
Cas A



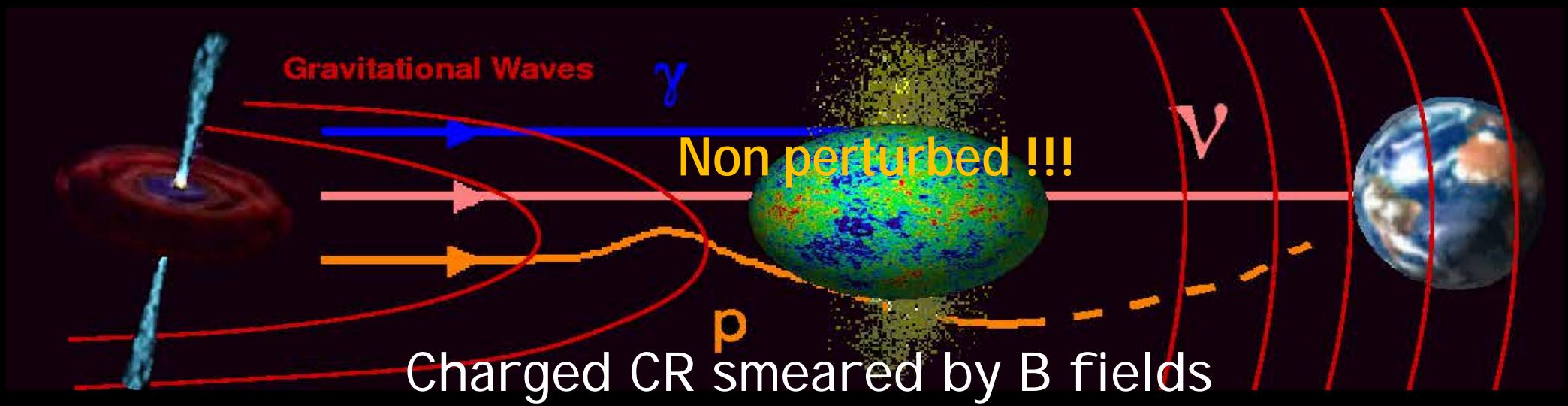
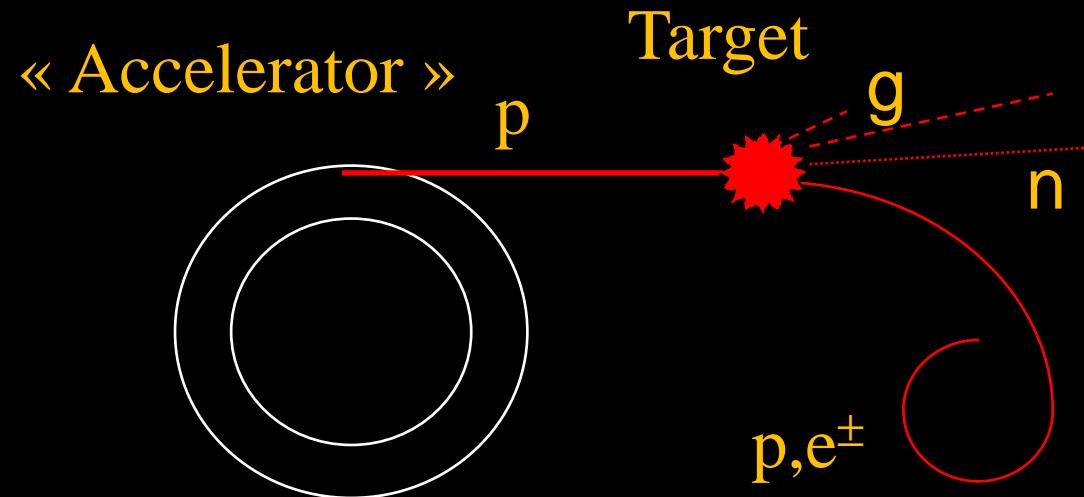
SN1006



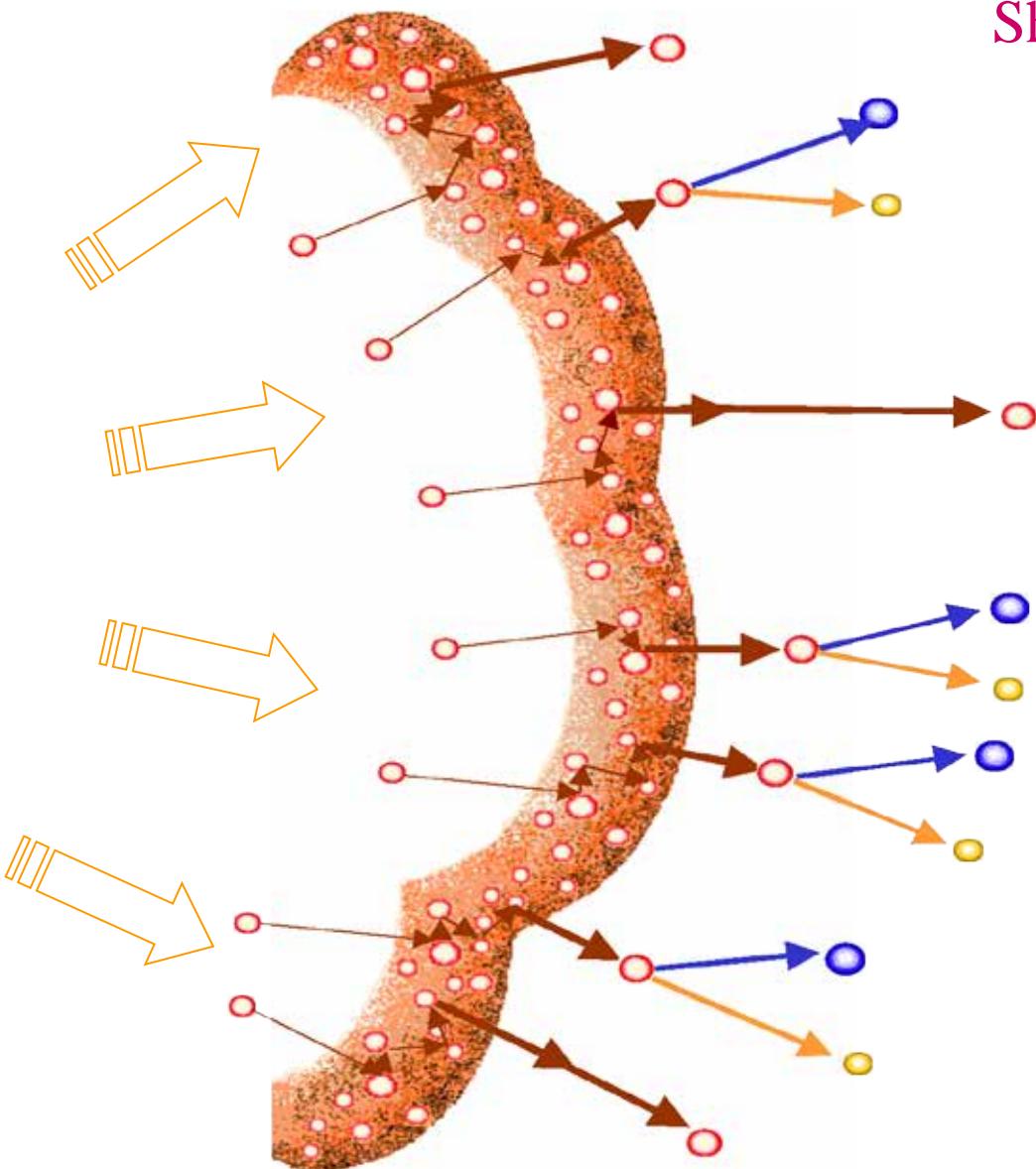
Cosmic Accelerators: (Hillas Plot)



Cosmic Neutrino Beam



Fermi acceleration and UHE neutrino production



Shock wave proton of nuclei acceleration:

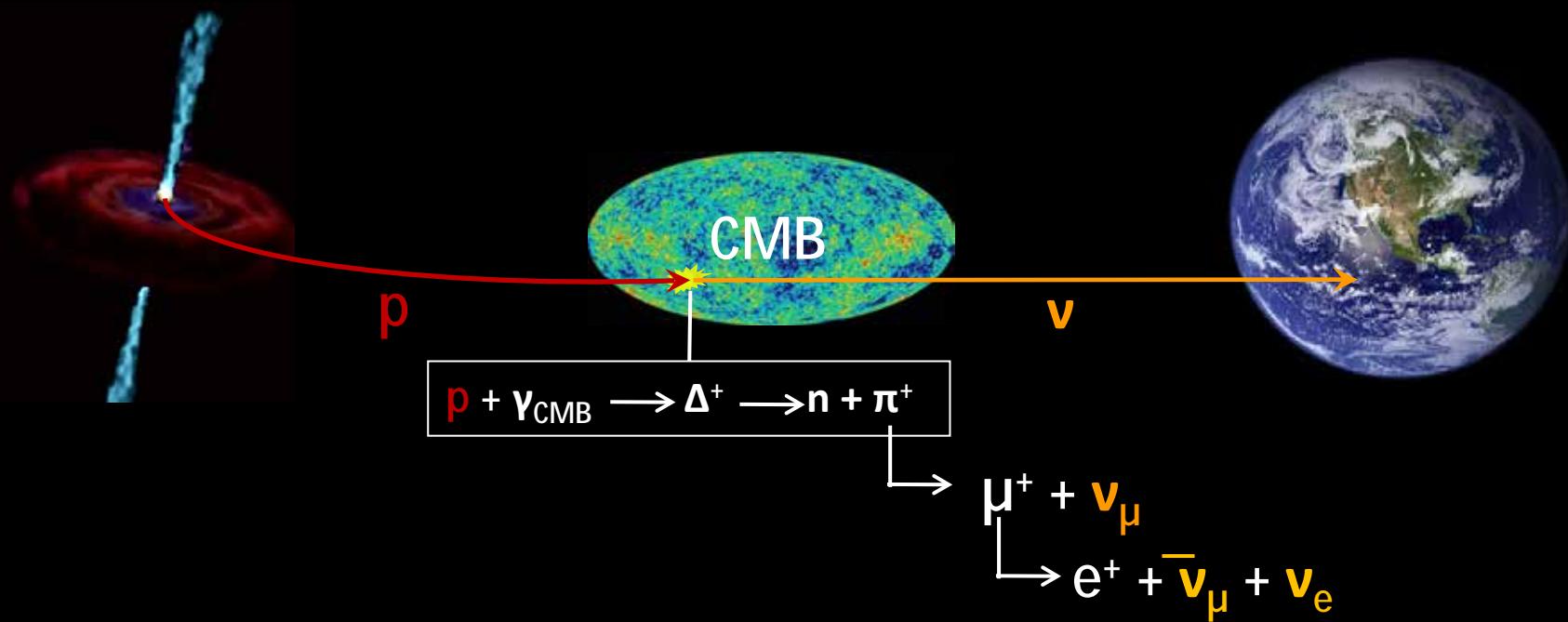
p/N interact with the ISM and produce pionic cascades



COSMOLOGICAL SOURCES

"Cosmogenic" or "GZK" Neutrinos

- Produced by the propagation of UHECR in the IGM.
- Main assumptions:
 - UHECR are protons or nuclei (makes a big difference!)
 - The sources distribution is following a given redshift distribution
 - Spectra are known and flux are generic
 - The target density is well known (CMB photons)
- This flux can be predicted in a relatively robust manner.

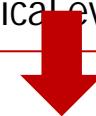


Cosmogenic or "GZK" Neutrinos

Flux attendus

I - Sources

- Composition at the source
- Spectral index + E_{\max}
- Source density + cosmological evol.



Many uncertainties

II - Fonds diffus

- Backgrounds (CMB, CIB...)
- Photon spectrum evol vs z



CMB very well known.
CIB and other bg less known

III - Interaction $\bar{\nu}\gamma$

- Cross Section
- Secondaries energy distribution
- Energy Distribution



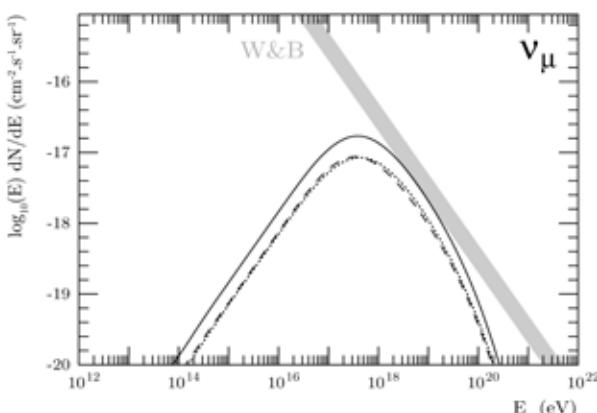
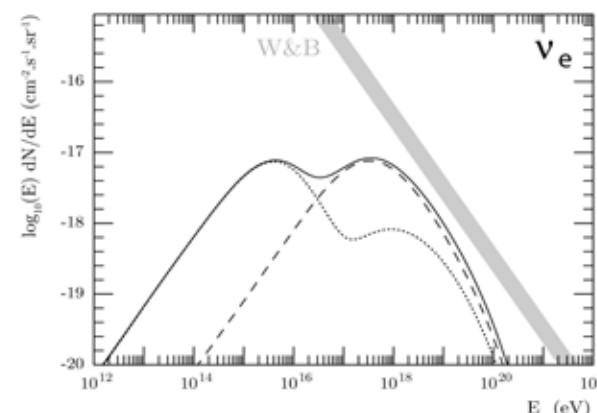
Very well known

I + II + III $\bar{\nu}_e$

$\bar{\nu}_\mu : \bar{\nu}_e : \bar{\nu}_\tau = 2:1:0$ (source)

oscillations

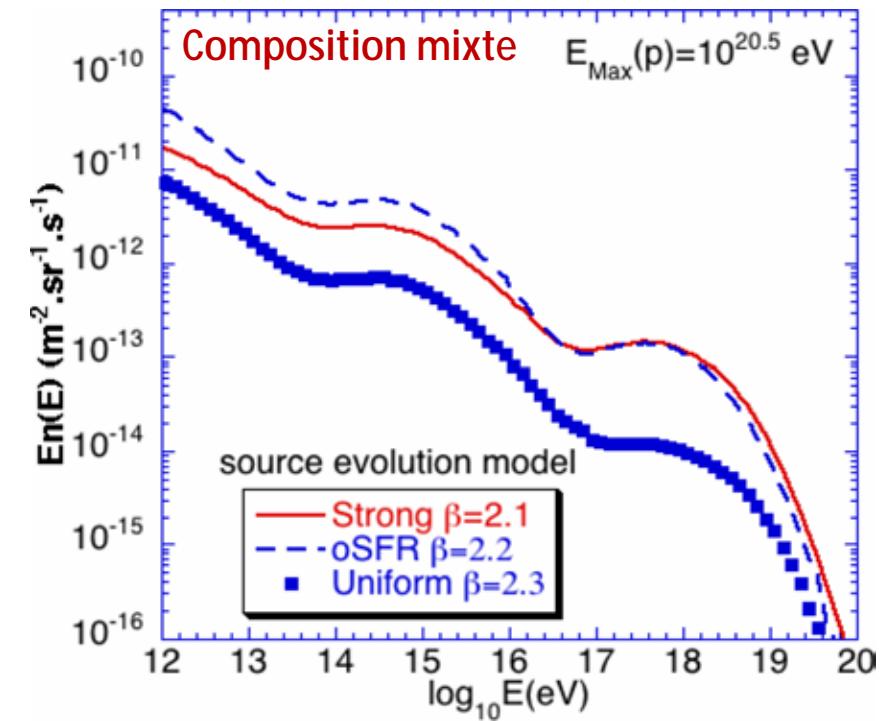
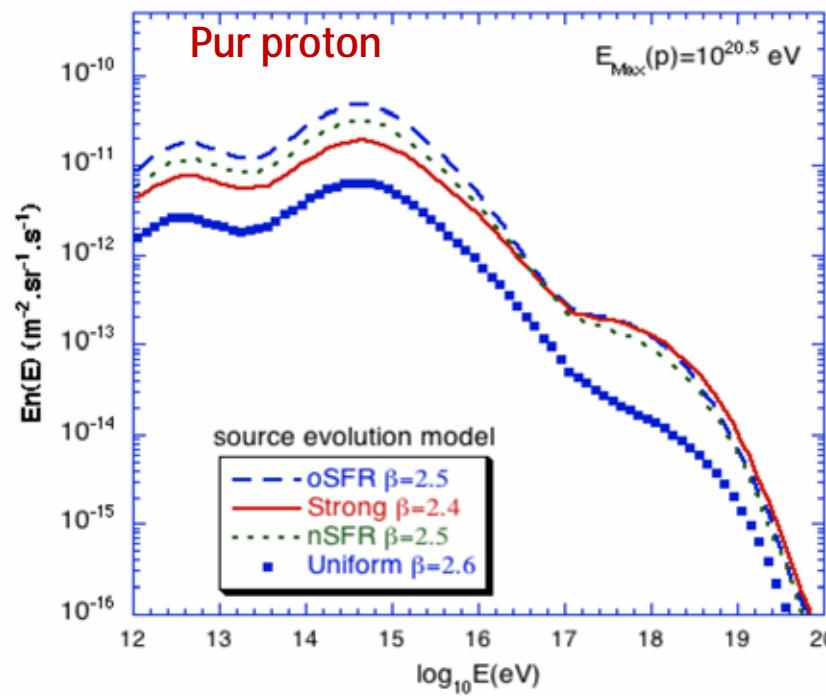
1:1:1 (earth)



Flux of "GZK" neutrinos

Influence of

- Composition of UHECR,
- Sources distribution and evolution.

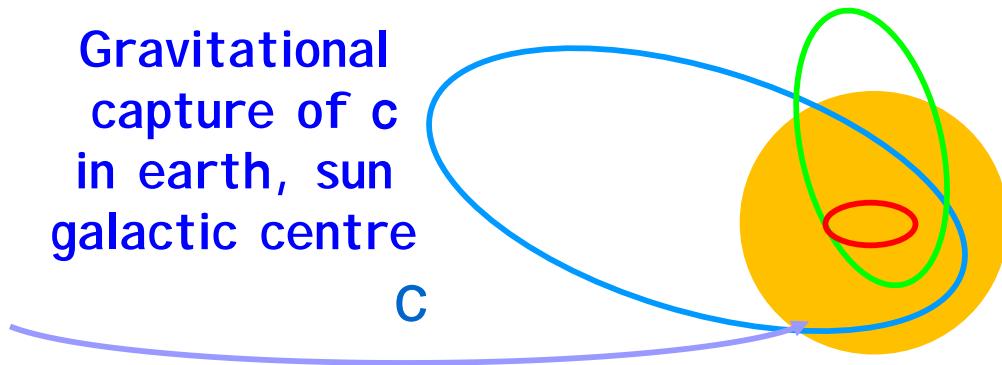


Dark Matter as a source of neutrinos

Annihilation in Halo, Earth, Sun or Galactic Centre

Signature	Experiment
Halo Positron, Antiproton Gamma rays $c\ c\ \textcircled{R}\ z\ g\ gg$	BESS, CAPRICE, AMS, ... HESS, GLAST, MILAGRO, ...
Earth, Sun, GC Neutrino $c\ c\ \textcircled{R}\ WW, ff$ $W, f\ \textcircled{R}\ nX$	SuperK, Baksan, IMB, MACRO AMANDA, ANTARES, Baikal, ...

Gravitational capture of c in earth, sun galactic centre



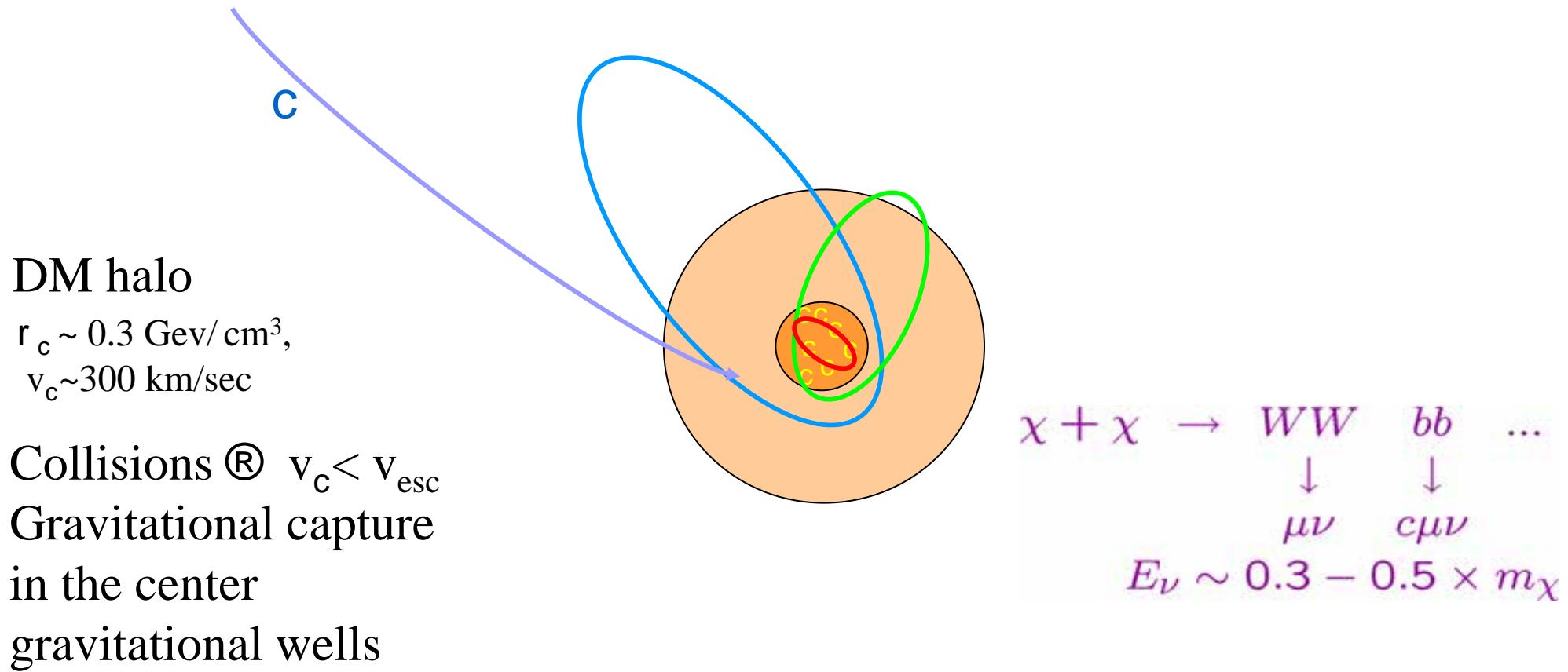
WI MP loses energy by elastic interaction

↳ if $v < v_{escape}$, capture

capture + annihilation balance

↳ constant density in core

Dark Matter as a source of neutrinos



► Indirect search for Neutralinos
 by neutrino telescopes

Propagation

The general problematic

- Thermal speeds ! RCUHE (few 10^{20} eV)



- From top to bottom (decay...)
- From bottom to top (acceleration)



- Energy losses (Synch., IC, p, pairs...)
- Destruction (photo-dissociation...)
- Escape probabilities
- Propagation in ISM and IGM (mag fields: deflection, confinement...)
- Re-acceleration

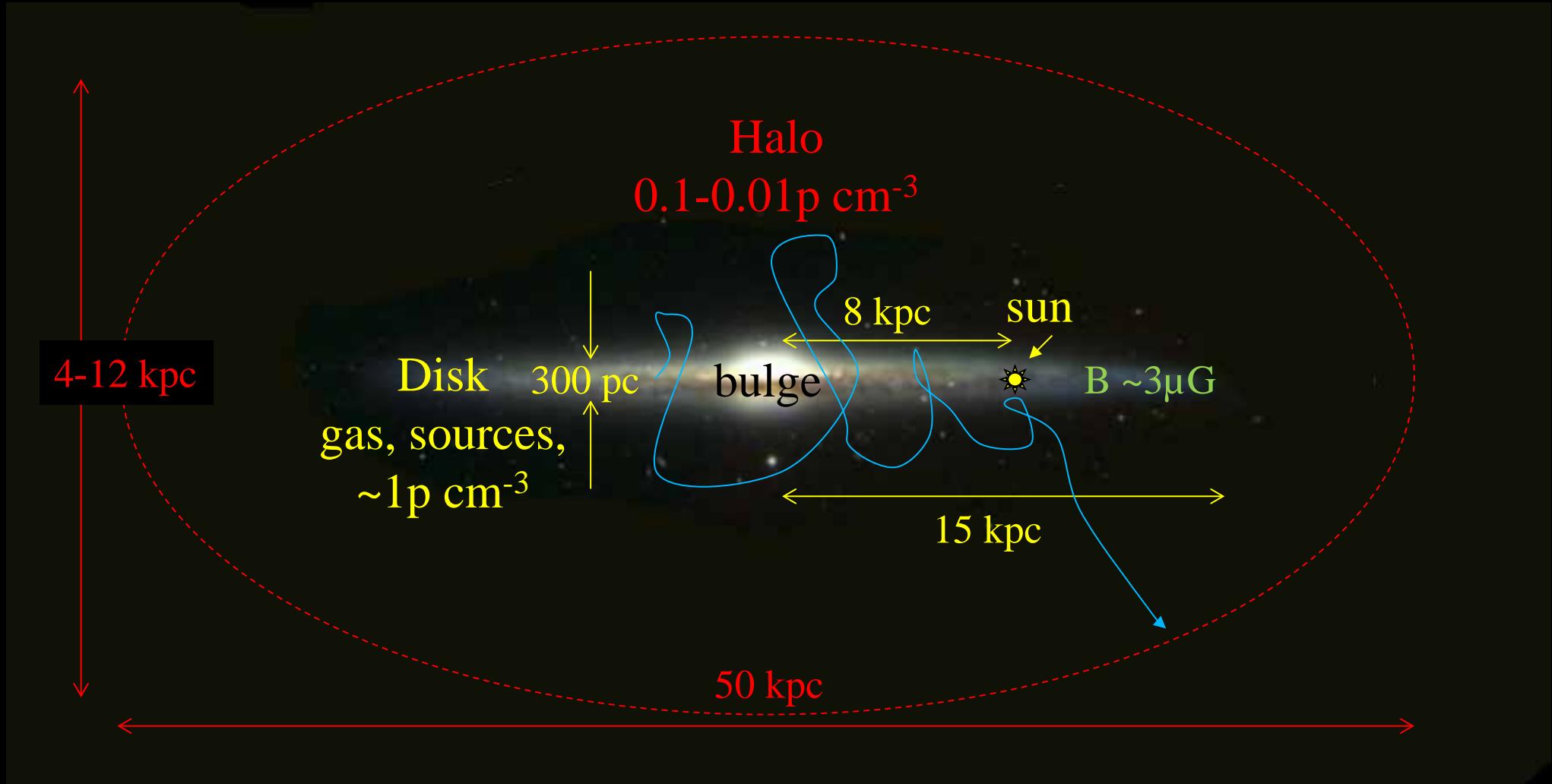


- Balloons, satellites...
- Air showers...
 - Cherenkov telescopes
 - Surface & Fluorescence Detectors



Dimensions of the Milky Way

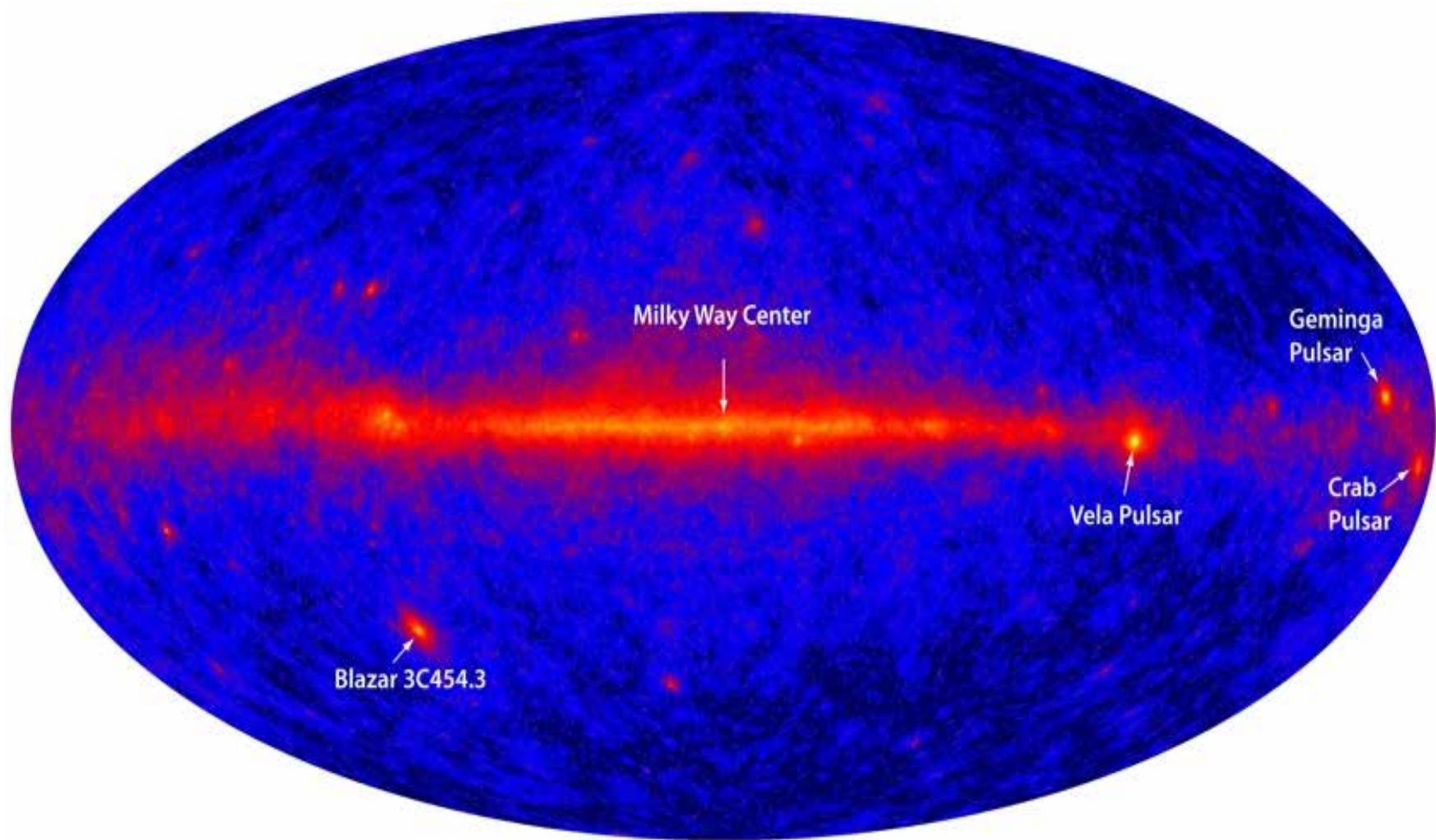
$$1 \text{ pc} \approx 3 \text{ l.y.} \approx 3 \times 10^{16} \text{ m}$$



Propagation des RC dans la Galaxie : Leaky box model

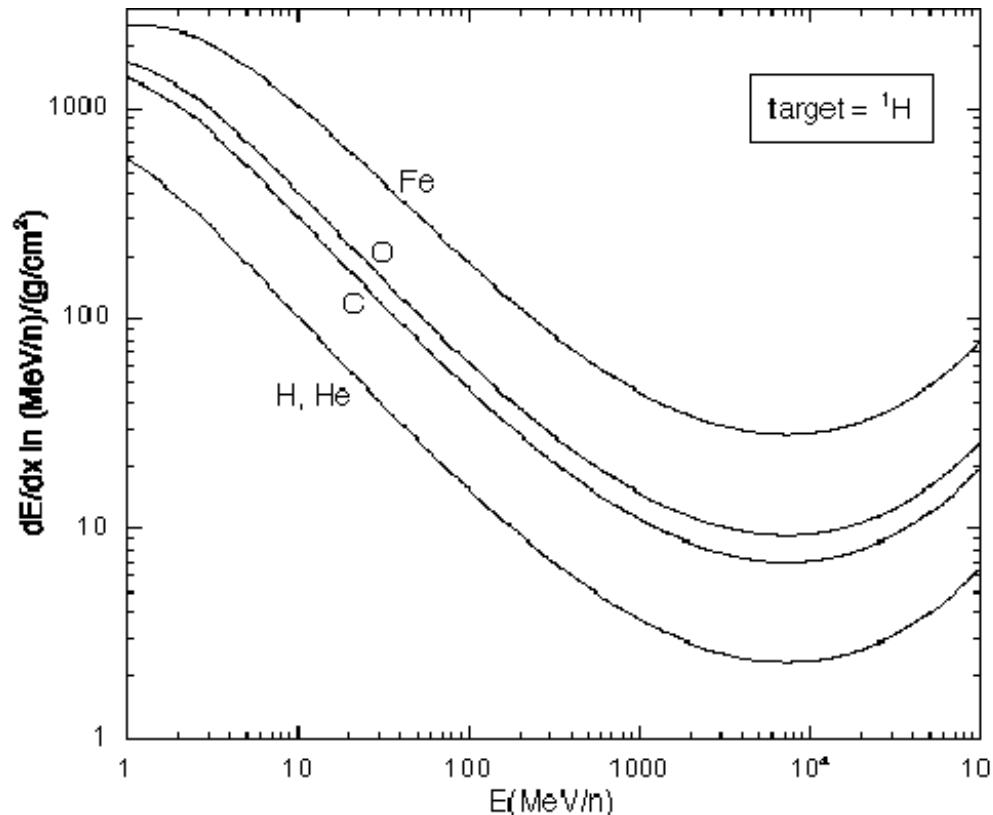
A thick target

- Diffuse gamma-ray emission from galactic CR interaction with matter (mostly molecular H clouds).

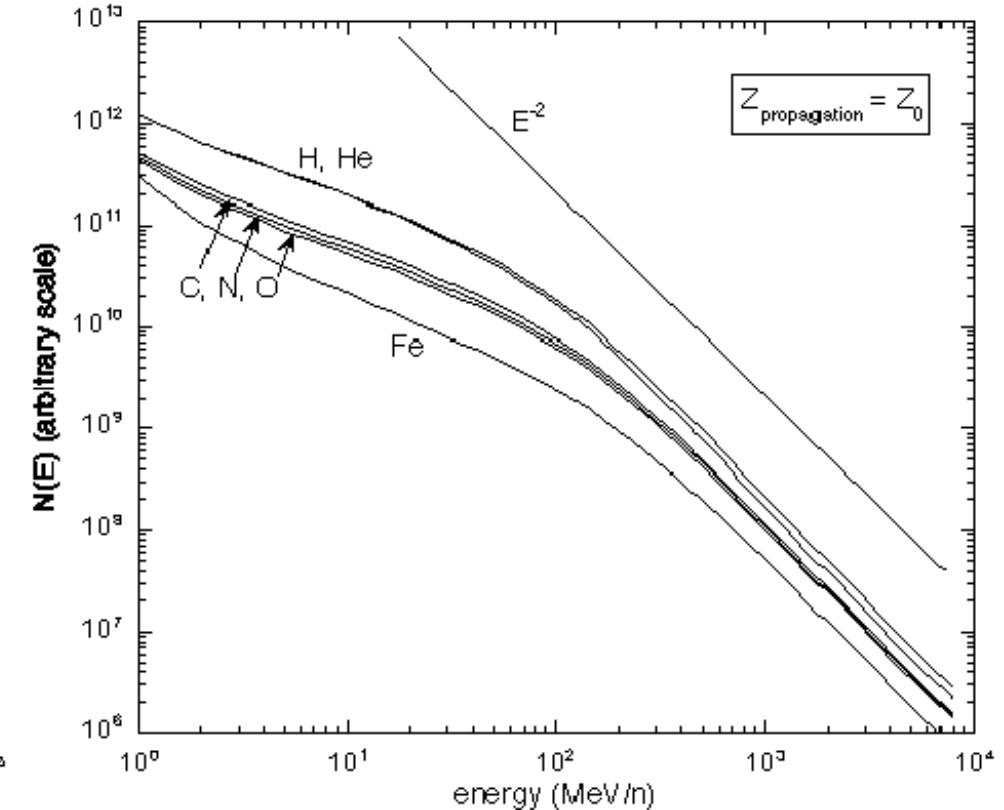


Cosmic rays transport

- Propagation in the interstellar medium



Energy loss: ionization,
Coulombian interactions



Propagated spectra
ionization losses only
(thick target)

Grammage

- Column density or quantity of matter traversed by the CR from its production site to earth (in kg cm^{-2} or g cm^{-2})
- Given the diffusion time (known from cosmic clocks see below) the measurement of grammage allows the understanding of the diffusion extension zone.
- The ratio secondary/primary allows estimating the grammage traversed:

$$\frac{dN_S}{dx} = -\frac{\sigma_P}{m} N_P$$

donc $N_S = N_P \exp -\frac{\sigma_P}{m} x$

et $x = -\frac{m}{\sigma_P} \log(S/P)$

$$B/C \approx 35\% \Rightarrow x = -\frac{m}{\sigma_P} \log(B/C) \approx 60 \text{ kg.m}^{-2}$$

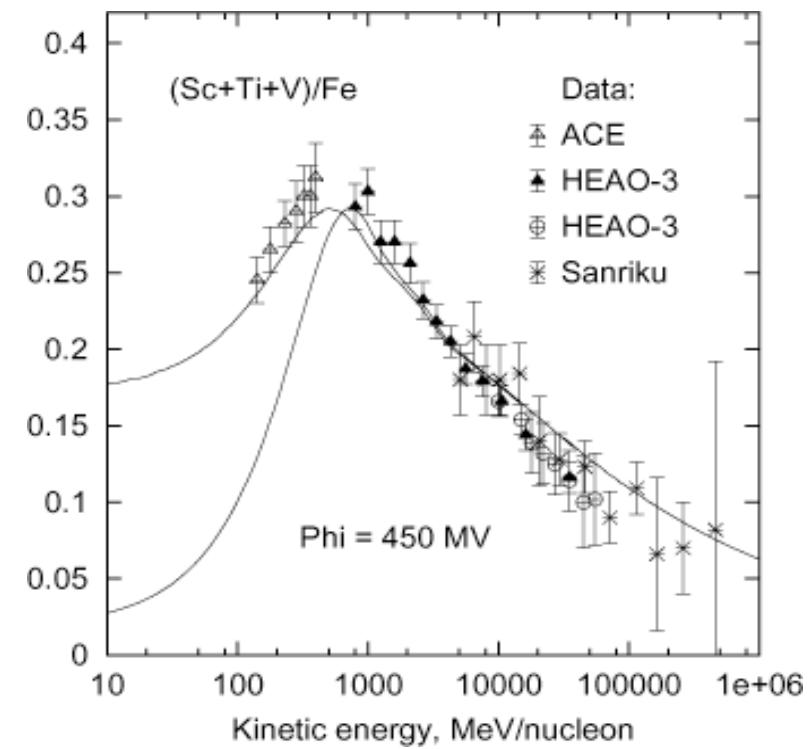
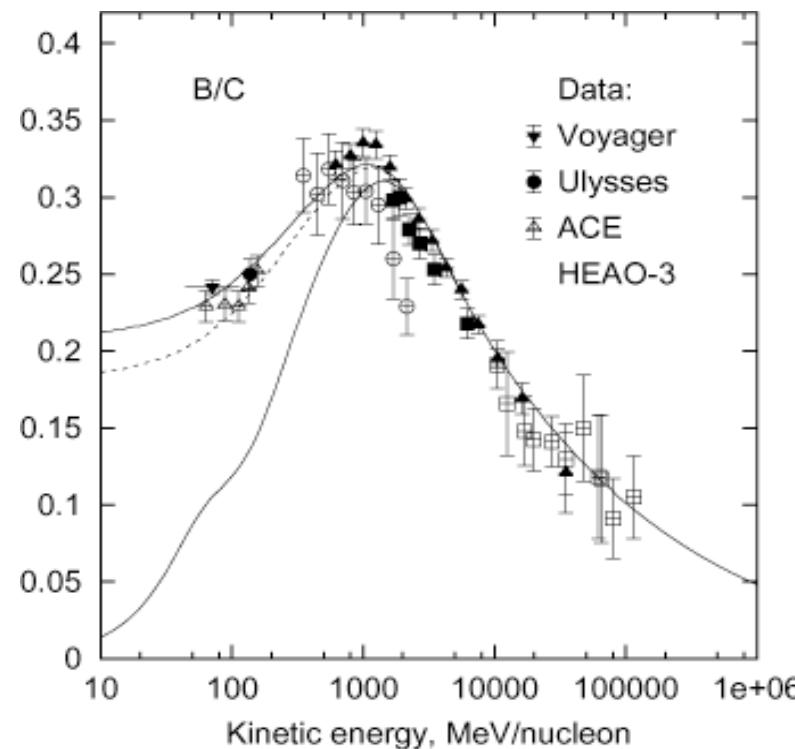
if $Br(C + P \rightarrow B + X) \approx 100\%$

Secondary / Primary ratio

The grammage depend on the parent nucleus:

- $(\text{Li}+\text{Be}+\text{B}) / (\text{C}+\text{N}+\text{O})$ ↗ mean grammage of 50 kg.m^{-2}
- $(\text{Sc}+\text{Ti}+\text{V})/\text{Fe}$ ↗ mean grammage of 20 kg.m^{-2}

and Primary/Secondary ratio (thus grammage) depends on the energy as well:



Secondary / Primary ratio

- A complete CR transport model
 - The secondary to primary ratio can be expressed by:

$$\frac{N_S}{\tau_{esc}} + \frac{N_S}{\tau_{spallation}} = \frac{N_P}{\tau_{P \rightarrow S}}$$
$$\Rightarrow \frac{N_S}{\tau_{esc}} + n\beta c \sigma_S N_S = n\beta c \sigma_{P \rightarrow S} N_P$$



$$\frac{N_S}{N_P} = \frac{\sigma_{P \rightarrow S}}{\sigma_S + 1/\lambda_{esc}}$$

with $\lambda_{esc} = n\beta c \tau_{esc}$

Cosmic clocks

Unstable nuclei with lifetimes comparable to the escape time $T_{1/2} \approx \tau^{esc}$ can be used as cosmic clocks.

The ratio unstable/stable isotope helps disentangling density and escape time.

$$^{10}Be \rightarrow \tau = 2.17 Myr$$

$$^{26}Al \rightarrow \tau = 1.31 Myr$$

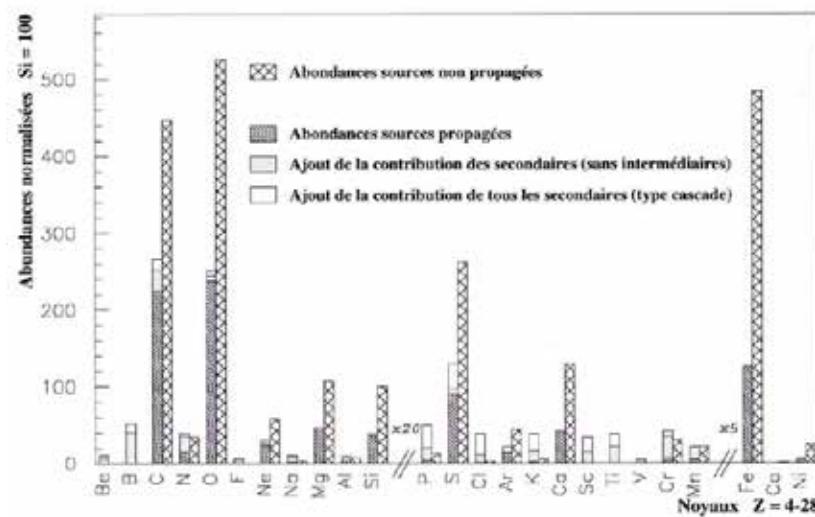
$$^{36}Cr \rightarrow \tau = 0.44 Myr$$

$$\frac{N_j}{\tau_{esc}} + \frac{N_j}{\tau_{rad}} + \frac{N_j}{\tau_{spallation}} = Q_j + \sum_{k>j} \frac{N_j}{\tau_{k \rightarrow j}}$$

Si $\tau_{rad} \ll \tau_e$ et $\tau_{rad} \ll \tau_{spallation}$:

Measure isotopic ratio

$$\frac{N_{rad}}{N_{stable}} = \frac{\tau_{rad}}{\tau_{esc}} + \frac{\tau_{rad}}{\tau_{spallation}}$$



Estimate escape time.

On gets $\tau_e \approx 20 Myr$

Cosmic clocks and halo size

- Radioactive decay:

$$\frac{N_j}{\tau_{esc}} + \frac{N_j}{\tau_{rad}} + \frac{N_j}{\tau_{spallation}} = Q_j + \sum_{k>j} \frac{N_j}{\tau_{k \rightarrow j}}$$

Measure isotopic ratios

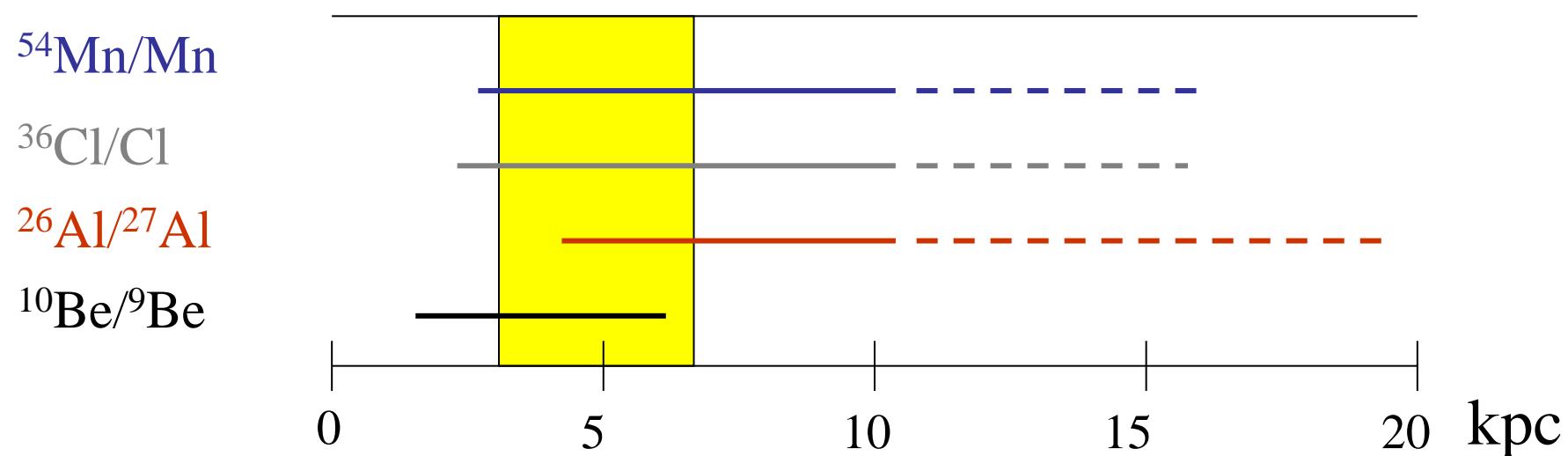
$$\frac{N_{rad}}{N_{stable}} = \frac{\tau_{rad}}{\tau_{esc}} + \frac{\tau_{rad}}{\tau_{spallation}}$$

Estimate escape time

- $^{12}\text{C} + \text{H} \xrightarrow{\text{R}} {}^9\text{Be}$ (stable secondary nucleus)
 $^{12}\text{C} + \text{H} \xrightarrow{\text{R}} {}^{10}\text{Be}$ (unstable secondary nucleus: $\sim 4 \times 10^8$ years)
- The ratio ${}^{10}\text{Be} / {}^9\text{Be}$ depends on secondaries history (and on cross sections).
 - Link between quantity of matter traversed and diffusion time.

Cosmic clocks and halo size

- $^{12}\text{C} + \text{H} \xrightarrow{\gamma} {}^9\text{Be}$ (stable secondary nucleus)
 $^{12}\text{C} + \text{H} \xrightarrow{\gamma} {}^{10}\text{Be}$ (unstable secondary nucleus: $\sim 4 \times 10^8$ years)
- The ratio ${}^{10}\text{Be} / {}^9\text{Be}$ depends on secondaries **history** (and on cross sections).
 - Link between quantity of matter traversed and diffusion time.
- Diffusion parameters adjustments (excursion in the less dense galactic halo)
 - ▷ determination of the CR confinement zone



Confinement and escape

- The average measured grammage is $x = 50 \text{ kg.m}^{-2}$
- Associated lengths:

$$\lambda_{esc} = x/\rho \approx 750 \text{ kpc},$$

with $\rho = 1.4 n_H m_p \approx 2.2 \times 10^{-21} \text{ kg.m}^{-3}$

- $\lambda_{esc} \gg R = 20 \text{ kpc} \Rightarrow \text{CR are confined}$
- $\lambda_{esc} \ll \lambda_{pp} = (n_H \sigma)^{-1} \approx 6 \text{ Mpc} \Rightarrow \text{CR can escape}$
- Long lived radioactive secondaries (cosmic clocks) indicate $\tau_{esc} \approx 20 \text{ Myr}$
- Average density scanned by CR:

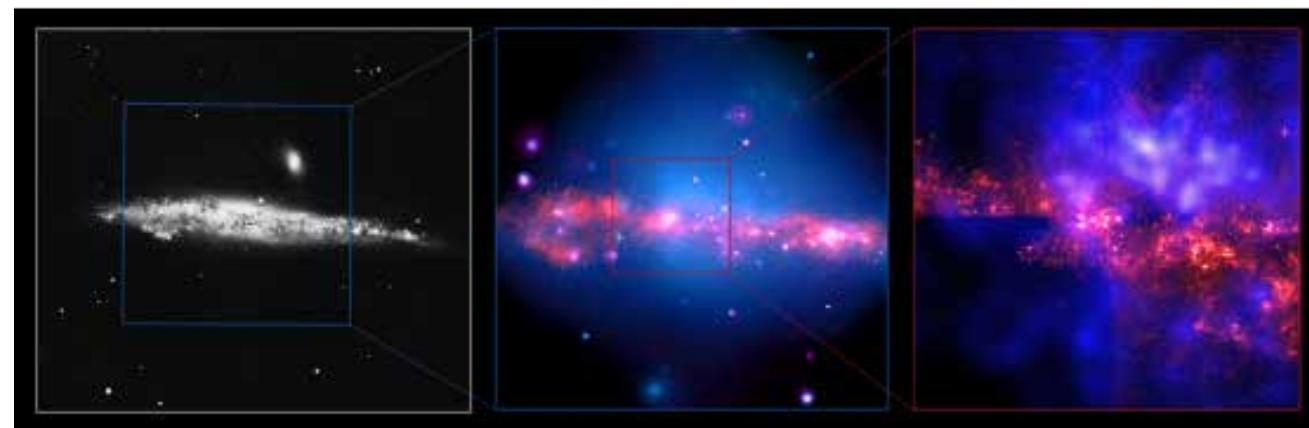
$$n_H = \lambda_{esc}/c \tau_{esc} m_p \approx 0.3 \text{ cm}^{-3} < n_{disk} = 1 \text{ cm}^{-3}$$

$\Rightarrow \text{CR diffuse in a thinner region: the Halo}$

Disk & Halo

- CR can wander out of the disk in a magnetized halo of hot ionized matter

$$T = 10^6 \text{ K} \quad \text{et} \quad n = 10^{-3} \text{ cm}^{-3}$$



- NGC 4631 galaxy and its halo of hot ionized matter emitting X-rays as seen by Chandra

The « leaky box » model

- Diffusive approximation...

$$\frac{\partial N_i}{\partial t} + \frac{\partial}{\partial E} [b(\vec{r}, E)N_i(\vec{r}, E, t)] = Q_i(\vec{r}, E) - \frac{N_i}{\tau_{tot}(\vec{r}, E)} + D(\vec{r}, E)\nabla^2 N_i$$

The diagram shows the leaky box model equation with several annotations:

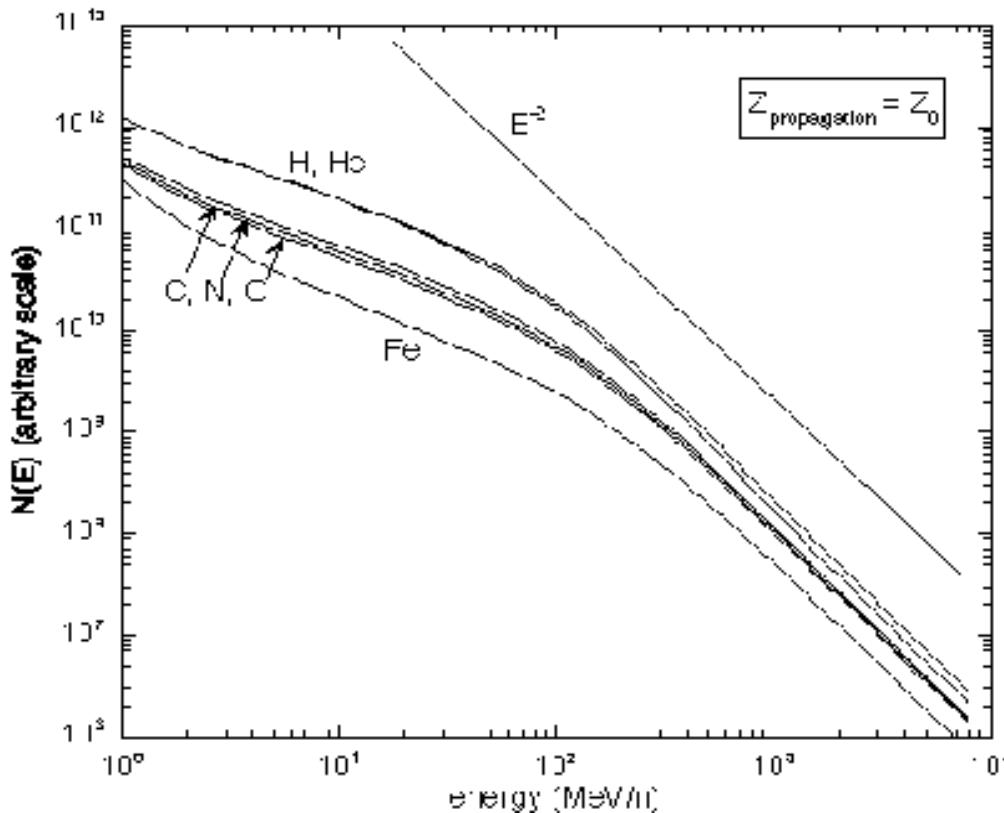
- A purple arrow points to the term $\frac{\partial}{\partial E} [b(\vec{r}, E)N_i(\vec{r}, E, t)]$ with the label "flux in energy space (losses + acceleration)".
- A purple arrow points to the term $Q_i(\vec{r}, E)$ with the label "injection, 'in flight' production".
- A purple arrow points to the term $\frac{N_i}{\tau_{tot}(\vec{r}, E)}$ with the label "destruction, decay, escape".
- A purple arrow points to the term $D(\vec{r}, E)\nabla^2 N_i$ with the label "diffusion".

- Re-acceleration...

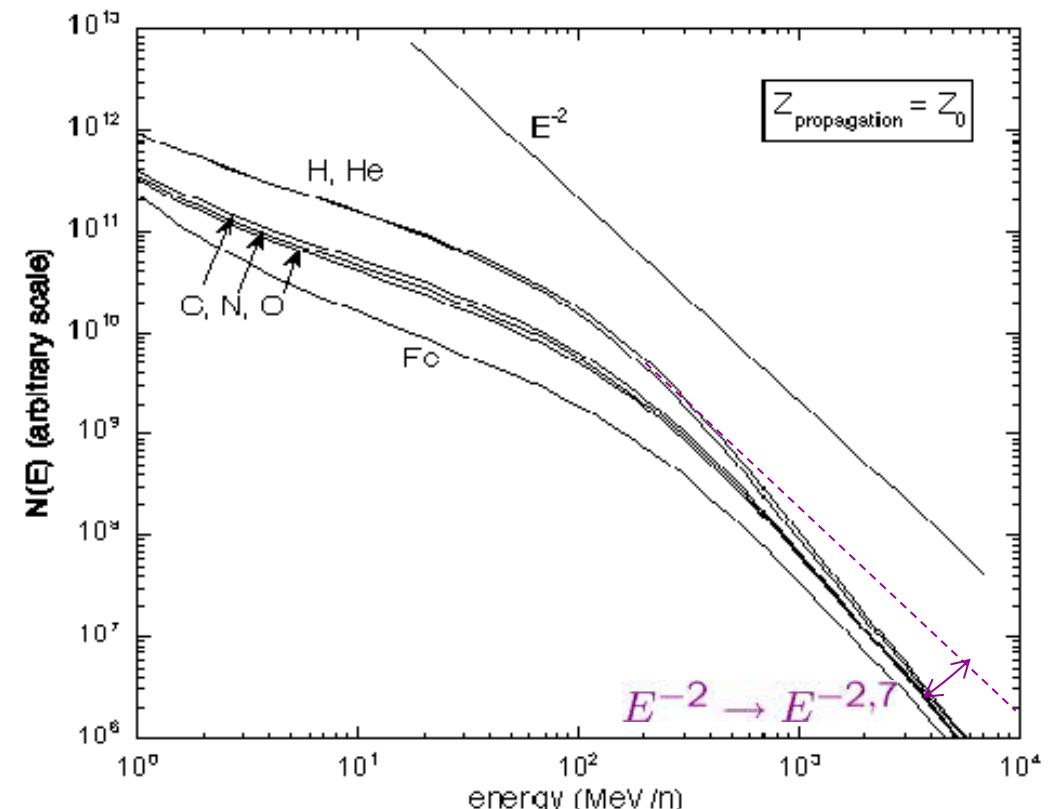
$$+ \frac{1}{2} \frac{\partial^2}{\partial E^2} [c(\vec{r}, E)N_i(\vec{r}, E, t)]$$

Slope of the propagated spectrum

- Escape out of the confinement zone
 - Confinement (escape probability) decrease with E



Without escape
(thick target)



$t_{\text{conf}} \mu E^{-0.7}$
④ $E^{-2.7}$ spectrum

CR confinement

- Escape depends on E
 - Diffusion on magnetic inhomogeneities
 - When $E \ll r_g$ thus interaction with inhomogeneities with larger wavelengths.

- $D(E)$ is an increasing function

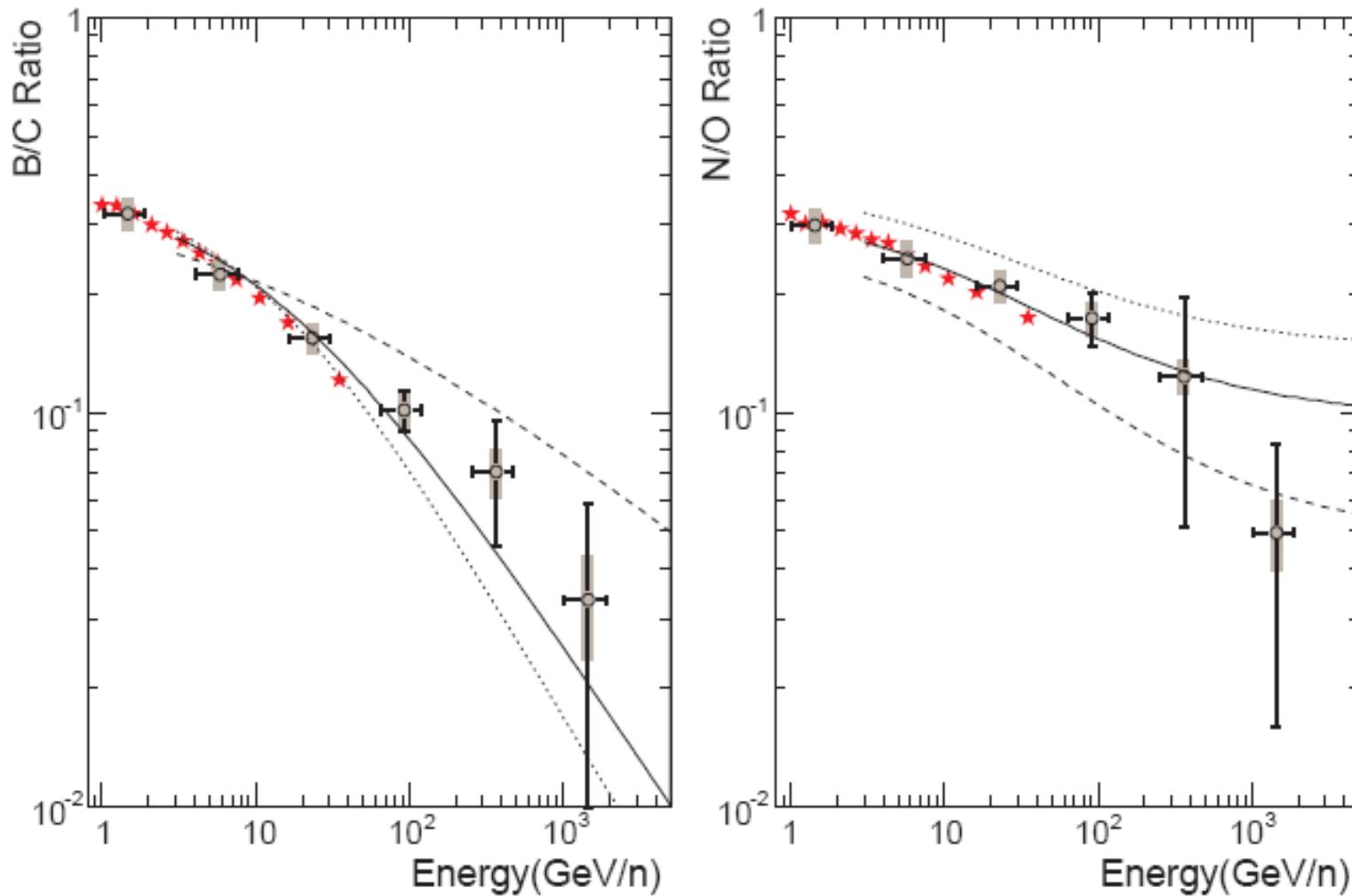
$$D = \beta D_0 \left(\frac{\rho}{\rho_0} \right)^x \text{ where } \rho \text{ is the particle rigidity}$$
$$\rightarrow \tau_{conf}(E) \propto E^{-x}$$

- Kolmogorov spectrum $\rightarrow \tau_{conf}(E) \propto E^{-1/3}$
 - $x - 2 = 1/3 < 0,7$... clearly not enough but...
 - ISM perturbations ? Diffusion-convection, MHD ?
- Determination of $\tau_{conf}(E)$ a posteriori :

$$2,7 - 2 = 0,7 !!!$$

The CREAM results

Secondary/Primary ratio



Compatible with propagation models with
escape terms in $E^{0.6}$

CR anisotropy

- Very weak observed anisotropy (first angular harmonic)
 - $\delta \sim (0,5 - 1)10^{-3}$ de 10^{12} à 10^{14} eV
 - $\delta \sim 1$ à 2% à 10^{17} eV
 - Upper limit : 5% à 10^{18} eV, 10% à $10^{18,5}$ eV, 30% à 10^{19} eV
 - Constraints on $D(E)$
 - In the diffusive approximation, one can compute d taking into account the position of the solar system wrt the Galaxy.
- Observed anisotropy incompatible with sources distributed like SNR and a diffusion term $D(E)\mu E^{0.6}$ below 10 GeV...
... but compatible with $D(E) \propto E^{1/3}$!

Sources Distribution

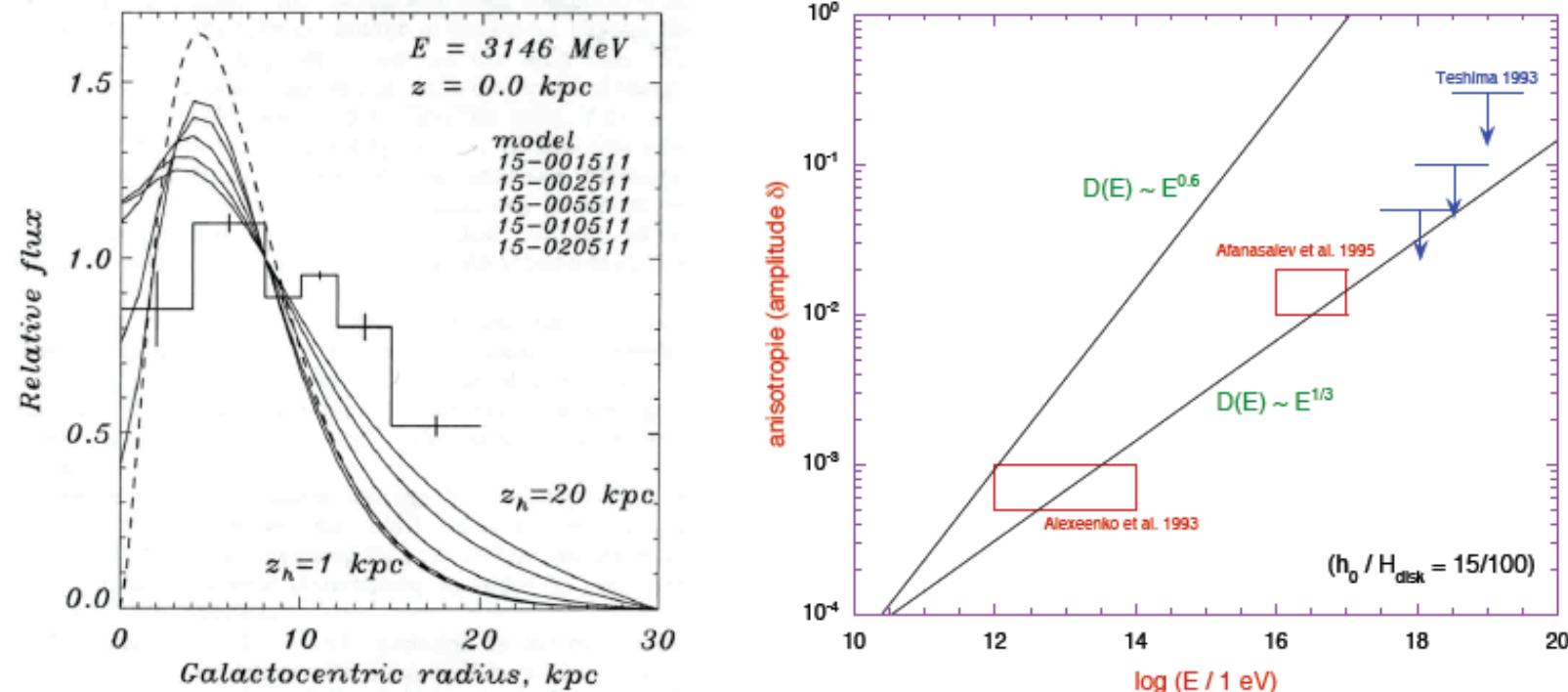


FIG. 5.1 – À gauche : distribution galactocentrique des flux de rayons cosmiques prédicts par les modèles de propagations pour différentes tailles du halo de confinement (courbes en trait continu), comparée à la ditribution déduite des observations d'EGRET. La courbe en pointillé montre la distribution source supposée, conforme à la distribution des SNRs, piquée sur l'anneau moléculaire à 4 kpc. À droite : évolution de l'anisotropie du rayonnement cosmique prédicté pour un coefficient de diffusion en $E^{0.6}$ (hypothèse SNR) ou en $E^{1/3}$ (comme attendu théoriquement et suggéré par les rapports d'abondance du rayonnement cosmique), comparée aux contraintes observationnelles. La seconde solution est seule en accord avec les données.

④ Source distribution is more uniform...

Full transport equation

$$\frac{\partial \psi(\vec{r}, p, t)}{\partial t} = q(\vec{r}, p) \quad \text{sources (SNR, nuclear reactions...)}$$

diffusion + $\vec{\nabla} \cdot [D_{\chi\chi} \vec{\nabla} \psi - \vec{V} \psi]$ convection

diffusive reacceleration + $\frac{\partial}{\partial p} \left[p^2 D_{pp} \frac{\partial}{\partial p} \left(\frac{\psi}{p^2} \right) \right]$

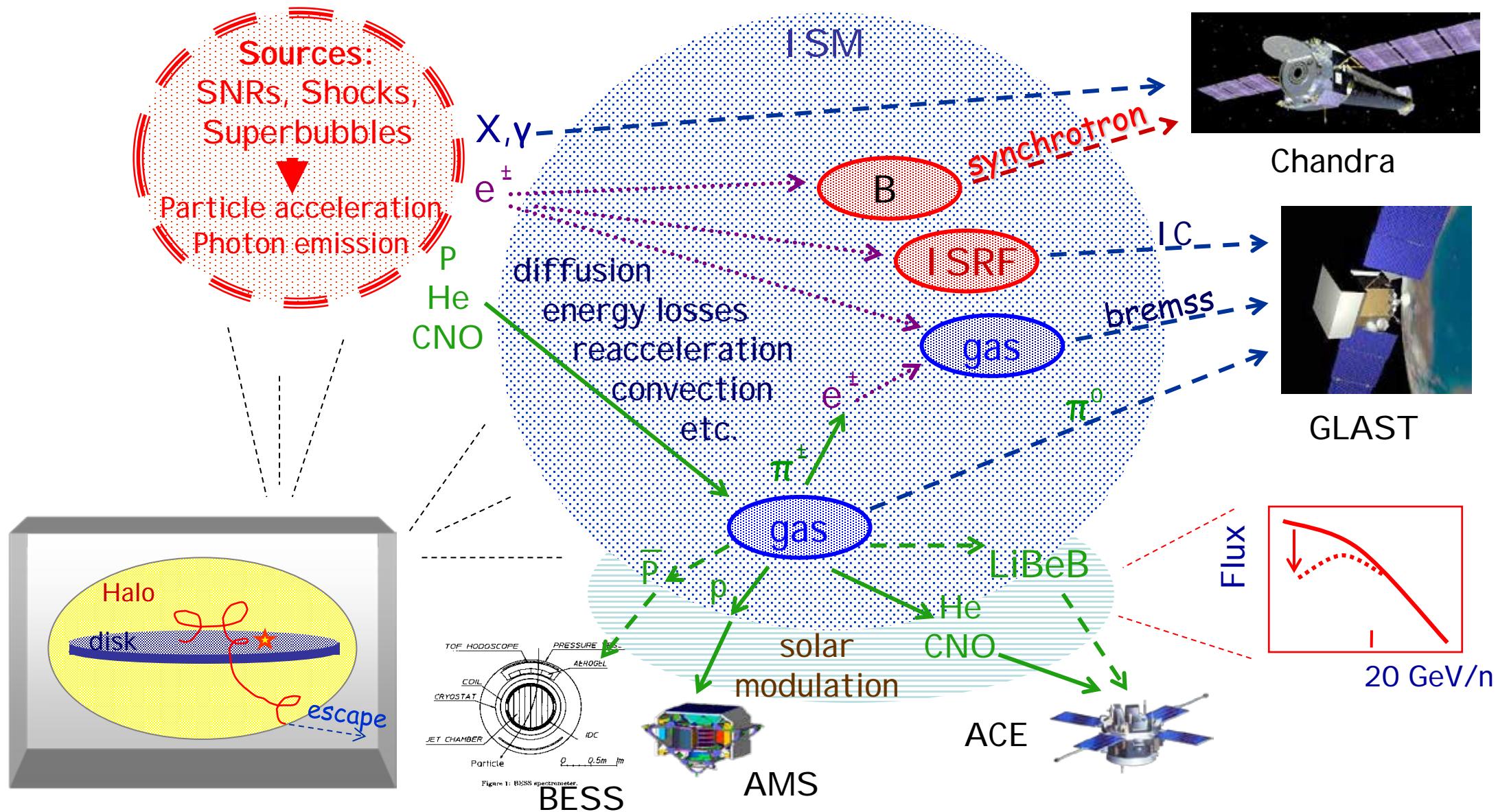
E-loss - $\frac{\partial}{\partial p} \left[\frac{dp}{dt} \psi - \frac{1}{3} p \vec{\nabla} \cdot \vec{V} \psi \right]$ convection

fragmentation - $\frac{\psi}{\tau_f} - \frac{\psi}{\tau_d}$ Radioactive decay

$\psi(\vec{r}, p, t)$ – momentum density

Propagation in the ISM et observational contrains

2010-2011



Summary for galactic CR

Everything works fairly well...

- Propagation in the ISM:
 - Complete theory with energy losses, diffusion, in flight nuclear reactions, CR escape, reacceleration, ... impressive results.
(see for example GLAPROP model, A. Strong et I Moskalenko)
- Secondaries / Primaries
- Cosmic clock
- Anisotropies
- Theoretical expectations (~ Kolmogorov spectrum : $D(E) \propto E^{0.36}$)

...except naïve acceleration models!

- Observation + models require source spectra / $E^{-2.35}$ (high energy spectral shape and $I_{\gamma}^{\text{aires}}/I_{\gamma}^{\text{radio}}$ ratio "best fit")
- "Softer" (steeper) than standard spectra for strong shocks f (E^{-2})

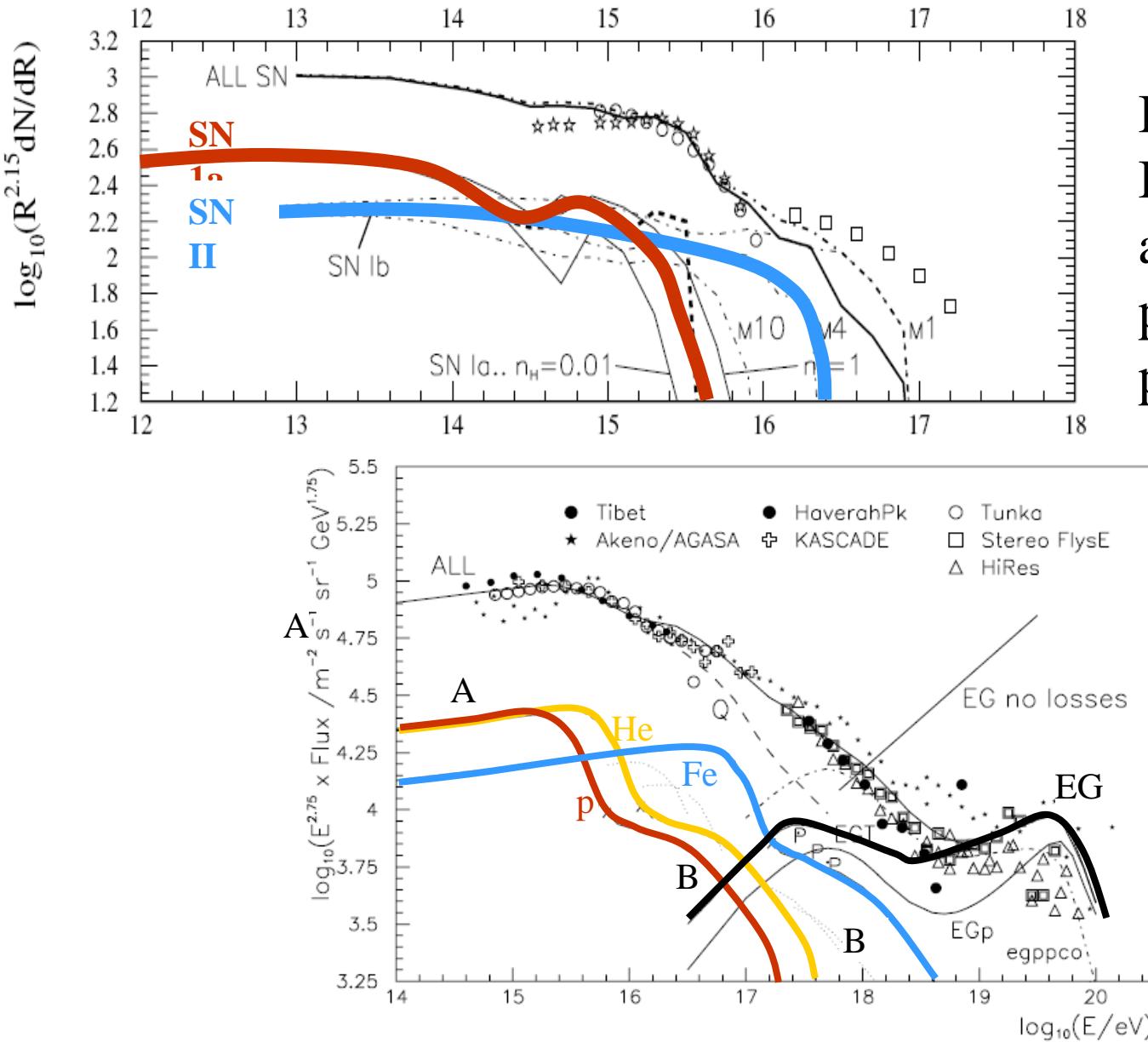
It is possible to find an agreement between diffusive propagation models and standard SNR models,

- Cut off energy, knee, non-linearities, γ -ray emission by SNR, source distribution...

Many parameters) need many observational constrains.

A conciliation exemple CR SM & diffusion

M.Hillas J. Phys., Conf. Ser. 47 (2006) 168 (<http://iopscience.iop.org/1742-6596/47/1/021>)

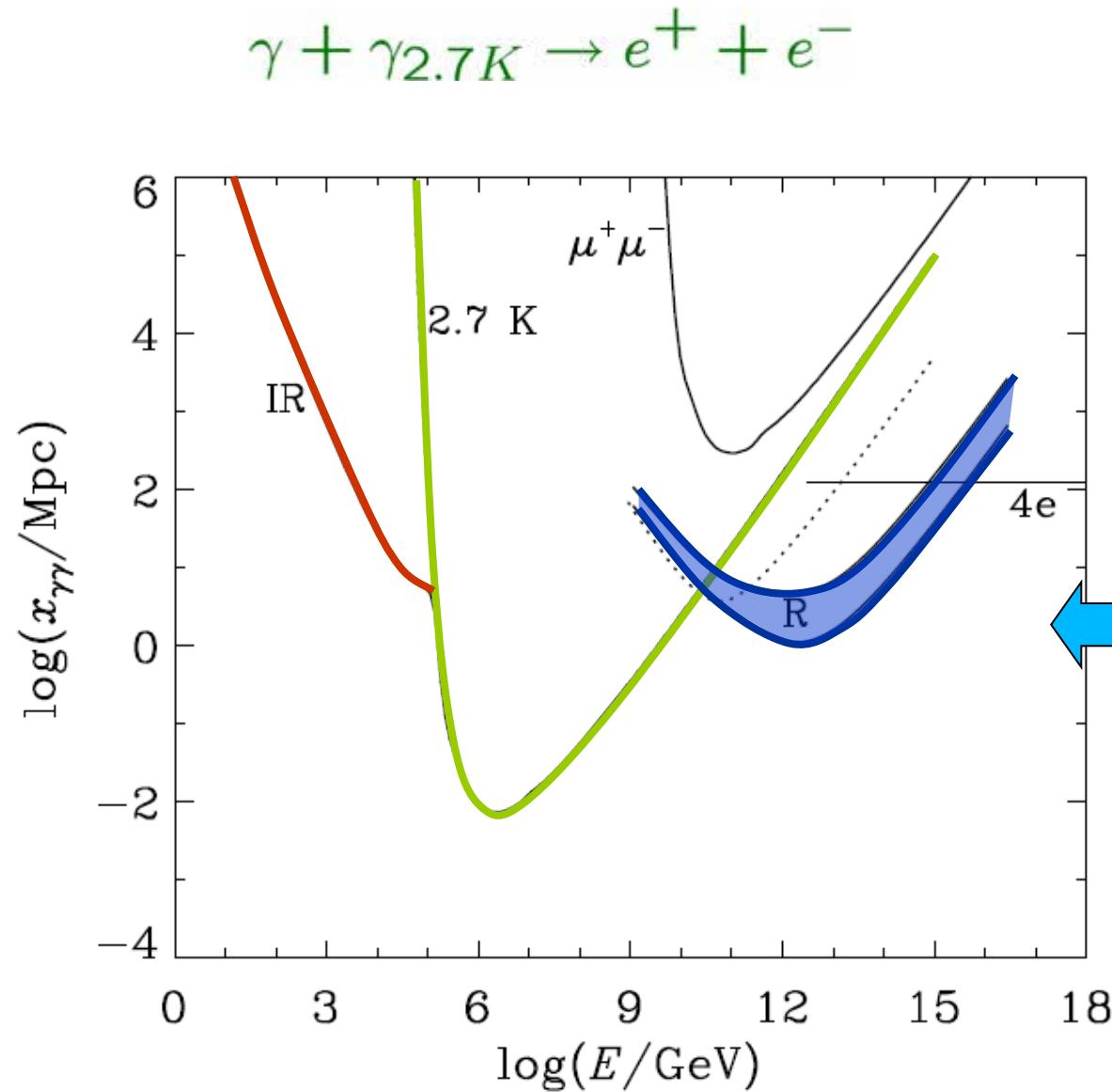


Based on former work by Bell et Lucek on amplification and self-production of Bfield perturbations by CR.

GAMMA-RAY PROPAGATION

Photon attenuation at VHE by intergalactic photon backgrounds

2010-2011



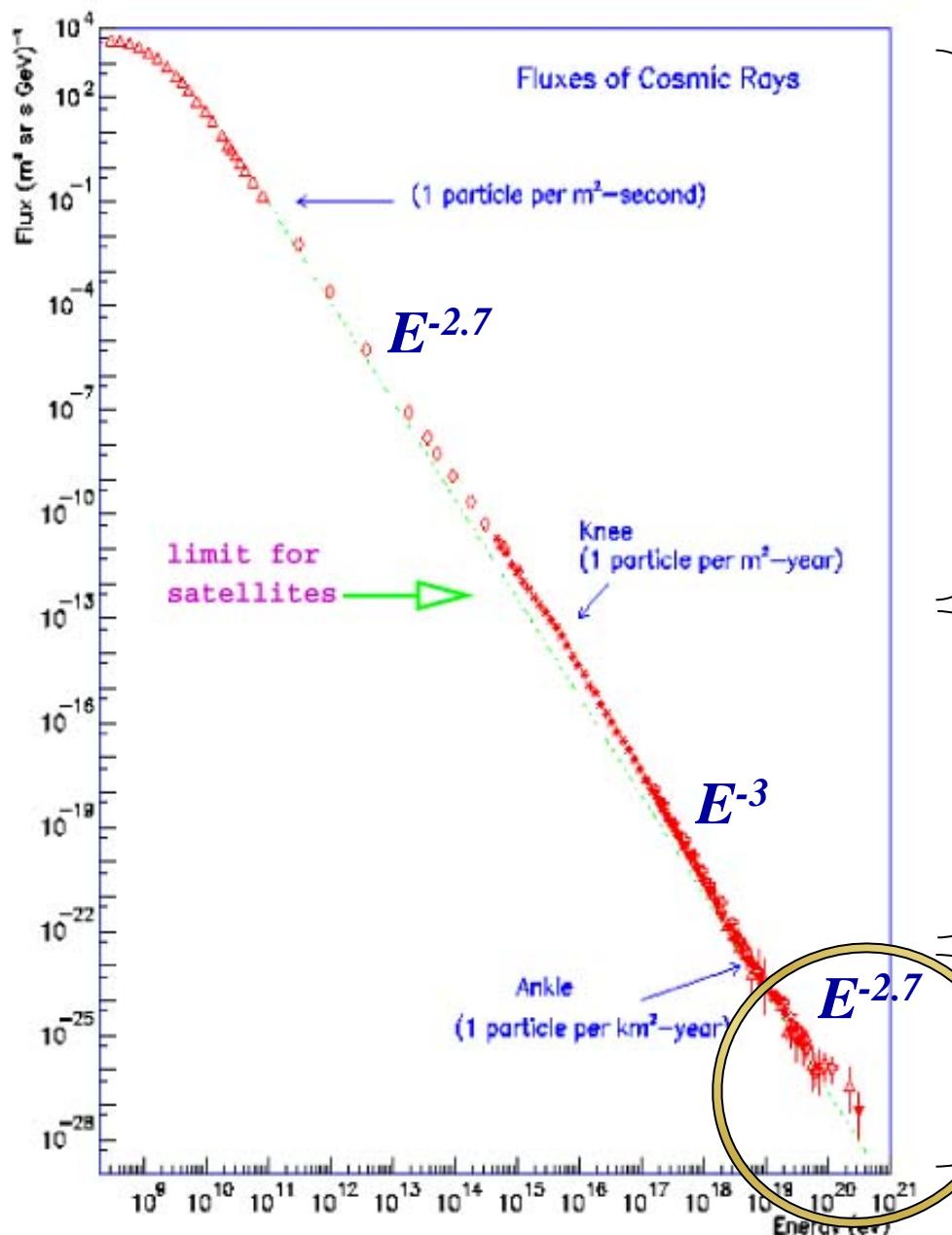
Effective ghorizon:
100Mpc at 1TeV
1Mpc at 10 TeV
and above

UHECR PROPAGATION

The CR spectrum

2010-2011

F.Montanet Cours Astroparticules M2R PSA & AMD Grenoble



Galactic CR :
Supernovae, MIS,
but no source pointing!

Galactic ?
SuperNovae? Superbubbles?
reacceleration?
Heavier nuclei ® protons ?

Extragalactic ?
source ?, composition ?

UHECR, terra incognita

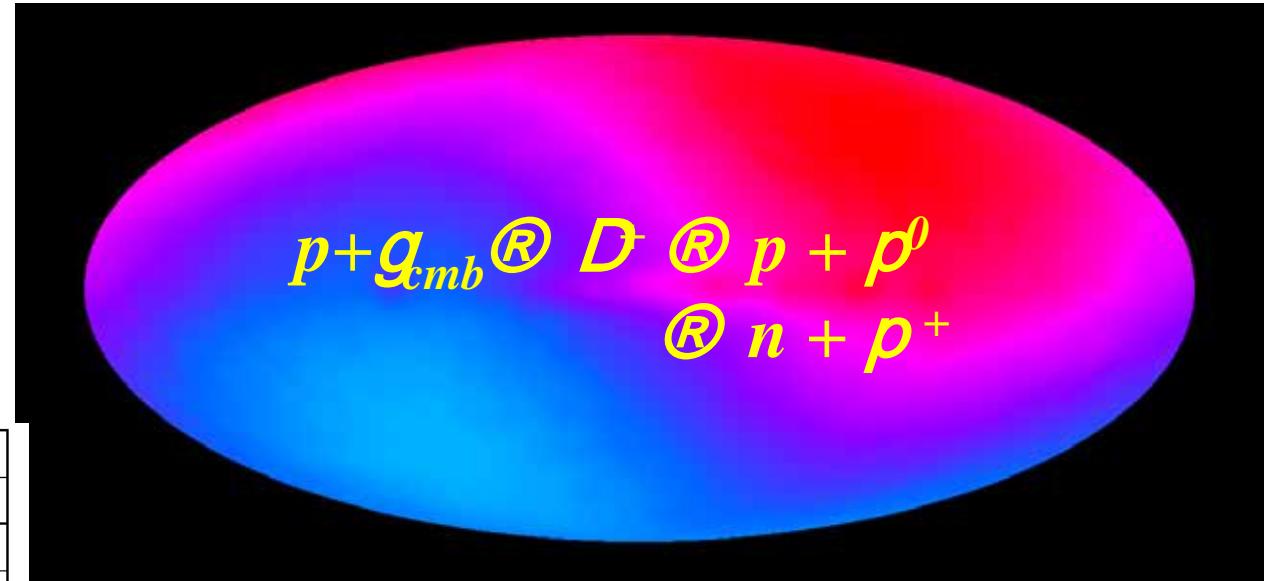
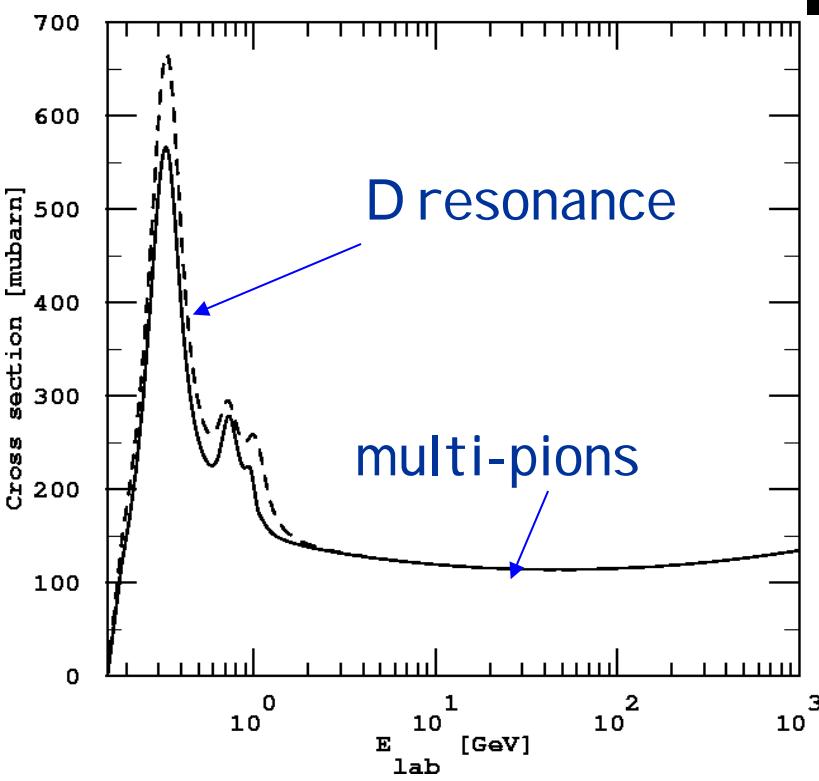
AUGER

UHECR propagation

3 essential effects :

- Energy losses: modify the spectral shape
- Particle confinement
(escape depending on energy)
- Spatial and angular diffusion
due to magnetic fields.
(regular or fluctuating, inhomogeneities, waves)

An extrem case of relativistic kinematics !!!

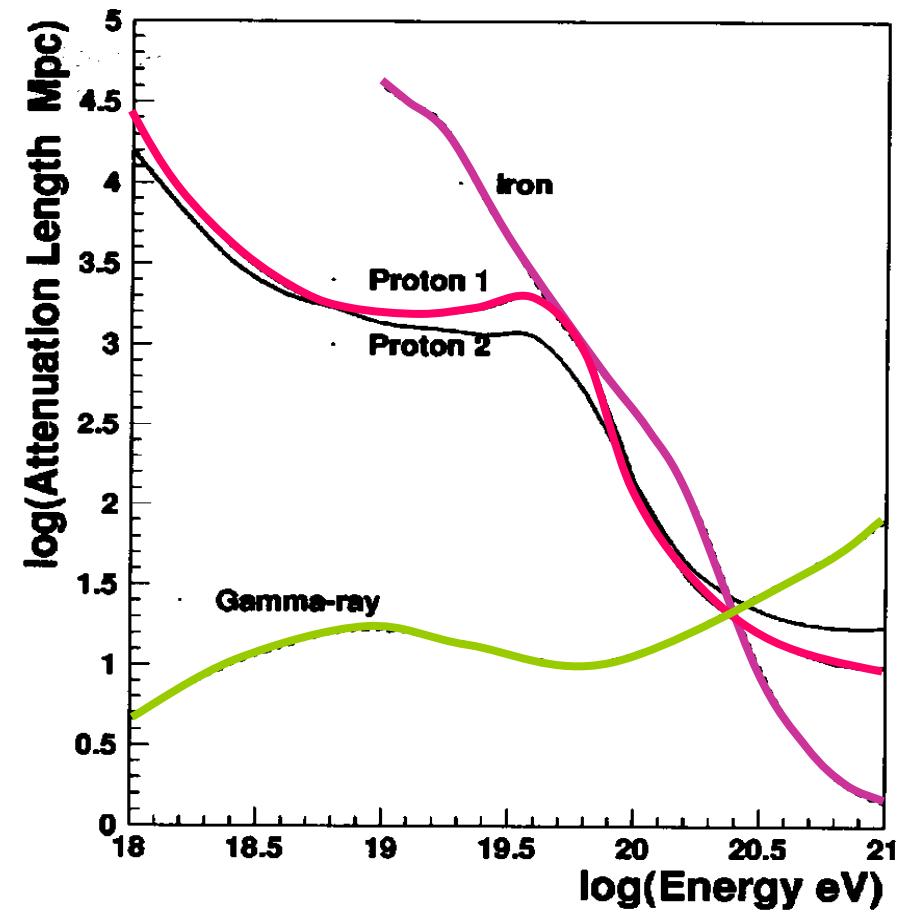
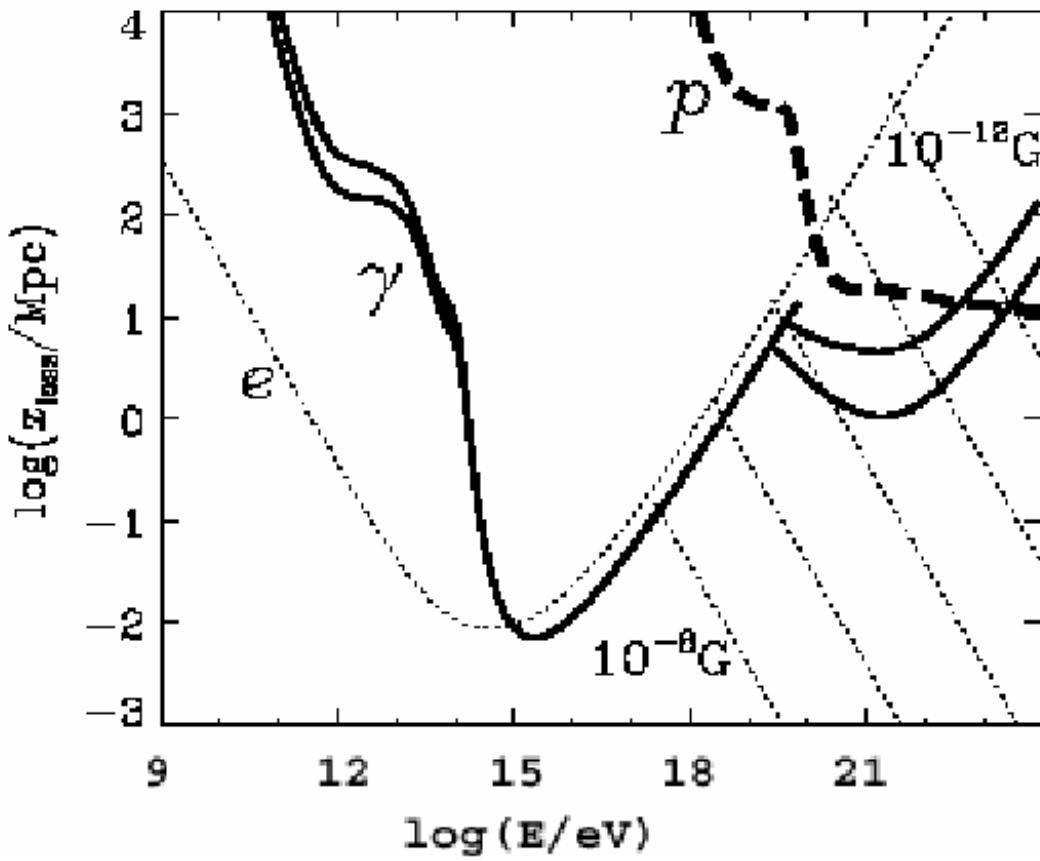


GZK
"cutoff"

Greisen '66, Zatsepin & Kuzmin '66

Energy losses

- $p + \gamma_{2.7K} \rightarrow n + \pi^+; p + \pi^0; p + e^+ + e^-$
- $A + \gamma_{2.7K} \rightarrow (A - 1) + N; (A - 2) + 2N; A + e^+ + e^-$
- $\gamma + \gamma_{2.7K} \rightarrow e^+ + e^-$



Consequences on spectral shape

Protons:

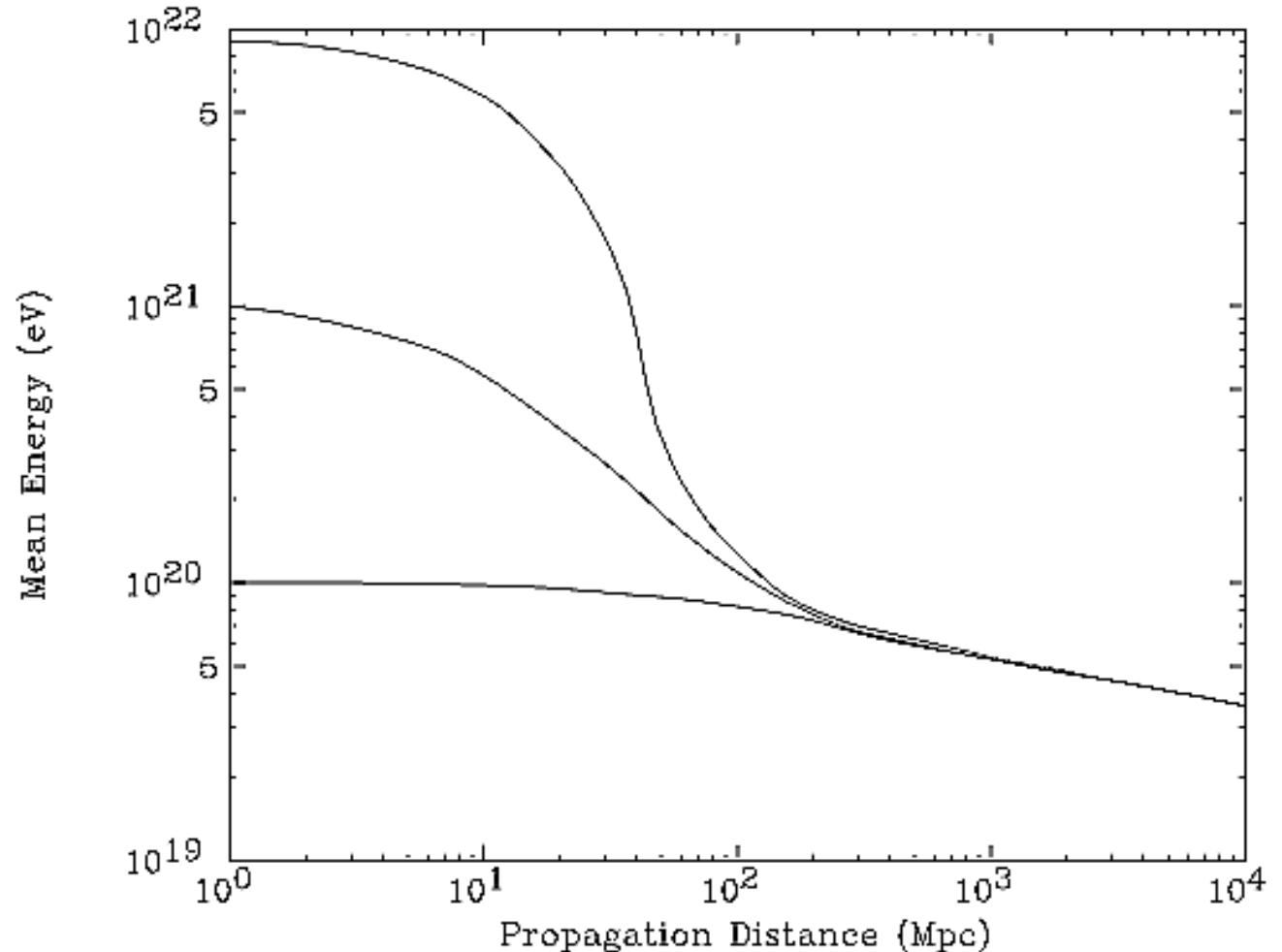
*Photo Pion production
CMB photons*

$p + g_{\text{cmb}} \xrightarrow{\text{R}} D \xrightarrow{\text{R}} p/n + p$

Fe:

*photo-dissociation
on IR bg and CMB*

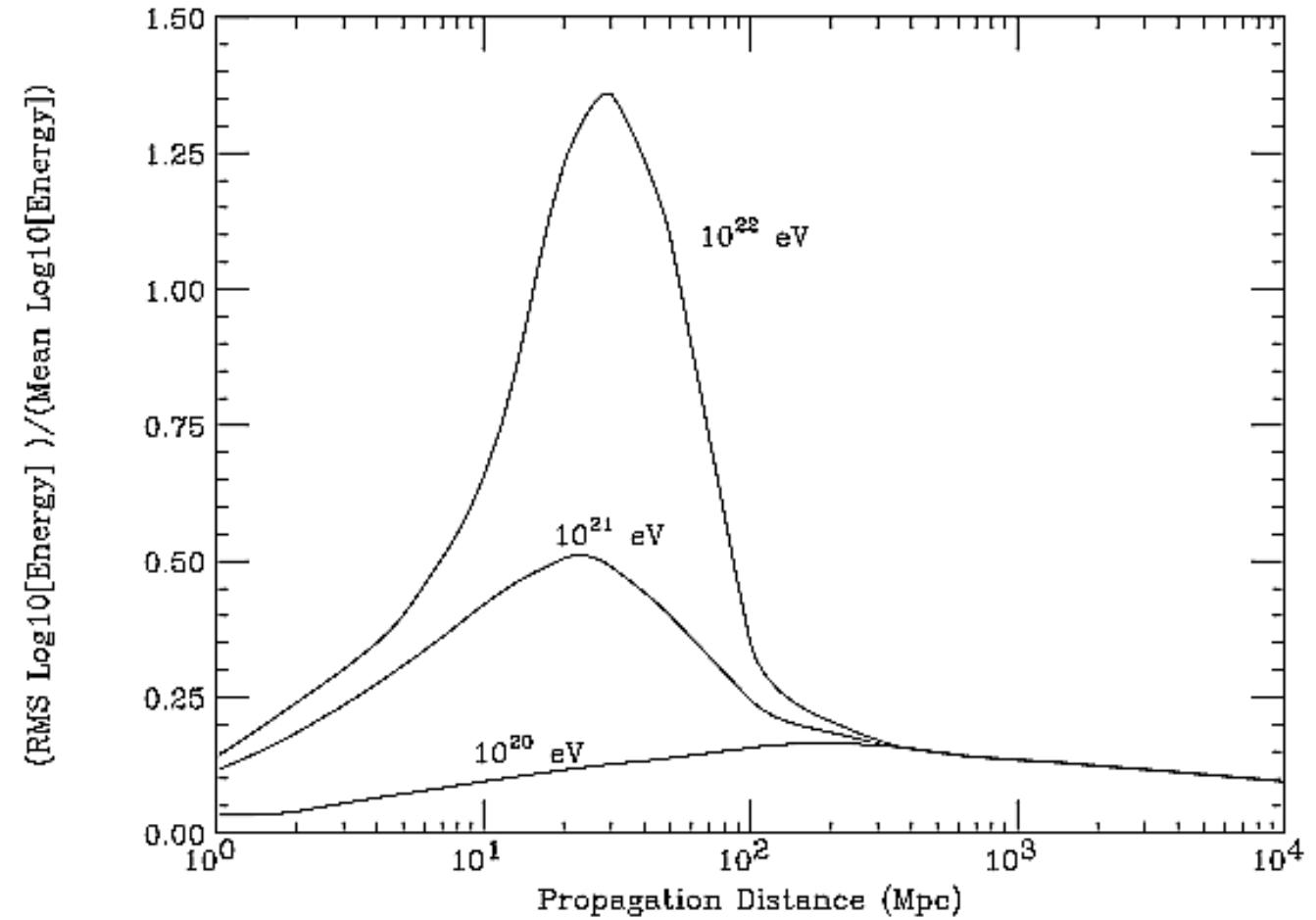
*Photons: paires e^+e^-
on radio bg*



Protons energy vs. distance (J. Cronin)
Energy loss on CMB

UHE Extragalactic Particles

Fluctuation dues to multiple scattering



Protons energy vs. distance (J. Cronin)
Energy loss on CMB

UHE Extragalactic Particles

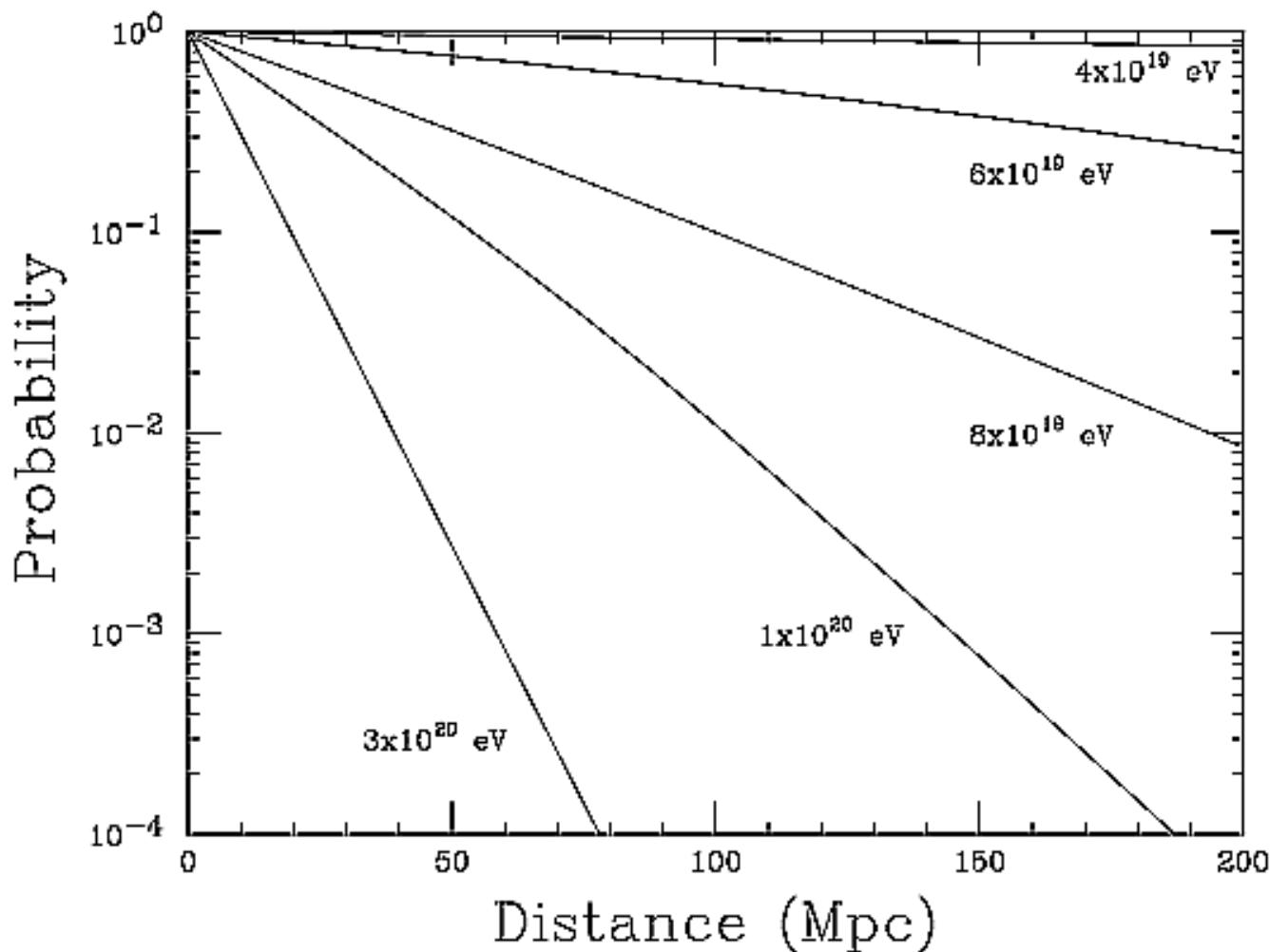
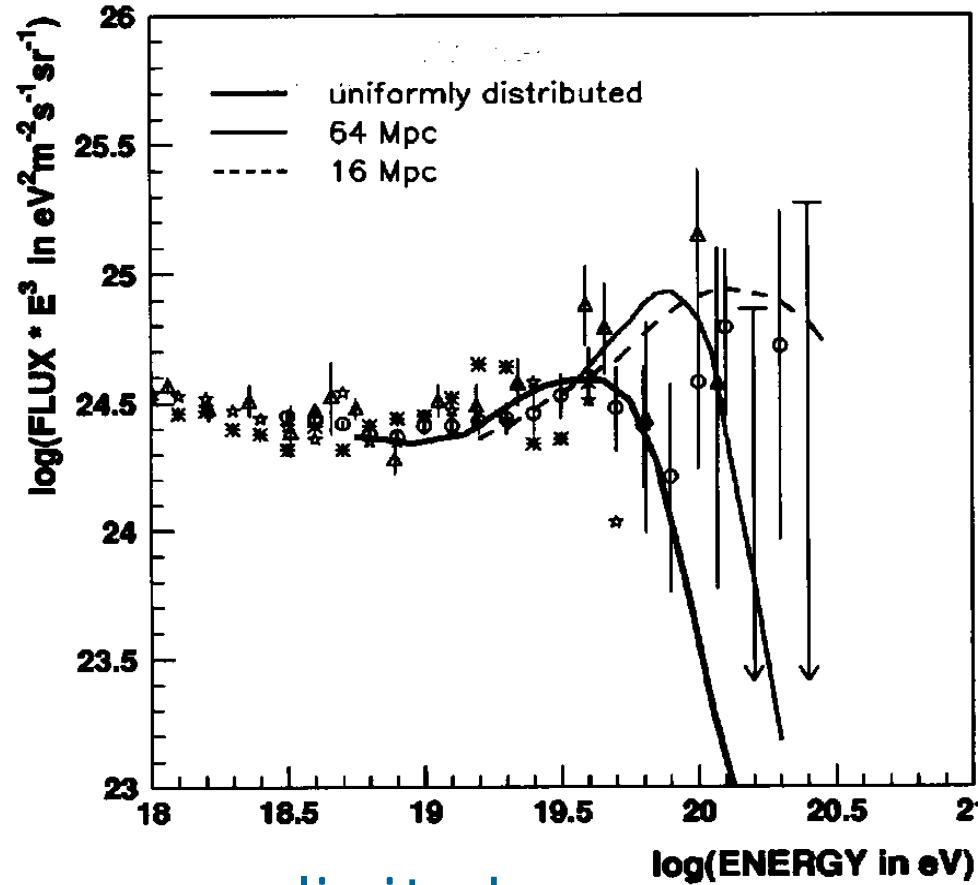


Figure 6: Probability that an observed event at a given energy has its source at a distance greater than the indicated distance. A source spectrum proportional to $E^{-2.5}$ is assumed. Figure provided Paul Sommers, University of Utah.

Coupure GZK

- Greisen-Zatsepin-Kuz'min

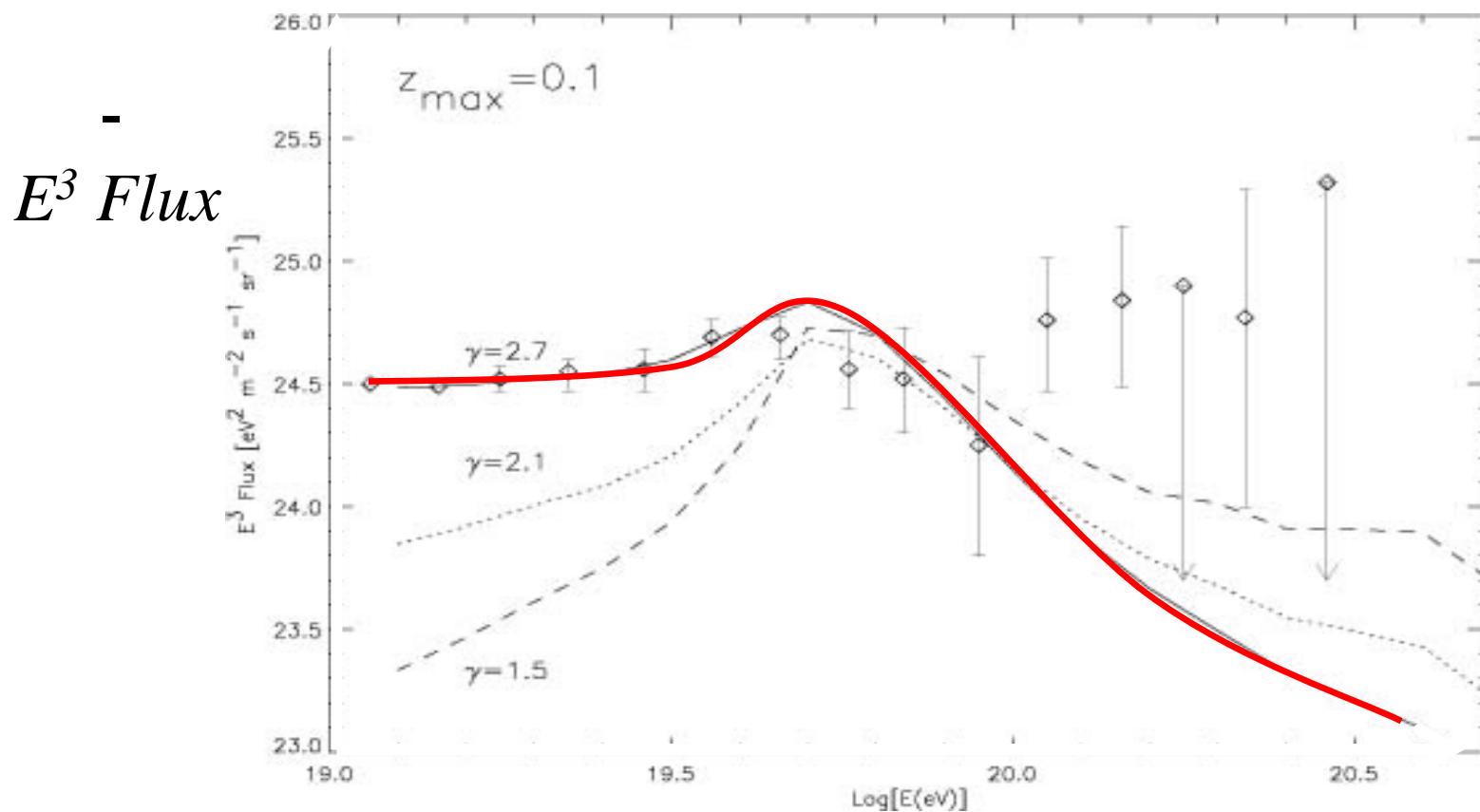


- Distance to the source limited to 10^{20} eV protons, and 15 Mpc for $3 \cdot 10^{20}$!
- Actually even worse if particles are deflected ($D_{\text{effectif}} > D_{\text{linear}}$)

Spectrum above GZK cutoff

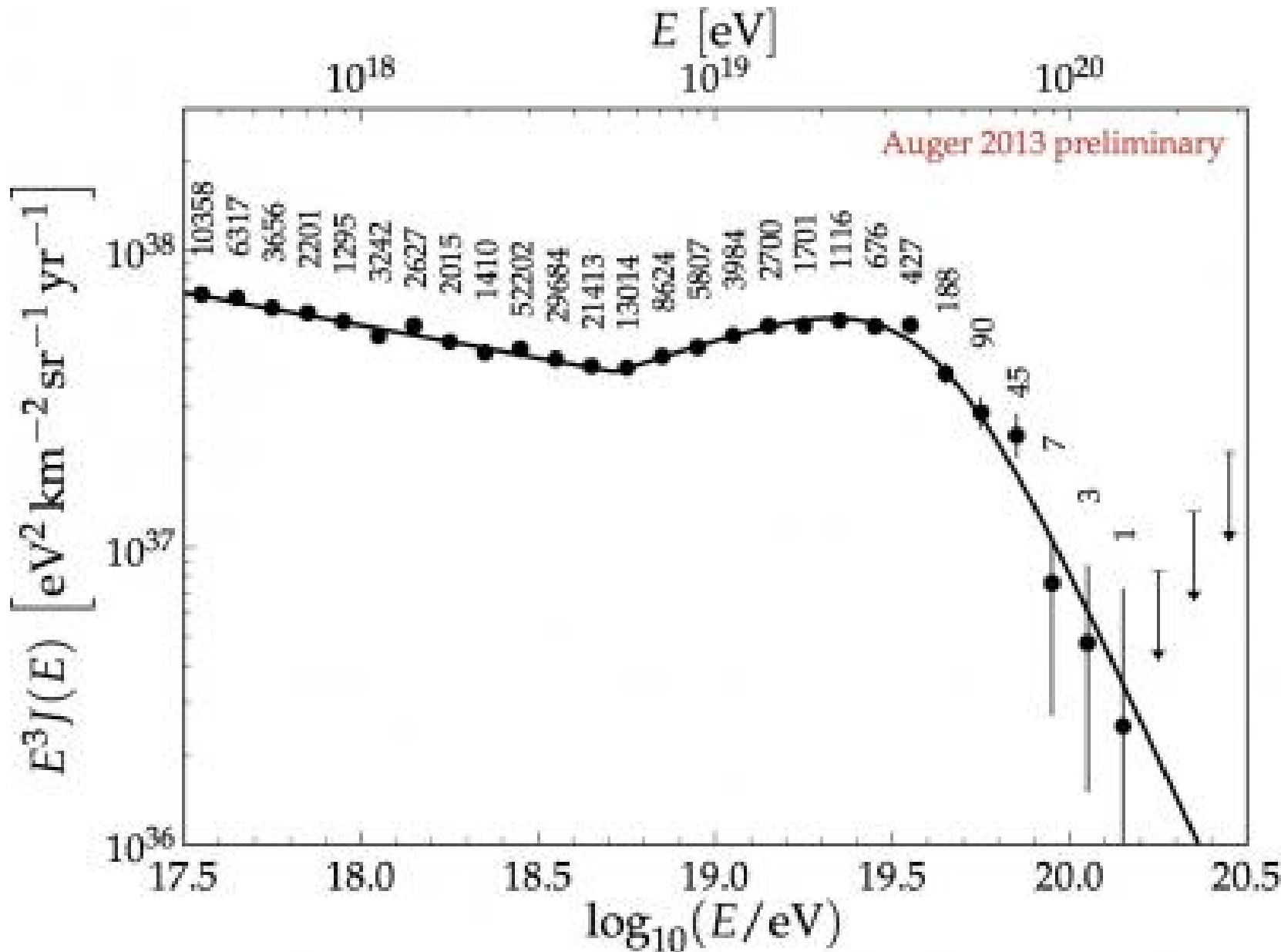
Depends among other things:

- ━ From the injection spectrum
- ━ From the sources cosmological distribution (evolution)
- ━ From IG magnetic fields

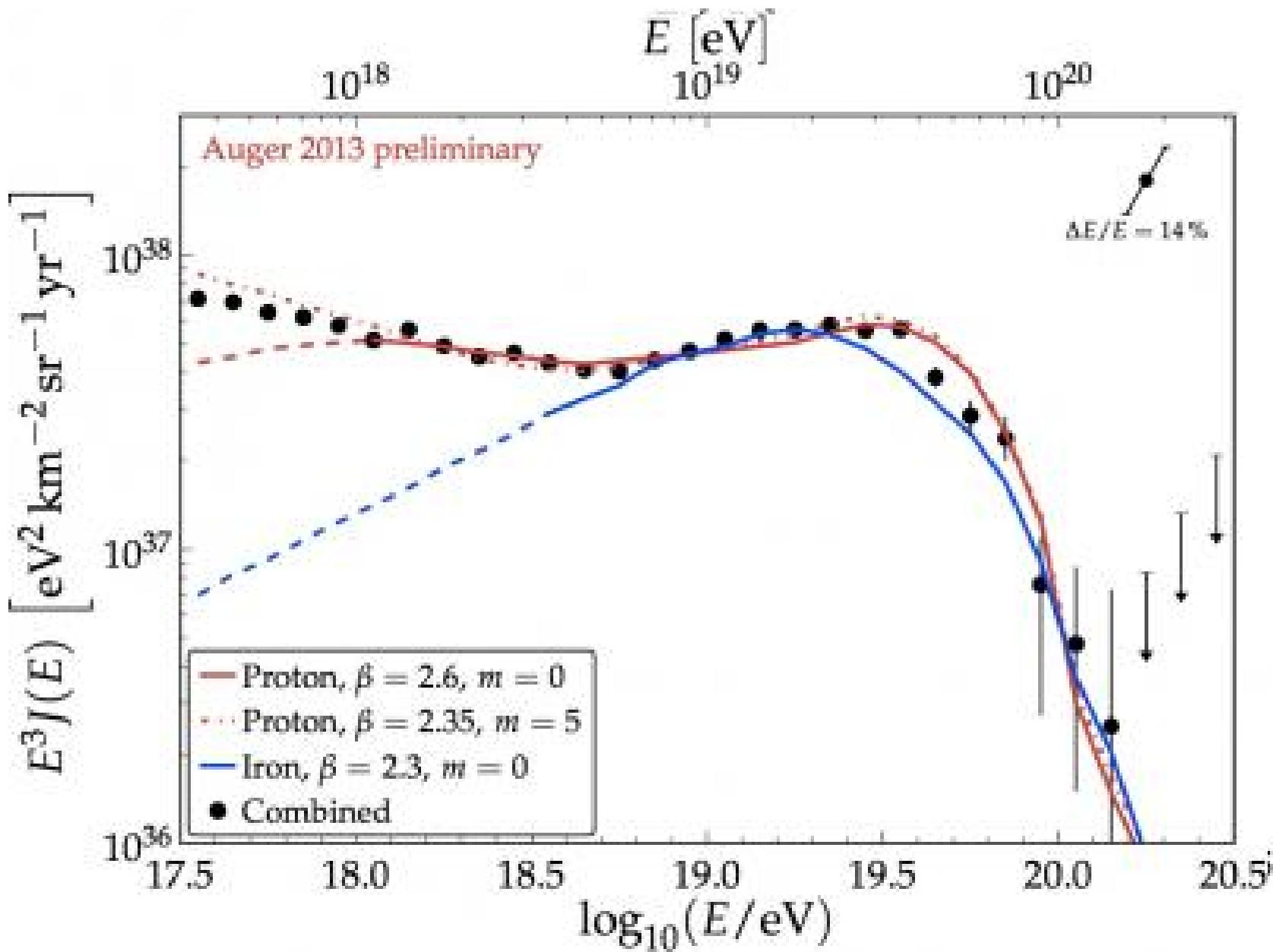


Blasi, Olinto. '01

GZK like suppression (Auger)



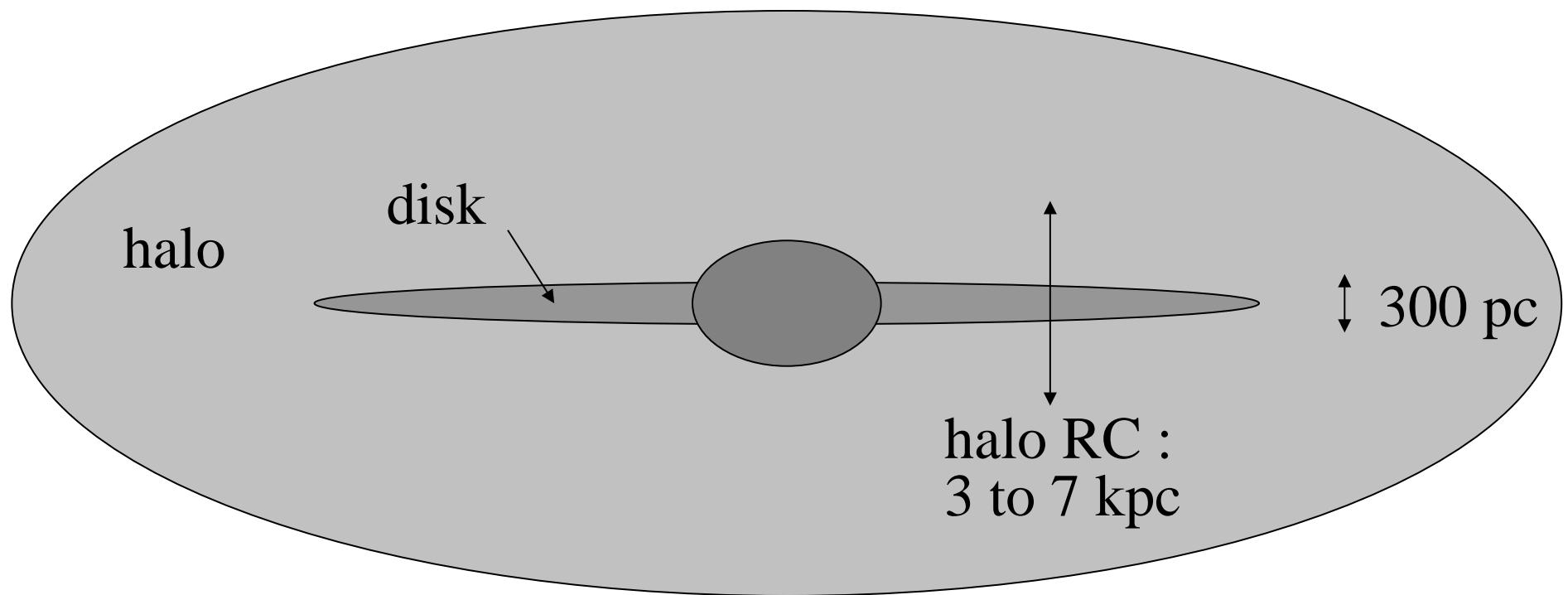
GZK like suppression (Auger)



MAGNETIC DEFLECTIONS

Galactic magnetic deflection

- 10^{18} eV proton in a $B = 3 \mu\text{G}$ field $\rightarrow r_g \sim 370 \text{ pc}$



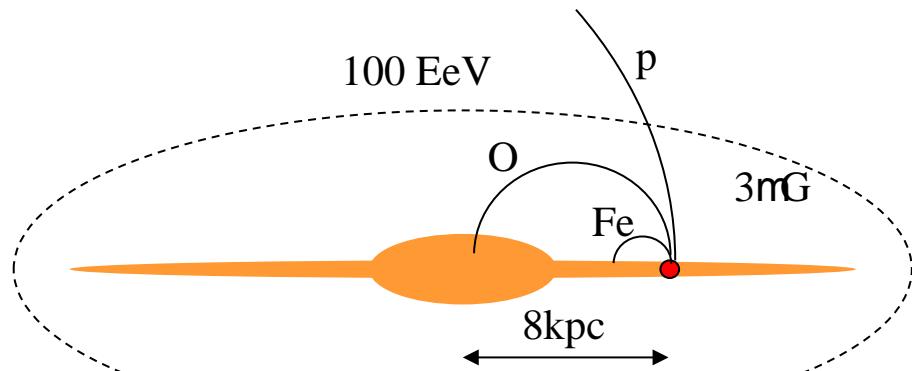
- $2 \cdot 10^{19}$ eV proton in $B = 3 \mu\text{G}$ $\rightarrow r_g \sim 7 \text{ kpc}$
- $5 \cdot 10^{20}$ eV Fe in $B = 3 \mu\text{G}$ $\rightarrow r_g \sim 7 \text{ kpc}$

Propagation the Galaxy

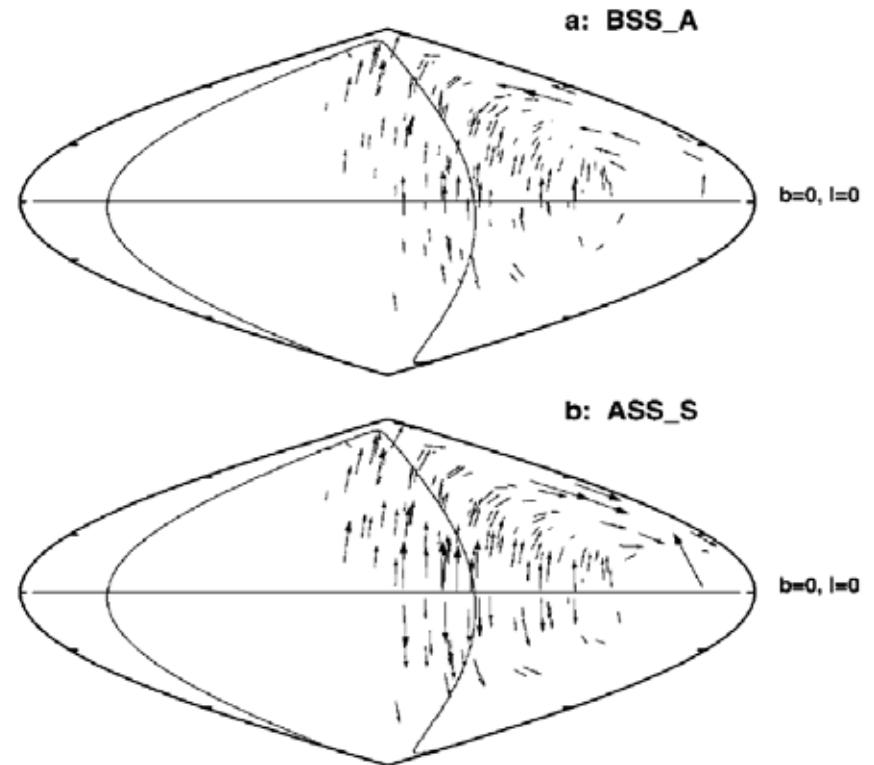
2010-2011

- Galactic magnetic field model

$$\frac{eR_{Larmor}}{kpc} = \frac{eZ}{e^2 m G} \frac{E}{1EeV} \frac{B}{mG}$$



- Possible galactic confinement of $10^{20} eV$ nuclei
- $10^{18} eV$ neutrons decay length
 $\gg 10 kpc \gg$ galactic distances

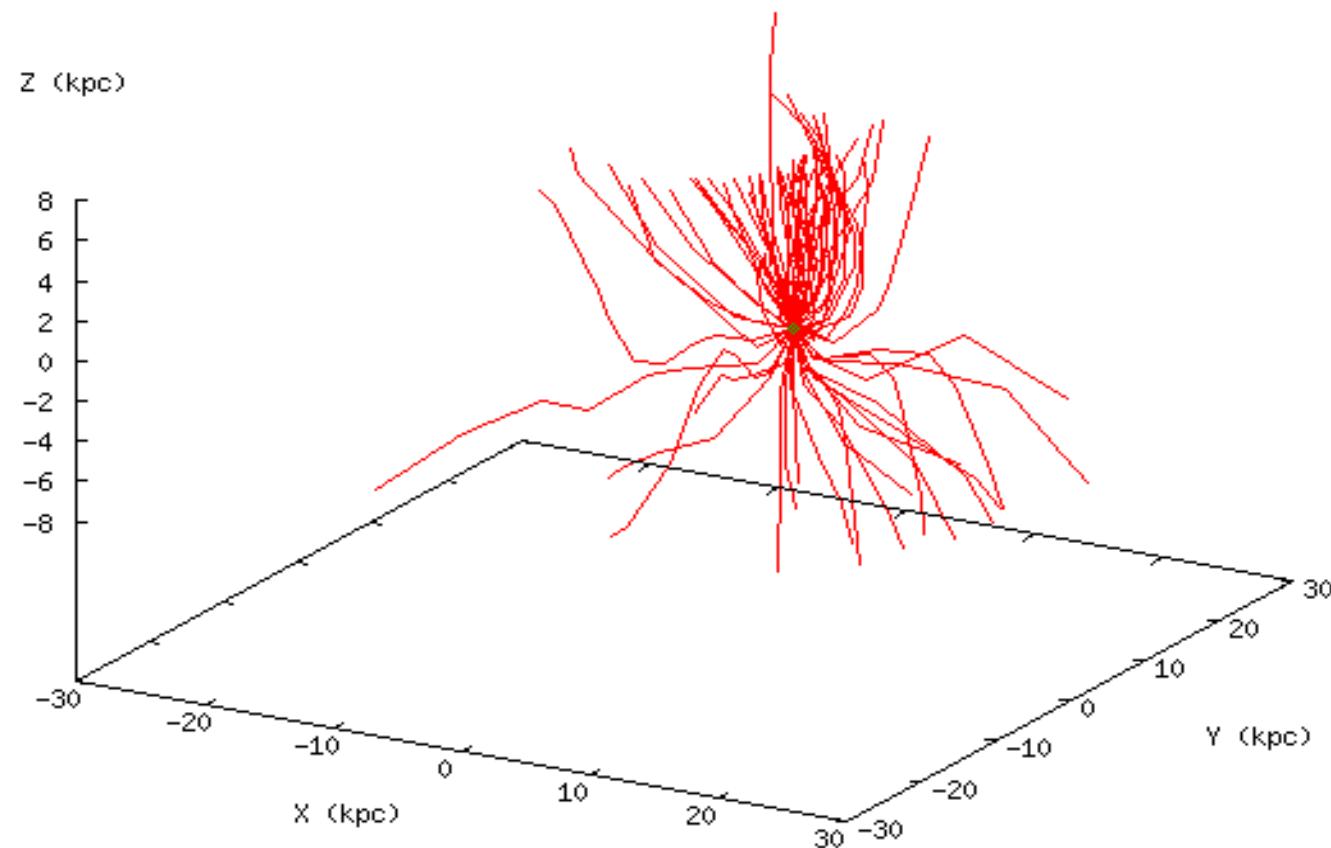


Tracking back direction of proton events
 $>4 \times 10^{19}$ out of the Galaxy, two different field hypothesis [Stanev97]

Pointing at UHECR sources?

2010-2011

100 EeV Iron Nucleus Distribution Under the Influence of Regular Galactic Field and Galactic Wind Field



O'Neil, Olinto, Blasi '01

Pointing at UHECR sources?

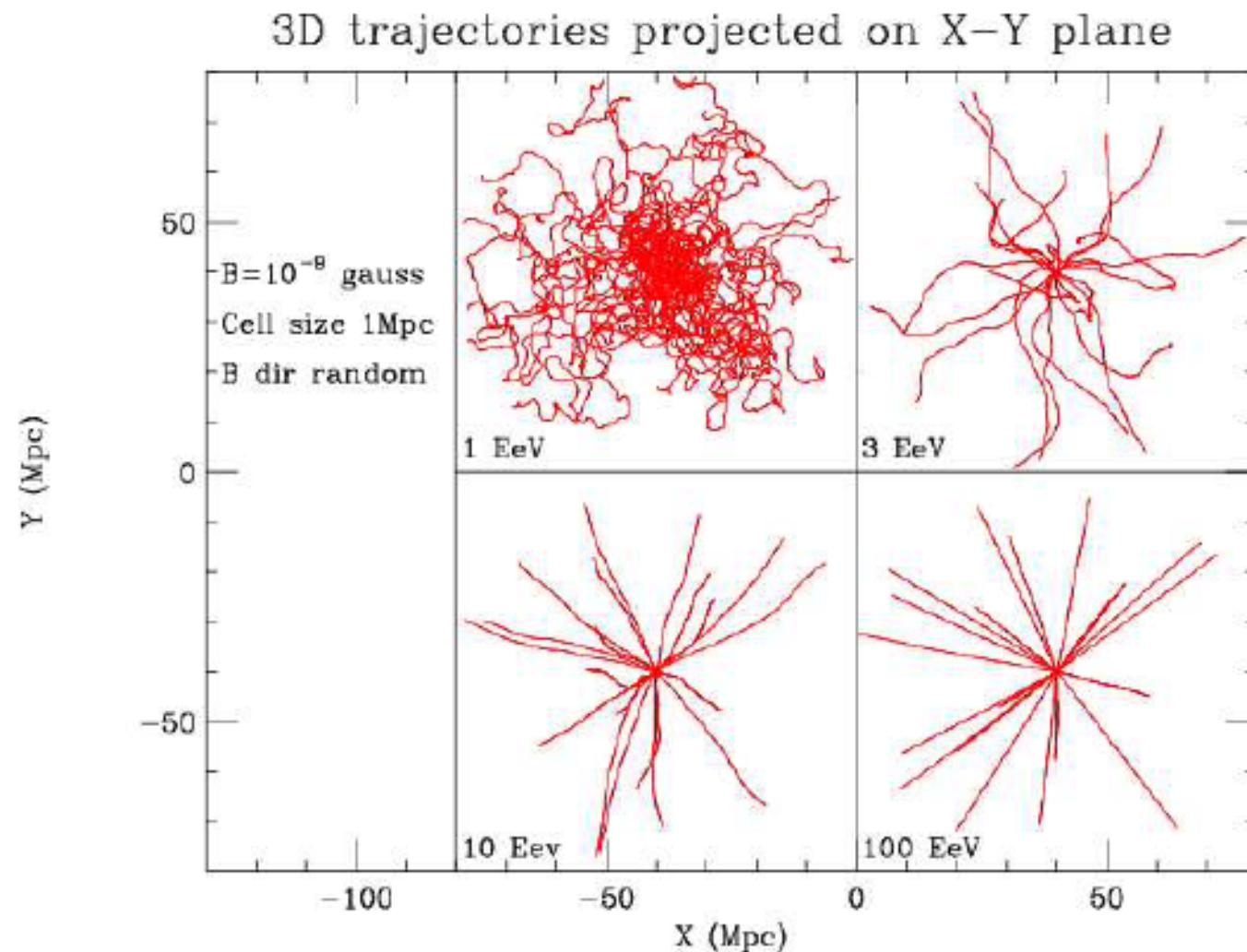


Figure 7: Projected view of 20 trajectories of proton primaries emanating from a point source for several energies. Trajectories are plotted until they reach a physical distance from the source of 40Mpc. See text for details.

Mapping IG fields with UHECR?

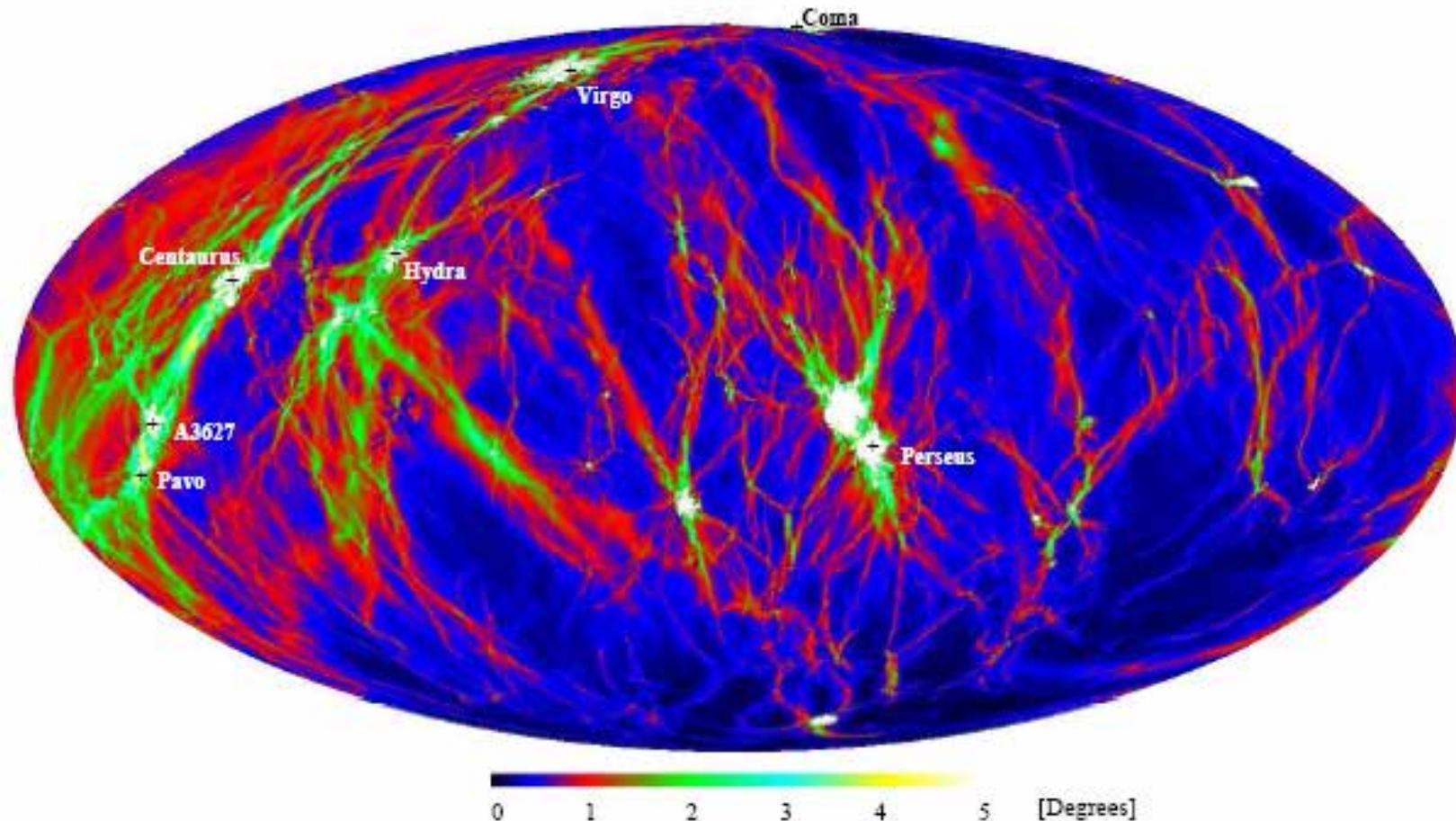


FIG. 1. Full sky map (area preserving projection) of deflection angles for UHECRs with energy 4×10^{19} eV using a linear color scale. All structure within a radius of 107 Mpc around the position of the Galaxy was used. The coordinate system is galactic, with the galactic anti-center in the middle of the map. Positions of identified clusters are marked using the locations of the corresponding halos in the simulation. Note that deflections internal to the Milky Way have not been included.

Diffusion in the Universe

- If they are protons,
arrival direction \approx source direction

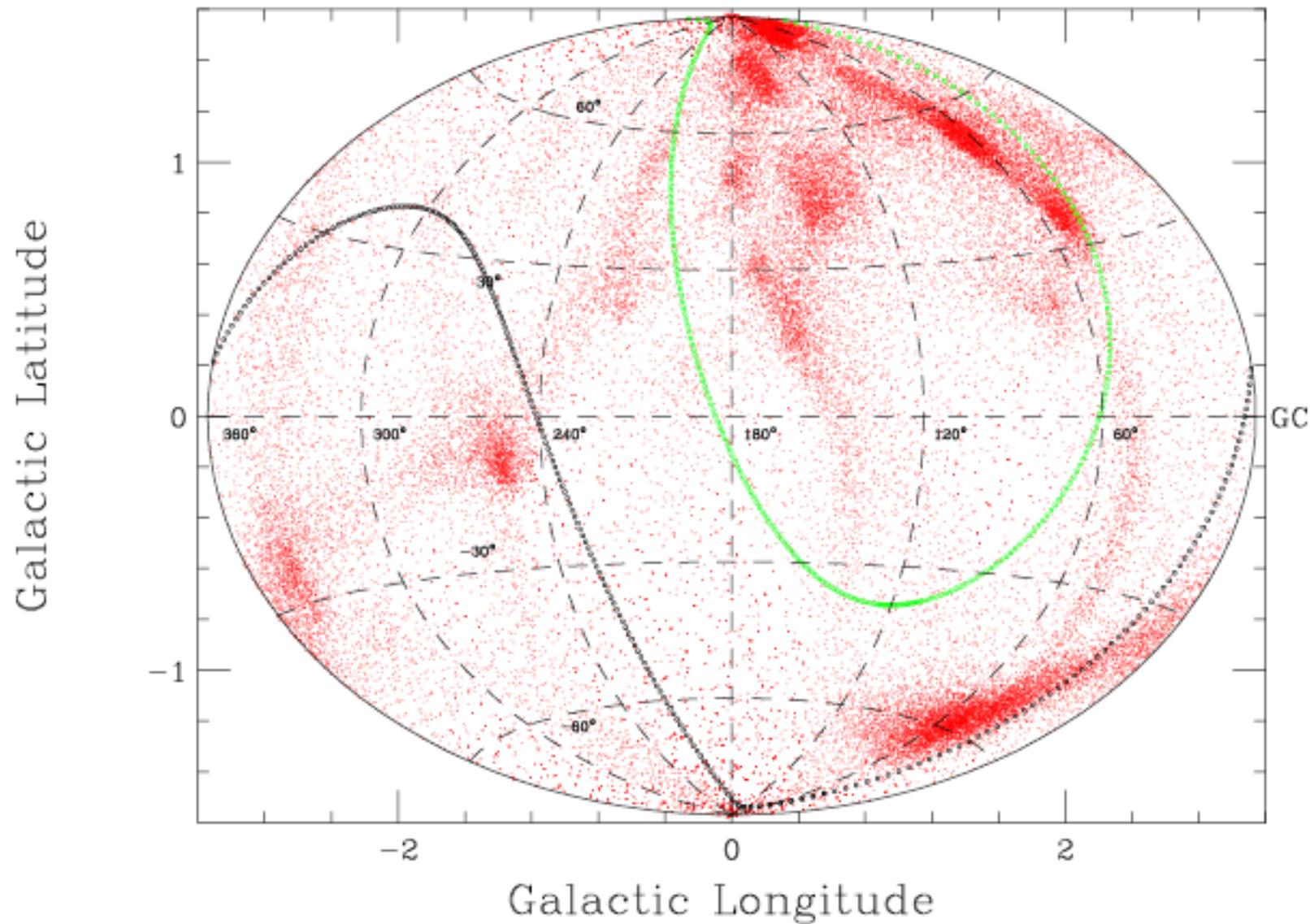
$$r_L \approx 100 \text{ kpc} \times Z \times (E/10^{20} \text{ eV}) \times (B/10^{-6} \text{ G})^{-1}$$

$\delta\theta \sim \lambda_B/r_L$ deviation per field correlation length $\Rightarrow \Delta\theta \sim \sqrt{D\lambda_B}/r_L$
® Proton astronomy !

- Correlations between arrival directions and sources:
UHECR distribution is NOT ISOTROPI C!!! (AUGER 2007)
 - Confirmation of a GZK limited horizon
 - Few sources in the GZK sphere ® anisotropie
 - Astrophysical origine is confirmed!
 - Arrival time delay
 - $\Delta t \sim \Delta\theta^2 d/c \sim D^2 \lambda_B / r_L^2 c$
 - If eruptive or transient sources (GRBs, TDs), they must overlap in time (otherwise E(t)!!)
 - Multiplets of events from same direction observed but no significant ordering in E or deviation.
- ® Correlation must be confirmed! (statistics...)

Matter distribution in the GZK sphere

Matter distribution 7–21 Mpc. Exclusion zones; north array (black), south array (green)



Observables & Observations

PRIMARY RC DETECTION (ON TOP OF ATMOSPHERE)

How to characterize the primary particle?

- Mass m
- Electric charge Ze
- Velocity $v = \beta c$
- Lorentz Facteur $\gamma = E/mc^2$
- Momentum $p = mc\beta\gamma$
- Kinetic energy $T = mc^2(\gamma - 1)$

How to characterize the primary particle?

Detector	Observable	Link with the particle
Magnetic spectrometer	Rigidity & Sign of Z	pc/Ze
Time of flight	Velocity/c	β
Proportionnal counters Scintillators Ionisation chamber	Ionisation	$dE/dx = Z^2 f(\beta)$
Čerenkov effect	Č photons density	$dN/dx = Z^2 g(\beta)$
Transition radiation	Number of photons X	$N = Z^2 h(\gamma)$
Calorimeter	Deposited energie	$mc^2(\gamma - 1)$

Two important radiations for particle identification

Two effects of the **polarization** induced by charged particles in dielectric medium

Proportionnal to Z^2

- **Čerenkov radiation** : si $v > c/n$
Sensitive to $\beta = v/c$
- **Transition radiation** : at the interface of \neq dielectric media
Sensitive to $\gamma = E/(mc^2)$

Cherenkov Radiation

- Emitted on a cone whose axis is along the particle trajectory and of half opening angle θ_C tel que $\cos \theta_C = 1/(\beta n(\omega))$
- Threshold defined by the condition $\cos \theta_C < 1$
- Emission at all frequencies with $n(\omega) > 1$ (from UV to radio), flat in $h\nu$ in photon yield Generally detected from near UV to visible light.

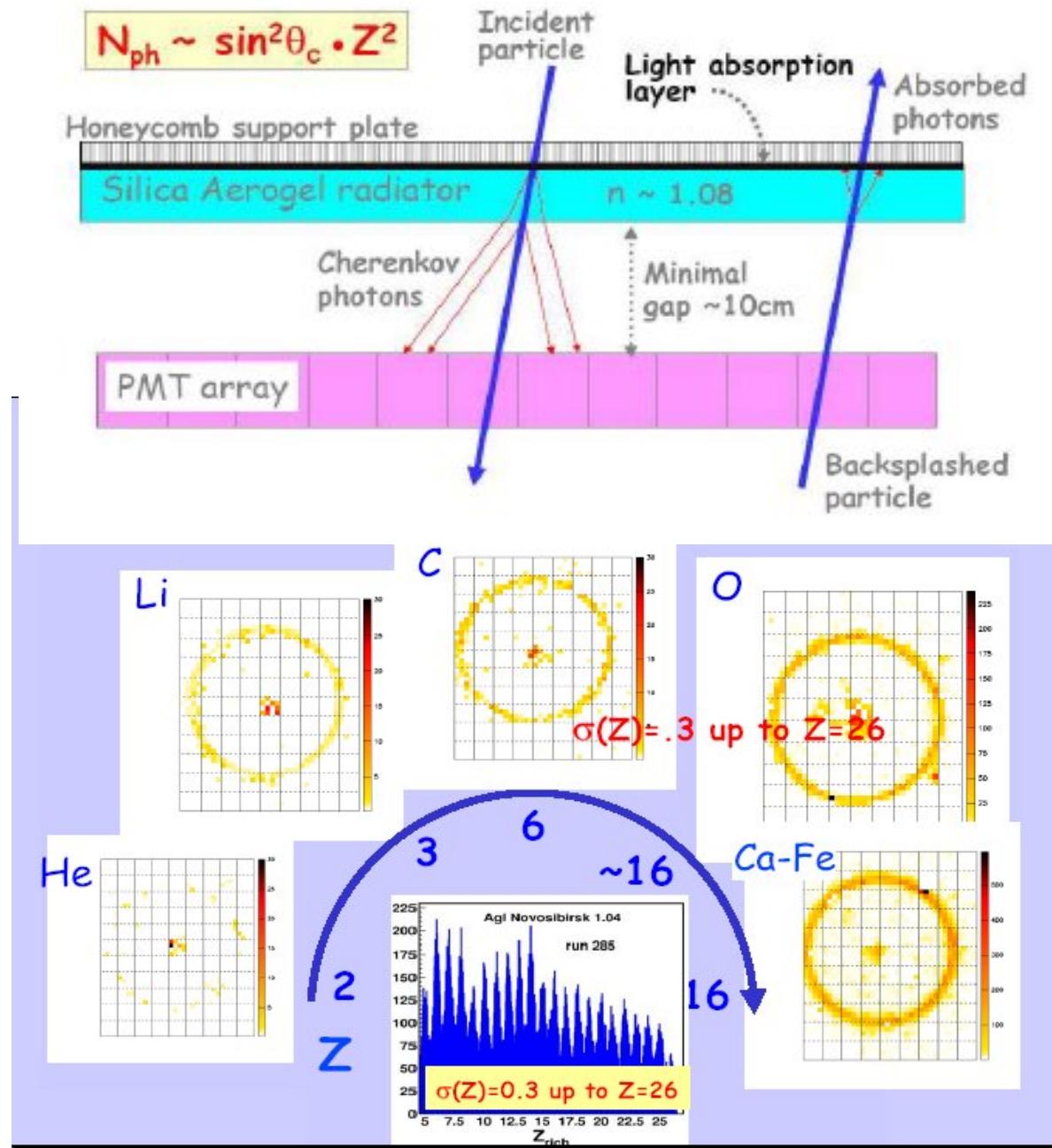
$$\frac{dN^2}{dxd\omega} = \frac{\alpha}{c} Z^2 \sin^2 \theta_C$$

- Allow discrimination between particles with same momentum and \neq masses.
(électrons / protons / noyaux) up to energies of the order of 10 GeV/nucleon if $\Delta\beta/\beta \approx 10^{-3}$
- Allow determining the orientation of the particle direction.
- "Ring Imaging Cherenkov" detector or RICH : allow a precise measurement of the charge.

Cherenkov imaging (RICH) and charge measurement

**RICH
principle →**

**AMS 2
Prototype
→**



Transition radiation

- Origine : if a particle traverse the boundary between 2 \neq dielectric media, the solution corresponding to each medium do not satisfy the conditions at the boundary.
→ nécessité d'une onde «libre» supplémentaire.
- A dielectric medium is characterised by its "plasma" frequency ω_p (oscillation frequency of free-like electrons)

$$\omega_p = \sqrt{\frac{n_e e^2}{\epsilon_0 m_e}} \text{ or } \hbar \omega_p = 2E_R \sqrt{4\pi n_e a_B^3} \approx 20 \text{ eV}$$

n_e = electron density ;
 E_R = Rydberg energy = 13.6 eV ;
 a_B = Bohr radius
- Roughly half of the energy is emitted in the frequency domaine $0.1\gamma\omega_p < \omega < \gamma\omega_p$ for a Lorentz factor $\gamma \approx 1000$, This is the X-ray domaine (2 to 20 keV)
- Energy emitted by the interface if $I = \alpha z^2 \gamma \hbar \omega_p / 3$

Transition radiation (cont)

- Angular distribution peaked at small angles around particle direction:
 $\theta \approx 1/\gamma$
- Small yield of X-photons per interface : $N \approx \alpha Z^2 \approx 10^{-2}Z^2$
 - multiply the number of interfaces
 - stack plastic sheets or fibers
- X-ray detection by photo-electric effect: proportional tubes
- Discriminates between particles with same energy and \neq masses at high energies (100 GeV to 1 TeV) (instrumental detection threshold for X-rays).
- Can measure the Lorentz factor γ up to 10^5 . In this case, choose material adequately to have a progressive threshold.

Charge particles and cosmic antimatter

Satellite born experiment : AMS-1, PAMELA, AMS-2

- First satellite based experiments on cosmic rays : HEAO-C, Ariel-VI (1979) → relatively low energy (up to a few 10 GeV/nucleon)
- First satellite based magnetic spectrometer en satellite : AMS-1 on the space shuttle « Discovery » (1998)
- New generation of experiments:
PAMELA (since Juin 2006)
AMS-2 (since May 2011)
on the International Space Station → data up to ~TeV
and precise measurements of flux of cosmic antiparticles

AMS-2 On Board ISS

2014

Mission Number: STS-134

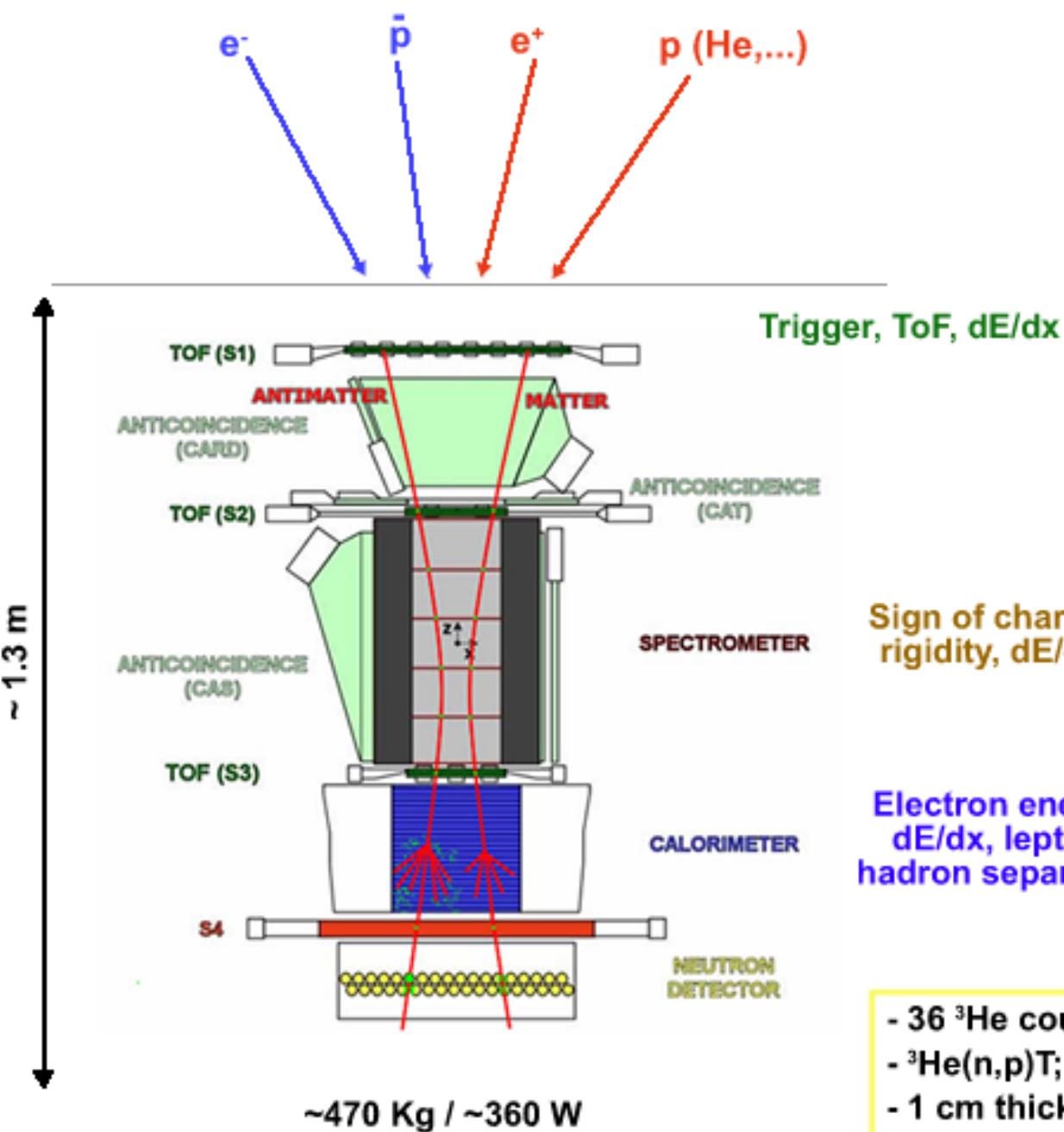
Launch: May 19, 2011

Orbiter: Endeavour



Space spectrometers

	AMS-1 (June 1998)	PAMELA (June 2006 - ...)	AMS-2 (May 2011 - ...)
Spectrometer Acceptance	0.82 m ² sr	20.5 cm ² sr	0.82 m ² sr
Spectrometer	Aimant permanent Nd Fe B 0.15 T $BL^2 = 0.15 \text{ T m}^2$ 6 plans (Si)	Aimant permanent Nd Fe B 0.48 T $BL^2 = 0.10 \text{ T m}^2$ 6 plans (Si)	Aimant permanent Nd Fe B 0.15 T $BL^2 = 0.15 \text{ T m}^2$ 6 plans (Si)
Time of Flight	yes	yes	yes
Cherenkov	Aerogel (threshold)	-	Ring Imaging Ch.
Transition rad	-	yes	yes
Neutrons det.	-	³ He	-
Anticoincidence	-	yes	yes
Calorimeter	-	16,3 X ₀ W+22 plans (Si)	16 X ₀ Pb+fibers sc.



PAMELA

- S1, S2, S3; double layers, x-y
- plastic scintillator (8mm)
- ToF resolution ~ 300 ps (S1-3 ToF > 3 ns)
- lepton-hadron separation < 1 GeV/c
- S1.S2.S3 (low rate) / S2.S3 (high rate)

- Permanent magnet, 0.43 T
- $21.5 \text{ cm}^2 \text{ sr}$
- 6 planes double-sided silicon strip detectors (300 μ m)
- 3 μ m resolution in bending view \rightarrow MDR ~ 800 GV (6 plane) ~ 500 GV (5 plane)

Sign of charge,
rigidity, dE/dx

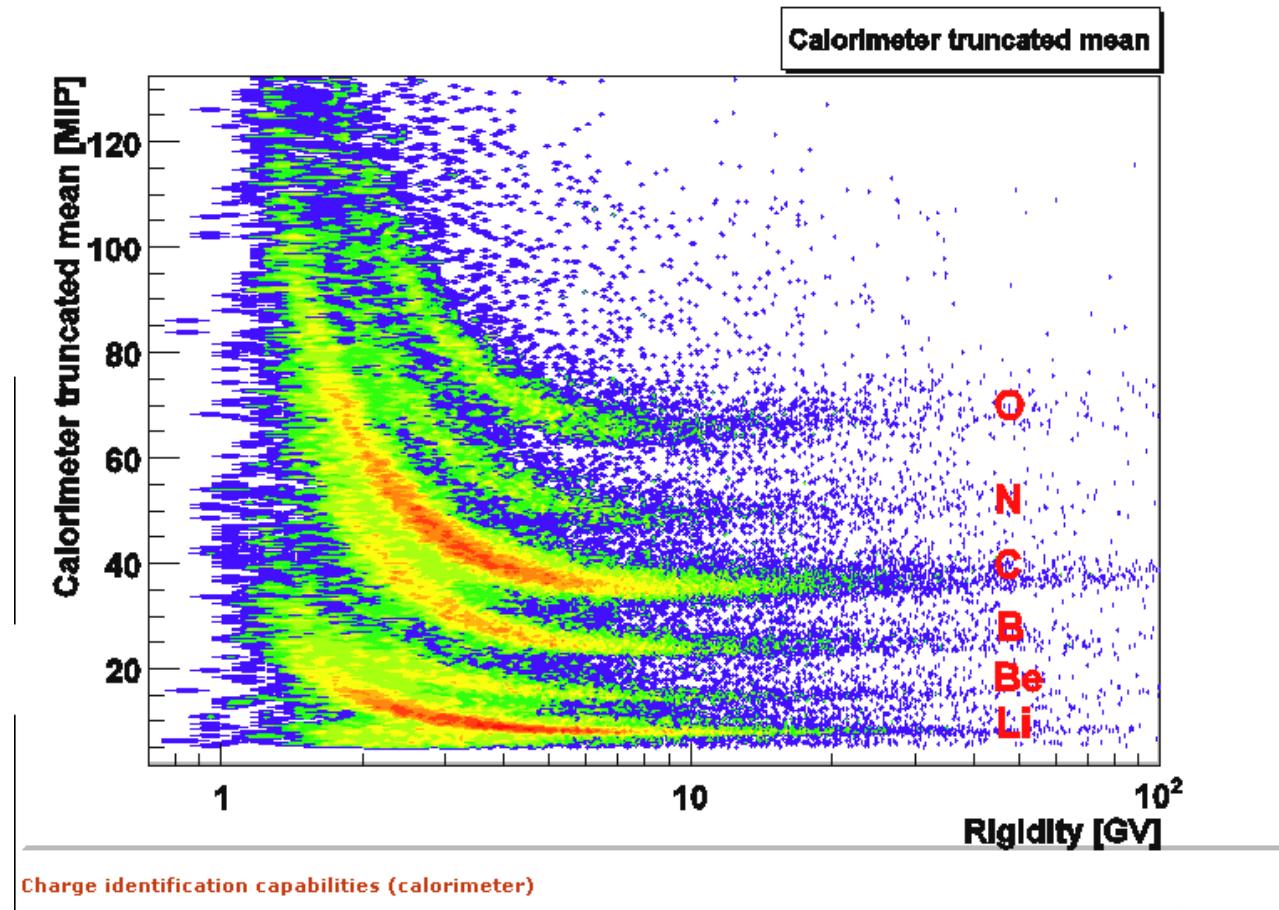
Electron energy,
 dE/dx , lepton-
hadron separation

- 44 Si-x / W / Si-y planes (380)
- 16.3 X0 / 0.6 L
- $dE/E \sim 5.5\%$ (10 - 300 GeV)
- Self trigger > 300 GeV / $600 \text{ cm}^2 \text{ sr}$

- 36 ^3He counters
- $^3\text{He}(n,p)\text{T}$; $E_p = 780$ keV
- 1 cm thick poly + Cd moderator
- 200 μ s collection

PAMELA

Ionisation in the
calorimeter ®
Z identification of
nuclei

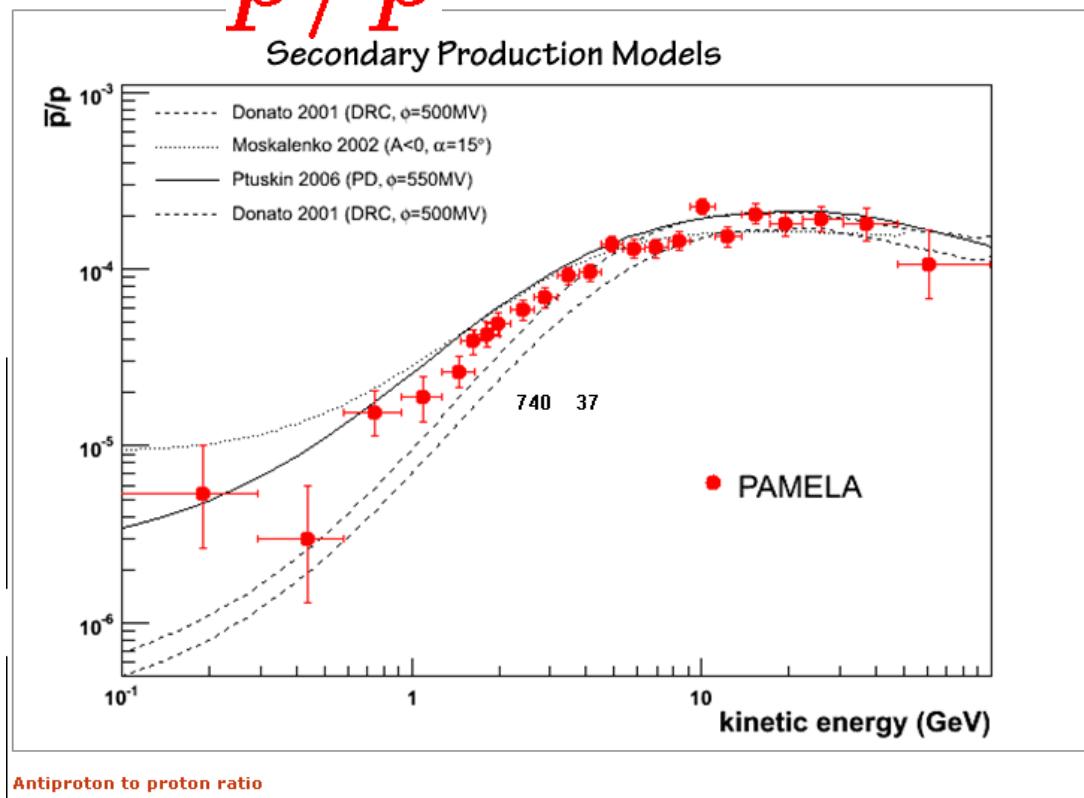


PAMELA

Antimatter ?

$$\bar{p}/p$$

Secondary Production Models



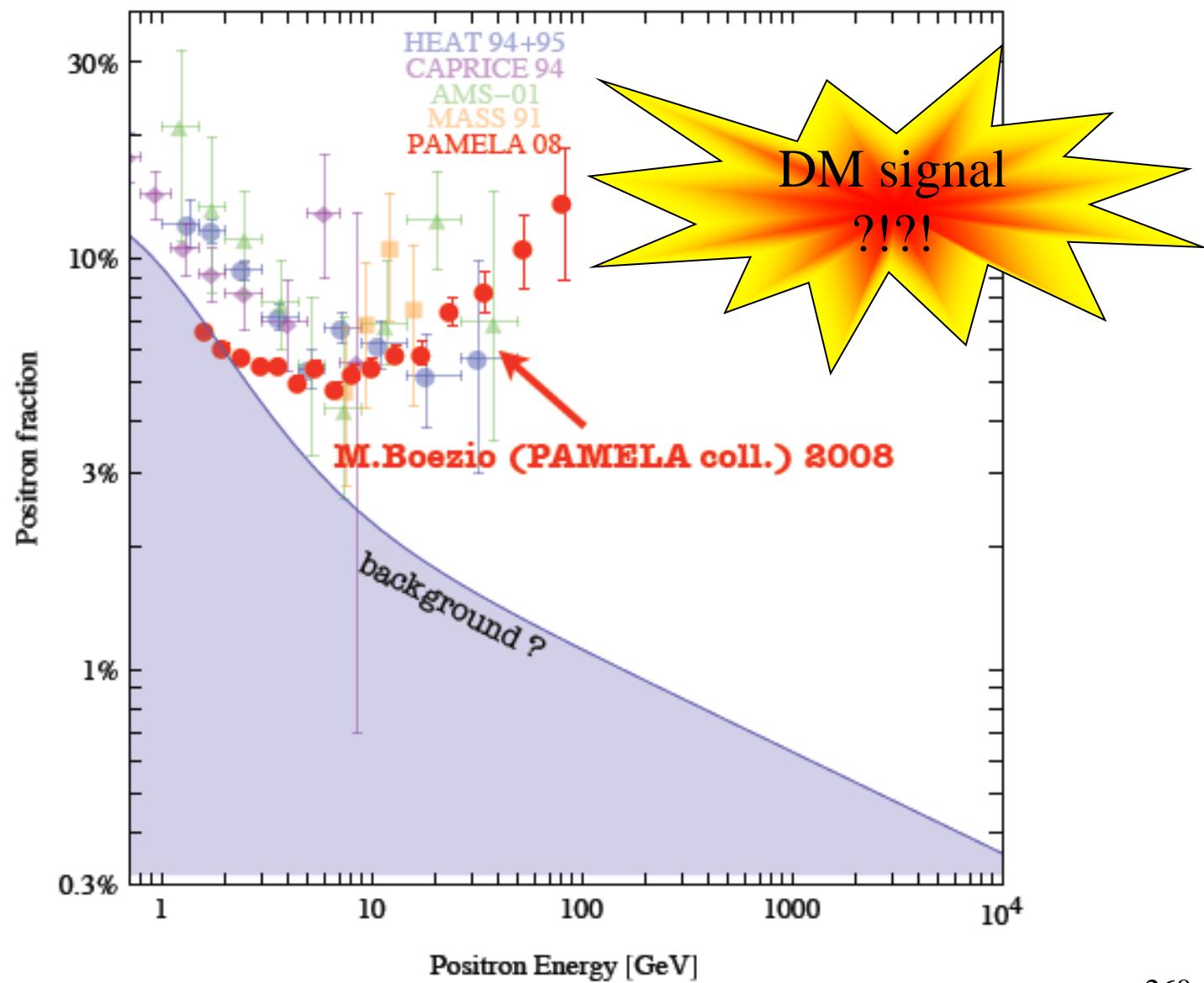
b-ph 0810.4994

Antimatter ?

$$\frac{e^+}{e^- + e^+}$$

strong excess
above 10 GeV
Important flux
confirmed by AMS-2

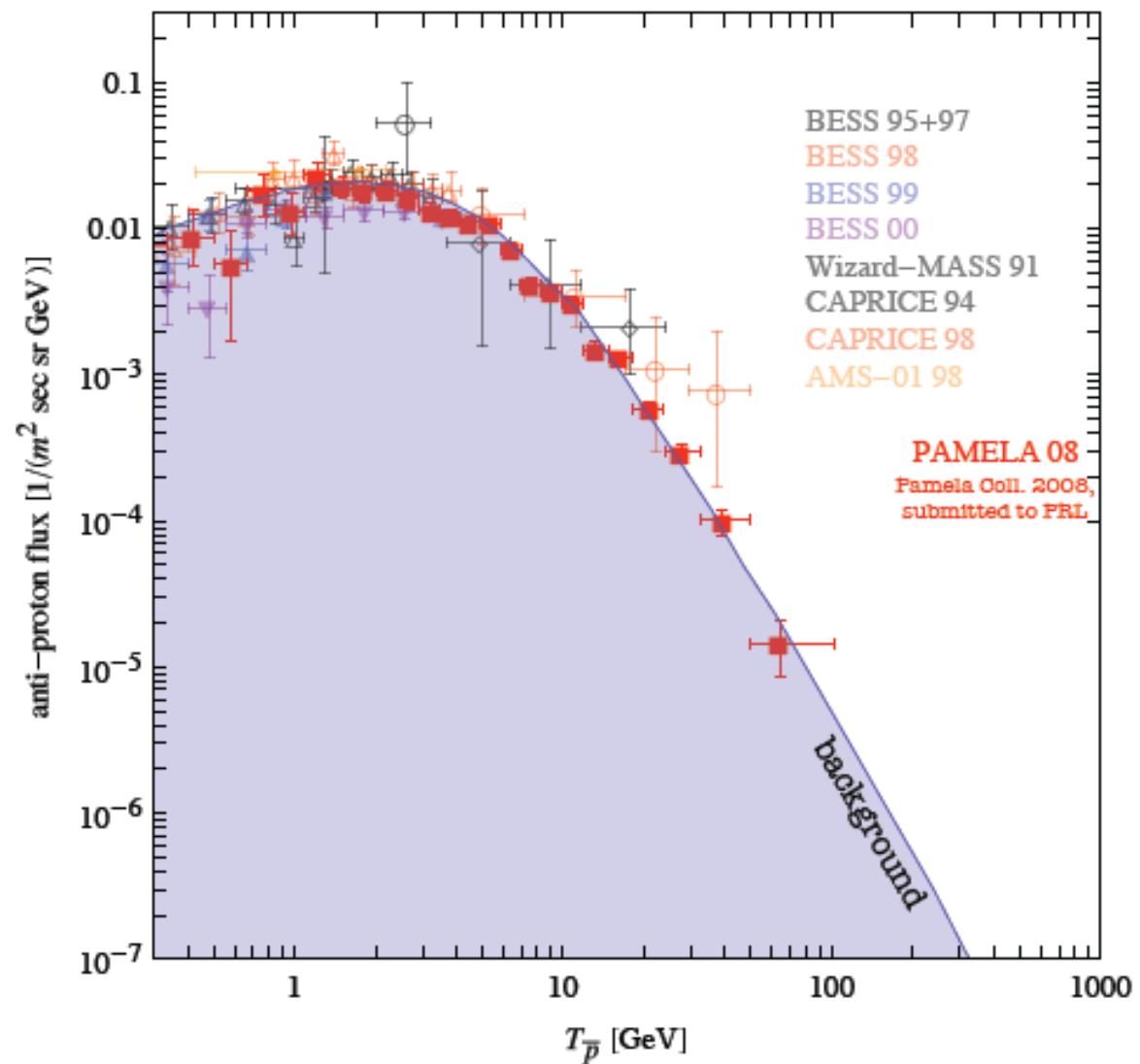
PAMELA



PAMELA

Antimatter ?

\bar{p}/p
Cistant
with "background"

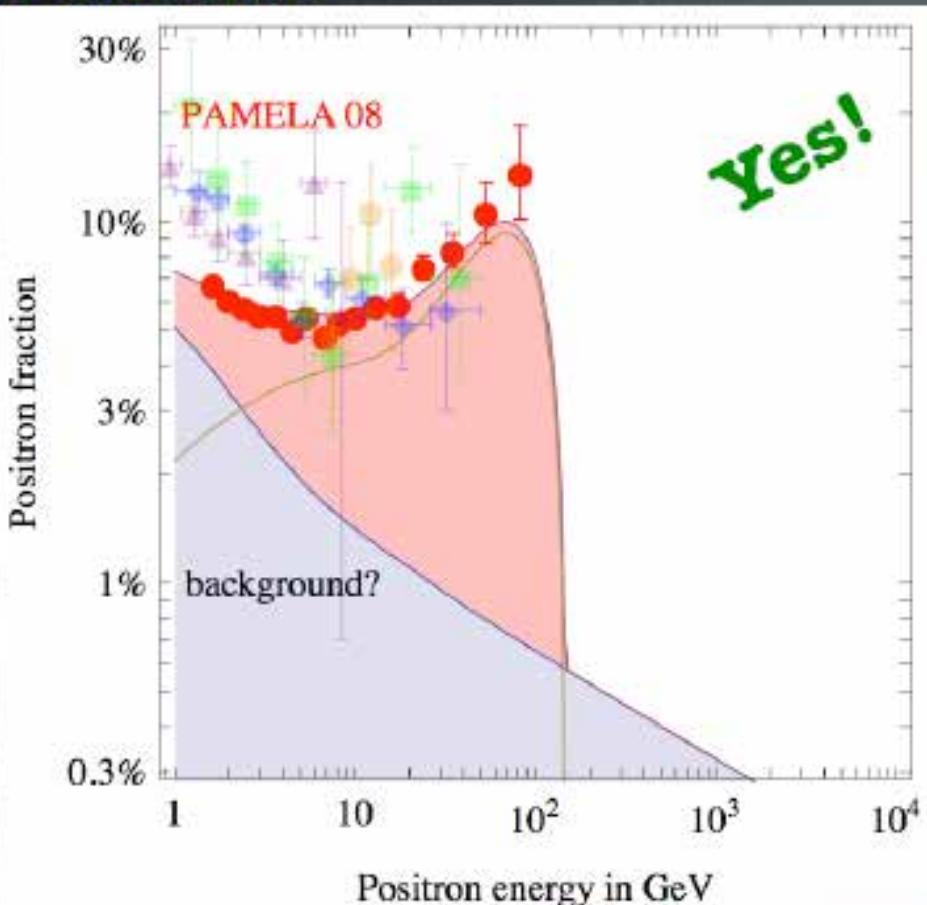


Results

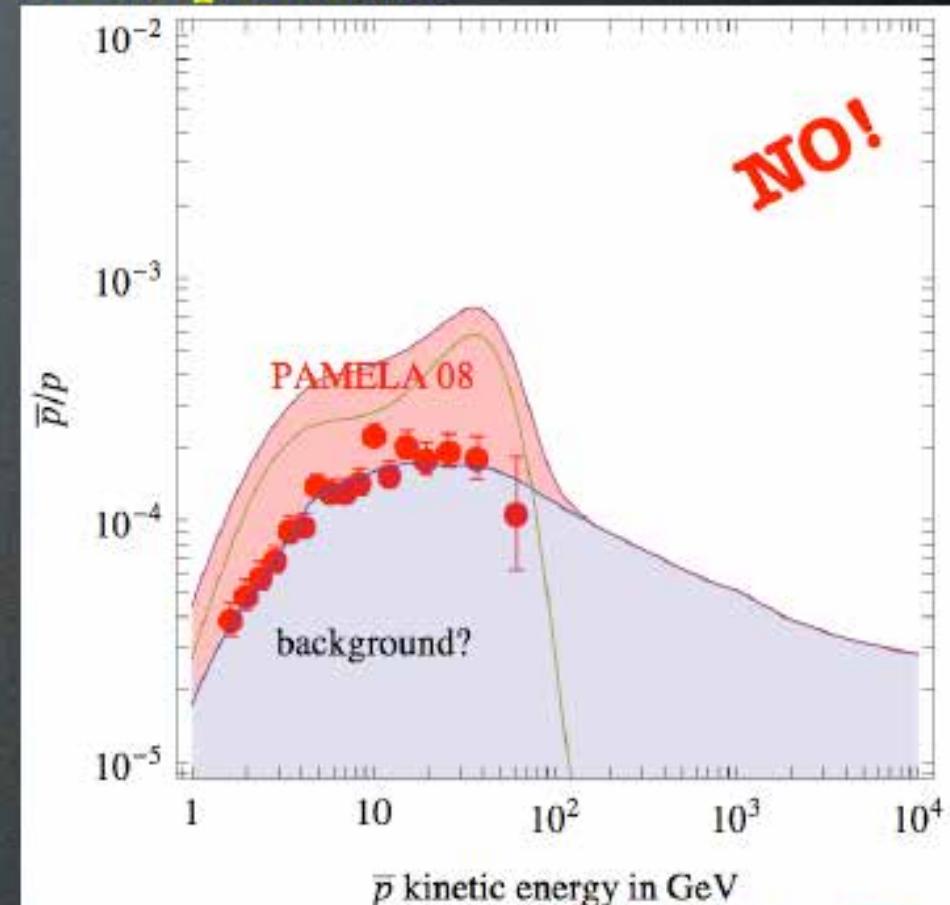
Which DM spectra can fit the data?

E.g. a DM with: -mass $M_{\text{DM}} = 150 \text{ GeV}$
 -annihilation $\text{DM DM} \rightarrow W^+W^-$
 (a possible SuperSymmetric candidate: wino)

Positrons:



Anti-protons:



[insisting on Winos]

Results

M.Cirelli, seminar @LPSC
December 2009

Which DM spectra can fit the data?

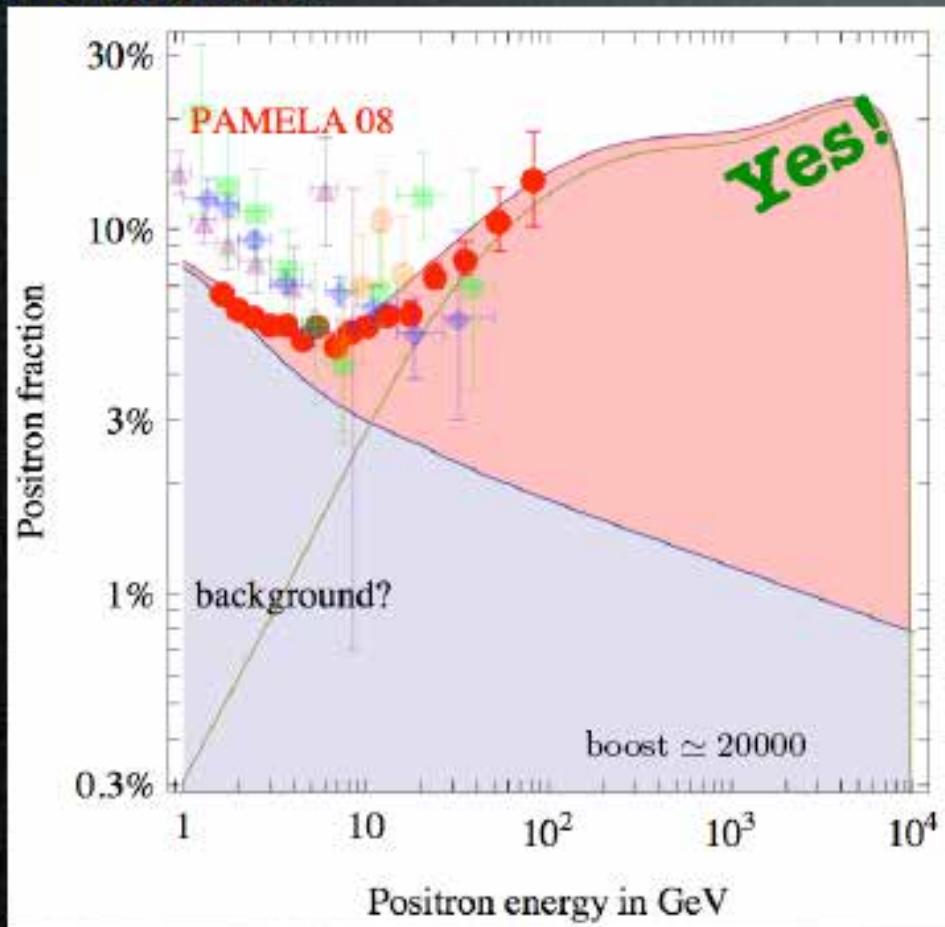
E.g. a DM with: -mass $M_{\text{DM}} = 10 \text{ TeV}$

-annihilation $\text{DM DM} \rightarrow W^+W^-$

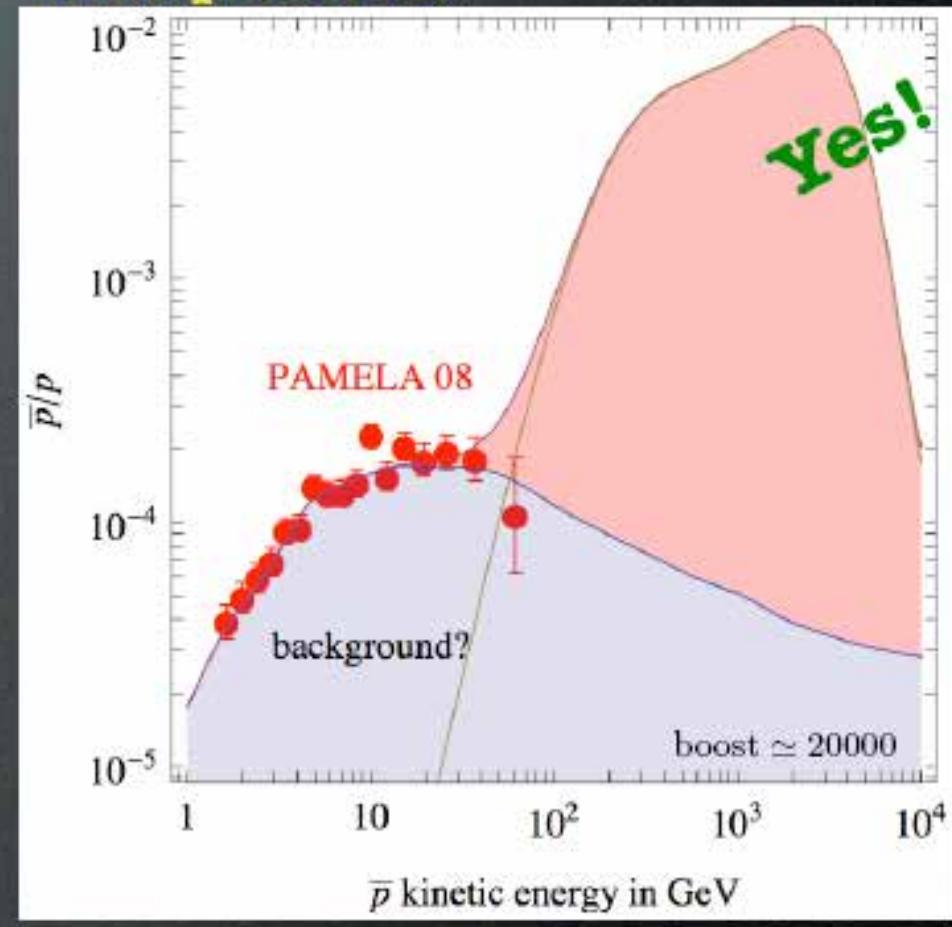
but...: -cross sec $\sigma_{\text{ann}}v = 6 \cdot 10^{-22} \text{ cm}^3/\text{sec}$

Mmm...

Positrons:



Anti-protons:



Background "measurements"

e^+ background:
mainly spallation of CR on IS gas.

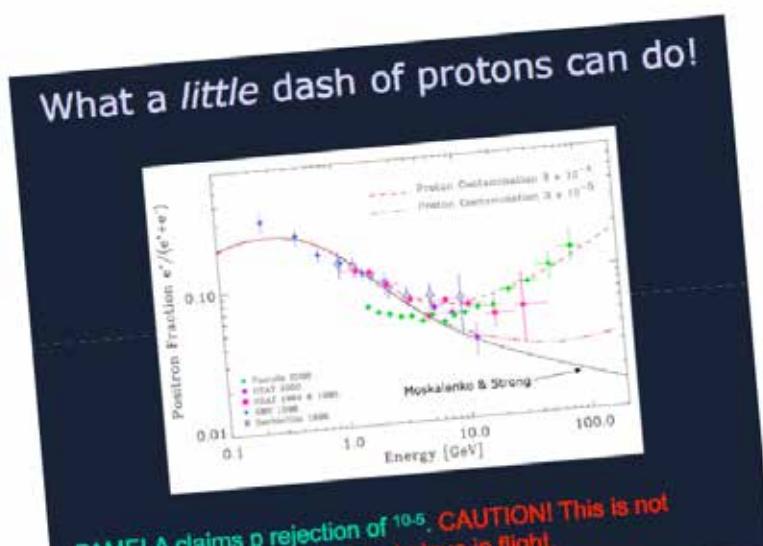
Baltz, Edsjo 1999

Moskalenko, Strong 1998

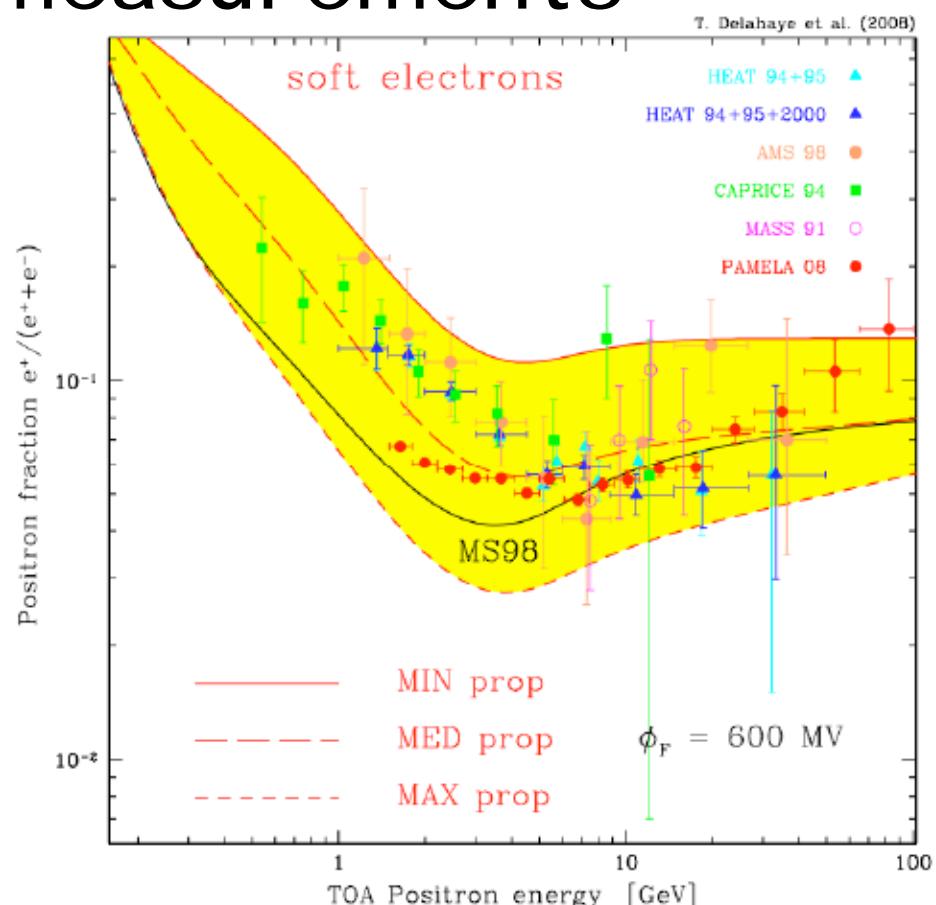
Delahaye et al., 0809.5268

P.Salati, Cargese 2007

Control of instrumental effects ?

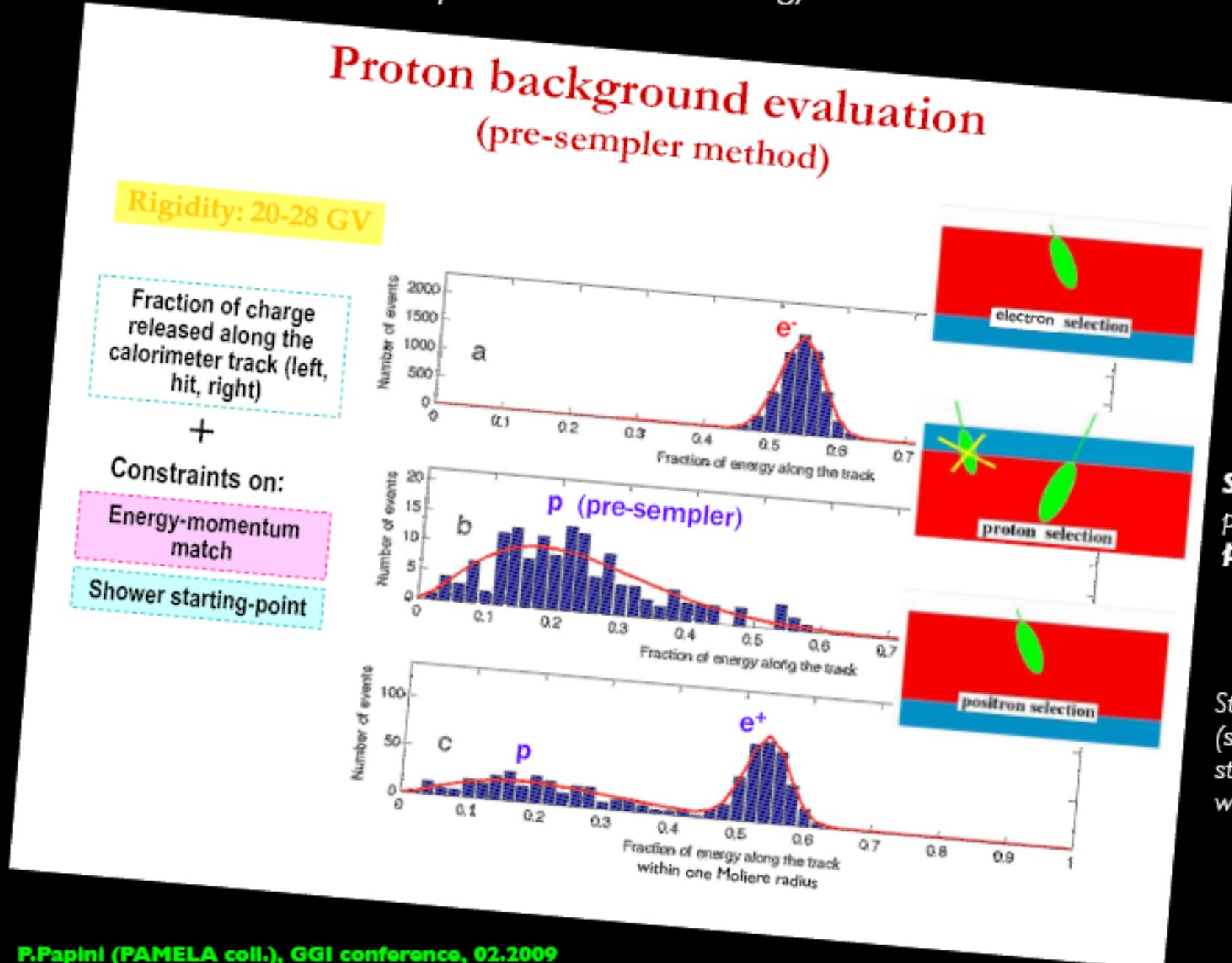


PAMELA claims p rejection of 10^{-5} . CAUTION! This is not verified using independent technique in flight.



“PAMELA did do in-flight checks of the p rejection rate”

Method: in the calorimeter, leptons leave all their energy and on the top; protons leave little energy and in the bottom.



A nearby Pulsar ?

Mature pulsars: ($0.01 < T < 1$ Myr)

e^\pm are confined to the pulsar wind nebula until it merges with the ISM. Merger process is fast so that pulsars can be treated as burst-like sources of e^\pm

Contribution could be from a few local pulsars

- Geminga: $d = 160$ pc, $T = 0.37$ Myr
- Monogem: $d = 290$ pc, $T = 0.11$ Myr

but pair conversion efficiency needs to be high (30 - 40%)

Contribution could be from a large number of pulsars, distant and local, with an assumed continuum distribution and injection spectrum

$$\frac{dN_{e^\pm}}{dE} \propto E^{-1.5} e^{-E/E_p}$$

A precision, multipurpose spectrometer up to TeV

TRD

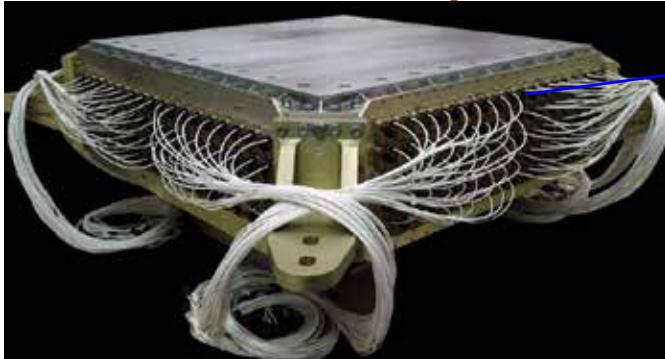
Identify e^+ , e^-



Silicon Tracker
 Z, P



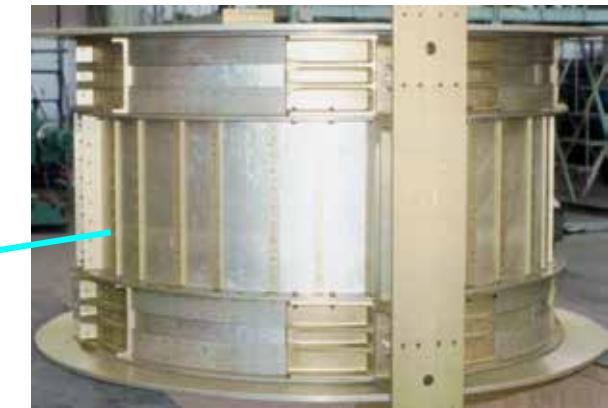
ECAL
 E of e^+ , e^- , γ



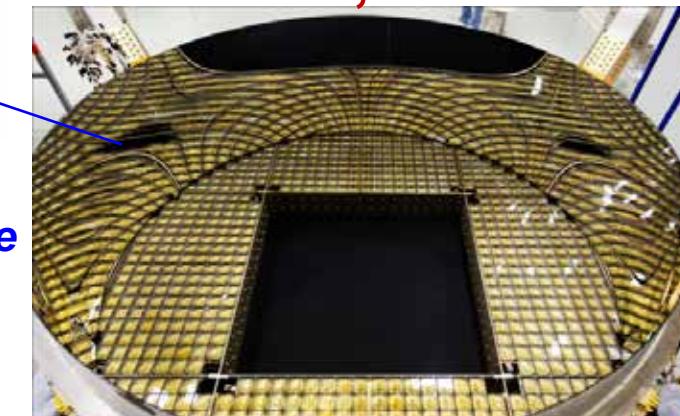
TOF
 Z, E



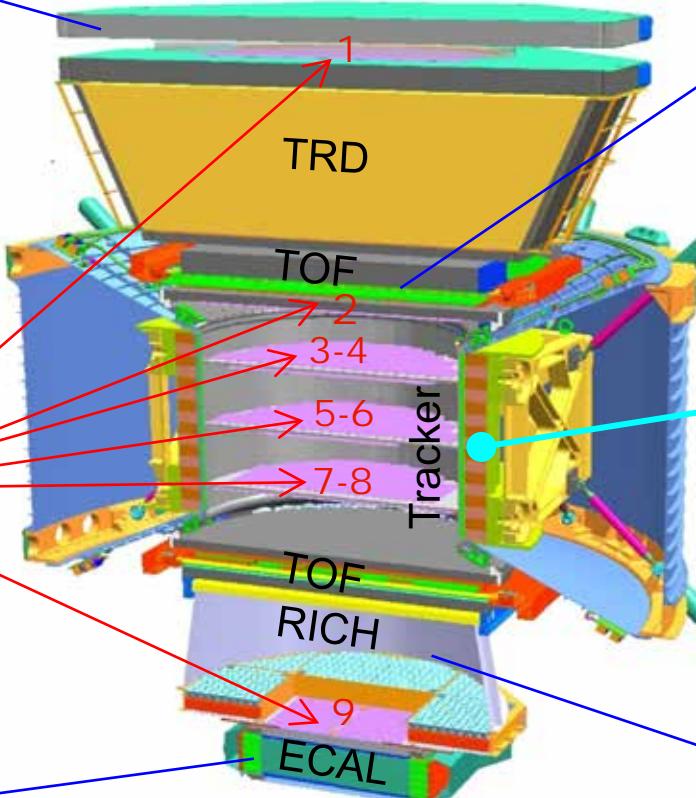
Magnet
 $\pm Z$



RICH
 Z, E

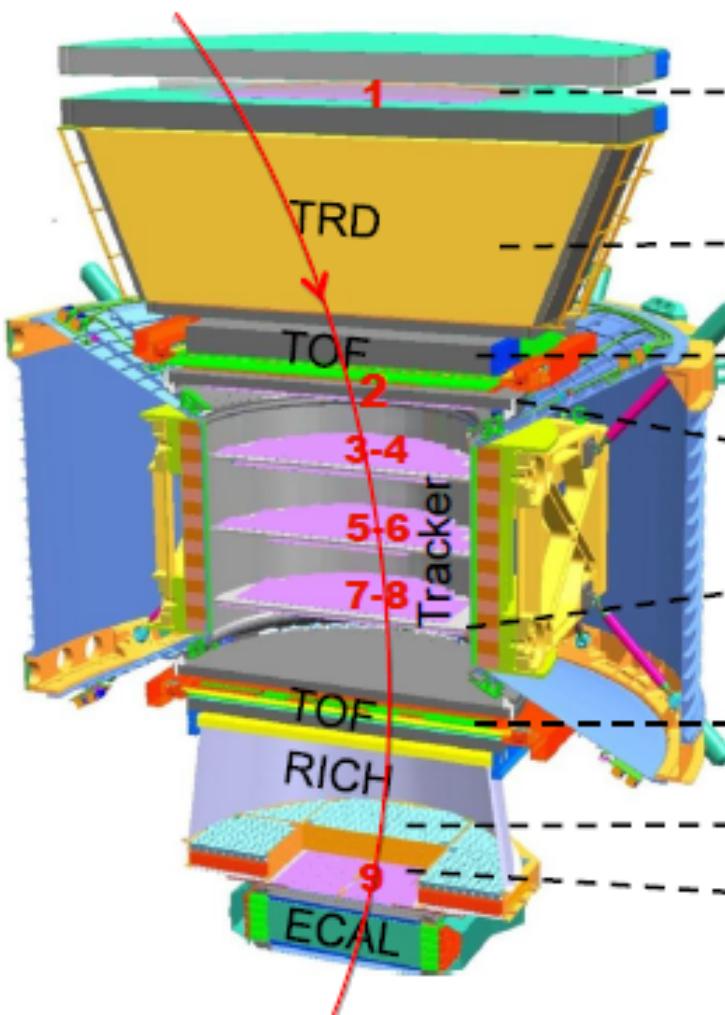


**Z, P are measured independently by the
Tracker, RICH, TOF and ECAL**



AMS charge identification

AMS: Multiple Independent Measurements
of the Charge ($|Z|$)

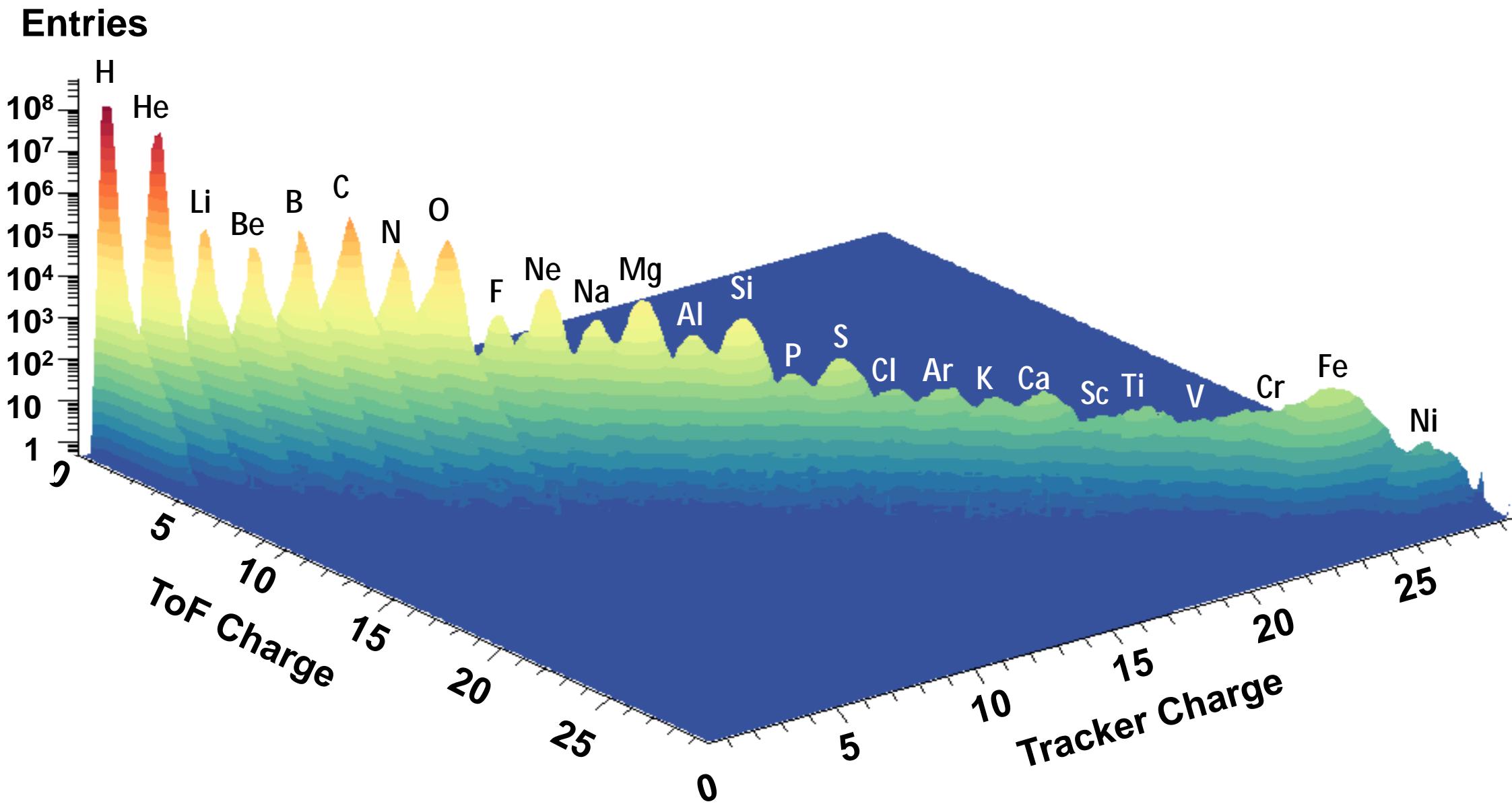


Carbon ($Z=6$)	ΔZ (cu)
1. Tracker Plane 1	0.30
2. TRD	0.33
3. Upper TOF (1 counter)	0.16
4. Tracker Planes 2-8	0.12
5. Lower TOF (1 counter)	0.16
6. RICH	0.32
7. Tracker Plane 9	0.30

Full coverage of anti-matter & CR physics

	e^-	P	He,Li,Be,..Fe	g	e^+	\bar{P}, \bar{D}	\bar{He}, \bar{C}
TRD							
TOF							
Tracker							
RICH							
ECAL							
Physics example	Cosmic Ray Physics				Dark matter		Antimatter

AMS Nuclei Measurement on ISS



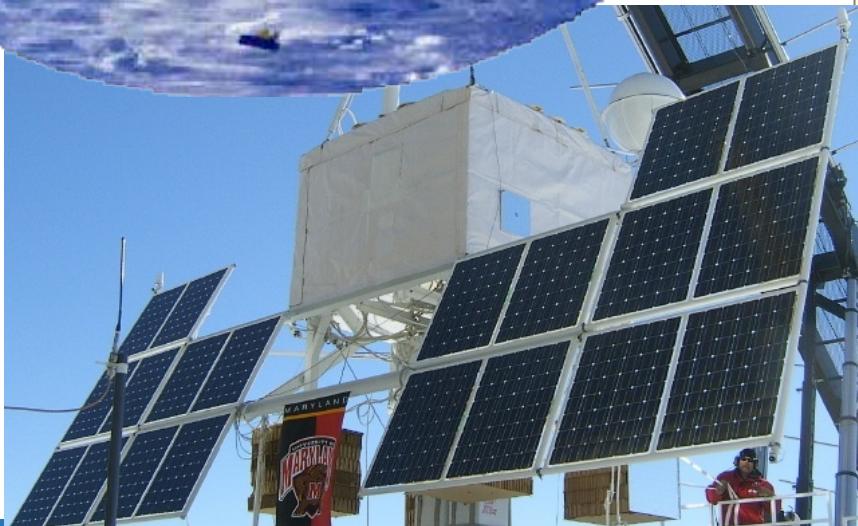
Charged particles (1 TeV → few 100 TeV)

Balloon Experiments

- Very high altitude flights: 40 km ↵ 3,9 g cm⁻²
- Load: up to ≈ 270 kg and 10 m³
- Many former flights:
JACEE (USA + Japon) ; RUNJOB (Russie + Japon) etc.
using emulsion chambers !
- Recently: Ultra Long Duration Balloon (ULDB) Flights
60-100 jours (NASA, Antarctic) → the CREAM experiment
(Cosmic Ray Energetics and Mass)



CREAM



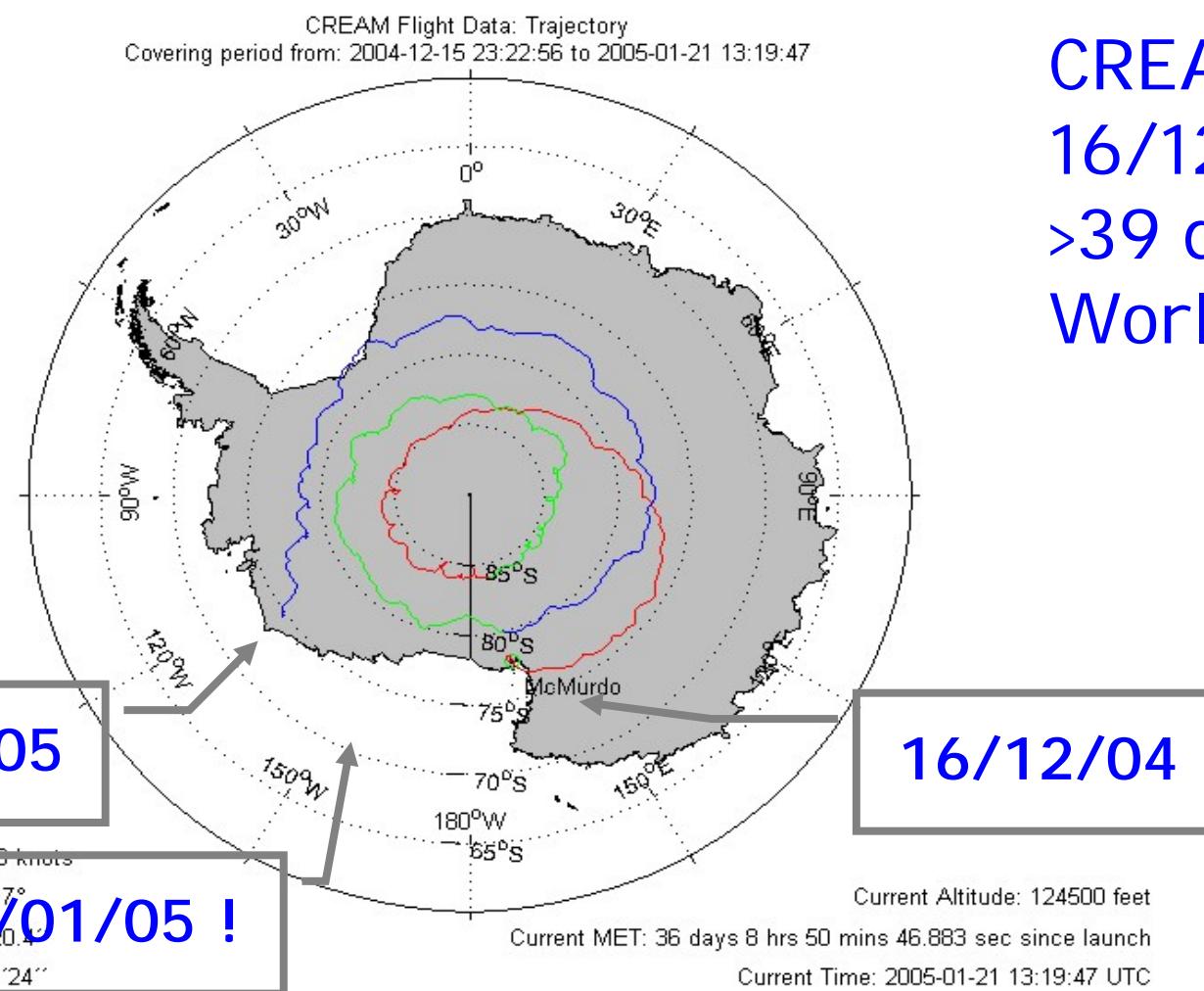
Ultra Long Duration Balloon

ULDB Proj., Adv.Sp.Res33,1633(2004) :
NASA project to develop

- Flight of < 100 days
- Payload . 2 tons
- Alt 33000 meter
- CREAM n° 1 : 2006 (2005/LDB)



CREAM 2004

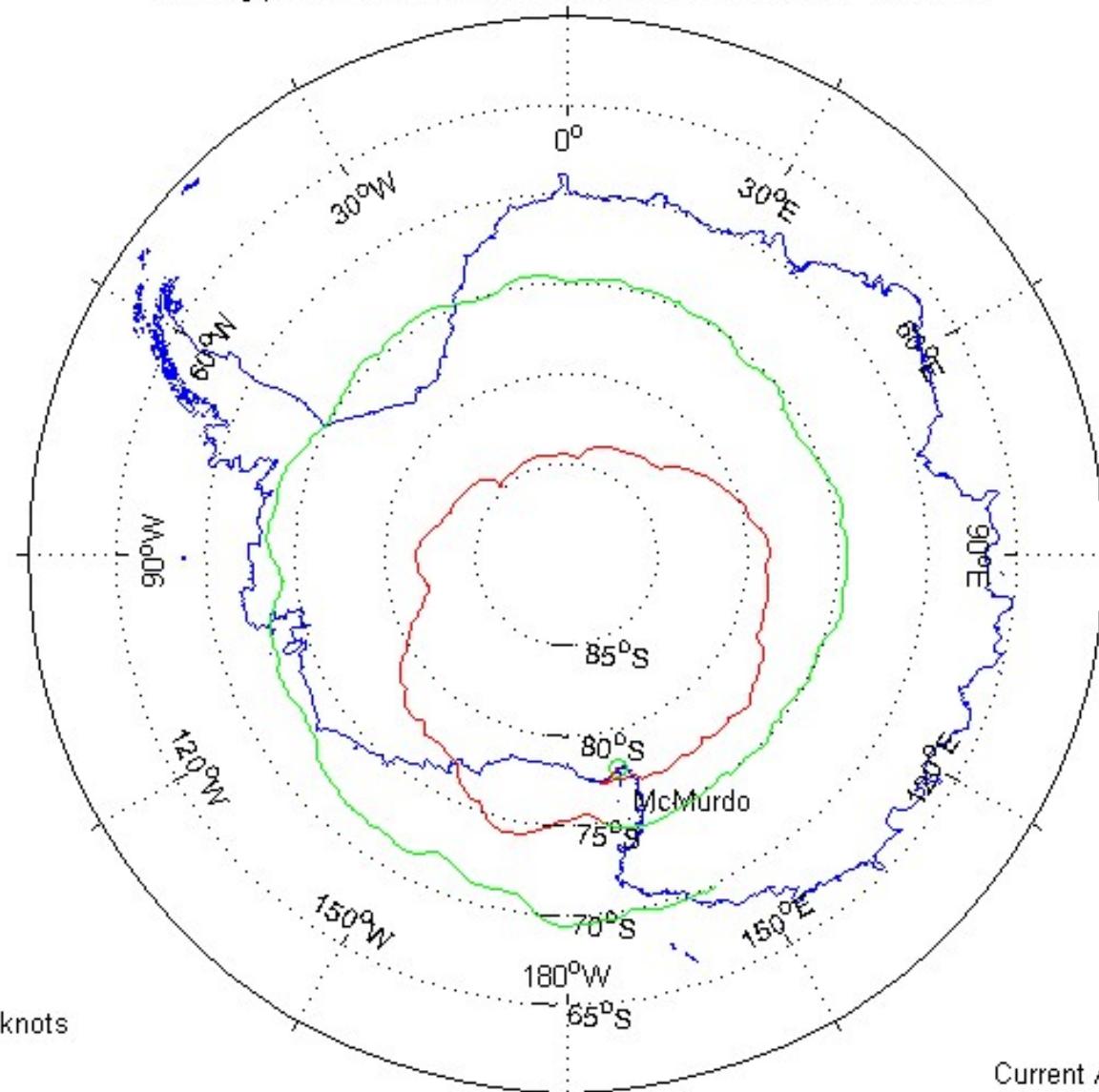


CREAM I:
16/12/04 - **/01/05
>39 days flight
World Record

CREAM 2008

CREAM Flight Data: Trajectory

Covering period from: 2008-12-16 21:08:56 to 2009-01-07 11:29:46



Current Speed: 0.01 knots

Current Course: 0°

Current Lat: -69°48'13.2"

Current Lon: 155°51'54"

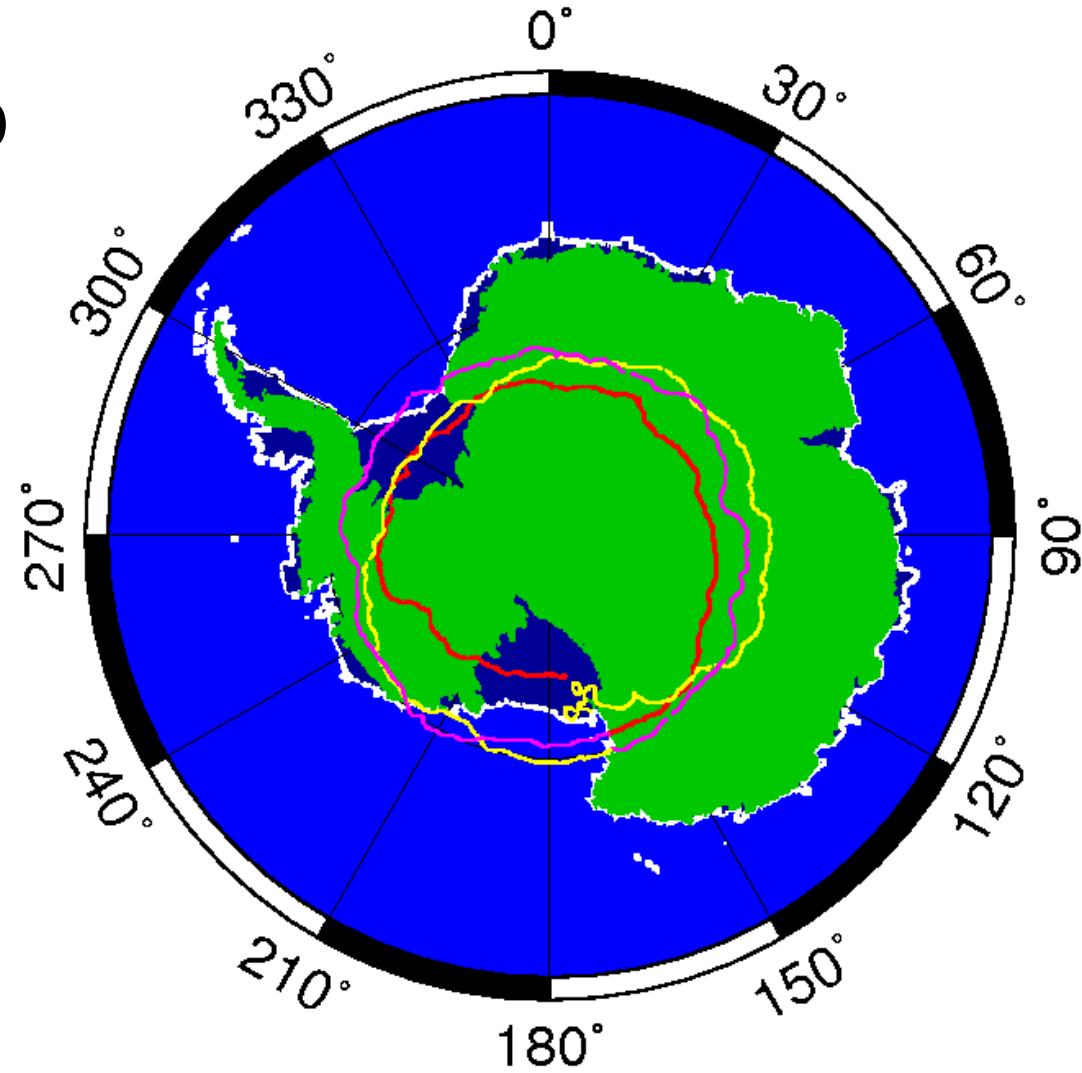
Current Altitude: 4720.1444 feet

Current MET: 19 days 13 hrs 35 mins 0.65 sec since launch

Current Time: 2009-01-07 11:29:46 UTC

1 décembre 2009
au 8 janvier 2010

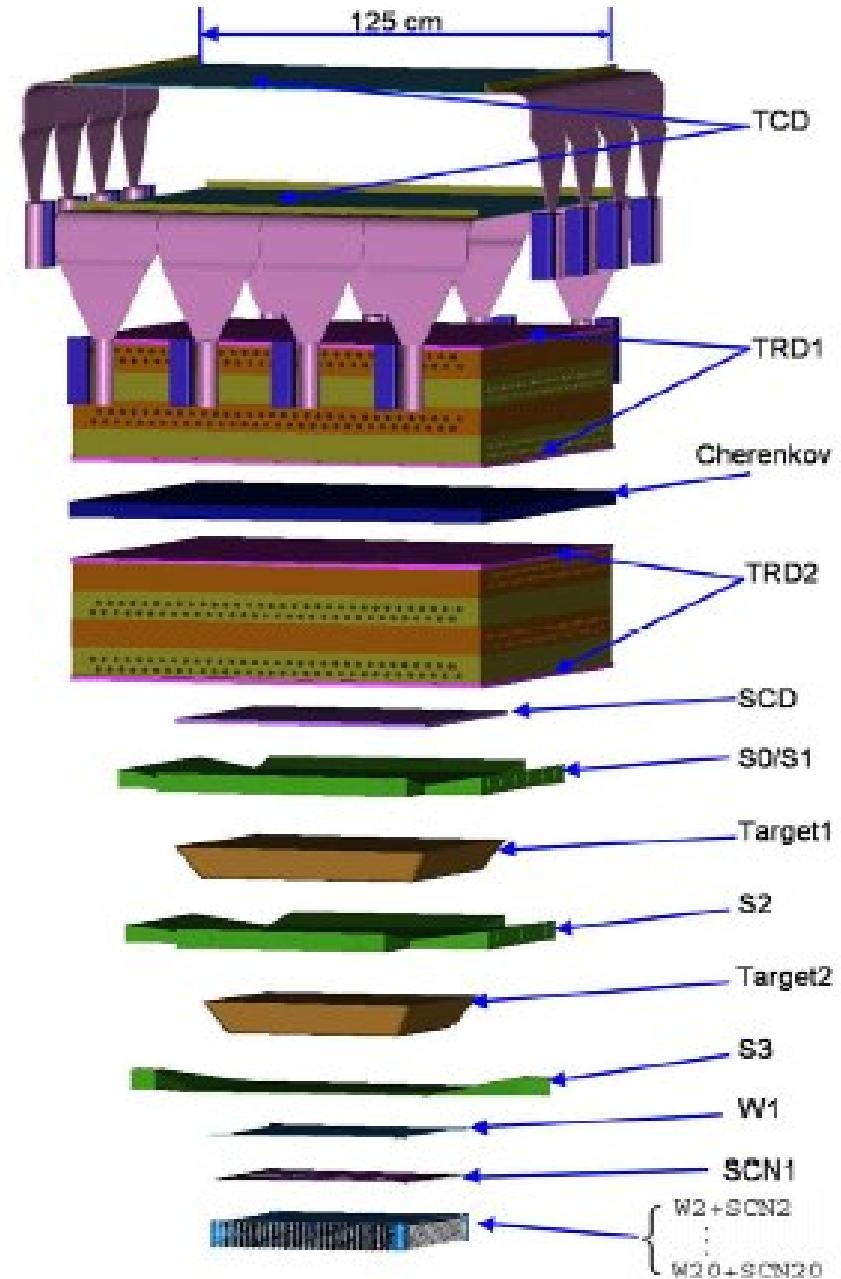
CREAM 2009



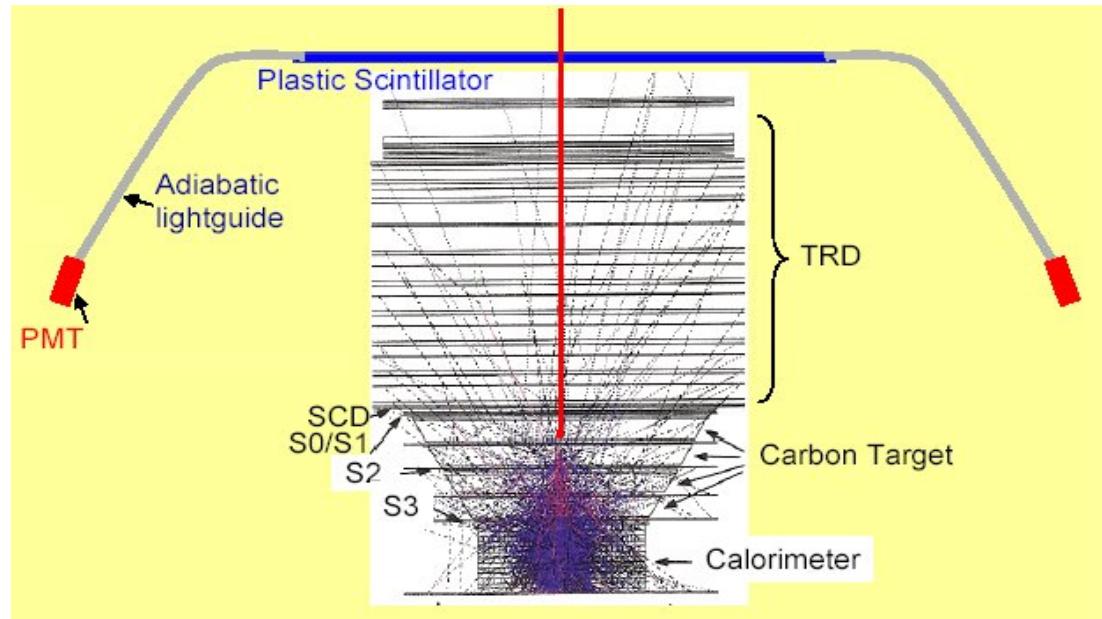
CREAM

Cosmic Ray Energetics and Mass

- **Objectives :**
CR composition and spectrum of the different elements
(from TeV to ~500 TeV)
- **Acceptance :** $2,2 \text{ m}^2 \text{ sr}$
- **Energy measurement:**
 - Calorimeter $20 X_0$ (W + scint. fibres)
 - Transition Radiation Detector
- **Identification :**
 - TRD
 - Cherenkov detector "CHERCAM"
similar to AMS-2

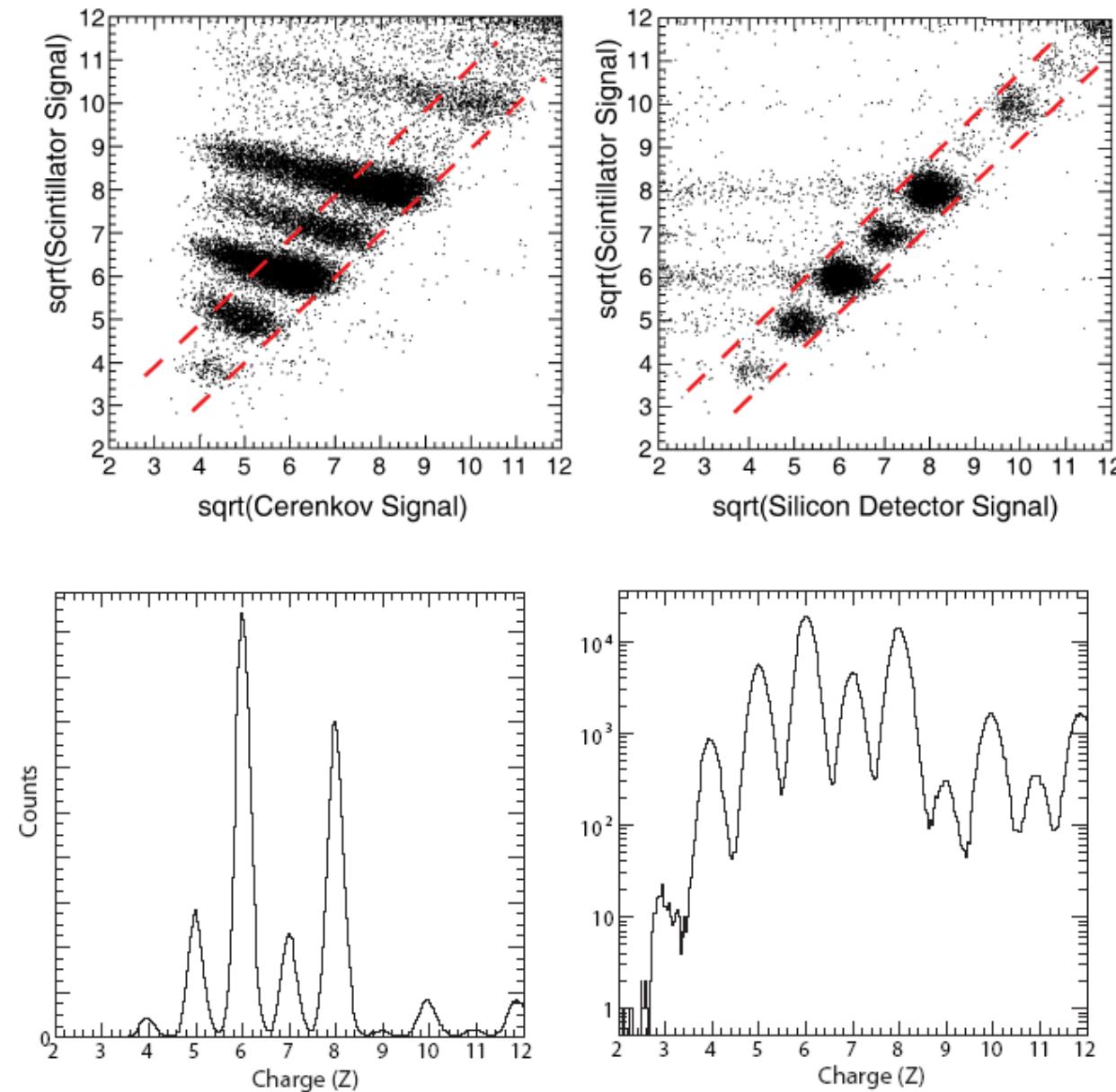


CREAM experiment

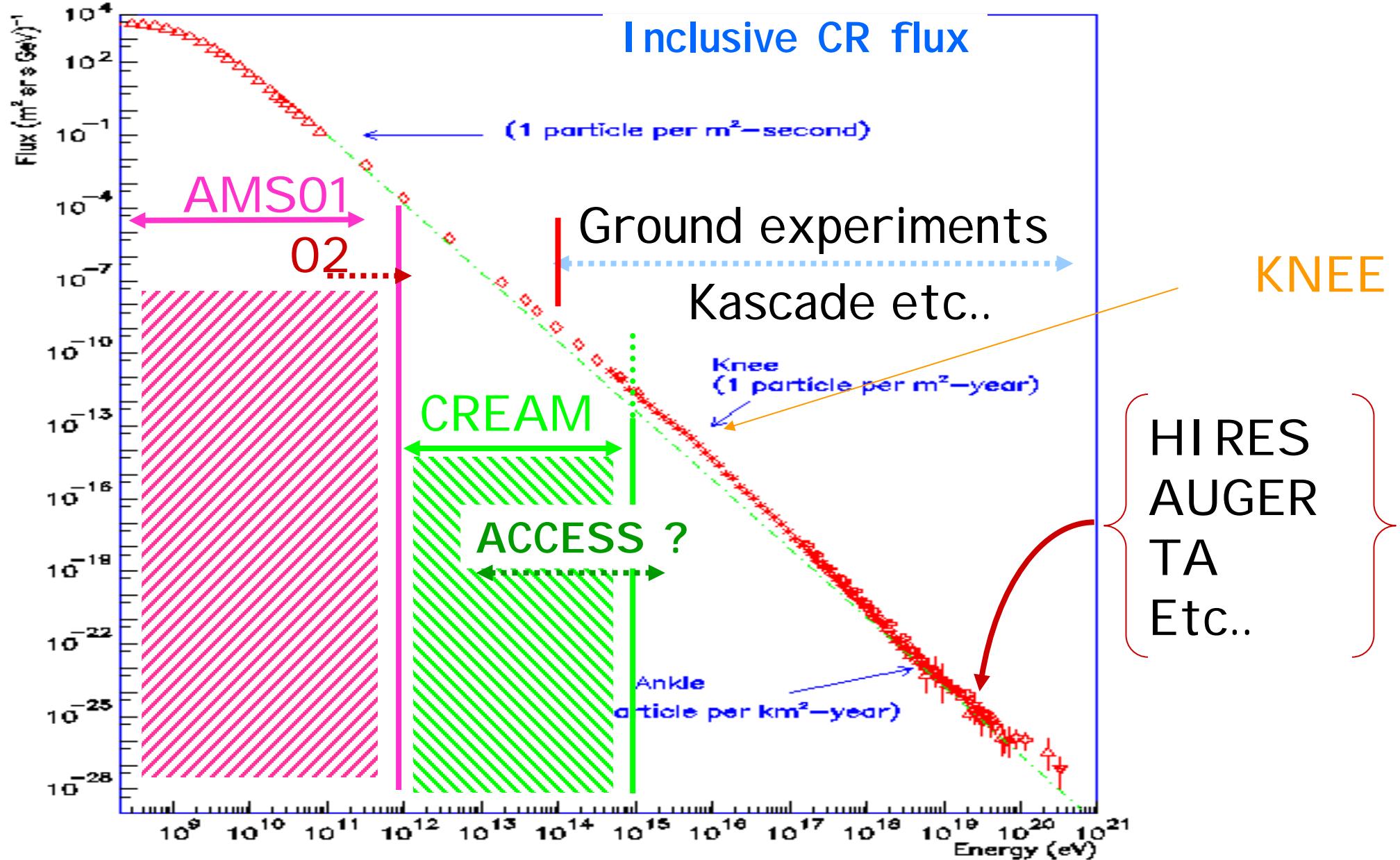


- At TeV energies, the interaction of CR in the calorimeter induces many backscattered secondary particles that one have to veto.
- The "CHERCAM" cherenkov solves this problem by measuring accurately the time of any through going particle as well as achieving a precise charge measurement ($\pm 0,3$ e)

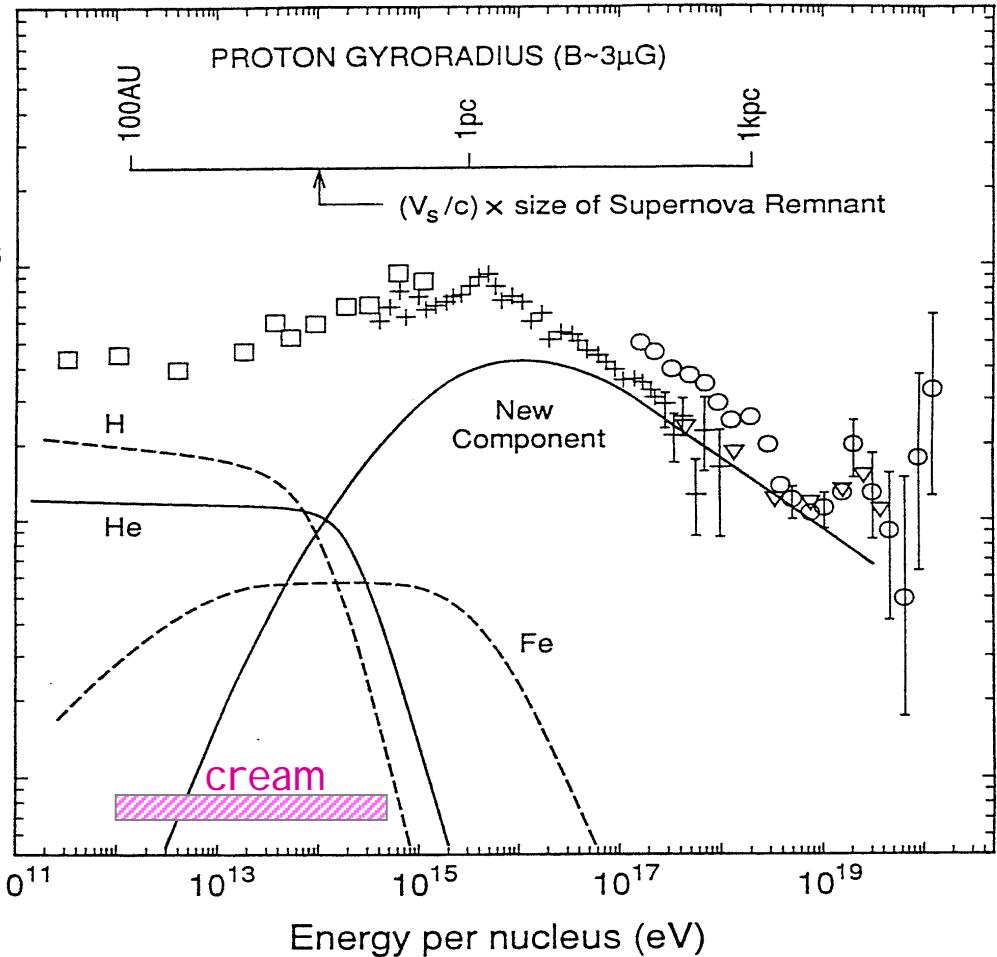
CREAM experiment



Contexte expérimental



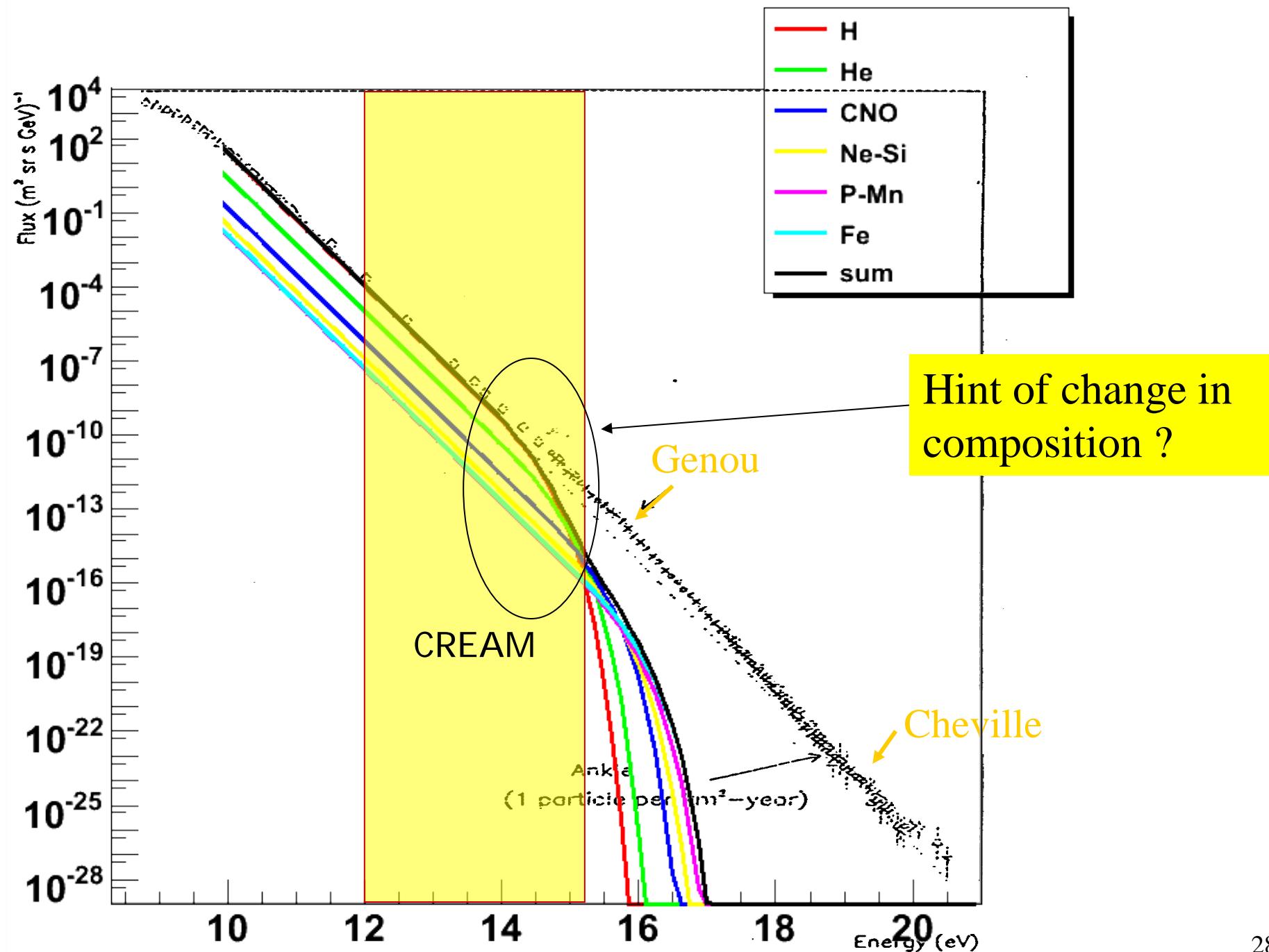
The knee



- Is the knee due to:
 - Acceleration mechanisms or to changes :
 - in propagation?
 - in CR sources?
 - in interaction properties (threshold) ?
- A diffuse SNR shock acceleration with E_{\max} implies a change in composition around $\sim 10^{14}$ eV.

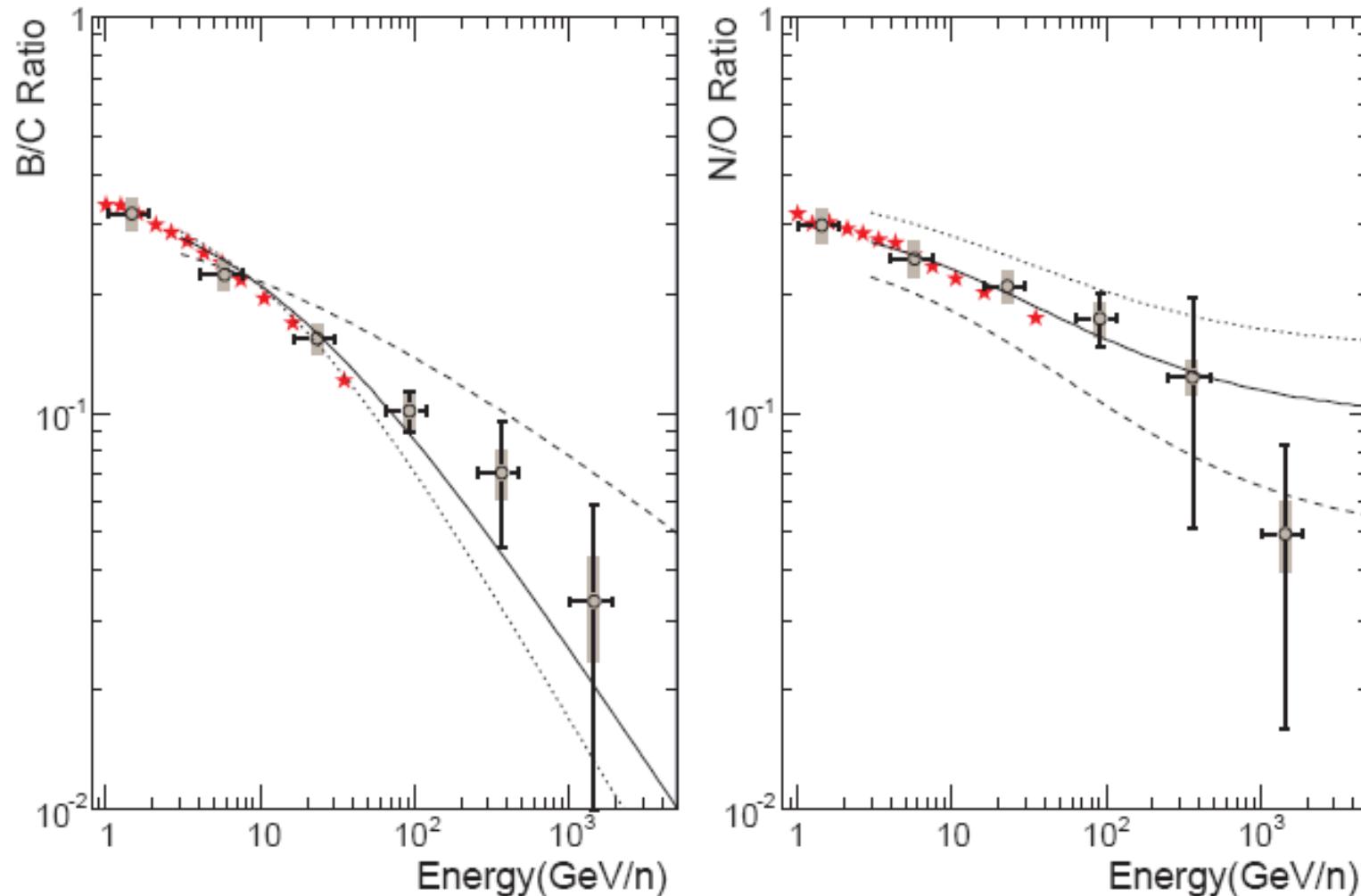
SNR energy limit: $E_{\max} \sim Z \cdot 10^{14}$ eV

CR SM with Emax / Z



Galactic CR propagation with CREAM

Secondary/Primary ratio



Compatible with propagation models with escape terme in $E^{0.6}$

AIR SHOWERS DEVELOPMENT MODELS

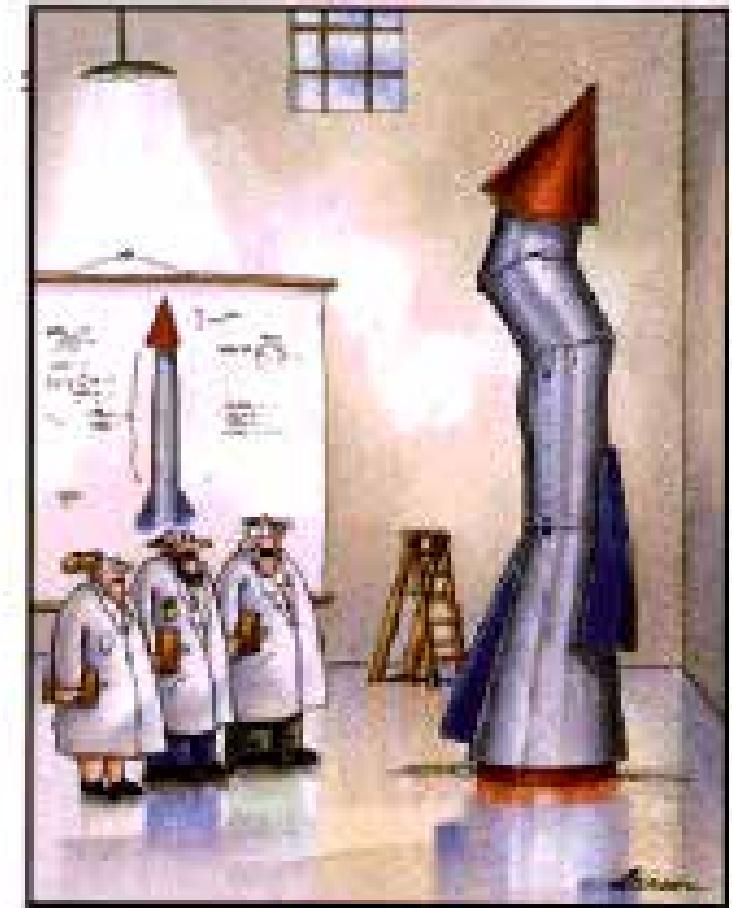
A peek above the knee !

To measure the inclusive spectrum at the knee, one needs a 10m^2 exposed during 10 years !

The realistic experimental limits are:

- For satellites $\sim 1\text{m}^2$ (sr) during \sim few years
- For balloons, $\sim 10\text{m}^2$ (sr) during ~ 30 jours

$$\textcircled{R} \quad E < 10^{15} \text{ eV}$$



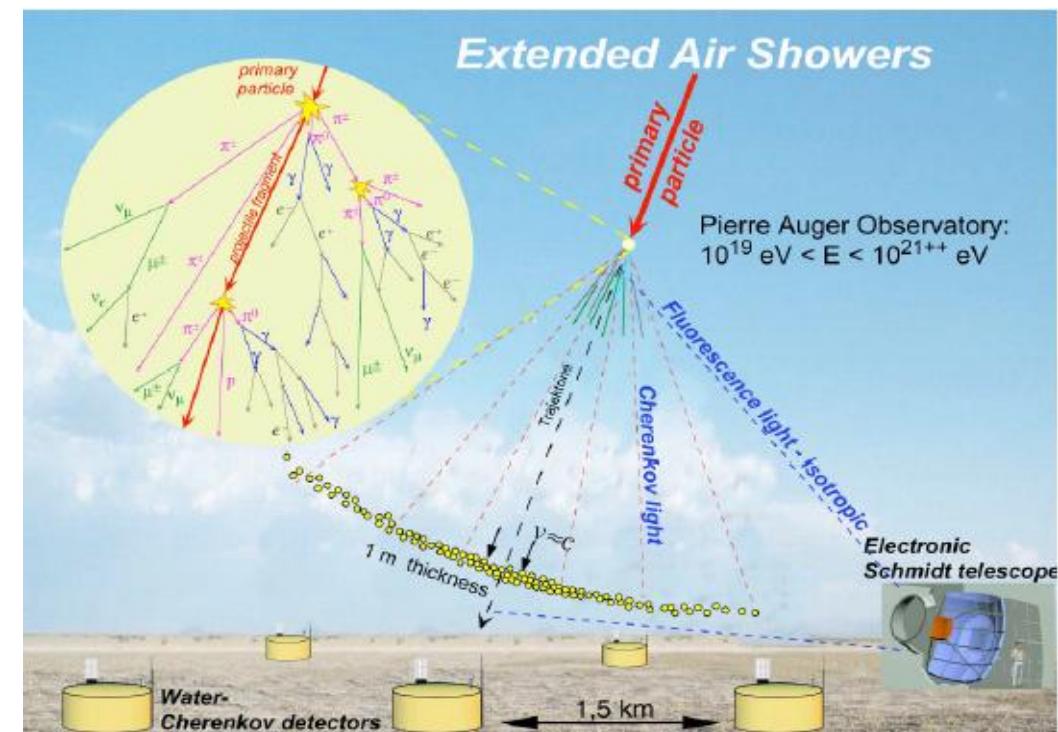
It's time we face reality my friends,
we should keep to ground detectors !

Extensive Air Showers: the phenomenon and the observables

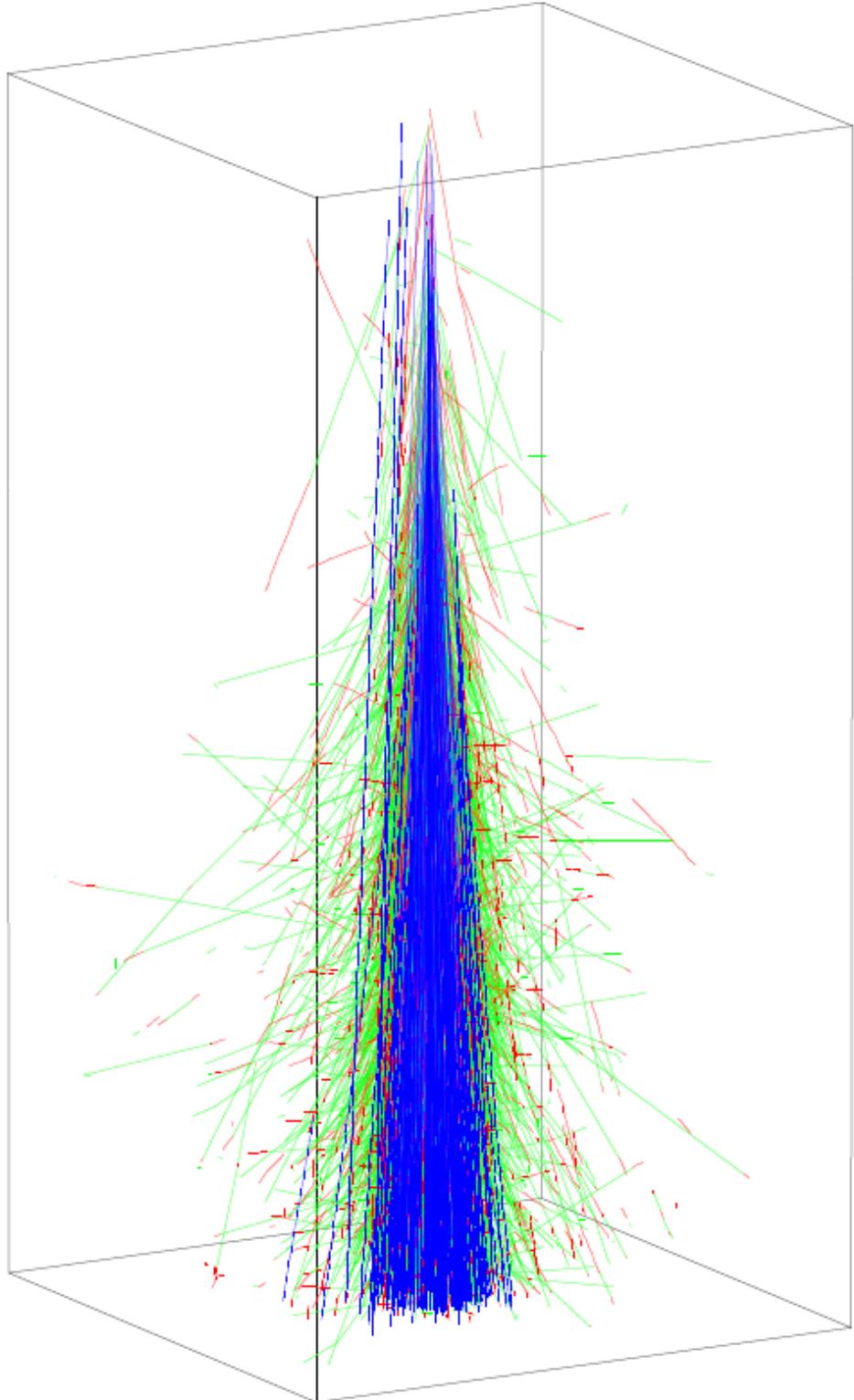
- The large shower of secondary particles induced by the interaction of a primary CR in the upper atmosphere can be detected on an extensive area
→ large effective surfaces to fight against low flux at $E \geq 1000$ TeV
- Atmosphere used as an calorimeter
(~ 1000 g cm⁻² at sea level for a vertical shower)
- From the observables, one aims at measuring:
 - Incident direction;
 - Primary energy E_0 ;and if possible, get access to the nature of the primary particle :
 - distinction γ -hadron ;
 - distinction light nuclei (p, He) - heavy nuclei(Fe)

p or nucleus + N or O nucleus → hadronic cascade

- Hadronic component: nuclear fragments, nucleons, mesons π , K , etc.
- Electromagnetic component: induced by $\pi^0 \rightarrow \gamma\gamma$ and other radiative decays
- Muonic component: induced by decays of π^\pm and K^\pm
- Atmospheric Neutrinos issued from π^\pm K^\pm and μ^\pm decays



Primary electrons and γ induce an **electromagnetic shower** consisting mainly of secondary electrons, positrons and γ (muon poor)



Shower development

« des giboulées d'électrons »

**Rayon cosmiques
par Pierre Auger
1941 PUF**

A 10^{19} eV shower

10^{11} particles
at sea level

Photons + electrons (99%),
muons (1%)

Ground observables

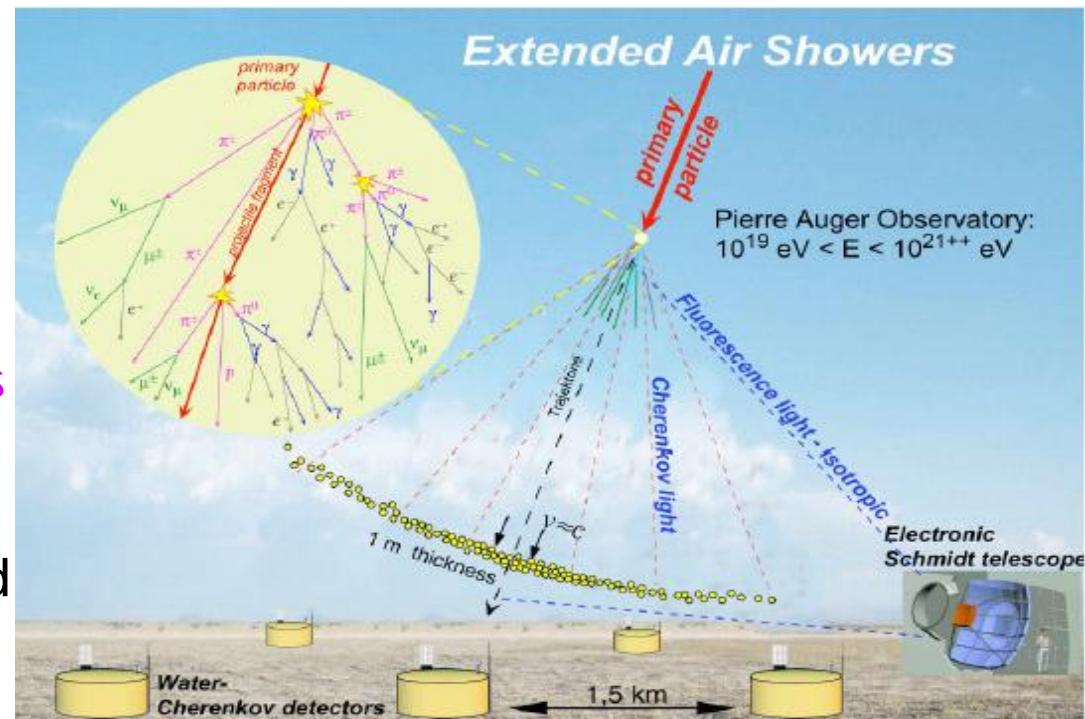
- **Secondary particles reaching ground**

As a function of the primary energy and of the altitude:

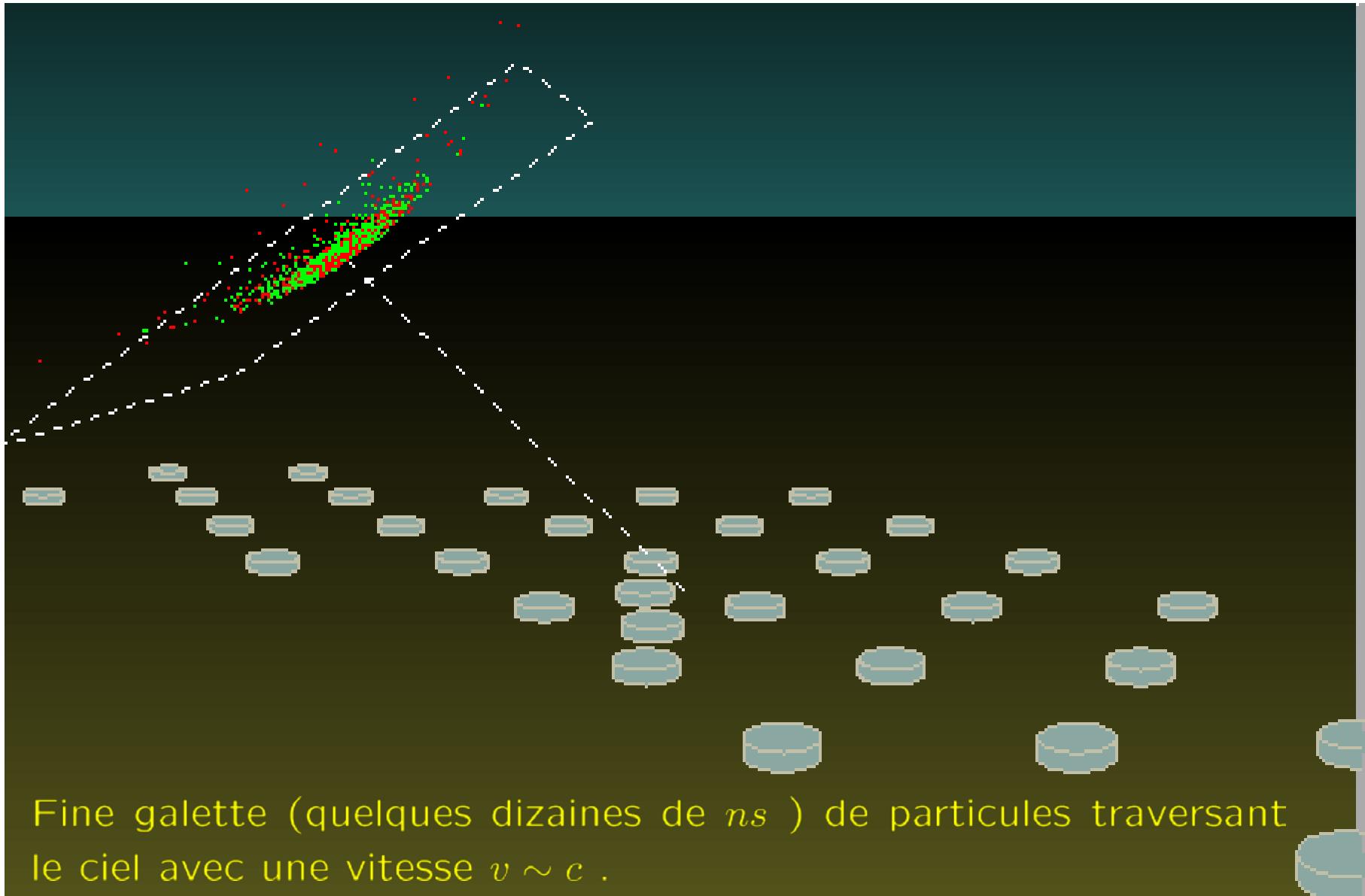
- Residual Hadrons (nuclear fragments): not numerous ($>11 \text{ } \mu\text{m}_{\text{int}}$).
 - e^\pm : the more numerous at shower development maximum.
 - μ^\pm : most reach ground and may penetrate deep underground.
 - ${}^{\circ}\text{ secondaries}$: may be detected at ground level via e^+e^- pair conversion (e.g. Cherenkov effect in water).
-
- Photons (visible, UV) emitted along the trajectories of charged particles (Cherenkov effect, N_2 fluorescence) during the shower development
→ Calorimetric 3D information !
-
- Radio emission by the shower particles in the geomagnetic field or by the induced plasma.

Temporal aspects

- A light speed moving "pancake" of charged particles.
- This front is more or less curved depending on the shower development stage.
- The front thickness (~ 10 m) induce as signal time spread in each detector.
- The arrival time differences at ground on the sampling detectors \rightarrow arrival direction ($\Delta\theta \approx 1^\circ$).
- The Cherenkov light front (forward emission) is thinner (\sim m) than the charged particle front \rightarrow well defined timing.



Time structure



Fine galette (quelques dizaines de ns) de particules traversant le ciel avec une vitesse $v \sim c$.

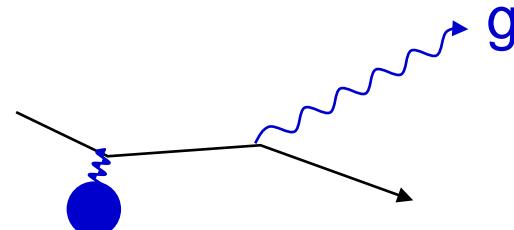
EM shower

Longitudinal development

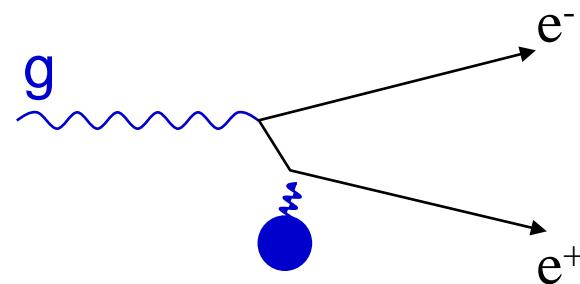
- Mean number of particles (e^+, e^- or γ) crossing a plan \perp to the shower axis after a slant depth t (in units of X_0).
- As long as the ionization losses are small wrt radiation losses (bremsstrahlung and pair prod) the number of particle increase exponentially.
- When the mean energy per particle decreases below the critical energy ($E_c \approx 84,2$ MeV in air), the number of particle decreases (shower extinction phase).
- At the transition between the two phases, (maximal development), the mean energy is equal to the critical energy.

Radiative processes ($E > E_c$)

Bremsstrahlung :



Pairs production :

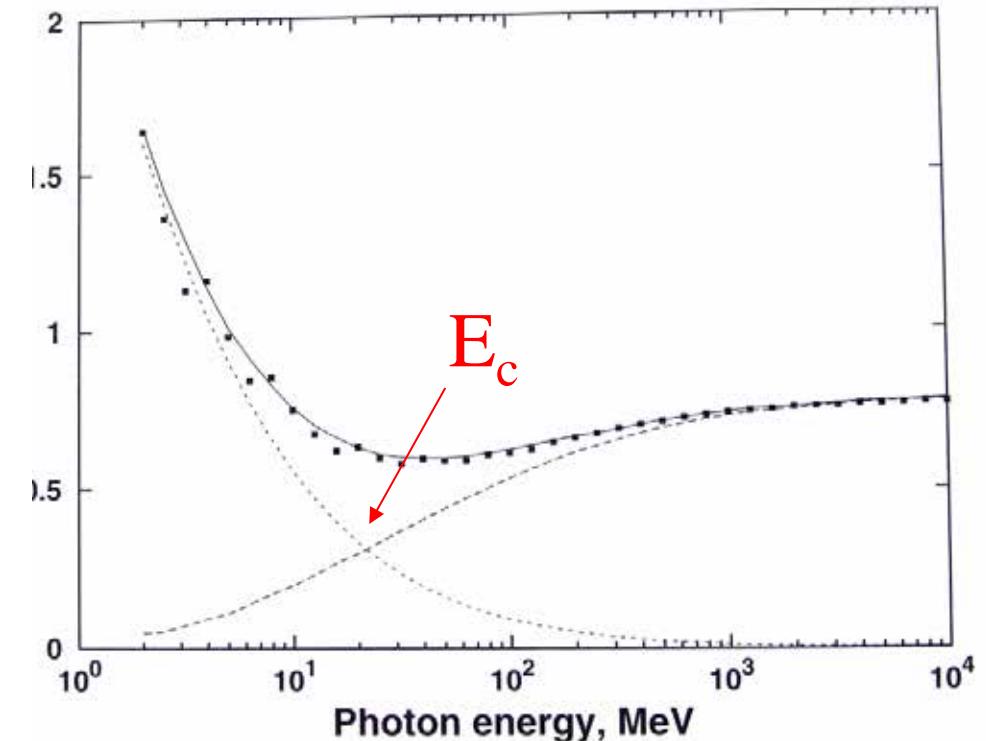
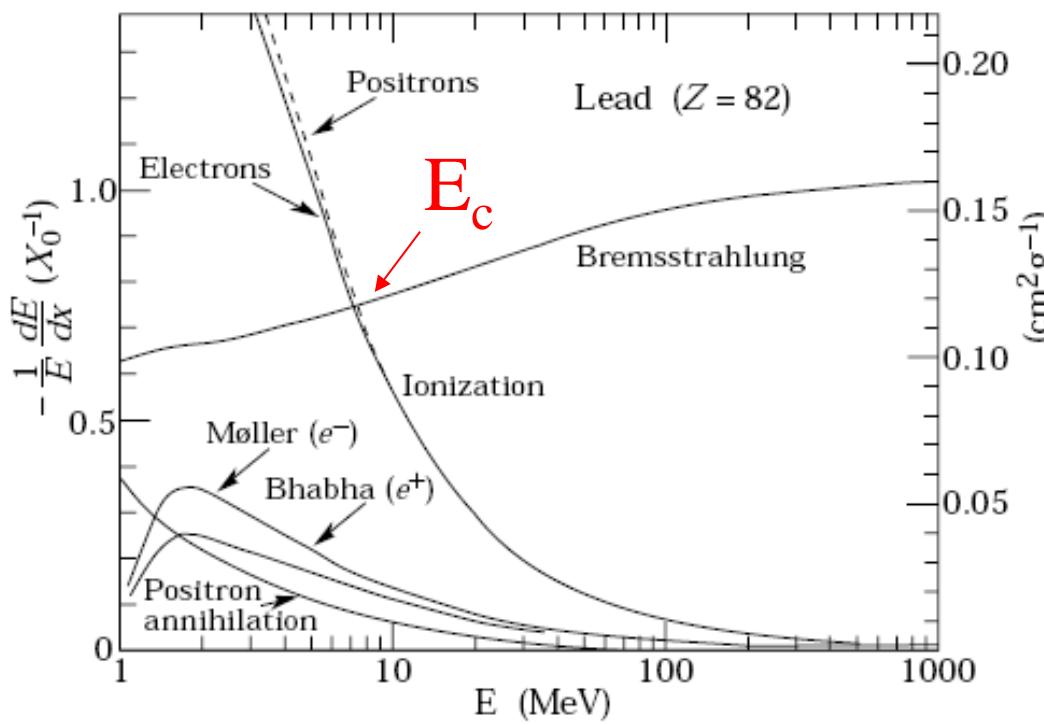


Radiation length X_0 :

- energy loss = $1/e$ due to bremsstrahlung
- $7/9$ of the range of a γ due to pair production.

In air : $X_0 = 36.7 \text{ g/cm}^2$

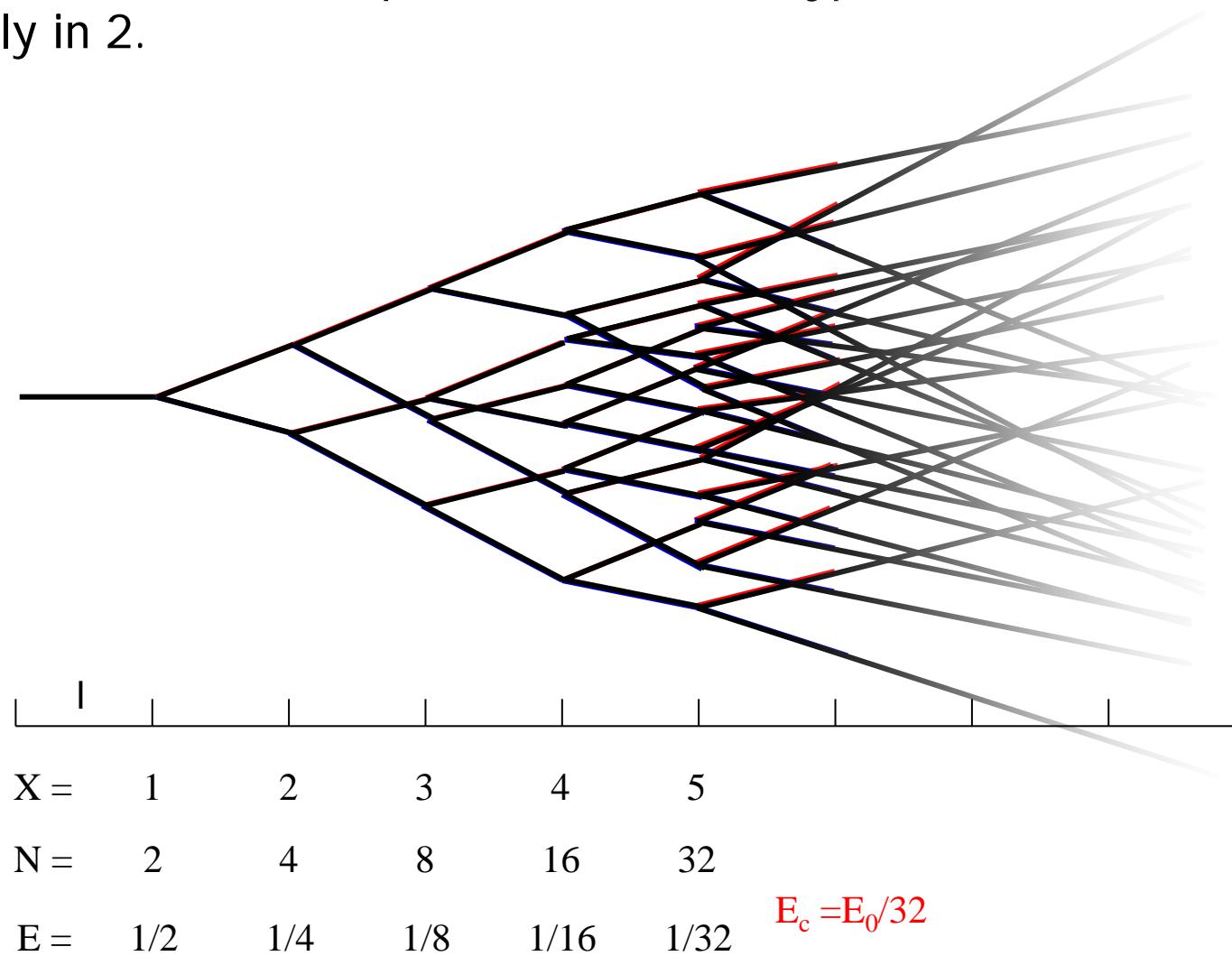
EM cascades (Rossi & Greisen)



Critical energy: below this energy, ionization losses dominate.

Simplified development model (Heitler)

- Cascade consisting of only one type of particles having an interaction length l .
- At each interaction, 2 particles of same type are emitted sharing the energy exactly in 2.



Longitudinal development

- After t radiation length, there are 2^t particles with energy

$$E = E_0/2^t$$

$$\text{soit : } t \ln 2 = \ln(E_0/E)$$

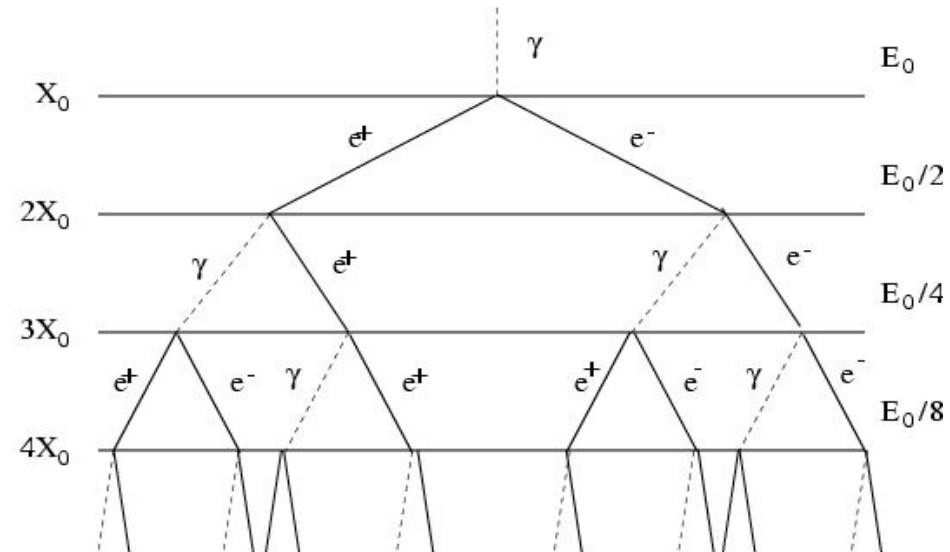
- The particles of energy E are produced at thickness:

$$t(E) \approx \ln(E_0/E)$$

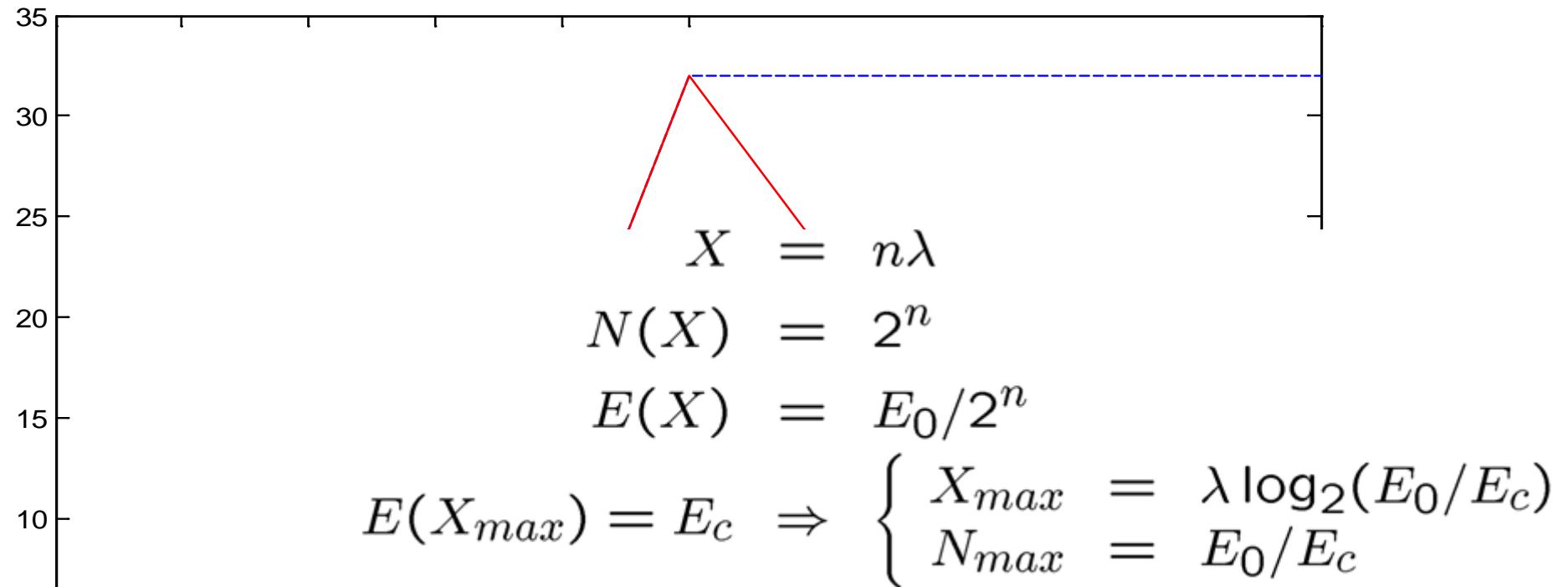
- The maximal development of the shower is reach for a thickness:

$$t_{max}(E_0) \approx \ln(E_0/E_c)$$

- More realistic models agree with this rough estimate.



Simplified development model (Heitler)



$$X = 1 \quad 2 \quad 3 \quad 4 \quad 5$$

$$N = 2 \quad 4 \quad 8 \quad 16 \quad 32$$

$$E = 1/2 \quad 1/4 \quad 1/8 \quad 1/16 \quad 1/32$$

$$E_c = E_0/32$$

Longitudinal development: Approximation "A" (B. Rossi, K. Greisen)

- Approximation "A" describes the shower development phase where only bremsstrahlung and pair creation are in action.
- From Bethe-Heitler theory, one obtains des integro-differential linear and coupled equations leading to:
 - $\Pi(E, t)dE$ = average number of e^\pm with energy $\in [E, E + dE]$, at tX_0 depth
 - $\Gamma(W, t)dW$ = average number of γ with energy $\in [W, W + dW]$, at tX_0 depth
- The simplifying factor is the absence of any energy scale.

Approximation A (cont)

- $\Pi(E, t)dE$ = average number of e^\pm with energy $\in [E, E + dE]$, at tX_0 depth
- $\Gamma(W, t)dW$ = average number of γ with energy $\in [W, W + dW]$, at tX_0 depth
- Initial condition :
 - If the primary particle is a γ : $\Gamma(W, 0) = \delta(E - E_0)$
 - If the primary particle is an e^\pm : $\Pi(E, 0) = \delta(E - E_0)$
- Obvious special solutions:
 $\Gamma(W, t) = f(t)/W^{s+1}$ et $\Pi(E, t) = g(t)/E^{s+1}$
(absence of energy scale)
... but they don't satisfy the initial conditions!

Approximation A (suite)

- The obvious solutions (power-law spectra, therefore scale invariant) correspond to an initial condition interesting in itself: an incident beam with a power law spectrum with an integral spectral index s .
- These special solutions form a base and a solution that fulfills the initial condition (photon or electron with an energy E_0) is obtained from a superposition of $1/E^{s+1}$ spectra (Mellin transformation, analogue to Fourier or Laplace transforms).
- Result : for a given value of t , the particle spectrum is very close to a power law $1/E^{s+1}$ with a value of s that varies with t and $y = \ln(E_0/E)$ following:

$$s = \frac{3t - 1}{t + 2y}$$

- The number of particle with energy E is maximal for $s = 1$

Taking into account ionisation energy losses: the "age" parameter

- Approximation A is not valid anymore when the electron mean energy energie is close to the critical energy E_c .
- One can modify the above results:

$$y = \ln\left(\frac{E_0}{E_c}\right) \text{ et } s = \frac{3t}{t + 2y}$$

- Semi empirical formula given by Greisen for an incident γ , for the mean number of electrons after traversing t radiation length:

$$\bar{N}_t = \frac{0.31}{\sqrt{y}} \exp\left[t\left(1 - \frac{3}{2} \ln s\right)\right]$$

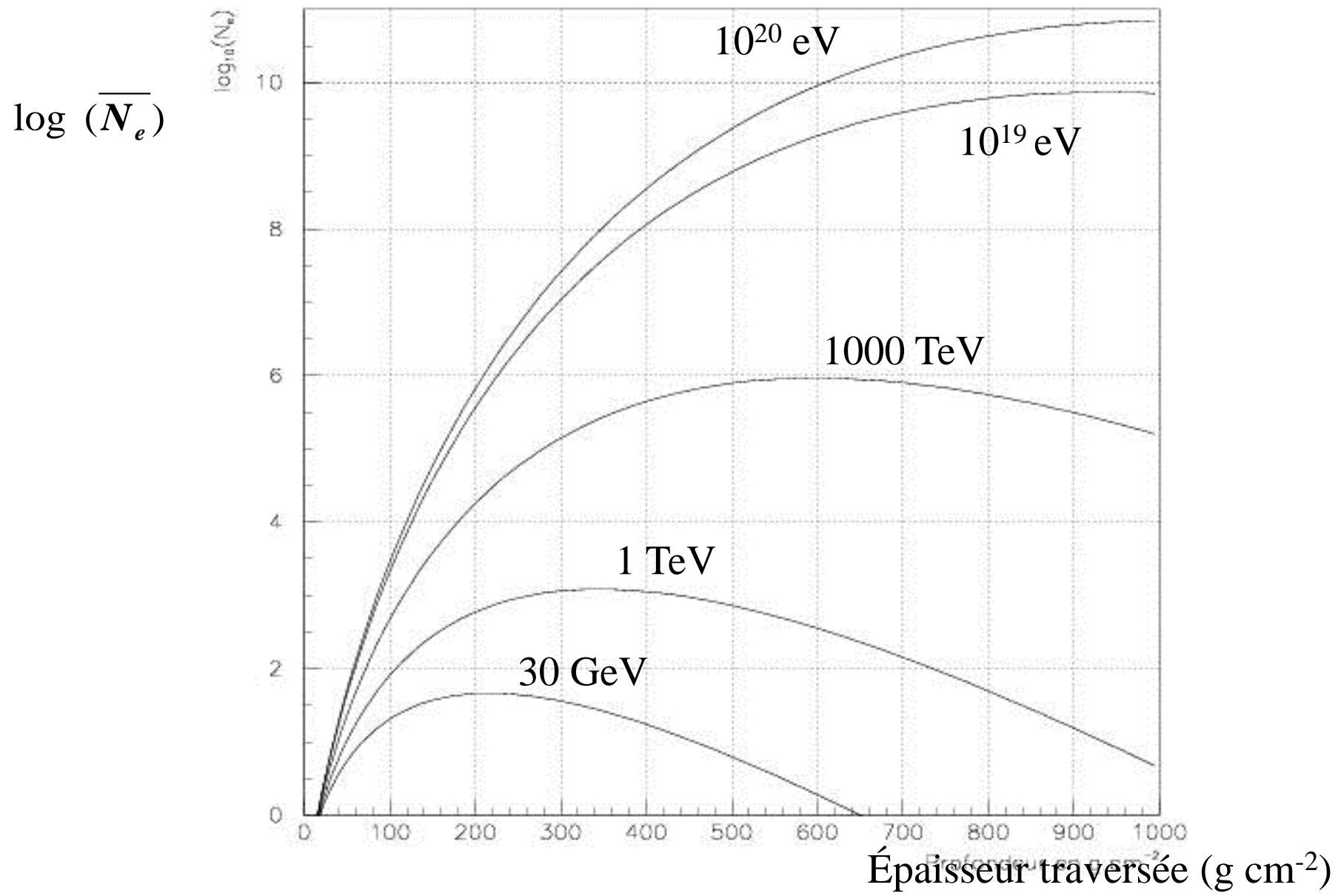
- The parameter s increase with t . It is < 1 during the development phase, reaches 1 at the maximal development stage for $t_{max} = y = \ln(E_0/E_c)$ and is > 1 during the extinction phase.
- s is called the "age".

EM showers : some orders of magnitude

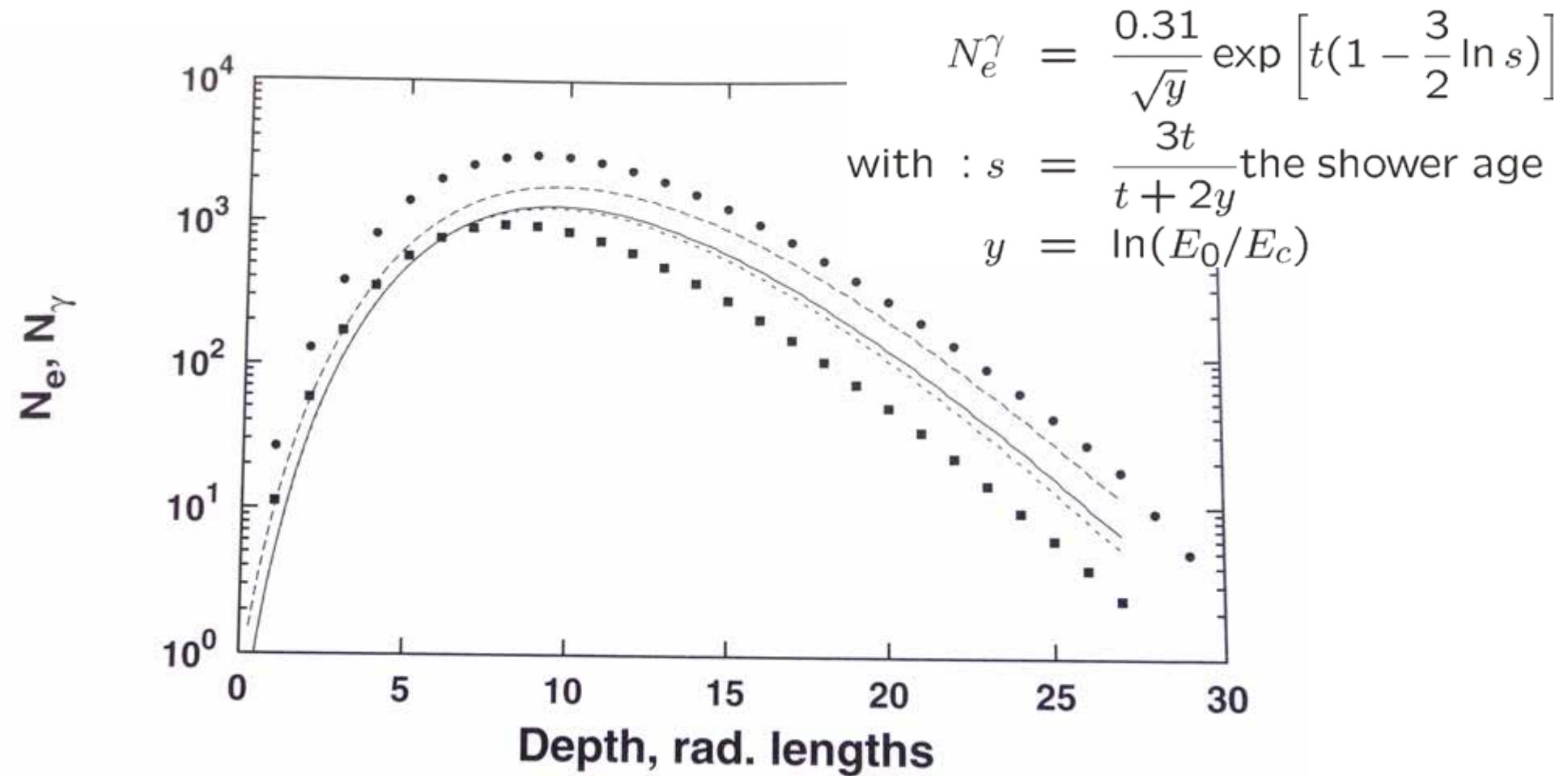
Primary g energy E_0	Thickness traverse $t_{\max} X_0$ (g cm $^{-2}$)	Altitude (m)	$N_e(t_{\max})$
30 GeV	216	12000	50
1 TeV	345	8000	1200
1000 TeV	600	4400	$0,9 \times 10^6$
10^{19} eV	936	1200	$7,4 \times 10^9$
10^{20} eV	1021	0	$7,0 \times 10^{10}$

γ

EM shower average profiles



EM cascades (Rossi & Greisen)



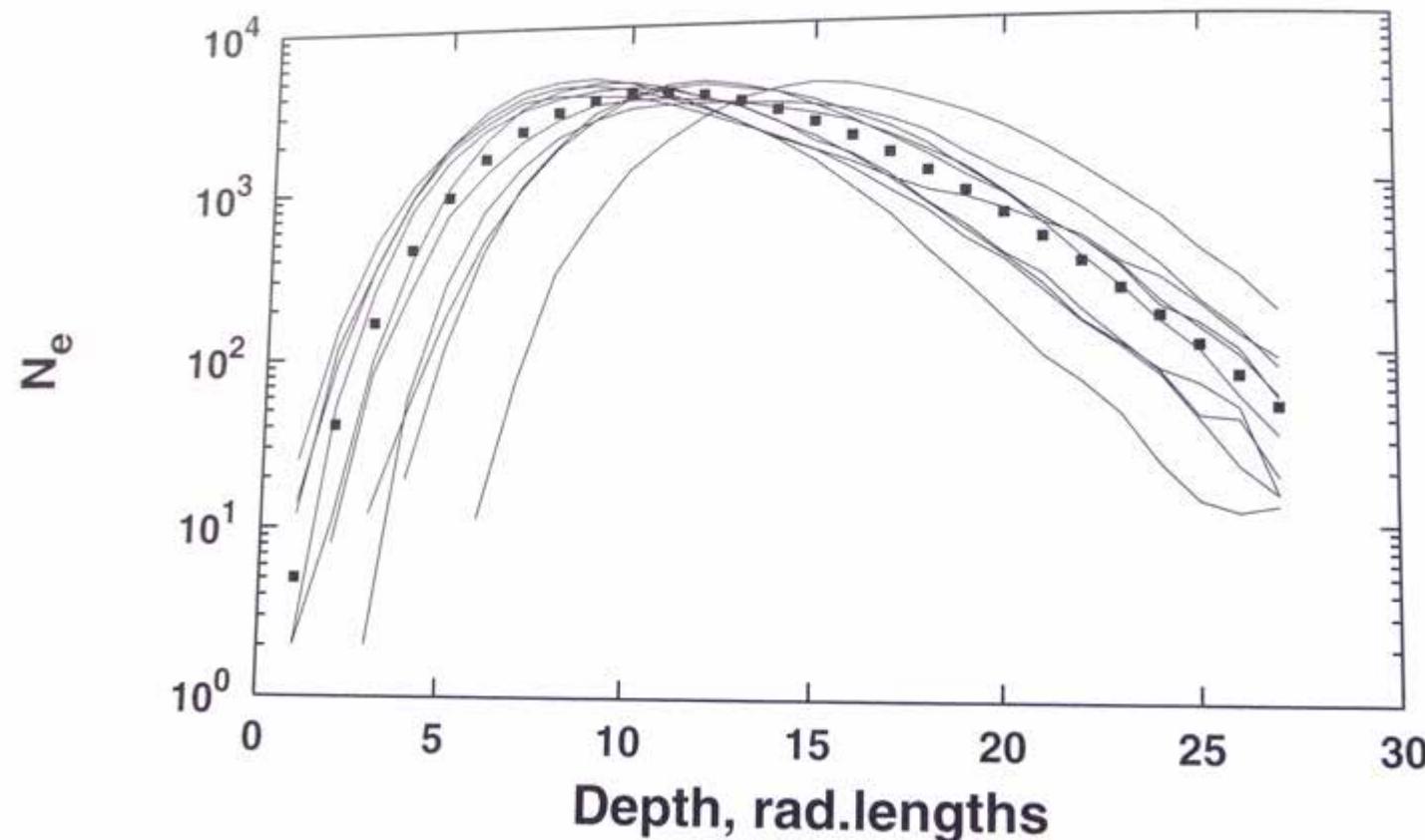
Parametrisation (Greisen) and Monte Carlo (EGS4)
photons 1 TeV, $E_c = 10$ MeV

Shower size

i.e. number of electrons at ground level
as an energy estimator

- At maximal development level, the mean number of electrons quasi proportionnal to the primary energy ($y = \ln(E_0/E_c)$).
- Fluctuations on N_e :
 - Fluctuations on the depth of first interaction (exponential law)
 - Fluctuations in the shower development (approximately log-normal because of the multiplicative behaviour)
 - Sampling fluctuations (depends on the type of detectors, their arrangement on the ground etc.)
- If the altitude of the maximal development is known (direct optical measurement), or if one can estimate the age independently (from lateral distribution of the electrons) one can avoid the first kind of fluctuations.
- Fluctuations are minimal at the maximum of development.

Cascades EM (Rossi & Greisen)

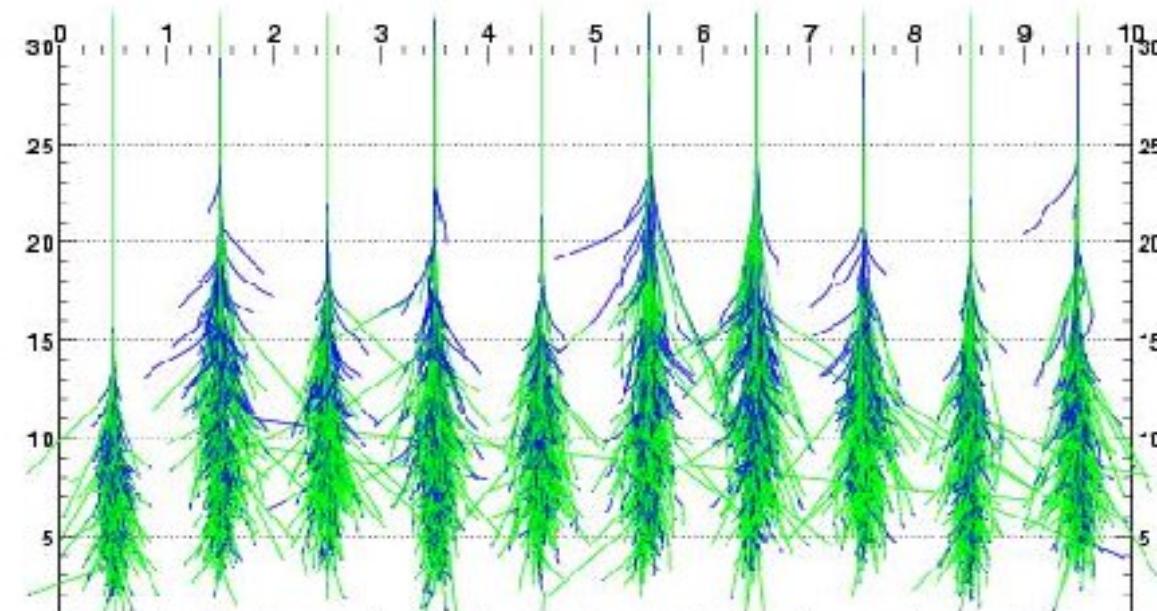


Shower to shower fluctuations

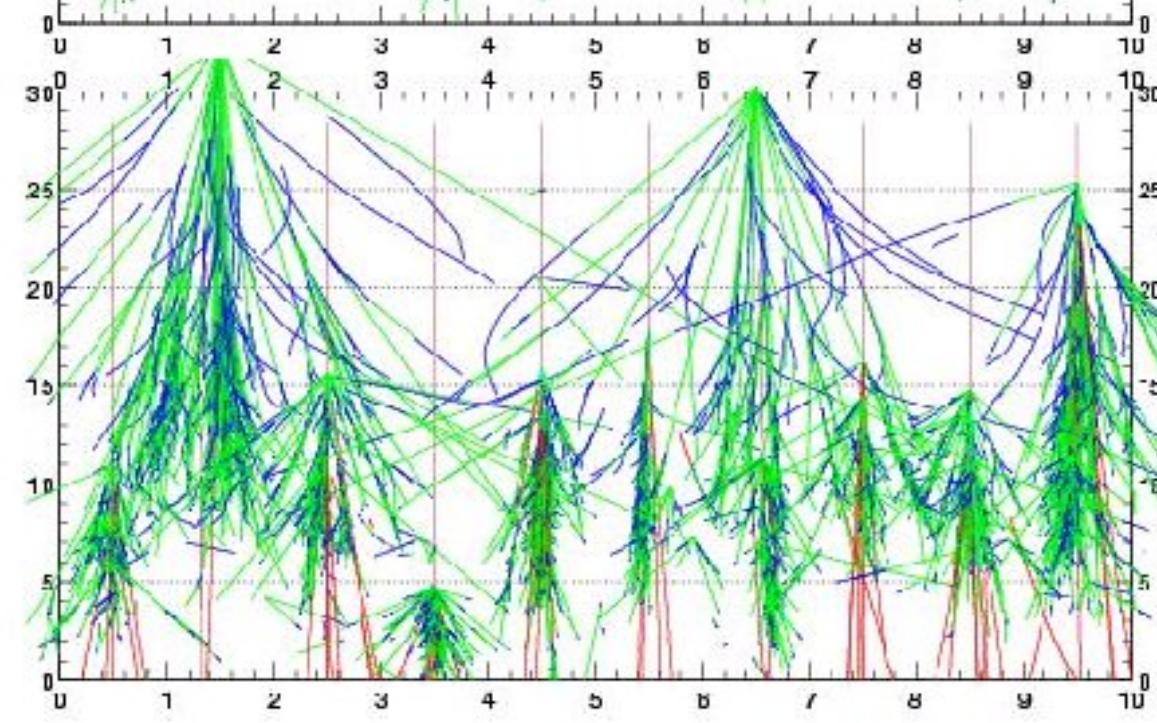
10 showers at 10^{14} eV compared to the average of 100 showers.

GAMMA-RAY (EM) INDUCED SHOWERS

10 γ
300 GeV



10 protons
300 GeV



*Simulations de
M. de Naurois*

Electromagnetic showers (e^\pm or γ primary)

Dominating phenomena

- Radiation processes:
 - Bremsstrahlung of e^\pm
 - Pair production ($> \text{MeV}$) pairs e^+e^-
- Multiple scattering
(small angular deflexions) of e^\pm
- Energy losses by e^\pm
 - par ionisation
 - excitation des atomes



In the coulombian
field of nuclei

γ induced
shower 300 GeV

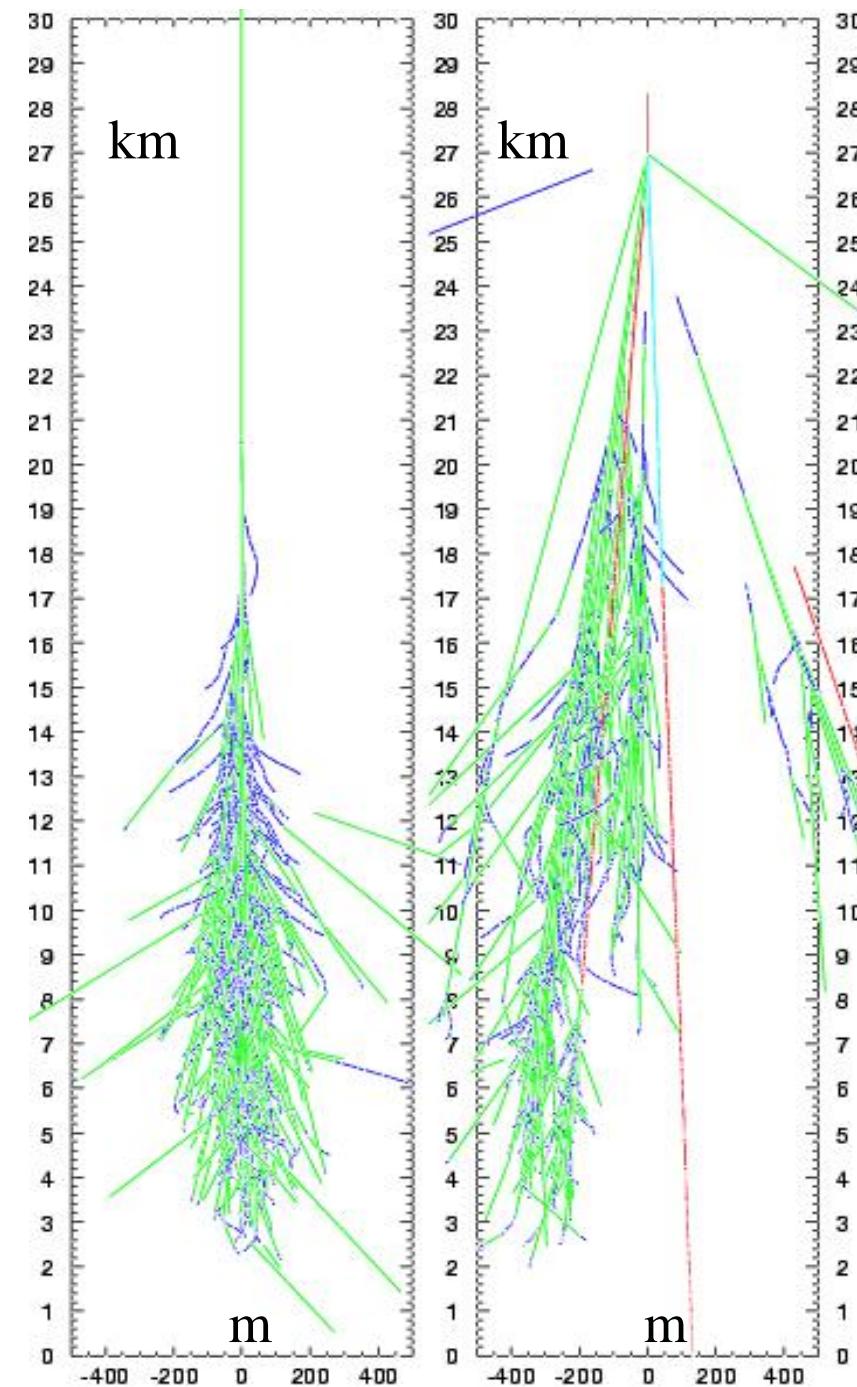
Roughly
symmetric
around the axis

Small transverse
dispersion
(multiple scattering)

(almost) no muons
...

(unless $E_0 > 1$ PeV)

Essentially
 e^+ e^- and γ
secondaries



proton induced
shower 300 GeV

Large transverse
momentum

Muon component
(from mesons decays)

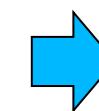
A hadronic shower
does contain
EM sub-showers

Optical photon emission by showers

2010-2011

F.Montanet Cours Astroparticles M2R PSA & AMD Grenoble

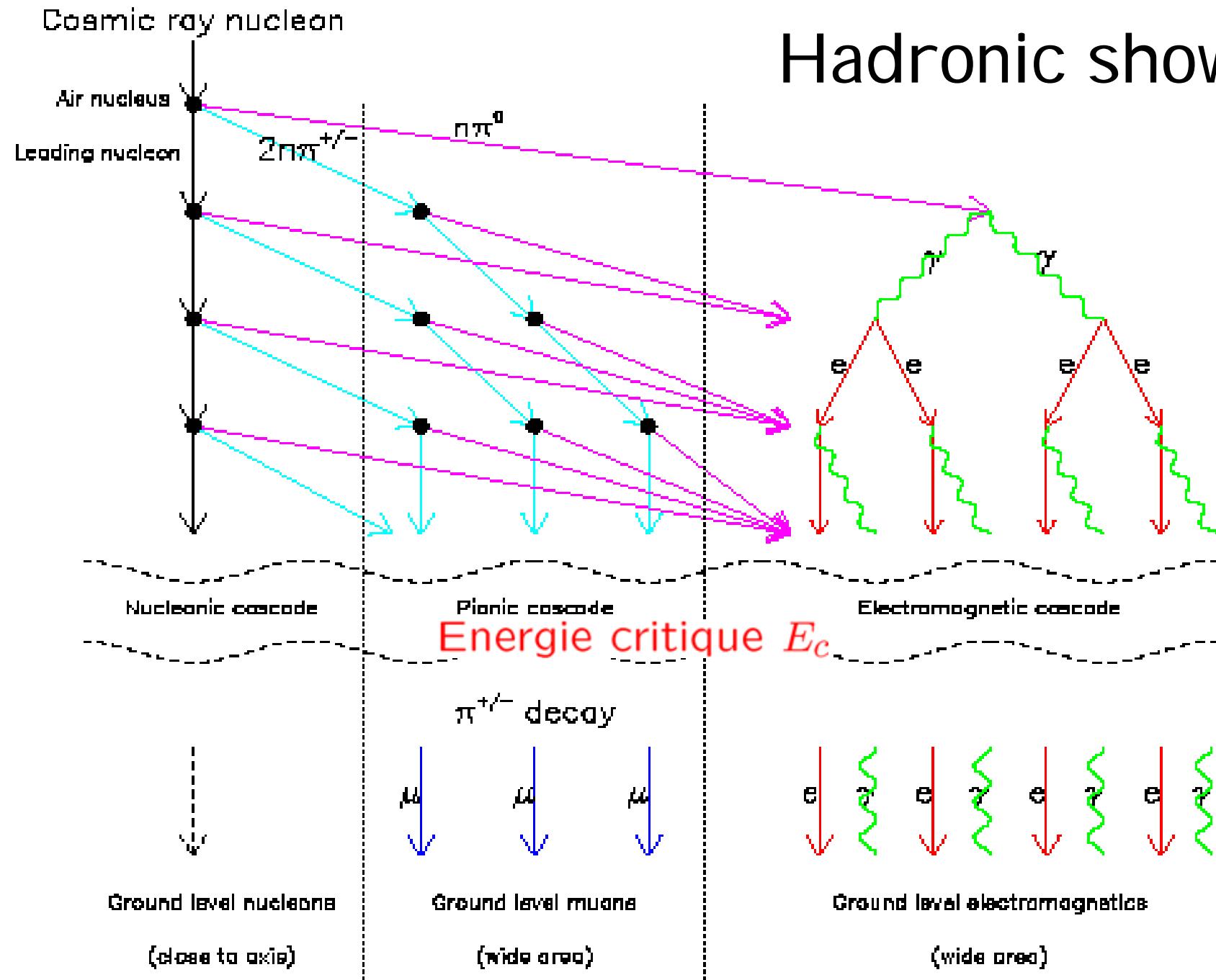
- Showers charged particles emit light:
 - **Cherenkov light** : very collimated along the shower axis (Cherenkov angle at 1 Atm. $\approx 1^\circ$) threshold depending on the altitude : at ground 22 MeV for e^\pm et 4.5 GeV for μ^\pm
(20 photons per m per $\beta \approx 1$ charged particle at 1 atm)
Essentially used for gamma-ray astronomy
 - **Nitrogen fluorescence**: isotropic emission
($\frac{1}{4}$ 4 photons per electron per m)
Essentially used at UHE $\geq 10^{18}$ eV.
- This light detected by ground telescopes gives us very rich information on the **3D development of the showers**. It give a quasi calorimetric reliable measurement of the energy.
- ... but optical detectors can only work during moonless clear sky nights ($\approx 10\%$ duty cycle).



Lecture on
Imaging & Cherenkov
Detectors

HADRONIC SHOWERS MODELS AND DETECTION

Hadronic showers

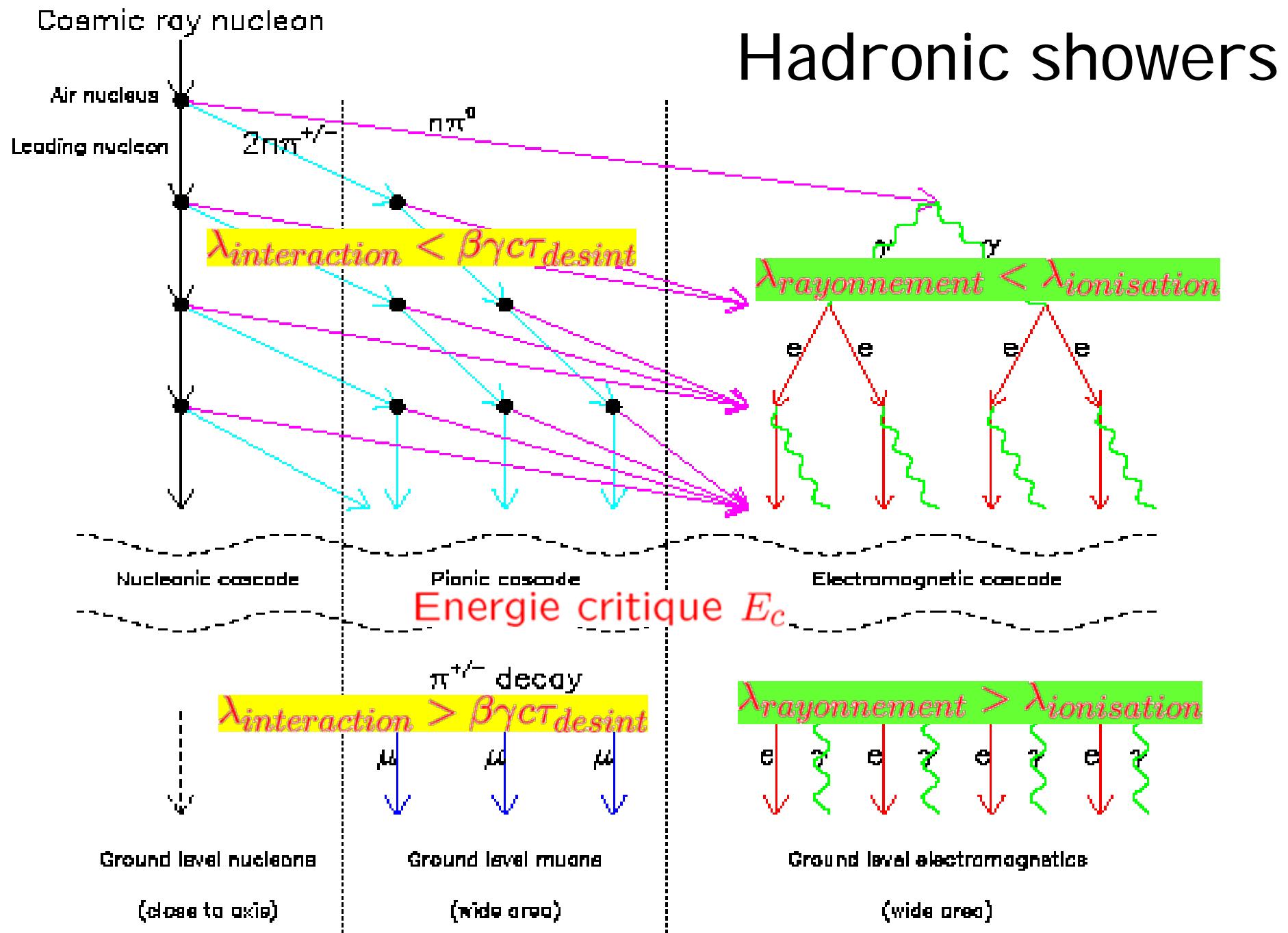


"Hadronic" showers (proton ou noyau primaire)

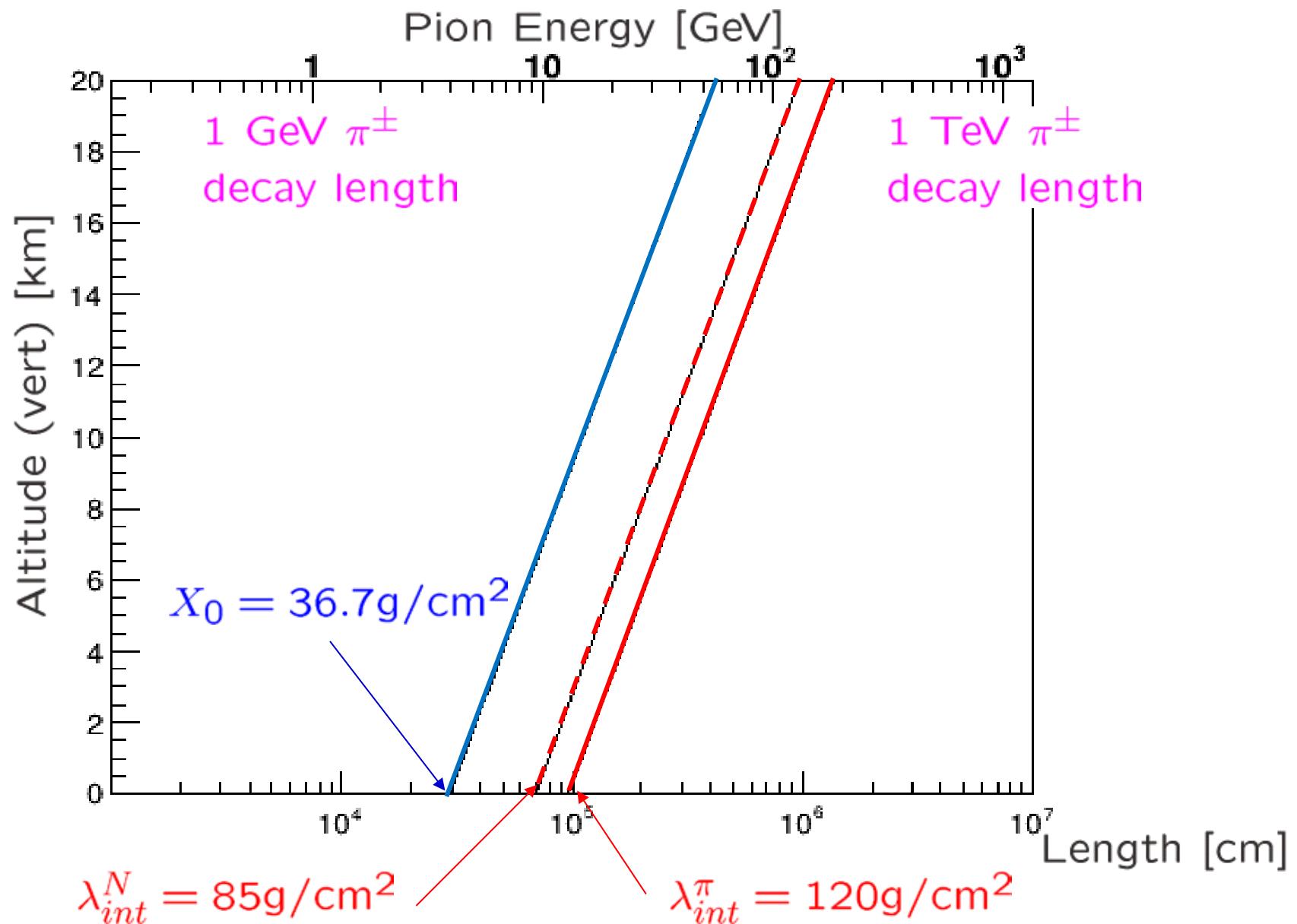
- Great complexity implying the use of numerical simulations:
 - Many length scales : nucleon interaction length, pion interaction length, EM radiation length, atmosphere density height scale...
 - Superposition of a nuclear cascade, a pionic cascade and an electromagnetic cascade (the later from π^0 decay γ).
 - Large fluctuations in the multiplicity of secondaries.
- But simulations are subject to many uncertainties:
 - p+N or N+N interactions: sensitivity to nuclear models.
 - Energy range unexplored by accelerators and colliders : sensitivity to nucleon structure functions (parton distributions) and fragmentation functions extrapolated far from the measured regions.
 - The inelasticity and in general the very forward diffractive physics is not well measured in fixe target experiment (even worse at colliders). Still, the main behaviour observed on EM showers remains valid.

From EM to Hadronic showers

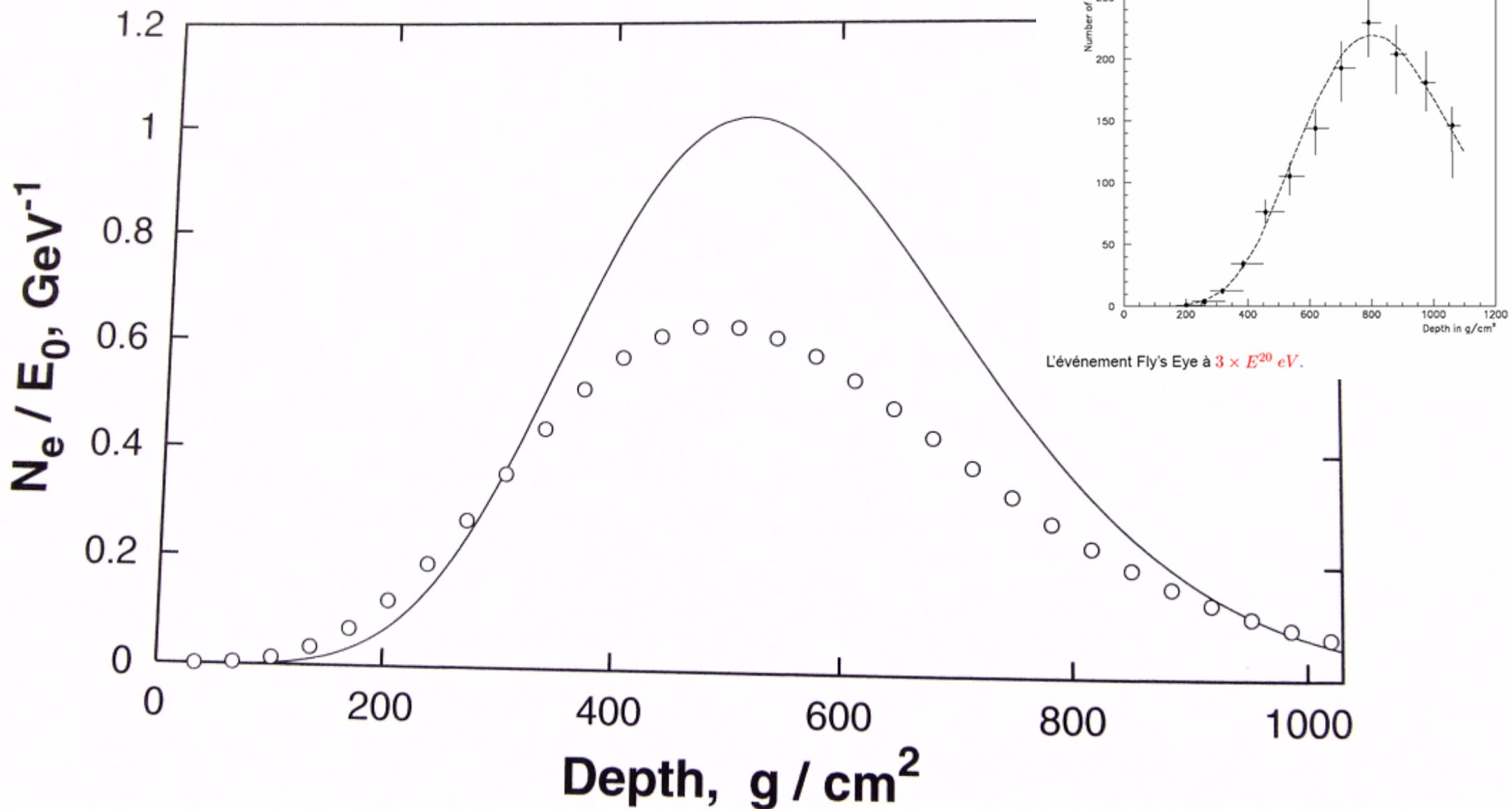
- The main observables are the same:
 - Number of electrons, gamma **but also muons** at ground and their lateral distributions.
 - Longitudinal profil and maximal dev. altitude (optical detectors).
 - Number of muons at ground level and lateral distribution of muons.
- Feynman scaling is rather well verified in the fragmentation: it plays an role analogue to that of Bethe-Bloch formulae for EM showers (absence of mass/energy scale).
- Simulations have allowed to establish empirical formulae inspired by EM showers useful to to quite estimates (*T.K. Gaisser, A.M. Hillas*)



Interaction and radiation lengths in atmosphere



Development of Hadronic vs EM showers



Gaisser longitudinal Parametrisation

Gaiser Hillas formulae :

$$N_e(X - X_1) = N_e^{max} e^p \left(\frac{X - X_1}{X_{max} - \lambda} \right)^p \exp - \left(\frac{X - X_1}{\lambda} \right)$$

avec $p = \frac{X_{max} - \lambda}{\lambda}$

Averaging on X_1 depth of 1st interaction :

$$\bar{N}_e(X) = N_e^{max} \frac{p}{p+1} e^p \left(\frac{X}{X_{max} - \lambda} \right)^{p+1} \exp - \left(\frac{-X}{\lambda} \right)$$

$$X_{max} = X_0 \log \left(\frac{E_0}{\epsilon_0} \right)$$

$$N_e^{max} = \frac{E_0}{\omega}$$

Radiation length : $\approx 36.7 \text{ g/cm}^2$

Critical energy : $\epsilon_0 \approx 74 \text{ eV}$

Empirical relation between size and energy: $\omega \approx 1.7 \text{ GeV}$

Incident nucleus interaction length (of energy E_0) $\lambda_N \approx 70 \text{ g/cm}^2$

Longitudinal developpement

Xmax and energie :

$$X_{max} \approx X_0 \log \left(\frac{E_0}{\epsilon_0} \right)$$
$$\Rightarrow 80 \text{ g/cm}^2 \text{ per energie decade}$$

Nuclei :

Superposition principle : a nucleus A_N is equivalent to A protons.

Thus :

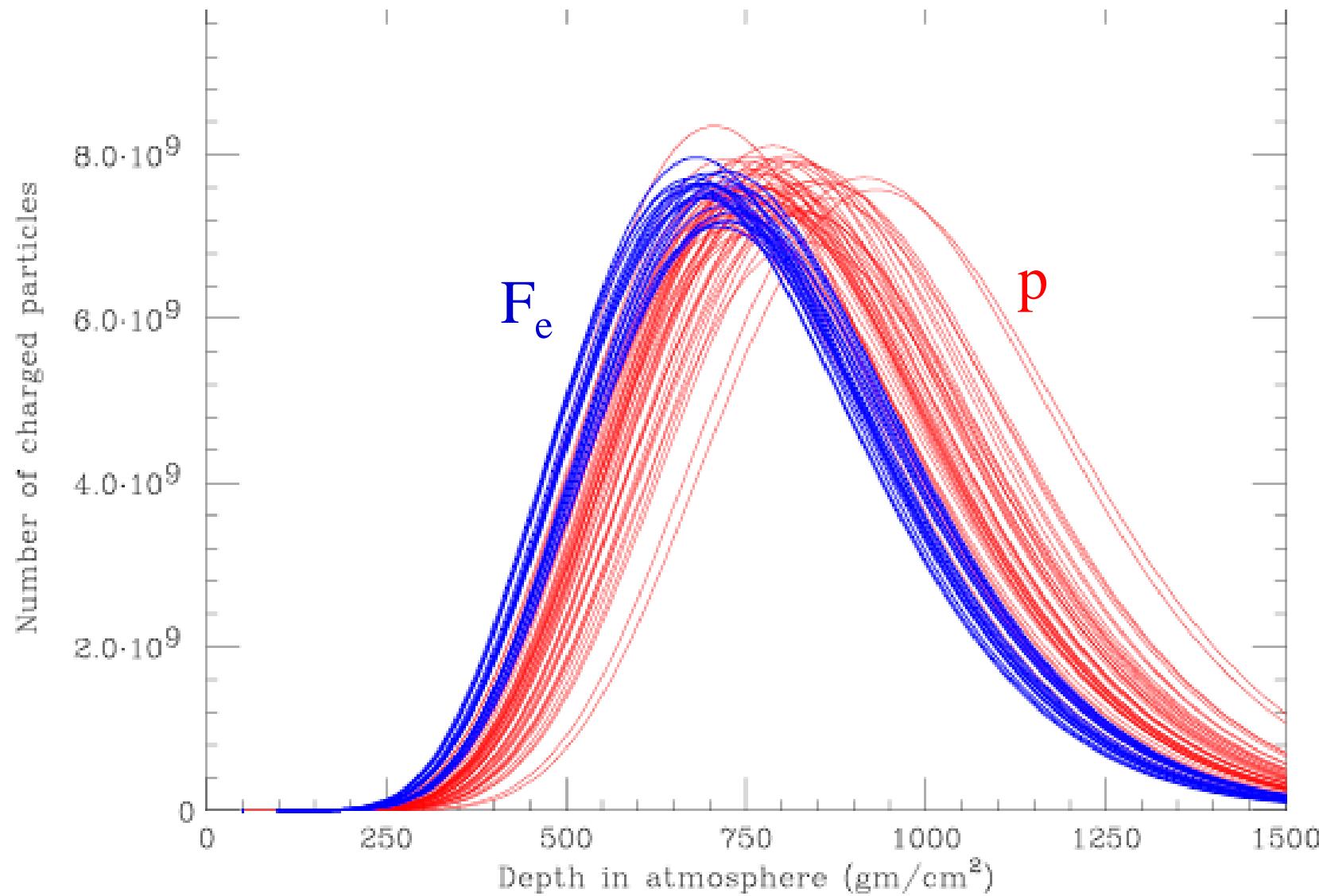
$$X_A^{max} = X_0 \log \left(\frac{E_0}{A\epsilon_0} \right)$$
$$= X_p^{max} - X_0 \log(A)$$

For example iron/proton $A = 56$:

$$X_0 \log(A) = 36.7 \log(56) = 148 \text{ g/cm}^2$$

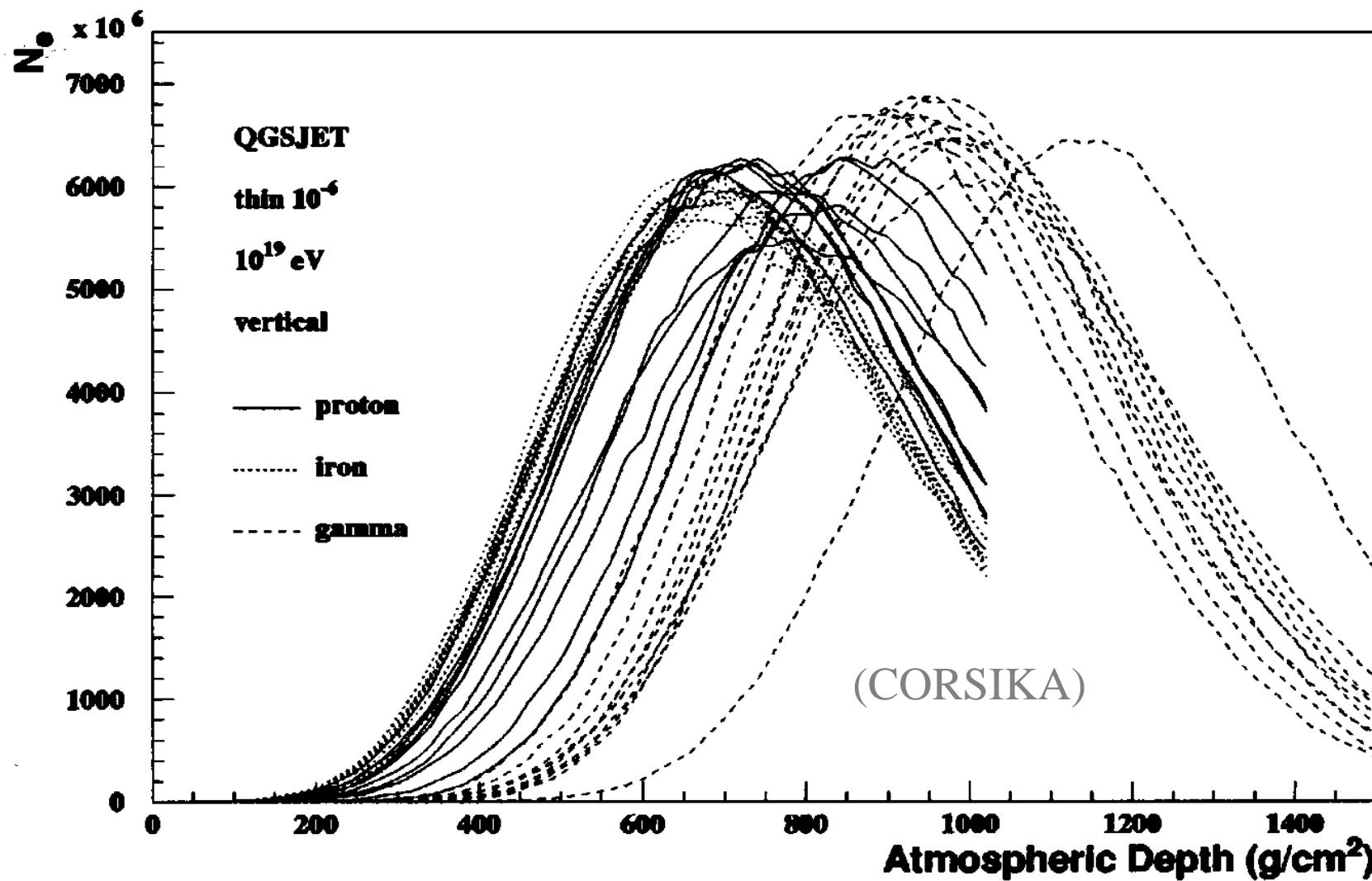
Structure in space

Shower to shower fluctuations largely due to the depth of the first interaction.



Primary identification

- Requires a good statistics and a good knowledge of the initial energy, the shower angle (+ systematic corrections because of atmospheric attenuation)



Radial extention

The radial distribution is determined by **the mean transverse momentum (P_T)** from hadronic interactions and by **multiple scattering**. In air, the Molière radius is $\approx 75\text{m}$.

Molière radius (~1/4 of the radiation length) :

$$\begin{aligned}\langle \delta\theta^2 \rangle &= \left(\frac{21\text{MeV}}{E} \right) \delta X \\ r_1 &= \left(\frac{E_s}{E_c} \right) X \approx 9.3 \text{ g/cm}^2\end{aligned}$$

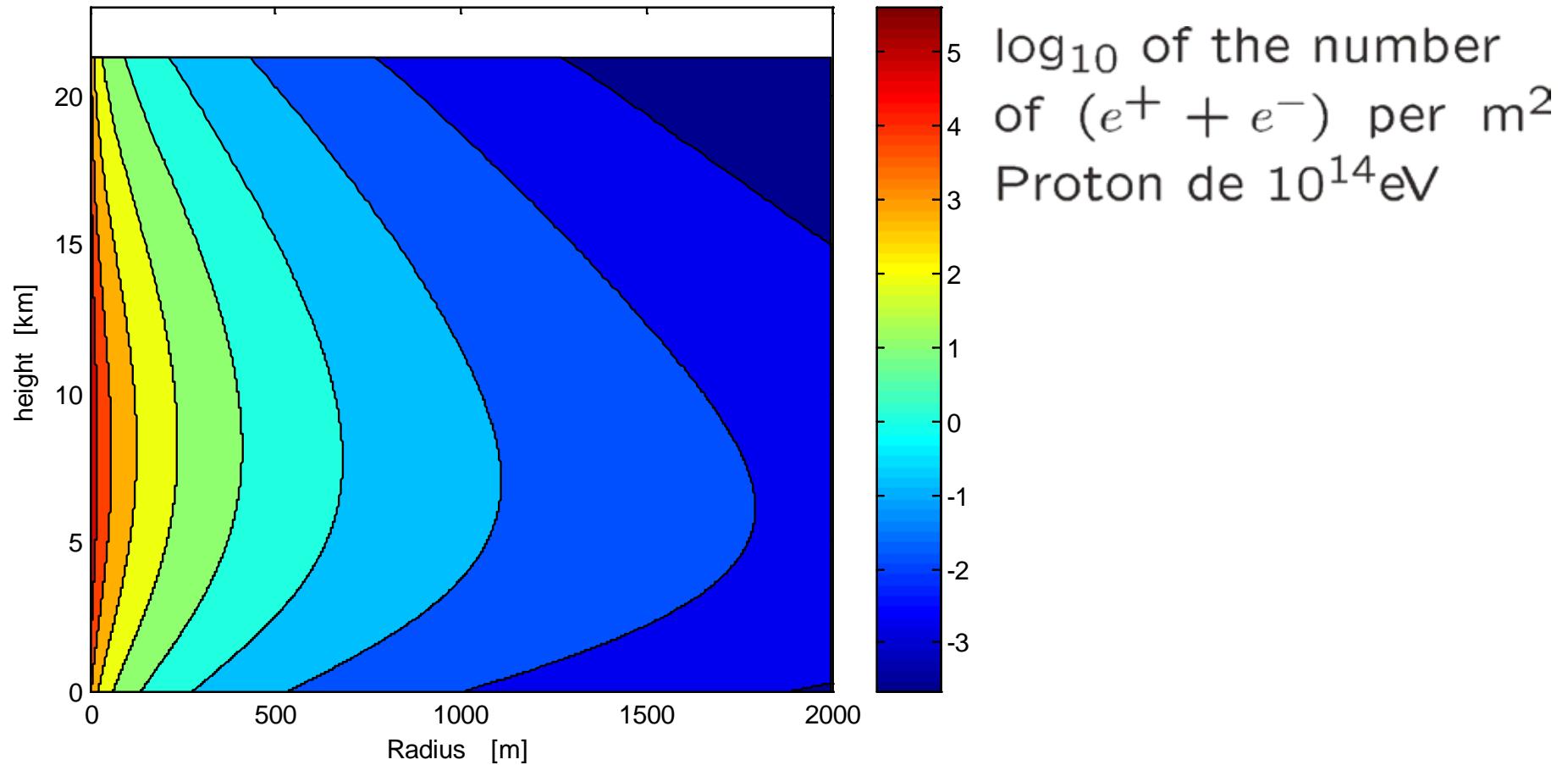
Nishimura, Kamata, Greisen :
multiple scattering + transverse momentum

$$\begin{aligned}xf(x) &= C(s)x^{(s-1)}(1+x)^{(s-4.5)} \\ \text{with : } x &= \frac{r}{r_1}\end{aligned}$$

normalization tel que :

$$2\pi \int_0^\infty xf(x)dx = 1$$

$e^+ + e^-$ lateral density



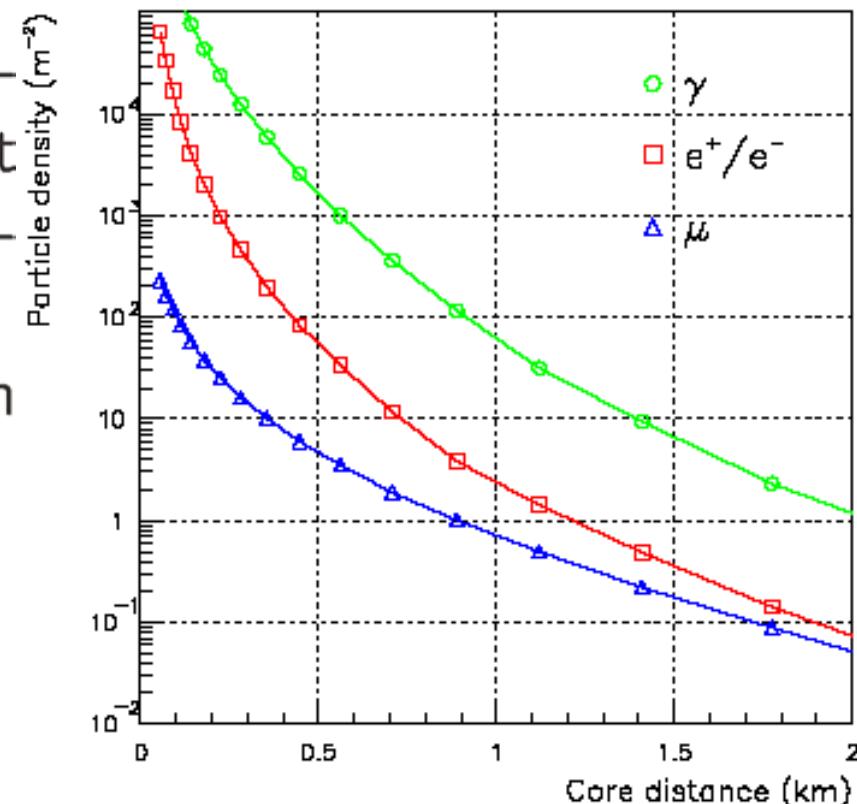
Lateral evolution

The density as a function of the distance to the center of the shower is characterized by a **lateral desity function (LDF)**

$$\rho(r) \propto k \times r^{-[\eta+f(r)]}$$

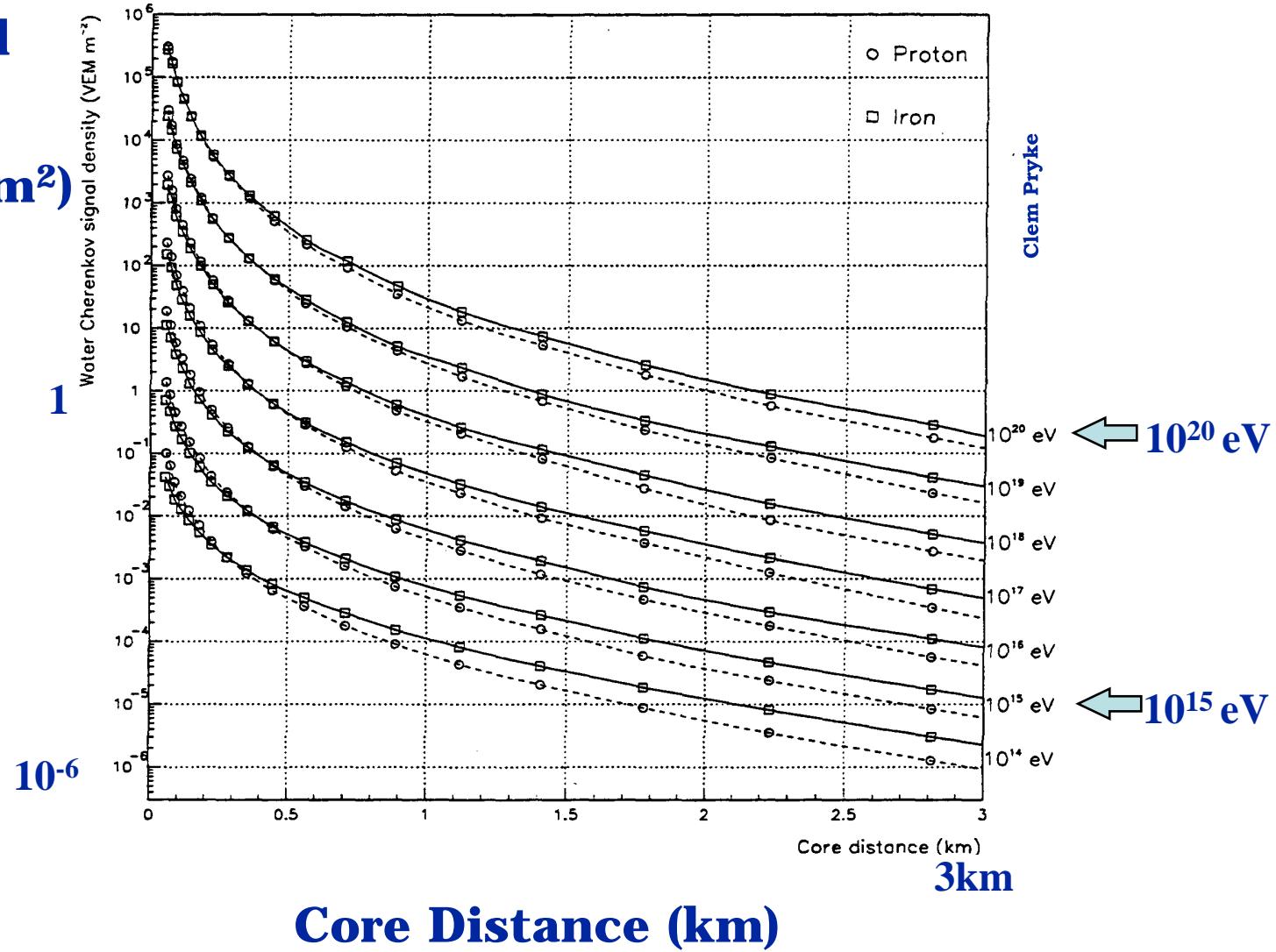
where f et k depends on the type of detectors used where η depends on the incident angle of the shower and the primary energy.

For $r > 800\text{m}$ this (empirical) expression must be modify as $(r/800)^{1.03}$



Shower Density Lateral Distribution (simulation)

**Detector Signal
Density**
(equiv.muons/m²)



Particle energy distribution

Rossi Greisen :

$$\frac{dN}{d(\log E)} \approx \frac{1}{E^{s+1}}$$

E.Nerling (thesis) :

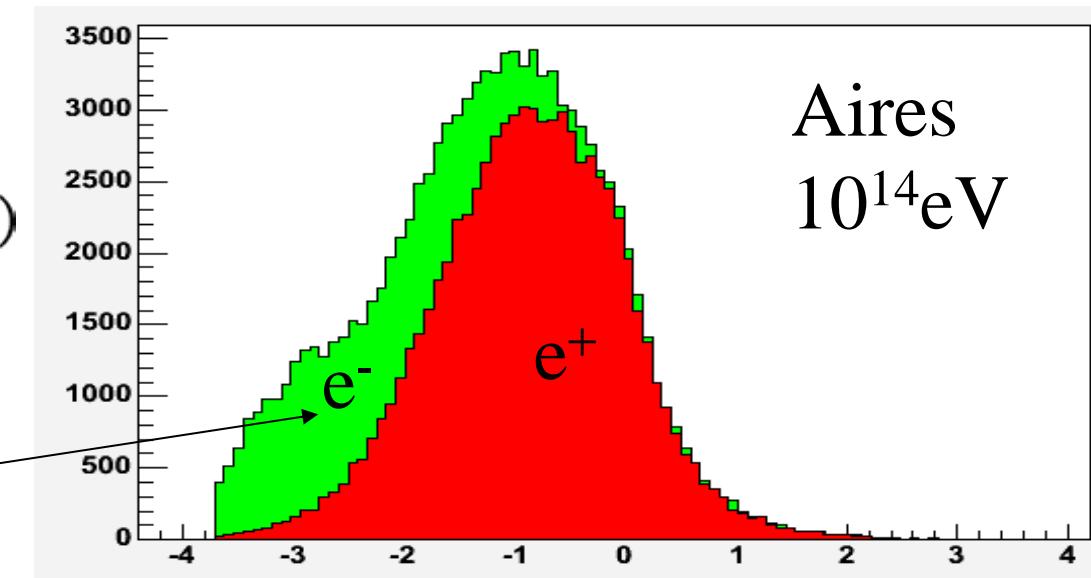
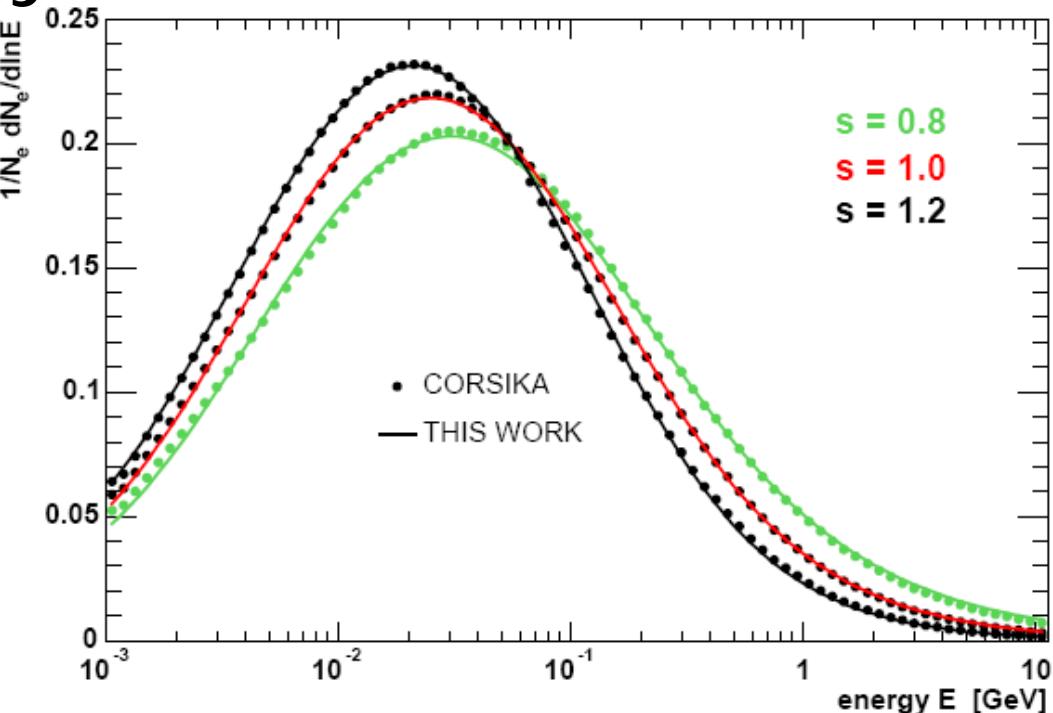
$$f_e(E, s) = a_0 \frac{E}{(E + a_1)(E + a_2)^s}$$

$$a_1 = 6.42522 - 1.53183.s$$

$$a_2 = 168.168 - 42.1368.s$$

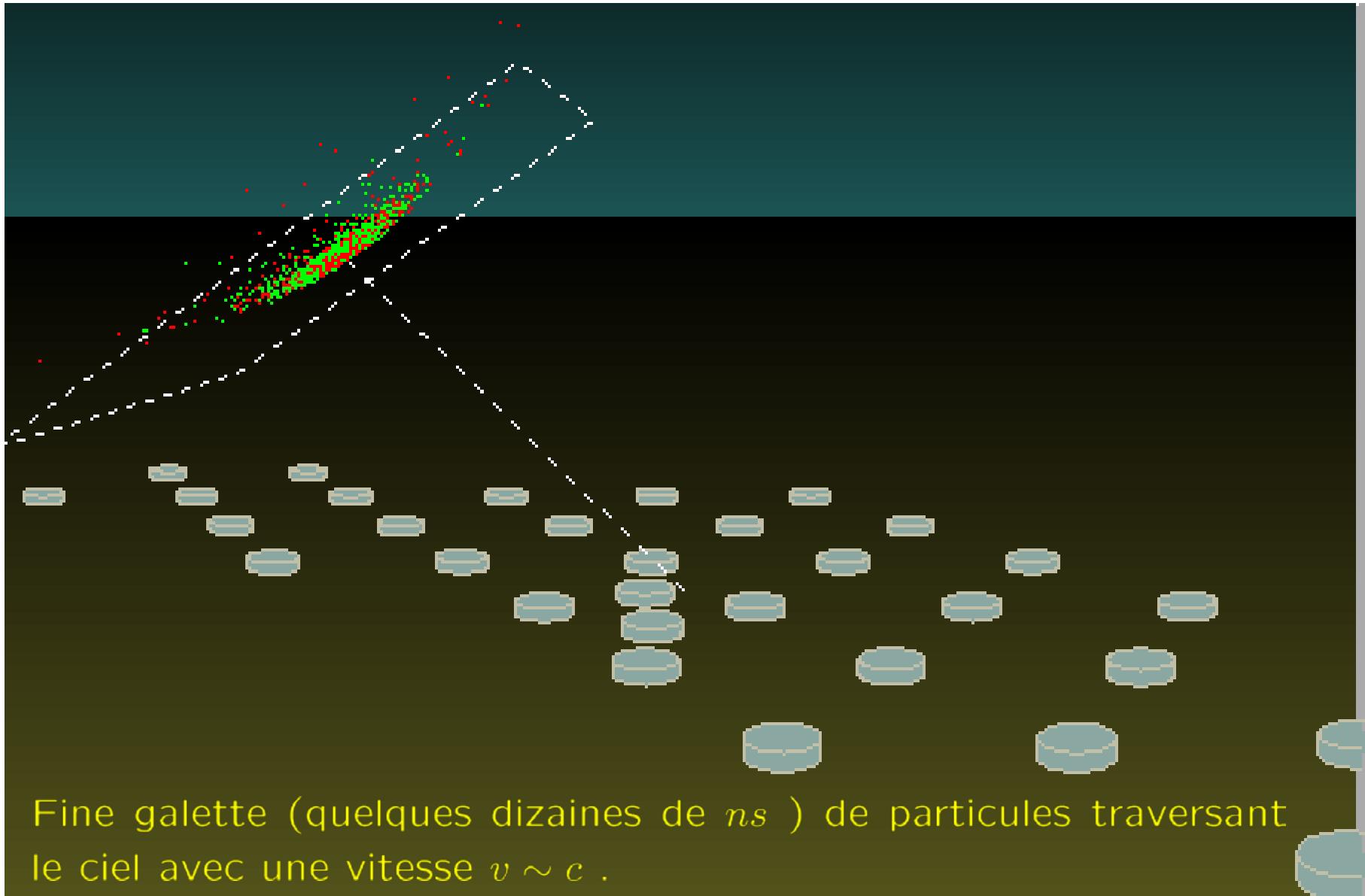
E en MeV

$$a_0^{-1} = \int_{\log E_{cut}} f(E, s) d(\log E)$$



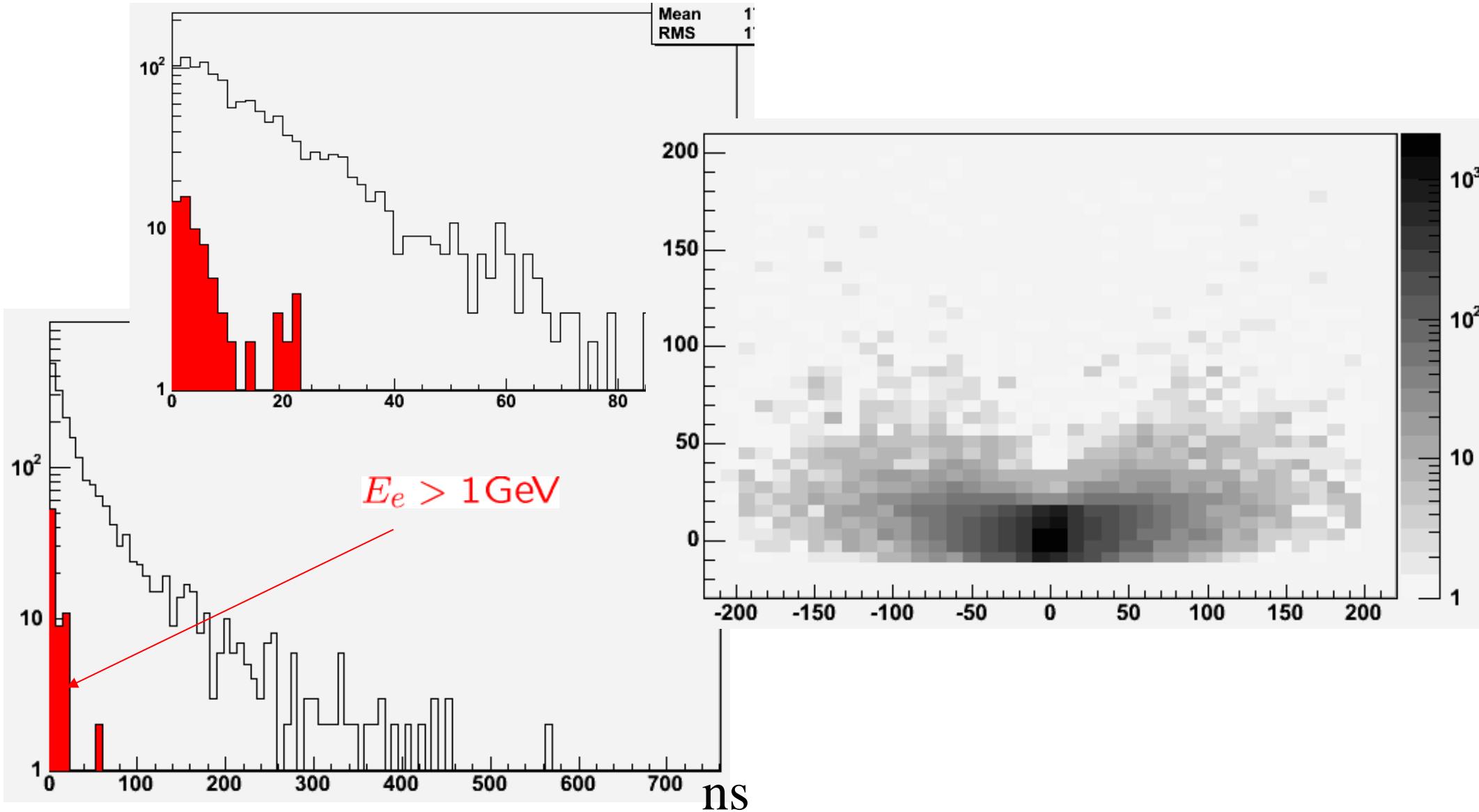
Excess e⁻ at low energy (ionization)

Time structure

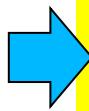


Fine galette (quelques dizaines de ns) de particules traversant le ciel avec une vitesse $v \sim c$.

Time structure



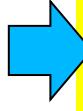
UHECR detection

 Lecture on
Imaging & Cherenkov
Detectors

Neutrino Physics with astroparticules

2014

F.Montanet Experimental Astroparticle Physics ESIPAP

 Lecture on
Imaging & Cherenkov
Detectors