



Radiative B meson decays at LHCb

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Introduction

- Radiative $b \rightarrow s\gamma$ decays proceed via an effective FCNC (loop penguin diagrams)
 - Sensitive probe of physics at high mass scales
 - Good testing ground in search of new physics
- FCNC processes are described by an effective Hamiltonian in the form of Operator Product Expansion :

$$H_{eff} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \left[\underbrace{C_i(\mu) O_i(\mu)}_{\text{left-handed part}} + \underbrace{C'_i(\mu) O'_i(\mu)}_{\text{right-handed part suppressed in SM}} \right] + \sum \frac{c}{\Lambda_{NP}^2} O_{NP}$$

$i = 1, 2$	Tree
$i = 3 - 6, 8$	Gluon penguin
$i = 7$	Photon penguin
$i = 9, 10$	Electroweak penguin
$i = S$	Higgs (scalar) penguin
$i = P$	Pseudoscalar penguin

Introduction

- NP can modify the Wilson coefficients (C_i) affecting observable quantities as (in case of radiative decays) photon polarization (C'_7)
 - Observation of the photon polarization in $B^+ \rightarrow K^+ \pi^- \pi^+ \gamma$ decays
 - Measurement of the photon polarization in $B_s \rightarrow \varphi \gamma$ decays
- The left-handed nature of the weak interactions $SU(2)_L$ predicts:

W boson couples to left handed quarks



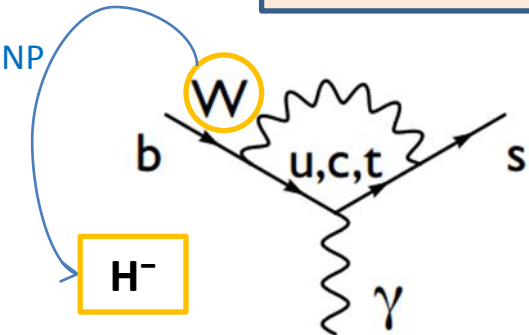
The photon is left handedly polarized

$b \rightarrow s \gamma_L$ (left-handed polarization)

$\bar{b} \rightarrow \bar{s} \gamma_R$ (right-handed polarization)

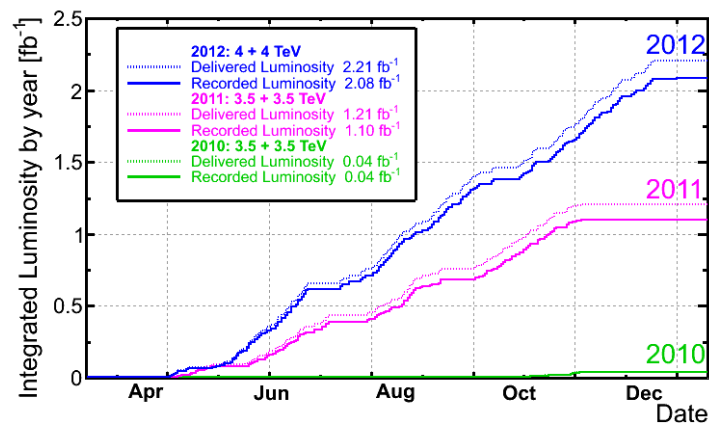
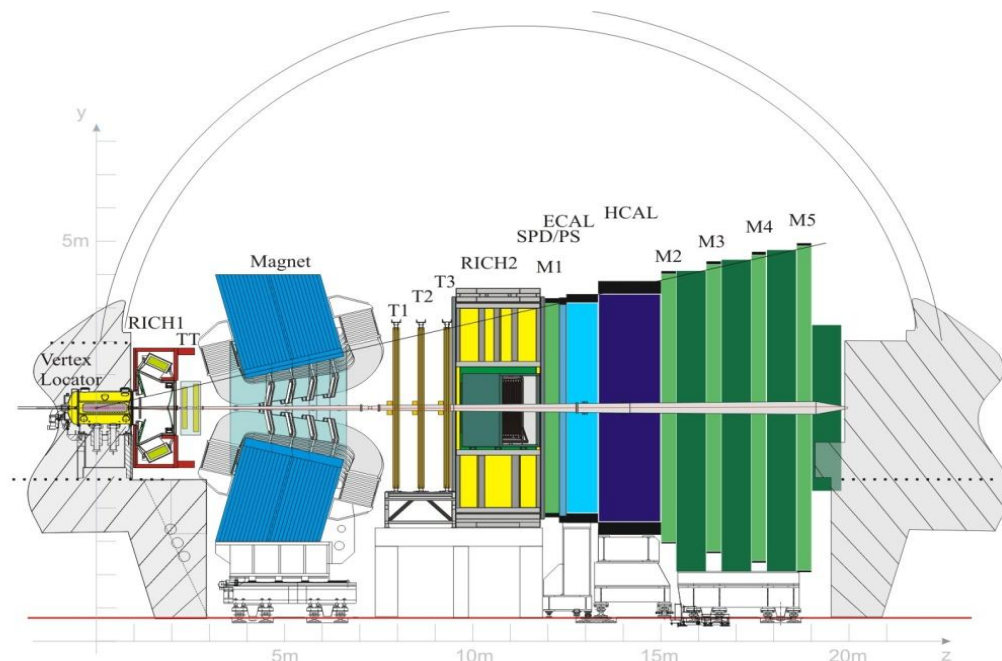
The opposite polarization is suppressed by m_s/m_b

Pheno 2014



The LHCb detector

- Single-arm forward spectrometer ($2 < \eta < 5$)
- Vast program of heavy flavor studies
- Challenge:
 - Precision measurements in a hadronic environment
 - High background level from large multiplicities
 - 30 particles for hard PP collision
 - High statistics from 30 KHz rate
 - Access to all b species:
 - Bd, Bu, Bs, Ab, Ξ b,
- Excellent performance:
 - $\sigma(m) \sim 10\text{-}25 \text{ MeV}/c^2$ (except for radiative $\sim 90 \text{ MeV}/c^2$)
 - $\Delta p/p \sim 0.4\text{-}0.6 \%$ ($5\text{-}100 \text{ GeV}/c$)
 - High proper time resolution $\sigma(t) = 60 \text{ fs}$
 - $\sigma_E/E = 10\%/ \sqrt{E} \oplus 1\%$
 - $\epsilon(K \rightarrow K) \sim 95\%$ for $\epsilon(\pi \rightarrow K) \sim 5\%$



Rad. Decays at LHCb: $B^0 \rightarrow K^{*0} \gamma$ and $B_s \rightarrow \phi \gamma$

2011 Data [Nucl. Phys. B 867(2012) 1-18]

- Mass resolution $\sim 90 \text{ MeV}/c^2$ dominated by LHCb Electromagnetic CALorimeter (ECAL) resolution.
- Many background contributions

$$\text{BR}(B^0 \rightarrow K^{*0} \gamma) = (4.33 \pm 0.15) \times 10^{-5}$$

$$\text{BR}(B_s \rightarrow \phi \gamma) = (5.7 + 2.1 - 1.8) \times 10^{-5}$$

HFAG 2010

- World best measurement of the ratio of branching fractions

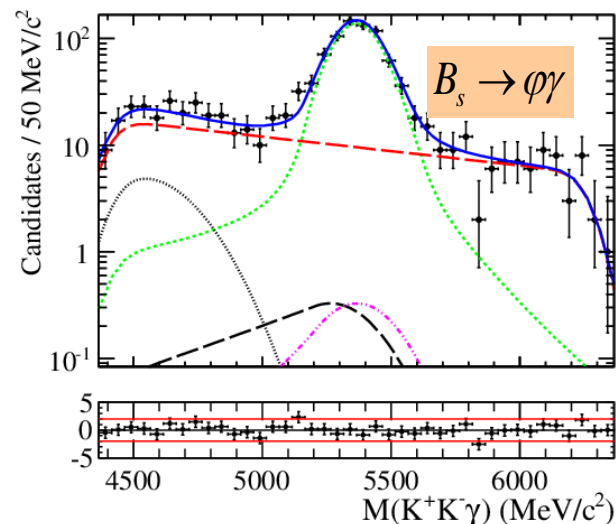
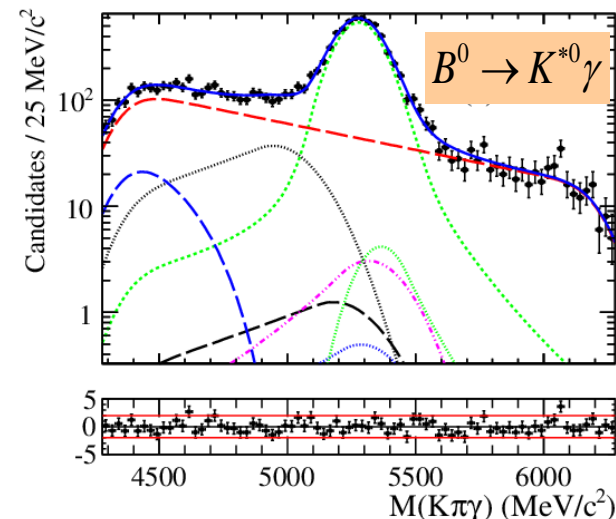
$$\frac{\mathcal{B}(B^0 \rightarrow K^{*0} \gamma)}{\mathcal{B}(B_s^0 \rightarrow \phi \gamma)} = 1.23 \pm 0.06 (\text{stat.}) \pm 0.04 (\text{syst.}) \pm 0.10 (f_s/f_d)$$

Theory predictions:

$$\frac{\mathcal{B}(B^0 \rightarrow K^{*0} \gamma)}{\mathcal{B}(B_s^0 \rightarrow \phi \gamma)} = 1.0 \pm 0.2$$

- And $\text{BR}(B_s \rightarrow \phi \gamma)$

$$\mathcal{B}(B_s^0 \rightarrow \phi \gamma) = (3.5 \pm 0.4) \times 10^{-5}$$



Rad. Decays at LHCb: $B^0 \rightarrow K^* \gamma$

2011 Data [Nucl. Phys. B 867(2012) 1-18]

- $B^0 \rightarrow K^* \gamma$

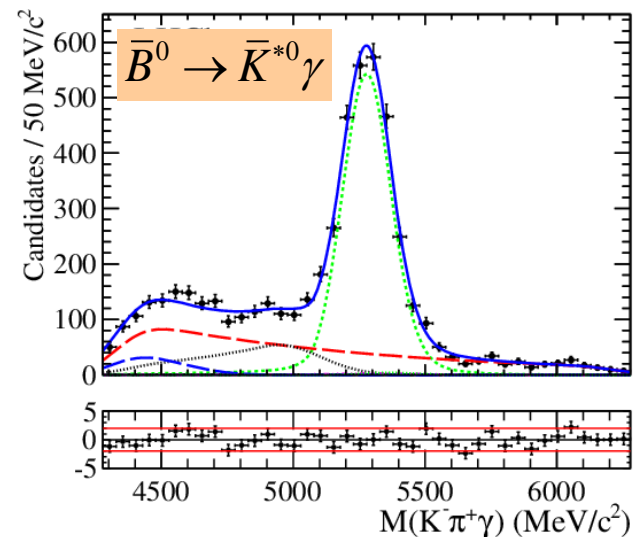
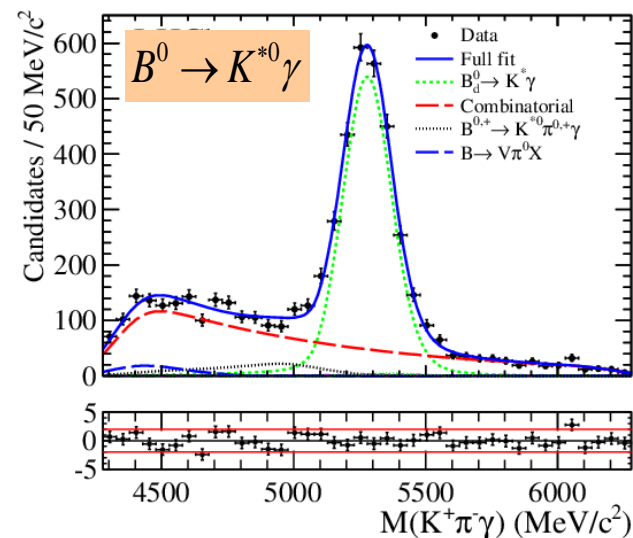
$$N_{B_0} + N_{\bar{B}_0} = 5300 \pm 100$$

- World best measurement of the direct CP asymmetry in $B_0 \rightarrow K^* \gamma$

$$\mathcal{A}_{CP}(B^0 \rightarrow K^{*0} \gamma) = (0.8 \pm 1.7 \text{ (stat.)} \pm 0.9 \text{ (syst.)})\%$$

Theory predictions:

$$\mathcal{A}_{CP} = -0.0061 \pm 0.0043$$



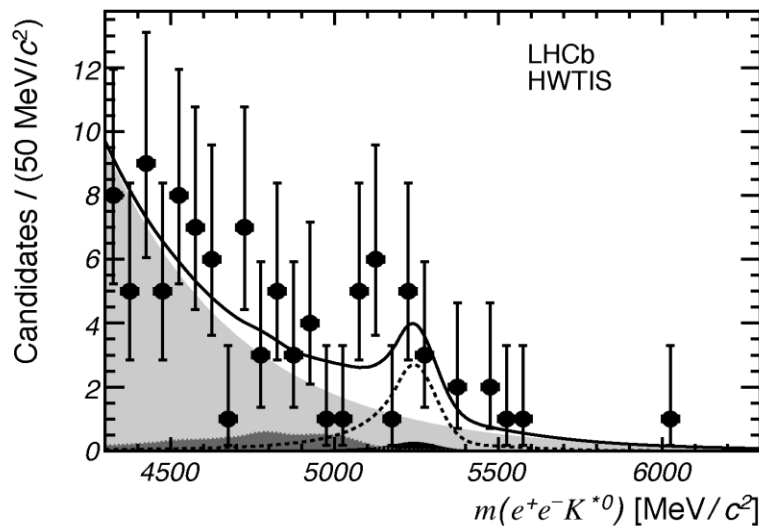
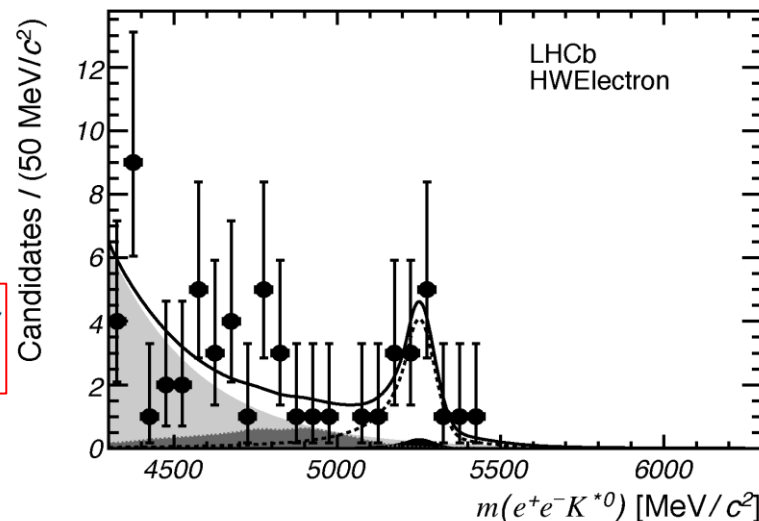
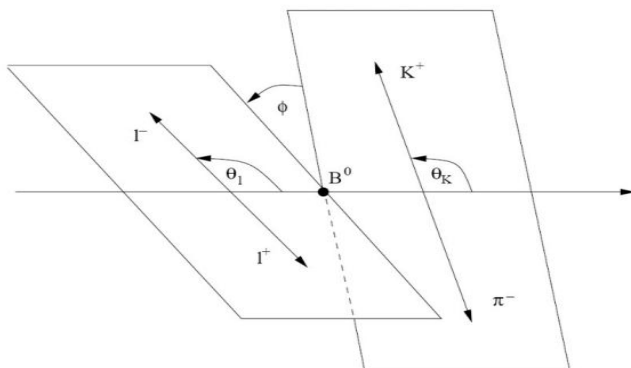
RD at LHCb: $B^0 \rightarrow K^* \gamma^* (e^+ e^-)$

2011 Data [J. High Energy Phys. 05 (2013) 159]

- The world best measurement of the BR of $B^0 \rightarrow K^* e^+ e^-$

$$\mathcal{B}(B^0 \rightarrow K^{*0} e^+ e^-)^{30-1000 \text{ MeV}/c^2} = (3.1^{+0.9}_{-0.8} {}^{+0.2}_{-0.3} \pm 0.2) \times 10^{-7}$$

- An update with 2012 data is ongoing
 - Perform an angular analysis to extract a virtual photon polarization



Photon polarization from $B \rightarrow K\pi\pi\gamma$

- The main goal is to extract photon polarization from $B \rightarrow K_{\text{res}}\gamma$ where $K_{\text{res}} \rightarrow K\pi\pi$
- The decay amplitude is given by

$$|A(B \rightarrow K_{\text{res}}\gamma \rightarrow P_1 P_2 P_3 \gamma)|^2 = |c_R|^2 |\mathcal{M}_R|^2 + |c_L|^2 |\mathcal{M}_L|^2$$

- The photon polarization

$$\lambda_\gamma \equiv \frac{|c_R|^2 - |c_L|^2}{|c_R|^2 + |c_L|^2}$$

- With $\lambda_\gamma = -1$ for B^- decays and $+1$ for B^+ decays
- This 3-body decay allow to access the γ polarization through the up-down asymmetry defined by the γ direction w.r.t the plane defined by the daughters of the kaon resonance

Photon polarization from $B \rightarrow K\pi\pi\gamma$

Kou et al, PRD83 (2011) 094007

Gronau et al, PRL88 (2002) 051802

- The differential decay rate, for a given intermediate resonance is given as:

$$\frac{d\Gamma(B \rightarrow K_{\text{res}}\gamma \rightarrow P_1 P_2 P_3 \gamma)}{ds ds_{13} ds_{23} d\cos\theta} \propto |\vec{\mathcal{J}}|^2 (1 + \cos^2\theta) + \lambda_\gamma 2 \text{Im} [\vec{n} \cdot (\vec{\mathcal{J}} \times \vec{\mathcal{J}}^*)] \cos\theta$$

Helicity amplitude

Photon direction

- And for multiple intermediate resonances, we have

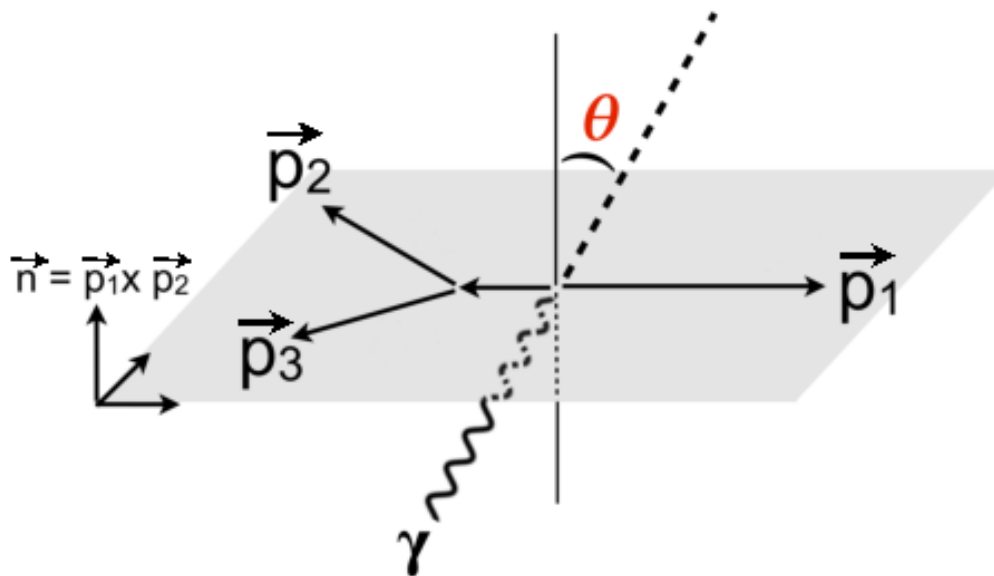
$$\frac{d\Gamma(\sum B \rightarrow K_{\text{res},i}\gamma \rightarrow P_1 P_2 P_3 \gamma)}{ds ds_{13} ds_{23} d\cos\theta} \propto \sum_{j=0,\text{even}}^4 a_j(s_{13}, s_{23}) \cos^j\theta + \sum_{j=1,\text{odd}}^3 \lambda_\gamma a_j(s_{13}, s_{23}) \cos^j\theta$$

Eq. 1

- To extract an up-down asymmetry, the photon direction should be well defined

Photon polarization

The photon direction



- We define:

$$\cos\theta \equiv -\frac{\vec{p}_\gamma}{|\vec{p}_\gamma|} \cdot \hat{n}$$

- With:

$$\hat{n} \equiv \frac{\vec{p}_1 \times \vec{p}_2}{|\vec{p}_1 \times \vec{p}_2|}$$

The up-down asymmetry: formalism

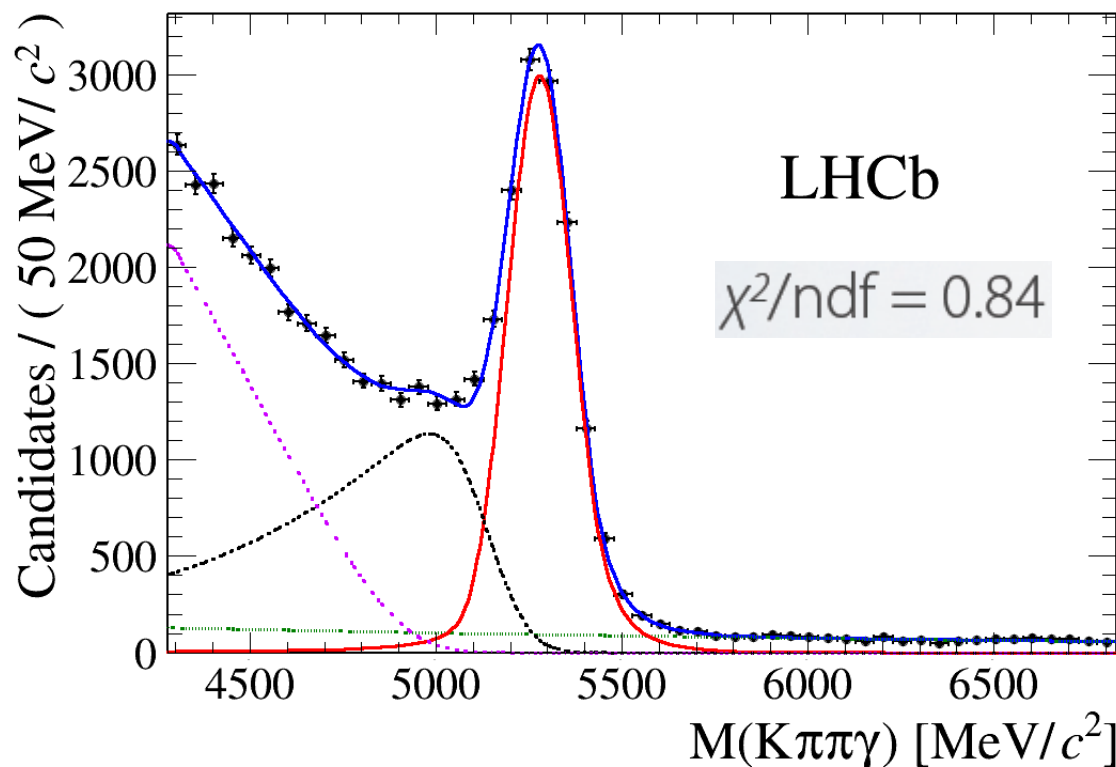
- The up-down asymmetry:
 - The number of events having the photon above the plane defined by the daughter of the Kaon resonance subtracted from those with the photon below the plane

$$\mathcal{A}_{UD} \equiv \frac{\int_0^1 d\cos\theta \frac{d\Gamma}{d\cos\theta} - \int_{-1}^0 d\cos\theta \frac{d\Gamma}{d\cos\theta}}{\int_{-1}^1 d\cos\theta \frac{d\Gamma}{d\cos\theta}} \stackrel{\text{single resonance}}{=} \frac{3}{4} \lambda_\gamma \frac{\int ds ds_{13} ds_{23} \text{Im} [\vec{n} \cdot (\vec{\mathcal{J}} \times \vec{\mathcal{J}}^*)]}{\int ds ds_{13} ds_{23} |\mathcal{J}|^2}$$

- \mathcal{A}_{UD} is proportional to λ_γ
- If J is known, i.e. \mathcal{A}_{UD} is known for a single resonance, the up-down asymmetry would allow to compute the photon polarization λ_γ

The $B \rightarrow K\pi\pi\gamma$ invariant mass spectrum

2011+2012 Data [Phys. Rev. Lett. **112**, 161801]



- Fit components:

Signal

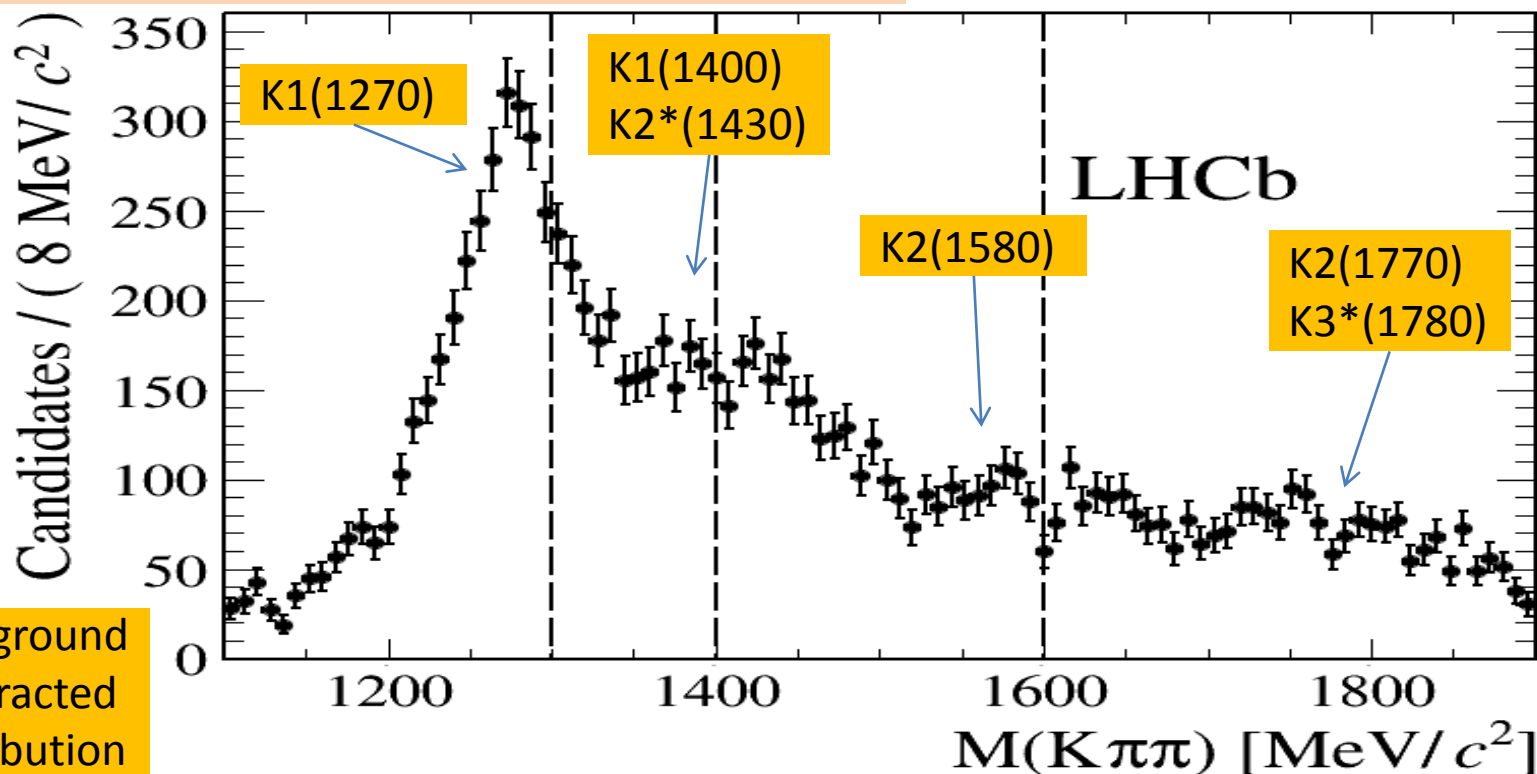
Part' reco' Missing π

Combinatorial

Partially reconstructed

The $K\pi\pi$ invariant mass Spectrum

2011+2012 Data [Phys. Rev. Lett. **112**, 161801]



- Can't isolate individual components without amplitude analysis
 - Asymmetry measurement need to be inclusive
- Up-down asymmetry can't be converted to photon polarization, only a significance w.r.t no-polarization is extracted (in each mass region)

The up-down asymmetry: measurement

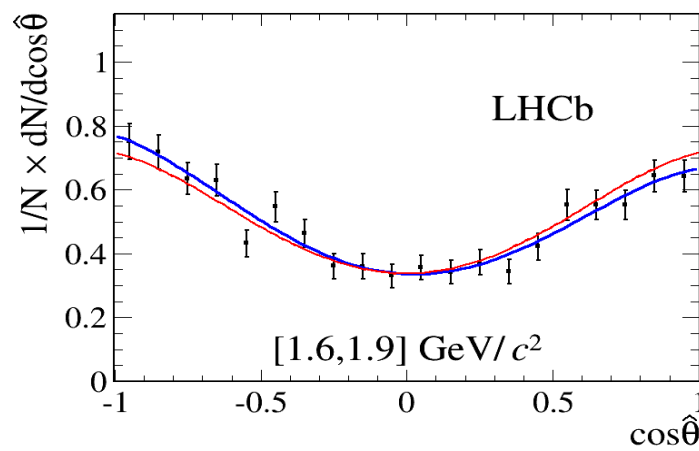
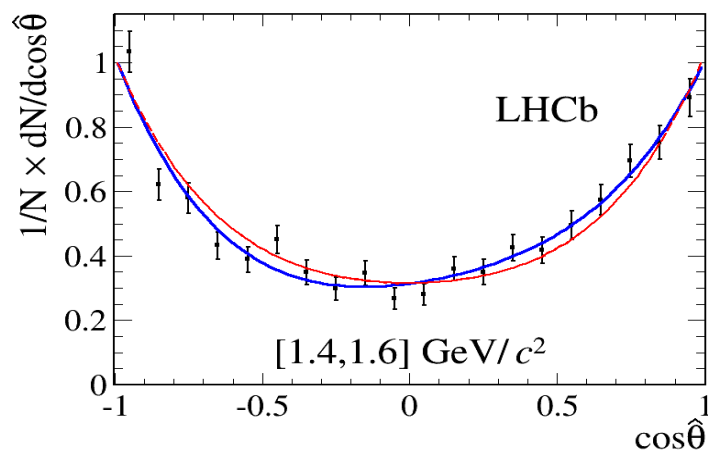
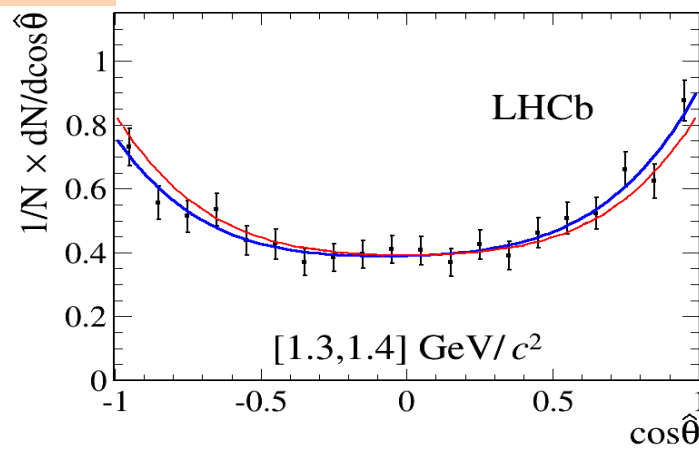
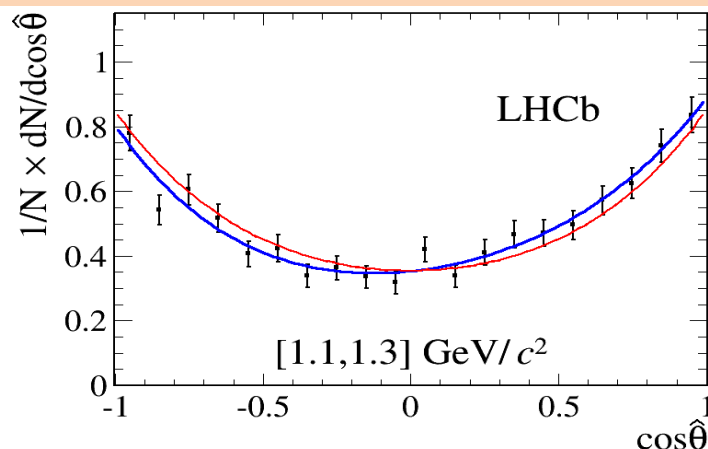
- A background-subtracted $\cos(\theta)$ distribution is extracted from each $K\pi\pi$ invariant mass bin
- $\cos(\theta)$ distribution is fitted (unbinned ML fit) with 4-th order polynomial normalized to unit area

$$f(\cos \theta; c_0 = 0.5, c_1, c_2, c_3, c_4) = \sum_{i=0}^4 c_i L_i(\cos \theta)$$

Eq. 2

cos (θ) distribution

2011+2012 Data [Phys. Rev. Lett. 112, 161801]



-Odd terms are put to zero: assume no photon polarization
-Fit of odd and even components

5.2 σ significance: taking into account correlations between errors

Up-down asymmetry: results

2011+2012 Data [Phys. Rev. Lett. **112, 161801]**

- The asymmetry is then given by: [from Eq. 1(slide 9) and Eq. 2 (slide 14)]

$$A_{UD} = c_1 - \frac{c_3}{4}$$

Obtained from the odd terms of the poly. fit

	[1.1, 1.3]	[1.3, 1.4]	[1.4, 1.6]	[1.6, 1.9]
c_1	6.3 ± 1.7	5.4 ± 2.0	4.3 ± 1.9	-4.6 ± 1.8
c_2	31.6 ± 2.2	27.0 ± 2.6	43.1 ± 2.3	28.0 ± 2.3
c_3	-2.1 ± 2.6	2.0 ± 3.1	-5.2 ± 2.8	-0.6 ± 2.7
c_4	3.0 ± 3.0	6.8 ± 3.6	8.1 ± 3.1	-6.2 ± 3.2
\mathcal{A}_{ud}	6.9 ± 1.7	4.9 ± 2.0	5.6 ± 1.8	-4.5 ± 1.9

- The quoted uncertainties contain systematic and statistical contributions

Conclusions I

- Radiative decays at LHCb: many important results
 - Direct CP asymmetry in $B^0 \rightarrow K^* \gamma$
 - Ratio of BR of $B^0 \rightarrow K^* \gamma$ / $B_s \rightarrow \phi \gamma$
 - The BR of $B^0 \rightarrow K^* e^+ e^-$
- Observation of photon polarization in $B \rightarrow K \pi \pi \gamma$
 - Up-down asymmetry
- And more to come..

Conclusions II

- LHCb has, for the first time, an evidence of the polarization of the photon with a 5.2σ significance
 - A dedicated amplitude analysis will help to translate this result to a value for the photon polarization

Or

- Input from theory is needed to translate this result into a value of the photon polarization λ_γ
 - How to derive the polarization over all the mass range
- Constrain effects of physics beyond the SM in $b \rightarrow s\gamma$ sector

Perspectives

- $B_s \rightarrow \phi \gamma$ time dependent analysis:
- In the time dependent decay rate, the photon polarization appears through two parameters, S and \mathcal{A}^Δ

$$\Gamma_B(t) \propto |A|^2 e^{-\Gamma t} \left[\cosh(\Delta\Gamma t / 2) - \mathcal{A}^\Delta \sinh(\Delta\Gamma t / 2) \pm C \cos(\Delta m t) \mp S \sin(\Delta m t) \right]$$

Defining:

$$\tan \psi = \frac{\overline{B} \rightarrow f^{CP} \gamma_R}{\overline{B} \rightarrow f^{CP} \gamma_L}$$

$$\mathcal{A}^\Delta = \sin(2\psi) \cos \phi$$

$$S = \sin(2\psi) \sin \phi$$

Where ϕ is the B mixing phase

- Which channels?
- $B_d \rightarrow K^*(K_S \pi^0) \gamma$
 - $\Delta\Gamma \sim 0 \rightarrow$ the sinh term cancels and we have only access to S
 - Done at Babar but not possible at LHCb $S = 0.9 \pm 1.0 \pm 0.2$ (Babar, $1.1 \text{ GeV} < m_{K_S \pi^0} < 1.8 \text{ GeV}$)
- $B_s \rightarrow \phi(K^+ K^-) \gamma$
 - $\Delta\Gamma_s$ is not negligible \rightarrow the dominant term is the sinh $\rightarrow \mathcal{A}^\Delta$ can be measured
 - Work in progress: untagged time dependent analysis to extract \mathcal{A}^Δ

AM I AN UNCLEAR
COMMUNICATOR?



SIX
O'CLOCK.



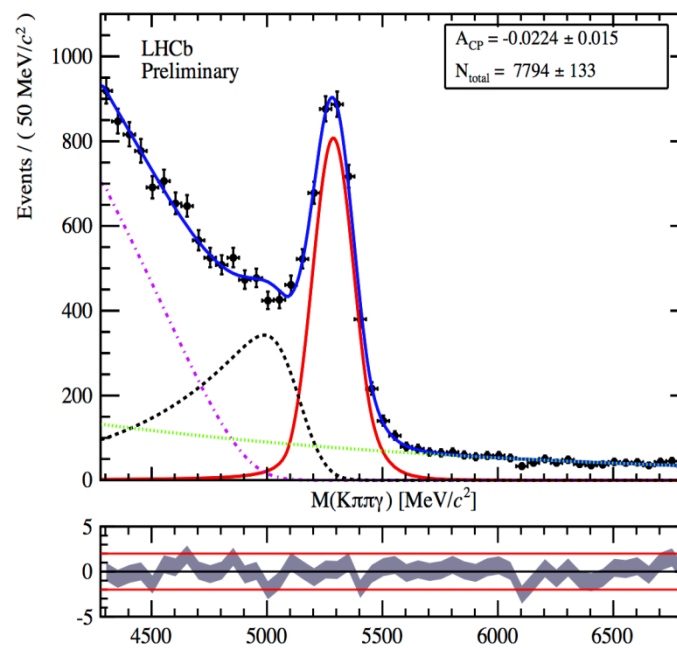
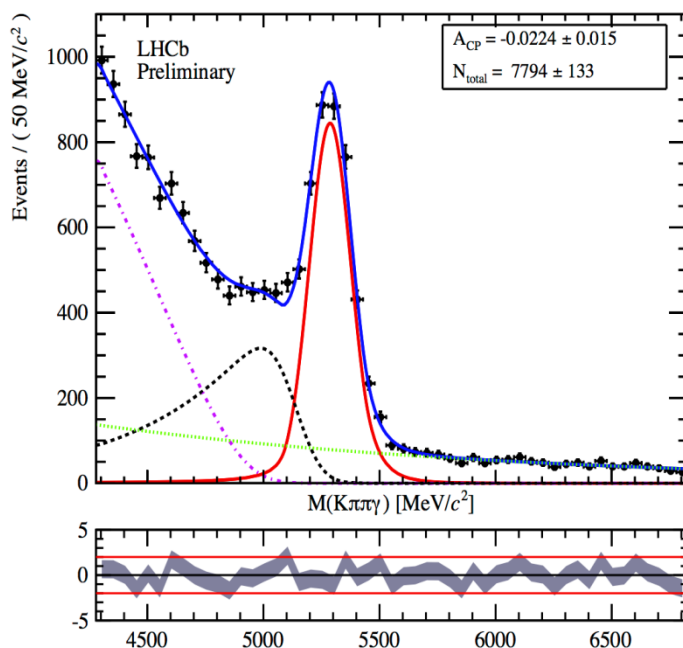
spares

CP asymmetry in $B \rightarrow K\pi\pi\gamma$

2011+2012 Data [LHCb-PAPER-2014-001]

- A CP asymmetry measurement can also be conducted:

$$A_{CP}^{\text{raw}} = -0.022 \pm 0.015$$



CP asymmetry in $B \rightarrow K\pi\pi\gamma$

- A CP asymmetry measurement can also be conducted:

$$\mathcal{A}_{CP}^{\text{raw}} = -0.022 \pm 0.015$$

- The raw ACP obtained from the fit is corrected to obtain the physical CP asymmetry
 - Charged B meson production asymmetry
 - Particle interaction with matter (cross-section) asymmetry (K^+ vs K^-)
 - Geometrical detection asymmetries
- Corrections are extracted from control channels and from data corresponding to different magnet polarities

CP asymmetry in $B \rightarrow K_{\text{res}} \gamma$: results

2011+2012 Data [LHCb-PAPER-2014-001]

- $B \rightarrow K \pi \pi \gamma$ has been observed by LHCb: world largest sample

$$N = 13876 \pm 153 \text{ events}$$

- CP asymmetry (consistent with 0) has been measured for the first time in $B \rightarrow K \pi \pi \gamma$

$$A_{CP} = -0.007 \pm 0.0155(\text{stat.})^{+0.012}_{-0.011}(\text{syst.})$$