MONO-HIGGS DETECTION OF DARK MATTER AT THE LHC

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In collaboration with Tongyan Lin and Lian-Tao Wang arXiv:1402.7074 [hep-ph]



(Also, see Carpentar et al., arXiv:1312.2592 and Petrov et al., arXiv:1311.1511









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We consider decay to bb or $\gamma\gamma$.

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Effective Field Theory (EFT)

- Dim. 6 : Direct Z+DM coupling. Direct detection and invisible Z width are very constraining.
 - <u>Dim. 7</u>: Mild constraints on the cutoff. Softer momentum transfer dependence.
- Dim. 8: Stronger constraints on the cutoff. Harder momentum transfer dependence.







<u>95% CL</u> <u>Constraints on</u> <u>Cutoff from</u> <u>Mono-Higgs</u>

<u>Red Lines</u>: Lower bound from naive requirement.

<u>Blue Regions</u>: Excluded regions after "unitarizing".

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Small mixing Z-Z' angle. Slightly alters Z mass.

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Yet another form of Higgs + MET is Z+h production via Z' resonance!

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Experimental Constraints



$$\rho_0 = 1 + \epsilon^2 \left(\frac{M_{Z'}^2 - M_Z^2}{M_Z^2} \right) \le 1.0009$$

Experimental Constraints











$h \rightarrow b\bar{b} + MET$

- □ Large branching fraction.
- **Dominant irreducible background = Z+h production.**
- \Box Dominant reducible background = tt and Zbb.
- $\Box tt$ dominates at large energies.
- □ Consider MET cuts from 120-250 GeV.

 $h \rightarrow \gamma \gamma + MET$

- Very small branching fraction.
- **Dominant irreducible background = Z+h production.**
- \Box Dominant reducible background = $Z\gamma\gamma$.
- □ Increased sensitivity at 14 TeV.
- □ Consider MET cuts from 100-250 GeV.

Effective Field Theory (EFT)

 $\frac{\mathbf{Dim. 7:}}{\frac{1}{\Lambda^3}} \bar{X} \gamma^{\mu\nu} X \times i \left[(D_{\mu}H)^{\dagger} D_{\nu}H - \text{h.c.} \right] \square$

Dim. 8: $\frac{1}{\Lambda^4} \frac{1}{2} (\phi^{\dagger} \partial^{\mu} \phi + \text{h.c.}) (B_{\nu\mu} H^{\dagger} D^{\nu} H + \text{h.c.})$ $\frac{1}{\Lambda^4} \bar{X} \gamma^{\mu} X \left(W^a_{\nu\mu} H^{\dagger} t^a D^{\nu} H + \text{h.c.} \right)$



Rate is smaller when DM couples to 2 Higgs

These same operators give Mono-Z/W. This is weaker due to schannel Higgs production.



EFT Unitarity

Mono-Higgs signal comes from high momentum transfer region.

- Constraints on cutoff often of order 100 GeV (if even). Can see that momentum transfer exceeds this easily.
- We should worry about the unitarity of our EFT description!



Momentum Transfer at 8 TeV, for10 GeV (gray) or 100 (black) GeV DM, from Mono-Higgs process of Dim.-8 fermionic operator. Zh production (red) is the irreducible SM background.

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(the naive limit)

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Don't take too literally. Only illustrates (conservatively) error in assuming 1 operator is relevant.



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Additional source of Higgs+MET: $\Gamma_{Z' \to hZ} = (g_z \sin^2 \beta)^2 \frac{|p|}{24\pi} \left(\frac{|p|^2}{M_{Z'}^2} + 3\frac{M_Z^2}{M_{Z'}^2}\right)$

Mono-Higgs Signal

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 $\frac{\text{For small tan }\beta:}{\text{For large tan }\beta:} \text{ rate } \sim \tan^4 \beta$ $\frac{\text{For large tan }\beta:}{\text{Increases for smaller }M_{Z'}}$

Dark Matter Coupling

- **Want order 1 BF of pseudoscalar to DM**
- Singlet-Doublet Fermionic Dark Matter (Cheung and Sanford arXiv 1311.5896 [hep-ph])
- □ Introduce two SU(2) doublet Weyls with hypercharge $\pm 1/2$, D₂ and D₁, respectively
- Introduce total gauge singlet Weyl S
- Assume all are neutral under $U(1)_{Z'}$

 $-\mathcal{L} \supset \frac{1}{2}M_S^2 S^2 + M_D D_1 D_2 + y_1 S D_1 \Phi_d + y_2 S \Phi_d^{\dagger} D_2 + \text{h.c.}$

