

THE HIGGS BOSON THE MEASUREMENT OF ITS MASS, WIDTH, AND SPIN

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Representing the CMS and ATLAS collaborations



Is it really
Lincoln?



Can the study of
the properties of
the Higgs boson
reveal cracks in
the SM?

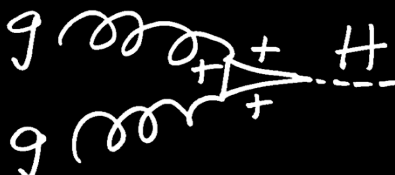
Higgs boson
mass, width, and
spin



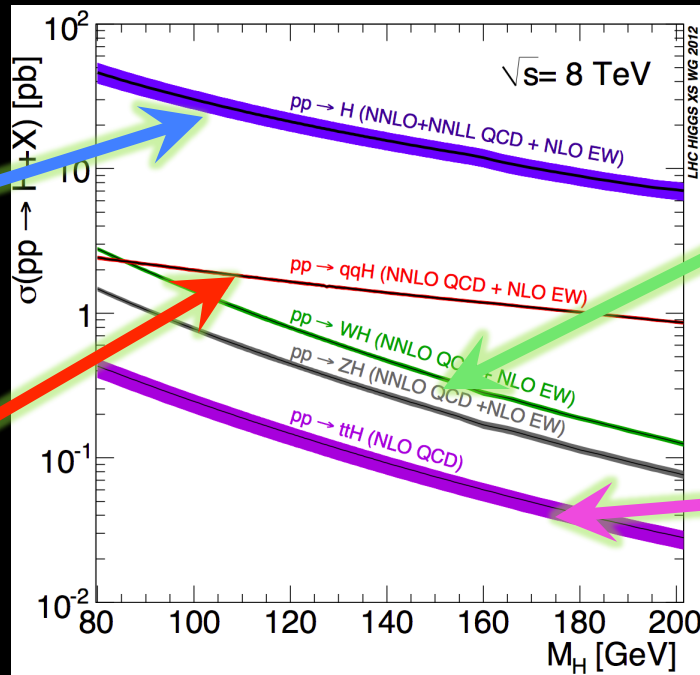
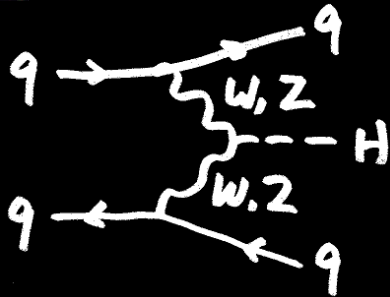
Higgs Production and decay



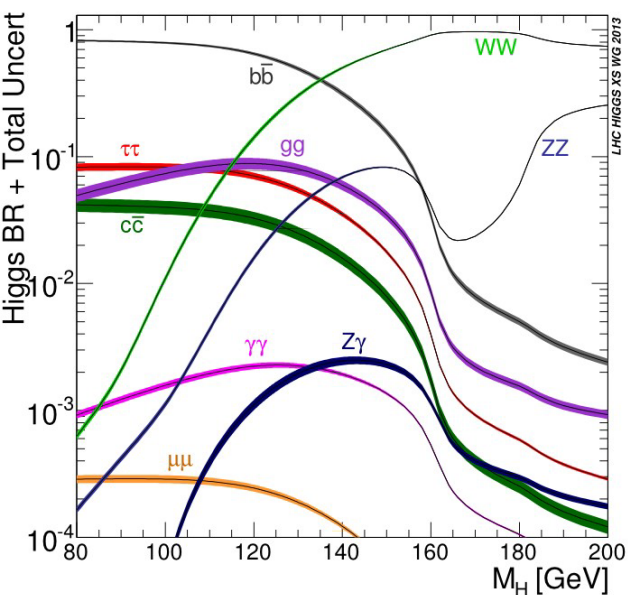
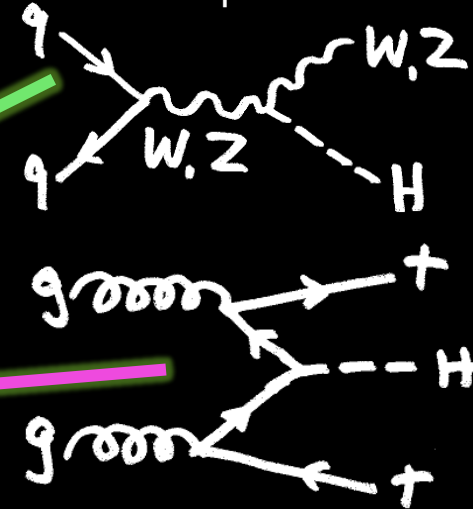
Gluon fusion



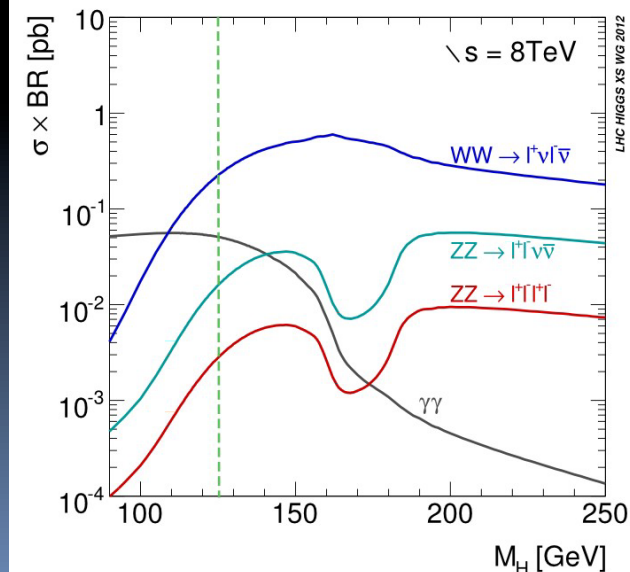
Vector Boson Fusion



Associated production



- The LHC has measured both fermionic ($\tau\tau, b\bar{b}$) and bosonic ($\gamma\gamma, WW, ZZ$) Higgs final states.
- Bosonic final states are currently providing mass, spin and width measurements





Higgs mass measurement

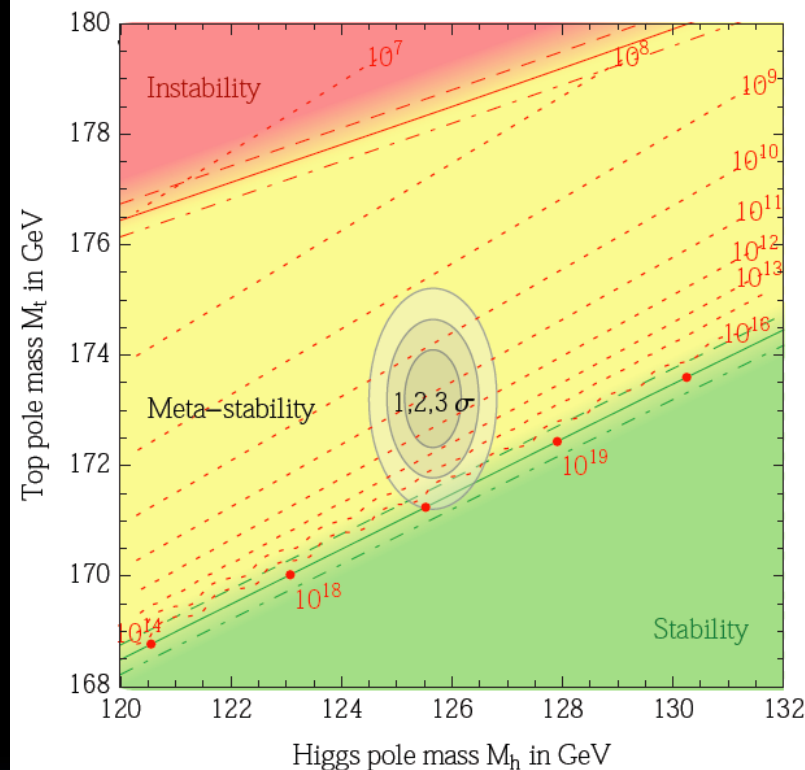


- M_H is a fundamental parameter not predicted by the SM
- Precision measurements of the Higgs mass provide important constraints.
- In the SM

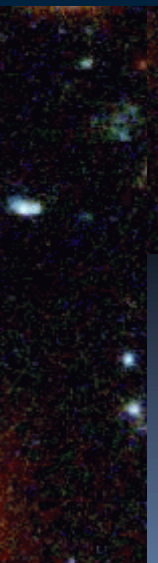
$$M_H^2 = 2\lambda_{SM} v^2$$

We must check this relation by measuring LHS and RHS independently

- Implication on the stability of the Higgs potential



Buttazzo, Degrandi, Giardino,
Giudice, Salab, Salvio, Strumia





Higgs mass measurement

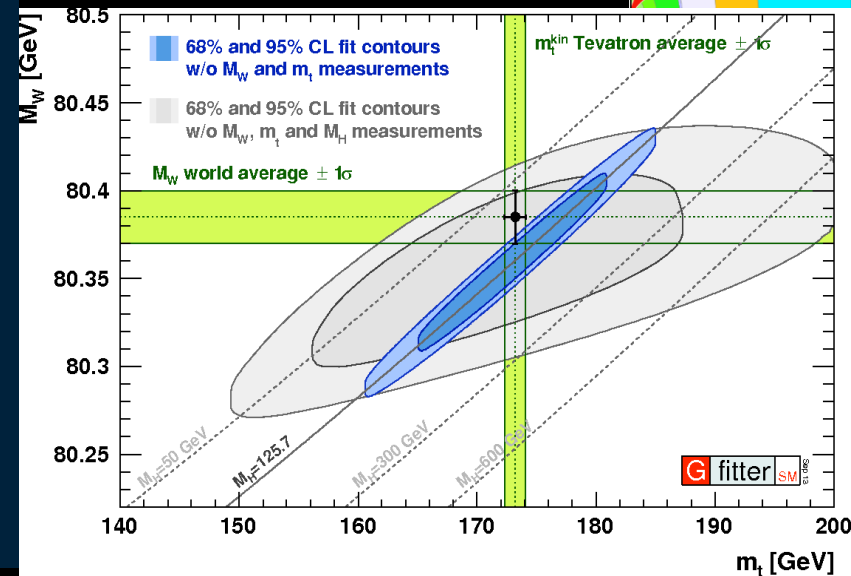


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- In the SM

$$M_H^2 = 2\lambda_{SM} v^2$$

We must check this relation by measuring LHS and RHS independently

- Implication on the stability of the Higgs potential
- Over constraining the EW fits
- In SUSY



$$m_h^2 = m_Z^2 \cos^2 2\beta + \frac{3G_F m_t^4}{\sqrt{2}\pi^2} \ln \frac{\Delta_S^2}{m_t^2},$$

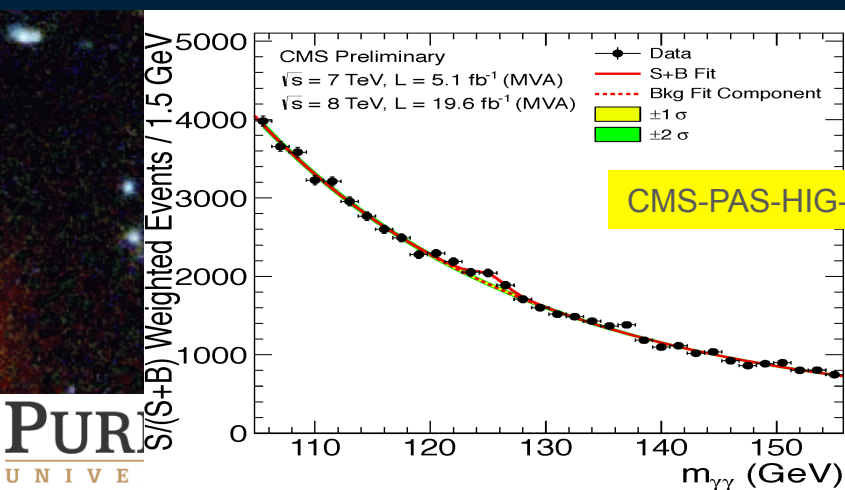
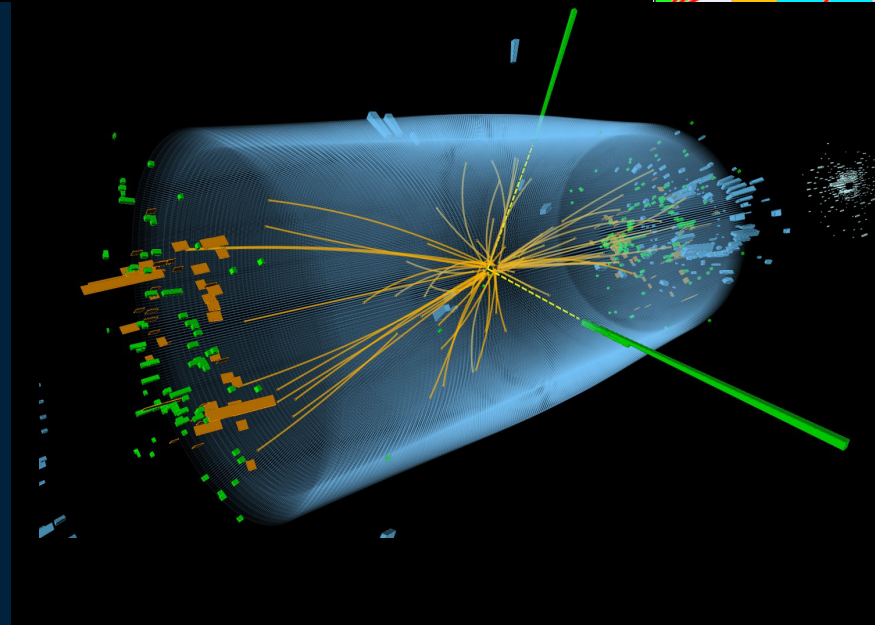
- LHC Measurement strategy
 - Use high resolution channels: $H \rightarrow \gamma\gamma$ and $H \rightarrow 4l$ which have a mass resolution of few %
- Challenges :
 - Maintain excellent mass-resolution in high-pile-up environment



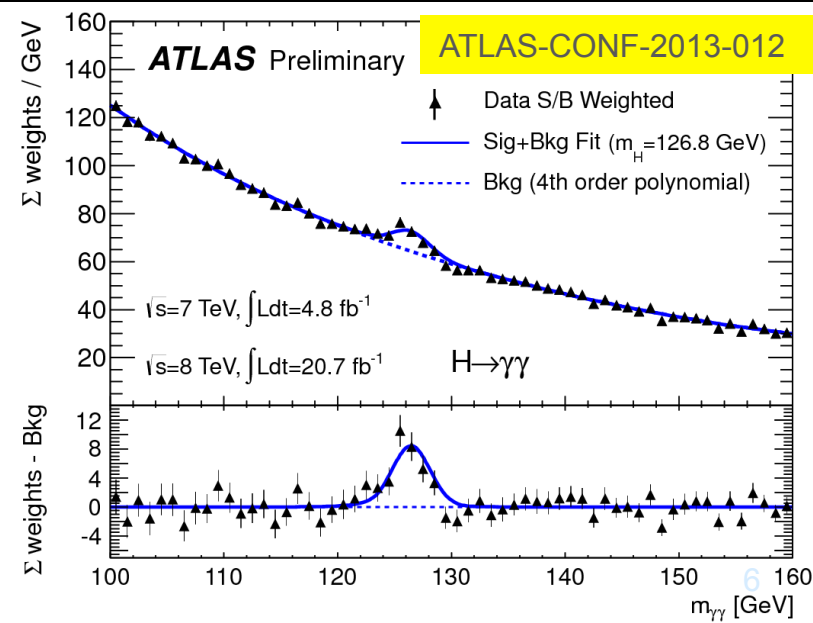
H $\rightarrow\gamma\gamma$ Mass Measurement



- Clean signature:
 - Two isolated, high p_T photons.
 - Small branching-fraction: $\sim 0.2\%$.
 - Excellent mass resolution: 1-2%.
 - Large background from QCD processes: $S/B \sim 1/1 \div 1/20$
- Strategy for analysis
 - Events categorized according to photon resolution and kinematics.
 - Exclusive channels targeting VBF and associated production.
 - Background modeled with polynomials or falling power-law or exponentials



2014





H $\rightarrow\gamma\gamma$ Mass Measurement

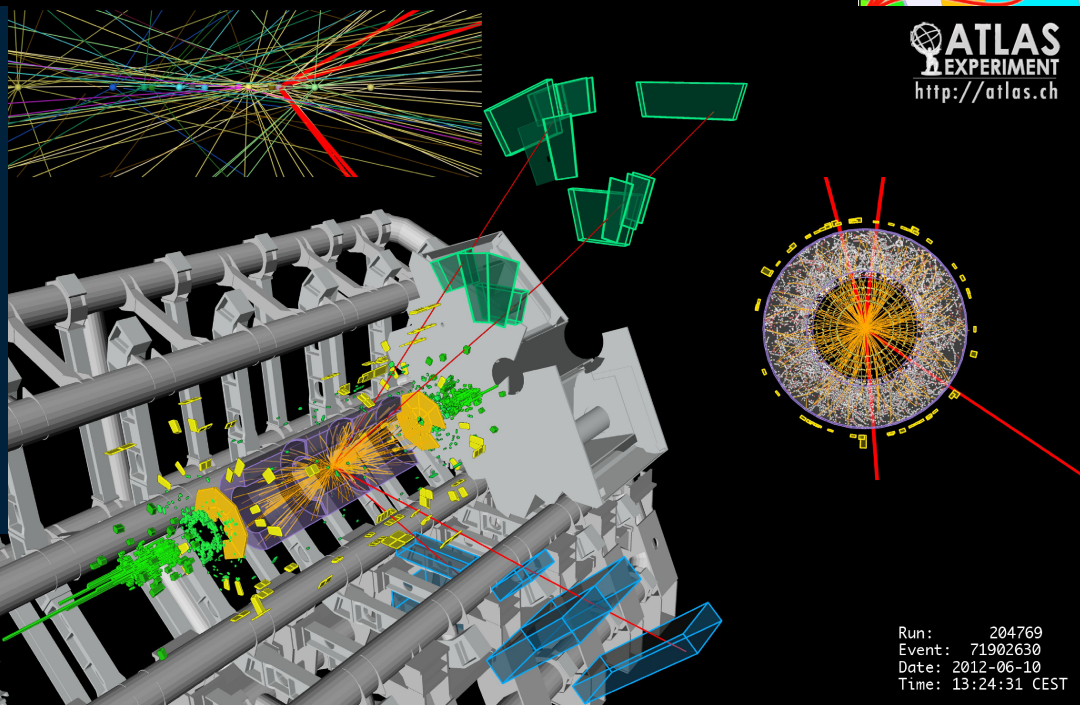


- Energy measurement:
 - CMS: e/γ energy estimated using multivariate regression
 - ATLAS: weighted sum of energy deposits in the different calorimeter layer
 - Energy scale and resolution validated with $Z \rightarrow ee$, $W \rightarrow e\nu$, J/ψ , and Upsilon resonances and corrected with MC input for e/γ differences
 - Additional smearing applied to MC to match the resolution in data
- Systematic uncertainty on mass scale currently $\sim 0.5\%$ but expected to improve considerably for legacy Run 1 results.

	ATLAS		CMS	
Mass [GeV]	$126.8 \pm 0.2(stat) \pm 0.7(sys)$		$125.4 \pm 0.5(stat) \pm 0.6(sys)$	
	Z $\rightarrow ee$ E-scale	0.3%	Residual non-linearity from Z $\rightarrow ee$ to H $\rightarrow \gamma\gamma$	0.4%
	Material modeling	0.3%	Material modeling	0.24%
	Pre-sampler scale	0.1%		
	Other	0.32%		
	Total	0.55%		0.47%



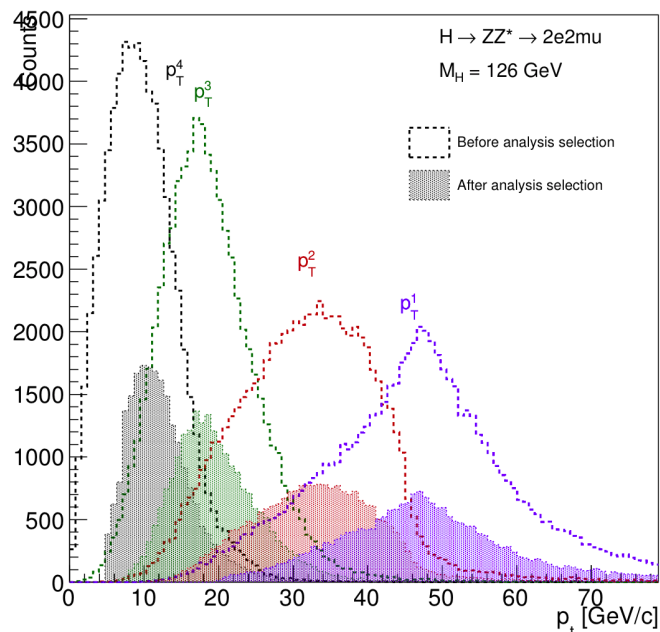
H → 4l Mass Measurement



Run: 204769
Event: 71902630
Date: 2012-06-10
Time: 13:24:31 CEST

- Golden channel:
 - Four isolated leptons.
 - Small branching fraction: $\sim 10^{-4}$
- Extremely pure: S/B $\sim 2:1$
- Challenges :
 - Maximize acceptance for low p_T leptons
- Precise calibration of lepton p_T scale

CMS Simulation, $\sqrt{s} = 8$ TeV

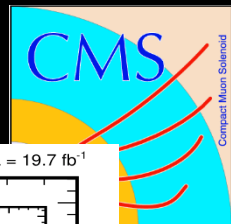


Analysis strategy:

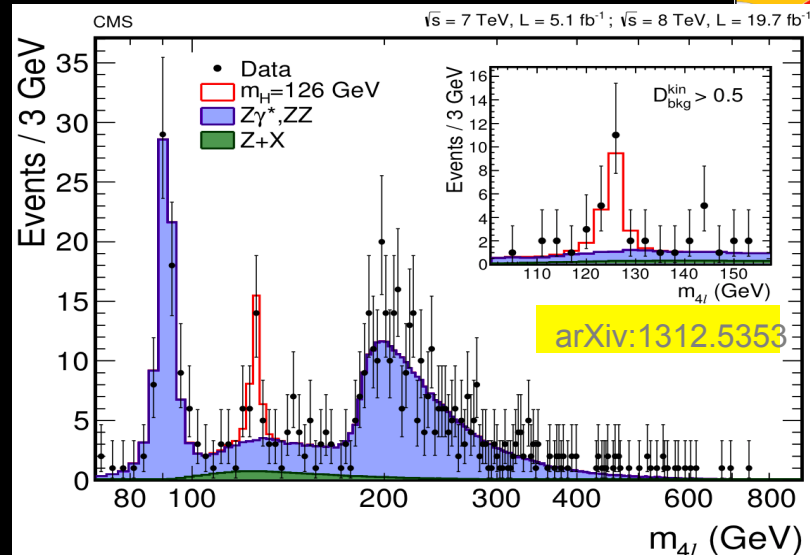
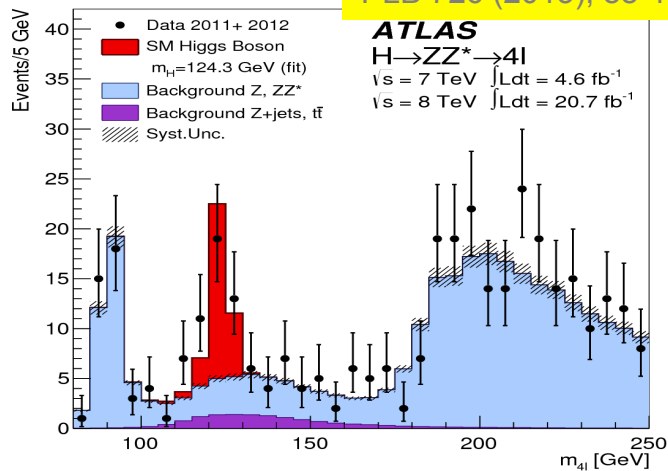
- CMS
 - Use $m(4l)$ vs kin. discriminant (KD) for S/B separation and event-by-event mass errors (estimated from lepton momentum errors)
- ATLAS:
 - Use $m(4l)$ for S/B separation and categorize events into VBF-like, VH-like, and untagged



H → 4l Mass Measurement

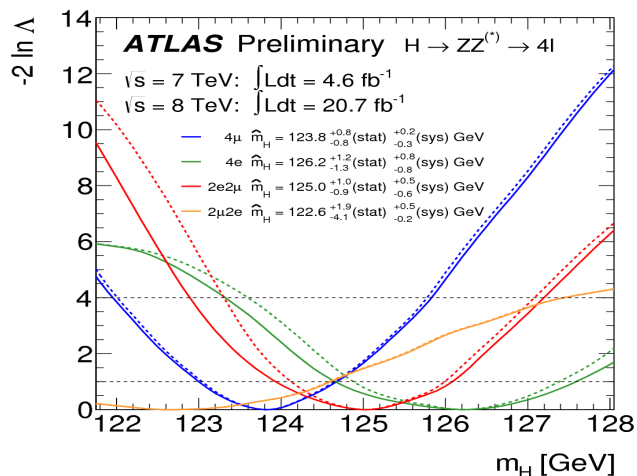


PLB 726 (2013), 88-119

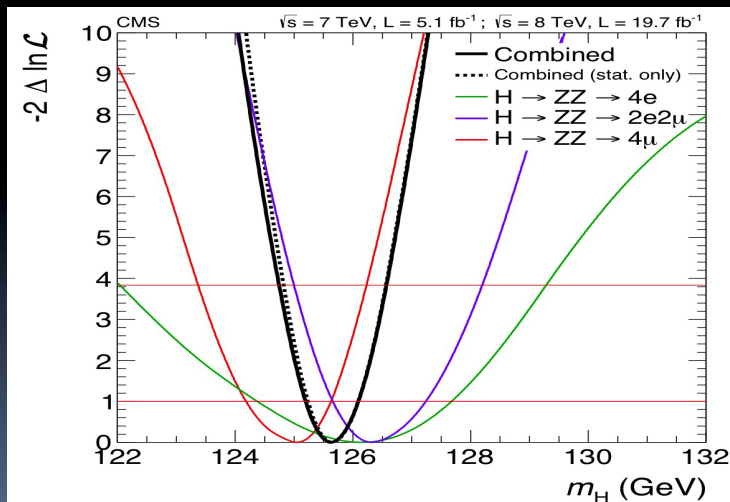


- ATLAS: 1D fit to m_{4l} . Using kinematic constraint to Z_1 candidate.

- CMS: 3D fit to m_{4l} , KD, and $\sigma(m_{4l})/m$



$$124.3^{+0.6}_{-0.5}(\text{stat}) \pm 0.7^{+0.5}_{-0.3}(\text{sys})$$



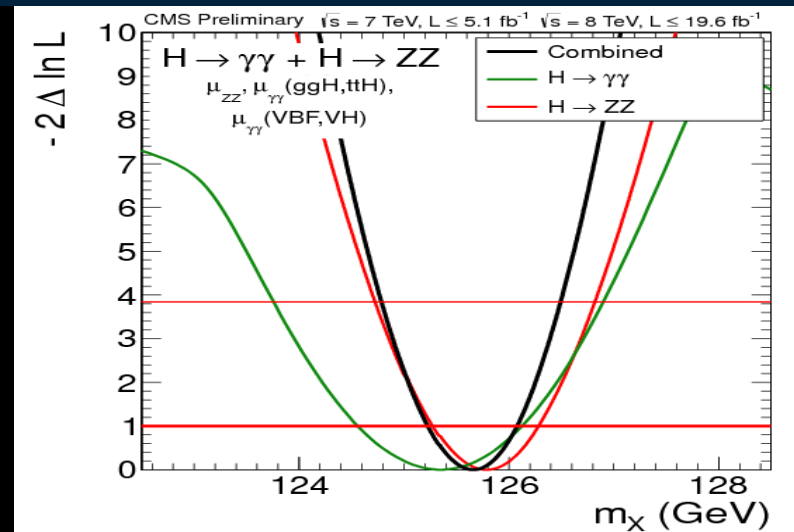
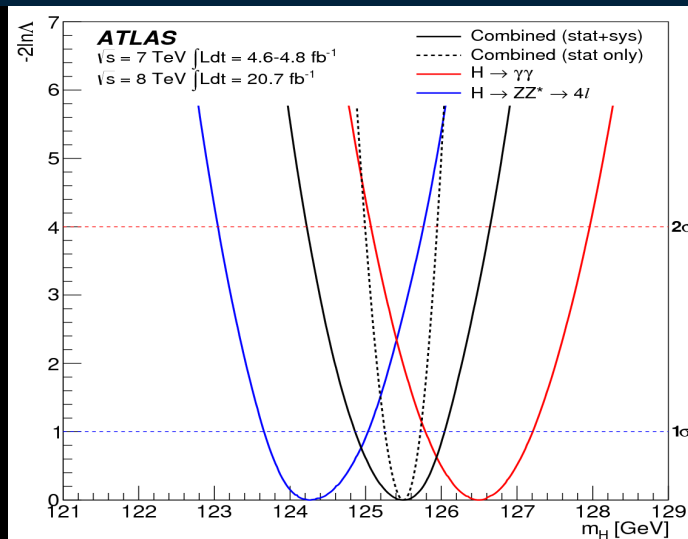
$$125.6 \pm 0.4(\text{stat}) \pm 0.2(\text{sys})$$



Mass Measurement: Combination



- Mass measurements $\gamma\gamma$ and $4l$ channels can be combined (under the hypothesis that the same state decays in both modes).

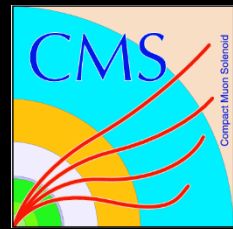


	ATLAS	CMS (new $4l$ measurement not used)
Mass [GeV]	$125.5 \pm 0.2(stat) \pm_{-0.6}^{+0.5}(sys)$	$125.6 \pm 0.3(stat) \pm 0.3(sys)$

- Precision expected to improve since both CMS&ATLAS will use updated calibrations and improved understanding of energy/momentum scale uncertainties for Run 1 legacy papers
- LHC combination expected later this year with precision possibly reaching $\approx 0.2\%$.



Spin and Parity



- In the Standard Model the Higgs boson is a 0^+ state.
 - This must be verified experimentally
- Most general form of spin 1 and 2 scattering amplitudes have many terms

$$\begin{aligned}
 A(X \rightarrow V_1 V_2) = \Lambda^{-1} & \left[2g_1^{(2)} t_{\mu\nu} f^{*(1)\mu\alpha} f^{*(2)\nu\alpha} + 2g_2^{(2)} t_{\mu\nu} \frac{q_\alpha q_\beta}{\Lambda^2} f^{*(1)\mu\alpha} f^{*(2)\nu\beta} + g_3^{(2)} \frac{\tilde{q}^\beta \tilde{q}^\alpha}{\Lambda^2} t_{\beta\nu} \left(f^{*(1)\mu\nu} f_{\mu\alpha}^{*(2)} + f^{*(2)\mu\nu} f_{\mu\alpha}^{*(1)} \right) \right. \\
 & + g_4^{(2)} \frac{\tilde{q}^\nu \tilde{q}^\mu}{\Lambda^2} t_{\mu\nu} f^{*(1)\alpha\beta} f_{\alpha\beta}^{*(2)} + m_V^2 \left(2g_5^{(2)} t_{\mu\nu} \epsilon_1^{*\mu} \epsilon_2^{*\nu} + 2g_6^{(2)} \frac{\tilde{q}^\mu q_\alpha}{\Lambda^2} t_{\mu\nu} (\epsilon_1^{*\nu} \epsilon_2^{*\alpha} - \epsilon_1^{*\alpha} \epsilon_2^{*\nu}) + g_7^{(2)} \frac{\tilde{q}^\mu \tilde{q}^\nu}{\Lambda^2} t_{\mu\nu} \epsilon_1^* \epsilon_2^* \right) \\
 & \left. + g_8^{(2)} \frac{\tilde{q}_\mu \tilde{q}_\nu}{\Lambda^2} t_{\mu\nu} f^{*(1)\alpha\beta} \tilde{f}_{\alpha\beta}^{*(2)} + m_V^2 \left(g_9^{(2)} \frac{t_{\mu\alpha} \tilde{q}^\alpha}{\Lambda^2} \epsilon_{\mu\nu\rho\sigma} \epsilon_1^{*\nu} \epsilon_2^{*\rho} q^\sigma + \frac{g_{10}^{(2)} t_{\mu\alpha} \tilde{q}^\alpha}{\Lambda^4} \epsilon_{\mu\nu\rho\sigma} q^\rho \tilde{q}^\sigma (\epsilon_1^{*\nu} (q\epsilon_2^*) + \epsilon_2^{*\nu} (q\epsilon_1^*)) \right) \right], \quad (18)
 \end{aligned}$$

- Until there is enough data, ATLAS and CMS are testing alternative hypothesis against SMH (0^+).
- Keep only dim-4 terms ($g_1 = g_5 \neq 0$): Graviton-like “couplings” (2^+m).
- Most sensitive channels:
 - $H \rightarrow ZZ \rightarrow 4l$, $H \rightarrow WW \rightarrow 2l2\nu$ and $H \rightarrow \gamma\gamma$
- $H \rightarrow ZZ \rightarrow 4l$ Extremely rich channel in terms of angular information.
 - Best suited to study spin and parity



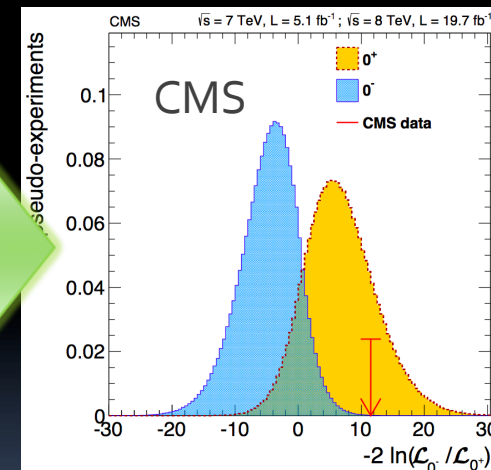
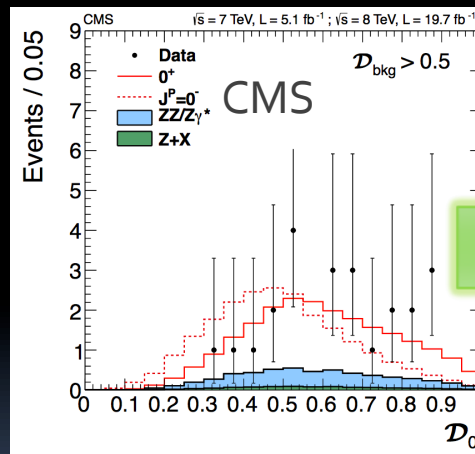
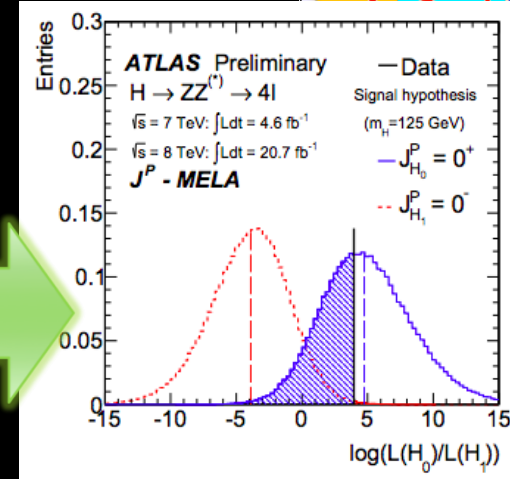
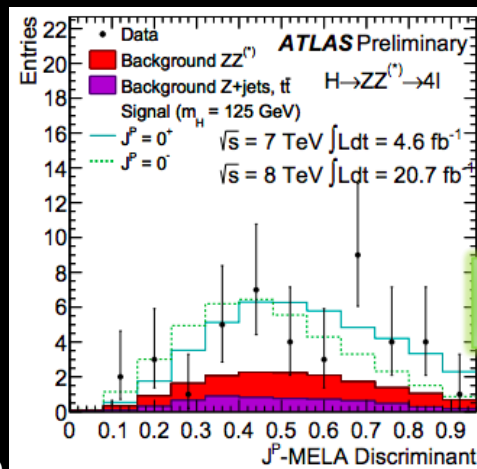
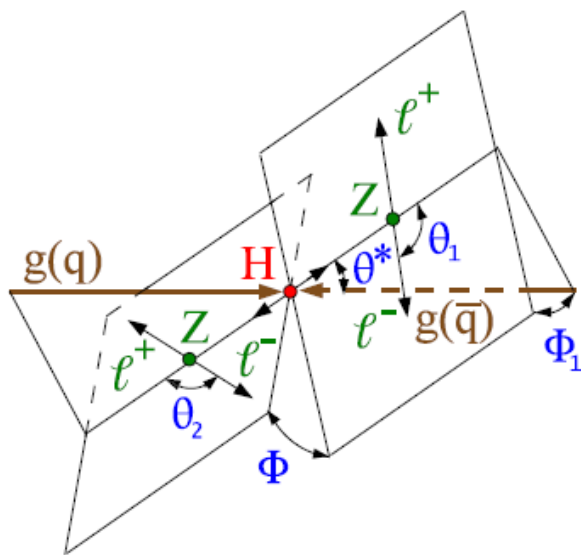
J^P in $H \rightarrow ZZ \rightarrow 4\ell$



- Use discriminants from decay angles and invariant masses

$$\mathcal{D}_{J^P} = \left[1 + \frac{\mathcal{P}_{J^P}^{\text{kin}}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell})}{\mathcal{P}_{0^+}^{\text{kin}}(m_{Z_1}, m_{Z_2}, \vec{\Omega} | m_{4\ell})} \right]^{-1}$$

- Profiled likelihood ratio test statistic.
 - ATLAS: template fit of BDT score distribution.
 - CMS: 2D fit of superKD($m_{4\ell} | x_{KD}$) vs KD(J^P)



	ATLAS	CMS
CL_s	0.37%	0.09%
$P(0^+)$	0.2σ	-0.9σ
$P(0^-)$	2.8σ	3.6σ

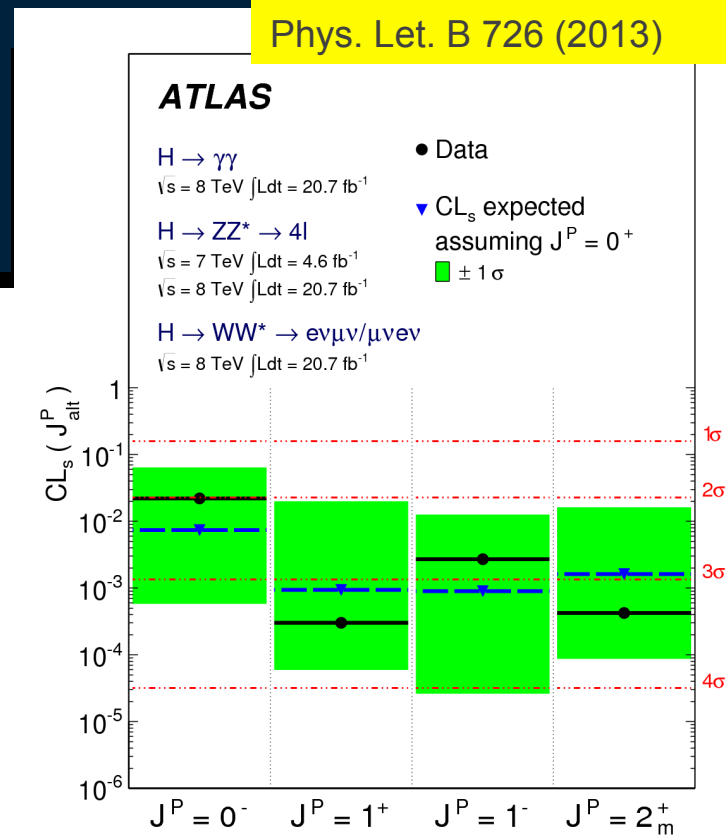
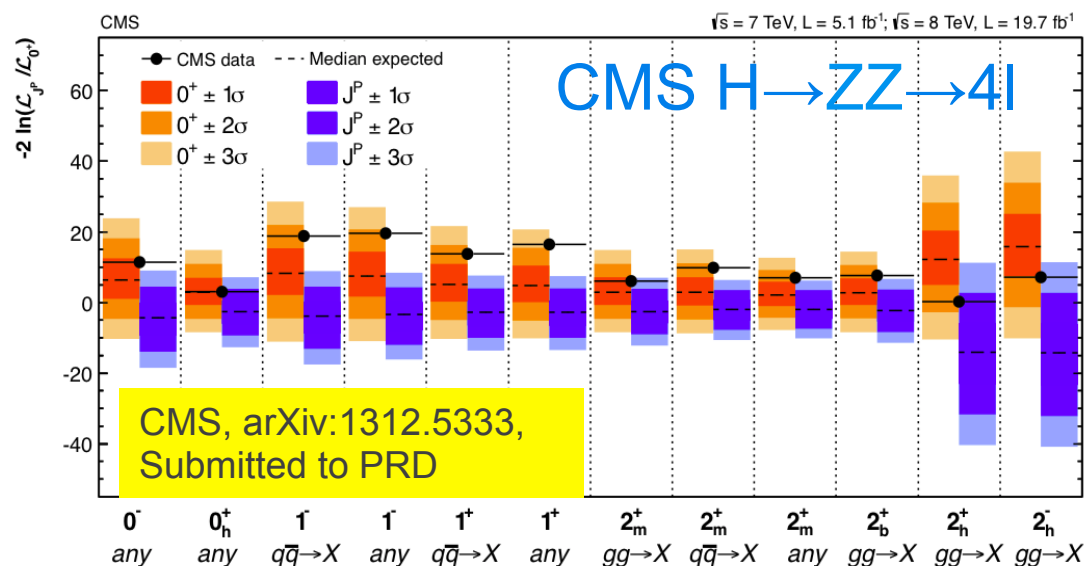
$$CL_s(J_{\text{alt}}^P) = \frac{p_0(J_{\text{alt}}^P)}{1 - p_0(0^+)}$$



Summary J^{PC} measurements

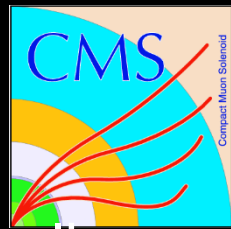


- CMS: 8 alternative J^{CP} hypotheses tested, for spin 0, 1 and 2, including production-independent analysis.
- ATLAS: 4 alternative J^{CP} hypotheses tested for spin 0, 1 and 2, considering qq, gg initiated productions (and mixtures).
- Data favors SM hypothesis.
 - Exclusion of spin 1 and graviton like resonance at $>3\sigma$
 - 0^- excluded at $> 2\sigma$ level





J^P in $H \rightarrow ZZ \rightarrow 4l$



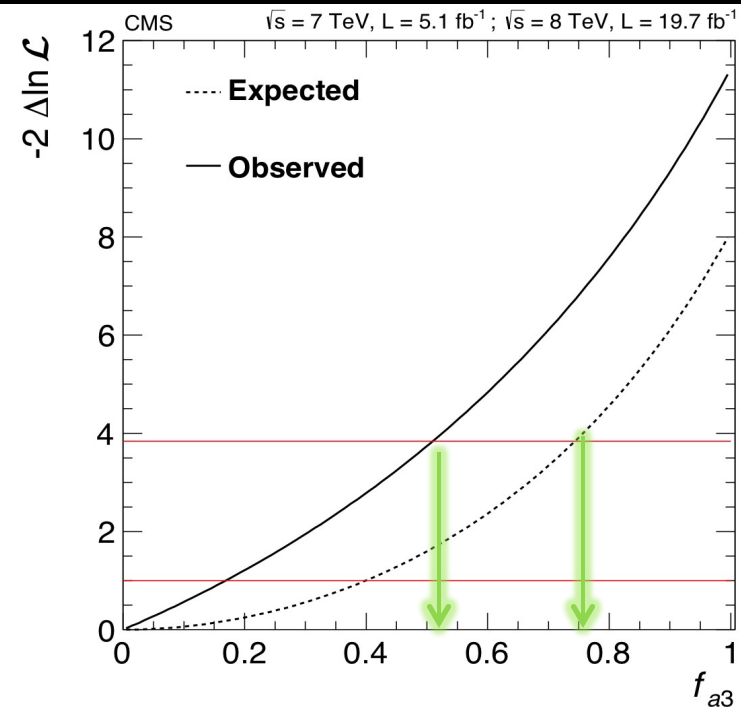
- Spin/parity studies will eventually constrain anomalous coupling parameters in the Higgs sector.
- First example, in the CMS $H \rightarrow ZZ \rightarrow 4l$ analysis.

$$A(H \rightarrow ZZ) = v^{-1} \left(\underbrace{a_1 m_Z^2 \epsilon_1^* \epsilon_2^*}_{\text{SM}} + \underbrace{a_2 f_{\mu\nu}^{*(1)} f^{*(2),\mu\nu}}_{\text{A-CP even}} + \underbrace{a_3 f_{\mu\nu}^{*(1)} \tilde{f}^{*(2),\mu\nu}}_{\text{A-CP odd}} \right)$$

- Effective CP odd fraction extracted from data, re-parameterizing the likelihood fit used in the 0^+ vs 0^- test as a function of:

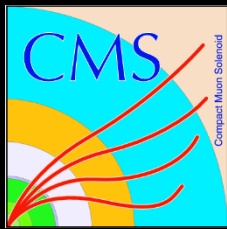
$$f_{a3} = \frac{|a_3|^2 \sigma_3}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3}$$

- Observed (expected) 95% C.L. Bound: $f_{a3} < 0.5$ (0.7).



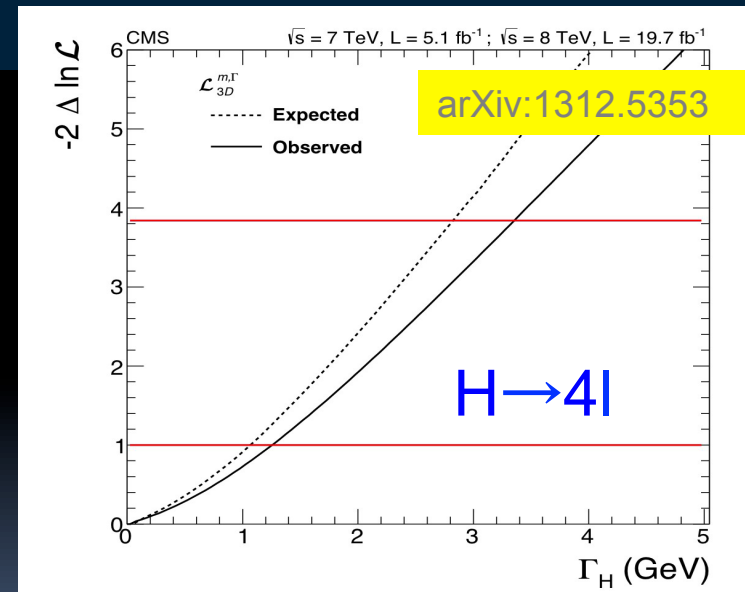
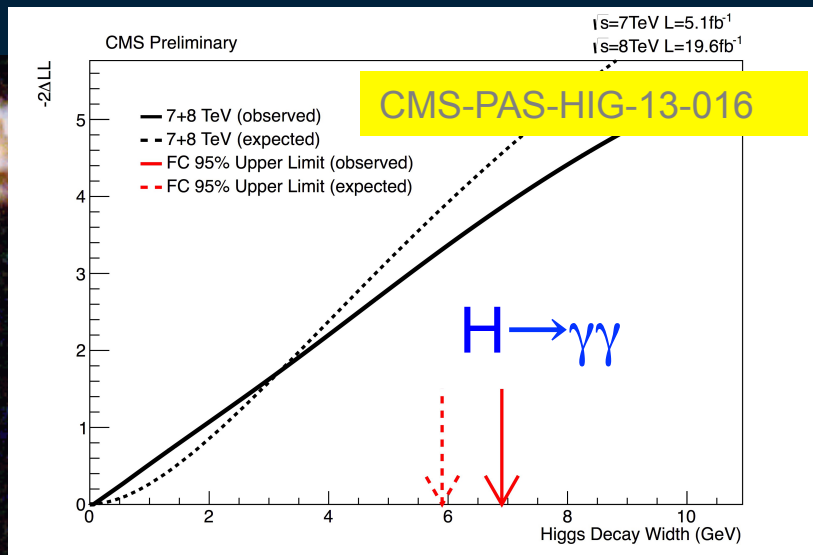


Higgs boson width



- Lack of knowledge of Γ_H introduces degeneracy in the knowledge of the Higgs couplings
- In the SM $\Gamma_H(125 \text{ GeV}) \sim 4 \text{ MeV}$
- Direct measurements limited by experimental resolution ($\approx \text{GeV}$).
 - Current upper limits 3.4(7) GeV from $4l(\gamma\gamma)$ decay modes.

$$\sigma_{i \rightarrow H \rightarrow f} \sim \frac{g_i^2 g_f^2}{\Gamma_H}$$



- New: Use interference between Higgs resonance in gluon fusion and the continuum back-ground to measure the Higgs width
 - Complements more direct measurements possible at lepton colliders



Total width: interferometry



- Off-shell $H^* \rightarrow VV$ ($V = W, Z$) enhances the $H(126)$ cross-section at high mass [$\sim 8\%$ of $\sigma(H \rightarrow ZZ)$ found in $m_{ZZ} > 2m_Z$]

$$\frac{d\sigma_{pp \rightarrow H \rightarrow ZZ}}{dM_{4l}^2} \sim \frac{g_{Hgg}^2 g_{HZZ}^2}{(M_{4l}^2 - m_H^2)^2 + m_H^2 \Gamma_H^2}$$

$$\sigma_{gg \rightarrow ZZ}^{\text{on-peak}} \propto \frac{g_{ggH}^2 g_{HZZ}^2}{\Gamma}$$

$$\sigma_{gg \rightarrow ZZ}^{\text{off-peak}} \propto g_{ggH}^2 g_{HZZ}^2$$

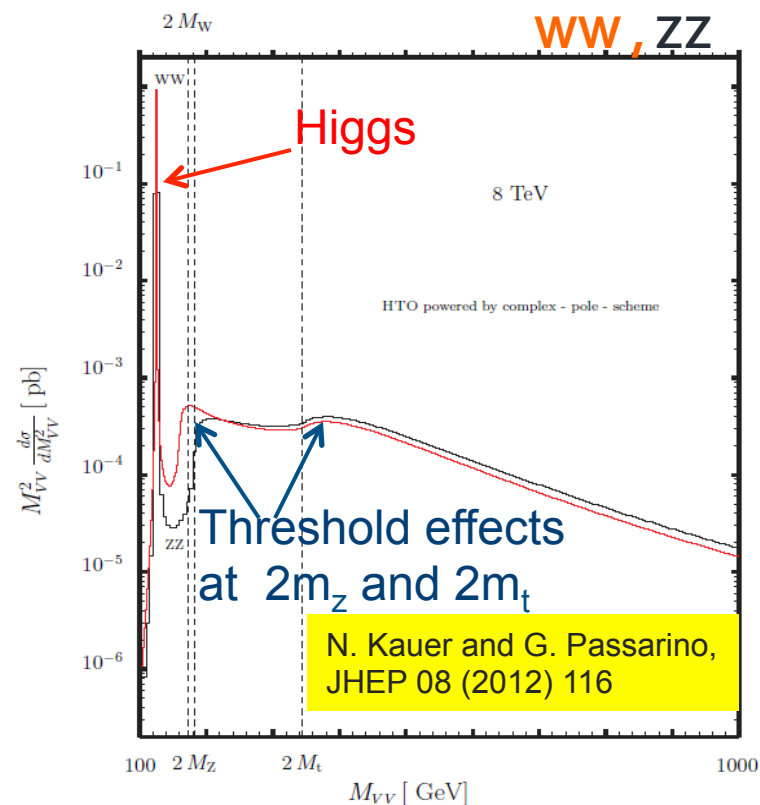
Measuring

$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{on-peak}}$$

$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{\text{off-peak}}$$

constrains Γ_H

gluon-gluon fusion production



F. Caola, K. Melnikov, Phys. Rev. D88 (2013) 054024
J. Campbell et al. arXiv:1311.3589



Total width: interferometry

Define

$$\kappa_g = g_{ggH} / g_{ggH}^{SM}$$

$$\kappa_Z = g_{HZZ} / g_{HZZ}^{SM}$$

$$r = \Gamma_H / \Gamma_H^{SM}$$

$$\sigma_{gg \rightarrow H \rightarrow ZZ}^{on-peak} = \frac{\kappa_g^2 \kappa_Z^2}{r} (\sigma \cdot BR)_{SM} \equiv \mu (\sigma \cdot BR)_{SM}$$

$$\frac{d\sigma_{gg \rightarrow H \rightarrow ZZ}^{off-peak}}{dm_{ZZ}} = \kappa_g^2 \kappa_Z^2 \cdot \frac{d\sigma_{gg \rightarrow H \rightarrow ZZ}^{off-peak, SM}}{dm_{ZZ}} = \mu r \frac{d\sigma_{gg \rightarrow H \rightarrow ZZ}^{off-peak, SM}}{dm_{ZZ}}$$

- If μ is known r can be determined. Use μ from $H \rightarrow ZZ \rightarrow 4l$

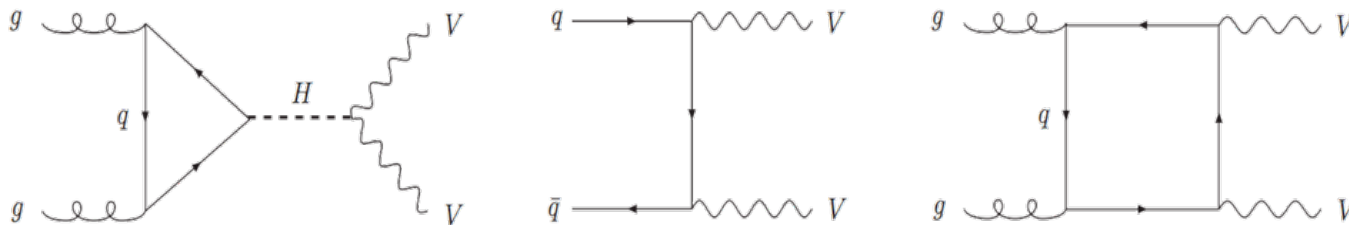
- Expected

$$\mu = 1.00^{+0.27}_{-0.24}$$

- Observed

$$\mu = 0.93^{+0.26}_{-0.24}$$

- Challenges: Interference $gg \rightarrow ZZ$ (gg2VV/MCFM)



- Effect in VBF $\approx 10\%$ at high mass (assume $\mu_F = \mu_V$, PHANTOM)

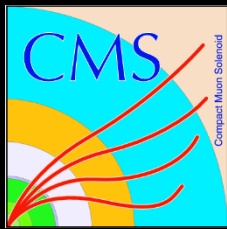
- Fit

$$P_{tot} = \mu r P_{sig} + \sqrt{\mu r} P_{int} + P_{bkg}$$

where P are MC- or data-derived templates



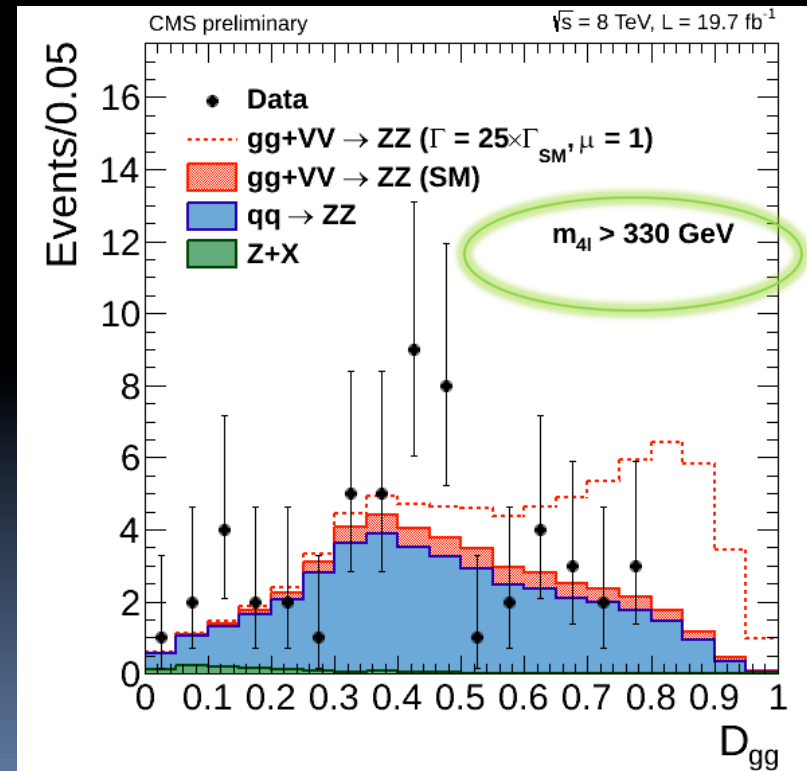
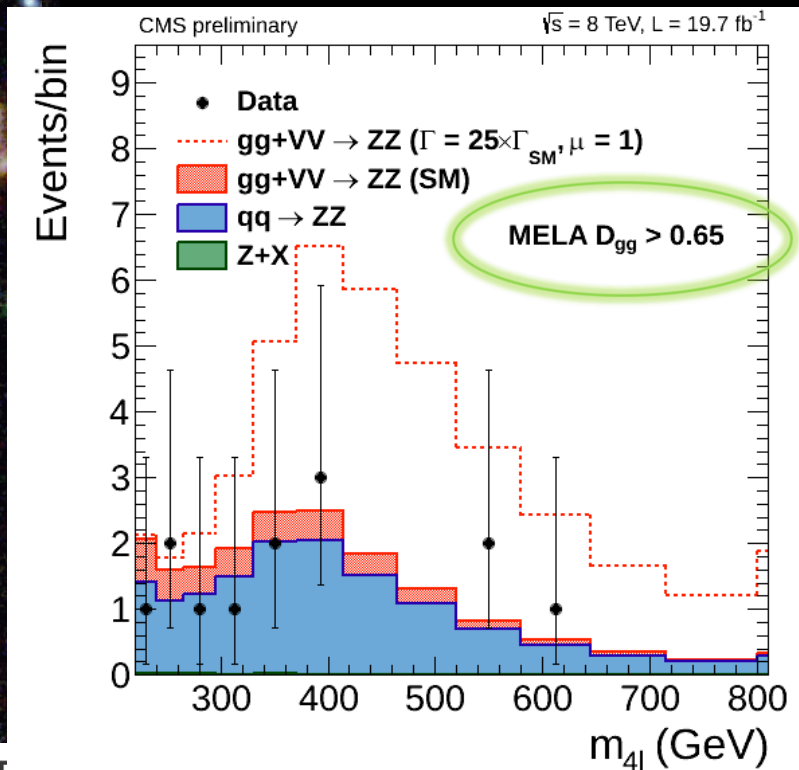
Width in $H \rightarrow ZZ \rightarrow 4l$



- $H \rightarrow ZZ \rightarrow 4l$ analysis (arXiv:1312.5333)
 - Require $m(4l) > 220 \text{ GeV}$
 - Develop NEW MELA discriminant using with 7 kinematics variables (m_{Z1}, m_{Z2} , five angles) to distinguish between $gg \rightarrow ZZ$ production (signal background and interference) vs $qq \rightarrow ZZ$

$$\mathcal{D}_{gg,a} = \frac{\mathcal{P}_{gg,a}}{\mathcal{P}_{gg,a} + \mathcal{P}_{q\bar{q}}}$$

$a=10$ [$a=1$ for SM]



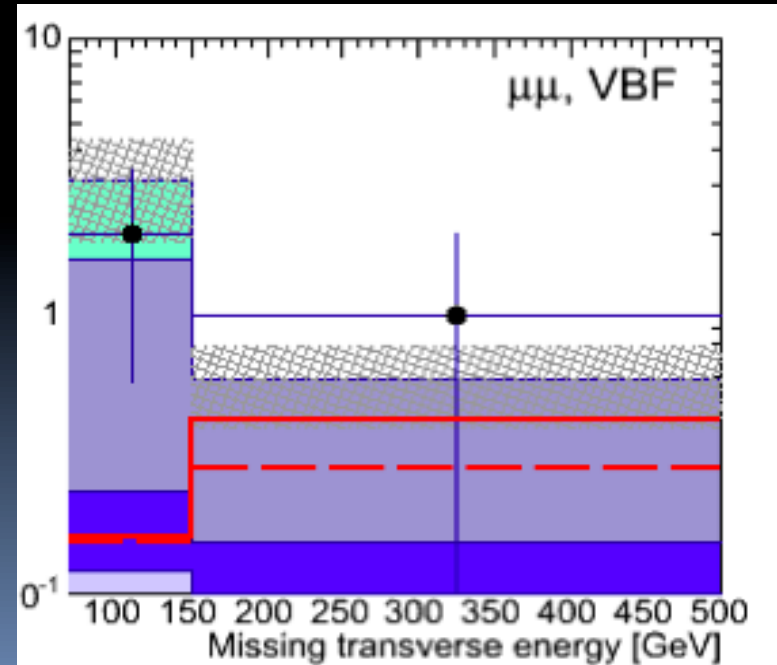
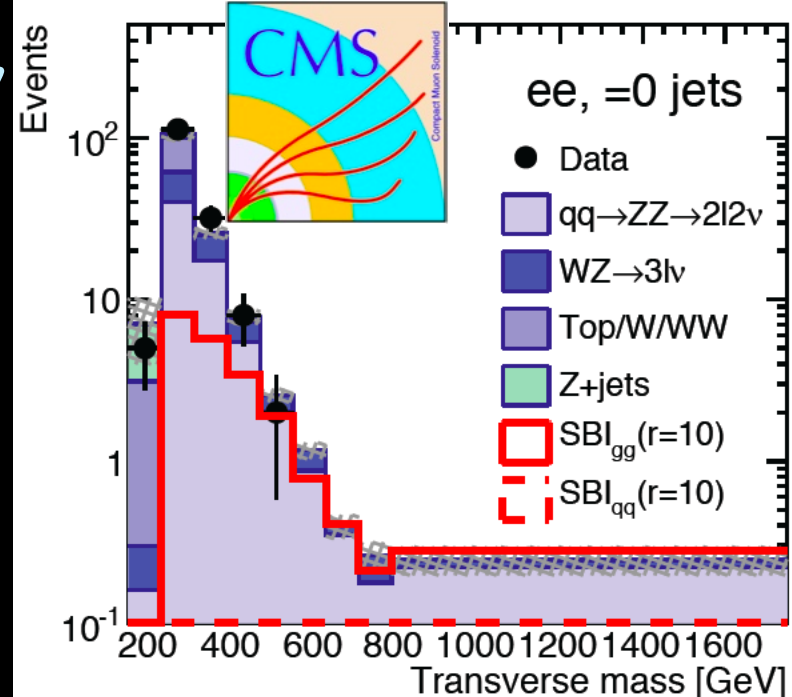


Width $H \rightarrow ZZ \rightarrow 2l2\nu$

- $H \rightarrow ZZ \rightarrow 2l2\nu$ analysis as described in Eur. Phys. J. C (2013) and CMS PAS-HIG-13-014
 - Require large $p_T(Z)$ and $E_{T,miss}$ (boosted)
 - Veto 3rd lepton and b-tagged jets
 - Split events in 3 categories according to number no. jets ($p_T > 30 \text{ GeV}$ and $|\eta| < 4.7$)
 - VBF-like: two jet with $m_{JJ} > 500 \text{ GeV}$ and $|\Delta\eta_{JJ}| > 4$
 - ≥ 1 jets: excluding VBF-like category
 - 0 jets
 - discriminating variables:
 - m_T for 0 and 1-jet category:

$$m_T^2 = \left[\sqrt{p_{T,\ell\ell}^2 + m_{\ell\ell}^2} + \sqrt{E_T^{miss^2} + m_{\ell\ell}^2} \right]^2 - \left[\vec{p}_{T,\ell\ell} + \vec{E}_T^{miss} \right]^2$$

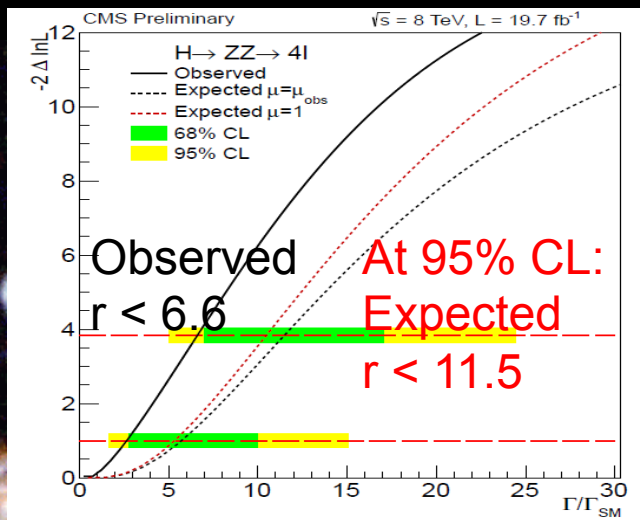
- $E_{T,miss}$ for VBF-like category



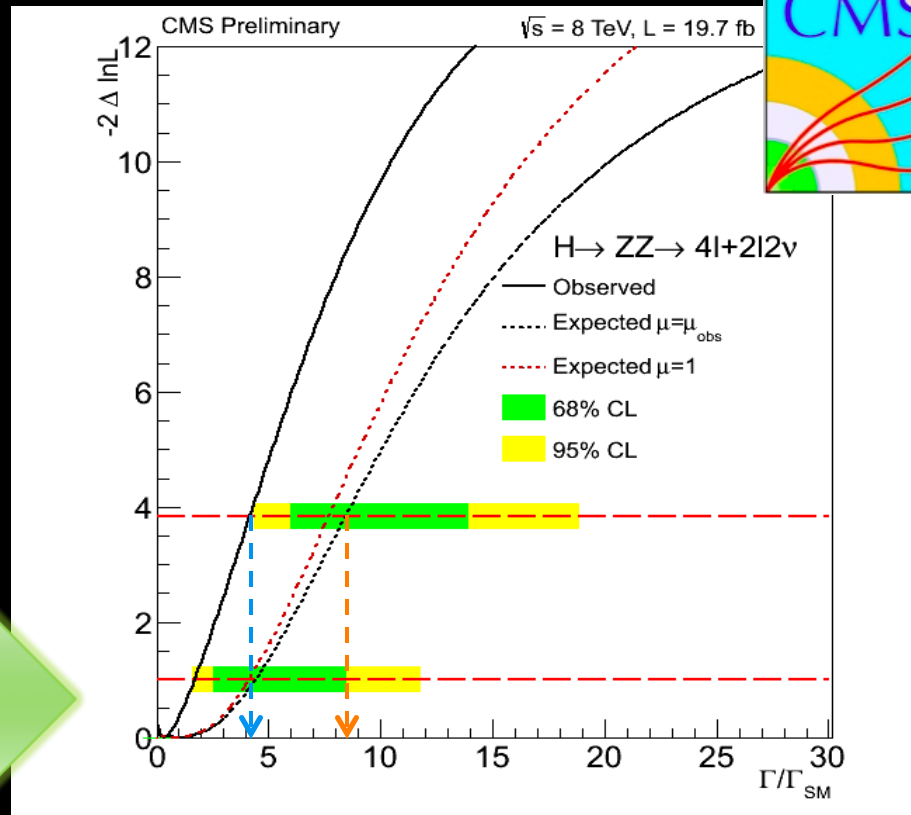
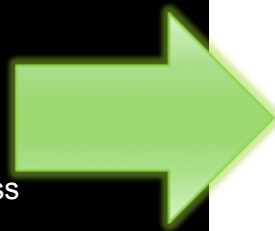
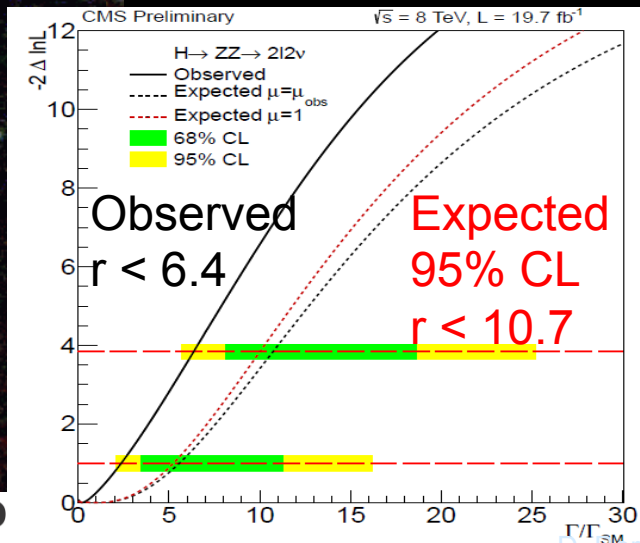


Width $H \rightarrow ZZ$

4l : 2D likelihood in m_{4l} and D_{gg}



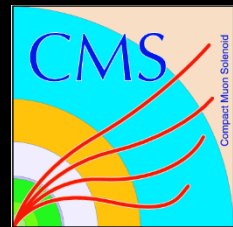
2l2v : 1D likelihood with m_T or E_{tmiss}



- Combined observed (expected) values
 - $r = \Gamma/\Gamma_{\text{SM}} < 4.2$ (8.5) @ 95% CL (p-value = 0.02)
 - $r = \Gamma/\Gamma_{\text{SM}} = 0.3+1.5-0.3$
- equivalent to:
 - < 17.4 (35.3) MeV @ 95% CL
 - $= (1.4+6.1-1.4) \text{ MeV}$



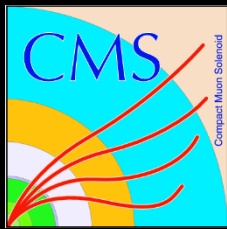
Conclusions



- Run1 of the LHC has brought the discovery of the new boson and the first measurements of its properties
 - Expect precision measurements of Higgs boson mass to $\approx 0.2\%$ from LHC Run 1 combination.
 - Spin/parity of new boson consistent with 0^+ .
 - A few alternative models have been tested and experiments are moving towards setting limits on anomalous couplings.
 - A new technique to constrain the Higgs boson width from non-resonant ZZ production has emerged
- The large luminosity of the HL-LHC upgrades will eventually
 - Allow the transition into precision Higgs physics More fully test the consistency of the new particle with the SM Higgs boson.
 - Measurements of rare Higgs boson decays
 - Study of the Higgs self coupling
- Work will be required to bring down the theory uncertainties as well



Conclusions



- Run the new boson
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BACKUP

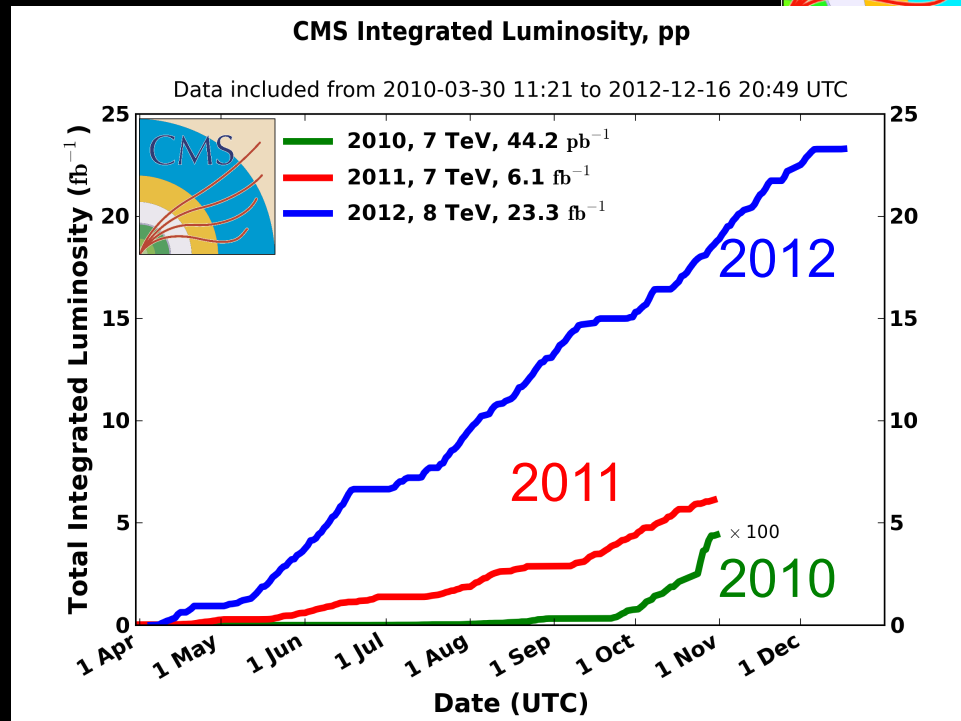
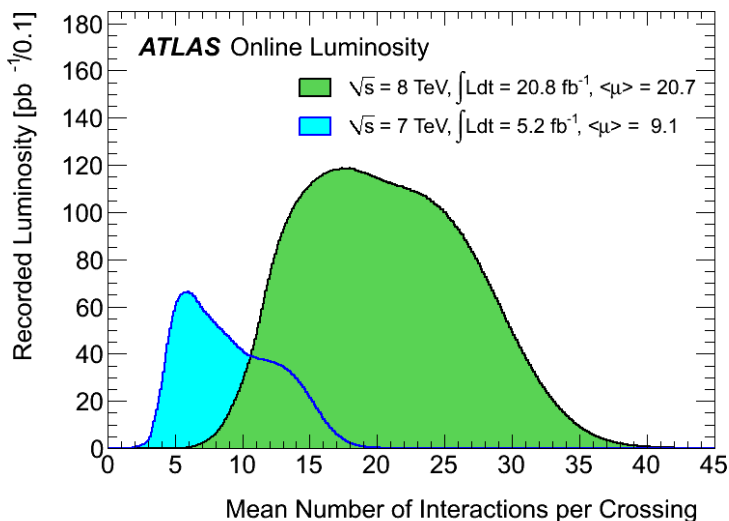




LHC Run 1 data set



- Spectacular performance of the machine and the detectors
- ~90% of the delivered data available for offline analysis.
- Available dataset:
~5fb⁻¹ \sqrt{s} = 7TeV +
~20fb⁻¹ \sqrt{s} = 8TeV



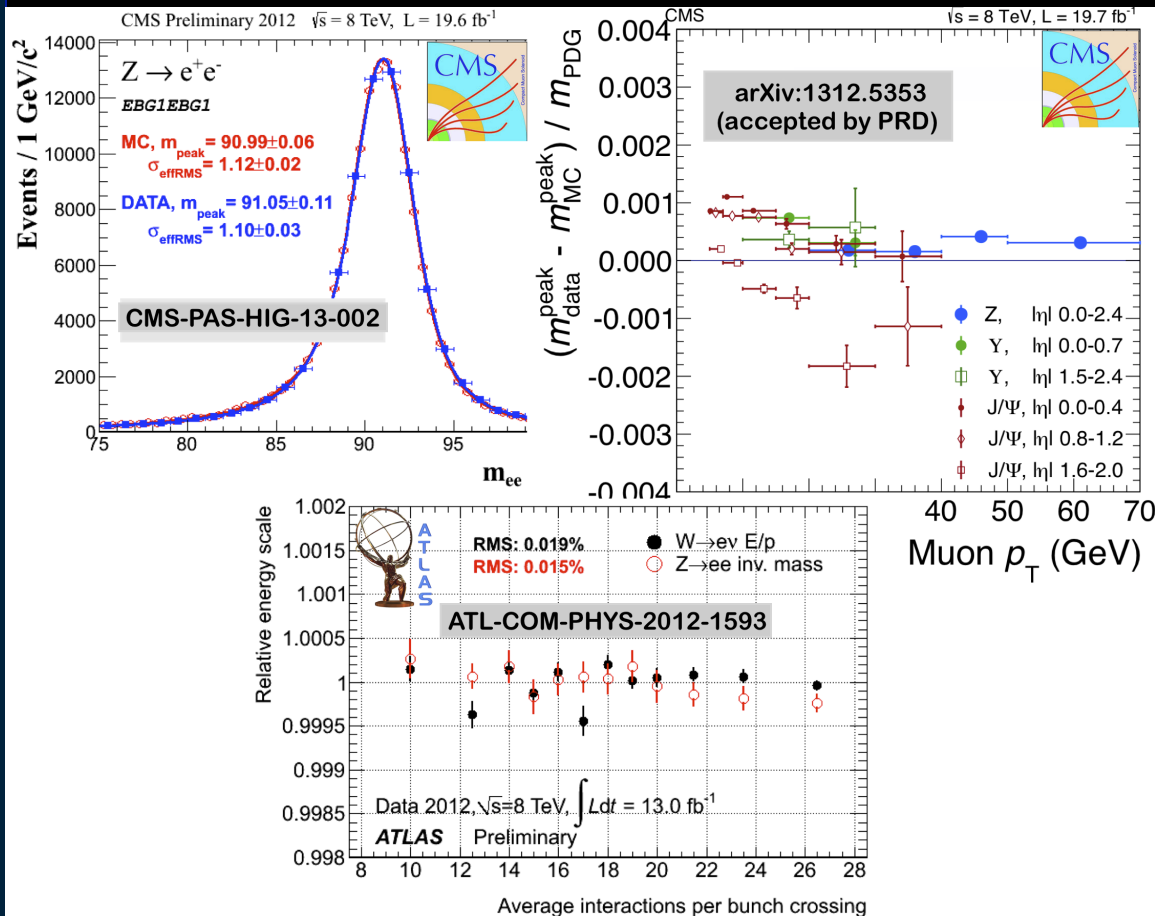
- Challenging pile-up conditions.
 - Up to 30 average interactions per bunch-crossing.
 - Development of pile-up safe analysis techniques



Mass resolution and scale uncertainties

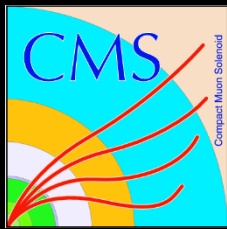


- CMS: e/γ energy estimated using multivariate regression
- ATLAS: weighted sum of energy deposits in the different calorimeter layers
- Scale and resolution is obtained from W ; Z ; J/ψ and Upsilon resonances
- Additional smearing is applied to MC to match the resolution in data
- Resulting systematic uncertainty on mass measurements is 0.5% per channel ($H \rightarrow \gamma\gamma$ and $H \rightarrow 4l$)





H \rightarrow 4l Mass Measurement



- Systematics
 - Electron energy corrections.
 - MC correction using multivariate regression (CMS) or weighted sum of energy deposits.
 - Data/MC corrections derived on $Z \rightarrow ee$ and checked with $Z \rightarrow ee$ and low-mass resonances.
 - Muon momentum scale corrections.
 - Data/MC corrections from $Z \rightarrow \mu\mu$ + check w/ low-mass resonances.

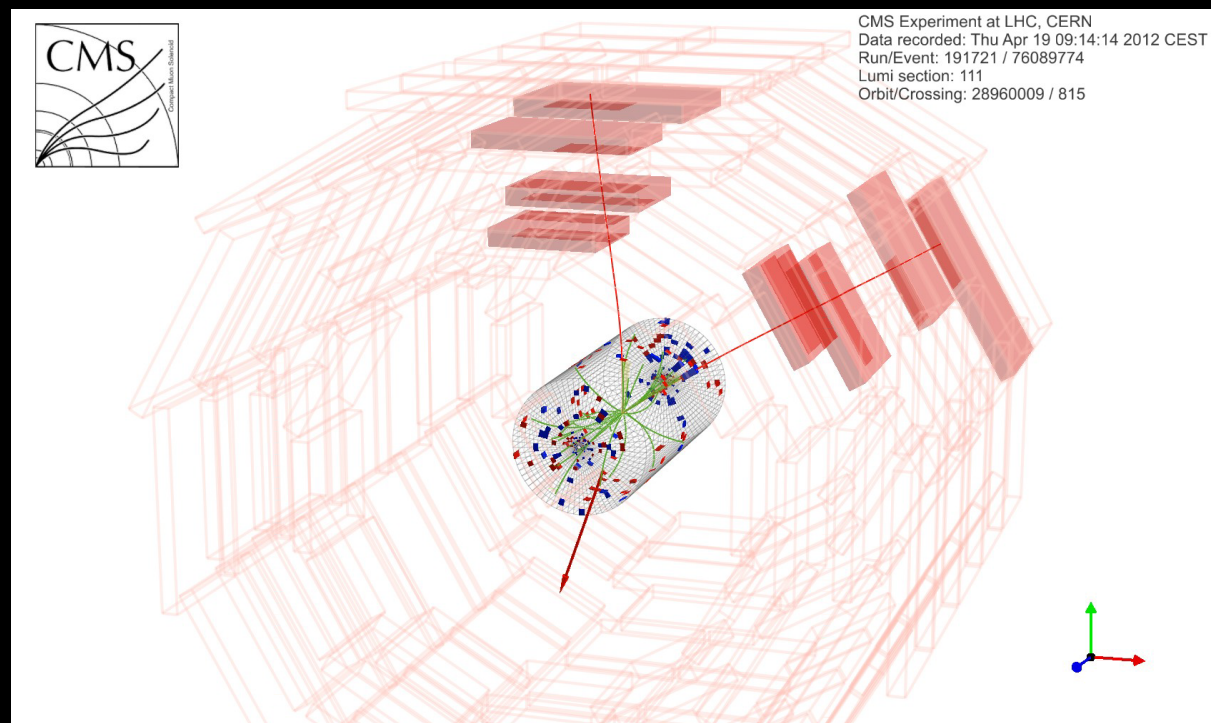
	ATLAS PLB 726 (2013), 88-119		CMS arXiv:1312.5353	
Mass [GeV]	$124.3^{+0.6}_{-0.5}(stat) \pm 0.7^{+0.5}_{-0.3}(sys)$		$125.6 \pm 0.4(stat) \pm 0.2(sys)$	
Sys. Uncertainty	Electron e/p scale	0.2-0.4%	Electron e/p scale	0.1-0.3%
	Muon p-scale	0.1-0.2%	Muon p-scale	0.1%



J^P in $H \rightarrow WW \rightarrow \ell\nu\ell\nu$



- Distinct signature:
 - Two high p_T leptons
 - Missing transverse energy.
- Large branching fraction.
- Poor mass resolution.
- Large backgrounds.
- Angular correlation between final state leptons provides information on the polarization of the resonance.

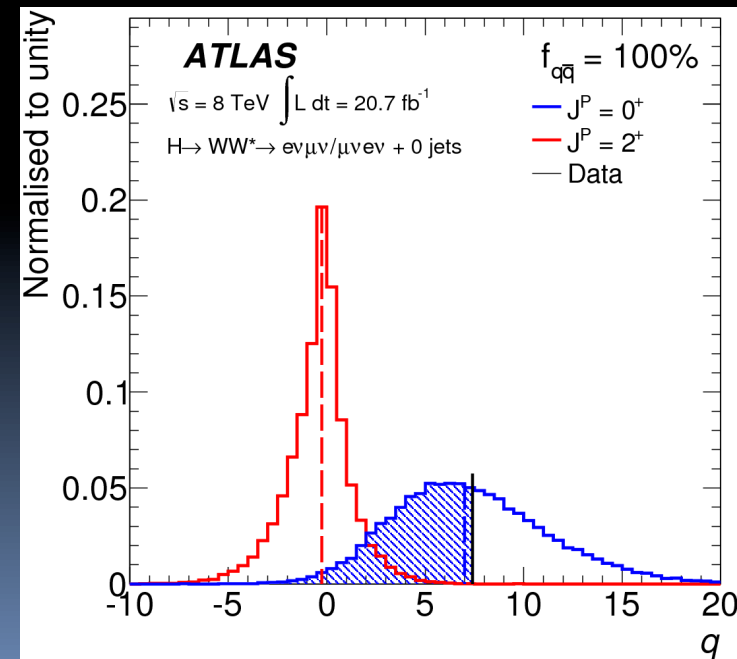
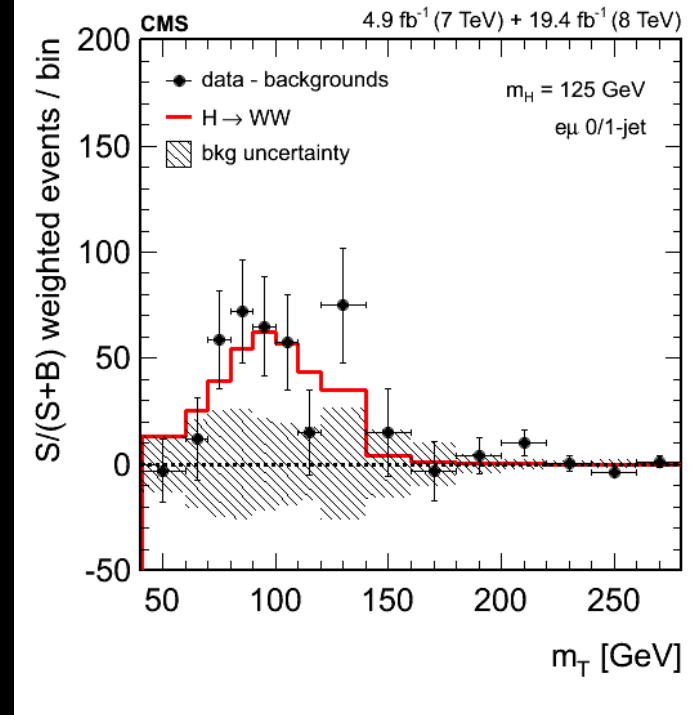


- Challenges:
 - Missing energy resolution.
 - Background modeling



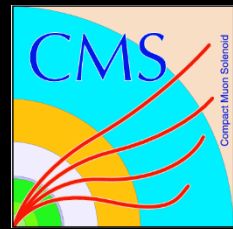
J^P in $H \rightarrow WW \rightarrow l\nu l\nu$

- Analysis strategy:
- Select two high p_T , different flavor leptons plus missing E_T .
 - CMS: categorize events in 0 and 1 jet bins
 - ATLAS: no categorization in number of jets.
- Hypothesis test from 2D template fit to data:
- CMS: m_{ll} vs m_T
- ATLAS: use two BDT discriminants ($\Delta\phi_{ll}$, m_{ll} , m_T)
 - DT0 (discriminate SM from background)
 - BDTalt (discriminate alternative hyp from background).
- Tested alternative models:
 - CMS: 2+m “graviton-like” and 0^- .
 - ATLAS: 2+ m “graviton-like”.
 - For 2+ m model both qq, gg production modes (and mixtures) are considered.

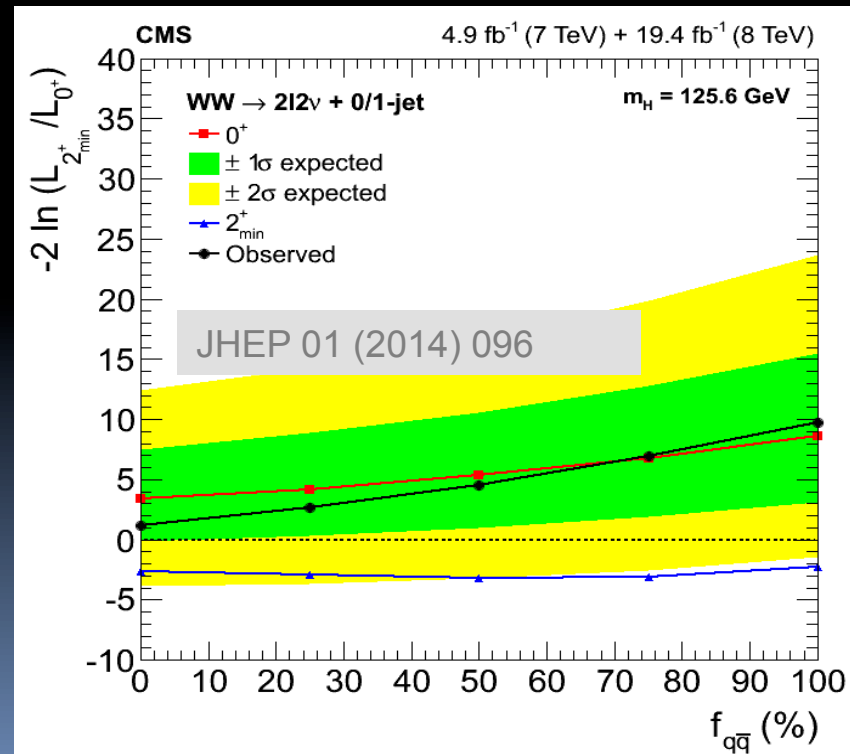
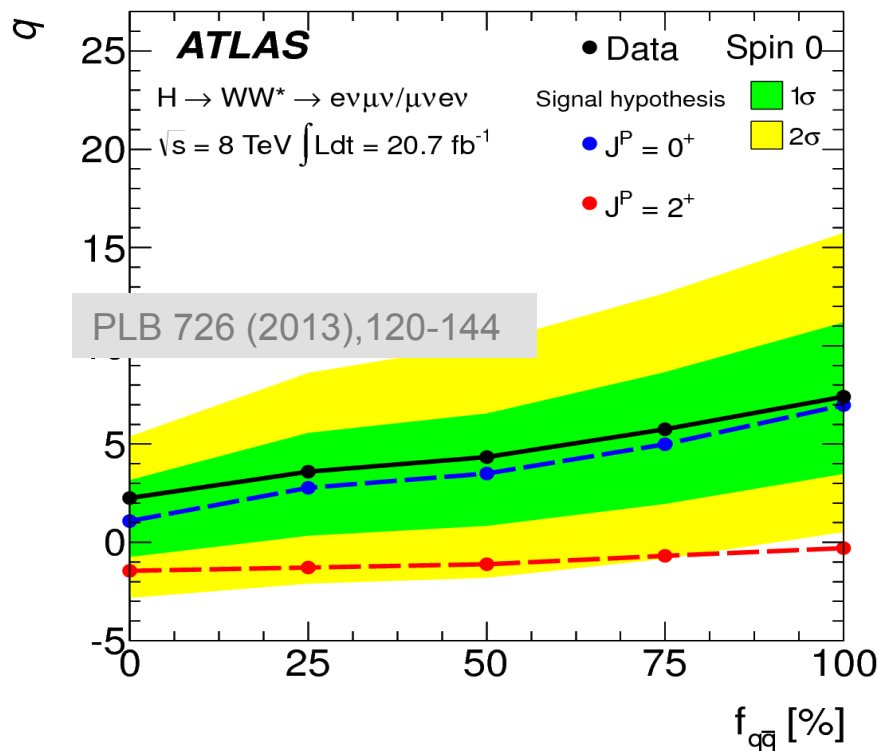




J^P in $H \rightarrow WW \rightarrow l\nu l\nu$



- Expected (post-fit) exclusion for 2+m model 1-CLs > 0.94 .
- In CMS, 0- expected (post-fit) exclusion 1-CLs = 0.72.
- Observed results favor SM hypothesis.





J^P in $H \rightarrow \gamma\gamma$



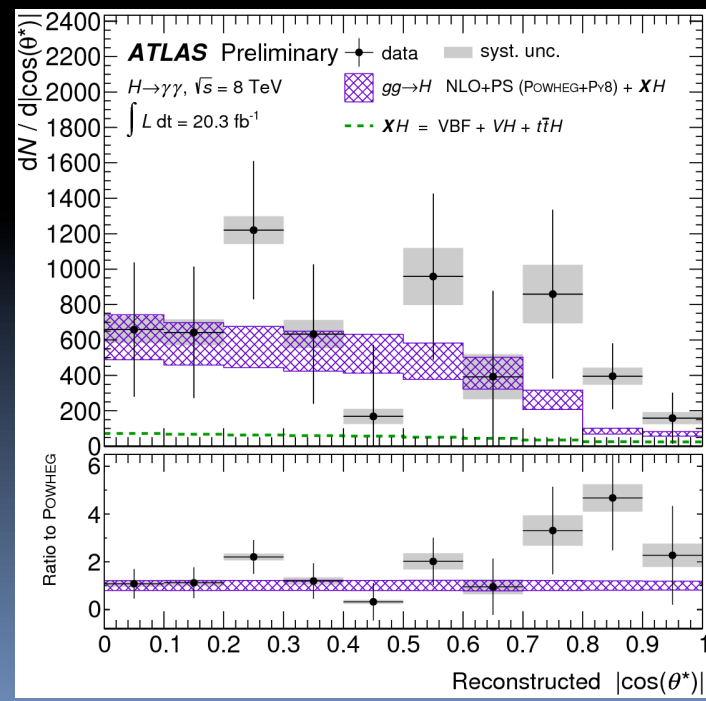
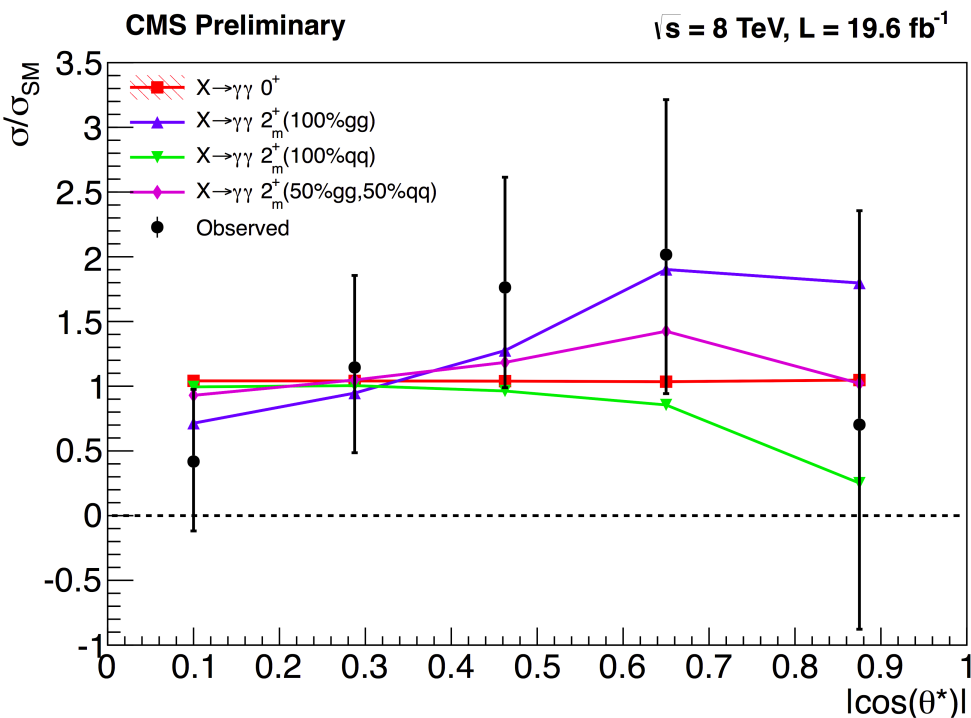
Distribution of production angle sensitive to spin/parity.

- Event selection similar to mass analysis.
- ATLAS: no categorization in photon kin. or resolution
- CMS: simple 4 categories cut-based categorization
- Hypothesis test:

$$\cos\theta^* = 2 \frac{E_2 p_{z1} - E_1 p_{z2}}{m_{\gamma\gamma} \sqrt{m_{\gamma\gamma}^2 + p_{T\gamma\gamma}^2}}$$

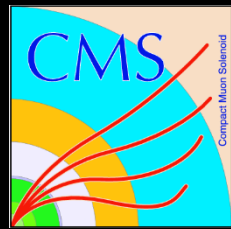
- CMS: simultaneous fit to $m_{\gamma\gamma}$ in 5 $\cos(\theta^*)$ bins

- ATLAS: 2D fit of $\cos(\theta^*)$ vs $m_{\gamma\gamma}$

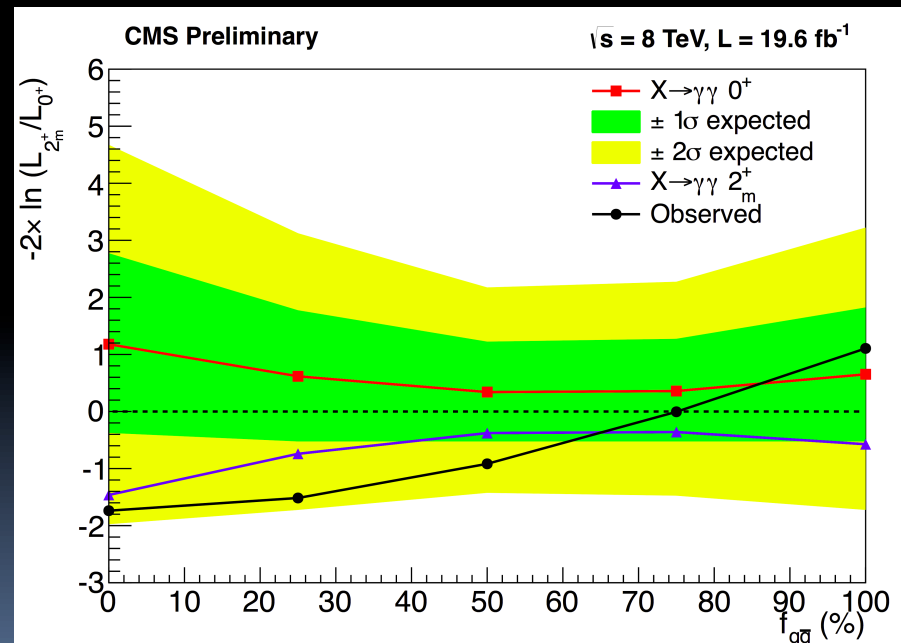
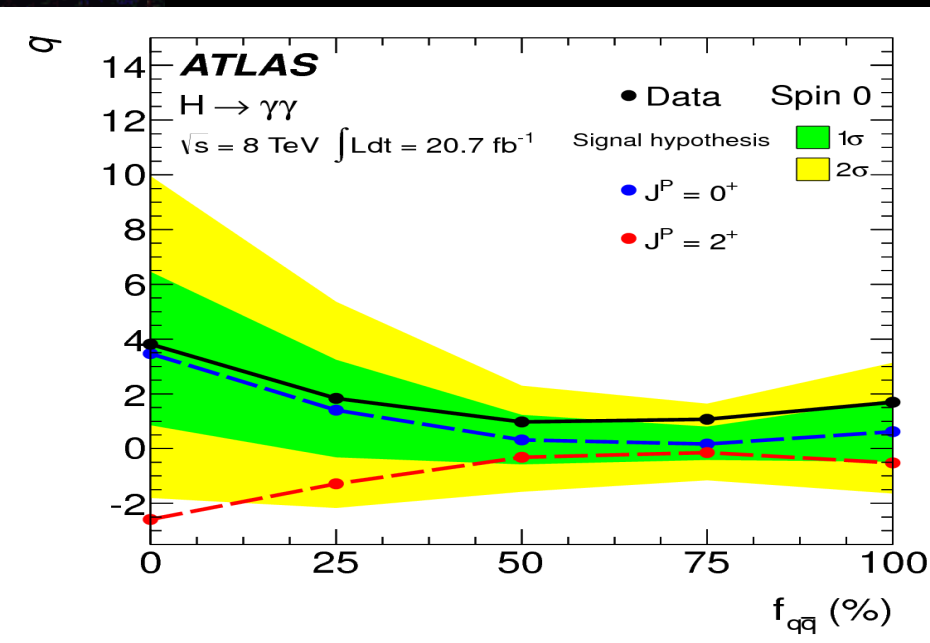




J^P in $H \rightarrow \gamma\gamma$



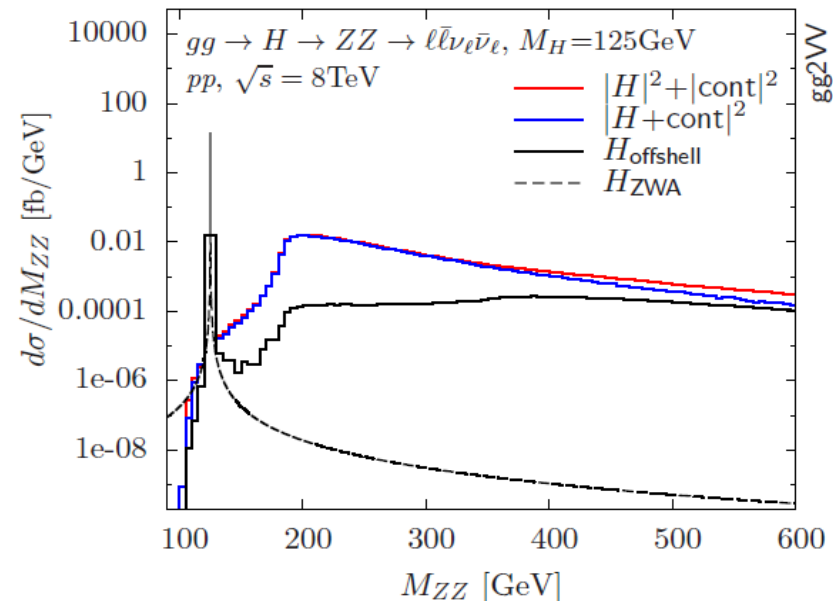
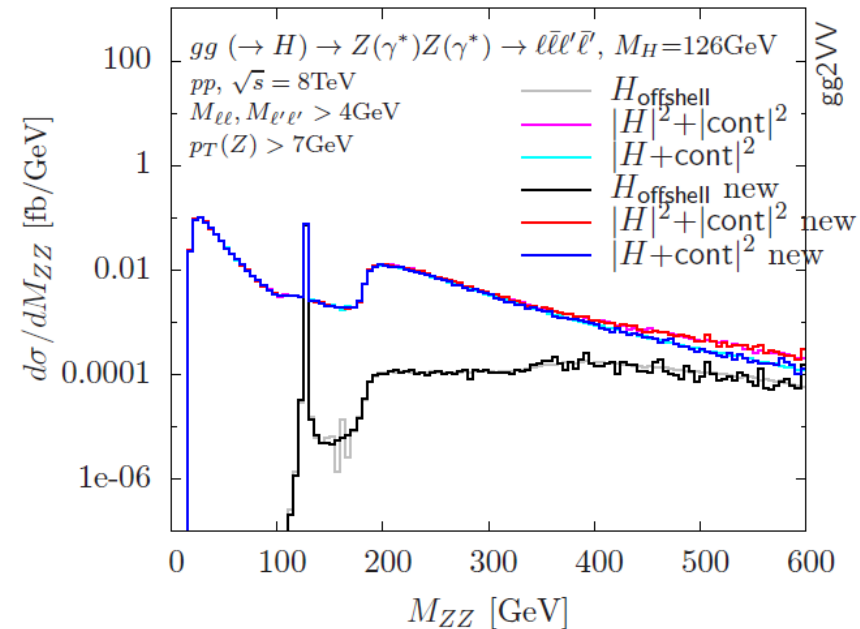
- Post-fit) Expected separation: 1-CLs > 17(55)-60(99)% for CMS
- (ATLAS).
- Better sensitivity for ATLAS analysis partially driven by higher observed excess.
- SM hypothesis generally favored in data.





The 4l and 2l2 ν final states

- 4l final state ($l = e, \mu$)
 - At high mass, basically only background is $qq \rightarrow ZZ$ (known at NLO, QCD uncertainties at the level of %)
 - Fully reconstructed state \rightarrow can use matrix element probabilities of lepton 4-vectors to distinguish between gg and qq production
- 2l2 ν final state ($l = e, \mu$)
 - Much larger BR (x6) but smaller acceptance (tight p_T selection)
 - Rely on transverse mass distributions



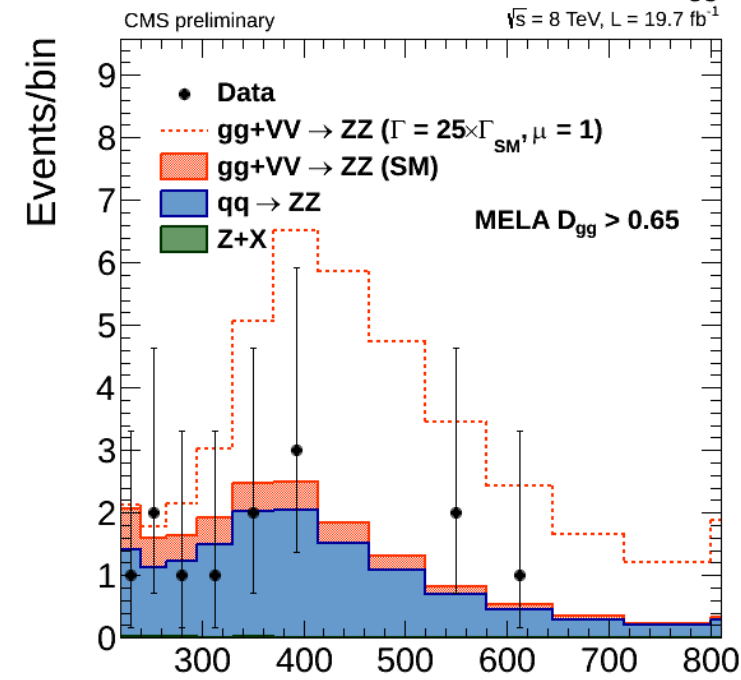
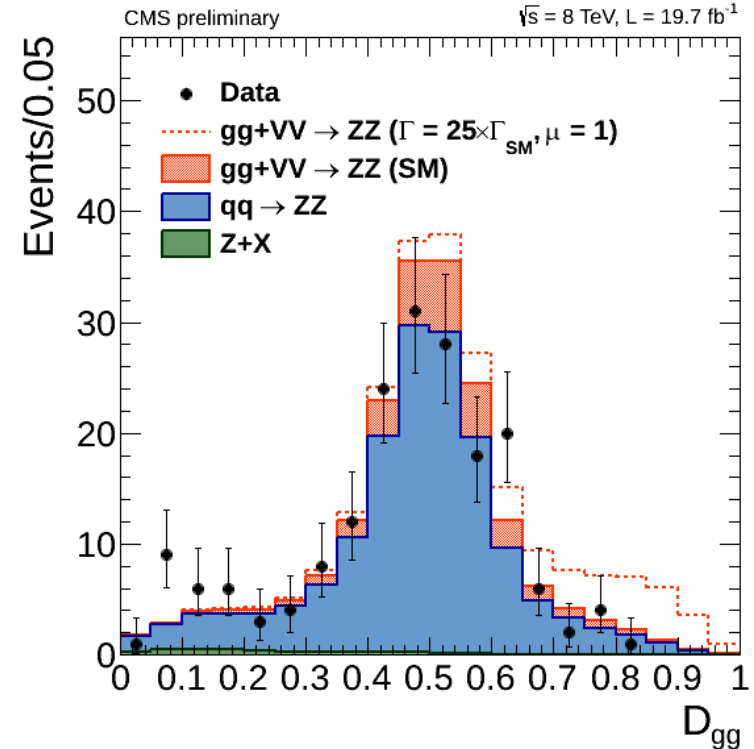


Width 4l analysis

- No changes in selection w.r.t. CMS collab. , arXiv:1312.5353
- Lepton p_T cuts, Z invariant masses, impact parameter significance, loose isolation
- In the matrix element likelihood approach (MELA), design a specific discriminant for $gg \rightarrow ZZ$ production:

$$D_{gg,a} = \frac{P_{gg,a}}{P_{gg,a} + P_{q\bar{q}}}$$

- Built with 7 variables completely describing kinematics (m_{Z1} , m_{Z2} , five angles)
- $P_{gg,(q\bar{q})}$ are joint probabilities for $gg \rightarrow ZZ$, signal + background + interference ($q\bar{q} \rightarrow ZZ$) from MCFM matrix elements

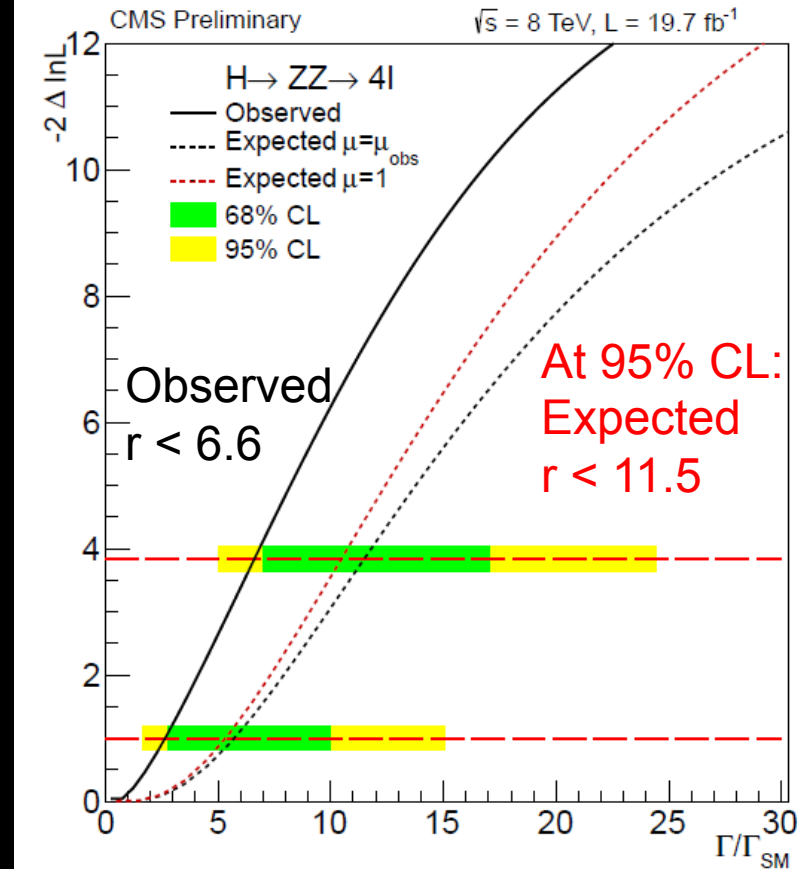




Width $H \rightarrow ZZ \rightarrow 4\ell$

Main systematic uncertainties:

- QCD scale and PDFs for $qq \rightarrow ZZ$ and $gg \rightarrow ZZ$
- μ uncertainties 4ℓ low-mass paper
- Uncertainty on k-factor approximation for $gg \rightarrow ZZ$ continuum
- Experimental uncertainties (lepton trigger/reconstruction efficiencies etc.)



	Full region	Signal-enriched region
$gg + \text{VBF} \rightarrow 4\ell$ (signal, $\Gamma_H/\Gamma_H^{\text{SM}} = 1$)	$2.22^{+0.15}_{-0.17}$	$1.20^{+0.08}_{-0.09}$
$gg + \text{VBF} \rightarrow 4\ell$ (background)	$31.1^{+3.0}_{-3.1}$	2.12 ± 0.21
(a) $gg + \text{VBF} \rightarrow 4\ell$ (total, $\Gamma_H/\Gamma_H^{\text{SM}} = 1$)	$29.6^{+2.8}_{-2.9}$	$1.73^{+0.16}_{-0.17}$
$gg + \text{VBF} \rightarrow 4\ell$ (total, $\Gamma_H/\Gamma_H^{\text{SM}} = 15$)	$51.8^{+4.9}_{-5.0}$	13.1 ± 1.1
(b) $q\bar{q}$	154.7 ± 7.4	8.6 ± 0.4
(c) Reducible background	3.7 ± 0.6	0.44 ± 0.08
(a+b+c) Total expected ($\Gamma_H/\Gamma_H^{\text{SM}} = 1$)	188.0 ± 7.9	10.8 ± 0.4
Observed	183	8



Width $H \rightarrow ZZ \rightarrow 2l2\nu$



- No changes in selection to CMS PAS-HIG-13-014
 - Large $p_T(Z)$ and $E_{T,miss}$
 - Vetoing 3rd lepton and b-tagged jets (removing Z+heavy-flavor jets)
 - Events split in three purity categories according to number of selected jets ($p_T > 30$ GeV and $|\eta| < 4.7$)
 - VBF-like: two jets with $m_{JJ} > 500$ GeV and $|\Delta\eta_{JJ}| > 4$
 - ≥ 1 jets: excluding events in VBF-like category
 - 0 jets
 - Data-derived estimation of reducible backgrounds (double and single top, WW, W+jets, Z+jets), $qq \rightarrow ZZ$ and WZ from MC
 - Fit the distribution of the transverse mass for 0 and 1-jet category

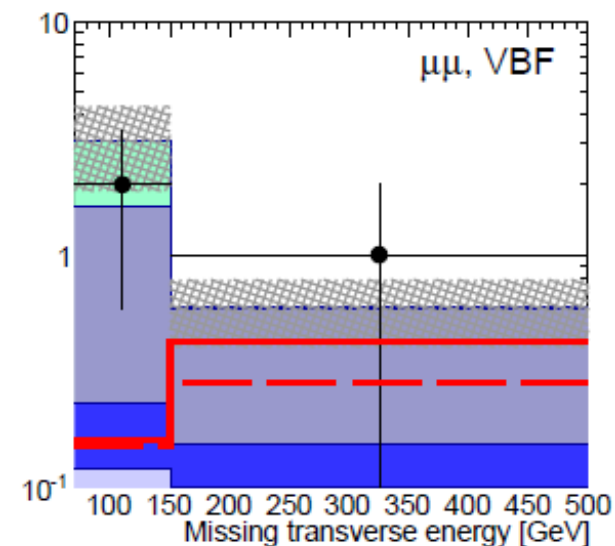
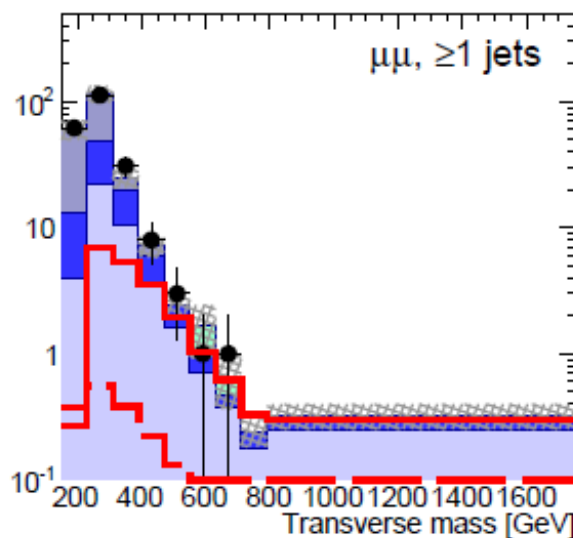
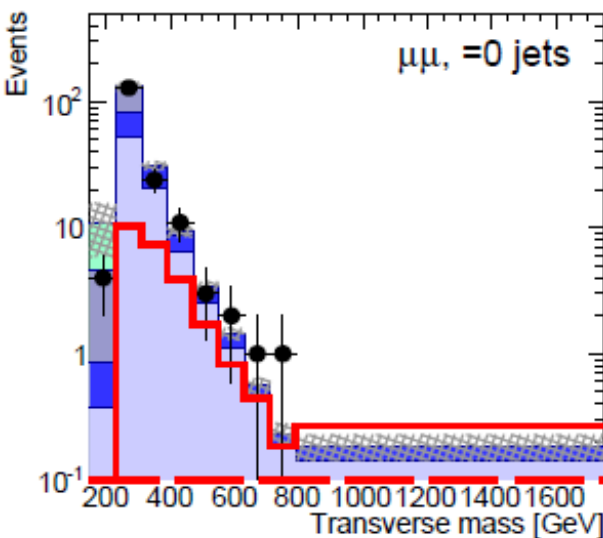
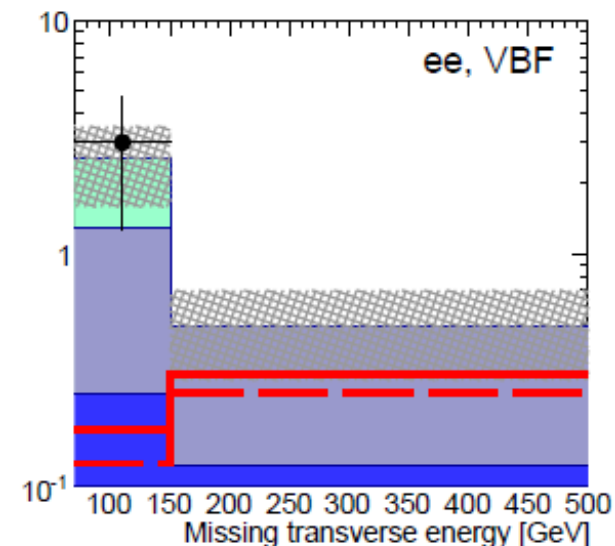
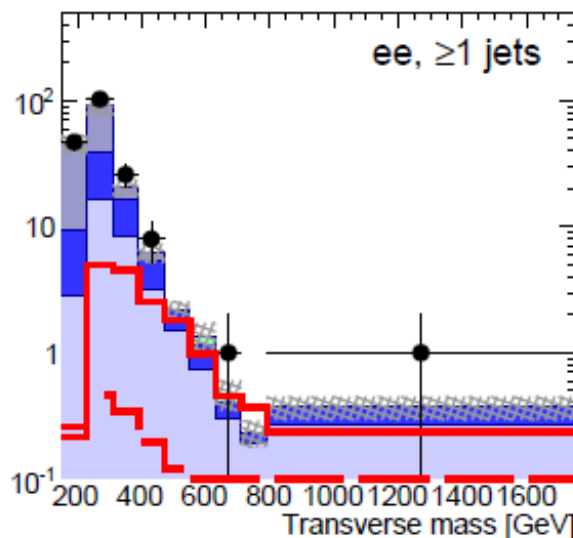
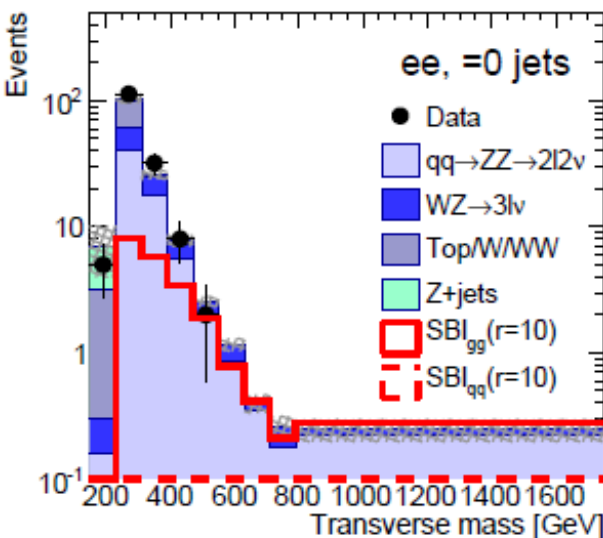
$$m_T^2 = \left[\sqrt{p_{T,ee}^2 + m_{ee}^2} + \sqrt{E_T^{miss^2} + m_{ee}^2} \right]^2 - \left[\vec{p}_{T,ee} + \vec{E}_T^{miss} \right]^2$$



Width $H \rightarrow ZZ \rightarrow 2l2\nu$

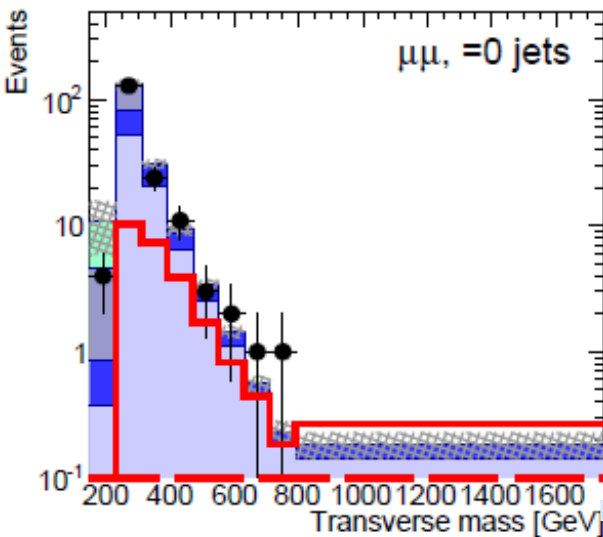
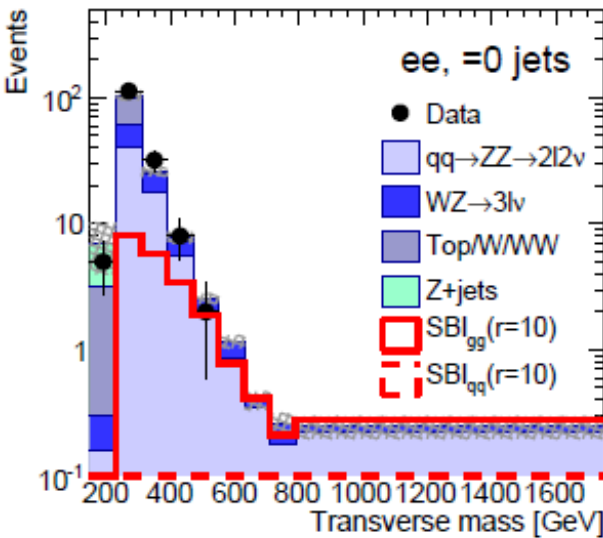
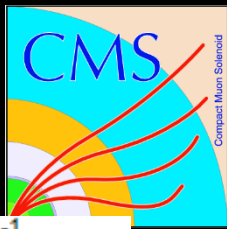


CMS preliminary, $\sqrt{s}=8.0$ TeV, $|\mathcal{L}|=19.7$ fb $^{-1}$





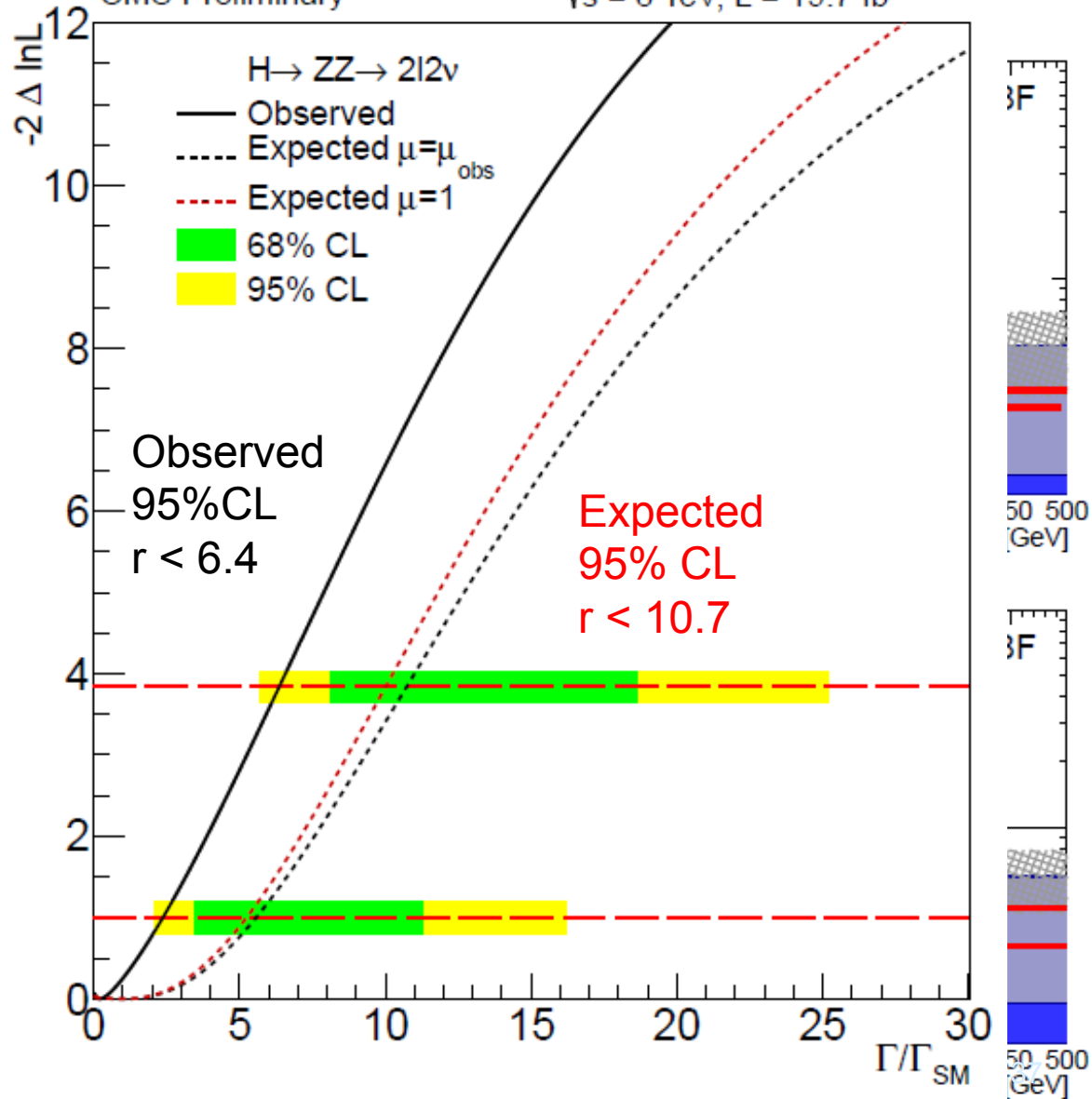
Width $H \rightarrow ZZ \rightarrow 2l2\nu$



CMS

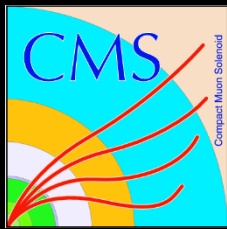
CMS Preliminary

$\sqrt{s} = 8 \text{ TeV}, L = 19.7 \text{ fb}^{-1}$

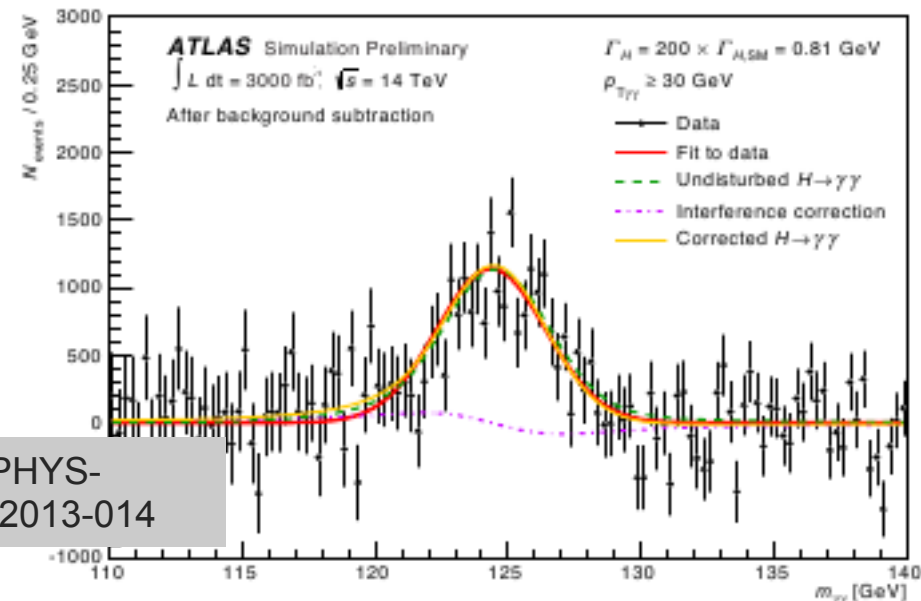
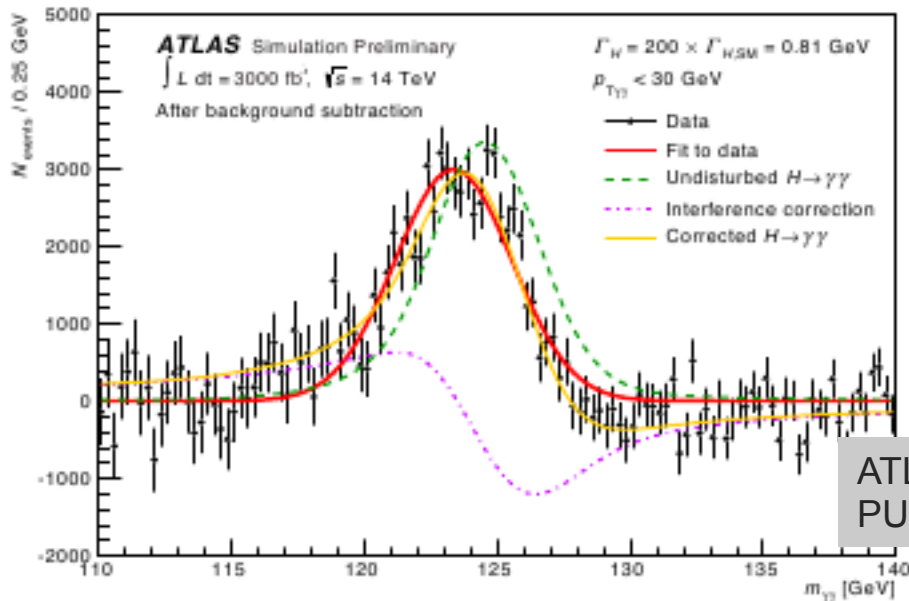




Interferometry - di-photon



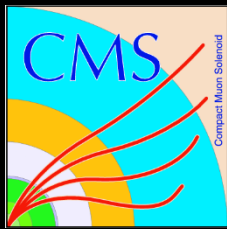
- Can also exploit destructive interference between $gg \rightarrow \gamma\gamma$ and $gg \rightarrow H \rightarrow \gamma\gamma$.
- Generate effective mass shift, which magnitude varies as a function of the boson p_T .
- Constraint of the width from measurement of m_H vs p_{TH} .
- Projected sensitivity for $3ab^{-1}$ $\Gamma < 30 \Gamma_{SM}$ (95% CL).



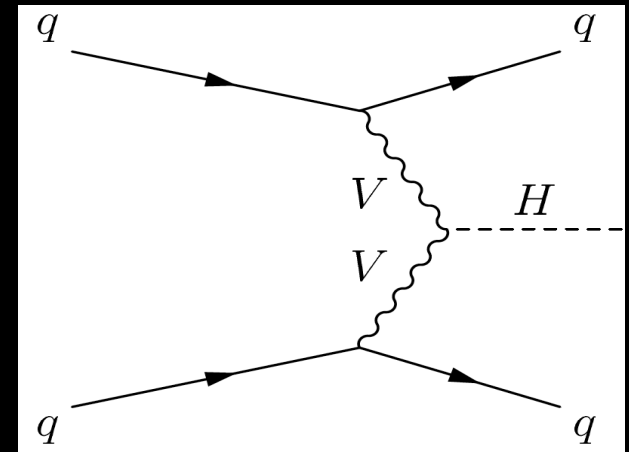
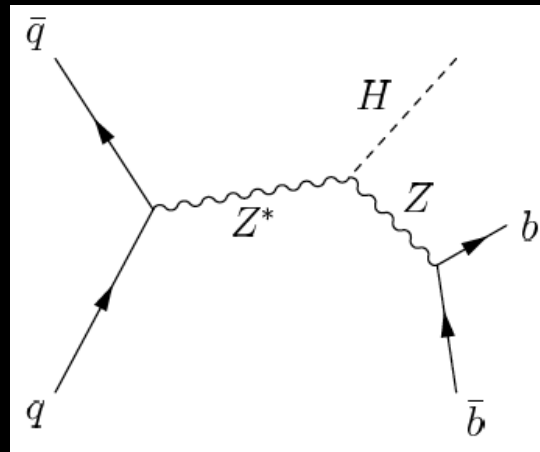
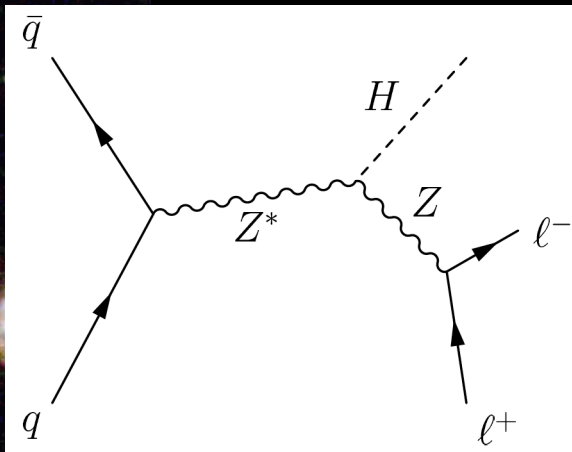
ATL-PHYS-
PUB-2013-014



Invisible Higgs decays



- Look for decays of Higgs boson to weakly interacting particles $B(H \rightarrow \text{inv})$ using VBF and ZH production.



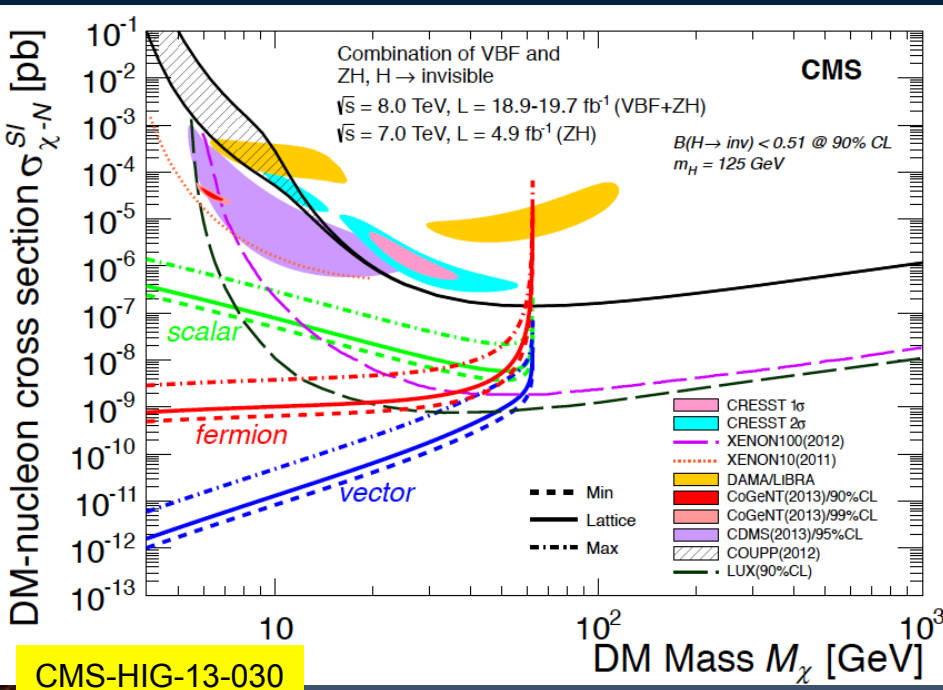
- Complementary to direct searches for dark matter.
 - Exploited channels:
 - ATLAS: $Z(\ell\ell) + \text{MET}$
 - CMS: $Z(\ell\ell, b\bar{b}) + \text{MET}$, VBF + MET
- Direct constraint on $B(H \rightarrow \text{inv})$ can also be obtained from global fit of measured decay modes



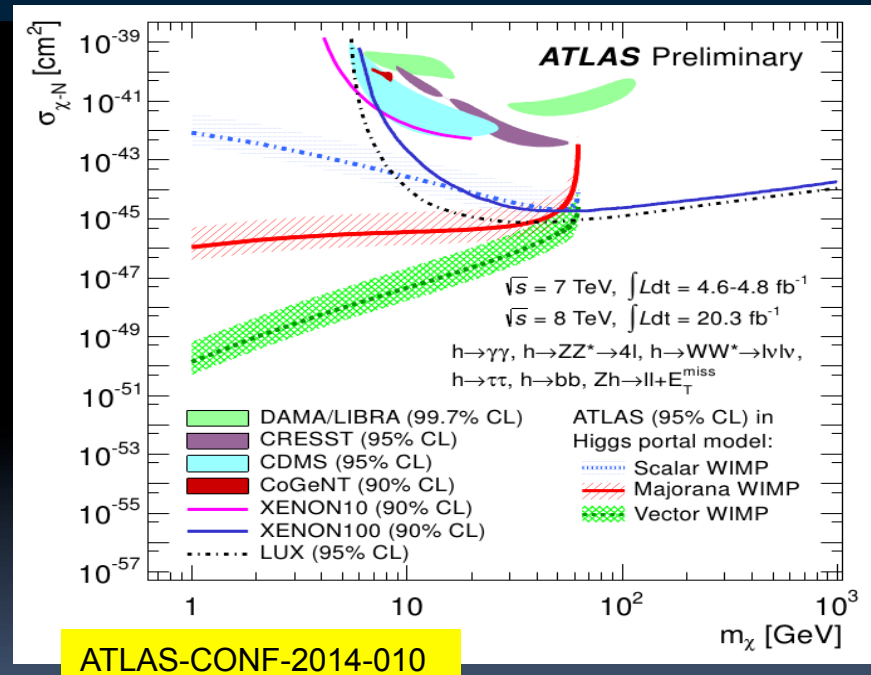
Invisible Higgs Decay



- CMS:
 - combination of Z(lℓ)+ MET, Z(bb)+MET and VBF + MET searches yields
 - Observed (expected) $B(H \rightarrow \text{inv})/\sigma_{\text{SM}} < 0.58 (< 0.44)$ @ 95% CL



- ATLAS:
 - Combination of direct and indirect results
 - Observed (expected) limit $B(H \rightarrow \text{inv}) < 0.41 (0.55)$ at 95% CL



- Re-Interpret $B(H \rightarrow \text{inv})$ limit in Higgs-portal model:
 - DM sector decoupled from SM, except for Higgs-mediated interactions with $m_\chi < m_{H/2}$

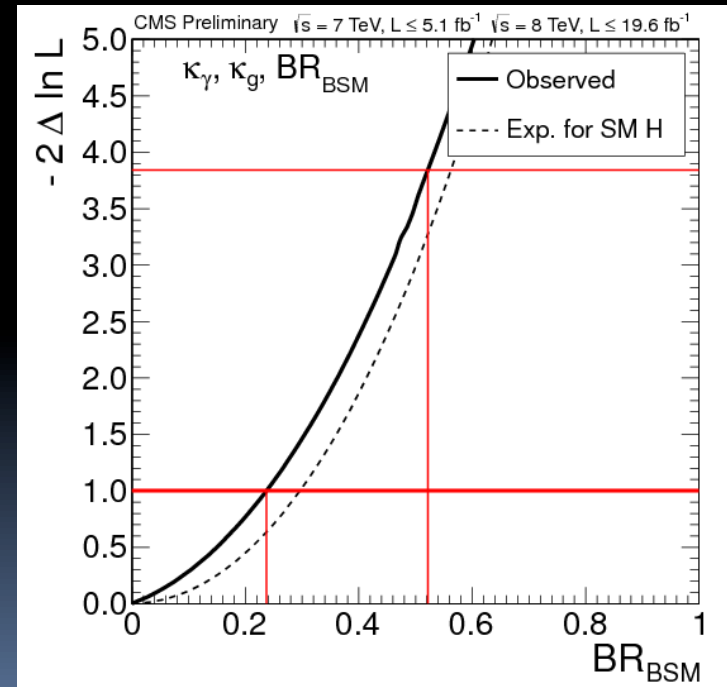
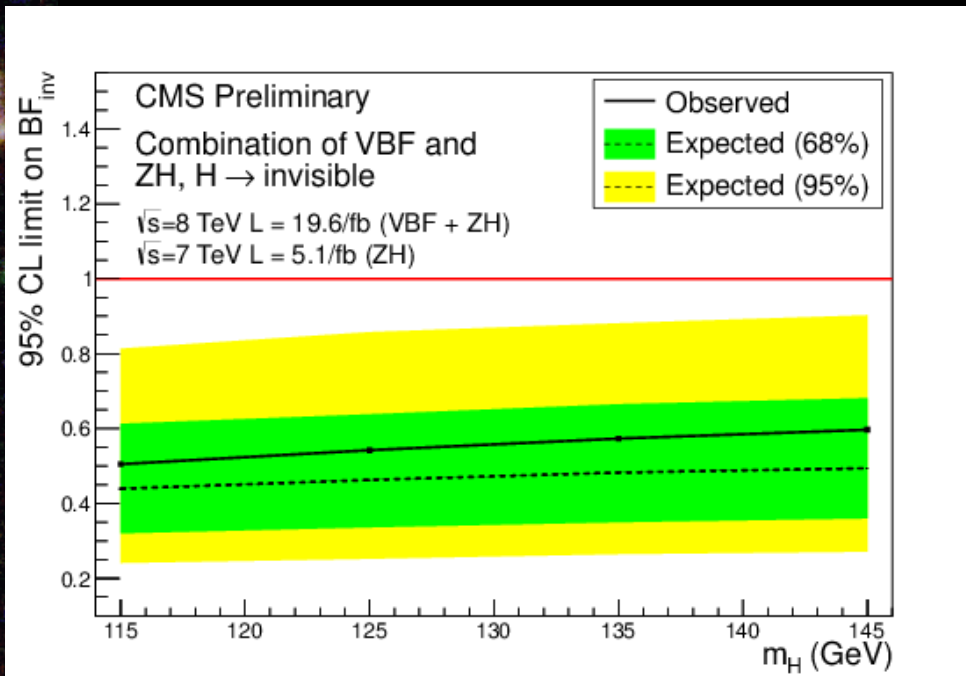
- Djouadi, A. Falkowski, Y. Mambrini, and J. Quevillon
- B. Patt and F. Wilczek



Direct and indirect constraints

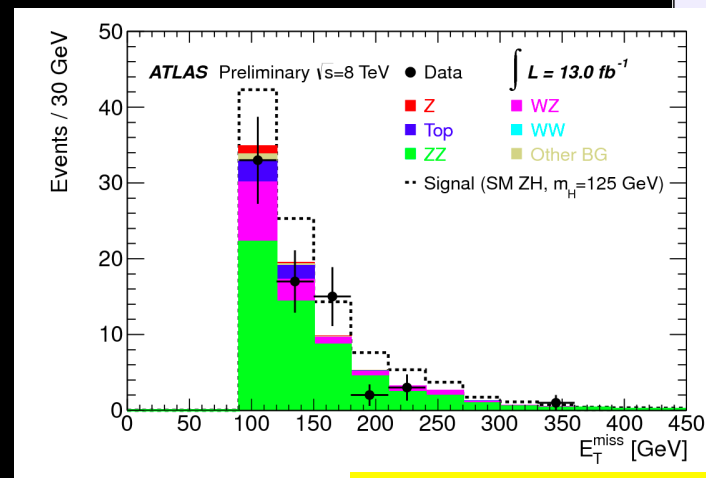
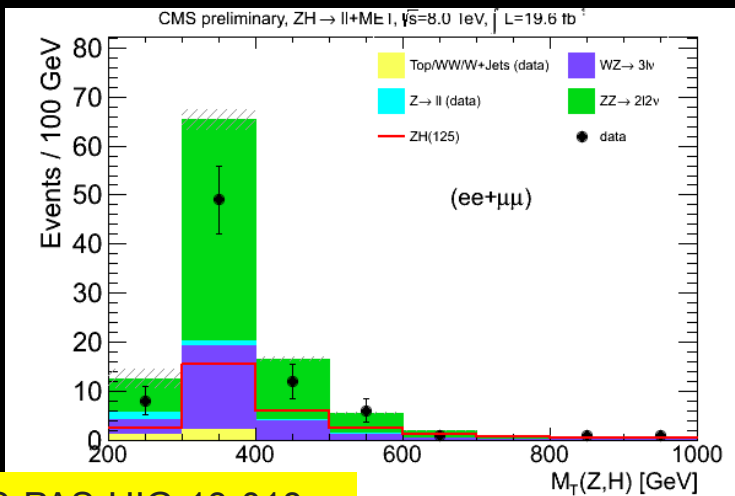
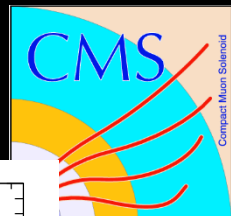


- Indirect constraint on BR_{inv} can also be obtained from global fit of measured decay modes.
- Fixing unmeasured modes to SM predictions and assuming $\kappa_V < 1$.
- Direct and indirect limits have comparable magnitudes



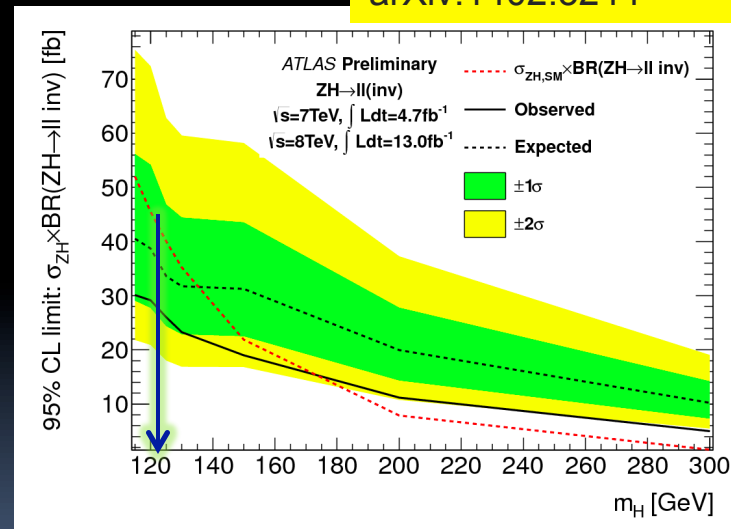
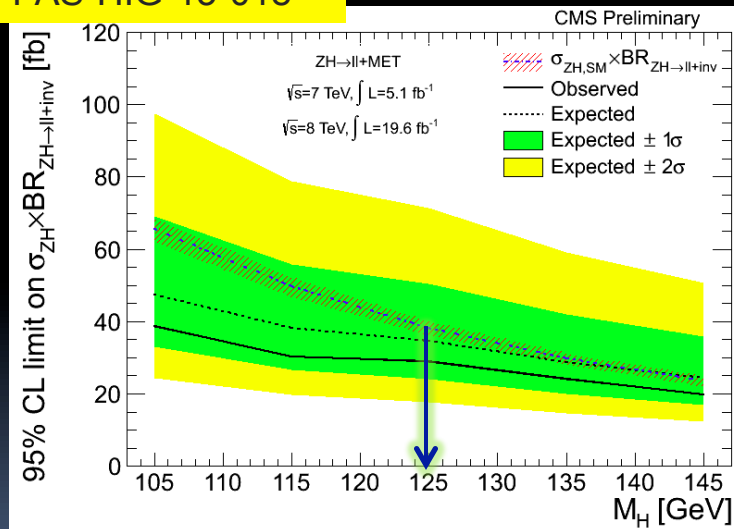


ZH \rightarrow ll+invisible



CMS-PAS-HIG-13-018

arXiv:1402.3244



Observed (expected)

CMS

ATLAS

B(H \rightarrow inv) at 125 GeV @95% CL

< 0.75 (0.91)

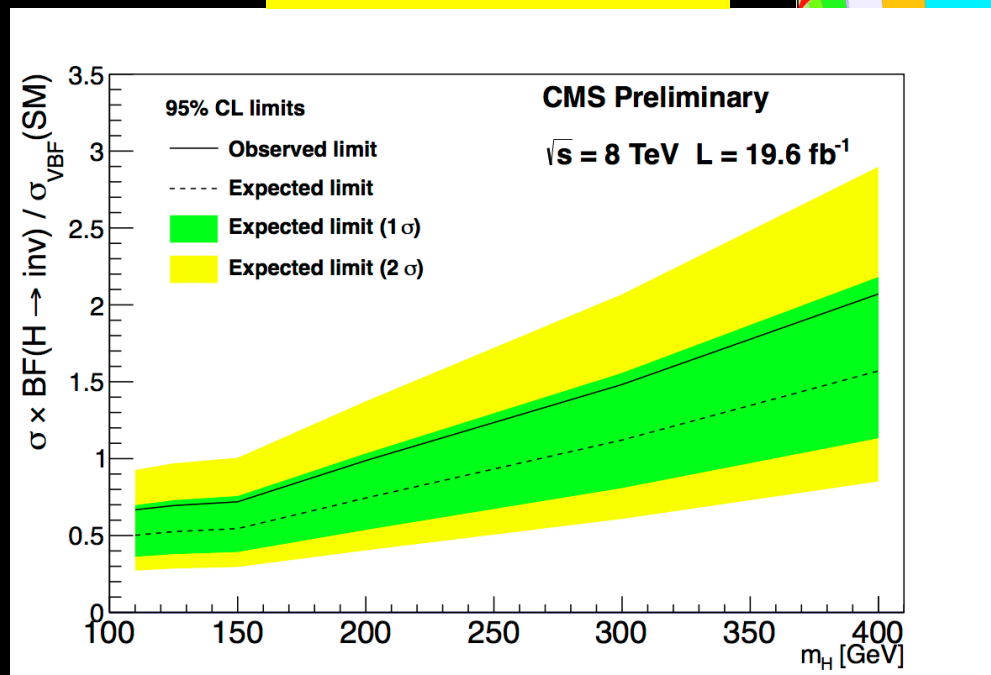
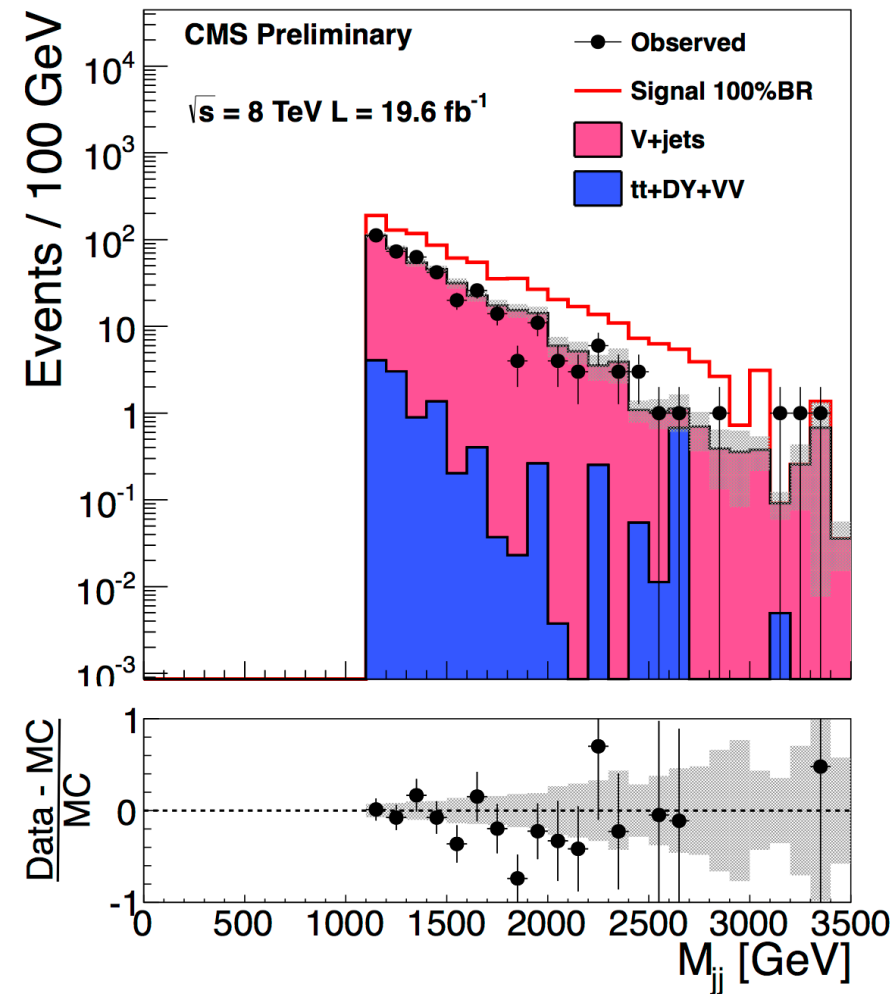
< 0.65 (0.84)



VBF $H \rightarrow$ Invisible



CMS-PAS-HIG-13-013

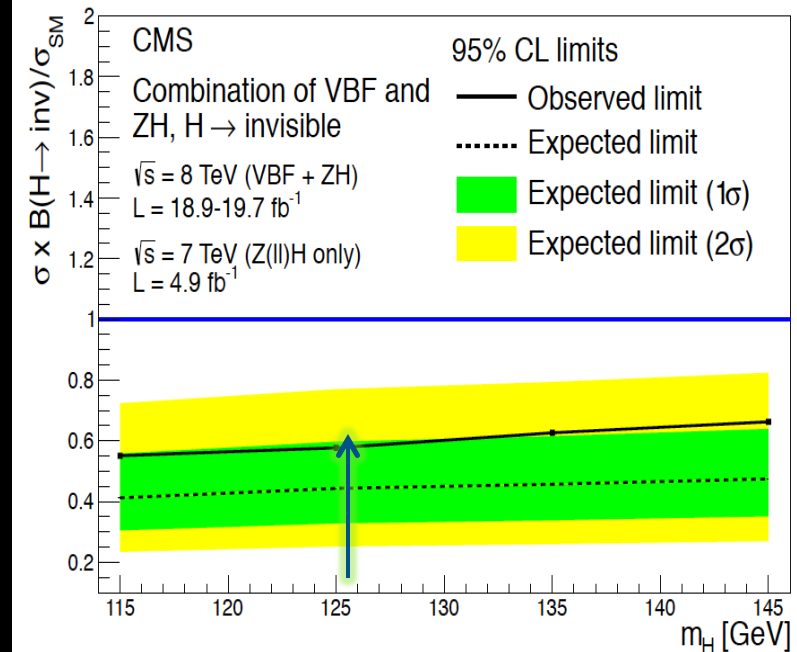


- At $m_H = 125 \text{ GeV}$ observed (expected) limits at 95% CL on $\sigma \times \text{B}(H \rightarrow \text{inv}) / \sigma_{\text{SM}} < 0.69$ (0.53) at 95%CLs

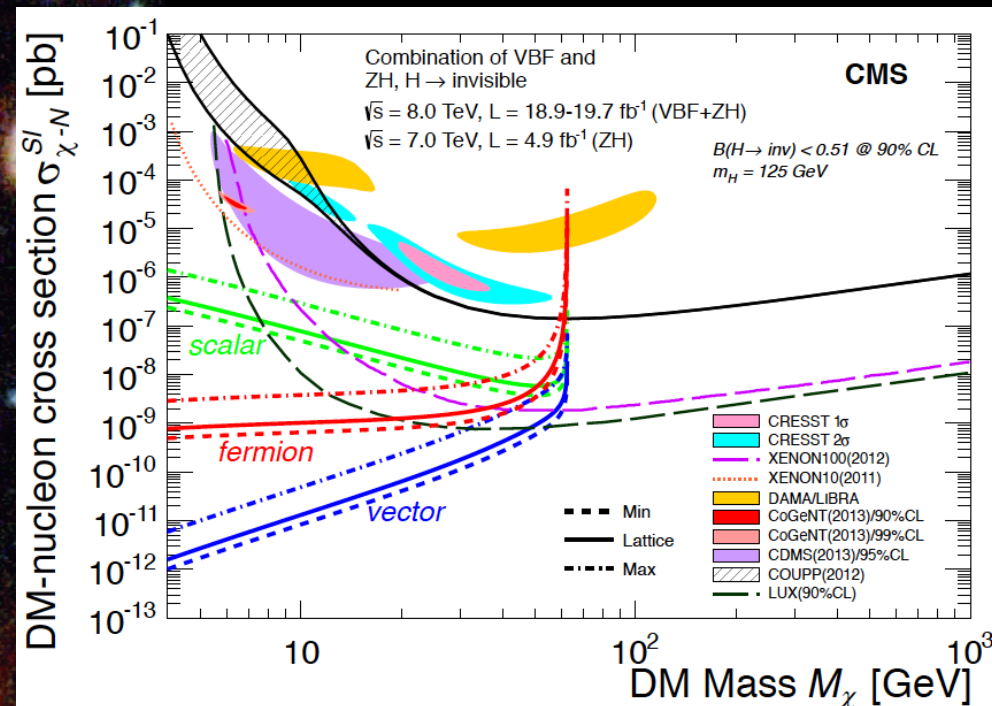


CMS combination

- Three individual CMS searches combined assuming the SM production cross section and acceptance
- $B(H \rightarrow \text{inv}) / \sigma_{\text{SM}} < 0.58$ (< 0.44) observed (expected) @ 95% CL



- $B(H \rightarrow \text{inv})$ for $m_H = 125 \text{ GeV}$ used to obtain upper limits at 90% CL on the DM-nucleon cross section as a function of the DM mass in Higgs-portal models of DM interactions



- Djouadi, A. Falkowski, Y. Mambrini, and J. Quevillon
- B. Patt and F. Wilczek

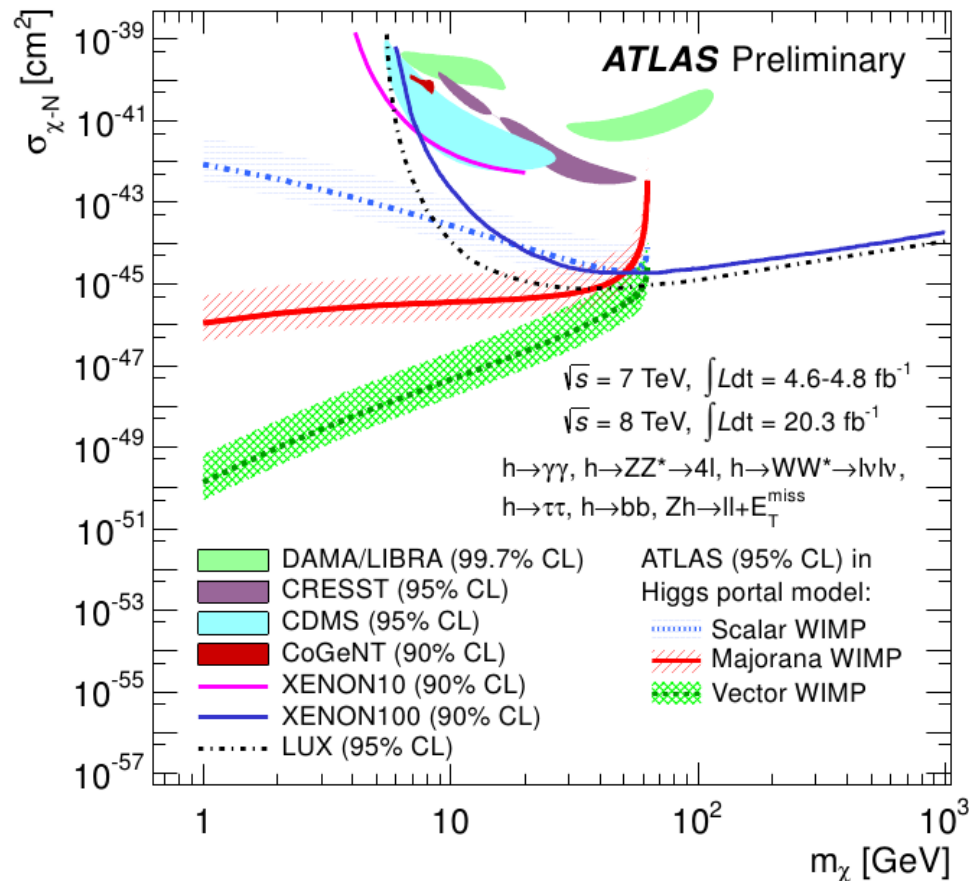
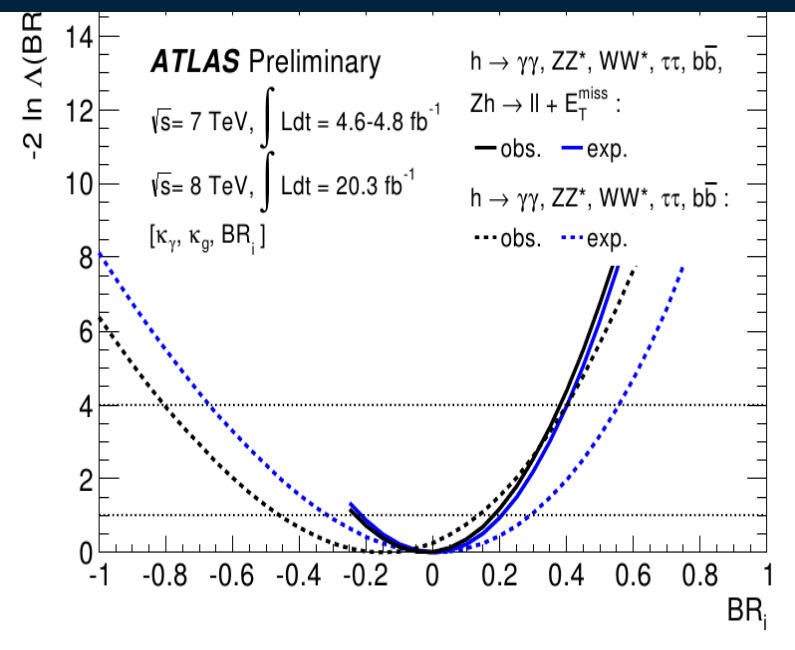


Invisible decays and dark matter



- Indirect constraint on $B(H \rightarrow \text{inv})$ can also be obtained from global fit of measured decay modes.
- Combination of direct and indirect results
- Observed (expected) limit $B(H \rightarrow \text{inv}) < 0.41(0.55)$ at 95% CL

- Interpret limit on $B(H \rightarrow \text{inv})$ as direct bounds on massive dark particles with $m_\chi < m_{H/2}$ coupling to the Higgs
- Interpretation in Higgs-portal model:
 - DM sector decoupled from SM, except for Higgs-mediated interactions



ATLAS-CONF-2014-011